

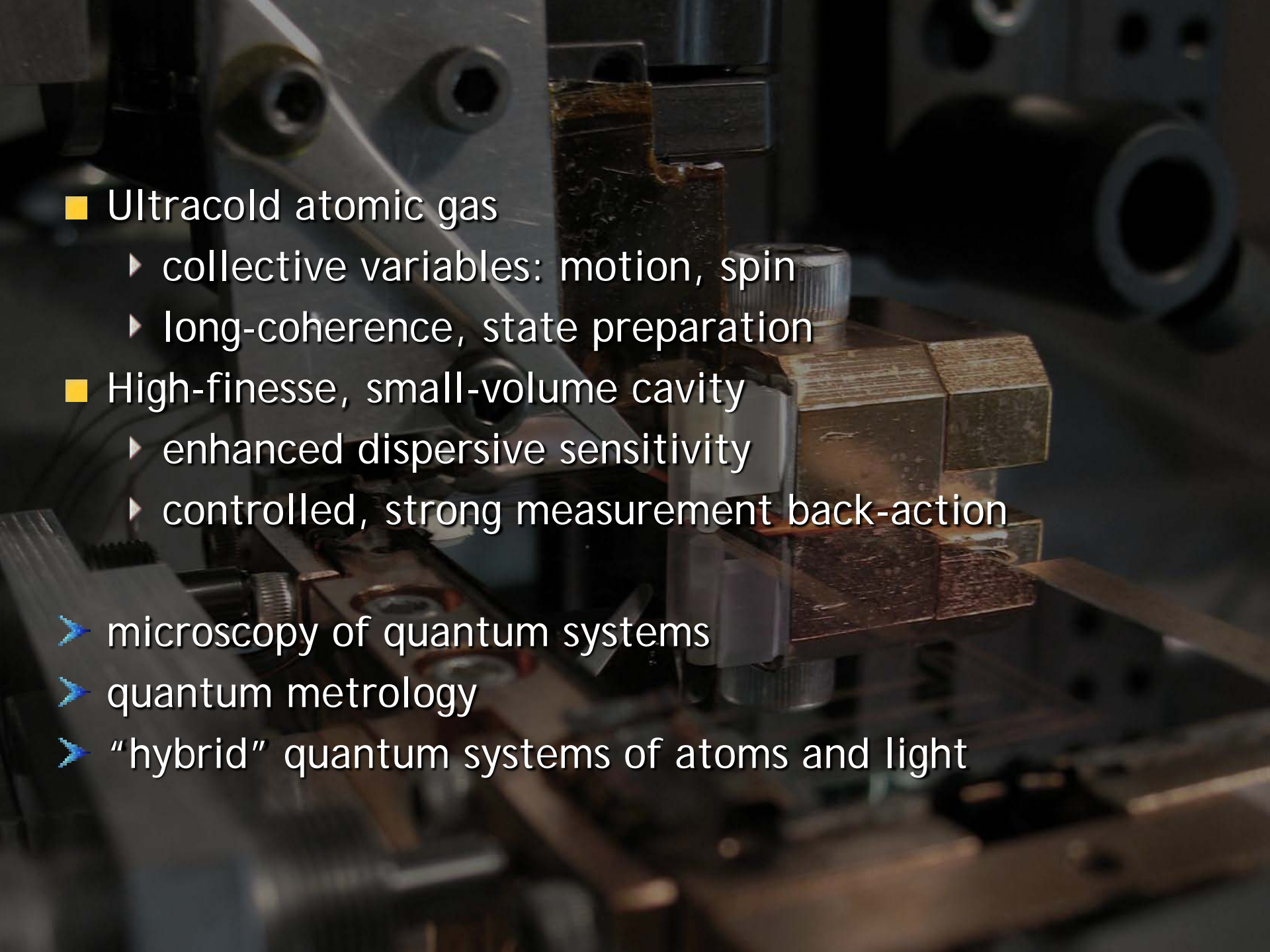


Optomechanics and spin dynamics of cold atoms in a cavity

Thierry Botter, Nathaniel Brahms, Daniel Brooks, Tom Purdy

Dan Stamper-Kurn
UC Berkeley

Lawrence Berkeley National Laboratory

- 
- Ultracold atomic gas
 - ▶ collective variables: motion, spin
 - ▶ long-coherence, state preparation
 - High-finesse, small-volume cavity
 - ▶ enhanced dispersive sensitivity
 - ▶ controlled, strong measurement back-action
 - microscopy of quantum systems
 - quantum metrology
 - “hybrid” quantum systems of atoms and light

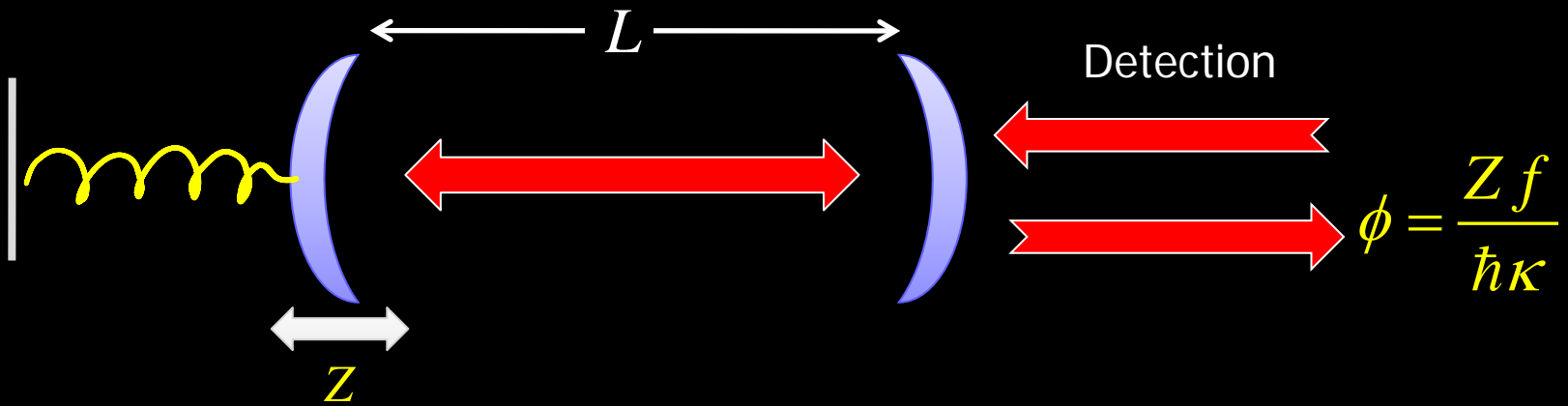
Why “beyond standard optical lattice” folks might care

- Optical probes of optical lattice systems
- Many-body physics in a Hamiltonian with backaction
- Many-body physics in a system under constant measurement
- Cavity-mediated interactions b/w mechanical elements, atomic spins, both: \neq super-exchange
- Motion \leftrightarrow spin
- Potential for using solid-state optomechanical systems as quantum simulators



Mechanics

One paradigm for cavity optomechanics



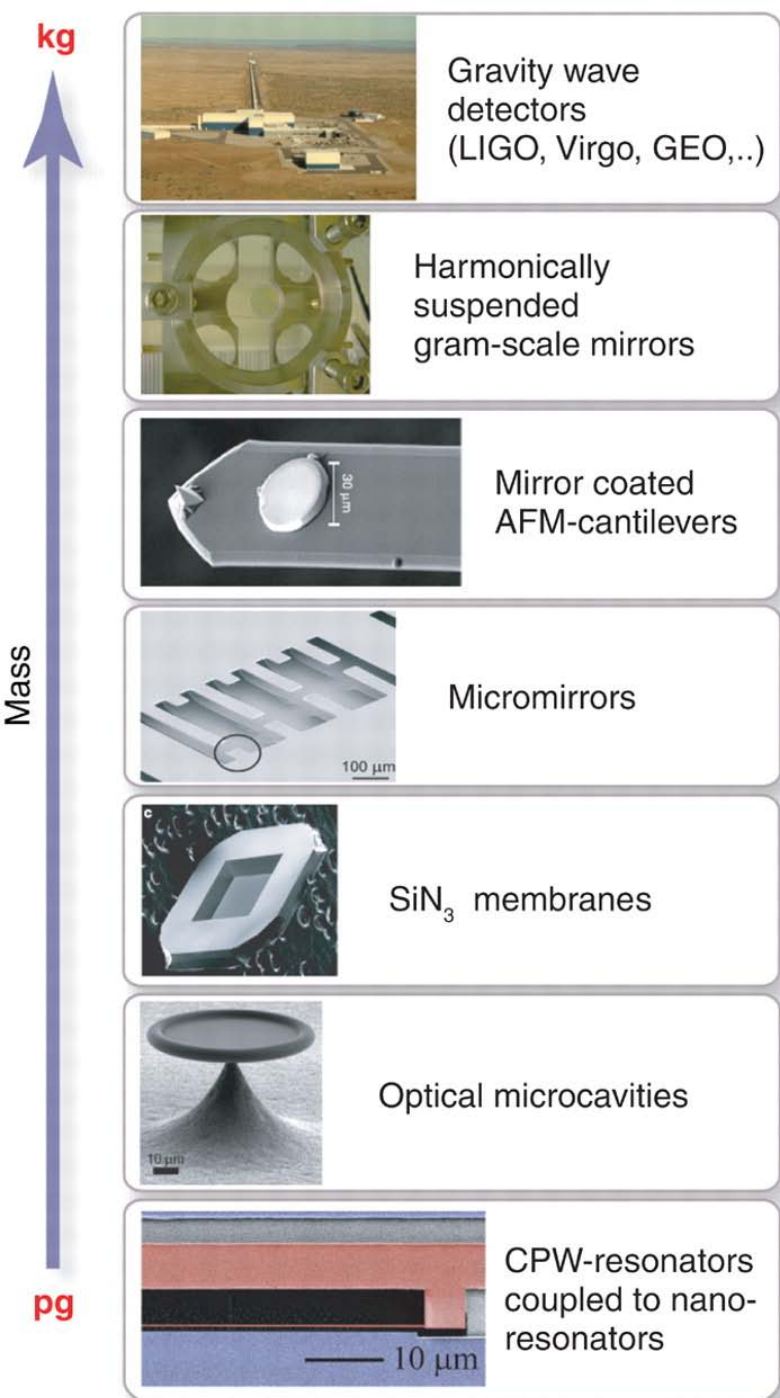
$$H = H_{osc} + H_{cav} + H_{in/out} - \boxed{Zfn}$$

force per photon:

$$f = \frac{hc}{\lambda L}$$

Cavity frequency shift due to oscillator displacement
(sensitivity to oscillator motion)

Optical force on oscillator
(optical back-action on oscillator)



Hz

Mechanical frequency ↓

MHz

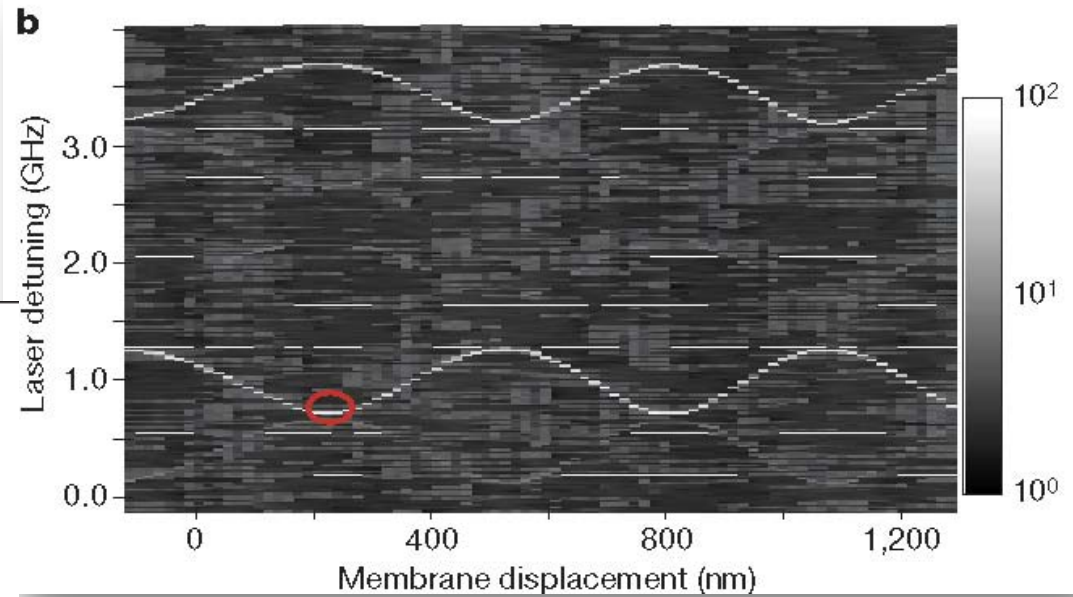
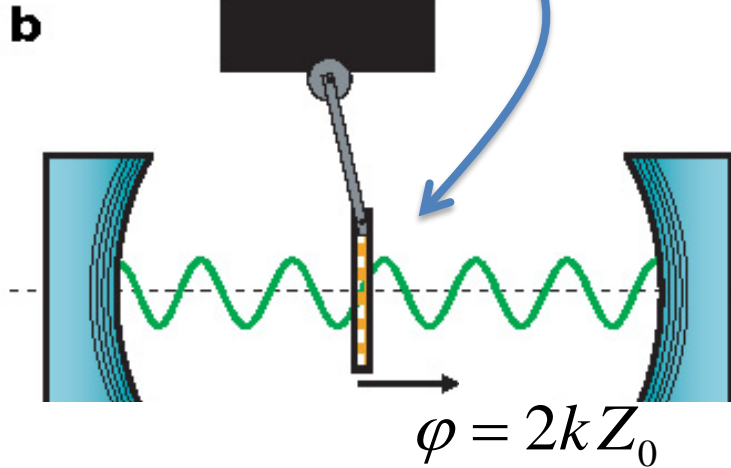
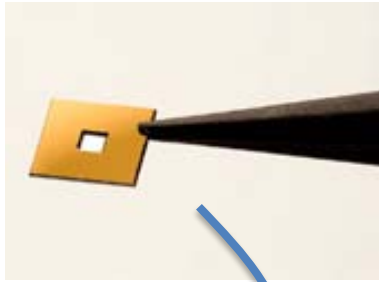
Common goals:

- Dominance of quantum fluctuations over thermal fluctuations
 - ◆ cooling mechanical oscillator to ground state
 - ◆ reaching quantum limits for sensitivity
- Study and use quantum effects
 - ◆ quantifying measurement backaction
 - ◆ squeezed light (pondermotive squeezing)
 - ◆ entanglement of macroscopic object with light

Common means:

- Better isolation from environment
- Colder starting points
- Stronger optomechanical coupling

Experimental Cavity Optomechanics: e.g. mechanical oscillator = membrane



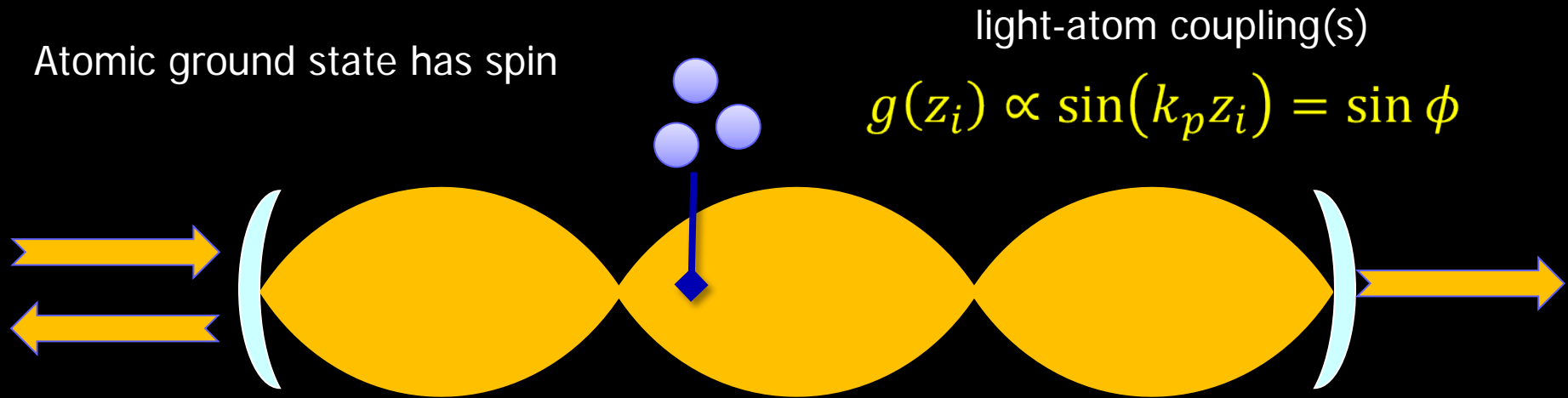
Harris group (Yale), Nature **452**, 72 (2008)

$$H_{om} = E(2\varphi)\hat{n} - F \sin(2\varphi)\hat{Z}_{CoM}\hat{n} - Fk \cos(2\varphi)\hat{Z}_{CoM}^2\hat{n} + \dots$$

linear coupling:
optical spring, bistability,
ponderomotive squeezing...

quadratic coupling:
phonon QND, ...

