

# Energy barriers and cell migration in densely packed tissues

**Dapeng (Max) Bi**  
*Syracuse University*

Active Matter: Cytoskeleton, Cells, Tissues and Flocks

KITP

Apr 2, 2014



10 um

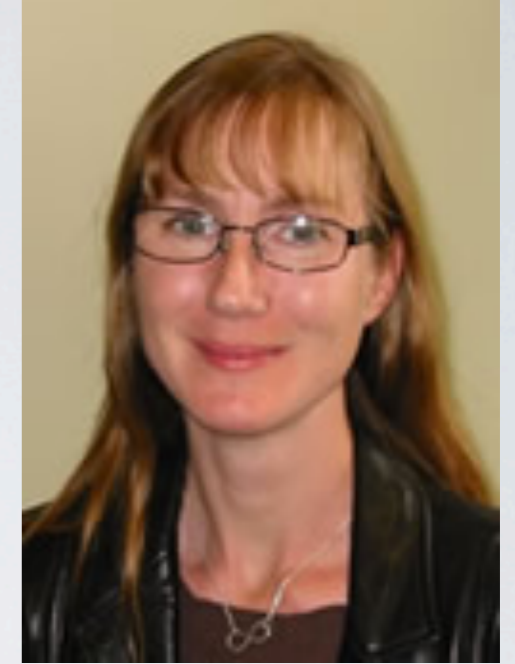
# Collaborators



M. Lisa Manning



Jorge Lopez



Jen Schwarz



Syracuse University Chancellor Fund  
NSF-DMR-0645373

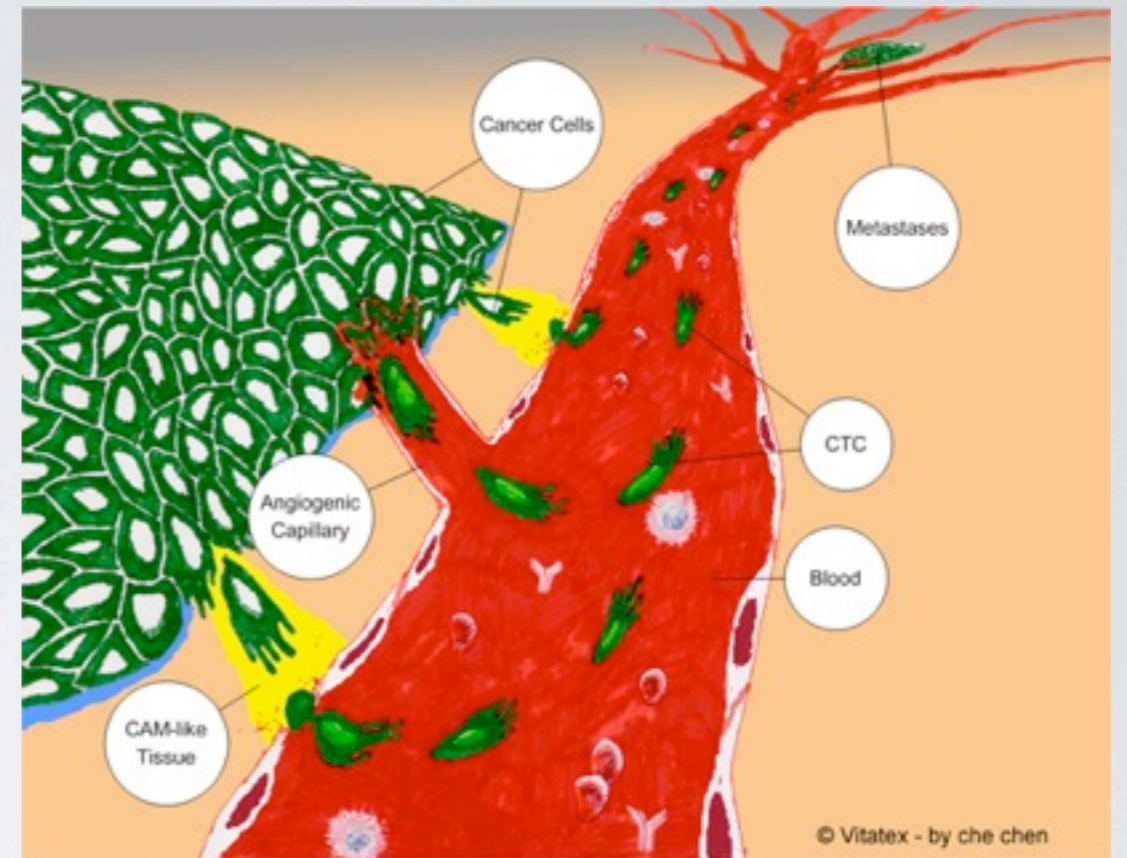
[bdpmax@gmail.com](mailto:bdpmax@gmail.com)

[WWW.DAPENGBI.COM](http://WWW.DAPENGBI.COM)

# Why study cell migration in dense tissues?

Wound healing assay

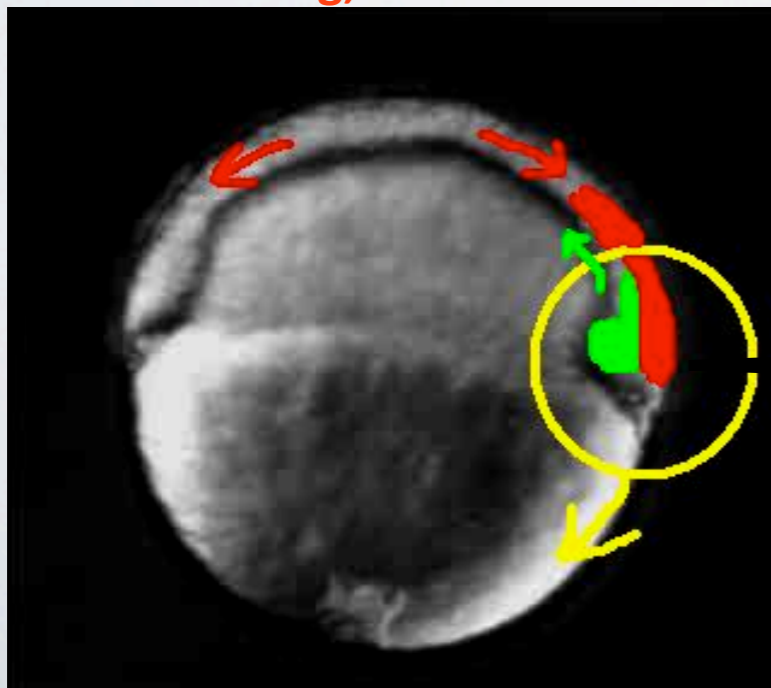
Tumor cell migration



[youtube.com/watch?v=v9xq\\_GiRXeE](https://www.youtube.com/watch?v=v9xq_GiRXeE)

During development of embryo

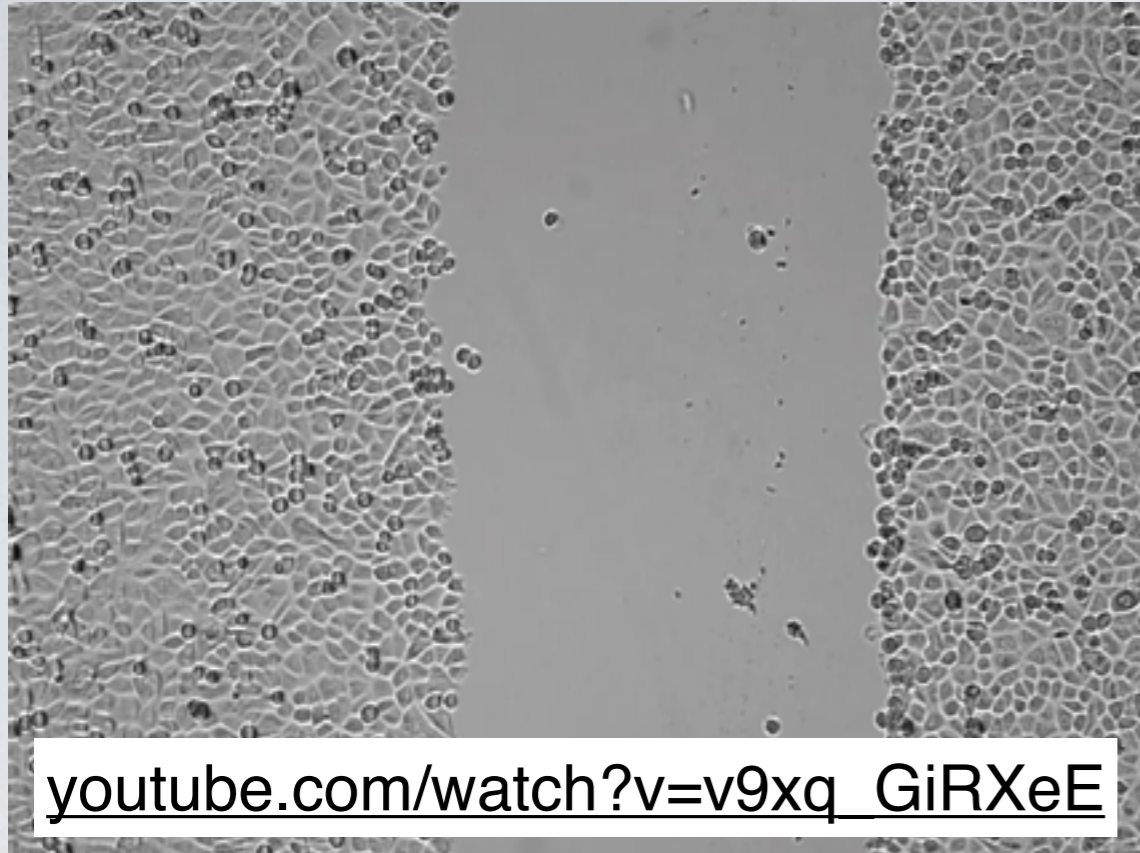
**Differential Adhesion Hypothesis:**  
Steinberg, Science 1962



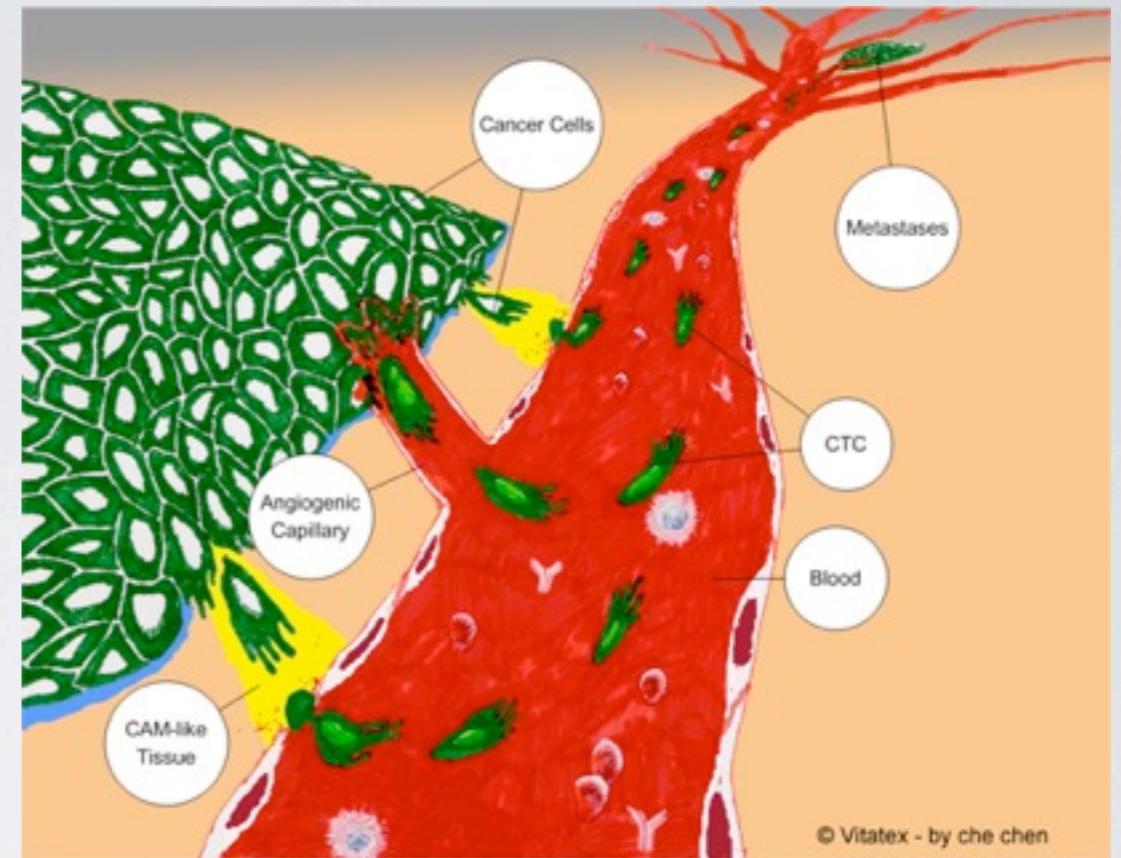
**E-M. Schöetz Thesis 2008**

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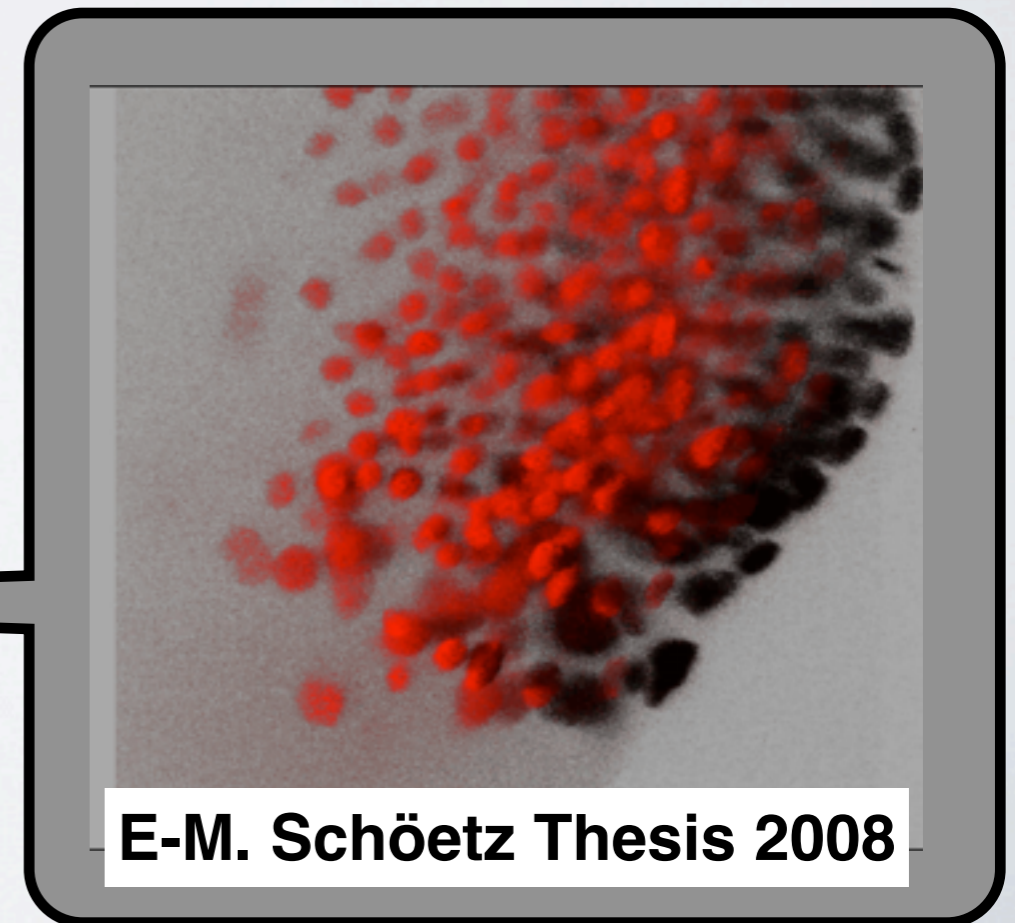
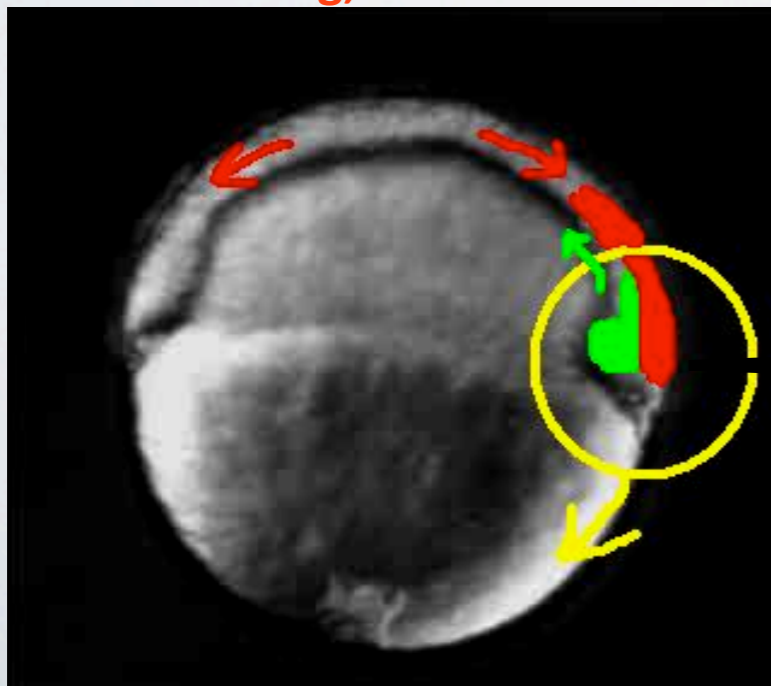


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# How do cells move?

