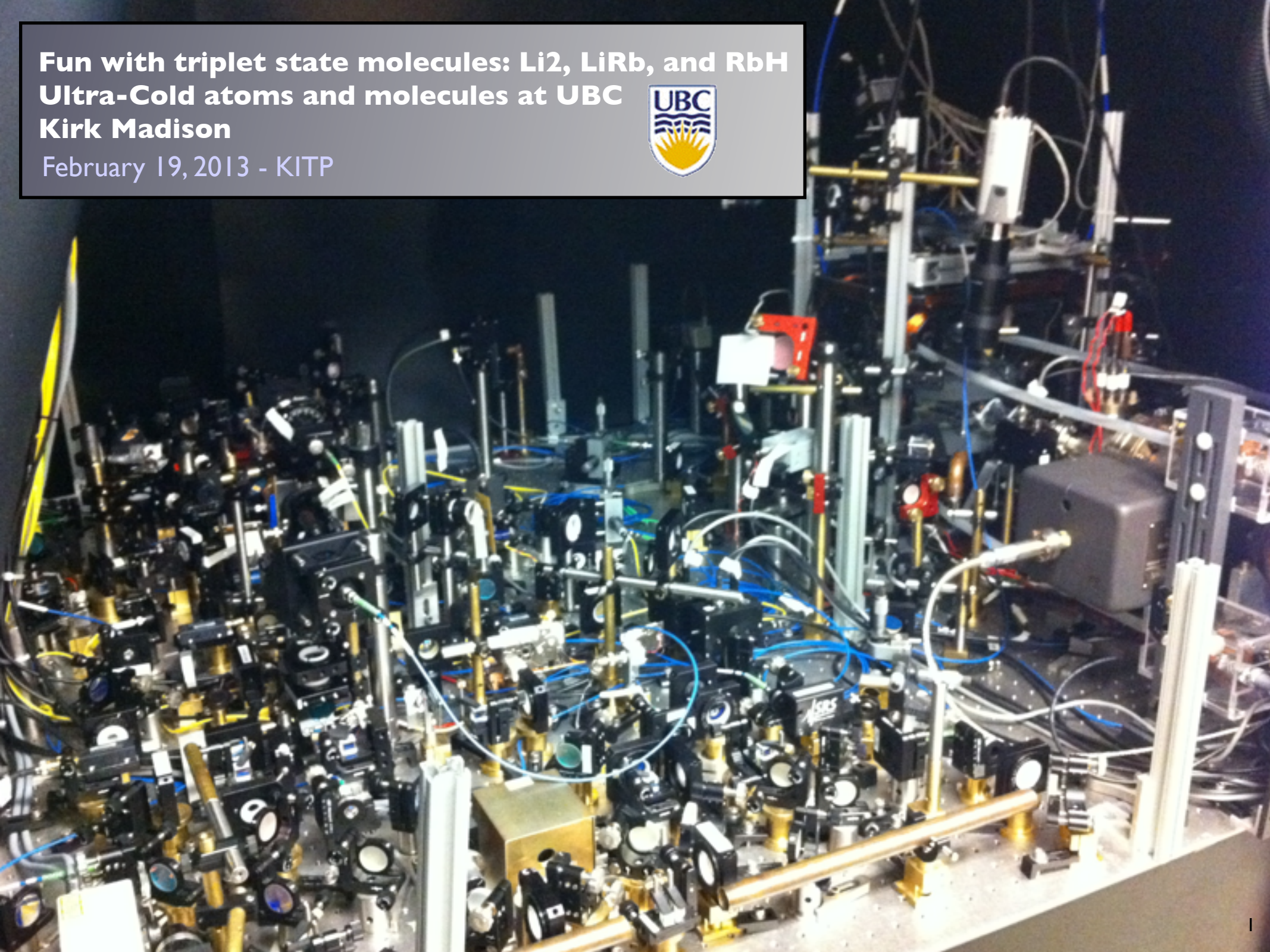


Fun with triplet state molecules: Li₂, LiRb, and RbH
Ultra-Cold atoms and molecules at UBC
Kirk Madison

February 19, 2013 - KITP



Li+Rb mixtures

Will
Gunton



Mariusz
Semczuk



Theory support for Li_2 from Xuan Li¹ and Nike Dattani²

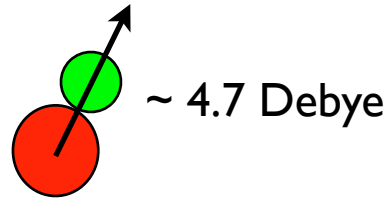
¹Chemical Science Division, Lawrence Berkeley National Laboratory, Berkeley, USA

²Department of Chemistry, University of Oxford, Oxford, UK



Long term goal: production of dense, ultracold ensembles of polar molecules

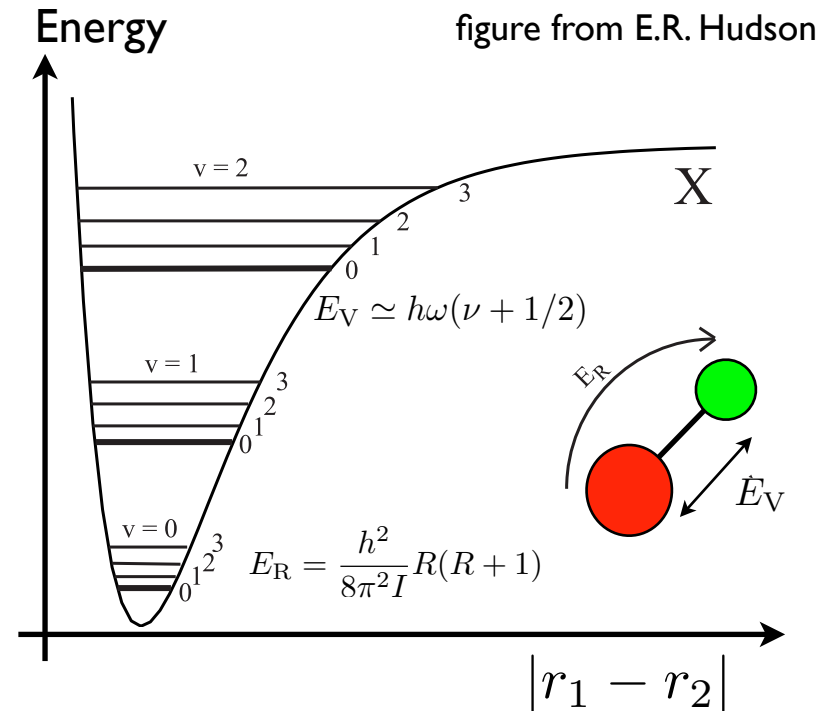
RbLi is a polar molecule



Polar molecules: scientific motivation

Additional **features** not available with atoms:

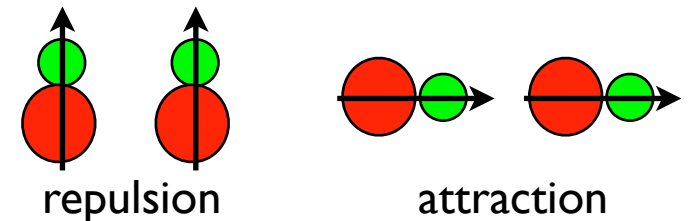
Rich internal structure:
- *rotational, vibrational*



Dipolar interactions:

- *large and long range* $\sim 1/r^3$
- *angular dependence (anisotropic)*

$$U_{dd}^{\text{electric}} = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{d}_1 \cdot \mathbf{d}_2 - 3(\mathbf{d}_1 \cdot \hat{r})(\mathbf{d}_2 \cdot \hat{r})}{r^3}$$



... and zero at the 'magic angle'

Polar molecules: scientific motivation

Dipolar interactions - *how strong?*

* fully polarized

Inter-particle interactions between RbLi dimers (4.7 Debye)*

~1000x larger than mean field interaction in a BEC

~10,000x larger than magnetic dipolar interactions in Cr

Dipolar interactions - *at what range?*

532 nm
optical
lattice

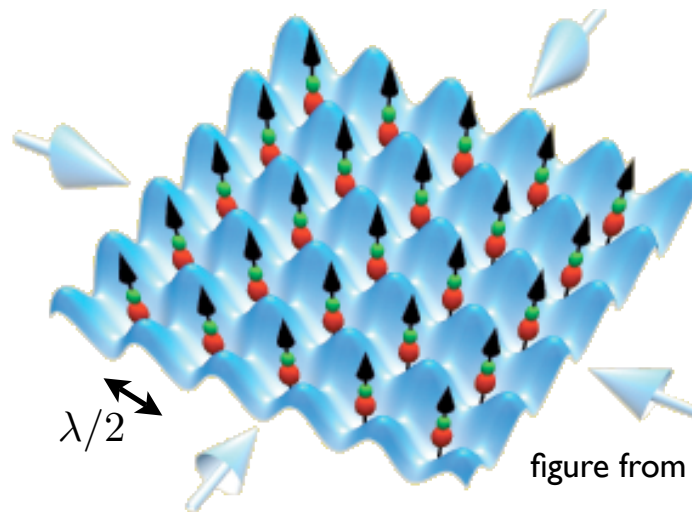


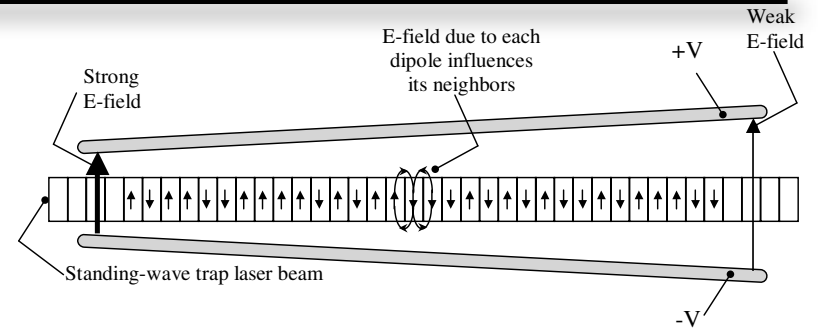
figure from Nature **2**, 341 (2006)

$$U_{\text{dd}} \sim \frac{(4.7 \text{ Debye})^2}{4\pi\epsilon_0 r^3} \Big|_{r=0.266\mu\text{m}} = k_B \times 8.5 \mu\text{K} = h \times 177 \text{ kHz}$$

Polar molecules: scientific motivation

few and many body QM

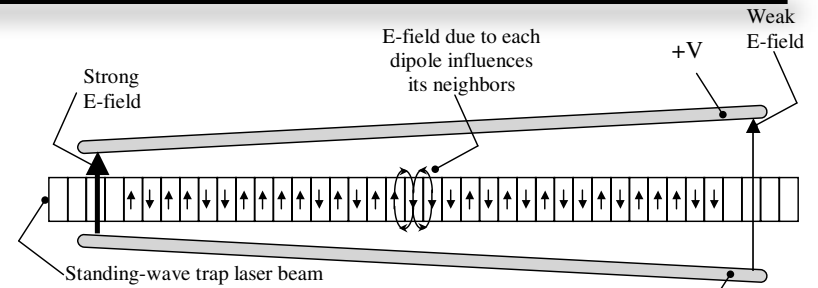
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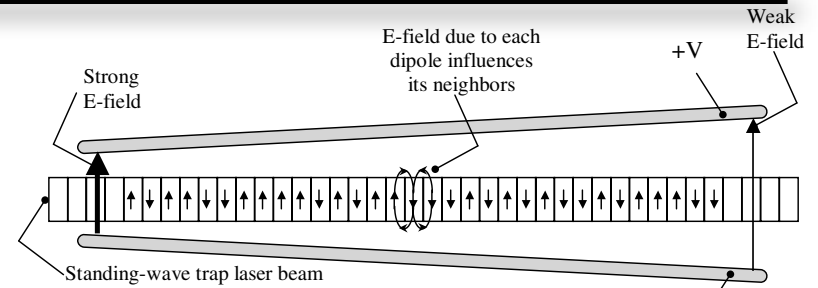
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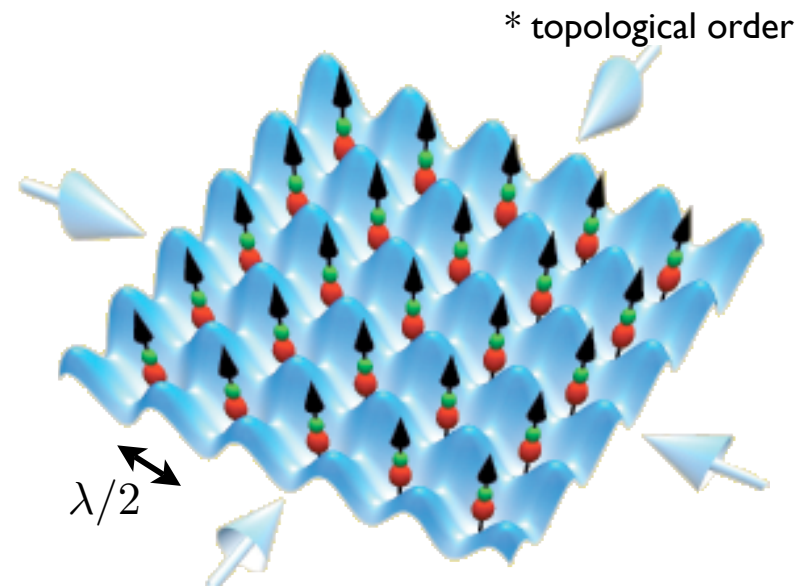
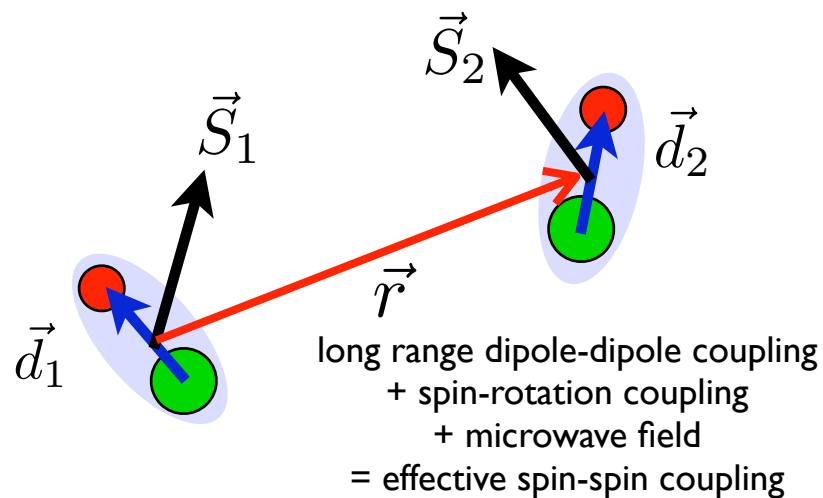
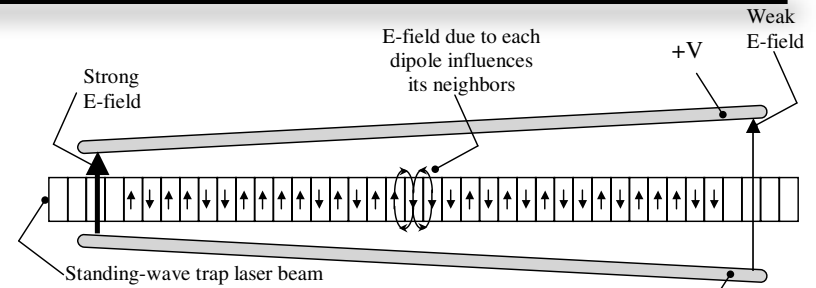
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“Inelastic collisions and chemical reactions of cold molecules in external fields”



Polar molecules: scientific motivation

few and many body QM

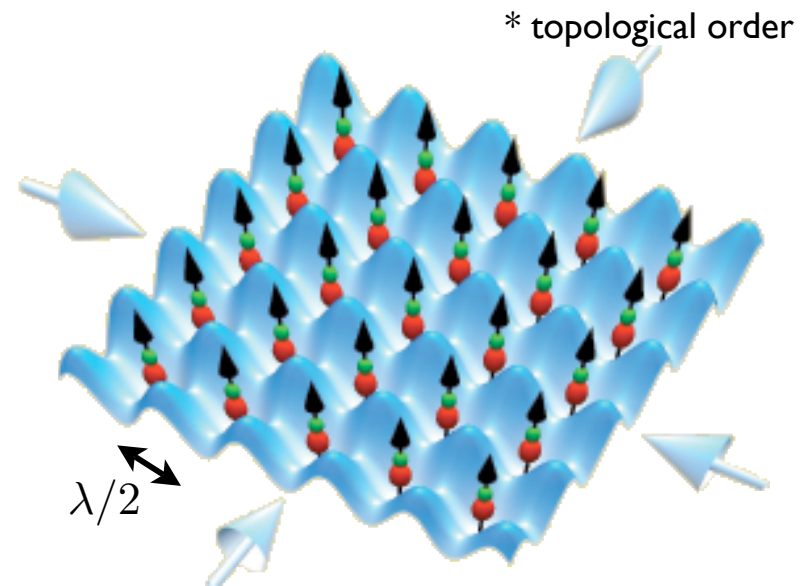
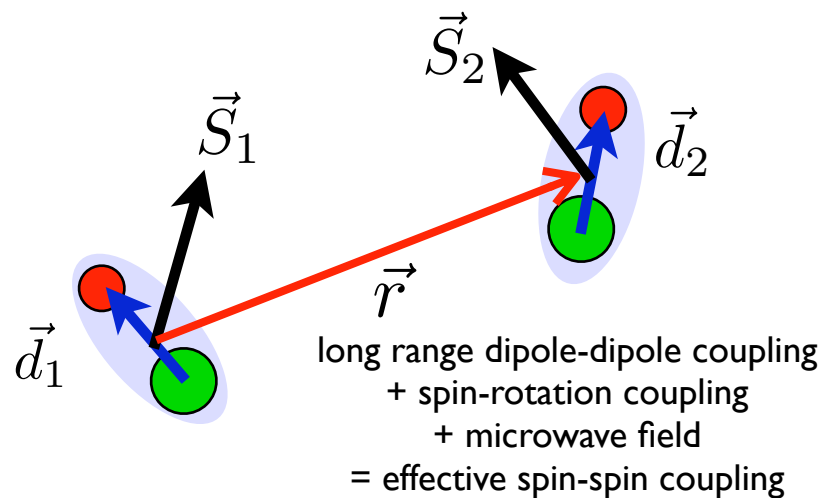
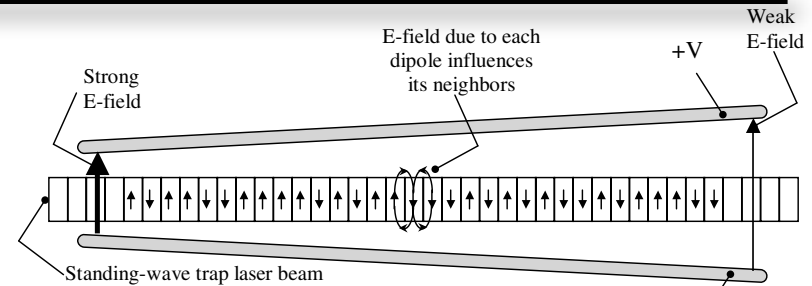
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with polar molecules in an optical lattice dressed with a microwave field, you can realize just about any spin lattice model : 1D xyz, 2D Ising, 3D Heisenberg, Kitaev model, etc...*



Polar molecules: scientific motivation

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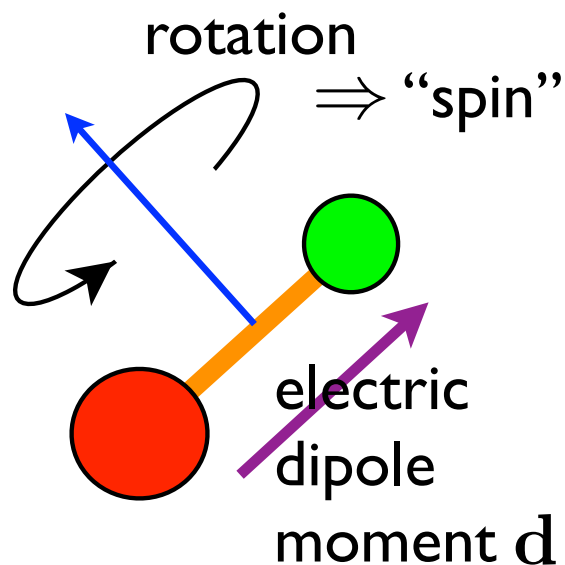
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for this, you need paramagnetic, polar molecules !

Polar molecules: scientific motivation

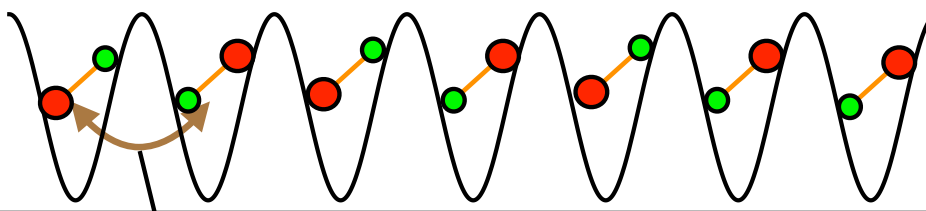
from Alexey V. Gorshkov's Jan. 24th talk, "Topological Phases in Polar-Molecule Quantum Magnets"



highly tunable exotic spin models

Our achievements:

- stronger interactions
- much higher tunability \Rightarrow exotic physics



dipole-dipole interactions \Rightarrow "spin-spin" interactions

Barnett et al., PRL (2006)
Micheli, Brennen, Zoller, Nat Phys (2006)
Brennen, Micheli, Zoller, NJP (2007)
Buchler, Micheli, Zoller, Nat Phys (2007)
Watanabe, PRA (2009)
Wall, Carr, NJP (2009)
Pupillo et al, in *Cold Molecules* (2009)

Wall, Carr, PRA (2010)
Schachenmayer et al., NJP (2010)
Perez-Rios, Herrera, Krems, NJP (2010)
Trefzger et al., NJP (2010)
Herrera, Litinskaya, Krems, PRA (2010)
Kestner et al., PRB (2011),
Lemeshko et al, PRL (2012), etc...

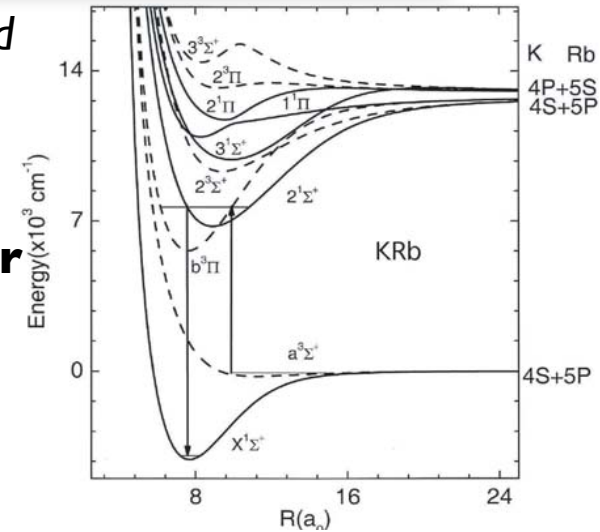
for some of these proposals, you need paramagnetic, polar molecules !

Quantum degenerate polar molecules from cold atoms: state of the art

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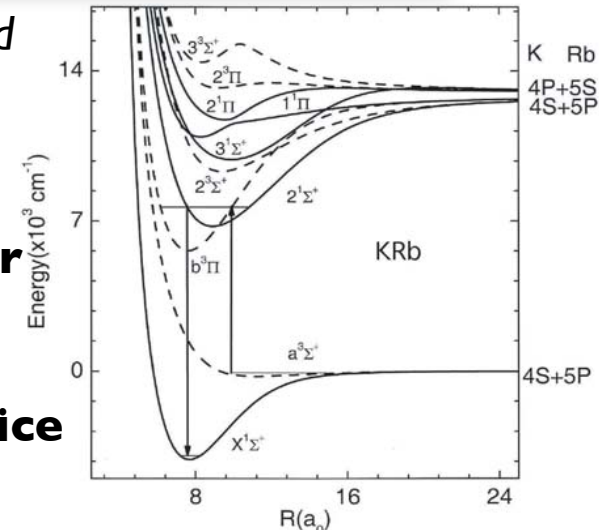
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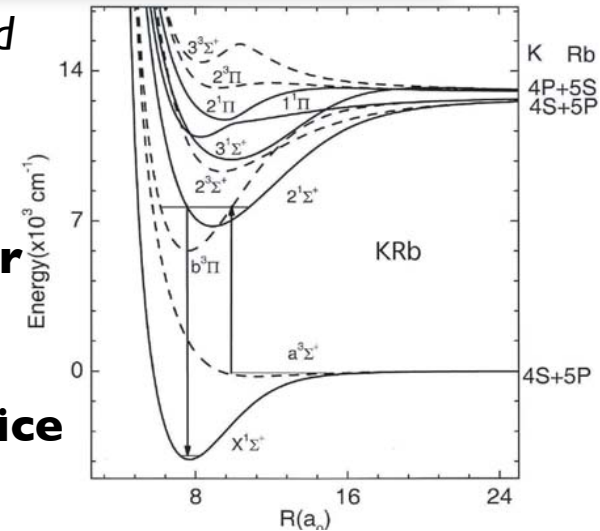
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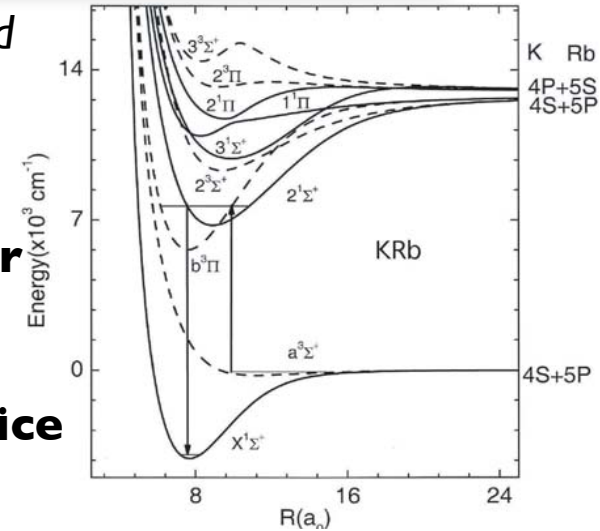


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Dipole moments (Debye)

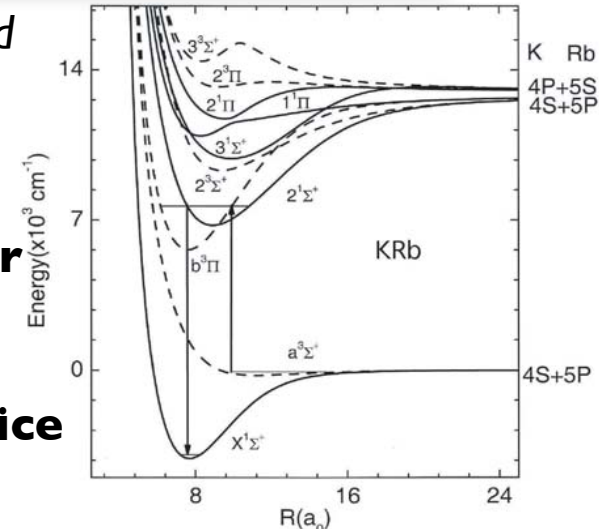
$\begin{matrix} Y \\ X \end{matrix}$	Na	K	Rb	Cs
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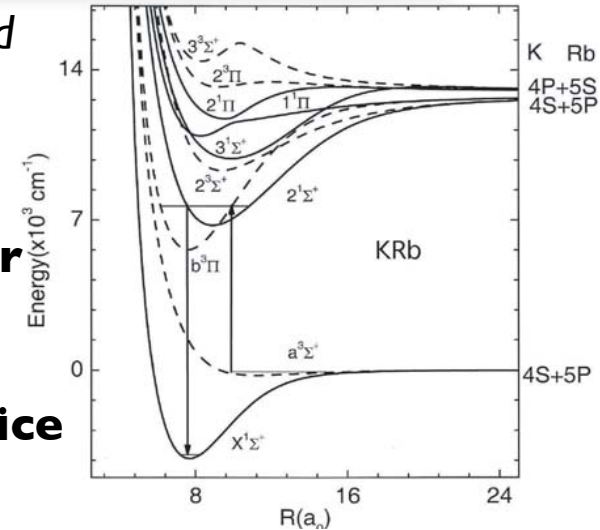
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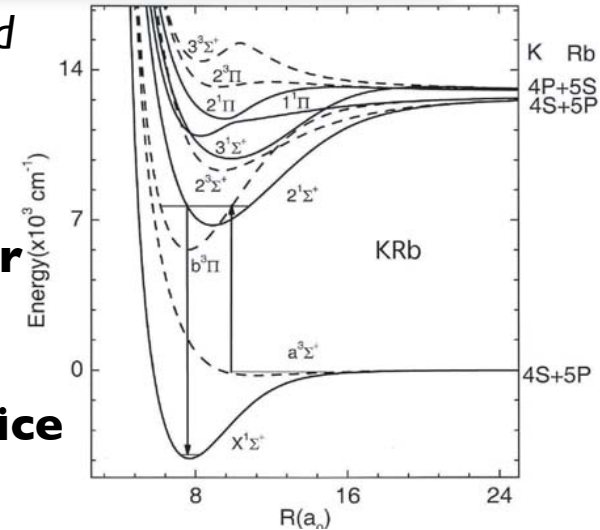
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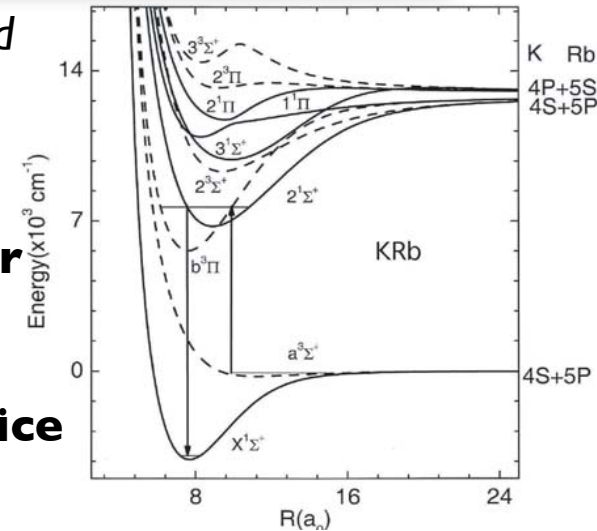
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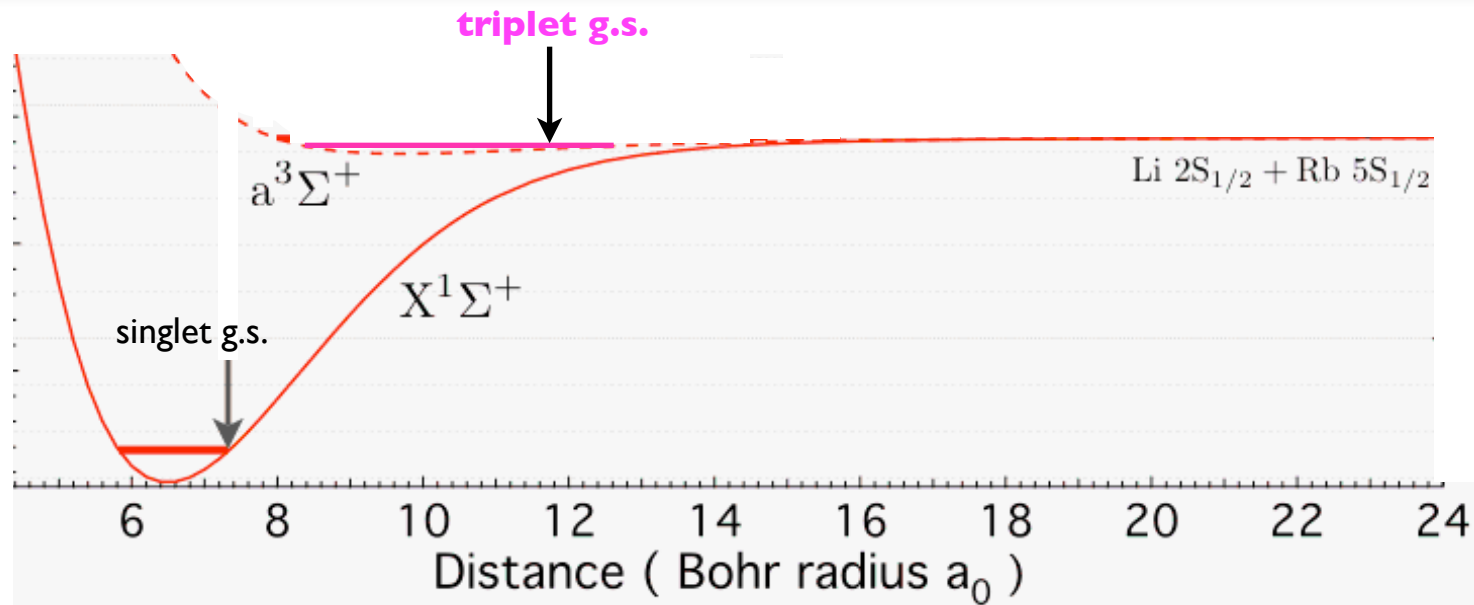
what about the triplet state of a bi-alkali molecule?

Dipole moments (Debye)

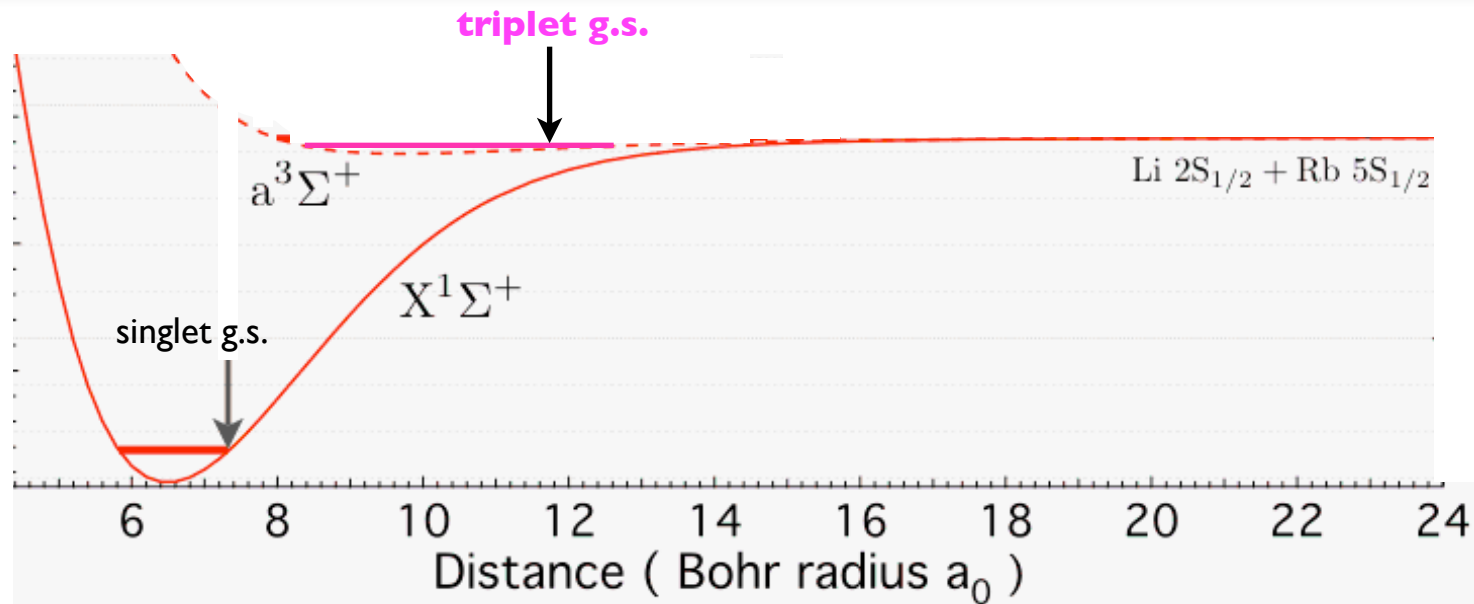
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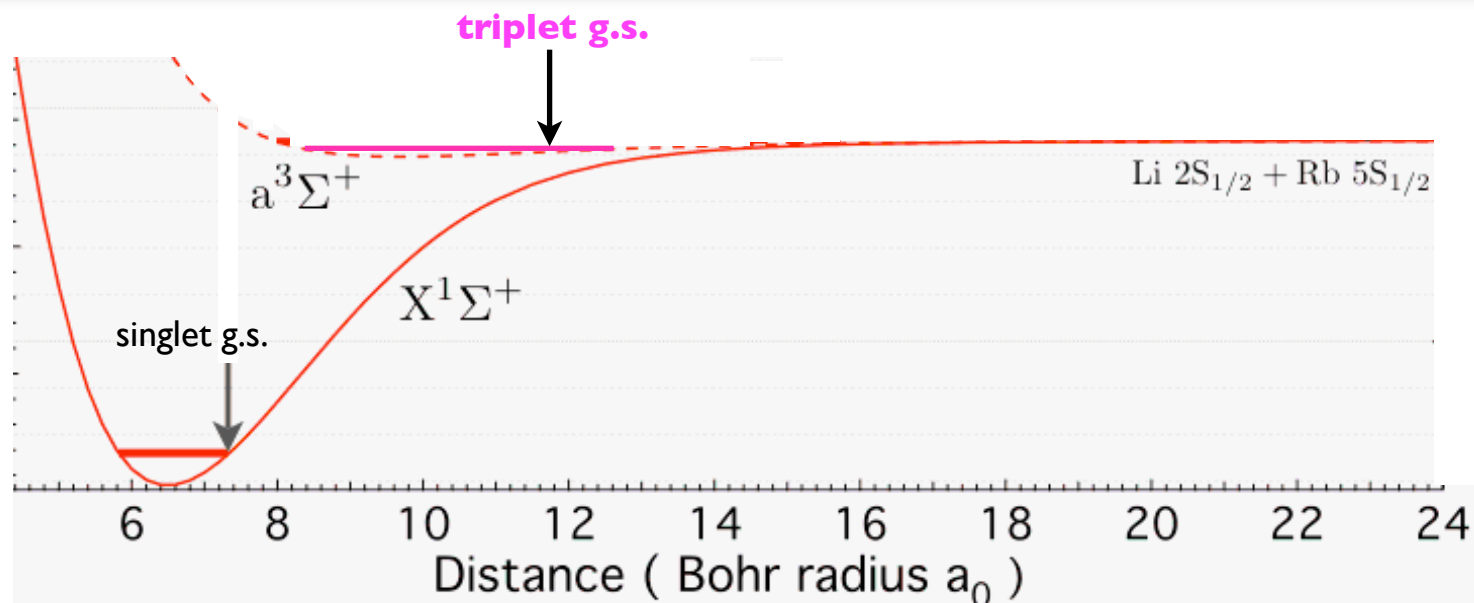


What about the lowest triplet state?



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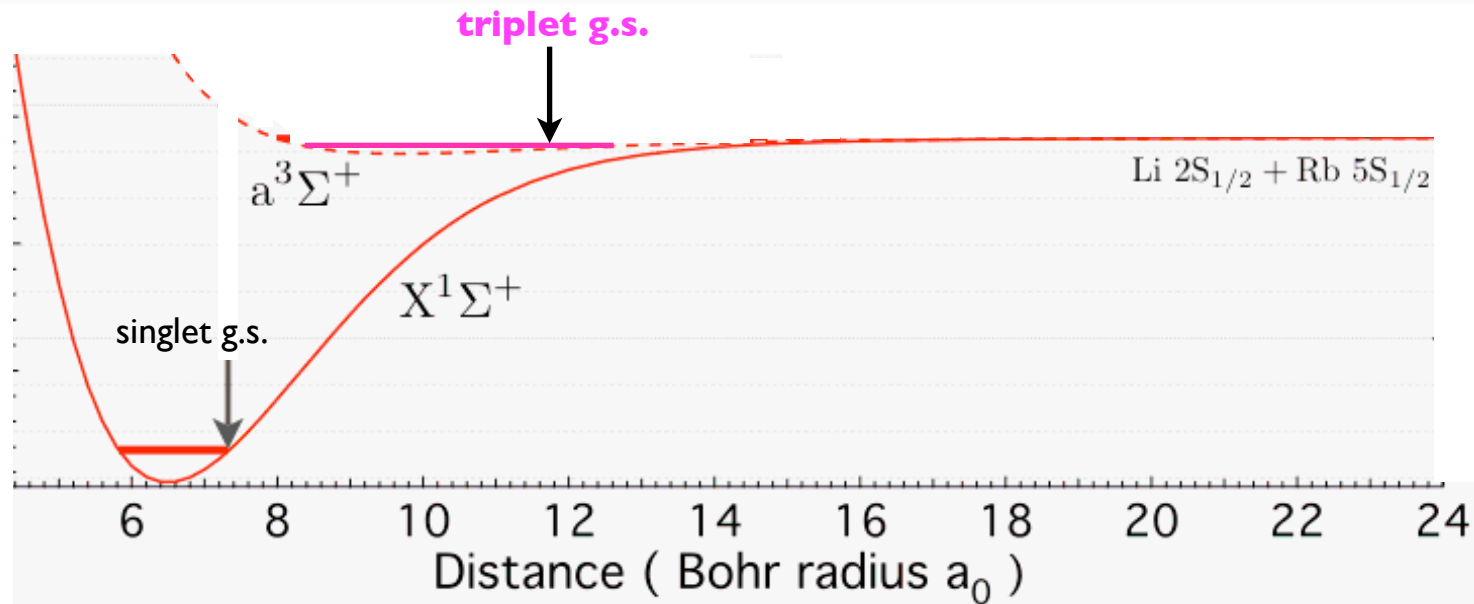
DM of triplet state LiRb is about 0.4 Debye
(DM of singlet state KRb is about 0.5 Debye)

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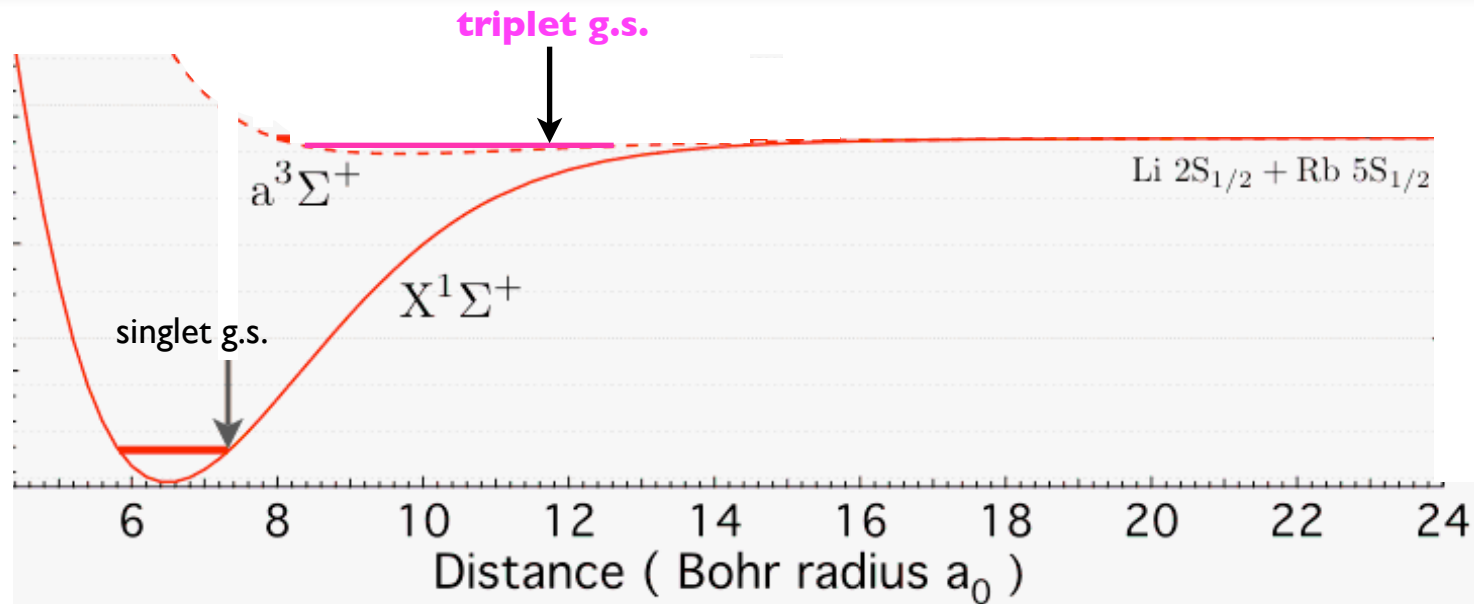
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- **chemical reactivity ?**
- **spontaneous emission ?**

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Is the lowest triplet state stable?

- spin relaxation collisions (triplet to singlet coupling) ?

- chemical reactivity ?

by some magic, is the triplet state stable?*

** or suitably physical reasons*

- spontaneous emission ?

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“Spin-orbit couplings are small in light systems. For Li_2 , one wouldn't expect fast relaxation. ... but will be more important in heavy systems (like LiRb), but how important?”

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Physical Chemistry Chemical Physics

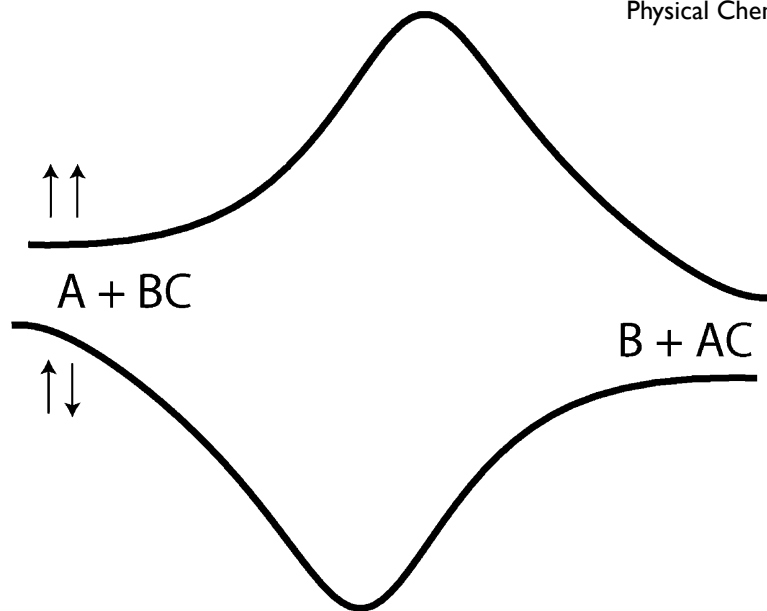


Fig. 4 Schematic illustration of minimum energy profiles for an $A(^2S) + BC(^2\Sigma)$ chemical reaction in the singlet-spin (lower curve) and triplet-spin (upper curve) electronic states. Electric fields may induce non-adiabatic transitions between the different spin states and modify the reaction mechanism.

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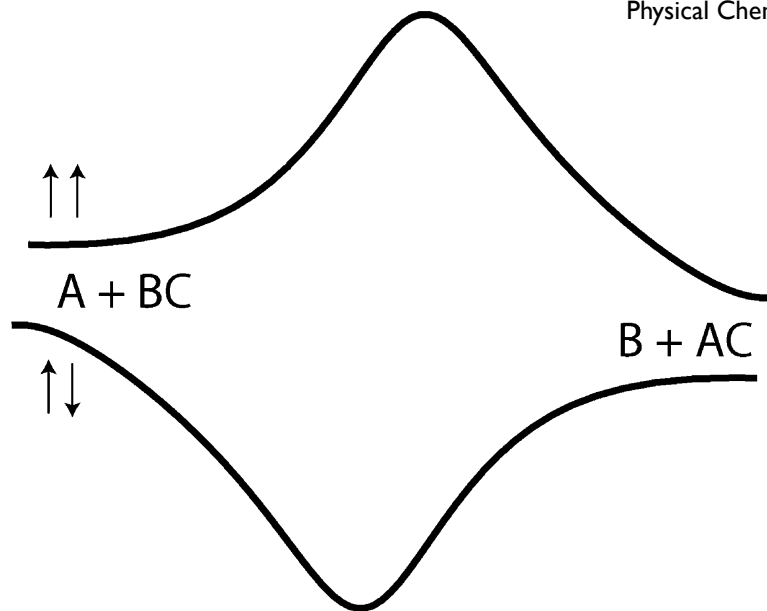


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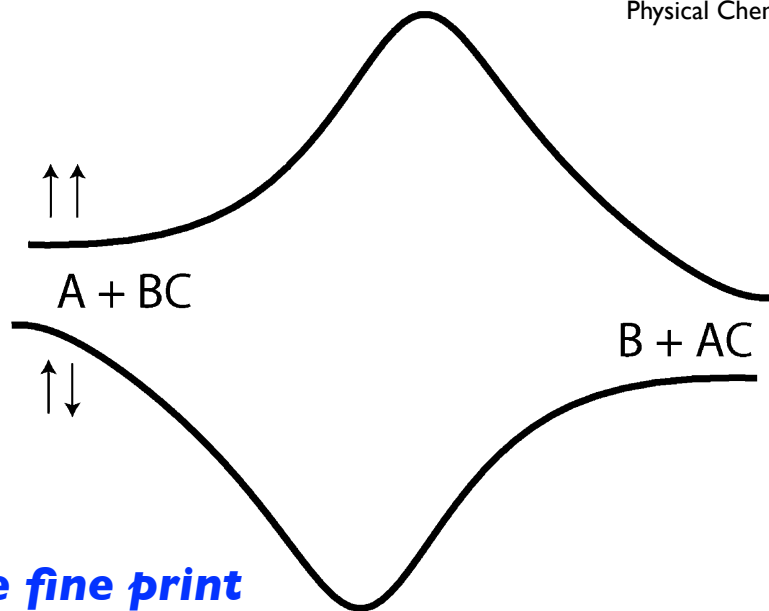


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the fine print

“... non-adiabatic coupling ... may be induced by the **spin-rotation interaction** and the **magnetic dipole–dipole** interaction. The latter is negligibly small and ... transitions are determined by the spin-rotation interaction in the open-shell molecule. The spin-rotation interaction can be effectively manipulated with an external electric field.”

