

COMPOSITE HIGGS, COMPOSITE MEDIATOR

Flip Tanedo

UC Riverside Particle Theory

with Javi Serra and Alex Wijancgo



30 April 2018

KITP CONFERENCE ON

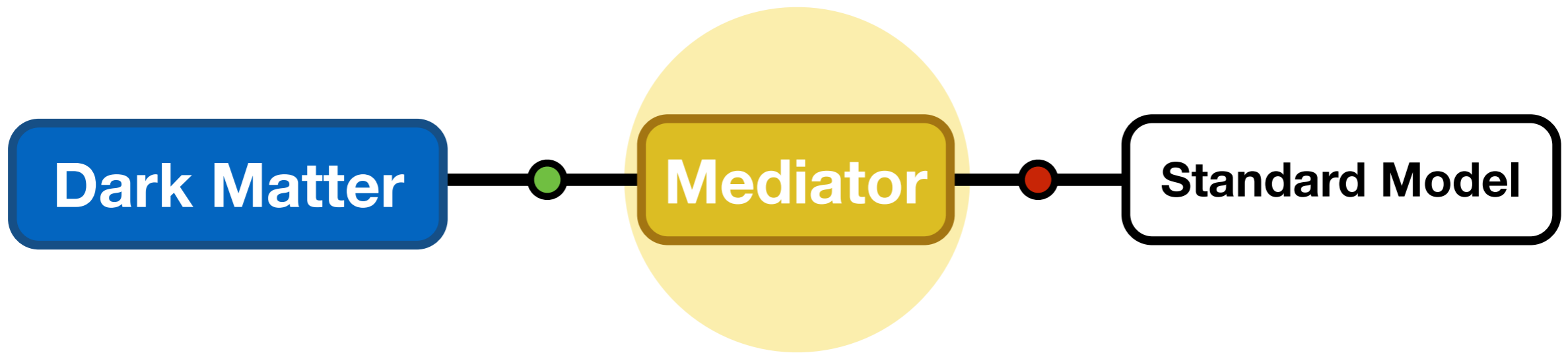
Dark matter detection and detectability:

paradigm confirmation or shift?



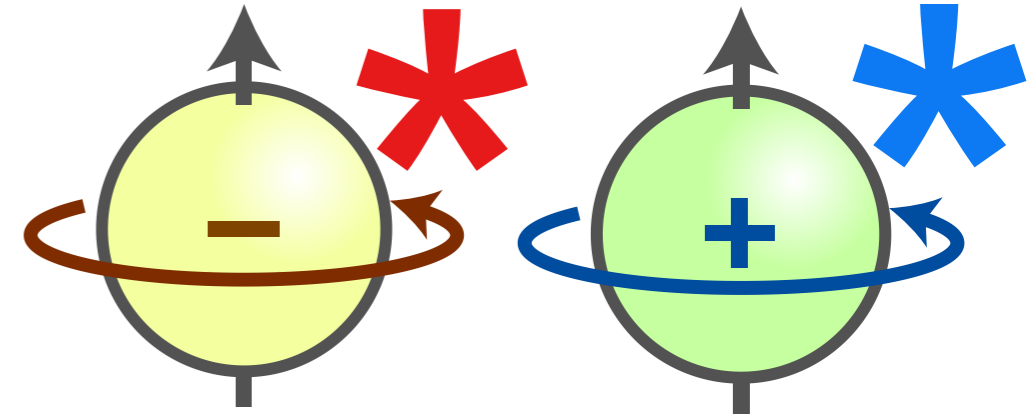
Julio Navarro

Dark Sectors with Light Mediators

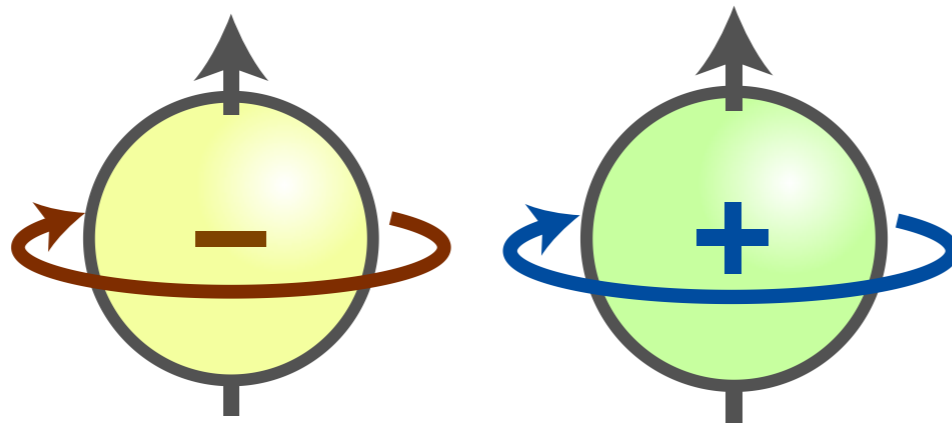


Gauge Invariance

In Standard Model, these are totally different fermions



$$\varphi \bar{f} i \gamma^5 f = i \varphi f_L f_R - i \varphi \bar{f}_L \bar{f}_R$$



See, e.g.

Bell et al. 1503.07874

Kahlhoefer et al. 1510.02110

Bell et al. 1612.03475

Ko et al. 1701.04131

+ many others

How to UV complete?

1. ≥ 2 Higgs Doublet Model

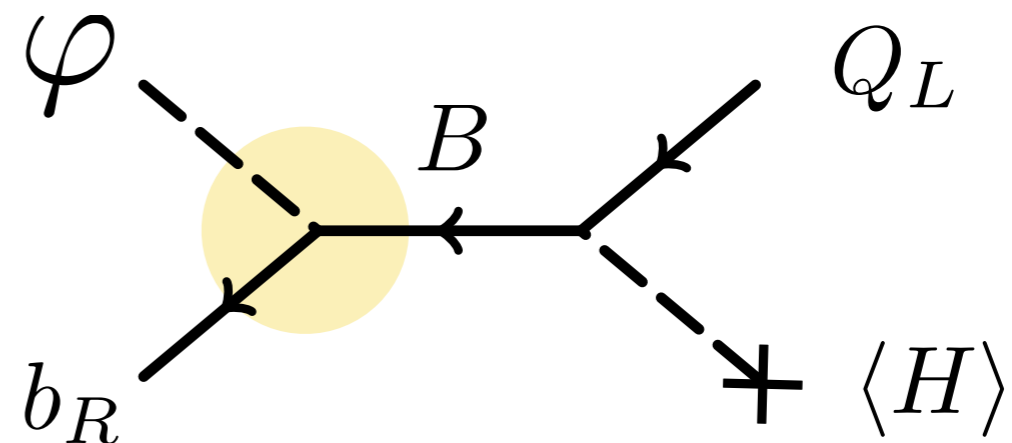
for example, the next-to-minimal SUSY SM

e.g. Ipek et al. 1404.3716, Berlin et al. 1502.06000

2. Heavy **Vector-like Fermions**

... then decouple them

Fan et al. 1507.06993



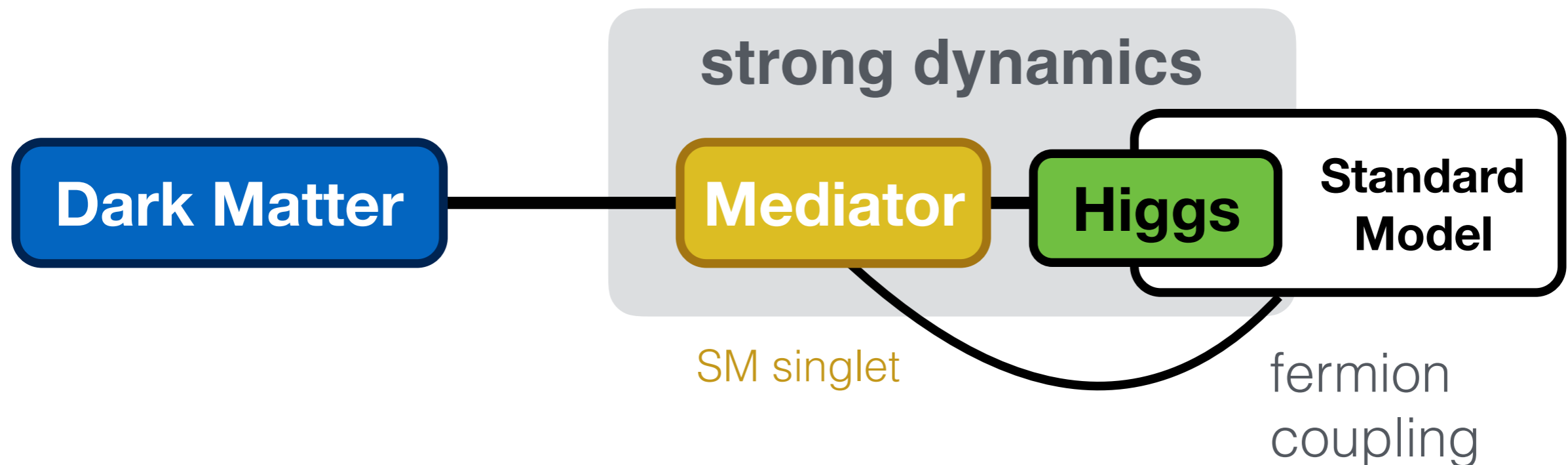
How to UV complete? This talk.

