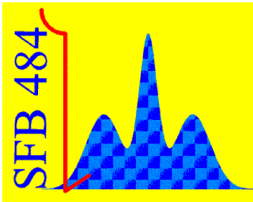


Center for
Electronic Correlations and Magnetism
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Towards a global phase diagram
for perovskite manganites

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and P. Majumdar (Allahabad)

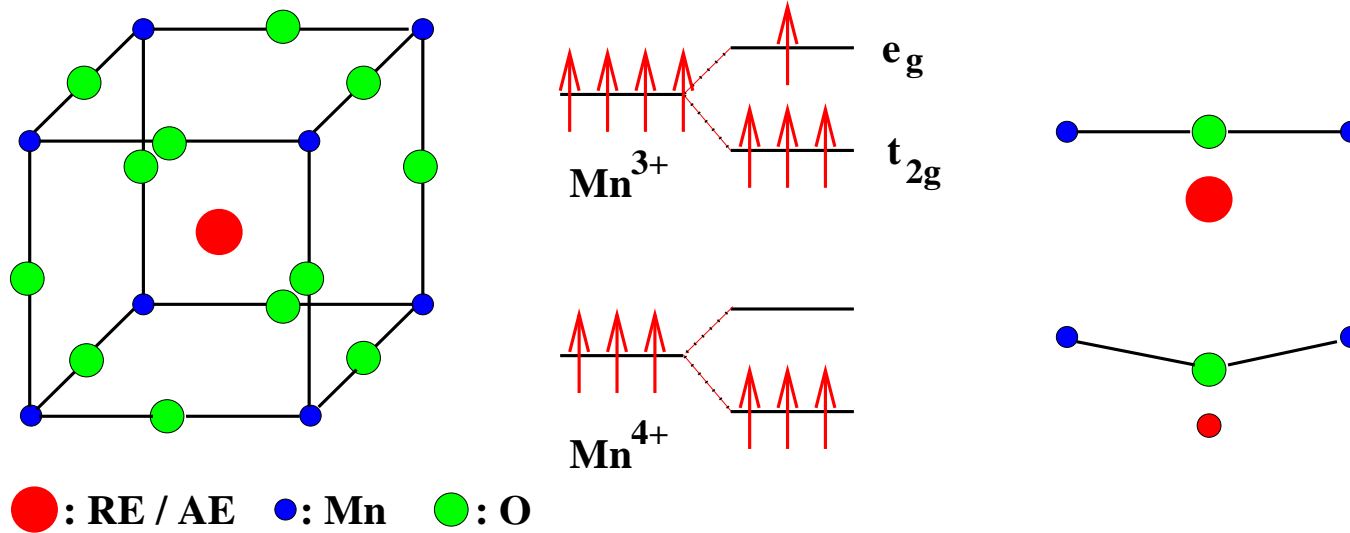


Sonderforschungsbereich 484
Kooperative Phänomene im Festkörper: Metall-Isolator-
Übergänge und Ordnung mikroskopischer Freiheitsgrade

Outline:

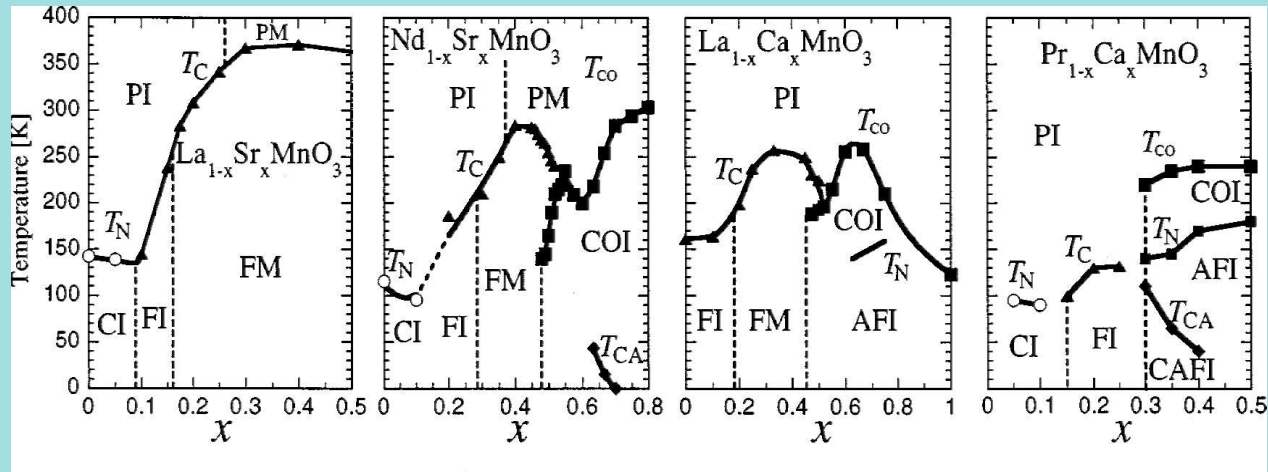
- perovskite manganites:
crystal structure, electronic configurations
- recent experiments:
controlling disorder and bandwidth
- model for manganites:
2 bands, Jahn-Teller and Hund's rule coupling
- traveling cluster approximation:
exact diagonalisation and Monte Carlo
- results:
phase diagrams and transport
- summary

