

# Dynamics of living liquid crystals: Bacteria in Lyotropic Chromonic Liquid Crystals

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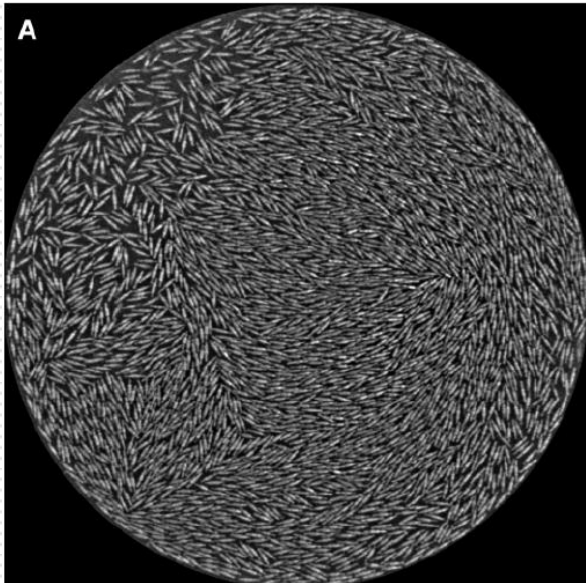


Support: NSF DMR

Active Processes in Living and Nonliving Matter, Feb 10-14, 2014

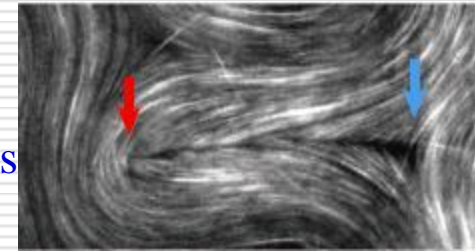
# Examples of active fluids

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Narayan, Ramaswamy, Menon, *Science* (2007): Active granular rods with disclinations and giant number fluctuations

Sanchez, Dogic et al, *Nature* (2012): Active microtubules with flows and disclinations



Review: Marchetti et al *RMP* (2013)

This work: Living Liquid Crystal (LLC), dispersion of rod-like self-propelling bacteria in lyotropic (water-based) chromonic liquid crystal

Motivation: The system allows one to control orientational order and activity separately

- Orientational order modifies bacterial behavior
- Bacterial activity modifies orientational order

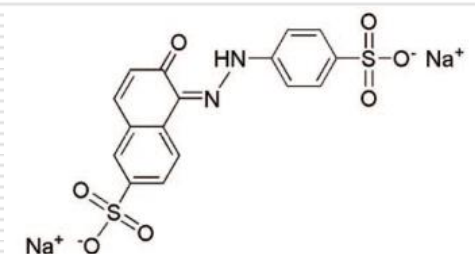
# Living Liquid Crystals: Content

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- Chromonics (inactive): Basic properties
  - Phase diagram
  - Viscoelasticity
- Living liquid crystals = *Chromonics* + *B. Subtilis*
  - Behavior of individual bacteria in a chromonic
  - Collective effects: activity-controlled transition from equilibrium uniformity to topological turbulence

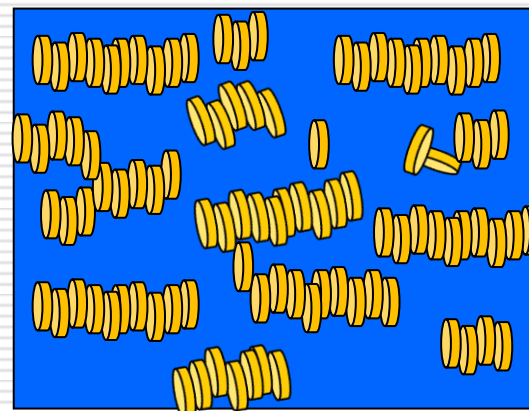
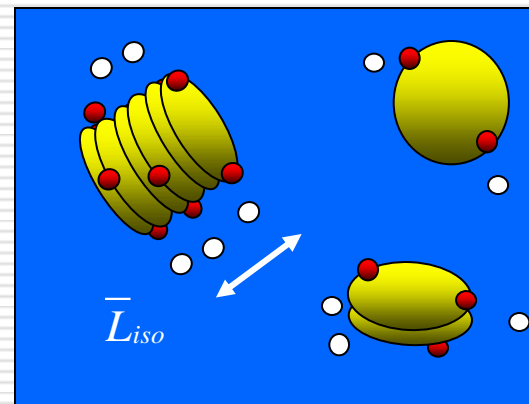
# Chromonics

- Chromonics are formed by rigid polyaromatic molecules that assemble face-to-face into elongated aggregates in water



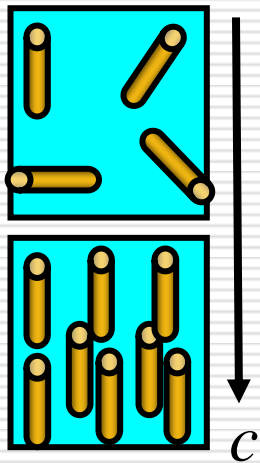
- As the concentration increases, the aggregates form a nematic
- Similar to the Onsager mechanism, but not exactly, since the system is also “thermal”: The aggregates are kept together by weak non-covalent interactions  $E \sim 10 k_B T$ , which makes their length strongly temperature dependent and concentration dependent:

$$\bar{L}_{iso} \approx L_0 \sqrt{\phi} \exp\left(\frac{E}{2k_B T}\right)$$

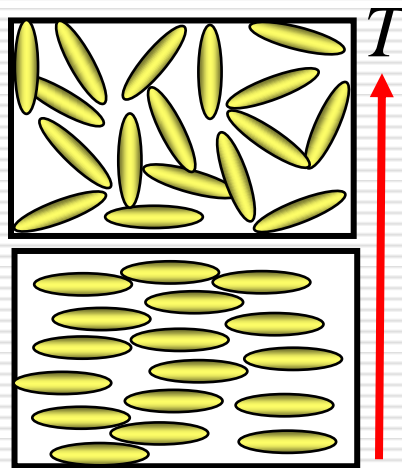


# Inactive chromonics as “living polymers”

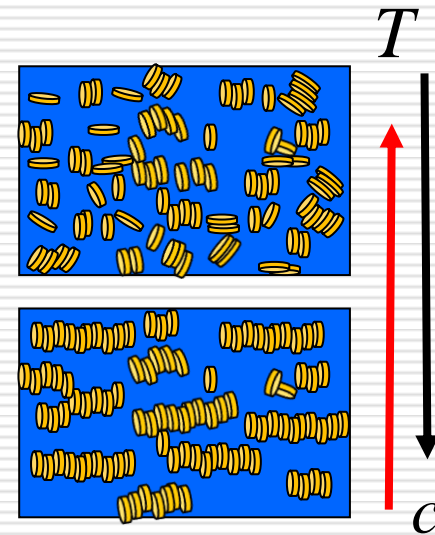
Onsager systems (TMV):  
The same rods, athermal, aligned at high  $c$



