

Moiré is Different:

Wigner crystallization in TBLG



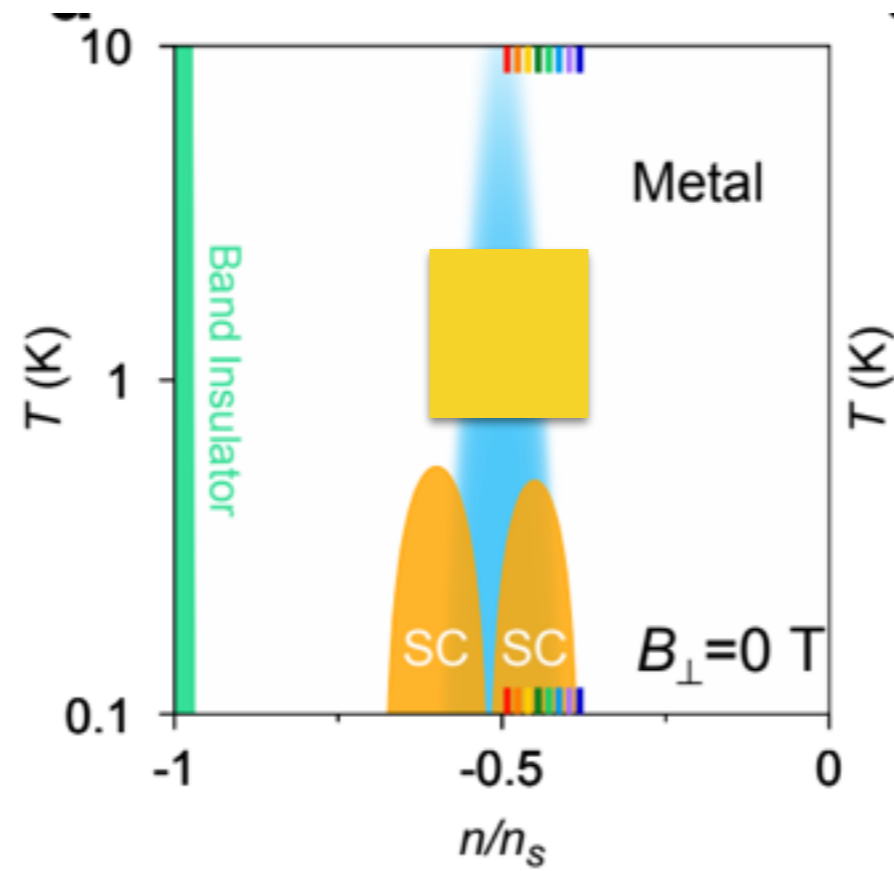
Bikash Padhi



Chandan Setty

Nano. Lett. 2018+
arXiv:1810.00884

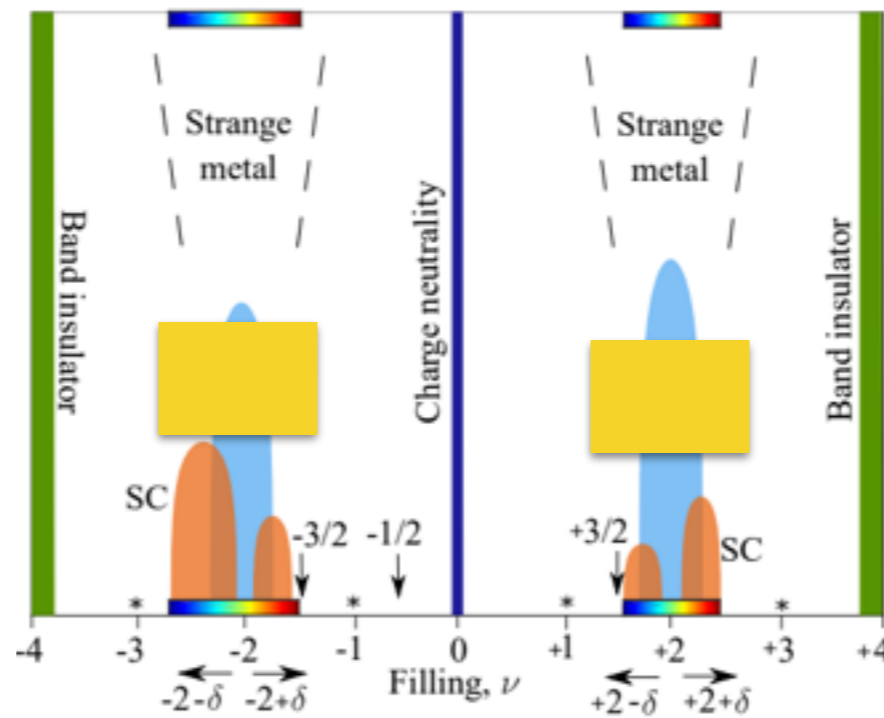
EFRC, NSF



2018

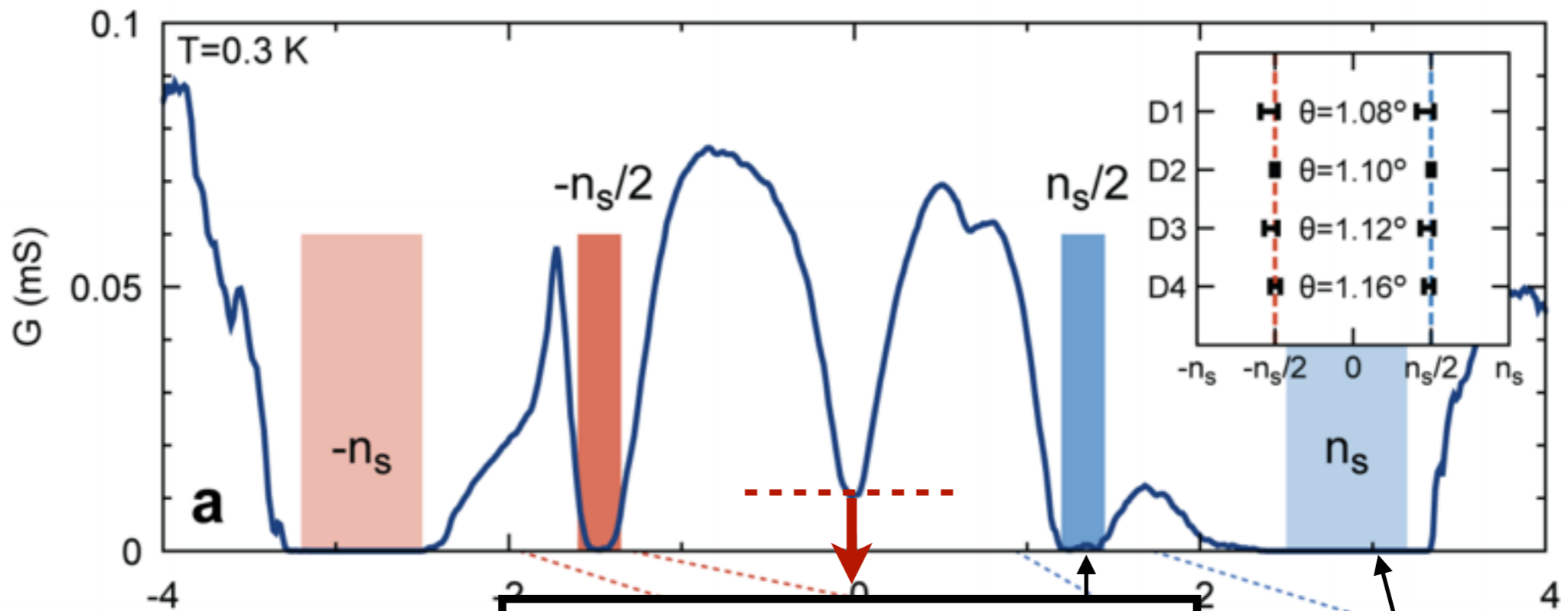
8months

2019



filling

$$\nu = n_e A_s$$



$$\sigma_{\min} = \frac{2e^2}{\pi^2 \hbar} \left(1 + \frac{1}{W} \right)$$

Koshino & Ando, PRB (2006)
Mott Insulator

$\nu = 4$

band insulator

Mottness in TBLG?

$$\nu = n_e A_s \quad \left\{ \begin{array}{l} \text{---} \\ \text{---} \end{array} \right\} \quad \nu_{\max} = 4$$

$$U = E^{\nu+1} + E^{\nu-1} - 2E^{\nu}$$

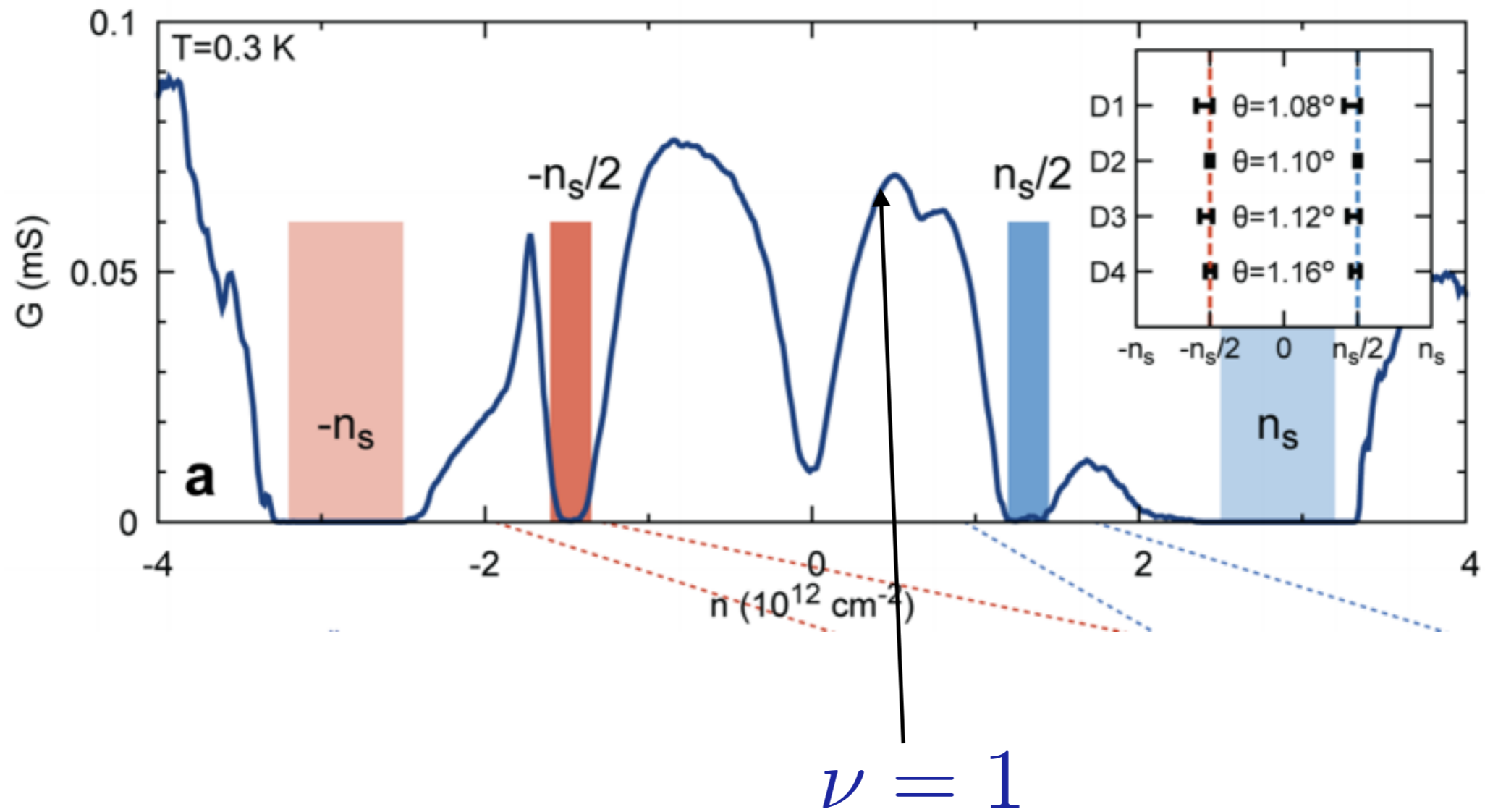
If Mott

$$\forall \nu < 4$$

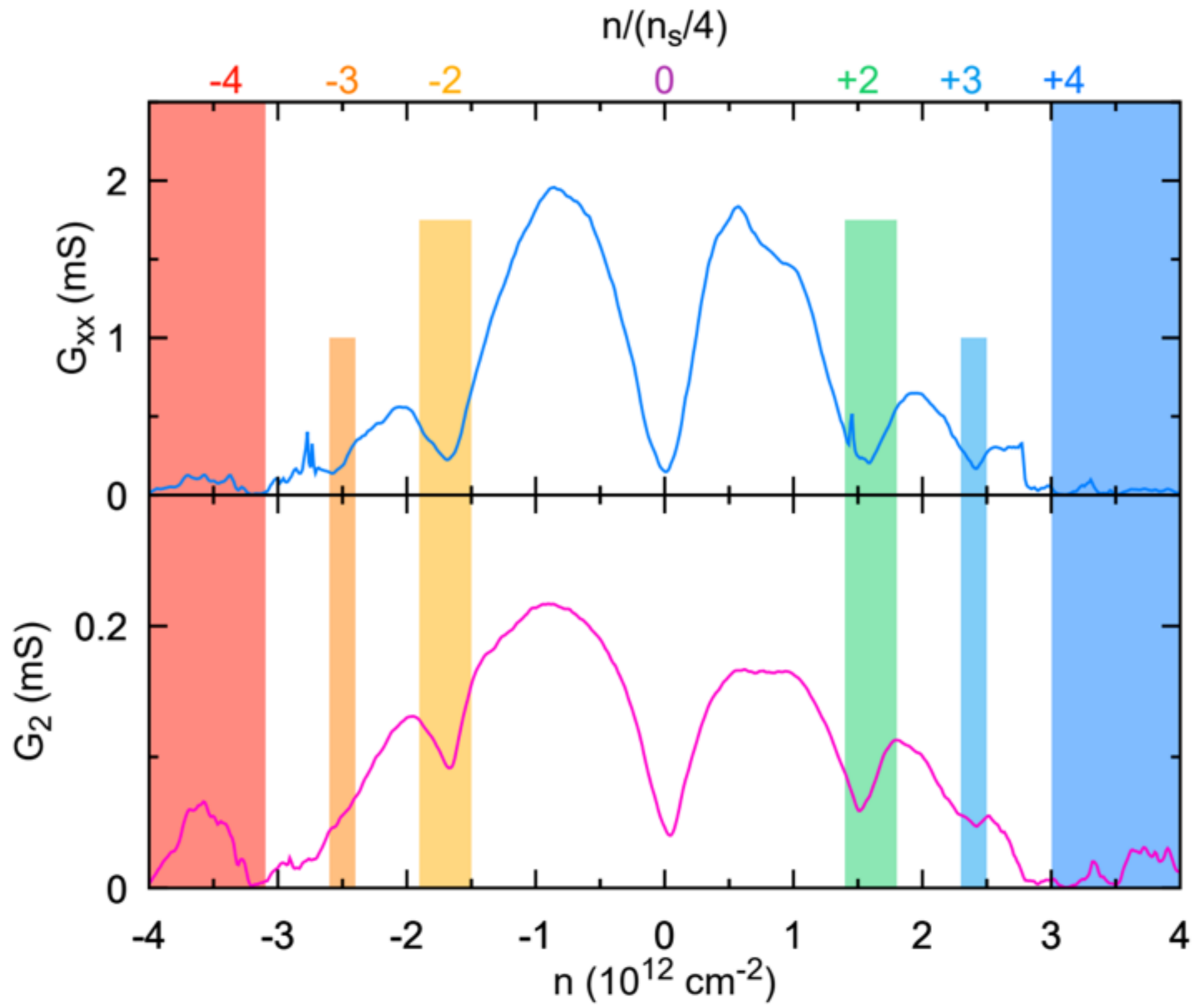


insulator

Where's the insulator at $\nu = 1$?



$\nu = 1 \neq$ insulator



Where's the insulator at $\nu = 1$?

