

Scaling laws for dynamos in rotating spherical shells

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in collaboration with

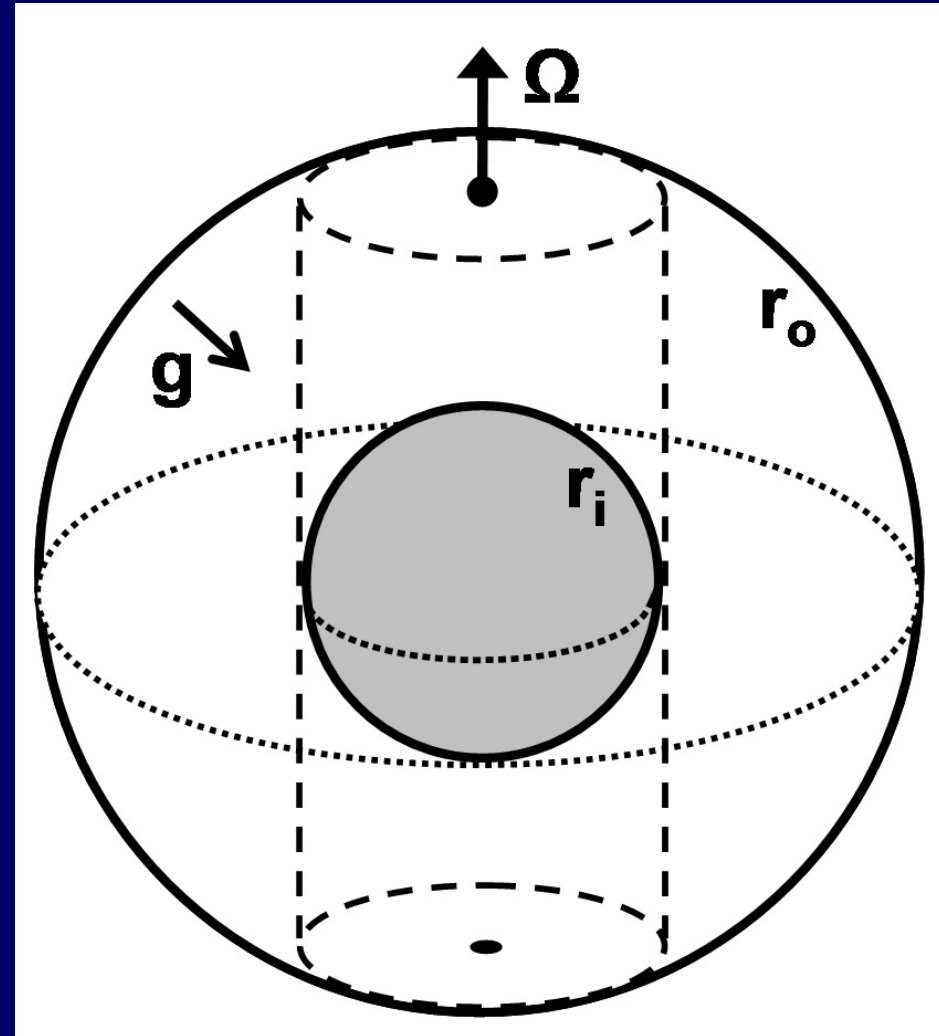
**Julien Aubert, Peter Olson,
Andreas Tilgner**

Introduction

- Geodynamo models successfully reproduced first-order properties of the geomagnetic field
 - This seems surprising, because several control parameters are far from Earth values (viscosity and thermal diffusivity too large, rotation too slow)
 - Pessimistic view: Models give right answer for wrong reasons
 - Optimistic assumption: Perhaps current models approach a regime in which diffusive processes do not play a first-order role in the force balance
- Use many case studies, covering decent range of control parameters, to derive scaling laws

Outline of dynamo models

- **Boussinesq equations for convection-driven MHD flow**
- **Rigid inner and outer boundary**
- $r_i / r_o = 0.35$
- **Fixed temperature contrast, no internal heat sources**



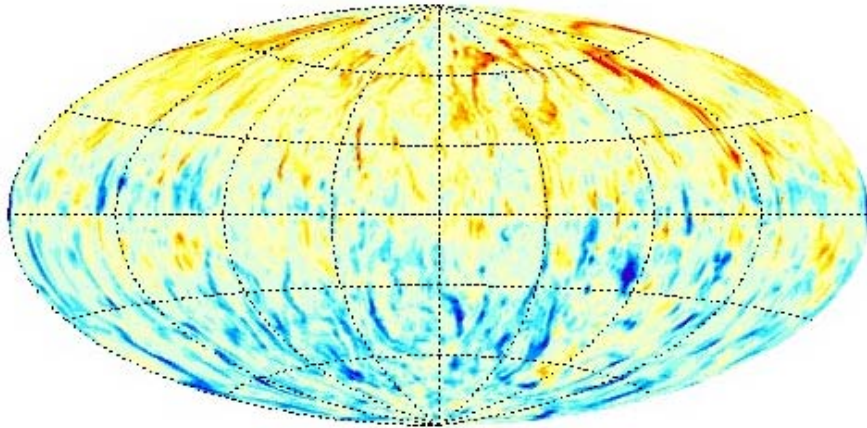
Control parameters

- Ekman number $E = \nu/(\Omega D^2)$ $10^{-6} \dots 10^{-3}$ (*)
- Prandtl number $Pr = \nu/\kappa$ $0.1 \dots 10$
- Magnetic Prandtl # $Pm = \nu/\lambda$ $0.06 \dots 20$
- Modified Rayleigh # $Ra^* = \alpha g_o \Delta T / \Omega^2 D$ $5 \dots 100 \times Ra_{crit}$

