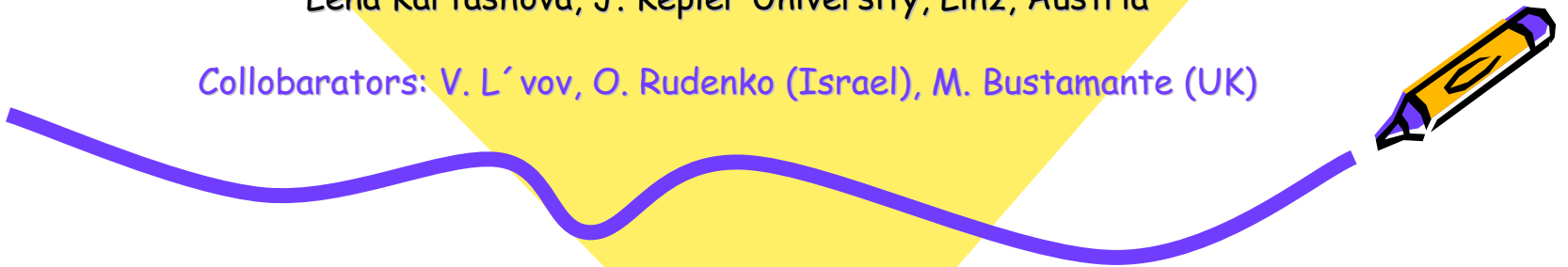




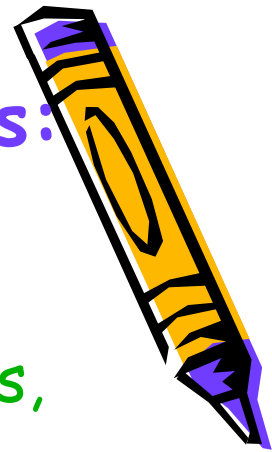
# Simple generic model of intra-seasonal oscillations in the Earth atmosphere

Lena Kartashova, J. Kepler University, Linz, Austria

Collobarators: V. L'vov, O. Rudenko (Israel), M. Bustamante (UK)



# Intra-seasonal oscillations - main facts:



- Periods: 10-100 days
- Location: everywhere (SH, NH; mid-latitudes, tropical zone)
- Data: AAM (atmospheric angular momentum, zonal winds, )
- Reason (NH, m.-l. - topography? SH, tropics - ?...)
- Modeling (no general model, particular numerical models, e.g. m.-l. IOs - reproduced with and without topography)
- Open questions - many (later)



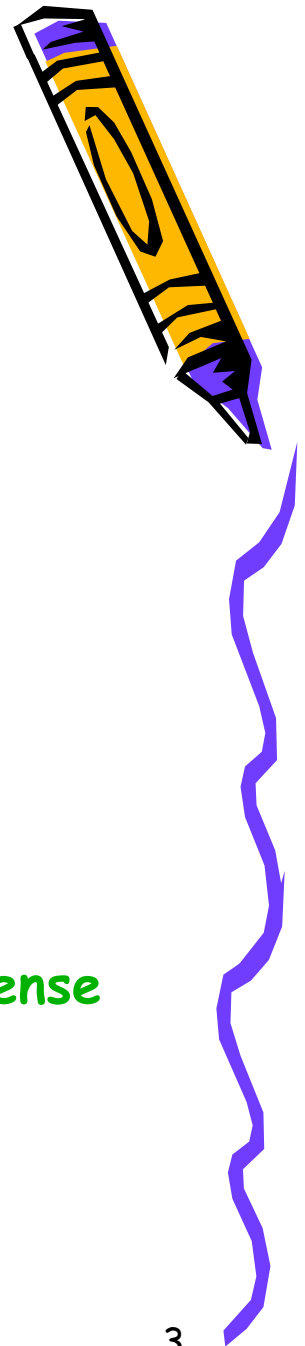
# Nonlinear resonance theory (NRT)

$$\begin{cases} \omega(\mathbf{k}_1) + \omega(\mathbf{k}_2) - \omega(\mathbf{k}_3) = \Omega, \\ \mathbf{k}_1 + \mathbf{k}_2 - \mathbf{k}_3 = 0, \end{cases} \quad \Omega \geq 0$$

- **Not all** modes interact resonantly
- Resonant modes form small **dynamically independent** clusters
- Form of clusters, if any, is defined by **boundary conditions**
- Resonant interactions are **not local** in Kolmogorov's sense
- **Quasi-resonances** possess similar properties



EK: **Amer.Math.Soc.** (1998); **PRL** (1994,2007)