3. Non-minimal **composite Higgs**

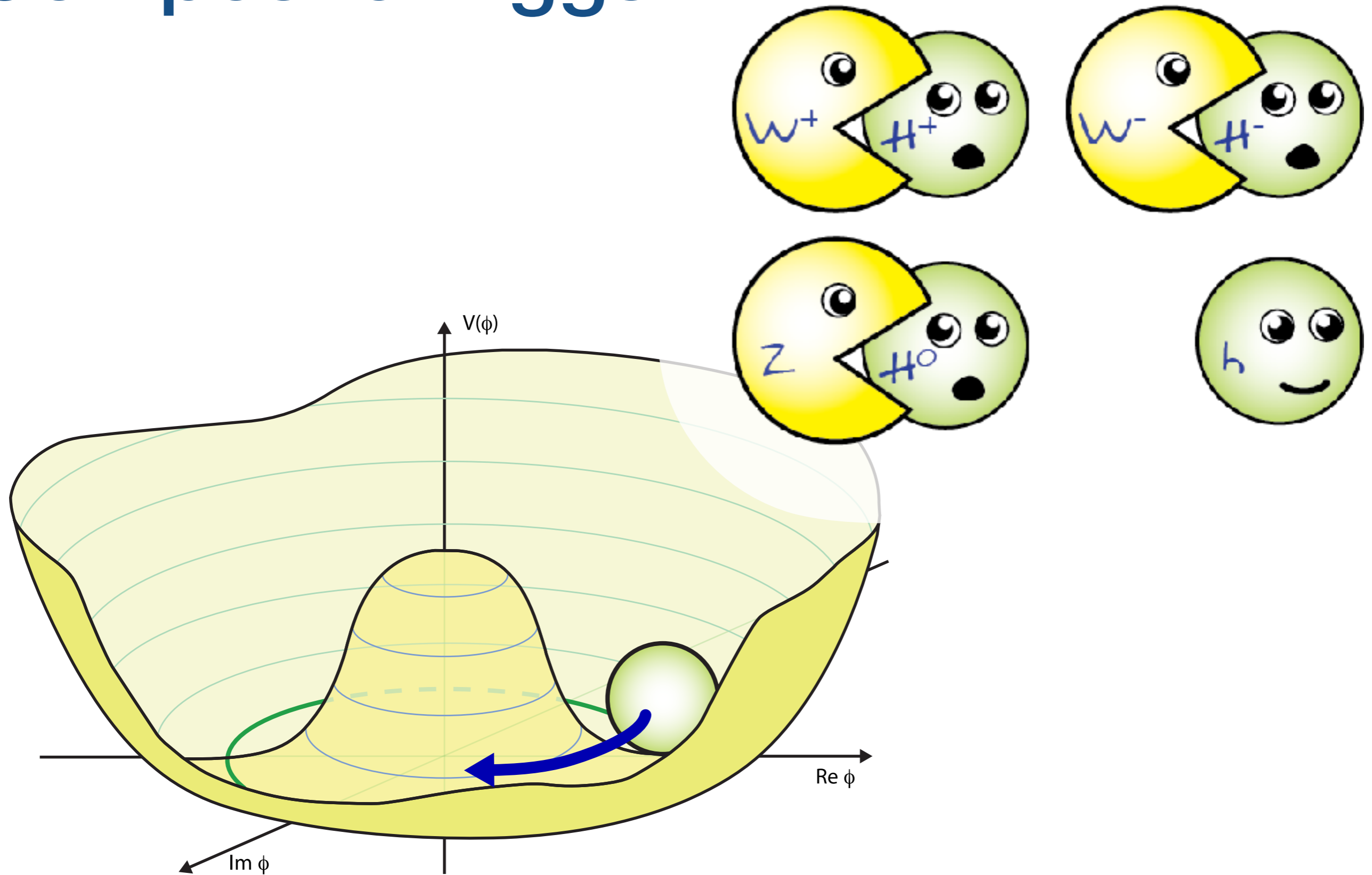
Gauge-singlet pseudoscalars

... predictive, but heavy states decouple

... connects to Hierarchy problem



Composite Higgs



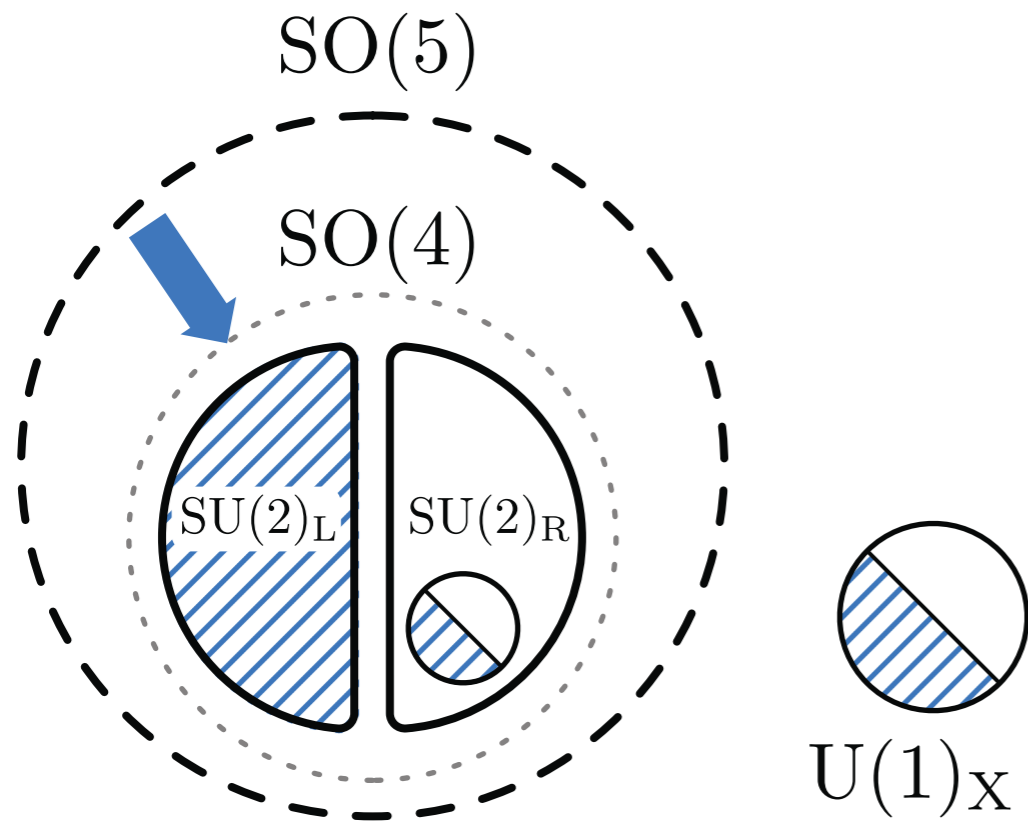
Kaplan & Georgi '84 (Phys. Lett B136 & Phys. Lett. B145)

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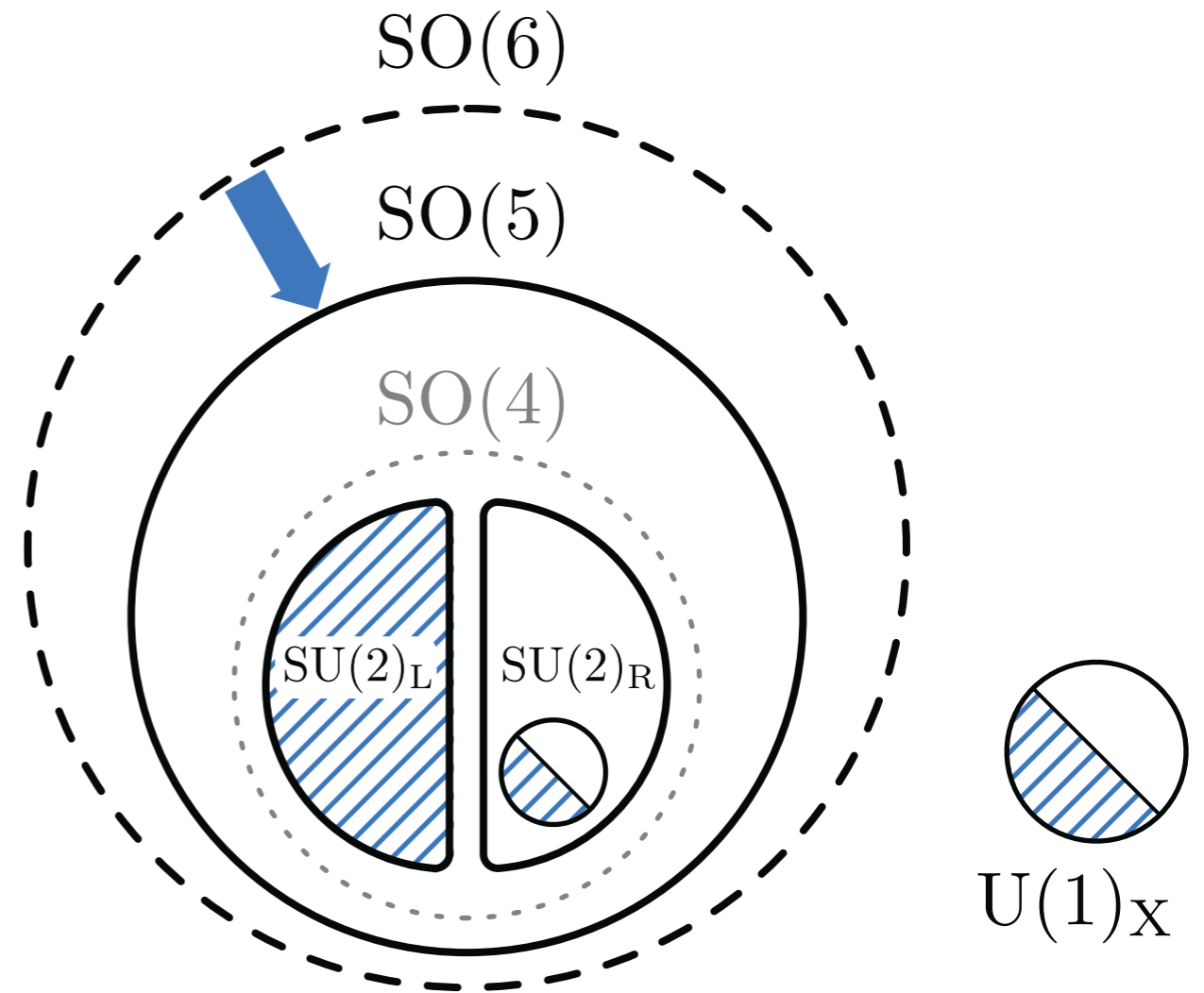
KITP: DM DETECTION & DETECTABILITY

