

Friction, Failure, and Dynamics of Granular Materials

KITP, Friction Fracture and Earthquakes

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OUTLINE

- Why granular materials?
- Where granular materials and molecular matter part company—open questions of relevant scales
- Review of models for force transmission
- Obtaining vector forces
- Force distributions and correlations
- Diffusion, dynamics, fluctuations in shear
- Dynamical transition: freezing by heating in a sheared and shaken system
- Force Transmission (via an older experimental approach)
- Conclusions

Examples of Granular Materials

- Earthquake gouge
- Avalanches and mudslides
- Food and other natural grains: wheat, rice,...
- Industrial materials: coal, ores,...
- Soils and sands
- Pharmaceutical powders
- Dust
- Chemical processing—e.g. fluidized beds

What are Granular Materials?

- Collections of macroscopic ‘hard’ (but not rigid) particles: interactions are dissipative
 - Classical $\hbar \rightarrow 0$
 - A-thermal $T \rightarrow 0$
 - Draw energy for fluctuations from macroscopic flow
 - Exist in phases: granular gases, fluids and solids
 - Large collective systems, but outside normal statistical physics

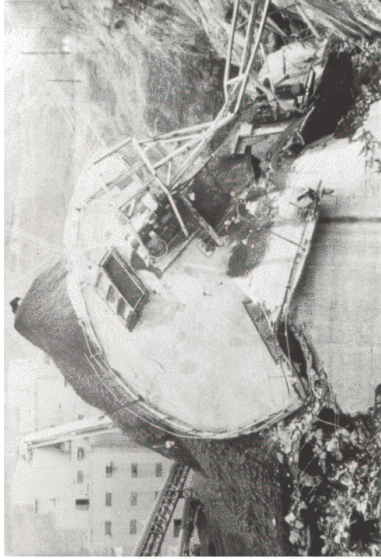
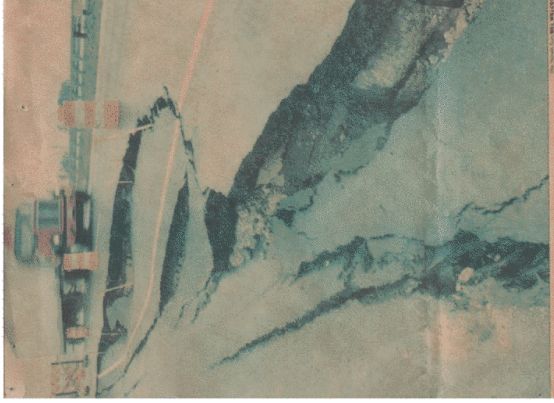
Questions

- Fascinating and deep statistical questions
 - What is the nature of granular friction?
 - What is the nature of granular fluctuations—what is their range?
 - Is there a granular temperature?
 - Phase transitions
 - Jamming and connections to other systems: e.g. colloids, foams, glasses,...
 - The continuum limit and ‘hydrodynamics—at what scales?
 - What are the relevant macroscopic variables?
 - Novel instabilities and pattern formation phenomena

Practical Issues

- o Massive financial costs Claim:
~\$1 Trillion/year in US for granular handling
- o Failures are frequent, typical facilities operate at only ~65% of design
- o Soil stability is difficult to predict/assess
- o How is stress/information transmitted in granular materials?

Problems closer to (my) home



....And a bit further from home...



Assessment of theoretical understanding

- Basic models for dilute granular systems are reasonably successful
- For dense granular states, theory is far from settled, and under intensive debate and scrutiny
- Current ability to predict for dense granular states is poor relative to other systems—e.g. fluids

Granular Material Phases-Dense Phases

Granular Solids and fluids much less well understood than granular gases

Forces are carried preferentially on force chains → multiscale phenomena

Deformation leads to large spatio-temporal fluctuations

Granular Material Phases-Dense Phases Continued

Friction and extra contacts → preparation history matters

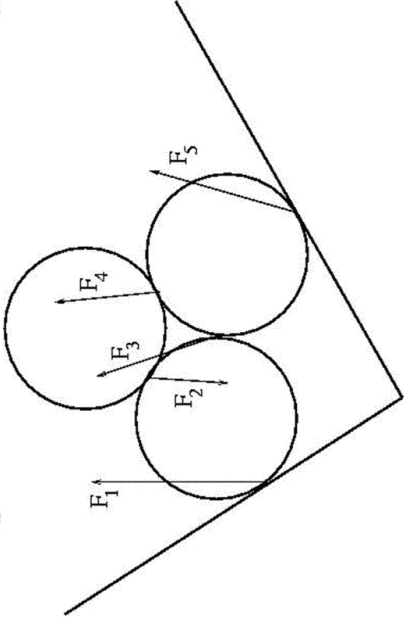
Jamming/glassy behavior near solid-fluid transition
(Liu, Nagle, O'Hern, Bouchaud et al.)

--interesting connections to plasticity in disordered solids (e.g. Falk, Langer, Lemaitre, Maloney...)

In most cases, a statistical approach may be the only possible description

A look at some properties that affect dense granular states

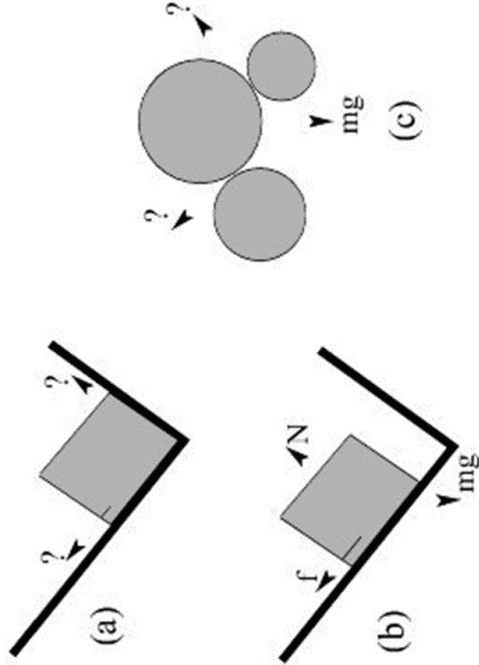
Multiple contacts => indeterminacy



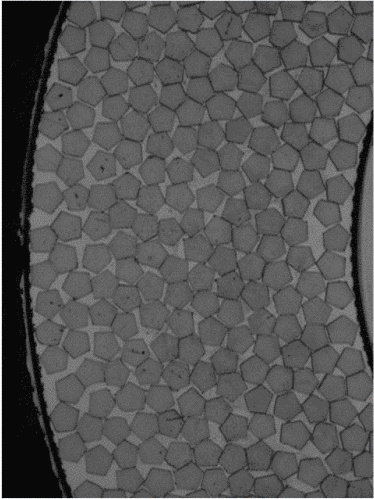
Note: 5 contacts => 10 unknown force components.

3 particles => 9 constraints

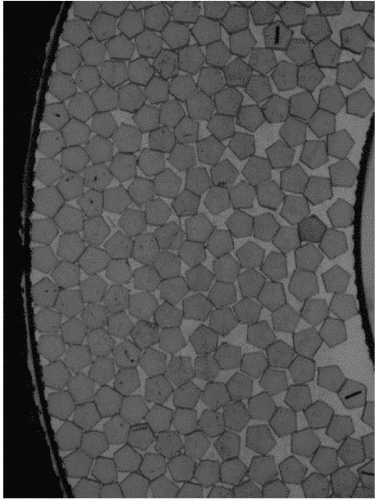
Frictional indeterminacy => history dependence



Dilation under shear

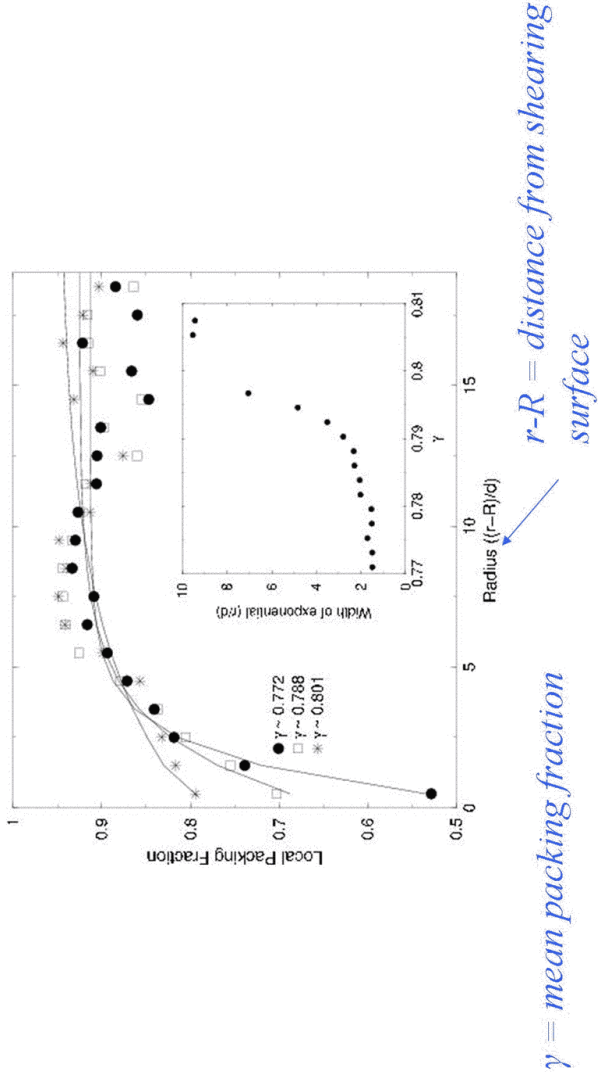


Before shearing



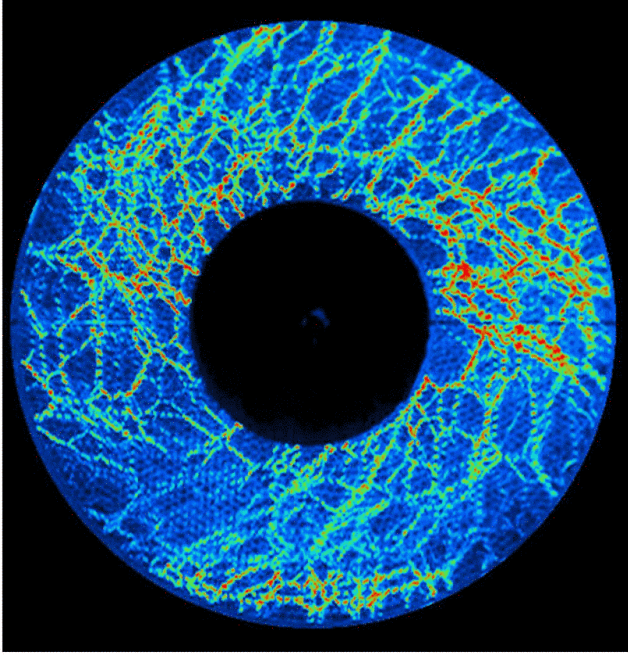
After sustained shearing

Density profiles following sustained shearing



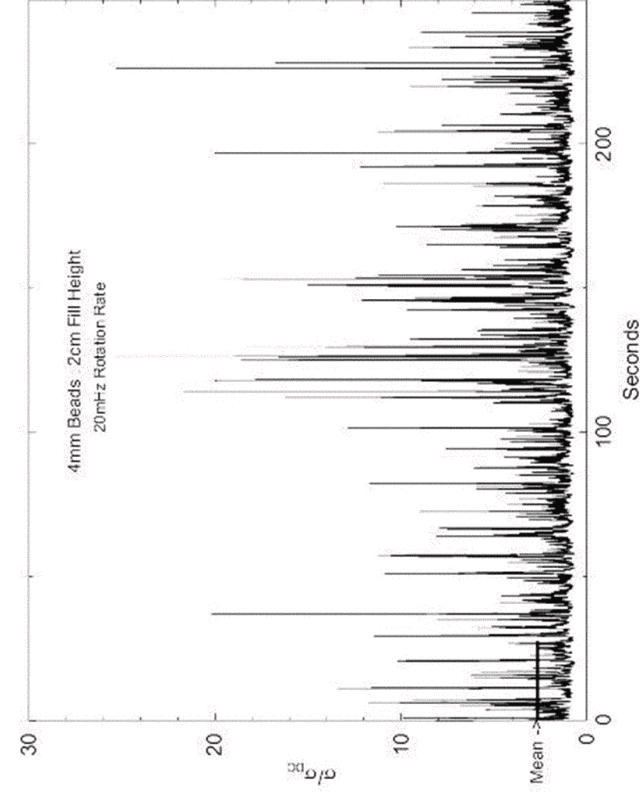
Example of Force Chains—Shear Experiment

Howell et al. PRL 82, 5241 (1999)



Stress Fluctuations in 3D Shear Flow

Miller et al. PRL 77, 3110 (1996)



Video of 2D shear flow



Understanding Static Stress Balance—Ideally from Micromechanics

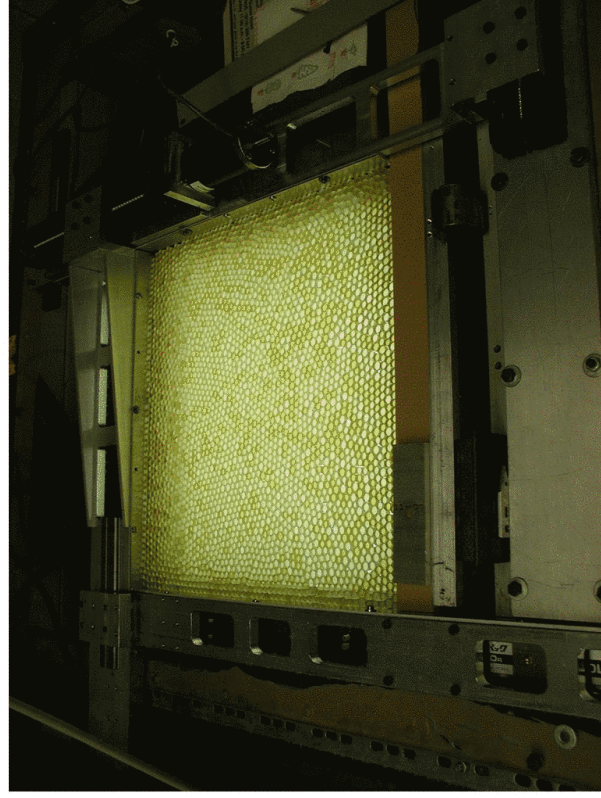
- Four unknown stress components (2D)
- Three balance equations
 - Horizontal forces
 - Vertical forces
 - Torques
- Need a constitutive equation

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xz}}{\partial z} = 0 \qquad \frac{\partial \sigma_{xz}}{\partial x} + \frac{\partial \sigma_{zz}}{\partial z} = 0 \qquad \sigma_{xz} = \sigma_{zx}$$

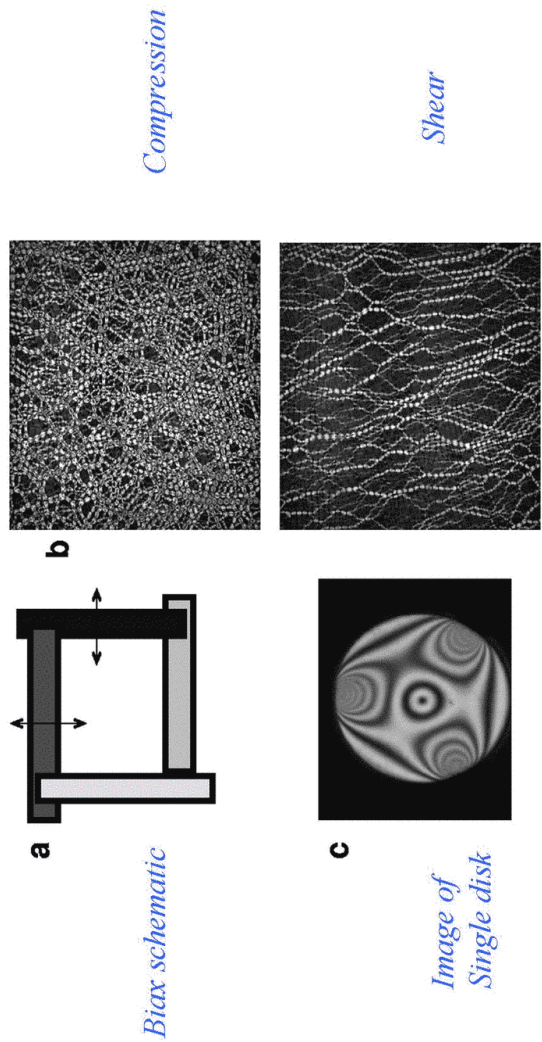
Some approaches to describing stresses

- Elasto-plastic models (Elliptic, then hyperbolic)
- Lattice models
 - Q-model (parabolic in continuum limit)
 - 3-leg model (hyperbolic (elliptic) in cont. limit)
 - Anisotropic elastic spring model
- OSL model (hyperbolic)
- Telegraph model (hyperbolic)
- Double-Y model (type not known in general)

