

BH Accretion, FB Efficiencies & Thermal Instability: a step to modeling AGN feedback (FB)

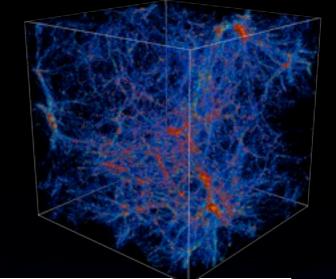
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Relevant papers:

Kurosawa, Proga, KN, 2009, ApJ, 707, 823 Barai, Proga, KN, 2011, MNRAS, 418, 591 (Paper I) Barai, Proga, KN, 2012, MNRAS, 424, 728 (Paper II) Moscibrodzka & Proga, 2013, ApJ, 767, 156 Kashi, Proga, KN, Barth, Greene 2013, arXiv:1306.1090

Outline

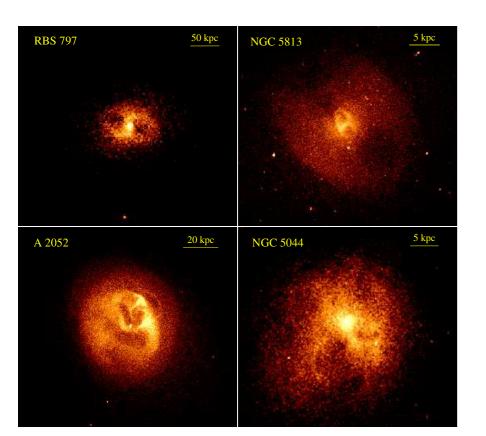


- Motivation: understand AGN feedback on smallscales and link to large-scales
- AGN feedback (FB) efficiencies
- Simulations: Bondi accretion (both SPH & Grid)
- Understanding Thermal Instability
- Implications for AGN feedback

Kurosawa, Proga, KN, 2009, ApJ, 707, 823 Barai, Proga, KN, 2011, MNRAS, 418, 591 (Paper I) Barai, Proga, KN, 2012, MNRAS, 424, 728 (Paper II) Moscibrodzka & Proga, 2013, ApJ, 767, 156 Kashi, Proga, KN, Barth, Greene 2013, arXiv:1306.1090

Many Signs of AGN Feedback (FB) in Galaxy Formation

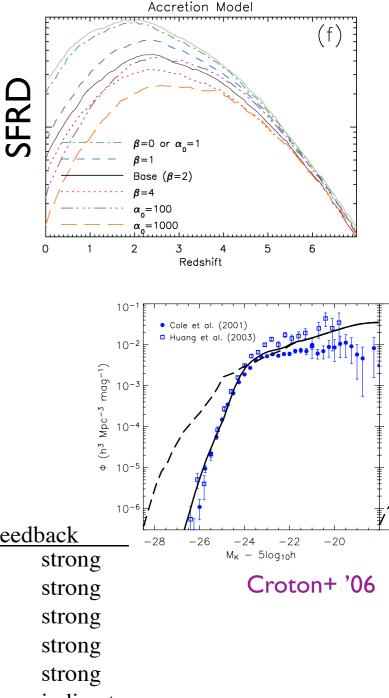
- AGN FB is considered to regulate star formation (quenching) in massive galaxies, causing the decline of cosmic SFRD at z≤1.
- Maybe it suppresses the massive-end of galaxy stellar mass func.
- Galaxy clusters show signs of central heating.



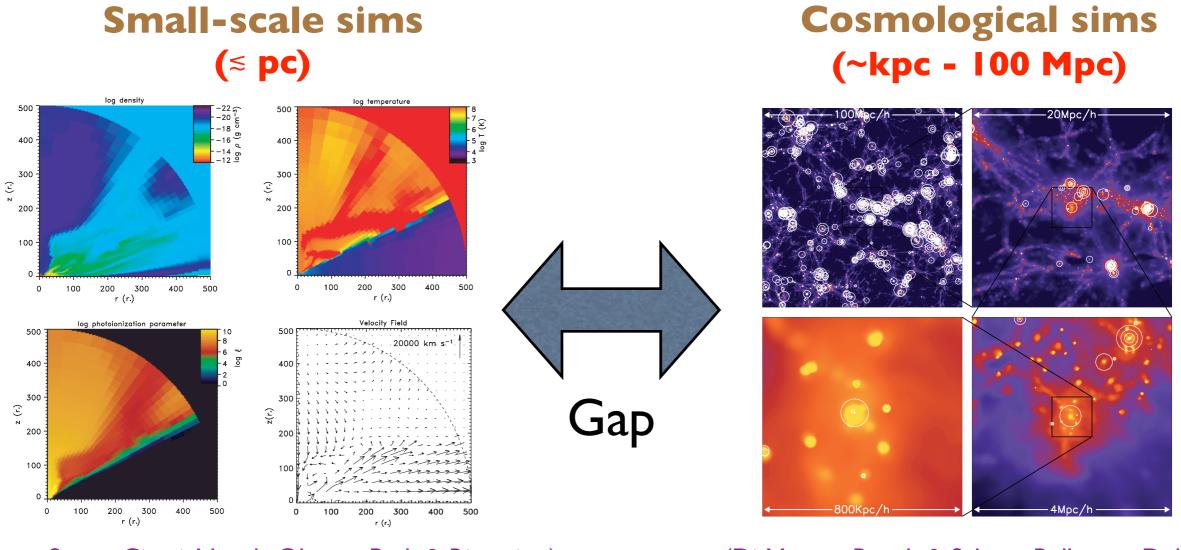
	10				
Table 1: Observational Evidence for AGN Feedback					
High velocity broad absorption lines in quasars	strong				
Strong winds in AGN	strong				
1000 km/s galactic outflows	strong				
Bubbles and ripples in BCGs	strong				
Giant radio galaxies	strong				
Lack of high SFR in cool cluster cores	indirect				
$M - \sigma$ relation	indirect				
Red and dead galaxies	indirect				
Lack of high lambda, moderate N_H , quasars	indirect				
Steep $L - T$ relation in low T clusters and groups	indirect				



