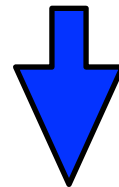


# Realizing Topological Phases with Dipolar Spins

Alexey V. Gorshkov

Institute for Quantum Information and Matter (IQIM),  
Caltech

IQIM



(summer 2013)

jqi

NIST

Joint Quantum Institute (JQI)  
NIST and University of Maryland



KITP conference  
New Science with Ultracold Molecules  
March 14, 2013

# Realizing Topological Phases with Dipolar Spins

- symmetry protected topological phases [[PRB 87, 081106\(R\) \(2013\)](#)]:  
S. Manmana, K. Hazzard, A. M. Rey - JILA [related: [arXiv:1301.5636](#)]  
M. Stoudenmire - UCI
- fractional Chern insulators [[PRL 109, 266804 \(2012\)](#) & [arXiv:1212.4839](#)]:  
N. Yao, C. Laumann, S. Bennett, E. Demler, M. Lukin - Harvard  
A. Läuchli, P. Zoller - Innsbruck  
J. Ye - JILA



# Realizing Topological Phases with Dipolar Spins

- symmetry protected topological phases [PRB 87, 081106(R) (2013)]:  
S. Manmana, K. Hazzard, A. M. Rey - JILA [related: arXiv:1301.5636]  
M. Stoudenmire - UCI
- fractional Chern insulators [PRL 109, 266804 (2012) & arXiv:1212.4839]:  
N. Yao, C. Laumann, S. Bennett, E. Demler, M. Lukin - Harvard  
A. Läuchli, P. Zoller - Innsbruck  
J. Ye - JILA

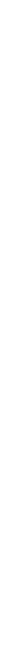


# Motivation

Interacting dipoles

Electric

Magnetic

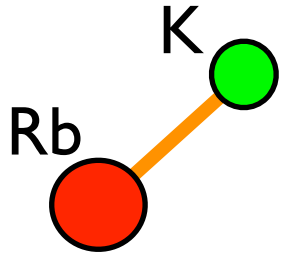


# Motivation

## Interacting dipoles

Electric

Magnetic



polar molecules

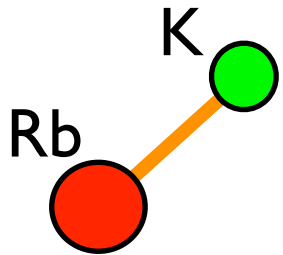
[ultracold, rovibrational  
ground state: Ye, Jin,  
Weidemuller, Inouye,  
Nagerl, etc...]

# Motivation

## Interacting dipoles

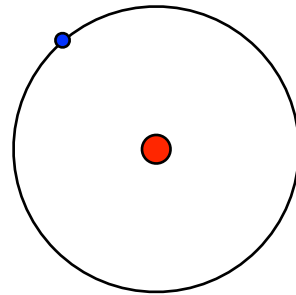
Electric

Magnetic



polar molecules

[ultracold, rovibrational ground state: Ye, Jin, Weidemuller, Inouye, Nagerl, etc...]



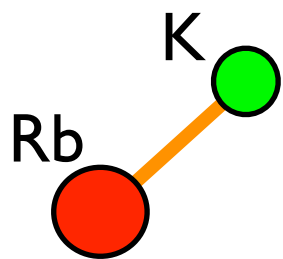
Rydberg atoms

[Grangier, Saffman, Pfau, Weidemuller, Bloch, etc...]

# Motivation

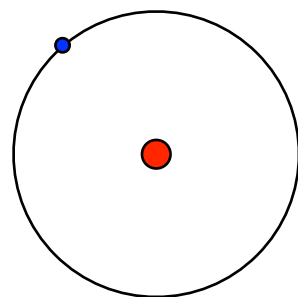
## Interacting dipoles

### Electric



polar molecules

[ultracold, rovibrational ground state: Ye, Jin, Weidemuller, Inouye, Nagerl, etc...]



Rydberg atoms

[Grangier, Saffman, Pfau, Weidemuller, Bloch, etc...]

### Magnetic

large  
electronic  
angular  
momentum  $J$

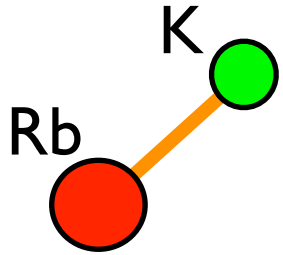
magnetic atoms  
(e.g. Dy, Er, Cr)

[Lev, Ferlaino, Pfau, Laburthe-Tolra, etc...]

# Motivation

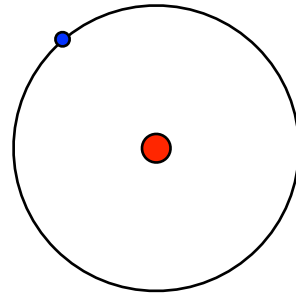
## Interacting dipoles

### Electric



polar molecules

[ultracold, rovibrational ground state: Ye, Jin, Weidemuller, Inouye, Nagerl, etc...]



Rydberg atoms

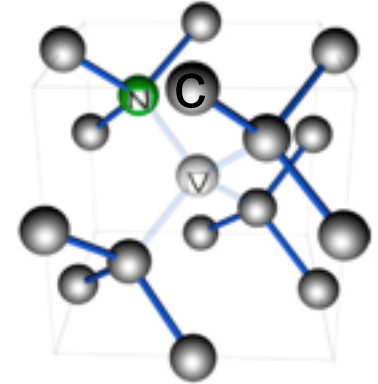
[Grangier, Saffman, Pfau, Weidemuller, Bloch, etc...]

### Magnetic

large  
electronic  
angular  
momentum  $J$

magnetic atoms  
(e.g. Dy, Er, Cr)

[Lev, Ferlaino, Pfau, Laburthe-Tolra, etc...]



NV centers

[Wrachtrup, Jelezko, Lukin, etc...]

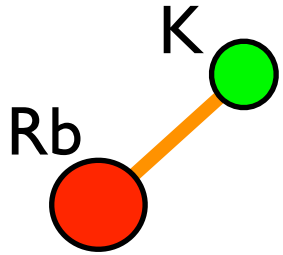
[drawn by Wrachtrup et al.]



# Motivation

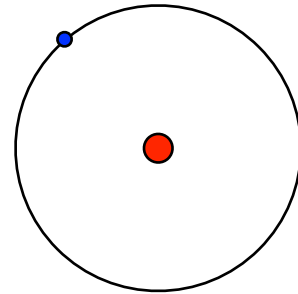
## Interacting dipoles

### Electric



polar molecules

[ultracold, rovibrational ground state: Ye, Jin, Weidemuller, Inouye, Nagerl, etc...]



Rydberg atoms

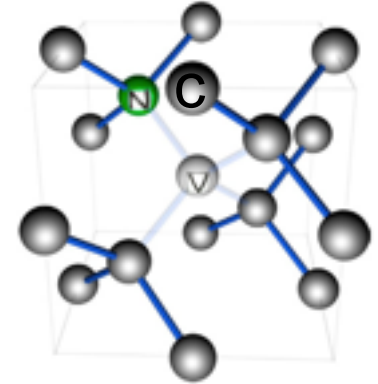
[Grangier, Saffman, Pfau, Weidemuller, Bloch, etc...]

### Magnetic

large electronic angular momentum  $J$

magnetic atoms (e.g. Dy, Er, Cr)

[Lev, Ferlino, Pfau, Laburthe-Tolra, etc...]



NV centers

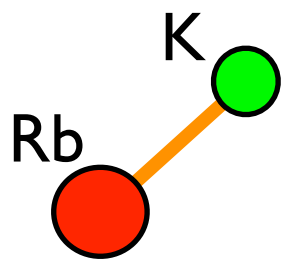
[Wrachtrup, Jelezko, Lukin, etc...]

[drawn by Wrachtrup et al.]

# Motivation

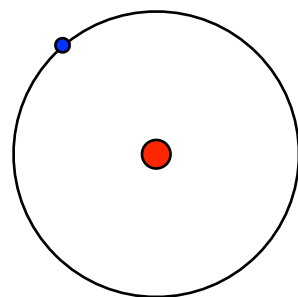
## Interacting dipoles

### Electric



polar molecules

[ultracold, rovibrational ground state: Ye, Jin, Weidemuller, Inouye, Nagerl, etc...]



Rydberg atoms

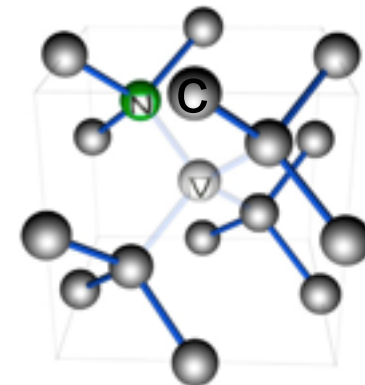
[Grangier, Saffman, Pfau, Weidemuller, Bloch, etc...]

### Magnetic

large electronic angular momentum  $J$

magnetic atoms (e.g. Dy, Er, Cr)

[Lev, Ferlaino, Pfau, Laburthe-Tolra, etc...]



NV centers

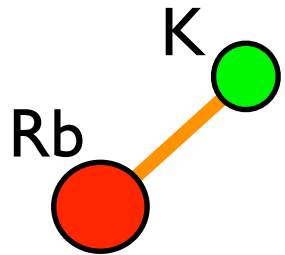
[Wrachtrup, Jelezko, Lukin, etc...]

[drawn by Wrachtrup et al.]

# Motivation

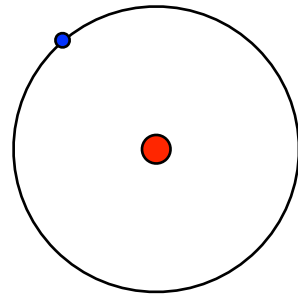
## Interacting dipoles

### Electric



polar molecules

[ultracold, rovibrational ground state: Ye, Jin, Weidemuller, Inouye, Nagerl, etc...]



Rydberg atoms

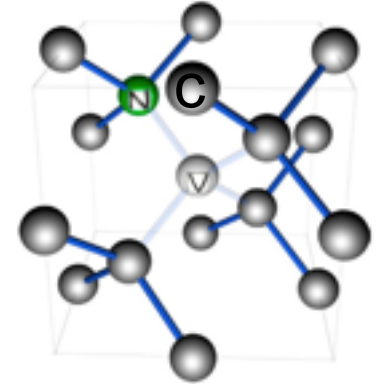
[Grangier, Saffman, Pfau, Weidemuller, Bloch, etc...]

### Magnetic

large electronic angular momentum  $J$

magnetic atoms (e.g. Dy, Er, Cr)

[Lev, Ferlaino, Pfau, Laburthe-Tolra, etc...]



NV centers

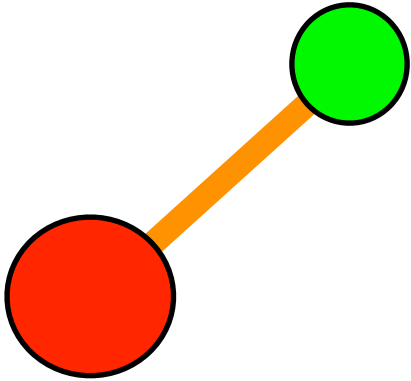
[Wrachtrup, Jelezko, Lukin, etc...]

[drawn by Wrachtrup et al.]

- KRb already loaded into a 3D optical lattice at JILA!

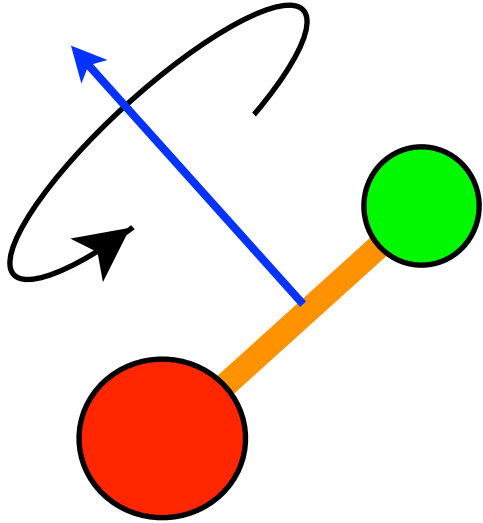
[Jun Ye's talk; Chotia et al, PRL (2012)]

# Motivation

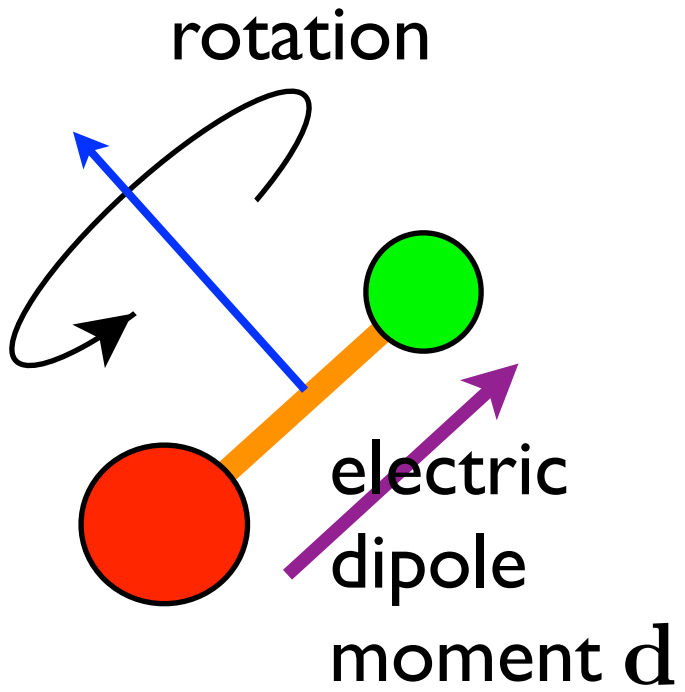


# Motivation

rotation

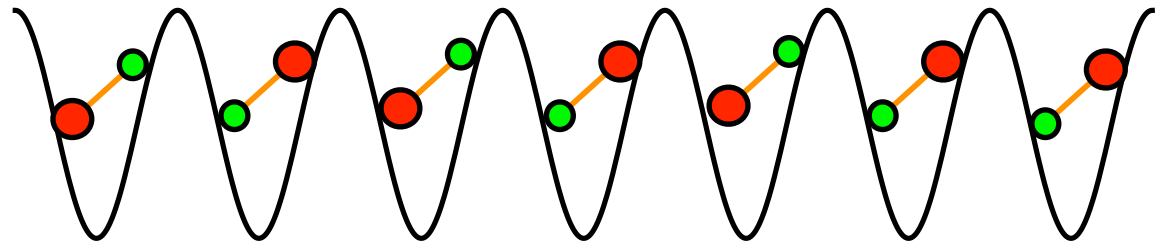
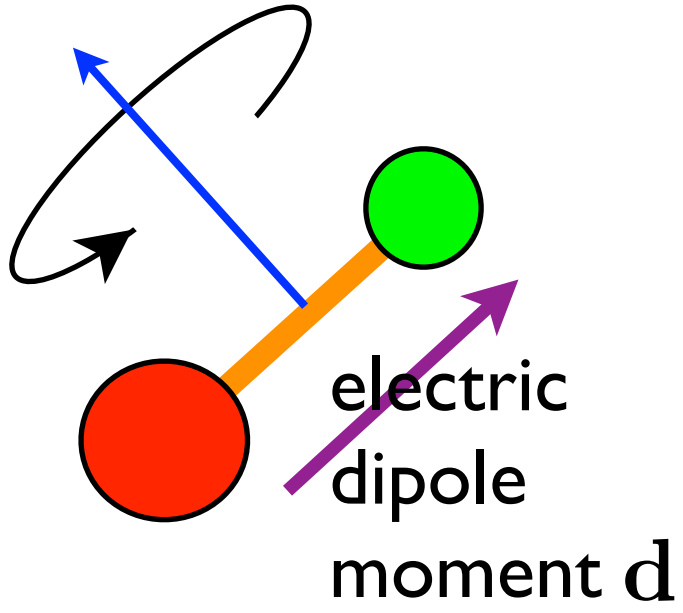


# Motivation

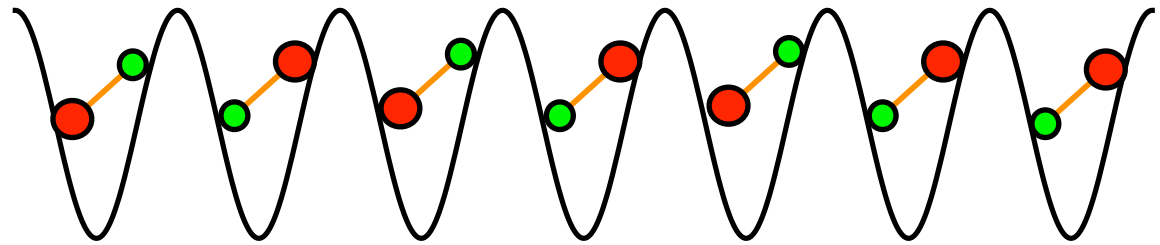
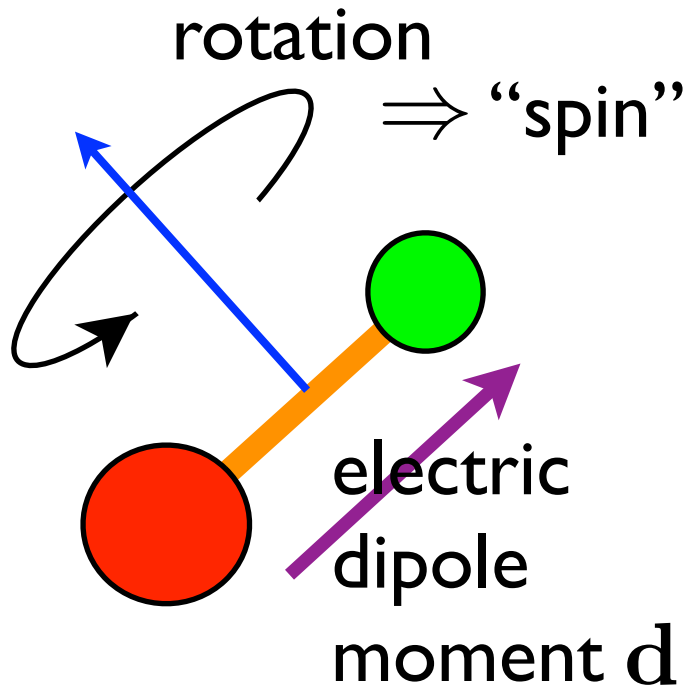


# Motivation

rotation

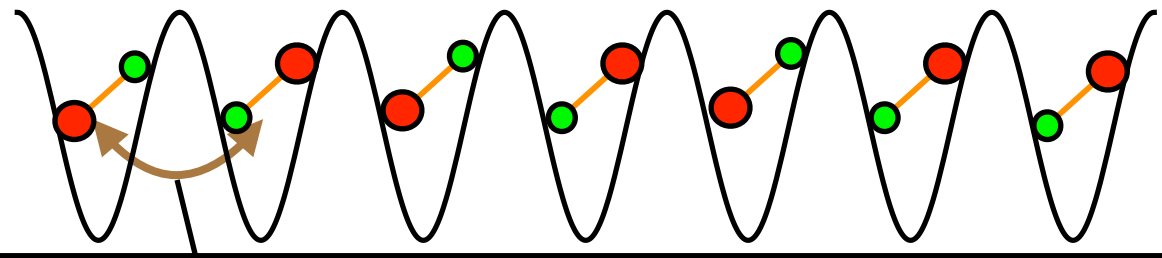
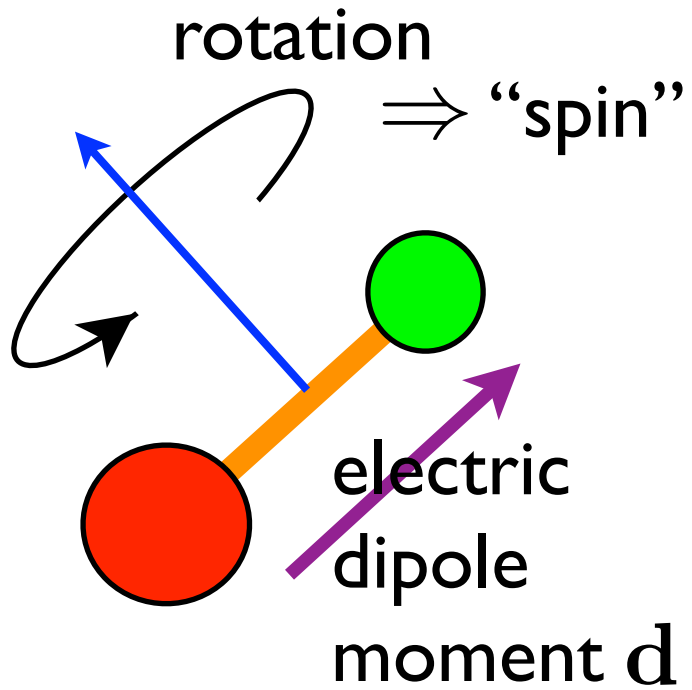


# Motivation





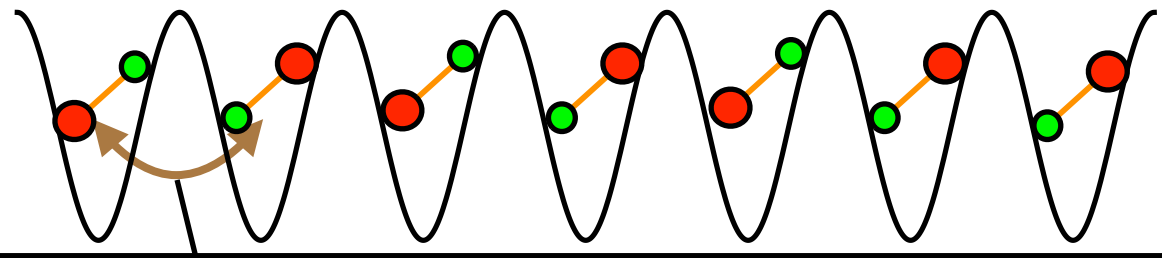
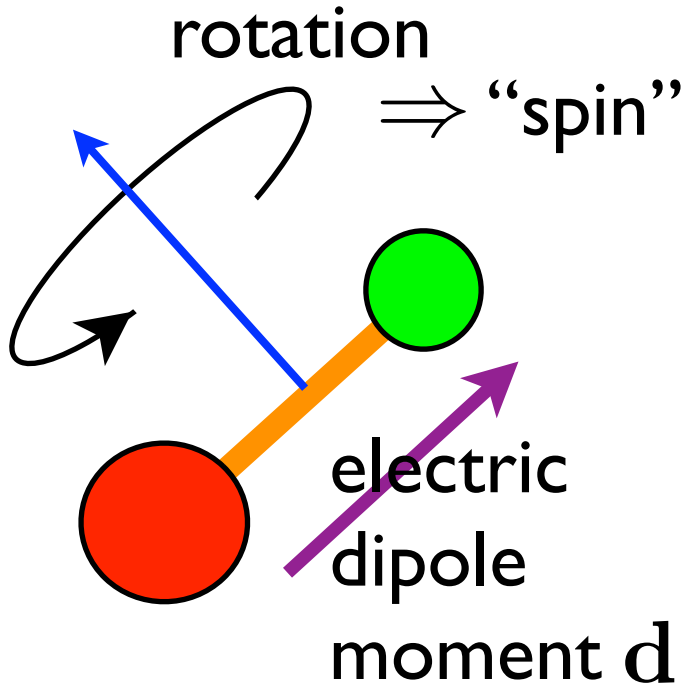
# Motivation



dipole-dipole interactions  $\Rightarrow$  "spin-spin" interactions

# Motivation

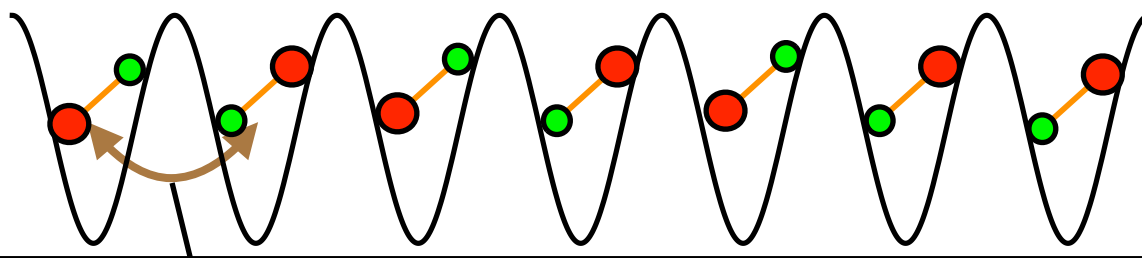
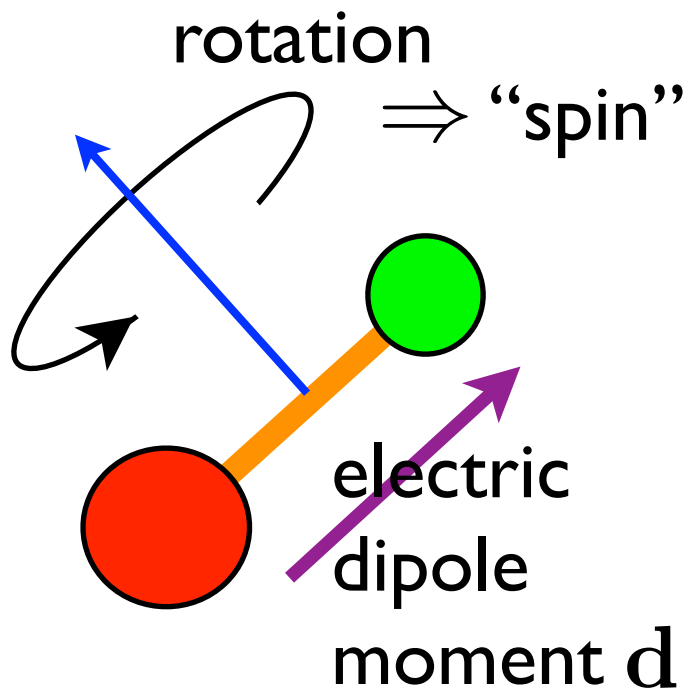
highly tunable exotic spin models



dipole-dipole interactions  $\Rightarrow$  "spin-spin" interactions

# Motivation

highly tunable exotic spin models



dipole-dipole interactions  $\Rightarrow$  "spin-spin" interactions

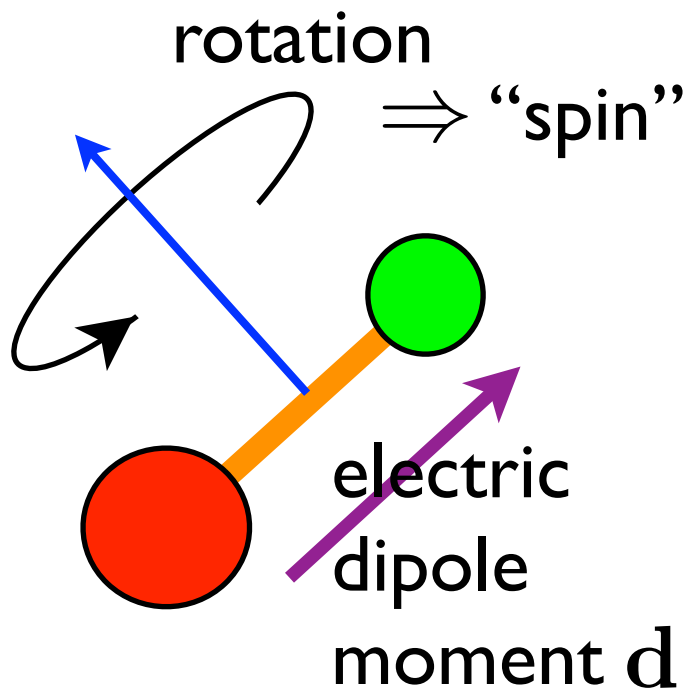
Seminal: Barnett, Petrov, Lukin, Demler PRL (2006); Micheli, Brennen, Zoller, Nat Phys (2006); Brennen, Micheli, Zoller, NJP (2007); Buchler, Micheli, Zoller, Nat Phys (2007)

- Wall, Carr, Schachenmayer, Daley, Perez-Rios, Herrera, Litinskaya, Krems, Lewenstein, Trefzger et al, Kestner, Das Sarma, et al, Lemeshko, Friedrich, Weimer, Ortner, Pupillo, Micheli, Zoller, Santos, Cooper, Shlyapnikov, Watanabe, Baranov, Rabl, Dalmonte, Zhou, Bekaroglu, Buchler, Peter et al, etc...

Review: Baranov, Dalmonte, Pupillo, and Zoller, Chem. Rev. 112, 5012 (2012).

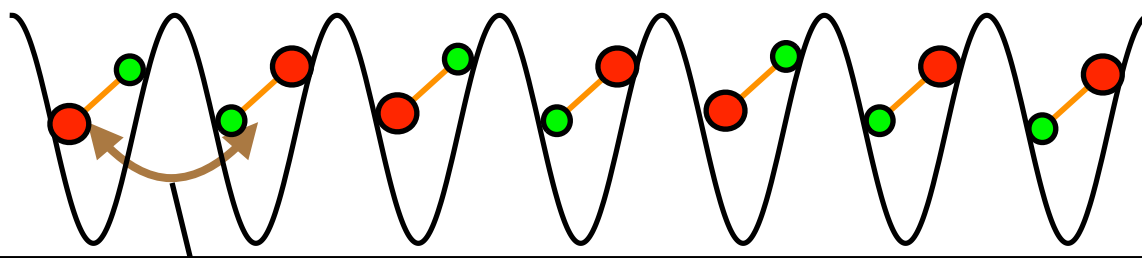
# Motivation

highly tunable exotic spin models



## Our achievements:

- stronger interactions
- much higher tunability



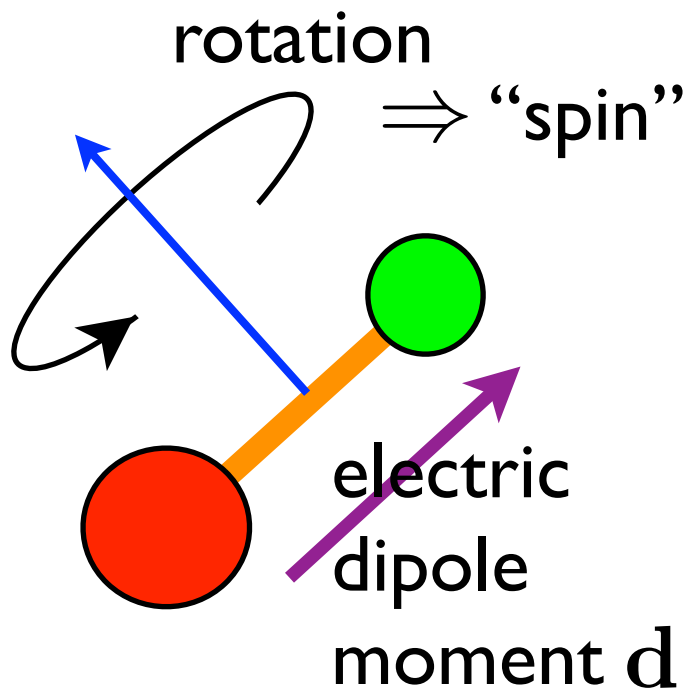
dipole-dipole interactions  $\Rightarrow$  "spin-spin" interactions

Seminal: Barnett, Petrov, Lukin, Demler PRL (2006); Micheli, Brennen, Zoller, Nat Phys (2006); Brennen, Micheli, Zoller, NJP (2007); Buchler, Micheli, Zoller, Nat Phys (2007)

- Wall, Carr, Schachenmayer, Daley, Perez-Rios, Herrera, Litinskaya, Krems, Lewenstein, Trefzger et al, Kestner, Das Sarma, et al, Lemeshko, Friedrich, Weimer, Ortner, Pupillo, Micheli, Zoller, Santos, Cooper, Shlyapnikov, Watanabe, Baranov, Rabl, Dalmonte, Zhou, Bekaroglu, Buchler, Peter et al, etc...

Review: Baranov, Dalmonte, Pupillo, and Zoller, Chem. Rev. 112, 5012 (2012).

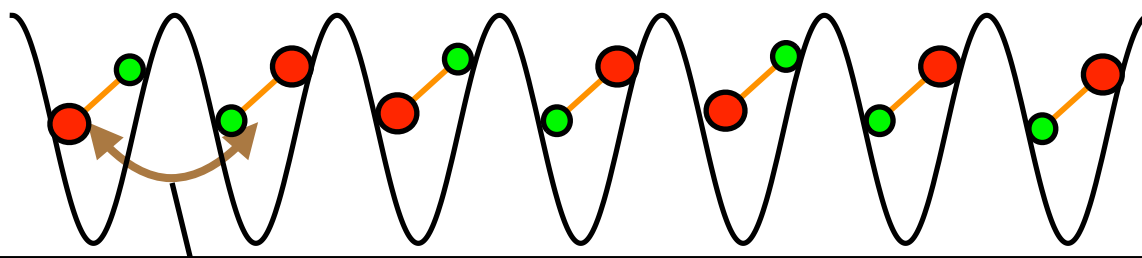
# Motivation



highly tunable exotic spin models

## Our achievements:

- stronger interactions
- much higher tunability  $\Rightarrow$  exotic physics



dipole-dipole interactions  $\Rightarrow$  "spin-spin" interactions

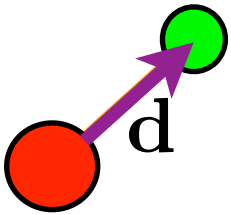
Seminal: Barnett, Petrov, Lukin, Demler PRL (2006); Micheli, Brennen, Zoller, Nat Phys (2006); Brennen, Micheli, Zoller, NJP (2007); Buchler, Micheli, Zoller, Nat Phys (2007)

- Wall, Carr, Schachenmayer, Daley, Perez-Rios, Herrera, Litinskaya, Krems, Lewenstein, Trefzger et al, Kestner, Das Sarma, et al, Lemeshko, Friedrich, Weimer, Ortner, Pupillo, Micheli, Zoller, Santos, Cooper, Shlyapnikov, Watanabe, Baranov, Rabl, Dalmonte, Zhou, Bekaroglu, Buchler, Peter et al, etc...

Review: Baranov, Dalmonte, Pupillo, and Zoller, Chem. Rev. 112, 5012 (2012).

# Rigid rotor

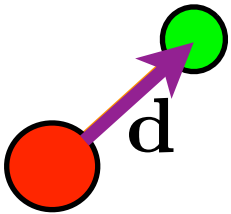
$$H_0 = B\mathbf{N}^2$$



# Rigid rotor

$$H_0 = B\mathbf{N}^2$$

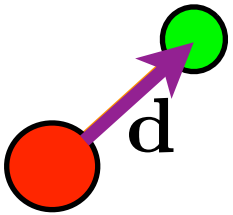
$$\mathbf{N}^2|N, N_z\rangle = N(N + 1)|N, N_z\rangle$$



# Rigid rotor

$$H_0 = B\mathbf{N}^2$$

$$\mathbf{N}^2|N, N_z\rangle = N(N+1)|N, N_z\rangle$$

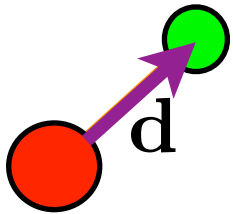


$$\begin{array}{cccc} N_z = & -1 & 0 & 1 \\ 2B \left\{ \begin{array}{l} \text{---} \quad \text{---} \quad \text{---} \\ \text{---} \end{array} \right. & & & \begin{array}{l} N = 1 \\ \\ N = 0 \end{array} \end{array}$$



# Rigid rotor

$$\mathbf{E} = E\hat{z}$$



$$H_0 = B\mathbf{N}^2$$

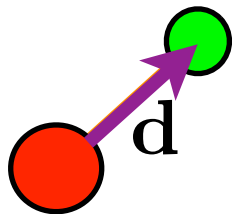
$$\mathbf{N}^2|N, N_z\rangle = N(N+1)|N, N_z\rangle$$

$$N_z = -1 \quad 0 \quad 1$$

$$2B \left\{ \begin{array}{ccc} \text{---} & \text{---} & \text{---} \\ & \text{---} & \\ \end{array} \right. \begin{array}{l} N = 1 \\ \\ N = 0 \end{array}$$

# Rigid rotor

$$\mathbf{E} = E\hat{z}$$



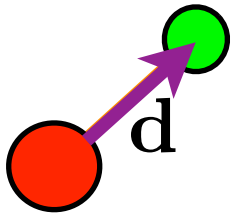
$$H_0 = B\mathbf{N}^2 - d^z E$$

$$\mathbf{N}^2|N, N_z\rangle = N(N+1)|N, N_z\rangle$$

$$\begin{array}{ccc} N_z = & -1 & 0 & 1 \\ 2B \left\{ \begin{array}{l} \text{---} & \text{---} & \text{---} & N = 1 \\ & \text{---} & & N = 0 \end{array} \right. \end{array}$$

# Rigid rotor

$$\mathbf{E} = E\hat{z}$$



$$H_0 = B\mathbf{N}^2 - d^z E$$

$$\mathbf{N}^2|N, N_z\rangle = N(N+1)|N, N_z\rangle$$

$N_z =$	-1	0	1	
		—		
{	—		—	$N = 1$
		—		$N = 0$

# Rigid rotor

$$\mathbf{E} = E\hat{\mathbf{z}}$$



$$N_z = \begin{array}{ccc} -1 & 0 & 1 \\ \hline & & \hline & & \\ & \hline & & \hline \end{array}$$

# Rigid rotor

$$\mathbf{E} = E\hat{z}$$



$$N_z = -1 \quad 0 \quad 1$$



# Simplest spin Hamiltonian

$$\mathbf{E} = E\hat{\mathbf{z}}$$

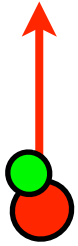


$$N_z = -1 \quad 0 \quad 1$$



# Simplest spin Hamiltonian

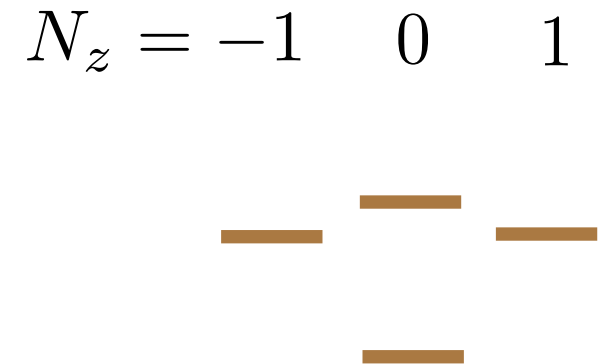
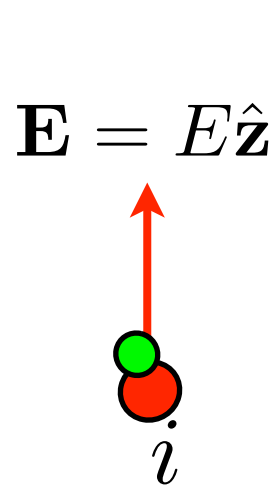
$$\mathbf{E} = E\hat{\mathbf{z}}$$



$$N_z = -1 \quad 0 \quad 1$$

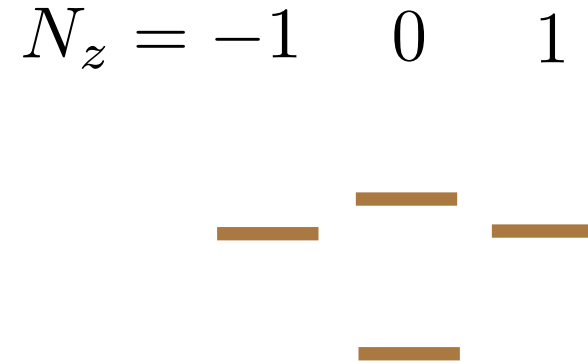
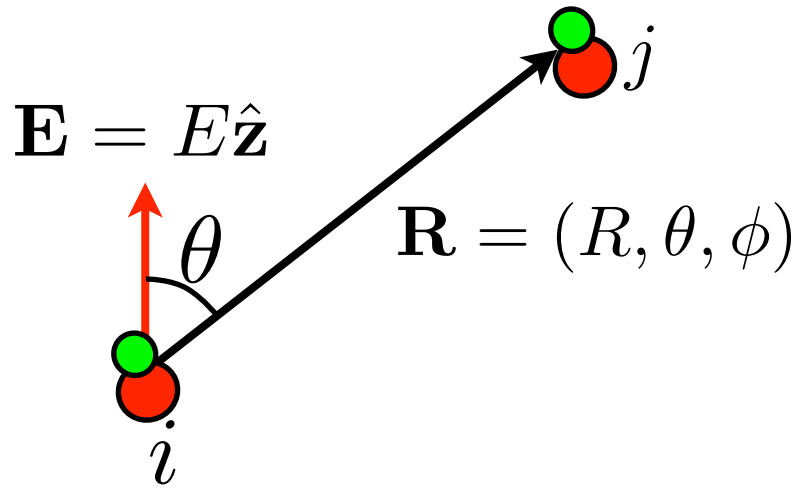