Thermotropic LC (in your displays):  
The same molecules,  $c = \text{const}$ , aligned at low  $T$

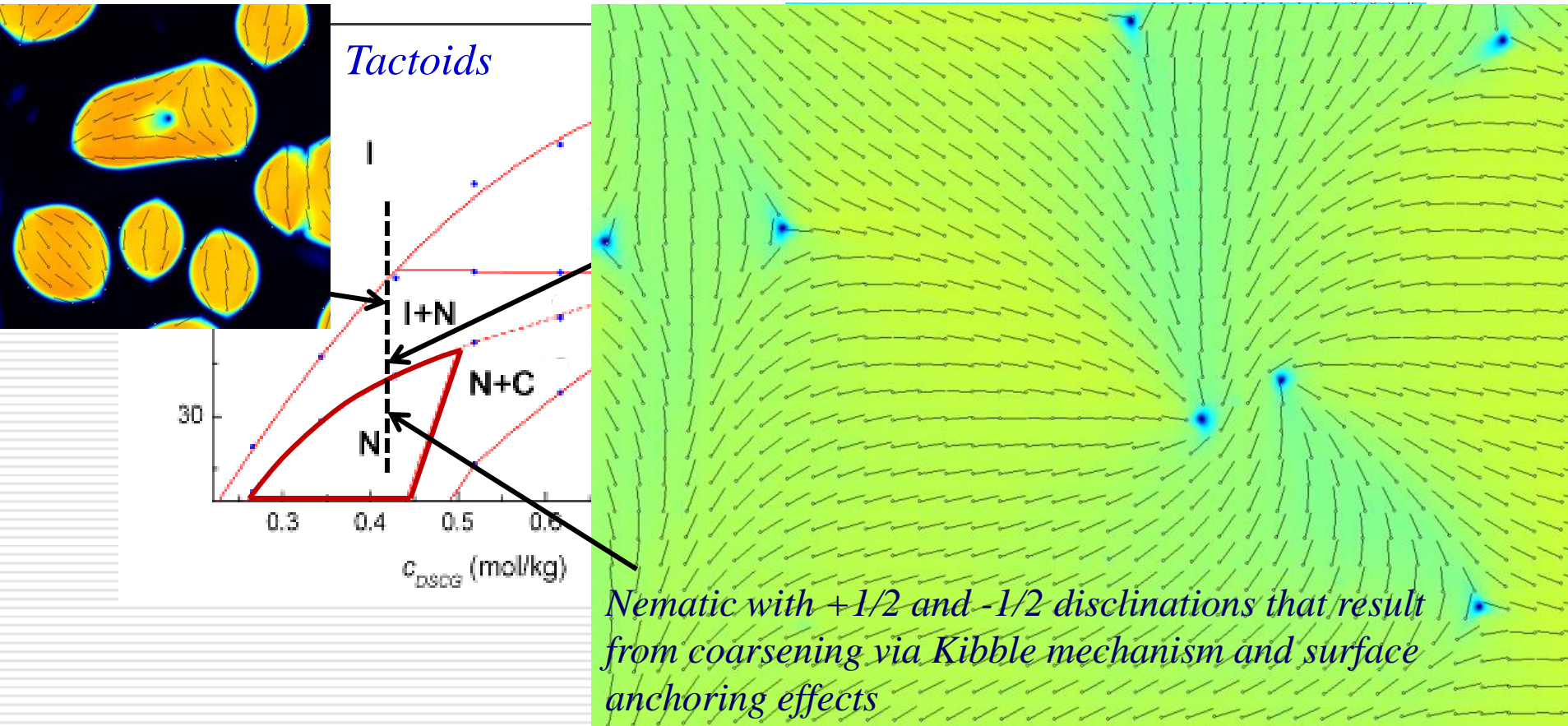


Chromonics (close to “living polymers” and “worm-like micelles”): phase diagram is controlled by both  $c$  and  $T$



