

DARK MATTER SEARCHES WITH ICECUBE

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KAVLI INSTITUTE, 11 JUNE, 2013

THANKS IN ABSENTIA TO C. ROTT, M. DANNINGER, J. MILLER, R. REIMANN

- **ICECUBE IS COMPLETE:** TAKING DATA WITH 86 STRINGS SINCE MAY 2011
- RESULTS FROM 2010/11 (79/86 STRINGS) ARE BEING RELEASED
- ANALYSES ON 86 STRINGS IN THE PIPELINE
- THE LOW-ENERGY EXTENSION, DEEPCORE, USEFUL TO LOWER ENERGY THRESHOLD TO ~ 10 GeV

ALREADY USEFUL IN CONVERTING ICECUBE TO AN ALL-SKY DETECTOR:

ICECUBE REACHES THE SOUTHERN SKY (GALACTIC CENTER)

DOUBLES EXPOSURE OF THE DETECTOR FOR SOLAR DM SEARCHES

- COMPETITIVE LIMITS ON DARK MATTER SEARCHES

- ON THE HIGH-ENERGY SIDE: ν_e ATMOSPHERIC NEUTRINOS

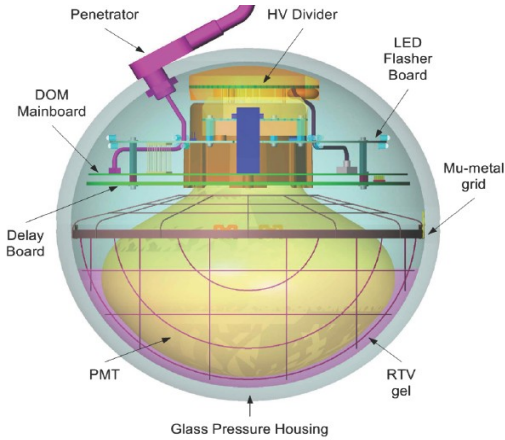
OBSERVATION OF $\mathcal{O}(100)$ TeV-PeV EVENTS

SENSITIVITY TO POINT SOURCES REACHING THE LEVEL WHERE A NON-OBSERVATION

CONSTRAINS MODELS (GRBs FOR EXAMPLE)

- SEVERAL “TARGET OF OPPORTUNITY” PROGRAMS WITH SATELLITES, CTAs AND GW DETECTORS WORKING

THE ICECUBE NEUTRINO TELESCOPE

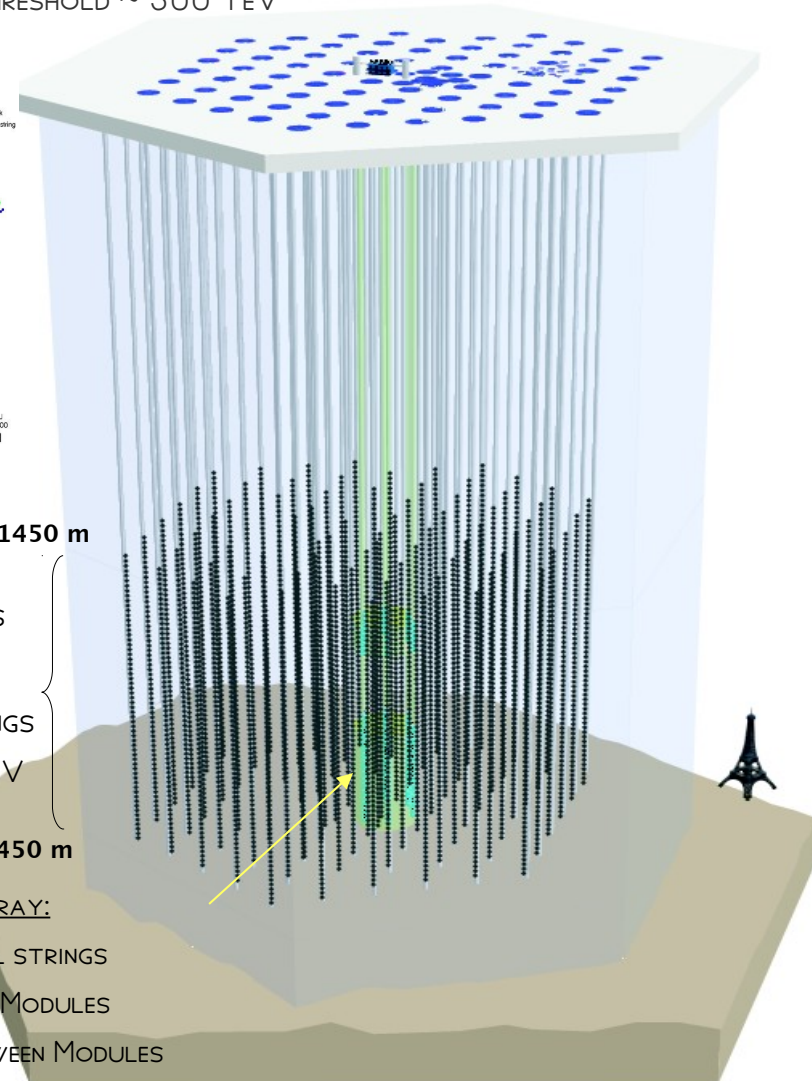
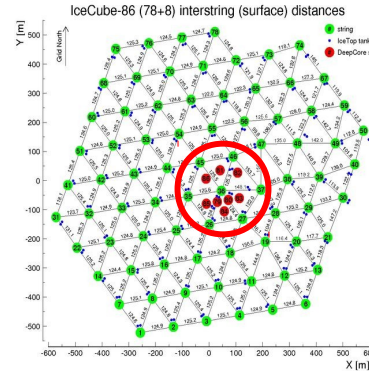


- PMT: HAMAMATSU, 10''
- DIGITIZERS:
 - ATWD: 3 CHANNELS. SAMPLING 300MHZ, CAPTURE 400 NS
 - FADC: SAMPLING 40 MHZ, CAPTURE 6.4 μ s DYNAMIC RANGE 500PE/15 NSEC, 25000 PE/6.4 μ s
- **FLASHER BOARD**:
 - 12 controllable LEDs at 0° or 45°
- DARK NOISE RATE \sim 400 Hz
- LOCAL COINCIDENCE RATE \sim 15 Hz
- DEADTIME < 1%
- TIMING RESOLUTION \leq 2-3 NS
- POWER CONSUMPTION: 3W

CLOCK STABILITY: $10^{-10} \approx 0.1$ NSEC / SEC
 SYNCHRONIZED TO GPS TIME EVERY ≈ 5 SEC
 AT 2 NS PRECISION

ICE TOP: AIR SHOWER DETECTOR

80 STATIONS/2 TANKS EACH
 THRESHOLD \sim 300 TEV



INICE ARRAY:

- 80 STRINGS
- 60 OPTICAL MODULES
- 17 M BETWEEN MODU
- 125 M BETWEEN STRINGS
- ν THRESHOLD \lesssim 100 GeV

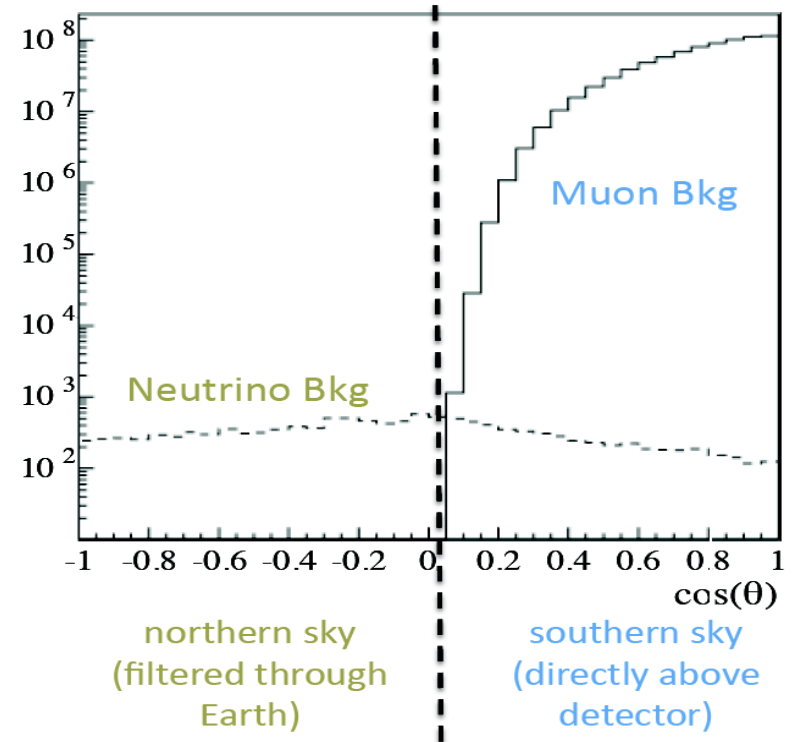
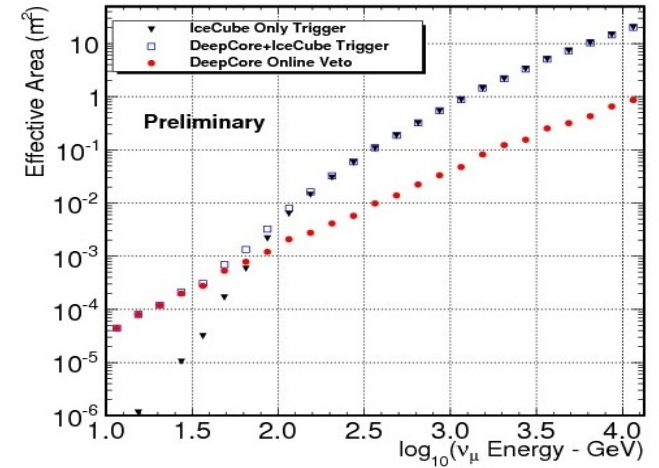
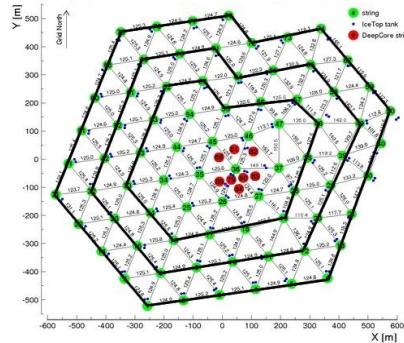
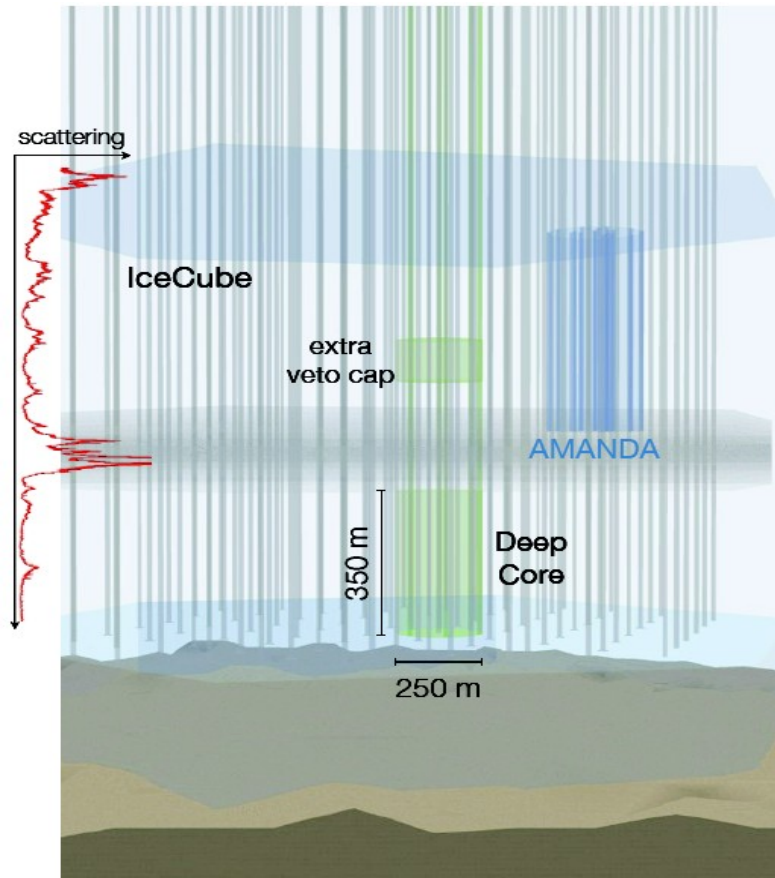
DEEPCORE ARRAY:

- 6 ADDITIONAL STRINGS
- 60 OPTICAL MODULES
- 7/10 M BETWEEN MODULES
- 72 M BETWEEN STRINGS
- ν THRESHOLD \sim 10 GeV

FULL DETECTOR VETO CAPABILITIES

Complete detector \rightarrow **full sky sensitivity**
using IceCube outer strings as a veto