Booth & Schaye '09



Bridging the Gap btw Large & Small Scale

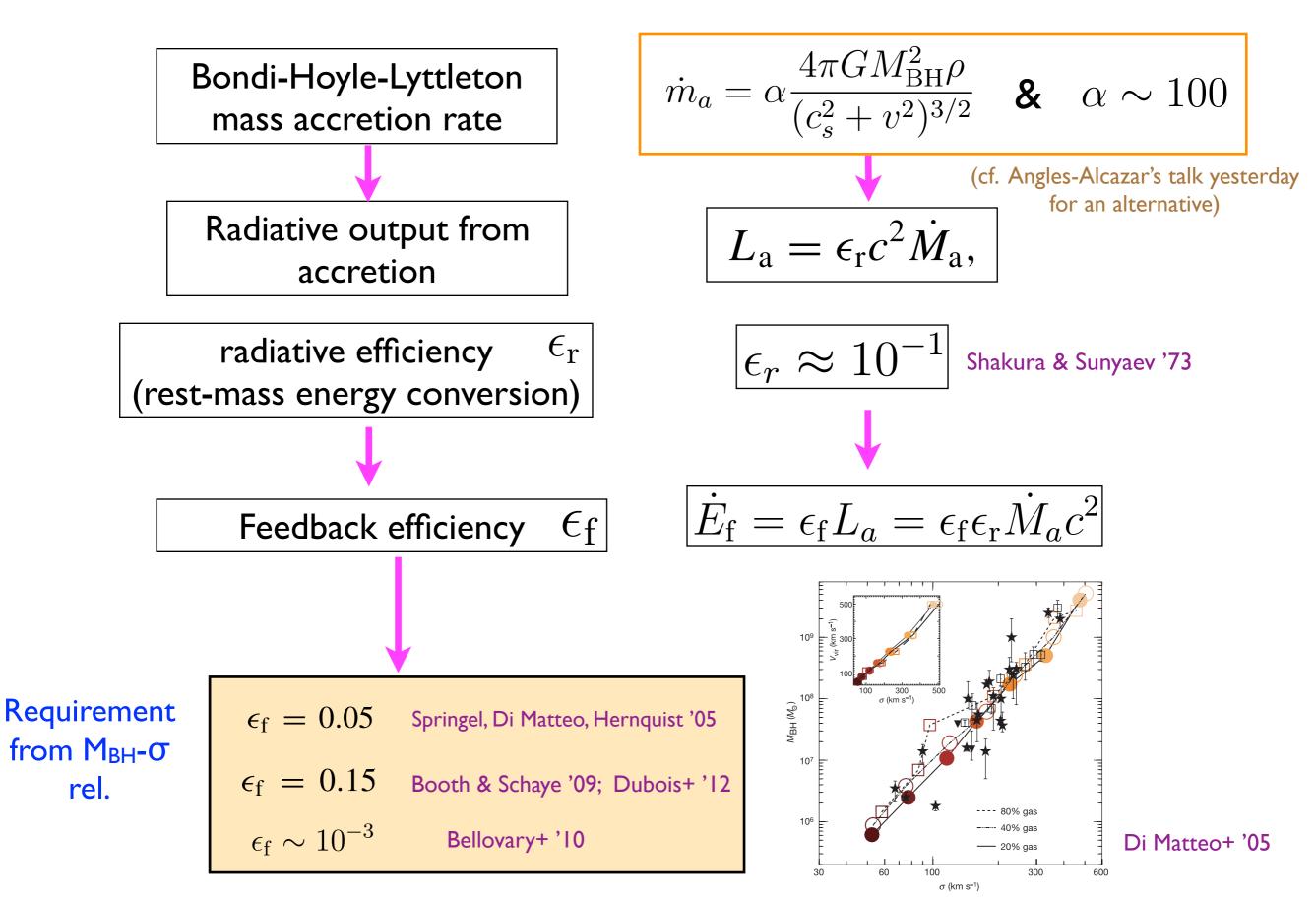


(Proga, Stone, Ciotti, Novak, Ohsuga, Park & Ricotti, ...)

(Di Matteo, Booth & Schaye, Bellovary, Dubois, ...)

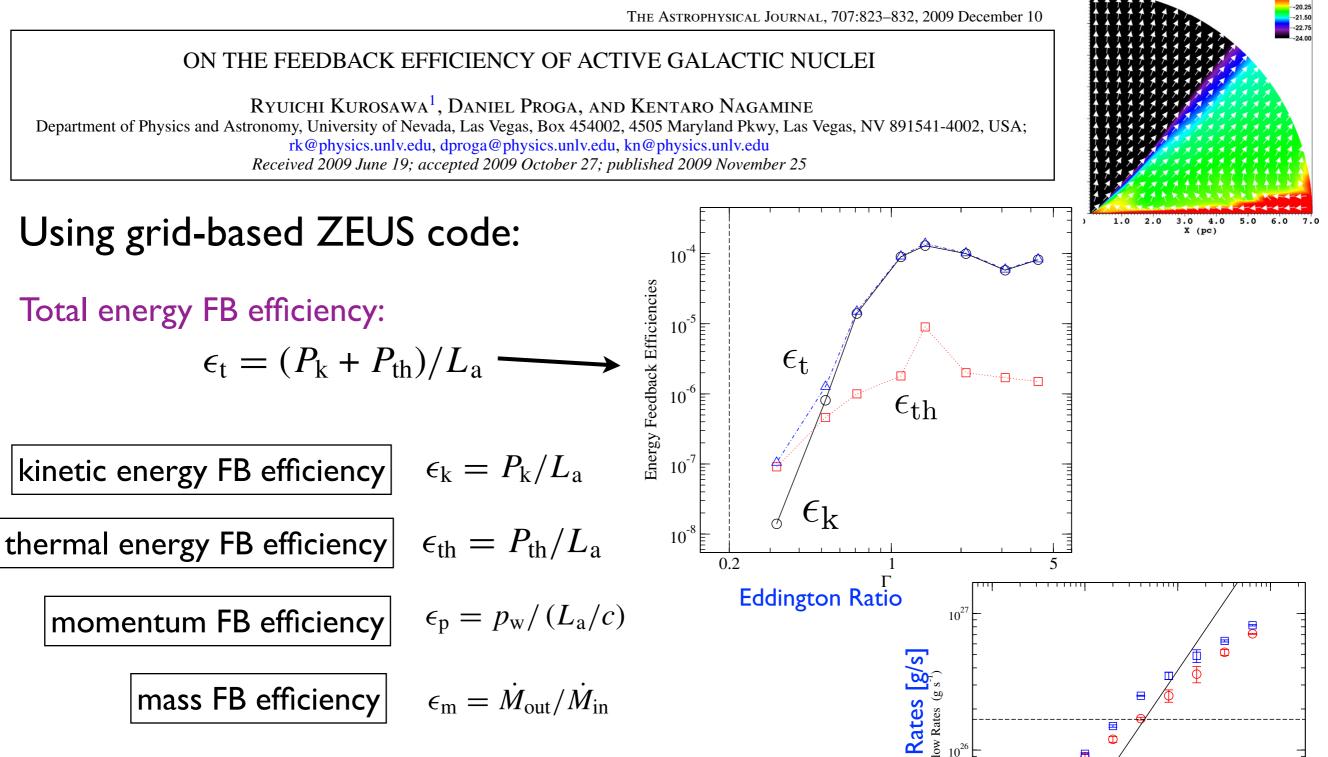
- Still a large gap btw small-scale sims & cosmological sims.
- Cosmo. sims use ad-hoc BH accretion models as "sub-grid" physics.

