

# Towards a First Principles Understanding of Black Hole Variability

Chris Reynolds

Drew Hogg

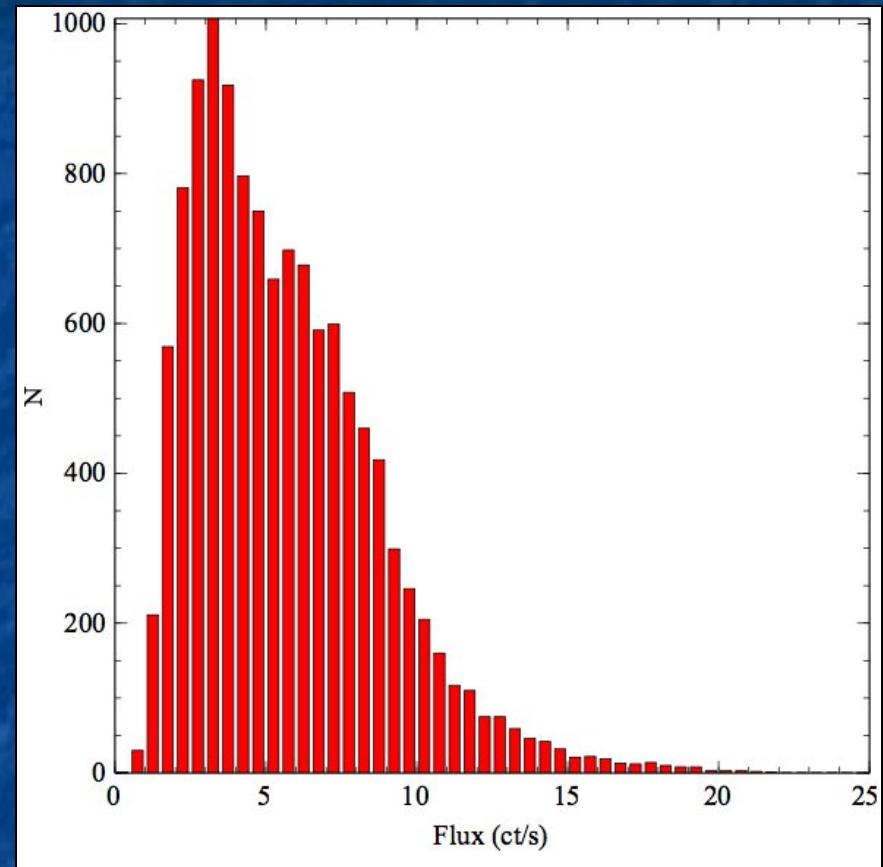
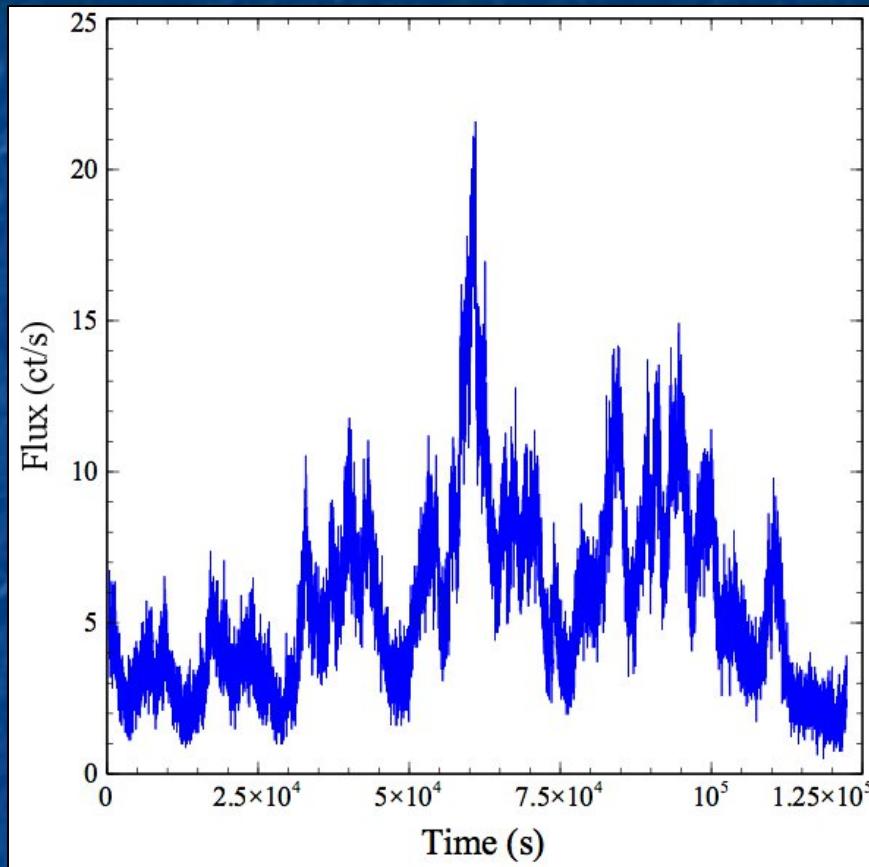
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USA*

# Outline

- The structure of the aperiodic “noise”
  - Propagating fluctuation model
  - Role of dynamo processes
- State transitions
  - The thin-disk/RIAF paradigm
  - First look at structure of truncated disks
- Goal : Use of variability as probe of the underlying accretion physics

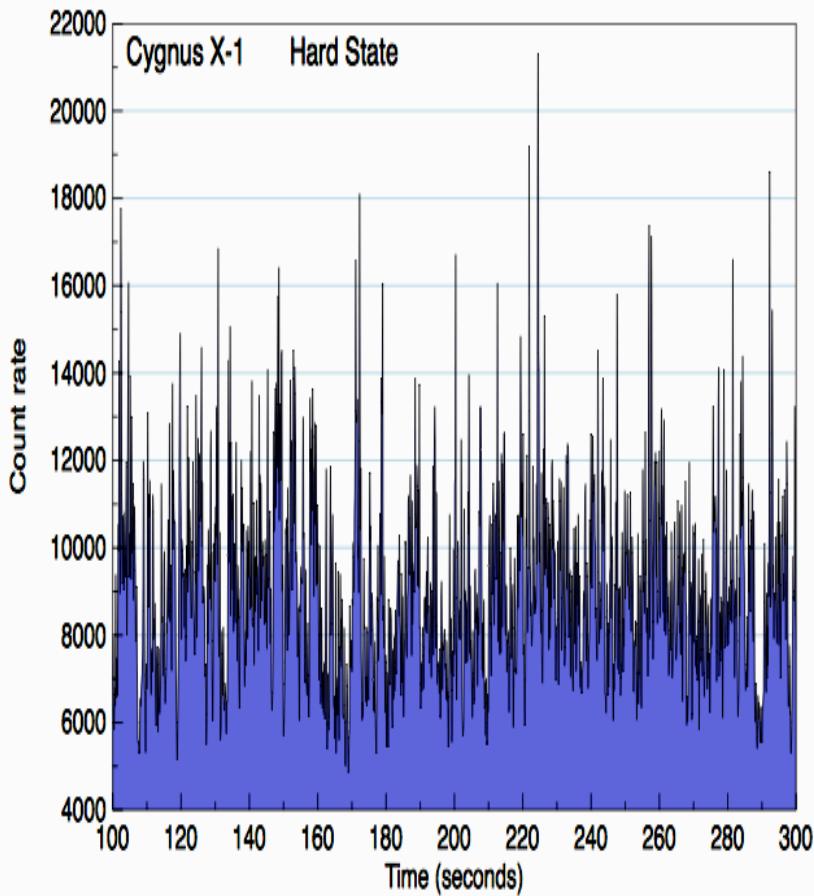
# I : Rapid aperiodic variability

# Seyfert galaxy 1H0707-495 (XMM-Newton)

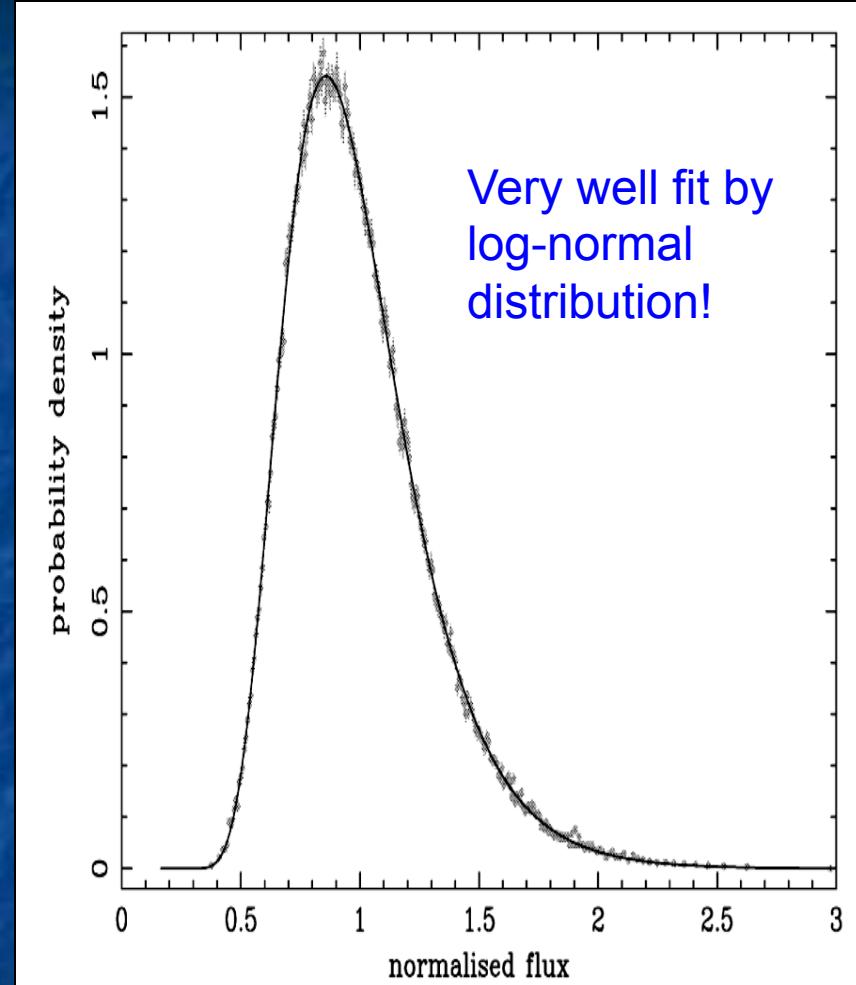


Light curve is “flary”... distribution function of flux measurements skewed to the high end

# Cygnus X-1 (RXTE)

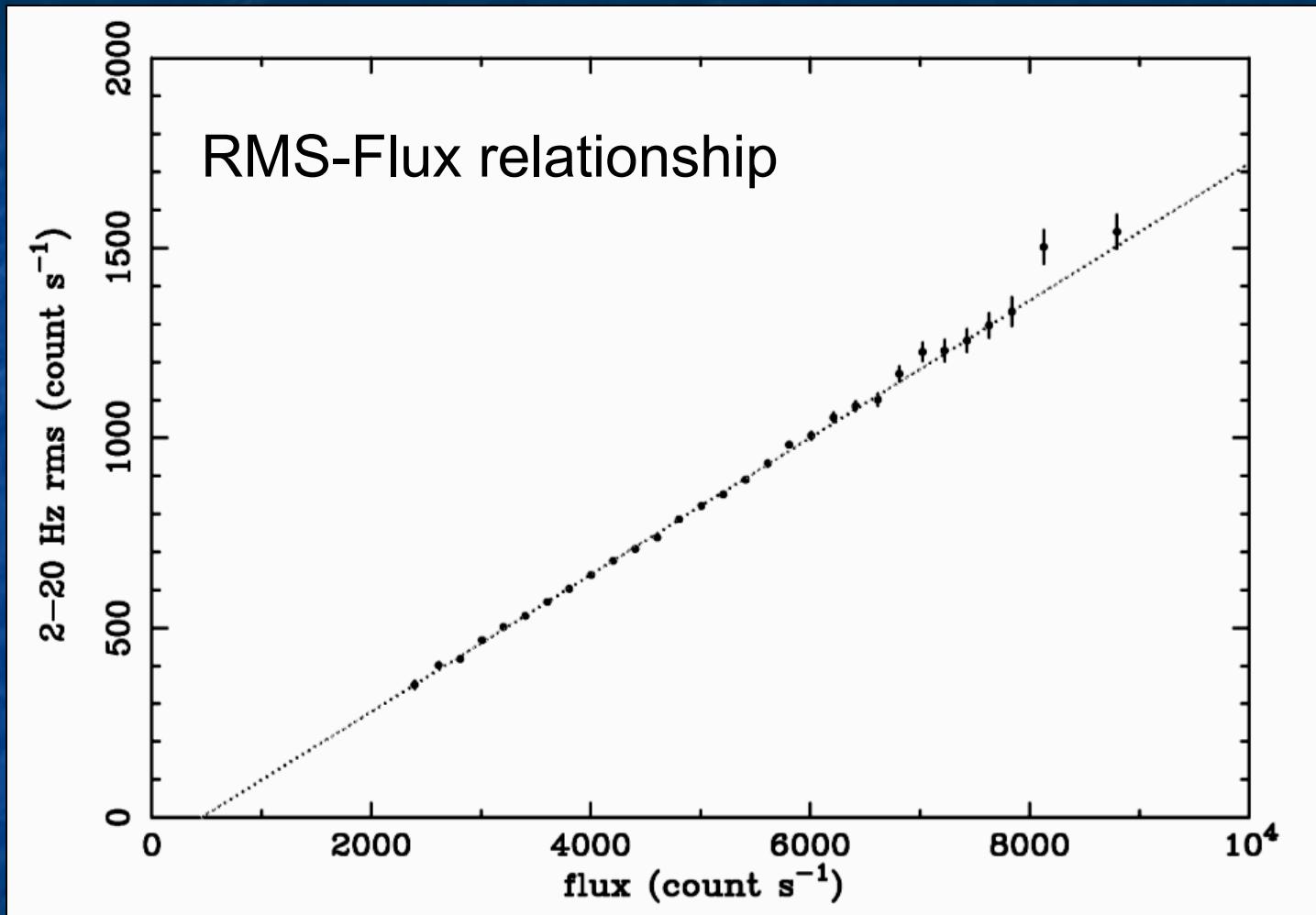


Belloni et al. (2010)



