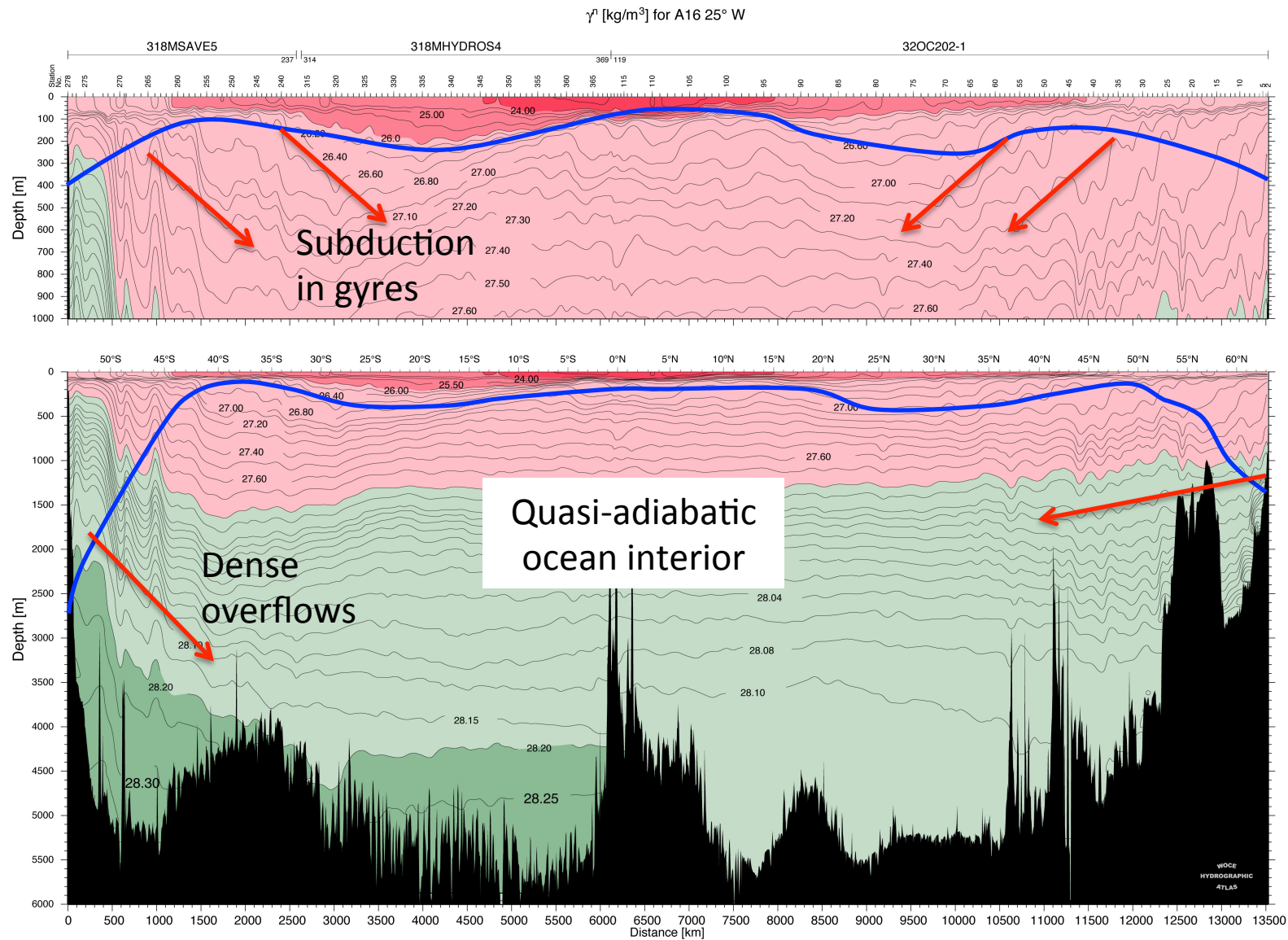
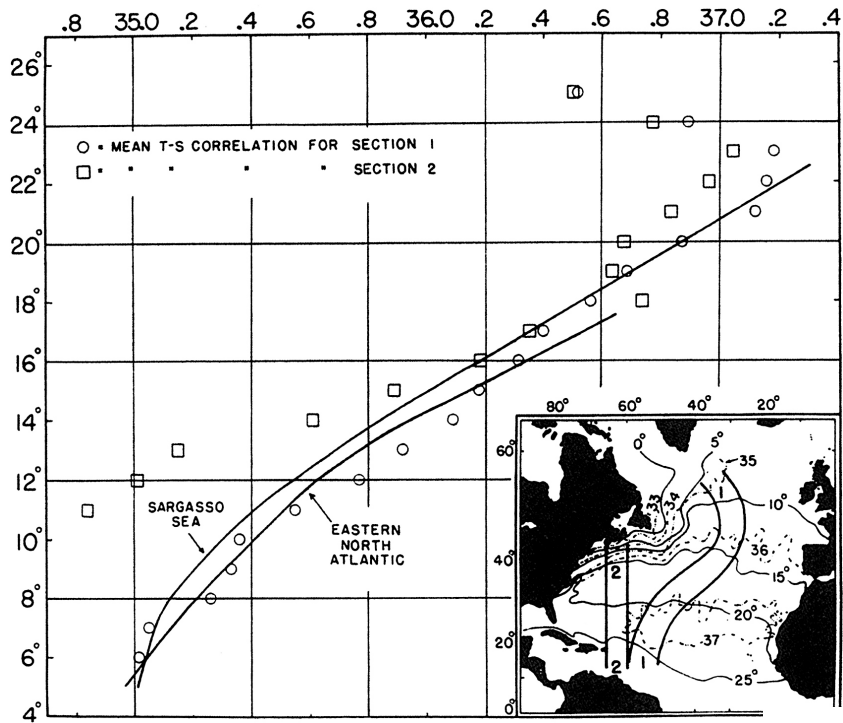


The oceanic boundary layers in models: our (unfortunate) reliance on parameterizations



Bulk of ocean properties are set in the boundary layers





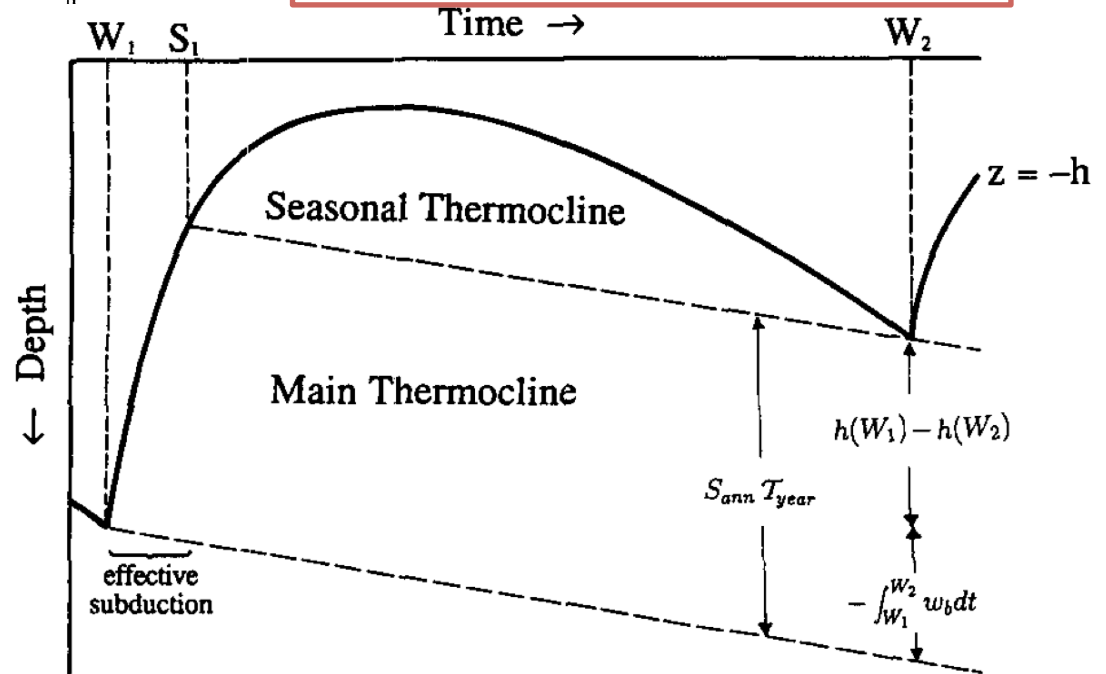
Temperature-salinity along surface swaths in the North Atlantic (dots and squares), and in the vertical (solid curves) at stations in the western North Atlantic (Sargasso Sea) and eastern North Atlantic. Source: From Iselin (1939).

