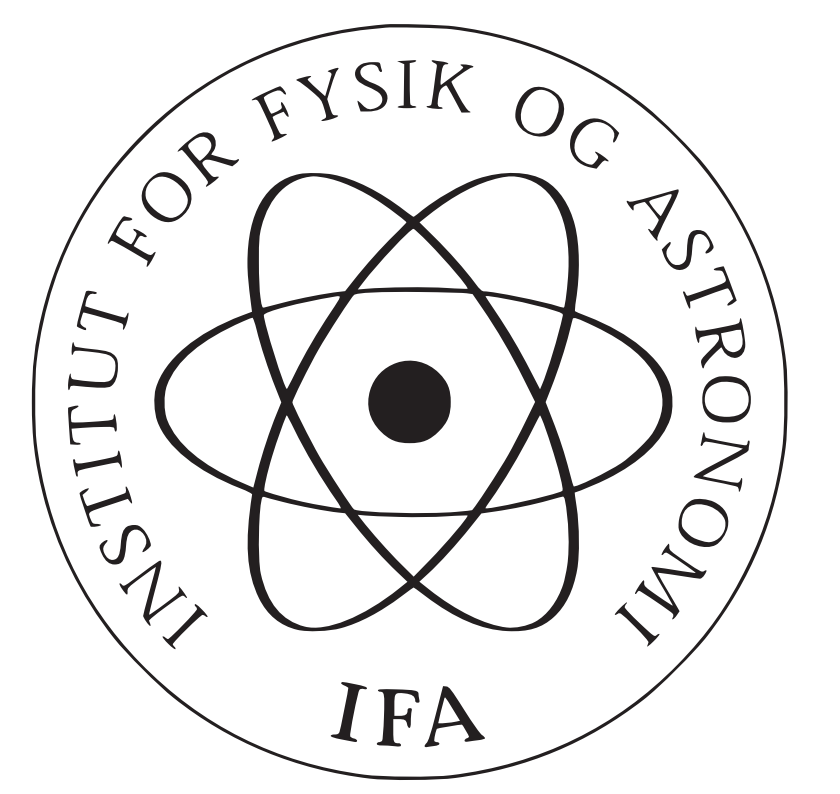


Dipolar fermions in a two-dimensional lattice at nonzero temperature



Eidgenössische Technische Hochschule Zürich
Swiss Federal Institute of Technology Zurich

Anne-Louise Gadsbølle (1,2) and Georg M. Bruun (1)

1) Department of Physics and Astronomy, University of Aarhus, DK 8000 Aarhus C, Denmark

