

**High Resolution High Photon Energy
Small Photon Spot
Bulk Sensitive Photoemission Study of the
Metal-Insulator Transition
in $(V_{1-x}Cr_x)_2O_3$:
Comparison to LDA + DMFT Theory**

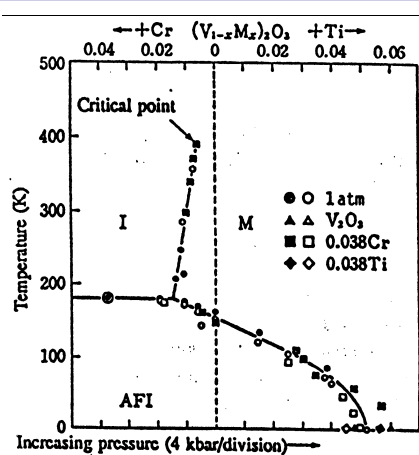
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K. Held (MPI Stuttgart), V.I. Anisimov (IFMLRS)

Work at UM Supported by the U.S. NSF and the U.S. DoE.

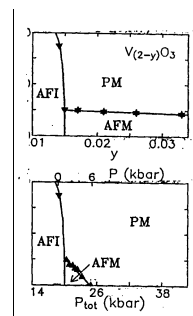
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V_2O_3 phase diagram



McWhan et al. PRL '69, PRB '73



S.A. Carter et al. PRL '91
New phase boundary

Interpreted as Mott transition
Brinkman-Rice for 1-band Hubbard

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Some standard history for V_2O_3

- McWhan, Rice *et al.* (PRL, PRB 1969 – 1973)
Vertical pair of V ions stabilize a_{1g} singlet $\rightarrow a_{1g}e_g$, $S = 1/2$
Ignore orbital degeneracy, use 1-band Hubbard
- Castellani *et al.* (PRB 1978).
Inclusion of orbital ordering for the AFI phase
Keep a_{1g} singlet & $a_{1g}e_g$ & $S = 1/2$
- Kotliar, George, and co-workers (e.g. Rev. Mod. Phys 1996)
DMFT for 1-band Hubbard
Analysis of optical spectroscopy
UV photoemission

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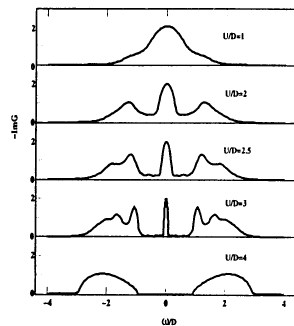
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DMFT for metal-insulator transition

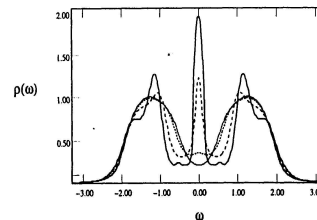
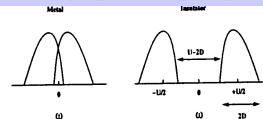
**1-band Hubbard, $d \rightarrow \infty$
k-independent self energy**

A. George, G. Kotliar, W. Krauth, M. Rozenberg Rev. Mod. Phys '96

**Spectral function for varying U
"Kondo-like" picture**



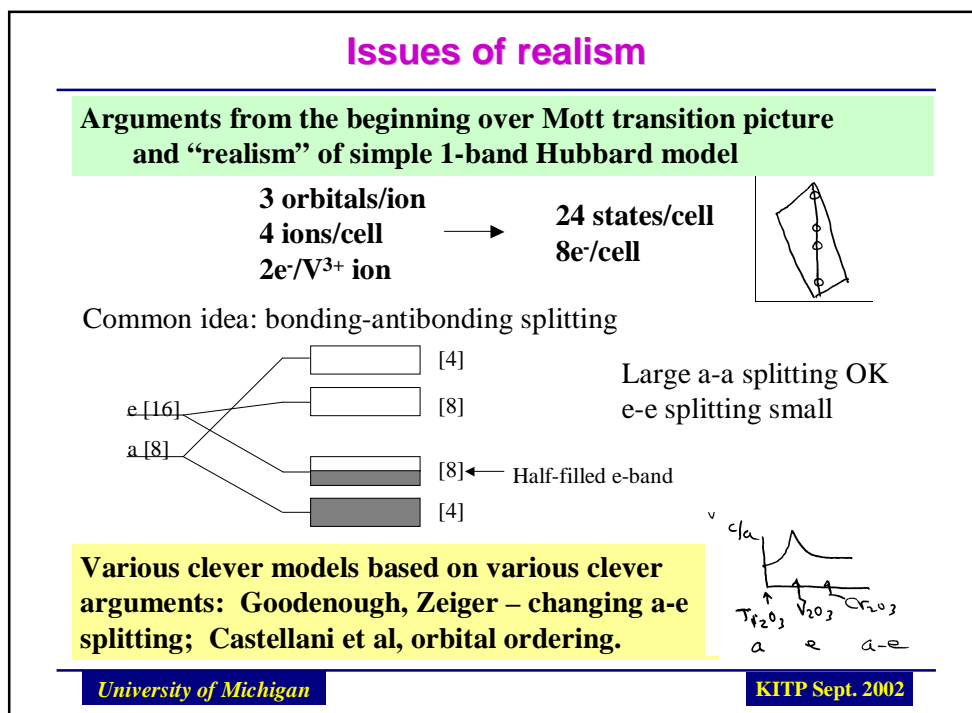
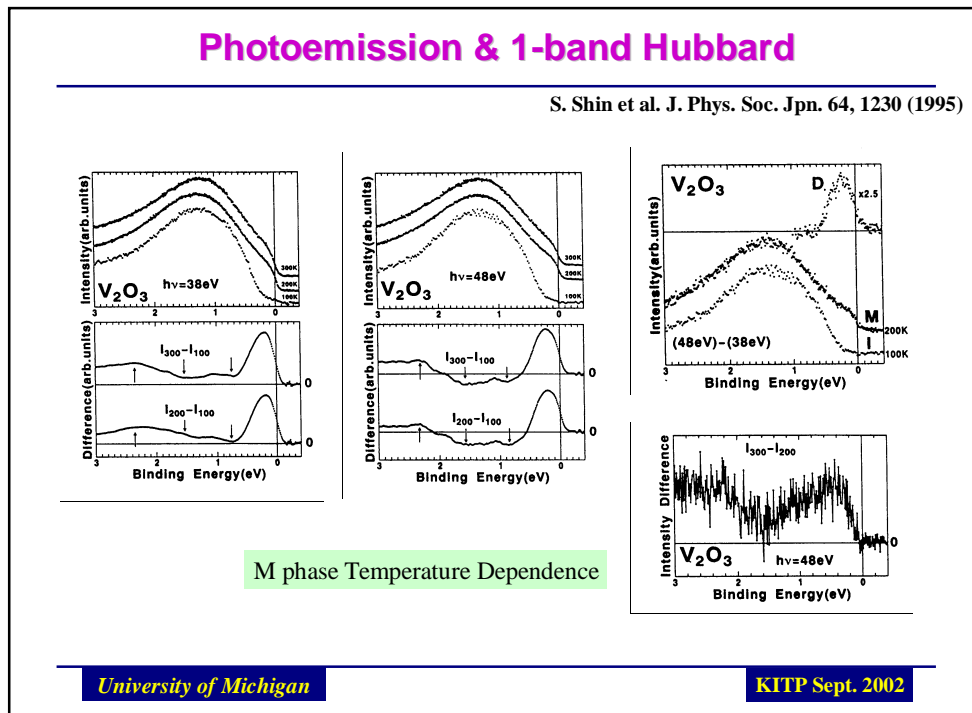
Old gap collapse picture



**Increase T in metallic phase
- lose quasi-particle peak**

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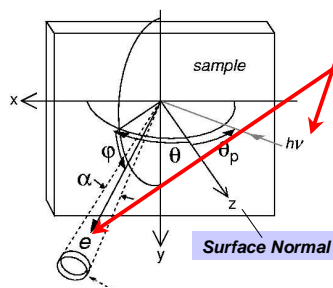
Newer work--multiband models needed for V₂O₃

- LDA+U: Ezhov *et al.* (PRL 1999).
only e occupation, S=1 on site and no e site-ordering
- Paolasini *et al.* (PRL 1999).
resonant x-ray scattering claims to see e-site ordering
- Park *et al.* (PRB 2000).
Polarization-dependent x-ray absorption
S=1 and ee:ea = 2:1 in AFI phase
- Mila *et al.* (PRL 2000) and Di Matteo *et al.* (cond-mat/2001)
2 different correlated models of c-axis pair states for AFI phase
dynamic mix of ee and ea with S=1 on sites
- Held *et al.* (PRL 2001) LDA+DMFT
multiband many body realism for PM and PI phases and
single particle spectra—COMPARE TO PHOTOEMISSION

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Angle resolved photoemission spectroscopy (ARPES) to measure $\rho(k, \omega)$



- **Photon In**
- **(Photo-)Electron Out**
- Electron KE, $h\nu$, \Rightarrow bind. en. ω
- Angles $\theta, \phi \rightarrow$ k-par, cons. at surf.
- k-perp -- not conserved, must model surface potential
- Electron Energy Distribution (ω)
= $\rho(k, \omega) \times$ Fermi function
 \times (ARPES cross-section)

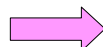
MDC (fix ω , scan k)

EDC (fix k, scan ω)

"FS" map ($\omega=E_F$, scan k region)

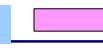
Cross-section resonances specific to atom and ang. mom.
for $h\nu$ at core level level absorption edges----- "RESPES"

Varying angles, move on spherical
k-space surface, radius fixed by $h\nu$



k-perp changes with $h\nu$

Larger KE, larger escape depth

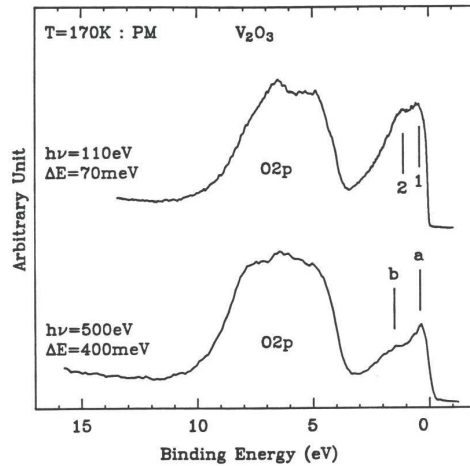


Higher $h\nu$, more bulk

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Early evidence of bulk/surface difference



