

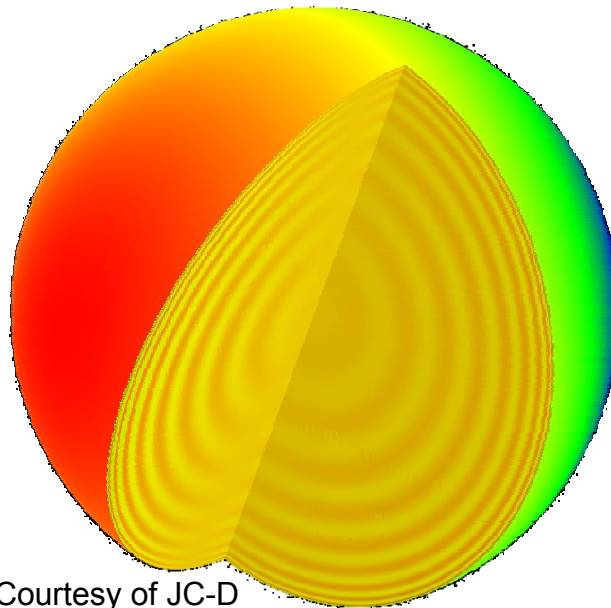


Stellar pulsations as convection probes

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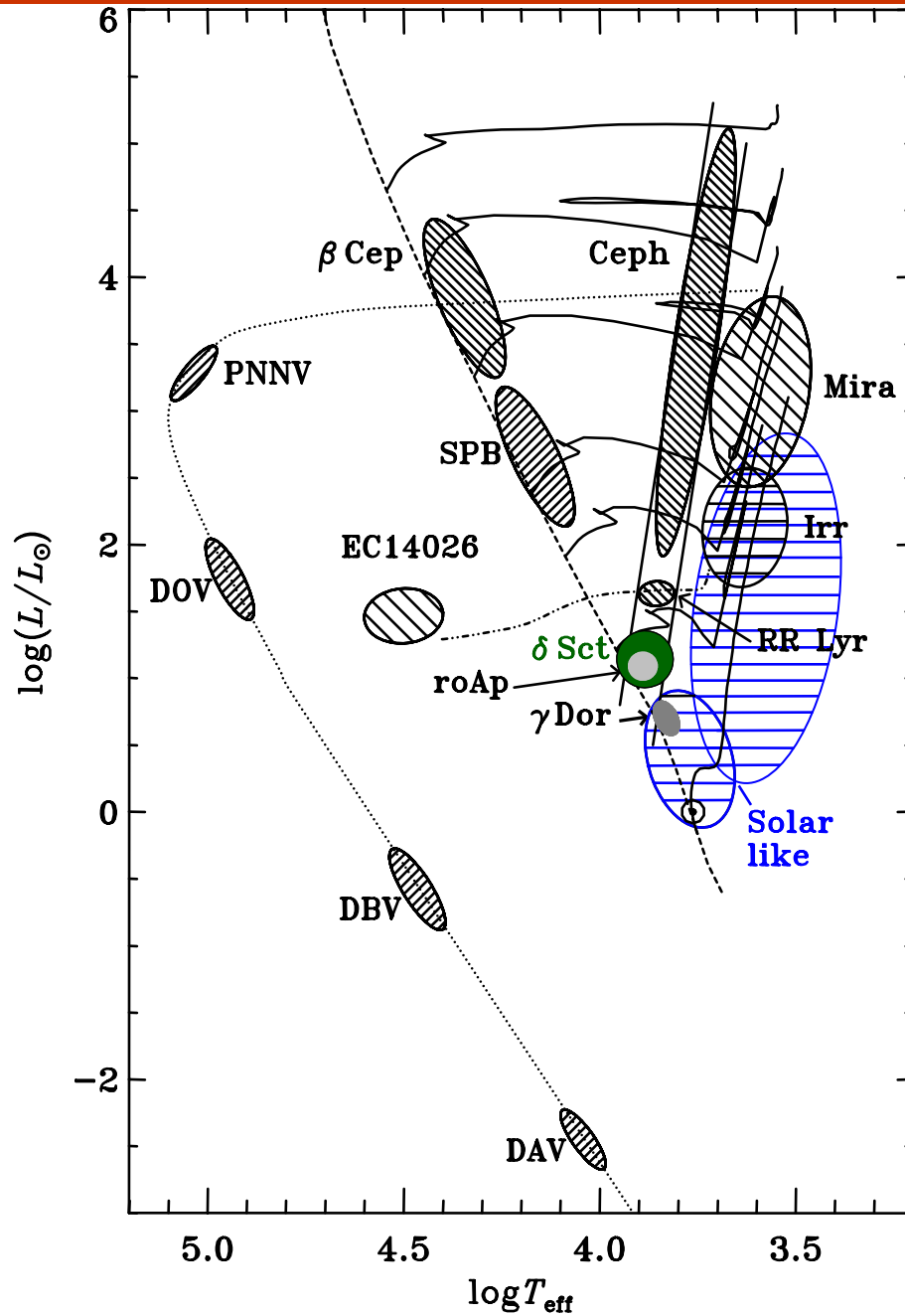


Courtesy of JC-D

The impact of asteroseismology across stellar astrophysics

Santa Barbara – October 2011

Pulsating H-R diagram



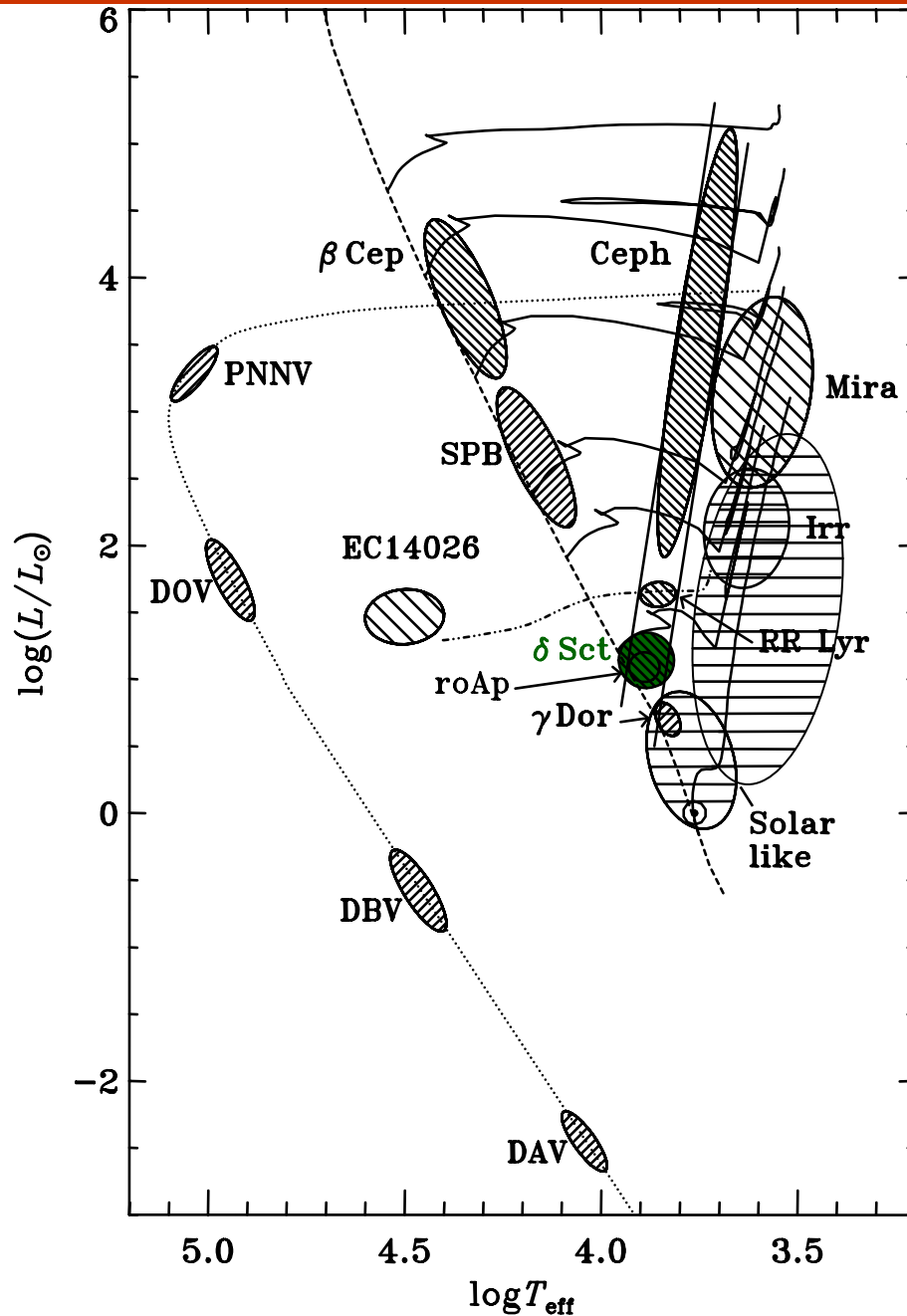
δ Scuti stars

- ▶ location of cool edge of IS
- ▶ acoustic radiation

Solar-like pulsators

- ▶ surface effects
- ▶ amplitudes

δ Scuti stars



- ▶ Central or shell hydrogen burning phase
- ▶ $1.5 M_{\odot} < M < 2.5 M_{\odot}$
- ▶ $0.6 < \log L/L_{\odot} < 2.0$
- ▶ Multiperiodic behaviour (e.g. FG Vir: 79 modes)
- ▶ Amplitudes from $10^{-3} \rightarrow 0.1$ mag
- ▶ low-order p modes
- ▶ Periods between 18 min and 8 h
- ▶ Driven by the κ mechanism in the HeII ionization zone

δ Scuti stars

Overview of selected **time-dependent convection models** (applications by)

Unno (1967, 89) \rightarrow Gabriel (1998) \rightarrow Grigahcène et al. (2004) : Dupret; Théado

Nonlinear mixing-length equations for a **Boussinesq fluid**
Infinite lifetime of fluid elements

Gough (1965, 77) \rightarrow Balmforth (1992) : Baker; Balmforth; Cunha; Gough; GH

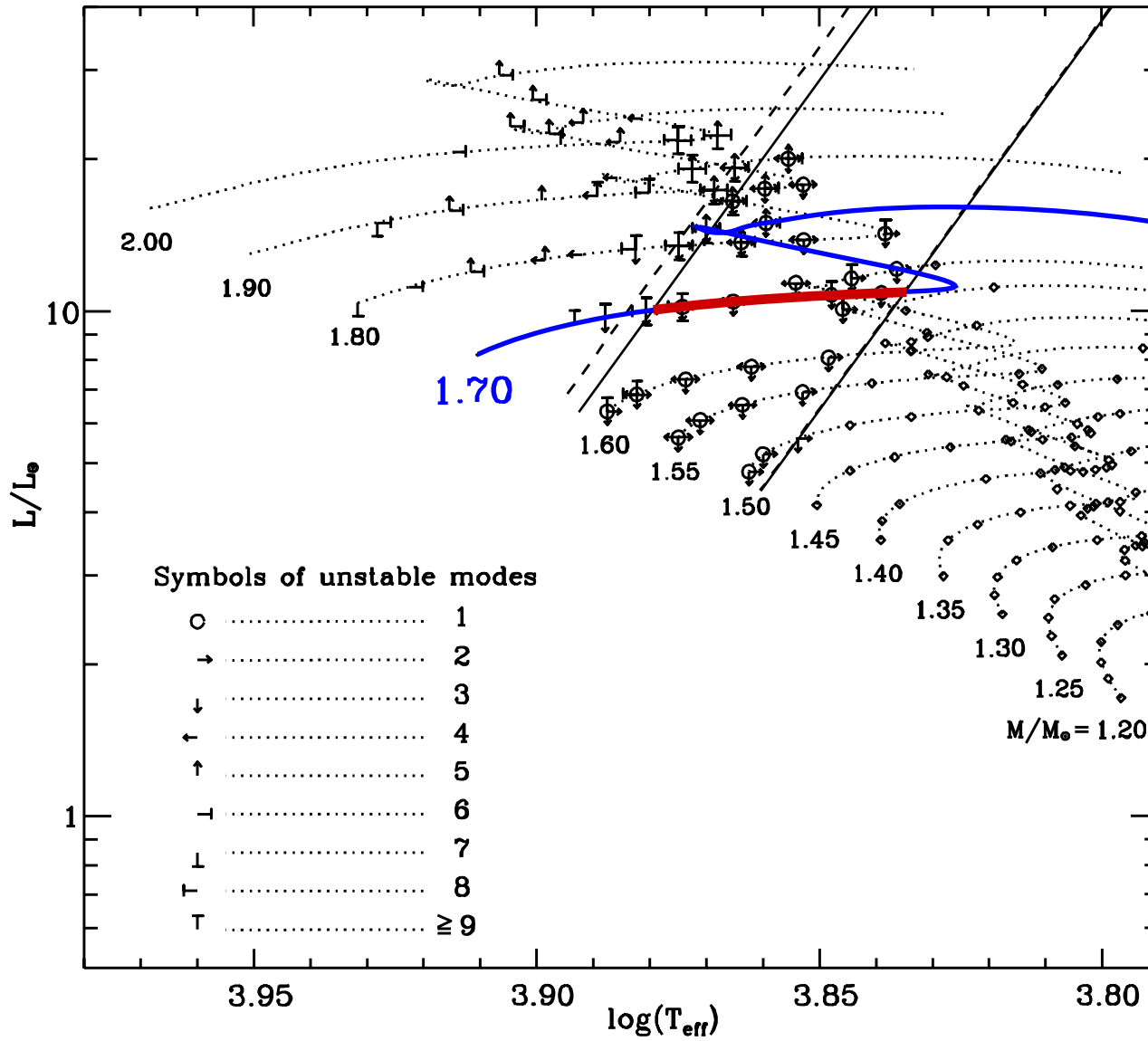
Linearized mixing-length equations for a **Boussinesq fluid**
Finite lifetime of fluid elements (linear growth rates)

