

Many-Body Physics with Up and Coming Molecules in Optical Lattices

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Dynamics



What sorts of molecules?

$^1\Sigma$

- E.g. KRb, RbCs (talks this week)
- “Assembled” from ultracold alkali gases

$^2\Sigma$

- E.g. SrF, YO via laser cooling
- Assemble from other ultracold species

$^2\Pi$

- E.g. OH, evaporative cooling

Symmetric Tops

- E.g. CH₃F
- Possibility of efficient Stark deceleration, sympathetic cooling

<http://coldmoles13.wikispaces.com/>

What is new with molecules?

- Large electric dipole moments!
 - Resonant dipole-dipole interactions

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repulsive

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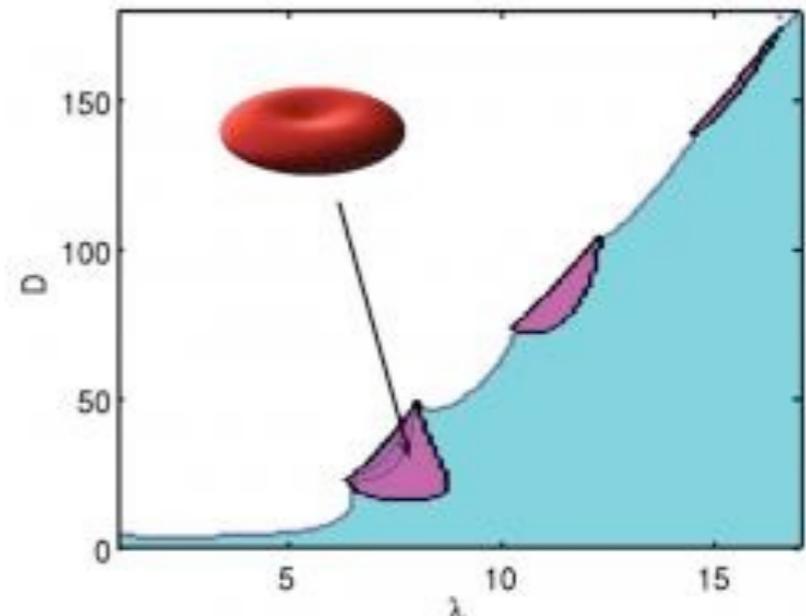
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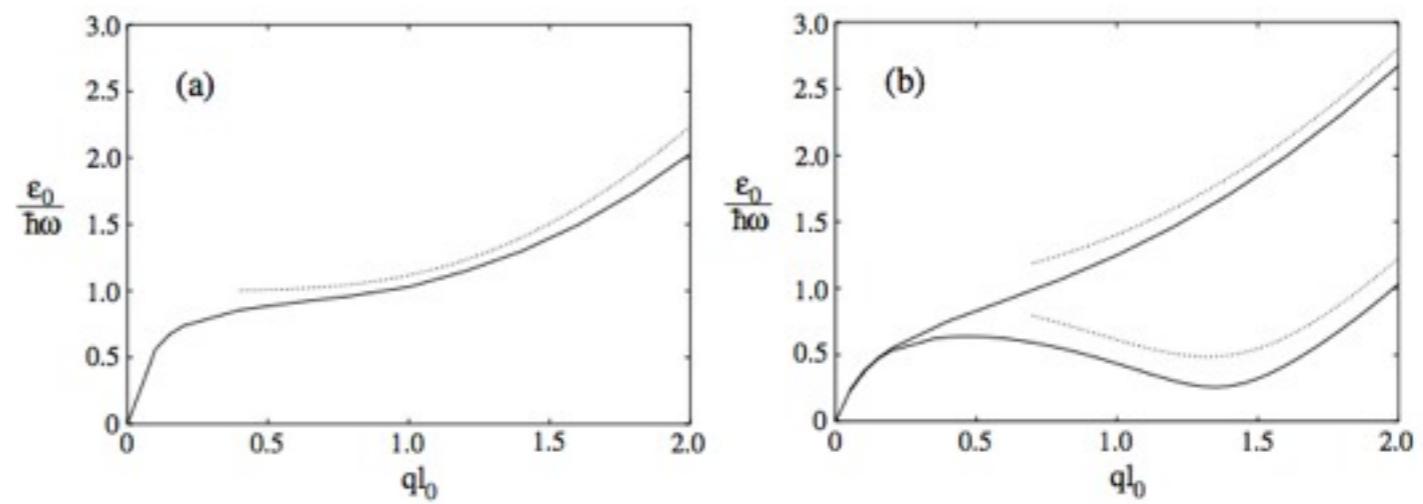
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Ronen et al. PRL 98 030406

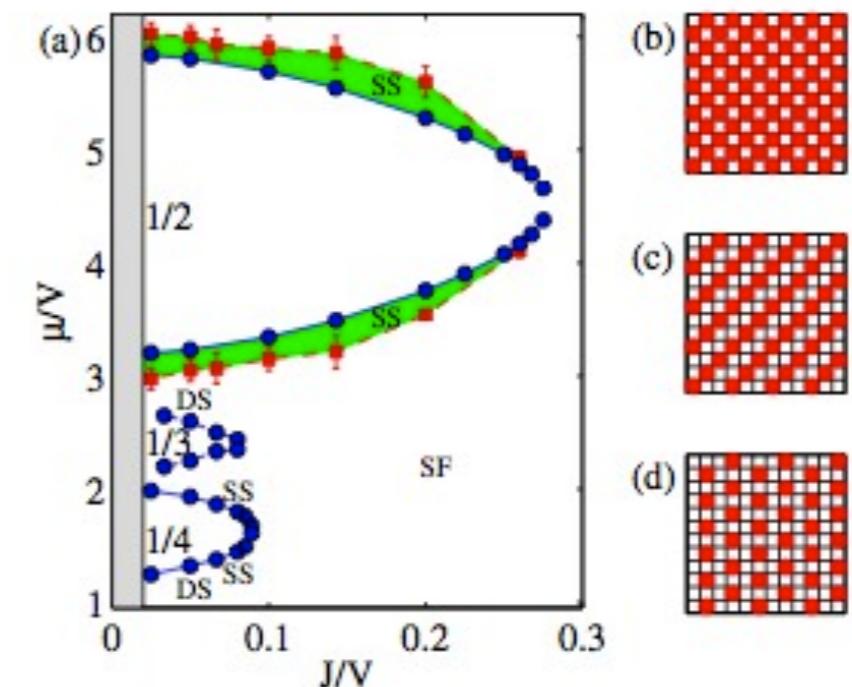
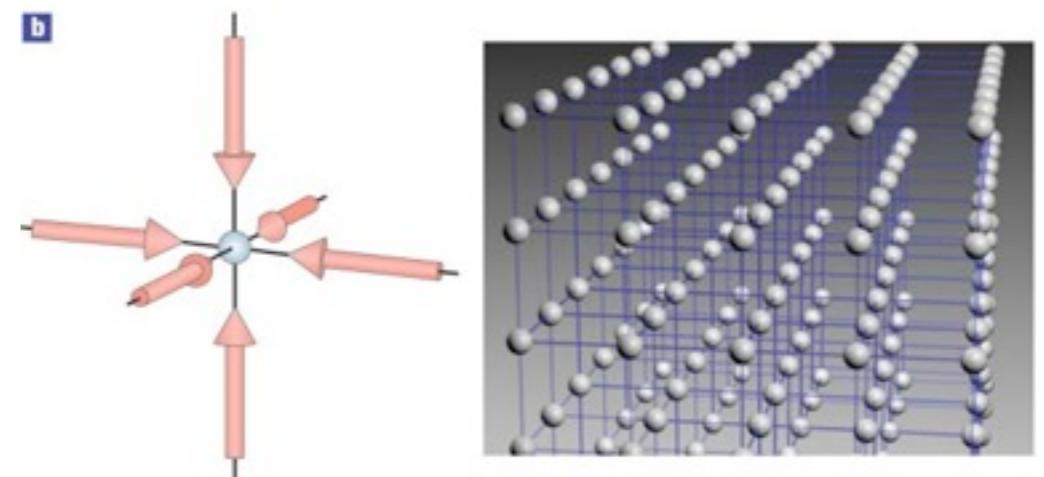
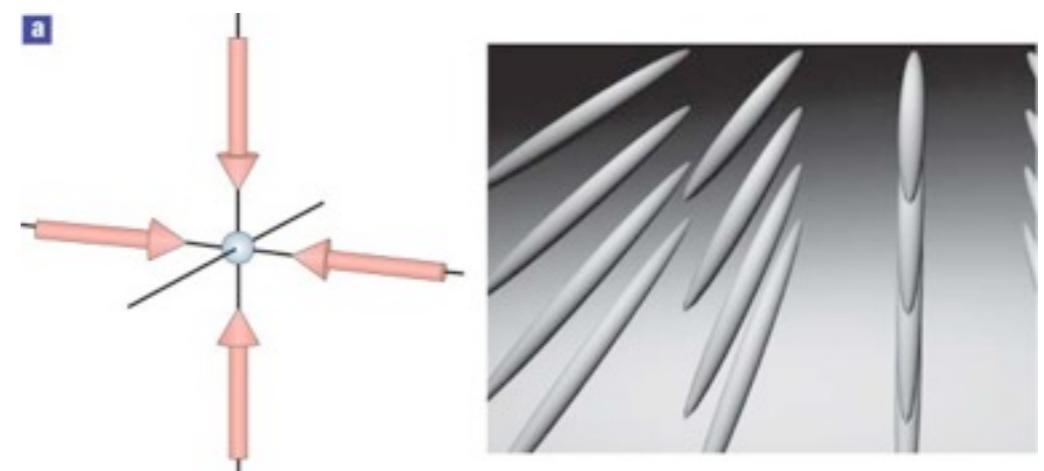
- Several novel predictions for structureless dipoles
- Rotons, “Red blood cell” BEC
- Anisotropic Fermi liquid
- Novel BCS-type pairing



Santos et al. PRL 90 250403

Dipoles in lattices

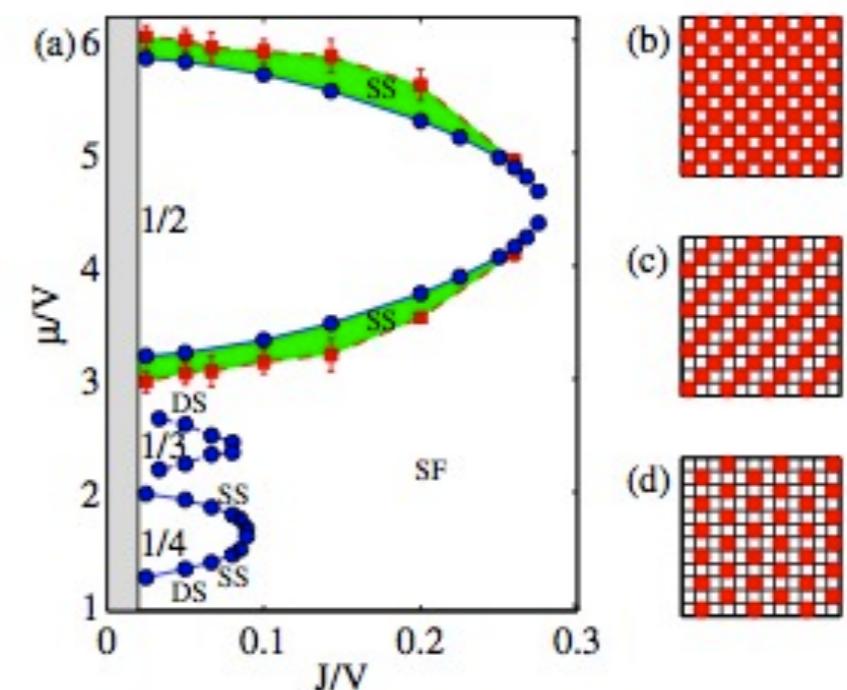
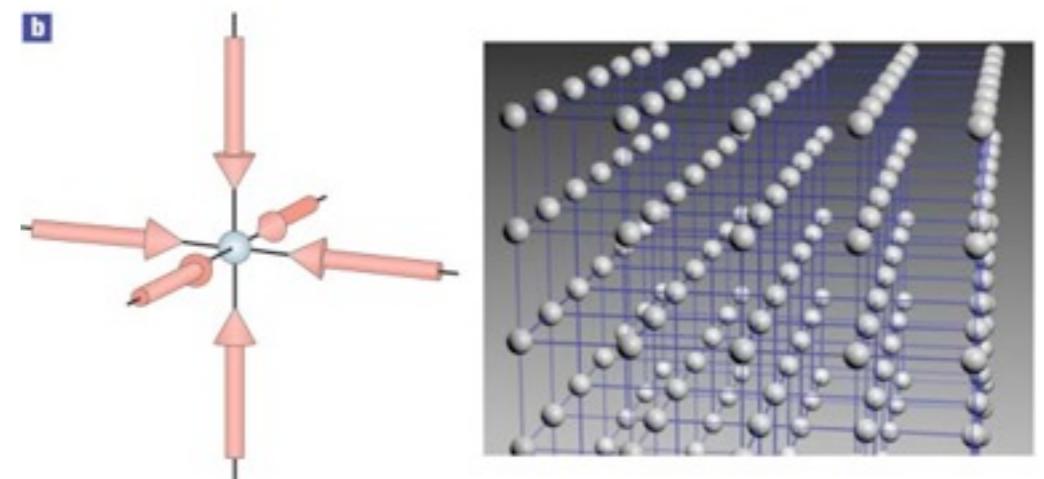
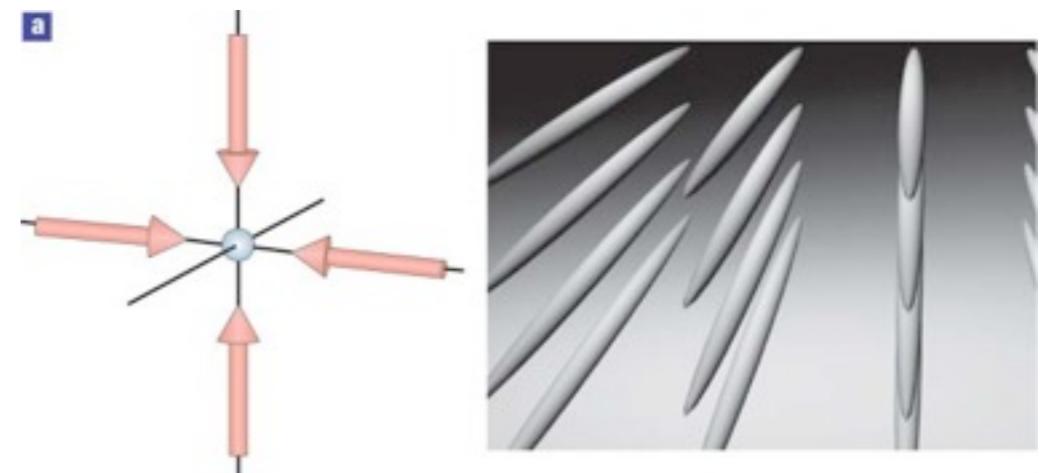
- Standing wave “Optical lattice”
- “Devil’s staircase” of insulating phases
- Supersolidity (Pupillo,...)



Capogrosso-Sansone et al. PRL 104 125301

Dipoles in lattices

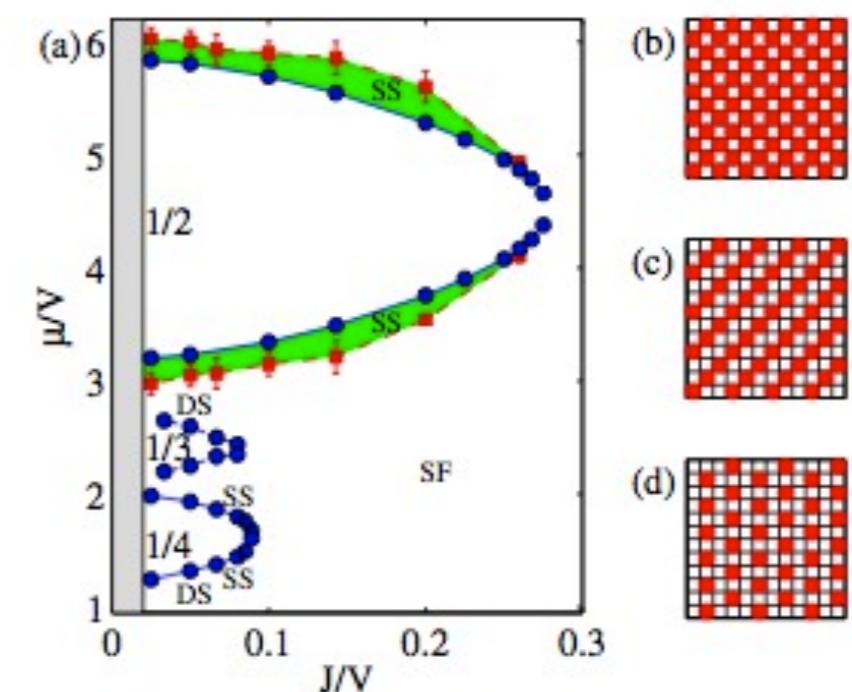
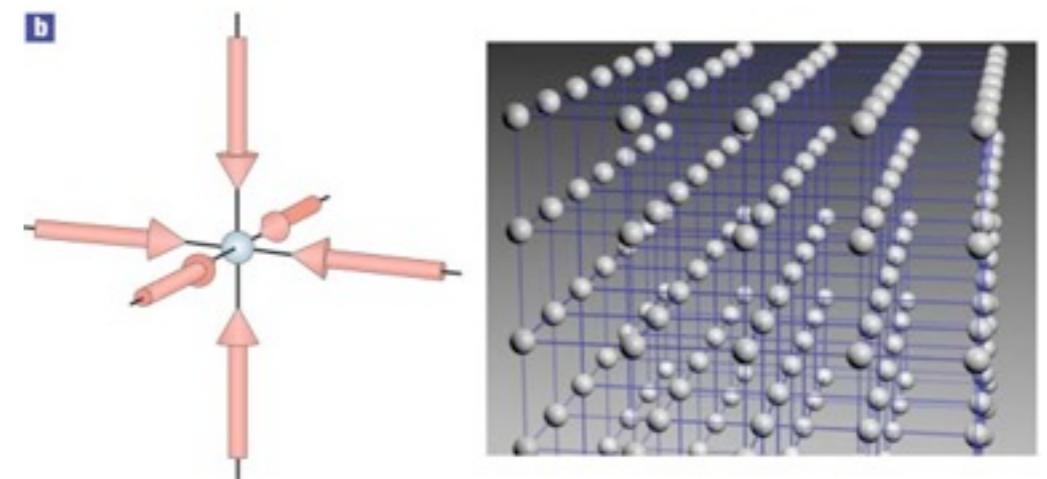
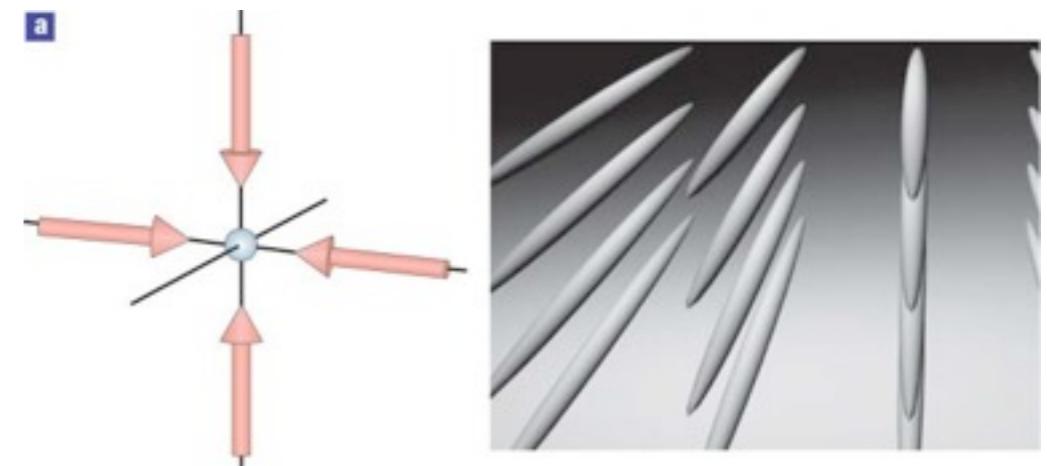
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- Are strong dipoles everything molecules have to offer for many-body physics?



Capogrosso-Sansone et al. PRL 104 125301

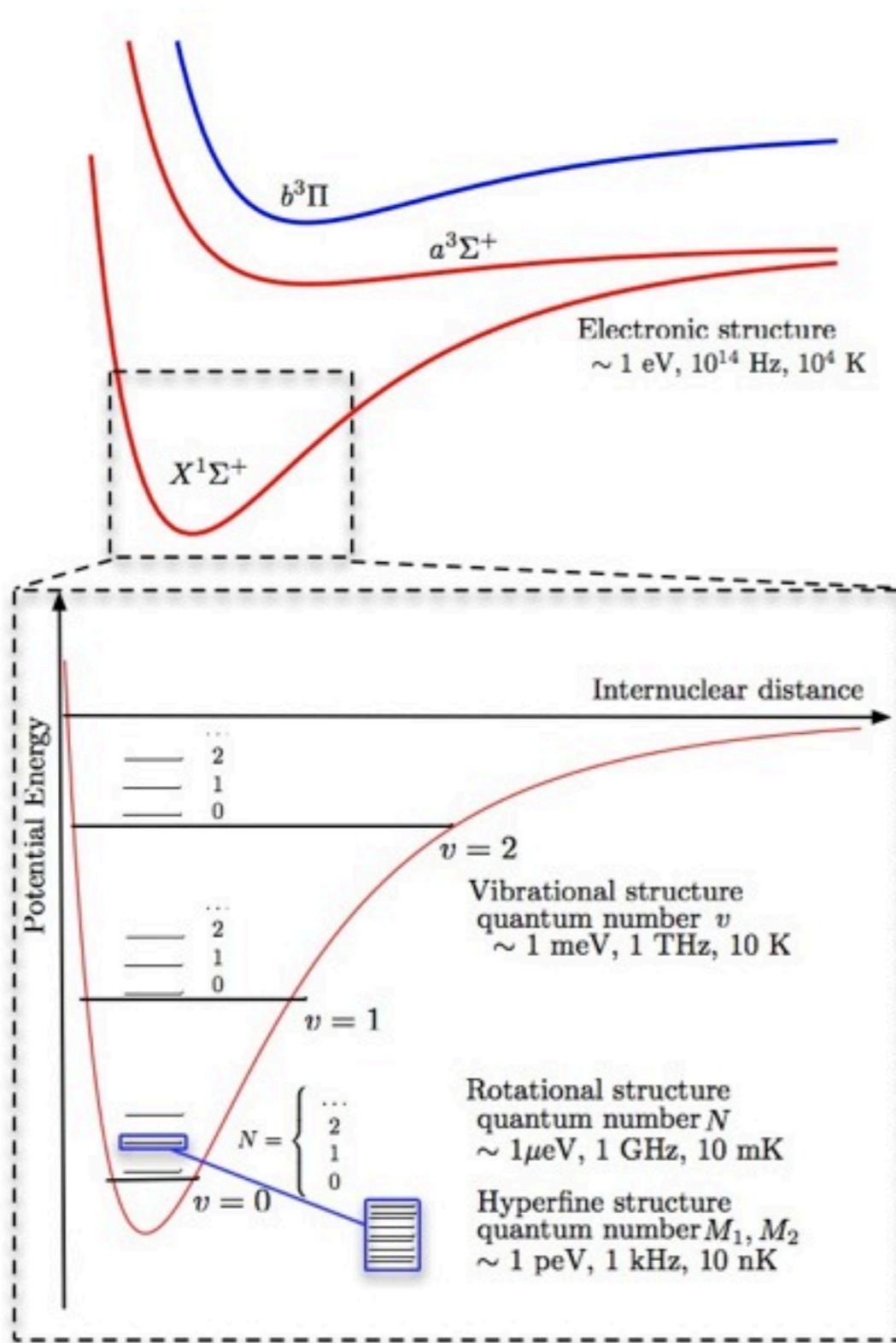
Dipoles in lattices

- Standing wave “Optical lattice”
- “Devil’s staircase” of insulating phases
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- Are strong dipoles everything molecules have to offer for many-body physics?
- Molecules are complex, tunable, multi-scale objects!
- Very different opportunities from solid state systems!



Capogrosso-Sansone et al. PRL 104 125301

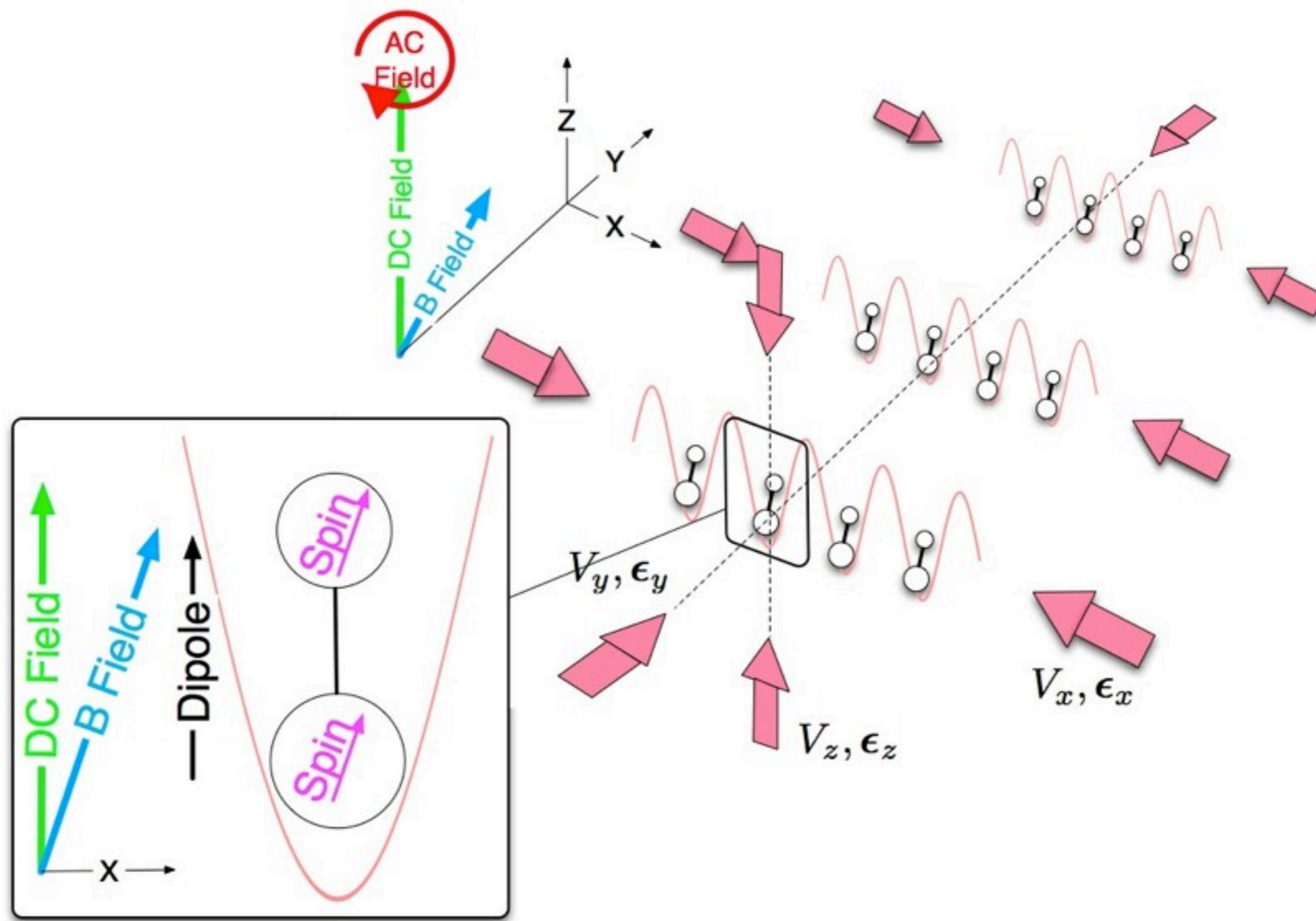
Characterization of ${}^1\Sigma$ molecules



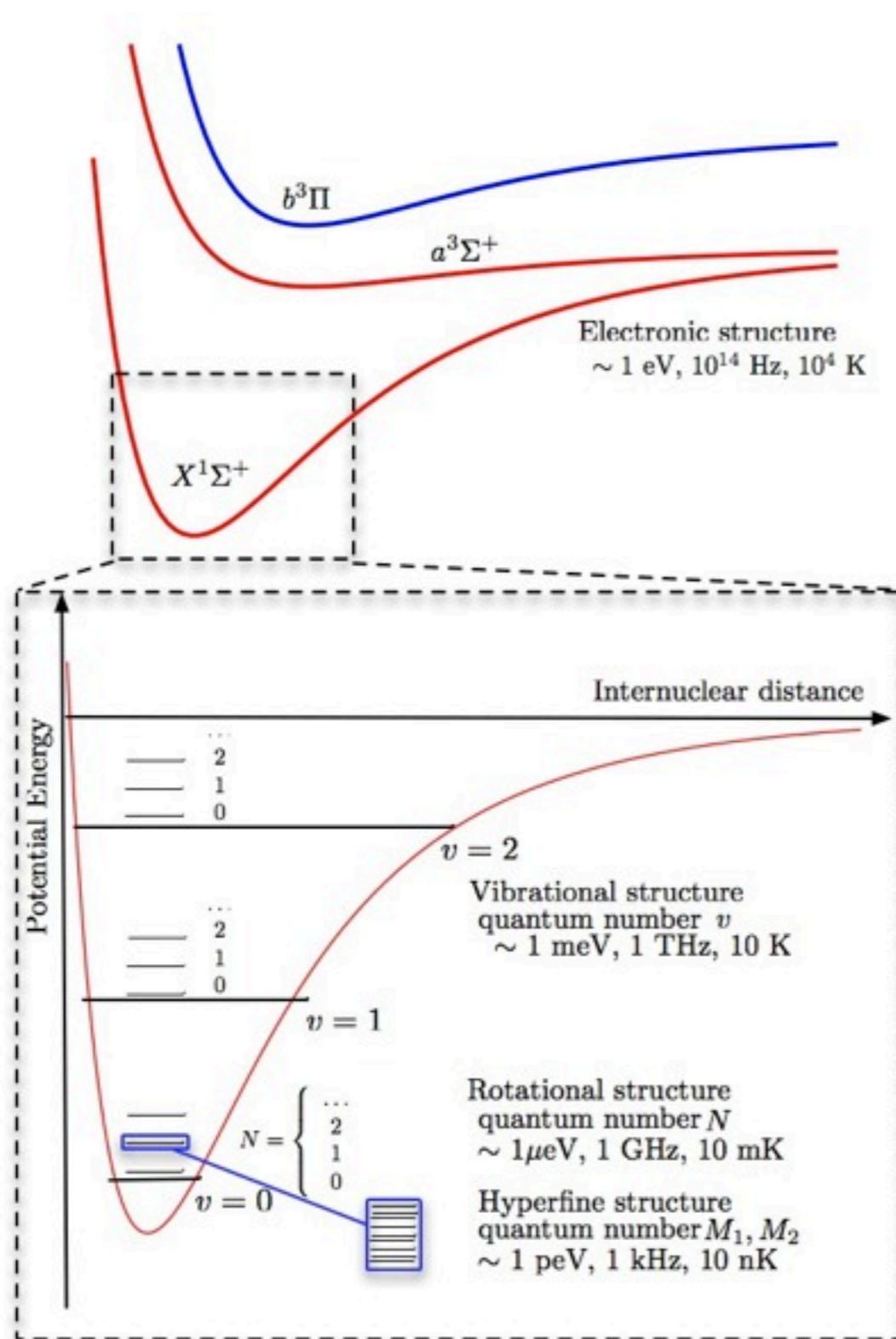
Multi-scale problem

- Electronic
 - Coupling to optical lattice
 - Anisotropic polarizability
- Rotational
 - Coupling to DC E-field
 - Dipole moments
 - Accessible with microwave field
- Hyperfine
 - Coupling to DC B-field
 - Quadrupole mixing with rotation
 - Large number of accessible internal states