**In isolation**

**vs.**

**inside dense tissue**

**Human bone cancer  
cell on fibronectin**

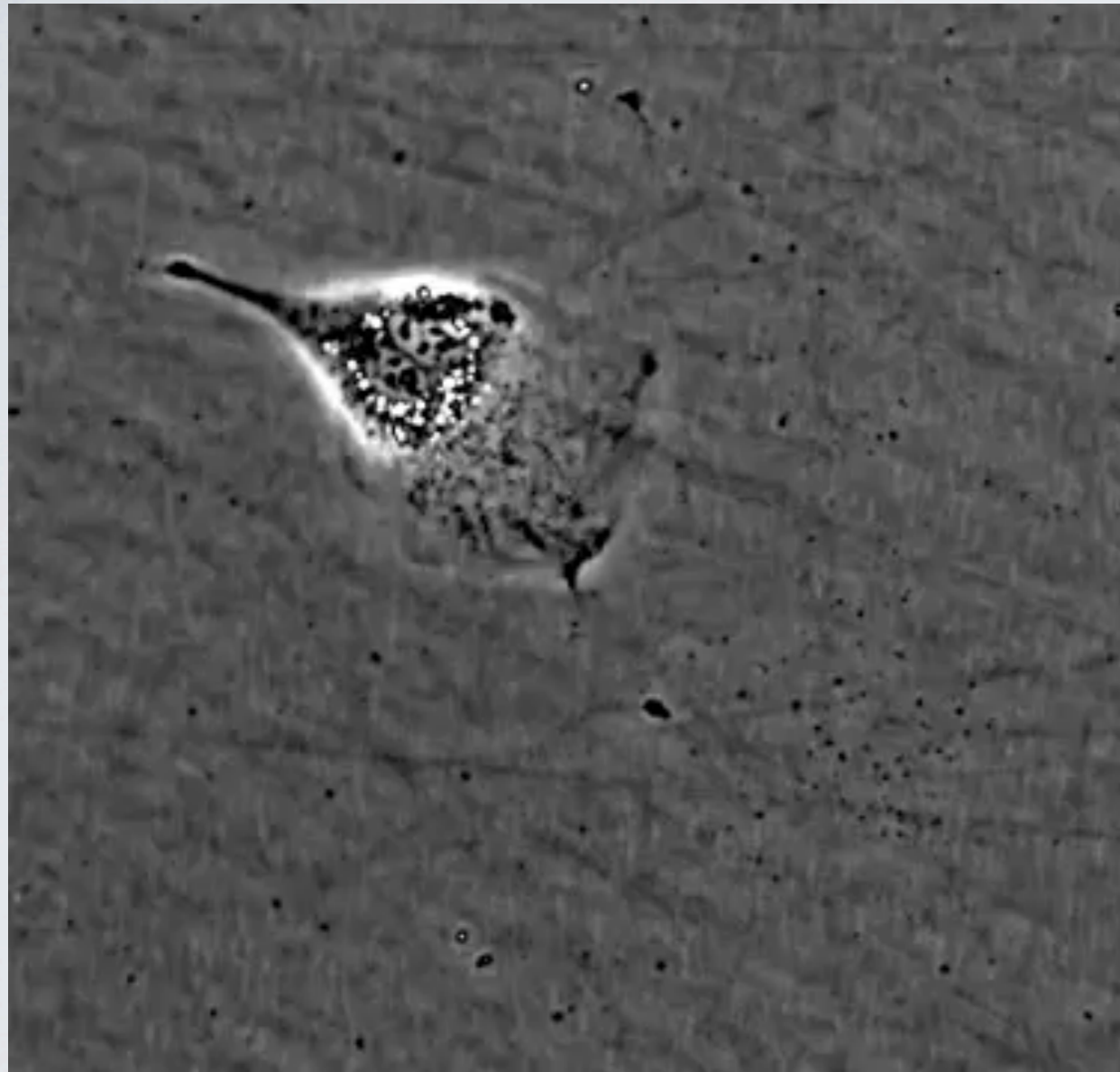
**Zebrafish embryo  
Schöetz Lab, UCSD**

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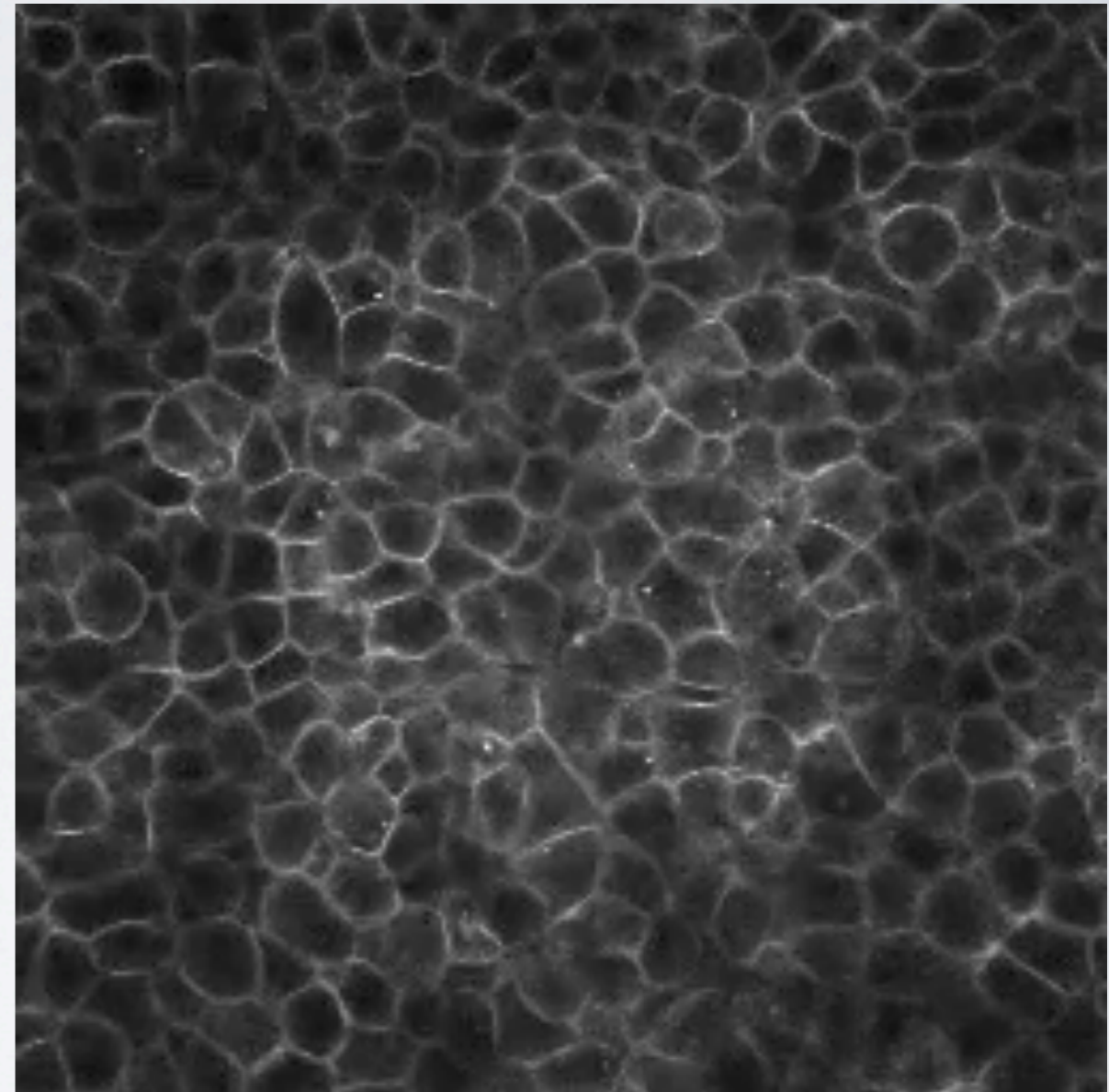
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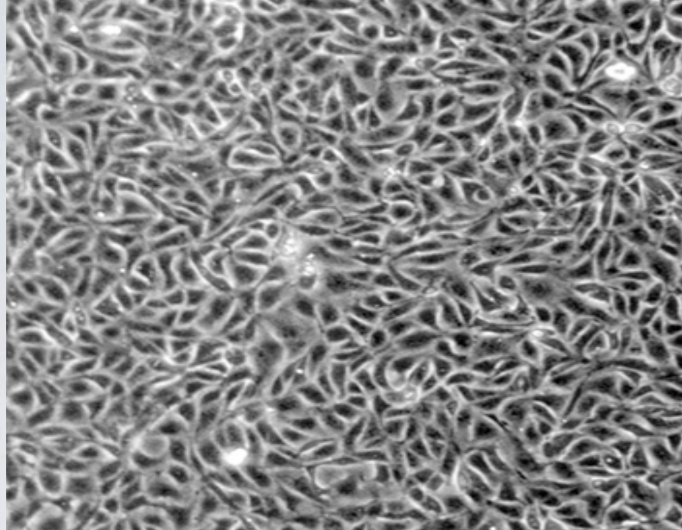


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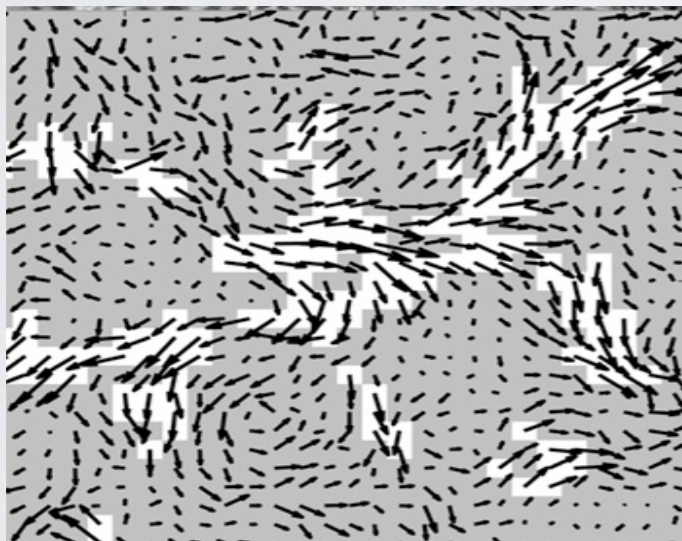
# Similarity to glassy materials

Living and Active

## Energy injection by cell



Madin-Darby canine kidney (MDCK) cells forming a 2-d confluent layer.



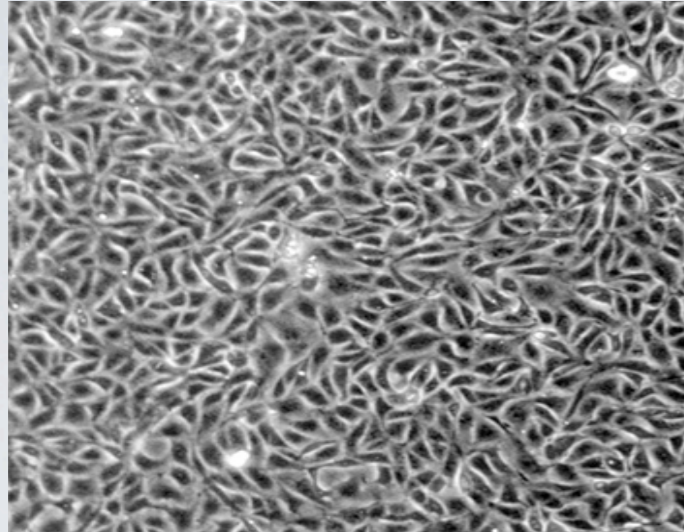
Velocity profile of cells show the spatially heterogeneous pattern in MDCK tissue.

**Angelini et al PNAS 2010**

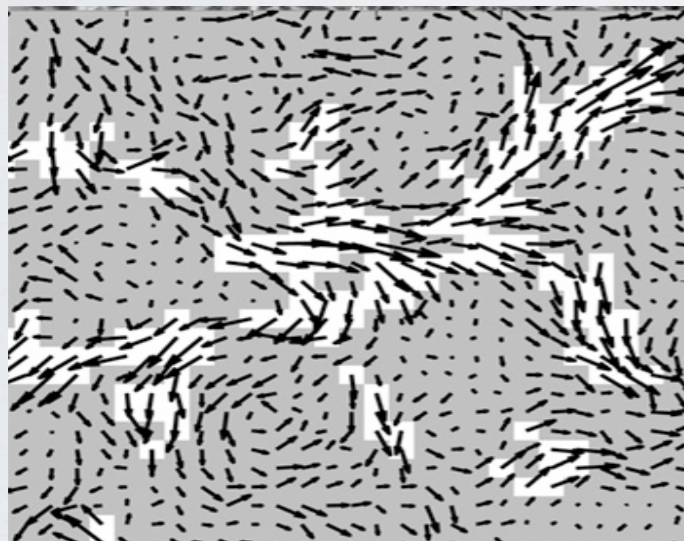
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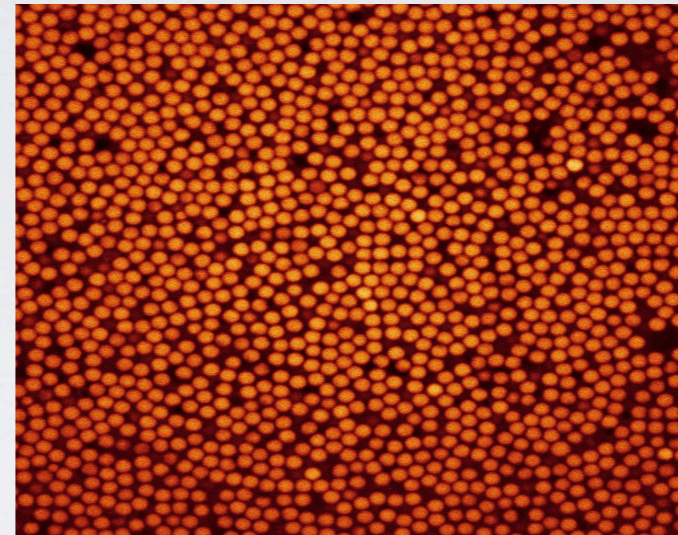


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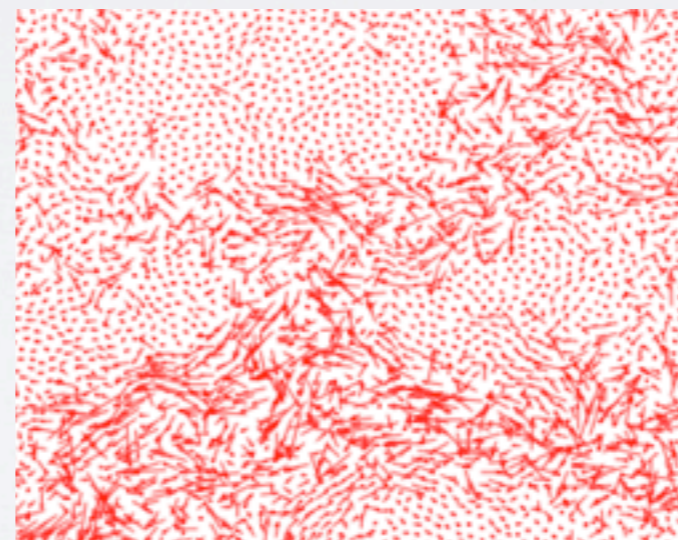
**Angelini et al PNAS 2010**

Non-living

## Thermal



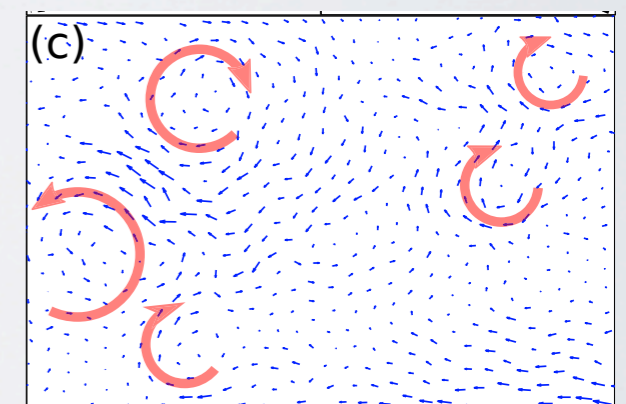
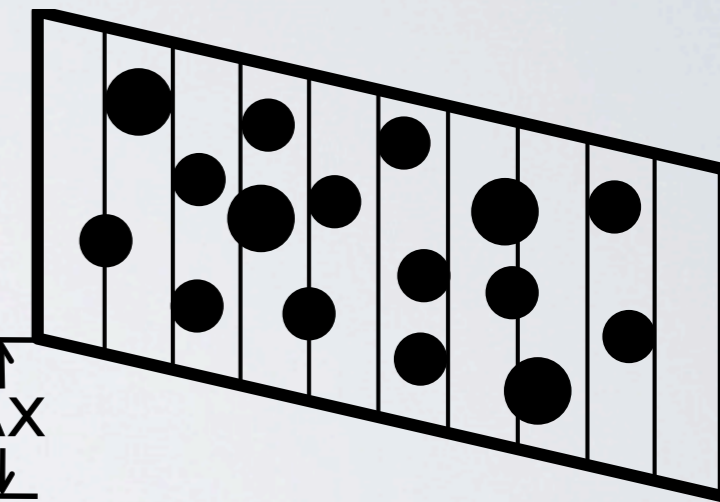
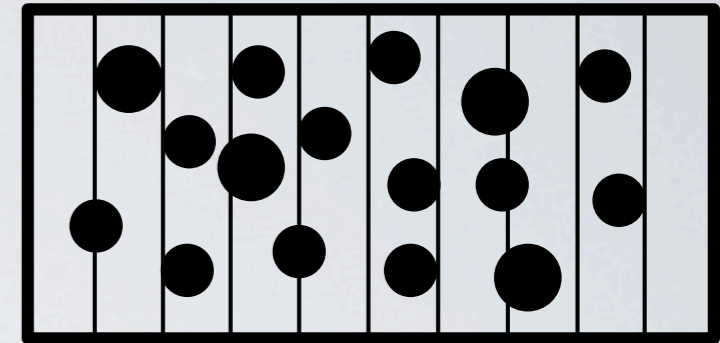
A colloidal glass.



Displacement profile in simulation of a 2-d glass former.