Xiong (1977, 1989) : Cheng; Deng; Xiong

Reynold's transport equations for a **Boussinesq fluid**
Third-order moments approximated with diffusion-like expressions
using parametrized length scales (closure coefficients)

\longrightarrow **estimates of turbulent flux perturbations:** δF_c ; δp_t

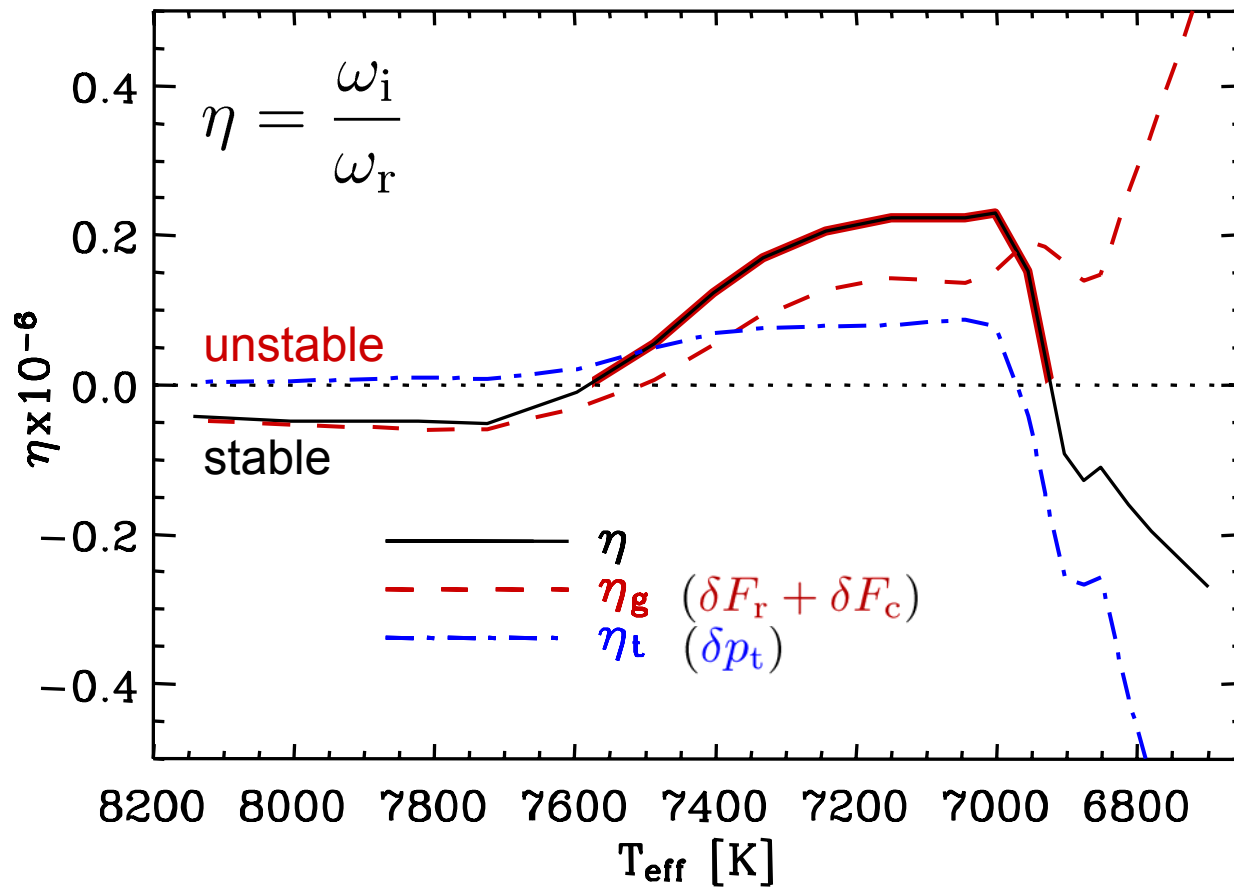
δ Scuti stars



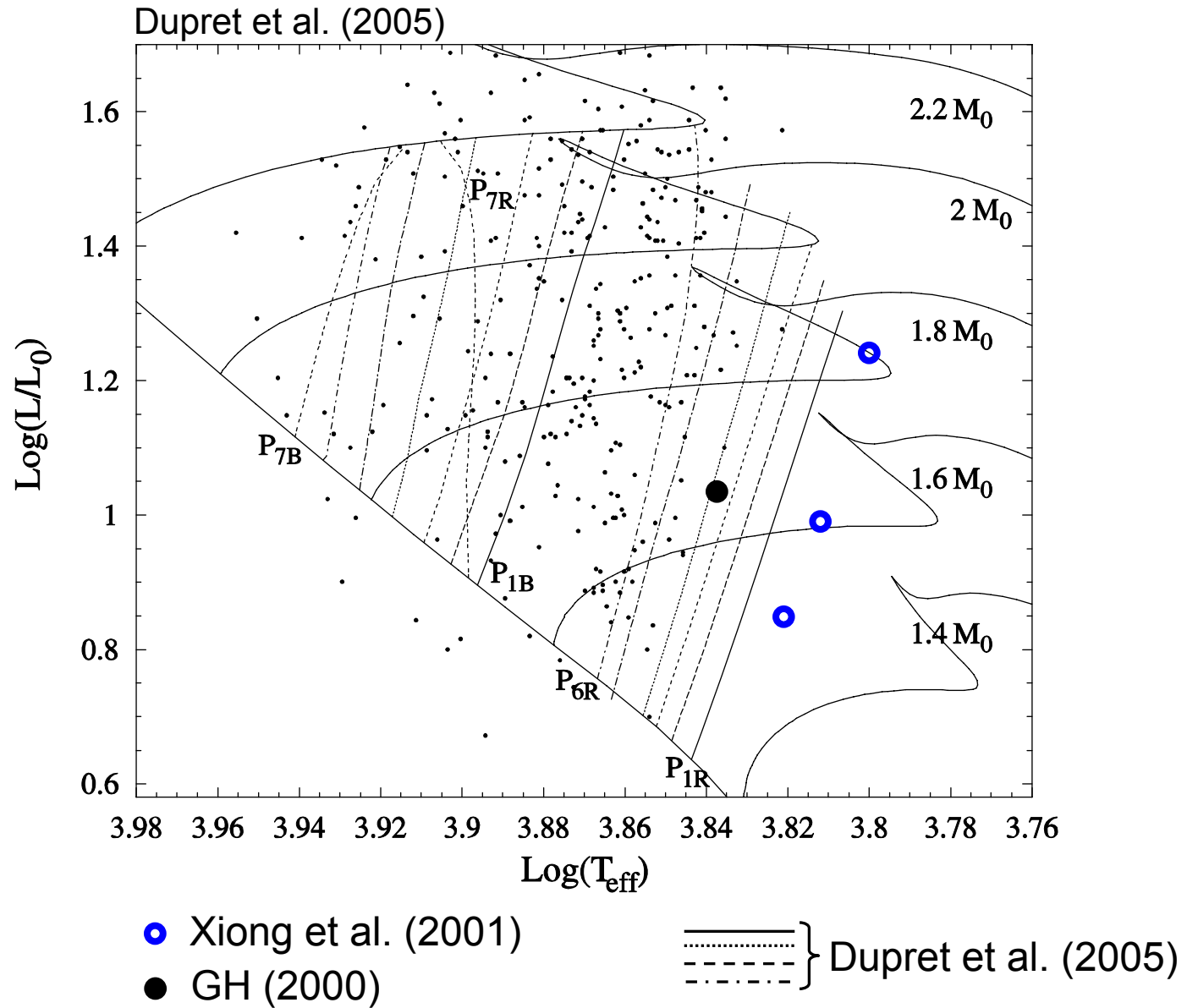
GH et al. (1999)