Experiments to determine vector contact forces
(Trush Majmudar and RPB, Nature, June 23, 2005)

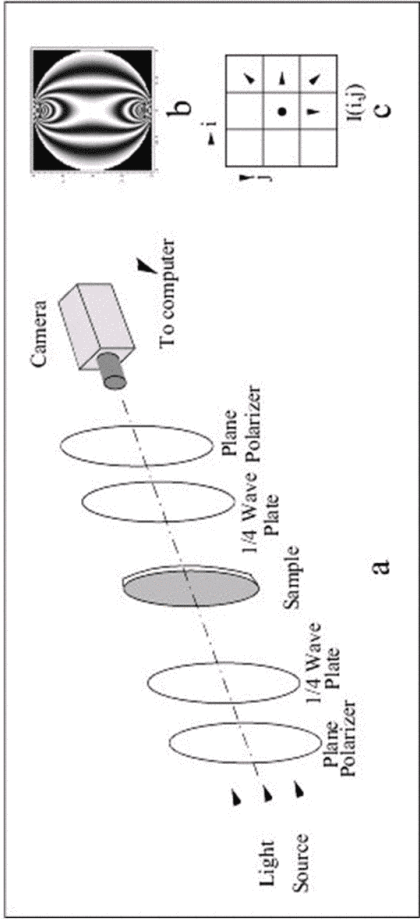


Experiments Use Photoelasticity:



~ 2500 particles, *bi-disperse*, $d_L=0.9\text{cm}$, $d_S=0.8\text{cm}$, $N_S/N_L=4$

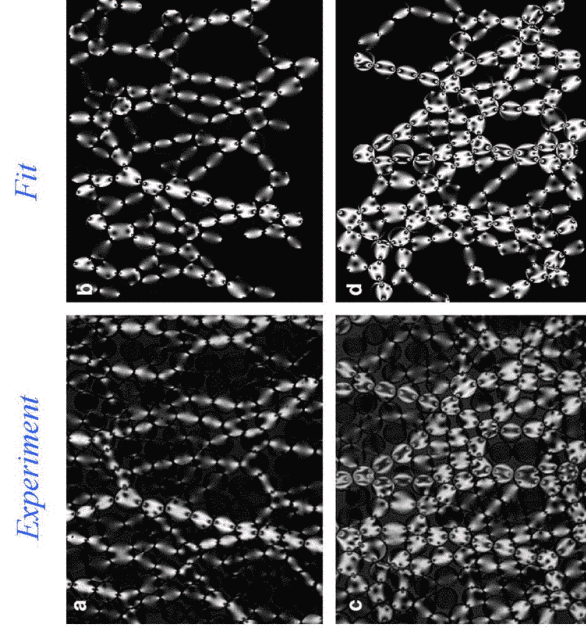
Measuring forces by photoelasticity



Basic principles of technique

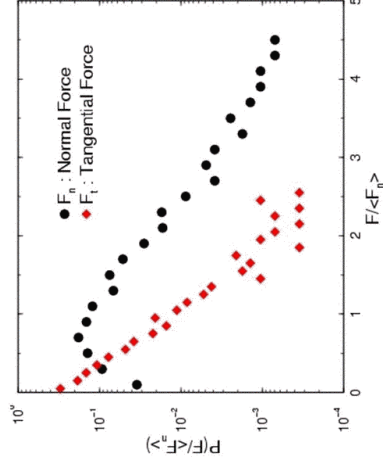
- Process images to obtain particle centers and contacts
- Invoke exact solution of stresses within a disk subject to localized forces at circumference
- Make a nonlinear fit to photoelastic pattern using contact forces as fit parameters
- $I = I_0 \sin^2[(\sigma_2 - \sigma_1)CT/\lambda]$
- In the previous step, invoke force and torque balance
- Newton's 3d law provides error checking

Examples of Experimental and 'Fitted' Images



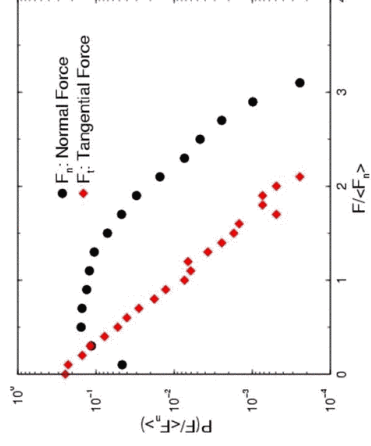
Force distributions for shear and compression

Shear



$$\epsilon_{xx} = -\epsilon_{yy} = 0.04; \quad Z_{avg} = 3.1$$

Compression

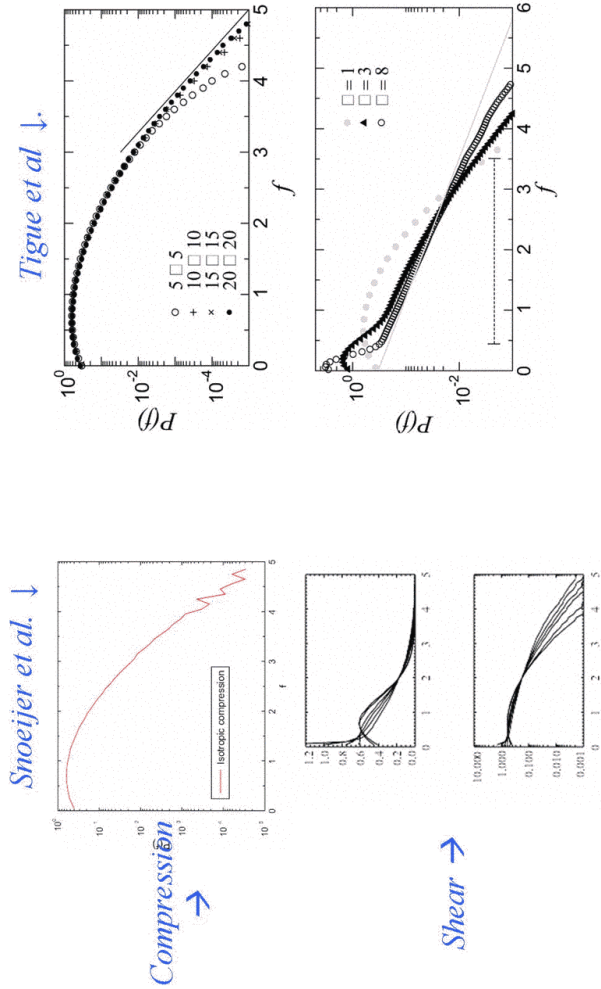


$$\epsilon_{xx} = -\epsilon_{yy} = 0.016; \quad Z_{avg} = 3.7$$

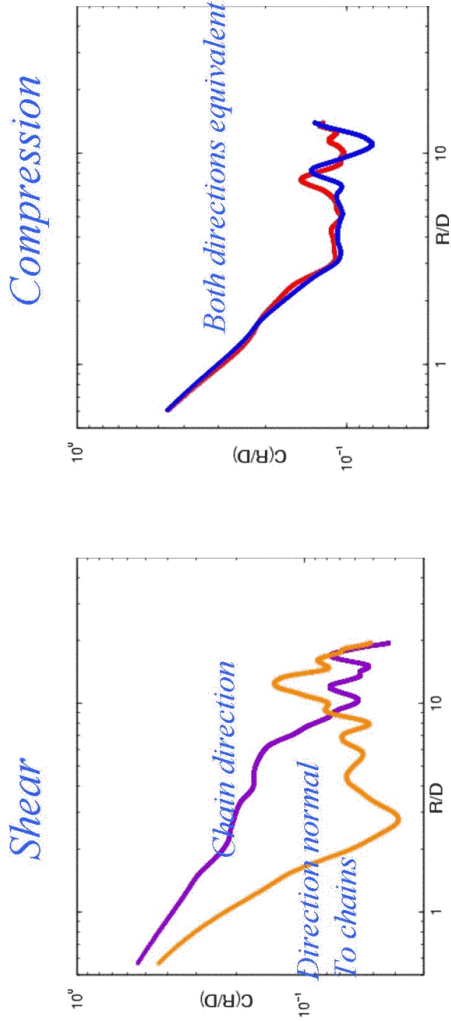
Edwards Entropy-Inspired Models for $P(f)$

- Consider all possible states consistent with applied forces
- Compute Fraction where at least one contact force has value $f \rightarrow P(f)$
- E.g. Snoeier et al. PRL 92, 054302 (2004)
- Tighe et al. preprint (Duke University)

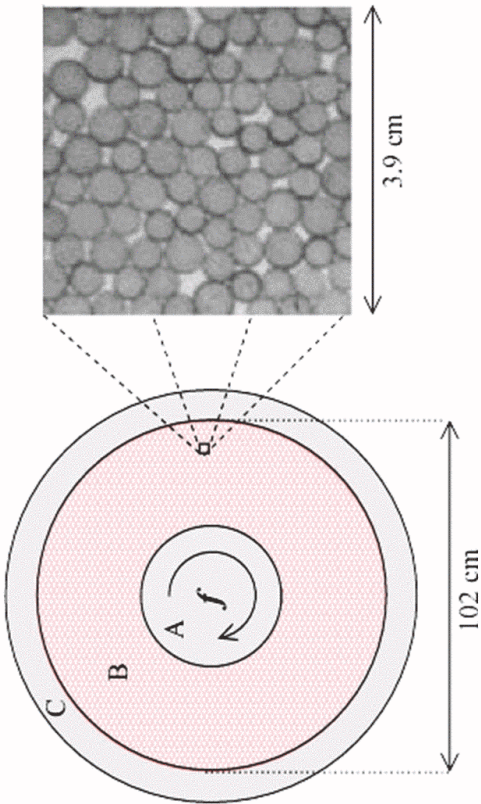
Some Typical Cases—**isotropic compression and shear**



Spatial correlations of forces—angle dependent

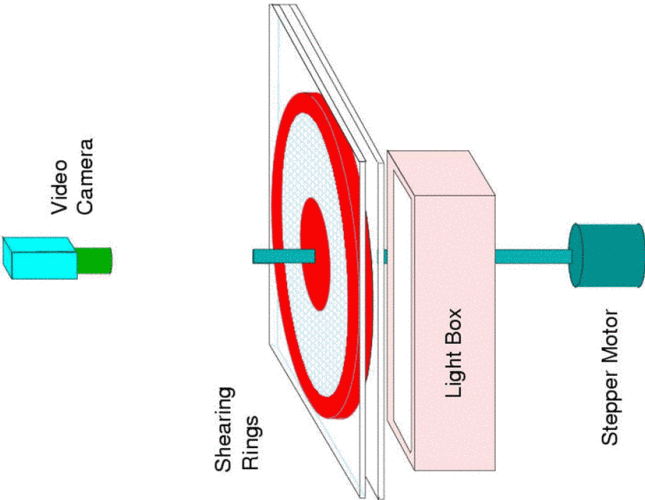


Granular friction and dynamics in a 2D
sheared system

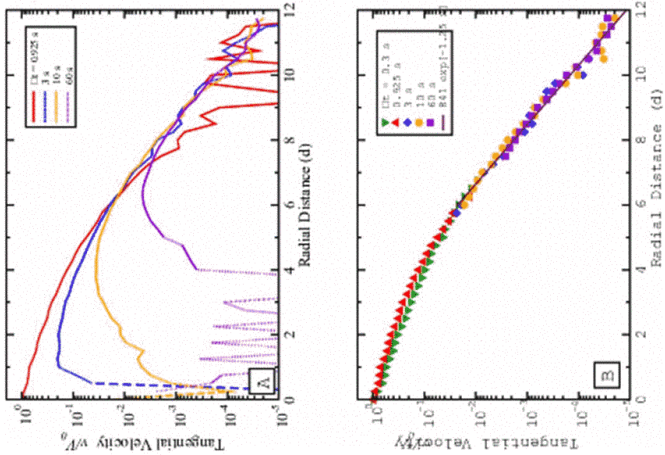


B.Utter and RPB PRE 69, 031308 (2004)
Eur. Phys. J. E 14, 373 (2004)

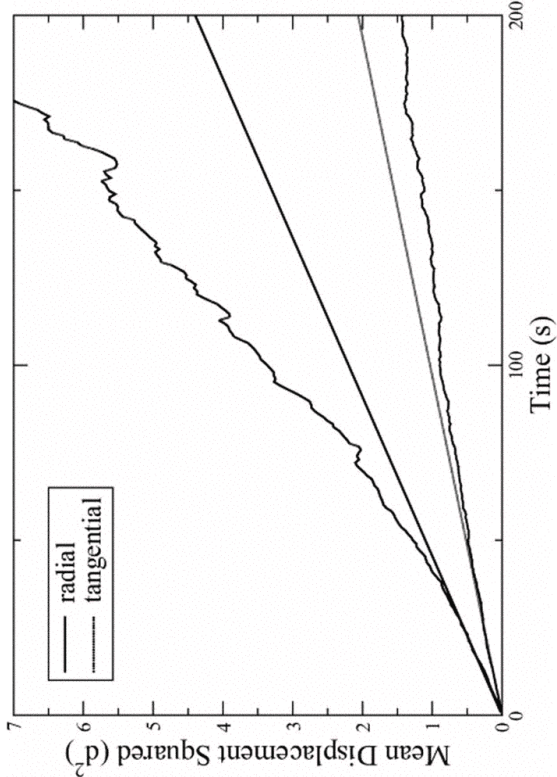
Schematic of apparatus



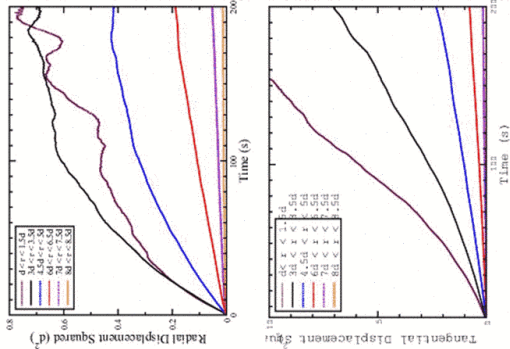
Typical near-exponential velocity profile



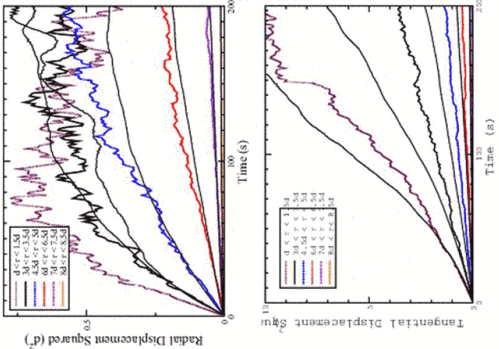
Variances in radial and tangential direction
yield diffusivities



Diffusivities only appear sub- or super-diffusive
due to Taylor-like dispersion and rigid
boundary

