Spin Liquids in Triangular Lattice Organic Compounds: a Nodal Fractionalized State?

Tarun Grover, Nandini Trivedi, T. Senthil, Patrick Lee

Tarun Grover, Nandini Trivedi, T. Senthil, Patrick Lee Spin Liquids in Triangular Lattice Organic Compounds: a Noda

Basics: Phases of Matter

- Most Phases of Matter can be classified by identifying an 'order parameter' O.
- Order parameter \Rightarrow Classical description.

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'Truly' Quantum Phases of Matter

- Zero Kelvin, no symmetry breaking, no classical description.
- Quantum hall phases, Fermi liquids, Insulators with one-electron per unit cell and no symmetry breaking.
- Spin-liquids ⇒ Oshikawa-Hastings argument gaurantees an interesting outcome e.g. emergence of topological order, artificial photons ...

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Outline

κ-ET Organic Superconductors and Insulators Spin-liquids on the Triangular Lattice Consequences of d-wave Spin-Liquid State Summary and Questions

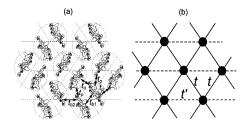
- 1 κ -ET Organic Superconductors and Insulators
 - Basics
 - Phenomenology of insulating region
 - Phenomenology of superconducting region
 - Summary
- 2 Spin-liquids on the Triangular Lattice
 - Ring-exchange Hamiltonian
 - Variational wave-functions
 - Result
- Onsequences of d-wave Spin-Liquid State
- 4 Summary and Questions

Basics

Phenomenology of insulating region Phenomenology of superconducting region Summary

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Structure



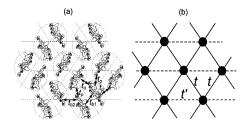
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- A moderate pressure can change t, t'!

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Material	Ground State	t'/t
$\kappa - (ET)_2 Cu_2 (CN)_3$	Unordered Mott insulator	1.06
$\kappa - (ET)_2 Cu[N(CN)_2]CI$	Ordered Mott insulator	0.75

• $\kappa - (ET)_2 Cu[N(CN)_2]Cl$ orders antiferromagnetically at 27 K.

• $\kappa - (ET)_2 Cu_2(CN)_3$ doesn't till lowest observed temperatures ($\approx 20 \text{ mK}$)!

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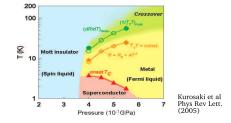
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Basics

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Phase diagram of κCN



• Pressure \Rightarrow Mott Insulator \rightarrow Superconductor \rightarrow Fermi liquid.

Superconductivity proximate to Mott insulator (coincidence?).

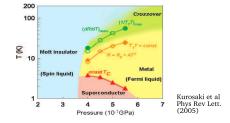
Mott insulator suspected to be a spin-liquid.

Basics

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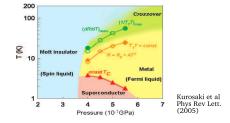
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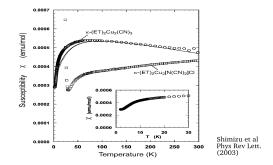
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Insulator: Spin Susceptibility



No signature of phase transition till 1.9 K (J ≈ 250 K!).
χ(T → 0) is finite.

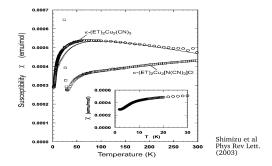
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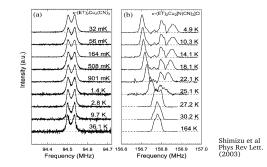
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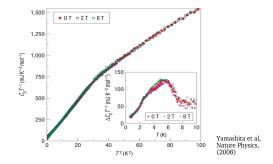
Insulator: NMR spectra



• No appreciable shift \Rightarrow No local magnetization.

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Insulator: Specific Heat



• *T*-linear specific heat at low temperatures!

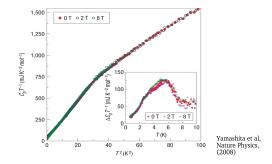
• Insensitive to external magnetic field.