# Simplest spin Hamiltonian

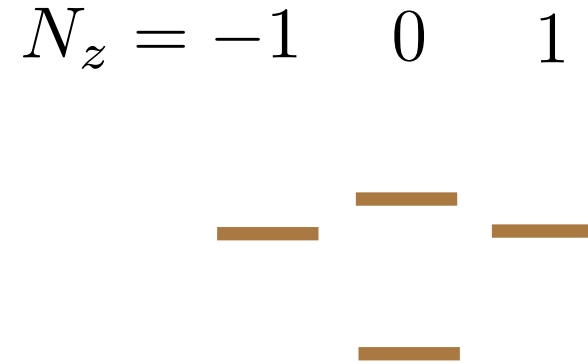
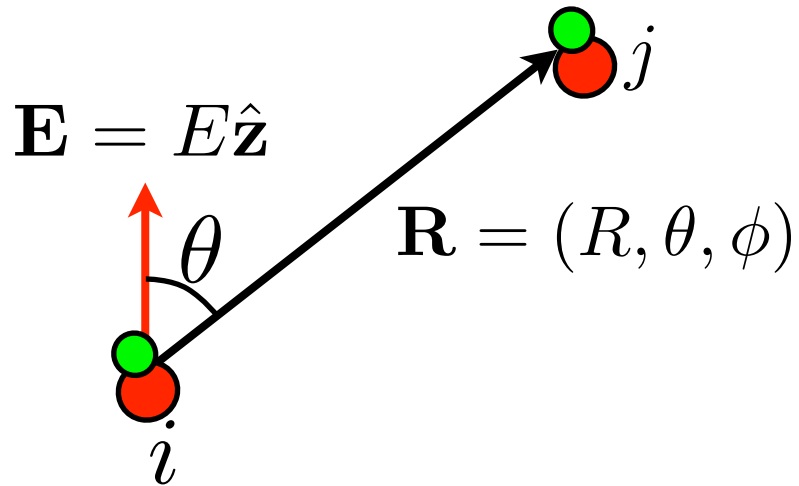




# Simplest spin Hamiltonian

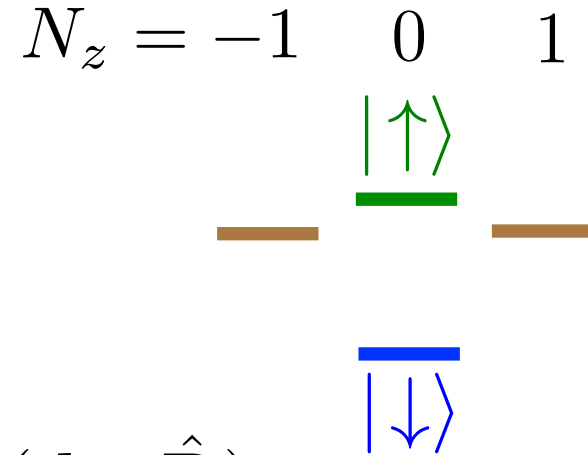
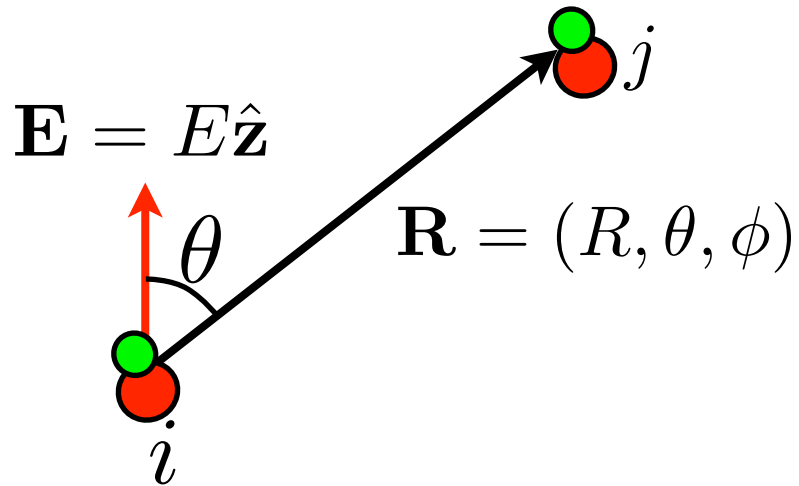


# Simplest spin Hamiltonian



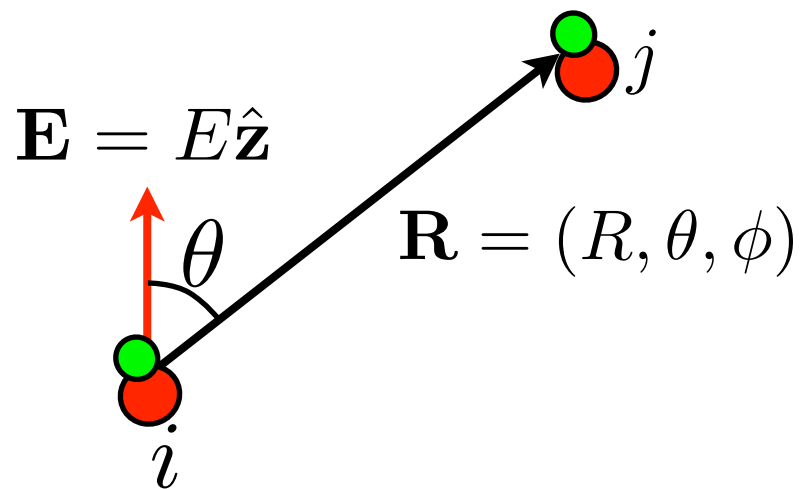
$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

# Simplest spin Hamiltonian



$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

# Simplest spin Hamiltonian



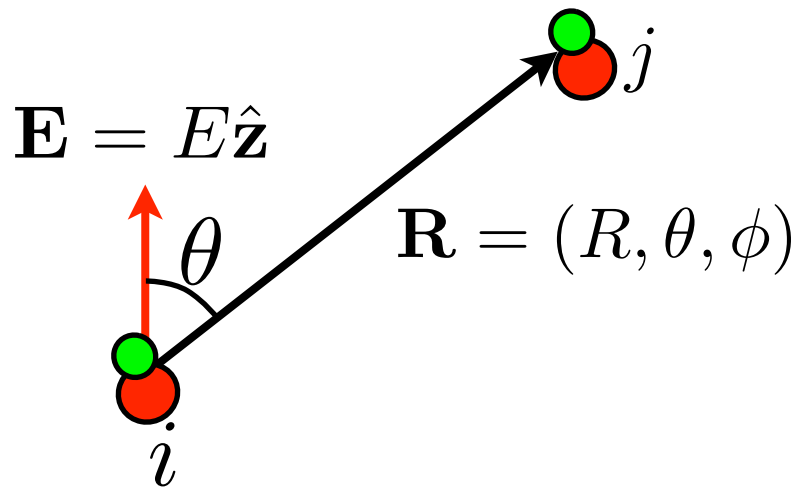
$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

$$N_z = -1 \quad 0 \quad 1$$

$$|\downarrow\rangle$$

$H_{dd}$   
weak

# Simplest spin Hamiltonian



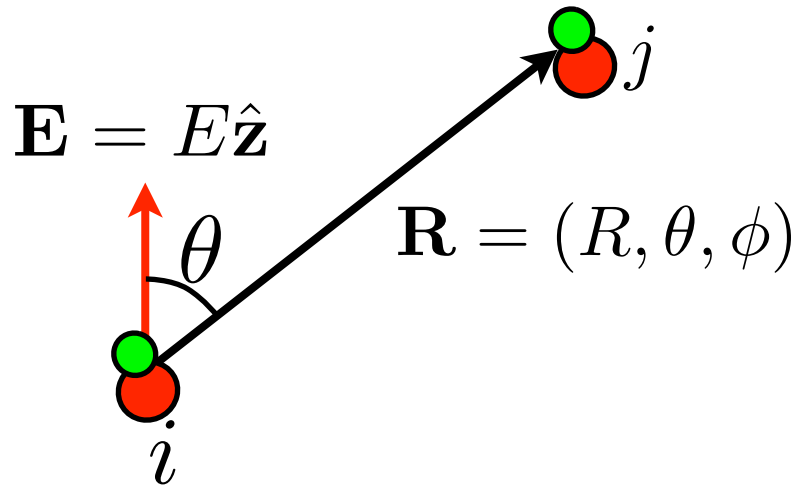
$$N_z = -1 \quad 0 \quad 1$$

$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

$H_{dd}$   
weak

- project on  $|\uparrow\rangle, |\downarrow\rangle$

# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$

$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

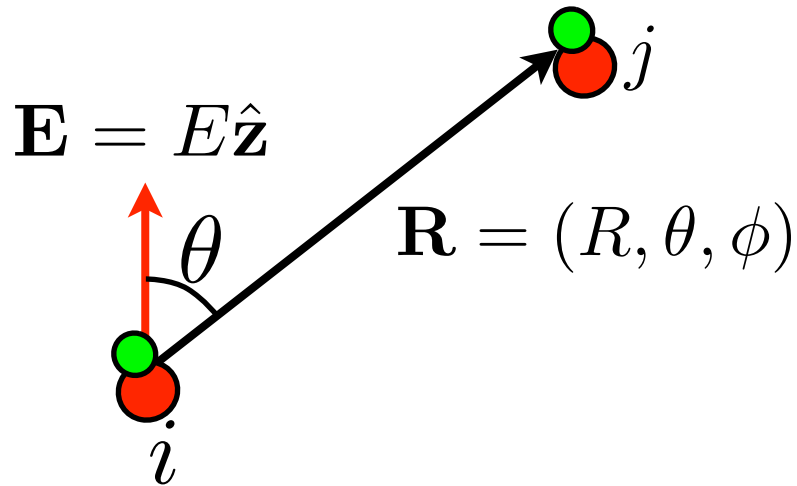
$H_{dd}$   
weak

- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + J_{xy} (S_i^x S_j^x + S_i^y S_j^y)]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$

$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

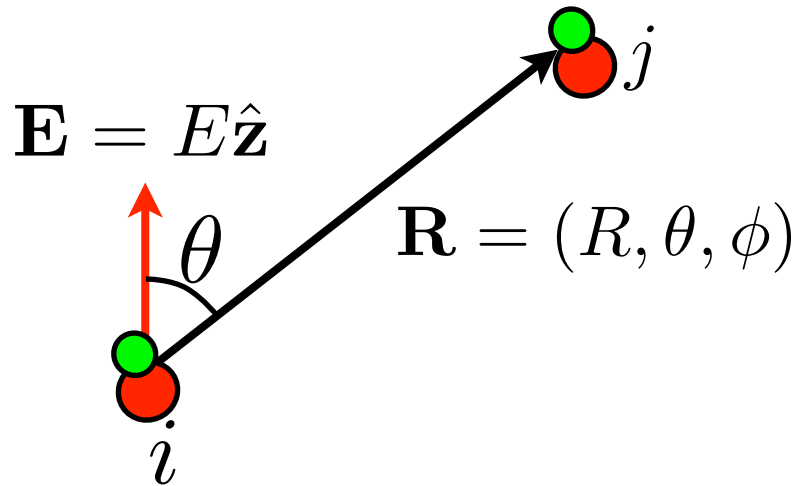
$H_{dd}$   
weak

- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z (S_i^z S_j^z + S_i^x S_j^x + S_i^y S_j^y)]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$

$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

$H_{dd}$   
weak

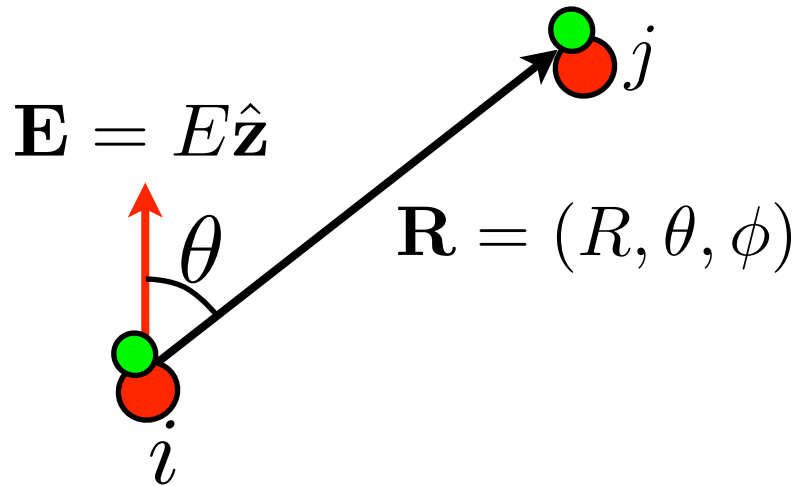
- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + J_{xy} (S_i^x S_j^x + S_i^y S_j^y)]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$



# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$

—
—
—

$$|\downarrow\rangle$$

$H_{dd}$   
 weak

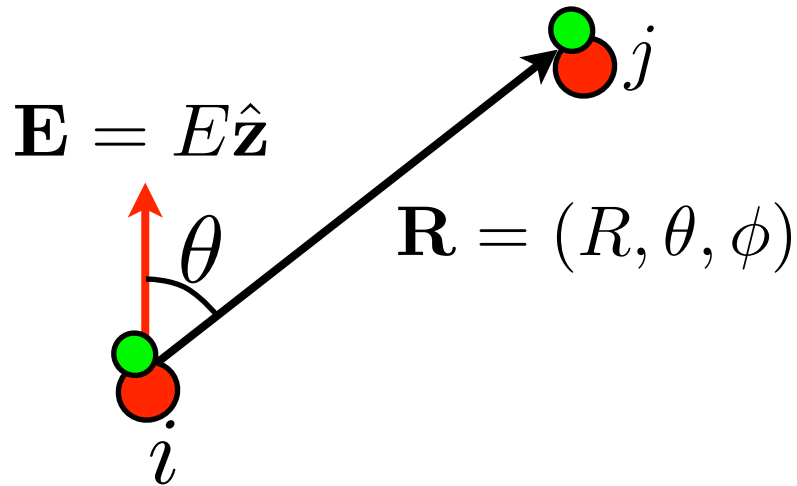
$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + J_{xy} (S_i^x S_j^x + S_i^y S_j^y)]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$

Energy level diagram showing the  $N_z$  quantum numbers. The  $N_z = 0$  level is split into  $|\uparrow\rangle$  (green) and  $|\downarrow\rangle$  (blue) states. The  $N_z = -1$  and  $N_z = 1$  levels are shown as single lines.

$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

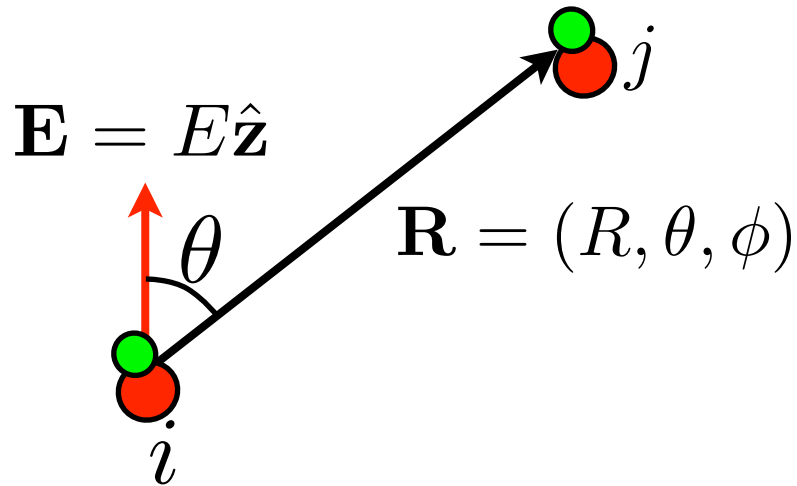
$H_{dd}$   
weak

- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + J_{xy} (S_i^x S_j^x + S_i^y S_j^y)]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$

—
—
—

$|\downarrow\rangle$

$H_{dd}$   
weak

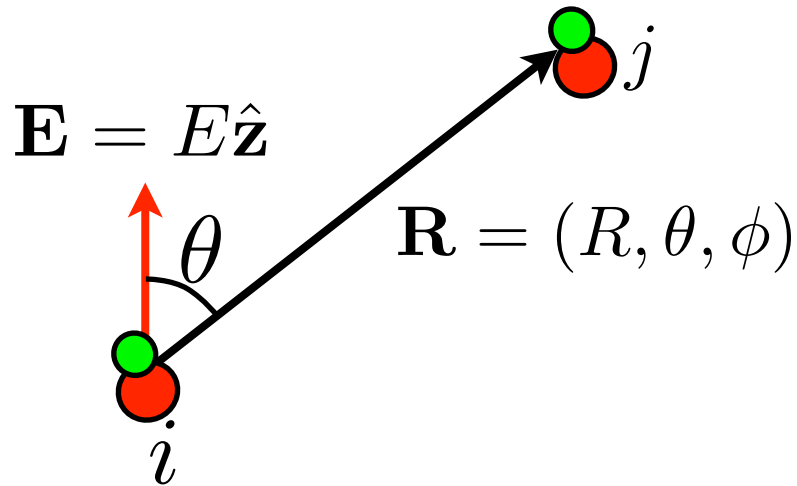
$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

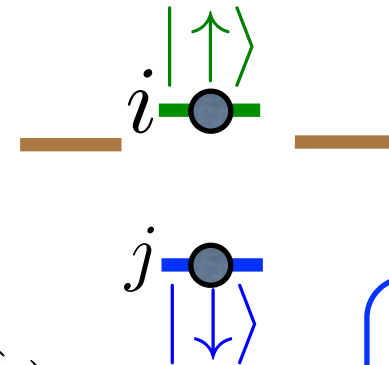
$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + \underbrace{J_{xy} (S_i^x S_j^x + S_i^y S_j^y)}_{\text{circled}}]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$



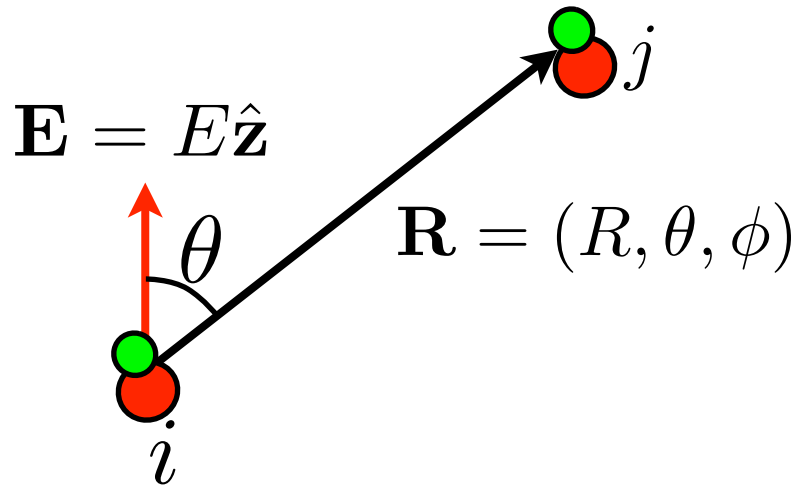
$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

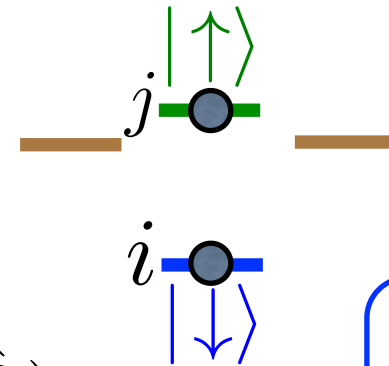
$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + \underbrace{J_{xy} (S_i^x S_j^x + S_i^y S_j^y)}_{\text{circled in red}}]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$



$H_{dd}$   
weak

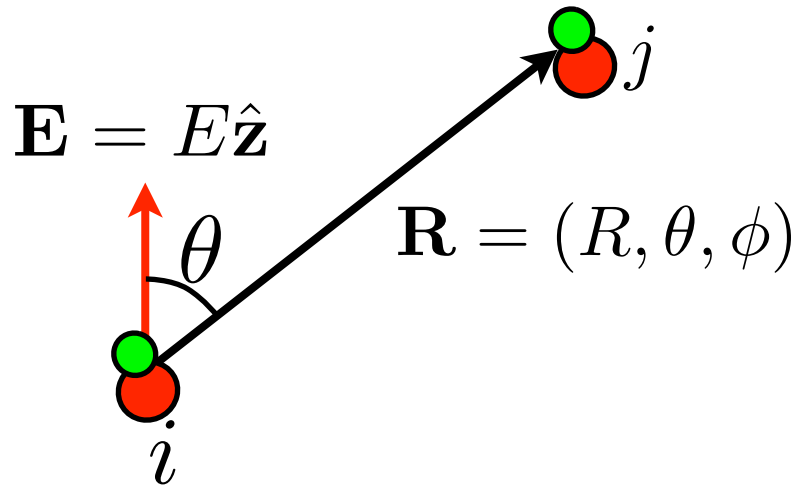
$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

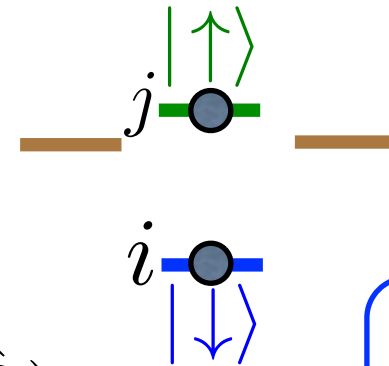
$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + \underbrace{J_{xy} (S_i^x S_j^x + S_i^y S_j^y)}_{\text{circled}}]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$



$H_{dd}$   
weak

$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

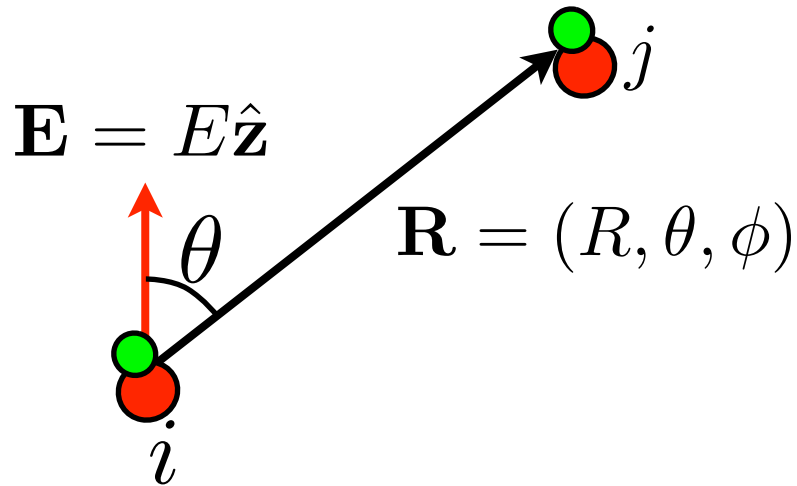
- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + J_{xy} (\overbrace{S_i^x S_j^x + S_i^y S_j^y}^{S_i^+ S_j^- + S_i^- S_j^+})]$$

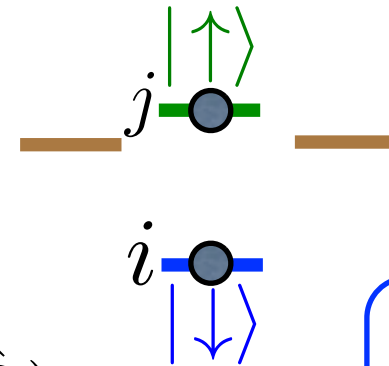
$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

- (• if add hopping  $\Rightarrow$  t-J model  $\Rightarrow$  d-wave superfluidity?)

# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$



$H_{dd}$   
weak

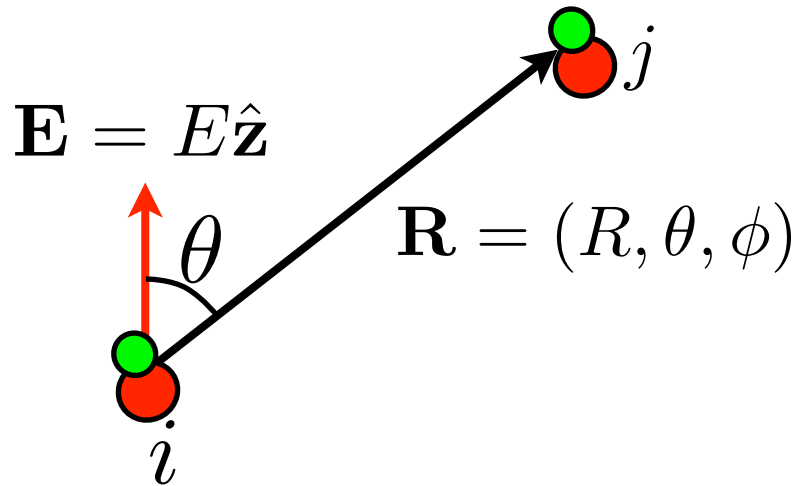
$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

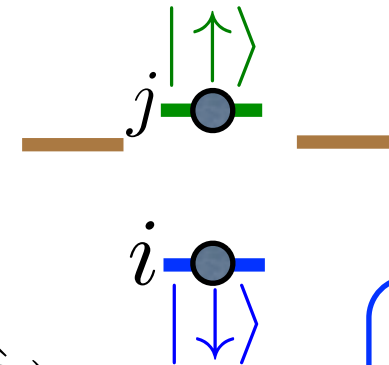
$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + J_{xy} (\overbrace{S_i^x S_j^x + S_i^y S_j^y}^{S_i^+ S_j^- + S_i^- S_j^+})]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$



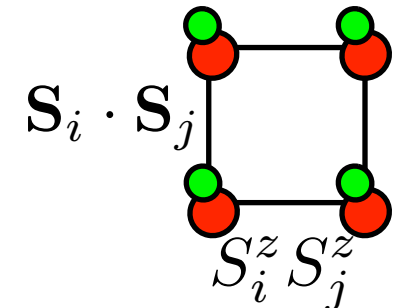
$H_{dd}$   
weak

$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

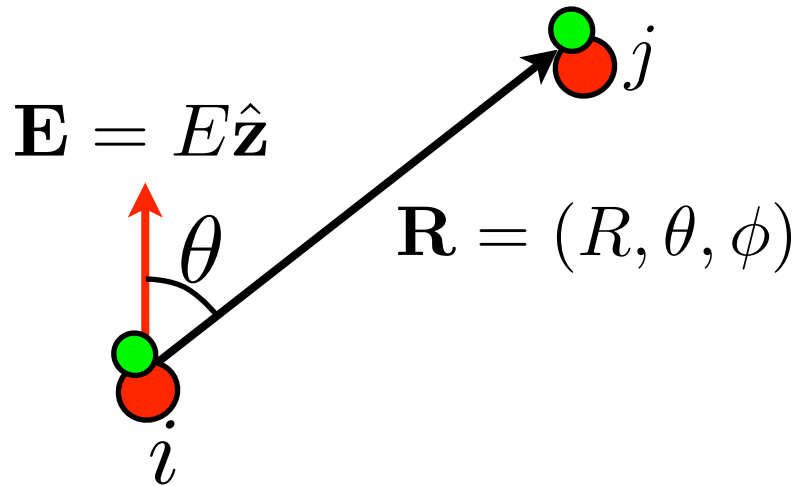
$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + J_{xy} (S_i^x S_j^x + S_i^y S_j^y)]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

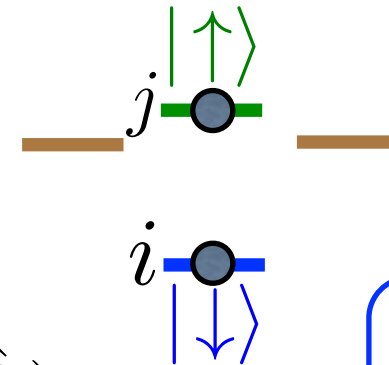




# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$



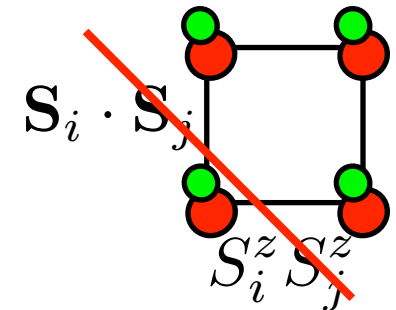
$H_{dd}$   
weak

$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

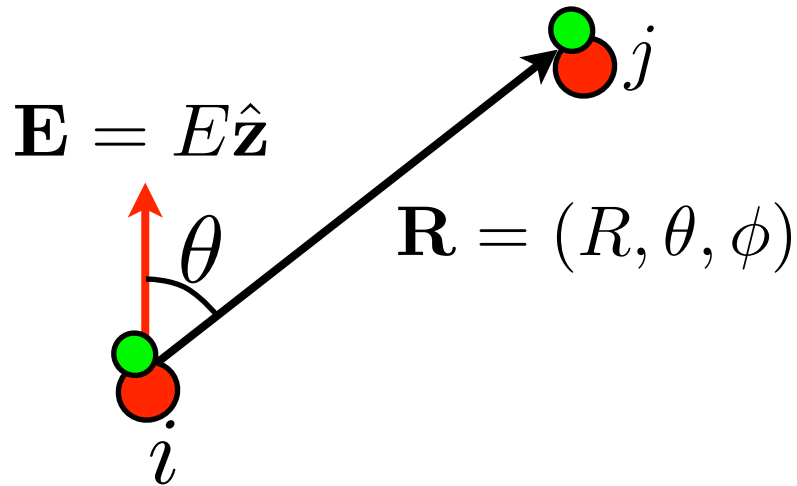
- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + J_{xy} (S_i^x S_j^x + S_i^y S_j^y)]$$

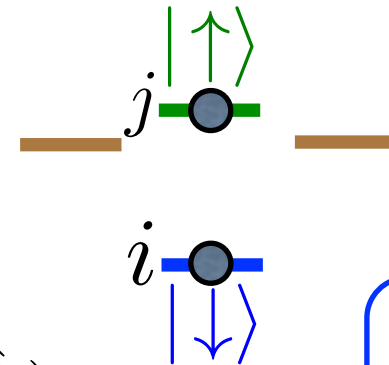
$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$



# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$



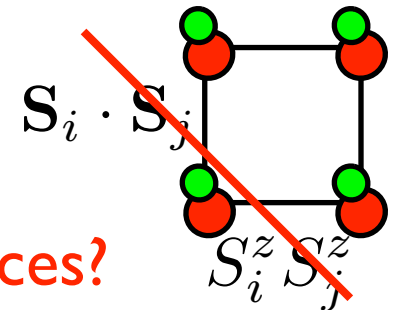
$H_{dd}$   
weak

$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

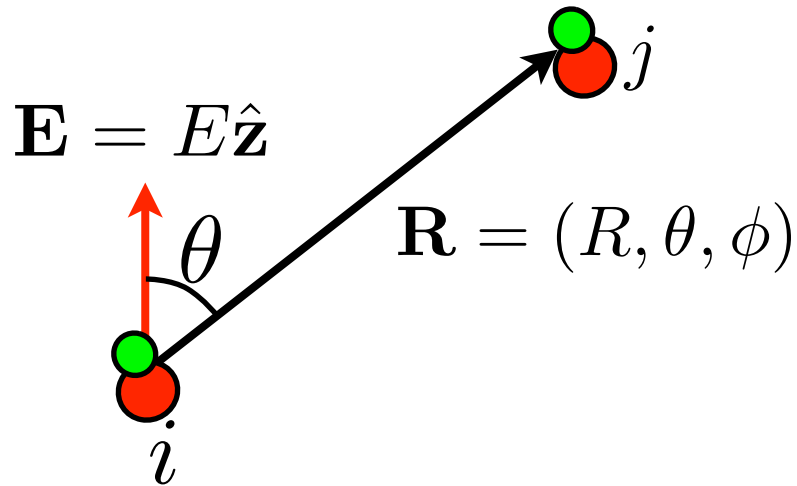
$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + J_{xy} (S_i^x S_j^x + S_i^y S_j^y)]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

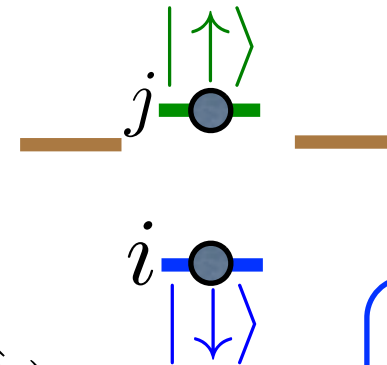


- can  $J_z$  and  $J_{xy}$  come with different angular dependences?

# Simplest spin Hamiltonian



$$N_z = -1 \quad 0 \quad 1$$



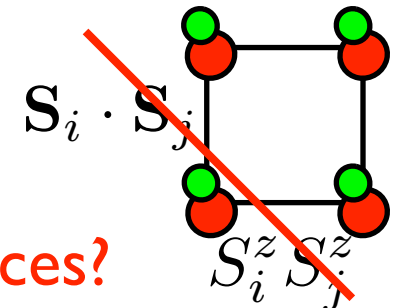
$H_{dd}$   
weak

$$R^3 H_{dd} = \mathbf{d}_i \cdot \mathbf{d}_j - 3(\mathbf{d}_i \cdot \hat{\mathbf{R}})(\mathbf{d}_j \cdot \hat{\mathbf{R}})$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = Y_{2,0}(\theta) [J_z S_i^z S_j^z + J_{xy} (S_i^x S_j^x + S_i^y S_j^y)]$$

$$Y_{2,0} \propto 1 - 3 \cos^2 \theta$$

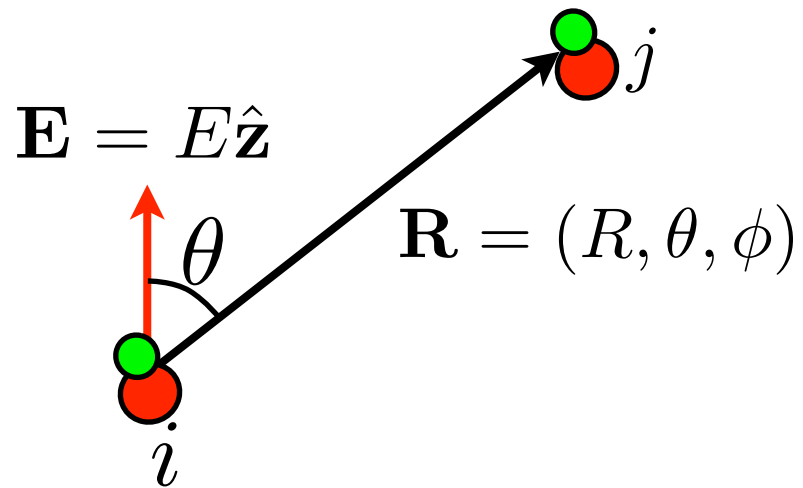


- can  $J_z$  and  $J_{xy}$  come with different angular dependences?

