

Kaluza-Klein Dark Matter

Jonathan Feng

UC Irvine

KITP, Santa Barbara

5 December 2002

Dark Matter

- Our best evidence for new particle physics
- We live in interesting times
 - we know how much there is ($\Omega_{\text{DM}} \sim 0.3$)
 - but not what it is (non-baryonic, cold)
- WIMPs are attractive
 - predicted in many particle theories (EWSB)
 - naturally give thermal relic density $\Omega_{\text{DM}} \sim \text{O}(1)$
 - $\Omega_{\text{DM}} < 1 \Rightarrow \chi\chi \rightarrow ff$ not small $\Rightarrow \chi f \rightarrow \chi f$ not small, so testable: promising for direct, indirect detection

Candidates from Particle Physics

- Supersymmetry
 - Neutralinos – partners of γ , Z , W , h
 - Requirements:
 - high supersymmetry-breaking scale (supergravity)
 - R -parity conservation
- Extra Dimensions
 - Kaluza-Klein particles – partners of γ , Z , W , h , $G_{\mu\nu}, \dots$
 - Requirements:
 - universal extra dimensions

Cheng, Feng, Matchev (2002)
Feng, Rajaraman, Takayama

Universal Extra Dimensions



- Kaluza (1921) and Klein (1926) considered $D=5$, with 5th dimension compactified on circle S^1 of radius R :

$D=5$ gravity \rightarrow $D=4$ gravity + EM + scalar

$$G_{MN} \rightarrow G_{\mu\nu} + G_{\mu 5} + G_{55}$$



- Kaluza: “virtually unsurpassed formal unity...which could not amount to the mere alluring play of a capricious accident.”

- Problem: gravity is weak
- Solution: introduce extra 5D fields: G_{MN} , V_M , etc.
- New problem: many extra 4D fields; some with mass n/R , but some are massless! E.g., 5D gauge field:

$$V_\mu(x^\mu, y) = \underbrace{V_\mu(x^\mu)}_{\text{good}} + \sum_n V_\mu^n(x^\mu) \cos(ny/R) + \sum_m V_\mu^m(x^\mu) \sin(my/R)$$

$$V_5(x^\mu, y) = \underbrace{V_5(x^\mu)}_{\text{bad}} + \sum_n V_5^n(x^\mu) \cos(ny/R) + \sum_m V_5^m(x^\mu) \sin(my/R)$$

- A new solution...

- Compactify on S^1/Z_2 instead (orbifold); require

$$y \rightarrow -y : \quad V_\mu \rightarrow V_\mu \quad V_5 \rightarrow -V_5$$

- Unwanted scalar is projected out:

$$V_\mu(x^\mu, y) = \underbrace{V_\mu(x^\mu)}_{\text{good}} + \sum_n V_\mu^n(x^\mu) \cos(ny/R) + \cancel{\sum_m V_\mu^m(x^\mu) \sin(my/R)}$$

$$V_5(x^\mu, y) = \underbrace{V_5(x^\mu)}_{\text{bad}} + \cancel{\sum_n V_5^n(x^\mu) \cos(ny/R)} + \sum_m V_5^m(x^\mu) \sin(my/R)$$

- Similar projection on fermions \rightarrow 4D chiral theory, ...
- Very simple (requires UV completion at $\Lambda \gg R^{-1}$)

Appelquist, Cheng, Dobrescu (2001)

KK-Parity

- An immediate consequence: conserved KK-parity $(-1)^{KK}$
Interactions require an even number of odd KK modes

- 1st KK modes must be pair-produced at colliders

Macesanu, McMullen, Nandi (2002)

- weak bounds: $R^{-1} > 200 \text{ GeV}$

Appelquist, Yee (2002)

- LKP (lightest KK particle) is stable – dark matter

Kolb, Slansky (1984)

Saito (1987)

Other Extra Dimension Models

- SM on brane; gravity in bulk (brane world)
 - Requires localization mechanism
 - No concrete dark matter candidate
- fermions on brane; bosons and gravity in bulk
 - Requires localization mechanism
 - $R^{-1} > \text{few TeV}$ from $f\bar{f} \rightarrow V_\mu^1 \rightarrow f\bar{f}$
 - No concrete dark matter candidate
- everything in bulk (UED)
 - No localization mechanism required
 - Natural dark matter candidate – LKP

UED and SUSY

Similarities:

- Superpartners \rightarrow KK partners
- R-parity \rightarrow KK-parity
- LSP \rightarrow LKP
- Bino dark matter $\rightarrow B^1$ dark matter
- Sneutrino dark matter $\rightarrow \nu^1$ dark matter