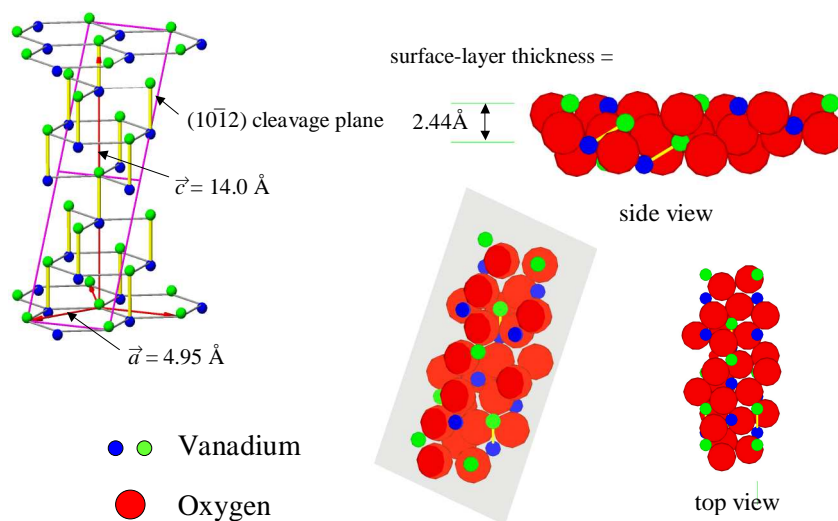
**J.-H. Park thesis
Univ. of Michigan, 1994**

Systematic reduction of near E_F peak in metallic phase for low photon energy relative to high photon energy implies surface effect but resolution not good at high photon energy at that time.

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Crystal structure and surface layer



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High photon energy high resolution bulk sensitive resonant photoemission at SPring8

- Samples from Purdue University: single crystals of $(V_{1-x}R_x)_2O_3$
 $x=0$; $R=Cr$ ($x=1.2\%$, 1.8% , 2.8%); $R=Ti$ ($x=1.0\%$)
- Spectrometer: beamline BL25SU, SPring-8, Japan with SCIENTA SES200
- Base pressure: low 10^{-10} Torr.
- Sample cleaning: cleaved to expose a (10-12) plane
- Energy resolution: ~ 90 meV.
- Fermi level position: Fermi edge of Pd metal electrically connected with a sample.

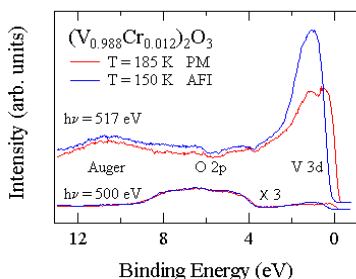


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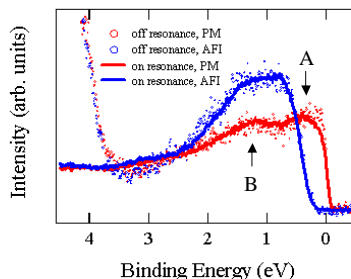
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Initial results, slide 1: RESPES of $(V_{0.988}Cr_{0.012})_2O_3$ at V 2p \rightarrow 3d edge

- On- and off-resonance spectra



- V 3d spectra



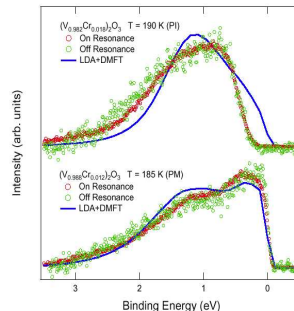
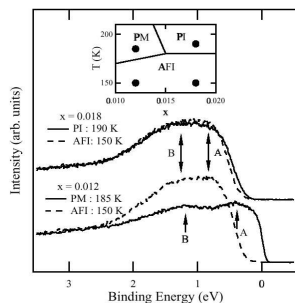
- Strong enhancement of V 3d signal at $2p \rightarrow 3d$ edge.
- Within signal to noise of off-resonance spectrum, on- and off-resonance spectra of V 3d essentially identical $\Rightarrow 2p \rightarrow 3d$ RESPES provides V 3d spectral function.

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Prominent Metal Phase Quasi-particle and High Temperature Correlation Gap Filling in Photoemission Spectra of $(V_{1-x}Cr_x)_2O_3$: Comparison to LDA+DMFT Theory

Initial results, slide 2



Compare to LDA+DMFT(QMC) 1100K from Held et al, PRL 86, 5345 (2001).

Fair agreement for PM phase— E_F feature not strong, with maximum ≈ 0.3 eV below E_F . Prominent central peak smeared out due to multiband realism?

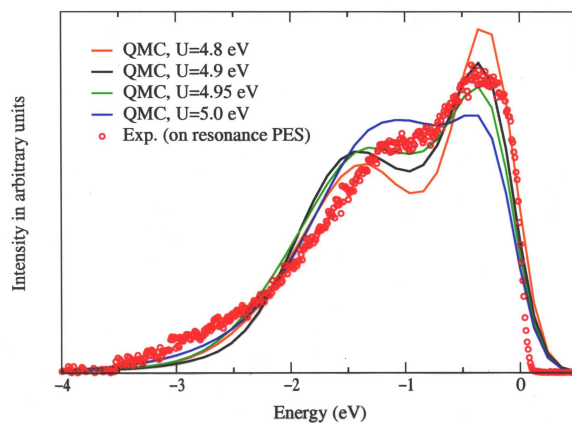
Less good agreement for PI phase—theory shows gap filling at high T of calculation, and other spectral differences.

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Initial results, slide 3: LDA+DMFT (1100K) – for various U values & initial data

Relative magnitudes of coherent/incoherent peaks change rapidly with U in 1100K theory—U 4.95 eV OK

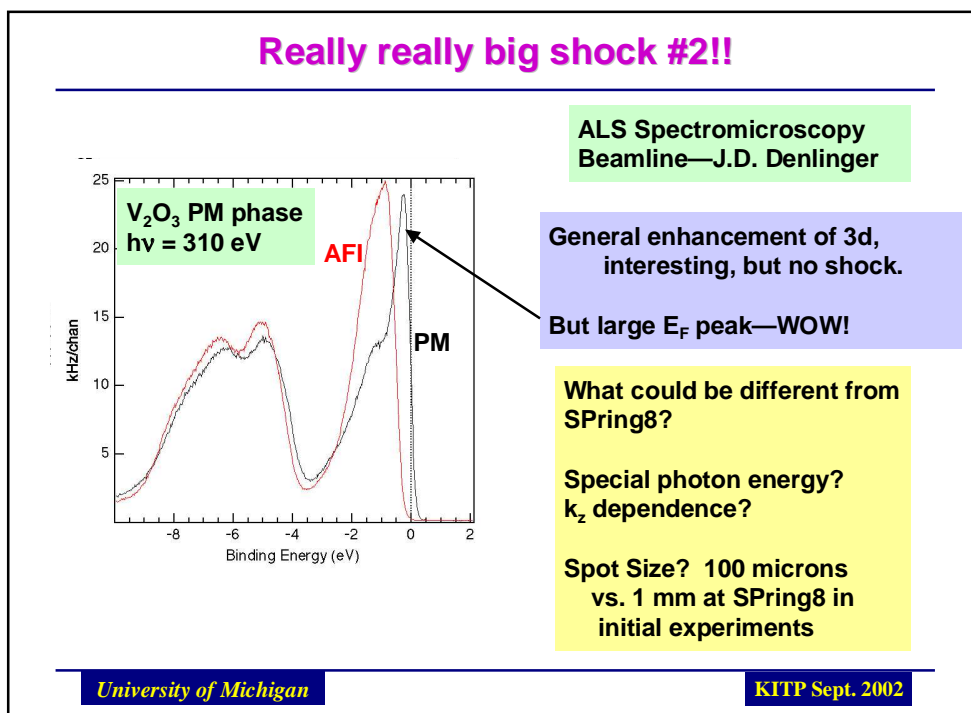
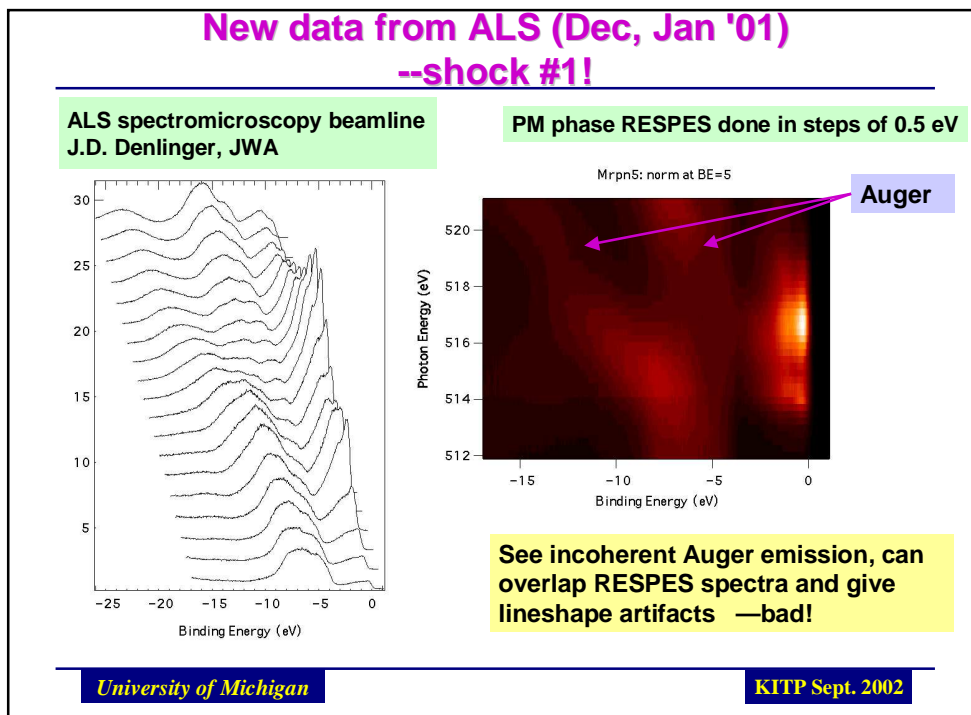