**YES!**

AVG et al, PRL, PRA 2011; Kuns et al, PRA 2011

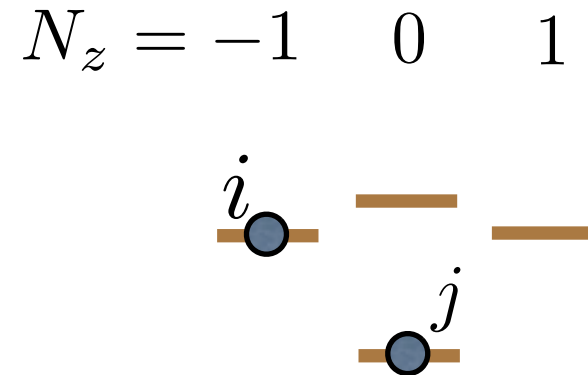
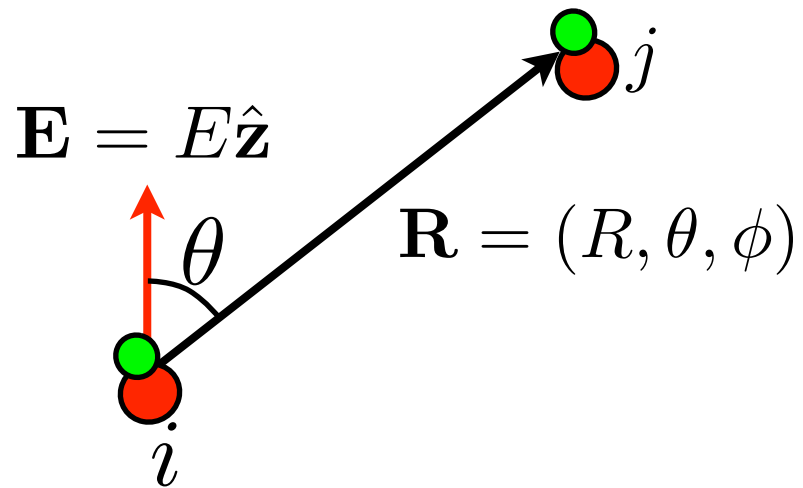
# Spin Hamiltonian from dressed states



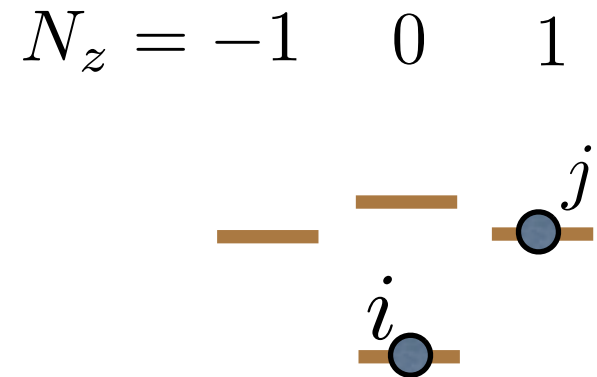
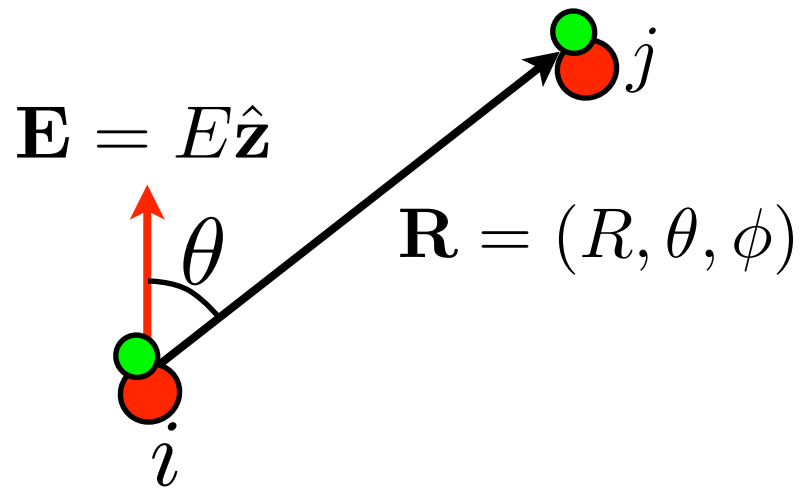
$$N_z = -1 \quad 0 \quad 1$$



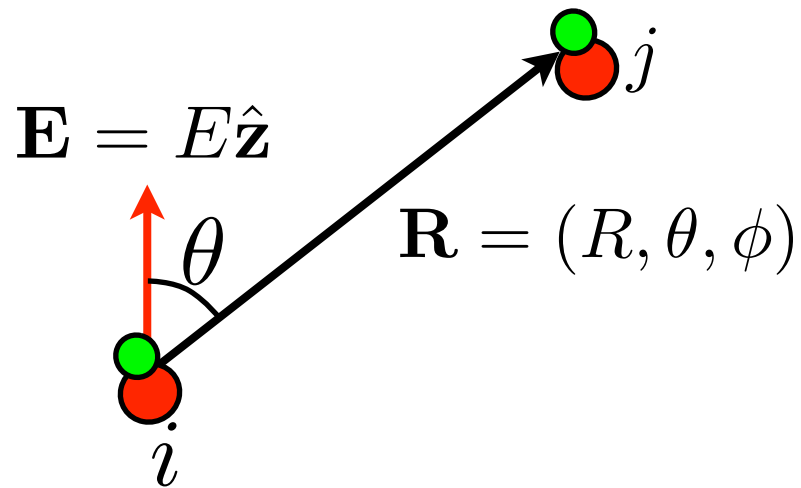
# Spin Hamiltonian from dressed states



# Spin Hamiltonian from dressed states

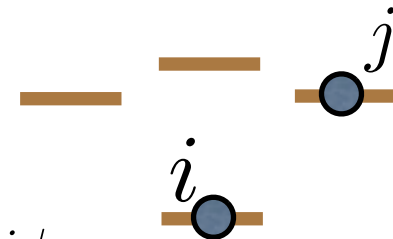


# Spin Hamiltonian from dressed states

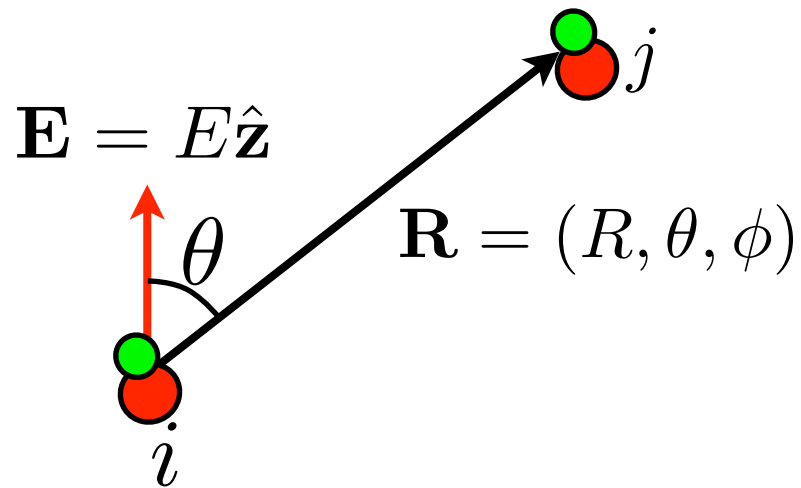


$$N_z = -1 \quad 0 \quad 1$$

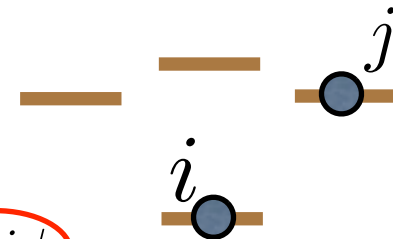
$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi}$$



# Spin Hamiltonian from dressed states



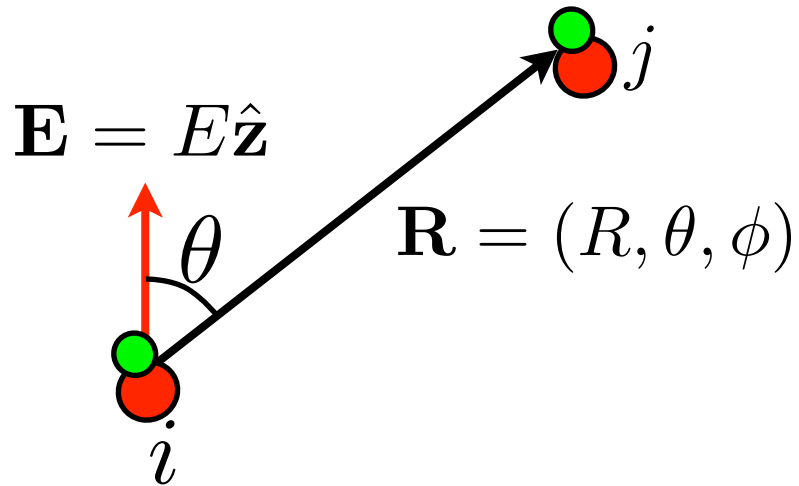
$$N_z = -1 \quad 0 \quad 1$$



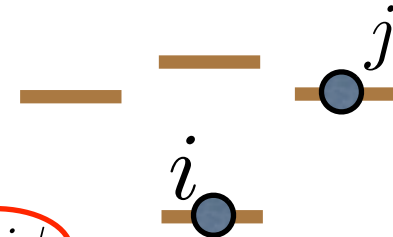
$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi}$$



# Spin Hamiltonian from dressed states



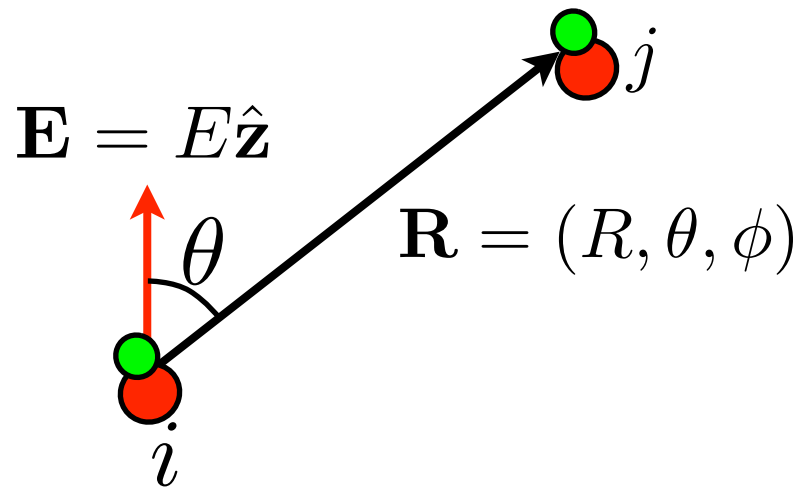
$$N_z = -1 \quad 0 \quad 1$$



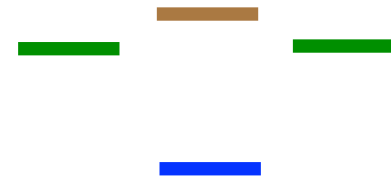
$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi}$$

[Einstein, de Haas effect]

# Spin Hamiltonian from dressed states



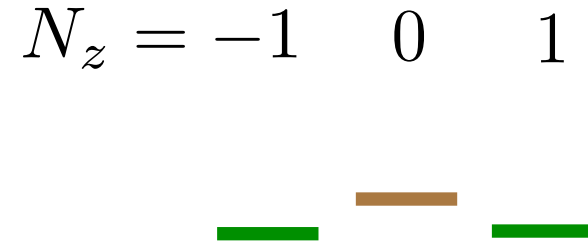
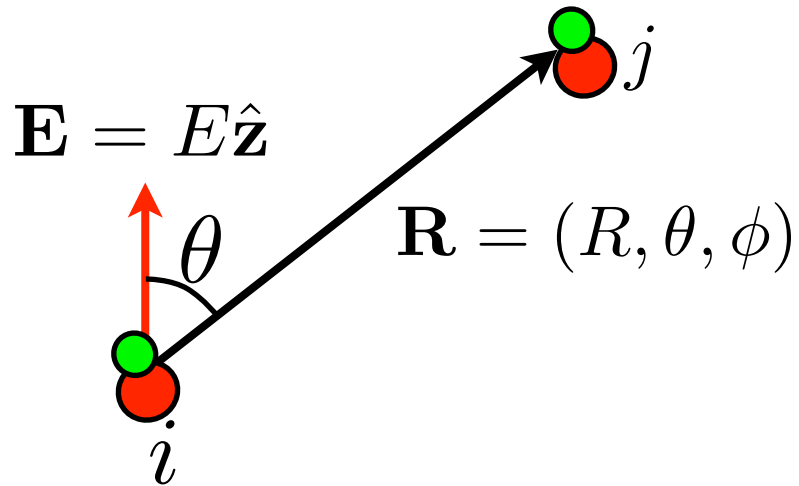
$$N_z = -1 \quad 0 \quad 1$$



$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi}$$

[Einstein, de Haas effect]

# Spin Hamiltonian from dressed states

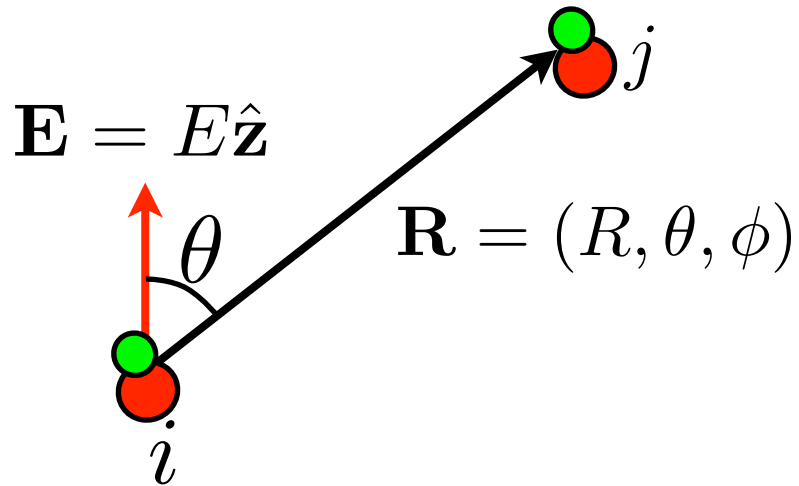


$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi}$$



[Einstein, de Haas effect]

# Spin Hamiltonian from dressed states



$$N_z = -1 \quad 0 \quad 1$$

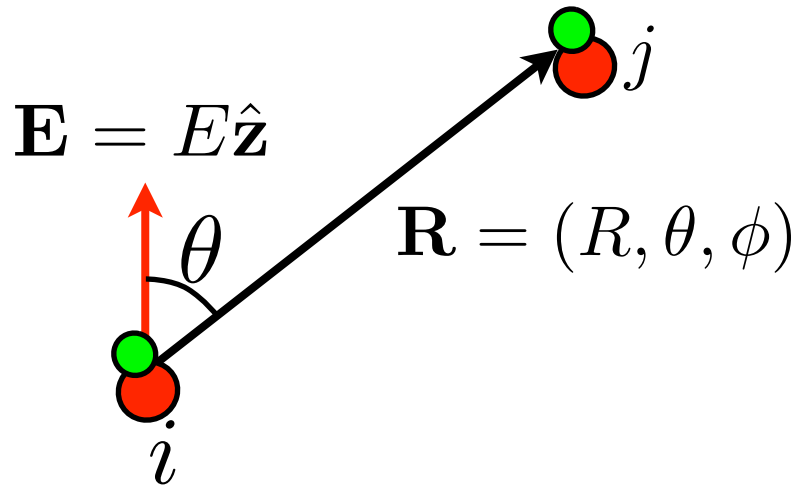
$$|1\rangle \text{---} \text{---} \text{---} |2\rangle$$

$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi}$$

$$|\downarrow\rangle$$

[Einstein, de Haas effect]

# Spin Hamiltonian from dressed states



$$|\uparrow\rangle = \Omega_2|1\rangle - \Omega_1|2\rangle$$

$$N_z = -1 \quad 0 \quad 1$$

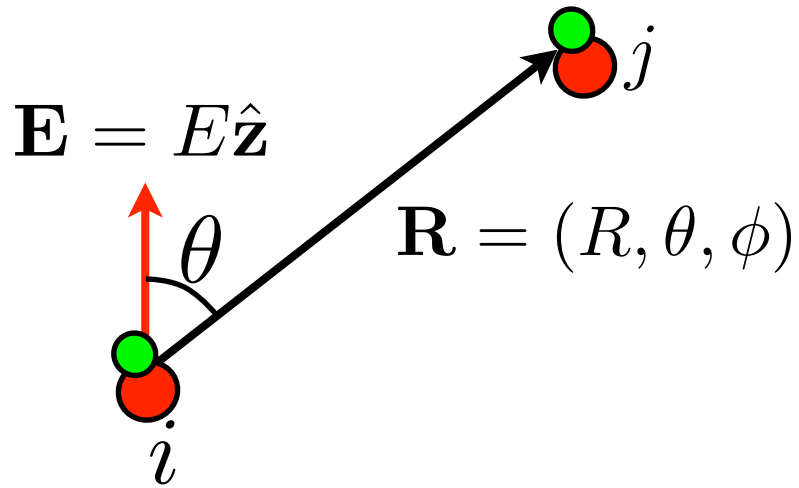
$$|1\rangle \text{---} \text{---} \text{---} |2\rangle$$

$$Y_{2,-2} \propto \sin^2\theta e^{-2i\phi}$$



[Einstein, de Haas effect]

# Spin Hamiltonian from dressed states



$$|\uparrow\rangle = \Omega_2|1\rangle - \Omega_1|2\rangle$$

$$N_z = -1 \quad 0 \quad 1$$

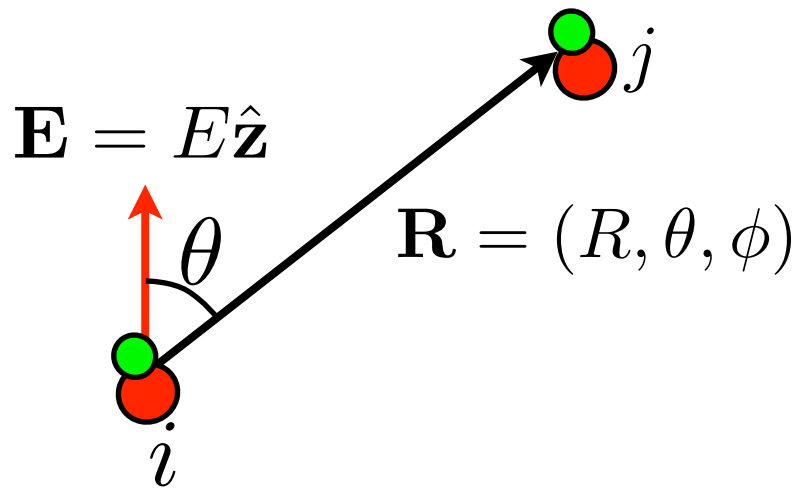
$$|1\rangle \text{---} \text{---} \text{---} |2\rangle$$

$$Y_{2,-2} \propto \sin^2\theta e^{-2i\phi}$$

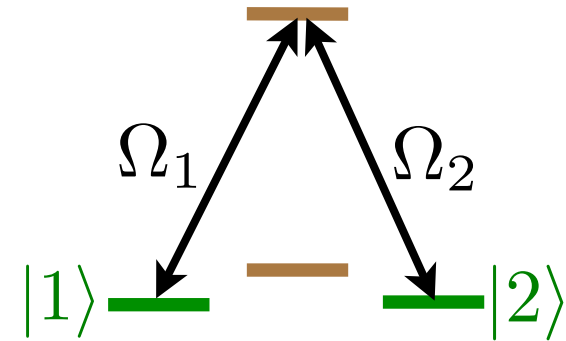


[Einstein, de Haas effect]

# Spin Hamiltonian from dressed states



$$|\uparrow\rangle = \Omega_2|1\rangle - \Omega_1|2\rangle$$

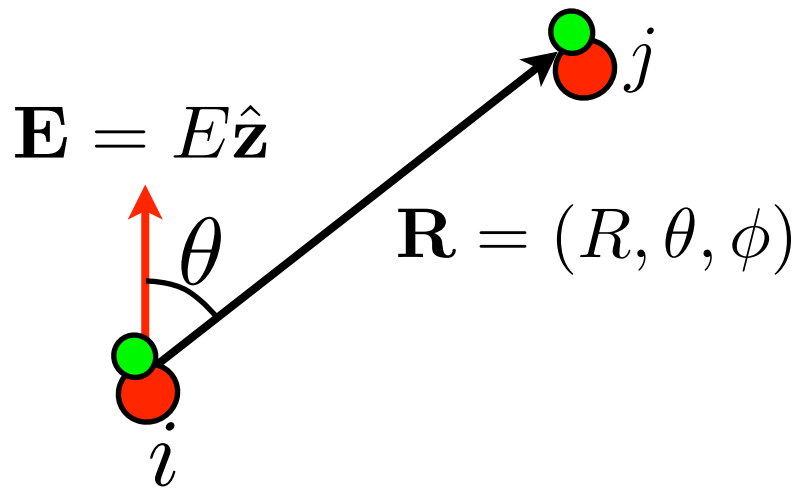


$$Y_{2,-2} \propto \sin^2\theta e^{-2i\phi}$$

$\overline{|\downarrow\rangle}$

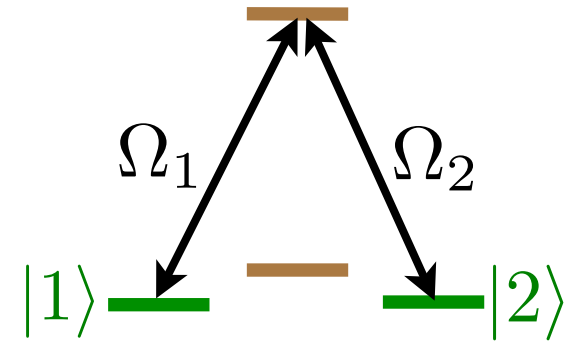
[Einstein, de Haas effect]

# Spin Hamiltonian from dressed states



$$|\uparrow\rangle = \Omega_2|1\rangle - \Omega_1|2\rangle$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

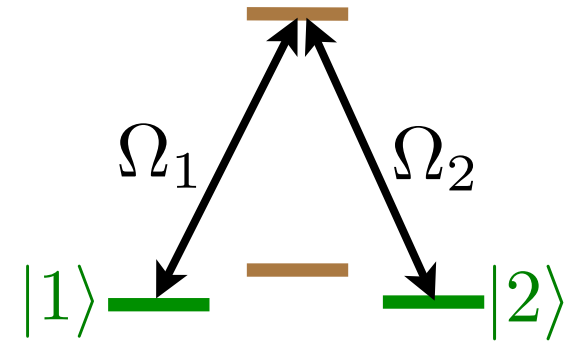
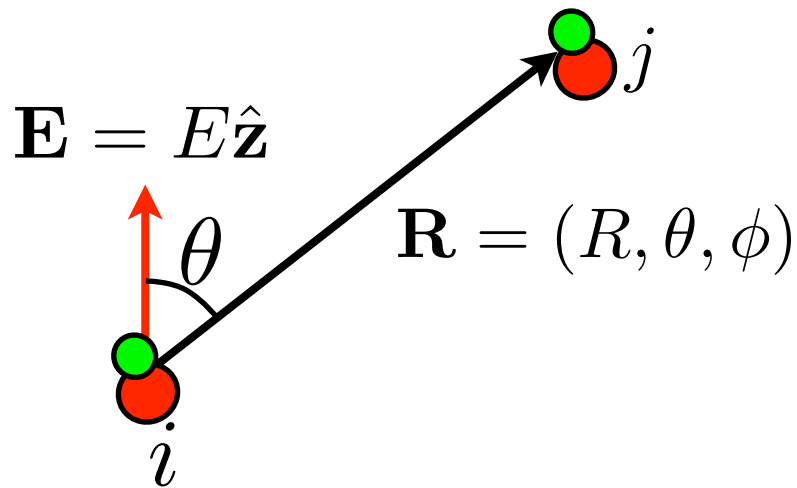


$$Y_{2,-2} \propto \sin^2\theta e^{-2i\phi} \quad \overline{|\downarrow\rangle}$$

[Einstein, de Haas effect]



# Spin Hamiltonian from dressed states



$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi} \quad \overline{|\downarrow\rangle}$$

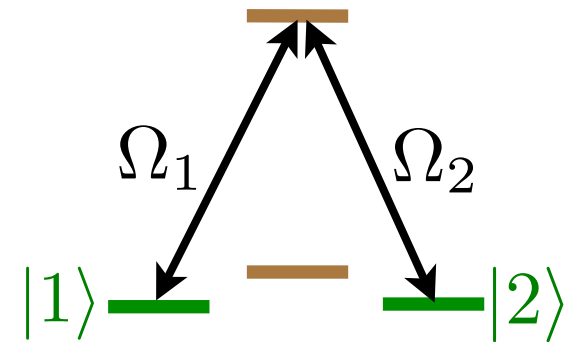
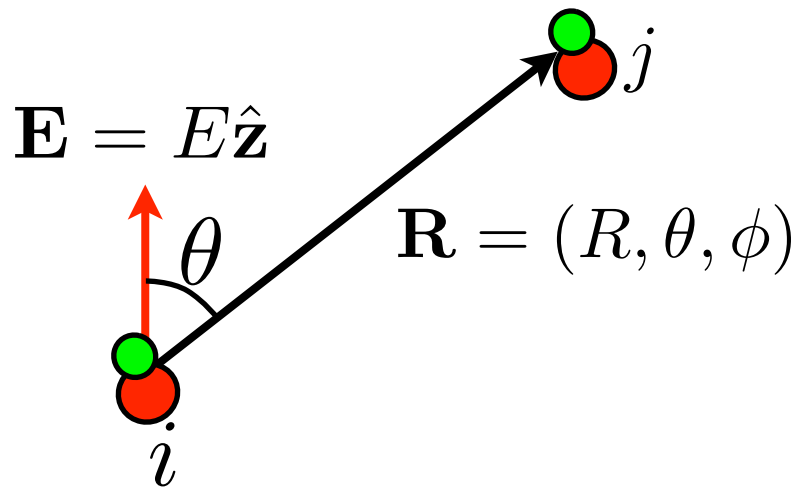
[Einstein, de Haas effect]

$$|\uparrow\rangle = \Omega_2 |1\rangle - \Omega_1 |2\rangle$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = J_z Y_{2,0} S_i^z S_j^z + (J_{xy} Y_{2,0} + \text{Re}[J'_{xy} Y_{2,-2}]) \frac{1}{2} (S_i^+ S_j^- + S_i^- S_j^+)$$

# Spin Hamiltonian from dressed states



$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi} \quad \overline{|\downarrow\rangle}$$

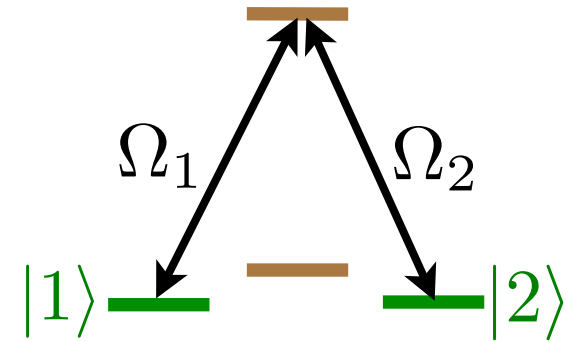
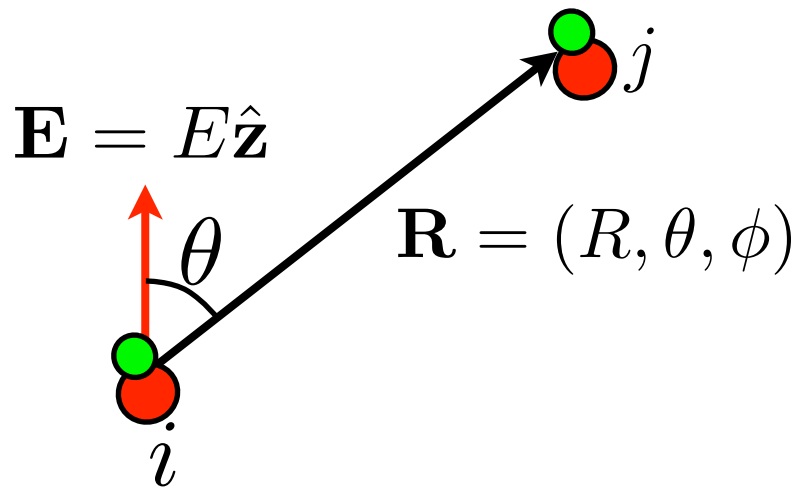
[Einstein, de Haas effect]

$$|\uparrow\rangle = \Omega_2 |1\rangle - \Omega_1 |2\rangle$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = \boxed{J_z Y_{2,0} S_i^z S_j^z} + \boxed{J_{xy} Y_{2,0}} + \text{Re}[J'_{xy} Y_{2,-2}] \frac{1}{2} (S_i^+ S_j^- + S_i^- S_j^+)$$

# Spin Hamiltonian from dressed states



$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi} \quad \overline{|\downarrow\rangle}$$

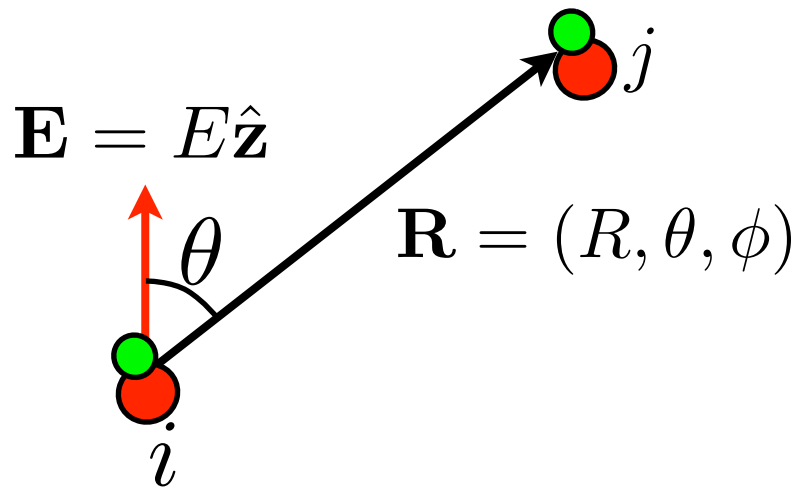
[Einstein, de Haas effect]

$$|\uparrow\rangle = \Omega_2 |1\rangle - \Omega_1 |2\rangle$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = J_z Y_{2,0} S_i^z S_j^z + (J_{xy} Y_{2,0} + \text{Re}[J'_{xy} Y_{2,-2}]) \frac{1}{2} (S_i^+ S_j^- + S_i^- S_j^+)$$

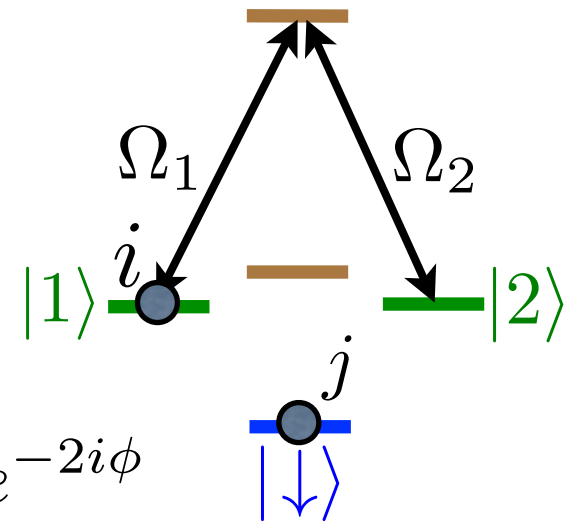
# Spin Hamiltonian from dressed states



$$|\uparrow\rangle = \Omega_2|1\rangle - \Omega_1|2\rangle$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

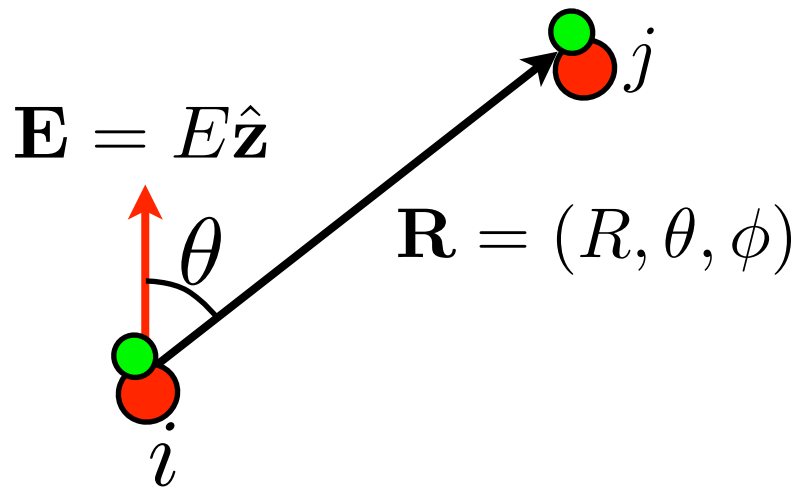
$$R^3 H_{dd} = J_z Y_{2,0} S_i^z S_j^z + (J_{xy} Y_{2,0} + \text{Re}[J'_{xy} Y_{2,-2}]) \frac{1}{2} (S_i^+ S_j^- + S_i^- S_j^+)$$



$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi}$$

[Einstein, de Haas effect]

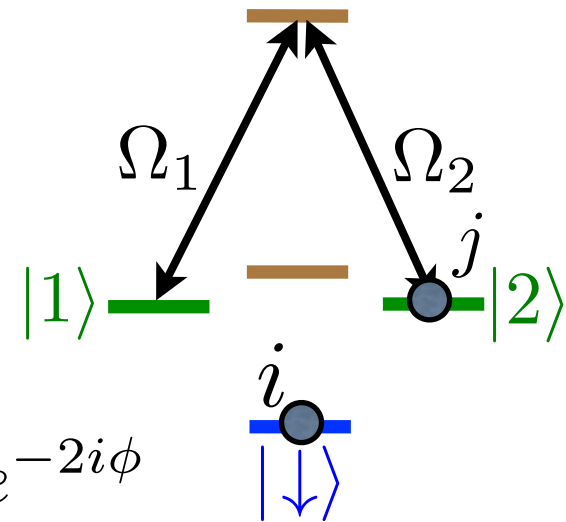
# Spin Hamiltonian from dressed states



$$|\uparrow\rangle = \Omega_2|1\rangle - \Omega_1|2\rangle$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

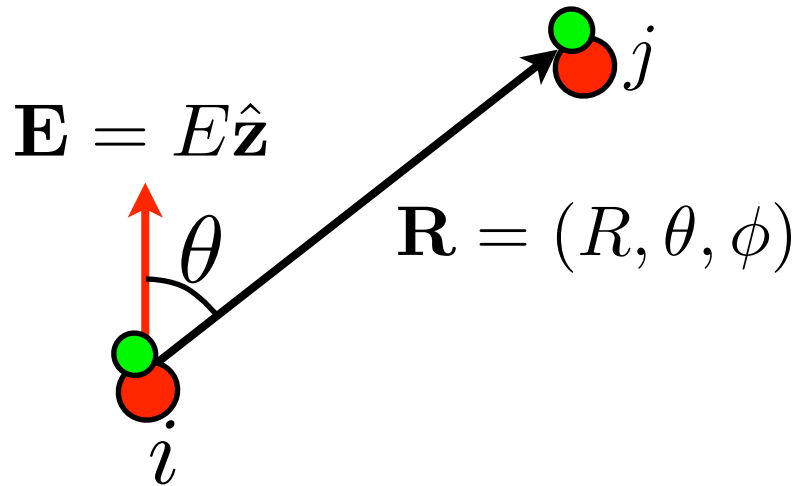
$$R^3 H_{dd} = J_z Y_{2,0} S_i^z S_j^z + (J_{xy} Y_{2,0} + \text{Re}[J'_{xy} Y_{2,-2}]) \frac{1}{2} (S_i^+ S_j^- + S_i^- S_j^+)$$



$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi}$$

[Einstein, de Haas effect]

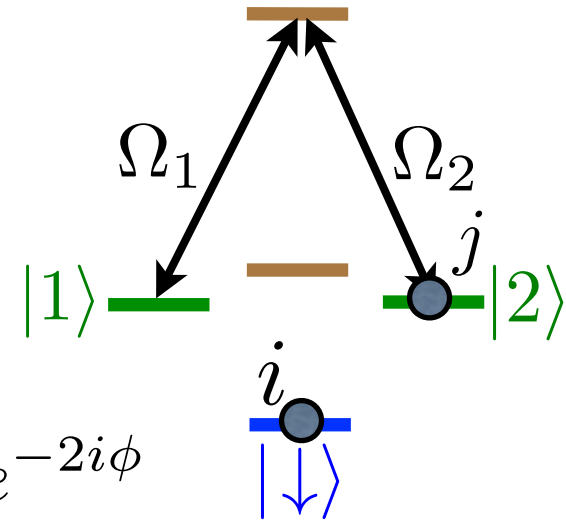
# Spin Hamiltonian from dressed states



$$|\uparrow\rangle = \Omega_2|1\rangle - \Omega_1|2\rangle$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

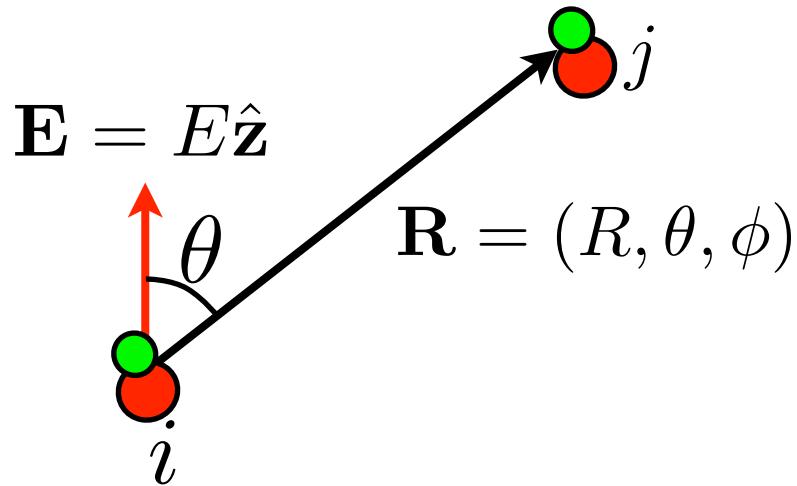
$$R^3 H_{dd} = J_z Y_{2,0} S_i^z S_j^z + (J_{xy} Y_{2,0} + \text{Re}[J'_{xy} Y_{2,-2}]) \frac{1}{2} \underbrace{(S_i^+ S_j^- + S_i^- S_j^+)}_{S_i^x S_j^x + S_i^y S_j^y}$$



$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi}$$

[Einstein, de Haas effect]

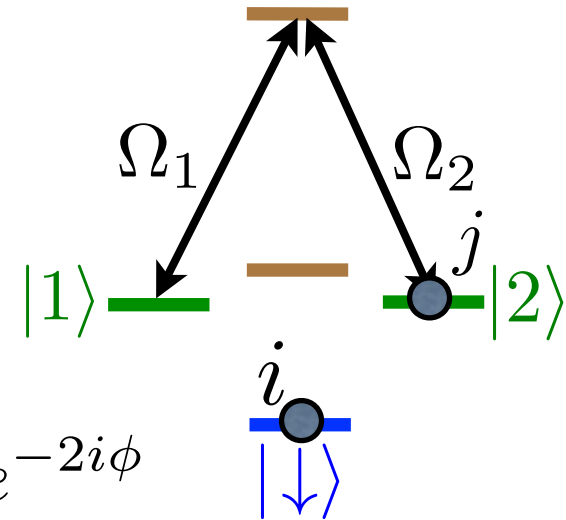
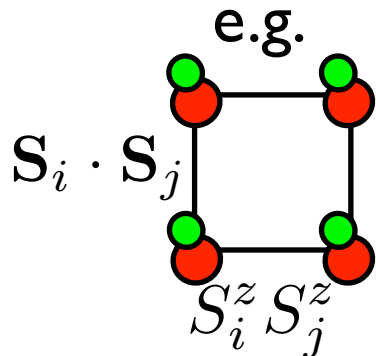
# Spin Hamiltonian from dressed states



$$|\uparrow\rangle = \Omega_2 |1\rangle - \Omega_1 |2\rangle$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

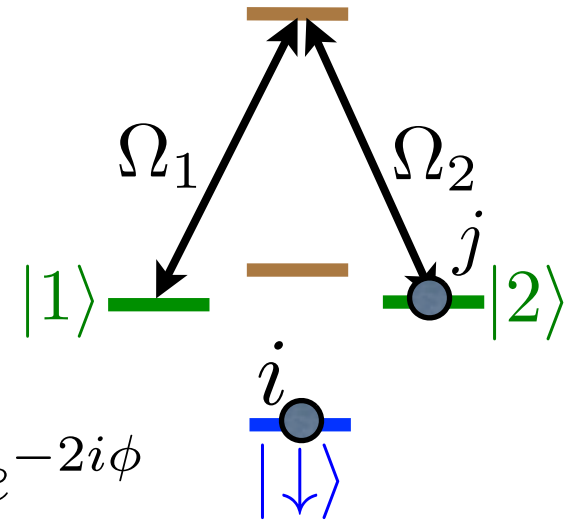
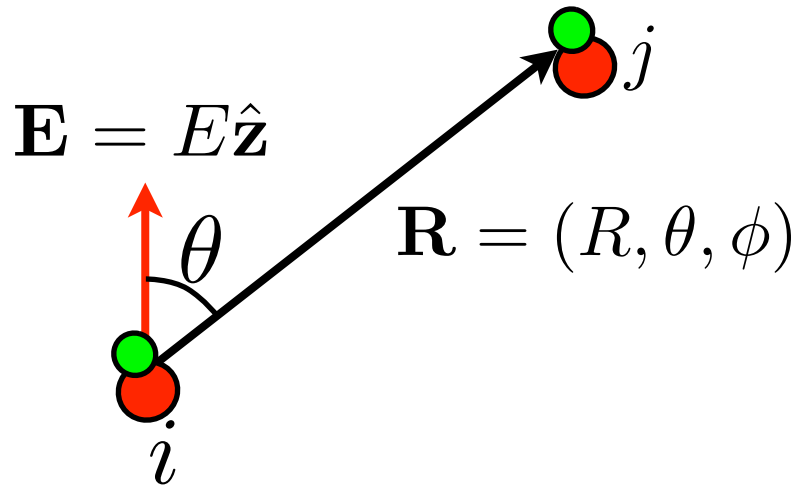
$$R^3 H_{dd} = J_z Y_{2,0} S_i^z S_j^z + (J_{xy} Y_{2,0} + \text{Re}[J'_{xy} Y_{2,-2}]) \frac{1}{2} \underbrace{(S_i^+ S_j^- + S_i^- S_j^+)}_{S_i^x S_j^x + S_i^y S_j^y}$$



$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi}$$

[Einstein, de Haas effect]

# Spin Hamiltonian from dressed states



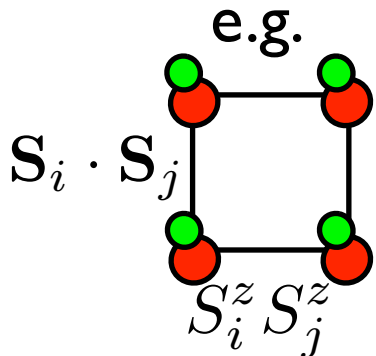
$$Y_{2,-2} \propto \sin^2 \theta e^{-2i\phi}$$

[Einstein, de Haas effect]

$$|\uparrow\rangle = \Omega_2 |1\rangle - \Omega_1 |2\rangle$$

- project on  $|\uparrow\rangle, |\downarrow\rangle$

$$R^3 H_{dd} = J_z Y_{2,0} S_i^z S_j^z + (J_{xy} Y_{2,0} + \text{Re}[J'_{xy} Y_{2,-2}]) \frac{1}{2} \underbrace{(S_i^+ S_j^- + S_i^- S_j^+)}_{S_i^x S_j^x + S_i^y S_j^y}$$



- what can we implement with this?



