

Synthetic dimensions in integrated photonics:

From optical isolation to
4D quantum Hall physics

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Many-Body Physics with Light
KITP, November 11th 2015



Outline

- Background on synthetic dimensions
- Synthetic dimensions in integrated photonics: how to do it?
- Topological physics in a **1D** ring-resonator chain:
 - Edge states for optical isolation
 - **2D** quantum Hall effect in photon transport
- Topological physics in a in a **3D** resonator array:
 - **4D** quantum Hall effect in photon transport

Synthetic dimensions in integrated photonics: arXiv:1510.03910

Tomoki Ozawa, HMP, Nathan Goldman, Oded Zilberberg & Iacopo Carusotto



*INO-CNR BEC Center &
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*Université Libre de
Bruxelles*



ETH Zurich



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4D quantum Hall effect with ultracold atoms: Phys. Rev. Lett. 115, 195303 (2015)

HMP, Oded Zilberberg, Tomoki Ozawa, Iacopo Carusotto & Nathan Goldman

Topological Physics for Photons

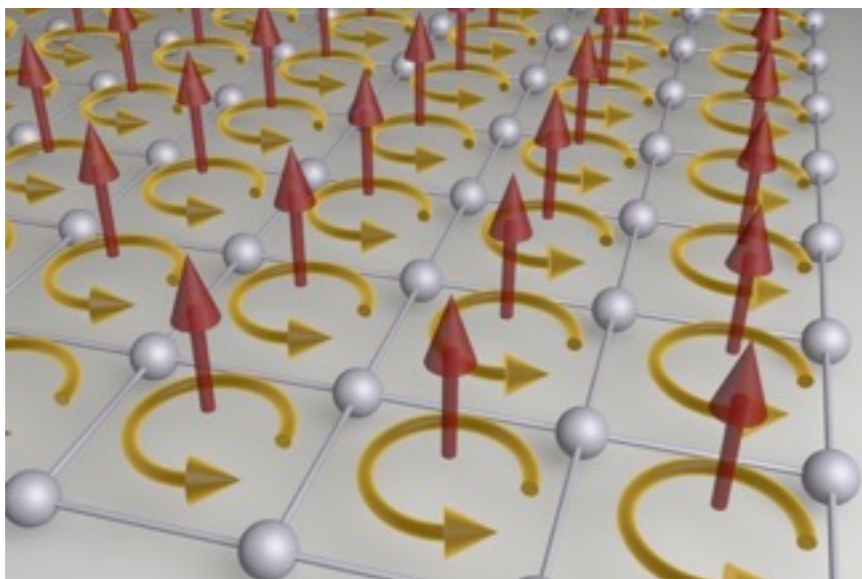
- Energy bands with topological invariants
This Talk: *1st and 2nd Chern numbers*
- Topological edge states for unidirectional propagation of light
- New topological physics: 4D quantum Hall effect

Focus on variants of Hofstadter model for charged particle:

Hofstadter, PRB 14, 2239, (1976)

$$H = -J \sum_{x,y} \left(\hat{b}_{x+a,y}^\dagger \hat{b}_{x,y} + e^{i2\pi\Phi x/a} \hat{b}_{x,y+a}^\dagger \hat{b}_{x,y} + \text{h.c.} \right)$$

in Landau gauge



uniform perpendicular
magnetic flux per plaquette of $2\pi\Phi$

energy bands with
non-zero 1st Chern numbers

Topological Physics for Photons

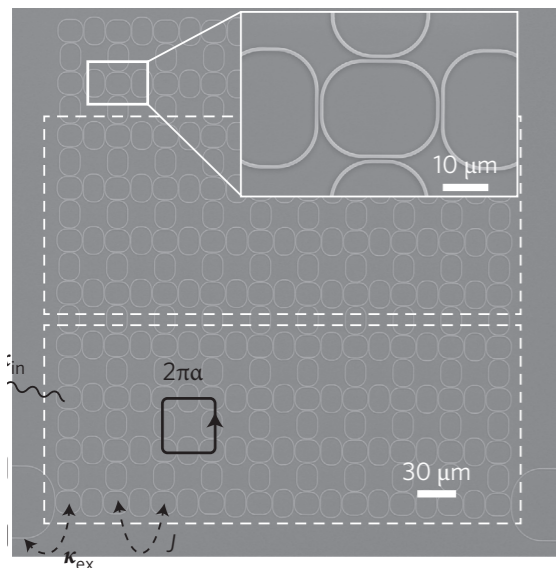
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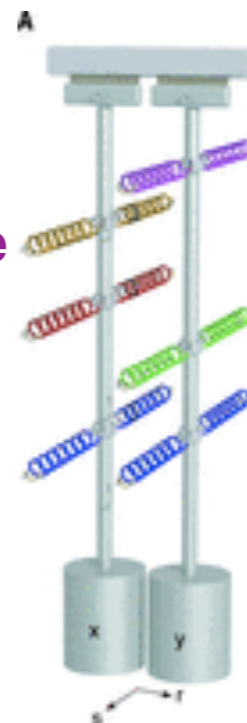
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in silicon photonics



Hafezi et al, Nat. Photon.
7, 1001, (2013)

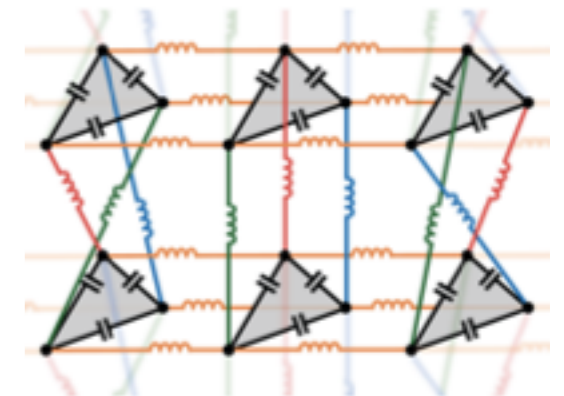
**Also in solid-state
 superlattices,
 ultracold atoms...**



pendula:
Susstrunk & Huber,
Science 349, 47
(2015)

& time-modulation proposal:
Salerno, Ozawa, HMP,
Carusotto,
arXiv:1510.04697

in Landau gauge



circuits:

Expt: **Ningyuan et al., PRX**
5, 021031, (2015)
 Theory: **V. V. Albert, et al**
PRL 114, 173902 (2015).

Synthetic dimensions in ultracold atoms

Ingredients:

1. Choose degrees of freedom \rightarrow site indices in synthetic dimension
2. Couple these degrees of freedom \rightarrow "hopping"

1. Spin states
2. Raman laser coupling

Theory:
 Boada et al., PRL, 108, 133001 (2012)
 Celi et al., PRL, 112, 043001 (2014)
 ... HMP et al, PRL 115, 195303 (2015)

Experiments:
 Mancini et al, Science, 349, 1510 (2015)
 Stuhl et al. Science, 349, 1514 (2015)

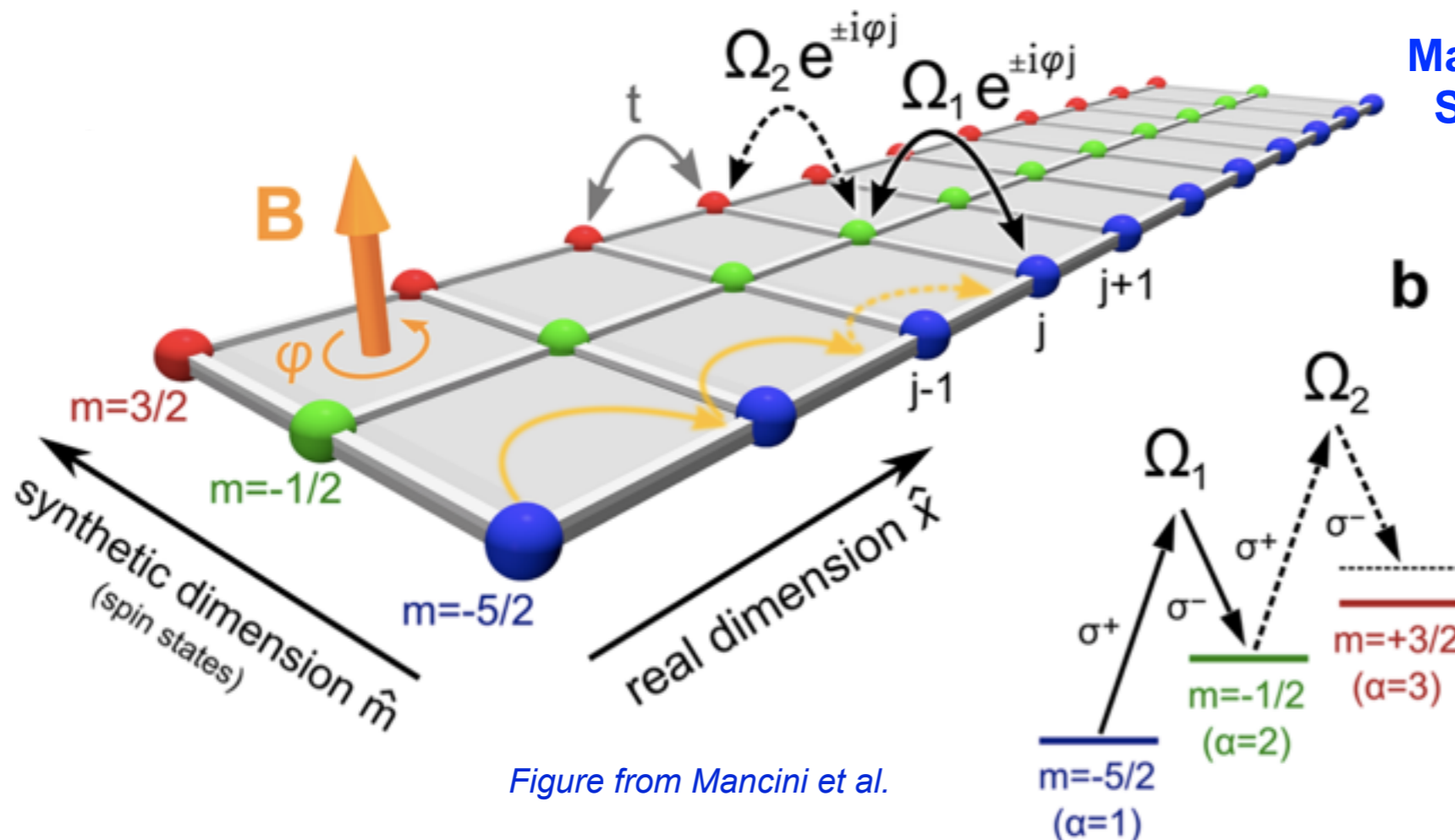


Figure from Mancini et al.