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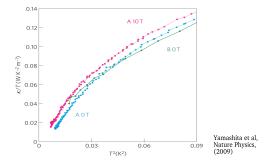
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Insulator: Thermal Conductivity



κ/T(T→0) extrapolates to different values for different T.
For T > 1K, κ/T(T→0) is non-zero.

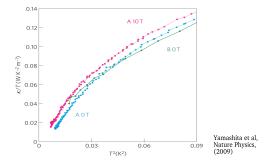
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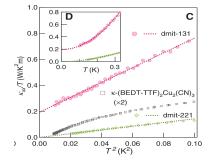


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Insulator: Thermal Conductivity for $EtMe_3Sb[Pd(dmit)_2]_2$



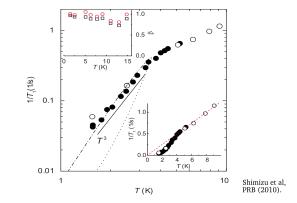
Yamashita et al, Science 2010.

• $\kappa/T(T \to 0)$ extrapolates to non-zero value as $T \to 0$.

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Superconductor: NMR Relaxation



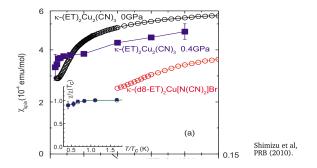
• $1/T_1T \rightarrow T^2$ for $T \lesssim 3.5$ K.

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Superconductor: Suscptibility



Almost no change in suscptibility across the Mott transition!

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Summary

Insulating Phase

- No magnetic ordering down to 32 mK ($\approx 10^{-4}J$) \Rightarrow Spin-liquid at T = 0?
- Specific heat $C_P \sim T$ at low $T \Rightarrow$ Gapless excitations?

- $1/T_1T \sim T^2$ at low $T \Rightarrow$ Nodes in SC gap?
- $\chi(T \rightarrow 0)$ is finite \Rightarrow Gapless excitations?

Basics Phenomenology of insulating region Phenomenology of superconducting region Summary

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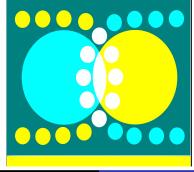
Summary

Topological and Quantum Order

Long Lost English Edition

Course of Theoretical Physics Volume 11

L.D. Landau, E.M. Lifshitz and X.G. Wen



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Ring-exchange Hamiltonian /ariational wave-functions Result

Gapless spin liquids and fractional particles

• κCN as gapless fractionalized Mott insulator.

- Basic idea: electron creation operator $c_{\sigma}^{\dagger} = f_{\sigma}^{\dagger} b$.
- Spinon f_{σ} carries spin while chargeon b carries charge.
- Emergent gauge fields.

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Ring-exchange Hamiltonian Variational wave-functions Result

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Hamiltonian for κ -ET insulators?

- Charge gap <u>BUT</u> significant charge fluctuations (proximity to metal-insulator transition).
- \Rightarrow Multiple exchange spin-model.

$$H = 2J_2 \sum_{\langle rr' \rangle} \vec{S}_{r'} \cdot \vec{S}_{r'} + J_4 \sum_{\Box} (P_{1234} + h.c) + \dots \quad (1)$$

Motrunich, Phys. Rev. B, 2005

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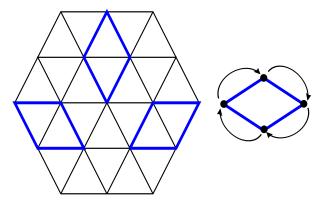
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Ring Exchange



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Ring-exchange Hamiltonian Variational wave-functions Result

• Construct variational states $|\Psi\rangle$ motivated by κCN .

- Minimize $E = \langle \Psi | H | \Psi \rangle$ for candidate $| \Psi \rangle$'s.
- Propose $|\Psi
 angle$ with the minimum *E* as the ground state.