Perovskite structure & electronic configurations



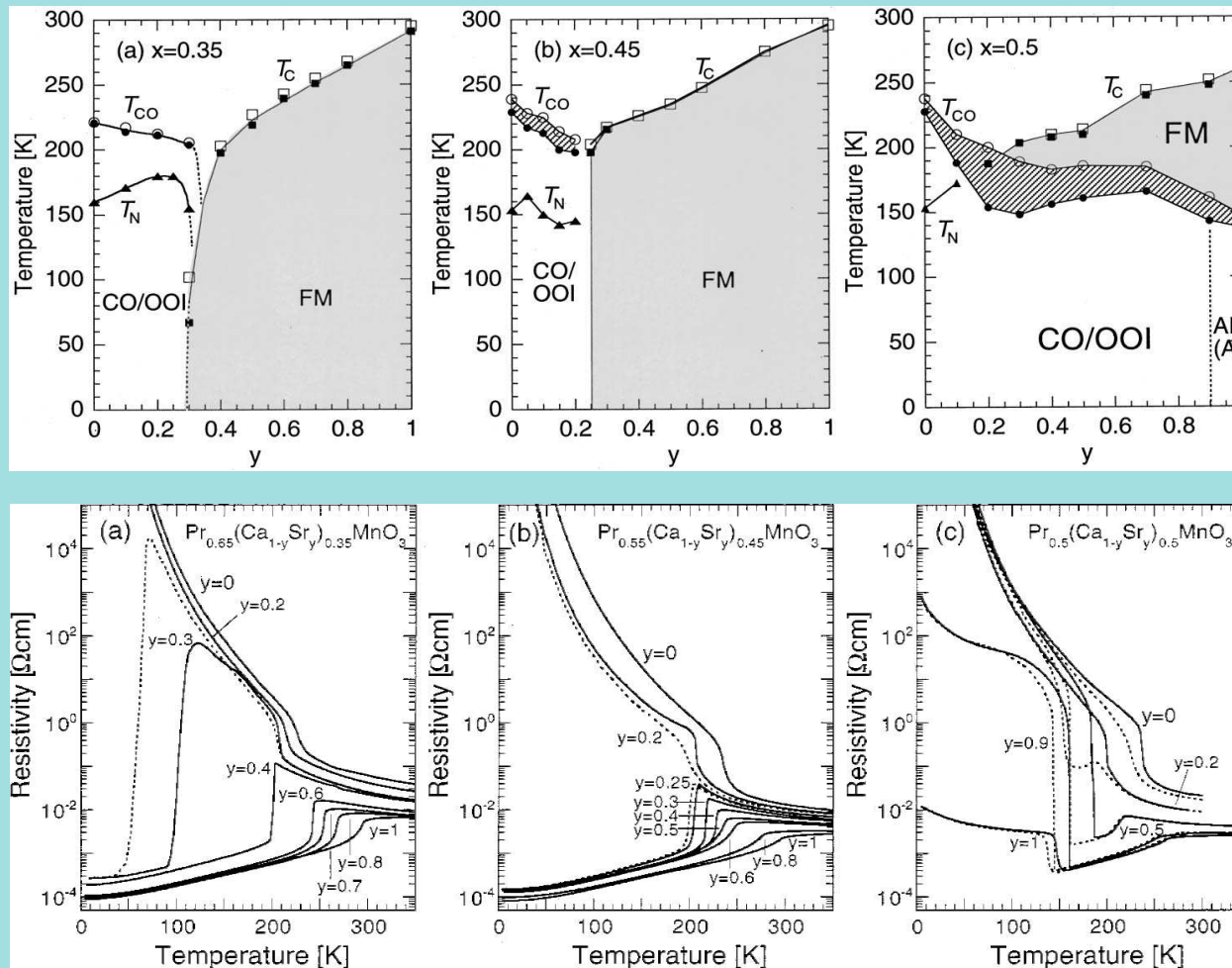
- parent compounds (e.g. LaMnO_3)
A-type AFM, Mott insulator, distorted octahedra
- doped manganites (e.g. $\text{La}_{1-x}\text{Ca}_x\text{MnO}_3$)
 - RE^{3+} replaced by AE^{2+} \longrightarrow hole doping
 - CaMnO_3 : G-type AFM, band insulator

T-x Phase Diagrams. Imada *et al.* (1998)



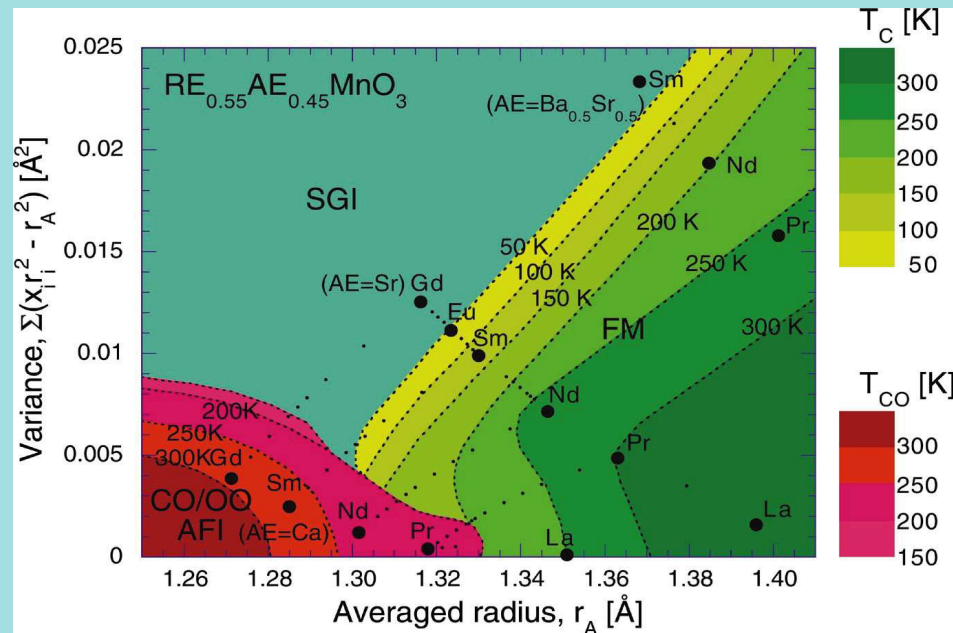
- Large bandwidth: FM-Metal with large T_c
- Intermediate bandwidth: reduction in T_c , metal insulator transition, colossal magnetoresistance(CMR).
- Small bandwidth: Only insulating phases.