δ Scuti stars

1.7 M_⊙ δ Scuti model evolving during H-core burning phase



δ Scuti stars

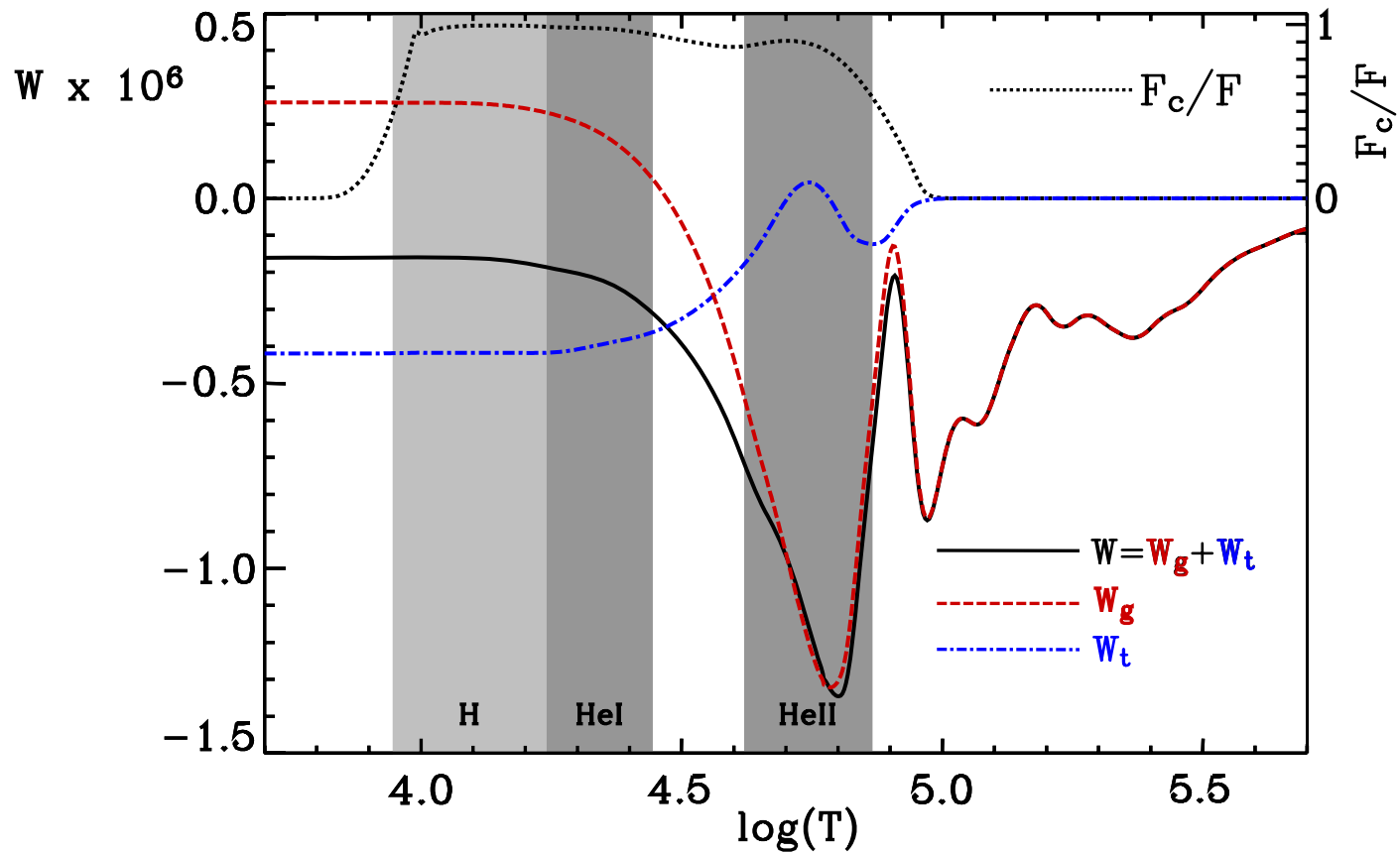


δ Scuti stars

Accumulated work integral

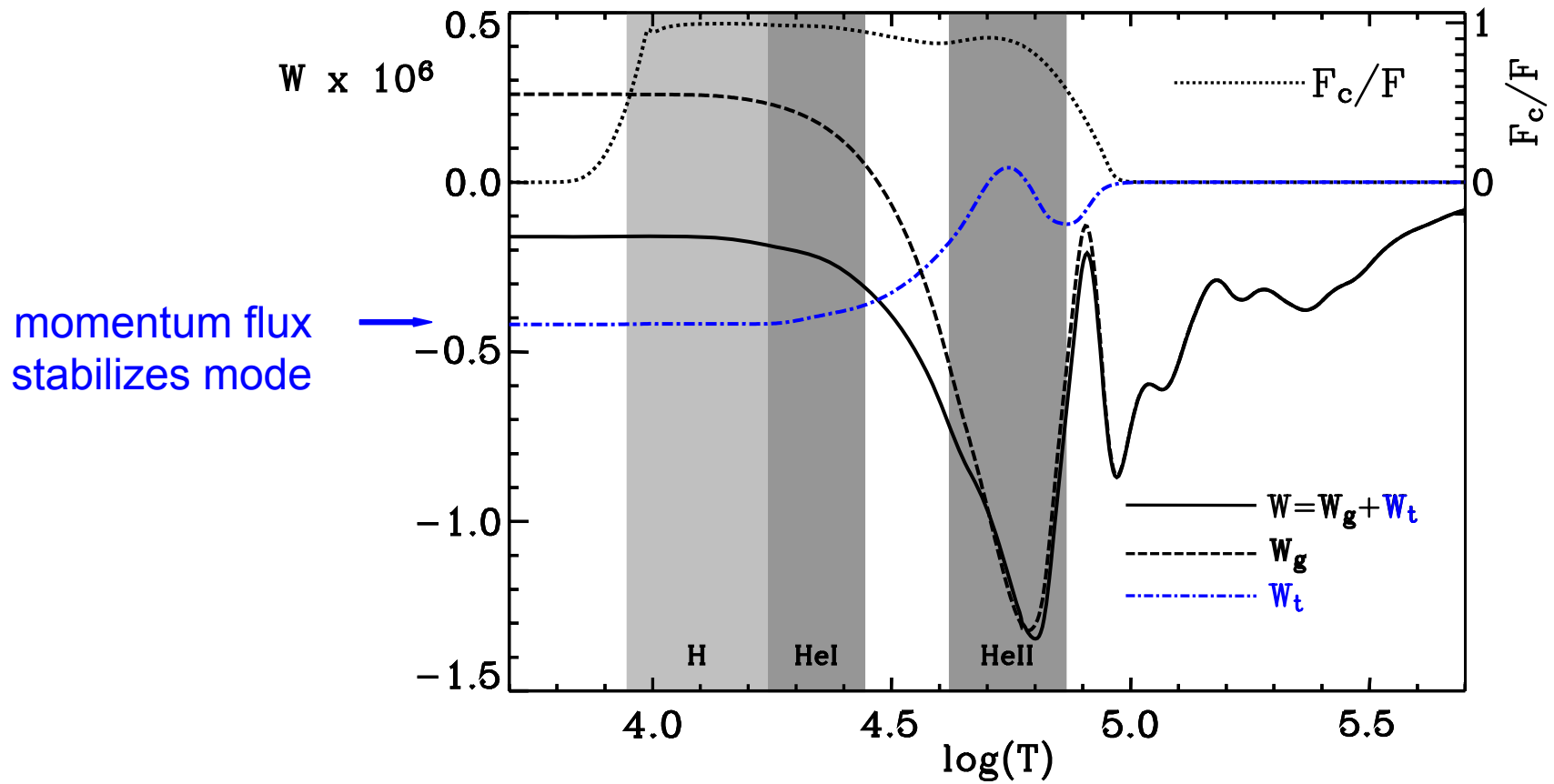
$$W_g = -\text{Im} \left[\int_M \delta\rho^* \delta p_g \rho^{-2} dm \right]$$

$$W_t \simeq -\text{Im} \left[\int_M \delta\rho^* \delta p_t \rho^{-2} dm \right]$$



δ Scuti stars

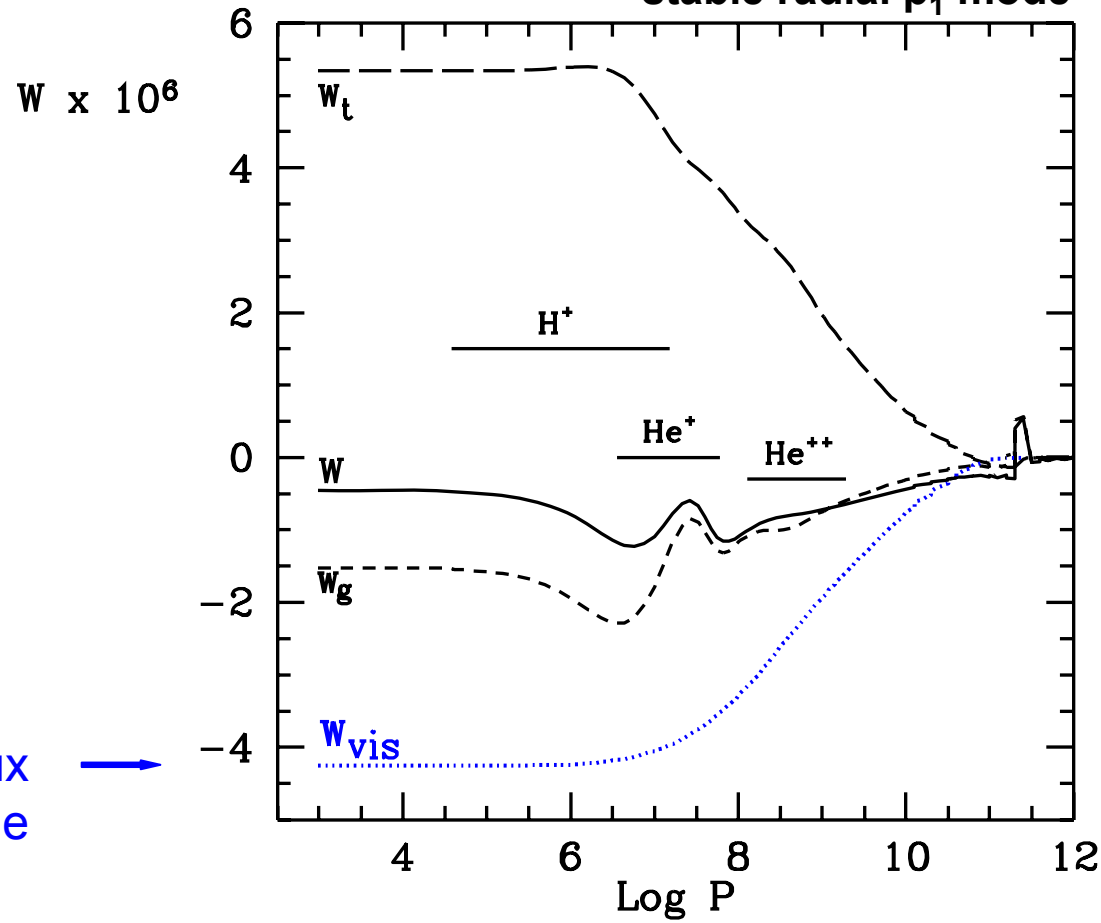
Accumulated work integral



δ Scuti stars

Accumulated work integral

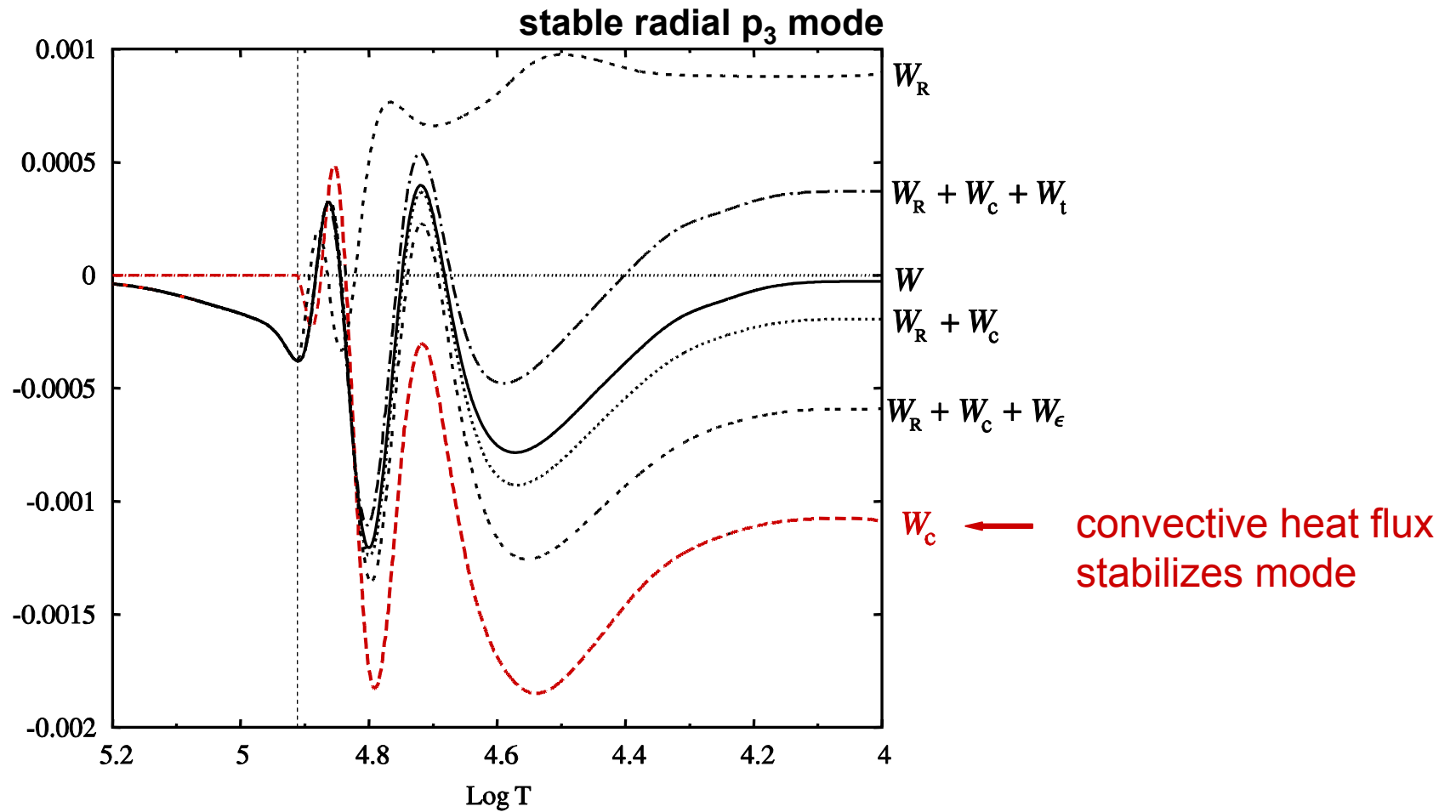
stable radial p_1 mode



viscous flux
stabilizes mode

δ Scuti stars

Accumulated work integral



Dupret et al. (2005)

