

Upper limit on spontaneous supercurrents in Sr_2RuO_4

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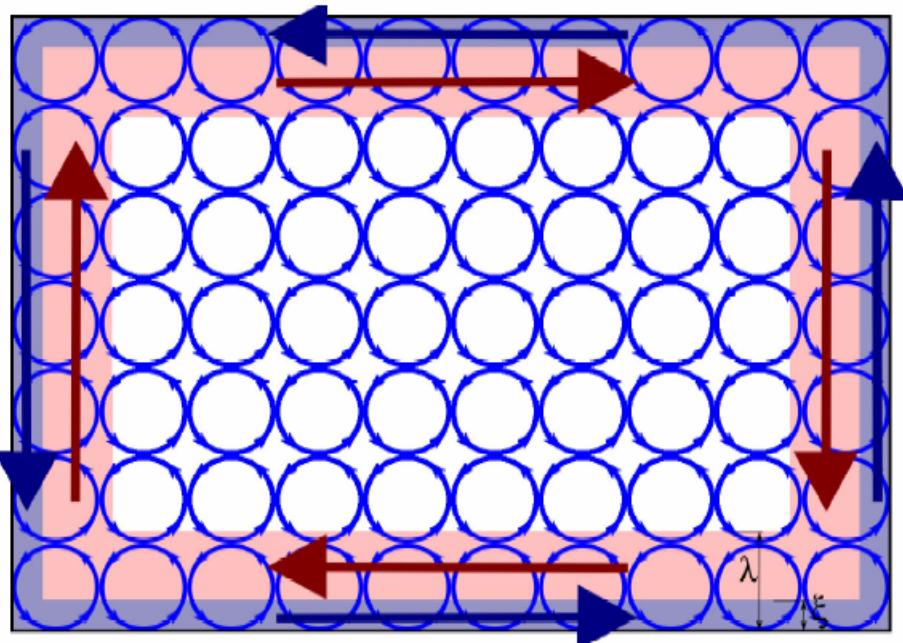
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Spontaneous supercurrents predicted to occur at edges of p_x+ip_y superconductors

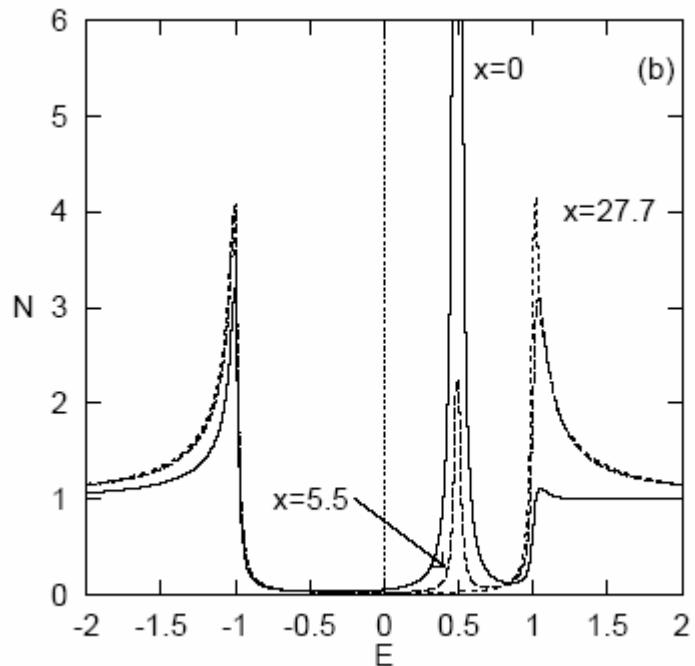
each Cooper pair has angular momentum- edges not cancelled out



for Sr_2RuO_4 - $2.8\mu\text{A}/\text{plane}$

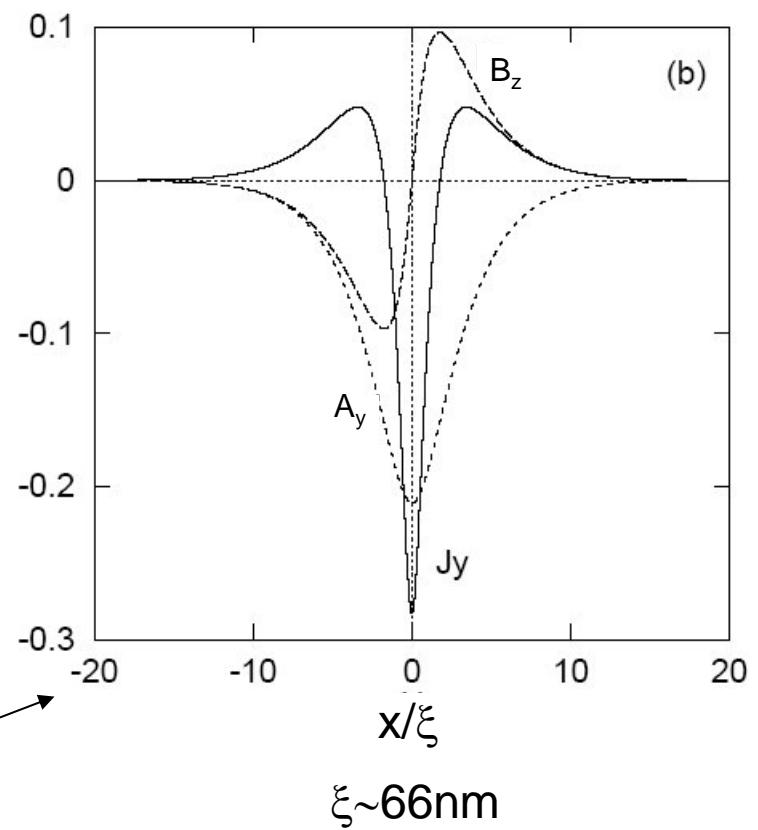
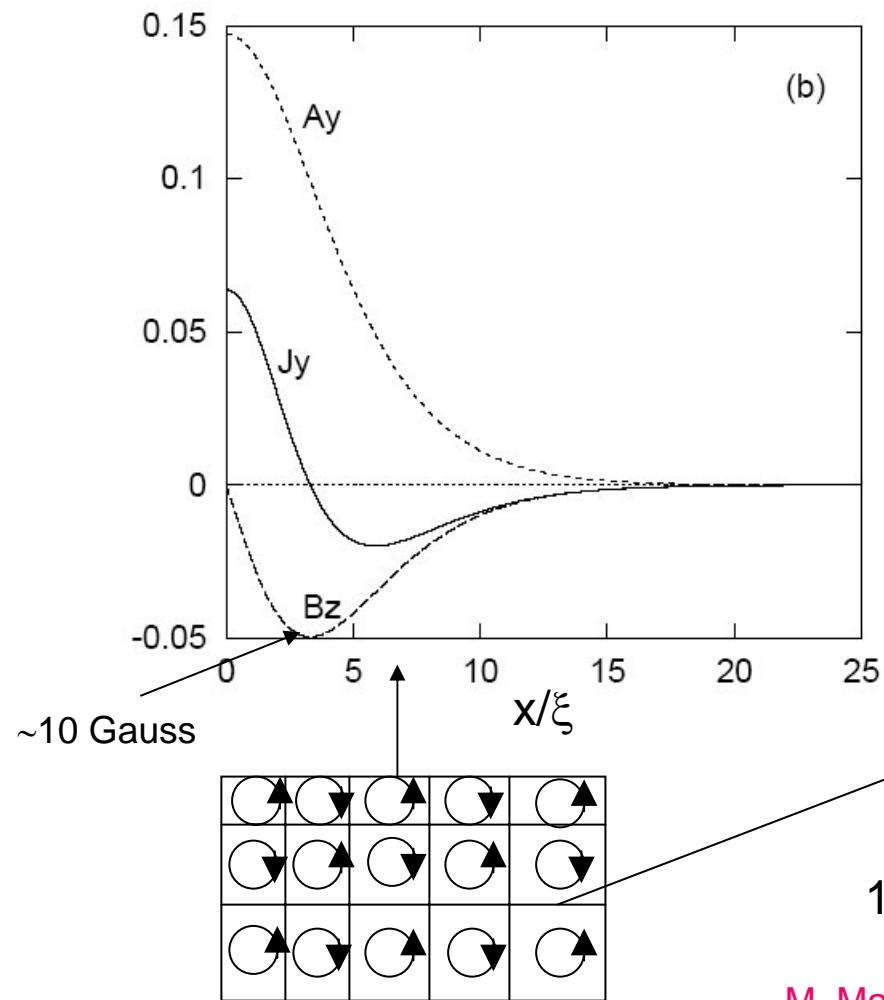
alternate view:

sub-gap surface states not at zero energy- carry current



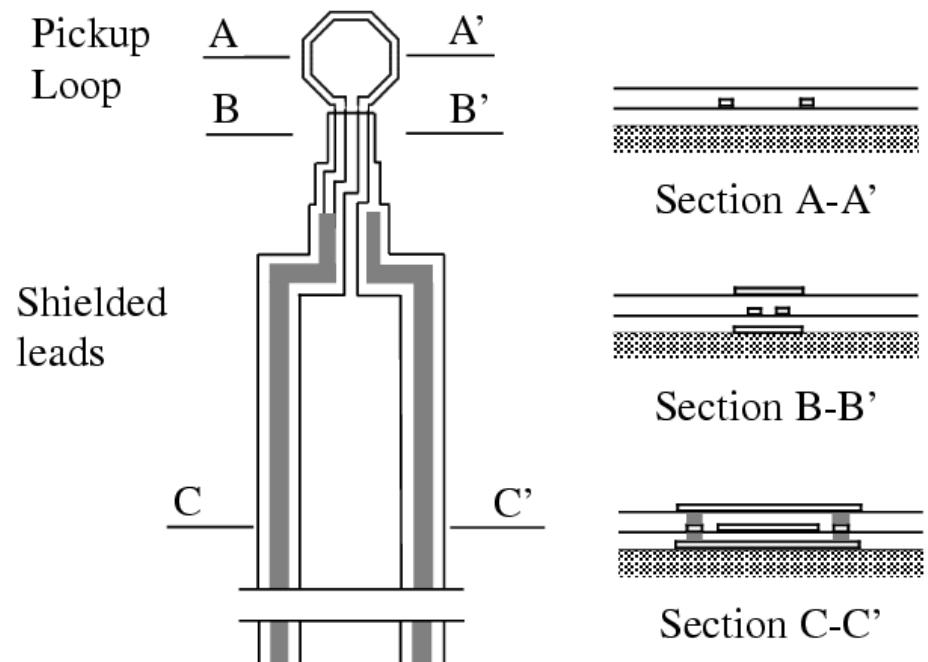
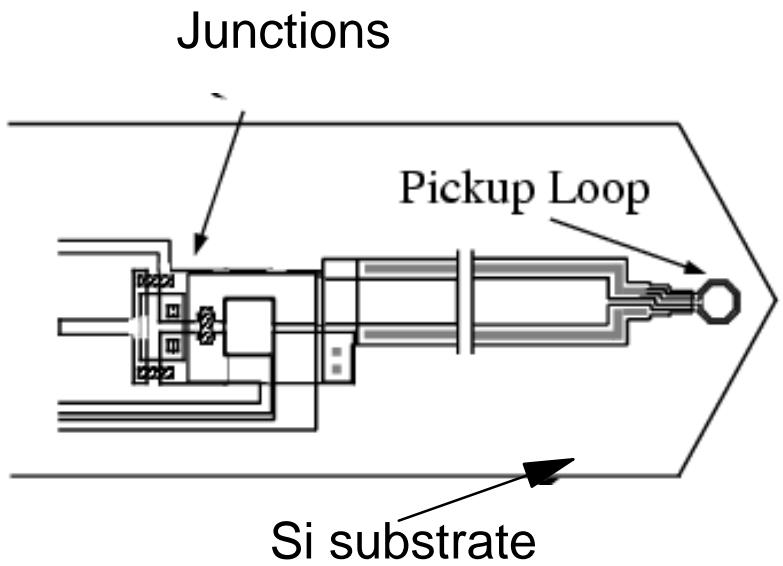
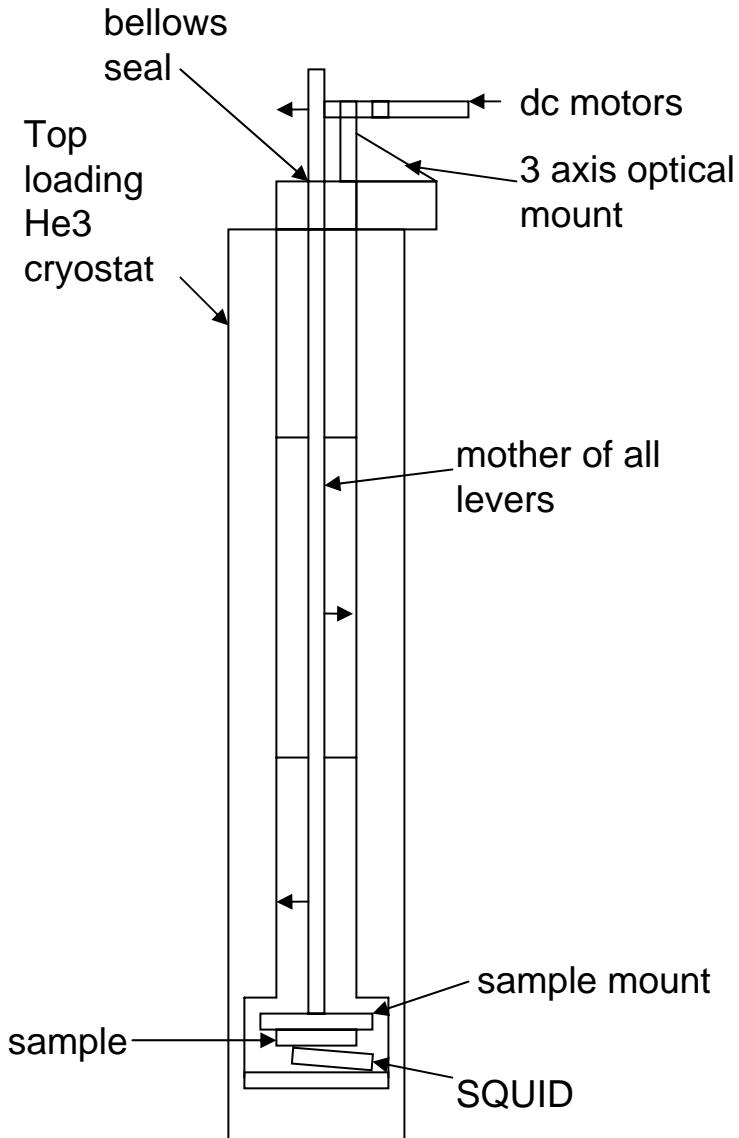
Spontaneous currents and fields at edge of $p_x \pm ip_y$ superconductor

Spontaneous currents and fields at domain boundaries of $p_x \pm ip_y$ superconductor