Phase Diagrams for $Pr_{1-x}(Ca_{1-y}Sr_y)_xMnO_3$.



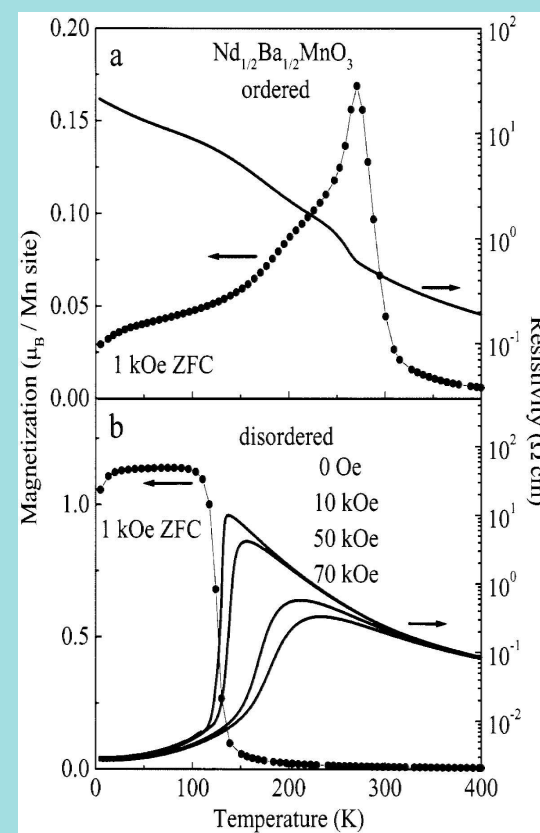
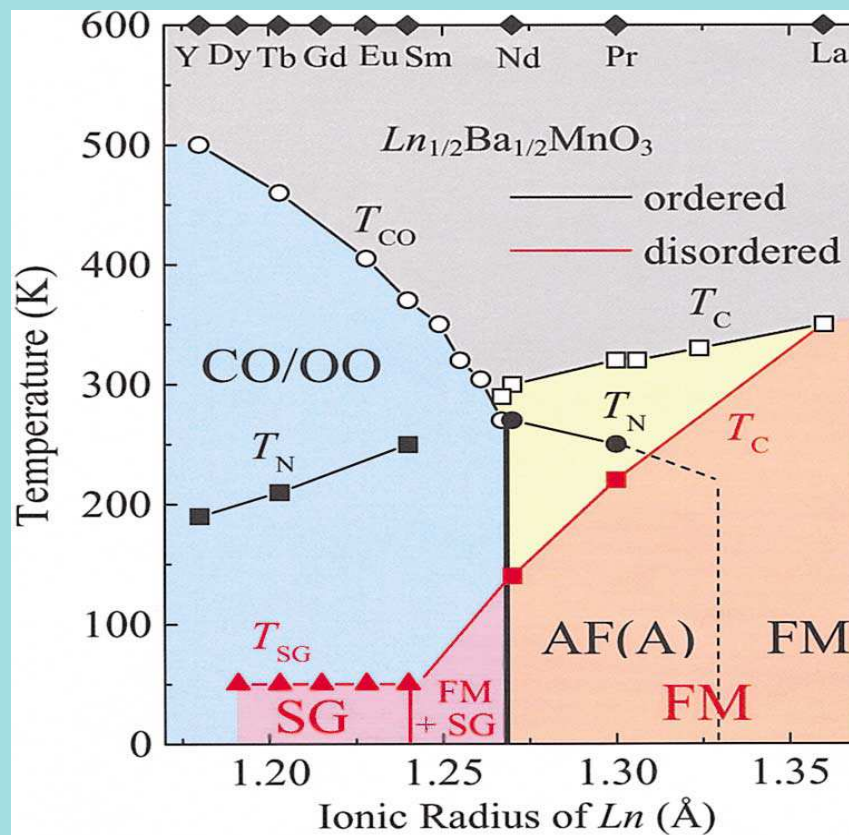
Tomioka *et al.* (2002)

Global Phase Diagram. Tomioka *et al.* (2004)

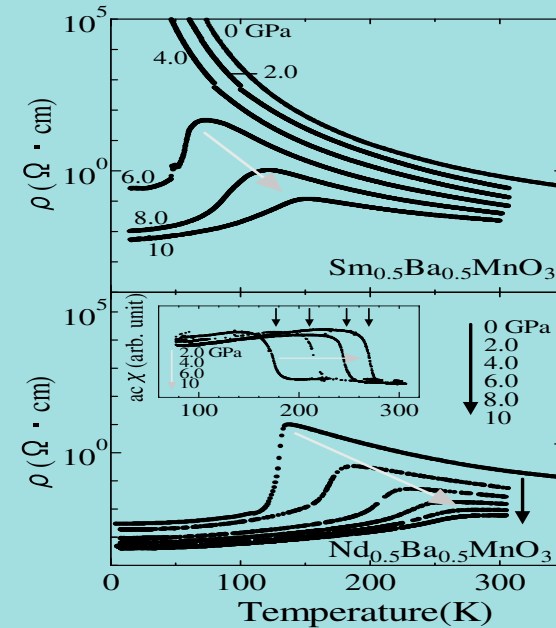
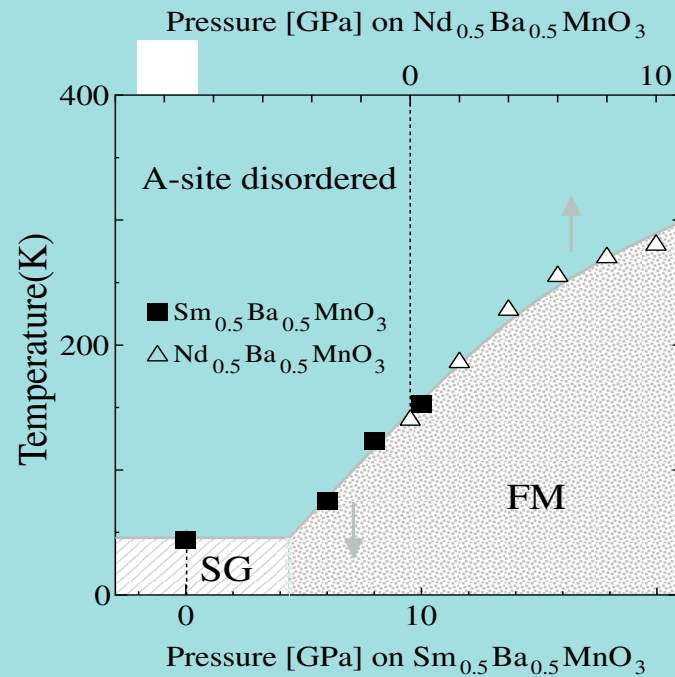


