

Elasticity-based mechanism for collective motion in natural and artificial swarms

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Featuring: Yael Katz, Christos Ioannou, Kolbjørn Tunstrøm, Iain Couzin

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Background: Biological motivation

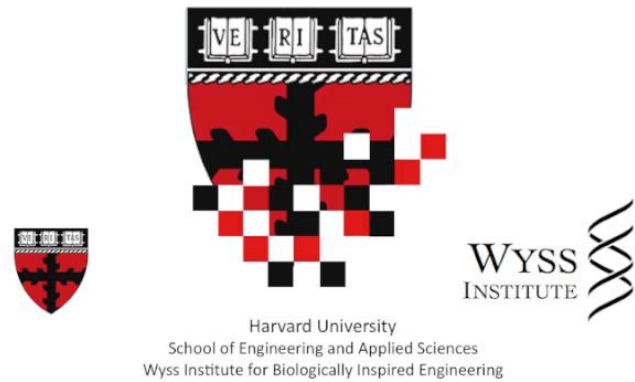
- ◆ Collective motion is observed in diverse animal species: From bacteria to humans
- ◆ Fish schools & bird flocks can have from a few to several thousand individuals
- ◆ Locust swarms can contain 10^9 insects traveling thousands of kilometers

Background: Robotics

- ❑ Groups of robots effective for:
 - Deploying sensor networks
 - Parallel tasks
 - Micro-robotics
- ❑ Control algorithms for groups of autonomous robots must be:
 - **Decentralized, Scalable, Robust**
- ❑ Additional constrains
 - Small processing power & bandwidth
 - Limited communication (nearby neighbors, line-of-sight, direct contact)



Self-Organizing Systems Research Group



Background: Random thoughts

❑ My first works on Collective Motion (before it was cool):

- *Phase transitions in self-driven many-particle systems and related non-equilibrium models: A network approach.*
M. Aldana and C. Huepe – *J. Stat. Phys.* 112 (1-2), 135-153 (**2003**).
- *Intermittency and clustering in a system of self-driven particles*
C. Huepe & M. Aldana – *Phys. Rev. Lett.* 92 (16)168701 (**2004**).

❑ Current research on active matter (non extensive list):

- Hydrodynamic descriptions
 - self-propelled agents on fluids or fluids of self-propelled agents
 - fluids with *aligning spins* or *self-energized advection*
 - active complex fluids, - etc.
- Agent-based
 - explicit/implicit alignment,
 - escape-pursuit, - etc.

❑ Most fascinating for me: how active systems self-organize

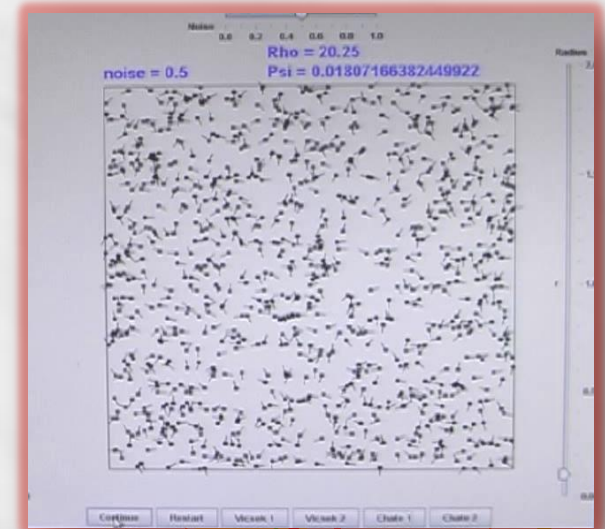
Background: Vicsek model

❑ Motivation

- ◆ Non-equilibrium self-organizing statistical system
- ◆ Energy injected at smallest scales

❑ The Vicsek model (1995 results)

$$\theta_i(t + \Delta t) = \text{Ang} \left\{ \sum_{Z_i(t)} \vec{V}el[\theta_j(t)] \right\} + \eta \xi_i(t)$$
$$\vec{x}_i(t + \Delta t) = \vec{x}_i(t) + \vec{V}el[\theta_i(t)] \Delta t.$$

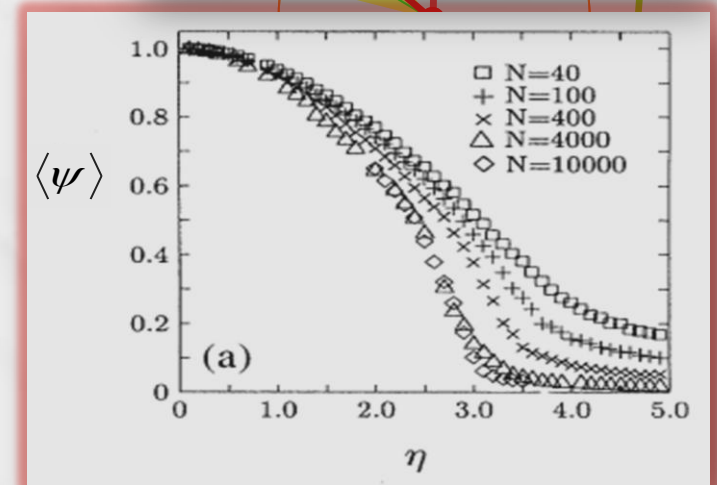


❑ Order parameter (alignment/magnetization)

$$\psi(t) = \frac{1}{Ns} \left| \sum_{i=1}^N \vec{V}el[\theta_i(t)] \right|$$

❑ Main Result

- ◆ “Novel type” of phase transition





Dr. Yael Katz



Prof. Kolbjørn Tunstrøm



Prof. Christos Ioannou



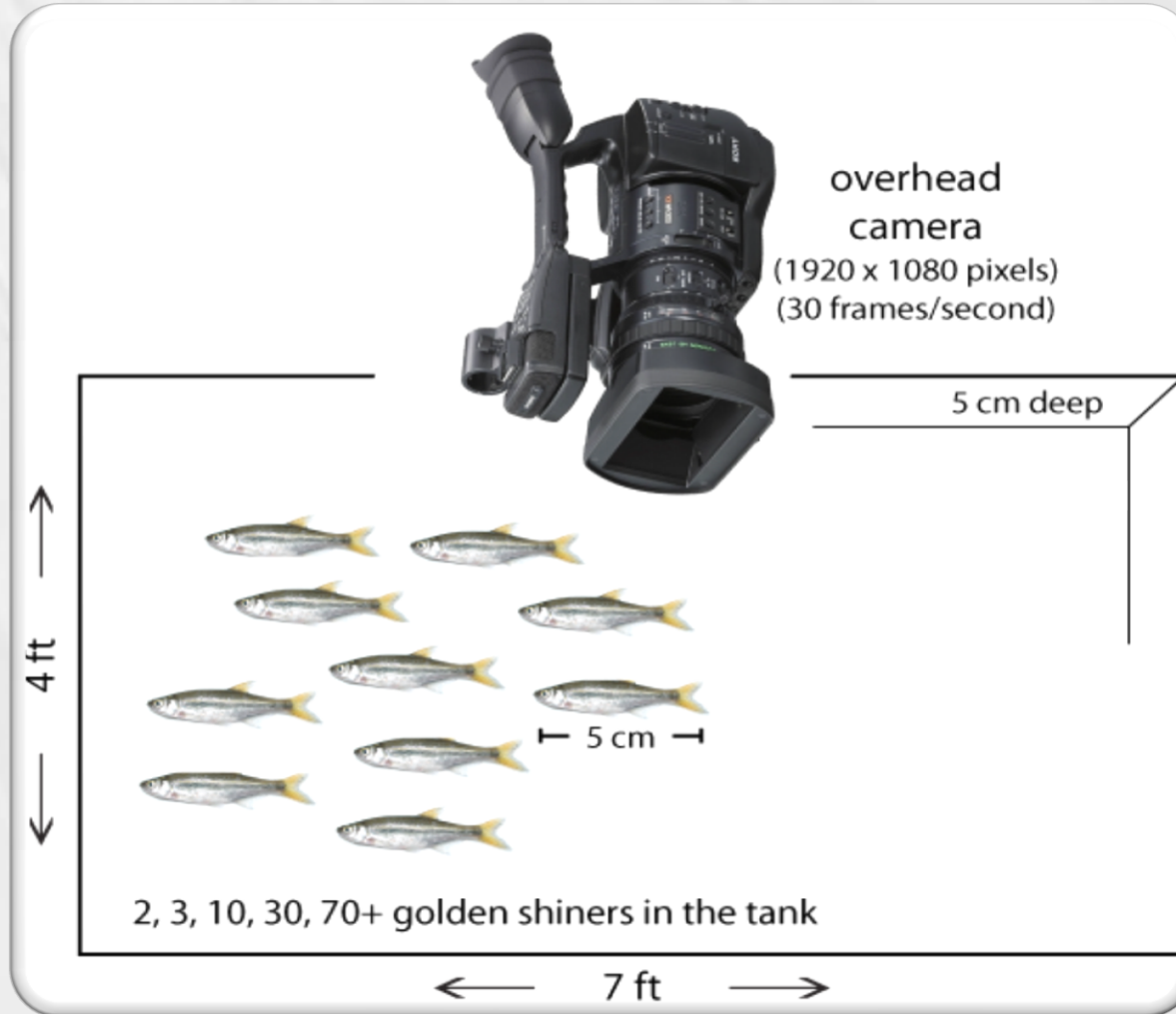
Prof. Iain Couzin

- Data-driven analysis of fish schooling experiments
(Inspiration for this work)

Publications:

- *Collective states, multistability and transitional behavior in schooling fish*, PLoS Computational Biology Feb. 2013 [9(2), e1002915]
- *Inferring the structure and dynamics of interactions in schooling fish*, PNAS July 2011 [doi:10.1073/pnas.1107583108]