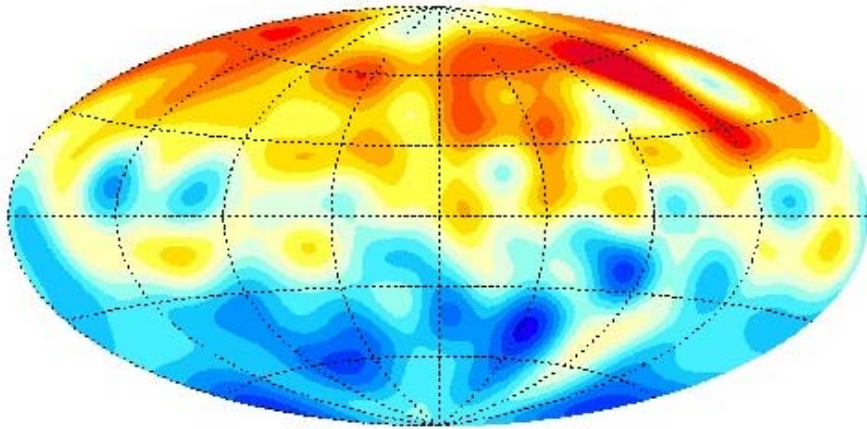
Modified Rayleigh number does not depend on diffusivity.

(*) With definition of Kono & Roberts [$E = \nu/(2\Omega r_o^2)$], the range is $2 \times 10^{-7} \dots 2 \times 10^{-4}$.