Composite Higgs + Singlet

Gauged symmetry in blue



Minimal Composite Higgs

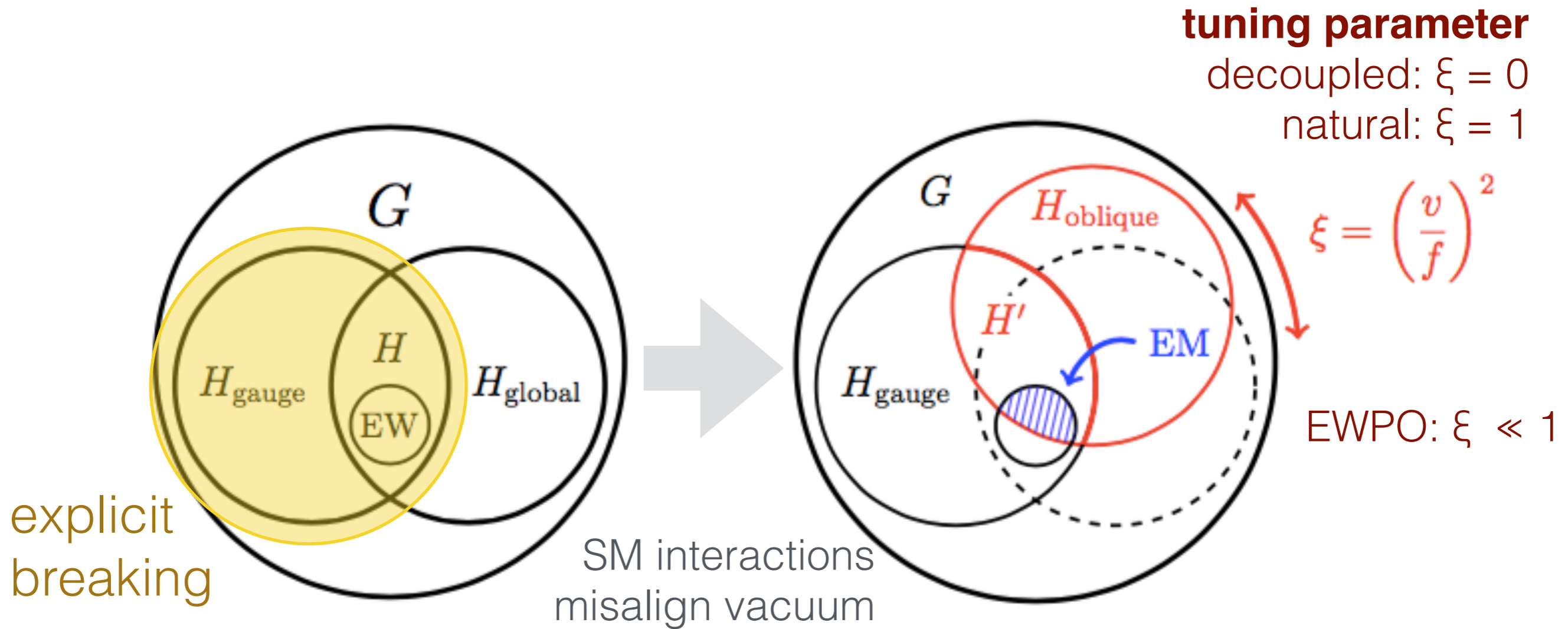


Next-to-Minimal Composite Higgs

Singlet as DM: Frigerio et al. 1204.2808, Marzocca et al. 1404.7419, Fonseca et al. 1501.05957, Carmona et al. 1504.00332, Antipin et al. 1503.08749, **Related:** Poland & Thaler 0808.1290; Asano and Kitano 1406.6374

Kaplan & Georgi '84 (Phys. Lett B136 & Phys. Lett. B145)

Explicit & Electroweak Breaking



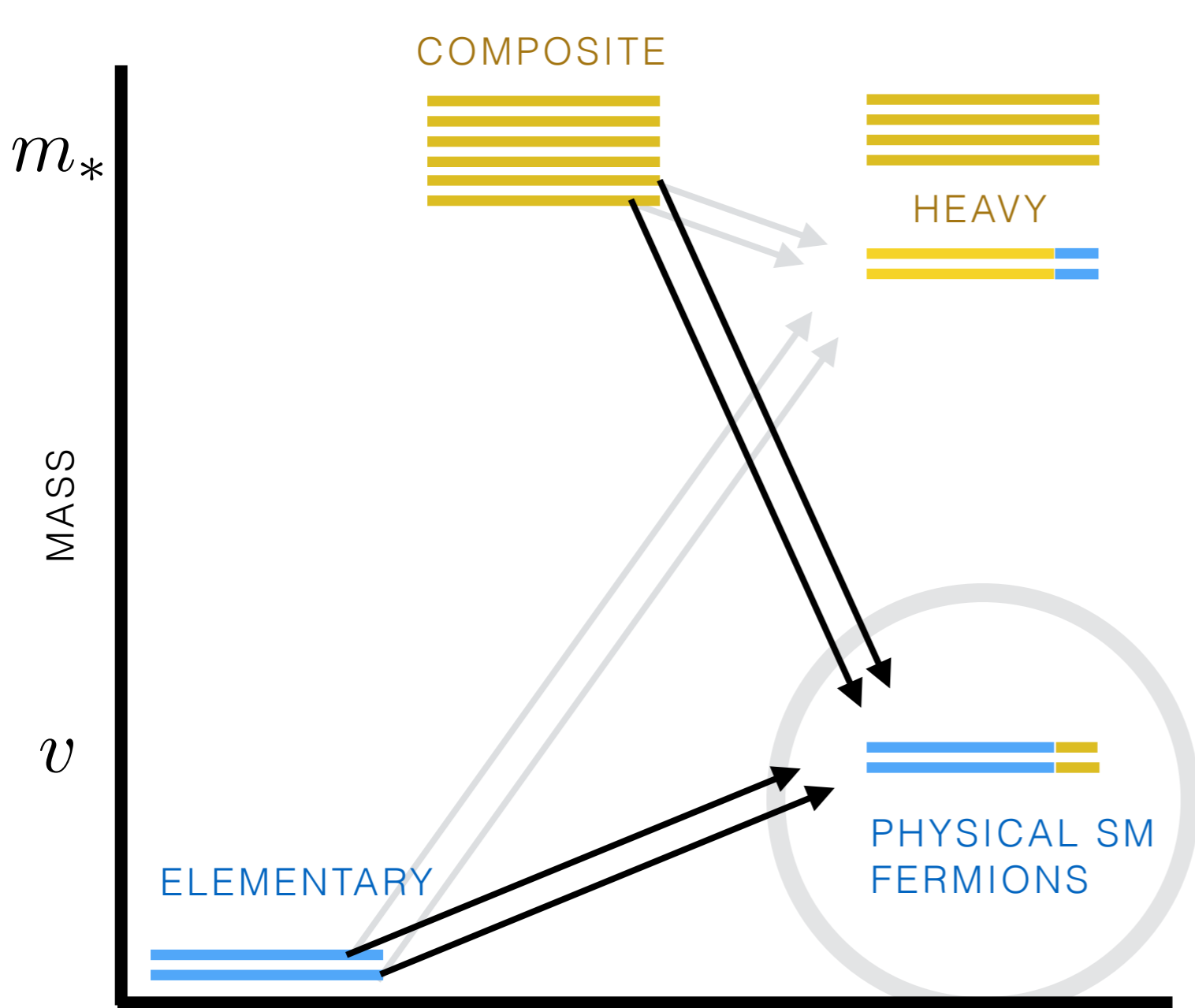
Loops of gauge bosons, fermions generate electroweak-breaking Higgs potential

Partial Compositeness

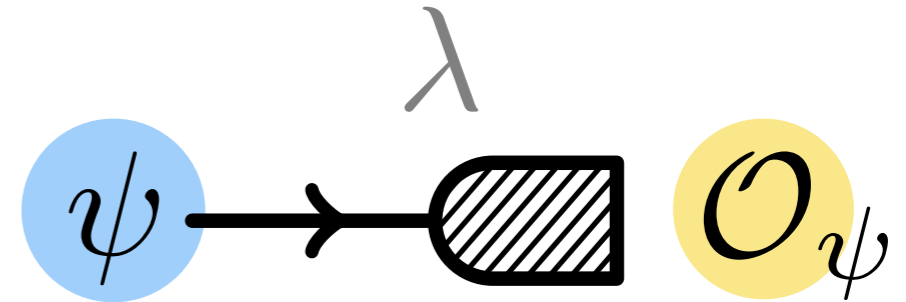
$SU(2)_L \ SU(2)_R$

$$6 = 4 + 1 + 1 = (\mathbf{2}, \mathbf{2}) + (\mathbf{1}, \mathbf{1}) + (\mathbf{1}, \mathbf{1})$$