- Experimental System



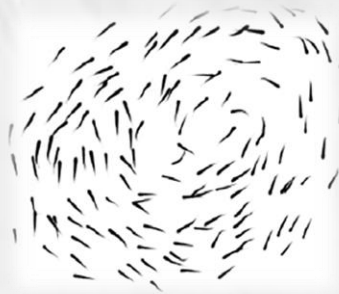
2, 3, 10, 30, 70+ golden shiners in the tank

- 1000 fish dynamics

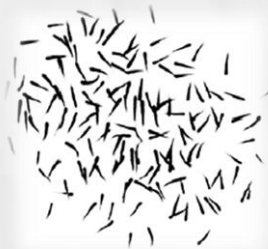
Polarized state



Rotating state



Swarming state



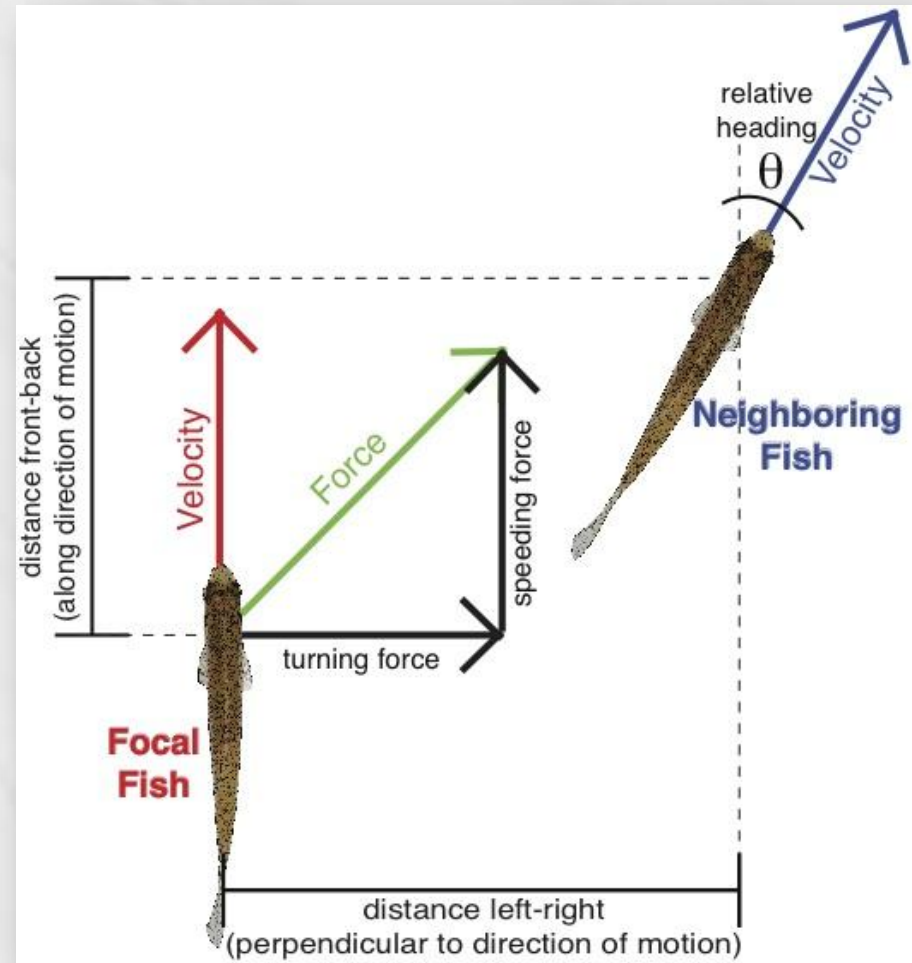
■ The two-fish system

❖ Effective-force approach

- ◆ Measure mean *effective (social) forces* on 2-fish & 3-fish systems
- ◆ Use *large dataset*: 14 experiments of 56 minutes each

❖ Goals

- ◆ "*Model-free*" approach on clear mathematical grounds
- ◆ *Gain intuition* over possible non-standard interactions

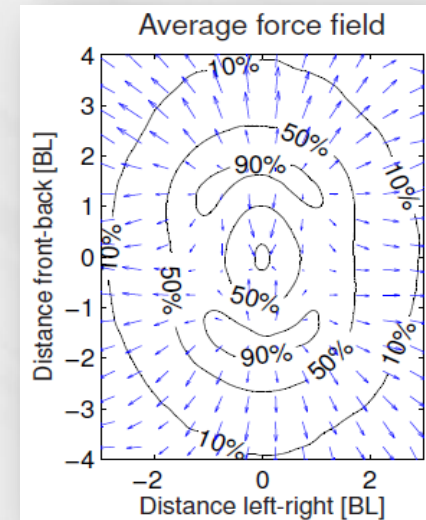
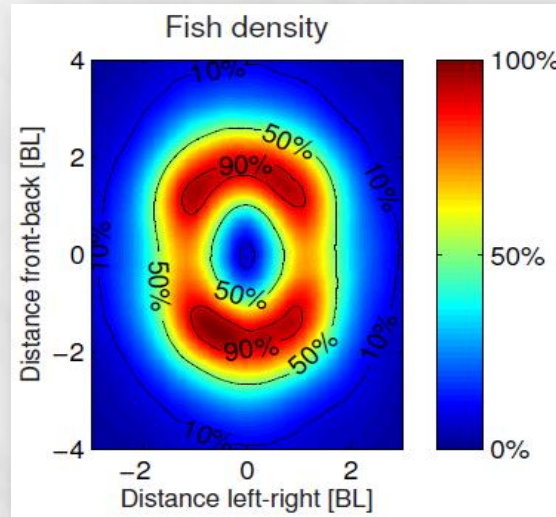


The two-fish system

□ Mean effective forces (social) on 2-fish & 3-fish systems

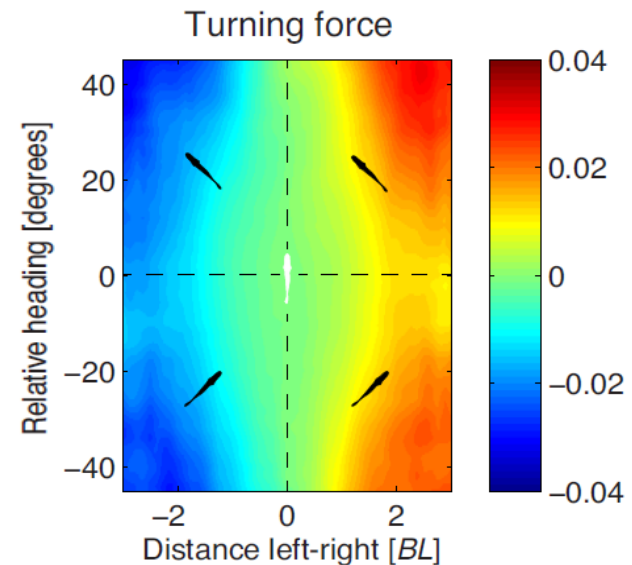
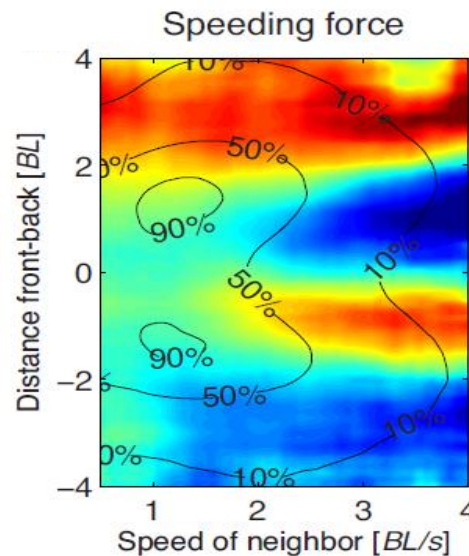
□ 14 experiments of 56 minutes each

• Zero force \Leftrightarrow high density



• Higher speed \Rightarrow larger front-back distance

• Higher left-right distance \Rightarrow larger turning force



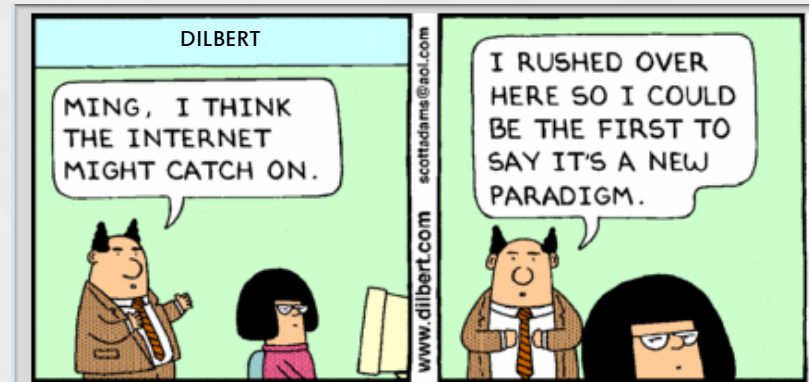
□ Elasticity-based mechanism for collective motion in active solids



Dr. Eliseo Ferrante



Prof. Ali-Emre Turgut



Publications:

- *Elasticity-Based Mechanism for the Collective Motion of Self-Propelled Particles with Springlike Interactions: A Model System for Natural and Artificial Swarms.* **Phys. Rev. Lett.**, 111, 268302 (2013)
- *Collective motion dynamics of active solids and active crystals.* **New J. Phys.** 15 095011 (2013)

- ❖ **Consider limit opposite to Vicsek model**
 - ◆ No explicit aligning interaction
 - ◆ Interactions based on relative positions (not relative orientations)
 - ◆ Fixed interacting neighbors
(Following Mermin-Wagner theorem condition)

