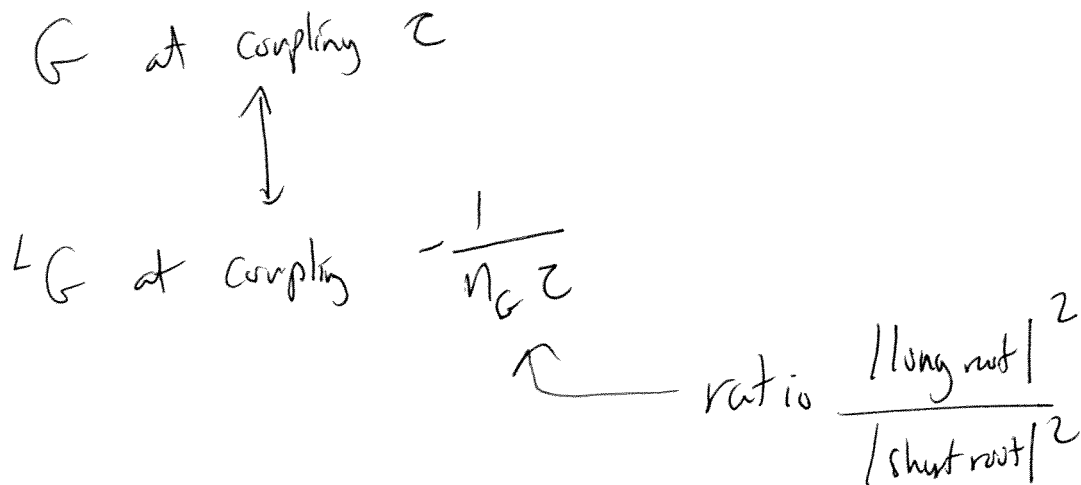


12 August 2010
Y. Tachikawa

The work of Argyres and Seiberg

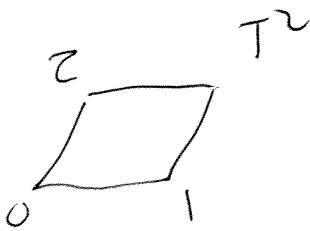
$\mathcal{N}=4$ SUSY Yang-Mills

S-duality Montonen-Olive (1977)



6d theory of A_n, D_n, E_n

Vafa (1999)



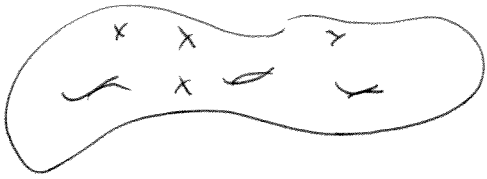
$\mathcal{N}=2$ SUSY theories

$\mathcal{N}=2$ $SU(2)$ 4 flavors \leftarrow self dual (1994)

Argyres-Seiberg (2007)
Argyres-Wittig



$\mathcal{N}=2, SU(3)$ with 6 flavors
 \leftarrow very strange dual theory

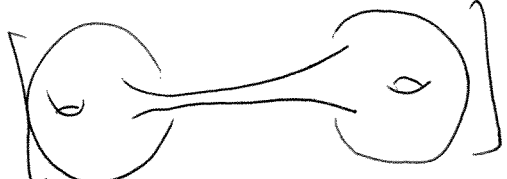


6d theory $g = A, D, E$




Gaiotto (2009)

6d theory $g = A, D, E$

4d SCFT_g :  \longrightarrow 4d SCFT_g []
(can associated 4d SCFT)

SCFT_g [] = SCFT_g [] + SCFT_g []