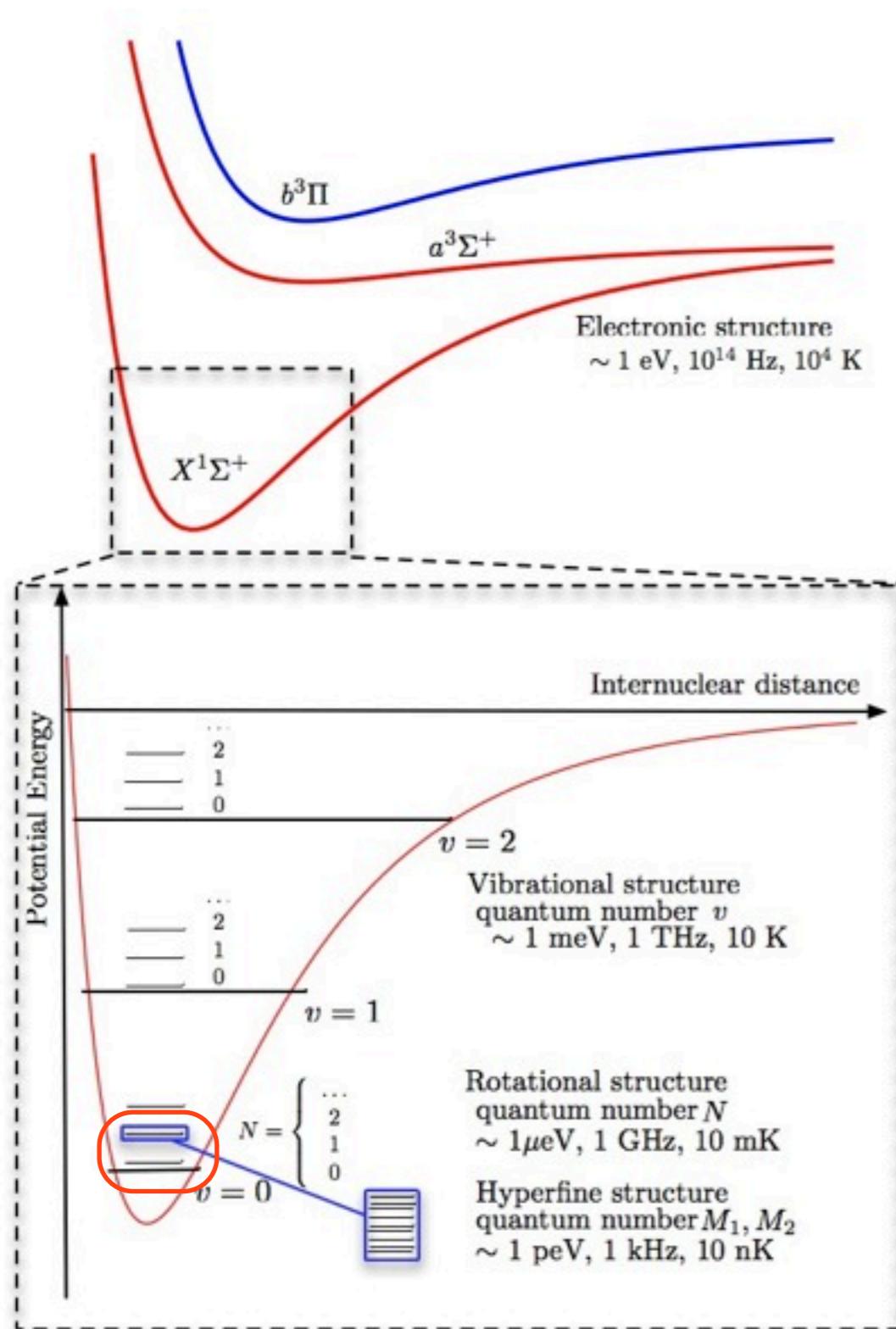
Experimental setup



Field Regimes

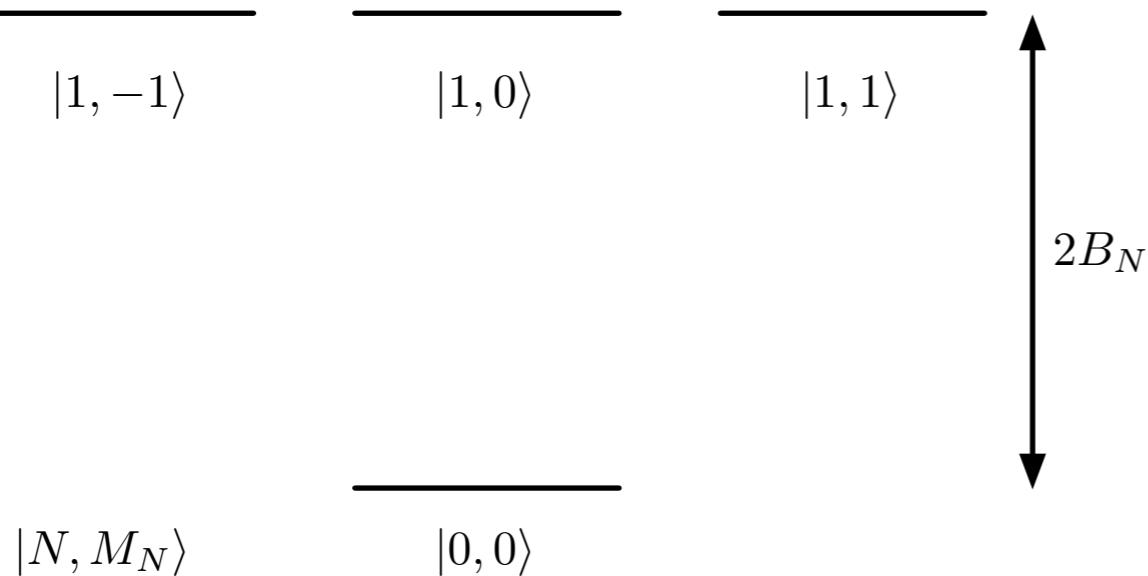


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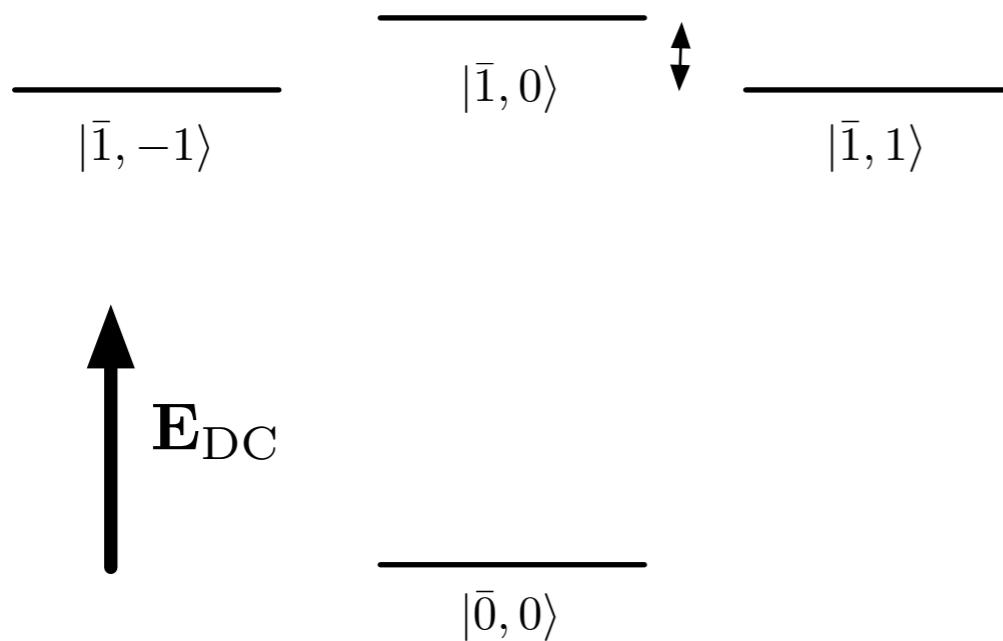
Field Regimes

- DC electric field
 - Splits $N=1$ degeneracy
 - Tunable dipole moments

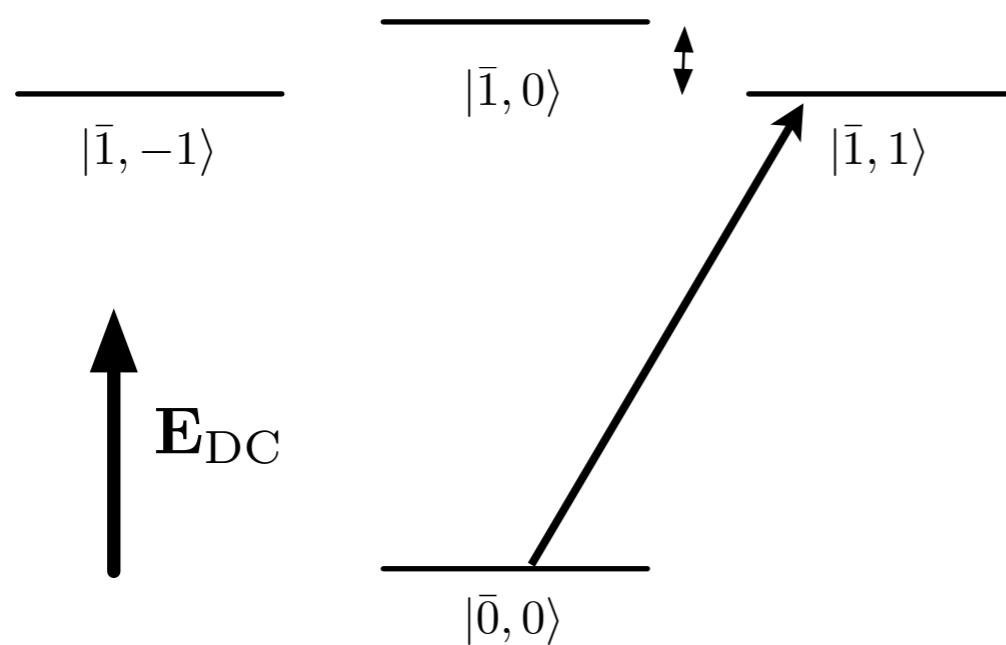


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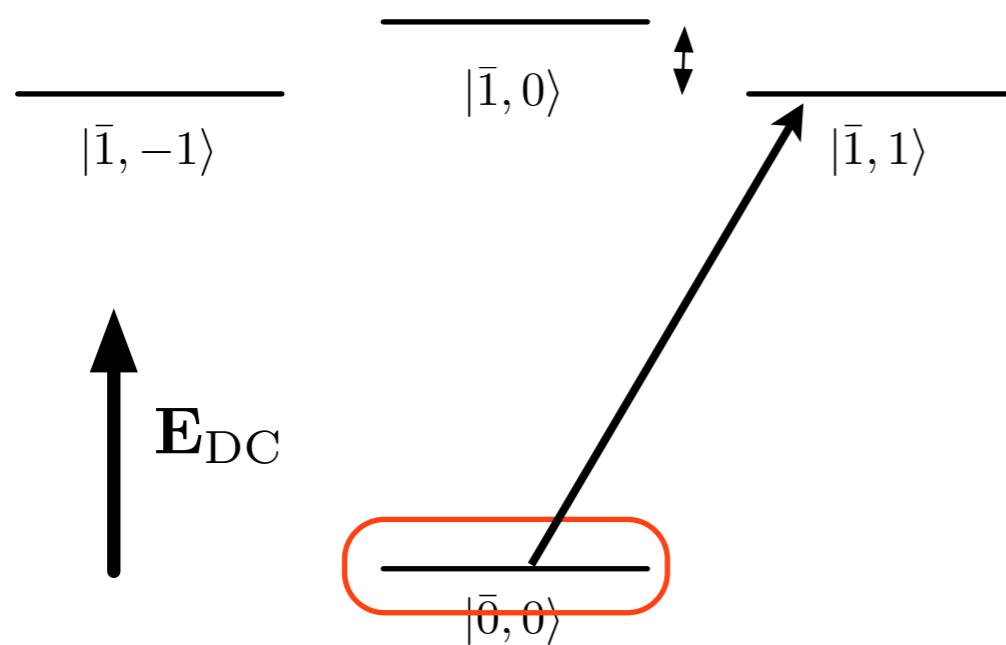


Field Regimes



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 - Tunable dipole moments
- AC microwave field
 - Tunable access to internal states

Field Regimes



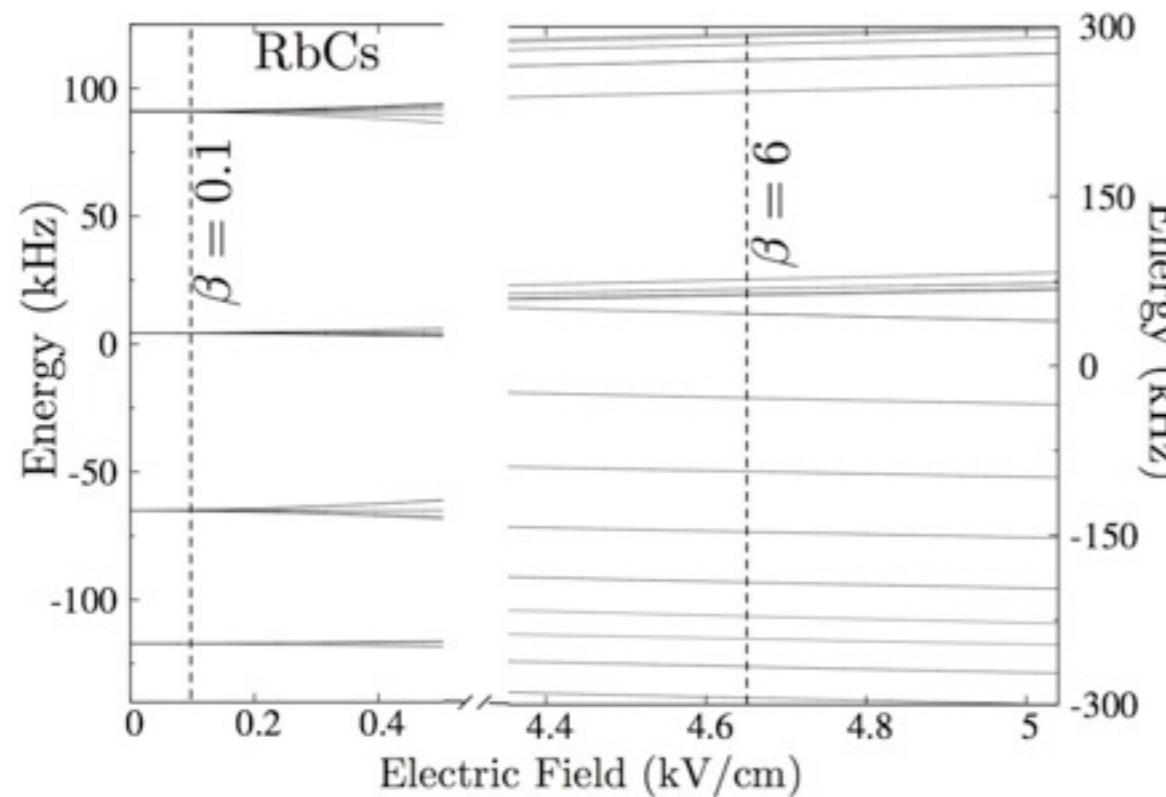
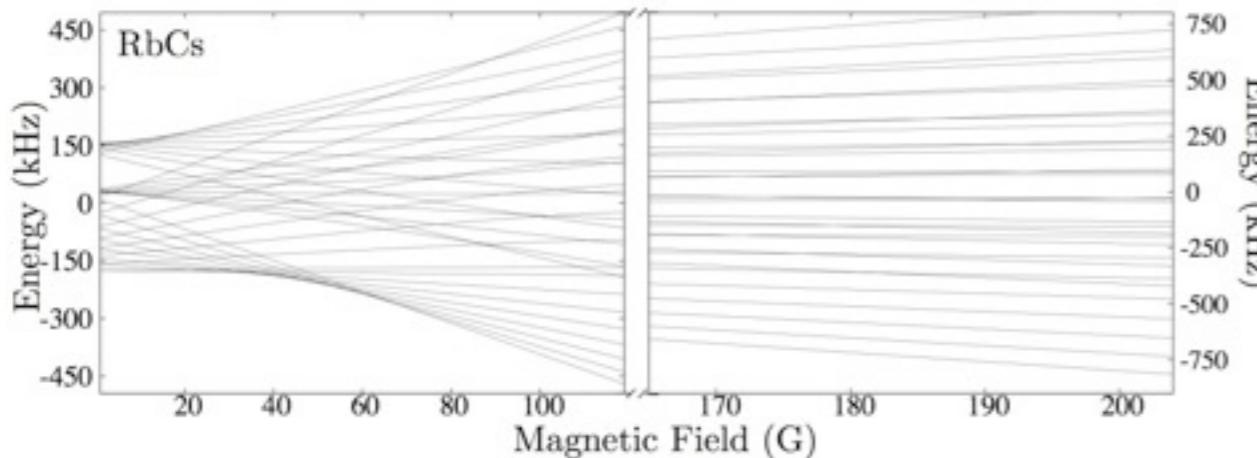
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- DC magnetic field
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 - Large fields present for magneto-association

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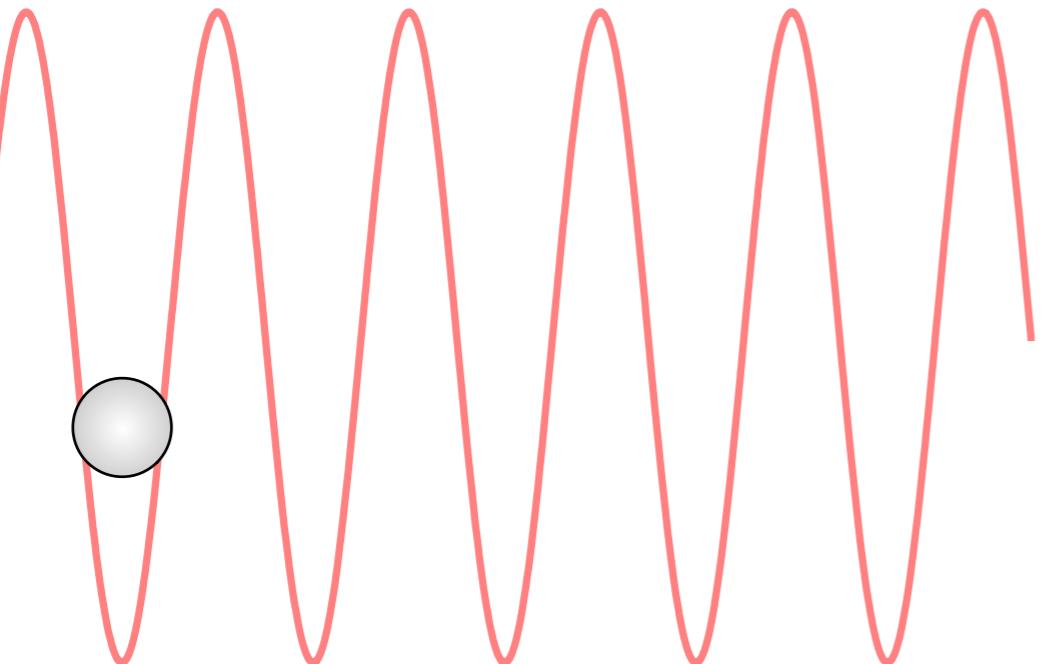


Atoms in optical lattices

$$\hat{H} = -t \sum_{\langle i,j \rangle} [\hat{b}_i^\dagger \hat{b}_j + \text{h.c.}] + \frac{U}{2} \sum_i \hat{n}_i (\hat{n}_i - 1)$$

Atoms in optical lattices

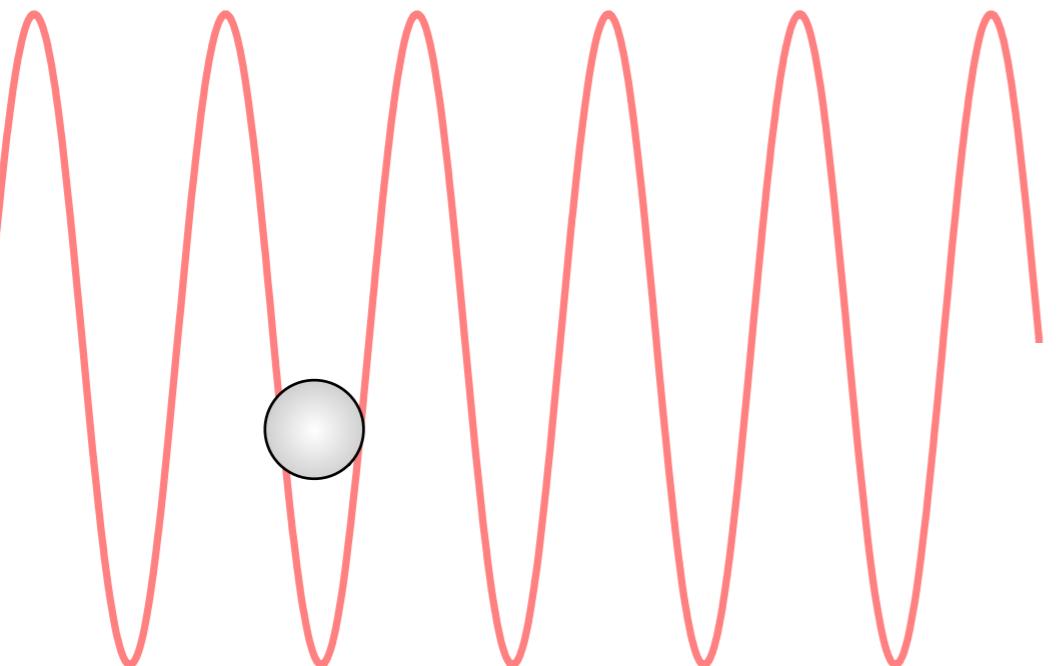
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Atoms in optical lattices

tunneling

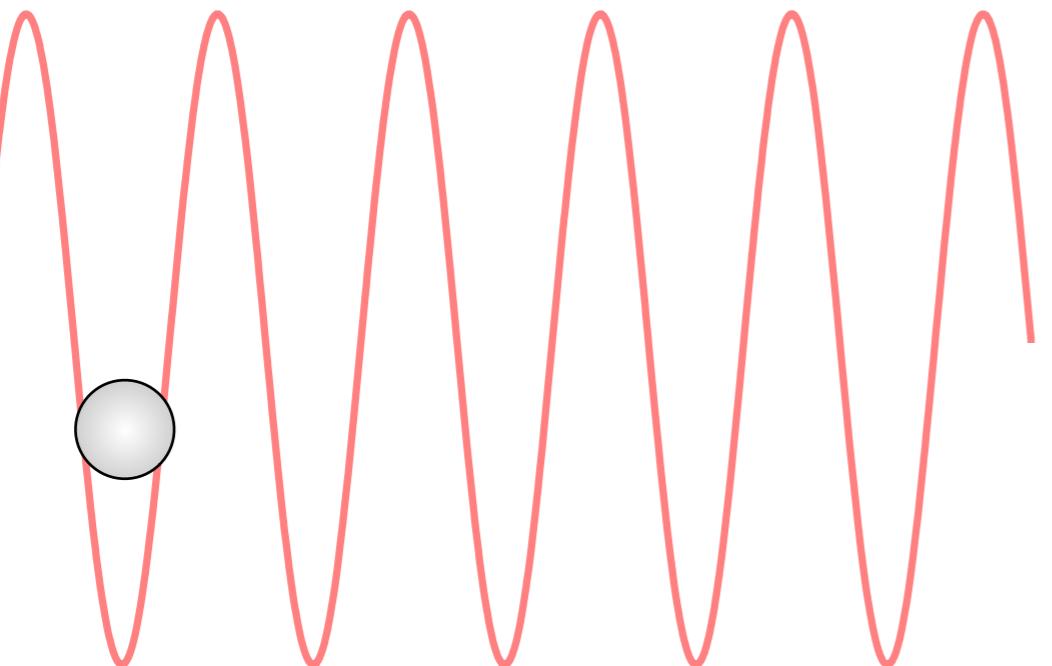
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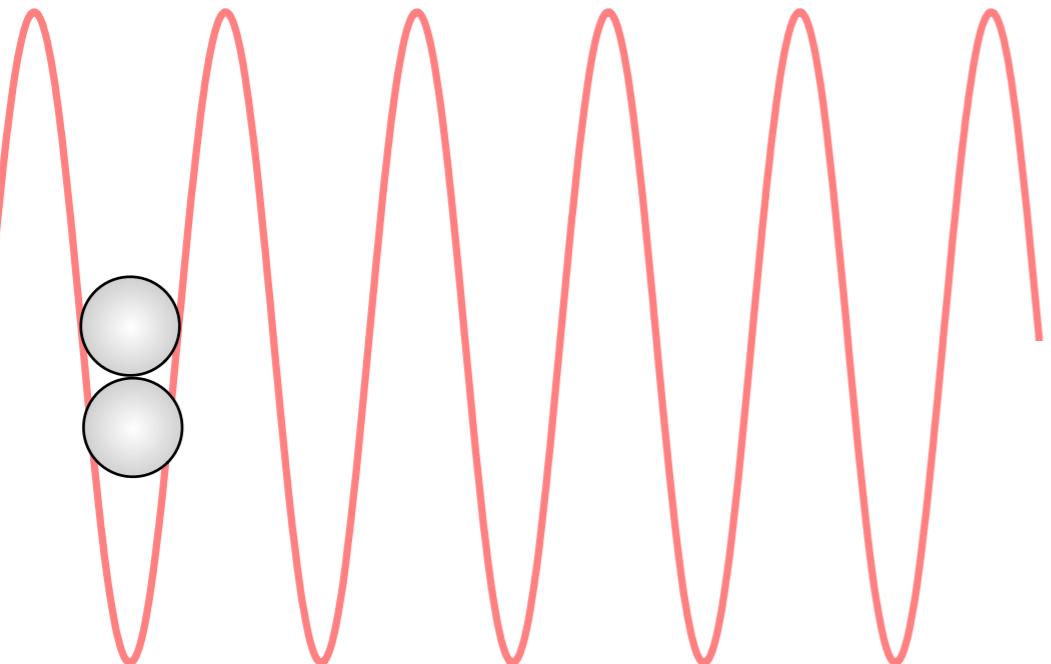


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s-wave scattering

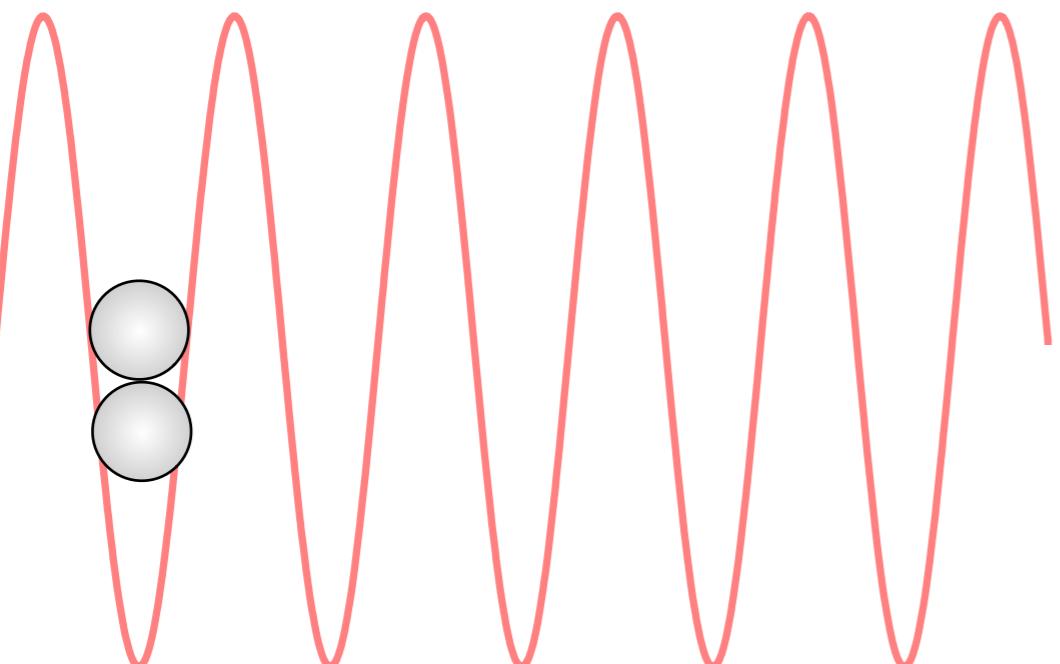


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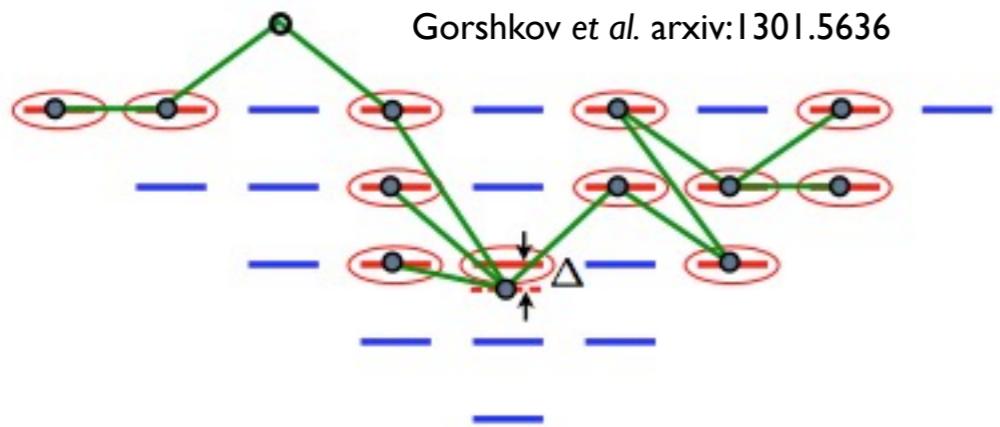
- “Quantum simulator”
- Also natural Hamiltonian for ultracold atoms
- What about molecules?