- universal parameters: Average cation size (r_A) and variance in cation size (σ^2).
- more than 50 different manganites follow a universal trend in terms of r_A and σ^2 .

Order vs. Disorder. Akahoshi *et al.* (2003)



Bandwidth controlled by external pressure.



Takeshita *et al.* (2004)

- *SBMO* at finite pressure \sim *NBMO* at zero pressure.
- *NBMO* at finite pressure \sim *PBMO* at zero pressure.

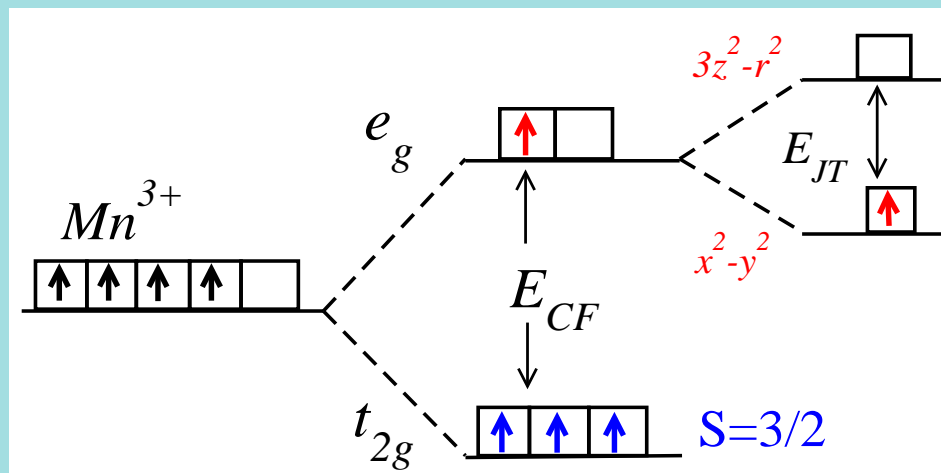
Summary of selected experiments

- special charge and orbital ordering phenomena at $x=1/2$
- insulating phases close to $x=0$ and $x=1$
- interesting doping regime for CMR: $0.25 < x < 0.45$
- 2 overall parameters: bandwidth and disorder,
independent of the specific chemical composition

Focus of present work:

understanding global trends away from commensurate fillings

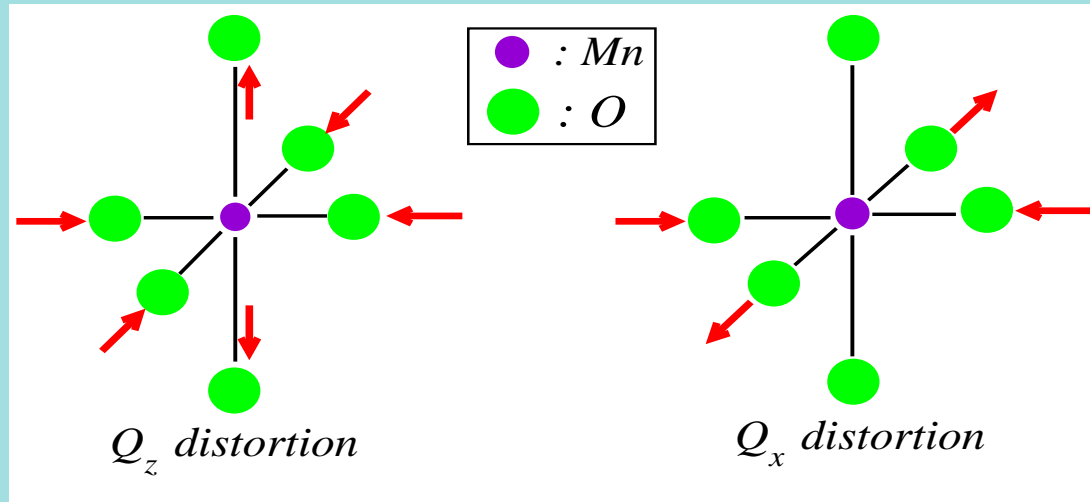
Model for Manganites



Main ingredients:

- Kinetic energy; complicated hopping structure ($t_{ij}^{\alpha\beta}$).
- Hund's rule coupling of t_{2g} spins ($S=3/2$) to e_g spin.
- Octahedra distortions; Jahn-Teller effect for e_g orbitals.
- Hubbard interactions: inter & intra orbital.

Jahn-Teller distortions



- Origin of JT effect: lifting of orbital degeneracy due to electron lattice coupling.