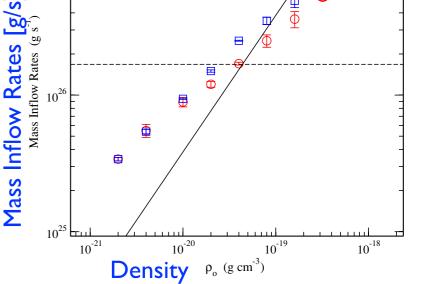
Assumptions in Cosmo. Sims.



FB Efficiency in small scale sim

THE ASTROPHYSICAL JOURNAL, 707:823–832, 2009 December 10





Taking Simple Steps

- First, check the code (Gadget SPH): the simplest case
 -- spherical Bondi accretion
- Include radiative cooling / heating -- only thermal feedback by X-rays
- Emergence of non-spherical accretion inflow/outflow, fragmentation due to thermal instability

Our goal is to test the simplest possible setup and see how the accretion flow reacts.

Barai, Proga, KN 2011, 2012

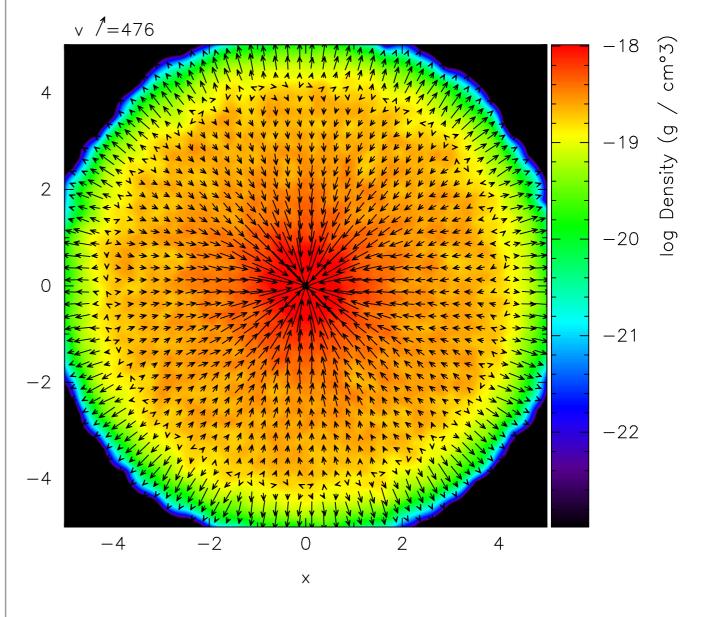
<u>Simplest Case: Spherical Bondi</u> <u>Accretion Flow onto a SMBH</u>

- GADGET-3 SPH (Springel '05)
- Central SMBH 10⁸ M_☉ represented by a pseudo-Newtonian Paczynsky & Wiita '80 potential
- r_{out}=5-20 pc, N_{ptcl}=64³-128³
- IC: uniform/spherical Bondi flow w/ γ =1.01, ρ_{∞} =10⁻¹⁹ g/ cm³, T_{\infty}=10⁷K, T_{init}=T_{\infty}
- $R_B=3pc, R_{sonic}=1.5pc, t_B=7.9e3yr$
- All runs: **r**_{in}=**0.1pc**, γ=1.01
- DM halo potential is minimal

			ρ∞				
			rin		r _{out}		
					T∞		
Run No.	$r_{ m out} \ [m pc]$	$N^{ m b}$	IC	$M_{ m tot,IC}$ ^c $[M_{\odot}]$	${M_{ m part}}^{ m d}$ $[M_{\odot}]$	${t_{ m end}}^{ m e} [10^4 \ { m yr}]$	
$egin{array}{c} 1 \\ 2 \\ 3 \end{array}$	$5 \\ 10 \\ 50$	${64}^3 \\ {64}^3 \\ {128}^3$	Uniform ⁱ Uniform Uniform	3.96×10^{5} 6.19×10^{6} 7.73×10^{8}	$1.51 \\ 23.61 \\ 368.60$	$3 \\ 7.2 \\ 20$	
4	5	64^{3}	Bondi ^j	1.81×10^{6}	6.89	2	
5 6	10 10	64^{3} 128^{3}	Bondi Bondi	9.76×10^{6} 9.76×10^{6}	$\begin{array}{c} 37.23\\ 4.65\end{array}$	8 8	
7	20	128^{3}	Bondi	6.24×10^7	29.75	8	
$7a^{ m \ k}$ $7b^{ m \ l}$	$\begin{array}{c} 20\\ 20 \end{array}$	128^{3} 128^{3}	Bondi Bondi	6.24×10^{7} 6.24×10^{7}	$29.75 \\ 29.75$	80 100	
8	50	128^{3}	Bondi	8.48×10^8	404.35	16	
9 10 11	20 20 20	128^3 128^3 128^3	$ ho_B, v_{ m init} = 0$ Uniform Hernquist ^m	6.24×10^{7} 4.95×10^{7} 6.24×10^{7}	29.75 23.60 29.75	$8\\8\\7.2$	
12 ⁿ	20	128^{3}	Bondi	6.24×10^7	29.75	8	

Spherical Bondi Accretion Flow

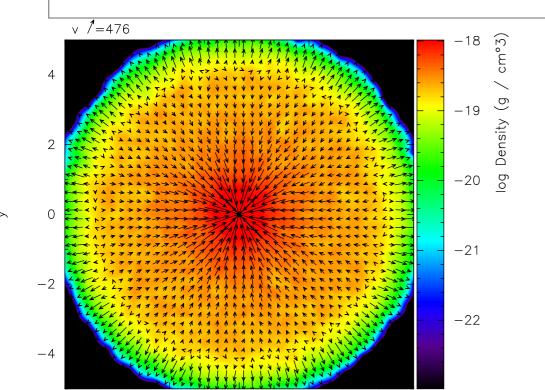
- Cross section, showing ρ_{gas}
- Run #4: $r_{out}=5 \text{ pc}$, $N_{ptcl}=64^3$
- At $t = 0.25t_B = 2e3$ yr
- Smooth, spherical Bondi accretion > flow is reproduced.
- Artificial outflow near the outer boundary due to vacuum outer B.C.
- Greater sim. volume reduces this effect on mass inflow.