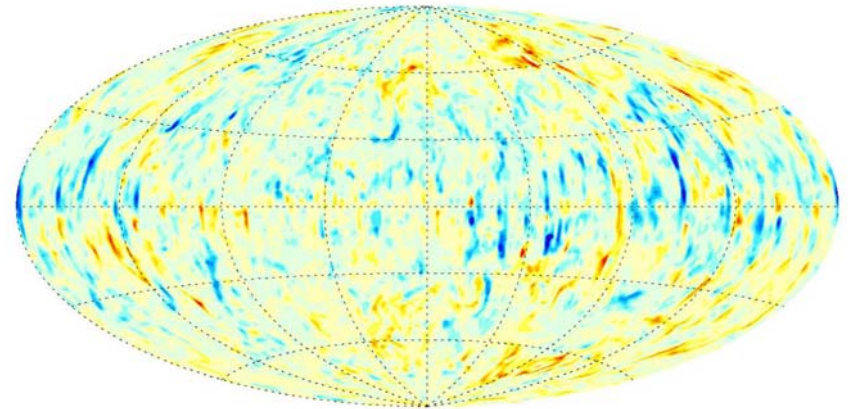
Dynamo classes



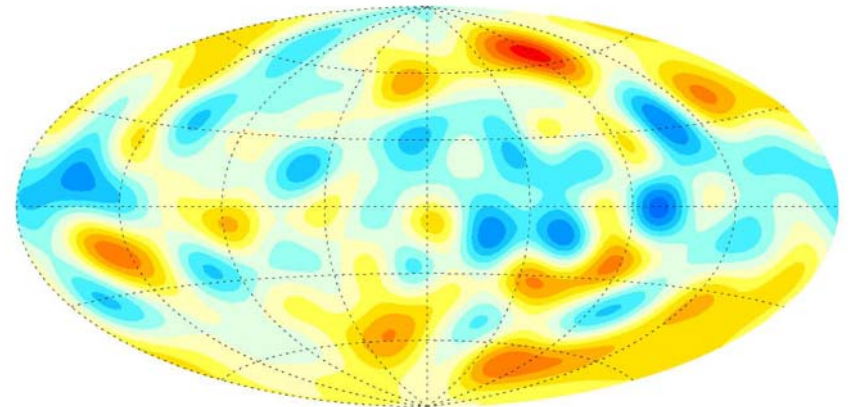
$Ra^*=0.12$ $E=10^{-5}$ $Pm=0.8$



Dipolar

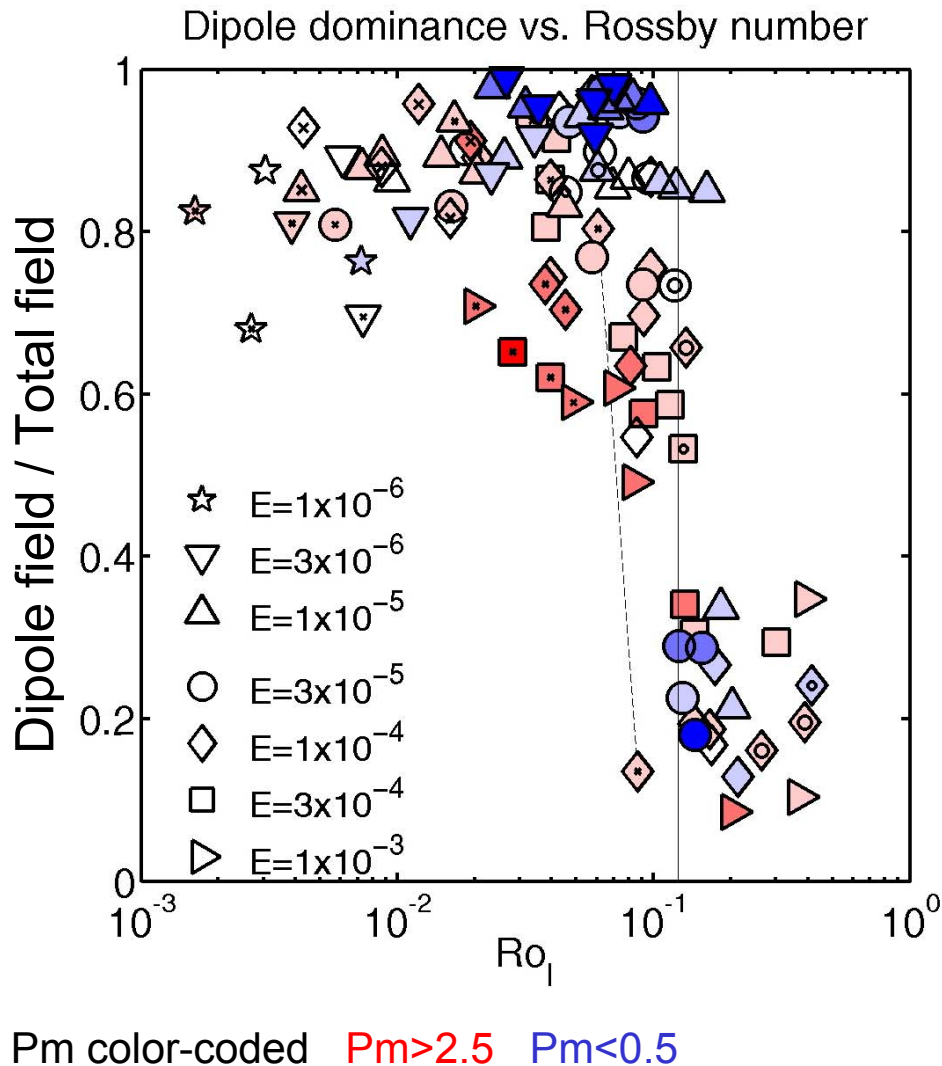


$Ra^*=0.17$ $E=10^{-5}$ $Pm=0.5$



Multipolar

Dynamo regimes



Inertial vs. Coriolis force:

Rossby number Ro_ℓ
calculated with mean
length scale ℓ in the
kinetic energy spectrum

$$Ro_\ell = U/\Omega\ell$$

Regime boundary at Ro_ℓ
 ≈ 0.12 (depends on
heating mode, b.c.)

Scaling analysis: data basis and case selection

- ~160 model cases (125 dynamos, 35 failed dynamos)
- Each run for at least 50 advection times (some much longer)
- Symmetry in longitude assumed for $E \leq 10^{-5}$

Selection criteria for scaling analysis:

- Self-sustained dynamo
- Dipole-dominated magnetic field ($f_{\text{dip}} > 0.35$)
- Fully developed convection ($Nu > 2.0$)

→ 87 model cases pass these criteria

Scaling and parameters II

- Ro Rossby number $Ro = U / \Omega D$
- Lo Lorentz number $Lo = B / (\rho\mu)^{1/2}\Omega D$
- Nu^* Modified Nusselt number $Nu^* = Q_{adv} / (4\pi r_o r_i \rho c_p \Delta T \Omega D)$

Relation to conventional parameters

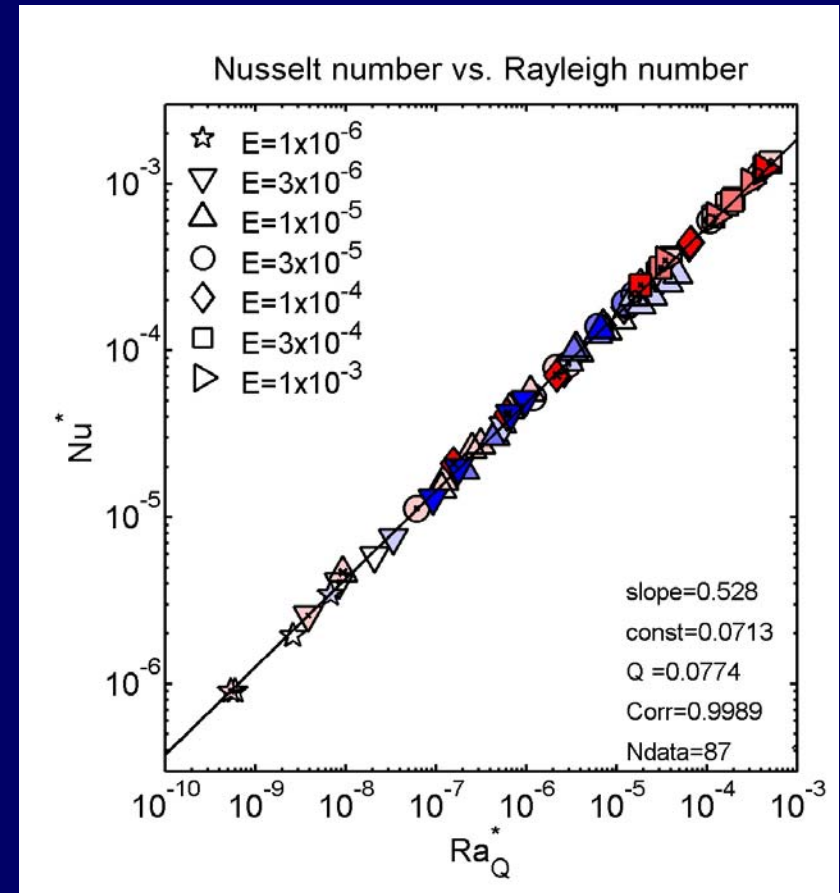
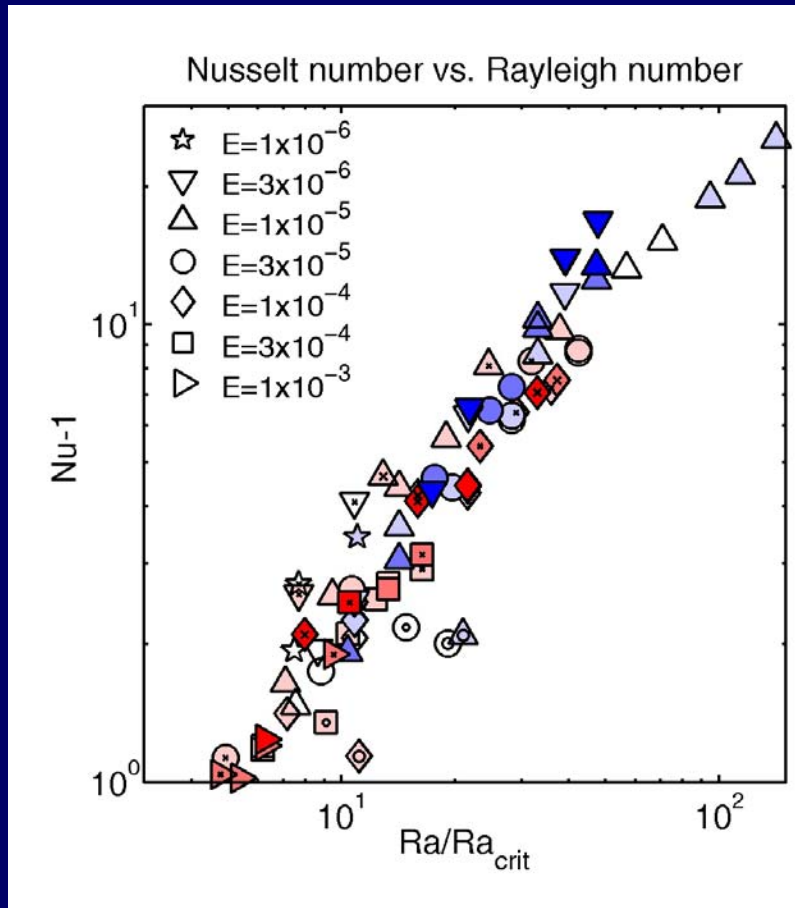
- Rayleigh number Ra $Ra^* = Ra E^2 Pr^{-1}$
- Nusselt number Nu $Nu^* = (Nu-1) E Pr^{-1}$
- Elsasser number Λ $Lo = (\Lambda E Pr^{-1})^{1/2}$