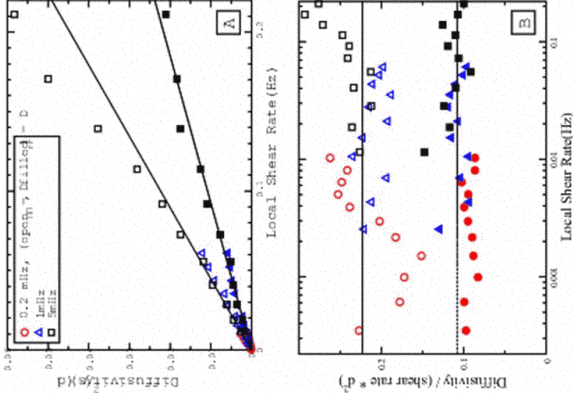


Experiment



Simulations of random
walk, with velocity profile, etc

Local shear rate determines diffusivities

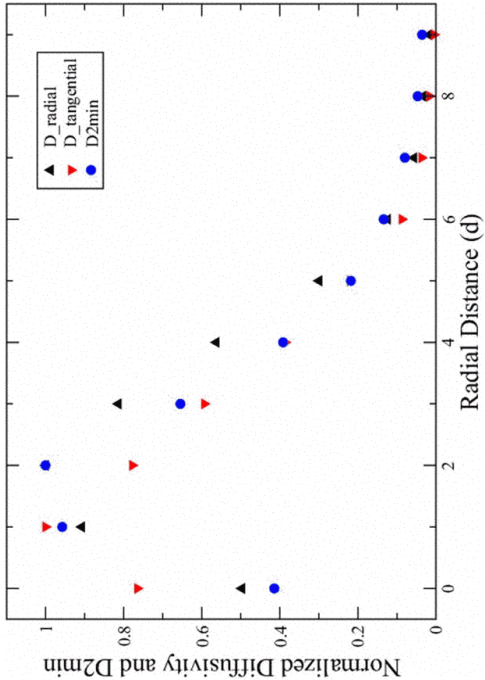


Parameterization of non-affine motion—after Falk and Langer



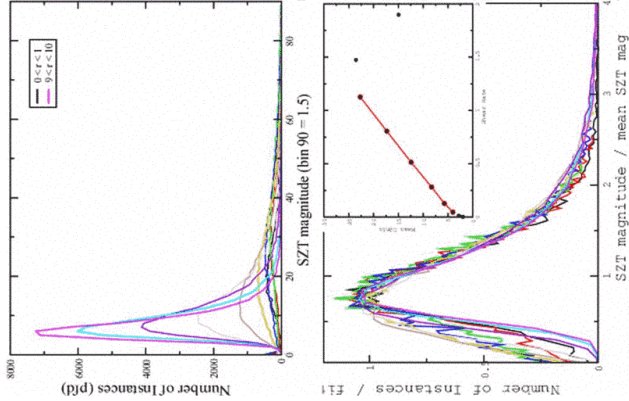
D^2_{min} characterizes local non-affine part of deformation

Relation to D^2_{min} of Falk and Langer



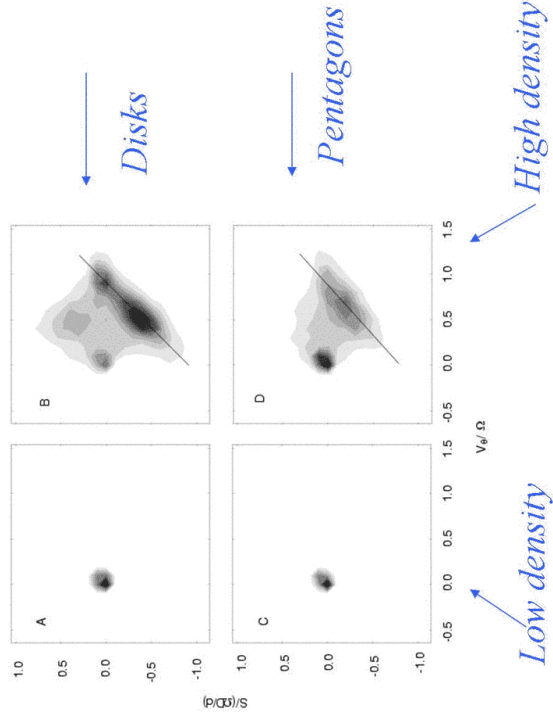
D^2_{min} measures non-affine particle motion

Distributions of D^2_{\min} for different distances from shearing wheel

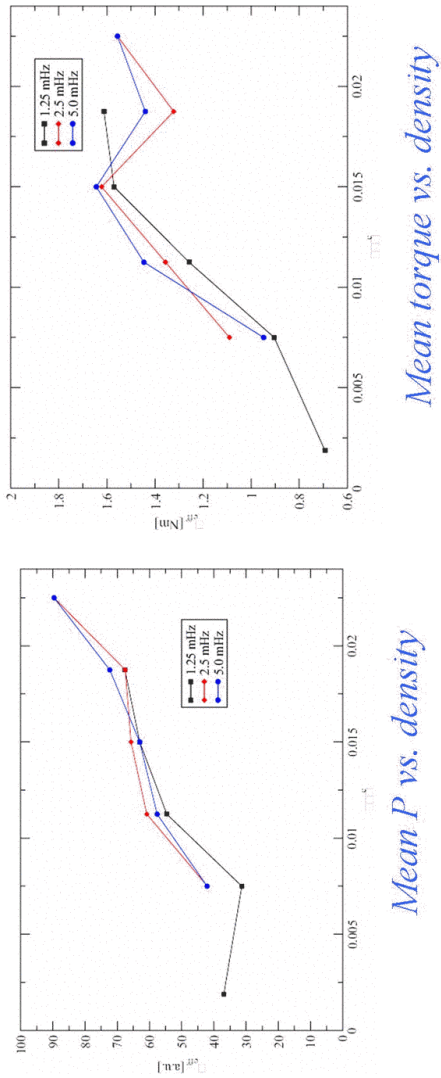


Useful candidate
for measure of
disorder?

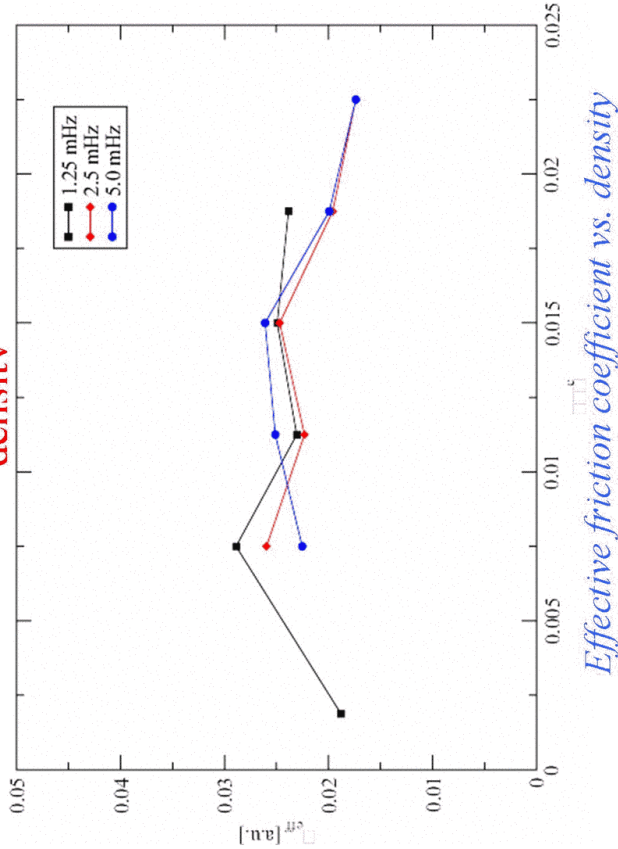
What role does particle rotation play in granular friction?
See greyscale representations of velocity and angular velocity



Measure pressure and shear stress, compare ratio, for different densities

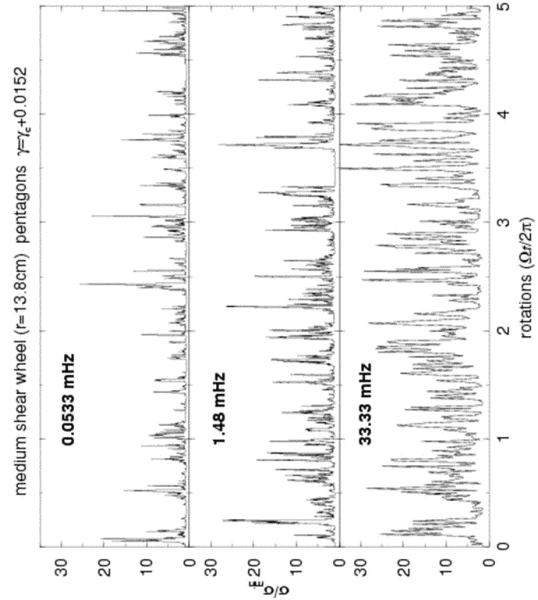


Effective friction coefficient does not change with density



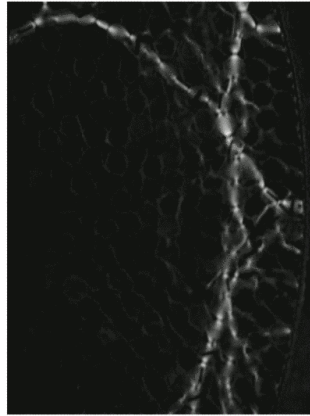
Rate Dependence in Granular Shear

γ gives density

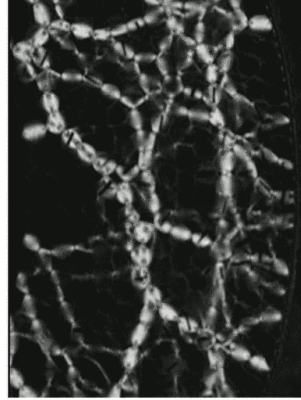


Stress fluctuations vs. time for different shear rates

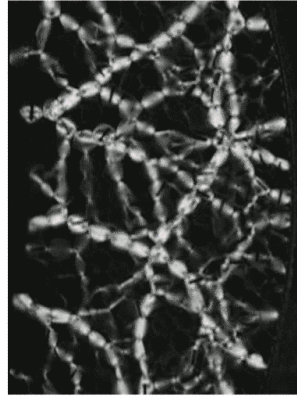
Videos at different shear rates



$\gamma = 0.0027\text{Hz}$

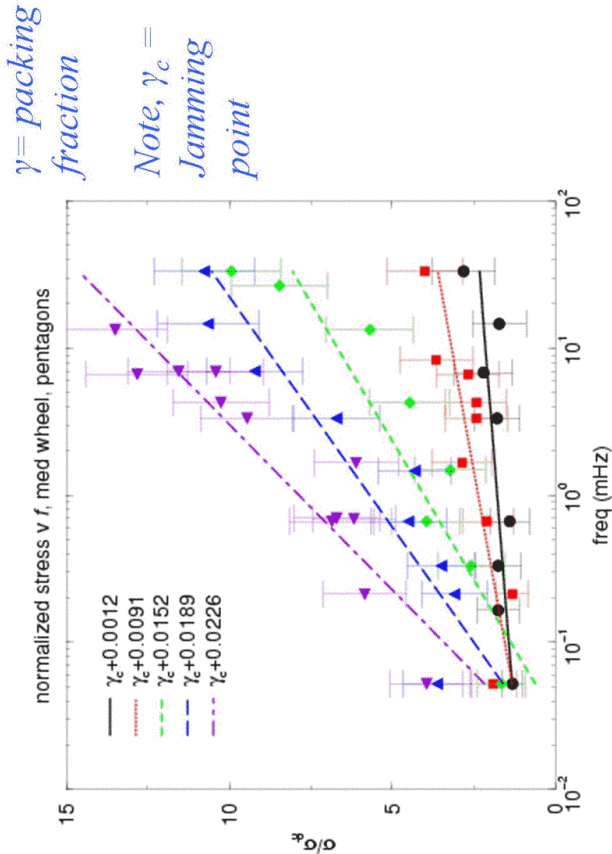


$\gamma = 0.027\text{Hz}$

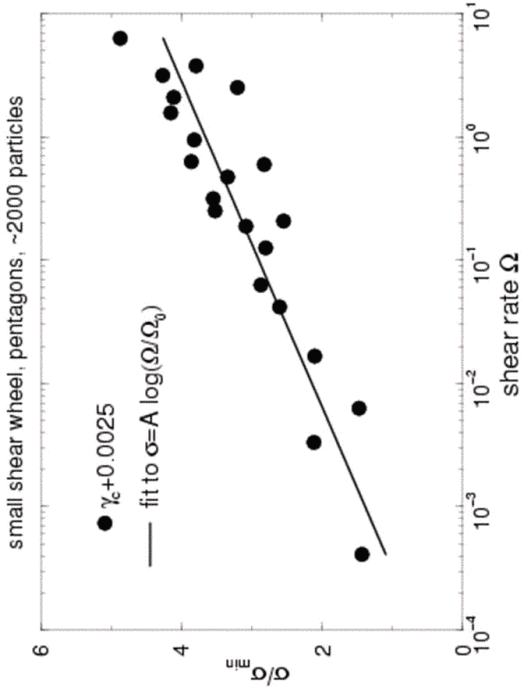


$\gamma = 0.27\text{Hz}$

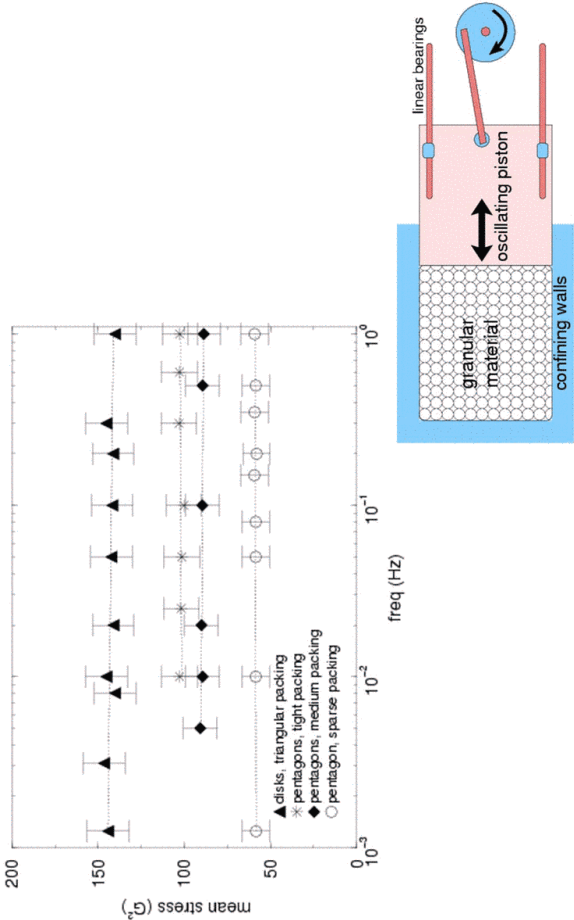
Time-averaged pressure vs. shear rate, and density



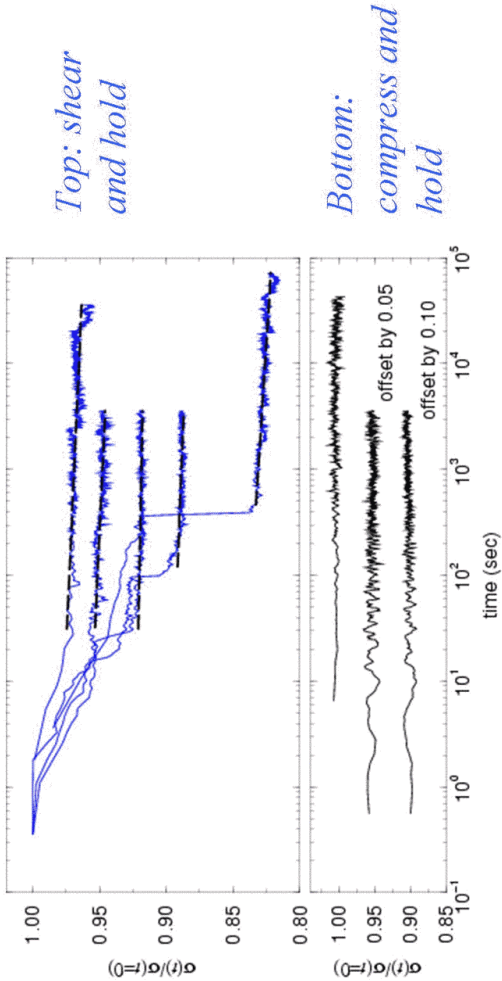
Time-averaged pressure vs. shear rate: 5 decades