b
 3-leg Hofstadter ladder
 for atoms in a 1D
 optical lattice

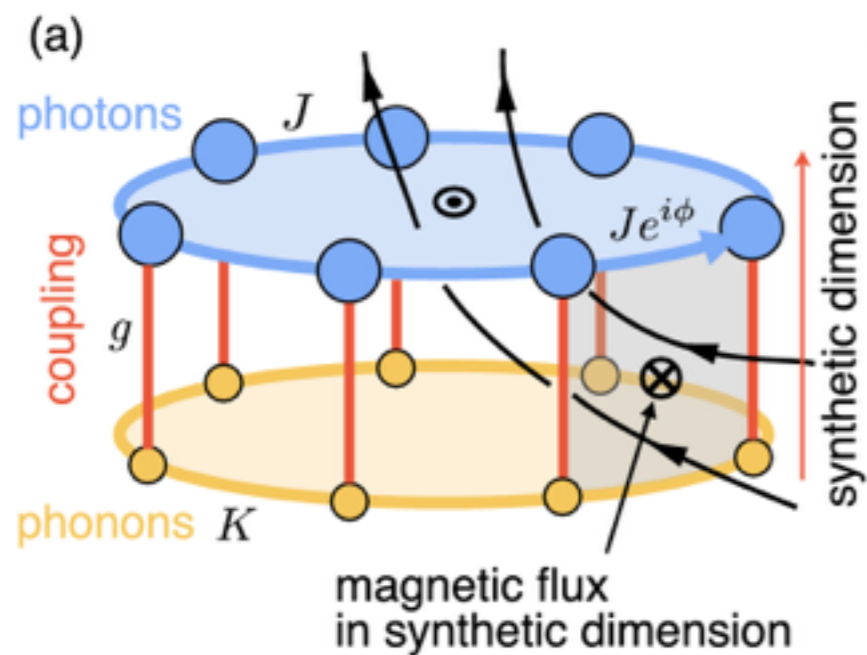
^{173}Yb
 $I = 5/2$

Synthetic dimensions in photonics so far...

Ingredients:

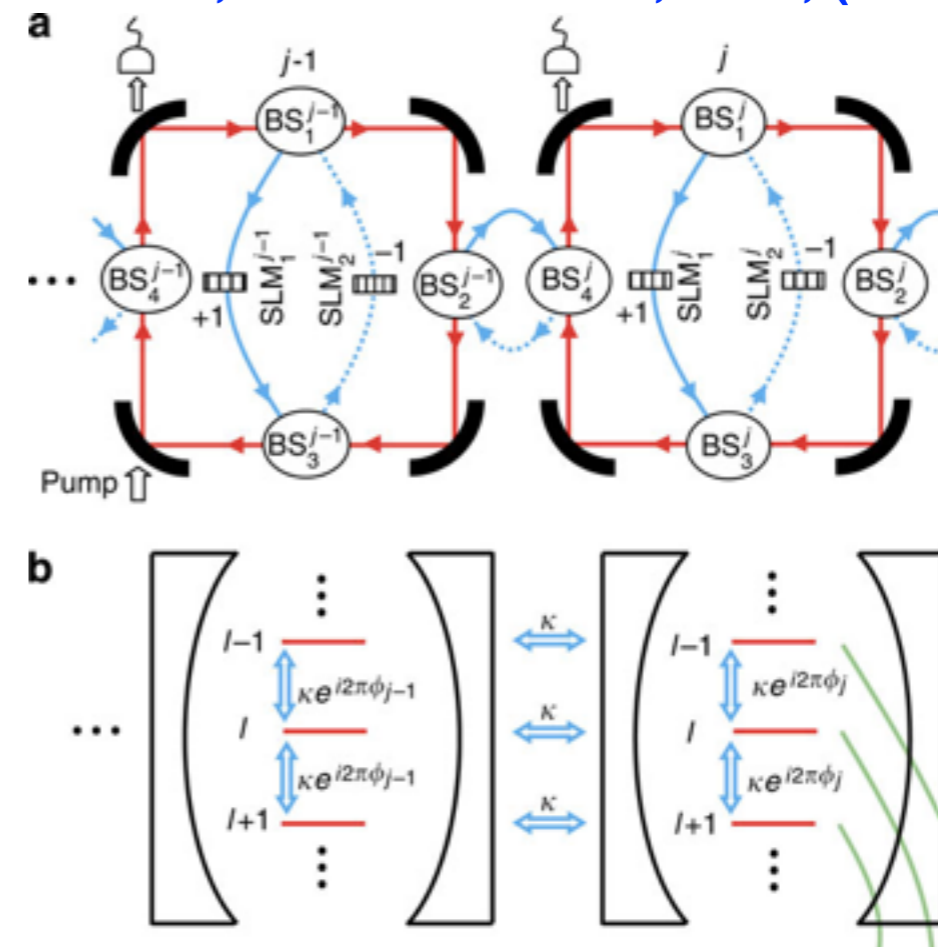
1. Choose degrees of freedom \rightarrow site indices in synthetic dimension
2. Couple these degrees of freedom \rightarrow "hopping"

Schmidt et al, *Optica* 2, 7, 635 (2015)



1. Photons & phonons
2. Optomechanical coupling

Luo et al, *Nature Comm.* 6, 7704, (2015)



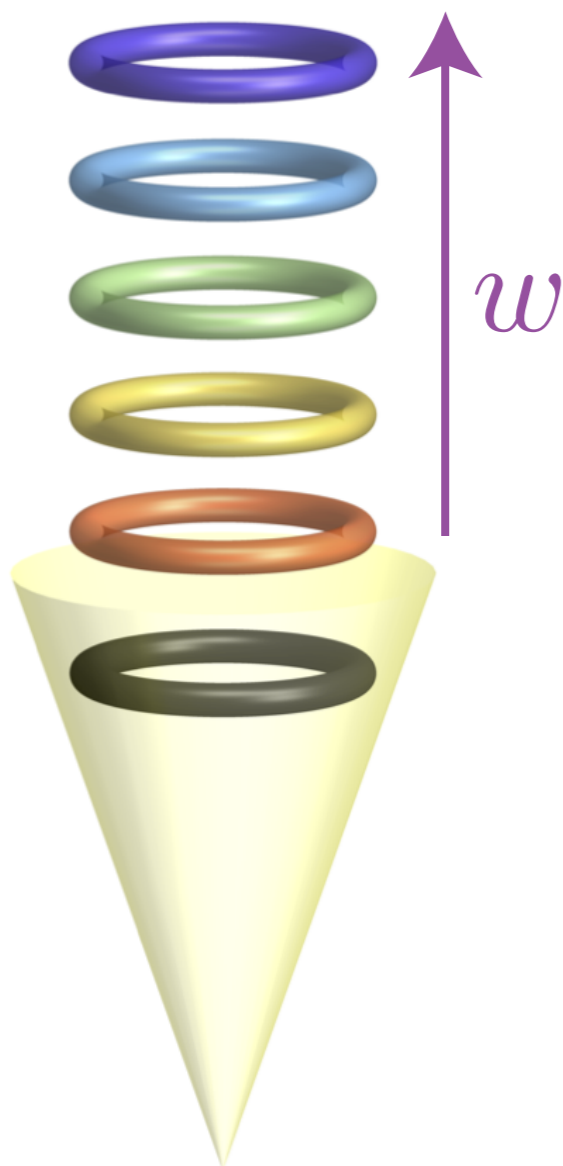
1. Orbital angular momentum of cavity modes
2. Spatial light modulators

...and also Schleier-Smith et al, KITP DenseLight conference

Synthetic dimensions in integrated photonics

Ingredients:

1. Choose degrees of freedom \rightarrow site indices in synthetic dimension
2. Couple these degrees of freedom \rightarrow "hopping" ($\hbar = 1$)



1. Modes of a ring resonator

$$\omega_w = \omega_{w_0} + \Delta\omega (|w| - w_0) + \frac{D}{2} (|w| - w_0)^2 + \dots$$

free
spectral
range

$$\Delta\omega = 2\pi c / n_{\text{eff}} R$$

2. Strong beam(s) modulates ϵ_{ij} at $\omega_{\text{mod}} \approx \eta\delta\omega$ via optical nonlinearity

$$H_{\text{mod}} = - \sum_w \Omega_{\mathbf{r}}(t) \hat{a}_{\mathbf{r}, w+\eta}^\dagger \hat{a}_{\mathbf{r}, w} + \text{h.c.},$$

$$\Omega_{\mathbf{r}}(t) = \Omega_{\mathbf{r}}^o e^{-i\omega_{\text{mod}} t}$$

Synthetic dimensions in integrated photonics

To get static inter-mode coupling from:

$$H_{\text{mod}} = - \sum_w \Omega_{\mathbf{r}}(t) \hat{a}_{\mathbf{r},w+\eta}^\dagger \hat{a}_{\mathbf{r},w} + \text{h.c.},$$

$$\Omega_{\mathbf{r}}(t) = \Omega_{\mathbf{r}}^o e^{-i\omega_{\text{mod}} t}$$

Go to rotating frame:

$$\hat{b}_{\mathbf{r},w}(t) \equiv \hat{a}_{\mathbf{r},w}(t) e^{i[\omega_{w_0} + (w-w_0)\omega_{\text{mod}}/\eta]t},$$

Can be complex:
easily tuneable synthetic
gauge field

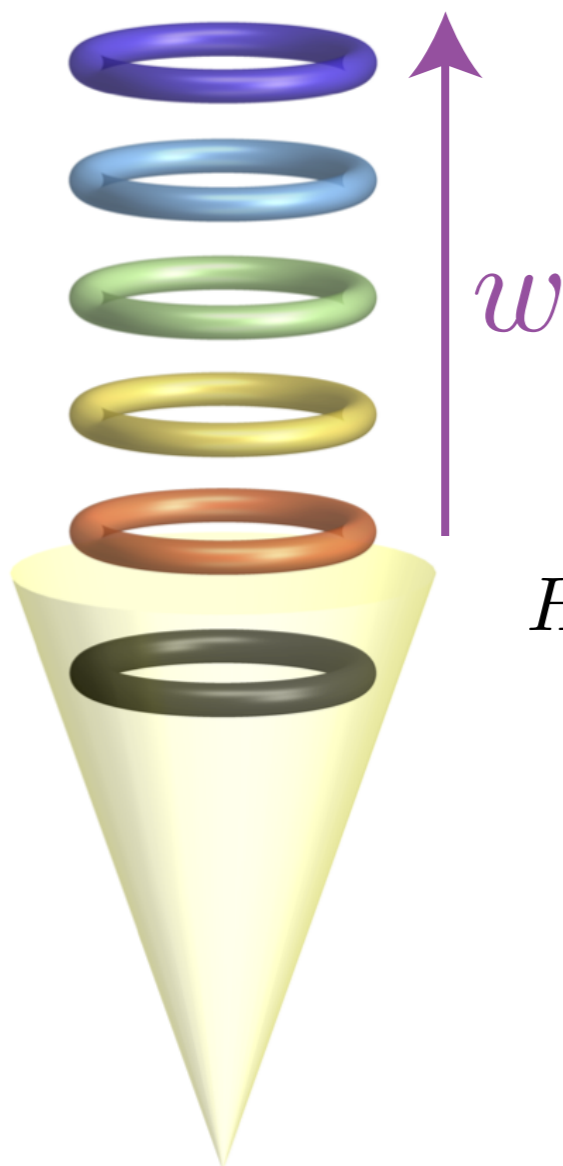
“hopping” in synthetic
dimension

$$H_{\Omega} = \sum_{w,\mathbf{r}} \left[-\Omega_{\mathbf{r}}^o \hat{b}_{\mathbf{r},w+\eta}^\dagger \hat{b}_{\mathbf{r},w} + \text{h.c.} \right.$$

$$\left. + \left((\Delta\omega - \omega_{\text{mod}}/\eta) (w - w_0) + \frac{D}{2} (w - w_0)^2 \right) \hat{b}_{\mathbf{r},w}^\dagger \hat{b}_{\mathbf{r},w} \right].$$

