# Gravito-turbulence in Irradiated Protoplanetary Disks 

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Confronting MHD Theories of Accretion Disks with Observations KITP

## Outline

1. Gravitational instability (GI) in accretion disks
2. Shearing box simulations of Gl at 50 AU in irradiated protoplanetary disks (SH and Shi, in review)
3. Radial dependence of GI in irradiated protoplanetary disks (SH and Shi, in preparation)

## Angular momentum transport in accretion disks

- The evolution equation of surface density $\Sigma(r, t)$ is written as

$$
\frac{\partial \Sigma}{\partial t}-\frac{1}{r} \frac{\partial}{\partial r}\left(\frac{2}{r \Omega} \frac{\partial}{\partial r}\left(r^{2} W_{r \phi}\right)\right)=0
$$

- If shear stress is locally determined as $W_{r \phi}(\Sigma ; r)$, the evolution equation can be solved as a diffusion equation.



## Origin of shear stress in accretion disks

$$
\begin{array}{cc}
W_{r \phi}= & \int\left\langle-B_{r} B_{\phi}\right\rangle d z+\int\left\langle\rho v_{r} \delta v_{\phi}\right\rangle d z+\int\left\langle\frac{g_{r} g_{\phi}}{4 \pi G}\right\rangle d z \\
\text { Maxwell stress } & \text { Reynolds stress } \\
\text { magneto-rotational } & \text { gertical shear }
\end{array} \text { gravitational stress }
$$



Kratter and Lodato (2016)

## Gravitational instability in accretion disks

- Instability condition for the axisymmetric mode in infinitesimally thin disks (Toomre 1964)

$$
Q \equiv \frac{c_{\mathrm{s}} \Omega}{\pi G \Sigma}<1
$$

- $Q$ tends to be self-regulated to $\sim 1$ (Paczynski 1978)

1. $Q<1$ : Gl and heating sets in
2. $T$ increases if heating overcomes cooling
3. $Q>1$ : Gl and heating is quenched
4. $T$ decreases
5. $\rightarrow 1$


## Angular mom. transport by GI depends on $\mathrm{M}_{\text {disk }} / \mathrm{M}_{\text {star }}$

Kratter and Lodato (2016)


| M $_{\text {disk }} / M_{\text {star }}$ | large | small |
| :---: | :---: | :---: |
| angular mom. <br> transport | non-local | local |

## Locality of angular momentum transport

- Transport is local when $M_{\text {disk }} / M_{\text {star }} \leq 0.25$ (Lodato \& Rice 2005)

$\mathbf{M}_{\text {disk }} / M_{\text {star }}=\mathbf{0 . 2 5}$



## Gravito-turbulence vs. fragmentation

- Using a simple cooling function with $\beta$ being constant cooling time

$$
\frac{\partial e}{\partial t}=-\frac{e}{\beta \Omega^{-1}},
$$

Gammie (2001) found for 2D (razor thin) disks of $\gamma=2$

$$
\begin{cases}\beta>3 & \text { gravito-turbulence with } Q \sim 1 \\ \beta<3 & \text { fragmentation. }\end{cases}
$$


$\beta=2$
(strong cooling)

$\beta=10$
(weak cooling)

## Fragmentation criterion in terms of $\alpha$

- In thermal equilibrium where the $\beta$ cooling equals the $\alpha$ dissipation,

$$
\alpha=\frac{4}{9} \frac{1}{\gamma(\gamma-1) \beta}
$$

- Gravito-turbulence of $\alpha>\alpha_{\max } \sim 0.06$ cannot be sustained (Rice, Lodato \& Armitage 2005).



## No solid fragmentation criterion?

- Fragmentation is stochastic since shock heating is very localized.
- Fragmentation becomes increasingly rare for larger $\beta$.



## Key parameters in GI in accretion disks

- Toomre parameter $Q \equiv \frac{c_{\mathrm{s}} \Omega}{\pi G \Sigma}$
- self-regulated to $Q \sim 1$
- disk-star mass ratio $M_{\text {disk }} / M_{\text {star }}$ and angular momentum transport
- local for small $M_{\text {disk }} / M_{\text {star }}$ and global for large $M_{\text {disk }} / M_{\text {star }}$
- cooling time $\beta \equiv t_{\text {cool }} \Omega$
- fragmentation criterion: $\beta_{\text {min }}=3$ for $\gamma=2\left(\alpha_{\max }=0.06\right)$
- no solid criterion due to stochastic fragmentation?