Many atom cavity QED



Dispersive regime (detuning Δ_{ca} is large):

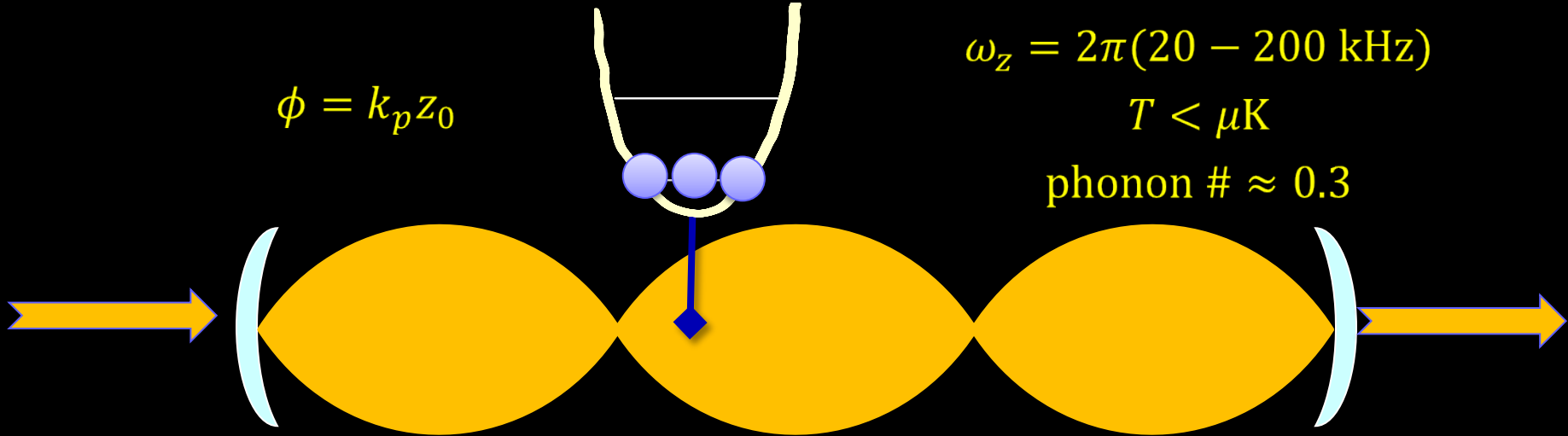
Cavity frequency shift per atom:

$$\frac{g_s^2(z_i)}{\Delta_{ca}} (1 + \epsilon \hat{k} \cdot \vec{s}_i)$$

Coupling to collective position variables
(atoms = membrane)

Coupling to collective spin variables
(circular birefringence, magneto-optics)

Tunable cavity optomechanics with cold atoms



$$H_{om} = \sum_{atoms} \frac{\hbar g^2(\hat{z}_i)}{\Delta_{ca}} \hat{n} \simeq \hbar \Delta_N^0 \hat{n} - F \sin(2\phi) \hat{Z}_{COM} \hat{n} - F k_p \cos(2\phi) [\hat{Z}_{COM}^2 + \hat{\sigma}^2] \hat{n}$$

- Tunability of optomechanical coupling (strength, type)
- Immediately in the quantum regime (ultracold)
- Dominated by quantum radiation pressure fluctuations (thermally isolated)
- Connected directly to basic theory (quantum optics, atomic physics)

Granular regime of optomechanics

define dimensionless granularity parameter:

$$\varepsilon = \frac{Z_{SQL}}{\delta Z = \hbar \kappa / F}$$

zero-point position spread
measurement uncertainty from single photon

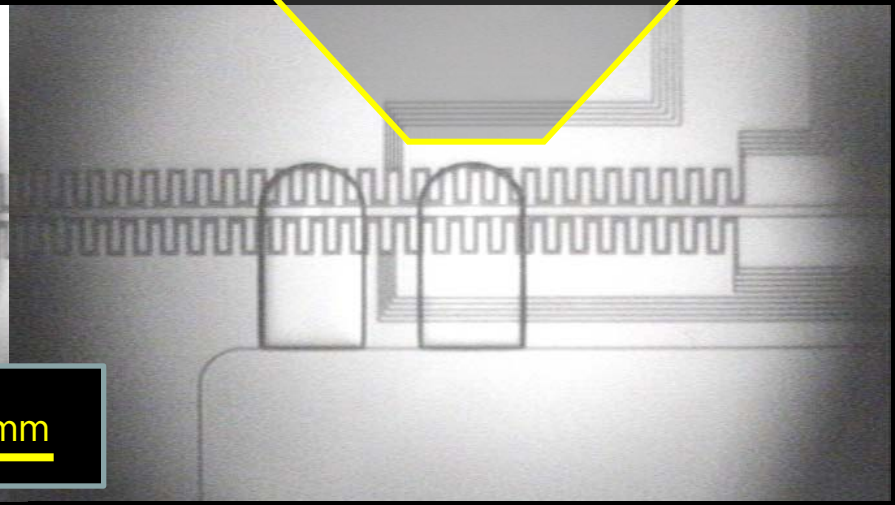
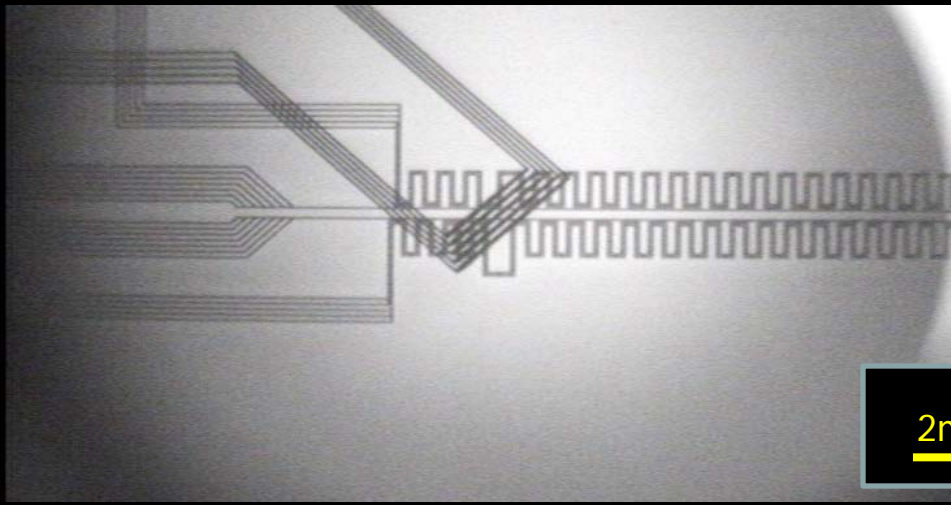
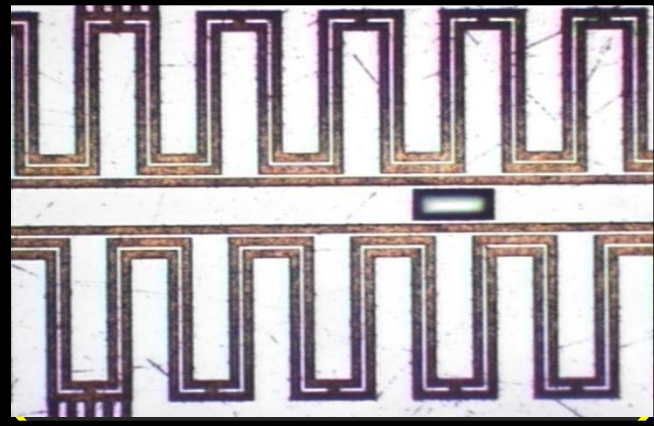
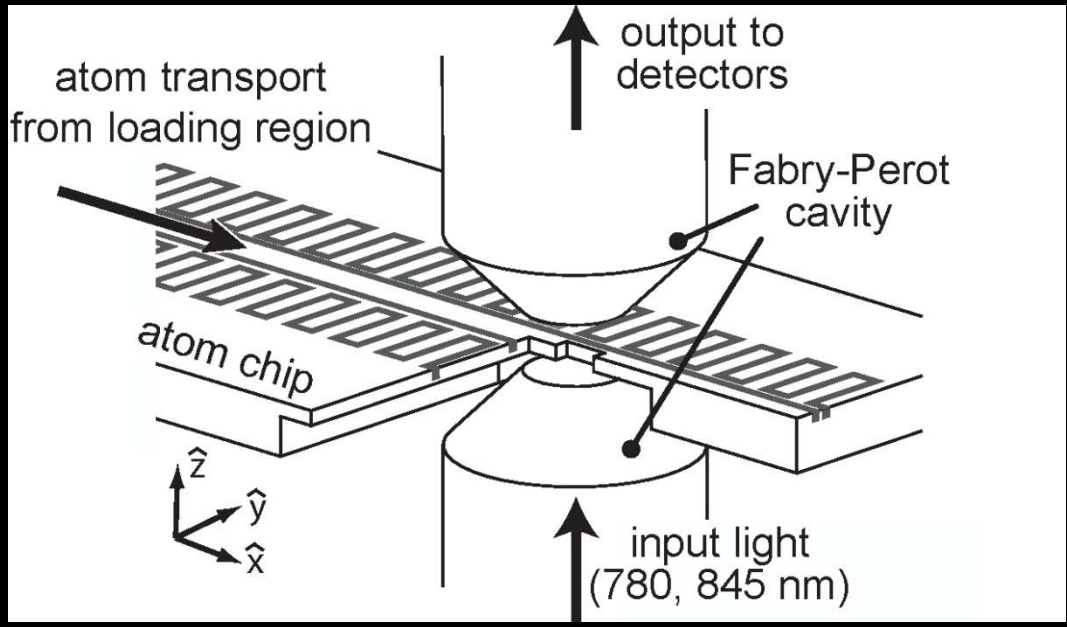
Does a single photon measure the cantilever to better than the SQL?

$$\varepsilon = \frac{F \times \frac{1}{\kappa}}{\hbar / Z_{SQL}}$$

(single-photon force) x (residence time of photon)
zero-point momentum spread

Does a single photon's kick significantly perturb the cantilever?

- ◆ Cantilever-based optomechanics: $\varepsilon = 10^{-7} - 10^{-5}$
- ◆ Atoms-based optomechanics: $\varepsilon = 0.01 - 10$

