
KITP – UCSB, 3 February 2004

*Automated resummation of QCD
final state observables*

Giulia Zanderighi

☞ In collaboration with

A. Banfi (Amsterdam) and G. Salam (Paris)



Fermilab

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$$T \equiv \frac{1}{Q} \max_{\vec{n}_T} \sum_i |\vec{p}_i \cdot \vec{n}_T| = \frac{1}{Q} \sum_i |p_{iz}|$$

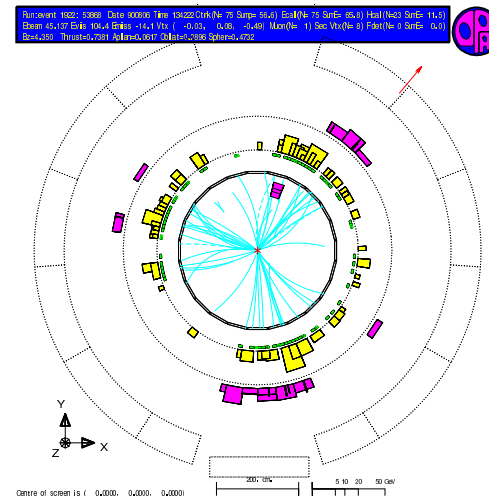
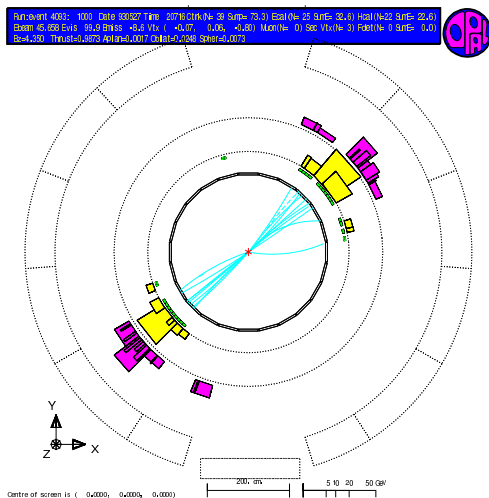
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Pencil-like event: $\tau \equiv 1 - T \ll 1$

Planar event: $T \simeq 2/3$



Event shapes are a **good compromise** between

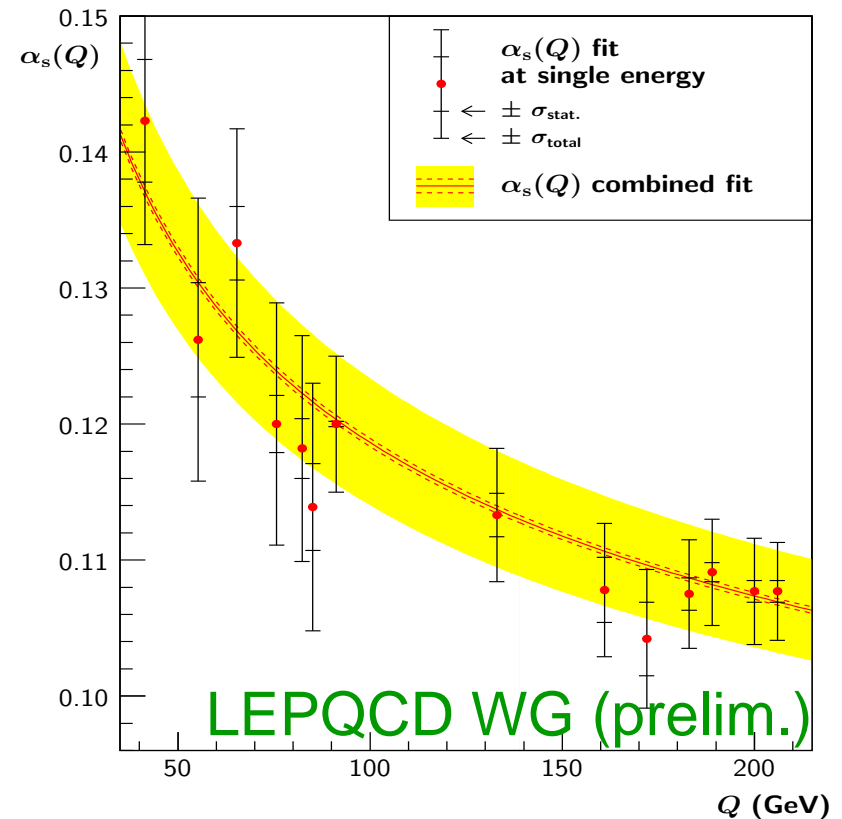
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- **sensitivity** to properties of QCD radiation

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Provide a wealth of information, e.g.:

- Measurements of the coupling α_s and its **renormalization group running**

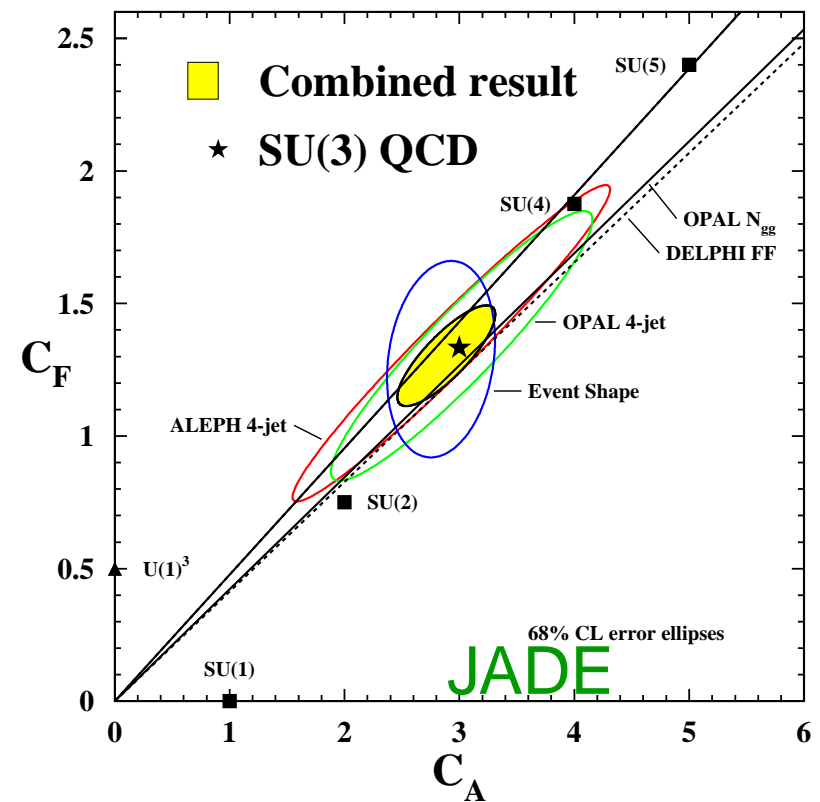


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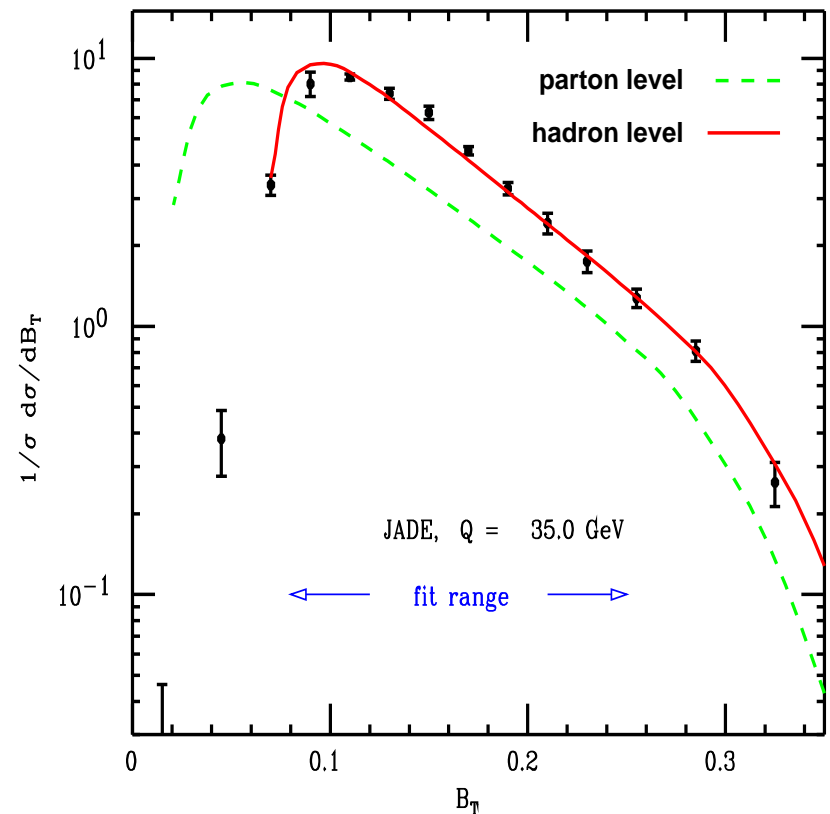


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Provide a wealth of information, e.g.:

- Measurements of the coupling α_s and its **renormalization group running**
- Measurements/cross checks of the values of the **colour factors** of QCD
- Studies of connection between **parton-level** (perturbative description of quarks and gluons) and **hadron-level** (the real)



Given an event shape (or a jet-rate) V then

$V \ll 1 \equiv$ going to the **Born limit** \equiv **forbidding** gluon radiation

