

BLACK HOLES, QUANTUM INFORMATION TRANSFER, AND HILBERT-SPACE NETWORKS

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Black Holes and Information
KITP

May 21, 2012

Refs: SBG 1108.2015, 1201.1037;
SBG and Y. Shi, 1205.xxxx

The information “paradox,” in a nutshell:

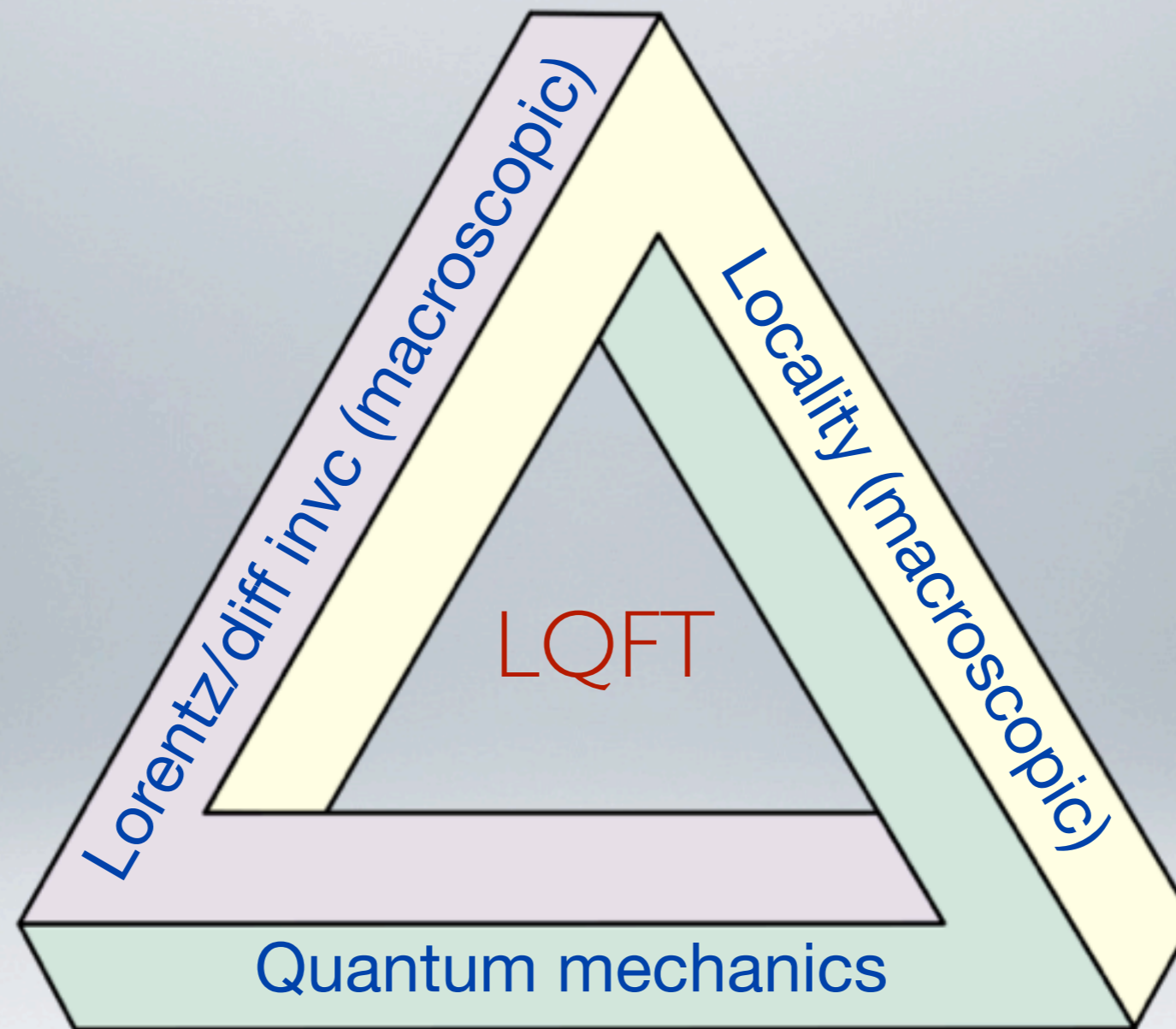
Information cast into a black hole

- can't get out
 - can't be destroyed
 - can't be left in remnant
- locality
QM; energy conserv.
catastrophic instabilities

Stupid mistake?

Harbinger of new physics?

A VIEWPOINT: REPRESENTS A FUNDAMENTAL CONFLICT



- QM, LI: hard to modify (consistency, observation). Locality?
- It's not about singularities, renormalizability? **Long distance.**

A SEEMINGLY SIMILAR CRISES:

Classical atom

CM fails here



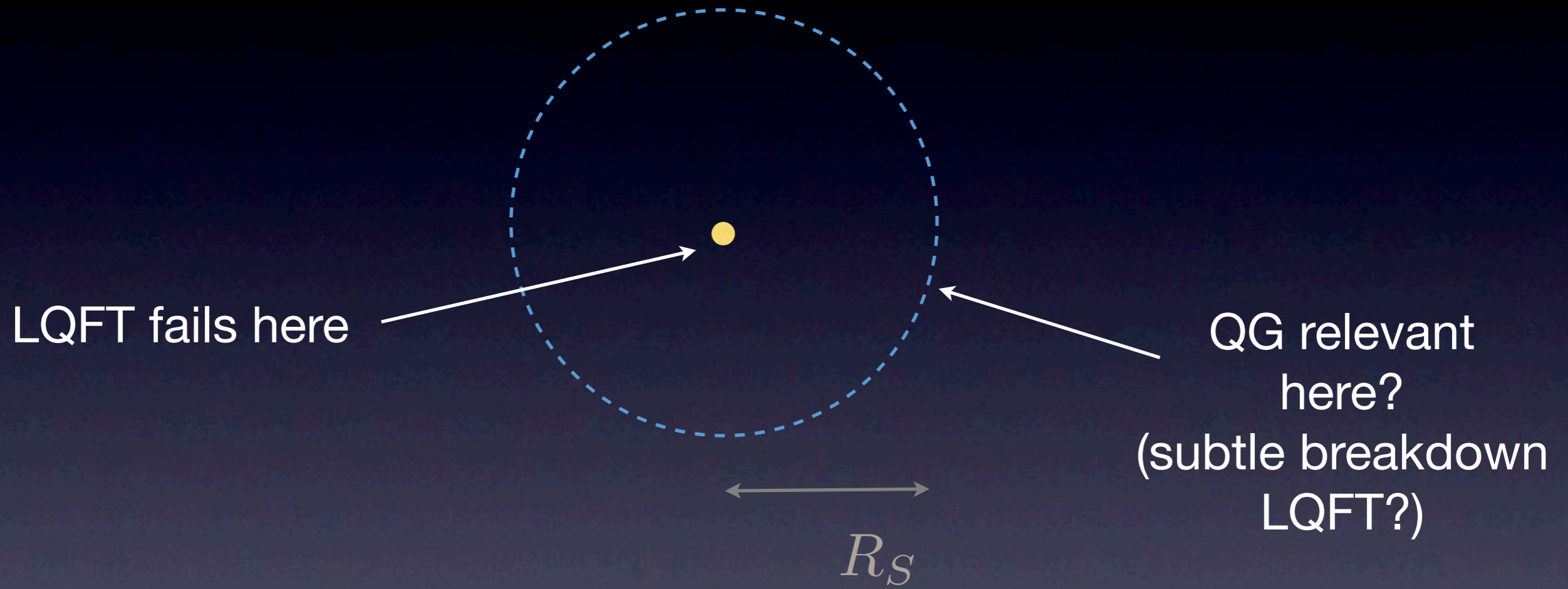
QM takes over here
(CM irrelevant)

New physics was needed:

Uncertainty principle
Wave mechanics...

A SEEMINGLY SIMILAR CRISES:

Black hole



New physics ~~was~~ needed:

is

~~Uncertainty principle~~

~~Wave mechanics...~~

???

“CLASSICAL INSTABILITY PARADOX”



“BLACK HOLE INFORMATION PARADOX”

The information problem is likely an important guide to understanding new principles/mechanisms.

(As was the stability problem of the atom)

What other guides do we have?

AdS/CFT??

Microstates??

workshop/conference

Cosmological comparisons; tests?

workshop/conference

Locality/ local observables

e.g. hep-th/0512200,
w/ Marolf and Hartle

S-matrix

Erice lectures: 1105.2036

Amplitude magic??

previous workshop;
nonperturbative?

Correspondence

Quantum info. transfer from BHs

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Comments on AdS/CFT (More detail: 1105.6359)

- many regard as resolution -

A question: can we recover from the boundary theory a sufficiently fine-grained bulk description, e.g., of evolution of a small ($\ll R$) BH, and of observers falling in?