DOUBLET



$$\mathcal{L} \supset \lambda \bar{\psi} \mathcal{O}_\psi + \text{h.c.}$$



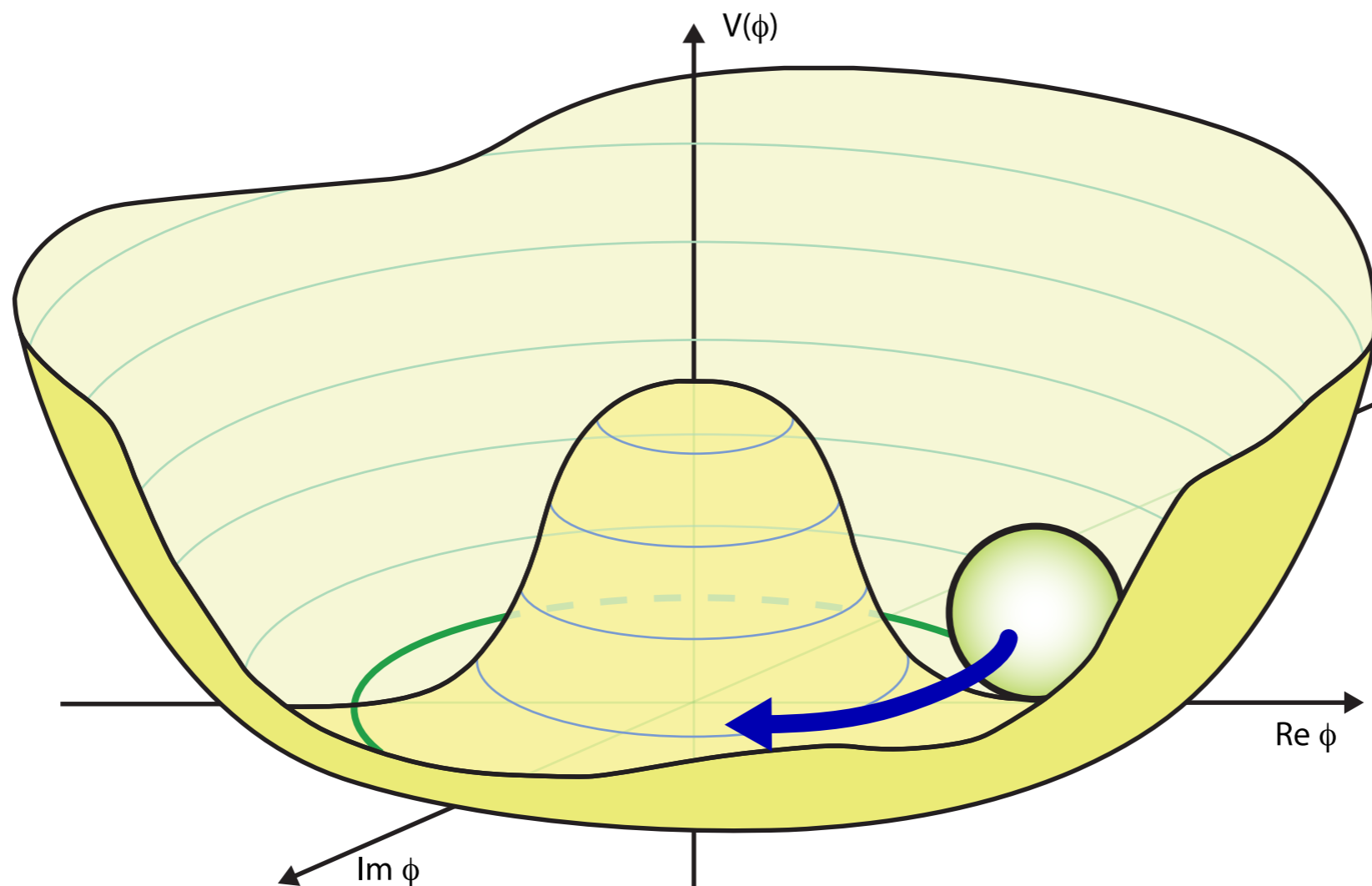
t_L & b_L
only

6 fermionic
composites

Elementary fields are not $SO(6)$ multiplets, mix with composite operators.

Non-linear realization

$$\Sigma = e^{iT^A} \pi^A(x) / f \Sigma_0$$



An $SO(6)$ basis

VEV

$$\Sigma_0 = (0, 0, 0, 0, 0, 1)^T$$

$$\Sigma = \frac{1}{f} (0 \quad 0 \quad 0 \quad h \quad \eta \quad \sqrt{f^2 - h^2 - \eta^2})$$

BROKEN

$$T_2^\alpha = \frac{i}{\sqrt{2}} \begin{pmatrix} & & & & -\delta_1^\alpha \\ & & & & -\delta_2^\alpha \\ & & & & -\delta_3^\alpha \\ & & & & -\delta_4^\alpha \\ \hline & & & & 0 \\ \delta_1^\alpha & \delta_2^\alpha & \delta_3^\alpha & \delta_4^\alpha & 0 \end{pmatrix}$$

$$T_\eta = \frac{i}{\sqrt{2}} \begin{pmatrix} 0 & & & & \\ & 0 & & & \\ & & 0 & & \\ & & & 0 & \\ \hline & & & & 1 \\ & & & & -1 \end{pmatrix}$$

MEDIATOR
U(1) SYM

Reminder: $SO(6) \supset SO(4) \simeq SU(2)_L \times SU(2)_R$ (custodial symmetry)

What the theory looks like

Use Σ and spurions to write invariants

dark matter embedding

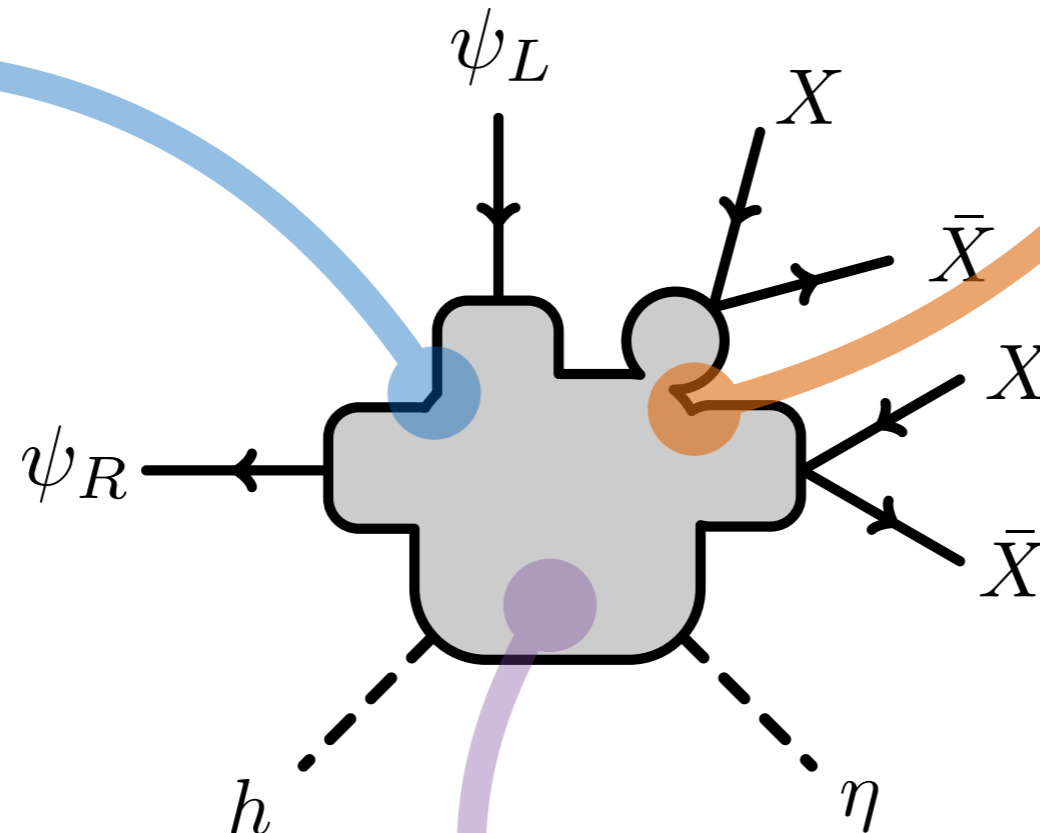
CHOICES LEAD TO MEDIATOR POTENTIAL



OUTPUTS:
MEDIATOR MASS AS A
FUNCTION OF DISCRETE
CHOICES AND DM MASS

visible matter embedding

CHOICES LEAD TO
- HIGGS POTENTIAL
- MEDIATOR POTENTIAL



COMPOSITE HIGGS

SPECIFY BREAKING SCALE f
GIVES MEASURE OF TUNING

$$\xi = \left(\frac{v}{f}\right)^2$$

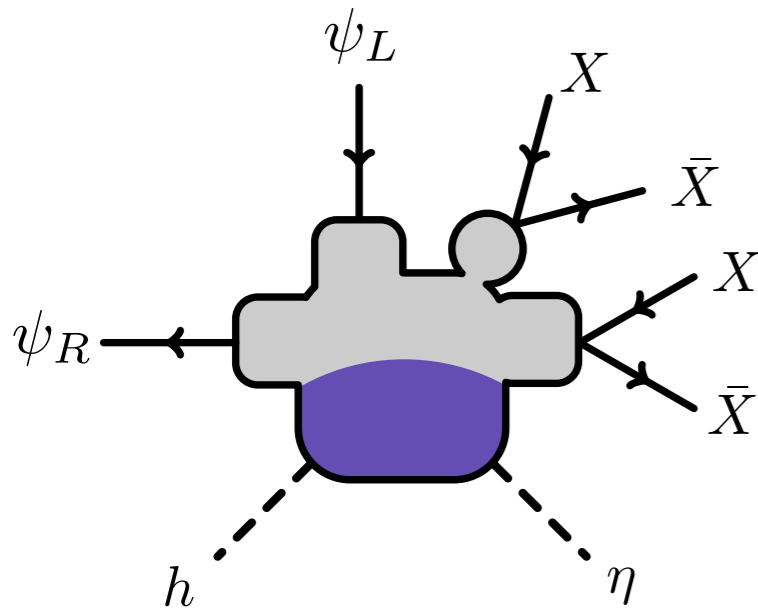
NL Σ IRREDUCIBLE INTERACTIONS

Non-linear Σ Model

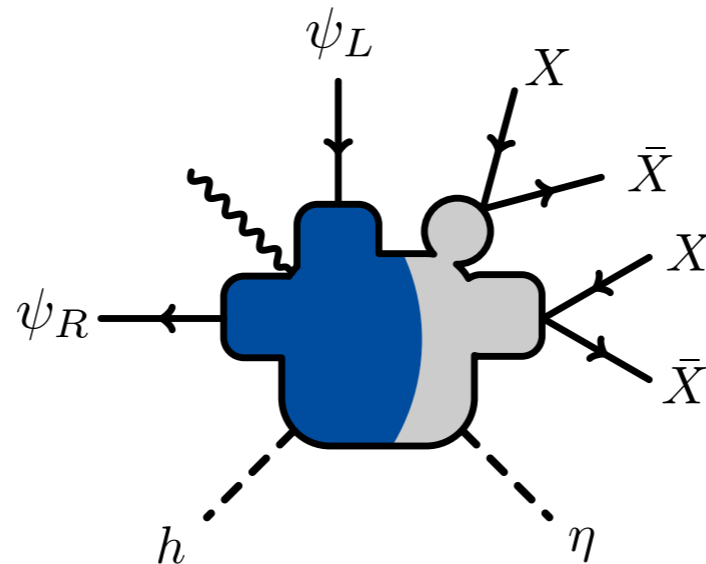
$$\frac{f^2}{2} |D_\mu \Sigma|^2 = \frac{(\partial_\mu h)^2}{2} + \frac{(\partial_\mu \eta)^2}{2}$$

$$+ \frac{1}{2} \frac{(h\partial_\mu h + \eta\partial_\mu \eta)^2}{f^2 - h^2 - \eta^2}$$

$$+ \frac{g^2}{4} h^2 \left(W_\mu^+ W^{\mu-} + \frac{1}{2 \cos^2 \theta_W} Z_\mu Z^\mu \right)$$