Simplifications:

- Mn ion is always at the centre of the octahedron.
- O displacements are along the $Mn - O - Mn$ bond.

The Hamiltonian

$$H = H_{kin} + H_S + H_{el-latt} + H_{dis}$$

$$H_{kin} = \sum_{\langle ij \rangle \sigma}^{\alpha\beta} \{t^{\alpha\beta} c_{i\alpha\sigma}^\dagger c_{j\beta\sigma} + h.c.\} \quad ; \quad H_{dis} = \sum_i \epsilon_i n_i$$

$$H_S = -J_H \sum_i \mathbf{S}_i \cdot \hat{\sigma}_i + J_s \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

$$H_{el-latt} = \lambda \sum_i \{Q_i^x \tau_i^x + Q_i^z \tau_i^z\} + \frac{K}{2} \sum_i \{(Q_i^x)^2 + (Q_i^z)^2\}$$

$$\tau_i^x = \sum_\sigma (c_{i1\sigma}^\dagger c_{i2\sigma} + c_{i2\sigma}^\dagger c_{i1\sigma}) \quad ; \quad \tau_i^z = \sum_\sigma (c_{i1\sigma}^\dagger c_{i1\sigma} - c_{i2\sigma}^\dagger c_{i2\sigma})$$

$\epsilon_i = \pm\Delta$; **binary disorder.**

Simplifying assumptions

- Treat $\{Q^x, Q^z\}$ in the adiabatic limit.
- Treat $\{S\}$ as classical variables.
- Justification: Large t_{2g} spin, large O mass.
- For a given set of $\{Q^x, Q^z\}$ and $\{S\}$, the Hamiltonian is bilinear in electronic operators.

Still a difficult problem to solve !

$\{Q^x, Q^z\}$ and $\{S\}$ have to be determined self consistently.

Parameter space $\lambda/t, \Delta/t, J_H/t, J_s/t, n$ and T/t .

Q: How to determine the classical configurations?

Double Exchange Limit

$J_H/t \rightarrow \infty$; e_g spins are slaved along the t_{2g} spin directions.



$$H_{kin} = \sum_{\langle ij \rangle}^{\alpha\beta} \{t_{ij}^{\alpha\beta} \gamma_{i\alpha}^\dagger \gamma_{j\beta} + h.c.\}$$

$$t_{ij}^{\alpha\beta} = t^{\alpha\beta} (\cos(\theta_i/2) \cos(\theta_j/2) + \sin(\theta_i/2) \sin(\theta_j/2) e^{i(\phi_i - \phi_j)})$$

θ_i, ϕ_i : polar and azimuthal angles for spin S_i

• hopping matrix specific to the perovskite manganites

$$t^{11} = t, \quad t^{12} = t^{21} = \mp \frac{1}{\sqrt{3}} t, \quad t^{22} = \frac{1}{3} t.$$

Combining ED with Monte Carlo

Formally integrating out the fermions

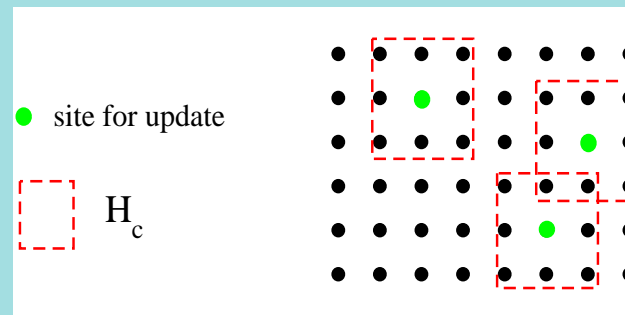
$$Z = \int \mathcal{D}S \mathcal{D}Q \operatorname{Tr} e^{-\beta H(\{Q, \mathbf{S}\})} \equiv \int \mathcal{D}S \mathcal{D}Q e^{-\beta H_{eff}(\{Q, \mathbf{S}\})}$$

$$H_{eff}(\{Q, \mathbf{S}\}) = -\frac{1}{\beta} \log \operatorname{Tr} e^{-\beta H(\{Q, \mathbf{S}\})}$$

- In practice one needs to know the eigenvalues of $H(\{Q, \mathbf{S}\})$ in order to evaluate the above trace.
- Straightforward method: combination of exact diagonalisation and classical Monte-Carlo (EDMC).
- Too costly in terms of cpu time. $\tau_N \sim N^4$

Travelling Cluster Approximation(TCA)

- Construct a Hamiltonian on a cluster of sites around the site to be updated during Monte-Carlo.
- Estimate ΔE for old and new configurations using the eigenvalues of the cluster Hamiltonian.



- $\tau_N \sim NN_c^3$. Achievable size $N \sim 1000$ for $N_c \sim 100$
- To compute electronic properties we use exact eigenfunctions of the full Hamiltonian in the TCA generated classical configurations.
- Detailed comparisons, [cond-mat/0406082](https://arxiv.org/abs/cond-mat/0406082)