Peak below E_F in spectrum agrees with theory

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Prominent Metal Phase Quasi-particle and High Temperature Correlation Gap Filling in Photoemission Spectra of (V_{1-x}Cr_x)₂O₃: Comparison to LDA+DMFT Theory



**Back to SPring8, April ' 02
---Pray we can sort it out!**

**Planned to focus on large hv ARPES--
did more PES instead**

Good fortune—

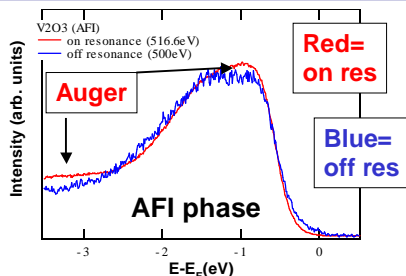
**Beam line had been modified to decrease photon spot size
from 1 mm to 100 μm, same as ALS**

**Spot size decrease also increased photon intensity so could
measure off resonance with good S/N and the full high
resolution capability that makes this beamline unique.**

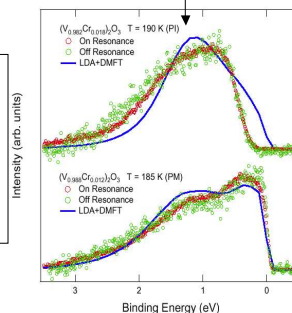
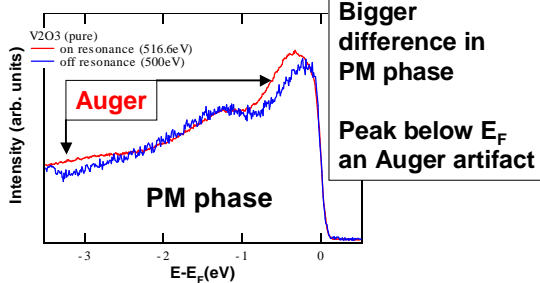
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Yes, find Auger in RESPES lineshape



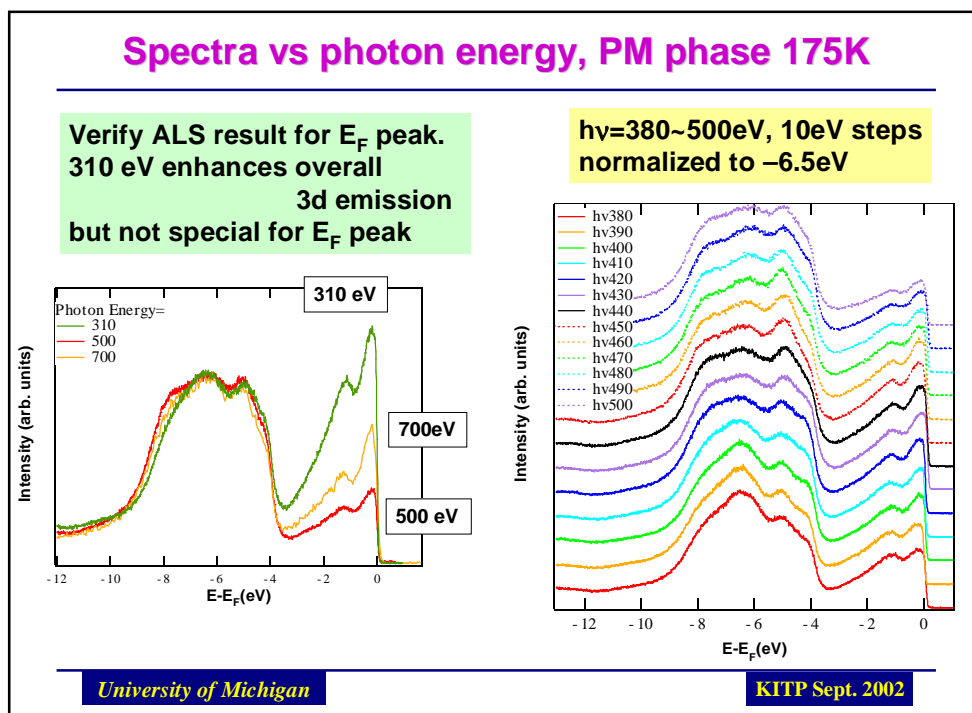
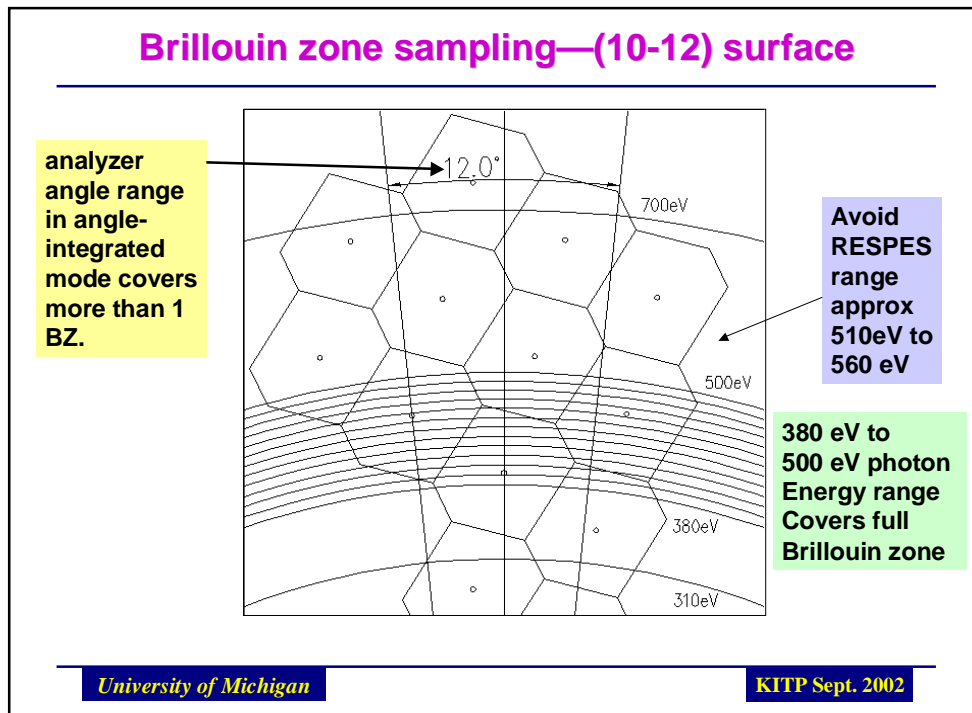
**With good resolution and good
S/N can see on-res/off-res
difference previously
hidden in poor statistics of initial
data for off-resonance spectra.**

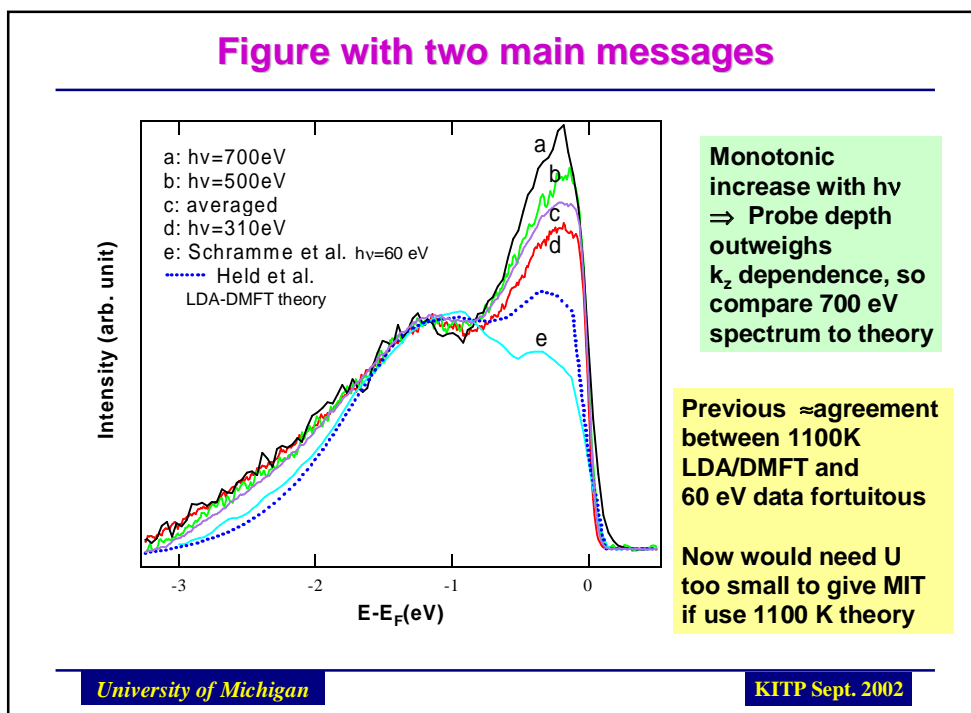
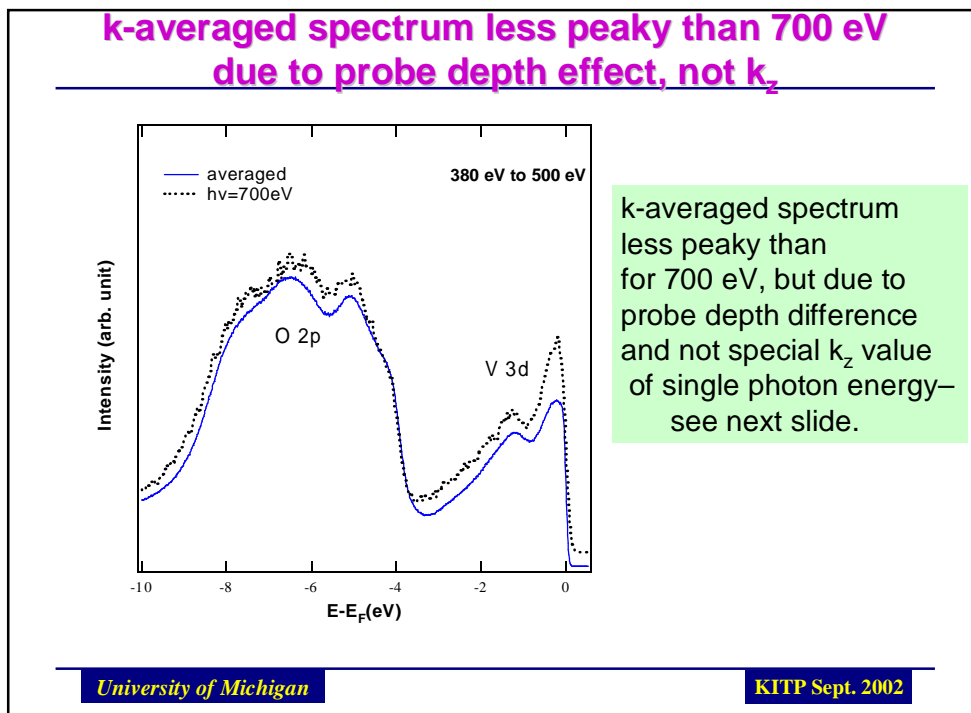