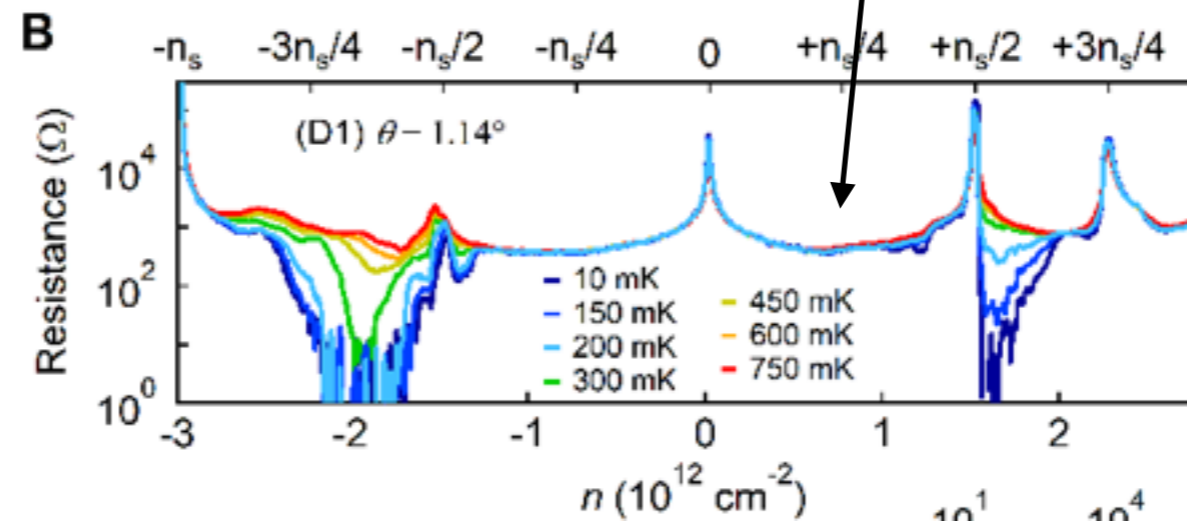
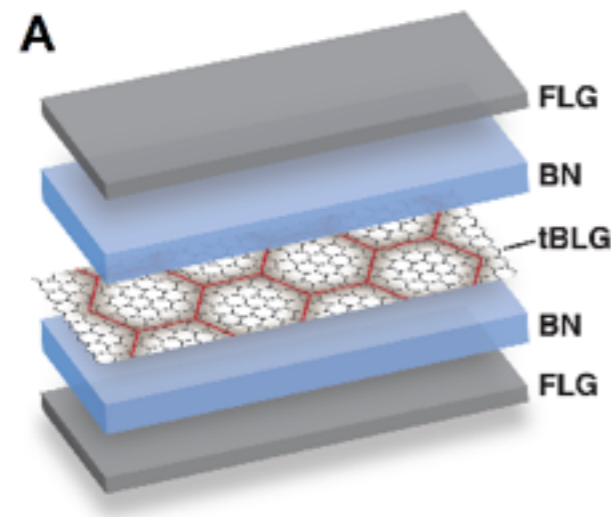
Tuning superconductivity in twisted bilayer graphene

Matthew Yankowitz^{1*}, Shaowen Chen^{1,2*}, Hryhoriy Polshyn^{3*}, K. Watanabe⁴,
T. Taniguchi⁴, David Graf⁵, Andrea F. Young^{3†}, and Cory R. Dean^{1†}

¹*Department of Physics, Columbia University, New York, NY, USA*

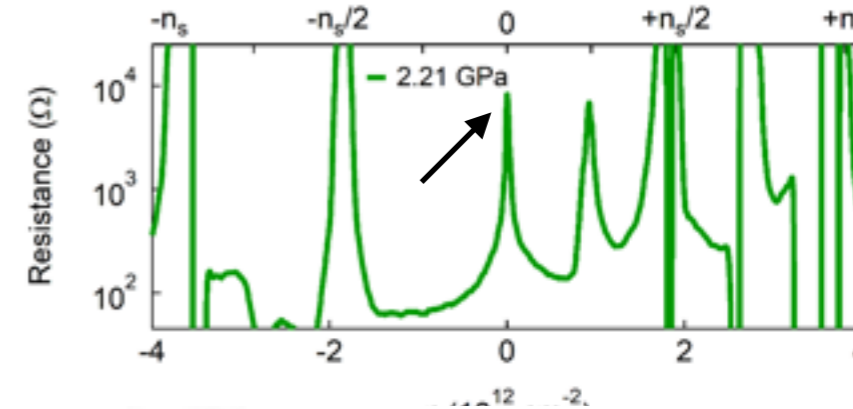
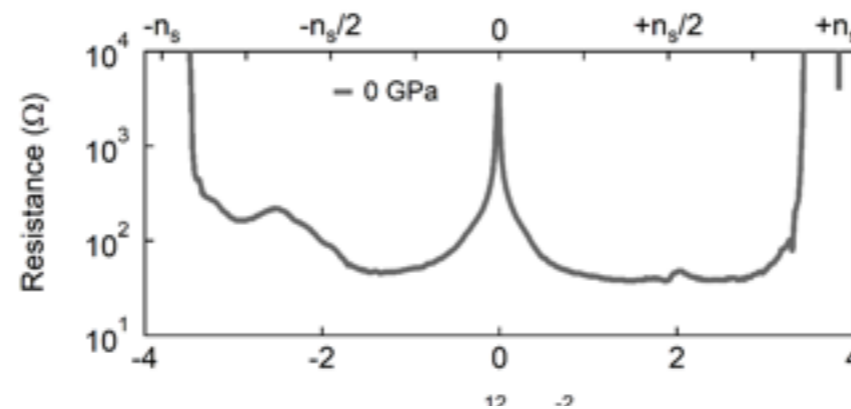
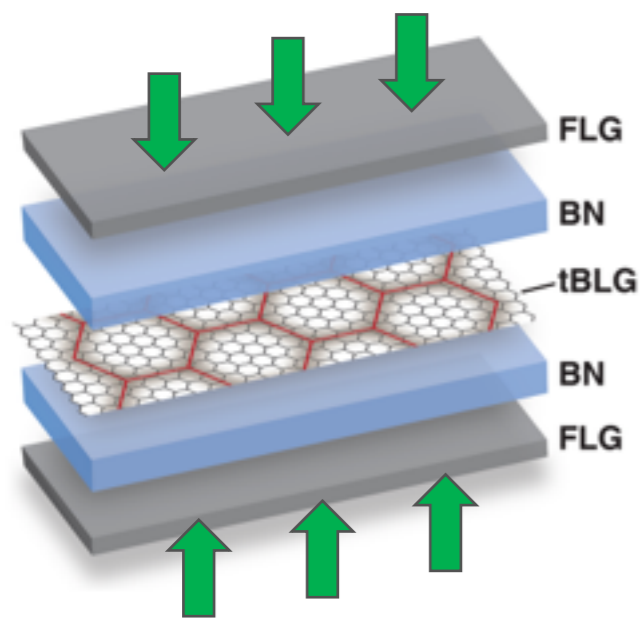
²*Department of Applied Physics and Applied Mathematics, Columbia University, New York, NY, USA*