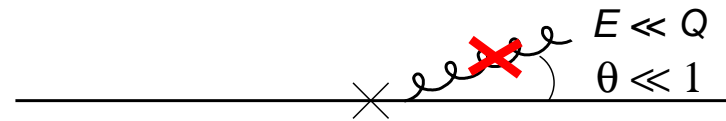
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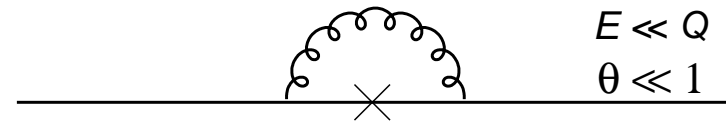
- **hard, large angle** emissions *forbidden* (no additional jet)
- **soft and collinear** real emissions are *constrained*

~~$$\frac{dE}{E} \frac{d\theta}{\theta} \alpha_s(\theta E)$$~~



- **virtual** corrections are *unaffected*

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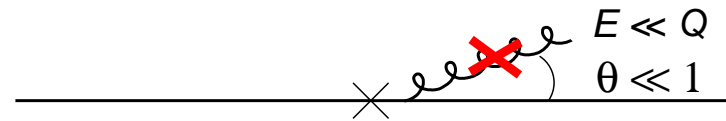
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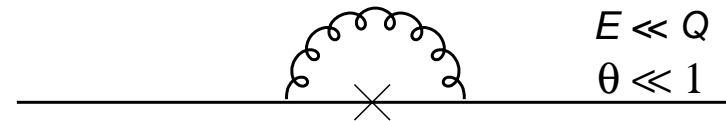
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Imbalance leads to **large logarithms** in distributions

$$\text{Prob}(V < v) \simeq 1 - \frac{\# \alpha_s C_F}{2\pi} \ln^2 v + \dots \quad [v \ll 1 \Rightarrow \frac{\alpha_s C_F}{2\pi} \ln^2 v = \mathcal{O}(1)]$$

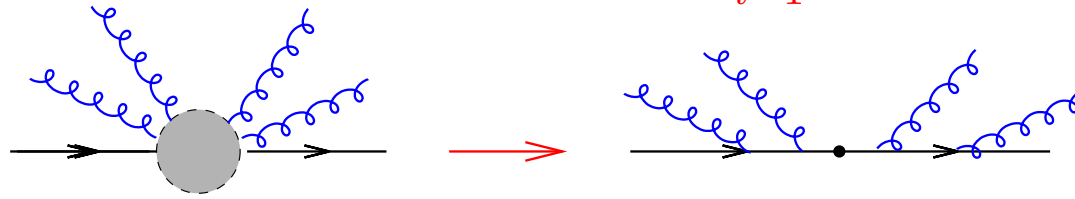
which need to be **resummed to all orders**

Basics of resummation: factorization

First half of the history: **Matrix elements and phase space**

exploit *angular ordering* \Rightarrow soft *independent emissions* (\Rightarrow QED)

e.g. $e^+e^- \rightarrow 2 \text{ jets} \Rightarrow w_{p\bar{p}}(k_1, \dots, k_n) = \frac{1}{n!} \prod_{i=1}^n w_{p\bar{p}}(k_i) \sim \frac{1}{n!} \prod_{i=1}^n \frac{dE}{E} \frac{d\theta}{\theta}$

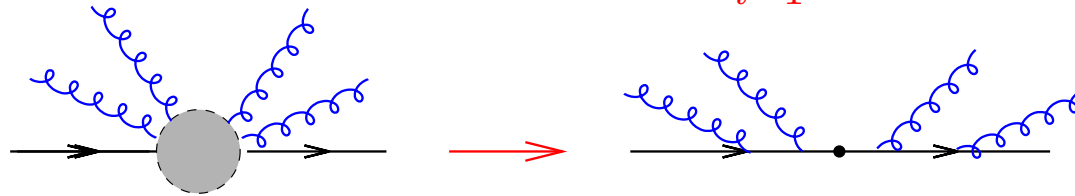


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Second half of the history: **The observable definition**

analyse the observable & use Mellin transforms

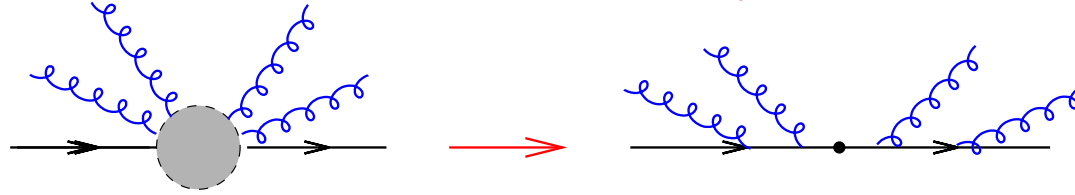
$$1 - T \simeq \frac{1}{Q} \sum_{i=1}^n \frac{E_i \theta_i^2}{2} \quad \longrightarrow \quad \Theta(1 - T < \tau) = \int \frac{d\nu}{2\pi i \nu} e^{\nu\tau} \prod_{i=1}^n e^{-\nu \frac{E_i \theta_i^2}{2Q}}$$

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THE ANSWER

$$\Sigma(\tau) \int \frac{d\nu}{2\pi i \nu} e^{\nu\tau} \exp \left[\int \frac{d\theta}{\theta} \frac{dE}{E} \alpha_s(E\theta) \left(e^{-\nu \frac{E_i \theta_i^2}{2Q}} - 1 \right) \right]$$