“uniform force” along
synthetic dimension

“harmonic potential” along
synthetic dimension



Extending to a lattice of resonators

$$H = H_J + H_\Omega$$

spatial tunnelling

coupling in
synthetic dimension

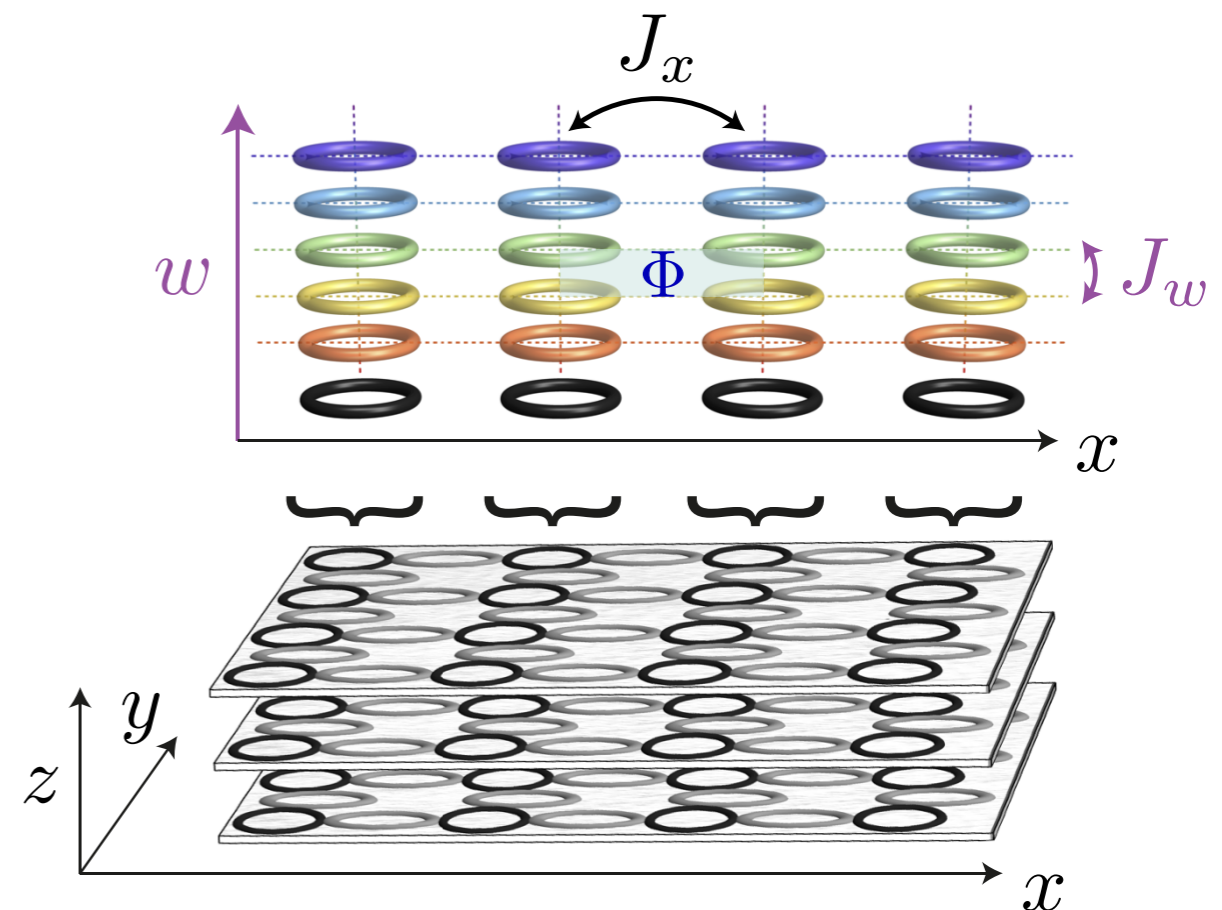
$$H_J = - \sum_{w, \mathbf{r}, j} J_j \hat{b}_{\mathbf{r}+\mathbf{a}_j, w}^\dagger \hat{b}_{\mathbf{r}, w} + \text{h.c.}$$

$$H_\Omega \approx - \sum_{w, \mathbf{r}} \Omega_{\mathbf{r}}^o \hat{b}_{\mathbf{r}, w+\eta}^\dagger \hat{b}_{\mathbf{r}, w} + \text{h.c.}$$

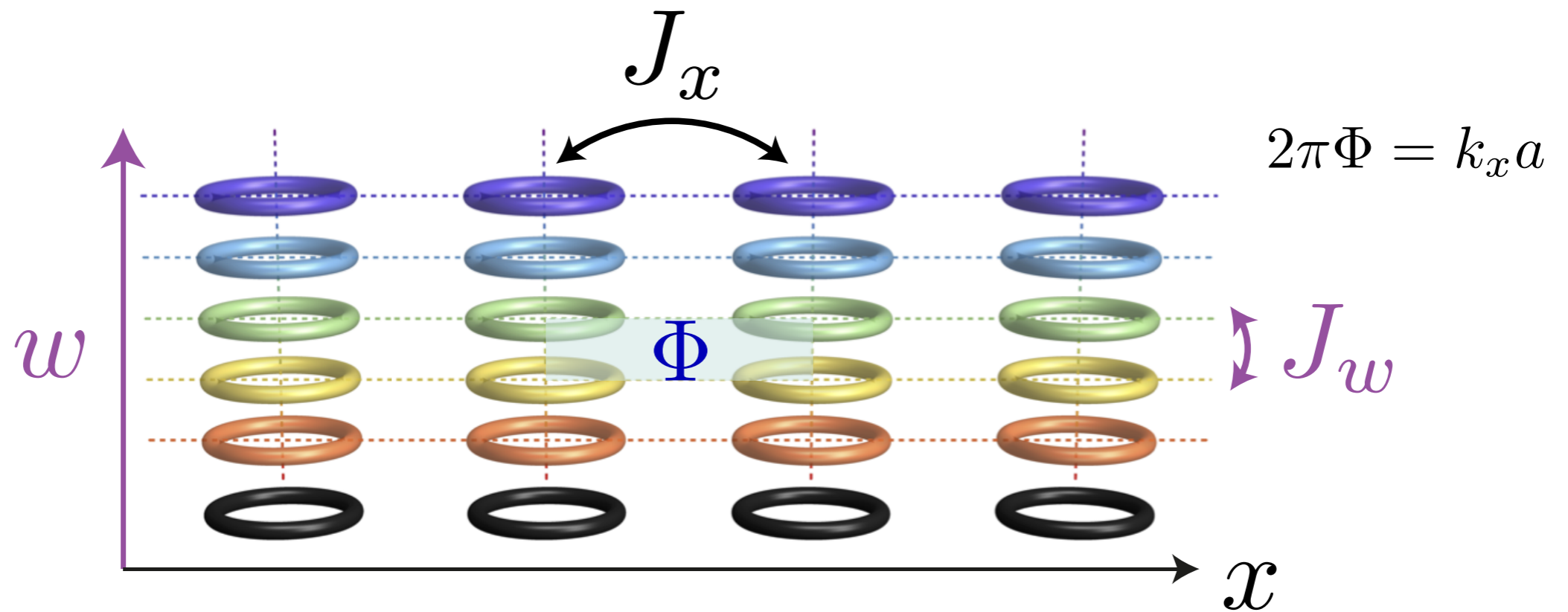
Now to specific systems:

1. 1D chain \rightarrow effective 2D lattice

2. 3D array \rightarrow effective 4D lattice



1D resonator chain with the synthetic dimension



$$H = H_J + H_\Omega$$

Choose explicitly: $\Omega_x^o = |\Omega_x^o| e^{ik_x x}$

$$H = - \sum_{x,w} \left(J_x \hat{b}_{x+a,w}^\dagger \hat{b}_{x,w} + |\Omega_x^o| e^{ik_x x} \hat{b}_{x,w+\eta}^\dagger \hat{b}_{x,w} + \text{h.c.} \right) \quad \text{Hofstadter Model}$$

Ultracold atoms: Celi et al., PRL 112, 043001 (2014) (Theory)

Mancini et al, Science, 349, 1510 (2015) (Expt)

Stuhl et al. Science, 349, 1514 (2015) (Expt)

Optical cavities: Luo et al, Nature Comm. 6, 7704, (2015) (Theory)

A few more details about the Hofstadter model

$$H = - \sum_{x,w} \left(J_x \hat{b}_{x+a,w}^\dagger \hat{b}_{x,w} + |\Omega_x^0| e^{ik_x x} \hat{b}_{x,w+\eta}^\dagger \hat{b}_{x,w} + \text{h.c.} \right) \quad \begin{aligned} 2\pi\Phi &= k_x a \\ \Phi &= p/q \end{aligned}$$