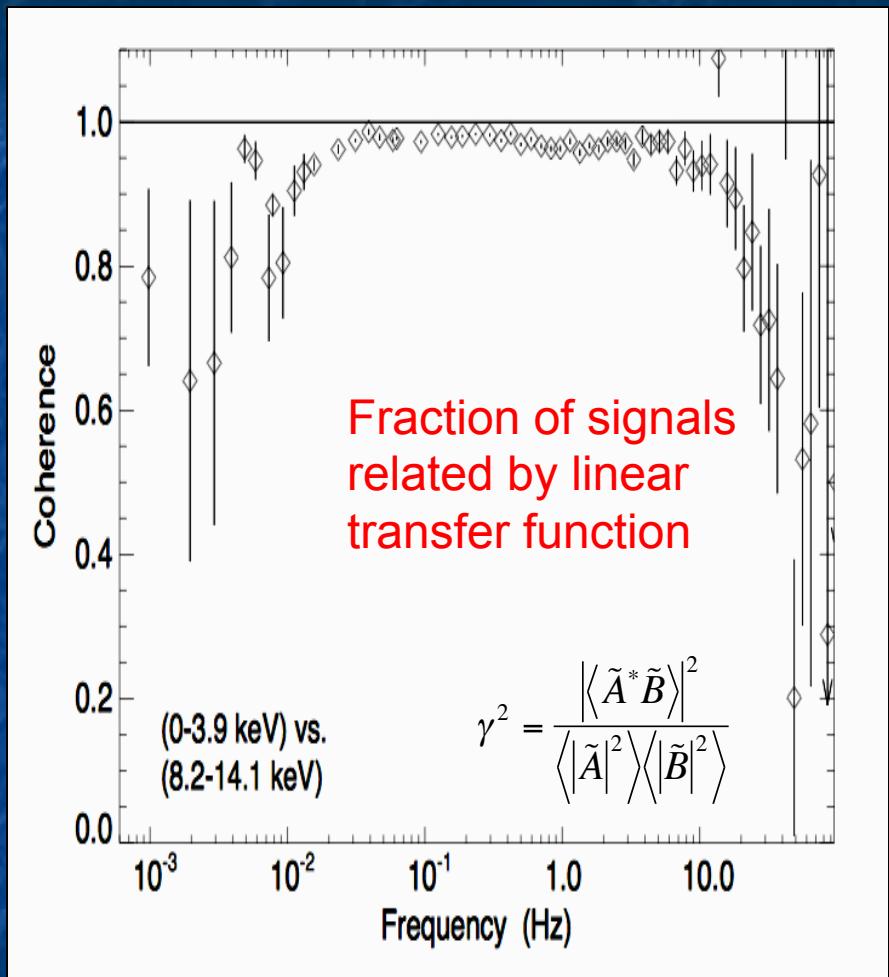
Uttley, McHardy & Vaughan (2005)

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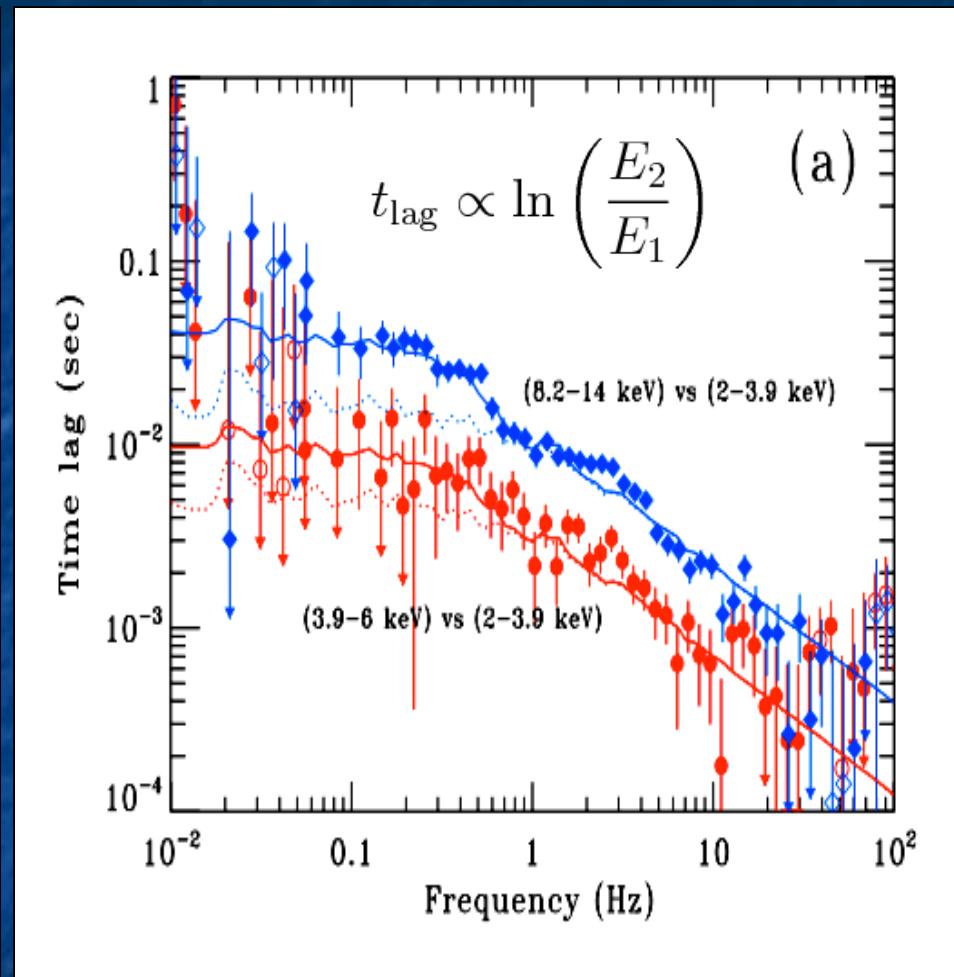


Uttley, McHardy & Vaughan (2005)

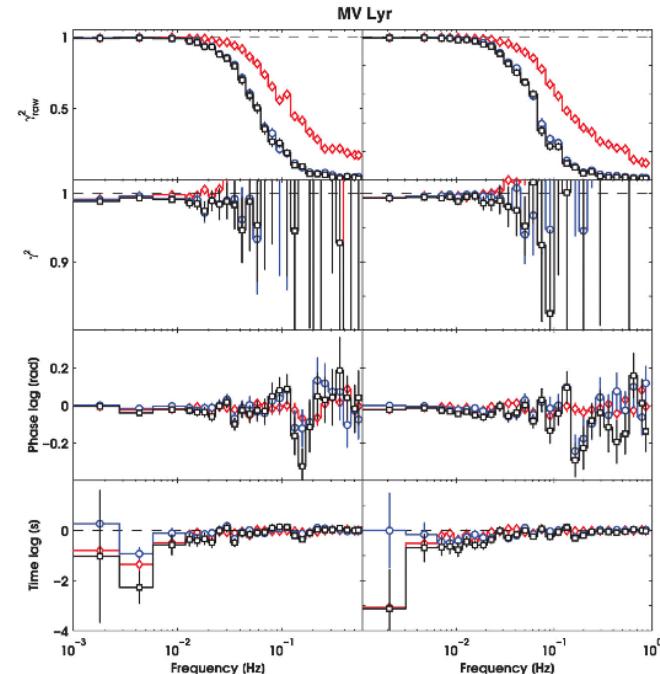
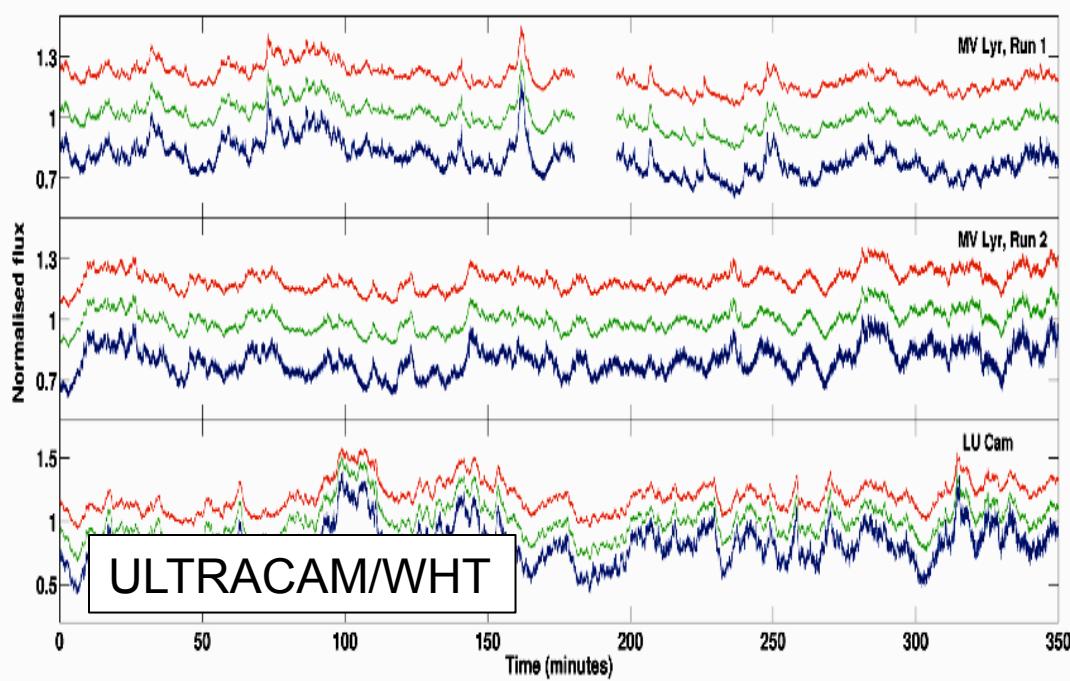
# Cygnus X-1 (RXTE)



Nowak et al. (1999)



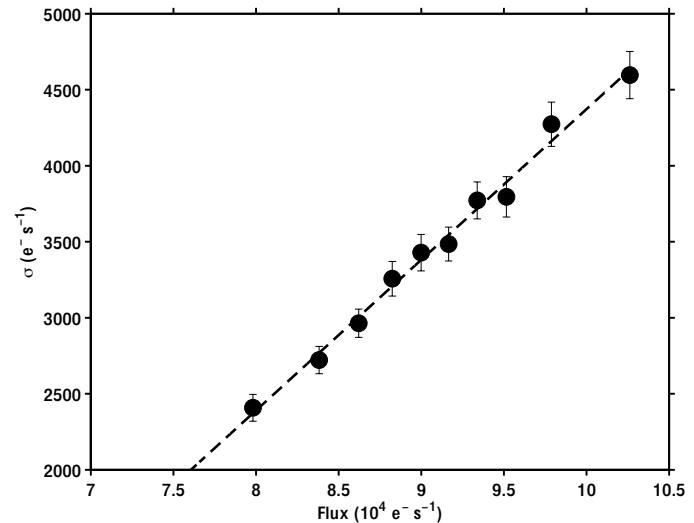
Poutanen (2001)



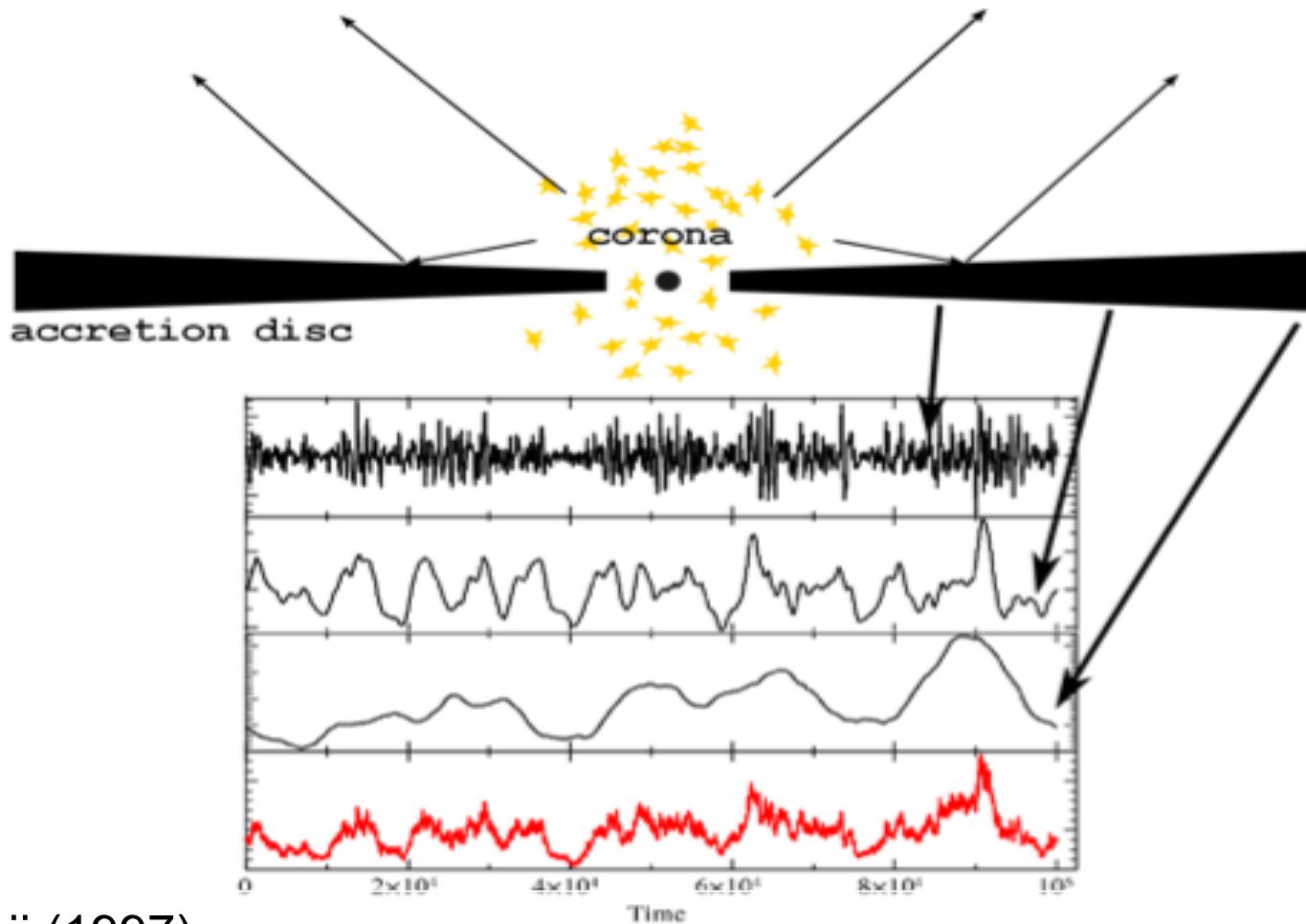
## White dwarfs do it too!