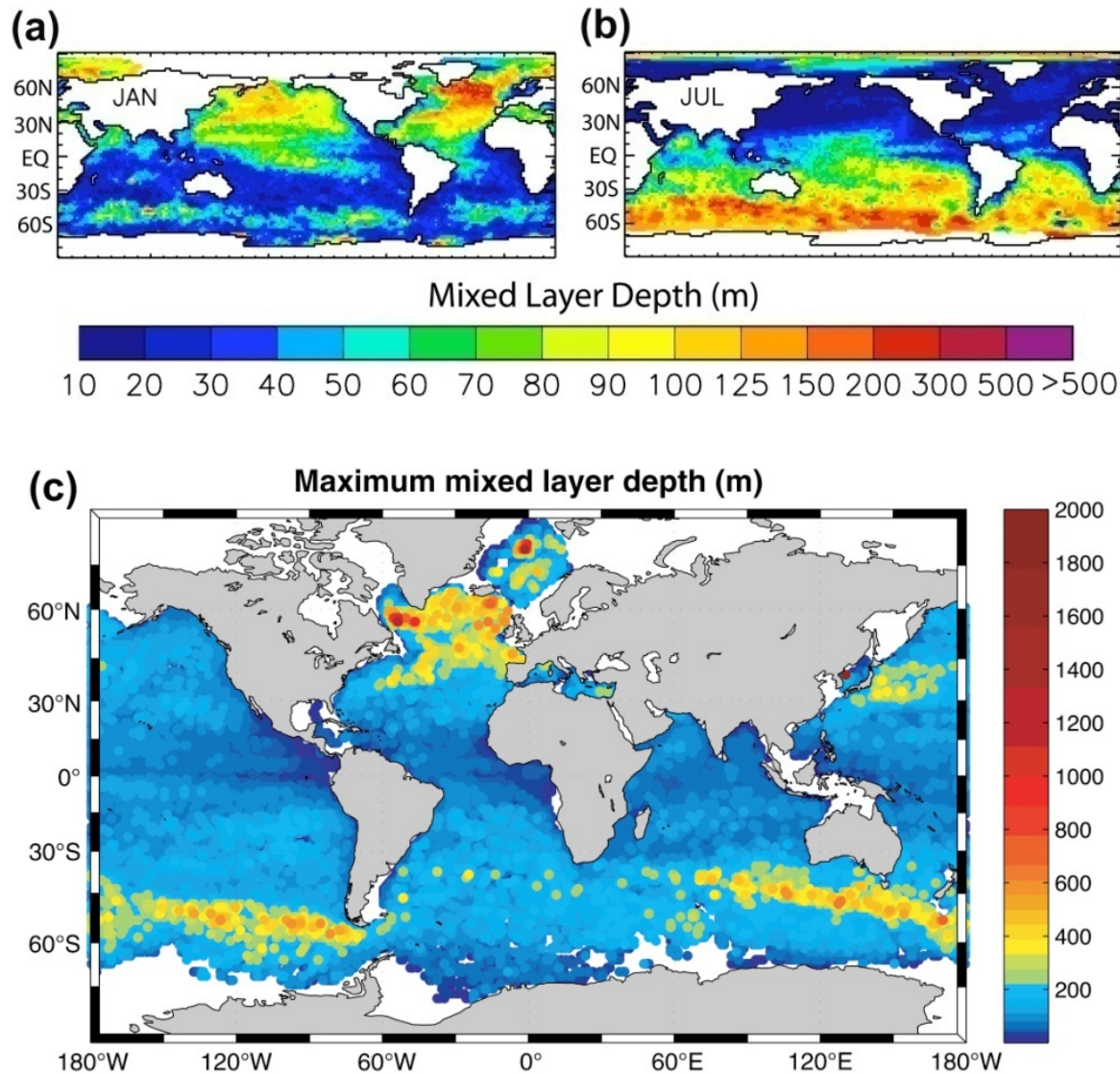
Ocean interior properties are set by the late winter ML properties

TALLEY
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FIGURE 4.6

Williams et al. (1994)

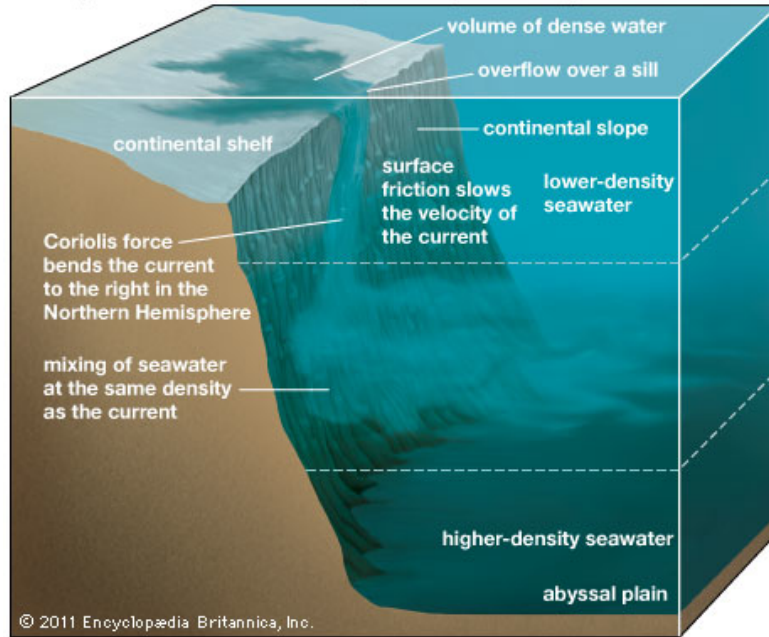




Mixed layer depth in (a) January and (b) July, based on a temperature difference of 0.2°C from the near-surface temperature. *Source: From deBoyer Montégut et al. (2004).* (c) Averaged maximum mixed layer depth, using the 5 deepest mixed layers in $1^{\circ} \times 1^{\circ}$ bins from the Argo profiling float data set (2000-2009) and fitting the mixed layer structure as in Holte and Talley (2009). This figure can also be found in the color insert.

