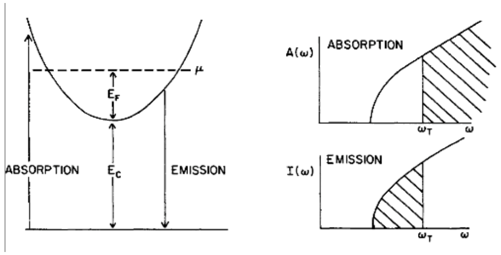
Non-Equilibrium Quantum Many-Body Systems: Universal Aspects of Weak Quenches

Stefan Kehrein
Ludwig-Maximilians-Universität München

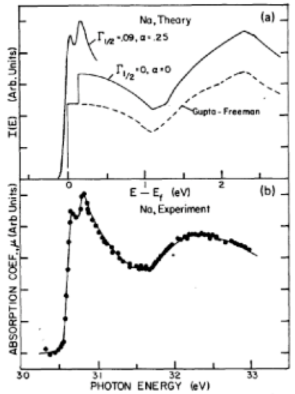
ARNOLD SOMMERFELD
CENTER FOR THEORETICAL PHYSICS

X-Ray Edge Problem



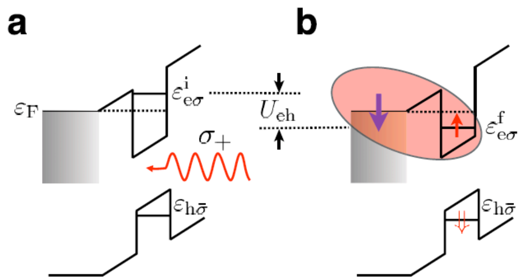
$$H_i = \sum_k \varepsilon_k c_k^\dagger c_k$$

$$H_f = \sum_k \varepsilon_k c_k^\dagger c_k + \sum_{k,k'} V_{kk'} c_{k'}^\dagger c_k$$



Mahan,
Many-Particle Physics

Kondo Excitons

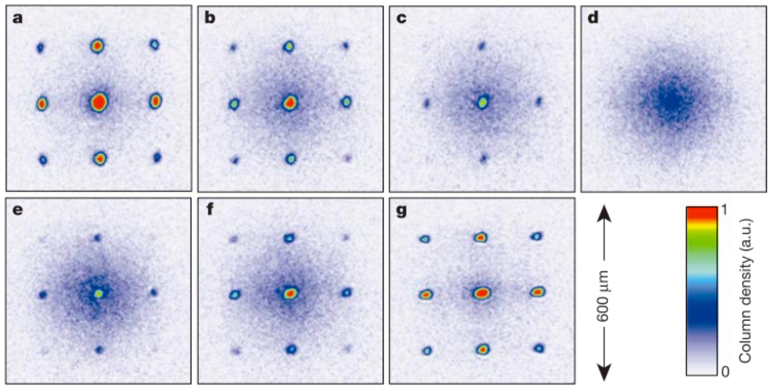


Türeci et al., arXiv:0907.3854

Initial Hamiltonian H_i: defines state $|\Psi_i\rangle$

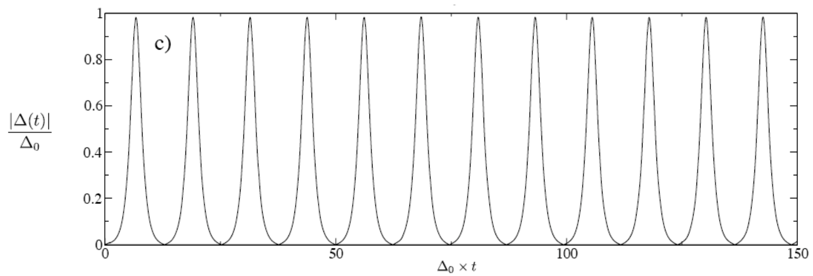
Final Hamiltonian H_f: defines time evolution $|\Psi(t)\rangle = \exp(-iH_f t)\,|\Psi_i\rangle$

Collapse and Revival



Greiner et al., Nature 419, 51 (2002)

Non-Equilibrium Cooper Pairing



Yuzbashyan et al., PRL 96 (2006)

H_i , $|\Psi_i\rangle$

Quench

H_f , $|\Psi(t)\rangle = \exp(-iH_f t)\,|\Psi_i\rangle$

Work performed on system

Excitation energy in system

$$E_{ex} = \langle \Psi_i | H_f | \Psi_i \rangle - \langle \Psi_f | H_f | \Psi_f \rangle > 0$$

Long time limit $t \to \infty$

- Memory of initial state?
- Thermalization?

- Time-independent overlap:

$$\forall t \quad \langle \Psi_f | \Psi(t) \rangle = \langle \Psi_f | \Psi_i \rangle$$

- Pure states remains pure

Universal Aspects of Weak Quenches

Working definition:
“weak” = amenable to some (suitable) perturbation theory

Factor 2 enhancement

- **I) Ferromagnetic Kondo model (d=0)**
- **II) Quenched Fermi gas (d>1)**
  (or: Beyond Landau adiabacity)
- **III) What about d=1?**

"Factor 2" for discrete Hamiltonians

- Perturbative Hamiltonian $H = H_0 + g\, H_{\text{int}}$
- Observable O(t) with (i) $O|\Omega_0\rangle = 0$ (ii) $[O, H_0] = 0$

- Ground states
  of H_0 $|\Omega_0\rangle$
  of H $|\Omega\rangle$

M. Moeckel and S. K., Ann. Phys. (NY) 324, 2146 (2009)

In second order perturbation theory:

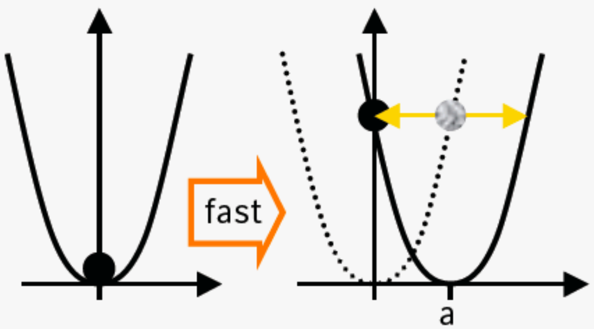
$$\lim_{T\to\infty} \frac{1}{T}\int_0^T dt\, \langle \Omega_0(t)|O|\Omega_0(t)\rangle = 2\, \langle \Omega|O|\Omega\rangle$$

Expectation value of excited state (nonequilibrium)

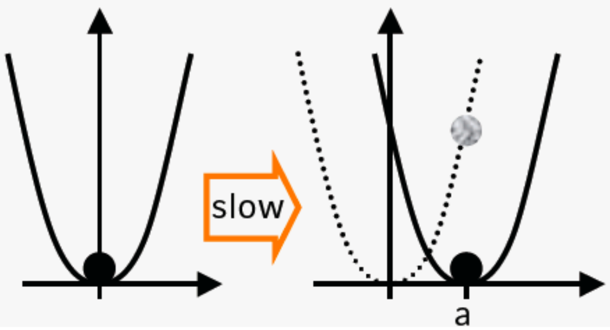
Ground state expectation value (equilibrium)

$$V^{(0)} = \frac{1}{2}x^2 \;\rightarrow\; V^{(S)} = \frac{1}{2}(x-a)^2$$

Sudden shift



Adiabatic shift

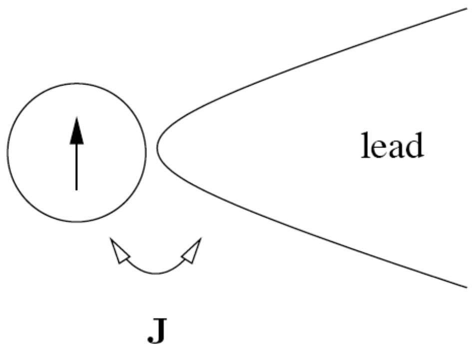


Ferromagnetic Kondo Model

A. Hackl, D. Roosen, S. K., W. Hofstetter, Phys. Rev. Lett. 102, 196601 (2009)
A. Hackl, M. Vojta and S. K., Phys. Rev. B 80, 195117 (2009)

$$H_i = \sum_{k,\sigma} \varepsilon_k c^\dagger_{k\alpha} c_{k\alpha} - 0^+ S_z$$

infinitesimal magnetic field



⇒ Product initial state: $|\Psi_i\rangle = |\uparrow\rangle \otimes |FS\rangle$

$$H_f = \sum_{k,\sigma} \varepsilon_k c^\dagger_{k\alpha} c_{k\alpha} - 0^+ S_z + J \vec{S} \cdot \sum c_{k'\alpha} \vec{\sigma}_{\alpha\beta} c_{k\beta}$$

ferromagnetic coupling (J<0): Coupling constant flows to zero
⇒ Expansion becomes better (asymptotically exact) for long times

Nonequilibrium spin expectation value:

$$\langle S_z(t) \rangle = \frac{1}{2}\left(\frac{1}{\ln(t) - \frac{1}{\rho J}} + 1 + \rho J + O(J^2) \right).$$

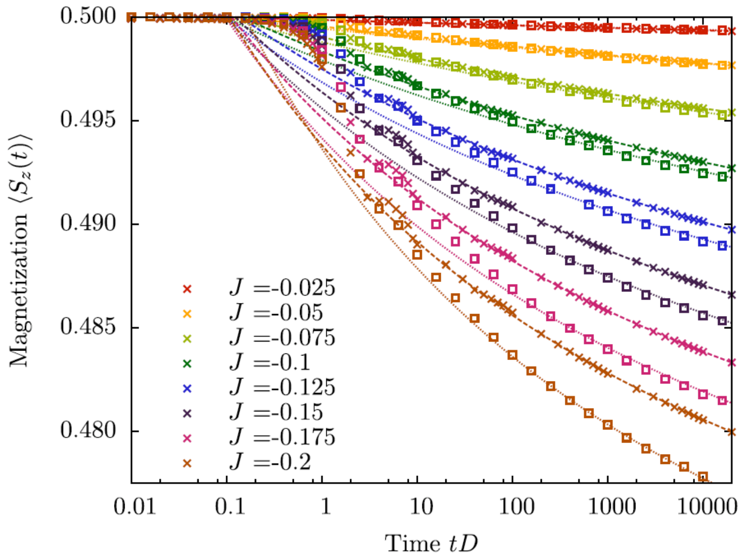
Equilibrium:

$$\langle S_z \rangle_{eq} = \frac{1}{2}\left(1 + \frac{\rho J}{2} + O(J^2)\right)$$

Observable:

$$O = S_z - \frac{1}{2}$$

Comparison with TD-NRG:
(Hackl et al., PRL 102)