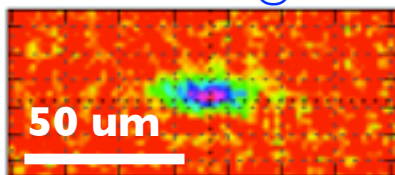
- **spontaneous emission ?**

Experimental goals and results

Li₂

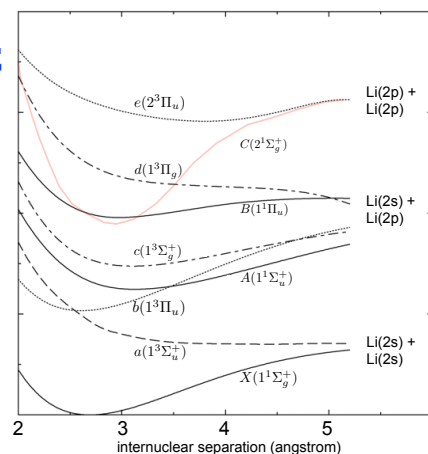
- ✓ Feshbach resonances and FR molecule production

BEC of Li dimers @ 300 nK



- ✓ excited triplet spectroscopy (1-color PA)

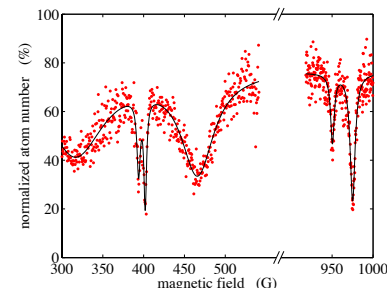
- ✓ ground triplet spectroscopy (2-color PA)



STIRAP production of triplet Li₂ from FR molecules

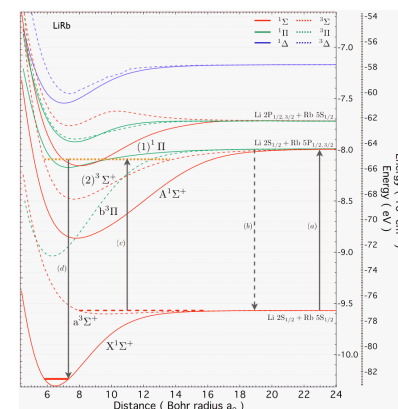
LiRb

- ✓ Feshbach resonances and FR molecule production



- excited triplet spectroscopy (1-color PA)

- ground triplet spectroscopy (2-color PA)



STIRAP production of triplet LiRb from FR molecules

Experimental goals and results

table from Eberhard Tiemann's Jan. 25th talk, "Molecular Properties and Potentials"

Overview of results

FR: Feshbach resonances observed

a: potential for $a^3\Sigma^+$ determined



	Li	Na	K	Rb	Cs
Li	FR				
Na	FR	FR (a)			
K	FR a	FR a	FR a		
Rb	FR (a)	FR a	FR a	FR a	
Cs	FR (a)	a	a	FR a	FR (a)

Experimental goals and results

table from Eberhard Tiemann's Jan. 25th talk, "Molecular Properties and Potentials"

Overview of results

FR: Feshbach resonances observed

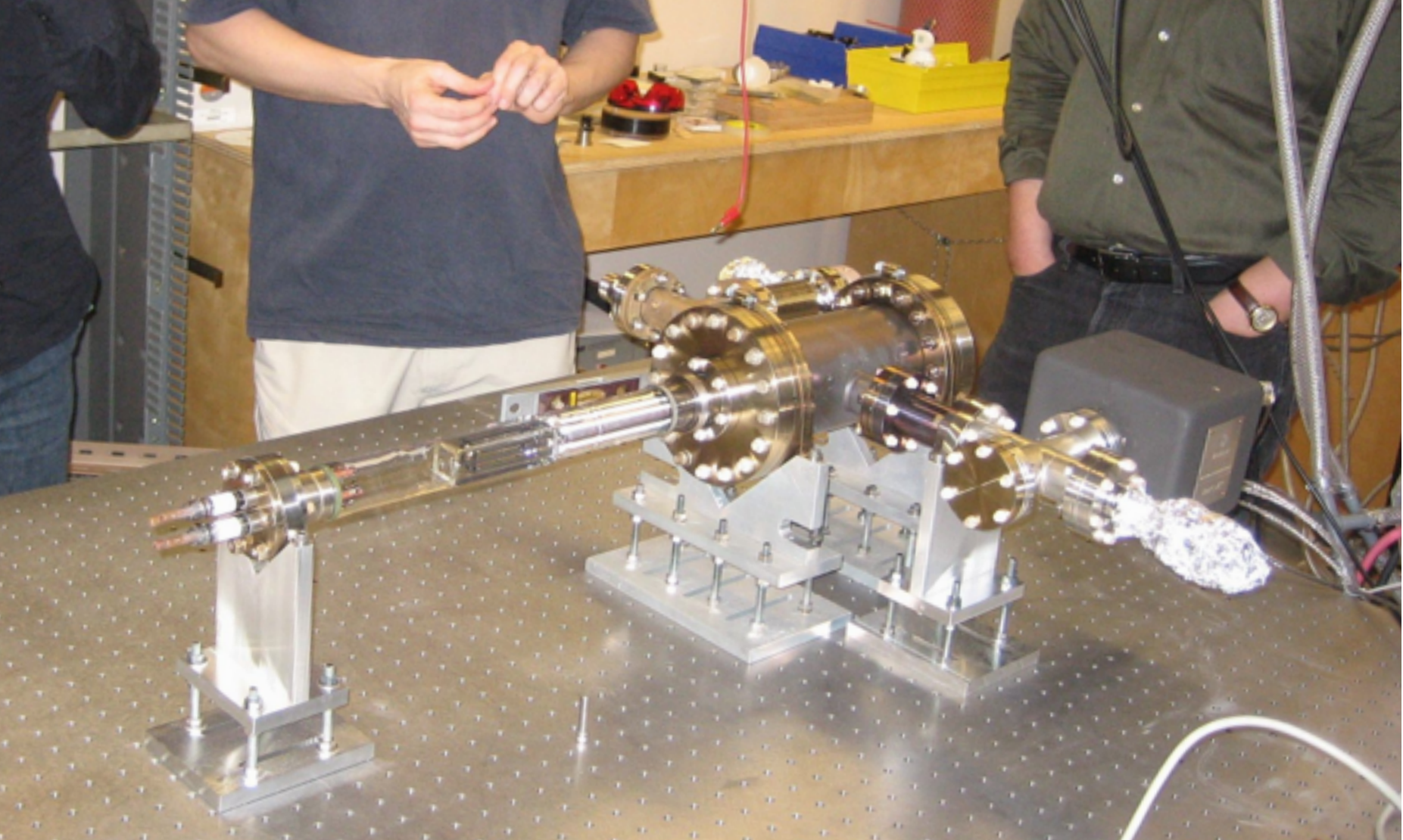
a: potential for $a^3\Sigma^+$ determined



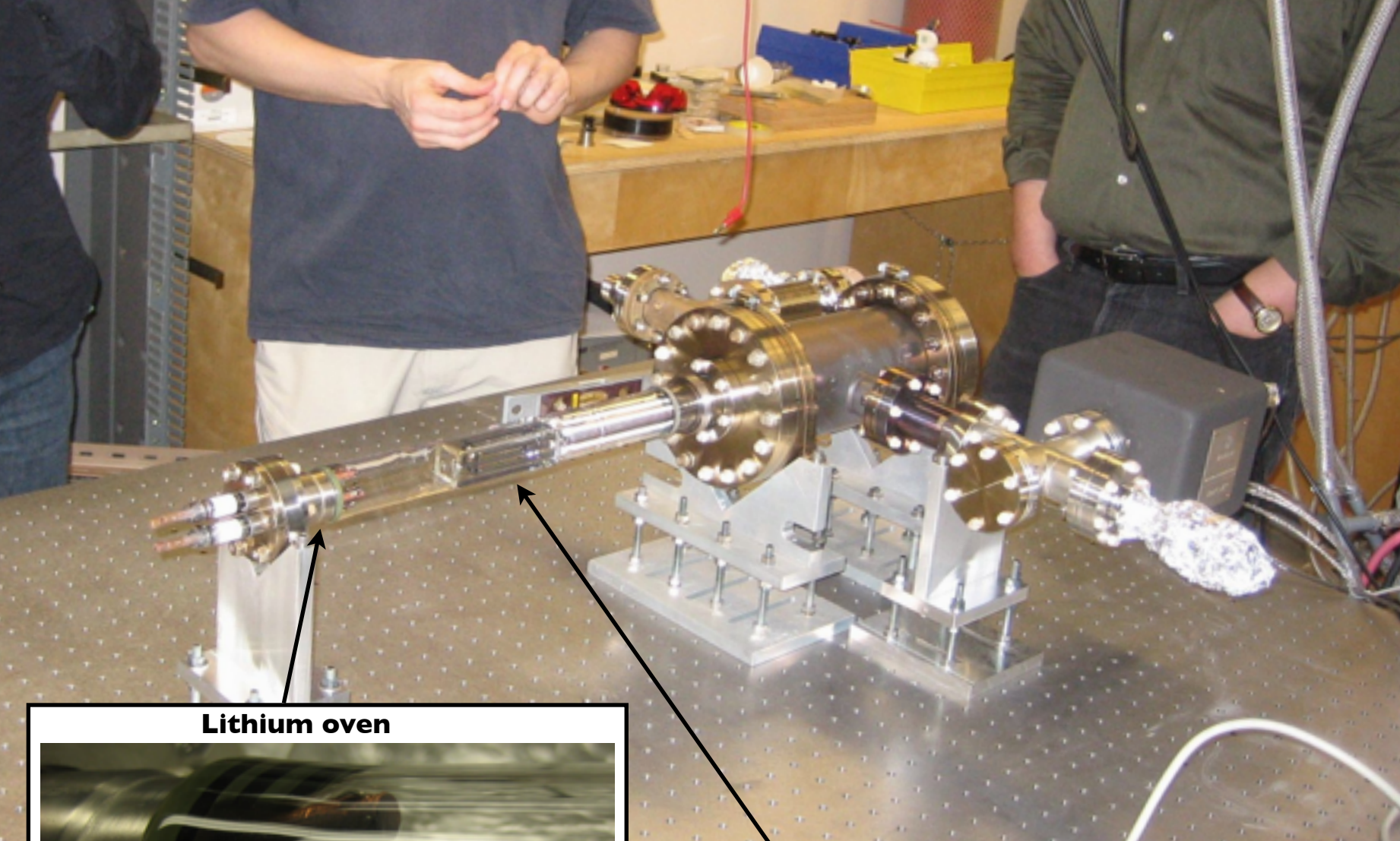
	Li	Na	K	Rb	Cs
Li	FR a				
Na	FR	FR (a)			
K	FR a	FR a	FR a		
Rb	FR (a) FR a	FR a	FR a	FR a	
Cs	FR (a)	a	a	FR a	FR (a)

our contributions in blue

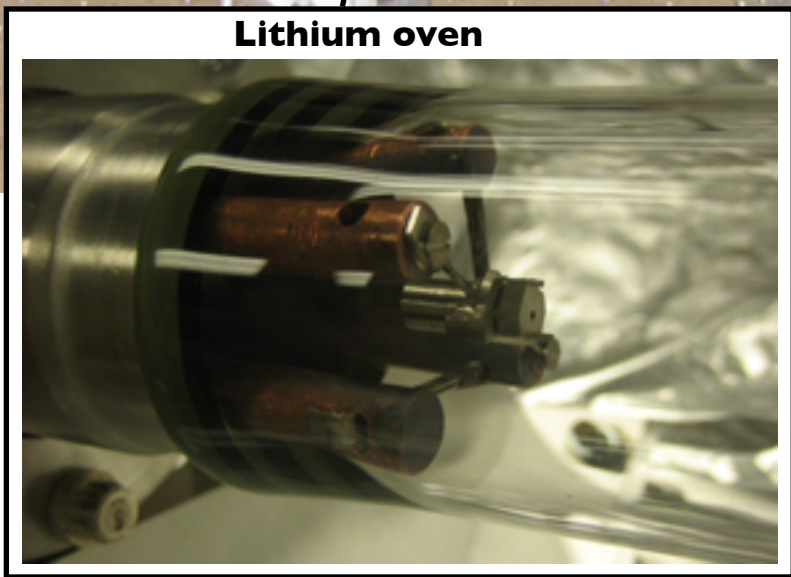
pending contributions in red



ultra-cold Li+Rb mixtures

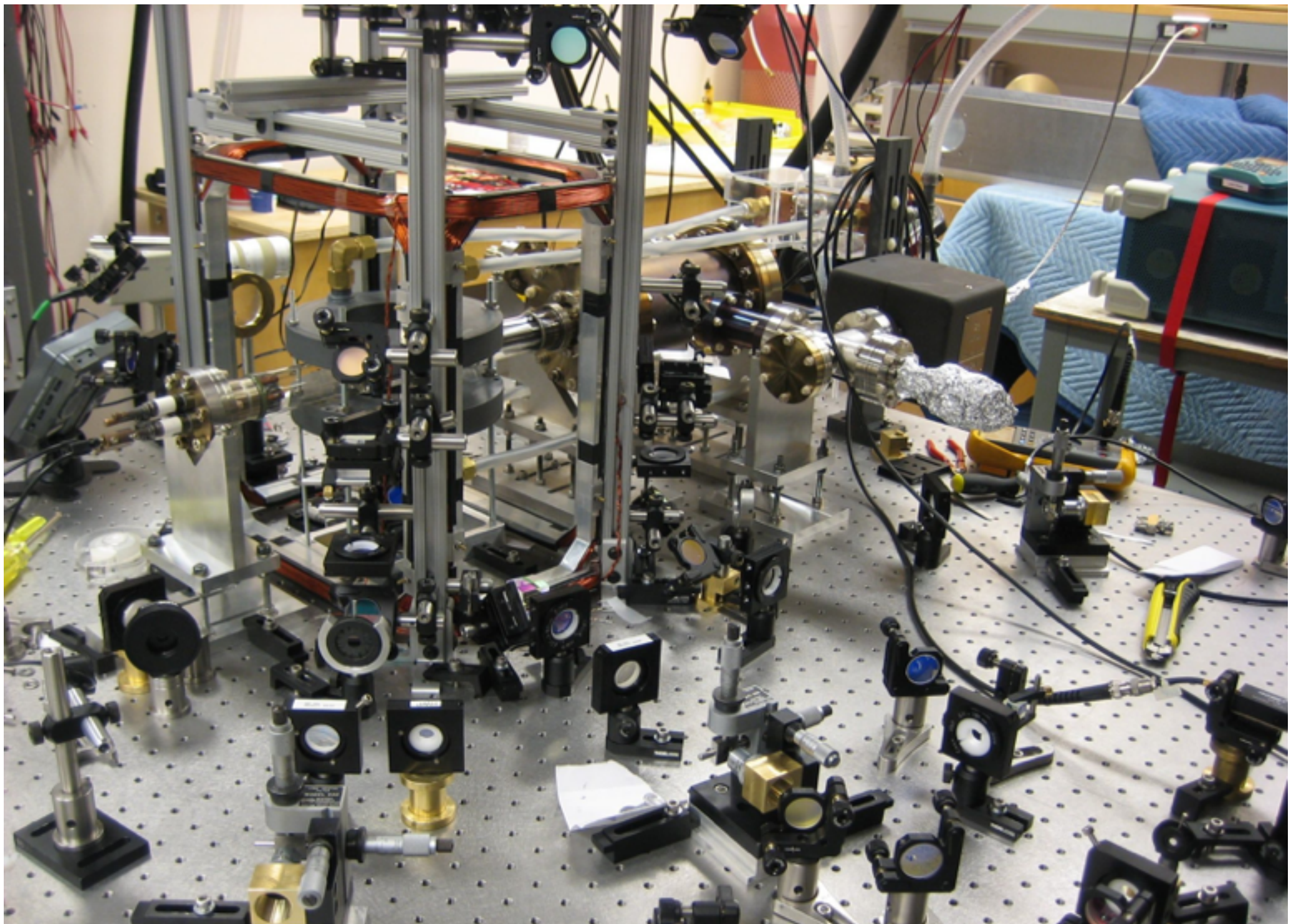


Lithium oven



MOT cell

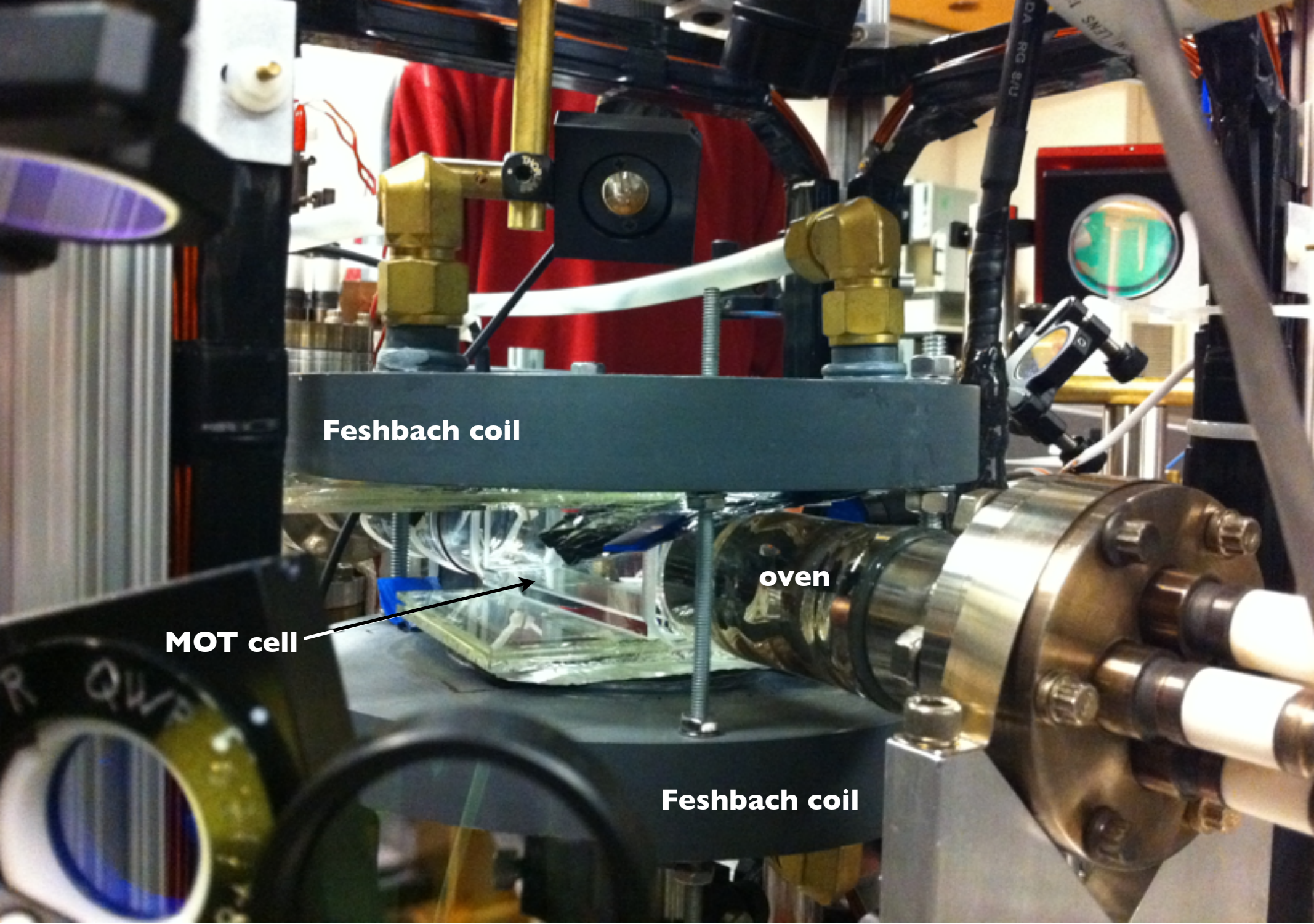
ultra-cold Li+Rb mixtures



ultra-cold Li+Rb mixtures



ultra-cold Li+Rb mixtures



Feshbach coil

oven

MOT cell

Feshbach coil

ultra-cold Li+Rb mixtures

Feshbach coil

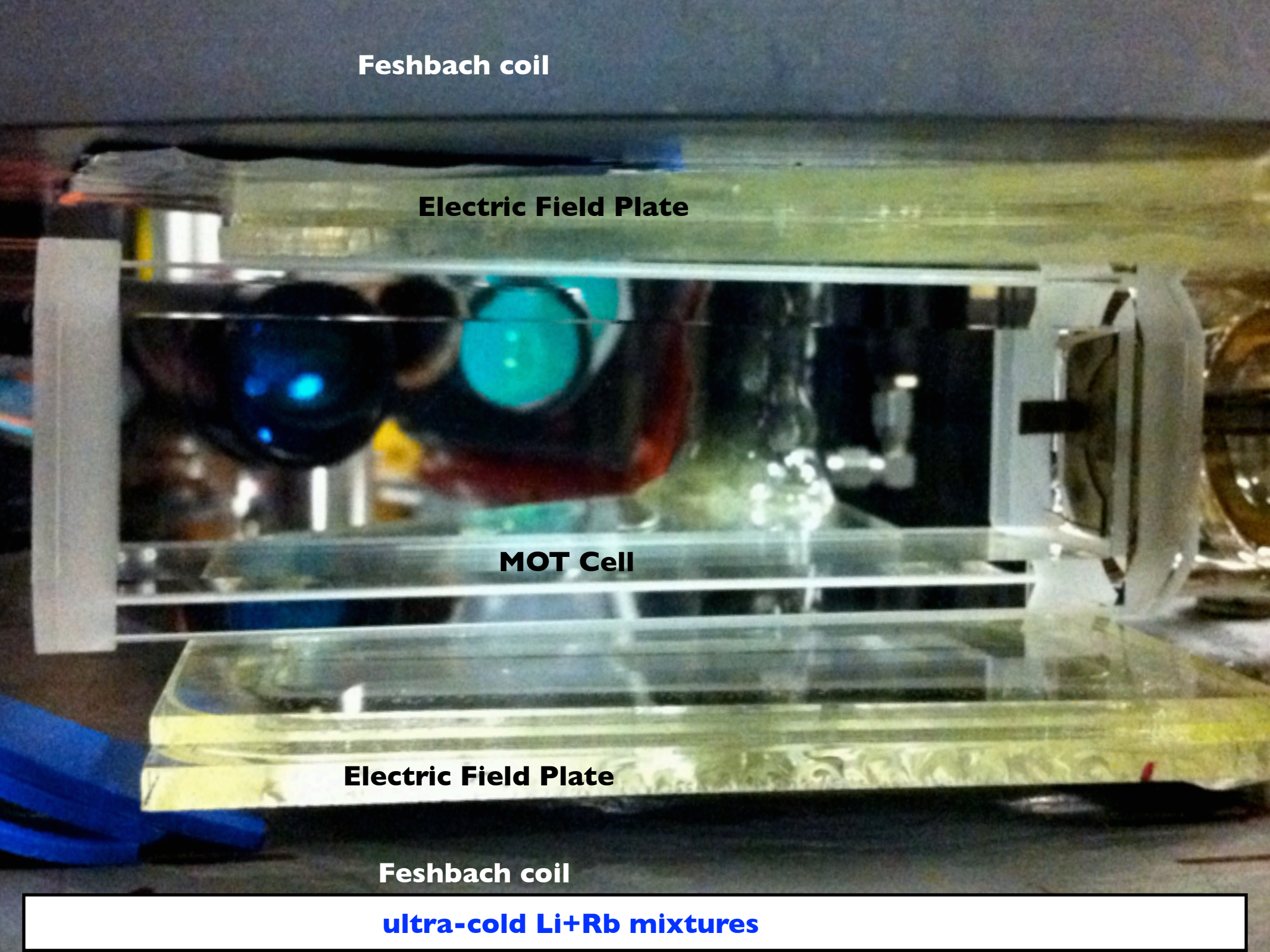
Electric Field Plate

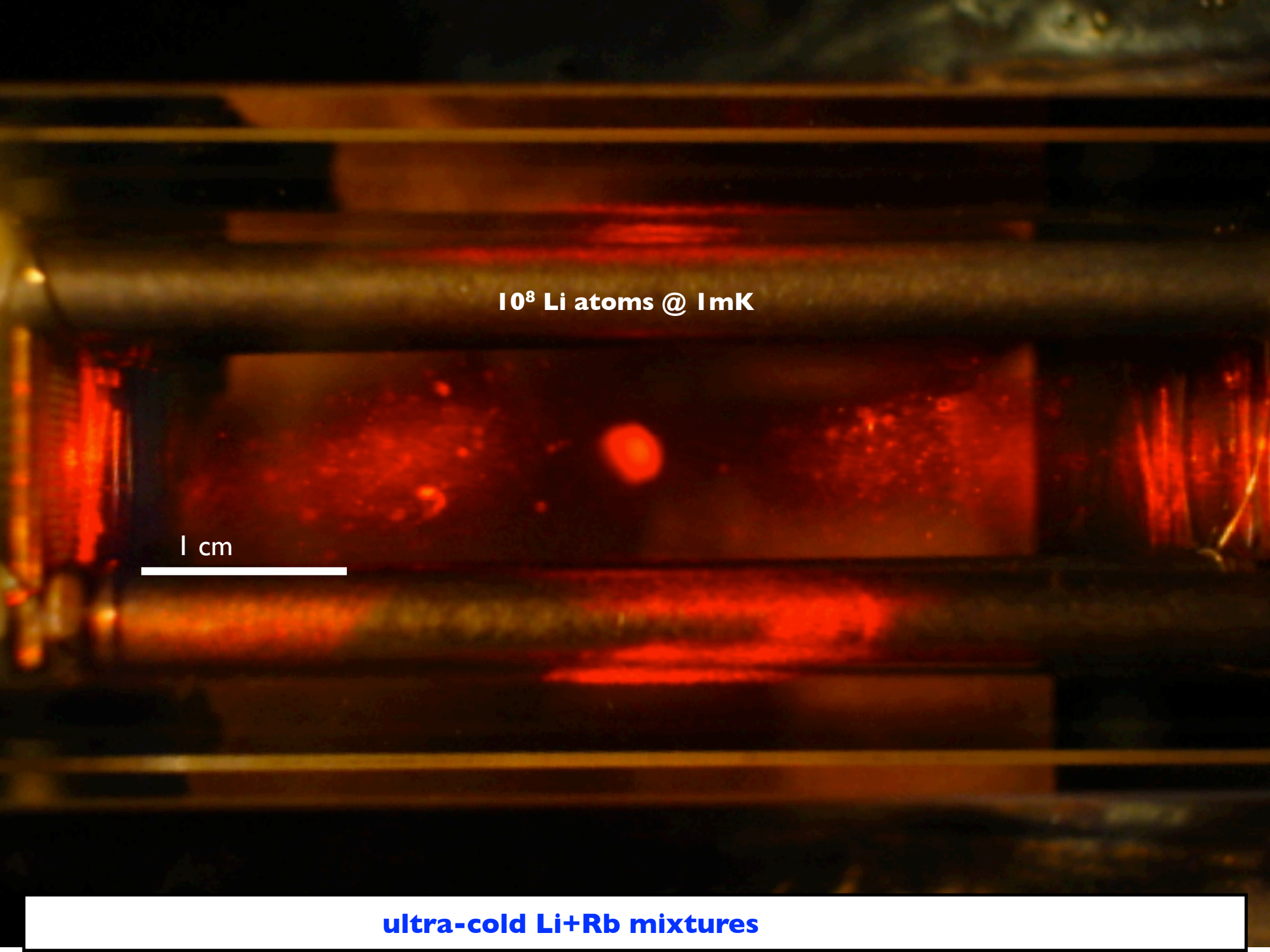
MOT Cell

Electric Field Plate

Feshbach coil

ultra-cold Li+Rb mixtures



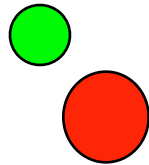


10^8 Li atoms @ 1 mK

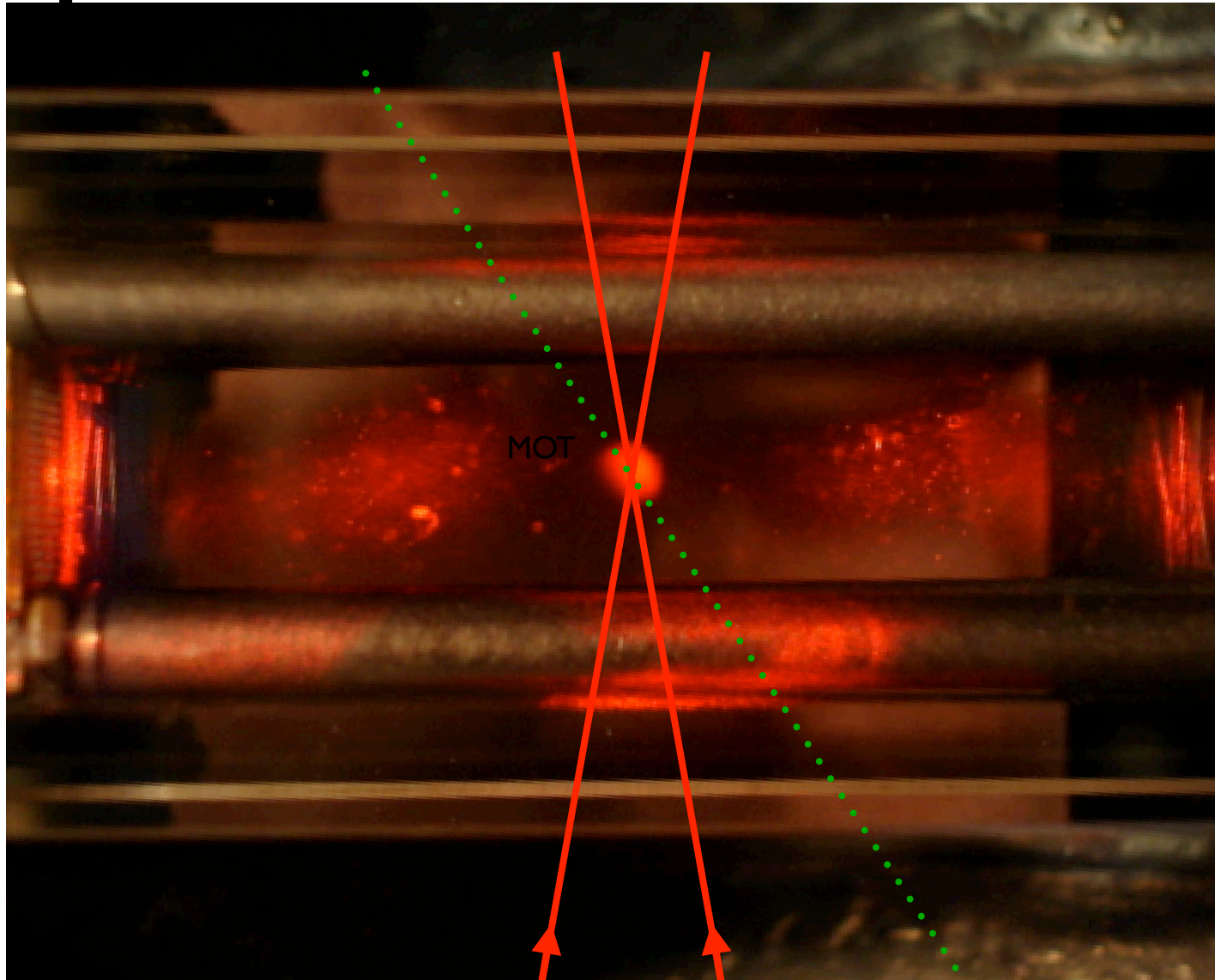
1 cm

ultra-cold Li+Rb mixtures

observing Rb+Li Feshbach resonances



Trap Schematic



MOT

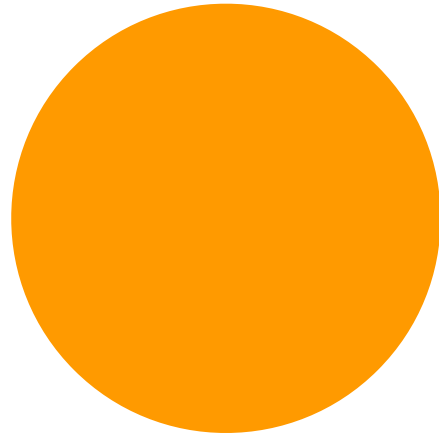
20 W

20 W

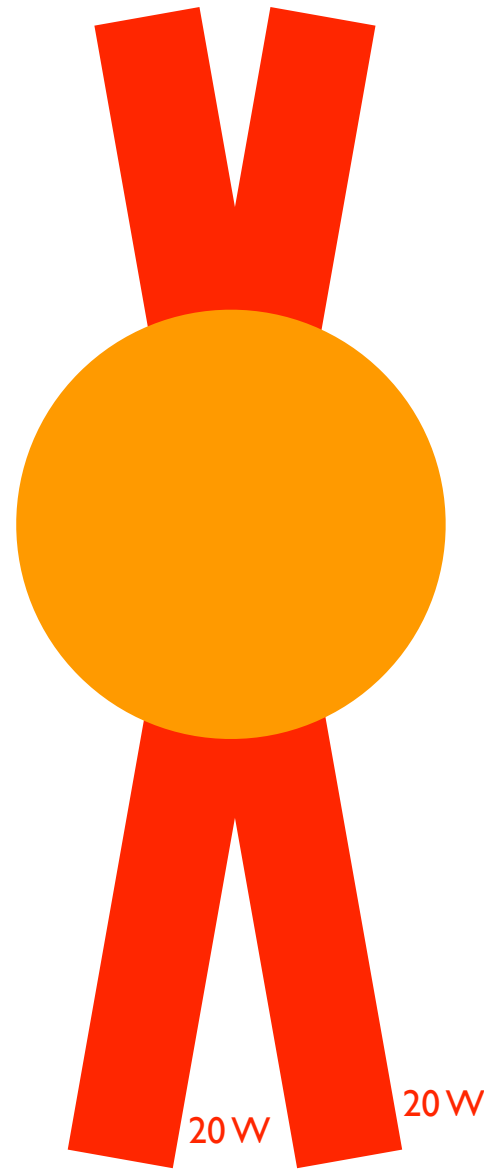
Crossed ODT
(optical dipole trap)

imaging axis

Load Lithium MOT

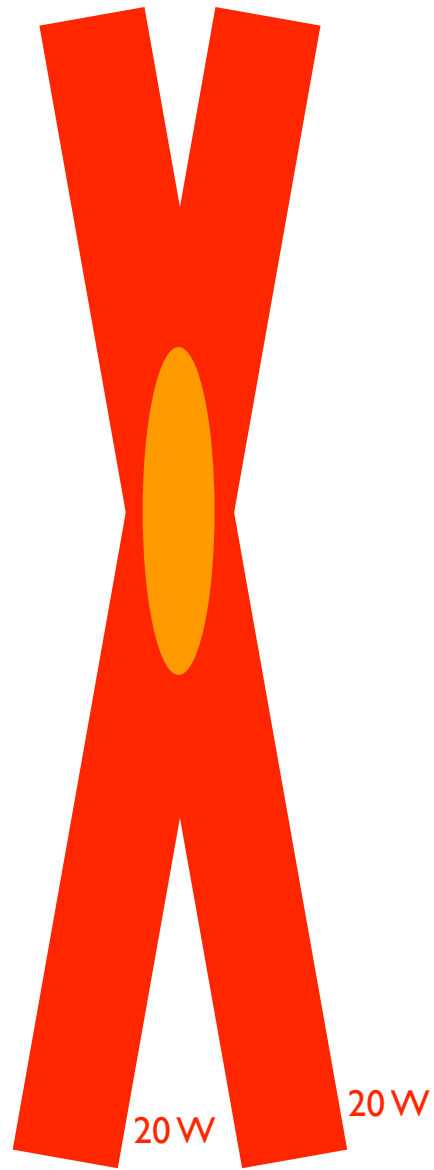


Transfer Li from MOT to crossed trap



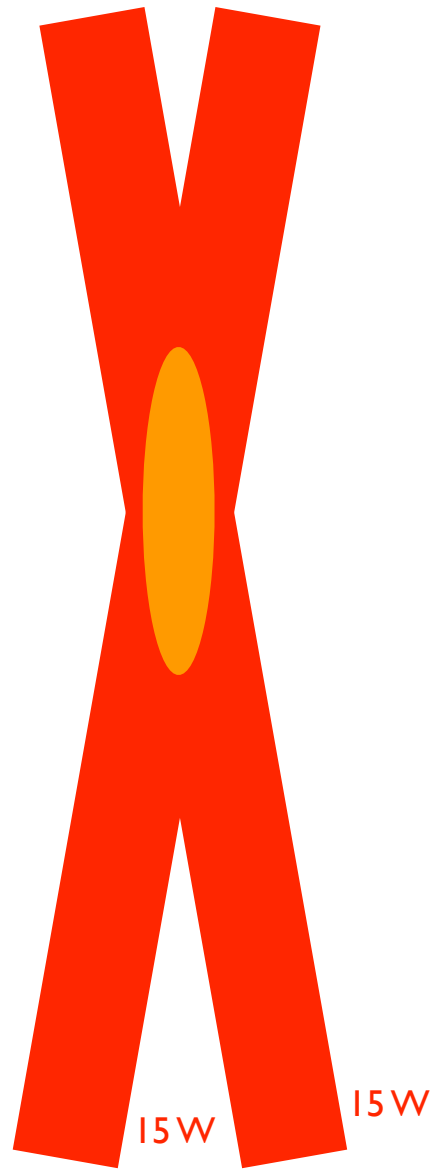
Crossed ODT

Transfer Li from MOT to crossed trap



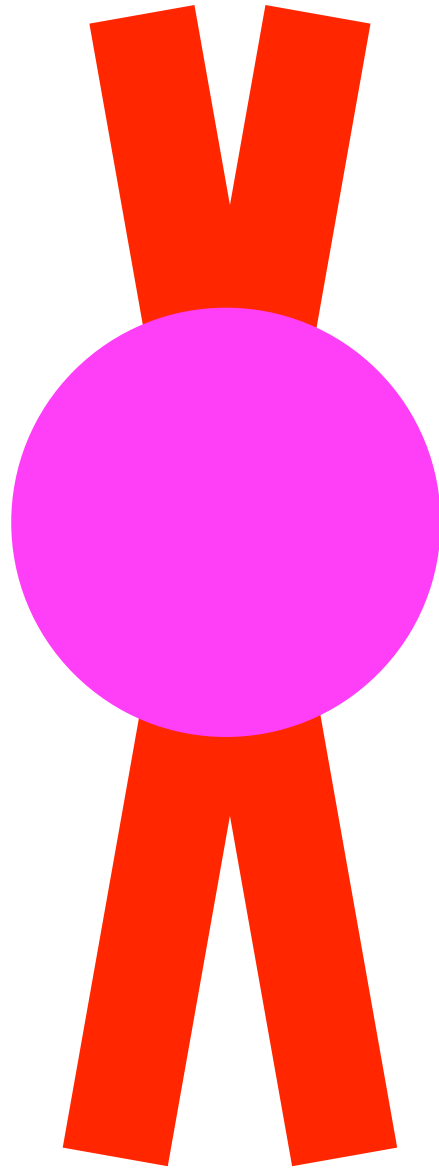
Crossed ODT

Evaporate Li from crossed trap



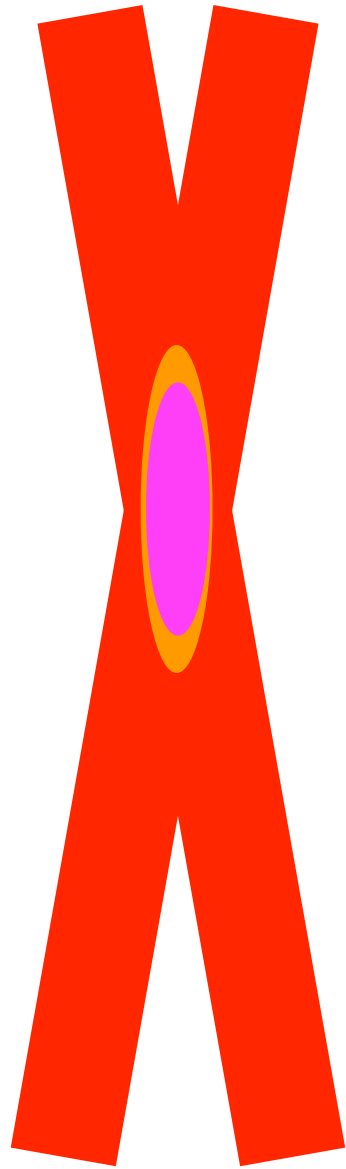
Crossed ODT

Load Rb MOT while holding Li



Crossed ODT

Transfer Rb to crossed trap

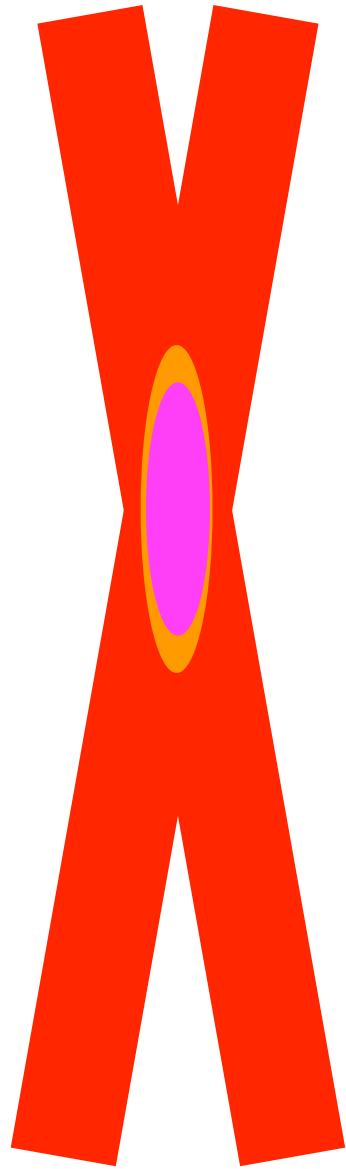


Rb MOT turned off:

In trap: Rubidium and Lithium

Crossed ODT

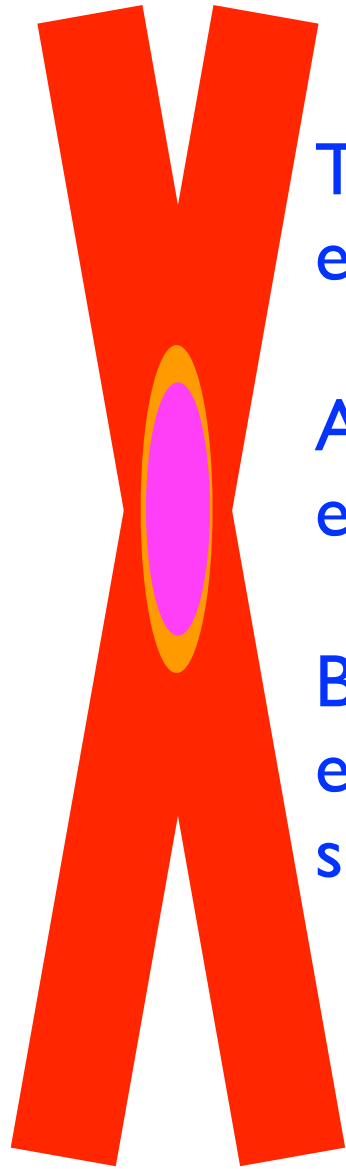
Both Rb and Li in trap: *trap depths different*



Trap for Li half as deep for Rb

Crossed ODT

Simultaneous evaporation of Li and Rb



Trap power lowered to force evaporation losses

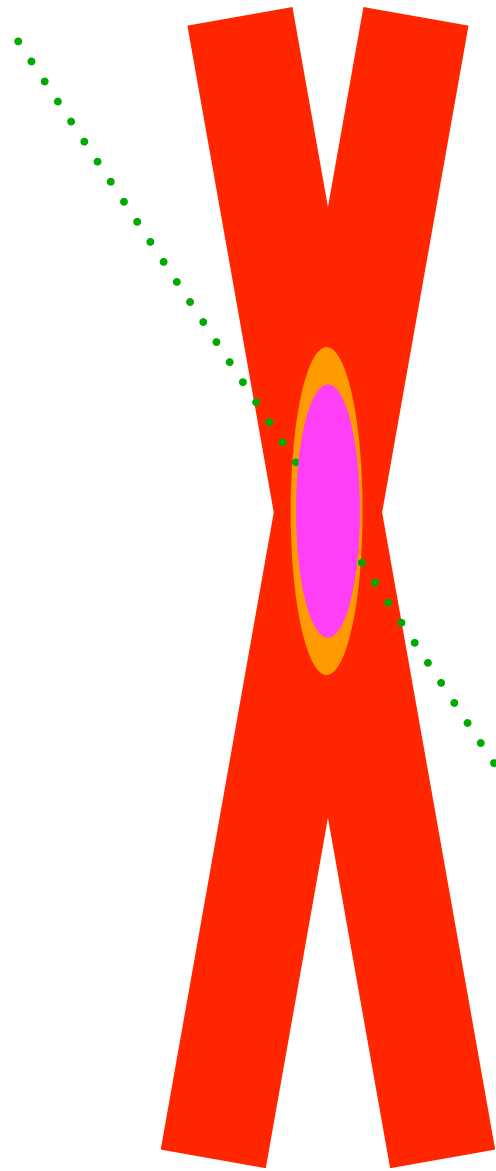
A) If Li+Rb reach thermal equilibrium, Li leaves trap quickly

B) If Li+Rb decoupled, Li evaporatively cools and leaves trap slowly

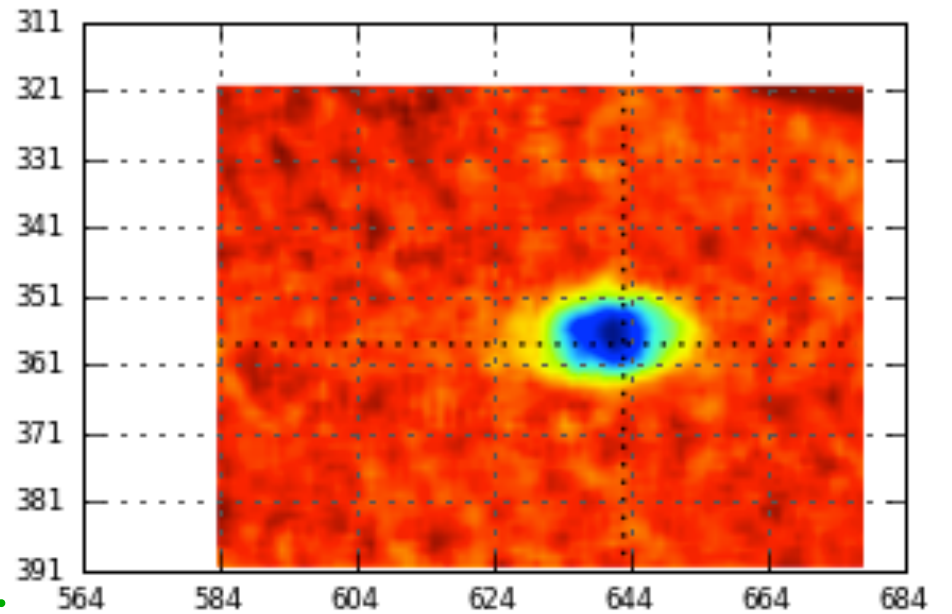
Trap for Li half as deep for Rb

Crossed ODT

Image Li or Rb

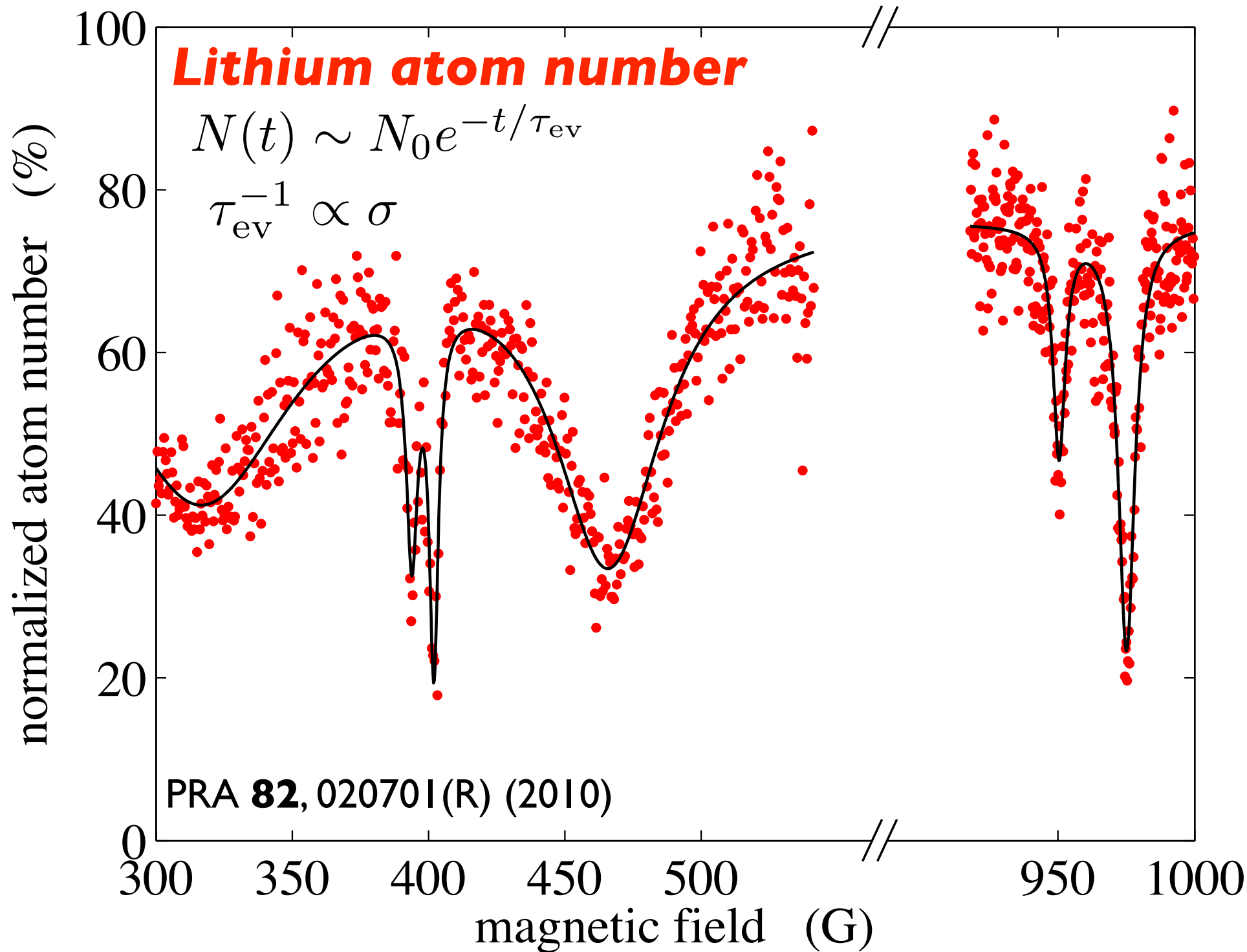


Crossed ODT

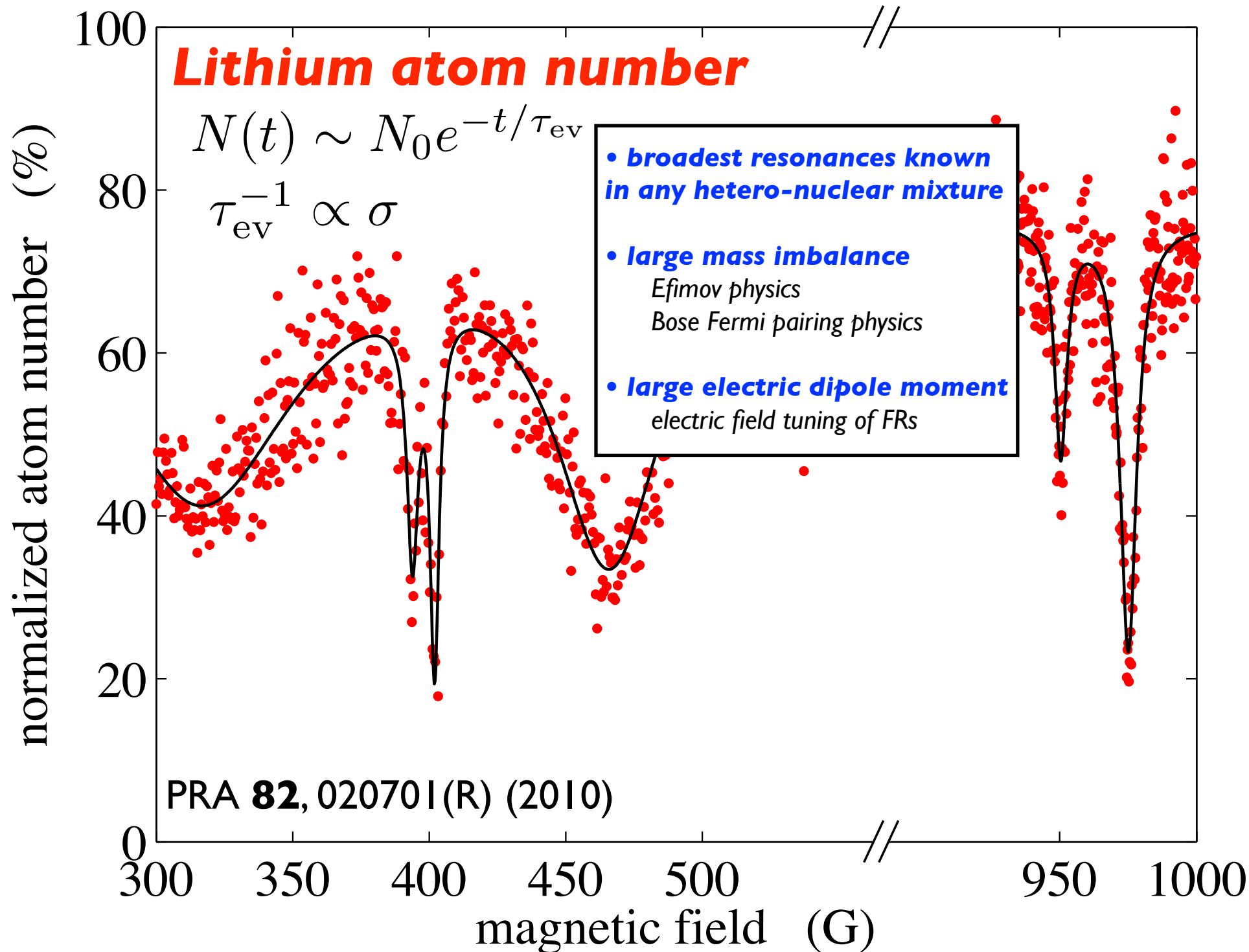


imaging axis

Feshbach resonances in ${}^6\text{Li}+{}^{85}\text{Rb}$ mixtures



Feshbach resonances in ${}^6\text{Li}+{}^{85}\text{Rb}$ mixtures



For photo-association, just add additional light

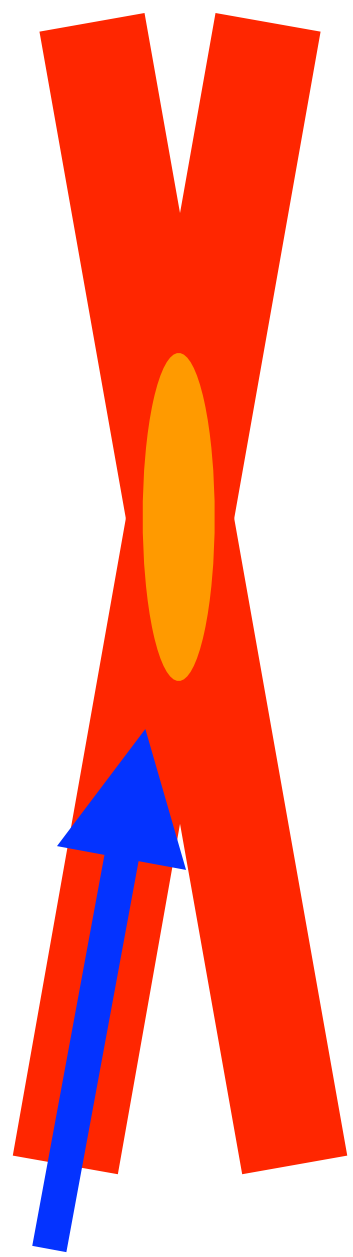
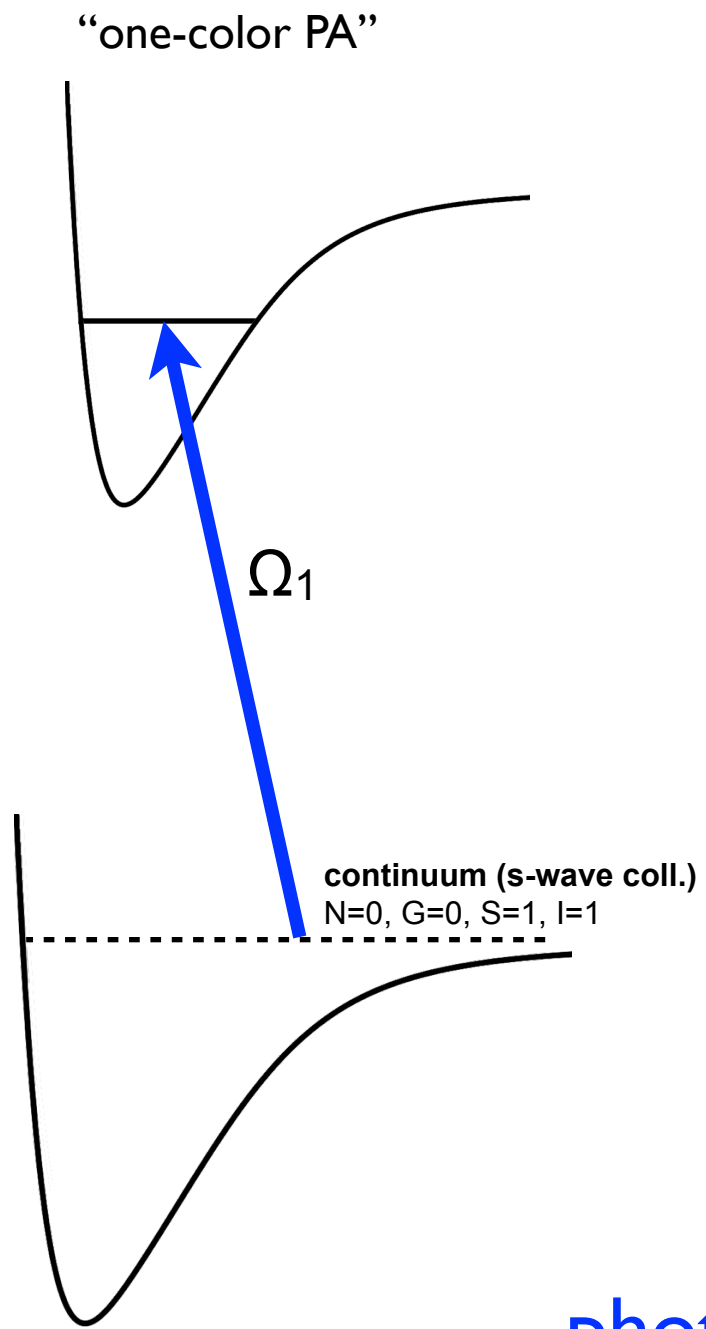