2) Theoretische Physik, ETH Zürich, Zürich CH-8093, Switzerland

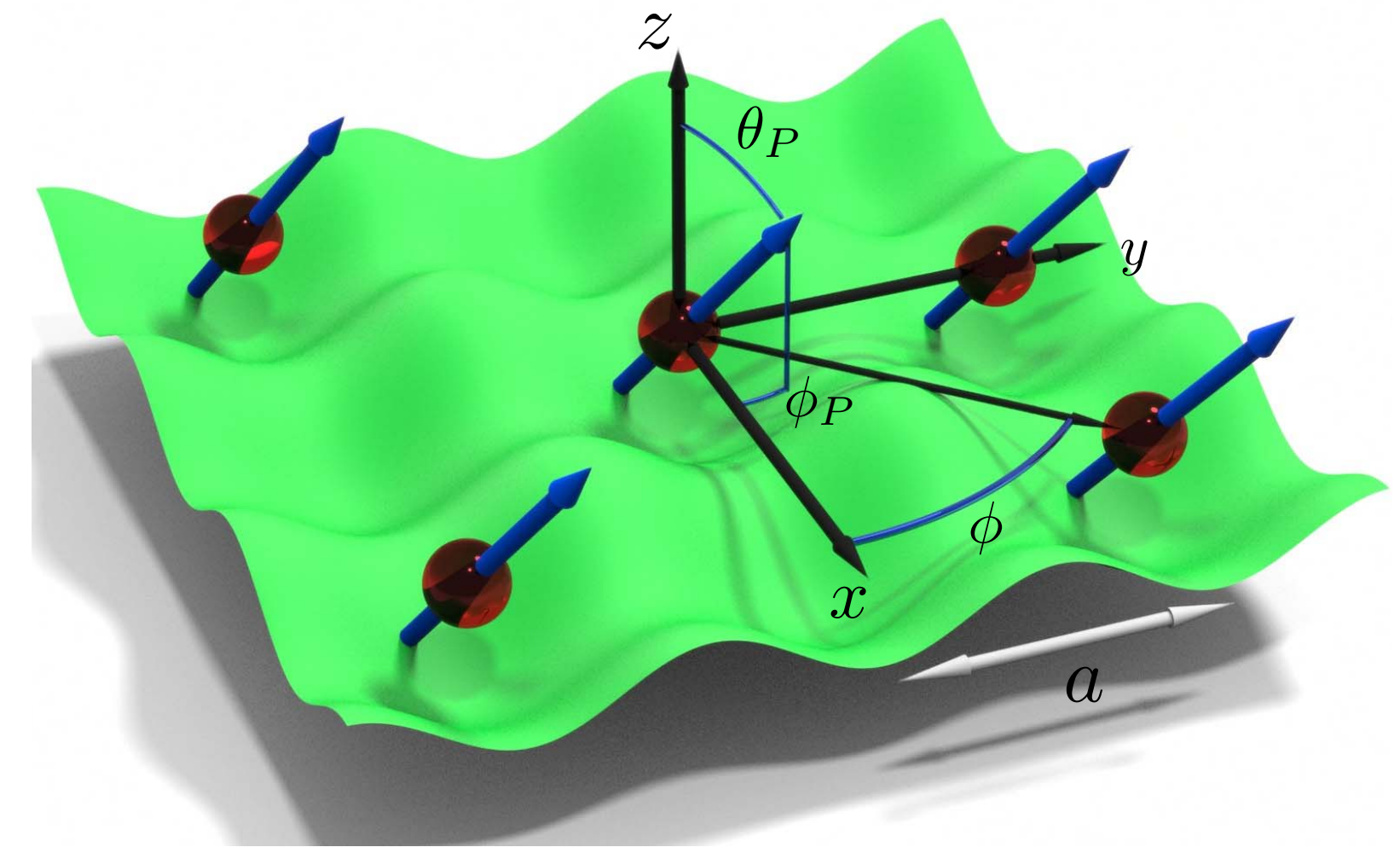
A.-L. Gadsbølle and G. M. Bruun, Phys. Rev. A **85** (R), 021604 (2012)
A.-L. Gadsbølle and G. M. Bruun, Phys. Rev. A **86**, 033623 (2012)

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Basic formalism

$$\hat{H} = -t \sum_{\langle ij \rangle} (\hat{c}_i^\dagger \hat{c}_j + h.c.) + \sum_i \left(\frac{1}{2} m \omega^2 r_i^2 - \mu \right) \hat{n}_i + \frac{1}{2} \sum_{i \neq j} V_D(\mathbf{r}_{ij}) \hat{n}_i \hat{n}_j$$

$$V_D(\mathbf{r}) = \frac{D^2}{r^3} [1 - 3 \cos^2(\phi_P - \phi) \sin^2(\theta_P)]$$



$$g = D^2/a^3$$

The Bogoliubov-de Gennes equations

$$\sum_j \begin{pmatrix} L_{ij} & \Delta_{ij} \\ \Delta_{ji}^* & -L_{ij} \end{pmatrix} \begin{pmatrix} u_\eta^j \\ v_\eta^j \end{pmatrix} = E_\eta \begin{pmatrix} u_\eta^i \\ v_\eta^i \end{pmatrix},$$