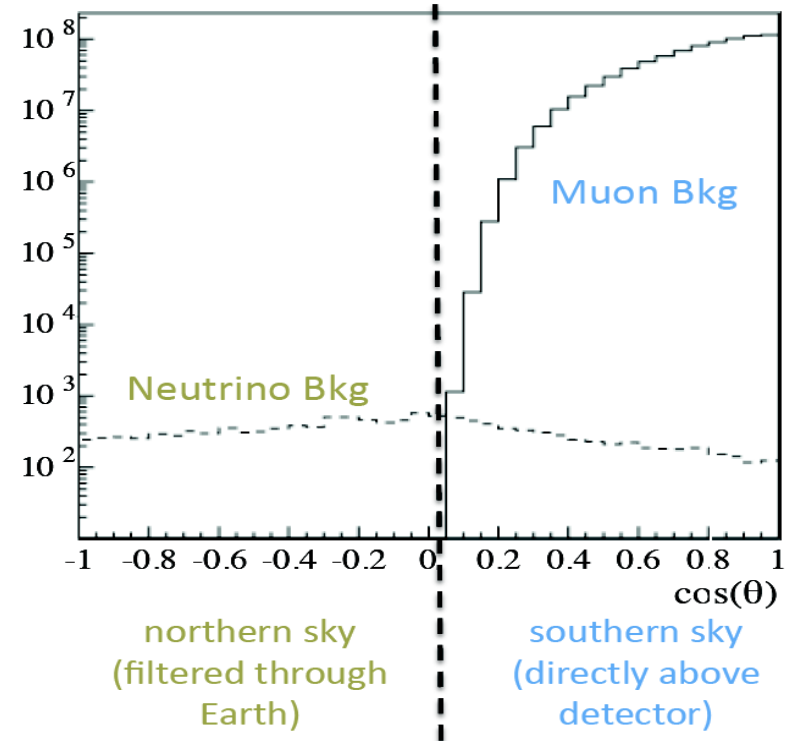
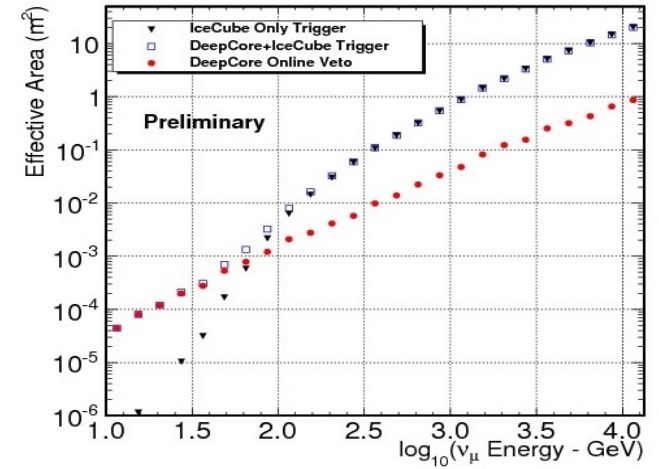
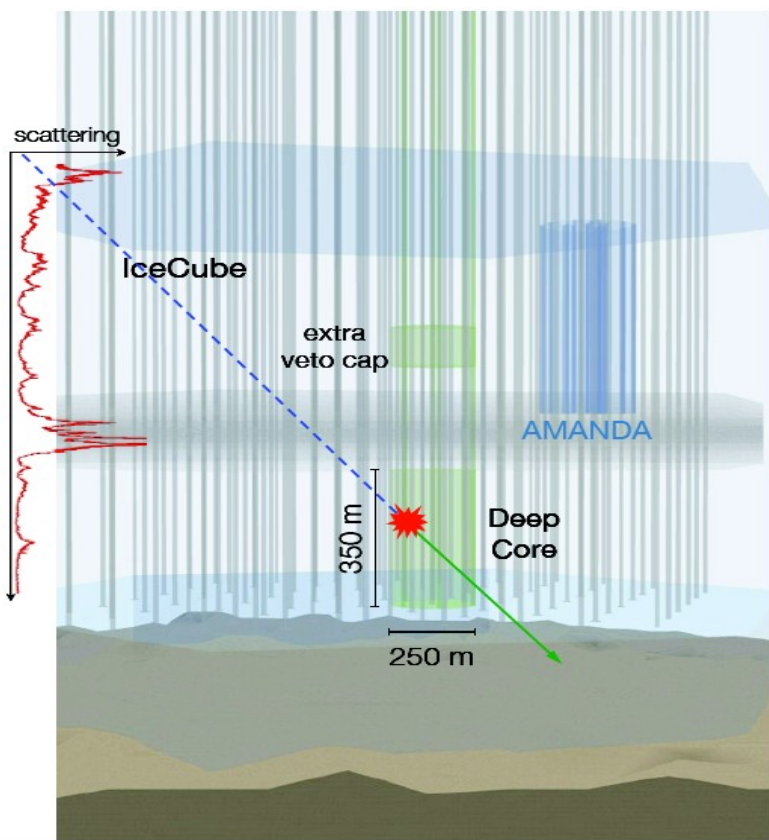
Define starting tracks == neutrinos



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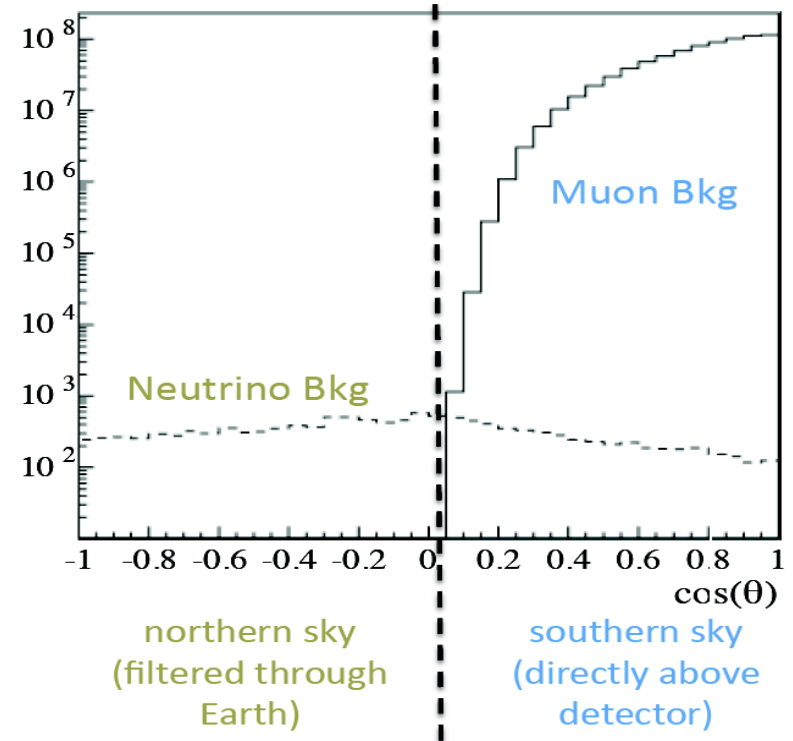
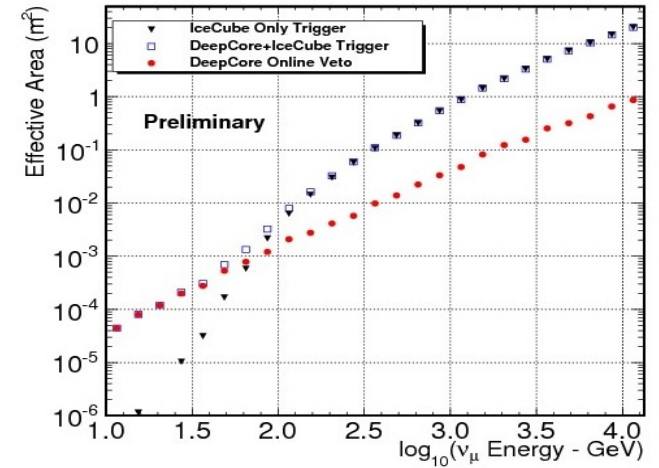
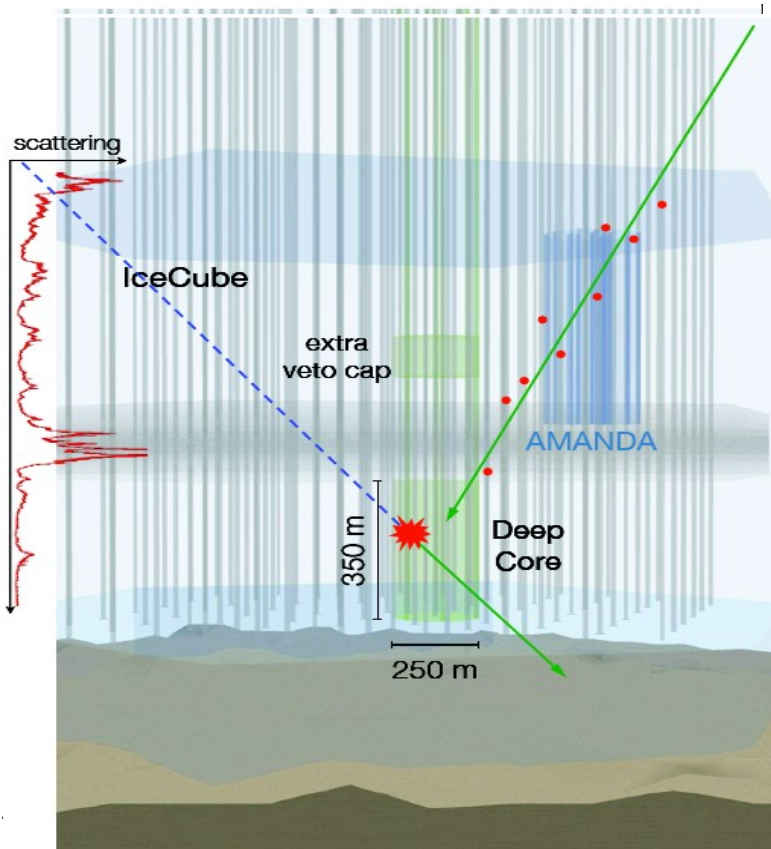
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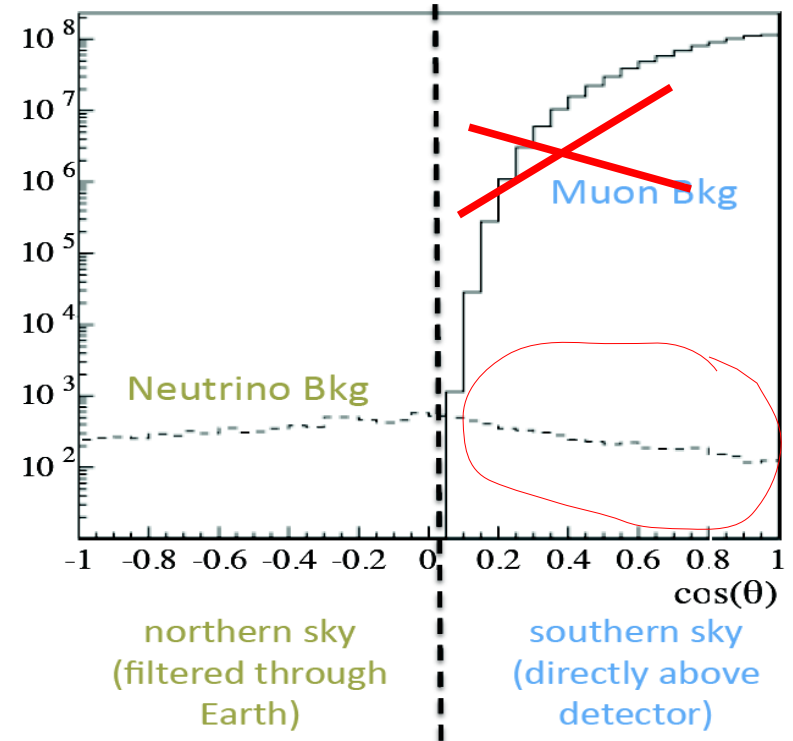
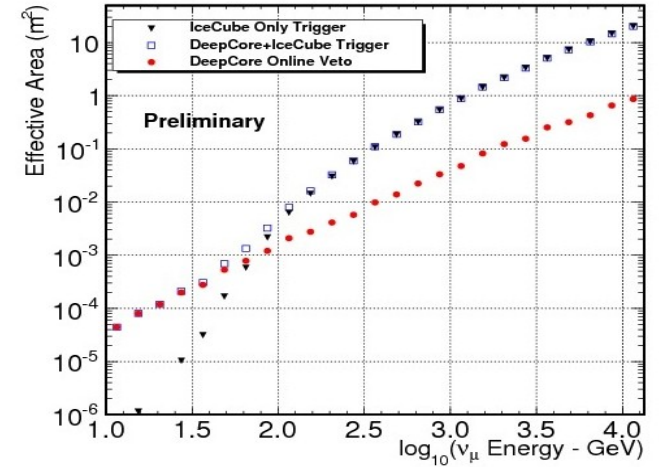
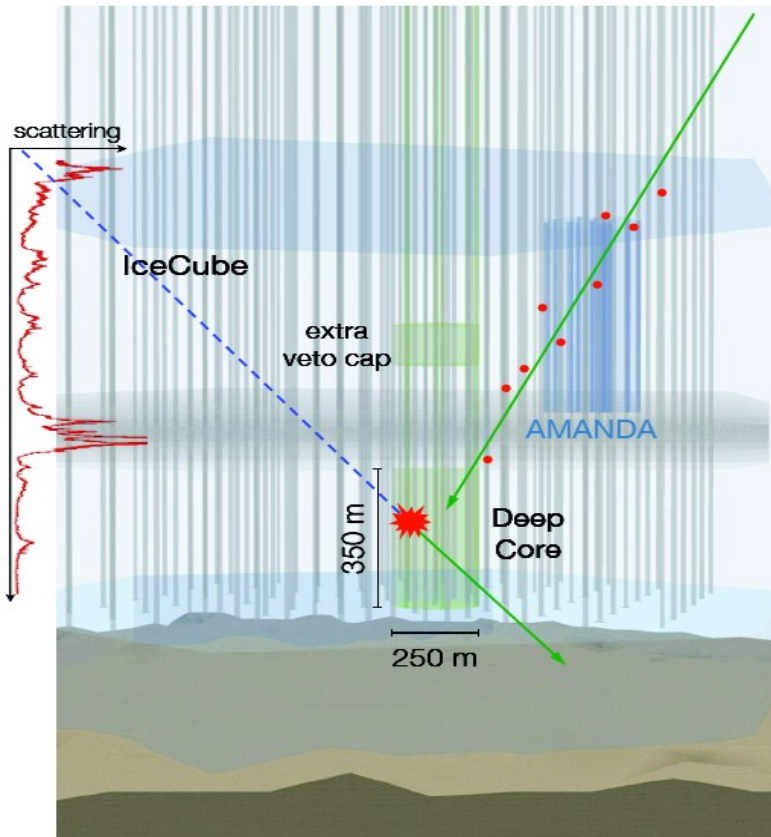


Complete detector → **full sky sensitivity**
using IceCube outer strings as a veto

Define starting tracks == neutrinos

→ access to southern hemisphere, galactic center
and all-year Sun visibility

IceCube is a 4π detector





WIMPS

- ARISE IN EXTENSIONS OF THE STANDARD MODEL
- ASSUMED TO BE STABLE: RELICS FROM THE BIG BANG
- WEAK-TYPE XSECTION GIVES NEEDED RELIC DENSITY

$$\Omega_\delta h^2 \approx \frac{10^{-27}}{\langle \sigma_{ann} v \rangle_{fr}} \text{ cm}^3 \text{ s}^{-1}$$

- MASS FROM FEW GeV TO FEW TeV
- MSSM CANDIDATE: LIGHTEST NEUTRALINO,

$$\chi_1^0 = N_1 \mathbf{B} + N_2 \mathbf{W}^3 + N_3 \mathbf{H}_1^0 + N_4 \mathbf{H}_2^0$$
- UED: LIGHTEST 'RUNG' IN THE KALUZA-KLEIN LADDER

SIMPZILLAS

- NON-THERMAL, NON-WEAKLY INTERACTING STABLE RELICS

A wealth of candidates from different theoretical models:

- dark baryons (primordial nucleosynthesis constraints)
- MACHOs – BHs, neutron stars, white/brown dwarfs... (microlensing constraints)
- neutrinos (mass constraint)
- primordial Black Holes (cosmological constraints)

Weakly Interacting Massive Particles

(LSPs from “x”MSSM, Kaluza-Klein modes...)