photo-association light

For photo-association, just add additional light

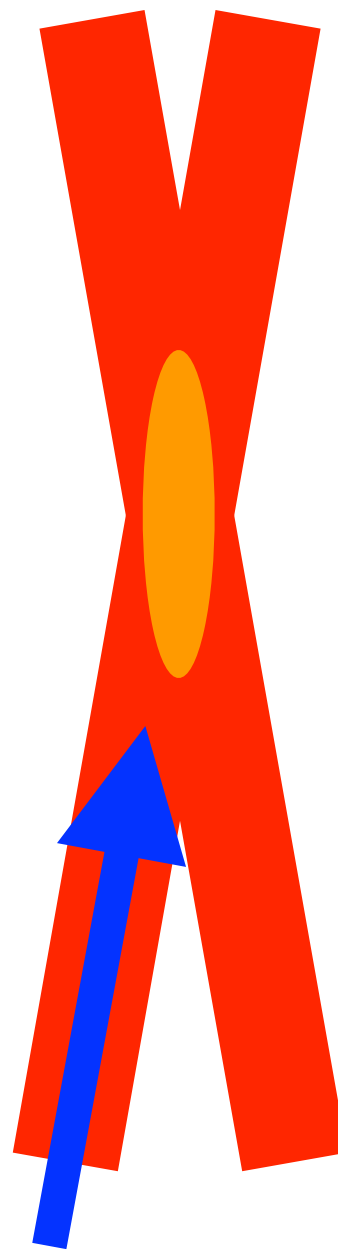
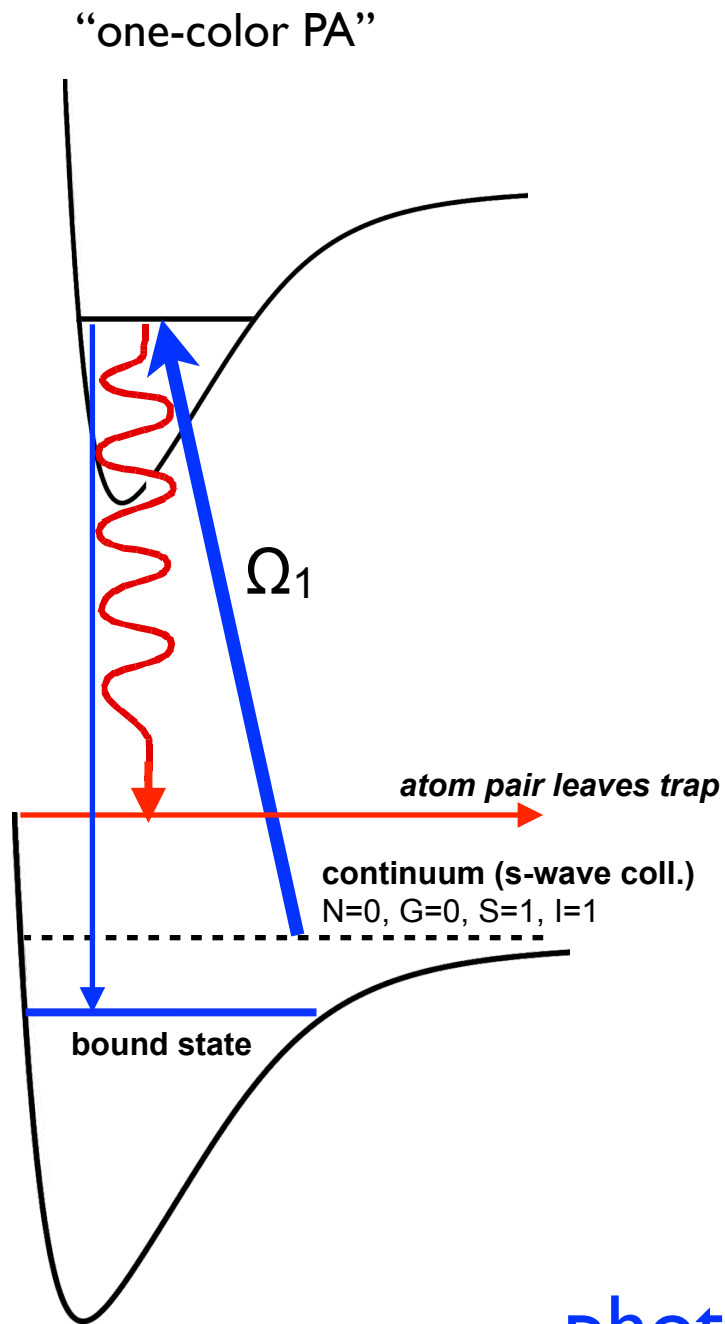
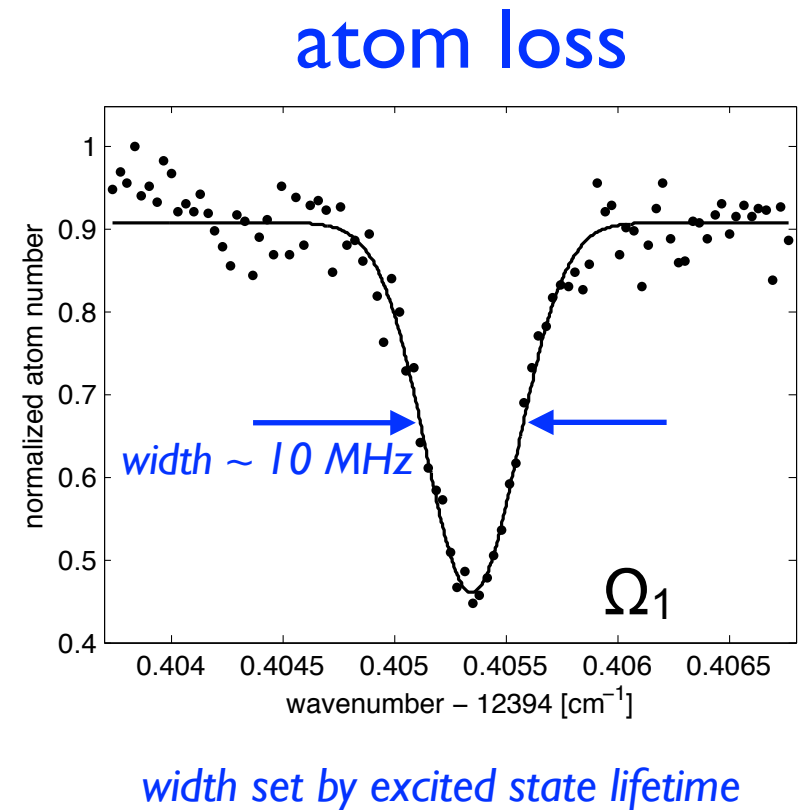
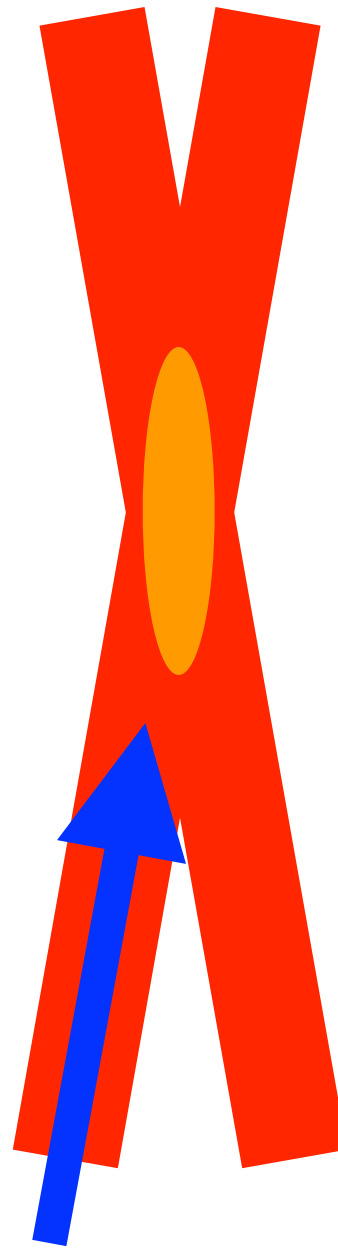
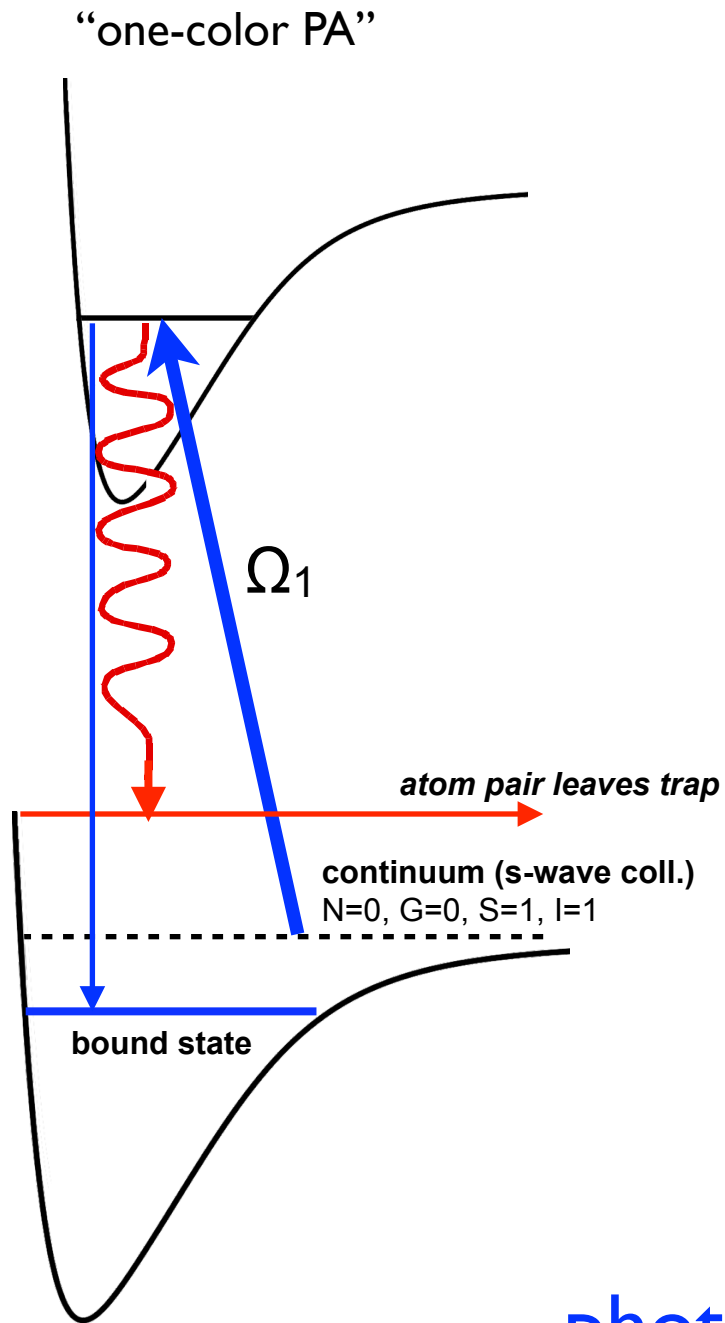


photo-association light

For photo-association, just add additional light



For photo-association, just add additional light

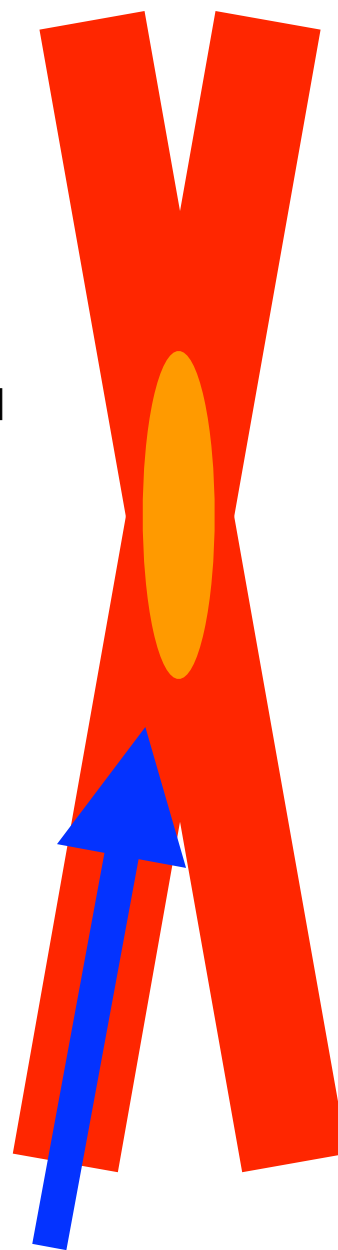
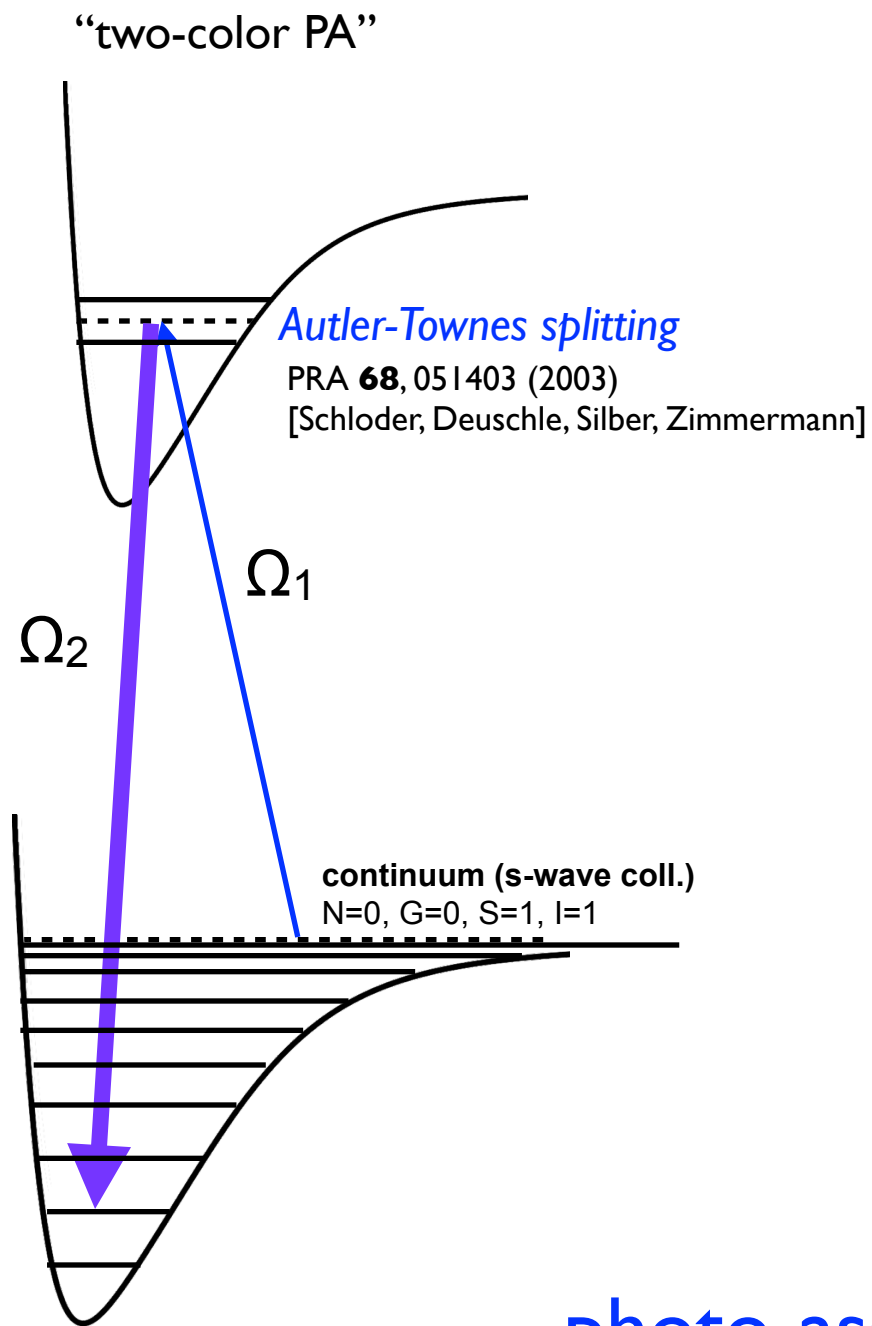
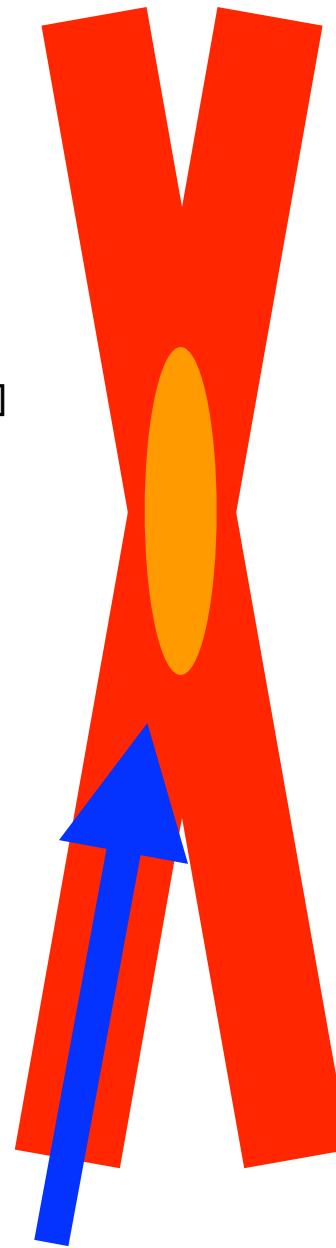
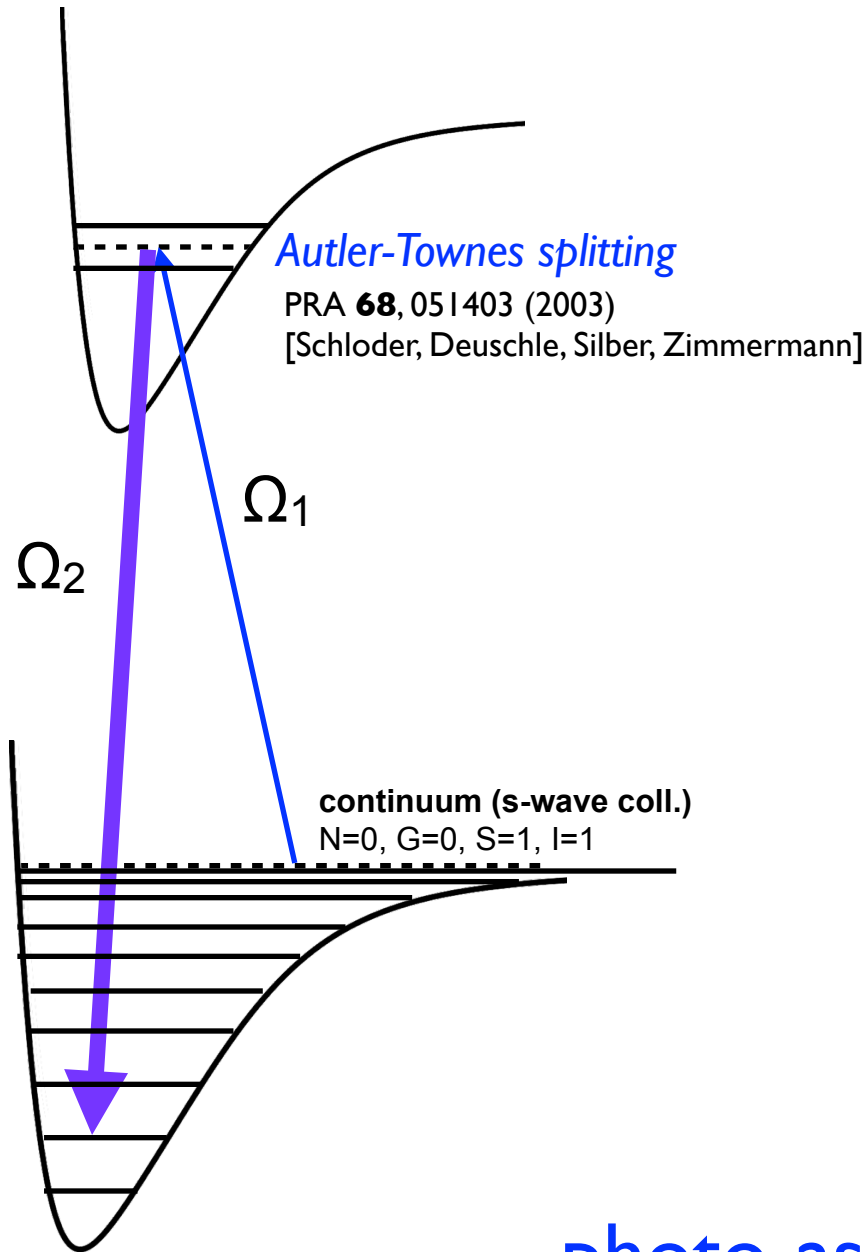


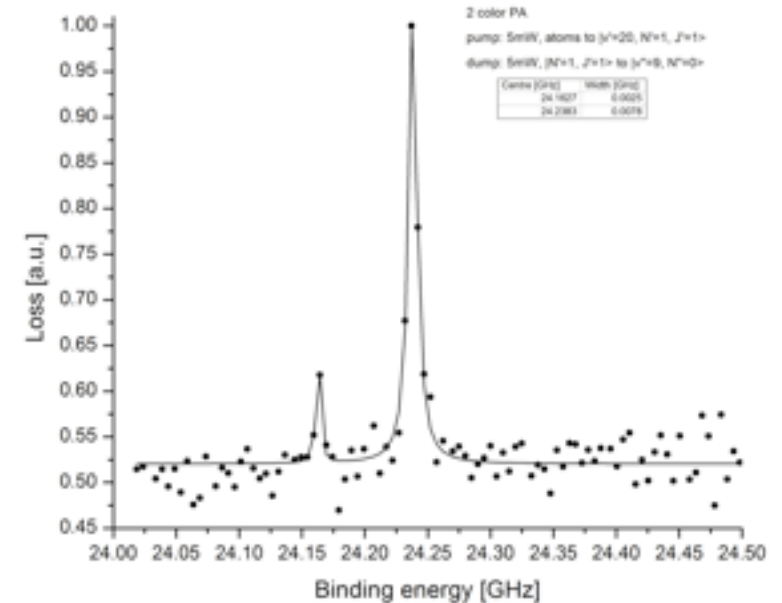
photo-association light

For photo-association, just add additional light

“two-color PA”



atom loss suppressed



$\Omega_2 - \Omega_1$

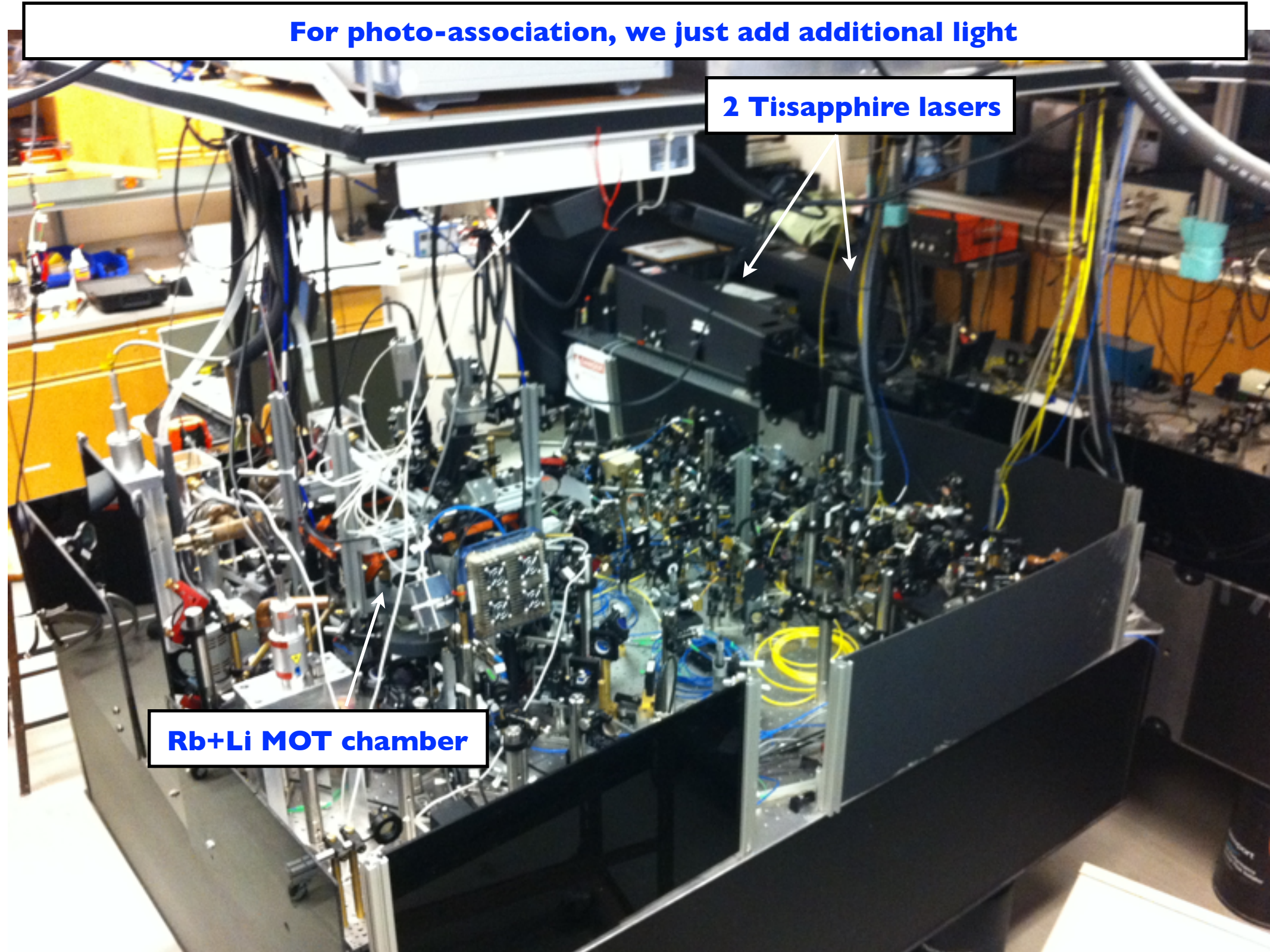
width set by coupling induced by Ω_2

photo-association light

For photo-association, we just add additional light

2 Ti:sapphire lasers

Rb+Li MOT chamber



For photo-association, we just add additional light

frequency comb

2 Ti:sapphire lasers

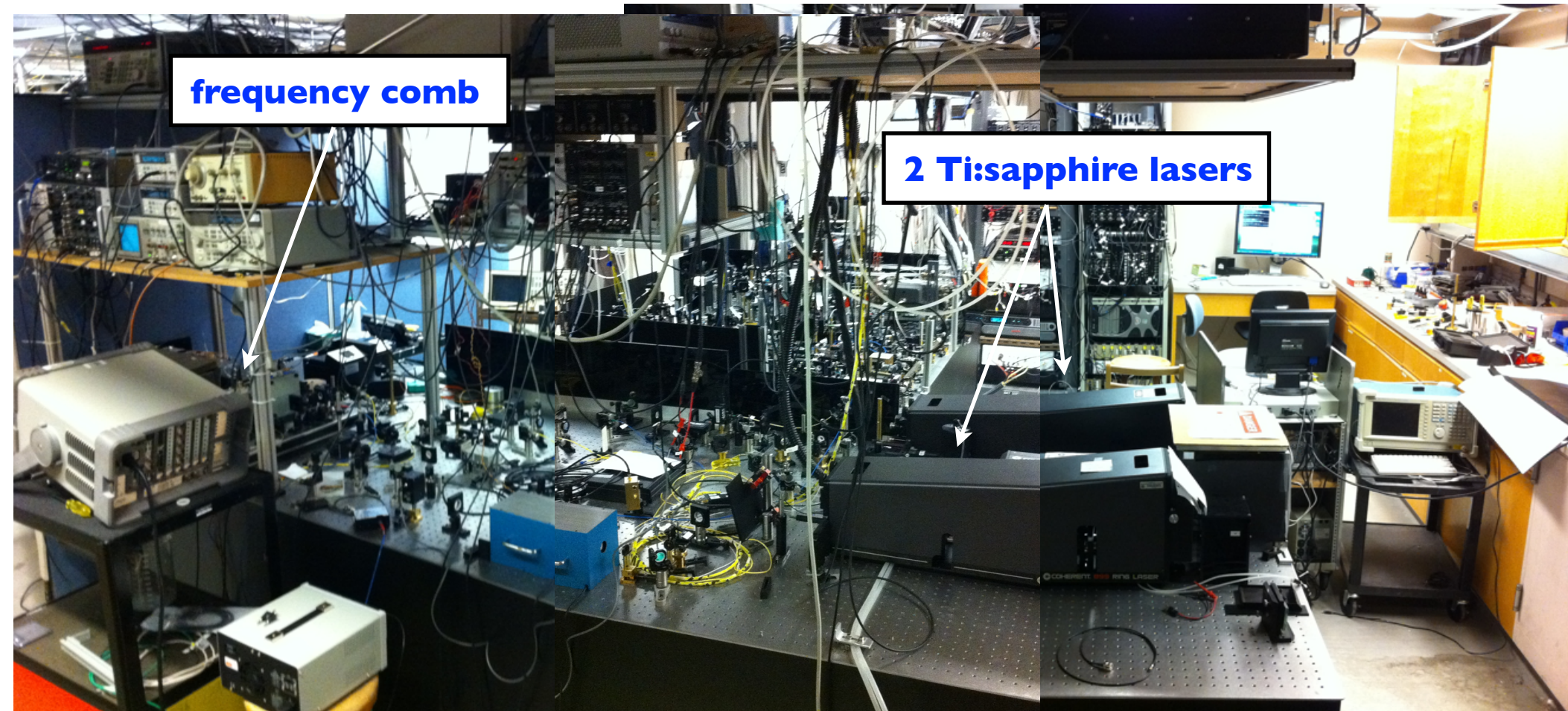
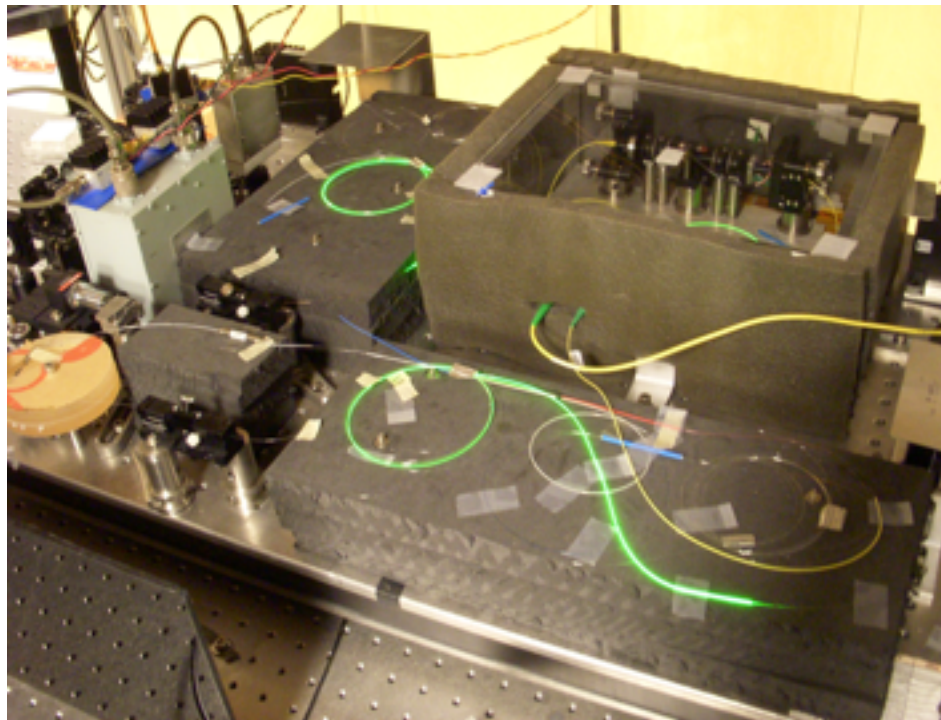
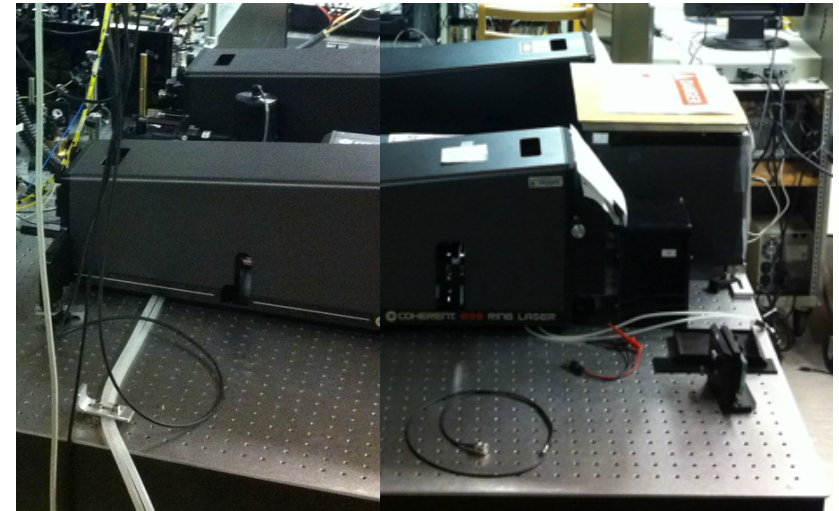


Photo-association (PA) laser system

frequency comb

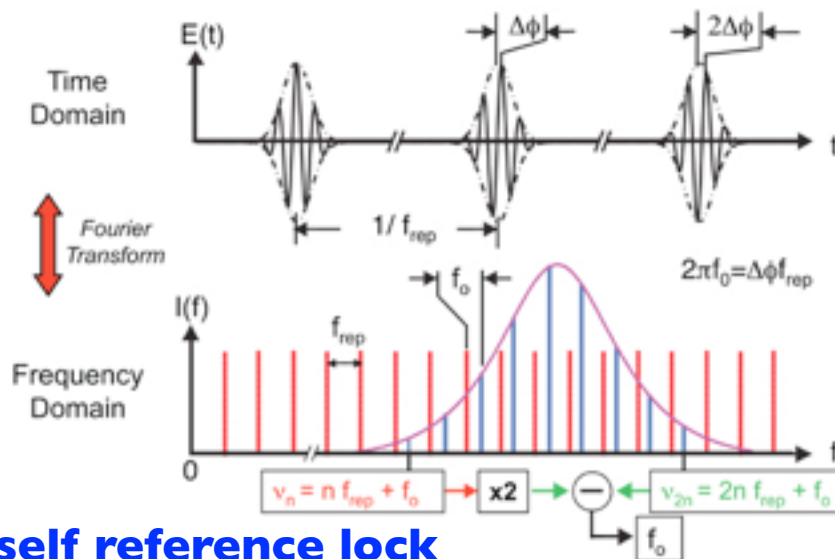


2 Ti:sapphire lasers



When locked to comb:

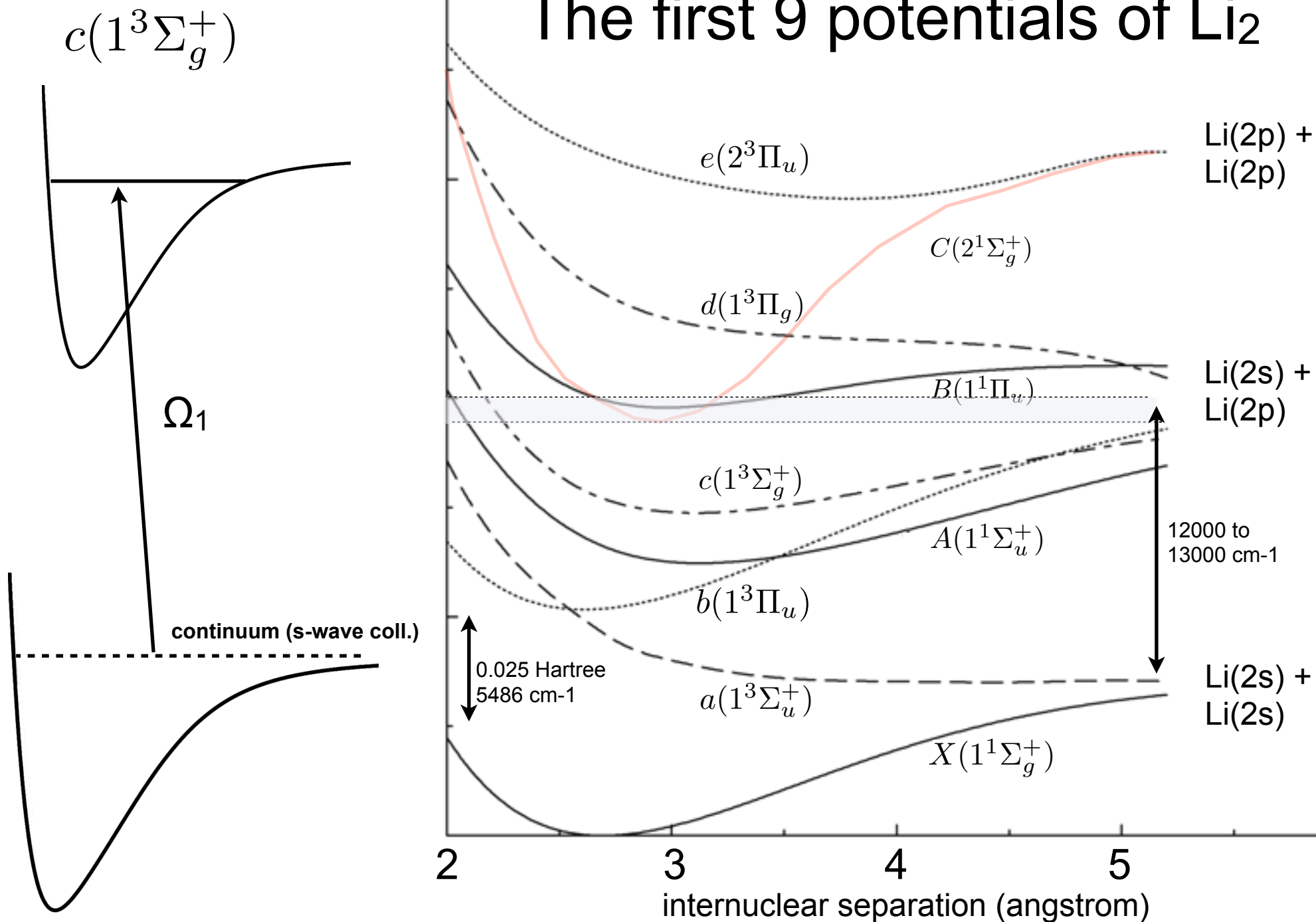
- 1) uncertainty on frequency difference < 10 kHz.
- 2) line width of each Ti:sapphire: ~ 100 kHz
(verified by an independent heterodyne measurement)



self reference lock

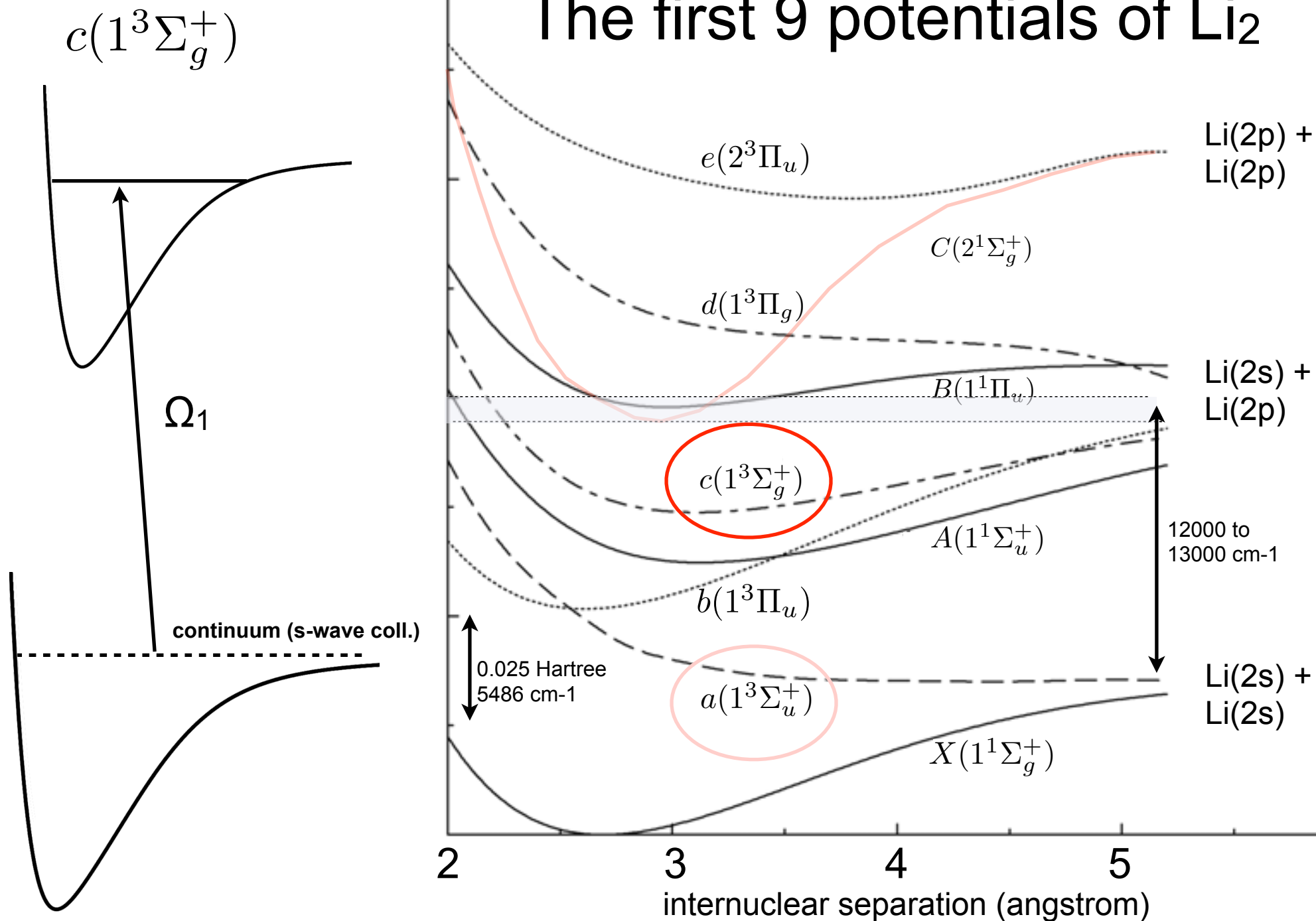
1-color PA spectroscopy results

The first 9 potentials of Li₂

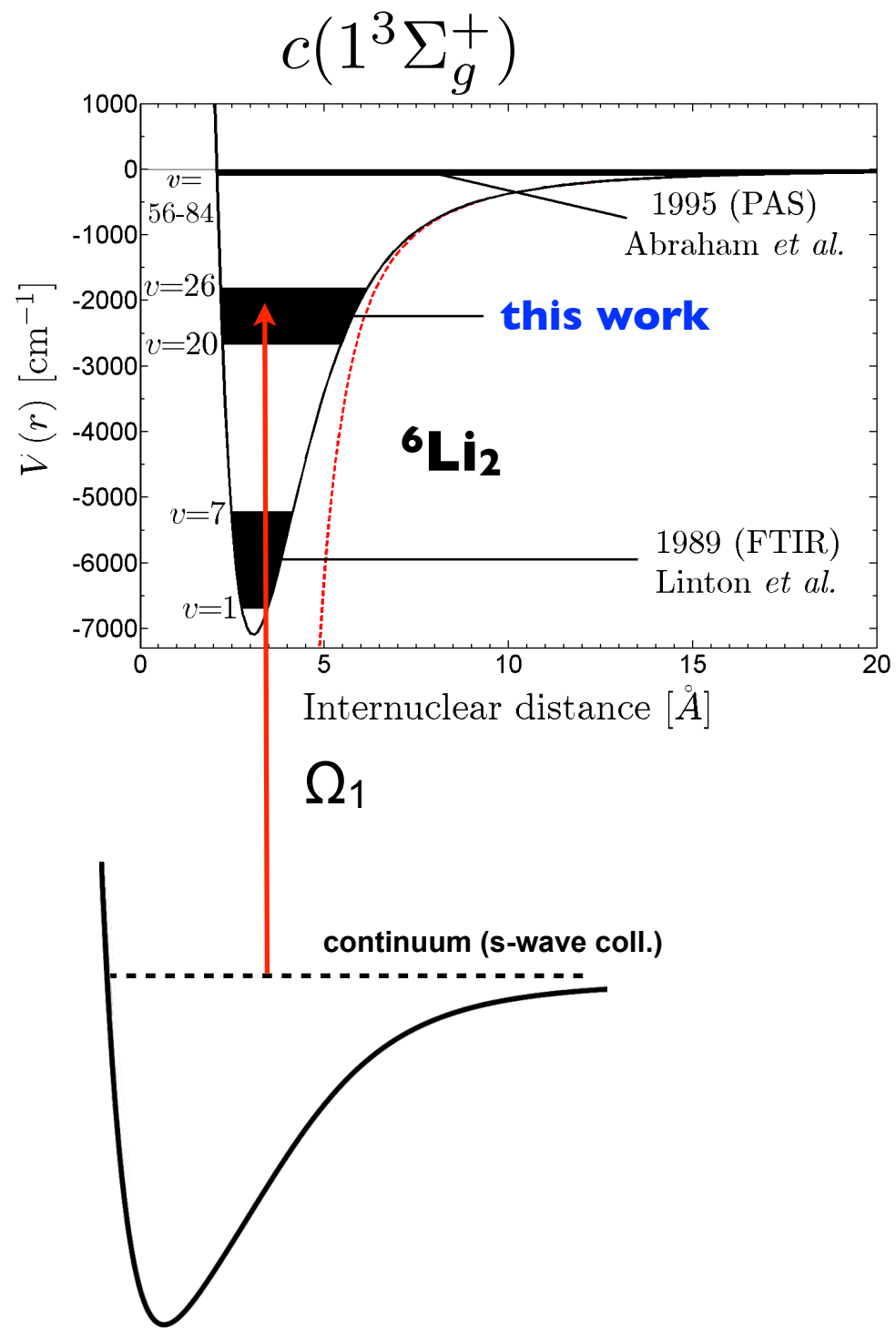


1-color PA spectroscopy results

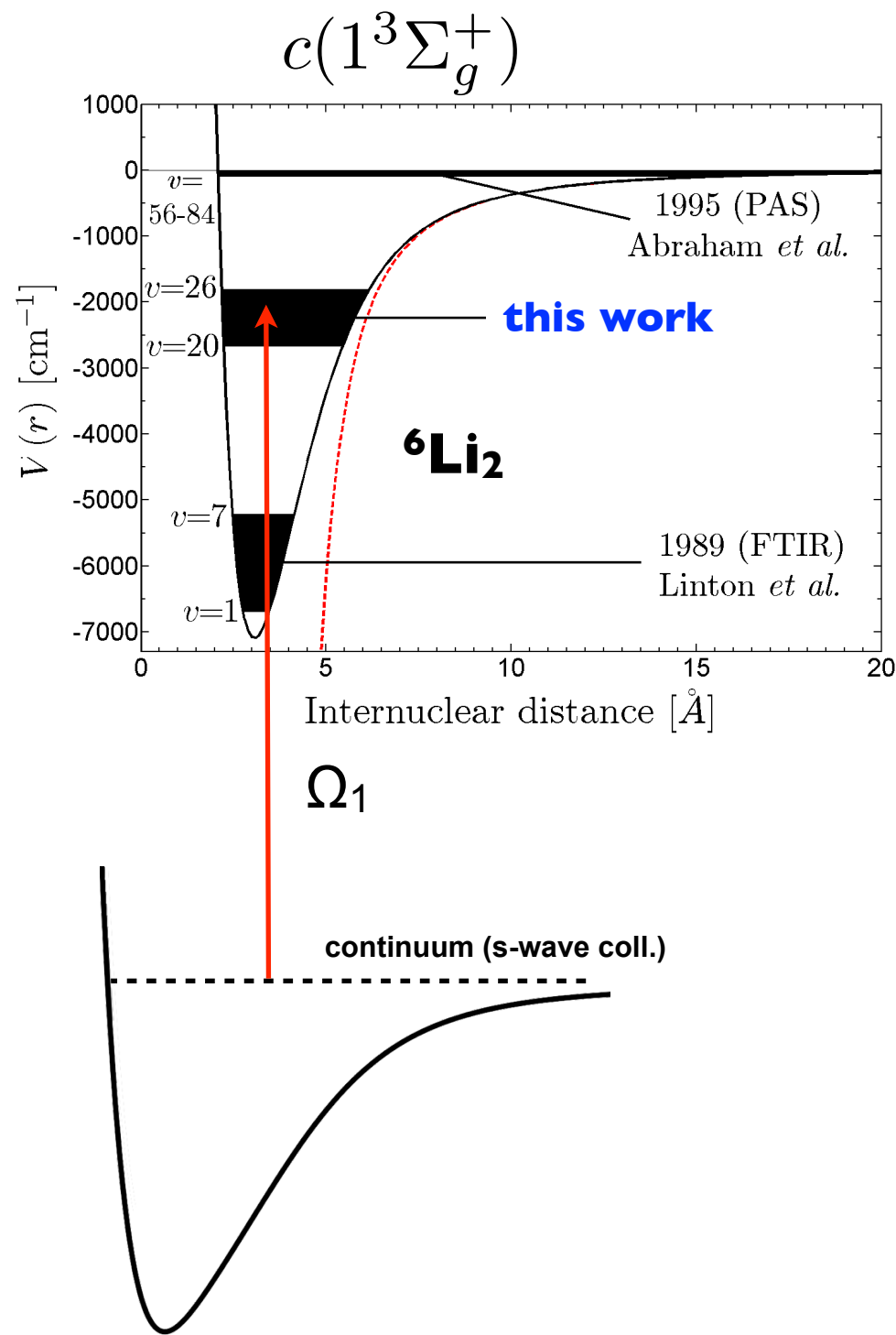
The first 9 potentials of Li₂



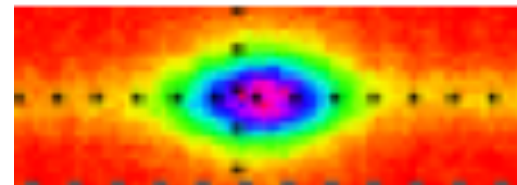
1-color PA spectroscopy results



1-color PA spectroscopy results



0) To minimize broadening and systematic shifts, we evaporate to very low trap powers and photo-associate a 2-component quantum degenerate Fermi gas ($T/T_F \sim 0.4$)



1-color PA spectroscopy results

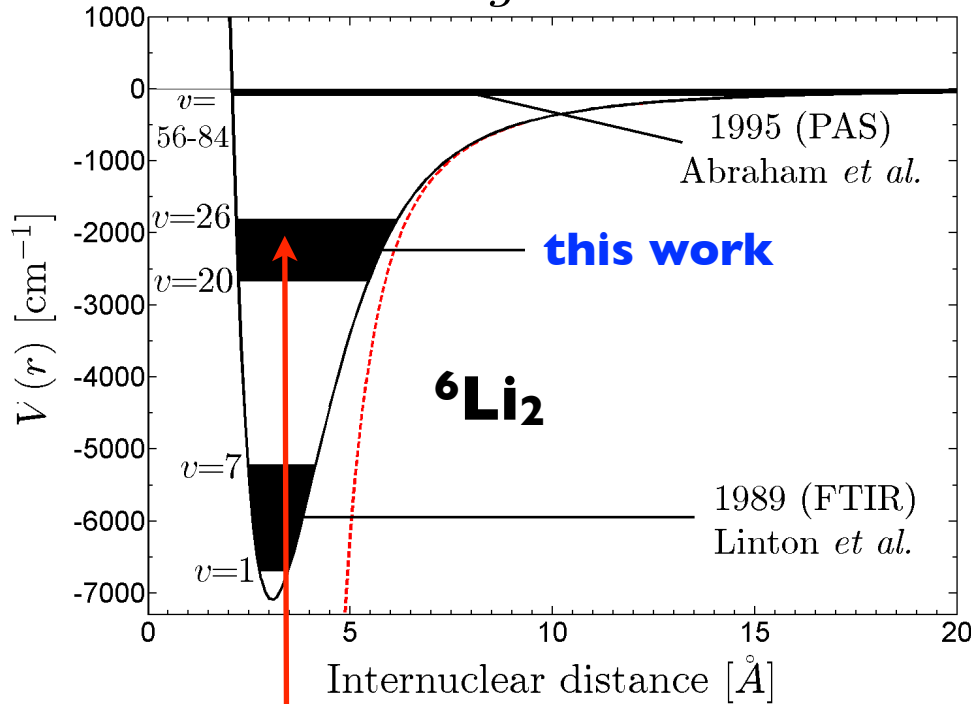
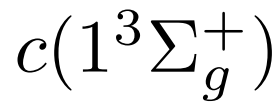
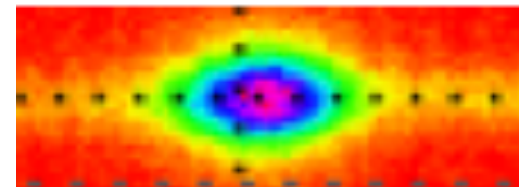
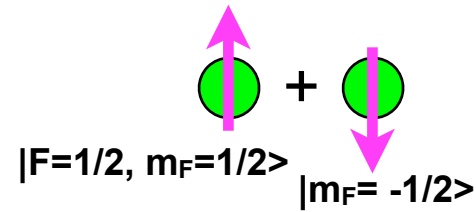
0) To minimize broadening and systematic shifts, we evaporate to very low trap powers and photo-associate a 2-component quantum degenerate Fermi gas ($T/T_F \sim 0.4$)

Initial state:

S-wave collision: $N=0$

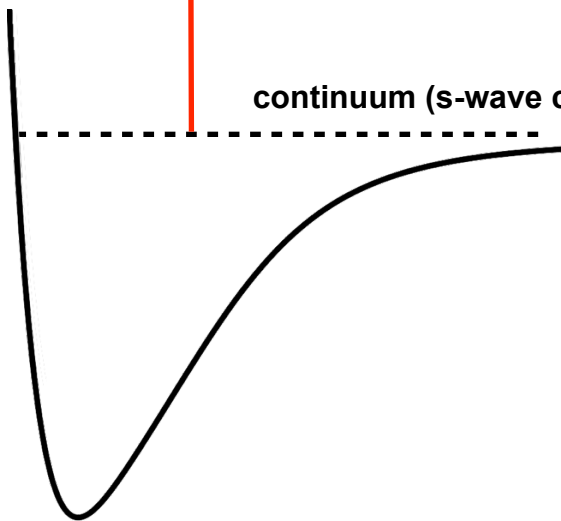
Total spin = 0 ($G = S+I = 0$)

electronic spin = 0, 1

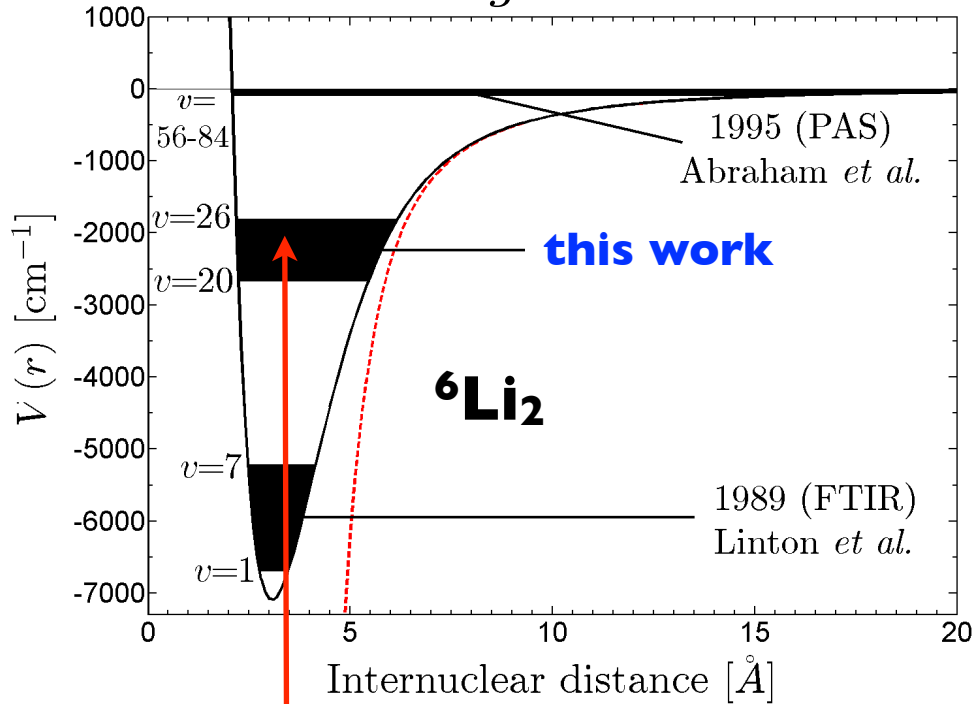
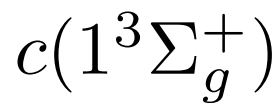


Ω_1

continuum (s-wave coll.)