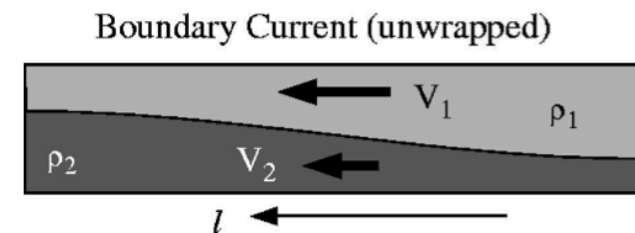
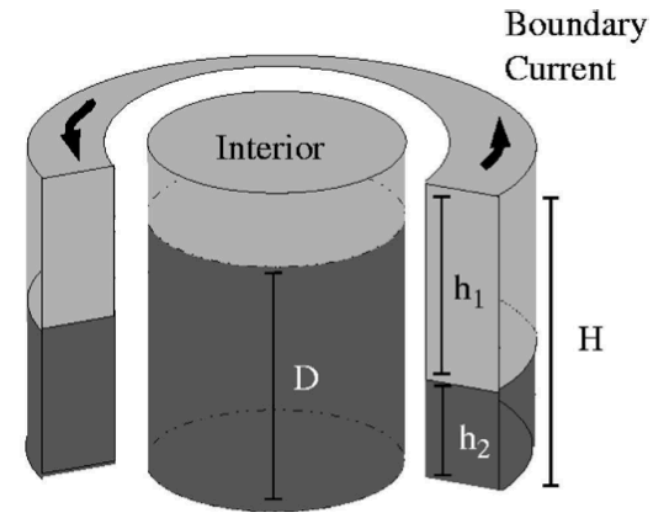
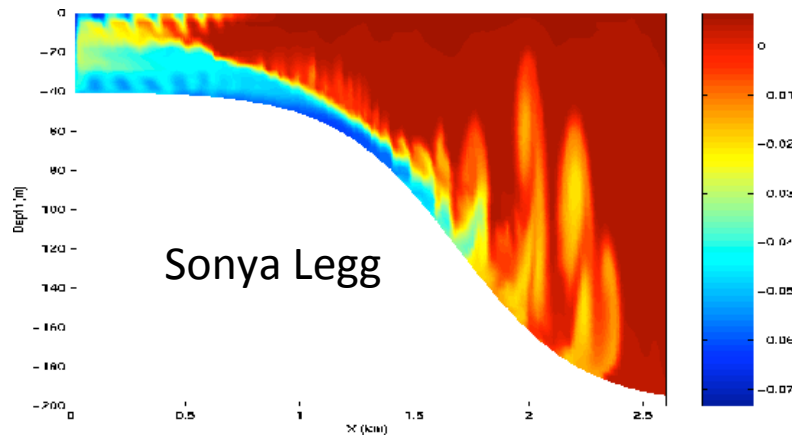
FIGURE 4.4

Density current: descent to a layer of equal density



At high latitudes, final properties of water set in:

- Overflow from continental shelves or over oceanic ridges
- Exchanges between boundary currents and open ocean



Straneo, 2008

How much boundary layer schemes influence the outcome of climate simulations?

→ Focus on the surface boundary layers

Compare:

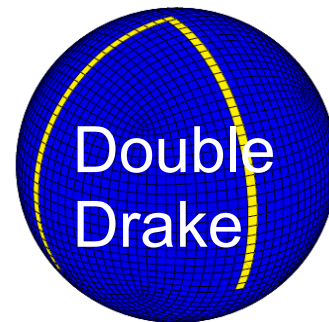
Simulation with/without a boundary layer

Simulations with different boundary layers

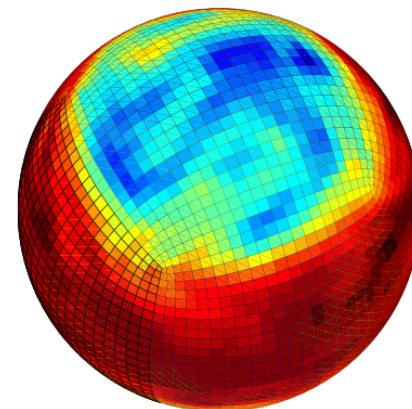
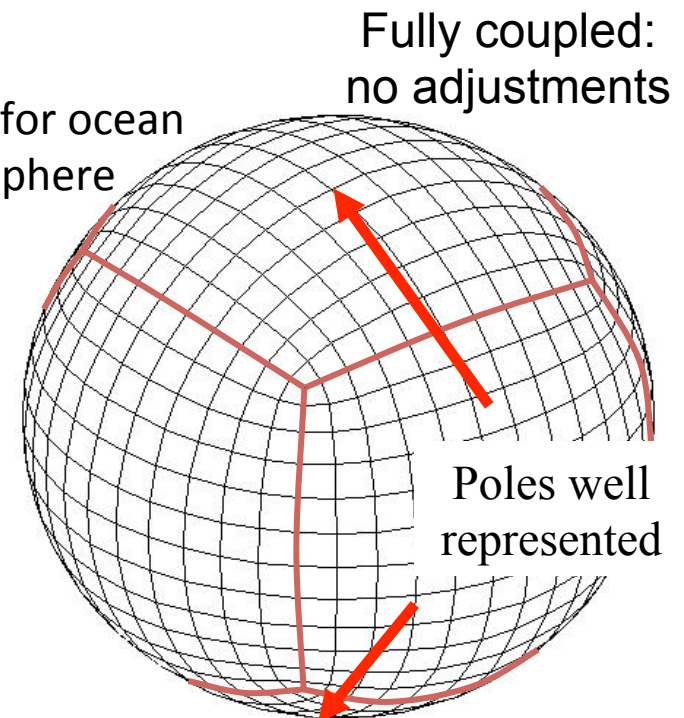
No attempt to tune, take schemes as “out of box”

2 MIT GCM: Coupled Ocean-Atmosphere-Sea ice:

- Primitive equation models,
- Cube-sphere grid: $\sim 3.75^\circ$,
- Synoptic scale eddies in the atmosphere,
- Gent and McWilliams eddy parameterization in the ocean,
- Simplified atmospheric physics (SPEEDY, Molteni 2003),
- Conservation to numerical precision (Campin et al. 2008)

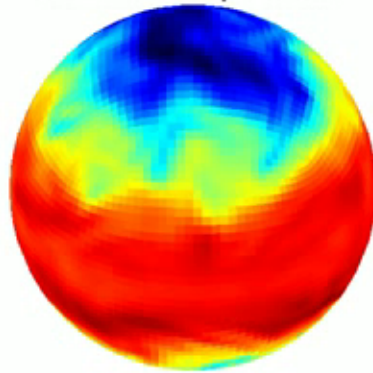


Same grid for ocean and atmosphere

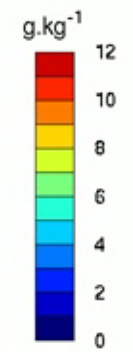
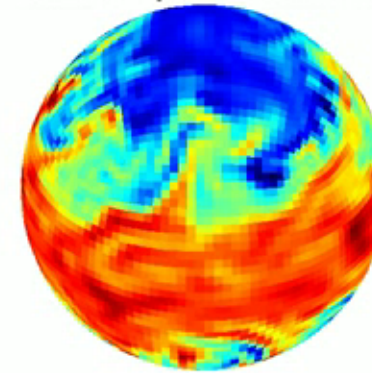


Temperature
snap-shot at
500 mb.