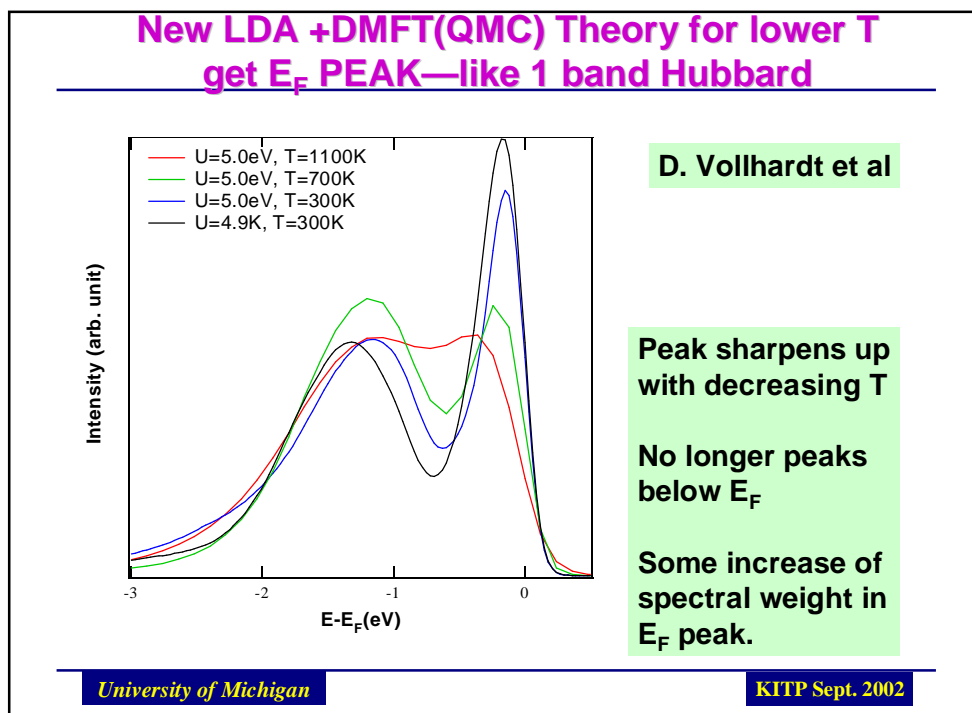
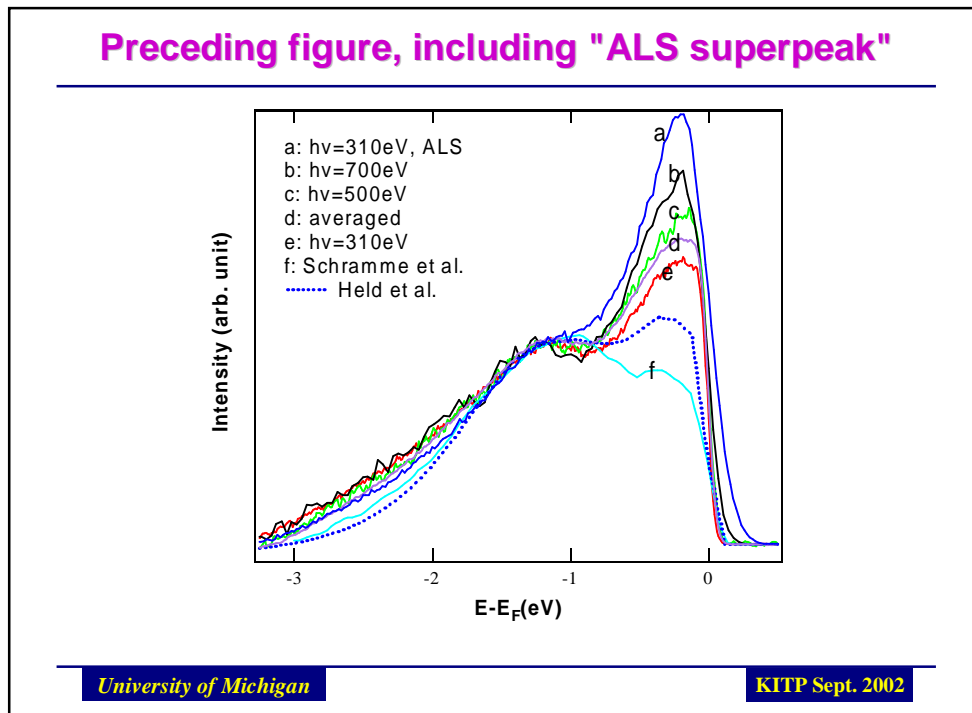
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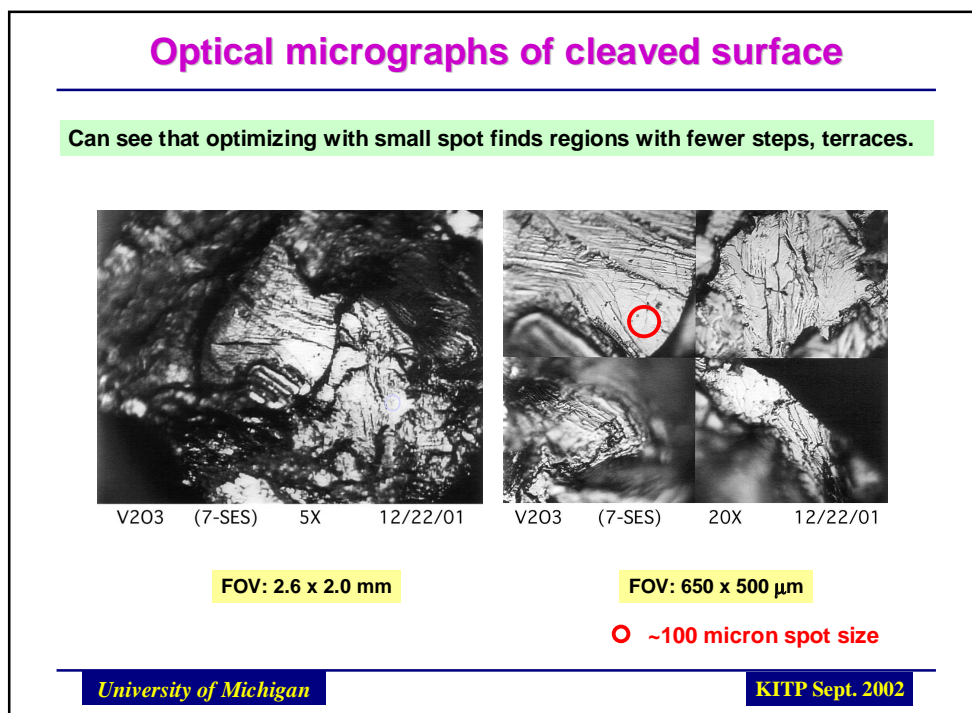
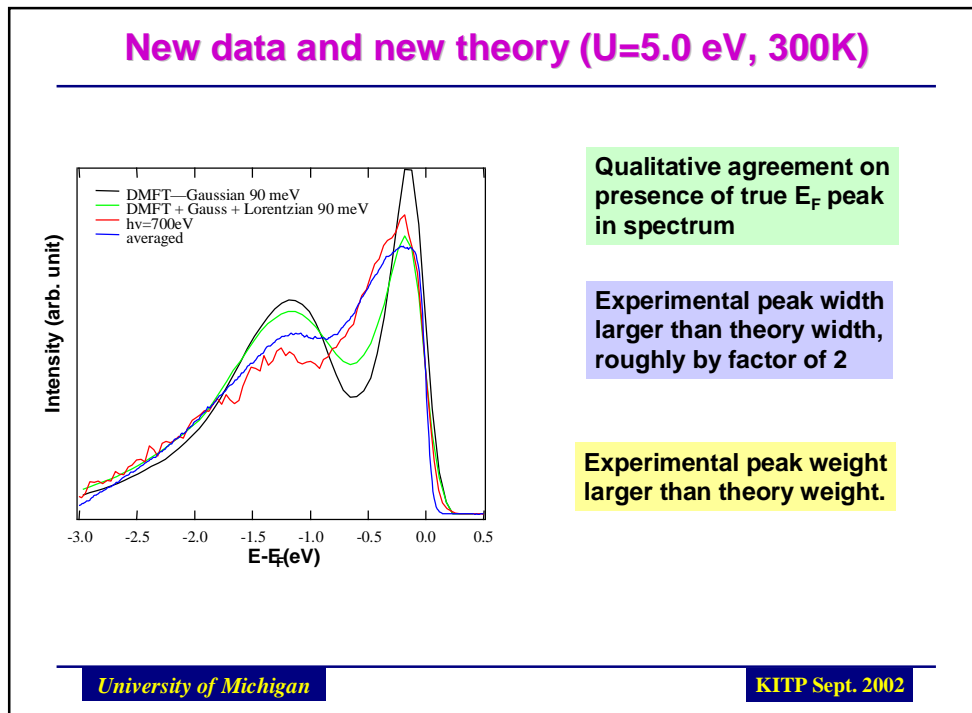
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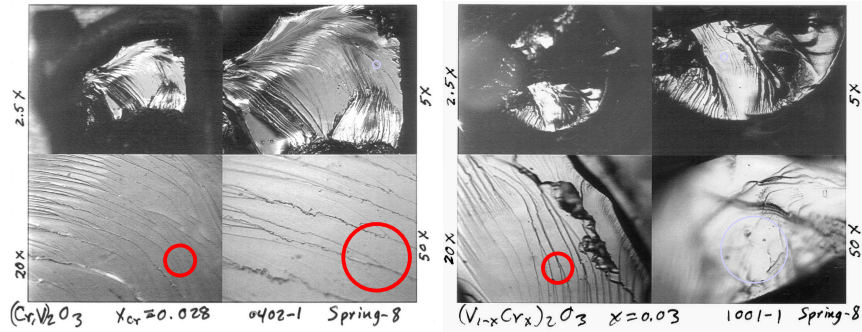






Optical micrographs of cleaved surfaces

Cr doped samples seem to give nicer cleaves, smoother, fewer terraces.