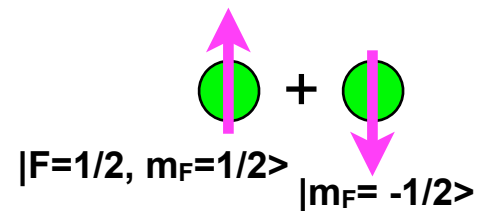
1-color PA spectroscopy results



Ω_1

the relevant initial state is:
 $|S=1, N=0, J=1, l=1\rangle$

0) To minimize broadening and systematic shifts, we evaporate to very low trap powers and photo-associate a 2-component quantum degenerate Fermi gas ($T/T_F \sim 0.4$)



Initial state:

S-wave collision: $N=0$

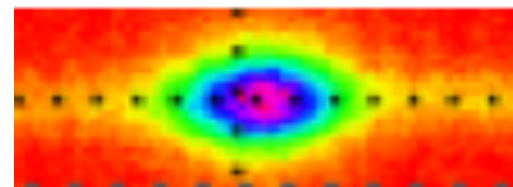
Total spin = 0 ($G = S+I = 0$)

electronic spin = 0, 1

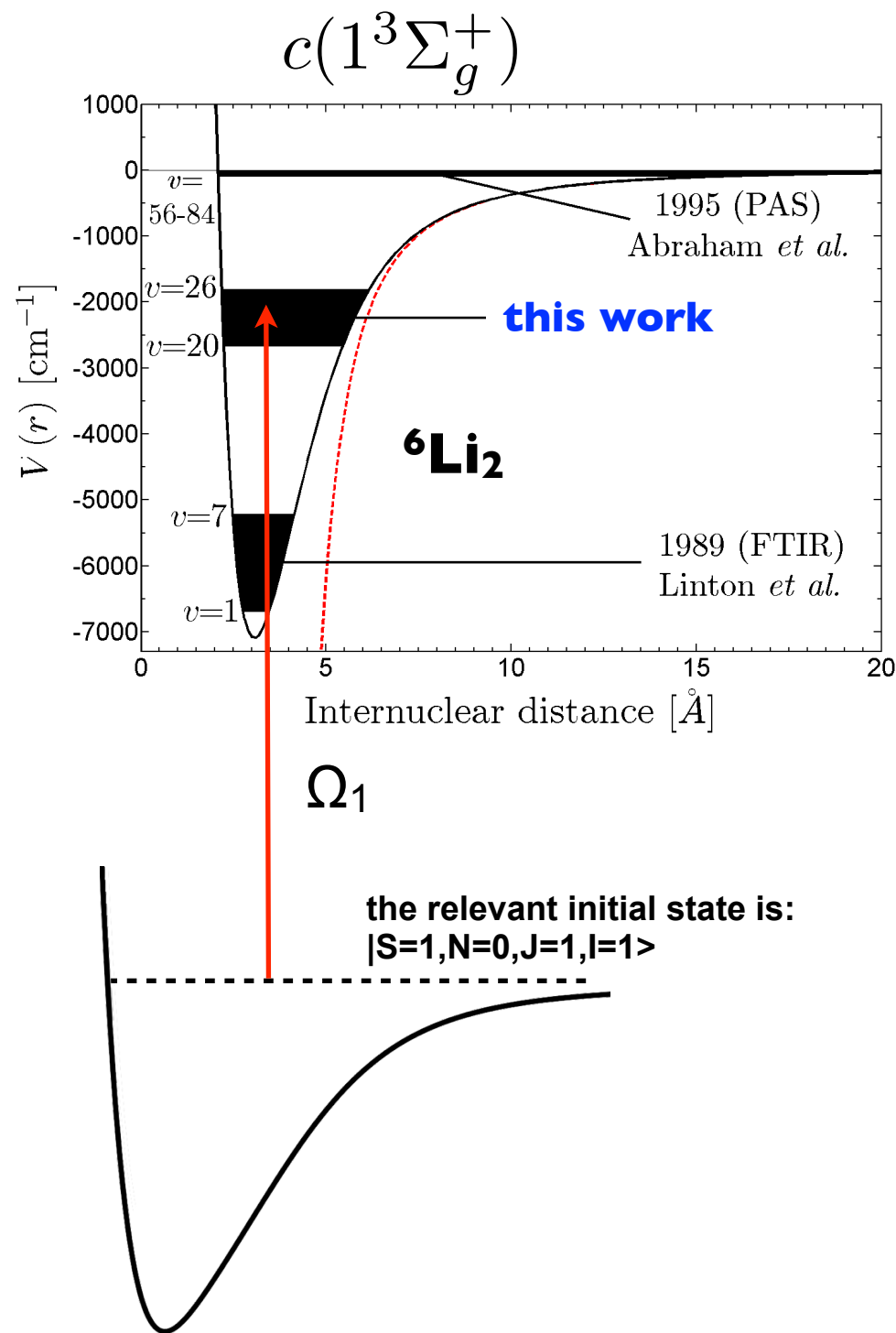
two states are:

$|S=0, N=0, J=0, l=0\rangle$

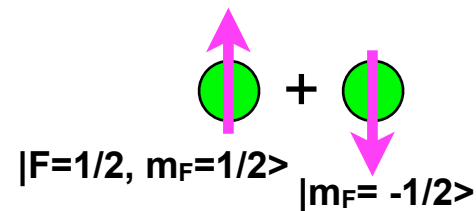
$|S=1, N=0, J=1, l=1\rangle$



1-color PA spectroscopy results



0) To minimize broadening and systematic shifts, we evaporate to very low trap powers and photo-associate a 2-component quantum degenerate Fermi gas ($T/T_F \sim 0.4$)



Initial state:

S-wave collision: $N=0$

Total spin = 0 ($G = S+I = 0$)

electronic spin = 0, 1

two states are:

$|S=0, N=0, J=0, l=0\rangle$

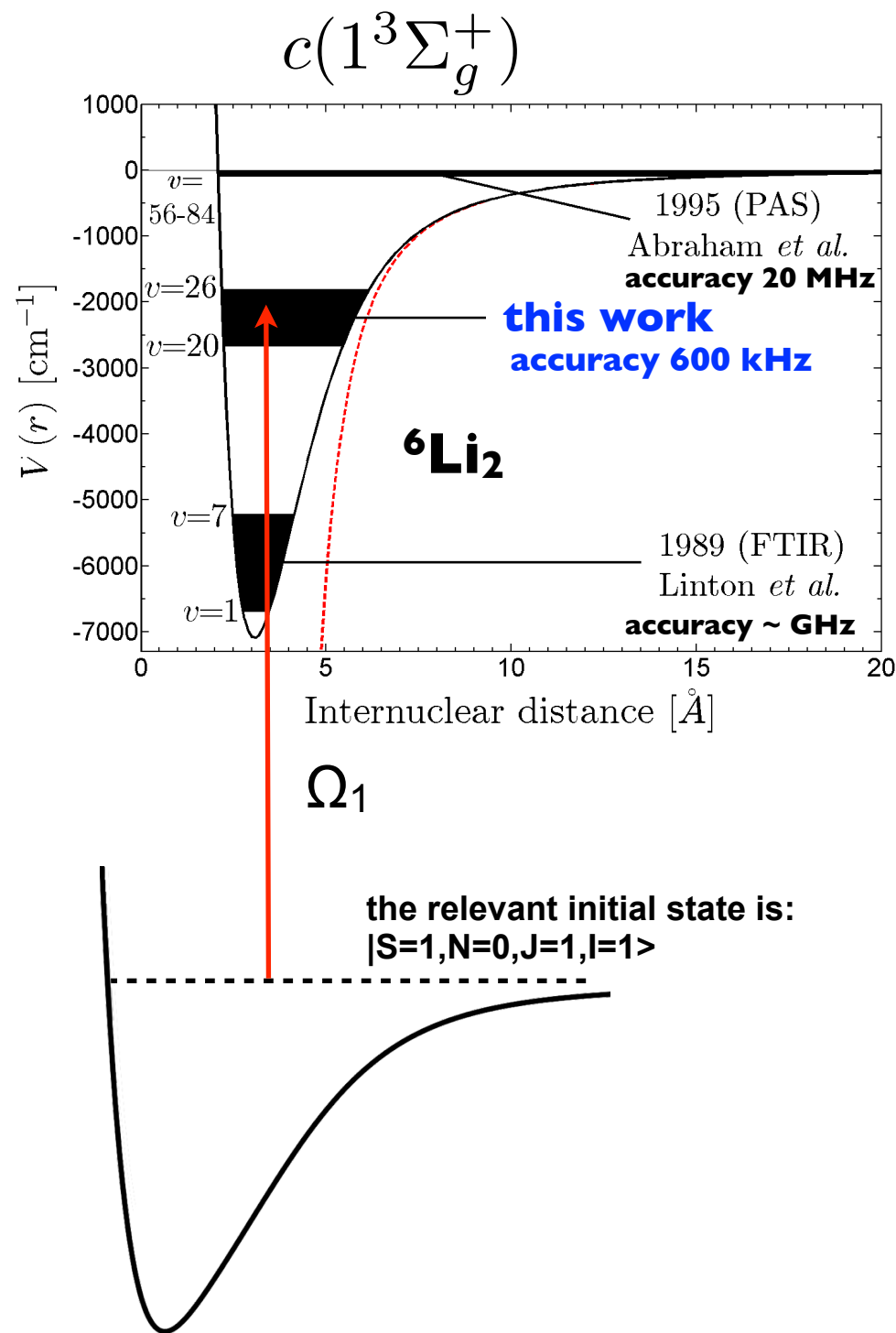
$|S=1, N=0, J=1, l=1\rangle$

1) We measure 7 levels from $v'=20$ to $v'=26$ with an absolute accuracy of ± 600 kHz (± 0.00002 cm^{-1})

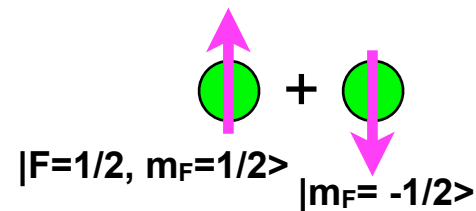
2) For each vibrational level, we observe the $N'=1$ level AND using a p-wave FR at 187.4 G, we observe the $N'=0$ and $N'=2$ rotational levels

3) We also characterized the systematic shifts due to magnetic field, and the residual AC Stark shifts of the PA laser and the trapping laser.

1-color PA spectroscopy results



0) To minimize broadening and systematic shifts, we evaporate to very low trap powers and photo-associate a 2-component quantum degenerate Fermi gas ($T/T_F \sim 0.4$)



Initial state:

S-wave collision: $N=0$

Total spin = 0 ($G = S+I = 0$)

electronic spin = 0, 1

two states are:

$|S=0, N=0, J=0, l=0\rangle$

$|S=1, N=0, J=1, l=1\rangle$

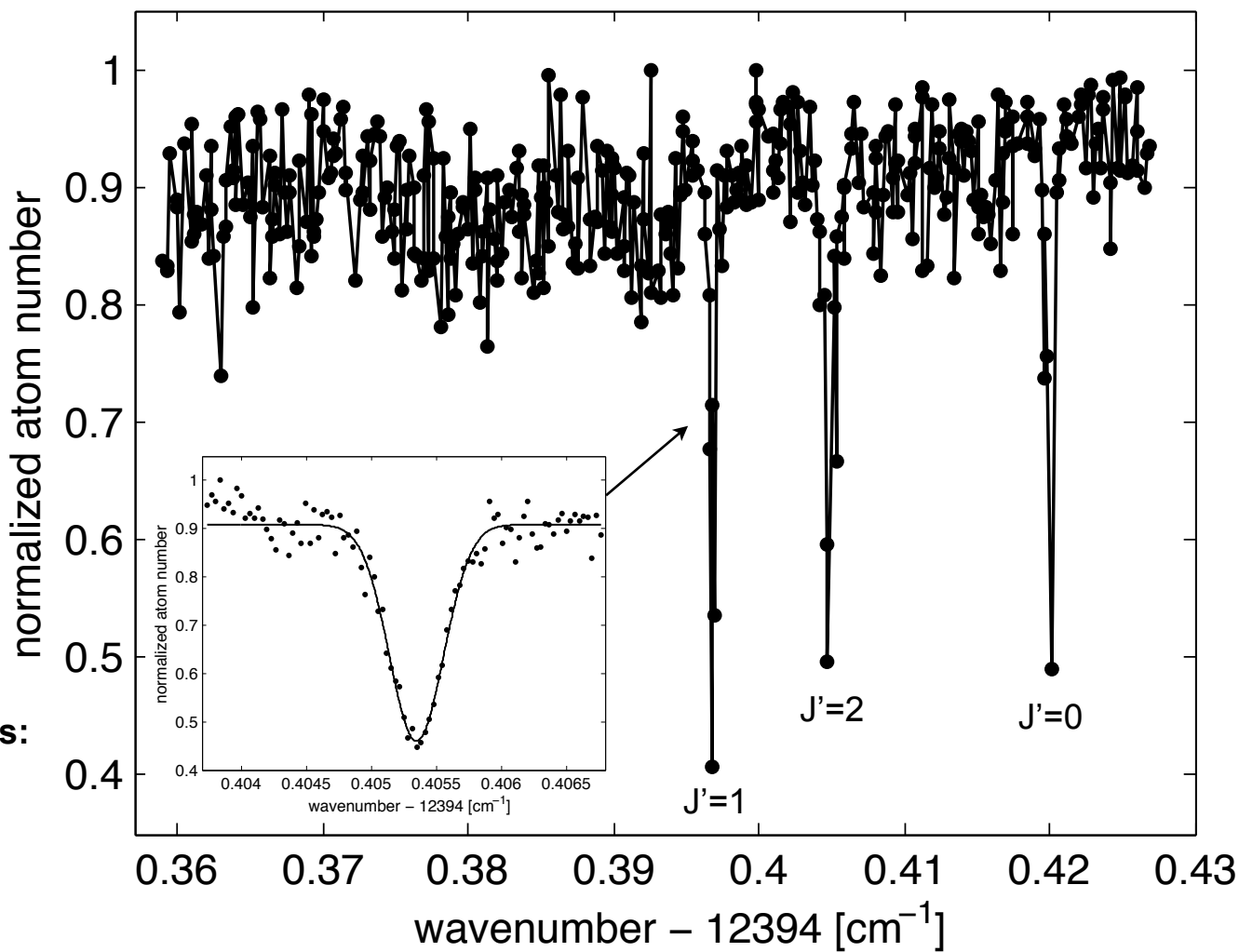
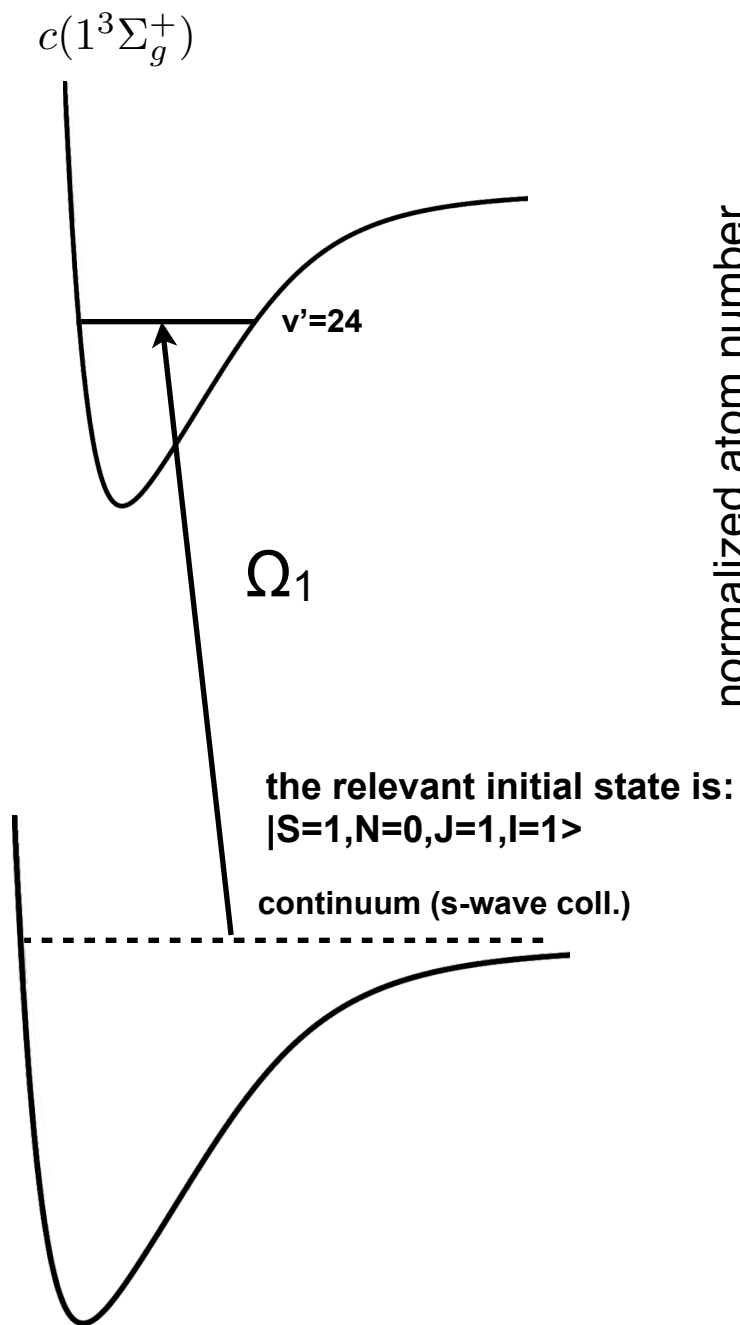
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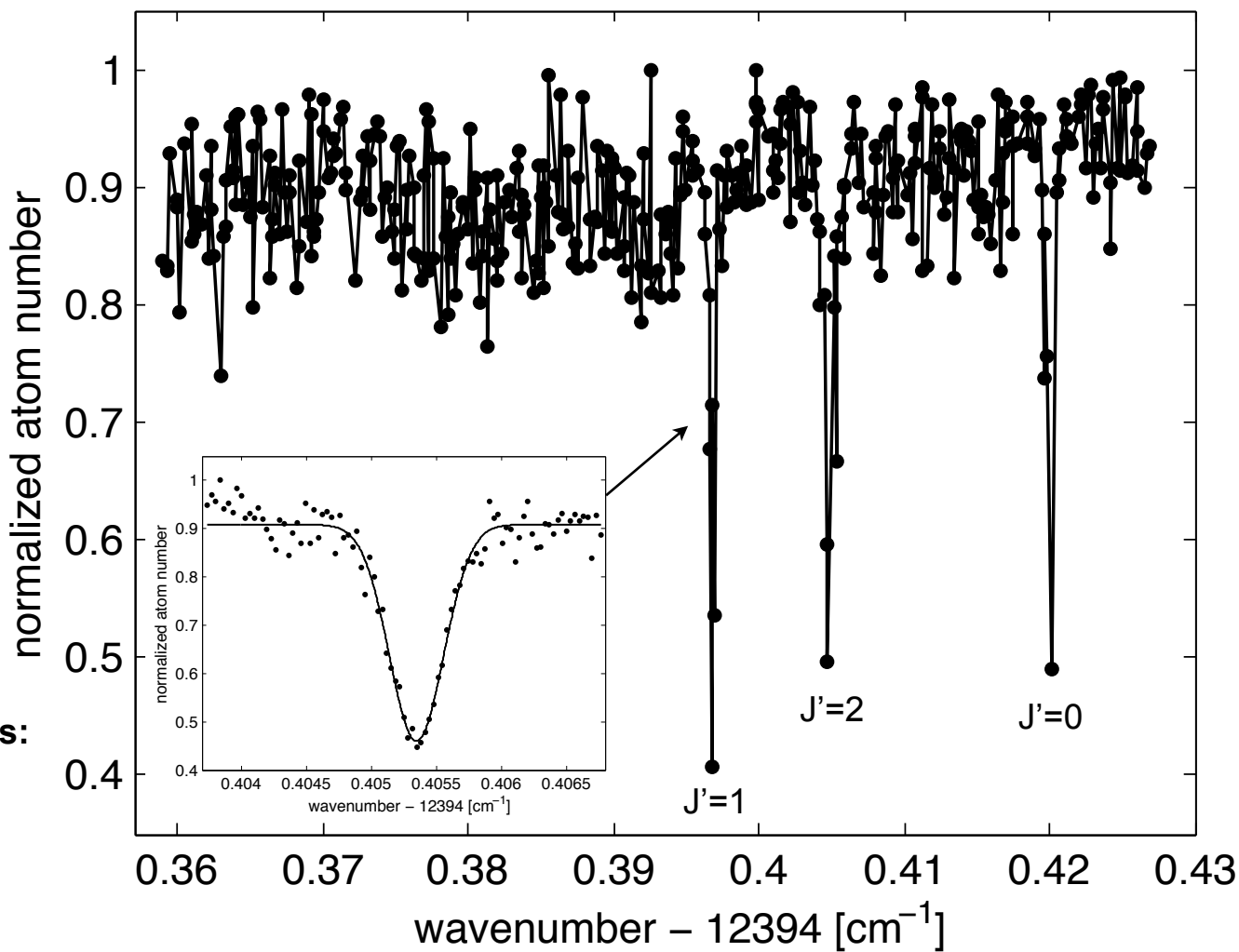
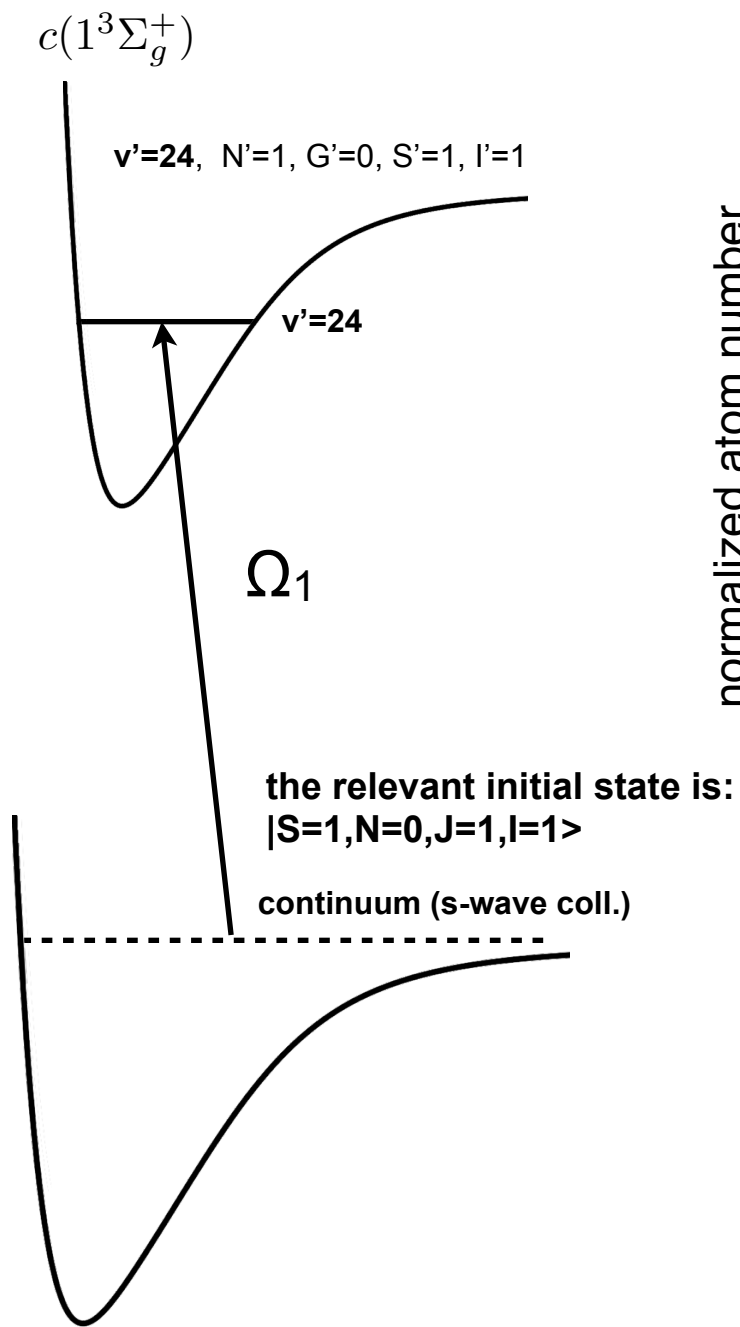
1-color PA spectroscopy results

We resolve three features for each vibrational level



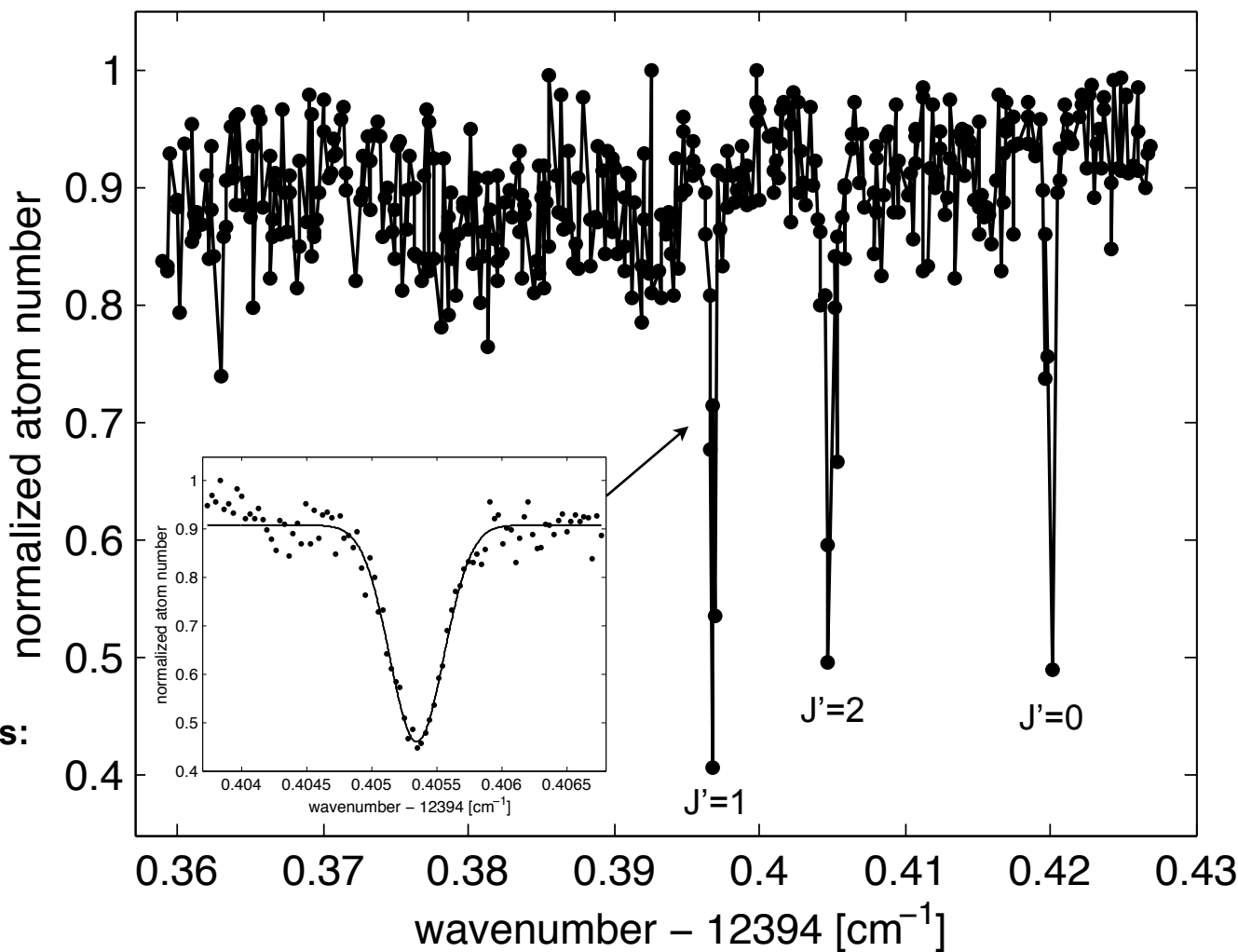
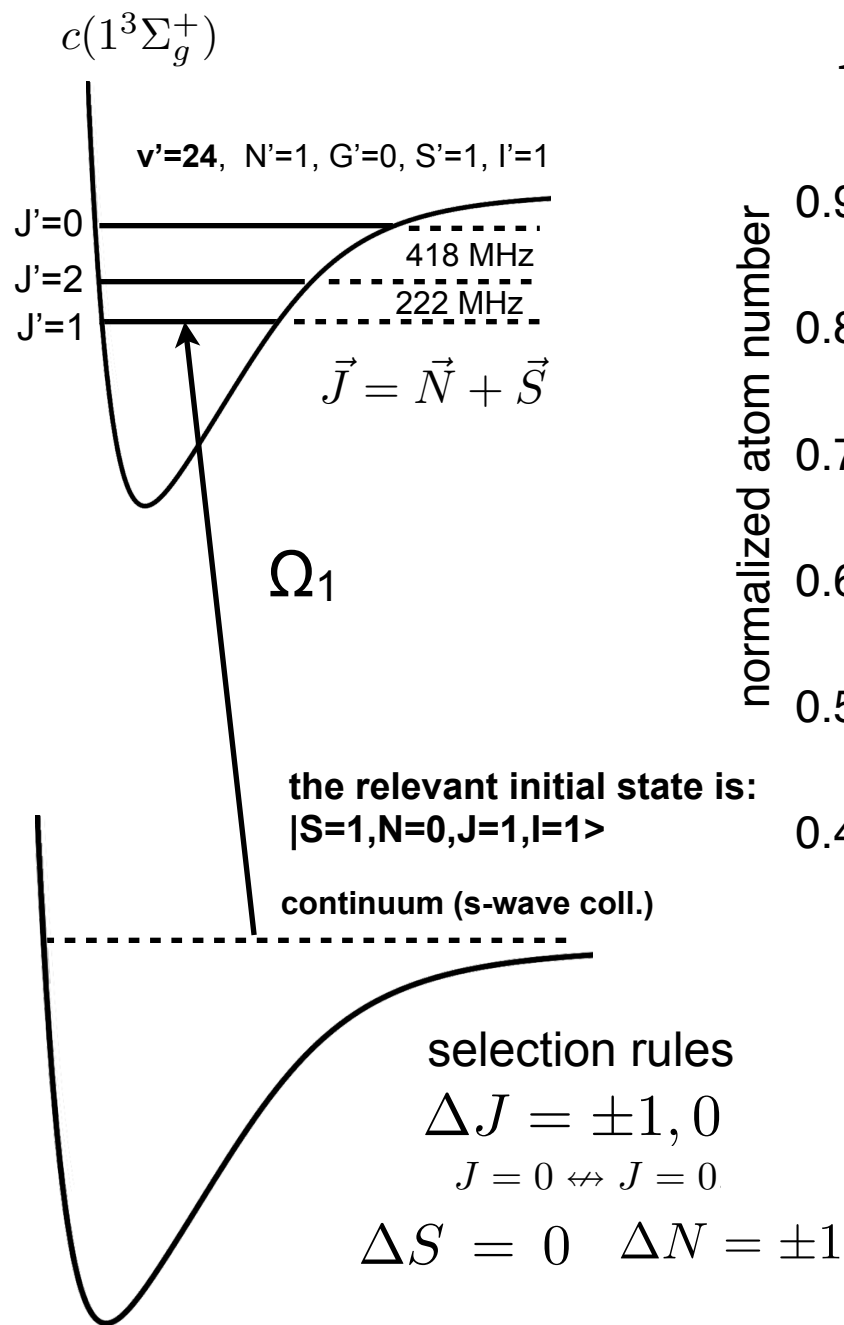
1-color PA spectroscopy results

We resolve three features for each vibrational level



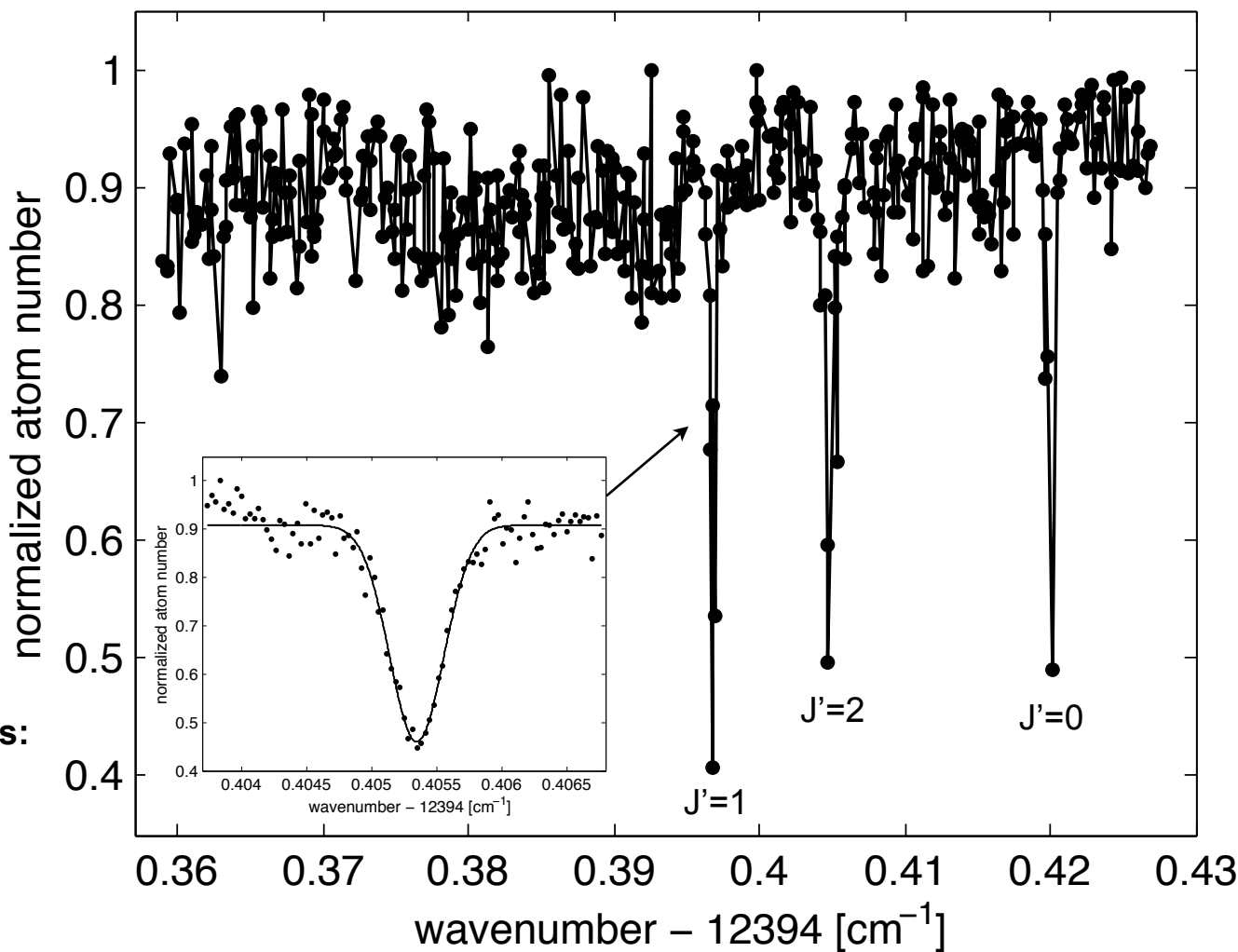
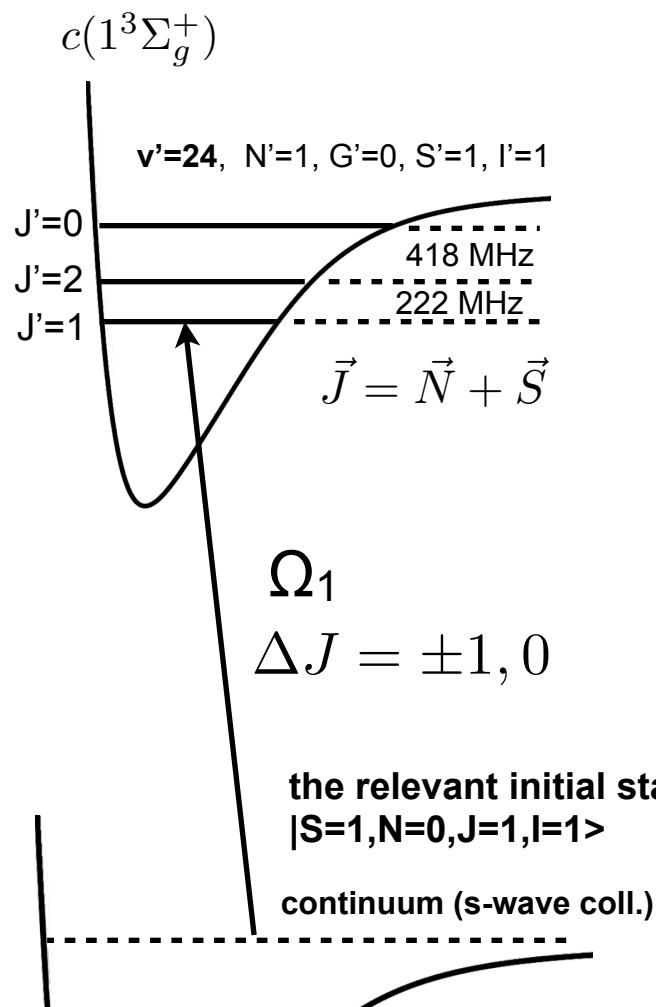
1-color PA spectroscopy results

We resolve three features for each vibrational level split by spin-spin and spin-rotation coupling



1-color PA spectroscopy results

We resolve three features for each vibrational level split by spin-spin and spin-rotation coupling

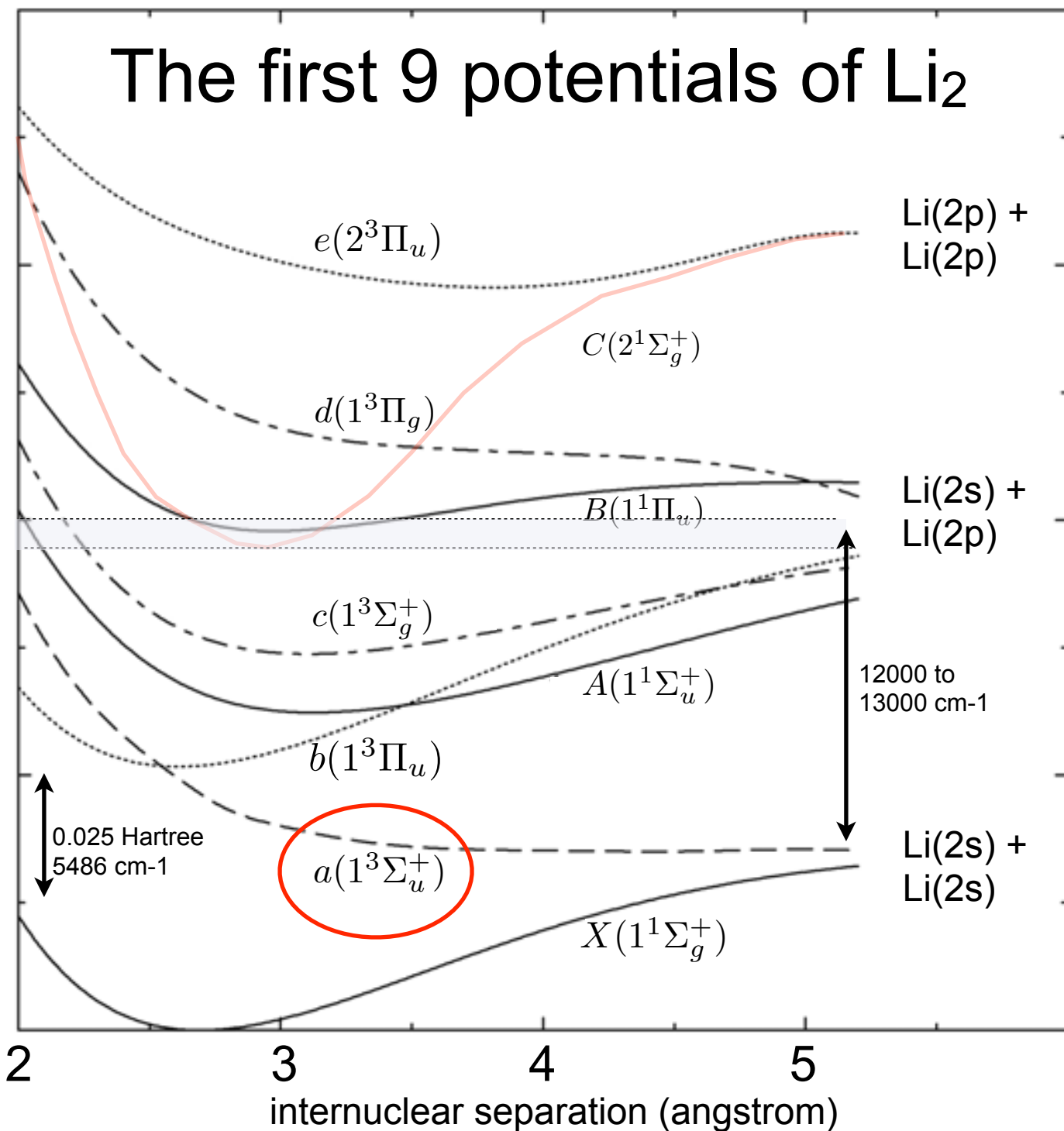
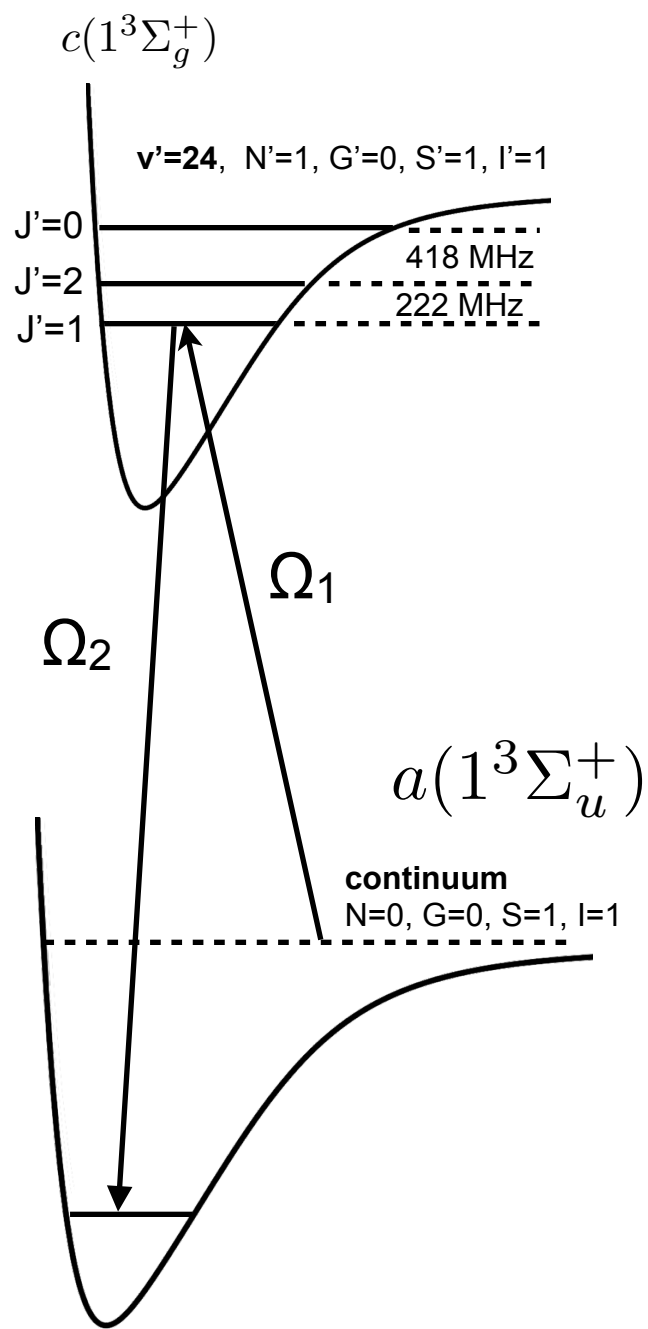


$$\hat{H}_{\text{spin-spin}} = 2\lambda_v [\hat{S}_z^2 - \hat{S}^2/3]$$

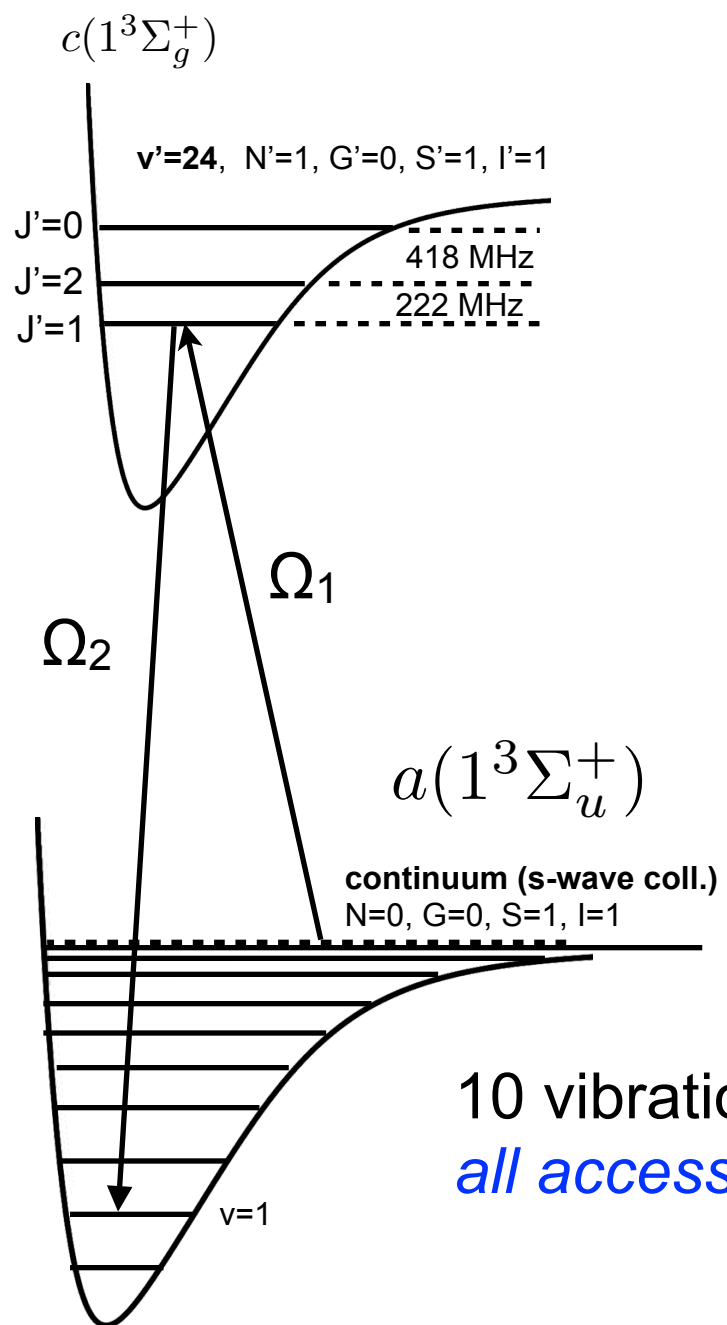
$$\hat{H}_{\text{spin-rot}} = \gamma_v \vec{N} \cdot \vec{S}$$

v'	λ_v (MHz)	γ_v (MHz)
20	-348.2	-14.5
21	-339.4	-14.5
22	-331.1	-14.7
23	-321.7	-14.2
24	-312.2	-14.4
25	-309.6	-14.0
26	-294.3	-14.3

2-color PA spectroscopy : ground state spectroscopy



Triplet ground state Li₂



Motivations:

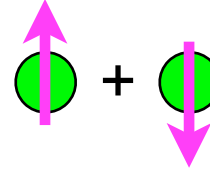
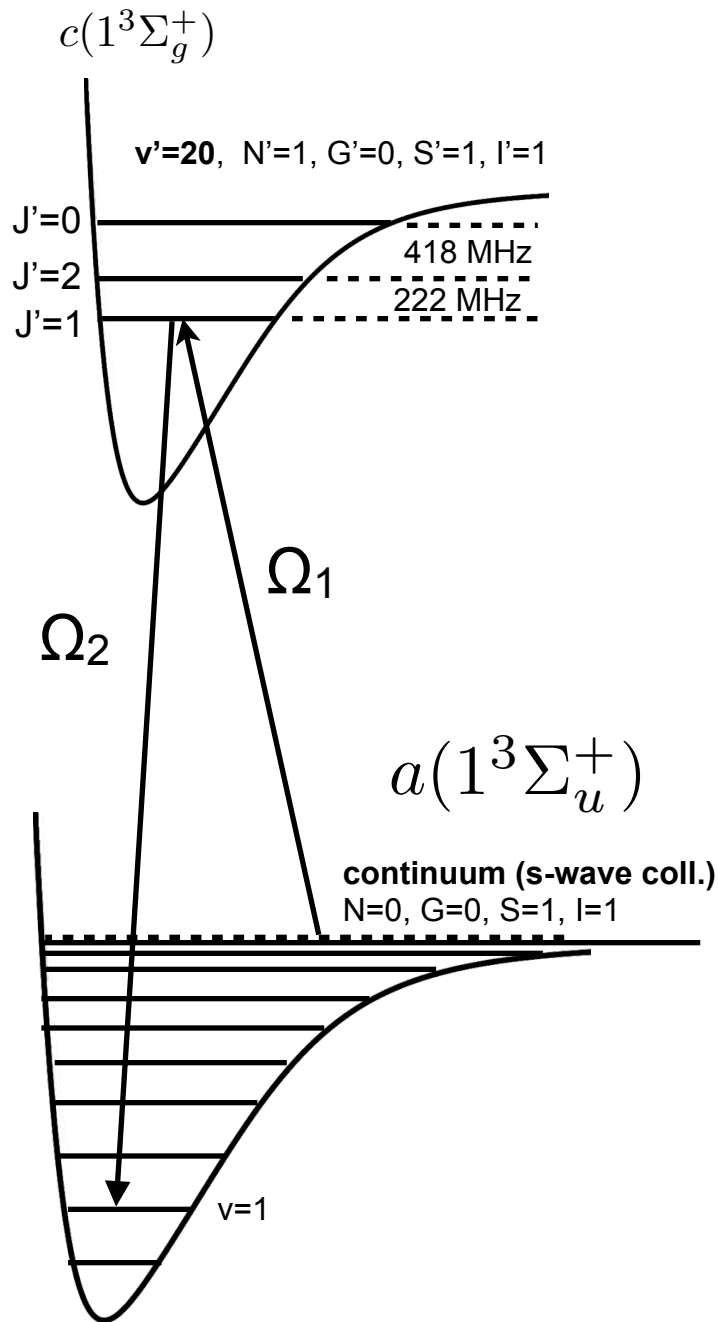
- 1) Technical: we can access with our lasers.
- 2) It has a magnetic moment (Molecular FRs !)
- 3) Make a BEC of ground state molecules
- 4) Measure the “Spin blockade”
- 5) Study collision properties of ultra-cold superrotors (collaboration with Valery Milner)
- 6) Stepping stone for making triplet LiRb

10 vibrational levels

all accessible with our laser system

2-color PA spectroscopy results : here Ω_1 is fixed and Ω_2 is scanned

In the data that follows, the loss of atoms due to photo-association induced by Ω_1 is **suppressed** by the AC Stark shift of the excited state levels induced by Ω_2 . We specify the excited state v' number and the J' number and we specify the ground state v'' number and N'' number.

