

Quantum matter & black hole ringing

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Motivation – unconventional phases at finite density

- 1 Low temperature and finite density (bosons and fermions)
- 2 Two uses of magnetic fields
- 3 Experimental examples
- 4 Free fermions and bosons

Strongly coupled theories with gravity duals

- 1 The normal state
- 2 Large N magnetic susceptibility
- 3 Black hole instabilities (superconducting instabilities)
- 4 $1/N$ corrections to the free energy
- 5 Black hole ringing
- 6 Quantum oscillations

Motivation – unconventional phases at finite density

- 1 Low temperature and finite density (bosons and fermions)
- 2 Two uses of magnetic fields
- 3 Quantum oscillations in High - T_c superconductors
- 4 Quantum criticality under the dome in High - T_c superconductors
- 5 Free fermions
- 6 Free bosons

Low temperature and finite density

- Effective field theories in condensed matter physics often have a finite charge density.
- Weak coupling intuition at low temperatures and finite density:
 - Charged fermions: **Fermi surface** is built up.
 - Charged bosons: **condensation instabilities** (e.g. superconductivity).
- Weakly interacting low energy excitations about a condensate or Fermi surface are very well characterised.
- There seem to be materials where these descriptions do not work.
- Perspective of this talk: AdS/CFT gives a tractable theory with an exotic finite density ground state.

Two uses of magnetic fields

- Magnetic fields useful for probing both fermions and bosons.
- de Haas - van Alphen effect: a Fermi surface leads to oscillations in the magnetic susceptibility as a function of $1/B$.
 - In a magnetic field

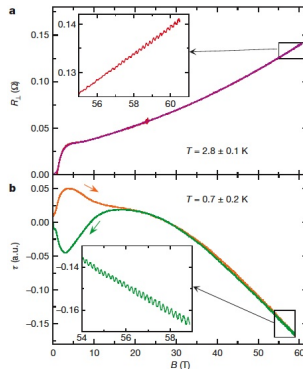
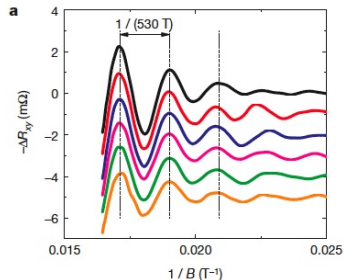
$$[P_x, P_y] \sim iB \quad \Rightarrow \quad \oint P_x dP_y \sim 2\pi(l + \frac{1}{2})B.$$

- When the area of the orbit is a cross section of the Fermi surface there is a sharp response. I.e. at $1/B \sim l/A_F \sim l/k_F^2 \sim l/\mu^2$.
- Large magnetic field will suppress superconducting instabilities.

Quantum oscillations in High - T_c superconductors

Doiron-Leyraud et al. 2007 (Nature), Vignolle et al. 2008 (Nature).

- de Haas - van Alphen oscillations in underdoped and overdoped cuprates.

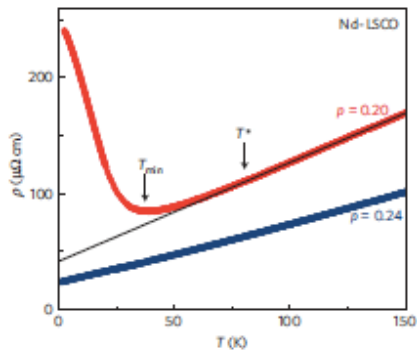


- In underdoped region, carrier density much lower than naïve expectation: “small Fermi surface”.

Criticality under the dome in High - T_c superconductors

Daou et al. 2008 (Nature Physics)

- Resistivity in 'normal phase' linear in temperature (anomalous).
- Applying a large magnetic field shows persistence down to $T = 0$ at critical doping.



Free fermions

- Free bosons or fermions in magnetic fields have Landau levels

$$\varepsilon_\ell = \sqrt{m^2 + 2|qB|(\ell + \frac{1}{2})}.$$

- Free energy for fermions (D=2+1)

$$\Omega = -\frac{|qB|AT}{2\pi} \sum_\ell \sum_{\pm} \log \left(1 + e^{-(\varepsilon_\ell \pm q\mu)/T} \right).$$

- Zero temperature limit

$$\lim_{T \rightarrow 0} \Omega = -\frac{|qB|A}{2\pi} \sum_\ell (q\mu - \varepsilon_\ell) \theta(q\mu - \varepsilon_\ell).$$

- Magnetic susceptibility has oscillations

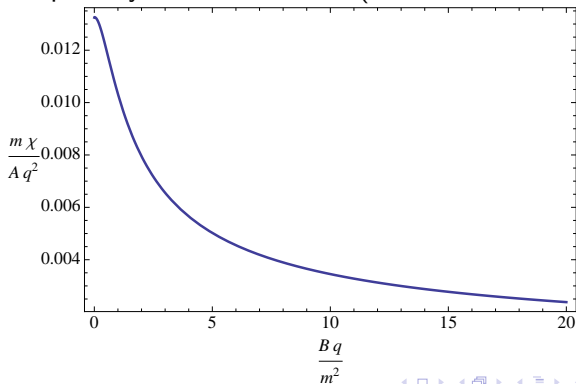
$$\chi \equiv \frac{\partial^2 \Omega}{\partial B^2} = \frac{|qB|A}{2\pi} \sum_\ell \frac{q^2(\ell + \frac{1}{2})^2}{\varepsilon_\ell^2} \delta(q\mu - \varepsilon_\ell) + \dots,$$