Also show rms-flux relation  
and frequency-dependent  
time-lags

(Scaringi et al. 2013)



Propagating fluctuations model : Inward propagating multiplicative fluctuations caused by stochasticity in the angular momentum transport mechanism



Lyubarskii (1997);

Kotov, Churazov & Gilfanov (2001);

Uttley, McHardy & Vaughan (2005)

# 1d viscous disk model with stochastic viscosity

Lyubarskii (1997) linearized problem → got rid of behavior we seek!

Cowperthwaite & Reynolds (2015):

Integrate non-linear disk evolution equation for surface density  $\Sigma$ :

$$\frac{\partial \Sigma}{\partial t} = \frac{3}{R} \frac{\partial}{\partial R} \left[ R^{1/2} \frac{\partial}{\partial R} (\nu \Sigma R^{1/2}) \right], \quad (\text{Pringle 1981})$$

Treat viscosity as stochastic parameter with characteristic frequency  $\omega_0$  (mimicking turbulent fluctuations):

$$\nu = \nu_0(1 + \beta)$$

$$d\beta(t) = -\omega_0(\beta(t) - \mu)dt + \xi dW(t)$$

driving  
frequency

mean  
(=0)

Normalized  
Gaussian  
white noise

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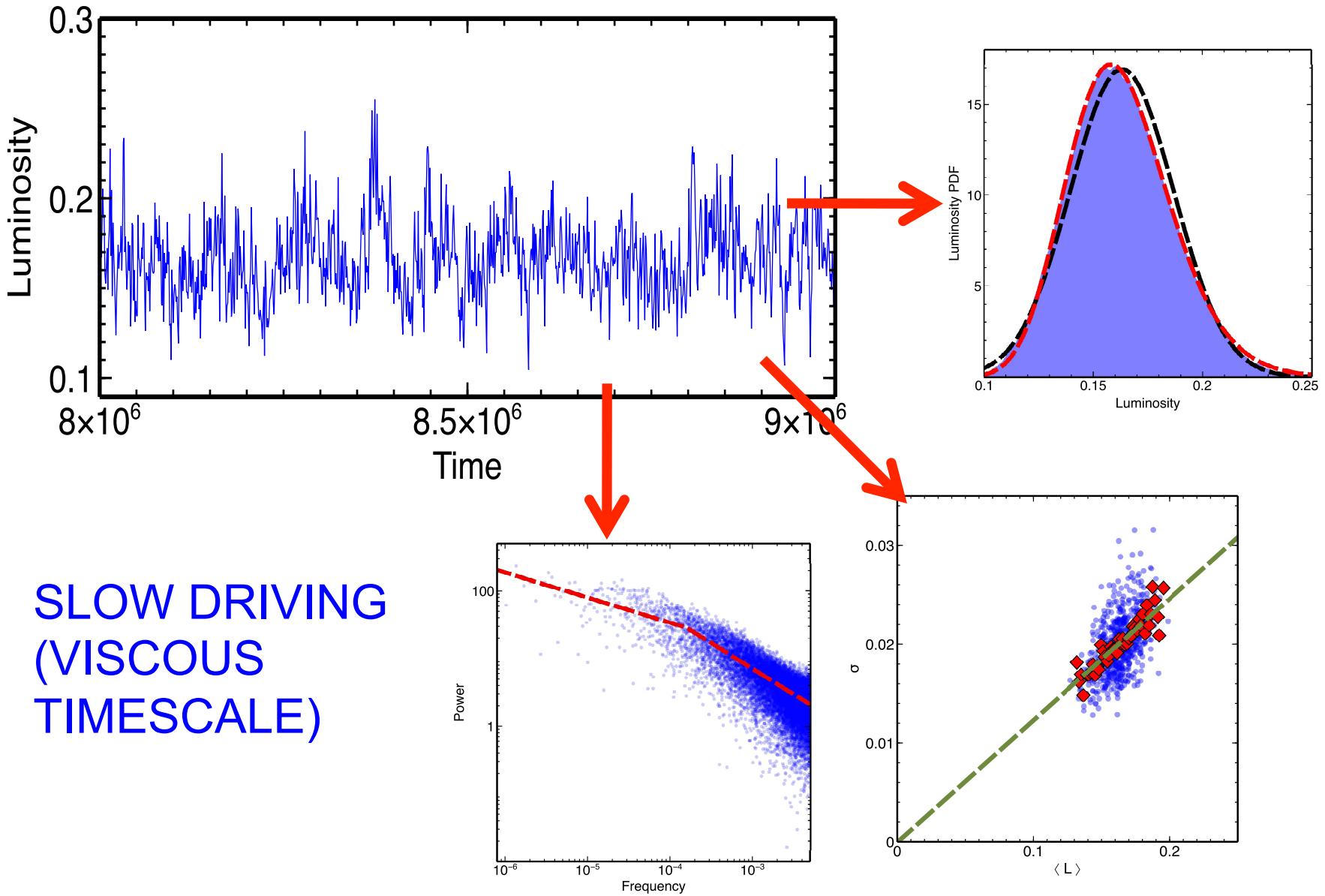
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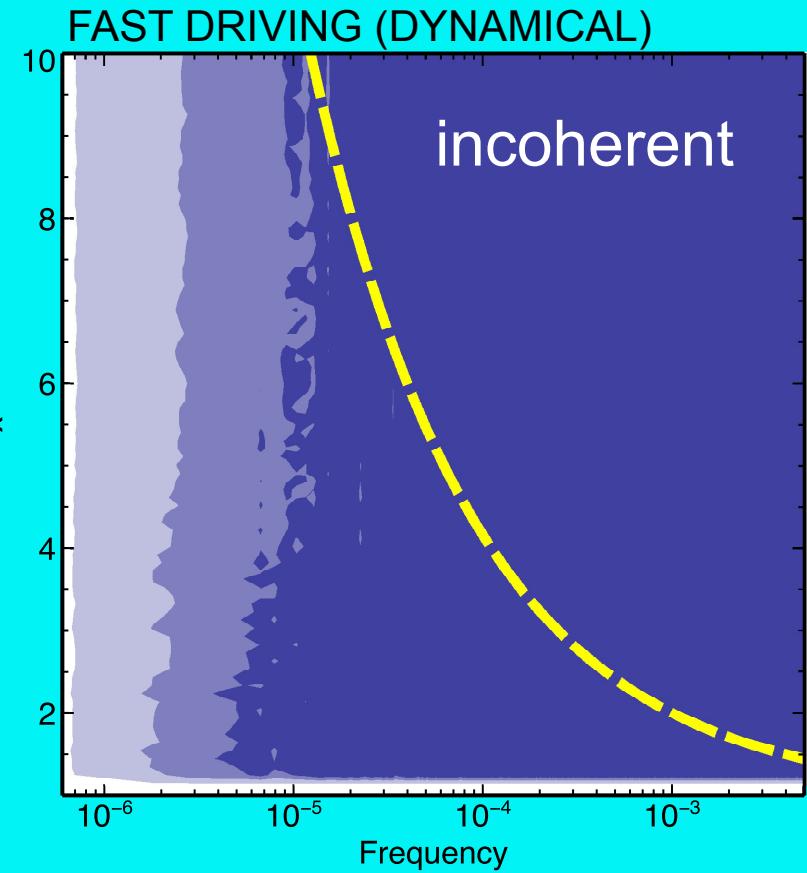
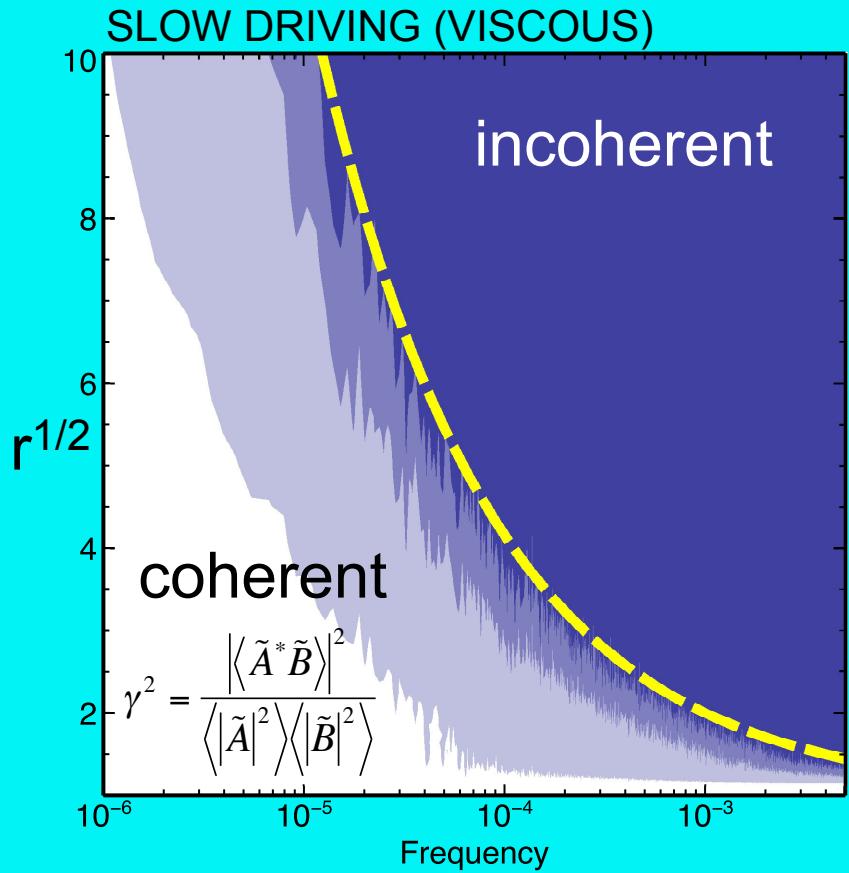
## Findings:

- If fluctuations are slow enough, then recover log-normal light curve, rms-flux relationship, coherence/lags between radii
- But, dynamical timescale fluctuations are too fast... no coherent propagation and no non-linearity in the light curve!



# COHERENCE FUNCTIONS

Cowperthwaite & Reynolds (2015)



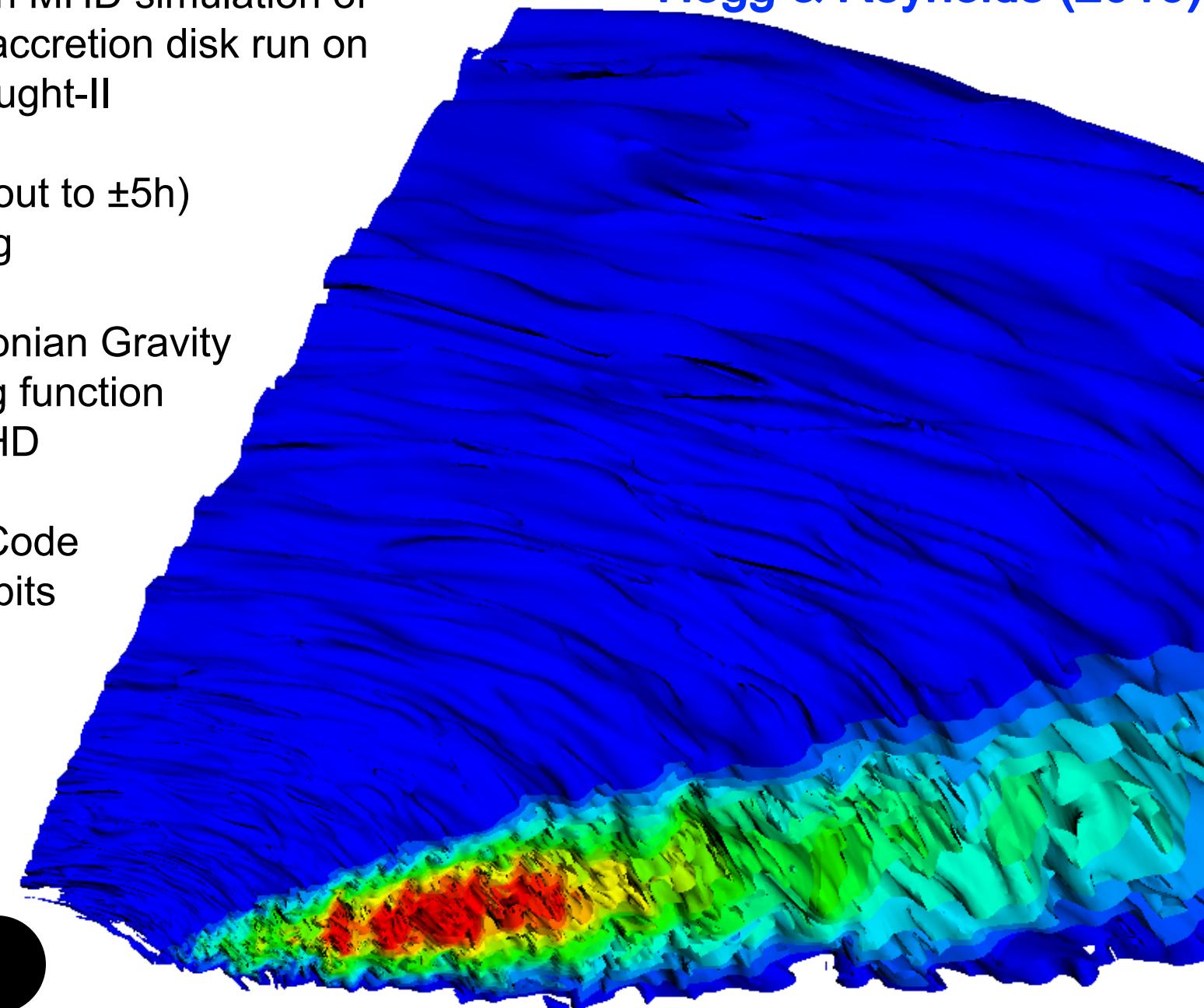
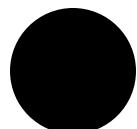
Radii are coherent on timescales longer than viscous timescale

