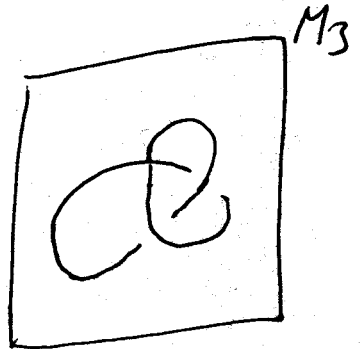


17 August 2000
E. Witten

Fivebranes and Knots

Chern-Simons in $d=3$

$$I = \frac{k}{4\pi} \int \text{Tr}(A \wedge A + \frac{2}{3} A \wedge A \wedge A)$$



$$Z_R(q) = \int \mathcal{D}A e^{iI} \text{Tr}_R \text{Pexp} \oint A$$

$M_3 = \mathbb{R}^3$

When you compute, one finds

$$Z_R(q; c) = \sum_{n=-\infty}^{\infty} a_n q^n$$

$$q = \exp \frac{2\pi i}{k+h}$$

$$a_n \in \mathbb{Z}$$

I. Frenkel, Khovanov

There is a "Hilbert" space \mathcal{H} , $\mathbb{Z} \times \mathbb{Z}$ graded

(two conserved quantities H F , $\sum u(i) \times u(i)$)
Hamiltonian fermion #

$$\text{Tr}_{\mathcal{H}} \mathcal{Z}^H (-1)^F = Z_R(q; c)$$

$$\mathcal{H} = \mathcal{H}(R; c)$$

$$\text{Tr}_{\mathcal{H}} e^{-\beta H} y^F$$

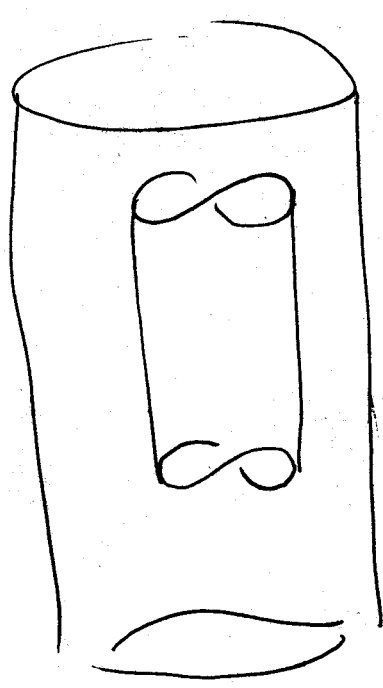
$$\text{Expect } M_3 \rightarrow M_3 \times \mathbb{R} \quad \uparrow \text{time}$$