**Berthier PRL 2011**

## Boundary Driven



Sheared granular material  
**Jie Ren thesis 2013,  
Duke Univ.**

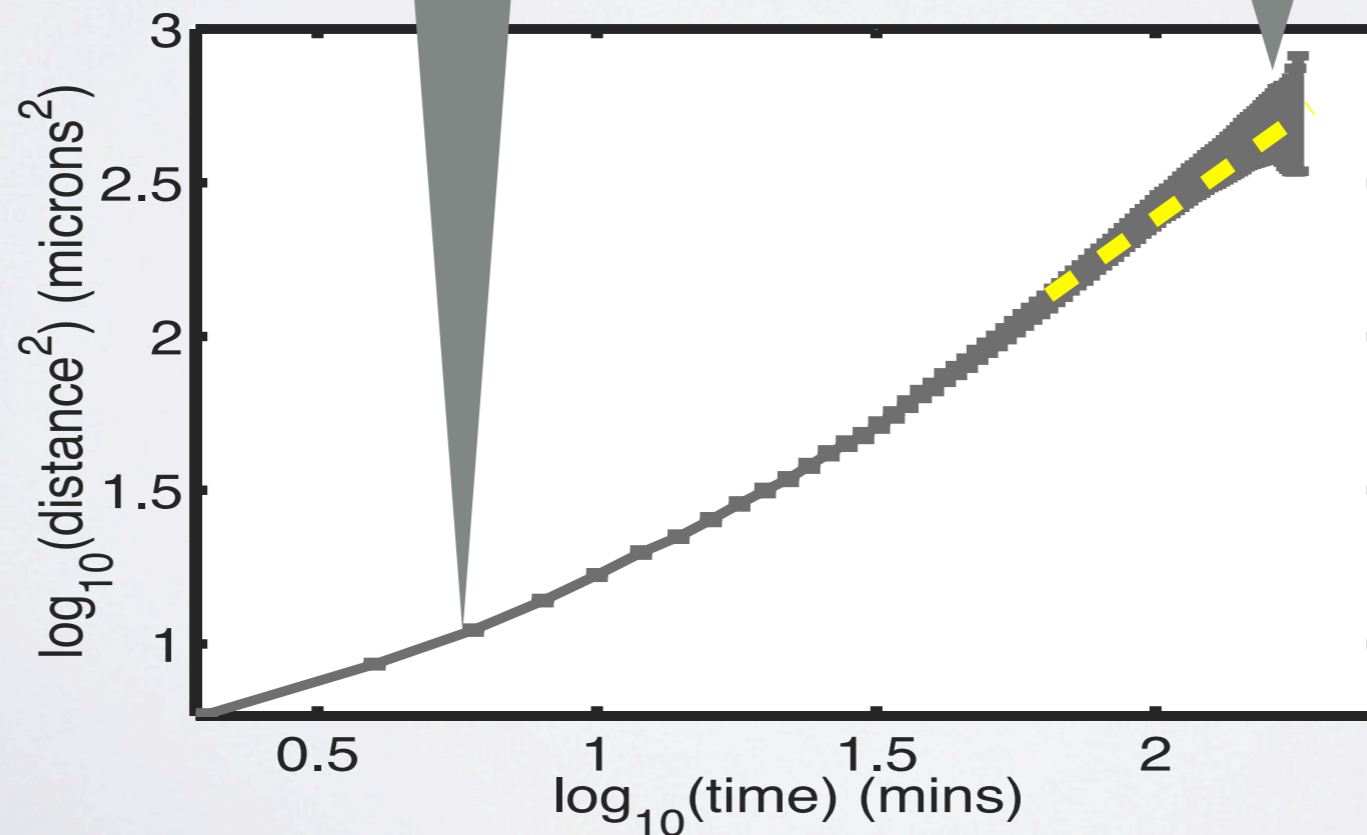


# Temporal behavior: viscoelasticity

timescale ~ minutes

about 5000 cells

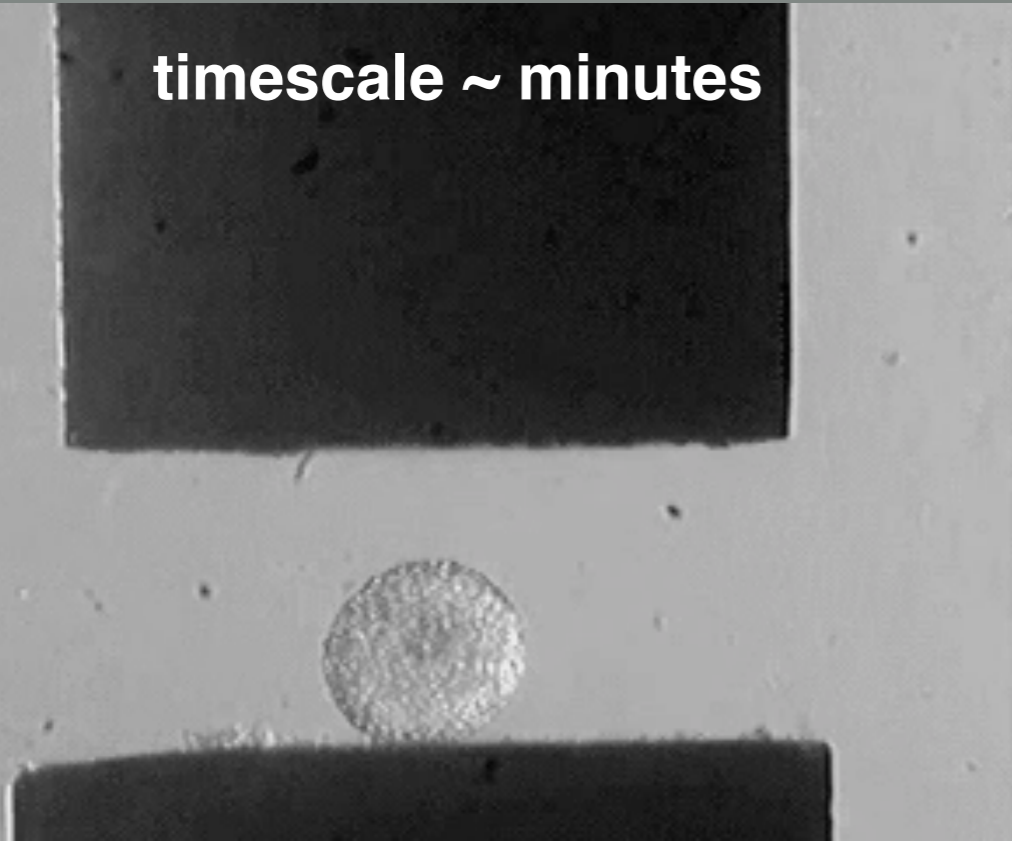
timescale ~ hours



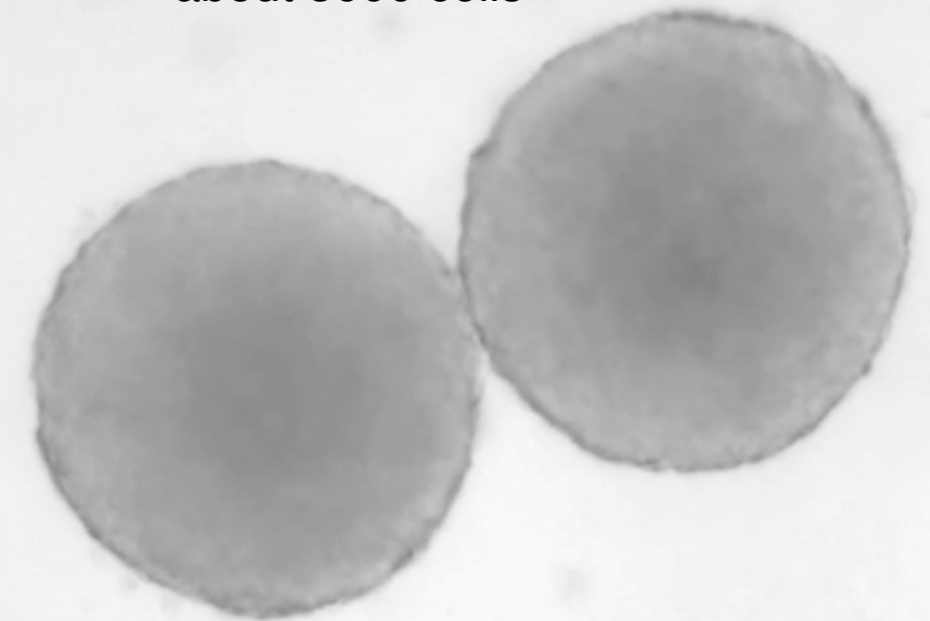
Schöetz et al  
J. R. Soc. Interface 2013

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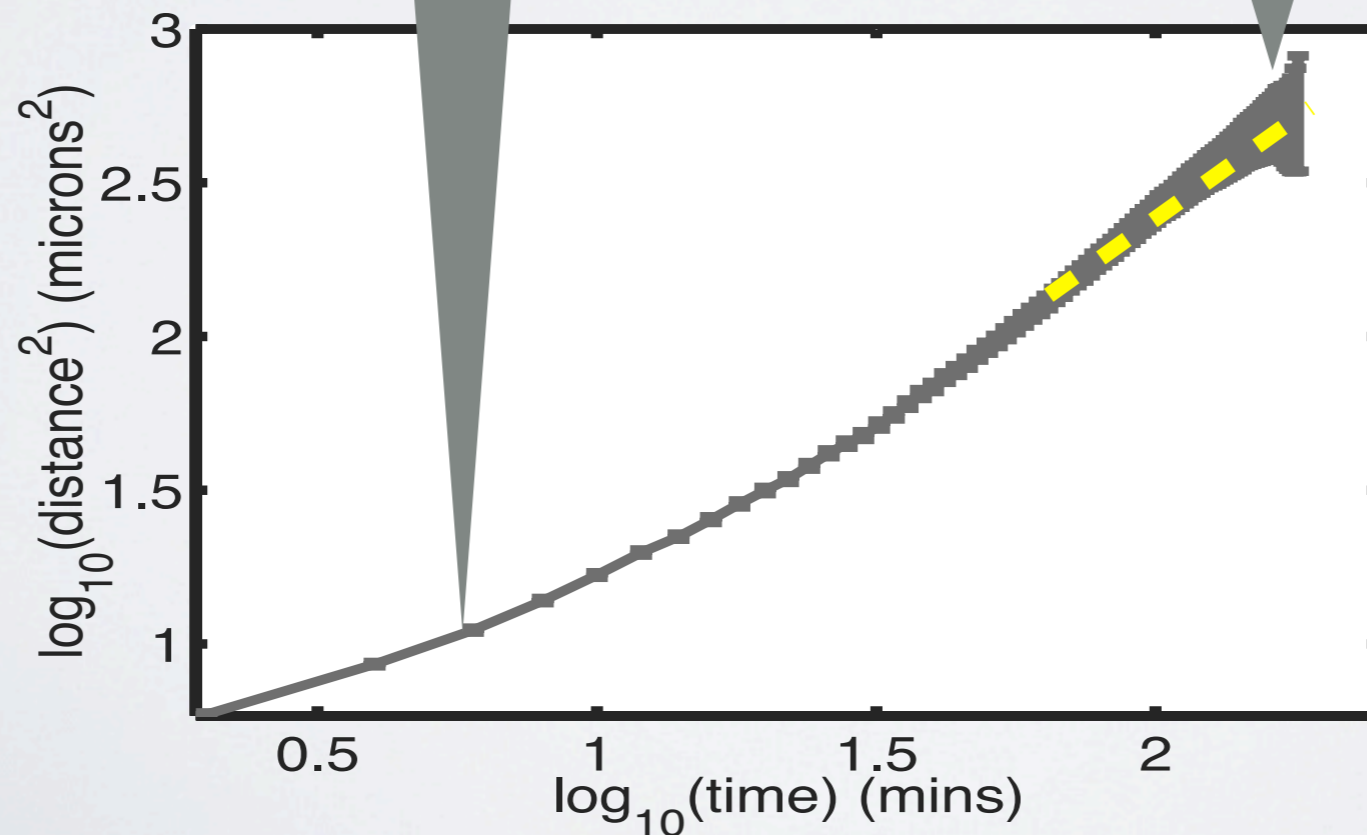
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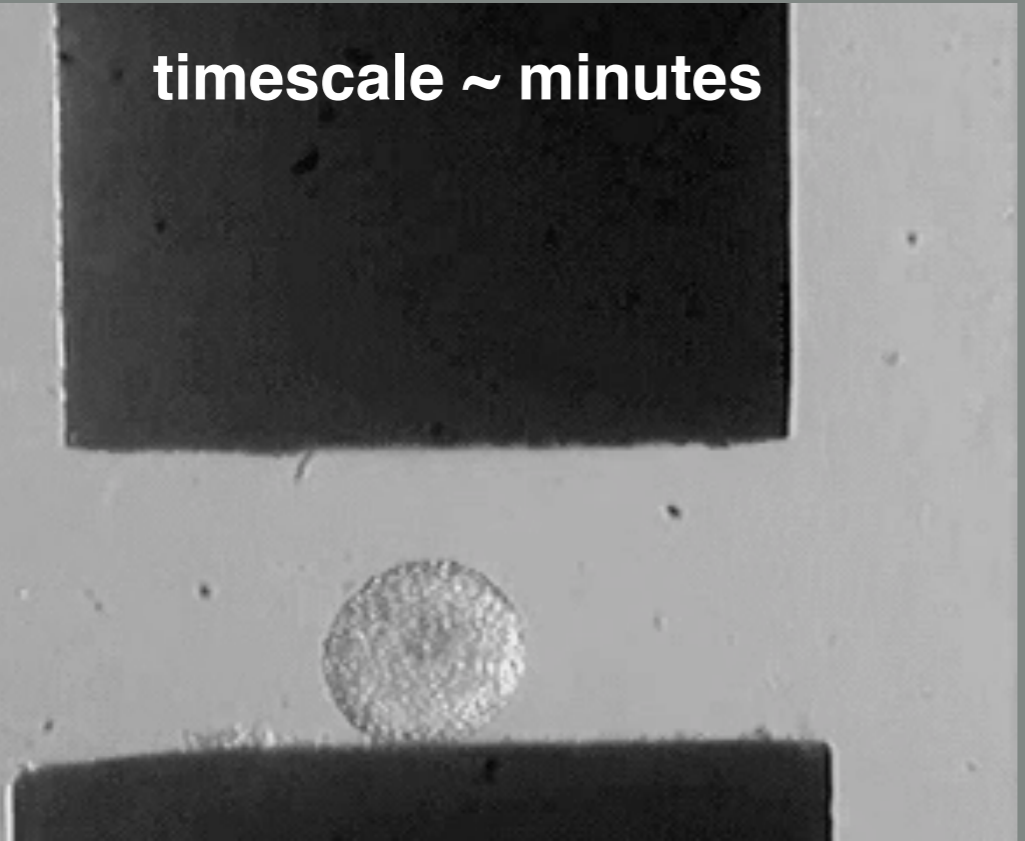
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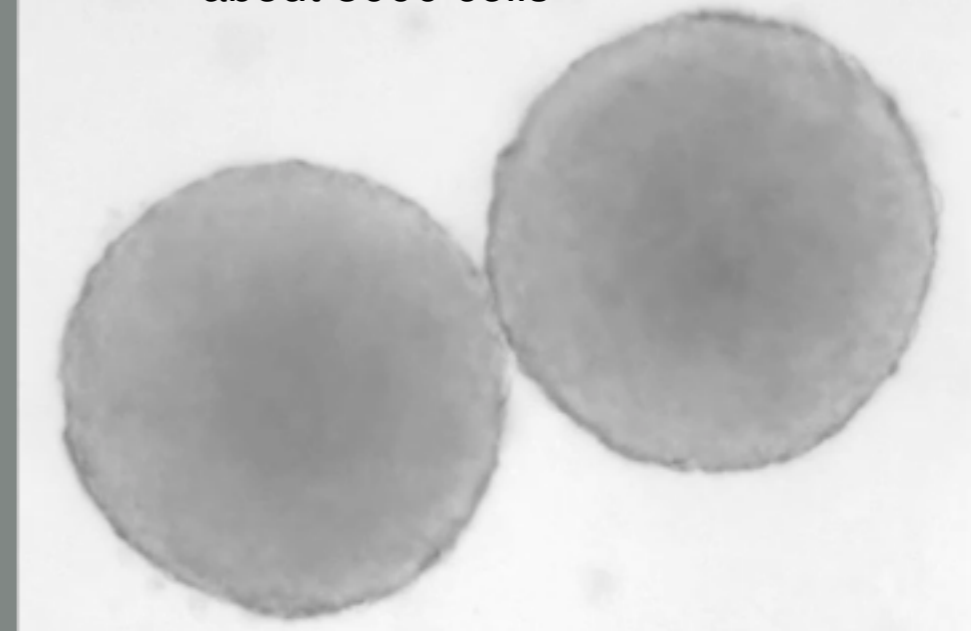
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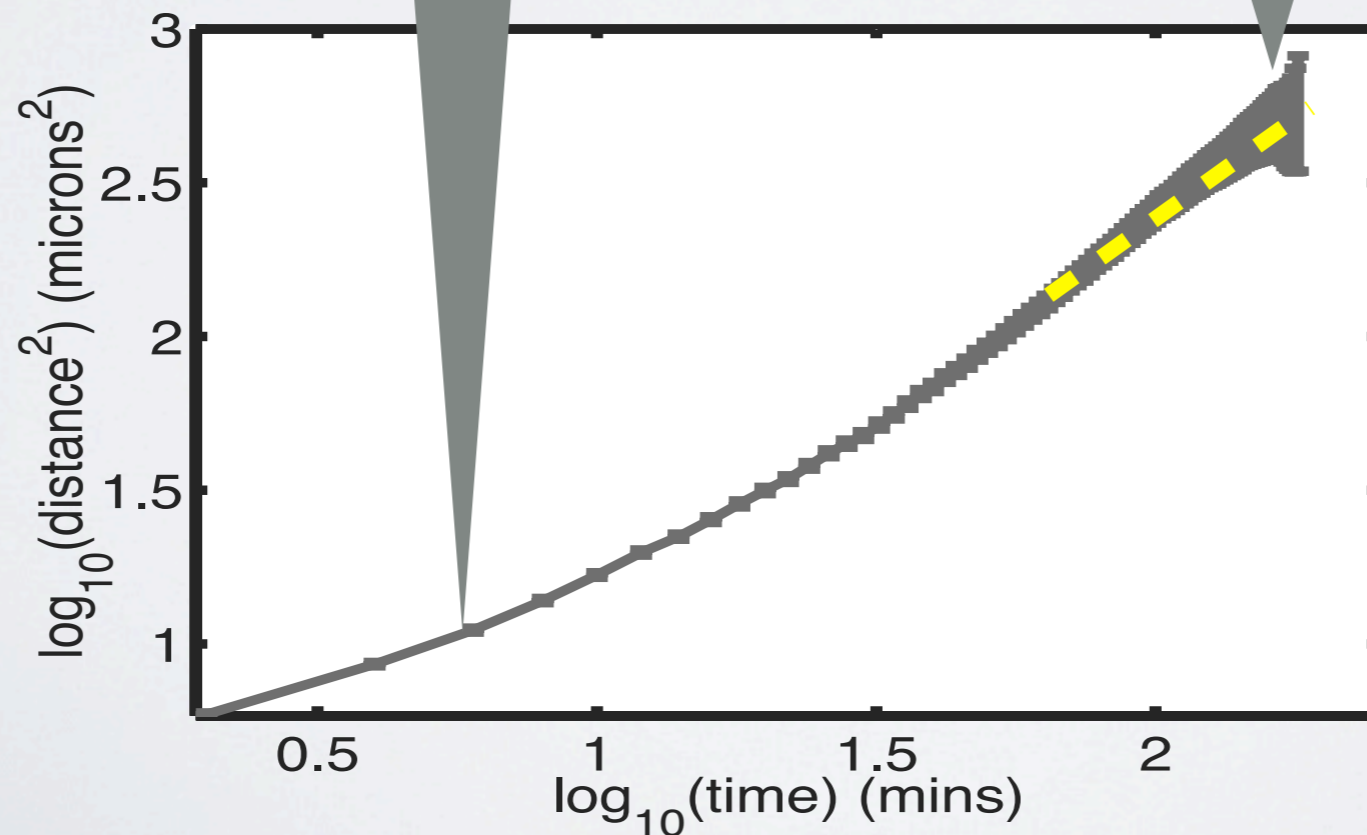
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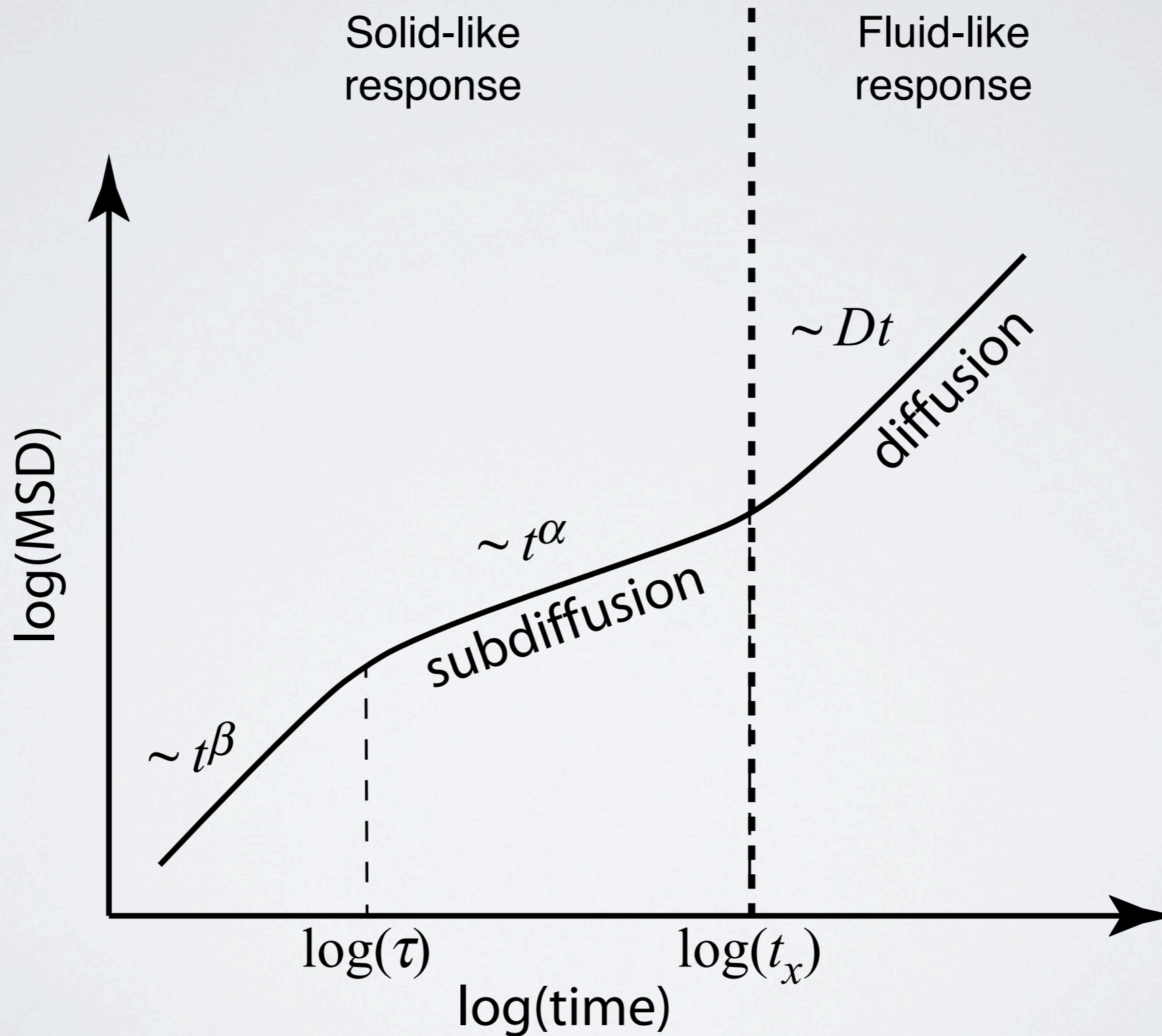