Need:

$$\begin{array}{ccc} & M : \mathcal{H}_{\text{bulk}} \rightarrow \mathcal{H}_{\text{bdy}} & \text{interacting} \\ & \nearrow \quad \text{1-1, unitary} \quad \nwarrow & \\ U_{\text{bulk}} = M^\dagger U M & & U \end{array}$$

Unitary bulk evolution?

Approaches:

(~) local bulk observables

challenge in QG

much discussion at workshop

general “relational” approach: e.g SBG, Marolf,
Hartle hep-th/0512200; used in inflation

S-matrix (flat space limit)

Problem: construct scattering states from
boundary data; extract fine-grained S-matrix

not much discussion at workshop

there are challenges

(more discussion:
1106.3553 w/ M. Gary)

More general approaches:

1. Investigate correspondence boundary
2. Quantum information transfer
3. The gravitational S-matrix

Correspondence

Existing theory: LQFT, semiclassical background

various proposals: planckian curvature, modified/string uncertainty, modified dispersion, holographic bounds ...

Configurations: $\phi_{x,k} \phi_{x',k'} |0\rangle$
(min uncertainty wavepackets)

Where description fails?



$\phi_{x,k}$



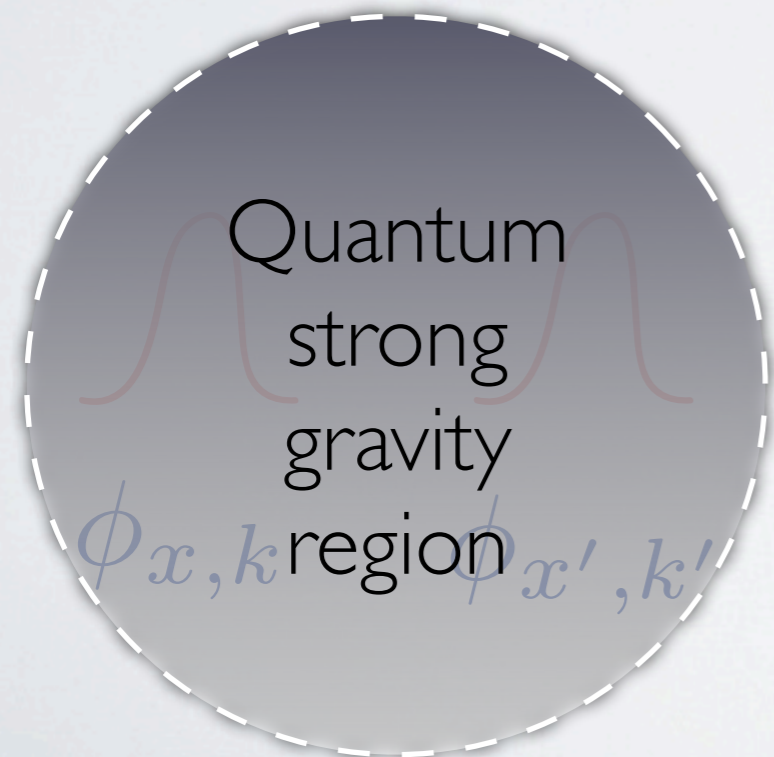
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Where description fails?

$$|x - x'|^{D-3} < (\hbar G) |k - k'|$$

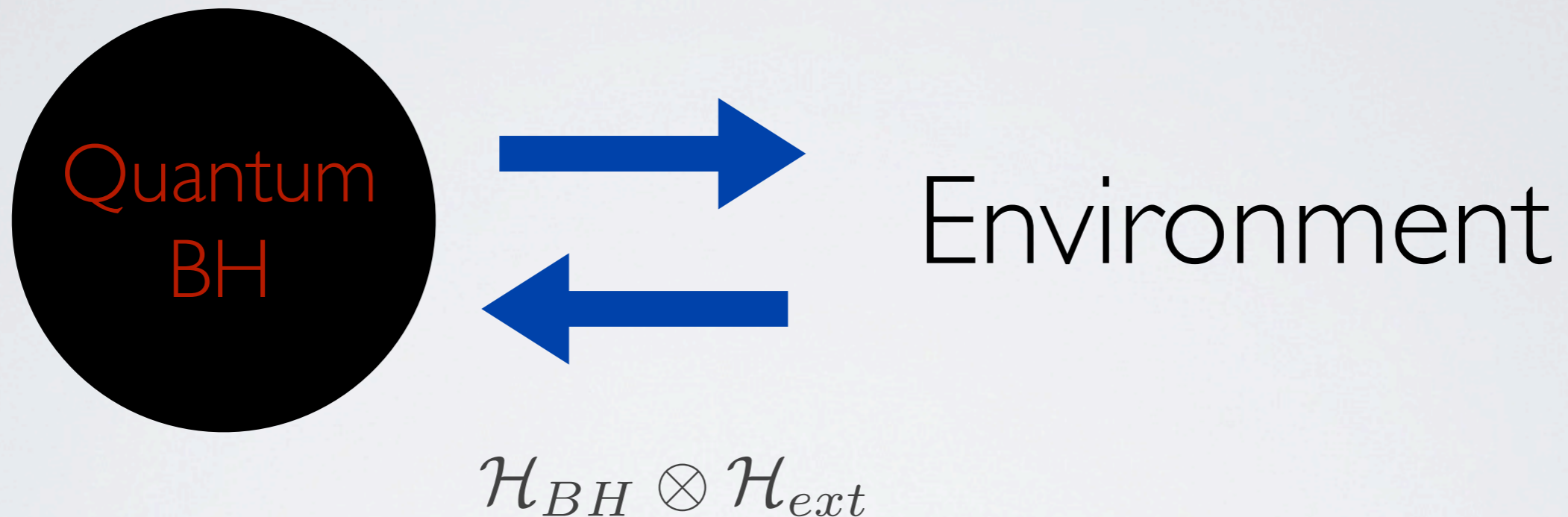
“locality bound”

~ Heisenberg microscope
(multiparticle generalizations)

We seek guides to this regime.

Assume: quantum mechanics (take seriously)

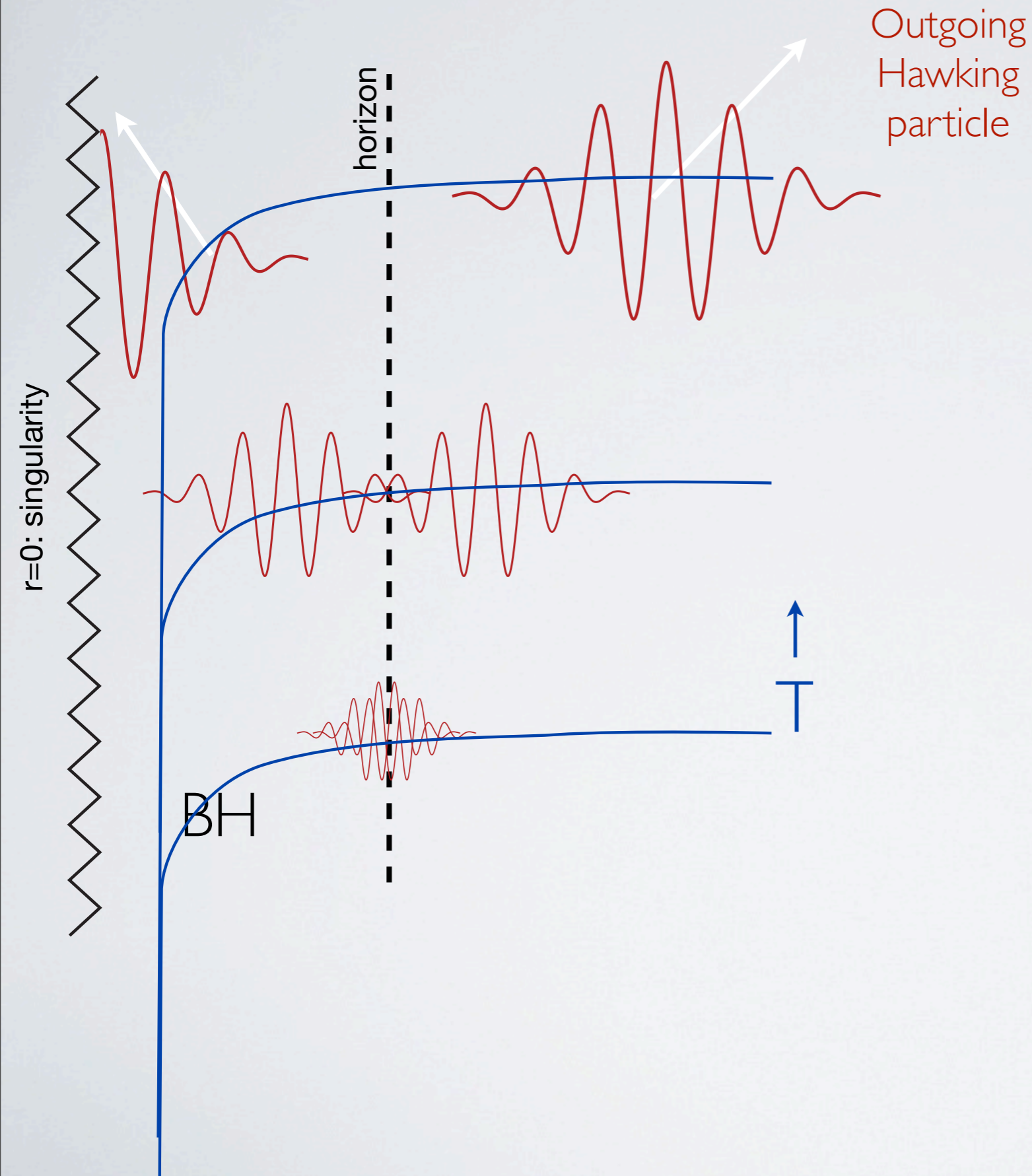
Also assume: quantum subsystems



Information exchange; unitary

“Effective quantum theory;” and more
(recall: “QM+locality+Poincare \Rightarrow QFT”)