- ❖ **Active Elastic Sheet model**
 - ◆ Only attraction-repulsion forces
 - ◆ Spring-mass model of elastic membrane
 - ◆ Overdamped & self-propelled dynamics

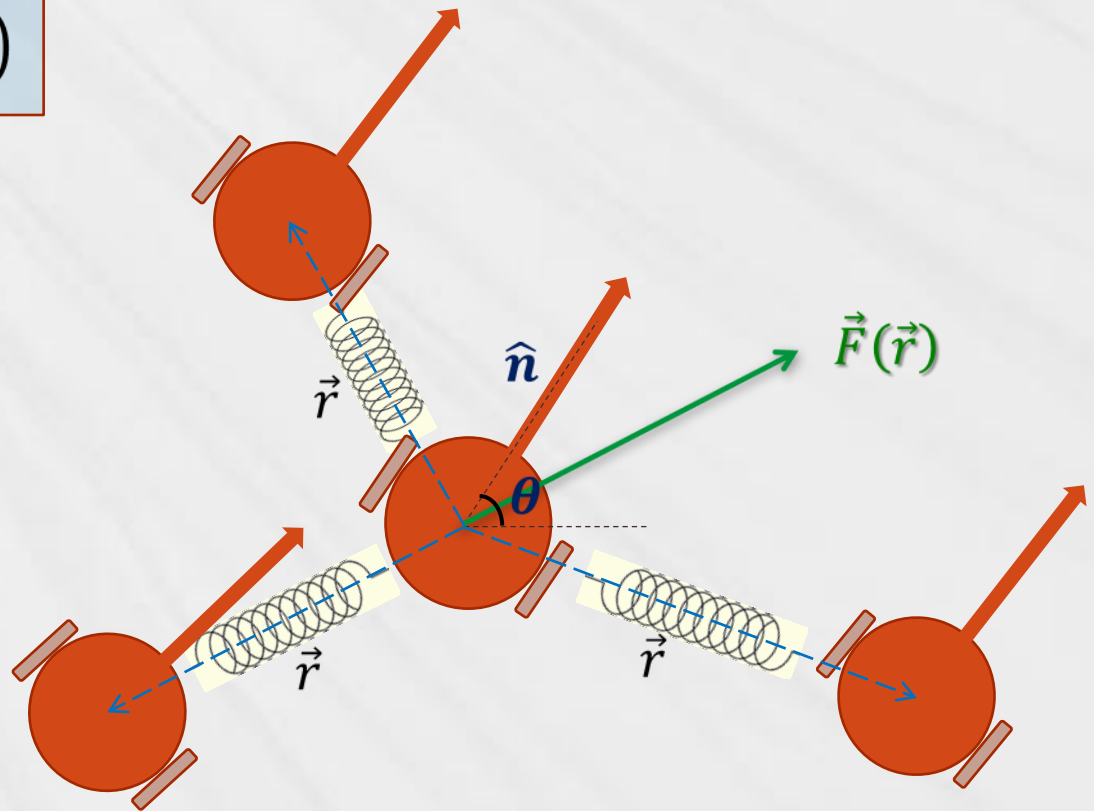
– Toy model with only attraction-repulsion interactions:
Active Elastic Sheet model

✘ Forward speed:

$$v = v_0 + \alpha(\vec{F} \cdot \hat{n})$$

□ Turning rate:

$$\dot{\theta} = \beta(\vec{F} \cdot \hat{n}_\perp)$$



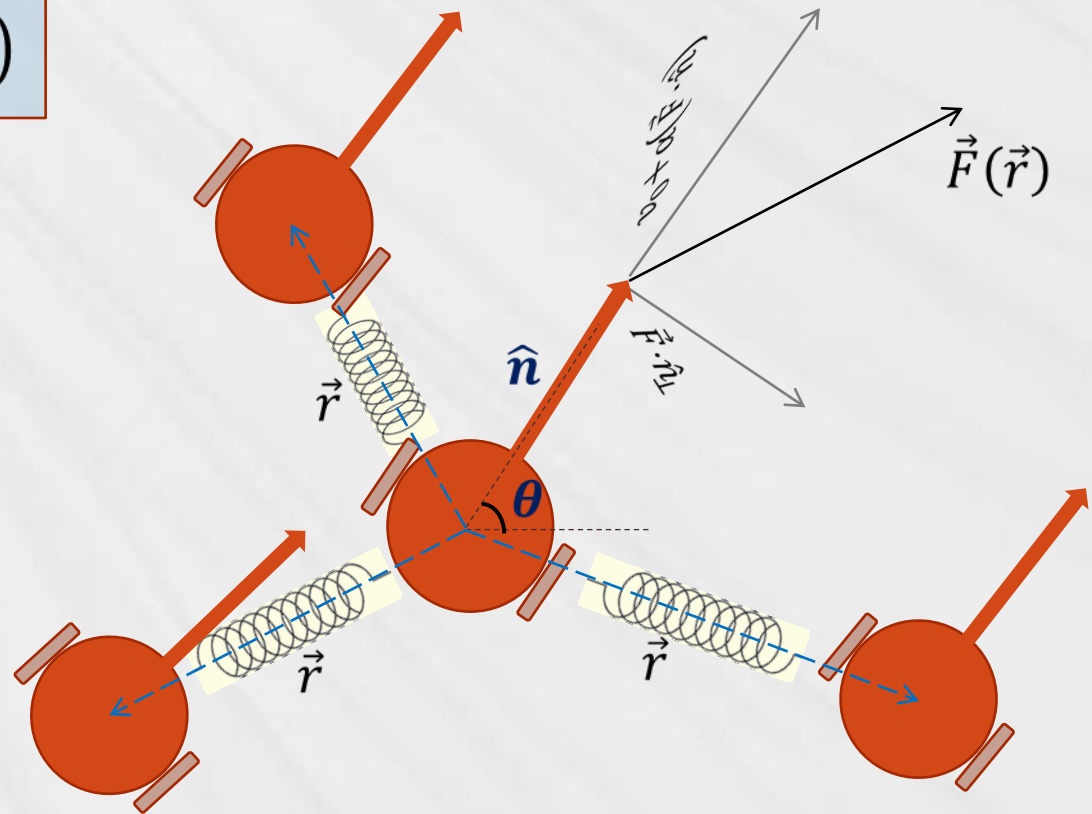
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– Toy model with only attraction-repulsion interactions:
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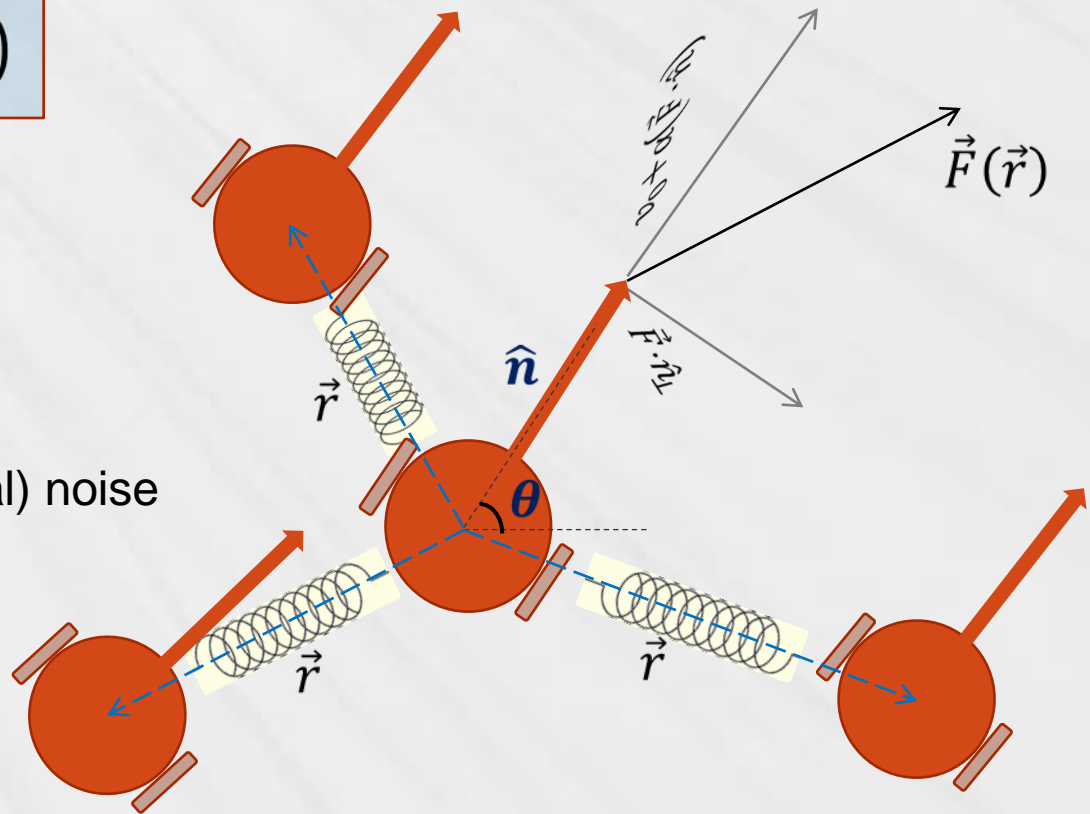
$$v = v_0 + \alpha(\vec{F} \cdot \hat{n})$$

□ Turning rate:

$$\dot{\theta} = \beta(\vec{F} \cdot \hat{n}_\perp)$$

- Sensing (vectorial) noise

- Actuation (angular) noise



– Toy model with only attraction-repulsion interactions: Active Elastic Sheet model

- Linear spring-like forces over each agent:

$$\vec{F} = \sum_{i \in S} \vec{f}_i + D_r \hat{\xi}_r = - \sum_{i \in S} k (|\vec{r}_i| - l) \frac{\vec{r}_i}{|\vec{r}_i|} + D_r \hat{\xi}_r$$