Non-weakly Interacting Supermassive particles
(Simpzillas)

- axions (too light+astrophysical constrains)
- many others
- ... + (alternative gravity theories)

DM-induced SM particles:

$$\kappa\kappa, \chi\chi, SS \rightarrow \left\{ \begin{array}{l} q\bar{q} \\ \ell^+\ell^- \\ W, Z, H \\ \dots \end{array} \right\} \rightarrow \nu, \gamma, e^+e^-, \bar{p}$$

Kaluza-Klein modes an additional useful channel:

$$\kappa\kappa \rightarrow \nu\nu$$

signature:

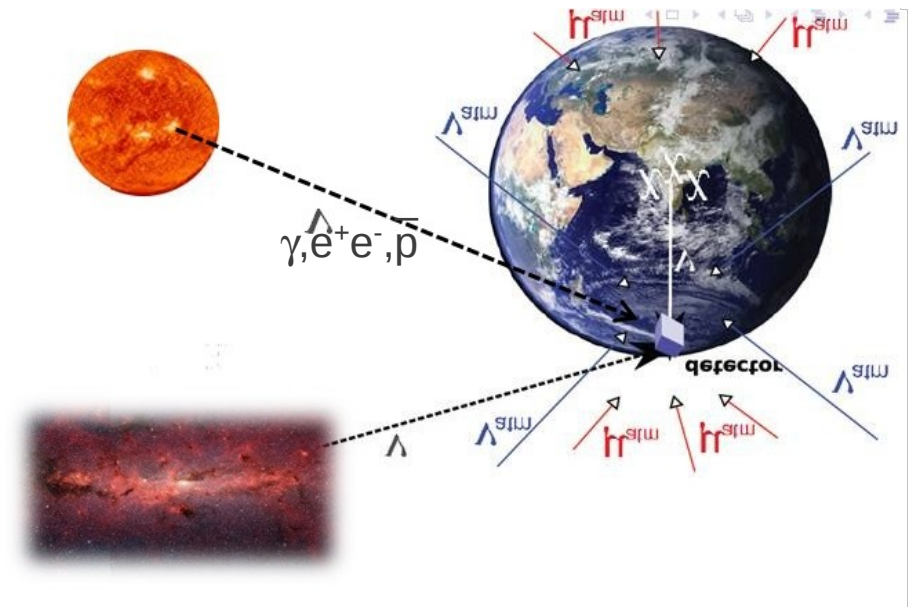
ν excess over background from Sun/Earth/Galactic Halo/nearby galaxies

A lot of physics uncertainties involved:

- relic density calculations
- DM distribution in the halo
- velocity distribution
- χ, K, S properties (MSSM/UED...)
- interaction of χ, K, S with matter (capture)
- self interaction (annihilation)

Look at objects where dark matter can have accumulated gravitationally over the evolution of the Universe

Sun, Earth, Galactic Halo/Center, dwarf spheroids



Atmospheric muons $\sim O(10^{10})$ events/year (downwards)

Atmospheric neutrinos $\sim O(10^4)$ events/year (all directions)

Solar Dark Matter Search with IceCube

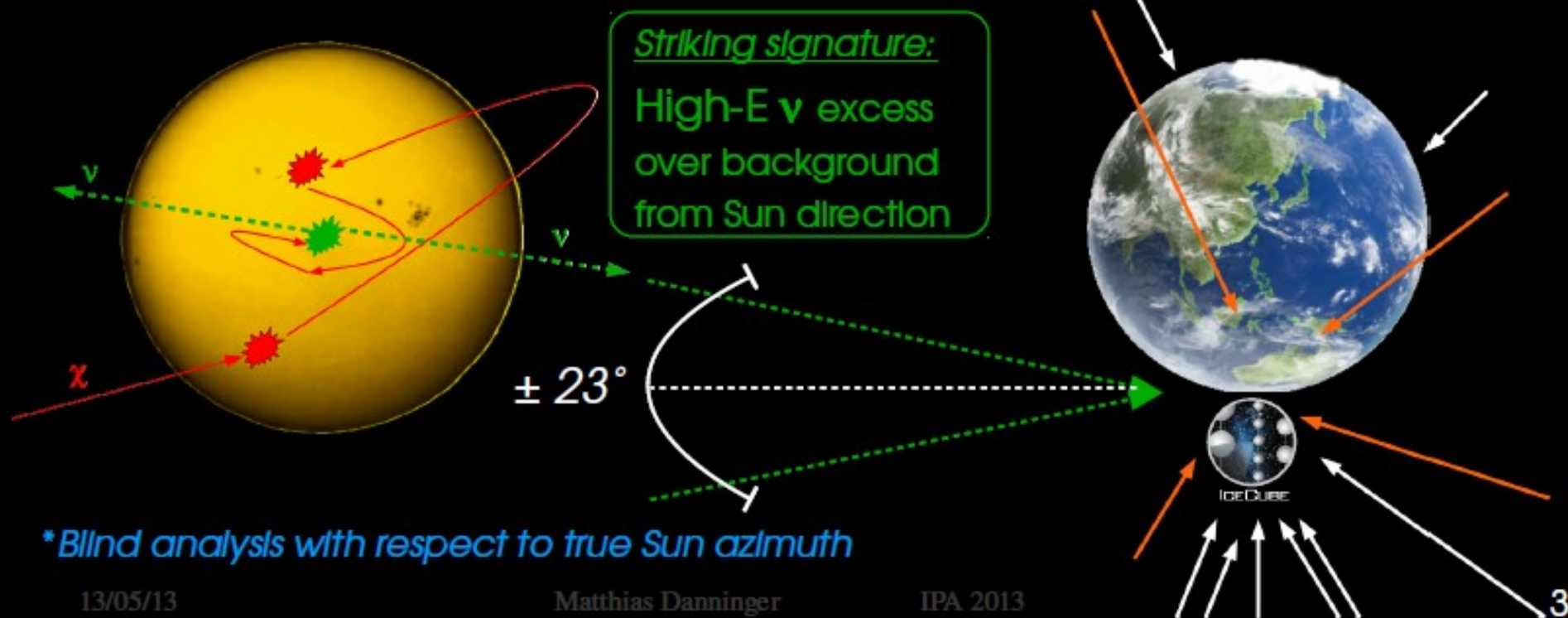


- All processes depend on WIMP mass
- Annihilation channel (branching ratios)
- Annihilation cross-section
- Capture (scattering)
→ Scattering cross-sections (SI & SD)

main analysis backgrounds:

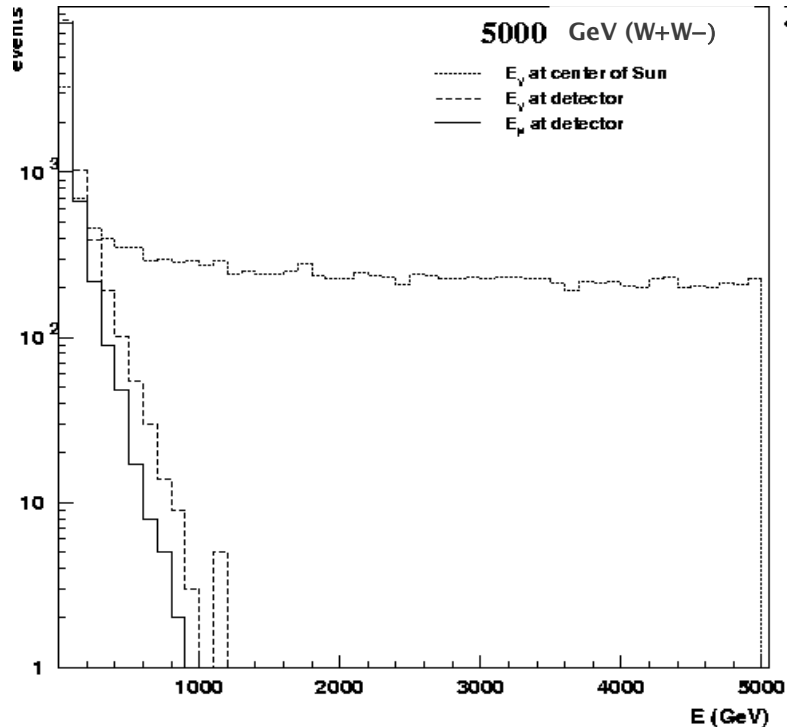
atm. $\nu \sim O(10^3 \text{ triggering events/day})$

atm. $\mu \sim O(10^8 \text{ triggering events/day})$

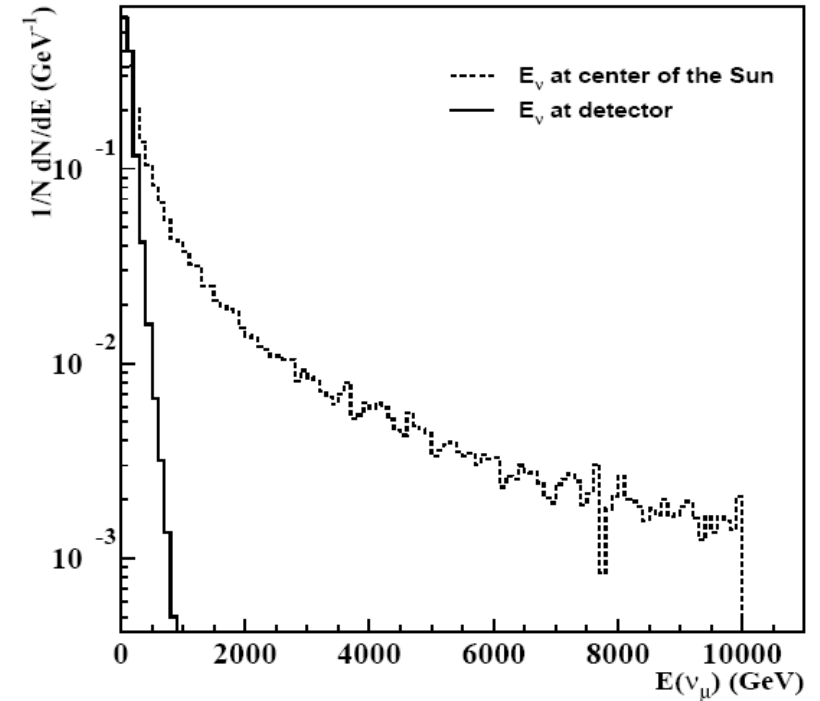


searches from the Sun: neutrino energies at the detector

5000 GeV Neutralino \rightarrow WW @ Sun

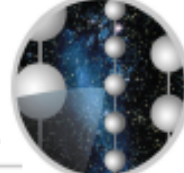


Simpzilla \rightarrow t \bar{t} @ Sun



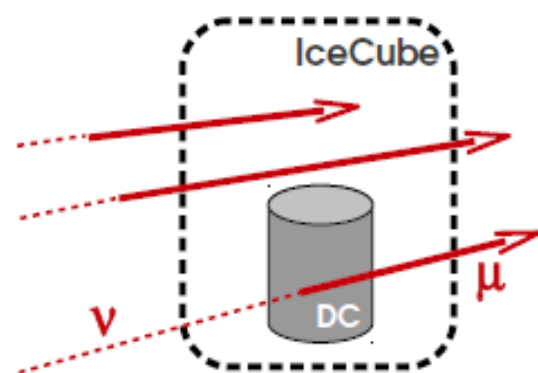
: Indirect dark matter searches from the **Sun** are a low-energy analysis in neutrino telescopes: even for the highest DM masses, we do not get muons above few 100 GeV

Not such effect for the Earth and Halo (no ν energy losses in dense medium)

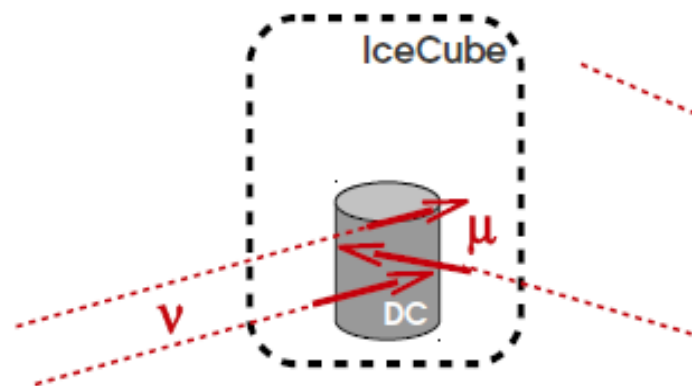


- Analysis for the **whole year!** Used 317 days livetime (151 days austral winter & 166 days austral summer)
- more than **60 billion** recorded events
- At final level **~25000** signal-like events in 3 independent samples
- With DeepCore, analysis reaches neutrino energies of **10-20 GeV**

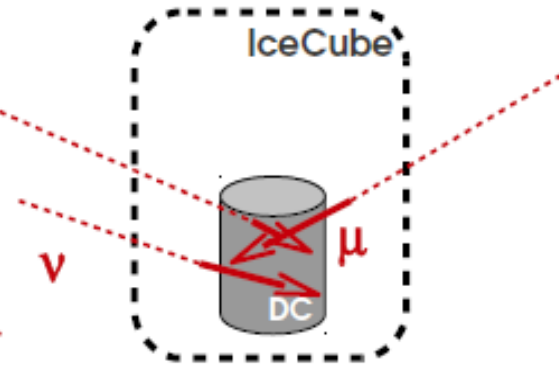
- **Up-going** ①
- No containment



- **Up-going** ②
- strong containment



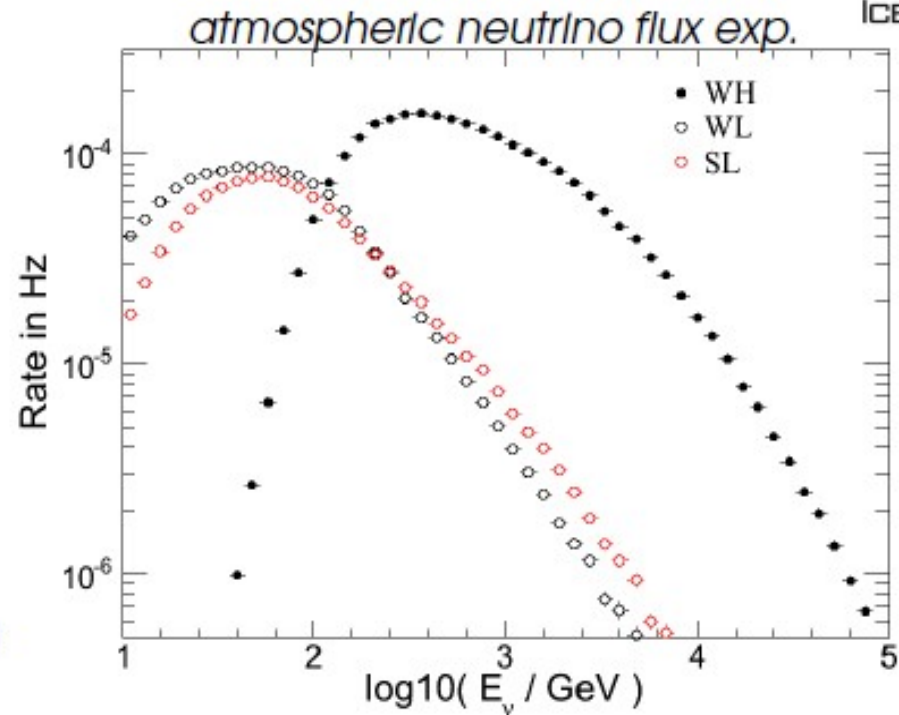
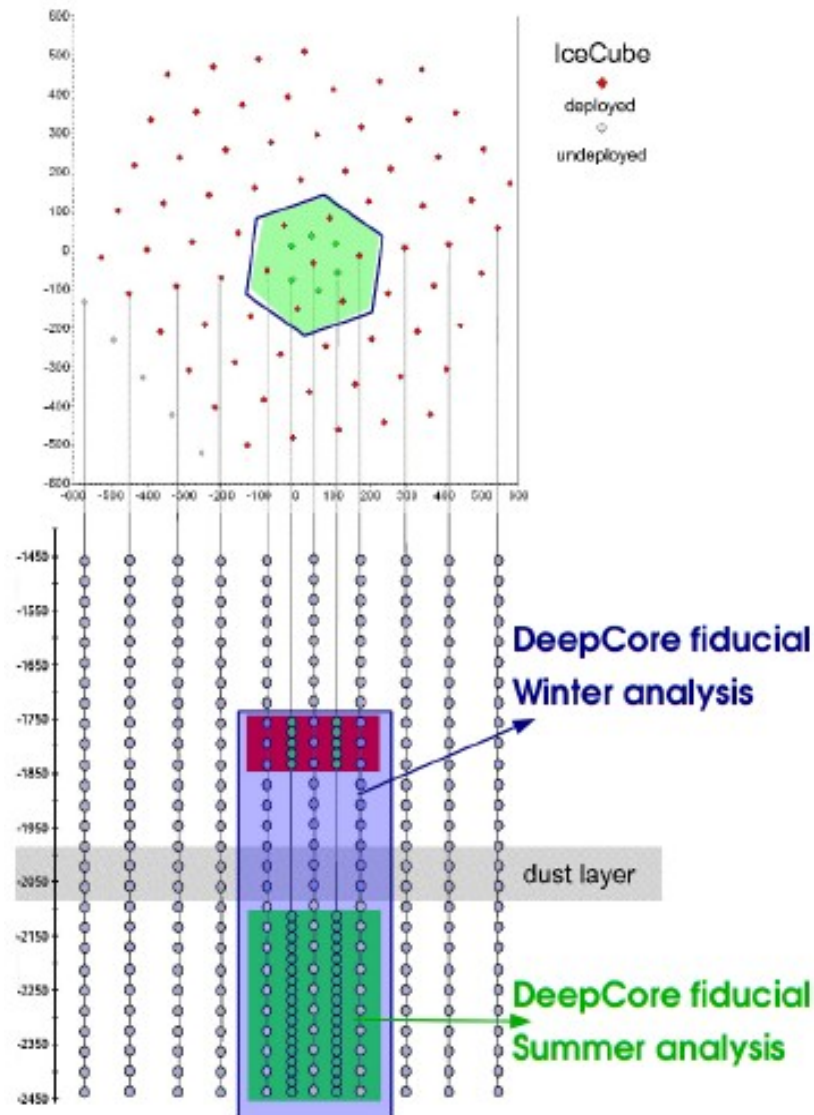
- **Down-going** ③
- strong containment