$$\text{or } M_4 = M_3 \times S^1$$

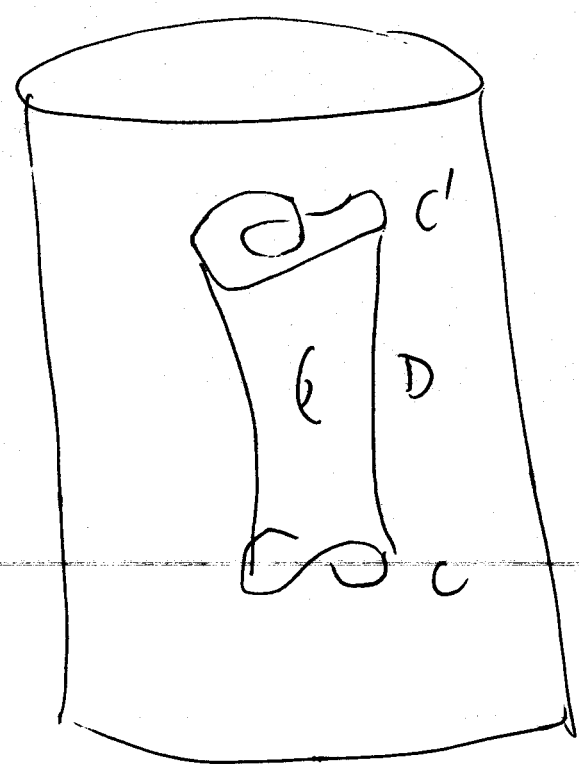
Path integral
⇒ number

Space of physical states

path integral
⇒ number
 \mathbb{Z}_2 graded trace
Supertrace in \mathcal{H}



World volume
= Support of a surface operator



$$\mathbb{F}_D: \mathcal{H}_c \rightarrow \mathcal{H}_{c'}$$

2004: Gubser-Vafa-Schwartz

⋮

1998: Ooguri-Vafa

$$\int DA \exp \frac{ik}{4\pi} \int \text{Tr}(ADA + \frac{2}{3}A^3) \text{Tr}_n \text{Perp} \int_C A$$

$$A \rightarrow a$$

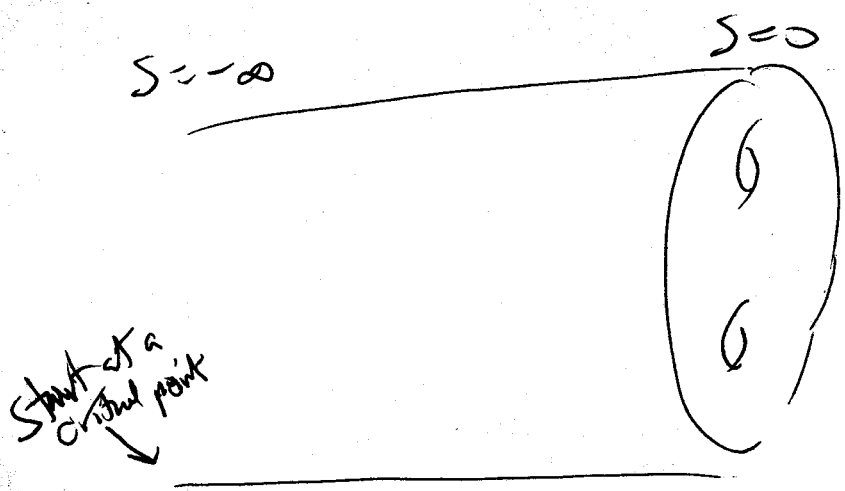
$$\int_D DA \exp \frac{ik}{4\pi} \int \text{Tr}(ada + \frac{2}{3}aaa) \text{Tr}_n \text{Perp} \int_C a$$

h = Real Part(---)

$$ds^2 = \int_{M_3} \text{Tr} \delta a_i \delta \bar{a}_i$$

$$= \int_{M_3} d^3x \sqrt{g} g \text{Tr} \delta a_i \delta \bar{a}_j$$

$$\frac{da}{ds} = - \frac{\delta h}{\delta a}$$



$$M_3 \times (-\infty, 0]$$

s

Cycle Γ consists of all values at $s=0$ of solutions of flow eqn in $\mathbb{R}_+ \times M_3$

If $M_3 = \mathbb{R}^3$, the only critical point is $a=0$ (up to gauge)

$$A = \sum c_\lambda A_{(\lambda)}$$

$$\int dc_\lambda \exp i \sum M_\lambda c_\lambda^2$$

$$\exp(i) (ada + \dots)$$

$$\exp(h + i\Phi)$$

Flow:

$$\frac{dA_i}{ds} = - \frac{\delta}{\delta A_i} \int \text{Tr}(A \dot{A} + \frac{2}{3} A^3)$$

$$\frac{dA_i}{ds} = \frac{1}{2} \underbrace{\epsilon_{ijk}}_{\vec{B}_i} F_{jk} \quad F^t = 0$$

$$A_5 = 0 \quad \vec{E}_i = \vec{B}_i$$

A real-valued connection

$$\frac{da}{ds} = - g^{aa} \frac{\delta h}{\delta a^a}$$

$$\phi \in \Omega^1(M_3 \times \mathbb{R}_+, \text{ad}(\mathfrak{g}))$$

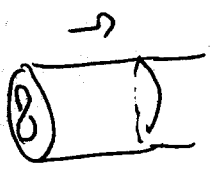
Flow equations are

$$F - \phi \wedge \phi = * D\phi$$

$$D*\phi = 0$$

↑ Kapustin-Witten parameter t

$$A = A + i \Phi \quad (t=i)$$



The theory on $M_3 \times \mathbb{R}_t$ is $\mathcal{N}=4$ SYM with the twist related to geometric Langlands (GL twist)