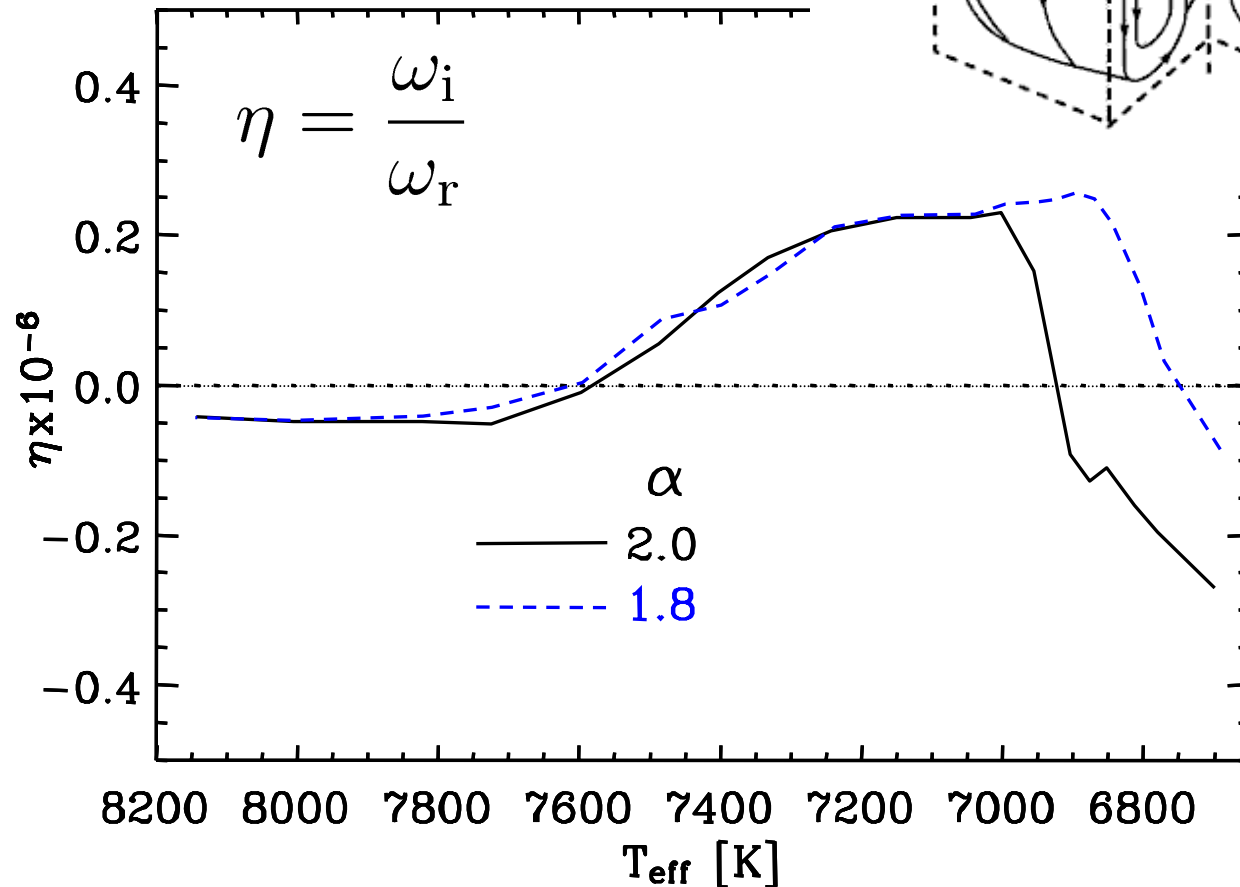
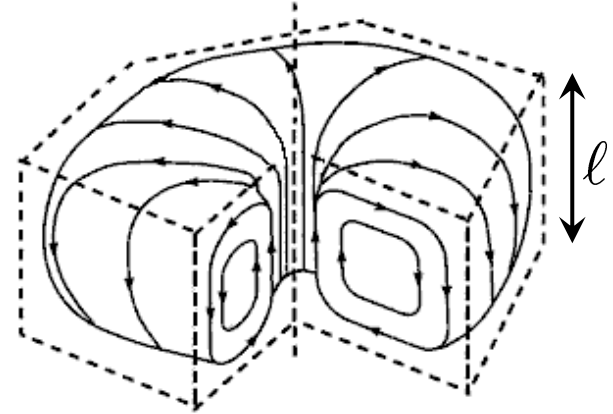
δ Scuti stars

Acoustic radiation in the equilibrium model

(Houdek & Gough 1998; Houdek 2000; Samadi, Goupil, Houdek 2002)

overturning eddy:
$$\frac{2w^2}{\ell} = g \frac{\tilde{\delta}}{T} T' - \frac{P_{ac}}{\rho w}$$

acoustic drag



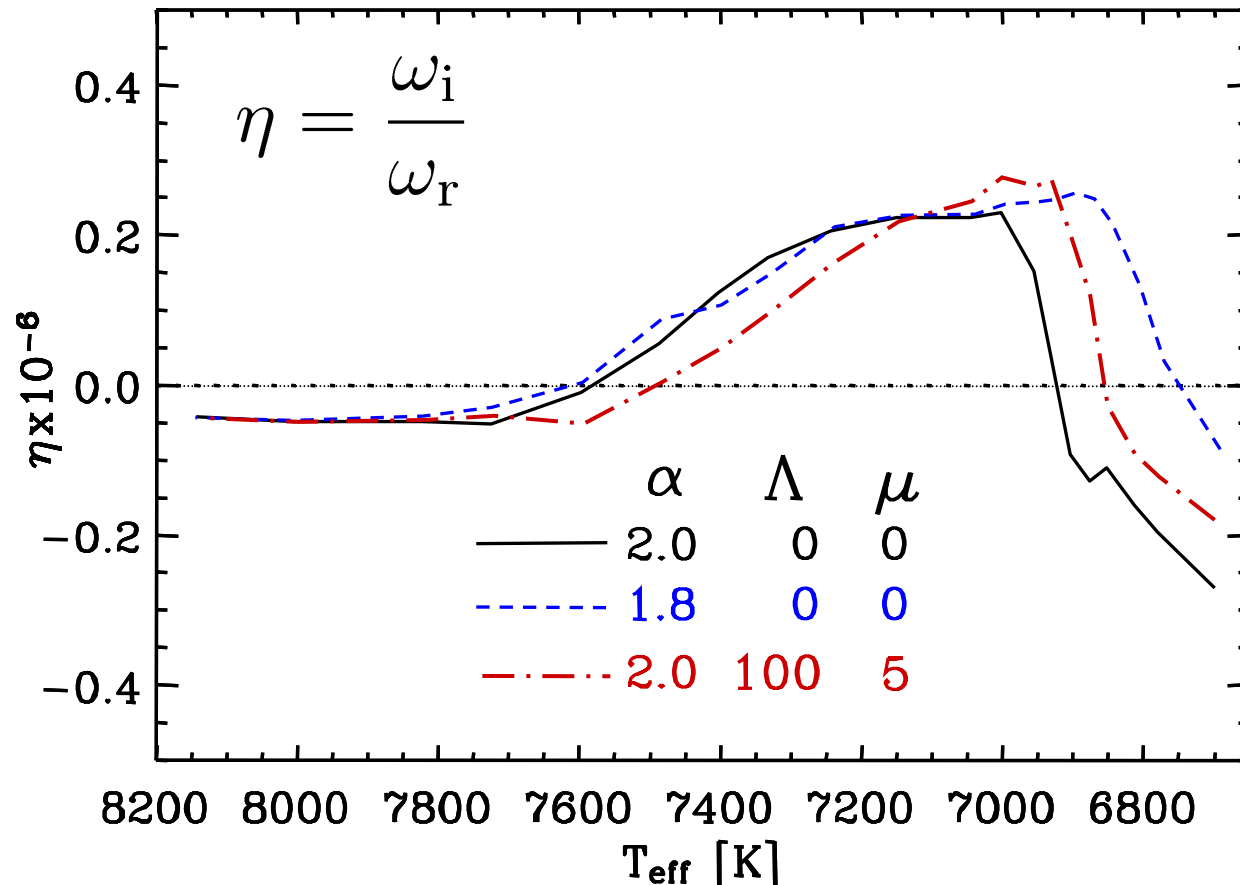
δ Scuti stars

Acoustic radiation in the equilibrium model

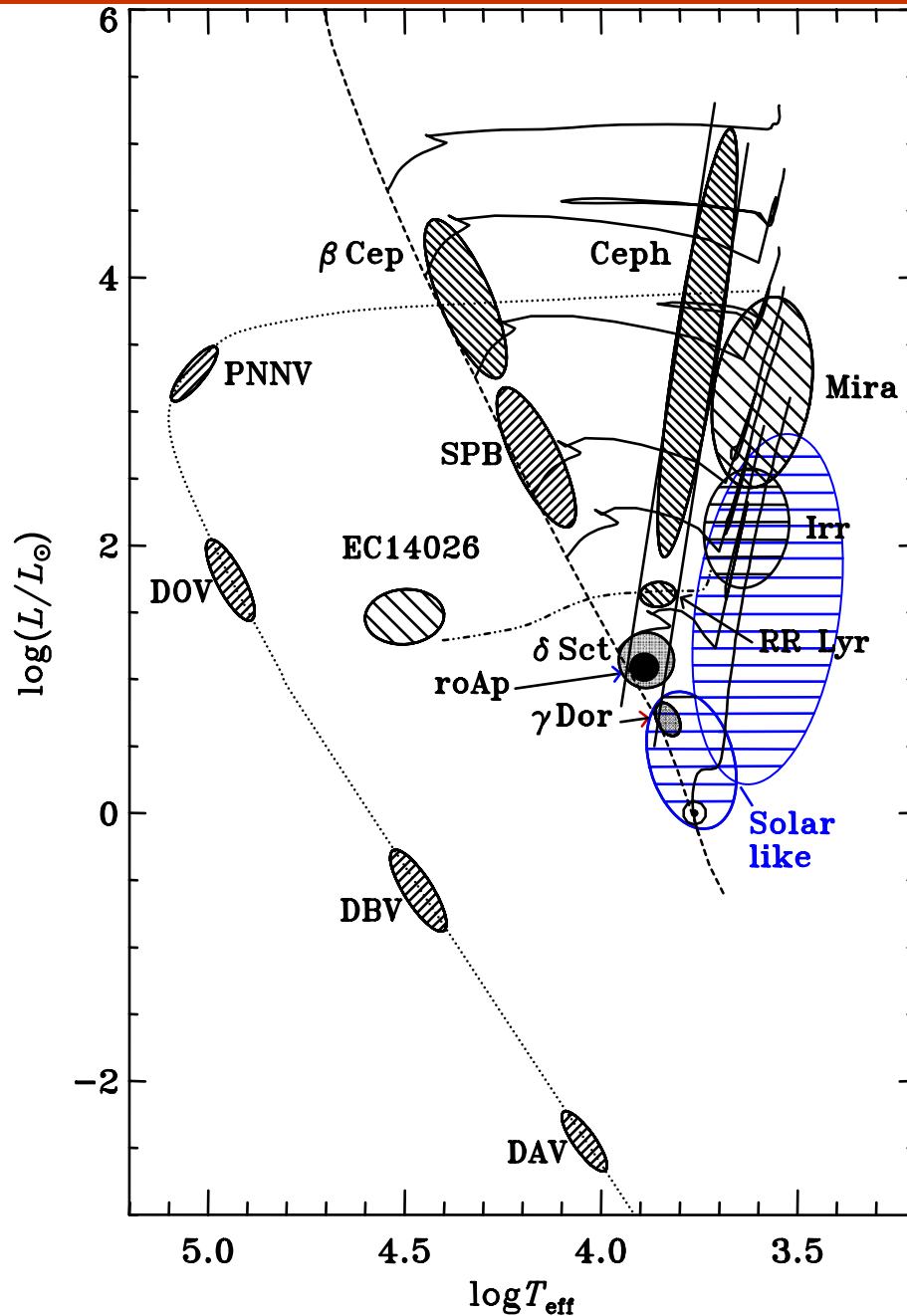
(Houdek & Gough 1998; Houdek 2000; Samadi, Goupil, Houdek 2002)

overturning eddy: $\frac{2w^2}{\ell} = g \frac{\tilde{\delta}}{T} T' - \frac{P_{ac}}{\rho w}$

Lighthill-Proudman: $P_{ac} = \Lambda \frac{\rho w^3}{\ell} \left(\frac{w}{c}\right)^\mu$



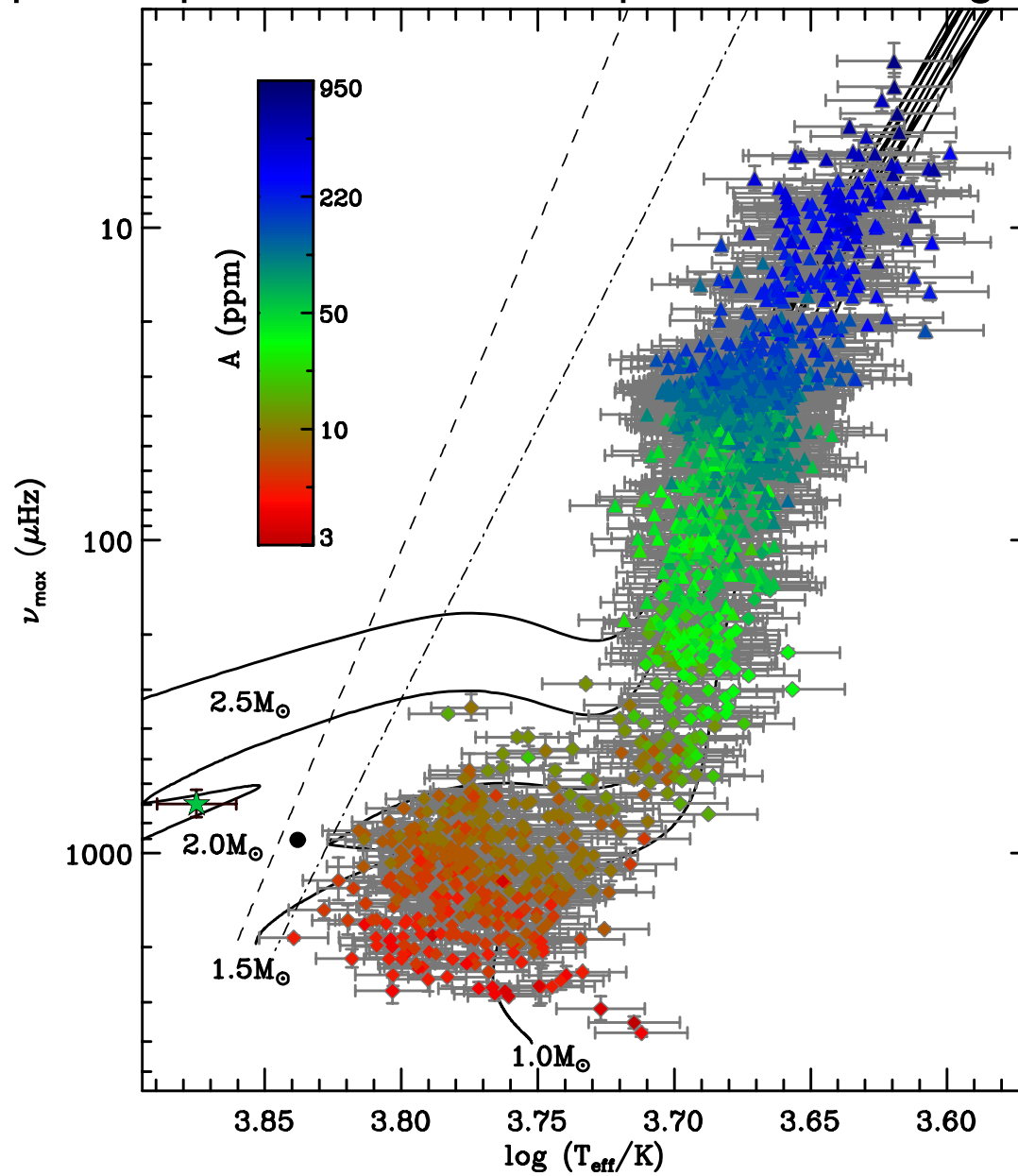
Solar-like pulsator



- ▶ Main-sequence/post main-sequence phase
- ▶ In the Sun up to 10 million p modes
- ▶ Amplitudes: $\text{cms}^{-1} \longrightarrow \text{ms}^{-1}$
few ppm \longrightarrow hundreds of ppm
- ▶ Driven stochastically by the turbulence in the outer stellar layers