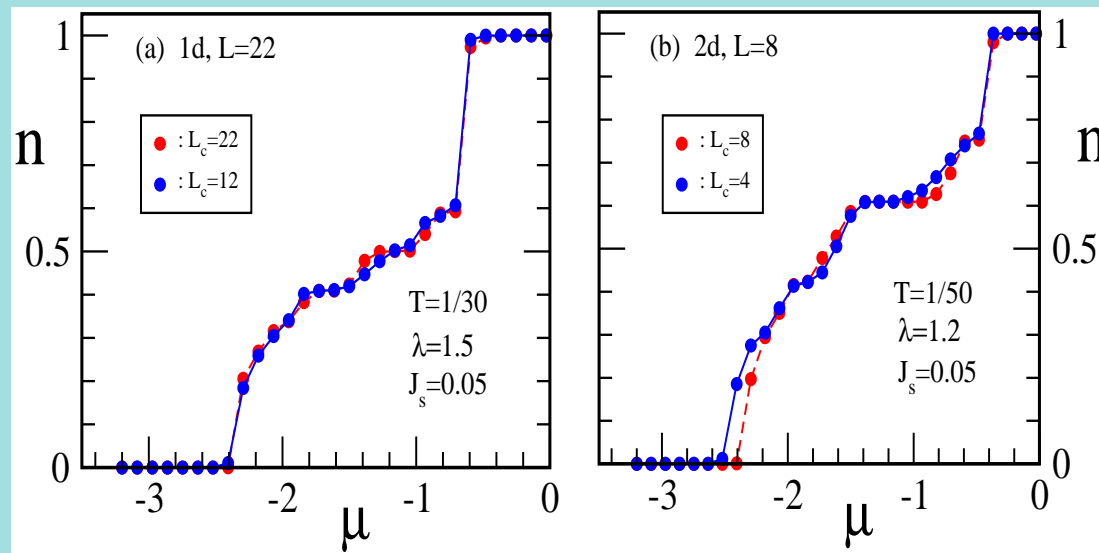
Calculated Quantities

1. spin structure factor: $S_{\vec{q}} = N^{-2} \sum_{ij} \mathbf{S}_i \cdot \mathbf{S}_j e^{i\vec{q} \cdot (\vec{r}_i - \vec{r}_j)}$
2. orbital structure factor: $\tau_{\vec{q}} = N^{-2} \sum_{ij} \tau_i \cdot \tau_j e^{i\vec{q} \cdot (\vec{r}_i - \vec{r}_j)}$
3. electronic density of states: $N(\omega) = \frac{1}{N} \sum_k \delta(\omega - \epsilon_k)$
4. local electronic density: $\langle n_i \rangle = \langle \gamma_{i1}^\dagger \gamma_{i1} + \gamma_{i2}^\dagger \gamma_{i2} \rangle$
5. conductivity using Kubo-Greenwood formula:

$$\sigma(\omega) = \frac{\pi e^2}{hN} \sum_{ab} |\langle \psi_a | \hat{j} | \psi_b \rangle|^2 \frac{(f_b - f_a)}{(\epsilon_a - \epsilon_b)} \delta(\omega - (\epsilon_a - \epsilon_b))$$

$$\hat{j} = i \sum_{\langle ij \rangle}^{\alpha\beta} (t_{ij}^{\alpha\beta} \gamma_i^\dagger \gamma_j - t_{ji}^{\alpha\beta} \gamma_j^\dagger \gamma_i)$$

Comparison with exact results

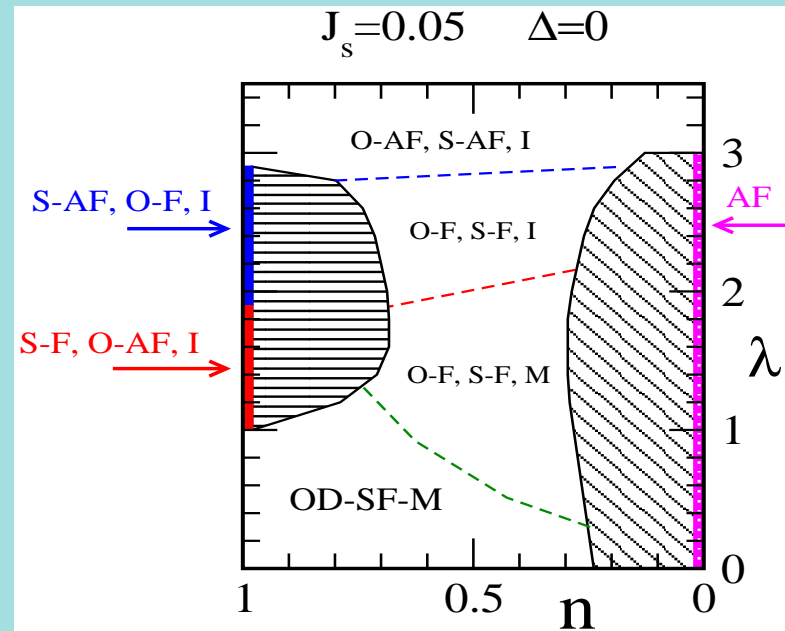


red dots: ED results from Yunoki *et al.* PRL(1998).

blue dots: results using the cluster method.

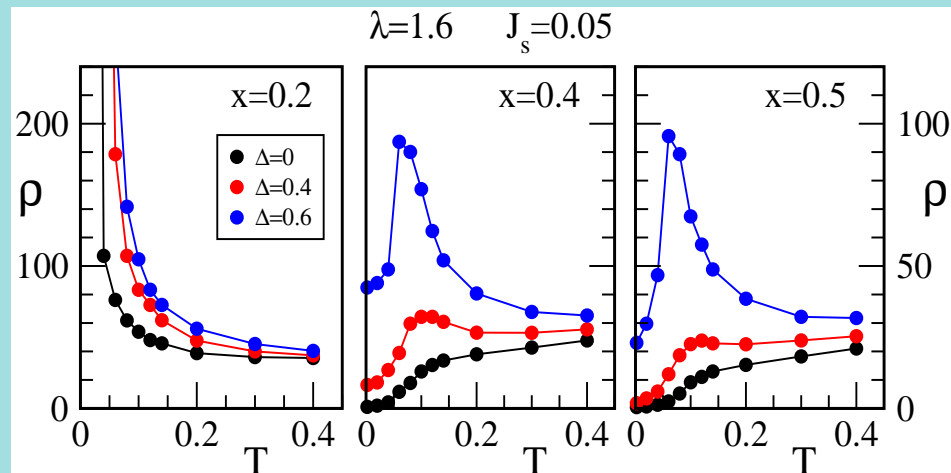
- cpu constraints in 3d, minimum 4^3 cluster required.