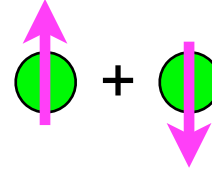
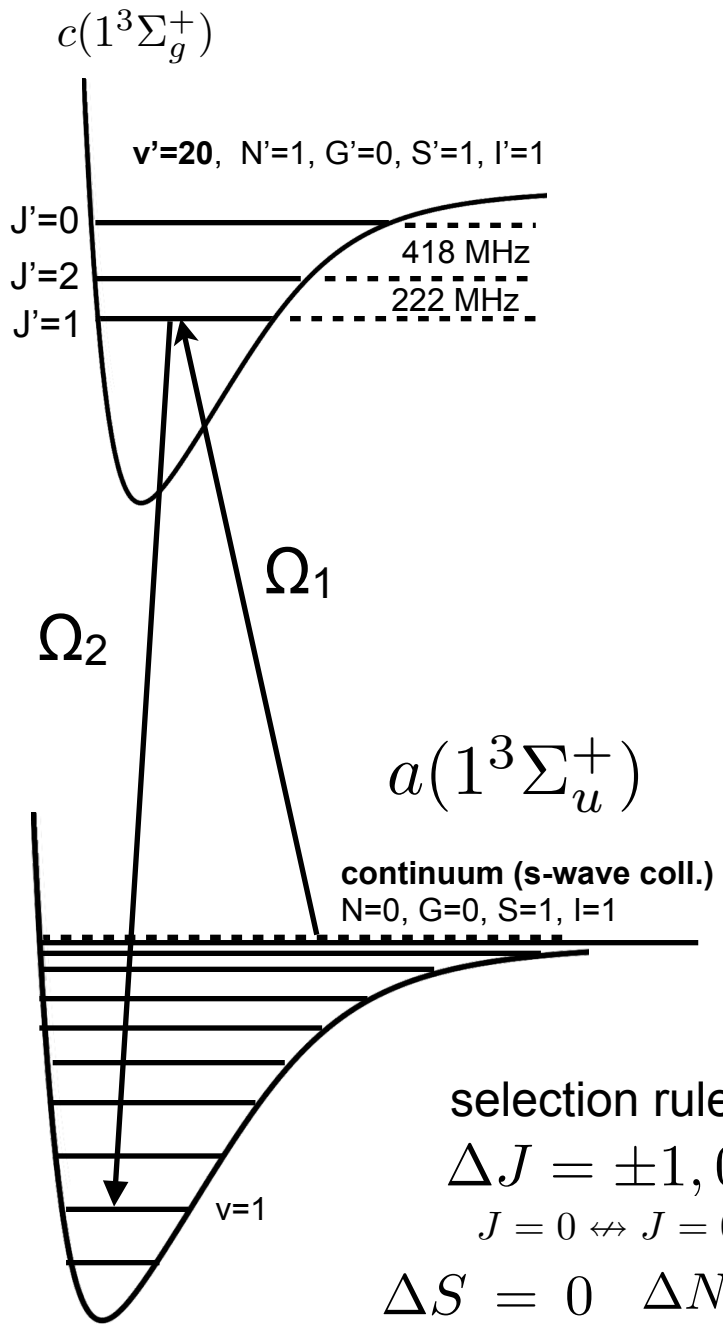


the relevant initial state is:
 $|S=1, N=0, J=1, l=1\rangle$

the intermediate state is:
 $|S=1, N=1, J=1, l=1\rangle$

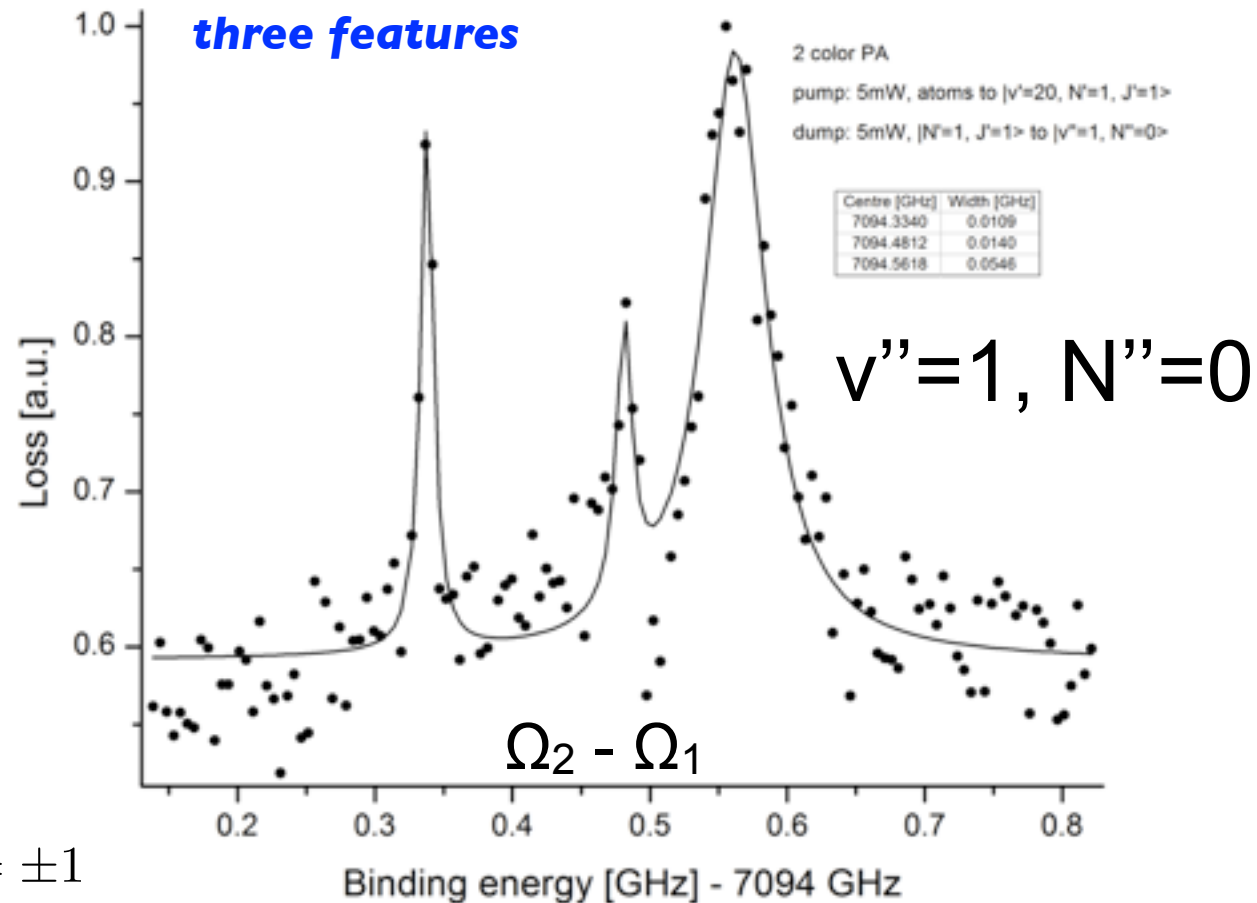
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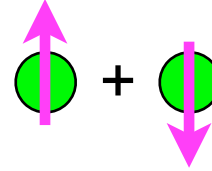
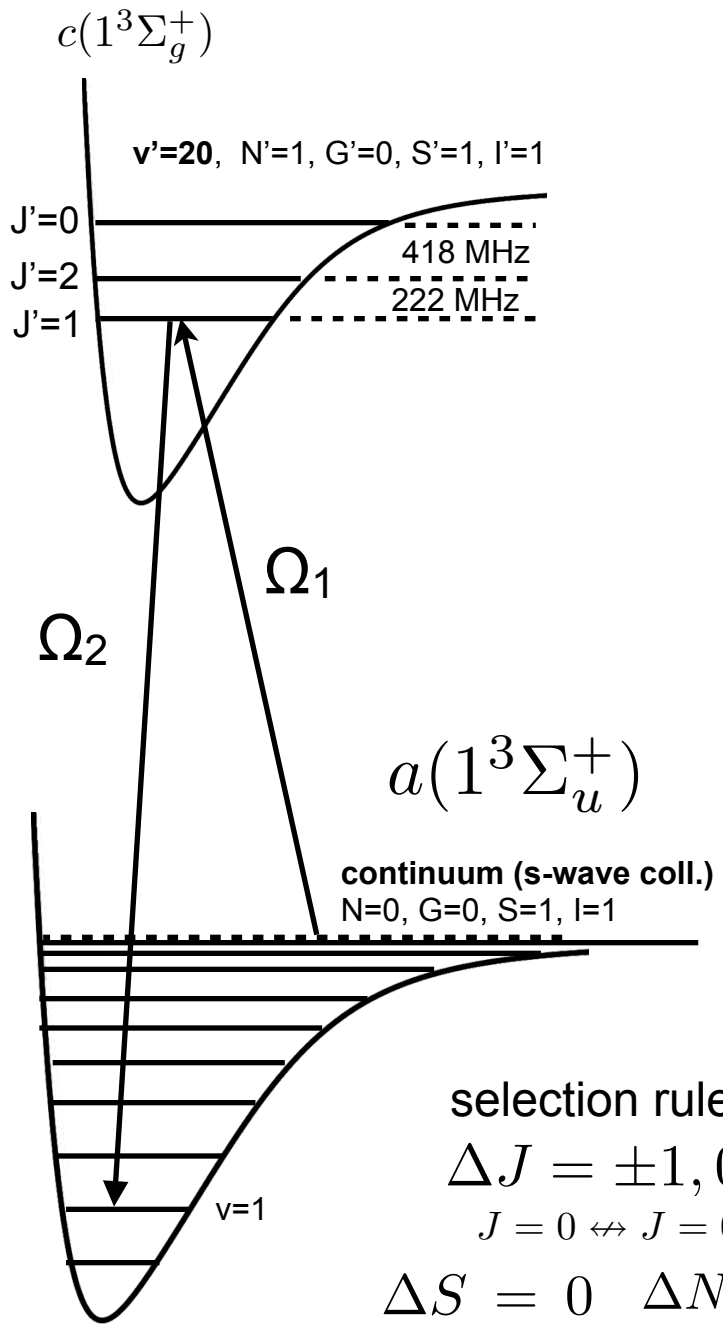
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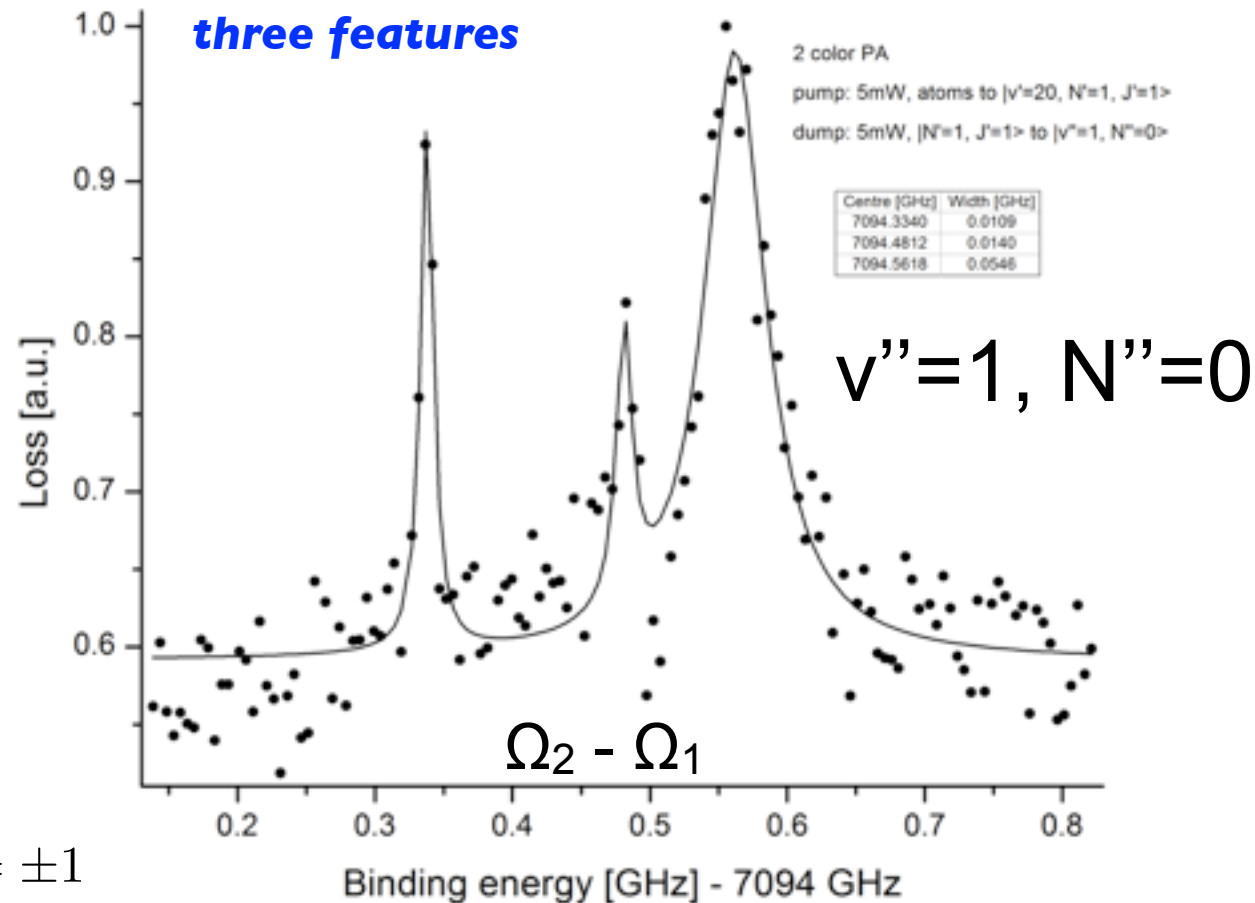
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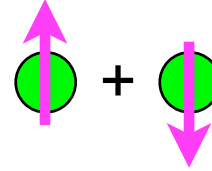
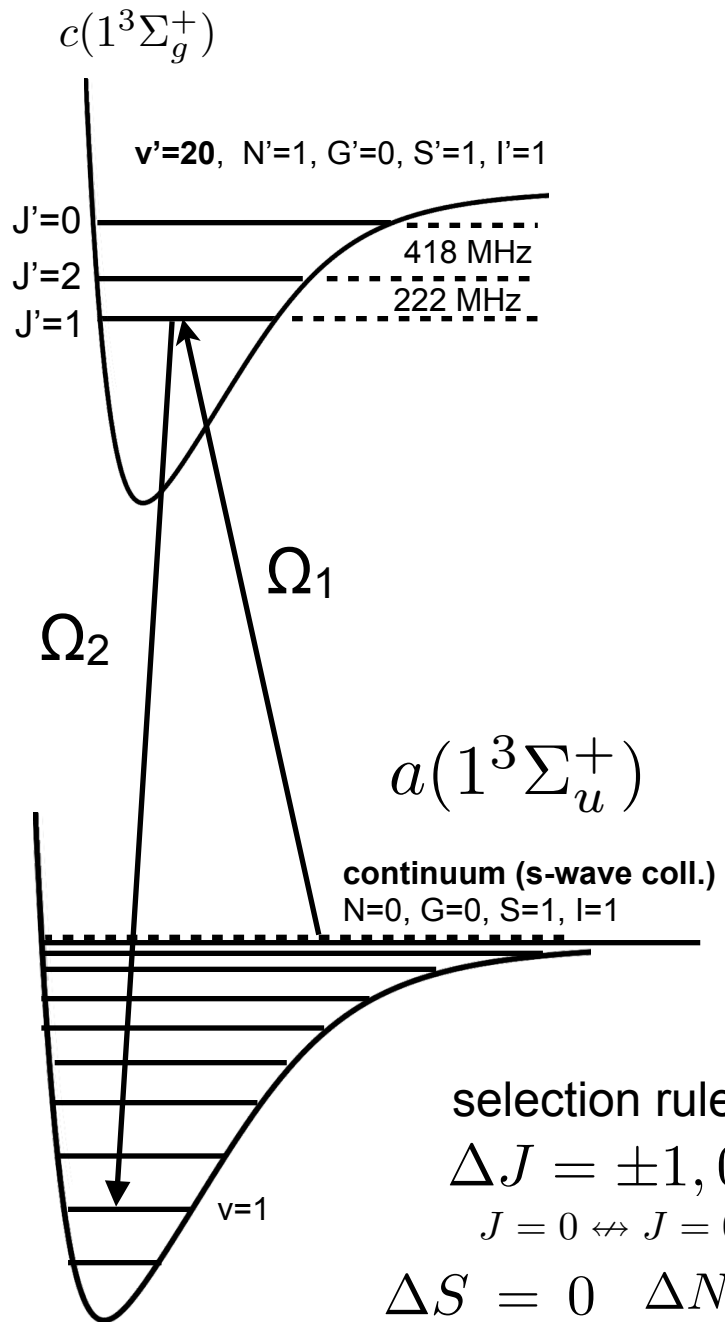
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the final state is:
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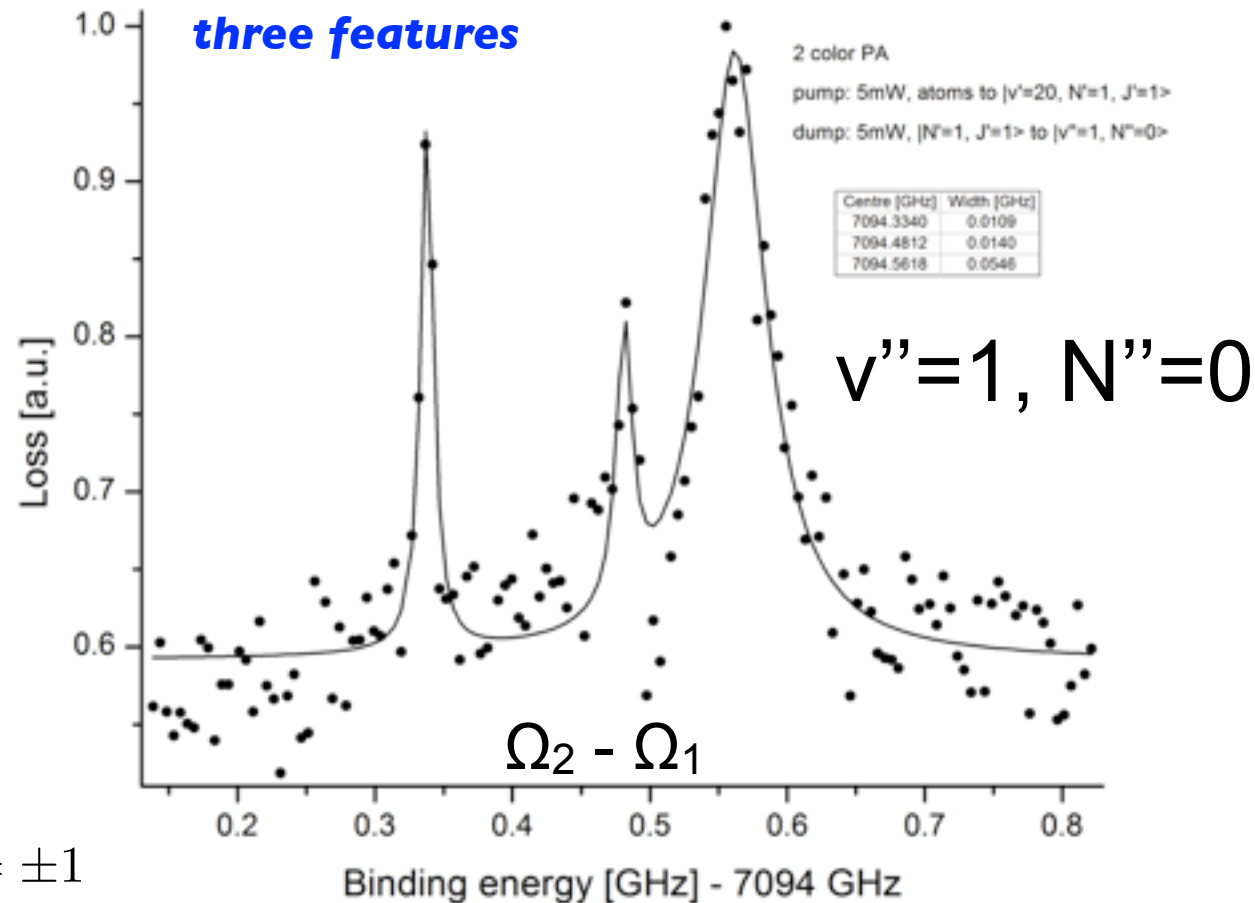
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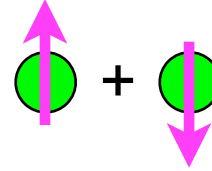
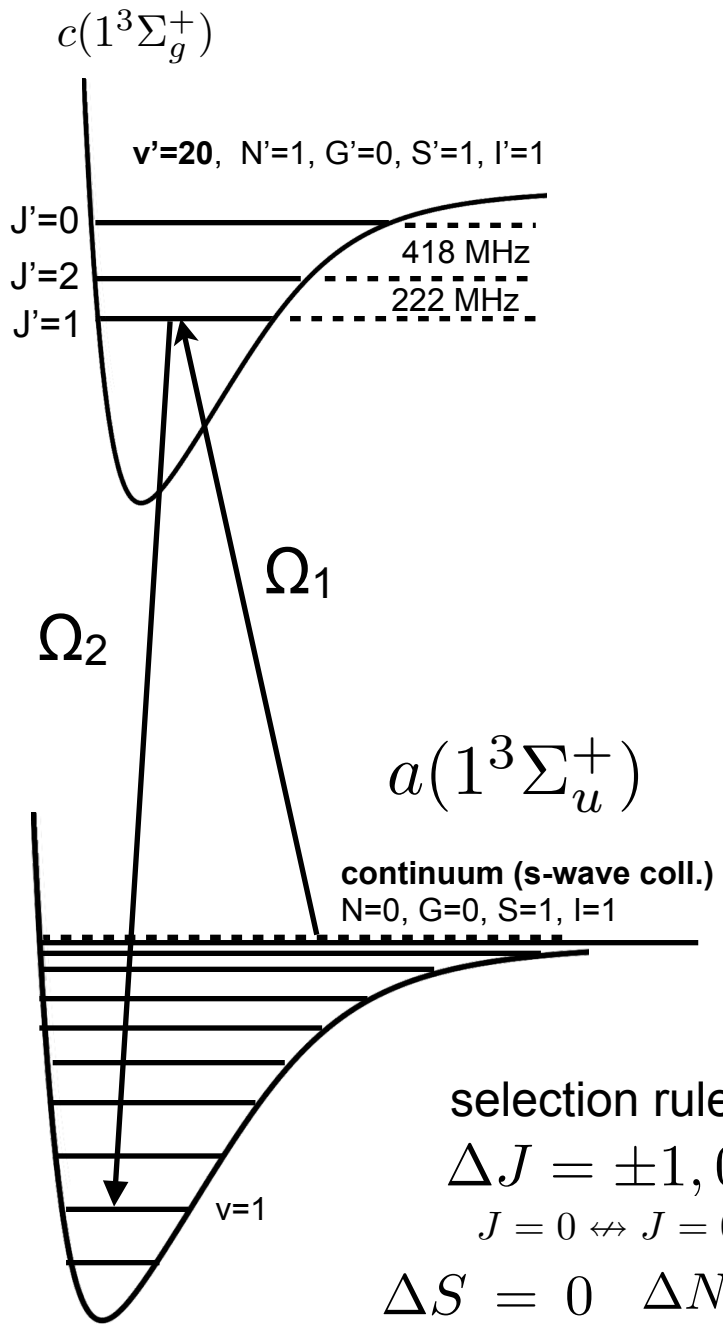
$F = I + S = 0, 1, 2$ (expect three hyperfine levels)

three features



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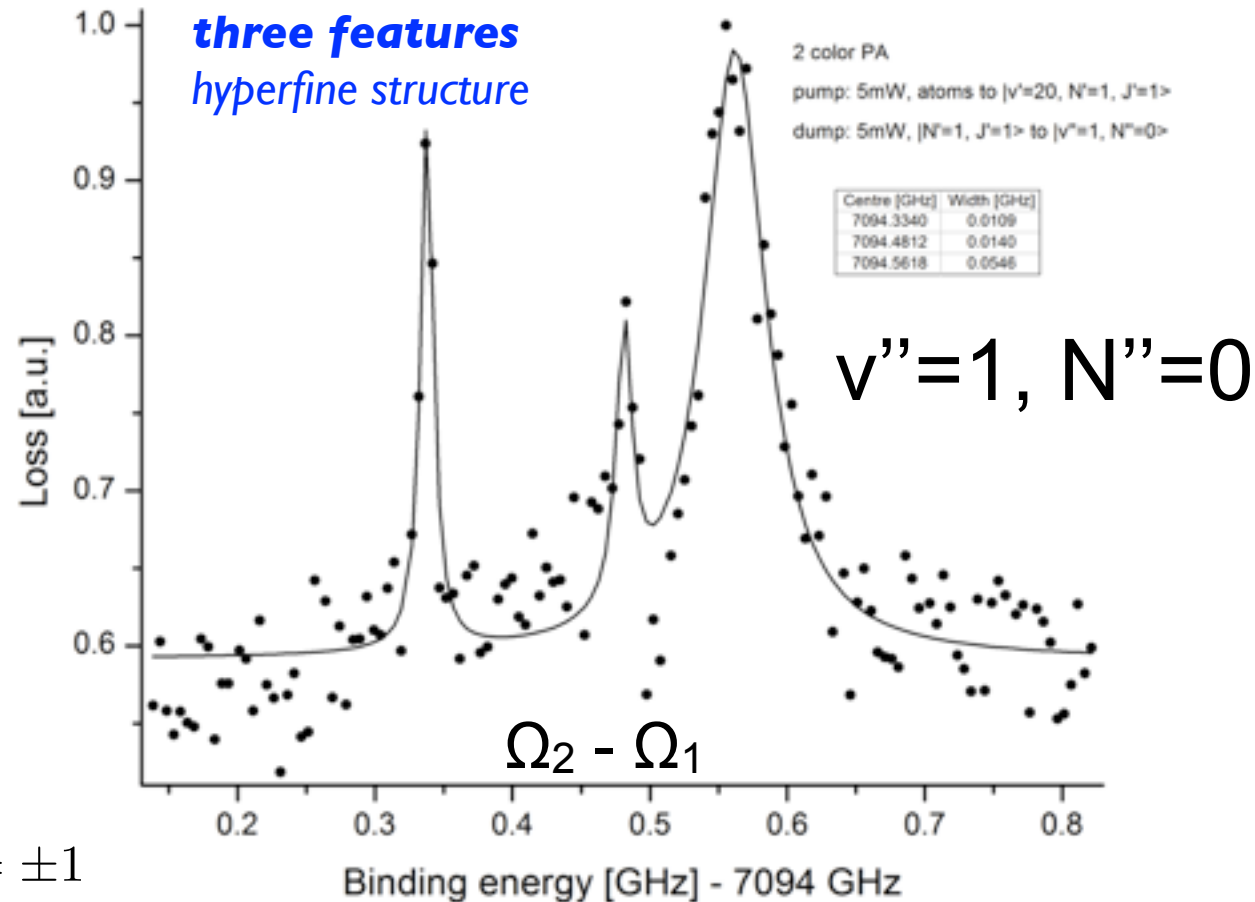


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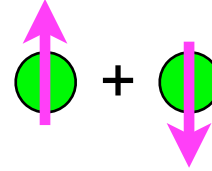
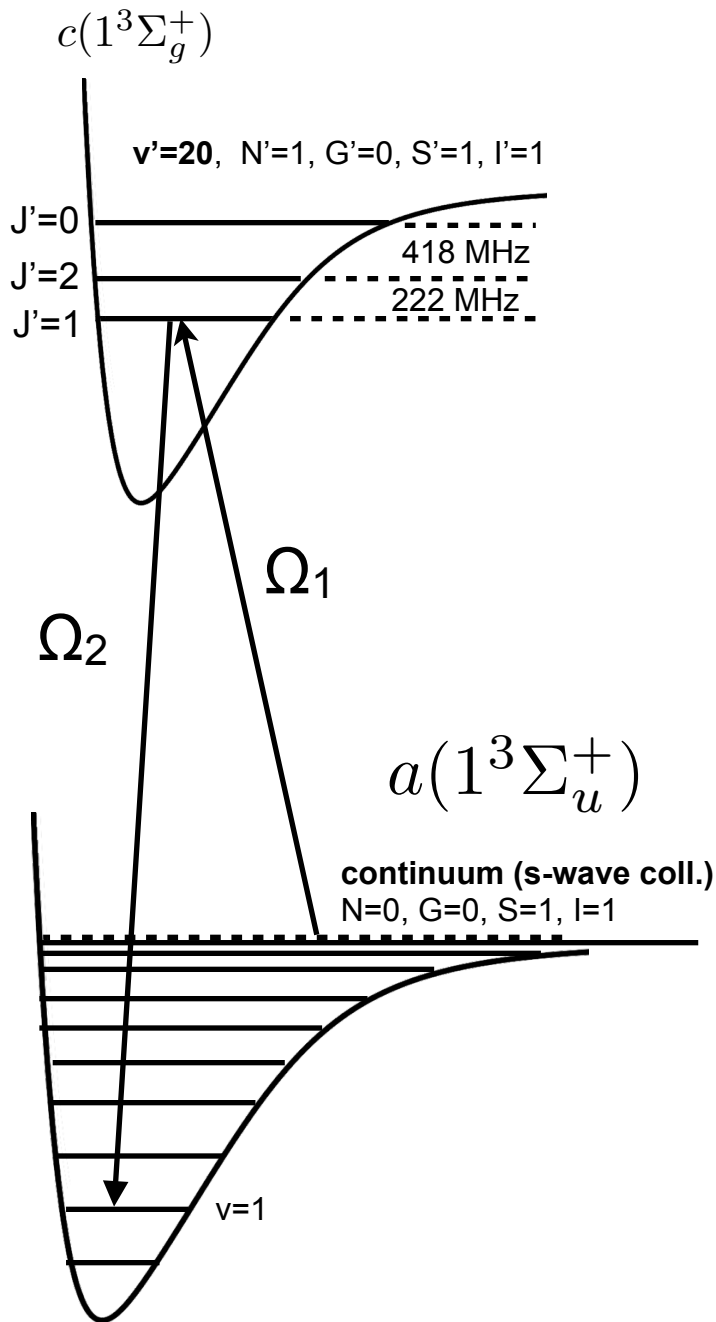
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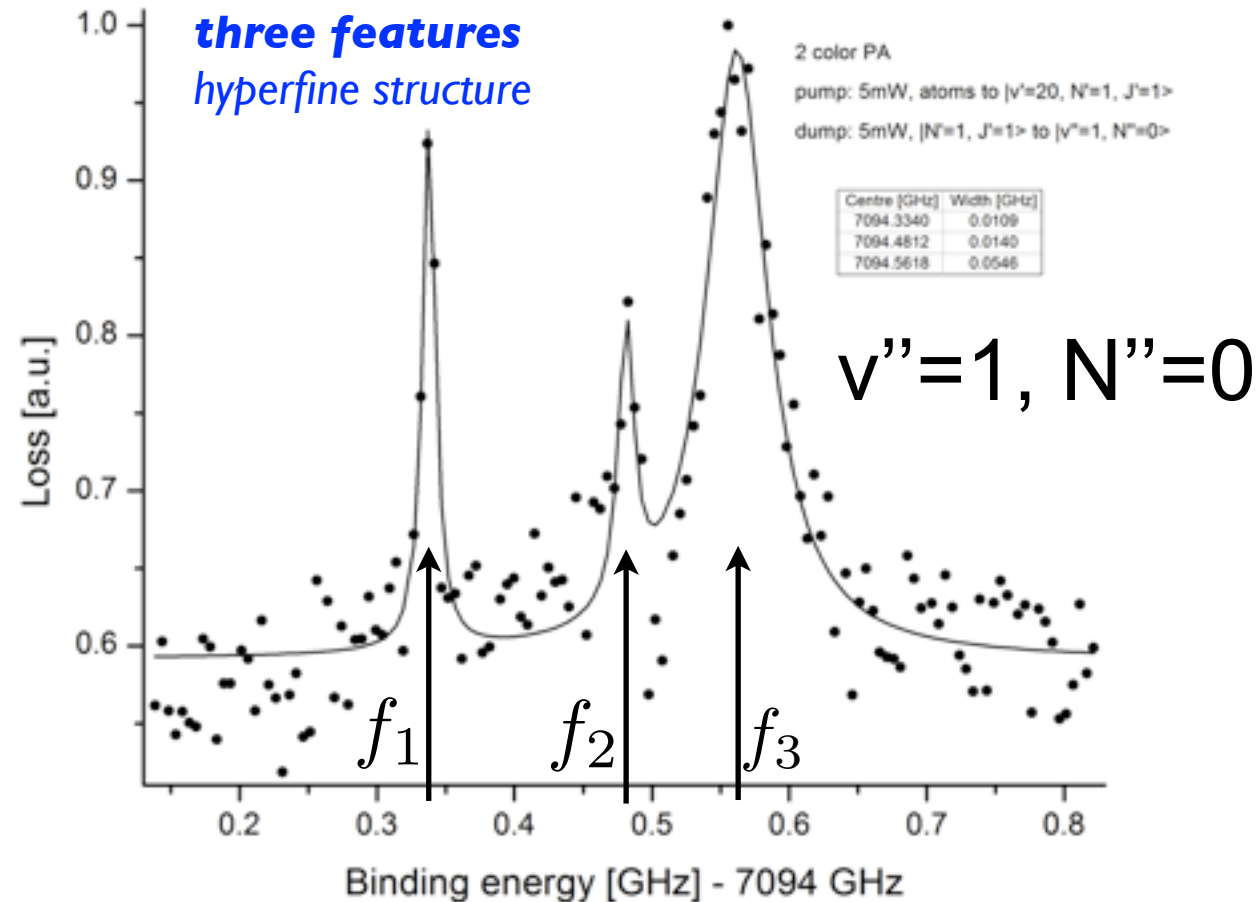
$|S=1, N=1, J=1, l=1\rangle$

the final state is:

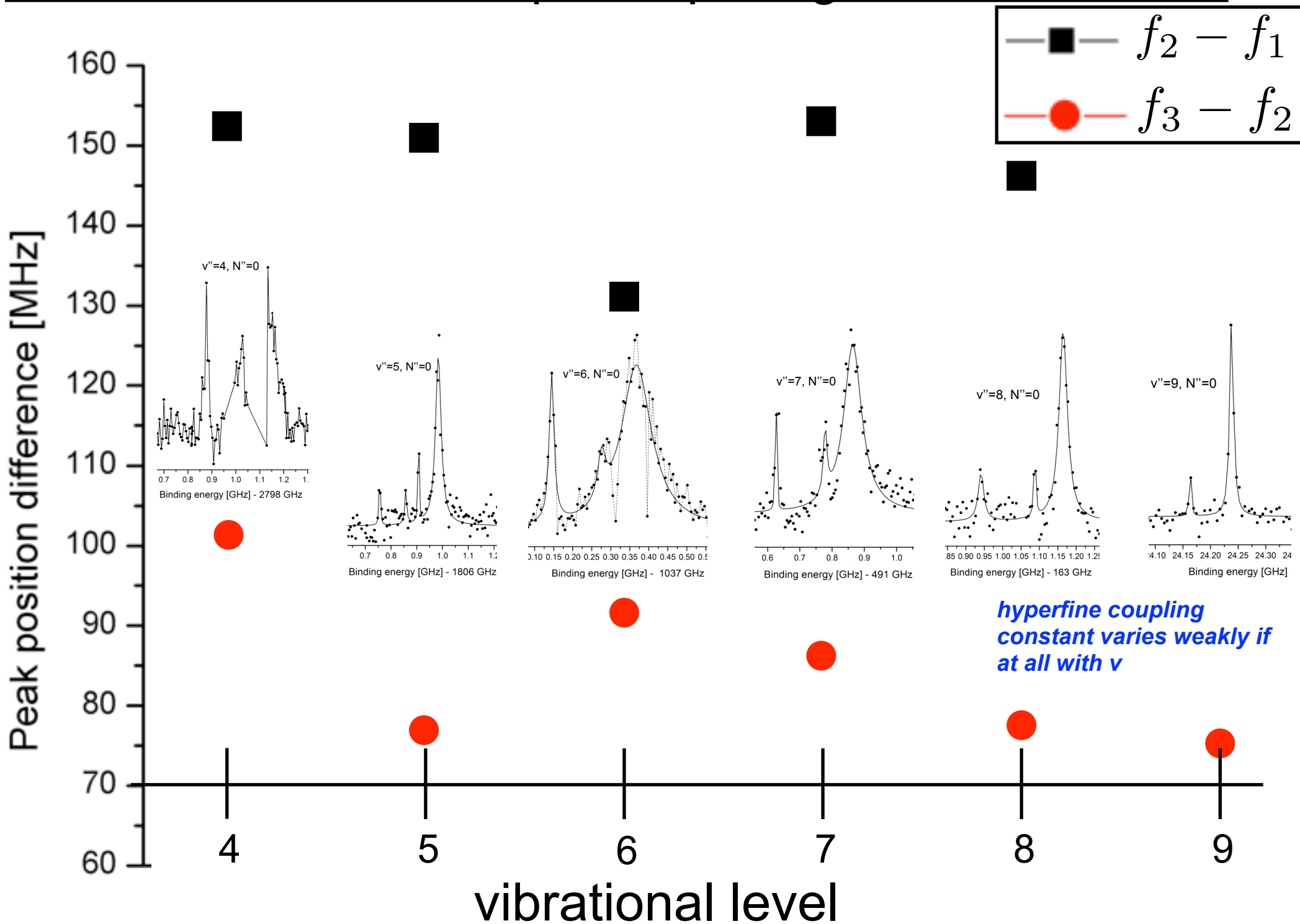
$|S=1, N=0, J=1, l=1\rangle$

only one value of J

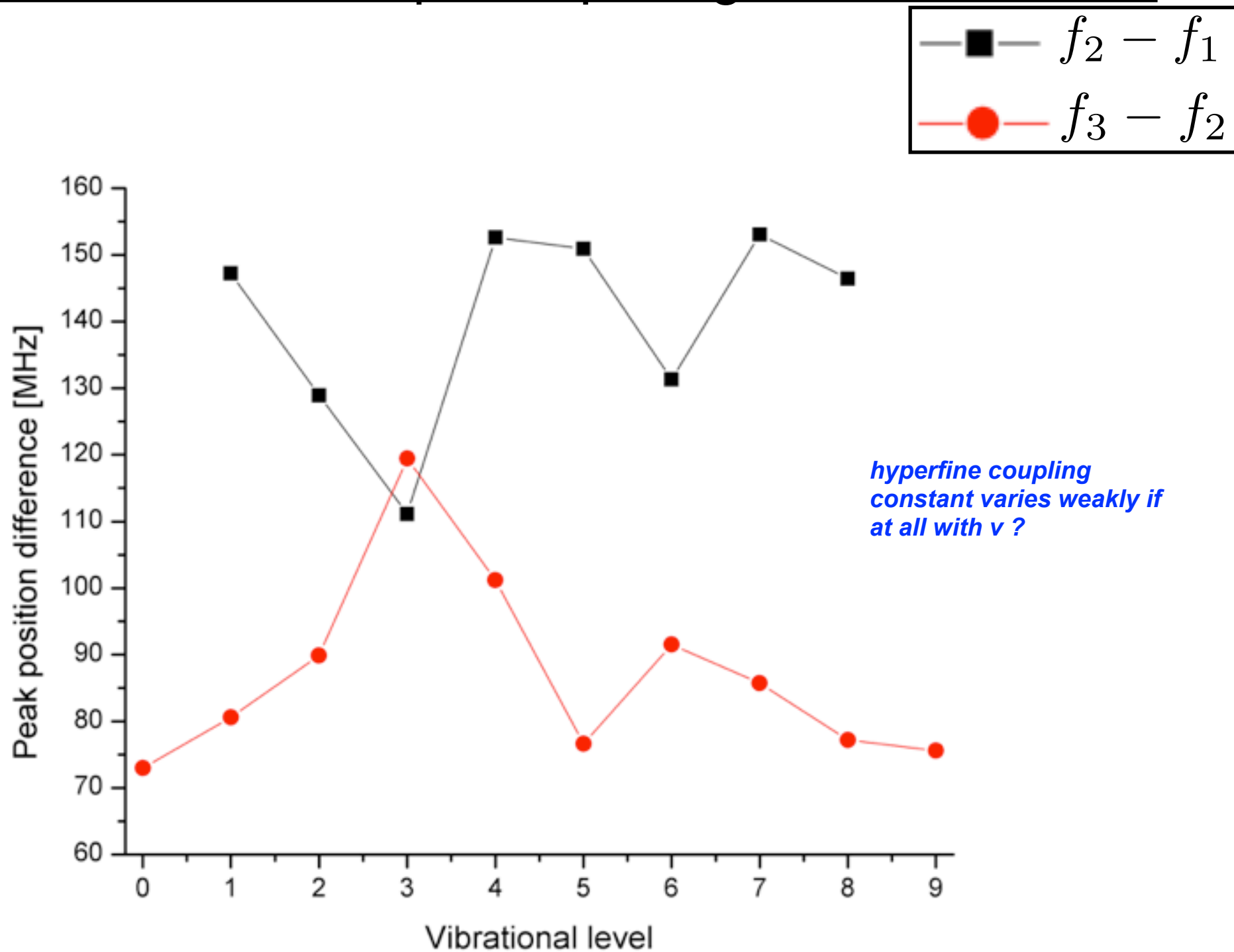
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N=0 peak splittings

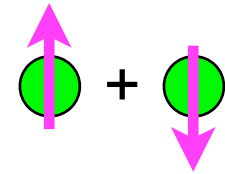
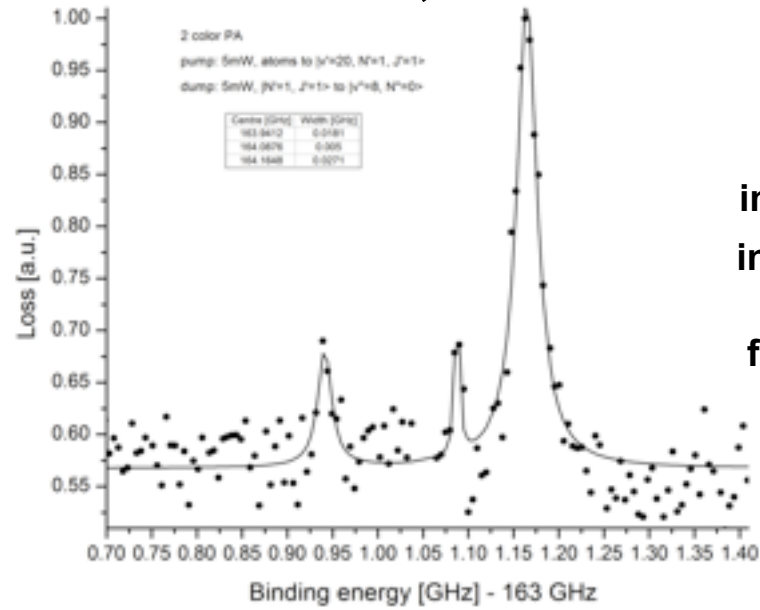
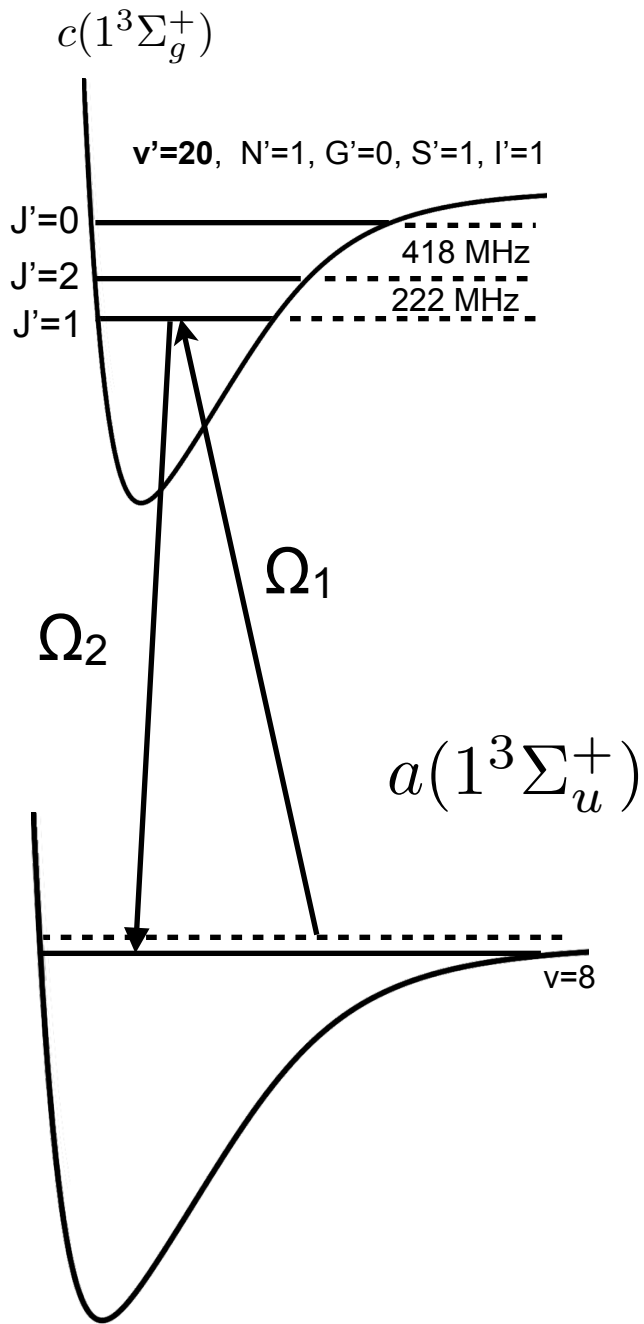