Molecules in optical lattices

- Two complementary questions:

Molecules in optical lattices

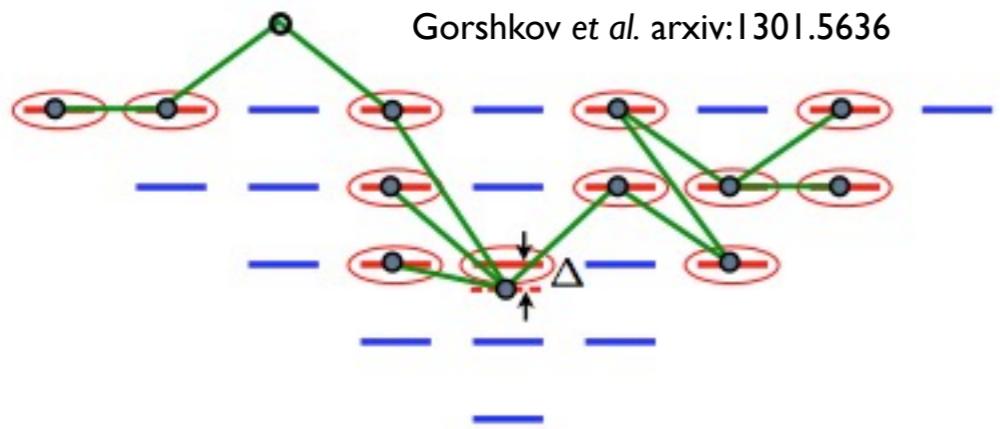
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- QI: What can we quantum simulate with ultracold molecules in optical lattices?
- AI: Quantum magnetism, topological phases, fractional Chern insulators,...

Molecules in optical lattices

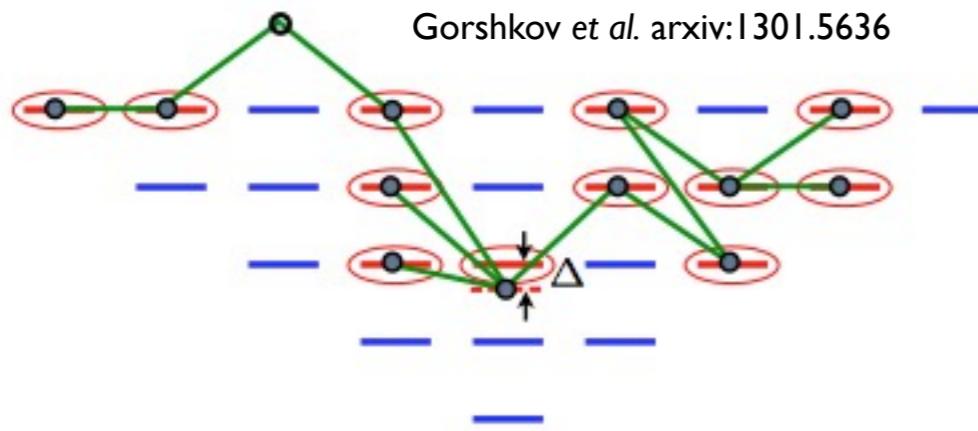
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- Q1: What can we quantum simulate with ultracold molecules in optical lattices?
- A1: Quantum magnetism, topological phases, fractional Chern insulators,...
- Q2: What is the natural many-body Hamiltonian describing up-and-coming ultracold molecules in optical lattices?
- A2: The molecular Hubbard Hamiltonian

Molecules in optical lattices

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- Q1: What can we quantum simulate with ultracold molecules in optical lattices?
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- A2: The molecular Hubbard Hamiltonian

Experimental
Controls

Case Studies

Different molecules
give different physics!

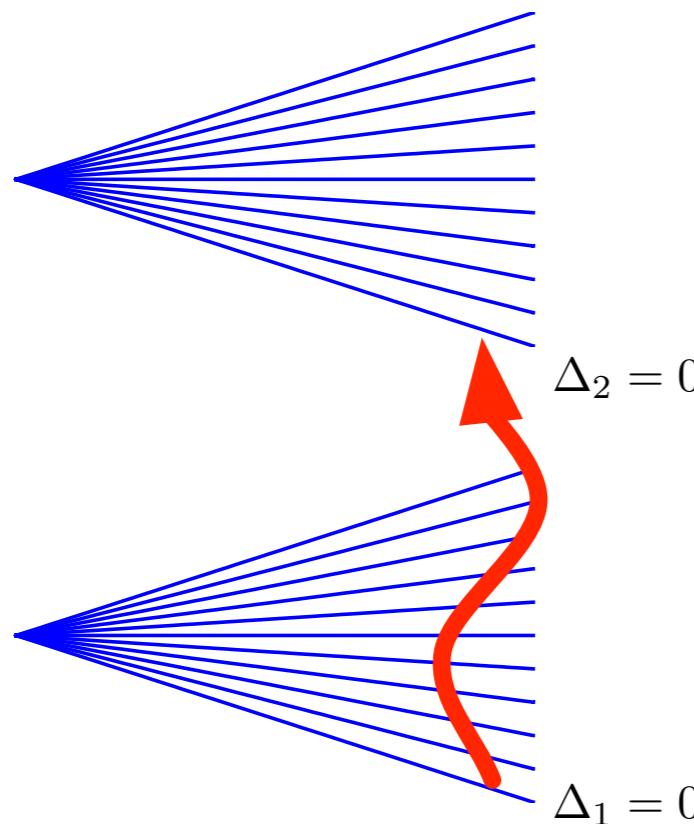
The molecular Hubbard Hamiltonian

$$\begin{aligned}\hat{H} = & \sum_{i\sigma} \Delta_\sigma \hat{n}_{i,\sigma} - \sum_{i \neq j; r_t} \sum_{\sigma} t_{j-i,\sigma} \hat{a}_{i,\sigma}^\dagger \hat{a}_{j,\sigma} + \frac{1}{2} \sum_{i \neq j; r_U} \sum_{\sigma\sigma'} U_{j-i,\sigma\sigma'} \hat{n}_{i,\sigma} \hat{n}_{j,\sigma'} \\ & - \sum_{\sigma\sigma'} \Omega_{\sigma\sigma'} \sum_i \hat{a}_{i,\sigma}^\dagger \hat{a}_{i,\sigma'} + \frac{1}{4} \sum_{i \neq j; r_E} \sum_{\sigma_1\sigma_2\sigma'_2\sigma'_1} E_{j-i,\sigma_1\sigma_2\sigma'_2\sigma'_1} \hat{a}_{i,\sigma_1}^\dagger \hat{a}_{j,\sigma_2}^\dagger \hat{a}_{j,\sigma'_2} \hat{a}_{i,\sigma'_1}\end{aligned}$$

The molecular Hubbard Hamiltonian

Detuning

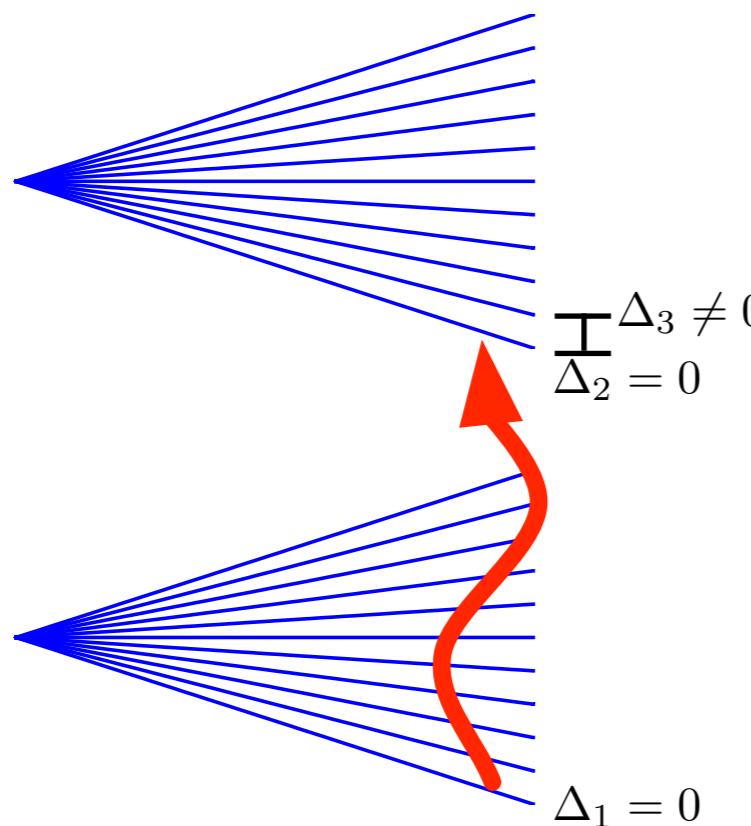
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The molecular Hubbard Hamiltonian

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The molecular Hubbard Hamiltonian

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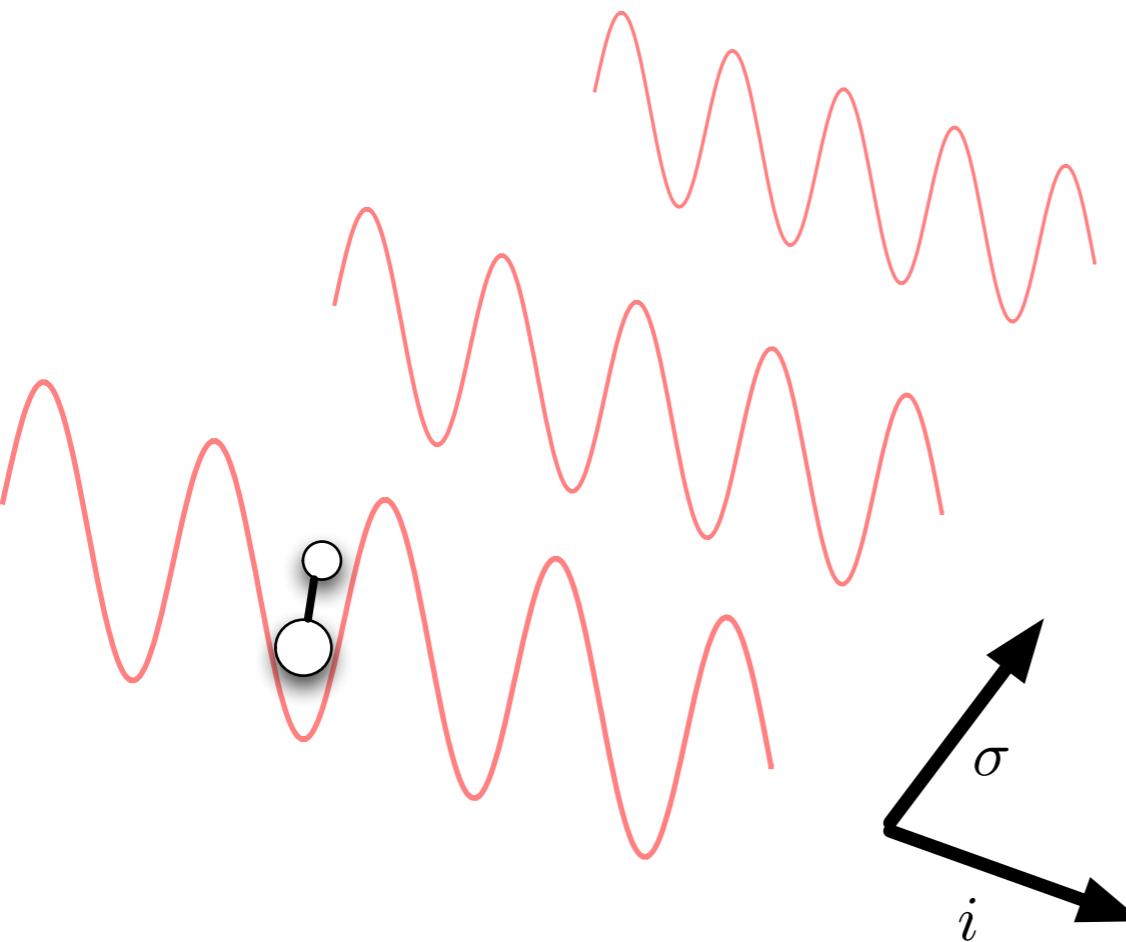
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$$\Delta_2 = 0$$

The molecular Hubbard Hamiltonian

Tunneling

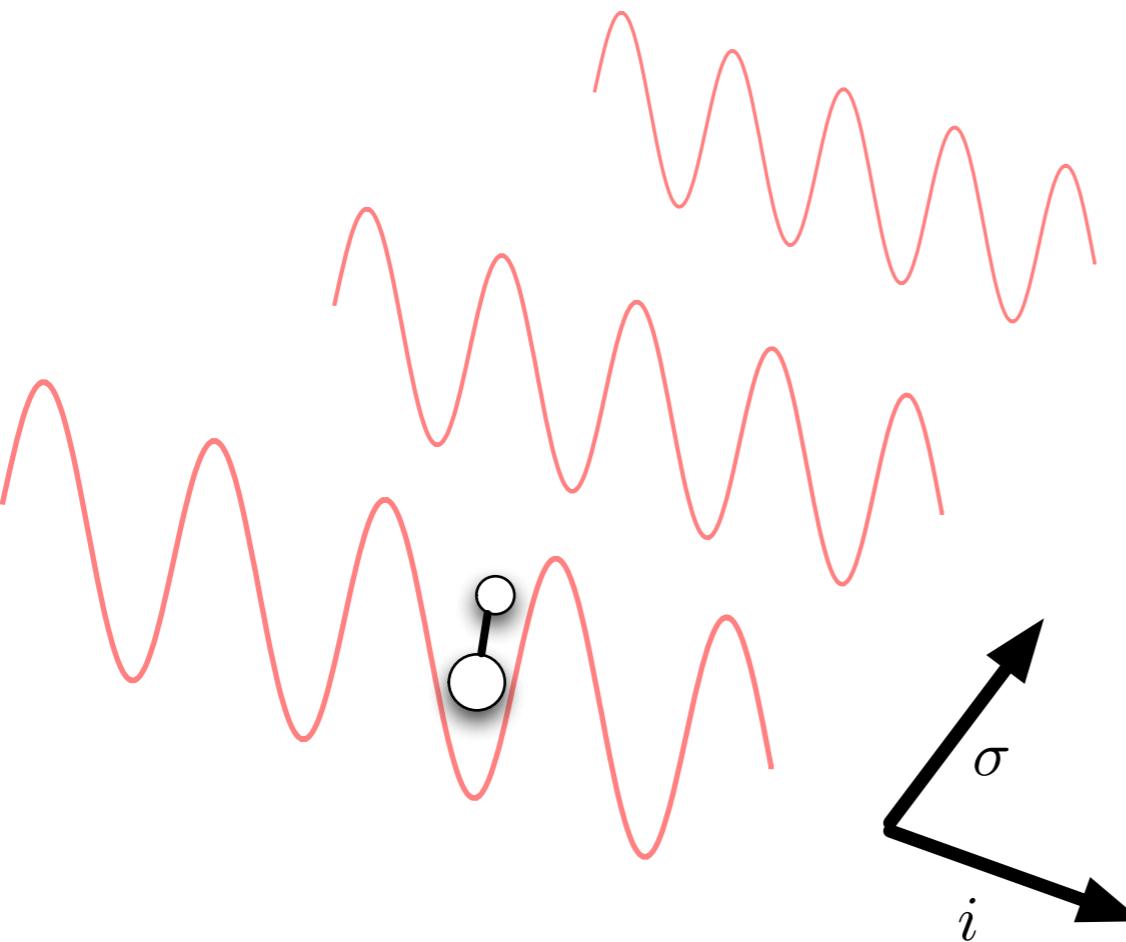
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The molecular Hubbard Hamiltonian

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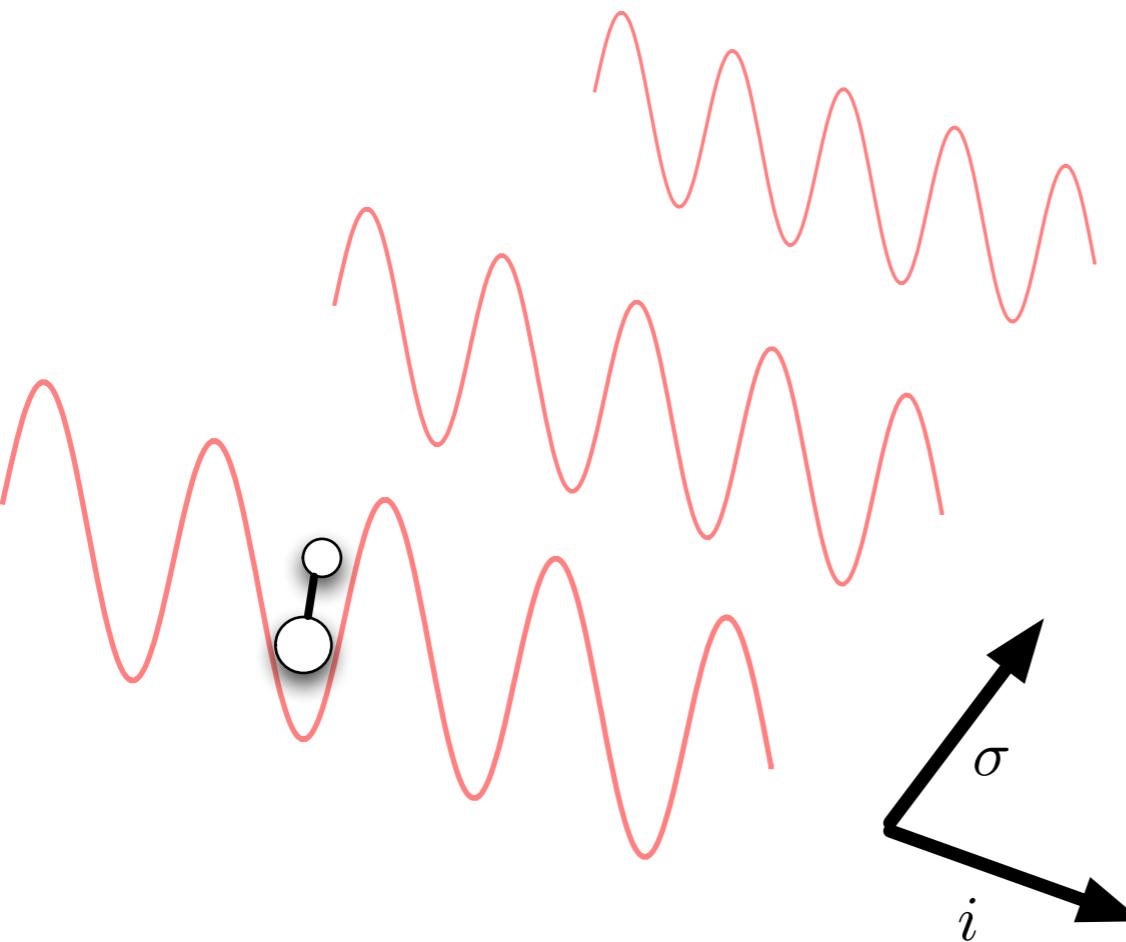
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The molecular Hubbard Hamiltonian

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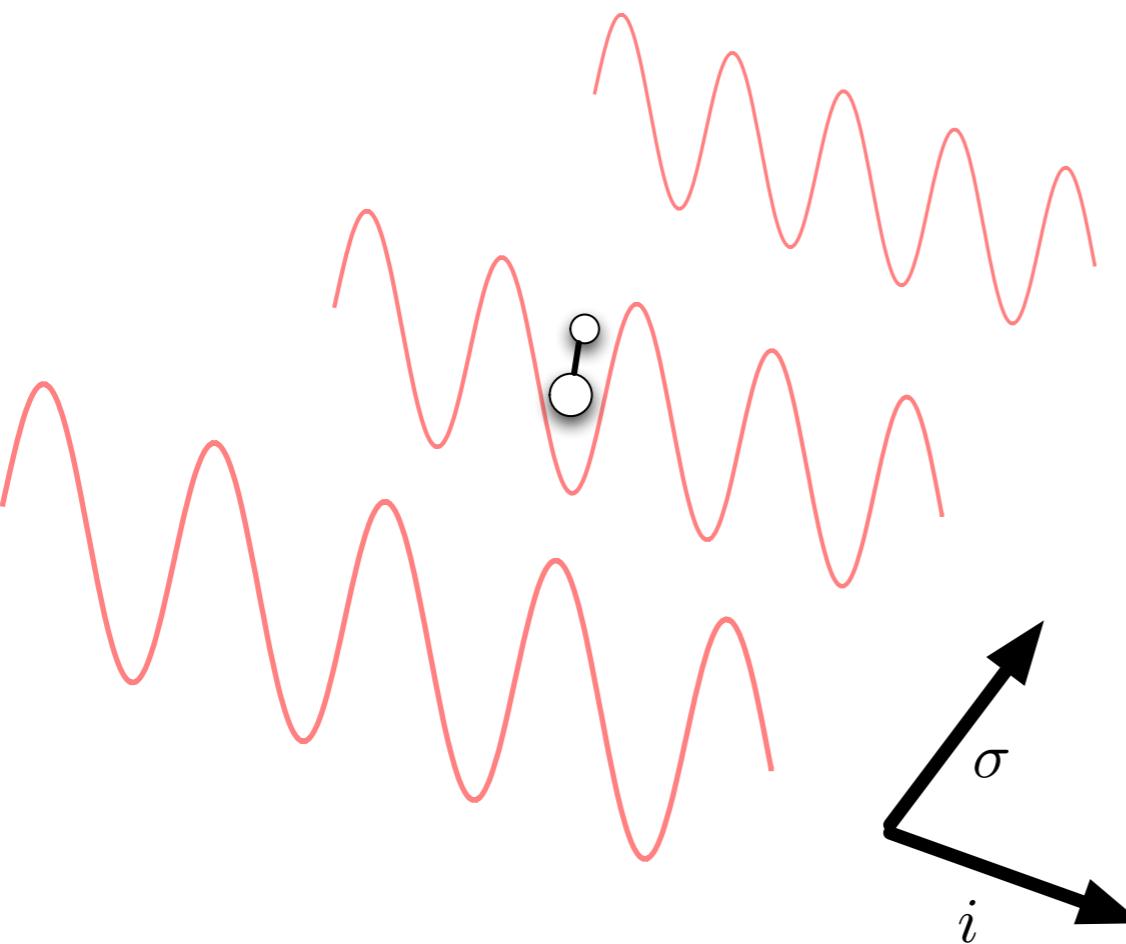
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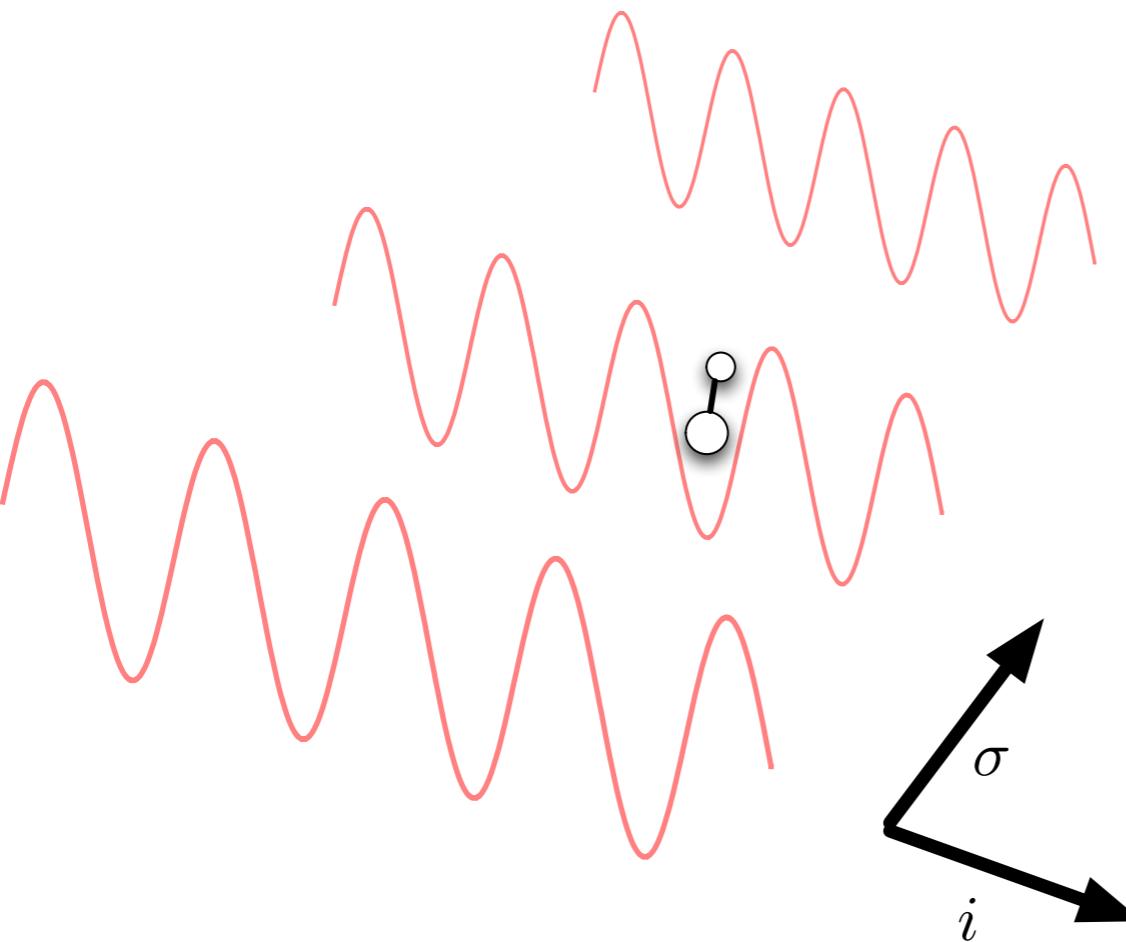
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The molecular Hubbard Hamiltonian