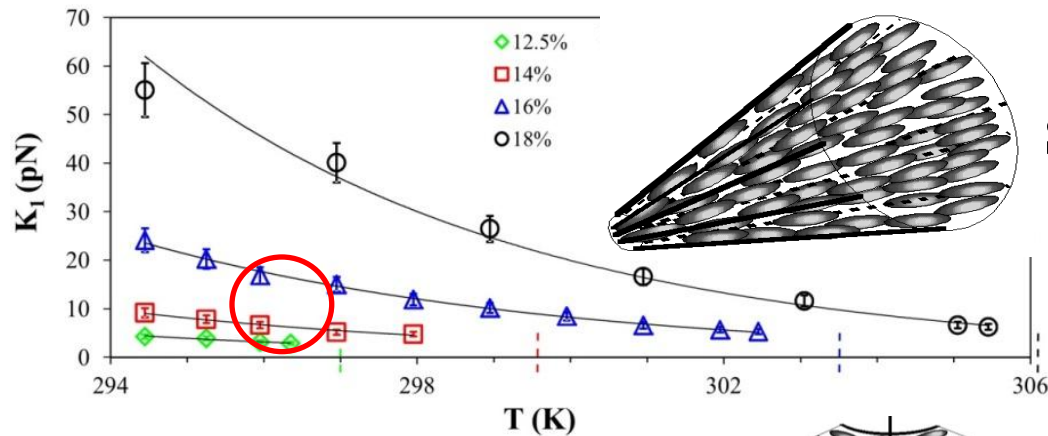


# Chromonics: $c$ and $T$ dependent phases

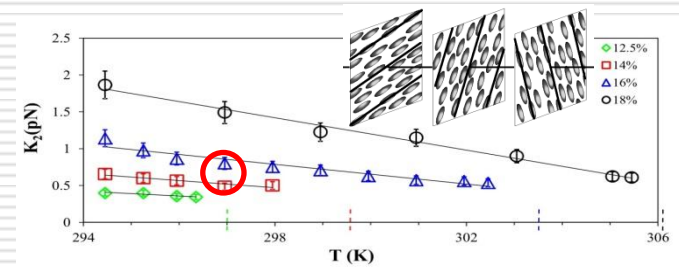


*Director field of real samples, reconstructed by LC PolScope (no simulations)*

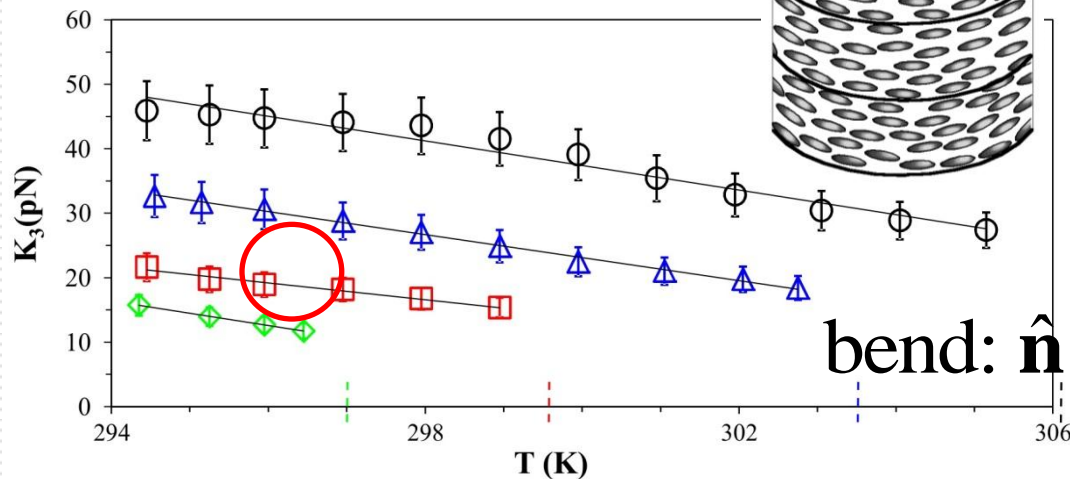
# $T, c$ dependent elastic constants



splay:  $\text{div} \hat{\mathbf{n}} \neq 0$

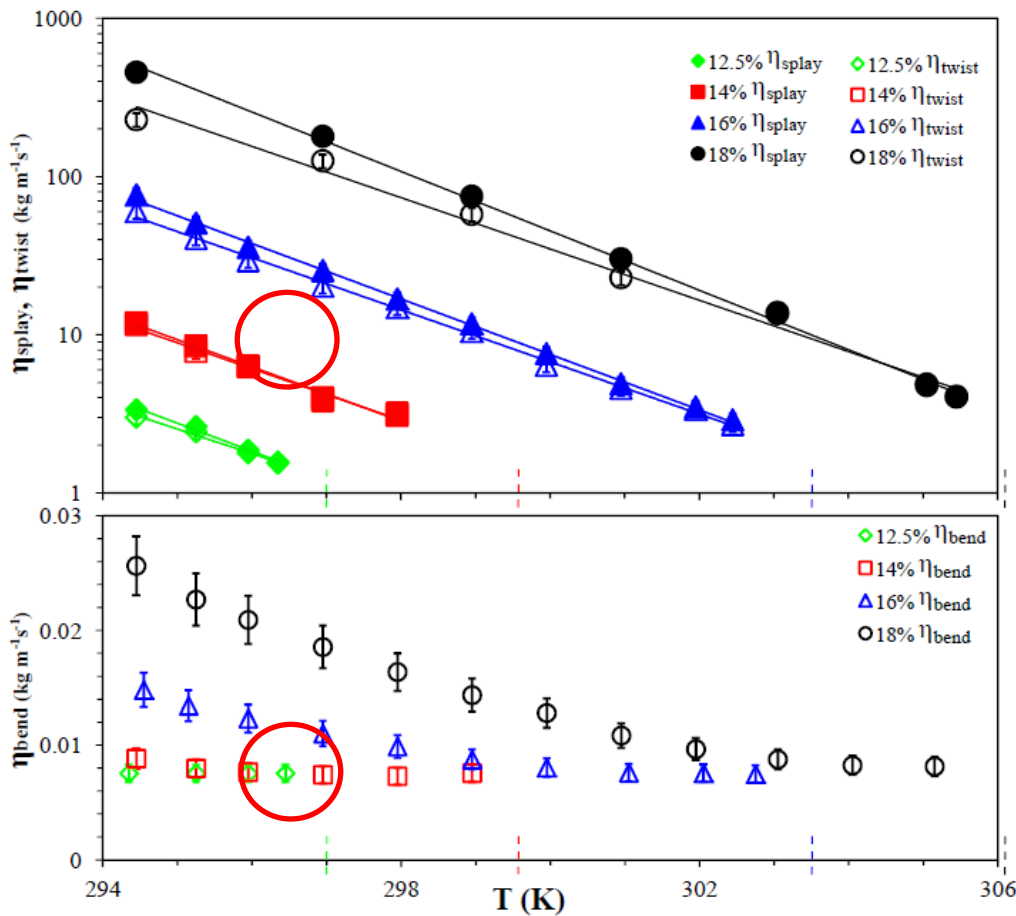


twist:  $\hat{\mathbf{n}} \cdot \text{curl} \hat{\mathbf{n}} \neq 0$



bend:  $\hat{\mathbf{n}} \times \text{curl} \hat{\mathbf{n}} \neq 0$

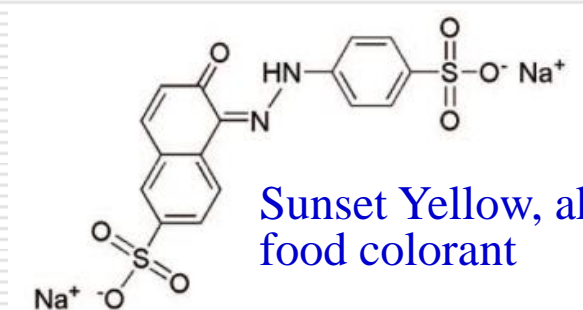
# $T, c$ dependent viscosities



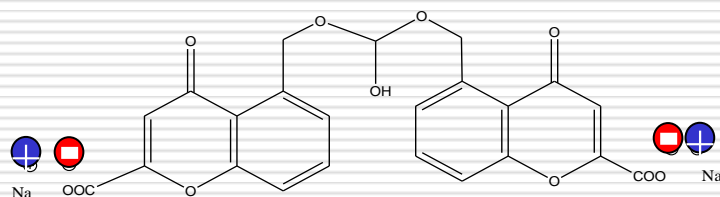


# Chemistry and biocompatibility of chromonics

- Chromonics are different from the surfactant based lyotropics: molecules have no aliphatic tails, thus chromonics are friendly to biological systems



Sunset Yellow, aka Yellow #6, food colorant



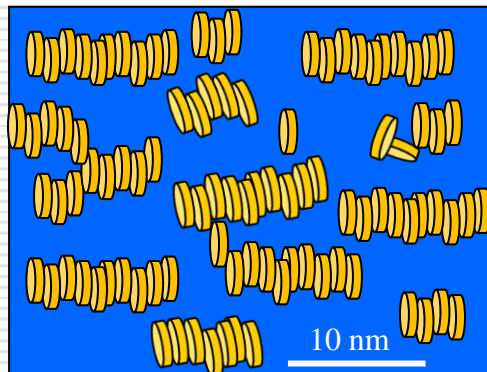
This work: DSCG, a.k.a. cromolyn, anti-asthma drug, concentration 14-17 wt%