³*Department of Physics, University of California, Santa Barbara, CA 93106*

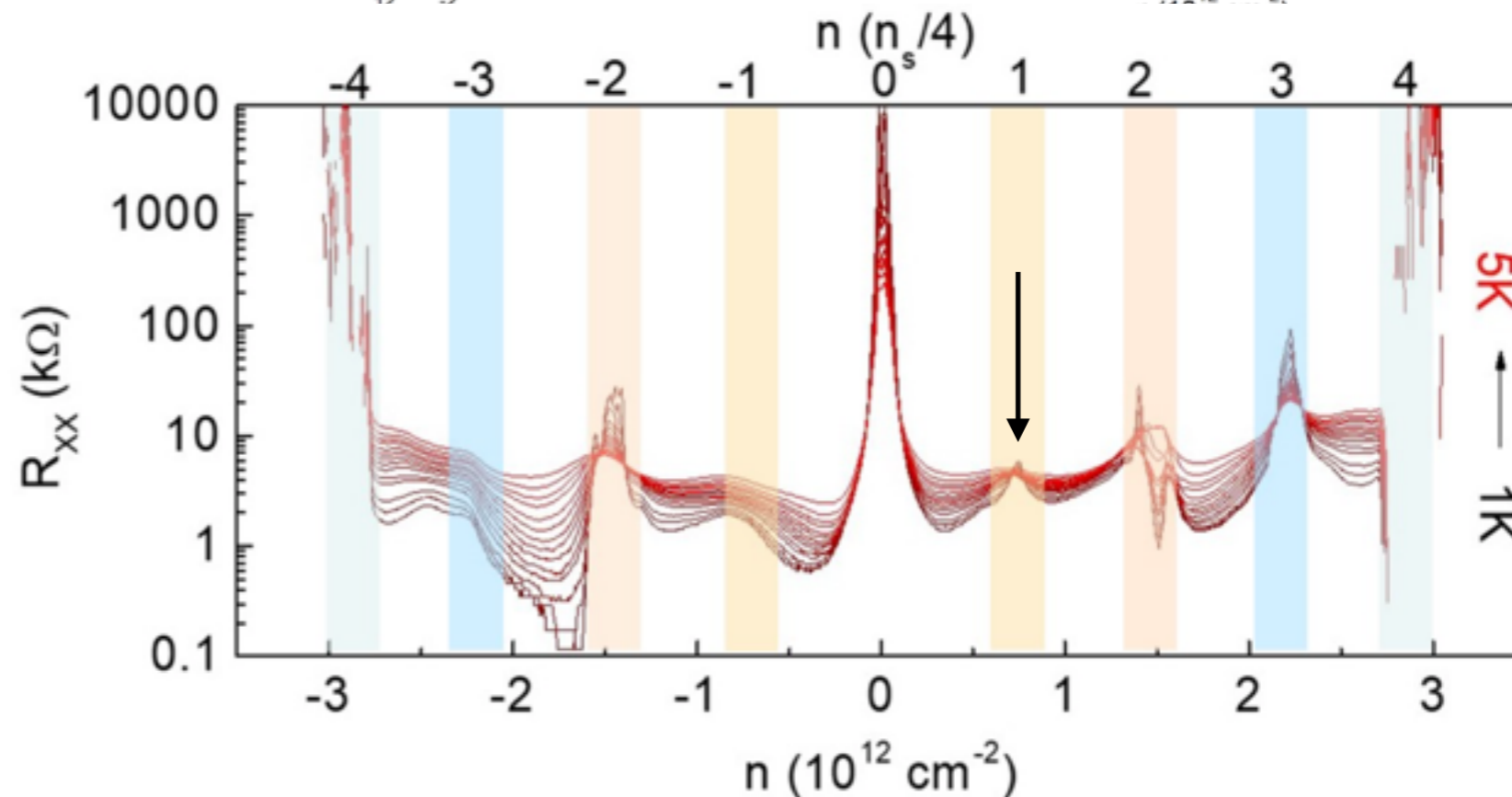


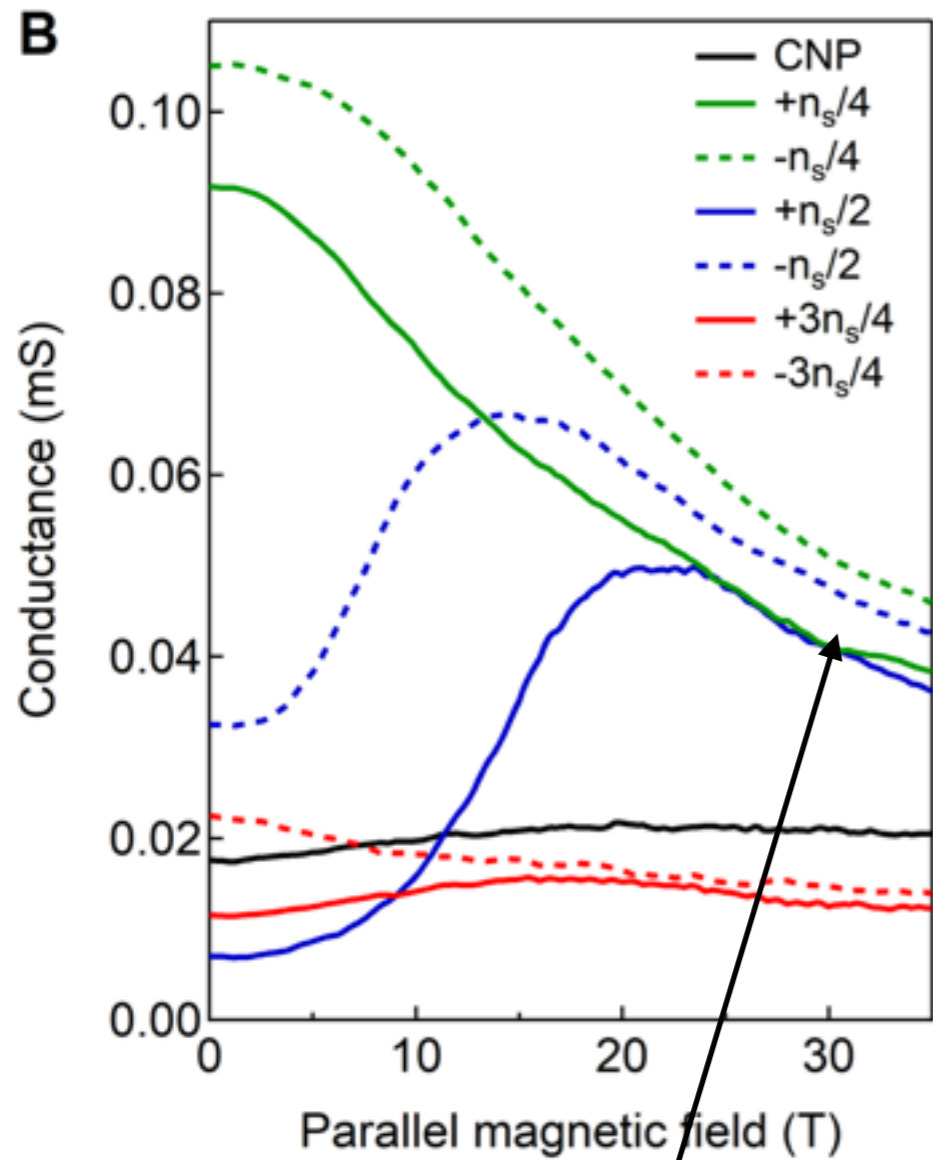
metallic

Modifying metallic-state resistivity with pressure

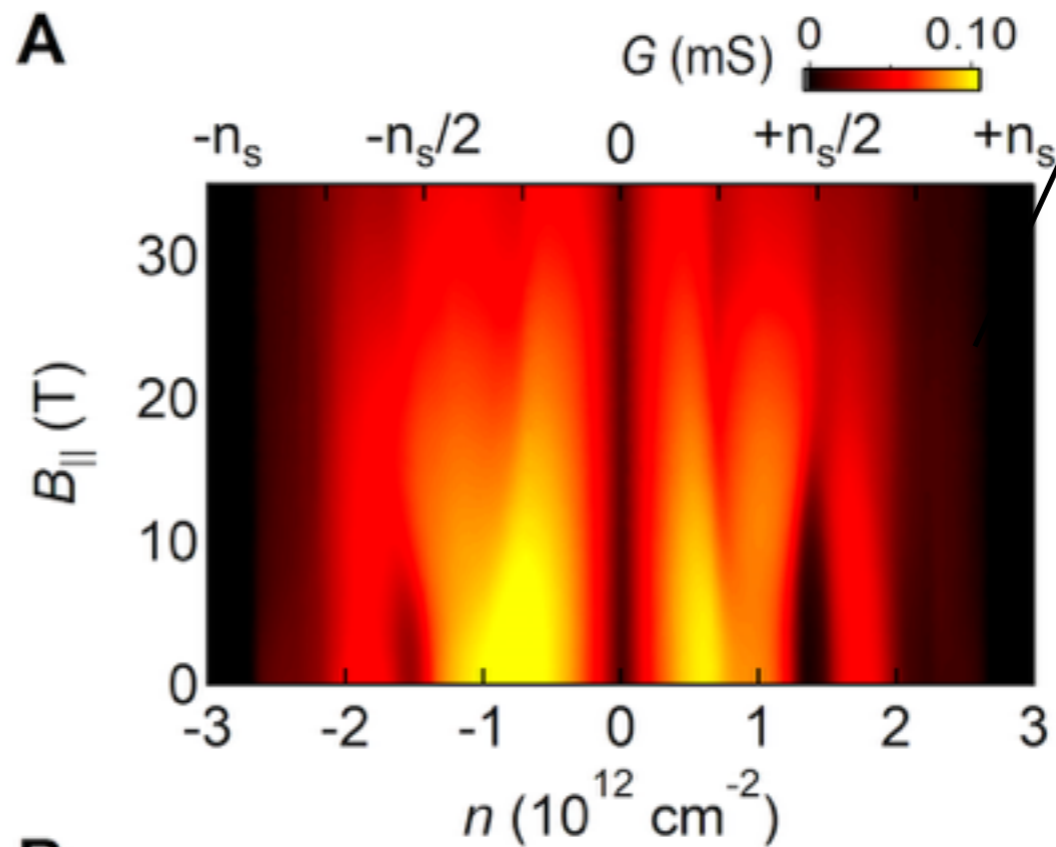


D. Efetov

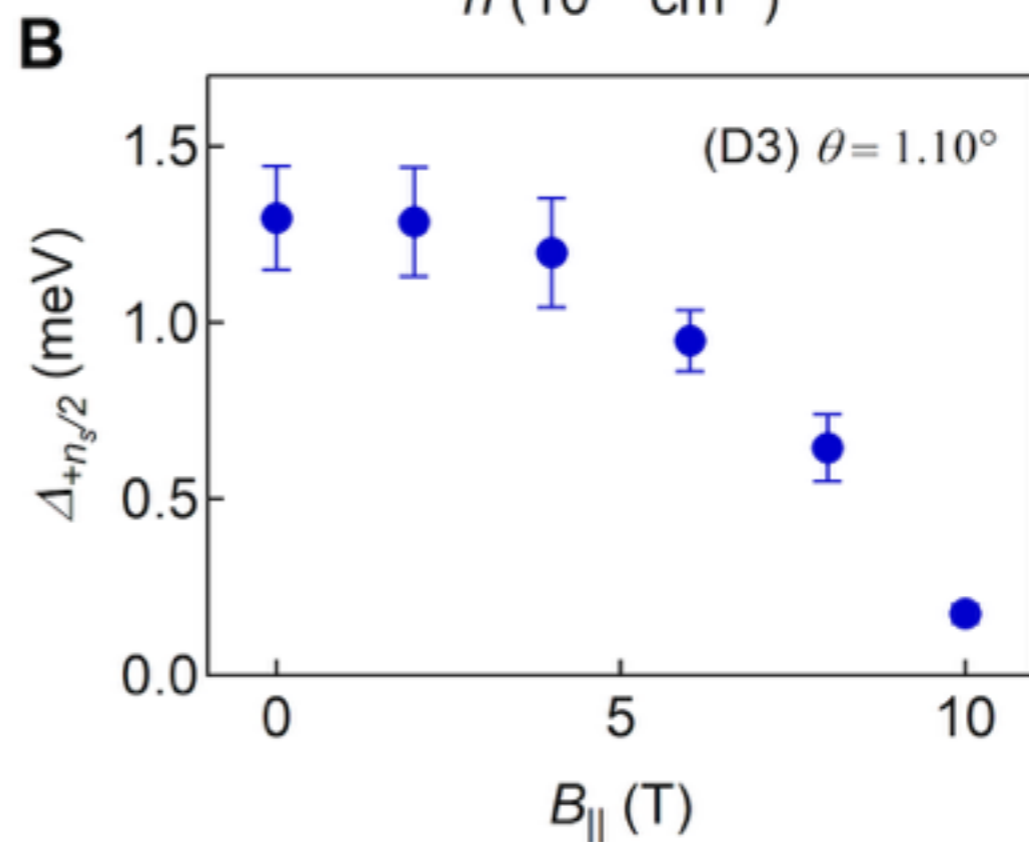




More
insulating

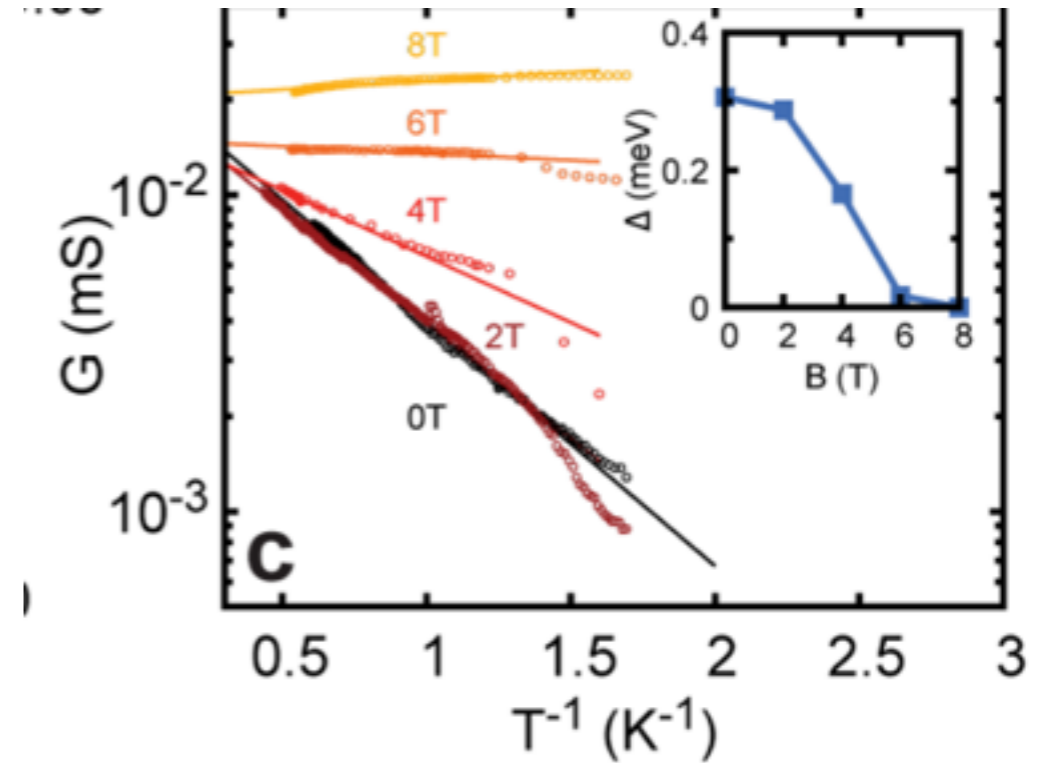
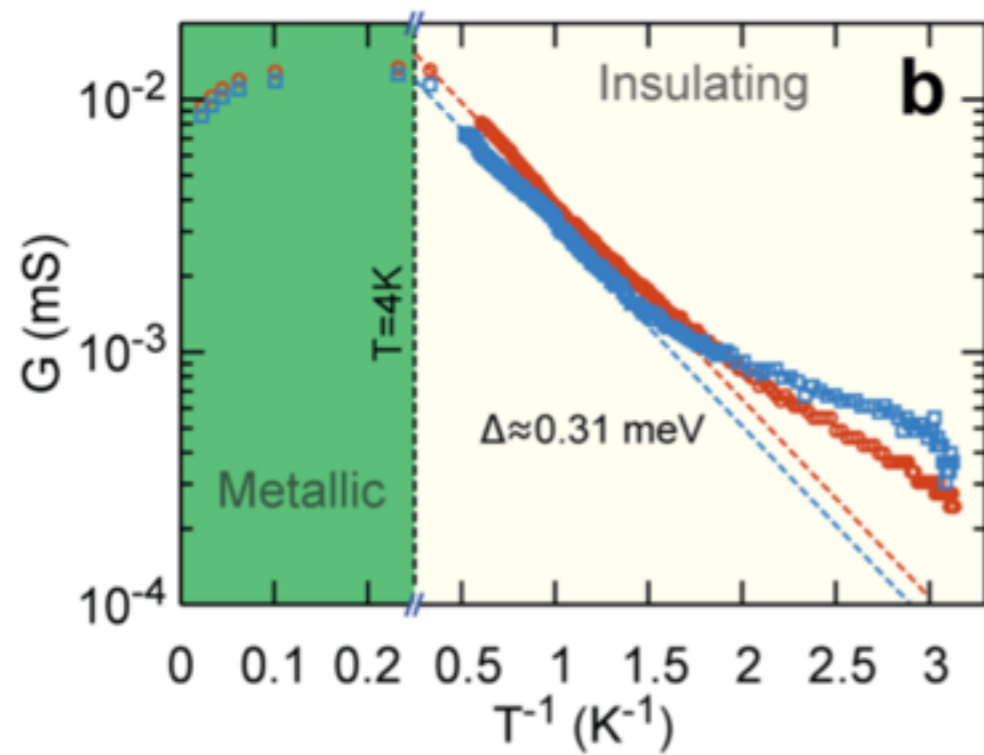


No effect



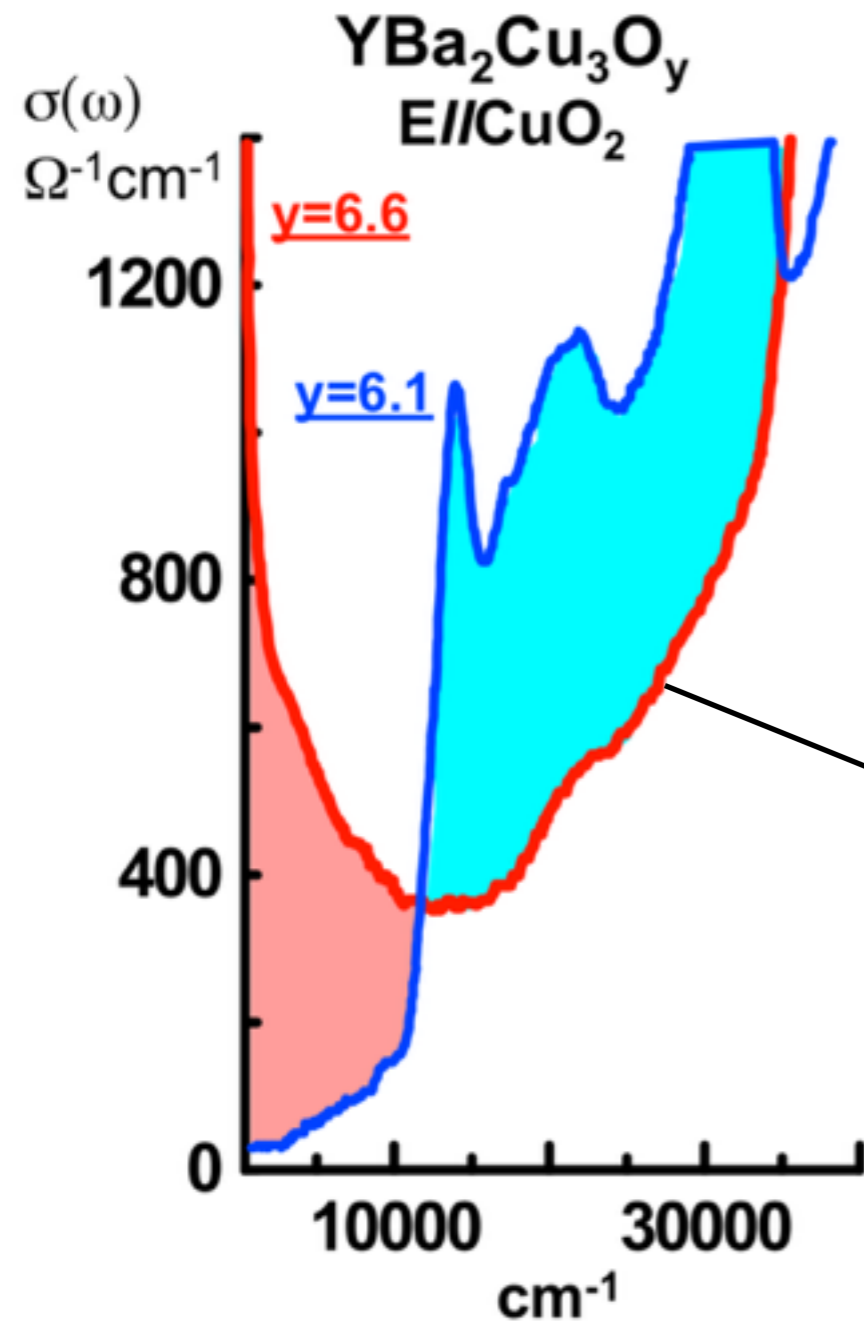
experiment

$$\Delta = .31\text{meV} \approx 4\text{K} \approx \mu_B B_c$$

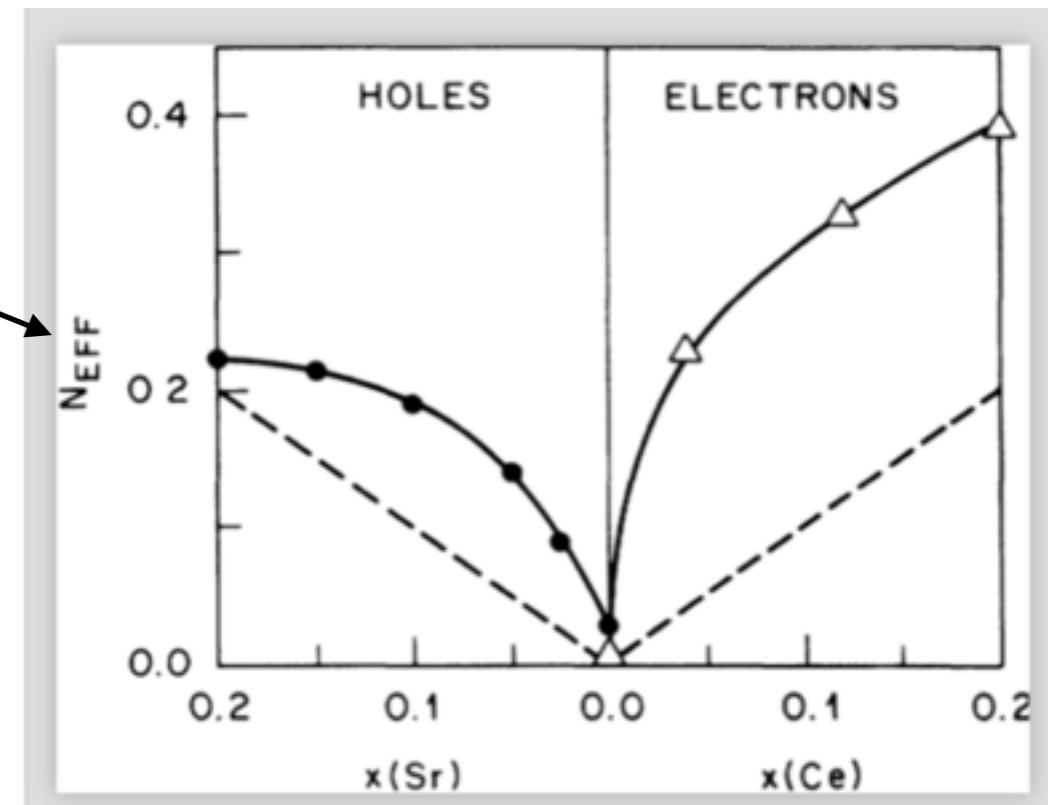
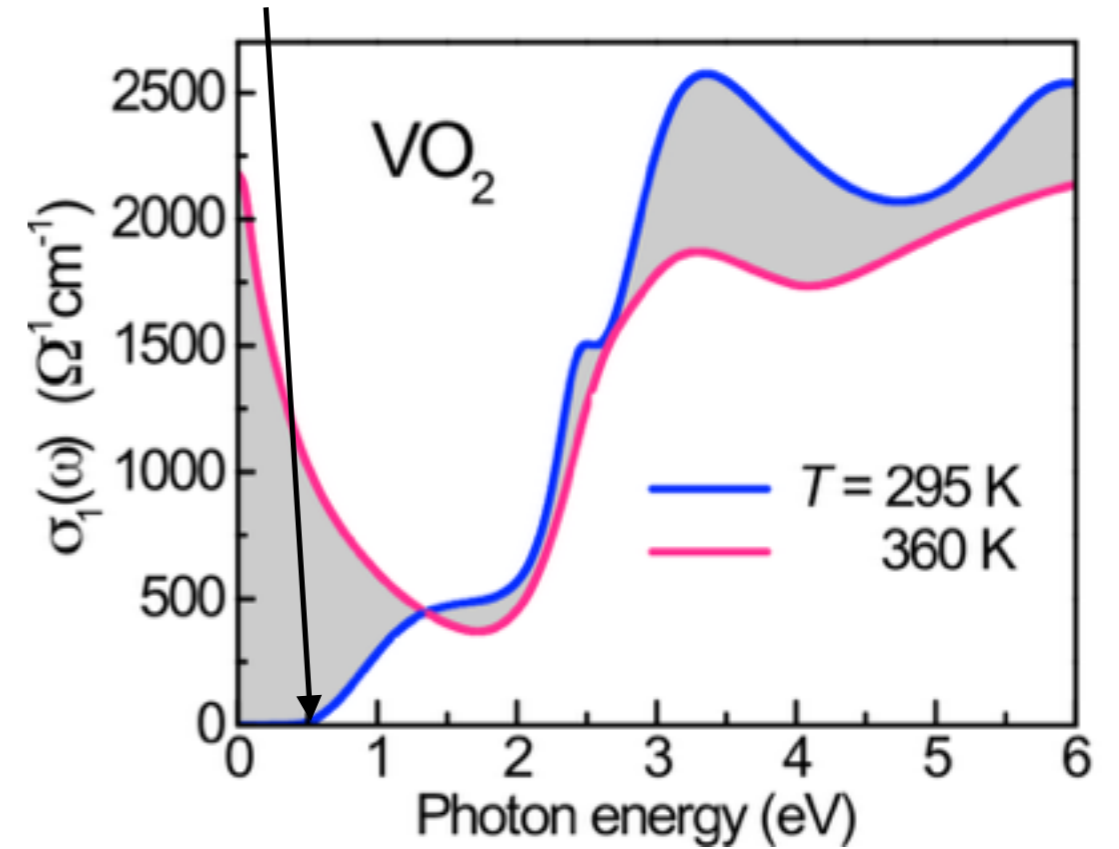