Solar-type pulsator

Kepler amplitudes of main-sequence and red-giant stars

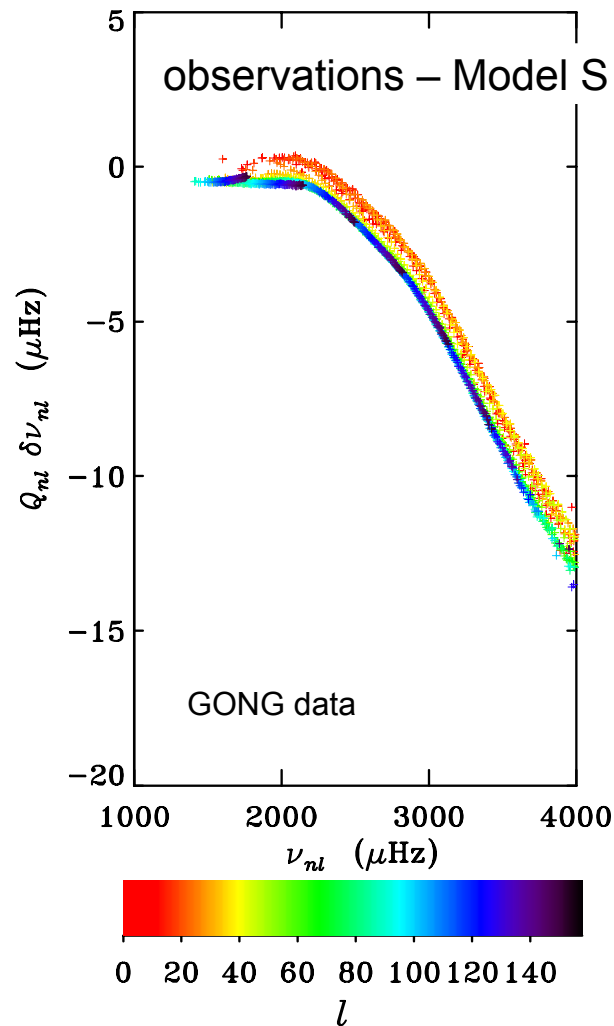


Huber et al. (2011)

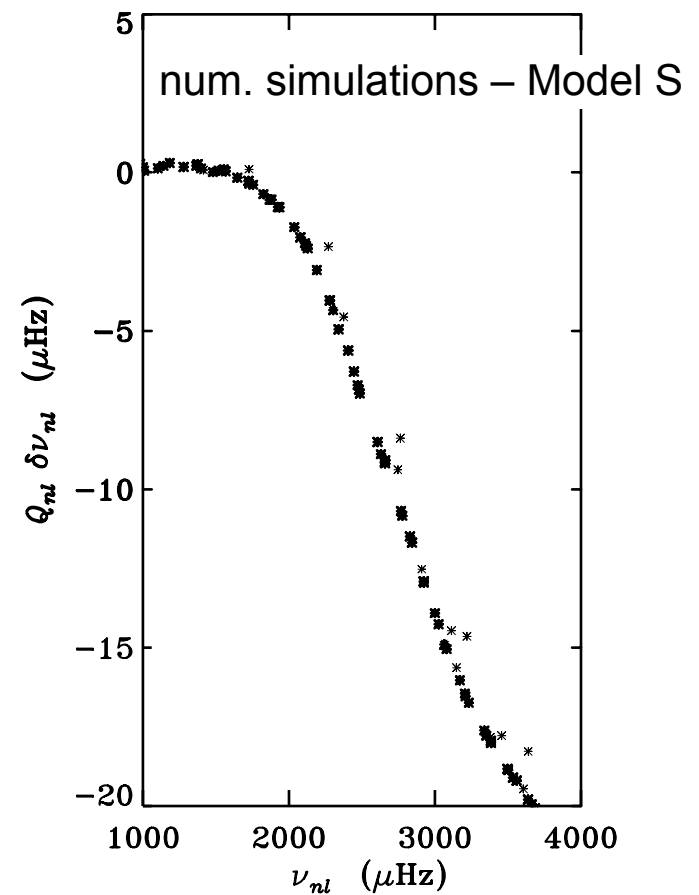
Solar-like pulsator: surface effects

Solar observations – **adiabatic** calculations

Christensen-Dalsgaard et al. (1996)



Rosenthal et al. (1995)

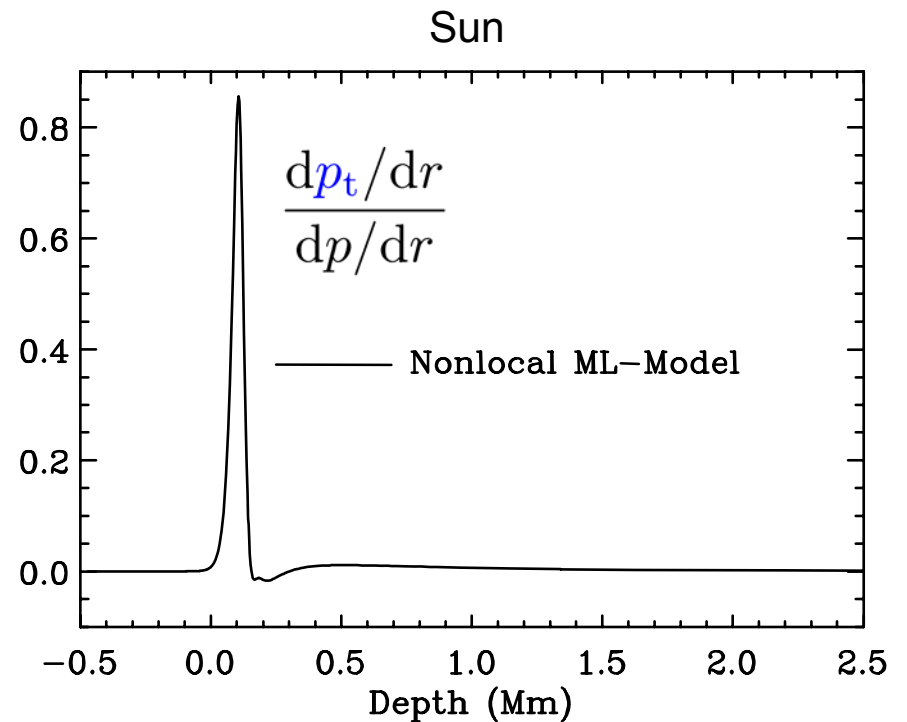
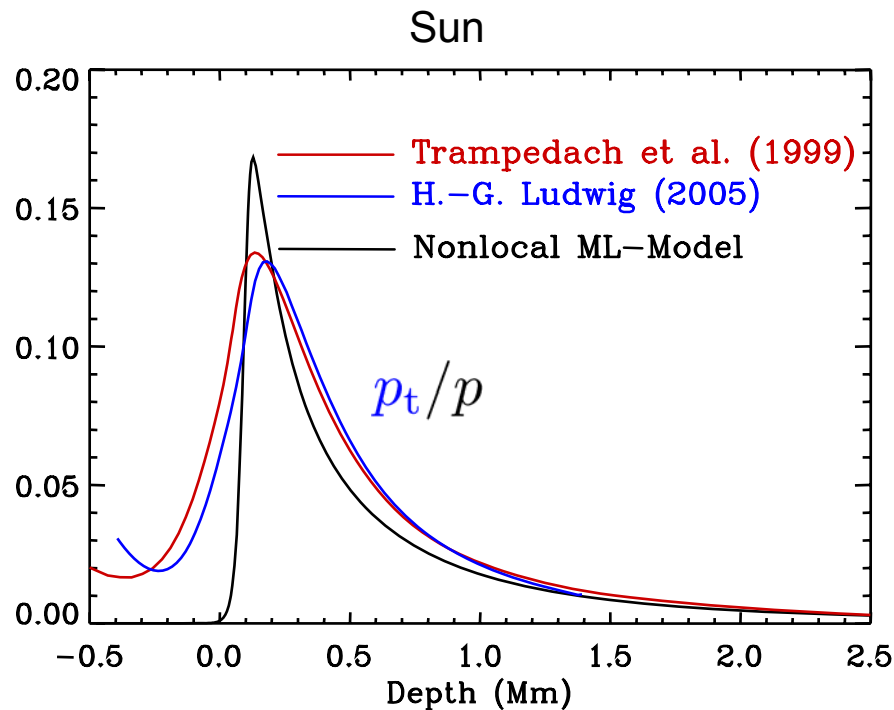


Solar-like pulsator: surface effects

Momentum equation of stellar (envelope) structure:

$$\frac{\partial}{\partial m} \underbrace{(p_g + p_t)}_p + (3 - \Phi) \frac{p_t}{4\pi r^3 \rho} = -\frac{1}{4\pi r^2} \left(\frac{Gm}{r^2} + \frac{\partial^2 r}{\partial t^2} \right)$$

(mean) turbulent momentum flux (turbulent pressure): $p_t = \overline{\rho w^2}$



GH (2006)

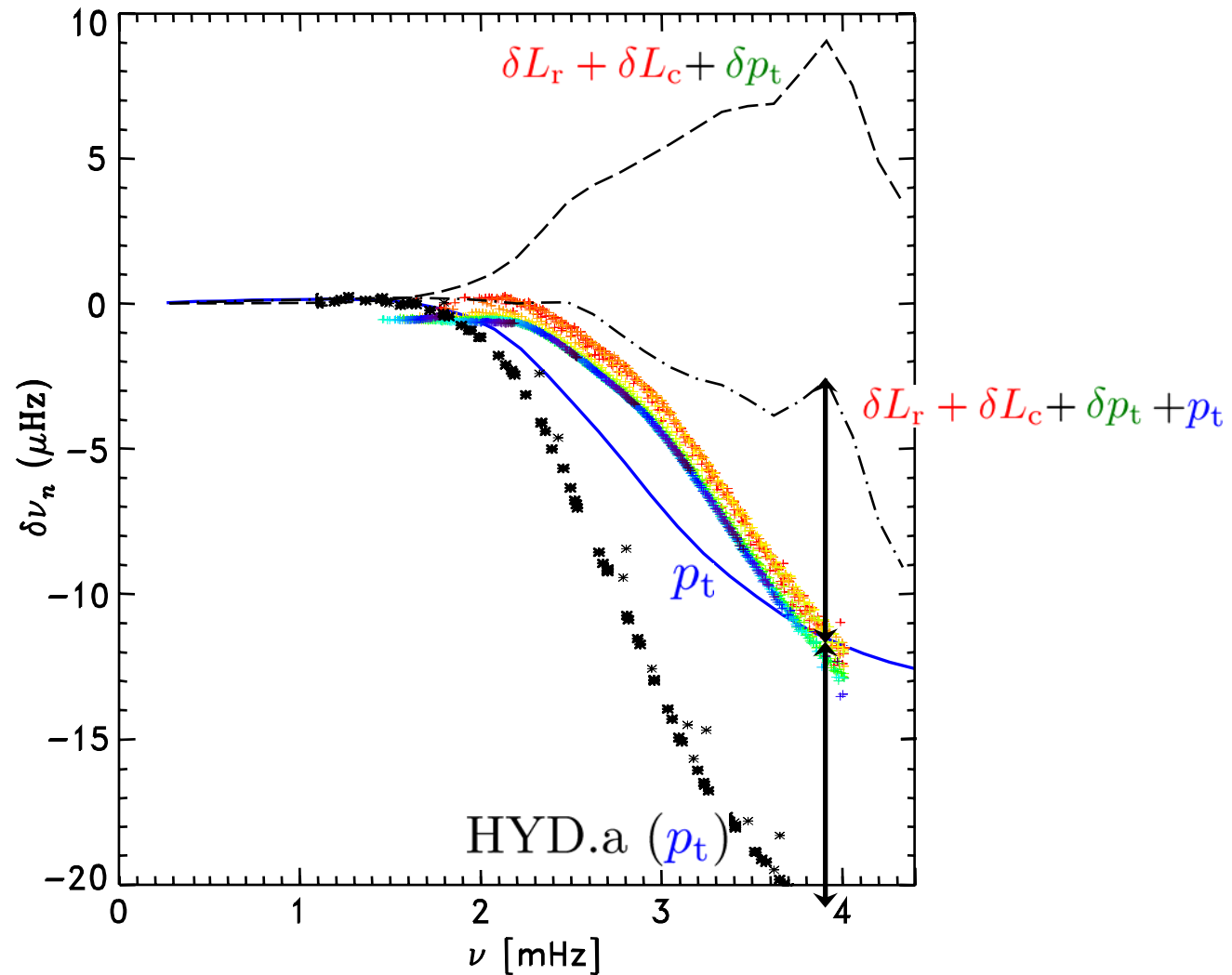
Solar-like pulsator: surface effects

Nonadiabaticity:

$$\frac{\partial}{\partial m} \left(\frac{\delta(L_r + L_c)}{L} \right) = -i\omega \frac{c_p T}{L} \left(\frac{\delta T}{T} - \nabla_{\text{ad}} \frac{\delta p}{p} \right) ;$$