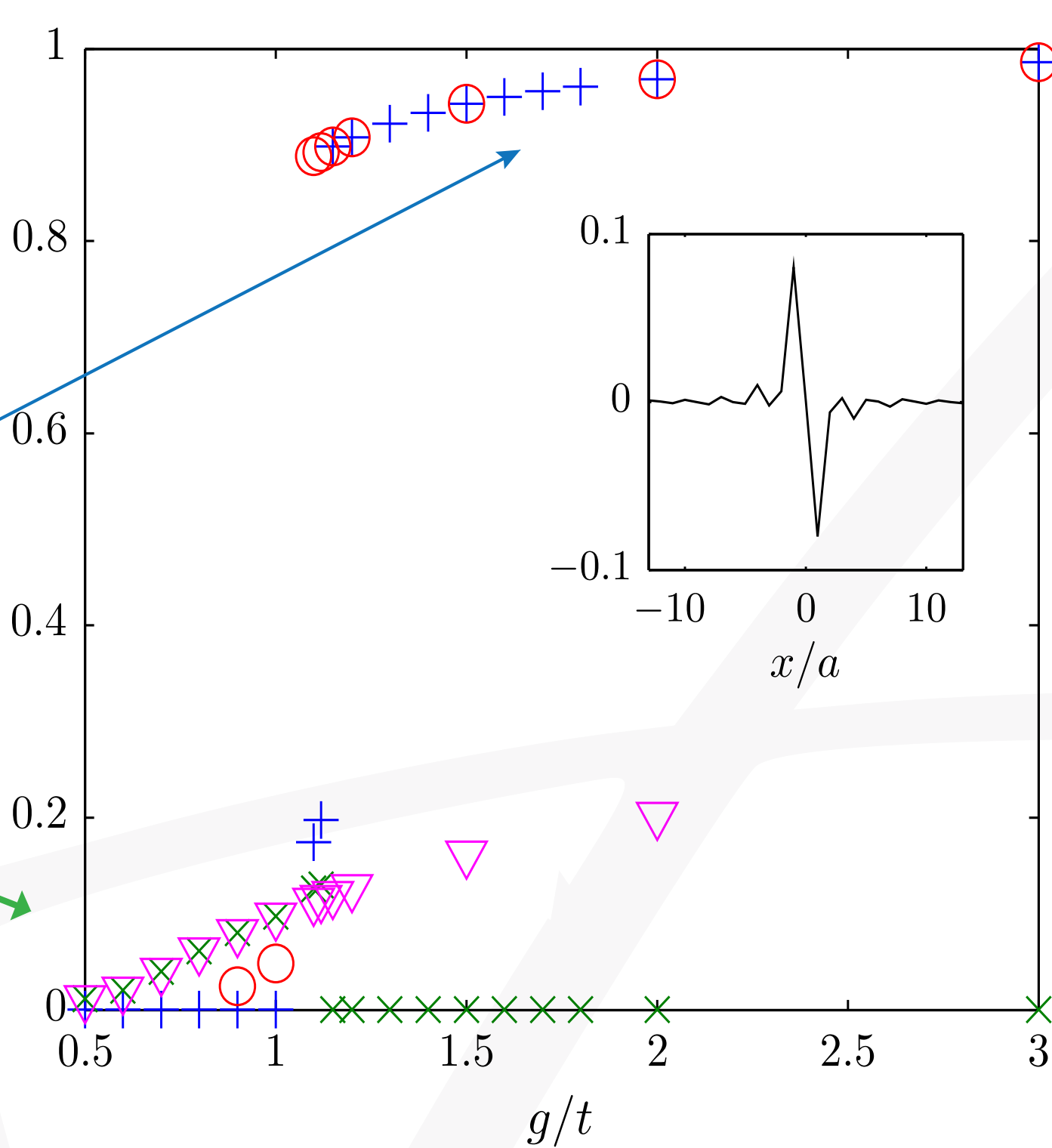
with $\Delta_{ij} = V_D(\mathbf{r}_{ij}) \langle \hat{c}_j \hat{c}_i \rangle$ and $L_{ij} = -t \delta_{ij} + (\sum_k V_D(\mathbf{r}_{ik}) \langle \hat{n}_k \rangle + \frac{m}{2} \omega^2 r_i^2 - \mu) \delta_{ij}$

The untrapped system at T=0:

p-wave pairing and striped density order for $f=1/3$

$$(\theta_P, \phi_P) = \left(\frac{\pi}{2}, 0 \right)$$

+: density order $M = \langle \hat{n}_i - \hat{n}_{i+e_y} \rangle$
x: *p*-wave pairing $\langle \hat{c}_i \hat{c}_{i+e_x} \rangle$



$$\theta_P > \arcsin\left(\frac{1}{\sqrt{3}}\right)$$

Melting of density order:

Self-consistency for stripes: $1 = \frac{1}{N_L} \sum_{k_y > 0} \frac{\tilde{V}_D(0, \pi/a) (f(E_{1\mathbf{k}}) - f(E_{2\mathbf{k}}))}{\sqrt{4t^2 \cos^2 k_y a + \left(\frac{\tilde{V}_D(0, \pi/a)}{2} M\right)^2}}$

Critical temperature: $T_c = -\frac{1}{4} \begin{cases} \tilde{V}_D(0, \pi/a) & \text{Stripes} \\ \tilde{V}_D(\pi/a, \pi/a) & \text{Checkerboard} \end{cases}$

$$T_c \propto g$$

Superfluid melting:

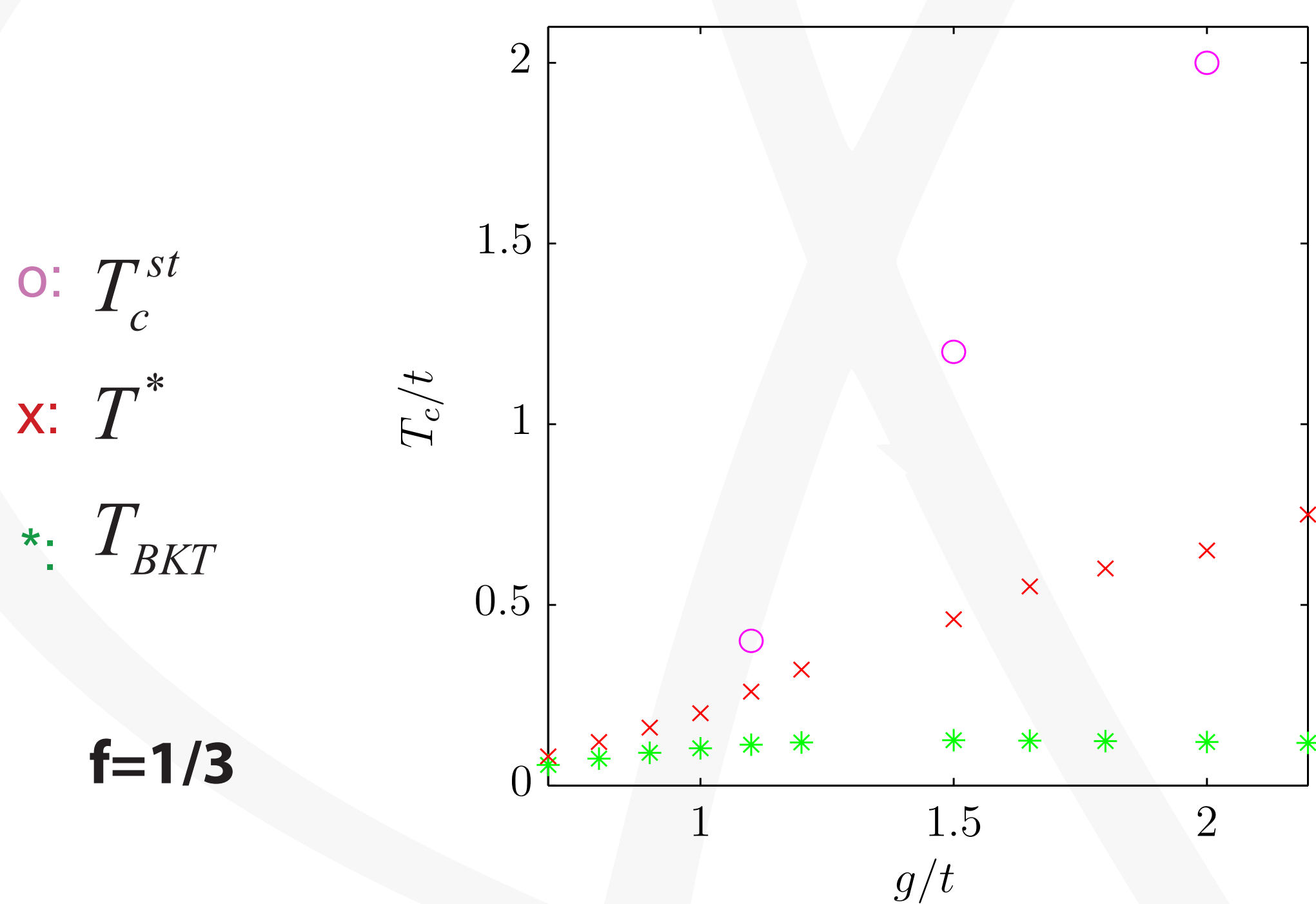
Phase-twist on order parameter: $\Delta_{ij} \rightarrow \Delta_{ij} e^{i(x_i+x_j)\delta\theta/a}$ $v_s = \frac{\hbar \delta\theta}{2m^* a}$