Flux-based modified Rayleigh number

$$Ra^*_Q = Ra^* Nu^* = Ra E^{-3} Pr^{-2}$$

Ra^*_Q is a measure for the work done by buoyancy forces

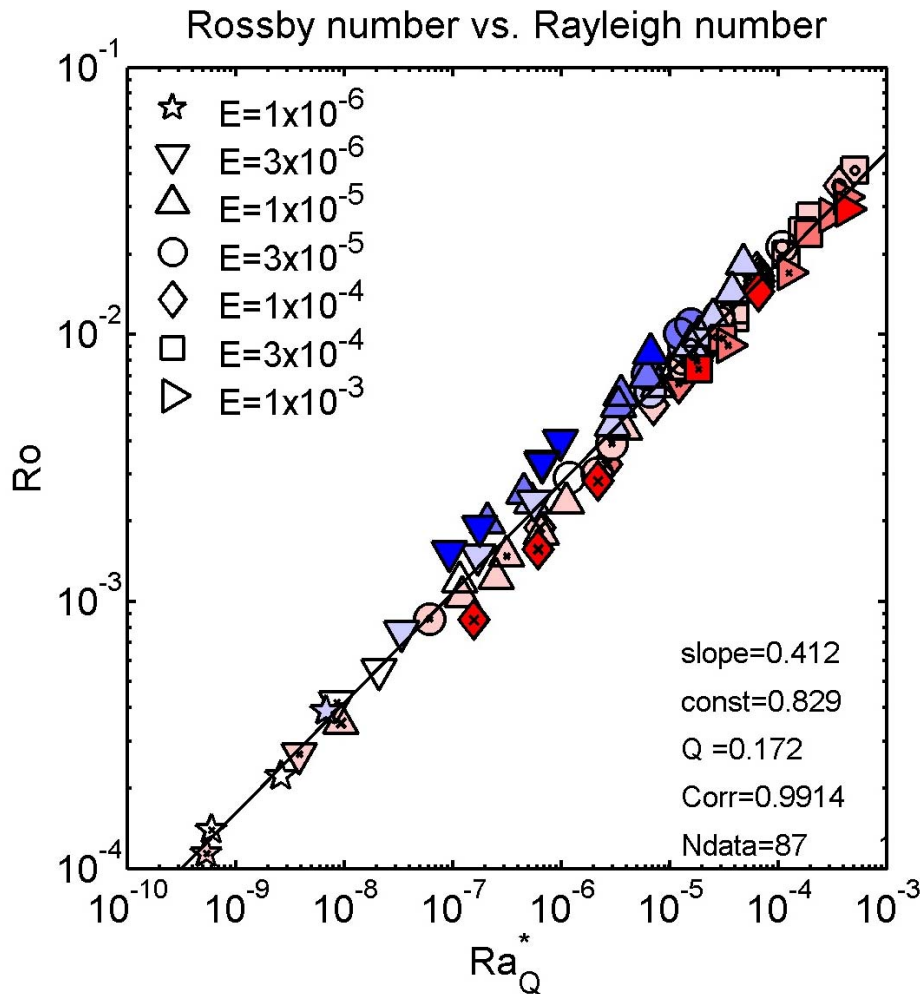
Scaling of Nusselt number



Use of modified „diffusionless“ parameters allows to collapse the data and express the dependence by a single power-law.

Compared to non-rotating convection, the exponent is very large (≈ 0.53).