# Planetary waves on a sphere

$$\partial \Delta \psi / \partial t + 2 \partial \psi / \partial \lambda = -J(\psi, \Delta \psi)$$

$$\omega_j = -2m_j / l_j (l_j + 1)$$

$$\begin{cases} N_1 \dot{A}_1 = 2Z(N_3 - N_2)A_3 A_2^*, \\ N_2 \dot{A}_2 = 2Z(N_1 - N_3)A_1^* A_3, \\ N_3 \dot{A}_3 = 2Z(N_2 - N_1)A_1 A_2, \end{cases}$$

All clusters in  $m, l \leq 21$

4 triads,

3 butterflies,

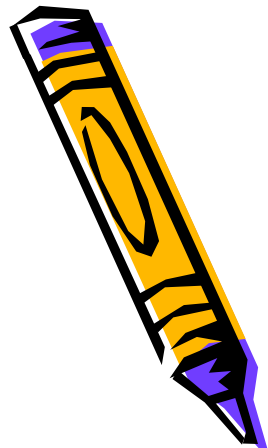
1 cluster of 6 connected triads

## Maximum Z at „interaction latitude“:

We can interpret the latitude  $\varphi_0$  as follows. The overlap of three wave functions in a triad  $Z(\lambda, \varphi)_{\ell_1}^{m_1} \times (\lambda, \varphi) Y_{\ell_2}^{m_2}(\lambda, \varphi) Y_{\ell_3}^{m_3}(\lambda, \varphi)$  shows a contribution to the interaction coefficient  $Z \propto \int Z(\lambda, \varphi) d\lambda d\varphi$  from a particular location on the sphere. The overlap  $Z(\lambda, \varphi)$  has a maximum at a particular latitude  $\varphi_0$ , and a narrow latitudinal belt around  $\varphi_0$  gives the main contribution to the global interaction amplitude  $Z$ . That is why  $\varphi_0$  can be understood as the interaction latitude.

EK, L.Piterbarg, G.Reznik, *Oceanology* (1990)

EK, V.L'vov, *PRL* (2007);

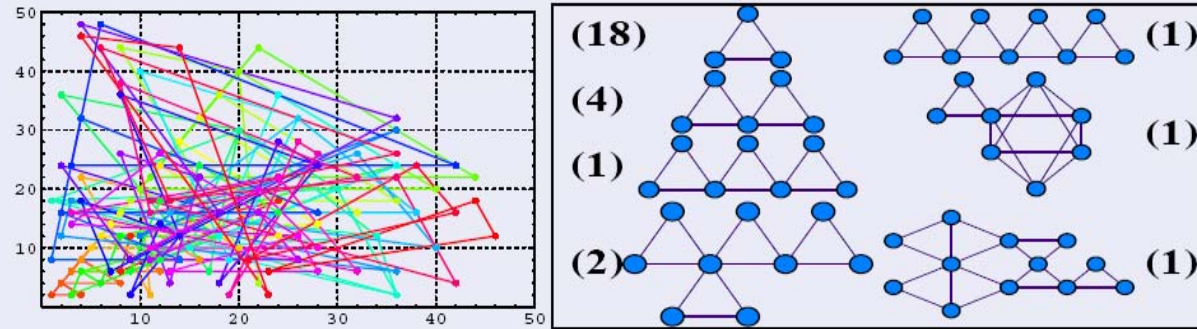


# Main idea: IOs are due to **resonance clusters** of spherical Rossby waves

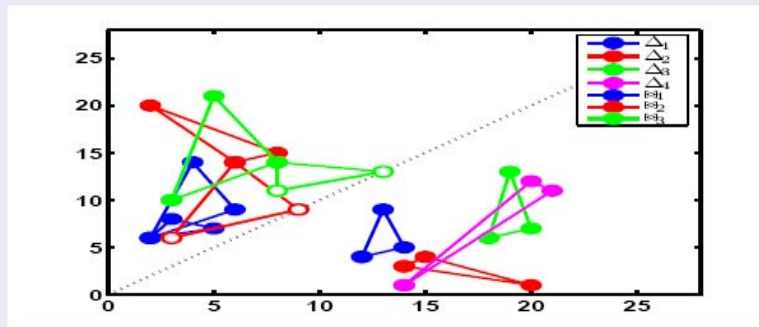
Galerkin:  
2500 Fourier  
harmonics  
( $m, n \leq 50$ )