10 Gauss \times 8 μm pickup loop $\rightarrow 30\Phi_0$!

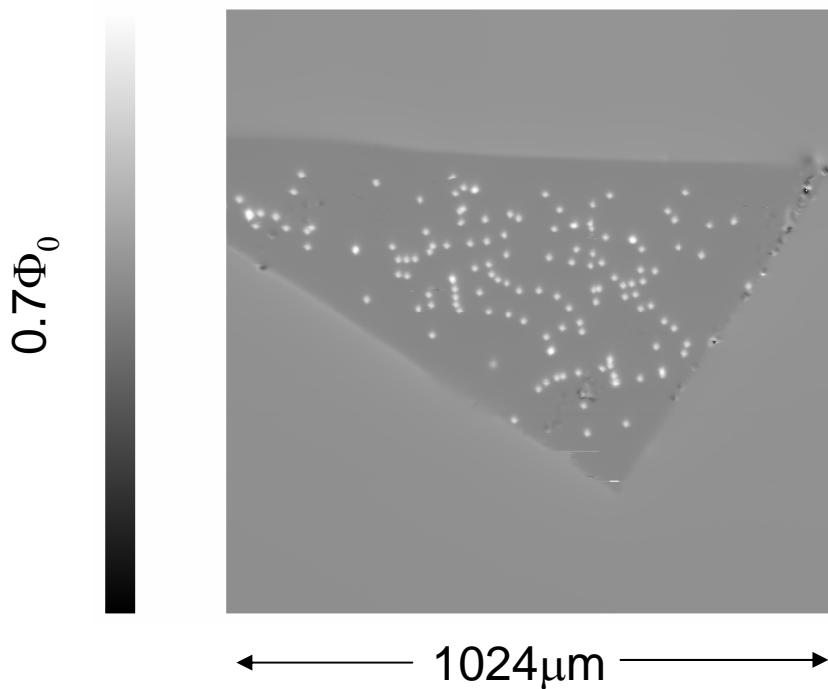
He3 Scanning SQUID microscope



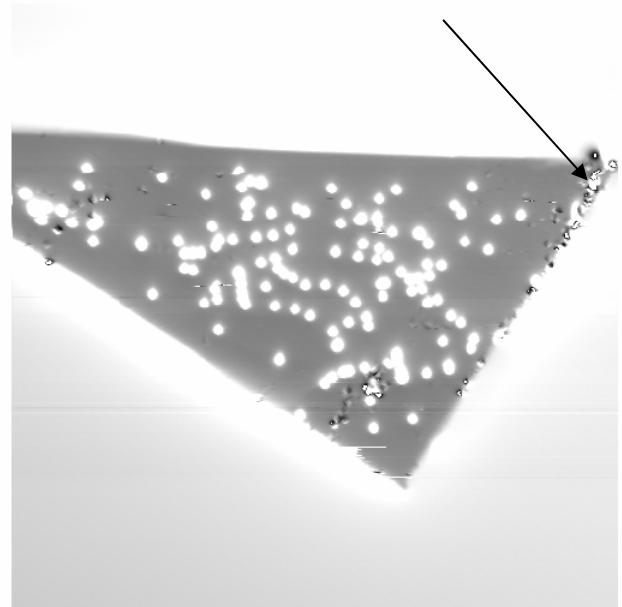
Scanning SQUID microscopy of "conventional" d-wave superconductor - YBCO

11.2 μm pickup loop

T=4.2K



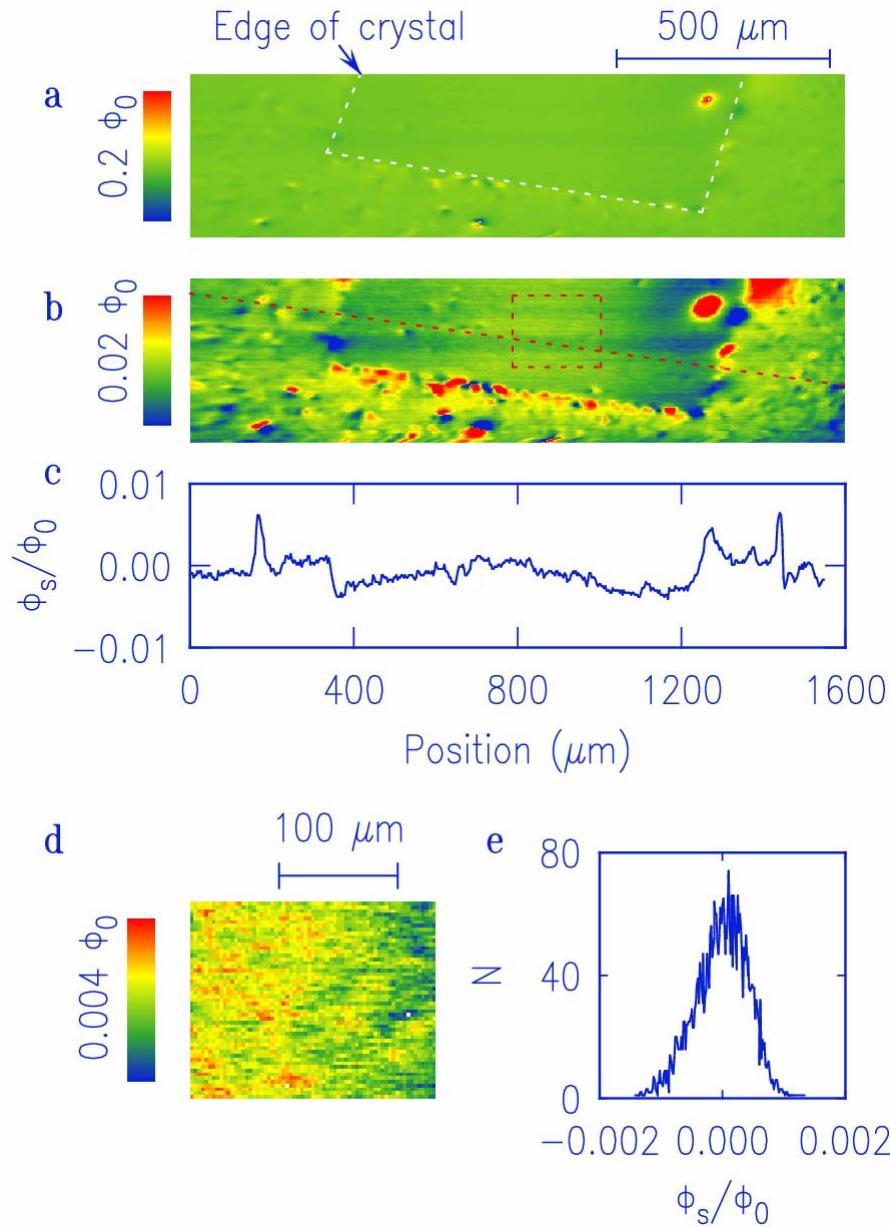
0.07 Φ_0



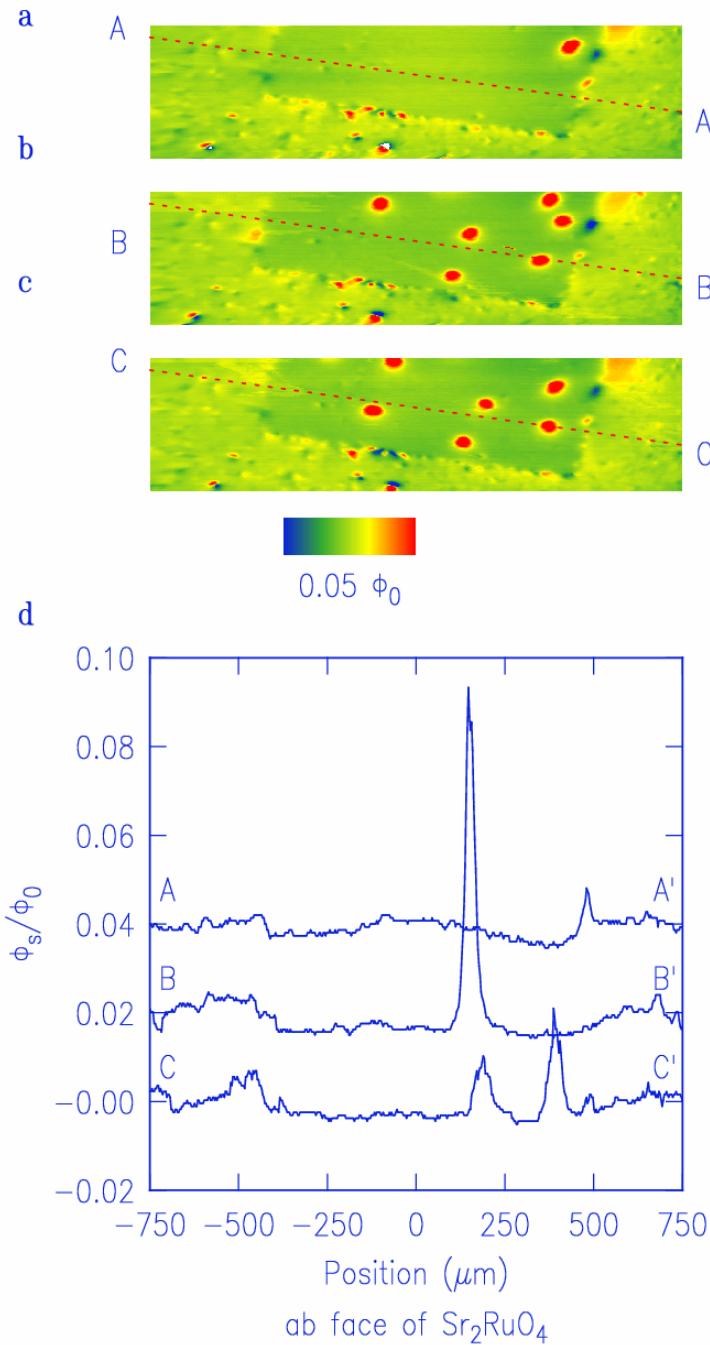
magnetic signals
at edges

K.A. Moler et al. PRB 55, 12753 (1997)

