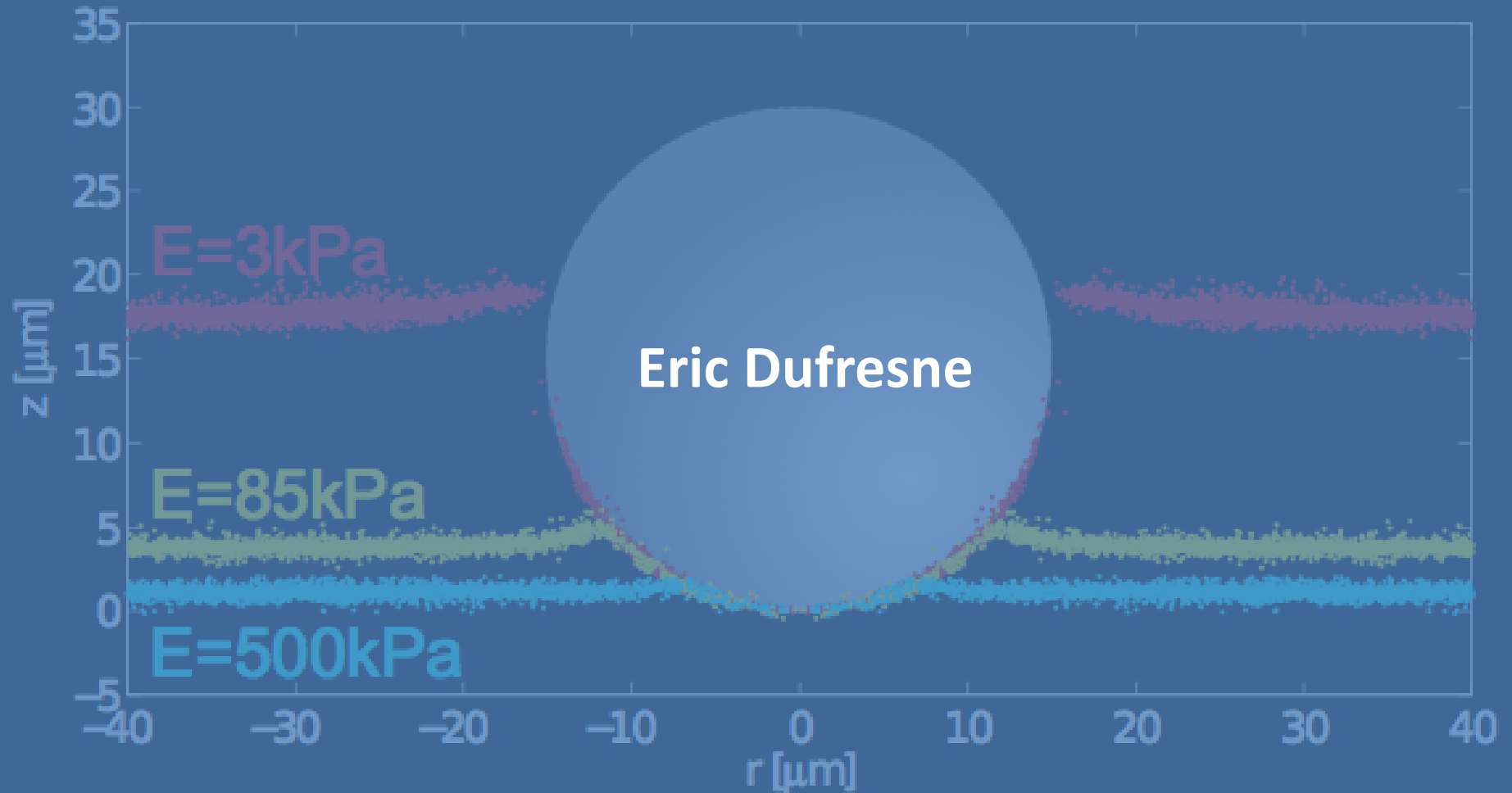
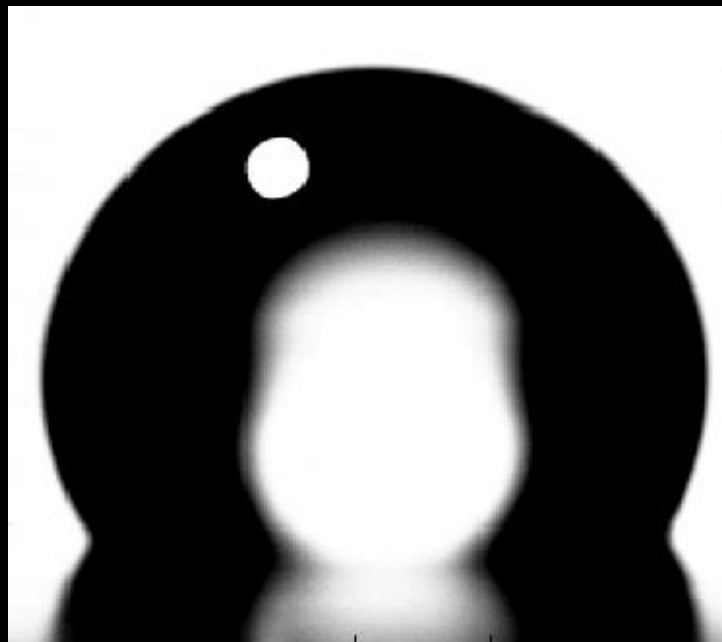


Wetting & Adhesion on Soft Surfaces

“Young’s Law is dead...long live Young’s Law”

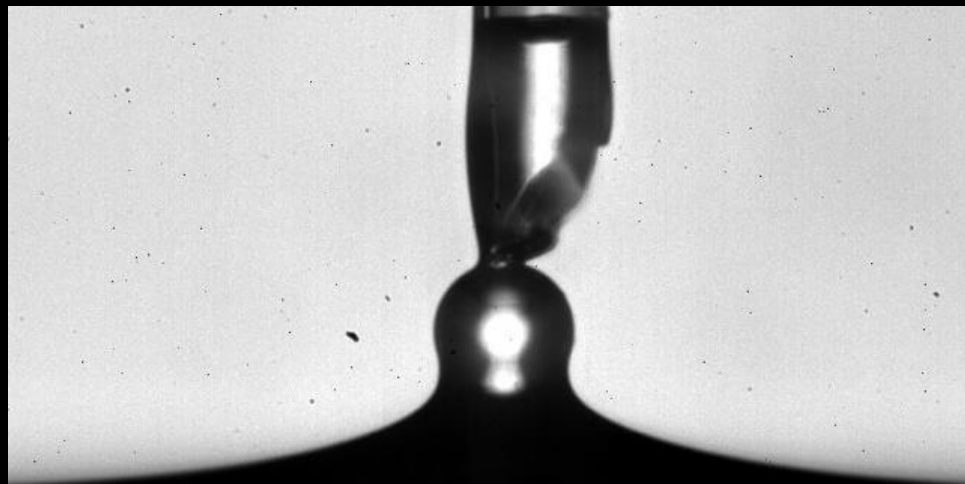


Wetting



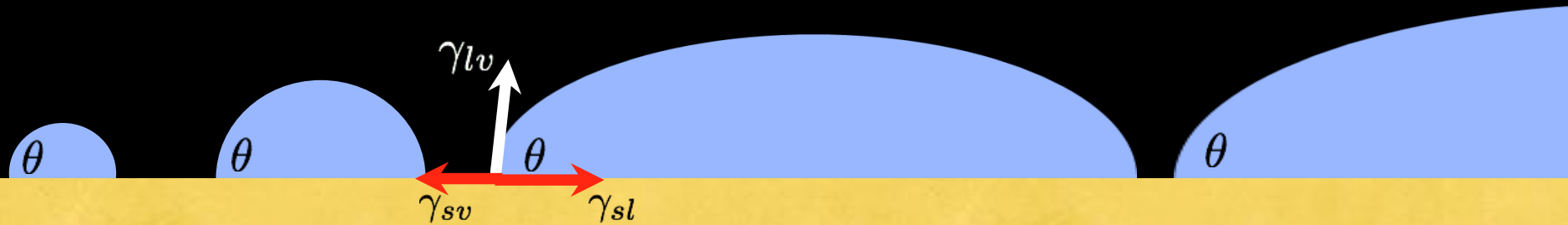
3mm

Adhesion



30 μm

Young's Law relates contact line geometry and material properties in equilibrium



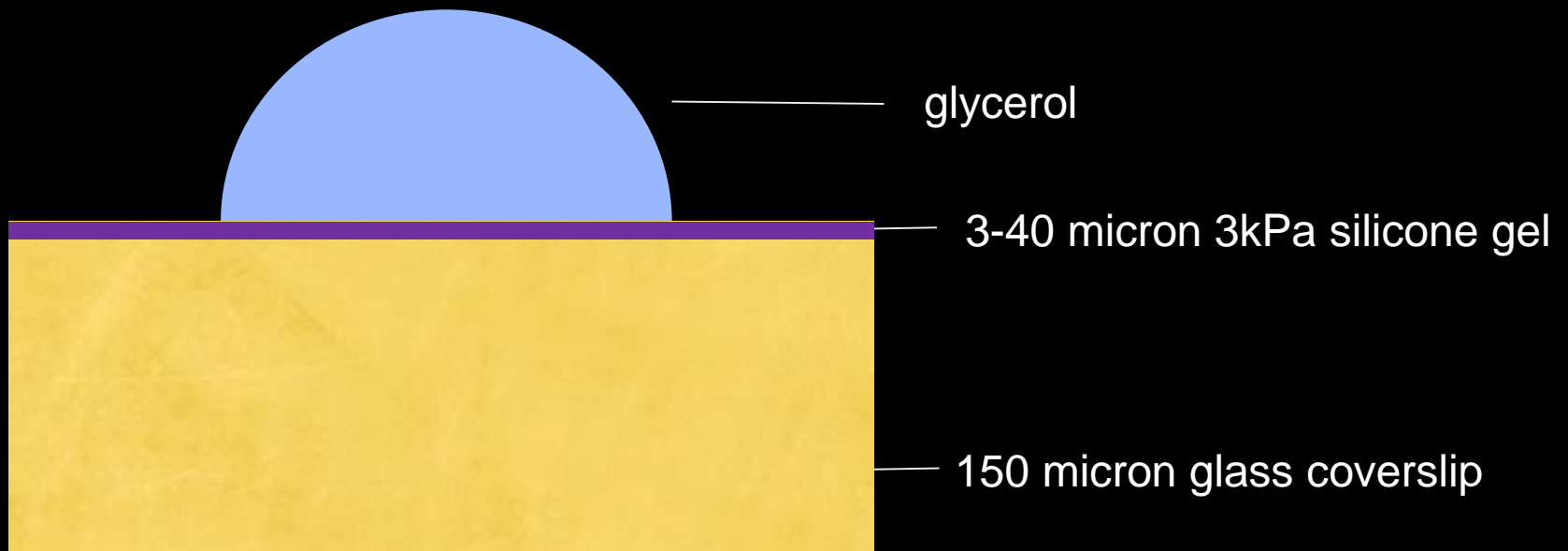
Thomas Young
1773-1829

$$\cos \theta = \frac{\gamma_{sv} - \gamma_{sl}}{\gamma_{lv}}$$

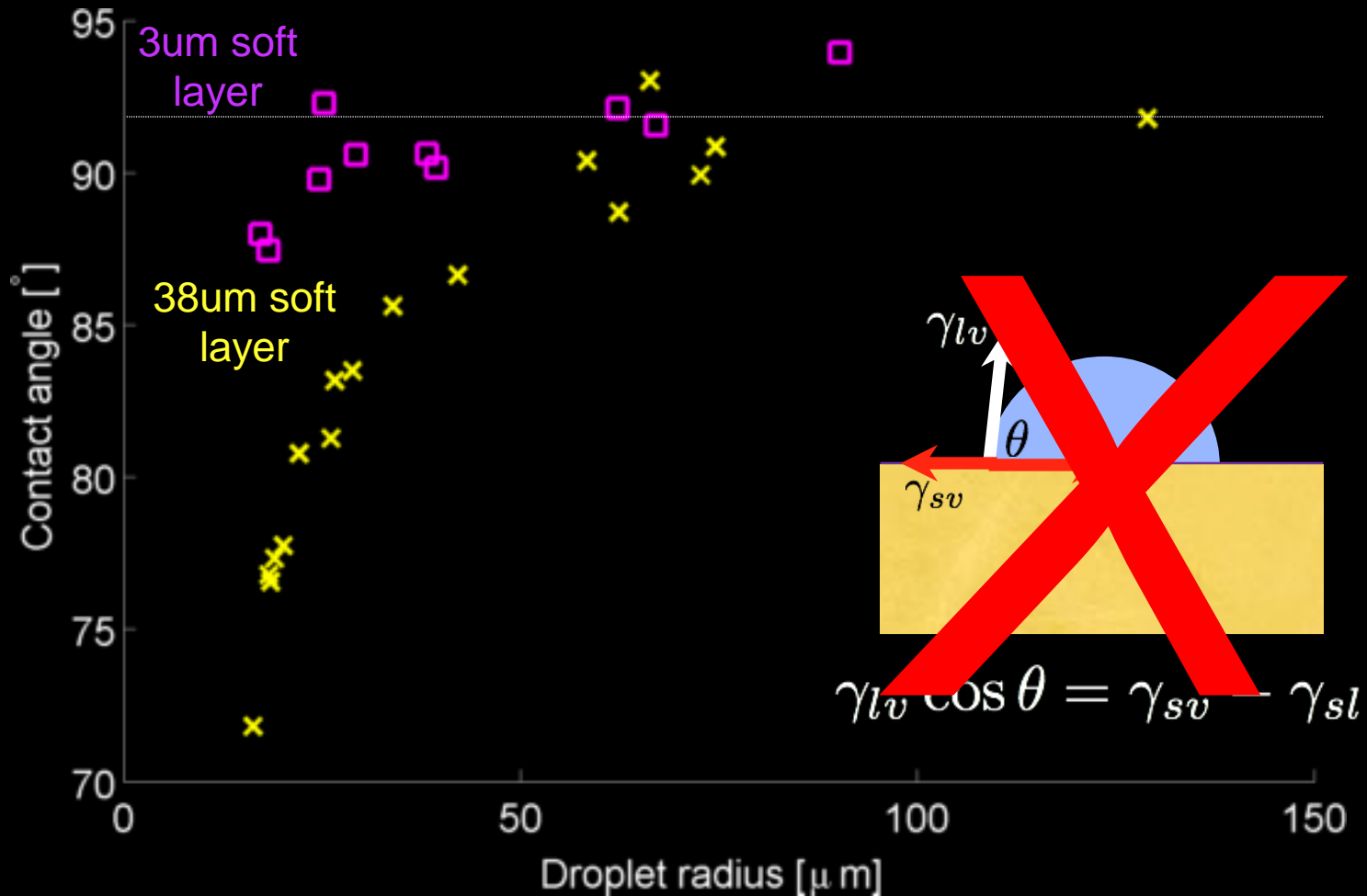
What happens when the solid surface is soft?