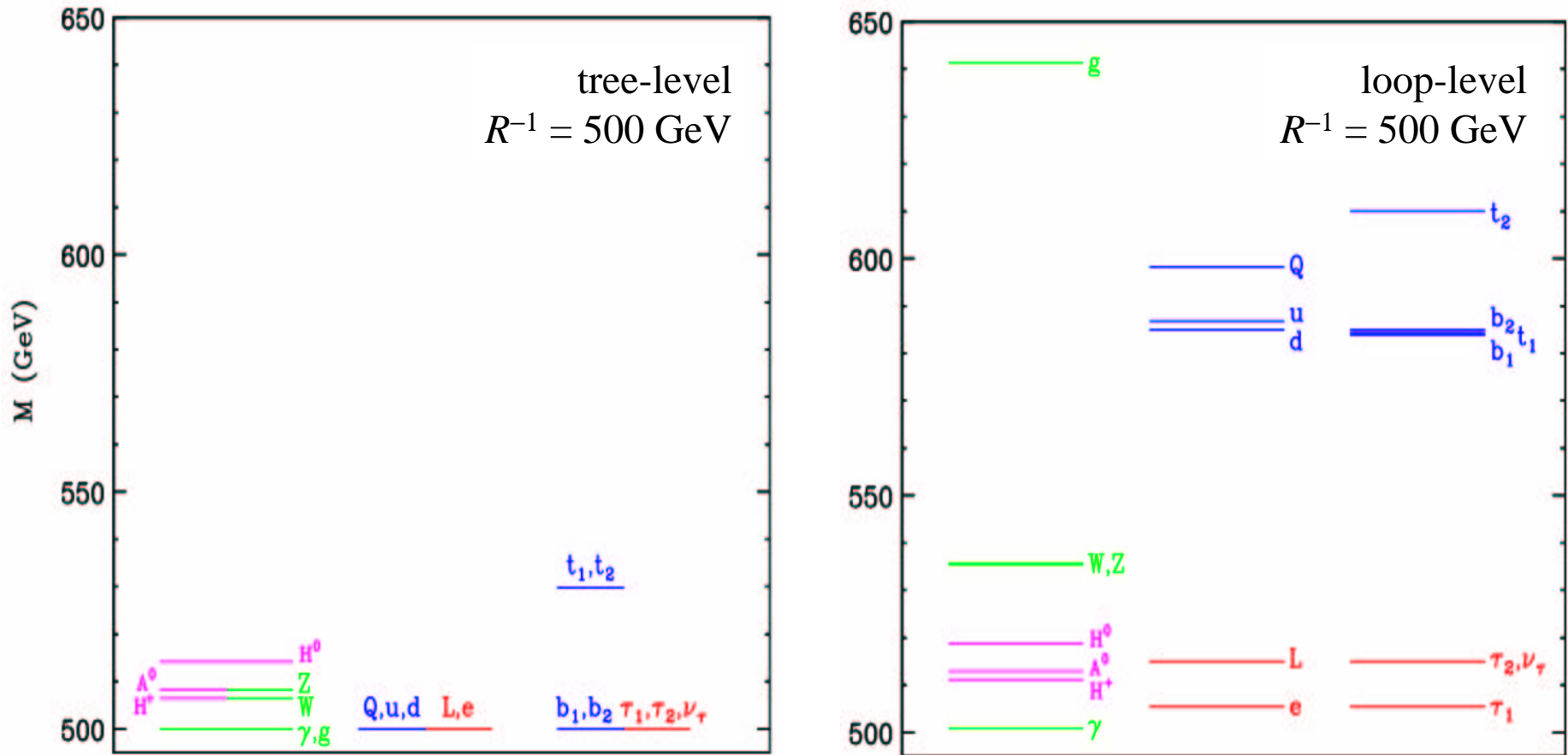
...

Not surprising: SUSY is also an extra (fermionic) dimension theory

Differences:

- KK modes highly degenerate, split by EWSB and loops
- Fermions \rightarrow Bosons