Free bosons

- Free energy for bosons – unstable if $\varepsilon_0 < |q\mu|$

$$\Omega = \frac{|qB|A}{2\pi} \sum_{\ell} \sum_{\pm} \log \left(1 - e^{-(\varepsilon_{\ell} \pm q\mu)/T} \right) + \Omega|_{T=0} .$$

- Magnetic susceptibility at $T=0$ if stable (Hurwitz zeta function)



Strongly coupled theories with gravity duals

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The normal state

- The minimal ingredient is Einstein-Maxwell theory

$$S_E[A, g] = \int d^4x \sqrt{g} \left[-\frac{1}{2\kappa^2} \left(R + \frac{6}{L^2} \right) + \frac{1}{4g^2} F^2 \right].$$

- The 'normal state' is dual to a dyonic black hole

$$ds^2 = \frac{L^2}{r^2} \left(f(r) d\tau^2 + \frac{dr^2}{f(r)} + dx^i dx^i \right),$$

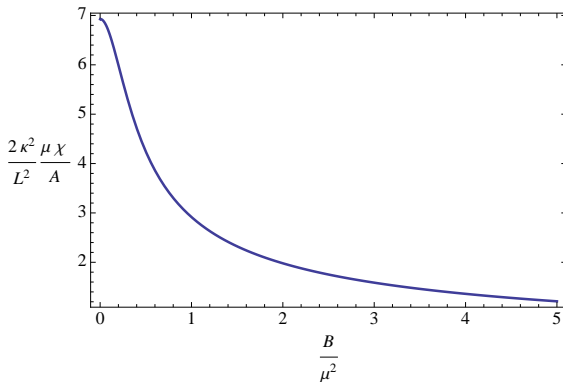
$$A = i\mu \left[1 - \frac{r}{r_+} \right] d\tau + B \times dy.$$

- Free energy is the action evaluated on shell

$$\Omega_0 = -\frac{AL^2}{2\kappa^2 r_+^3} \left(1 + \frac{r_+^2 \mu^2}{\gamma^2} - \frac{3r_+^4 B^2}{\gamma^2} \right).$$

Large N magnetic susceptibility

- Easy to compute $\chi \equiv \frac{\partial^2 \Omega_0}{\partial B^2}$
- Plot result:



- Looks just like free bosons.... (but massless!)

Black hole superconducting instabilities

- Add some matter to the bulk to make things more interesting...
- Charged bosons:

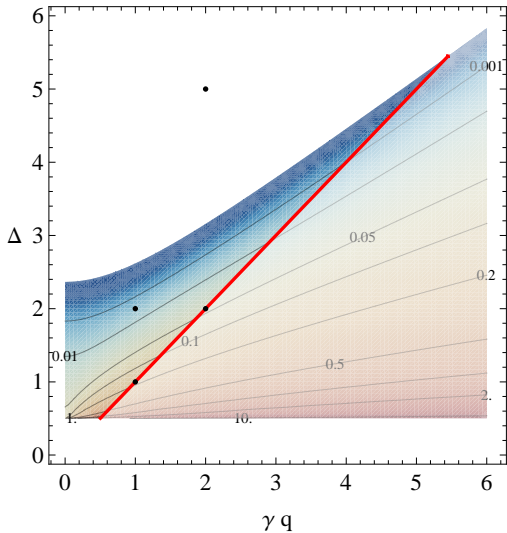
$$S_E[\phi] = \int d^4x \sqrt{g} \left[|\nabla\phi - iqA\phi|^2 + m^2|\phi|^2 \right].$$

- Charged fermions:

$$S_E[\psi] = \int d^4x \sqrt{g} \left[\bar{\psi} \Gamma \cdot \left(\partial + \frac{1}{4} \omega_{ab} \Gamma^{ab} - iqA \right) \psi + m \bar{\psi} \psi \right].$$

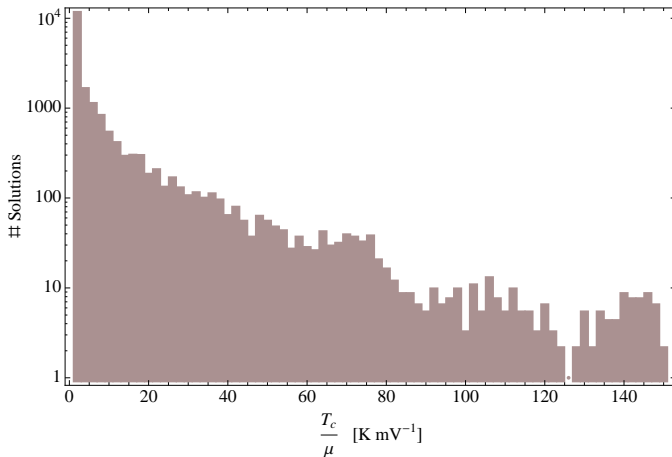
- Bosons: Criterion for homogeneous instability at $T = 0$

$$q^2 \gamma^2 \geq 3 + 2m^2, \quad \gamma^2 = \frac{2g^2 L^2}{\kappa^2}.$$



[Denef-SAH '09]