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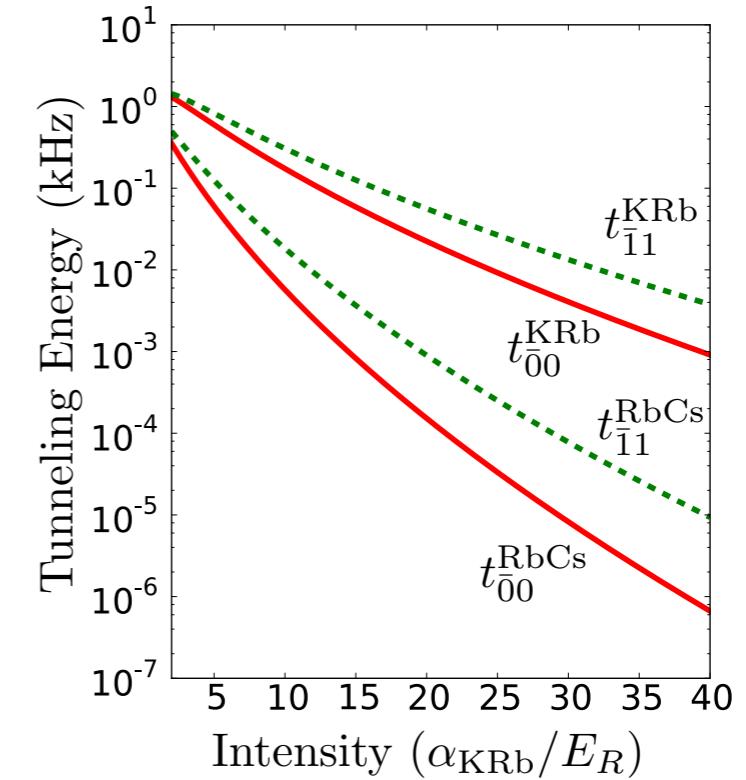
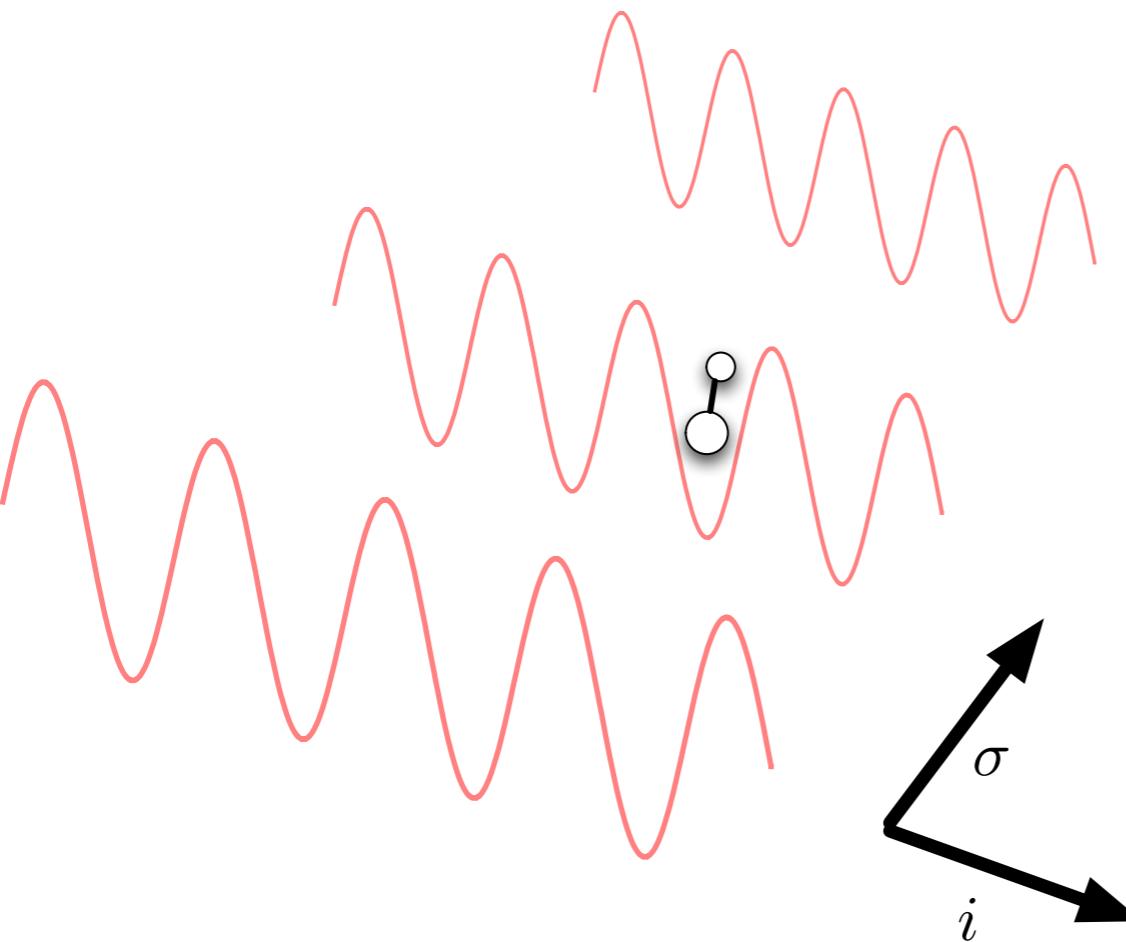
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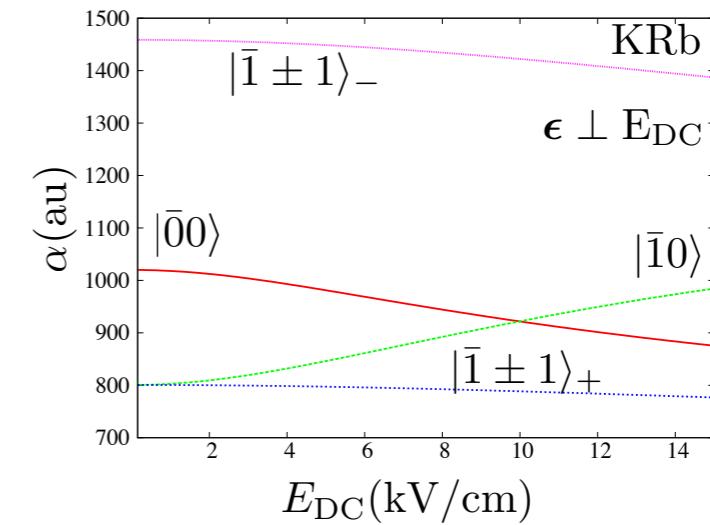
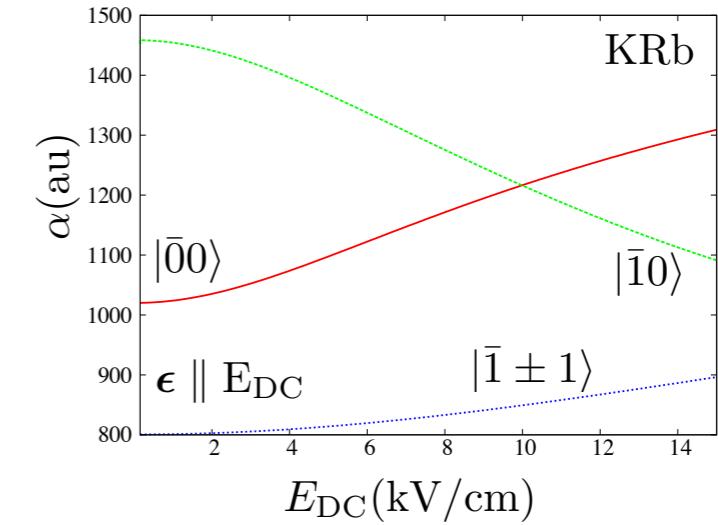
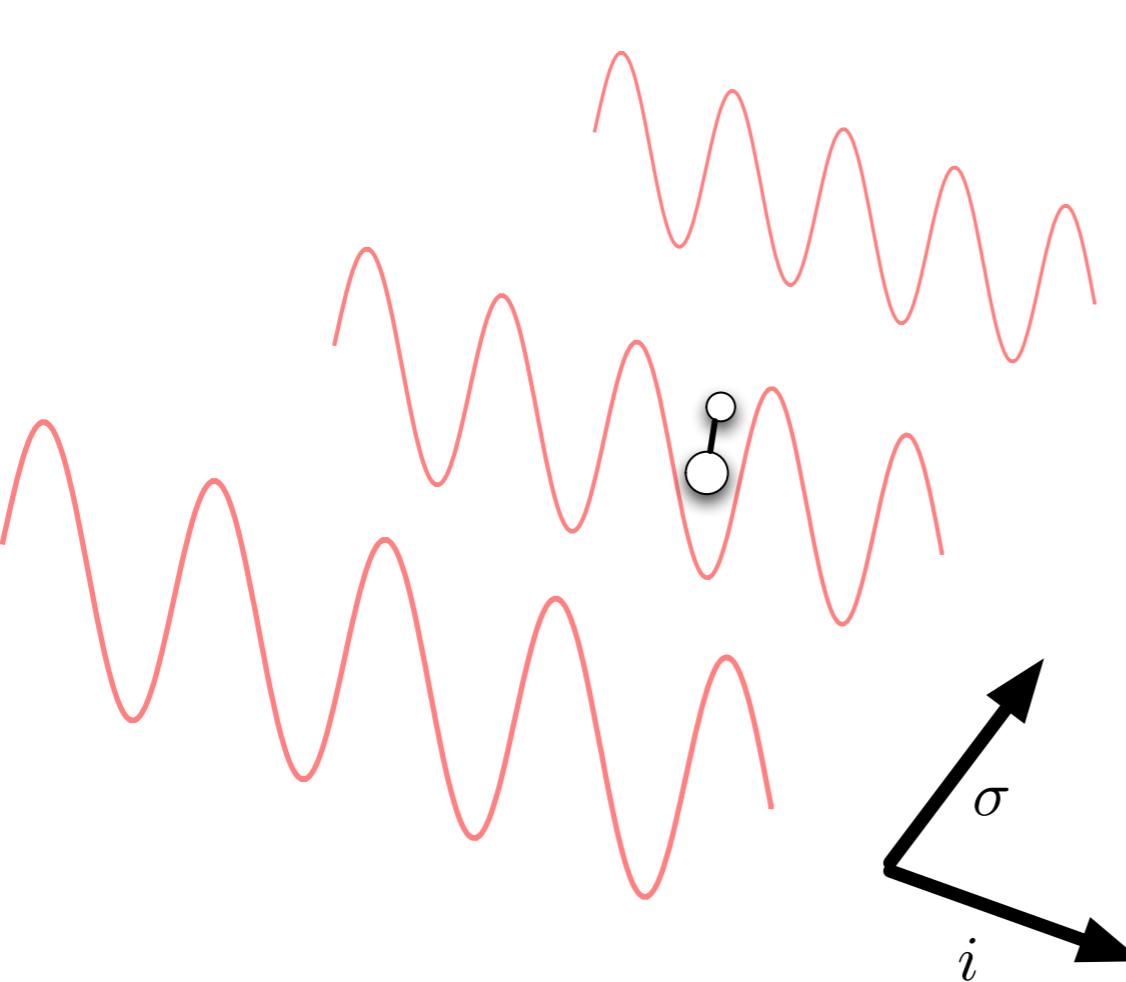
$$\hat{H} = \sum_{i\sigma} \Delta_\sigma \hat{n}_{i,\sigma} - \boxed{\sum_{i \neq j; r_t} \sum_{\sigma} t_{j-i,\sigma} \hat{a}_{i,\sigma}^\dagger \hat{a}_{j,\sigma}} + \frac{1}{2} \sum_{i \neq j; r_U} \sum_{\sigma\sigma'} U_{j-i,\sigma\sigma'} \hat{n}_{i,\sigma} \hat{n}_{j,\sigma'} \\ - \sum_{\sigma\sigma'} \Omega_{\sigma\sigma'} \sum_i \hat{a}_{i,\sigma}^\dagger \hat{a}_{i,\sigma'} + \frac{1}{4} \sum_{i \neq j; r_E} \sum_{\sigma_1 \sigma_2 \sigma'_2 \sigma'_1} E_{j-i,\sigma_1 \sigma_2 \sigma'_2 \sigma'_1} \hat{a}_{i,\sigma_1}^\dagger \hat{a}_{j,\sigma_2}^\dagger \hat{a}_{j,\sigma'_2} \hat{a}_{i,\sigma'_1}$$



The molecular Hubbard Hamiltonian

Tunneling

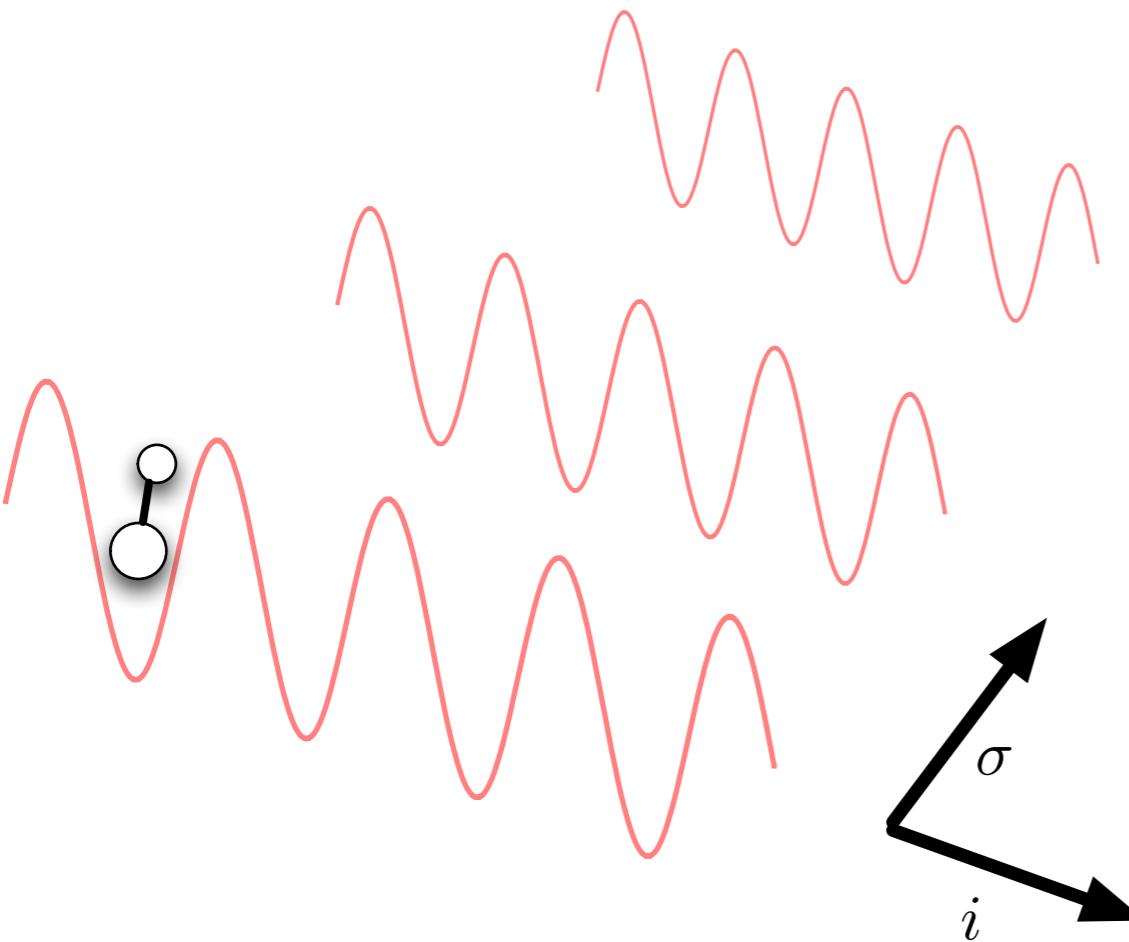
$$\hat{H} = \sum_{i\sigma} \Delta_\sigma \hat{n}_{i,\sigma} - \boxed{\sum_{i \neq j; r_t} \sum_{\sigma} t_{j-i,\sigma} \hat{a}_{i,\sigma}^\dagger \hat{a}_{j,\sigma}} + \frac{1}{2} \sum_{i \neq j; r_U} \sum_{\sigma\sigma'} U_{j-i,\sigma\sigma'} \hat{n}_{i,\sigma} \hat{n}_{j,\sigma'} \\ - \sum_{\sigma\sigma'} \Omega_{\sigma\sigma'} \sum_i \hat{a}_{i,\sigma}^\dagger \hat{a}_{i,\sigma'} + \frac{1}{4} \sum_{i \neq j; r_E} \sum_{\sigma_1\sigma_2\sigma'_2\sigma'_1} E_{j-i,\sigma_1\sigma_2\sigma'_2\sigma'_1} \hat{a}_{i,\sigma_1}^\dagger \hat{a}_{j,\sigma_2}^\dagger \hat{a}_{j,\sigma'_2} \hat{a}_{i,\sigma'_1}$$



The molecular Hubbard Hamiltonian

$$\hat{H} = \sum_{i\sigma} \Delta_\sigma \hat{n}_{i,\sigma} - \sum_{i \neq j; r_t} \sum_{\sigma} t_{j-i,\sigma} \hat{a}_{i,\sigma}^\dagger \hat{a}_{j,\sigma} + \frac{1}{2} \sum_{i \neq j; r_U} \sum_{\sigma\sigma'} U_{j-i,\sigma\sigma'} \hat{n}_{i,\sigma} \hat{n}_{j,\sigma'} \\ - \sum_{\sigma\sigma'} \Omega_{\sigma\sigma'} \sum_i \hat{a}_{i,\sigma}^\dagger \hat{a}_{i,\sigma'} + \frac{1}{4} \sum_{i \neq j; r_E} \sum_{\sigma_1 \sigma_2 \sigma'_2 \sigma'_1} E_{j-i,\sigma_1 \sigma_2 \sigma'_2 \sigma'_1} \hat{a}_{i,\sigma_1}^\dagger \hat{a}_{j,\sigma_2}^\dagger \hat{a}_{j,\sigma'_2} \hat{a}_{i,\sigma'_1}$$

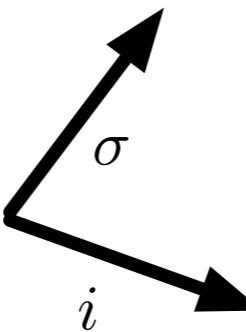
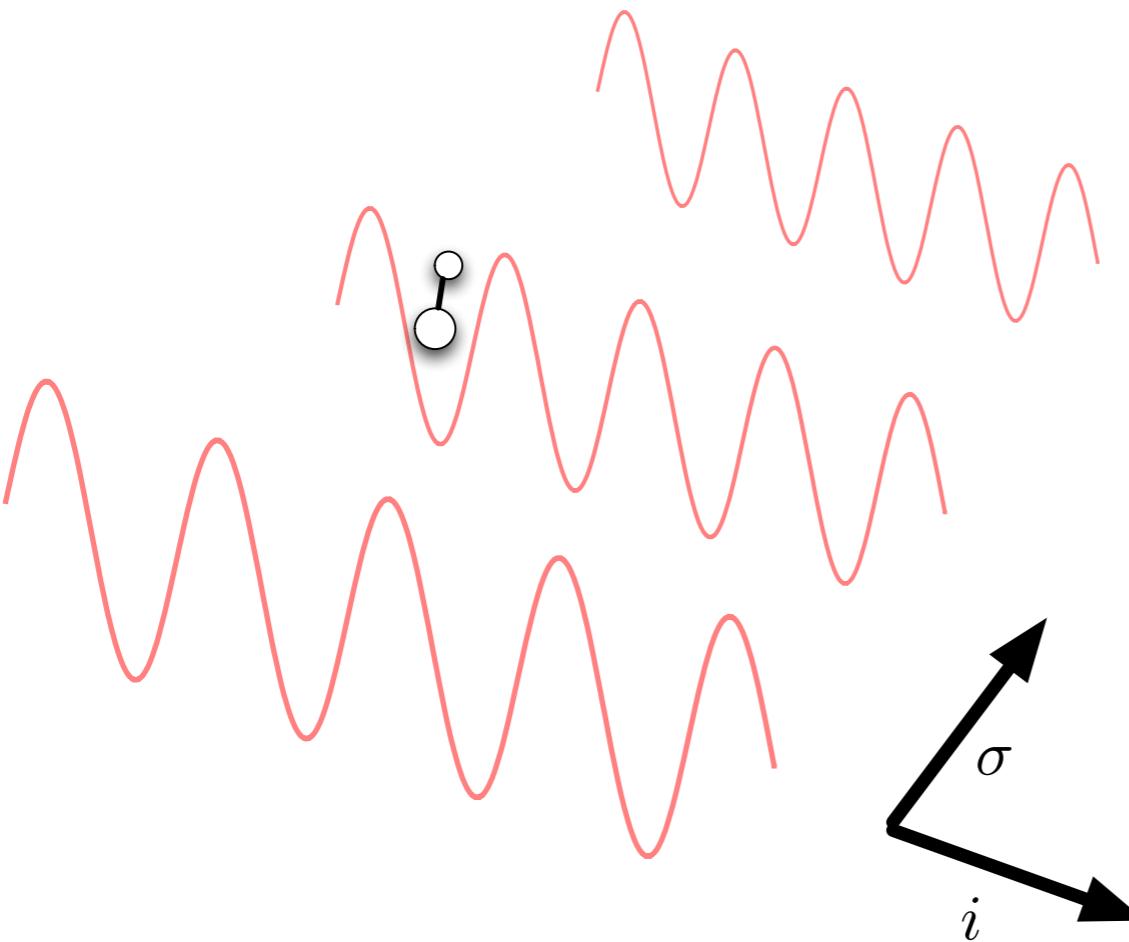
Transitions



The molecular Hubbard Hamiltonian

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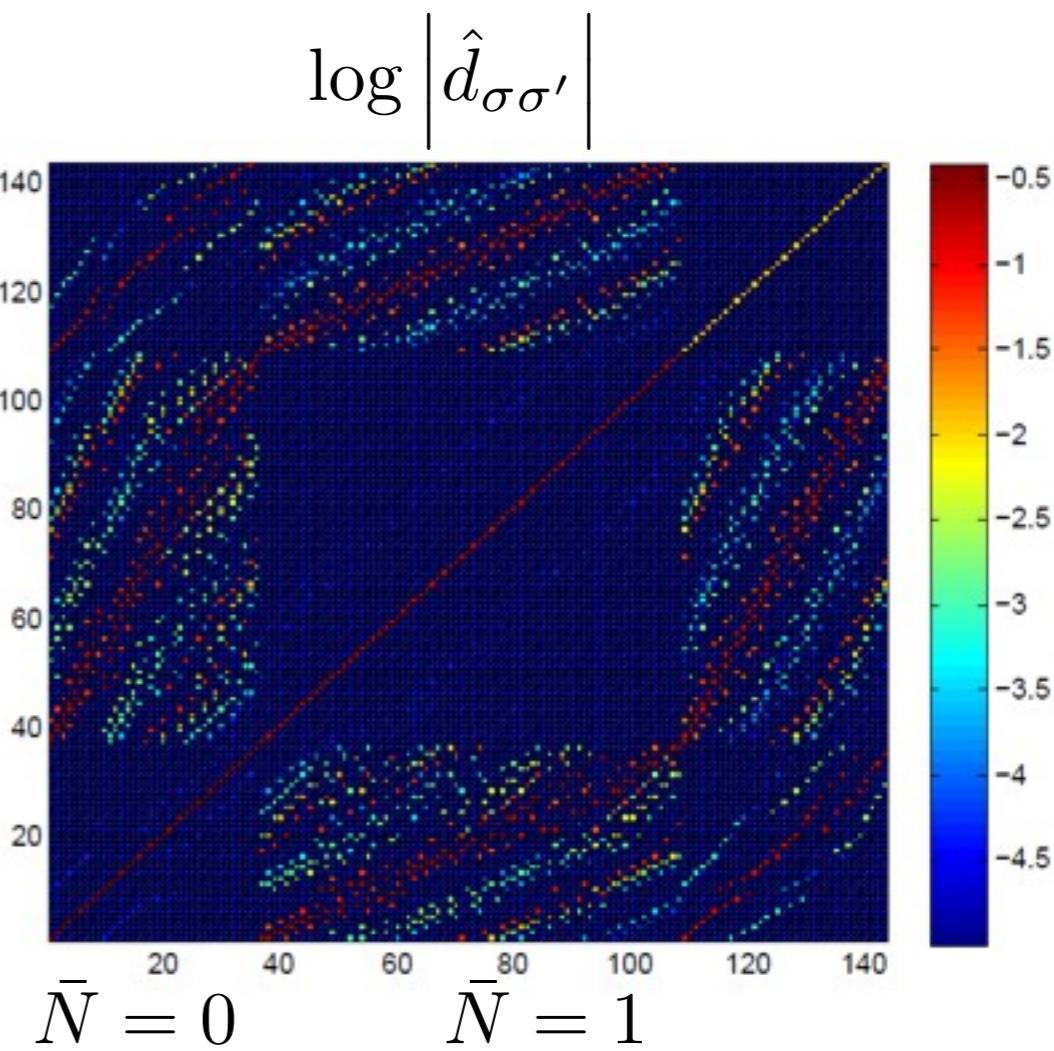
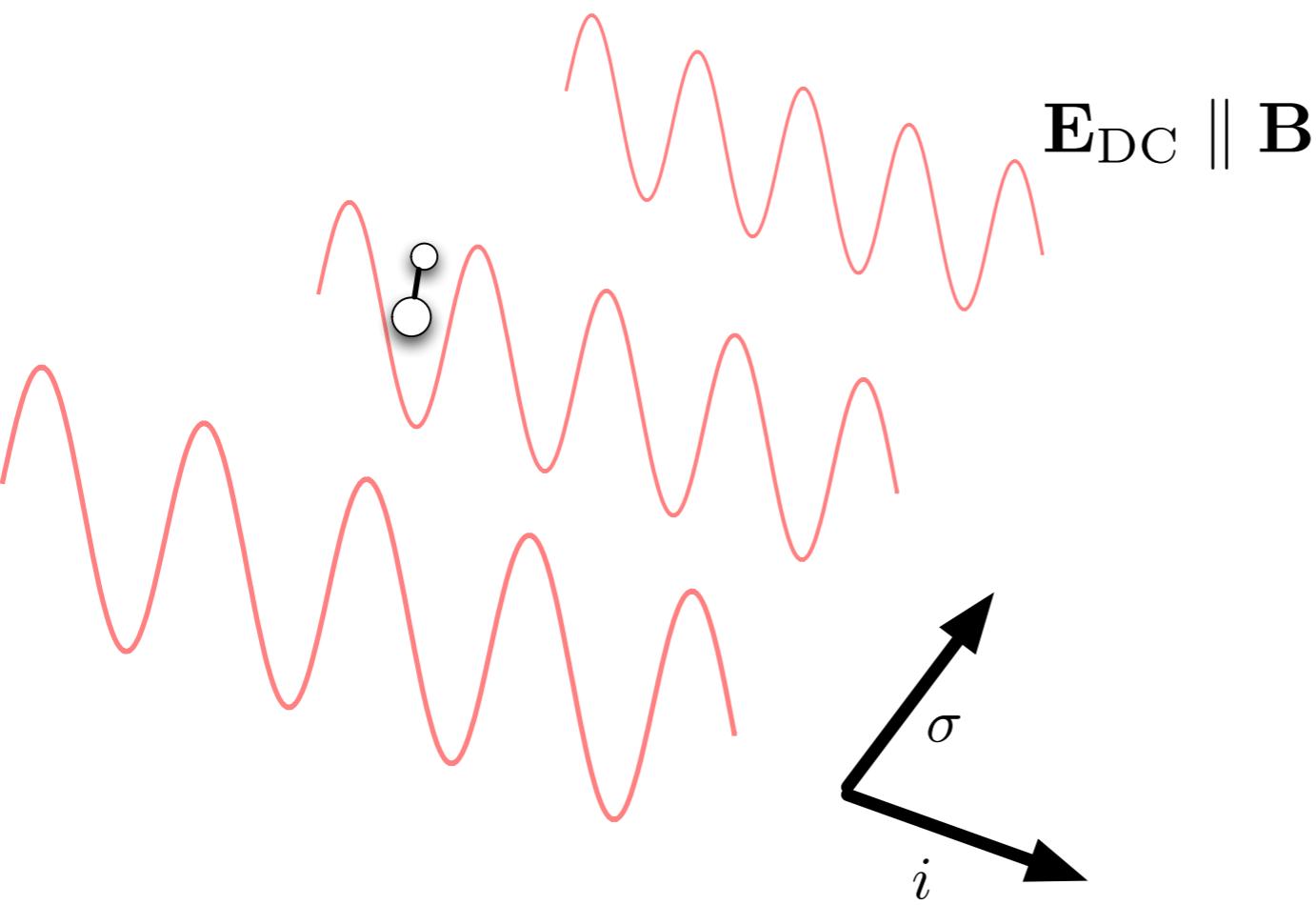
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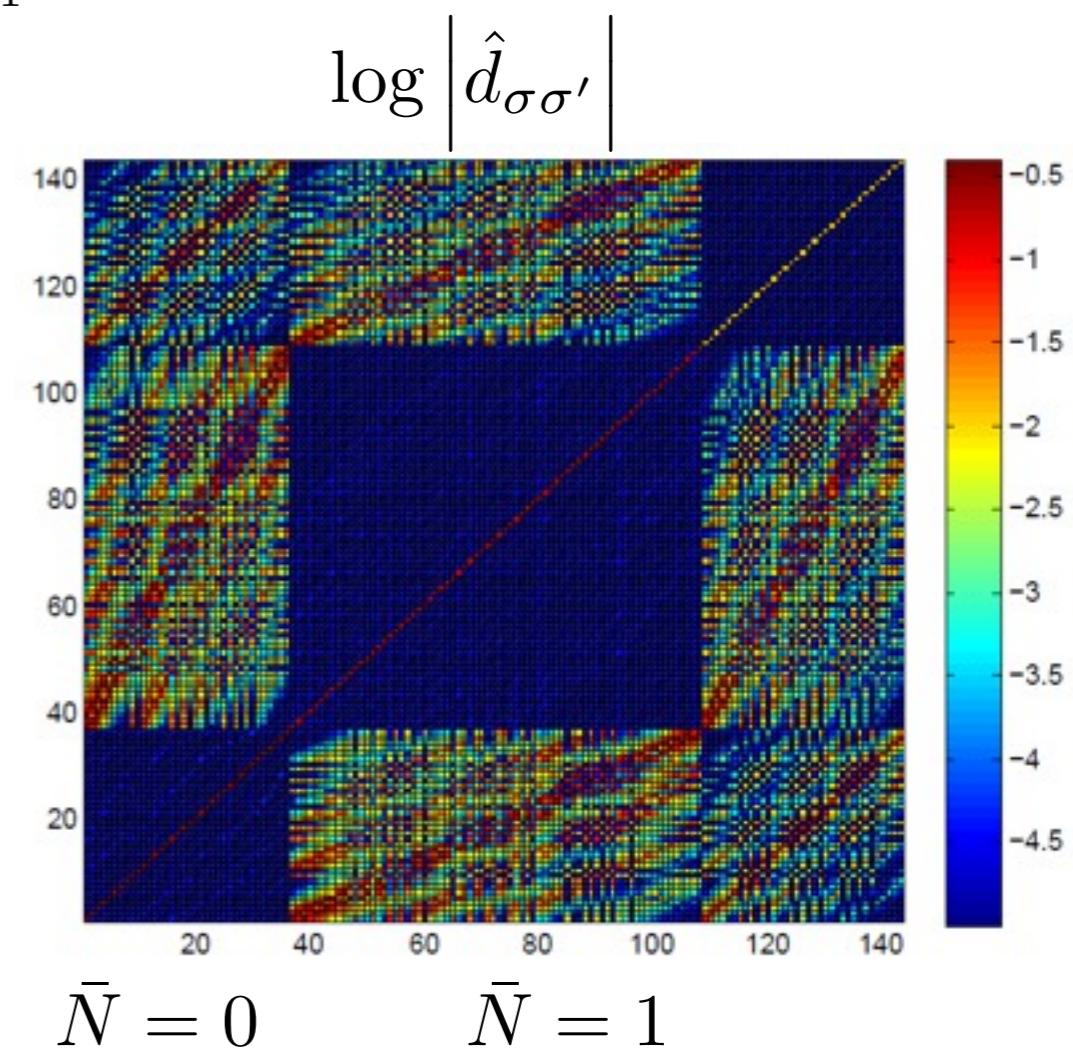
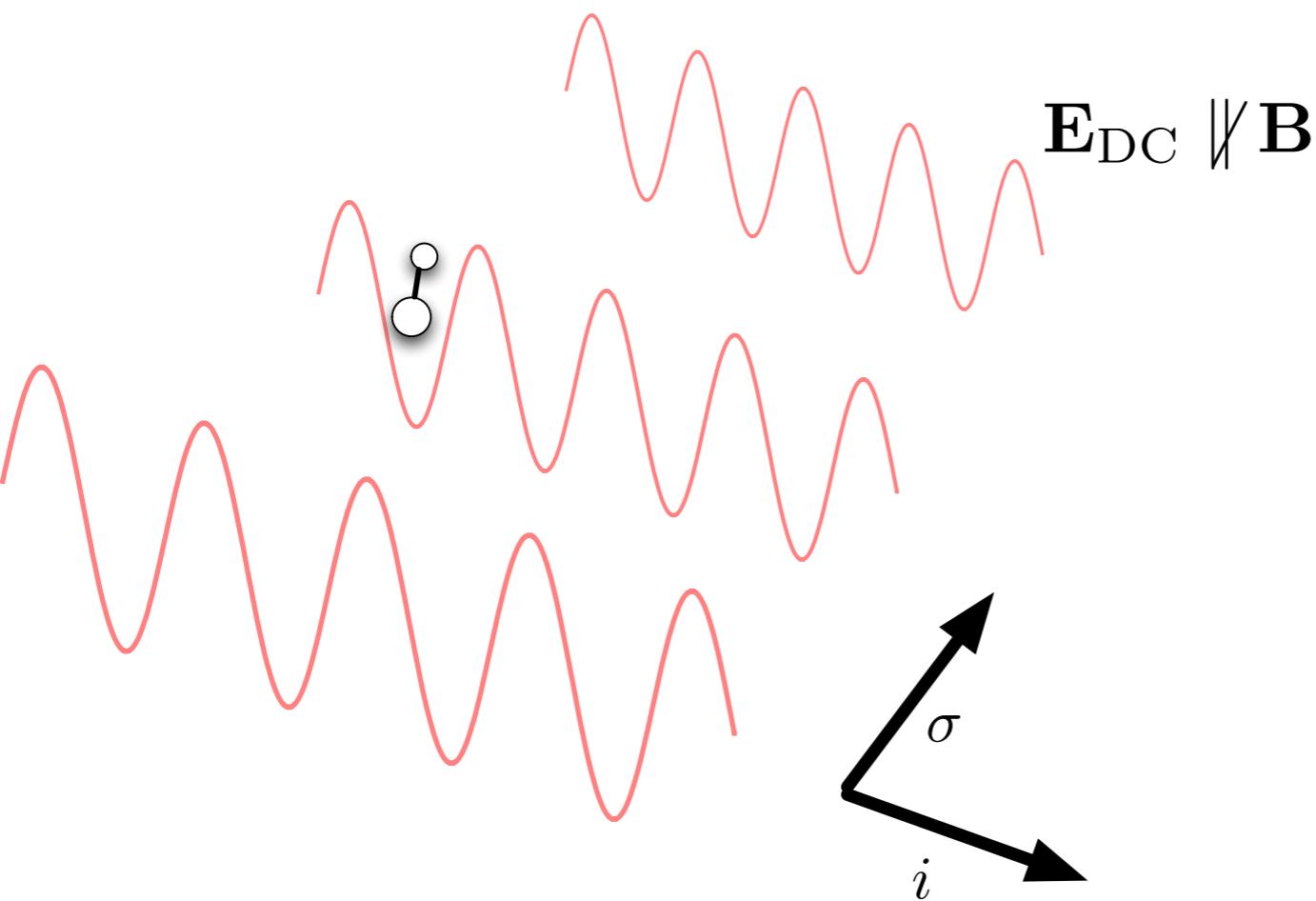
Transitions