Approved by FDA for intake by humans:



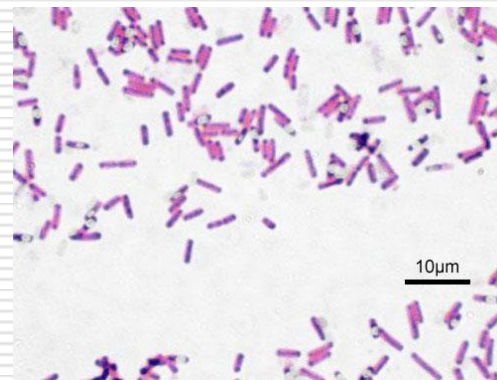
# Living Liquid Crystals

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- Living liquid crystals = *Chromonics* + *B. Subtilis*
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*Chromonics*

+

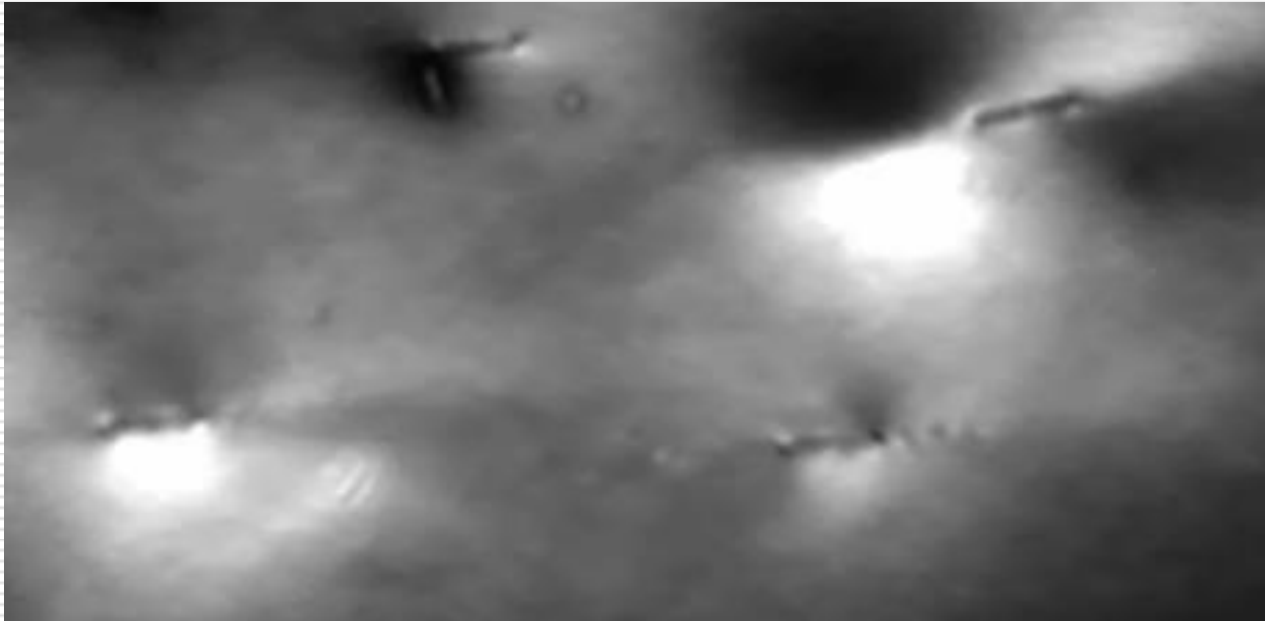


*Bacillus subtilis*

# Living LC: Individual bacteria follow $\hat{n} = \text{const}$

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Low concentration of bacteria ( $c < 10^8 / \text{cm}^3$ )



$\frac{P}{A}$

Rate  $\frac{1}{4}$

Isolated bacteria follow the uniform director, similar to observations by T. Galstian et al, Mol. Cryst. Liq. Cryst (2013) and N. Abbott et al, Soft Matter (2014)

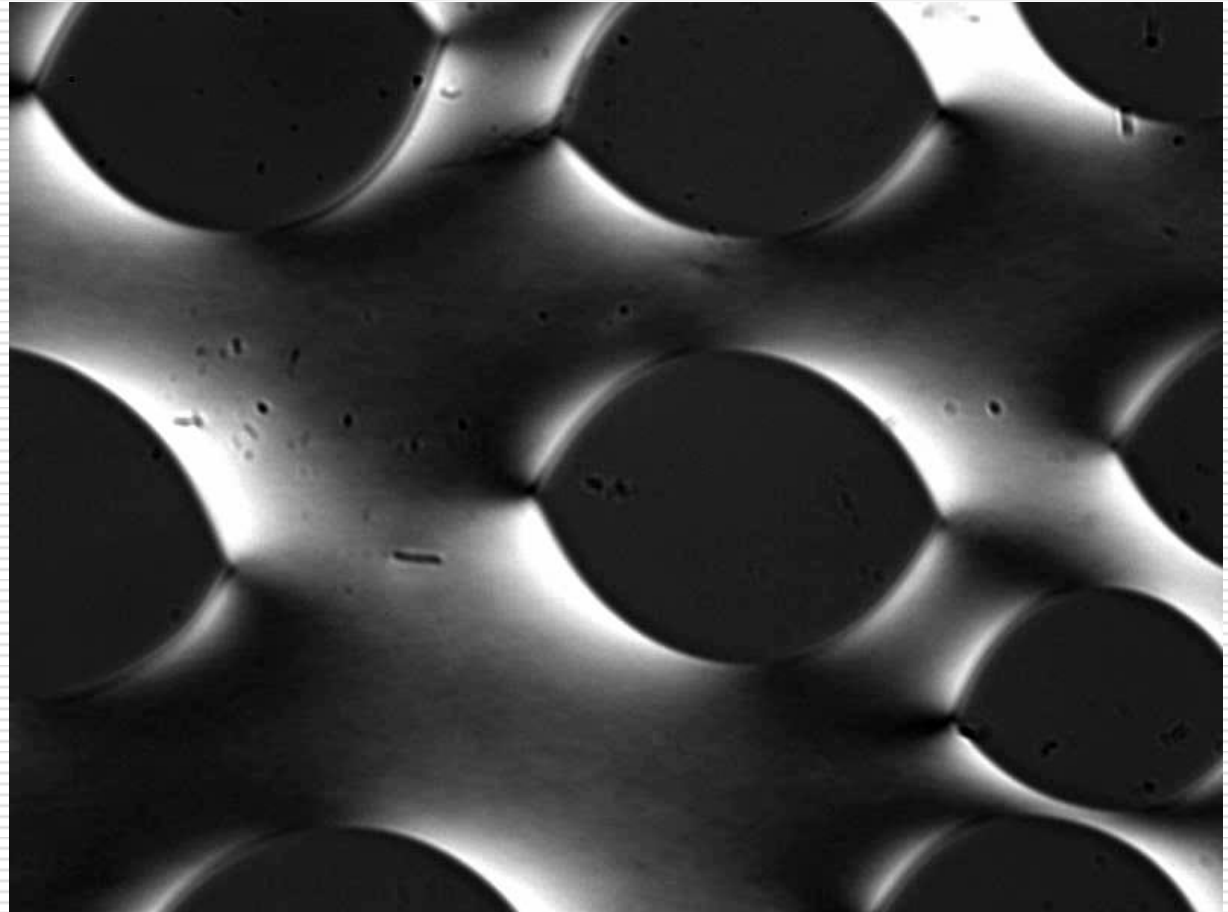
# Living LC: Individual bacteria follow $\hat{\mathbf{n}}(\hat{\mathbf{r}})$