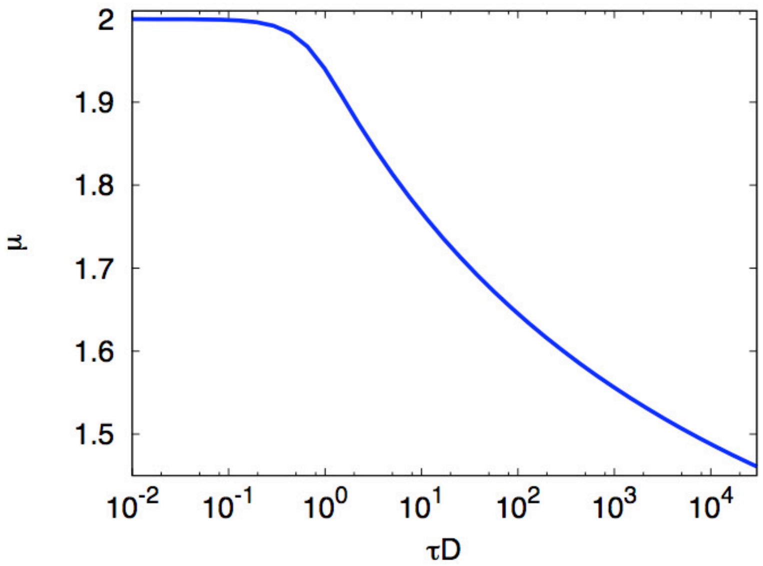
⇒ System remembers initial quantum state for all times

Crossover from adiabatic to instantaneous quenching:

Coupling J switched on on timescale τ

Measure of non-adiabacity:

$$\mu \stackrel{def}{=} \frac{\lim_{t\to\infty}\langle O(t)\rangle_{neq} - \langle O\rangle_0}{\langle O\rangle_{eq} - \langle O\rangle_0}$$



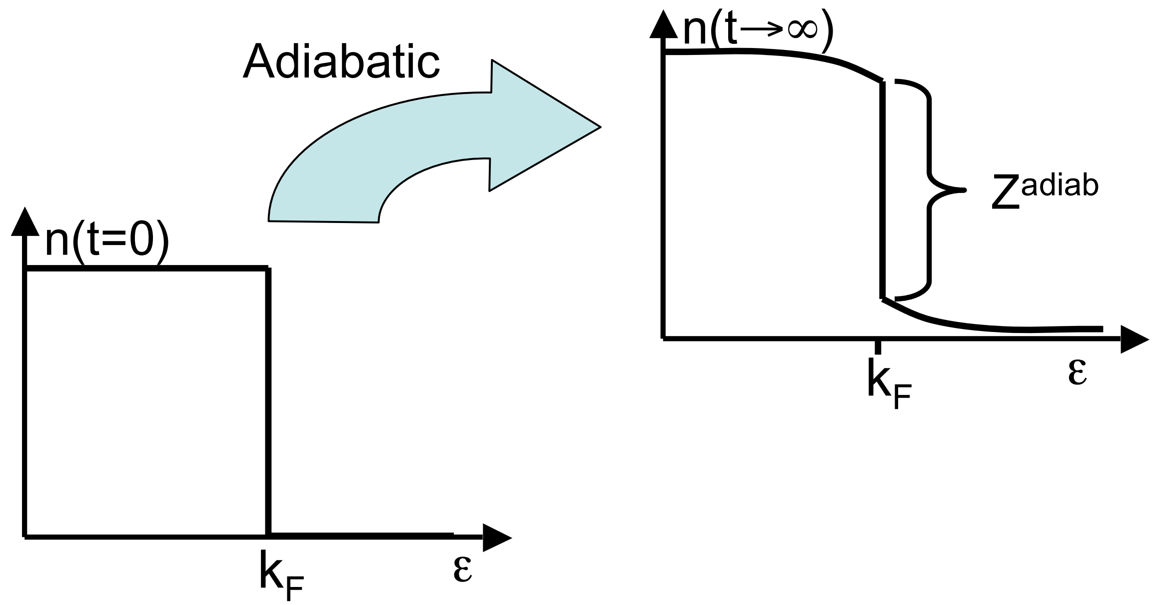
⇒ Crossover timescale nonperturbative (exponentially large) due to RG flow

[C. Tomaras, S. K., Preprint this week]

Sudden Interaction Fermi Liquid

Landau Fermi liquid theory:

Adiabatic switching on of interaction
→ 1 to 1 correspondence between physical electrons and quasiparticles