Energy cost: $F_\theta - F_0 \approx \frac{J^{xy}}{2} \sum_i \delta\theta^2 = \frac{N}{2} \rho_{s,x} m^* v_s^2$

Equivalent to gauge transformation: $\hat{H}_\Theta = e^{-i\delta\theta \sum_j \hat{x}_j/a} \hat{H} e^{i\delta\theta \sum_j \hat{x}_j/a}$

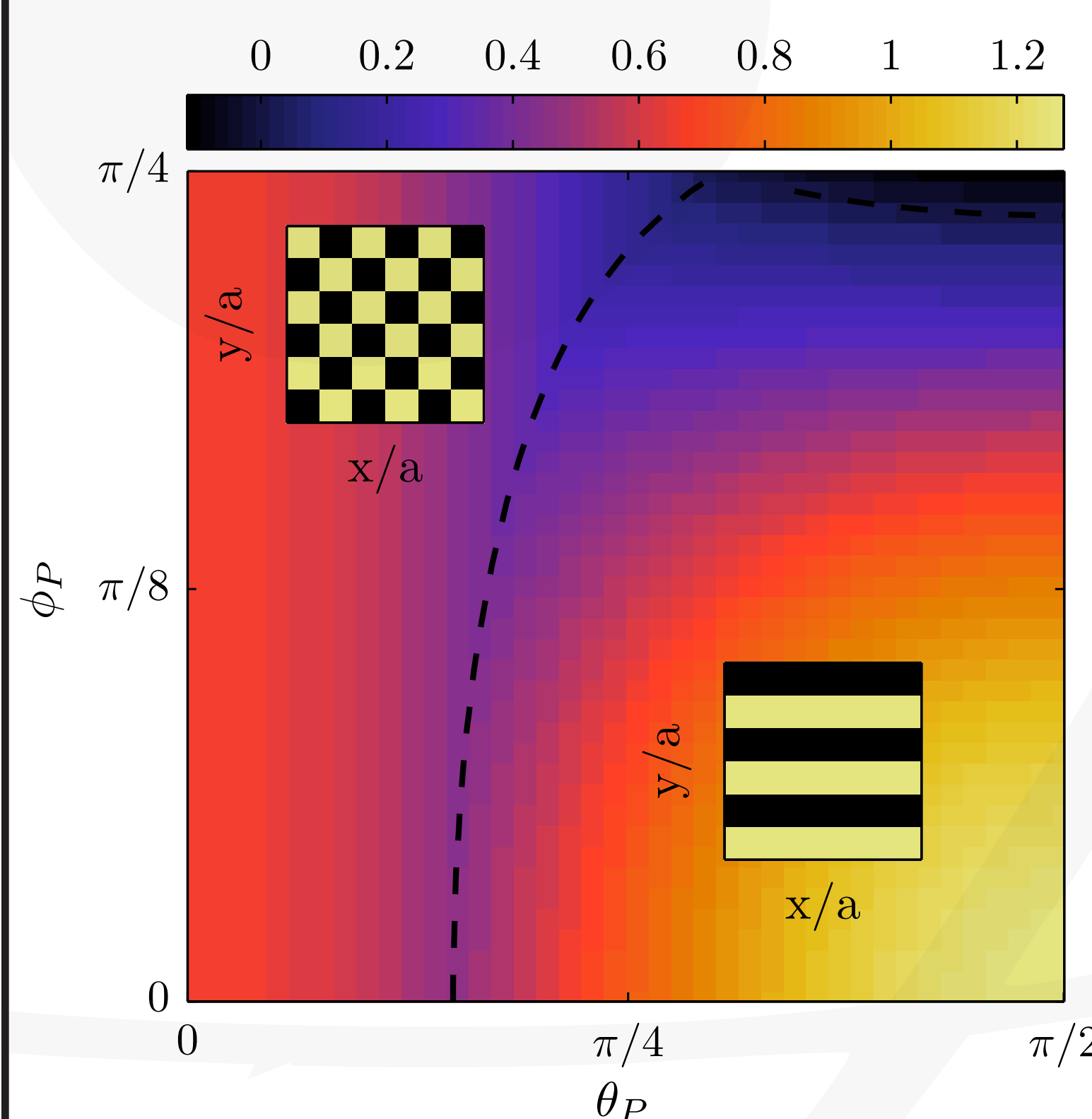
$$T_{BKT} = \frac{\pi \bar{J}}{2} = \frac{\pi}{4} \frac{N}{N_L} \bar{\rho}_s t$$