# Symmetry protected topological (SPT) phases

Shnyder et al 2008, Kitaev 2009, Gu & Wen 2009

# Symmetry protected topological (SPT) phases

- **topological**  $\approx$  no local order parameter, exotic

local order  
parameter:  
e.g.  $\uparrow \downarrow \uparrow \downarrow \uparrow \downarrow$

# Symmetry protected topological (SPT) phases

- **topological**  $\approx$  no local order parameter, exotic

local order  
parameter:  
e.g.  $\uparrow \downarrow \uparrow \downarrow \uparrow \downarrow$

- **symmetry protected** = distinct from other phases only in the presence of symmetry

# Symmetry protected topological (SPT) phases

- **topological**  $\approx$  no local order parameter, exotic

local order  
parameter:  
e.g.  $\uparrow \downarrow \uparrow \downarrow \uparrow \downarrow$

- **symmetry protected** = distinct from other phases only in the presence of symmetry

phase  
diagram:



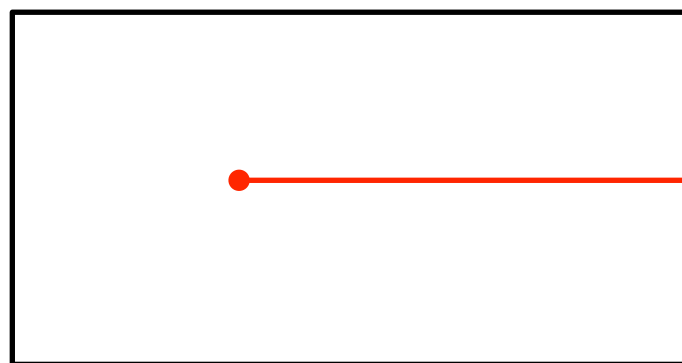
# Symmetry protected topological (SPT) phases

- **topological**  $\approx$  no local order parameter, exotic

local order  
parameter:  
e.g.  $\uparrow \downarrow \uparrow \downarrow \uparrow \downarrow$

- **symmetry protected** = distinct from other phases only in the presence of symmetry

phase  
diagram:



← line of phase  
transitions

Shnyder et al 2008, Kitaev 2009, Gu & Wen 2009

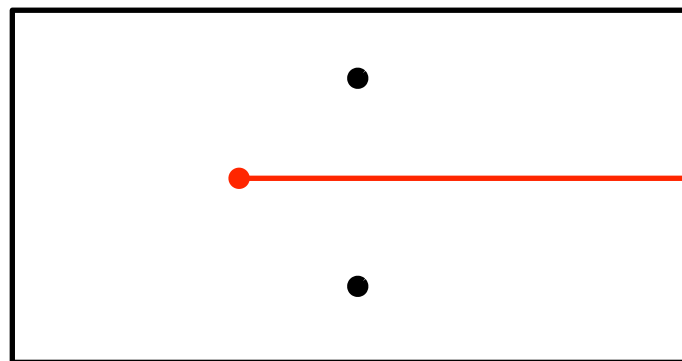
# Symmetry protected topological (SPT) phases

- **topological**  $\approx$  no local order parameter, exotic

local order  
parameter:  
e.g.  $\uparrow \downarrow \uparrow \downarrow \uparrow \downarrow$

- **symmetry protected** = distinct from other phases only in the presence of symmetry

phase  
diagram:



← line of phase  
transitions

Shnyder et al 2008, Kitaev 2009, Gu & Wen 2009

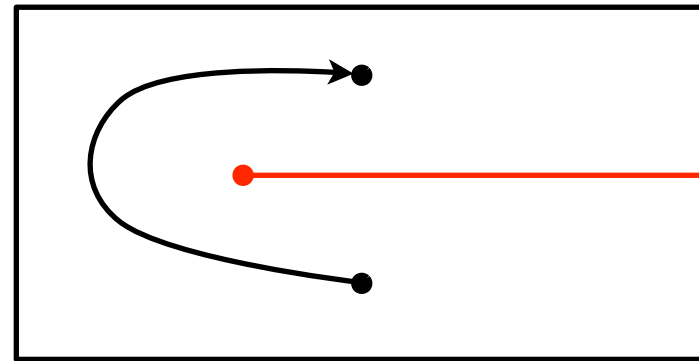
# Symmetry protected topological (SPT) phases

- **topological**  $\approx$  no local order parameter, exotic

local order  
parameter:  
e.g.  $\uparrow \downarrow \uparrow \downarrow \uparrow \downarrow$

- **symmetry protected** = distinct from other phases only in the presence of symmetry

phase  
diagram:



← line of phase  
transitions

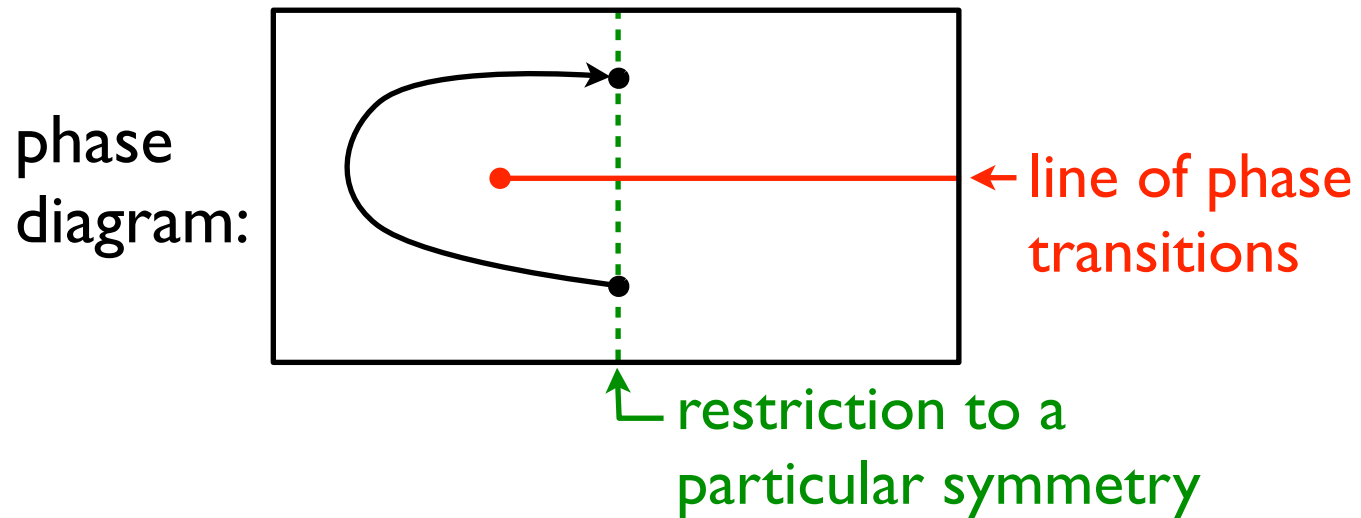
Shnyder et al 2008, Kitaev 2009, Gu & Wen 2009

# Symmetry protected topological (SPT) phases

- **topological**  $\approx$  no local order parameter, exotic

local order  
parameter:  
e.g.  $\uparrow \downarrow \uparrow \downarrow \uparrow \downarrow$

- **symmetry protected** = distinct from other phases only in the presence of symmetry



Shnyder et al 2008, Kitaev 2009, Gu & Wen 2009



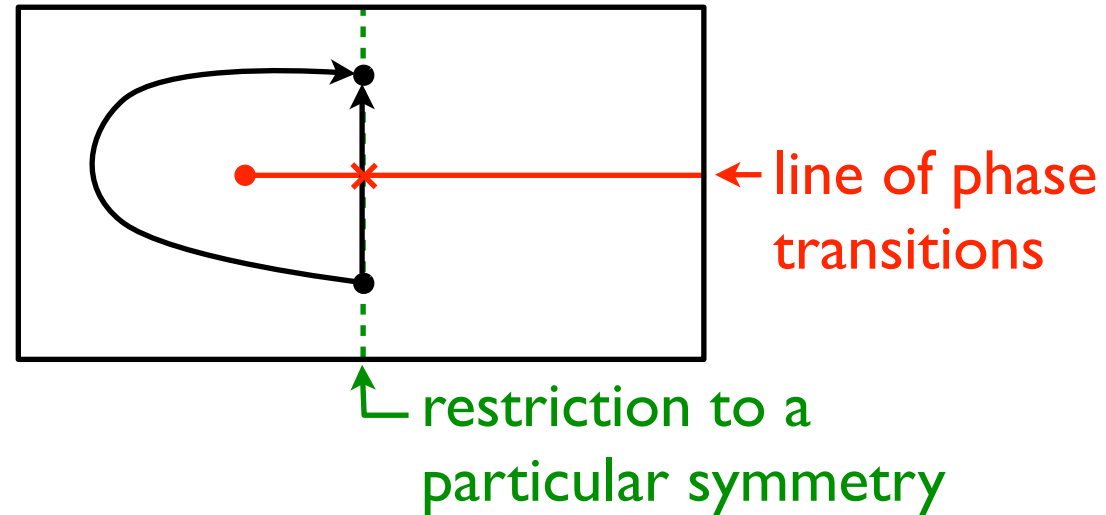
# Symmetry protected topological (SPT) phases

- **topological**  $\approx$  no local order parameter, exotic

local order  
parameter:  
e.g.  $\uparrow \downarrow \uparrow \downarrow \uparrow \downarrow$

- **symmetry protected** = distinct from other phases only in the presence of symmetry

phase  
diagram:



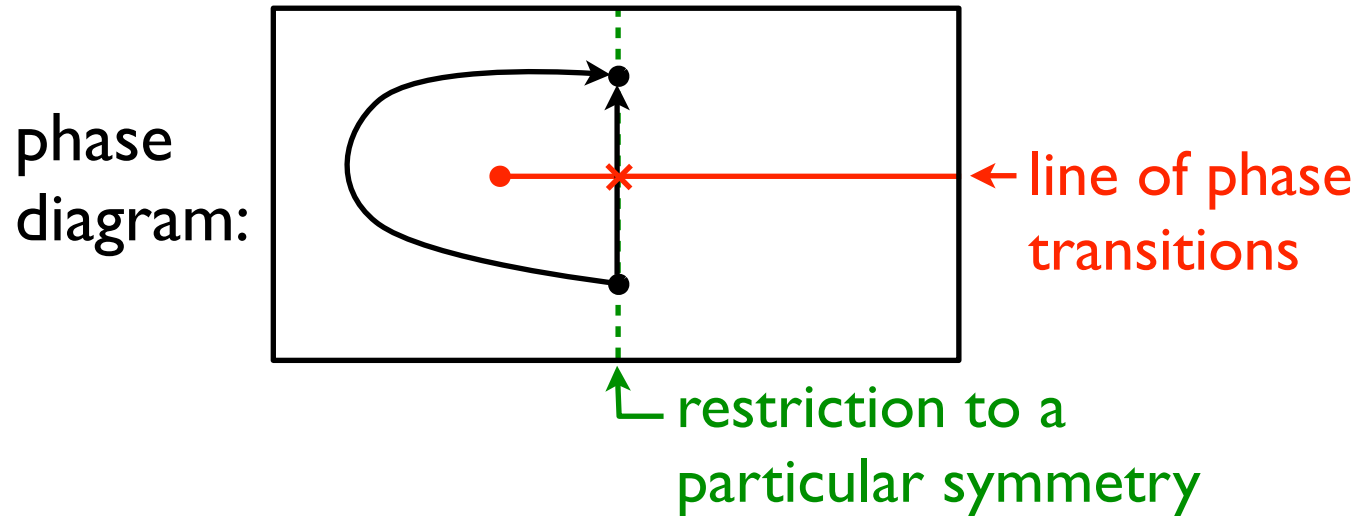
Shnyder et al 2008, Kitaev 2009, Gu & Wen 2009

# Symmetry protected topological (SPT) phases

- **topological**  $\approx$  no local order parameter, exotic

local order  
parameter:  
e.g.  $\uparrow \downarrow \uparrow \downarrow \uparrow \downarrow$

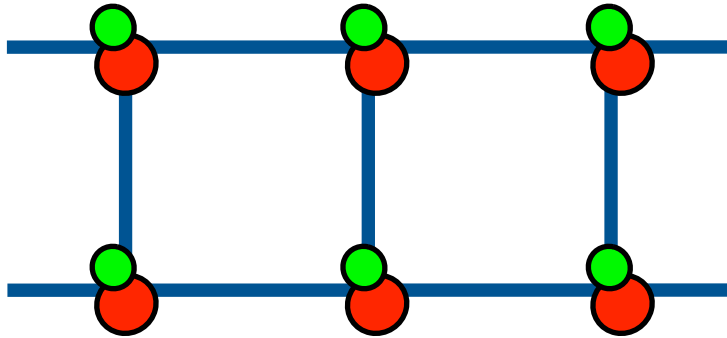
- **symmetry protected** = distinct from other phases only in the presence of symmetry



- **motivation:** quest to classify and understand all phases of matter

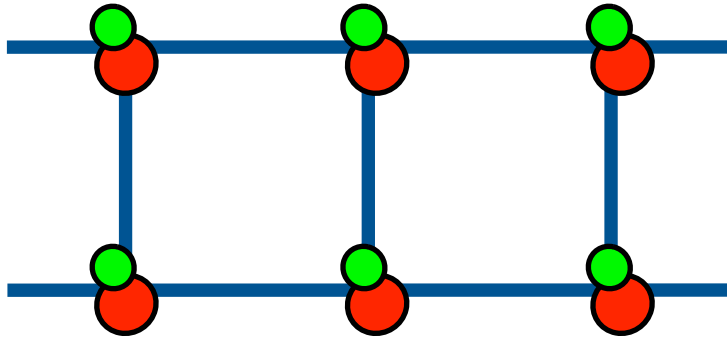
Shnyder et al 2008, Kitaev 2009, Gu & Wen 2009

# SPT phases in spin-1/2 ladders



Liu, Wen, et al, PRB 2012; Chen, Gu, Wen, PRB 2011; Pollman et al, PRB 2010

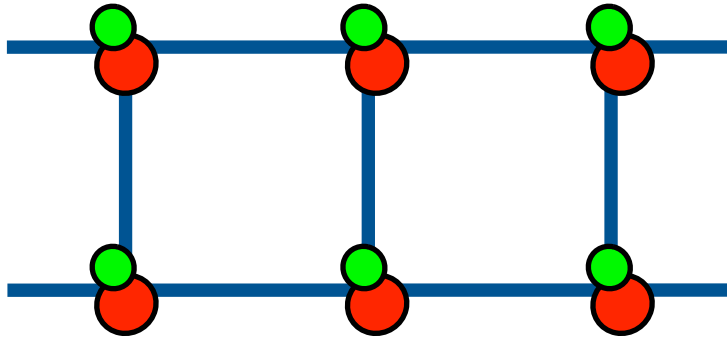
# SPT phases in spin-1/2 ladders



- symmetry:

$$D_2 \times \sigma$$

# SPT phases in spin-1/2 ladders

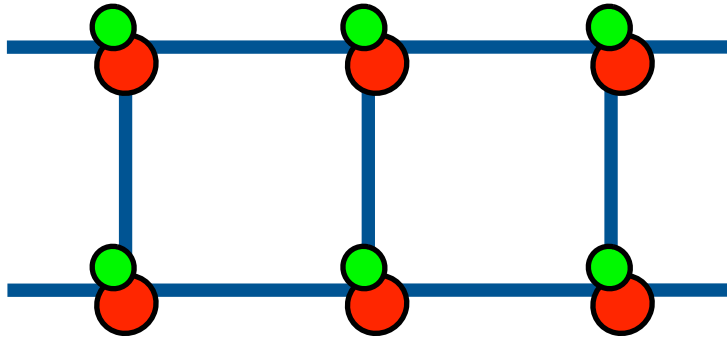


- symmetry:

$$D_2 \times \sigma$$

⏟ exchange of the two legs

# SPT phases in spin-1/2 ladders



- symmetry:

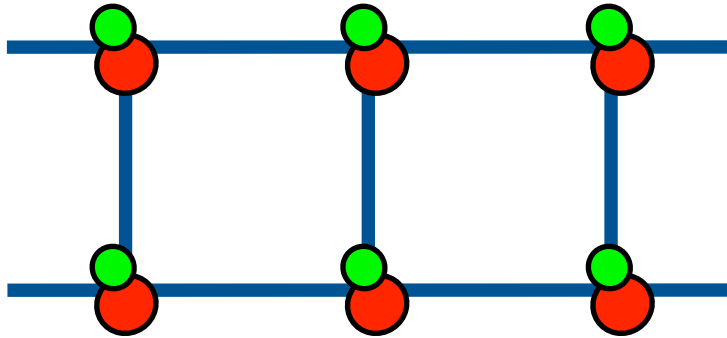
$$D_2 \times \sigma$$

exchange of the two legs

$$\{E, R_x, R_y, R_z\}$$

identity

# SPT phases in spin-1/2 ladders



- symmetry:

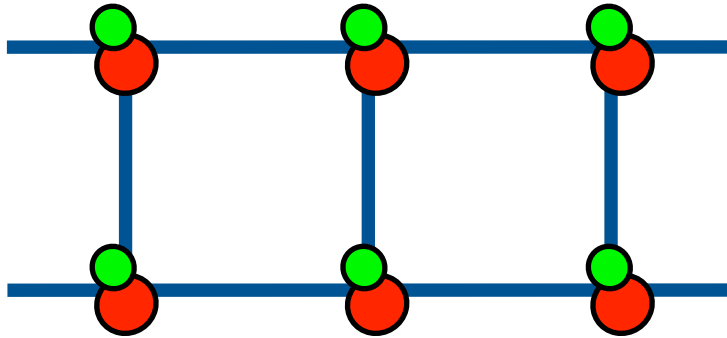
$$D_2 \times \sigma$$

exchange of the two legs

$$\{E, R_x, R_y, R_z\}$$

identity  $\pi$ -pulse around x-axis on all spins

# SPT phases in spin-1/2 ladders



- symmetry:

$$D_2 \times \sigma$$

exchange of the two legs

$$\{E, R_x, R_y, R_z\}$$

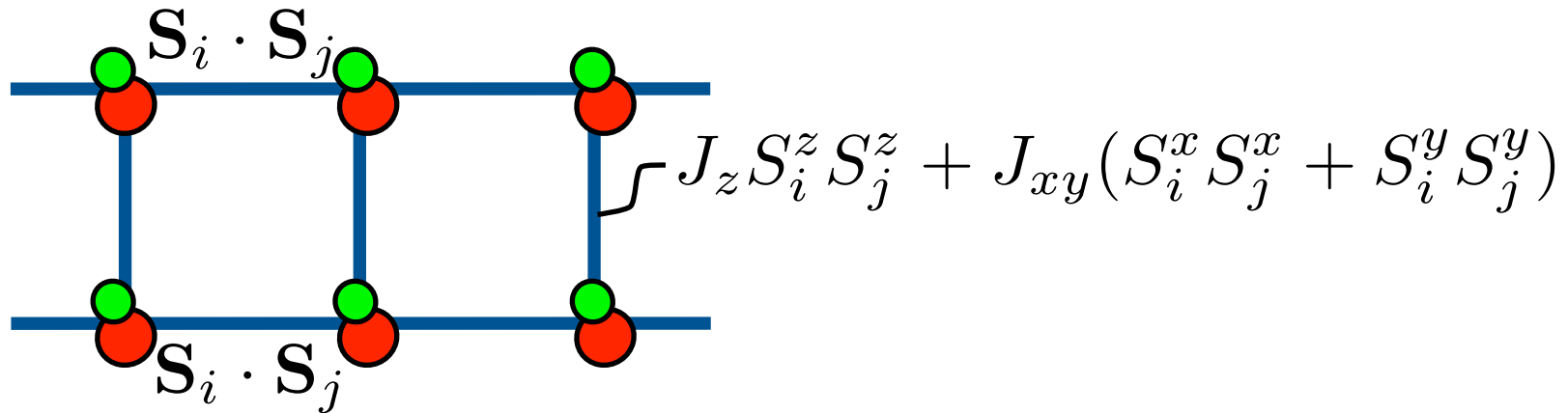
identity  $\pi$ -pulse around x-axis on all spins



7 non-trivial SPT phases



# SPT phases in spin-1/2 ladders



- symmetry:

$$D_2 \times \sigma$$

exchange of the two legs

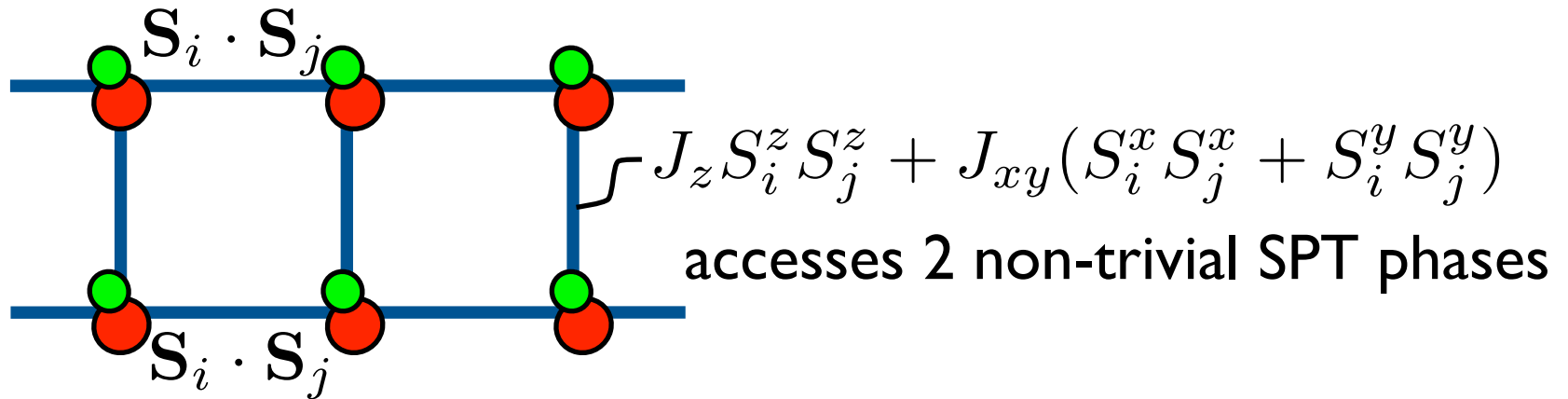
$$\{E, R_x, R_y, R_z\}$$

identity  $\pi$ -pulse around x-axis on all spins



7 non-trivial SPT phases

# SPT phases in spin-1/2 ladders



- symmetry:

$$D_2 \times \sigma$$

exchange of the two legs

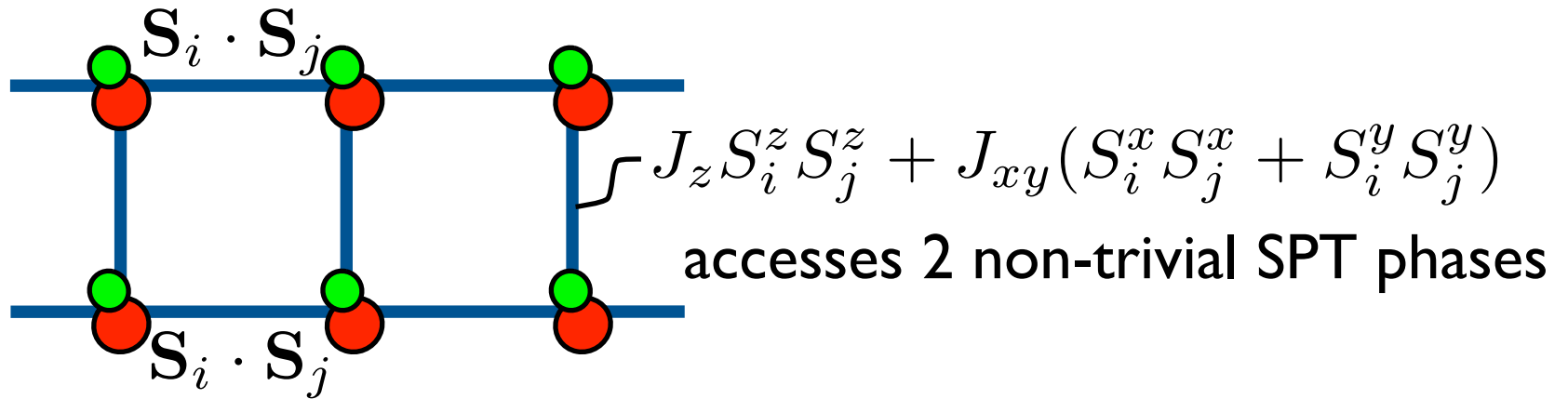
$$\{E, R_x, R_y, R_z\}$$

identity  $\pi$ -pulse around x-axis on all spins

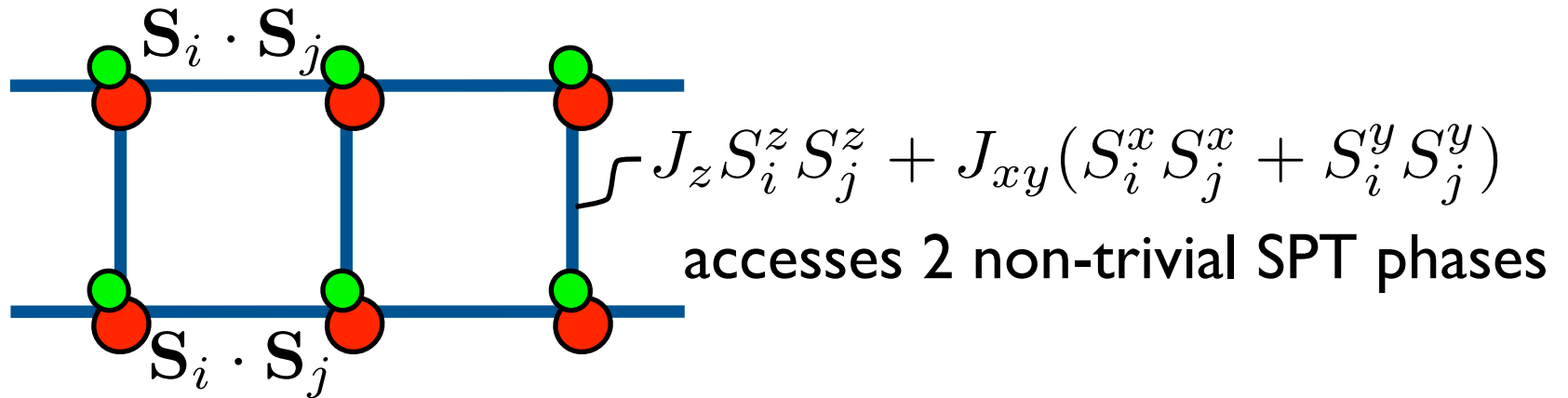


7 non-trivial SPT phases

# SPT phases in spin-1/2 ladders



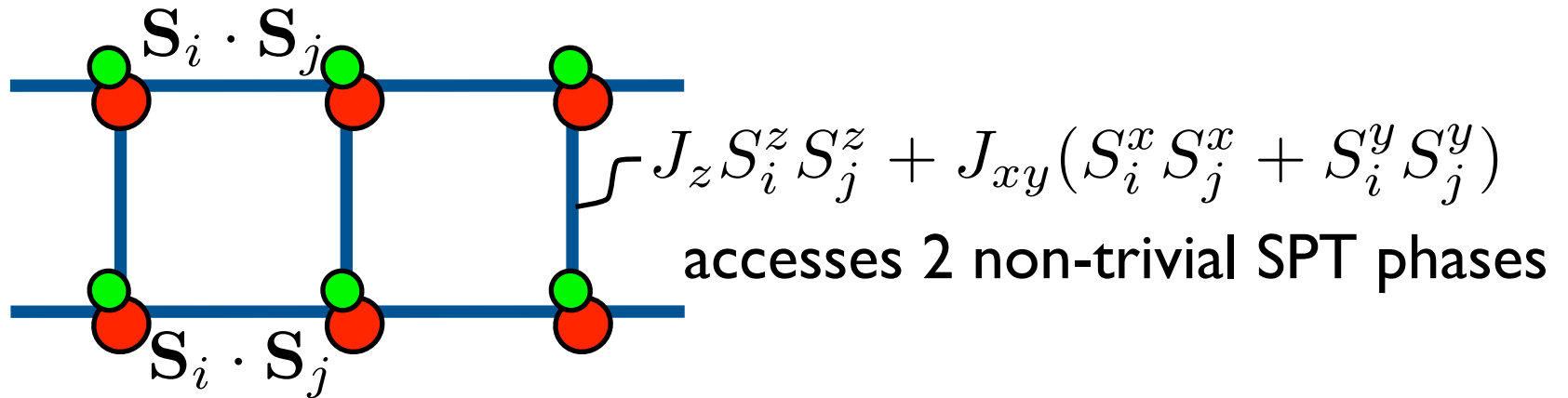
# SPT phases in spin-1/2 ladders



One SPT phase:

The other SPT phase:

# SPT phases in spin-1/2 ladders

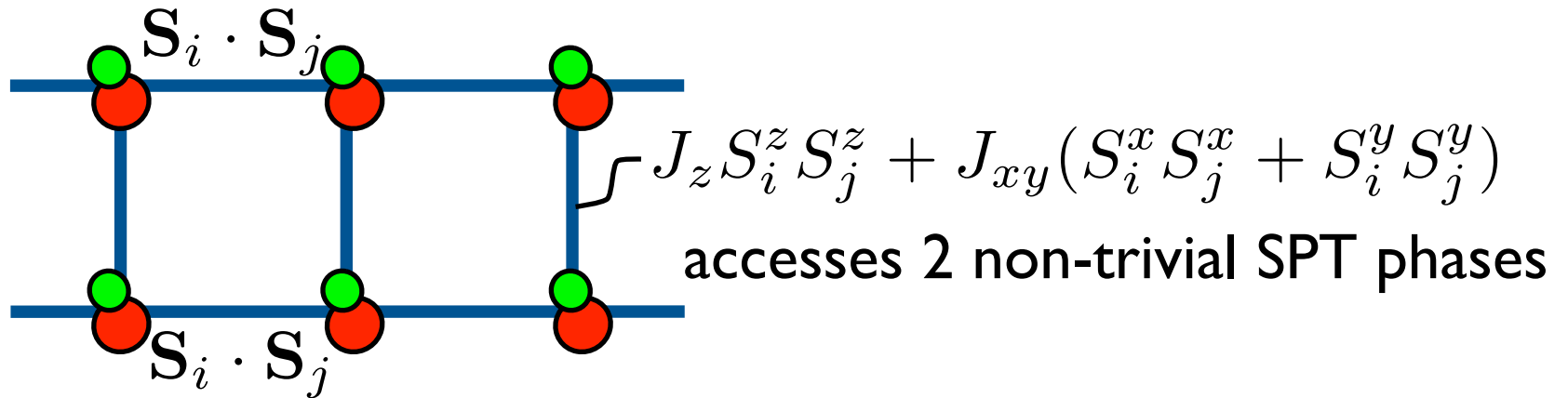


One SPT phase:

- $J_{xy} = J_z \ll -1$

The other SPT phase:

# SPT phases in spin-1/2 ladders

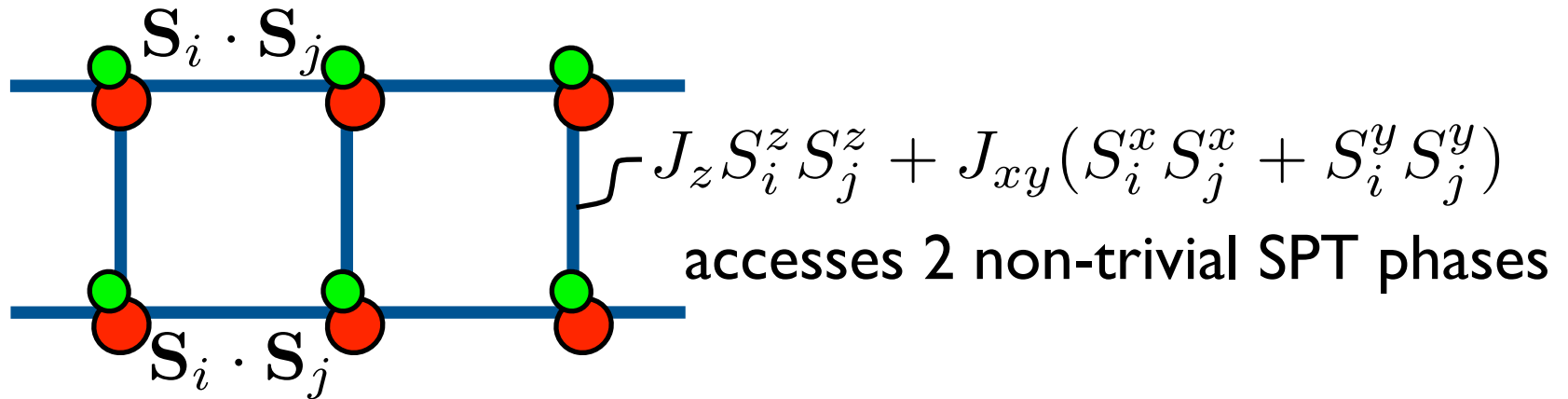


One SPT phase:

- $J_{xy} = J_z \ll -1$
- one rung:  $-|J| \mathbf{S}_i \cdot \mathbf{S}_j$

The other SPT phase:

# SPT phases in spin-1/2 ladders



One SPT phase:

- $J_{xy} = J_z \ll -1$
- one rung:  $-|J| \mathbf{S}_i \cdot \mathbf{S}_j$

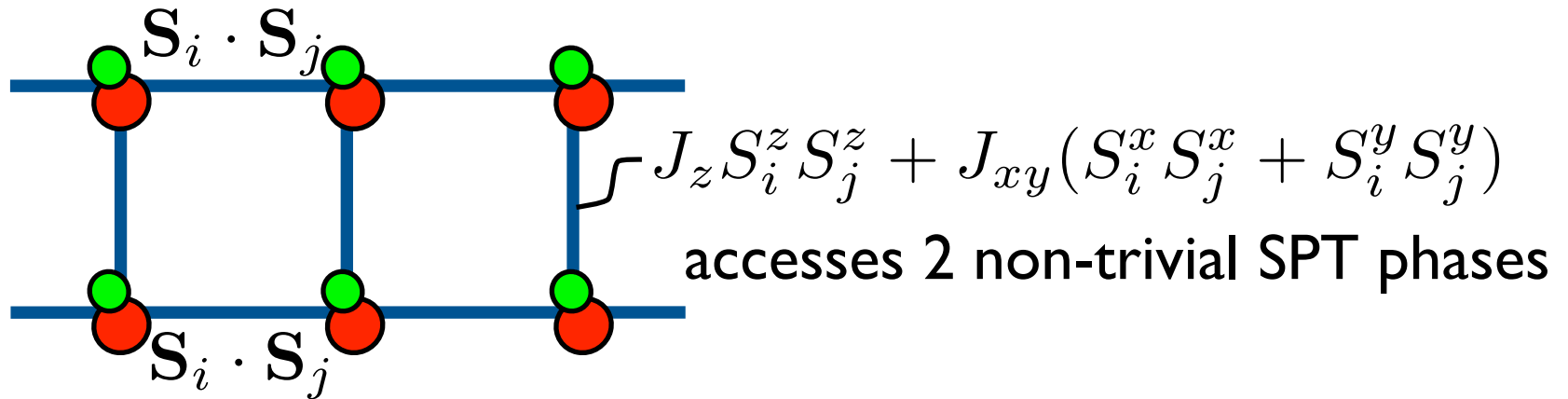
$$|\uparrow\uparrow\rangle$$

$$|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle$$

$$|\downarrow\downarrow\rangle$$

The other SPT phase:

# SPT phases in spin-1/2 ladders



## One SPT phase:

- $J_{xy} = J_z \ll -1$
- one rung:  $-|J| \mathbf{S}_i \cdot \mathbf{S}_j$

$$|\uparrow\uparrow\rangle$$

$$|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle$$

$$|\downarrow\downarrow\rangle$$

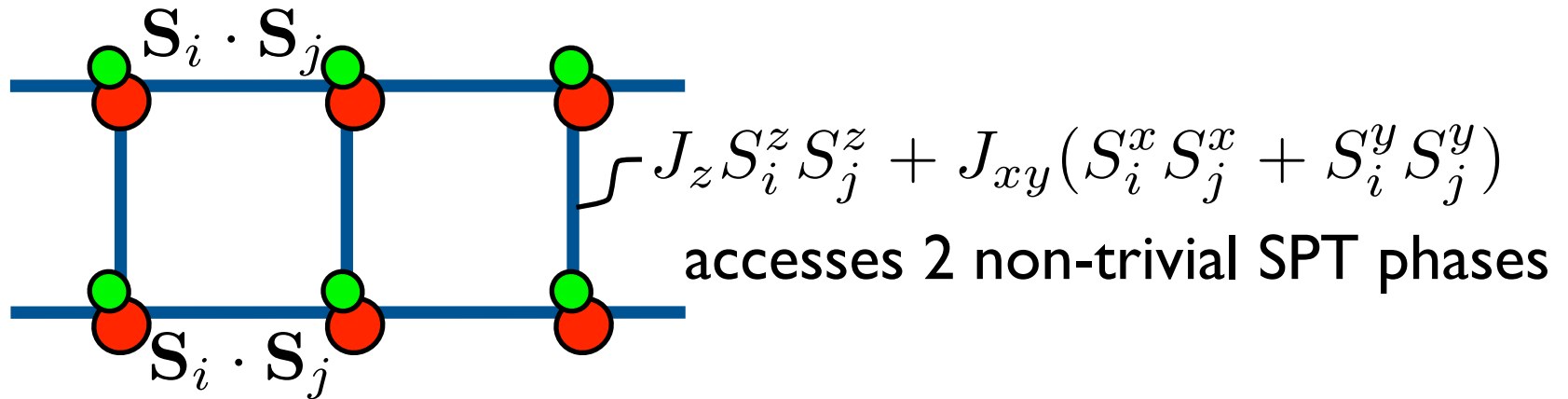
- spin-1 antiferromagnet  
 $\Rightarrow$  Haldane phase

## The other SPT phase:

Liu, Wen, et al, PRB 2012; Chen, Gu, Wen, PRB 2011; Pollman et al, PRB 2010



# SPT phases in spin-1/2 ladders



## One SPT phase:

- $J_{xy} = J_z \ll -1$
- one rung:  $-|J| \mathbf{S}_i \cdot \mathbf{S}_j$

$$|\uparrow\uparrow\rangle$$

$$|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle$$

$$|\downarrow\downarrow\rangle$$

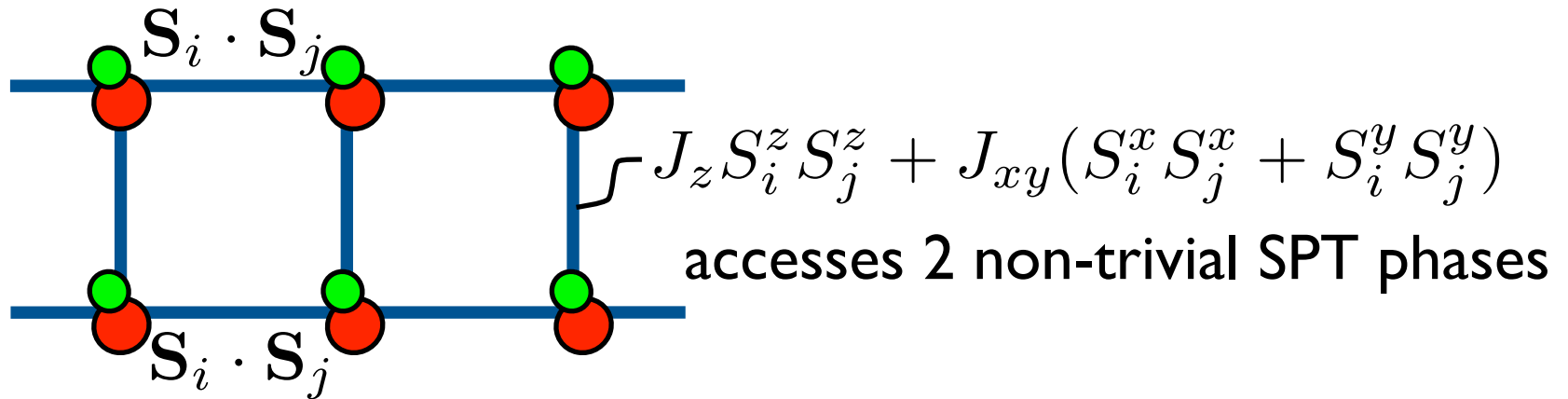
- spin-1 antiferromagnet  
 $\Rightarrow$  Haldane phase

## The other SPT phase:

- $J_{xy} \rightarrow -J_{xy}$

Liu, Wen, et al, PRB 2012; Chen, Gu, Wen, PRB 2011; Pollman et al, PRB 2010

# SPT phases in spin-1/2 ladders



## One SPT phase:

- $J_{xy} = J_z \ll -1$
- one rung:  $-|J| \mathbf{S}_i \cdot \mathbf{S}_j$

$$|\uparrow\uparrow\rangle$$

$$|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle$$

$$|\downarrow\downarrow\rangle$$

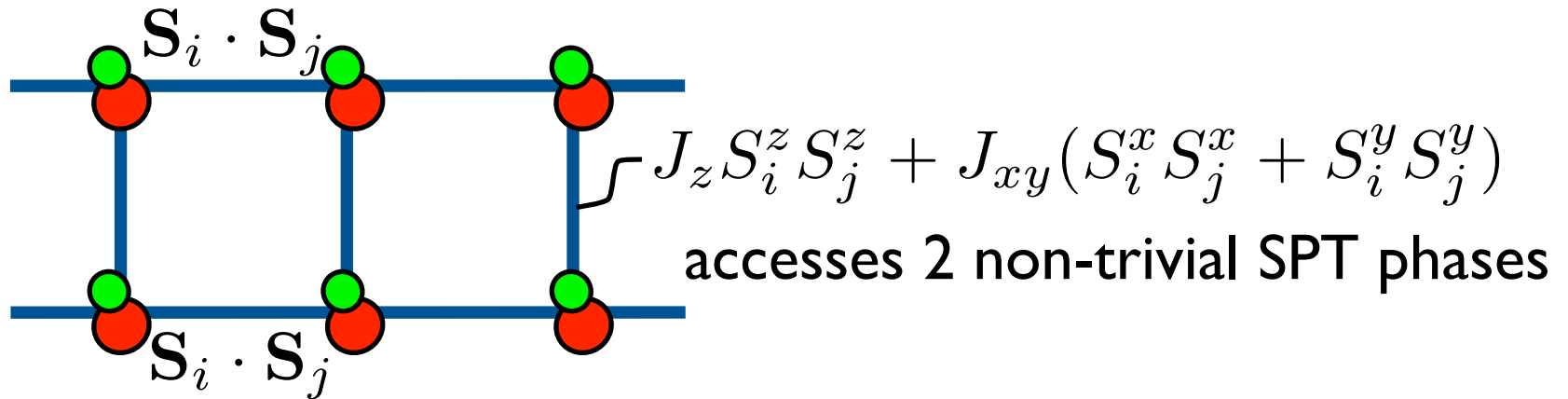
- spin-1 antiferromagnet  
 $\Rightarrow$  Haldane phase