A selection of analytical NLL predictions

$e^+e^- \rightarrow 2$ jets

- ☛ S. Catani, G. Turnock, B. R. Webber and L. Trentadue, *Thrust distribution in e^+e^- annihilation*, Phys. Lett. B **263** (1991) 491.
- ☛ S. Catani, G. Turnock and B. R. Webber, *Heavy jet mass distribution in e^+e^- annihilation*, Phys. Lett. B **272** (1991) 368.
- ☛ S. Catani, Yu. L. Dokshitzer, M. Olsson, G. Turnock and B. R. Webber, *New clustering algorithm for multi-jet cross-sections in e^+e^- annihilation*, Phys. Lett. B **269** (1991) 432.
- ☛ S. Catani, L. Trentadue, G. Turnock and B. R. Webber, *Resummation of large logarithms in e^+e^- event shape distributions*, Nucl. Phys. B **407** (1993) 3.
- ☛ S. Catani, G. Turnock and B. R. Webber, *Jet broadening measures in e^+e^- annihilation*, Phys. Lett. B **295** (1992) 269.
- ☛ G. Dissertori and M. Schmelling, *An Improved theoretical prediction for the two jet rate in e^+e^- annihilation*, Phys. Lett. B **361** (1995) 167.
- ☛ Y. L. Dokshitzer, A. Lucenti, G. Marchesini and G. Salam, *On the QCD analysis of jet broadening*, JHEP **9801** (1998) 011
- ☛ S. Catani and B. R. Webber, *Resummed C-parameter distribution in e^+e^- annihilation*, Phys. Lett. B **427** (1998) 377
- ☛ S. J. Burby and E. W. Glover, *Resumming the light hemisphere mass and narrow jet broadening distributions in e^+e^- annihilation*, JHEP **0104** (2001) 029
- ☛ M. Dasgupta and G. Salam, *Resummation of non-global QCD observables*, Phys. Lett. B **512** (2001) 323
- ☛ C. F. Berger, T. Kucs and G. Sterman, *Event shape / energy flow correlations*, Phys. Rev. D **68** (2003) 014012

DIS 1+1 jet

- ☛ V. Antonelli, M. Dasgupta and G. Salam, *Resummation of thrust distributions in DIS*, JHEP **0002** (2000) 001
- ☛ M. Dasgupta and G. Salam, *Resummation of the jet broadening in DIS*, Eur. Phys. J. C **24** (2002) 213
- ☛ M. Dasgupta and G. Salam, *Resummed event-shape variables in DIS*, JHEP **0208** (2002) 032

e^+e^- , DY, DIS 3 jets

- ☛ A. Banfi, G. Marchesini, Y. L. Dokshitzer and GZ, *QCD analysis of near-to-planar 3-jet events*, JHEP **0007** (2000) 002
- ☛ A. Banfi, Y. L. Dokshitzer, G. Marchesini and GZ, *Near-to-planar 3-jet events in and beyond QCD perturbation theory*, Phys. Lett. B **508** (2001) 269
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- ☛ A. Banfi, G. Marchesini, G. Smye and GZ, *Out-of-plane QCD radiation in DIS with high p(t) jets*, JHEP **0111** (2001) 066
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~ 1 observable per article

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Error prone business [✗ \sim 30%!]

Automated resummed predictions

Our goal: develop a computer code which resums final state observables at NLL accuracy in an automated way – as for fixed order calculations.

The user just

- ✗ fixes the Born process and the number of hard jets (legs)
- ✗ provides the definition of the observable in the form of a computer routine

➡ To achieve this one needs to understand the origin of all NLL terms in observable distributions in a general way.

The master formula

$$\Sigma(v) =_{NLL} \sum_{\text{sub.}} \int [d\Phi]_{\text{hard}} \Sigma_s(v) \cdot \mathcal{F}(R')$$

Banfi , Salam, GZ hep-ph/0304148

- ✓ Analytical resummation for the “easy” Σ_s : *pure LL and NLL terms*

$$\Sigma_s(v) = \prod_{\ell=1}^{n_{inc}} \underbrace{f_{\ell}(v^{\frac{2}{a+b_{\ell}}} \mu_F^2)}_{\text{pdfs}} \otimes \prod_{\ell=1}^N \underbrace{J_{\ell}(L)}_{\text{jet function}} \cdot \underbrace{S(T(L/a))}_{\text{soft}}$$

- ✚ soft and collinear emission \Rightarrow **jet function $J_{\ell}(L)$**
(all LL Sudakov suppression and some NLL terms)
- ✚ hard collinear splitting \Rightarrow **evolution of the pdfs**
- ✚ soft large angle \Rightarrow **QCD coherence and geometry dependence in S**

- ✓ the “difficult” \mathcal{F} is computed numerically but is **by construction a pure NLL function**



IDEA: Define a simpler observable with the same double logs but factorizes trivially

$$V(k_1, \dots, k_n) \Rightarrow V_s \equiv \max[V(k_1), \dots, V(k_n)]$$

➡ Simple factorization $\Theta(V_s - v) = \prod_i \Theta(V_i - v) \Rightarrow$ analytical resummation straightforward!