$g = \text{"length"}$ eg. 4-punctured sphere 

Coupled to G at $\tau = \log g$

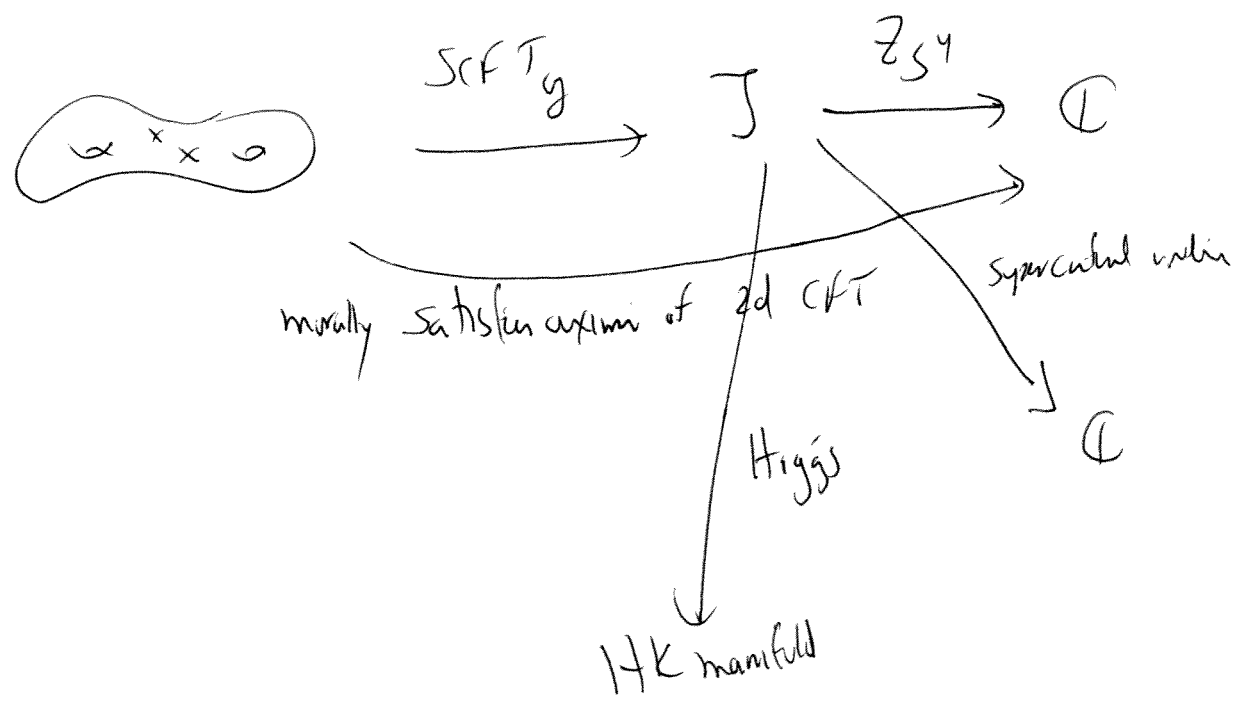
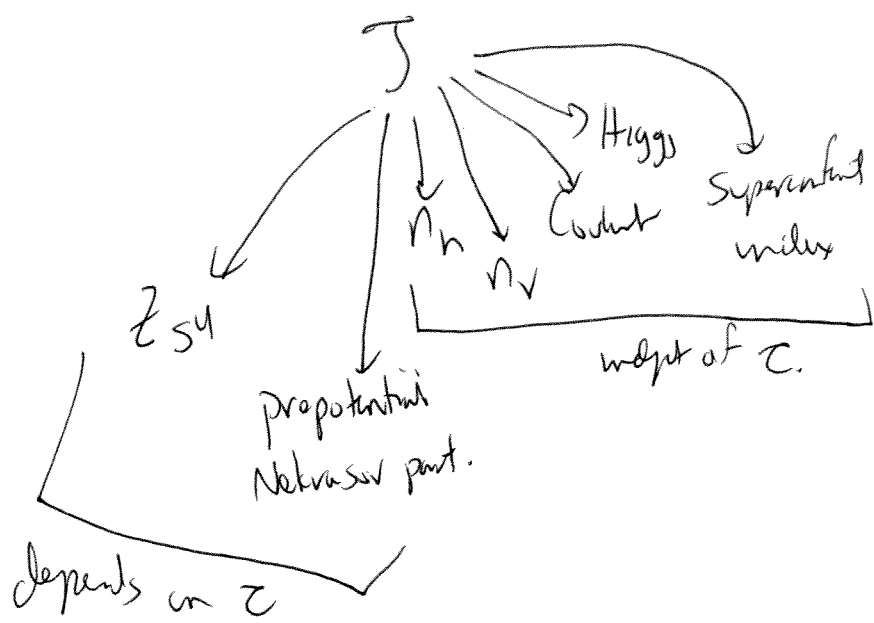
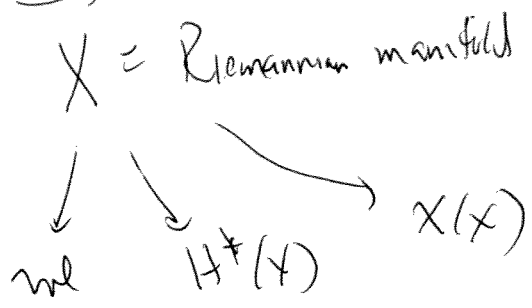
$\text{Im } \tau \gg 0$ weakly coupled

2d TQFT
CFT

Z :  \longrightarrow $\mathbb{1}$

$$Z[\text{circle } \alpha \text{ --- circle } \beta] = \sum_a Z[\text{circle } \alpha \times] Z[\text{circle } \times \beta]$$

analogy:



$$\text{Higgs}(T_1) \times \text{Higgs}(T_2) \equiv G$$

• Lagrangian vs. non-Lagrangian SCFT
 ↓
 has a path integral description → does not have a path integral

N=2 Lagrangian theory

$G =$ compact Lie group

$\tilde{\rho} =$ pseudo-real rep. of G .

$$\mathbb{R}^4 \quad \delta^M (\partial_\mu + \tilde{\rho}(A_\mu)) \not\chi = 0$$

$$\begin{matrix} \uparrow \\ \mathbb{D} \otimes \tilde{\rho} \end{matrix} \quad \swarrow \text{Spin}(4) \simeq \text{SU}(2) \oplus \text{SU}(2)$$

$$C\psi = \not\chi$$

\mathbb{D} pseudo real } strictly real.
 $\tilde{\rho}$ pseudo real }

conj. linear map $C: \not\chi \rightarrow \not\chi, C^2 = 1$.