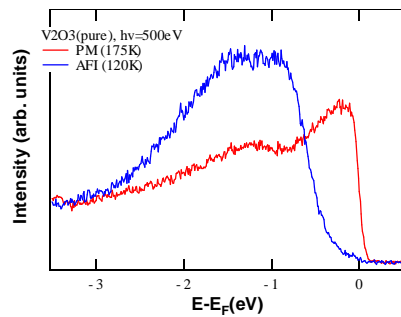
Field of View: 2.5X (5.2 x 4.0 mm) 5X (2.6 x 2.0 mm)
20X (650 x 500 μm) 50X (260 x 200 μm)

○ ~100 micron spot size

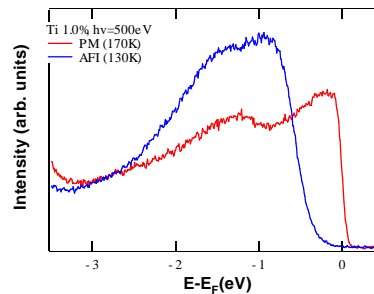
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PM/AFI transition – pure V_2O_3 and Ti 1.0%



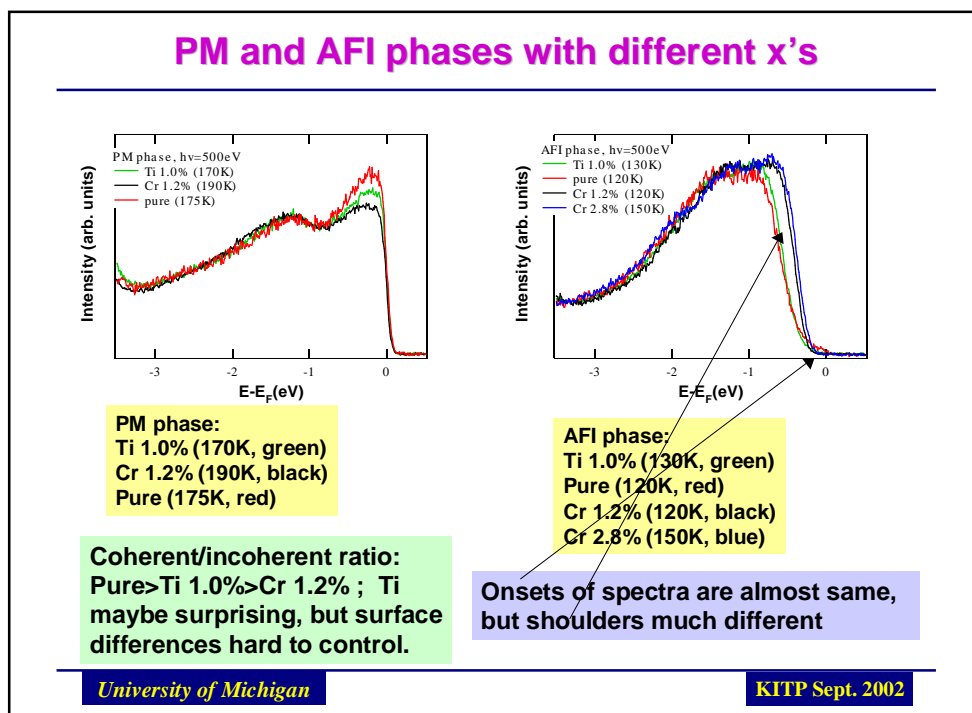
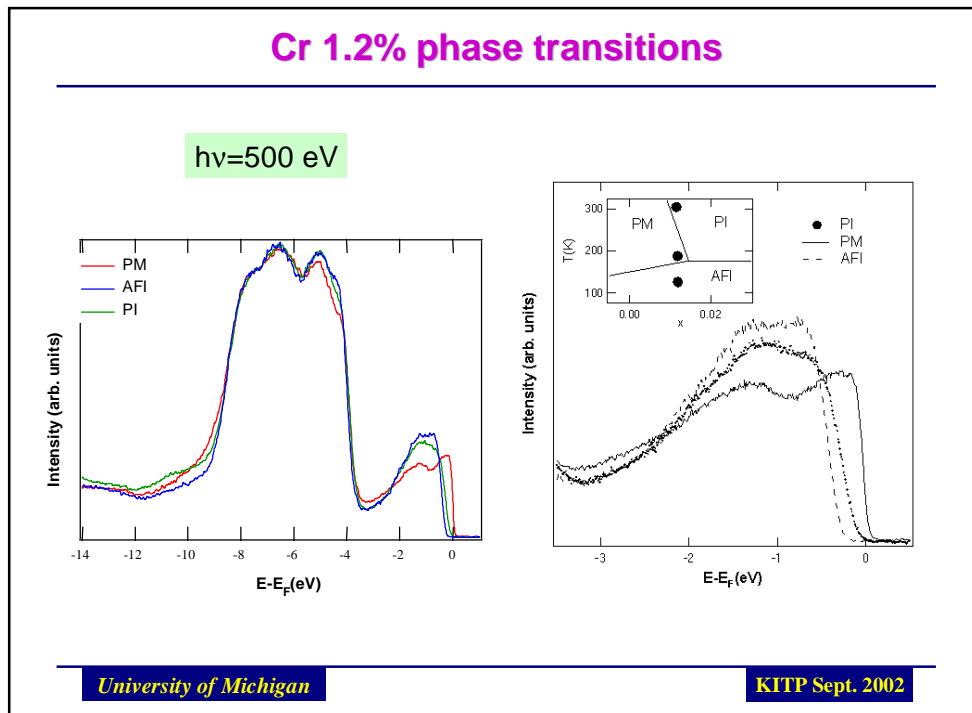
Pure V_2O_3 PM (175K) red
and AFI (120K) blue

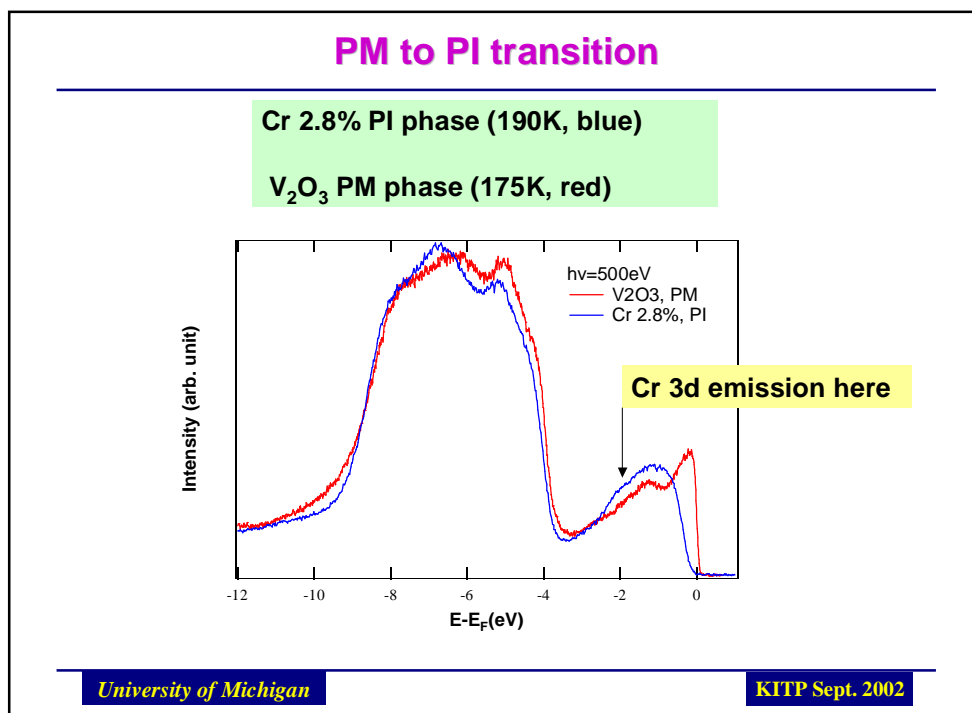
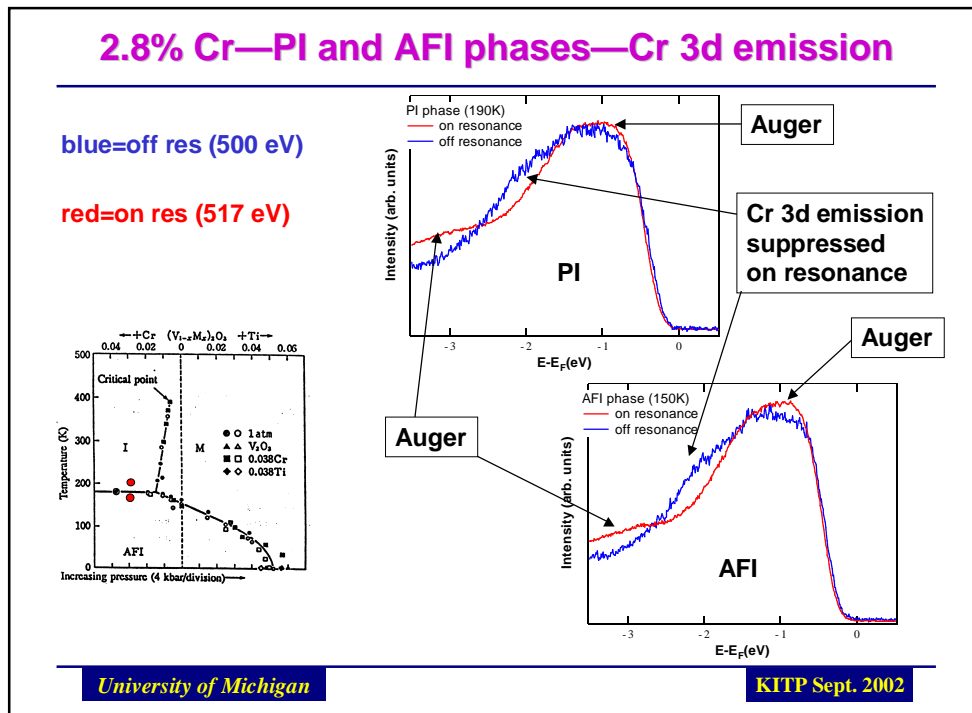