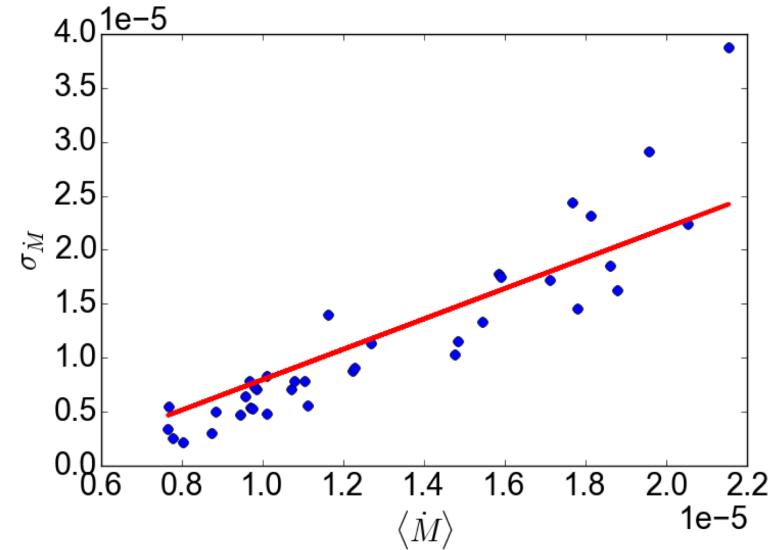
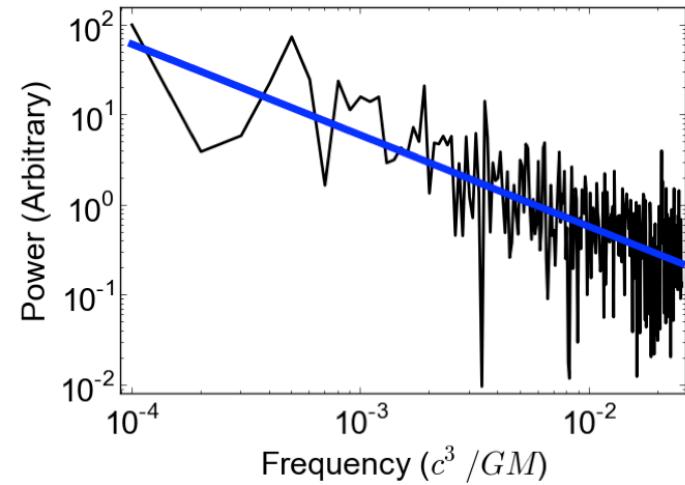
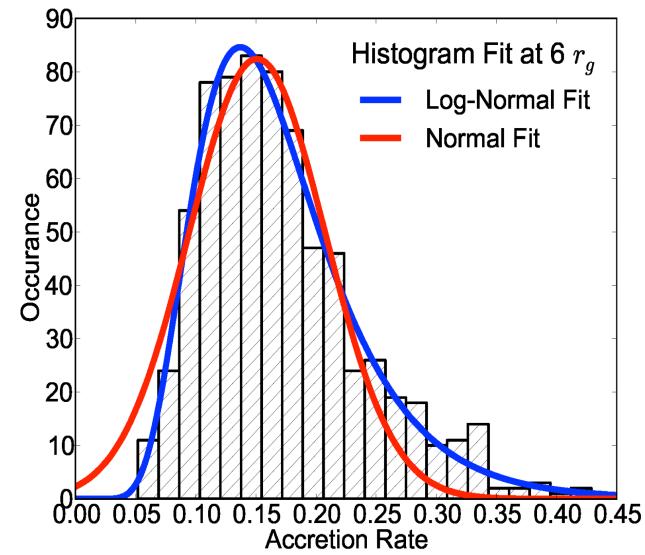
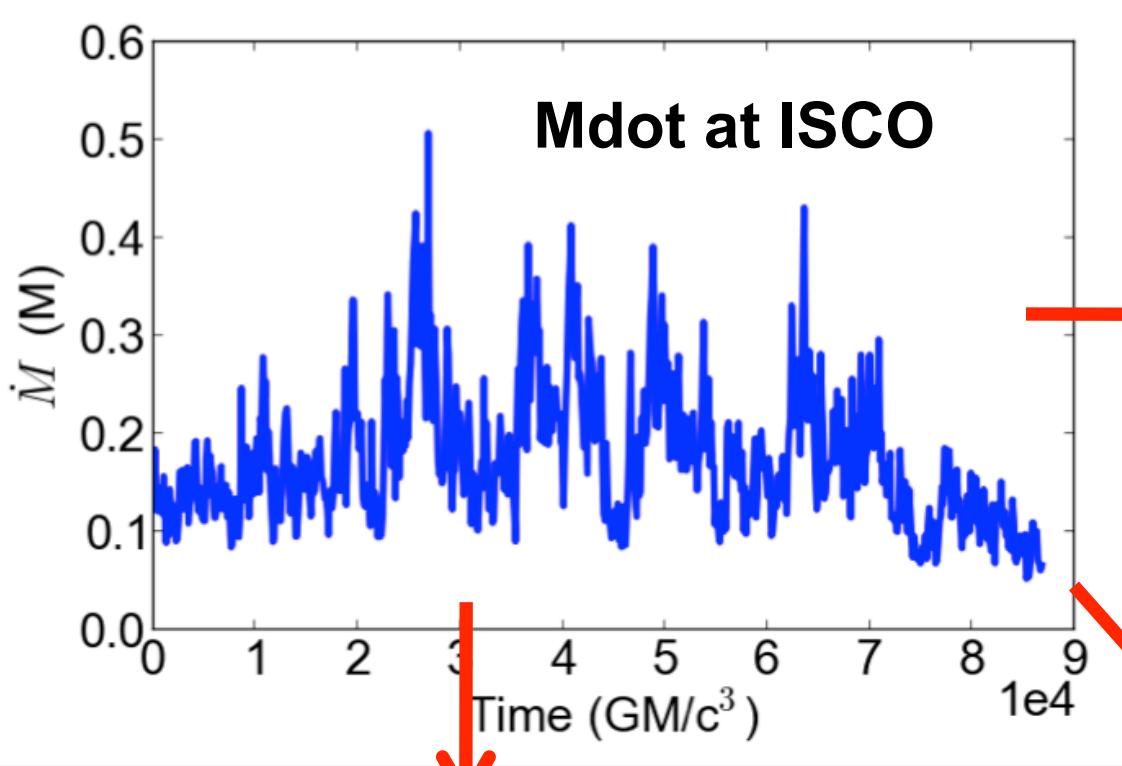
High-resolution MHD simulation of  
thin ( $h/r=0.1$ ) accretion disk run on  
UMd Deepthought-II

30-zones /  $h$  (out to  $\pm 5h$ )  
 $R : 4rg \rightarrow 145rg$

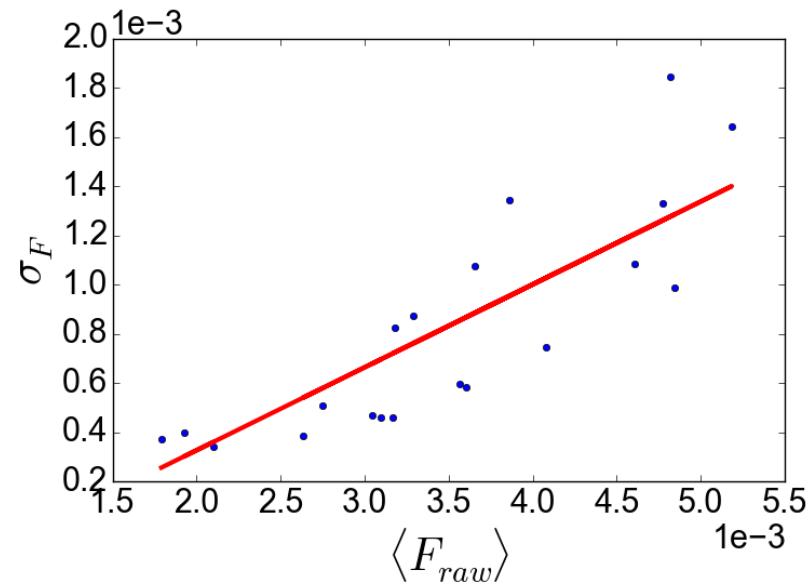
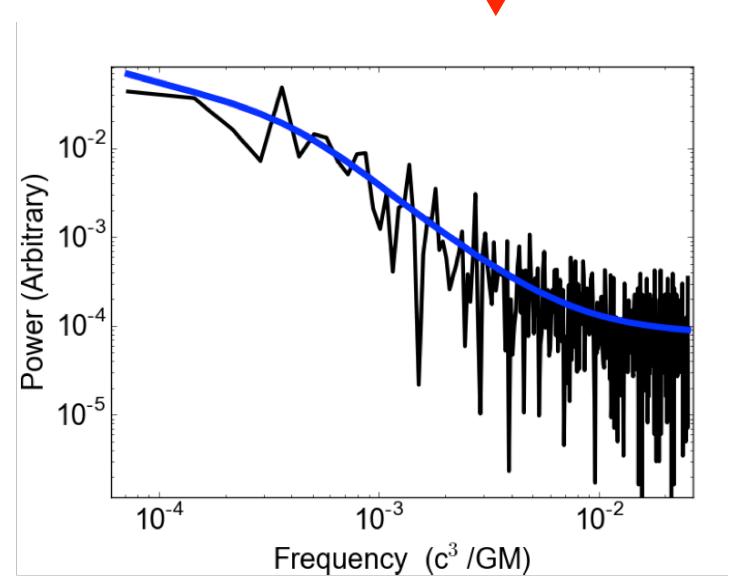
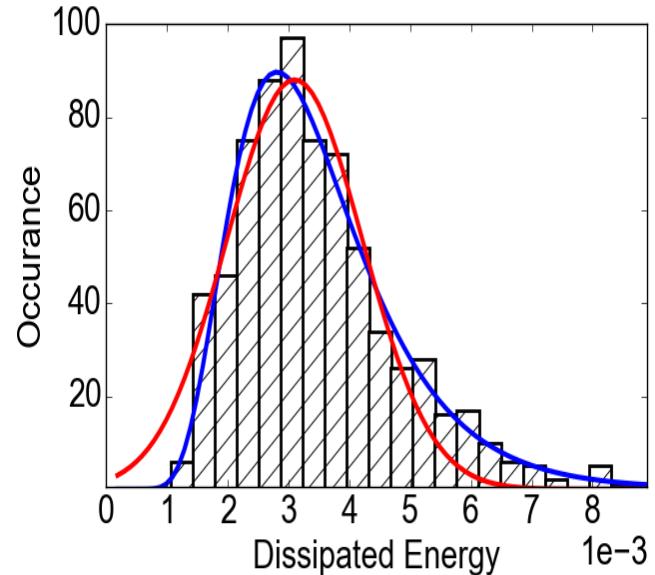
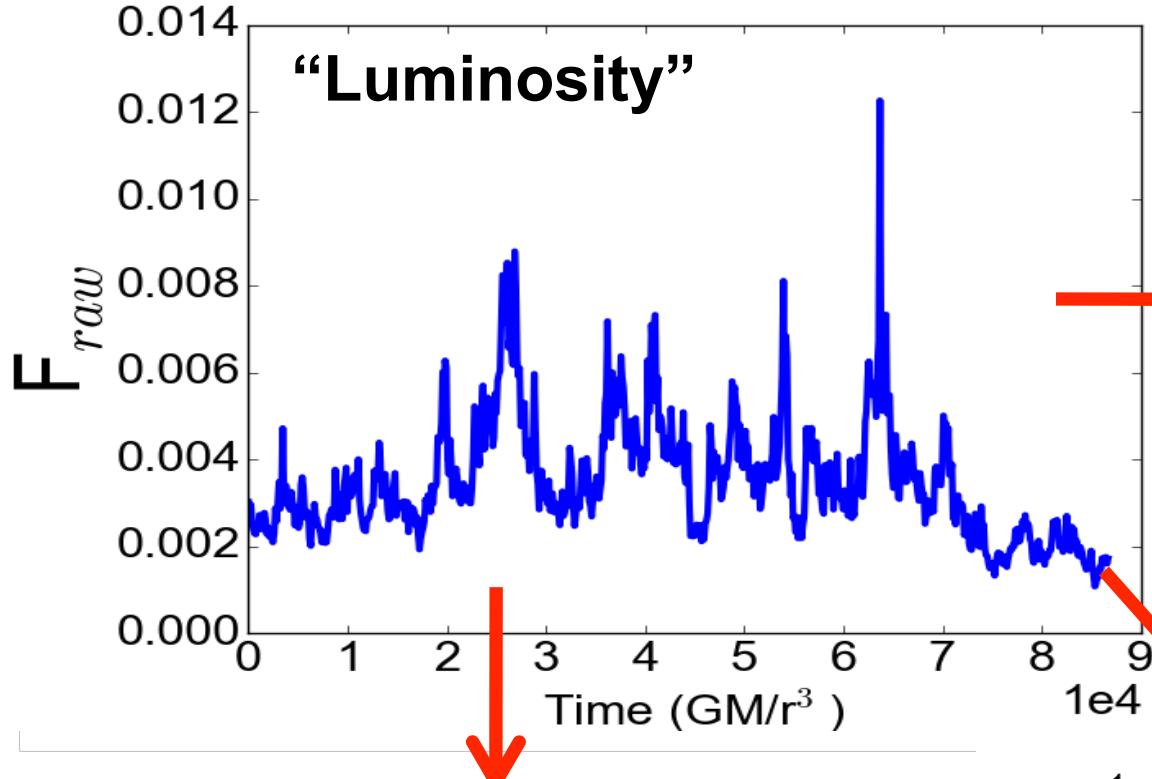
Pseudo-Newtonian Gravity  
Simple cooling function  
Newtonian MHD

