

# Electrons in narrow-band moire graphene: two-dimensional Bloch oscillations

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Correlations in Moire Flat Bands, KITP  
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# Dynamics of a Bloch electron in 1D solids

Bloch waves in a periodic lattice:  $H_0 = \frac{p^2}{2m} + U(x)$ :

$$\psi_n(x) = u_{n,k}(x)e^{ikx}, \quad \epsilon_n(k)\psi_n(x) = H_0\psi_n(x)$$

with band dispersion  $\epsilon_n(k)$  periodic in  $k$ , and the envelop functions  $u_{n,k}(x)$  periodic in  $x$ .

The  $k$ -space form of Newton's law for Bloch waves in the presence of an external field,  $H = H_0 - eEx$ :

$$\hbar \frac{d\mathbf{k}}{dt} = e\mathbf{E} \quad (1)$$

Follows from adiabatic apprx for one band, or rigorously for a multiband dynamics.

# Oscillatory movement in a DC electric field (video)

- Electron wavepacket for  $n$ th band:

$\tilde{\psi}_{n,k}(x) \sim \int dk' e^{-\alpha(k'-k)^2} \psi_{n,k'}(x)$ . The velocity is the group velocity  $v = \frac{1}{\hbar} \frac{\partial \epsilon(k)}{\partial k}$

- Sign-changing  $v(t)$ , current oscillations!
- Bloch oscillator **DC**  $\rightarrow$  **AC**

Quasiclassical explanation:  $\hbar \frac{dk}{dt} = eE$  gives

$k(t) = k_0 + eEt/\hbar$  that sweeps the entire Brillouin zone. After reaching the right end point  $k = \frac{\pi}{a}$  electron Bragg-reflects and continues from the left end point  $k = -\frac{\pi}{a}$ , since Brillouin zone is a circle (in 1D). Sign-changing velocity  $v(t)$ , time period  $T = (2\pi/a)\hbar/eE = h/eEa$ .

Bloch oscillation frequency  $\nu = eEa/h$ .

# Experiments: an ongoing quest

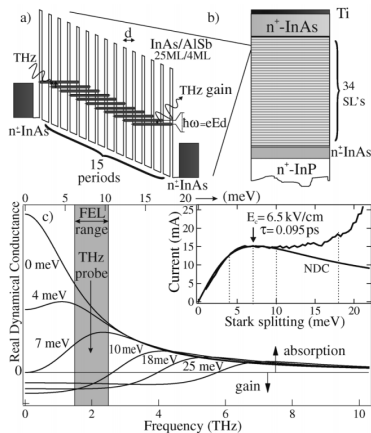
- **Requirements:** narrow bandwidth  $J < \omega_{\text{ph}}$ , and low disorder  $\gamma_{\text{dis}} < \nu$
- Large lattice constant, to achieve high  $\nu = eEa/h$  at lower  $E$

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THz Loss and Gain in a Bloch Oscillating InAs/AISb Superlattice

[Savvidis et al, PRL 92, 196802

(2004)]

## Experiments: an ongoing quest

“The THz part of the electromagnetic spectrum is marked by a lack of commercial technology. Bloch oscillating superlattices have the potential to provide broad band gain at THz frequencies and may be the basis of a technology that will help fill the gap in solid state THz fundamental oscillators.” (2004)

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- Oscillation at a frequency controlled by the external resonator or circuit but only at frequencies below the Bloch frequency
- Promising, but plagued by charge instabilities due to negative differential conductivity
- Coherent THz emission?
- Proof-of-principle demonstration with **cold atoms in optical lattices** (Weld @ UCSB)



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Other promising systems?

# Narrow bands in moire graphene

- High optical phonon frequency  $\omega_{\text{ph}} \sim 200\text{meV}$
- Condition  $J < \omega_{\text{ph}}$  easy to fulfill
- Large lattice constant  $a \sim 10\text{nm}$
- Opportunity: demonstrate BO in moire bands at  $eE < J/a$ , avoiding Wannier-Stark ladder localization and charge instabilities

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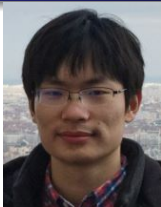
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Feasible?