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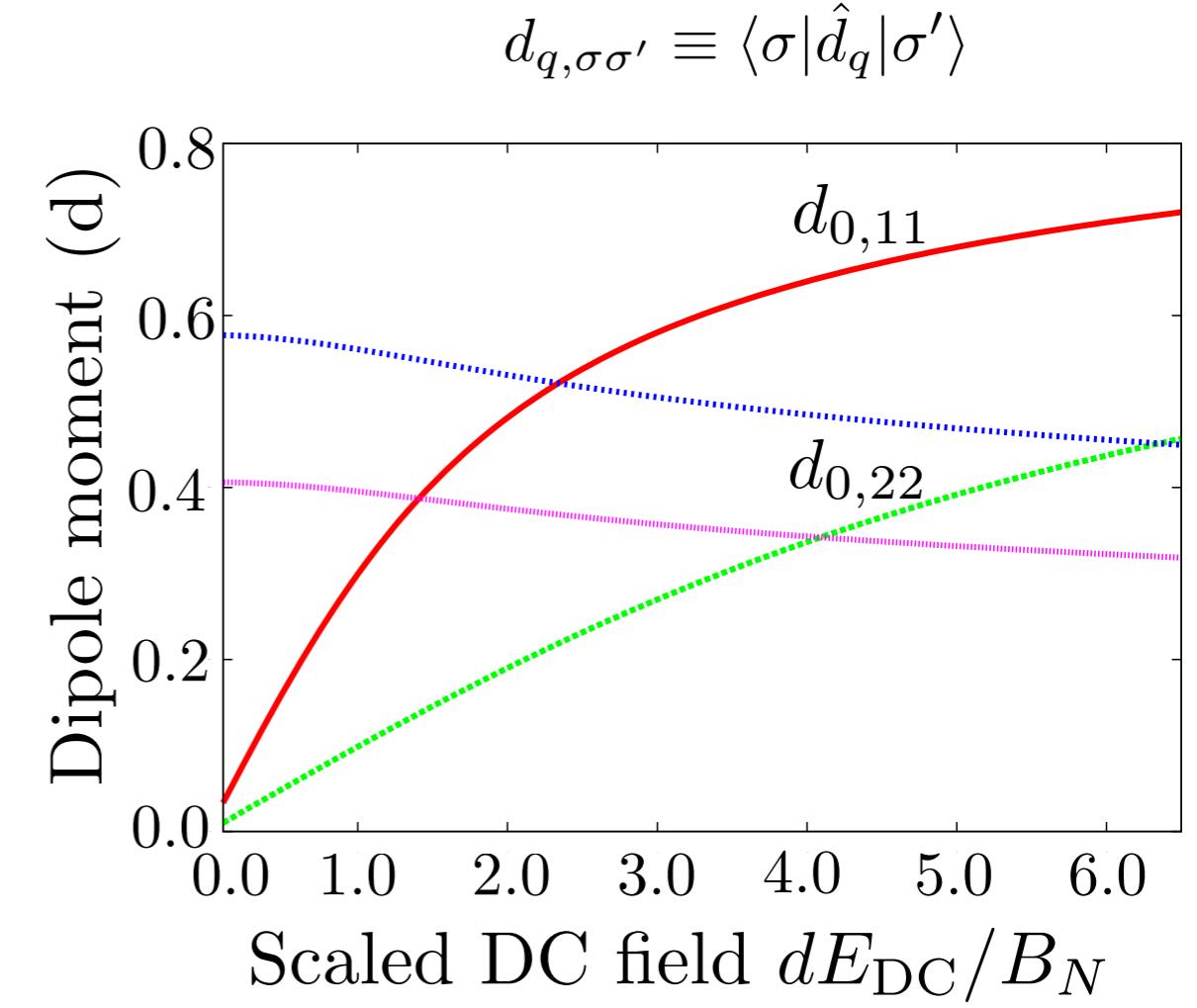
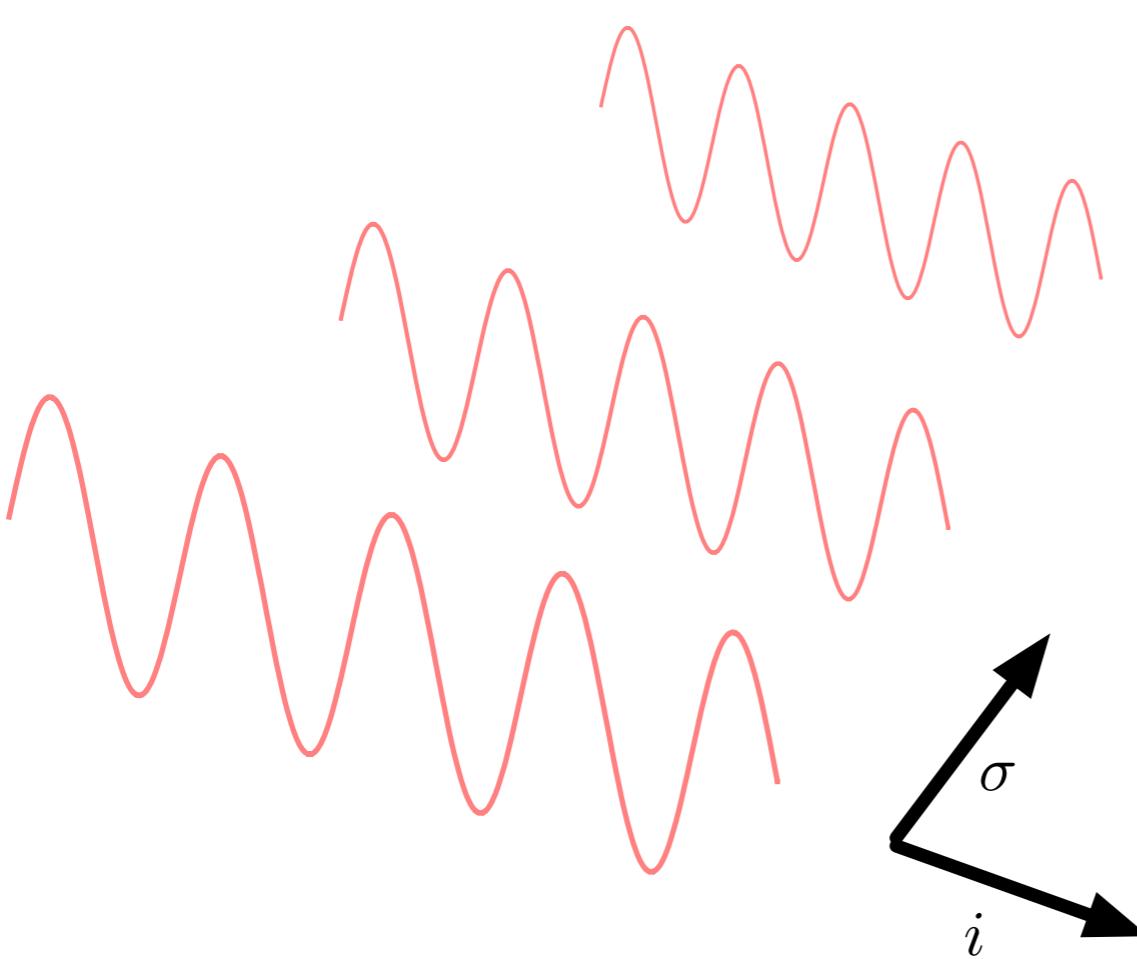
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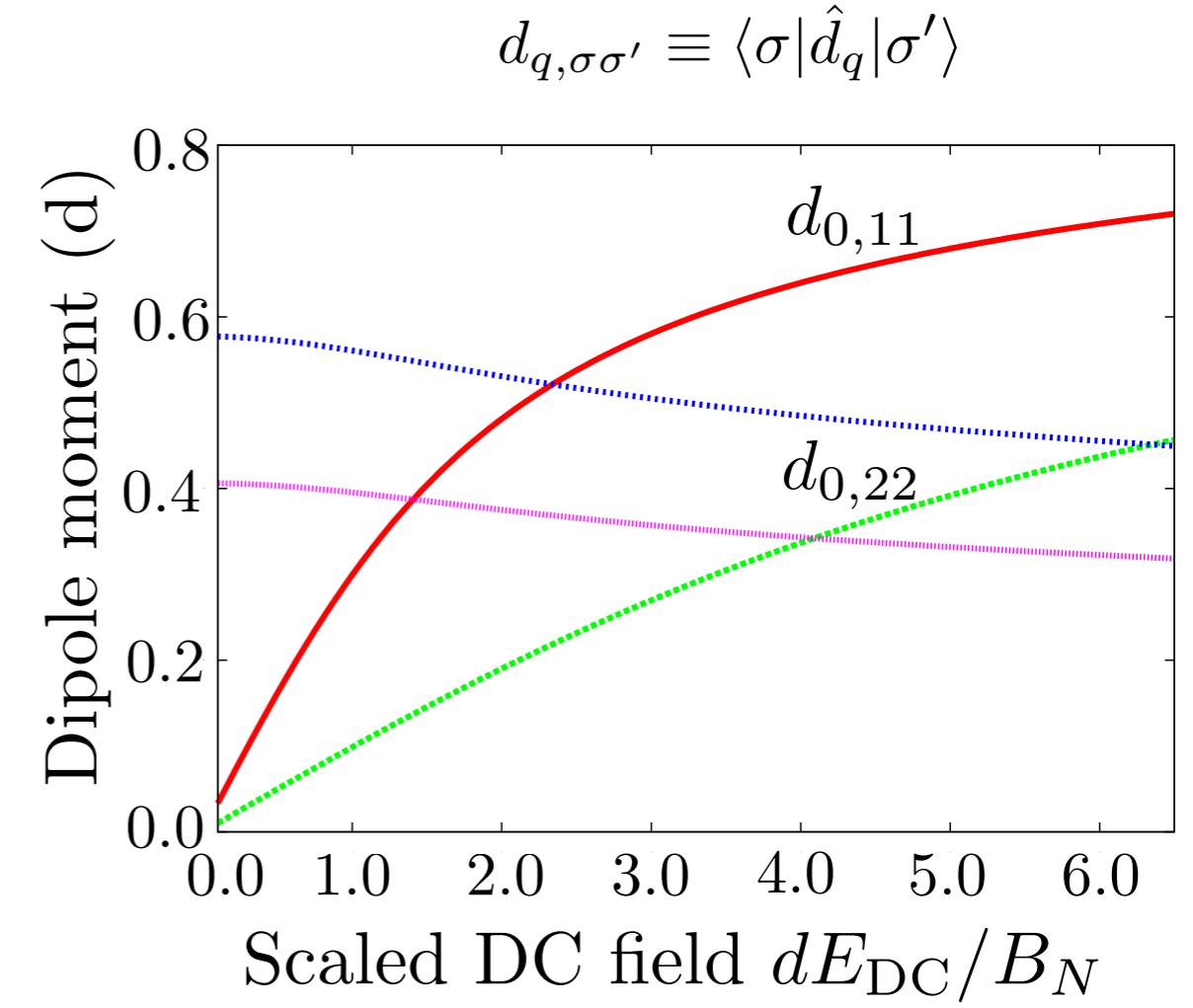
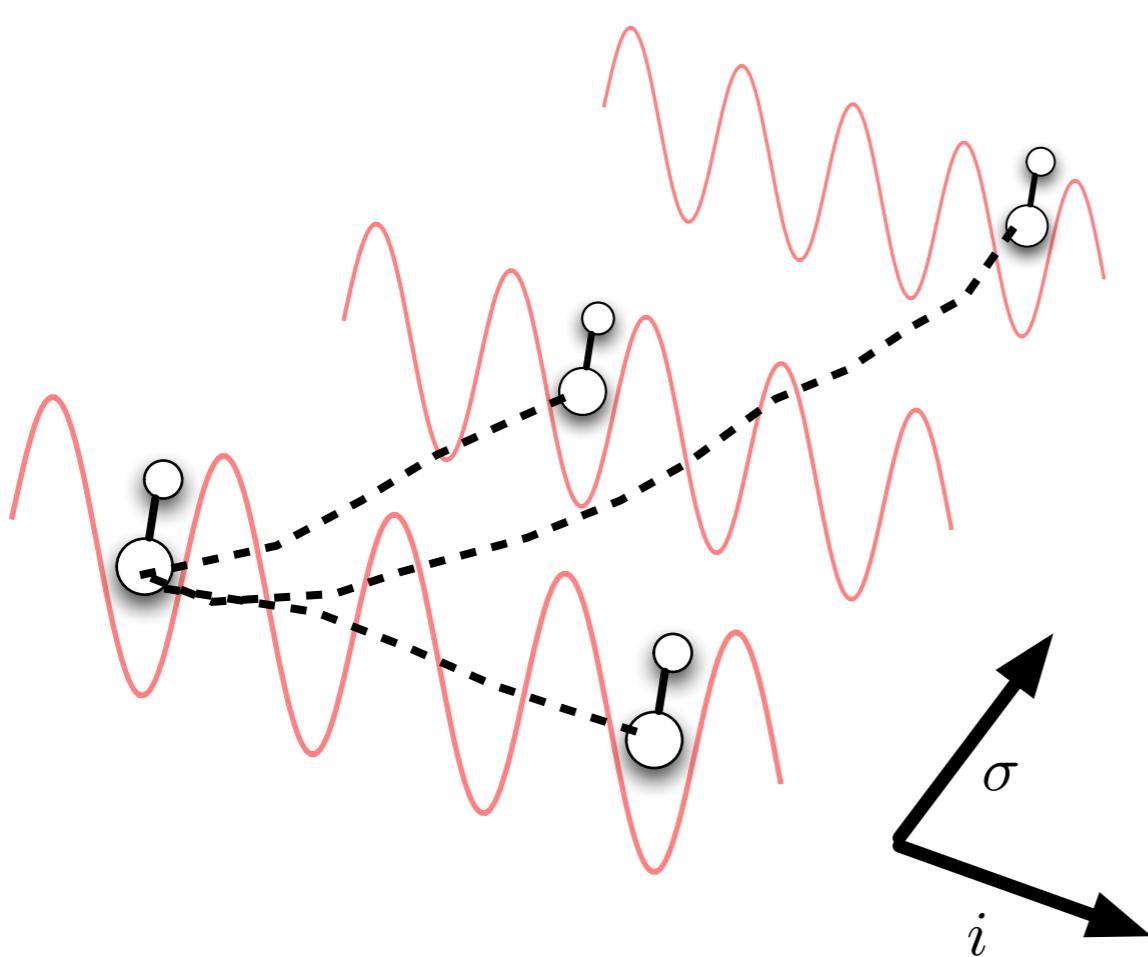
Direct interaction



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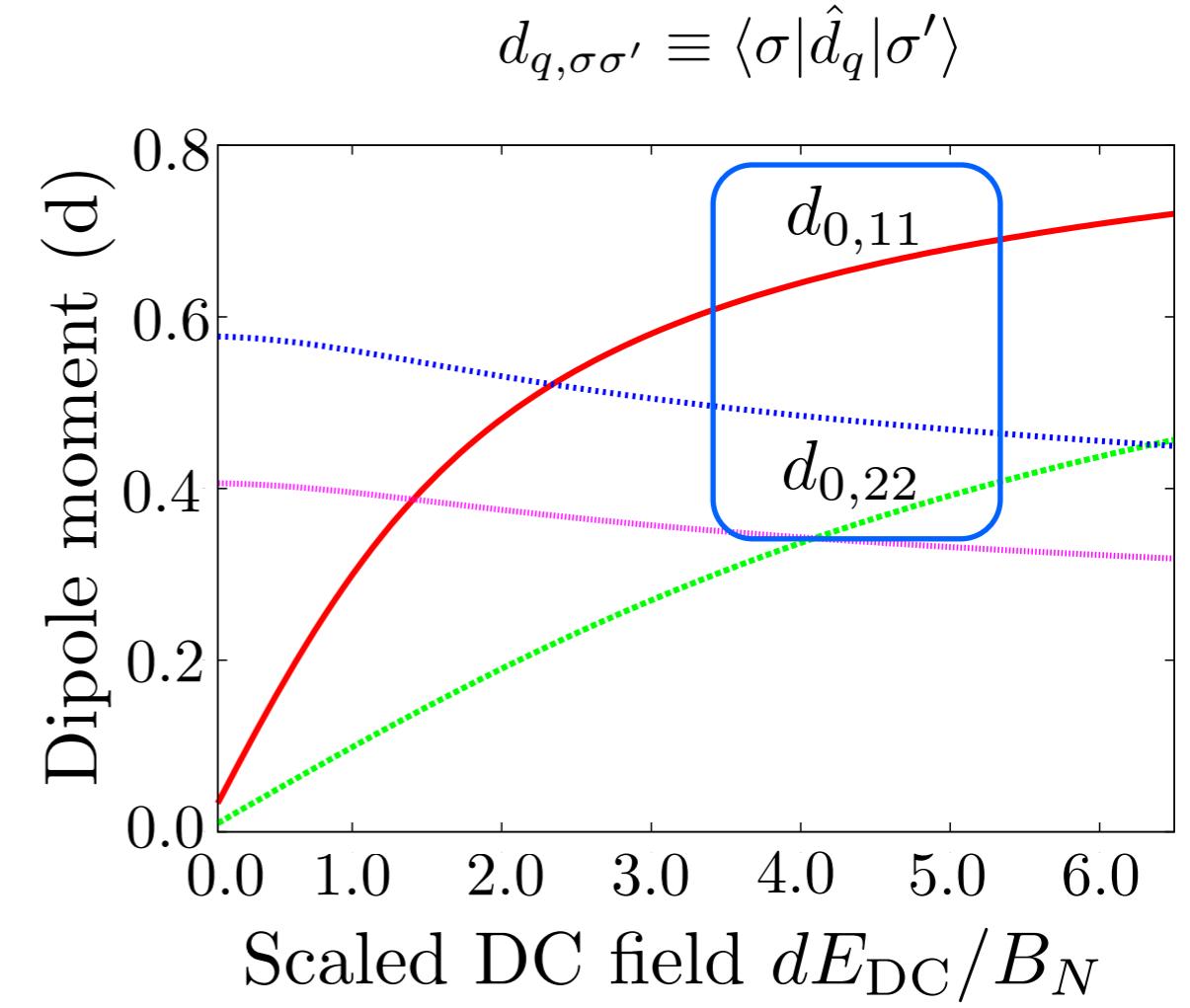
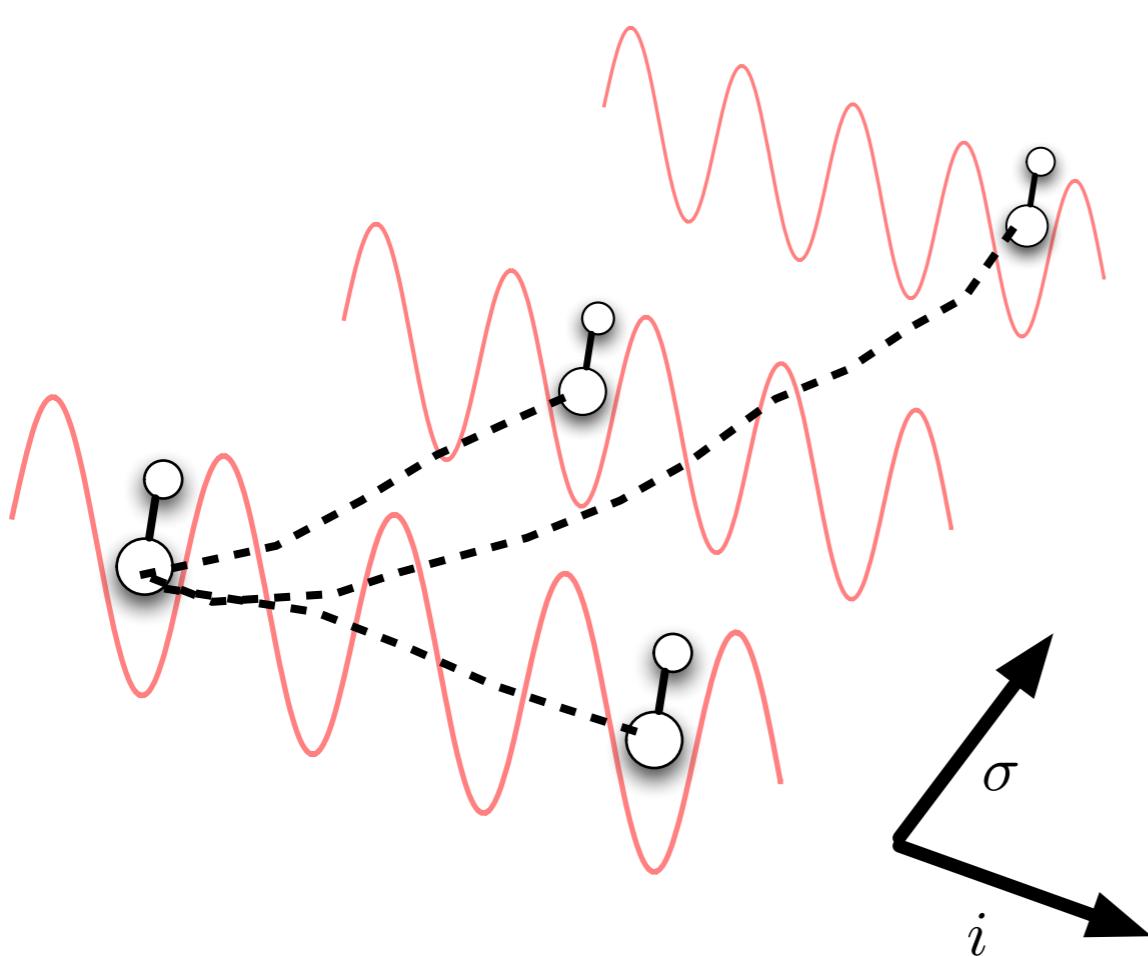
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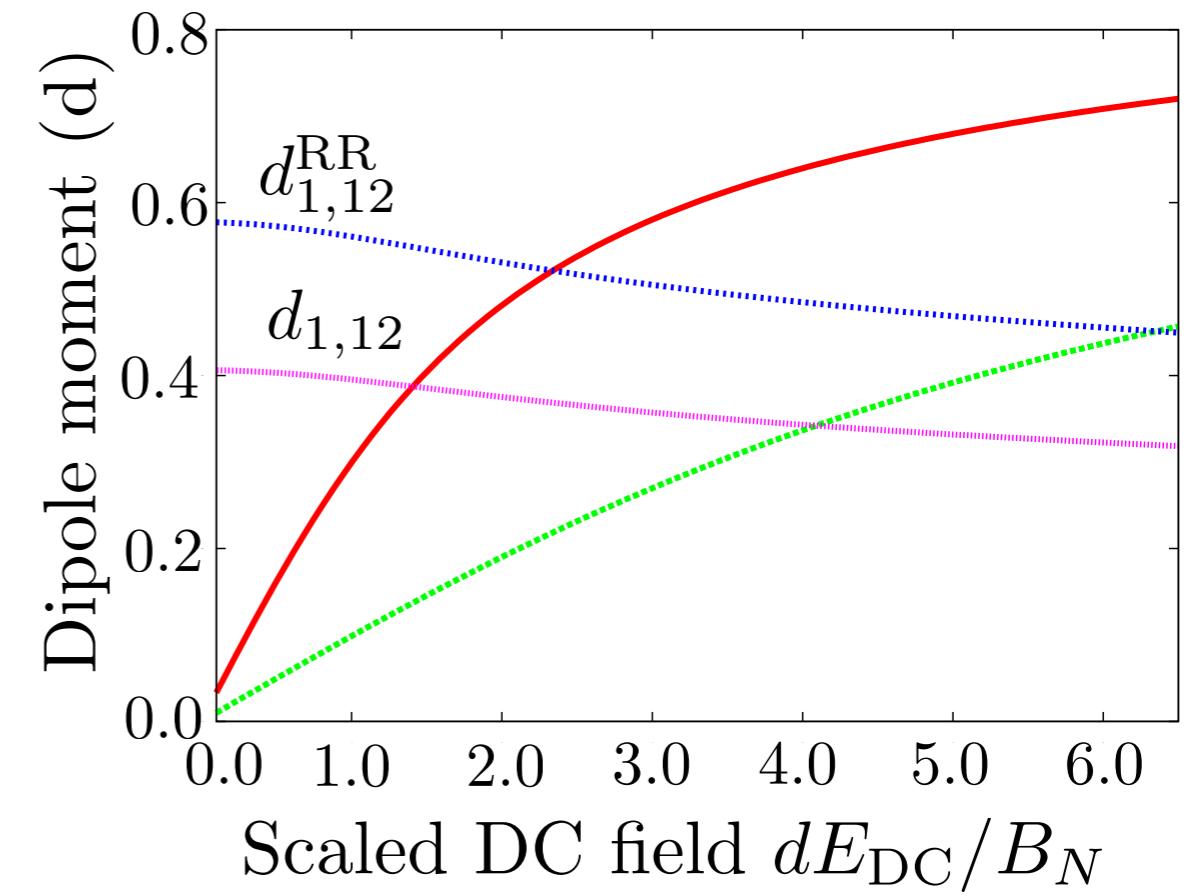
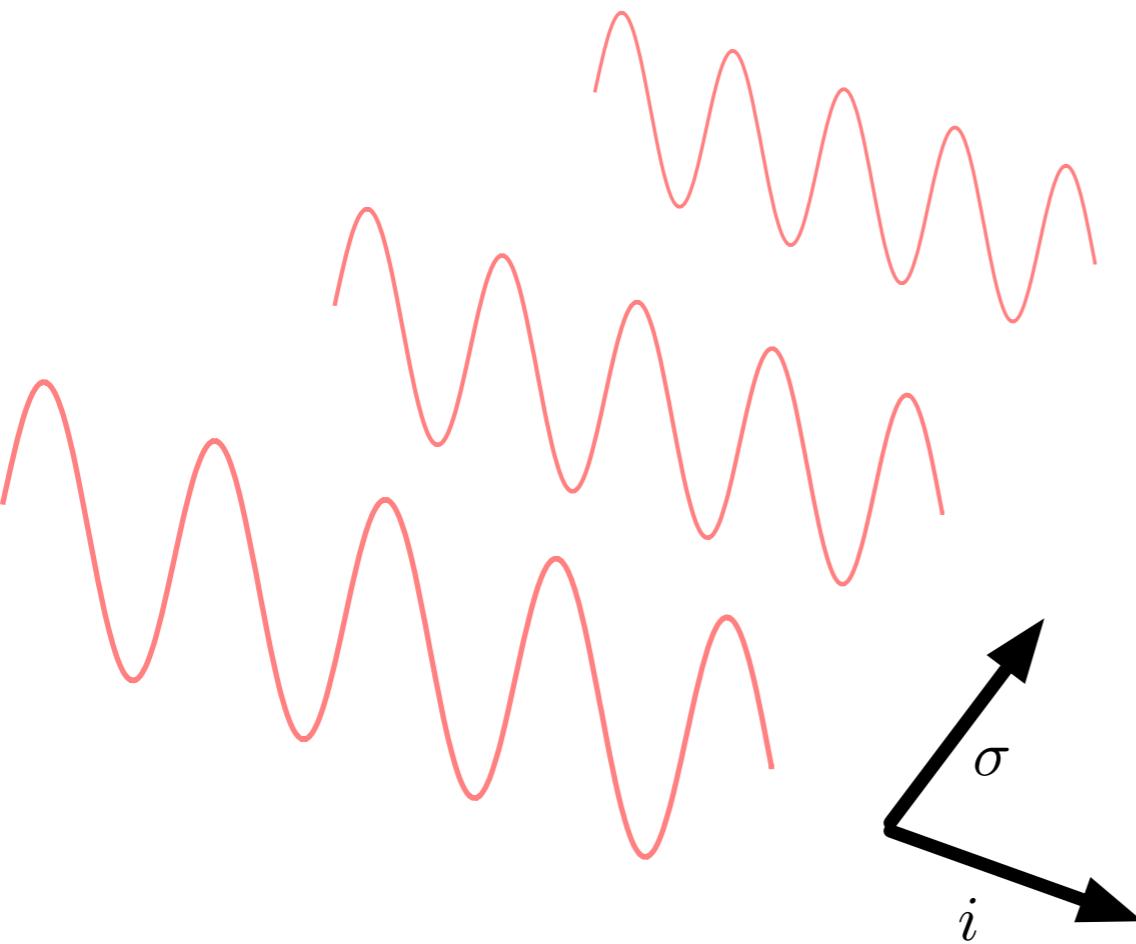
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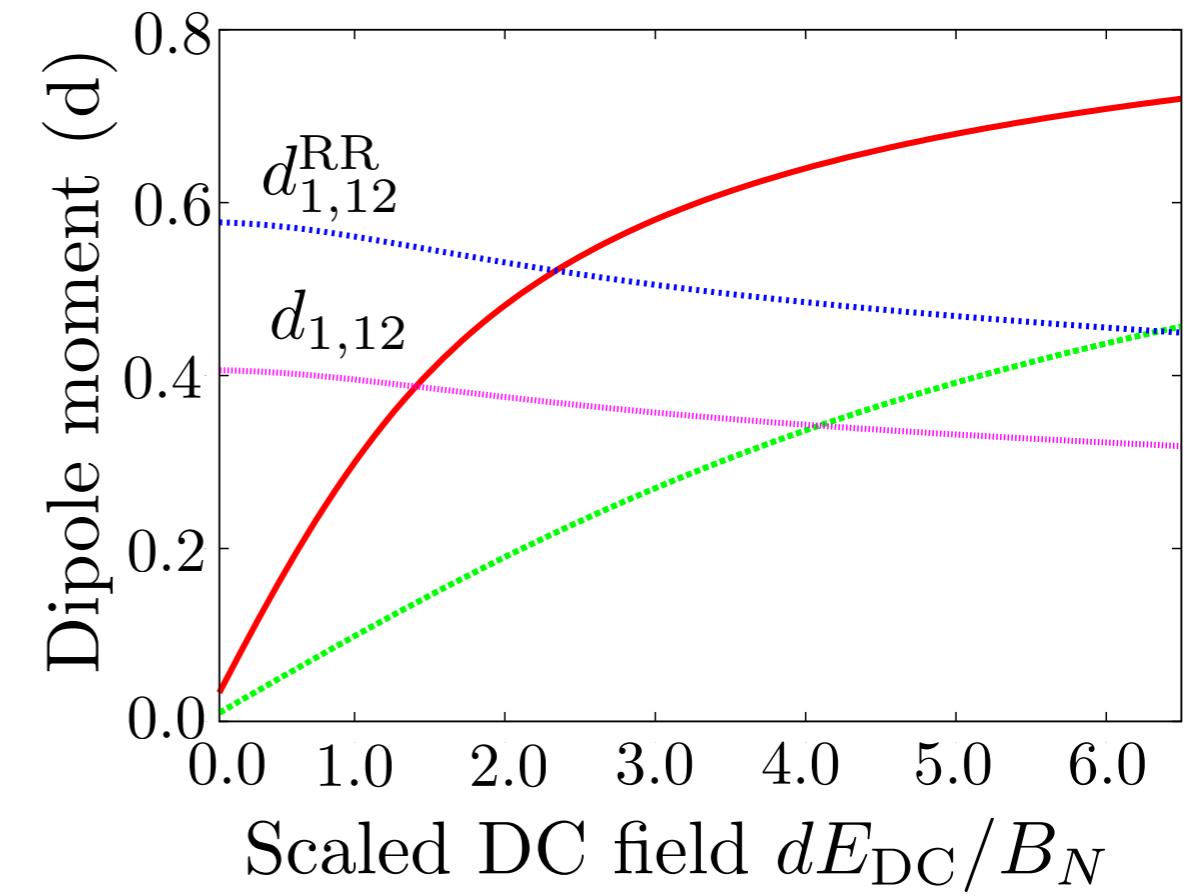
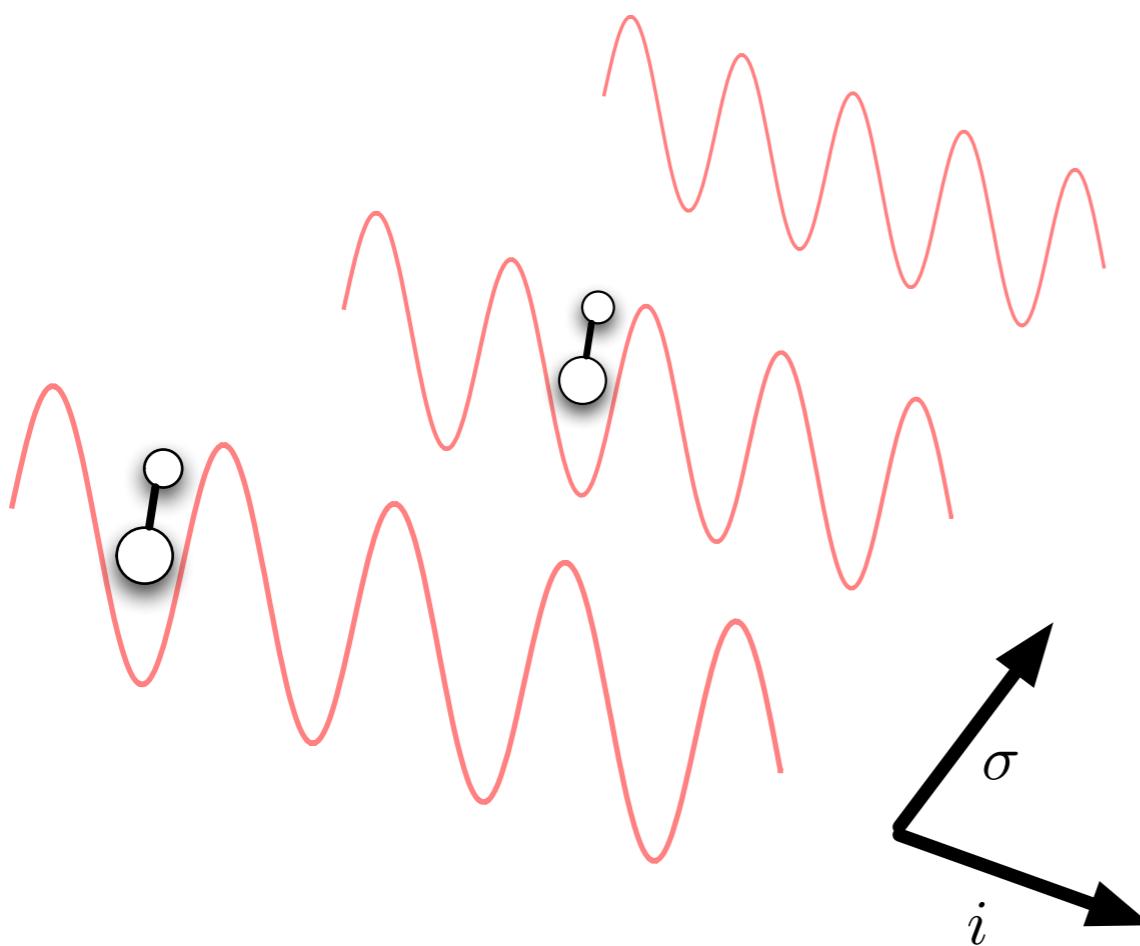
Exchange interaction