Dr. Robert Style

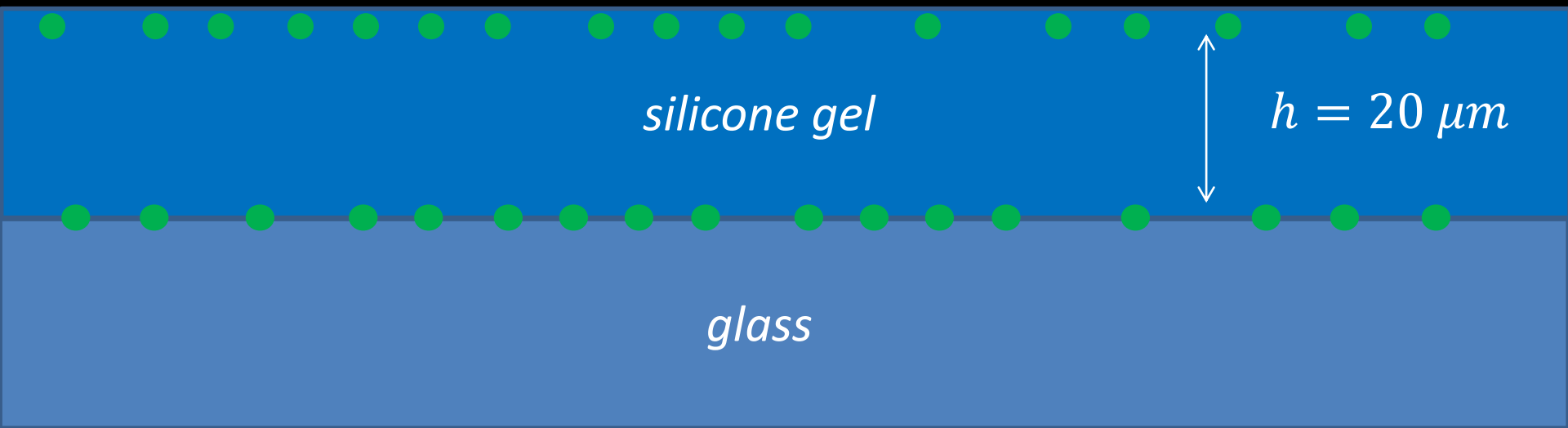


On a soft substrate, apparent contact angle depends on droplet size and thickness of soft layer

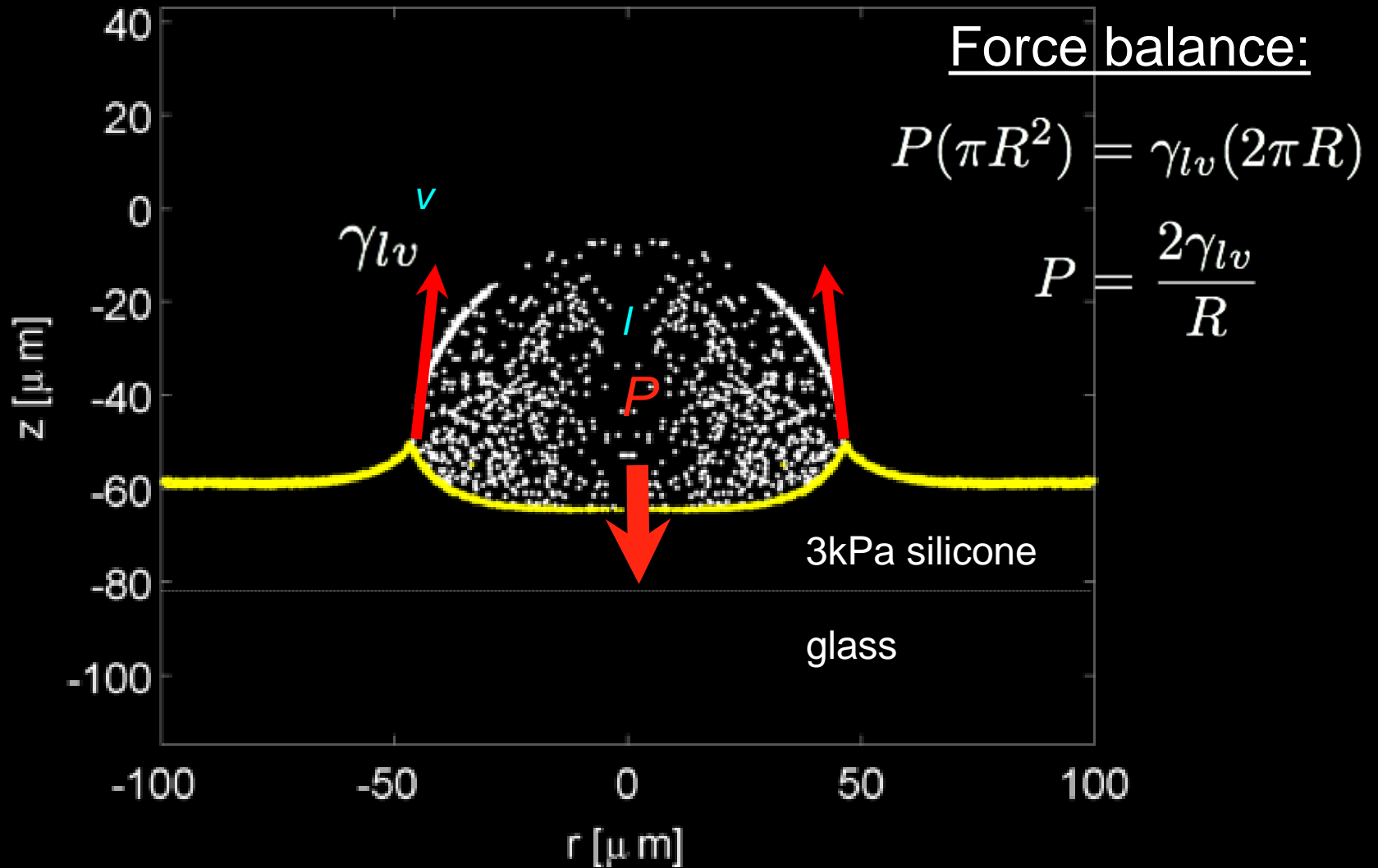


Glycerol drops on silicone ($E=3\text{kPa}$)
Zygo surface profilometer

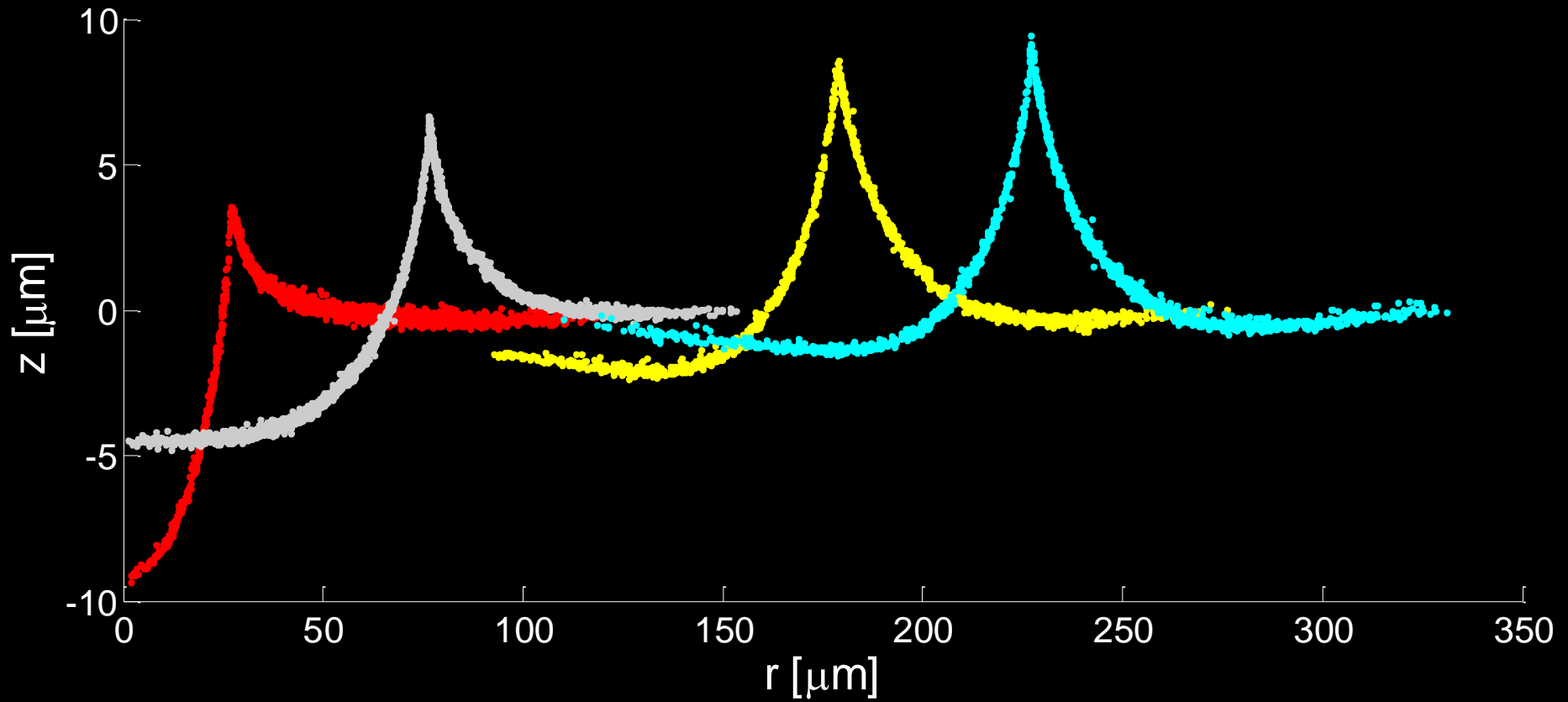
Confocal Imaging Near the Contact Line



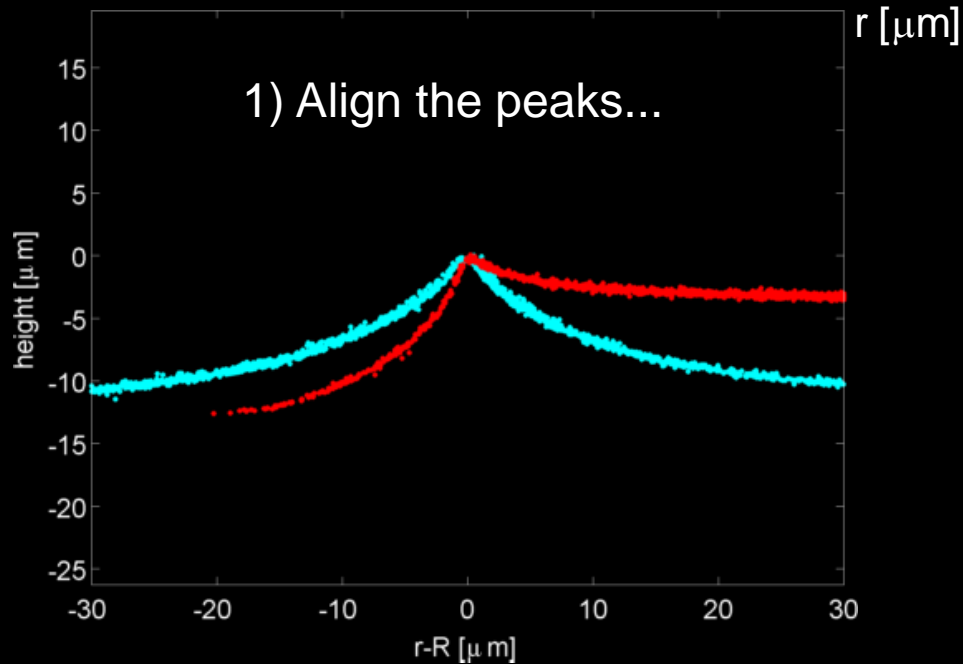
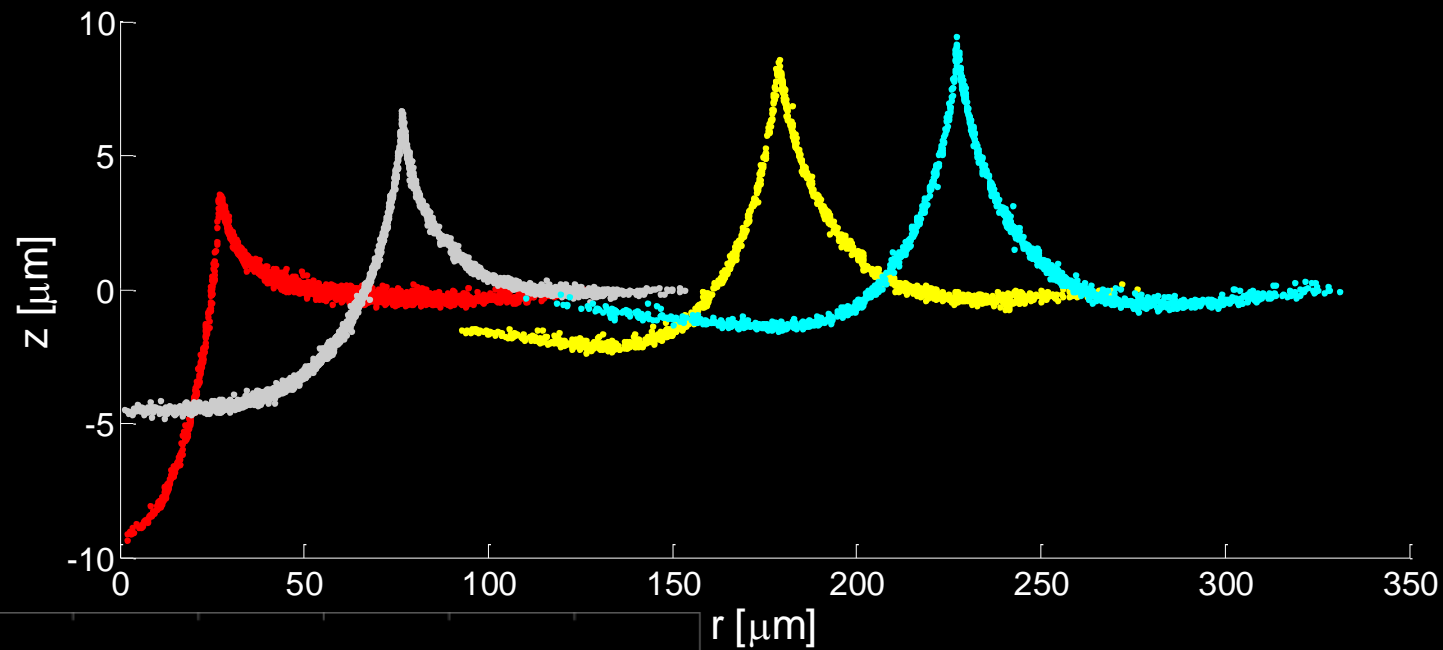
Droplets Deform Soft Substrates