timescale ~ hours

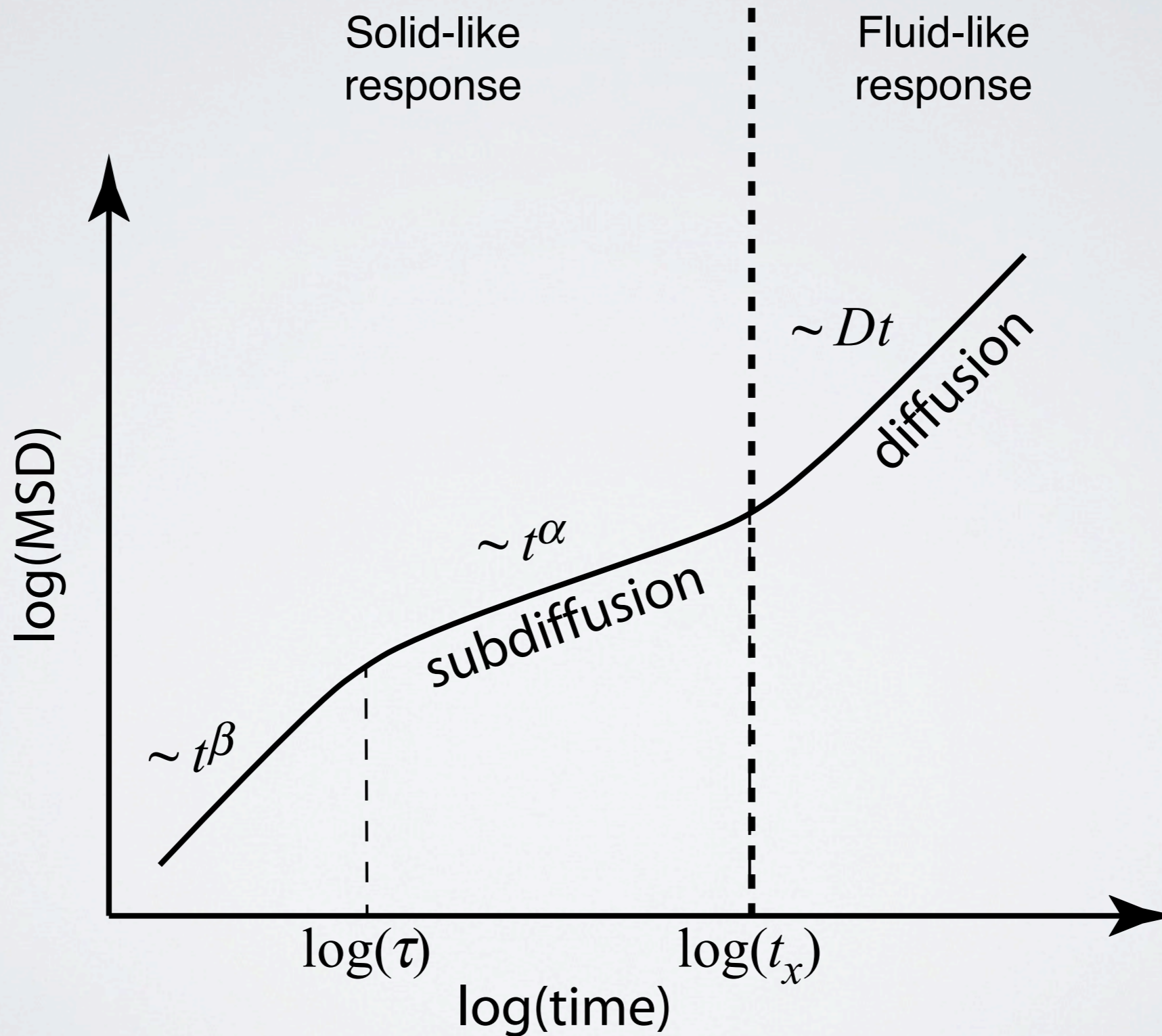


Schöetz et al  
J. R. Soc. Interface 2013

# Similarity to glassy materials



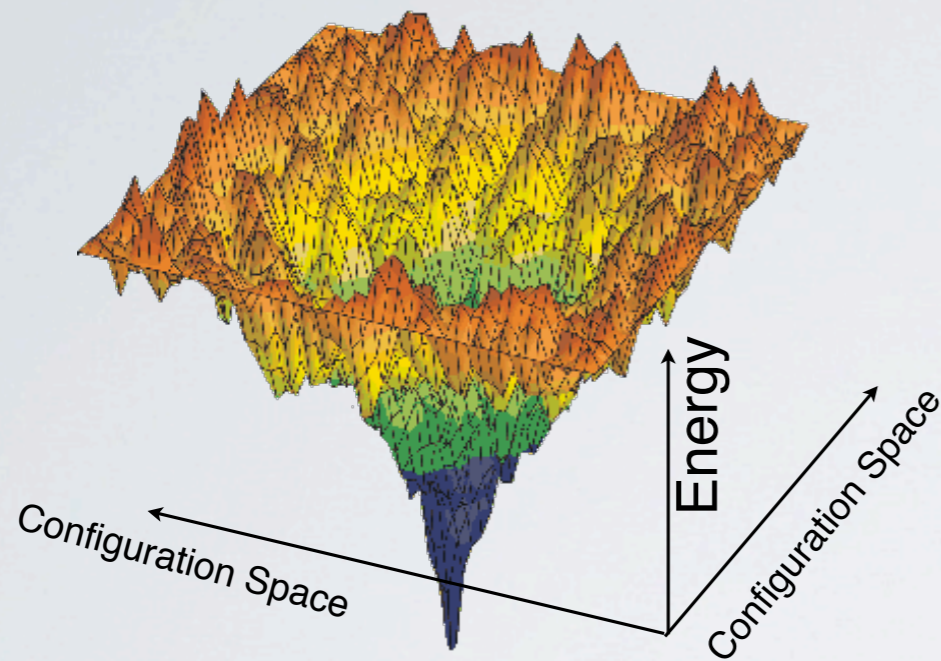
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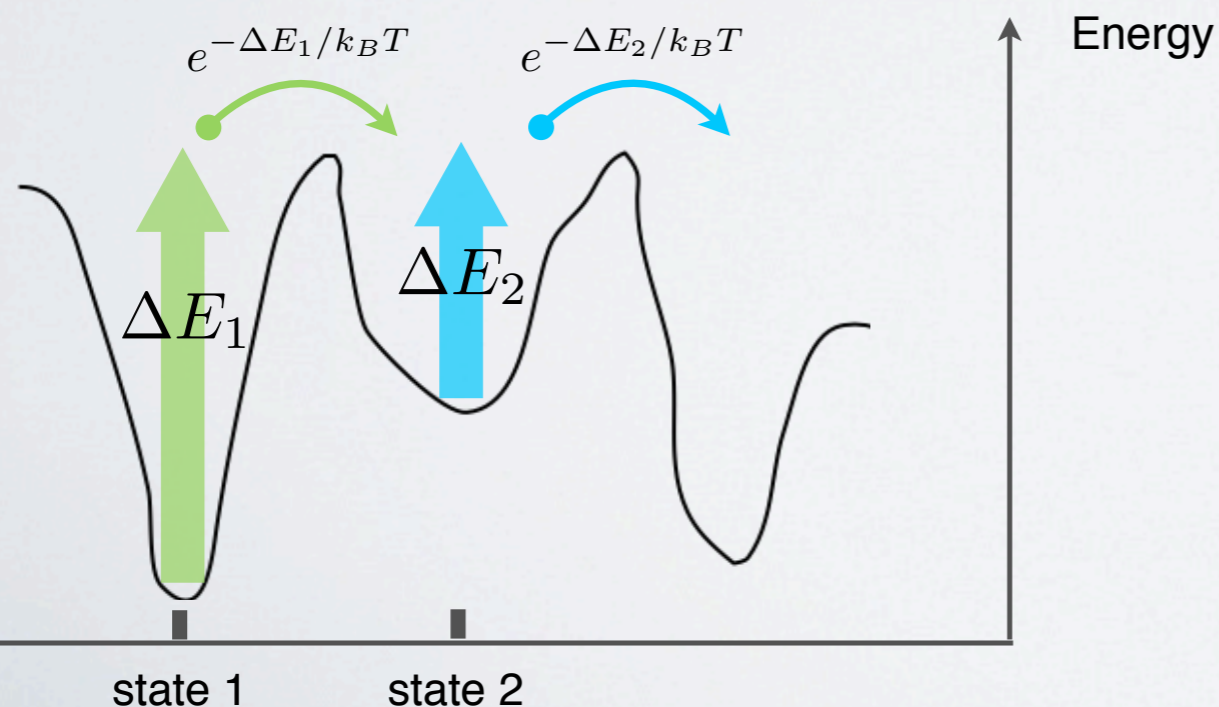
# Why might tissues generically be close to a glass transition?

- For wound healing, embryonic development, cancer invasion:
  - \* Initially need large scale flows (i.e. fluid-like rheology)
  - \* Subsequently need to support forces and shear stresses (i.e. a solid-like rheology)

# What is the microscopic origin of glassy behavior in non-active materials?



- \* The existence of a complex potential energy landscape
- \* System close to energy landscape surface
- \* System is trapped in metastable states, needs energy fluctuations to escape



## Trap Model for glassy dynamics

C. Monthus & J.-P Bouchaud  
J. Phys. A 29 3847 (1996)

## Soft Glassy Rheology

P. Sollich et al  
PRL 78 2020 (1997)

# Exploring the Potential Energy Landscape

**Sheared foam**

**Cells in tissues**

**Alexandre J Kabla**

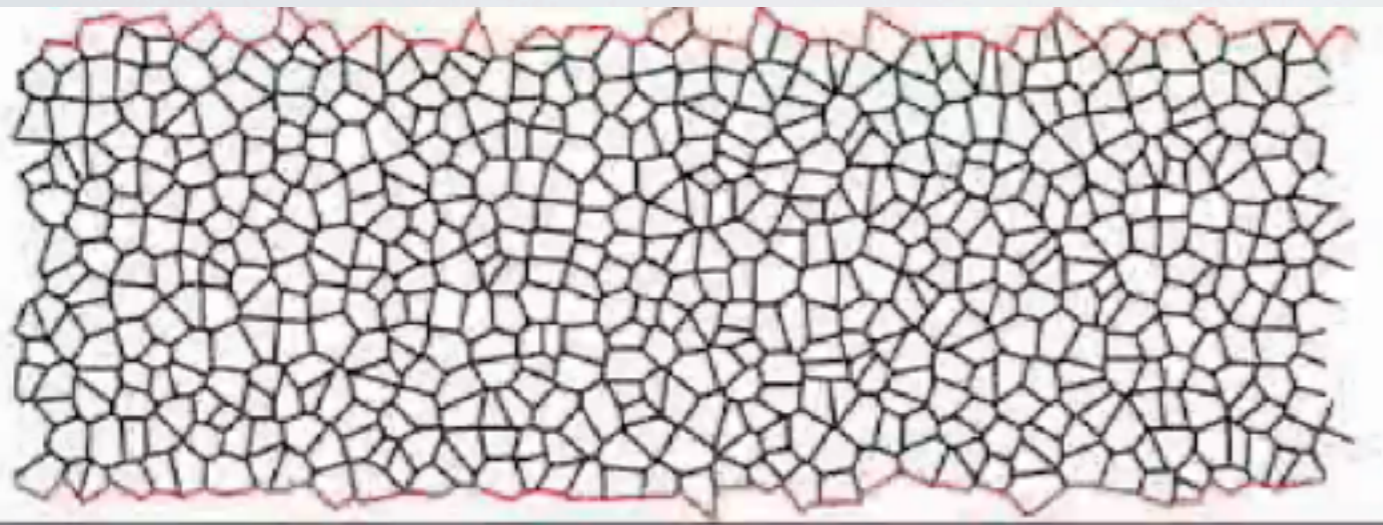
**Global injection of energy**

**Energy injection by cell**



# Exploring the Potential Energy Landscape

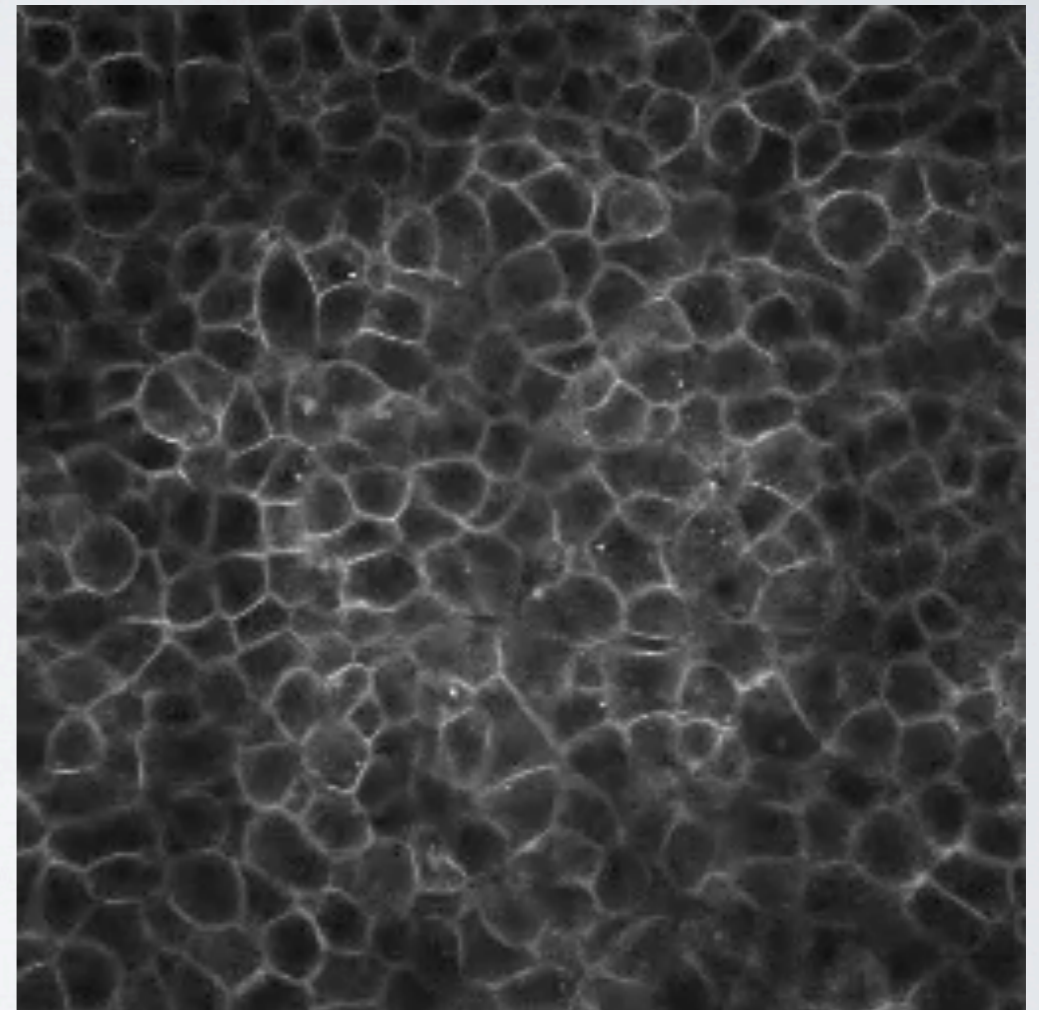
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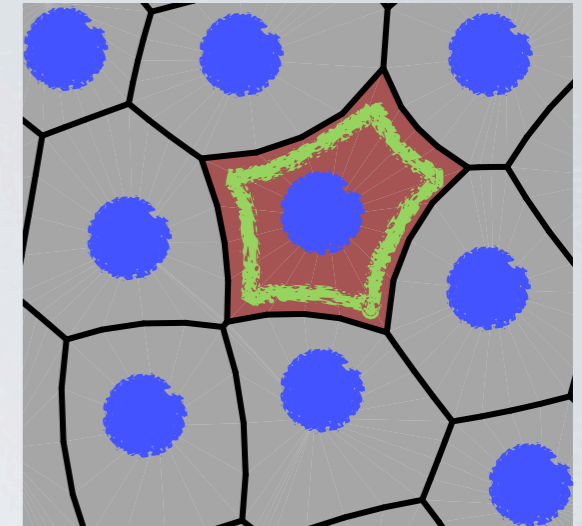
**Global injection of energy**

**Cells in tissues**



**Energy injection by cell**

# The Shape Equilibrium Model / Vertex Model



$$E_{cell} = k_A(A - A_0)^2 + k_P(P - P_0)^2$$
$$= k_A(A - A_0)^2 + k_P(P^2 - 2P_0P + P_0^2)$$

**Bulk elasticity term**

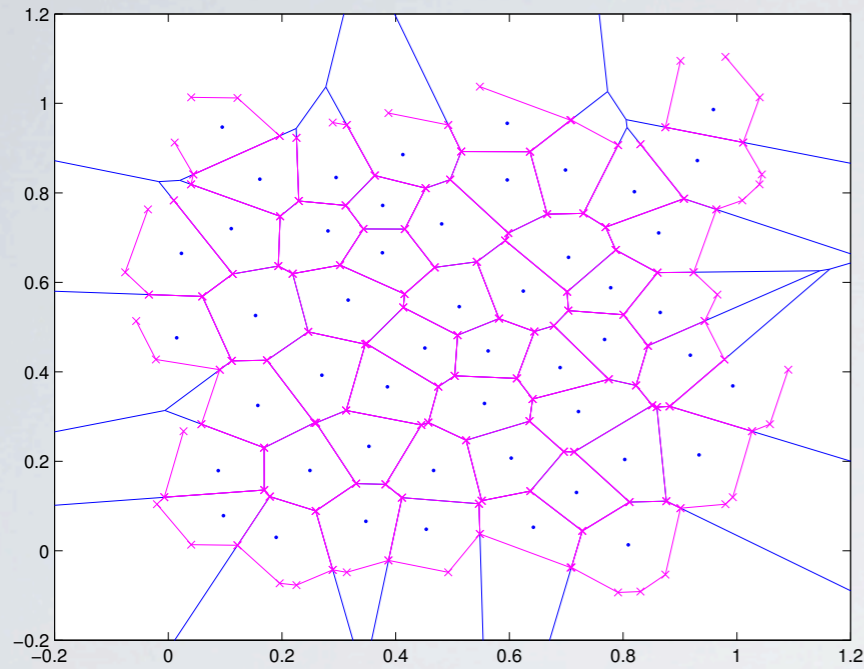
**Contractility**

**Line tension and Adhesion**

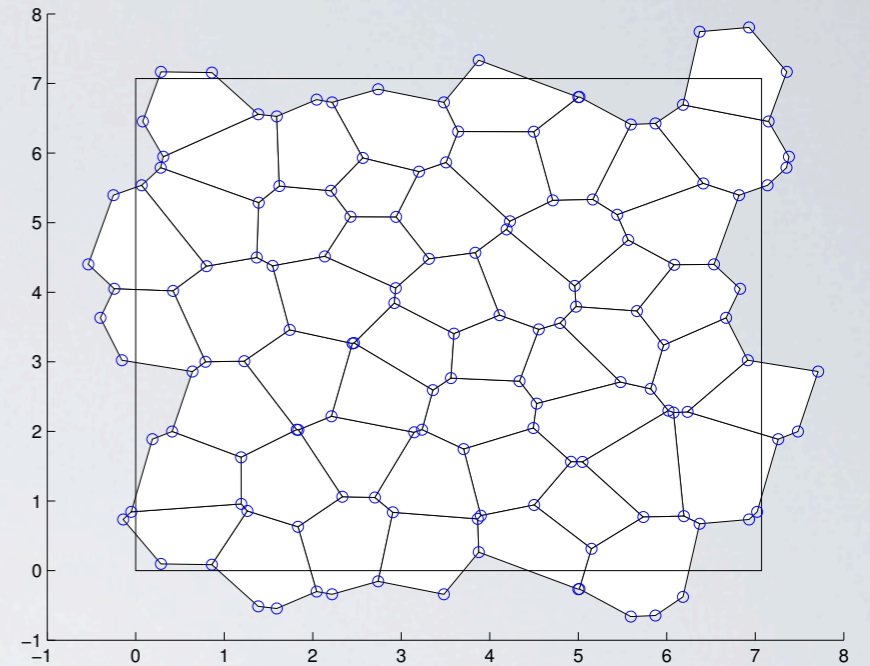
- Developed more than 10 years ago and well-studied
- Generalization of a foam model
- Reasonably good agreement with experimentally observed cell shapes

**Nagai & Honda Phil. Mag. B vol. 81 (7) (2001)**  
**Hufnagel et al, PNAS vol. 104 (10) pp. 3835 (2007)**  
**Farhadifar et al, Current Biology (2007)**  
**Jülicher et al Phys. Rep. (2007)**  
**Manning et al, PNAS (2010)**  
**Staple et al EPJE 33 (2) 117 (2010)**  
**Chiou et al PLOS Comp Bio 8 (5) e1002512 (2012)**