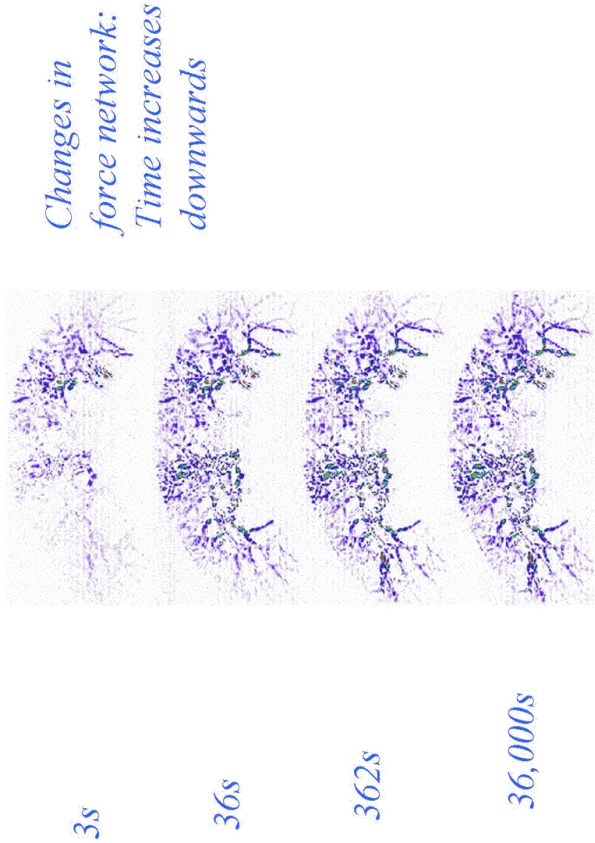
No rate dependence for periodic compression



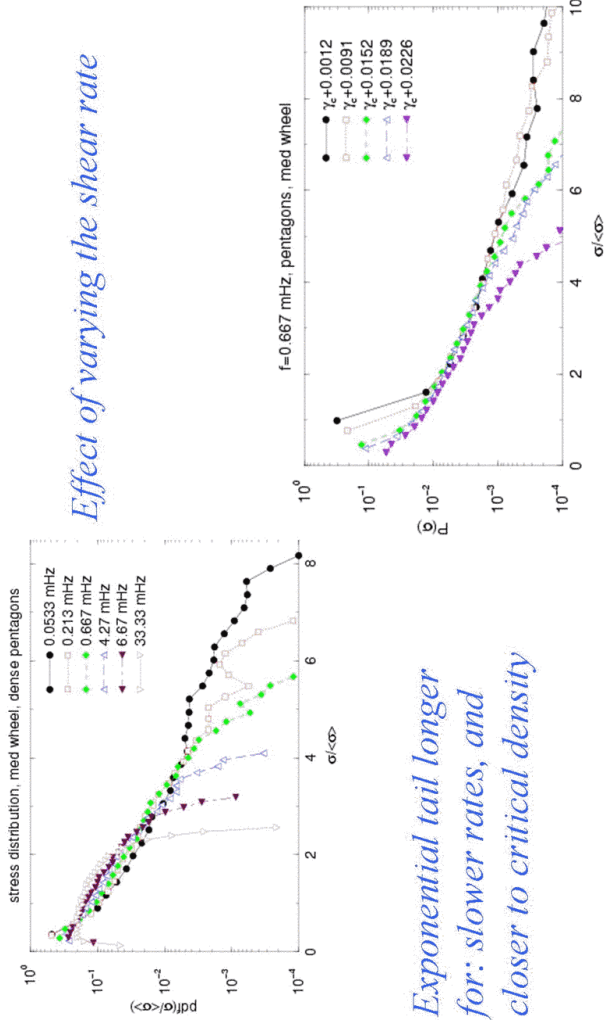
(Un?)Related Creep effects for sheared samples



Visible relaxation in sheared sample



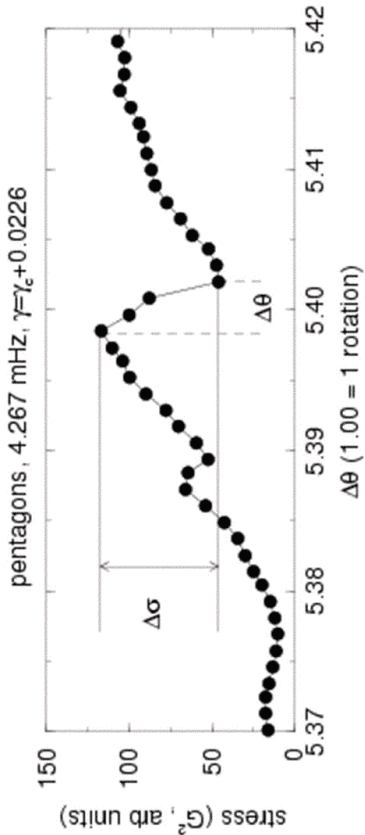
Stress statistics, rate-dependence and ‘avalanches’



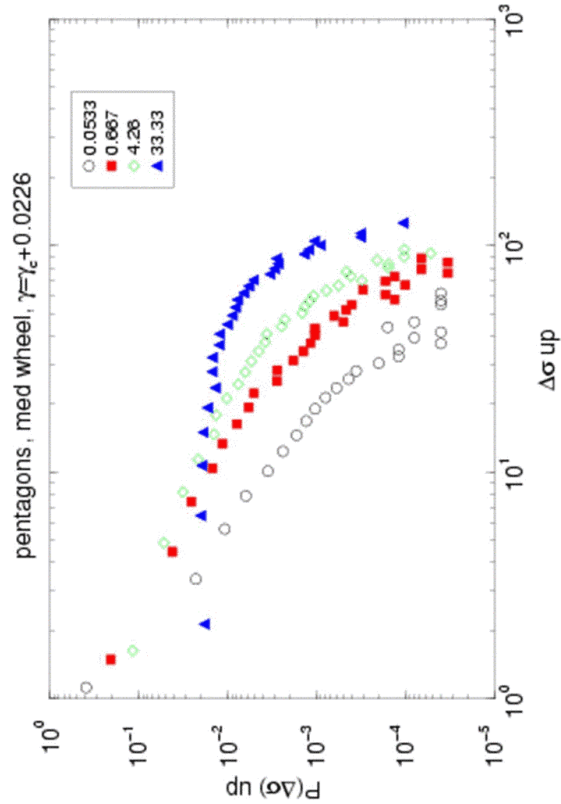
Exponential tail longer
for: slower rates, and
closer to critical density

Effect of varying the density

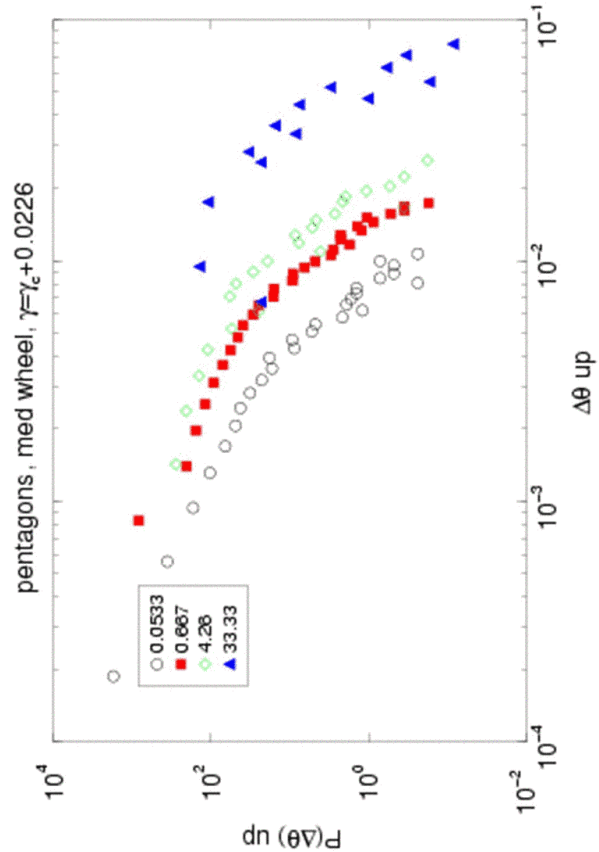
Stress Avalanches



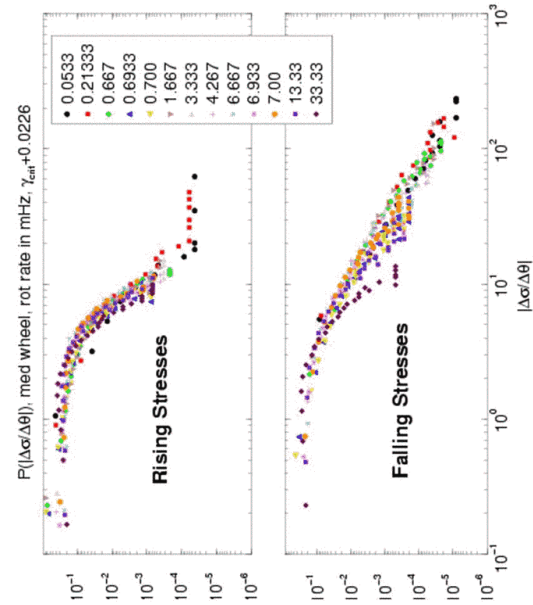
Pdf's for stress changes



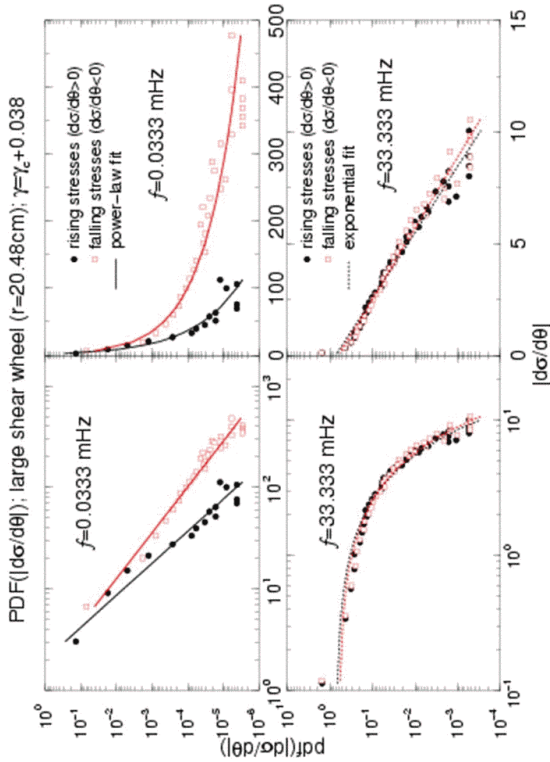
Pdf's for $\Delta\theta$



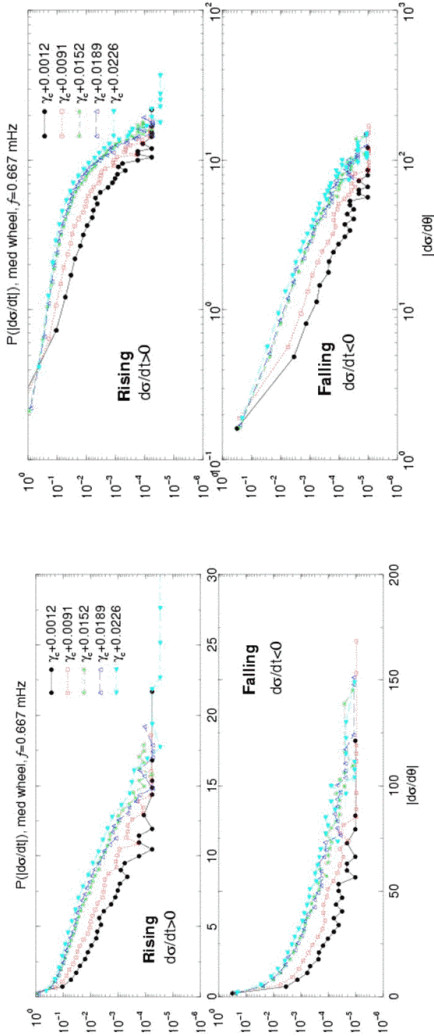
Pdf's for $d\sigma/d\theta$: for various rates



Pdf's for $d\sigma/d\theta$: for various rates



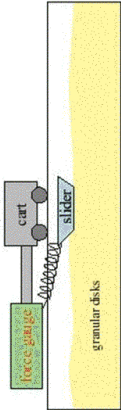
Pdf's for $d\sigma/d\theta$: for various densities



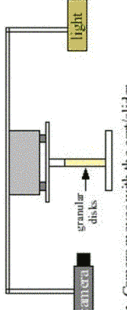
Granular Rheology—a slider experiment

Experimental Apparatus

side view:




end view:

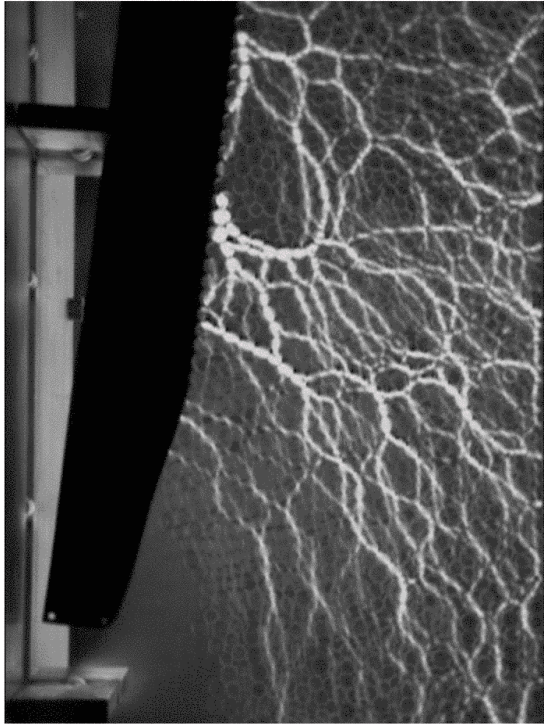


- Cart and force gauge move at constant speed v .
- Slider exhibits stick-slip motion on granular bed.
- Camera moves with the cart/slider.

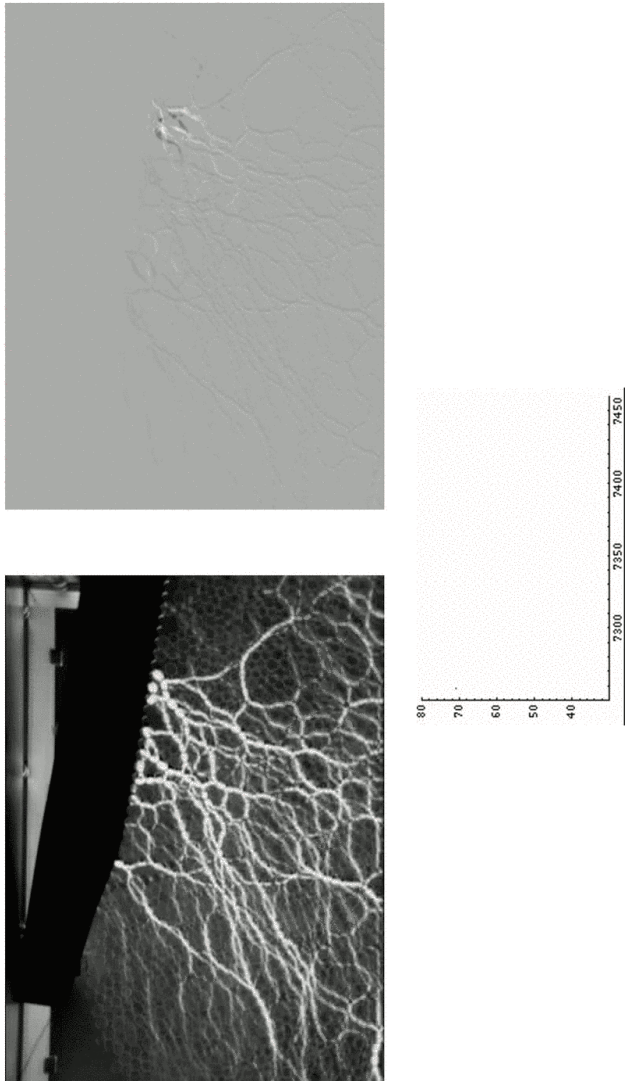
Granular bed = 500 d \times 20 d deep, $d = 0.41$ cm and 0.51 cm bidisperse photoelastic disks.
Typical speeds = 0.1-2 d/s. Slider length = 30-40 d.
Dragging force = 0-100 grams (0-1 Newtons).