Ti 1.0% PM(170K, red)
and AFI (130K, blue)

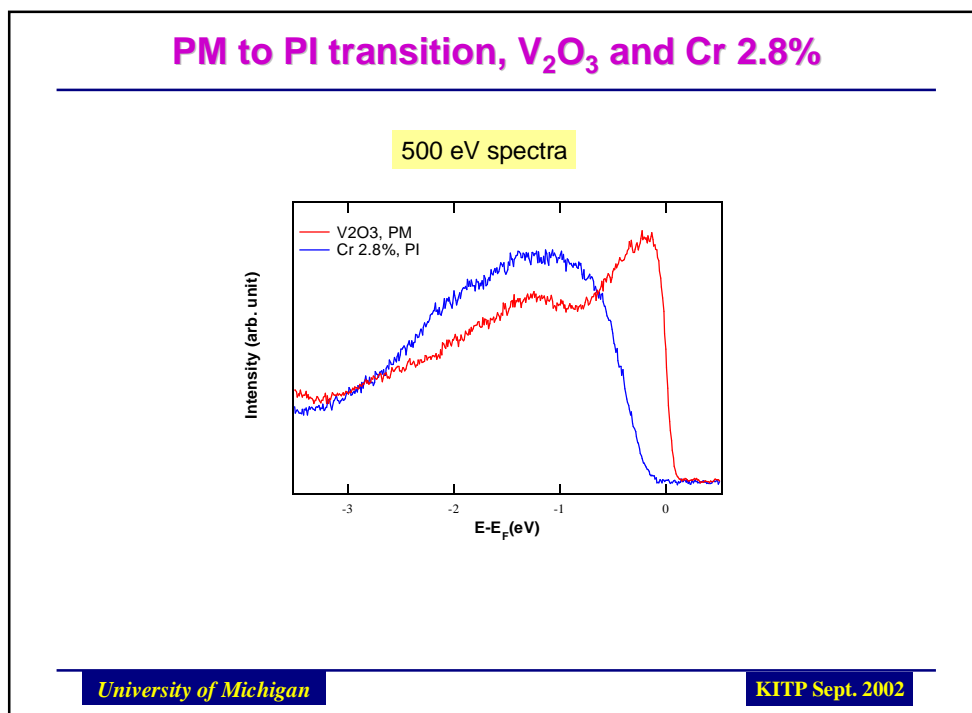
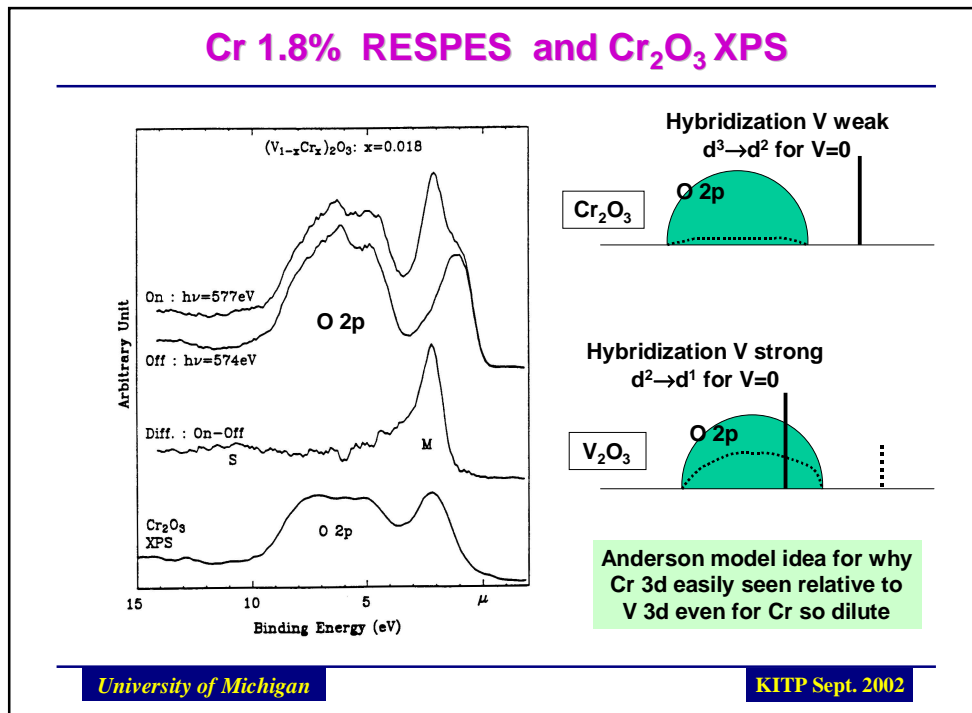
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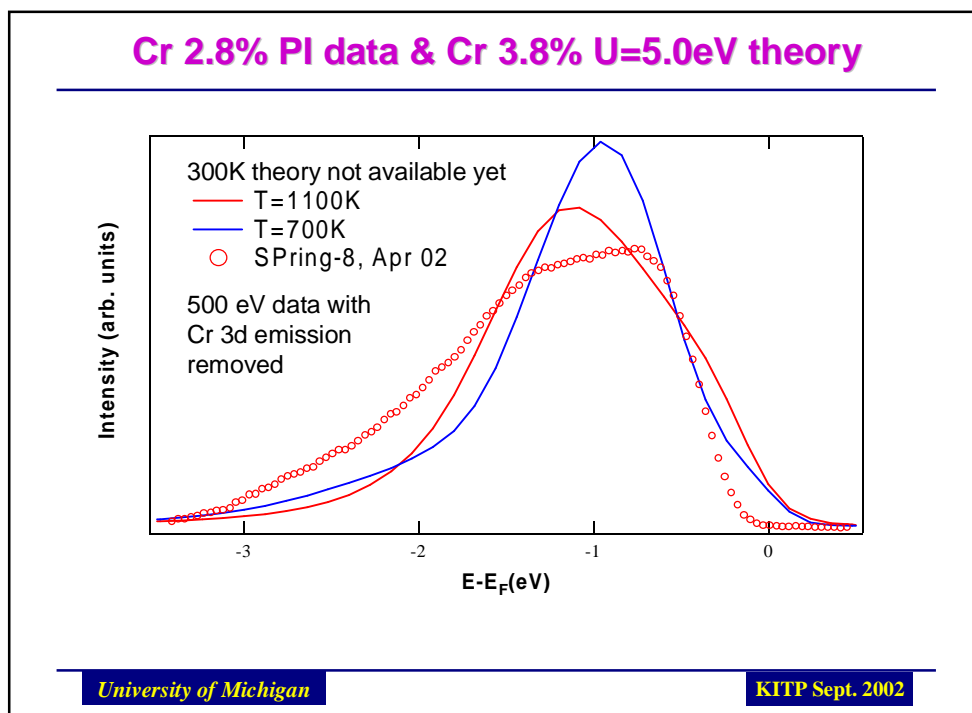
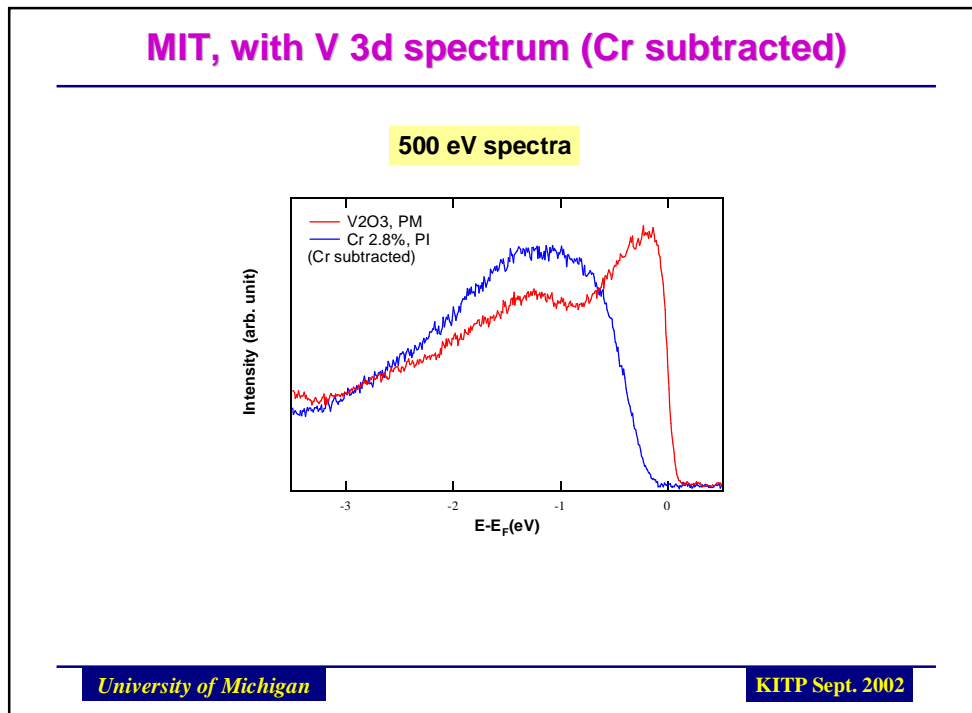
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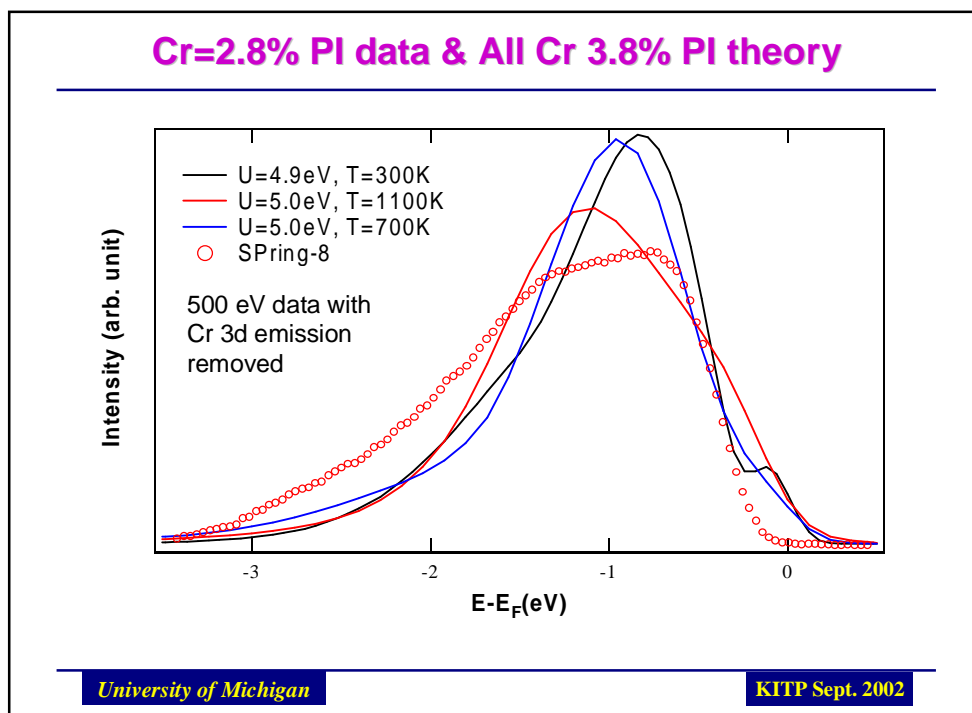
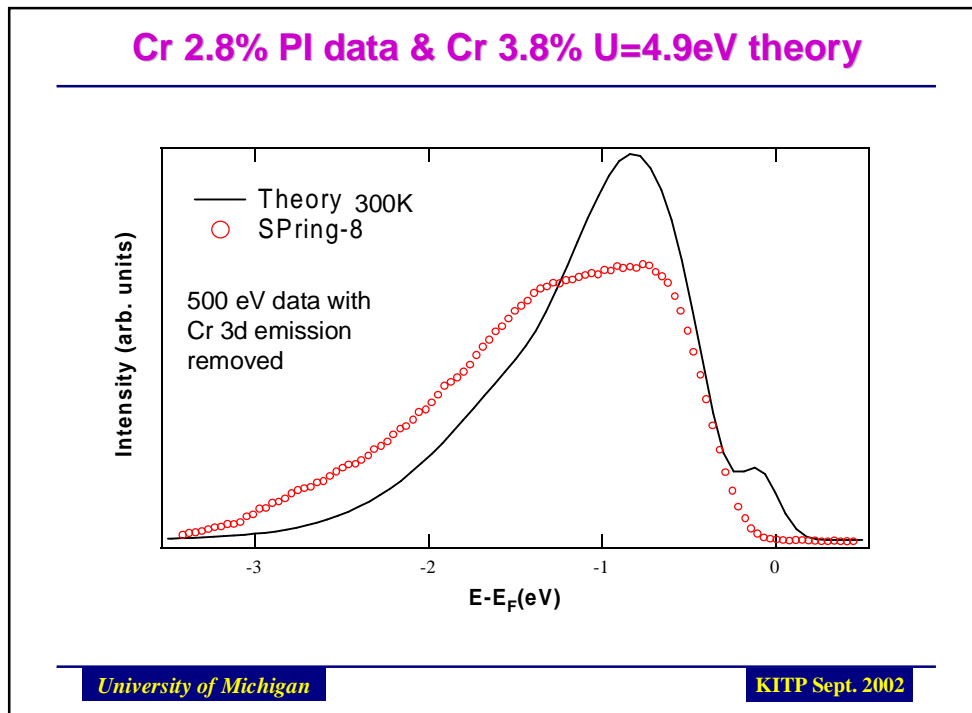