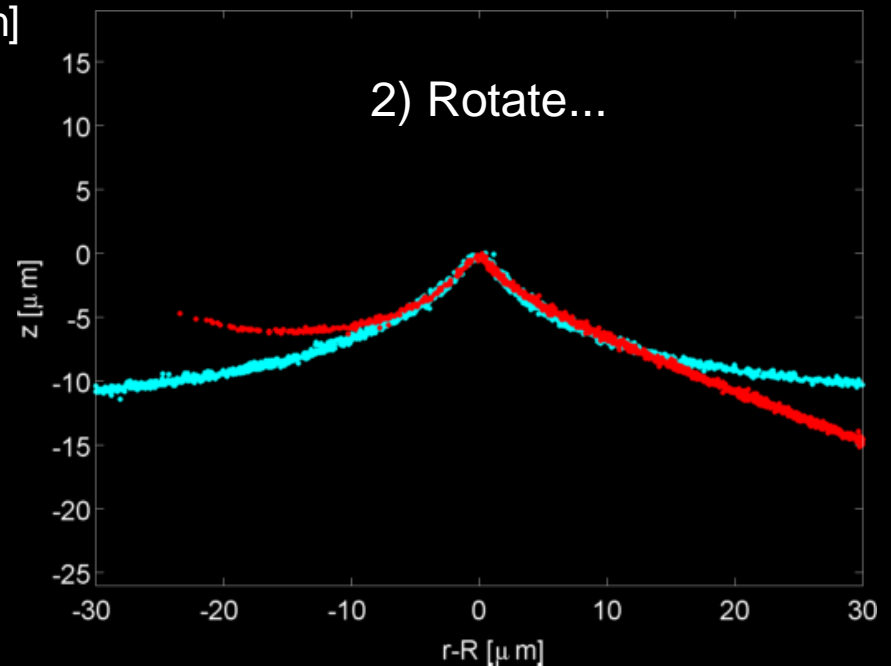
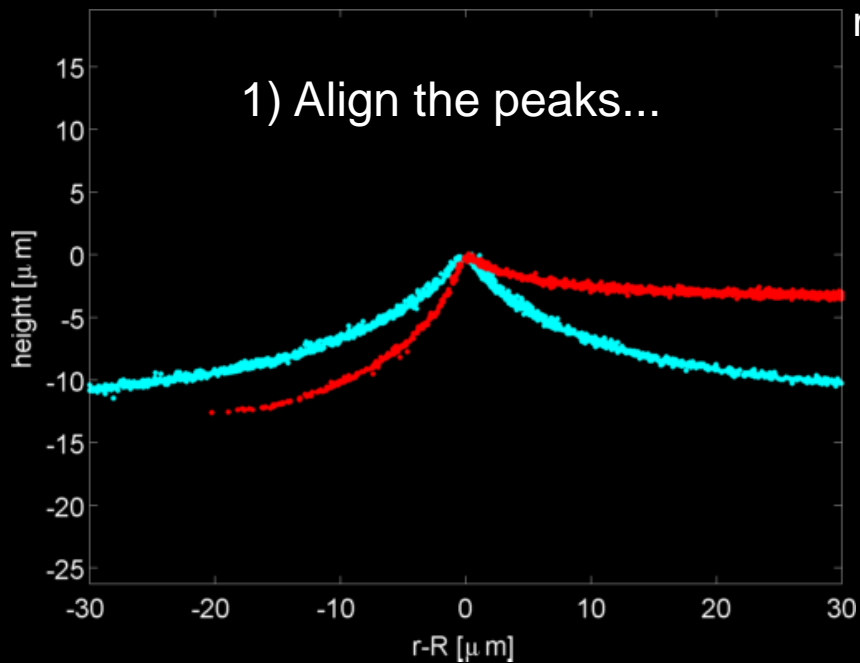
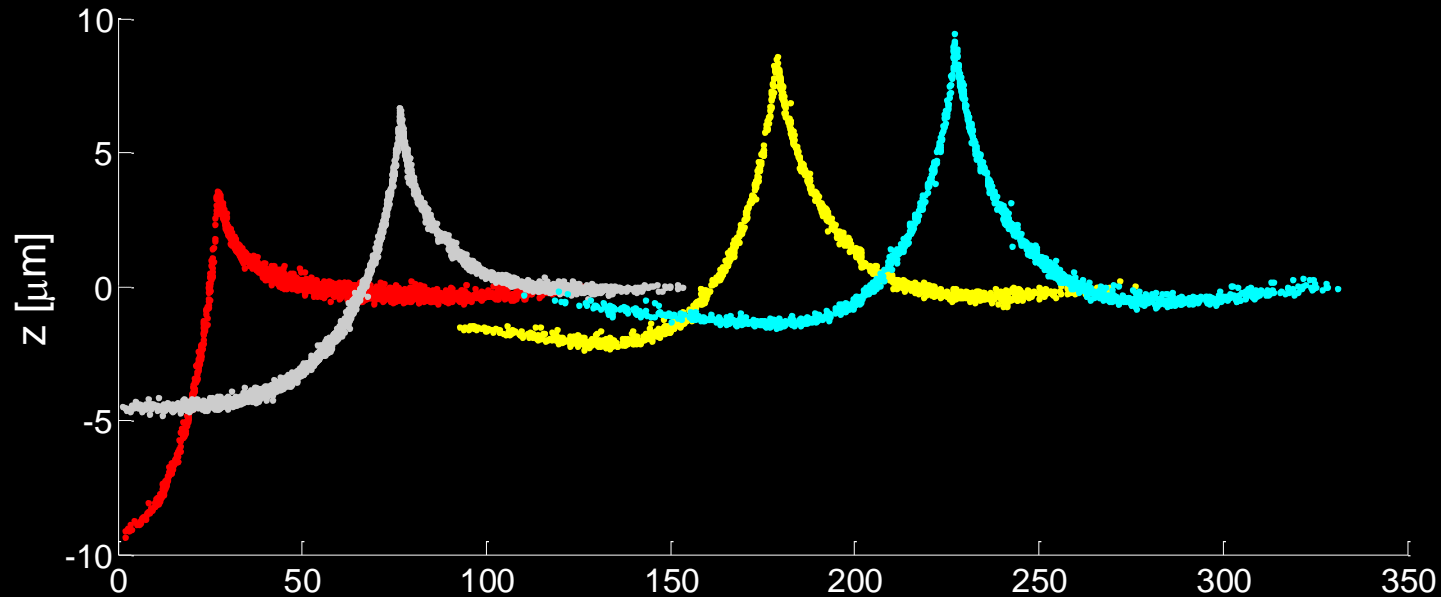
Profiles change dramatically with droplet size



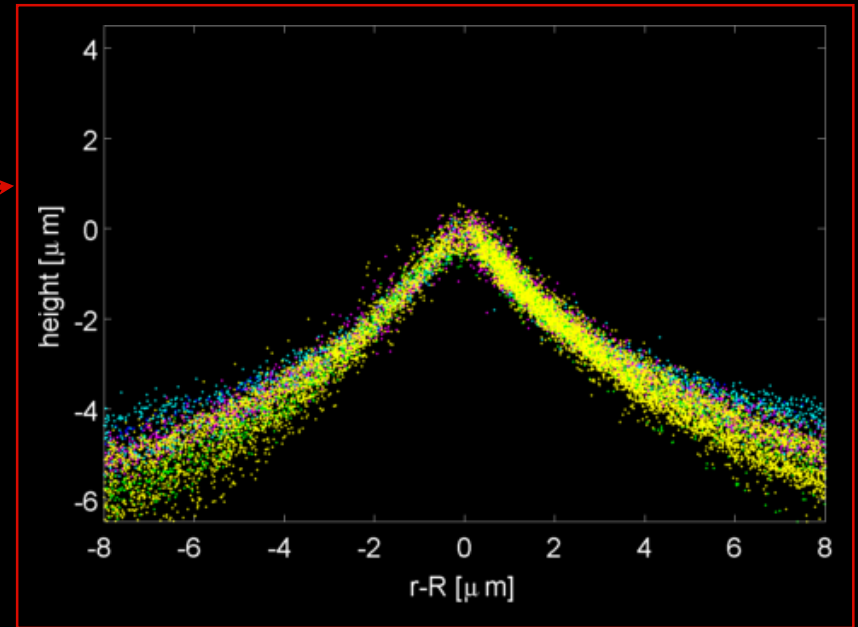
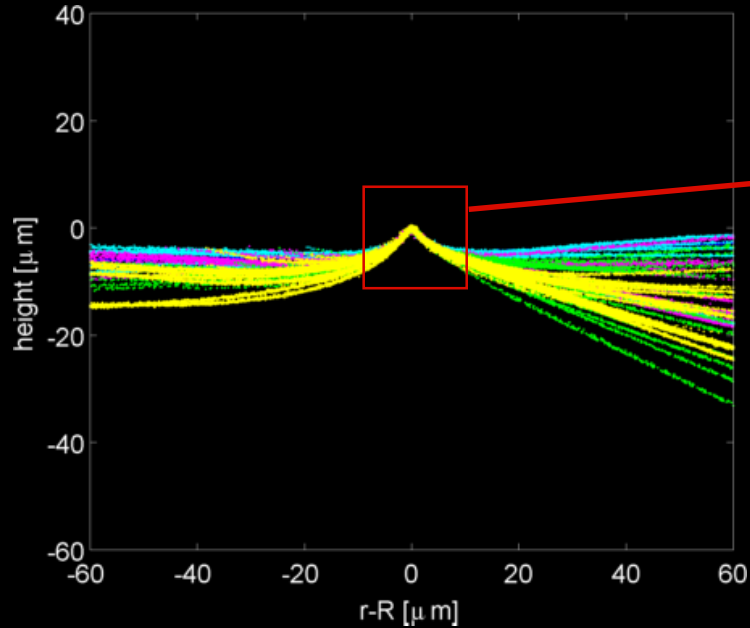
Overall profiles change – but how about the cusp?



Overall profiles change – but how about the cusp?

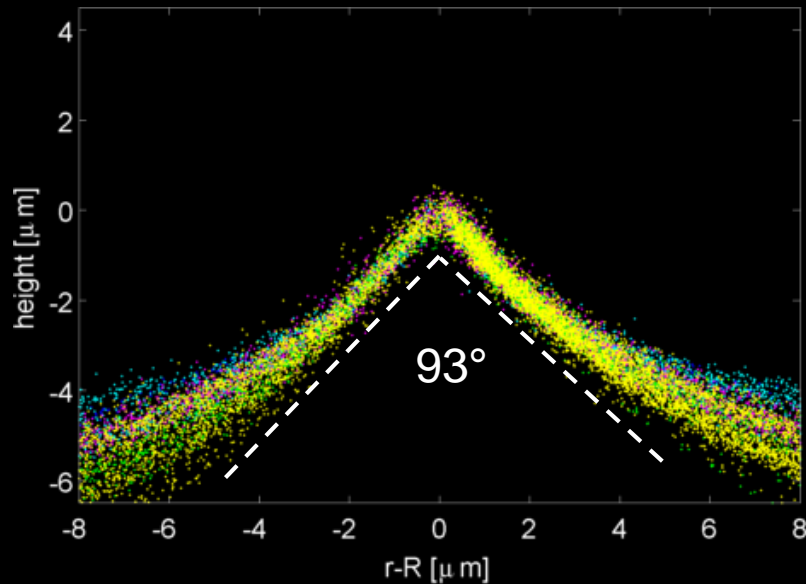


Cusp shape is universal near contact line

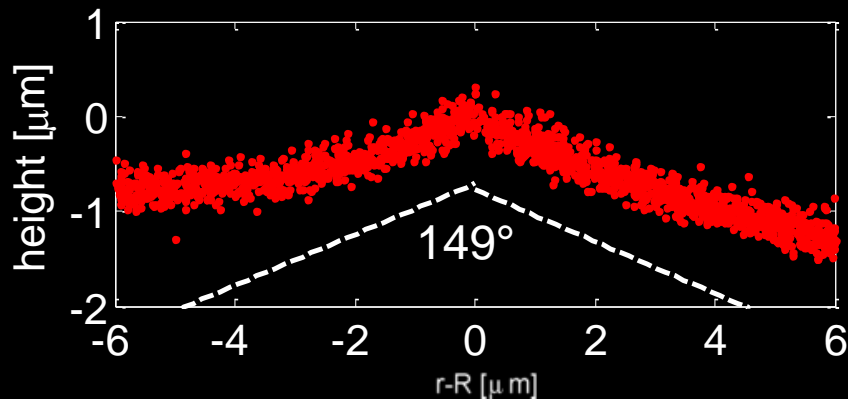


61 glycerol drops
radii: 18 μm - 1000 μm
Four different substrates: 13.5 - 50 μm thick

Contact line geometry depends on the wetting fluid

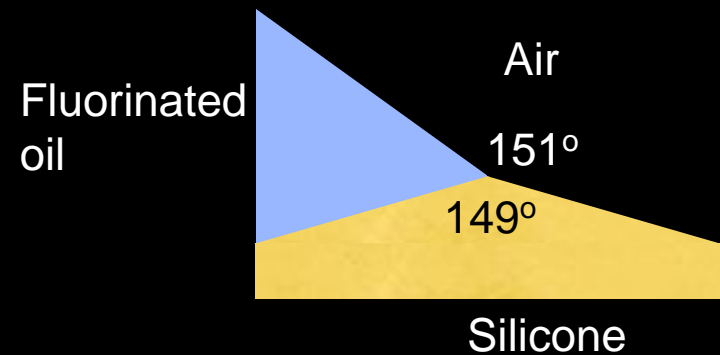
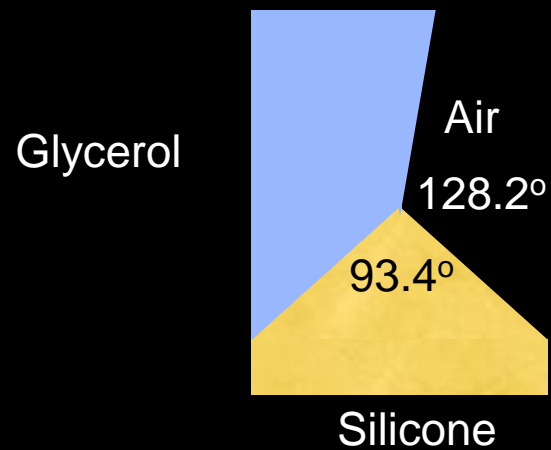


glycerol
61 drops
radii: 18 μm - 1000 μm
substrates: 13.5 - 50 μm thick

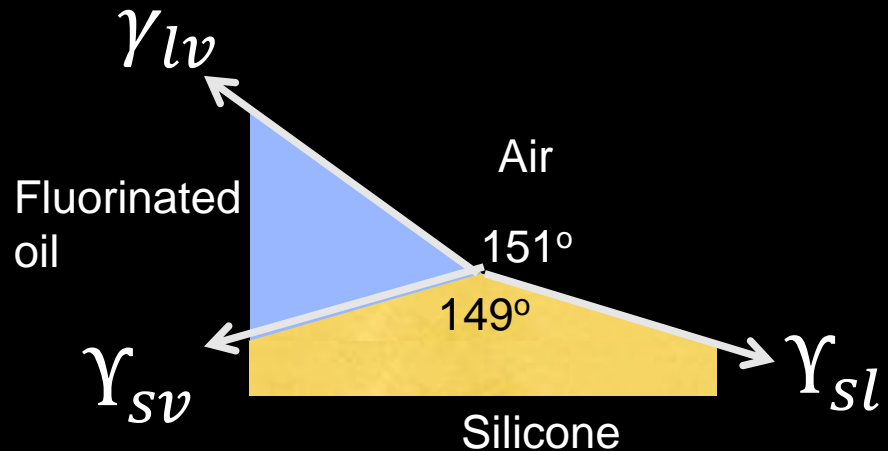
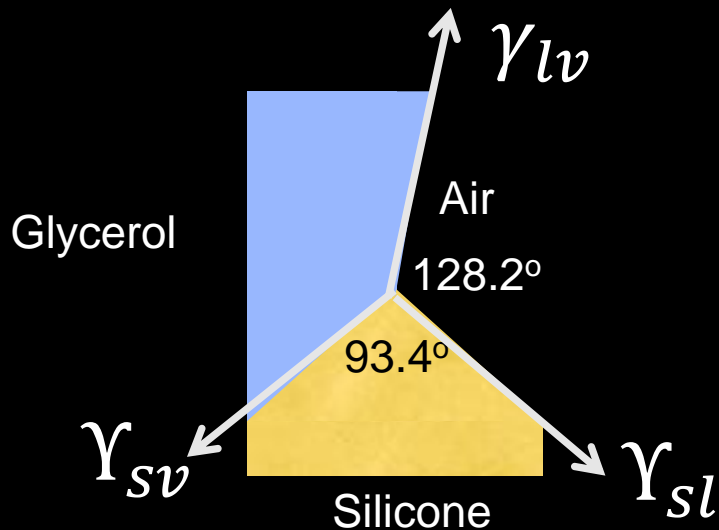


fluorinert fc-70
14 drops
radii: 140 μm - 270 μm
substrate: 23 μm thick

While apparent contact angle depends on boundary conditions...
microscopic configuration of interfaces is universal



Hypothesis: geometry at contact line is determined by a vector balance of surface tensions – a la Neumann