Another guide: "as close as possible" to LQFT



nice-slice description
 -- sharpens tension

X^+ , X^- : null, Kruskal

$$X^+(X^- + e^{-2T} X^+) = R_c^2$$

“Hawking state”

(explicit in 2d:
SBG and Nelson, 1992)

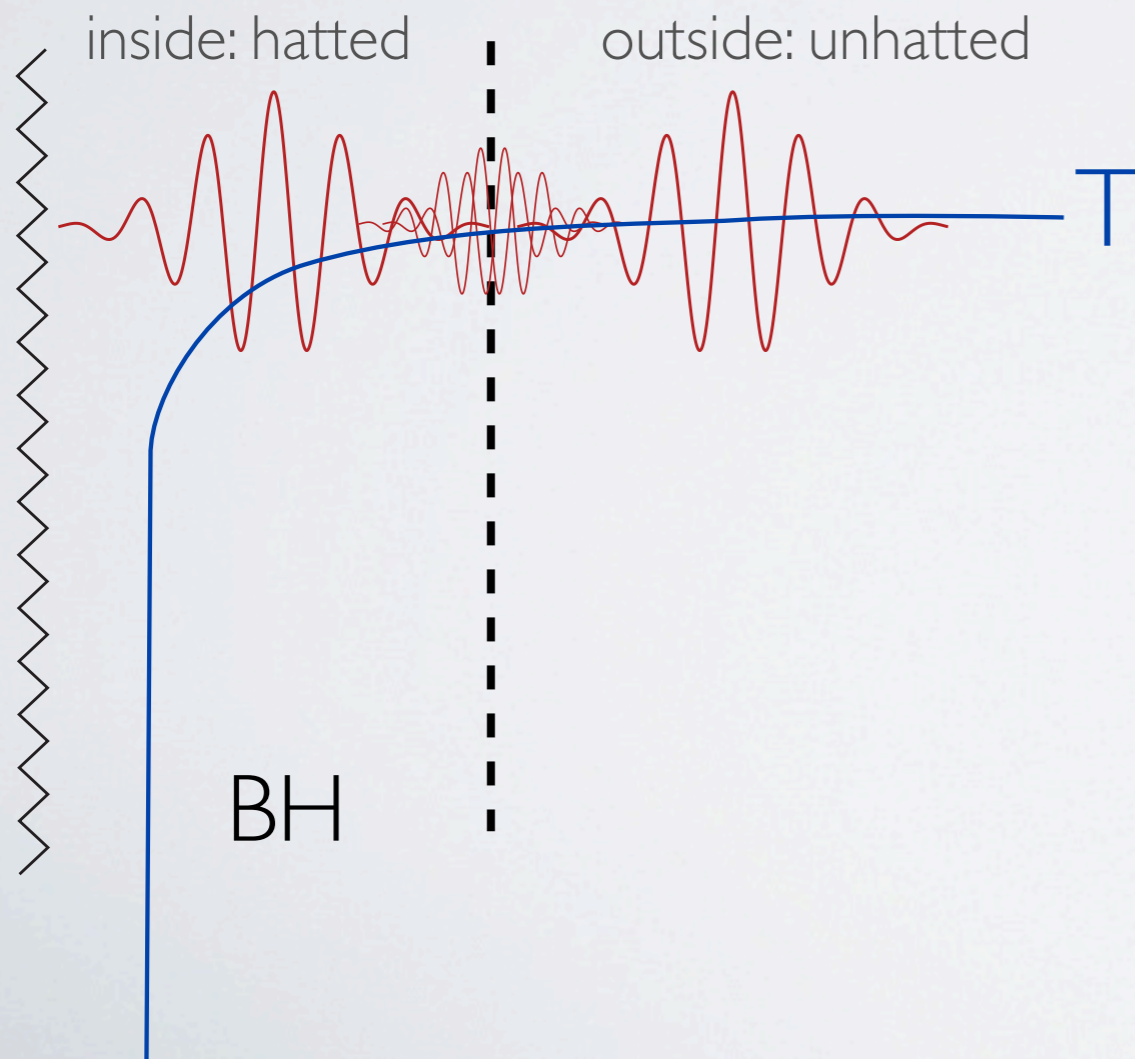
$$|\psi\rangle_{\text{Hawk}} = \prod_{jn} S_{jn} |\hat{0}\rangle |0\rangle$$

$j \sim$ asymp. frequency

$n \sim$ position along slice

$$\exp \left\{ \tanh^{-1} [e^{-\omega_j/2T}] \left(\hat{b}_{jn}^\dagger b_{jn}^\dagger - \text{h.c.} \right) \right\}$$

$$= \frac{1}{Z} \sum_{\{n_{jn}\}} e^{-\beta H/2} |\widehat{\{n_{jn}\}}\rangle |\{n_{jn}\}\rangle$$



For given T :

$$n \rightarrow \infty \leftrightarrow \rightarrow \text{horizon}$$

Regularize:

$$n < N(T) \sim \lambda > L$$

Shorter modes “look like” vacuum

“Hawking state” (cont’d):

So rewrite:

$$|\psi\rangle_{\text{Hawk}} = \prod_{jn} S_{jn} |\hat{0}\rangle |0\rangle = |0\rangle_N \prod_j \prod_n^{N-1} (S |\hat{0}\rangle |0\rangle)_{jn}$$

$\in \mathcal{H}_{BH} \otimes \mathcal{H}_{ext}$

As discussed

- $|0\rangle_N$ can either go with BH, or “ancillary”
- Hilbert spaces effectively finite dim. (if finite time)

Evolution: e.g 2d: $U = 1$

$$|\psi\rangle_{\text{Hawk}} = |0\rangle_N \prod_j \prod_n^{N-1} (S|\hat{0}\rangle|0\rangle)_{jn}$$

$$= \frac{|0\rangle_N}{Z} \sum_{\{n_{jn}\}, n < N} e^{-\beta H/2} |\widehat{\{n_{jn}\}}\rangle |\{n_{jn}\}\rangle$$

more generally: U_{LQFT}
 (also, can include infalling matter)

and: $N \rightarrow N + 1$

Note: “generalized” unitary transform: dimensions of $\mathcal{H}_{BH}, \mathcal{H}_{ext}$ change

Cartoon: timestep $\sim R$ (See also Mathur)

$$|0\rangle_N \rightarrow |0\rangle_{N+1} (|\hat{0}0\rangle + |\hat{1}1\rangle)_{n=N}$$