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Fix a Born event and emit a soft gluon k collinear to a given hard leg ℓ .

We parametrize

$$V(k) \simeq d_\ell \left(\frac{k_t}{Q} \right)^{a_\ell} e^{-b_\ell \eta} g_\ell(\phi)$$

$k_t \Rightarrow$ transverse momentum

$\eta \Rightarrow$ rapidity

$\phi \Rightarrow$ azimuth

✓ Σ_s known given the (automatically determined) quantities $a_\ell, b_\ell, d_\ell, g_\ell(\phi)$

Single emission properties



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To resum $V(k_1 \dots k_n)$ one needs to account for the observable specific mismatch between $V(k_1, \dots, k_n)$ and $V_s \Rightarrow$ multiple emission effects

Multiple emission effects

The function \mathcal{F} which encodes the information on how precisely the observable depends on multiple emissions, e. g.

- if $V(k_1, \dots, k_n) = \max\{V(k_1), \dots, V(k_n)\} \implies \mathcal{F} = 1$ [$y_3^{\text{Cam.}}$]
- if $V(k_1, \dots, k_n) = V(k_1) + \dots + V(k_n) \implies \mathcal{F} = \frac{e^{-\gamma_E R'}}{\Gamma(1 + R')}$ [τ]
- in general, compute \mathcal{F} via Monte Carlo event samples targeted to be observable

$$\mathcal{F} = \left\langle \exp \left\{ -R' \ln \frac{V(k_1, \dots, k_n)}{\max\{V(k_1), \dots, V(k_n)\}} \right\} \right\rangle$$

↔ Notation: $R' \equiv -dR/dL$ with $R(v)$ the LL Sudakov exponent $\Sigma_s(v) = e^{-R(v)}$

↪ R' and so \mathcal{F} are pure NLL functions!

Requirements on the observable

For the observable to be resummed automatically it should

- ✗ vanish in the Born limit and be positive defined
- ✗ behave as $V(k) \simeq d_\ell \left(\frac{k_t}{Q}\right)^{a_\ell} e^{-b_\ell \eta} g_\ell(\phi)$ for 1 SC gluon along leg ℓ
- ✗ be infrared and collinear safe
- ✗ be continuously global ($a_\ell = a \quad \forall$ hard legs ℓ)
- ✗ exponentiate (no JADE)

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- ✗ practically the limiting condition is the requirement of globalness
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- ✗ be **infrared and collinear safe**
- ✗ be **continuously global** ($a_\ell = a \quad \forall$ hard legs ℓ)
- ✗ **exponentiate** (no JADE)

While this might seem a long list

- ➡ practically the **limiting condition** is the requirement of **globalness**
(all other conditions are satisfied by all observables resummed so far)
- ➡ the essential feature of the program is the **ability to perform all checks automatically** and to resum the observable only when **correctness of the result is guaranteed at NLL**

Some observables have exponentiating double (and single) logs

$$P(v) = 1 - X \frac{\alpha_s C_F}{\pi} \ln^2 v + \frac{1}{2} X^2 \left(\frac{\alpha_s C_F}{\pi} \right)^2 \ln^4 v + \dots$$

Some observables have exponentiating double logs, others do not, e.g. Jade-algorithm jet rates:

$$P_{\text{Jade2-jet}}(y_{\text{cut}}) = 1 - \frac{\alpha_s C_F}{\pi} \ln^2 y_{\text{cut}} + \frac{1}{2} \cdot \frac{5}{6} \left(\frac{\alpha_s C_F}{\pi} \right)^2 \ln^4 y_{\text{cut}} + \dots$$

Brown and Strling, Phys.Lett.B 252 (1990)

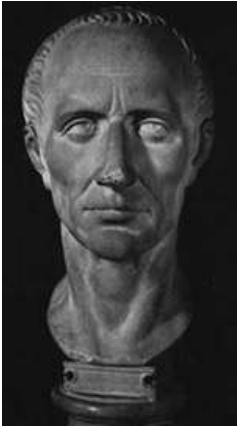
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Brown and Strling, Phys.Lett.B 252 (1990)

- ➡ No one jet knows how to resum Double Logs, let alone what matrix-element ingredients are needed to achieve NLL accuracy!

