

Inertial effects in crackling noise: magnetic materials and granular matter

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Granular media:

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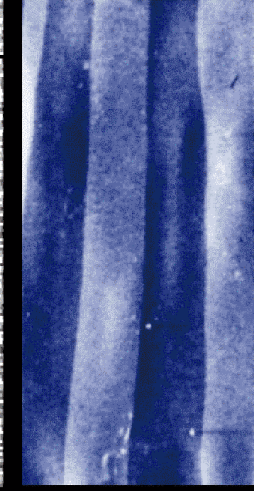
Luciano Pietronero

Stefano Zapperi

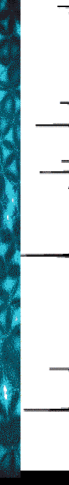
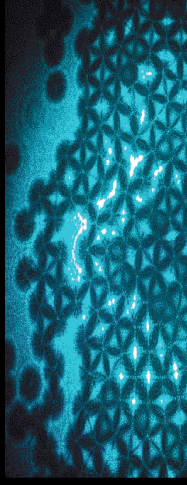
INFN, CNR

Roma

Crackling noise: from atomic to tectonic



Sethna et al Nature 410, 242 (2001)



outline & motivation

Crackling noise:

broad (power law)
amplitude distributions,
 $1/f^\alpha$ power spectra,
avalanches.

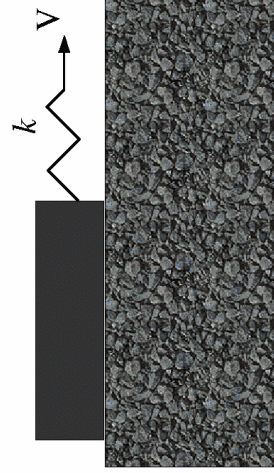
Questions:

can we explain the phenomenology?
is there universality?
can we walk across different
scales and phenomena?

Outline:

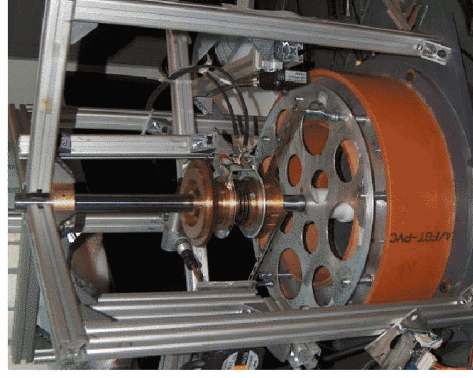
Stick-slip in granular friction
Analogy with magnetic Barkhausen noise
combined analysis of experiments and models

Granular friction



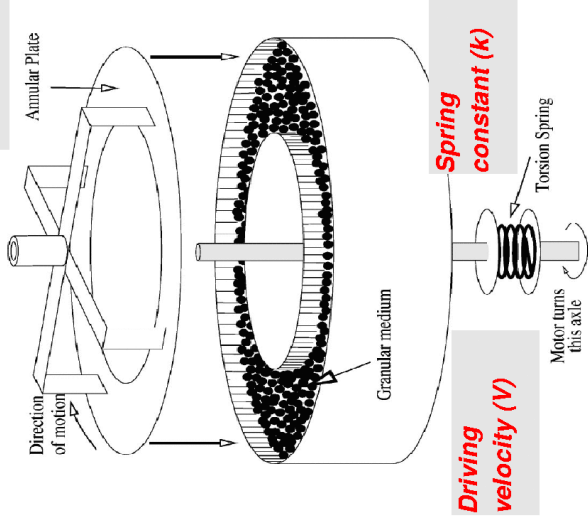
- Macroscopic contacts
- Dilatancy
- Strong internal stress fluctuations

Relevant for earth sciences
(see previous talks in this program)



Sheared granular medium

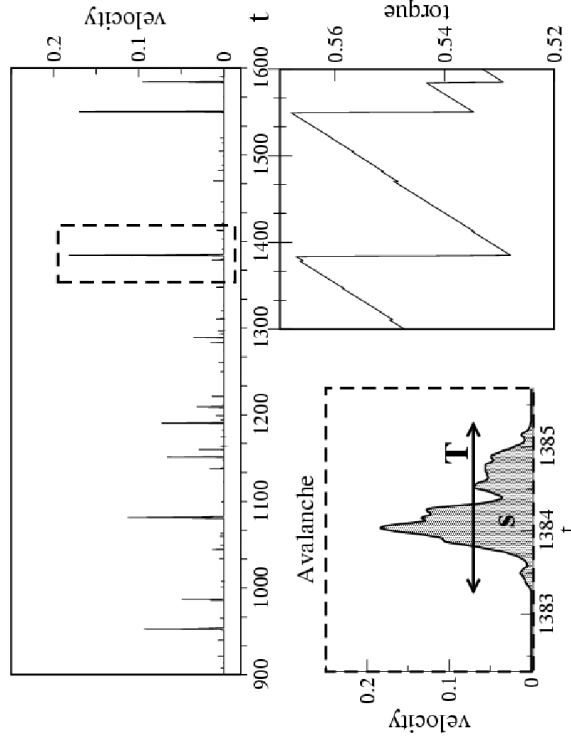
Disk Inertia (I)