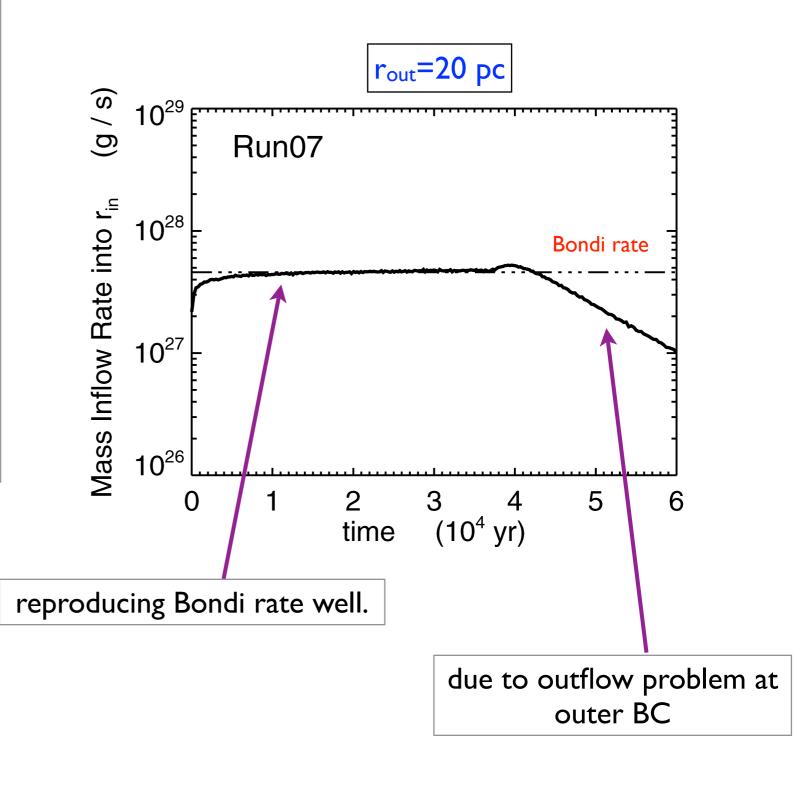


Barai, Proga, KN 'I I

Mass Inflow Rates at rin

- the larger r_{out}, the longer duration of Bondi inflow rate
- If started from a Bondi flow, Bondi rate is achieved quickly.
- After a while, the inflow rate decreases due to the artificial outflow at the outer boundary.
- Greater sim. volume reduces this effect on mass inflow.





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Radiative Heating & Cooling

• Xray emitting corona irradiates the accretion flow

$$L_X = f_X L_{Edd}, \quad L_{Edd} = \frac{4\pi c G m_p M_{BH}}{\sigma_e}, \qquad Flux: F_X = \frac{L_X}{4\pi r^2}.$$

(f_x=0.05 in KP09)

 Approx. analytic heating/cooling rates from Blondin '94; opt-thin gas illuminated by a 10 keV bremsstrahlung.

net rate:

$$\rho \mathcal{L} = n^2 \left(\underline{G_{Compton}} + \underline{G_X} - \underline{L}_{b,l} \right)$$
 [erg cm⁻³ s⁻¹],

 Compton h/c rate:
 $G_{Compton} = 8.9 \times 10^{-36} \xi (T_X - 4T)$ [erg cm³ s⁻¹].

 Net Xray photoioniz, heating and recomb, cooling rate:
 $G_X = 1.5 \times 10^{-21} \xi^{1/4} T^{-1/2} \left(1 - \frac{T}{T_X}\right)$ [erg cm³ s⁻¹].

 Brems, and line cooling rate:
 $L_{b,l} = 3.3 \times 10^{-27} T^{1/2} + \left[1.7 \times 10^{-18} \exp\left(-1.3 \times 10^5/T\right) \xi^{-1} T^{-1/2} + 10^{-24}\right] \delta$ [erg cm³ s⁻¹].

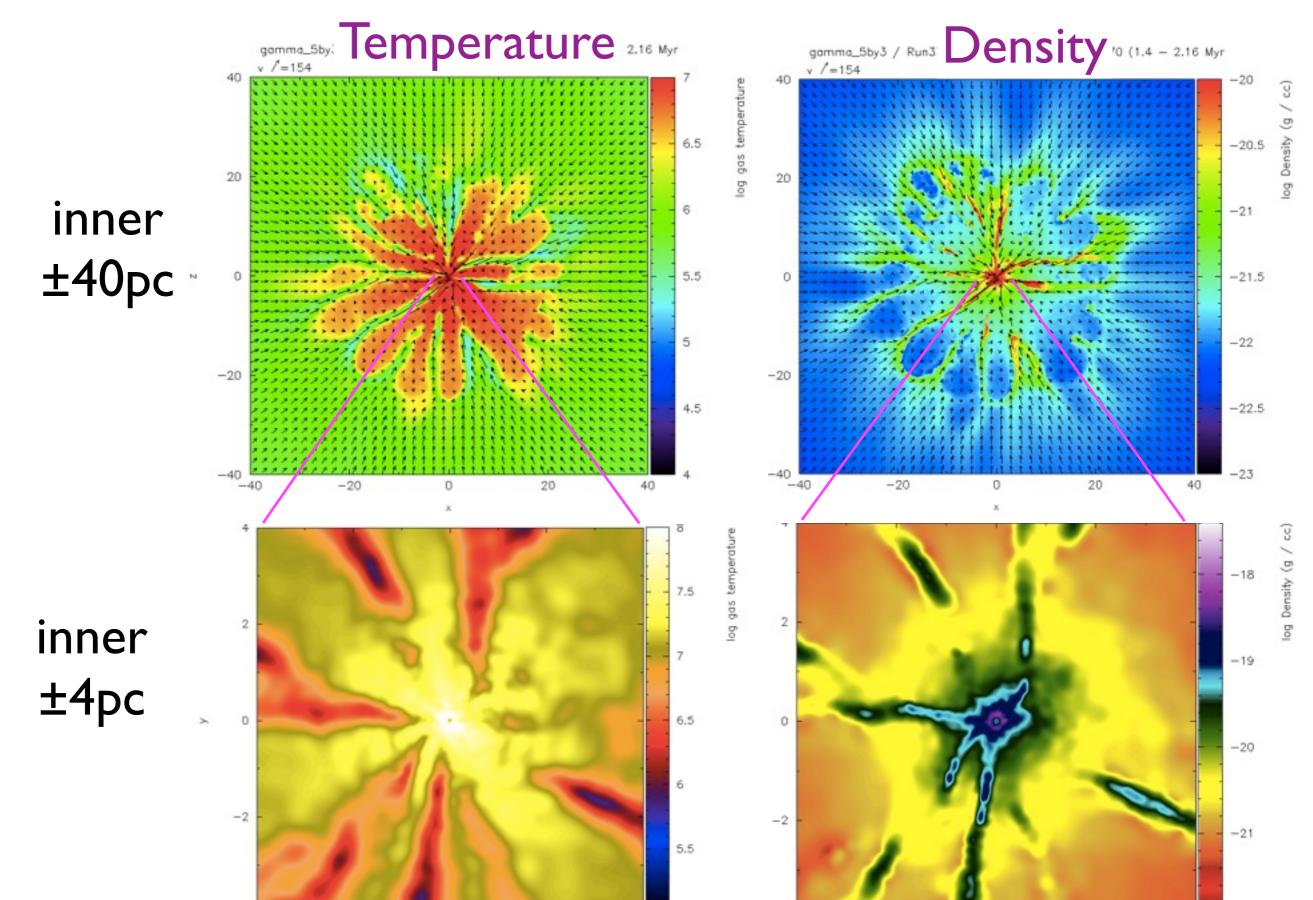
 $T_X=1.16 \times 10^8 \text{ K}$ (=10keV, Blondin '94)