Any automated approach to NLL resummation has better be able to establish whether an observables exponentiates



Computer Automated Expert Semi-Analytical Resummation

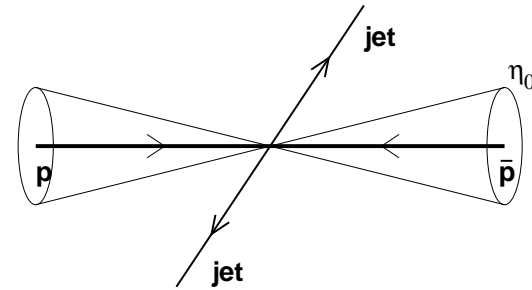


- currently limited to global observables
- tested against all known global, exponentiable event shapes
- results from an early version used by the LEP-QCD-WG for fits of α_s
- can be applied to
 - 2 & 3 jets in e^+e^-
 - [1+1] & [1+2] jets in DIS
 - Drell-Yan + 1 jet
 - hadron-hadron dijet events [\Leftarrow first resummations]

Observables in hadronic dijet production

Cut around the beam $|\eta| < \eta_0$

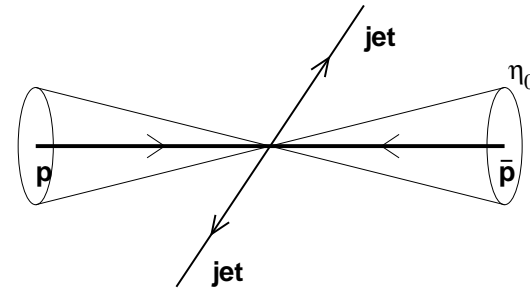
→ Problems with **globalness** 



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Directly global observables: $\eta_0 > 1$

X Transverse thrust

$$T_T = \frac{1}{E_T} \max_{\vec{n}_T} \sum_i |\vec{p}_{ti} \cdot \vec{n}_T|$$

X Thrust minor

$$T_m = \frac{1}{E_T} \sum_i |p_i^{out}|$$

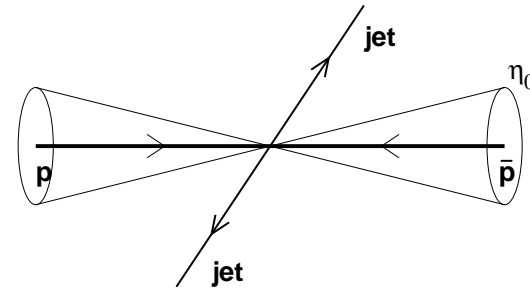
Predictions valid as long as

$$|\log v| < (a + b_\ell) |\eta_0|$$

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$$T_T = \frac{1}{E_T} \max_{\vec{n}_T} \sum_i |\vec{p}_{ti} \cdot \vec{n}_T|$$

✗ Thrust minor

$$T_m = \frac{1}{E_T} \sum_i |p_i^{out}|$$

Predictions valid as long as

$$|\log v| < (a + b_\ell)|\eta_0|$$

Indirectly global observables: $\eta_0 = \mathcal{O}(1)$

✗ Transverse thrust

$$T_T = \frac{1}{E_{T,\eta_0}} \left(\max_{\vec{n}_T} \sum_{|\eta_i| < \eta_0} |\vec{p}_{ti} \cdot \vec{n}_T| - \left| \sum_{|\eta_i| < \eta_0} \vec{p}_{ti} \right| \right)$$

✗ Thrust minor

$$T_m = \frac{1}{E_{T,\eta_0}} \left(\sum_{|\eta_i| < \eta_0} |p_i^{out}| + \left| \sum_{|\eta_i| < \eta_0} \vec{p}_{ti} \right| \right)$$

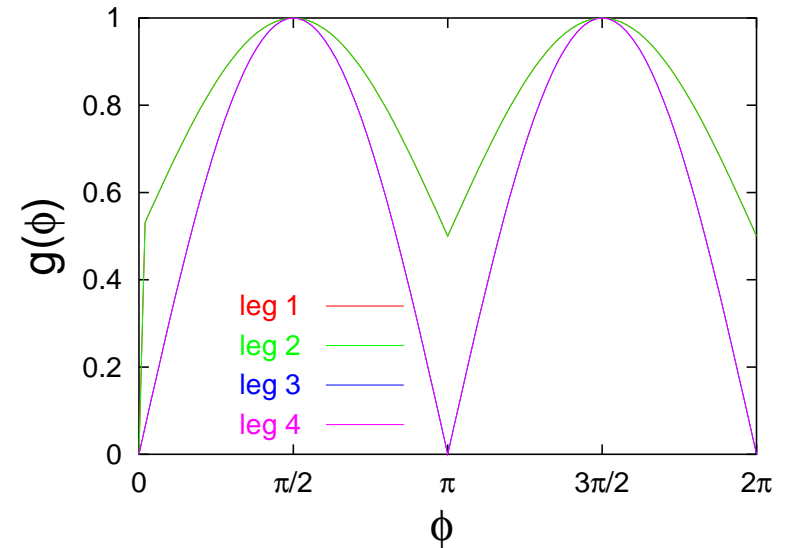
Predictions valid as usual,

but \mathcal{F} diverges at $R' = R'_c$

Sample output: the indirectly global thrust minor

✗ Tests on the observable

Test	result
check number of jets	T
all legs positive	T
global	T
continuously global	T
additive	F
exponentiate	T
eliminate subleading effects	T
opt. probe region exists	T



