

Slow relaxation in quasi-periodically driven random spin chains

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Phases in Isolated Quantum Systems

- Interested in non-equilibrium properties of highly excited many-body states
- New Quantum Dynamic Phases:
unitary evolution + coherent external fields
- Many recent experiments on systems which are well isolated from dissipation:
 Ultracold Atoms, Trapped Ions,
 Superconducting Qubits, NV Centers

Thermalization of Quantum Systems

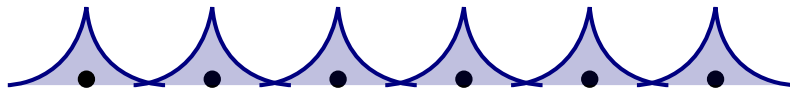
- ETH for Static Unitary Evolution:
pure state stay pure, but subsystems thermal
[Srednicki; Deutsch]



- Driven Quantum Systems:
Thermalize to infinite temperature, governed by
(entanglement) Hydrodynamics
[Nahum, Ruhman, Vijay, Haah; von Keyserlingk, Rakovszky, Pollmann, Sondhi]

Many-Body Localization in 1D (Static)

- Complete failure to thermalize
- Eigenstates are product states of local integrals of motion (LIOM / L-bits)




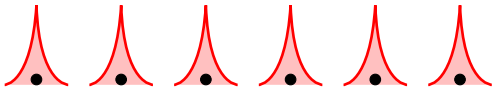
The diagram shows a horizontal line representing a 1D chain. Six black dots are placed along the line, representing sites. Above each dot is a blue, bell-shaped curve that is localized around that site. The curves are separated by small gaps, indicating that the states are localized and do not overlap significantly.

$$|\Psi\rangle = |\tau_1^z \tau_2^z \cdots \tau_L^z\rangle$$

- Excited states \sim equilibrium ground states
 - Area Law Entanglement: $S_{bp} \sim \text{const.}$
 - Have Phase Transitions [ETH-MBL / MBL-MBL]
 - Support symmetry breaking / topological phases

MBL with Periodic Driving

- Rapid, periodic, local driving keeps system MBL
[Ponte, Papić, Huveneers, Abanin; Lazarides, Das, Moessner]

H_+  H_- 

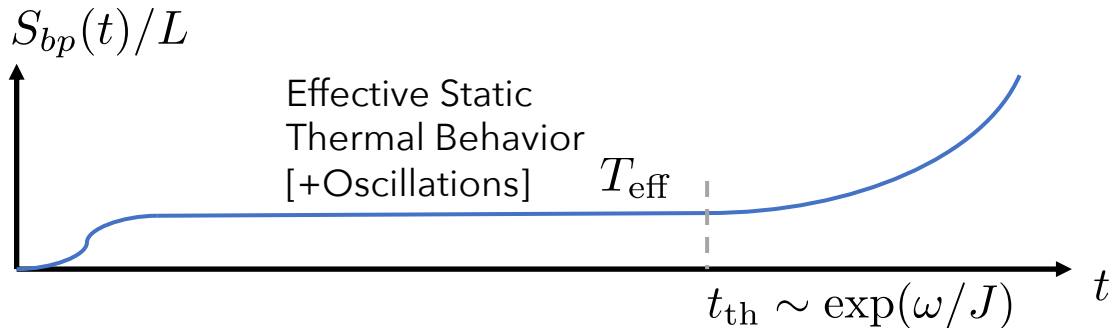
$U_F = e^{-i\lambda H_+} e^{-i\lambda H_-}$
 $= e^{-i(2\lambda)H_F}$

- Stability: can there be many-body resonances?
 - Single frequency (perturbative): exponential matrix element suppression, finite energy difference
 - No energy absorption; Modified L-bit structure
- Actual resonances with random driving or at low frequencies (Many-Body Landau-Zener)

Clean Pre-Thermal Floquet Phases

- No disorder, but heating rate can be exponentially suppressed with fast driving

[Abanin, De Roeck, Ho, Huveneers; Else, Bauer, Nayak]