γ_{lv} : l-v surface tension

γ_{sv} : s-v surface tension

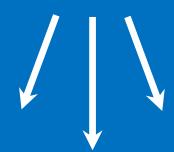
γ_{sl} : s-l surface tension

What about solid elasticity?

vapor

liquid

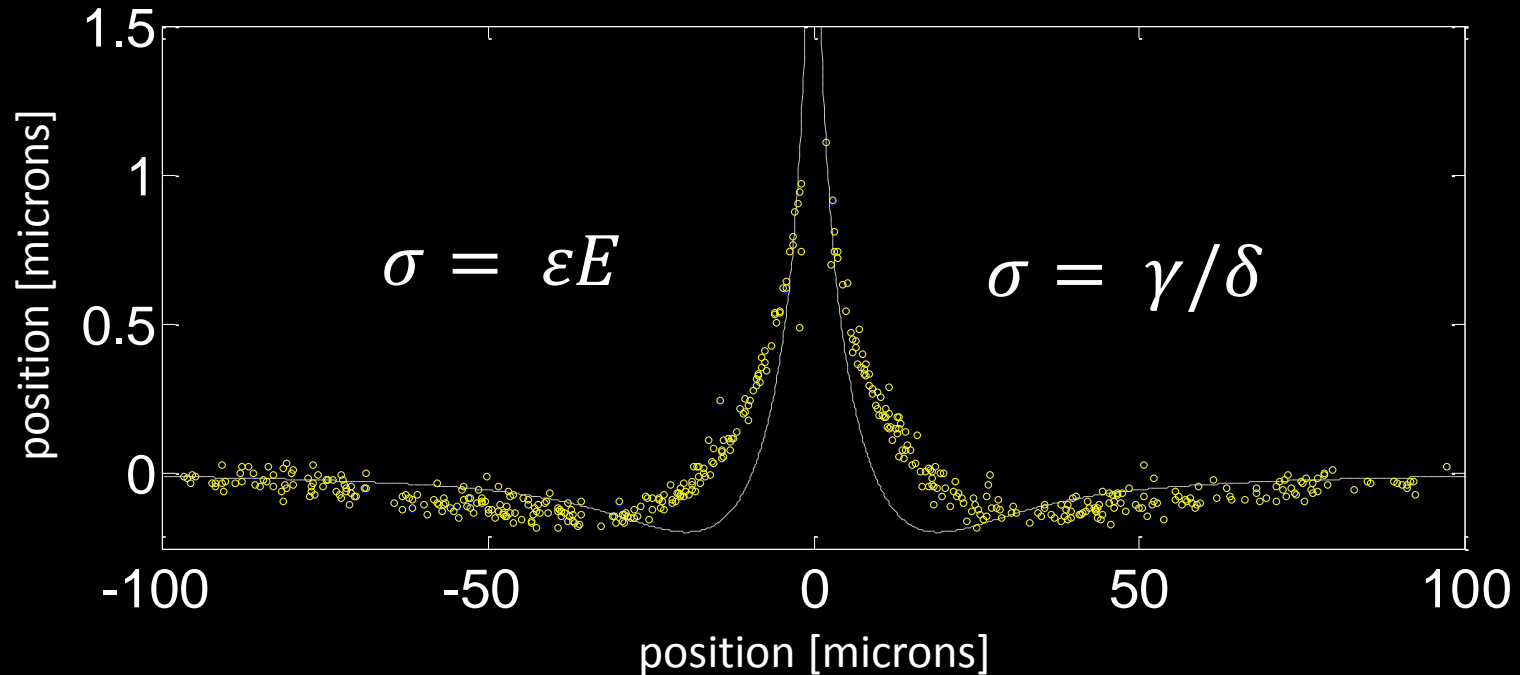
$\gamma_{LV} \sin \theta$



elastic restoring force?

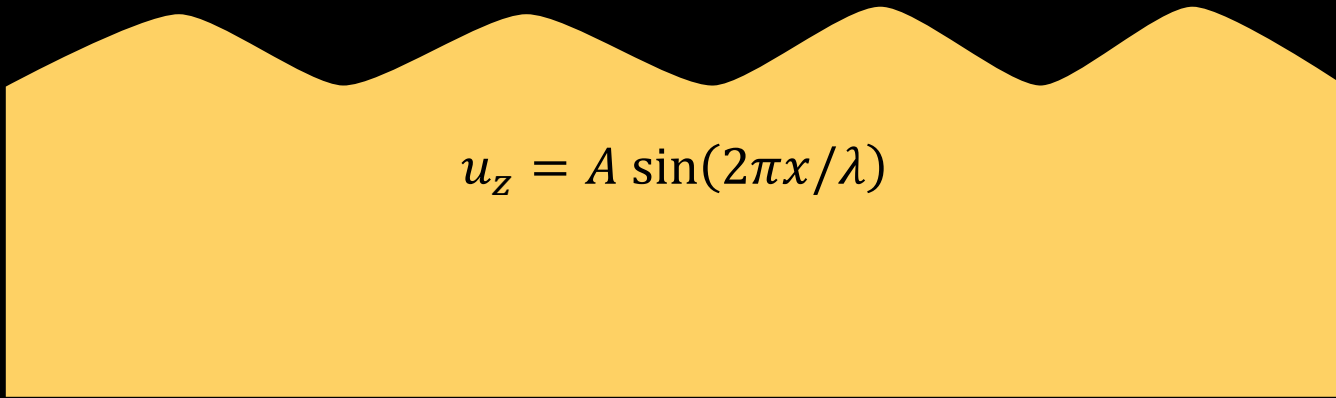
solid

Elastic Theories Cannot Balance Contact Line Forces



Reasonable estimates for contact line width lead to unreasonable strains and displacements

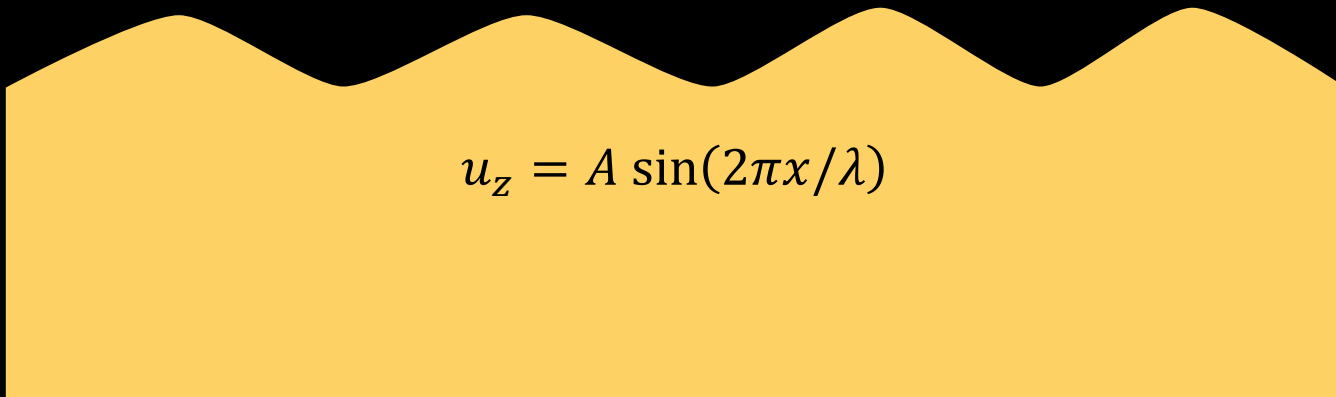
Flattening of a linear elastic solid



Elastic restoring force: $\sigma_E = \varepsilon E \sim AE/\lambda$

Flattening of a linear elastic solid by surface tension

Long Ajdari 1996, Jerison Dufresne 2011, Jagota 2012

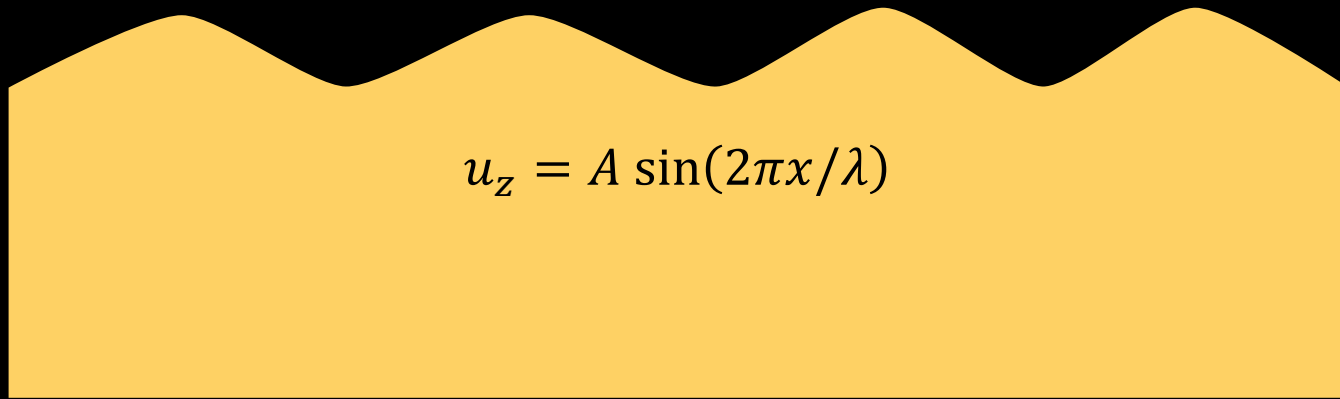


Elastic restoring force: $\sigma_E = \varepsilon E \sim AE/\lambda$

Capillary force: $\sigma_\gamma \sim \Upsilon A/\lambda^2$

Υ : solid surface tension

Balance of Elasticity and Capillarity Defines a Length scale



$$\frac{\sigma_E}{\sigma_\gamma} \sim \frac{\lambda}{l}$$

$$l = \gamma/E$$

Elastocapillary Length

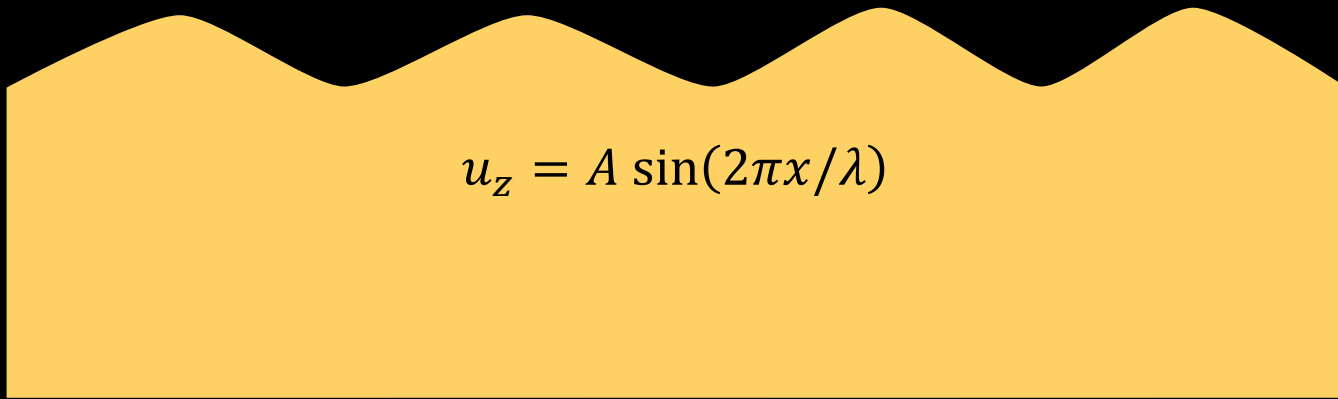
$$\lambda/l \gg 1$$

Elasticity Dominates

$$\lambda/l \ll 1$$

Surface Tension Dominates

Capillarity Dominates at Short Length Scales on Soft Materials



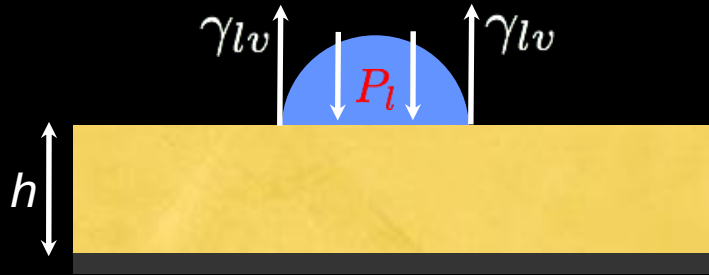
$$\gamma = 0.03 \text{ N/m}$$

$$l = 0.1 \text{ \AA} \text{ for } E = 3 \text{ GPa}$$

$$l = 10 \text{ nm for } E = 3 \text{ MPa}$$

$$l = 10 \text{ } \mu\text{m for } E = 3 \text{ kPa}$$

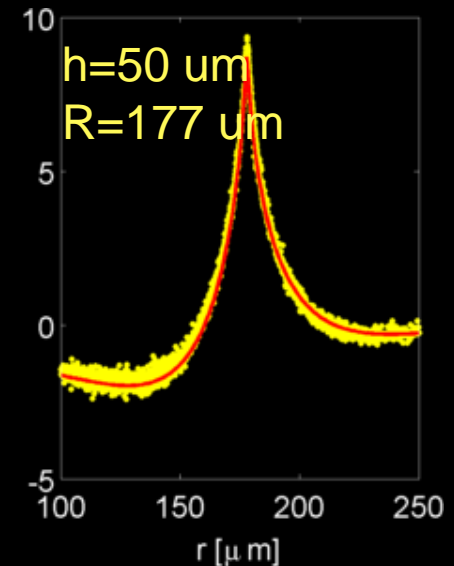
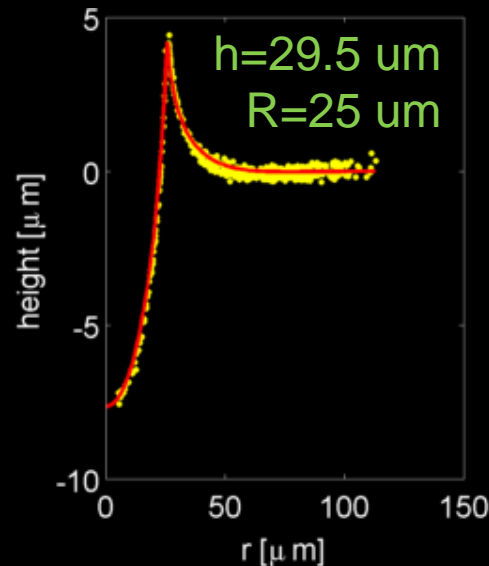
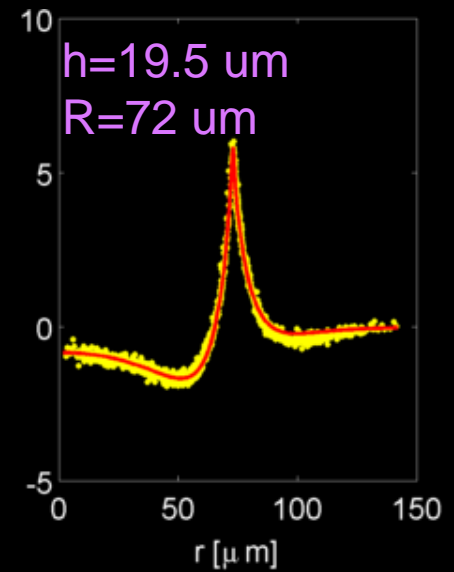
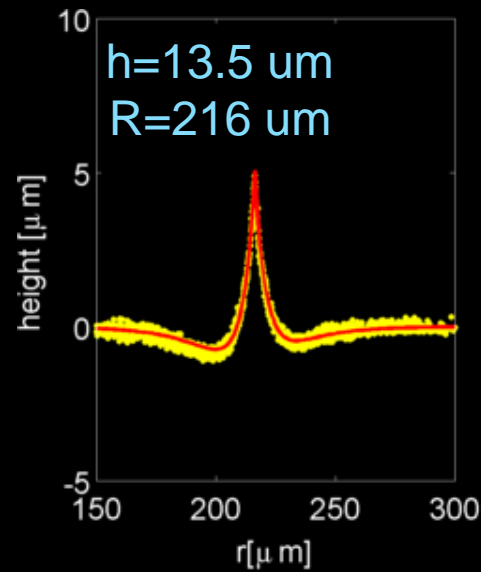
Linear Elasticity Plus Solid Surface Tension Captures Profiles