500 mb Temperature

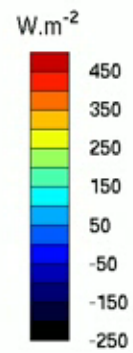
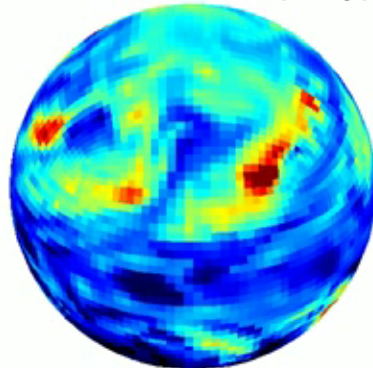


Surface Specific Humidity

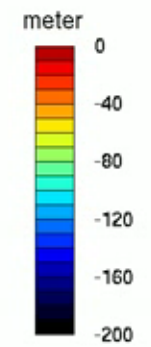
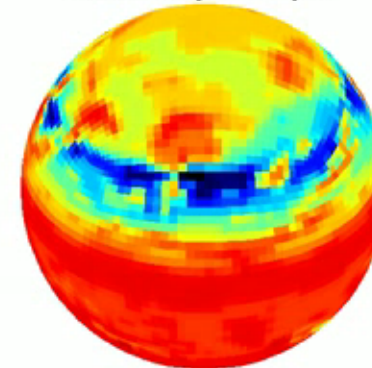


Day 1

Air-sea Heat Flux (+ = up)



Mixed-layer Depth



No boundary layer scheme (“Reference” run)

- First grid level: 30 m thick
 - Mixed layer is always at least 30 m thick
- Convective adjustment: intense mixing for $N^2 < 0$
 - Applies to T and S and other passive tracers
 - Nothing of momentum

Non-local K profile parameterization (KPP), Large et al. (1994)

$$\overline{wx}(d) = -K_x(\partial_z X - \gamma_x)$$

Local term
(although non-local too)

Non-local term depend on forcing

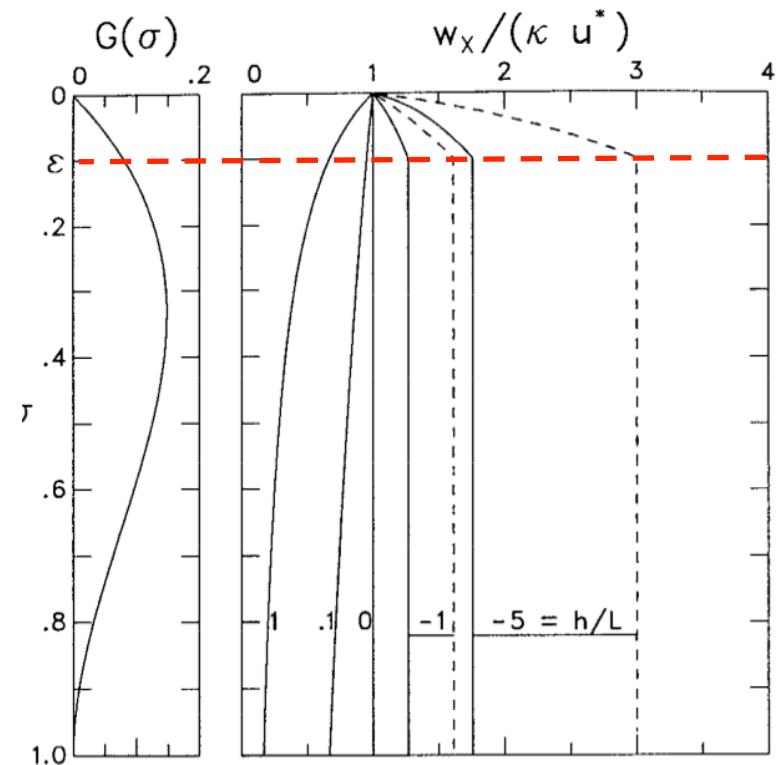
$$K_x(\sigma) = hw_x(\sigma)G(\sigma)$$

h boundary layer depth

$$\gamma_x = 0 \quad \zeta \geq 0$$

$$\gamma_m = 0 \quad \gamma_s = C_s \frac{\overline{ws_0}}{w_s(\sigma)h}$$

$$\gamma_\theta = C_s \frac{(\overline{w\theta_0} + \overline{w\theta_R})}{w_s(\sigma)h} \quad \zeta < 0$$



TKE model Gaspard et al. (1990), Blanke and Delecluse (1993)

Second-order moment closure

TKE prognostic equation:

$$\partial_t \bar{e} = -\partial_z (\overline{ew'} + \rho_0^{-1} \overline{p'w'}) - \overline{\mathbf{u}'_h w'} \cdot \partial_z \mathbf{U}_h + \overline{b'w'} - \epsilon,$$

~ Transport
Source of TKE
Sink of TKE
↑
dissipation

Closure

$$-\overline{\mathbf{u}'_h w'} = K_m \partial_z \mathbf{U}_h,$$

$$-\overline{b'w'} = -g \rho_0^{-1} K_\rho \partial_z \rho = K_\rho N^2$$

$$-\overline{(ew' + \rho_0^{-1} p'w')} = K_e \partial_z \bar{e}$$

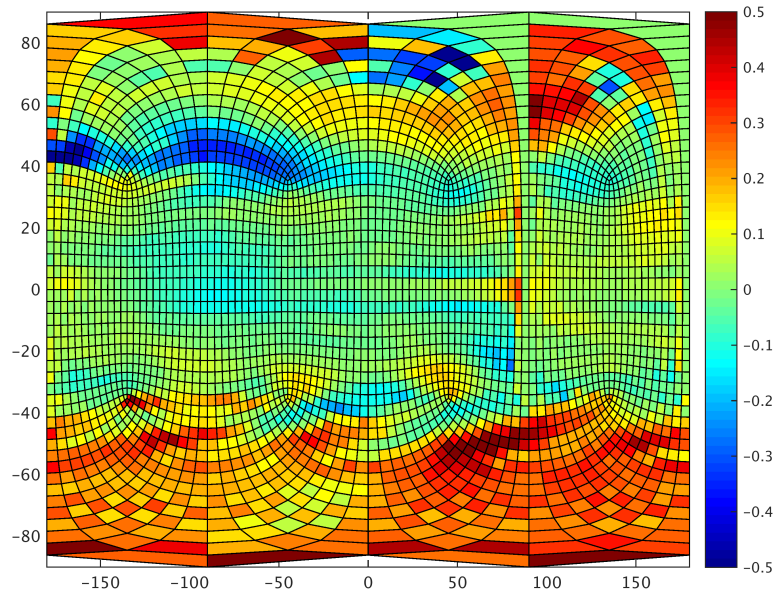
$$\epsilon = C_\epsilon \bar{e}^{3/2} l_\epsilon^{-1},$$