Runs with radiative cooling/heating

Run No.	$r_{ m out} \ [m pc]$	Ν	$M_{ m tot,IC}$ $[M_{\odot}]$	$M_{ m part}$ $[M_{\odot}]$	$\gamma_{ m init}$	T_{∞} [K]	R_B [pc]	$ ho_{\infty}$ [g/cm ³]	T_{init}	L_X $[L_{ m Edd}]$	$t_{ m end}$ $[10^5 m yr]$
13	20	128^{3}	5.81×10^5	0.277	1.4	10^{7}	2.19	10^{-21}	T_{∞}	0.5	1.0
14	50	128^{3}	8.23×10^6	3.92	1.4	10^{7}	2.19	10^{-21}	T_{∞}	0.5	2.9
$\frac{15}{16}$	20 20	128^{3} 256^{3}	5.81×10^{-1} 5.81×10^{-1}	2.77×10^{-7} 3.46×10^{-8}	$1.4\\1.4$	$\frac{10^7}{10^7}$	$\begin{array}{c} 2.19 \\ 2.19 \end{array}$	10^{-27} 10^{-27}	T_{∞} T_{∞}	$\begin{array}{c} 0.5\\ 5\times10^{-4}\end{array}$	$\begin{array}{c} 1.0 \\ 1.9 \end{array}$
$\frac{17}{18}$	$\begin{array}{c} 20\\ 20\end{array}$	$\frac{128^3}{128^3}$	$5.81 imes 10^5 \ 5.65 imes 10^5$	$0.277 \\ 0.269$	$1.4 \\ 5/3$	$\frac{10^7}{10^7}$	$\begin{array}{c} 2.19 \\ 1.84 \end{array}$	10^{-21} 10^{-21}	${T_{ m rad}}^{ m b} T_{ m rad}$	5×10^{-4} 5×10^{-4}	$\begin{array}{c} 2.9\\ 3.0 \end{array}$
19	20	128^{3}	1.47×10^{7}	7.0	5/3	10^{5}	183.9	10^{-21}	$T_{\rm rad}$	5×10^{-4}	1.5
$20\\21$	$\begin{array}{c} 200\\ 200 \end{array}$	256^{3} 256^{3}	1.33×10^9 4.95×10^8	$79.09 \\ 29.50$	$5/3 \\ 5/3$	$\frac{10^5}{10^7}$	$\begin{array}{c} 183.9\\ 1.84 \end{array}$	10^{-21} 10^{-21}	$T_{ m rad} \ T_{ m rad}$	5×10^{-4} 5×10^{-4}	$\begin{array}{c} 6.5 \\ 8.7 \end{array}$
22	200	128 ³	1.33×10^{7}	6.33	5/3	10^{5}	183.9	10^{-23} 10^{-23}	$T_{\rm rad}$	5×10^{-4}	70
23	200	256^{3}	1.33×10^7	0.791	5/3	10^{5}	183.9	10 -0	$T_{\rm rad}$	5×10^{-4}	20
24 ^c	200	1.24×10^{7}	9.77×10^{6}	0.791	5/3	10^{5}	183.9	10^{-23}	$T_{ m Run23}$	5×10^{-5}	19
25	200	1.24×10^7	9.77×10^{6}	0.791	5/3	10^{5}	183.9	10^{-23}	T_{Run23}	5×10^{-3}	21
26	200	1.24×10^{7}	9.77×10^{6}	0.791	5/3	10^{5}	183.9	10^{-23} 10^{-23}	$T_{ m Run23}$	1×10^{-2}	22
27 28	$200 \\ 200$	$\begin{array}{c} 1.24 \times 10^7 \\ 1.24 \times 10^7 \end{array}$	9.77×10^{6} 9.77×10^{6}	$0.791 \\ 0.791$	$5/3 \\ 5/3$	10^{5} 10^{5}	$183.9\\183.9$	10^{-23} 10^{-23}	$T_{ m Run23}$ $T_{ m Run23}$	2×10^{-2} 5×10^{-2}	$\begin{array}{c} 25 \\ 50 \end{array}$

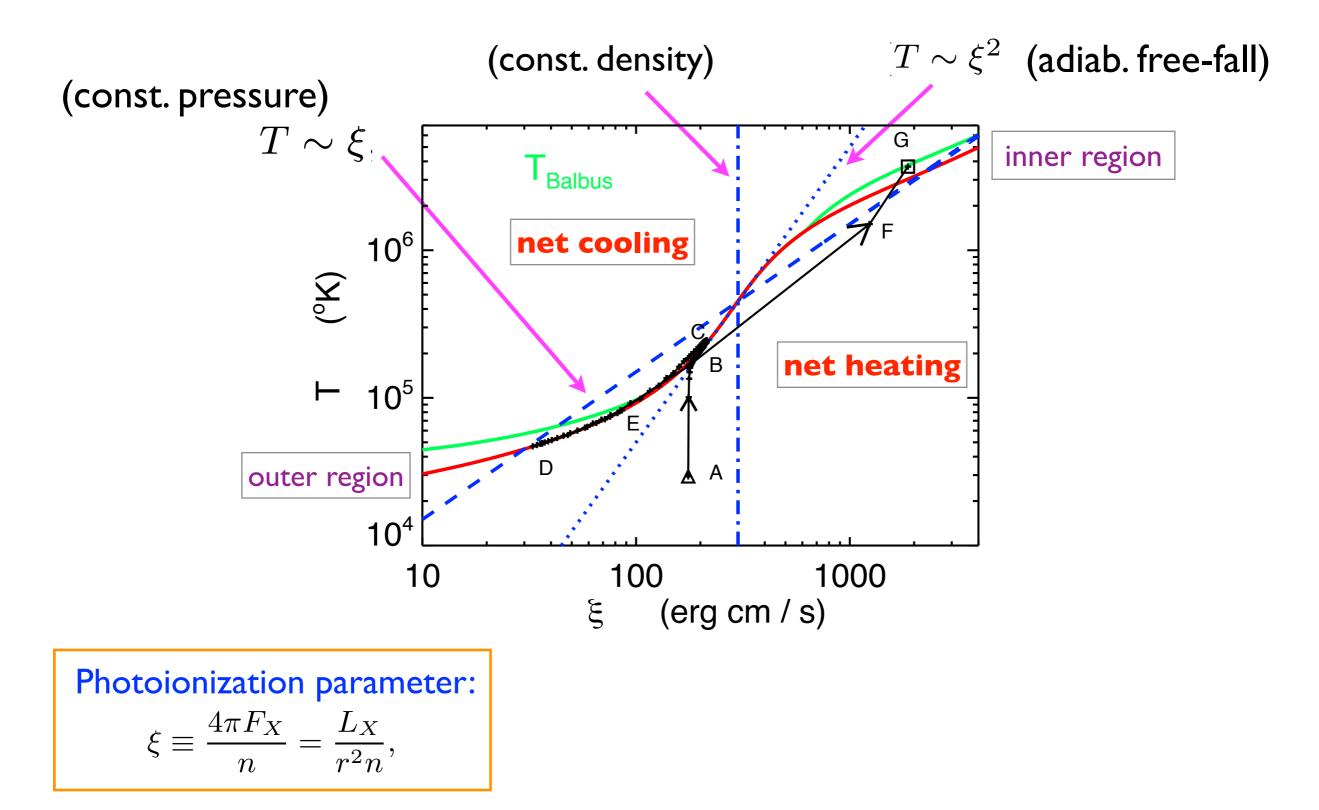
Non-spherical outflow: Run 26: rout=200pc, Lx/L_{Edd}=0.01