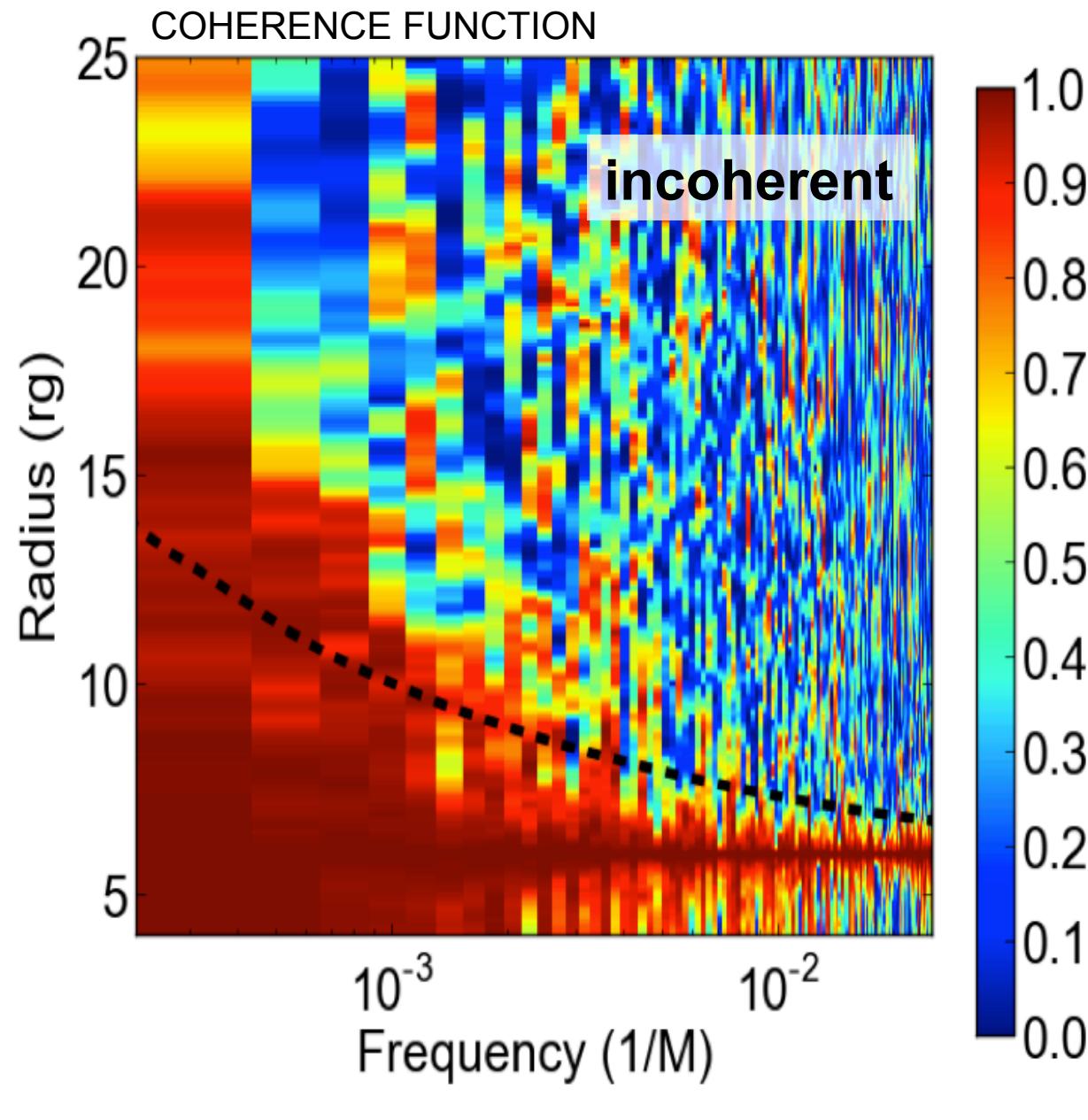
ZEUS-MPV2 Code  
2000 ISCO orbits

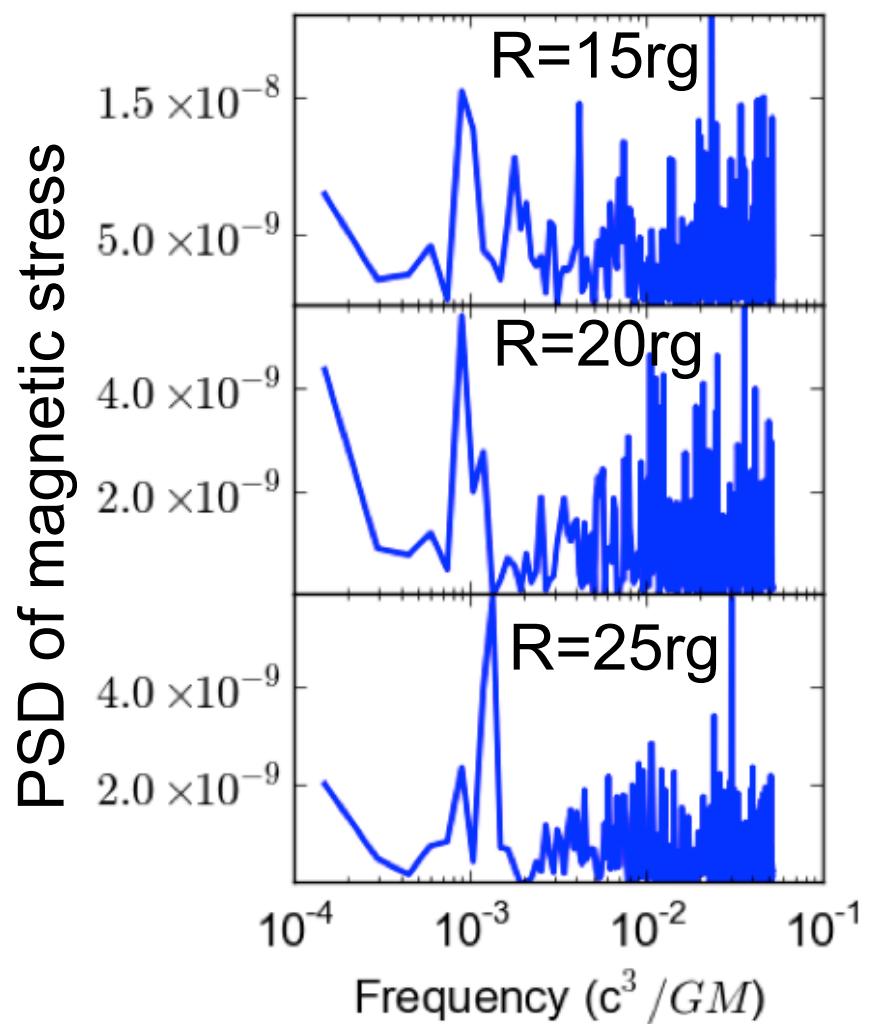


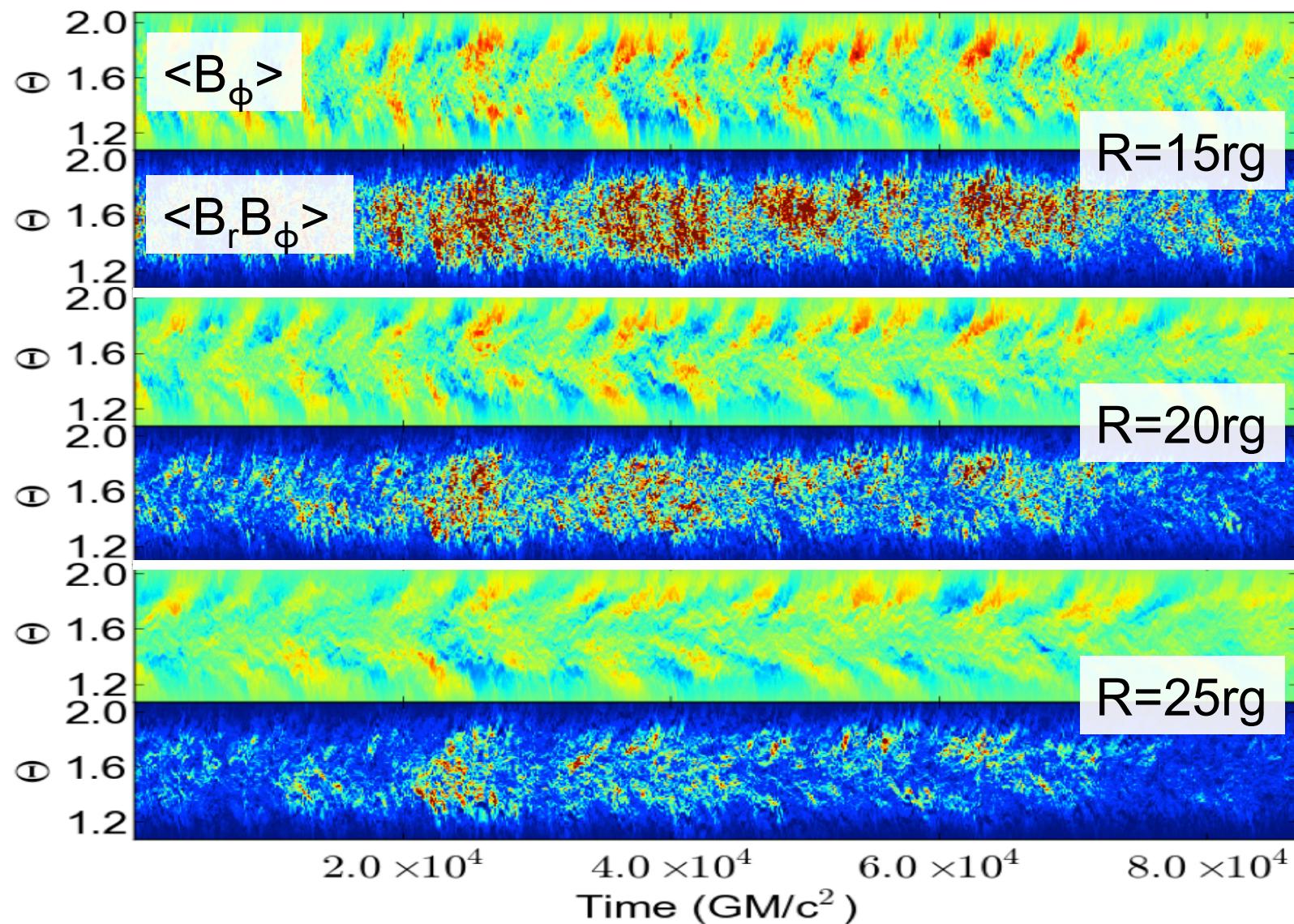


Hogg & Reynolds (2016)