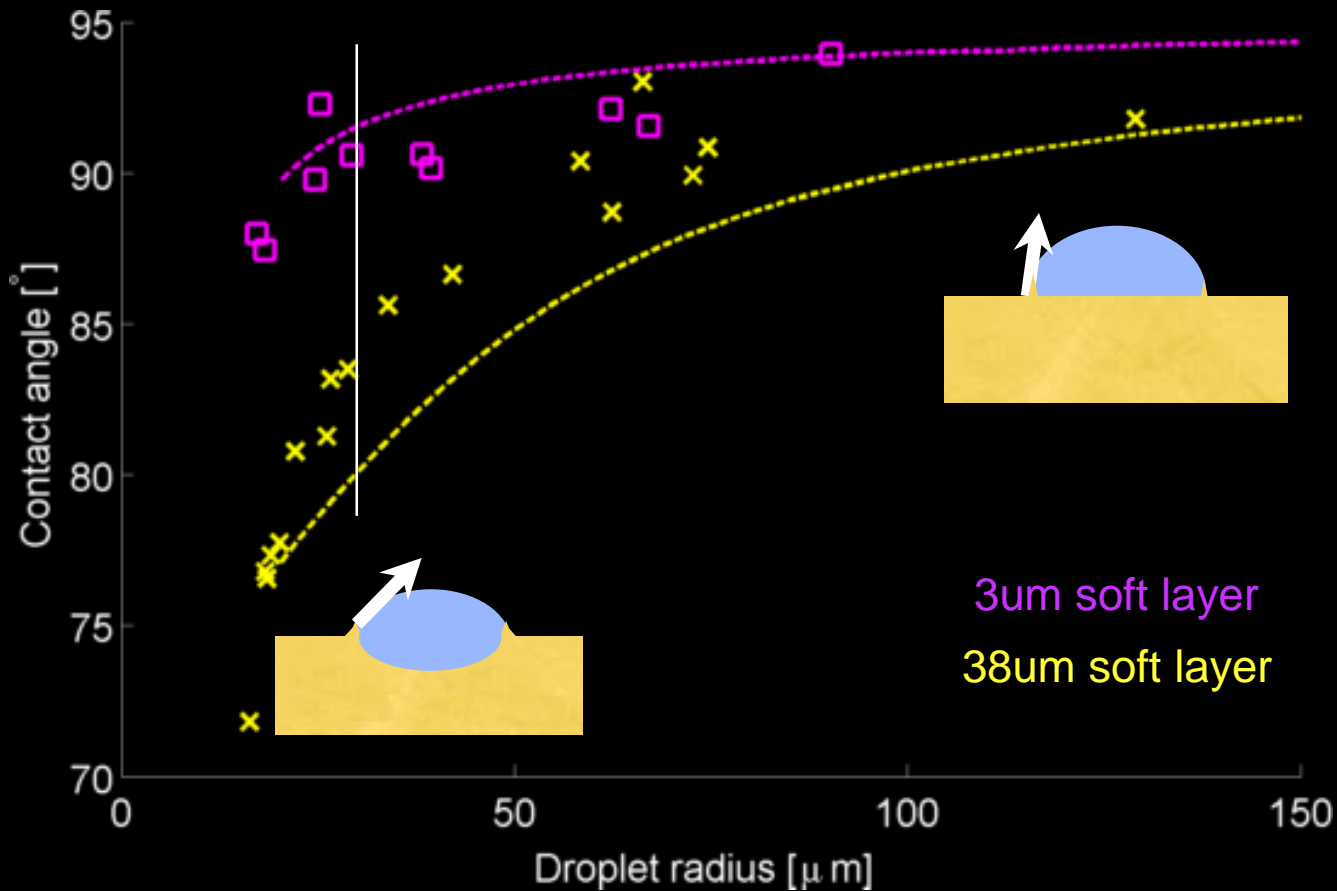
Linear elastic solid with surface tension

Sharp features near contact line are controlled by surface tension.

Far-field determined by elasticity

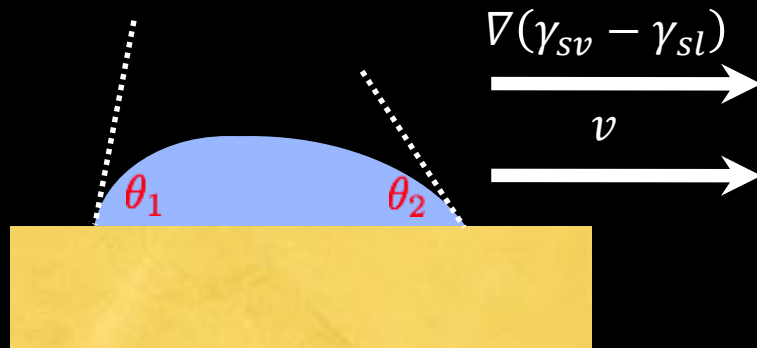
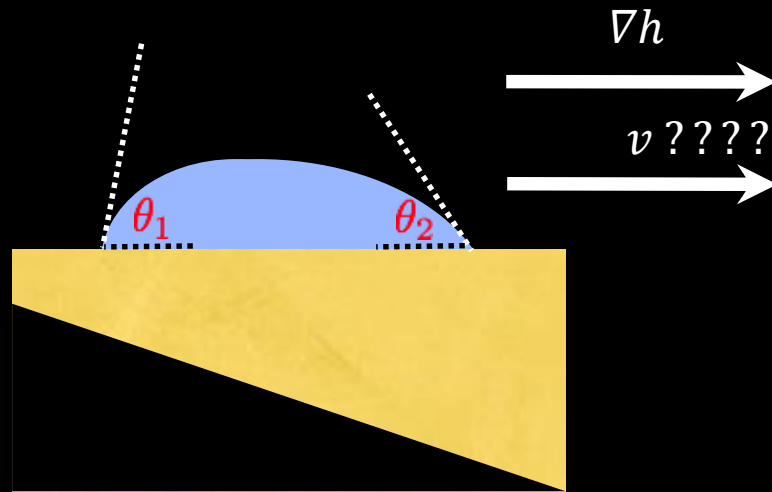


Linear Elasticity Plus Solid Surface Tension Predicts Change in Apparent Contact Angle



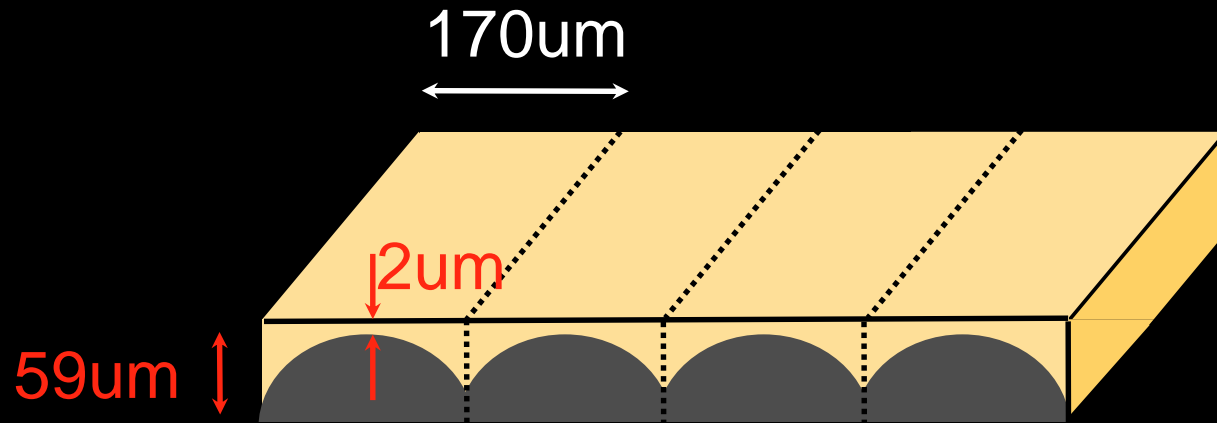
Glycerol drops on silicone ($E=3\text{kPa}$)

Differences in contact angle drive droplet motion?

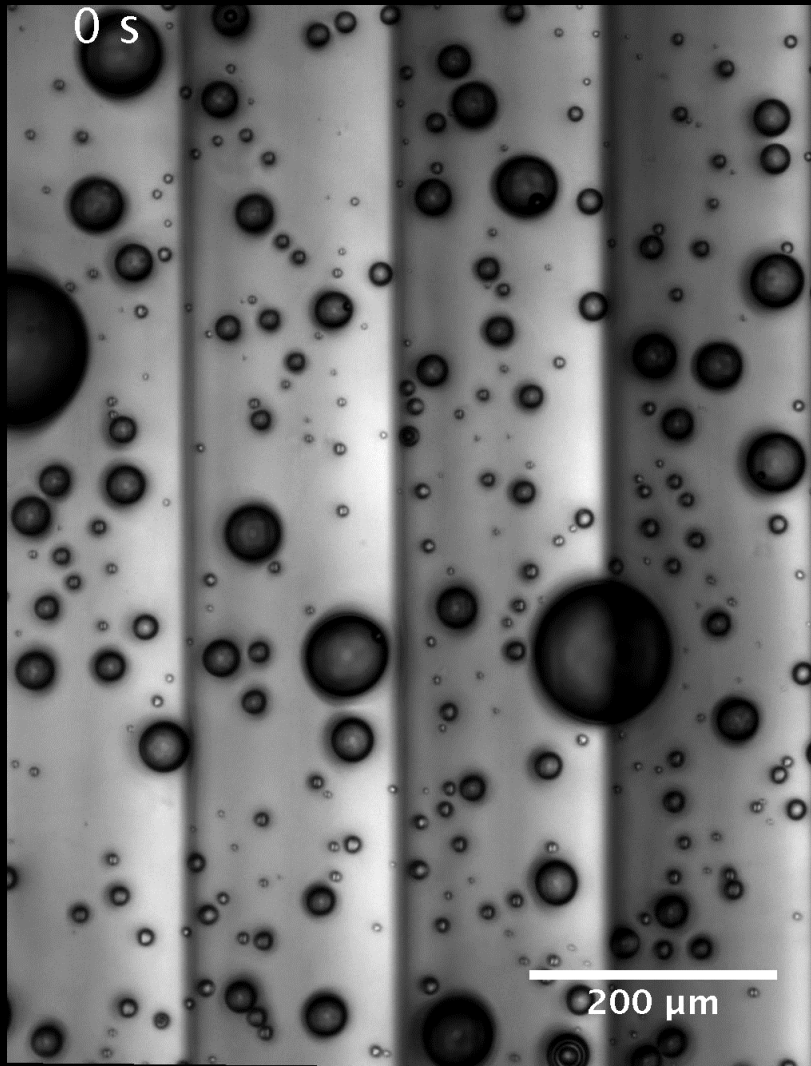


see e.g. chaudhury,
whitesides, troian

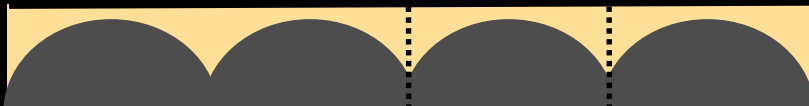
Making a thickness gradient



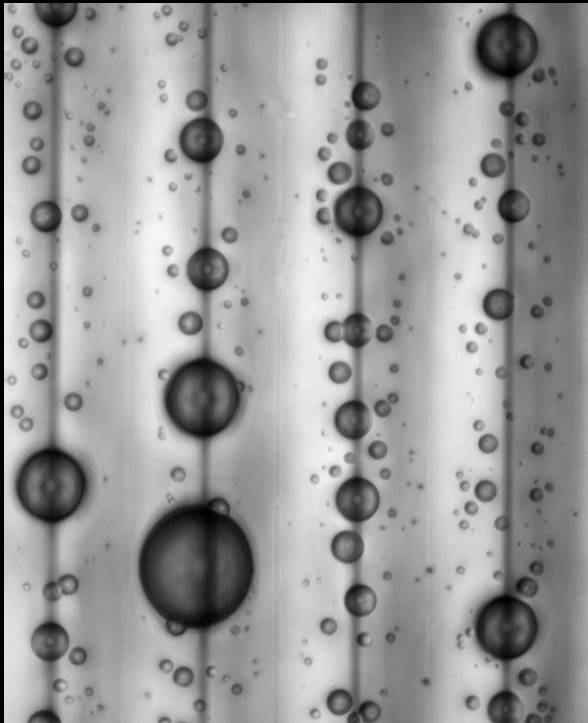
3kPa silicone gel coated on hard substrate
with close-packed cylindrical ridges



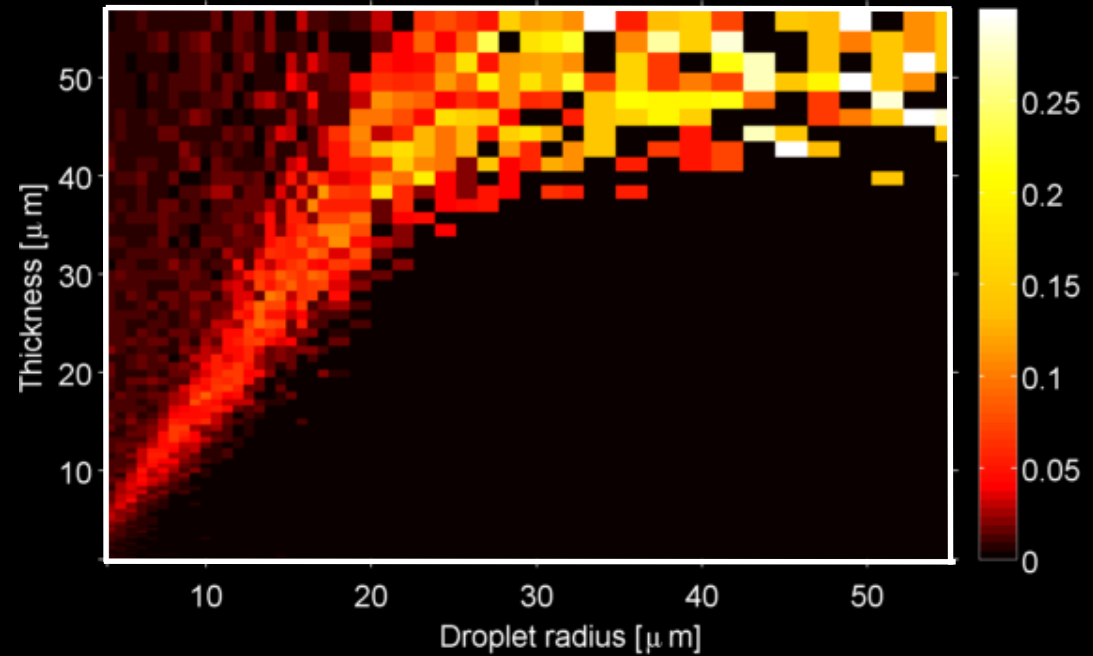
Atomized spray of
glycerol on substrate with
thickness/stiffness gradient