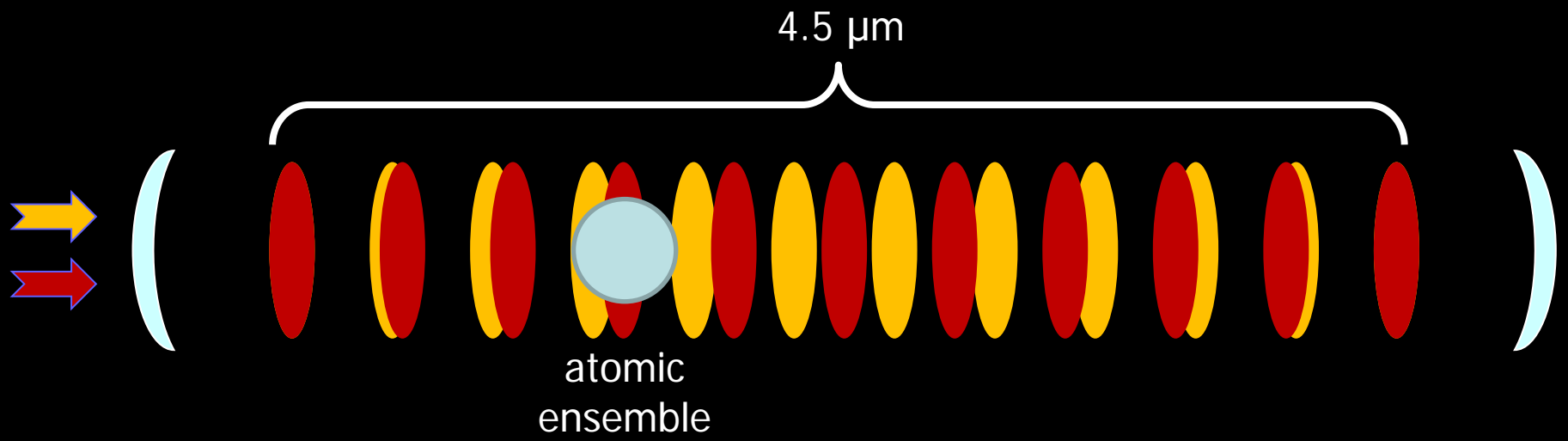


2mm

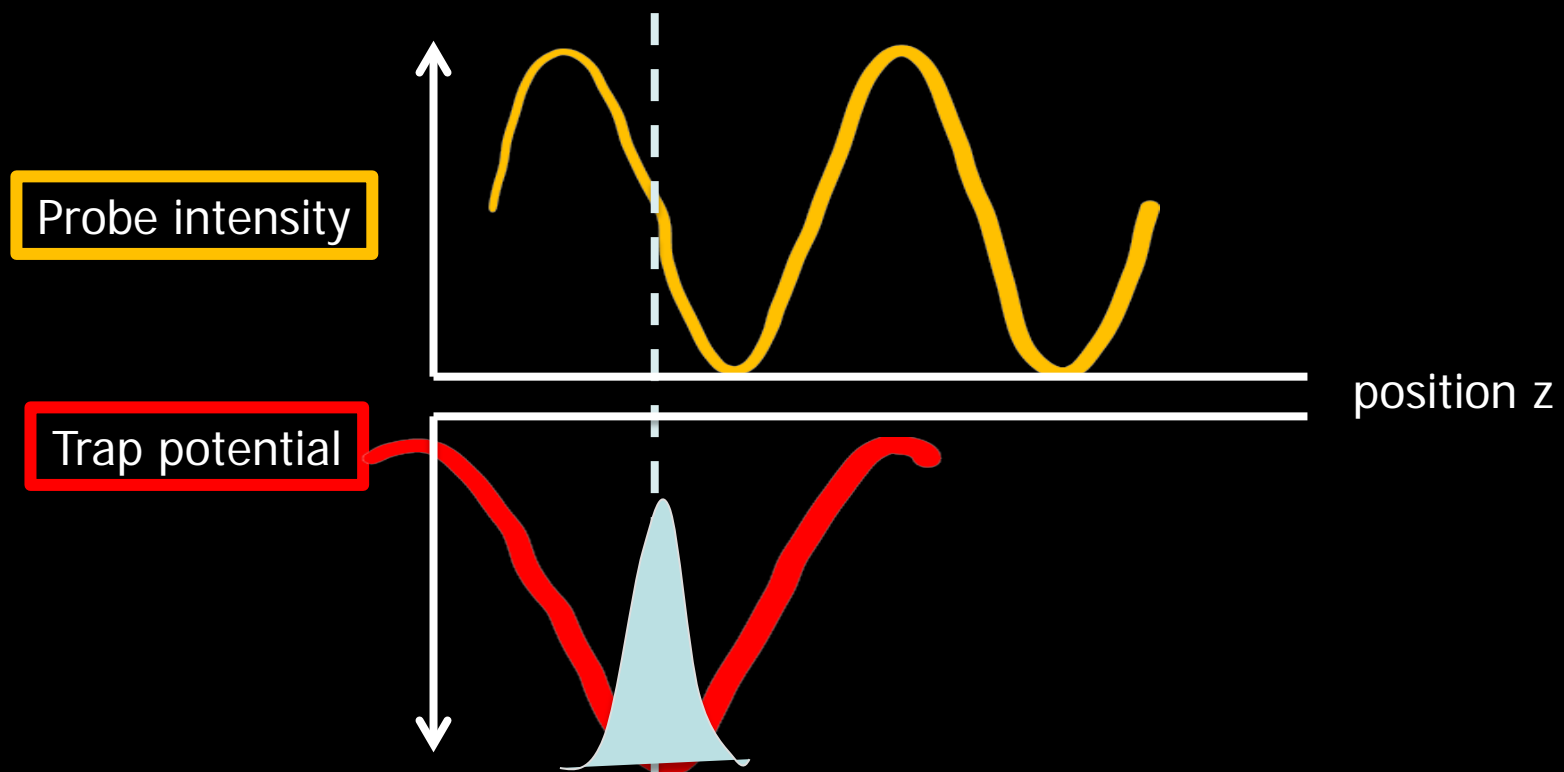
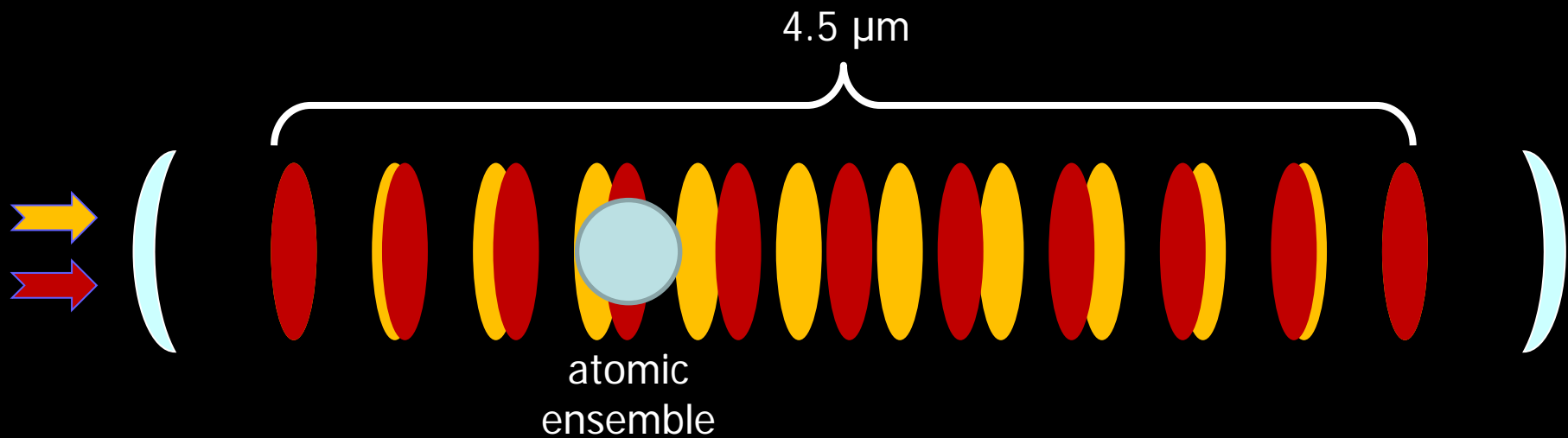
MOT Loading

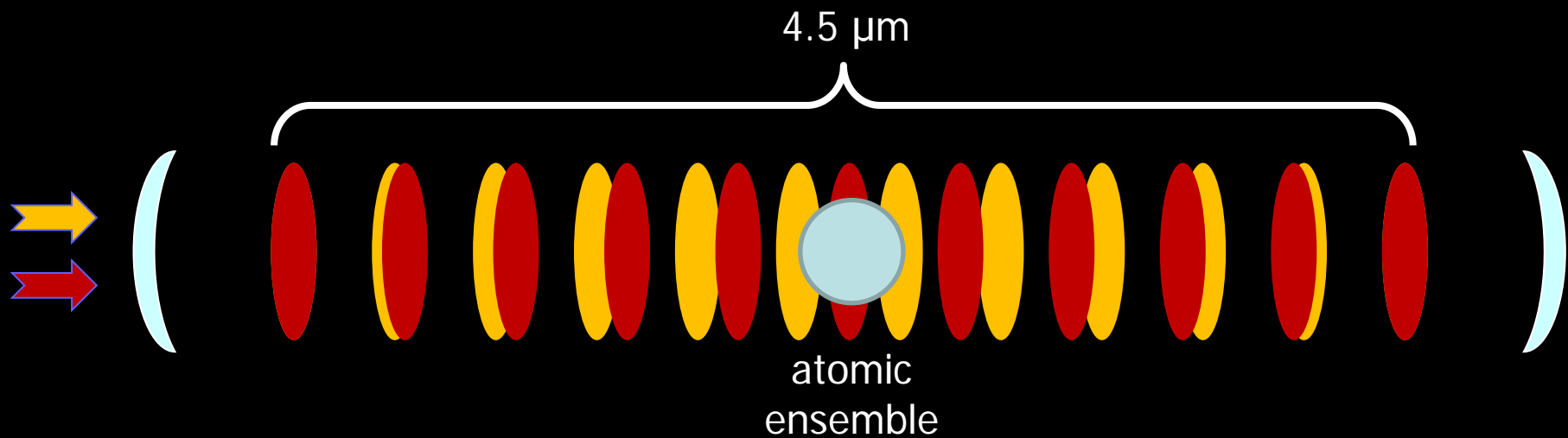
Conveyor Belt

Cavity Locations

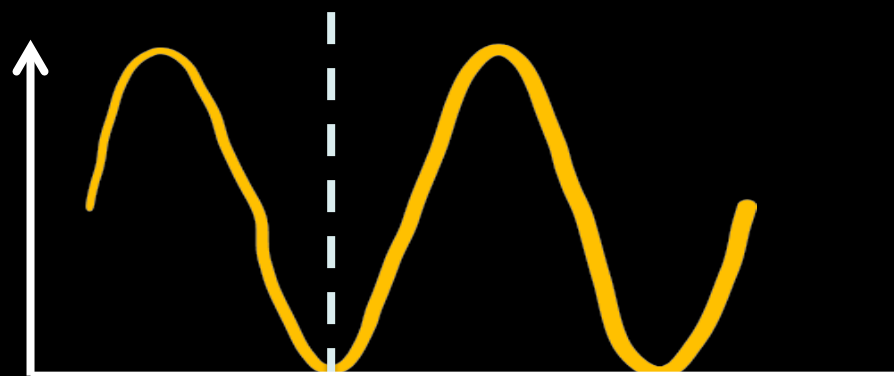


- pre-position with magnetic trap
- load into 845 nm optical trap
- probe with 780 nm light
(either red [attractive] or blue [repulsive] detuned)



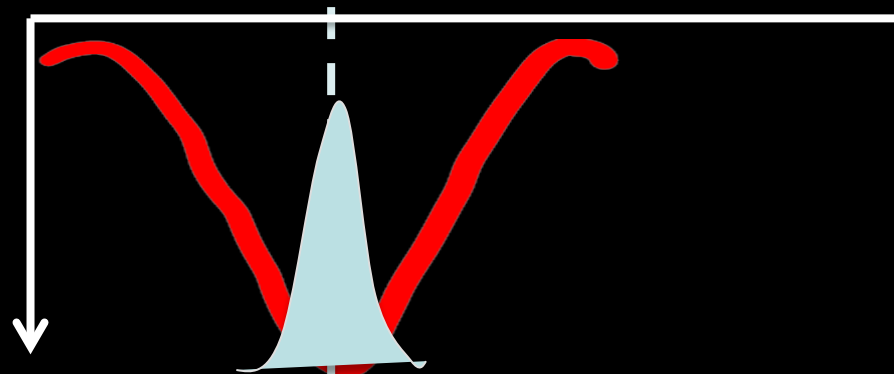


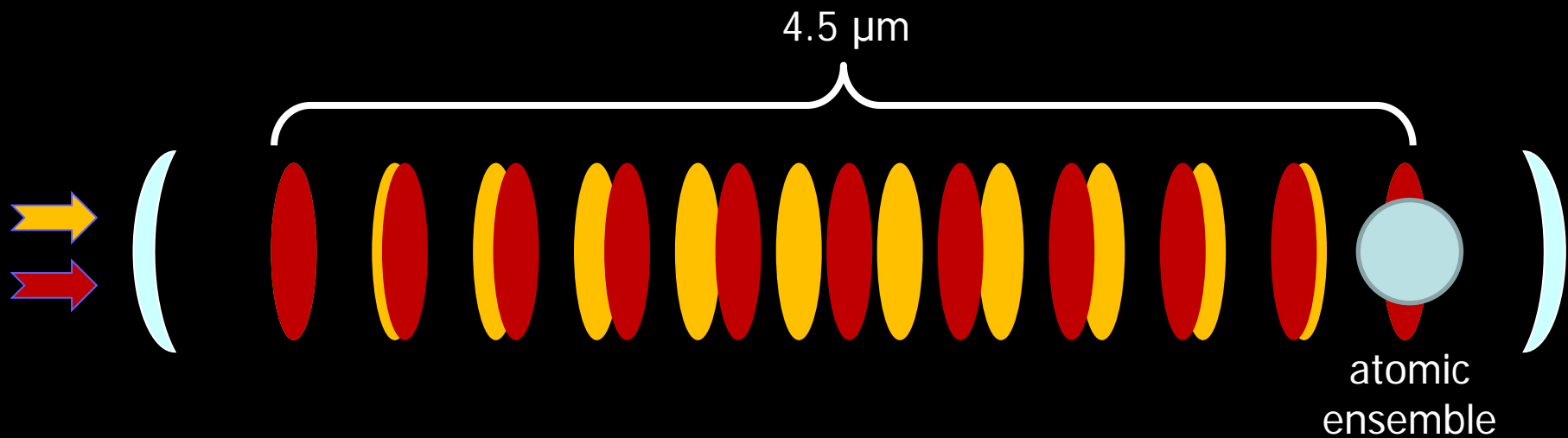
Probe intensity



position z

Trap potential



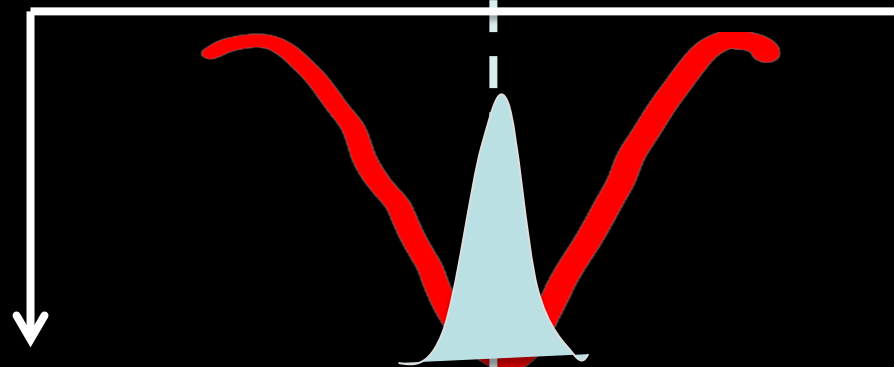


Probe intensity

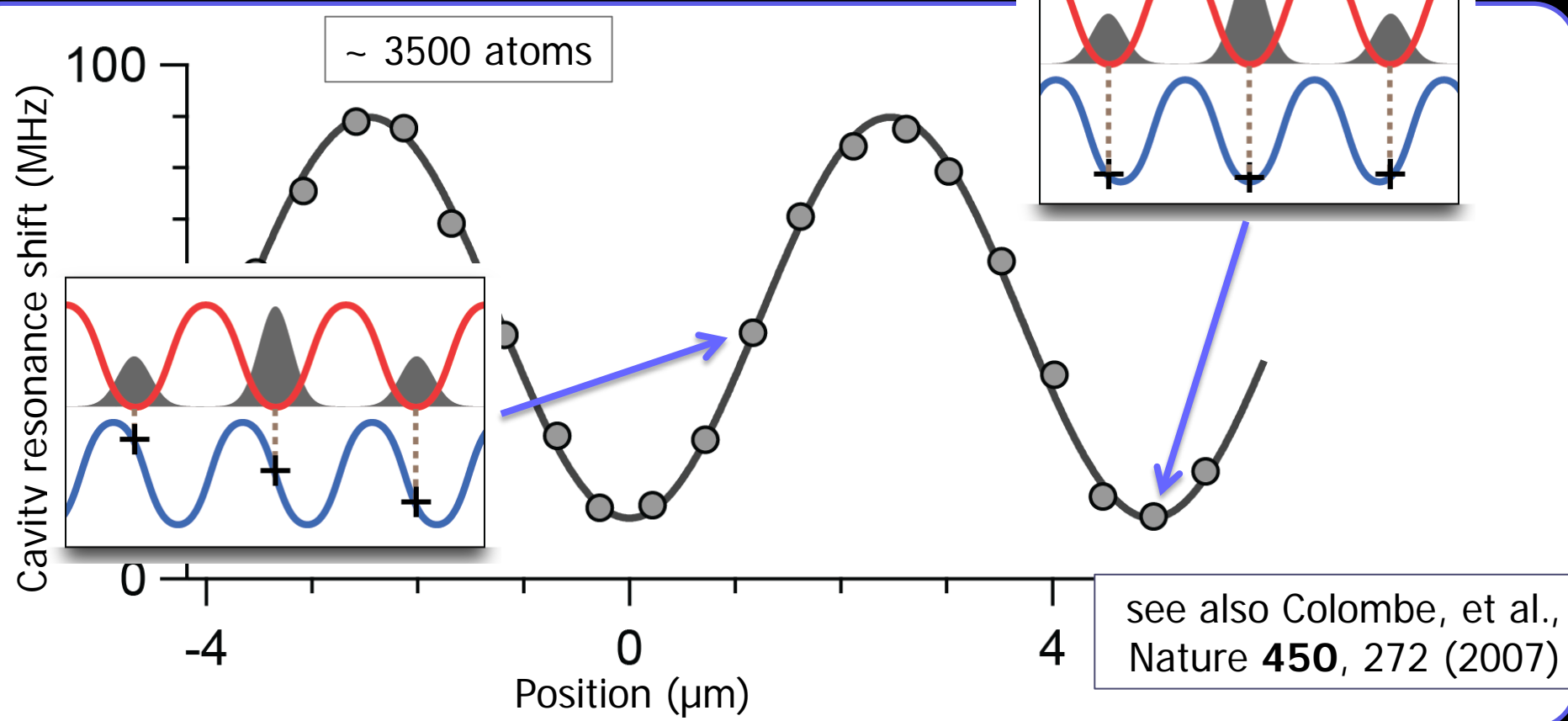
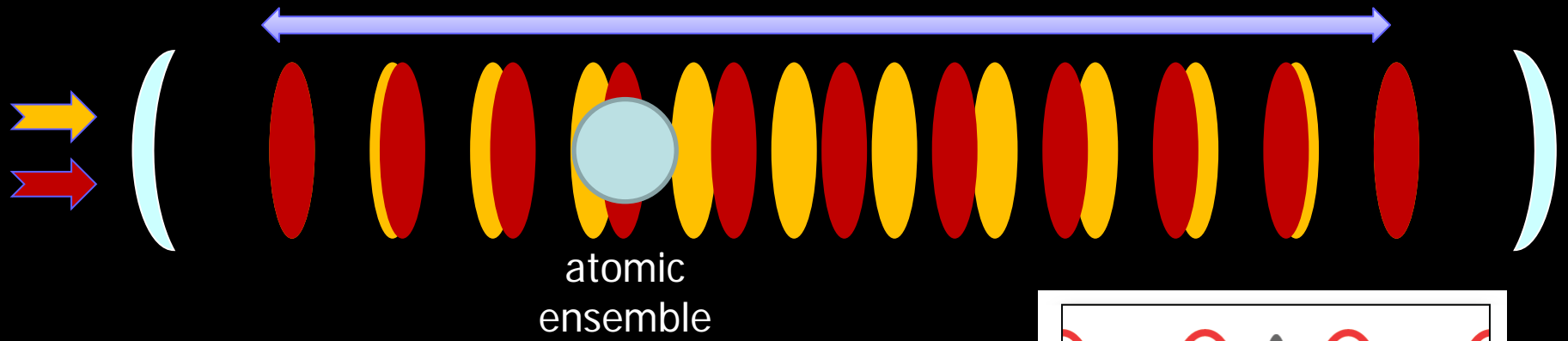