Clipping:  
Only 128  
are important

Structure of resonances in spectral space  $\leq 50$



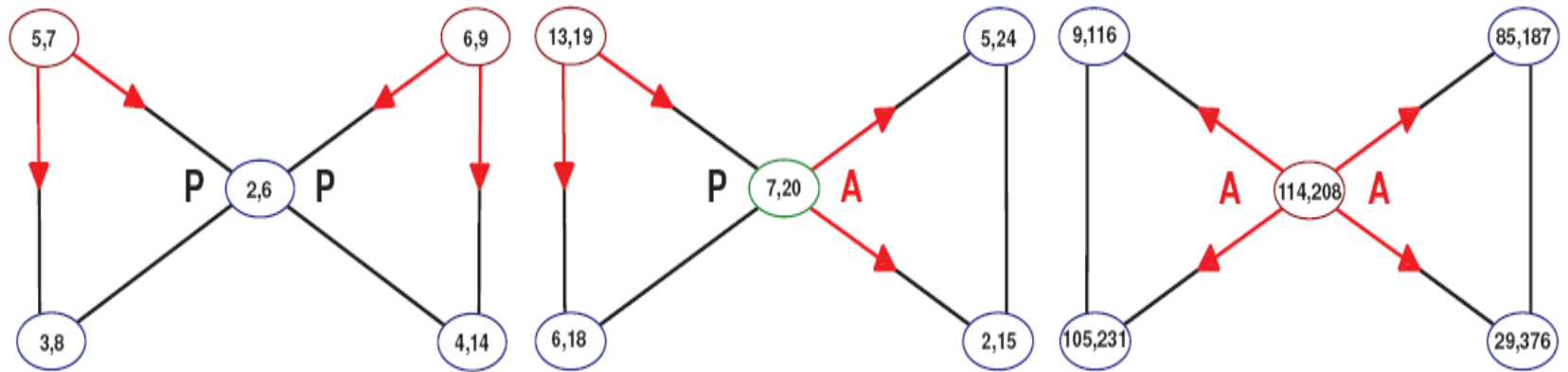
Spectral domain  $\leq 21$



and 1 group of 6 connected triads.