IceCube-79 string analysis details



ICECUBE



Signal event topology very different for low & high WIMP mass
→ find *geometrical cut* to split dataset into 2 non-overlapping datasets

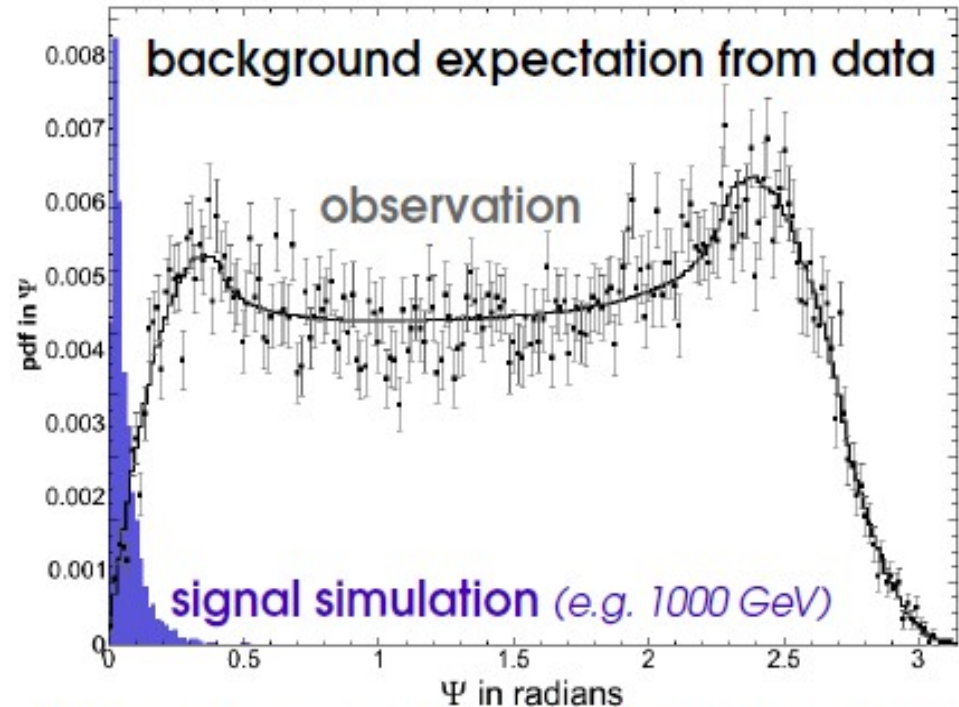


The observed angle to the Sun
is fitted with *signal* and
background pdf:s

Evaluate shape fit with log-likelihood rank (Feldman-Cousins) to construct confidence regions for the number of signal events μ_s

$$R(\mu) = \frac{\mathcal{L}(\mu)}{\mathcal{L}(\hat{\mu})}$$

where \mathcal{L} is the pdf product over the final sample



(Angle between event track & direction from the Sun)

$$f(\psi|\mu_s) = \frac{\mu_s}{n_{\text{obs}}} f_S(\psi) + \left(1 - \frac{\mu_s}{n_{\text{obs}}}\right) f_B(\psi)$$

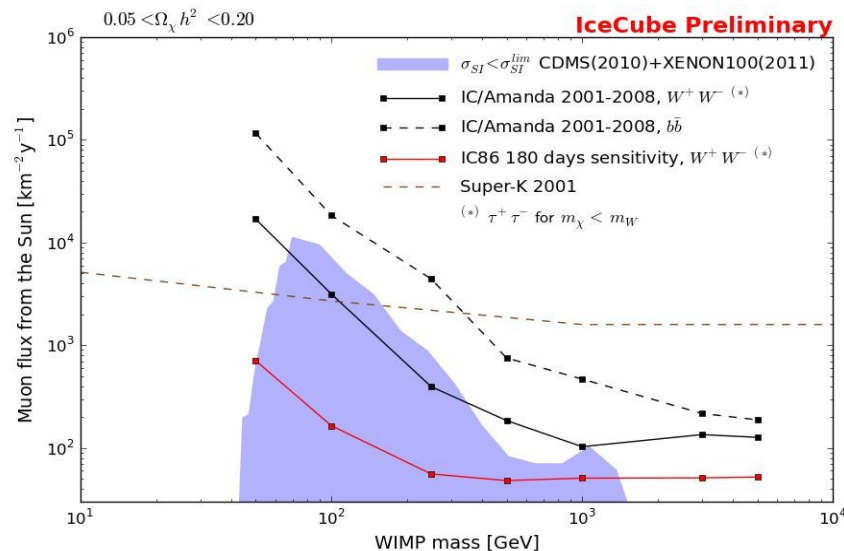
$$\mathcal{L}(\mu_s) = \prod_{i=1}^{n_{\text{obs}}} f(\psi_i|\mu_s)$$

$$\Psi_{\text{data}}, \Psi_{\text{bck}} \rightarrow N_{90} \longrightarrow \Gamma_{\nu\mu} \leq \frac{N_{90}}{V_{\text{eff}} \cdot t}$$

Experimentally obtained quantity:
allowed number of signal events still
compatible with background, at 90%
confidence level

Use model to convert
to a muon flux

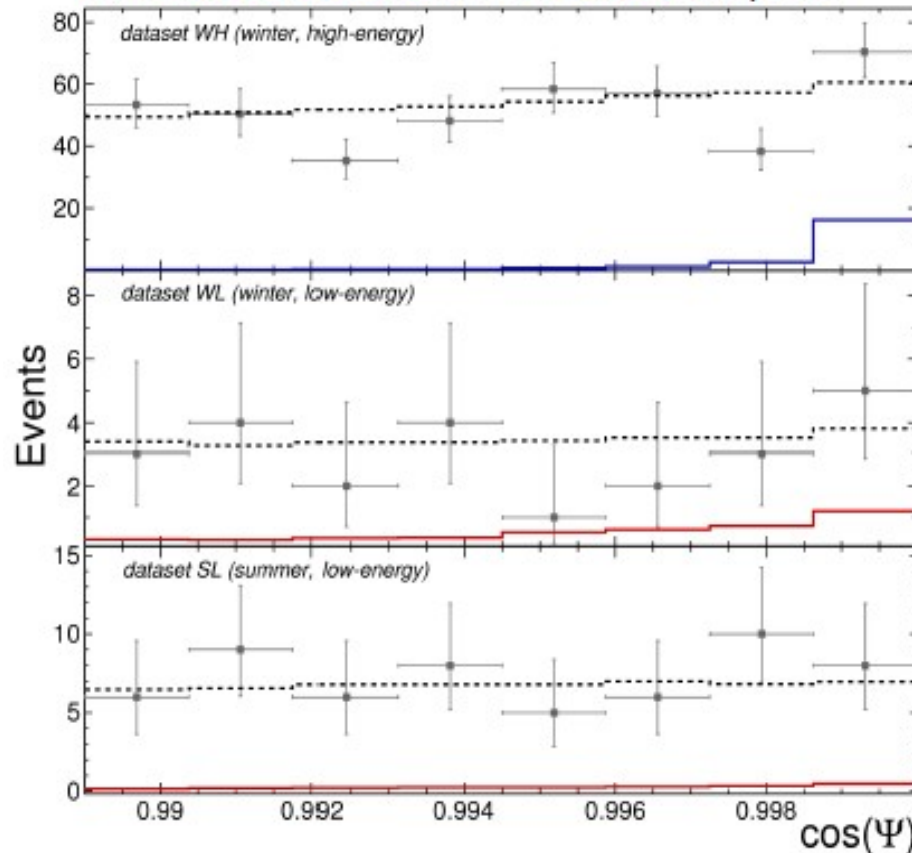
$$\Gamma_{\nu\mu}(m_\chi) = \Gamma_A \cdot \frac{1}{4\pi R_\oplus^2} \int_0^{m_\chi} \sum B_{\chi\bar{\chi} \rightarrow X} \left(\frac{dN_\nu}{dE_\nu} \right) \times \sigma_{\nu+N \rightarrow \mu + \dots}(E_\nu | E_\mu \geq E_{\text{thr}}) \rho_N dE_\nu \longrightarrow \phi_\mu(E_\mu \geq E_{\text{thr}}) = \frac{\Gamma_A}{4\pi D_\odot^2} \int_{E_{\text{thr}}}^\infty dE_\mu \frac{dN_\mu}{dE_\mu}$$



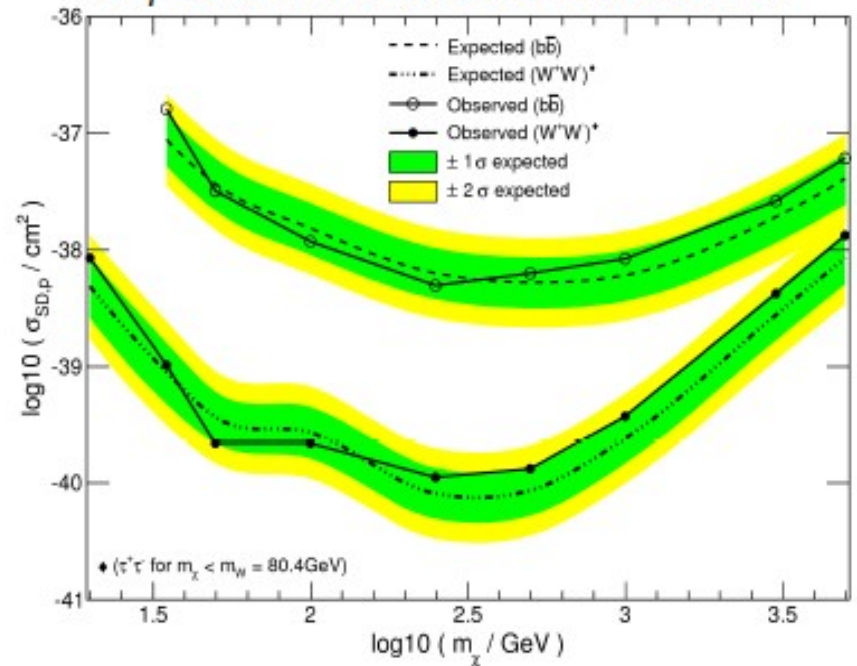
Solar WIMP search results (observed events)



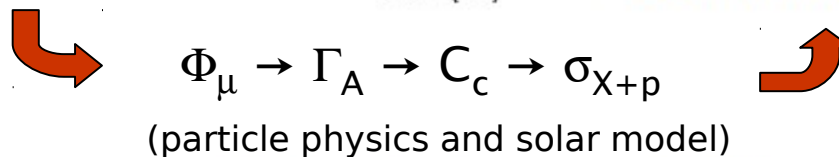
Unblinded events in different samples



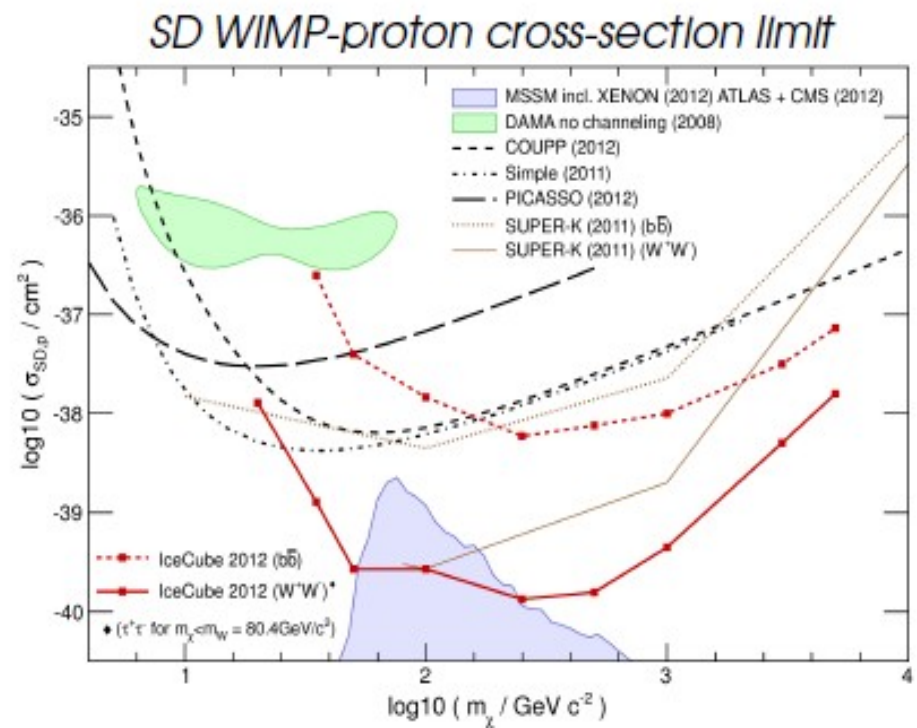
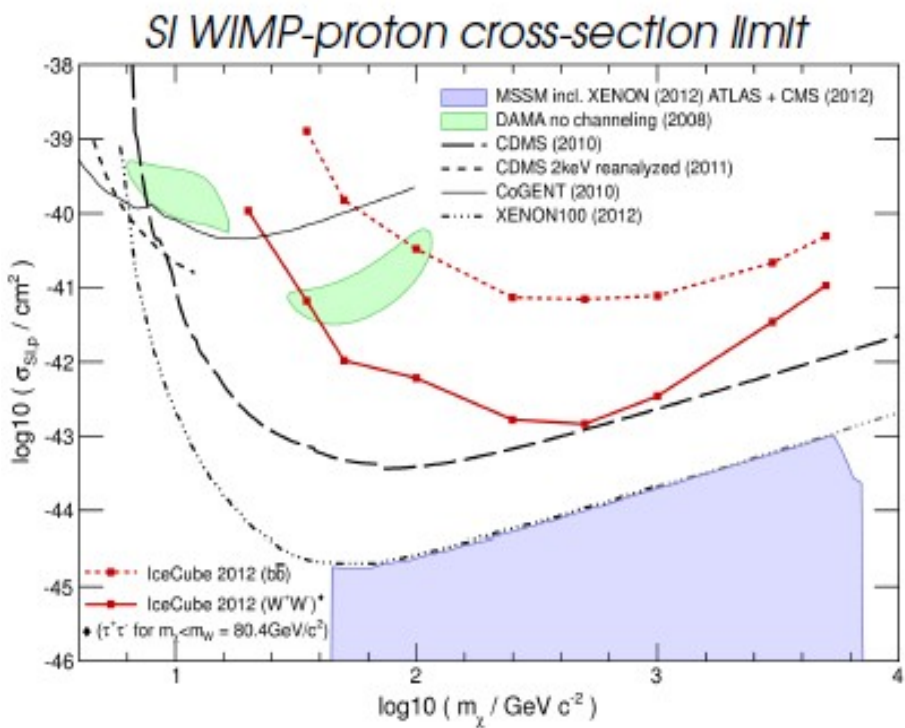
Expected sens. vs. observed result



details on systematic uncertainties,
see *PRL* **110** (2013) 131302



Solar WIMP search results (cross-section limit)