Consider the case $G = U(n)$.

$$S(F - \varphi_{\text{rep}} = *D\phi)$$

$\mathcal{N}=4$ SYM theory
 $M \times \mathbb{R}_t$

n parallel D3-branes

Here, the D3s end on an NS5

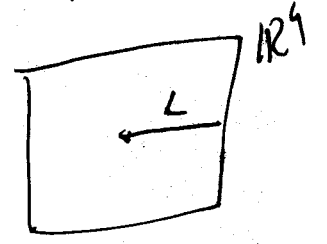
Bulk theory has $\theta \neq 0$, in fact $\theta = 2\pi i k$.

Geometry: 10 dim'l spacetime is $T^*M_3 \times \mathbb{R}^4$

NS5 supported on
 $T^*M_3 \times \{0\}$

D3 supported on
 $M_3 \times \mathbb{R}$
 \cap
 T^*M_3

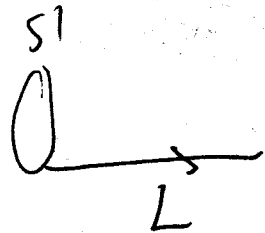
Half line from $\{0\}$ to ∞



$$\mathbb{R}^4 \rightarrow \mathbb{R}^3 \times S^1$$

$$T^*M_3 \times \mathbb{R}^3 \times S^1$$

$$L \in \mathbb{R}^3 \times \{p\}^{S^1}$$



(7)

(D3)

T-duality on S^1 :

$$D3 \rightarrow D4$$

$$NS5 \rightarrow (\dots)$$

$$T^*M_3 \times \mathbb{R}^3 \times S^1 \text{ asymptotically}$$

$$D4 \text{ on } M_3 \times L \times S^1$$

Event

$$T^*M_3 \times TN$$

Tach-Net



TN hyperkähler

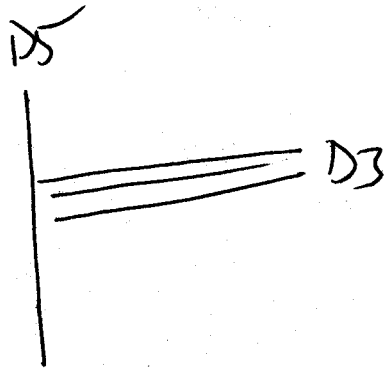
$$TN \rightarrow \mathbb{R}^3$$

$$\begin{array}{c} \text{genus } 1 \\ \uparrow \\ S^1 \end{array}$$

Actually, D4 on $M_3 \times$

Take the S-dual (before T-dual)

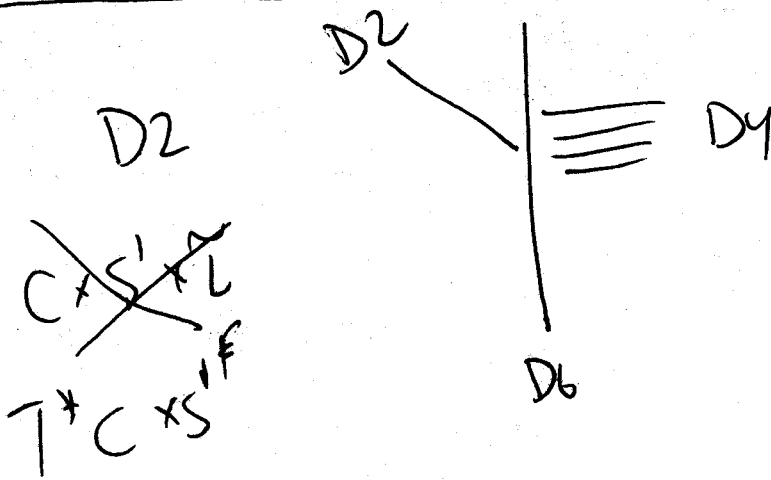
$$\mathbb{R}B: \quad NS5 \rightarrow D5$$



