



## Non-thermal X-ray Filaments in Cas A: Testing Plasma Turbulence

### David Lomiashvili

M. Lyutikov, M. Araya, C. Chang, W. Cui

Department of Physics, Purdue University



- Cosmic-Ray acceleration requires a strong turbulence, particle diffusion near astrophysical shocks
- Non-thermal X-ray filaments in shell type SNRs
- Estimate :
  - magnetic field near the shock
  - diffusion coefficient
  - level of turbulence
  - shock obliquity
- Test plasma turbulence near the shock

Non-thermal filaments

Cas



- 1 Ms Chandra observation of Cas A (Hwang et al. 2004)
- Obtain:
  - spatially resolved spectra
  - radial profiles at different energies
- Expect :
  - spectrum softer inside
  - larger width at lower photon energies

Araya et al. (in preparation)





- Non-thermal filaments modeled as thin spherical regions associated with forward shock
- Particles injected at the shock have a power-law distribution
- Advected downstream
- Diffusion
- Two different scenarios are considered:
  - Advection + Synchrotron losses "Advection-Loss"
  - Advection + Synchrotron losses + Diffusion "Advection Diffusion" ( with different types of energy dependence of diffusion coefficient )
- kinetic equation :







Magnetic turbulence



Particle diffusion due to the scattering on magnetic irregularities

Spectrum of magnetic Turbulence –  $W(k) \propto k^{-q}$ 



Spatial diffusion coefficient 
$$D\propto\gamma^{2-q}~(q-1\equiv\delta)$$

We parameterize diffusion coefficient as :

$$D(\gamma) = \kappa D_0 \gamma^{1-\delta}$$

$$D_0 = \left(mc^3/3eB\right)$$

Type of Turbulence	q	δ
White noise	1	0 ( <i>Bohm</i> )
Kraichnan	1/2	1/2
Kolmogorov	5/3	2/3



• Solution of the kinetic equation (*added Advection*):

$$n(\gamma, r) = \frac{1}{\beta B^2 \gamma^2} \int_{\gamma}^{\gamma_0} d\gamma_0 \ Q(\gamma_0) \ \frac{exp\left[-\frac{\left(r - V_{adv}t(\gamma, \gamma_0)\right)^2}{4\lambda(\gamma, \gamma_0)}\right]}{\left(4\pi\lambda(\gamma, \gamma_0)\right)^{3/2}}\right]}{\left(4\pi\lambda(\gamma, \gamma_0)\right)^{3/2}}$$

$$\lambda(\gamma,\gamma_0) = \int_{\gamma}^{\gamma_0} d\gamma' \frac{D(\gamma')}{\beta B^2 \gamma'^2}$$
 - Syrovatskii parameter

$$Q(\gamma_0) = K \gamma_0^{-\Gamma}$$

 $\lambda(\gamma,\gamma_0)$  is equivalent to the squared distance traveled by the particle while "cooling" from  $\gamma_0$  to  $\gamma$ 

$$\begin{array}{c}
l_{adv} = V_{adv}t_{cool} \\
\hline l_{dif} \sim \sqrt{\lambda(\gamma, \gamma_0)} \\
\hline R \\
\end{array}$$

$$\begin{array}{c}
C_1 = \frac{l_{dif}(1keV)}{R} \\
\hline C_2 = \frac{l_{adv}(1keV)}{l_{dif}(1keV)} \\
\hline \end{array}$$

$$l_{dif}(1keV) = \left(\frac{\kappa D_0}{\beta B^2} \left(\frac{1.5h\nu_L}{1keV}\right)^{\delta/2}\right)^{1/2}$$
$$l_{adv}(1keV) = \frac{V_{adv}}{\beta B^2} \left(\frac{1.5h\nu_L}{1keV}\right)^{1/2}$$



Model

• Synchrotron emissivity integrated along the line of sight :

$$F_{\nu}(z) = \int_0^{\bar{l}_{max}(z)} d\bar{l} I_{\nu}(\rho(\bar{l}, z))$$

$$I_{\nu}(\rho) = \int_{\gamma} d\gamma P_{\nu}(\gamma) n(\gamma, \rho)$$

- Synchrotron emissivity

$$P_{\nu}(\gamma) = \frac{\sigma_T c B^2 \gamma^2}{6\pi} \delta(\nu - \nu_c)$$

- Specific power by single electron

$$F_E(z) \propto E_{keV}^{-\Gamma/2+3\delta/4} \int_{\bar{l}_{min}}^{\bar{l}_{max}} d\bar{l} \int_1^\infty d\epsilon \, \frac{\epsilon^{-\Gamma+3\delta/2}}{\left(\epsilon^{\delta}-1\right)^{3/2}}$$
$$\times exp\left[-\frac{\delta \cdot E_{keV}^{\delta/2}}{4\left(1-\epsilon^{-\delta}\right)} \left(\frac{\rho(\bar{l},z)}{C_1} - \frac{C_2}{E_{keV}^{1/2}} \left(1-\frac{1}{\epsilon}\right)\right)^2\right]$$

• Precursor (upstream) solution same as downstream except: ( $V_{adv}$ )<sub>down</sub>  $\iff$  (-4<sup>\*</sup>  $V_{adv}$ )<sub>up</sub>, (B)<sub>down</sub>  $\iff$  (B/4)<sub>up</sub>





 $C_1 = 0.007$   $C_2 = 5$   $\Gamma = 2.5$   $B = 112 \ \mu G$ 



- In *all* 9 filaments spectrum is harder *outside* than *inside* by about 1σ (both *inside* and *outside* in a downstream region)
- This supports the assumption about *particle cooling*

Filament widths and spectral data from *Araya et al. (in preparation)* 



Inclusion of the diffusion in the model is necessary !!!