NONLINEAR



HIGGS-GAUGE INTERACTIONS

HIGGS POTENTIAL GIVES:

$$h \rightarrow v + h\sqrt{1 - \xi}$$

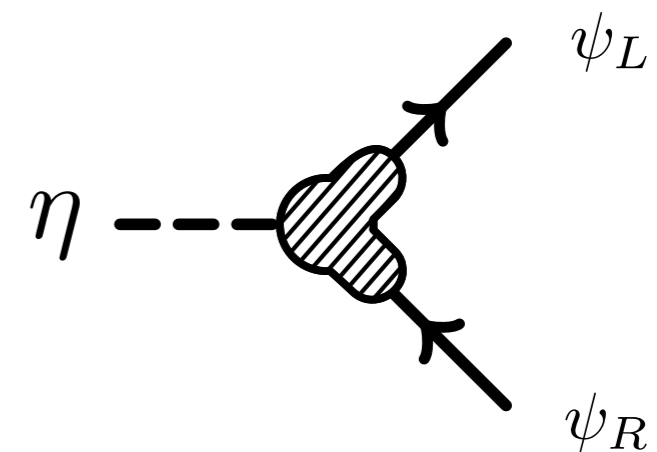
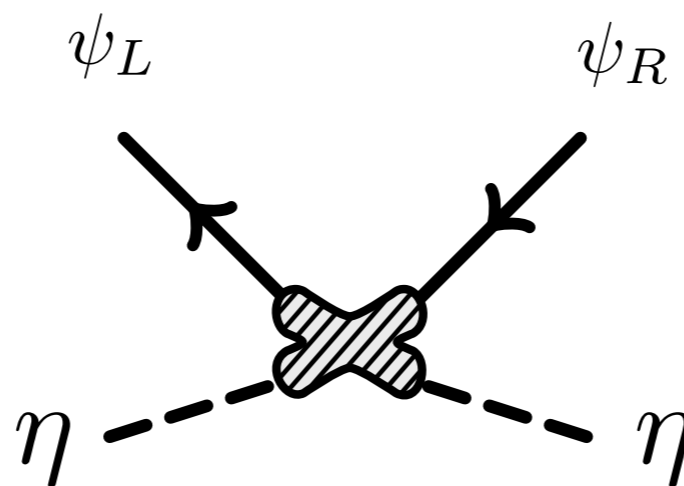
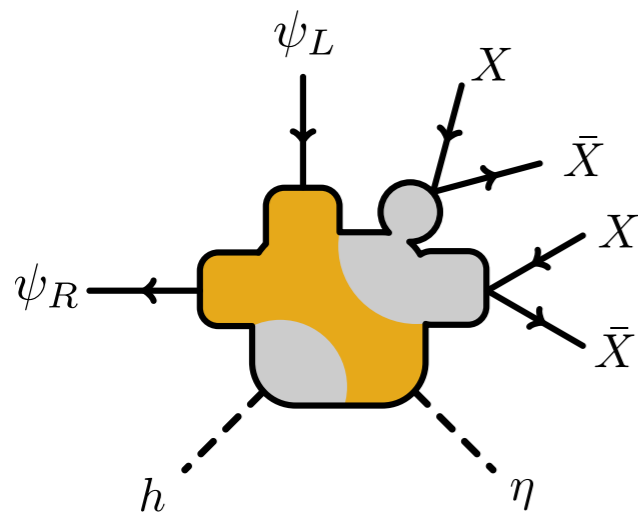
η -SM coupling: right-handed embed

$$(2, 2)_{\frac{2}{3}}: Q = \frac{1}{\sqrt{2}} (ib_L, b_L, it_L, -t_L, 0, 0)^T \quad \psi \rightarrow \overset{\lambda}{\text{[diagonal box]}} \mathcal{O}_\psi$$

$$(2, 2)_{\frac{2}{3}}: T = (0, 0, 0, 0, i\delta_t t_R, t_R)^T \quad \text{CHOICE: } \delta \text{ PARAMETER}$$

$$\Sigma = \frac{1}{f} (0 \quad 0 \quad 0 \quad h \quad \eta \quad \sqrt{f^2 - h^2 - \eta^2})$$

$$\mathcal{L} \supset y_t f (\bar{Q} \Sigma) (\Sigma^T T) = -\frac{y_t}{\sqrt{2}} h \bar{t}_L t_R \left(\sqrt{1 - \frac{h^2}{f^2} - \frac{\eta^2}{f^2}} + i\delta_t \frac{\eta}{f} \right)$$

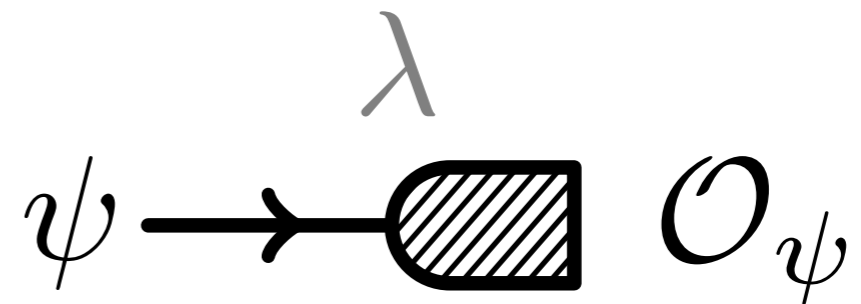


Dark matter: same thing again?

Introduce dark matter as elementary Dirac fermion

Why not use this again?

Why should such a particle be stable?



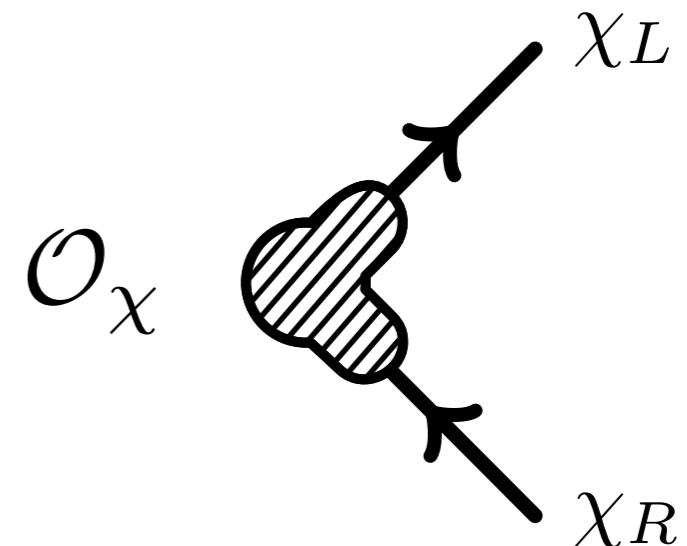
Better option: **bilinear** mixing

TECHNICOLOR-LIKE MASS GENERATION

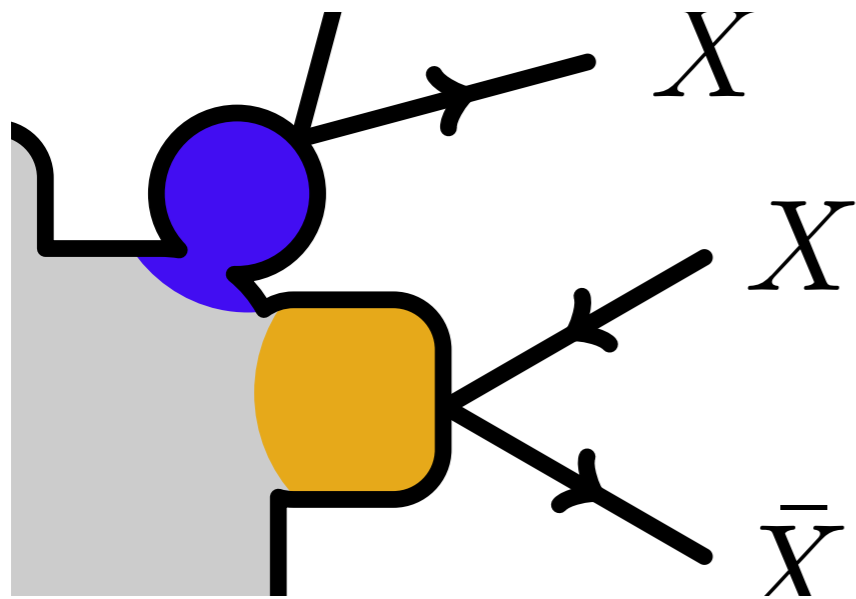
Doesn't work for visible matter (flavor)