Landscape of superconducting membranes



[Denef-SAH '09]

1/N corrections to the free energy

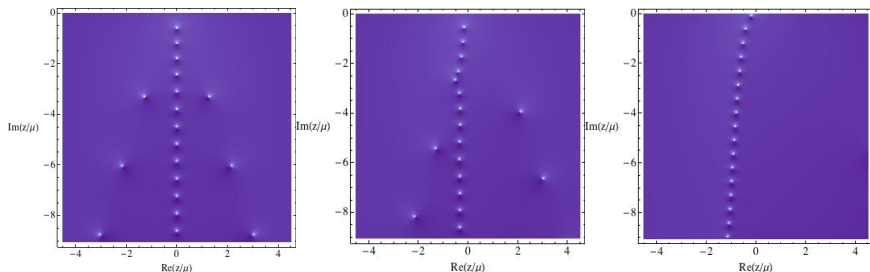
- Suppose no superconducting instability (or suppress with B)
- Nontrivial Landau-level structure subleading in $1/N$?
⇒ Quantum contribution from charged matter:

$$\Omega_{1\text{-loop}} = T \text{tr} \log \left[-\hat{\nabla}^2 + m^2 \right] - T \text{tr} \log \left[\Gamma \cdot \hat{D} + m \right] + \dots$$

- It is difficult to compute determinants in black hole backgrounds and it is hardly ever done...
- Reformulate the problem using quasinormal modes.

Black hole ringing

- Late times: a perturbed black hole ‘rings’ with characteristic frequencies.
- Quasinormal modes: poles of the retarded Green’s function (bulk or boundary).
- Some typical quasinormal for charged AdS black holes at low temperature (not easy to make these plots!)



The free energy and quasinormal modes

- We derived (new to my knowledge) formulae for the determinant as a sum over quasinormal modes $z_*(\ell)$ of the black hole

$$\Omega_{1\text{-loop, B}} = \frac{|qB|AT}{2\pi} \sum_{\ell} \sum_{z_*(\ell)} \log \left(\frac{|z_*(\ell)|}{2\pi T} \left| \Gamma \left(\frac{iz_*(\ell)}{2\pi T} \right) \right|^2 \right).$$

$$\Omega_{1\text{-loop, F}} = -\frac{|qB|AT}{2\pi} \sum_{\ell} \sum_{z_*(\ell)} \log \left(\left| \Gamma \left(\frac{iz_*(\ell)}{2\pi T} + \frac{1}{2} \right) \right|^2 \right).$$

- For the BTZ black hole we did the sum explicitly and checked agreement with the known result.

Quantum oscillations

- The power of these formulae is that if an individual quasinormal mode does something non-analytic, then this is directly identified.
- Faulkner-Liu-McGreevy-Vegh have shown that at zero temperature there is a fermion quasinormal mode whose trajectory in the complex frequency plane bounces off the real axis at $k = k_F$.
- At a finite magnetic field, this gives a bounce when $2B\ell = k_F^2$.
- The contribution of this mode to the free energy is

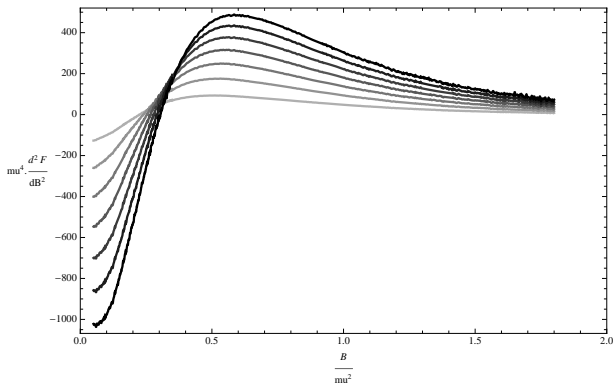
$$\Omega = -\frac{qBA}{2\pi^2} \log \frac{\mu}{T} \text{Im } z_*$$

- Taking two derivatives, delta functions in the magnetic susceptibility with period

$$\Delta \left(\frac{1}{B} \right) = \frac{2\pi q}{A_F} .$$

Preliminary comments on bosons

- It looks possible that stable bosons can also give oscillations if a zero temperature bosonic quasinormal mode crosses a branch cut.
- Summing numerically over the bosonic modes does not converge very quickly but may indicate some structure



Conclusions

- There exist systems with finite charge density that are described as neither conventional Fermi liquids or superfluids.
- AdS/CFT provides model exotic stable finite density systems.
- Magnetic fields are an essential experimental and theoretical tool for probing such systems.
- There is interesting structure at $1/N$ in AdS/CFT related to Landau levels for fermions and bosons.
- Found a method for computing determinants about black holes using quasinormal modes.
- Fermionic loops are shown to give de Haas - van Alphen oscillations.