# Simulating a cellular structure



Voronoi tessellation of a random point pattern



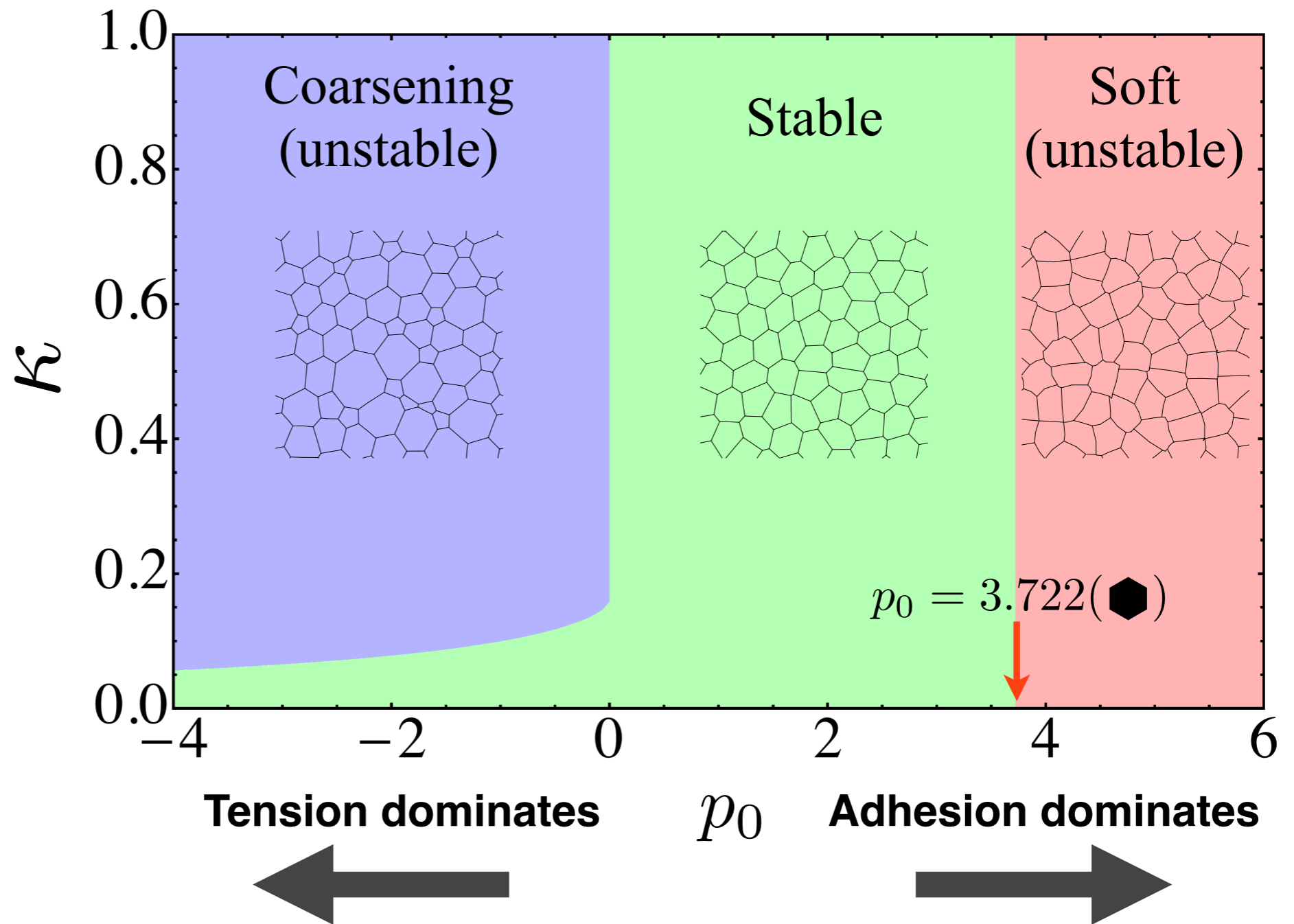
## \* Modeling tissues using **Surface Evolver**

- Using refined polygon tiling to represent 2-d confluent monolayer
- Configurations obtained by minimizing

$$E_{tot} = \sum_{cell} E_{cell}$$

# Phase space of disordered metastable states

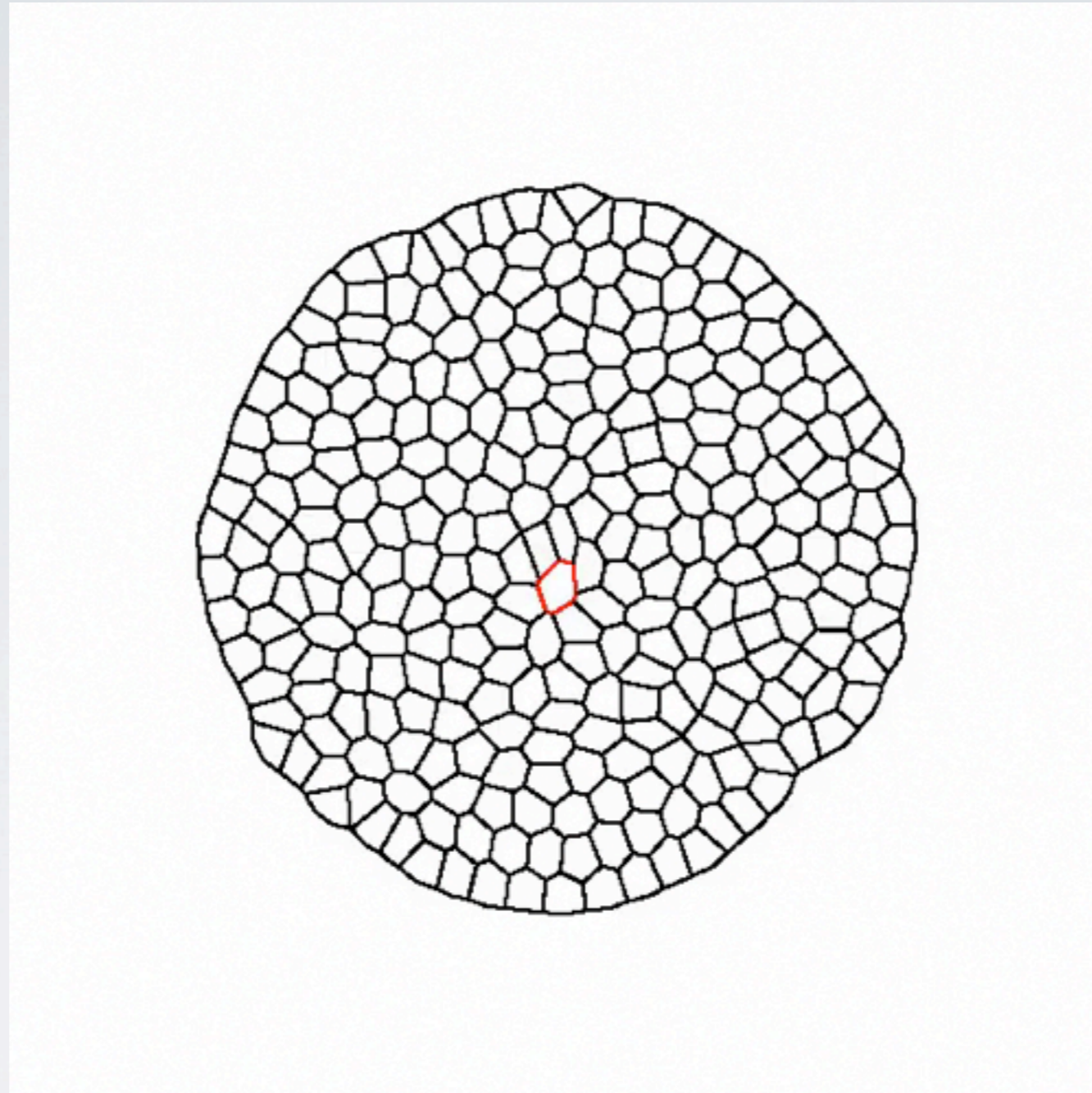
$$u_{tissue} = \sum_i u_i; \quad u_i = \kappa p_i^2 - 2\kappa p_0 p_i + (a_i - 1)^2$$



**Ground states** first studied by  
Staple et al, EPJ-E 33(2) 117 (2010)

D. Bi et al, unpublished

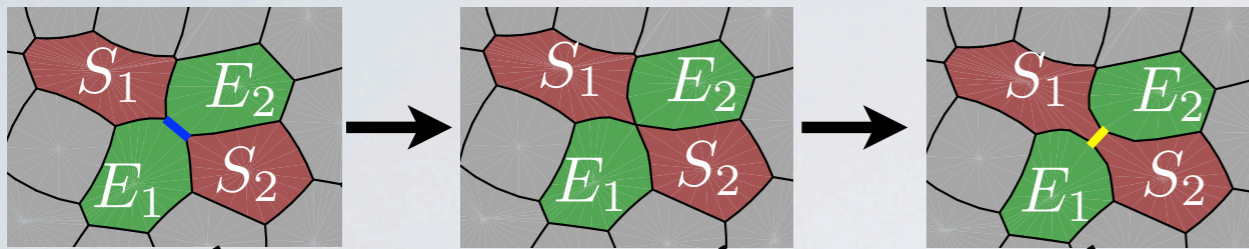
# Migration in 2D happens via T-1 transitions



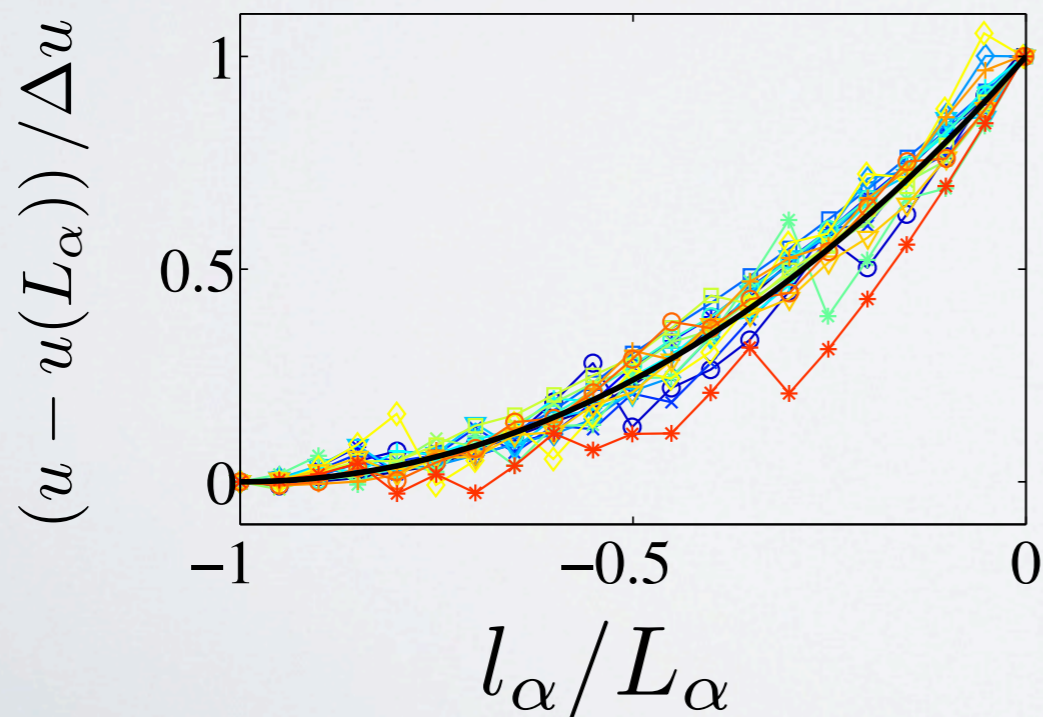
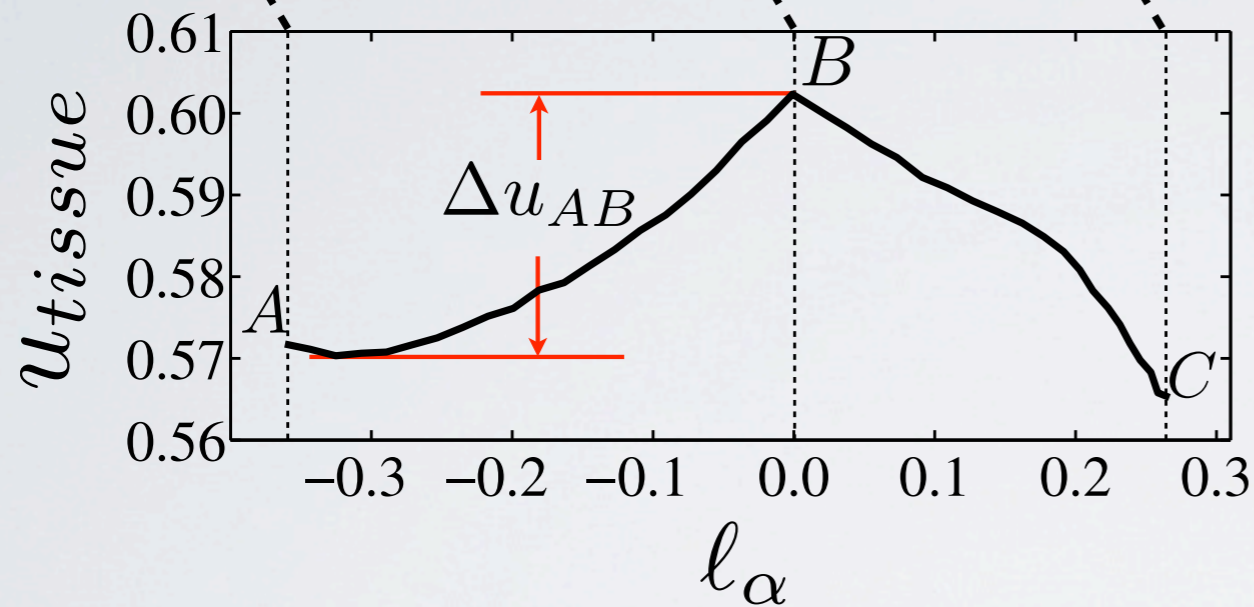
Cell divisions also and cell death (T2) also cause migrations, but it is not necessary for fluid behavior

**D. Bi et al, Soft Matter 2014**

# Energy trace for T-1 transitions

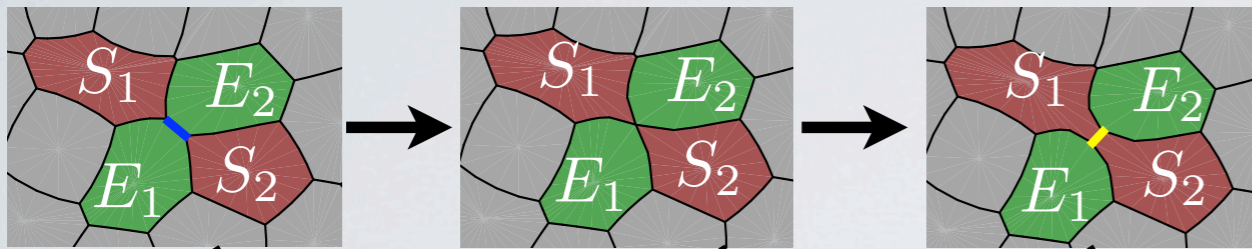


Testing all possible transition paths

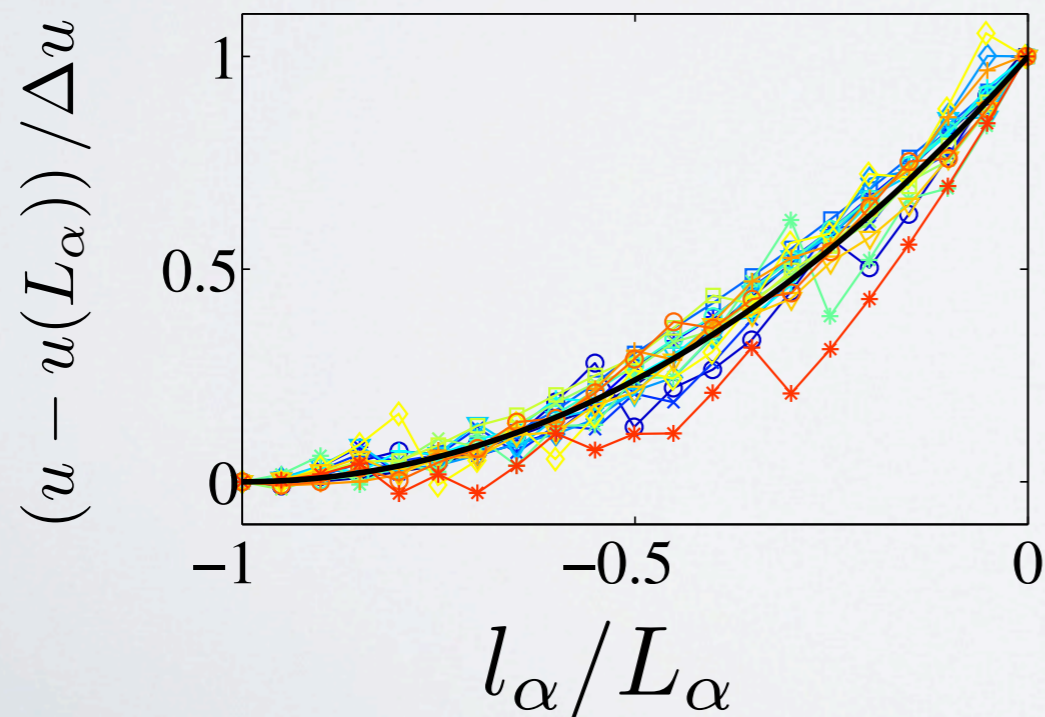
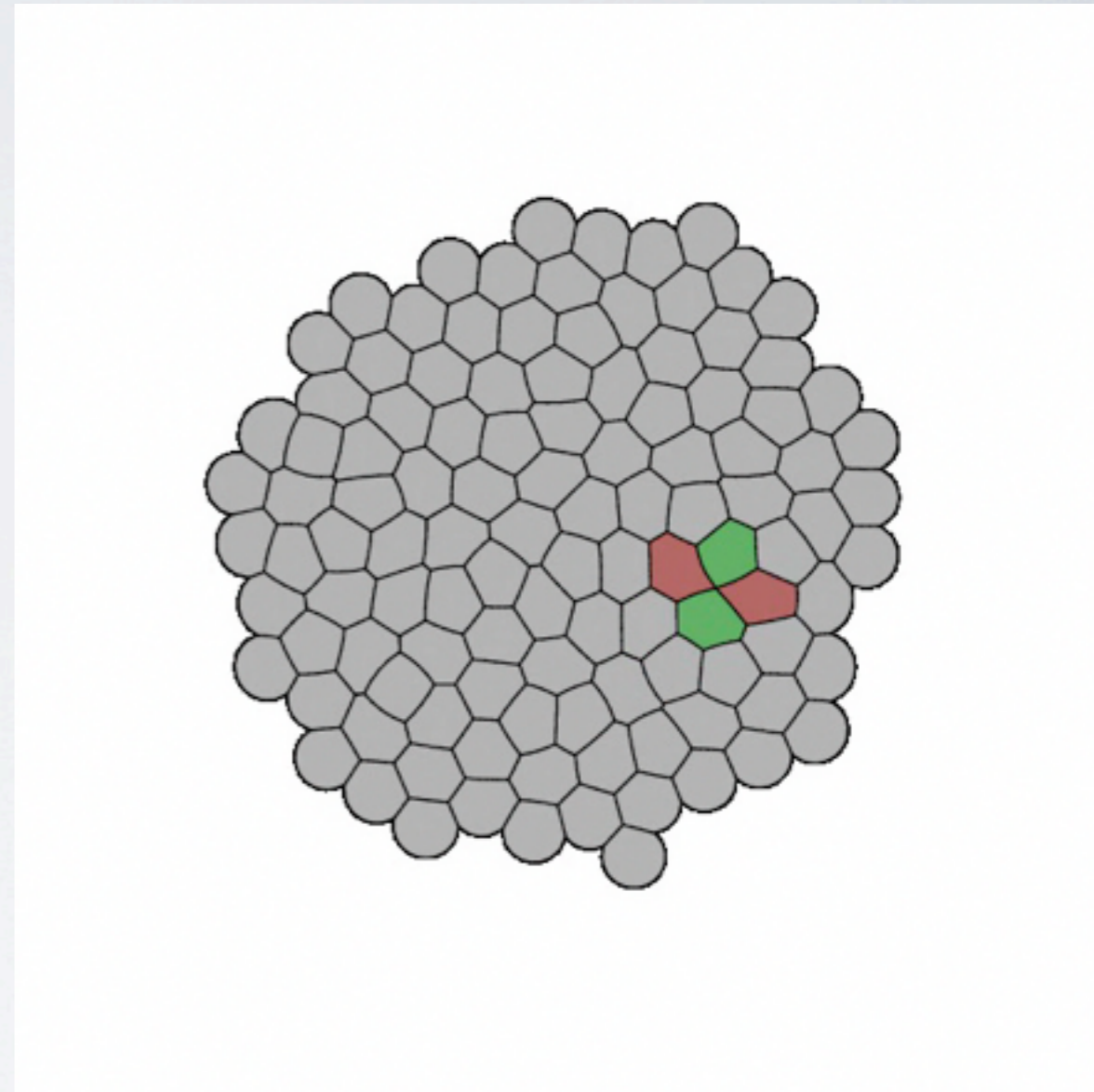
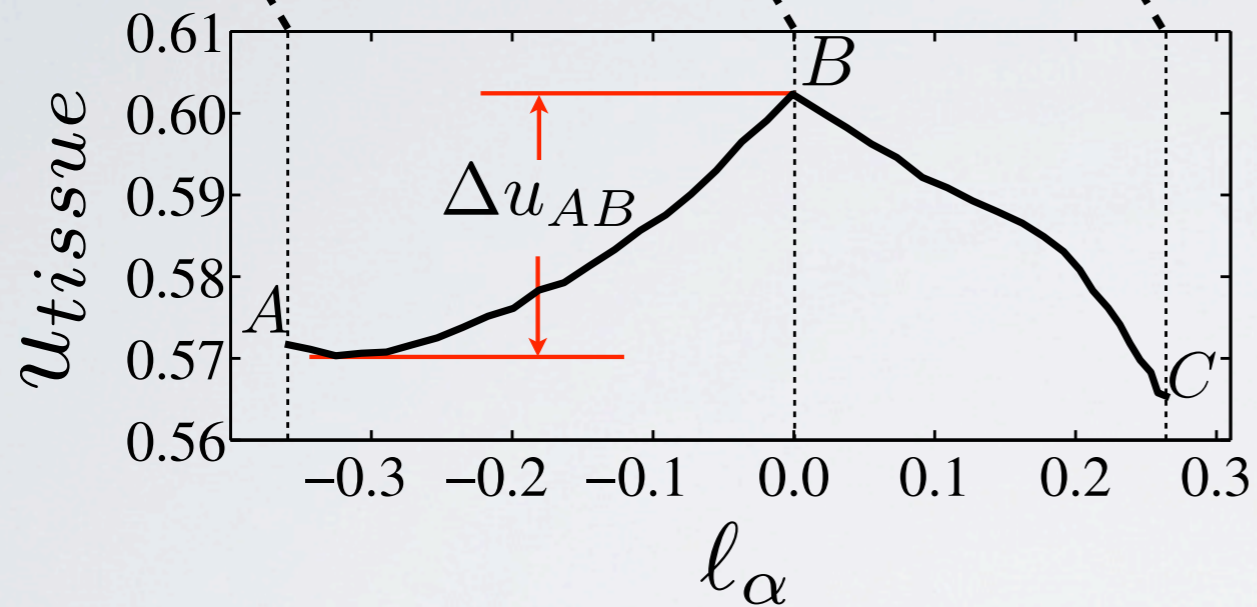