S : set of all agents connected to focal agent

\vec{r}_i : position of i with respect to focal agent

k & l : spring constants & natural lengths

- Equations of motion for focal agent

$$\text{Forward speed : } v = v_0 + \alpha(\vec{F} \cdot \hat{n})$$

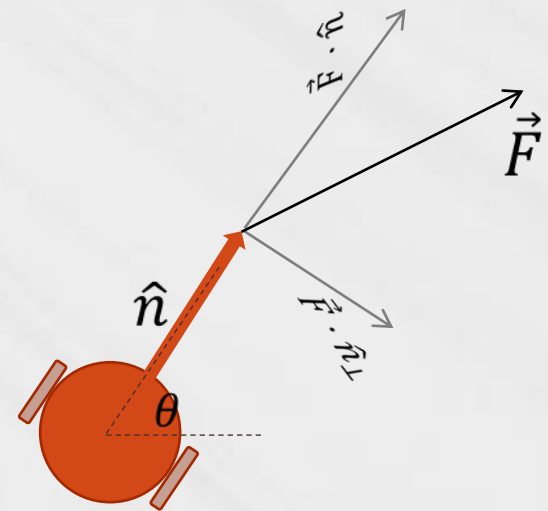
$$\text{Turning rate : } \dot{\theta} = \beta(\vec{F} \cdot \hat{n}_\perp) + D_\theta \xi_\theta$$

v_0 : preferred speed

α, β : inverse damping coefficients

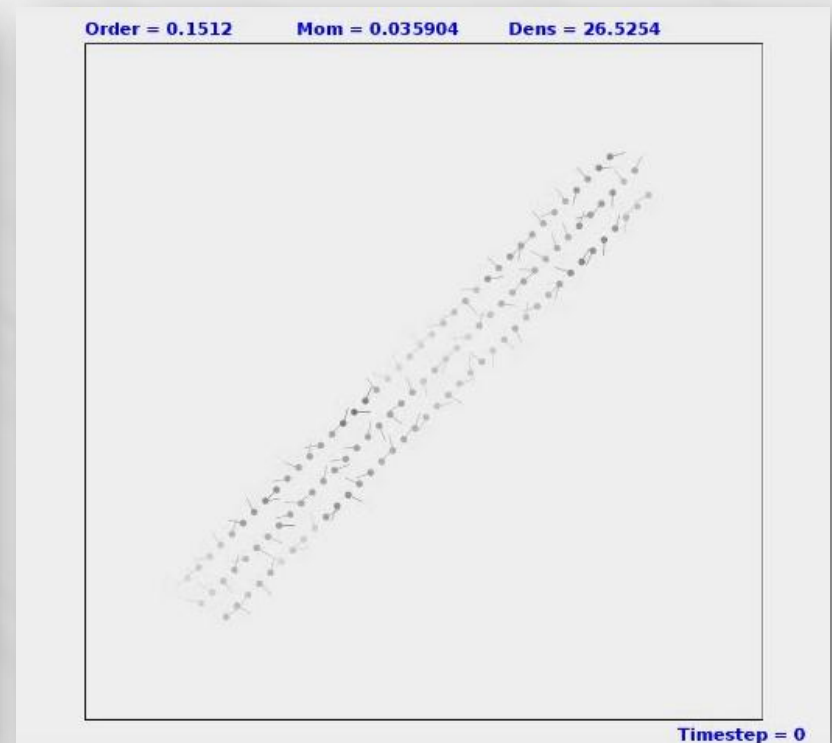
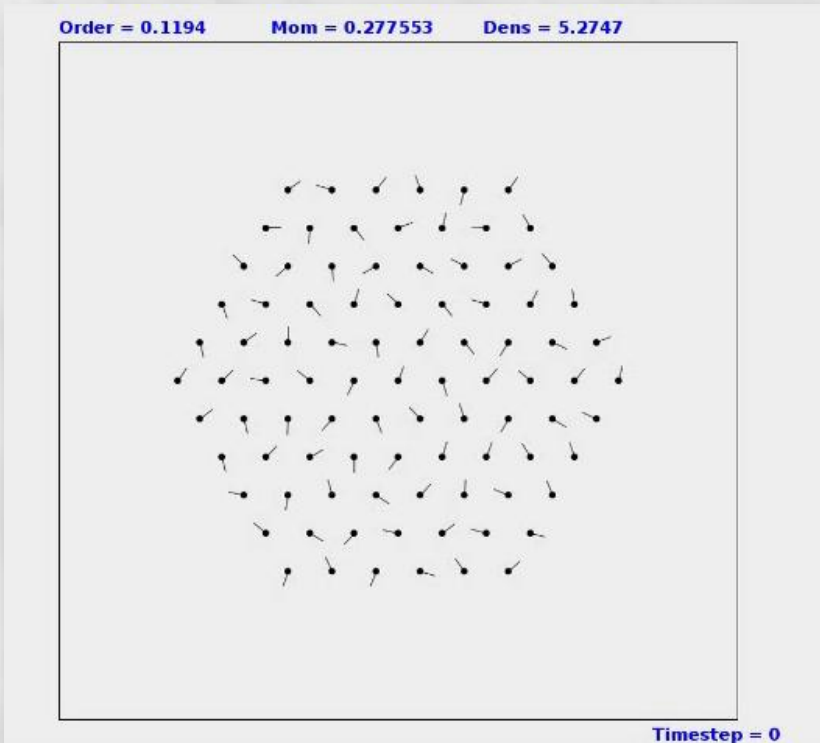
\hat{n} : unit heading vector

ξ : random variables



- Simulation dynamics

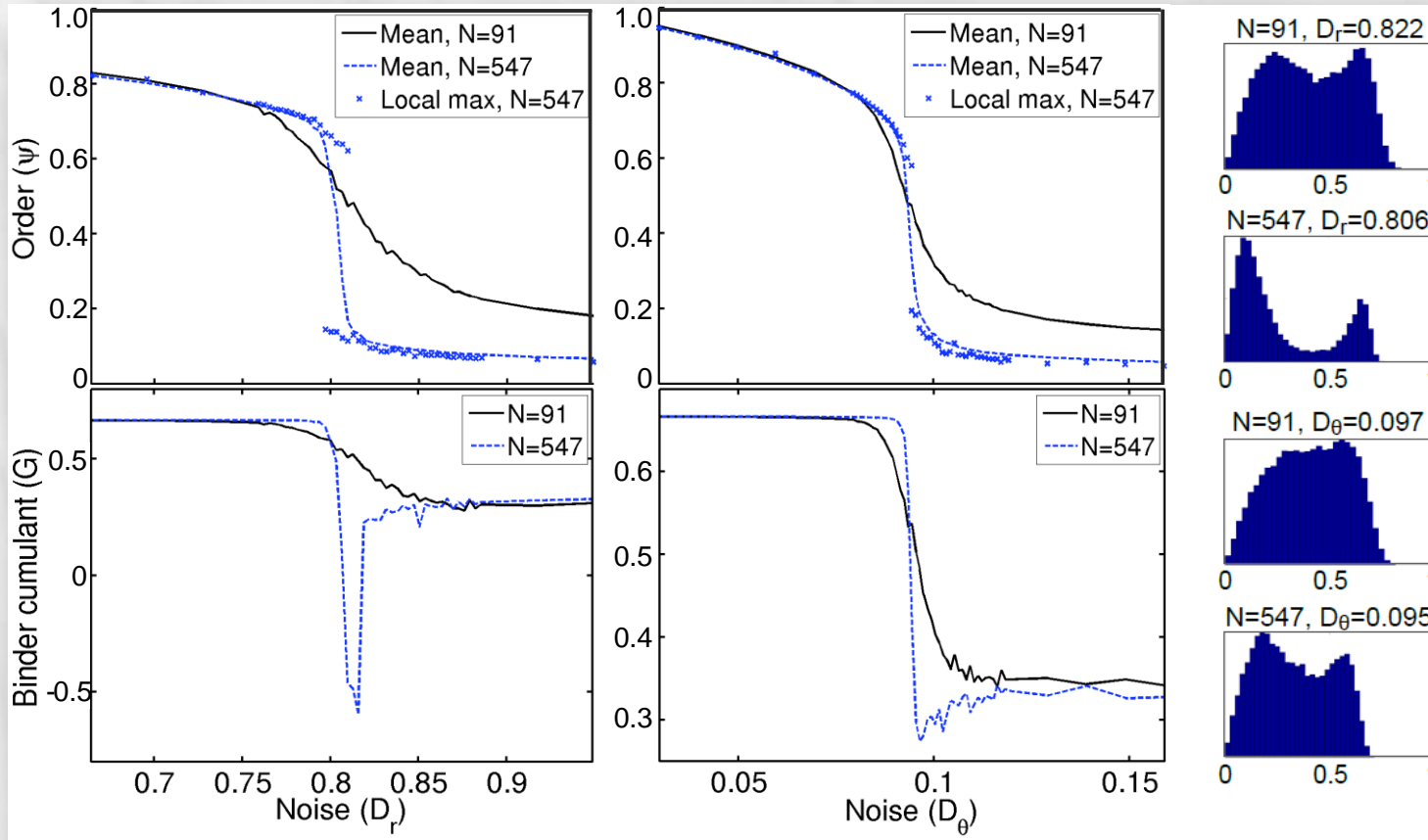
- *Spring-mass* model of elastic sheet
- Active crystals:



- ❑ **Rotational** ordered state (metastable)
- ❑ Ordered states develop below critical noise

Translational ordered state (preferred heading direction)