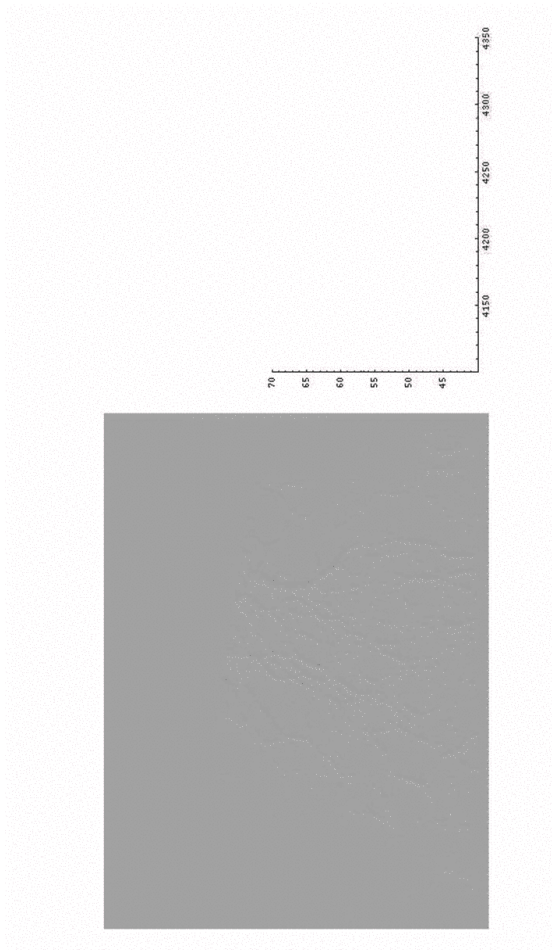
What is the relation between stick slip and granular force structure?



Videos of force evolution



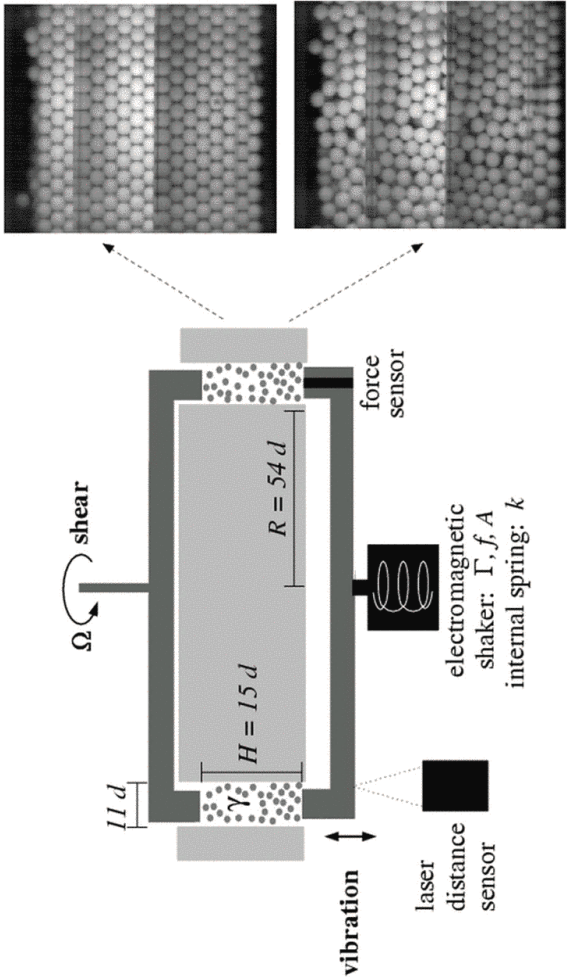
Composite force and photoelastic video: Relation between internal force network and force avalanches



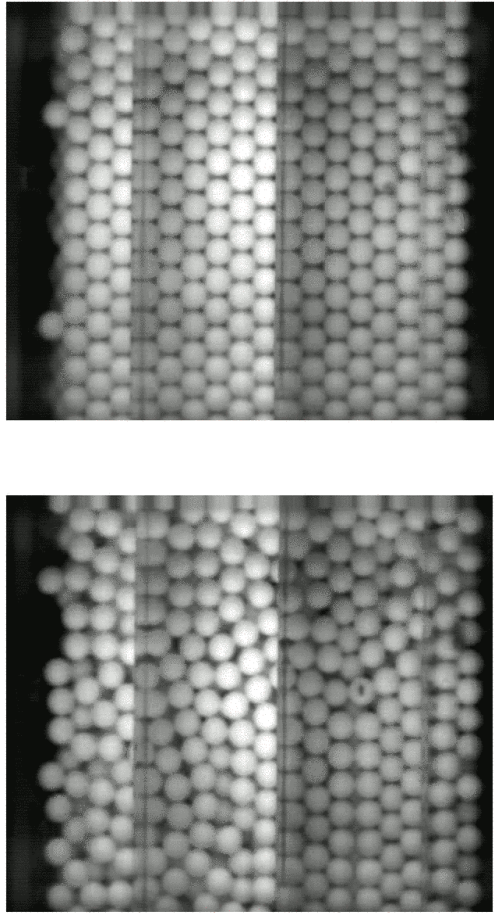
Order-disorder:
Transition from solid to dense fluid

Jamming/unjamming

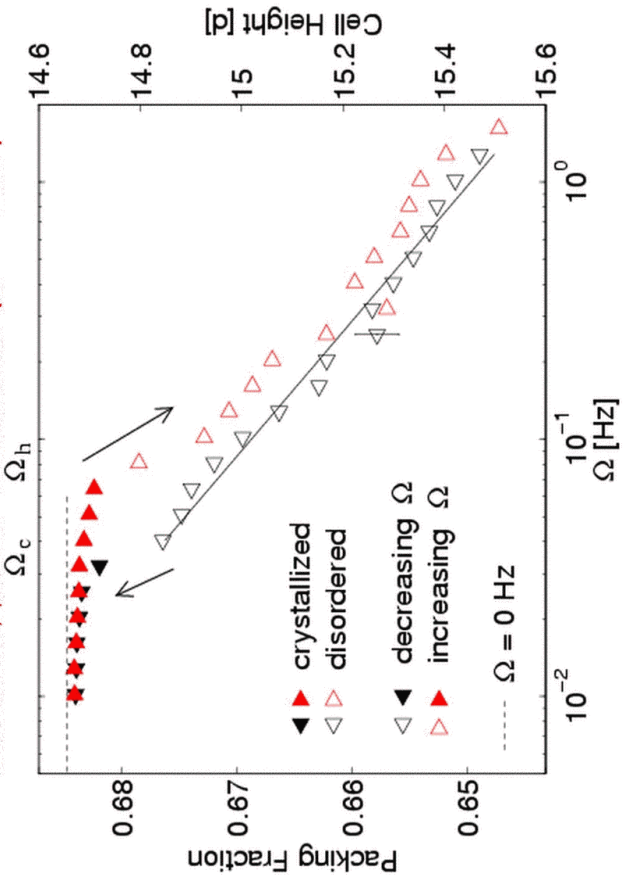
K. Daniels and RPB, PRL 94 168001 (2005)



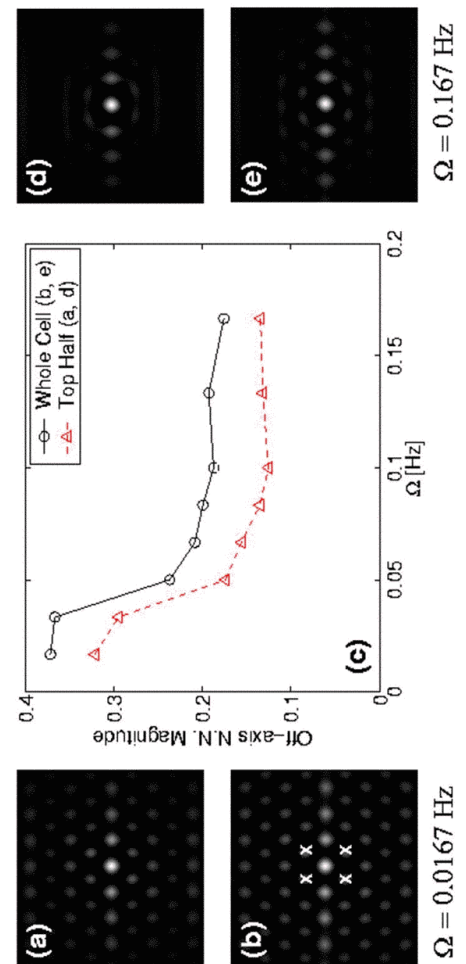
Videos of ordered/disordered states



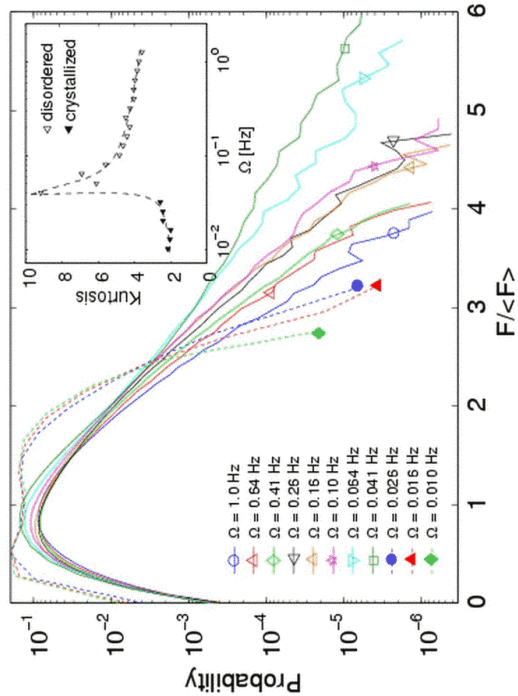
Freezing by Heating—Competition between shearing and vibration ($\Gamma = 2.0$)



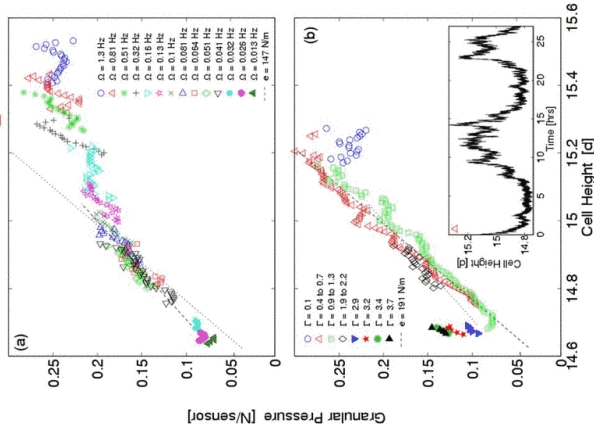
Spatial Autocorrelations show disorder with shear (a, b, d, e) and more quantitatively, c.



Force Probability Distributions: Singular behavior in the Kurtosis

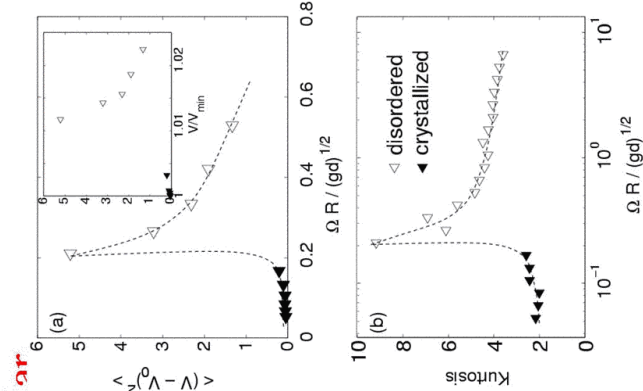


Large-scale Fluctuations, particularly near transition point

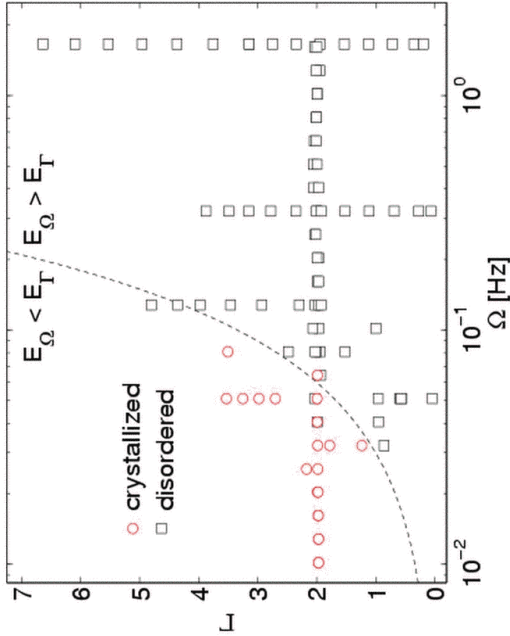


Kurtosis for force pdf and volume fluctuations are singular

Note: volume fluctuations in Edwards entropy picture are analogous to energy fluctuations in Canonical Ensemble → specific heat



Phase Diagram in Shear Rate (Ω) and Shaking Amplitude (Γ)



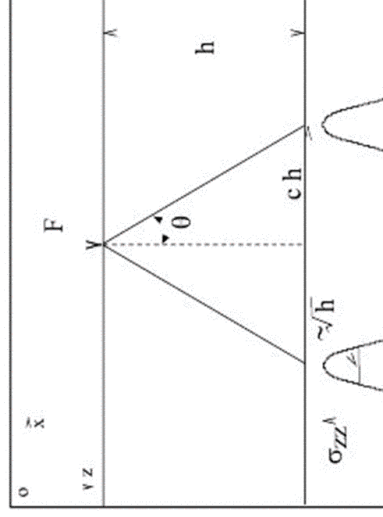
Some approaches to describing stresses

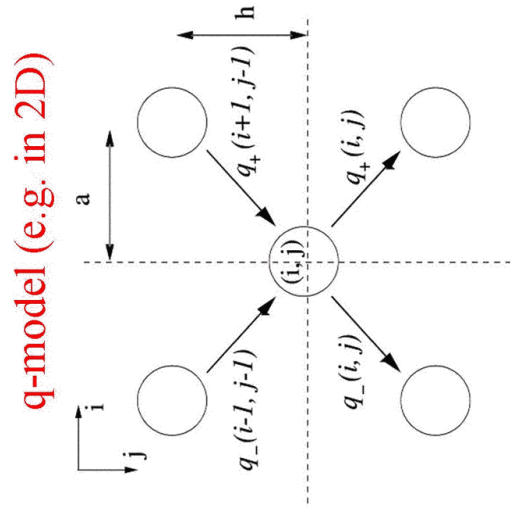
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OSL model—hyperbolic—wave equation

$$\sigma_{xx} = \eta \sigma_{zz} + \mu \sigma_{xz} \quad \eta, \mu: \text{phenomenological parameters}$$

$$\sigma_{zz}(x, z) = \frac{F}{2} [\delta(x + cz) + \delta(x - cz)]$$





q's chosen from uniform distribution on [0,1]

Predicts force distributions ~ exp(-F/F₀)