- Can realize Floquet phases
- Cool pre-thermal regime: quantum phases

Quasiperiodic Sequences

- Fibonacci Drive: $U_n = U_{n-2}U_{n-1}$

$$U_0$$

$$U_1$$

$$U_0U_1 = U_2$$

$$U_1U_0U_1 = U_3$$

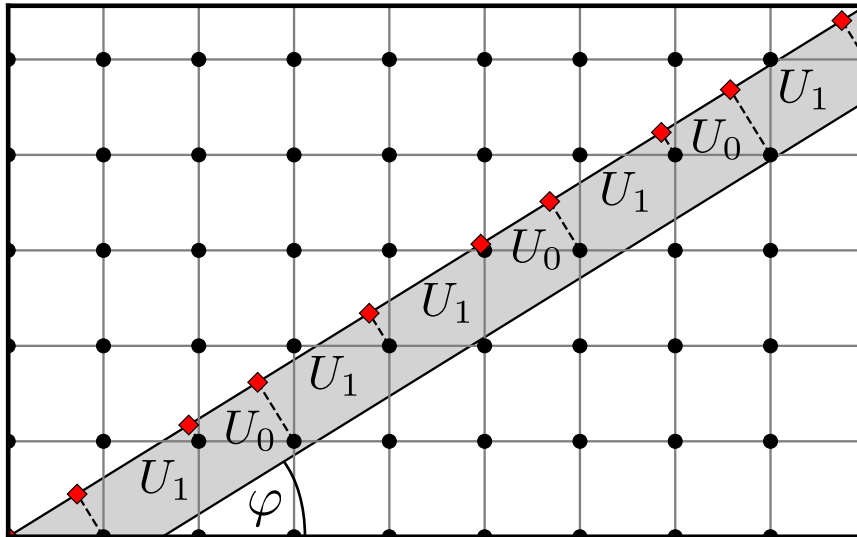
$$U_0U_1U_1U_0U_1 = U_4$$

$$U_1U_0U_1U_0U_1U_1U_0U_1 = U_5$$

- Efficient to simulate: need $\sim \log n$ multiplications

Quasiperiodicity From Projection

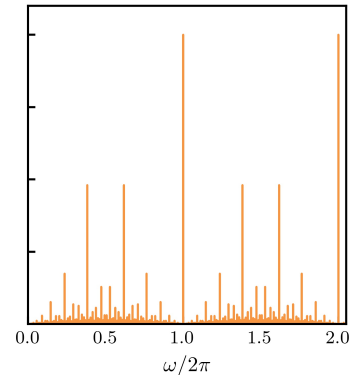
- Project 2D lattice at irrational angle φ :



$$U_5 = U_1 U_0 U_1 U_0 U_1 U_1 U_0 U_1$$

Quasiperiodic Spectrum

- Incommensurate Frequencies: $1, \varphi$
- Spectrum is dense, but discrete



- What happens to MBL?
All transitions possible but suppressed

[Levine-Steinhardt]

$$\Delta \sim |E - n - m\varphi| \quad \text{vs} \quad \Gamma \sim (\gamma_1/J)^n (\gamma_2/J)^m$$

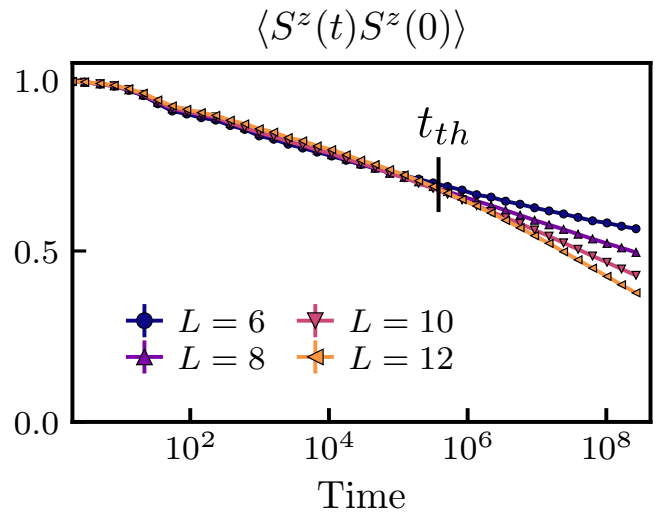
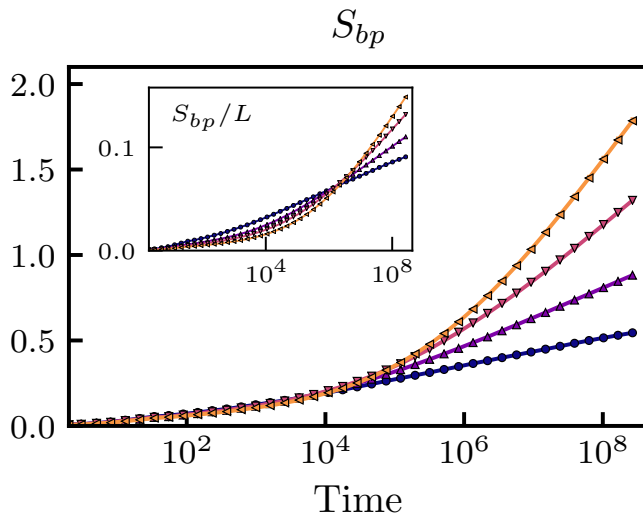
- Protocol:

$$U_0 = e^{-i\lambda H_+}, U_1 = e^{-i\lambda H_-}$$

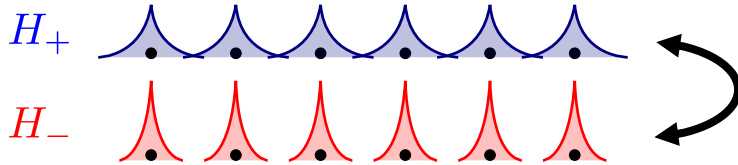
$$H_{\pm} = \pm\delta J \sum_i \mathbf{S}_i \cdot \mathbf{S}_{i+1} + \sum_i h_i S_i^z, \quad h_i \in [-2\pi, 2\pi)$$

Slow Thermalization

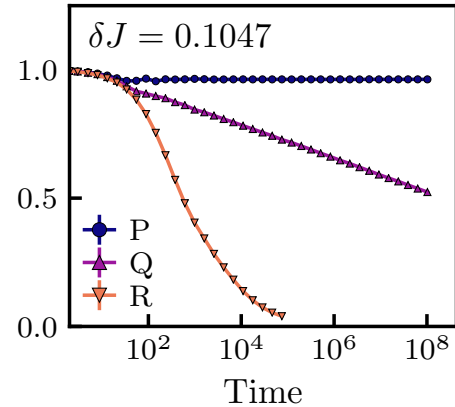
- Three Regimes:
Initially Coherent, Glassy Decay, Fast Thermalization



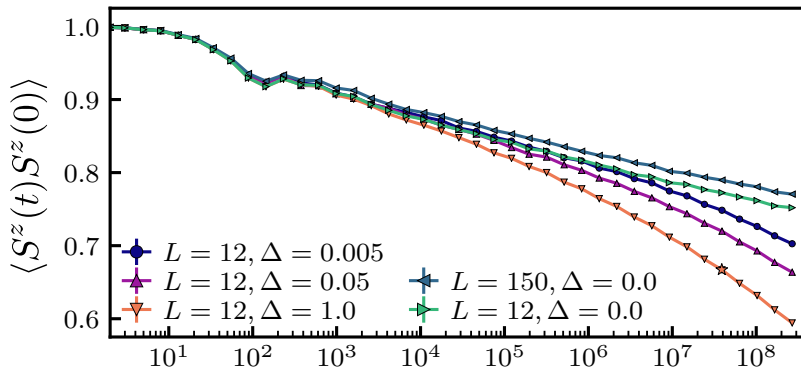
Are there L-Bits?