Velocity Scaling

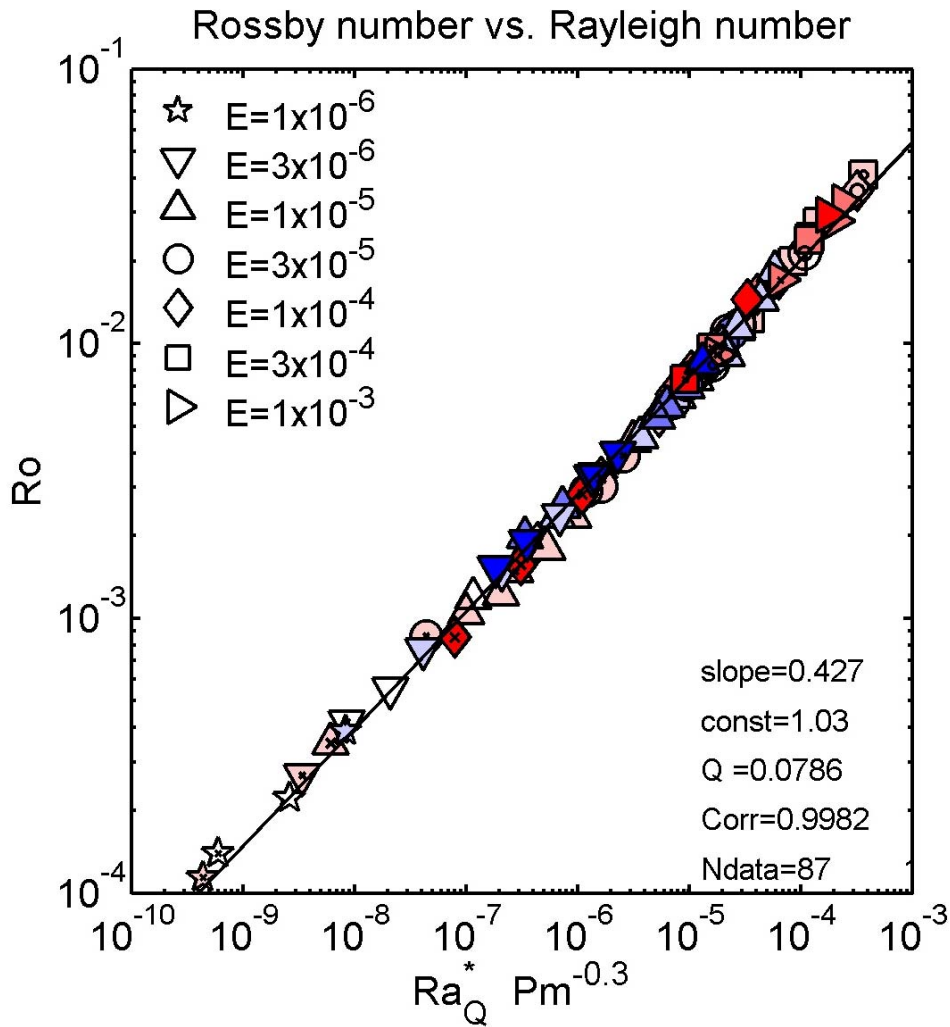


$$Ro \sim Ra_Q^*{}^{0.41}$$

Exponent 0.4 is predicted from a triple balance
Coriolis – Inertia - Buoyancy

Dependence on magnetic Prandtl number ?