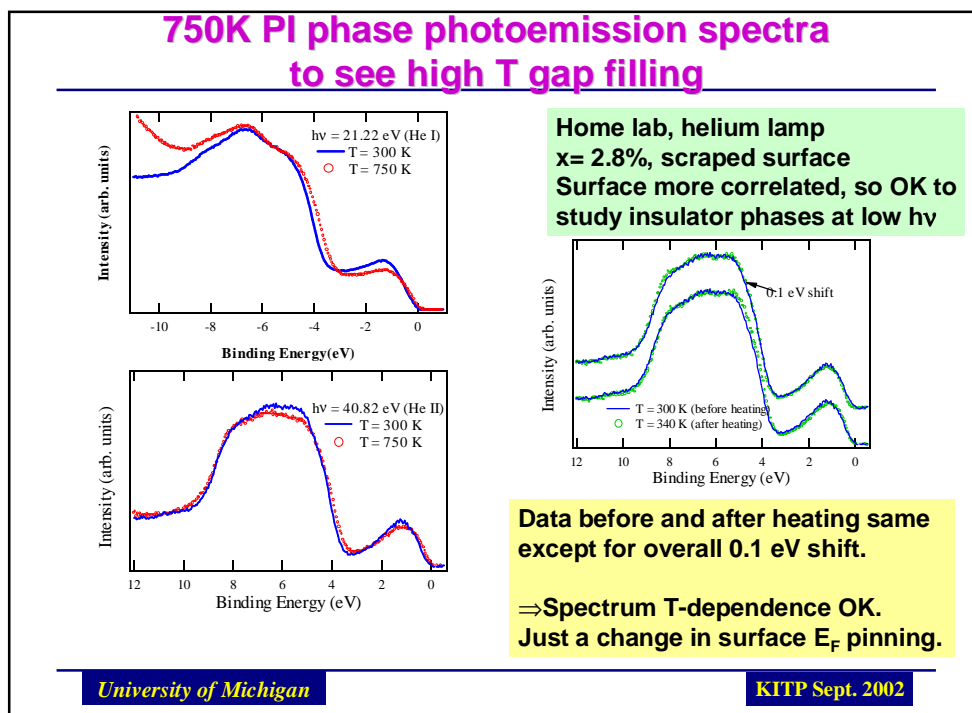
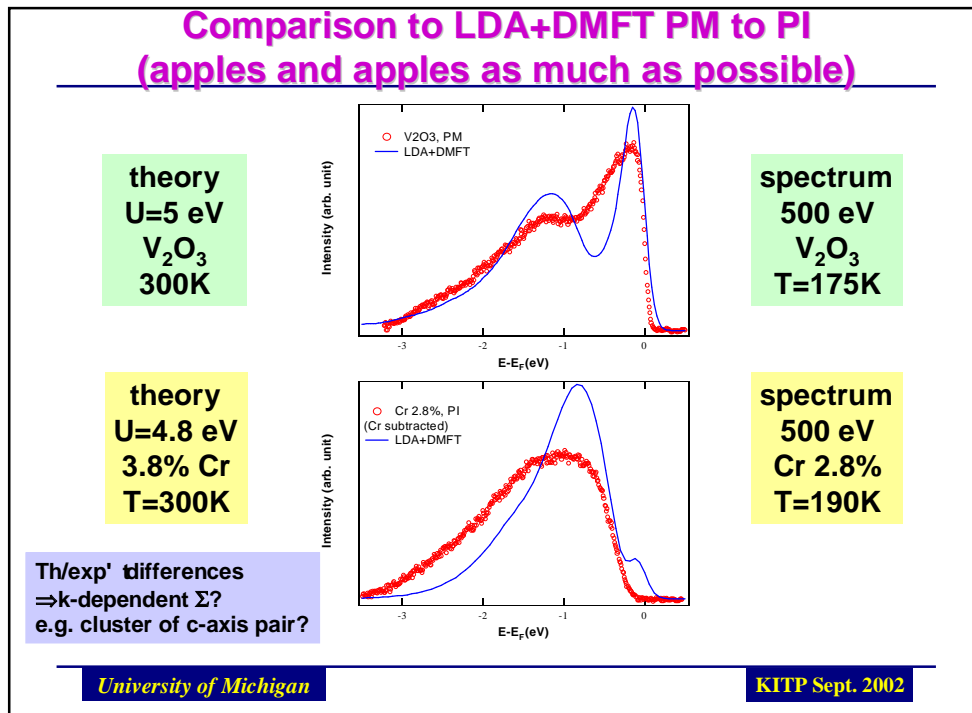
Prominent Metal Phase Quasi-particle and High Temperature Correlation Gap Filling in Photoemission Spectra of $(V_{1-x}Cr_x)_2O_3$: Comparison to LDA+DMFT Theory

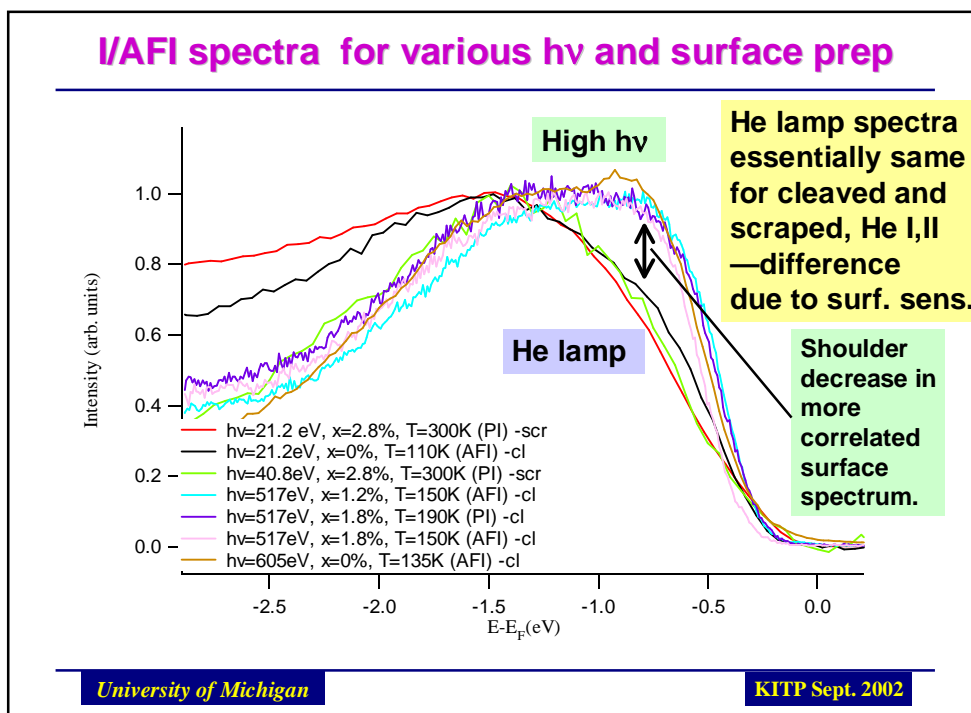
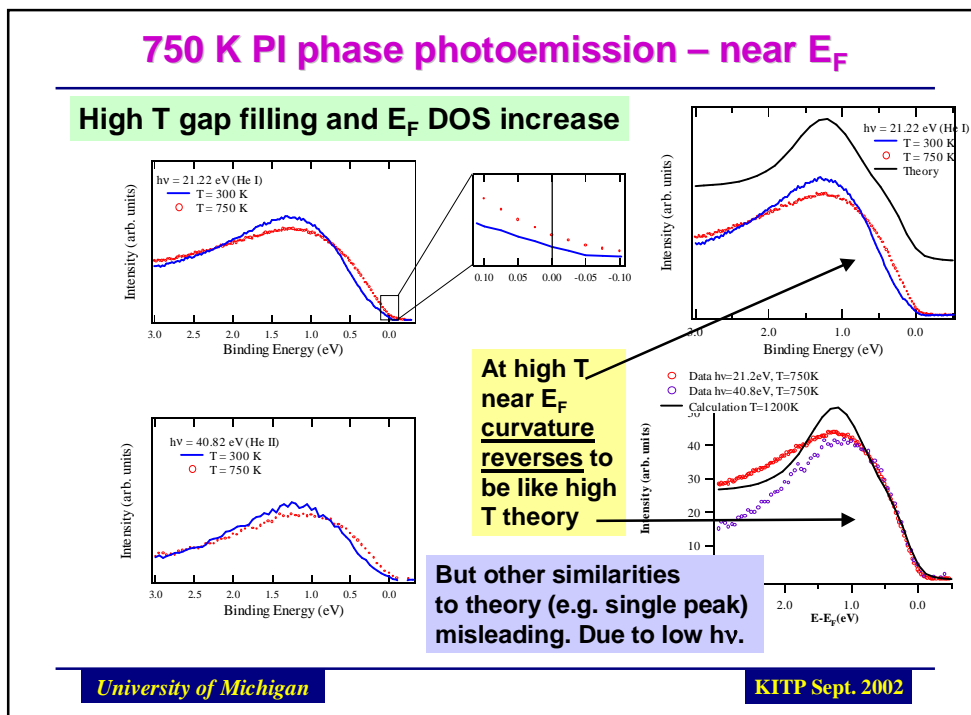