Minimal UED KK Spectrum

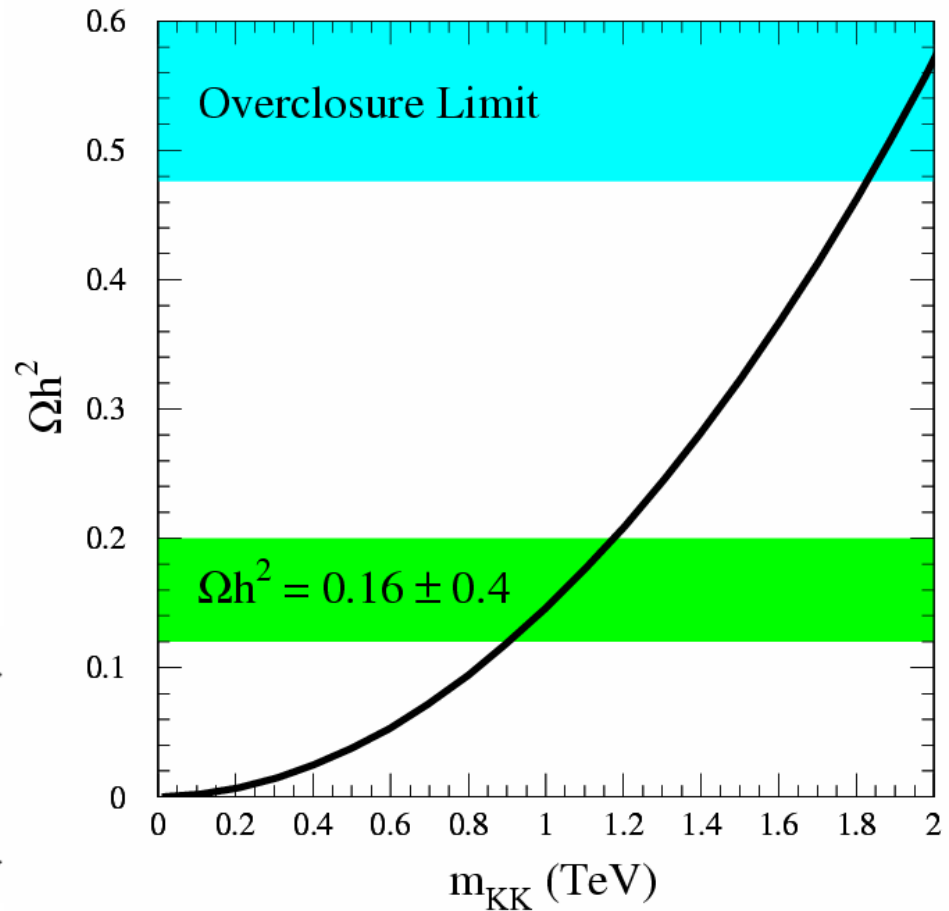
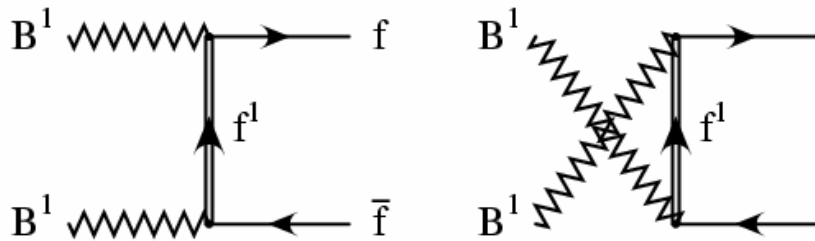


Cheng, Matchev, Schmaltz (2002)

B^1 Dark Matter

- LKP is nearly pure B^1 in minimal model (more generally, a B^1 - W^1 mixture)

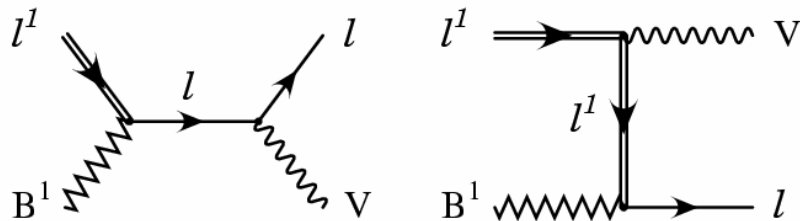
- Relic density:



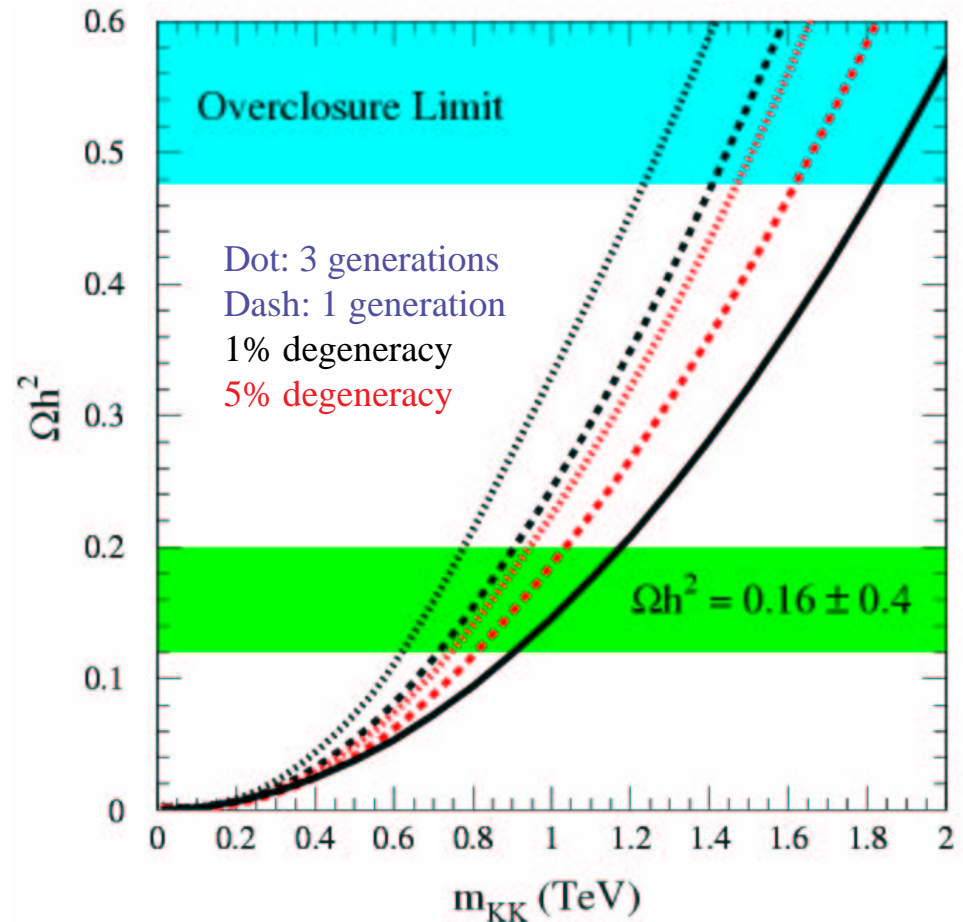
Servant, Tait (2002)

Co-annihilation

- But degeneracy \rightarrow co-annihilations important
- Co-annihilation processes:



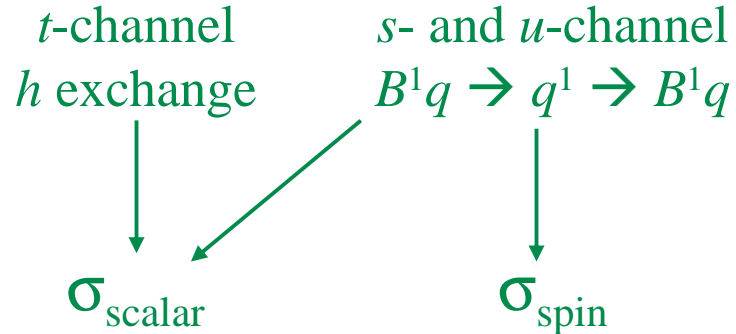
- Preferred m_{B^1} : l^1 lowers it, q^1 raises it; 100s of GeV to few TeV possible



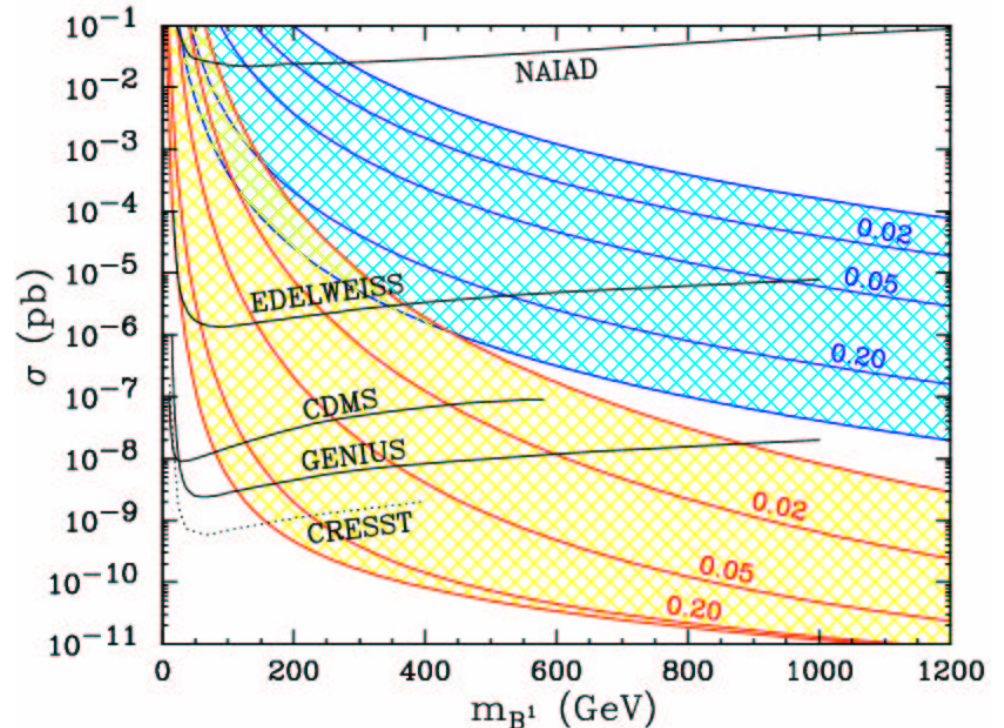
Servant, Tait (2002)