The molecular Hubbard Hamiltonian

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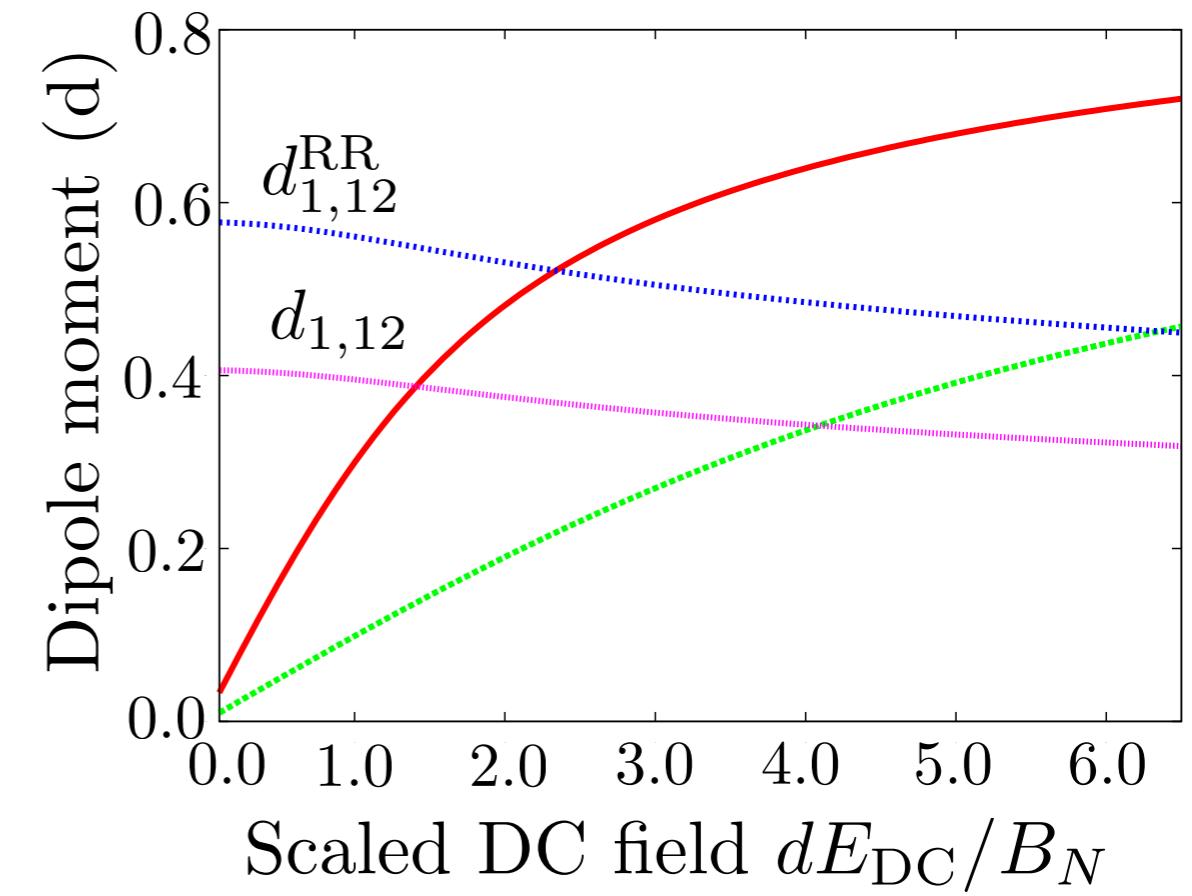
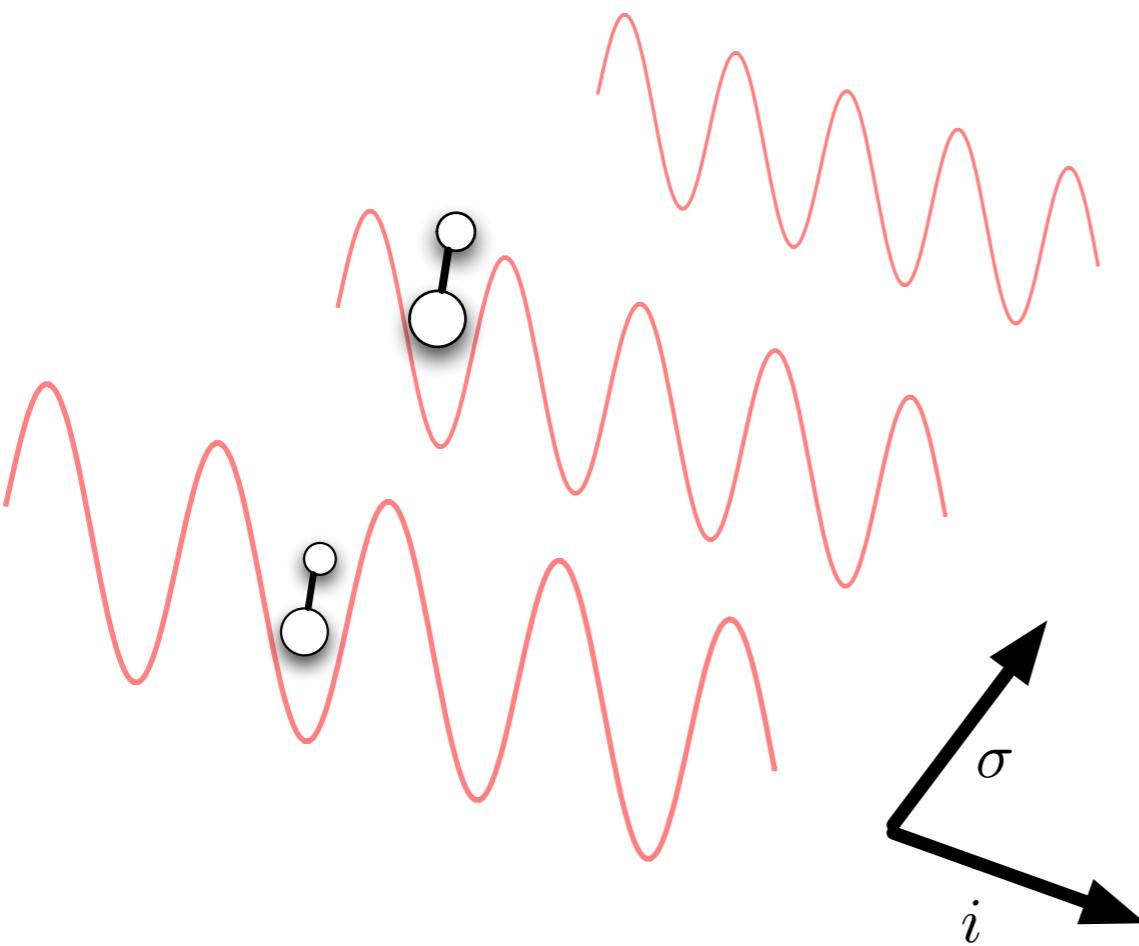
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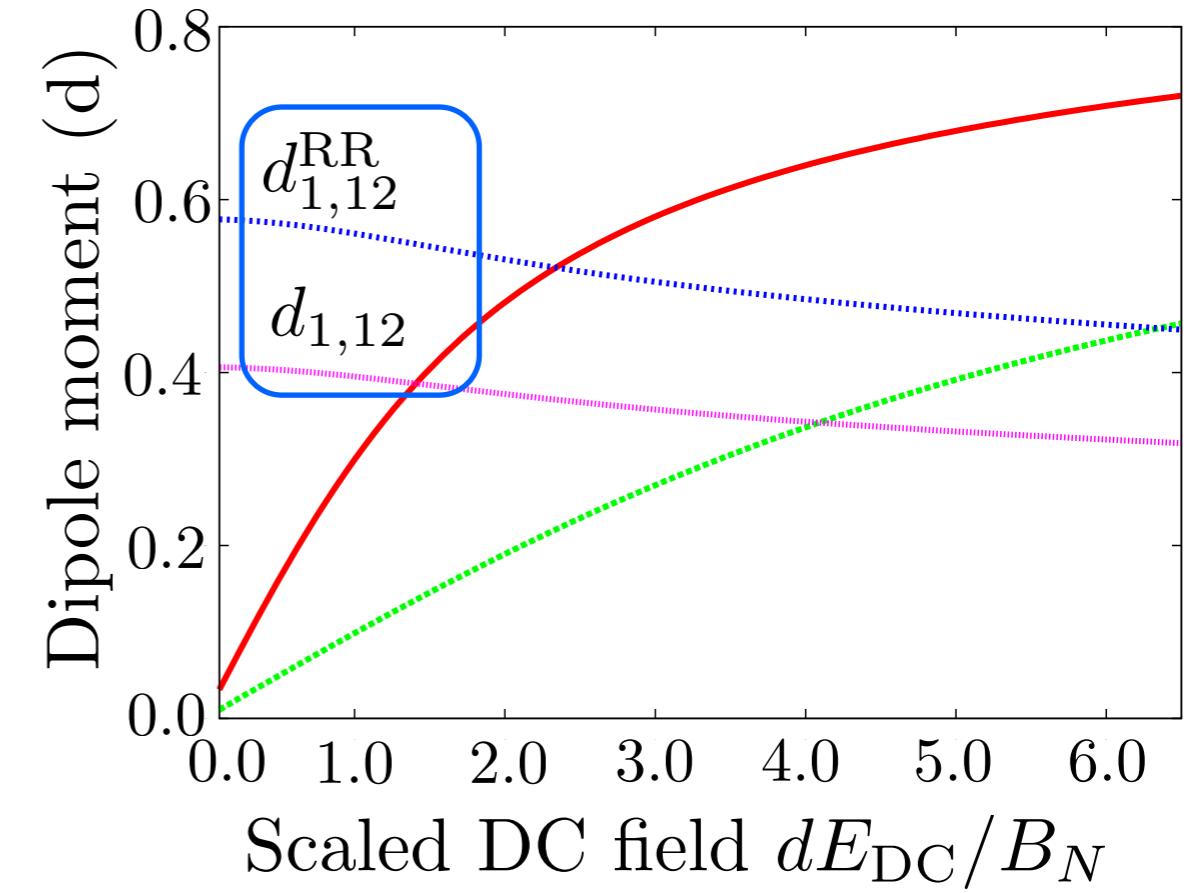
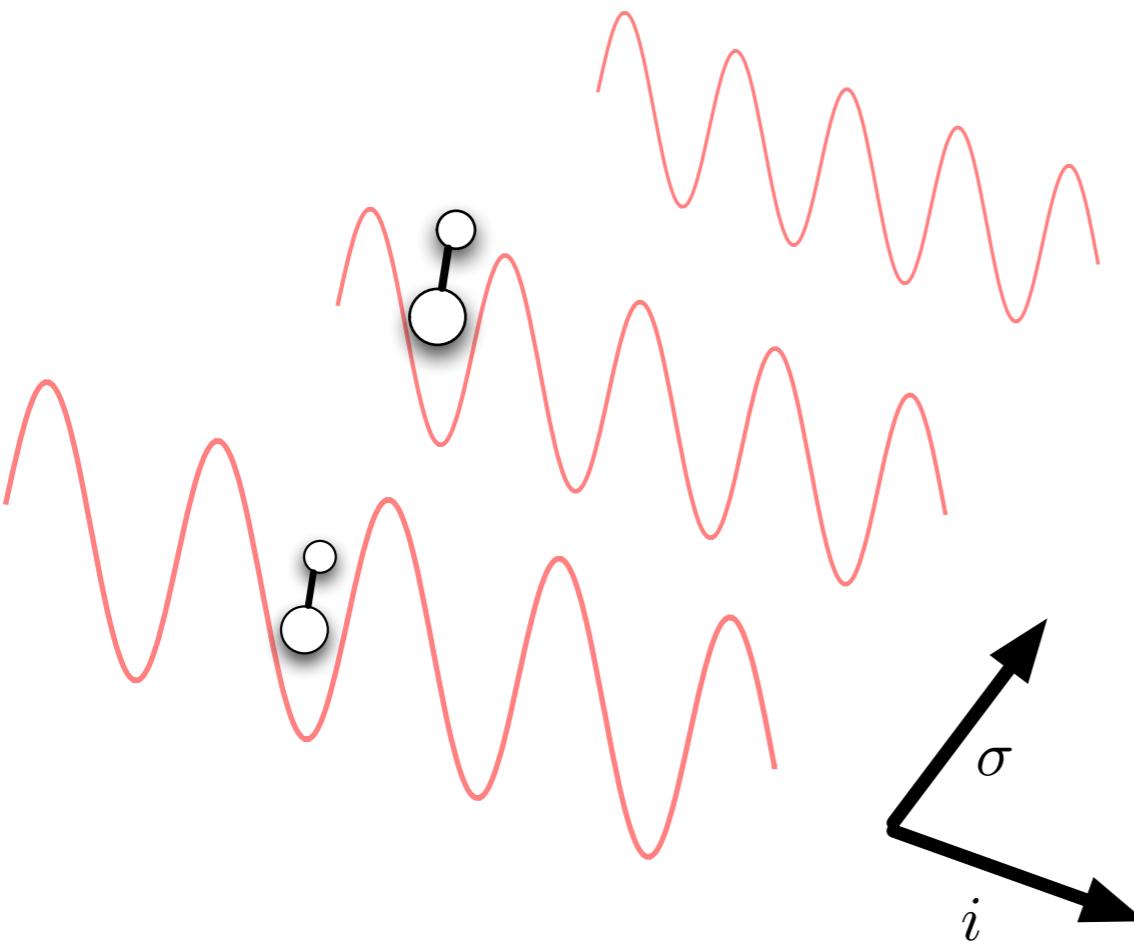
Exchange interaction



The molecular Hubbard Hamiltonian

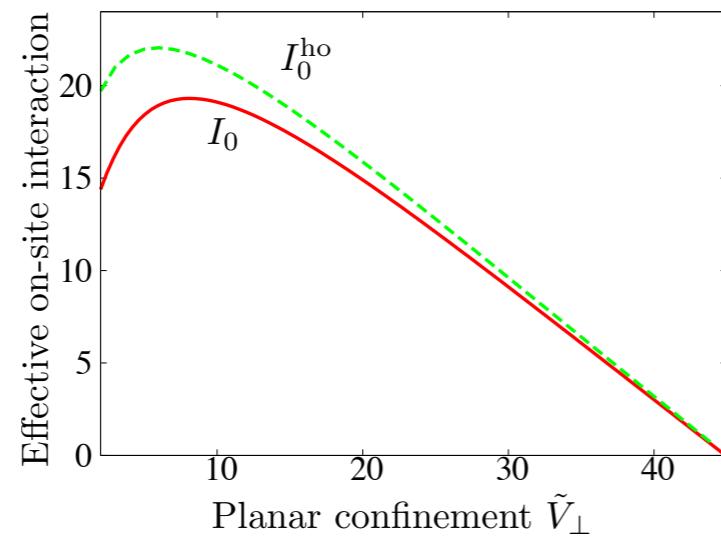
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Exchange interaction



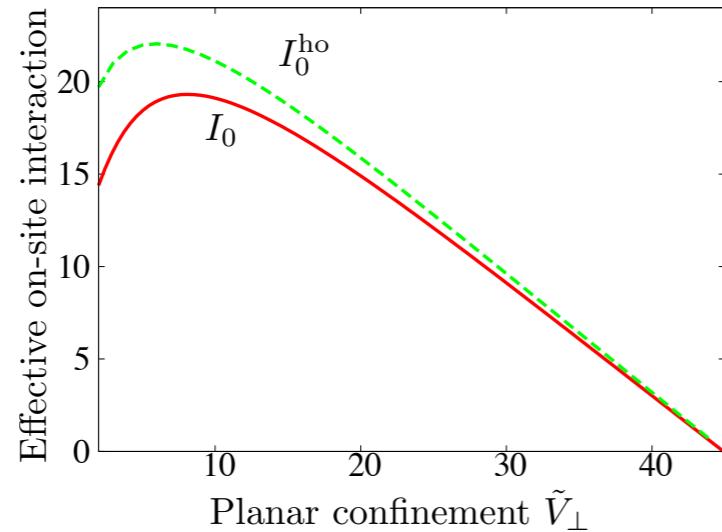
Confinement effects

- Harmonic oscillator: anisotropy affects energy



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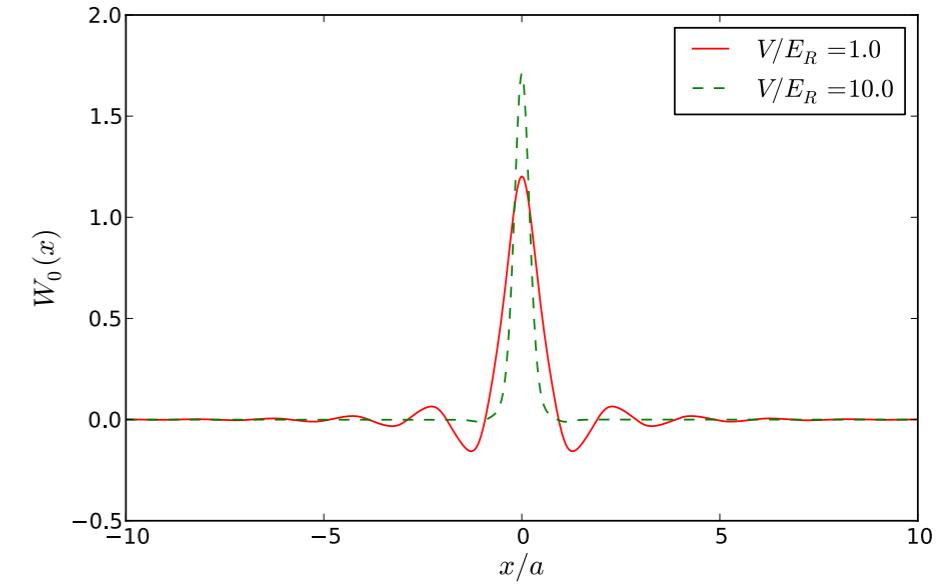
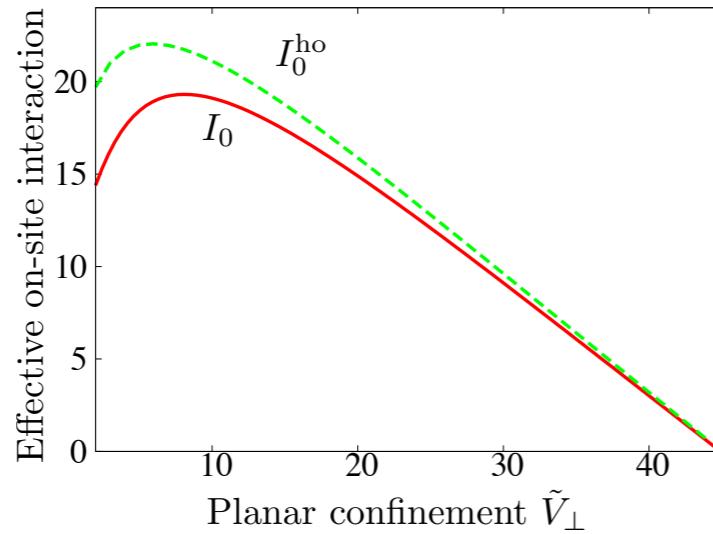
- Lattice: Long-ranged interactions also affected!
- Dipolar interactions are different in a lattice!

$$U_r \neq \frac{d^2}{a^3} \frac{1}{r^3} \quad U_r \sim \frac{d^2}{a^3} \left[\frac{1}{r^3} + a_e \exp(-b_e r) \right]$$

- 30-50% corrections at nearest-neighbor distance in quasi-low dimensional scenarios

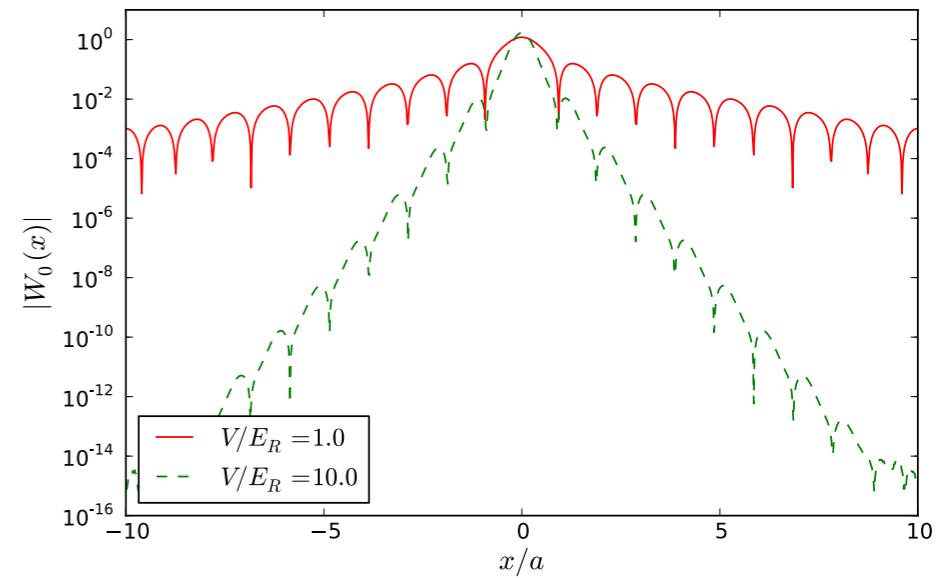
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Comparisons of molecules