with $K_m = C_\kappa I_\kappa \bar{e}^{1/2},$

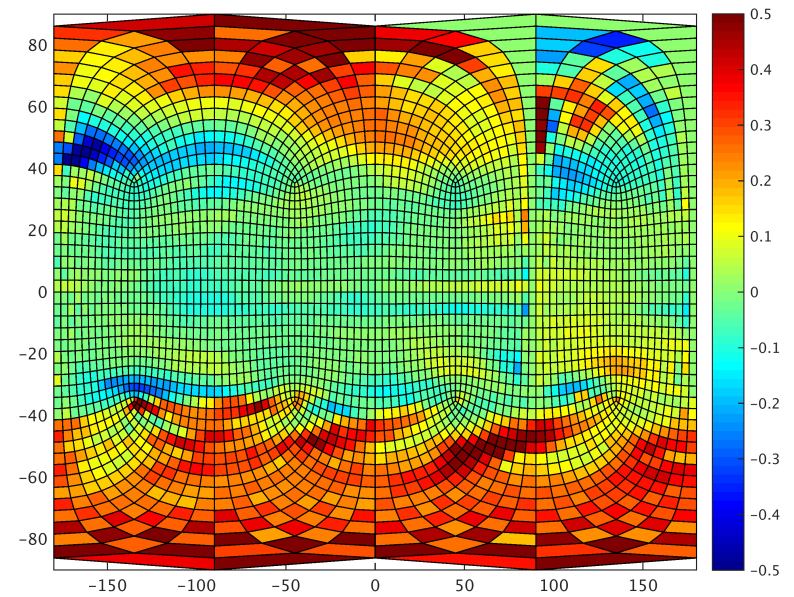
$$K_\rho = K_m / P_{rt},$$

Temperature

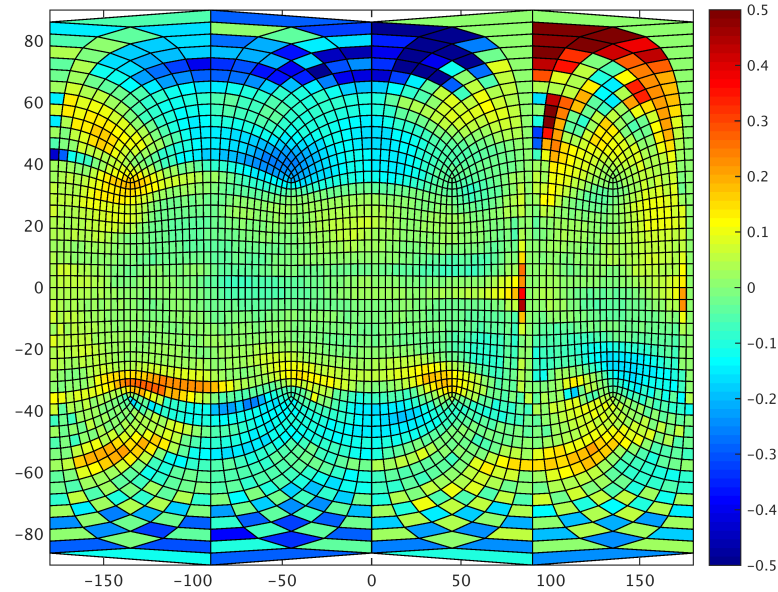
KPP minus Reference



GGL minus Reference

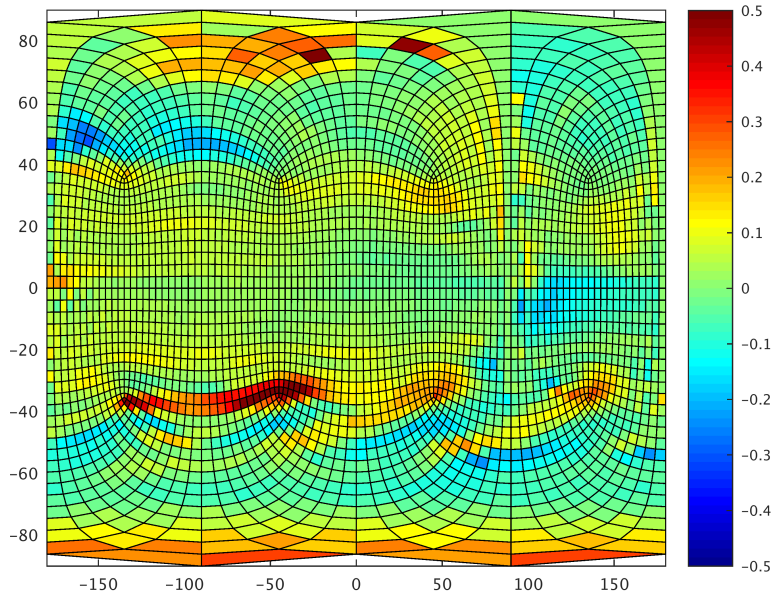


KPP minus GGL

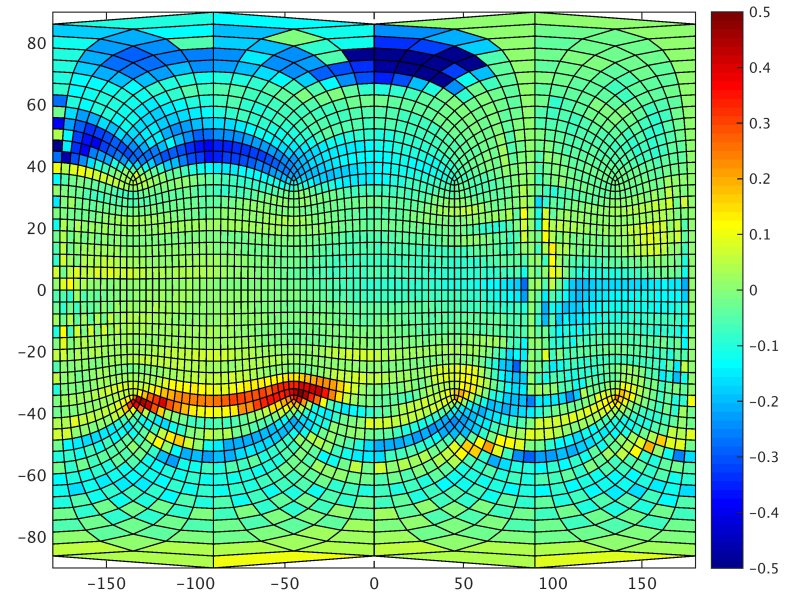


Salinity

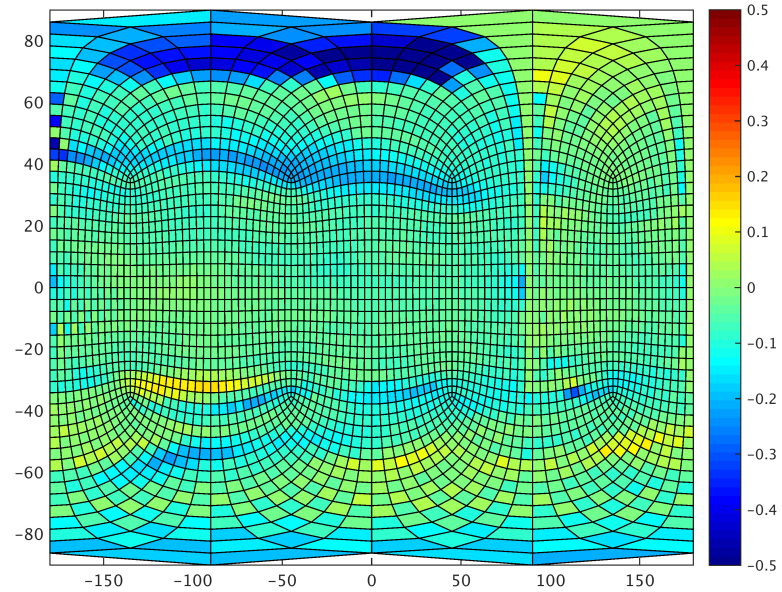
KPP - Reference



GGL - Reference



KPP - GGL



Interaction between eddy parameterization and mixed layer

Gent and McWilliams scheme: $\Psi^* = K_{GM} S_y$ S_y : Isopycnal slope

- Near the surface $S_y \rightarrow$ infinity (or close)
- Need for a “tapering” scheme to control stability
- Mostly ad-hoc, but see Ferrari et al. (08, 10), Ferreira and Marshall (05)