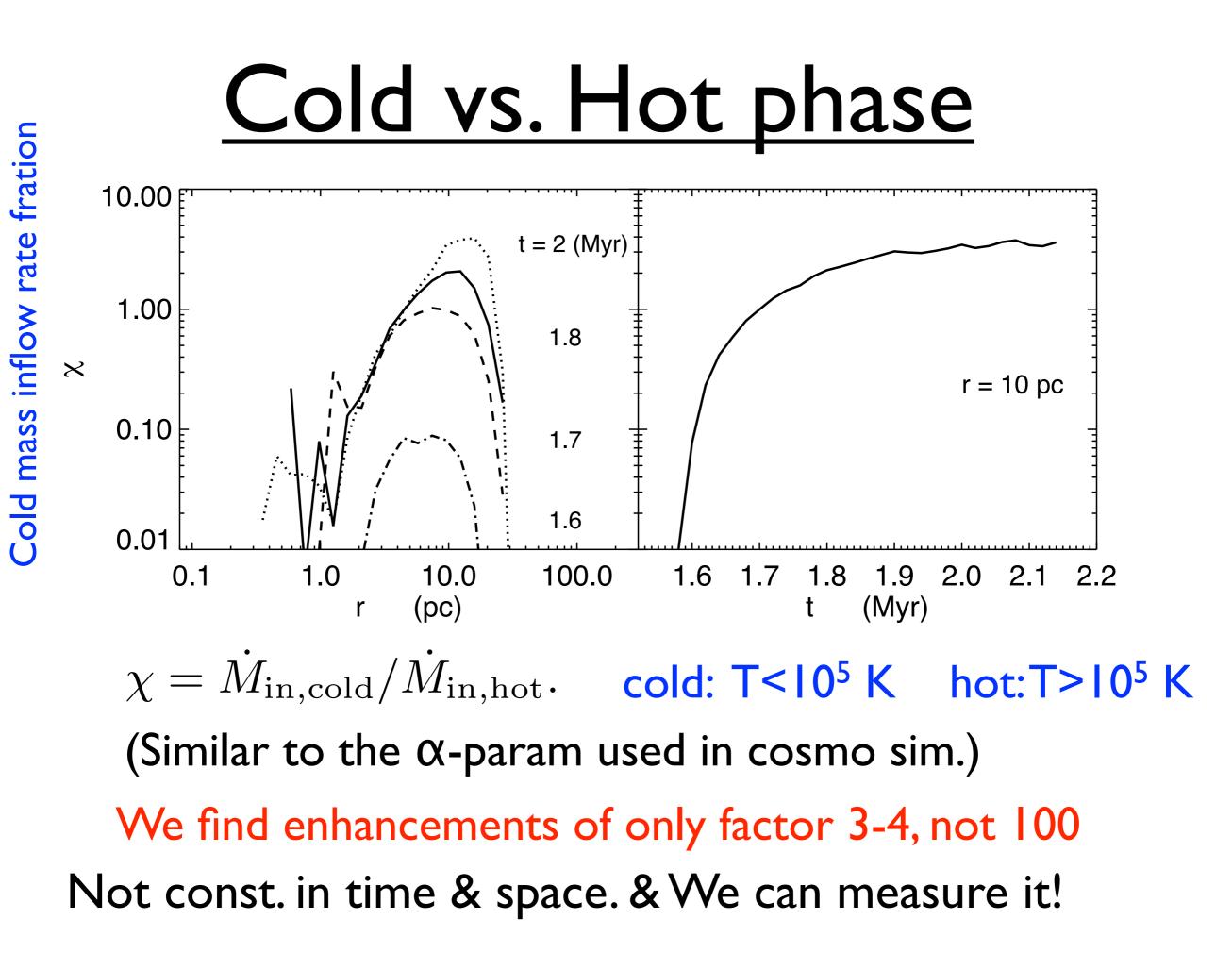
due to rad. feedback

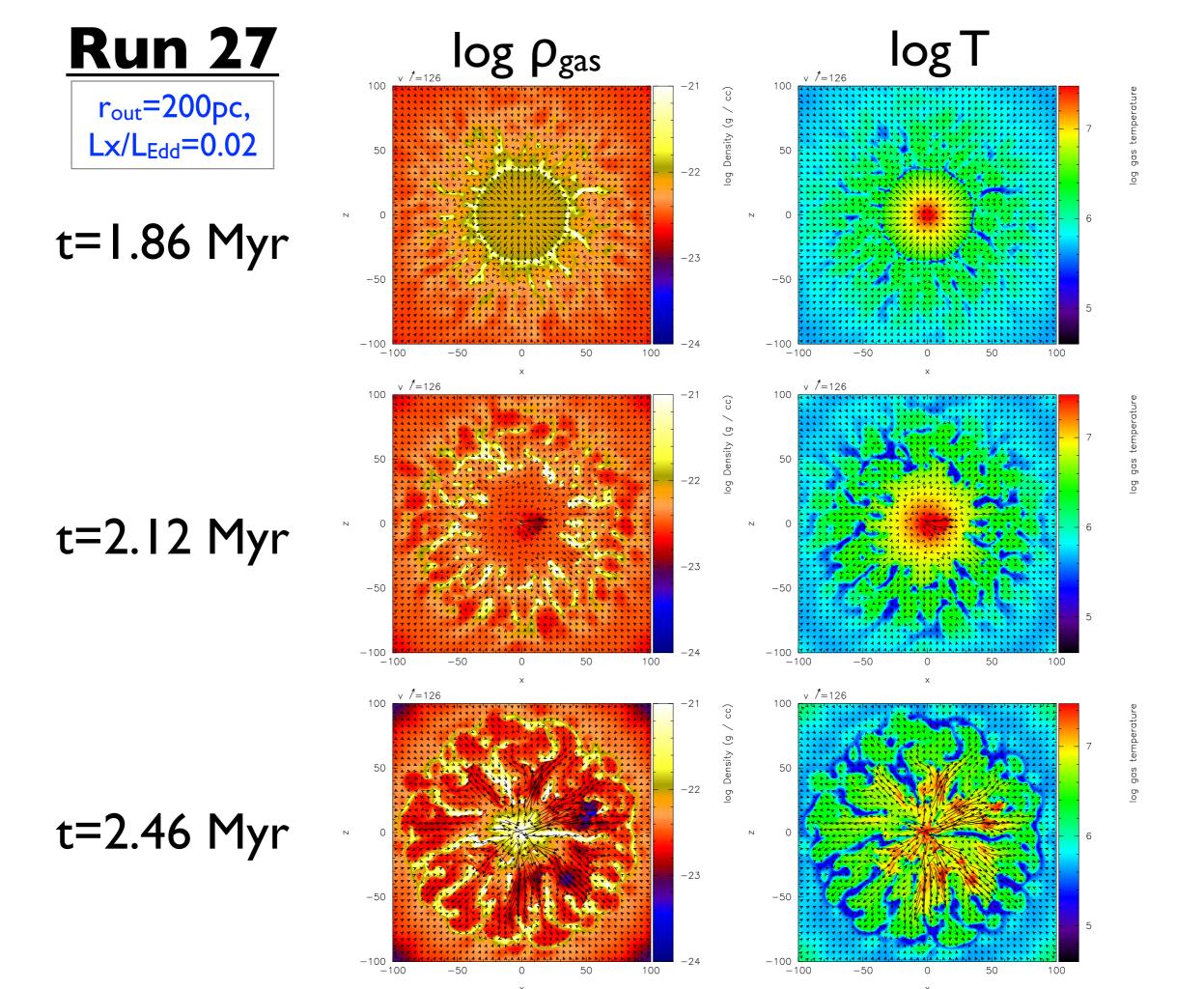


Indication of Thermal Instability

Run 26: r_{out}=200pc, Lx/L_{Edd}=0.01, t=2.0 Myr