# Take home message & WAGs

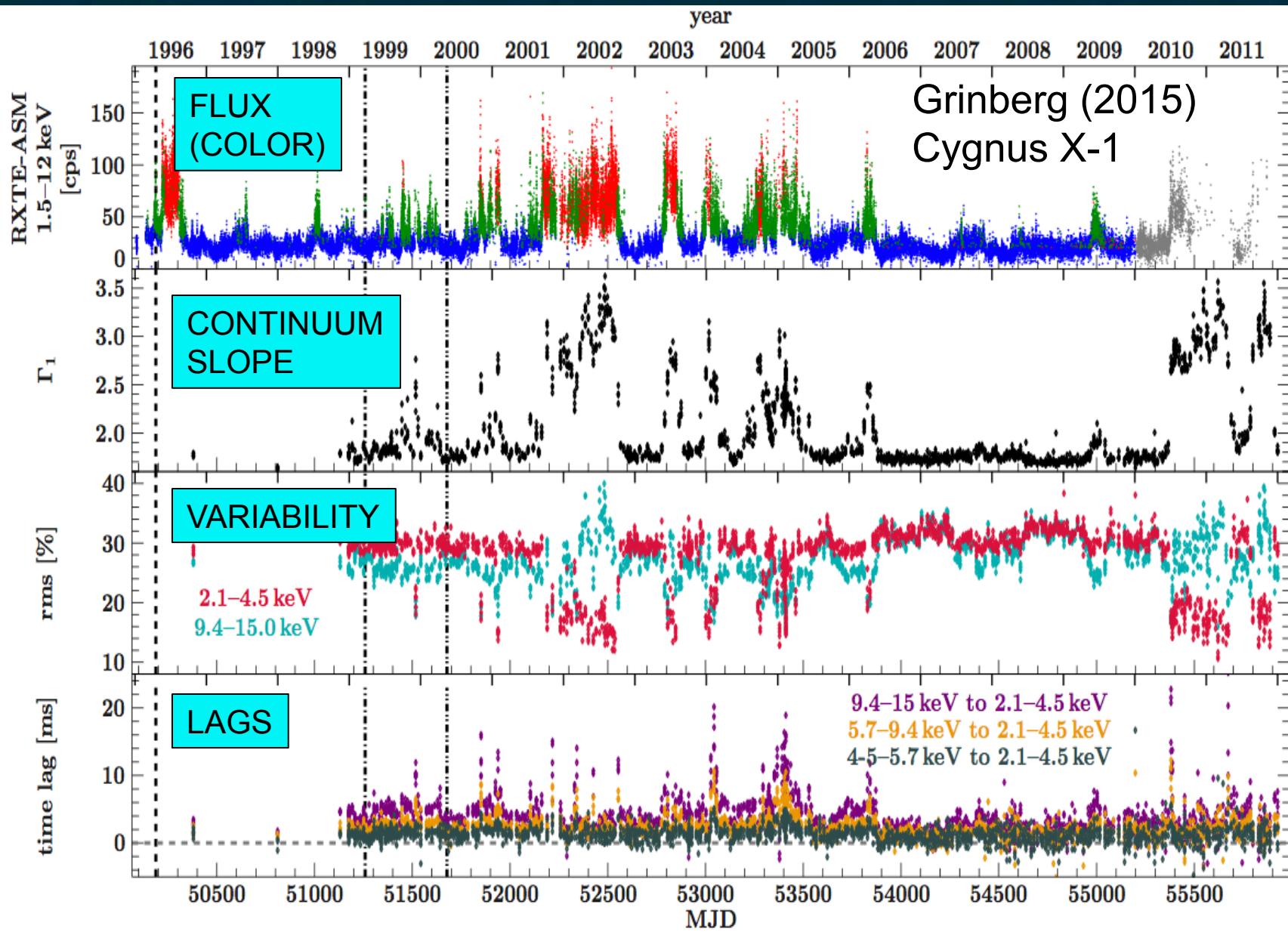
DYNAMO CYCLES → NON-LINEAR VARIABILITY

BUT... WWCS? {What Would Chris (Done) Say?}

High-soft state BHs have very little high-frequency variability at all!

Maybe dynamo is suppressed in the HSS; thermodynamics (maybe  $Pr_m$  instability), net magnetic flux...

## II : State transitions and the nature of truncated disks



Very High State



High State



Intermediate State



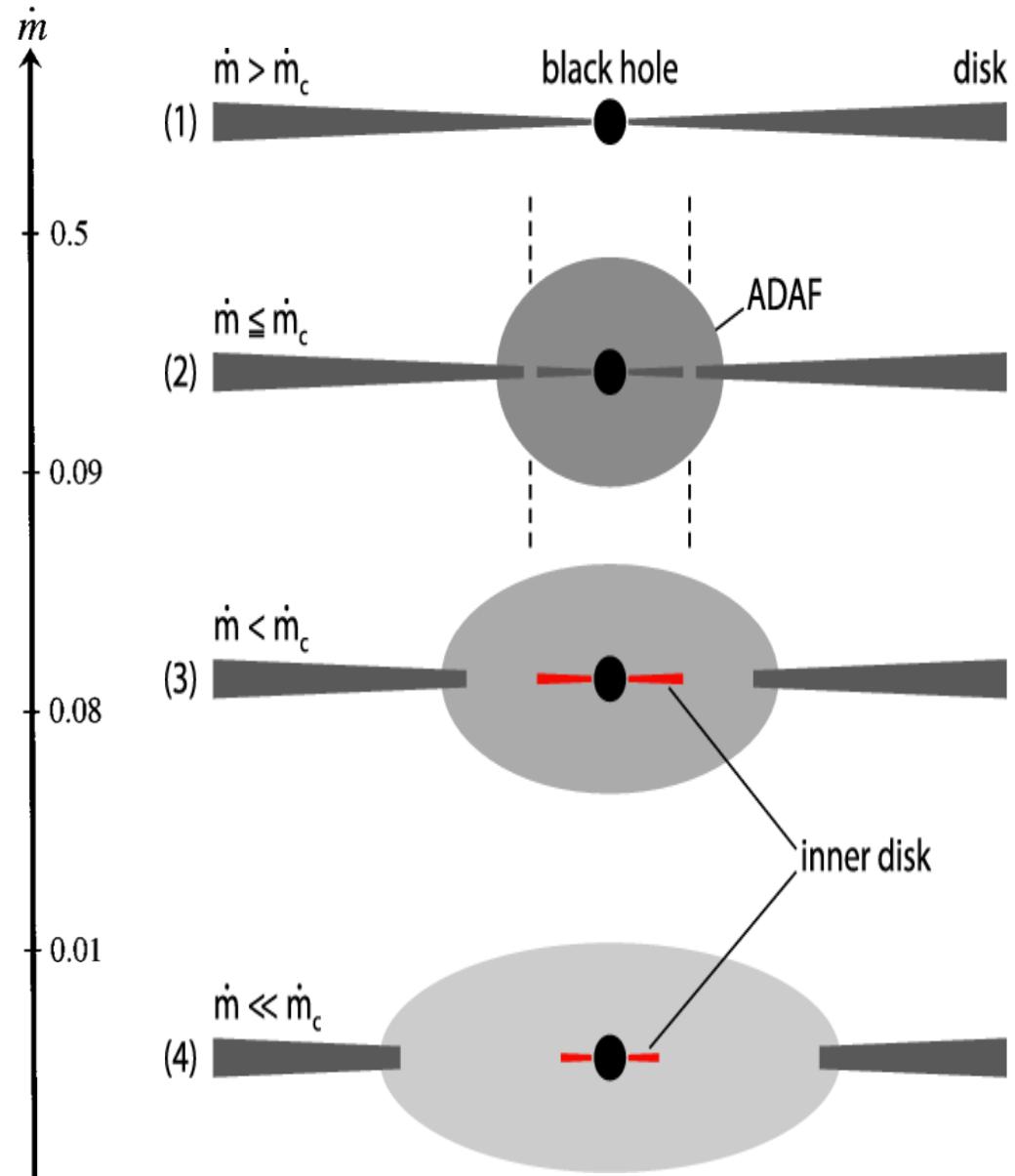
Low State



Quiescent State



Esin et al. (1997)



Meyer-Hofmeister et al. (2009)

At  $r=10rg$ ...

$$t_{\text{dyn}} = 1.5\text{ms}$$

$$t_{\text{th}} \sim 15\text{ms}$$

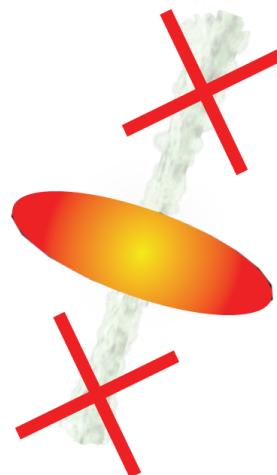
$$t_{\text{visc}} \sim 1.5\text{s}$$

At  $r=100rg$ ...

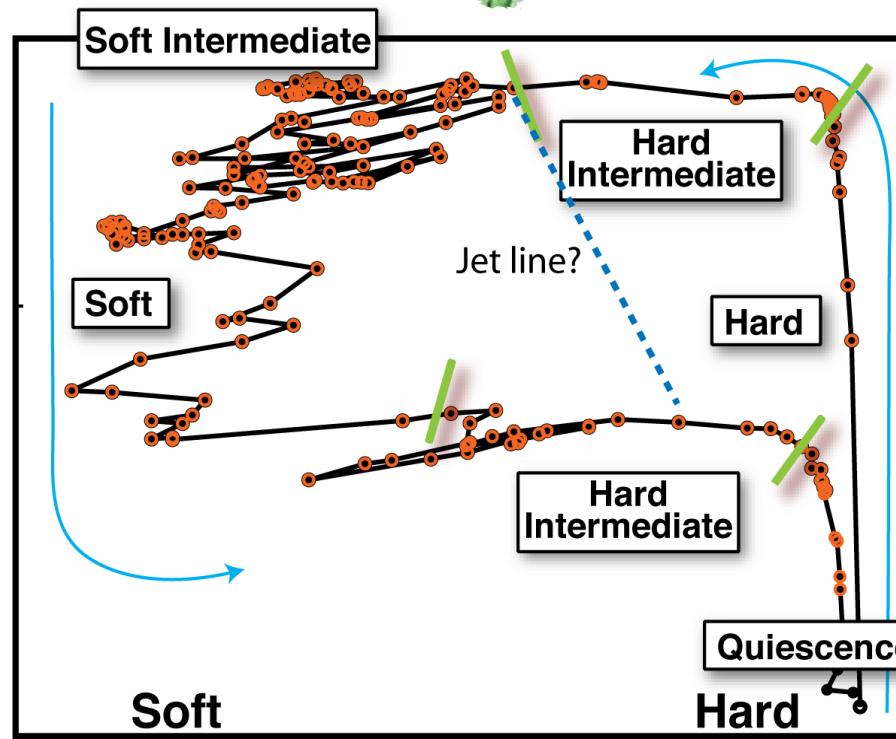
$$t_{\text{dyn}} = 50\text{ms}$$

$$t_{\text{th}} \sim 0.5\text{s}$$

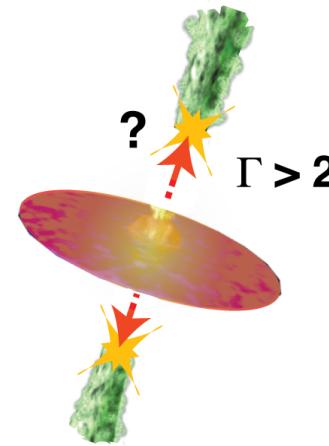
$$t_{\text{visc}} \sim 50\text{s}$$



X-ray Luminosity/Accretion rate



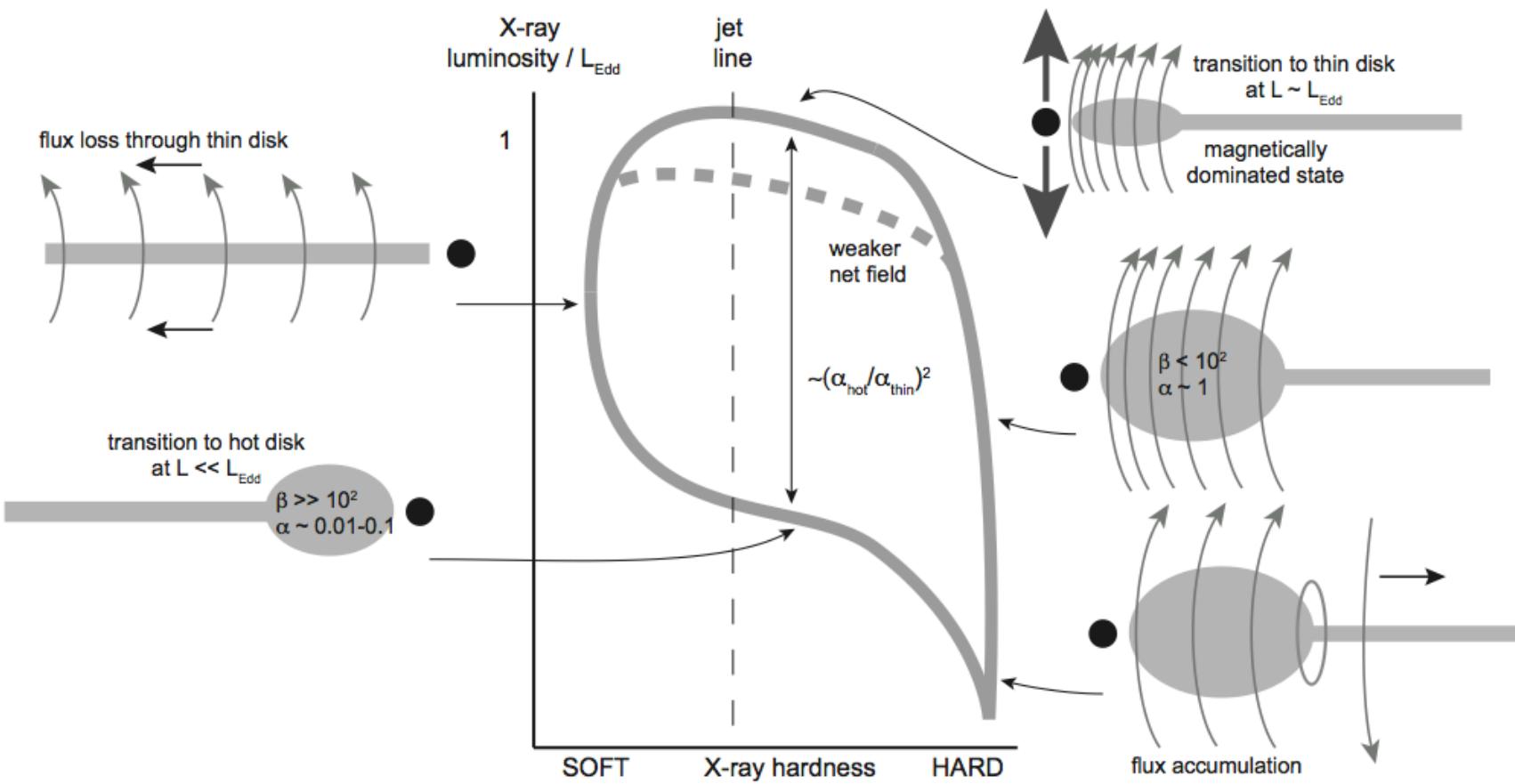
**HYSTERESIS!**



Courtesy of  
Sera Markoff

Spectral Hardness  
(soft=more thermal, hard=more nonthermal)

# Begelman & Armitage (2014) model



**Figure 1.** Illustration of the observed evolution of black hole binary disks in  $L_X$  versus spectral hardness (simplified from Fender, Belloni & Gallo [2004]), and how that relates to the assumed evolution of net magnetic fields in our model. The luminosity of the upper branch of the transition depends on the net accumulated flux, which can vary from cycle to cycle.