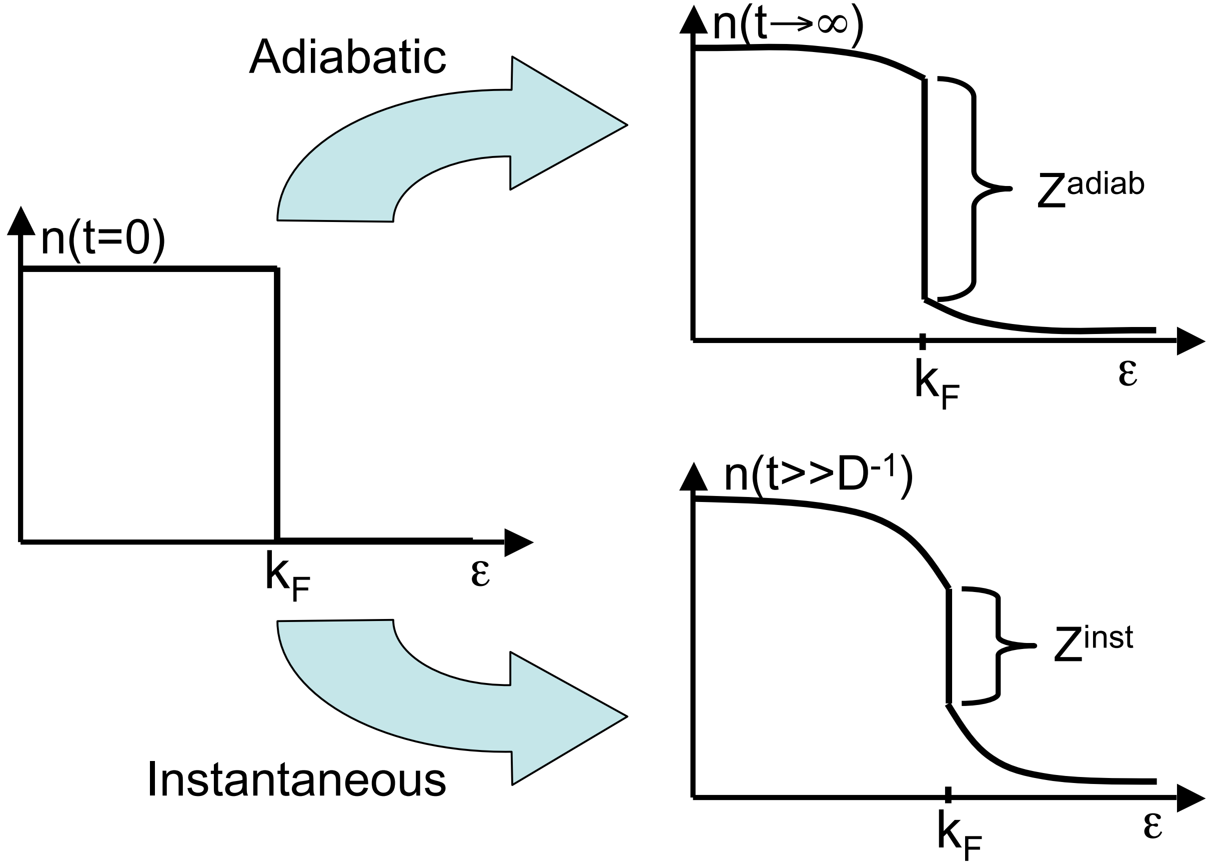


What happens for sudden switching?

Initial Hamiltonian: $$H_i = \sum_{k,\alpha} \varepsilon_k c^\dagger_{k\alpha} c_{k\alpha}$$

Observable $$O = \begin{cases} n_k & \text{for} \quad k > k_F \\ 1 - n_k & \text{for} \quad k < k_F \end{cases}$$

Calculation to order U^2



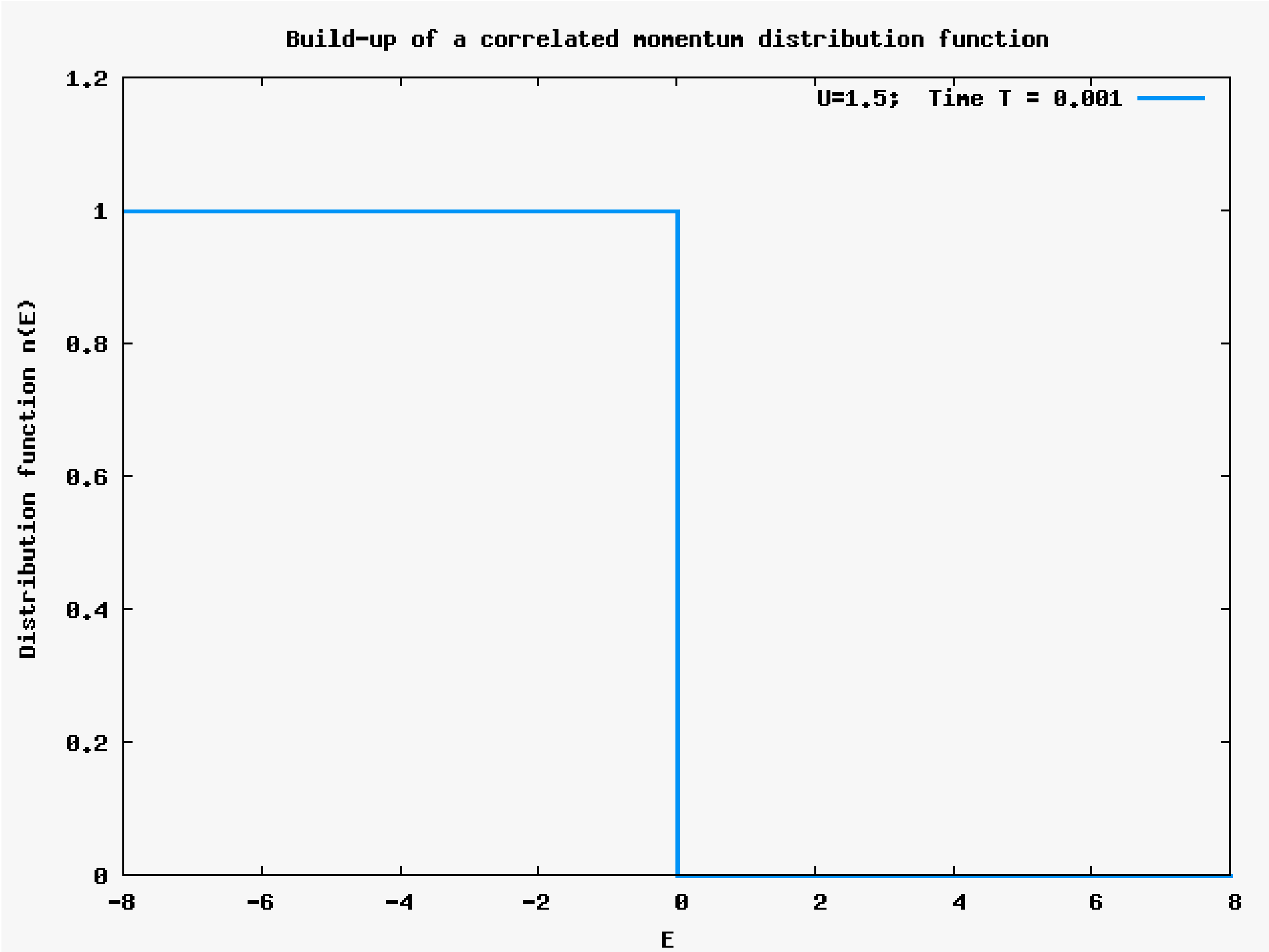
Sudden switching looks like T=0 Fermi liquid with "wrong" quasiparticle residue:

$$1-Z^{inst} = 2\,(1-Z^{adiab})$$

M. Möckel and S. K., Phys. Rev. Lett. 100, 175702 (2008); Ann. Phys. 324, 2146 (2009)

Hubbard model in d≥2 dimensions:

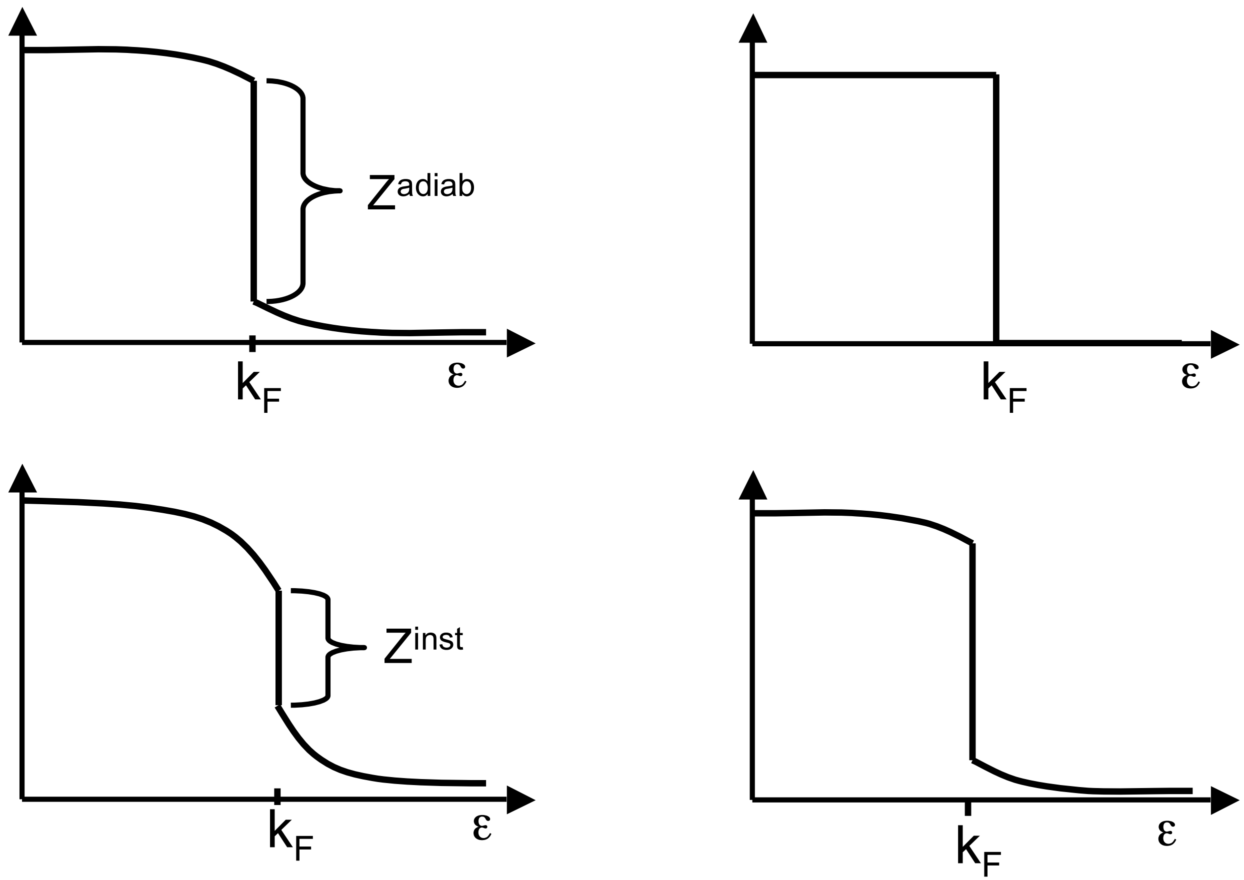
$$H = \sum_{k,\alpha} \varepsilon_k c_{k\alpha}^\dagger c_{k\alpha} + U\,\Theta(t) \sum_i \left(n_{i\uparrow} - \frac{1}{2}\right)\left(n_{i\downarrow} - \frac{1}{2}\right)$$



Physical electrons | Quasiparticles

Equilibrium (T=0)

Quench after $t \gg D^{-1}$



Nonthermal distribution function

⇒ unstable under Quantum Boltzmann equation dynamics

⇒ thermalization on timescale $t \propto U^{-4}$

Sudden quench (generic weak interaction g)