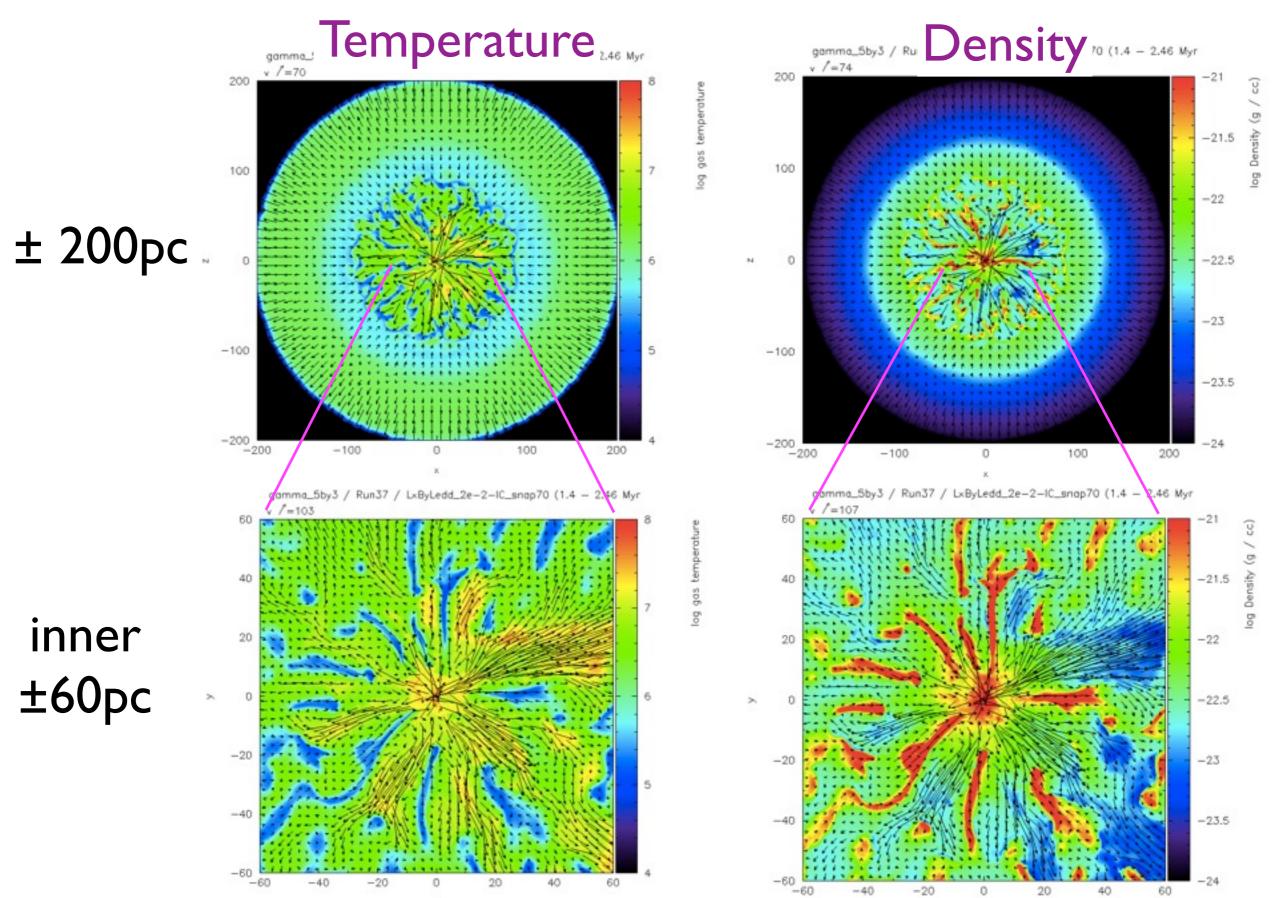






Non-spherical outflow: Run 27: rout=200pc, Lx/L_{Edd}=0.02

due to rad. feedback



Independent check of thermal instability with grid-based ZEUS code in 1,2,3-dimension