B^1 Dark Matter Detection

Direct Detection



- s -channel enhanced by B^1 - q^1 degeneracy



Cheng, Feng, Matchev (2002)

- Constructive interference: lower bound on both σ_{scalar} and σ_{spin}

B^1 Dark Matter Detection

- Indirect Detection:
 - Positrons from the galactic halo
 - Muons from neutrinos from the Sun and Earth
 - Gamma rays from the galactic center
- All rely on annihilation, very different from SUSY
 - For neutralinos (Majorana fermions), $\chi\chi \rightarrow ff$ is chirality suppressed
 - $B^1B^1 \rightarrow ff$ isn't; generically true for bosons

Positrons

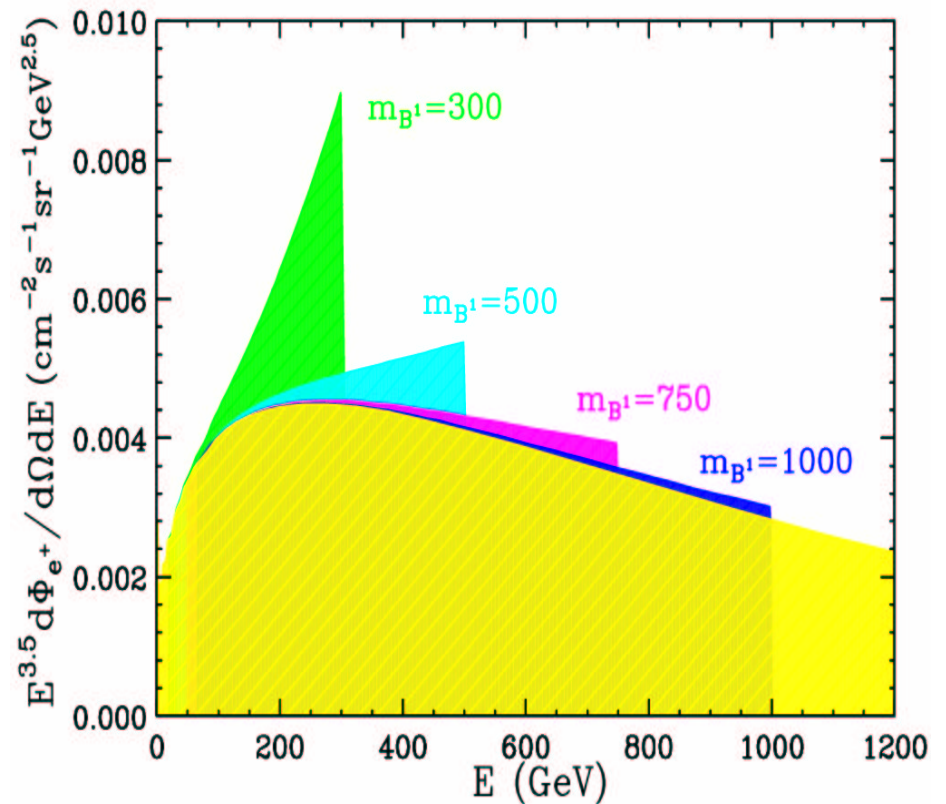
$$\frac{d\Phi_{e^+}}{d\Omega dE} = \frac{\rho^2}{m_{B^1}^2} \sum_i \langle \sigma_i v \rangle B_{e^+}^i \int dE_0 f_i(E_0) G(E_0, E)$$

where

$\langle \sigma_i v \rangle$ = the annihilation σ to channel i
 $B_{e^+}^i$ = e^+ branching fraction in channel i
 $f_i(E_0)$ = injection spectrum
 G = e^+ propagator in the galaxy