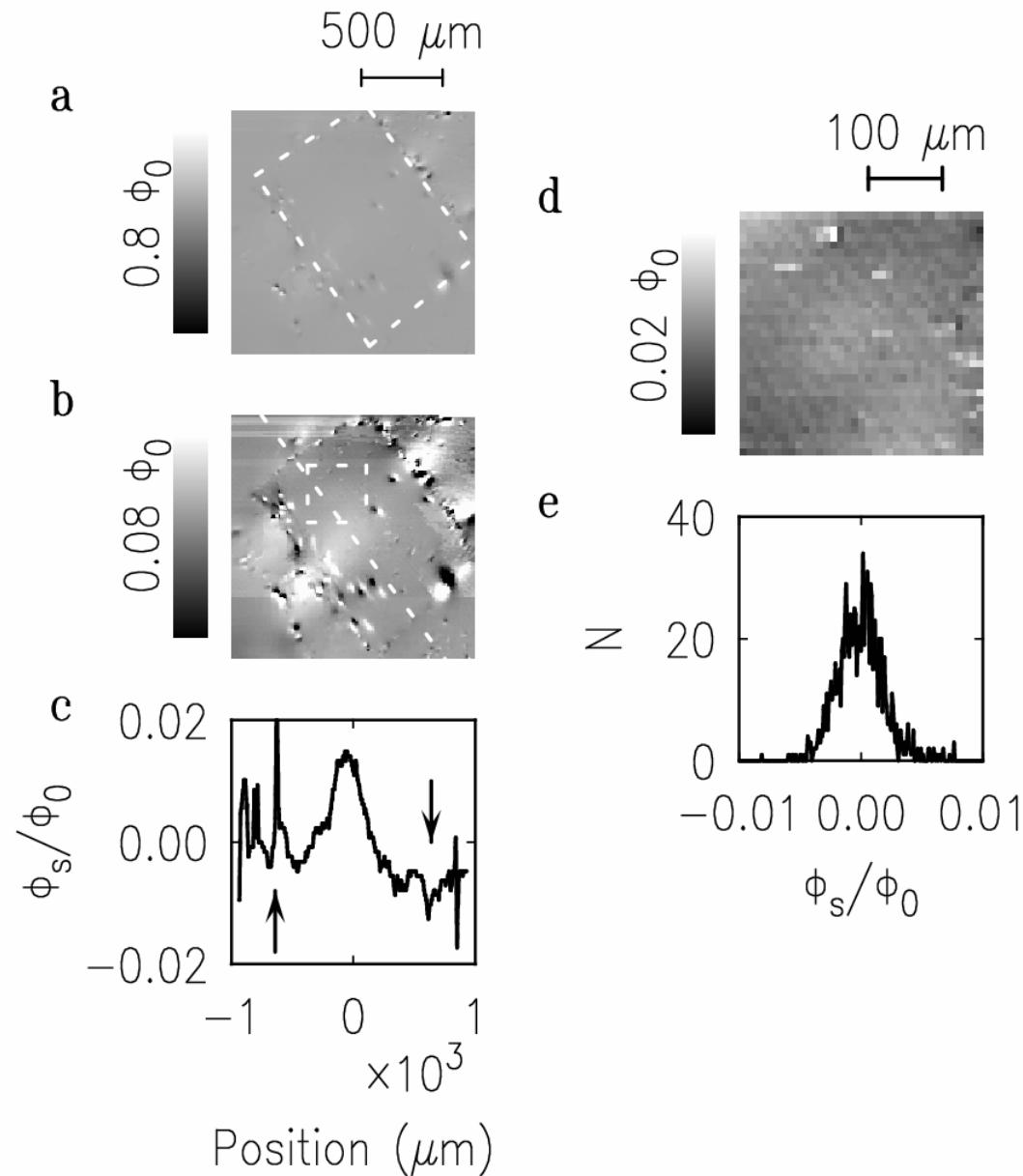
SQUID microscope image of ab face of Sr_2RuO_4 single crystal- $T=0.27\text{K}$
 $8\mu\text{m}$ pickup loop



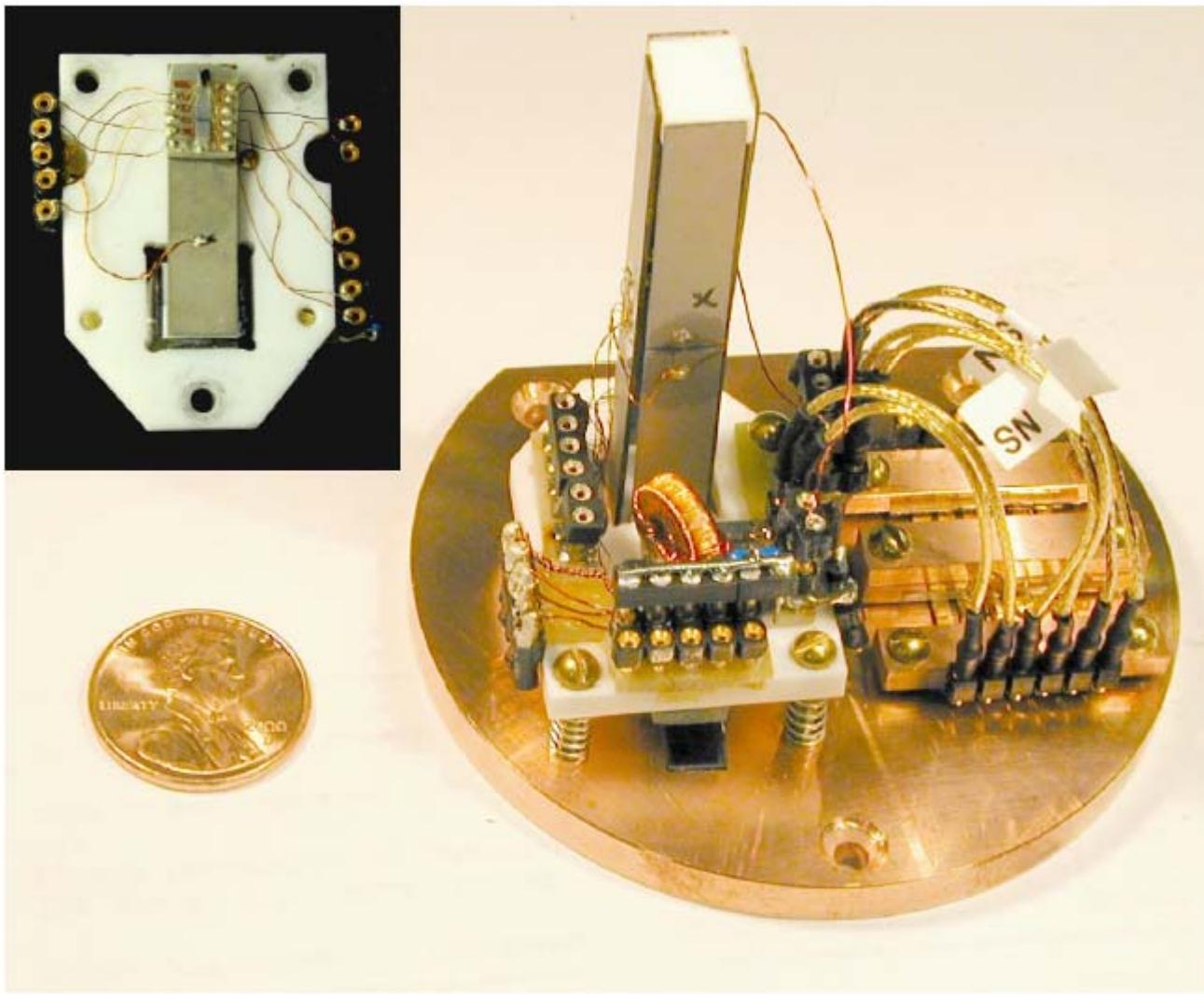
Three cooldowns of ab face in slightly difference fields