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Ring-exchange Hamiltonian Variational wave-functions Result

Wave-function for a gapless spin-liquid

- No symmetry breaking and gapless spinful excitations?
- Fermi liquid almost does it (!).
- Insulator \rightarrow Project out the charges.

$$PFL\rangle = \prod_{i} (1 - n_{i\uparrow} n_{i\downarrow}) |FL\rangle$$
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$$\equiv \text{ Fermi surface of spinons}$$
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Effective theory of $|PFL\rangle$

• Spinons coupled to U(1) gauge field.

$$S = \sum_{\langle ij \rangle} \left(f_i^{\dagger} f_j e^{i a_{ij}} + h.c. \right) + (\nabla \times a)^2$$
(4)

• Renormalized propagator $D(\omega, \vec{k})$ for photon:

$$D(\omega, \vec{k}) = \omega^2 + k^2 + \frac{|\omega|}{k}$$
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Ring-exchange Hamiltonian Variational wave-functions Result

Problems with $|PFL\rangle$ as the ground state of κCN ?

- U(1) gauge fluctuations $\Rightarrow C \sim T^{2/3}$ at low T.
- Superconductivity near the Mott transition?
- Many competitive PBCS states.

Motrunich, Phys. Rev. B, 2005 S.-S. Lee and P. A. Lee, Phys. Rev. Lett., 2005.

Consequences of discrete All Antices All A

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Ring-exchange Hamiltonian Variational wave-functions Result

More interesting spin-liquids

- Project charges out from the BCS wave-function.
- Instability of the spinon Fermi surface!

$$|PBCS\rangle = \prod_{i} (1 - n_{i\uparrow} n_{i\downarrow}) |BCS\rangle$$
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Ring-exchange Hamiltonian Variational wave-functions Result

Projected BCS states and κCN

• Explanation for superconductivity.



Ring-exchange Hamiltonian Variational wave-functions Result

Projected BCS states and κCN

- Explanation for superconductivity.
- Cooper pairs $\xrightarrow{\text{Projection}}$ Resonating dimers ('RVB').
- Or,

Superconductor $\xrightarrow{\text{Projection}}$ Spin-liquid.

Ring-exchange Hamiltonian Variational wave-functions Result

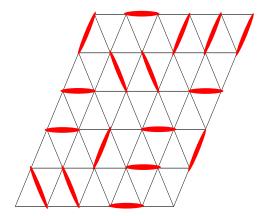
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Spin-liquid in pictures



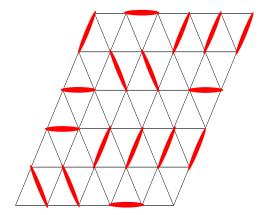
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Ring-exchange Hamiltonian Variational wave-functions Result

Spin-liquid in pictures

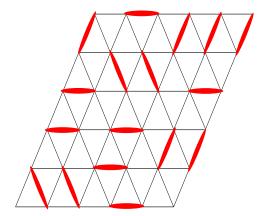


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Ring-exchange Hamiltonian Variational wave-functions Result

Spin-liquid in pictures

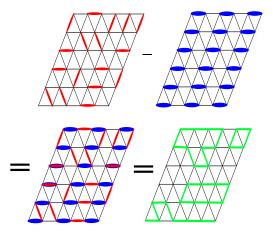


Tarun Grover, Nandini Trivedi, T. Senthil, Patrick Lee Spin Liquids in Triangular Lattice Organic Compounds: a Noda

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Spin Liquids in Triangular Lattice Organic Compounds: a Nodal

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Projected BCS states on triangular lattice

- Singlet $s, d_{xy} + id_{x^2-y^2}$ and triplet $f_{x^3-3xy^2}$ wave states fit triangular lattice.
- Anisotropic triangular lattice \rightarrow projected $d_{x^2-y^2}$ state has lower energy.
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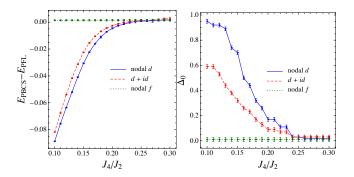
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Variational 'Answer' for Isotropic Triangular

Projected nodal $d_{x^2-y^2}$ has the lowest energy!