Long wavelength description is a diffusion equation

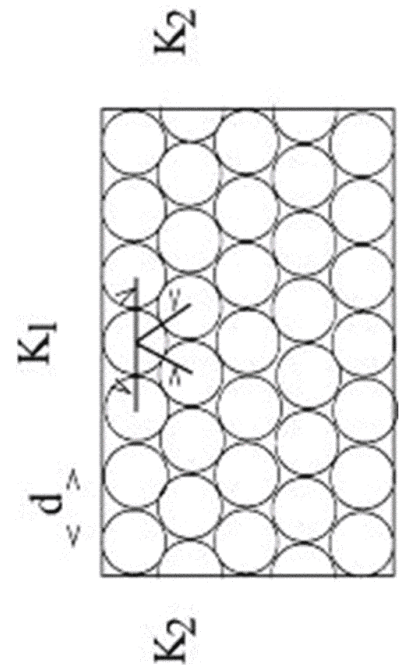
$$\frac{\partial w(z, j)}{\partial z} = \beta[w(z, j + 1) + w(z, j - 1) - 2w(z, j)]$$

$$\frac{\partial w}{\partial z} = D \frac{\partial^2 w}{\partial x^2}$$

Expected stress variation with depth

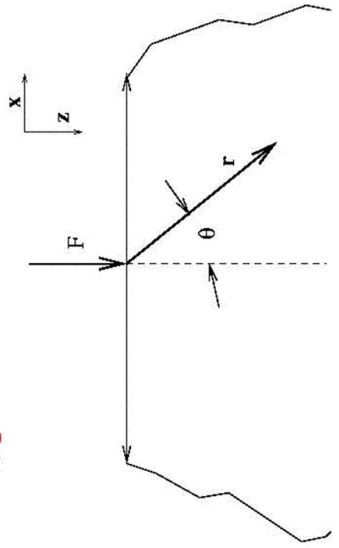
$$\sigma_{zz}(x, z) = \frac{F}{2\sqrt{\pi Dz}} \exp(-x^2 / 4Dz)$$

Anisotropic elastic lattice model



*Expect propagation along lattice directions
Linear widening with depth—e.g. Goldenberg
and Goldhirsch, Nature 435, 188 (2005)*

Elastic response, point force on a semi-infinite sheet



$$\sigma_{rr} = \frac{2F \cos \theta}{r\pi} \qquad \sigma_{r\theta} = \sigma_{\theta\theta} = 0$$

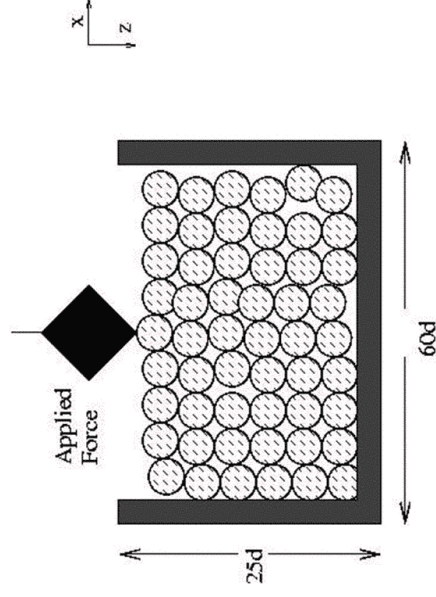
In Cartesian coordinates:

$$\sigma_{ii} = 1/[z(1 + (x/z)^2)]^p \qquad p = 1,2$$

Force response/transmission: what is the mechanical response to a small point force?

Physica D 182, 274 (2003)

Schematic of greens function apparatus



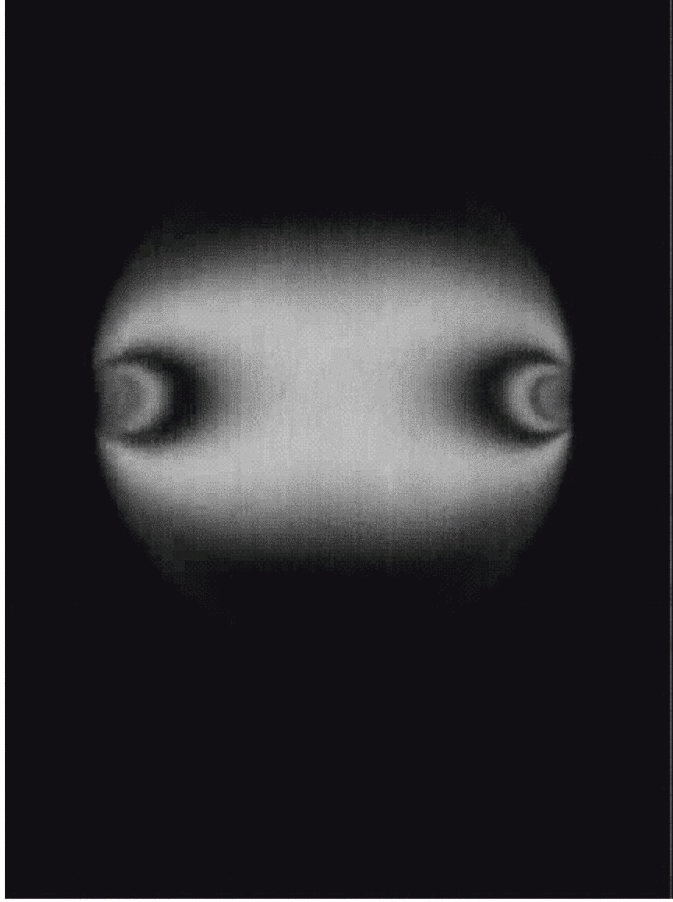
d = grain size

Use: 1) Monodisperse disks (spatially ordered)

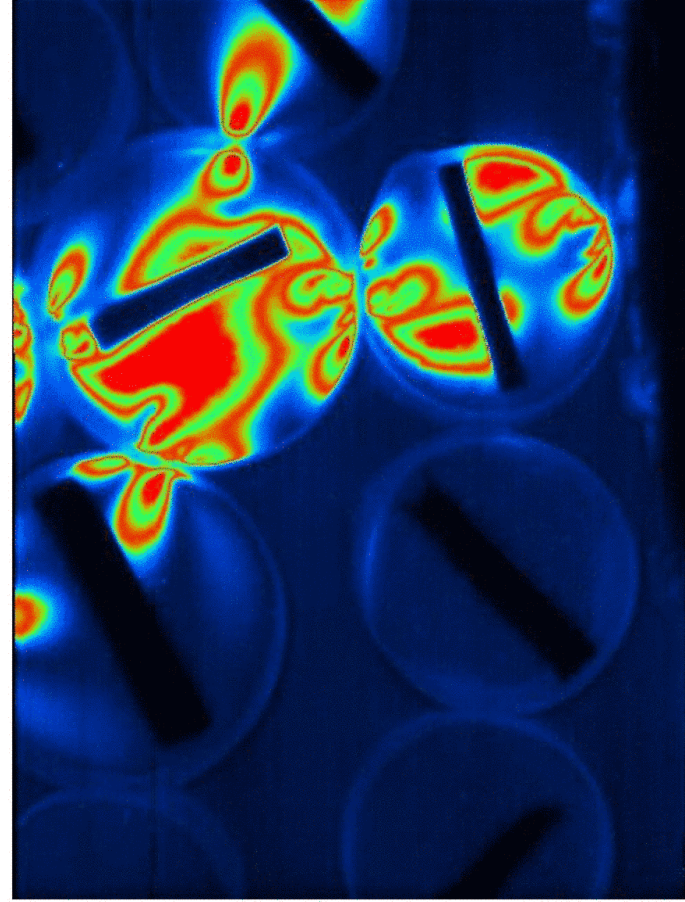
2) Bidisperse disks (weakly disordered)

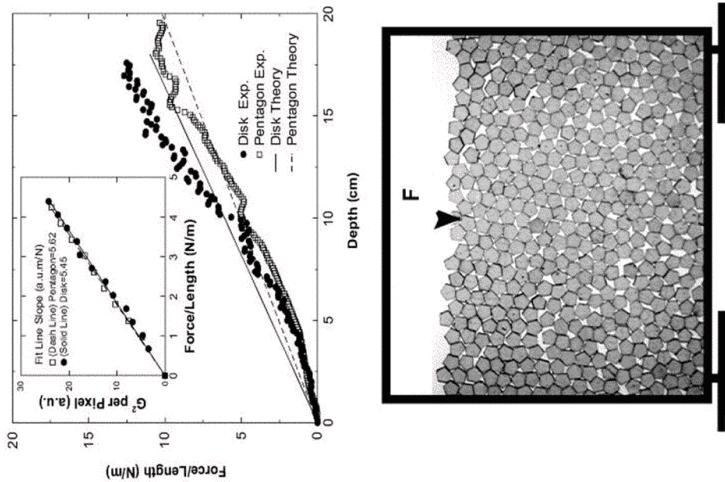
3) Pentagons (strongly disordered)

Diametrically opposed forces on a disk



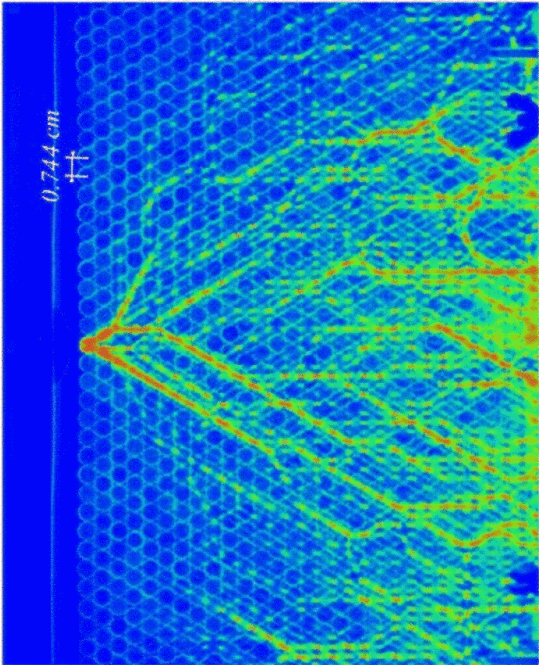
A gradient technique to obtain grain-scale forces



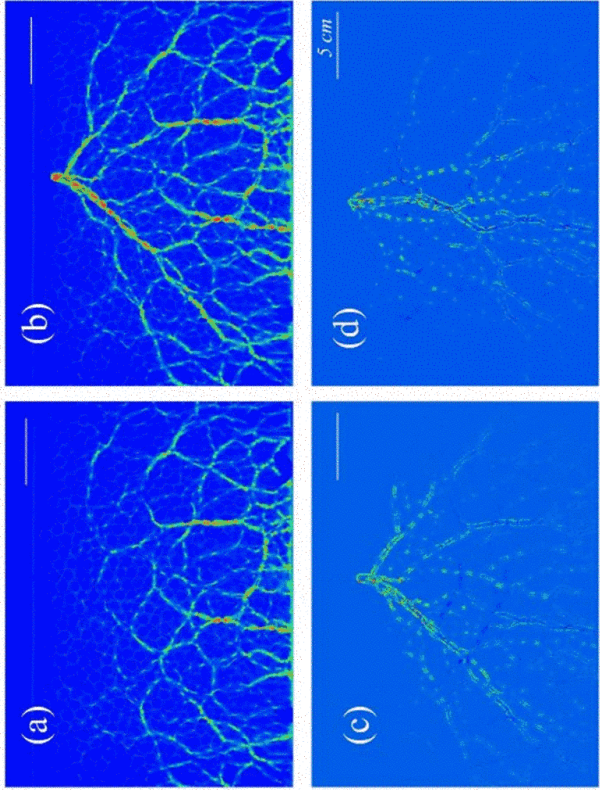


calibration

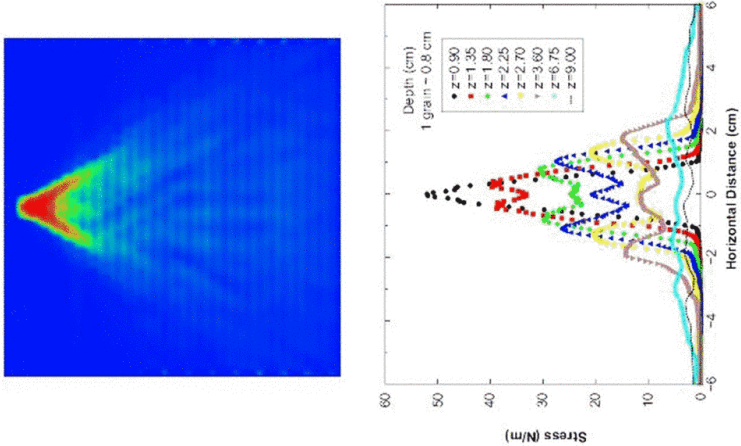
Disks-single response



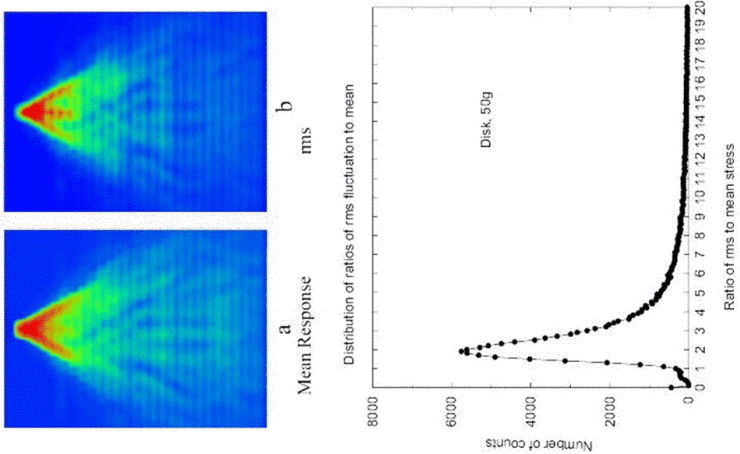
Before-after



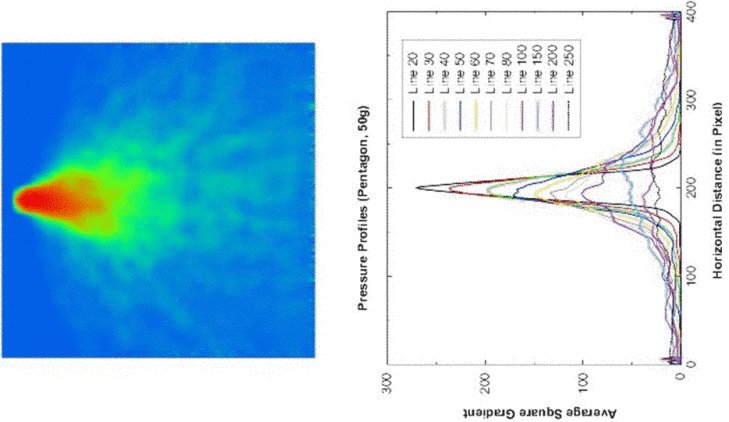
disk response mean



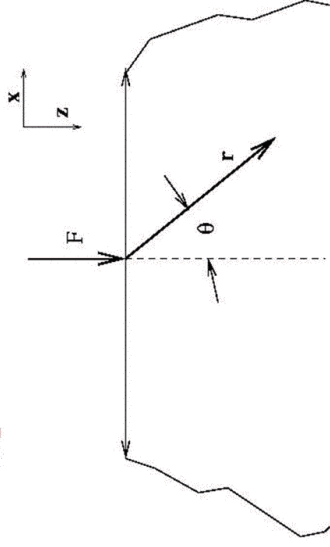
Large variance of distribution



Pentagon response



Elastic response, point force on a semi-infinite sheet

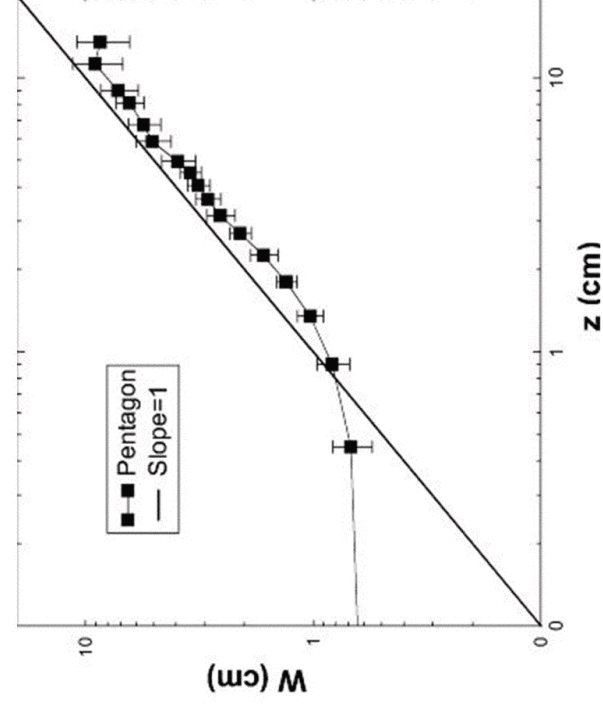


$$\sigma_{rr} = \frac{2F \cos \theta}{r\pi} \quad \sigma_{r\theta} = \sigma_{\theta\theta} = 0$$