...bacteria follow distorted director of the chromonic around isotropic tactoids

Rate  $\frac{1}{4}$

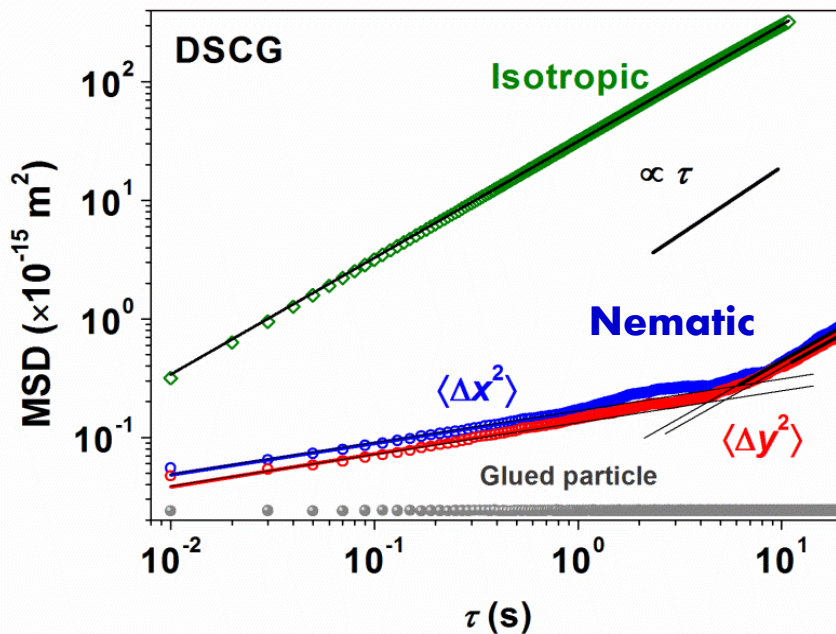


Isotropic tactoids



# Living LC: Individual bacteria follow $\hat{\mathbf{n}}(\hat{\mathbf{r}})$

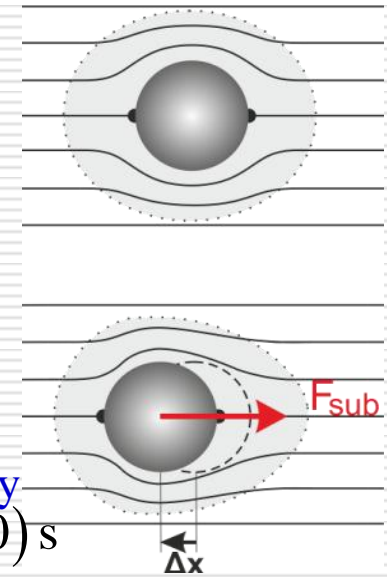
Both elastic and viscous effects favor swimming along the director lines



Diffusion of a  $5 \mu\text{m}$  sphere is easier parallel to the director as compared to perpendicular direction

$$D_{\parallel} = 6.8 \times 10^{-17} \text{ m}^2 \text{ s}^{-1}$$

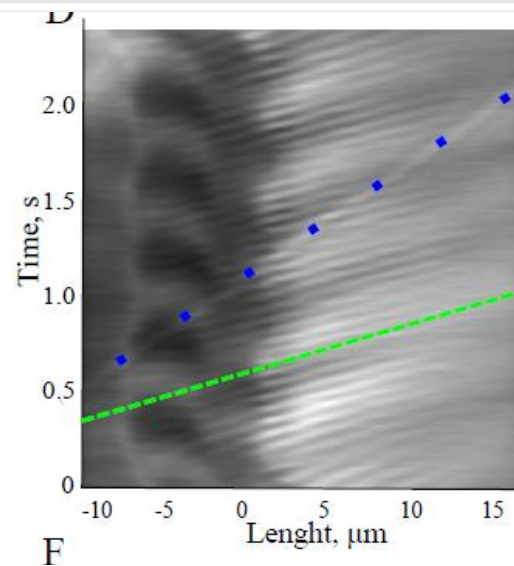
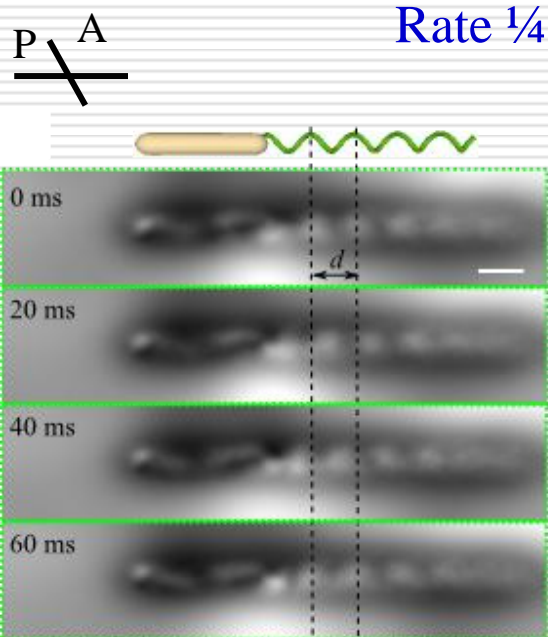
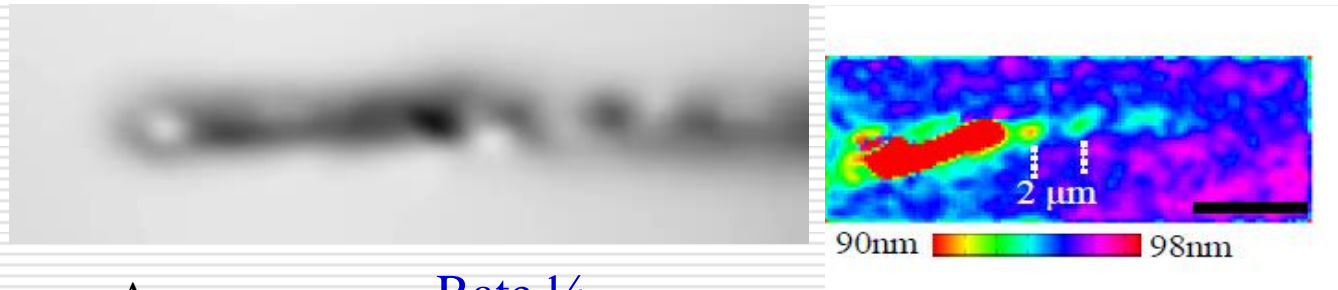
$$D_{\perp} = 4.5 \times 10^{-17} \text{ m}^2 \text{ s}^{-1}$$