# Connection types within a cluster

$$\omega_j = -2m_j/l_j(l_j + 1)$$



$$\frac{4}{14 \cdot 15} + \frac{2}{6 \cdot 7} = \frac{6}{9 \cdot 10}$$

$$\frac{4}{8 \cdot 9} + \frac{2}{6 \cdot 7} = \frac{5}{7 \cdot 8}$$

$$\frac{13}{19 \cdot 20} = \frac{7}{20 \cdot 21} + \frac{6}{18 \cdot 19}$$

$$\frac{7}{20 \cdot 21} = \frac{5}{24 \cdot 25} + \frac{2}{15 \cdot 16}$$

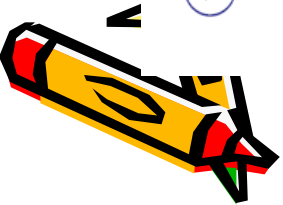
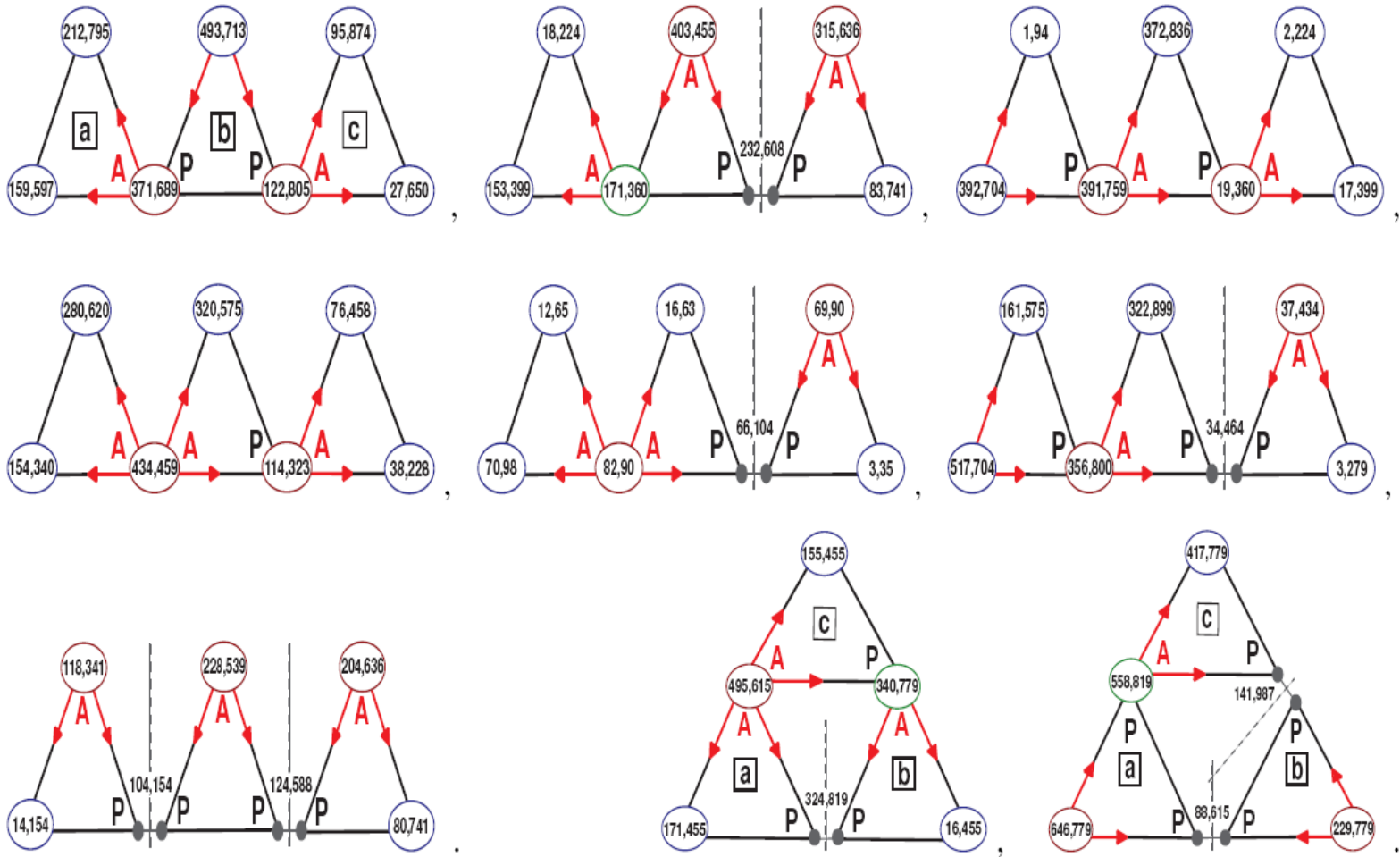
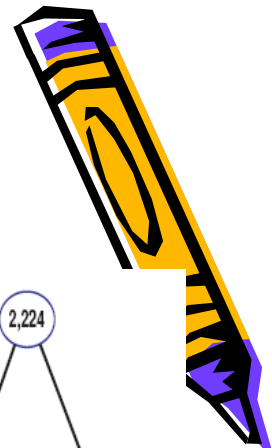
$$\frac{114}{208 \cdot 209} = \frac{9}{116 \cdot 117} + \frac{105}{231 \cdot 232}$$

$$\frac{114}{208 \cdot 209} = \frac{85}{187 \cdot 188} + \frac{29}{376 \cdot 377}$$