one energy scale

Mottness: 2-energy scales



$$\Delta = .5\text{eV} \gg 295\text{K}$$



Not Mott
Insulation

Experimental puzzle:
hierarchy of insulators

correlated 'insulator' $T < 4\text{K}$



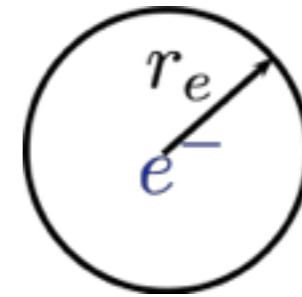
resistive but what?



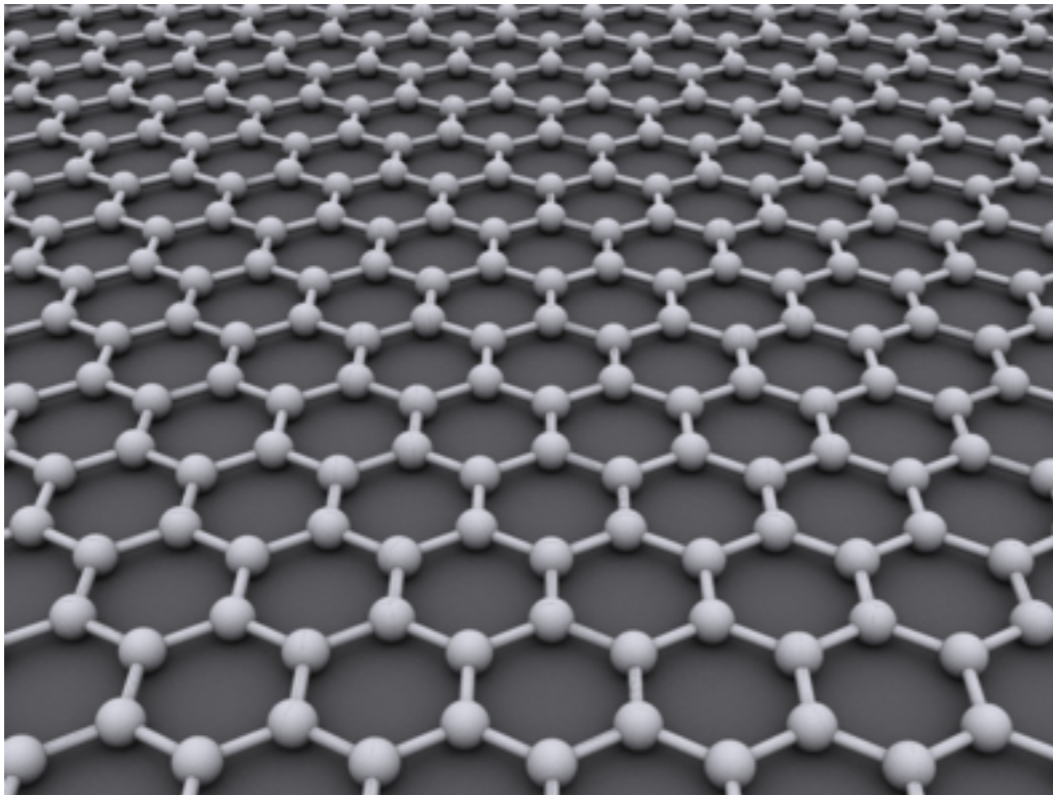
when are interactions important?

$$r_s = \frac{V_{ee}}{E_{\text{kin}}} = \frac{e^2}{\epsilon c_n} r_e^{n-1}$$

$c_n k^n$ → E_{kin}



Dirac spectrum
($n=1$)



density
independent

interactions
are irrelevant

$$H = v_F \boldsymbol{\sigma} \cdot (-i\boldsymbol{\partial} - \hat{\mathbf{A}})$$

non-abelian
gauge field
quenches kinetic energy

San-Jose, Gonzalez, Guinea PRL (2012)

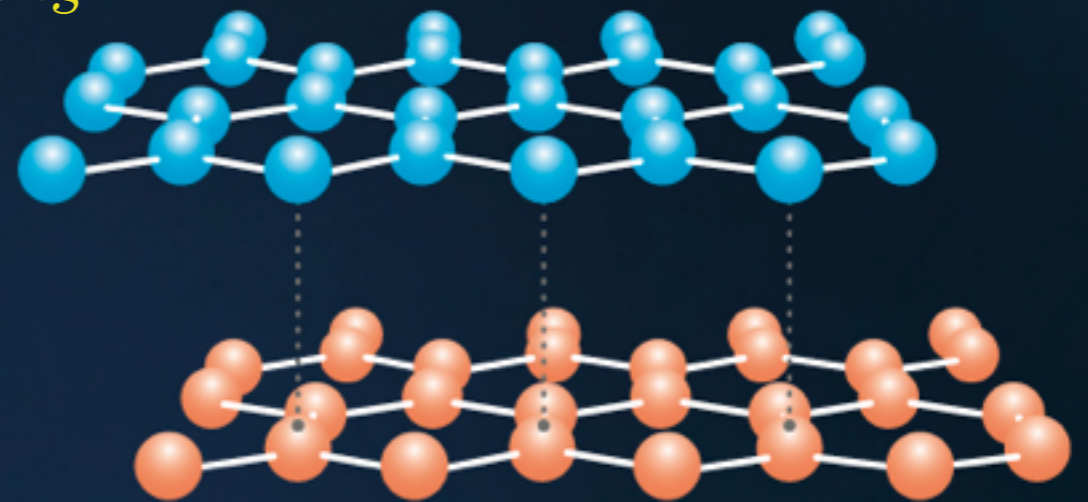
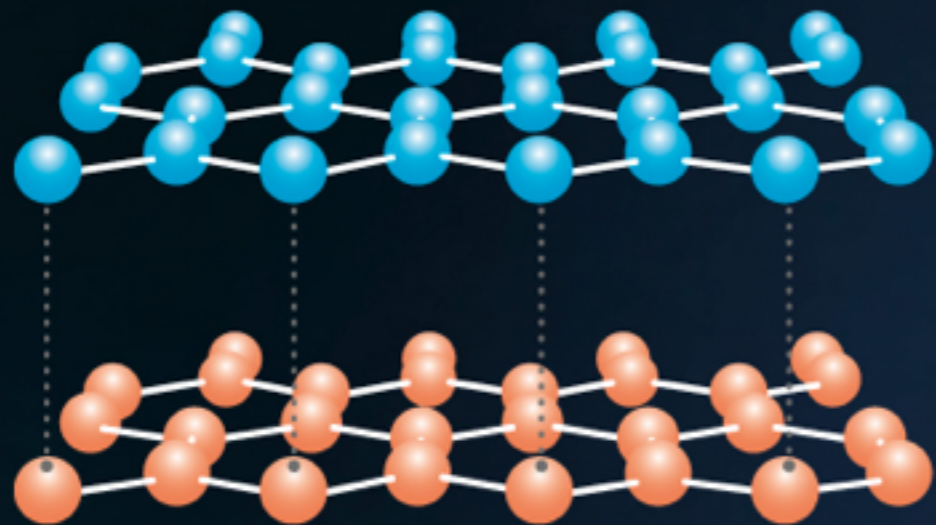
$$\ell_B \approx \lambda_s$$

3 Å

Wigner
crystallization
(Kindermann,...)

AA-stacking

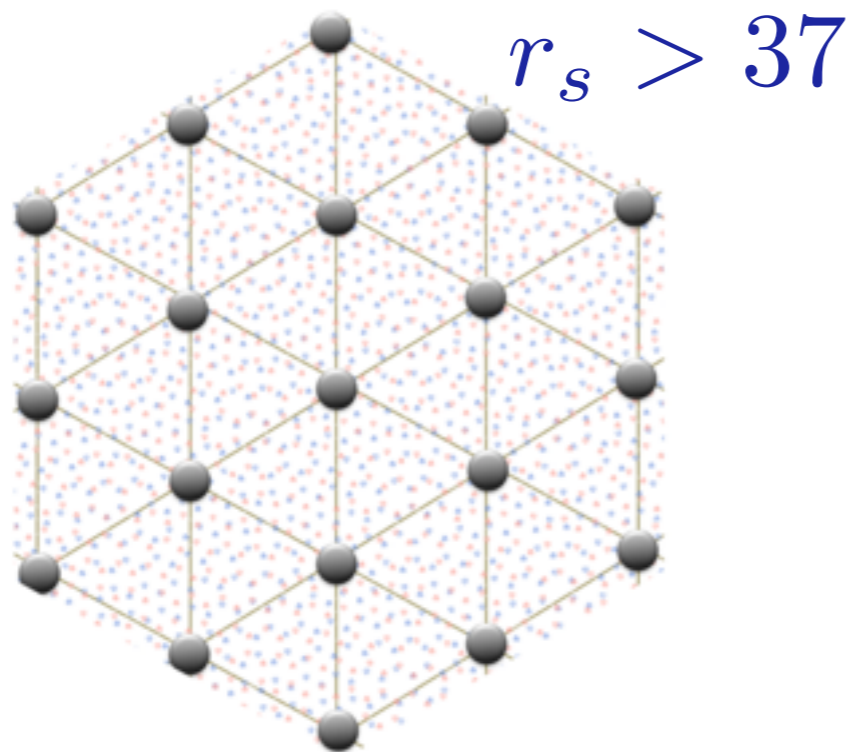
AB-stacking



Wigner crystallization

no underlying lattice

symmetry-broken state

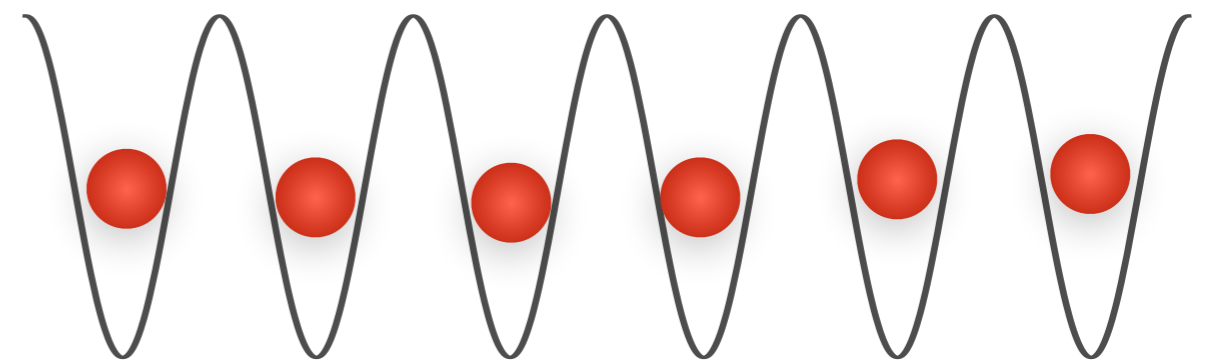


Mott Insulation

lattice needed

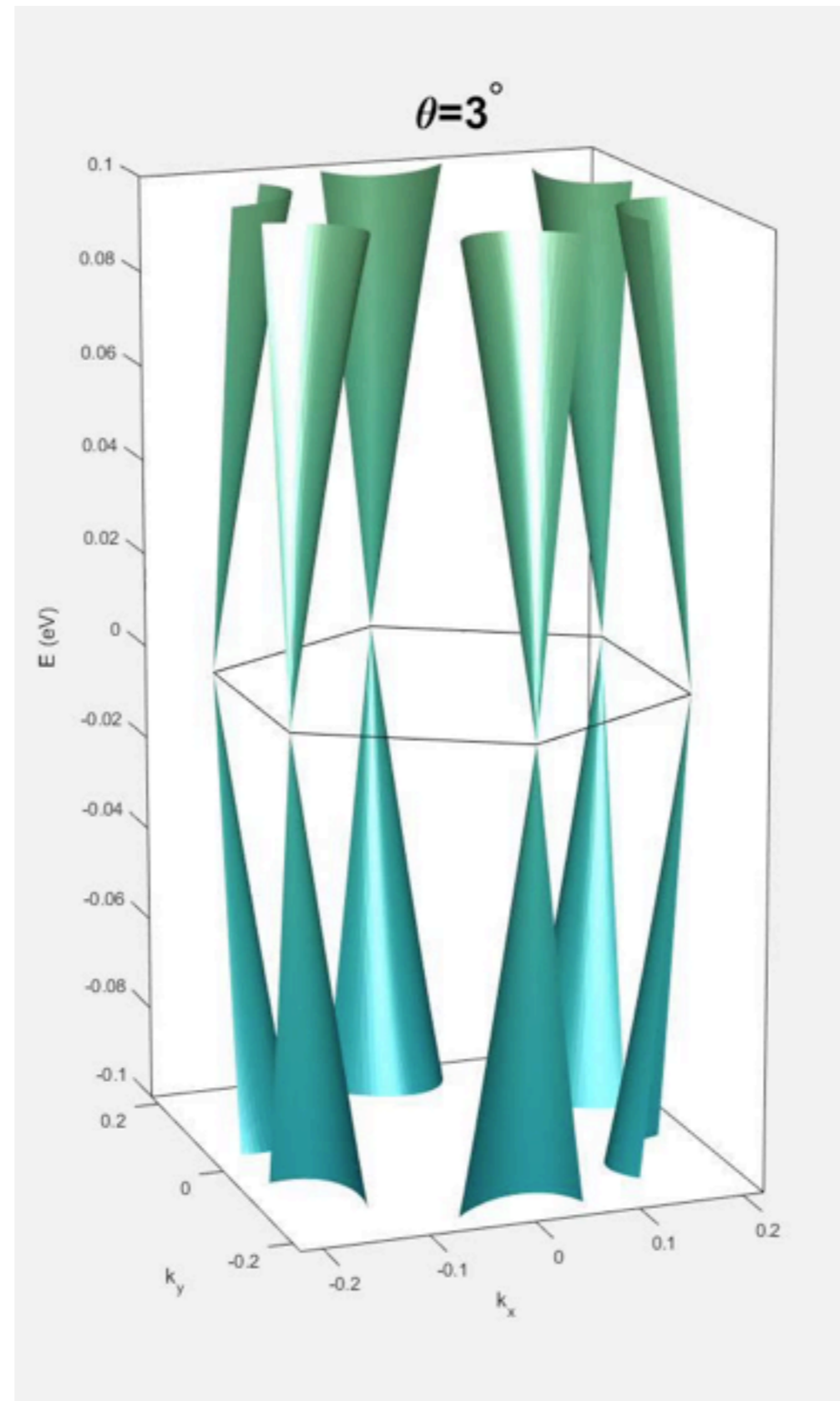
symmetry intact

$$U \gg W$$



$$n_e^{1/2} a_0^* = \frac{a_0^*}{r_e} = r_s^{-1} \approx O(1)$$

Evolution of Band Structure



recall

Dirac spectrum

$$r_s = \text{const}$$

TBLG: deviation from Dirac

$$H_\theta(\mathbf{k}) = \begin{pmatrix} 0 & v_F \mathbf{k}^\dagger + \frac{1}{2m} \mathbf{k}^2 \\ v_F \mathbf{k} + \frac{1}{2m} (\mathbf{k}^\dagger)^2 & 0 \end{pmatrix}$$