Time scale

$t \propto D^{-1}$ • Formation of quasiparticles

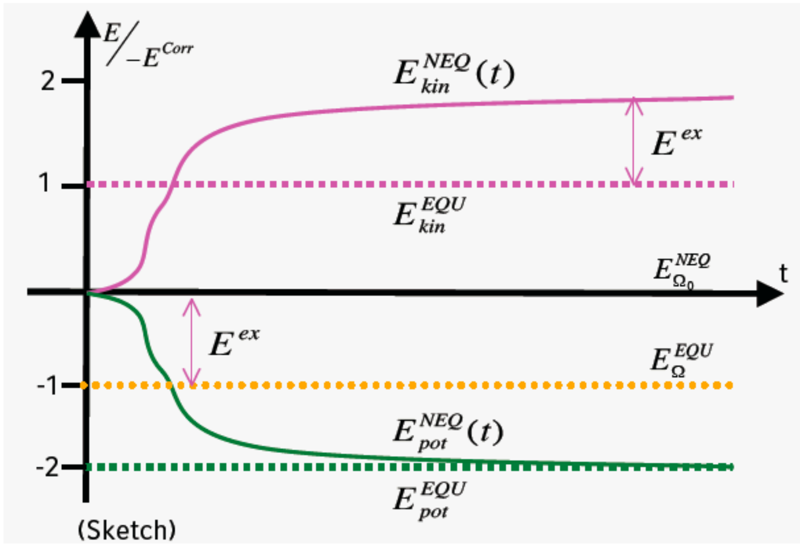
$D^{-1} \ll t \ll g^{-2}$ • T=0 Fermi liquid with "wrong" quasiparticle residue:

$$1-Z^{inst} = 2\,(1-Z^{adiab})$$

$D^{-1} \ll t \ll g^{-4}$

• Quasi-steady state

• Prethermalization (Berges et al. 2004)



$t \propto g^{-4}$ • Quantum Boltzmann equation (quasiparticles explore available phase space):

Thermalization with $T_{eff} \propto g$

Numerical Studies

M. Eckstein, M. Kollar and P. Werner, Phys. Rev. Lett. 103, 056403 (2009);
Phys. Rev. B 81, 115131 (2010)

Non-equilibrium DMFT with real time QMC for interaction quench in half-filled Hubbard model

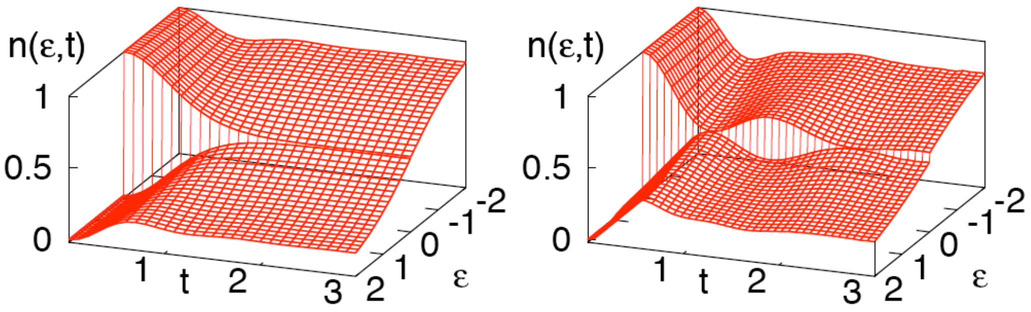
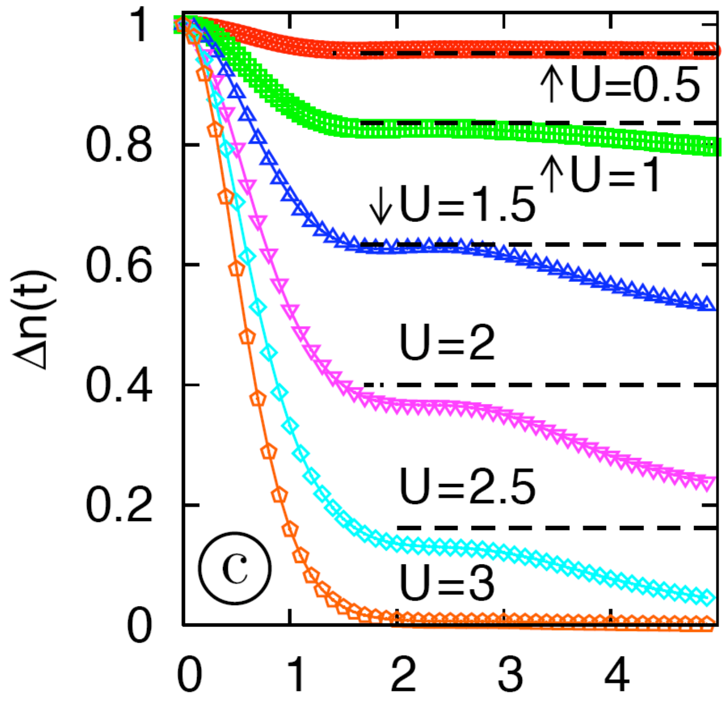


FIG. 1: Momentum distribution $n(\epsilon_{\boldsymbol{k}}, t)$ for quenches from $U = 0$ to $U = 3$ (left panel) and $U = 5$ (right panel).