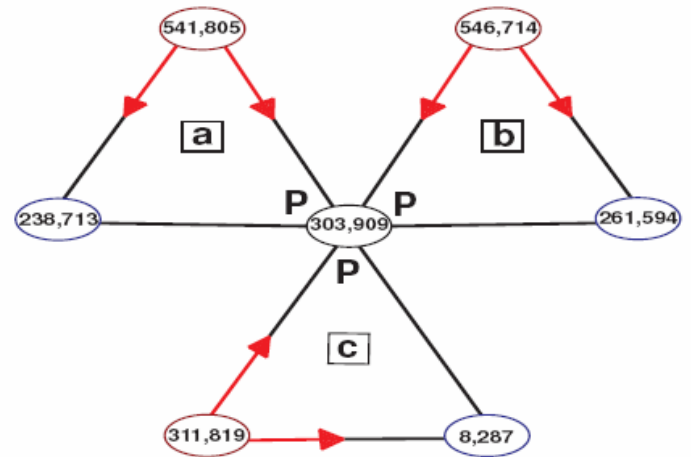
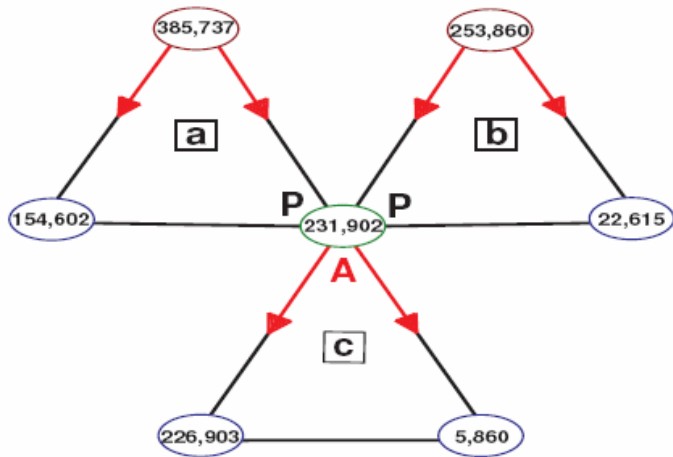
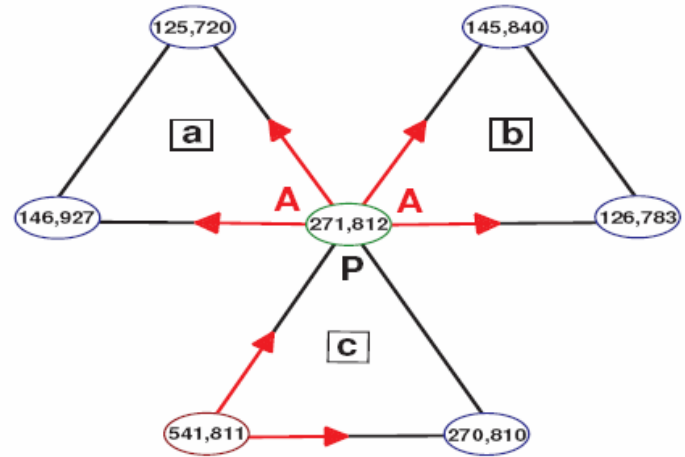
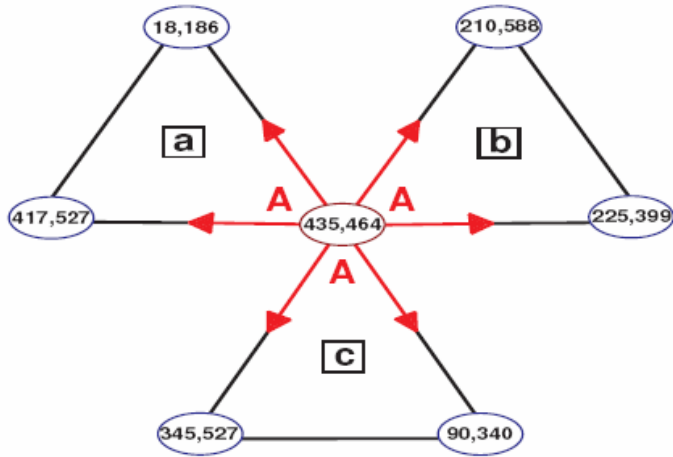
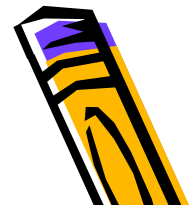


EK, V.L'vov, **EPL** (2008)