flat-band
condition

$$\tilde{v} = \frac{v_0}{v_F} = \frac{\partial_{\mathbf{k}} E_{\mathbf{k}}|_{\mathbf{k}=\Gamma}}{\partial_{\mathbf{k}} E_{\mathbf{k}}|_{\mathbf{k}=\mathbf{K}, \mathbf{K}'}} = \frac{1 + 6\kappa^2}{1 - 3\kappa^2}$$

$$\kappa = 1/\sqrt{3}$$

compute r_s

$r_s^{(1)}(\theta)$

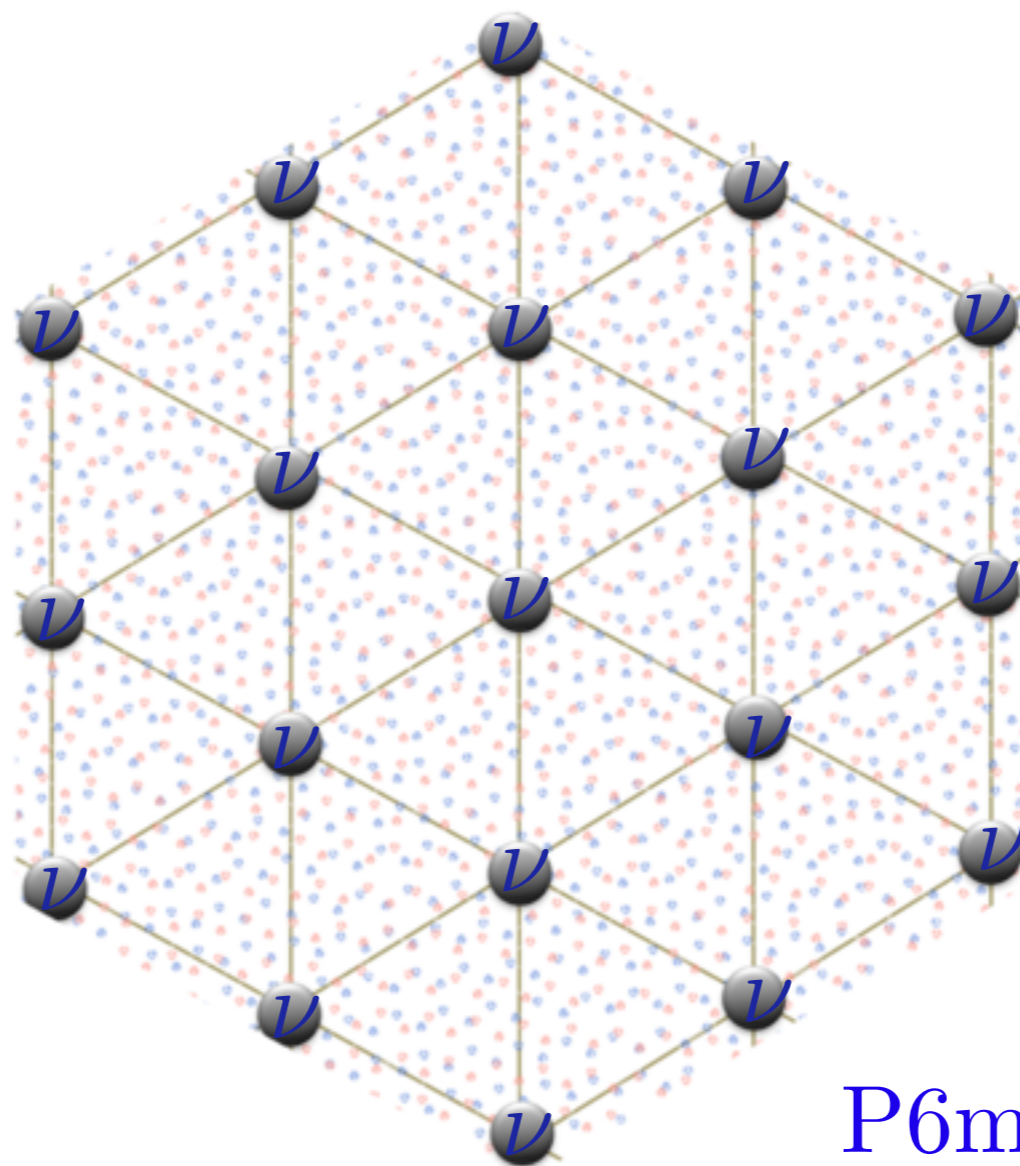
$$r_s(\theta, \nu) = \frac{\alpha}{\epsilon} \tilde{v}(\theta) \times \left(1 - \gamma + \frac{1}{2} \gamma^2\right)^{-1/2}$$

$$\alpha = \frac{e^2}{\hbar v_0} \sim 2$$

$$\gamma = \frac{\sqrt{2}\hbar}{m_* v_F} \frac{1}{2r_e} = \frac{\theta}{a} \frac{\hbar}{m_* v_F} \sqrt{\frac{\pi\nu}{\sqrt{3}}} \sim \frac{\theta^\circ}{100} \frac{\tilde{v}(\theta)}{\tilde{m}(\theta, \nu)} \sqrt{\nu}$$