Critical temperature in the untrapped system:



The critical temperature for the superfluid phase (*'s) and the striped phase (o's) for $(\theta_P, \phi_P) = (\pi/2, 0)$ as a function of coupling strength for one third filling obtained from a numerical calculation on a 27×27 lattice. The x's give the mean-field superfluid transition temperature T^* . For illustrative purposes, we plot the critical temperature of the superfluid phase even for $g/t > 1.15$, where stripe order suppresses superfluidity.

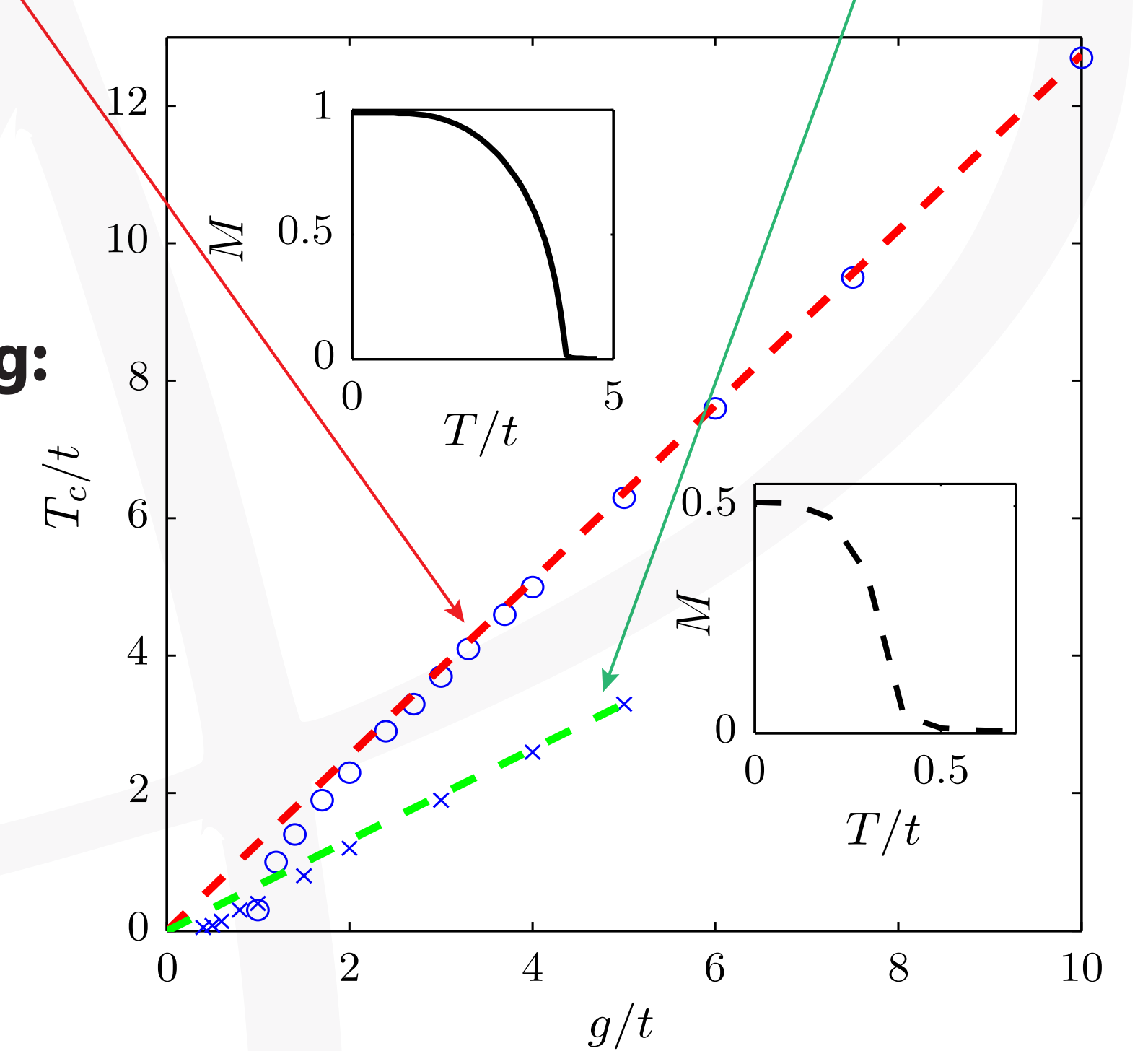
Critical temperature in the strong coupling regime:



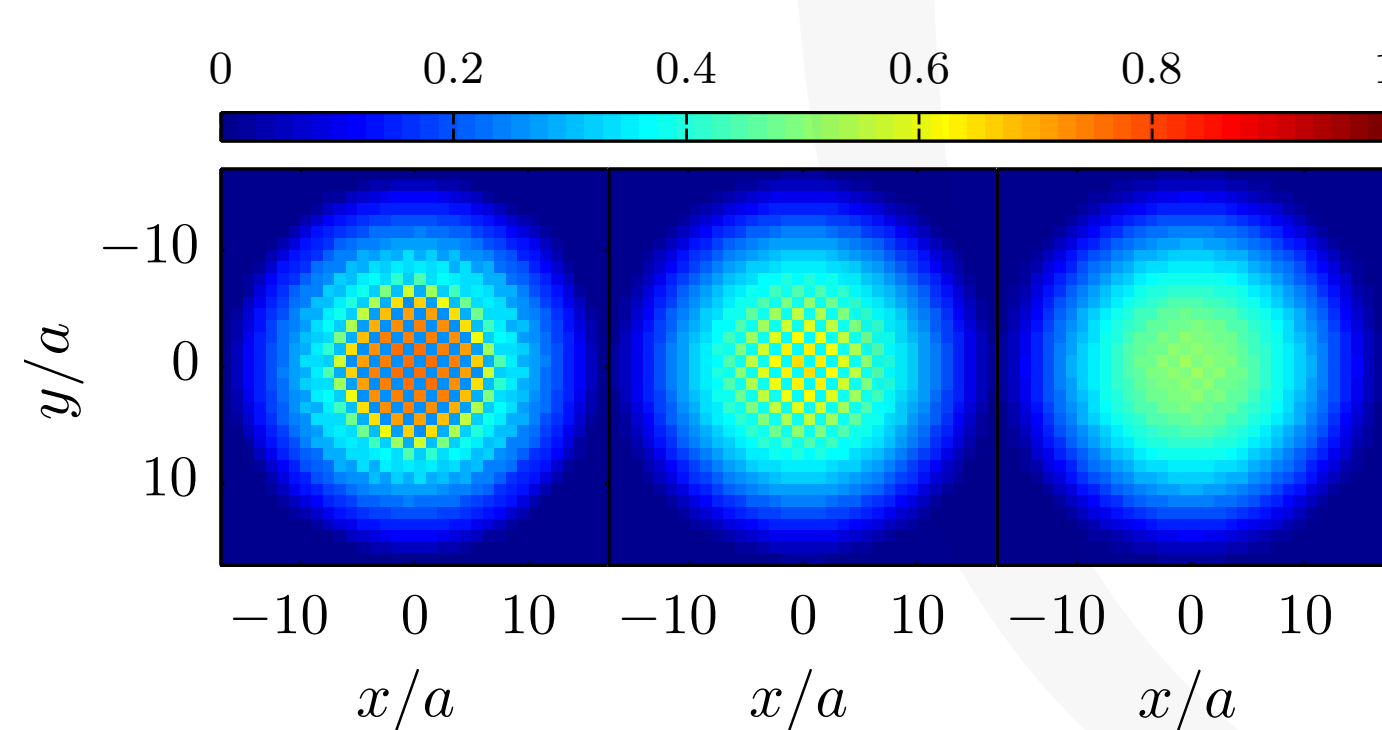
$$T_c^{st} = -\frac{1}{4} \tilde{V}_D(0, \pi/a)$$