Convection dynamics:

$$\frac{\partial}{\partial m} \left(\frac{\delta p}{p} \right) = f \left(\frac{\delta r}{r}, \frac{\delta T}{T}, \frac{\delta p}{p}, \frac{\delta p_t}{p}, \frac{\delta \Phi}{\Phi} \right)$$



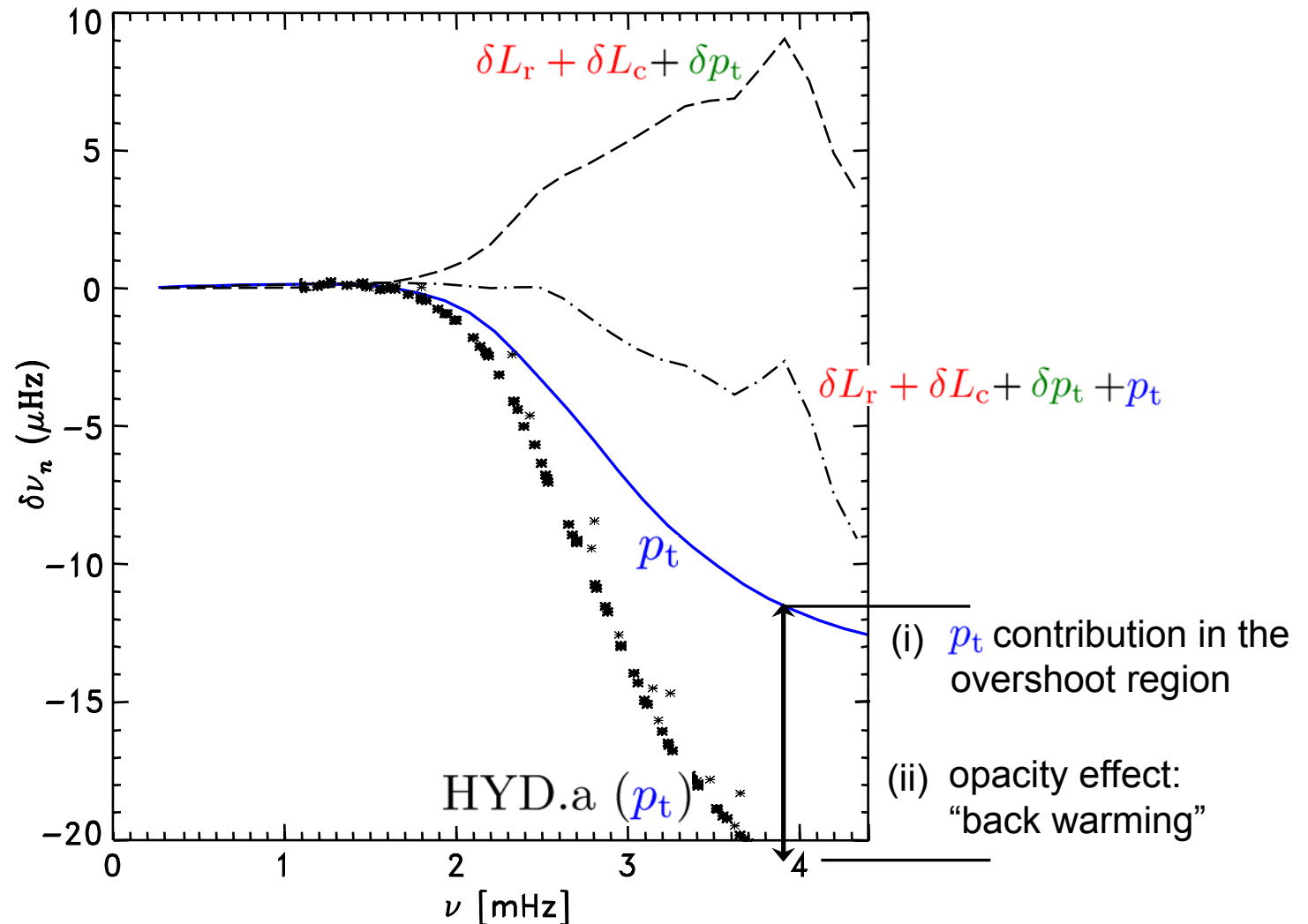
Solar-like pulsator: surface effects

Nonadiabaticity:

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Convection dynamics:

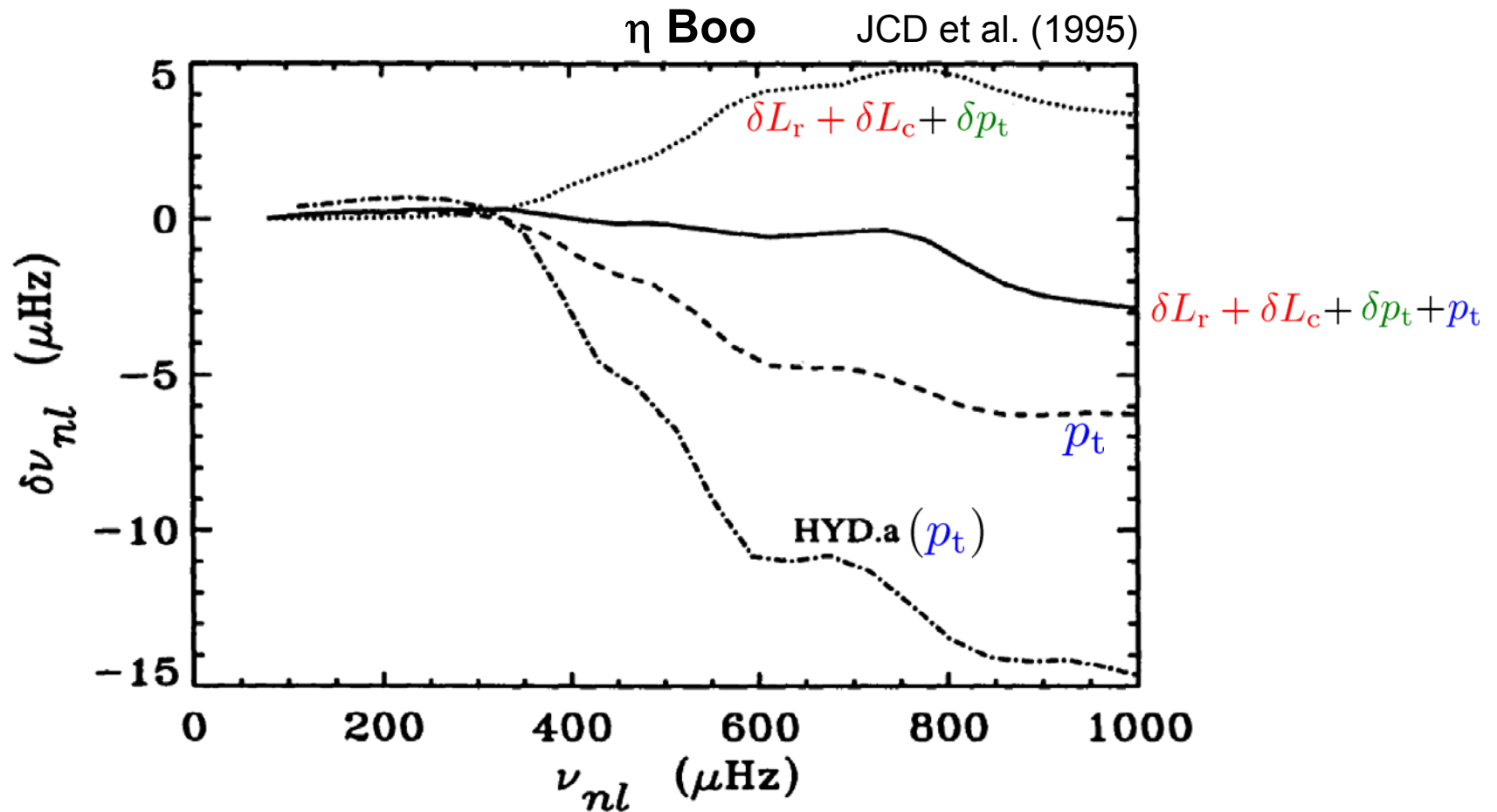
$$\frac{\partial}{\partial m} \left(\frac{\delta p}{p} \right) = f \left(\frac{\delta r}{r}, \frac{\delta T}{T}, \frac{\delta p}{p}, \frac{\delta p_t}{p}, \frac{\delta \Phi}{\Phi} \right)$$



Solar-like pulsator: surface effects

perturbed energy equation:
$$\frac{\partial}{\partial m} \left(\frac{\delta(L_r + L_c)}{L} \right) = -i\omega \frac{c_p T}{L} \left(\frac{\delta T}{T} - \nabla_{\text{ad}} \frac{\delta p}{p} \right)$$

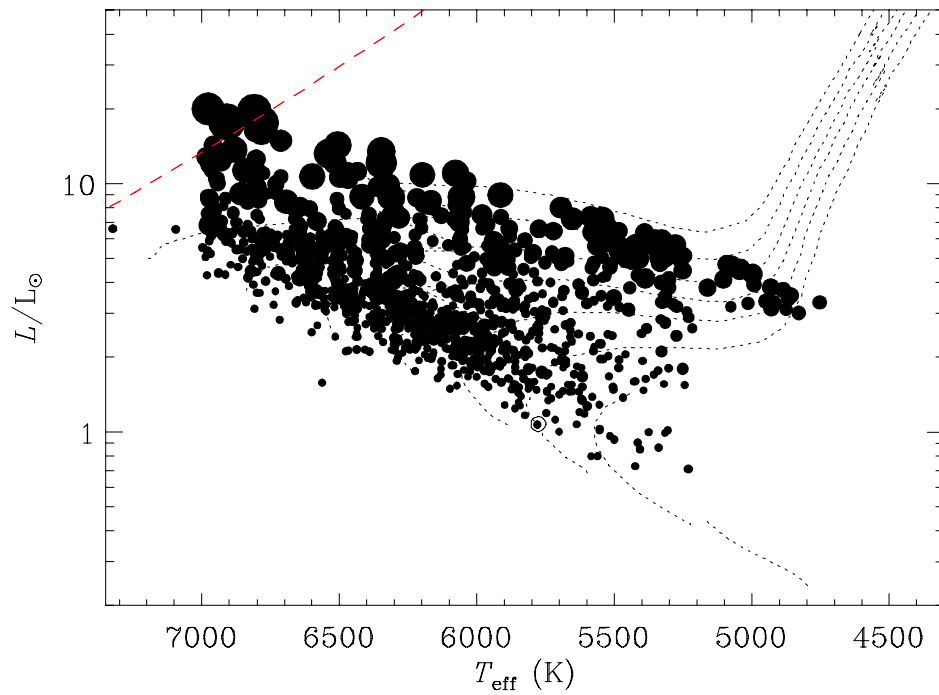
perturbed momentum equation:
$$\frac{\partial}{\partial m} \left(\frac{\delta p}{p} \right) = f \left(\frac{\delta r}{r}, \frac{\delta T}{T}, \frac{\delta p}{p}, \frac{\delta p_t}{p}, \frac{\delta \Phi}{\Phi} \right)$$