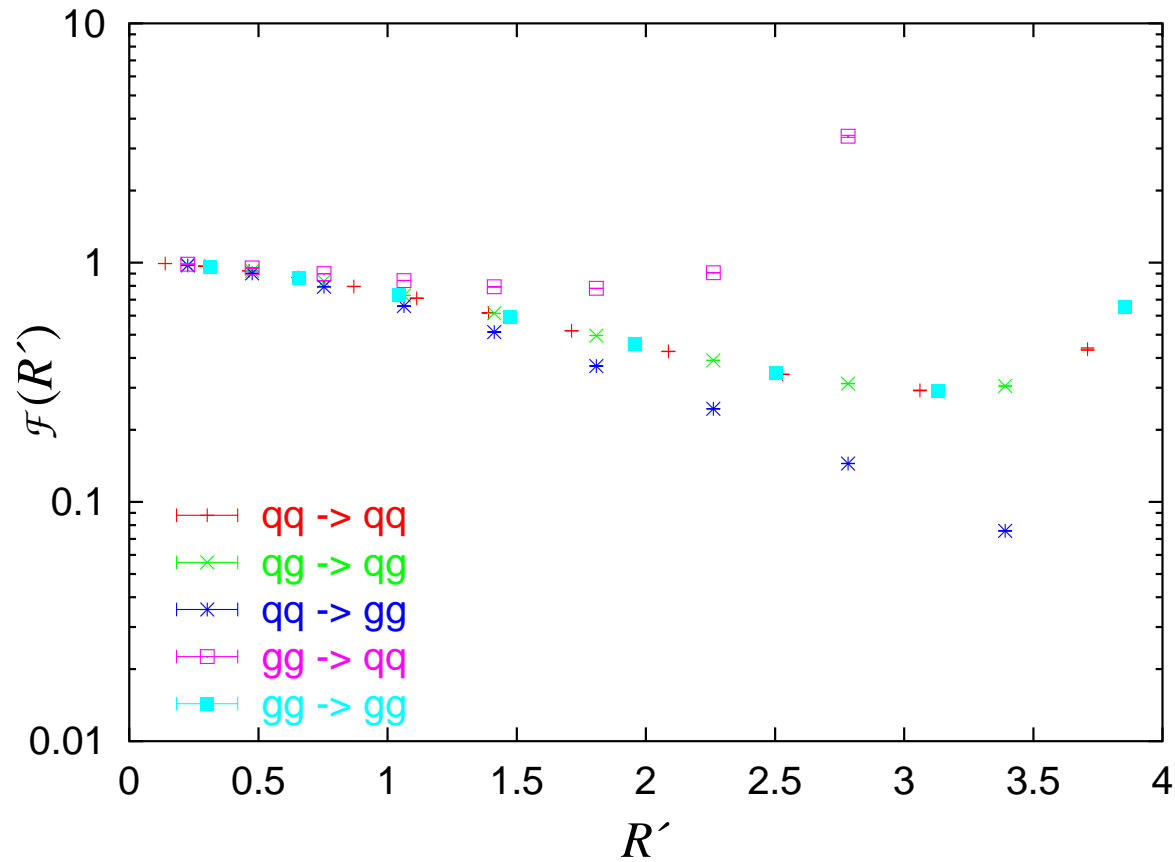
✗ Single emission properties

leg ℓ	a_ℓ	b_ℓ	$g_\ell(\phi)$	d_ℓ	$\langle \ln g_\ell(\phi) \rangle$
1	1	0	tabulated	2	-0.2201
2	1	0	tabulated	2	-0.2201
3	1	0	$\sin(\phi)$	2	$-\ln(2)$
4	1	0	$\sin(\phi)$	2	$-\ln(2)$