position z

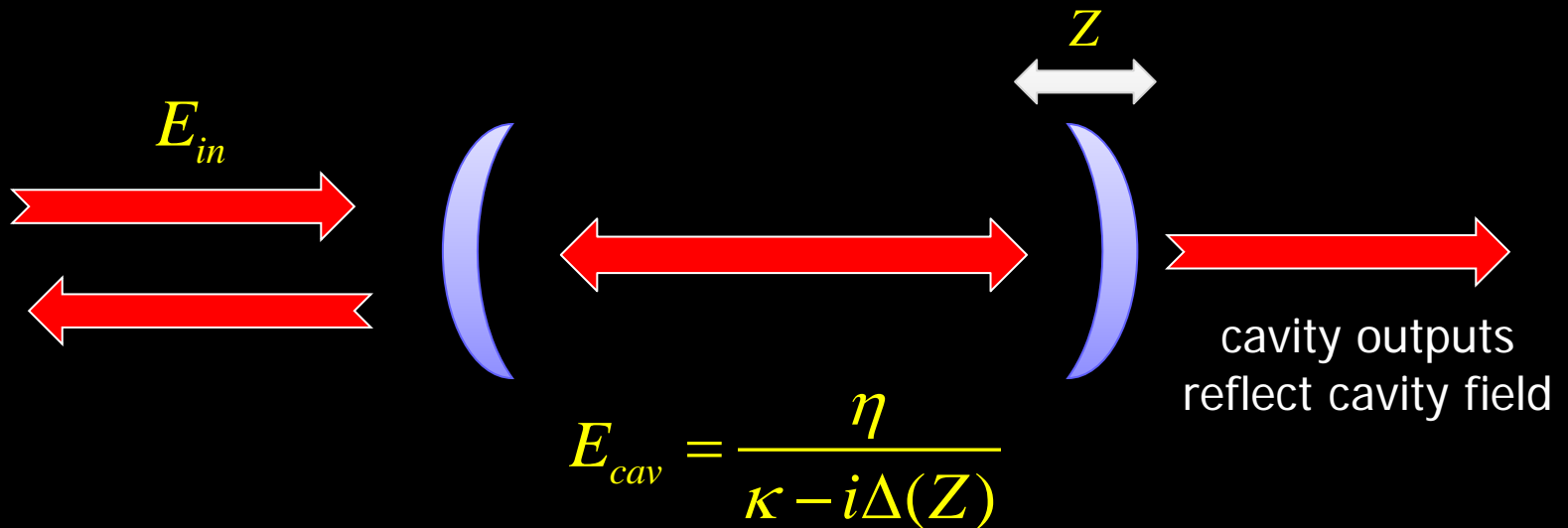
Trap potential



Sub-wavelength positioning within the cavity mode

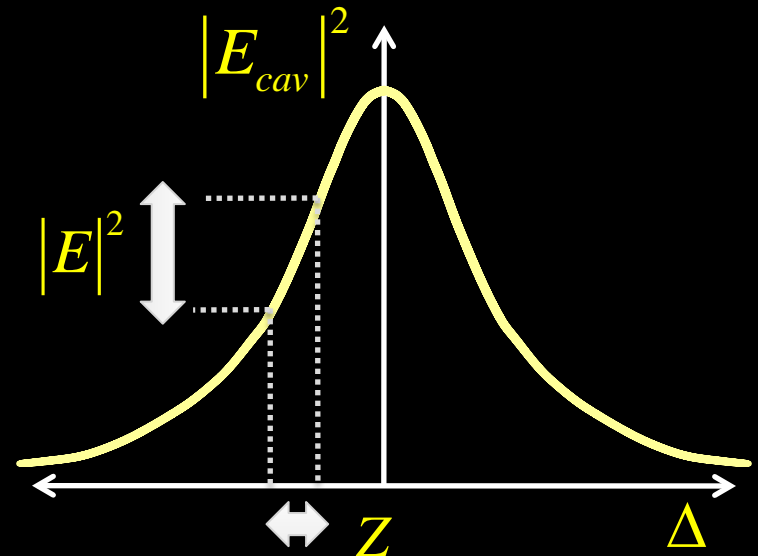


Motion detection in opto-mechanics

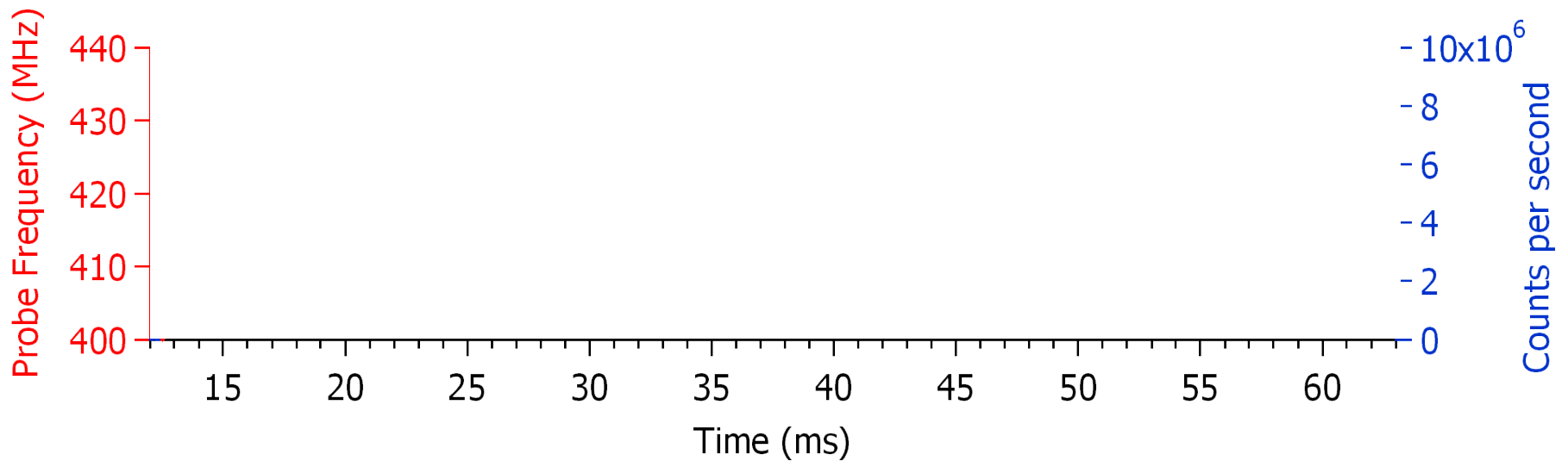
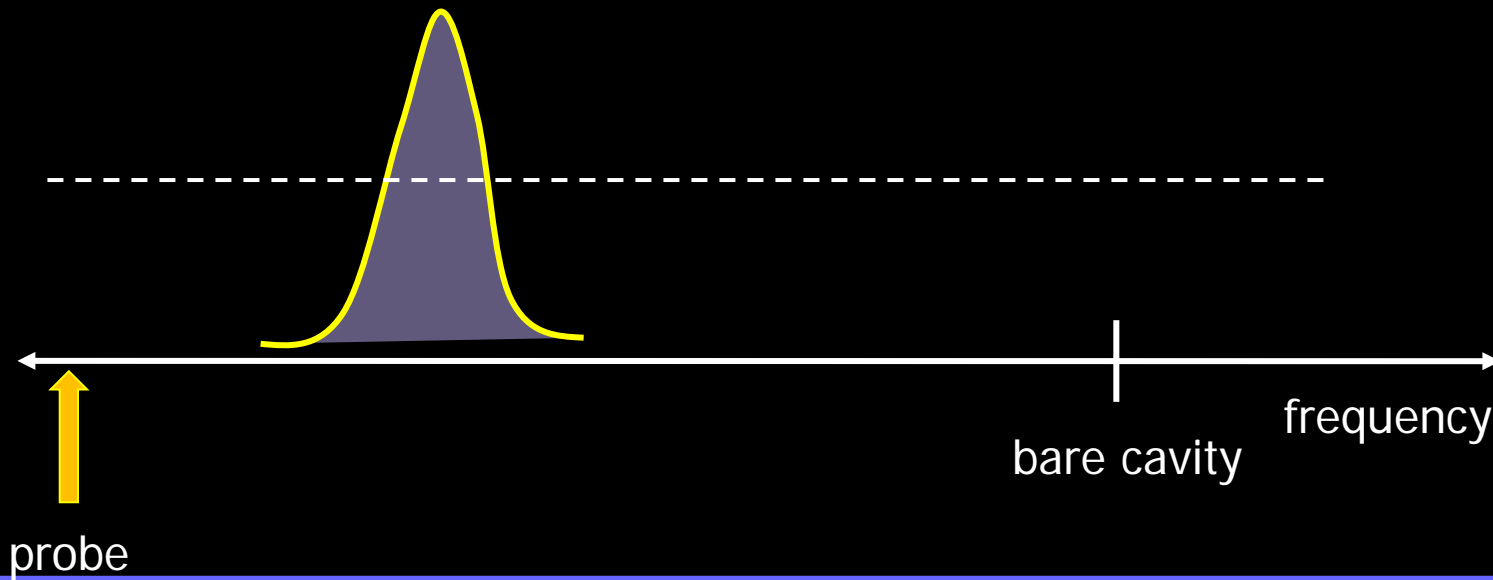


How to measure?

1. Optimal: Measure field quadrature
2. Sub-optimal: Measure intensity (transmission, fluorescence)

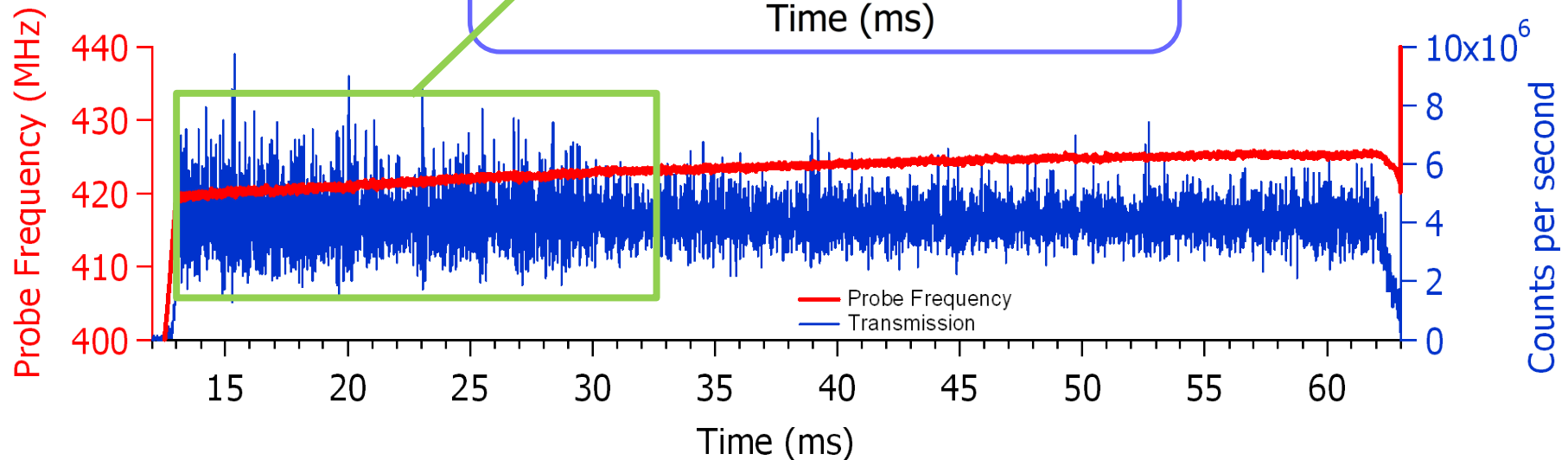
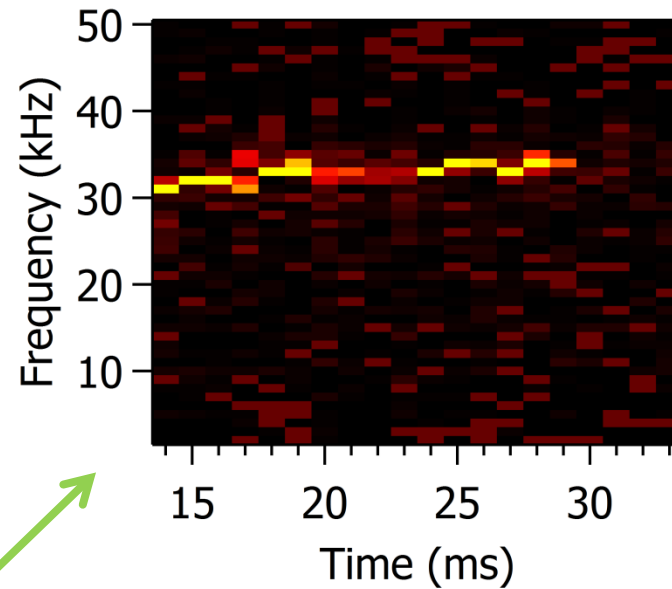


Direct observation of collective atomic motion: probe side-lock



Direct observation of collective atomic motion: probe side-lock

Spectral power in 1ms intervals



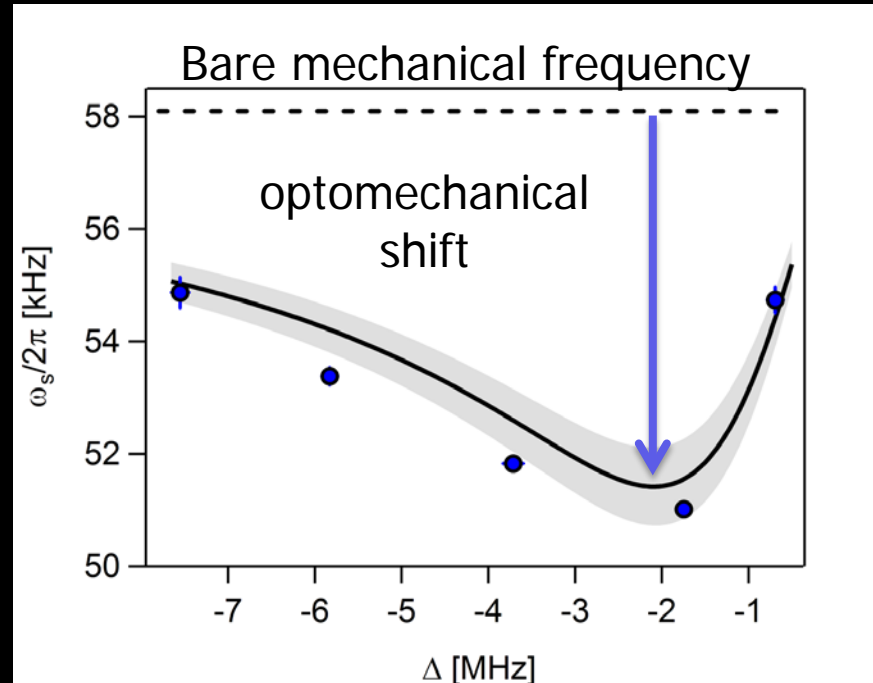
Detection is influential

Consider:

- Dynamical optomechanical frequency shift: the “optical spring”

force on collective
variable:

$$F_Z = Fn(Z) - M\omega_z^2 Z$$

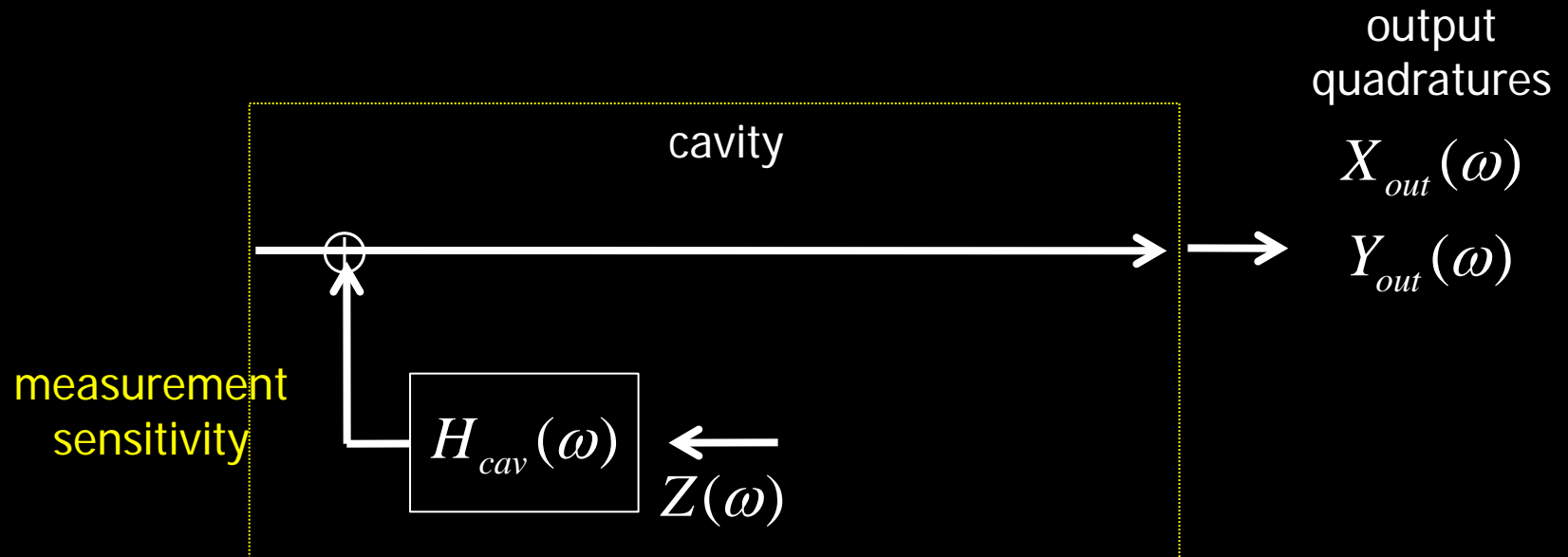


Purdy et al., PRL **105**, 133602 (2010)

- Q: Why is cantilever moving?
A: Radiation pressure fluctuations

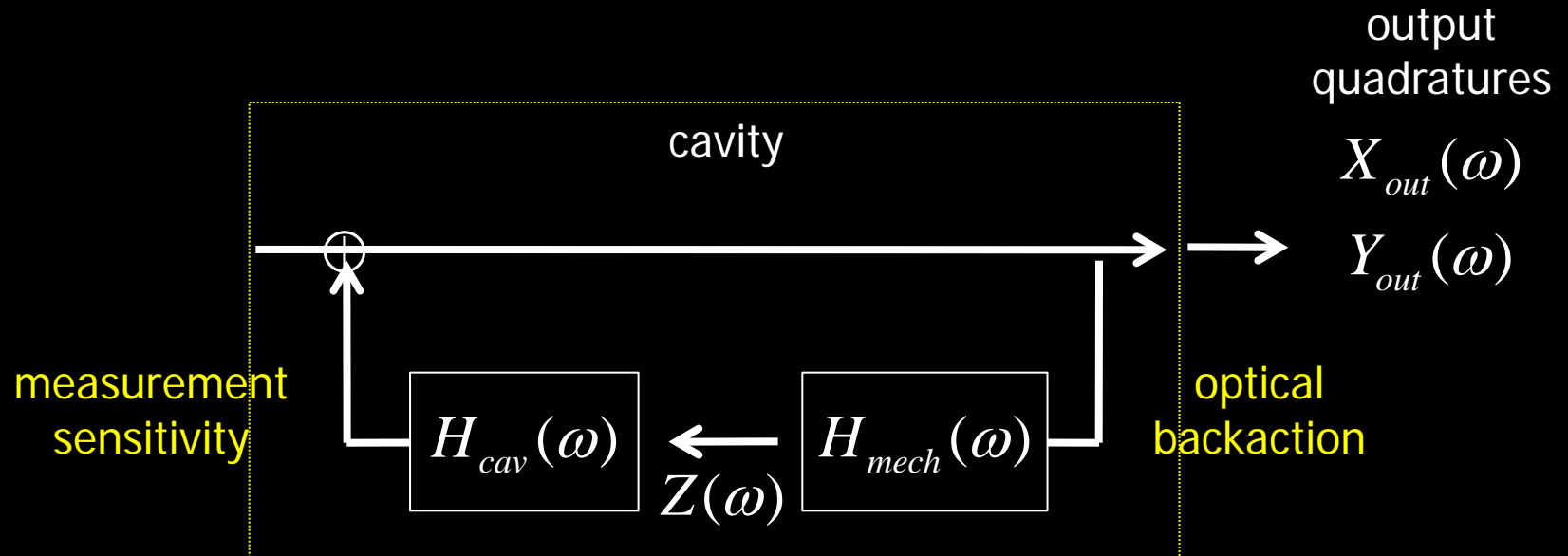
Measurement and feedback in cavity optomechanics

- Measurement scenario



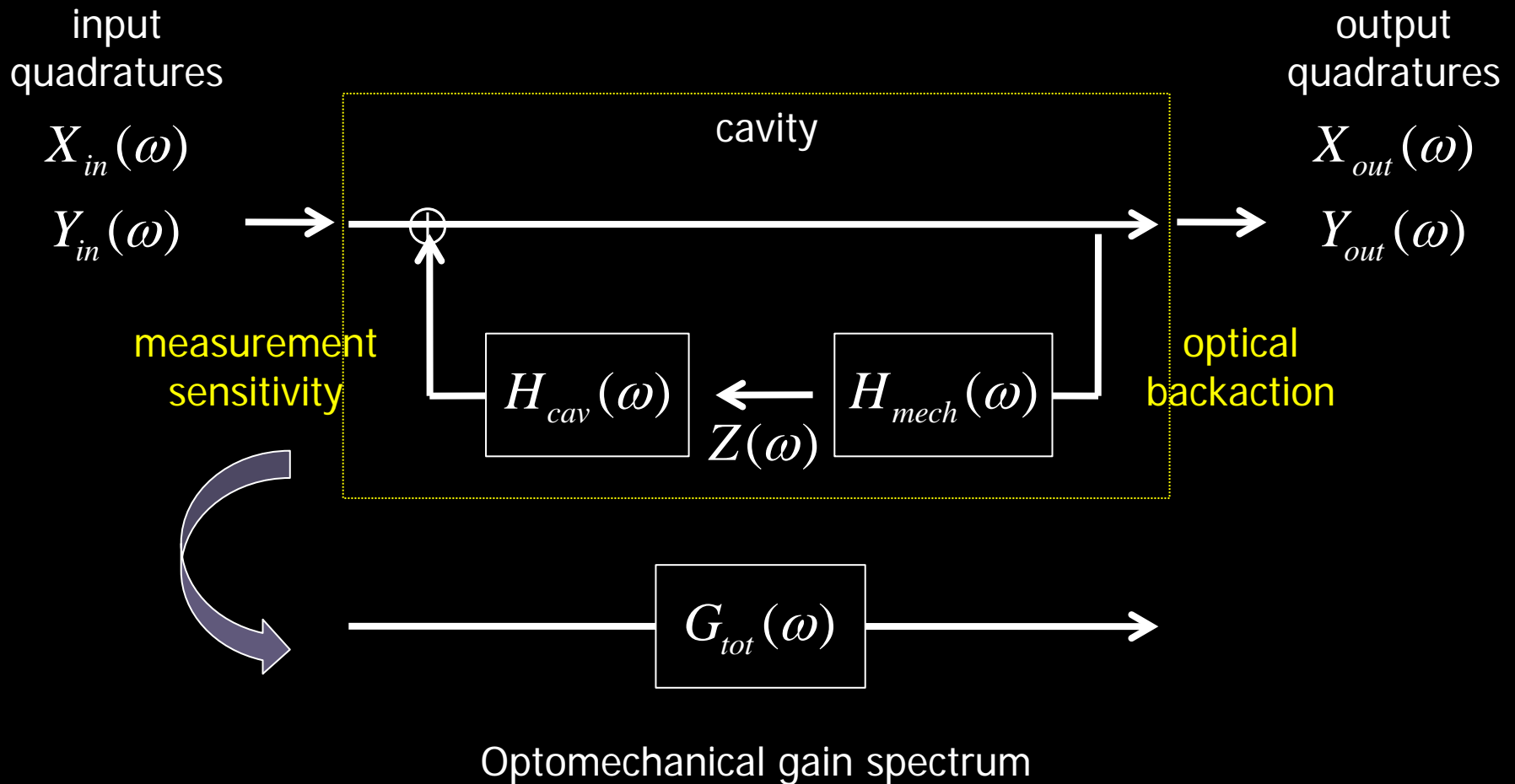
Measurement and feedback in cavity optomechanics

- Coherent backaction scenario
 - ◆ optomechanical frequency shift
 - ◆ cavity nonlinearity and bistability



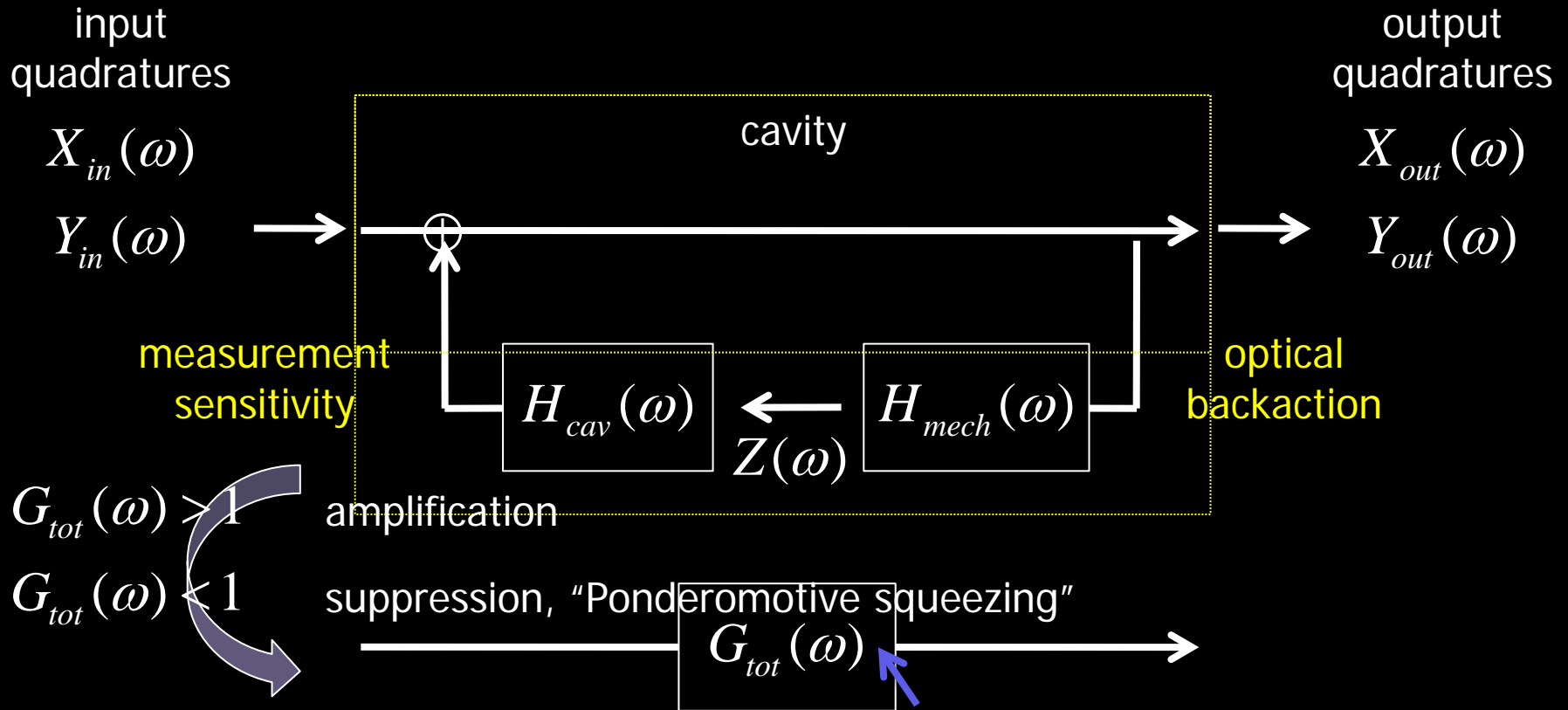
Measurement and feedback in cavity optomechanics

- Optomechanical amplification/squeezing of light



Measurement and feedback in cavity optomechanics

- Optomechanical amplification/squeezing of light

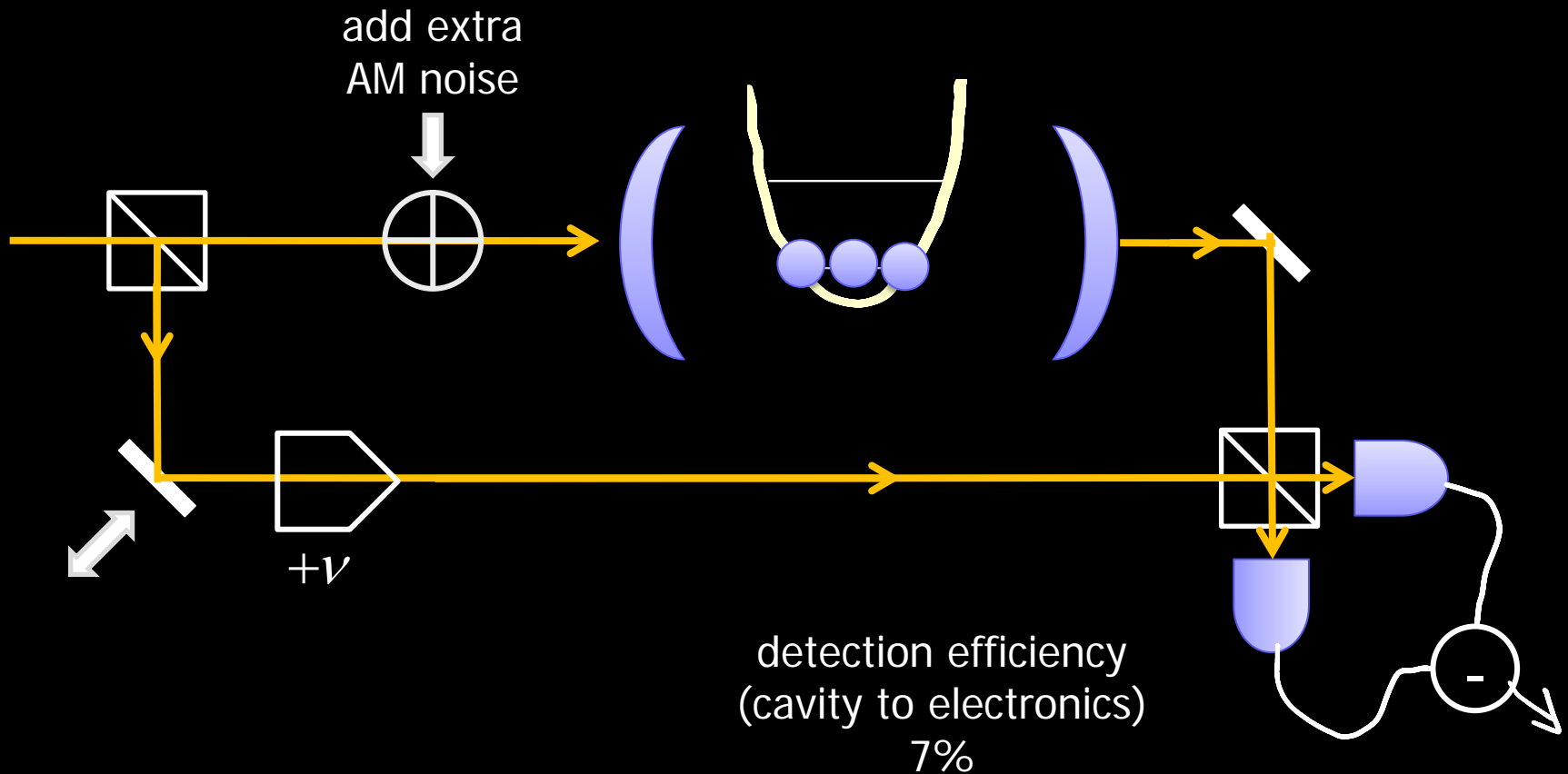


Key component in future of LIGO: from detector to observatory

See Kimble et al., PRD **62**, 022002 (2002).