SSM image of ac face of Sr_2RuO_4

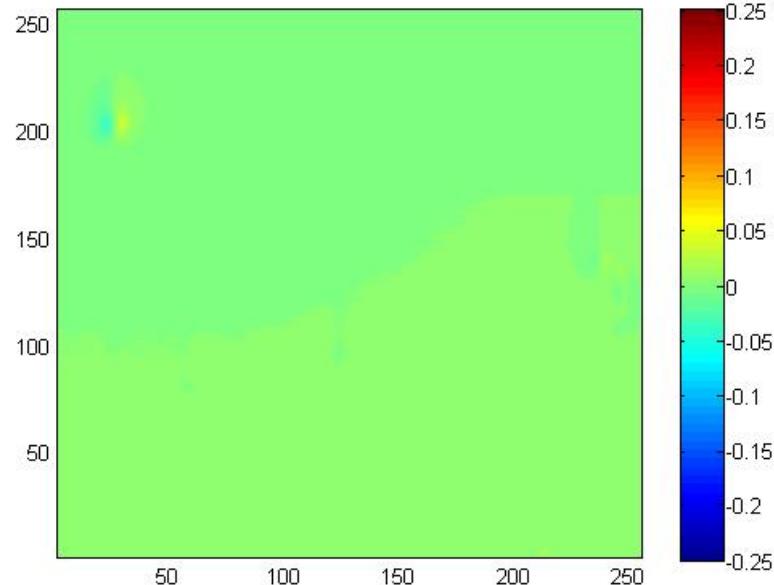


Stanford dilution refridgerator based scanning SQUID microscope

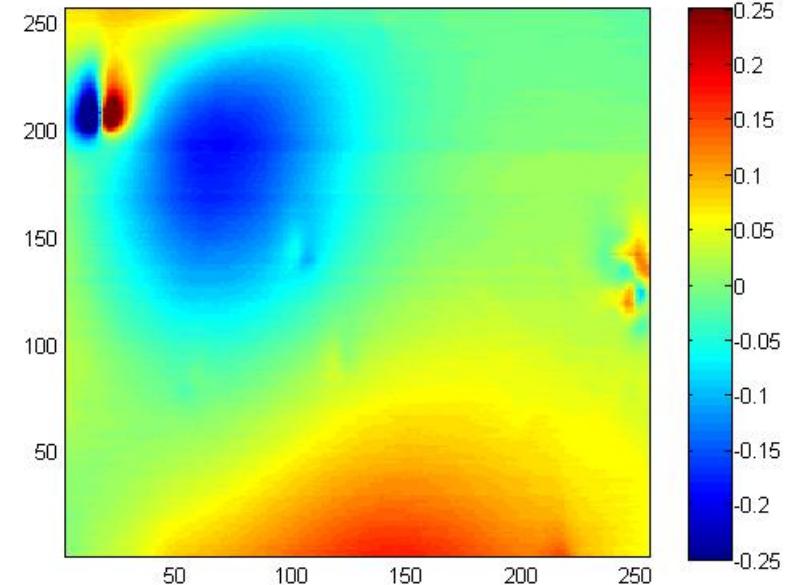


P. Bjornsson et al., RSI 72, 4153 (2001).

SQUID microscope image of
 Sr_2RuO_4 single crystal ac face

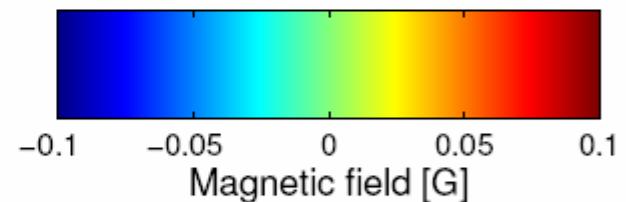
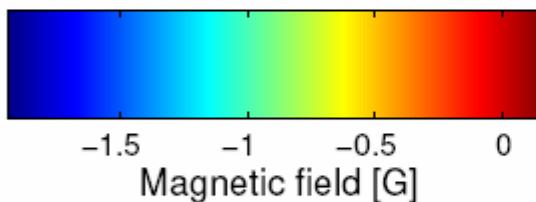
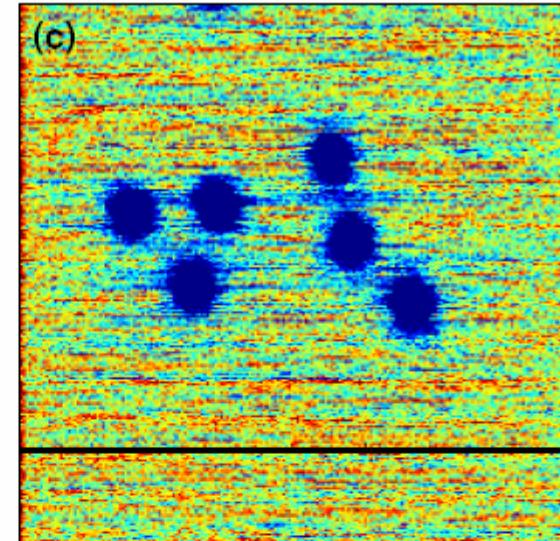
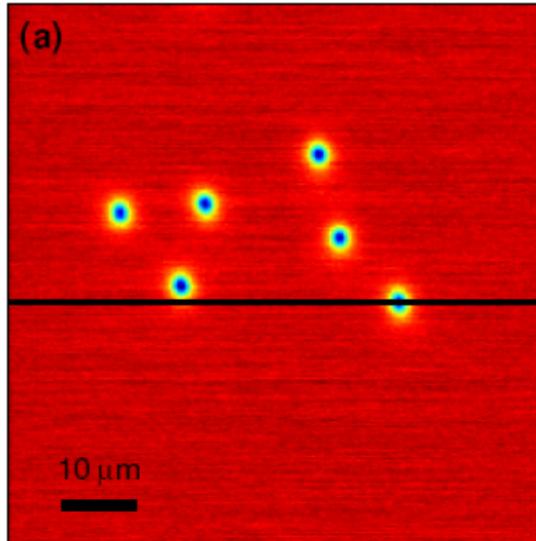
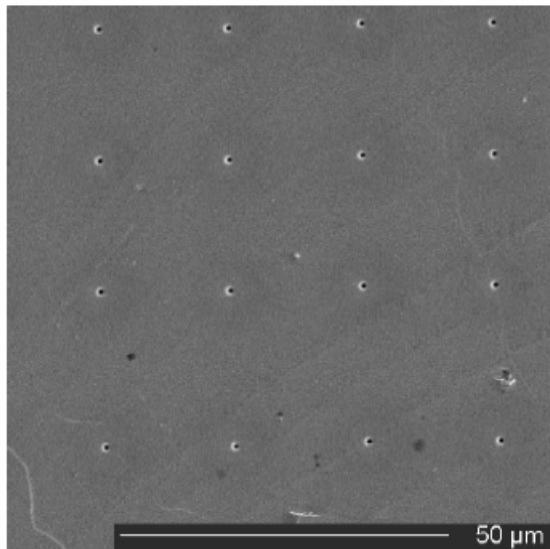