Moskalenko, Strong (1999)

- Here $f_i(E_0) \sim \delta(E_0 - m_{B^1})$, and the peak is not erased by propagation (cf. $\chi\chi \rightarrow W^+W^- \rightarrow e^+ \nu e^- \bar{\nu}$)
- AMS will have e^+/e^- separation at 1 TeV and see $\sim 1000 e^+$ above 500 GeV



Cheng, Feng, Matchev (2002)

Muons from Neutrinos

- Muon flux is

$$\Phi_\mu \propto \sum_{F,i} B_F \langle N z^2 \rangle_{F,i}$$

where $i = \nu, \bar{\nu}$, F labels final states, and $z \equiv E_\nu/E_{in}$.

Ritz, Seckel (1988)

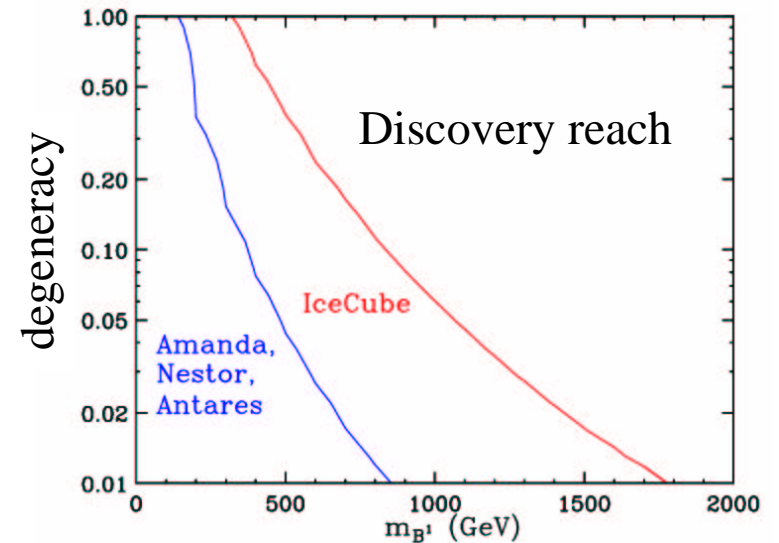
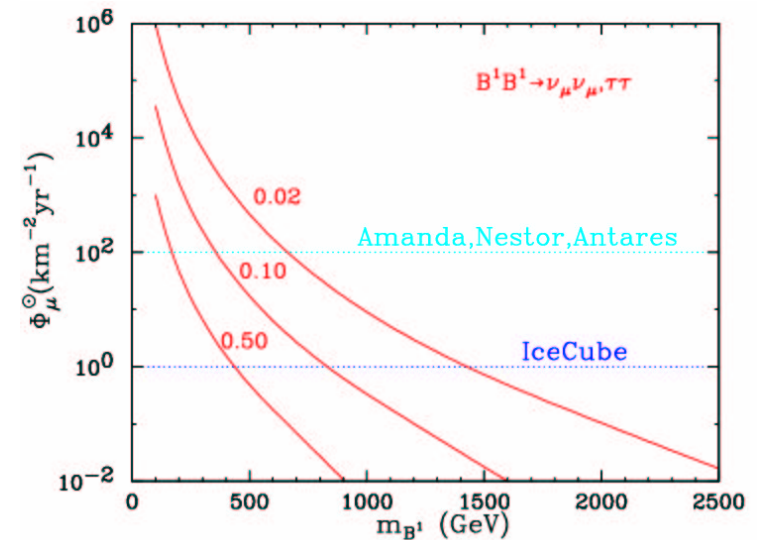
Jungman, Kamionkowski, Griest (1995)

- $B^1 B^1 \rightarrow \nu \nu$ is also unsuppressed, gives hard neutrinos, enhanced μ flux

Cheng, Feng, Matchev (2002)

Hooper, Kribs (2002)