The Astrophysical Journal, 767:156 (15pp), 2013 April 20

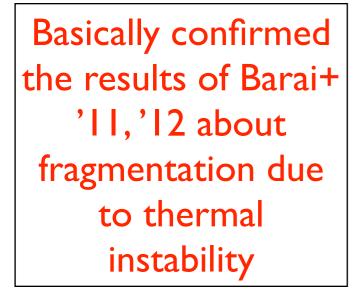
THERMAL AND DYNAMICAL PROPERTIES OF GAS ACCRETING ONTO A SUPERMASSIVE BLACK HOLE IN AN ACTIVE GALACTIC NUCLEUS

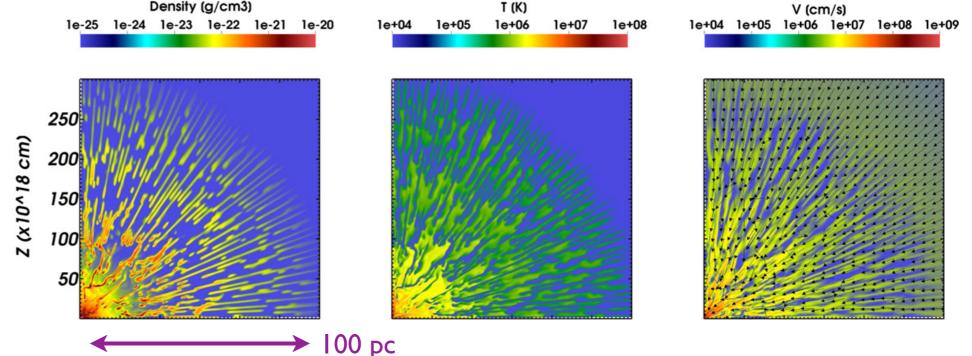
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$$M_{BH} = 10^8 M_{\odot}$$

$$L_X = 0.015 L_{\text{Edd}}$$
 -- right inbtw 1% and 2% L_{Edd}.

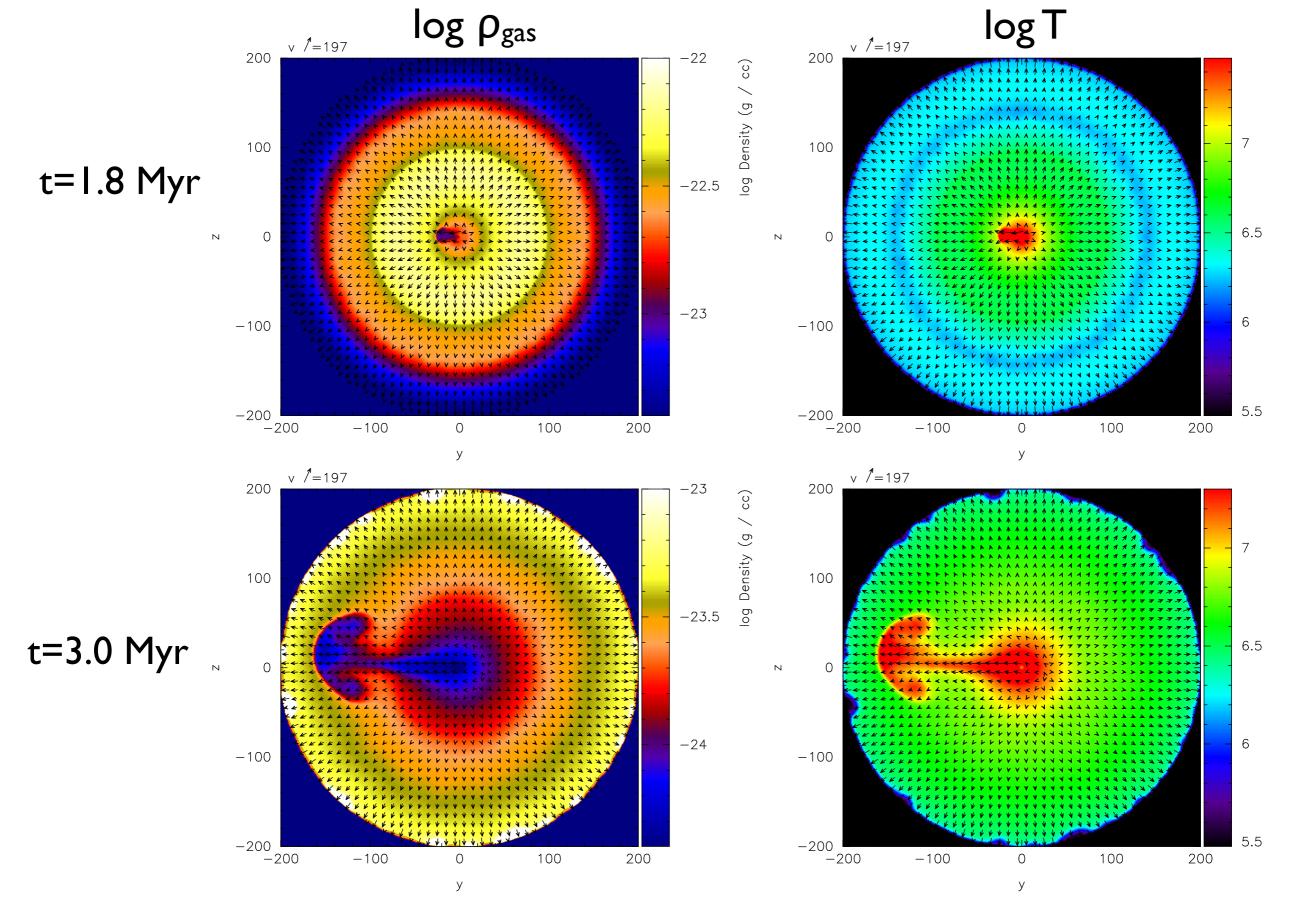




ZEUS does not suffer from artificial outflow problem at the outer boundary, and we can run it longer. Proga will show long-term evolution of the flow.

<u>Run 28</u>

r_{out}=200pc, Lx/L_{Edd}=0.05



Conclusions & Future Work

- GADGET SPH code *can* reproduce the *spherical Bondi accretion* properly w. some limitations. (repeat w. new SPH)
- Non-spherical inflow & outflow develops due to rad. feedback via thermal instability, even in the simplest case that we studied ---connection with NLR
- AGN FB efficiencies measured in our sims are LOW (≤10⁻⁴): cosmological implications? -- Is AGN FB really important for galaxy formation? (cf. Priya's talk, Ostriker+10)
- Future: include other physics; how do we apply these results to cosmo. sims? Maybe we should stop using Bondi formula in cosmo sims.