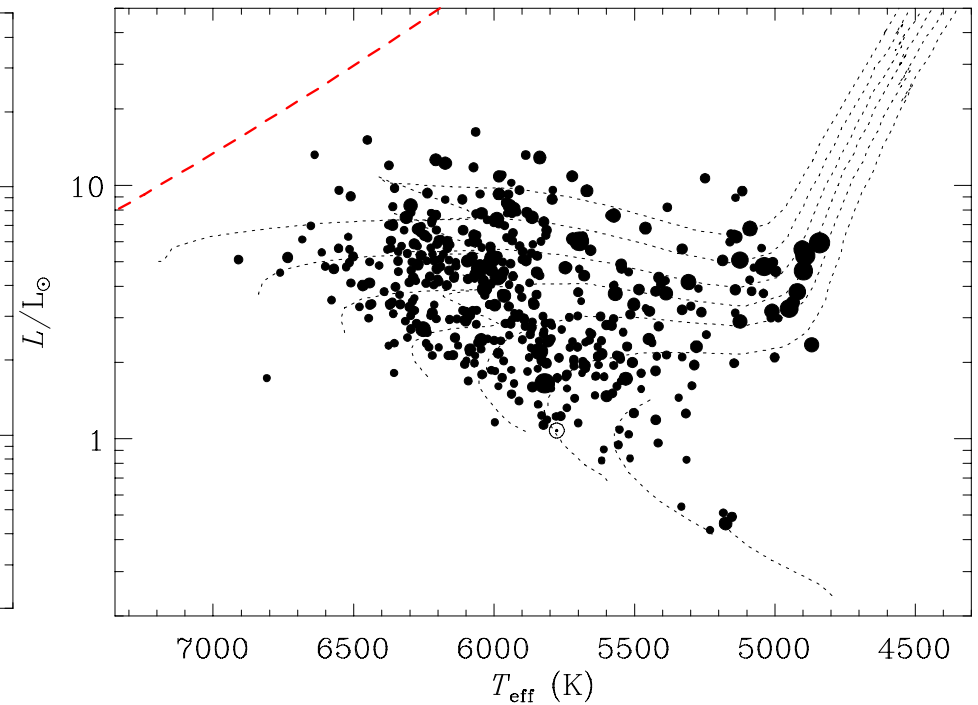
Solar-type pulsator

Pulsation amplitudes

Theoretical scaling: $V/V_{-} \sim (L/M)^S$

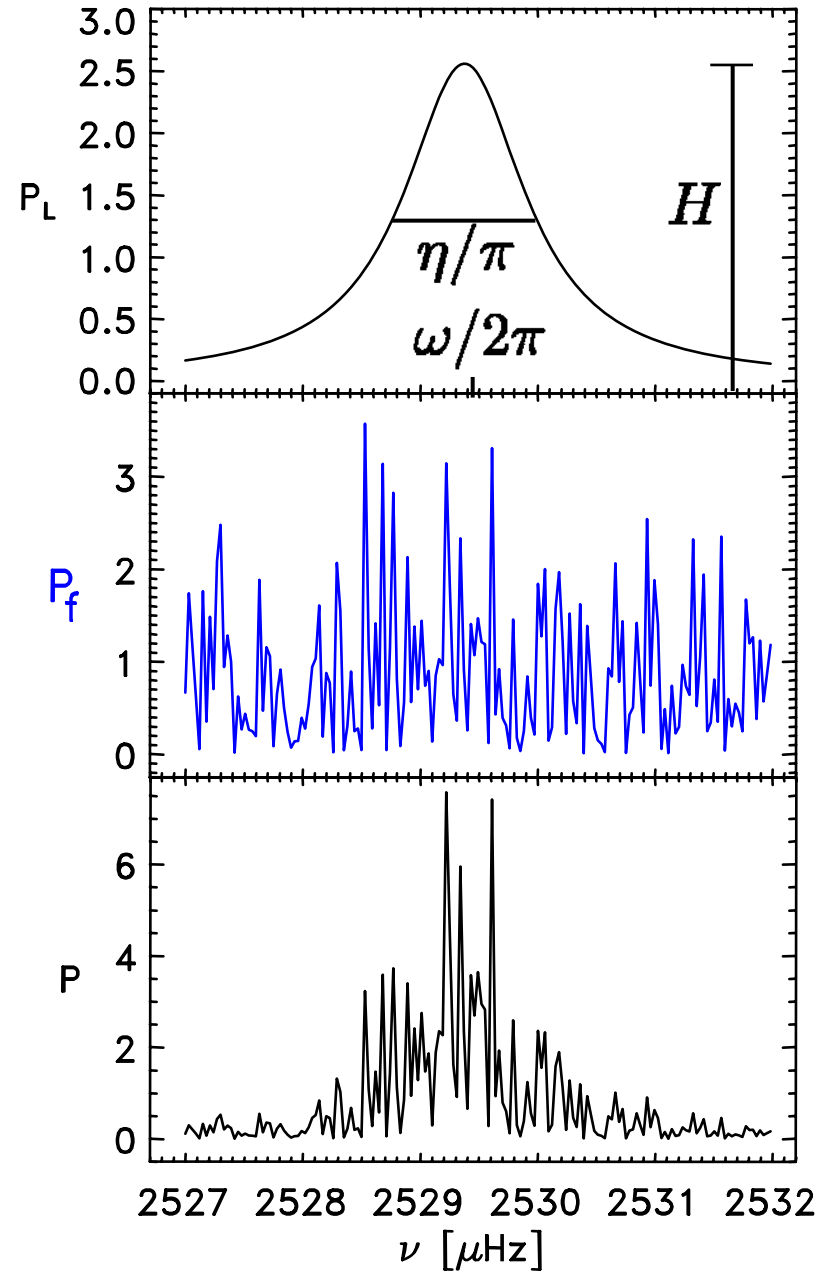
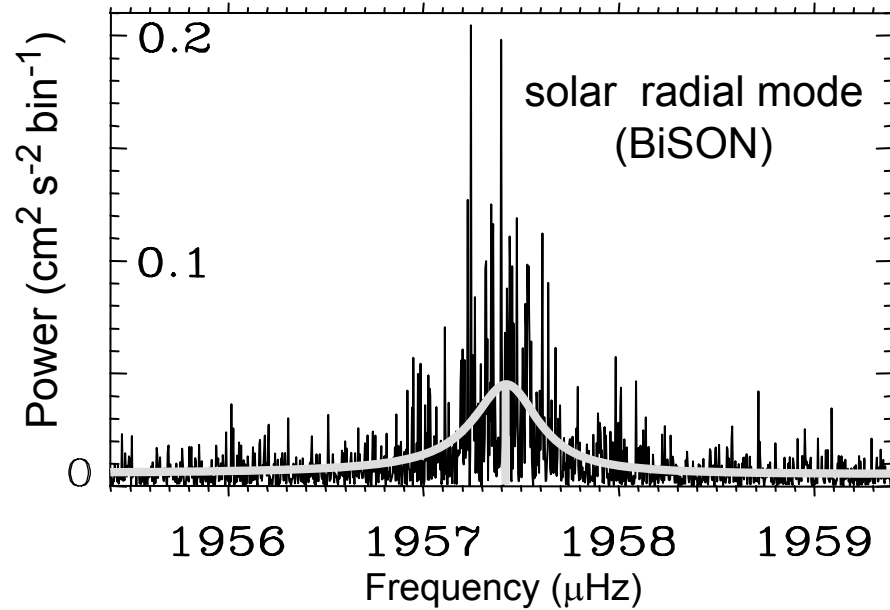


Kepler observations: SNR



(Chaplin et al. 2011)

Solar-type pulsator: excitation model



- Inhomogeneous wave equation:

$$\rho \left(\ddot{\xi} + 2\eta\dot{\xi} + \mathcal{L}\xi \right) = \mathcal{F}(u, t)$$

- Power (spectral density): $P \propto P_L P_f$

- Total integrated power (mean energy E):

$$E = IV^2 \propto \frac{P_f(\omega)}{\eta} \propto I\eta H$$

Solar-type pulsator: excitation model

(Goldreich & Keeley 1977, Balmforth 1992, Samadi et al. 2001, Chaplin et al. 2005)

Reynolds stress contribution

$$P \propto I^{-1} \int_0^R \ell^3 \left(r \frac{\partial \xi_{ir}}{\partial r} p_t \right)^2 \mathcal{S} dr$$

 **mean turbulent pressure**


$$H := P/\eta^2 I = 2V^2/\eta$$

Solar-type pulsator: excitation model

(Goldreich & Keeley 1977, Balmforth 1992, Samadi et al. 2001, Chaplin et al. 2005)

Reynolds stress contribution

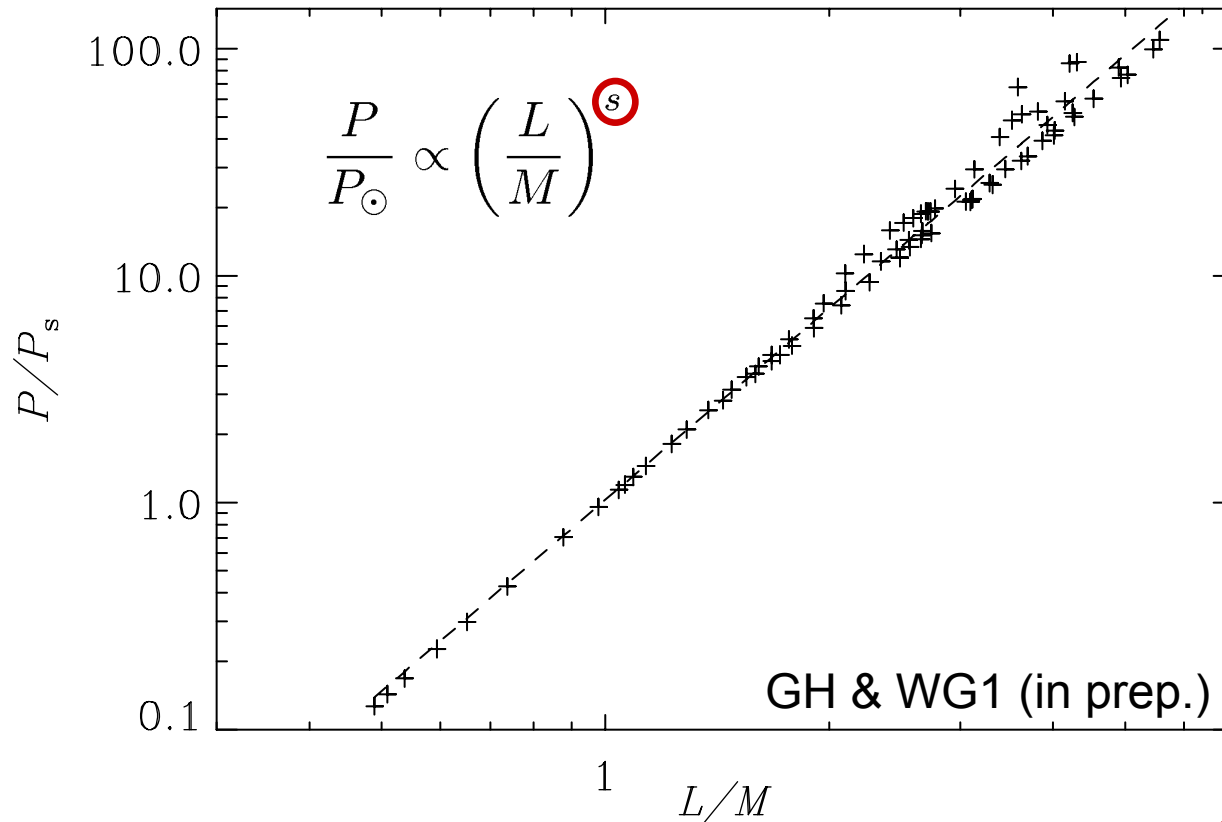
$$P \propto I^{-1} \int_0^R \ell^3 \left(r \frac{\partial \xi_{ir}}{\partial r} p_t \right)^2 \mathcal{S} dr$$

 **small-scale turbulence**

$$H := P/\eta^2 I = 2V^2/\eta$$

Solar-type pulsator: excitation model

Energy supply rate P for 83 models with $0.9 \leq M/M_{\odot} \leq 1.5$



Excitation model	Ω	s
Chaplin et al. (2005)	Gaussian: $\lambda = 1.0$	2.80
Samadi et al. (2003)	mod. Lorentzian: $\lambda = 1.5; \beta = 3.2$	2.87
Samadi et al. (2003)	mod. Lorentzian: $\lambda = 1.0; \beta = 3.2$ (Kaneda 1993)	2.72

Solar-type pulsator: excitation model

(Goldreich & Keeley 1977, Balmforth 1992, Samadi et al. 2001, Chaplin et al. 2005)

Reynolds stress contribution

$$P \propto I^{-1} \int_0^R \ell^3 \left(r \frac{\partial \xi_{ir}}{\partial r} p_t \right)^2 \mathcal{S} dr$$