Note sub-diffusive regime at short times: An intrinsic memory associated with the director relaxation  $\tau_d : \eta d^2 / K \sim (1-10) \text{ s}$



# Living LC: Individual bacteria distort $\hat{n}(\hat{r})$



Flagella rotation: 16 Hz  
 Body rotation:  $f=2.5$  Hz

The rotation is fast enough  
 to make the Ericksen  
 number larger than 1

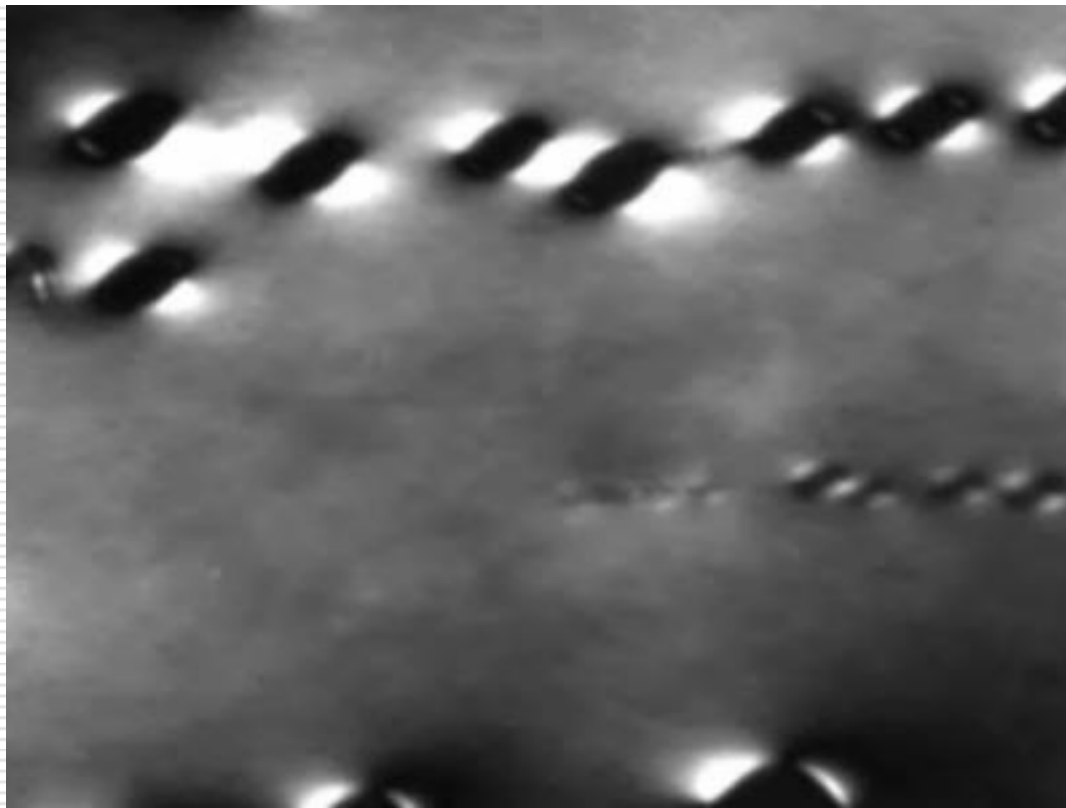
$$Er = \eta_{eff} f r h / K \sim 10$$

implying that the director is  
 distorted by moving flagella

# Living LC: Individual bacteria melt LC

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Moving bacterium can also change the scalar order parameter, melting the material and forming isotropic droplets-tactoids in its wake (“Wilson chamber”)



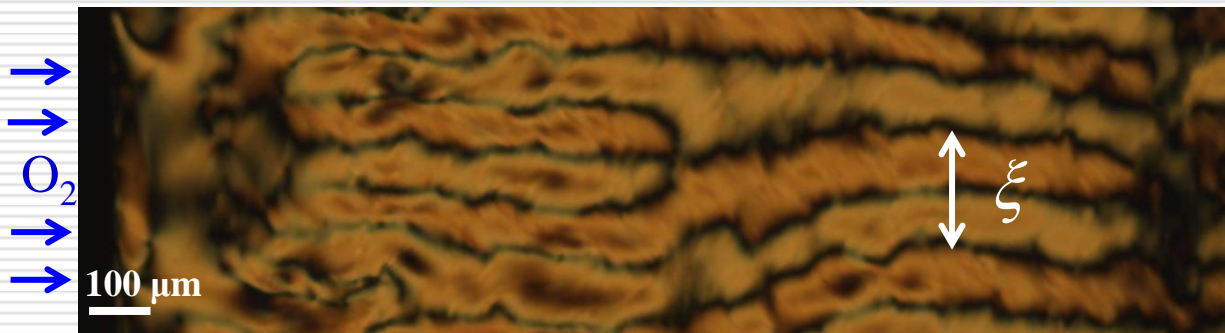
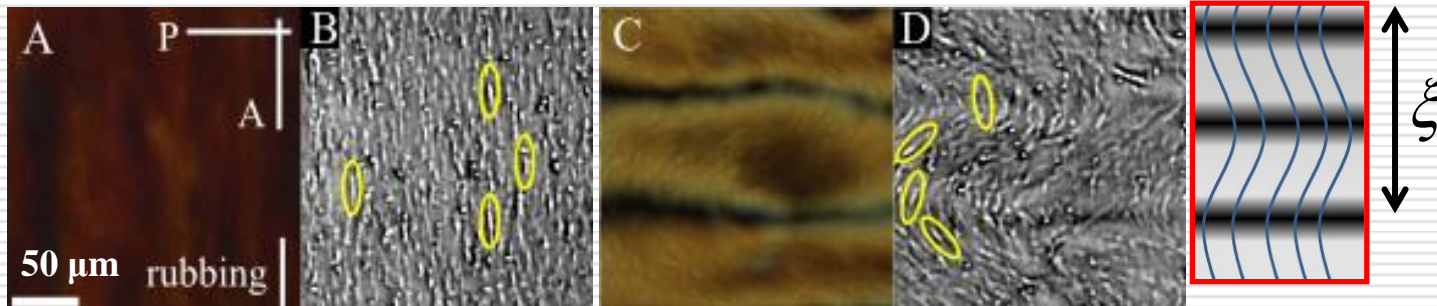
Rate  $\frac{1}{2}$