## The other SPT phase:

- $J_{xy} \rightarrow -J_{xy}$   
 $\Updownarrow$
- $|\downarrow\rangle \rightarrow -|\downarrow\rangle$  on one leg

Liu, Wen, et al, PRB 2012; Chen, Gu, Wen, PRB 2011; Pollman et al, PRB 2010

# SPT phases in spin-1/2 ladders



## One SPT phase:

- $J_{xy} = J_z \ll -1$
- one rung:  $-|J| \mathbf{S}_i \cdot \mathbf{S}_j$

$$|\uparrow\uparrow\rangle$$

$$|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle$$

$$|\downarrow\downarrow\rangle$$

- spin-1 antiferromagnet  
 $\Rightarrow$  Haldane phase

## The other SPT phase:

- $J_{xy} \rightarrow -J_{xy}$



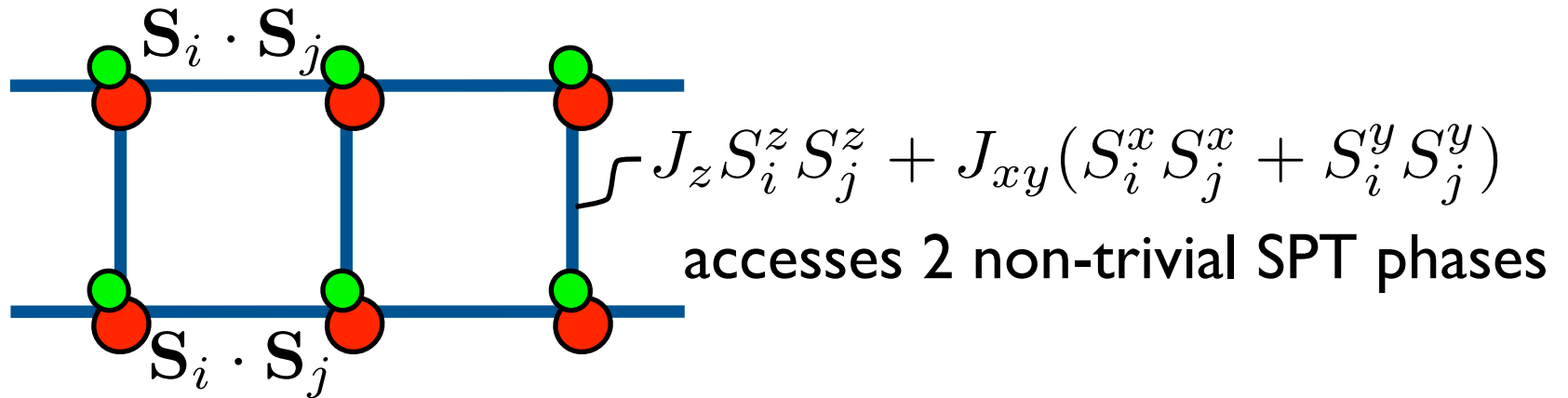
- $|\downarrow\rangle \rightarrow -|\downarrow\rangle$  on one leg

$$|\uparrow\uparrow\rangle$$

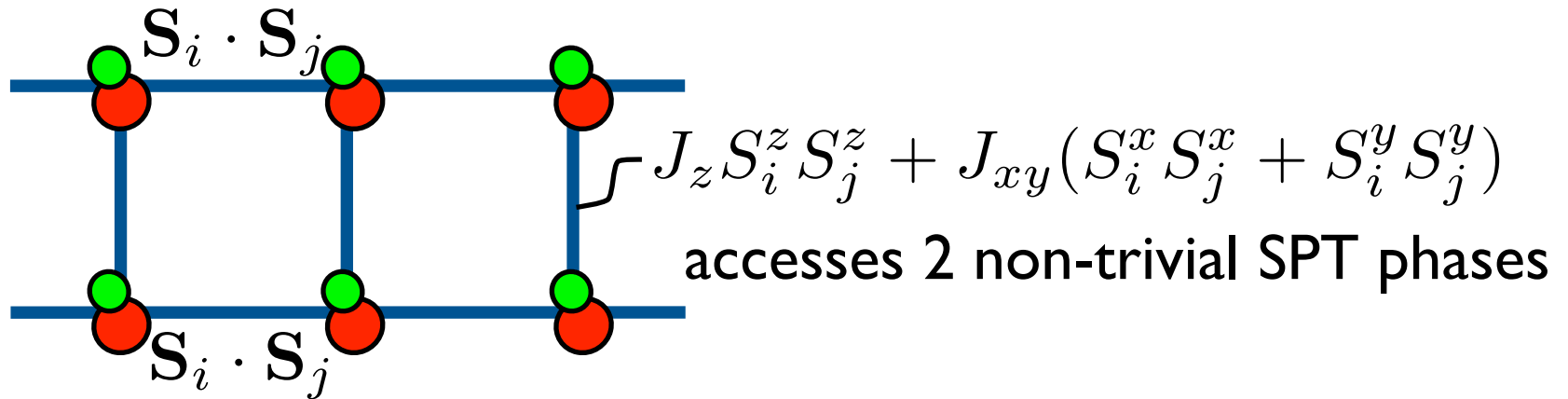
$$|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle$$

$$-|\downarrow\downarrow\rangle$$

# SPT phases in spin-1/2 ladders

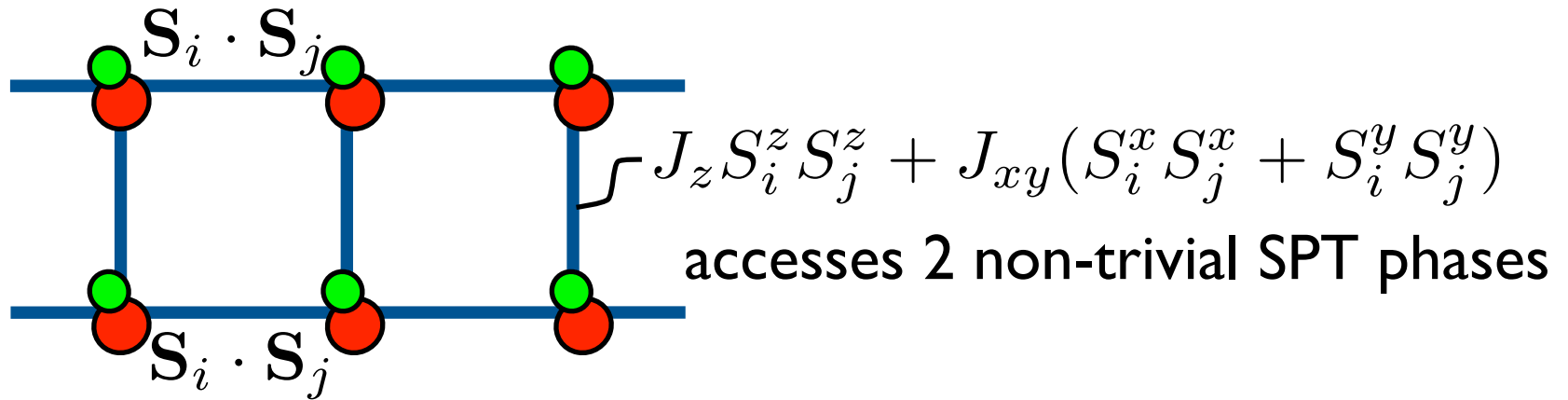


# SPT phases in spin-1/2 ladders



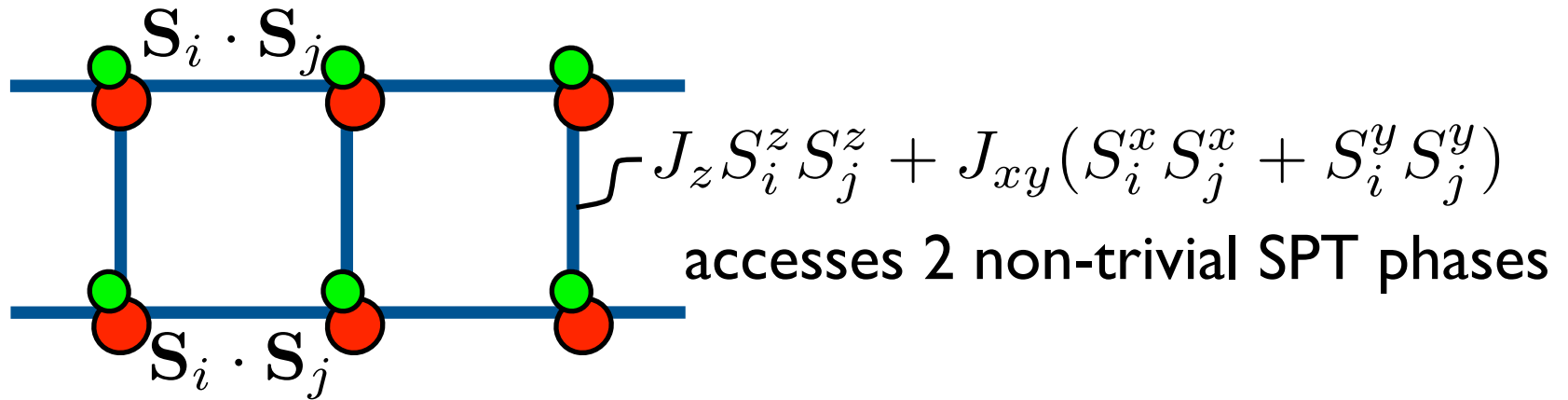
- symmetry-protected features:

# SPT phases in spin-1/2 ladders

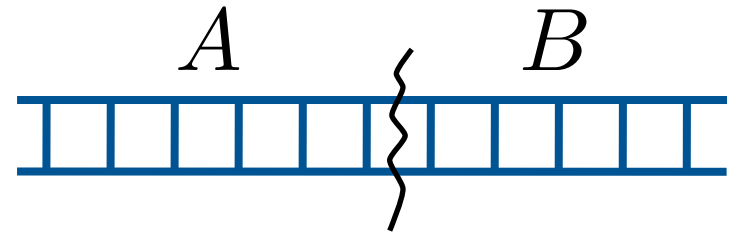


- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:

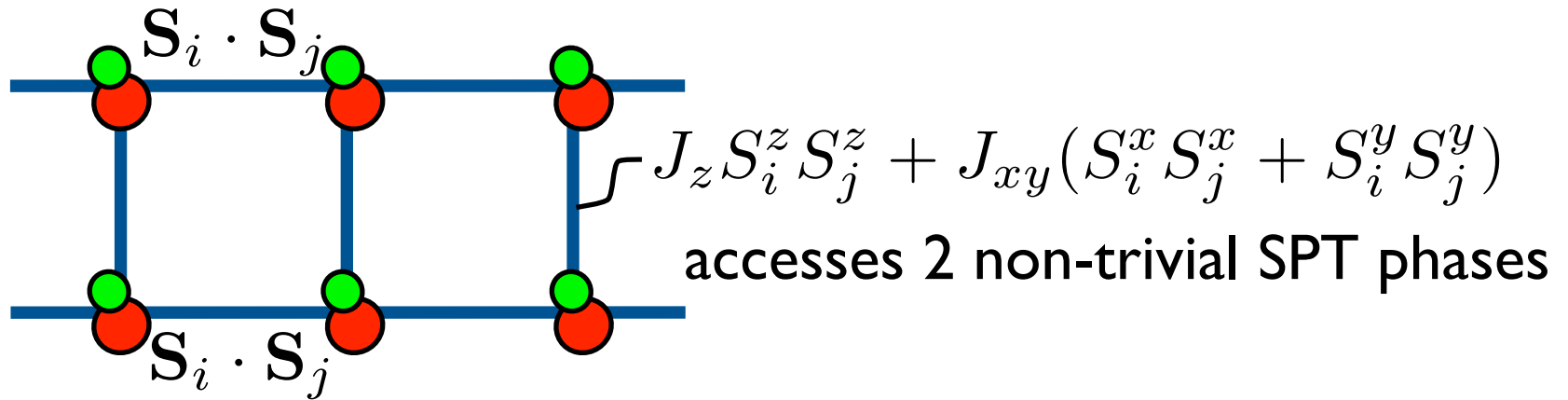
# SPT phases in spin-1/2 ladders



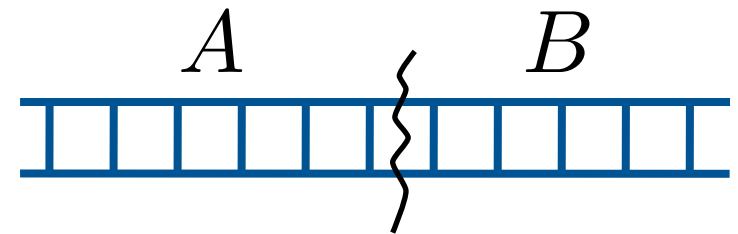
- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:



# SPT phases in spin-1/2 ladders



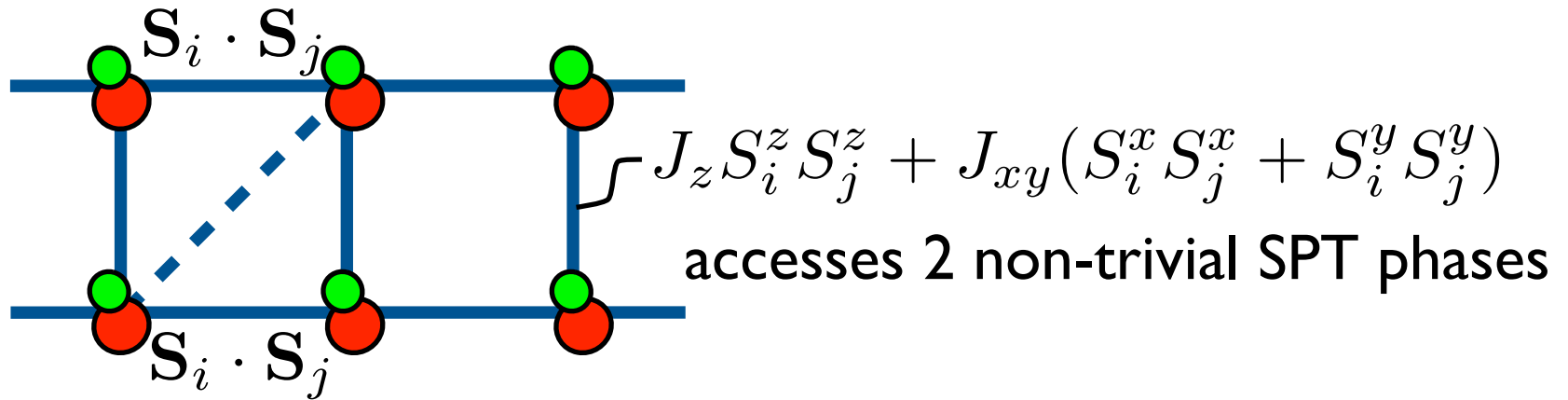
- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:



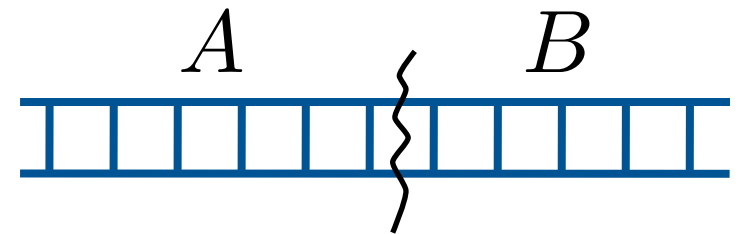
eigenvalues of  $\text{tr}_B [|\Psi\rangle\langle\Psi|]$  are doubly degenerate



# SPT phases in spin-1/2 ladders

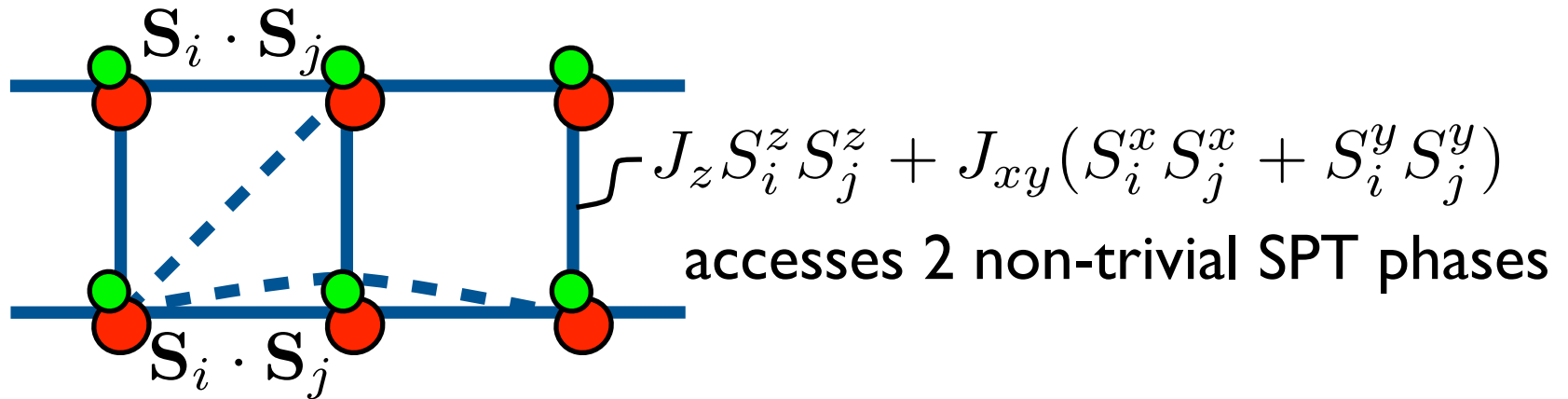


- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:

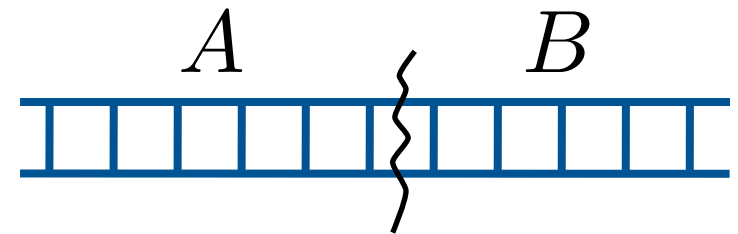


eigenvalues of  $\text{tr}_B [|\Psi\rangle\langle\Psi|]$   
 are doubly degenerate

# SPT phases in spin-1/2 ladders

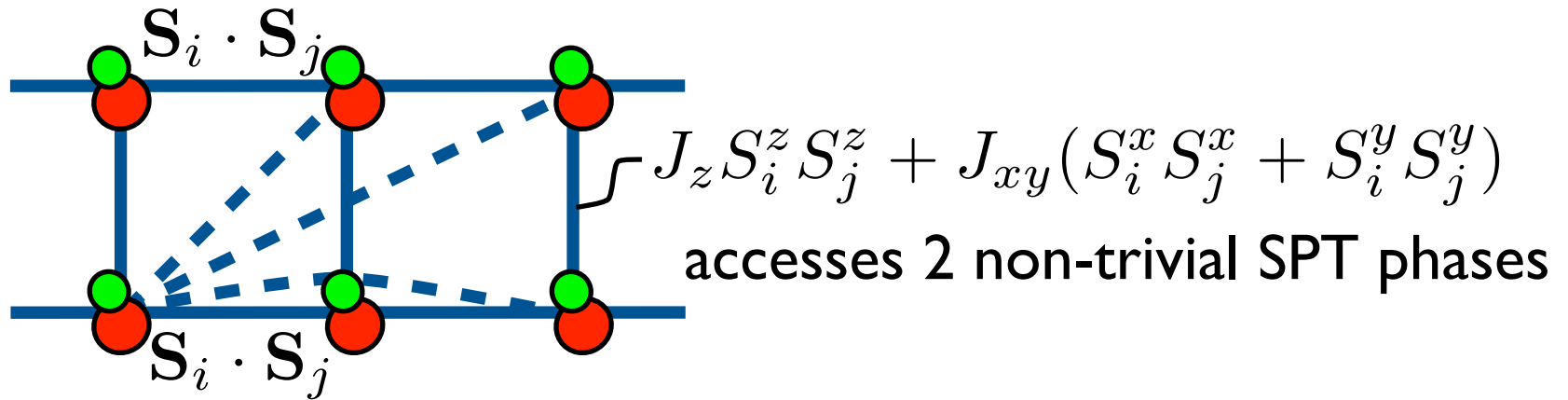


- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:

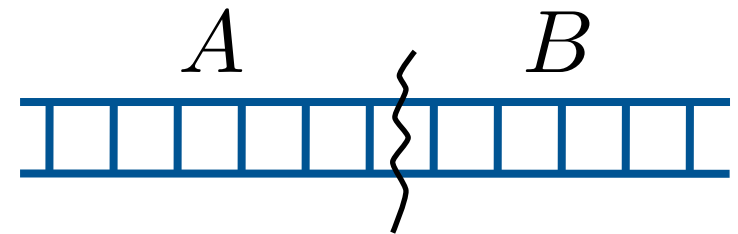


eigenvalues of  $\text{tr}_B [|\Psi\rangle\langle\Psi|]$   
 are doubly degenerate

# SPT phases in spin-1/2 ladders

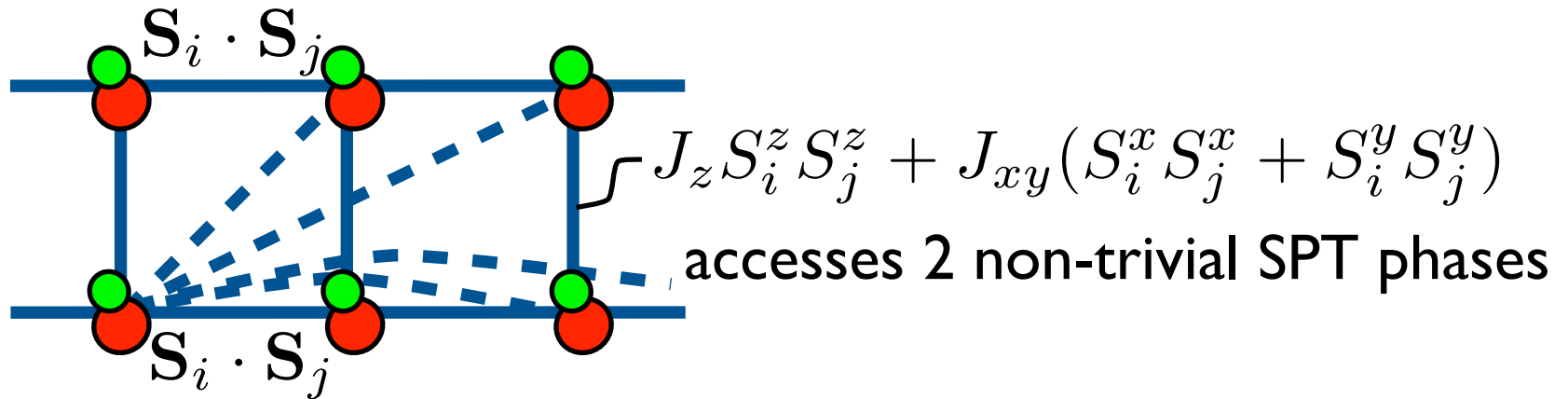


- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:

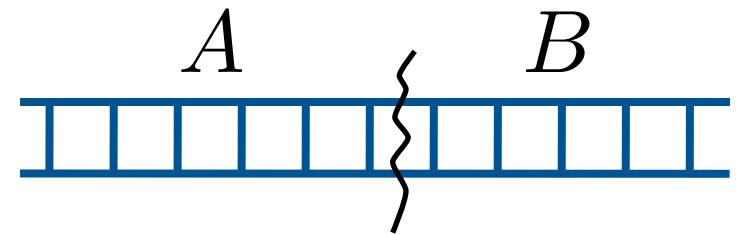


eigenvalues of  $\text{tr}_B [|\Psi\rangle\langle\Psi|]$   
are doubly degenerate

# SPT phases in spin-1/2 ladders

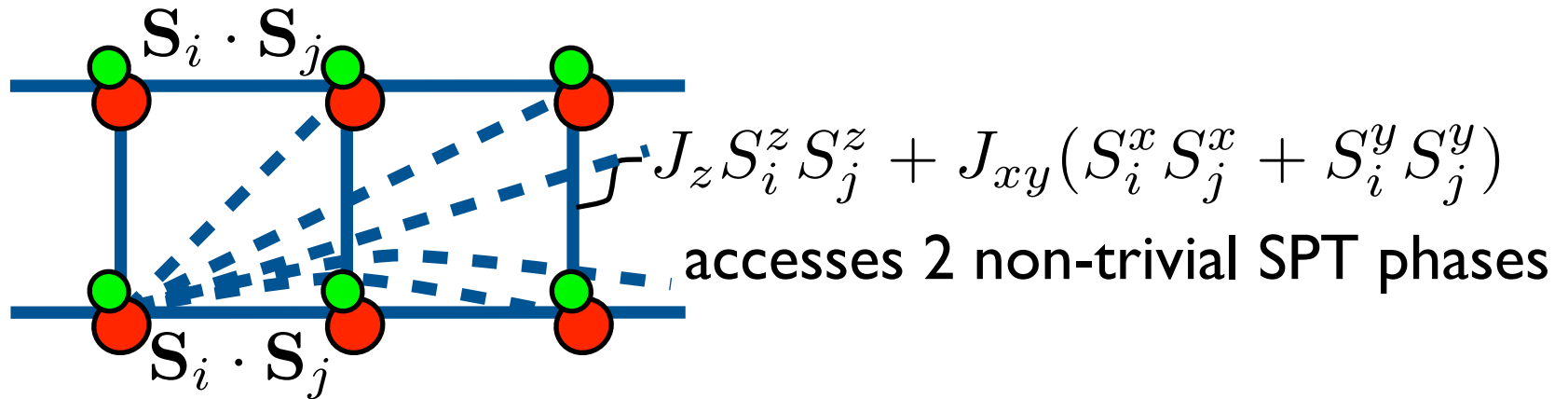


- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:

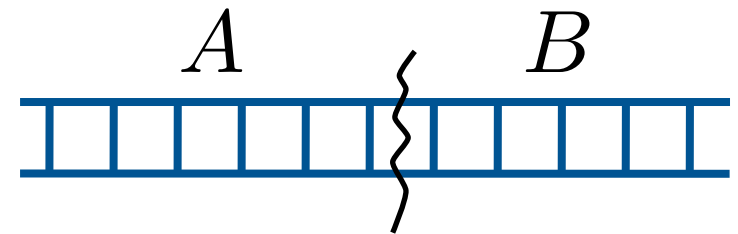


eigenvalues of  $\text{tr}_B [|\Psi\rangle\langle\Psi|]$   
are doubly degenerate

# SPT phases in spin-1/2 ladders

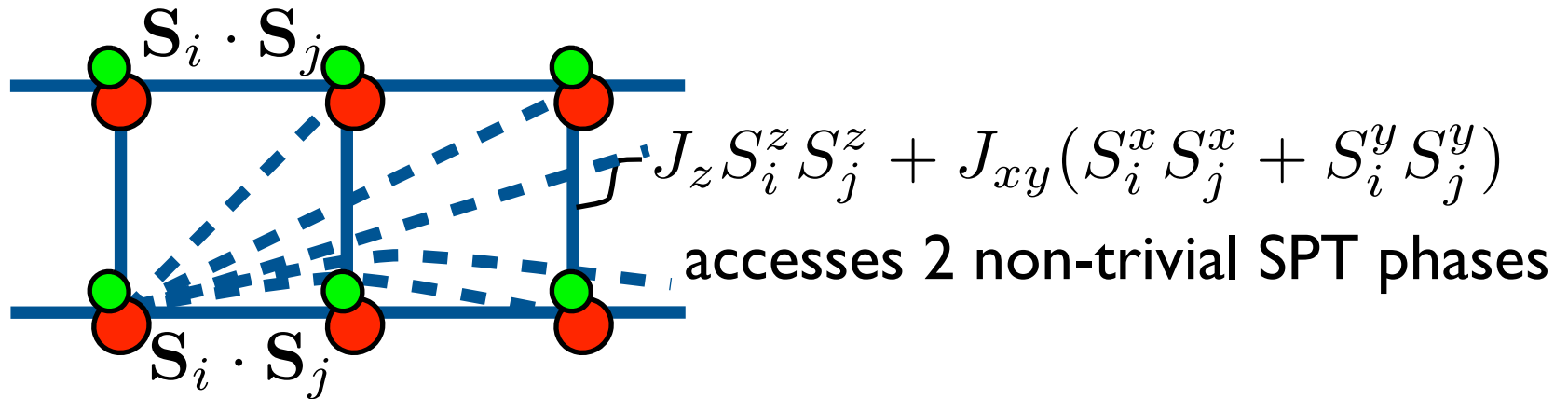


- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:



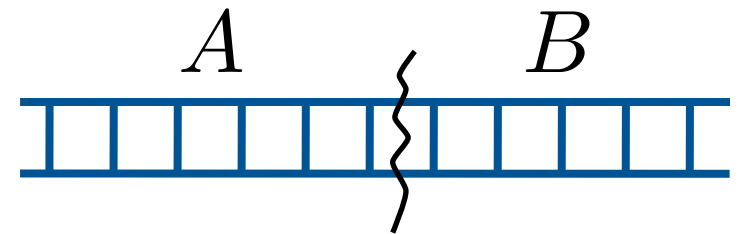
eigenvalues of  $\text{tr}_B [|\Psi\rangle\langle\Psi|]$   
 are doubly degenerate

# SPT phases in spin-1/2 ladders



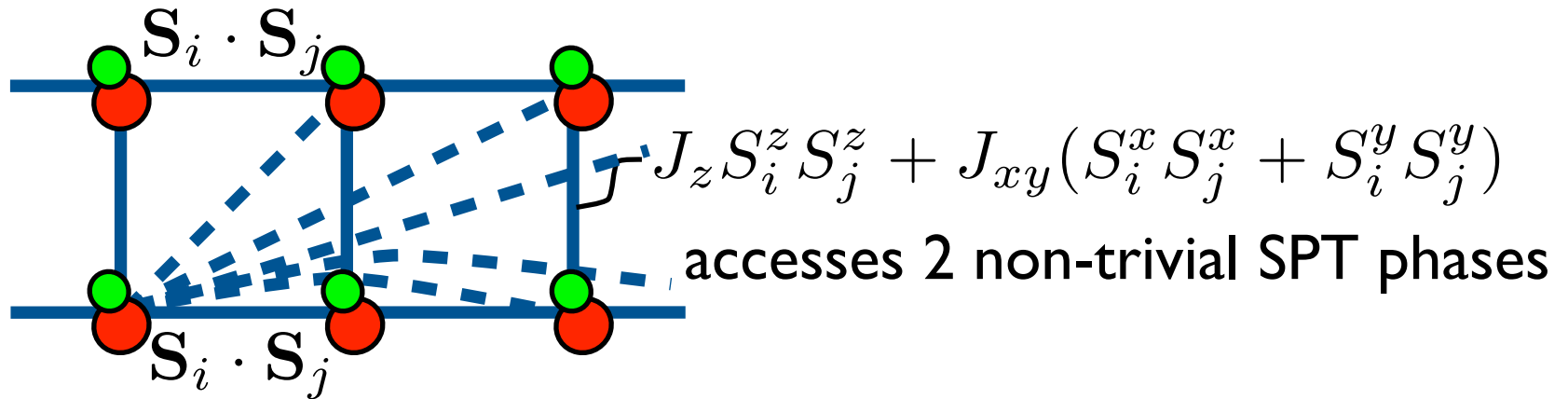
• do SPT phases survive?

- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:



eigenvalues of  $\text{tr}_B [|\Psi\rangle\langle\Psi|]$   
are doubly degenerate

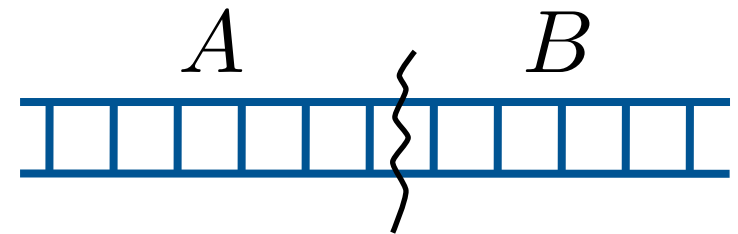
# SPT phases in spin-1/2 ladders



- do SPT phases survive?

**YES!**  
 [PRB 87, 081106(R) (2013)]

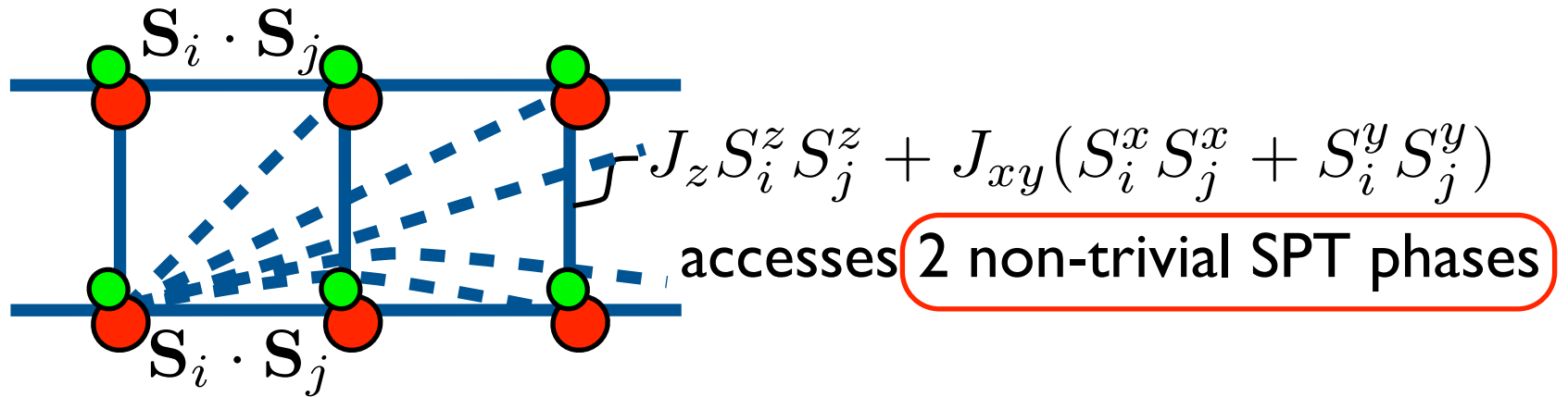
- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:



eigenvalues of  $\text{tr}_B [|\Psi\rangle\langle\Psi|]$   
 are doubly degenerate

Liu, Wen, et al, PRB 2012; Chen, Gu, Wen, PRB 2011; Pollman et al, PRB 2010

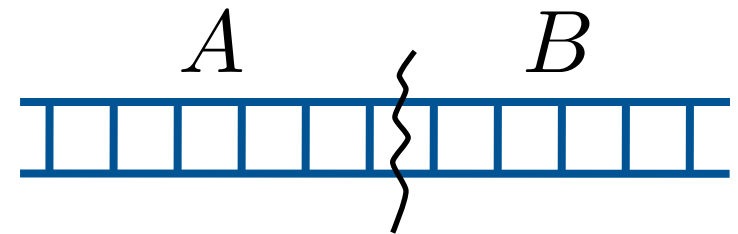
# SPT phases in spin-1/2 ladders



- do SPT phases survive?

**YES!**  
 [PRB 87, 081106(R) (2013)]

- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:

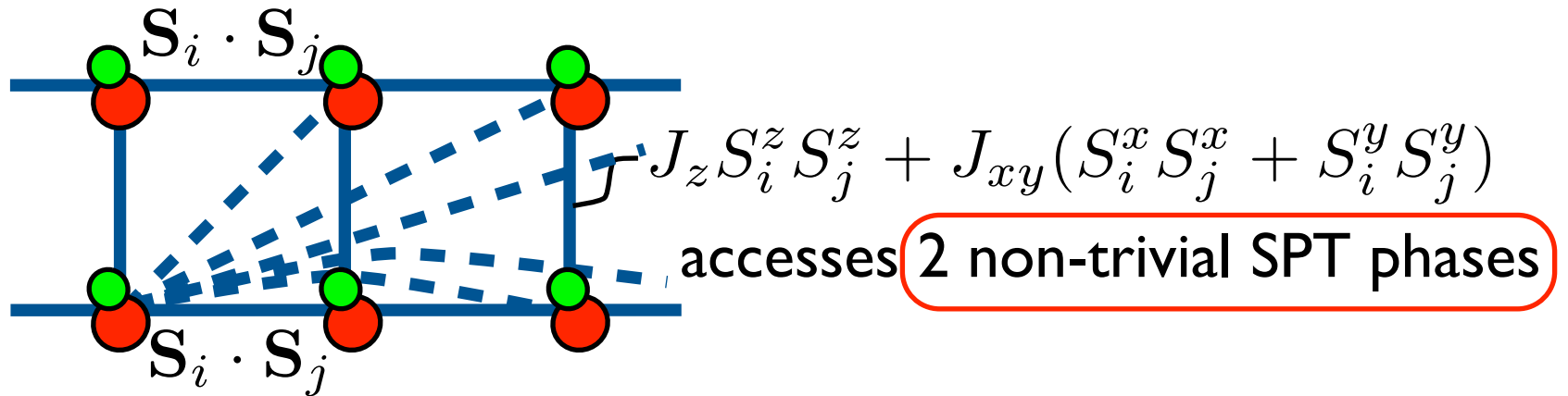


eigenvalues of  $\text{tr}_B [|\Psi\rangle\langle\Psi|]$   
 are doubly degenerate

Liu, Wen, et al, PRB 2012; Chen, Gu, Wen, PRB 2011; Pollman et al, PRB 2010



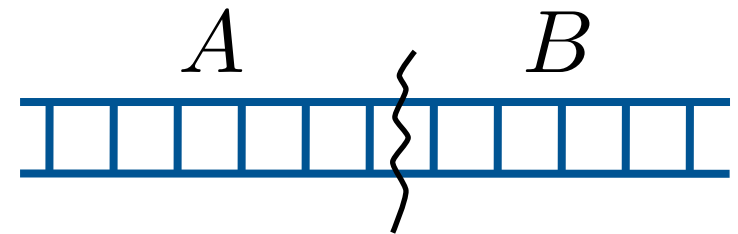
# SPT phases in spin-1/2 ladders



- do SPT phases survive?

**YES!**  
 [PRB 87, 081106(R) (2013)]

- symmetry-protected features:
  - 2-fold-degenerate edge states
  - 2-fold-degenerate entanglement spectrum:



eigenvalues of  $\text{tr}_B [|\Psi\rangle\langle\Psi|]$   
 are doubly degenerate

Liu, Wen, et al, PRB 2012; Chen, Gu, Wen, PRB 2011; Pollman et al, PRB 2010

# Conclusions and Outlook

# Conclusions and Outlook

- simple microwave dressing gives access to interacting SPT phases

# Conclusions and Outlook

- simple microwave dressing gives access to interacting SPT phases
- SPT phases survive in the presence of long-range interactions

# Conclusions and Outlook

- simple microwave dressing gives access to interacting SPT phases
- SPT phases survive in the presence of long-range interactions
- hope that same holds for more interesting examples  
[PRB 87, 081106(R) (2013) & arXiv:1301.5636]:
  - e.g. Kitaev honeycomb, bilinear-biquadratic spin-1 model, etc...

# Conclusions and Outlook

- simple microwave dressing gives access to interacting SPT phases
- SPT phases survive in the presence of long-range interactions
- hope that same holds for more interesting examples  
[PRB 87, 081106(R) (2013) & arXiv:1301.5636]:
  - e.g. Kitaev honeycomb, bilinear-biquadratic spin-1 model, etc...
- classification of SPT phases in the presence of long-range interactions?