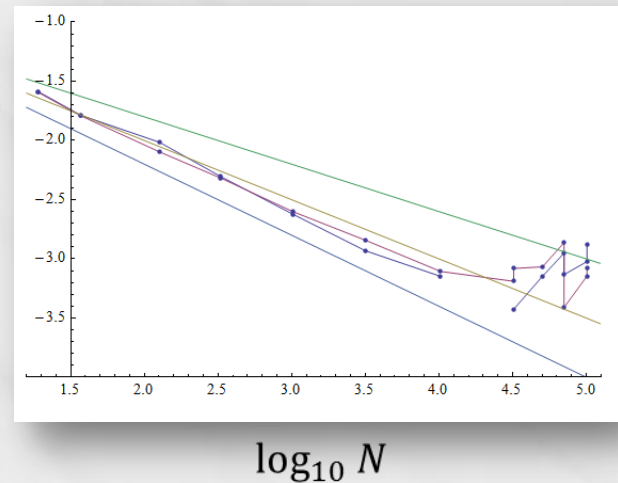
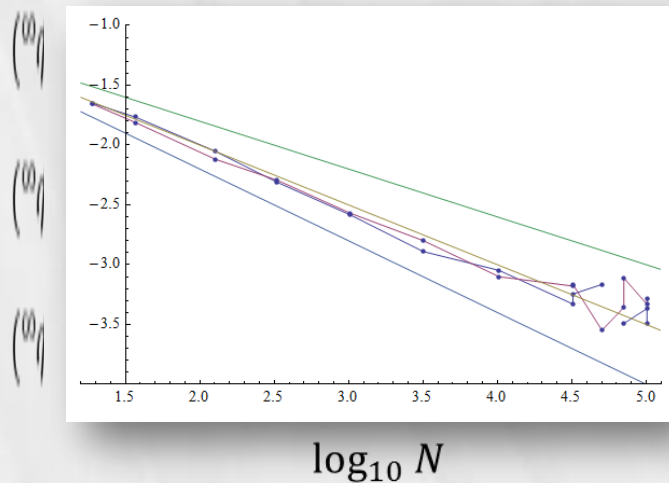
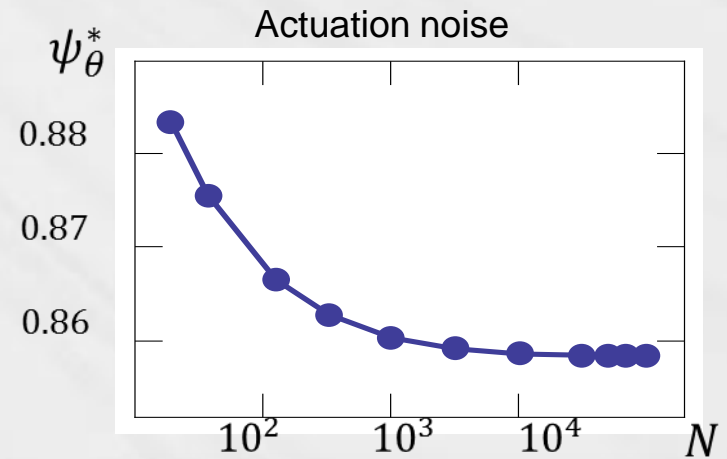
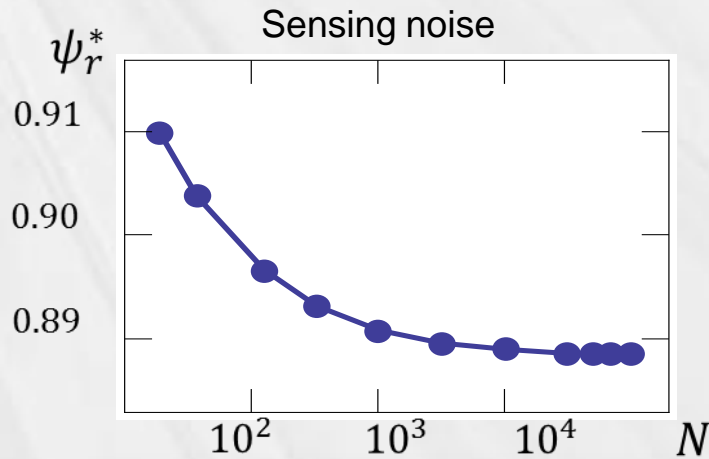
Phase transition



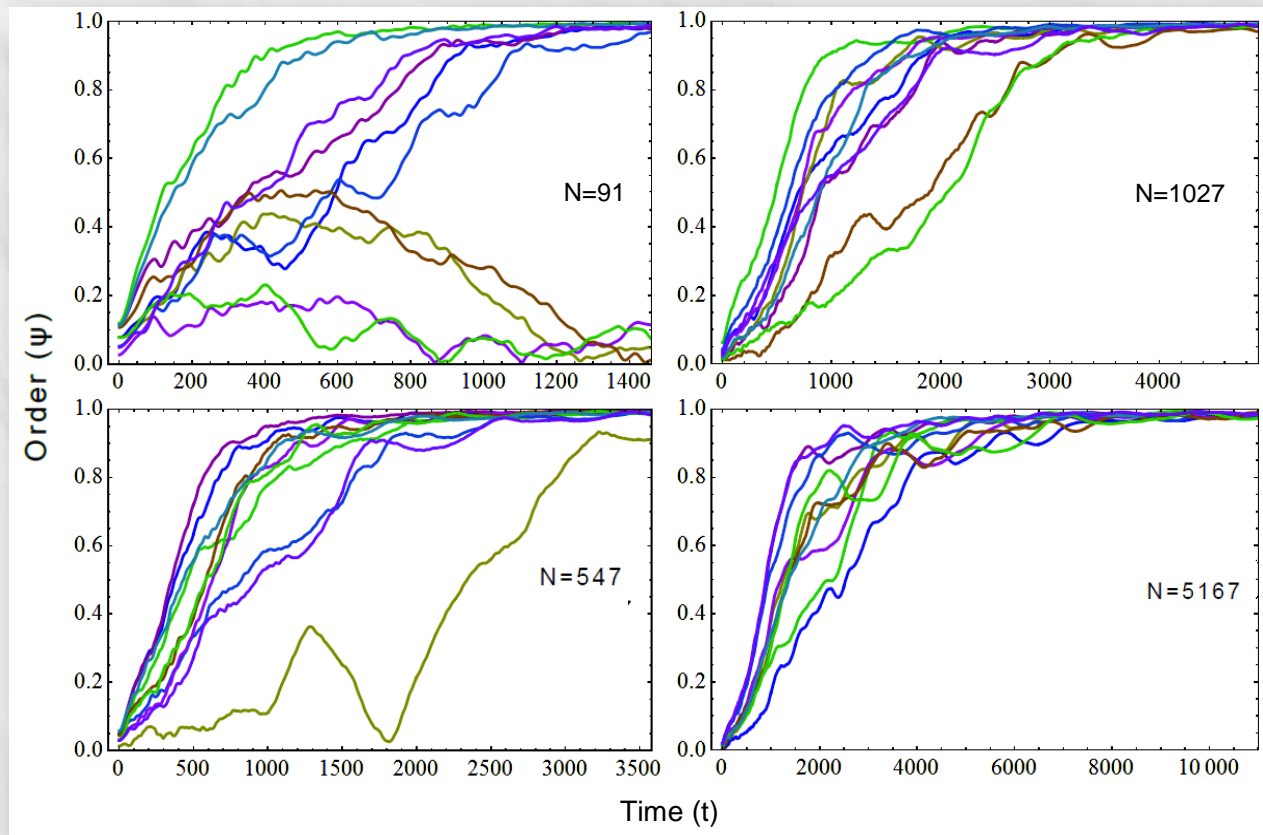
- Alignment order parameter, Binder cumulant, and bimodal distributions:
→ First order phase transition

Long-Range Order

- Long-range order verified through numerical finite scaling analysis
- Approximate scaling at $\sim 60\%$ of critical noise: $\psi = C N^{-1/2} + \psi_\infty$