Prominent Metal Phase Quasi-particle and High Temperature Correlation Gap Filling in Photoemission Spectra of (V_{1-x}Cr_x)₂O₃: Comparison to LDA+DMFT Theory





Some early optical studies

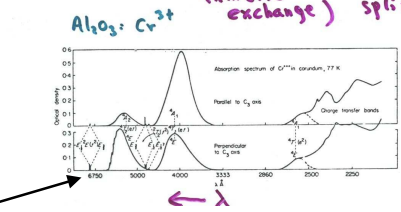
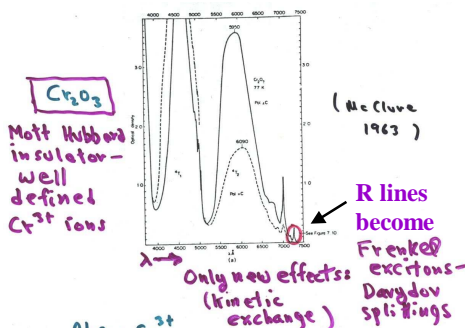
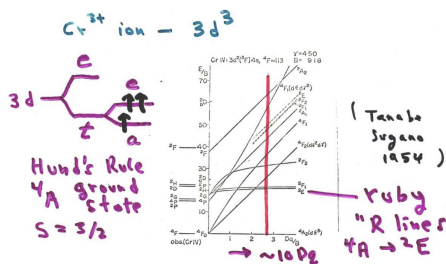
End by brief summary of evidence from optical spectra for important role of c-axis pairs in forming electronic structure of PI and AFI phases.

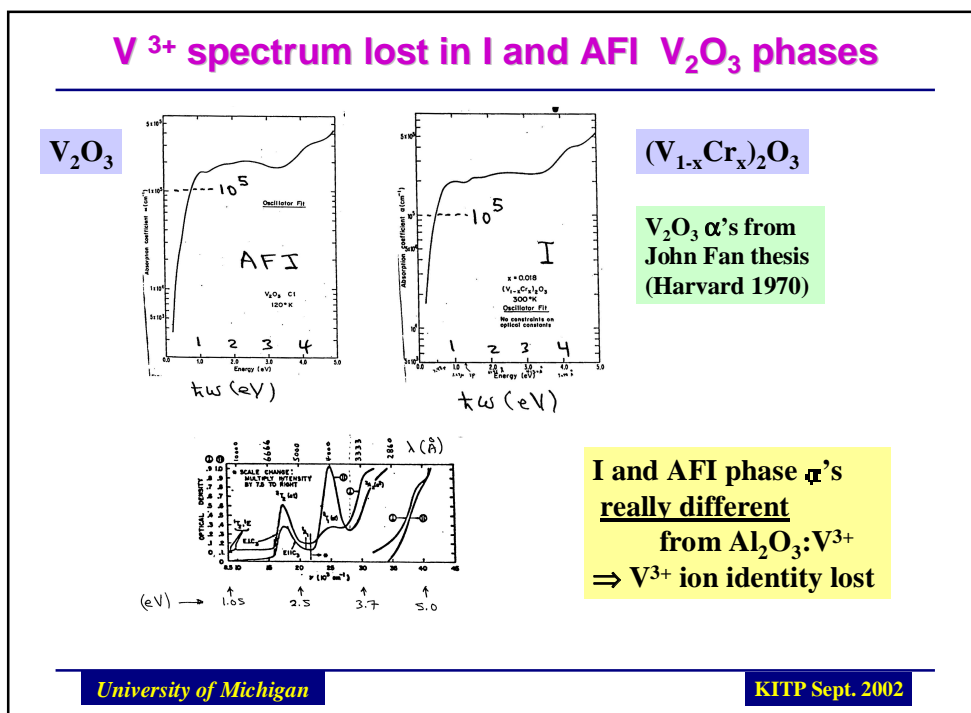
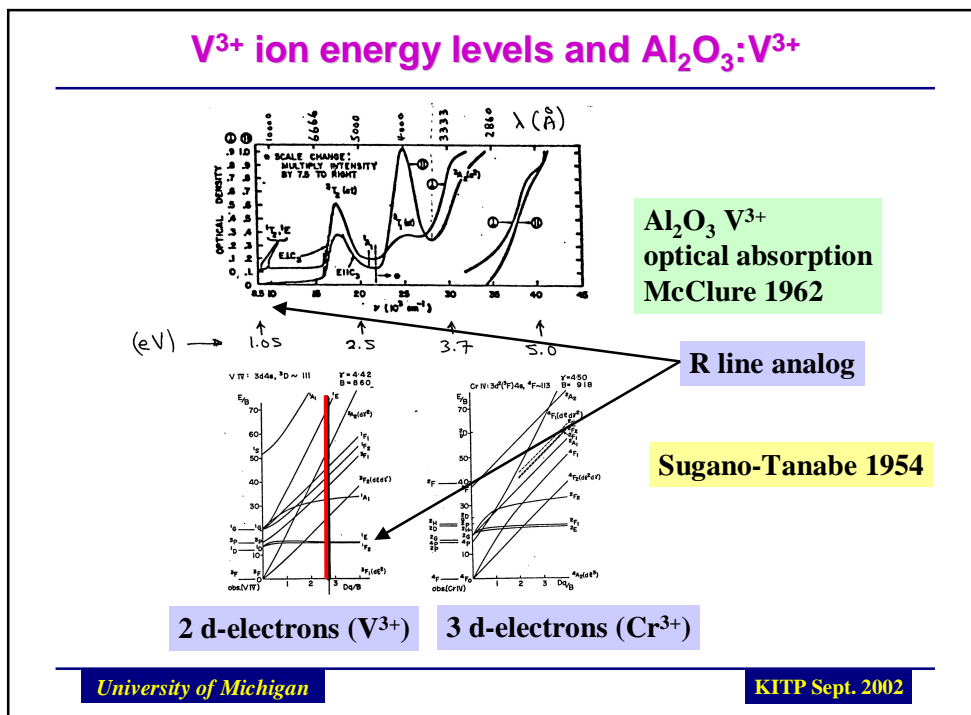
(The idea does not mean that other V-V couplings are unimportant, just that couplings for c-axis pairs are the largest.)

Cr³⁺ optical spectrum in ruby and Cr₂O₃

Cr³⁺ optical absorption spectrum of ruby persists in concentrated crystal—just weakly coupled Cr³⁺ ions

Cr₂O₃ a “real” Mott-Hubbard insulator. Same for NiO, etc.

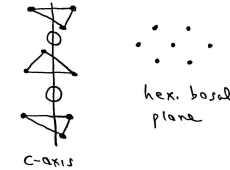




$(V_2O_3)_x(Al_2O_3)_{1-x}$ – early optical study

A. $x \sim 0.1\%$ crystal field spectrum
 $x \sim 100\%$ no V^{3+} spectrum

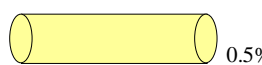
1. Evolution of optical properties with x
2. Do coupled pairs of V^{3+} ions retain basic ionic identity?



C-axis
hex. basal plane

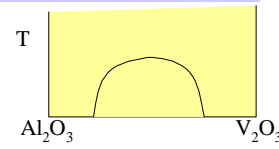
B. Samples (T. B. Reed)

1. Concentration gradient



10% 0.5%

2. V rich 2nd phase, $x > .08$



Al₂O₃ V₂O₃

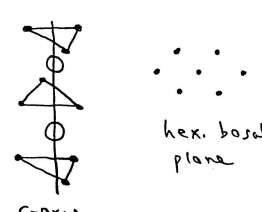
Neighbors

Sharing oxygens	#
1 st	1
2 nd	3
3 rd	3
4 th	6

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$(V_2O_3)_x(Al_2O_3)_{1-x}$ – doping statistics

C. Corundum crystal structure



C-axis
hex. basal plane

Neighbors	#
1 st	1
2 nd	3
3 rd	3
4 th	6

D. Statistics (assuming random distribution)

$P(\text{no neighbor}) = (1-x)^{13}$
 $P(\text{neighbor}) = 1 - (1-x)^{13}$
 $P(\text{1st neighbor}) = x = 1 - (1-x)^1$
 $P(\text{2nd neighbor}) = 1 - (1-x)^3$
 $P(\text{3rd neighbor}) = 1 - (1-x)^3$
 $P(\text{4th neighbor}) = 1 - (1-x)^6$

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