# Conclusions and Outlook

- simple microwave dressing gives access to interacting SPT phases
- SPT phases survive in the presence of long-range interactions
- hope that same holds for more interesting examples  
[PRB 87, 081106(R) (2013) & arXiv:1301.5636]:
  - e.g. Kitaev honeycomb, bilinear-biquadratic spin-1 model, etc...
- classification of SPT phases in the presence of long-range interactions?
- in general, peculiar effects of long-range interactions  
[see e.g. poster by H.P. Buchler & other talks and posters]

Topological flat bands  
and  
fractional Chern insulators

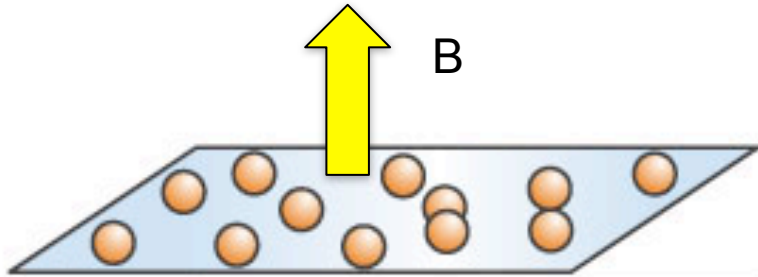


# Motivation

- fractional quantum Hall effect (FQHE):

# Motivation

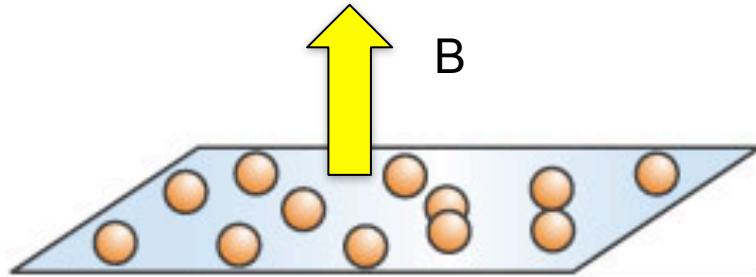
- fractional quantum Hall effect (FQHE):
  - 2D electron gas in a strong magnetic field  $B$



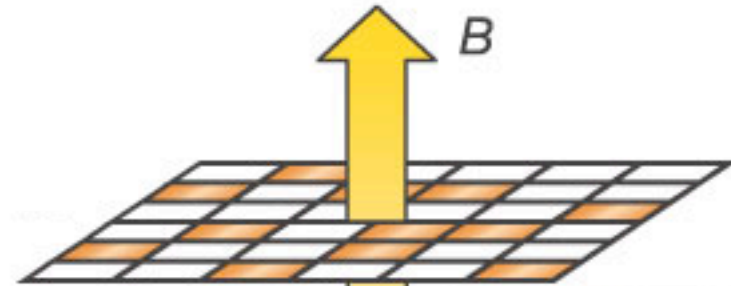
[Eisenstein & MacDonald, Nature (2004)]

# Motivation

- fractional quantum Hall effect (FQHE):
  - 2D electron gas in a strong magnetic field  $B$



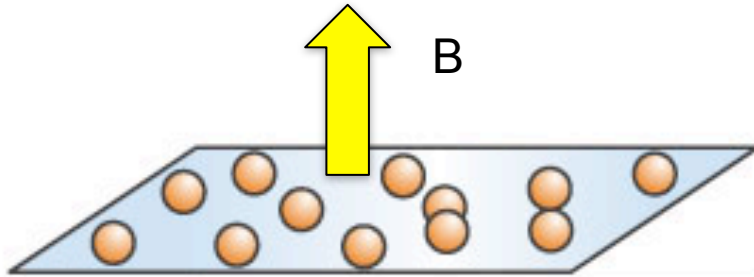
[Eisenstein & MacDonald, Nature (2004)]



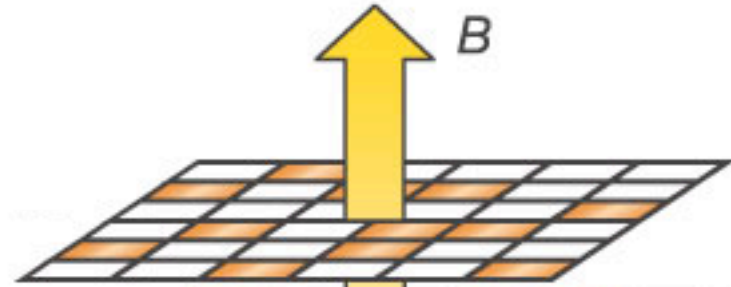
lowest Landau level

# Motivation

- fractional quantum Hall effect (FQHE):
  - 2D electron gas in a strong magnetic field  $B$



[Eisenstein & MacDonald, Nature (2004)]



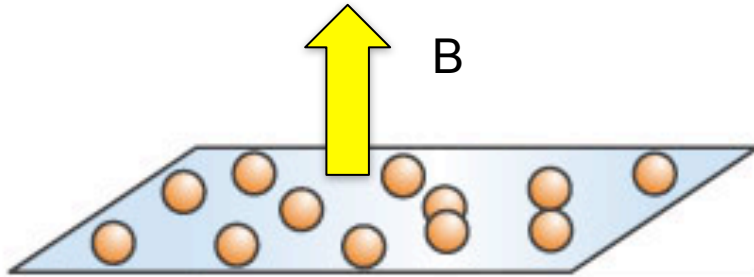
lowest Landau level

- huge number of degenerate  
cyclotron orbitals

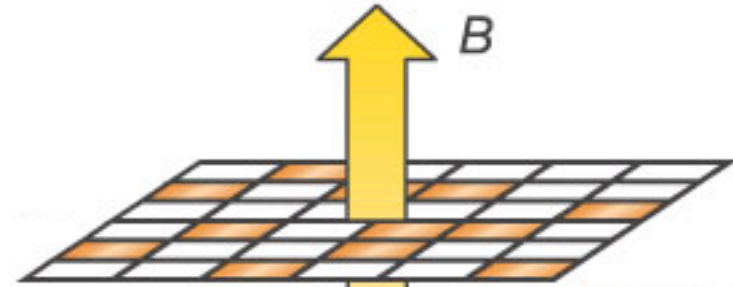
# Motivation

- fractional quantum Hall effect (FQHE):
  - 2D electron gas in a strong magnetic field  $B$

integer quantum  
Hall effect (IQHE)



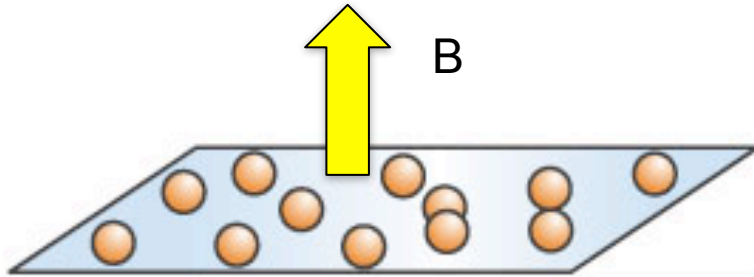
[Eisenstein & MacDonald, Nature (2004)]



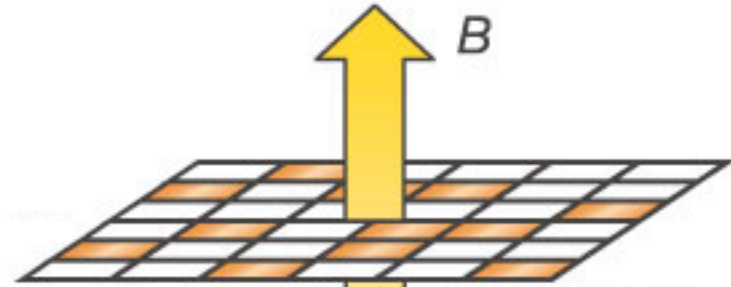
- lowest Landau level
- huge number of degenerate  
cyclotron orbitals

# Motivation

- fractional quantum Hall effect (FQHE):
  - 2D electron gas in a strong magnetic field  $B$



[Eisenstein & MacDonald, Nature (2004)]

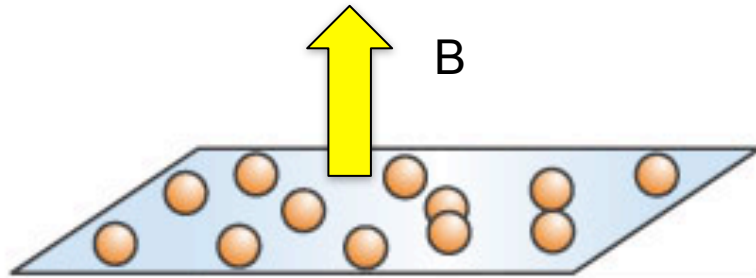


lowest Landau level

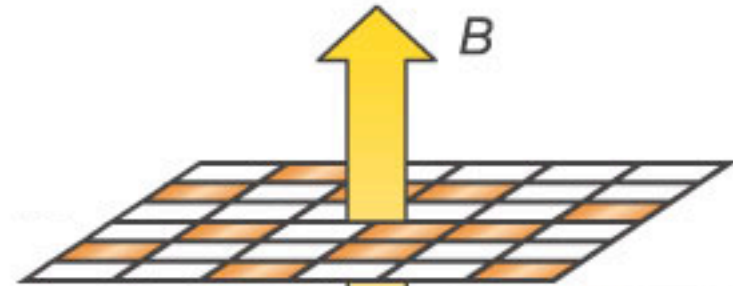
- huge number of degenerate  
cyclotron orbitals

# Motivation

- fractional quantum Hall effect (FQHE):
  - 2D electron gas in a strong magnetic field  $B$



[Eisenstein & MacDonald, Nature (2004)]



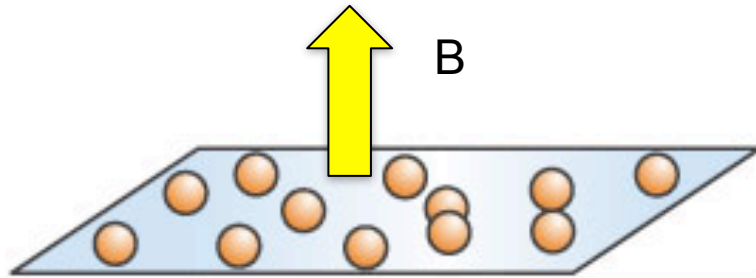
lowest Landau level

- huge number of degenerate **cyclotron orbitals**

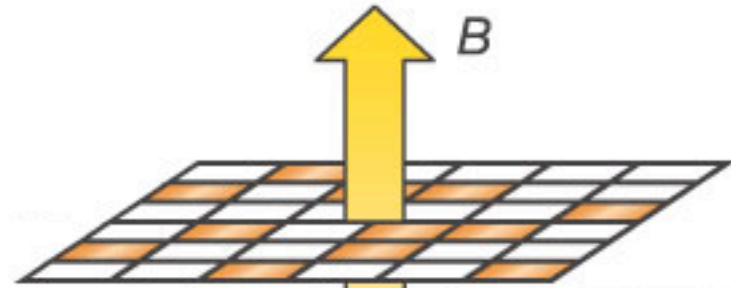
- Coulomb interactions break the degeneracy in a partially filled Landau level and can open a bulk gap  $\Rightarrow$  FQHE

# Motivation

- fractional quantum Hall effect (FQHE):
  - 2D electron gas in a strong magnetic field  $B$



[Eisenstein & MacDonald, Nature (2004)]



lowest Landau level

- huge number of degenerate **cyclotron orbitals**

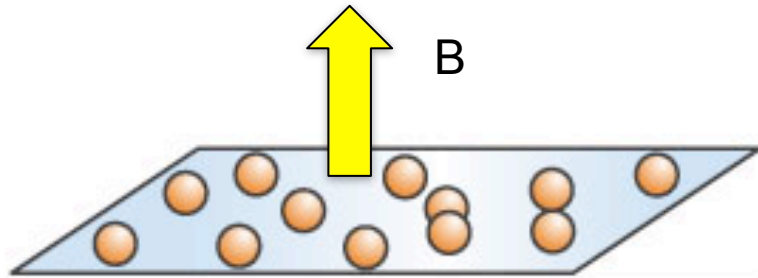
- Coulomb interactions break the degeneracy in a partially filled Landau level and can open a bulk gap  $\Rightarrow$  FQHE

- **key question:** can one find FQHE physics
  - on a lattice (not on a continuum) &
  - without a magnetic field?

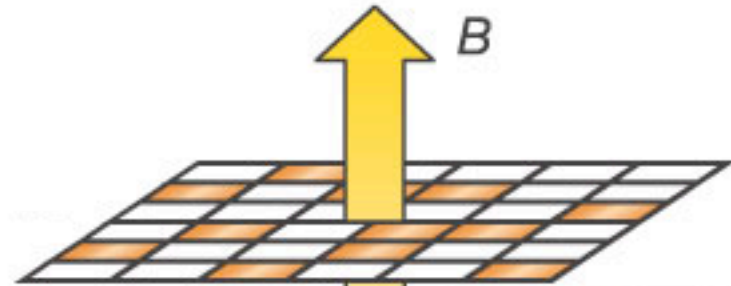


# Motivation

- fractional quantum Hall effect (FQHE):
  - 2D electron gas in a strong magnetic field  $B$



[Eisenstein & MacDonald, Nature (2004)]



lowest Landau level

- huge number of degenerate cyclotron orbitals

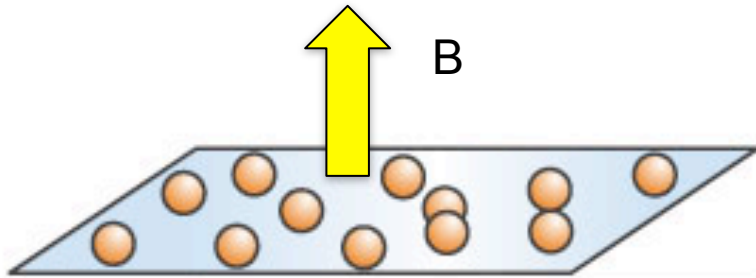
- Coulomb interactions break the degeneracy in a partially filled Landau level and can open a bulk gap  $\Rightarrow$  FQHE

- key question: can one find FQHE physics
  - on a lattice (not on a continuum) &
  - without a magnetic field?
- yes! Fractional Chern Insulators. Need:
  - flat single-particle energy bands [degeneracy]

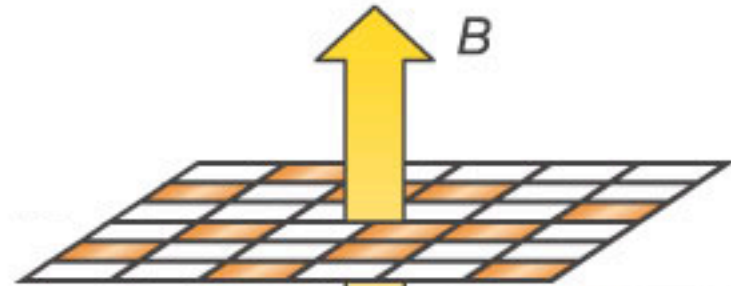
# Motivation

- fractional quantum Hall effect (FQHE):

- 2D electron gas in a strong magnetic field  $B$



[Eisenstein & MacDonald, Nature (2004)]



lowest Landau level

- huge number of degenerate cyclotron orbitals

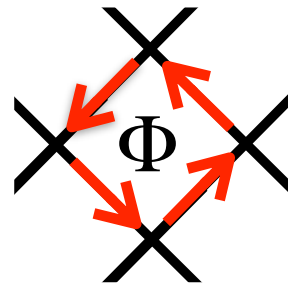
- Coulomb interactions break the degeneracy in a partially filled Landau level and can open a bulk gap  $\Rightarrow$  FQHE

- key question: can one find FQHE physics

- on a lattice (not on a continuum) &
- without a magnetic field?

$$\prod t_{ij} \propto e^{i\Phi}$$

[as in Spielman, etc...]



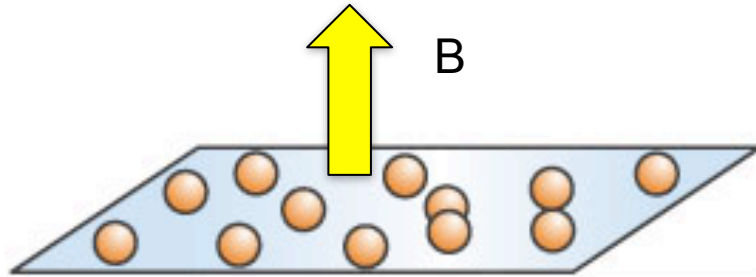
- yes! Fractional Chern Insulators. Need:

- flat single-particle energy bands [degeneracy]
- complex hopping amplitudes (like Aharonov-Bohm phases) [topology]

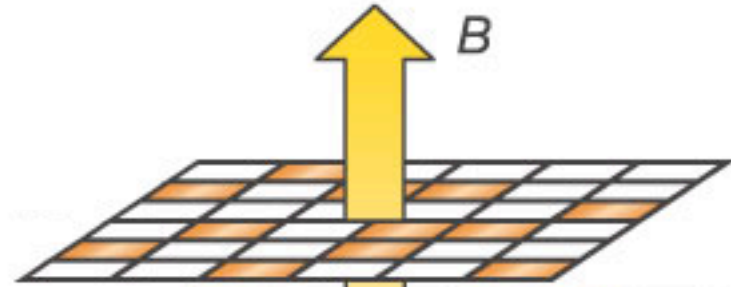
# Motivation

- fractional quantum Hall effect (FQHE):

- 2D electron gas in a strong magnetic field  $B$



[Eisenstein & MacDonald, Nature (2004)]



lowest Landau level

- huge number of degenerate **cyclotron orbitals**

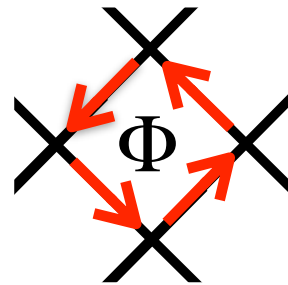
- Coulomb interactions break the degeneracy in a partially filled Landau level and can open a bulk gap  $\Rightarrow$  FQHE

- key question: can one find FQHE physics

- on a lattice (not on a continuum) &
- without a magnetic field?

$$\prod t_{ij} \propto e^{i\Phi}$$

[as in Spielman, etc...]



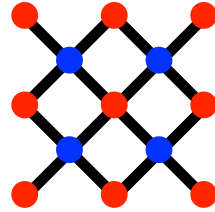
- yes! Fractional **Chern** Insulators. Need:

- flat single-particle energy bands [degeneracy]
- complex hopping amplitudes (like Aharonov-Bohm phases) **[topology]**

# Chern number and topology

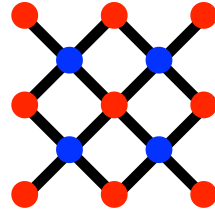
# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell



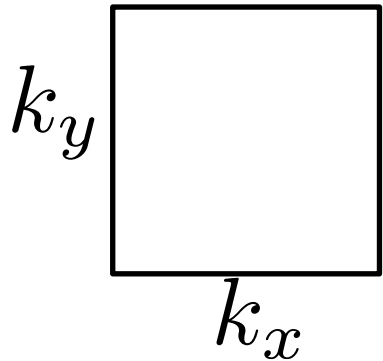
# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell



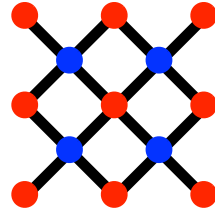
$$\mathbf{k} = (k_x, k_y)$$

Brillouin zone



# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell



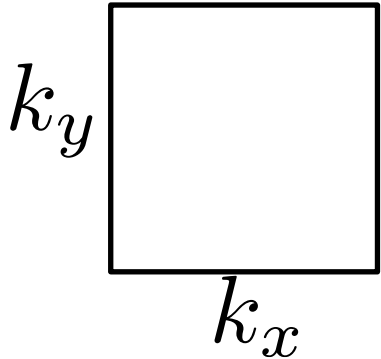
- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \sigma + f(\mathbf{k})$$

└ Pauli matrices

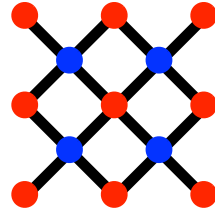
$$\mathbf{k} = (k_x, k_y)$$

Brillouin zone



# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell



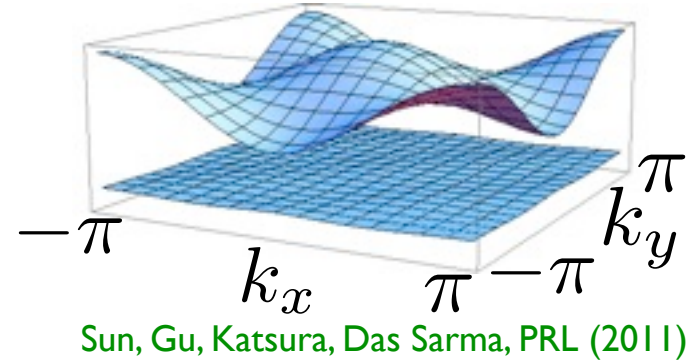
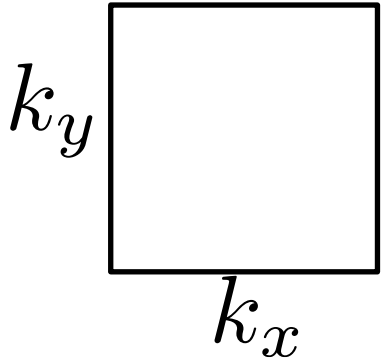
- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

( Pauli matrices

$$\mathbf{k} = (k_x, k_y)$$

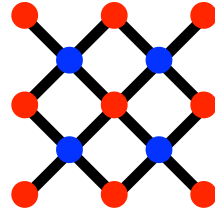
Brillouin zone





# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell



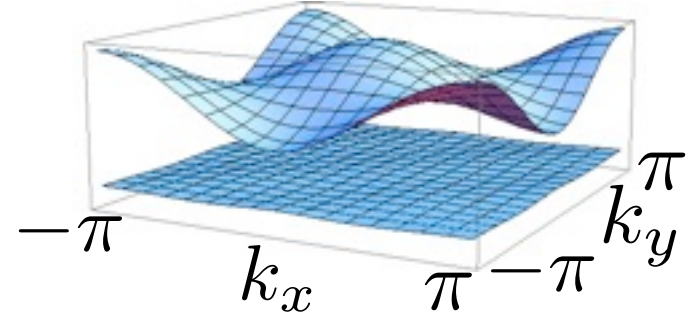
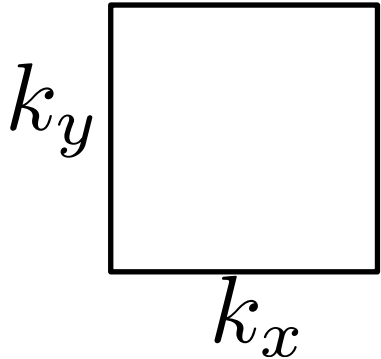
- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

( Pauli matrices

$$\mathbf{k} = (k_x, k_y)$$

Brillouin zone

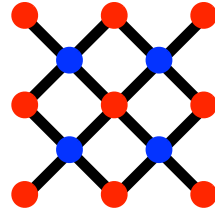


Sun, Gu, Katsura, Das Sarma, PRL (2011)

- topological flat band

# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell



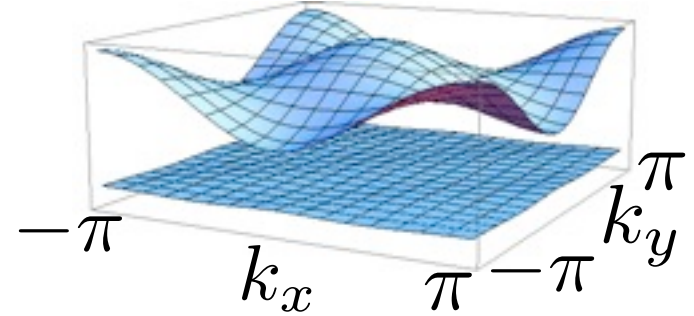
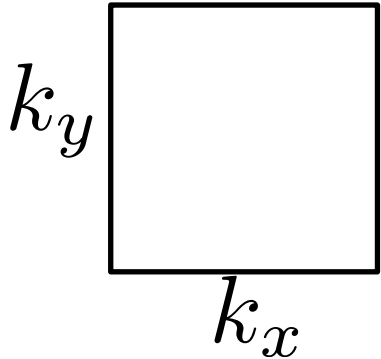
- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

( Pauli matrices

$$\mathbf{k} = (k_x, k_y)$$

Brillouin zone

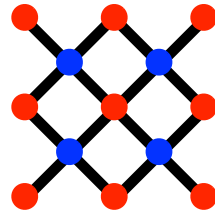


Sun, Gu, Katsura, Das Sarma, PRL (2011)

- topological **flat** band

# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell



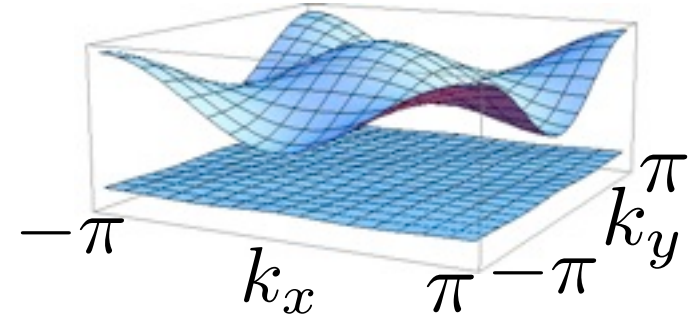
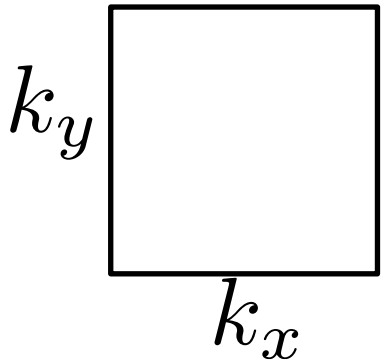
- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

( Pauli matrices

$$\mathbf{k} = (k_x, k_y)$$

Brillouin zone

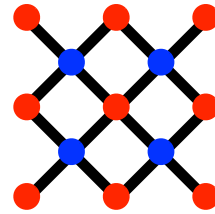


Sun, Gu, Katsura, Das Sarma, PRL (2011)

- **topological** flat band

# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

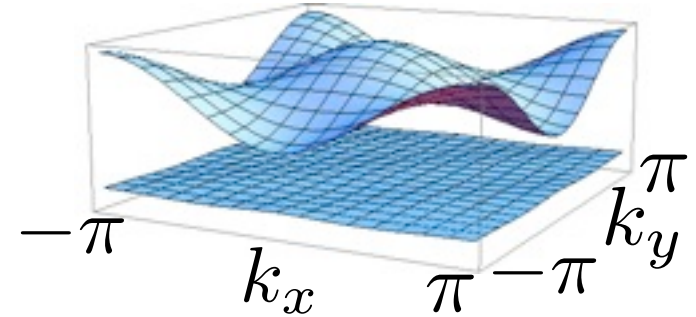


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

( Pauli matrices

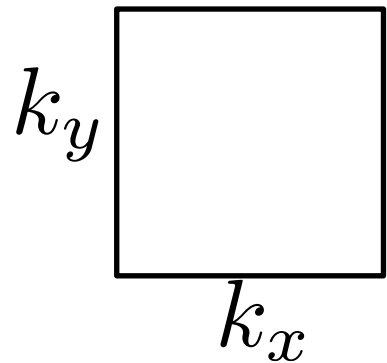
$$\mathbf{k} = (k_x, k_y)$$



Sun, Gu, Katsura, Das Sarma, PRL (2011)

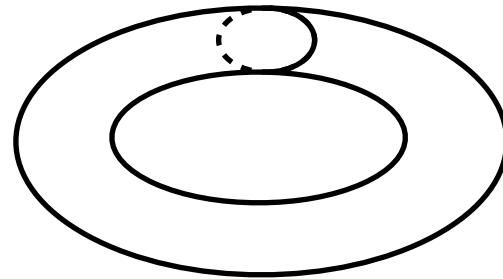
- topological flat band

Brillouin zone



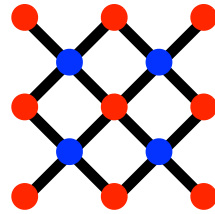
=

torus



# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

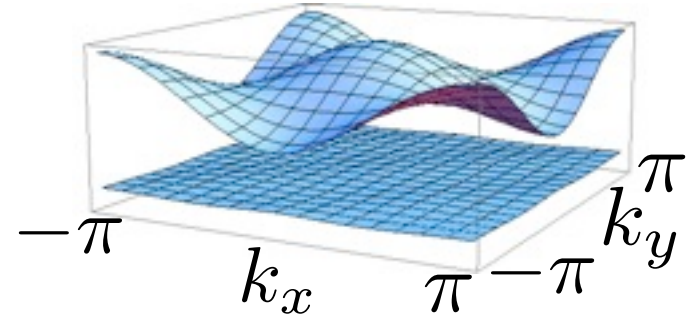


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

$\underbrace{\quad}_{\text{Pauli matrices}}$

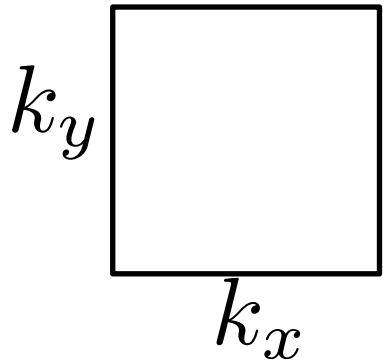
$$\mathbf{k} = (k_x, k_y)$$



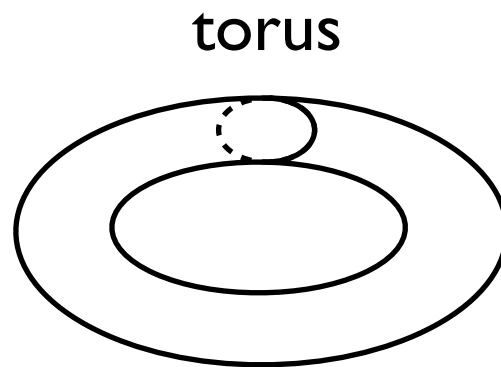
Sun, Gu, Katsura, Das Sarma, PRL (2011)

- topological flat band

Brillouin zone



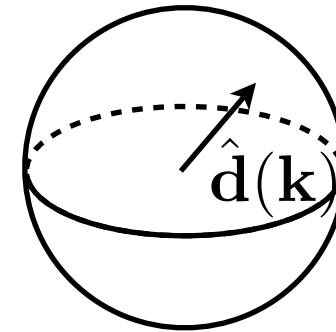
=



torus

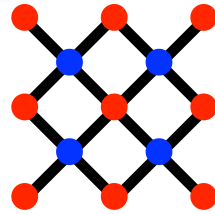


unit sphere



# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

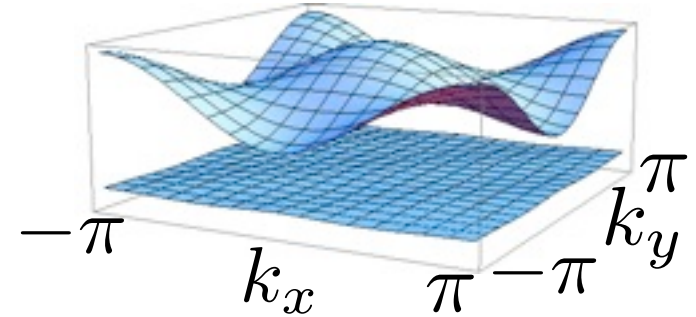


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

$\underbrace{\quad}_{\text{Pauli matrices}}$

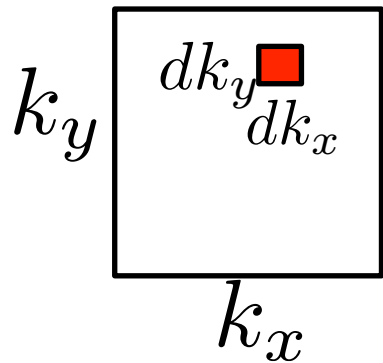
$$\mathbf{k} = (k_x, k_y)$$



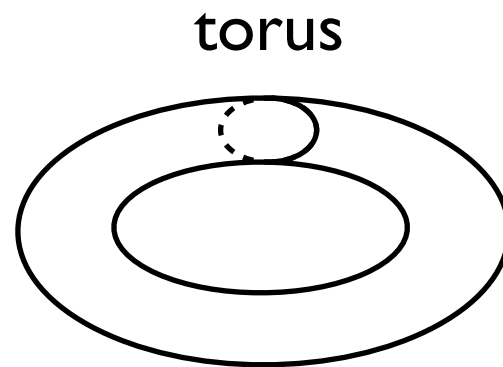
Sun, Gu, Katsura, Das Sarma, PRL (2011)

- topological flat band

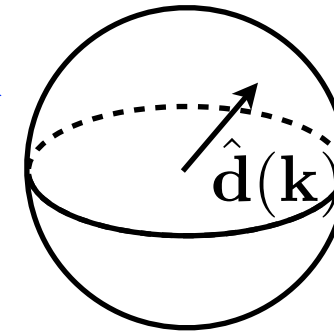
Brillouin zone



=

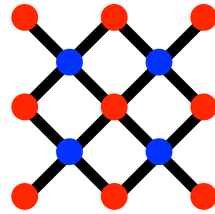


unit sphere



# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

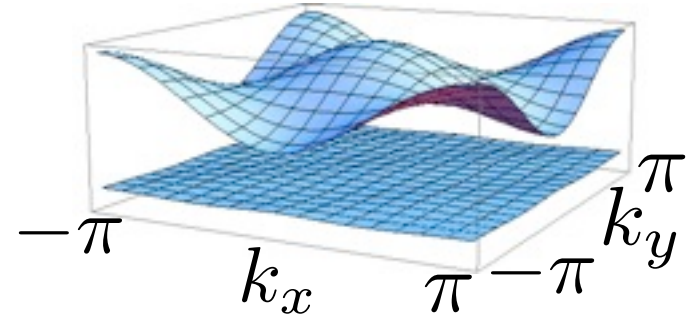


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

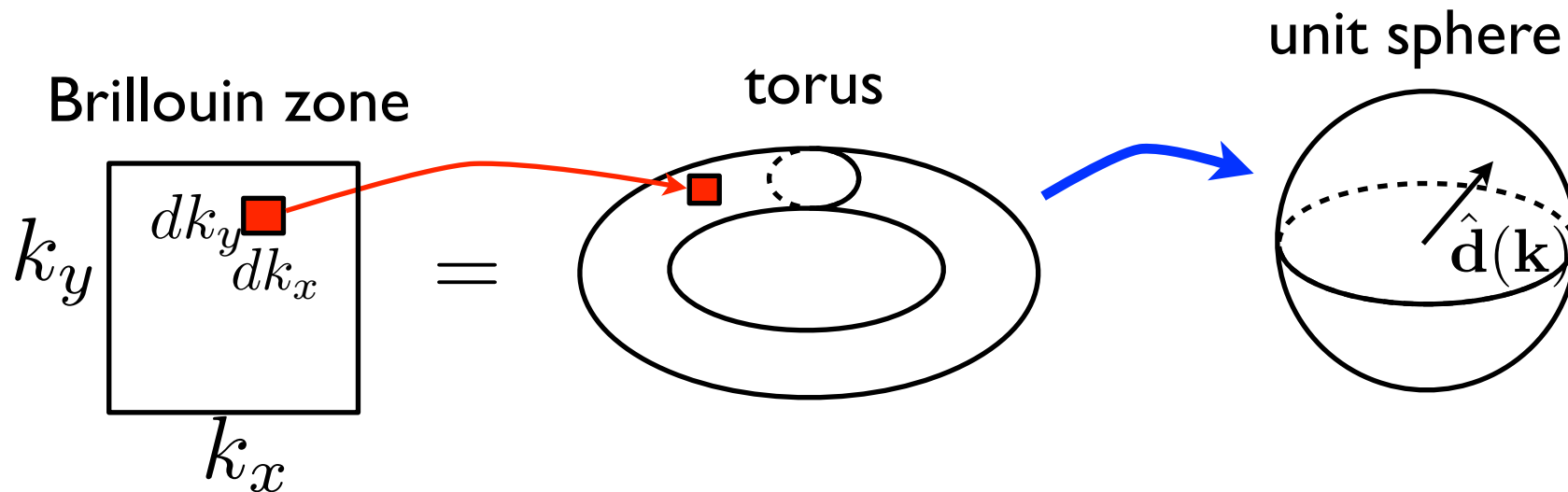
$\underbrace{\quad}_{\text{Pauli matrices}}$

$$\mathbf{k} = (k_x, k_y)$$



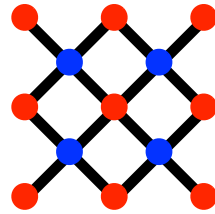
Sun, Gu, Katsura, Das Sarma, PRL (2011)

- topological flat band



# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

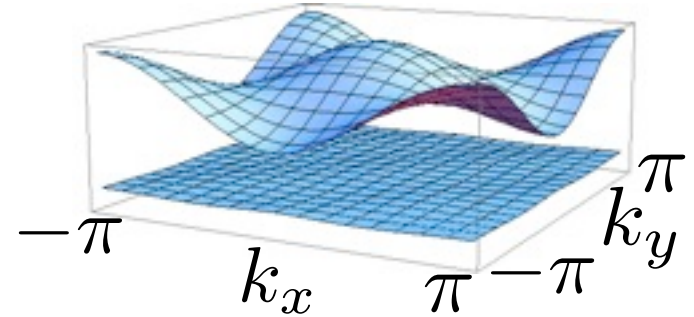


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

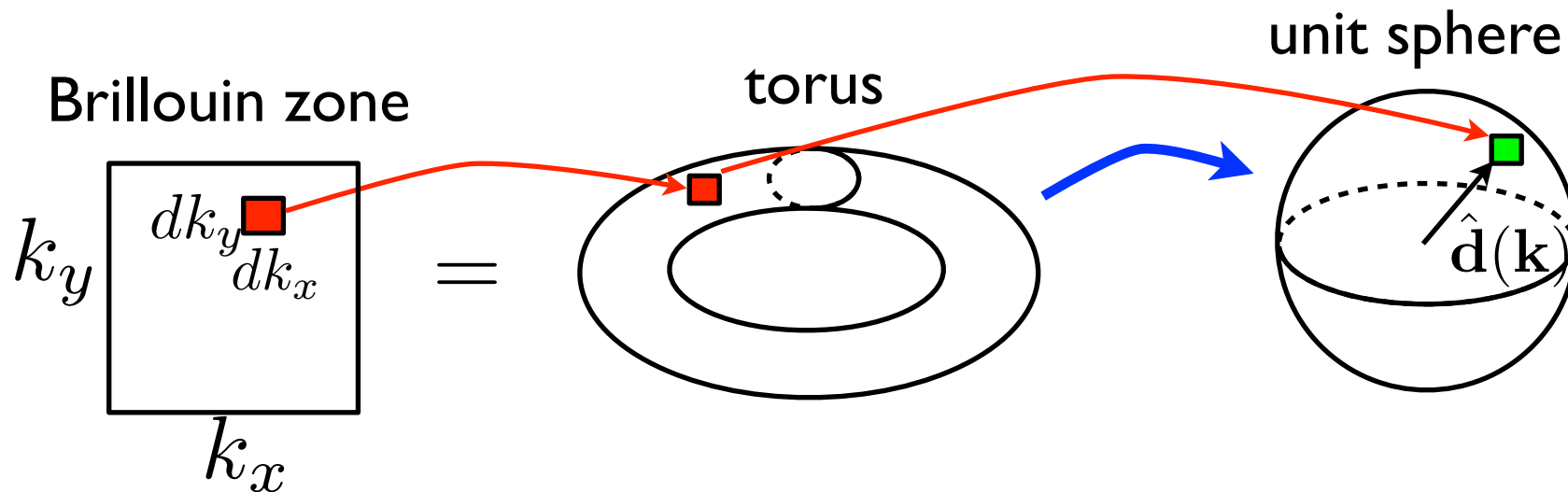
$\underbrace{\quad}_{\text{Pauli matrices}}$

$$\mathbf{k} = (k_x, k_y)$$



Sun, Gu, Katsura, Das Sarma, PRL (2011)