Two-parameter fit

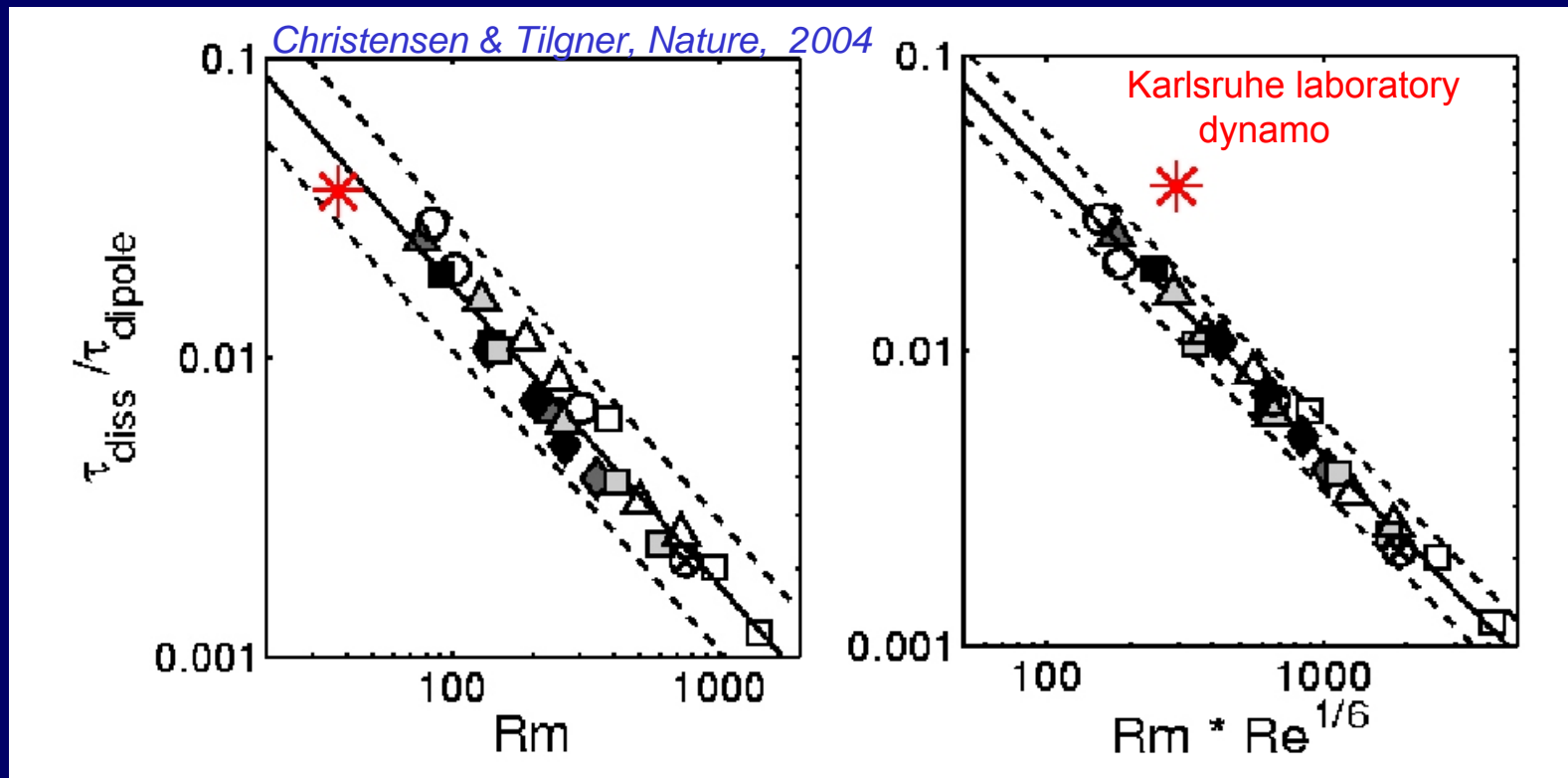


$$Ro \sim Ra_Q^{*0.43} Pm^{-0.13}$$

Two-parameter regression
reduces misfit by factor two

Scaling of magnetic dissipation time

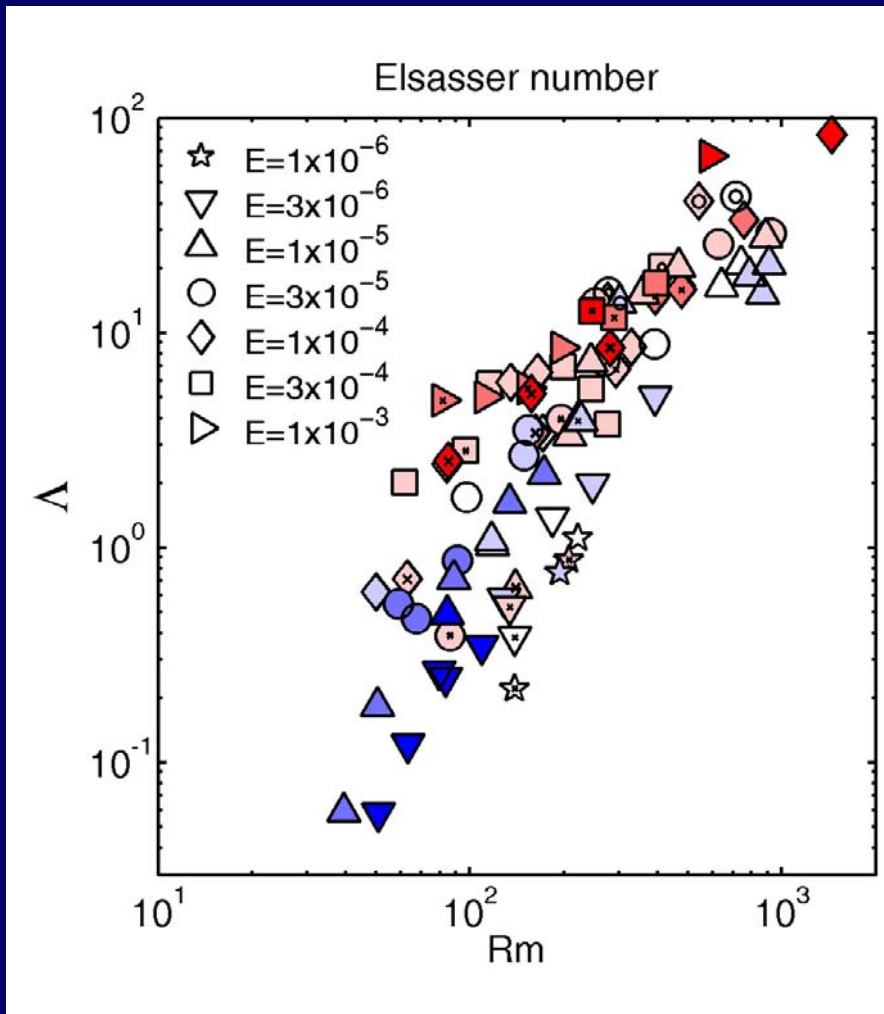
$$\tau_{\text{diss}} = E_{\text{mag}} / D_{\text{ohm}} \quad (\text{Magnetic energy} / \text{Ohmic Dissipation})$$



$$\tau_{\text{diss}} \sim Rm^{-1}$$

$$\tau_{\text{diss}} \sim Rm^{-7/6} Pm^{1/6}$$

What controls the strength of the magnetic field?



The presence of magnetostrophic balance is often associated with an Elsasser number $\Lambda = B^2/\mu\eta\rho\Omega$ of order one.

In the numerical models, the Elsasser number varies in the range 0.06 – 100.

→ Either the force balance is not magnetostrophic, or Λ not a good measure for magnetostrophy.

Alternative scaling:
Magnetic field strength
based on available power ?

Power-limited magnetic field strength

$$\frac{1}{2} L_0^2 = E_{\text{mag}} = D_{\text{ohm}} \tau_{\text{diss}} \sim D_{\text{ohm}} \text{Ro}^{-1} \text{E Pm}^{-1}$$