# Two-dimensional Bloch oscillations



with Zhiyu Dong

- Two independent frequencies (trajectories wind BZ in two different directions)
- Values tunable by the field polar angle  
 $\mathbf{E} = E(\cos \theta, \sin \theta)$
- Sensitive to band dispersion, a probe of carrier dynamics
- Interesting dynamics in magnetic field (phase transition to one-frequency oscillation)

## Illustrate for triangular lattice tight-binding band:

$$\epsilon(\mathbf{k}) = -J \sum_{j=0,1,2} \cos(\mathbf{k}\mathbf{a}_j), \quad \mathbf{a}_j = a(\cos \theta_j, \sin \theta_j)$$

$\theta_j = \frac{2\pi}{3}j$ . Combining with Newton's dynamics  $\hbar \frac{d\mathbf{k}}{dt} = e\mathbf{E}$ , gives trajectories  $\mathbf{k}(t) = e\mathbf{E}t + k_0$ . An oscillatory particle velocity time dependence

$$\mathbf{v}(t) = -J \sum_j \mathbf{a}_j \sin\left(\frac{e}{\hbar} \mathbf{E} \cdot \mathbf{a}_j t + \phi_0\right) \quad (2)$$

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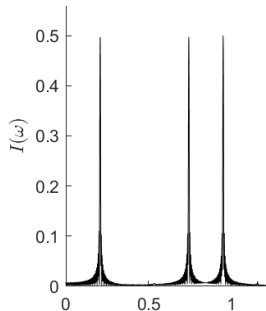
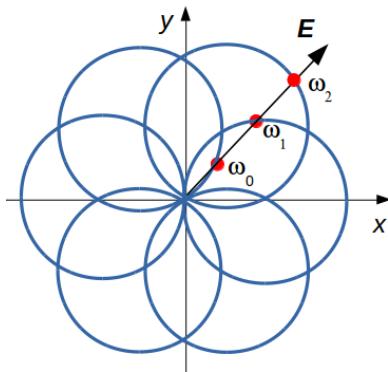
Three angle-dependent frequency values

$$\omega_j(\theta) = \frac{e}{\hbar} \mathbf{E} \cdot \mathbf{a}_j = \frac{e}{\hbar} E a \cos(\theta - \theta_j) \quad (3)$$

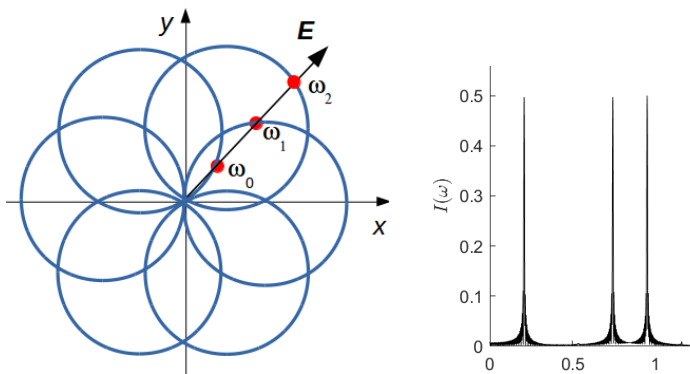
Only two frequencies are independent,  $\omega_2 = \omega_0 + \omega_1$  as expected.



# Three Bloch frequencies: angle dependence



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More frequencies for a more general dispersion.  
Tomography of the bandstructure?

## Adding magnetic field:

Newton's EOM for Bloch electron (use mks units):

$$\hbar \dot{\mathbf{k}} = e\mathbf{E} + e\mathbf{v} \times \mathbf{B} \quad (4)$$

Lorentz force with electron band velocity

$$\mathbf{v}(k) = \frac{1}{\hbar} \nabla_{\mathbf{k}} \epsilon(\mathbf{k}) \text{ (periodic in } k)$$

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Two distinct types of orbits arise depending on the relative strength of  $\mathbf{E}$  and  $B$ :

- 1) **confined orbits** (reside within BZ), B-dominated, one Bloch frequency
- 2) **deconfined orbits** (extended over infinitely many BZ's), E-dominated, two or more frequencies.

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Confined in  $k$  space means deconfined in  $r$  space, and vice versa.

# Geometric visualization of dynamics

Try to put EOM in a Hamiltonian form ( $k$  only).

$$\dot{k}_1 = \partial_2 H(k_1, k_2), \quad \dot{k}_2 = -\partial_1 H(k_1, k_2) \quad (5)$$

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$$H(k_1, k_2) = \frac{eB}{\hbar^2} \epsilon(\mathbf{k}) + \frac{e}{\hbar} \mathbf{E} \times \mathbf{k} \quad (6)$$