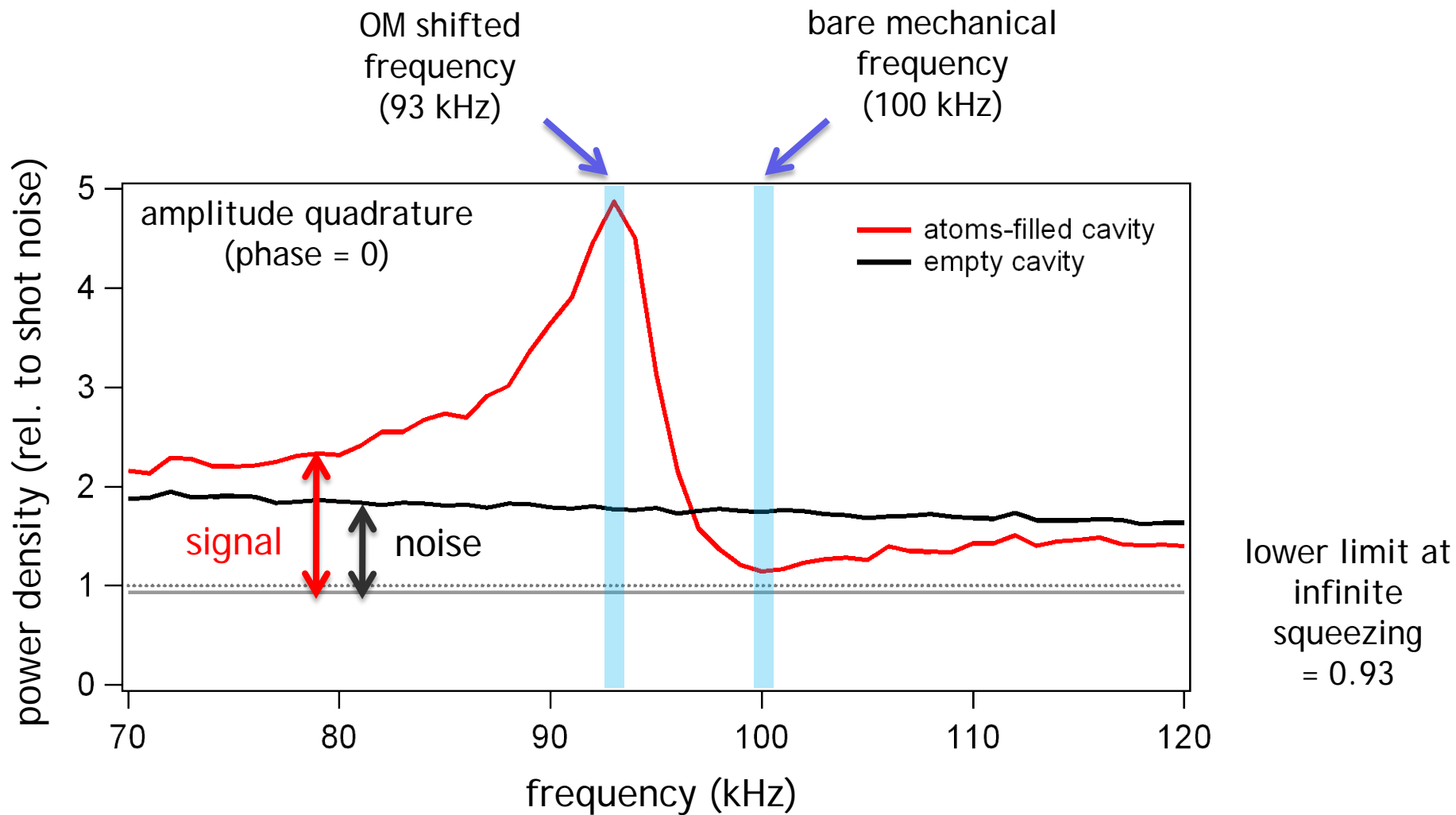
Measurements of optomechanical gain spectrum

see Marino et al., PRL **104**, 073601 (2010); Verlot et al., *ibid*, 133602.



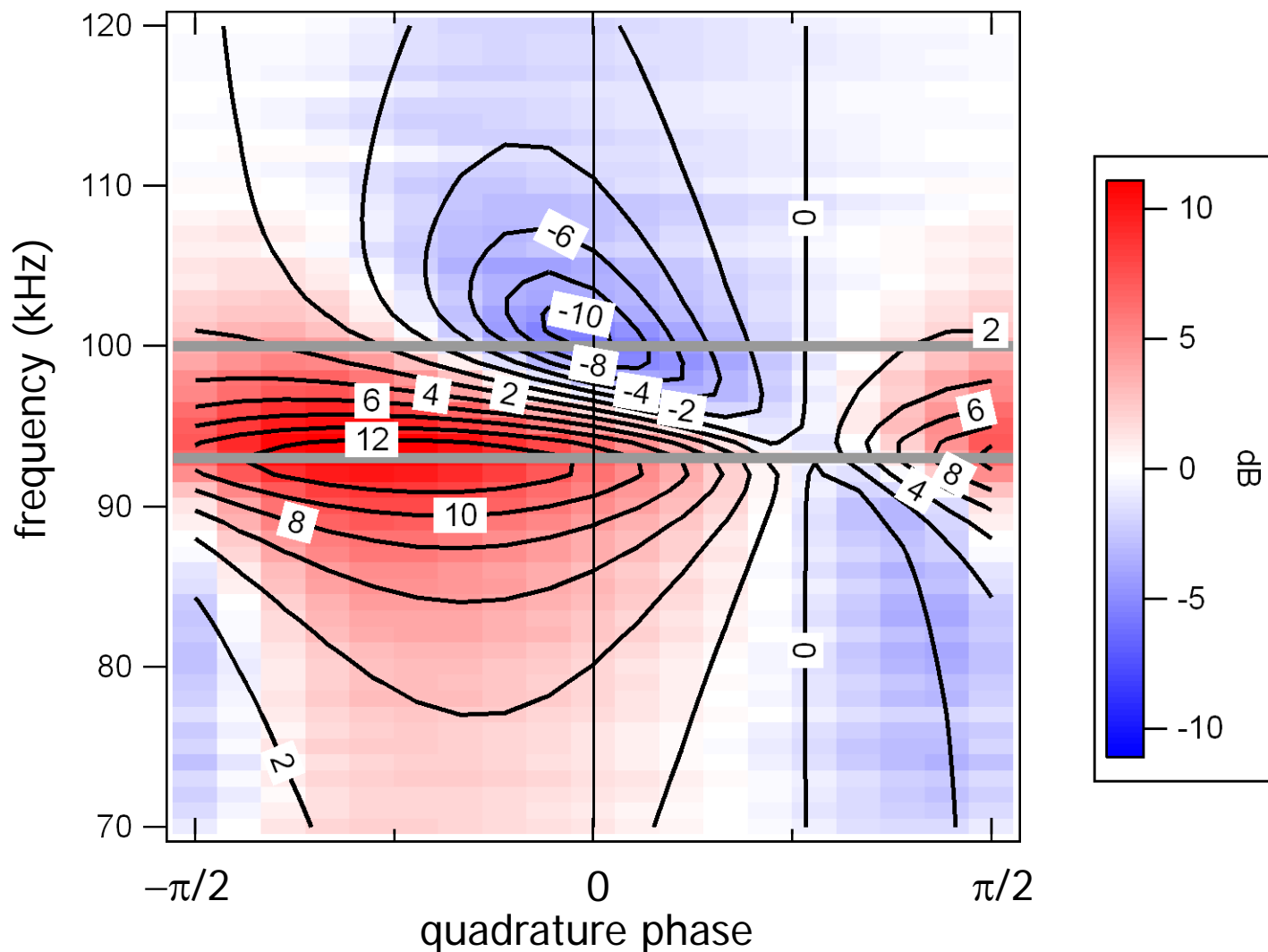
Spectral record of noise-driven atomic motion

■ added noise ~ 10x shot noise



Measurements of optomechanical gain spectrum

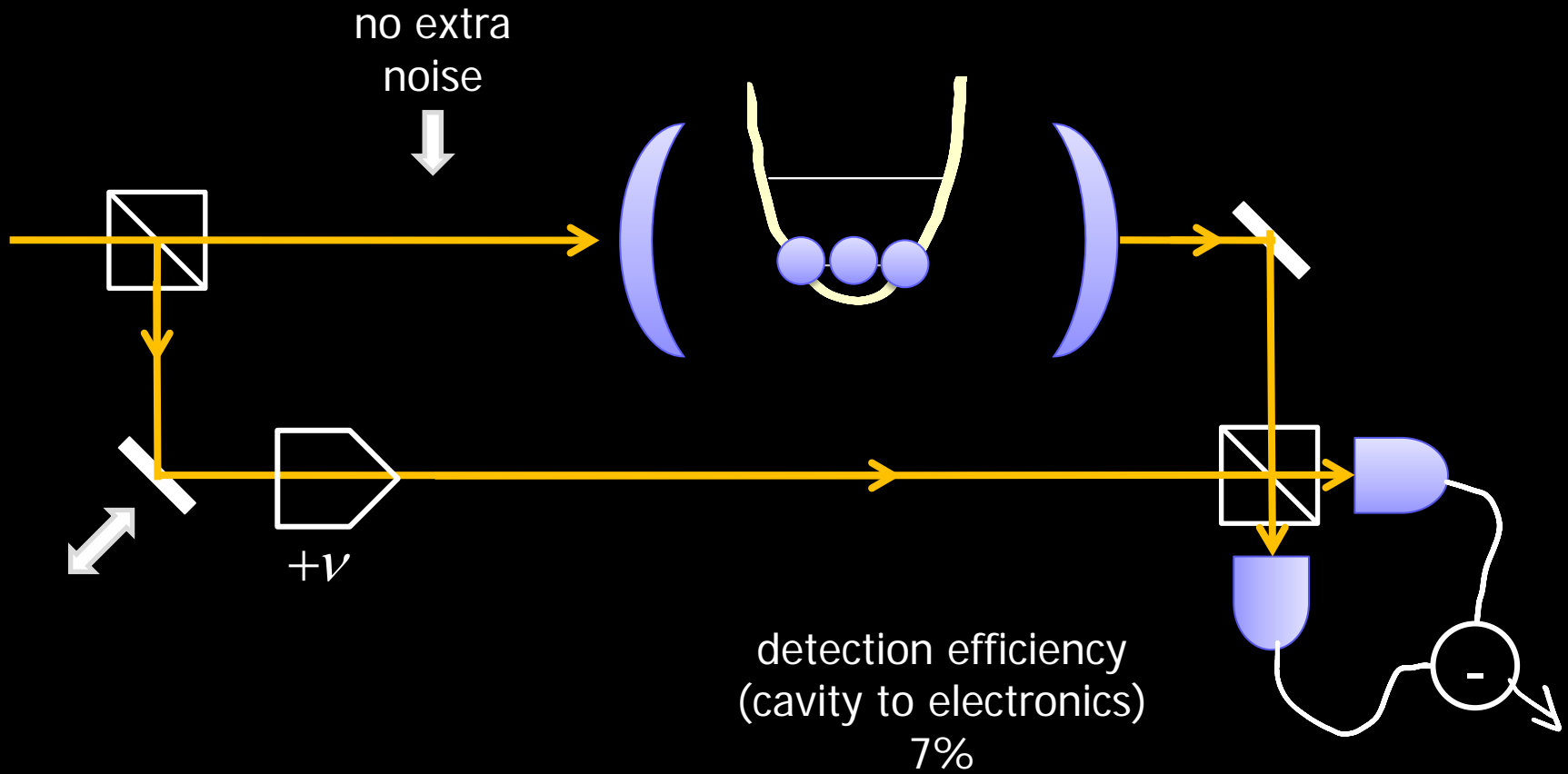
- adjusted for detector efficiency



settings:
 $\Delta_{ca} = -2$ GHz
5000 atoms
 $\langle n \rangle = 2$

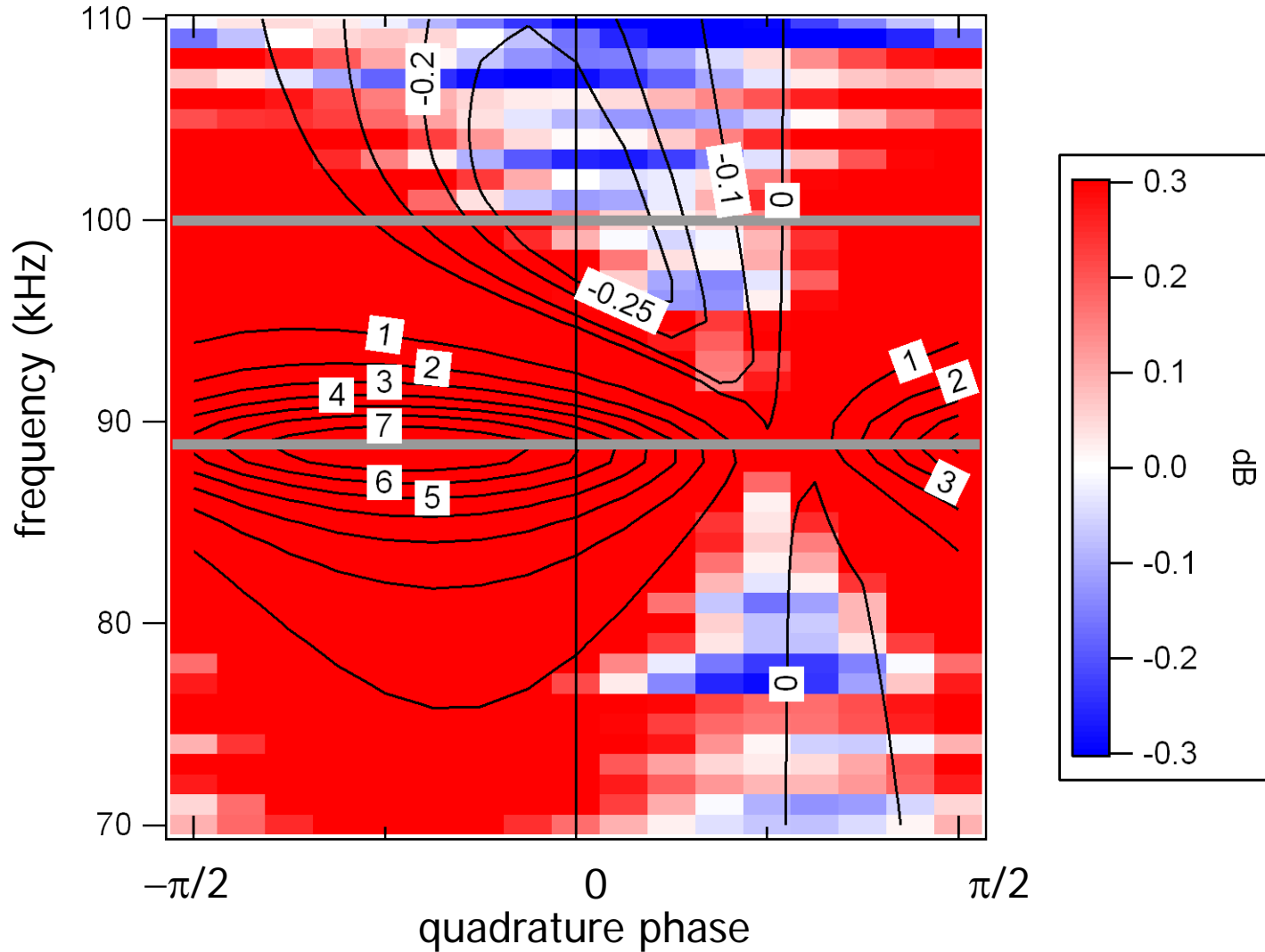
- Theory: semiclassical Langevin equations, one free parameter
Fabre et al., PRA **49**, 1337 (1994); Mancini and Tombesi, *ibid.*, 4055.