	${}^6\text{Li}{}^{133}\text{Cs}$	${}^{23}\text{Na}{}^{40}\text{K}$	${}^{87}\text{Rb}{}^{133}\text{Cs}$	${}^{40}\text{K}{}^{87}\text{Rb}$	${}^6\text{Li}{}^{23}\text{Na}$
d/Debye	5.52	2.76	1.25	0.566	0.56
I_1	1	$3/2$	$3/2$	4	1
I_2	$7/2$	4	$7/2$	$3/2$	$3/2$
B_N/GHz	6.5202	2.8268	0.504	1.114	12.7355
$\frac{(eQq)_1}{\text{MHz}}$	3×10^{-4}	-0.167	-0.872	0.45	8×10^{-4}
$\frac{(eQq)_2}{\text{MHz}}$	0.187	0.796	0.051	-1.41	0.687
g_1	0.822	1.479	1.834	-0.324	0.822
g_2	0.738	-0.324	0.738	1.834	1.479
c_1/Hz	15.2	110.6	98.4	-24.1	83.2
c_2/Hz	3476.0	-90.9	194.1	420.1	802.2
c_3/Hz	53.1	-46.5	192.4	-48.2	196.2
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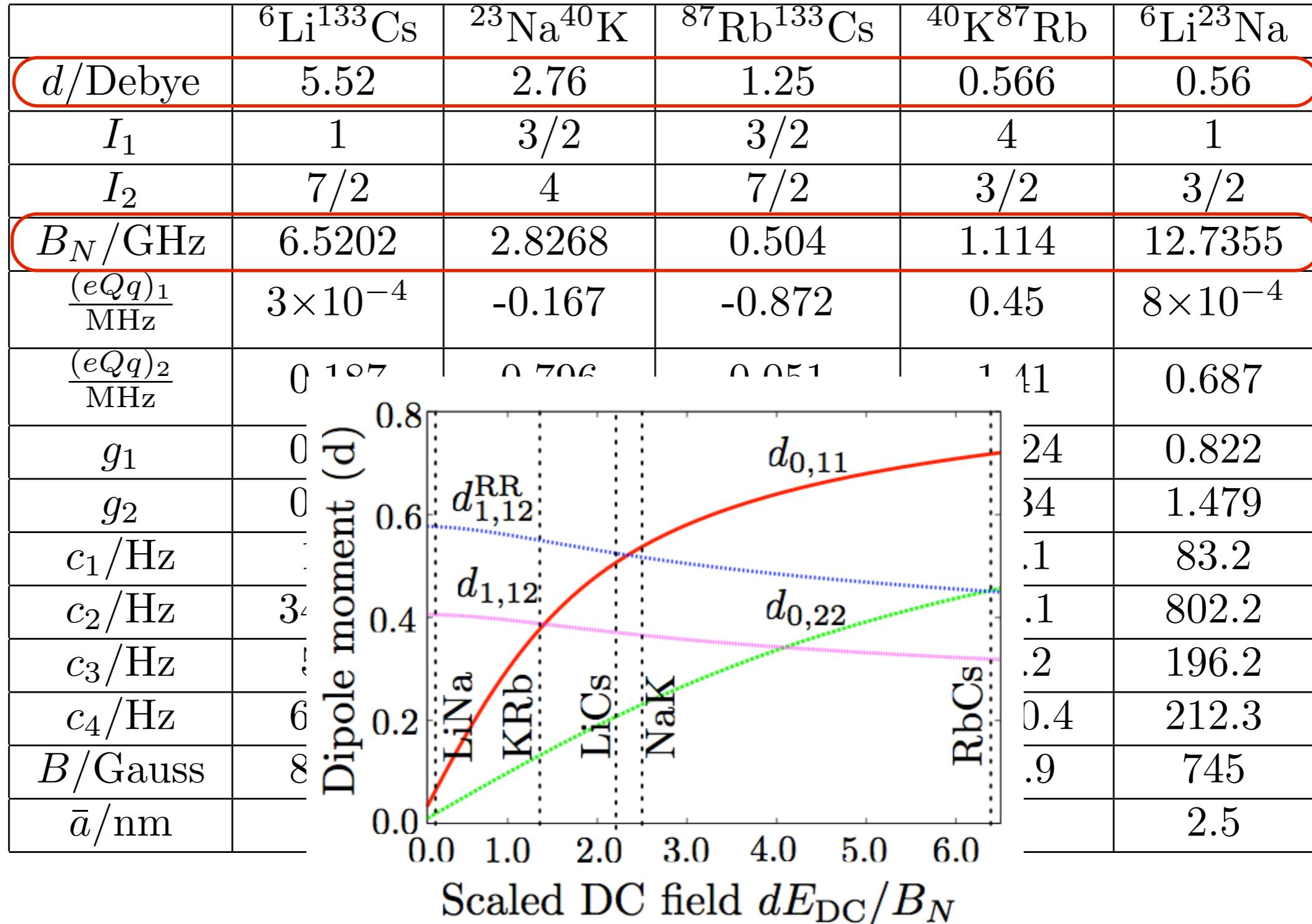
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Comparisons of molecules

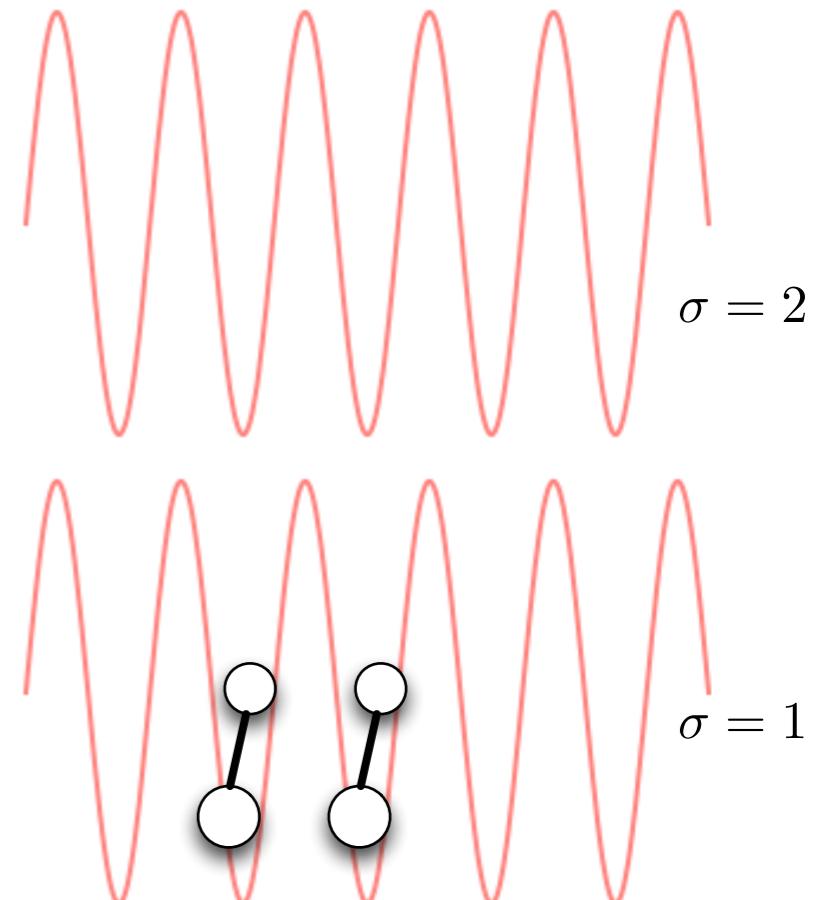
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I_1	1	3/2	3/2	4	1
I_2	7/2	4	7/2	3/2	3/2
B_N/GHz	6.5202	2.8268	0.504	1.114	12.7355
$\frac{(eQq)_1}{\text{MHz}}$	3×10^{-4}	-0.167	-0.872	0.45	8×10^{-4}
$\frac{(eQq)_2}{\text{MHz}}$	0.187	0.796	0.051	-1.41	0.687
g_1	0.822	1.479	1.834	-0.324	0.822
g_2	0.738	-0.324	0.738	1.834	1.479
c_1/Hz	15.2	110.6	98.4	-24.1	83.2
c_2/Hz	3476.0	-90.9	194.1	420.1	802.2
c_3/Hz	53.1	-46.5	192.4	-48.2	196.2
c_4/Hz	620.8	-466.2	17345.4	-2030.4	212.3
B/Gauss	843.5	139.7	181.7	545.9	745
\bar{a}/nm	25	23.5	23.5	6	2.5

Many-body case studies

	${}^6\text{Li}{}^{133}\text{Cs}$	${}^{23}\text{Na}{}^{40}\text{K}$	${}^{87}\text{Rb}{}^{133}\text{Cs}$	${}^{40}\text{K}{}^{87}\text{Rb}$	${}^6\text{Li}{}^{23}\text{Na}$
$E_{\text{DC}} = 1 \text{ kV/cm}$					
$U_{1,11}$ (kHz)	0.70	0.218	0.253	0.002	2×10^{-5}
$U_{1,22}/U_{1,11}$	0.084	0.085	0.108	0.082	0.08
$U_{1,12}/U_{1,11}$	0.28	0.29	0.32	0.28	0.28
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 $U_{r,\sigma\sigma'}$
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 $E_{r,\sigma_1\sigma_2\sigma'_2\sigma'_1}$

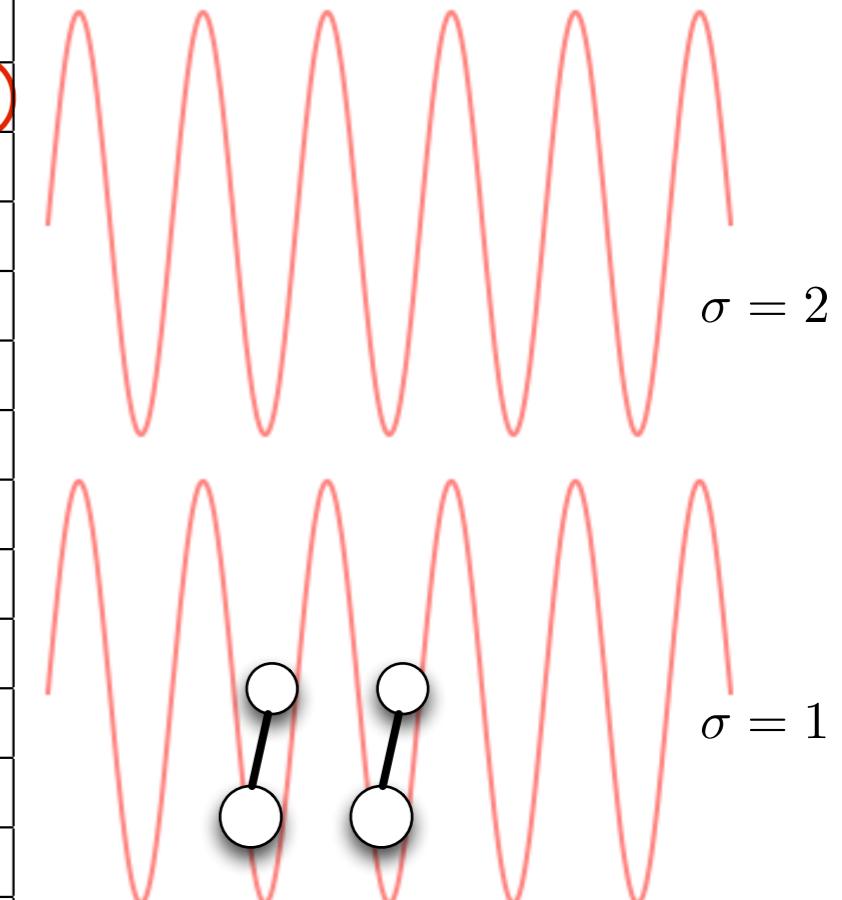


Many-body case studies

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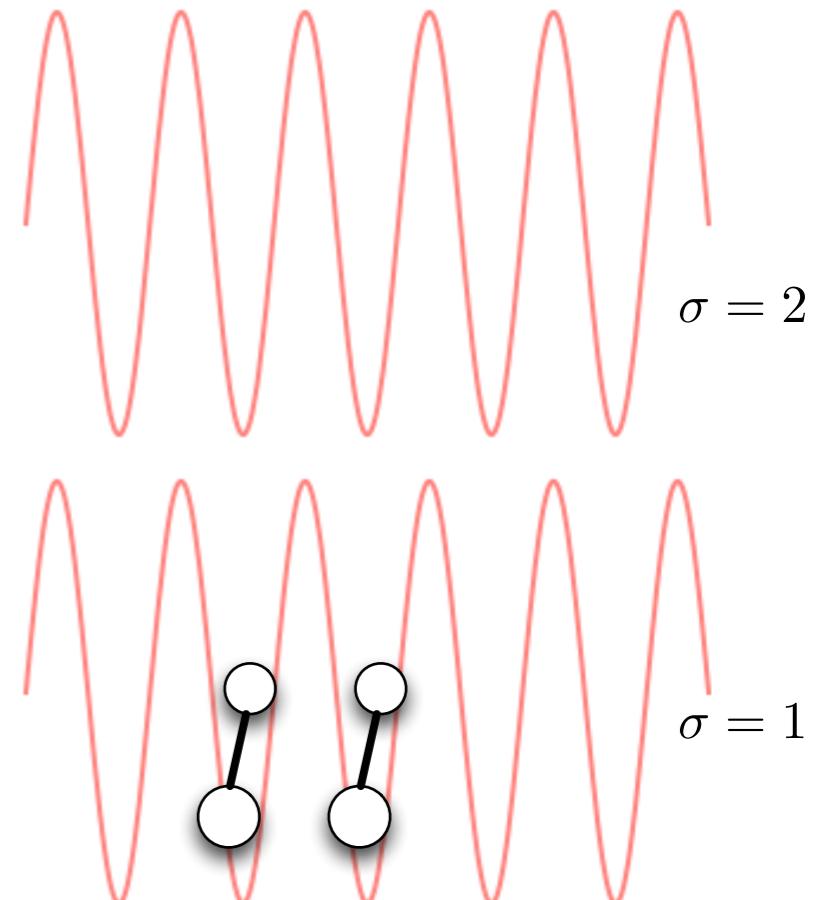


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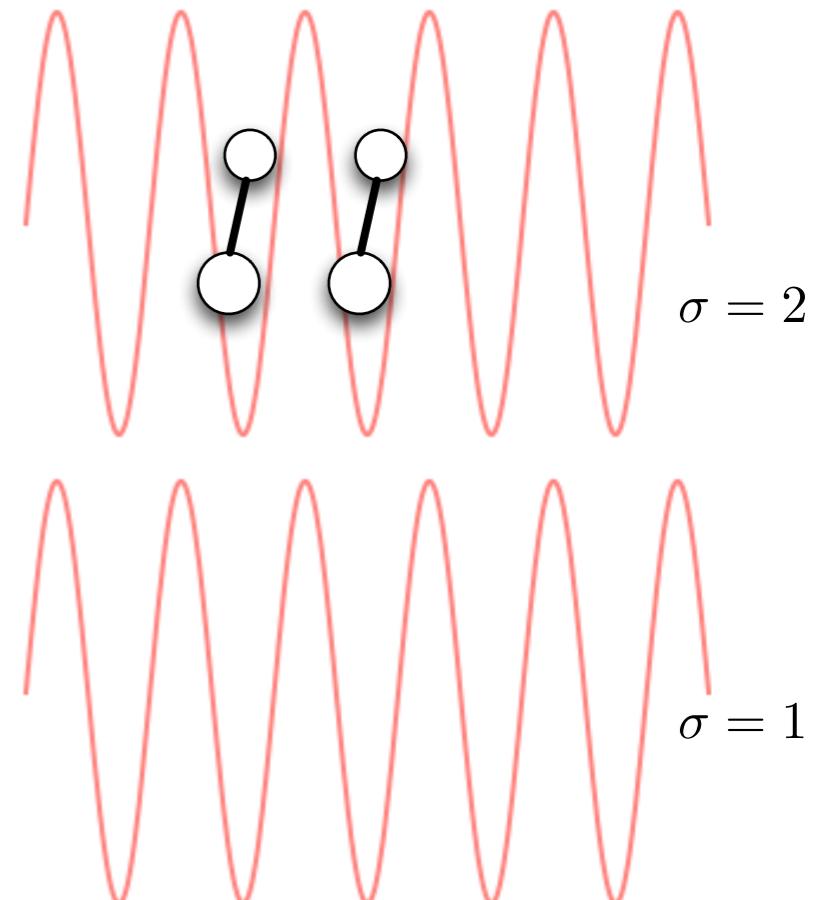


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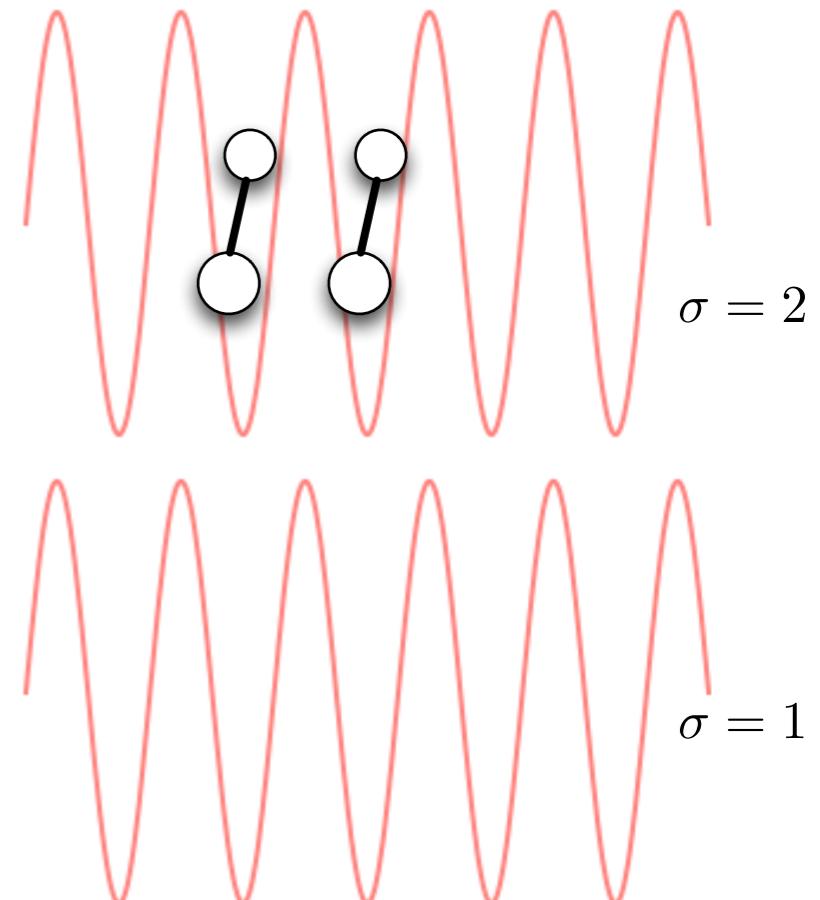


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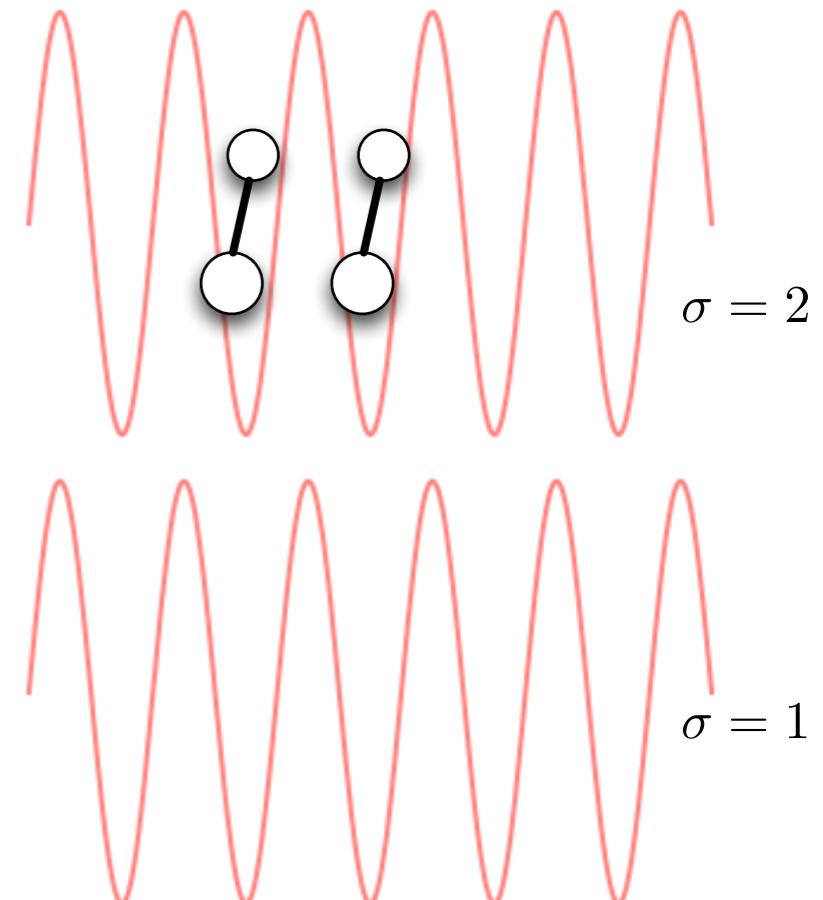


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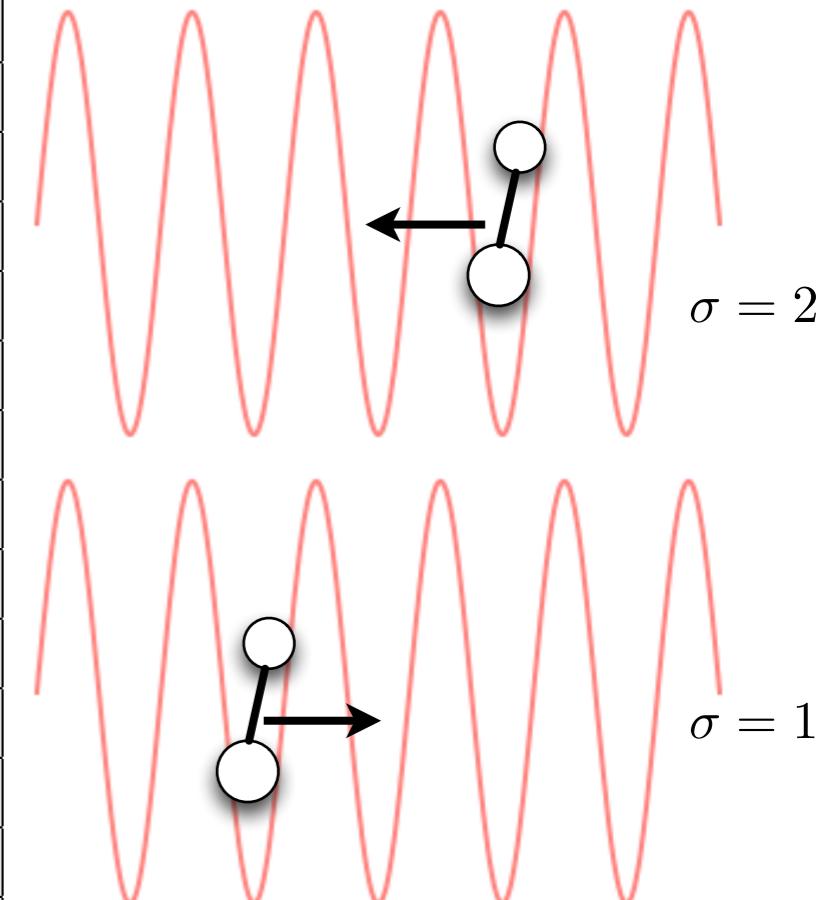


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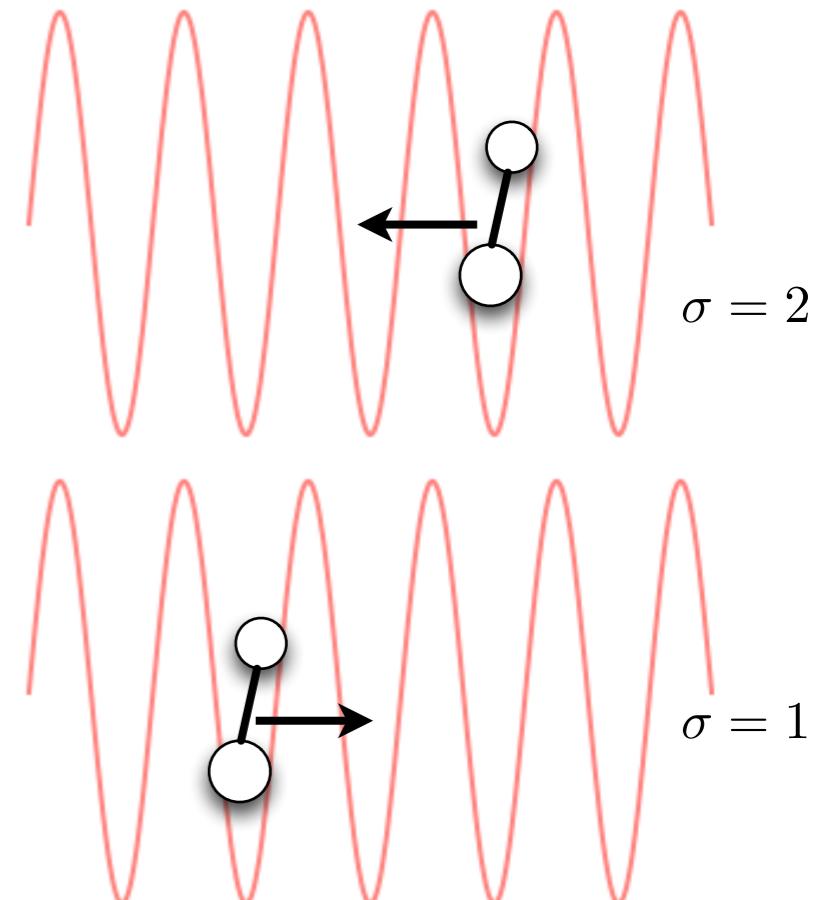


Many-body case studies

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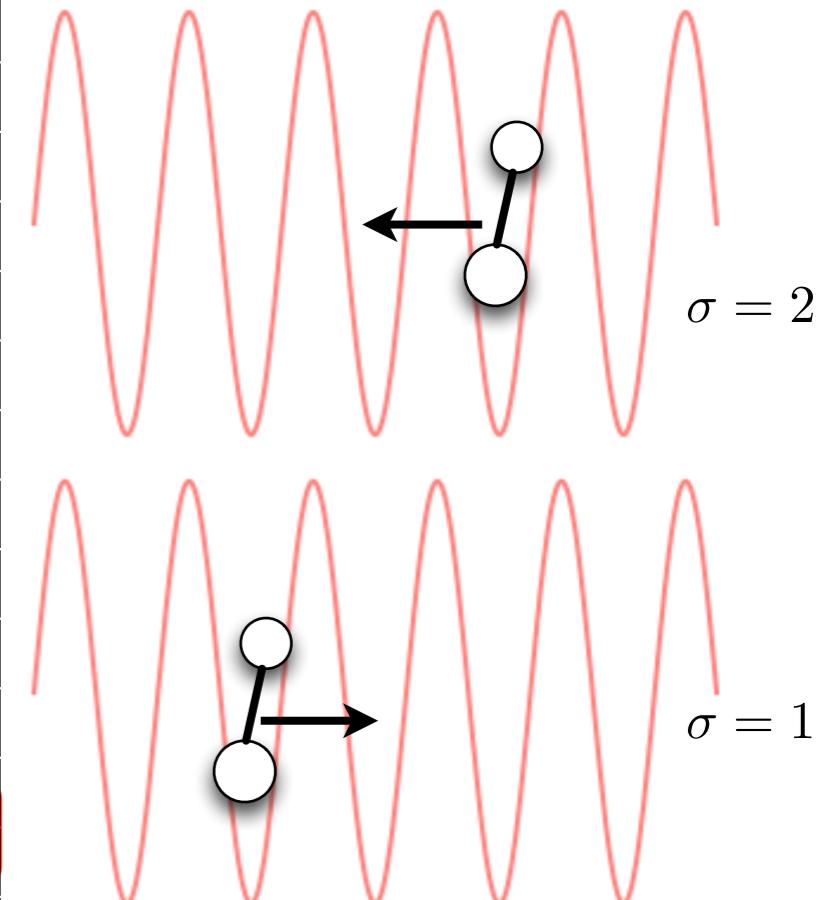


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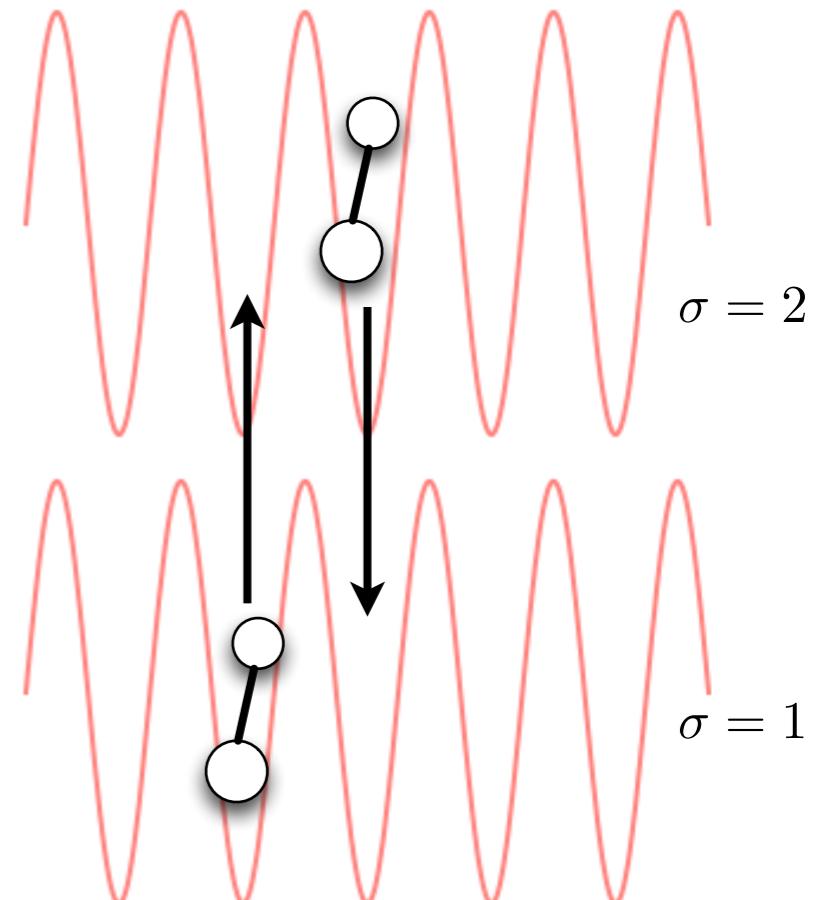