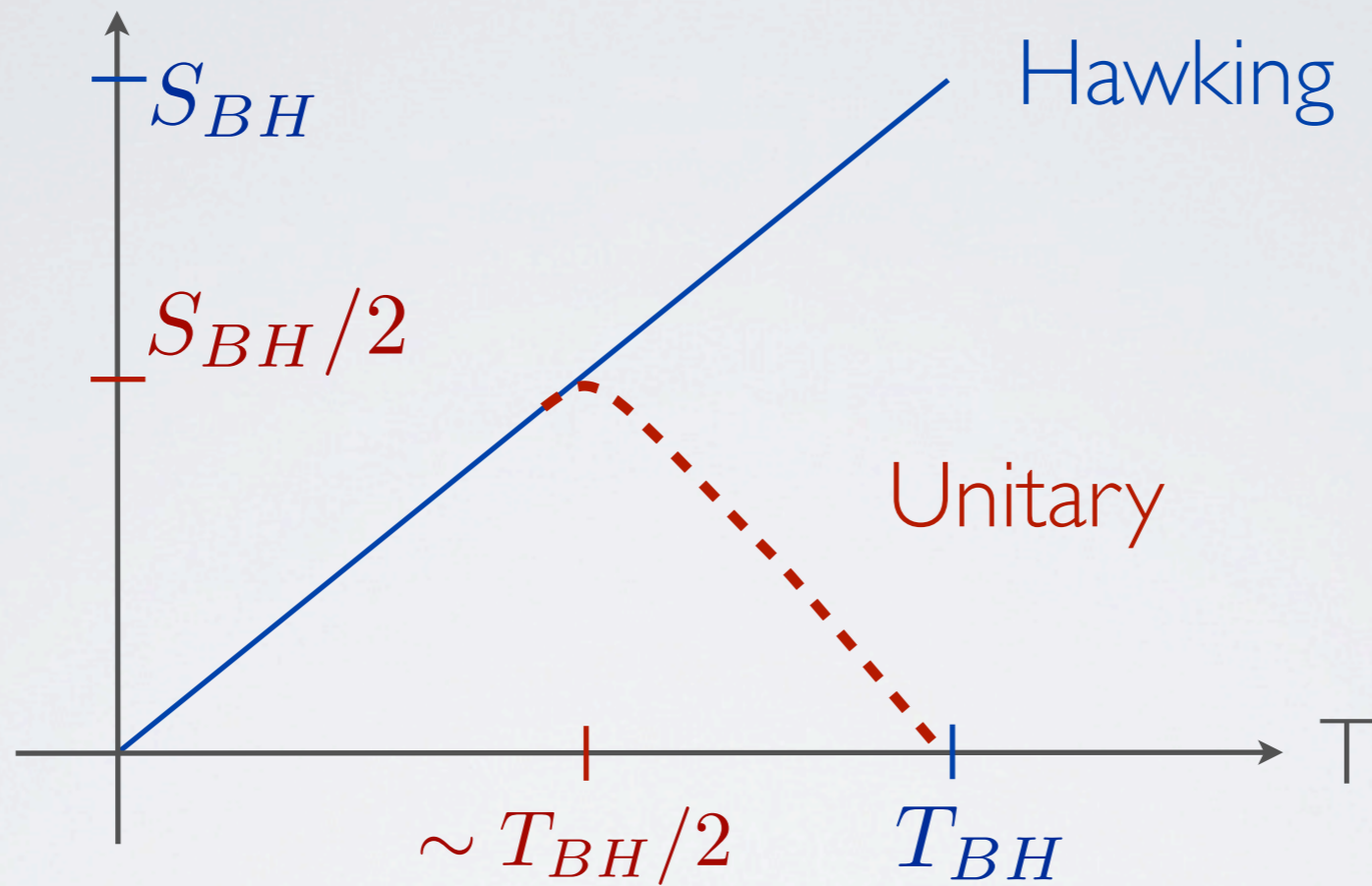
“qubit model”

↑
pairs produced

$(\omega_j \simeq 1/\beta)$

Nonunitarity $\rho_{ext} = \text{Tr}_{BH} (|\psi\rangle_{\text{Hawking}} \langle\psi|) \propto \sum_{\{n_{jn}\}, n < N} e^{-\beta H} |\{n_{jn}\}\rangle \langle\{n_{jn}\}|$

$$S_{ext} = -\text{Tr} \rho_{ext} \ln \rho_{ext}$$

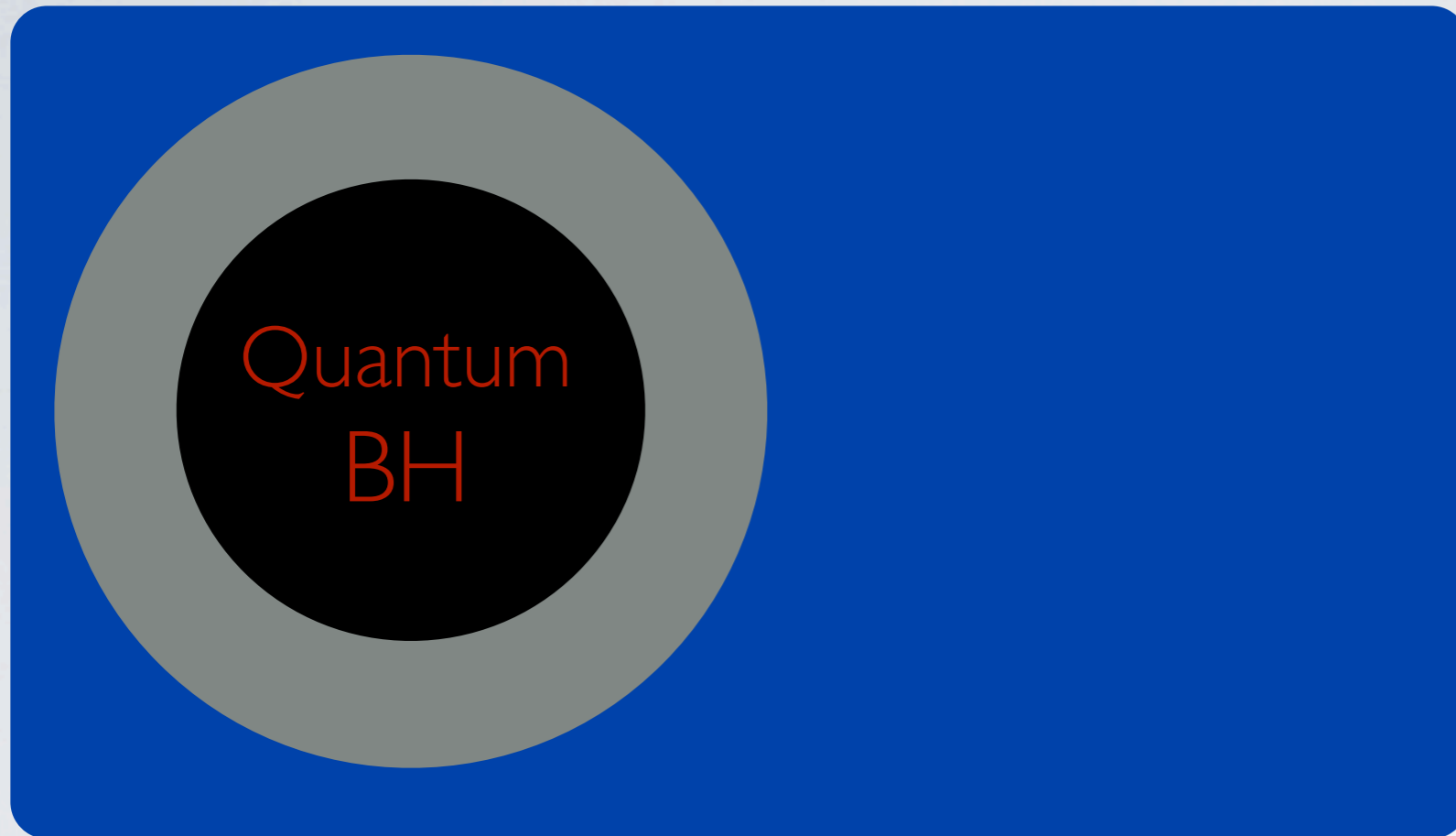


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arguments for failure of LQFT/nice-slice description
(hep-th/0703116, 0911.3395, etc.)

- problems making gauge invariant
- problems w/ perturbative quant.

But how do we describe presumably correct **unitary** evolution?



$$\mathcal{H}_{BH} \otimes \mathcal{H}_{near} \otimes \mathcal{H}_{far}$$

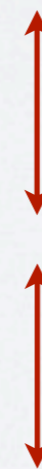
(nb: AdS/CFT could work this way)

Some expectations:

$$\mathcal{H}_{far} \approx LQFT$$

$$\mathcal{H}_{near} \sim LQFT$$

$$\mathcal{H}_{BH} \neq LQFT$$



$$U \approx LQFT$$

$$U \neq LQFT$$

\mathcal{H}_{BH} , \mathcal{H}_{near} finite dimensional; $\log \dim \mathcal{H}_{BH} = S_{BH}$?

Physical constraints on evolution:

A) $S \rightarrow 0$

B) Innocuous to infalling observers (?) (preserve “pain-free” horizon)

C) $dS/dt \sim 1/R$ (?)

D) Near-Hawking; \sim thermodynamic (?)

E) Correspondence limit: LQFT + GR

F) Energy conservation

G) Complete, consistent

...