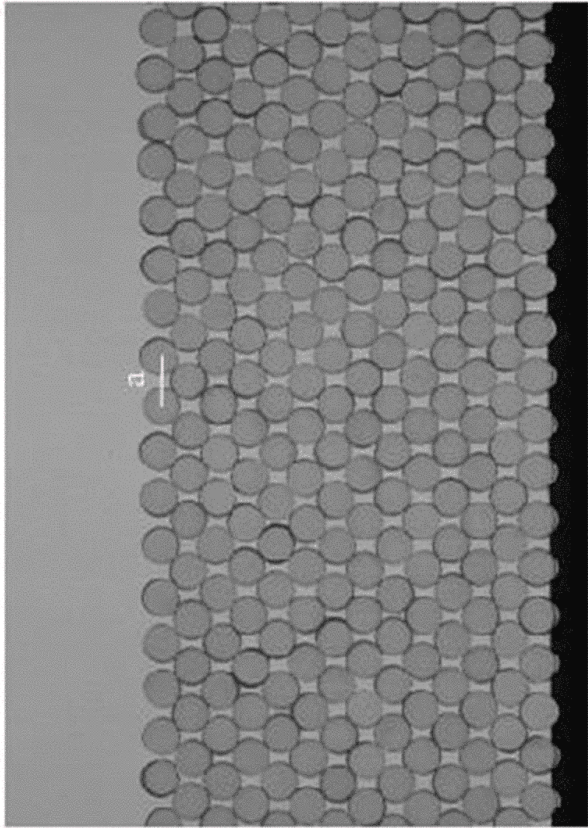
In Cartesian coordinates:

$$\sigma_{ii} = 1/[z(1 + (x/z)^2)]^p \quad p = 1, 2$$

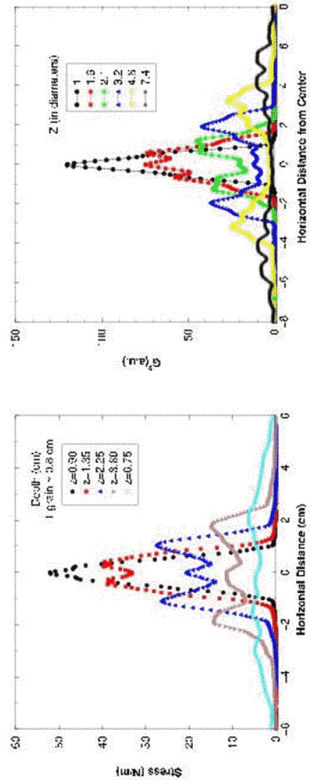
Pentagons, width vs. depth



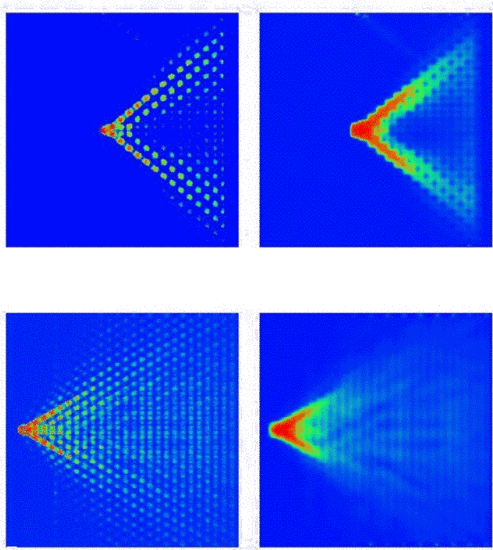
Rectangular packing reduces contact disorder



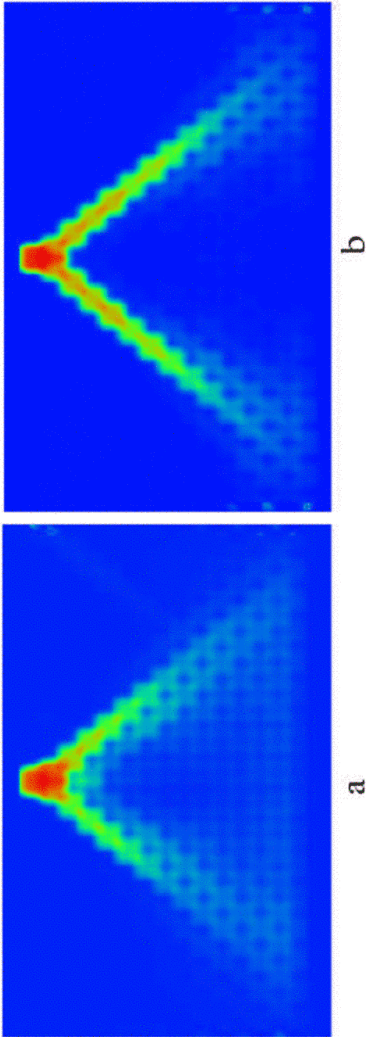
Hexagonal vs. square, data



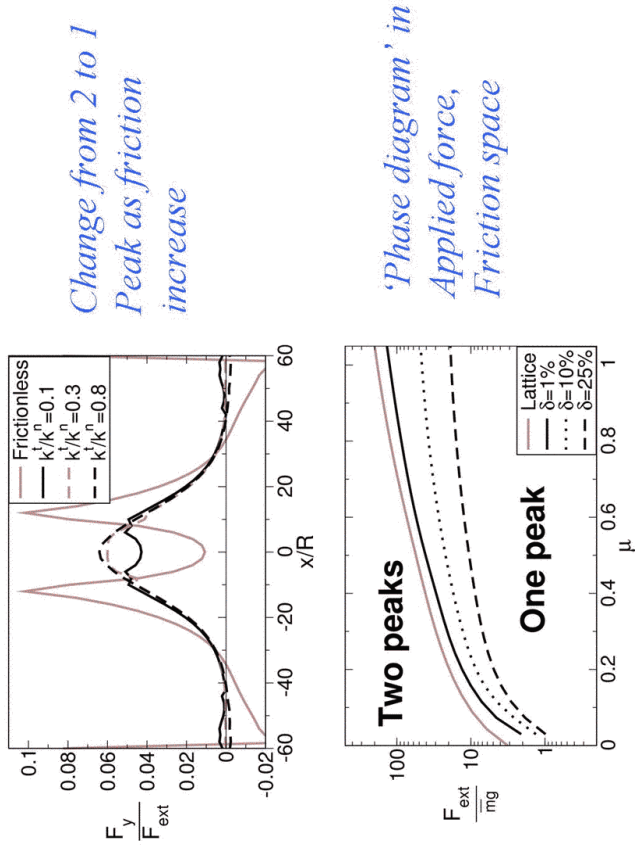
Hexagonal vs. square packing



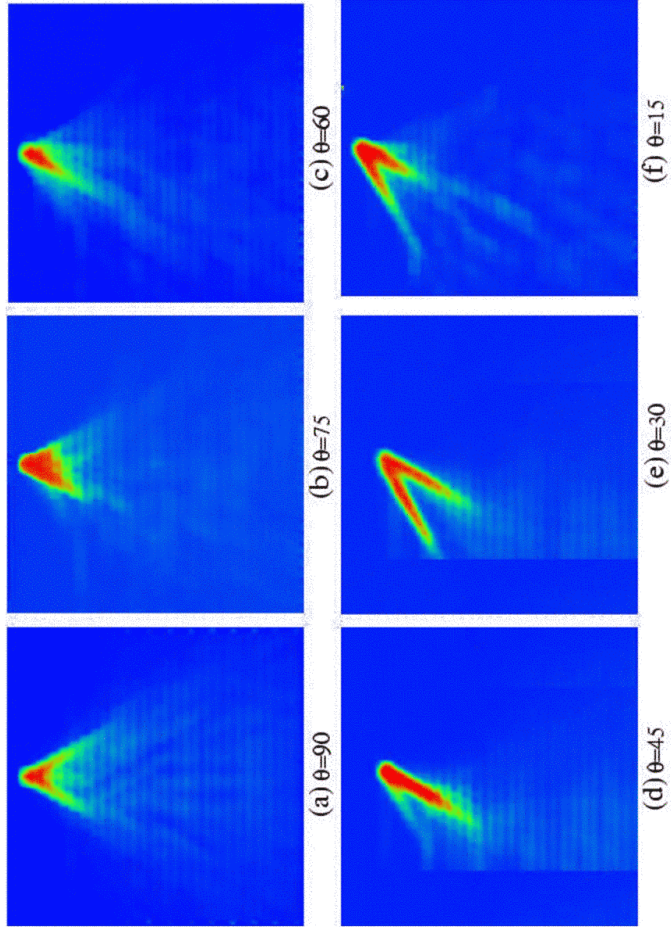
Square packs, varying friction



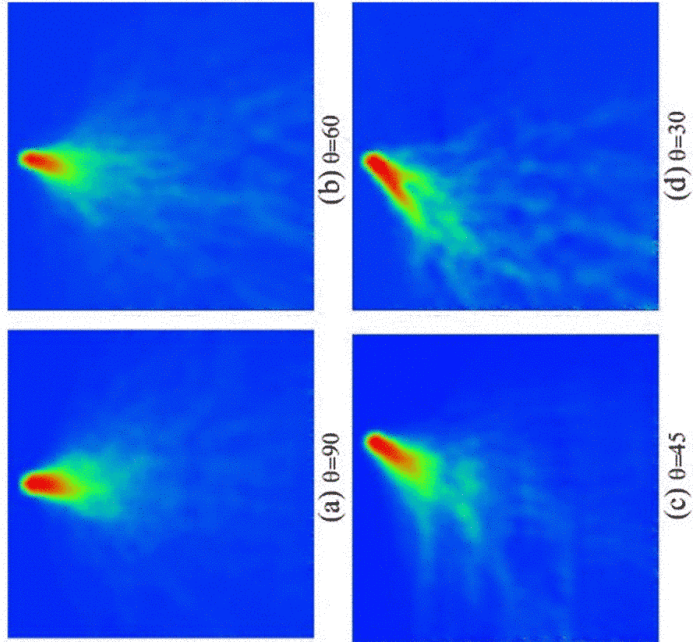
Friction ‘enhances’ elasticity—e.g. tends to produce single-peaked response (Goldenberg/Goldhirsch)



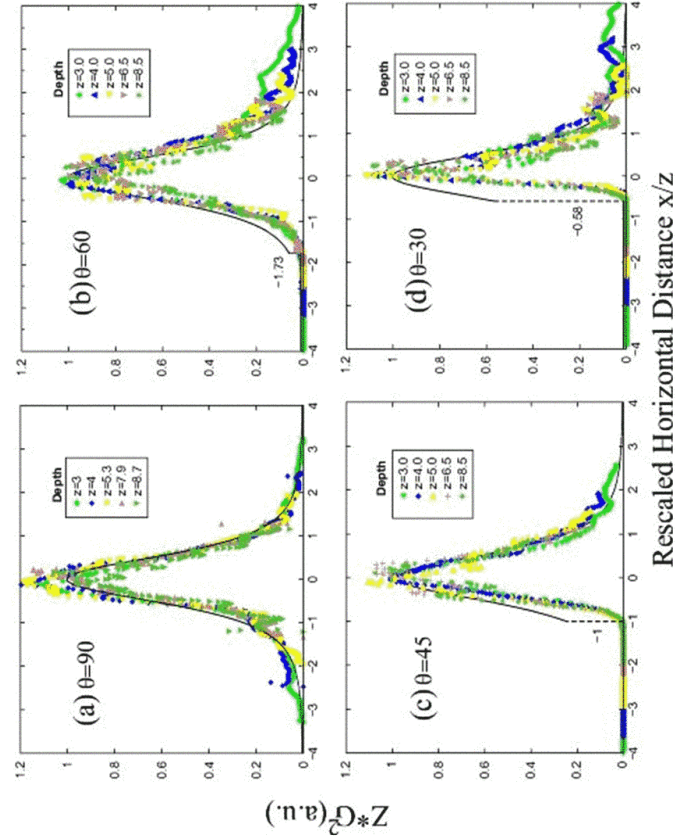
Non-normal response, disks, various angles



Non-normal responses, pentagons



Non-normal response, pentagons, rescaled



Conclusions

- Normal force distributions are sensitive to stress state
- Long-range correlations for forces in sheared systems—thus, force chains can be mesoscopic at least
- Diffusion in sheared systems: insights into microscopic statistics of driven granular materials
- Logarithmic rate dependence is seen in sheared granular systems
- Interesting connections to avalanches/earthquakes...
- Order-disorder transition—first order characterizes jamming-unjamming, contradictions notions of vibration \rightarrow temperature in granular systems
- Strong effects on transmission from order/disorder (spatial and force-contact)—overall response is mostly elastic

What are important questions? (Dense materials)

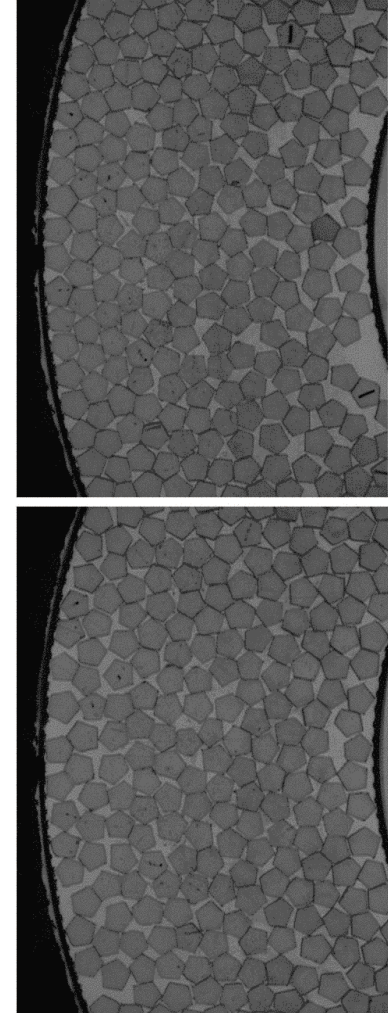
- What are statistical properties/variability of granular systems?
- What is the nature of spatio-temporal correlations/fluctuations?

The answer to this requires addressing the relevant multi-scale phenomena involved—something that is just now being considered

- Is there a universal description for stress, deformation, etc?

Extra Slides Follow

Reynolds Dilatancy



Jamming and fragility (Liu and Nagle, Bouchaud et al.)

