

Thoughts on 'Theory'

→ Is magnetostrophic turbulence as simple as it looks?

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→ initial stages → discussion, please!

→ background: see 3 papers on Wiki

- re: RGT (resistivity gradient driven turbulence), from Rippling Modes (aka FKR)

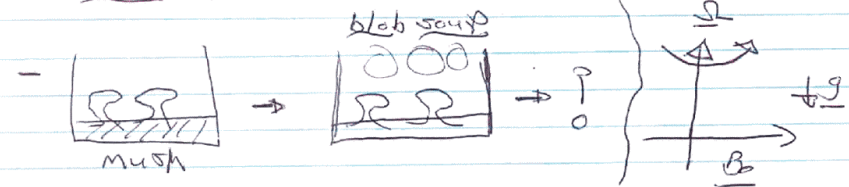
- structure similar to magnetostrophic (as is Vlasov turbulence)

- $\langle dT^2 \rangle$ dynamics central

→ see also: <http://physics.ucsd.edu/plasmatheorygroup/> (shameless advertising)

2.

① → What's going on? (or what we think is....)



- slugs/blobs erupt from mush at top of inner core

- slug $\left\{ \begin{array}{l} B \rightarrow \text{slices (St. Pierre)} \\ \Omega \rightarrow \text{columns} \end{array} \right. \Rightarrow$ symmetrical blob

- blobs → buoyant rise → flow → complex strain field ⇒ turbulence, cascade...

→ wakes → radiated MAC waves → ①

⇒ blob { fragments } during rise ↔ turbulence
but { radiates }

- $\langle \theta \rangle$ from erupting slugs ⇒ $\left\{ \begin{array}{l} \langle \theta(z) \rangle \text{ evolves on transport time scales} \\ \text{local } \partial_z \langle \theta \rangle \rightarrow \text{local instability?} \\ \text{aka 'Braginskii' } G_{\theta}^{\theta} \dots \end{array} \right.$

this suggests:

turbulent blob soup { structure collective modes }

2.

- also blobs \Rightarrow shear flow

$$\underline{u}^{(0)} = -\underline{\nabla} \cdot \underline{\tau} \times \underline{z} / \Omega, \quad \underline{\tau} \sim \langle \theta^2 \rangle$$

\rightarrow structure of shearing field? $\left\{ \begin{array}{l} \text{mean} \\ \text{zonal} \end{array} \right\}$
 \rightarrow impact on blob lifetime

\rightarrow A Few Equations

$$\left[\frac{\partial \theta}{\partial t} + \underline{u} \cdot \underline{\nabla} \theta - \chi \nabla^2 \theta = S \right] \quad \left\{ \begin{array}{l} \theta = \langle \theta \rangle + \delta \theta \\ \underline{u} = \langle \underline{u} \rangle + \delta \underline{u} \end{array} \right.$$

$$\theta_{y, \omega} \sim \alpha \sim \langle \delta \theta^2 \rangle_{y, \omega}$$

$$\delta u_{y, \omega} = \frac{A(k)}{d(k, \omega)} \delta \theta_{y, \omega}$$

$\rightarrow \left\{ \begin{array}{l} \beta, \beta, \beta \\ \omega, \eta k^2 \end{array} \right.$ $R_m \ll 1$ important
 $\rightarrow \eta$ is key damping

n.b.:
 $\langle u \rangle \sim \langle \delta \theta^2 \rangle \Rightarrow$
 $\langle u \cdot \nabla \theta \rangle \sim \delta \theta^3$

modes: $(-i\omega + \chi k^2) = \frac{-A(k) \cdot \nabla \langle \theta \rangle}{d(k, \omega)} \Rightarrow$ Braginskii's $\frac{6}{4} \dots$
 buoyancy modes

\Rightarrow revisit with shearing, etc. ...

Quantity of Interest: $\langle \delta \theta^2 \rangle_{y, \omega}$
 \downarrow
 Theatery spectrum

3.

(I)
 \rightarrow why? \Leftrightarrow What is Needed?
 \rightarrow A Statistical Theory of Theatery/Blobs

- statistically: "blob" \Leftrightarrow peak $\langle \delta \theta(x) \delta \theta(z) \rangle$

expect correlation structure in time:

$$\langle \theta(x) \theta(z, t) \rangle = |\theta_0|^2 e^{-i(\omega_0 - k \cdot v) t - \gamma / \tau_{ch}}$$

$\omega <$ frequency decay time \downarrow modes frozen \downarrow frozen flow \downarrow lifetime

- apart χ , $\left\{ \begin{array}{l} \frac{d\theta}{dt} = 0 \rightarrow \left\langle \frac{d\delta \theta^2}{dt} \right\rangle = -\frac{\partial \langle \theta \rangle^2}{\partial t} \\ \frac{\partial \langle \theta \rangle}{\partial t} = -\nabla \cdot \langle \underline{u} \delta \theta \rangle \end{array} \right.$

\Rightarrow $\left[\frac{\partial \langle \delta \theta(x) \delta \theta(z) \rangle}{\partial t} + \langle (\underline{u}(x) - \underline{u}(z)) \cdot \nabla \delta \theta(x) \delta \theta(z) \rangle + \dots \right]$ \uparrow shear + fragmentation \uparrow diffn

$$= -\langle \delta \underline{u}(x) \delta \theta(x) \rangle \cdot \nabla \langle \theta \rangle + [A \neq 2]$$

i.e. $\frac{\partial \langle \delta \theta^2 \rangle}{\partial t} + T_{1,2} \langle \delta \theta^2 \rangle = P$

$T_{1,2} \Rightarrow$ 2 pt. evolution \rightarrow buoyancy flux

$P \rightarrow$ production $\lim P = -\theta \cdot \nabla \langle \theta \rangle$
 $\rightarrow \sim$ extra production

4.