Ponderomotive squeezing?



Ponderomotive squeezing? (not clear yet)

- NOT adjusted for detector efficiency (max reduction to 0.93 = - 0.3 dB)

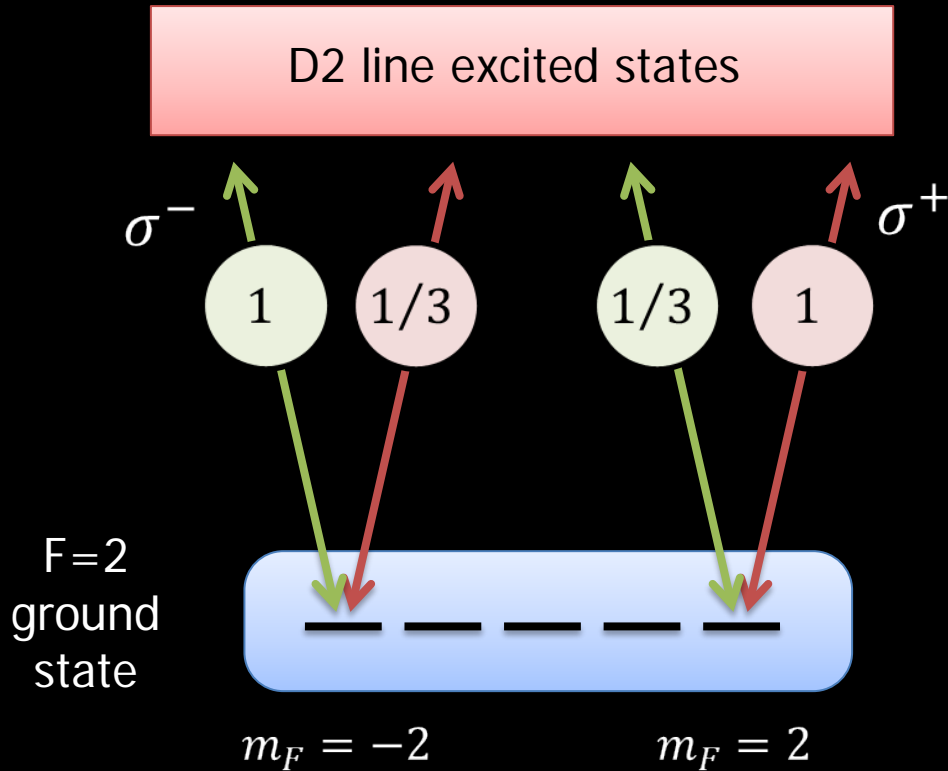


settings:
 $\Delta_{ca} = -2$ GHz
3700 atoms
 $\langle n \rangle = 2.5$



Spins

Cavity optical detection of spin ensembles

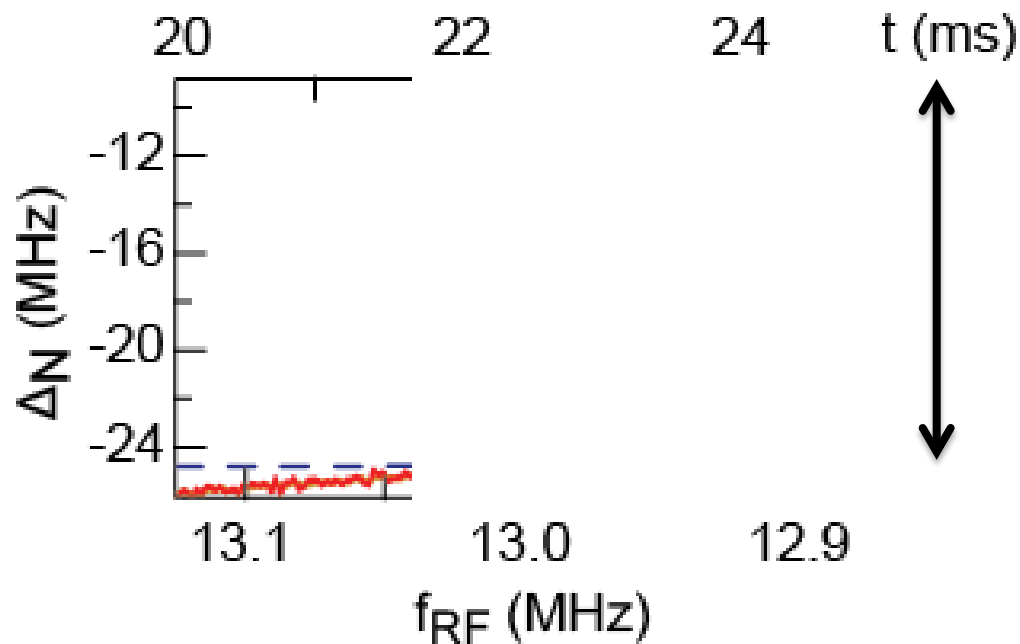
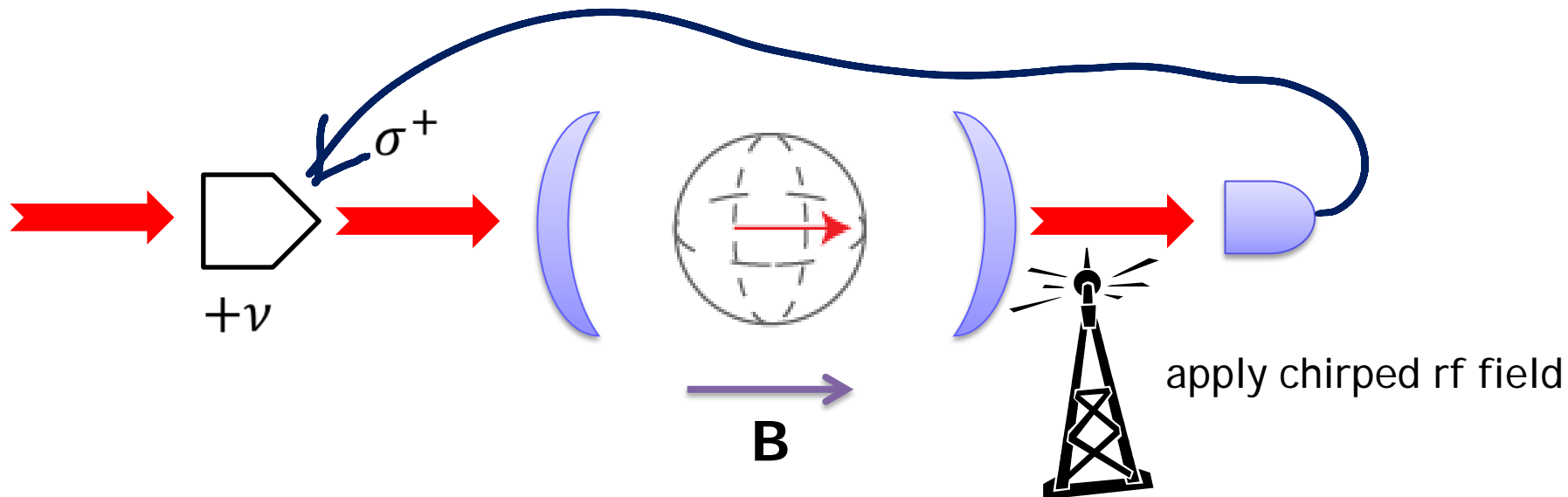


Circular birefringence:

Cavity resonance frequency depends linearly on projection of collective spin (measurement)

Polarized cavity photons produce effective magnetic field along cavity axis (back action)

Cavity optical detection of spin ensembles

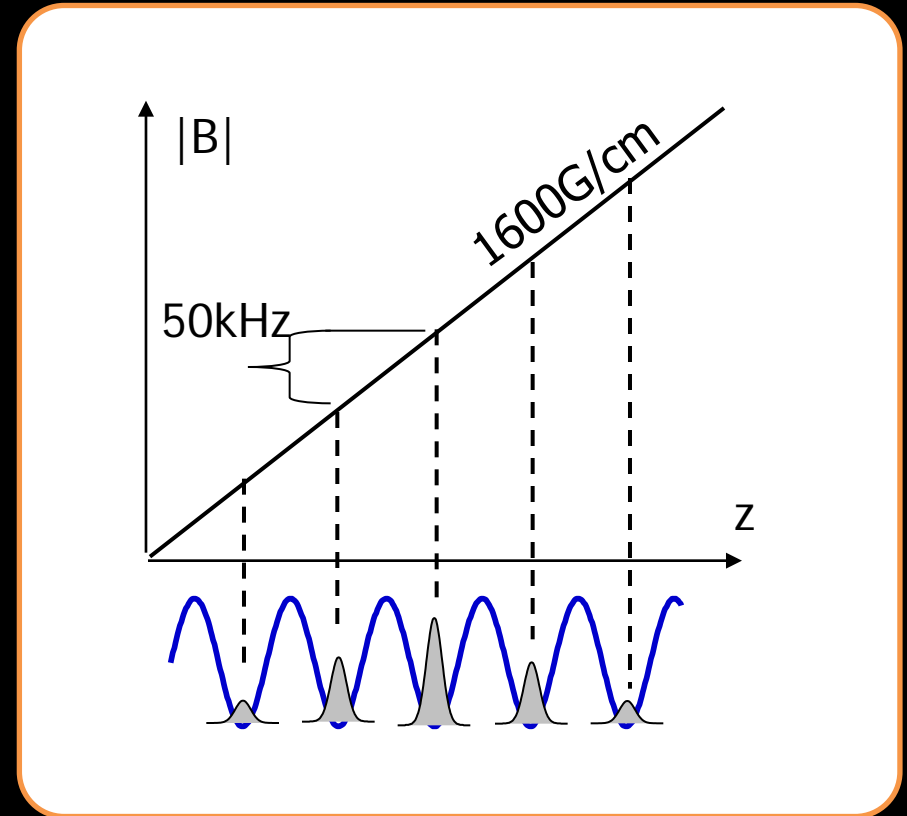


Freq. shift varies factor of 3
when spins are inverted.