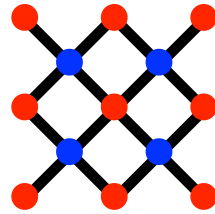
- topological flat band





# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

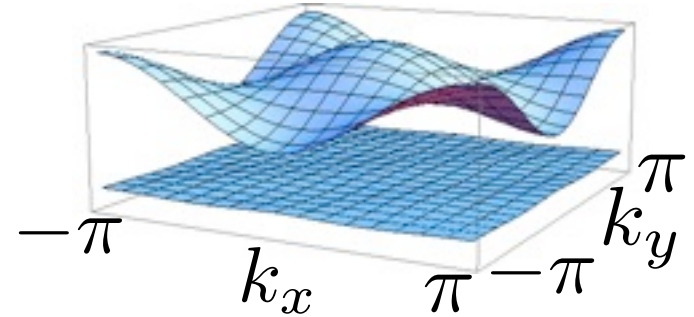


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

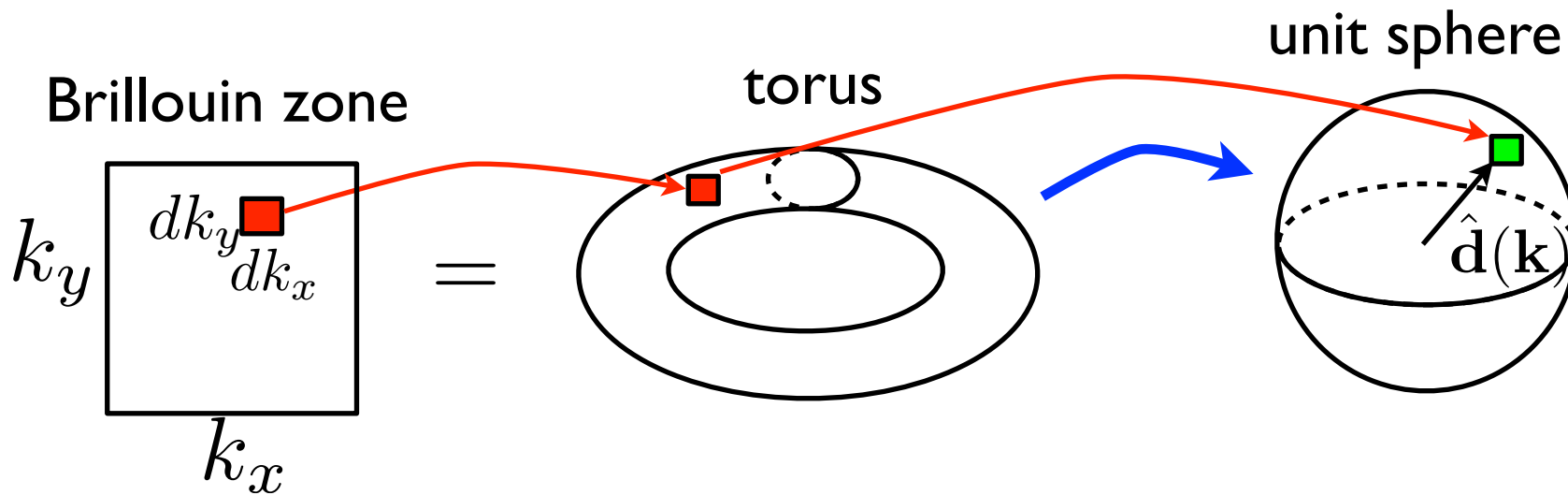
$\underbrace{\quad}_{\text{Pauli matrices}}$

$$\mathbf{k} = (k_x, k_y)$$



Sun, Gu, Katsura, Das Sarma, PRL (2011)

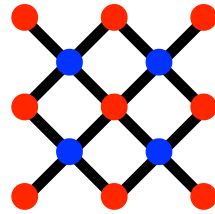
- topological flat band



$$\underbrace{dk_x dk_y \hat{\mathbf{d}} \cdot (\partial_{k_x} \hat{\mathbf{d}} \times \partial_{k_y} \hat{\mathbf{d}})}_{\text{area of } \blacksquare}$$

# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

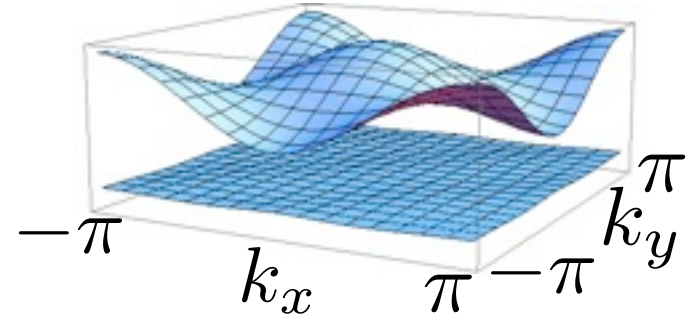


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

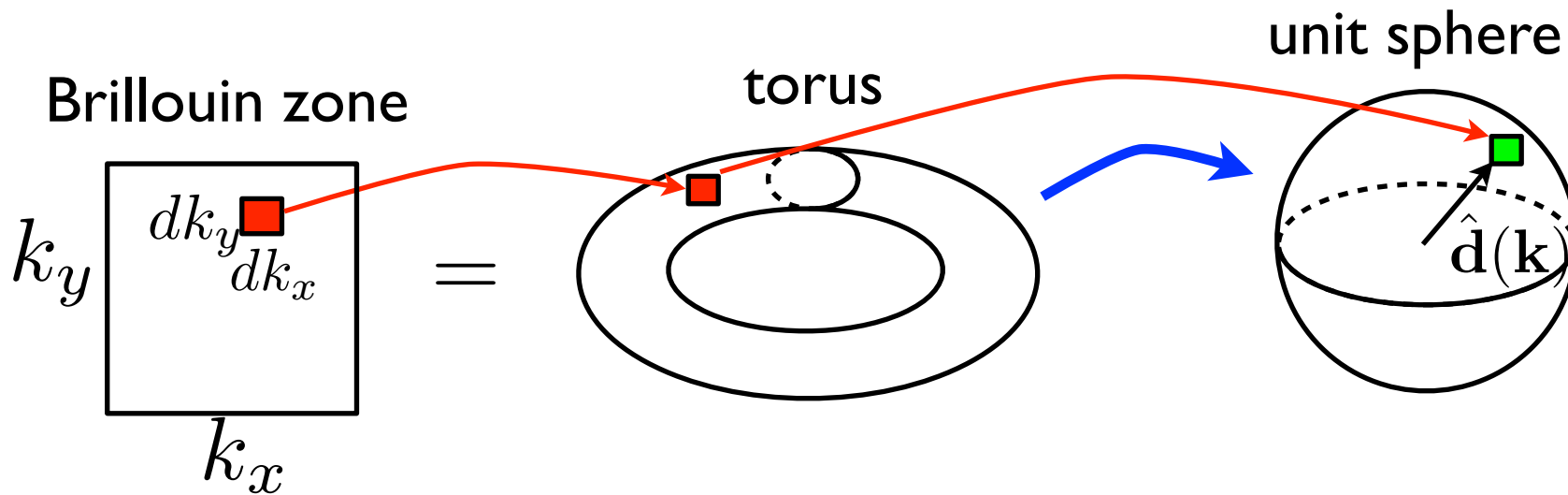
$\underbrace{\quad}_{\text{Pauli matrices}}$

$$\mathbf{k} = (k_x, k_y)$$



Sun, Gu, Katsura, Das Sarma, PRL (2011)

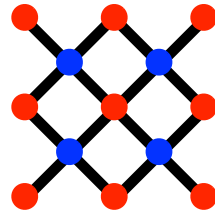
- topological flat band



$$C = \frac{1}{4\pi} \int dk_x dk_y \underbrace{\hat{\mathbf{d}} \cdot (\partial_{k_x} \hat{\mathbf{d}} \times \partial_{k_y} \hat{\mathbf{d}})}_{\text{area of } \blacksquare}$$

# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

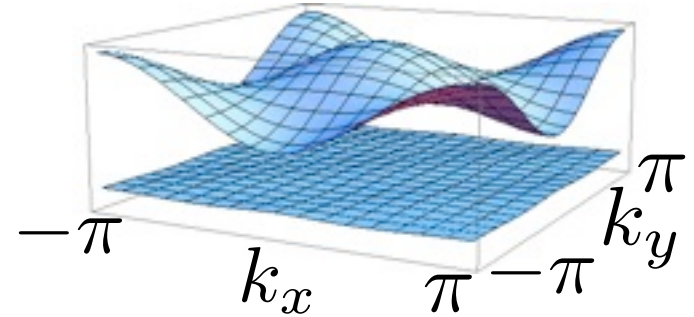


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

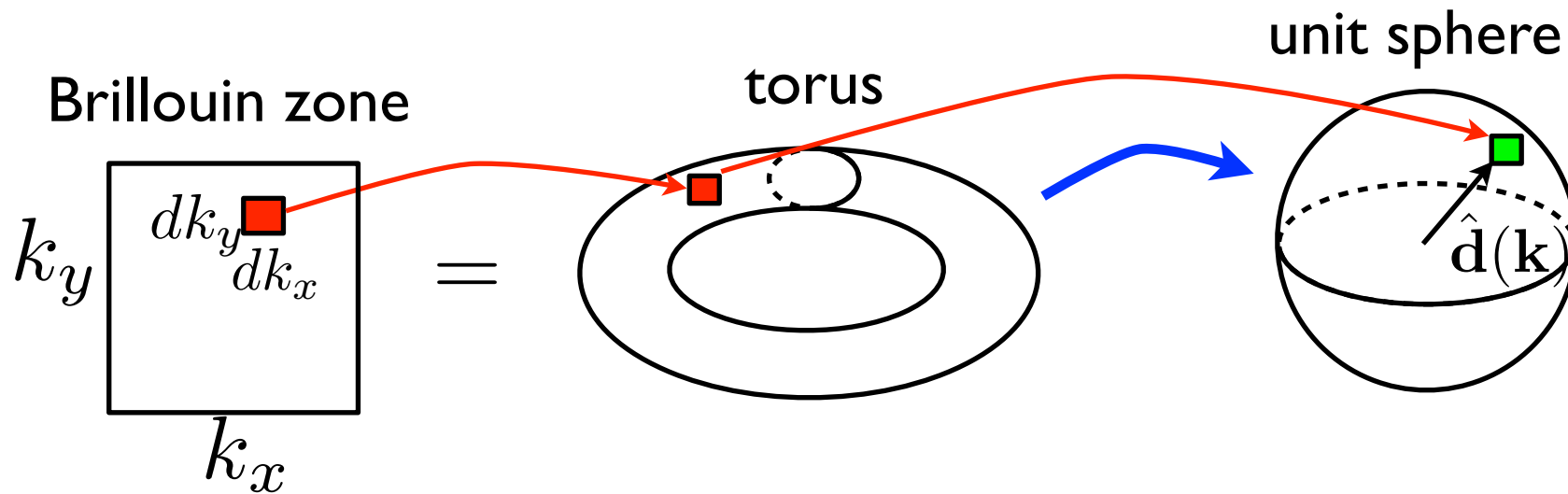
$\underbrace{\quad}_{\text{Pauli matrices}}$

$$\mathbf{k} = (k_x, k_y)$$



Sun, Gu, Katsura, Das Sarma, PRL (2011)

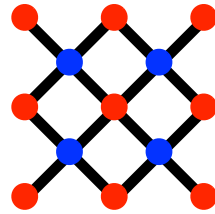
- topological flat band



$$C = \frac{1}{4\pi} \int \underbrace{dk_x dk_y \hat{\mathbf{d}} \cdot (\partial_{k_x} \hat{\mathbf{d}} \times \partial_{k_y} \hat{\mathbf{d}})}_{\text{area of } \blacksquare} = \text{number of times the torus wraps over the sphere}$$

# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

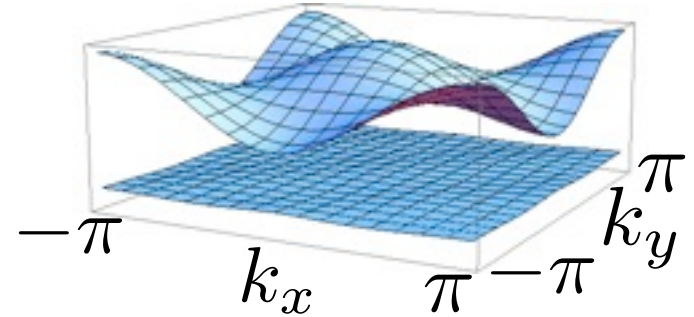


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

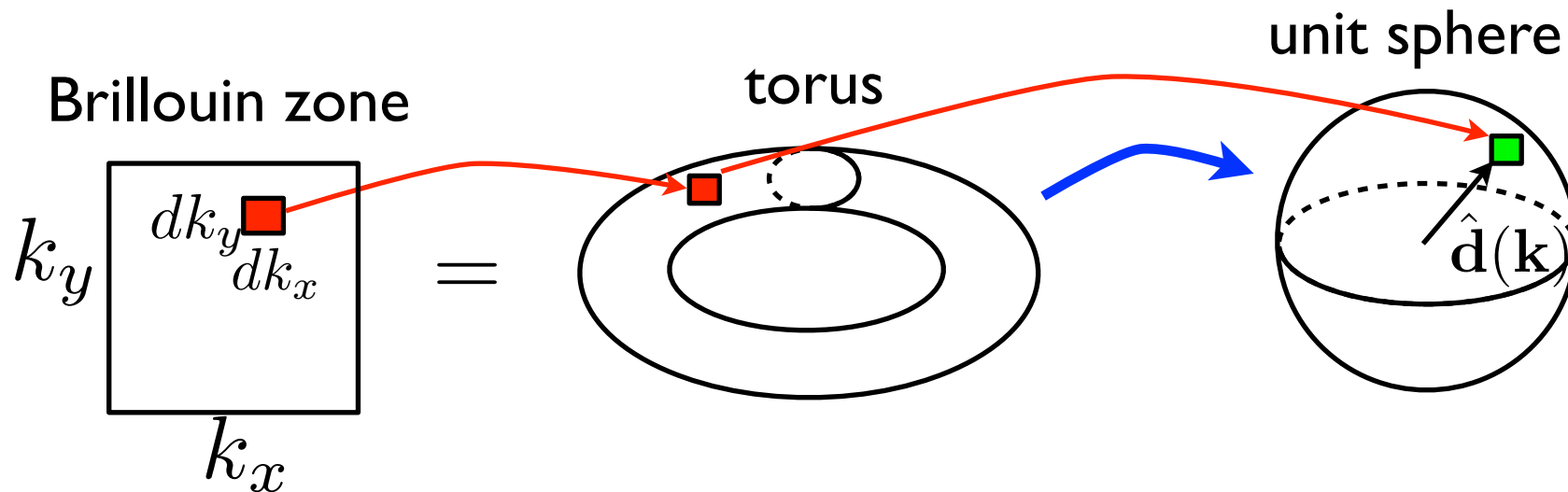
$\underbrace{\quad}_{\text{Pauli matrices}}$

$$\mathbf{k} = (k_x, k_y)$$



Sun, Gu, Katsura, Das Sarma, PRL (2011)

- topological flat band

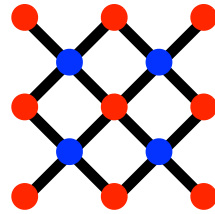


$$C = \frac{1}{4\pi} \int dk_x dk_y \underbrace{\hat{\mathbf{d}} \cdot (\partial_{k_x} \hat{\mathbf{d}} \times \partial_{k_y} \hat{\mathbf{d}})}_{\text{area of } \blacksquare} = \text{number of times the torus wraps over the sphere}$$

**Chern number**

# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

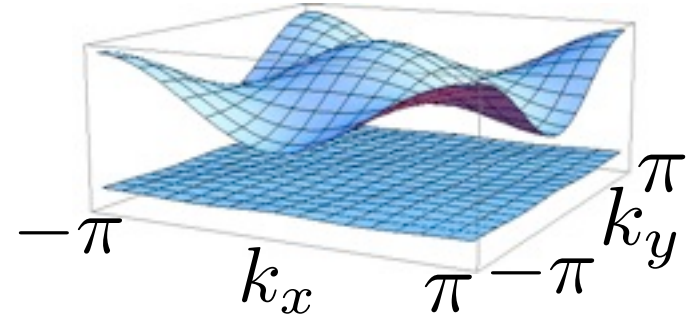


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

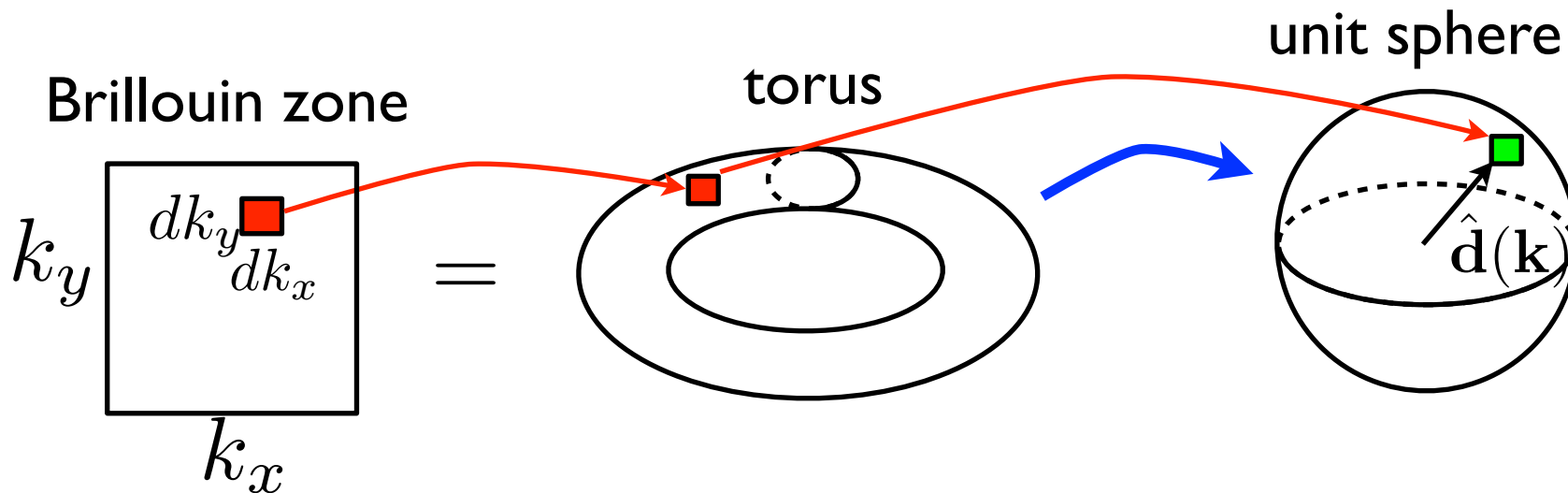
$\underbrace{\quad}_{\text{Pauli matrices}}$

$$\mathbf{k} = (k_x, k_y)$$



Sun, Gu, Katsura, Das Sarma, PRL (2011)

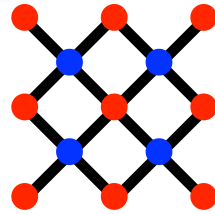
- topological** flat band



$$C = \frac{1}{4\pi} \int dk_x dk_y \underbrace{\hat{\mathbf{d}} \cdot (\partial_{k_x} \hat{\mathbf{d}} \times \partial_{k_y} \hat{\mathbf{d}})}_{\text{area of } \blacksquare} = \text{number of times the torus wraps over the sphere}$$

# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

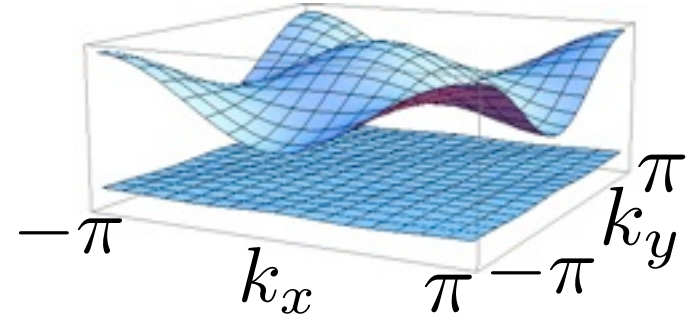


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

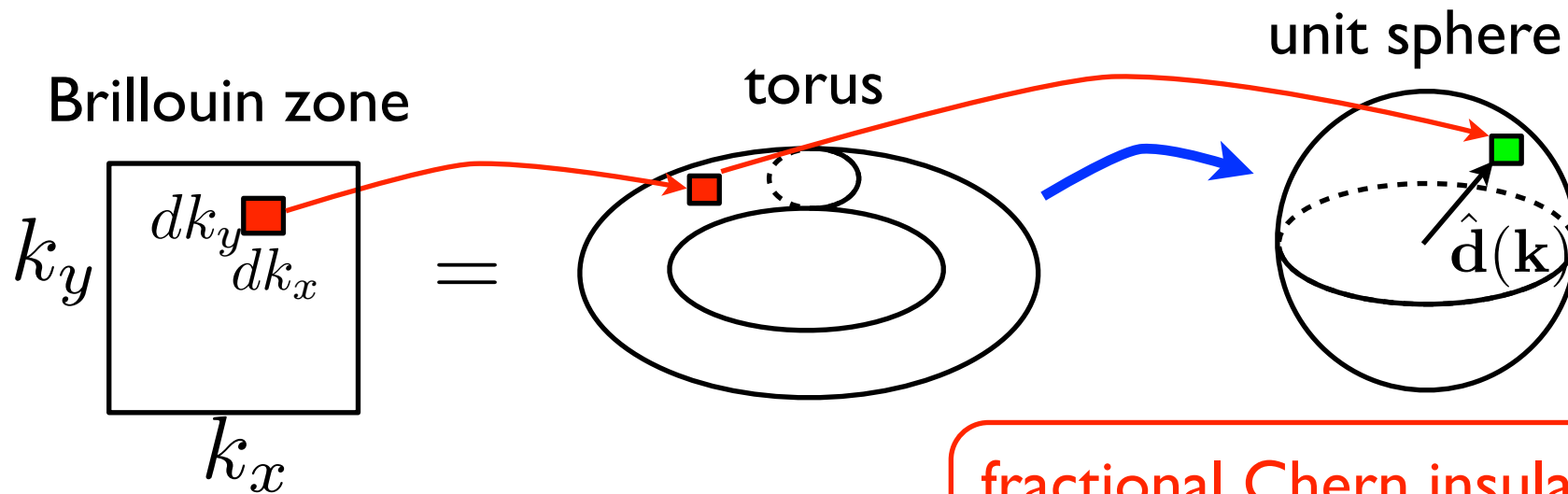
$\underbrace{\quad}_{\text{Pauli matrices}}$

$$\mathbf{k} = (k_x, k_y)$$



Sun, Gu, Katsura, Das Sarma, PRL (2011)

- topological flat band



fractional Chern insulator (FCI)

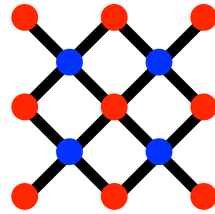
$$C = \frac{1}{4\pi} \int dk_x dk_y \hat{\mathbf{d}} \cdot (\partial_{k_x} \hat{\mathbf{d}} \times \partial_{k_y} \hat{\mathbf{d}})$$

= number of times the torus wraps over the sphere

Chern number area of ■

# Chern number and topology

- single particle hopping in a 2D square lattice with a two-site unit cell

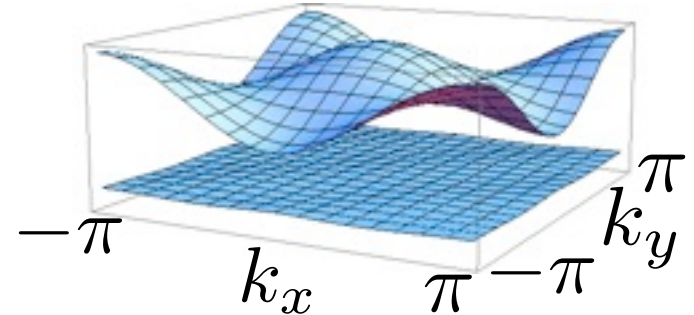


- 2x2 Bloch Hamiltonian

$$H(\mathbf{k}) = \mathbf{d}(\mathbf{k}) \cdot \boldsymbol{\sigma} + f(\mathbf{k})$$

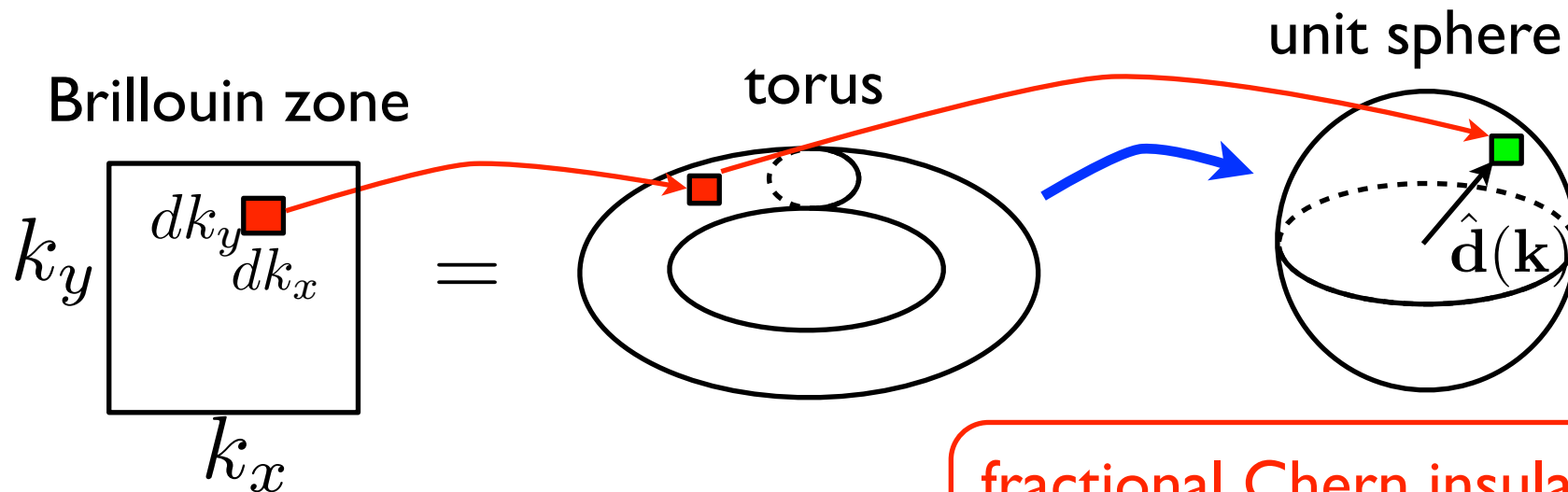
$\underbrace{\quad}_{\text{Pauli matrices}}$

$$\mathbf{k} = (k_x, k_y)$$



Sun, Gu, Katsura, Das Sarma, PRL (2011)

- **topological flat band**



**fractional Chern insulator (FCI)**

$$C = \frac{1}{4\pi} \int dk_x dk_y \underbrace{\hat{\mathbf{d}} \cdot (\partial_{k_x} \hat{\mathbf{d}} \times \partial_{k_y} \hat{\mathbf{d}})}_{\text{area of } \blacksquare}$$

= number of times the torus wraps over the sphere

**Chern number** area of  $\blacksquare$

# Topological flat bands and fractional Chern insulators

F. D. M. Haldane, PRL 61, 2015 (1988) [IQHE]

E. Tang, J-W. Mei, and X-G. Wen, PRL 106, 236802 (2011)

T. Neupert, L. Santos, C. Chamon, and C. Mudry, PRL 106, 236804 (2011)

K. Sun, Z. -C Gu, H. Katsura, and S. Das Sarma, PRL 106, 236803 (2011)

Y.-F. Wang, Z. -C Gu, C.-D. Gong, D. Sheng, PRL 107 146803 (2011)

N. Regnault & B.A. Bernevig, PRX 1 021014 (2011)

Z. Liu, E. J. Bergholtz, H. Fan, R. Moessner, A. M. Laeuchli PRL 109, 186805 (2012), arXiv: 1207.6094 (2012)

S. Yang, Z.-C. Gu, K. Sun, S. Das Sarma, arXiv:1205.5792 (2012)

... many more ...



# Topological flat bands and fractional Chern insulators

F. D. M. Haldane, PRL 61, 2015 (1988) [IQHE]

E. Tang, J-W. Mei, and X-G. Wen, PRL 106, 236802 (2011)

T. Neupert, L. Santos, C. Chamon, and C. Mudry, PRL 106, 236804 (2011)

K. Sun, Z. -C Gu, H. Katsura, and S. Das Sarma, PRL 106, 236803 (2011)

Y.-F. Wang, Z. -C Gu, C.-D. Gong, D. Sheng, PRL 107 146803 (2011)

N. Regnault & B.A. Bernevig, PRX 1 021014 (2011)

Z. Liu, E. J. Bergholtz, H. Fan, R. Moessner, A. M. Laeuchli PRL 109, 186805 (2012), arXiv: 1207.6094 (2012)

S. Yang, Z.-C. Gu, K. Sun, S. Das Sarma, arXiv:1205.5792 (2012)

... many more ...

- **realistic physical system missing** [recent exception: Cooper & Dalibard, 2012]

# Topological flat bands and fractional Chern insulators

F. D. M. Haldane, PRL 61, 2015 (1988) [IQHE]

E. Tang, J-W. Mei, and X-G. Wen, PRL 106, 236802 (2011)

T. Neupert, L. Santos, C. Chamon, and C. Mudry, PRL 106, 236804 (2011)

K. Sun, Z. -C Gu, H. Katsura, and S. Das Sarma, PRL 106, 236803 (2011)

Y.-F. Wang, Z. -C Gu, C.-D. Gong, D. Sheng, PRL 107 146803 (2011)

N. Regnault & B.A. Bernevig, PRX 1 021014 (2011)

Z. Liu, E. J. Bergholtz, H. Fan, R. Moessner, A. M. Laeuchli PRL 109, 186805 (2012), arXiv: 1207.6094 (2012)

S. Yang, Z.-C. Gu, K. Sun, S. Das Sarma, arXiv:1205.5792 (2012)

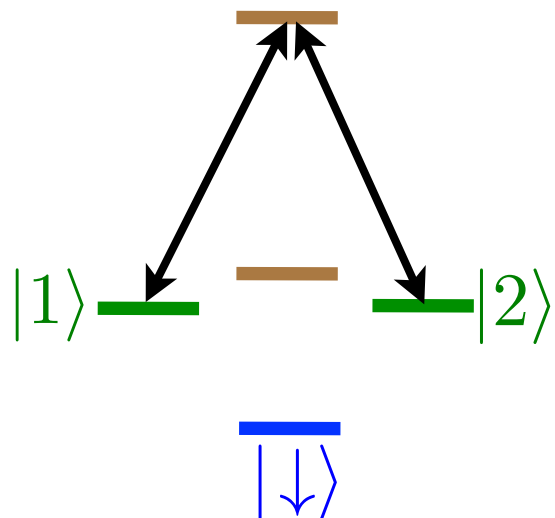
... many more ...

- **realistic physical system missing** [recent exception: Cooper & Dalibard, 2012]
- **dipolar spin systems (e.g. polar molecules) naturally admit topological flat bands and fractional Chern insulator ( $\sim$ FQHE) ground states!**

# Fractional Chern insulator with polar molecules

# Fractional Chern insulator with polar molecules

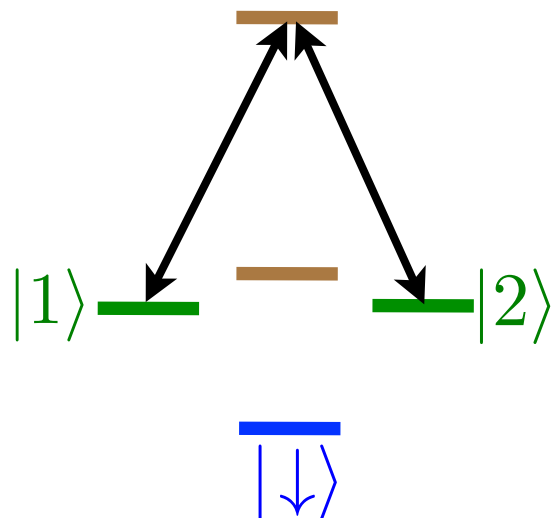
Recall SPT discussion:



$$|\uparrow\rangle = x|1\rangle + y|2\rangle$$

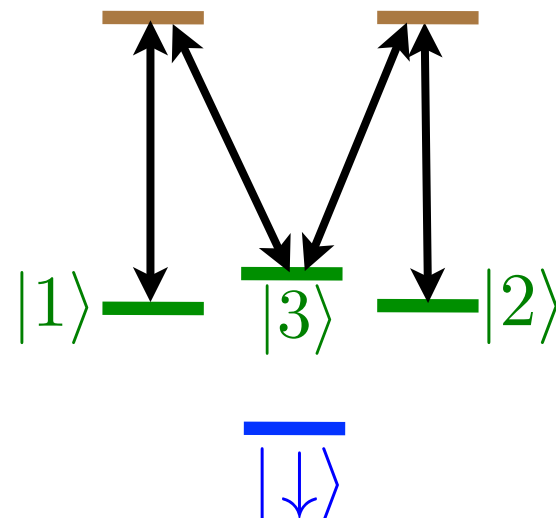
# Fractional Chern insulator with polar molecules

Recall SPT discussion:



$$|\uparrow\rangle = x|1\rangle + y|2\rangle$$

To get fractional Chern insulators:

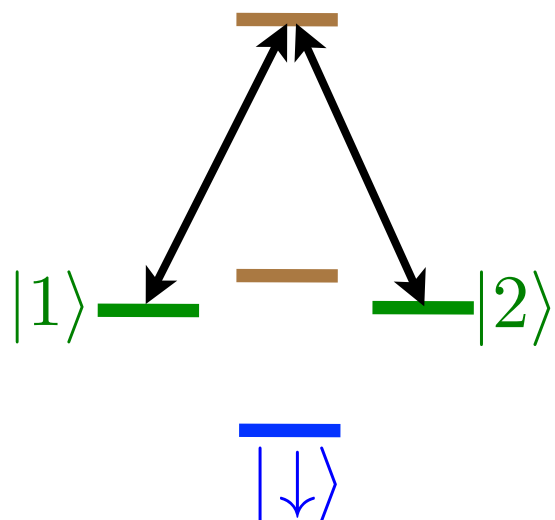


$$|\uparrow\rangle_i = x_i|1\rangle + y_i|2\rangle + z_i|3\rangle$$

[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

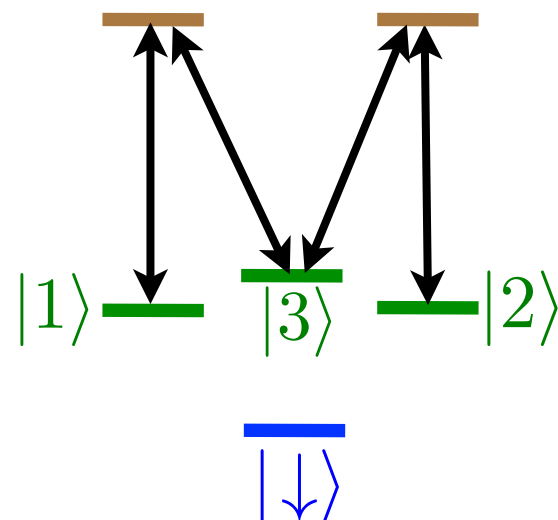
# Fractional Chern insulator with polar molecules

Recall SPT discussion:



$$|\uparrow\rangle = x|1\rangle + y|2\rangle$$

To get fractional Chern insulators:



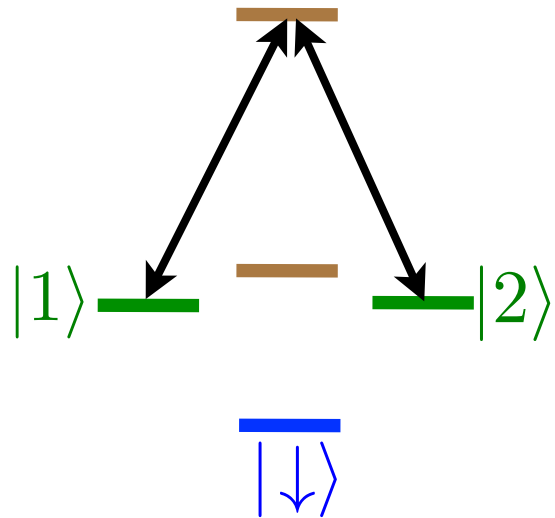
$$|\uparrow\rangle_i = x_i|1\rangle + y_i|2\rangle + z_i|3\rangle$$

Two differences:

[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

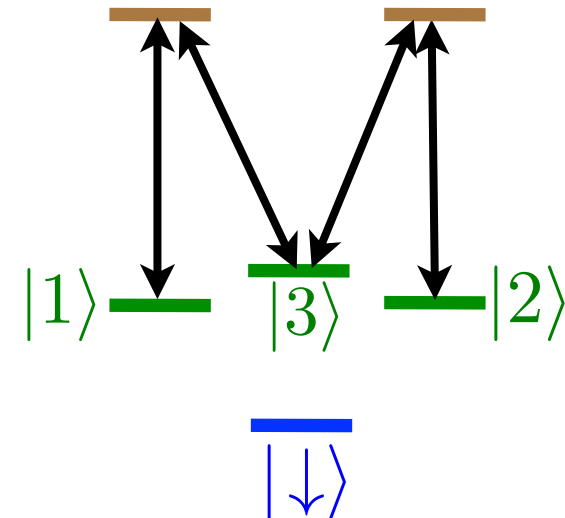
# Fractional Chern insulator with polar molecules

Recall SPT discussion:



$$|\uparrow\rangle = x|1\rangle + y|2\rangle$$

To get fractional Chern insulators:



$$|\uparrow\rangle_i = x_i|1\rangle + y_i|2\rangle + z_i|3\rangle$$

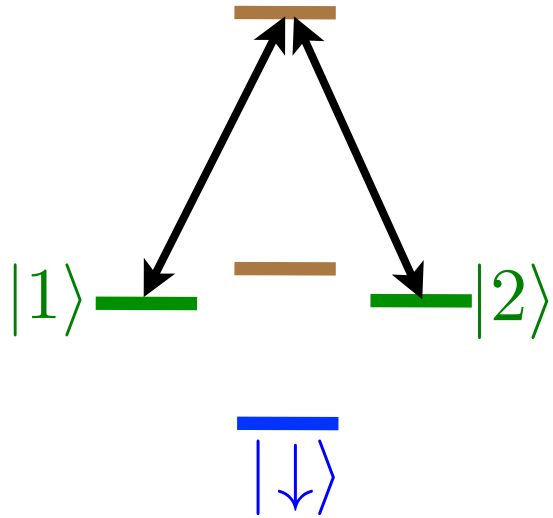
Two differences:

- use  $|3\rangle$

[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

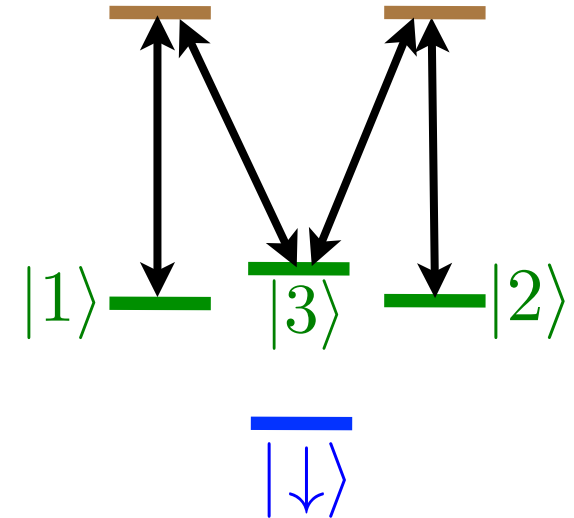
# Fractional Chern insulator with polar molecules

Recall SPT discussion:



$$|\uparrow\rangle = x|1\rangle + y|2\rangle$$

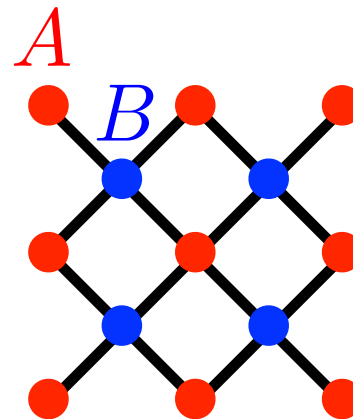
To get fractional Chern insulators:



$$|\uparrow\rangle_i = x_i|1\rangle + y_i|2\rangle + z_i|3\rangle$$

Two differences:

- use  $|3\rangle$

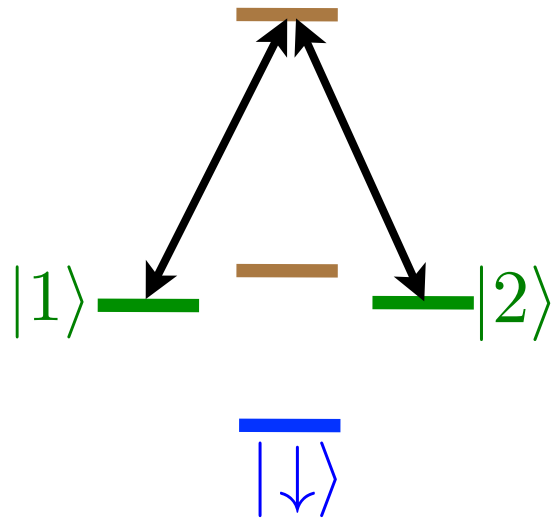


[PRL 109, 266804 (2012)  
& arXiv:1212.4839]