But preserves $\chi \rightarrow -\chi$

Sets the mass scale to f



Dark Matter: bilinear coupling



$$\mathcal{L} \supset -y'_\chi f \bar{\chi}_L \chi_R$$

SINGLET

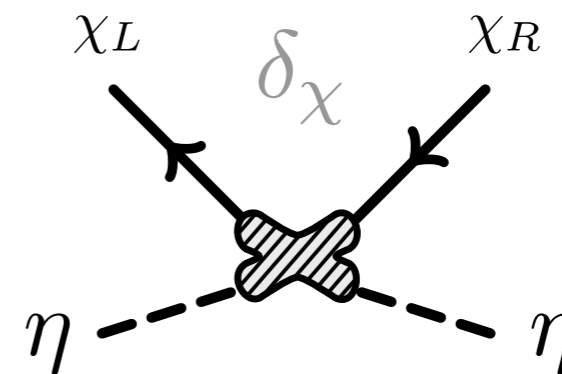
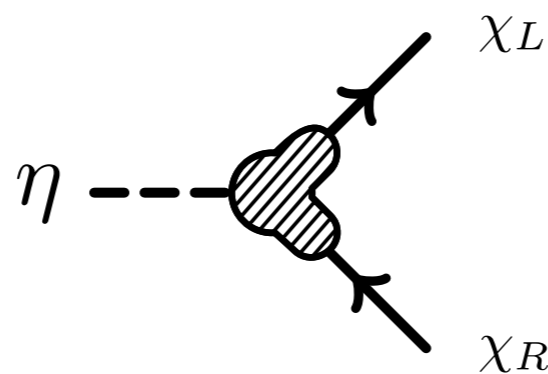
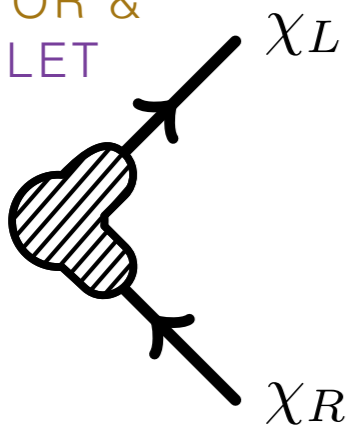
$$\mathcal{L} \supset -y_\chi f (\Sigma^T \mathbb{X})$$

VECTOR

$$\mathbb{X} = (0, 0, 0, 0, i \bar{\chi}_L \chi_R, \delta_\chi \bar{\chi}_L \chi_R)^T$$

$$\Sigma = \frac{1}{f} (0 \quad 0 \quad 0 \quad h \quad \eta \quad \sqrt{f^2 - h^2 - \eta^2})$$

VECTOR &
SINGLET



+ HIGGS COUPLINGS

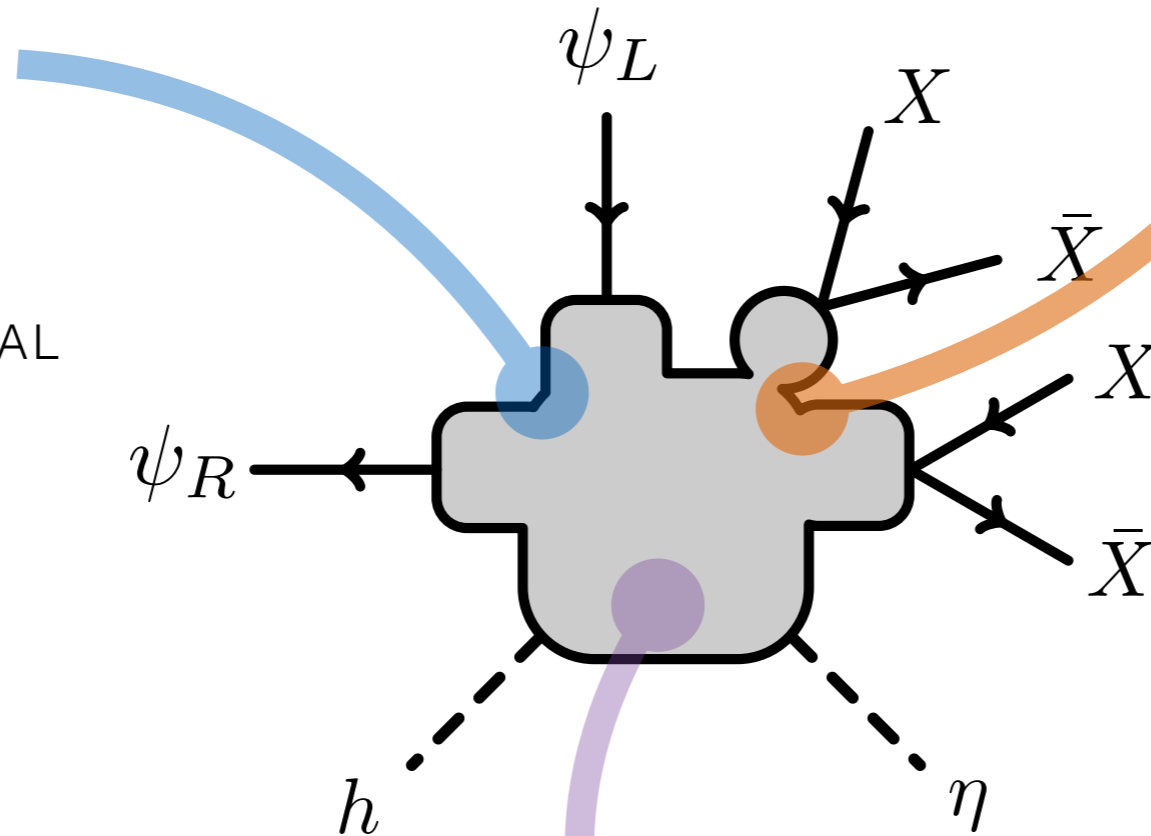
What does this buy us?