→ Fragments of the Theory $\sim \langle d\theta^2 \rangle$

ⓐ $T_{12}[\langle d\theta^2 \rangle] = \frac{\partial \langle d\theta^2 \rangle}{\partial t} + \underbrace{(\langle u_1 \rangle - \langle u_2 \rangle)}_{\text{shearing}} \cdot \nabla \langle d\theta^2 \rangle$

+ $\langle (\partial u_1 - \partial u_2) \cdot \nabla d\theta^2 \rangle$ + visc.
 $\sim \langle d\theta^3 \rangle \rightarrow$ closure needed....
 (1)

from painful experience, expect:

$T_{12} = \frac{\partial \langle d\theta^2 \rangle}{\partial t} + (\langle u_1 \rangle - \langle u_2 \rangle) \cdot \nabla \langle d\theta^2 \rangle - \nabla \cdot \underline{\underline{D_{rel}}} \cdot \nabla \langle d\theta^2 \rangle$

①: + →
 ②: + →

$\underline{\underline{D_{rel}}} = \sum_{k, \omega} \langle d\theta^2 \rangle_{k, \omega} \frac{AA}{|k|^2} \tau_{\omega} \sin(k \cdot x) \rightarrow \begin{matrix} 0 \\ (1 \rightarrow 2) \end{matrix}$

⇒ for blobs (lifetime): (scale l)

- D/l^2 relative diffusion
- $1/\tau_l \rightarrow$ eddy Richardson
- $(\frac{1}{l^2} \langle u \rangle^2 D)^{1/3} \rightarrow$ diffusion shearing

key issues:

- interaction locality in scale l
- importance of shear flow

5.

ⓐ → Production

→ not apparent that "inertial range" exists
 ↔ collective mode effects, $\nabla \langle \theta \rangle$, etc...

→ $P = -\underline{\underline{Q}}_0 \cdot \nabla \langle \theta \rangle = \langle \partial u \partial \theta \rangle \cdot \nabla \langle \theta \rangle$
 (P) → physical meaning

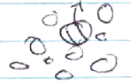
- can write directly, by $\partial u \leftrightarrow \partial \theta$

- "usual" $\partial \theta = \theta^c = -R_{\omega} \partial u_{\omega} \cdot \nabla \langle \theta \rangle$
 coherent response

⇒ $P = \nabla \langle \theta \rangle \cdot \underline{\underline{D}} \cdot \nabla \langle \theta \rangle$

→ mixing $\nabla \langle \theta \rangle$ feeds fluctuations
 → analogous to l^{-1} theory

What of wake, dynamical friction



→ condensing a long story:

- $\partial \theta = \theta^c + \tilde{\theta}$
 coherent incoherent
- $\tilde{\theta} \Rightarrow$ dynamical friction → F-P drag partic/cancellations
- θ^c screens $\tilde{\theta}$

6.

$$\rho = -\nabla \langle \theta \rangle \cdot \sum_{\mathbf{q}, \omega} \frac{A(\mathbf{q}, \omega)}{|\mathbf{d}(\mathbf{q}, \omega)|^2} \frac{\langle \tilde{\theta}^2 \rangle_{\mathbf{q}, \omega}}{|\mathbf{E}(\mathbf{q}, \omega)|^2}$$

buoyancy modes

and

$$\partial_t \langle \rho \theta^2 \rangle + T_{1,2} [\langle \rho \theta^2 \rangle] = \rho$$

$$\langle \rho \theta^2 \rangle \approx \tau_L \rho \rightarrow \dots \in \tau_{EM}, \tau_T$$

→ What do we gain from this approach?

- systematic statistical framework
- route to understanding wave-blob interaction ...
- lifetime \leftrightarrow { shearing, fragmentation }
- Production $\begin{cases} \rightarrow$ diffusion \rightarrow dynamical friction \rightarrow $\nabla \langle \theta \rangle$ \uparrow \downarrow $\nabla \langle \theta \rangle$

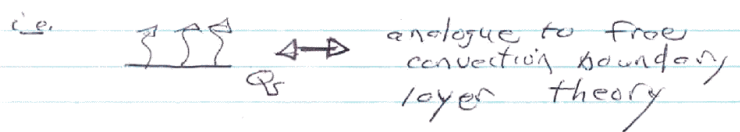
7.

III

→ Issues for Discussion

- basic picture ?
- "fragmentation" vs. "bubble competition"
- "cascade" vs. "coagulation"
- ie. large $\rho \theta$ overtake/engulf smaller ?
- \Rightarrow few blob state ?

- better off with macroscopic ?



- are collective modes important in story ?
- how best represent { local $\frac{\rho \theta}{\tau}$? }

8.

→ structure of shear flow }
 ↓

- smooth, zonal, both }
 ↓

- modulational stability of
 magnetostrophic turbulence }
 ↓

→ dominant interaction: local cascade
 vs. non-local interaction (shearing) }
 ↓

→ structure based blob kinetics model }
 ↓