D. Bi et al, Soft Matter 2014

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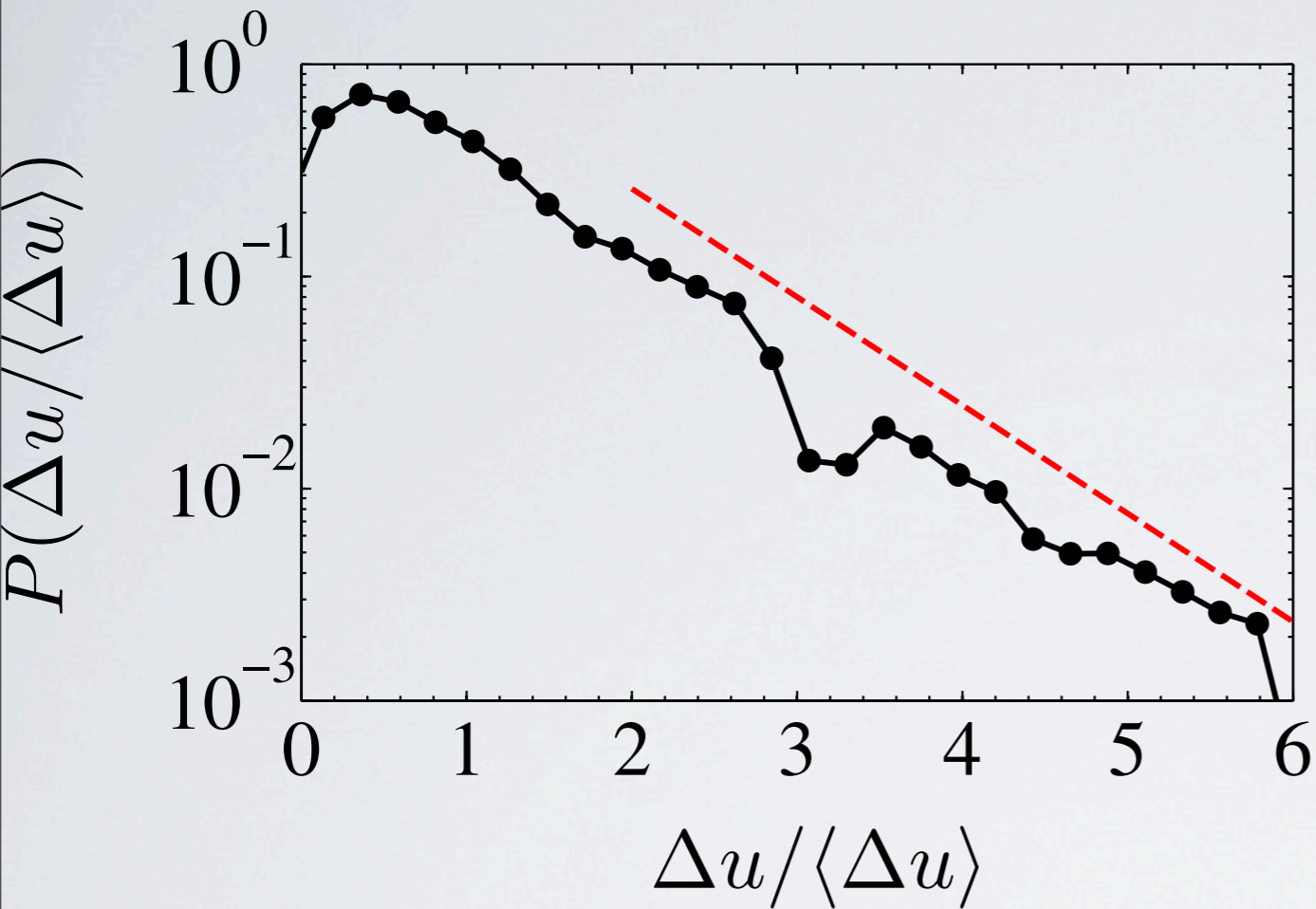
Testing all possible transition paths



D. Bi et al, Soft Matter 2014

# Energy barrier statistics

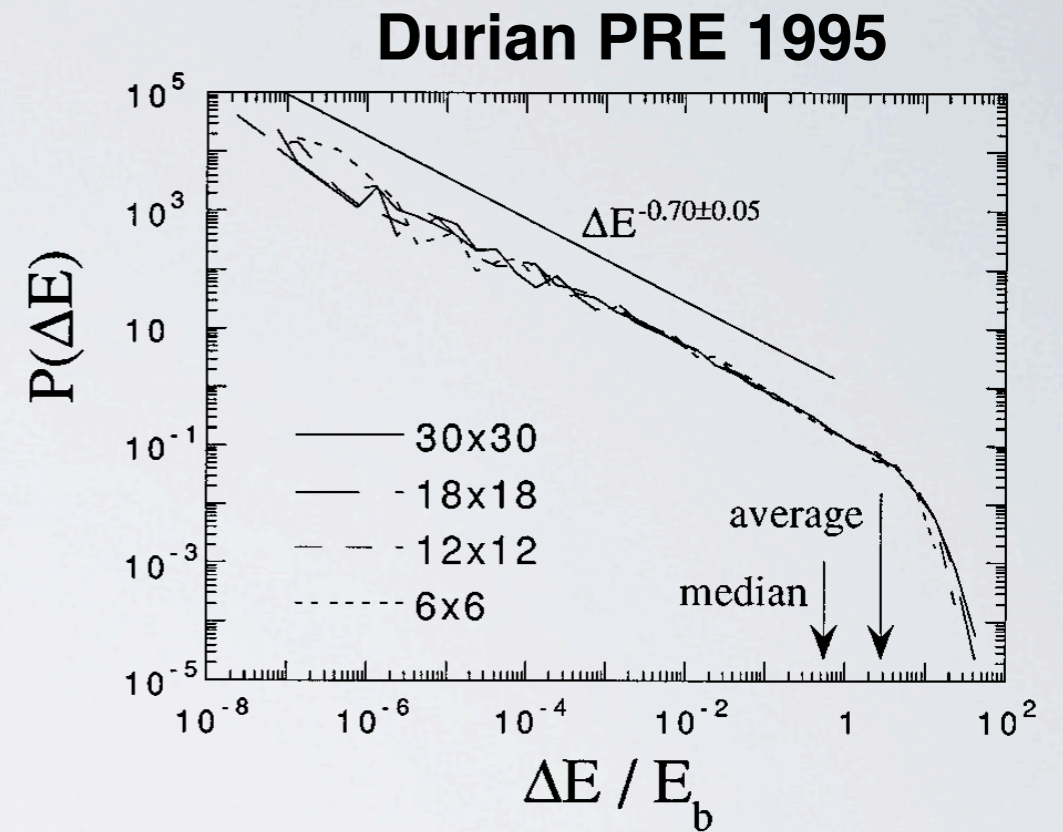
## Tissues



- \* Robust exponential tail
- \* Energy is injected locally, all sites are probed

D. Bi et al, *Soft Matter* 2014

## Sheared foams, grains



Also: Okuzono, Kawasaki and Ngai (1992)

- \* Power-law distributed, with exponential cut-off
- \* Energy is injected globally and failure occurs at a special soft spots in a material that is tuned near a critical point

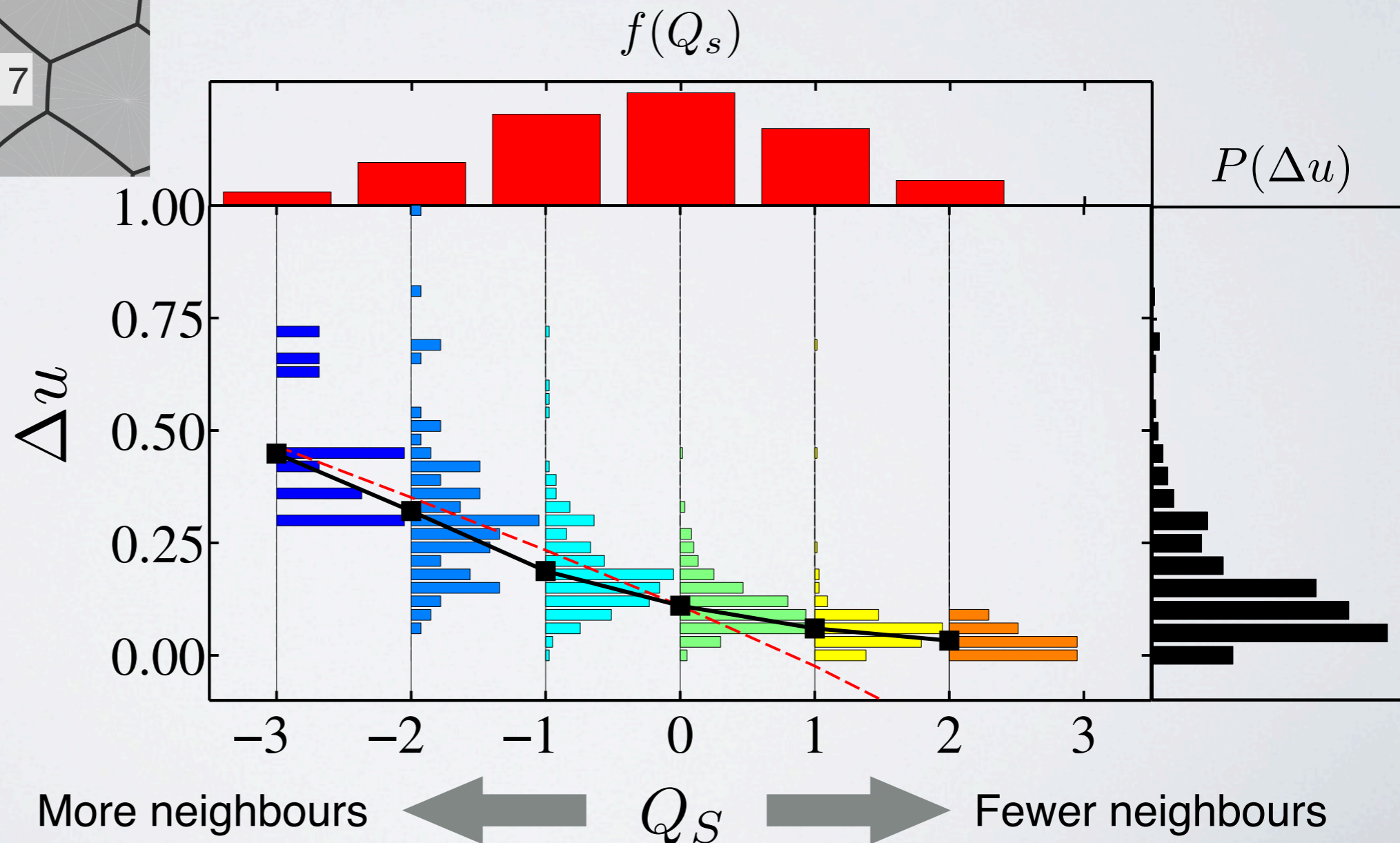
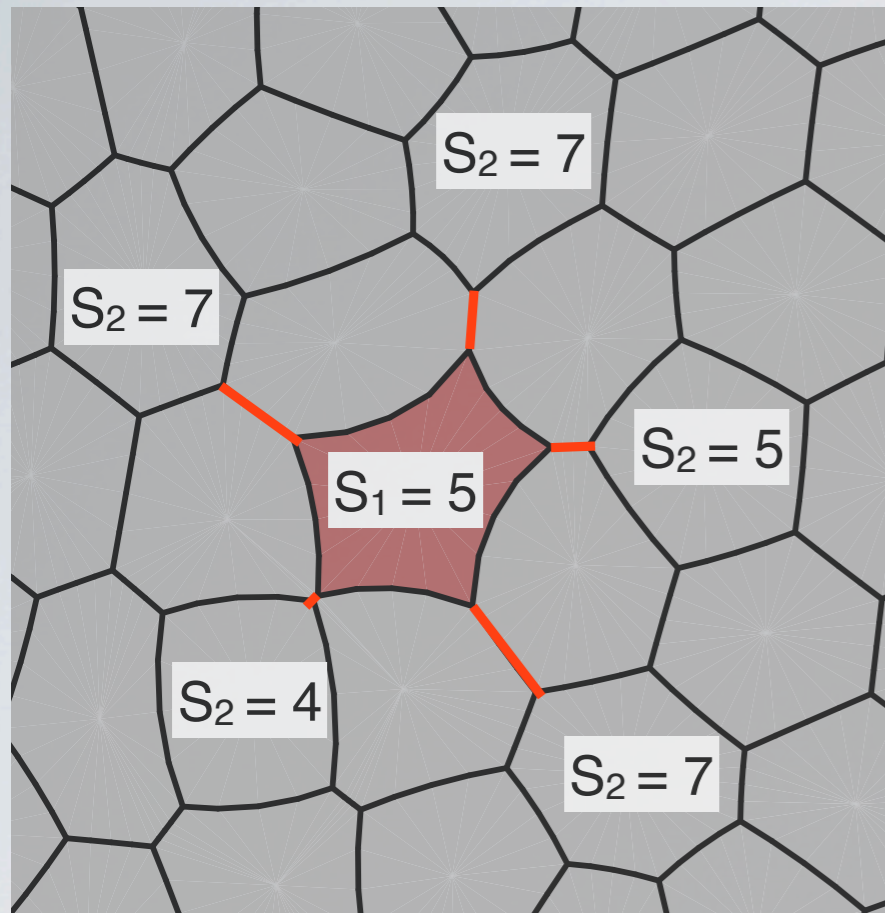


# Energy barrier depend on cell topology

$S_1$  and  $S_2$  are T-1 neighbours,  
Define a topological measure:

$$Q_S = (6 - S_1) + (6 - S_2)$$

D. Bi et al, Soft Matter 2014



# From energy barriers to cell migration

## A very minimal model:

- \* A cell is 'caged' by its neighbors
- \* Can escape cage by:

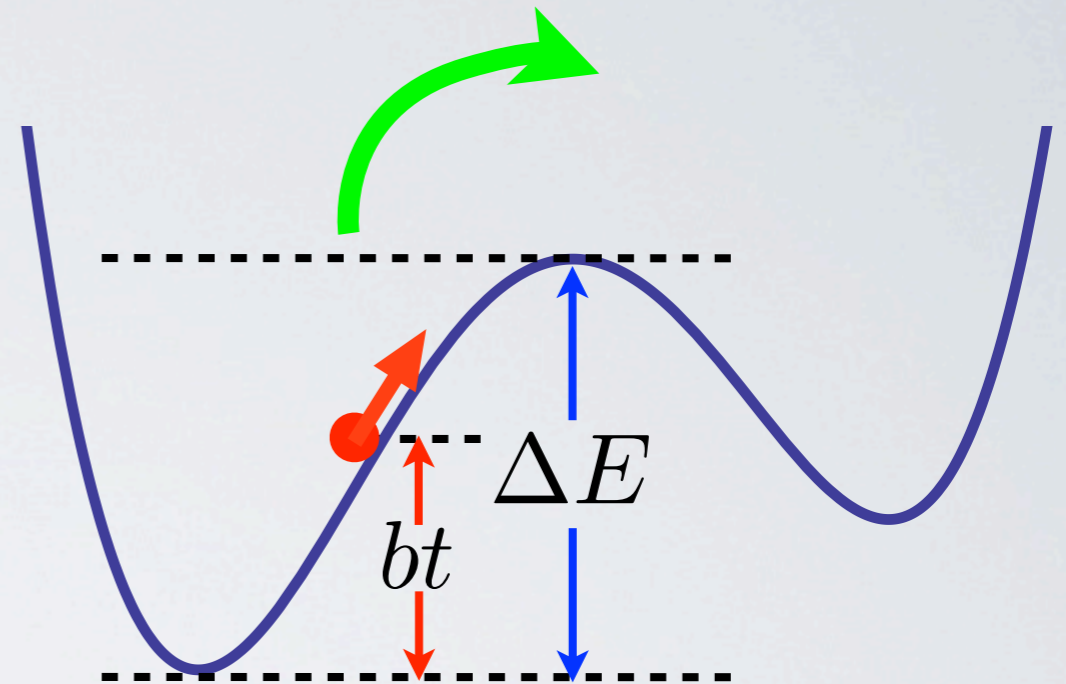
✓ Polarized cell motion **b**

+

✓ Active shape fluctuations  **$\epsilon$**

- \* Energy landscape has shape given by  $\rho(\Delta E) = e^{-\Delta E/\epsilon_0}$

$$R = \omega_0 \exp [ -(\Delta E - bt) / \epsilon ]$$



Master equation for dynamics



$$\frac{\partial}{\partial t} P(\Delta E, t) = -\omega_0 e^{-[\Delta E - bt]/\epsilon} P(\Delta E, t)$$

$$+ \varrho(\Delta E) \int d\Delta E' \omega_0 e^{-[\Delta E' - bt]/\epsilon} P(\Delta E', t)$$