← Tables and plots generated automatically by CAESAR

$\mathcal{F}(R')$ for the indirectly global thrust minor

The multiple emission function $\mathcal{F}(R')$

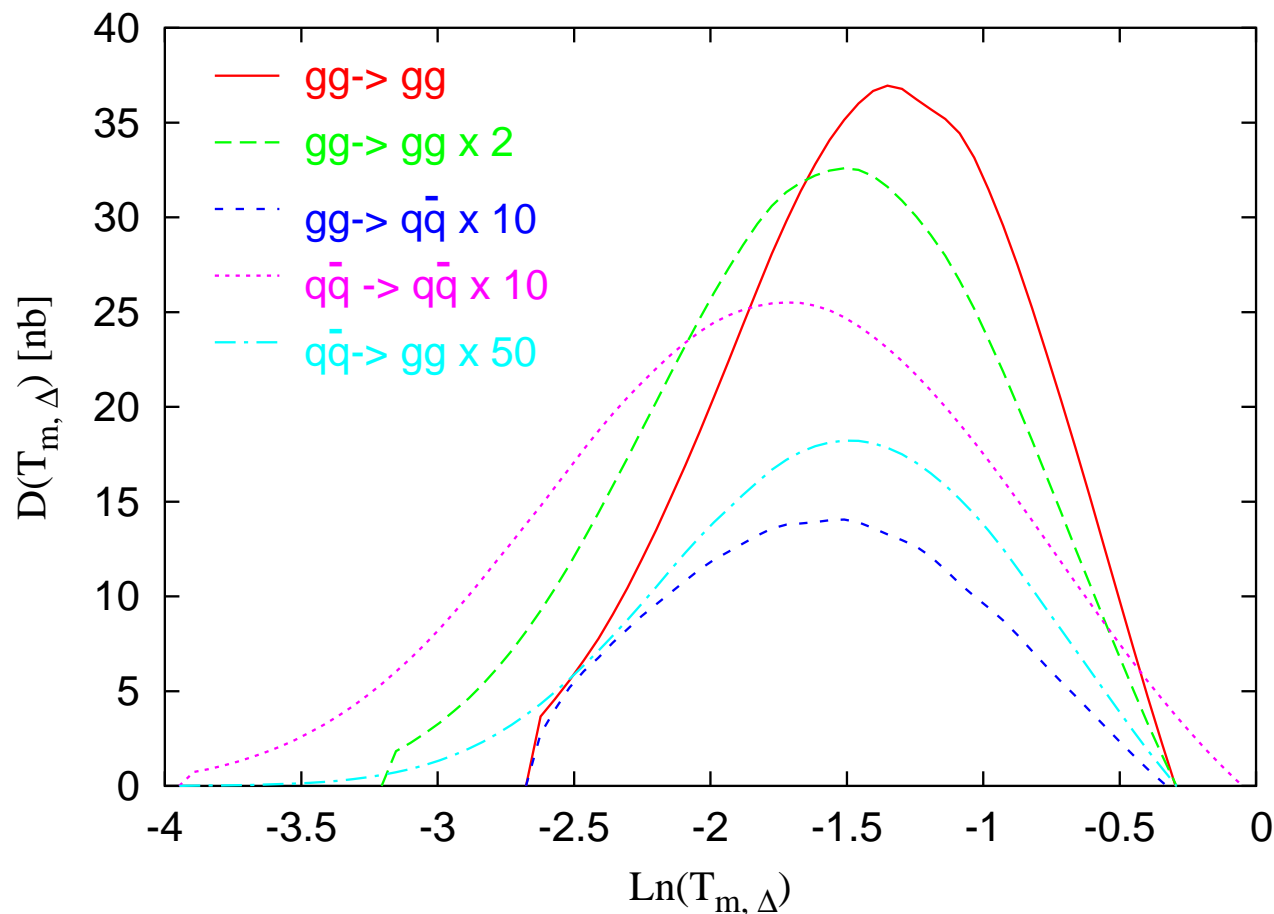


➡ Different result for different colour configurations

The indirectly global thrust minor

Dijets events at Tevatron run II regime

- ▶ run II regime $\sqrt{s} = 1.96 \text{ TeV}$
- ▶ cut on jet transverse energy $E_T > 50 \text{ GeV}$ and on rapidity $|\eta| < 1$



$$\alpha_s(M_Z) = 0.118$$

$$\mu_F = \mu_R = P_T$$

$$X_c = 1$$

PDFS: CTEQ6M

Physical/mathematical/technical content of CAESAR

✓ Born processes currently implemented

✌ e^+e^- -collisions: $e^+e^- \rightarrow 2 \text{ jets}$ $e^+e^- \rightarrow 3 \text{ jets}$

✌ DIS collision: $p e \rightarrow 2 \text{ jets}$ $p e \rightarrow 3 \text{ jets}$

✌ Drell Yan collision: $p_1 p_2 \rightarrow Z_0 + \text{jet}$

✌ Hadronic collisions: $p_1 p_2 \rightarrow 2 \text{ jets}$

$(p_i = q, \bar{q}, g)$

✓ Implementation of **exact analytical formulas** whenever possible

✓ Recoil in dipole method

Catani & Seymour, Nucl. Phys. B 485 (1997) 291

✓ Evolution of colour charge (soft radiation at large angle)

Kidonakis, Oderda & Sterman, Nucl. Phys. B 531 (1998) 365

✓ PDF evolution code

Dasgupta & Salam, Eur. Phys. J. C 24, 213 (2002)

✓ Extended **arbitrary precision arithmetic** package

Bailey, RNR Technical Report RNR-94-013

Conclusions & outlook

- ☞ In less inclusive regions fixed order calculations insufficient
 - ⇒ resummation of logarithmic enhanced terms mandatory
- ☞ the use of resummations limited by availability of analytical results

Main result: rigorous procedure to perform resummation semi-analytically

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Applications

- EX: first NLL predictions in **hh collisions (indirectly globalness)**
- TH: necessary and sufficient condition for **exponentiation**

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Work in progress

- release **CAESAR v1.0**

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To-do list

- automated matching of **NLL** with **NLO(JET++)**

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Wish-list

- extension non-global observables and inclusion of mass effects