Nonequilibrium Cooper pairing

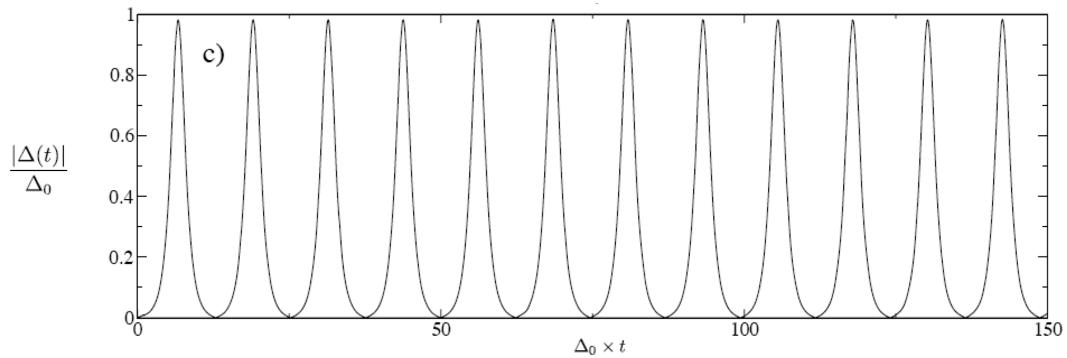
Barankov et al., PRL 93 (2004); Yuzbashyan et al., J. Phys. A 38 (2005)

Time-dependent BCS-Hamiltonian:

$$H = \sum_{k,\alpha} \varepsilon_k c^\dagger_{k\alpha} c_{k\alpha} - \Theta(t)\, g \sum_{k,k'} c^\dagger_{k\uparrow} c^\dagger_{-k\downarrow} c_{-k'\downarrow} c_{k'\uparrow}$$

BCS channel

Coherent order parameter oscillations for quench from normal state



Yuzbashyan et al., PRL 96 (2006)

Problem: For weak physical (short range) interaction

$$\Delta_{BCS} \propto \exp(-1/g) << T_{eff} \propto g$$

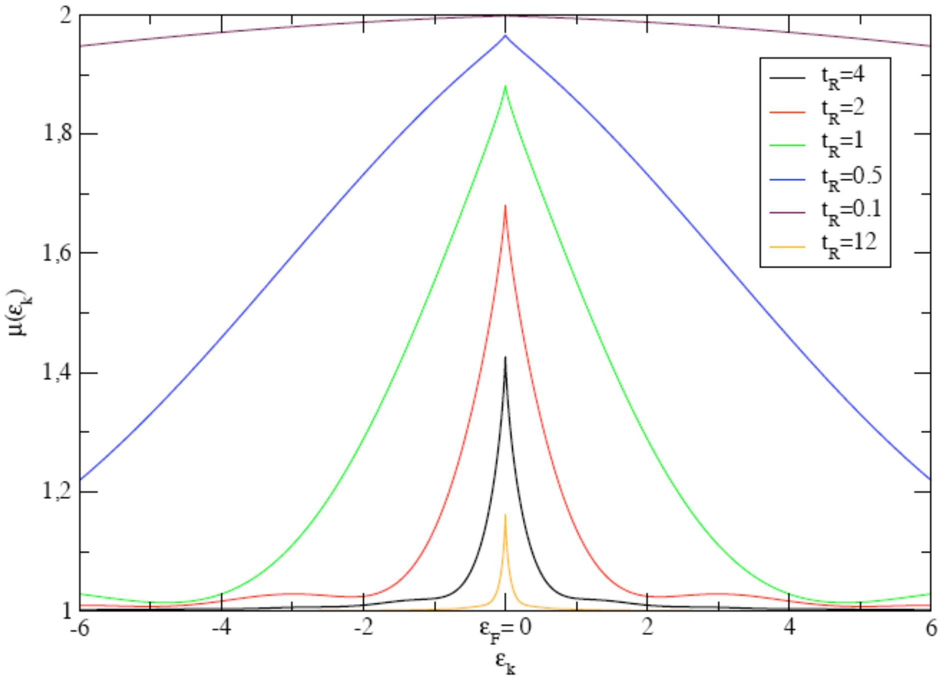
$\Rightarrow$ Thermalization destroys BCS physics

Crossover from adiabatic to instantaneous quenching

M. Moeckel and S. K., New J. Phys. 12, 055016 (2010)

Linear ramping on timescale t_R:

$$U(t) = U \begin{cases} 0 & t \leq 0 \\ t/t_R & 0 < t < t_R \\ 1 & t > t_R \end{cases}$$