Ring-exchange Hamiltonian Variational wave-functions Result

Why projected nodal *d*-wave is the best state?

- Ring-exchange (J_4) tends to delocalize the spinons.
- Heisenberg (J_2) drives BCS instability.
- For J₄/J₂ ≫ 1, mean-field gives projected Fermi liquid. (Motrunich 2005)
- \Rightarrow As J_4/J_2 decreases, best compromise between J_2 and $J_4 \rightarrow$ <u>nodal BCS</u> state.

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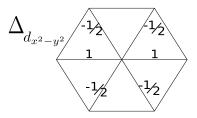
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Basic Properties of the $d_{x^2-y^2}$ spin-liquid



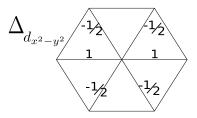
• Spinons coupled to a Z_2 gauge field.

$$S = \sum_{\langle ij \rangle} \sigma_{ij} \left(f_i^{\dagger} f_j + \Delta_{ij} f_i f_j + h.c. \right) + \prod_{\Box} \sigma_{ij} \sigma_{jk} \sigma_{kl} \sigma_{li}$$
(8)

- Spinons gapless at nodes.
- Breaks lattice orientational symmetry ⇒ 'Nematic spin-liquid'.

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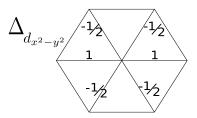
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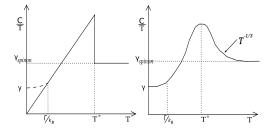


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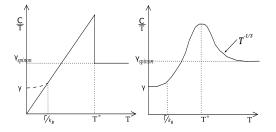
Nematic Spin Liquid: Specific heat and Spin Susceptibility



- No impurities $\Rightarrow C(T) \sim T^2$ at low T.
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- Issue with the magnetic field dependence of specific heat.

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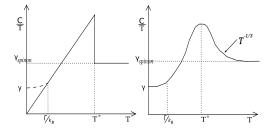
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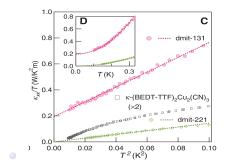


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Nematic Spin Liquid: Thermal Conductivity

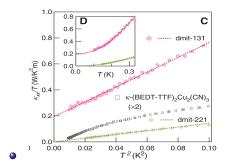
• Universal thermal conductivity $\kappa \sim T$



Yamashita et al, Science 2010.

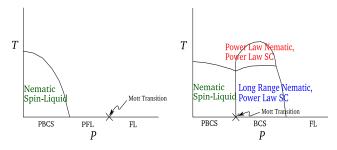
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Pressure-Temperature Phase Diagram



(a) Mott Transition After Pair-Breaking

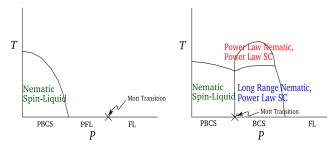
(b) Mott Transition Before Pair-Breaking

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• Case a and b relevant for two different materials.

• $a \rightarrow EtMe_3Sb[Pd(dmit)_2]_2$, $b \rightarrow \kappa CN$

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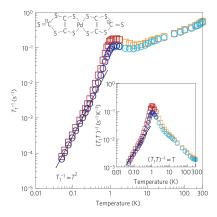


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Phase transition/cross-over in dmit-331?



T. Itou et al, Nature Physics 6, 2010.

Summary:

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 Projected nodal *d*-wave Z₂ spin liquid state is a promising candidate for *κCN*.

Questions:

- Experimental detection of the Z₂ topological order and/or spinon Fermi surface?
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- DMRG for J₂ J₄ model in quasi 1-d triangular lattice geometry (cf. arXiv:1009.1179, Matthew S. Block, D. N. Sheng, Olexei I. Motrunich, and Matthew P. A. Fisher).
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Acknowledgements

Special thanks to Matthew Fisher, Olexei Motrunich, Ying Ran, Ashvin Vishwanath for interesting discussions and suggestions.

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THANK YOU!

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