**Topological
invariant**

$$\mathcal{F}^{xw} = i \left(\langle \partial_{k_x} u | \partial_{k_w} u \rangle - \langle \partial_{k_w} u | \partial_{k_x} u \rangle \right)$$

geometrical Berry curvature

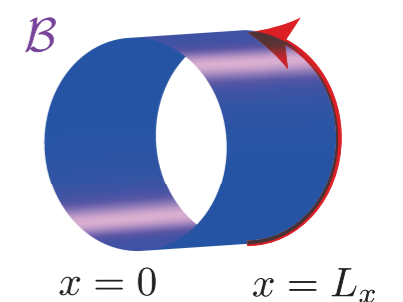
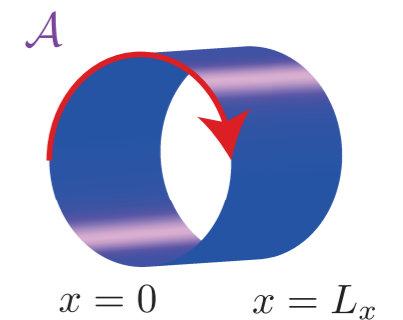
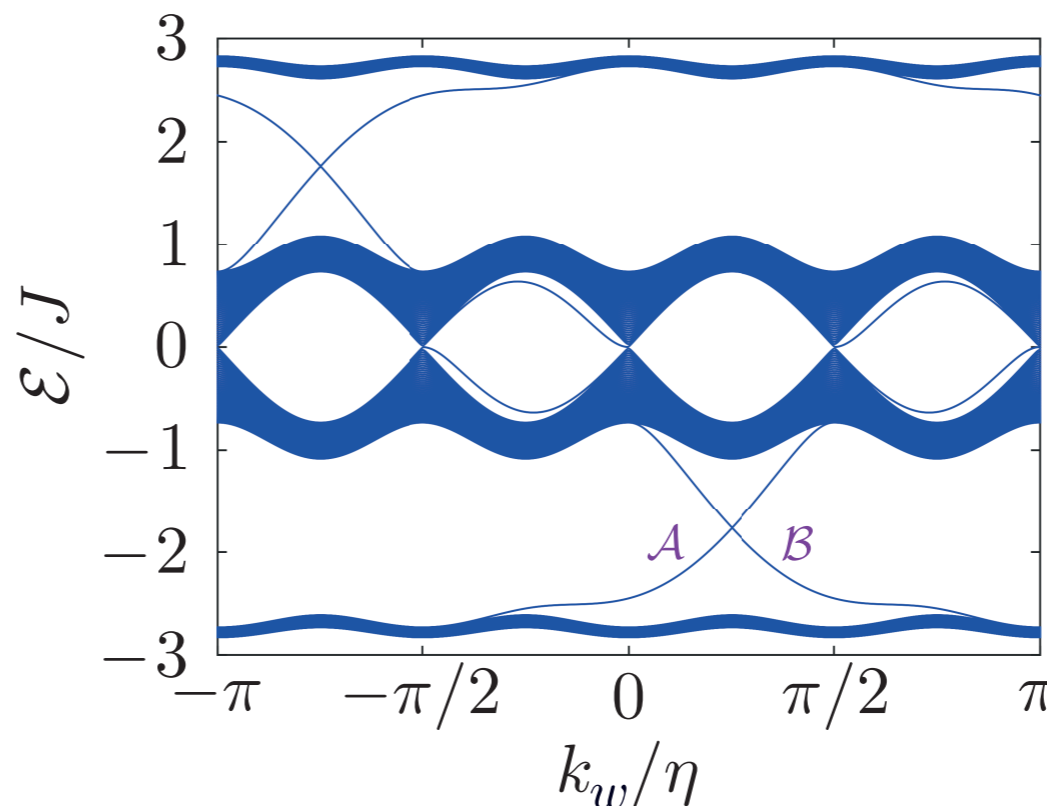
$$\nu_1^{xw} = \int_{BZ} \mathcal{F}^{xw} d^2k \in \mathbb{Z}$$

topological 1st Chern number

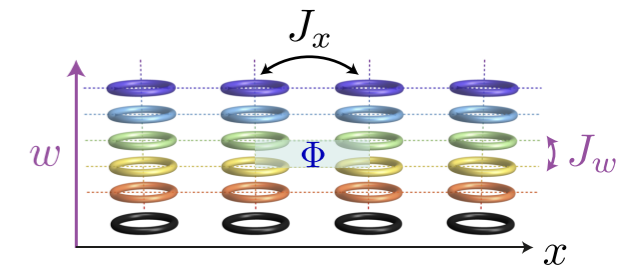
**Topological
edge States**

$$\Phi = 1/4$$

*Energy spectrum with
PBCs around w and
finite length along x*



Driven-dissipative experiments



Include **losses** $\gamma_{x,w}$

and **coherent driving**

**Ozawa & Carusotto,
PRL, 112, 133902, (2014)**

$$H_{\text{drive}} = \sum_{x,w} \left(f_{x,w}(t) \hat{b}_{x,w} + f_{x,w}^*(t) \hat{b}_{x,w}^\dagger \right)$$

assume $f_{x,w}(t) \propto e^{-i\omega_{\text{drive}} t}$

i.e. in non-rotating frame

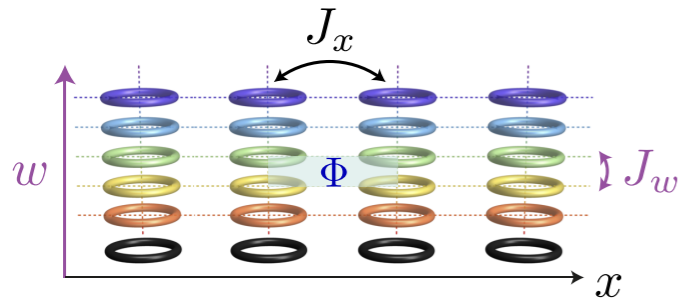
$$\omega_w + \omega_{\text{drive}}$$

Look for long-time **steady-state**

- Consider single-site driving, couples to all momenta in the Brillouin zone
- Vary position and frequency of drive to excite different supermodes

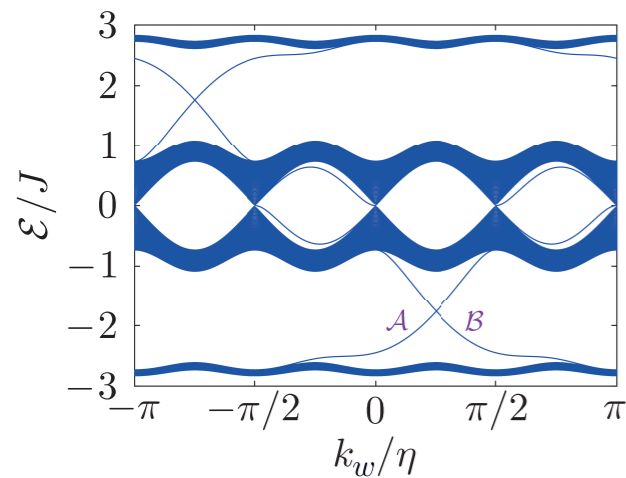
**Carusotto & Ciuit,
RMP, 85, 299, (2013)**

Topological edge states



Excite edge states:

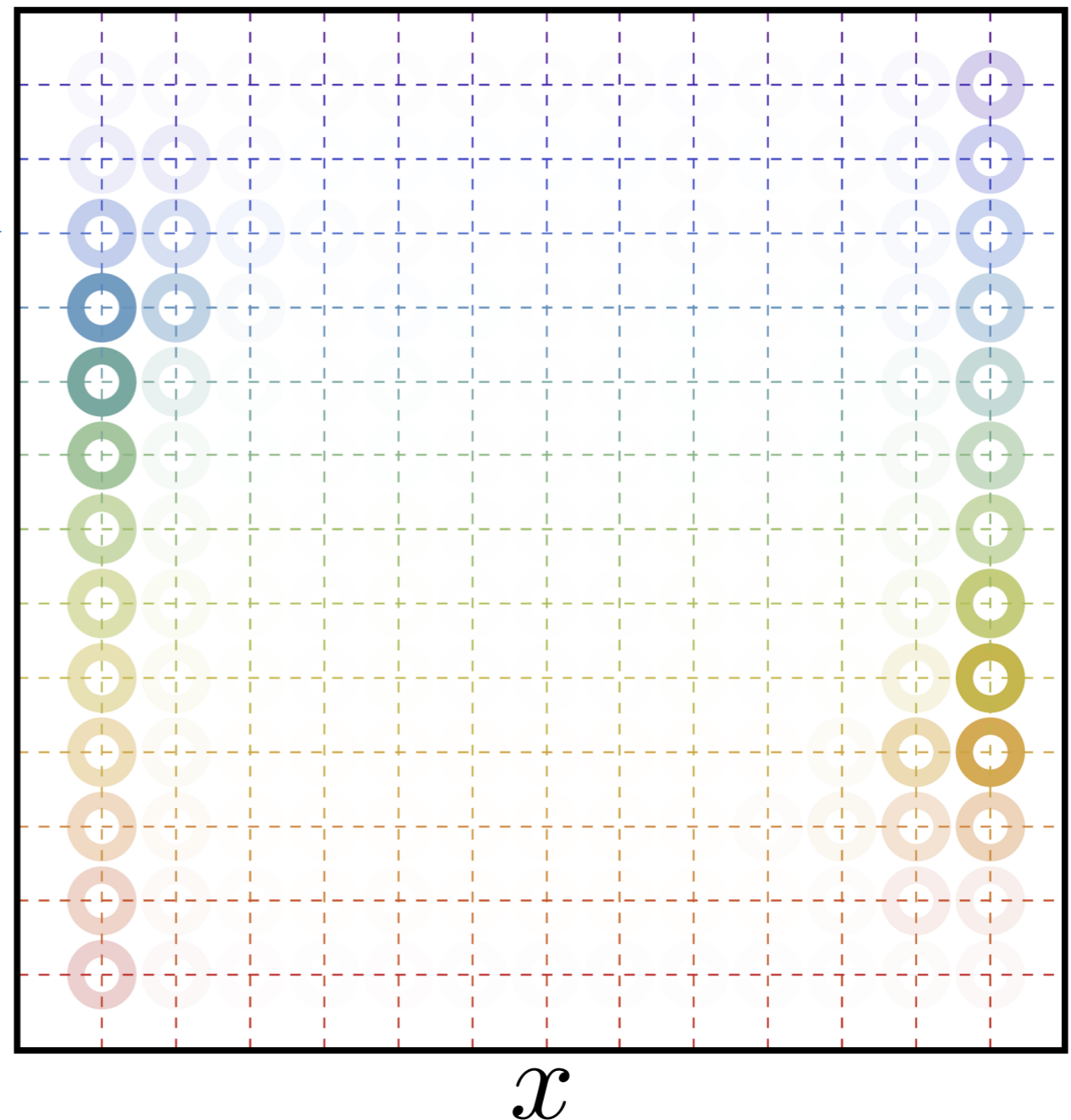
$$\omega_{\text{drive}} = -2J$$



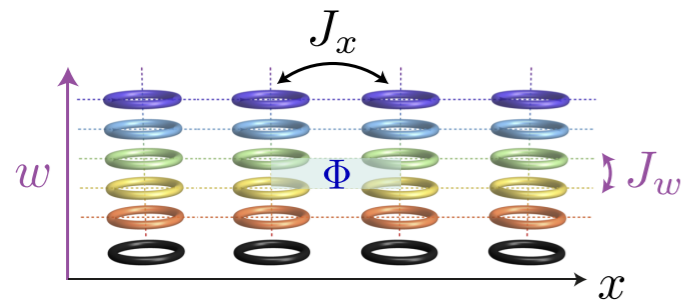
with uniform loss rate:

$$\gamma = 0.1J$$

(a)