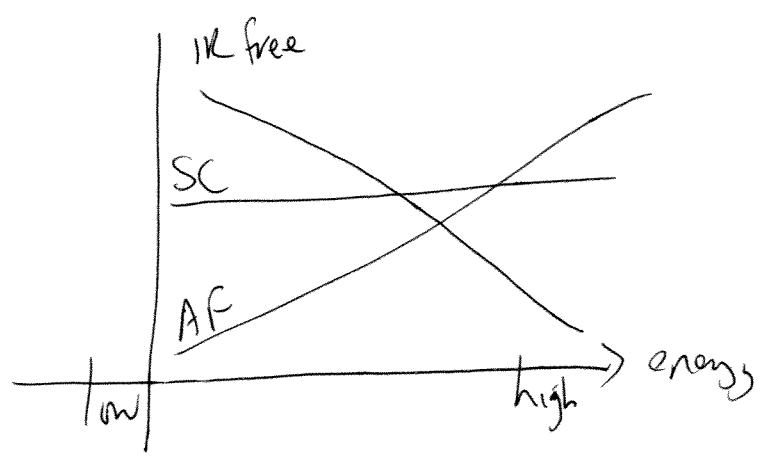
Special case: $\tilde{\rho} = \rho \oplus \rho^*$ $\gamma^M(\partial_{\mu} + \rho(A_{\mu}))\psi = 0$

$k_G(\rho \oplus \rho^*) = 2 \times \dim_{\mathbb{C}} \text{Index of } D_{A,\rho}$
 when A is an 1-instanton configuration

$k_G(\text{adj} \oplus \text{adj}) = 4h^v(G)$

$k_{\text{SU}(N)}(N \oplus \bar{N}) = 2$

- $k_G(\tilde{\rho}) > 4h^v(G)$ infrared free
- $= 4h^v(G)$ Superconformal
- $< 4h^v(G)$ asymptotically free



$$G > \tilde{\rho} = \text{adj} \oplus \text{adj}$$

$N=4$ SYM.

$$G = \text{SU}(N), \quad \tilde{\rho} = N \otimes \bar{2N} \oplus \bar{N} \otimes 2N$$

$\nearrow \text{SU}(N) \times \text{U}(2N)$

$N=2$ $\text{SU}(N)$ with $2N$ flavors

\nwarrow $f = \text{flavor symmetry}$

$$n_V(\tilde{\rho}/G) = \dim G$$

$$n_H(\tilde{\rho}/G) = \dim_{\mathbb{H}} \tilde{\rho} = \dim_{\mathbb{C}} \rho \quad \text{when } \tilde{\rho} = \rho \oplus \rho^*$$

$$k_F(\tilde{\rho}/G) = k_F(\tilde{\rho})$$

(Normalized quadratic Casimir)

$$C_{\text{Coulomb}}(\tilde{\rho}/G) = \bigwedge_{\text{Spec}} \mathbb{C}[\phi]^G$$

\nwarrow adj. of G .

$$= \bigwedge_{\text{Spec}} \mathbb{C}[u_1, u_2, \dots, u_r] \quad r = \text{rank } G$$

$$d(u_i) = e_{i+1}$$

\nwarrow exponent of G .

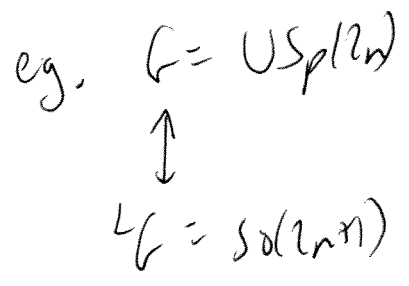
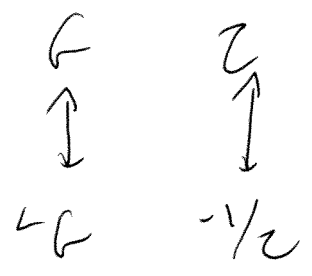
eg. $\text{SU}(N), \quad u_i = \text{tr } \phi^{i+1} \quad i=1, \dots, N-1$

$$\text{Higgs } (\tilde{\rho}/G) = \tilde{\rho} // G.$$

Things depending on τ :

- / prepotential
 - / Nekrasov partition function
 - / Z_{SYM}
 - / SW curves
- } to discuss next week

$N=4$ SYM.