$T=4\text{K}$



$T=100\text{mK}$

Hall bar images of Sr_2RuO_4 single crystal -ab face

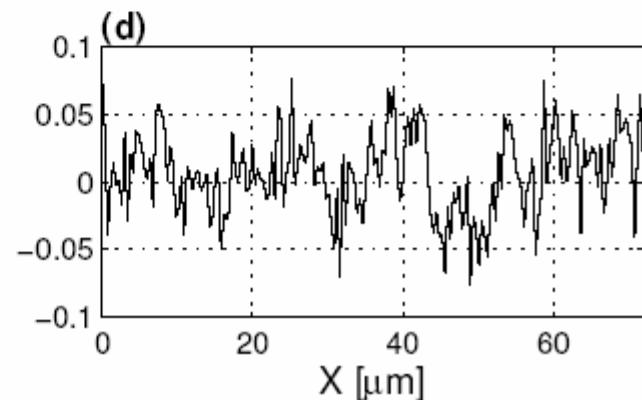
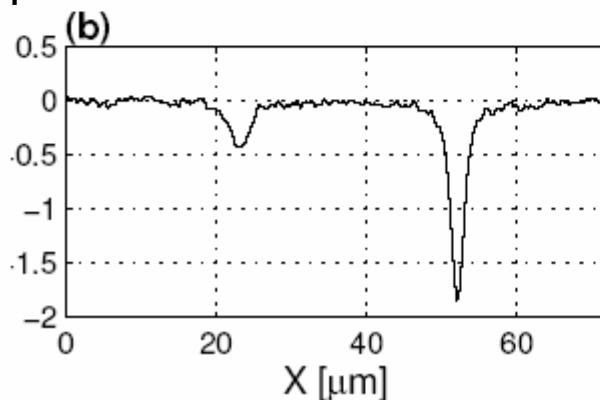


FIB etched holes $1\text{ }\mu\text{m}$ deep

$0.5\mu\text{m}$ Hall bar

$z=1.2\mu\text{m}$

$T=100\text{mK}$



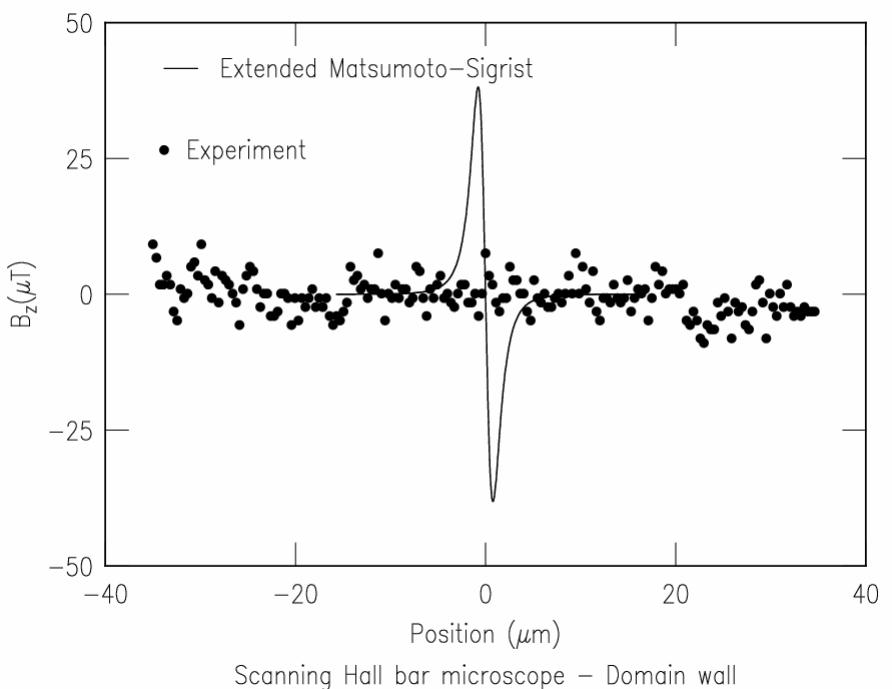
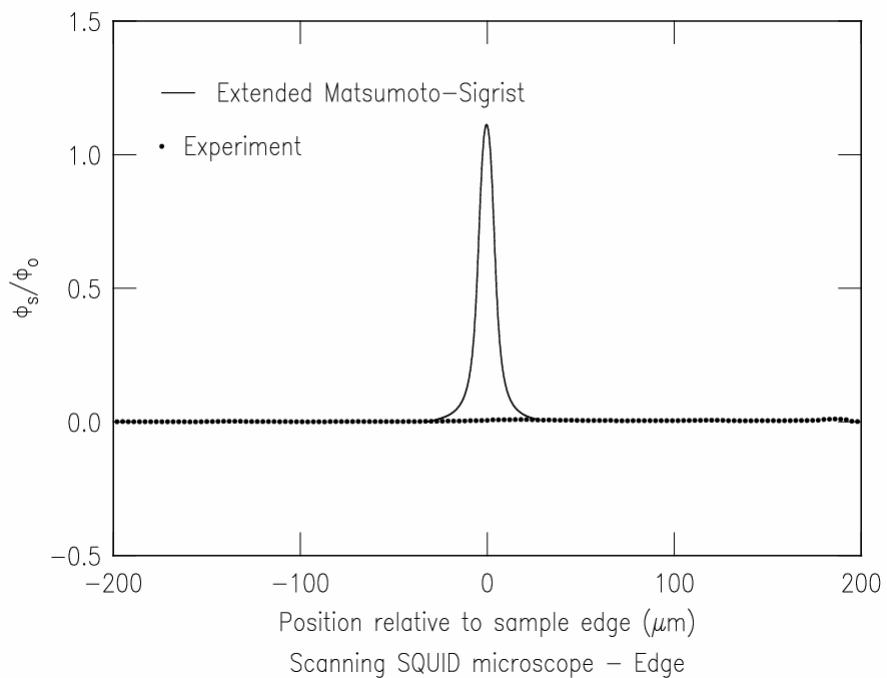
Predicted magnetic signal above $p_x \pm ip_y$ superconductor

(assuming B_z at surface same as B_z in bulk)

$$b_z(k_x, k_y, z) = b_z(k_x, k_y, 0) e^{-Kz} \quad K = (k_x^2 + k_y^2)^{1/2}$$

$$b_z(k_x, k_y, 0) = \frac{1}{(2\pi)^{1/2}} \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy e^{i(k_x x + k_y y)} B_x(x, y, 0)$$

$$\Phi = \int_{-L/2}^{L/2} dx \int_{-L/2}^{L/2} dy B_z(x, y, z)$$



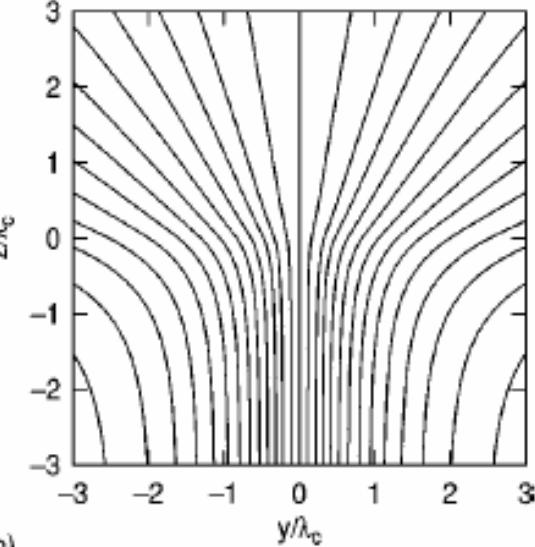
What about the effects of surface screening?