Class of systems that are constrained or jammed

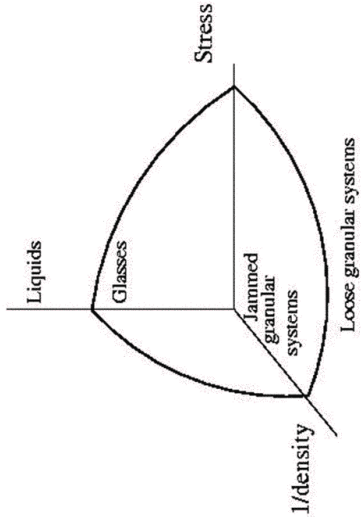
Granular Materials

Foams

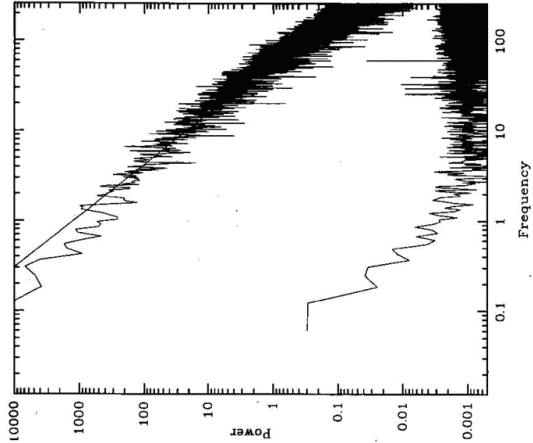
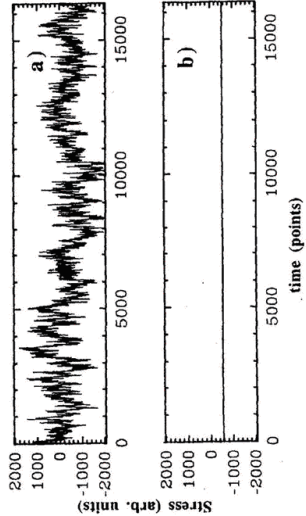
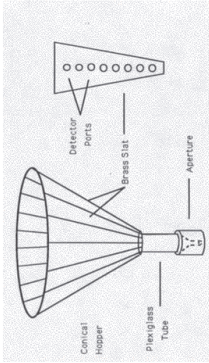
Colloids

Glasses

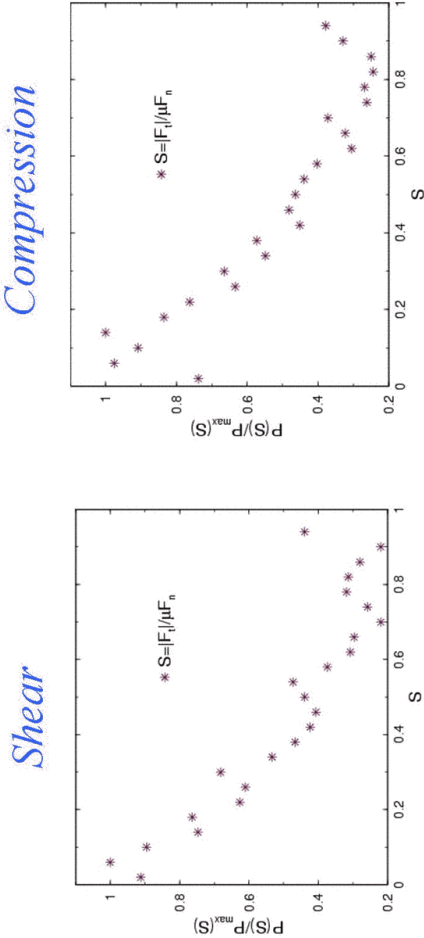
Temperature



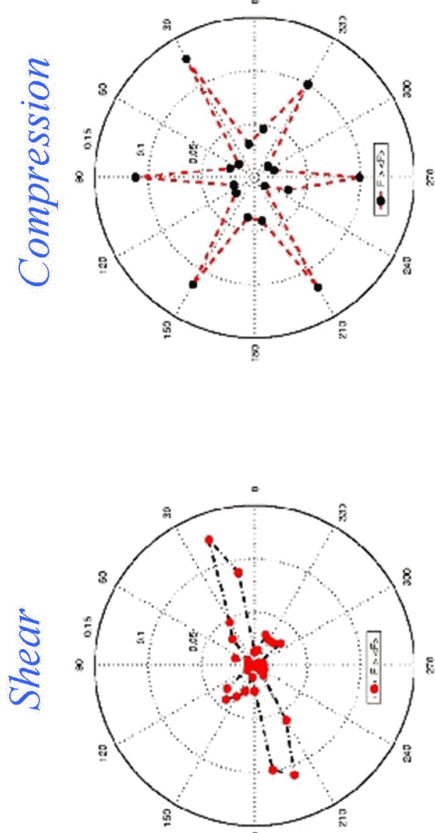
Fluctuations in a Simple Granular Flow
Baxter et al. Eur. J. Mech.B 1991



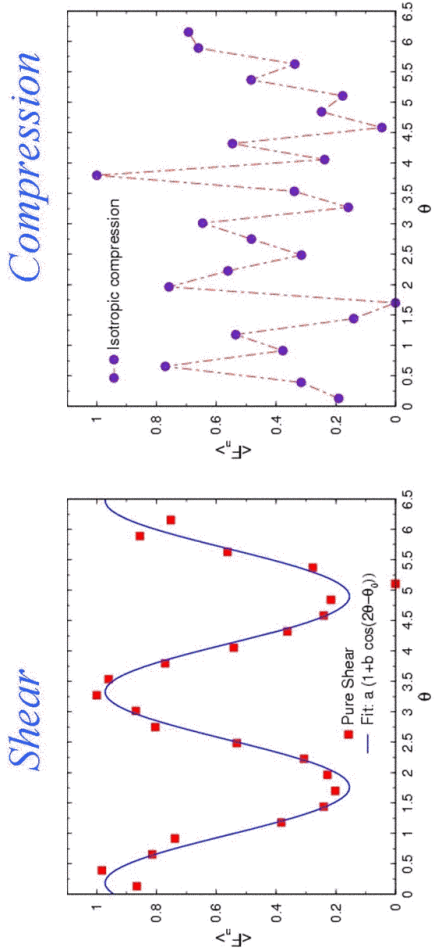
Distributions of frictional mobilization



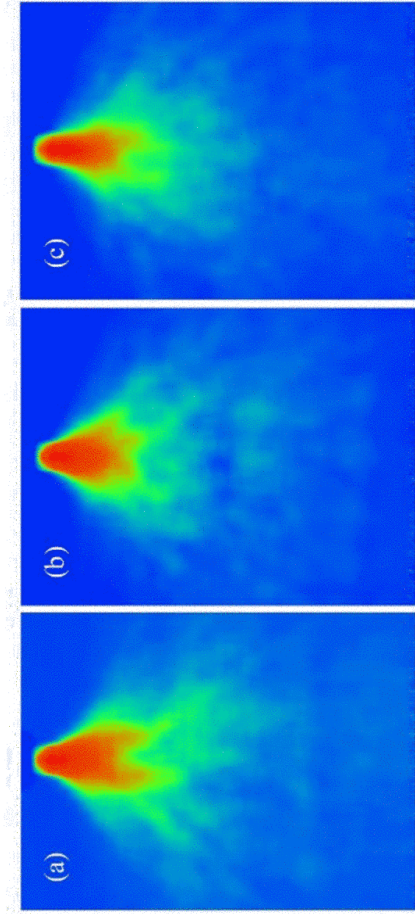
Distributions of contact angles for larger forces
(Geometric anisotropy)



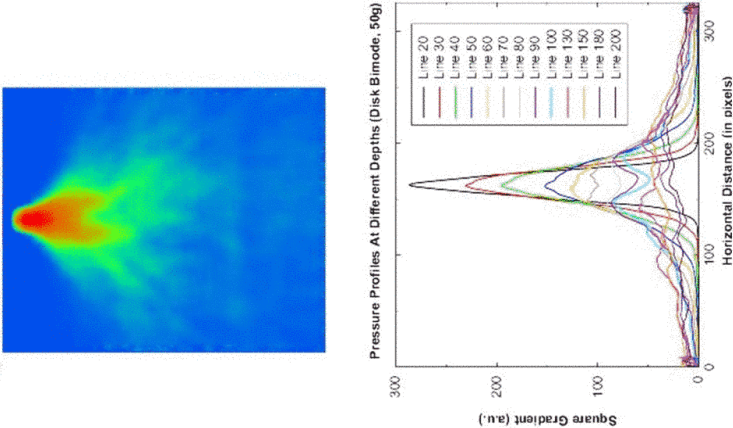
Variation with angle of mean normal force
(Force anisotropy)



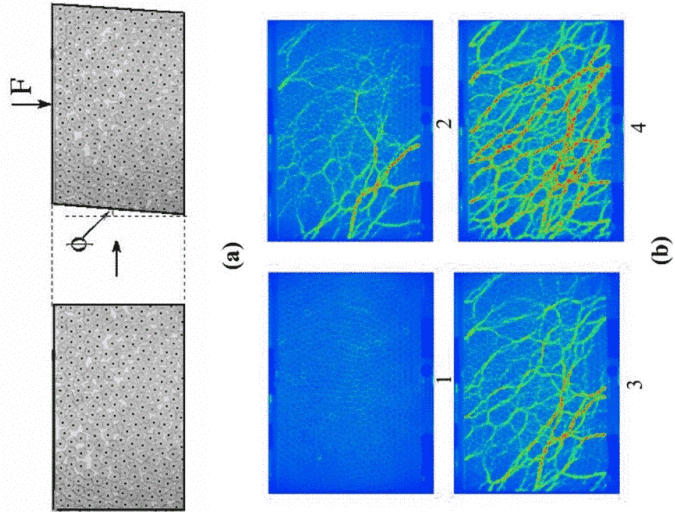
Bidisperse responses vs. A



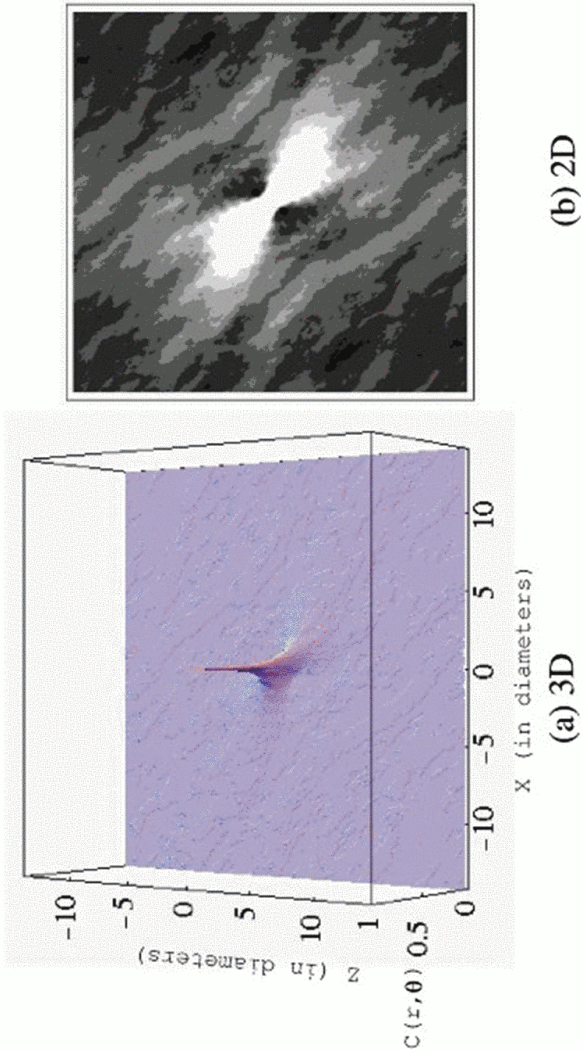
Weakly bi-disperse: two-peak structure remains



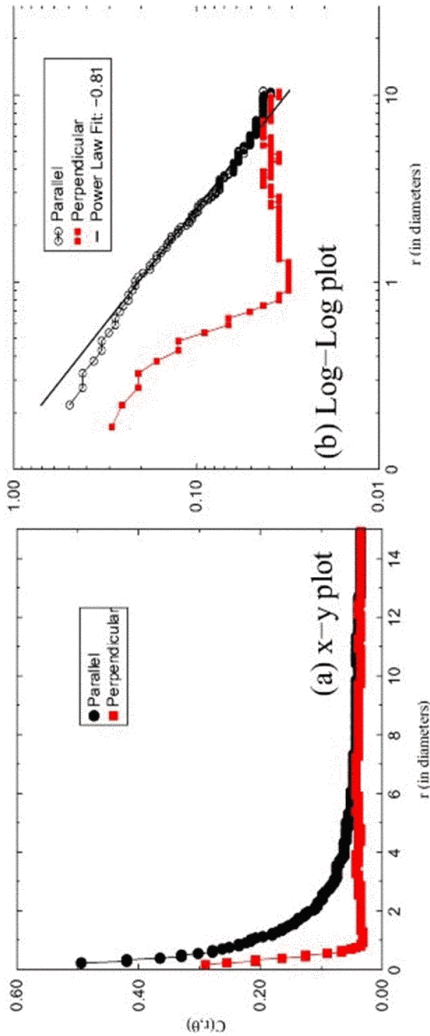
Creating a texture by shearing



Force correlation function



Correlation functions along specific directions



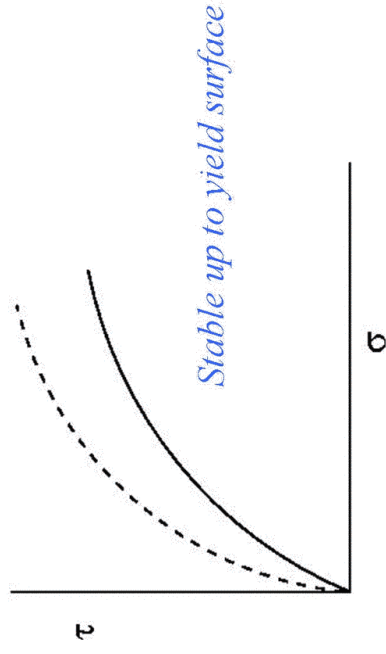
Features of elasto-plastic models

Conserve mass: $\partial \rho / \partial t + \partial_i (\rho v_i) = 0$

(Energy: lost by friction)

Conserve momentum: $\rho \frac{dv_i}{dt} = -\partial_j T_{ij}$

Concept of yield and rate-independence



$\tau \Rightarrow$ shear stress, $\sigma \Rightarrow$ normal stress

Example of stress-strain relationship for deformation

$$T_{ij} = P\delta_{ij} + kPV_{ij} / |V|$$

$$V_{ij} = -(\partial_j v_i + \partial_i v_j) / 2$$

(Strain rate tensor with minus)

$$|V|^2 = \Sigma V_{ij}^2 \quad |V| = \text{norm of } V$$

Contrast to a Newtonian fluid:

$$T_{ij} = P\delta_{ij} + 2\eta[V_{ij} - \text{Tr}(V)/3] + (2\zeta/3)\text{Tr}(V)$$

Convection-diffusion/3-leg model

Applies for weak disorder

$$O^+ O^- \sigma = 0$$

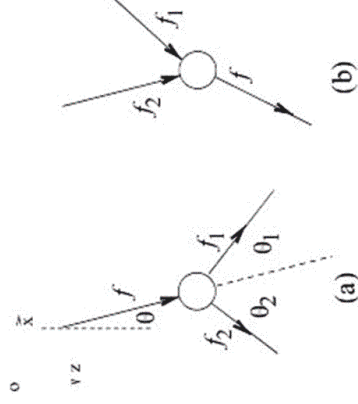
$$O^\pm = [\partial/\partial z \pm c\partial/\partial x - D\partial^2/\partial x^2]$$

Expected response to a point force:

$$\sigma_{zz} = \frac{F}{2} \frac{1}{\sqrt{4\pi Dz}} \{ \exp[-(x+cz)^2/4Dz] + \exp[-(x-cz)^2/4Dz] \}$$

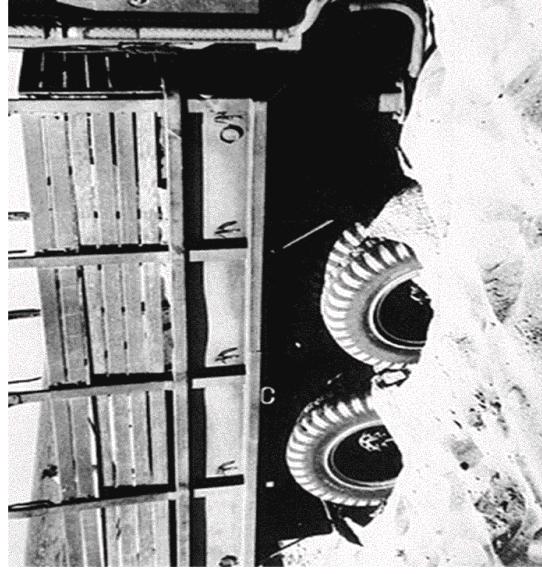
Double-Y model

Assumes Boltzmann equation for force chains



For shallow depths: One or two peaks
Intermediate depths: single peak-elastic-like
Largest depths: 2 peaks, propagative, with diffusive widening

A bit further from home



Granular Material Phases-Gases

Granular Gases:

Cool spontaneously, show clustering instability

$$T_g = (1/2)m\langle v^2 \rangle$$

Clustering in a Cooling Granular Gas (from work by S. Luding, H. Herrmann)

- Cooling