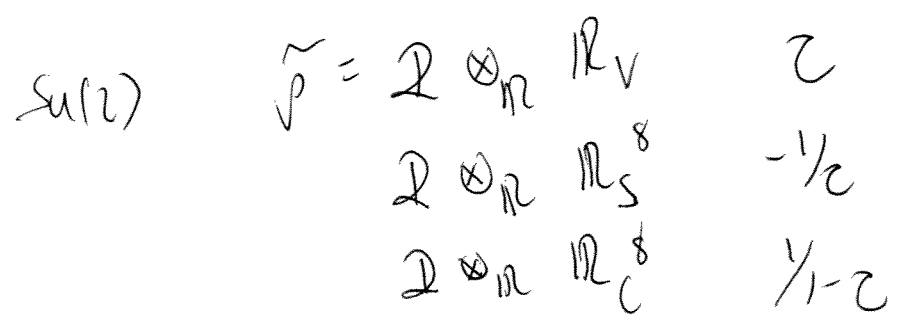


$N=2$ $SU(2)$ 4 flavors

$$\tilde{\rho} = \mathbb{Z} \otimes \bar{4} \oplus \bar{\mathbb{Z}} \otimes 4$$

$$= \mathbb{Z} \otimes_{\mathbb{R}} \mathbb{R}^8$$

$SU(2) \times SO(8)$



$\tilde{J} // \text{SU}(2)$ is 1-instanton moduli of $\text{SO}(8)$
 (depends on the adjoint of $\text{SO}(8)$)

non-Lagrangian theories

Mitsunaka + Nemeschansky in 1996:

$\text{MN}(E_{6,7,8})$

	n_V	n_H	Coulomb	Flavor symmetry
$\text{MN}(E_6)$	5	16	$\mathbb{C}[u]$ $d(u)=3$	E_6
$\text{MN}(E_7)$	7	24	$\mathbb{C}[u]$ $d(u)=4$	E_7
$\text{MN}(E_8)$	11	40	$\mathbb{C}[u]$ $d(u)=6$	E_6

Higgs branch.	$\dim_{\mathbb{H}} \text{higgs}$	k_E	family of SW curve
centered, framed 1-inst. moduli space of E_6	11	6	Kodaira
of E_7	17	8	
of E_8	29	12	

Given two $W=2$ SCFT A, B

(1) theory $A+B$.

n_v
 n_h
 ~~n_h~~ } additive

Coulomb
Higgs } direct product

If a theory A has a flavor symmetry $\tilde{F} \supset G$ st.

$$k_G(A) = 4h^v(G)$$

then "theory A/G at coupling τ " satisfies

$$n_v(A/G) = n_v(A) + \dim G$$

$$n_h(A/G) = n_h(A)$$

$$\text{Coulomb}(A/G) = \text{Coulomb}(A) \otimes \mathbb{C}[\phi]^G$$

deg $\phi = 1$
(adjoint of G)

$$\text{Flavor Sym}(A/G) = \text{centralizer of } G \text{ in } \text{Flavor}(A)$$

$$\text{Higgs}(A/G) = \text{Higgs}(A) // G$$

\tilde{P} = free theory

\tilde{P}/G (Lagrangian theory) is a special case.

Argyres-Seiberg Example

$G = SU(3)$

$$\tilde{P} = B \otimes \bar{6} \oplus \bar{B} \otimes 6$$

$\curvearrowright SU(3) \times U(6)$

dual

$G = SU(2)$

$$\tilde{P} = D \otimes_{\mathbb{R}} \mathbb{R}^2 \leftarrow SU(2) \times SU(3)$$
$$+ MN(E_6) \leftarrow E_6 \supset SU(2) \times SU(6)$$

$$k_{SU(2)}(MN(E_6))$$

" "

$$k_{E_6}(MN(E_6))$$

" "

6.

total $k_{SU(2)} = 8 = 4h^v(SU(2))$.

So we can couple via $SU(2)$.

Higgs branch:

$$\mathbb{B} \otimes \bar{\mathbb{6}} \oplus \bar{\mathbb{B}} \otimes \mathbb{6} // SU(3)$$

$$\dim 18 - 8 = 10$$

$$\mathbb{Z} \otimes_{\mathbb{R}} \mathbb{R}^2 \times W_{\min} \left(\begin{smallmatrix} \mathbb{E} \\ \mathbb{6} \end{smallmatrix} \right) // SU(2)$$

$$\dim 2 + 11 - 3 = 10$$

Gaiotto-Neitzke-Tachikawa: definition agrees