N=0 peak splittings



2-color PA spectroscopy results : here Ω_1 is fixed and Ω_2 is scanned

$v''=8, N''=0$



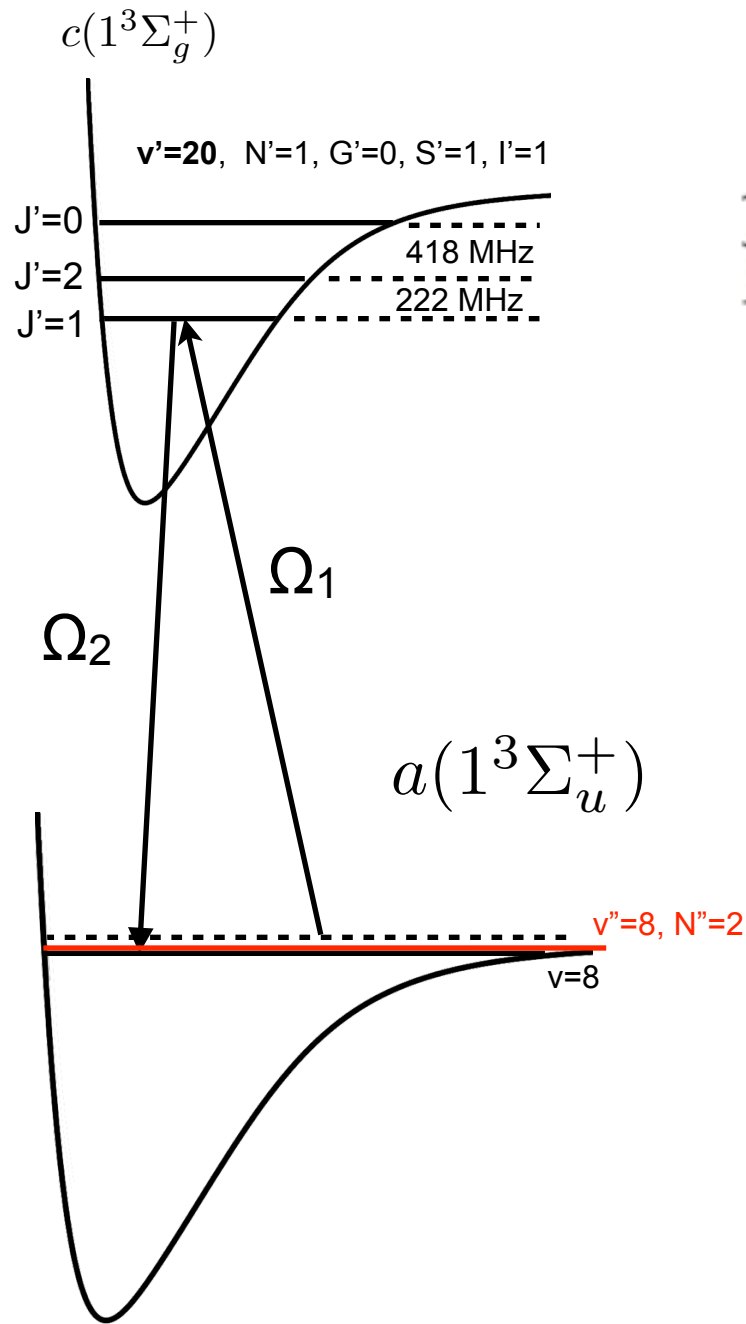
initial state $|S=1, N=0, J=1, l=1\rangle$

intermediate: $|S=1, N=1, J=1, l=1\rangle$

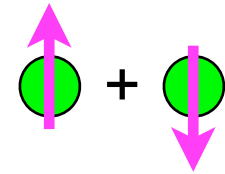
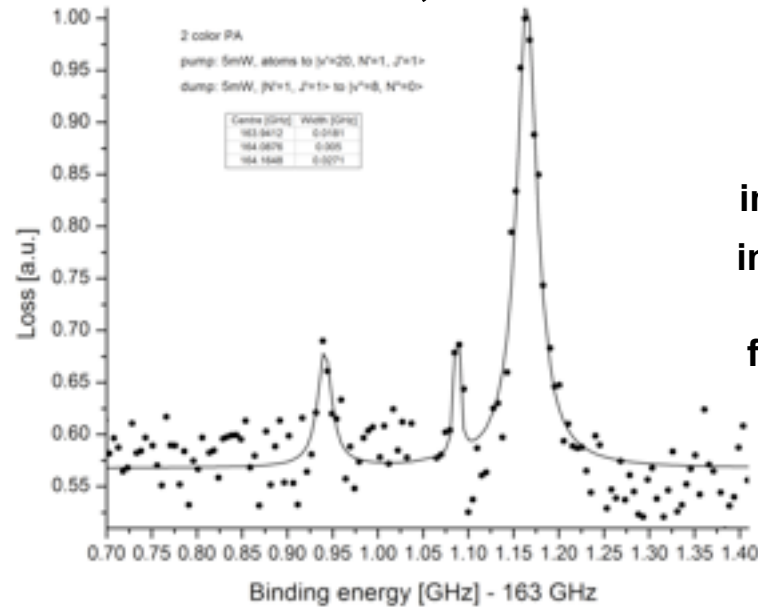
final: $|S=1, N=0, J=1, l=1\rangle$

three features

2-color PA spectroscopy results : here Ω_1 is fixed and Ω_2 is scanned



$v''=8, N''=0$



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 intermediate: $|S=1, N=1, J=1, l=1\rangle$

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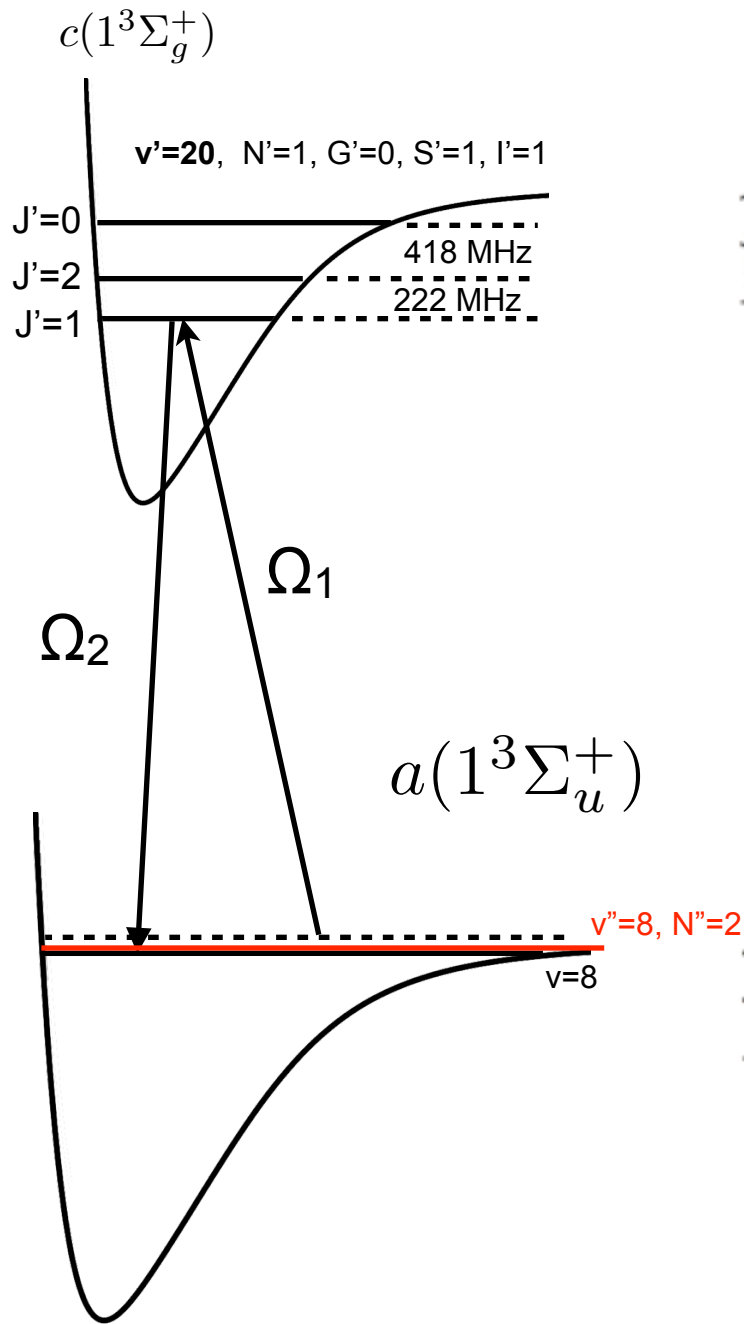
three features

$v''=8, N''=2$

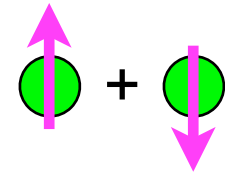
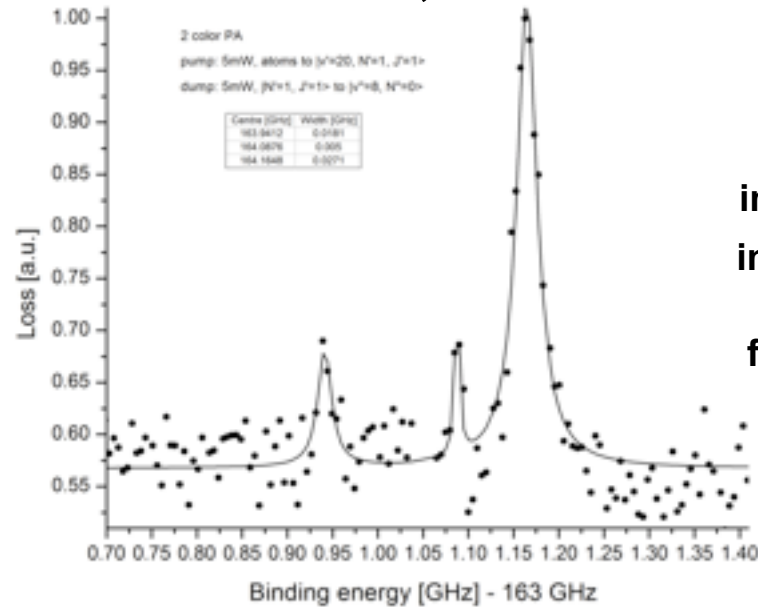
final state:
 $|S=1, N=2, J=1, l=1\rangle$
 OR
 $|S=1, N=2, J=2, l=1\rangle$

six features ?

2-color PA spectroscopy results : here Ω_1 is fixed and Ω_2 is scanned



$v''=8, N''=0$



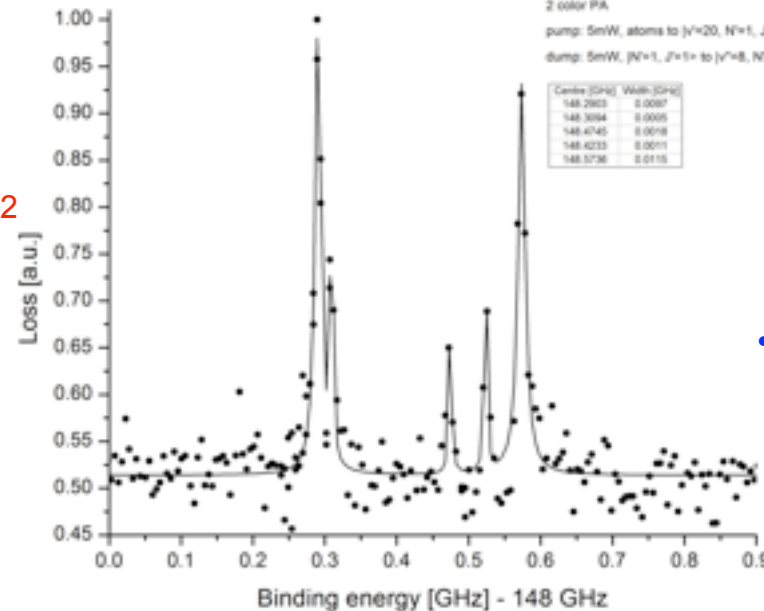
initial state $|S=1, N=0, J=1, l=1\rangle$

intermediate: $|S=1, N=1, J=1, l=1\rangle$

final: $|S=1, N=0, J=1, l=1\rangle$

three features

$v''=8, N''=2$



final state:
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OR

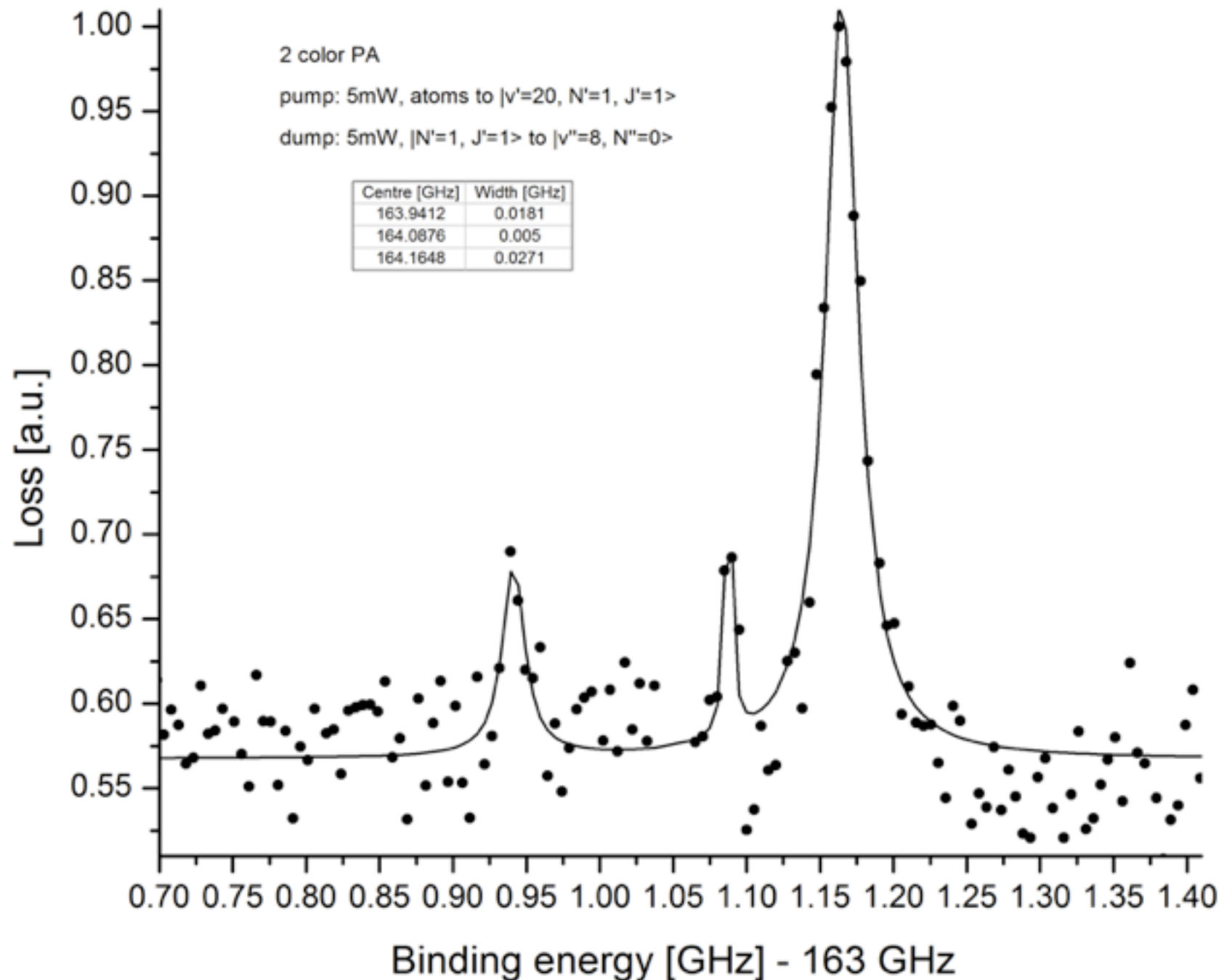
$|S=1, N=2, J=2, l=1\rangle$

six features ?

... we see five features

2-color PA spectroscopy results : here Ω_1 is fixed and Ω_2 is scanned

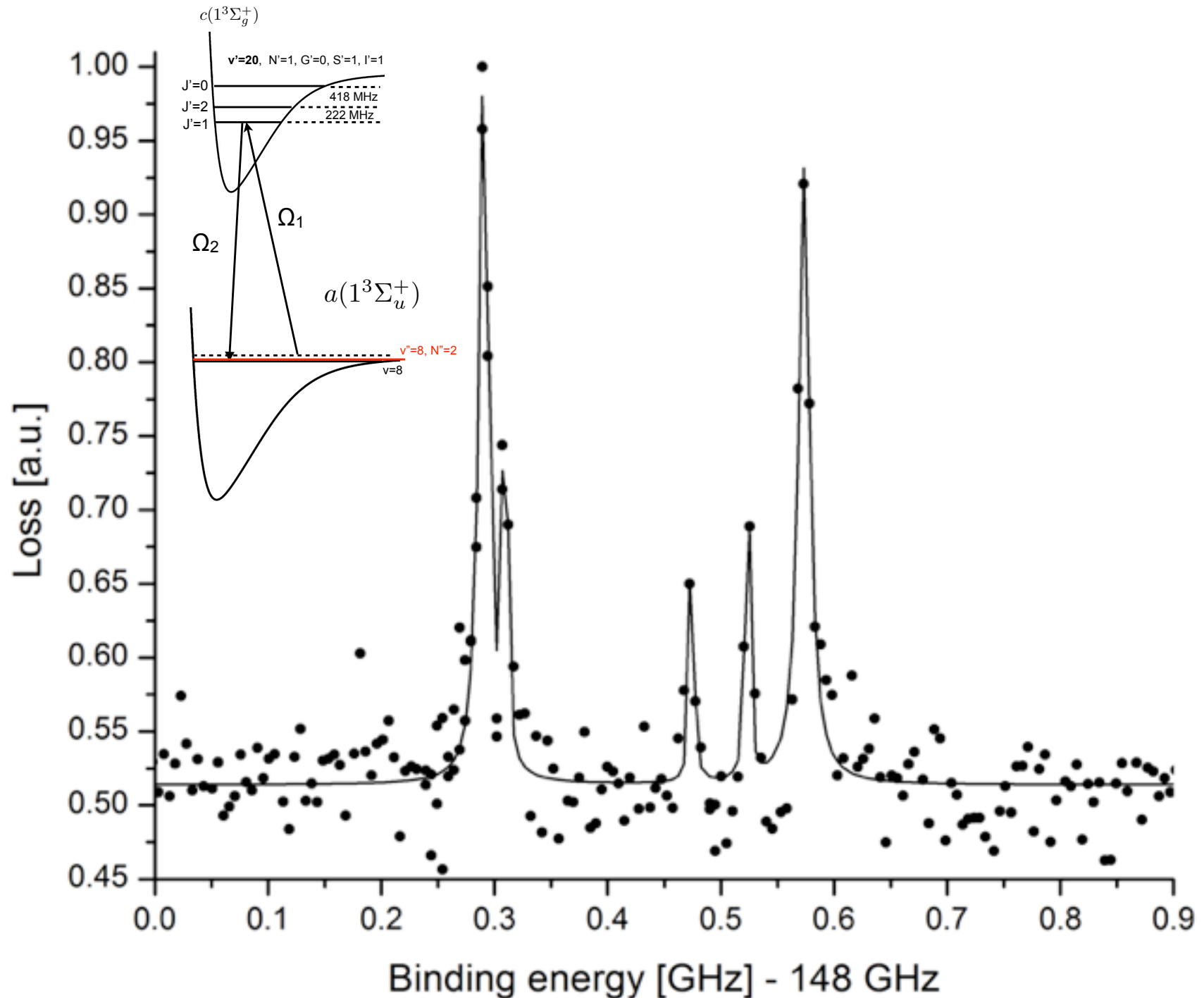
$$v''=8, N''=0$$



what follows is unchanged and what was presented, but after the talk Mariusz and Will sent me a correction : the data plotted in red were not obtained with Ω_1 tuned to the $J=0$ state, rather it was obtained with Ω_1 tuned to the $J=2$ state! So the conclusions are, as stated, incorrect.

2-color PA spectroscopy results : here Ω_1 is fixed and Ω_2 is scanned

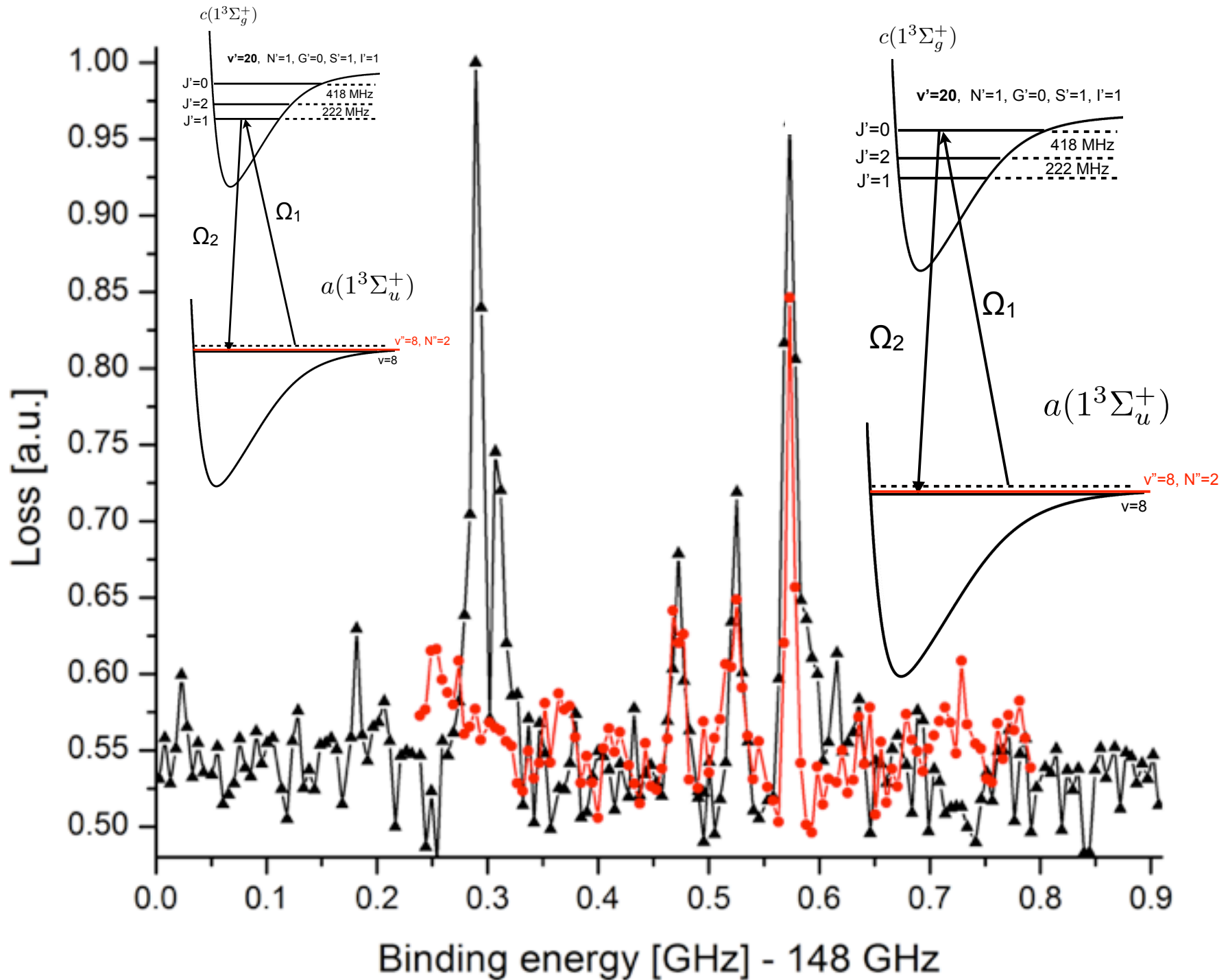
final: $|S=1, N=2, J=1 \text{ or } 2, l=1\rangle$



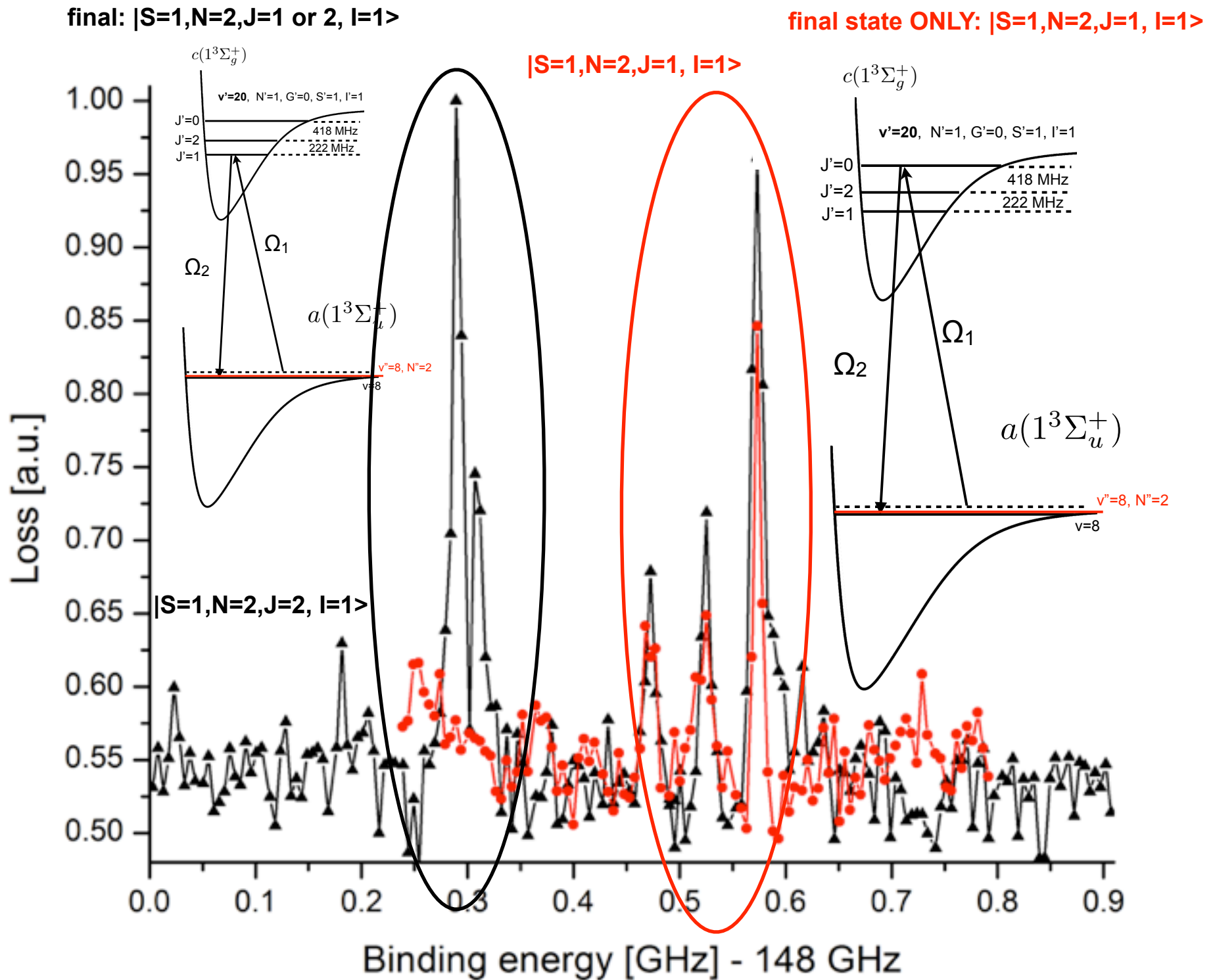
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final state ONLY: $|S=1, N=2, J=1, l=1\rangle$

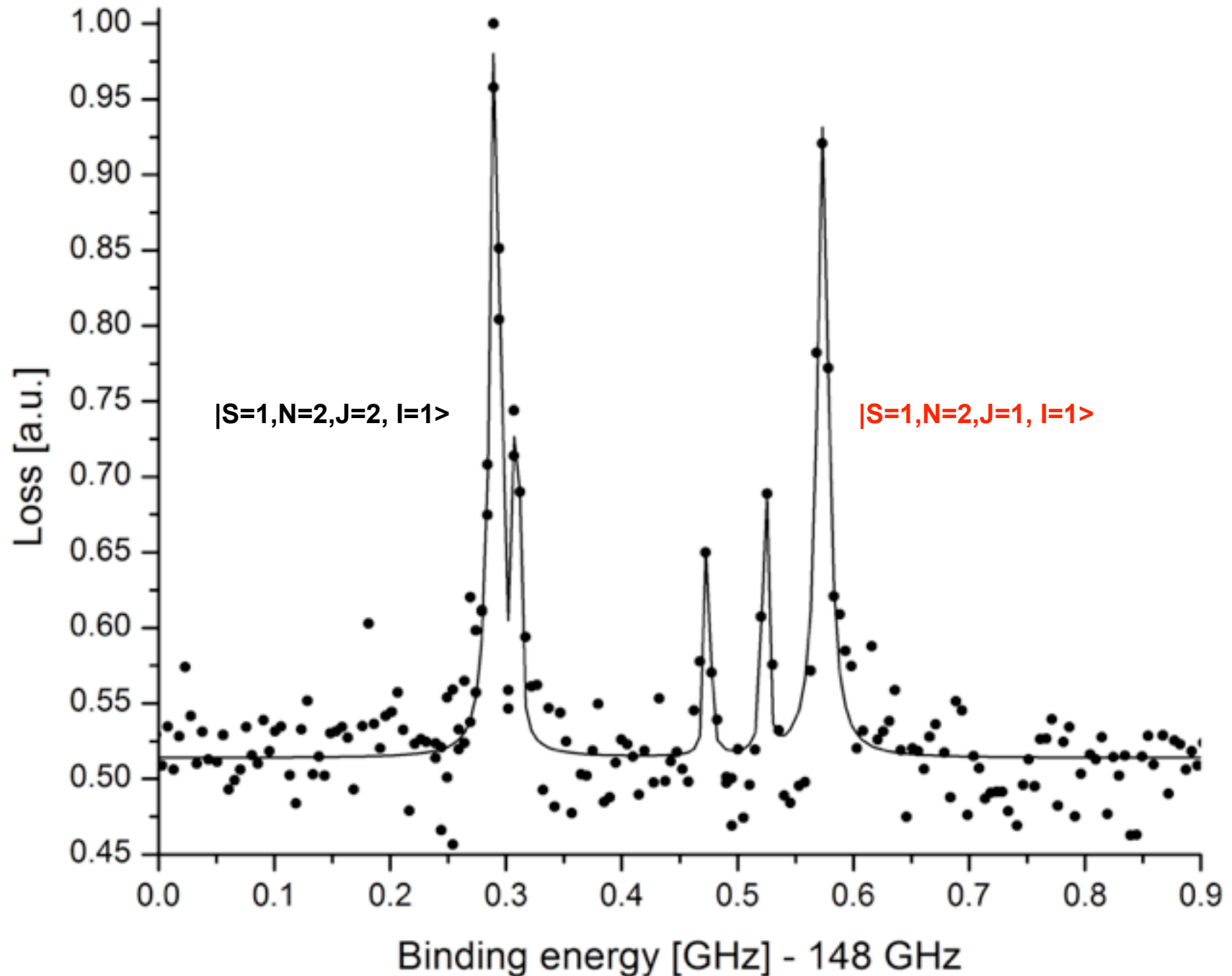


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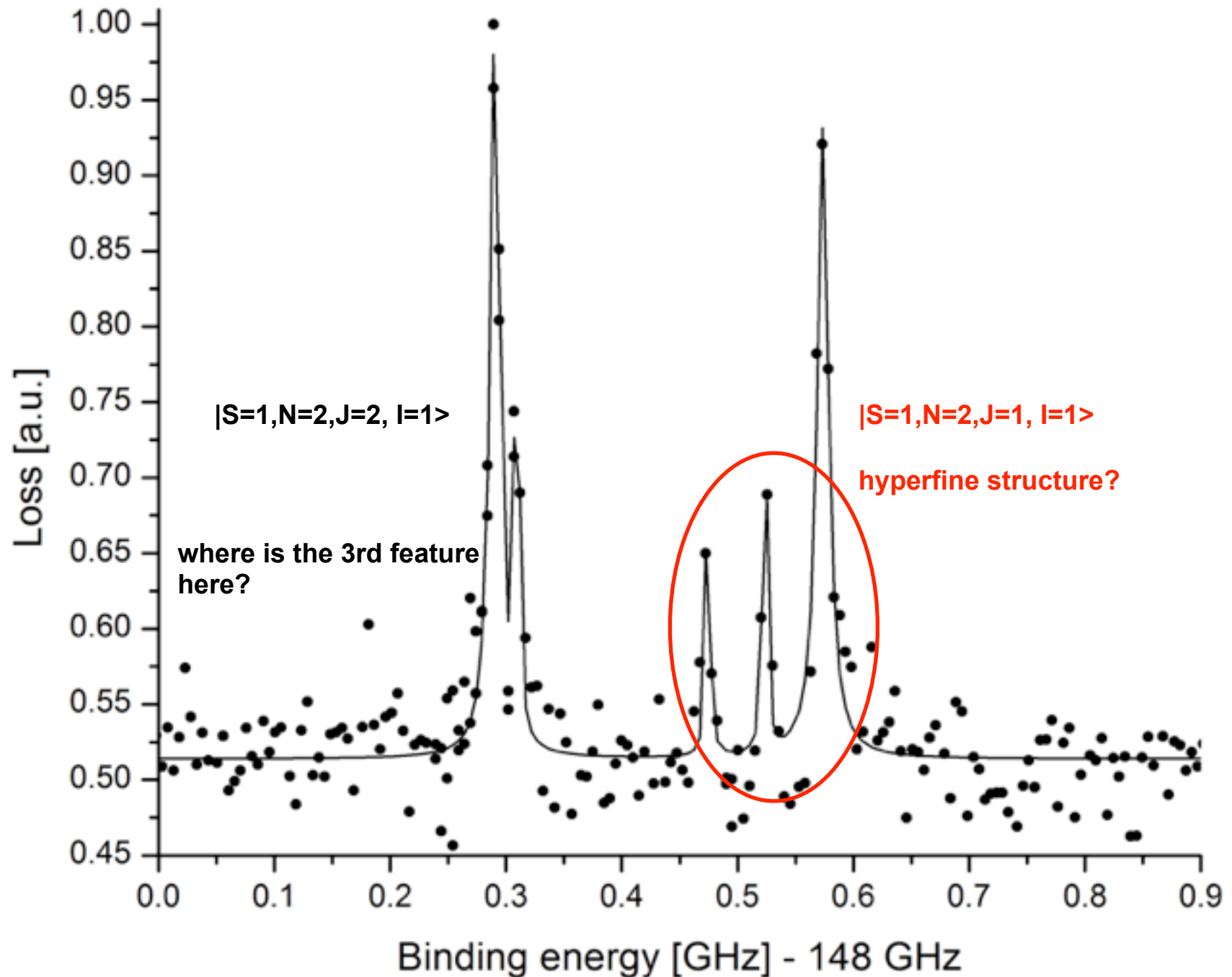
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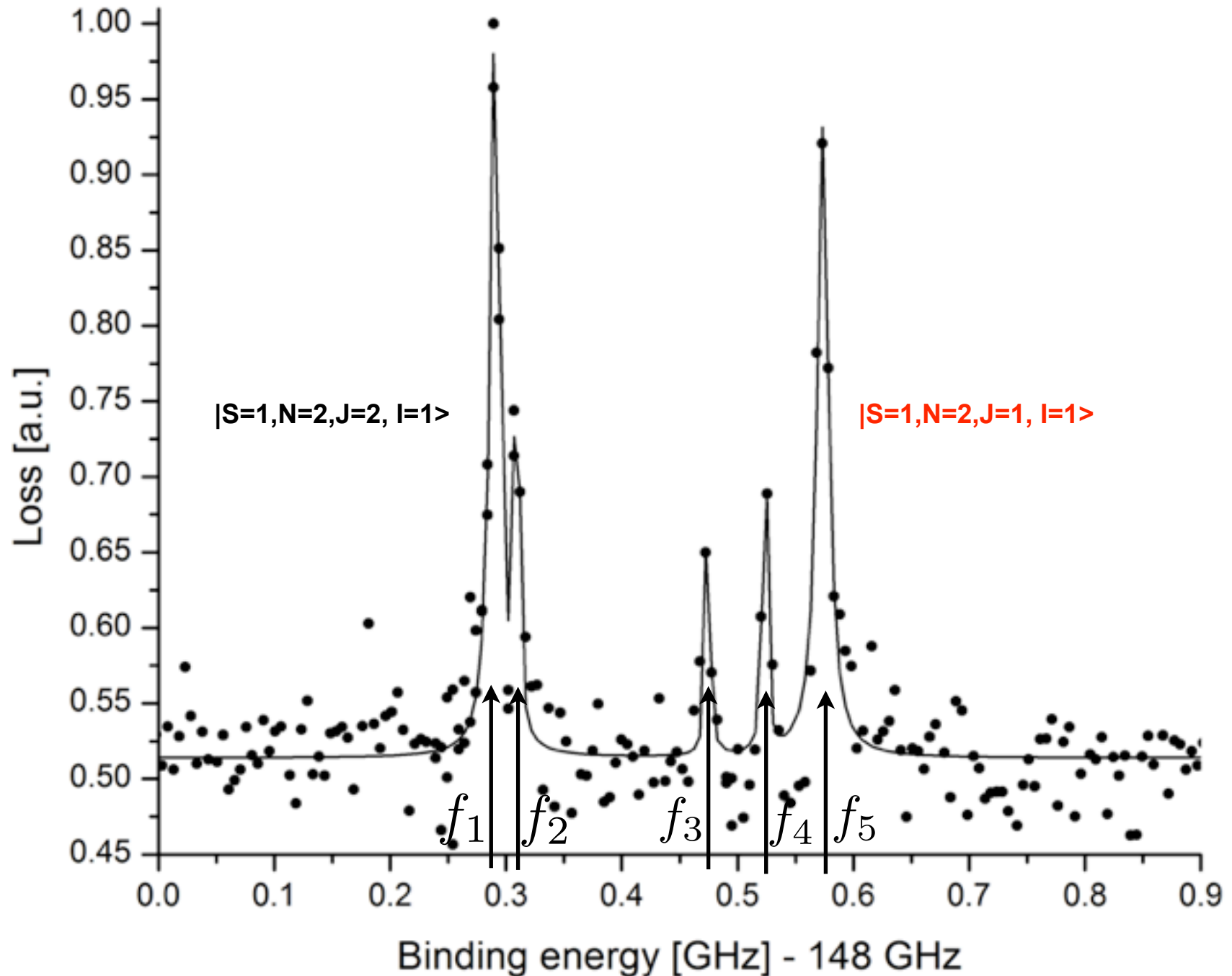
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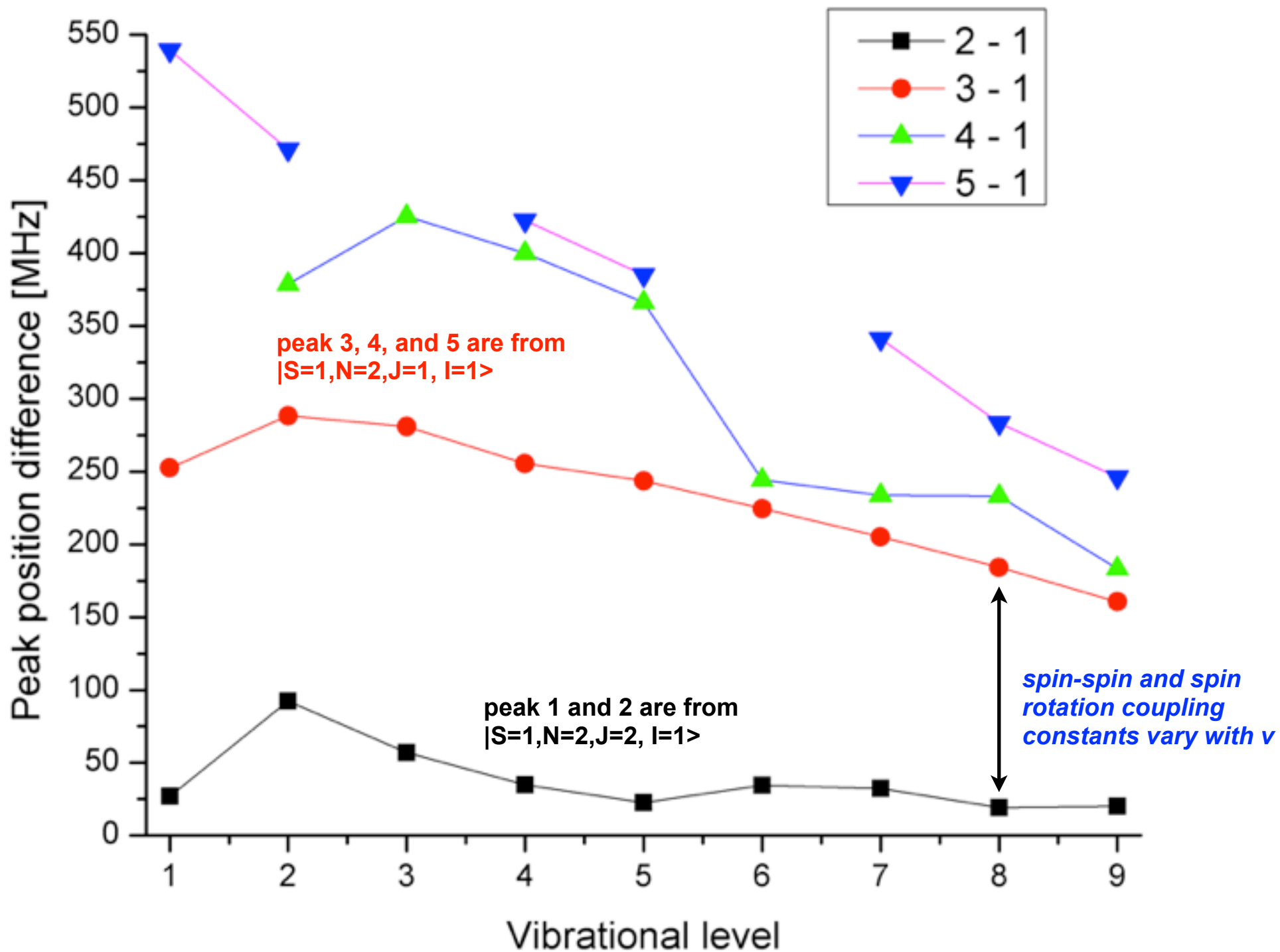
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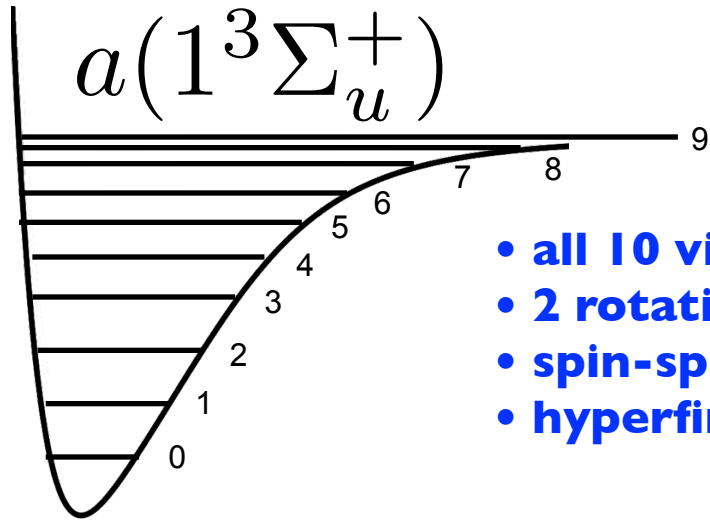
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Triplet ground state Li_2

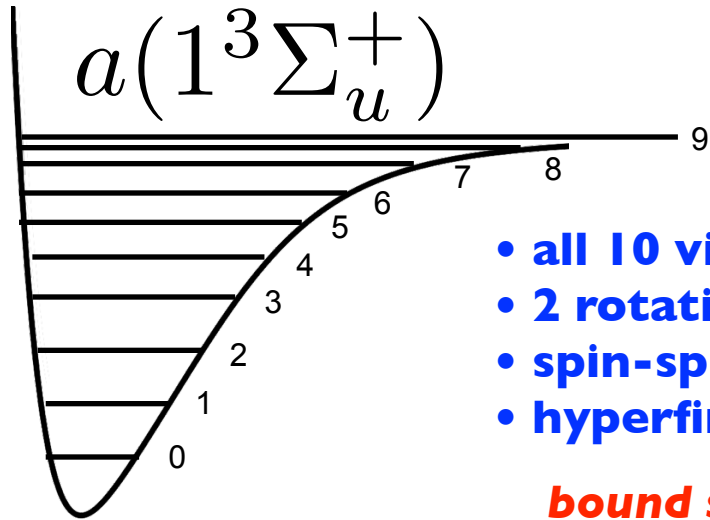
Full characterization of the bound levels of the ground triplet state of ${}^6\text{Li}_2$



- all 10 vibrational levels found
- 2 rotational levels for each vibrational level ($N''=0, N''=2$)
- spin-spin and spin-rotation splitting of J levels observed
- hyperfine structure resolved

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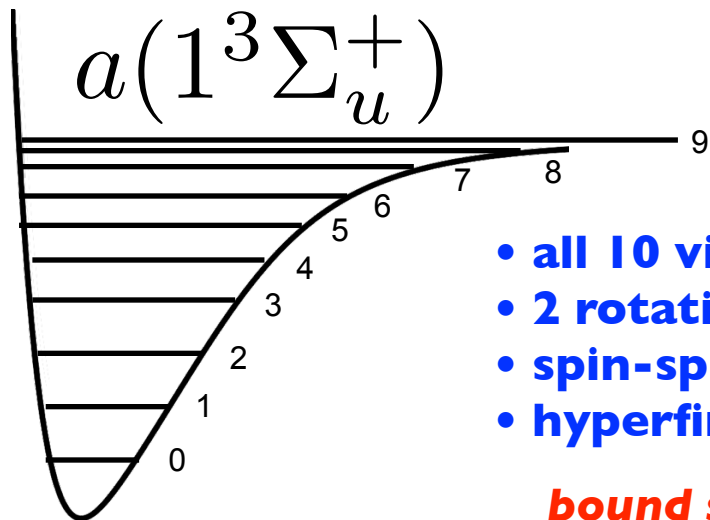


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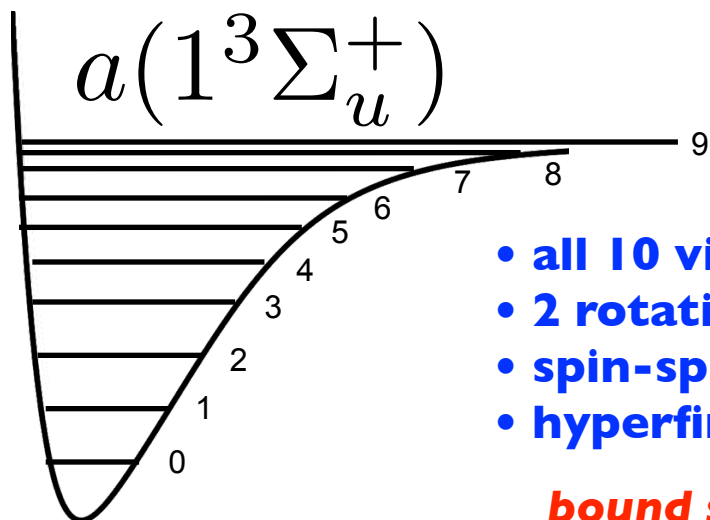
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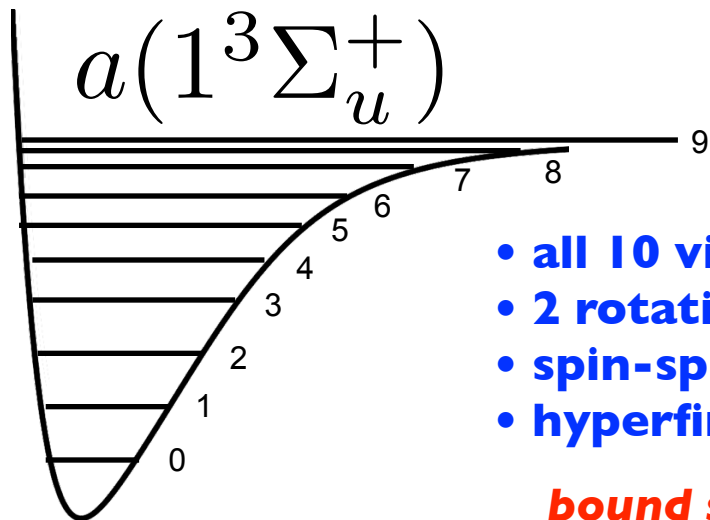
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It's the simplest molecule after to H_2

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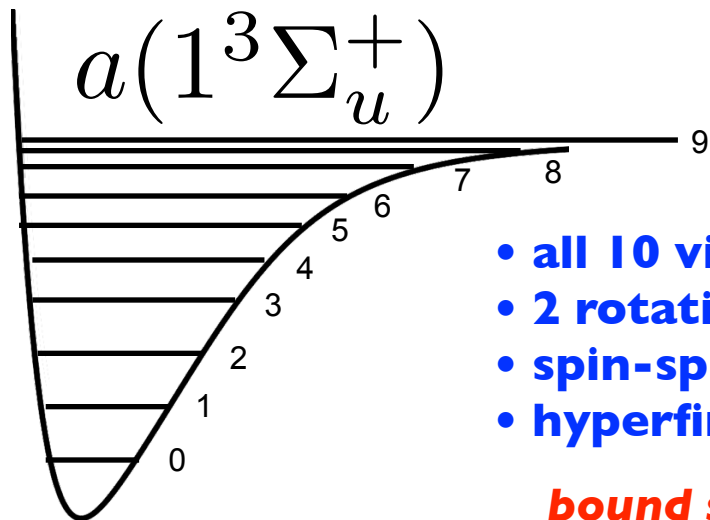
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Next step: STIRAP transfer FR molecules into these states and measure their lifetime

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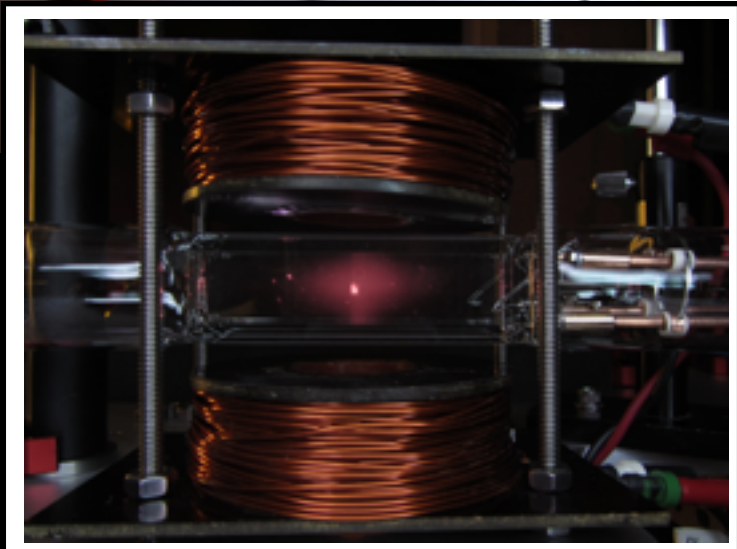
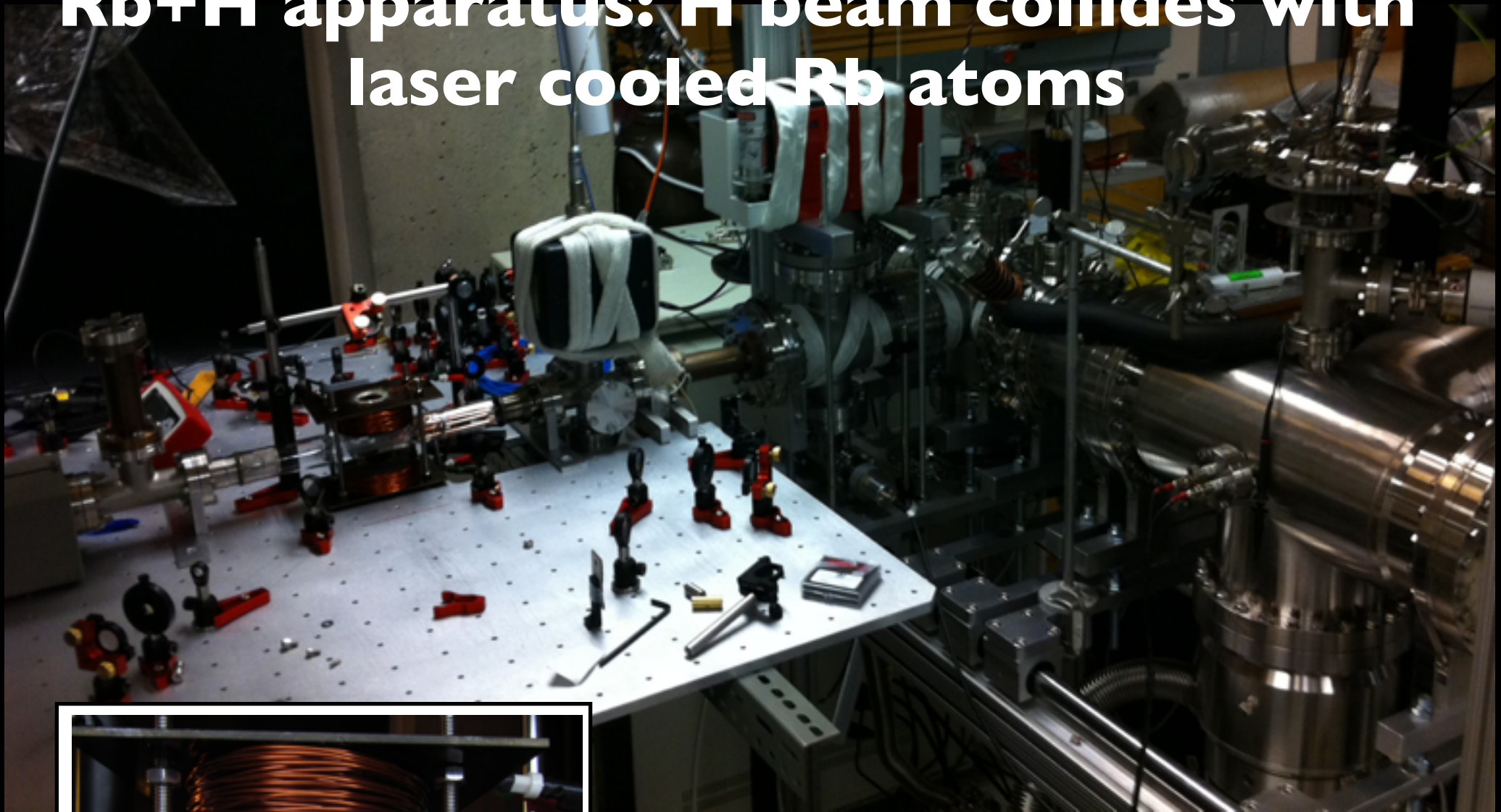
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fin