- "Advection + Synchrotron" model ruled out
- "Advection + Synchrotron + Diffusion" model (Bohm type diffusion):
  - Determine B-field
  - Determine  $\kappa$  (D=  $\kappa D_{Bohm}$ )
- Average magnetic field B ~ 60  $\mu$ G (30 - 160  $\mu$ G)
  Diffusion scaling factor (D =  $\kappa D_{Bohm}$ )  $\kappa \sim 0.1$  (most around 0.05)
  Power-law index of electrons  $\Gamma \sim 2.5$  (2.0 - 4.0)

$$B \approx 56 \mu G \quad \left(\frac{C_1}{0.02}\right)^{-2/3} \left(\frac{C_2}{5.0}\right)^{-2/3} \left(\frac{V_{adv}}{1.3 \times 10^8 cm \ s^{-1}}\right)^{2/3} \left(\frac{R}{10^{19} cm}\right)^{-2/3}$$

$$\kappa \approx 0.05 \ \left(\frac{C_2}{5.0}\right)^{-2} \ \left(\frac{V_{adv}}{1.3 \times 10^8 cm \ s^{-1}}\right)^2$$

#### Level of Turbulence and Shock Obliquity

For the scalar diffusion coefficient:

$$D = D_{\parallel} \cos^2 \theta + D_{\perp} \sin^2 \theta$$

$$D = D_{\parallel} \left[ \left( \cos^{2} \theta + \frac{\sin^{2} \theta}{1 + \eta^{2}} \right) \right]$$

$$D_{\parallel} = \eta D_{0} \gamma^{1 - \delta}$$

$$D = \kappa D_{0} \gamma^{1 - \delta}$$

$$D = \kappa D_{0} \gamma^{1 - \delta}$$

$$(Jokipii 1987)$$

$$\eta \equiv \left( \frac{\lambda_{mfp}}{r_{g}} \right) = (\delta B/B)_{res}^{-2}$$

$$1 \le \eta \quad \text{-gyrofactor}$$

$$\kappa = 0.05 \Rightarrow \left[ 88.5^{\circ} \le \theta \le 90^{\circ} \right] \Rightarrow \eta \ge 20 \Rightarrow \left[ \frac{\delta B}{B} \le 0.22 \right]$$

$$\kappa = 0.15 \Rightarrow \left[ 86^{\circ} \le \theta \le 90^{\circ} \right] \Rightarrow \eta \ge 10 \Rightarrow \left[ \frac{\delta B}{B} \le 0.4 \right]$$

$$\kappa = 0.3 \Rightarrow \left[ 81^{\circ} \le \theta \le 90^{\circ} \right] \Rightarrow \eta \ge 3 \Rightarrow \left[ \frac{\delta B}{B} \le 0.4 \right]$$

$$\kappa = 0.3 \Rightarrow \left[ 81^{\circ} \le \theta \le 90^{\circ} \right] \Rightarrow \eta \ge 3 \Rightarrow \left[ \frac{\delta B}{B} \le 0.4 \right]$$

$$Shocks are quasi-perpendicular ! consistent with wind environment$$

Maximum energy of accelerated electrons :

90

88

86

84

78

76

74 L

80

heta - obliquity angle

$$E_{max} \approx 330 \ \left(\frac{\kappa}{0.05}\right)^{-1/2} \left(\frac{B}{30\mu G}\right)^{-1/2} \left(\frac{V_{sh}}{5.2 \times 10^8 cm \ s^{-1}}\right) \ TeV$$



## Caveats

- We did  $\delta$ -function profile for specific synchrotron power of a single electron.
- More detailed calculations may be needed near the cut-off
- Detailed shapes of profiles are not well reproduced Half width @ half-maximum still relevant
- Possible solution precursor (upstream) emission
- Thin radio filaments found in some SNRs can't be explained by loss-limited approach
- Possible solution spatial structure of magnetic field ( Pohl, Yan, & Lazarian 2005 )
- Can't yet distinguish different types of turbulence (White noise, *Kolmogorov*, *Kraichnan*)
- Possible solution Fit the data for different values of  $\delta$





# Summary

- Important observational test for the shock acceleration mechanisms
- Diffusion is important in the downstream region:
  - Various levels of *spatial evolution of spectra*
  - *Weak* energy dependence of filament widths
- Magnetic field in filaments :  $B \approx 30 160 \ \mu G$
- Diffusion coefficient as a fraction of *Bohm* :  $\kappa \sim 0.05 0.35$
- Turbulence is moderately strong for all regions : δ**B/B** ~ 0.2-0.4
- For all studied filaments shocks are *quasi-perpendicular*
- Future work : i) Fit the data for different values of  $\delta$  to test the turbulence spectrum
  - ii) Understand the implications of our results on DSA mechanism
  - iii) Study other shell type SNRs