# Connection types (continuation)

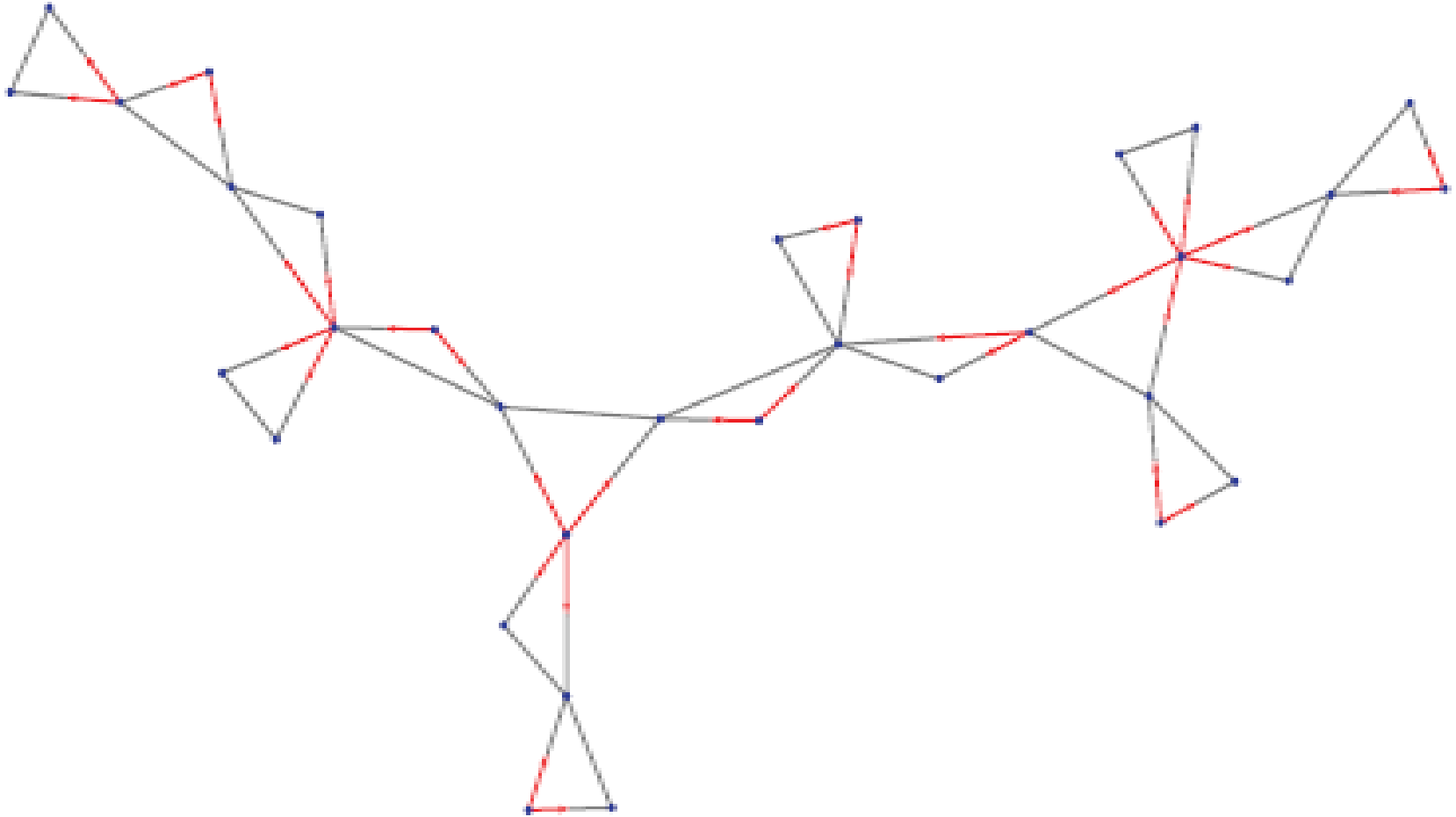
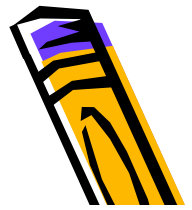


# Connection types (continuation)

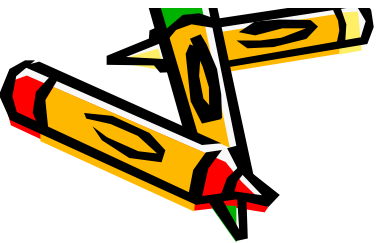




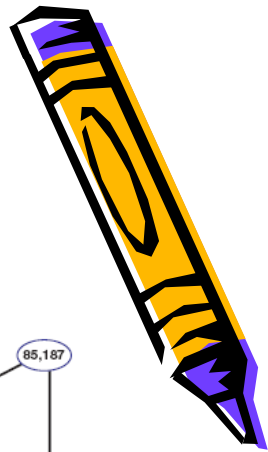
# Connection types (continuation)



16-triad's PP-reducible cluster



# Examples of dynamical systems

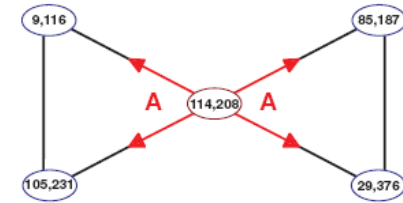


1.

$$\begin{cases} \dot{B}_{1a} = Z_a B_{2a}^* B_{3a}, & \dot{B}_{1b} = -Z_b B_{2b}^* B_{3a}, \\ \dot{B}_{2a} = Z_a B_{1a}^* B_{3a}, & \dot{B}_{2b} = Z_b B_{1b}^* B_{3a}, \\ \dot{B}_{3a} = -Z_a B_{1a} B_{2a} - Z_b B_{1b} B_{2b}. \end{cases}$$

$$\begin{cases} I_{12a} = |B_{1a}|^2 - |B_{2a}|^2, & I_{12b} = |B_{1b}|^2 - |B_{2b}|^2, \\ I_{ab} = |B_{1a}|^2 + |B_{3a}|^2 + |B_{3b}|^2. \end{cases}$$