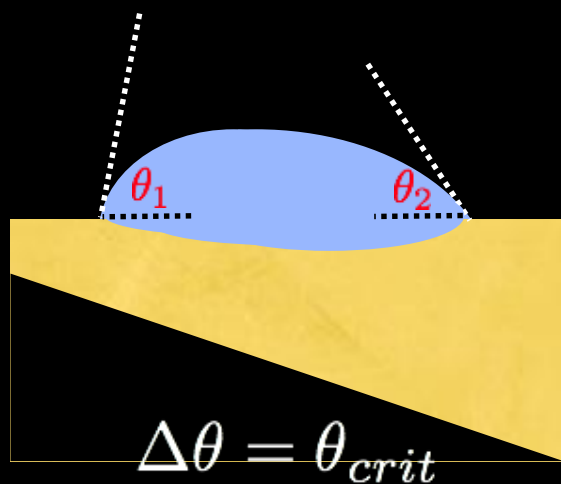
Resting depth depends on drop size



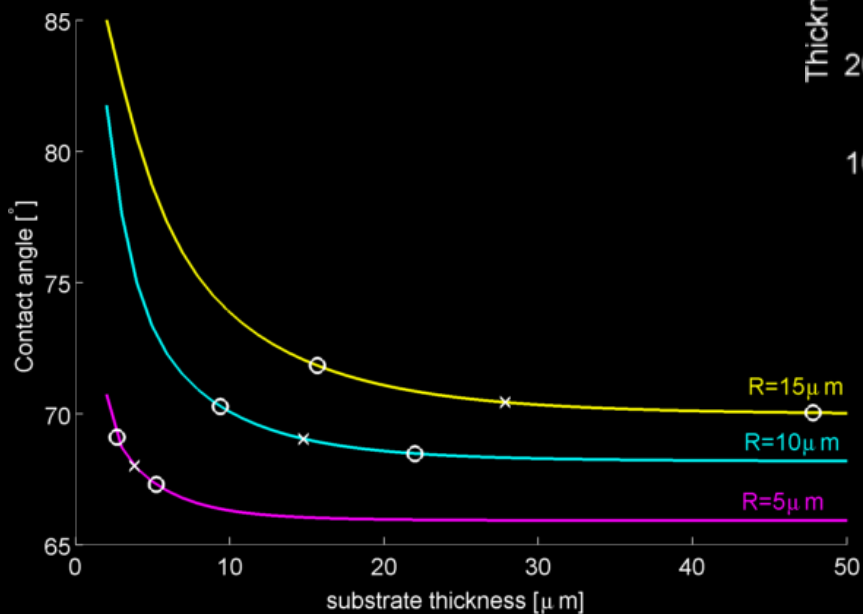
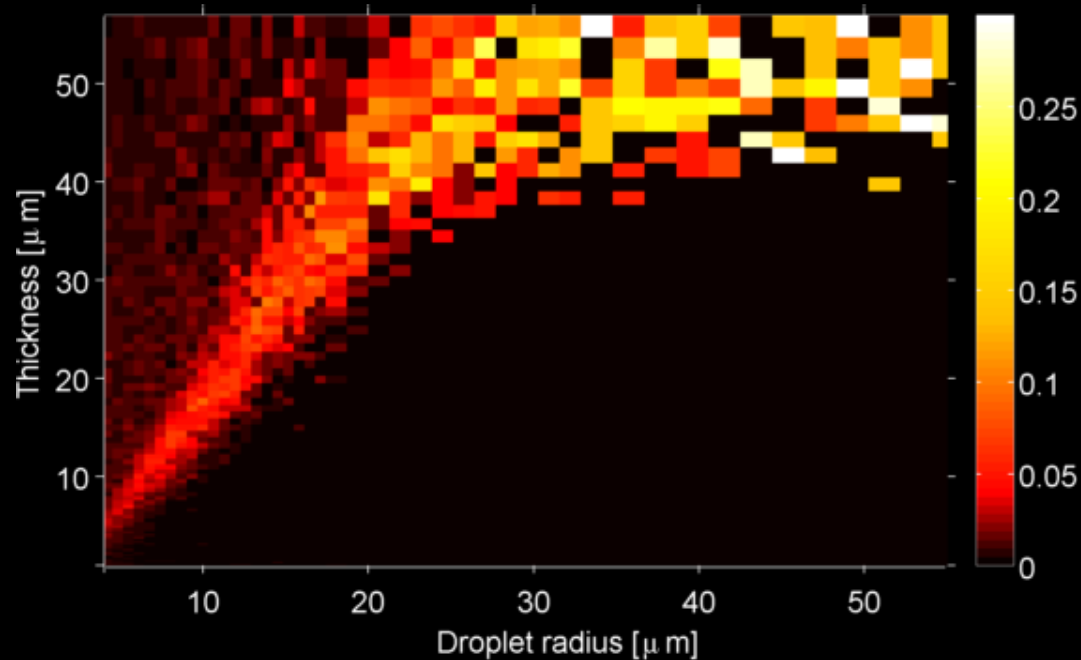
Probability density of 'final' drop position



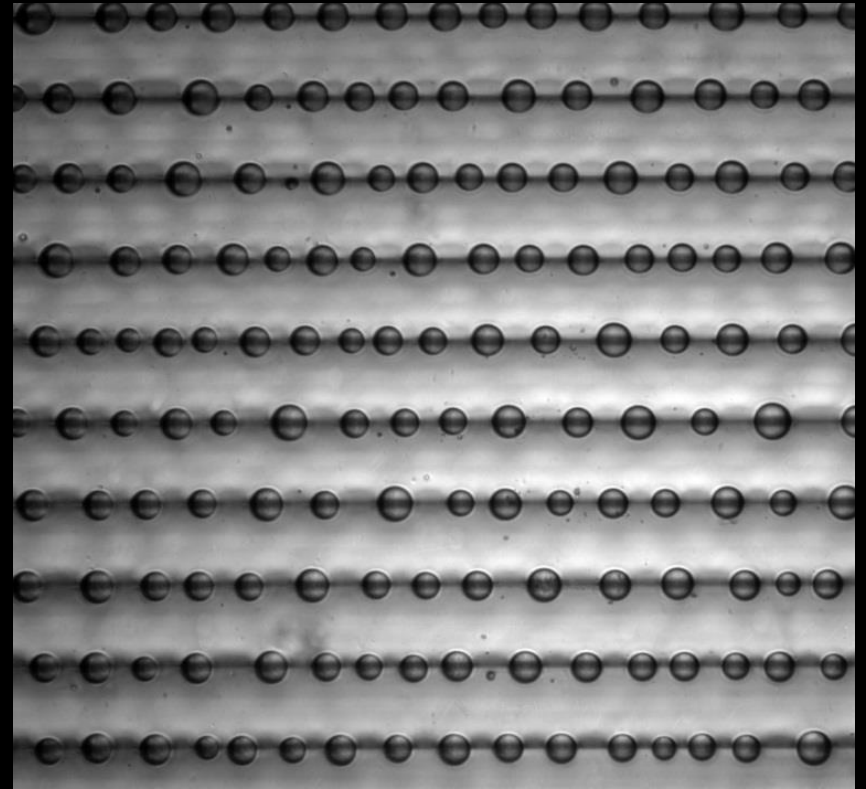
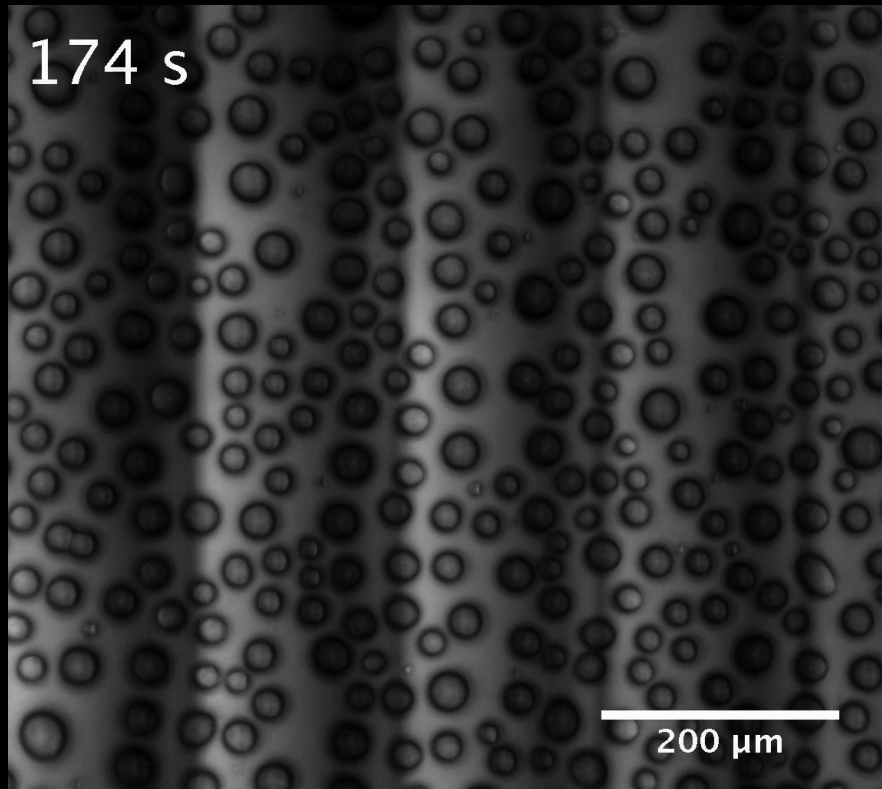
Simple theory for droplet motion on thickness gradient



Probability density of 'final' drop position

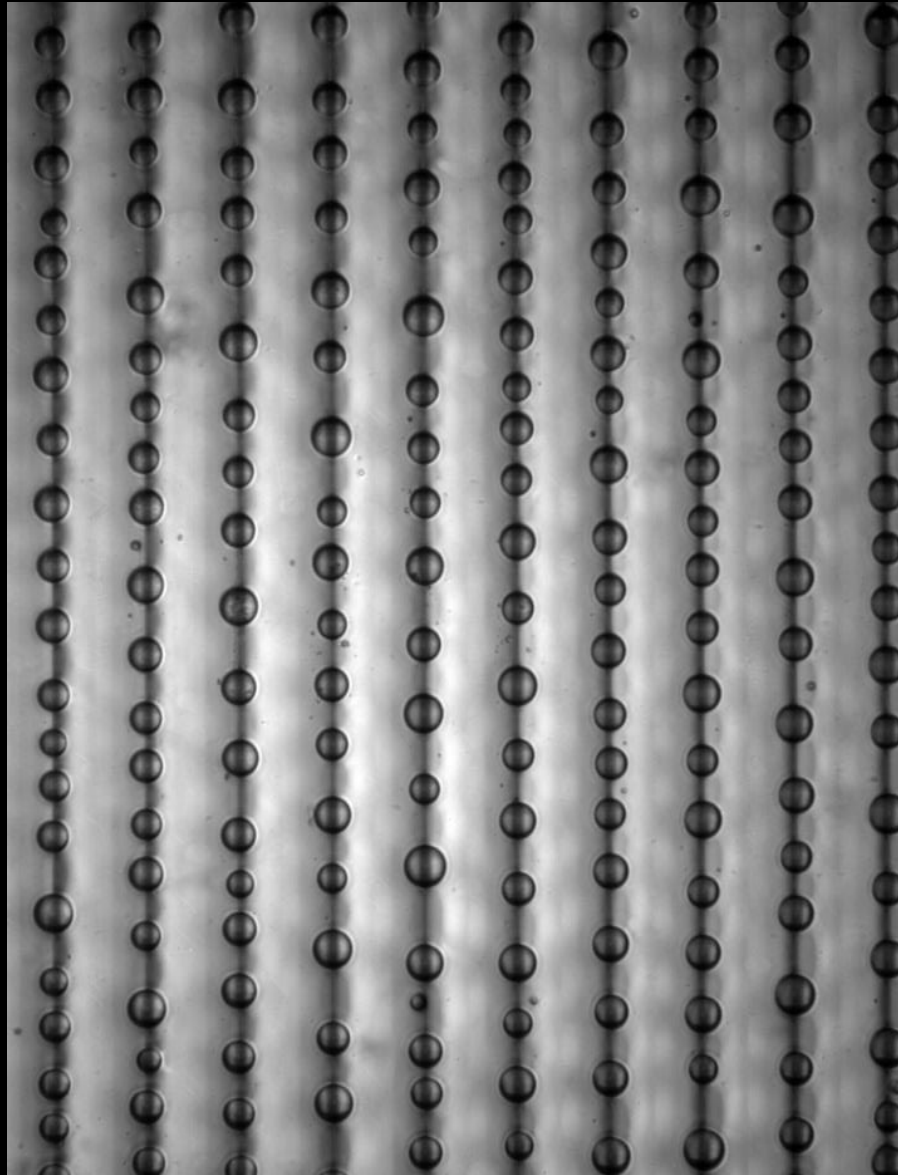


Condensation, coarsening and evaporation enhance patterning



0-600 s cooling by Peltier
>600 s Peltier off

Wetting Summary



Young's law fails on soft substrates as R approaches γ/E .

The shape of the cusp close to contact line is universal and determined by the surface tensions.

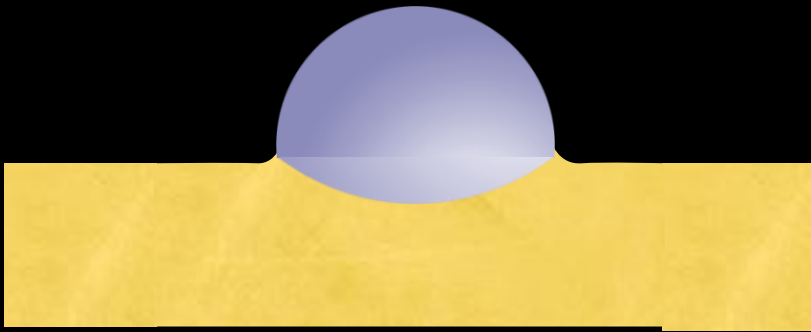
While the apparent contact angle is not universal it seems to drive droplet motion

Theory and preliminary expts:
Jerison et al Physical Review Letters 2011
Style et al Soft Matter 2012

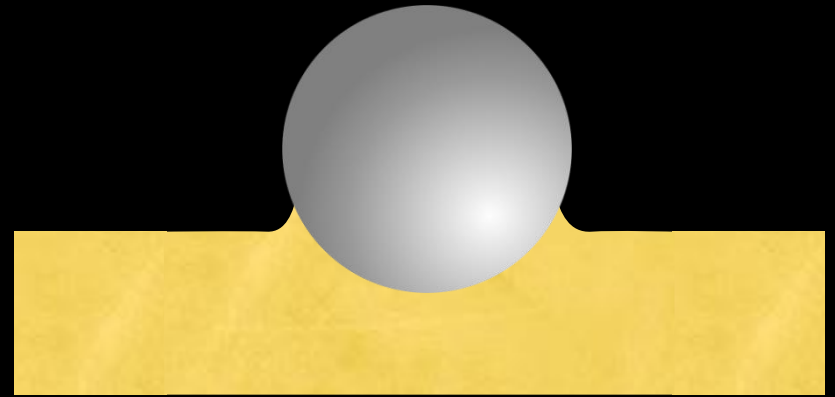
Breakdown of Young's Law
Style et al Physical Review Letters 2013