Zero Temperature Phase Diagram in 2D



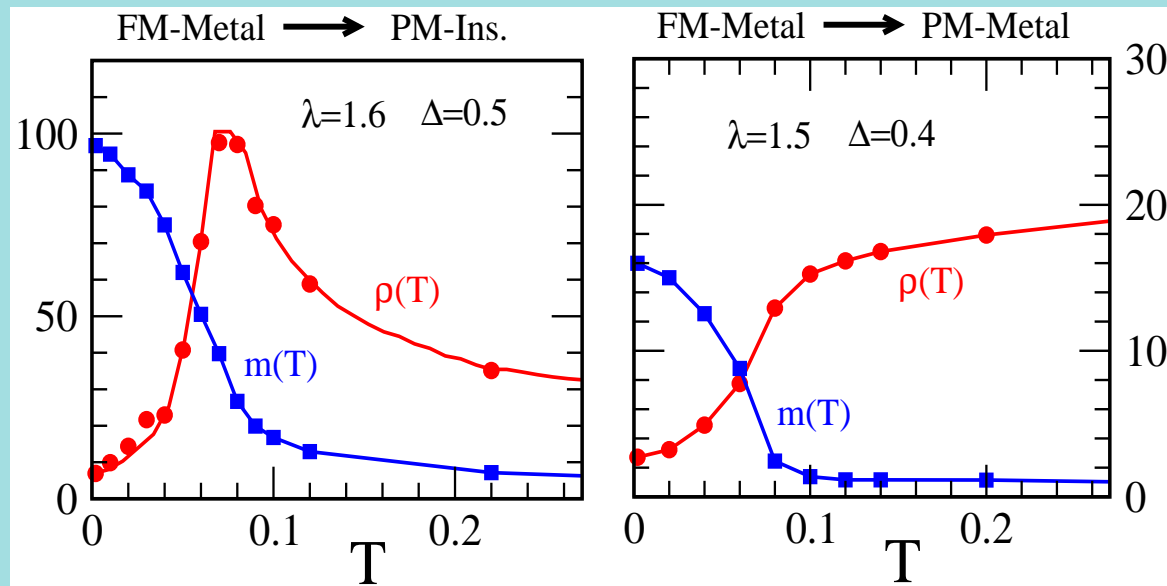
- insulators at $n = 0$ and $n = 1$. Phase separations.
- FM-I phase with reduced magnetisation.
- larger J_s induces CO-OO-I phase at $n = 0.5$.

Resistivity: variations in hole doping (x).



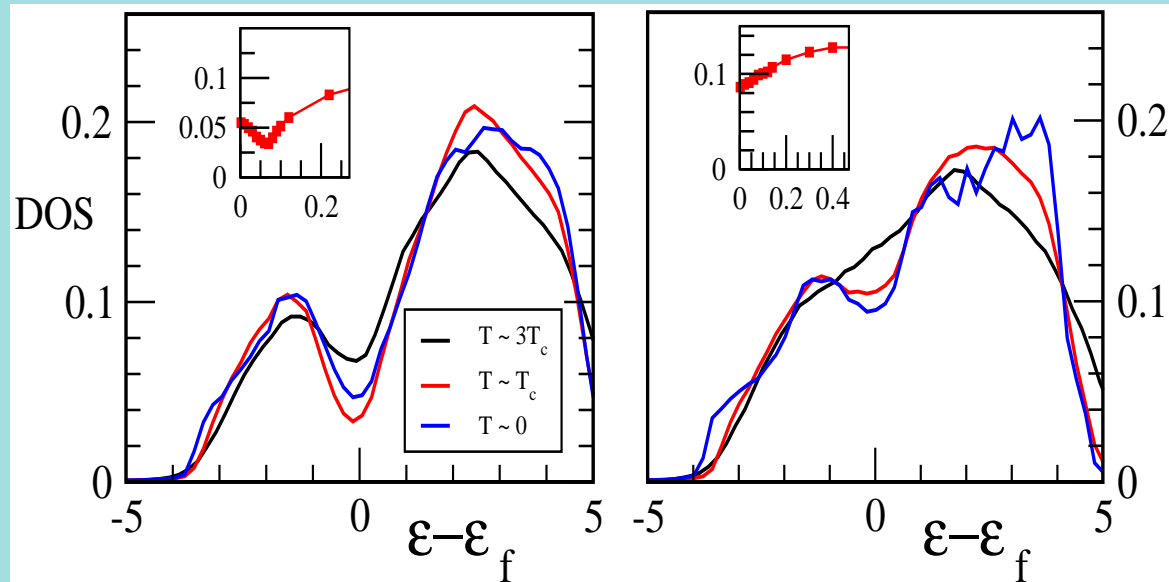
- Low doping \Rightarrow insulating phase.
- Disorder strongly affects the region near T_c . Inducing a metal-insulator transition.

Resistivity and Magnetisation



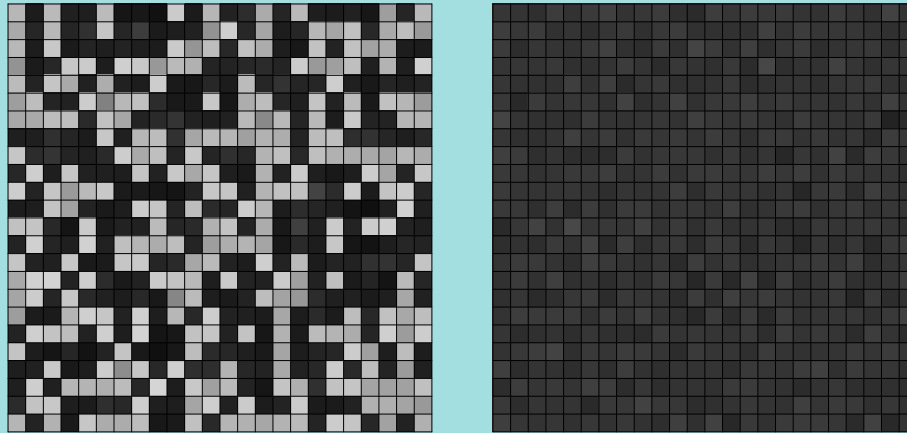
- correlation between FM-PM and M-I transitions
- weak λ, Δ has typical DE like behavior.
- small disorder induces polaron trapping, drastically changing $\rho(T)$.