Rb+H apparatus: H beam collides with laser cooled Rb atoms



Collaboration with
Takamasa Momose (UBC CHEM)

A new twist on Feshbach resonances: *polar* collision + *electric* field

PRL **96**, 123202 (2006)

PHYSICAL REVIEW LETTERS

week ending
31 MARCH 2006

Controlling Collisions of Ultracold Atoms with dc Electric Fields

R. V. Krems*

Department of Chemistry, University of British Columbia, Vancouver, British Columbia V6T 1Z1, Canada

PHYSICAL REVIEW A **75**, 032709 (2007)

Electric-field-induced Feshbach resonances in ultracold alkali-metal mixtures

Z. Li and R. V. Krems*

Department of Chemistry, University of British Columbia, Vancouver, British Columbia, Canada V6T 1Z1

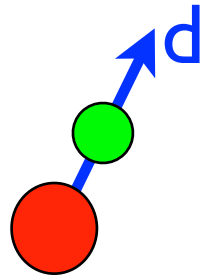
PHYSICAL REVIEW A **79**, 042711 (2009)

Effects of electric fields on heteronuclear Feshbach resonances in ultracold ${}^6\text{Li}$ - ${}^{87}\text{Rb}$ mixtures

Z. Li¹ and K. W. Madison²

... a new twist on Feshbach resonances

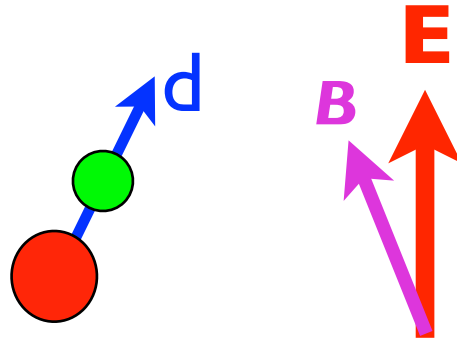
polar collision



A **heteronuclear** collision complex has an instantaneous **electric dipole moment**

... a new twist on Feshbach resonances

polar collision + electric field

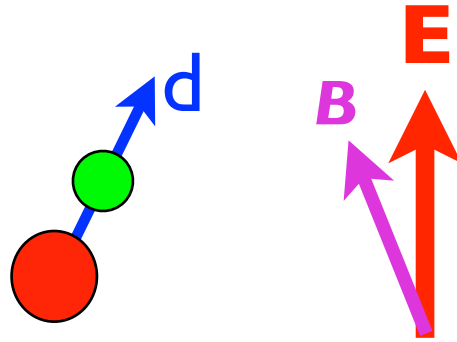


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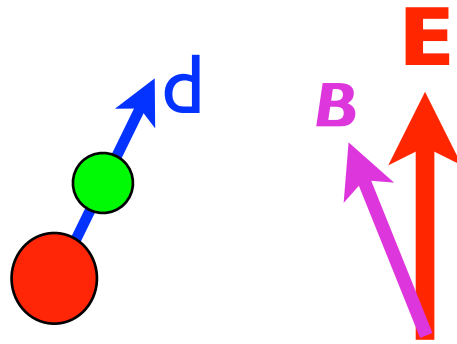
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$$\hat{V}_E(R) = -\vec{E} \cdot \vec{d} = -E(\hat{e}_E \cdot \hat{e}_d) \sum_S \sum_{M_S} |SM_S\rangle d_S(R) \langle SM_S|$$

this operator couples different partial wave states: $l = l' \pm 1$

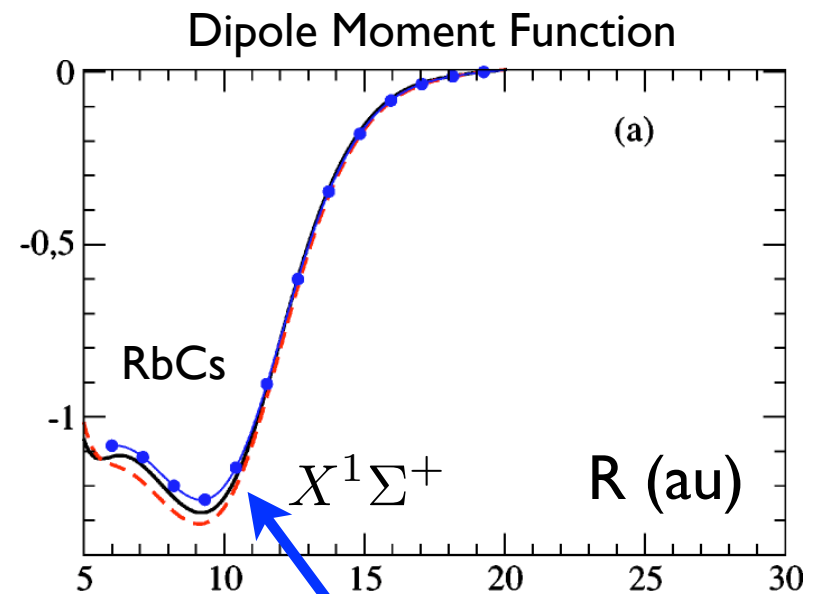
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polar collision + electric field



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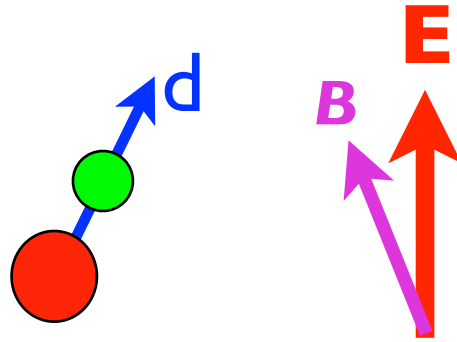


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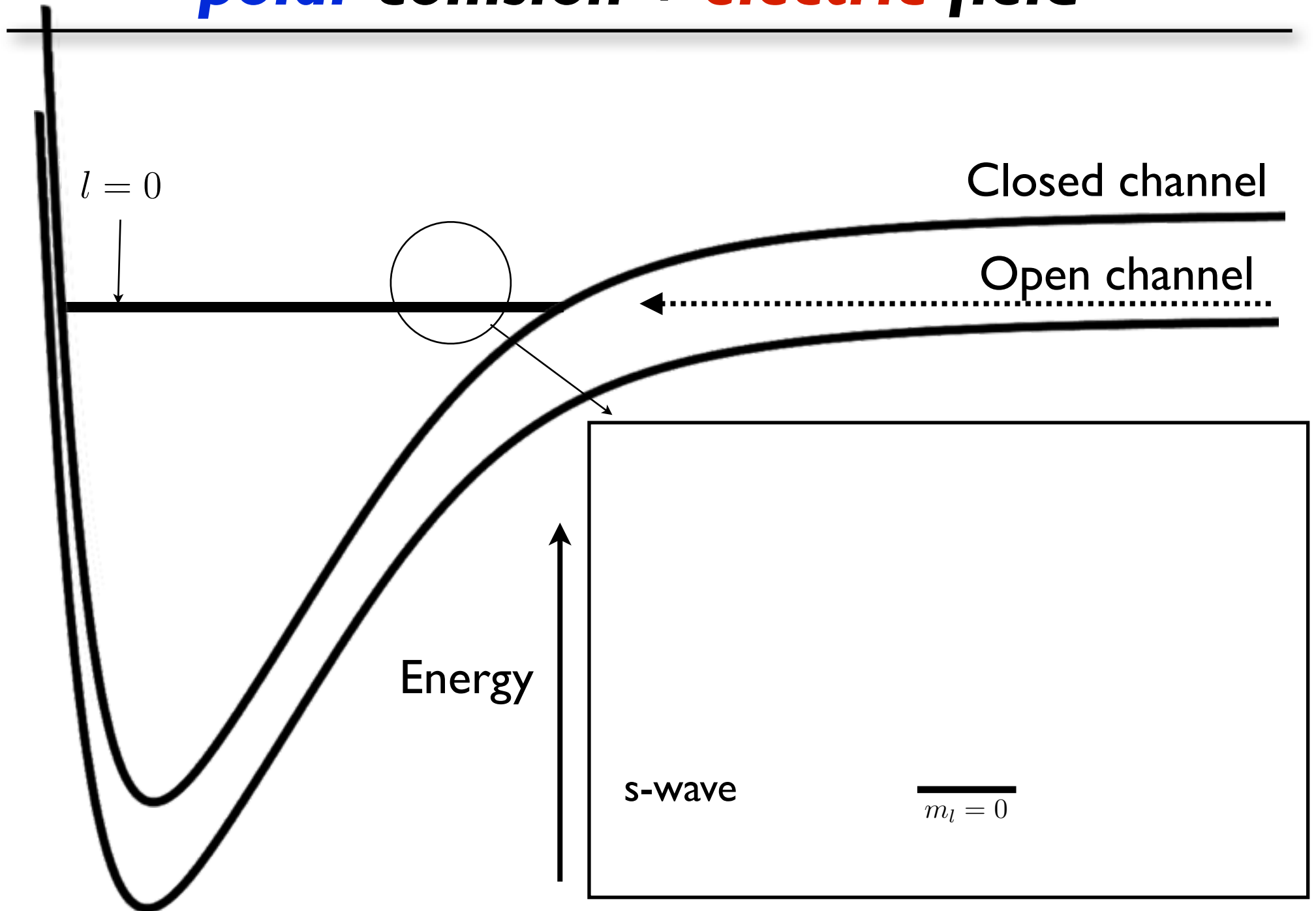
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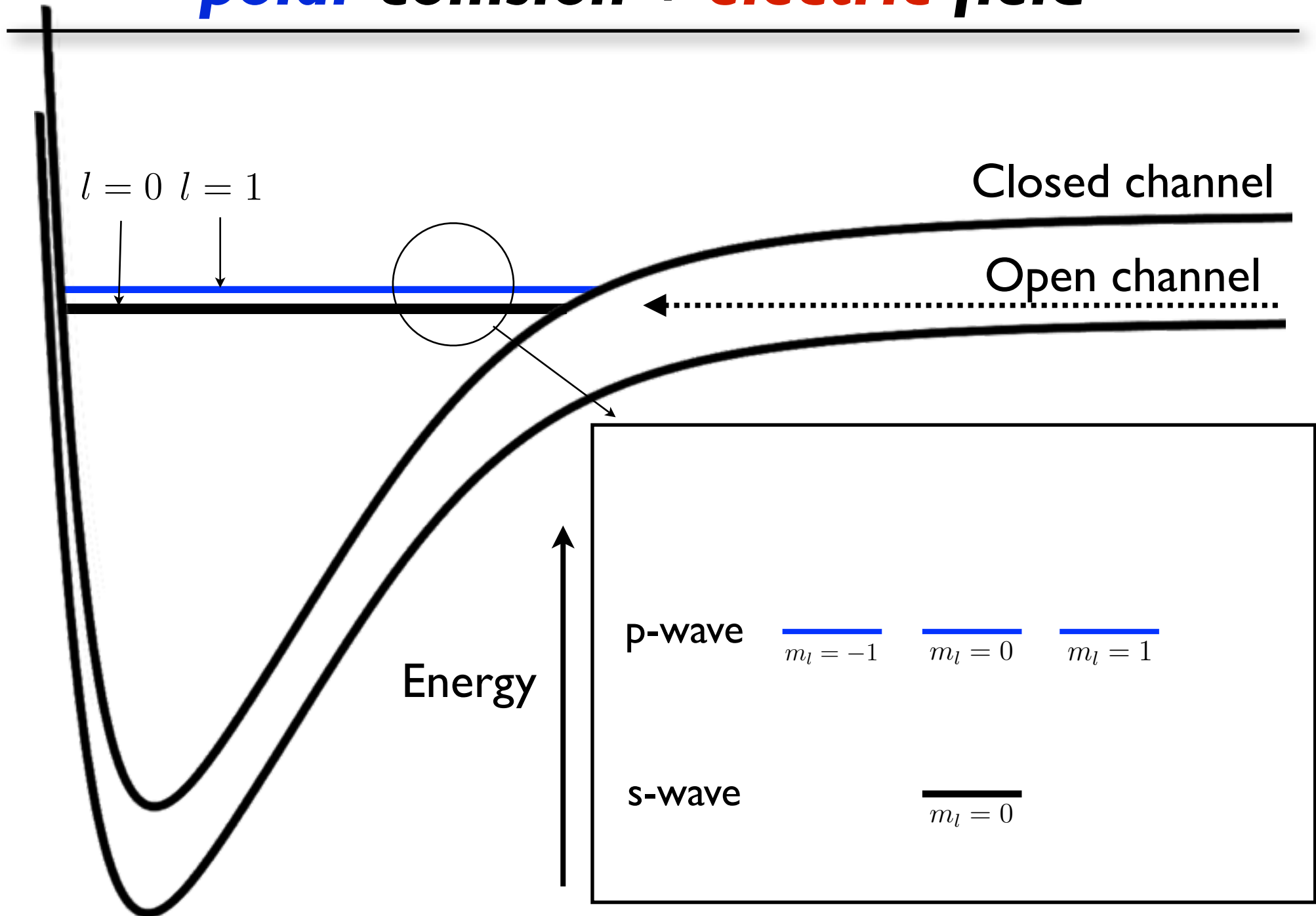


- induce new Feshbach resonances
mixing of partial waves
- shift position of intrinsic resonances
bound state level repulsion
- split resonances into $l+1$ multiplets
tunable anisotropic interactions

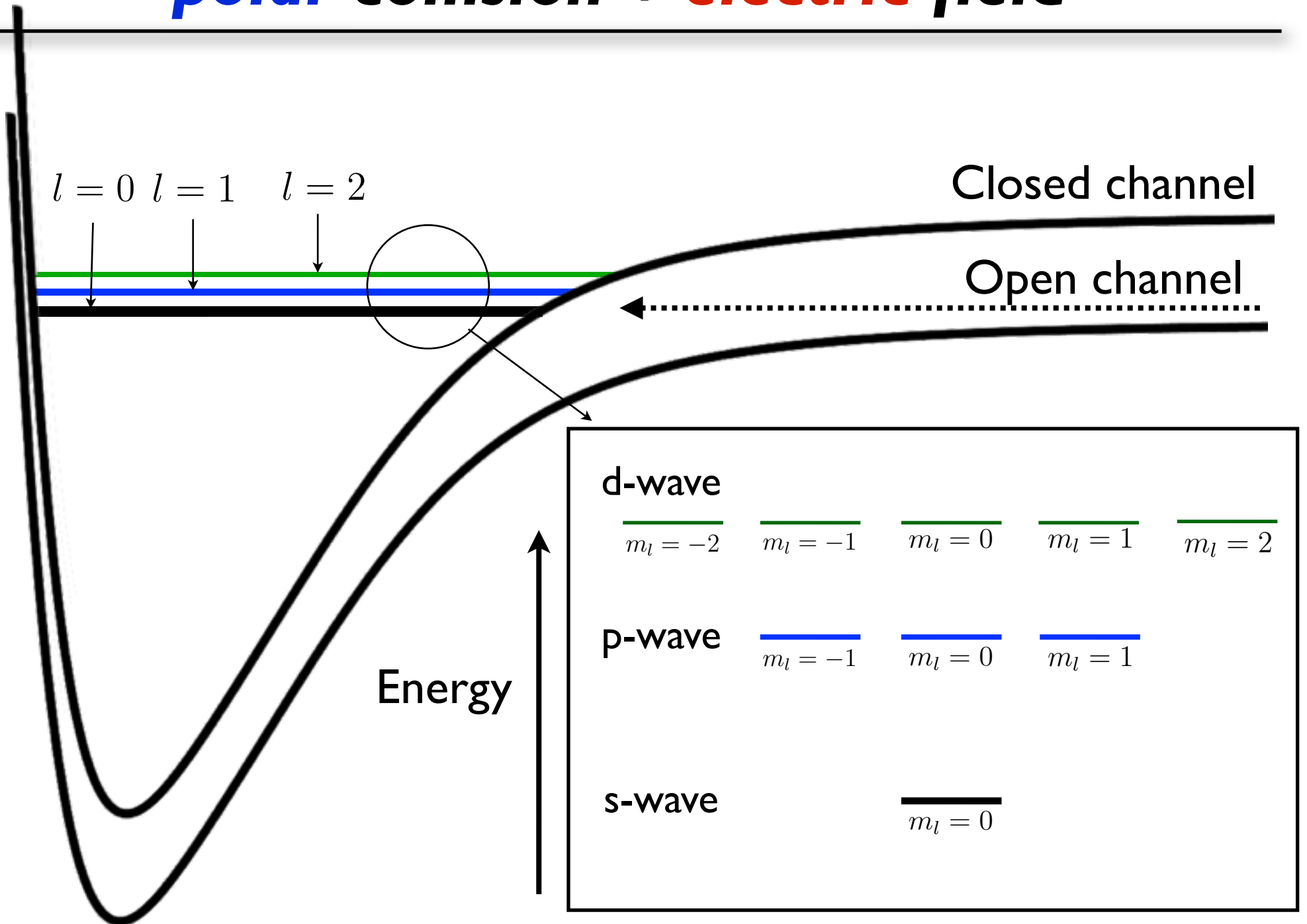
A new twist on Feshbach resonances: *polar* collision + *electric* field



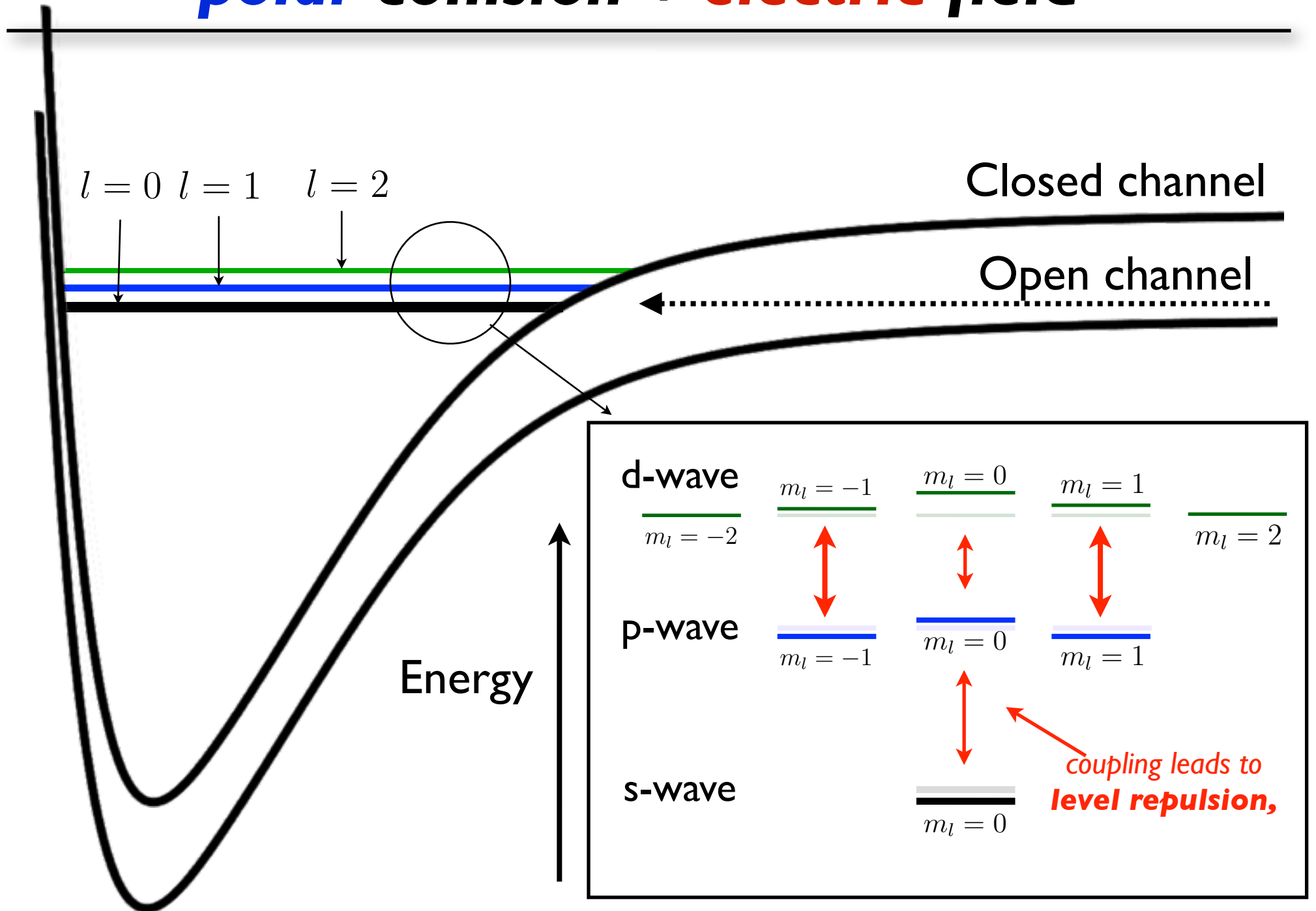
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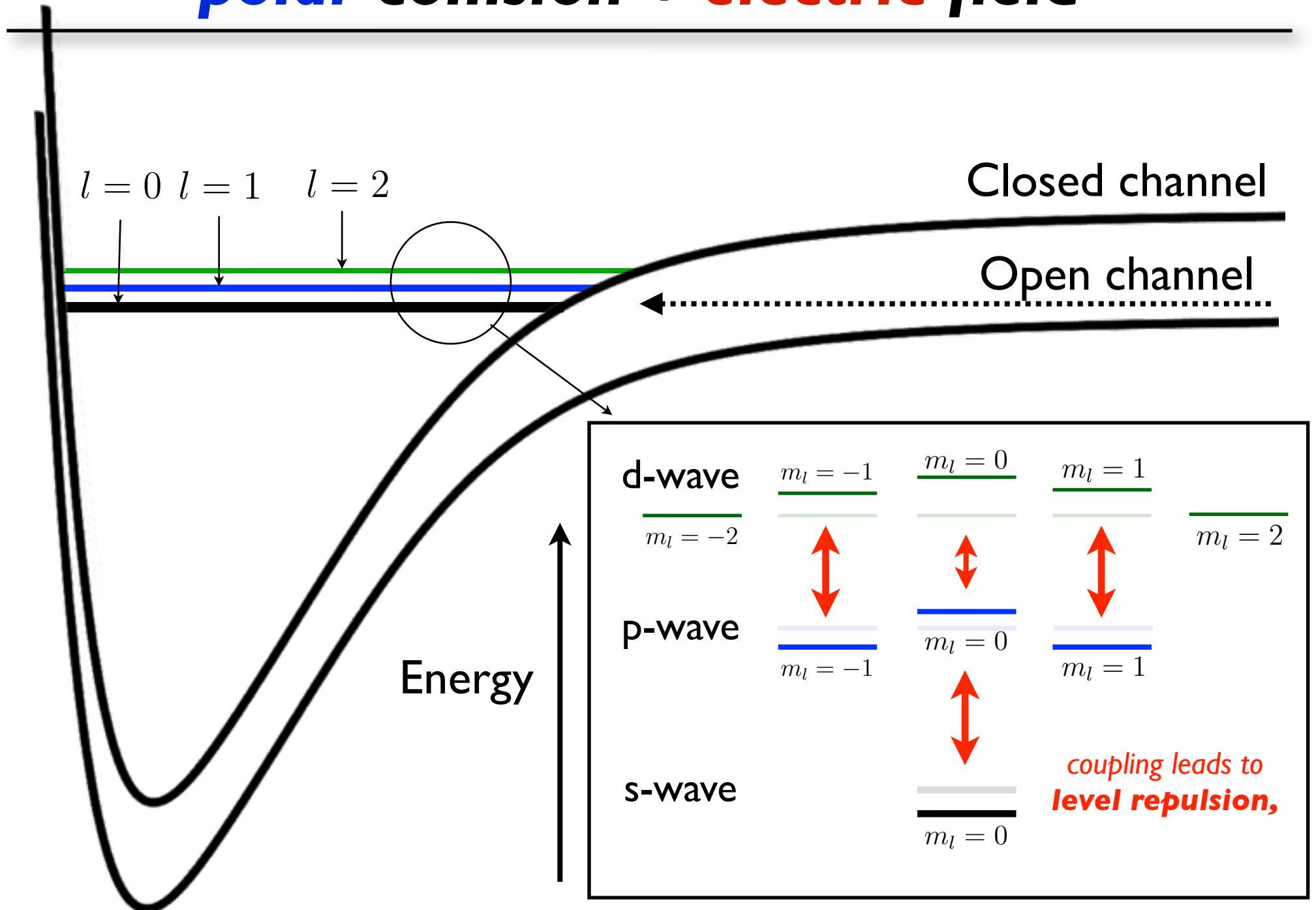
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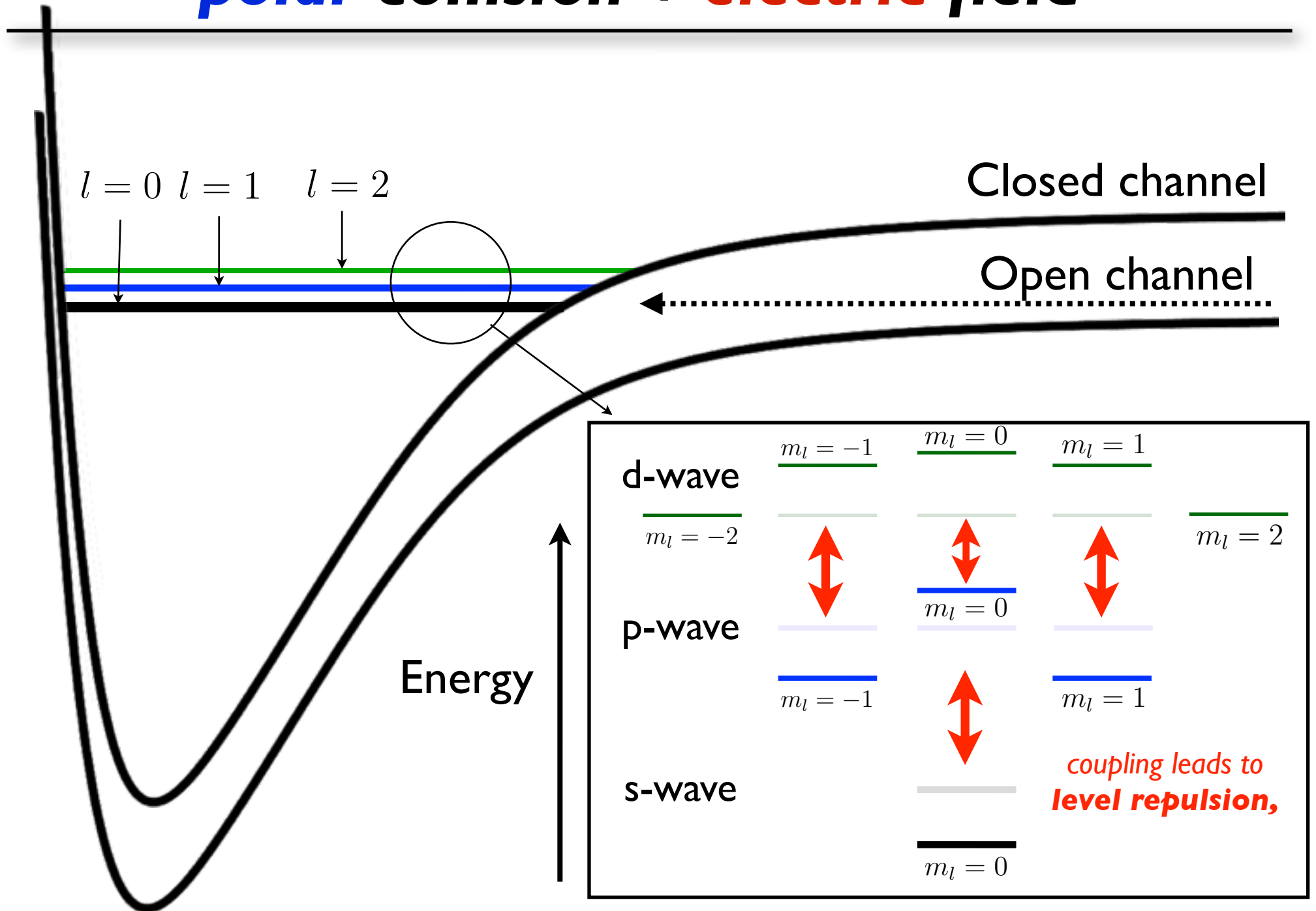
A new twist on Feshbach resonances: *polar* collision + *electric* field



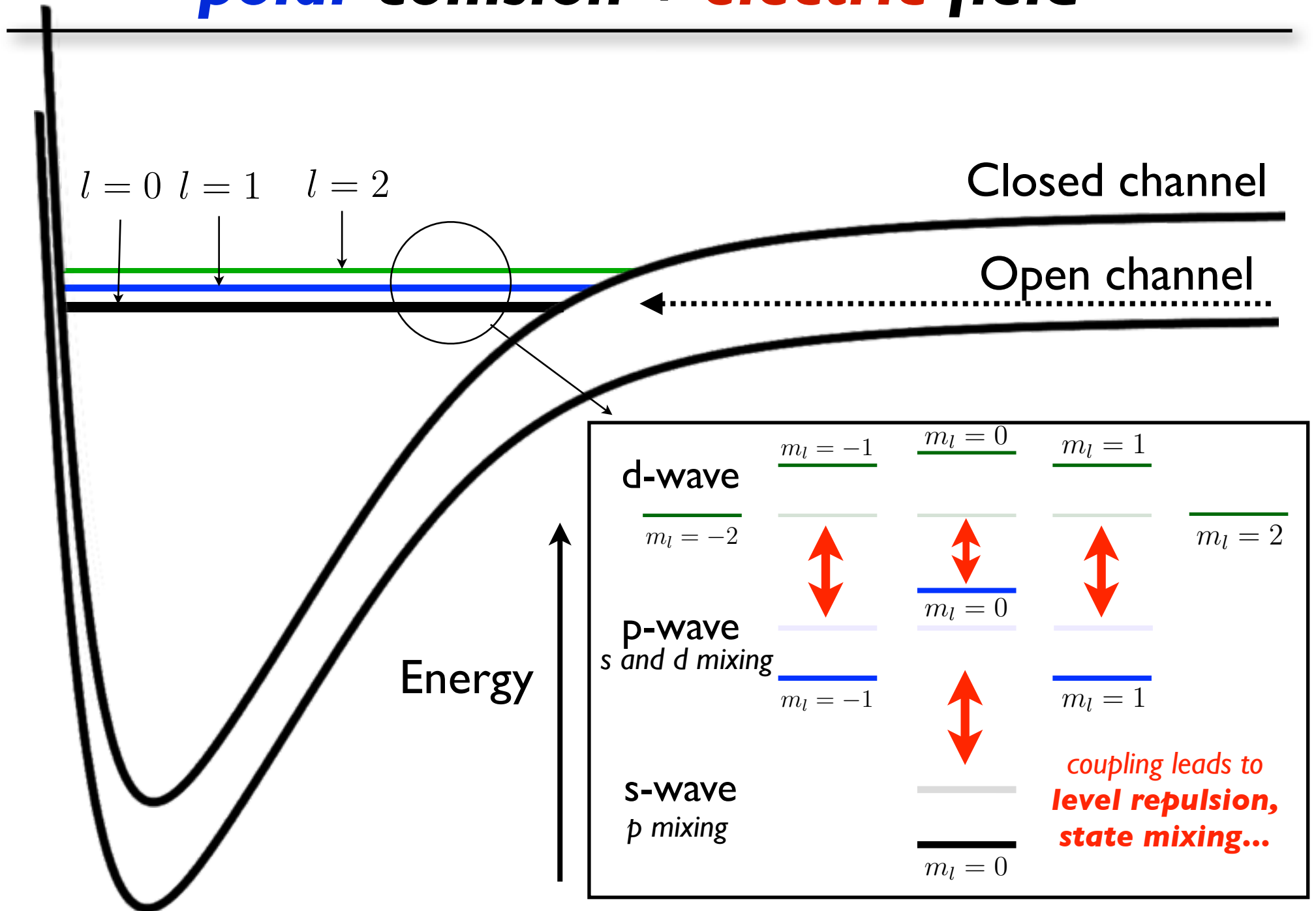
A new twist on Feshbach resonances: *polar collision* + *electric field*



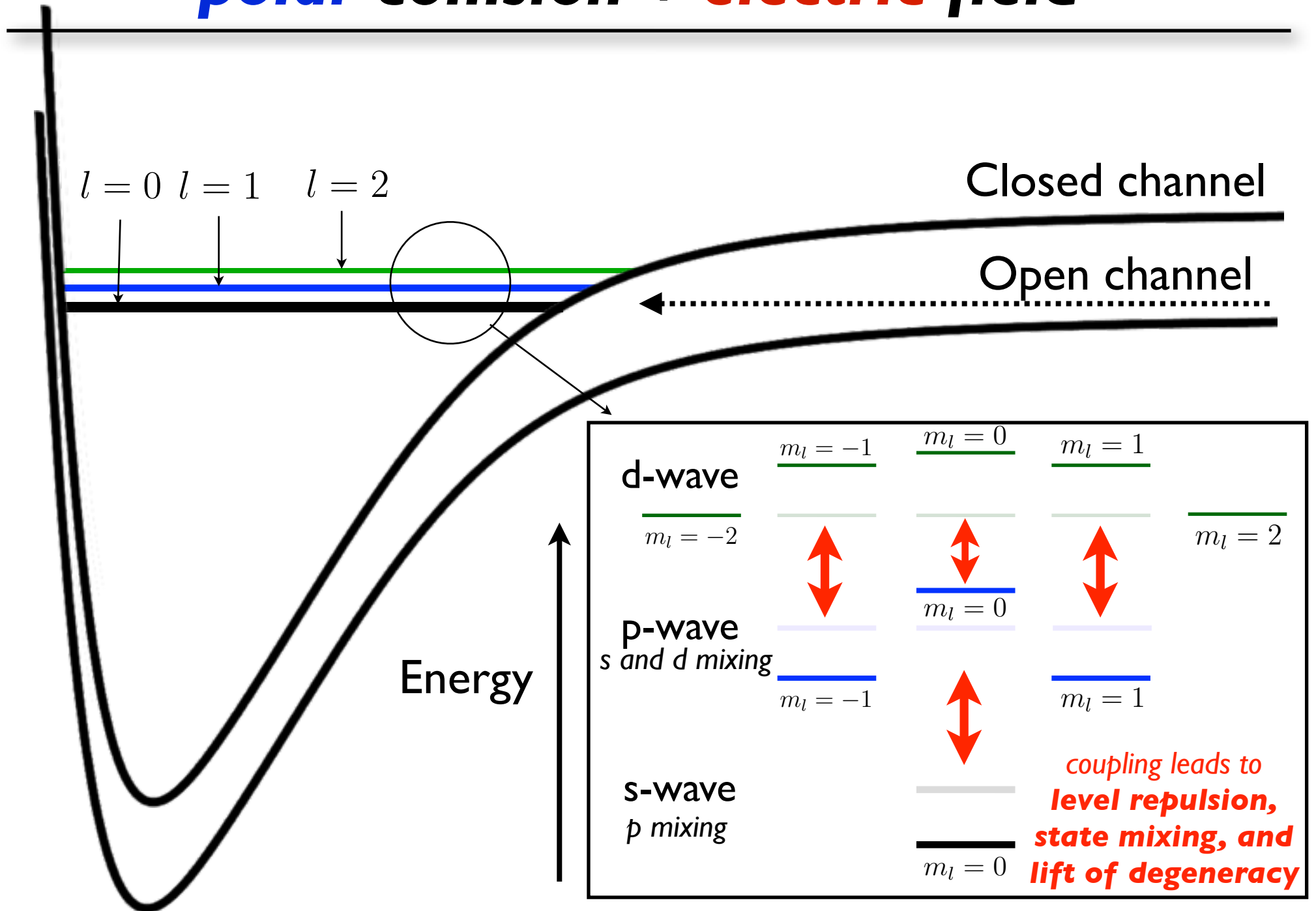
A new twist on Feshbach resonances: *polar* collision + *electric* field



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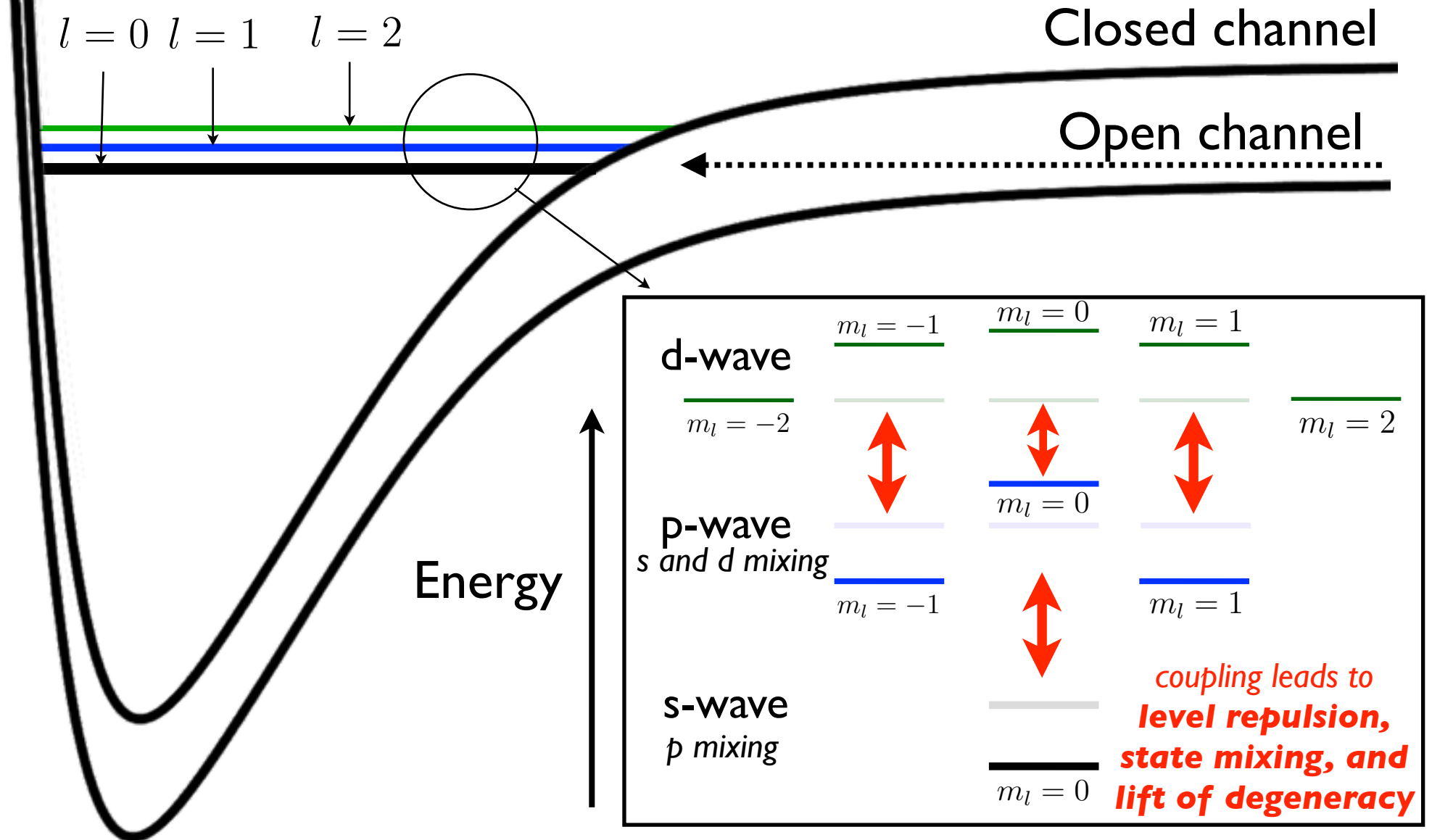


A new twist on Feshbach resonances: *polar collision* + *electric field*



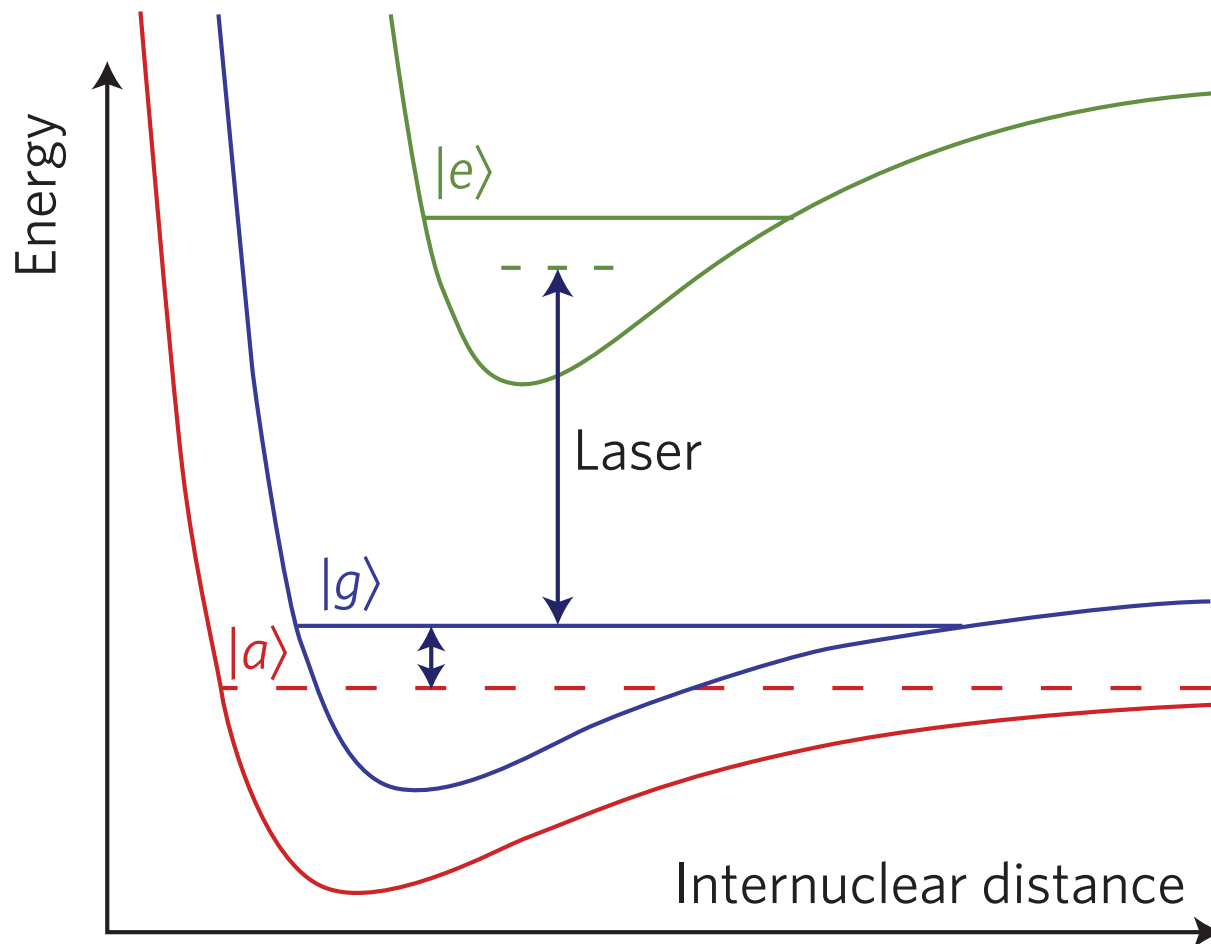
A new twist on Feshbach resonances: *polar* collision + *electric* field

Effects of electric fields on heteronuclear Feshbach resonances in ultracold $6\text{Li}-87\text{Rb}$ mixtures, PRA 79, 042711 (2009).



Control of a magnetic Feshbach resonance with laser light

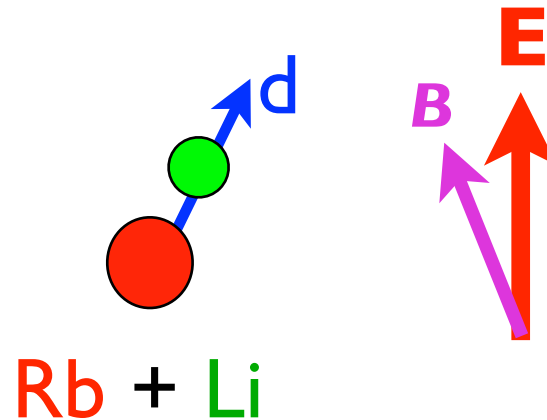
Dominik M. Bauer, Matthias Lettner, Christoph Vo, Gerhard Rempe and Stephan Dürr*



A new twist on Feshbach resonances:

Effects of electric fields on heteronuclear Feshbach resonances in ultracold 6Li - 87Rb mixtures, PRA 79, 042711 (2009).