$$T_c^{CB} = -\frac{1}{4} \tilde{V}_D(\pi/a, \pi/a)$$

Half-filling:
 $f=1/2$

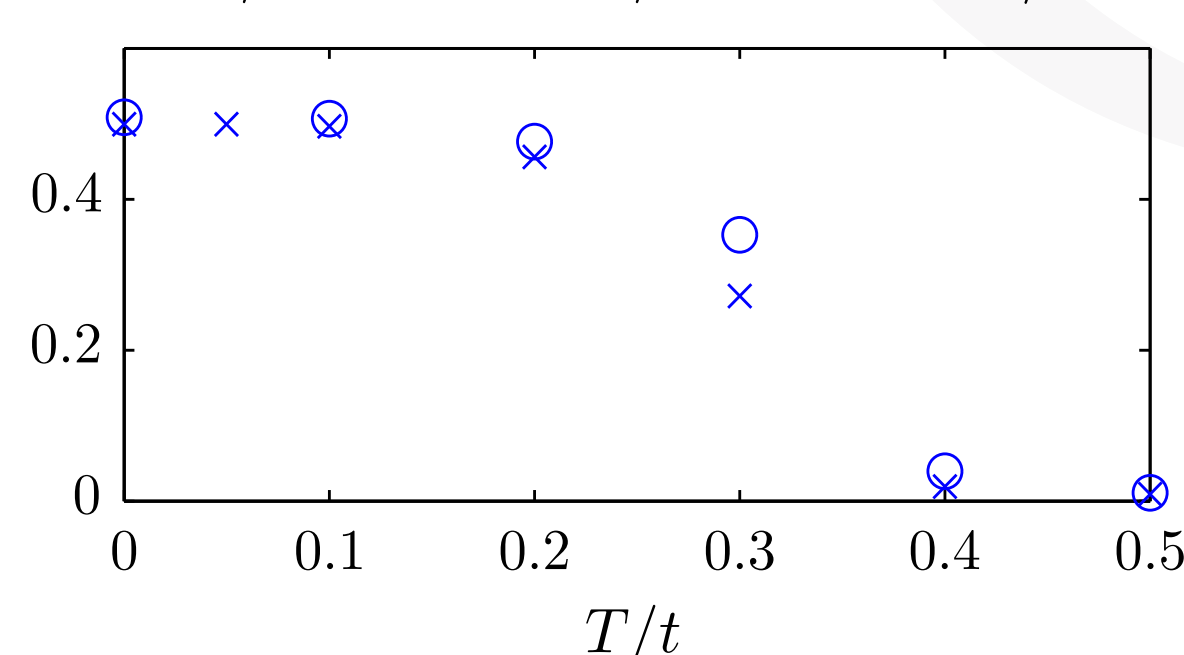


Checkerboard melting: $\theta_P = 0$



$$M = \langle \hat{n}_i - \hat{n}_{i+e_y} \rangle$$

$$\frac{1}{2} \left(\langle \hat{c}_{i+e_x} \hat{c}_i \rangle + \langle \hat{c}_i \hat{c}_{i+e_x} \rangle \right)$$



Stripe and superfluid melting: $(\theta_P, \phi_P) = (\pi/2, 0)$