Density of states



- pseudogap in the DOS.
- DOS at E_f has different behavior in the two cases.
- does provide hints for different resistivity behaviors.

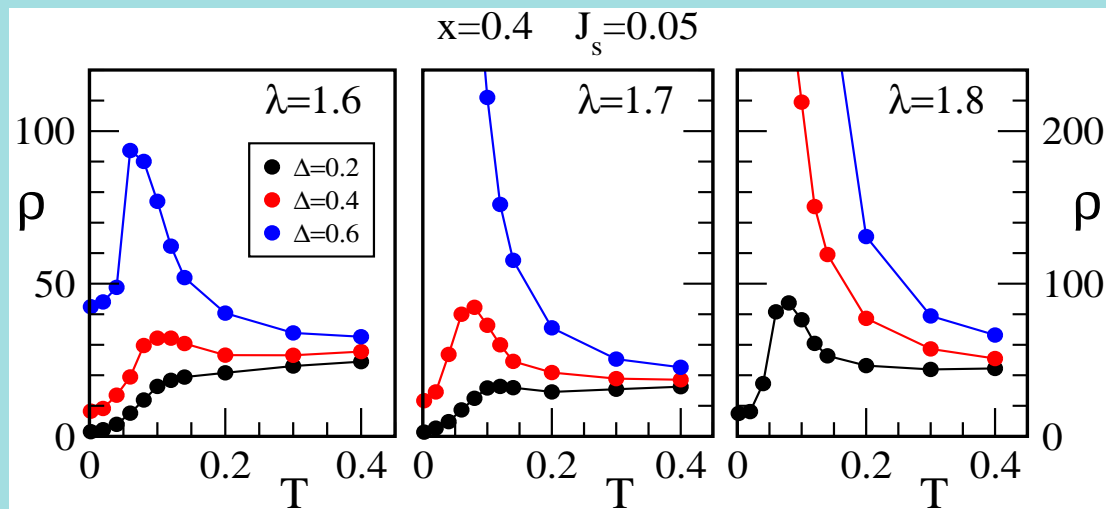
Real Space Density Profile



Understanding the MIT: Two type of states

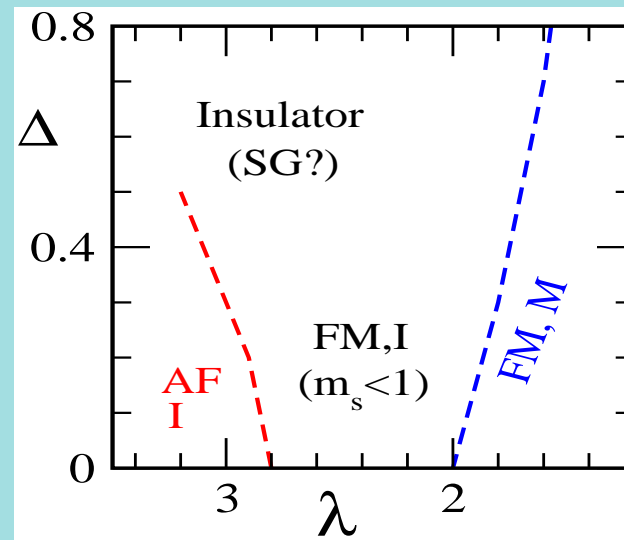
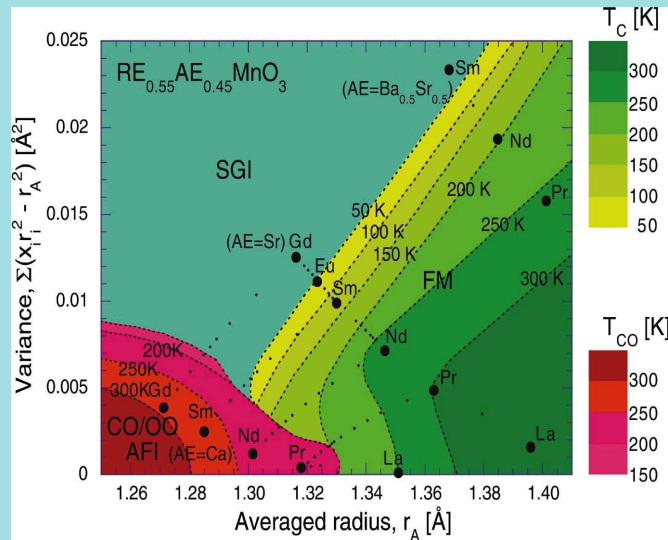
- Mixture of polaronic and extended states at $T = 0$.
- Rise in $\rho(T)$, scattering of mobile electrons.
- Additional spin disorder manages to trap extended states \Rightarrow insulating paramagnetic phase.

Resistivity: variations in bandwidth and disorder



- change in $t \Rightarrow$ change in λ/t , J_s/t and Δ/t .
- we explore the effect of changing λ/t and Δ/t .
- 3 transport regimes on increasing λ or Δ or both.

Summary



- Similar trends as seen in the experiments.
- spin glass phase in our model, yet to be confirmed.
- More realistic modeling: disorder in hoppings, changes in J_s/t .
- Work still in progress...