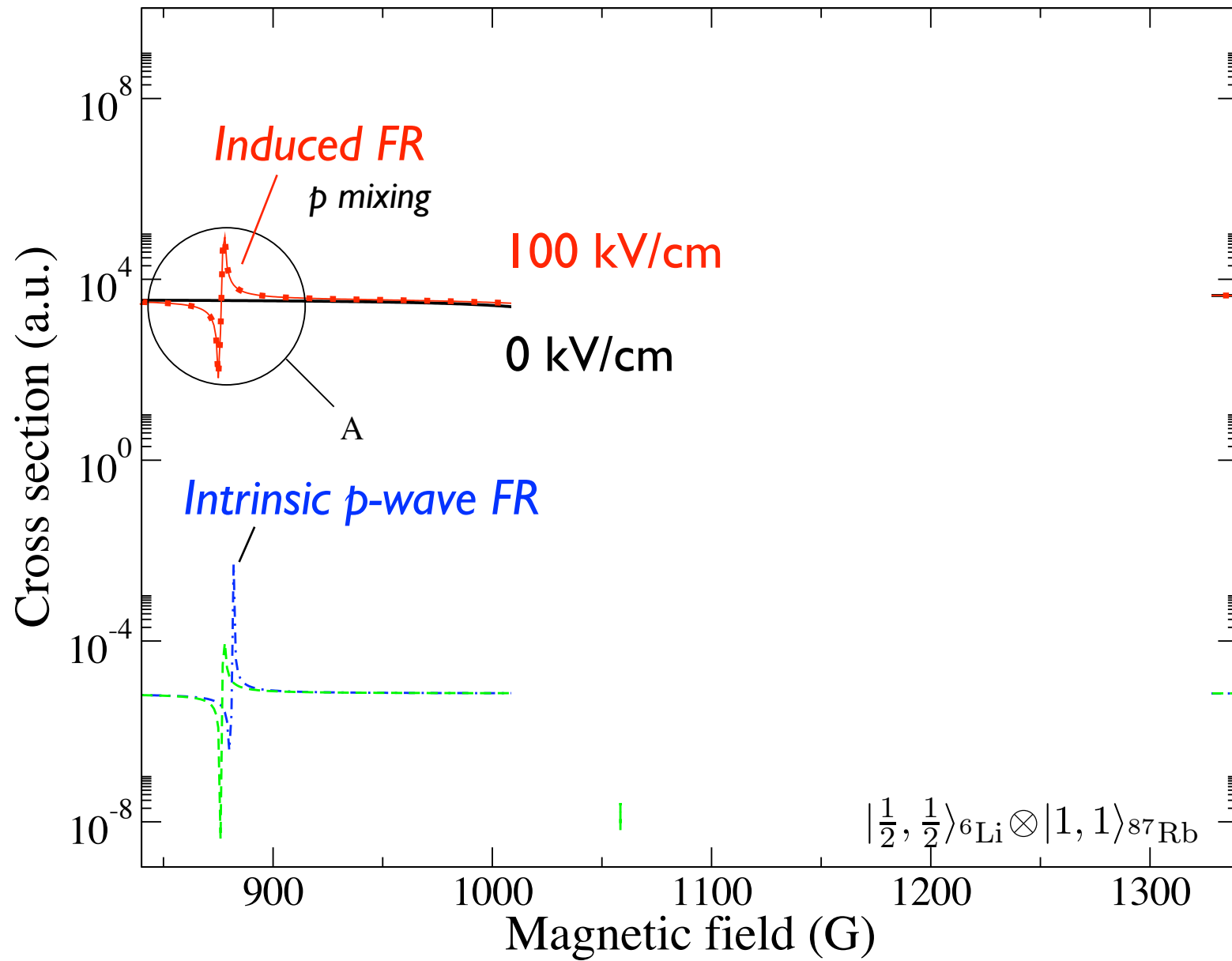
polar collision + electric field



- shift position of intrinsic Feshbach resonances
bound state level repulsion
- induce new Feshbach resonances
mixing of partial waves
- split resonances into $l+1$ multiplets
tunable anisotropic interactions

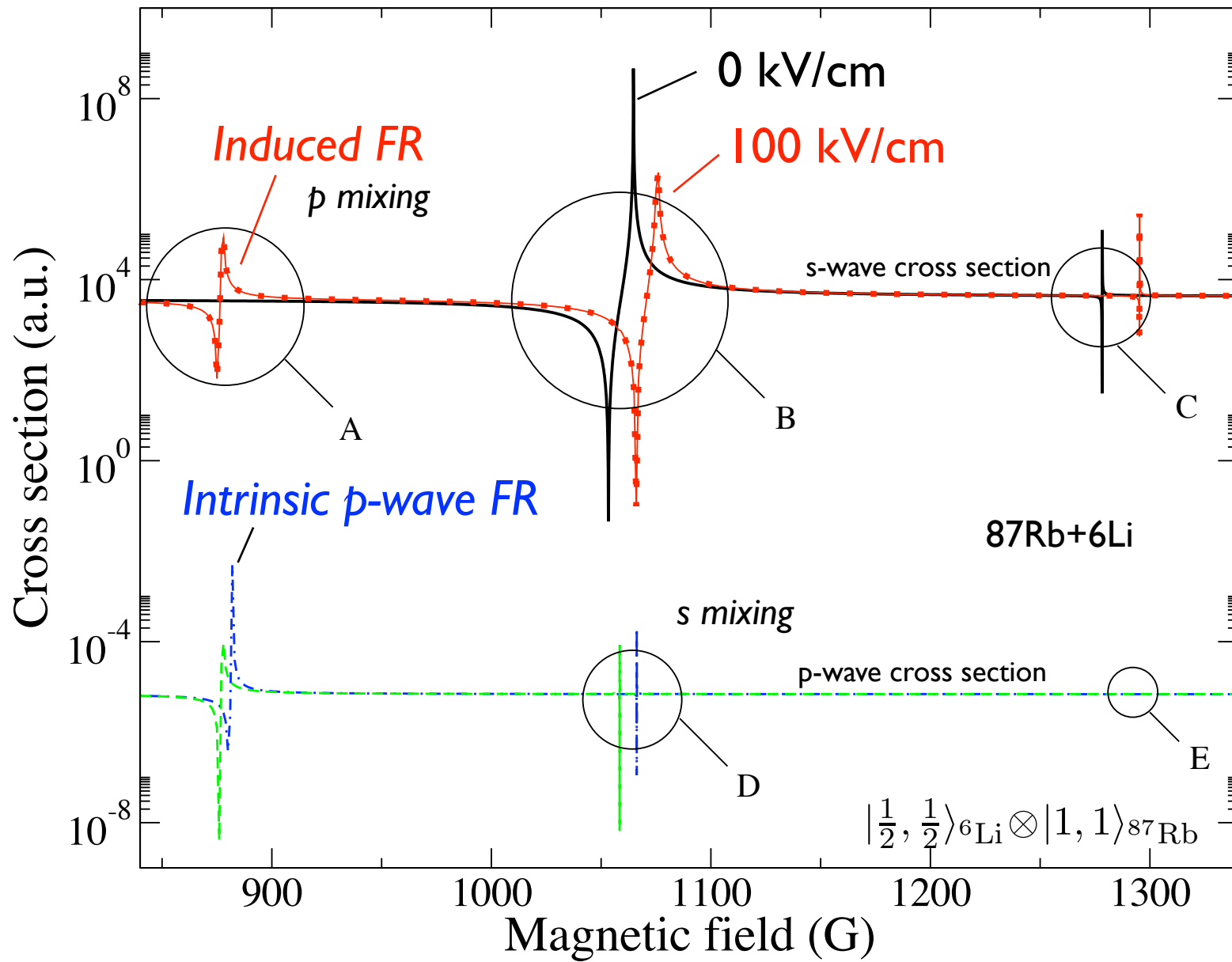
A new twist on Feshbach resonances: *Induced resonances*

new s-wave resonances appear at the location of intrinsic ($E=0$) p-wave resonances... and vice-versa



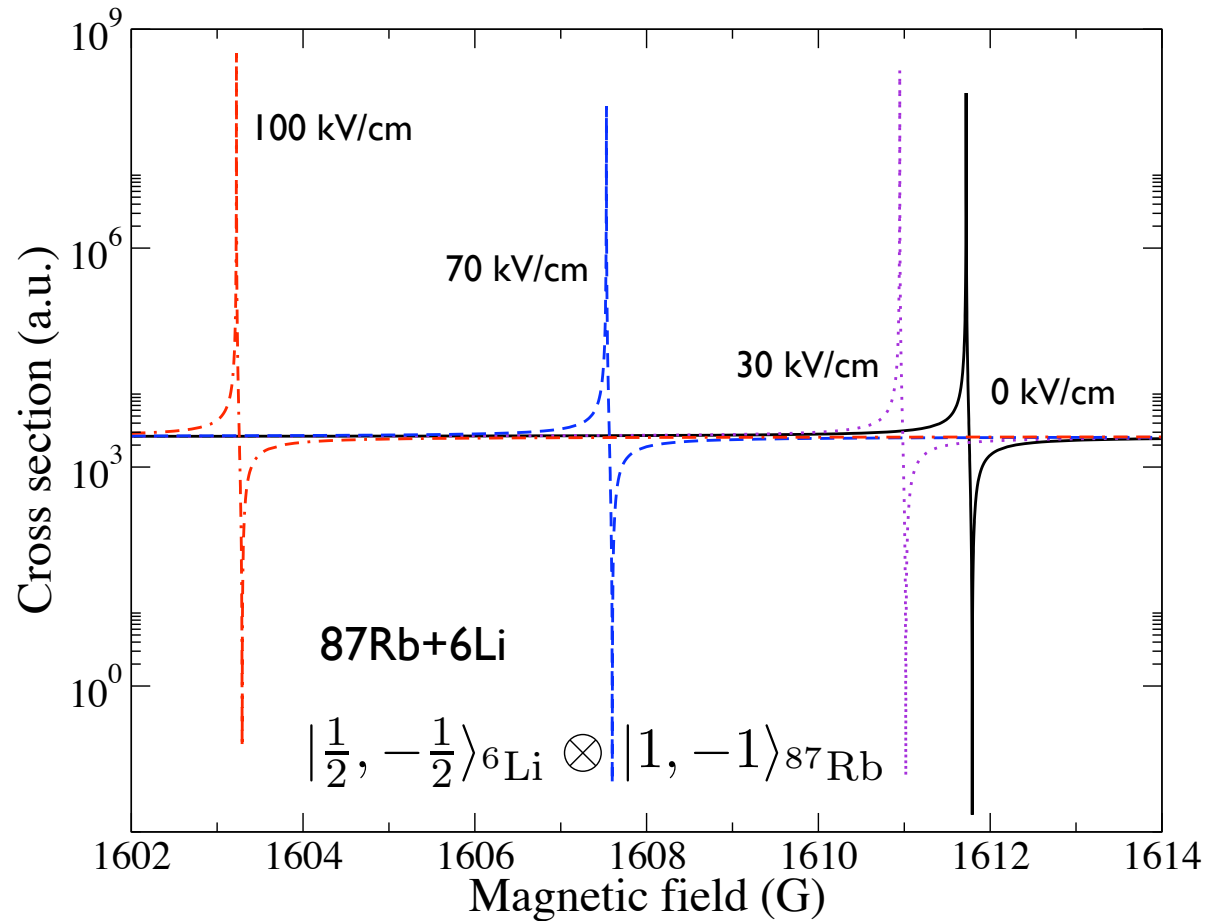
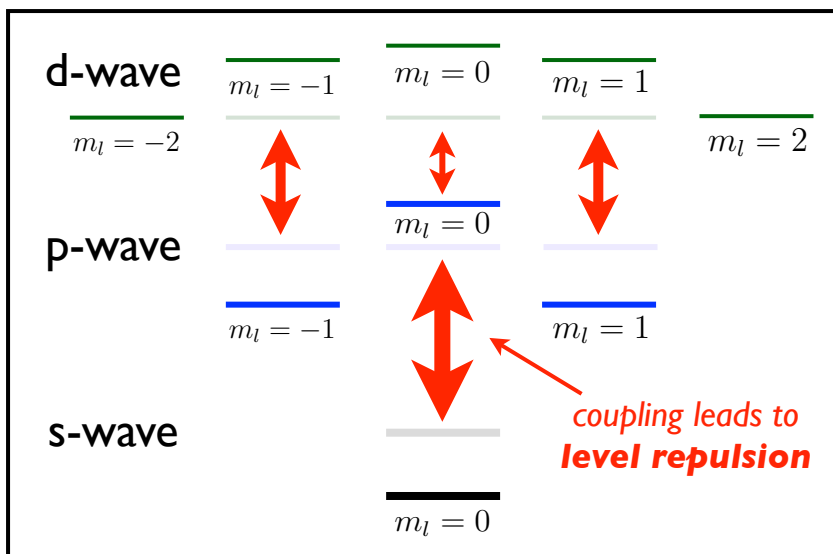
A new twist on Feshbach resonances: *shift of resonances*

intrinsic s-wave resonances shift with applied electric field



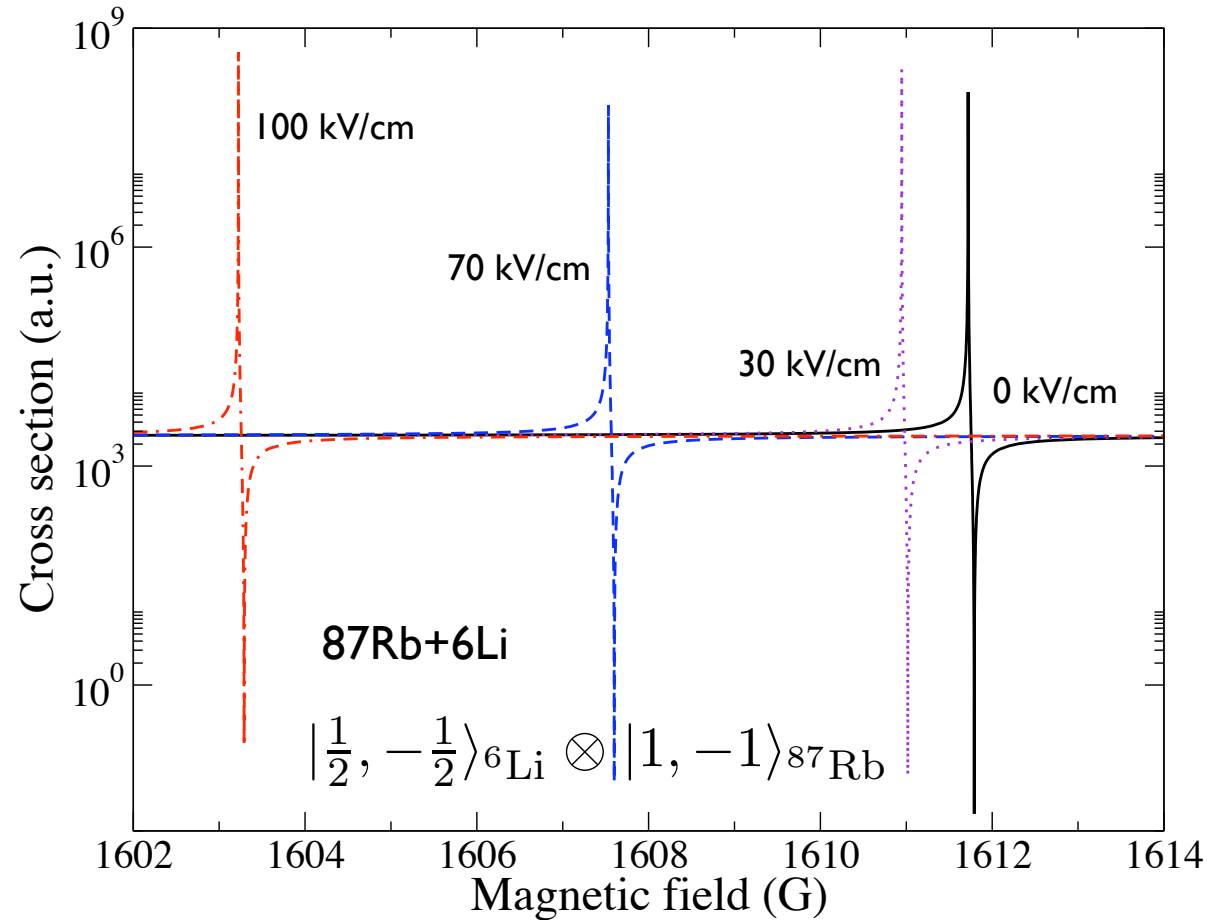
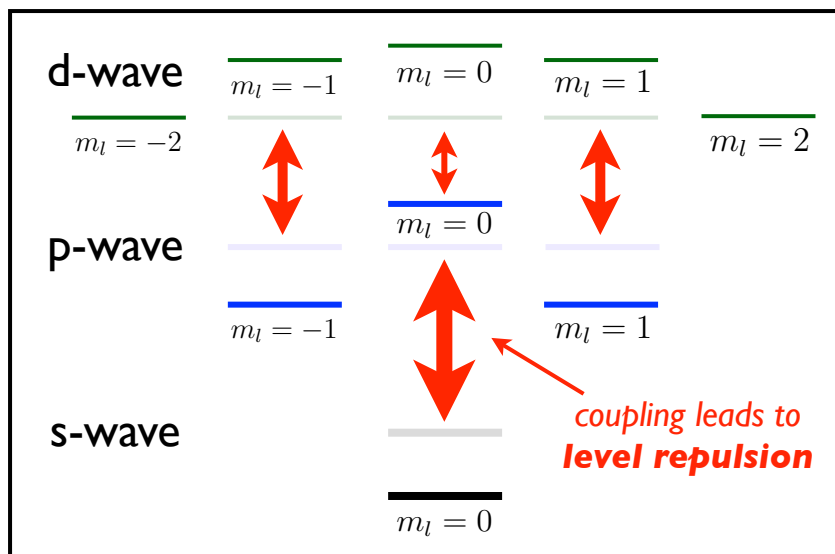
Shift of resonances:

intrinsic ($E=0$) resonances shift in position as magnitude of electric field increased

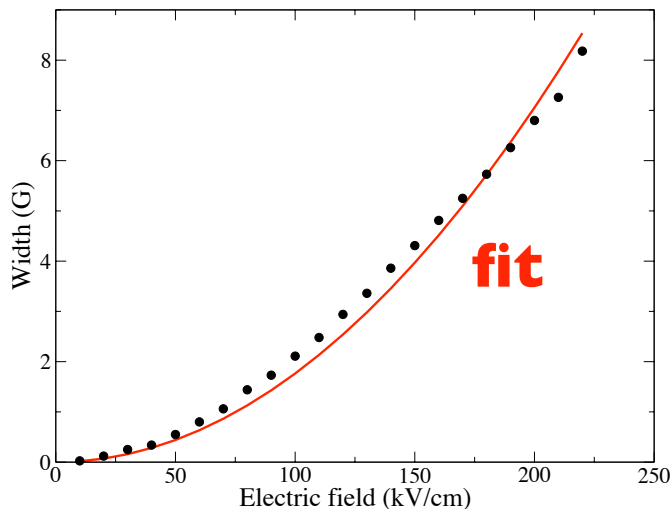


Shift of resonances:

intrinsic ($E=0$) resonances shift in position as magnitude of electric field increased



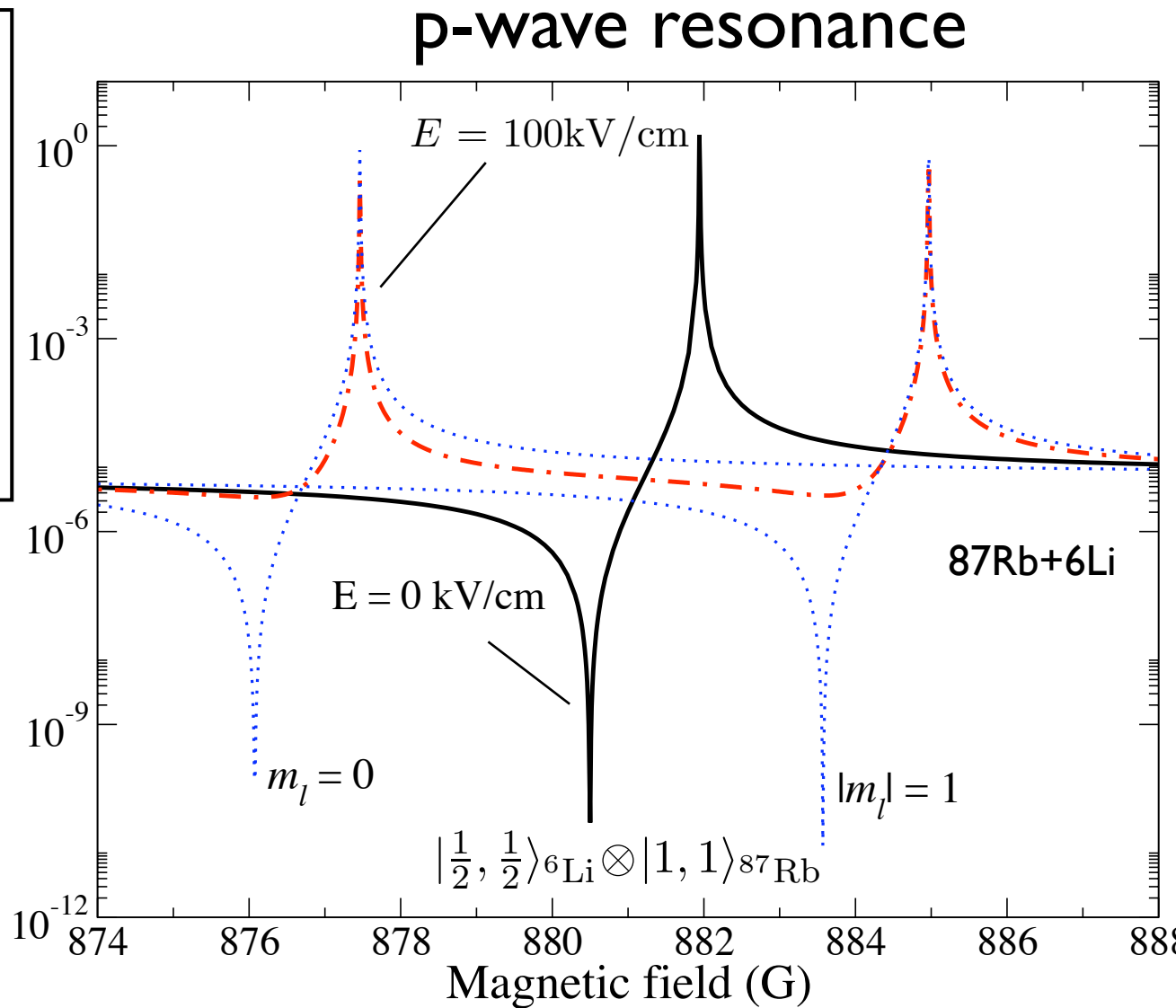
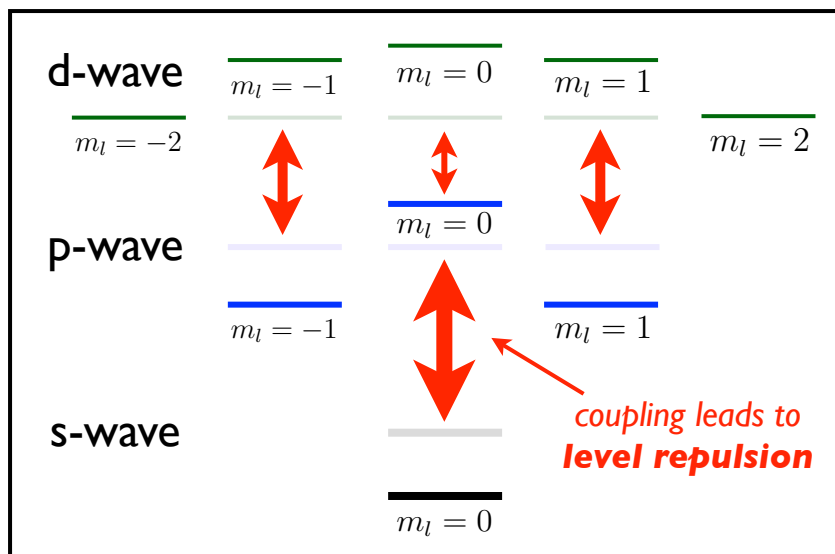
FR **shift** and induced FR **width** are both quadratic in E



FRs with tunable positions and widths

Splitting of resonances:

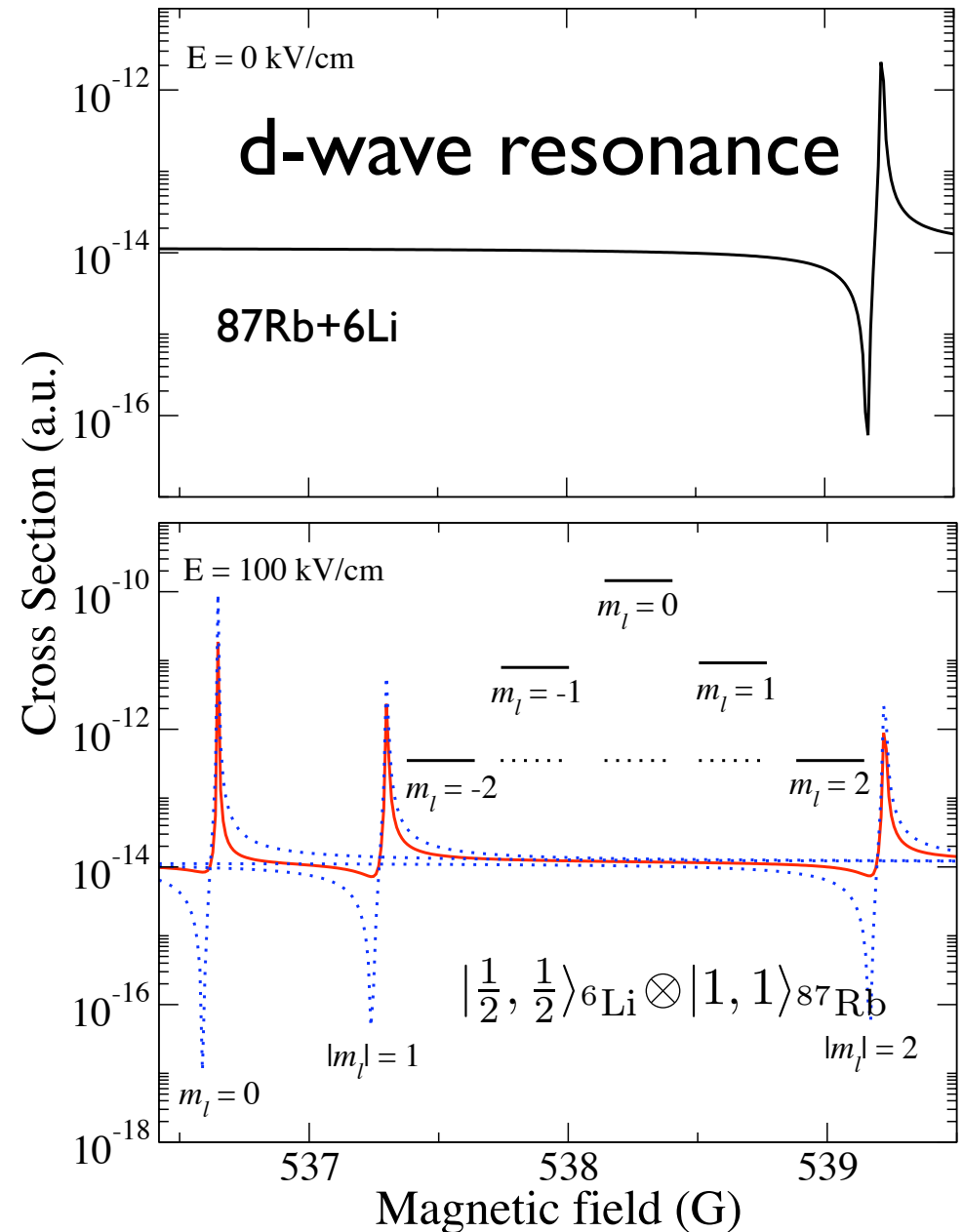
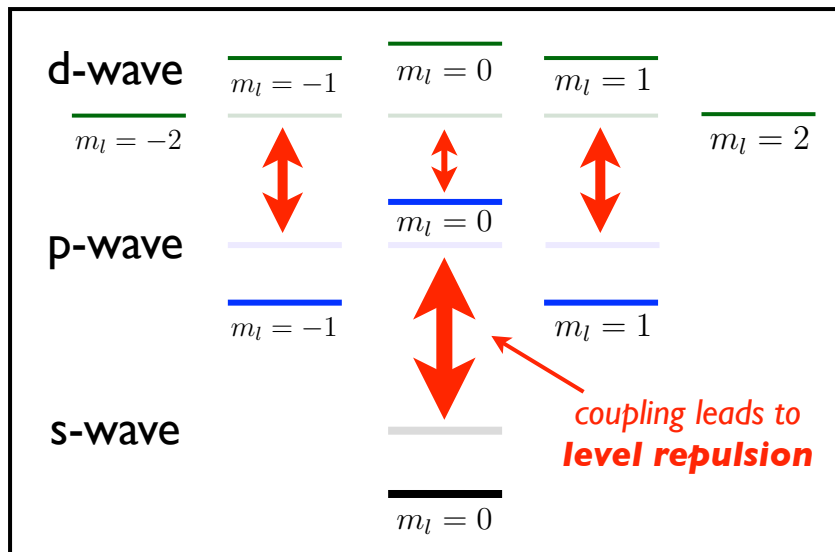
partial wave resonances split into $l+1$ distinct resonances

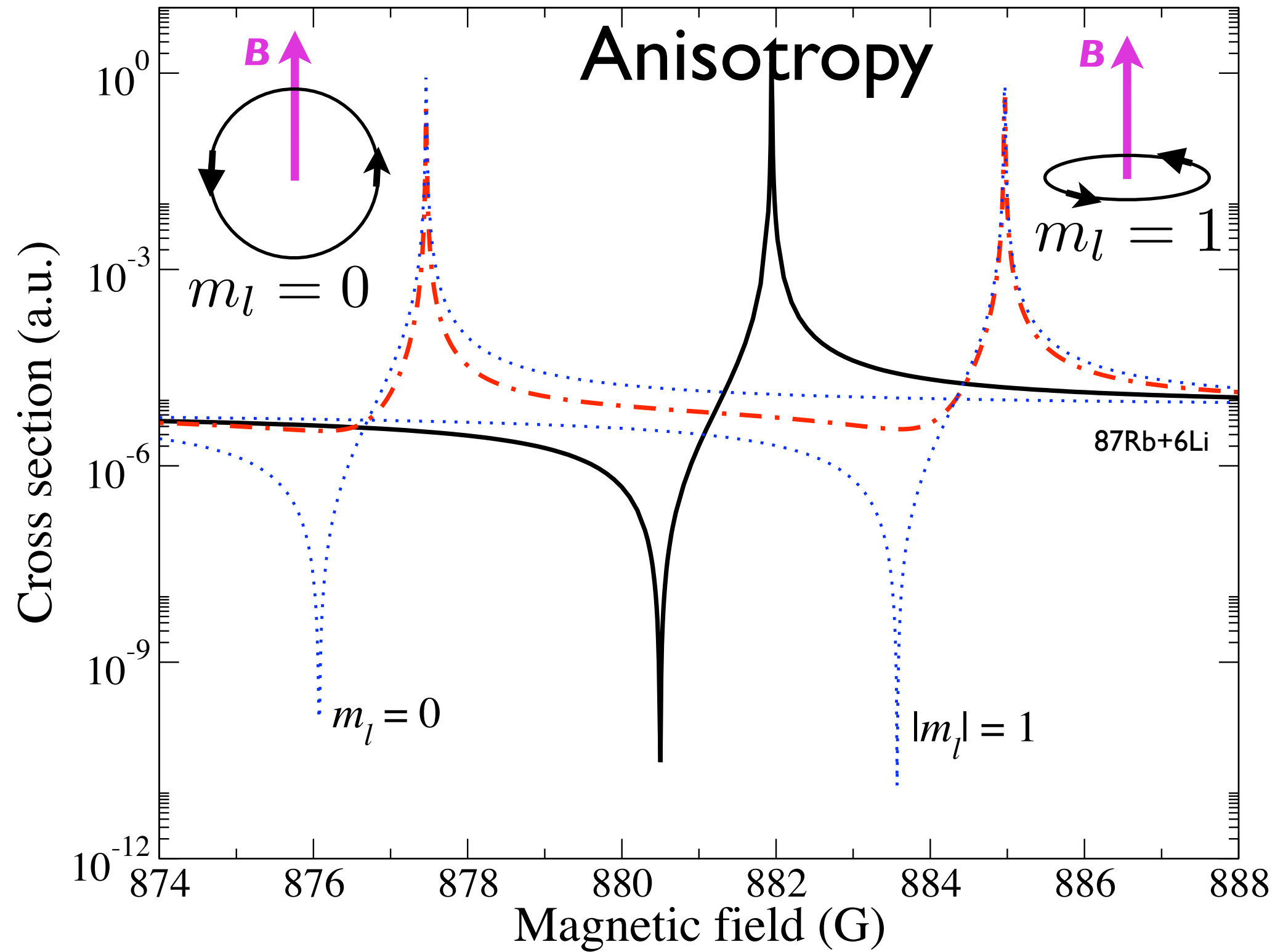


this is similar to splitting due to dipole-dipole coupling
 Regal, Ticknor, Bohn, Jin, PRL **90**, 053201 (2003), and PRA **69**, 042712 (2004).

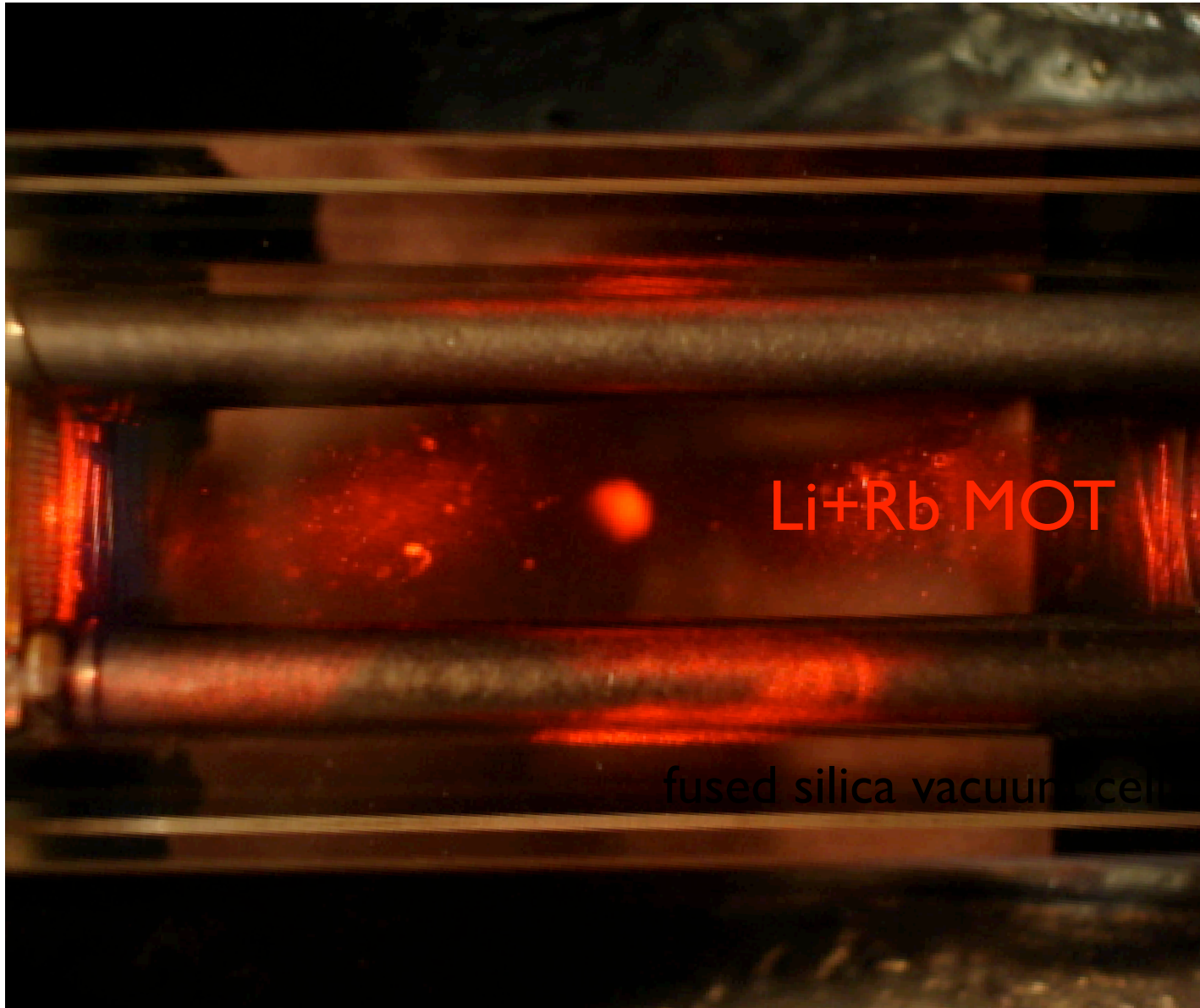
Splitting of resonances:

partial wave resonances split into
 $l+1$ distinct resonances





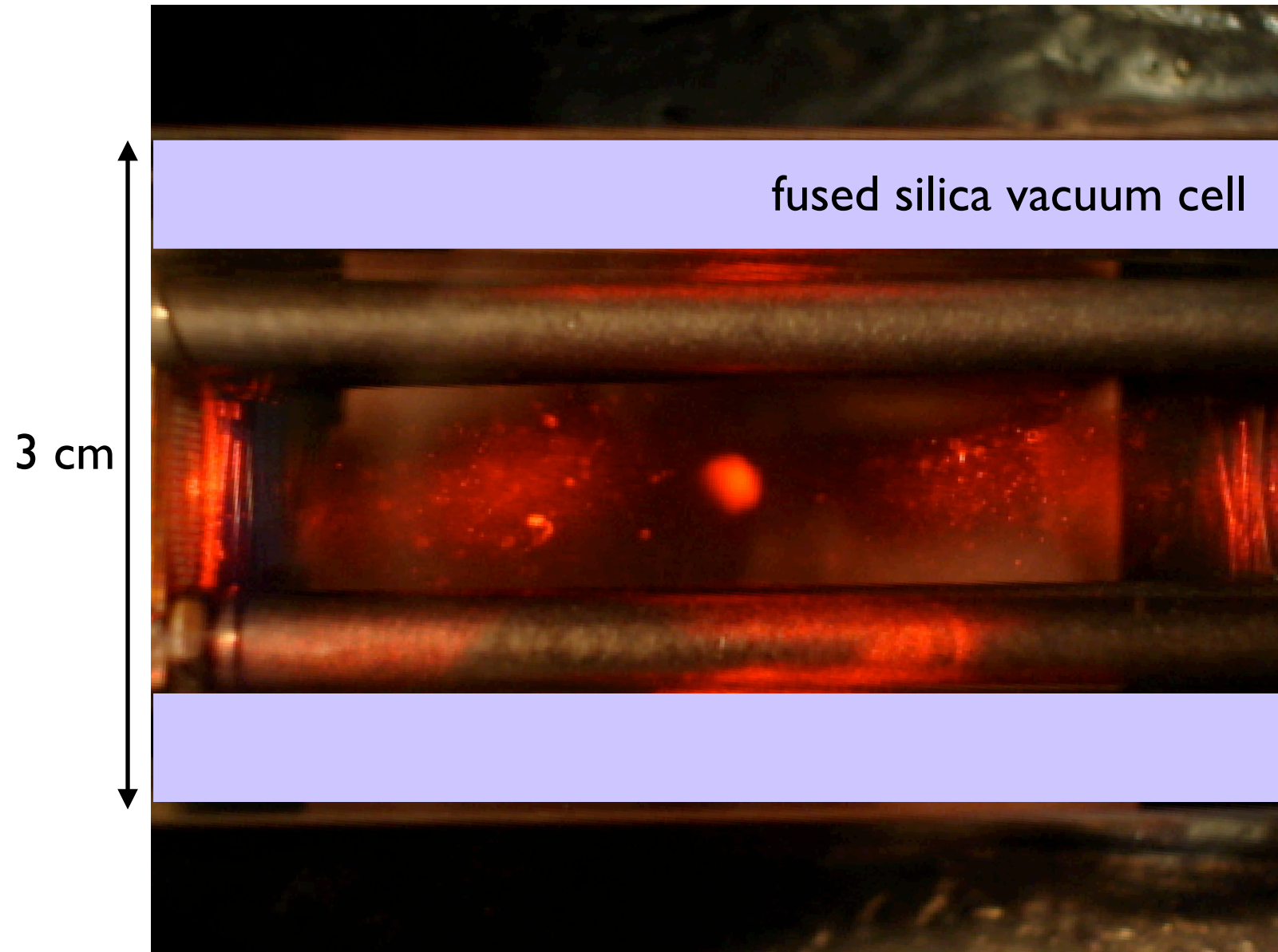
Experimental challenge: adding the E field



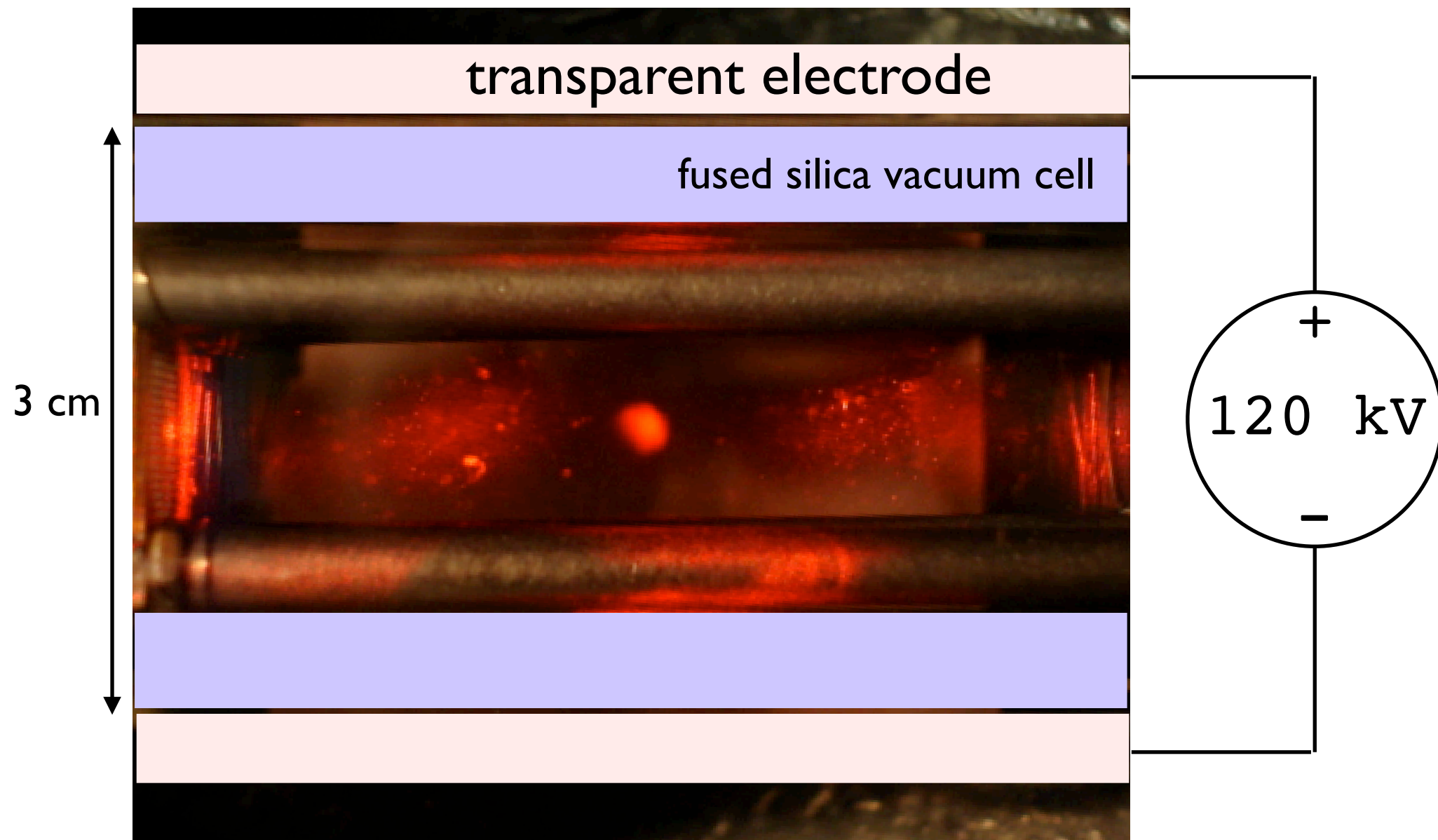
*we want a
field of about
30 kV/cm*

*how to add
electrodes
while keeping
good
optical access?*

Experimental challenge: adding the E field



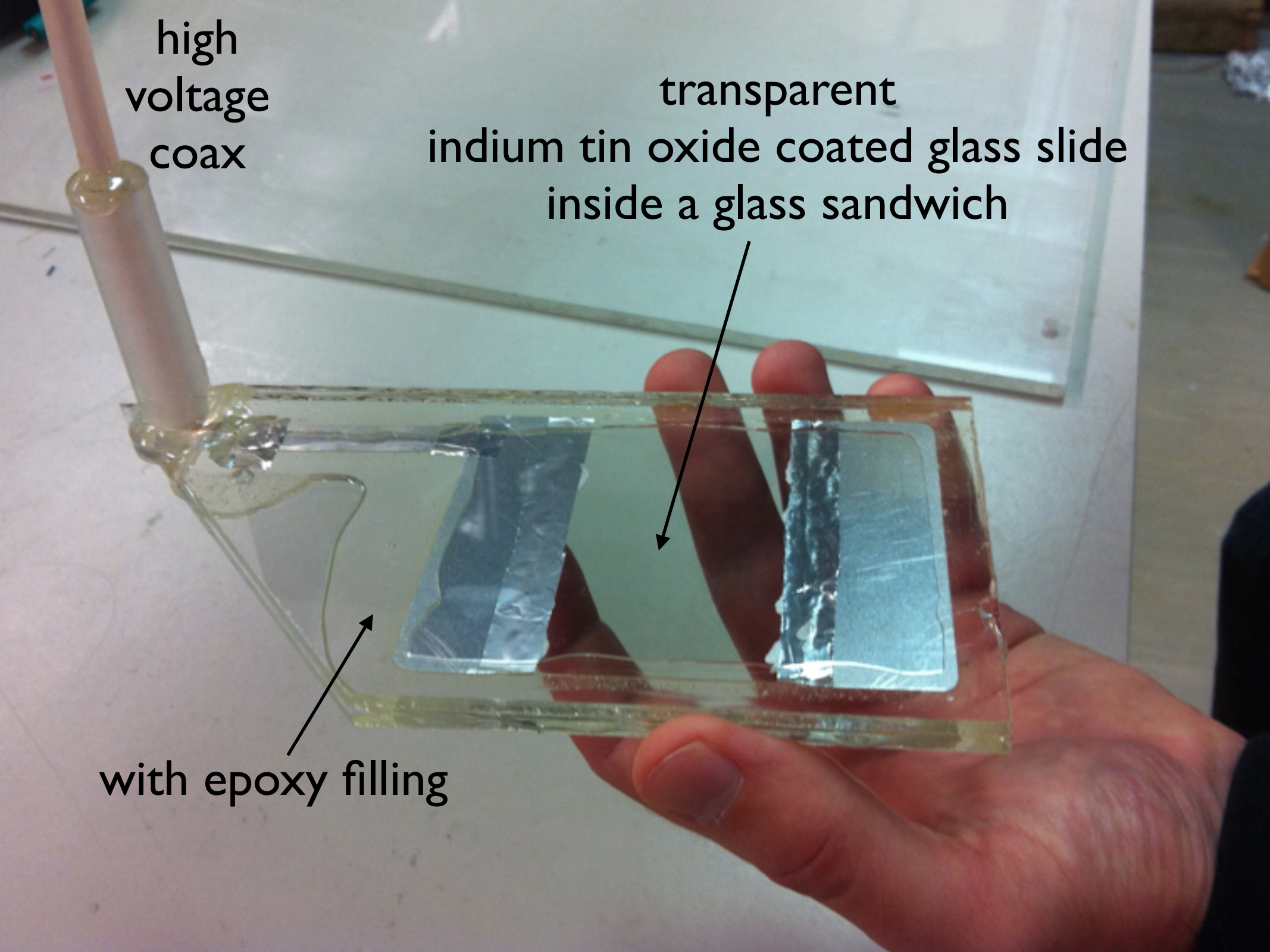
Experimental challenge: adding the E field



high
voltage
coax

transparent
indium tin oxide coated glass slide
inside a glass sandwich

with epoxy filling





Experimental challenges and mysteries

Status:

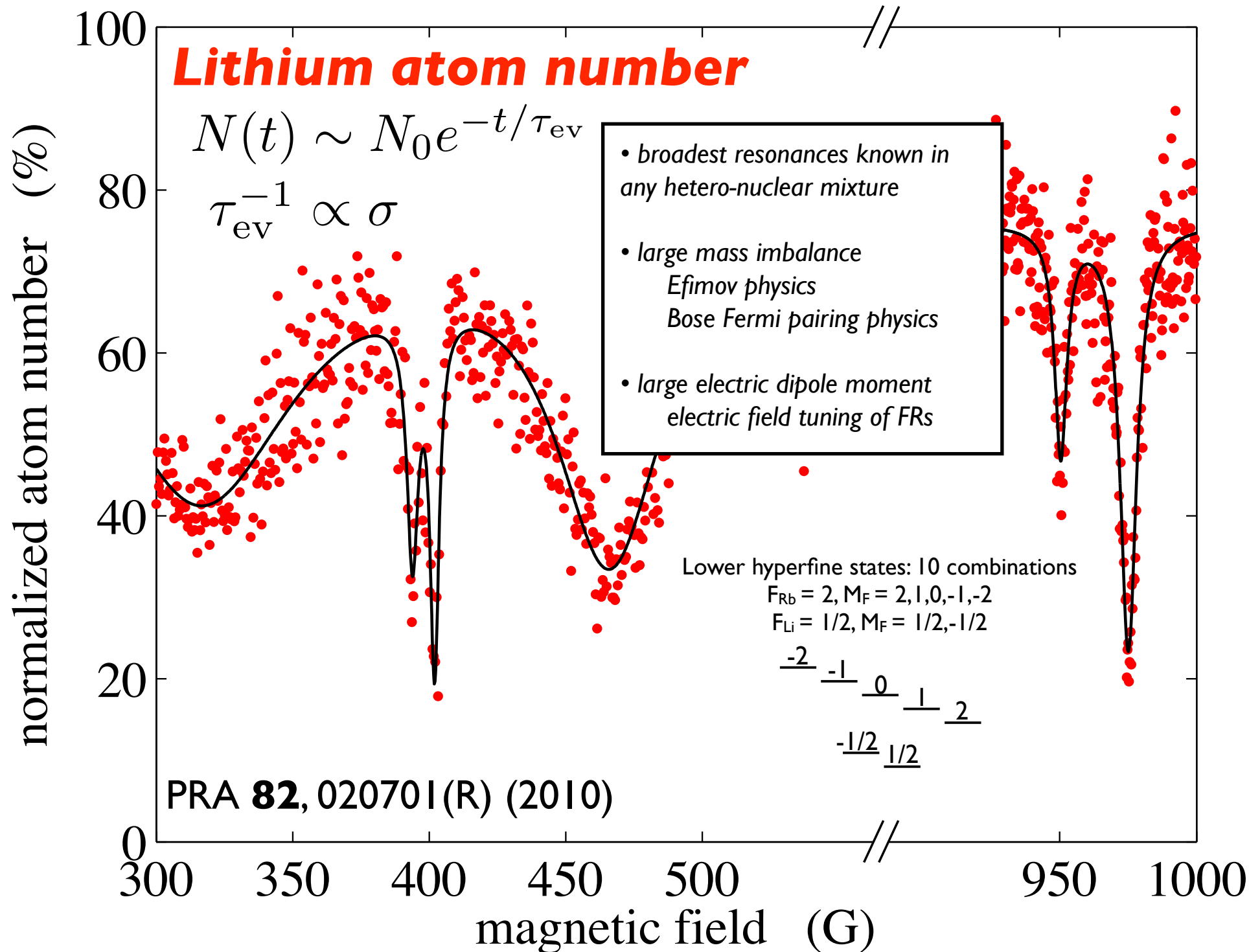
We have generated and verified fields of 18kV/cm and have another factor of 2 in voltage.

Mysteries:

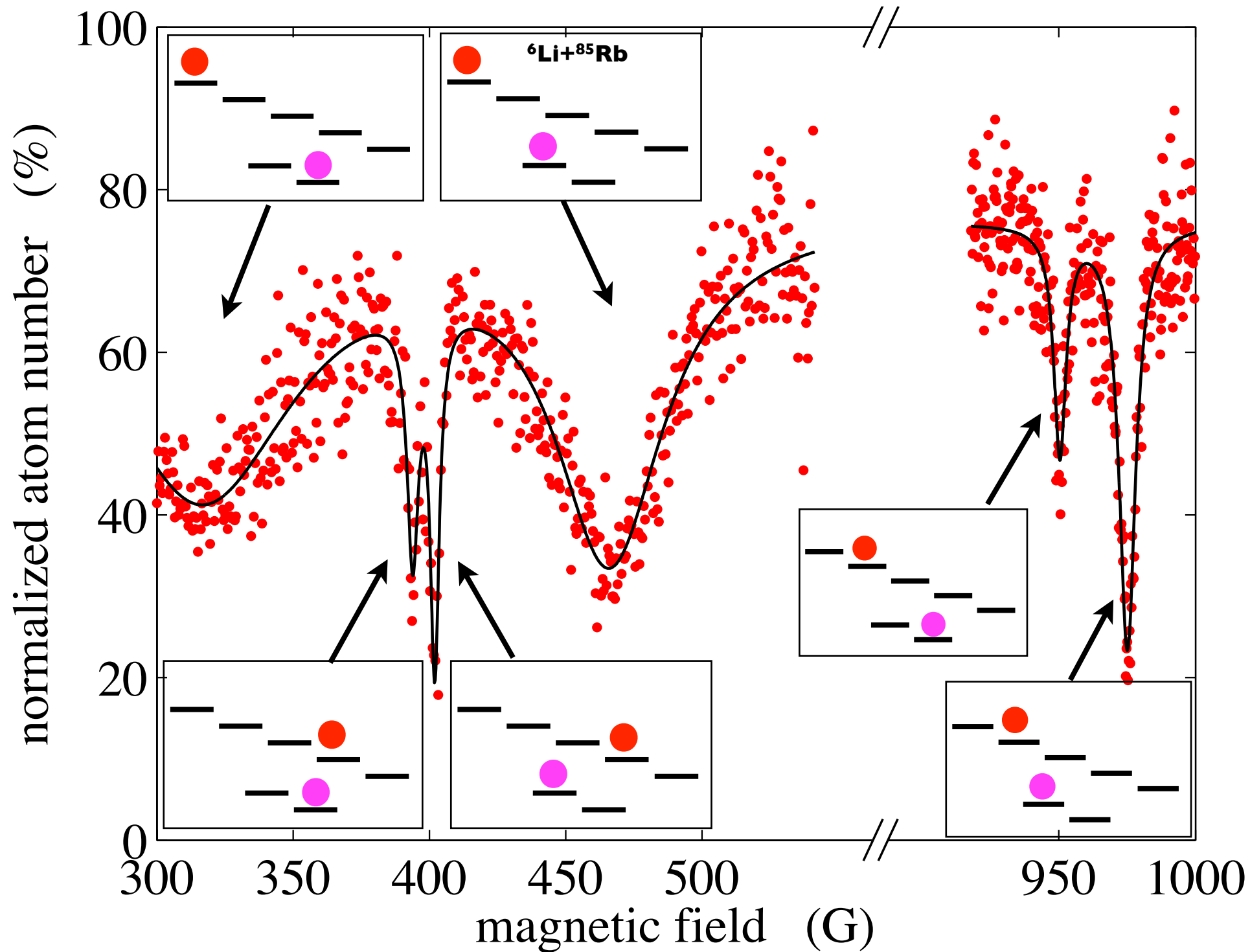
- 1) Field peculiarities (evidence of shielding)***
- 2) Trap loss associated with energizing plates***

**Feshbach resonances in ${}^6\text{Li}+{}^{85}\text{Rb}$ mixtures
the strength of resonances**

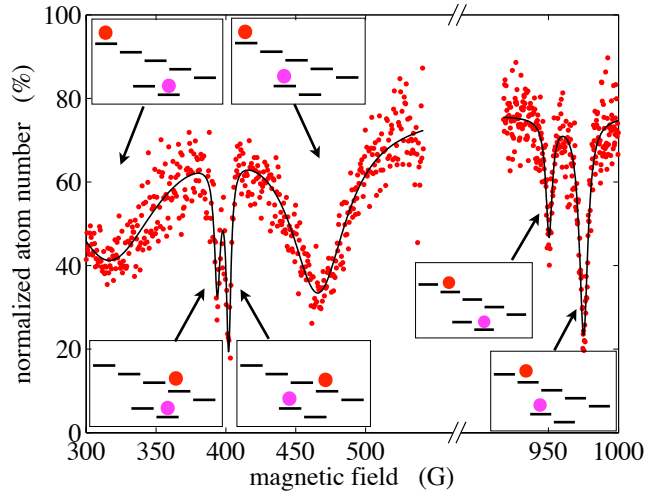
Feshbach resonances in ${}^6\text{Li}+{}^{85}\text{Rb}$ mixtures



Feshbach resonances in ${}^6\text{Li}+{}^{85}\text{Rb}$ mixtures



Feshbach resonances in ${}^6\text{Li}+{}^{85}\text{Rb}$ mixtures



$$\text{effective range: } r_e = \hbar^2 / (2m_R |a_{\text{bg}}| \mu_{\text{rel}} \Delta B)$$

$$\text{resonance strength: } E_{\text{res}} = \hbar^2 / (2m_R r_e^2)$$

TABLE II. The calculated characteristics of seven large FRs for stable Fermi-Bose mixtures of ${}^6\text{Li}$ - ${}^{85}\text{Rb}$ and ${}^6\text{Li}$ - ${}^{87}\text{Rb}$. The background scattering length a_{bg} and effective range of the resonance r_e are given in units of the Bohr radius, a_0 . The value for μ_{rel} , the difference in the magnetic moments of the closed channel (molecule) and the open-channel threshold, is given in units of the Bohr magneton $\mu_B = 9.27400915(23) \text{ J T}^{-1}$. For each FR, the magnetic fields at which the mixture is energetically stable with respect to two-body spin relaxation is provided. In some cases, because of nearby resonances, we can only provide an approximate lower bound on the resonance width.

Atomic states $ f, m_f\rangle \otimes f, m_f\rangle$	B_0 (G)	ΔB (G)	μ_{rel} (μ_B)	a_{bg} (a_0)	r_e (a_0)	Stability (G)
${}^6\text{Li}$ - ${}^{85}\text{Rb}$						
$ \frac{1}{2}, \frac{1}{2}\rangle \otimes 2, 2\rangle$	40.7	>40	1.66	-14.9	<231	ground state
$ \frac{1}{2}, -\frac{1}{2}\rangle \otimes 2, 1\rangle$	402.5	27.3	1.58	-14.9	358	≥ 149
$ \frac{1}{2}, -\frac{1}{2}\rangle \otimes 2, 0\rangle$	643.7	61.0	1.34	-14.9	189	≥ 141
$ \frac{1}{2}, -\frac{1}{2}\rangle \otimes 2, -1\rangle$	961.3	75.6	1.81	-14.7	113	≥ 133
$ \frac{1}{2}, -\frac{1}{2}\rangle \otimes 2, -2\rangle$	466.7	>100	0.58	-14.8	<264	≥ 0
${}^6\text{Li}$ - ${}^{87}\text{Rb}$						
$ \frac{1}{2}, \frac{1}{2}\rangle \otimes 1, 1\rangle$	1065.0	11.5	2.36	-19.0	442	ground state
$ \frac{1}{2}, -\frac{1}{2}\rangle \otimes 1, 1\rangle$	1108.6	11.0	2.36	-19.0	463	≥ 75