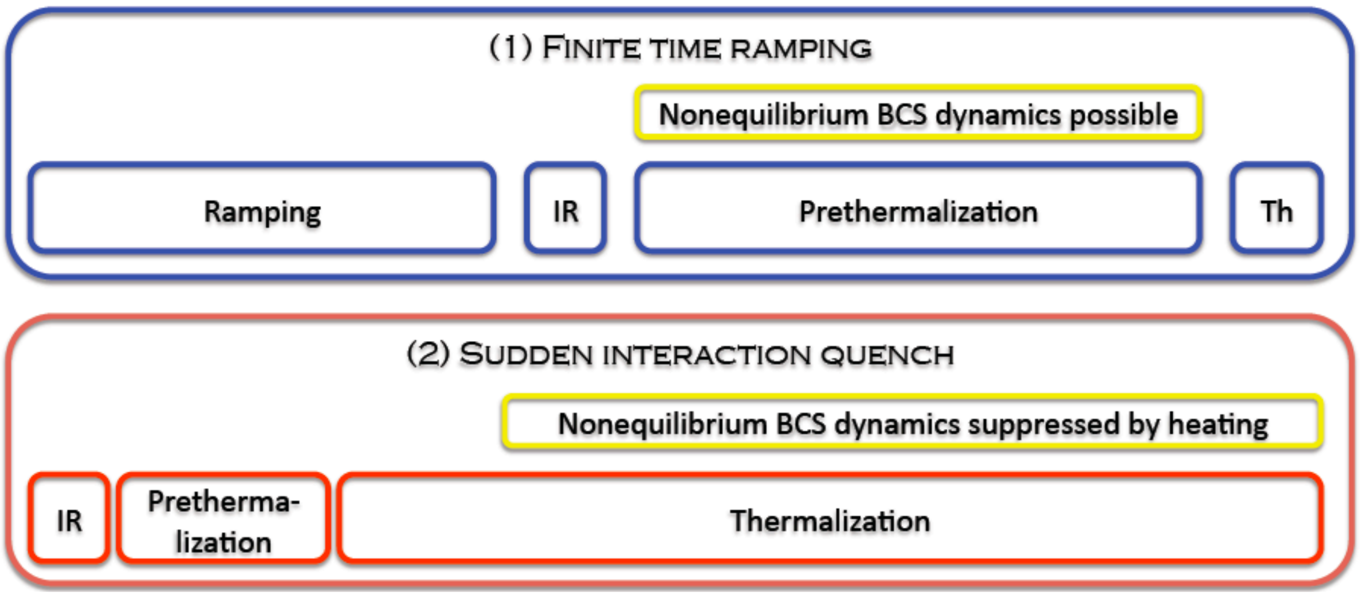


Uniform approach to equilibrium also at Fermi surface!

Effective temperature for linear ramping: $T_{eff} \propto t_R^{-1}$

Quench "instantaneous" for $t_R \leq \Delta_{BCS}^{-1}$

But: Excitation energy mainly in high-energy modes & relaxation bottleneck due to prethermalization plateau
$\Rightarrow$ Broadening of Fermi surface delayed to times $\propto t_R^2$

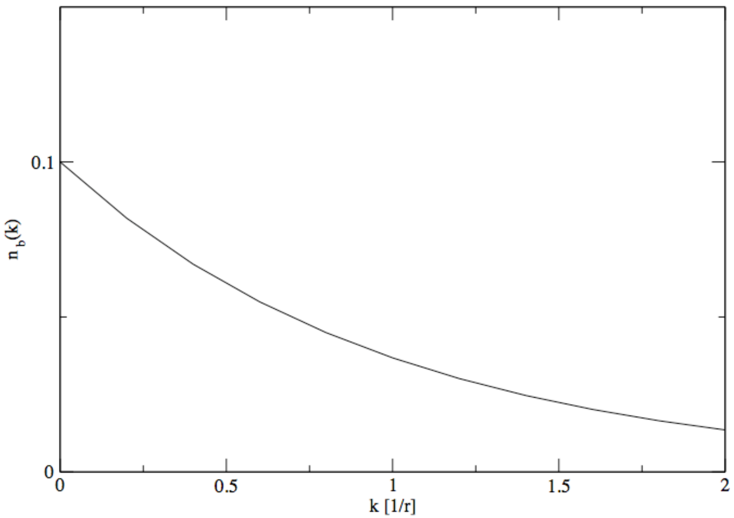


Interaction Quench in a Luttinger Liquid

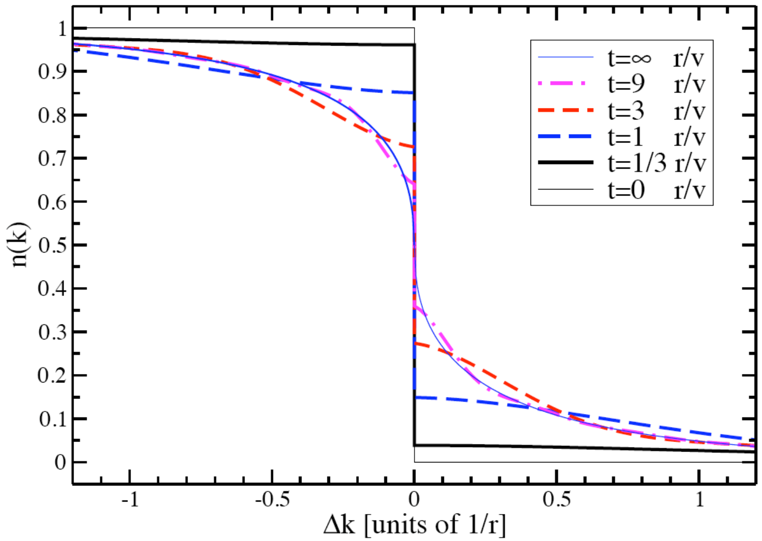
M. Cazalilla, Phys. Rev. Lett. 97 (2006):

Sudden forward scattering in 1d Fermi gas
⇒ Exact solution via Bogoliubov transformation to free bosons

No time evolution for free bosons (quasiparticles)



Time evolution for physical fermions (G. Uhrig, Phys. Rev. A 80)



Quasiparticle momentum distribution function time-invariant due to lack of quasiparticle interaction.

However, even Boltzmann dynamics from 2-quasiparticle interaction is generically ineffective in 1d!
(No time evolution beyond prethermalized regime.)

Systems with well-defined quasiparticles

No quasiparticle description possible

Prethermalization: Momentum-averaged quantities time-indep., but distribution over momentum modes non-thermal

- Weak quenches for d>1:
  - Fast prethermalization
  - Thermalization with Boltzmann dynamics

- Quench "to" Mott-Hubbard transition ($d=\infty$)
- Strong-coupling quantum critical points

- Weak quenches for d=1:
  - Fast prethermalization
  - Long time limit? (no 2-particle Boltzmann dynamics, constraints due to integrability, ...)