Example: Gerdes et al. 1991

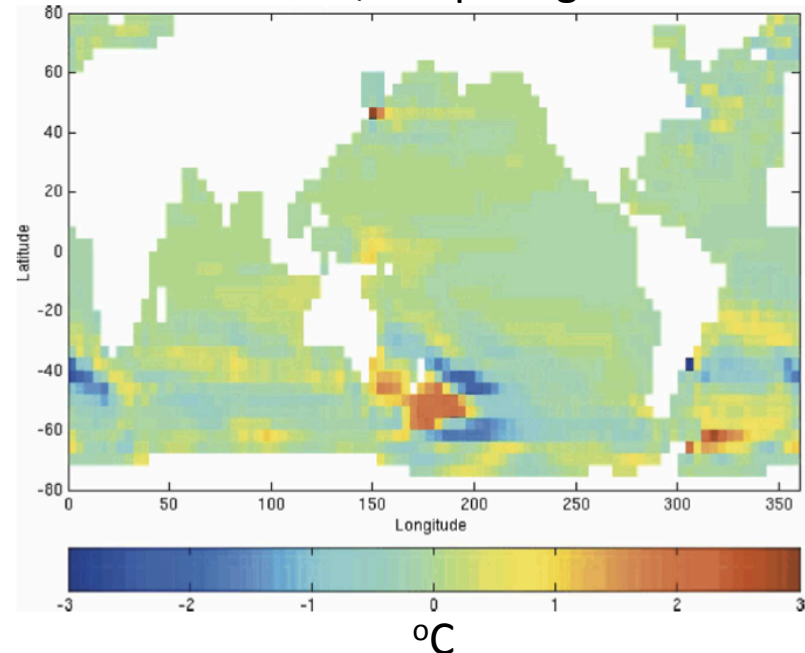
If $S_{big} > |S_y| > S_{max}$

$$\Psi^* = \frac{S_{max}^2}{S_y^2} K_{GM} S_y$$

If $|S_y| > S_{big}$

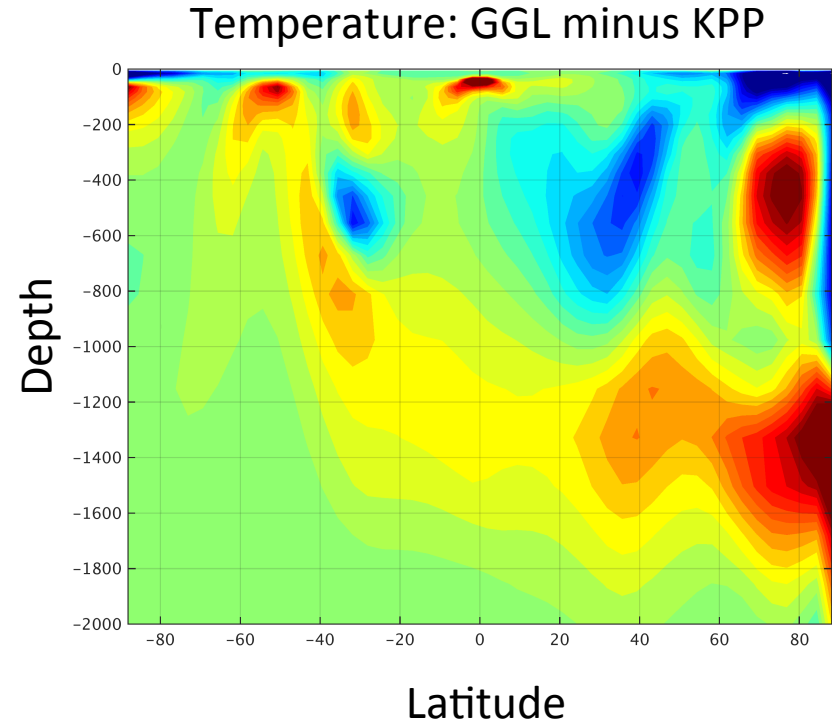
$$\Psi^* = 0$$

SST difference, 2 tapering schemes



Differences:

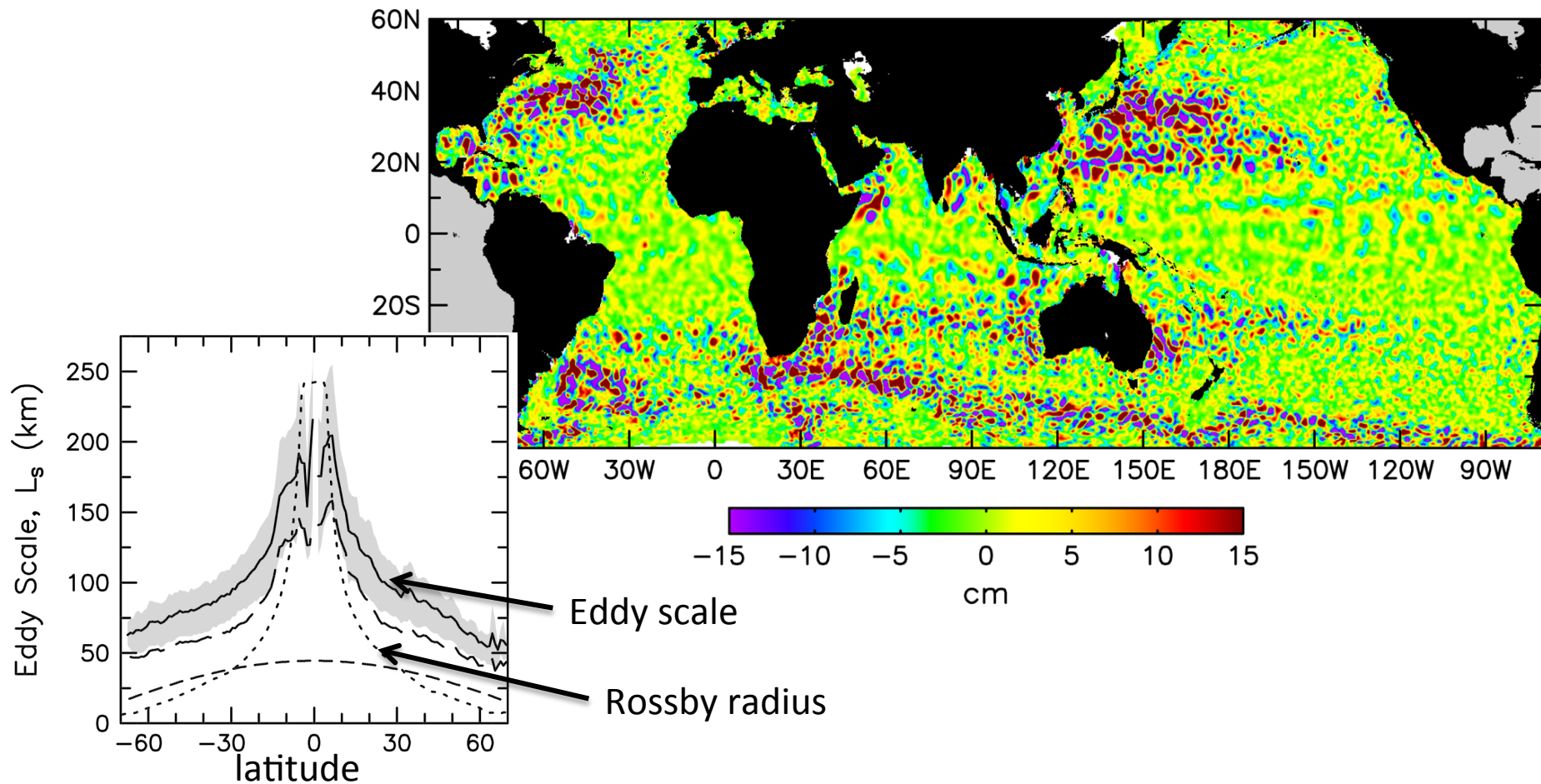
- Without/with BL scheme :
 - Locally: +/- 0.3-0.5 °C
 - Globally: + 0.2 °C
- Between BL schemes:
 - Locally: +/- 0.3-0.5 °C
 - Globally ~ 0 °C
- Can one tune GGL and KPP so they do the same thing?
 - Probably yes
- Can we always do it?
- If yes, it is good?
 - Possibly no.
 - i.e. there are too many “unknown” parameters
 - Not enough $O(1)$ constants



Ocean are full of mesoscale eddies

- Coherent vortices, radius of about 50-100 km
- Generated by baroclinic and barotropic instabilities
- ~215,000 eddies with 4 weeks or longer lifetime over 20 years (Chelton et al., 2011)
- They are everywhere