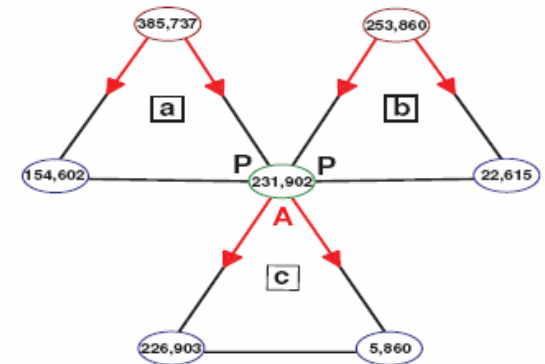
AA-butterfly



2.

$$\begin{cases} \dot{B}_{3a} = -Z_a B_{1a} B_{2a}, & \dot{B}_{2a} = Z_a B_{1a} B_{3a}, \\ \dot{B}_{3b} = -Z_b B_{1a} B_{2b}, & \dot{B}_{2b} = Z_a B_{1b} B_{3a}, \\ \dot{B}_{1c} = Z_c B_{2c}^* B_{1a}, & \dot{B}_{2c} = Z_c B_{1a} B_{3a}, \\ \dot{B}_{1a} = Z_a B_{2a}^* B_{3a} + Z_b B_{2b}^* B_{3b} - Z_c B_{1c} B_{2c}. \end{cases}$$

PPA-star

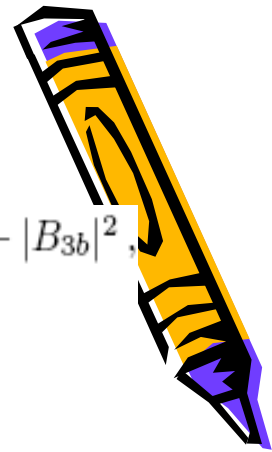


$$\begin{cases} \text{PPA-star:} & I_{12a} = |B_{1a}|^2 - |B_{2a}|^2, \\ I_{23b} = |B_{2b}|^2 + |B_{3b}|^2, & I_{23c} = |B_{2c}|^2 + |B_{3c}|^2, \\ I_{abc} = |B_{1a}|^2 + |B_{3a}|^2 + |B_{3b}|^2 + |B_{3c}|^2. \end{cases}$$

CLs constructed as combinations of known CLs for a triad



# New Conservation Laws, PP-butterfly :



$$\begin{cases} \dot{B}_{1a} = Z_a B_{2a}^* B_{3a} + Z_b B_{2b}^* B_{3b}, & B_{1a} = B_{1b}, \\ \dot{B}_{2a} = Z_a B_{1a}^* B_{3a}, & \dot{B}_{2b} = Z_b B_{1a}^* B_{3b}, \\ \dot{B}_{3a} = -Z_a B_{1a} B_{2a}, & \dot{B}_{3b} = -Z_b B_{1a} B_{2b}. \end{cases} \quad \begin{cases} I_{23a} = |B_{2a}|^2 + |B_{3a}|^2, & I_{23b} = |B_{2b}|^2 + |B_{3b}|^2, \\ I_{ab} = |B_{1a}|^2 + |B_{3a}|^2 + |B_{3b}|^2, \\ I_0 = \text{Im}(Z_a B_1 B_{2a} B_{3a}^* + Z_b B_1 B_{2b} B_{3b}^*), \end{cases}$$

Reduction: from 10 equations to 4:

$$\frac{dC_{3a}}{dt} = -Z_a C_1 C_{2a} \cos \varphi_a,$$

$$\frac{dC_{3b}}{dt} = -Z_b C_1 C_{2b} \cos \varphi_b,$$

$$\frac{d\varphi_a}{dt} = Z_a C_1 \left( \frac{C_{2a}}{C_{3a}} - \frac{C_{3a}}{C_{2a}} \right) \sin \varphi_a - \frac{I_0}{(C_1)^2},$$

$$\frac{d\varphi_b}{dt} = Z_b C_1 \left( \frac{C_{2b}}{C_{3b}} - \frac{C_{3b}}{C_{2b}} \right) \sin \varphi_b - \frac{I_0}{(C_1)^2}.$$