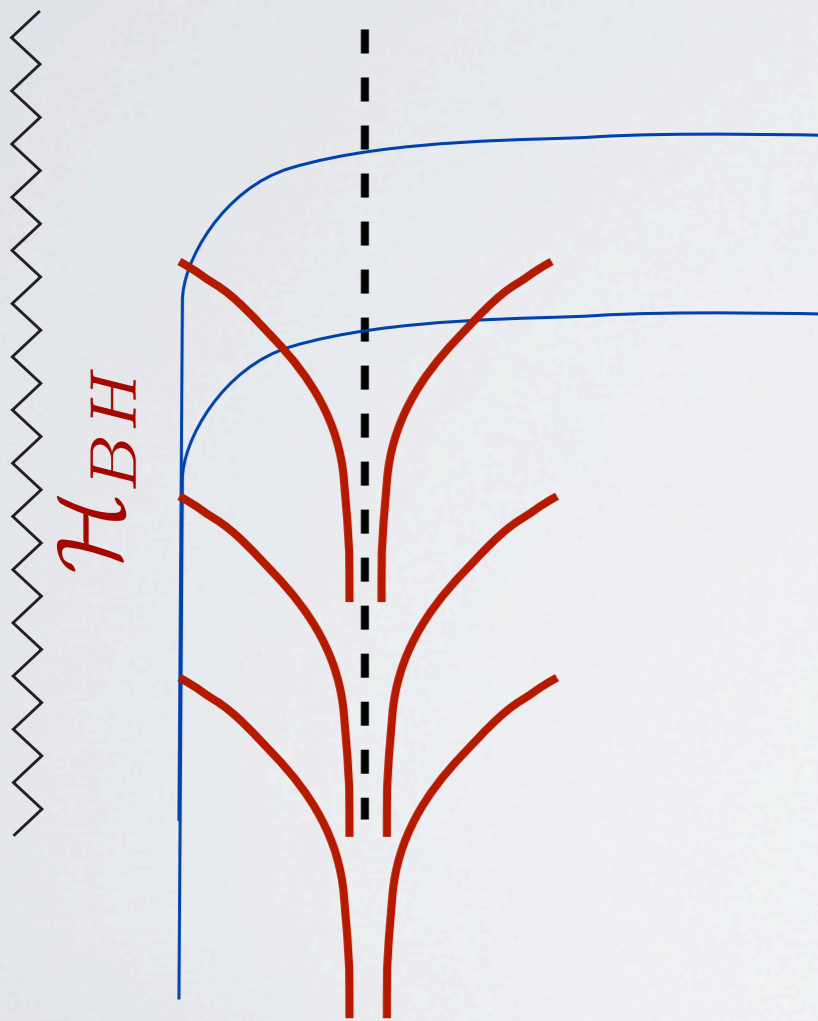
Basic guideline: be as conservative as possible!

What is “as close as possible to LQFT”?

Explore examples
(e.g. qubit models)

most conservative?

Recap:



Explore examples
(e.g. qubit models)

Not most conservative?

I A) e.g. qubit “q=0,1” in:

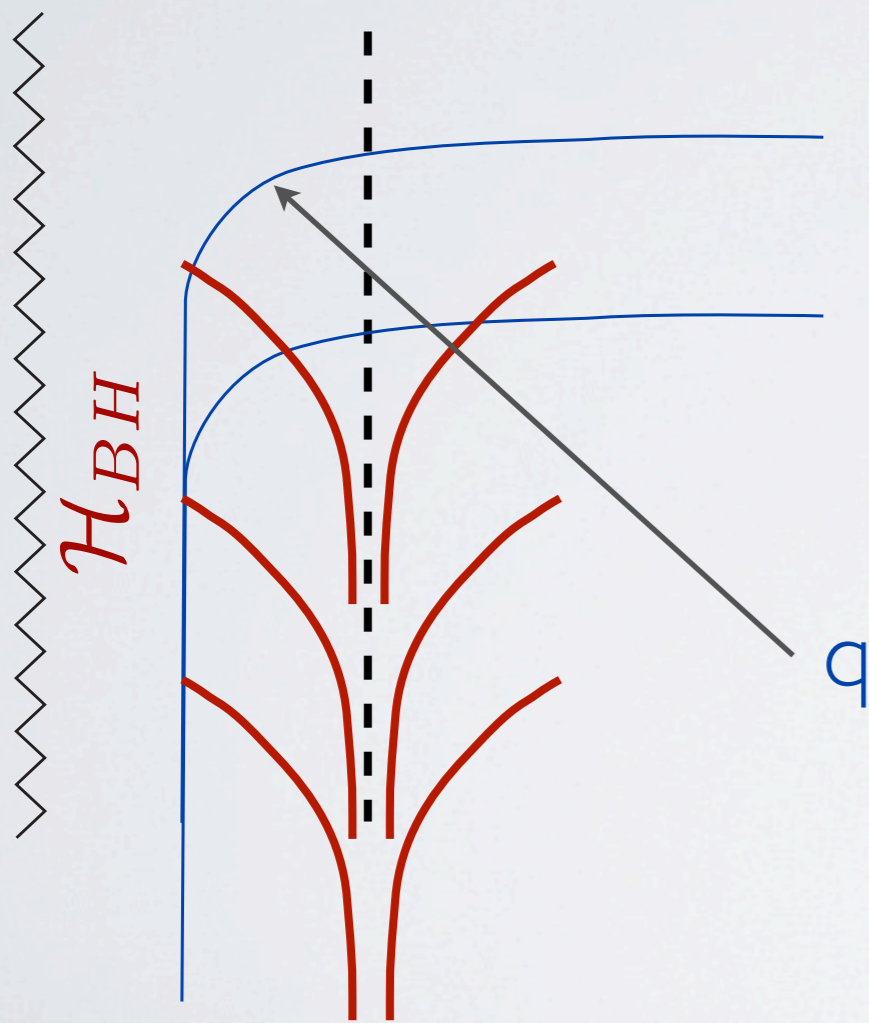
$$\rightarrow |\hat{a}\rangle |\hat{q}\rangle |\psi\rangle$$



$$\hat{U} (|\hat{a}\rangle |\hat{q}\rangle)$$

$$\Delta t \sim R$$

Recap:



“Fast scrambling”

(or, w/ $|\hat{00}\rangle + |\hat{11}\rangle$)

then:

$$|\hat{q}' \hat{a}'\rangle |\psi\rangle \rightarrow |\hat{a}'\rangle |\hat{q}'\rangle |\psi\rangle$$

(separate scrambling/transfer)

(these are representative of more general finite dim models)

Not most conservative?

Big departure from LQFT evolution

(\leftrightarrow complementarity?)

Indeed, Hayden/Preskill:

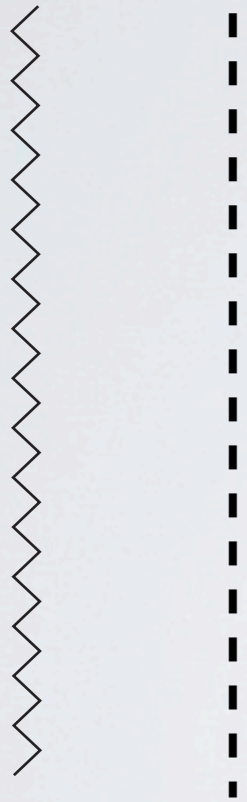
After sufficient evolution, a BH behaves as an information mirror on the scrambling time!

One classification:

	$T_{scramble}$	$T_{transfer}$
Susskind	$R \ln R$	$R \ln R$
Page		$< \mathcal{O}(RS)$
HR, nat. slice	$\sim R?$	∞
HR, nice slice	∞	∞

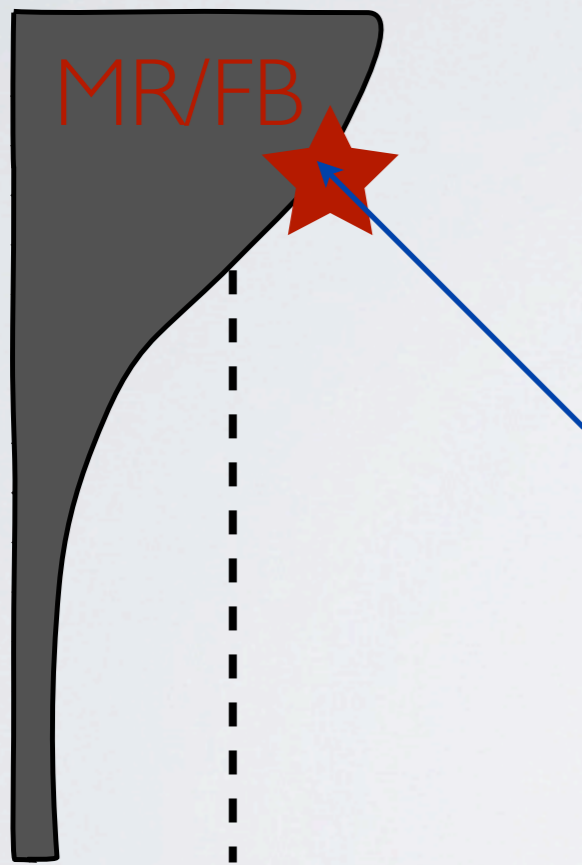
Not most conservative? (cont'd)