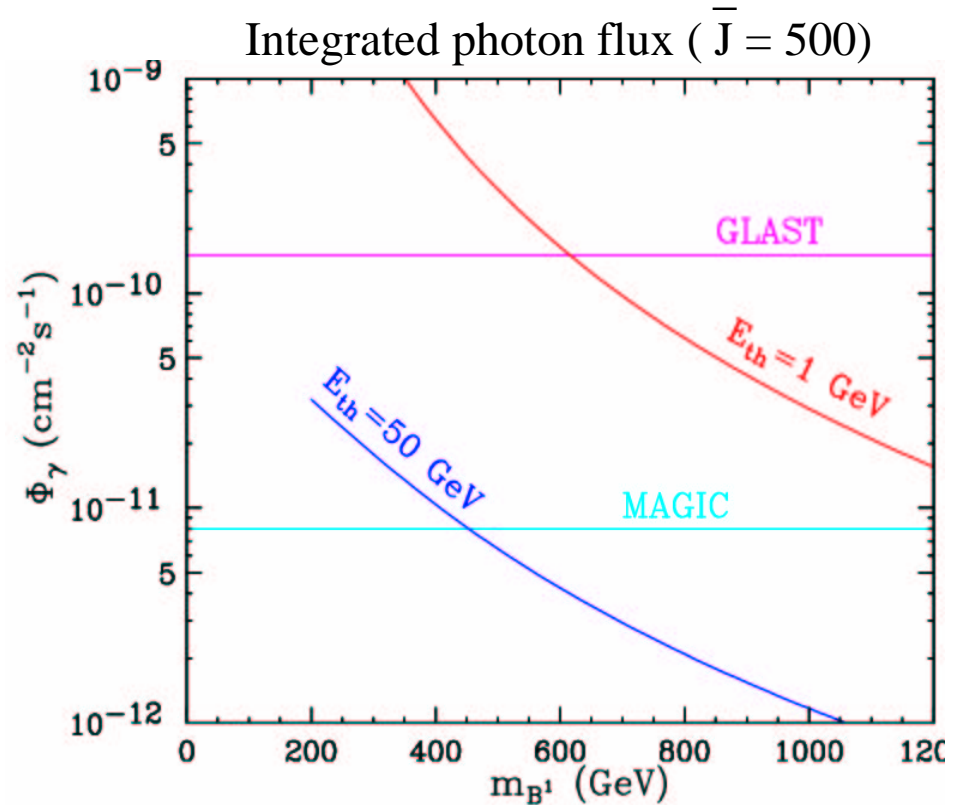
Bertone, Servant, Sigl (2002)



Gamma Rays

- $B^1 B^1 \rightarrow \gamma \gamma$ is loop-suppressed, but light quark fragmentation gives hardest photons, so absence of chirality suppression helps again
- Results sensitive to halo clumpiness; choose moderate value

Bergstrom, Ullio, Buckley (1998)

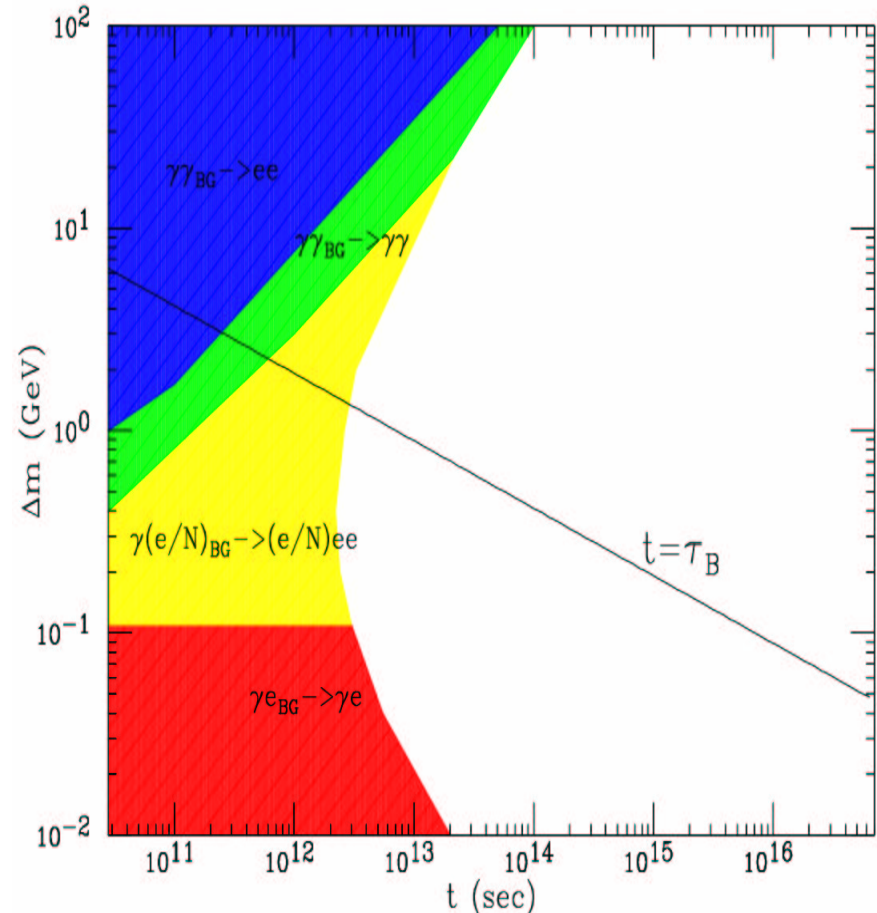


Cheng, Feng, Matchev (2002)