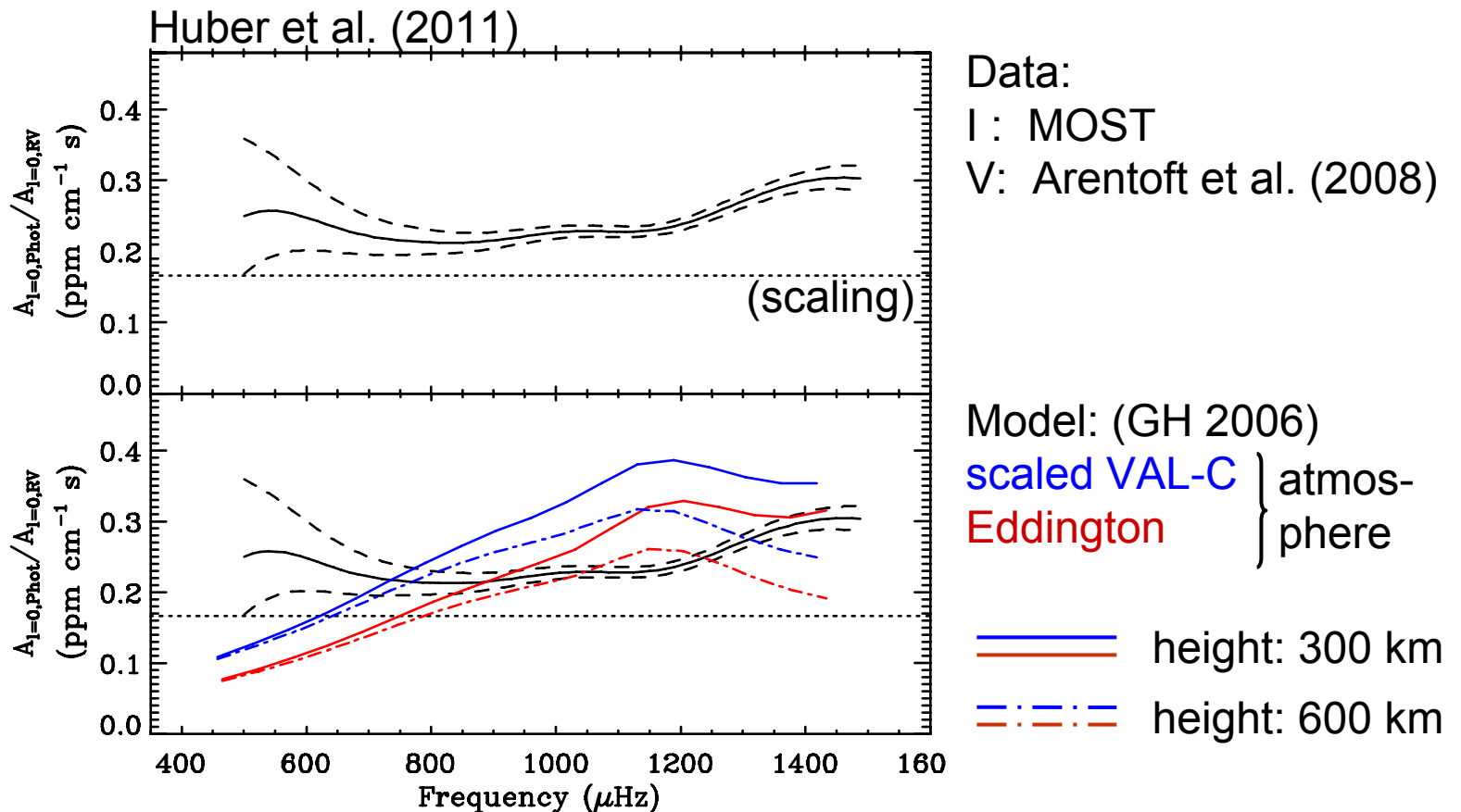
eigenfunction ξ

$$H := P/\eta^2 I = 2V^2/\eta$$

Solar-type pulsator: eigenfunctions ξ

Amplitude ratios in a model for Procyon A

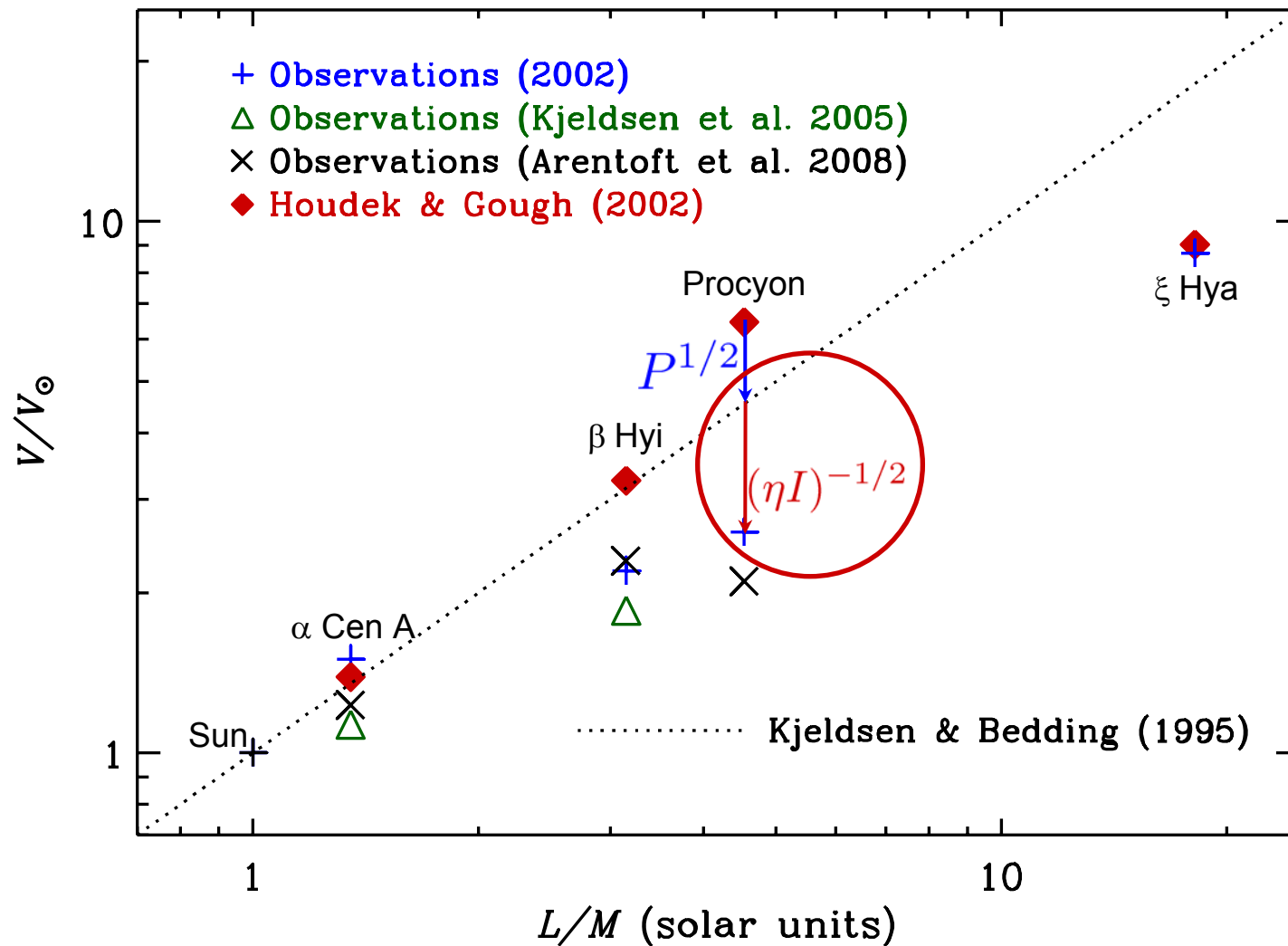
$$\text{modelled amplitude ratio} = \frac{\delta L/L}{\omega_r r \delta r/r}$$



Note: amplitude ratios do not depend on a stochastic excitation model !!

Solar-type pulsator: mode lifetimes τ

Velocity amplitudes



L/M (solar units)

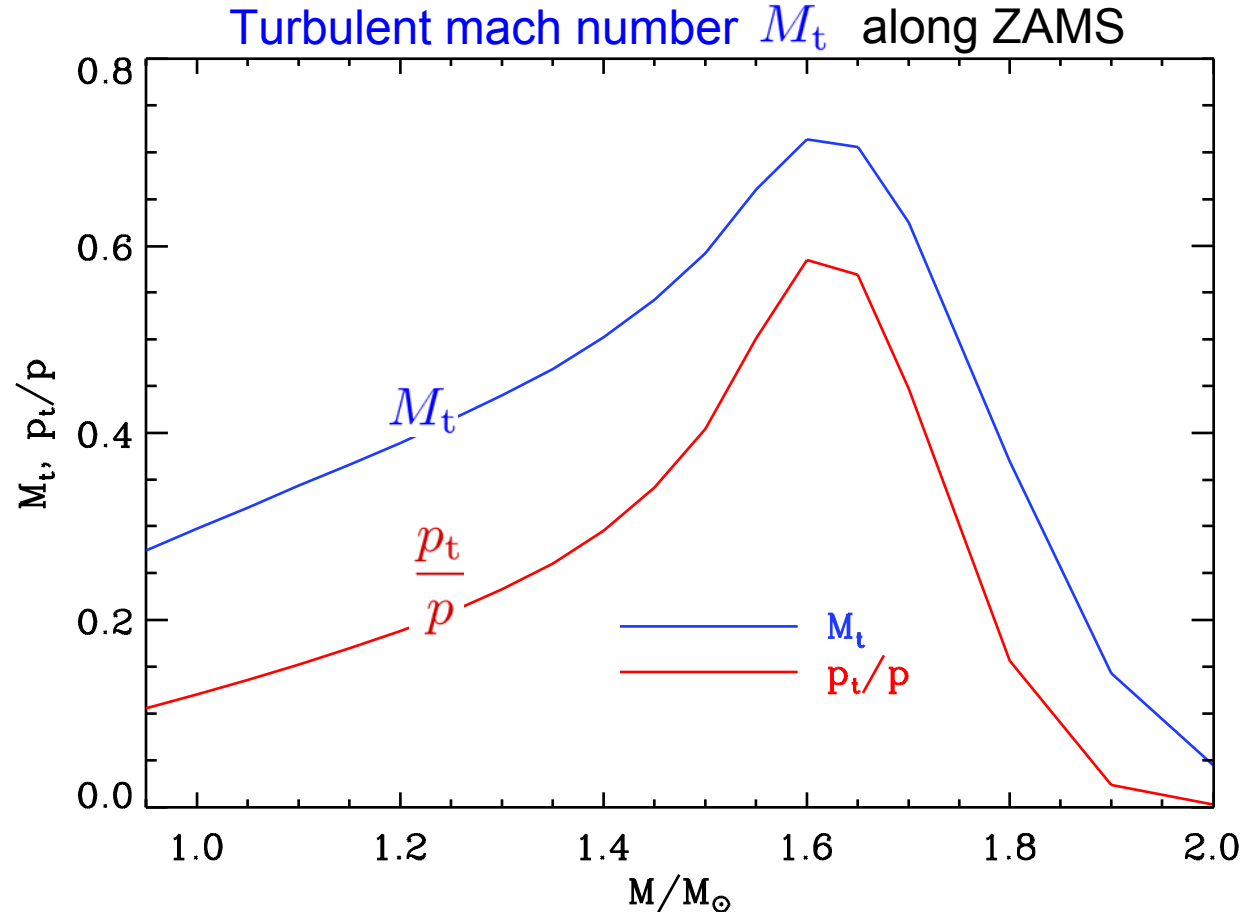
$$H := P/\eta^2 I$$

Solar-type pulsator: mode lifetimes τ

Scattering contribution (Goldreich, Murray 1994):

$$\eta_s \sim \frac{2\omega}{\pi(n+1)} M_t^2 \quad M_t := \frac{u}{c}$$

pulsation energy is scattered from low- l \rightarrow high- l modes
ultimate recipients: f and g modes



GH et al. (1999)

Summary/Conclusion

- ▶ **Turbulent convection** controls location of **cool edge of IS** in **δ Scuti stars**; the detailed mechanism (i.e. convective heat flux, momentum flux of turbulent viscosity) for return to stability, however, is still uncertain.
- ▶ **Acoustic radiation** in the equilibrium model affects both the **hot & cool edge of IS** in **δ Scuti stars**.
- ▶ **Mean turbulent pressure** appears to be dominating **surface effect** in explaining differences between observed and modelled frequencies of **solar-like oscillations**; **nonadiabaticity** and **convection dynamics** must also not be neglected.
- ▶ **Poor modelling** of pulsation **eigenfunctions** and mode **lifetimes** together with partial cancellation of acoustic **excitation sources** may explain disagreement between modelled and observed pulsation amplitudes of **solar-like oscillations** in hotter stars.