**These are sufficient to lead to glassy dynamics!**

**Trap model for glassy dynamics**

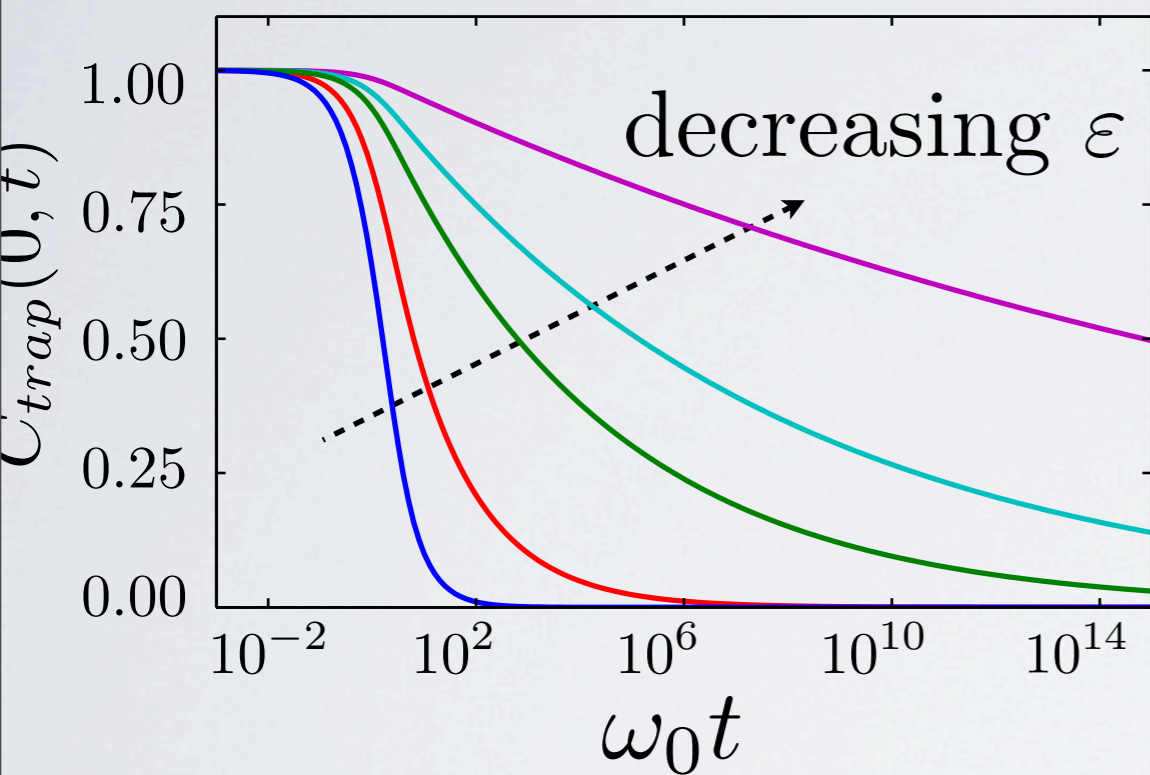
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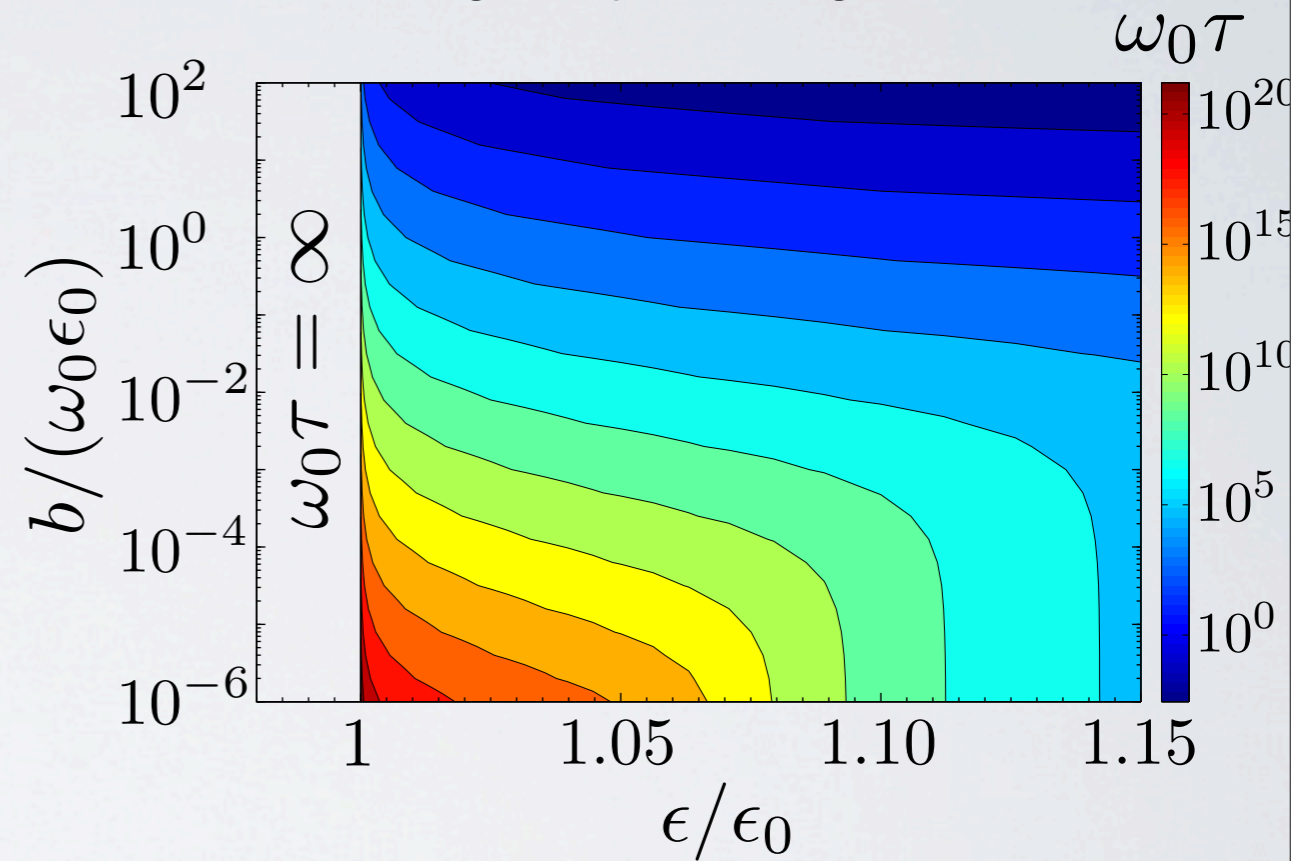
# From energy barriers to cell migration

Probability to remain in a 'cage'



↑  
More activity

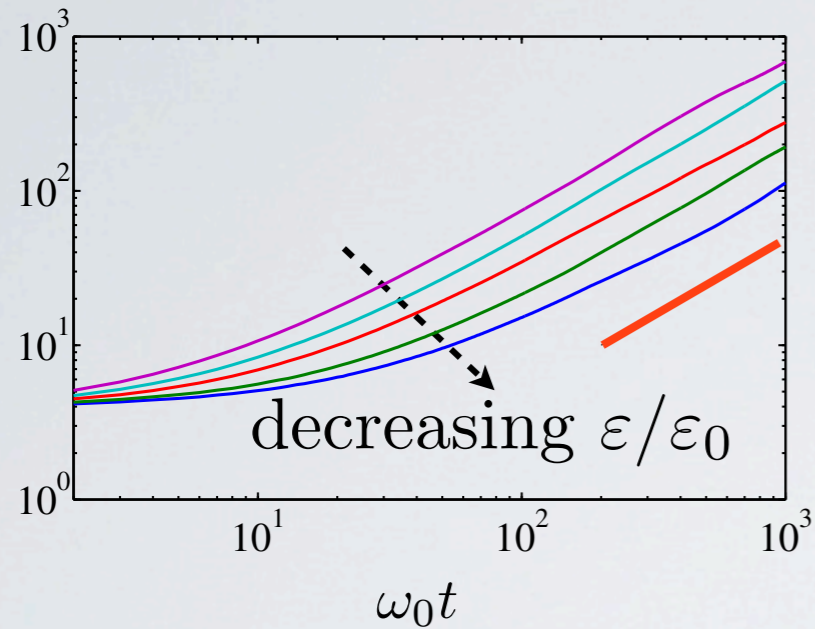
A glassy waiting time



$$\epsilon/\epsilon_0 = \frac{\text{shape fluctuation energy scale}}{\text{characteristic energy scale landscape}}$$

# From energy barriers to cell migration: experimentally accessible predictions

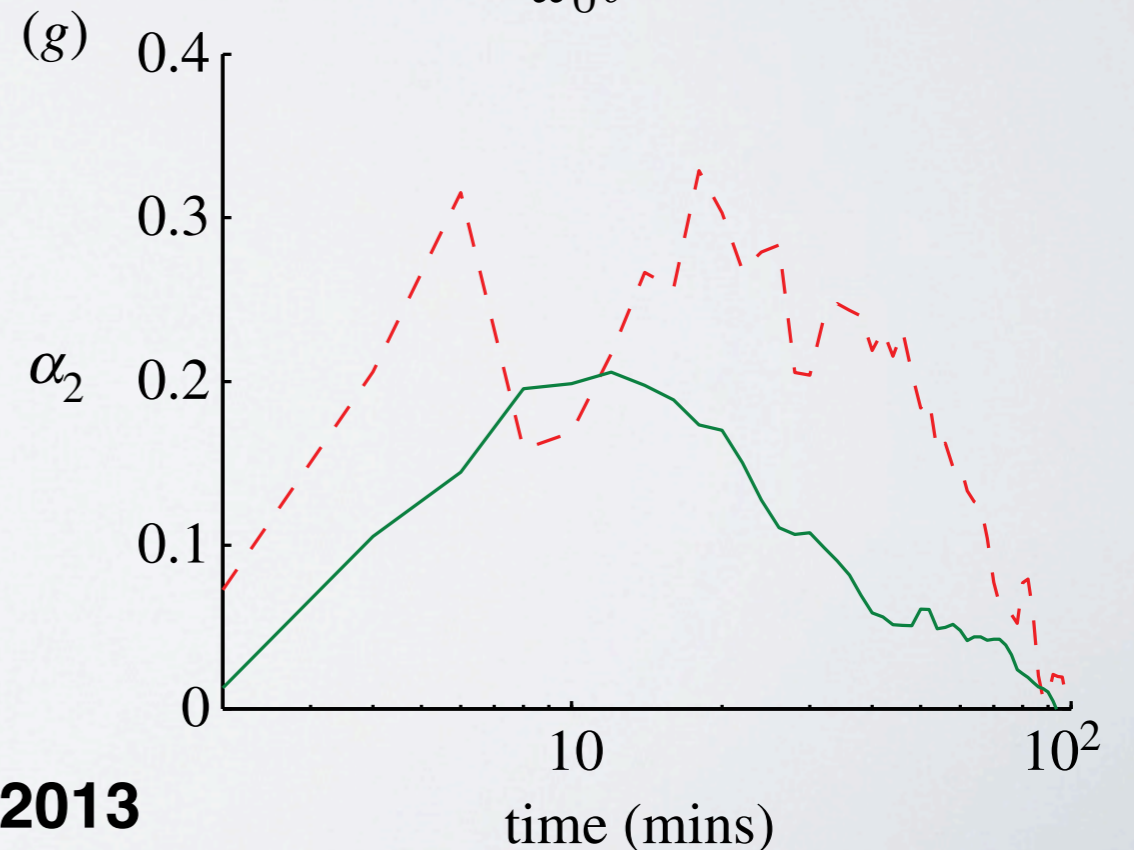
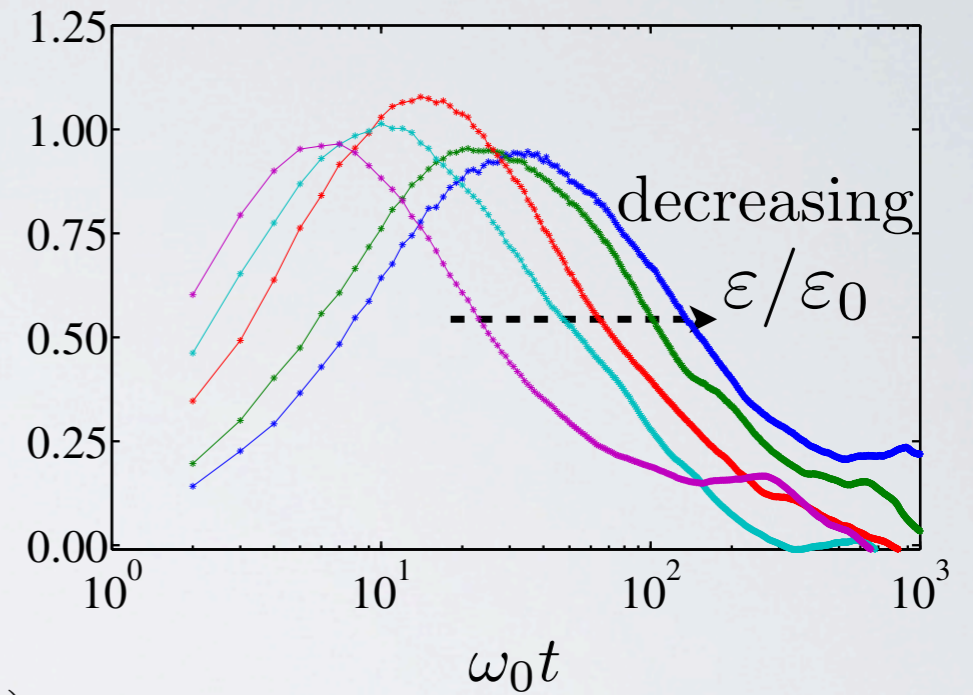
mean squared displacement