Drop movement
Style et al PNAS 2013

Wetting and adhesion don't look so different on a soft surface

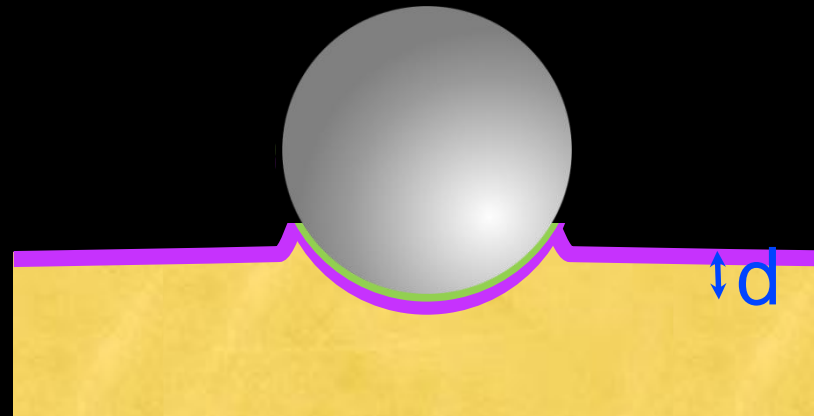


Wetting



Adhesion

Johnson, Kendall & Roberts (1971) – ‘JKR’

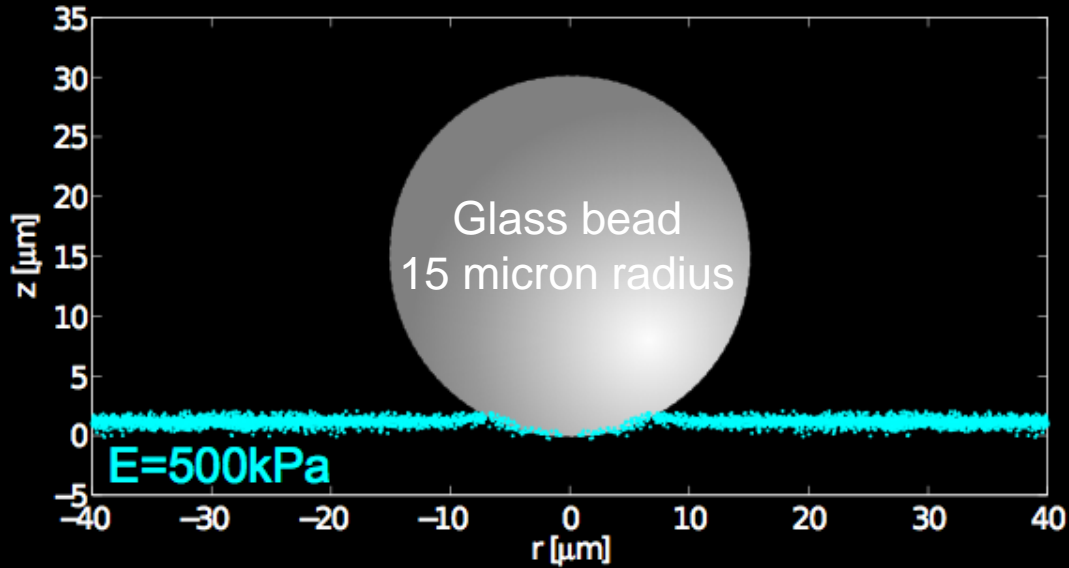


Adhesion Energy, $W = \gamma_{sp} - \gamma_{sv} - \gamma_{pv}$

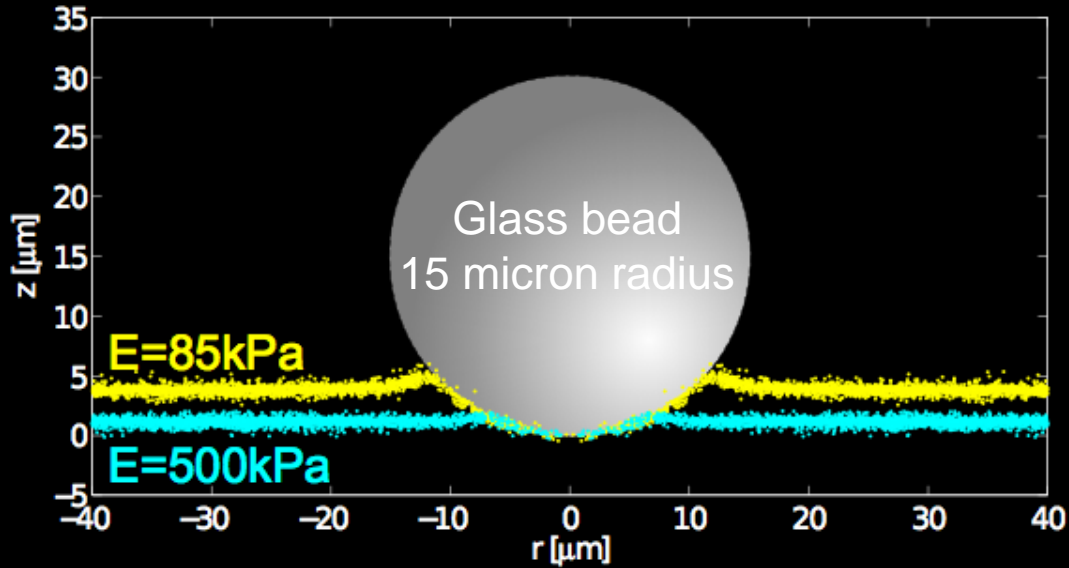
Substrate Elasticity, E

Substrate surface tension, $\gamma_{sv}, \gamma_{sp}?$

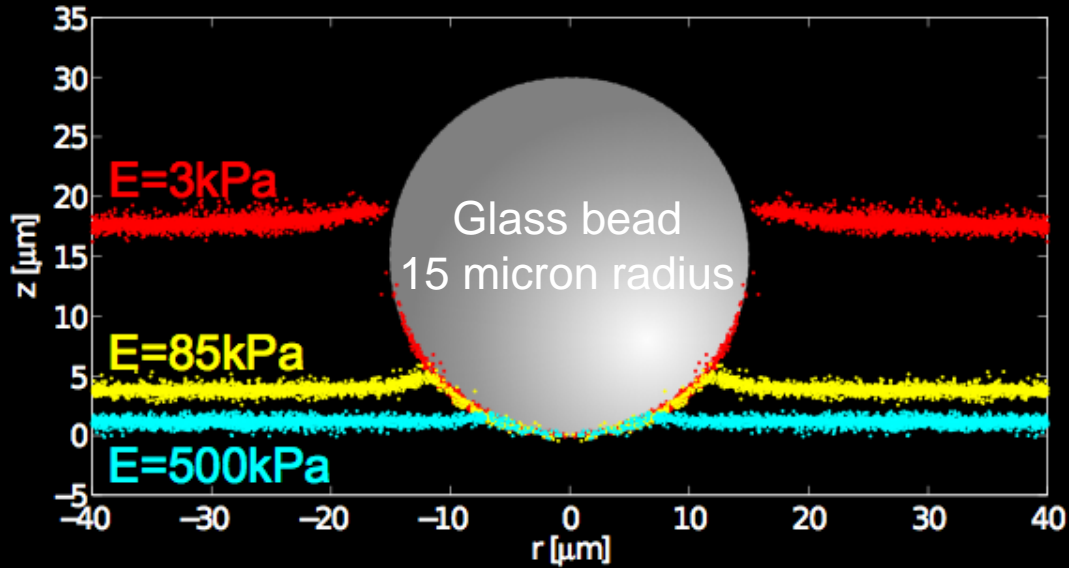
Surface profiles for different stiffness substrates



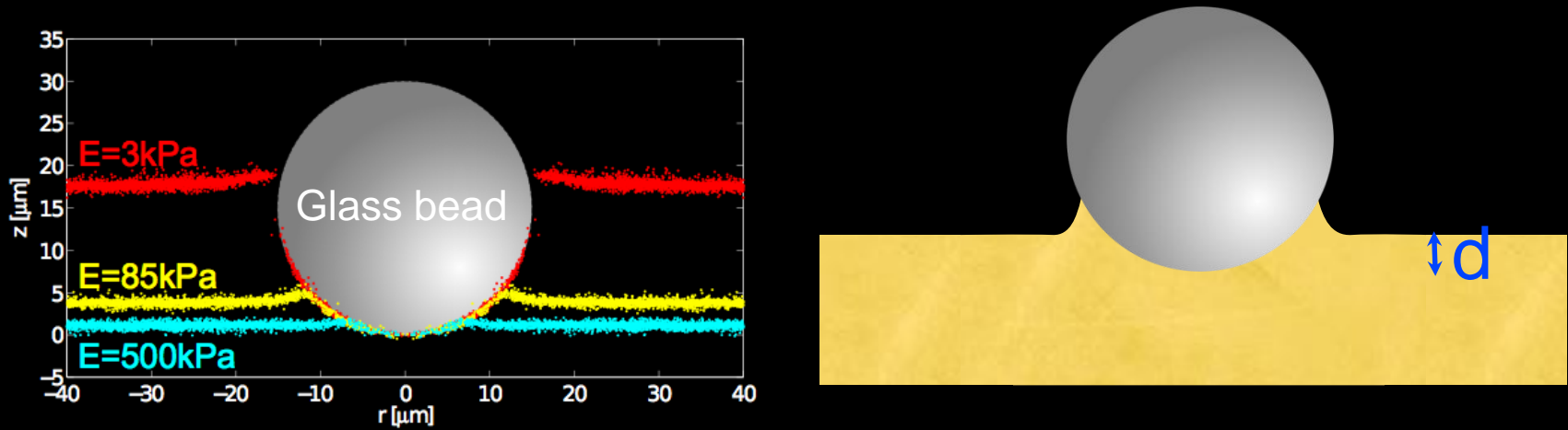
Surface profiles for different stiffness substrates



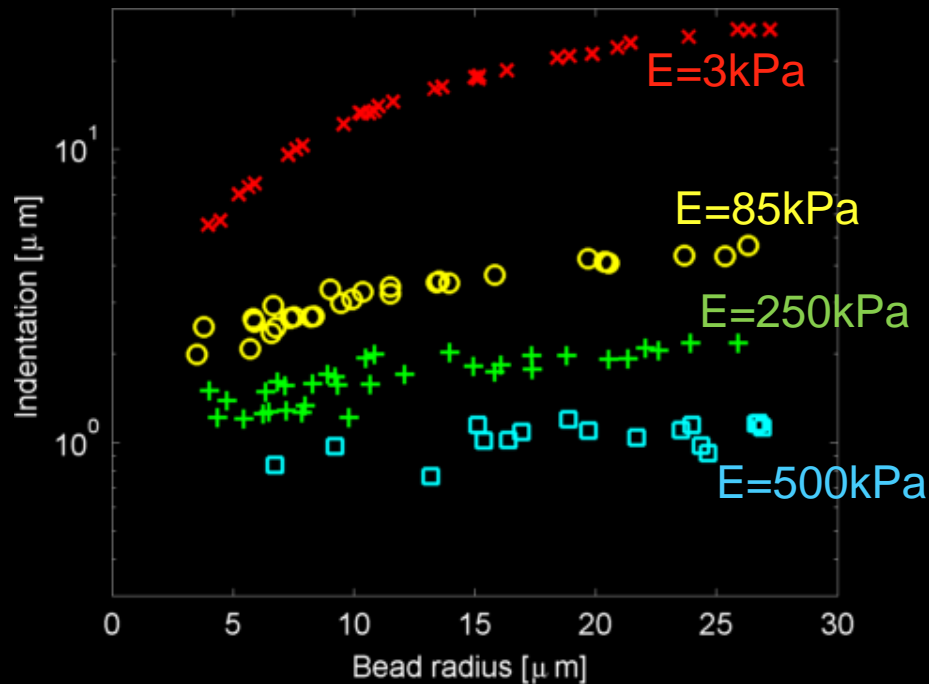
Surface profiles for different stiffness substrates