Convergence dynamics



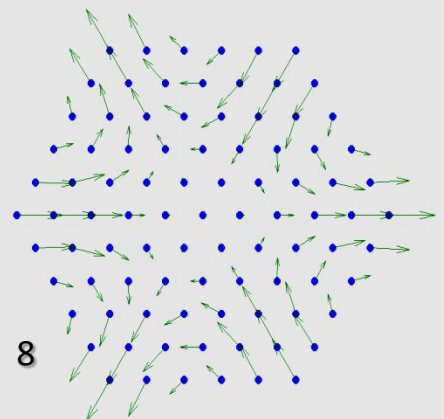
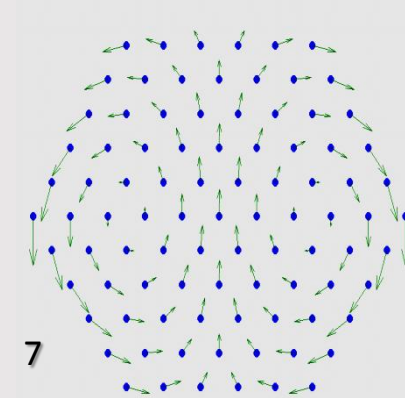
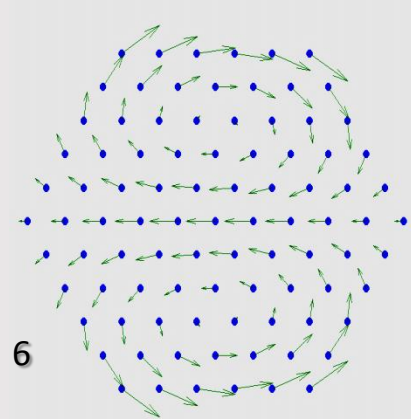
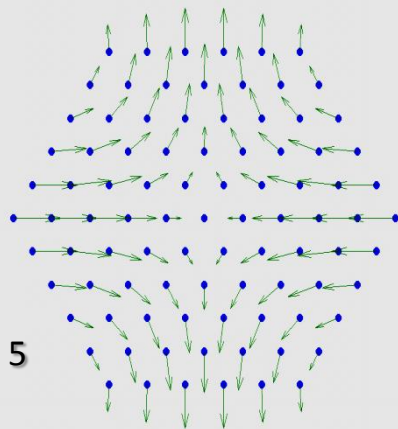
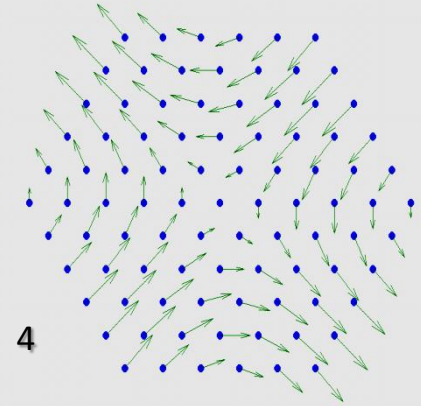
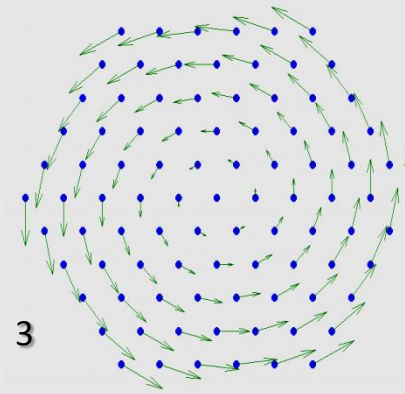
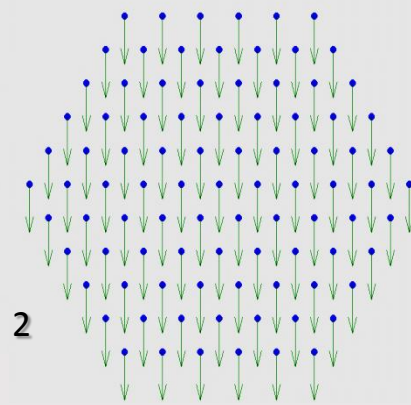
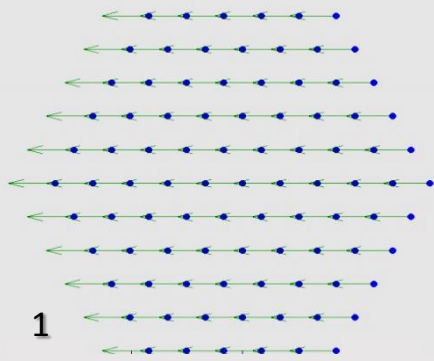
- ❑ Variability reflects complex dynamical landscape
- ❑ Very different amounts of time spent in metastable states.
- ❑ Typical convergence time grows here slower than $\sim N^{1/2}$

- Self-organizing Mechanism

- ❖ Why does the system self-organize into common heading direction?
 - ◆ Despite:
 - ◆ no ferromagnetic-like interaction
 - ◆ no clear consensus-based dynamics
 - ◆ fixed interacting neighbors
 - ◆ Elasticity-based mechanism?

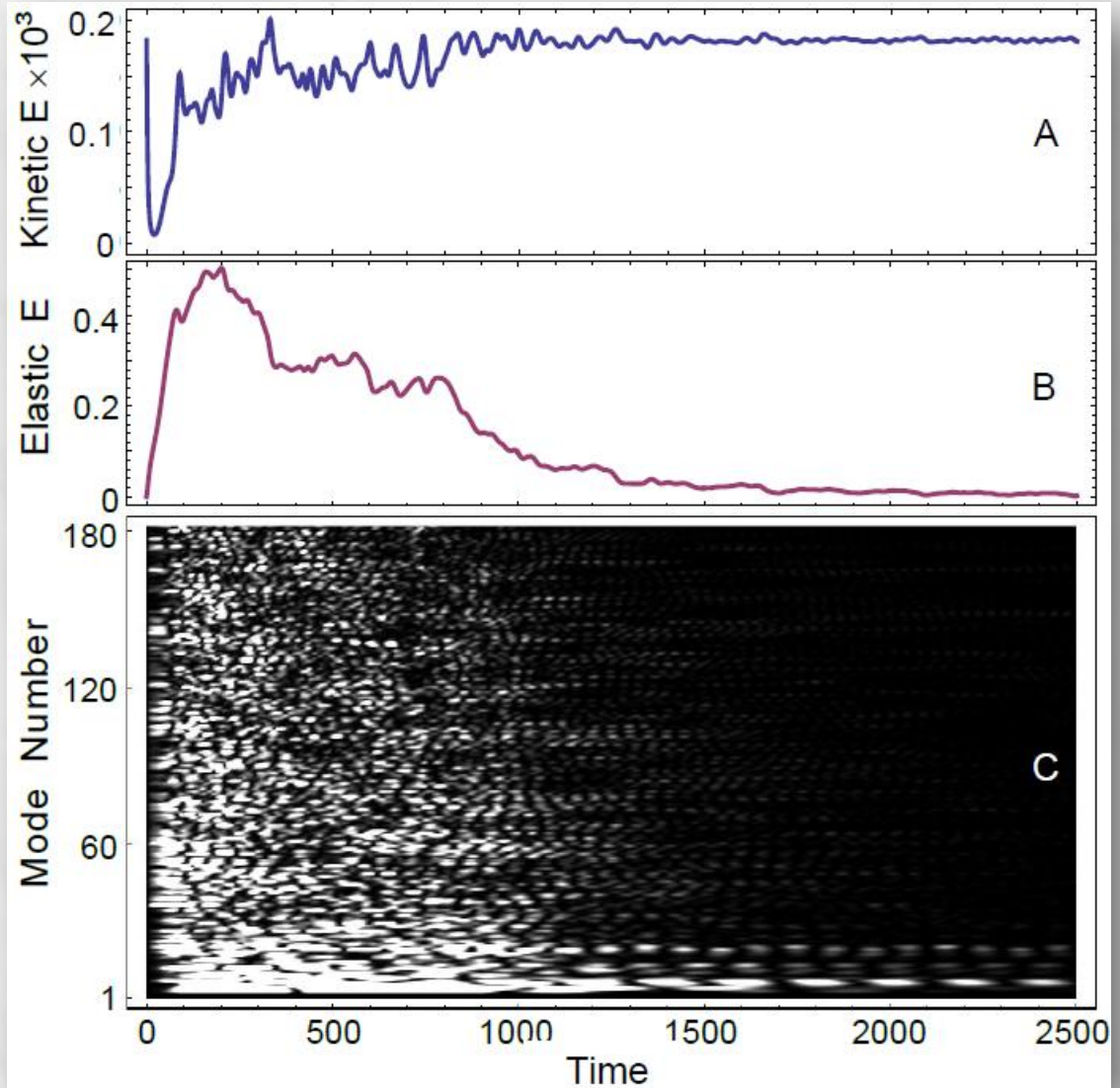
Unveiling the self-organizing mechanism

- We first decompose into normal elastic modes
 - E.g.: First 8 modes of hexagonal structure:



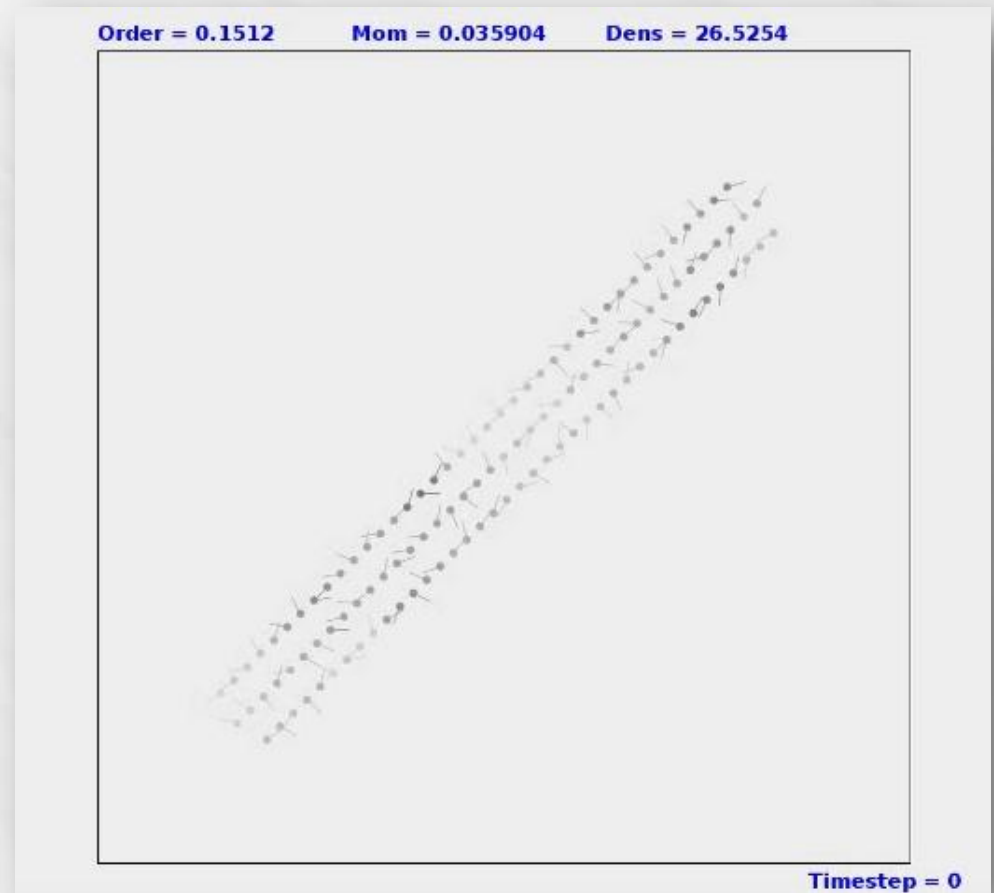
Spectral decomposition of elastic energy vs. time