# Living LC: Collective effects, Bend stripes

No oxygen; equilibrium state of uniform director

Added oxygen; director undulations

Higher concentration of bacteria ( $c_B \sim 10^9 / \text{cm}^3$ )



Oxygen supplied from the left hand side

# Living LC: Collective effects, Bend stripes

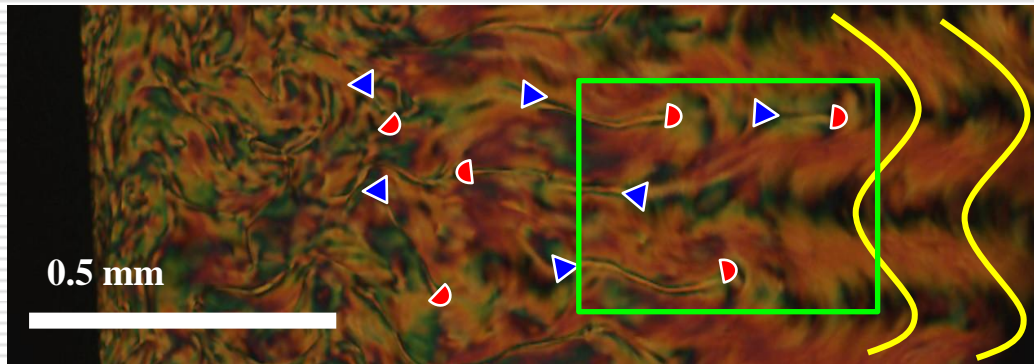
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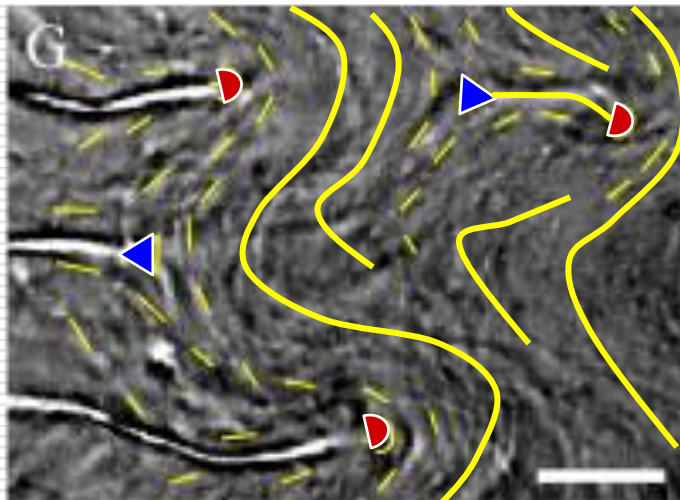
High concentration of bacteria, addition of oxygen: periodic undulations with a characteristic spatial scale that depends on  $c_B$ , amount of oxygen, etc.

Oxygen supplied from the left hand side, rate 100

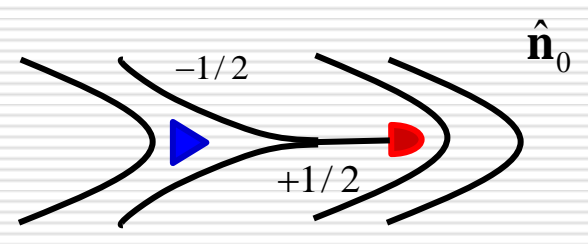
# Higher activity: Bend stripes replaced by disclination pairs



Nucleation of disclination pairs

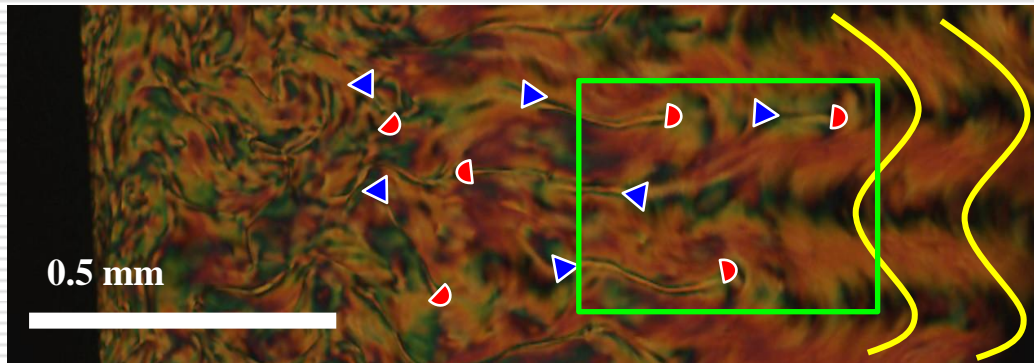


Director within the pair realigned by  $90^\circ$  w.r.t. the original director





# Bending vs disclination pairs: Elasticity and surface anchoring



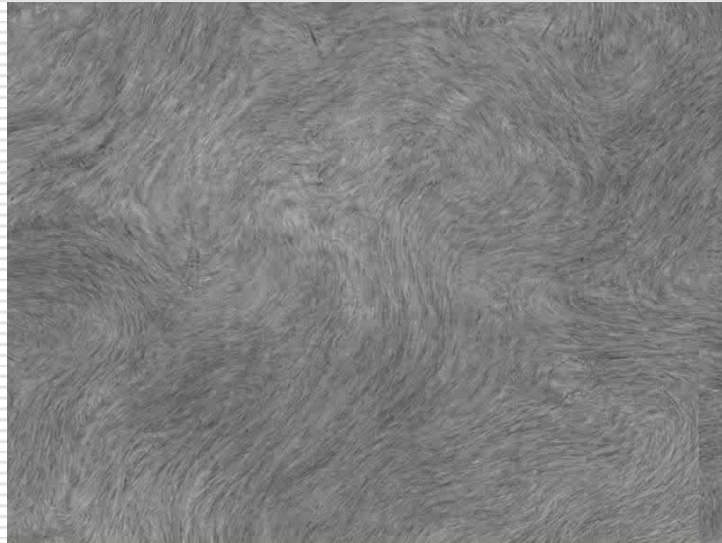
As activity increases, the uniform state (1) undergoes transition into periodic bend stripes, then (2) into the sea of nucleating and annihilating disclination pairs. Similar 2 stage scenario is seen in numerical simulations:

Thampi, Golestanyan, Yeomans, EPL (2014); poster at BIOACTER;  
M. Shelley et al, <http://arxiv.org/pdf/1401.8059v1.pdf>  
See also Shi, Ma, Nature Comm (2013)