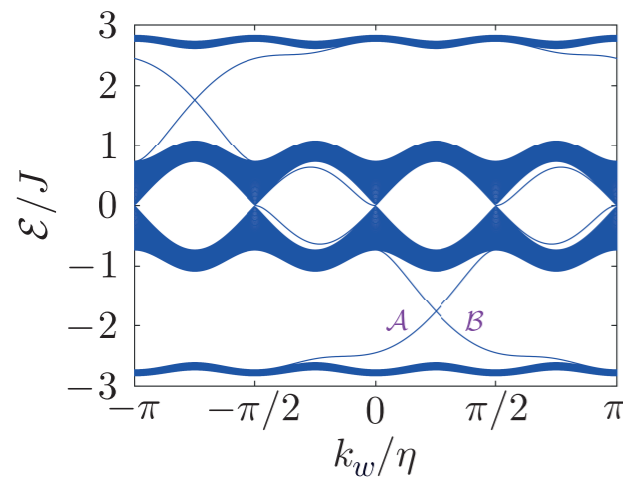


Topological edge states



Excite edge states:

$$\omega_{\text{drive}} = -2J$$

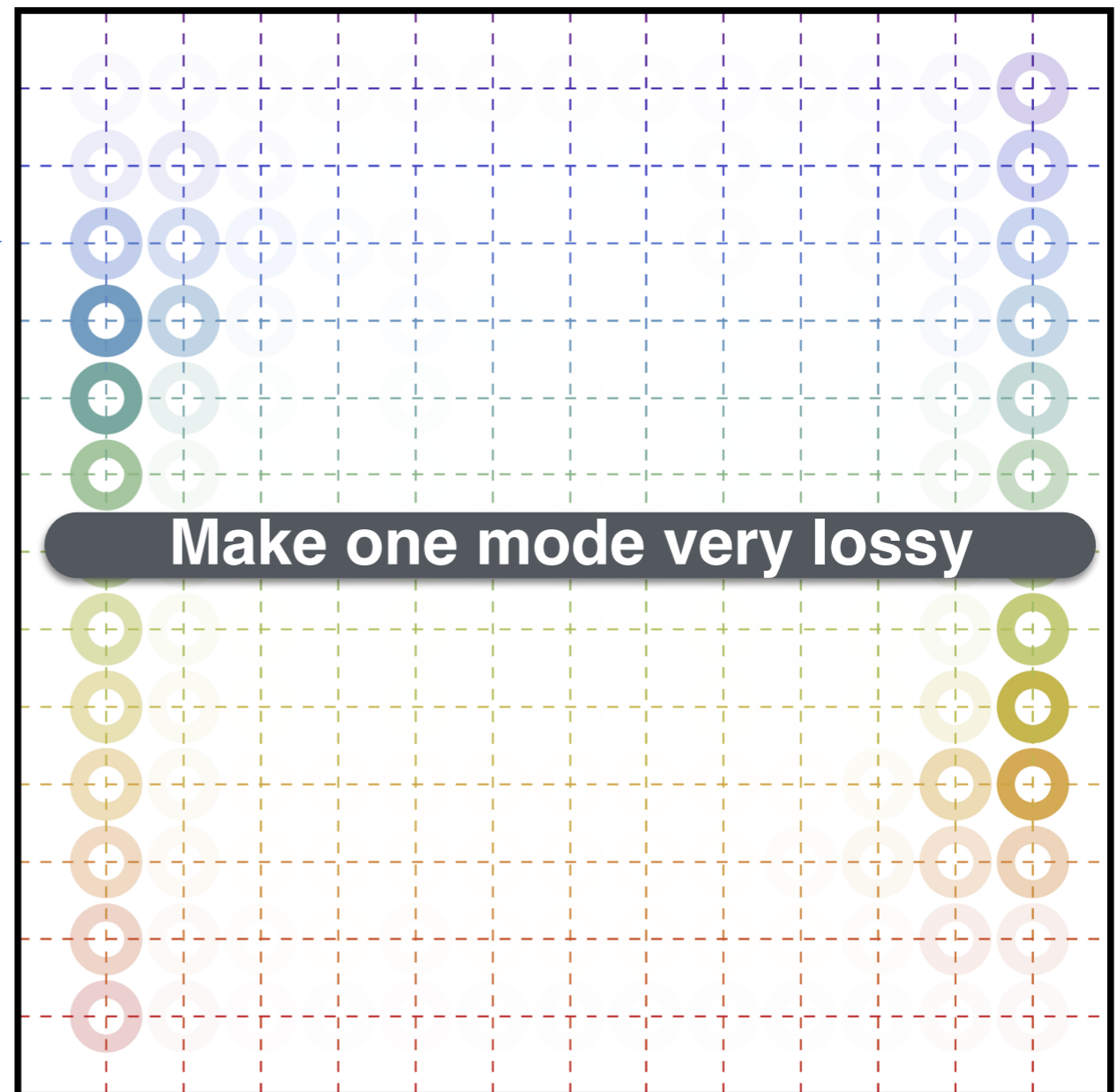


with uniform loss rate:

$$\gamma = 0.1J$$

How to make an effective boundary in the synthetic dimension?

(a)



the Zeno effect

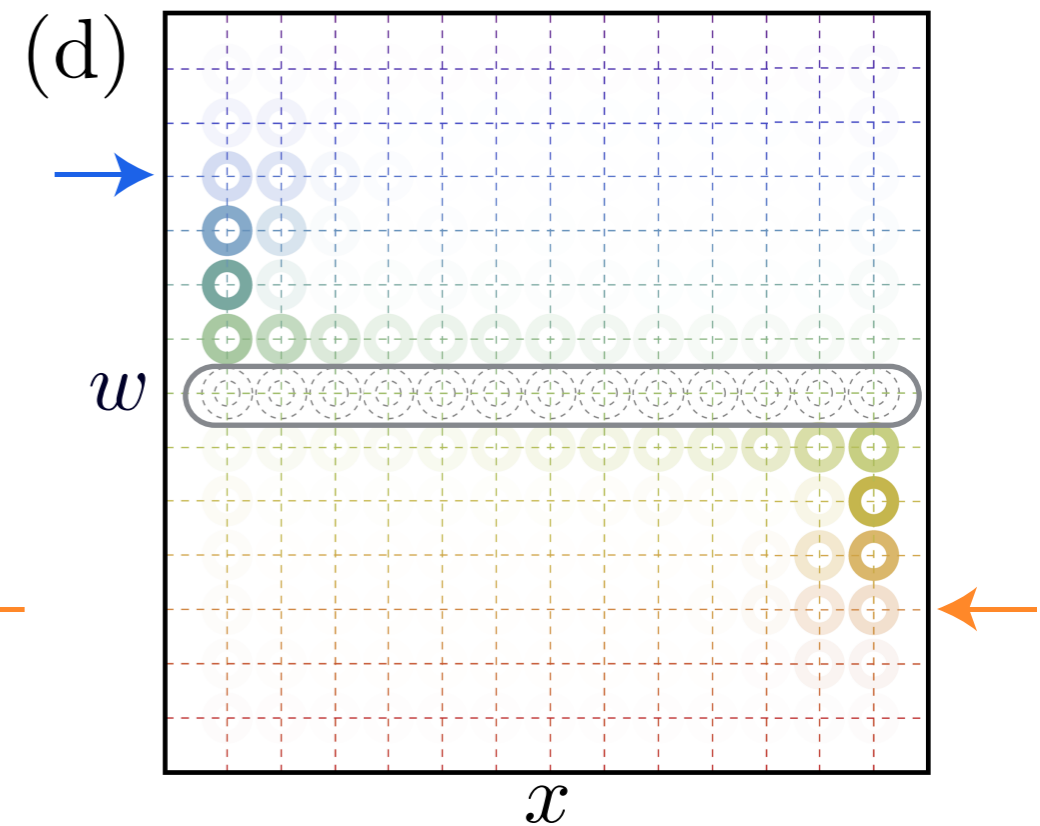
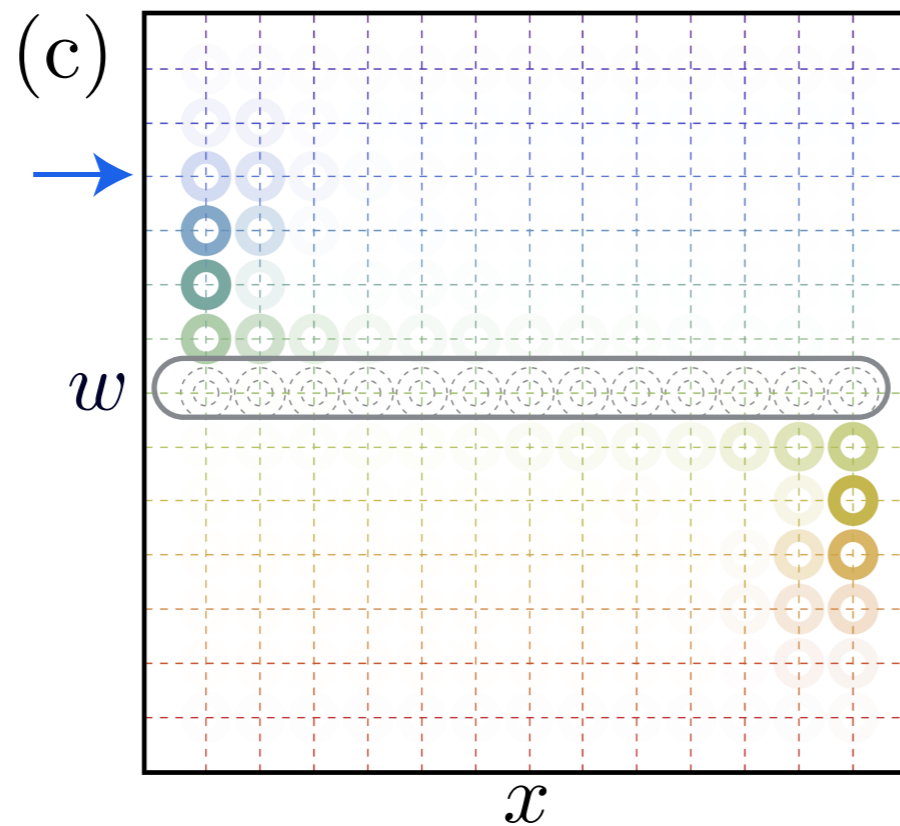
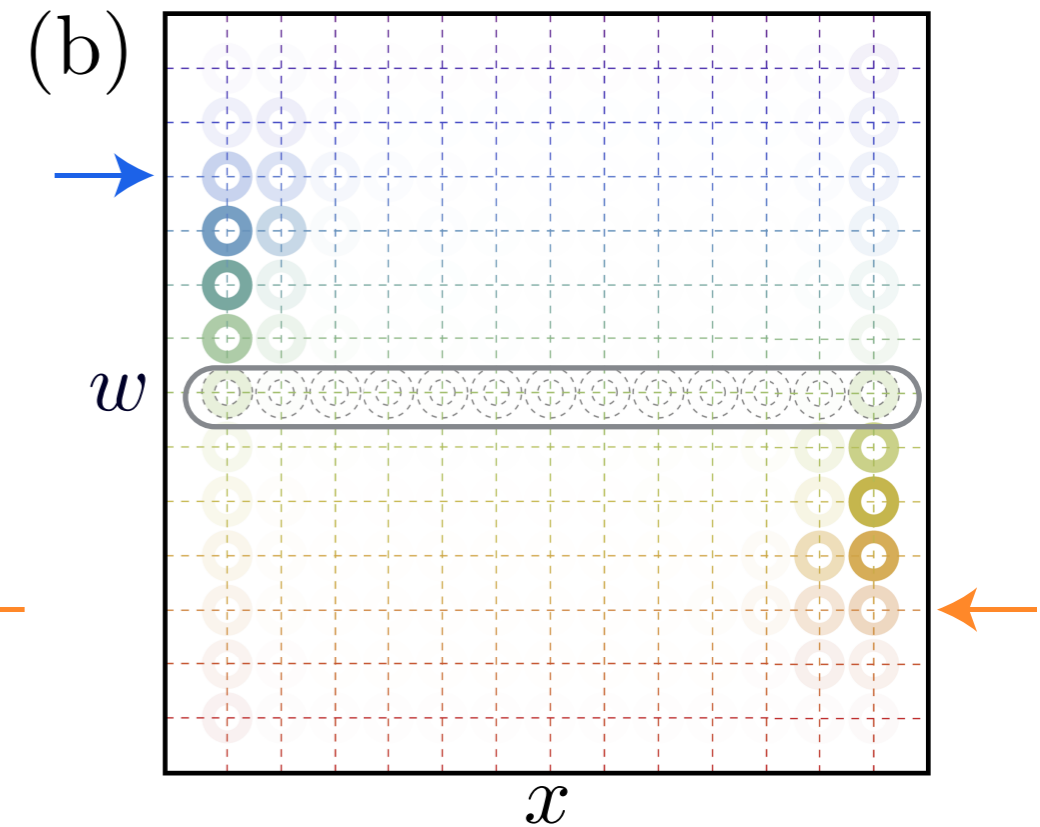
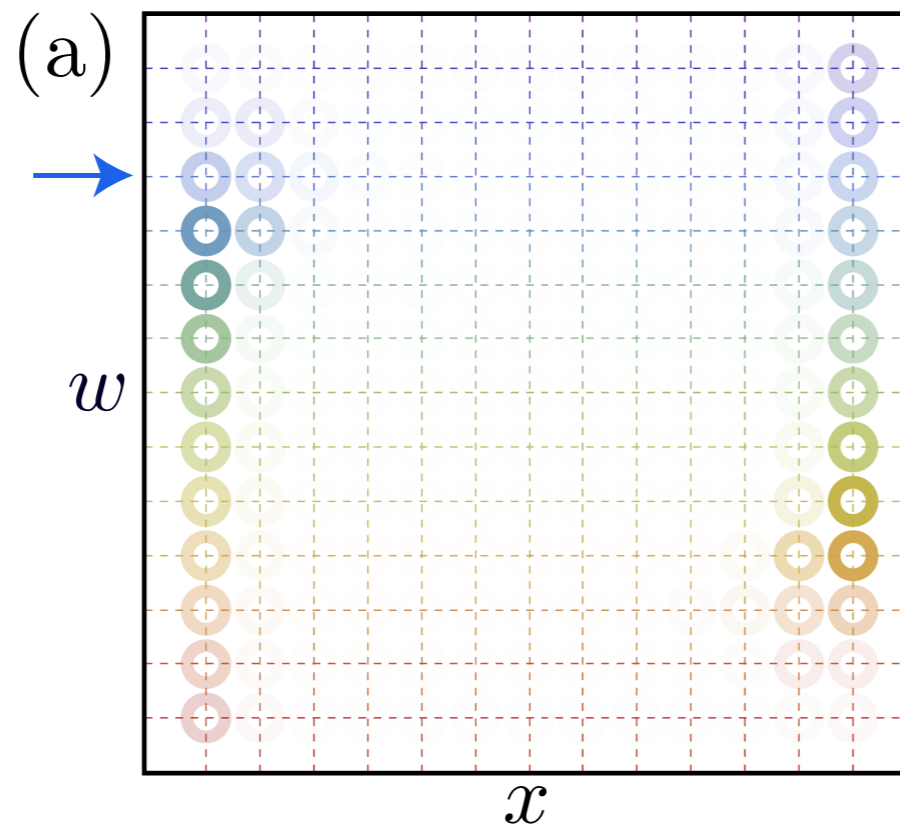
e.g. Barontini et al., PRL, 110, 035302 (2014)