- most stringent SD cross-section limit for most models
- complementary to direct detection search efforts
- Different astrophysical & nuclear form-factor uncertainties

SEARCHES FROM THE SUN: COMPLEMENTARITY WITH COLLIDER RESULTS

Assume (ie. model dependent) effective quark-DM interaction,

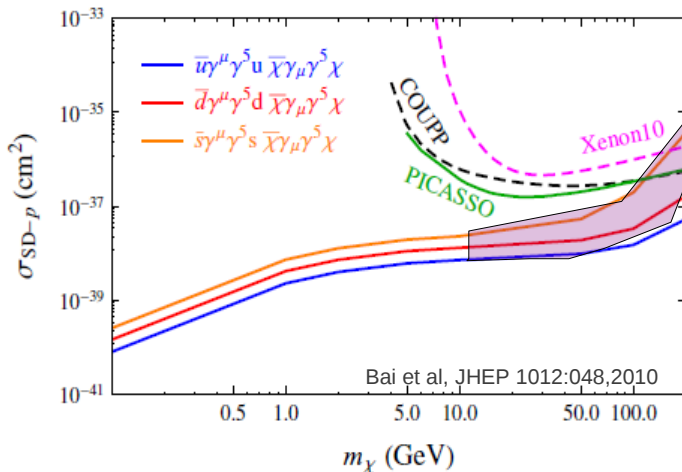
$$\lambda^2/\Lambda^2 (\bar{q}\gamma_5\gamma_\mu q)(\bar{\chi}\gamma_5\gamma^\mu\chi)$$

and look for monojets in $p\bar{p}$ collisions,

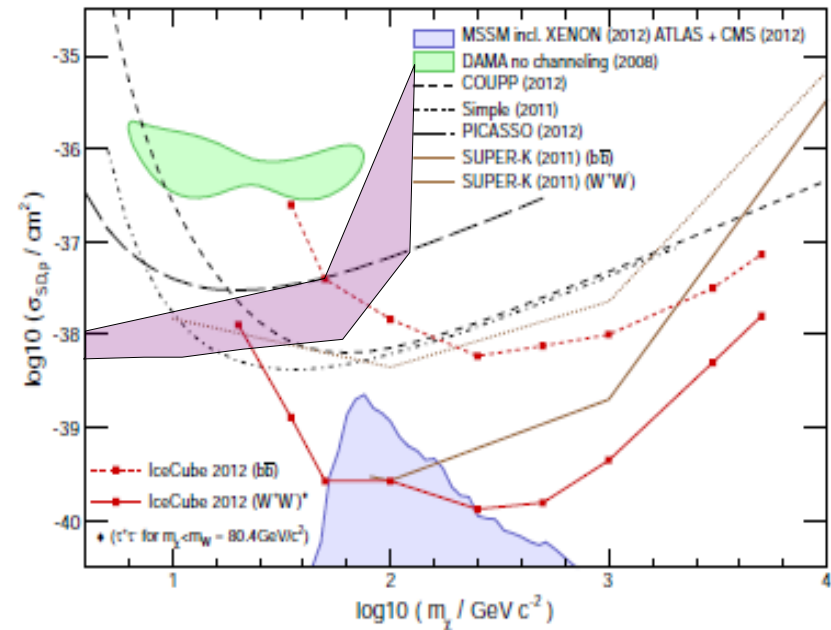
$$p\bar{p} \rightarrow \chi\bar{\chi} + \text{jet}$$

(as opposed to the SM process $p\bar{p} \rightarrow Z+\text{jet}$ and $p\bar{p} \rightarrow W+\text{jet}$)

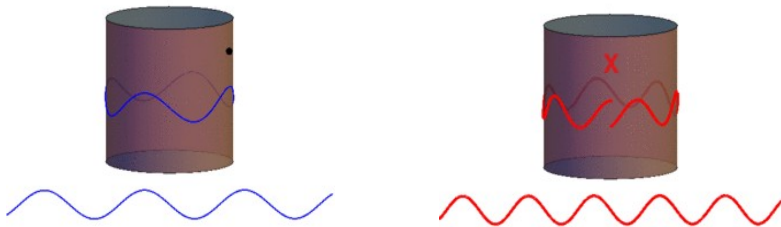
Constraints from monojet searches at the TeVatron:



90% CL **neutralino-p Xsection** limit



- Universal Extra Dimensions:**



$$n \frac{\lambda}{2} = 2\pi R, \quad n \frac{h}{2p} = 2\pi R \Rightarrow p = n \frac{h}{4\pi R}$$

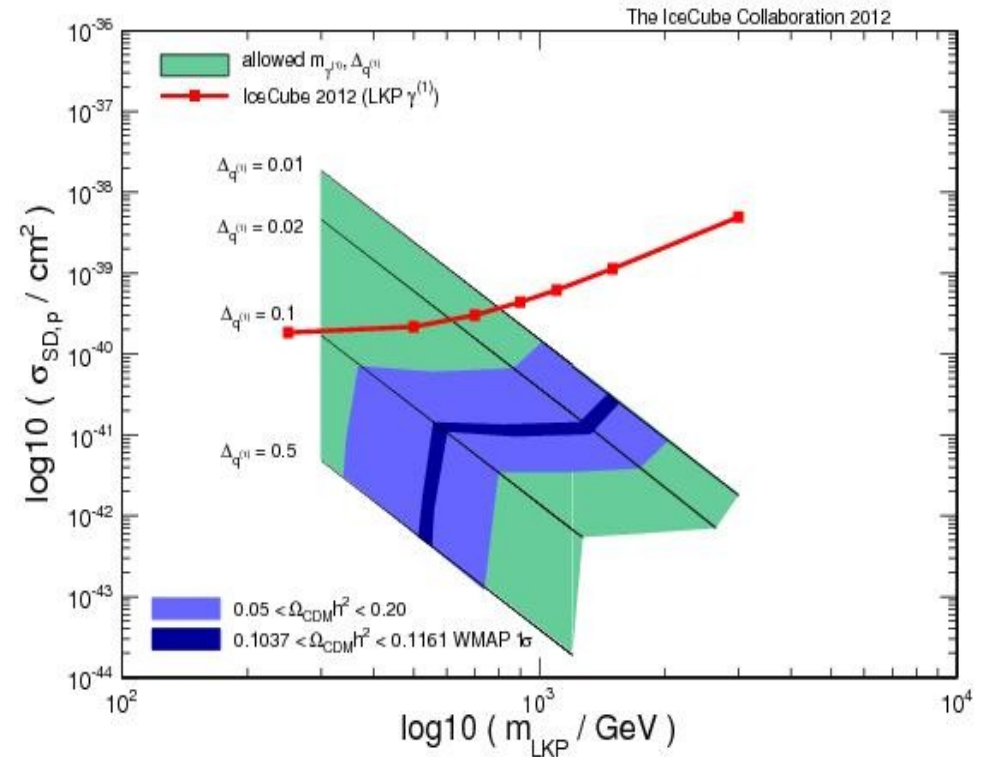
$$E^2 = p^2 c^2 + m_o^2 c^4 = n^2 \frac{1}{R^2} c^2 + m_o^2 c^4 = m_n^2 c^4$$

$$m_n^2 = \frac{n^2}{c^2 R^2} + m_o^2$$

$n=1 \rightarrow$ Lightest Kaluza-Klein mode, \mathbf{B}^1

good DM candidate

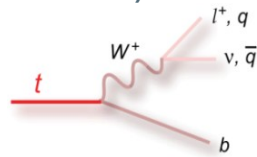
90% CL LKP-p Xsection limit vs LKP mass



SIMPZILLAS (Superheavy DM)

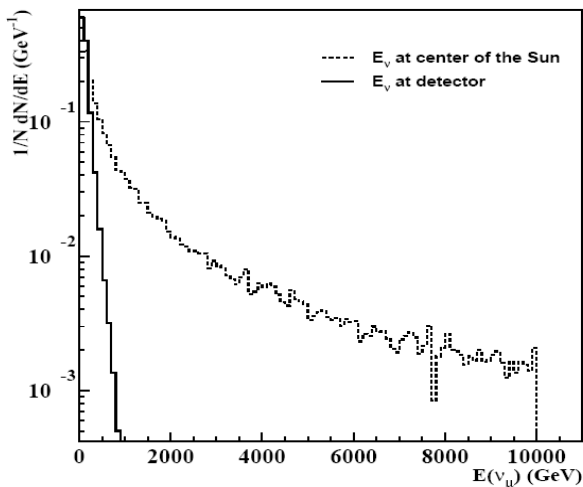
- Produced **non-thermally** at the end of inflation through vacuum quantum fluctuations or decay of the inflaton field
- strong Xsection (simply means non-weak in this context)
- m from $\sim 10^4$ GeV to 10^{18} GeV (no unitarity limit since production non thermal)

$$S+S \rightarrow t \bar{t}$$



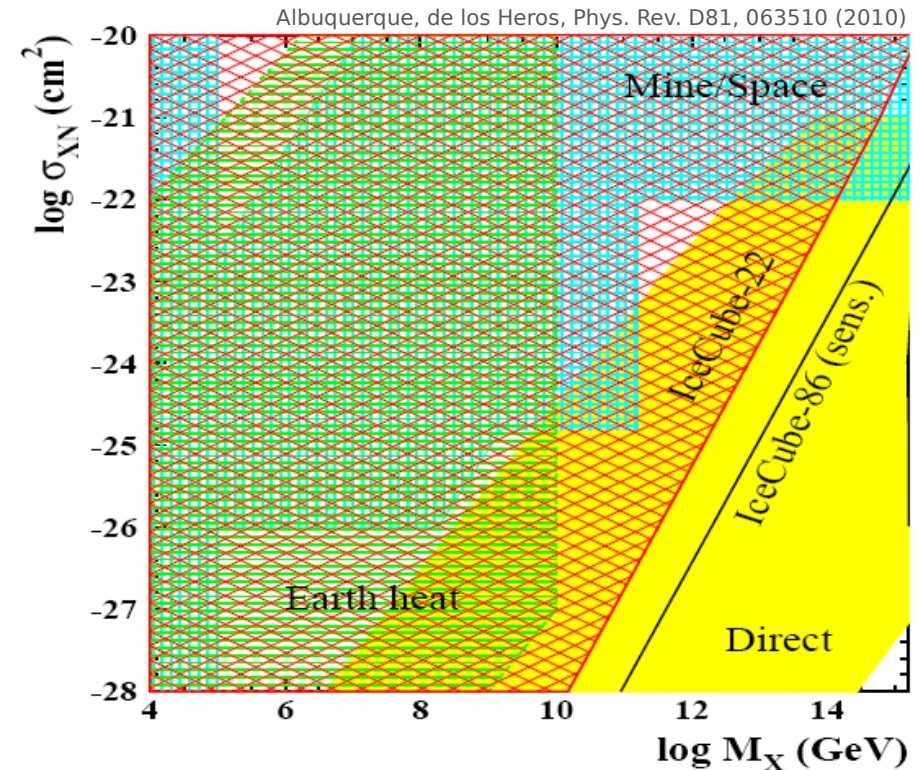
$2.8 \times 10^5 \sqrt{m_{X/12}}$ tops per annihilation

$$\frac{dN}{dE_\nu} \propto \frac{E_\nu + m_W}{\sqrt{(E_\nu + m_t)[(E_\nu + m_t)^2 - m_t^2][(E_\nu + m_W)^2 - m_W^2]}}$$

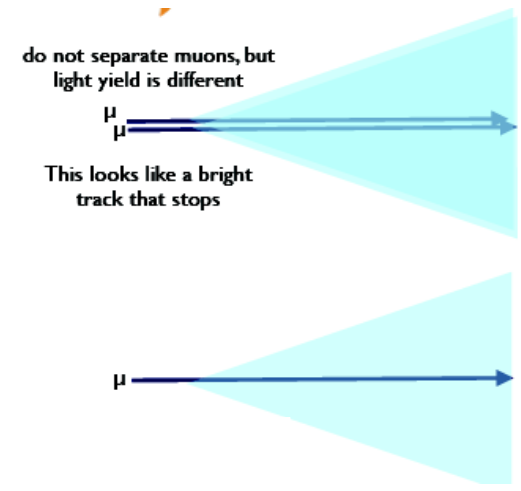
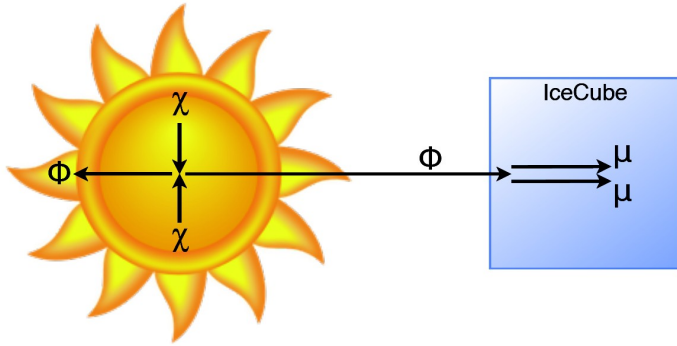


$$N_s(m_X, \sigma_{XN}) = N_t \cdot BR_W \cdot \Gamma_A(m_X, \sigma_{XN}) \cdot T \cdot \int \frac{dN_\nu}{dE} A_{eff} dE$$