filling dependence

Wigner Crystals for $\nu \neq 1$

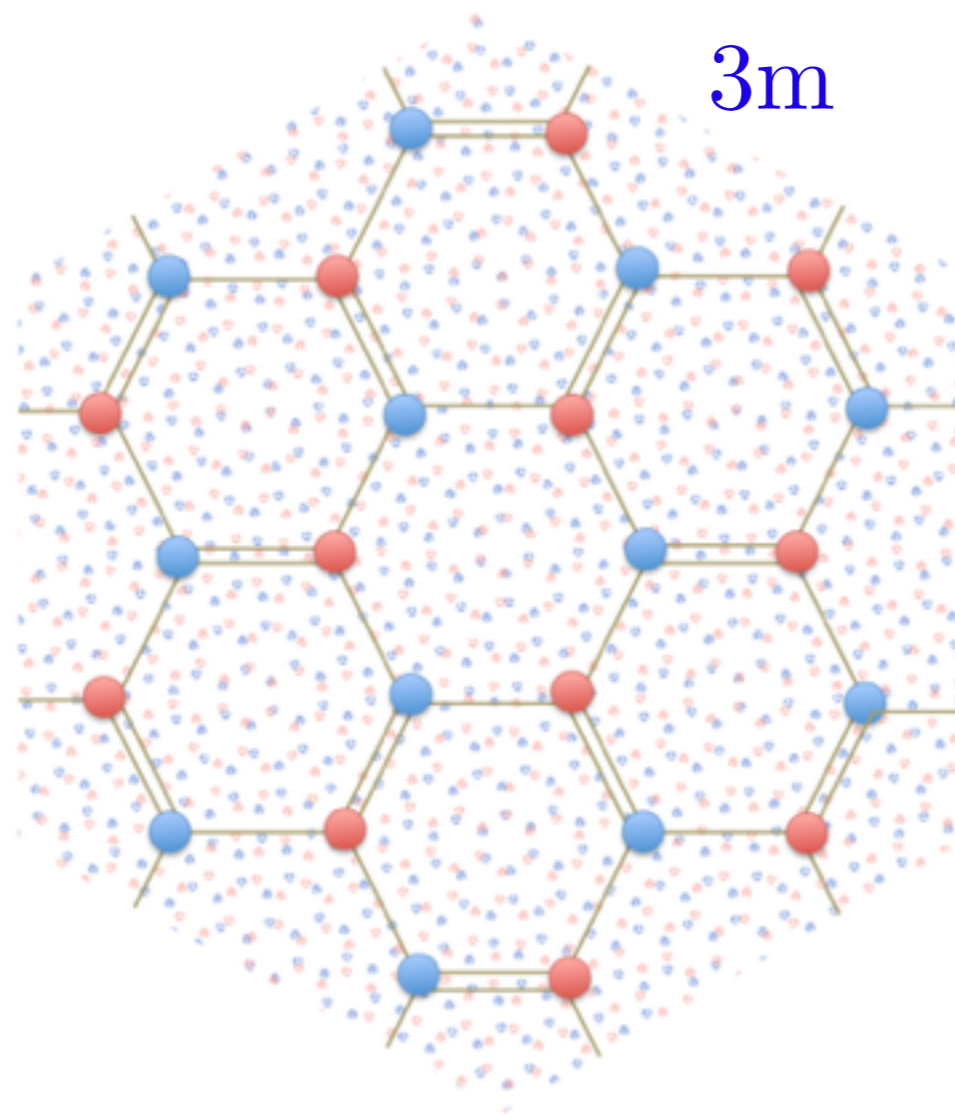


$$r_s^{\text{crit}} = \frac{37}{\nu^2}$$

lower bound

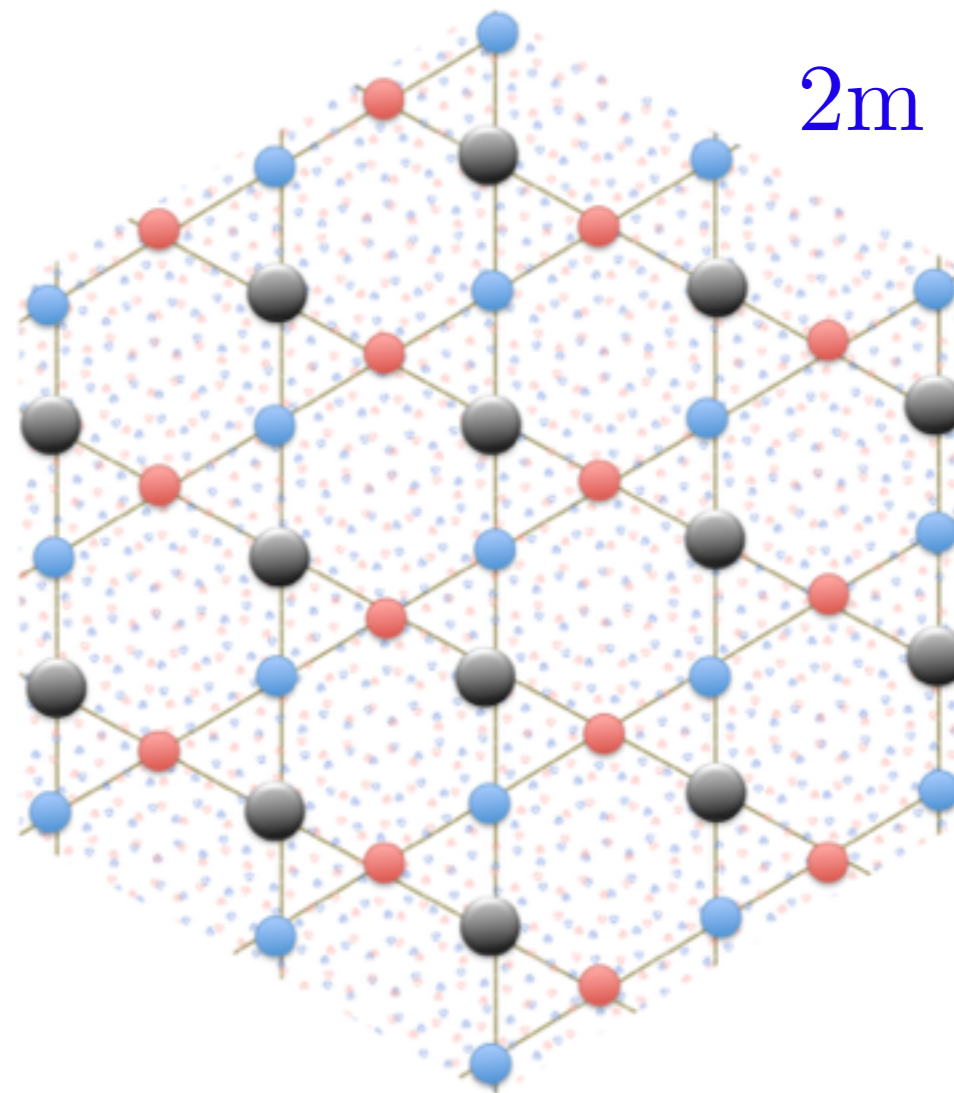
realistic lattices for $\nu > 1$

$\nu = 2$



$$\lambda_s \rightarrow \lambda_s / \sqrt{3}$$

$\nu = 3$



$$\lambda_s \rightarrow \lambda_s / 2$$

$$r_s = \frac{a}{a_0^*}$$



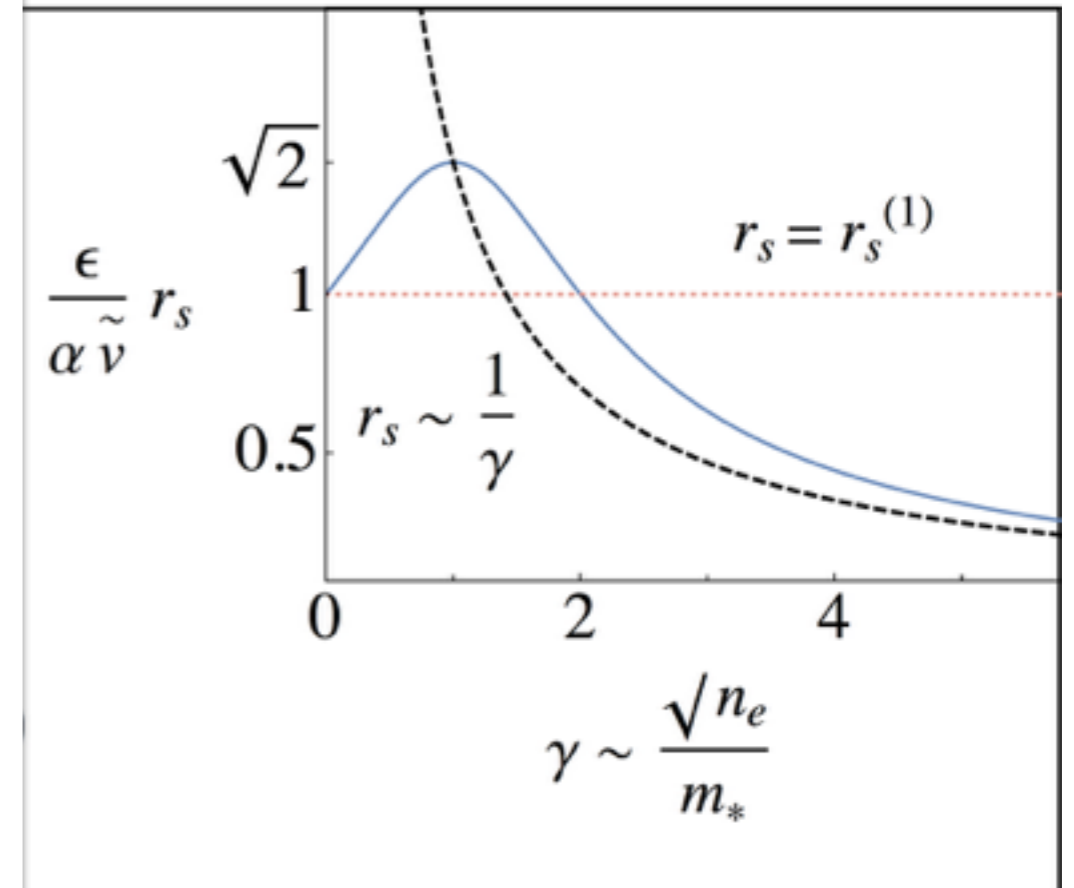
$$r_s = d_\nu 37$$

$$d_2 = 1/\sqrt{3}$$

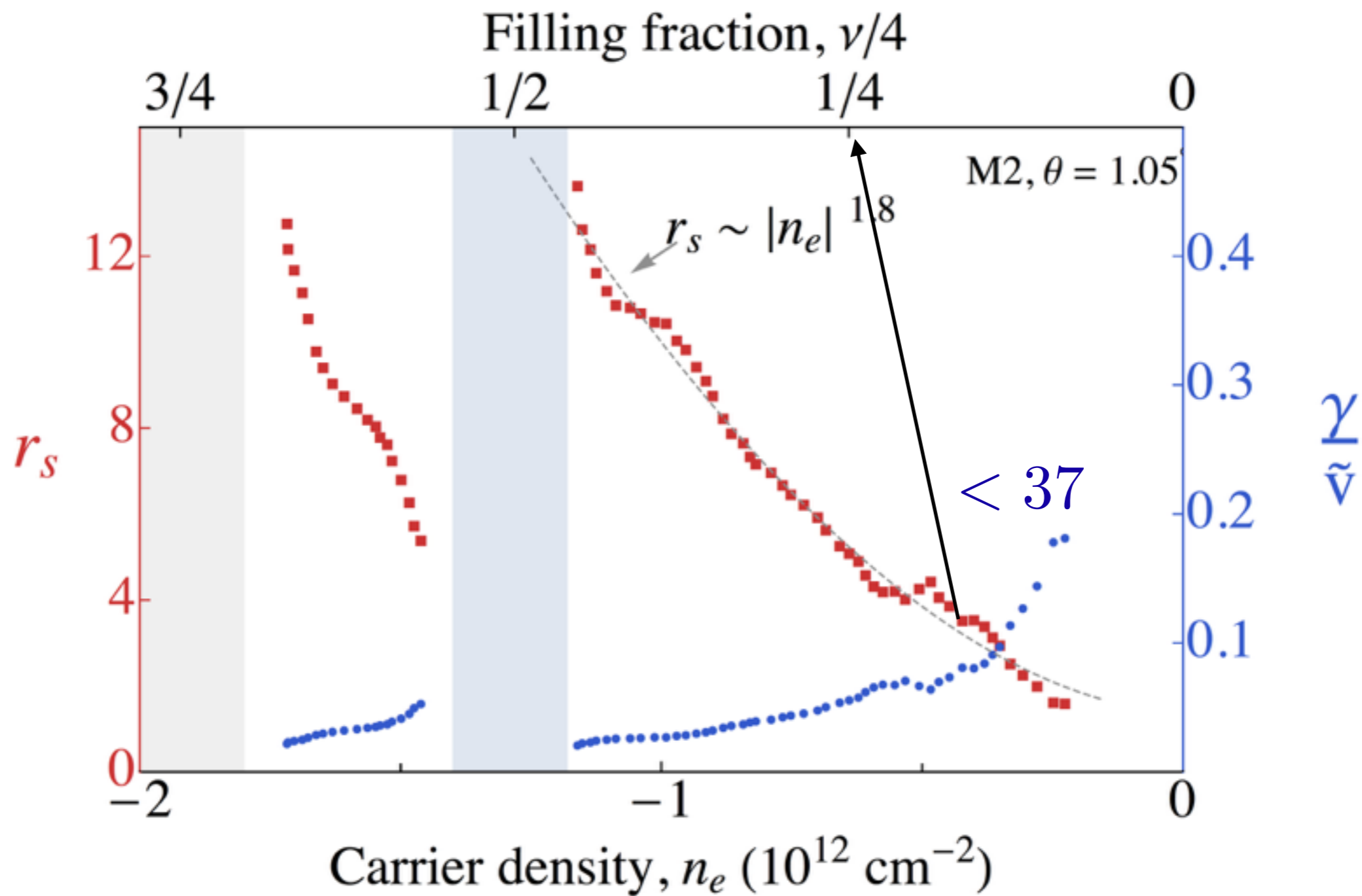
$$d_3 = 1/2$$

$$\frac{37}{\nu^2} < r_s^{\text{crit}} < d_\nu 37$$

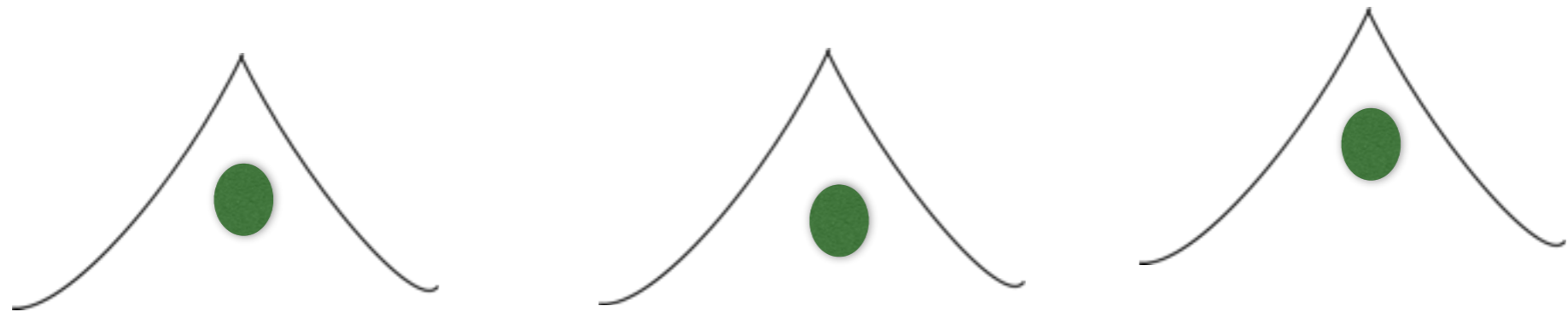
experiments



using experimental data for effective mass



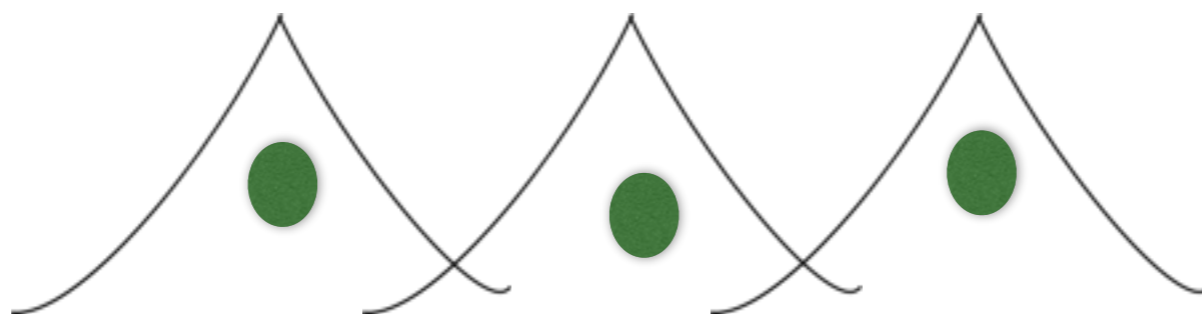
What happened to Motttness?



Mott Criterion

$$n_e^{1/2} a_0^* \approx 1$$

dynamics



$$m_*(\theta, \nu) = m_e \frac{(\gamma^2 - 2\gamma + 1)^{3/2}}{\gamma^3 - 3\gamma^2 + .45\gamma - 2}$$

Mott criterion

$$\frac{\lambda_s(\theta)}{a_0} \frac{m_*(\theta, \nu)}{m_e \sqrt{\nu}} \lesssim \epsilon$$

150 – 200

const

never satisfied

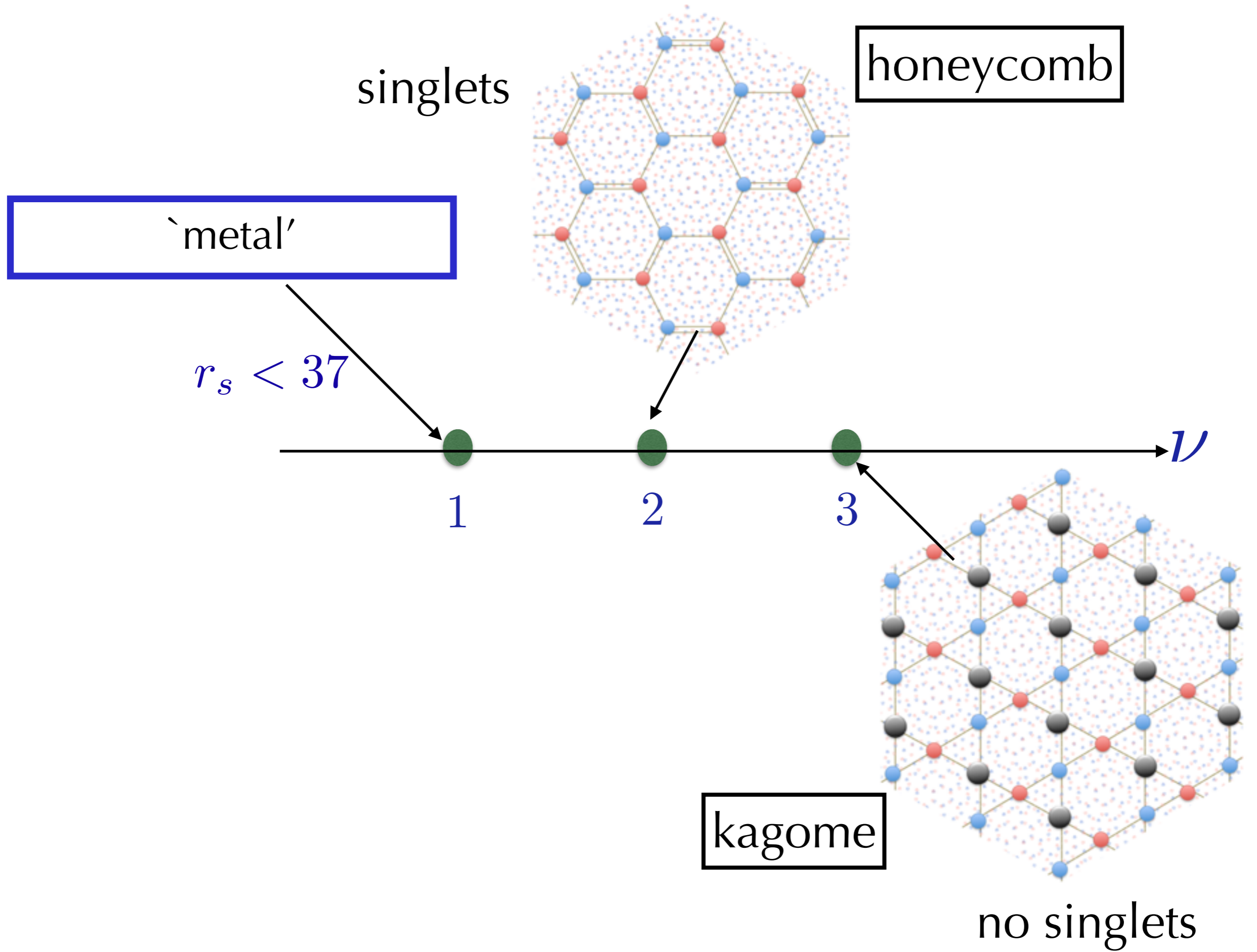
$$m_*/m_e = \tilde{v} \sqrt{h^2 n_e / 8\pi v_0^2 m_e^2}$$