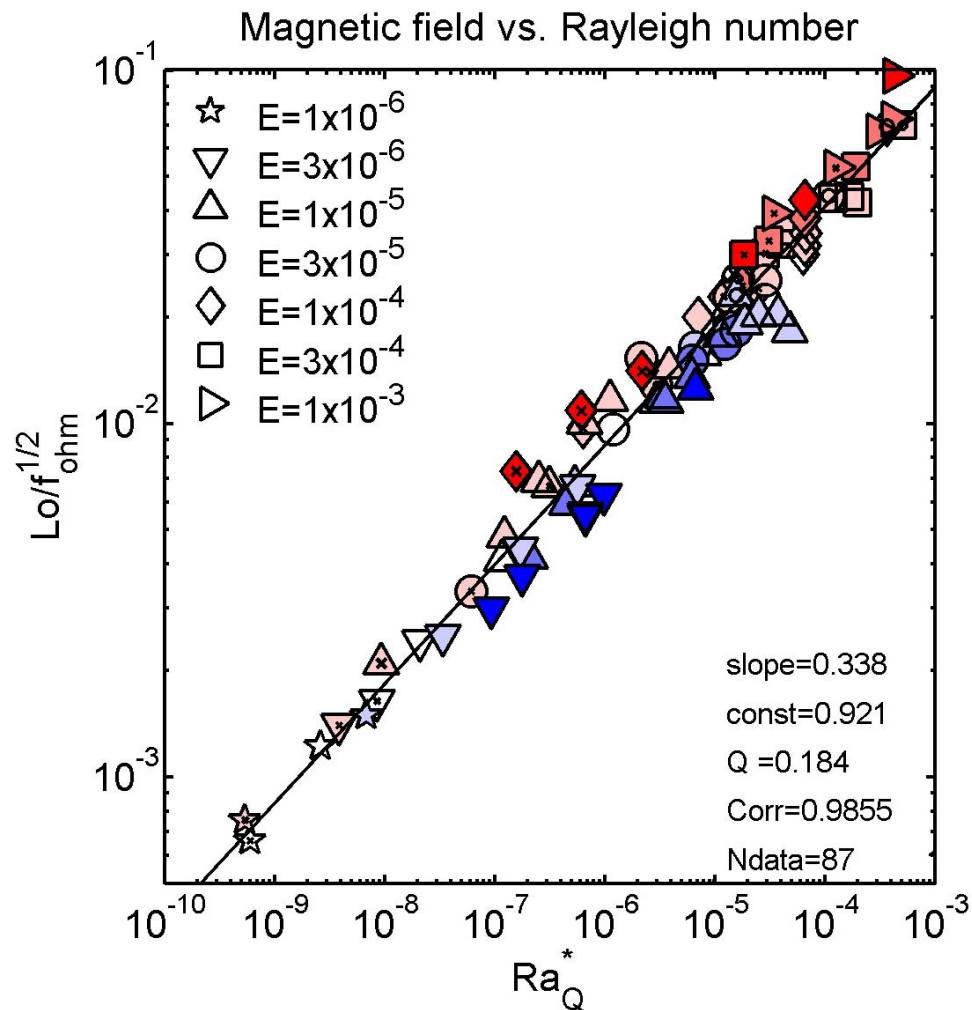
$$\text{Ra}^*_Q \sim \text{Power} = D_{\text{ohm}} / f_{\text{ohm}}$$

$$\text{Ro} \sim \text{Ra}^*_Q{}^{2/5}$$

$$\text{Prediction: } L_0 \sim \sqrt{f_{\text{ohm}}} \text{Ra}^*_Q{}^{3/10}$$

f_{ohm} : *fraction of ohmic dissipation*

Magnetic Field Scaling



$$Lo \sim Ra_Q^*{}^{1/3}$$

Empirical fit is close to
predicted dependence

Scaling in dimensional form

$$B \sim \mu^{1/2} \rho^{1/6} \left(\frac{P}{4\pi R^2} \right)^{1/3}$$

μ : magnetic permeability

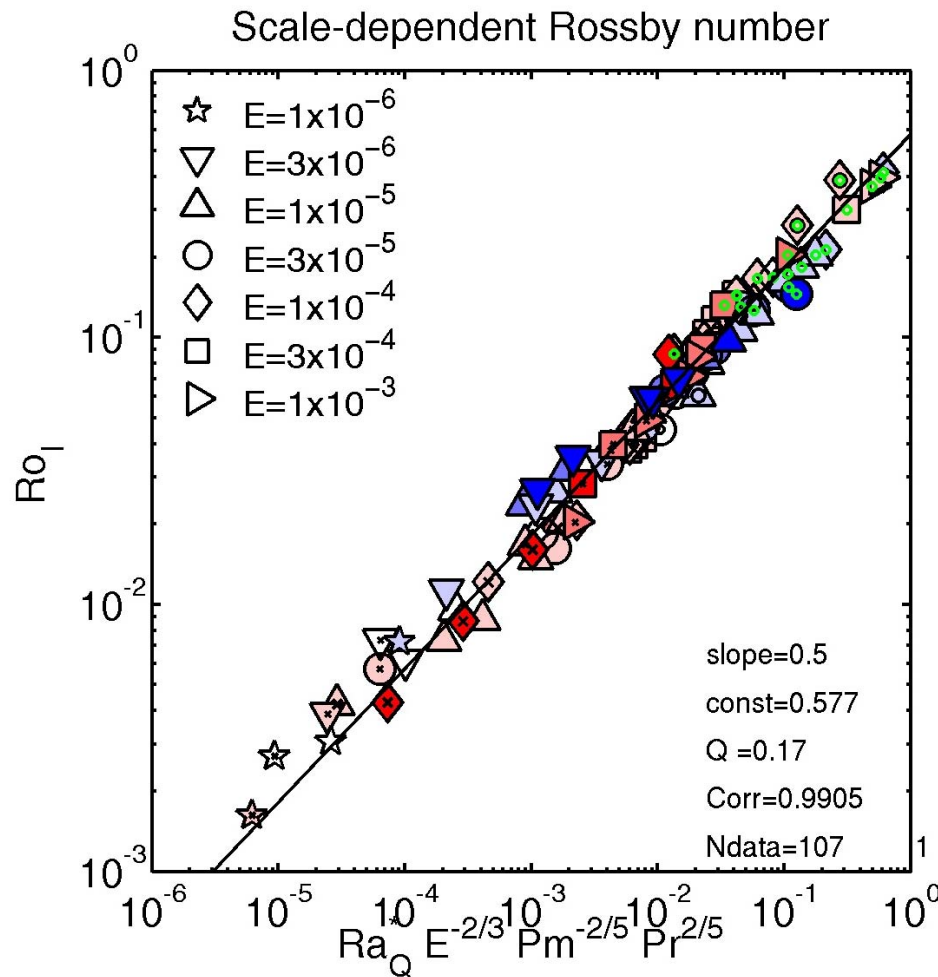
ρ : density

P: power

R: radius of dynamo layer

B independent of conductivity and rotation rate !