- ❑ Self-propelled agents **steer away** from high-energy modes
- ❑ Energy **cascades** towards **low-energy** modes
- ❑ Agents **self-organize** into growing regions of coherent motion



Going back to see the modes

- ❑ Darker agents: higher local alignment
- ❑ Visible cascade of bending modes
- ❑ Preferred direction of self-organized motion



Linear stability analysis (no noise)

Linear continuous approximation

$$\dot{u}_x = \alpha F_x$$

$$\dot{u}_y = v_0 \phi$$

$$\dot{\phi} = \beta F_y$$

$$F_x = (\lambda + \dots)$$

$$F_y = (\lambda + \dots)$$

Stability analysis in Fourier space

$$\text{Stability matrix} = \begin{pmatrix} -k^2 \alpha (\lambda + 3\mu) & -k^2 \alpha \\ 0 & -k^2 \beta (\lambda + \mu) \end{pmatrix}$$

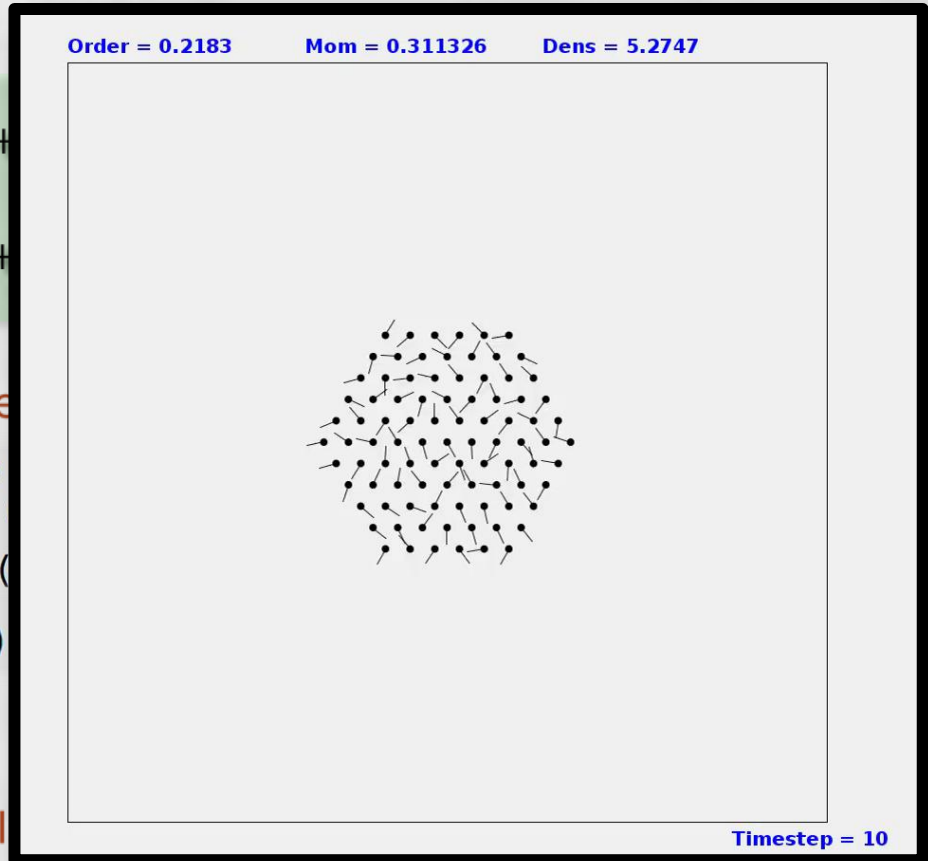
$$\text{Characteristic polynomial: } \Lambda^3 + \alpha (\lambda + 3\mu) k^2 \Lambda = 0$$

Routh's stability criterion: C0 C3 – C1 C2

Is the fixed speed case ($\alpha = 0$) all

$$\text{Characteristic polynomial: } \Lambda^3 + \beta v_0 (\lambda + 3\mu) k^2 \Lambda = 0$$

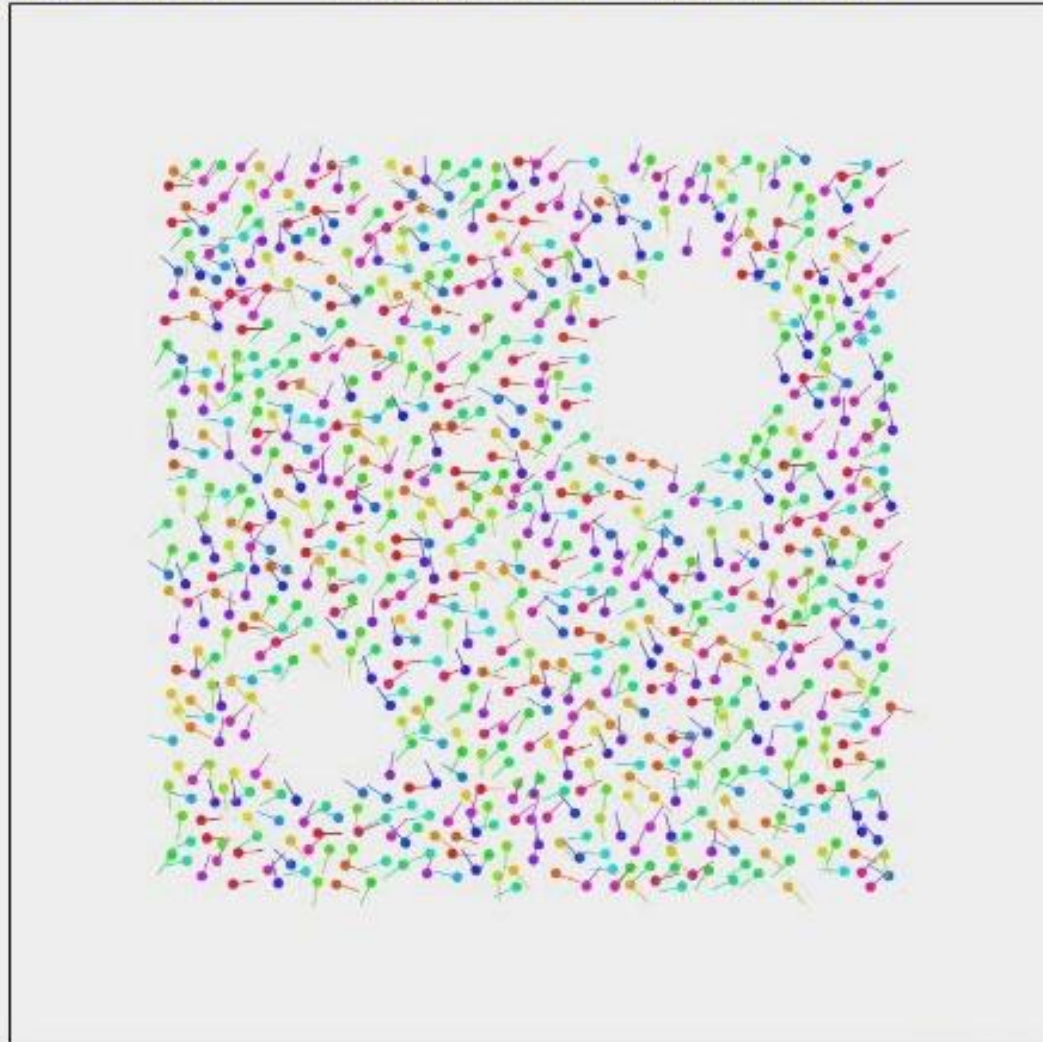
$$\Rightarrow \Lambda = 0 \quad \text{or} \quad \Lambda = \pm \sqrt{-k^2 v_0 \beta \lambda - 3 k^2 v_0 \beta \mu} \in i \mathbb{R}$$



Order = 0.0101

Mom = 0.003354

Dens = 13.2255



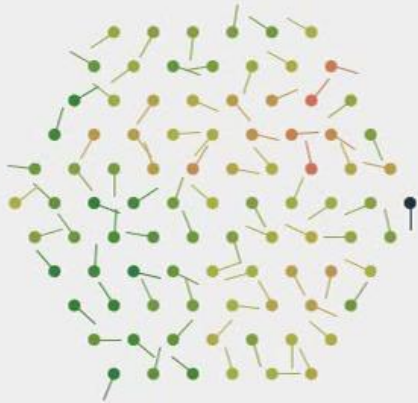
Timestep = 0

- ❑ Active solid setup
 - Random positions
 - Neighbors connected by linear springs

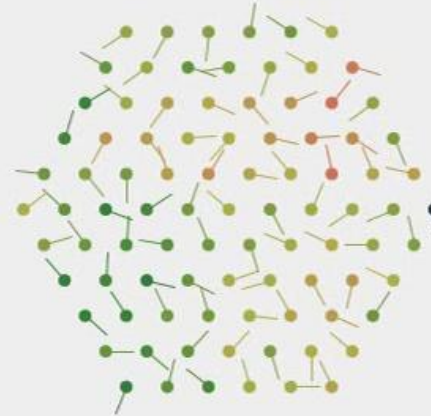
- ❑ Agents here colored by angle
 - Coherent regions grow
 - Energy flows to lower modes