Mott criterion

$$r_s^{(1)}(\theta) \lesssim \frac{1}{\sqrt{\pi}}$$

$$\theta - \theta_{\text{magic}} = .7$$

$$n_e \approx 10^4 n_s$$



Effects of Pressure

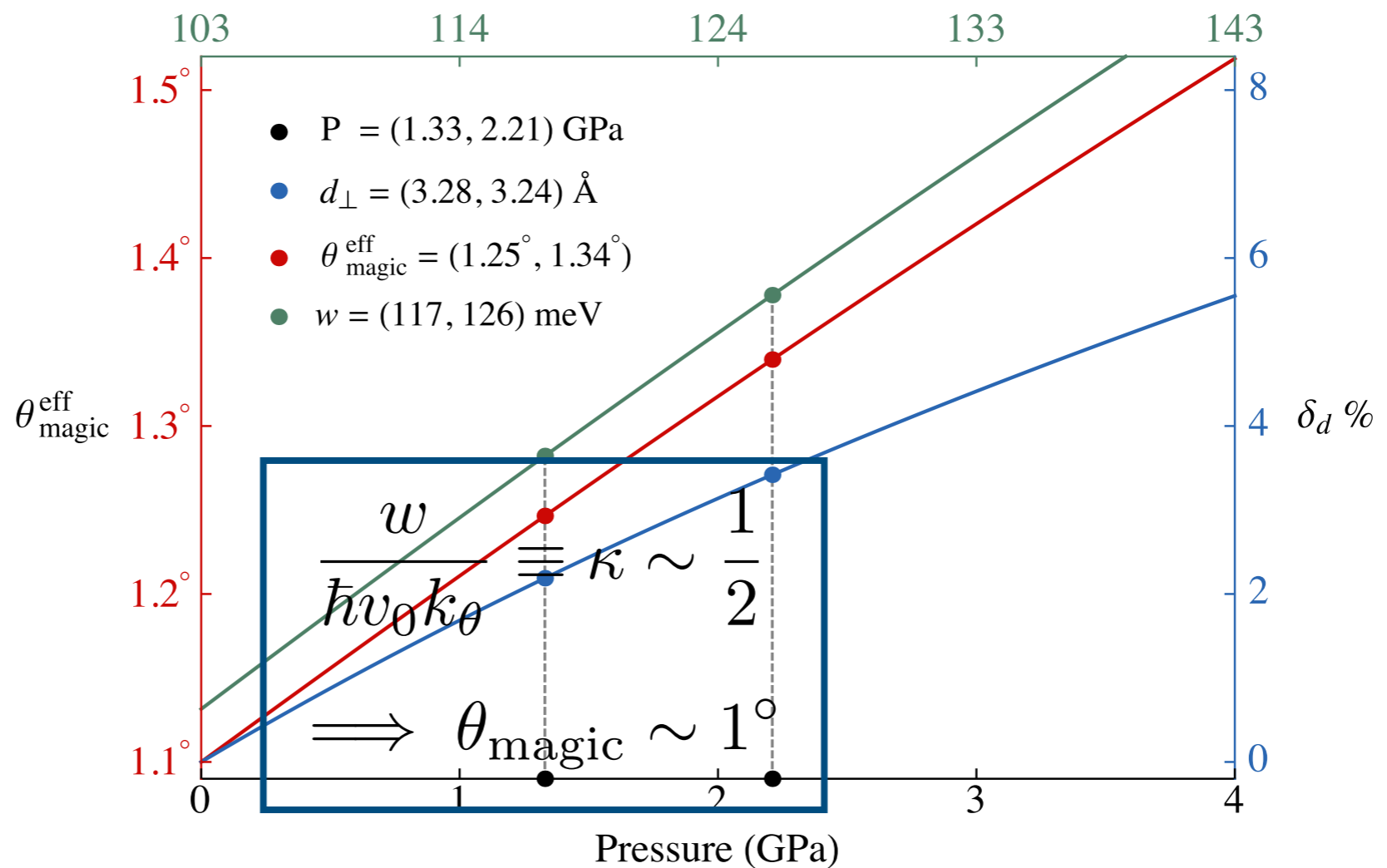
Inter-particle
Distance

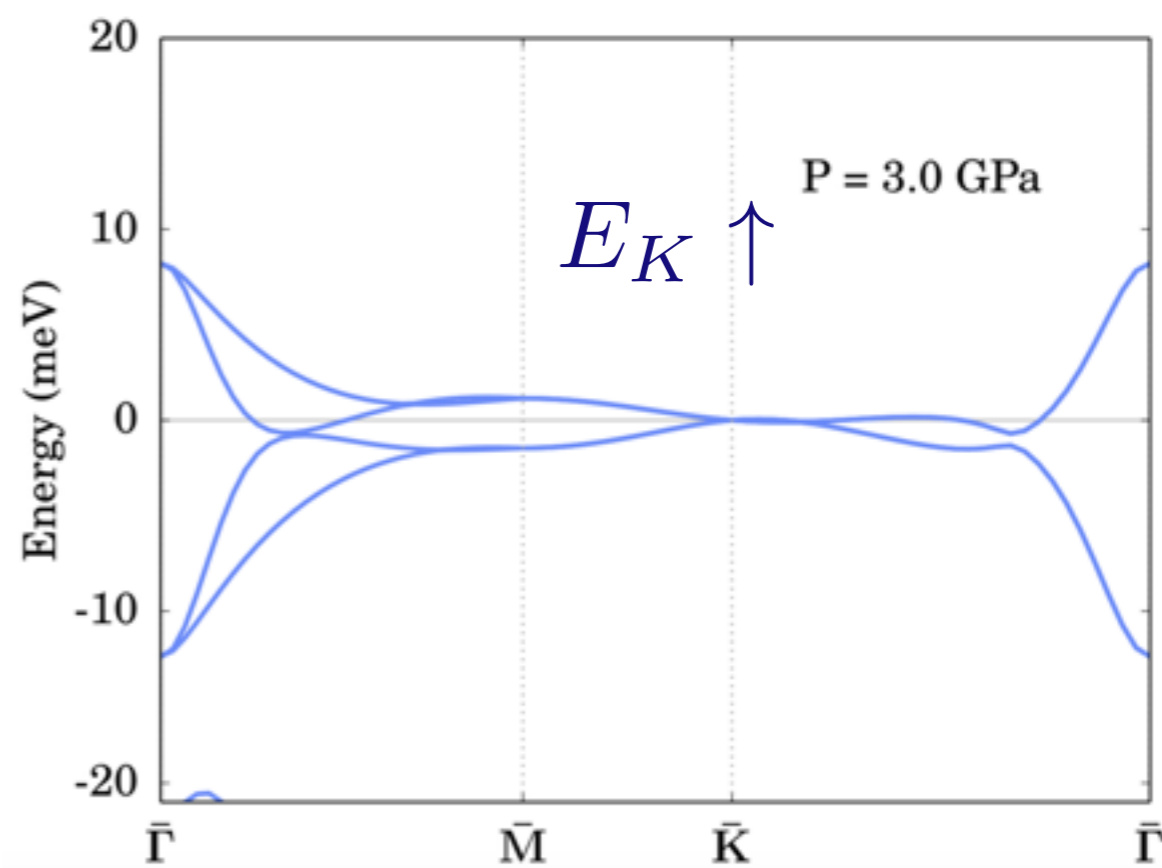
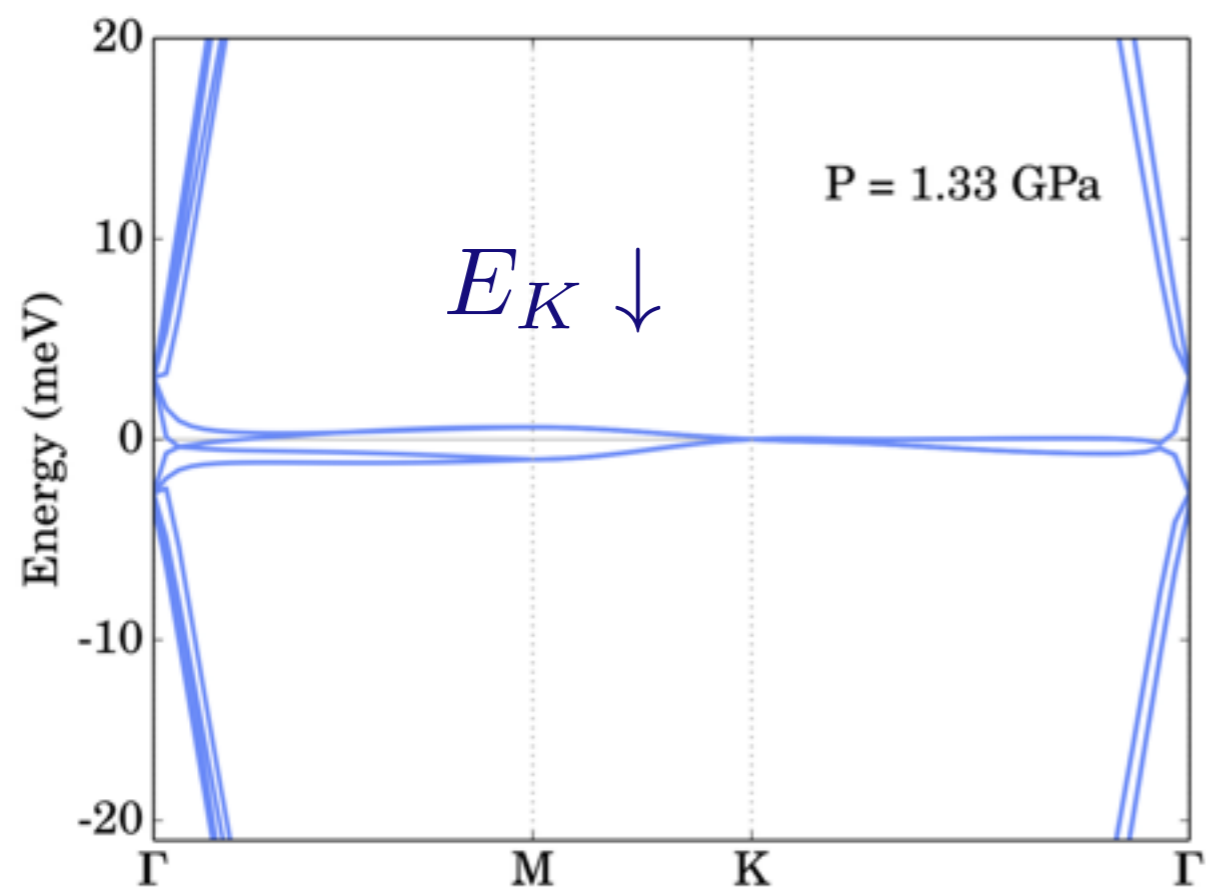
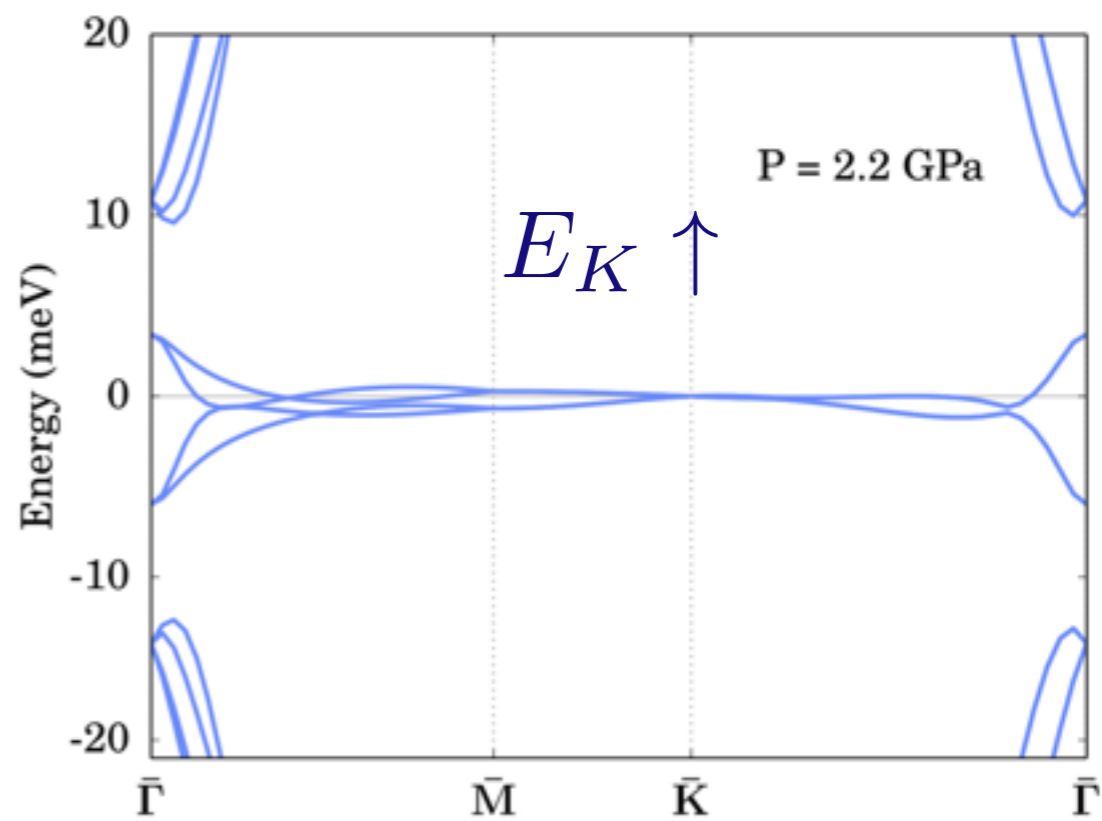
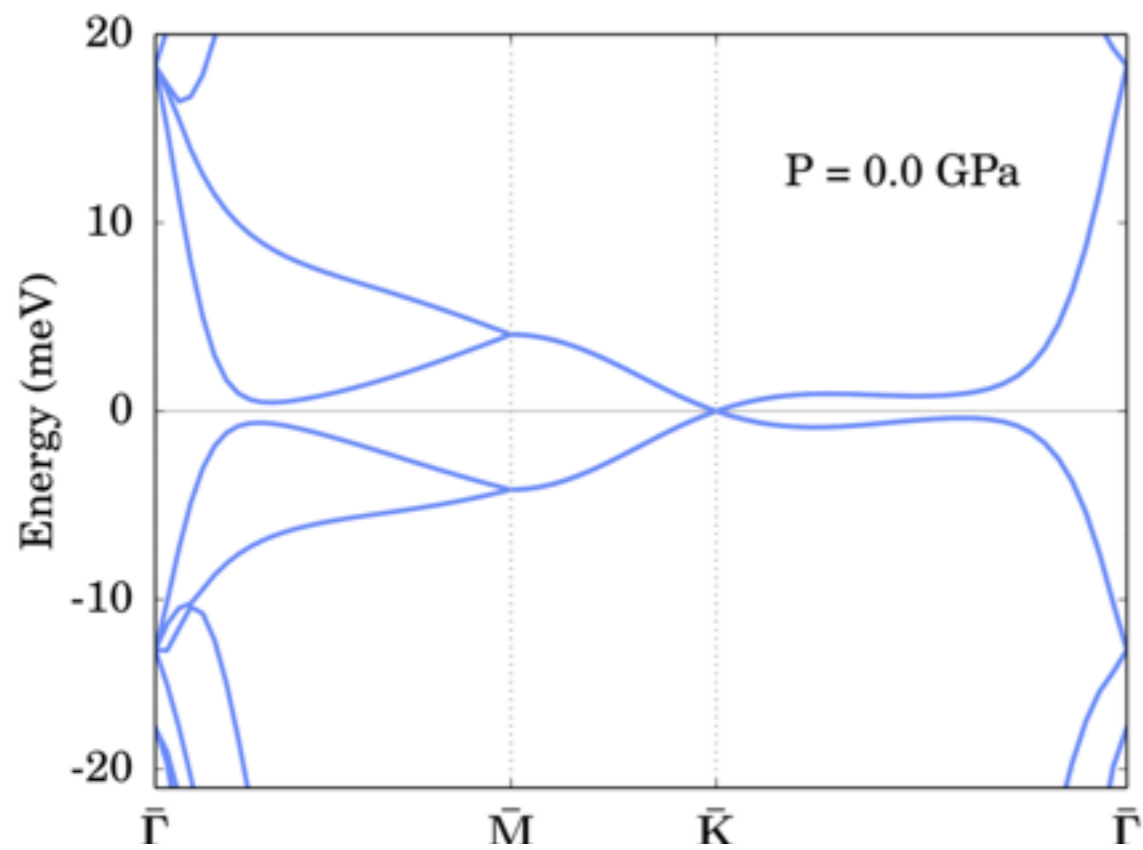
$$1 - \frac{d_{\perp}}{\tilde{d}} \equiv \delta_d = 10.48 \log \left(1 + \frac{P}{P_0} \right) \%$$

Interlayer Tunneling, w (meV)

Inter-layer
Tunneling

Effective
Magic-Ang





estimated

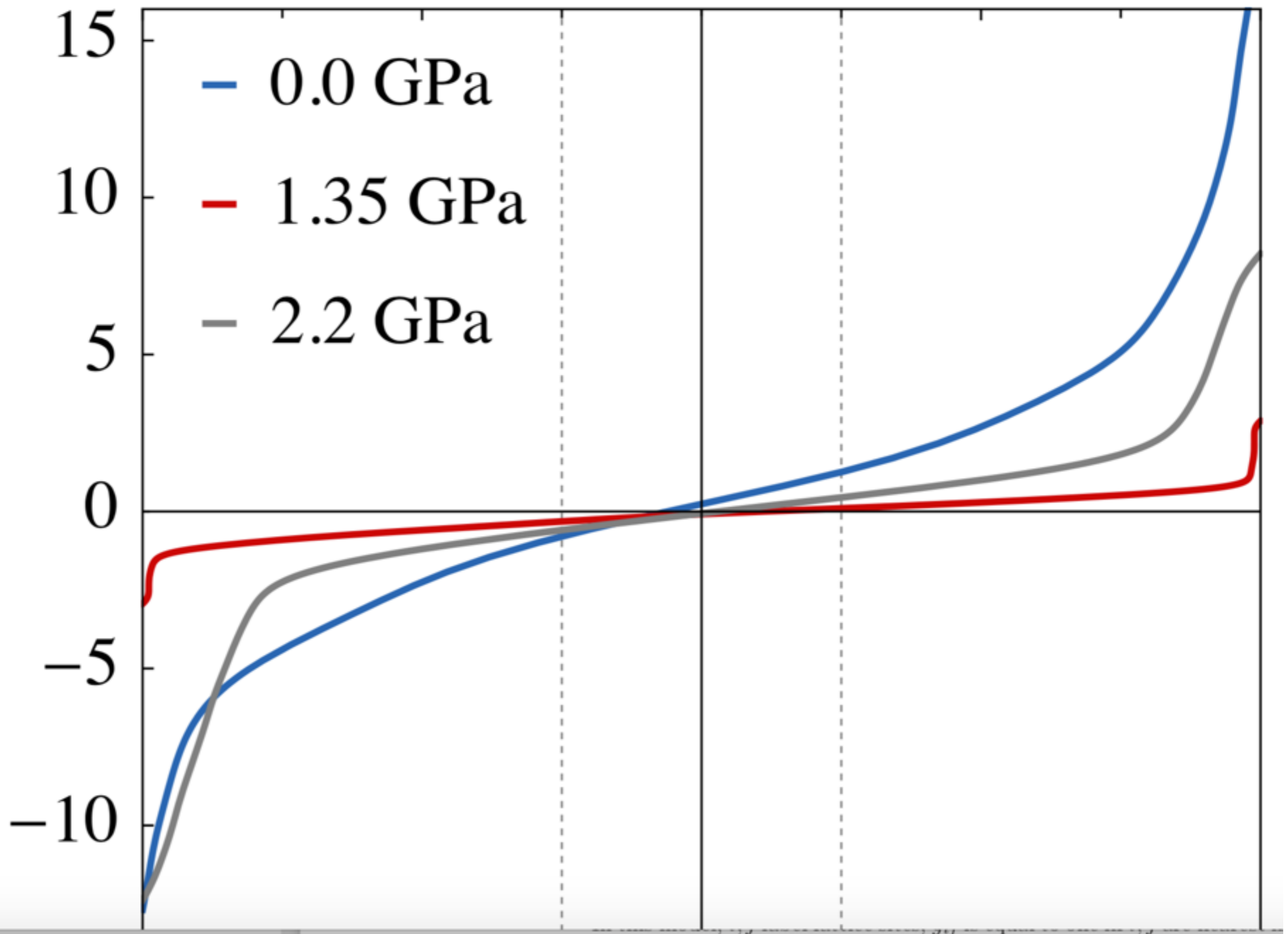
r_s

$$r_s \approx 15 \text{ meV} \frac{\theta^\circ}{E_K} \sqrt{\nu} \xrightarrow[\text{D2}]{\text{Device}} \frac{20 \text{ meV}}{E_K \text{ meV}} \sqrt{\nu}$$