Ro_ℓ vs. control parameters

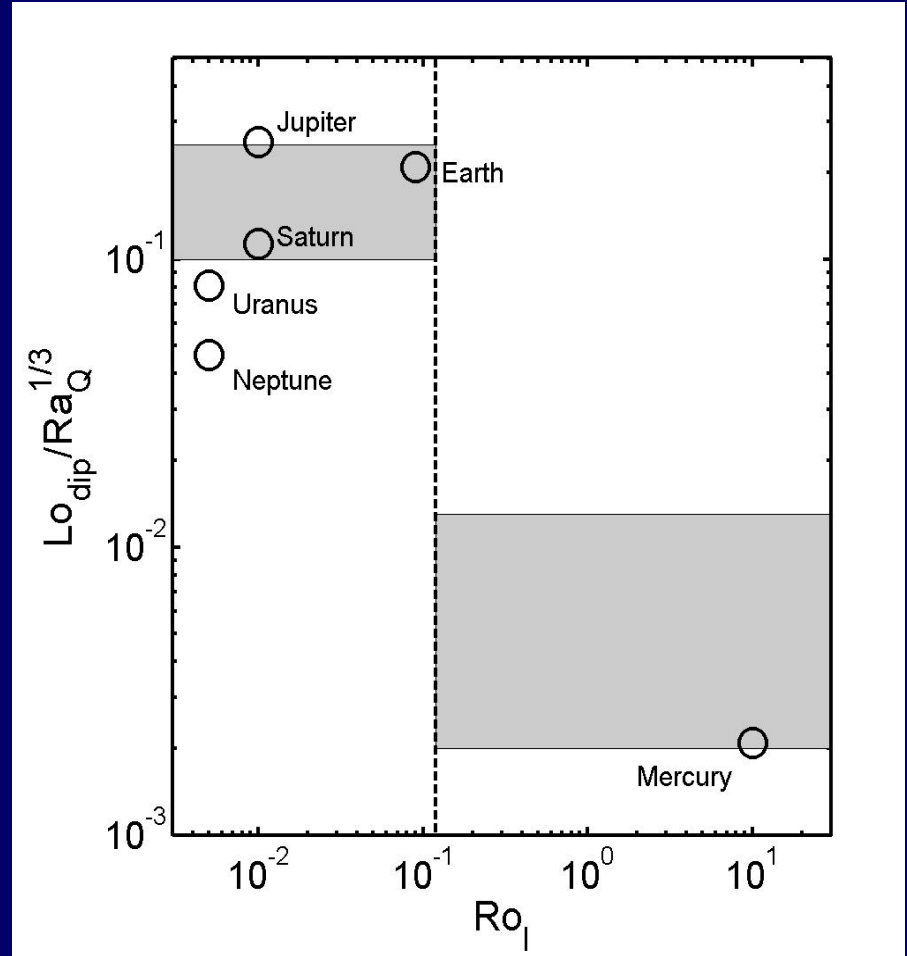
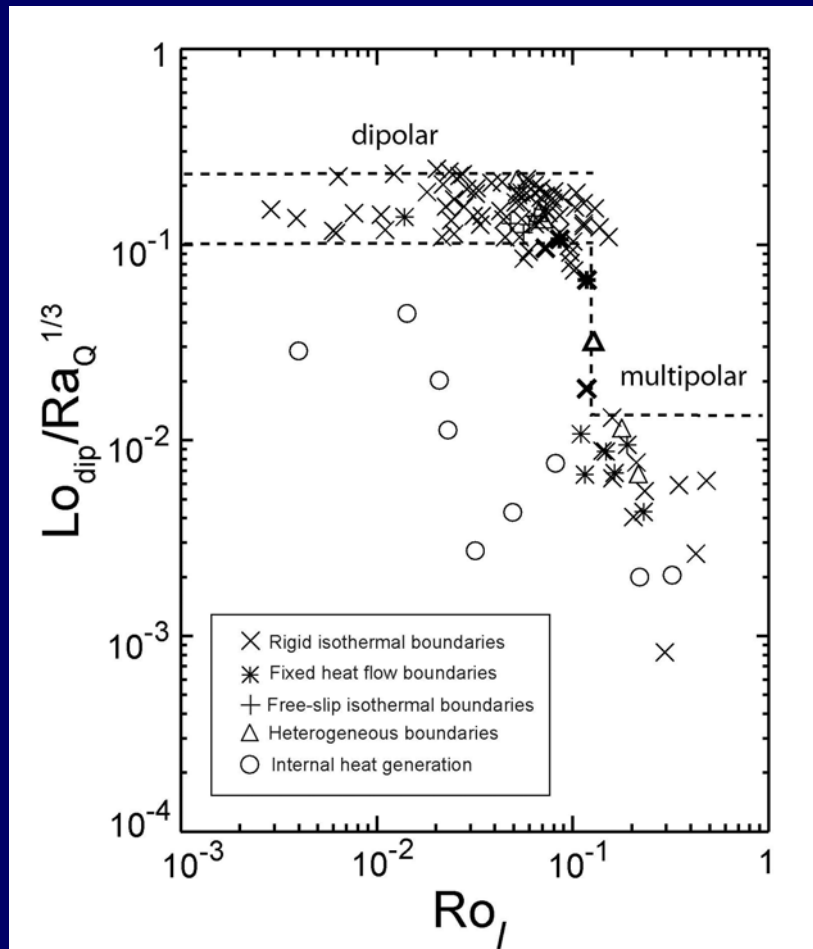


Fit involves all four control parameters

$$Ro_\ell \sim$$

$$Ra_Q^{*1/2} E^{-1/3} Pm^{-1/5} Pr^{1/5}$$

Dipole moment scaling



Earth predicted to lie close to transition dipolar - multipolar