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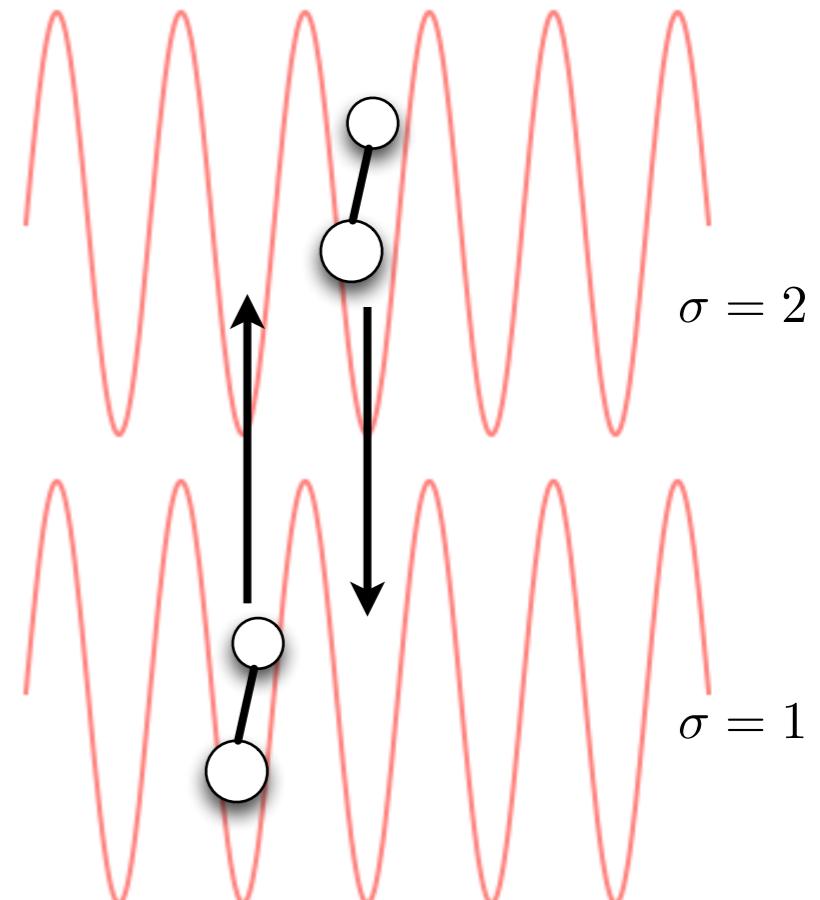


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$t_{1,1}/U_{1,11}$	0.026	0.20	0.15	5.13	2×10^3
$t_{1,2}/U_{1,22}$	0.068	0.48	0.14	19.75	1.2×10^5
$E_{1,1212}/U_{1,11}$	0.519	0.43	0.18	1.05	114

Quasi-1D, 2 internal states
 r , distance
 σ , internal state

Tunneling
 $t_{r,\sigma}$
 Direct dipole-dipole interaction
 $U_{r,\sigma\sigma'}$
 Exchange dipole-dipole interaction
 $E_{r,\sigma_1\sigma_2\sigma'_2\sigma'_1}$

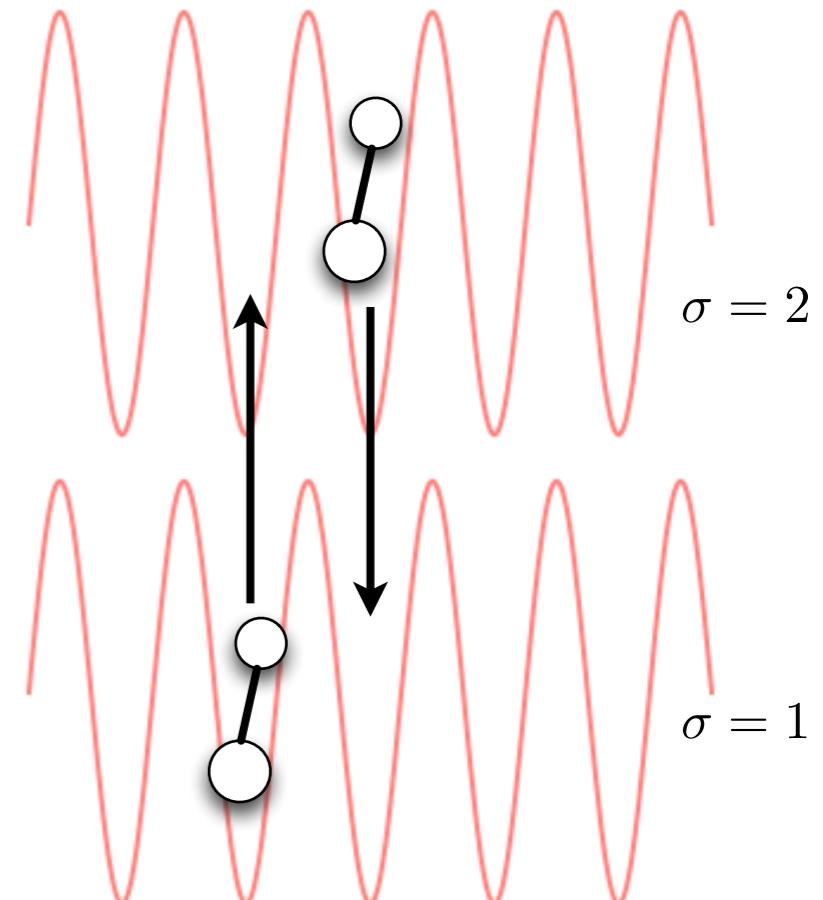


Many-body case studies

	${}^6\text{Li}{}^{133}\text{Cs}$	${}^{23}\text{Na}{}^{40}\text{K}$	${}^{87}\text{Rb}{}^{133}\text{Cs}$	${}^{40}\text{K}{}^{87}\text{Rb}$	${}^6\text{Li}{}^{23}\text{Na}$
$E_{\text{DC}} = 1 \text{ kV/cm}$					
$U_{1,11}$ (kHz)	0.70	0.218	0.253	0.002	2×10^{-5}
$U_{1,22}/U_{1,11}$	0.084	0.085	0.108	0.082	0.08
$U_{1,12}/U_{1,11}$	0.28	0.29	0.32	0.28	0.28
$t_{1,1}/U_{1,11}$	0.33	2.36	0.58	93.74	5×10^5
$t_{1,2}/U_{1,22}$	1.76	12.37	2.25	508.9	3×10^6
$E_{1,1212}/U_{1,11}$	7.93	6.05	0.99	20.97	2×10^3
$E_{\text{DC}} = 5 \text{ kV/cm}$					
$U_{1,11}$ (kHz)	8.97	2.48	0.970	0.049	5×10^{-4}
$U_{1,22}/U_{1,11}$	0.14	0.16	0.35	0.11	0.08
$U_{1,12}/U_{1,11}$	0.38	0.40	0.59	0.32	0.28
$t_{1,1}/U_{1,11}$	0.026	0.20	0.15	5.13	2×10^3
$t_{1,2}/U_{1,22}$	0.068	0.48	0.14	19.75	1.2×10^5
$E_{1,1212}/U_{1,11}$	0.519	0.43	0.18	1.05	114

Quasi-1D, 2 internal states
 r , distance
 σ , internal state

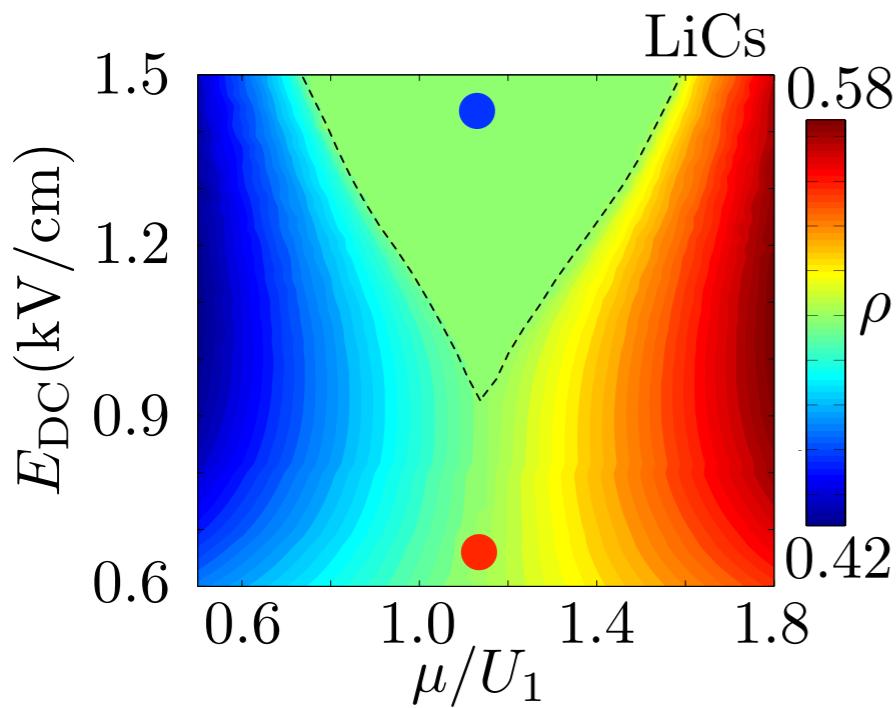
Tunneling
 $t_{r,\sigma}$
 Direct dipole-dipole interaction
 $U_{r,\sigma\sigma'}$
 Exchange dipole-dipole interaction
 $E_{r,\sigma_1\sigma_2\sigma'_2\sigma'_1}$



Numerical Methods

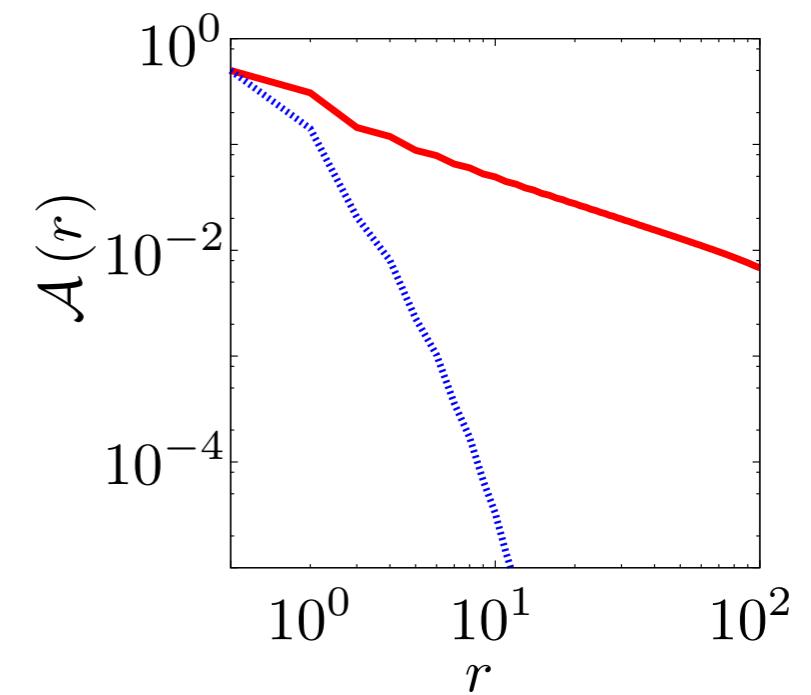
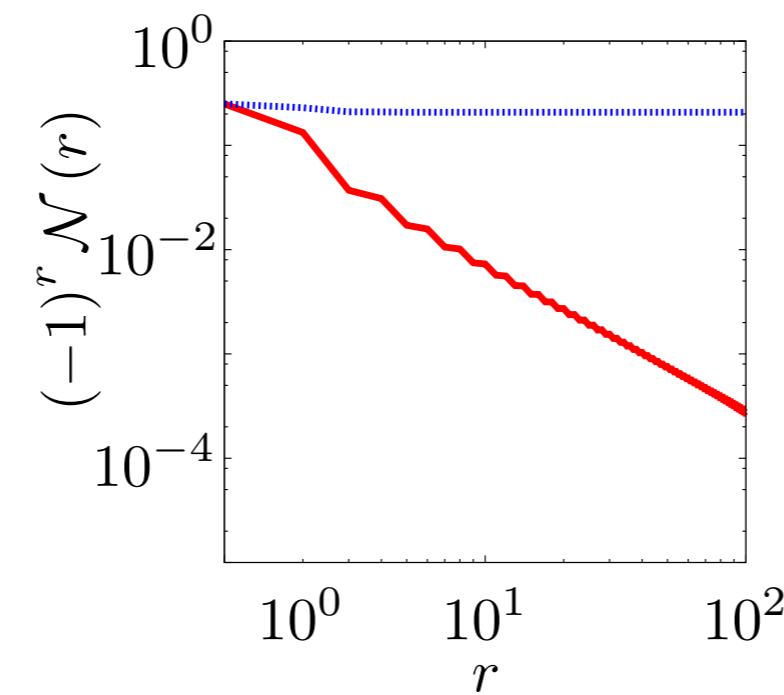
- Assume many-body state on infinite lattice has periodic structure
- Matrix product state ansatz for “unit cell”
- Recursive optimization of unit cell.
- Single convergence parameter; entanglement cutoff
- Straightforward extensions to long-range interactions
- Finite-system algorithms for excited states, time evolution

Many-body case studies

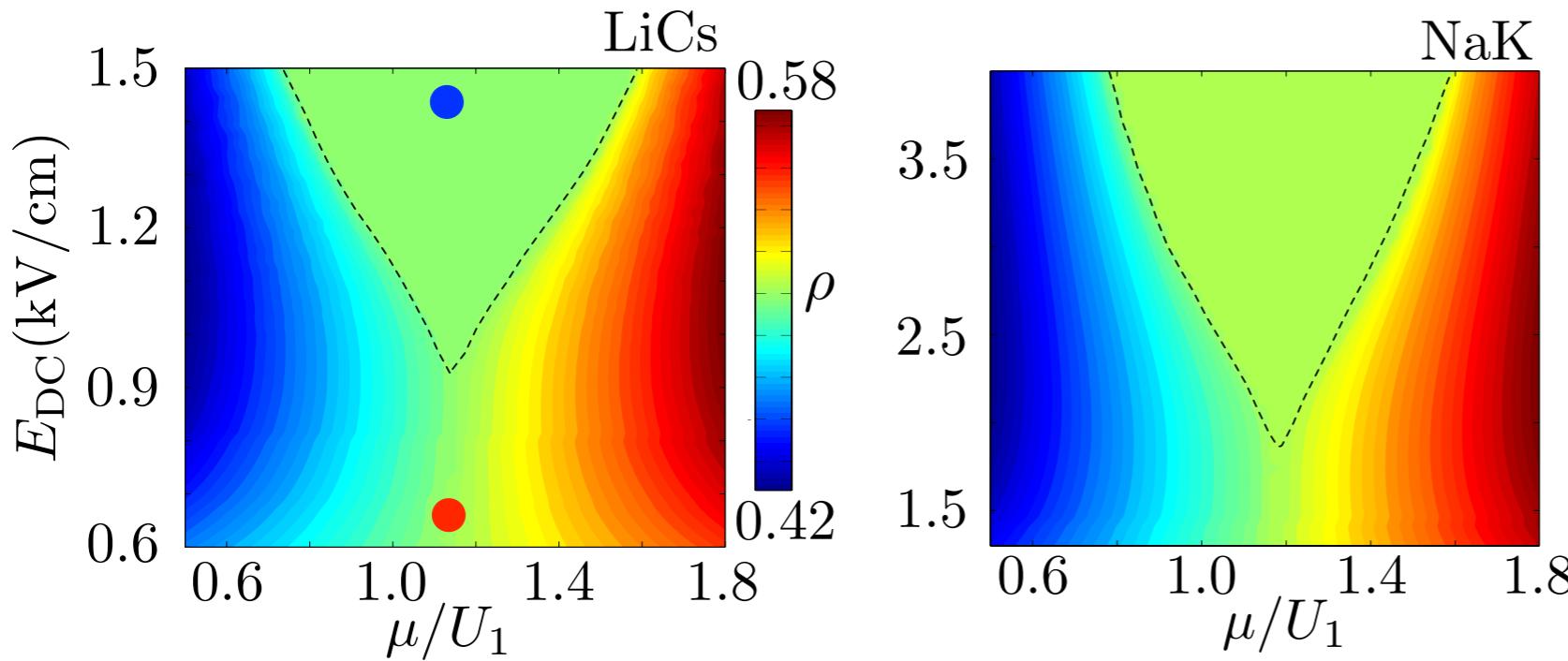


$$\mathcal{A}(r) = \langle \hat{a}_0^\dagger \hat{a}_r \rangle$$

$$\mathcal{N}(r) = \langle (\hat{n}_0 - \rho)(\hat{n}_r - \rho) \rangle$$

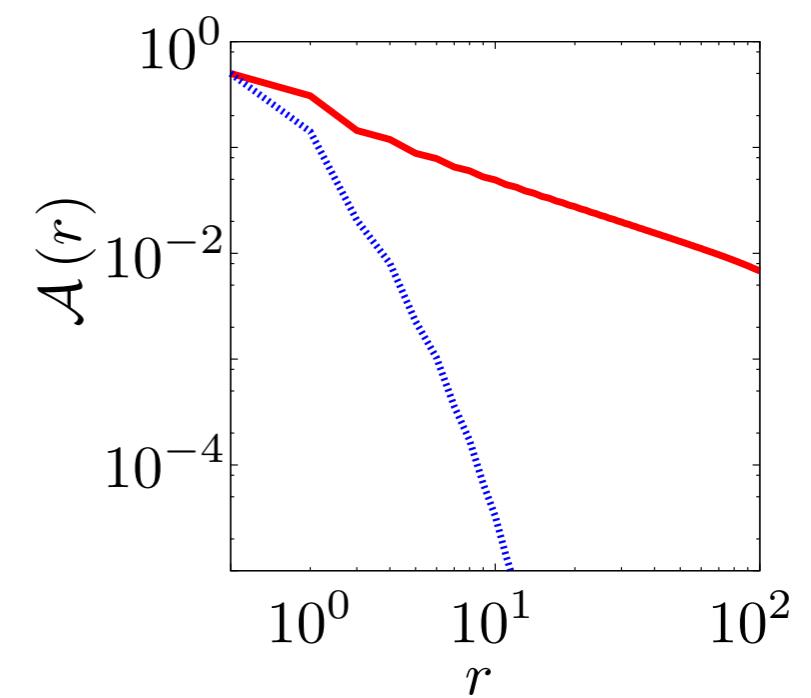
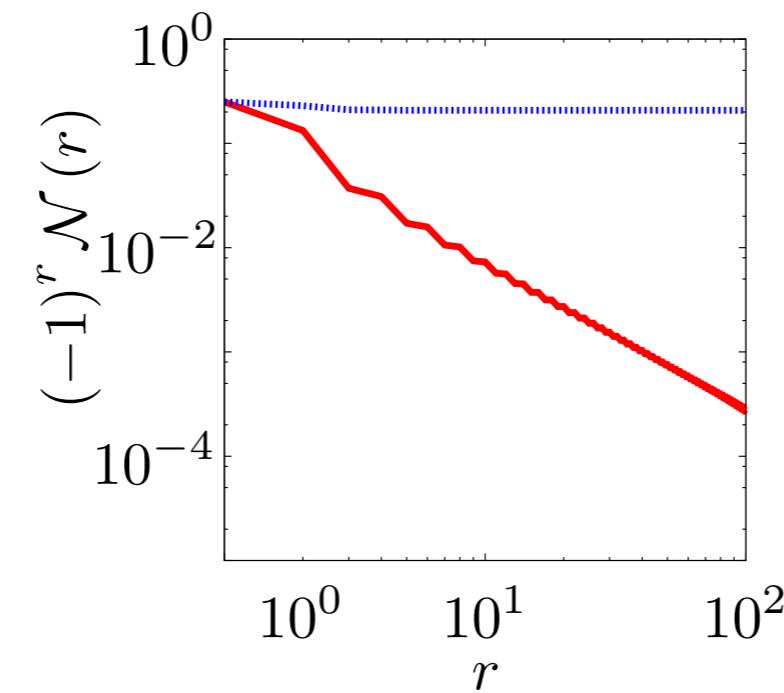


Many-body case studies

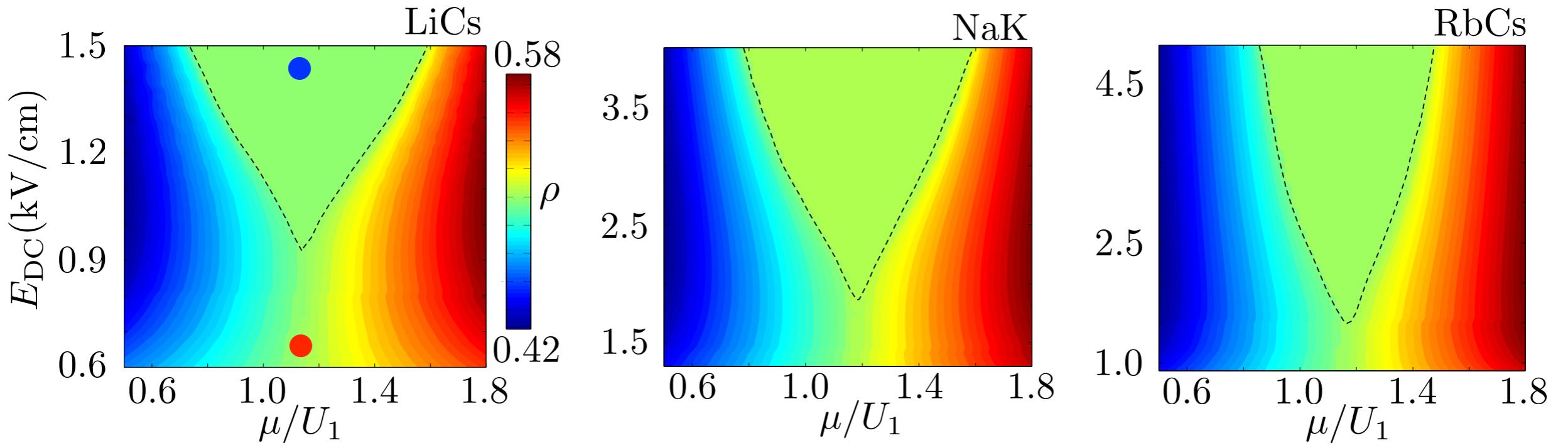


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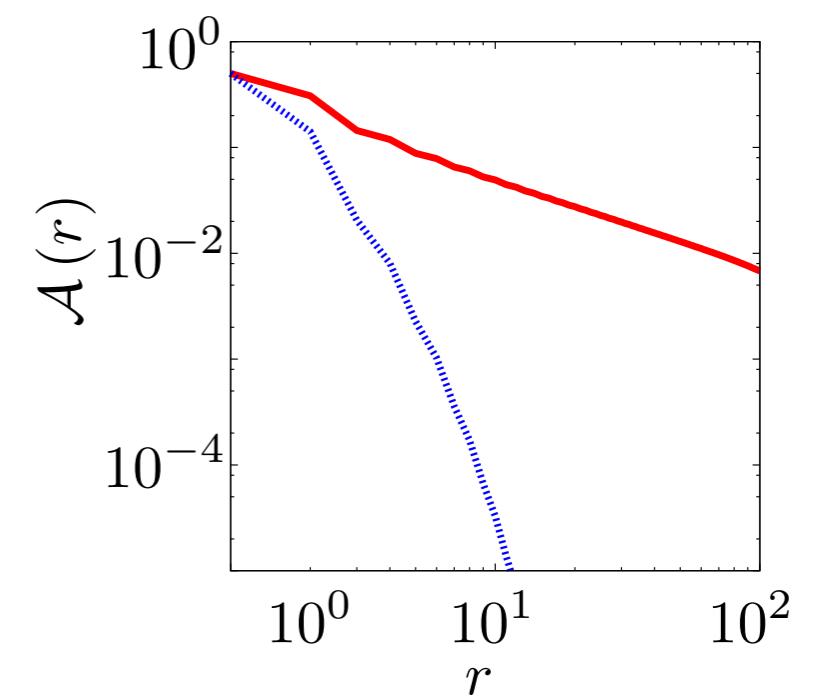
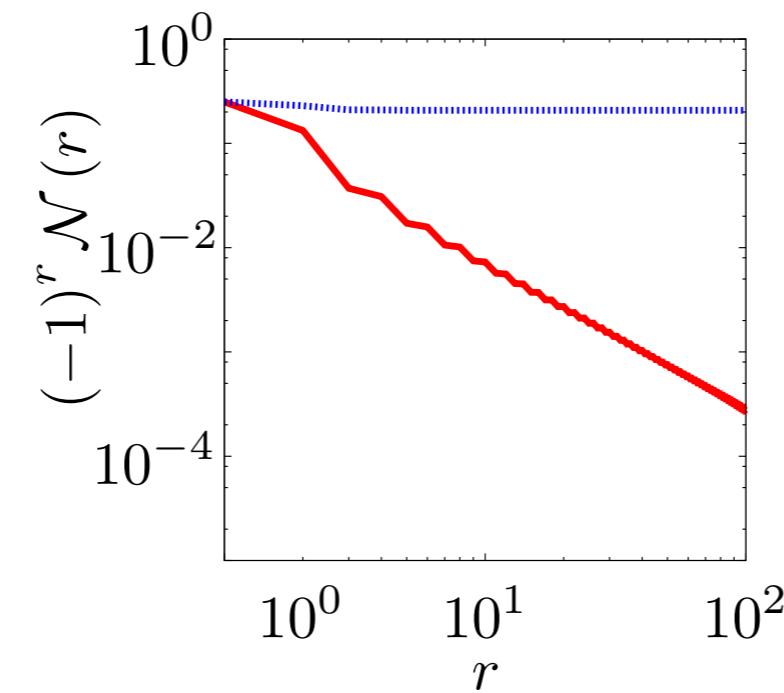


Many-body case studies

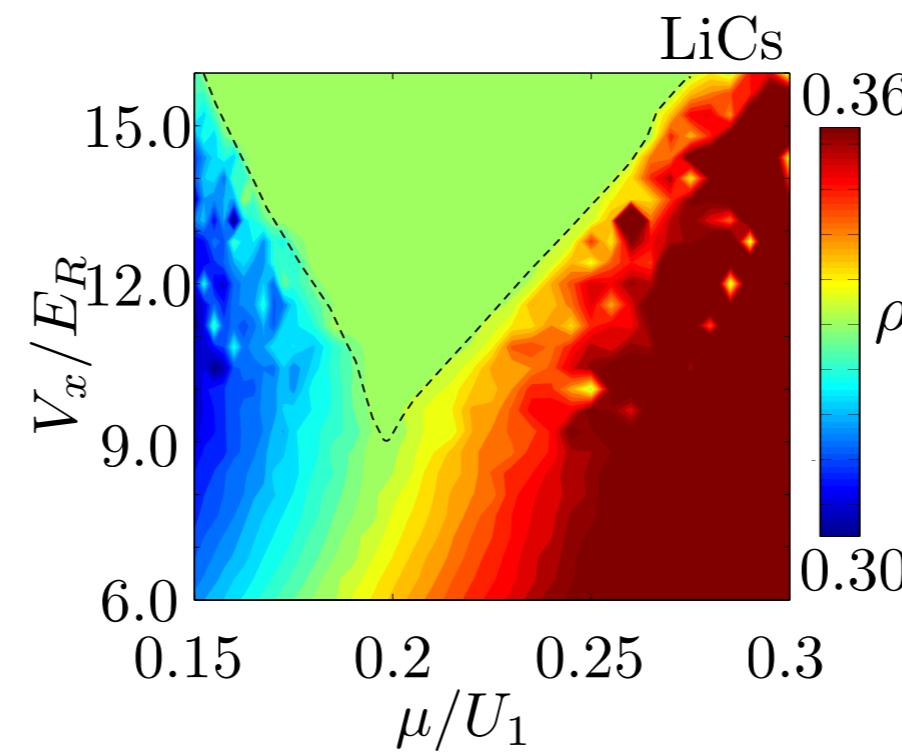
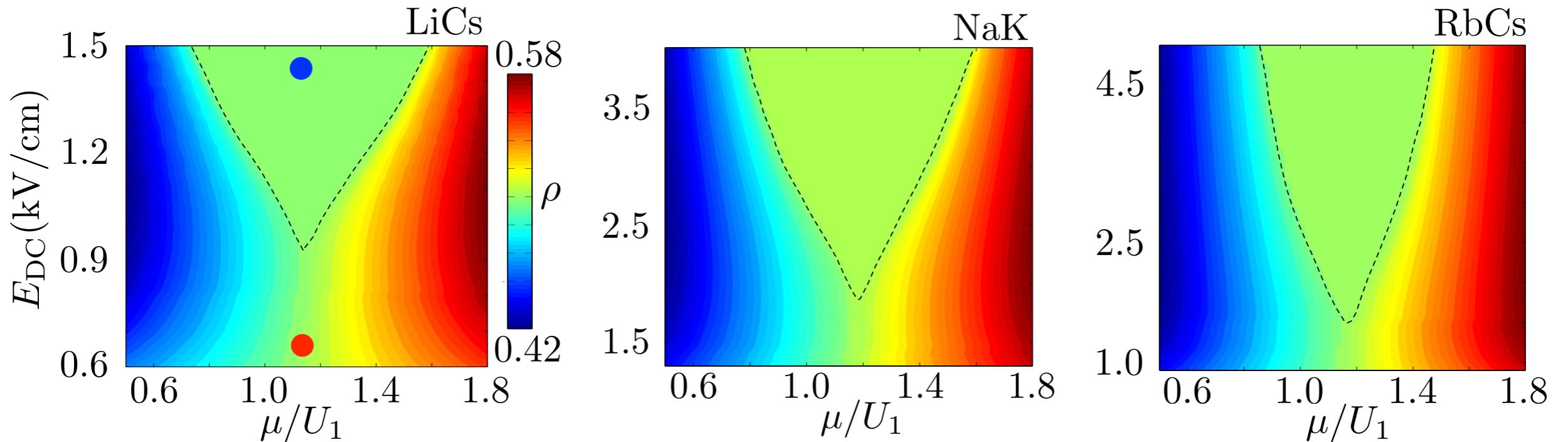


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$$\mathcal{N}(r) = \langle (\hat{n}_0 - \rho)(\hat{n}_r - \rho) \rangle$$



Many-body case studies



Summary

- Developed Molecular Hubbard Hamiltonian (MHH)
 - Treated¹ Σ case in detail, see also work of Gorshkov et al.
 - Different molecules access different regimes of MHH
- Wide array of experimental controls
 - Optical lattice intensity and polarization: tunneling and internal state character
 - Strength of DC field: ratio of direct and exchange dipolar interactions
 - AC field: number and timescale of internal states involved dynamically
 - B field: Coupling between hyperfine and rotation
 - Anisotropy in lattice confinement=> modified dipolar interactions
 - Species can also be considered an experimental control!
- New numerical many-body techniques
 - Able to handle long-range Hamiltonians, wide array of interactions
 - Used to explore phase diagrams of different molecules
 - Plan to open source code-contact us!

- MLW and L. D. Carr, New J. Phys. 11 055027 (2009)
- MLW and L. D. Carr Phys. Rev. A 82 013611 (2010)
- MLW, E. Bekaroglu, and L. D. Carr arxiv:1212.3042
- Confinement effects: MLW and L. D. Carr arxiv:1303.1230
- Algorithmic developments: MLW and L. D. Carr, New J. Phys. 14 125015 (2012)
- 4 posters in common room

Thank You!