# Need to understand truncation...

- Before we can understand state changes, we must understand the structure of flows with thin/efficient outer regions and thick/inefficient inner regions. Some issues
  - Can the thin→thick transition occur over  $\Delta r \ll r$  ?
  - Is a steady flow structure possible?
  - Nature of stresses/energy fluxes across the transition zone?
- Thermal physics driving the radial transition may be complex...
- So, make a toy system with a simple cooling function designed to drive a thermal instability; then study resulting thin→thick transition

# Toy cooling function

- **Design specs** : Want optically-thin cooling which is positive-definite, does not drive viscous instability, but can drive thermal instability. For steady-state and efficient disk, this will balance heating

$$\mathcal{H} = \frac{3}{2}\alpha\mathcal{R}\rho T\Omega_K$$

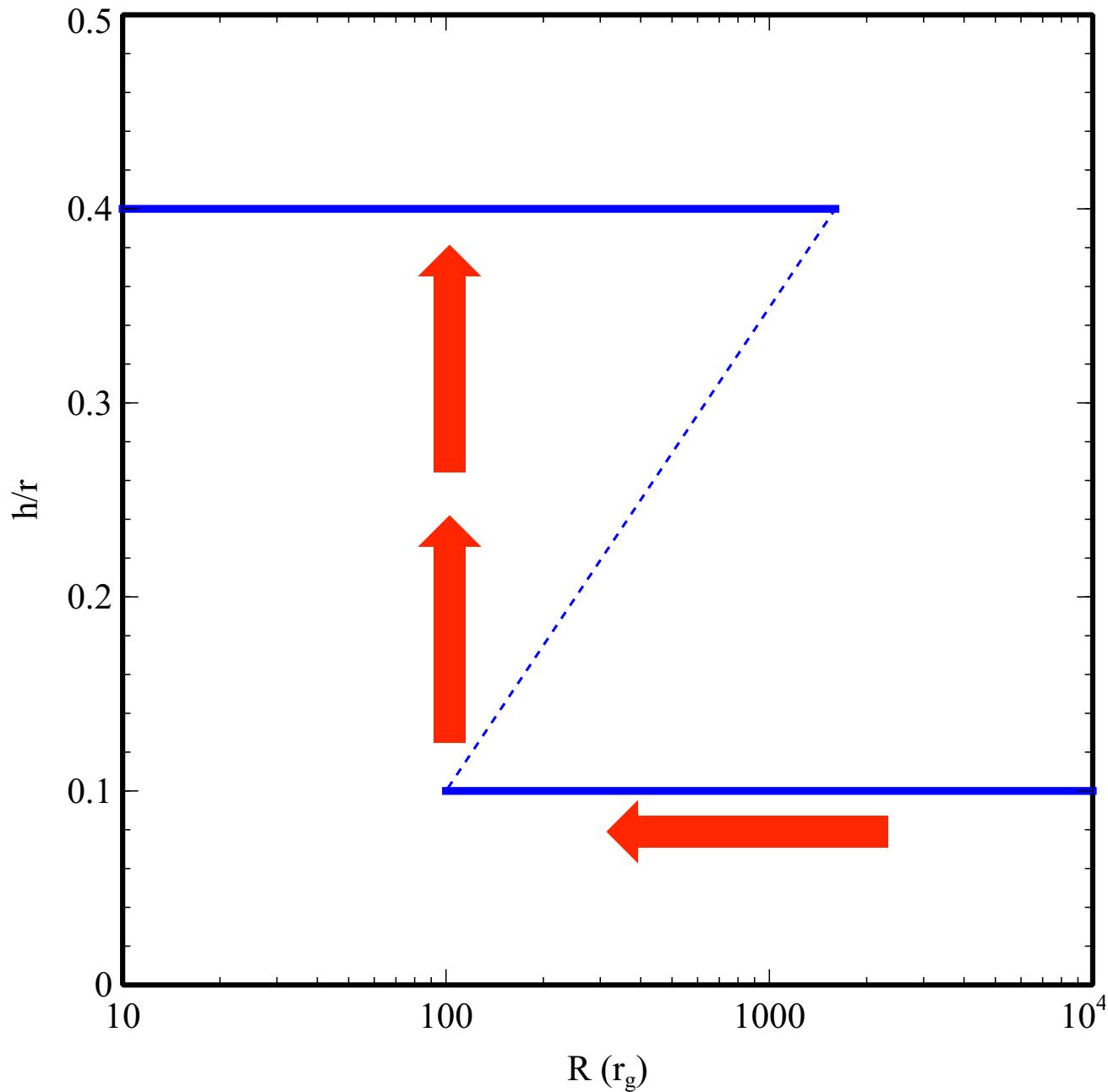
- Choose  $\Lambda = \Lambda_0(T)\rho T^2 f(r)$ ,

$$\mathcal{H} = \Lambda \rightarrow \frac{h}{r} = \sqrt{\frac{3\alpha}{2\Lambda_0(T)}},$$

NOT a simple Noble et al.  
style “target-h/r” cooling  
function!

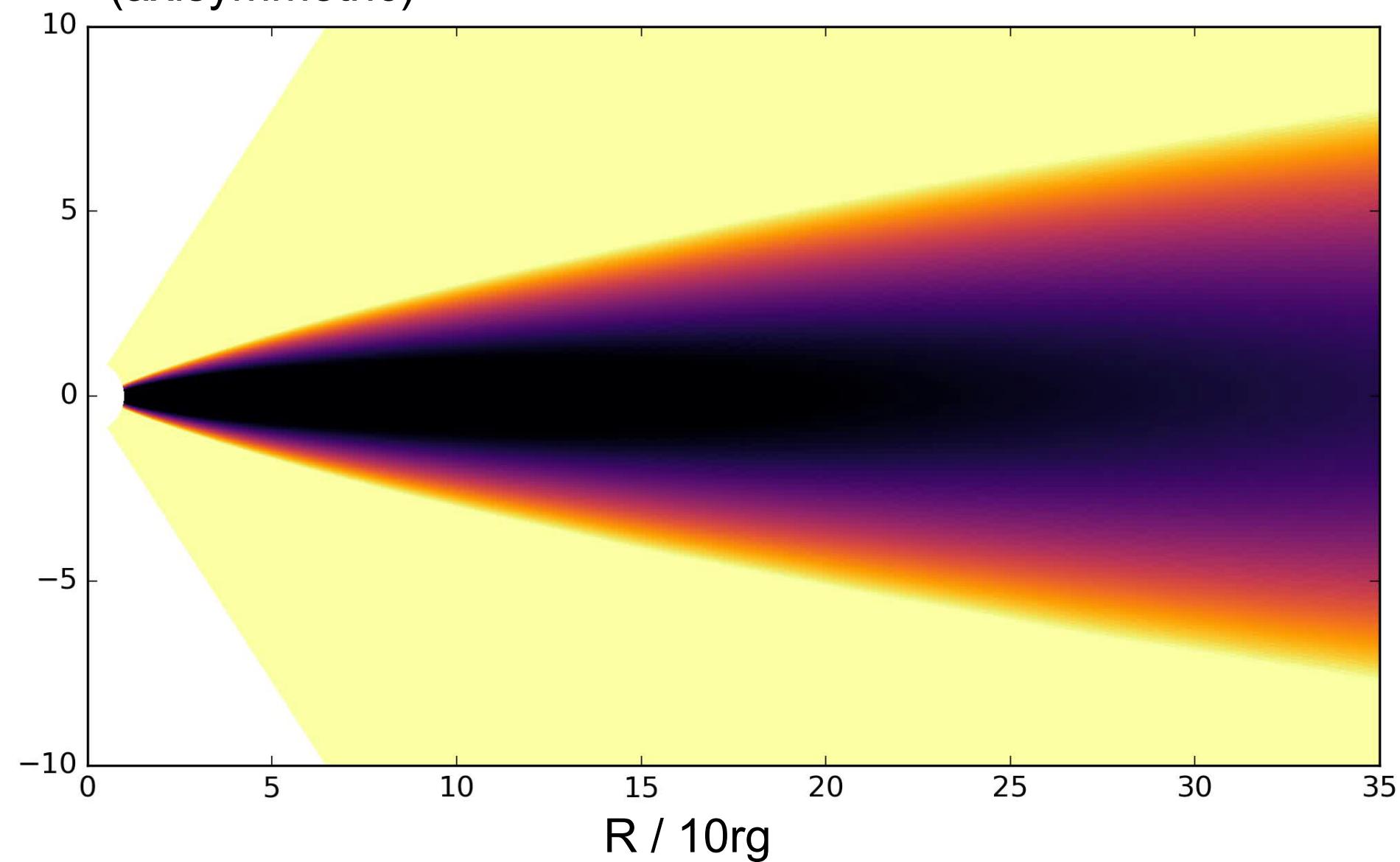
- Thermal instability introduced by T-dependent  $\Lambda_0$

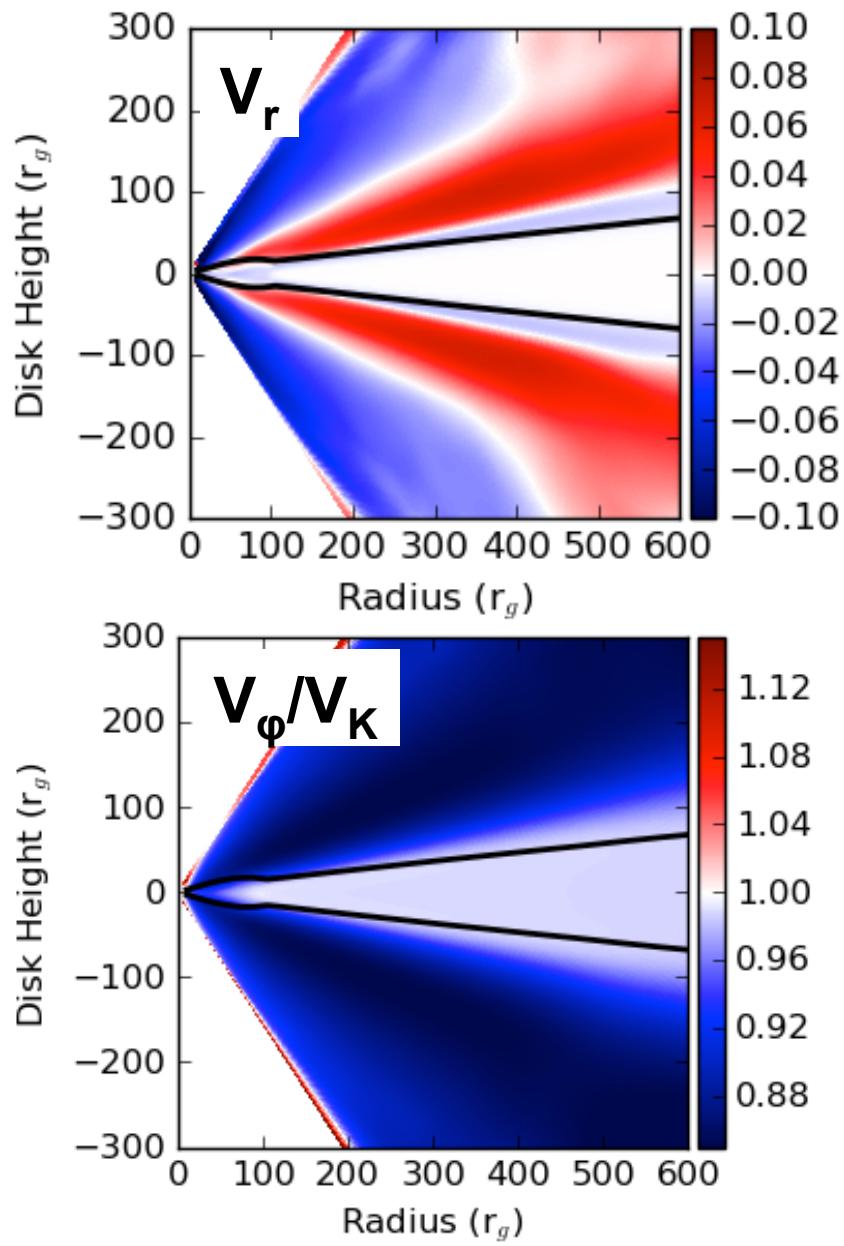
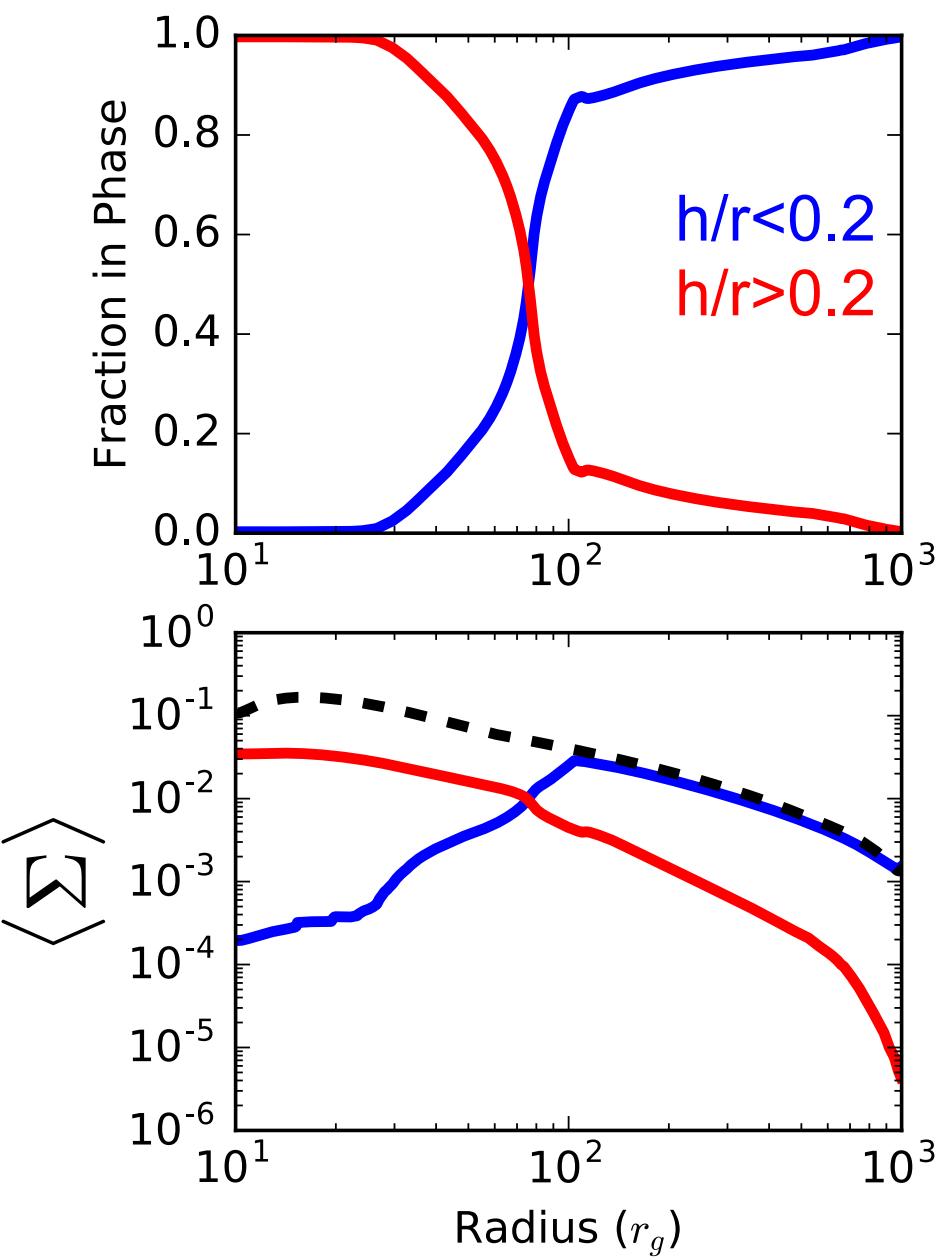
$$\Lambda_0(T) = \begin{cases} \Lambda_c & (T < T_T) \\ \Lambda_h & (T \geq T_T) \end{cases}$$