$$\tilde{B}_z(k, z) = \frac{K}{k + K} \tilde{B}_{0z}(k, 0) e^{-kz}$$

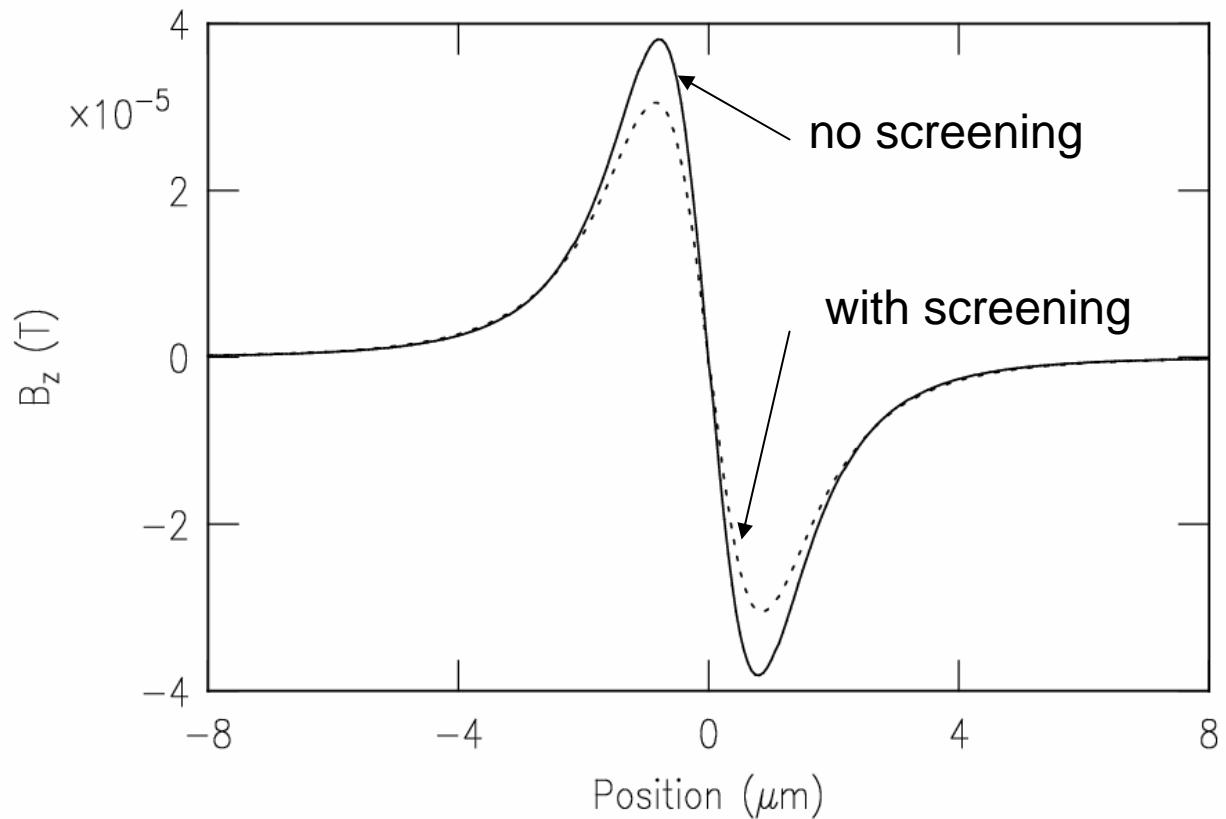
$\tilde{B}_z(k, z)$ - screened field
 $\tilde{B}_{0z}(k, 0)$ - bulk field

$$k = \sqrt{k_x^2 + k_y^2} \quad K = \sqrt{k^2 + 1/\lambda^2}$$

(a)

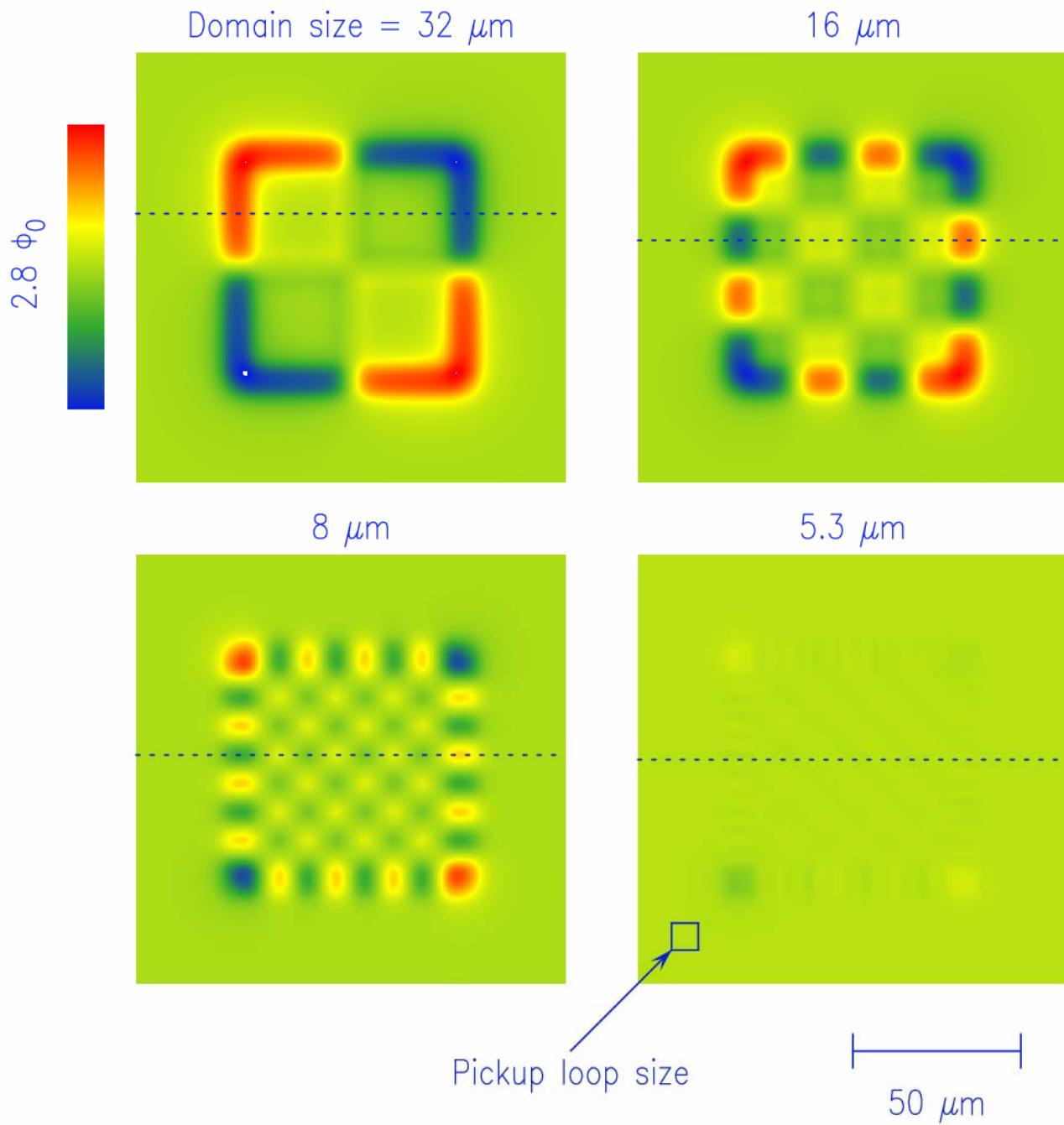


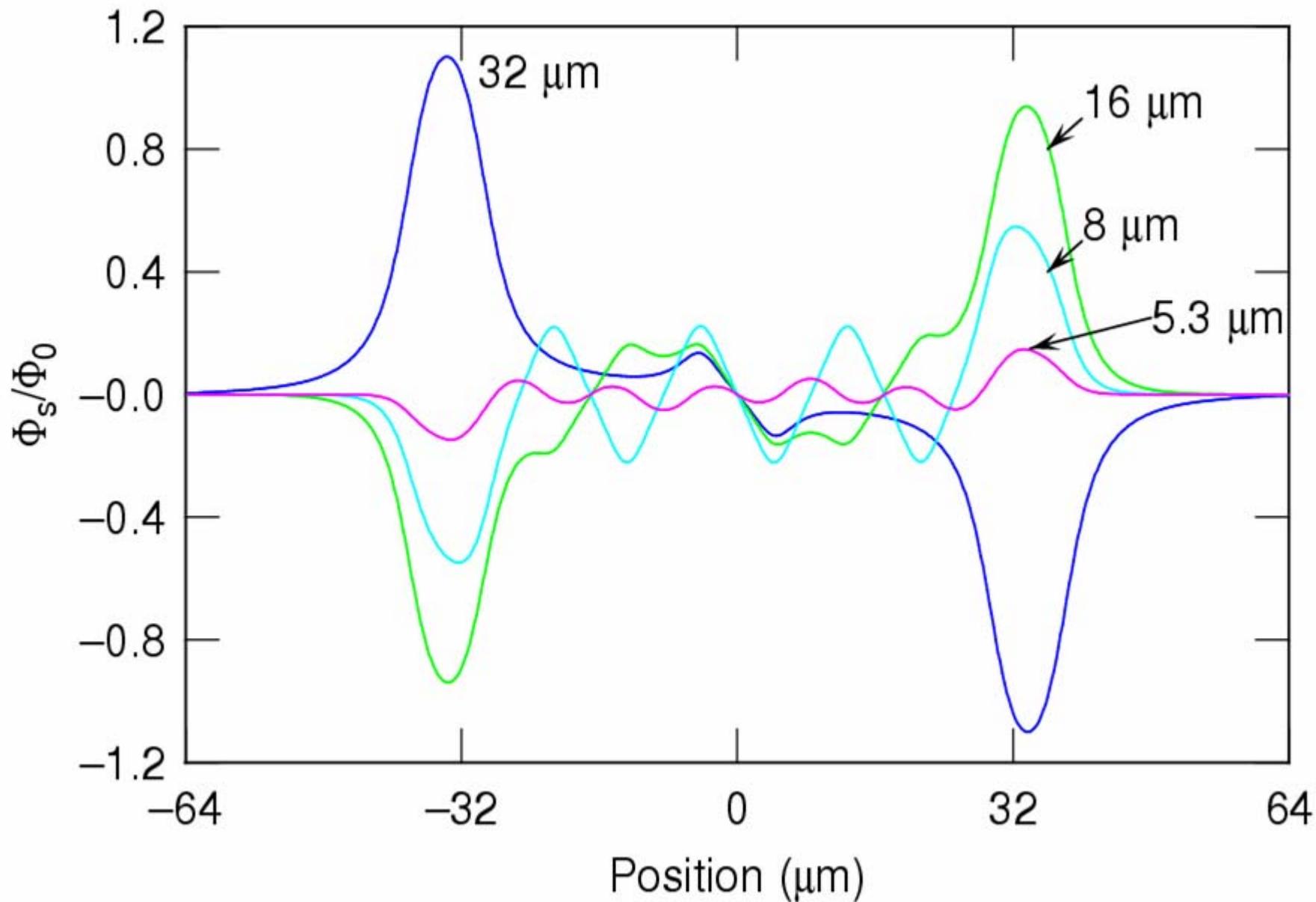
(b)