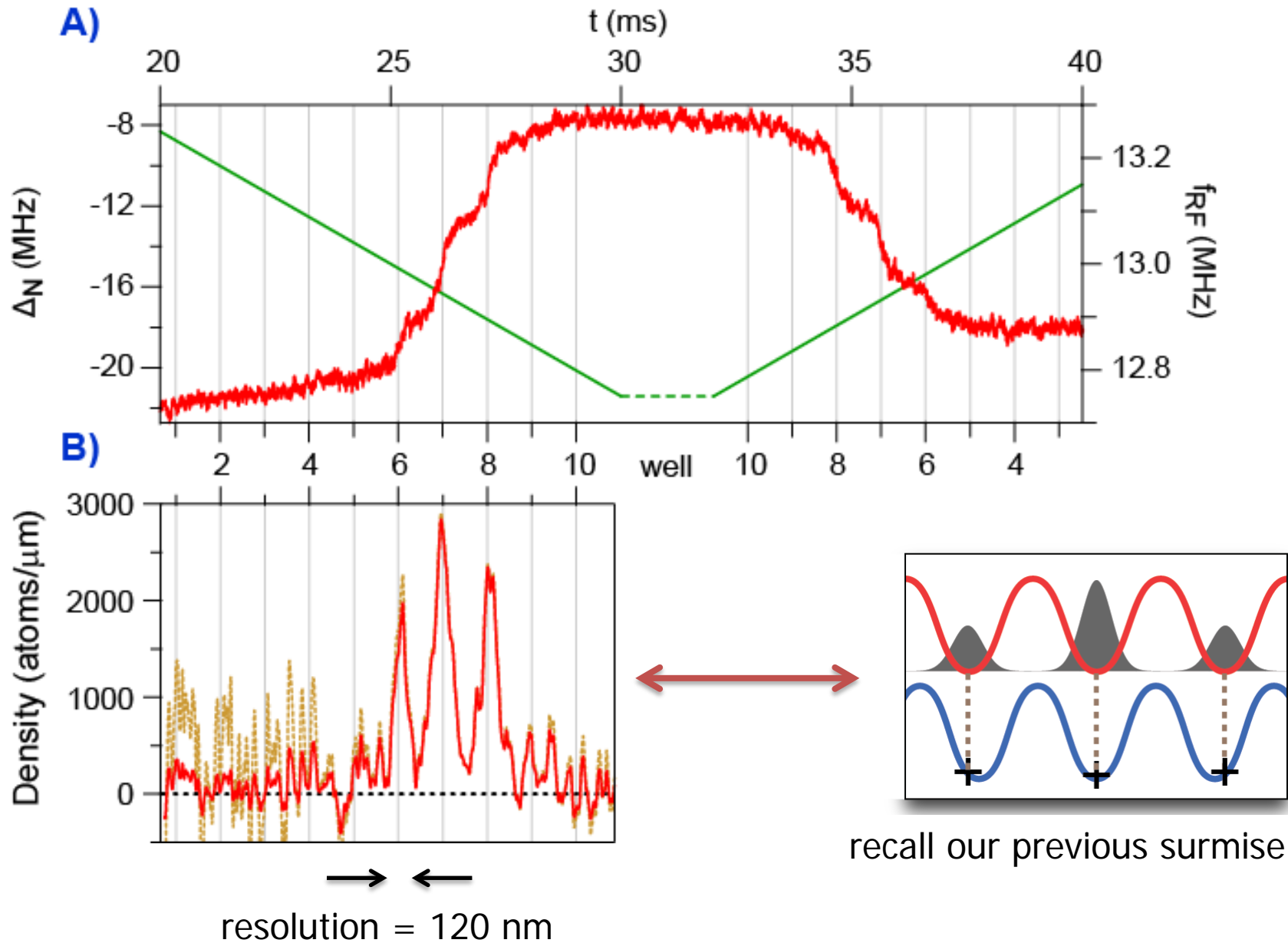
Step height measures
magnetization

Magnetic resonance imaging of atoms in a 1D lattice

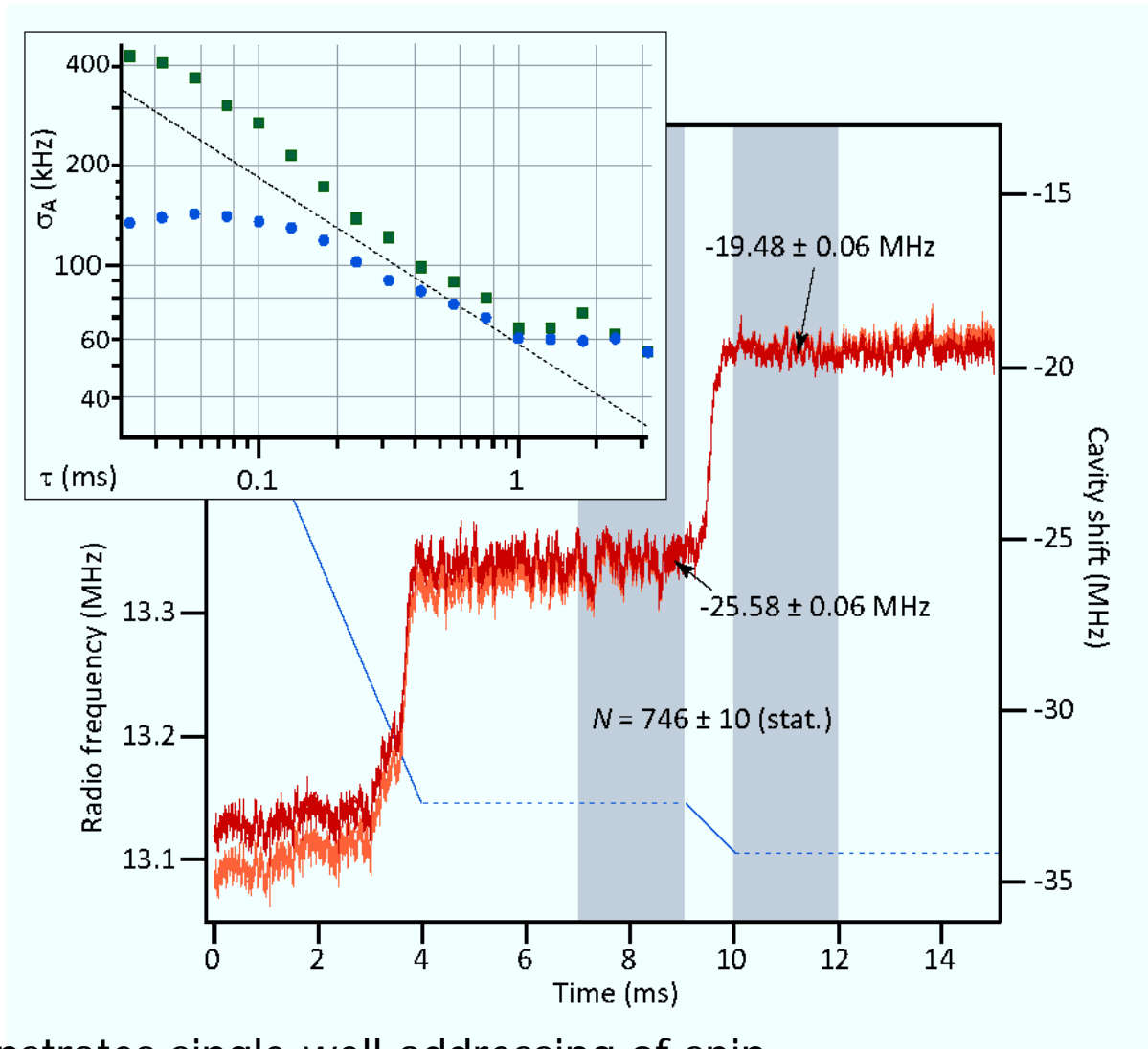
- Apply strong magnetic field to distinguish RF resonance at each lattice site
- Sweep frequency of applied RF drive
- Monitor cavity resonance frequency (using side lock)



Single shot MRI of atoms in a 425-nm-spacing optical lattice



Dynamic range and precision for a single well



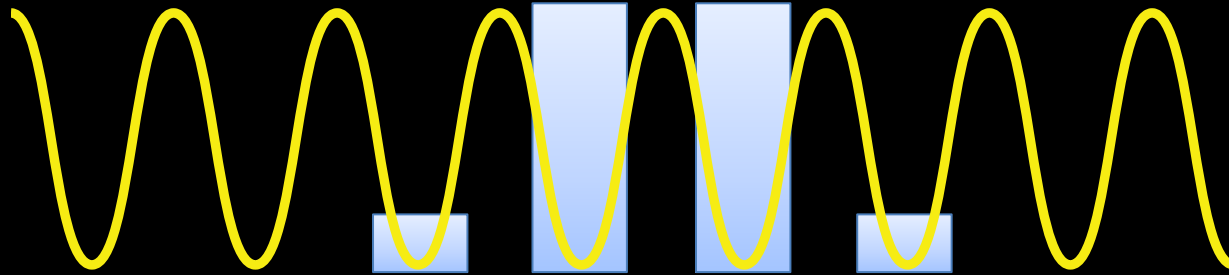
Statistical
uncertainty
below the
Poisson limit

Also demonstrates single-well addressing of spin

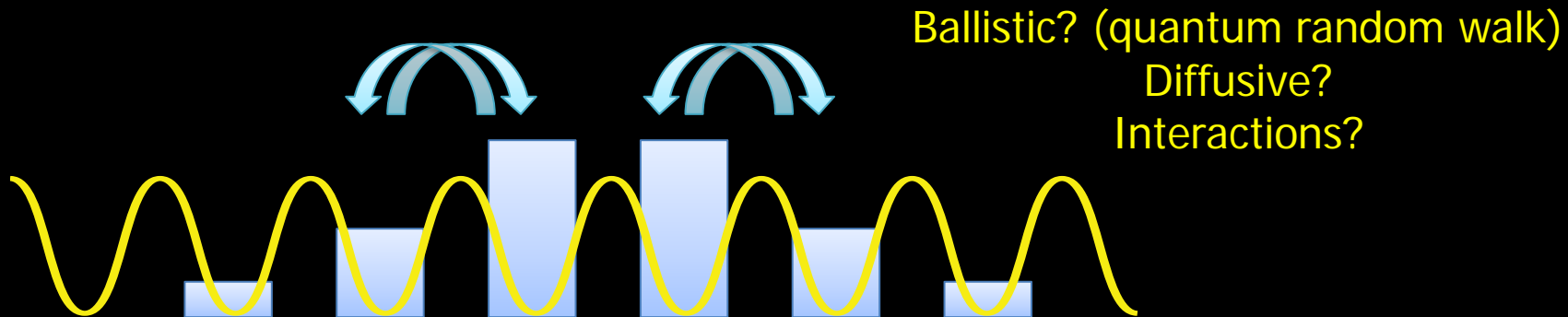
“Optical cavity-aided magnetic resonance imaging of atoms in an optical lattice,” submitted

app: microscopy of quantum transport in an optical lattice

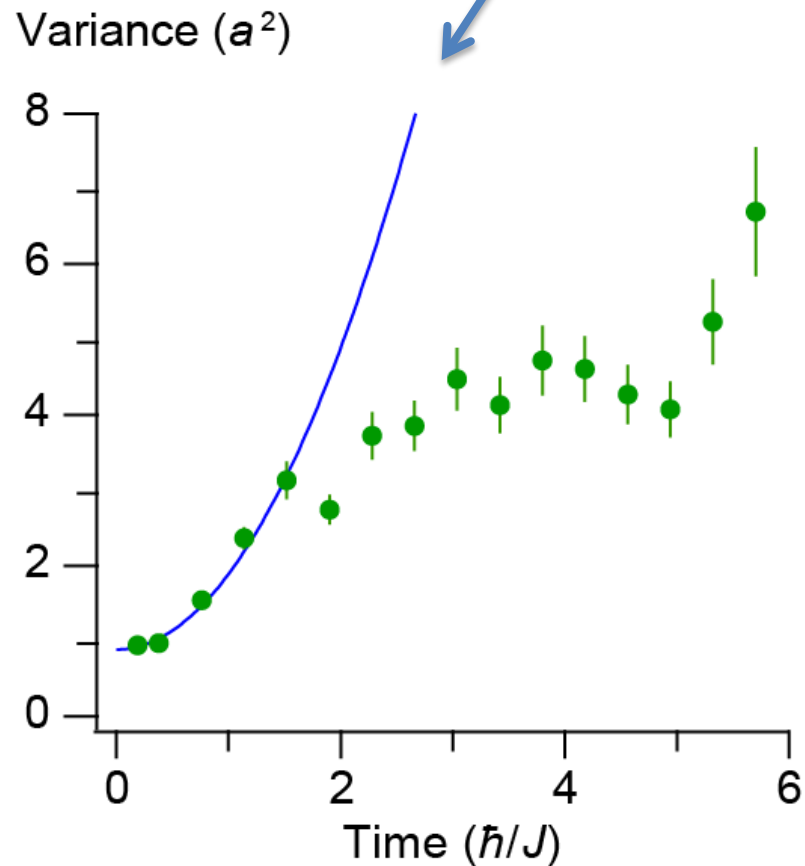
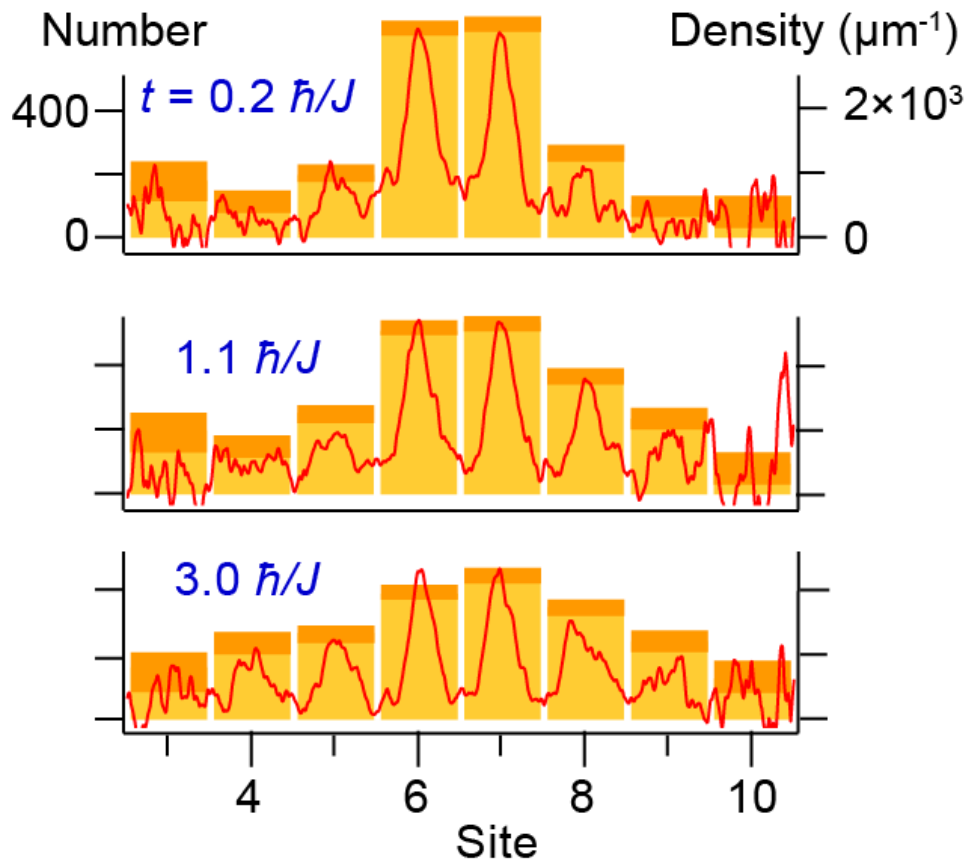
1. Deposit atoms localized in a few wells of the optical lattice



2. Lower lattice depth + tune to resonance to allow tunneling



3. Raise lattice depth + take MRI

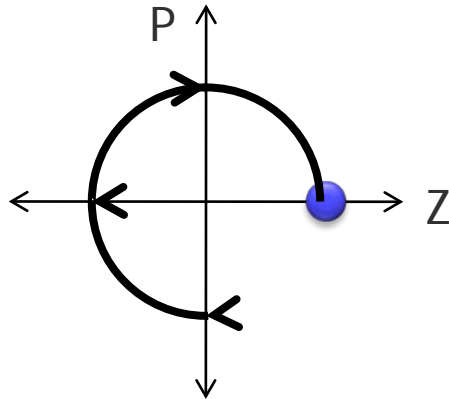


Theory help from Hazzard and Rey (JILA)

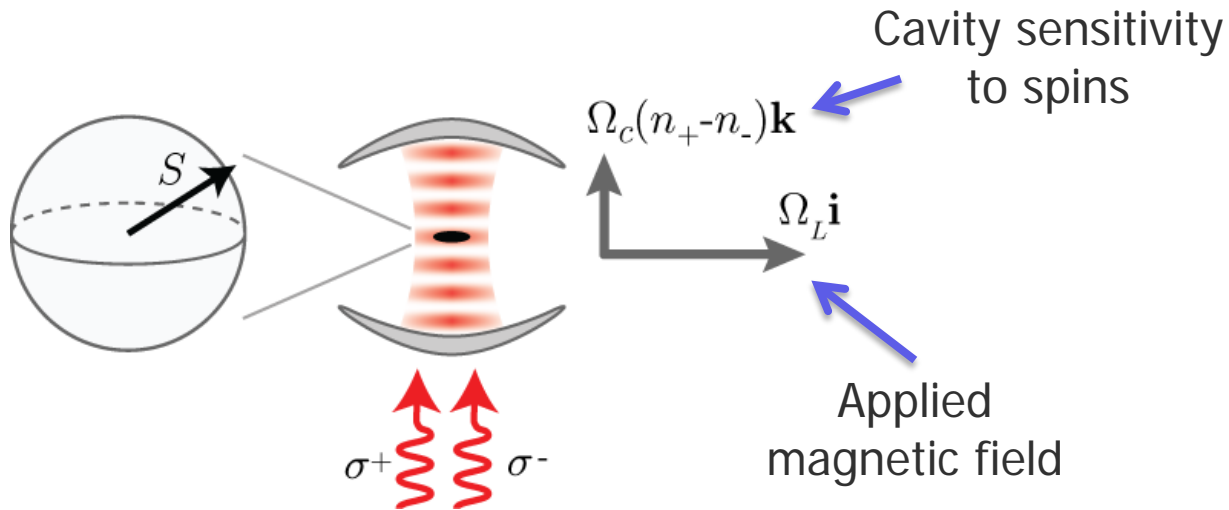
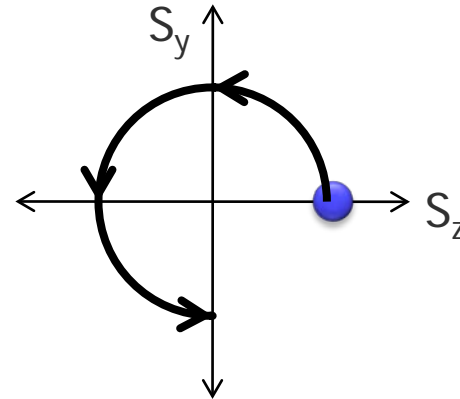
Cavity spin optodynamics

spins = "cantilever"

cantilever:



Larmor precession:



Cavity spin optodynamics

spins = "cantilever"

cavity optomechanics

optomechanical bistability



optomechanical frequency shift



cavity-induced cooling / amplification



ponderomotive squeezing



cavity spin optodynamics

cavity-spin bistability

Larmor precession shifts

coherent amplification and damping
of spin

spin optodynamical squeezing



Thierry Botter

Nathaniel Brahms

Daniel Brooks

Tom Purdy

Zhao-Yuan Ma

Supported generously by:
US National Science Foundation
US Air Force Office of Scientific Research