90% CL simpzilla-p Xsection limit vs simpzilla mass



DM SEARCHES FROM THE SUN: SECLUDED DARK MATTER



Secluded dark matter

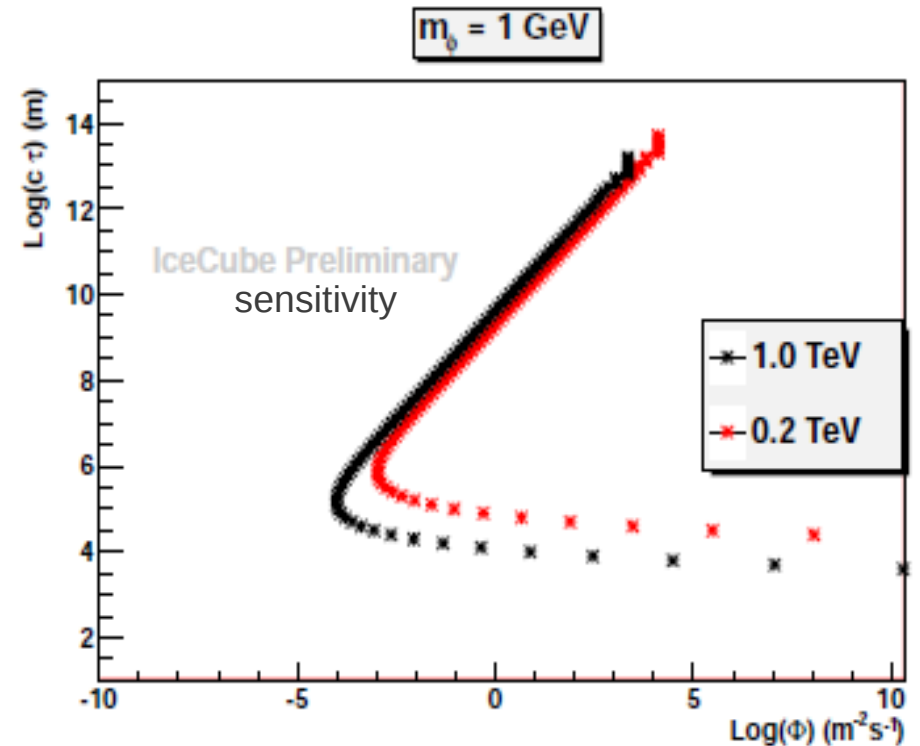
$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{DM}} + \mathcal{L}_{\text{mediator}}$$

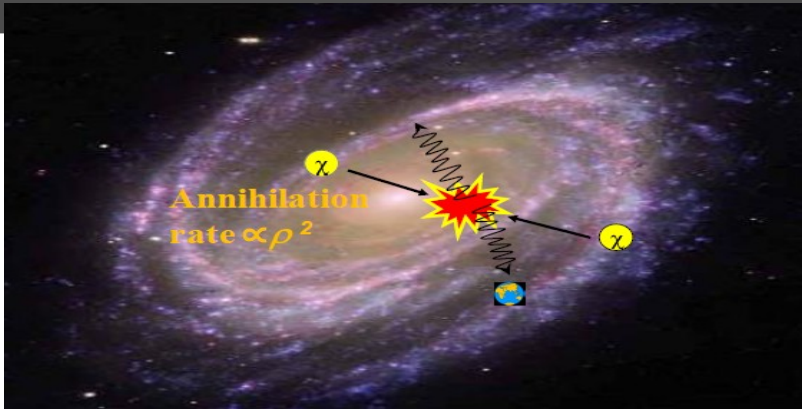
DM annihilates into mediator, $\chi\chi \rightarrow \phi\phi \rightarrow \text{SM}$
 with $m_\phi = \mathcal{O}(\text{GeV})$

ϕ is long lived, escapes the Sun and decays into $\mu^+\mu^-$ in or near the detector

→ **signature**: two closely separated muon tracks ($\sim 1\text{m}$)

look for stopping pairs of tracks in order to further reduce the background.





Look for an excess of events in the on-source region w.r.t. the off-source:

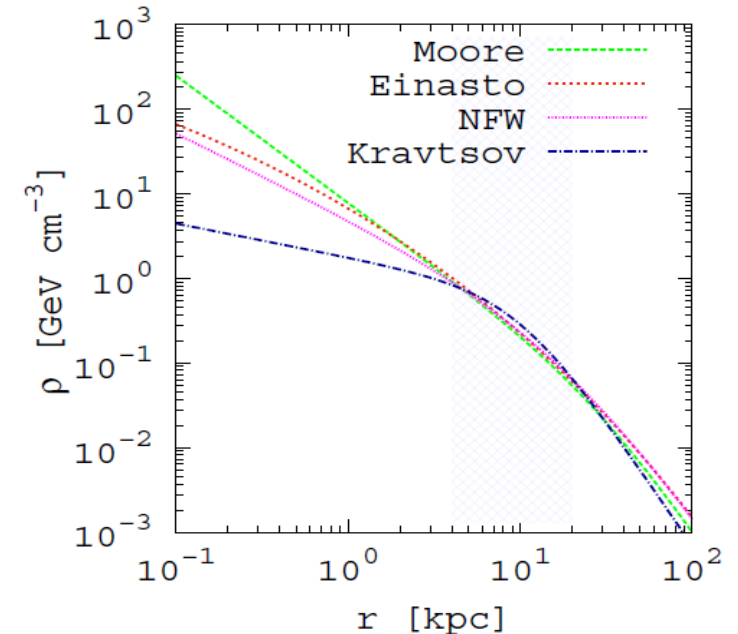
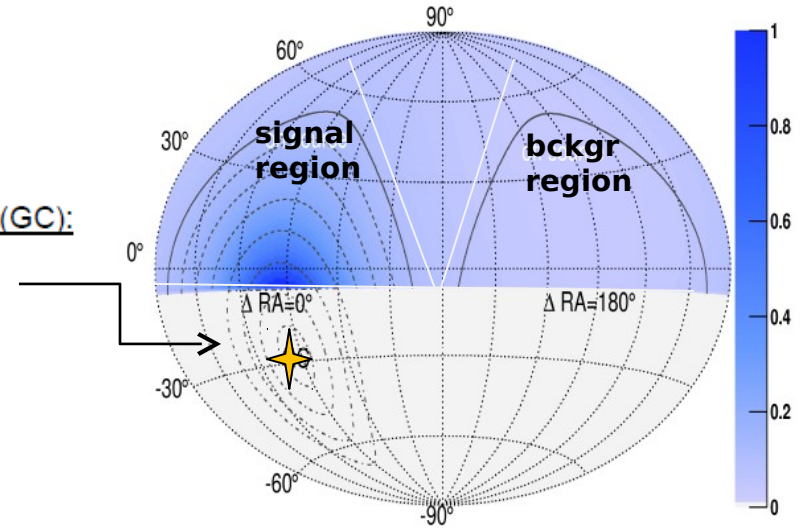
IC22: observed on-source: 1367 evts
observed off-source: 1389 evts

Need expected neutrino flux from SUSY and halo model. Limit on the self annihilation cross section:

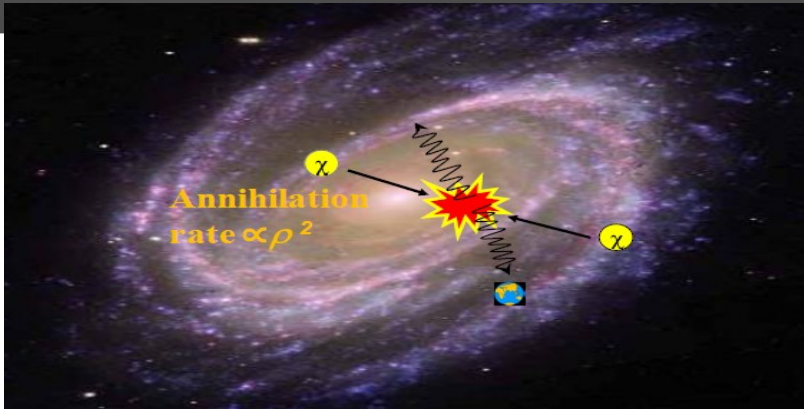
$$\phi_\nu = \frac{dN}{dE dA_{\text{eff}} dt d\Omega} = \frac{1}{2} \frac{1}{4\pi} \langle \sigma v \rangle J_\Omega R_{SC} \frac{\rho_{SC}^2}{m_\chi^2} \frac{dN_\nu}{dE}$$

Galactic Center (GC):

$$\begin{aligned} \text{R.A.} &= 277^\circ \\ \Theta &= -28^\circ \end{aligned}$$

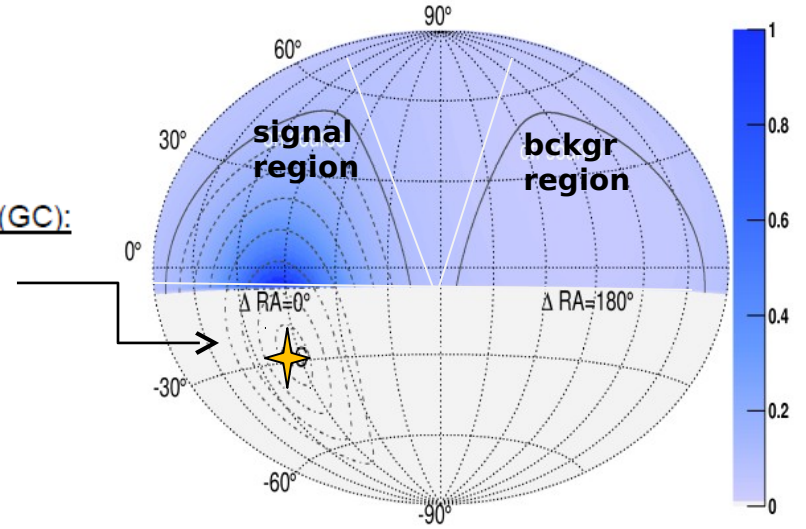


DM SEARCHES FROM THE GALACTIC HALO



Galactic Center (GC):

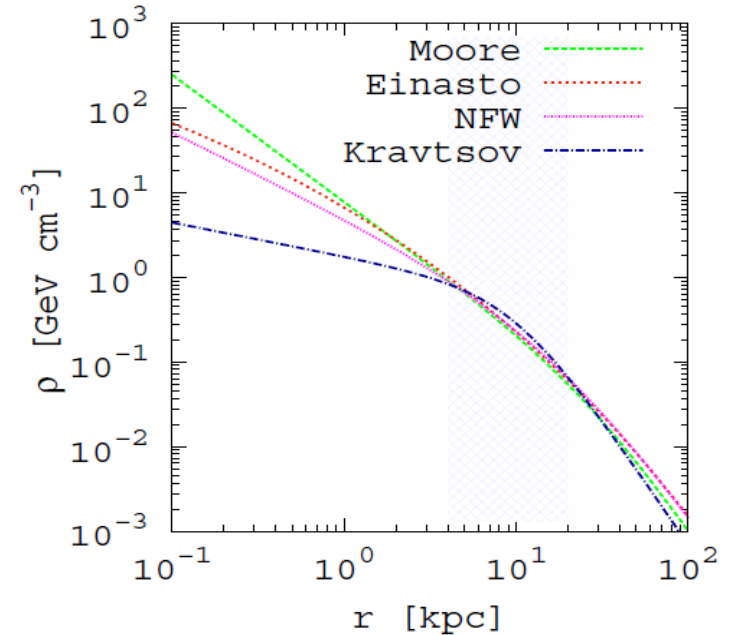
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measure **CONSTRAIN** halo model particle physics

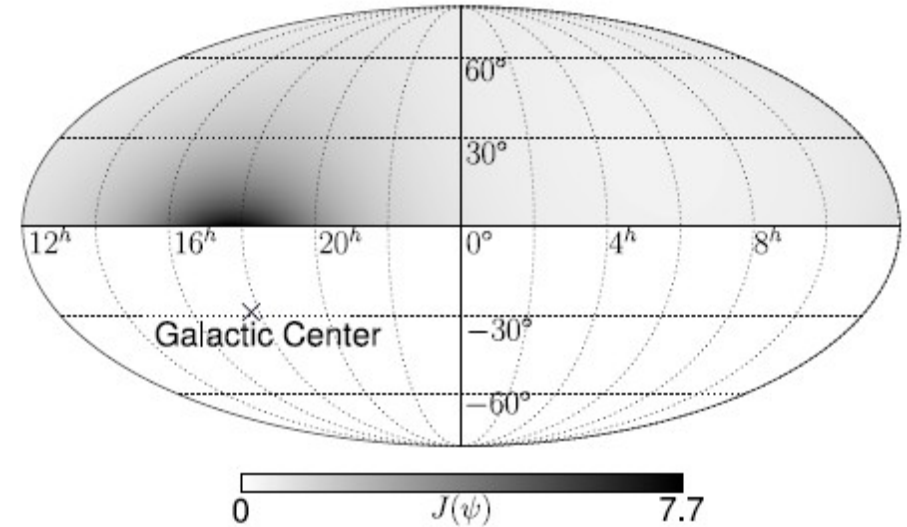
line of sight (los) integral

DM SEARCHES FROM THE GALACTIC HALO: MULTIPOLE ANALYSIS

$$Y_\ell^m(\theta, \phi) = \sqrt{\frac{(2\ell+1)(\ell-m)!}{4\pi(\ell+m)!}} P_\ell^m(\cos(\theta)) \exp(im\phi)$$

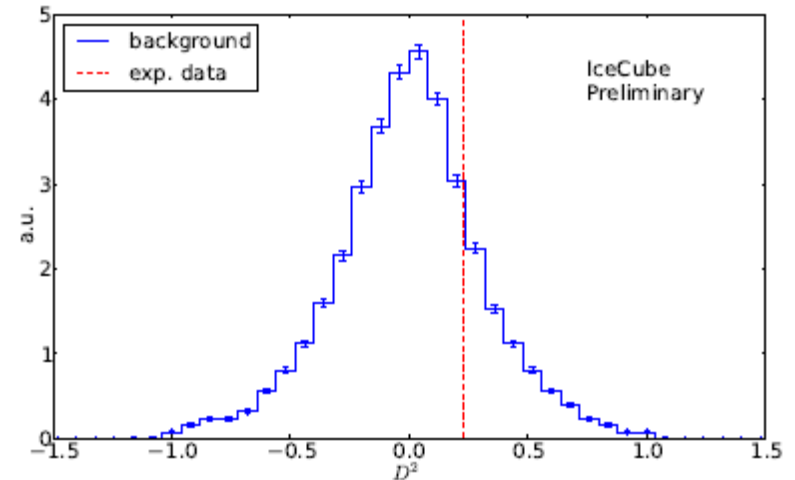
$$a_\ell^m = \int_\Omega d\Omega f(\theta, \phi) Y_\ell^{m*}(\theta, \phi)$$

$$f(\theta, \phi) = \sum_{i=1}^{N_V} \delta(\theta - \theta_i) \cdot \delta(\phi - \phi_i)$$