Charge per supercell (e^-)

4 2 0 -2 -4

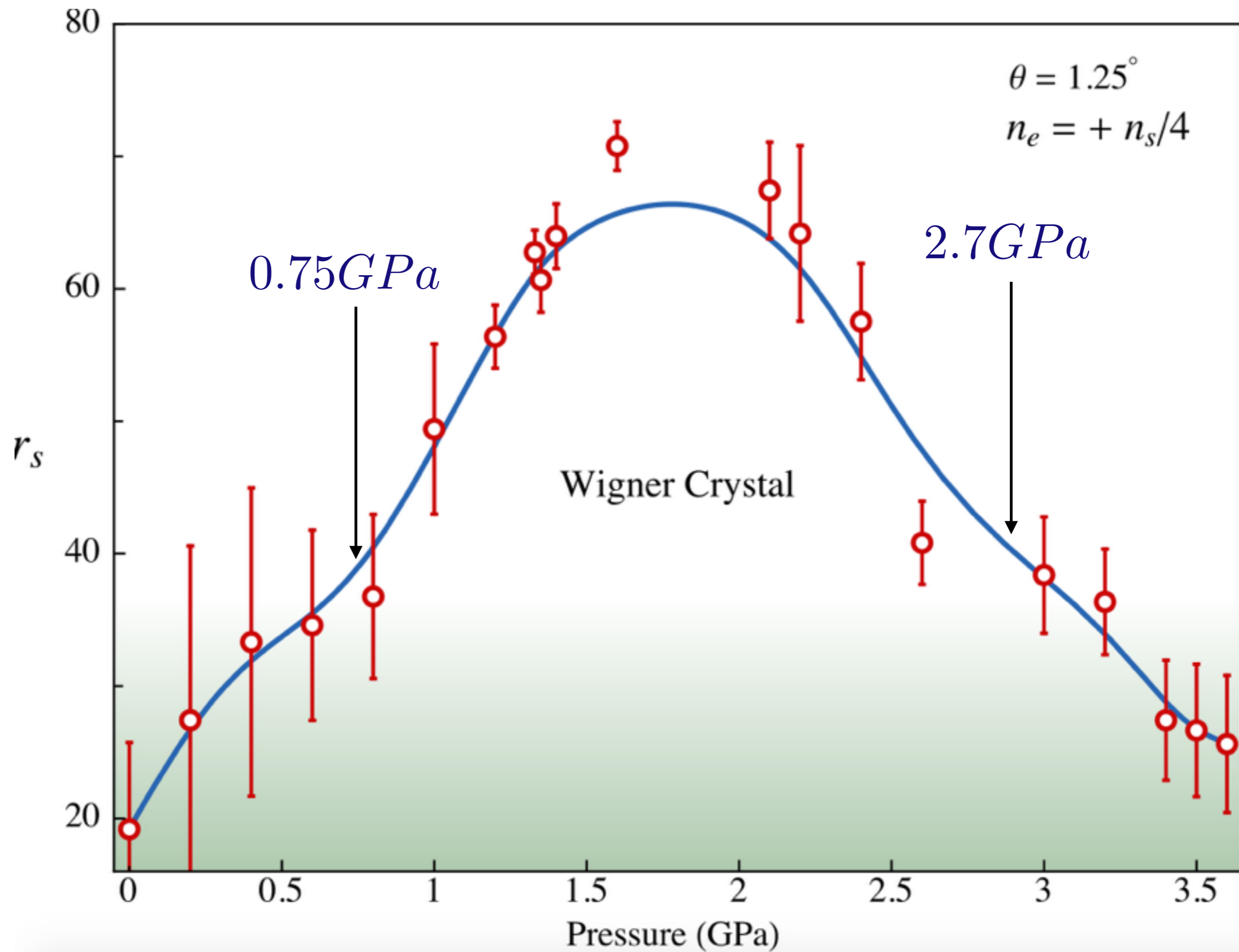
Chemical Potential, μ (meV)

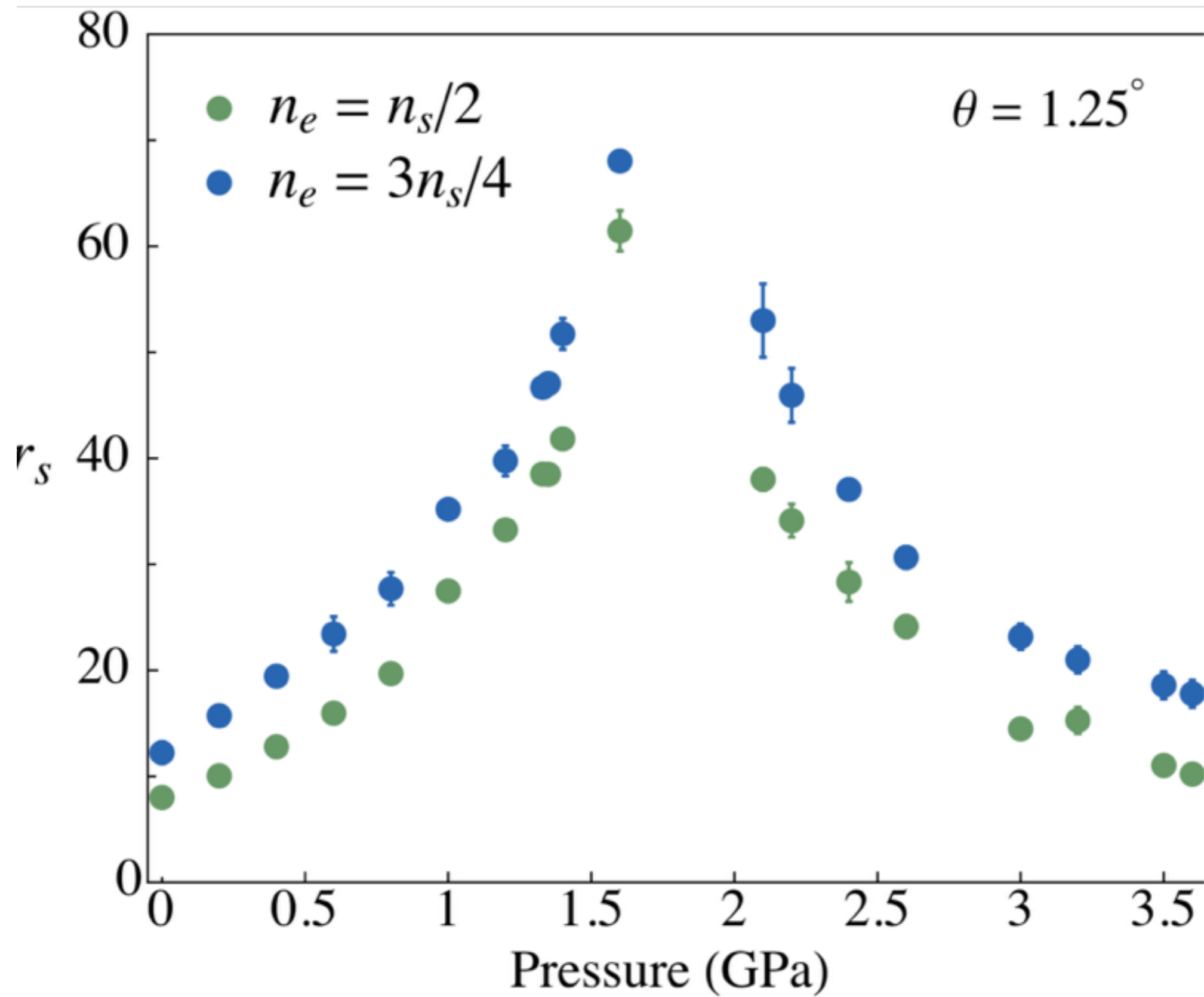


— 0.0 GPa

— 1.35 GPa

— 2.2 GPa





Prediction

$$0.75\text{GPa} < P < 2.75\text{GPa}$$

dome-shaped
phase diagram

Melting temperature

$$\frac{e^2}{4\pi\epsilon\lambda_s} \approx 30meV$$



$$T_{\text{melt}} \propto .01U_{\text{coul}}$$



$4K$

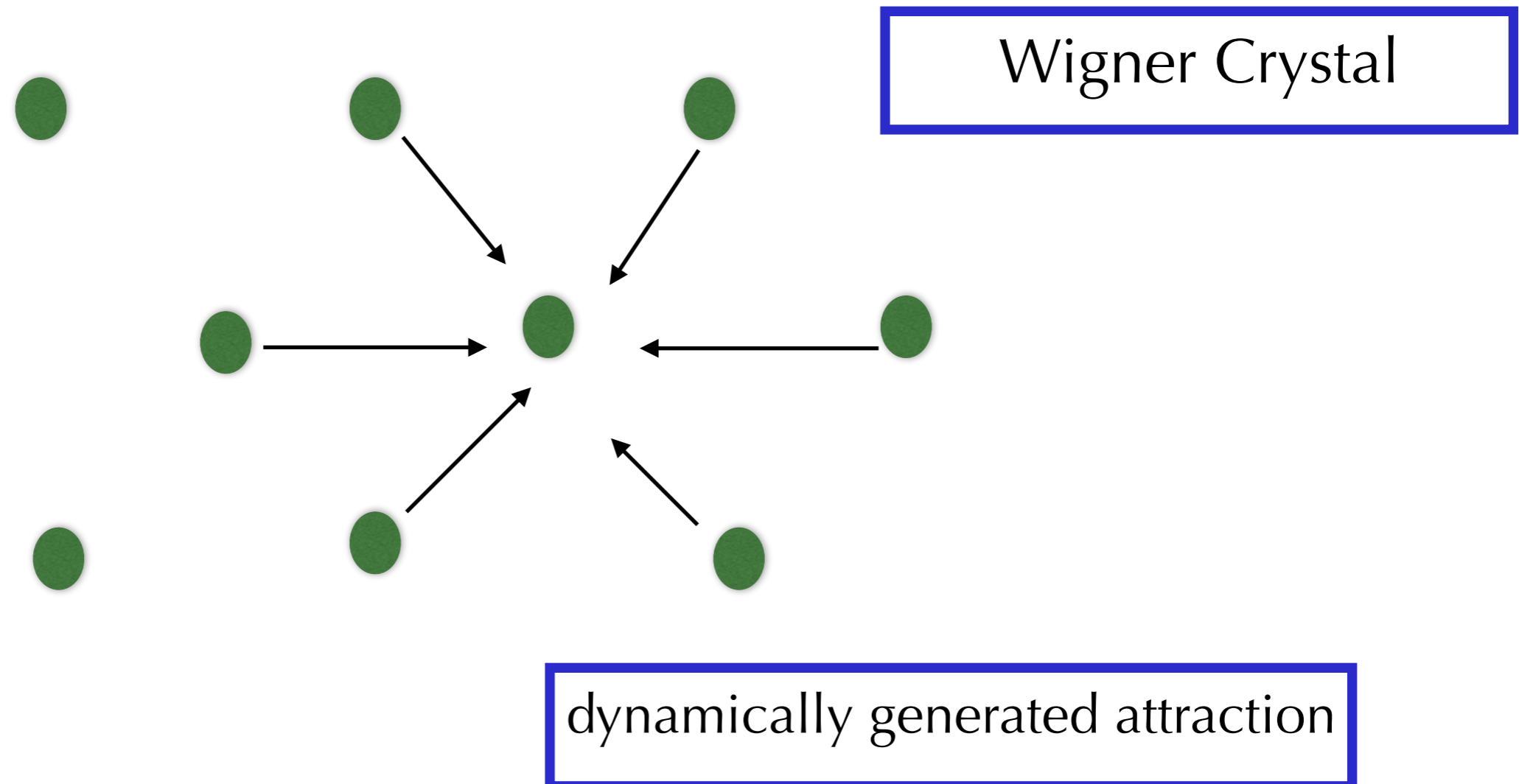
Letter

Superconductivity in a two-dimensional electron gas

Philip Phillips , Yi Wan, Ivar Martin, Sergey Knysh & Denis Dalidovich

Nature **395**, 253–257 (17 September 1998)

Received: 19 May 1998



Moiré is Different

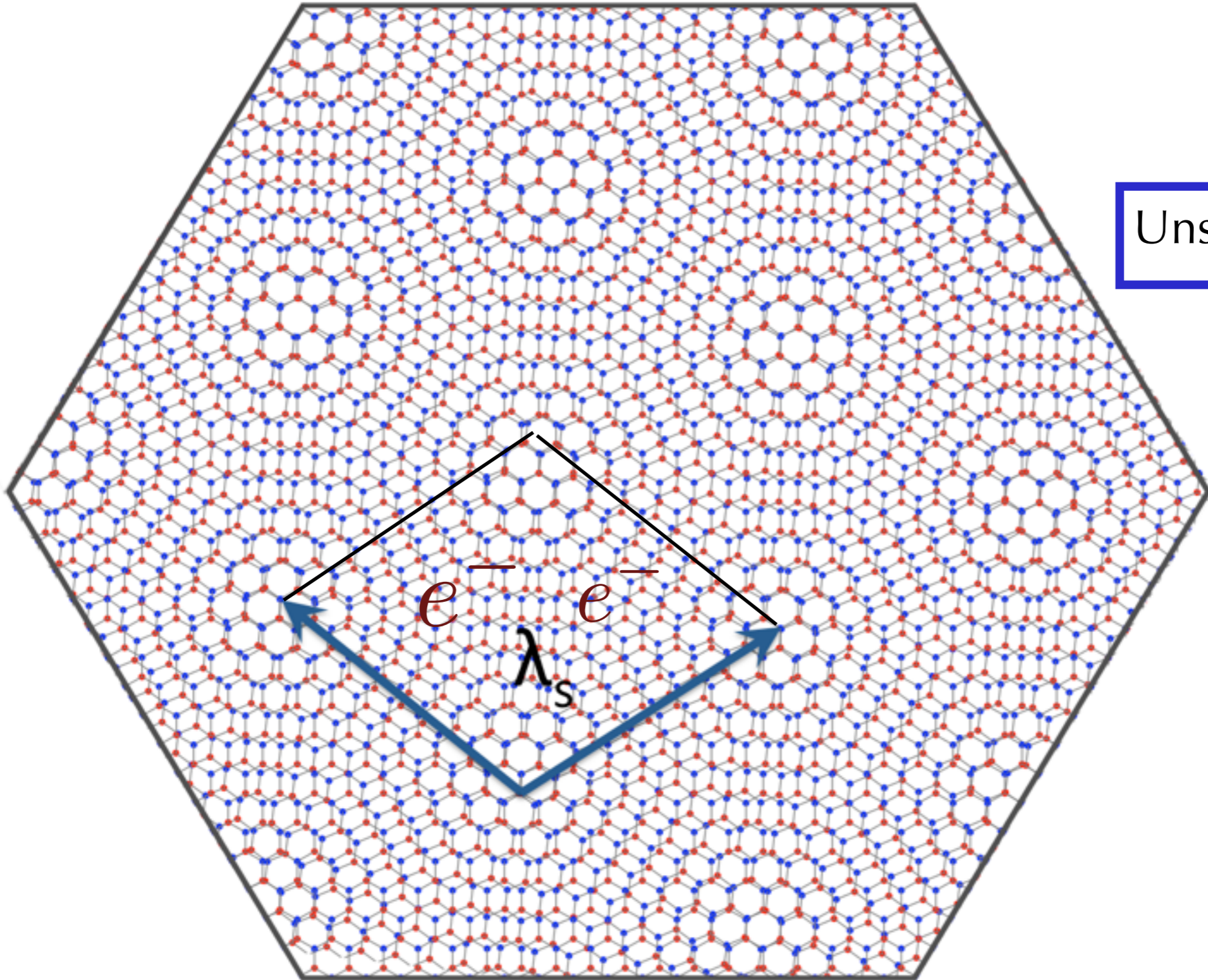
FL

WC

?



Unsolved Problem



$$r_s \gg 1$$