## Our approach to Gl in accretion disks

- Toomre parameter $Q \equiv \frac{c_{\mathrm{s}} \Omega}{\pi G \Sigma}$
- self-regulated to $Q \sim 1$
- disk-star mass ratio $M_{\text {disk }} / M_{\text {star }}$ and angular momentum transport
- local for small $M_{\text {disk }} / M_{\text {star }} \rightarrow$ shearing box simulations
- cooling time $\beta \equiv t_{\text {cool }} \Omega \rightarrow$ realistic thermodynamics with irradiation
- fragmentation criterion: $\beta_{\text {min }}=3$ for $\gamma=2\left(\alpha_{\max }=0.06\right)$
- no solid criterion due to stochastic fragmentation?
- 3D - vertical stratification


## Purpose of this study and simulation setup

- to explore nonlinear outcome of Gl in irradiated protoplaneatary disks
- nature of gravito-turbulence
- fragmentation criterion
- using 3D stratified shearing box simulations
- realistic thermodynamics (opacities and equation of state)
- irradiation



## Basic equations

$$
\begin{array}{ll}
\frac{\partial \rho}{\partial t}+\nabla \cdot(\rho \boldsymbol{v})=0, & \text { continuity eq. } \\
\frac{\partial(\rho \boldsymbol{v})}{\partial t}+\nabla \cdot(\rho \boldsymbol{v} \boldsymbol{v})=-\nabla p-\rho \nabla \Phi+\frac{\kappa_{\mathrm{R}} \rho}{c} \boldsymbol{F}, & \text { momentum eq. } \\
\frac{\partial e}{\partial t}+\nabla \cdot(e \boldsymbol{v})=-(\nabla \cdot \boldsymbol{v}) p-(4 \pi B(T)-c E) \kappa_{\mathrm{P}} \rho, & \text { gas energy eq. } \\
\frac{\partial E}{\partial t}+\nabla \cdot(E \boldsymbol{v})=-\nabla \boldsymbol{v}: \mathrm{P}+(4 \pi B(T)-c E) \kappa_{\mathrm{P}} \rho-\nabla \cdot \boldsymbol{F}, & \text { rad. energy eq. } \\
\boldsymbol{F}=-\frac{c \lambda}{\kappa_{\mathrm{R}} \rho} \nabla E, & \text { solved time-implicitly } \\
\nabla^{2} \Phi=4 \pi G \rho . & \text { FLD approx. } \\
\text { Poisson eq. }
\end{array}
$$

Fiducial run

$$
\begin{gathered}
\theta=0.02 \\
r=50 \mathrm{AU} \\
\Sigma=100 \mathrm{gcm}^{-3}
\end{gathered}
$$

( $\mathrm{M}_{\text {disk }} / \mathrm{M}_{\text {star }} \sim 0.09$ )

## Gravito-turbulence



## Gravito-turbulence

$t=20.00$ orbits


## Gravito-turbulence



## Gravito-turbulence is not usual "turbulence"

- There is no apparent turbulent cascade with energy dissipation occuring on the smallest scales.
- GI repeatedly excites density waves, which dissipate through shock waves and compressional heating.


Hydrostatic balance (gravitational accelerations)


Hydrostatic balance (pressure gradients)


Thermal balance (heating rates)
radiative
cooling


Thermal balance (cooling rates)



## Dependence on surface density $\boldsymbol{\Sigma}$

$$
\begin{gathered}
\theta=0.02 \\
r=50 \mathrm{AU} \\
30<\Sigma<300 \mathrm{gcm}^{-2}
\end{gathered}
$$

$$
\text { (0.03 < } \left.M_{\text {disk }} / M_{\text {star }}<0.26\right)
$$

$$
\Sigma=60 \mathrm{gcm}^{-3}
$$

$\Sigma=100 \mathrm{gcm}^{-3}$ (fiducial run)
$\Sigma=200 \mathrm{gcm}^{-3}$
$\Sigma=300 \mathrm{gcm}^{-3}$

laminar
gravitoturbulence
gravitoturbulence
fragmentation

Fragmentation run $\left(\Sigma=300 \mathrm{gcm}^{-3}\right)$


Fragmentation run $\left(\Sigma=300 \mathrm{gcm}^{-3}\right)$

"First core collapse"


Mass accretion rate as a function of $\Sigma$


## alpha as a function of $\boldsymbol{\Sigma}$



## Cooling time $\boldsymbol{\beta}$ as a function of $\boldsymbol{\Sigma}$



## Toomre $\mathbf{Q}$ as a function of $\boldsymbol{\Sigma}$



Dependence on radius $r$ and surface density $\Sigma$

$$
\theta=0.02
$$

## Nonlinear outcome of Gl in protoplanetary disks



## Phase diagram of nonlinear outcome of Gl



## Gravito-turbulence is sustained when $1>Q>0.7$



## Fragmentation criterion on Q



## Another fragmentation criterion on $\beta$ (or a)



## Summary

1. We investigated the nonlinear outcome of GI in irradiated protoplanetary disks using radiation hydrodynamics simulations.
2. At a fixed radius of 50AU, gravito-turbulence is sustained for a range of $\Sigma$ corresponding to $0.7 \leq Q \leq 1$, where $\beta$ tends to be constant.
3. In gravito-turbulence, density waves excited by GI dissipate through both shock waves and compressional heating.
4. Vertically diverging flows generated by collision of the density waves contribute to both hydrostatic and thermal balances.
5. From parameter survey on both $r$ and $\Sigma$, fragmentation seems to occur when either $Q \leq 0.7$ or $\beta \leq 2$ is satisfied.