x

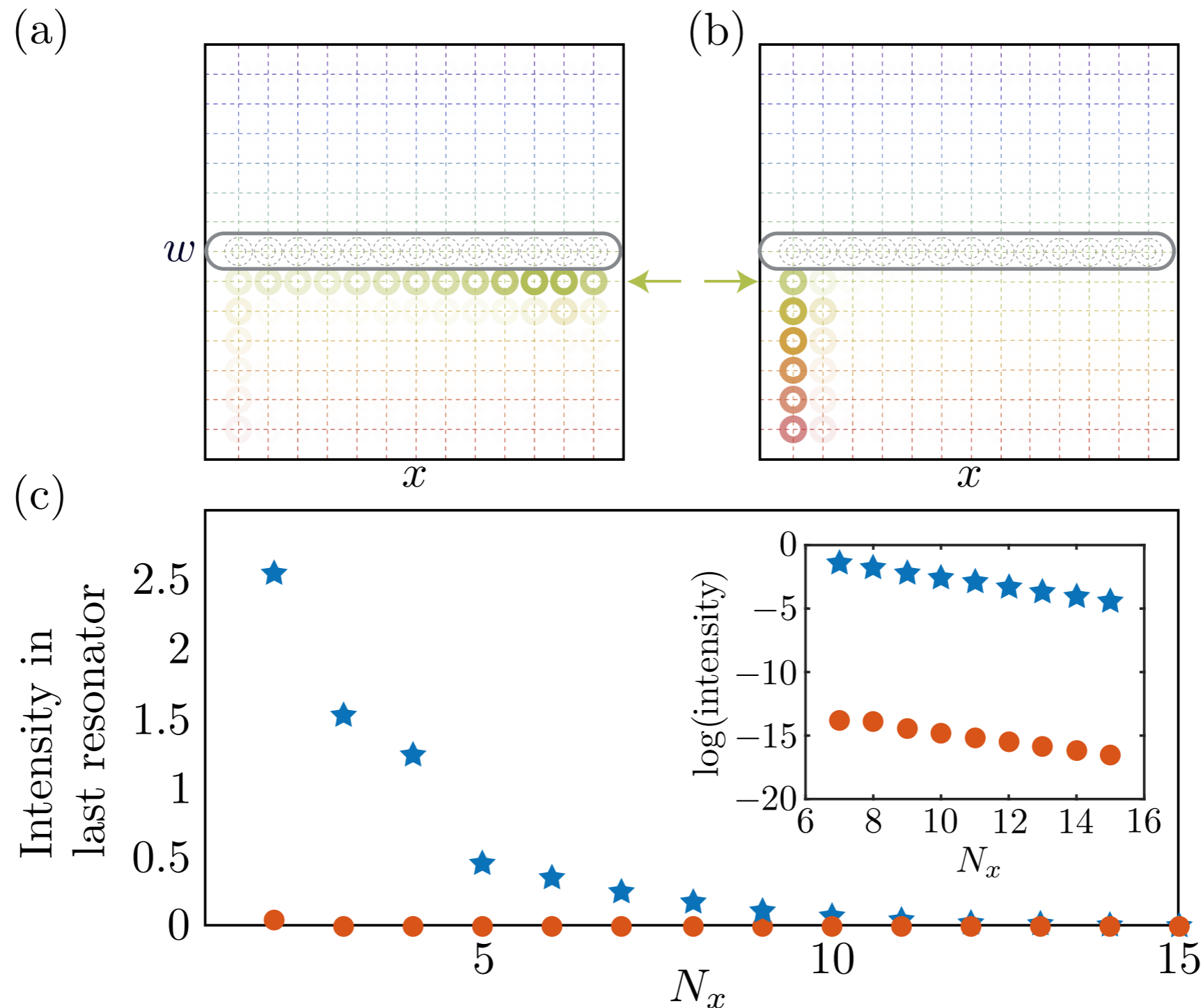
Topological edge states

Losses in the central row equal to:

- (a) γ
- (b) 10γ
- (c) 100γ
- (d) 1000γ



Topological edge states for optical isolation



- Extension of works on spatiotemporal modulation to include topological protection

Yu & Fan, Nat. Photon., 3, 91, (2009)
Lira et al., PRL 109, 033901 (2012)
Tzuan et al., Nat. Photon., 8, 701, (2014)

2D quantum Hall effect in photon transport in 1D chain

2D quantum Hall for a filled band of fermions:

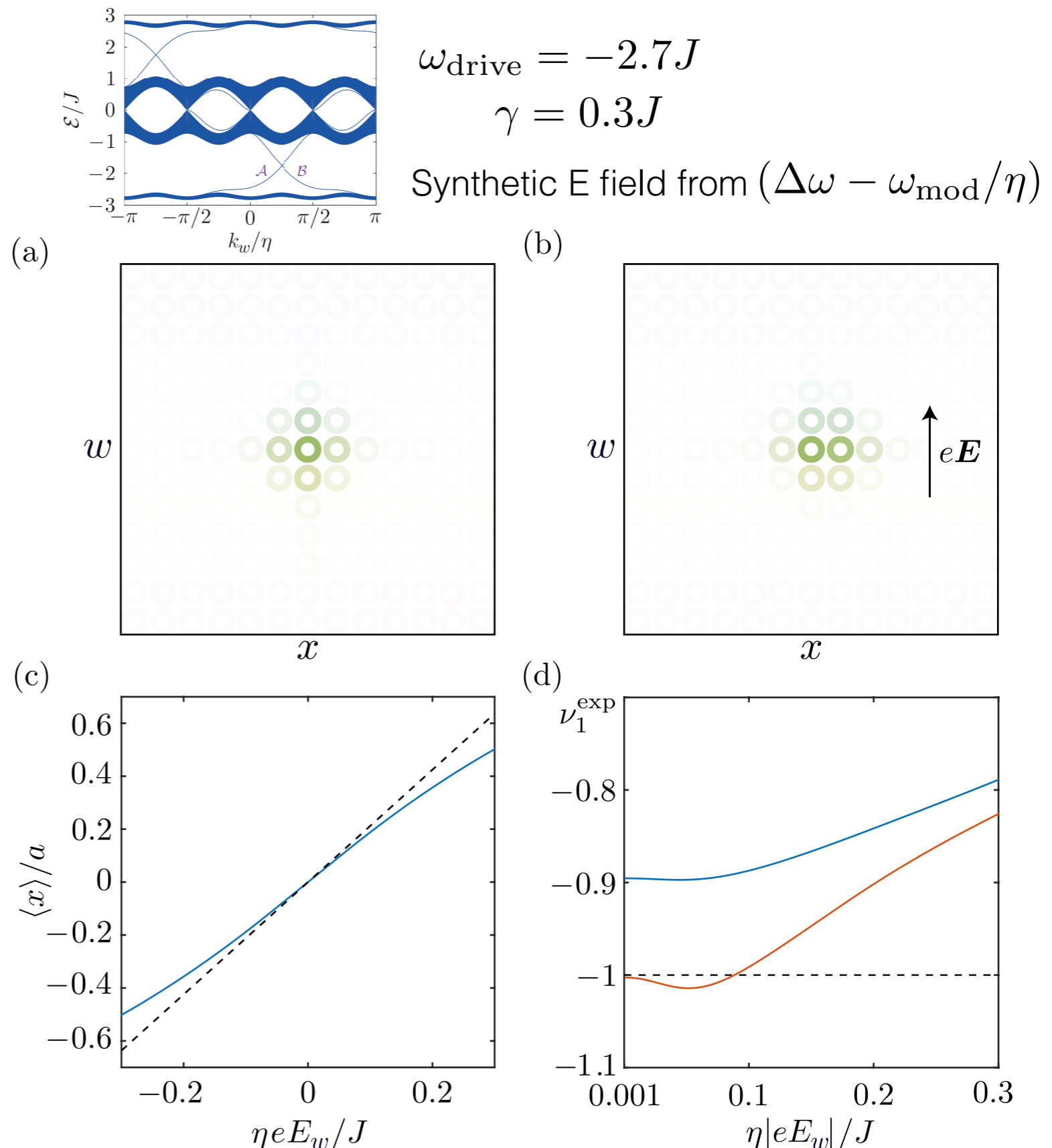
$$j^x = -\frac{eE_w}{2\pi} \nu_1^{xw}$$

Optical bosonic analogue when drive on resonance with a band and $\Delta\mathcal{E}_{\text{band}} < \gamma < \Delta\mathcal{E}_{\text{gap}}$

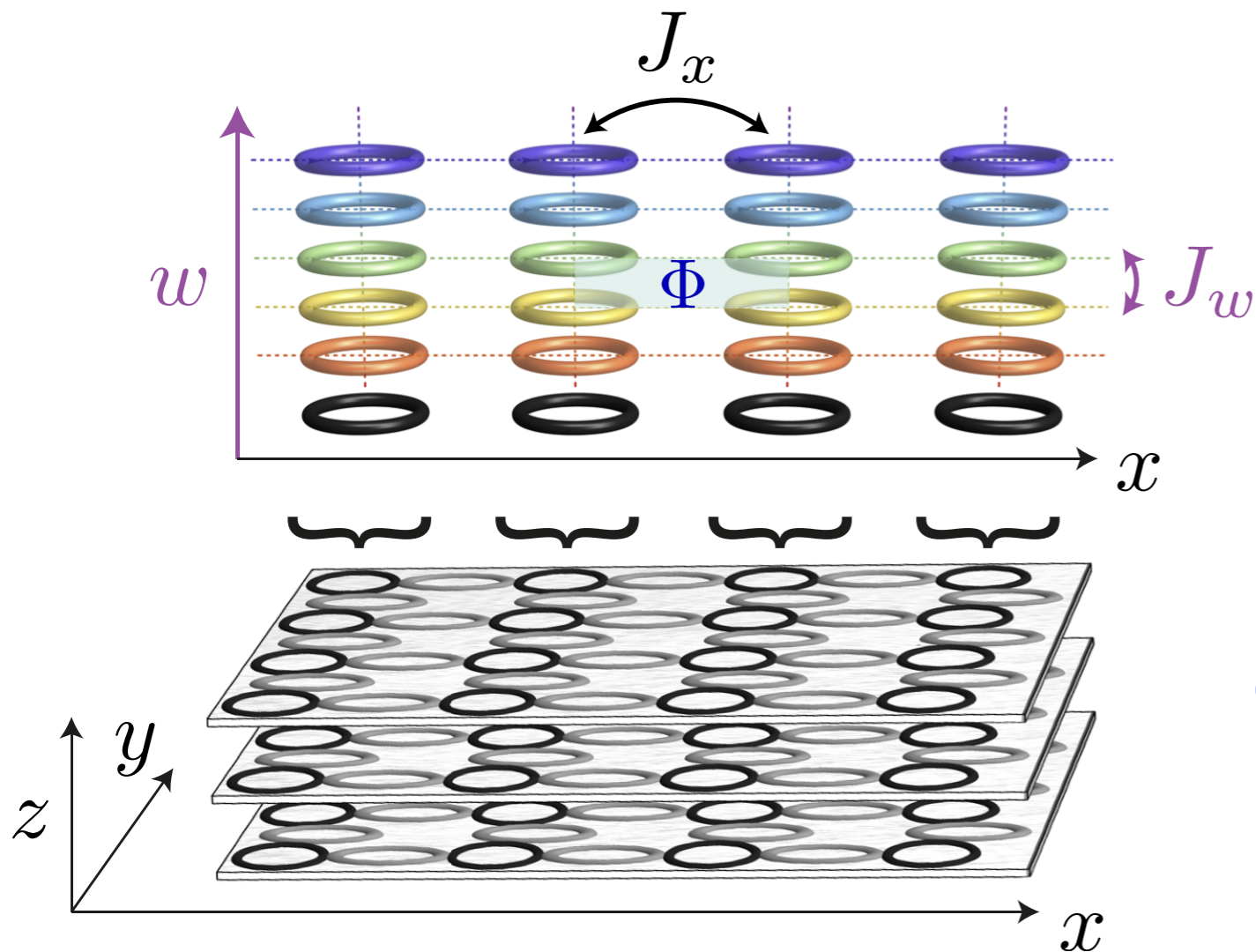
$$\langle x \rangle = -q(\eta a) \frac{eE_w}{2\pi\gamma} \nu_1^{xw} + \mathcal{O}(\gamma^0),$$

for $\Phi = p/q$

**Ozawa & Carusotto,
PRL, 112, 133902, (2014)**



3D resonator array with the synthetic dimension



**What about 4D
quantum Hall physics?**

S.-C.Zhang & J. Hu, *Science* 294, 823, (2001),
Qi, Hughes & Zhang, *PRB* 78, 195424, (2008)...

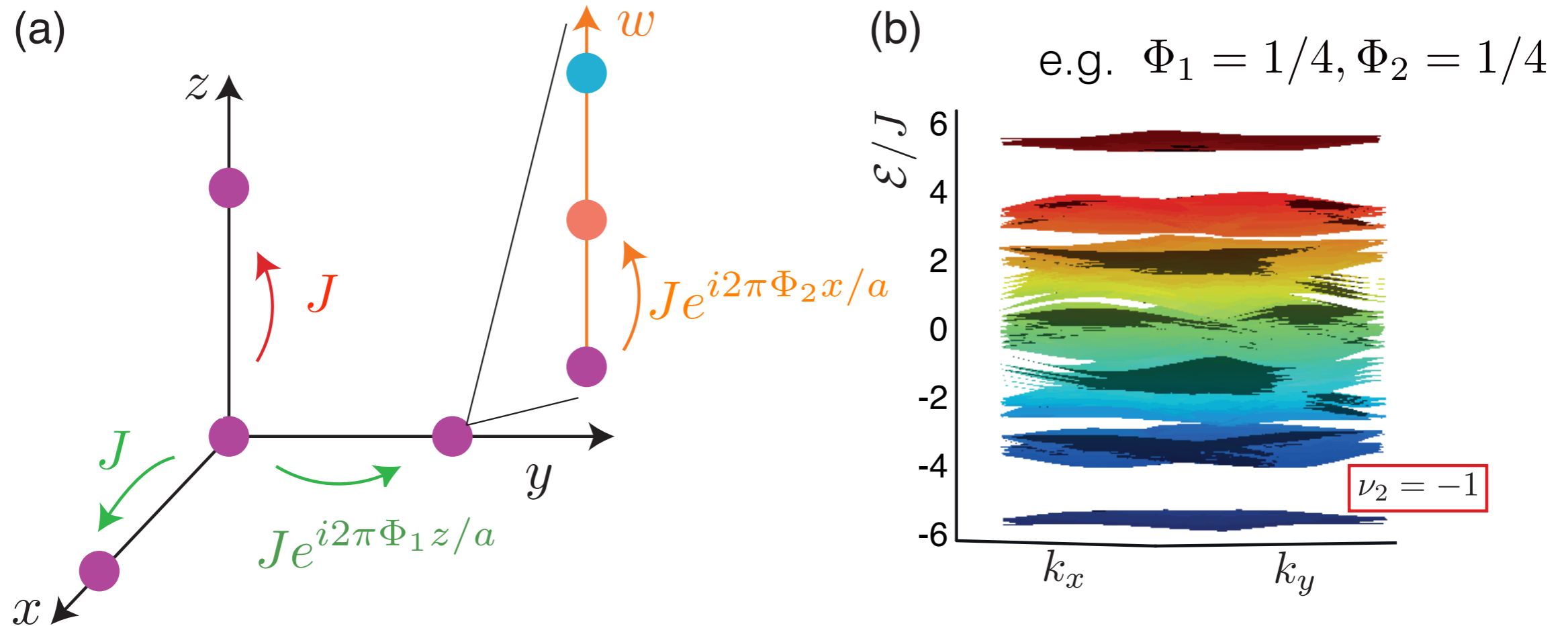
What do we need for interesting 4D topological physics?

$$\nu_2 \equiv \frac{1}{(2\pi)^2} \int_{\text{BZ}} (\mathcal{F}^{xy} \mathcal{F}^{zw} + \mathcal{F}^{wx} \mathcal{F}^{zy} + \mathcal{F}^{zx} \mathcal{F}^{yw}) d^4k \in \mathbb{Z},$$