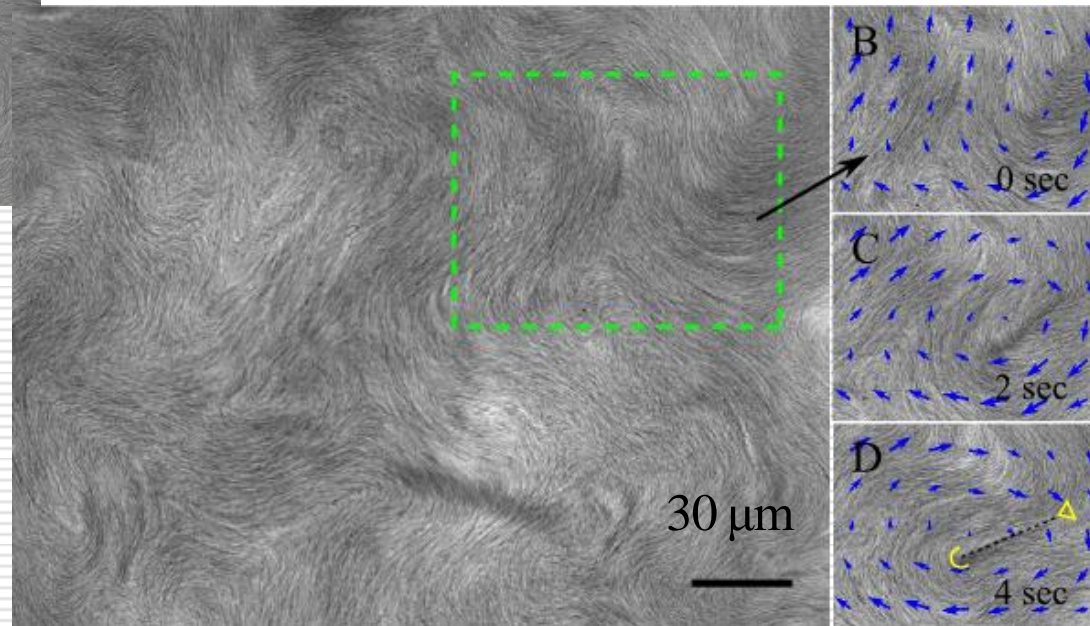
Why the walls are replaced by disclination pairs?

$$\text{Walls: } F_w \propto \frac{K}{\xi^2} \quad \text{Pairs: } F_w \propto \frac{K}{\xi_d^2} \ln \xi_d / r_c$$

# High activity: Topological turbulence

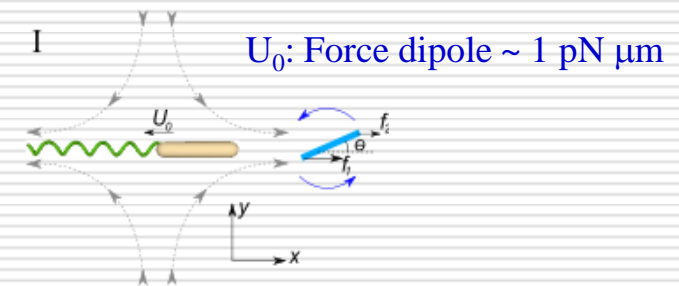
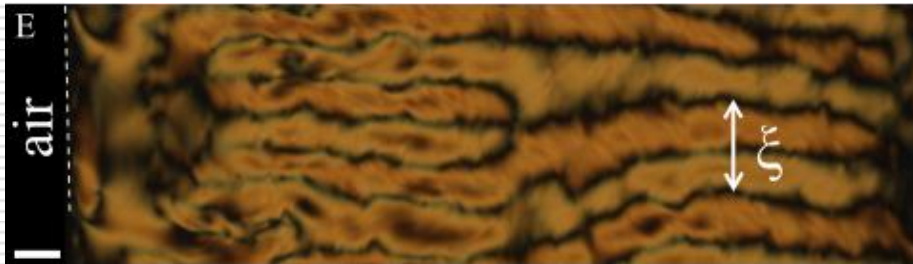


Blue arrows: LLC flow velocity field



# Bending: Activity vs Elasticity

Spatial scale: balances viscous shear (bacterial) and elastic (LC) torques

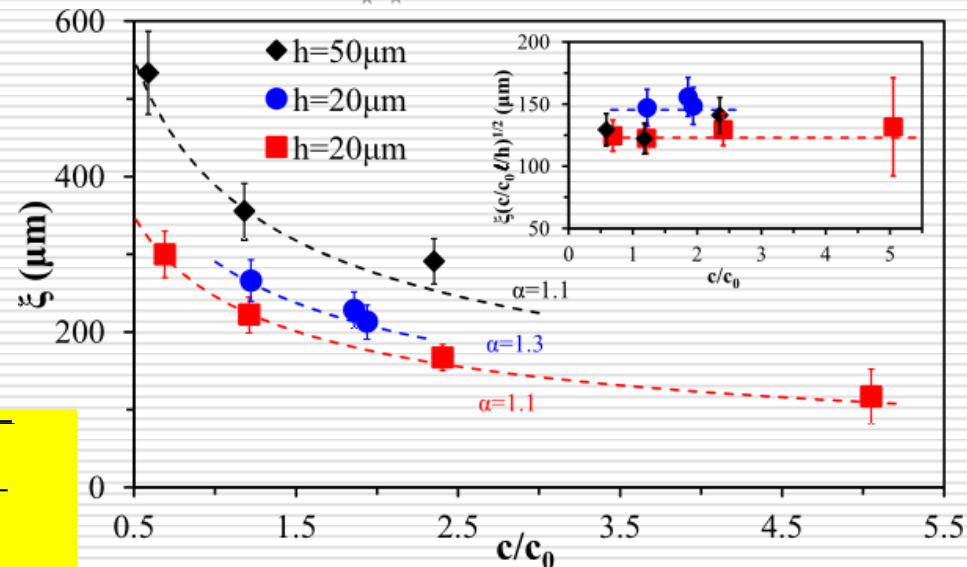


$$\left. \begin{aligned} \Gamma_{shear} &\sim \alpha c U_0 \theta \\ \Gamma_{elastic} &\sim K \frac{\partial^2 \theta}{\partial r^2} \end{aligned} \right\} \Rightarrow \xi = \sqrt{\frac{K}{\alpha c U_0}}$$

Similar to Ramaswamy Ann. Rev. CMP (2010)

Cell thickness correction (mass conservation)  $\alpha \rightarrow \alpha_0 l / h \Rightarrow$

$$\xi = \sqrt{\frac{Kh}{\alpha_0 l c U_0}}$$



# Living Liquid Crystals: Summary

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- Non-uniform director guides nonlinear trajectories of bacteria
- Birefringence-enhanced visualization of microflows generated by nanometers-thick flagella
- Bacteria produce local melting of chromonic LC
- Activity increase drives a transition from a uniformly ordered state to the topological turbulent state through two steps:
  - Formation of periodic bend, period  $\xi = \sqrt{K / \alpha c U_0}$
  - Nucleation and proliferation of disclination pairs