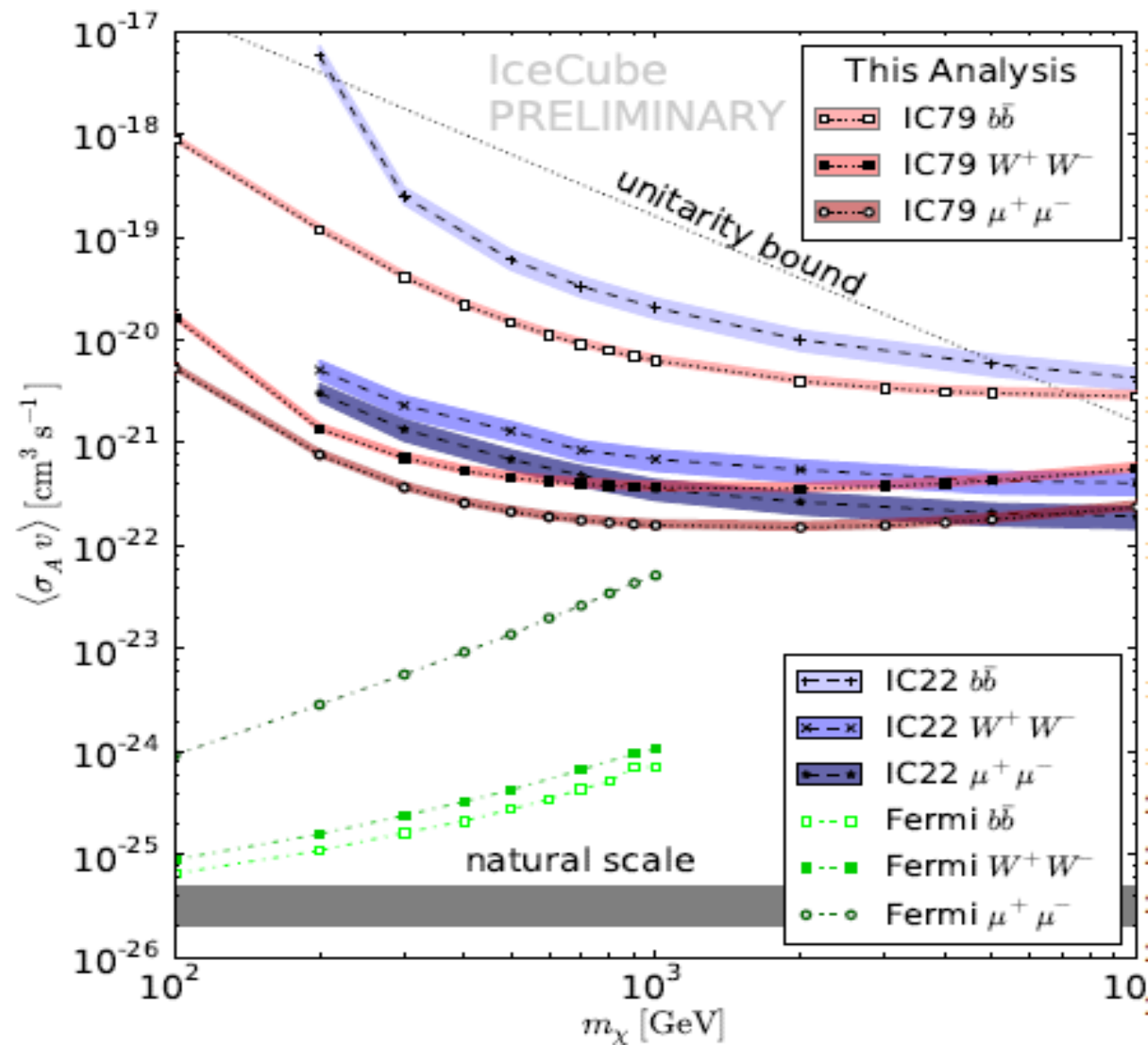


$$D^2 = \frac{1}{\sum w_\ell^m} \sum_{\ell=1}^{\ell_{\max}} \sum_{m=1}^{\ell} \text{sign}(\mathcal{A}_\ell^m) w_\ell^m \left(\frac{\mathcal{A}_\ell^m - \langle \mathcal{A}_{\ell, \text{atm}}^m \rangle}{\sigma(\mathcal{A}_{\ell, \text{atm}}^m)} \right)^2$$

$$\langle \sigma_A v \rangle_{90} = \frac{8\pi m_\chi^2}{R_{\text{SC}} \rho_{\text{SC}}^2} \frac{1}{T_{\text{live}}} \frac{1}{\int \int J(\psi) A_{\text{eff}} \frac{d\mathcal{N}_V}{dE} dE d\Omega} N_{90}$$

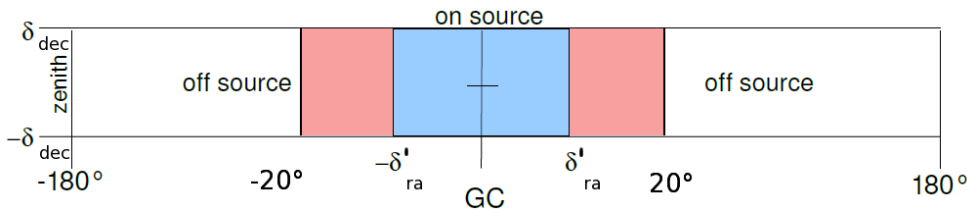


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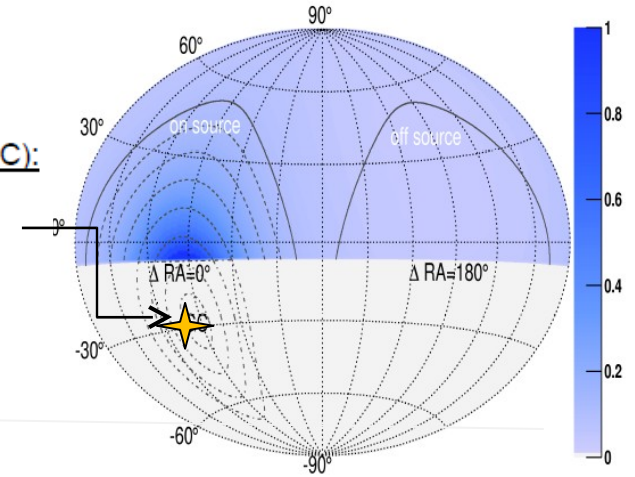


not to scale



Galactic Center (GC):

$$\begin{aligned} \text{R.A.} &= 277^\circ \\ \text{Dec} &= -28^\circ \end{aligned}$$

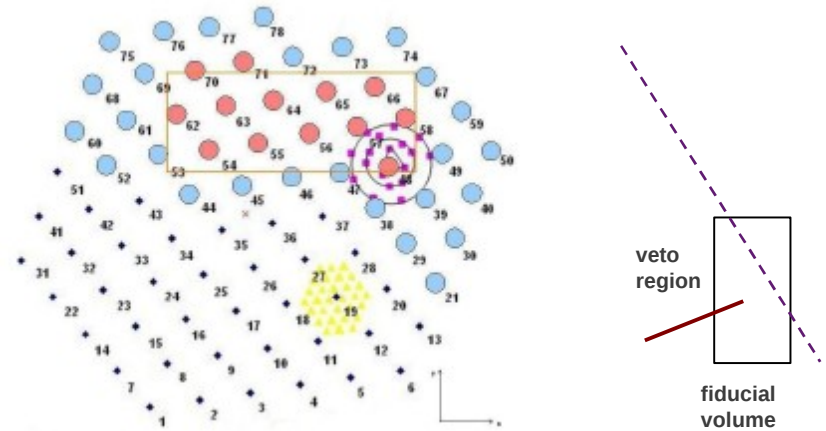


Look for an excess of events in the on-source region w.r.t. the off-source

on-source region below the horizon: need to veto downgoing μ s.

Use central strings of detector as fiducial volume, surrounding layers as veto. Only from IC40 this is possible.

IC40: observed on-source: 798842 evts
observed off-source: 798819 evts



$$\frac{d\Phi}{dE} = \frac{\langle \sigma_{AV} \rangle}{2} J(\psi) \frac{R_{sc}^2 \rho_{sc}^2}{4\pi m_\chi^2} \frac{dN}{dE}$$

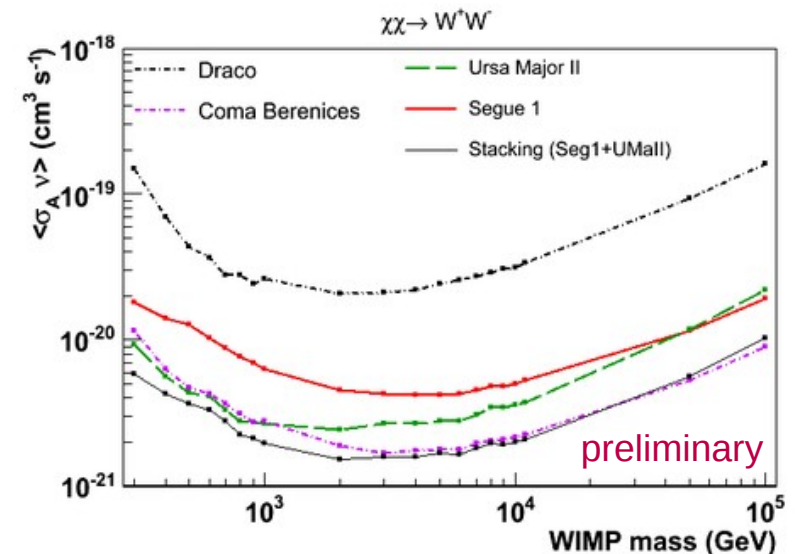
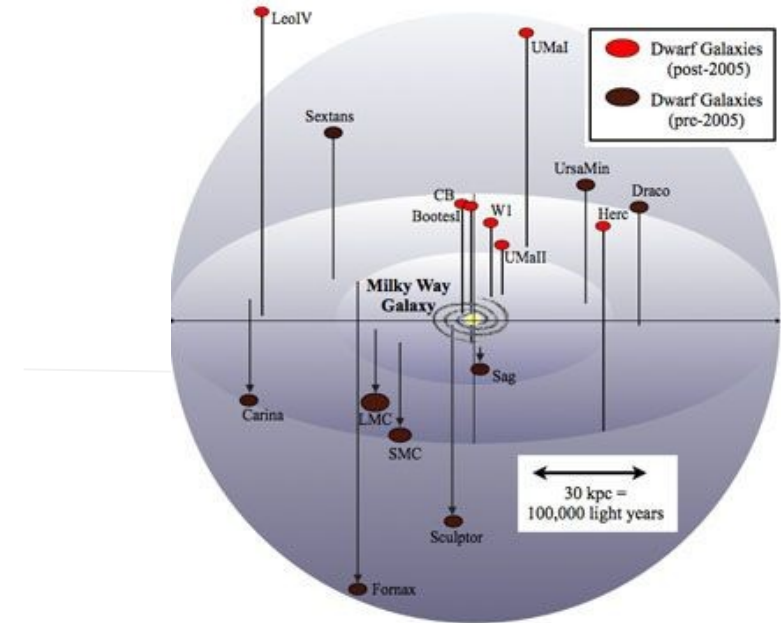
Measure Constrain Halo SUSY

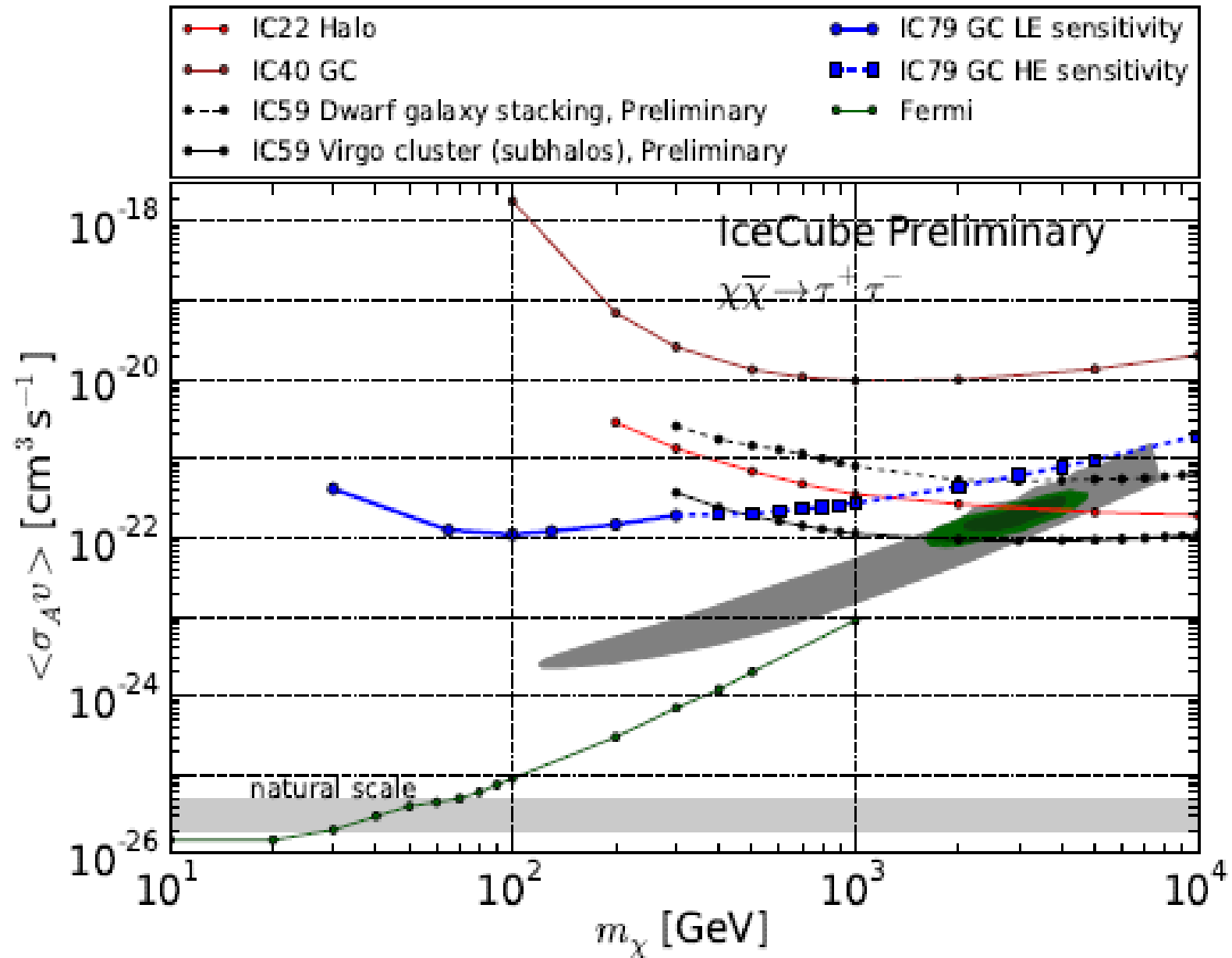
Same strategy as in the galactic halo analysis:

- dwarf galaxies: high mass/light ratio
- → high concentration of DM in the halos
- known location. Distributed both in the north and southern sky.
 - Point-like search techniques: stacking
 - known distance -> determination of absolute annihilation rate if a signal is detected
- same expected neutrino spectra as for the galactic center/halo
- IceCube results from various sources