I B) Massive remnant; fuzzball



Not most conservative? (cont'd)

I B) Massive remnant; fuzzball



expect:

big mods. to \mathcal{H}_{BH}

rapidly varying microstructure
outside horizon

rapid, more limited(?) scrambling

(~ "neutron star")

Unitary models: “more conservative”

2)

$$|\hat{0}\hat{0}\hat{a}\rangle|a\rangle \rightarrow \hat{U}|\hat{a}\rangle \frac{(|\hat{0}0\rangle + |\hat{1}1\rangle)}{\sqrt{2}} U|a\rangle$$

$$|\hat{0}\hat{1}\hat{a}\rangle|a\rangle \rightarrow \hat{U}|\hat{a}\rangle |\hat{0}1\rangle U|a\rangle$$

$$|\hat{1}\hat{0}\hat{a}\rangle|a\rangle \rightarrow \hat{U}|\hat{a}\rangle |\hat{1}0\rangle U|a\rangle$$

$$|\hat{1}\hat{1}\hat{a}\rangle|a\rangle \rightarrow \hat{U}|\hat{a}\rangle \frac{(|\hat{0}0\rangle - |\hat{1}1\rangle)}{\sqrt{2}} U|a\rangle$$

e.g. LQFT

e.g. “leftmost”
qubits

... representative of more general (e.g. multimode) models:

(and e.g. extra unitaries)

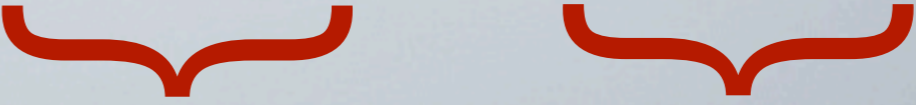
“Hawking-like”

“minimal” mods to evolution

- info imprinted in typical Hawking modes

$$- \langle T_r^0 \rangle = \langle T_r^0 \rangle_{Hawking}$$

$$3) \quad |\hat{q}_1 \hat{q}_2 \hat{a}\rangle |a\rangle \rightarrow \hat{U} |\hat{a}\rangle \frac{(|\hat{0}\rangle|0\rangle + |\hat{1}\rangle|1\rangle)}{\sqrt{2}} |\hat{0}'\hat{0}''\rangle |q'_1 q''_2\rangle U|a\rangle$$



Usual Hawking particles
Not typically occupied

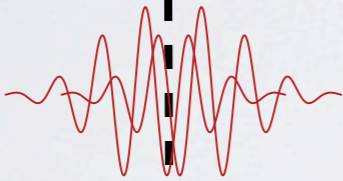
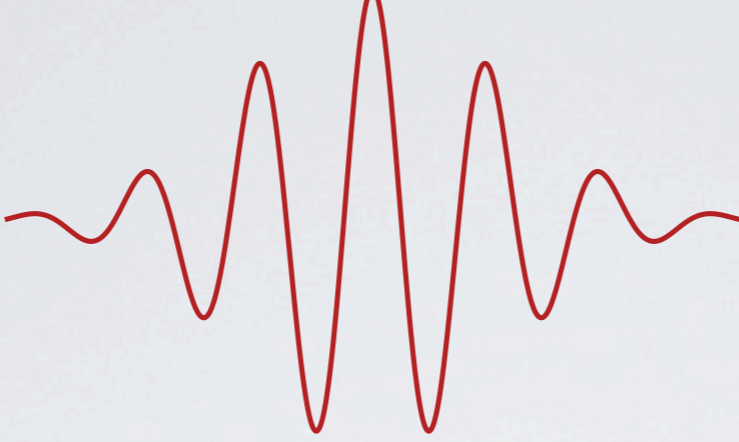
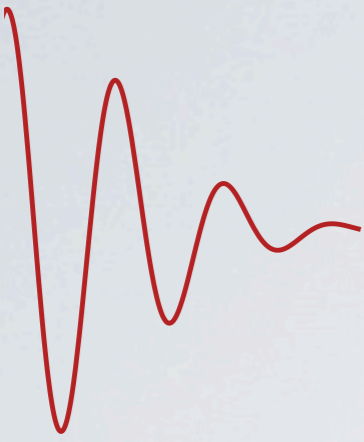
(again, modulo unitaries)

... representative of more general (e.g. multimode) models:

- more generic
- yield extra energy flux

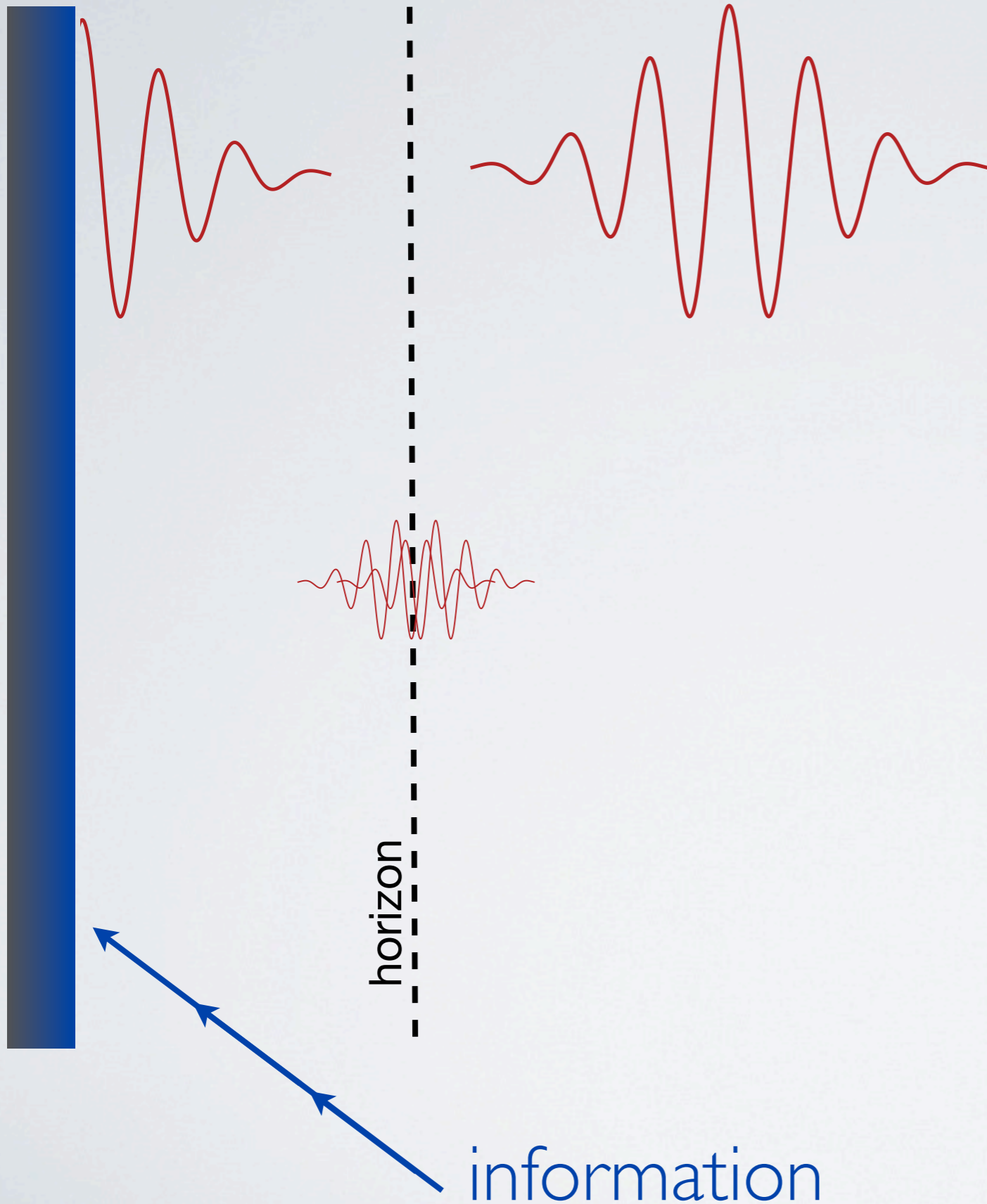
(But: $dE/dI > 1/R$?)

- This evolution is **not local** with respect to the semiclassical geometry
- “Pain-free” horizon?