Indentation Depth Changes with Size and Stiffness

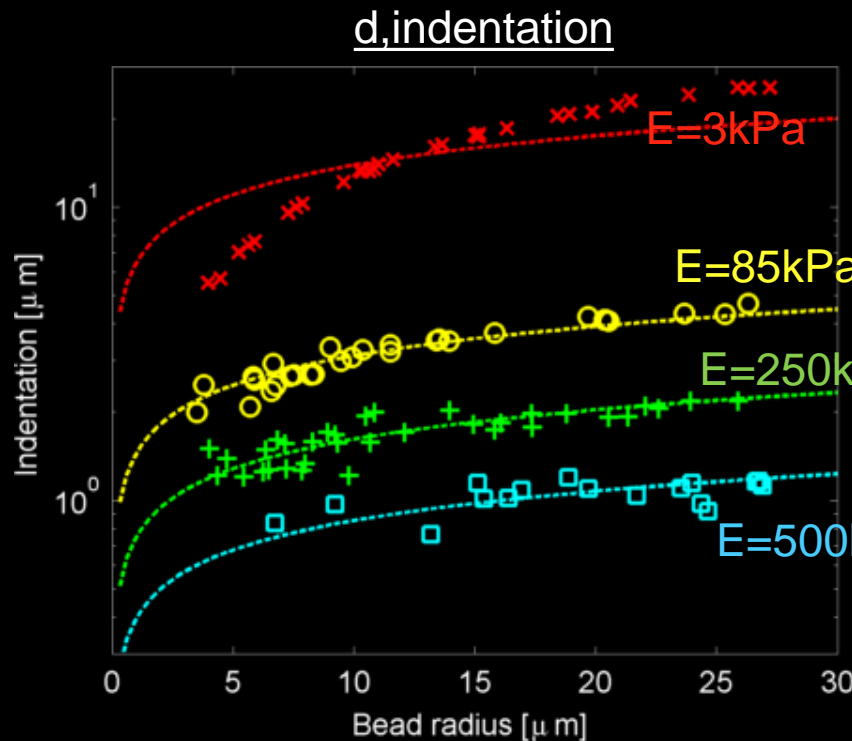


d , indentation



Comparing results with JKR

$$d = \left(\frac{\sqrt{3}W(1 - \nu^2)}{2E} \right)^{2/3} R^{1/3}$$

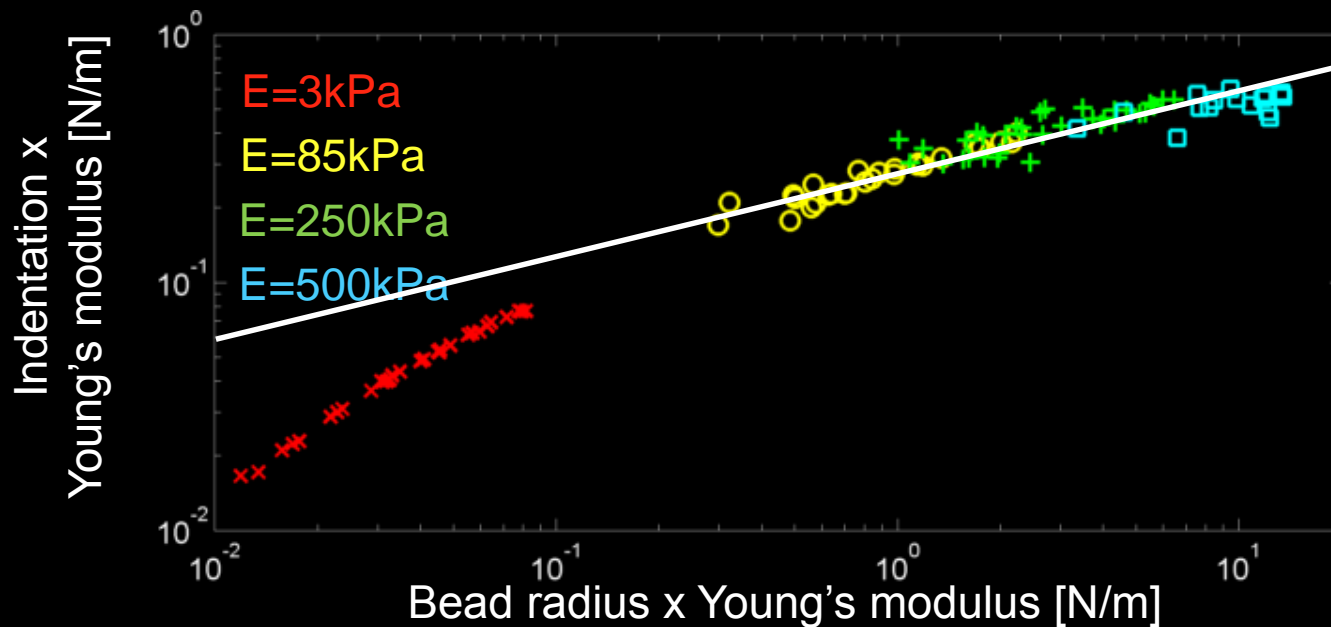


$W \sim 25\text{mN/m}$

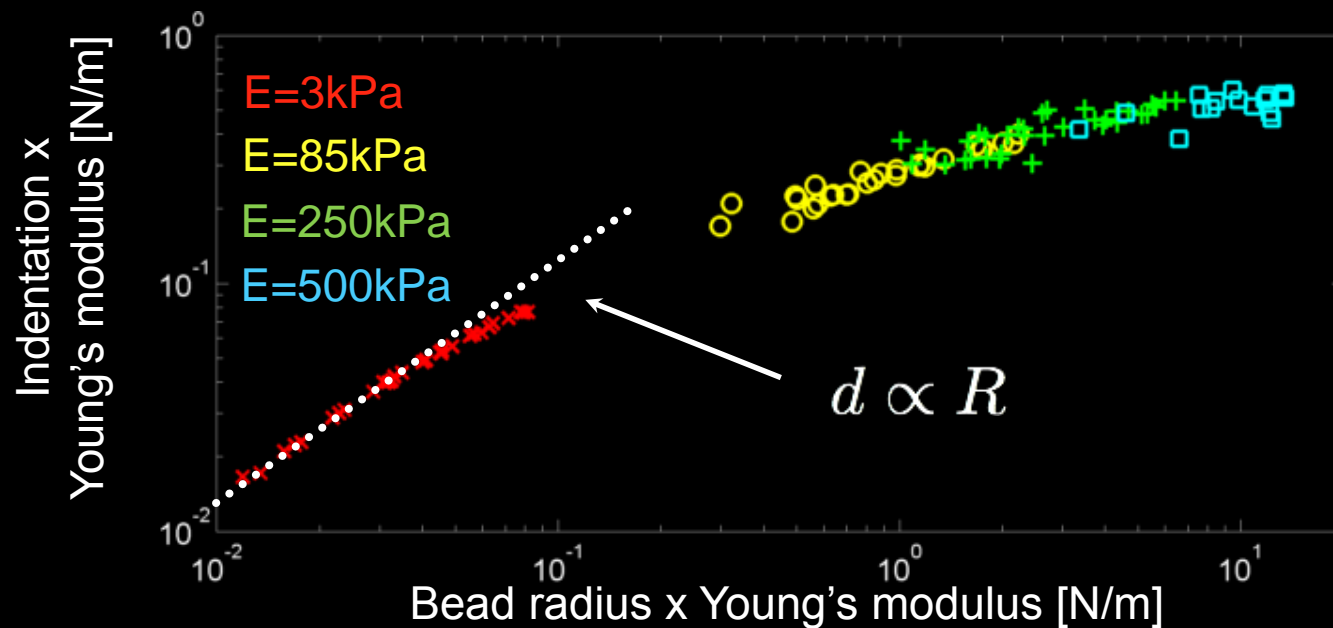
$W \sim 70\text{mN/m}$

Collapsing the data

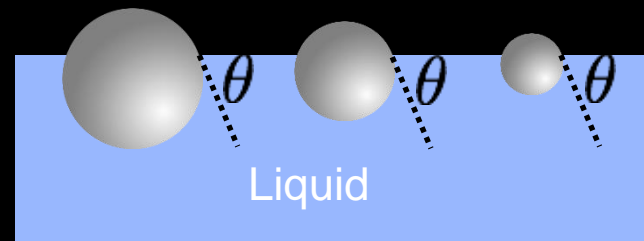
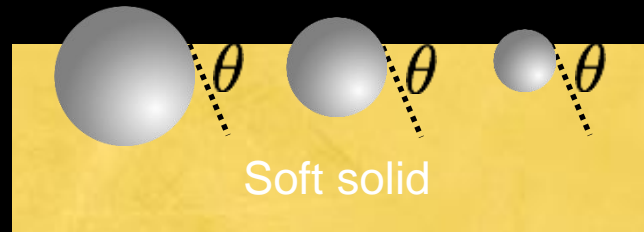
$$dE = \left(\frac{\sqrt{3}W(1 - \nu^2)}{2} \right)^{2/3} (RE)^{1/3}$$



Indentation proportional to bead radius for small beads



Indentation, $d \sim R$ implies constant contact angle



$$\gamma_{lv} \cos \theta = \gamma_{sv} - \gamma_{sl}$$

JKR plus surface tension fits data

Minimize total energy (a la Carillo/Raphael/Dobrynin 2010):

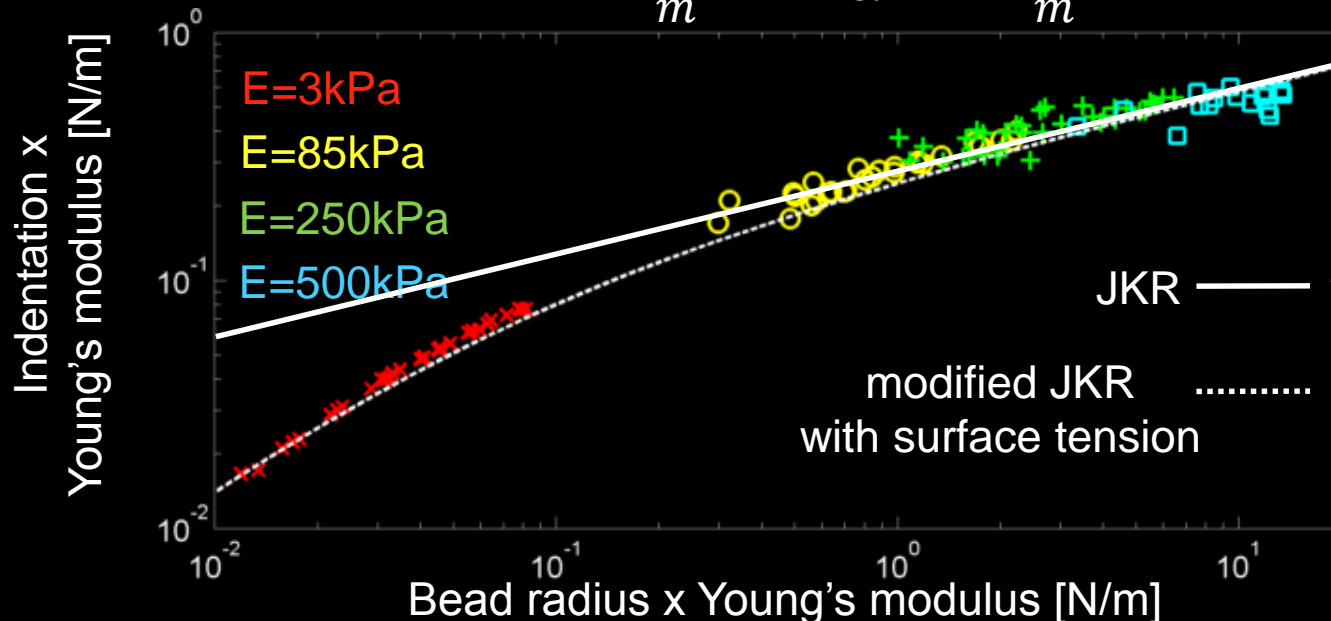
Elastic energy + Adhesion Energy + Surface Tension

(from JKR)

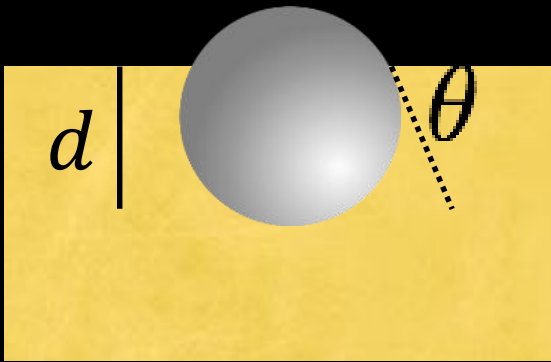
$$\Upsilon_{sv} \Delta A$$

$$W = 71 \frac{mN}{m},$$

$$\Upsilon_{sv} = 45 \frac{mN}{m}$$



Small contact limit of JKR plus surface tension...



$$d = WR/\Upsilon_{sv}$$

equivalently...

$$\Upsilon_{sv} \cos \theta = W - \Upsilon_{sv}$$

Smells like Young's Law

for $\Upsilon_{sv} = \gamma_{sv}$,

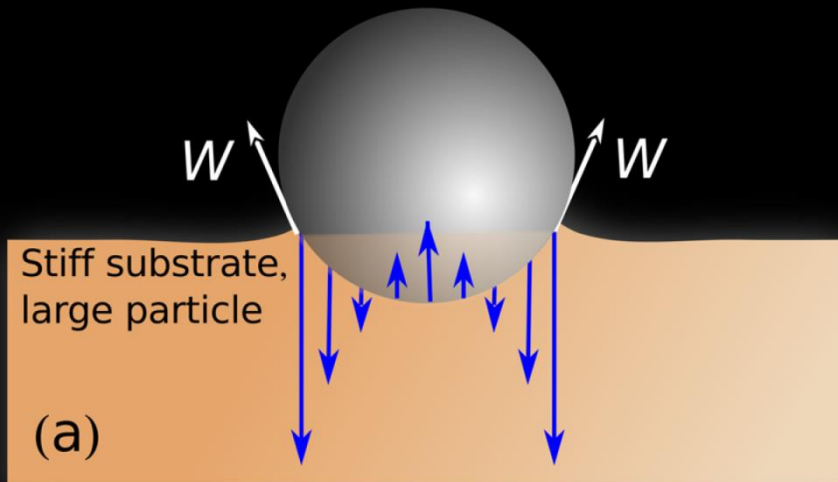
$$\gamma_{sv} \cos \theta = \gamma_{sp} - \gamma_{pv}$$

Young's Law with soft substrate in the place of the liquid

Two Regimes of Contact

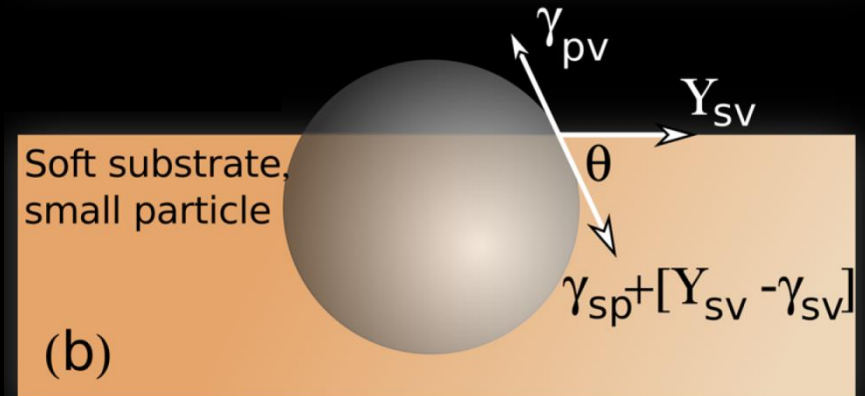
Style, Hyland, Dufresne *Nature Communications* 2013

$$a \gg Y_{sv}/E$$



JKR

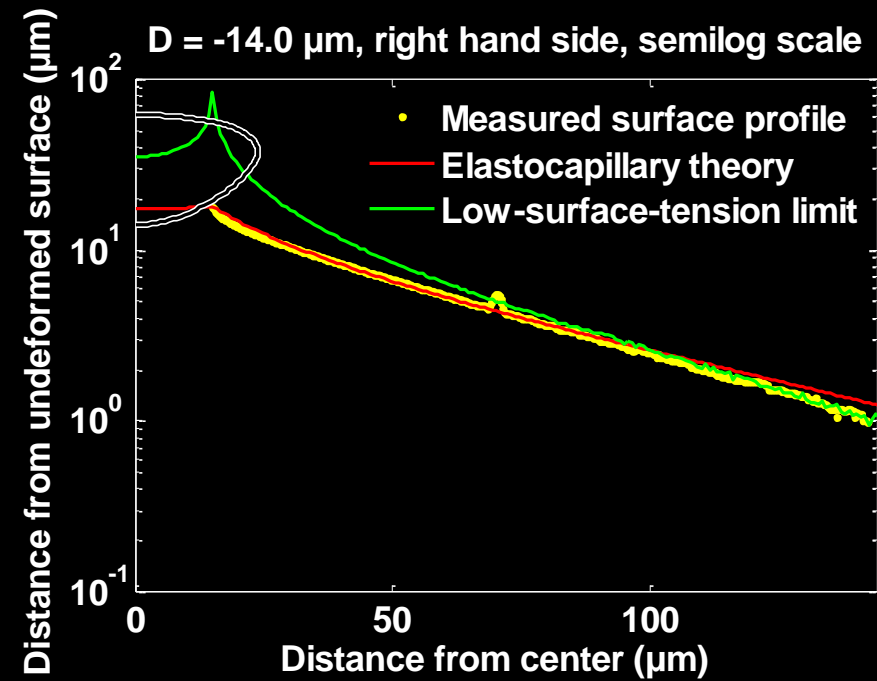
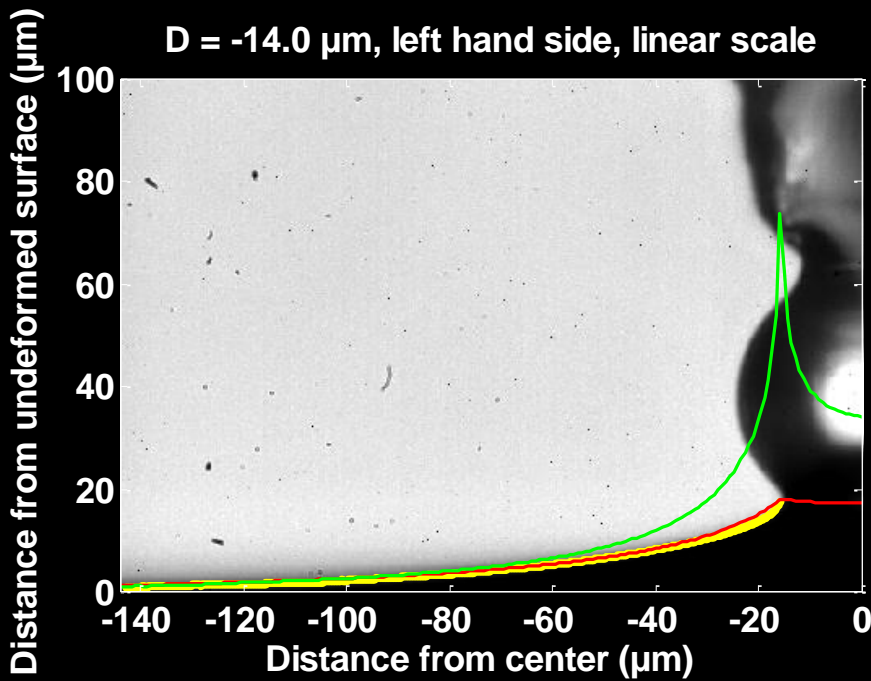
$$a \ll Y_{sv}/E$$



Young's Law

Next Step: How does surface tension modify forced-adhesion?

Kate Jensen



Key Findings

- Young's law doesn't describe wetting of small drops on soft substrates
- Young's law does describe adhesion of small particles to soft substrates
- Soft solids have a characteristic length scale $\frac{\gamma_{sv}}{E}$, below which elastic response to surface deformation is swamped by surface tension