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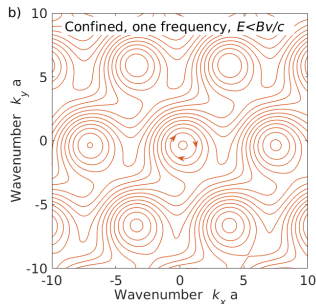
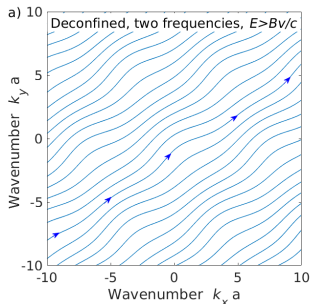
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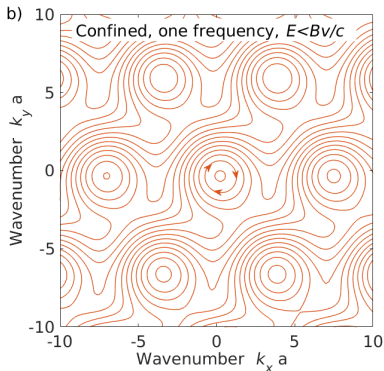
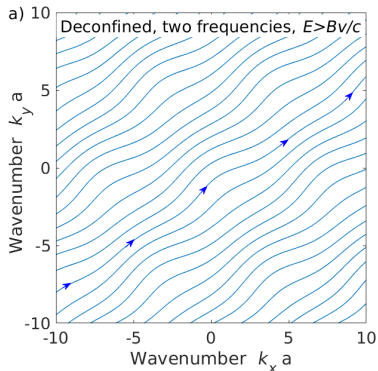
Trajectories are nothing but contours of  $H(k_1, k_2)$ !





# The confined-to-deconfined transition

Contours of a periodic function + a linear function:



# DC current at the transition

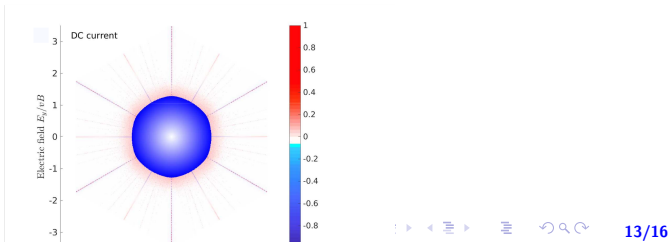
Find  $\mathbf{R}(t) = \int_0^t \mathbf{v}(t') dt'$  by integrating EOM:

$\hbar(\mathbf{k}(t) - \mathbf{k}(0)) = e\mathbf{E}t + e\mathbf{R}(t) \times \mathbf{B}$ . Therefore,

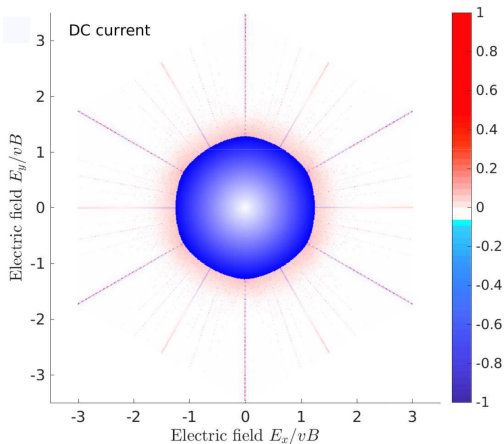
- In the confined (B-dominated) regime the drift velocity  $\mathbf{u} = \mathbf{R}(t)/t|_{t=nT}$  equals

$$\mathbf{u} = \mathbf{E} \times \mathbf{B} / B^2$$

- At transition to the deconfined (E-dominated) regime drift velocity abruptly drops to zero.



# DC current component perpendicular to E

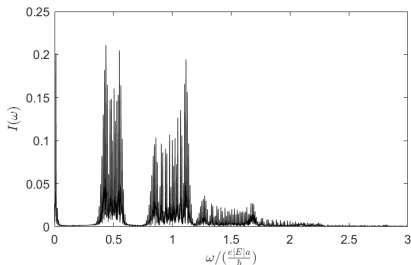
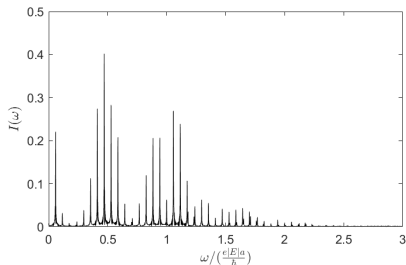


Abrupt drop to zero at transition; peculiar fine structure in the deconfined (E-dominated) state

# Narrow-band noise near transition

Observations:

- Prominent for  $\mathbf{E}$  near commensurability
- Turns into two-frequency spectrum at small  $B$
- Indicates complex dynamics at the transition



## 2D Bloch oscillation summary

- An exotic driven quantum system
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Future:

- Coherent THz radiation?
- Finite carrier density? Charge instabilities?
- DC + AC driving. Negative differential conductivity? Gain vs. loss?
- Nonstationary Quantum Hall state in the B-dominated regime?