Hydro with  $\alpha$ -viscosity  
(axisymmetric)

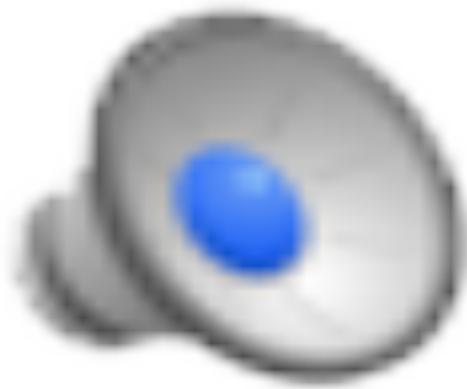
Hogg & Reynolds (submitted)  
Also, Drew's poster!

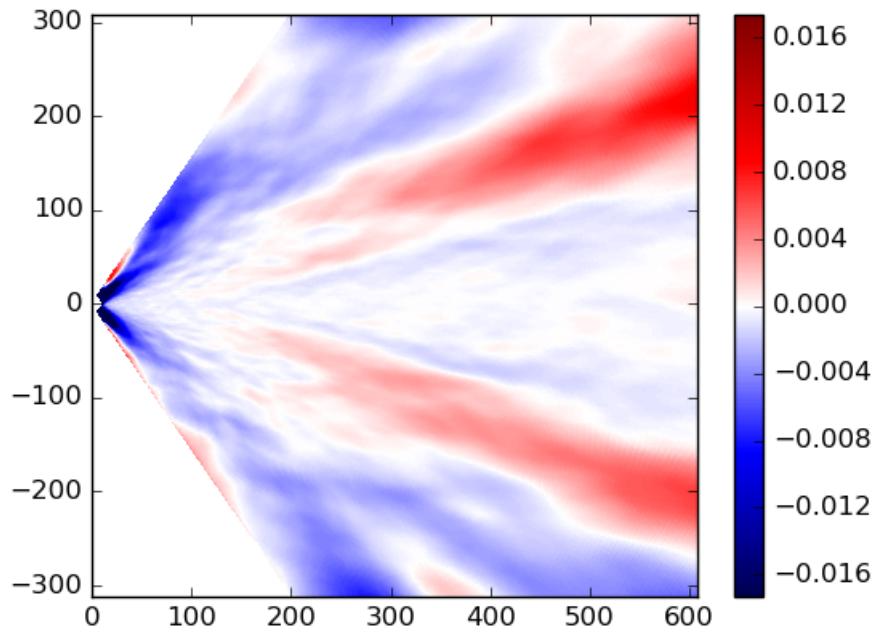
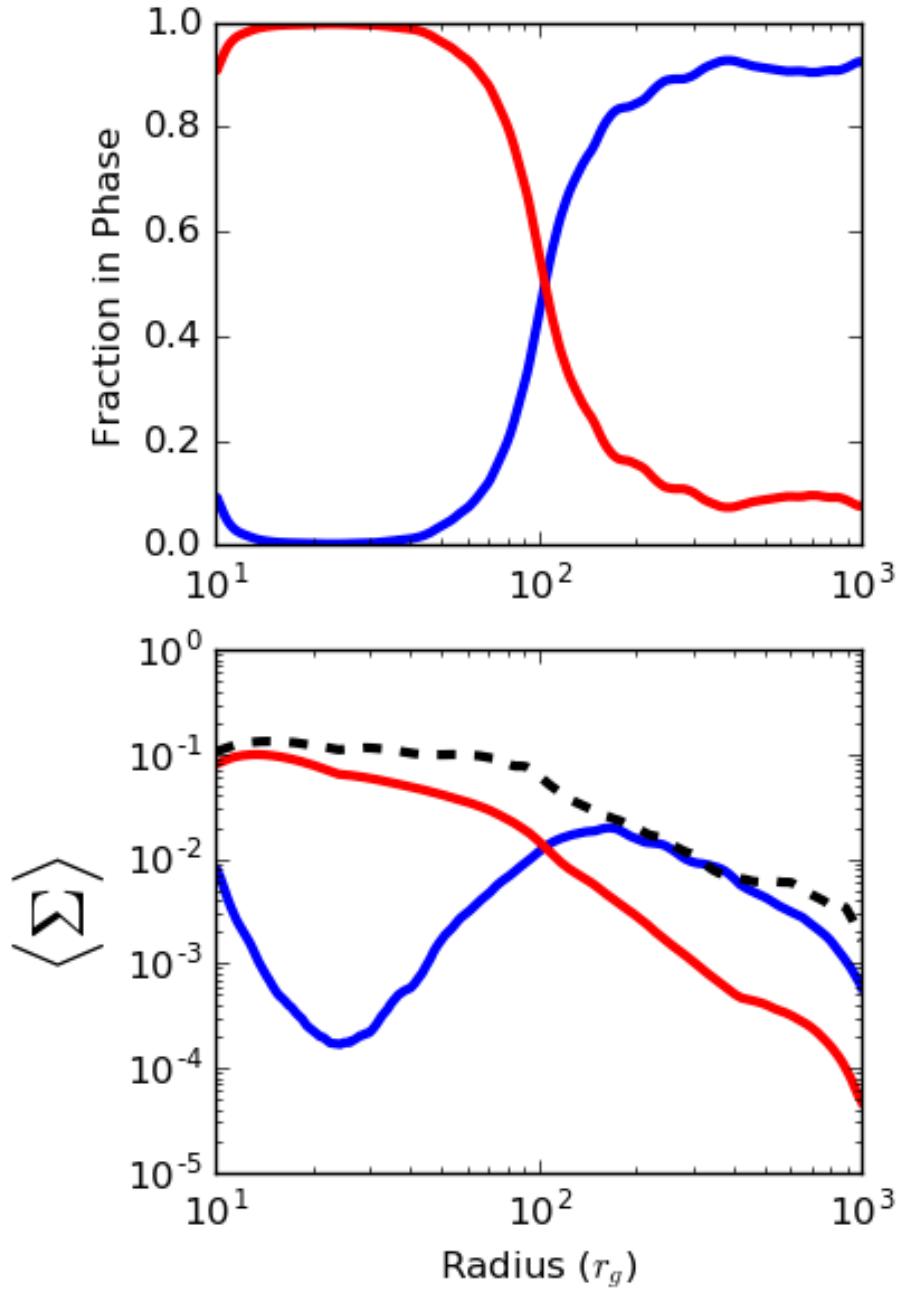


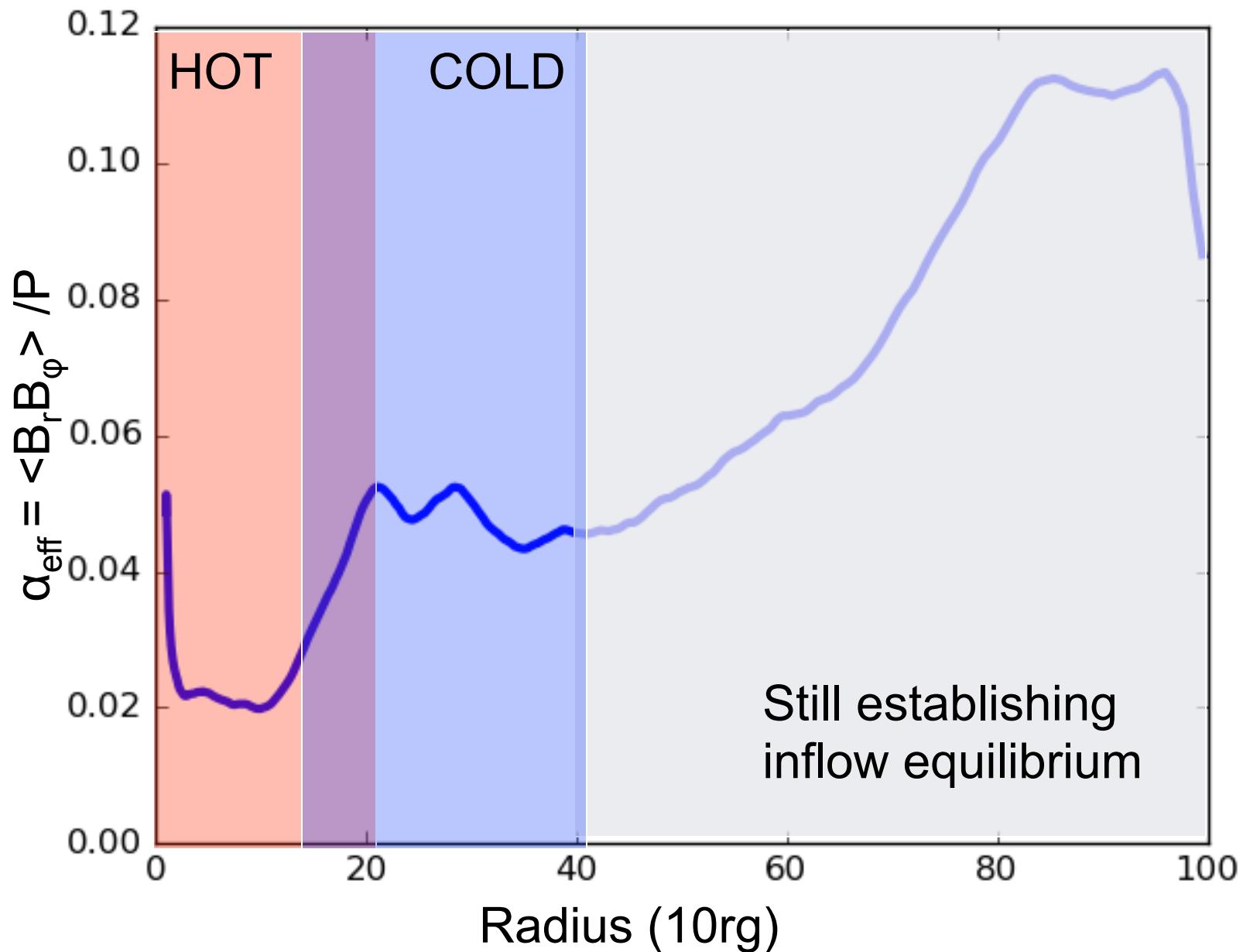


# Non-relativistic ideal MHD 3D ( $\pi/3$ wedge simulation)

**PRELIMINARY**







# Conclusions

- Structure of aperiodic variability
  - Propagating fluctuations driven by stochastic changes in the angular momentum transport
  - Dynamical timescale turbulence non-sufficient... slower dynamo drives variability
- Nature of truncated disks
  - Can recover thin→thick transition
  - Viscous HD model... thermal limit cycle, thermal breeze, and significant cold matter interior to transition
  - MHD model (preliminary)... no limit cycle, thermal breeze, little cold matter interior to transition.