visible matter embedding

CHOICES LEAD TO
- HIGGS POTENTIAL
- MEDIATOR POTENTIAL

COMPOSITE HIGGS
SPECIFY BREAKING SCALE f
GIVES MEASURE OF TUNING

$$\xi = \left(\frac{v}{f}\right)^2$$



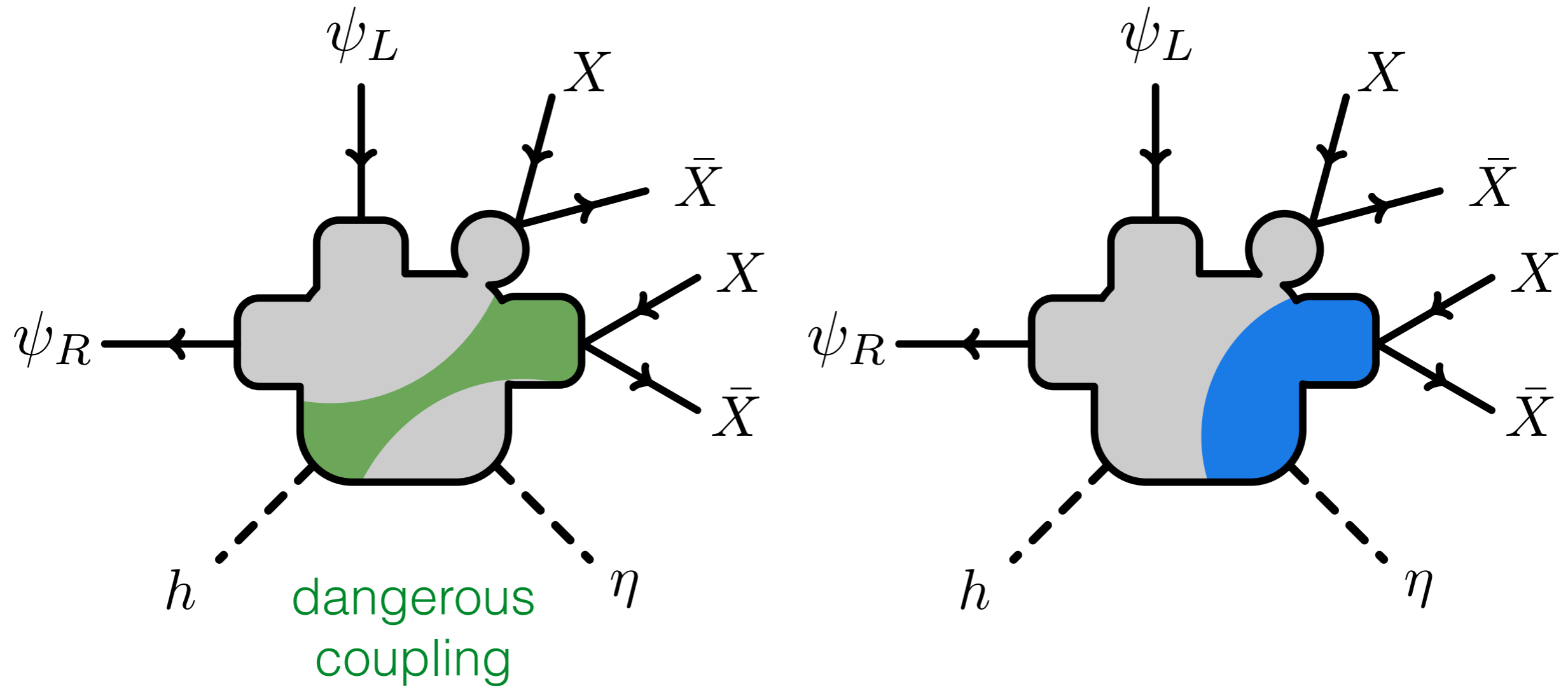
dark matter embedding

CHOICES LEAD TO
MEDIATOR POTENTIAL

OUTPUTS:
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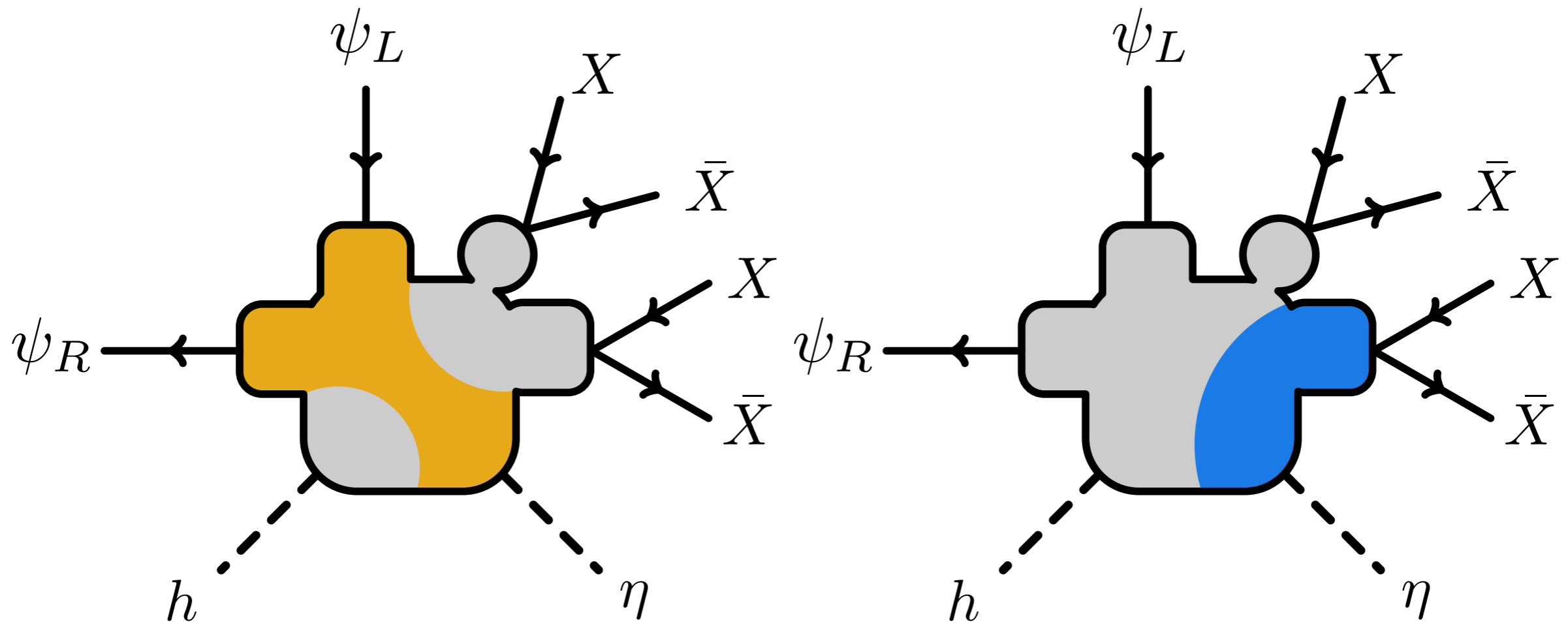
NLΣM IRREDUCIBLE INTERACTIONS

Dark Sector Couplings

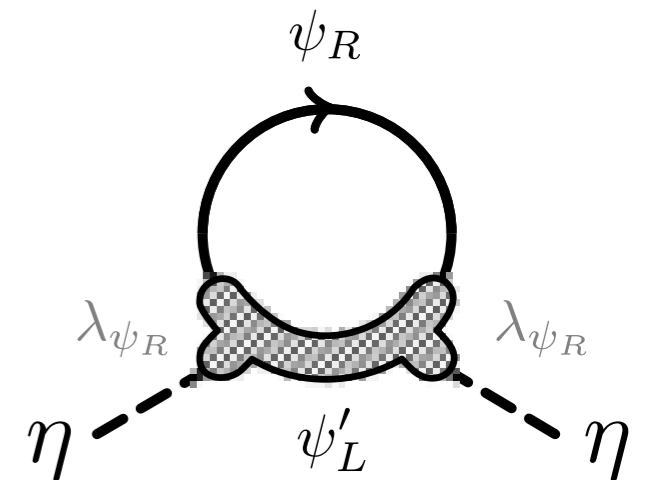


These are determined once you specify the order parameters of explicit breaking.

Mediator Coupling



Fermion loops give **mediator mass**
Interaction strength \iff **mediator mass**



So finding a thermal relic is ... kind of a **miracle**.

Mediator Mass

CHOOSE ONE NON-TRIVIAL REP

Standard Model

Approximate contribution to m_η^2 from fermion loops

$$\begin{aligned}
 t_R (\mathcal{O}_t) & \quad (\delta_t^2 - 1) \frac{3y_t}{8\pi^2} \frac{m_*^3}{f} \approx (1 \text{ TeV})^2 \left(\frac{m_*}{3 \text{ TeV}} \right)^3 \left(\frac{\text{TeV}}{f} \right) (\delta_t^2 - 1) \\
 b_R (\mathcal{O}_b) & \quad (\delta_b^2 - 1) \frac{3y_b}{8\pi^2} \frac{m_*^3}{f} \approx (130 \text{ GeV})^2 \left(\frac{m_*}{3 \text{ TeV}} \right)^3 \left(\frac{\text{TeV}}{f} \right) (\delta_b^2 - 1) \\
 \tau_R (\mathcal{O}_b) & \quad (\delta_\tau^2 - 1) \frac{y_\tau}{8\pi^2} \frac{m_*^3}{f} \approx (60 \text{ GeV})^2 \left(\frac{m_*}{3 \text{ TeV}} \right)^3 \left(\frac{\text{TeV}}{f} \right) (\delta_\tau^2 - 1)
 \end{aligned}$$

Dark Matter

Approximate contribution to m_η^2 from dark matter loops

$$\begin{aligned}
 \text{Vector } (\mathcal{O}_\chi) & \quad (1 - \delta_\chi^2) \frac{y_\chi^2 m_*^2}{8\pi^2} \approx (70 \text{ GeV})^2 \left(\frac{m_*}{3 \text{ TeV}} \right)^2 \left(\frac{y_\chi}{0.2} \right)^2 (1 - \delta_\chi^2) \\
 \text{Singlet } (\mathcal{O}'_\chi) & \quad \frac{m_\chi^{(1)} m_\chi^{(6)} m_*^2}{8\pi^2 f^2} \approx (10 \text{ GeV})^2 \left(\frac{m_*}{3 \text{ TeV}} \right)^2 \left(\frac{1 \text{ TeV}}{f} \right)^2 \left[\frac{m_\chi^{(1)} m_\chi^{(6)}}{(25 \text{ GeV})^2} \right]
 \end{aligned}$$

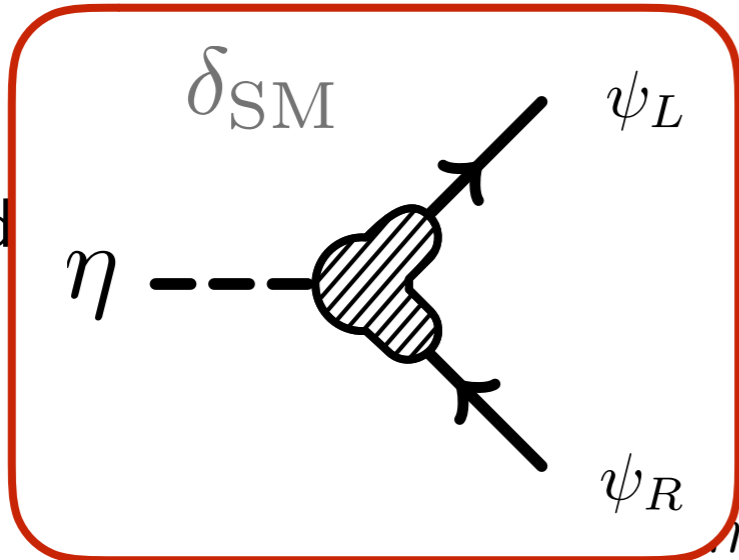
$\delta = 1$ corresponds to unbroken $U(1)_\eta$

Mediator Mass

$$\mathbb{X} = (0, 0, 0, 0, i \bar{\chi}_L \chi_R, \delta_\chi \bar{\chi}_L \chi_R)$$

$$T = (0, 0, 0, 0, i \delta_{t^c} t_R, t_R)^T$$

Standard



NON-TRIVIAL REP

contribution to m_η^2 from fermion loops

t_R (\mathcal{O}_t)

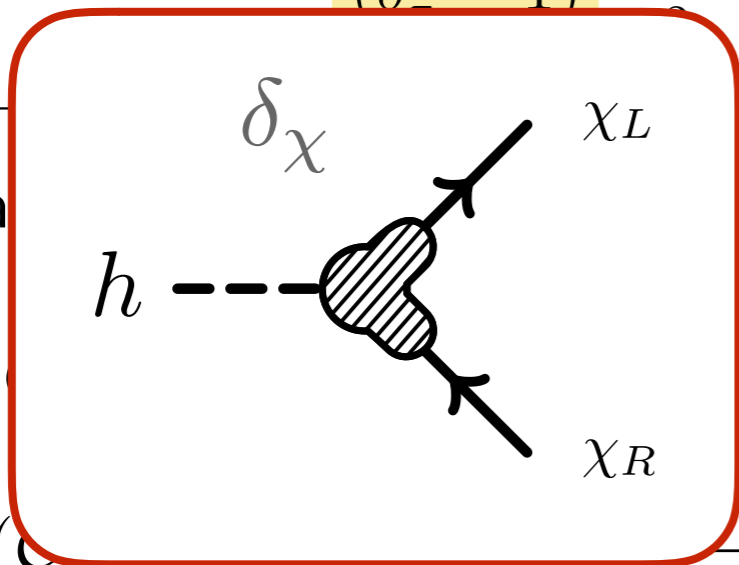
b_R (\mathcal{O}_b)

τ_R (\mathcal{O}_b)

Dark Ma

Vector (\mathcal{O}_χ)

Singlet (\mathcal{O}_χ)



contribution to m_η^2 from dark matter loops

$n_*^3 \approx$

$n_*^2 \approx$

$n_*^2 \approx$

\approx

\approx

\approx

$(1 \text{ TeV})^2$

$(130 \text{ GeV})^2$

$(60 \text{ GeV})^2$

$(70 \text{ GeV})^2$

$(10 \text{ GeV})^2$

$$\left(\frac{m_*}{3 \text{ TeV}}\right)^3 \left(\frac{\text{TeV}}{f}\right) (\delta_t^2 - 1)$$

$$\left(\frac{m_*}{3 \text{ TeV}}\right)^3 \left(\frac{\text{TeV}}{f}\right) (\delta_b^2 - 1)$$

$$\left(\frac{m_*}{3 \text{ TeV}}\right)^3 \left(\frac{\text{TeV}}{f}\right) (\delta_\tau^2 - 1)$$

$$\left(\frac{m_*}{3 \text{ TeV}}\right)^2 \left(\frac{y_\chi}{0.2}\right)^2 (1 - \delta_\chi^2)$$

$$\left(\frac{m_*}{3 \text{ TeV}}\right)^2 \left(\frac{1 \text{ TeV}}{f}\right)^2 \left[\frac{m_\chi^{(1)} m_\chi^{(6)}}{(25 \text{ GeV})^2} \right]$$