- $\langle S^z(t) S^z(0) \rangle$ does not saturate:
no conserved spin component



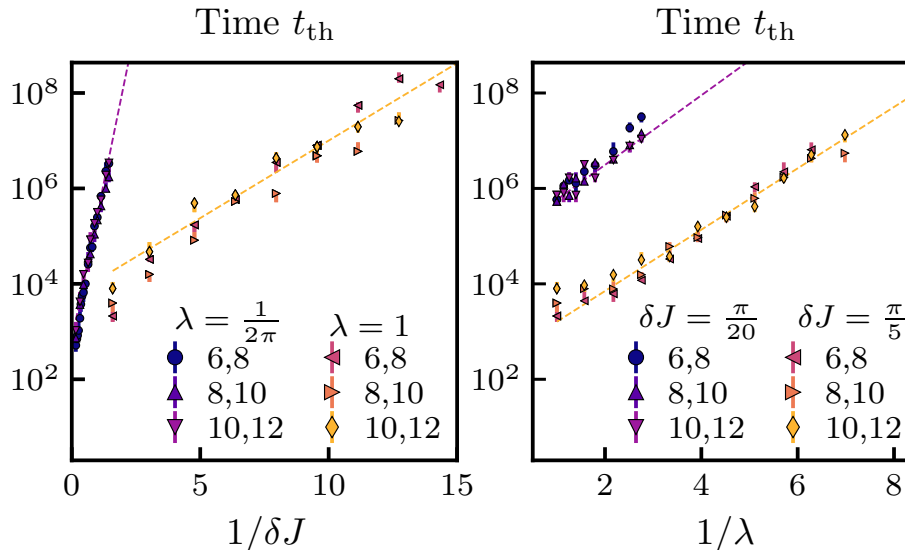
- But log decay in non-interacting limit:



Glassy Regime is Long-lived

- Thermalization time t_{th} depends exponentially on parameters

$$t_{th} \sim e^{1/\lambda}, t_{th} \sim e^{1/\delta J}$$



Effective Hamiltonian

- Effective Hamiltonian for the glassy regime?
- Periodic Driving: Floquet Operator

$$U_F = e^{\Omega_0} e^{\Omega_1} = e^{-iH_{\text{eff}}T},$$

$$U_F \sim e^{\Omega_0 + \Omega_1 + \frac{1}{2}[\Omega_0, \Omega_1] + \frac{1}{12}([\Omega_0, [\Omega_0, \Omega_1]] + [[\Omega_0, \Omega_1], \Omega_1]) + \dots}$$

- For Fibonacci Drive:

$$\left. \begin{array}{l} U_0 = U_0 \\ U_1 = U_1 \\ U_0 U_1 = U_2 \\ U_1 U_0 U_1 = U_3 \\ U_0 U_1 U_1 U_0 U_1 = U_4 \\ U_1 U_0 U_1 U_0 U_1 U_1 U_0 U_1 = U_5 \end{array} \right\} \begin{array}{l} t_n = F_n \sim \varphi^n \\ U_n = e^{\Omega_n} \rightarrow e^{-iH_{\text{eff}}\varphi^n} \end{array}$$

Recursive Magnus Expansion

- Fibonacci Drive has 'deflation' property.
Rule $U_0 \rightarrow U_1, U_1 \rightarrow U_0 U_1$ generates next term:

$$\Omega_{n+1}(\Omega_0, \Omega_1) = \Omega_n(\Omega_1, \Omega_0 * \Omega_1)$$

$$\Omega_0 * \Omega_1 = \log(\exp \Omega_0 \exp \Omega_1)$$

- Expand onto all possible commutators

$$\Omega_n = a_n \Omega_0 + b_n \Omega_1 + h_n [\Omega_0, \Omega_1]$$

$$+ f_n [\Omega_0, [\Omega_0, \Omega_1]] + g_n [[\Omega_0, \Omega_1], \Omega_1] + \dots$$

Order-by-order difference equations are solved recursively

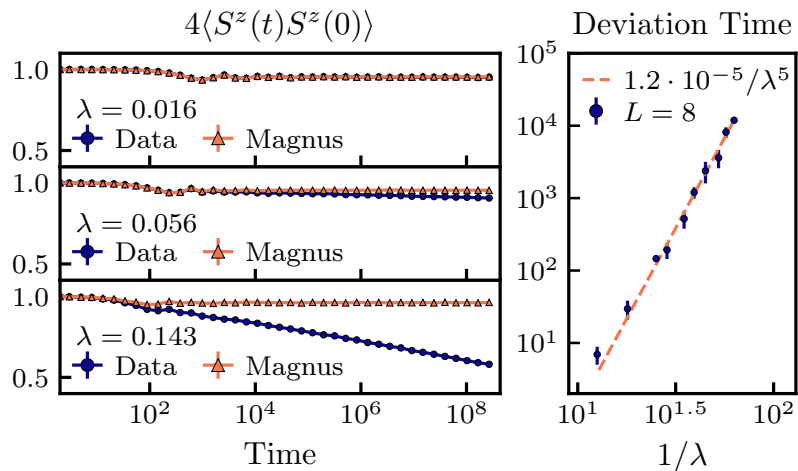
Breakdown of Magnus Expansion

- Asymptotic properties, order k of expansion

$$a_n \sim \varphi^{n|k-2|}$$

- Hamiltonian interpretation breaks down at $k=4$
- Up to $k=3$:

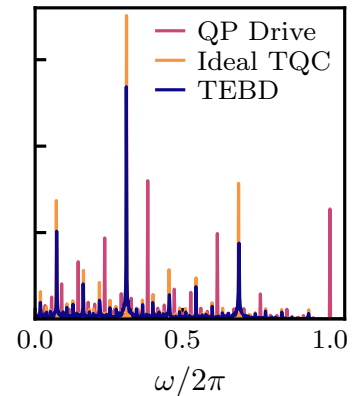
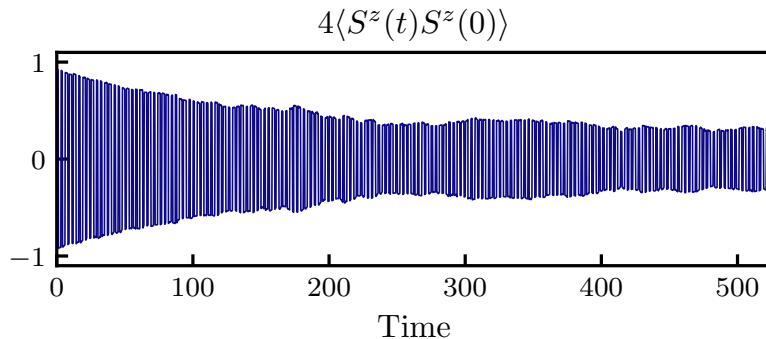
- Expansion works
- No Decay
- Resummation possible?



Time Quasicrystal

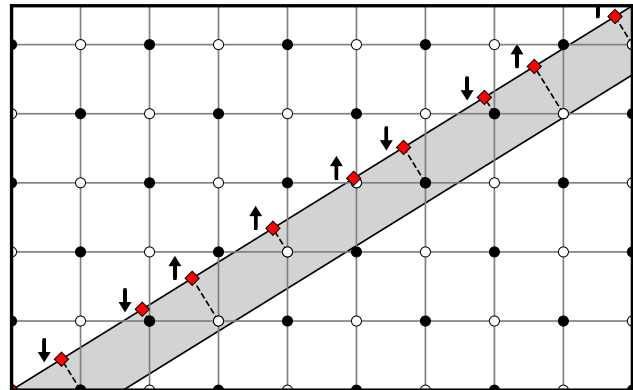
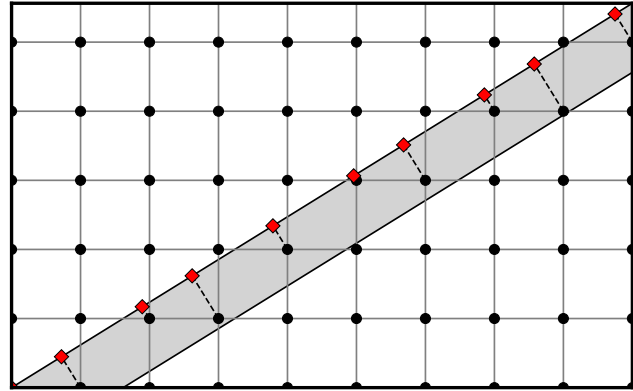
- Use glassy regime for new quantum phases [holds as good as log decay]
- Simplest case: break down a symmetry of the quasiperiodic time drive ['time quasicrystal']

$$U_0 = e^{-i\theta \sum_i S_i^x}, U_1 = e^{-i\lambda \sum_i (J_i S_i^z S_{i+1}^z + h_i^z S_i^z + h_i^x S_i^x)}.$$



Time Quasicrystal Symmetry

- Time Quasicrystal is symmetry breaking in the 2d space defining sequence
- Pattern survives the projection to 1d
- Dual to QP Majoranas [Peng, Refael]
- Generalize to other topological phases, but projection



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[arXiv:1708.00865]