T-dual is a a D4-D6 system.

$$D6 \text{ on } T^*M_3 \times \underbrace{\{0\}}_{\mathbb{R}^3} \times S^1$$

$$D4 \text{ on } M_3 \times L \times S^1$$



~~$$C \times S^1 \times L$$

$$T^*C \times S^1 \times F$$~~

$$C \subset M_3 \subset T^*M_3$$

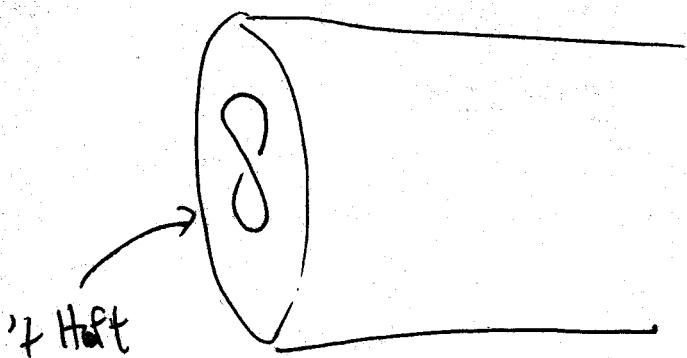
$$T^*C \subset T^*M_3$$

U(n) gauge theory on Dn's.

$$M_3 \times \mathbb{R}_4 \times \underbrace{\mathbb{R}}_{\text{time}}$$

$$\text{D2 end on } \underbrace{C}_{M_3} \times \underbrace{\{s=0\}}_{\mathbb{R}_+} \times \mathbb{R}$$

$$M_3 = 2(M_3 \times \mathbb{R}_+)$$



In general, could be on $M_4 \times \mathbb{R}_+$

$$M_4 = M_3 \times S^1$$

$$M_3 \times \mathbb{R}$$

$G = \text{Chern-Simons}$

$$M_4 \times \mathbb{R}_+$$

$$D = 2\text{-manifold}$$

$$D = \underbrace{D'}_{M_4} \times \underbrace{\{0\}}_{\mathbb{R}_+}$$

Fields are

$A = \text{gauge field}$
 $= \text{connection on } G^V \text{ bundle}$

$$E \rightarrow M_4 \times \mathbb{R}_+$$

$$B \in \Gamma(M_4 \times \mathbb{R}_+, \pi^* \Omega^{2,1}(M_4) \otimes \text{ad } E)$$

$$F^+ + B \times B = D_S B$$

$$A^2(\Omega^{3+}) \cong A^{2+}$$

$$*_{(14)} D B + L_{(2/25)} F = 0$$

$$L_{(2/25)} (* D B + F) =$$

$$\text{Instanton } \# = H = \int_{M_3 \times \mathbb{R}_+} \left(\frac{\text{Tr } F \wedge F}{8\pi^2} - \frac{\text{Tr}(P \wedge R)}{8\pi^2} \right)$$

G_n = "number" of sols for given $H=n$.

$$Z(M_4 \times \mathbb{R}_+) = \sum G_n q^n$$

$$M_4 = M_3 \times \mathbb{R} \text{ (time)}$$

- 1) chain groups
classical approach to ground states.
time-independent solutions

- 2) define a differential by "counting" time dependent solutions

gauge theory on $M_4 \times \mathbb{R}_+$

Good version in the (d, r, c) theory on $M_4 \times$