$\delta = 1$ corresponds to unbroken $U(1)_\eta$

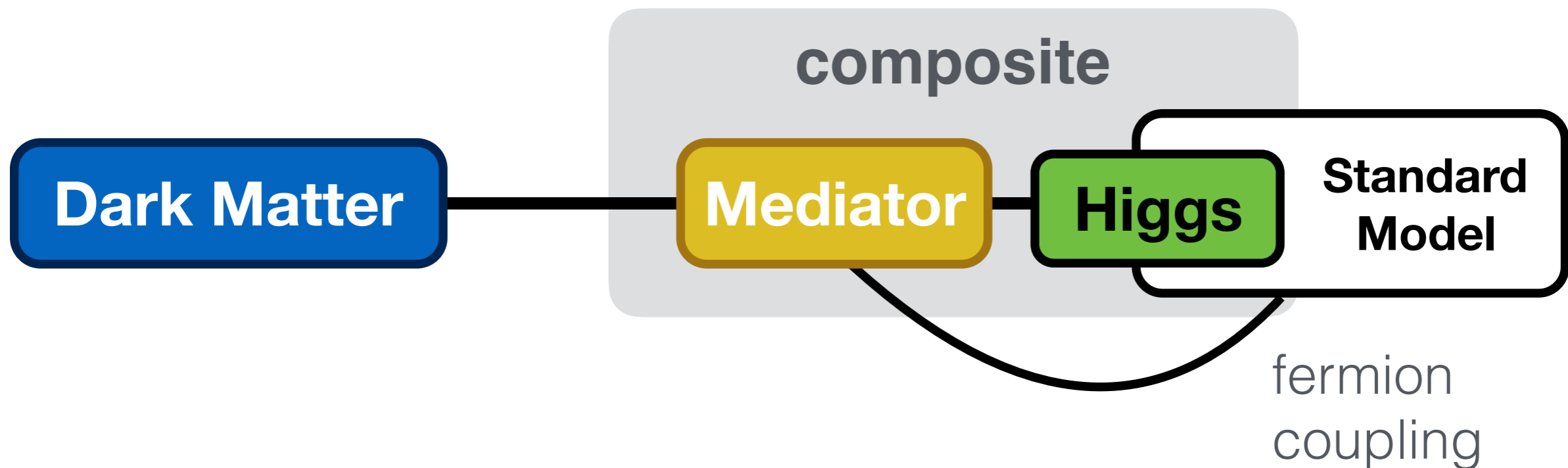
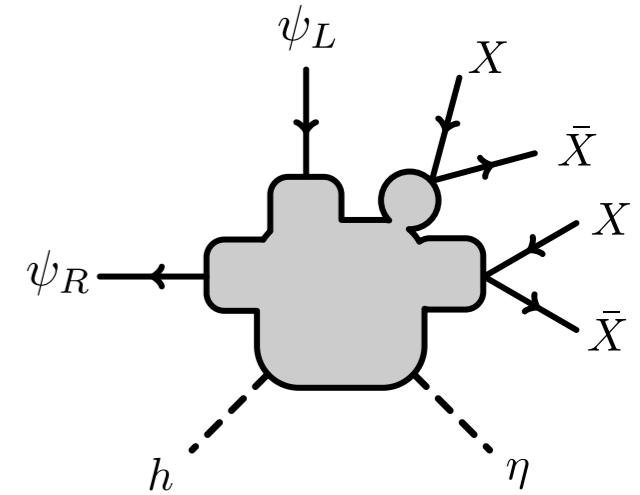
Punchline

New completion of pseudoscalar mediator

Connects **dark matter** and **Higgs naturalness**

Small number of parameters + discrete choices

In some sense, variant of “**WIMP**”



EXTRA SLIDES

Pseudoscalar Mediators

assume
Dirac X

$$(\bar{X} i\gamma^5 X)$$

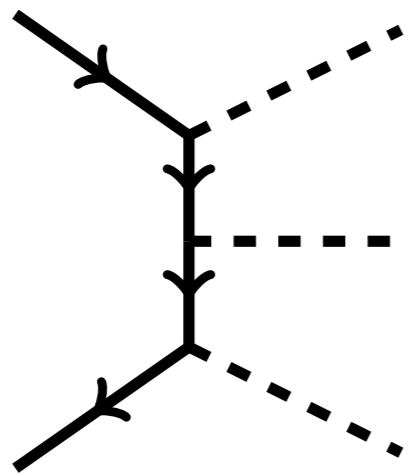
$$(\bar{q} i\gamma^5 q)$$

see, e.g.

Kumar & Marfatia 1305.1611

$\sigma_{\text{spin dependent}} \propto q^2 \times q^2$ suppressed
direct detection

$\sigma_{\text{spin independent}} = 0$



s-wave annihilation in s-channel or
to three on-shell pseudoscalars

Indirect detection:

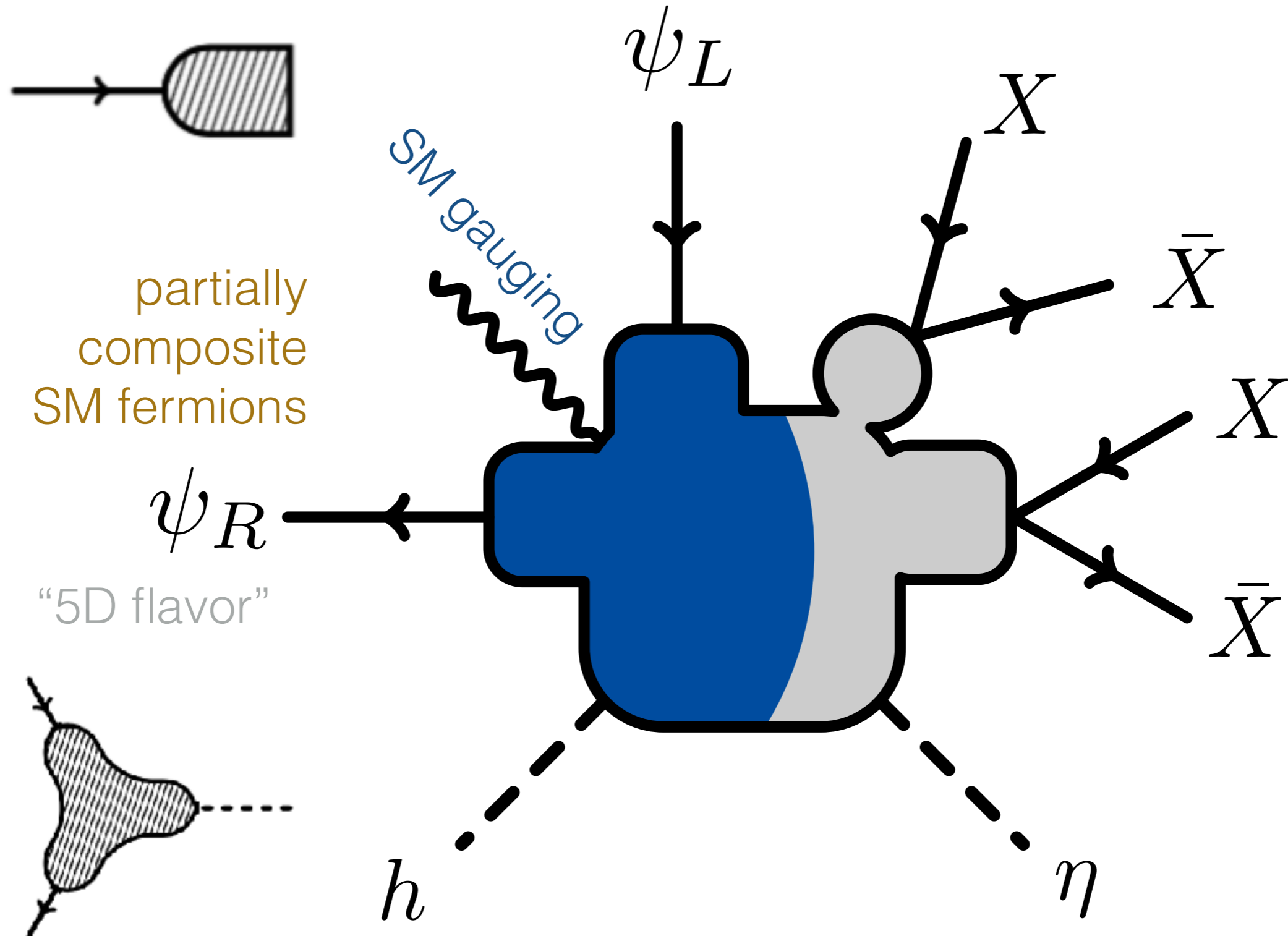
Boehm et al. 1401.6458

FT & UCI folks 1404.6528, 1503.05919

Berlin et al. 1502.06000

SO(5)/SO(4) & the Standard Model

described by usual minimal composite Higgs



Naive Dimensional Analysis

one scale, one coupling ansatz

mass scale of
new resonances

bosons

fermions

$$\mathcal{L} = \frac{m_*^4}{g_*^2} \hat{\mathcal{L}} \left(\frac{\partial}{m_*}, \frac{g_* \Phi}{m_*}, \frac{g_* \Psi}{m_*^{3/2}} \right)$$

characteristic
resonance coupling

$$g_* = \frac{m_*}{f} \text{ breaking scale}$$

scaling from: mass and \hbar dimensional analysis

Georgi & Manohar '84; see e.g. Panico & Wulzer 1506.01961 for a review