*topological
2nd Chern number
for TRS-breaking model*

Minimal 4D topological lattice model

Two copies of Hofstadter models in “disconnected” planes



$$H = -J \sum_{\mathbf{r}, w} \hat{b}_{\mathbf{r}+a\mathbf{e}_x, w}^\dagger \hat{b}_{\mathbf{r}, w} + \hat{b}_{\mathbf{r}+a\mathbf{e}_z, w}^\dagger \hat{b}_{\mathbf{r}, w} \\ + e^{i2\pi\Phi_1 z/a} \hat{b}_{\mathbf{r}+a\mathbf{e}_y, w}^\dagger \hat{b}_{\mathbf{r}, w} + e^{i2\pi\Phi_2 x/a} \hat{b}_{\mathbf{r}+\eta\mathbf{e}_w, w}^\dagger \hat{b}_{\mathbf{r}, w} + \text{h.c.}$$

4D quantum Hall effect in photon transport

4D quantum Hall for a filled (non-degenerate) band of fermions:

perturbing magnetic field

$$\delta B_{\rho\sigma} \equiv \partial_\rho A_\sigma - \partial_\sigma A_\rho$$

$$j^\mu = -eE_\nu \int_{\text{BZ}} \frac{d^4 k}{(2\pi)^4} \mathcal{F}^{\mu\nu}(\mathbf{k}) + e^2 E_\nu \delta B_{\rho\sigma} \frac{\nu_2}{(2\pi)^2} \epsilon^{\mu\nu\rho\sigma},$$

essentially 2D quantum
Hall physics

the genuine 4D quantum
Hall physics

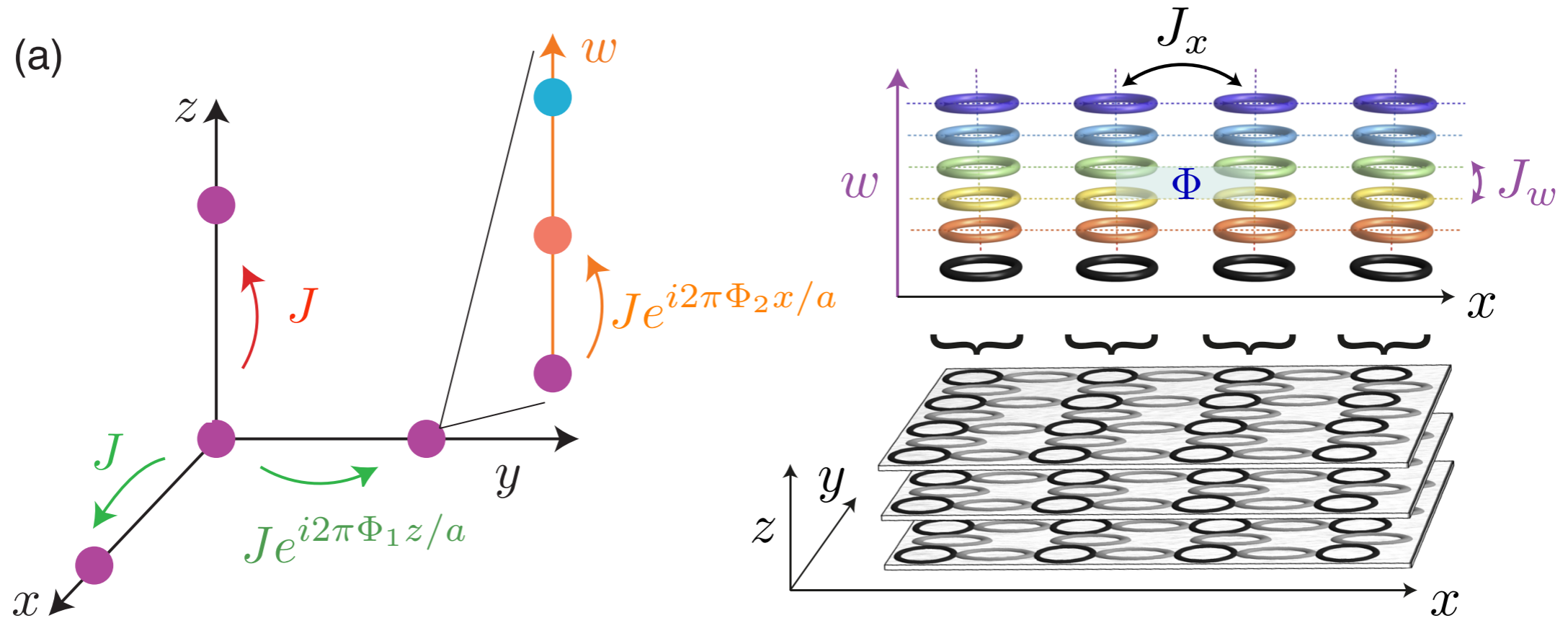
Generalising 2D optical analogue effect to 4D, again when on resonance and when $\Delta\mathcal{E}_{\text{band}} < \gamma < \Delta\mathcal{E}_{\text{gap}}$

$$\langle r^\mu \rangle = \frac{-eE_\nu \int_{\text{BZ}} \frac{d^4 k}{(2\pi)^4} \Omega^{\mu\nu}(\mathbf{k}) + e^2 E_\nu \delta B_{\rho\sigma} \frac{\nu_2}{(2\pi)^2} \epsilon^{\mu\nu\rho\sigma}}{\gamma \int_{\text{BZ}} d^4 k D(\mathbf{k})}.$$

Ozawa et al
arXiv:1510.03910

for case discussed here, modified density of states is simply $D(\mathbf{k}) = 1/(2\pi)^4$

4DQH in 3D array with a synthetic dimension



1. Synthetic flux in xw plane: as for the 1D chain of resonators through $\Omega_{\mathbf{r}}^o = |\Omega_{\mathbf{r}}^o| e^{ik_x x}$

2. Synthetic flux in yz plane: non-resonant links with a z -dependent displacement

like in Hafezi et al.,
Nat. Photon. 7, 1001, (2013)

3. Synthetic perturbing electric field E_x : e.g. let cavity size vary uniformly

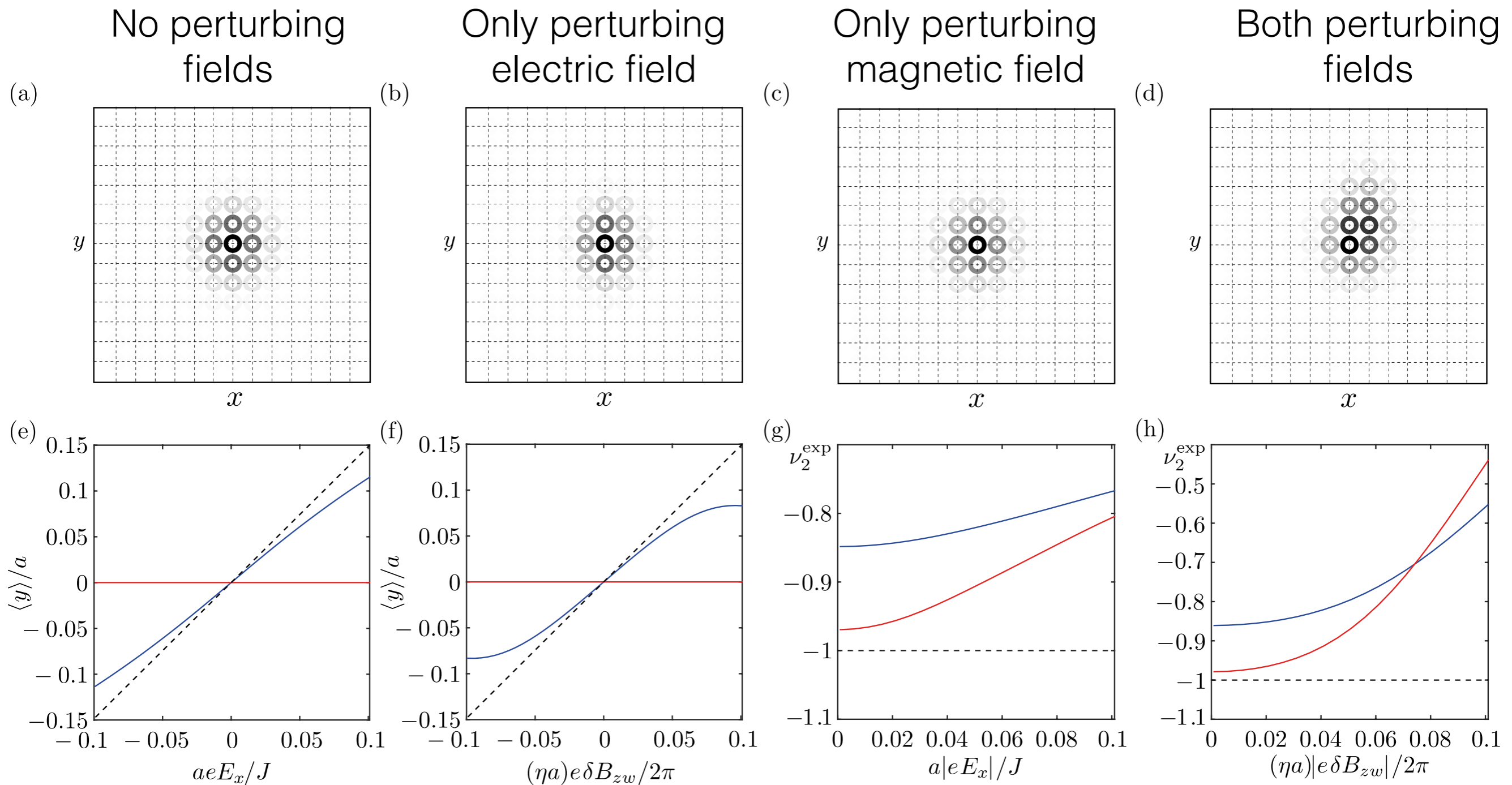
4. Synthetic perturbing magnetic field δB_{zw} : additional slow z -dependence

$$\Omega_{\mathbf{r}}^o = |\Omega_{\mathbf{r}}^o| e^{ik_x x} e^{ik_x z}$$

4DQH in 3D array with a synthetic dimension

$$\langle y \rangle = -(\eta a^3) e^2 \frac{q^{yz} q^{xw}}{\gamma} \frac{\nu_2}{(2\pi)^2} E_x \delta B_{zw} + \mathcal{O}(\gamma^0),$$

$$\begin{aligned} \omega_{\text{drive}} &= -6.3J \\ \nu_2 &= -1 \end{aligned}$$



In summary



- Different modes of a silicon ring-resonator can be exploited as an extra synthetic dimension for photons
- Topological physics in a **1D** ring-resonator chain:
 - Optical isolation with topological edge states
 - **2D** quantum Hall effect
- Topological physics in a in a **3D** resonator array:
 - **4D** quantum Hall effect