Same strategy as in the galactic halo analysis:

$$\phi_\nu = \frac{dN}{dE dA_{\text{eff}} dt d\Omega} = \frac{1}{2} \frac{1}{4\pi} \langle \sigma v \rangle J_\Omega R_{SC} \frac{\rho_{SC}^2}{m_\chi^2} \frac{dN_\nu}{dE}$$



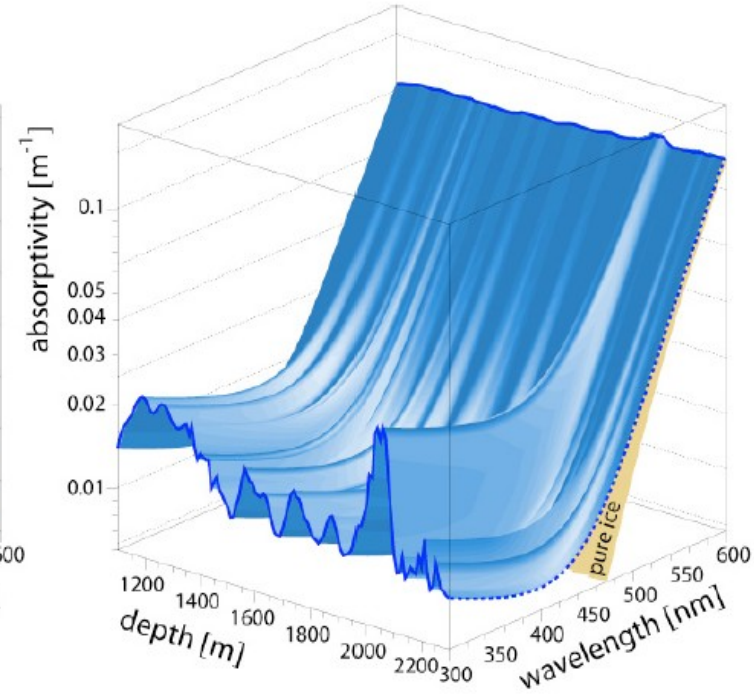
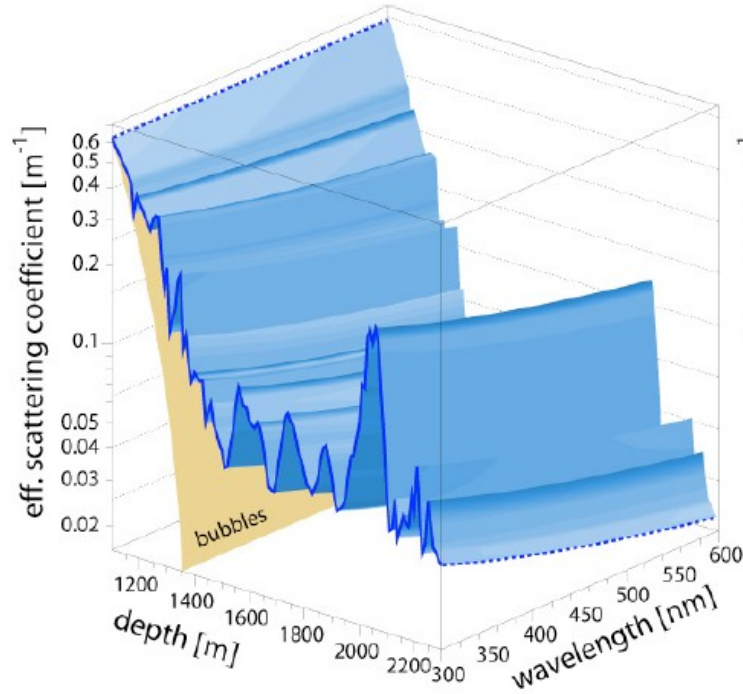


- we have 1 km³ of ice instrumented with optical modules
- we can detect flavours (muon tracks, e/ τ cascades)
- we can define through-going, starting and contained tracks
- we cover a wide neutrino energy range, from O(10) GeV to PeV
- we can look at all the sky (at once and continuously)

..... if you have a model of exotic physics that involves neutrinos, we can probe it



- Depth dependence of λ_{eff} and λ_{abs} from *in situ* LEDs
- Ice below 2100 m in DeepCore fiducial region very clear
 - $\langle \lambda_{\text{eff}} \rangle \sim 47 \text{ m}$, $\langle \lambda_{\text{abs}} \rangle \sim 155 \text{ m}$



- Constant temperature $\sim -35\text{C}$


Compare observed number of events n and predicted number θ for each model, taking into account error σ_ϵ on acceptance:

$$\mathcal{L}_{\text{num}}(n|\theta) = \frac{1}{\sqrt{2\pi}\sigma_\epsilon} \int_0^\infty \frac{(\epsilon\theta)^n e^{-\epsilon\theta}}{n!} \exp\left[-\frac{1}{2}\left(\frac{1-\epsilon}{\sigma_\epsilon}\right)^2\right] d\epsilon. \quad (1)$$

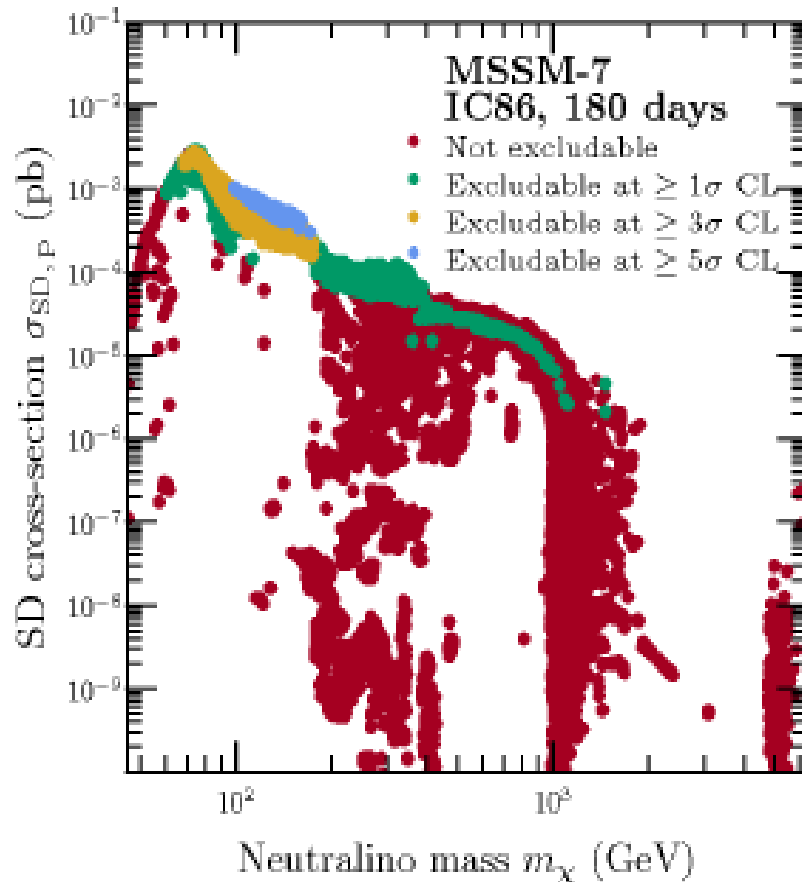
From this, construct a modified p -value as

$$p(n) = \frac{p_{\text{signal+BG}}(n)}{p_{\text{BG}}(n)} = \frac{\sum_{n_i \leq n} \mathcal{L}_{\text{num}}(n_i|\theta_{\text{signal+BG}})}{\sum_{n_i \leq n} \mathcal{L}_{\text{num}}(n_i|\theta_{\text{BG}})} \quad (2)$$

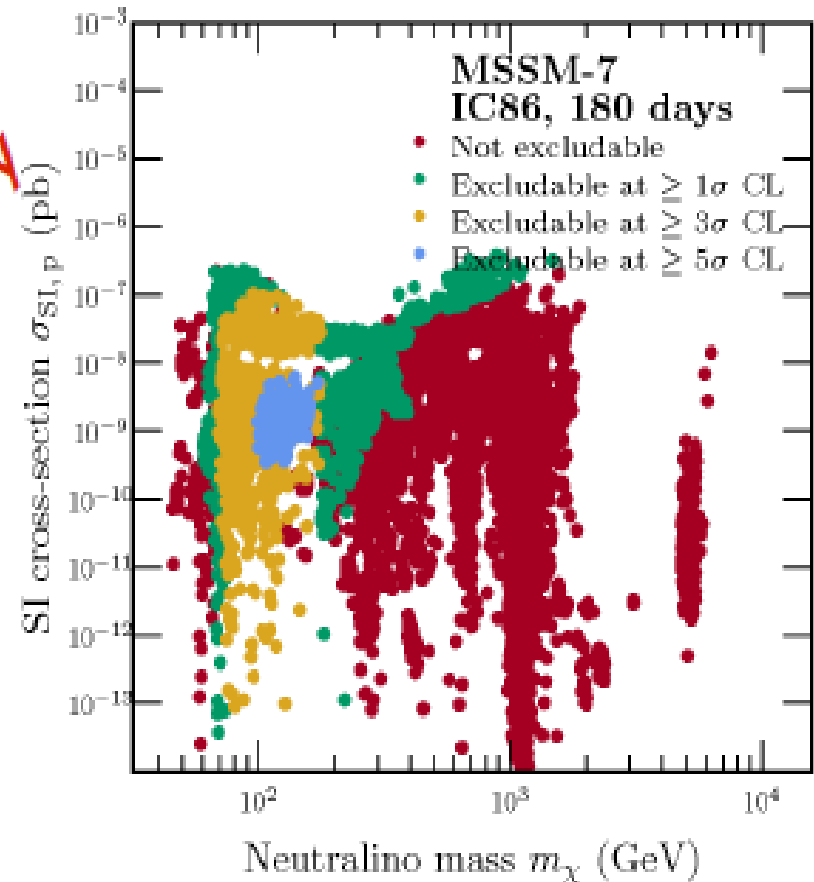
Can now say immediately, for a single point, that the point is excluded at a confidence level of $1 - p$.

\implies **model exclusion** exercise – IN/OUT analysis 

Assuming preliminary (conservative) estimate of IC-86 effective area




preliminary



- ✗ Only partial goodness of fit, no measure of convergence, no idea how to generalise to regions or whole space.
- ✗ Frequency/density of models in IN/OUT scans means essentially nothing.

Full unbinned likelihood with number (\mathcal{L}_{num}), spectral ($\mathcal{L}_{\text{spec}}$) and angular (\mathcal{L}_{ang}) parts

$$\mathcal{L} = \mathcal{L}_{\text{num}}(n|\theta_{\text{signal+BG}}) \prod_{i=1}^n \mathcal{L}_{\text{spec},i} \mathcal{L}_{\text{ang},i} \quad (3)$$

with **Number of lit channels (energy estimator)** theory 

$$\mathcal{L}_{\text{spec},i}(N_i, \Xi) = \frac{\theta_{\text{BG}}}{\theta_{\text{signal+BG}}} \frac{dP_{\text{BG}}}{dN_i}(N_i) + \frac{\theta_{\text{signal}}}{\theta_{\text{signal+BG}}} \int_0^{\infty} E_{\text{disp}}(N_i|E'_i) \frac{dP_{\text{signal}}}{dE'_i}(E'_i, \Xi) dE'_i \quad (4)$$

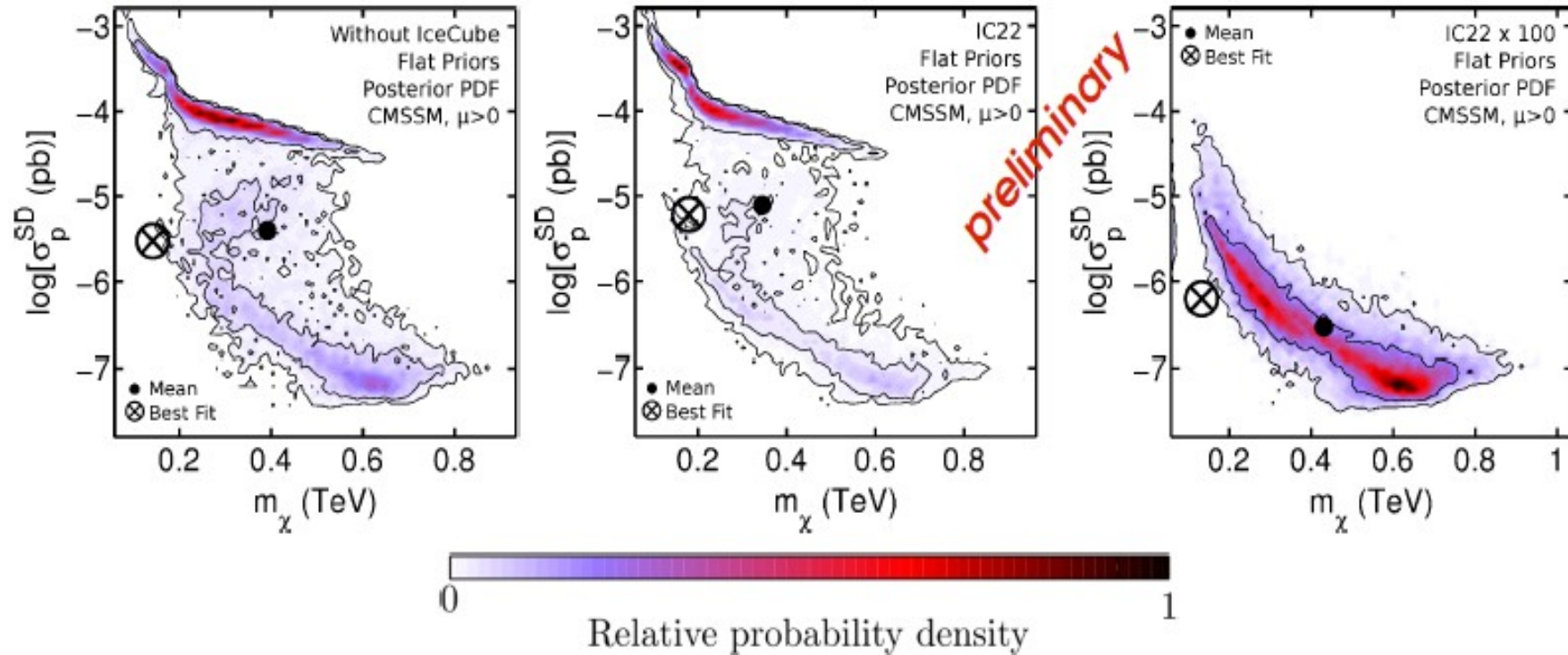
and **SUSY parameters**

$$\mathcal{L}_{\text{ang},i}(\cos \phi_i) = \frac{\theta_{\text{BG}}}{\theta_{\text{signal+BG}}} \frac{dP_{\text{BG}}}{d \cos \phi_i}(\cos \phi_i) + \frac{\theta_{\text{signal}}}{\theta_{\text{signal+BG}}} \text{PSF}(\cos \phi_i|1) \quad (5)$$

Event arrival angle

Example of what we want in the end (work for IC86 in progress):

SD nuclear scattering cross-section in the CMSSM with IceCube-22 events



- x Contours indicate 1σ and 2σ credible regions
- x Shading+contours indicate relative probability only, not overall goodness of fit
- x Scans performed with modified SuperBayes 1.5.1 and unreleased DarkSUSY