Parameter exploration

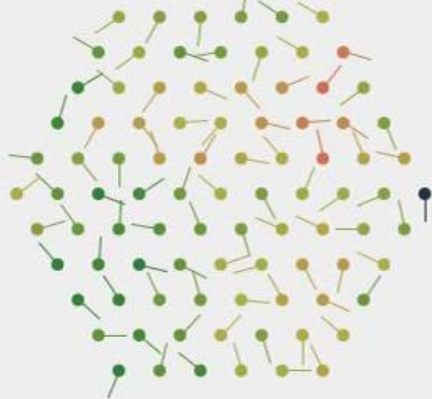
Hi α | Low v_0



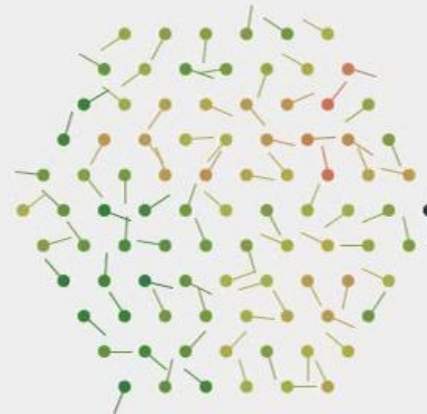
Low α



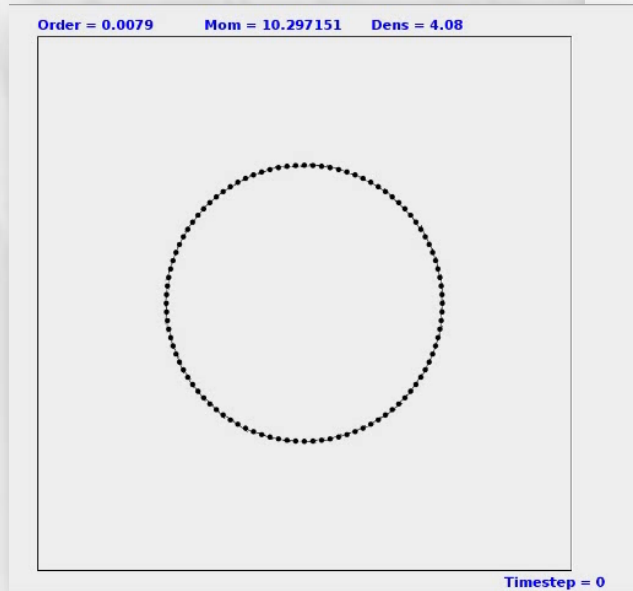
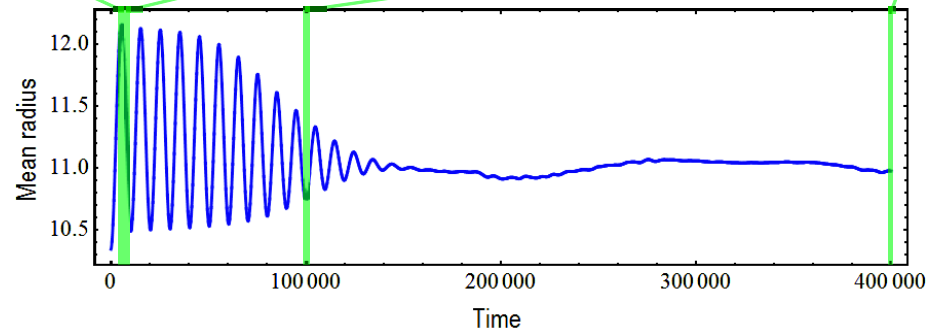
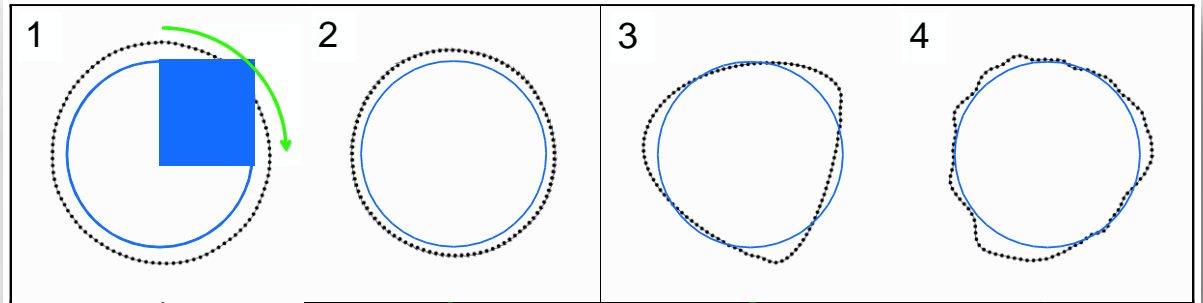
Hi β



Low β

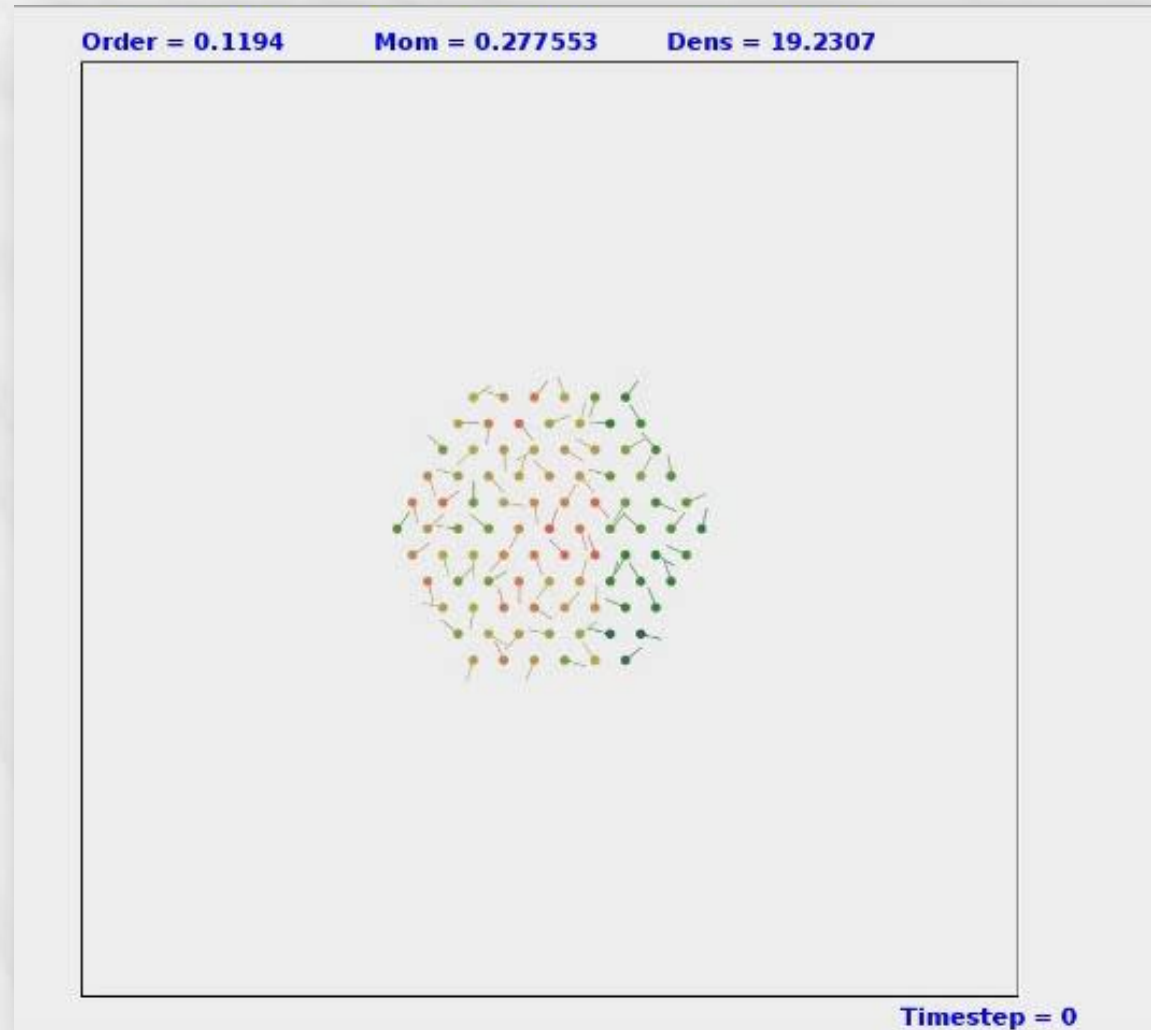


Ring dynamics



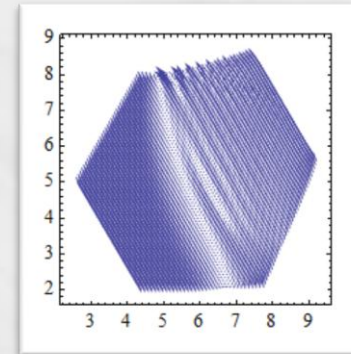
Heterogeneous self-propulsion speeds

- ❑ Randomly assigned self-propulsion speeds
- ❑ System self-organizes in much longer time

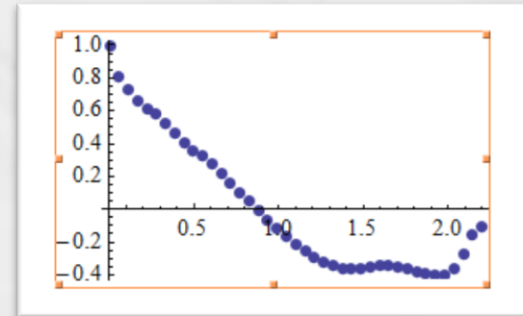


Final teaser – work in progress... lots to do...

- **Dispersion relation**
(with Naomi Oppenheimer – U of C)



- **Alternative explanation for flocking scale-free correlations**



- **Other elasticity-based self-organizing systems?**
- **Give me agents with different α , β , v_0 , k , l_0 and I'll build you a bicycle... or an 'amoeba'**

Recap & *The End*

- Active elastic membrane (crystal/solid structure)
- No alignment interactions required
- No mixing required
- Elasticity-based mechanism for self-organization
- Rich idealized limit-case model system

... *Fin*