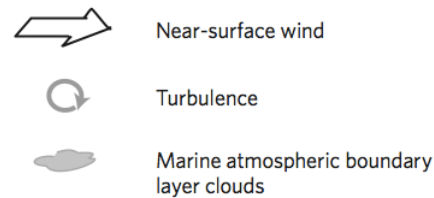
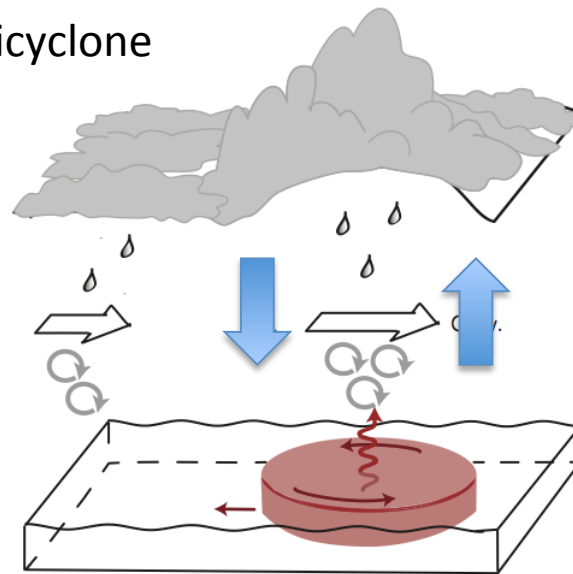
SSH anomalies 28 Aug 1996



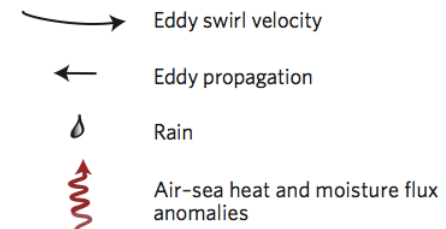
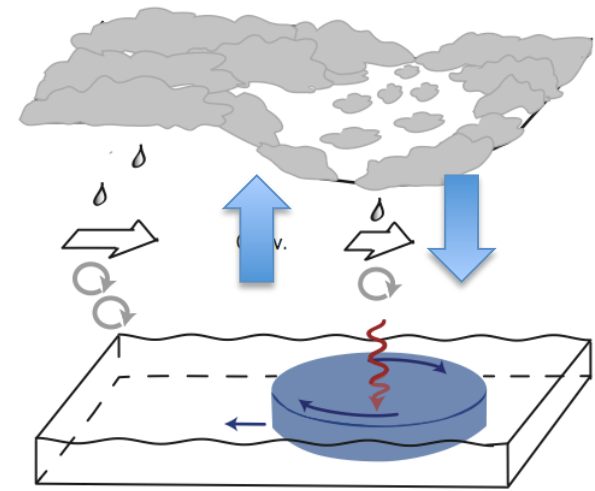
Mechanism

- Dominant mechanism: Modulation of the BL vertical mixing
- Warmer SST
 - Decreased vertical stability
 - Downward momentum flux
 - Decreased shear
 - Larger surface winds

Anticyclone



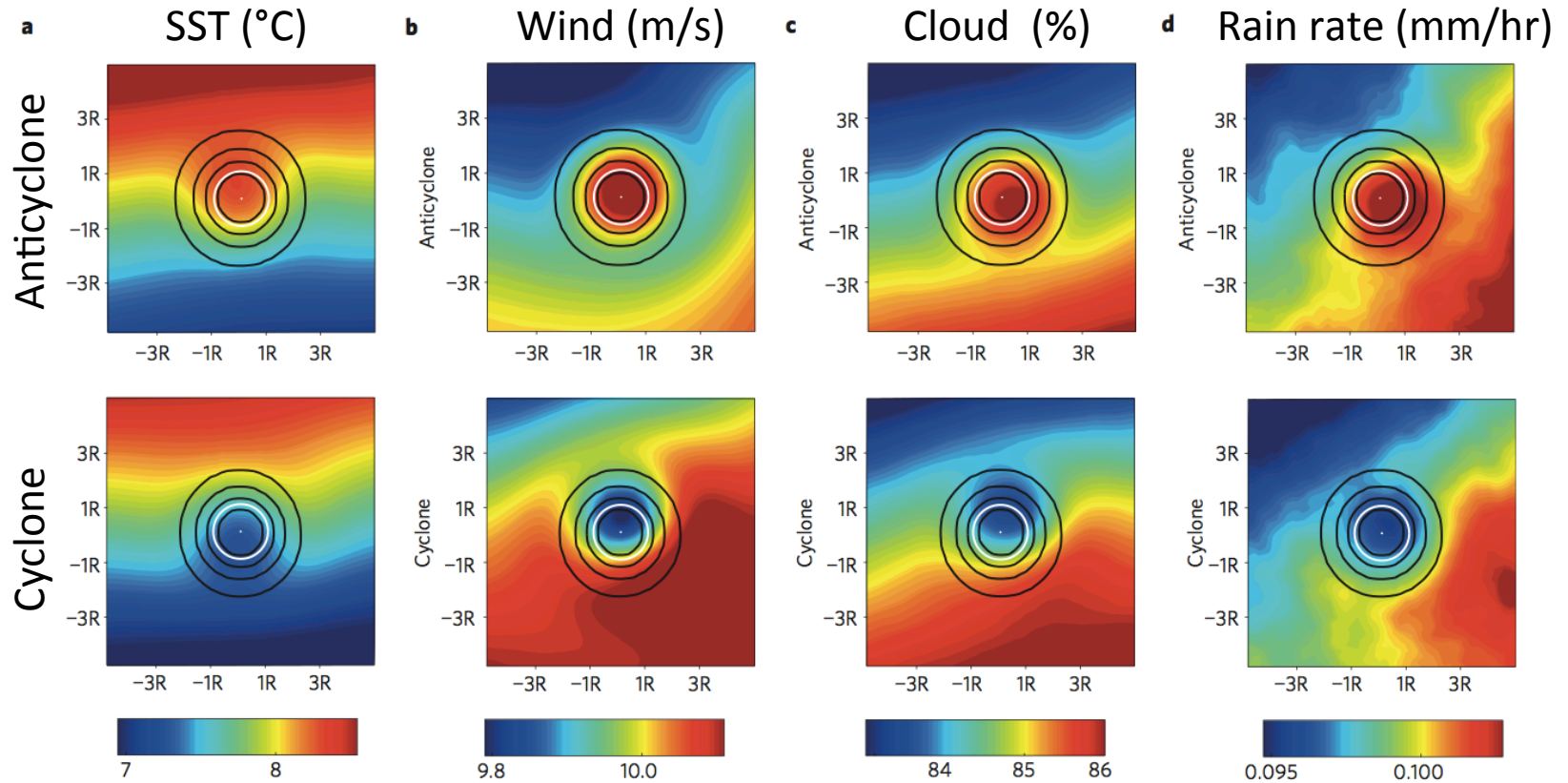
Cyclone



Pressure gradient adjustment does not appear important but it may depend on the scale.

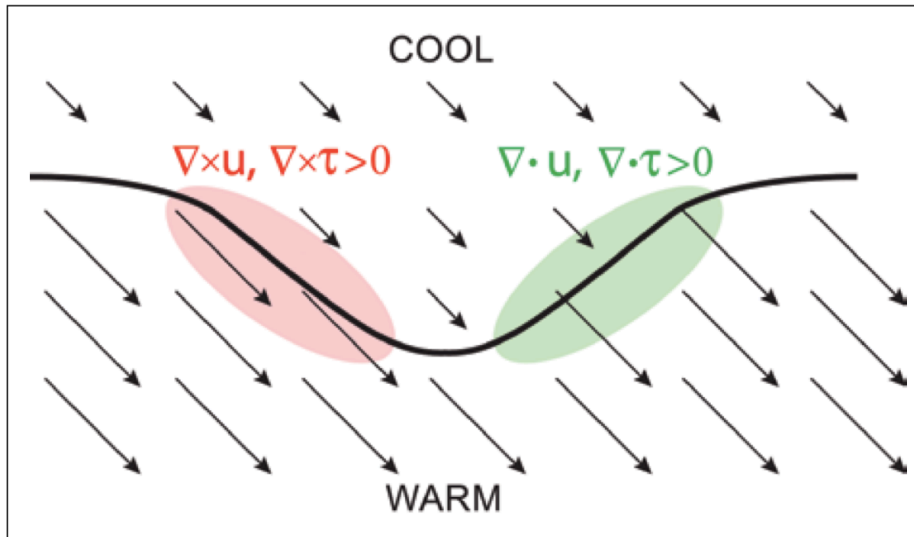
Observations: Mid-latitudes

Southern Ocean case (Frenger et al. 2013):



Wind direction →

See Bourras et al. (2004) for a North Atlantic example.