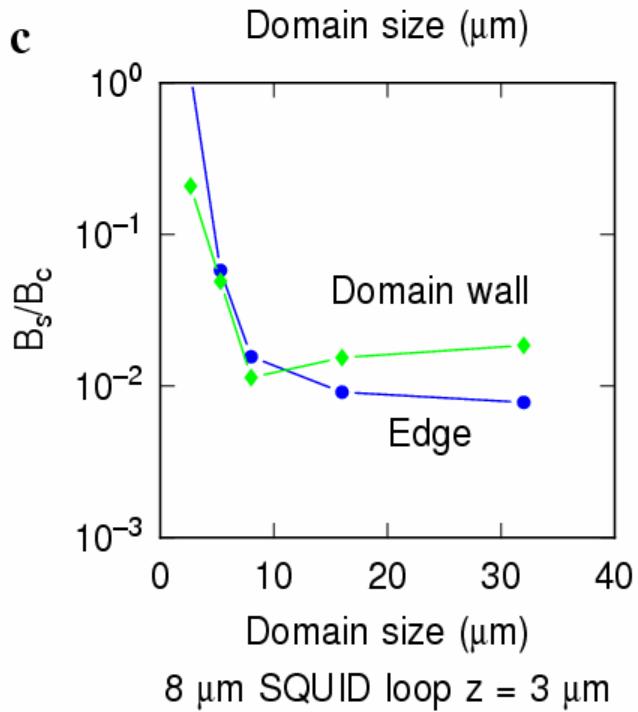
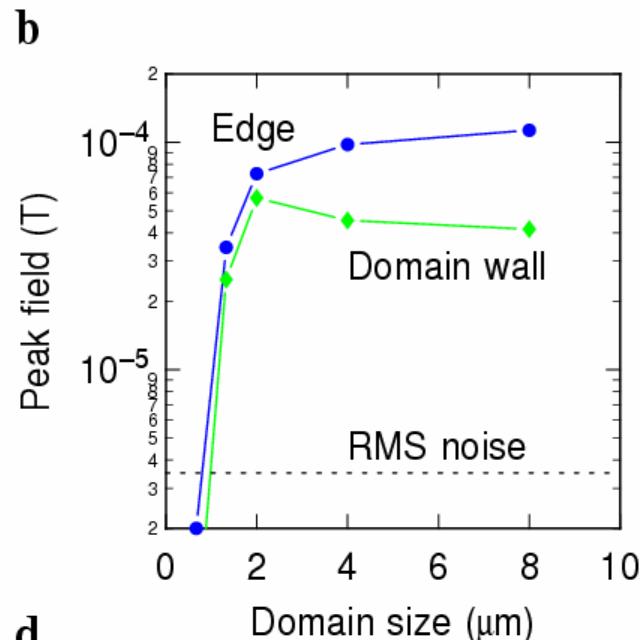
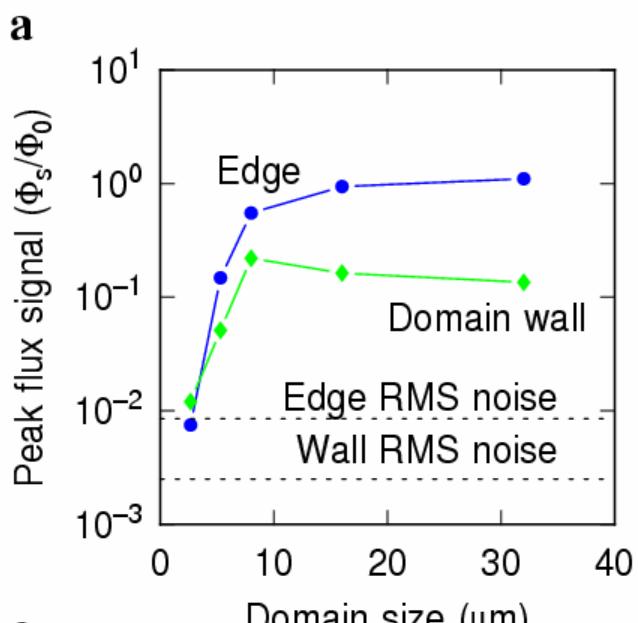


Kirtley et al., PRB 59, 4343 (1999)

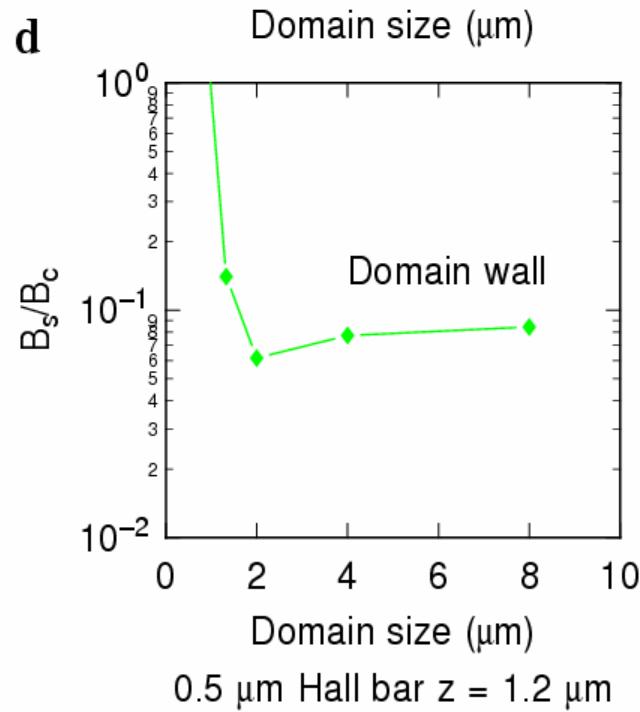
0.5 μm Hall bar 1.2 μm above the surface







8 μm SQUID loop $z = 3 \mu\text{m}$



0.5 μm Hall bar $z = 1.2 \mu\text{m}$

Conclusions:

1. Sharp, large spontaneous currents are not observed
2. Explanation A: Large spontaneous currents exist but the domain sizes are (conservatively) less than a few microns
to be compared with
 - 1 micron - van Harlingen - Illinois Urbana-Champaign
 - 50-100 microns - Kaptitulnik - Stanford
 - 1mm - Ying Liu - Penn State
3. Explanation B: Large spontaneous currents don't exist
 - a. Not (simple) $px \pm ipy$
 - b. ?