$$B_j = C_j \exp(i\theta_j)$$

$$\varphi_a = \theta_{1a} + \theta_{2a} - \theta_{3a}$$

$$\varphi_b = \theta_{1b} + \theta_{2b} - \theta_{3b};$$

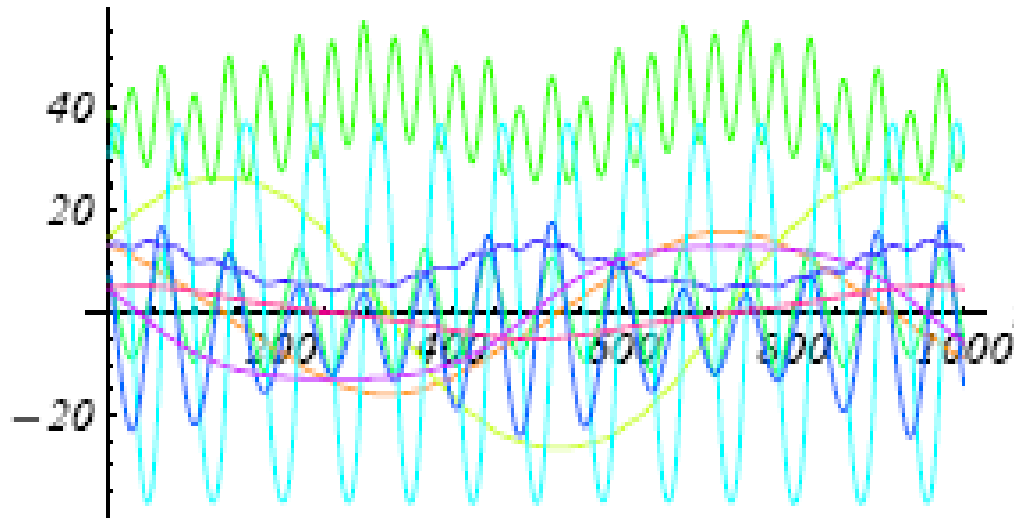
Integrable cases of butterflies:

1. Real amplitudes; 2.  $I_0 = 0$  ; 3.  $Z_a = Z_b$  ; 4. Some more



# 4-triad's chain, real amplitudes

$10^5 \times B_i, i=1, \dots, 9$



Amplitudes

$ba_1$	Orange
$ba_2$	Yellow-green
$ba_3$	Green
$ba_4$	Light green
$ba_5$	Cyan
$ba_6$	Blue
$ba_7$	Dark blue
$ba_8$	Purple
$ba_9$	Magenta



# Open questions (M. Ghil, 2004)



- Cause of IOs in the SH? *The same*
- NH oscillation: why by some researchers - 40 days, by others - 20-30 days. *Period is inversely proportional to the energy, more energy -> shorter period*
- Why IOs are better observable in winter data? *For larger energies appr. resonances may destroy the clusters*
- How do tropical and mid-latitude IOs interact? *Two mechanisms are possible:*
  - 1) *One cluster, different latitudinal belts;*
  - 2) *Via mode which is quasi-resonant for one triad and exact for another*
- How to predict IOs in data? *Obvious (spher. Harmonics with known wavenumbers)*



# Energy flux between latitudinal belts

Triad	Modes $[m, \ell]$	$Z$	$\varphi_0$	$K[\mu]$	$10^6 E_0$	$T_0$
$\Delta_1$	[4,12] [5,14] [9,13]	7.82	34	1.62	14.4	24
$\Delta_2$	[3,14] [1,20] [4,15]	37.46	19	1.14	5.4	5
$\Delta_3$	[6,18] [7,20] [13,19]	13.66	34	1.74	32.0	19
$\Delta_4$	[1,14][11,21][12,20]	47.67	28	1.21	0.58	13
$\boxtimes_1$	[2,6] [3,8] [5,7]	3.14	35	1.64	5.08	30
	[2,6] [4,14] [6,9]	14.63	37	1.61	0.395	10
$\boxtimes_2$	[6,14] [2,20] [8,15]	69.25	31	1.13	0.61	8
	[3,6] [6,14] [9,9]	11.31	—	1.17	0.360	13
$\boxtimes_3$	[3,10] [5,21] [8,14]	61.99	31	1.27	0.133	7
	[8,11] [5,21] [13,13]	8.71	—	1.36	0.784	24
$\boxtimes$	[1,6] [2,14] [3,9]	28.98	17	1.38	0.247	6
	[2,7] [11,20] [13,14]	2.77	42	1.08	1.78	26
	[1,6] [11,20] [12,15]	15.08	29	1.06	0.262	11
	[9,14] [3,20] [12,15]	74.93	50	1.36	0.487	8
	[3,9] [8,20] [11,14]	32.12	40	1.11	0.251	9
	[2,14][17,20][19,19]	11.05	—	1.05	3.33	24

the mode (13,19) is near-resonant for  $\Delta_4$  (with resonance discrepancy  $\delta = 0.16$ ) and is resonant for  $\Delta_3$ .

Thank you!