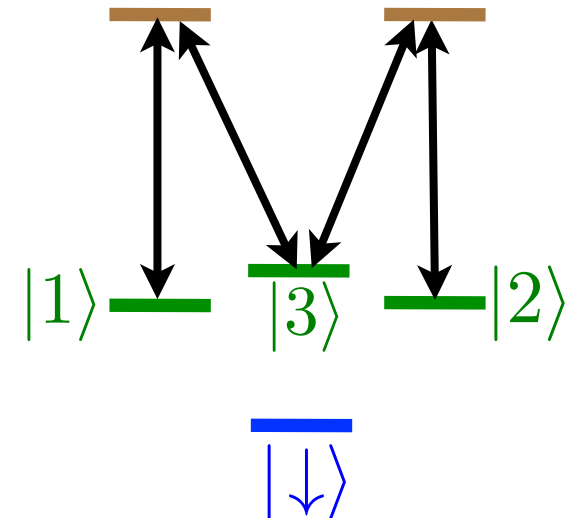
# Fractional Chern insulator with polar molecules

Recall SPT discussion:



$$|\uparrow\rangle = x|1\rangle + y|2\rangle$$

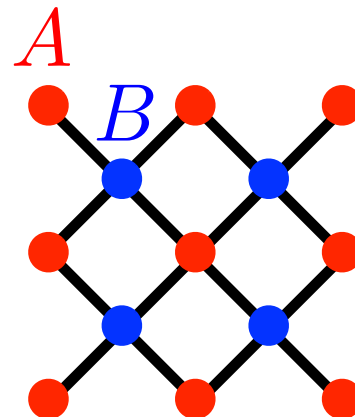
To get fractional Chern insulators:



$$|\uparrow\rangle_i = x_i|1\rangle + y_i|2\rangle + z_i|3\rangle$$

Two differences:

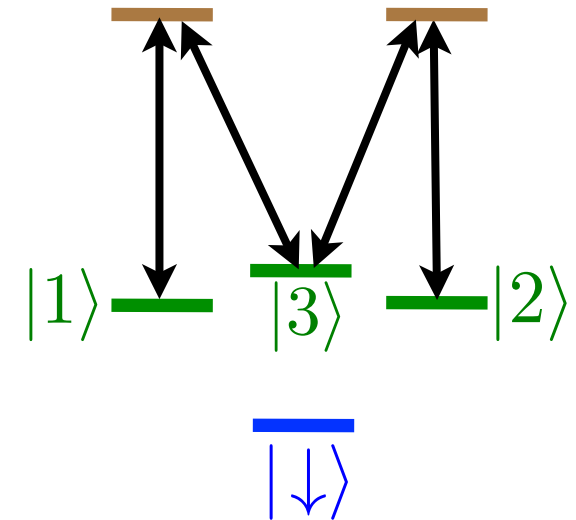
- use  $|3\rangle$
- $|\uparrow\rangle_i$  depends on  $i = A, B$



[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

# Fractional Chern insulator with polar molecules

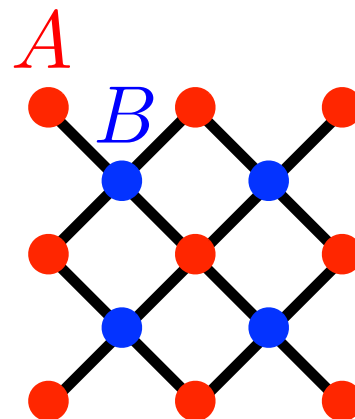
To get fractional Chern insulators:



$$|\uparrow\rangle_i = x_i|1\rangle + y_i|2\rangle + z_i|3\rangle$$

Two differences:

- use  $|3\rangle$
- $|\uparrow\rangle_i$  depends on  $i = A, B$

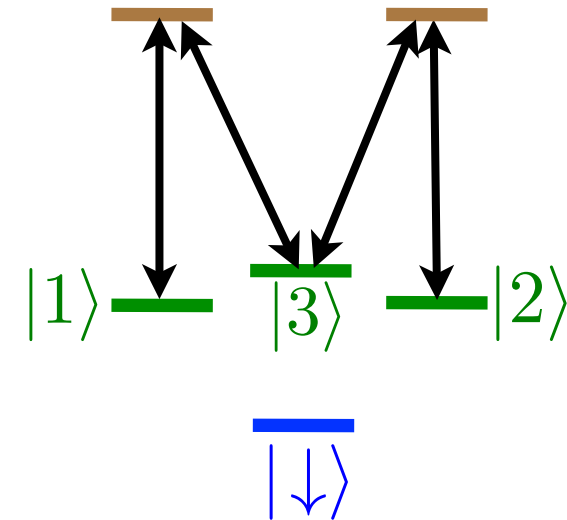


[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

# Fractional Chern insulator with polar molecules

- weak DC electric field

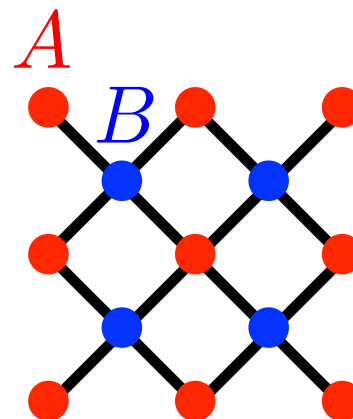
To get fractional Chern insulators:



$$|\uparrow\rangle_i = x_i|1\rangle + y_i|2\rangle + z_i|3\rangle$$

Two differences:

- use  $|3\rangle$
- $|\uparrow\rangle_i$  depends on  $i = A, B$



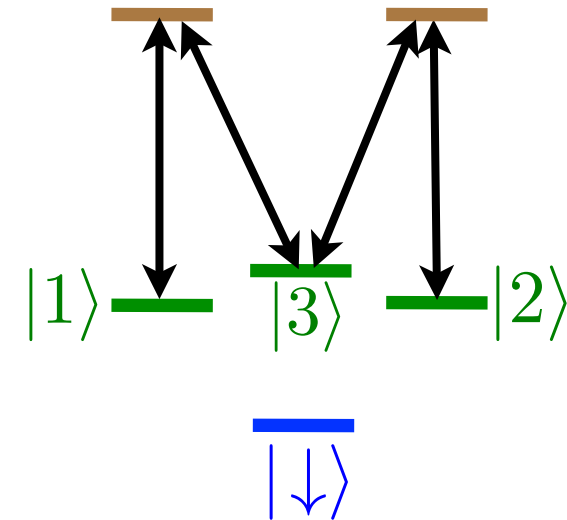
[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

# Fractional Chern insulator with polar molecules

- weak DC electric field

To get fractional Chern insulators:

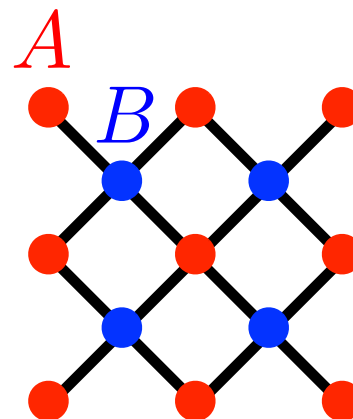
$$H = - \sum_{ij} t_{ij} S_i^+ S_j^-$$



$$|\uparrow\rangle_i = x_i |1\rangle + y_i |2\rangle + z_i |3\rangle$$

Two differences:

- use  $|3\rangle$
- $|\uparrow\rangle_i$  depends on  $i = A, B$



[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

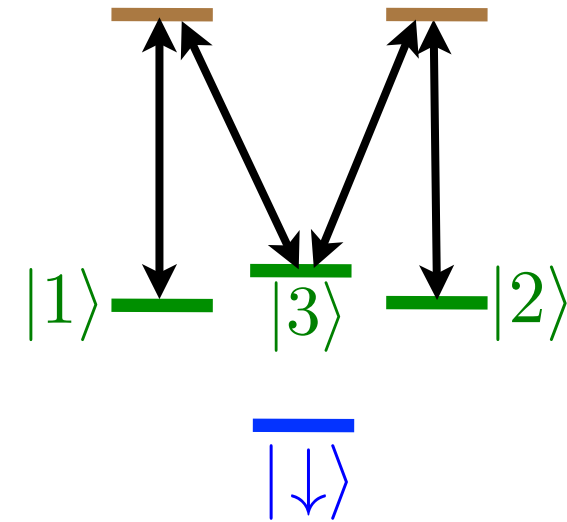
# Fractional Chern insulator with polar molecules

- weak DC electric field

To get fractional Chern insulators:

$$H = - \sum_{ij} t_{ij} S_i^+ S_j^-$$

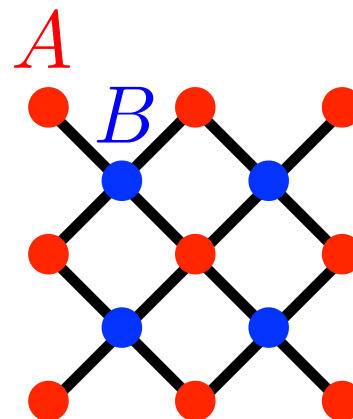
spin-1/2 = hardcore boson



$$|\uparrow\rangle_i = x_i |1\rangle + y_i |2\rangle + z_i |3\rangle$$

Two differences:

- use  $|3\rangle$
- $|\uparrow\rangle_i$  depends on  $i = A, B$



[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

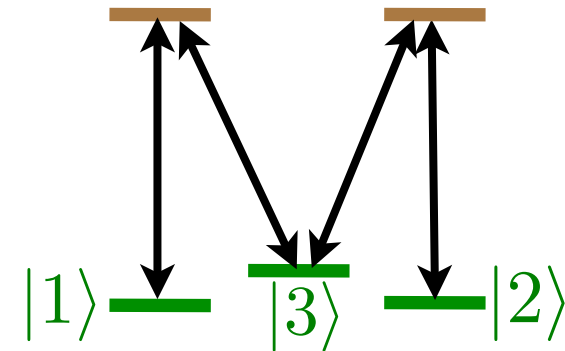
# Fractional Chern insulator with polar molecules

- weak DC electric field

To get fractional Chern insulators:

$$H = - \sum_{ij} t_{ij} S_i^+ S_j^-$$

spin-1/2 = hardcore boson

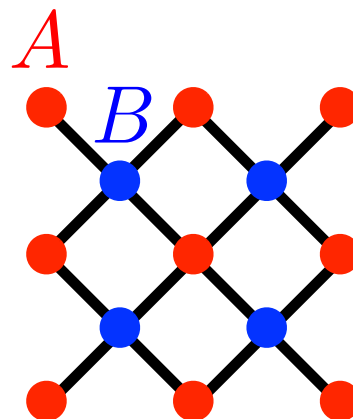


$$|0\rangle = |\downarrow\rangle$$

$$|\uparrow\rangle_i = x_i |1\rangle + y_i |2\rangle + z_i |3\rangle$$

Two differences:

- use  $|3\rangle$
- $|\uparrow\rangle_i$  depends on  $i = A, B$



[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

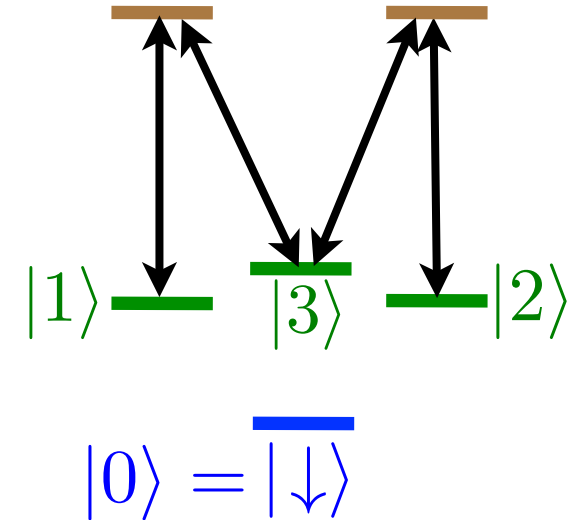
# Fractional Chern insulator with polar molecules

- weak DC electric field

To get fractional Chern insulators:

$$H = - \sum_{ij} t_{ij} S_i^+ S_j^-$$

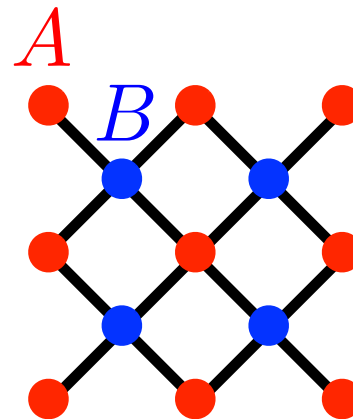
spin-1/2 = hardcore boson



$$a_i^\dagger |0\rangle = |\uparrow\rangle_i = x_i |1\rangle + y_i |2\rangle + z_i |3\rangle$$

Two differences:

- use  $|3\rangle$
- $|\uparrow\rangle_i$  depends on  $i = A, B$



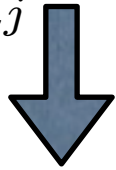
[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

# Fractional Chern insulator with polar molecules

- weak DC electric field

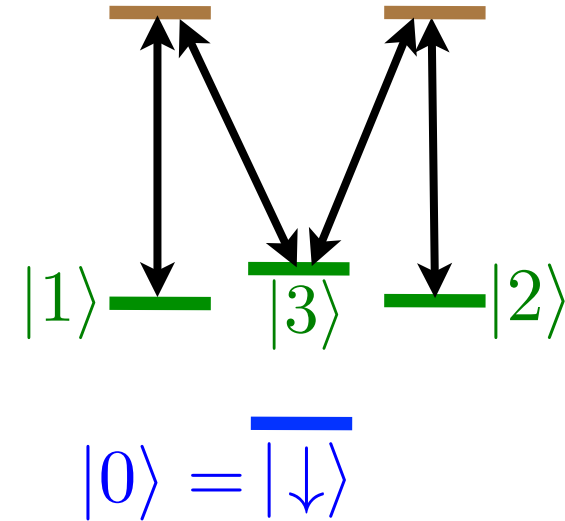
To get fractional Chern insulators:

$$H = - \sum_{ij} t_{ij} S_i^+ S_j^-$$



spin-1/2 = hardcore boson

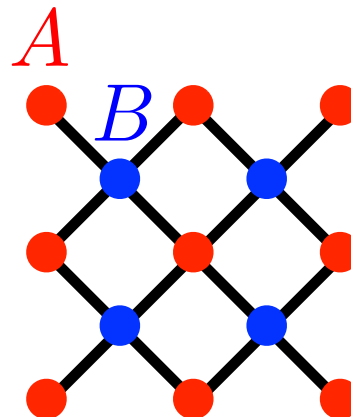
$$H = - \sum_{ij} t_{ij} a_i^\dagger a_j + (\text{hardcore constraint})$$



$$a_i^\dagger |0\rangle = |\uparrow\rangle_i = x_i |1\rangle + y_i |2\rangle + z_i |3\rangle$$

Two differences:

- use  $|3\rangle$
- $|\uparrow\rangle_i$  depends on  $i = A, B$



[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

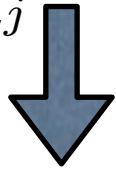


# Fractional Chern insulator with polar molecules

- weak DC electric field

To get fractional Chern insulators:

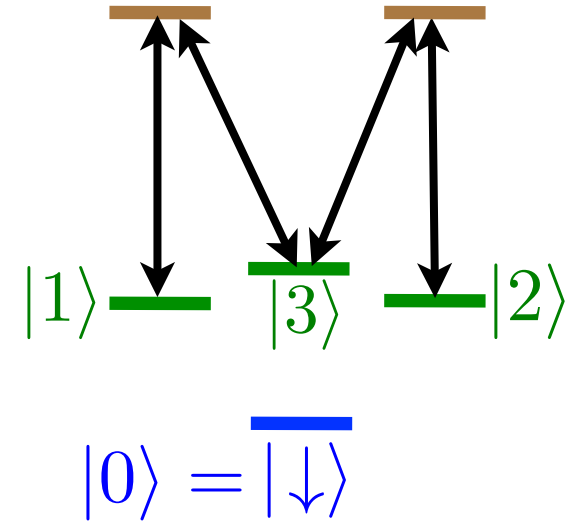
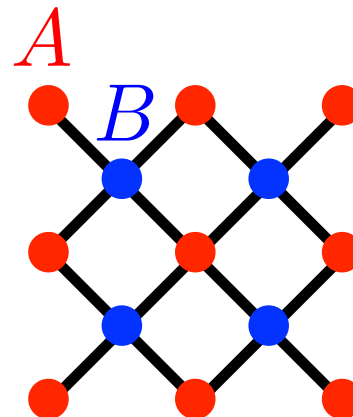
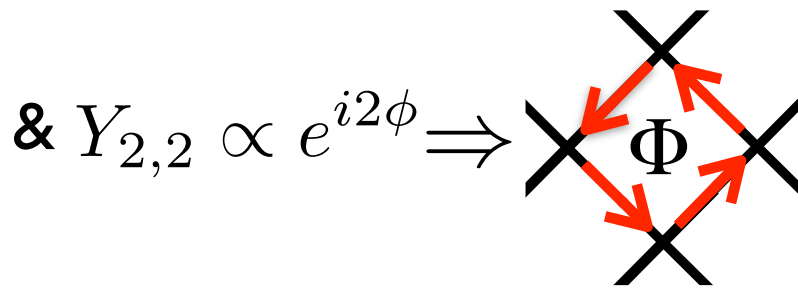
$$H = - \sum_{ij} t_{ij} S_i^+ S_j^-$$



spin-1/2 = hardcore boson

$$H = - \sum_{ij} t_{ij} a_i^\dagger a_j + (\text{hardcore constraint})$$

- sublattice-dependent dressing



$$a_i^\dagger |0\rangle = |\uparrow\rangle_i = x_i |1\rangle + y_i |2\rangle + z_i |3\rangle$$

Two differences:

- use  $|3\rangle$
- $|\uparrow\rangle_i$  depends on  $i = A, B$

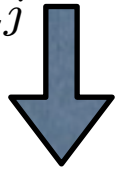
[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

# Fractional Chern insulator with polar molecules

- weak DC electric field

To get fractional Chern insulators:

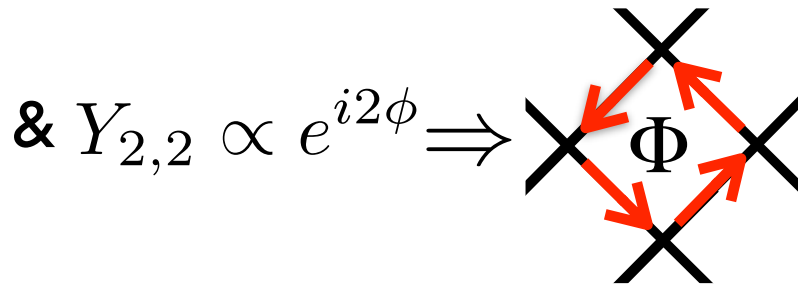
$$H = - \sum_{ij} t_{ij} S_i^+ S_j^-$$



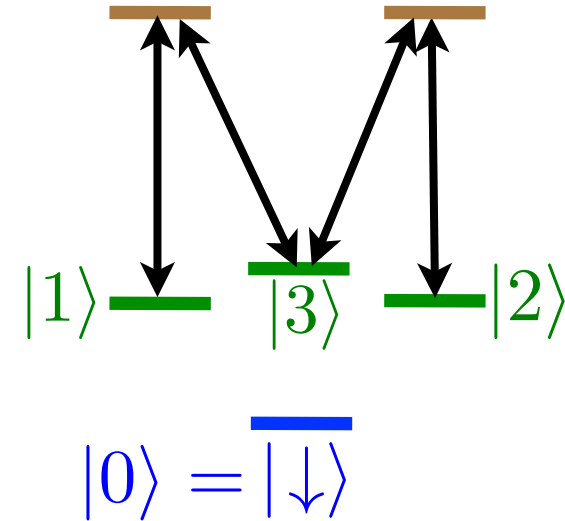
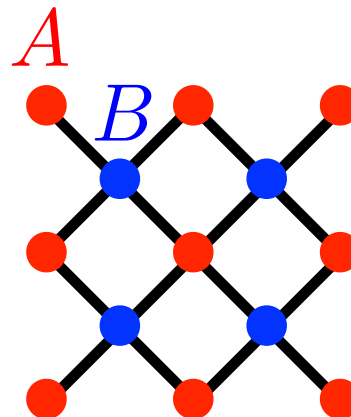
spin-1/2 = hardcore boson

$$H = - \sum_{ij} t_{ij} a_i^\dagger a_j + (\text{hardcore constraint})$$

- sublattice-dependent dressing



$\Rightarrow$  nonzero Chern number



$$a_i^\dagger |0\rangle = |\uparrow\rangle_i = x_i |1\rangle + y_i |2\rangle + z_i |3\rangle$$

Two differences:

- use  $|3\rangle$
- $|\uparrow\rangle_i$  depends on  $i = A, B$

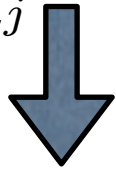
[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

# Fractional Chern insulator with polar molecules

- weak DC electric field

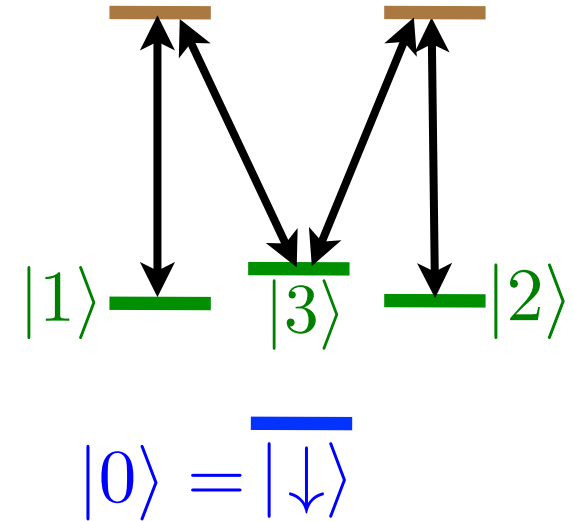
To get fractional Chern insulators:

$$H = - \sum_{ij} t_{ij} S_i^+ S_j^-$$

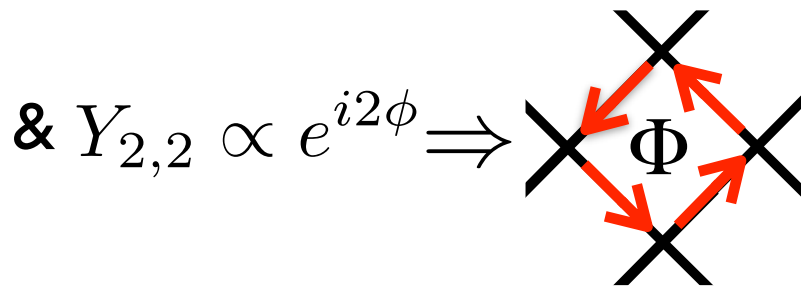


spin-1/2 = hardcore boson

$$H = - \sum_{ij} t_{ij} a_i^\dagger a_j + (\text{hardcore constraint})$$

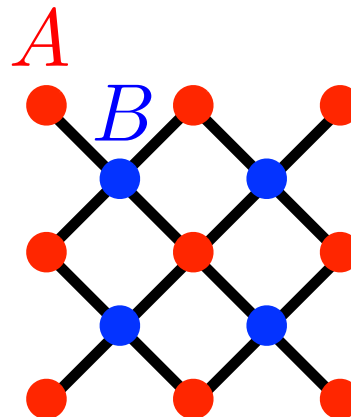


- sublattice-dependent dressing



$\Rightarrow$  nonzero Chern number

- anisotropic long-range hops interfere  $\Rightarrow$  flat bands



$$a_i^\dagger |0\rangle = |\uparrow\rangle_i = x_i |1\rangle + y_i |2\rangle + z_i |3\rangle$$

Two differences:

- use  $|3\rangle$
- $|\uparrow\rangle_i$  depends on  $i = A, B$

[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

# Fractional Chern insulator with polar molecules

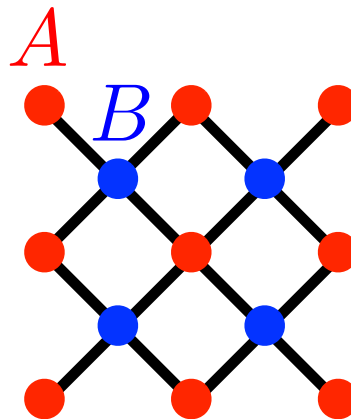
$$H = - \sum_{ij} t_{ij} a_i^\dagger a_j + \left( \begin{array}{l} \text{hardcore} \\ \text{constraint} \end{array} \right)$$

- sublattice-dependent dressing

$$\& Y_{2,2} \propto e^{i2\phi} \Rightarrow \begin{array}{c} \times \\ \nearrow \searrow \\ \Phi \\ \nwarrow \nearrow \\ \times \end{array}$$

$\Rightarrow$  nonzero Chern number

- anisotropic long-range hops interfere  $\Rightarrow$  flat bands

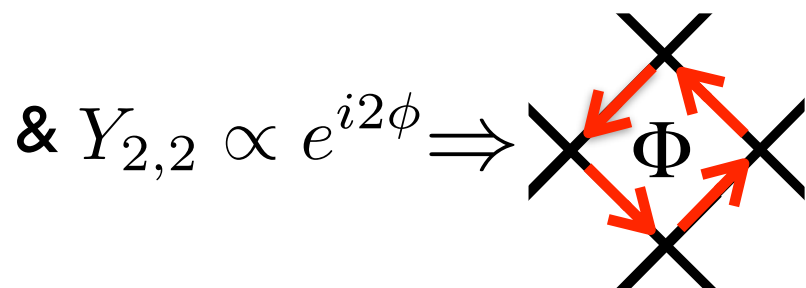


[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

# Fractional Chern insulator with polar molecules

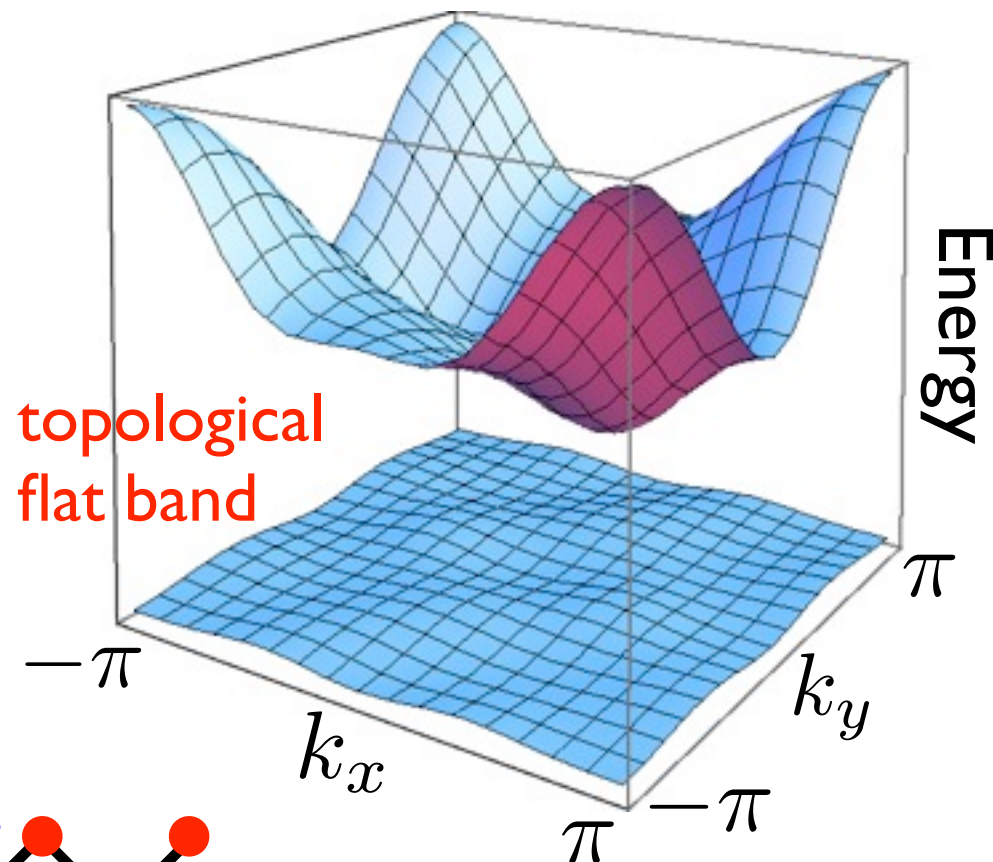
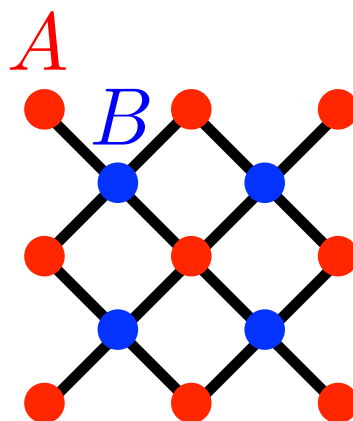
$$H = - \sum_{ij} t_{ij} a_i^\dagger a_j + (\text{hardcore constraint})$$

- sublattice-dependent dressing



$\Rightarrow$  nonzero Chern number

- anisotropic long-range hops interfere  $\Rightarrow$  flat bands

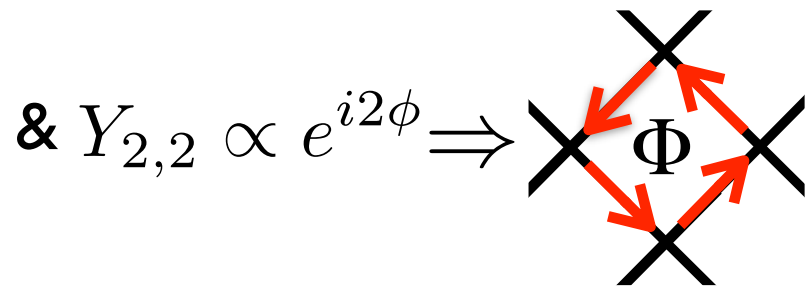


[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

# Fractional Chern insulator with polar molecules

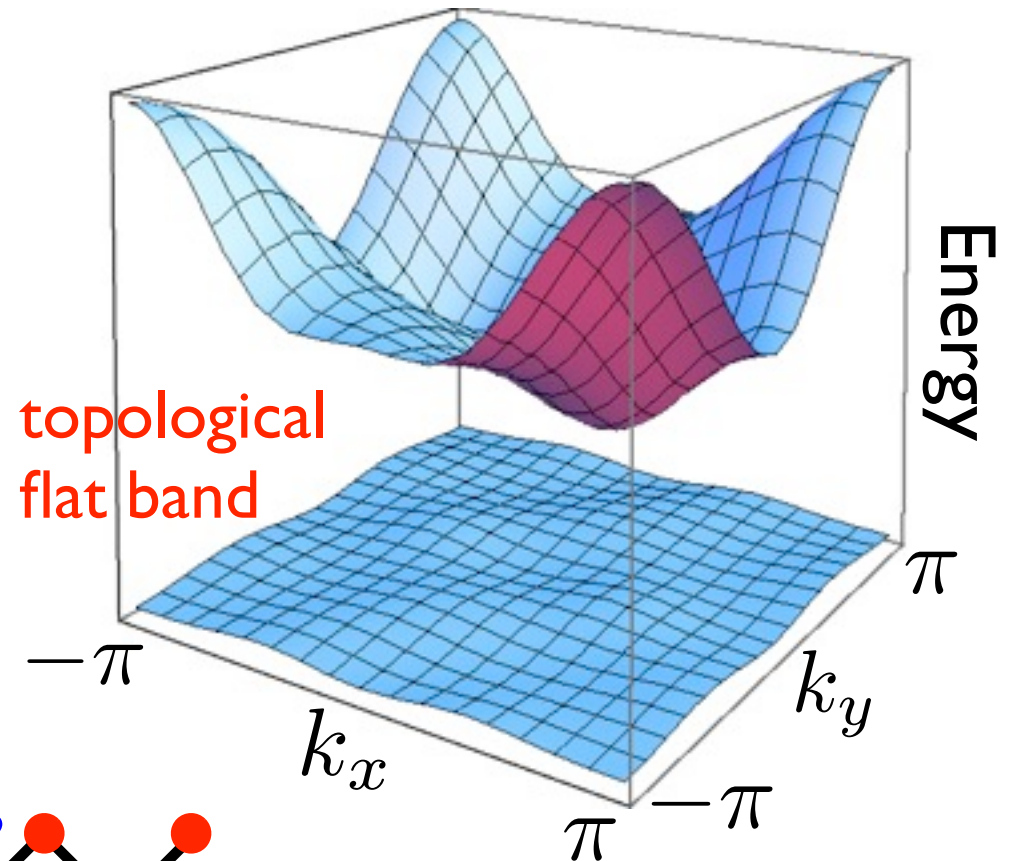
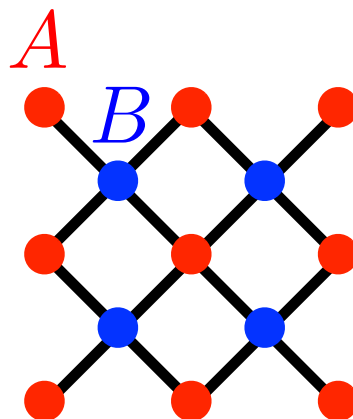
$$H = - \sum_{ij} t_{ij} a_i^\dagger a_j + \left( \text{hardcore constraint} \right)$$

- sublattice-dependent dressing



$\Rightarrow$  nonzero Chern number

- anisotropic long-range hops interfere  $\Rightarrow$  flat bands

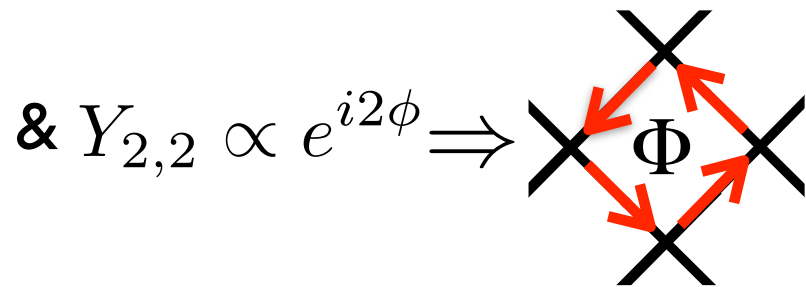


[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

# Fractional Chern insulator with polar molecules

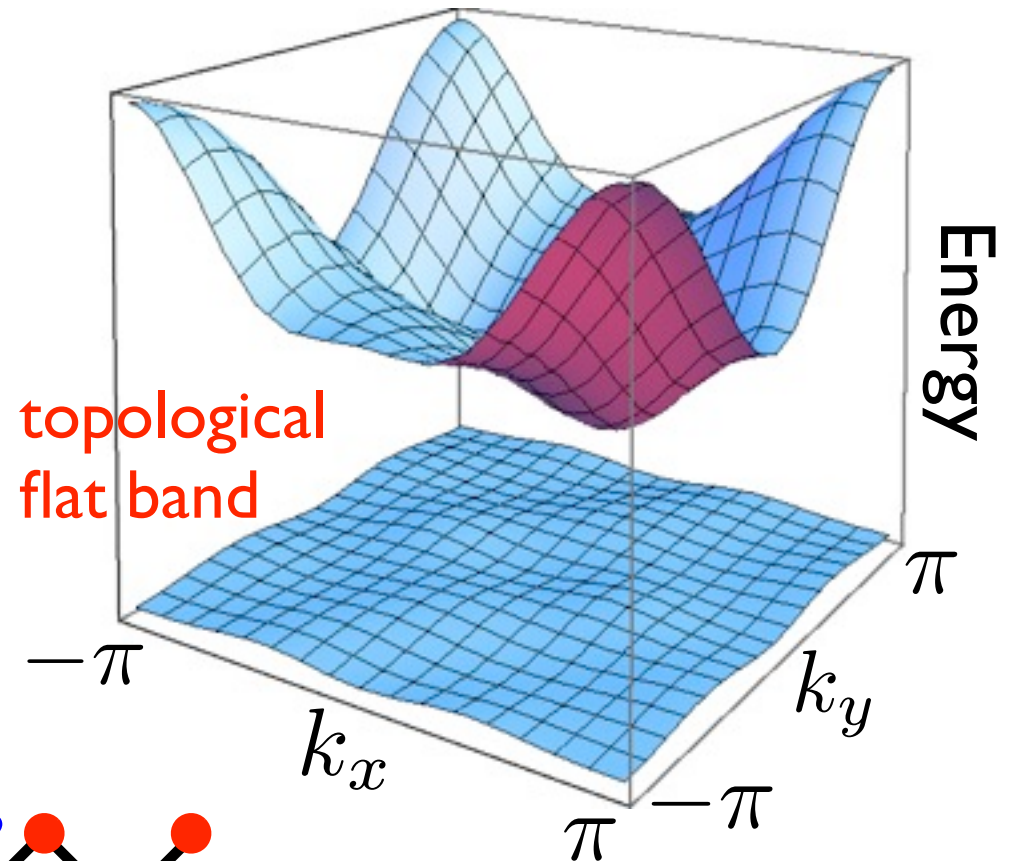
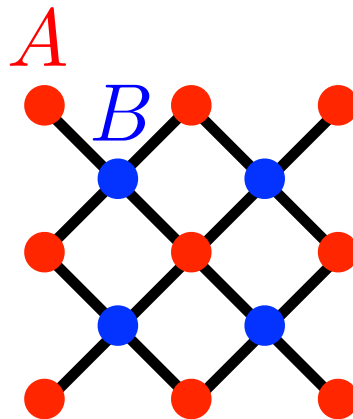
$$H = - \sum_{ij} t_{ij} a_i^\dagger a_j + (\text{hardcore constraint})$$

- sublattice-dependent dressing



$\Rightarrow$  nonzero Chern number

- anisotropic long-range hops interfere  $\Rightarrow$  flat bands



[PRL 109, 266804 (2012)  
& arXiv:1212.4839]

# Conclusion and Outlook



# Conclusion and Outlook

- first realistic implementation of fractional Chern insulators  
[see also Cooper & Dalibard 2012]

# Conclusion and Outlook

- first realistic implementation of fractional Chern insulators  
[see also Cooper & Dalibard 2012]
- $\nu = 1/2$  Laughlin state; Halperin (2,2,1) state

# Conclusion and Outlook

- first realistic implementation of fractional Chern insulators  
[see also Cooper & Dalibard 2012]
- $\nu = 1/2$  Laughlin state; Halperin (2,2,1) state
- phase transitions

# Conclusion and Outlook

- first realistic implementation of fractional Chern insulators  
[see also Cooper & Dalibard 2012]
- $\nu = 1/2$  Laughlin state; Halperin (2,2,1) state
- phase transitions
- preparation and detection

# Conclusion and Outlook

- first realistic implementation of fractional Chern insulators  
[see also Cooper & Dalibard 2012]
- $\nu = 1/2$  Laughlin state; Halperin (2,2,1) state
- phase transitions
- preparation and detection
- parafermions = fractionalized Majoranas  
[Clarke, Alicea, Shtengel 2012; Lindner, Berg, Refael, Stern, 2012]

# Conclusion and Outlook

- first realistic implementation of fractional Chern insulators  
[see also Cooper & Dalibard 2012]
- $\nu = 1/2$  Laughlin state; Halperin (2,2,1) state
- phase transitions
- preparation and detection
- parafermions = fractionalized Majoranas  
[Clarke, Alicea, Shtengel 2012; Lindner, Berg, Refael, Stern, 2012]
- exotic physics beyond fractional quantum Hall effect?

# Conclusion and Outlook

- first realistic implementation of fractional Chern insulators  
[see also Cooper & Dalibard 2012]
- $\nu = 1/2$  Laughlin state; Halperin (2,2,1) state
- phase transitions
- preparation and detection
- parafermions = fractionalized Majoranas  
[Clarke, Alicea, Shtengel 2012; Lindner, Berg, Refael, Stern, 2012]
- exotic physics beyond fractional quantum Hall effect?
- Rydberg atoms, magnetic atoms [Er, Dy, Cr], spins in solid state

# Conclusion and Outlook

- first realistic implementation of fractional Chern insulators  
[see also Cooper & Dalibard 2012]
- $\nu = 1/2$  Laughlin state; Halperin (2,2,1) state
- phase transitions
- preparation and detection
- parafermions = fractionalized Majoranas  
[Clarke, Alicea, Shtengel 2012; Lindner, Berg, Refael, Stern, 2012]
- exotic physics beyond fractional quantum Hall effect?
- Rydberg atoms, magnetic atoms [Er, Dy, Cr], spins in solid state

Thank You