Graviton Dark Matter

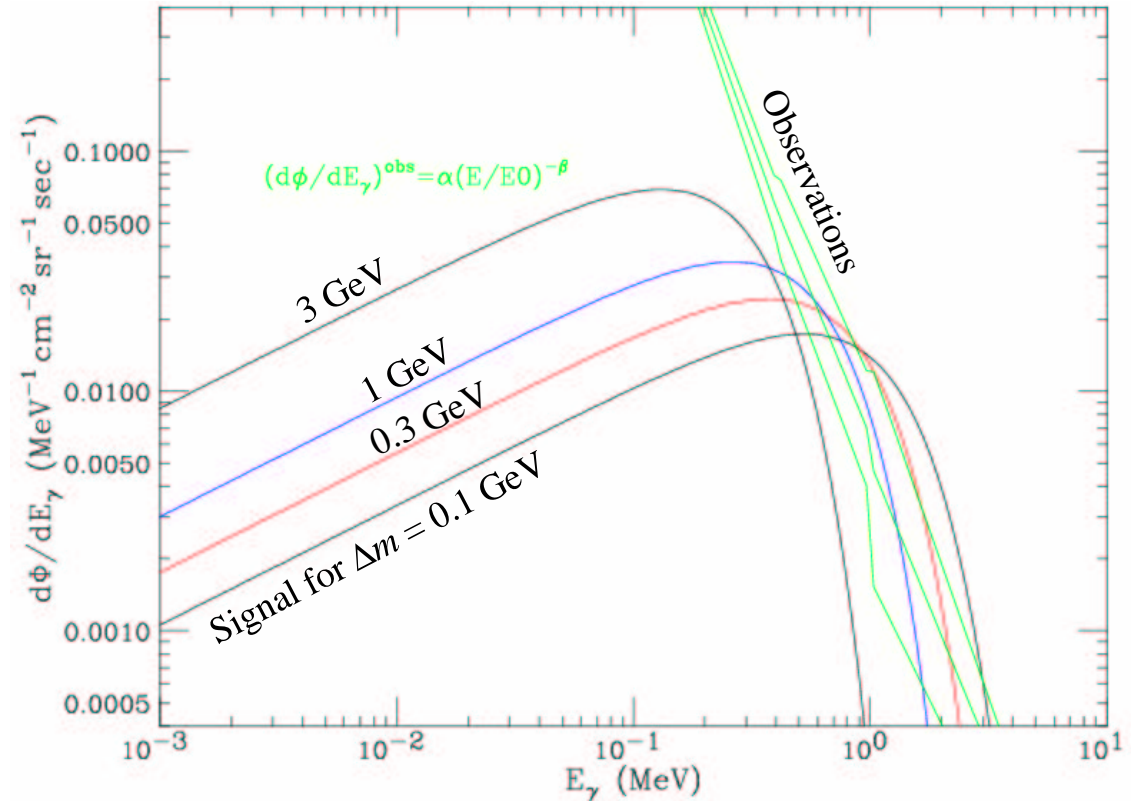
Feng, Rajaraman, Takayama

- LKP may be 1st KK graviton G^1
- If NLKP is B^1 , B^1 freezes out, then decays via $B^1 \rightarrow \gamma G^1$ with lifetime $\sim 10^{12}$ sec $[1 \text{ GeV} / \Delta m]^3$; get very late, very soft decays
- Evades BBN constraints for small Δm (or if NLKP is ν^1)
- G^1 DM retains WIMP virtues, but is undetectable by all conventional dark matter searches



Diffuse Photon Flux

- Late $B^1 \rightarrow \gamma G^1$ implies a novel WIMP signal: diffuse photon flux
- Large Δm implies larger initial energy, but also more red shifting; latter dominates
- Present flux peaks ~ 1 MeV, yields observable signal



Feng, Rajaraman, Takayama

Conclusions

- Extra Dimensions yield natural dark matter candidates
- Much work to be done: h^1 , W^1 , G^1 , non-minimal models, ..., but already several novel features:
 - s -channel enhancements from degeneracy
 - Annihilation not chirality suppressed
 - Graviton dark matter – naturally desired thermal relic density, but inaccessible to all conventional searches
- Direct detection, μ from ν , e^+ , γ rays, may all push sensitivity beyond collider reach
- **KKDM – escape from the tyranny of neutralino dark matter!**