On the oceanic mesoscale:

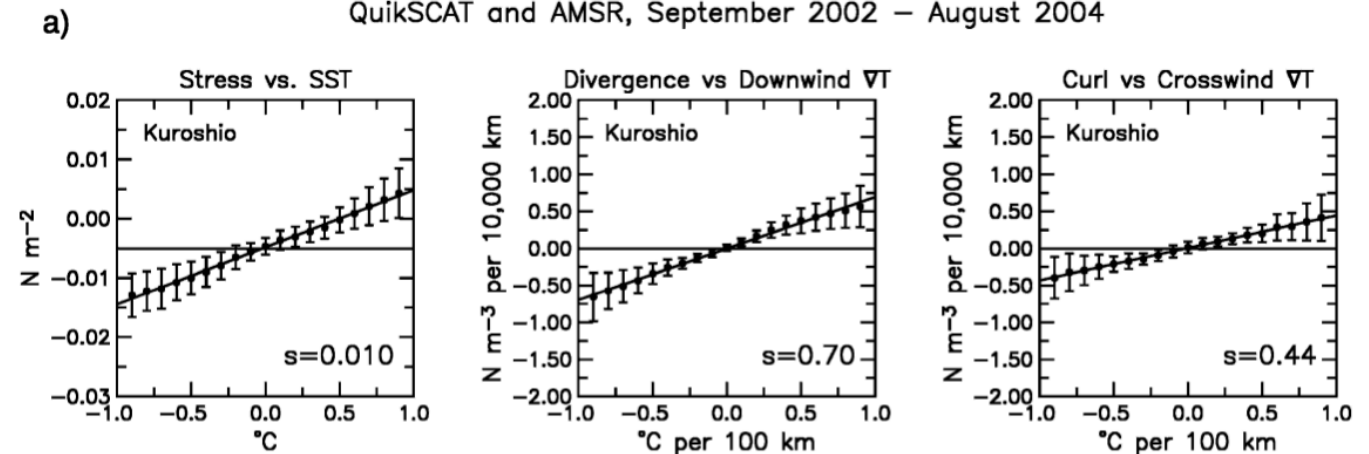
$$\bar{\mathbf{u}}' \propto T'$$

$$\nabla \times \bar{\mathbf{u}} \propto \text{crosswind } \nabla T$$

$$\nabla \cdot \bar{\mathbf{u}} \propto \text{downwind } \nabla T$$

Maloney and Chelton (2006)

QuikSCAT and AMSR, September 2002 – August 2004



Effects:

- $\vec{\tau}_a = \rho_a C_D |\vec{u}_a - \vec{u}_o| (\vec{u}_a - \vec{u}_o) \rightarrow$ eddies modify the wind stress they experience
- Surface fluxes modified by presence of eddies (latent, sensible, radiative)

Abyssal circulation, diapycnal mixing and Bottom boundary layer

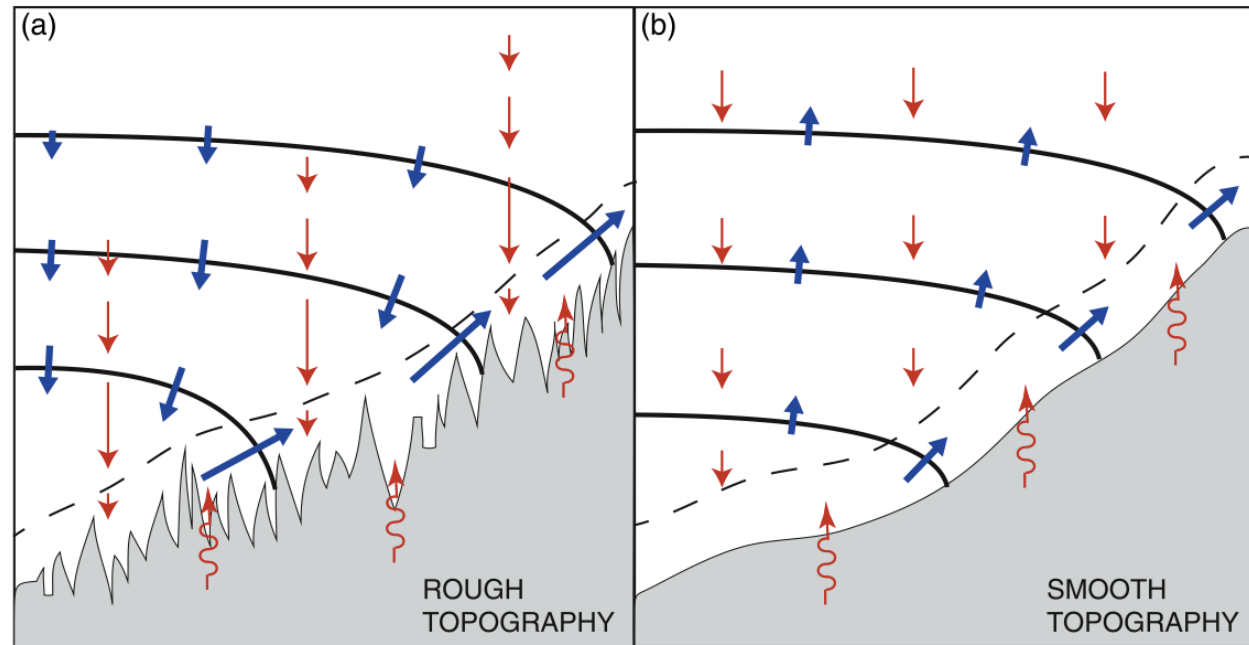
de Lavergne et al. 2016

1D advection-diffusion:

$$w^* N^2 = \frac{\partial}{\partial z} (k_V N^2)$$

Available energy for mixing

$$k_V N^2 = R_f \varepsilon_T$$



- Mixing-induced buoyancy fluxes
- Geothermal buoyancy fluxes
- Diapycnal transports

Observations: ε_T increases upward

→ Global overturning is closed within the BBL

$\varepsilon_T \downarrow$ with height
downwelling

$\varepsilon_T \uparrow$ with height
upwelling

But $k_V N^2 = 0$ @ bottom

- What do a ML scheme should do?
 - KPP does ocean interior mixing: Double-diffusion, shear mixing, etc...
 - GGL90 predicts TKE/Ks for the whole water column although in practice TKE is surface intensified
 - What about convective adjustment?
- Interaction between eddy parameterization and surface boundary layer
 - More generally interactions between subgrayscale processes (e.g. Eden & Olbers' work)
 - Horizontal v. vertical mixing in BL?
 - When do we introduce new parameterization or shut off a parameterization?
- Under sea ice/ice shelf?
 - Models use mixed layer schemes which have been developed/tested for open ocean conditions
- Bottom/benthic boundary layer? Downslope/overflows mixing scheme?
 - These are relatively unexplored
- Eddy-ABL interactions:
 - Does it matter ? Should it be parameterized in ocean-only and coarse climate models?