$\alpha_2$ : non-gaussian parameter

$\alpha_2 = 0$ : Brownian motion

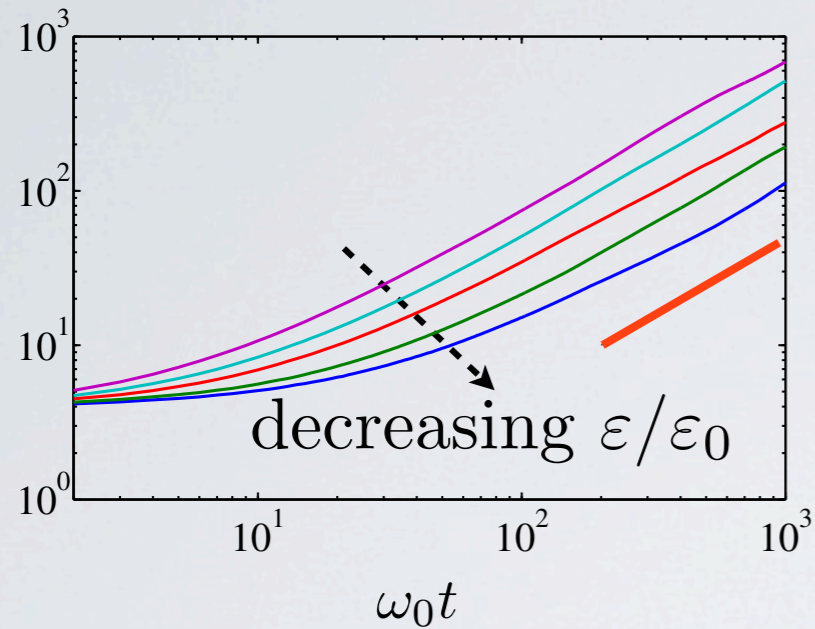
$\alpha_2 \sim 1$ : Cage breaking



Schöetz et al  
J. R. Soc. Interface 2013

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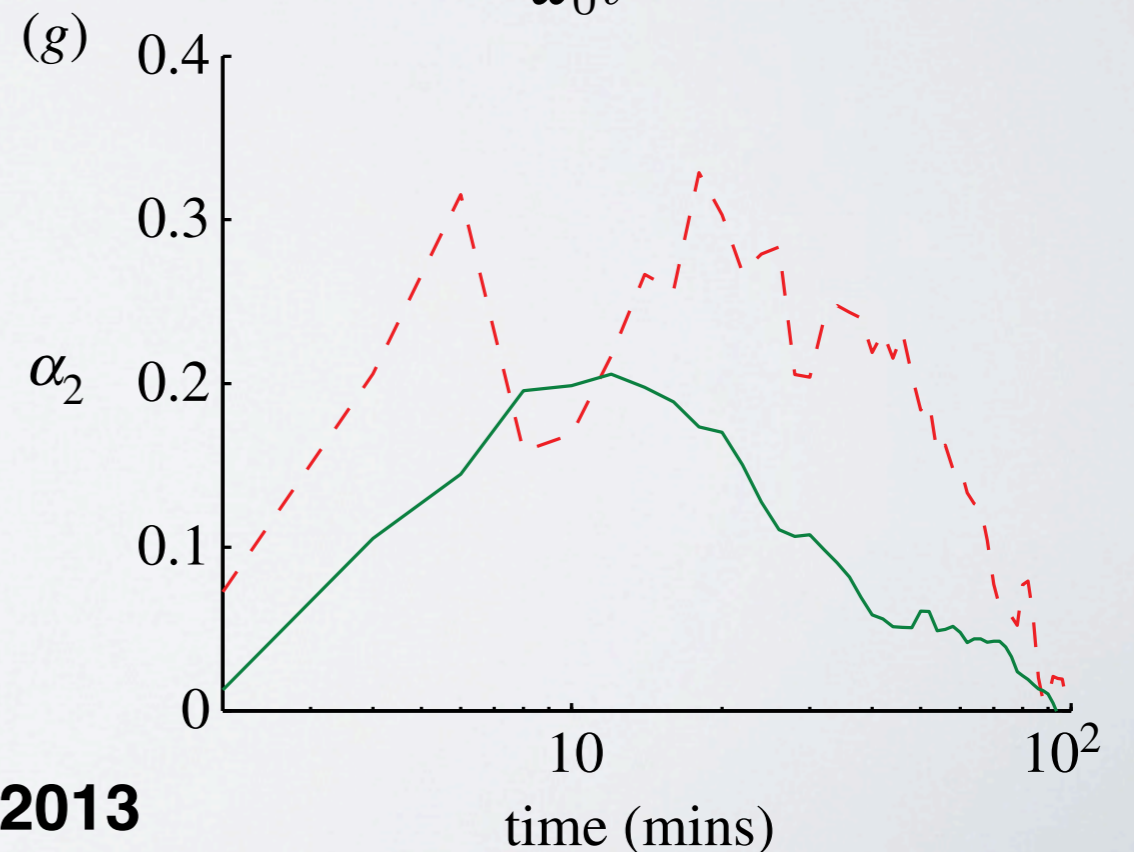
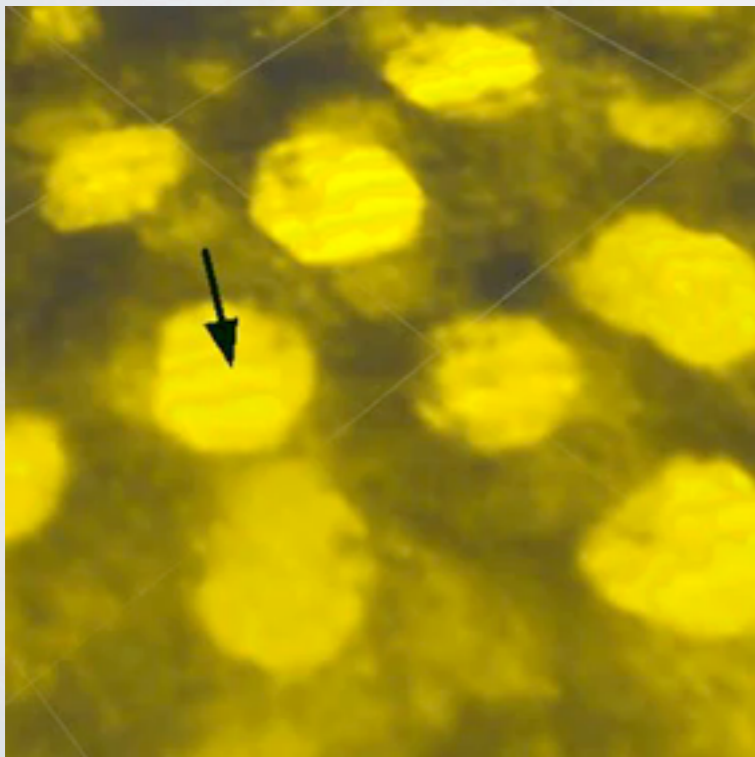
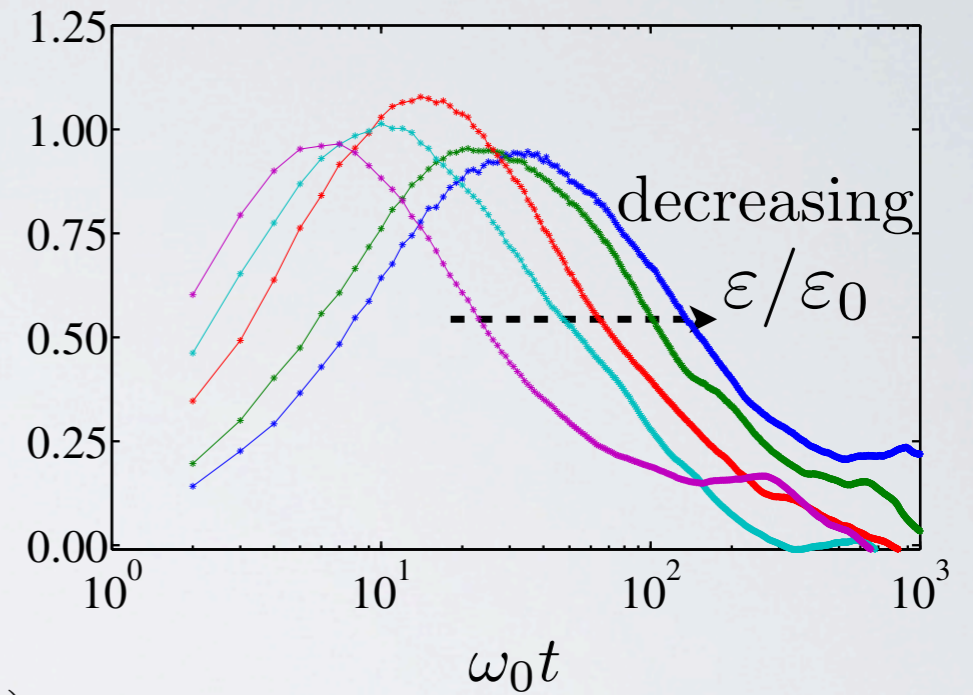
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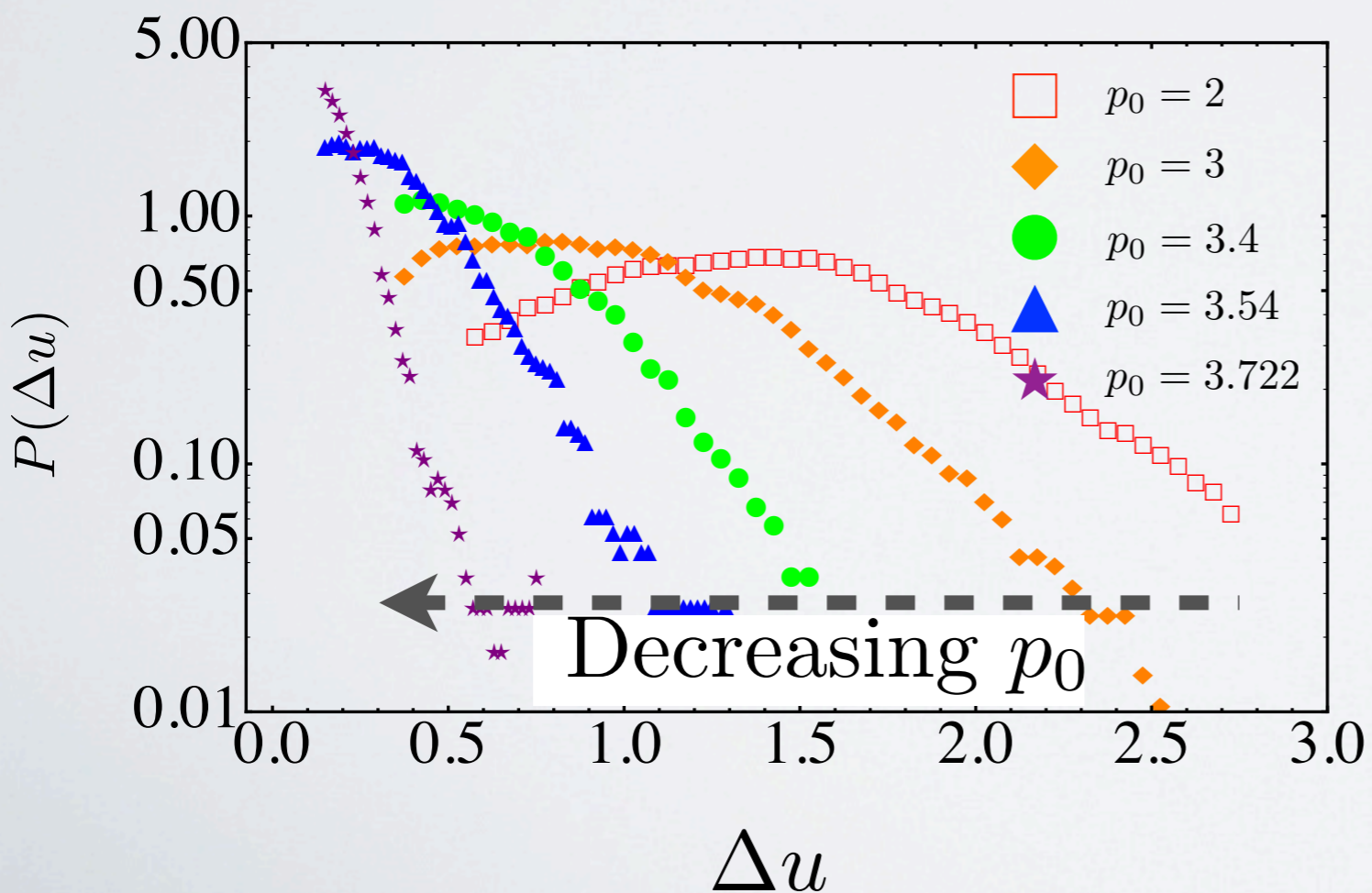
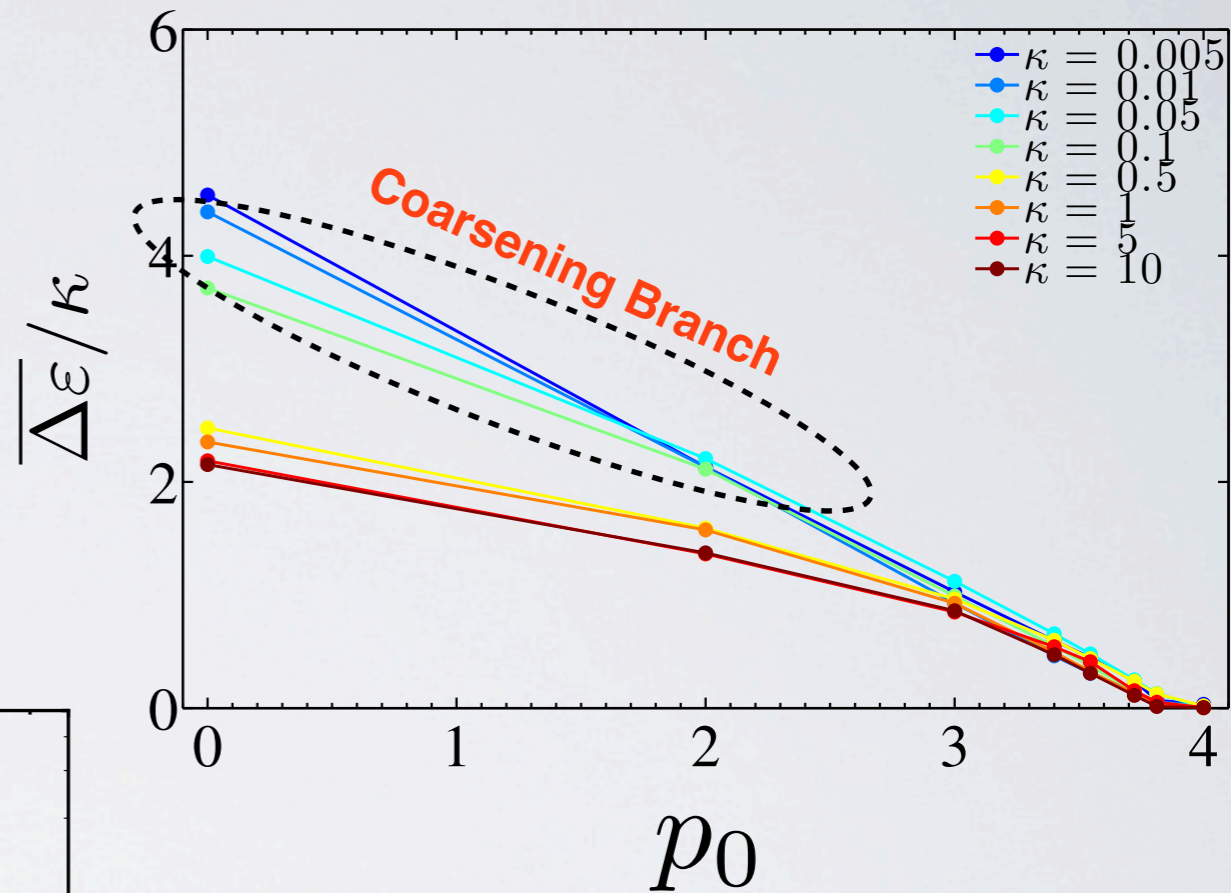
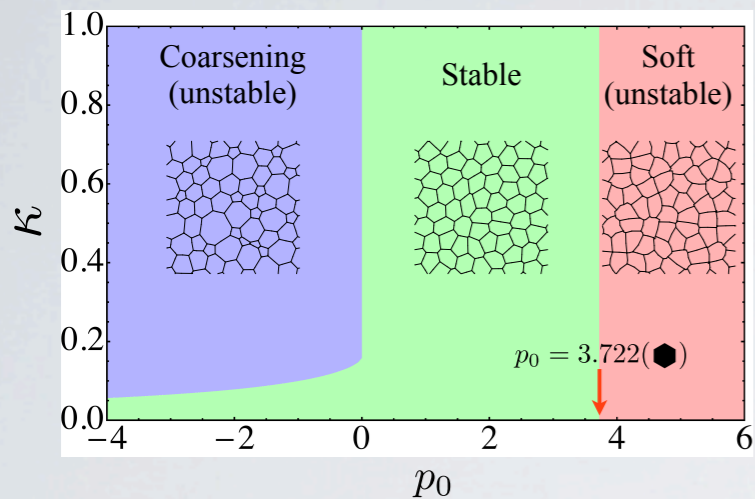
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Schöetz et al  
J. R. Soc. Interface 2013

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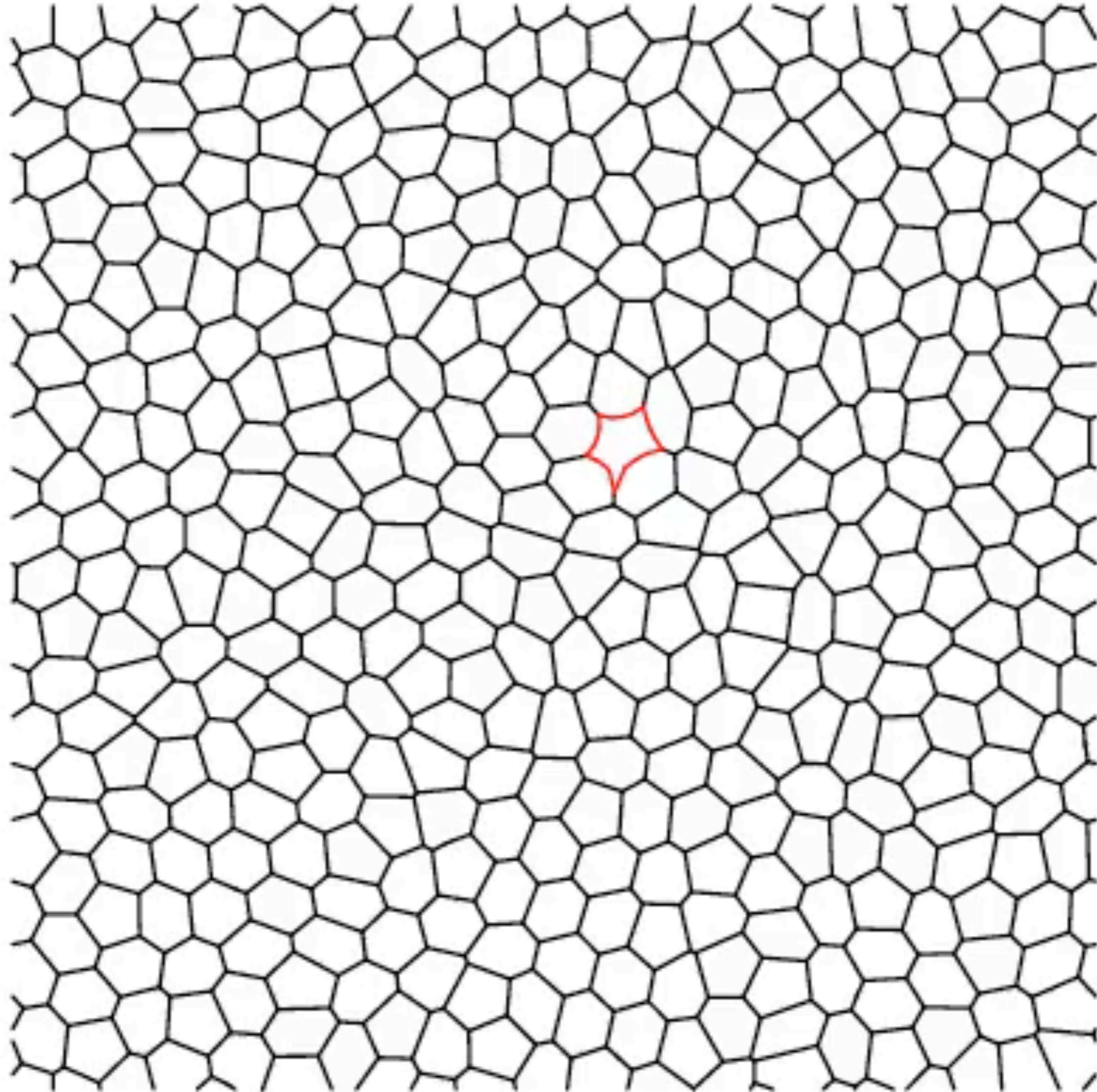
D. Bi et al, unpublished

**Work in progress**

**Migration of 'abnormal' cells embedded in healthy cells**

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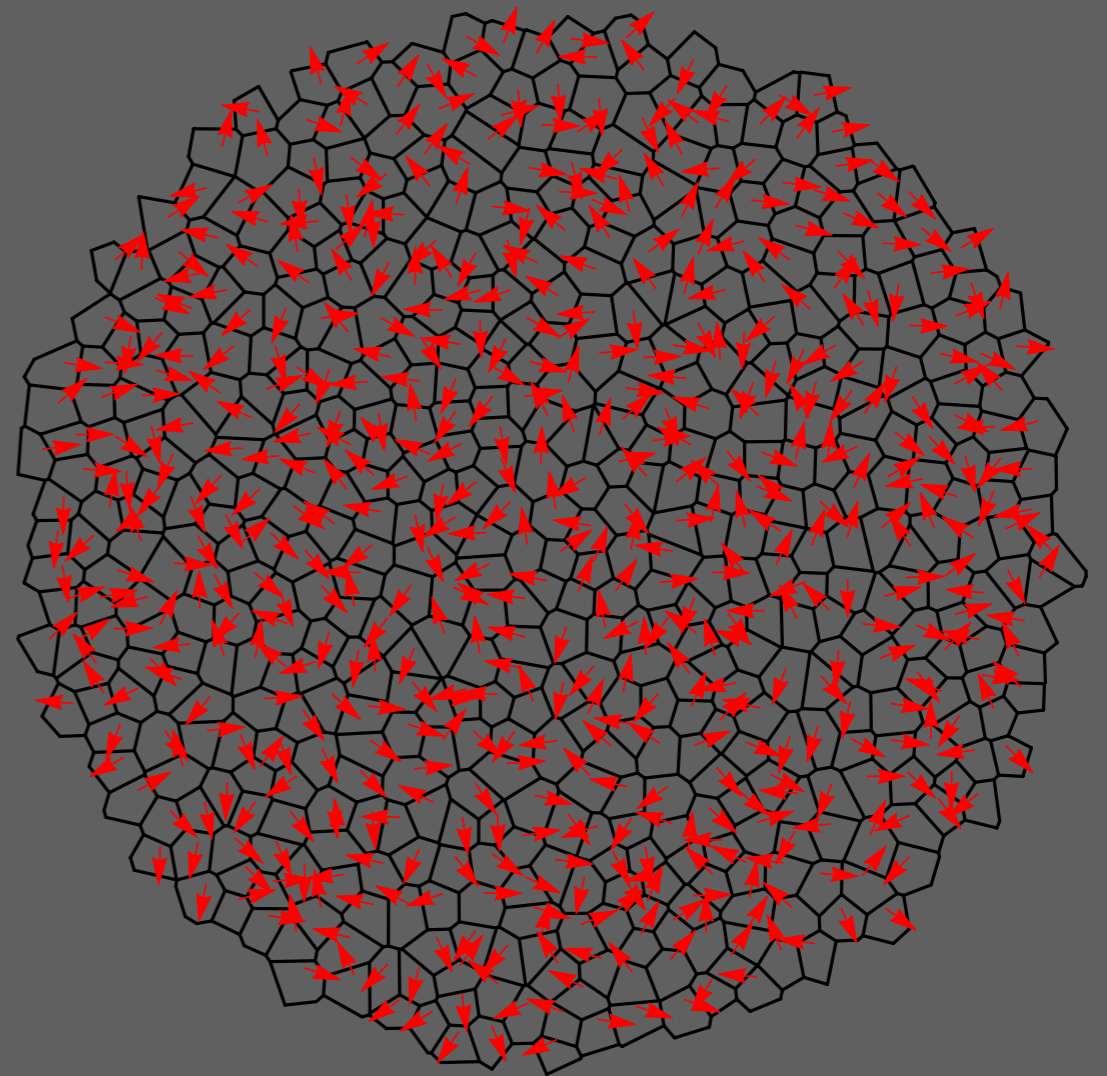
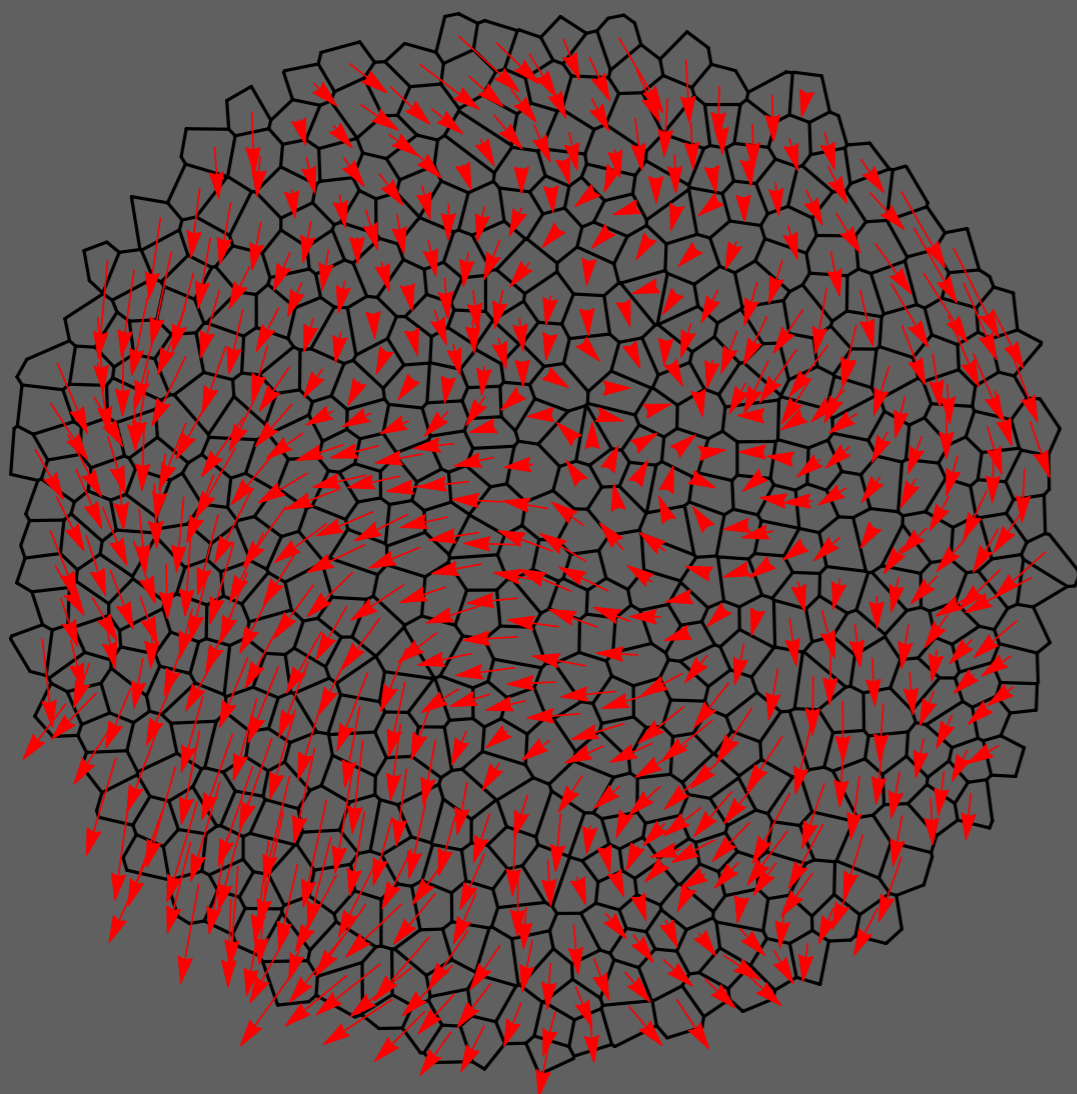
**Migration of 'abnormal' cells embedded in healthy cells**





Work in progress

# Dynamical simulation with active T-1 rearrangements



# Conclusion

- Many tissue types show evidence of being close to a glass transition (wound healing, embryonic development)
- For confluent (density  $\sim 1$ ) tissues, we have developed a framework for estimating energy barriers to cell rearrangements,
  - Self propelled particle models are perhaps less useful than shape equilibrium models at this density
  - energy barriers are exponentially distributed
  - energy barriers depend strongly on a cell's number of nearest neighbors (experimentally accessible predictions)
- Our model predicts how changing single-cell properties changes cell migration rates. Can apply cell sorting or cancer migration?

**D. Bi et al, Soft Matter 2014,10, 1885-1890**