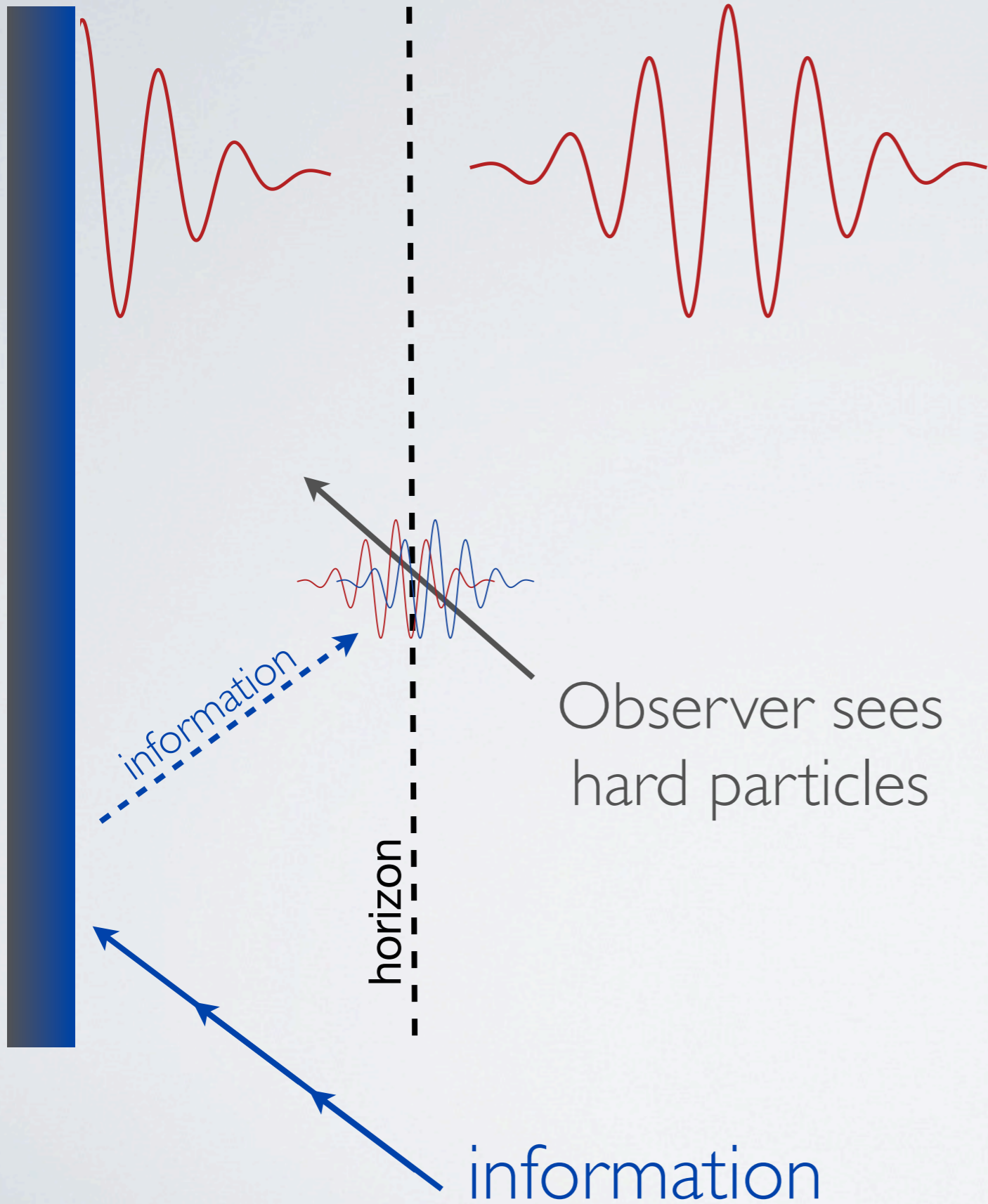
horizon





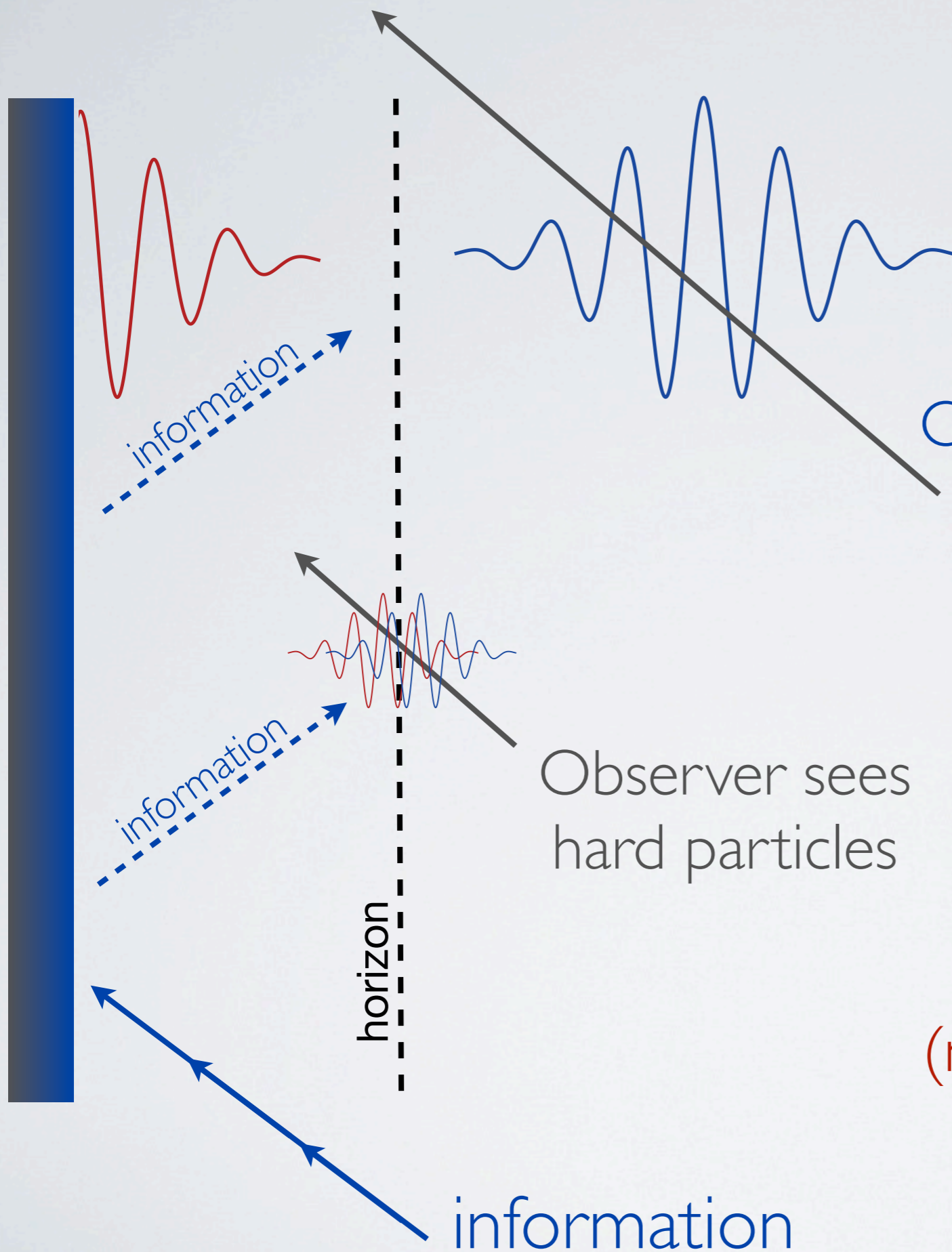
information

horizon



Observer sees
hard particles

information

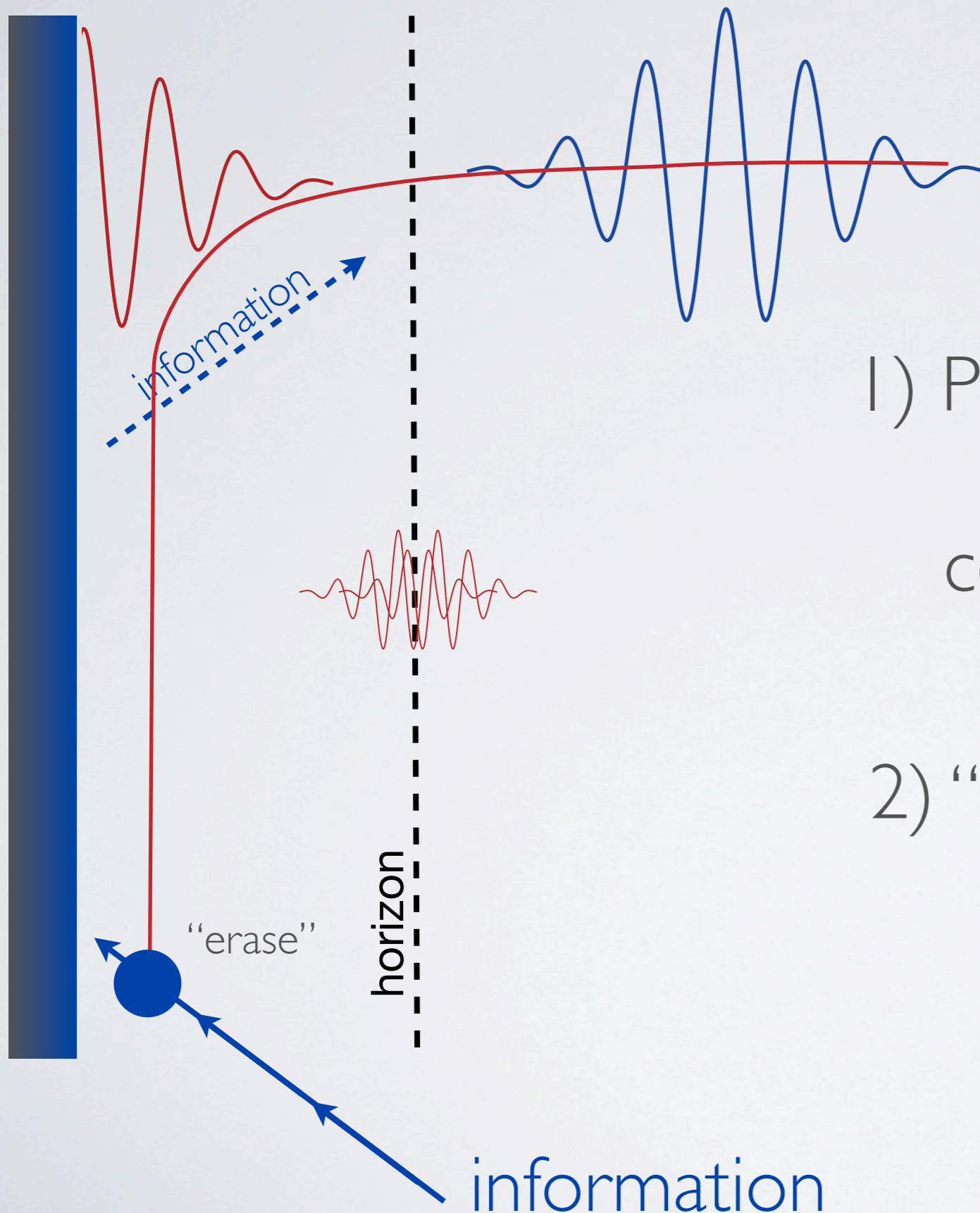


Observer sees ~ 1 extra quantum of energy $\sim 1/R$ per time R (\sim innocuous)

Observer sees hard particles

(n.b.: idea is: geometry not strictly correct picture of physics - spacetime=approx. !);

Two further comments:



- 1) Pain-free horizon:
small dim of \mathcal{H}_{near}
constraint on evolution
- 2) “Weak complementarity”
As long as $T \sim RS$
(e.g. nice slice descript.
not good at long times)

- Can explore other restrictions from physical + q. info-theoretic requirements (SBG & Shi | 205.xxxx & WIP)

E.g. characterize information transfer

Minimal -- Simplest, most efficient form:



A



B

- Can explore other restrictions from physical + q. info-theoretic requirements (SBG & Shi | 205.xxxx & WIP)

E.g. characterize information transfer

Minimal -- Simplest, most efficient form:



$$A = A_1 \otimes A_2$$



B

- Can explore other restrictions from physical + q. info-theoretic requirements (SBG & Shi | 205.xxxx & WIP)

E.g. characterize information transfer

Minimal -- Simplest, most efficient form:



$$A = A_1$$



$$A_2 \otimes B$$

“Subspace transfer”

Mod unitaries U_A, U_B etc. Saturates a subadd. condition

Contrast nonminimal:

$$\begin{array}{l} |\hat{0}\rangle|0\rangle \rightarrow |\hat{0}\rangle|0\rangle \\ |\hat{1}\rangle|0\rangle \rightarrow |\hat{1}\rangle|1\rangle \end{array}$$

Might allow you to measure dead or alive
Schrodinger's cat inside a black hole, but doesn't
transfer quantum information ...

Example 2 was minimal; Example 3 was not
(extra flux)

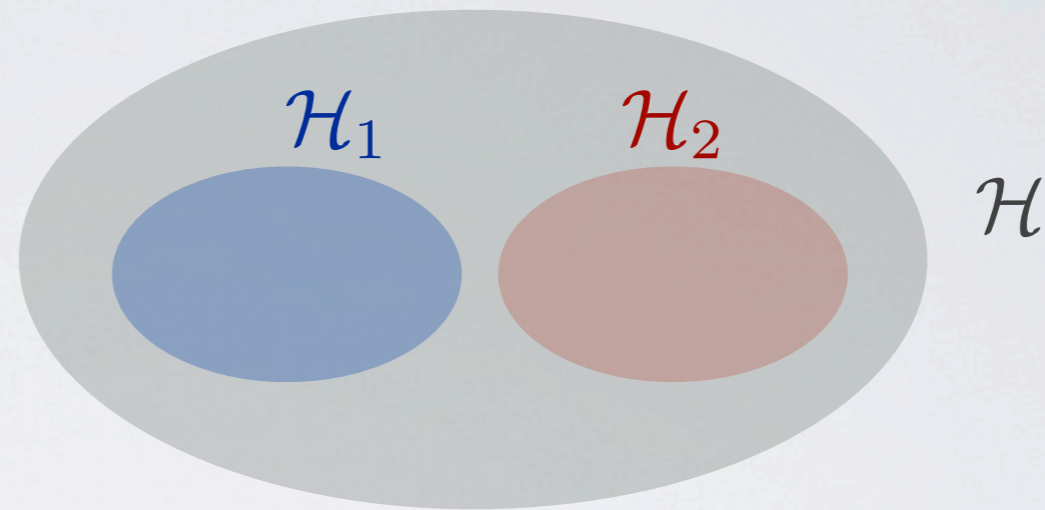
Possible reasons to expect (near-)minimal:

Weak coupling; \sim thermodynamic; small \mathcal{H}_{near}

A proposed broader picture:

Quantum states more basic than spacetime

Approximate
“localization”

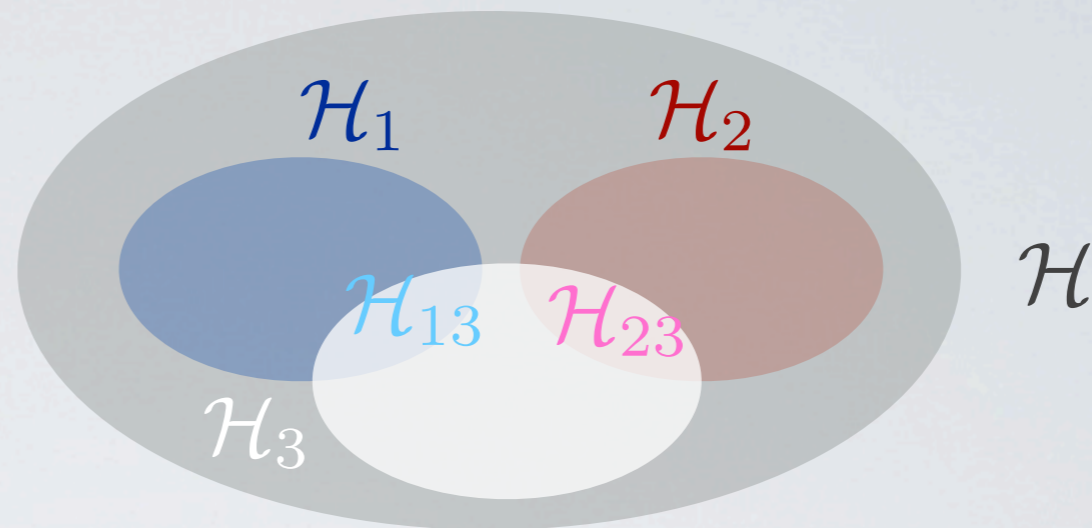


$$\mathcal{H}_1 \otimes \mathcal{H}_2 \in \mathcal{H}$$

Compare LQFT:

$$[\mathcal{O}(x), \mathcal{O}(y)] = 0 \quad , \quad (x - y)^2 > 0$$

~ quantum analog
of manifold:



(some common ideas w/ algebraic
QFT; also Banks “HST” -- though
important differences)

- Unitary evolution; ~local, LQFT

- Symmetries $\mathcal{H} \rightarrow S\mathcal{H}$ global
 $\mathcal{H} \rightarrow S_{loc}\mathcal{H} = S_1\mathcal{H}_1 \otimes S_2\mathcal{H}_2 \cdots$ local

Hilbert space networks: a possible framework for a
unitary theory of quantum gravity

Summary:

The information problem appears foundational.

A “most conservative” approach is to modify
macroscopic locality. (not QM, LI)

Unitarity can be restored through QI transfer from
BH subsystem ... **not LQFT** (but “nearly?”)

Proposal: approximate spacetime emerges from a
broader framework - Hilbert space network