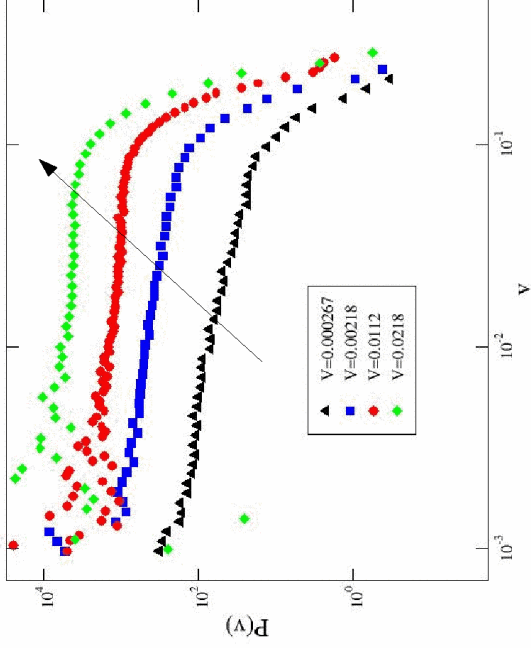
- ★ Annular shear cell, Couette geometry, driven by overhead plate via a torsion spring
- ★ 2mm glass beads
- ★ Max. humidity ~ 40%
- ★ Temperature 18-23°C
- ★ Free dilation of the bed

control parameters: I V k

Stick-slip and velocity fluctuations

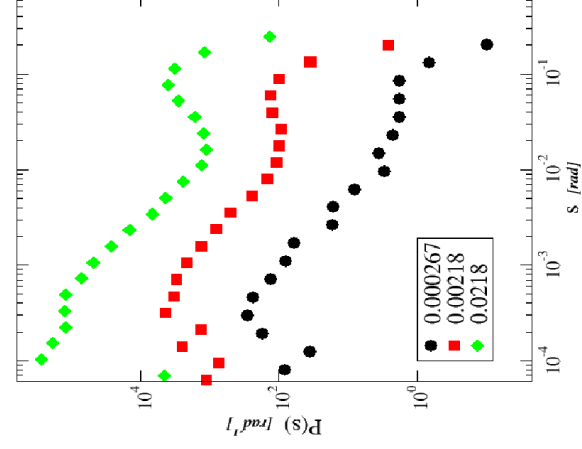
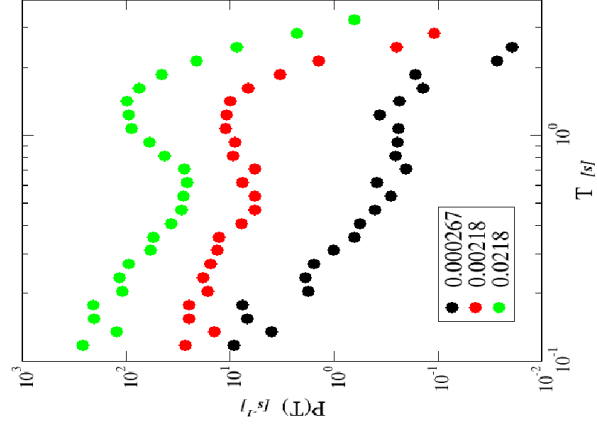


Velocity distributions



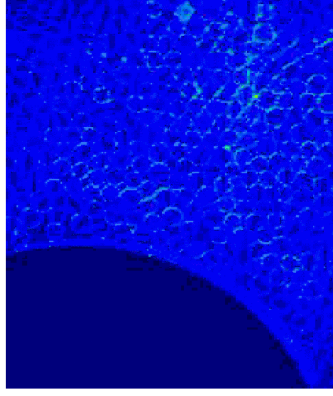
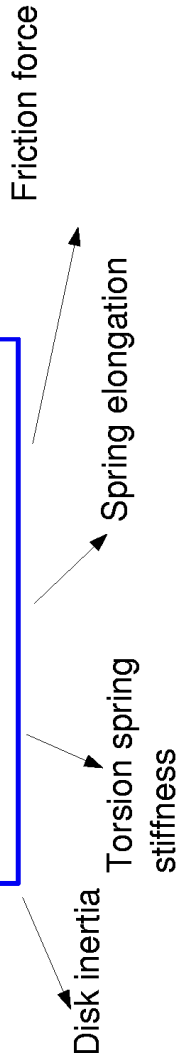
Exponent decreases as driving rate increases.

Slip duration and size distributions



Equation of motion

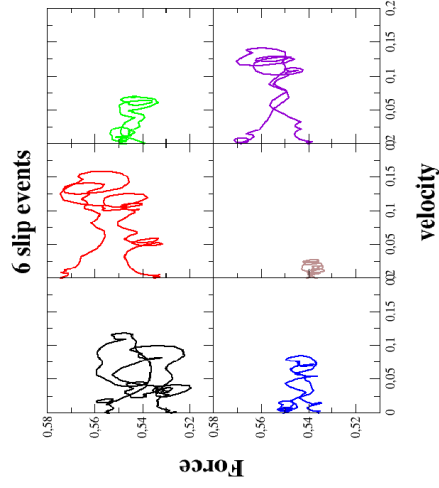
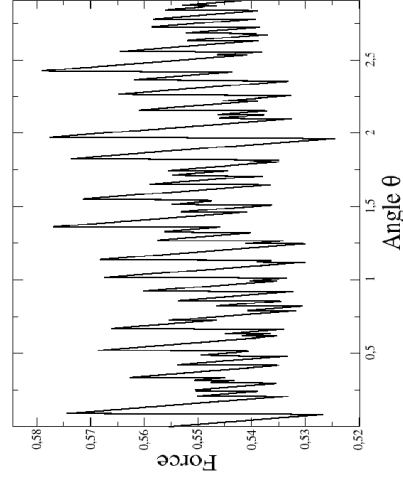
$$I\ddot{\theta} = k(Vt - \theta) - F(\theta, \dot{\theta}, \dots)$$



<http://www.phy.duke.edu/~dhowell/research.html>

Question: how to model the friction force?

Friction forces

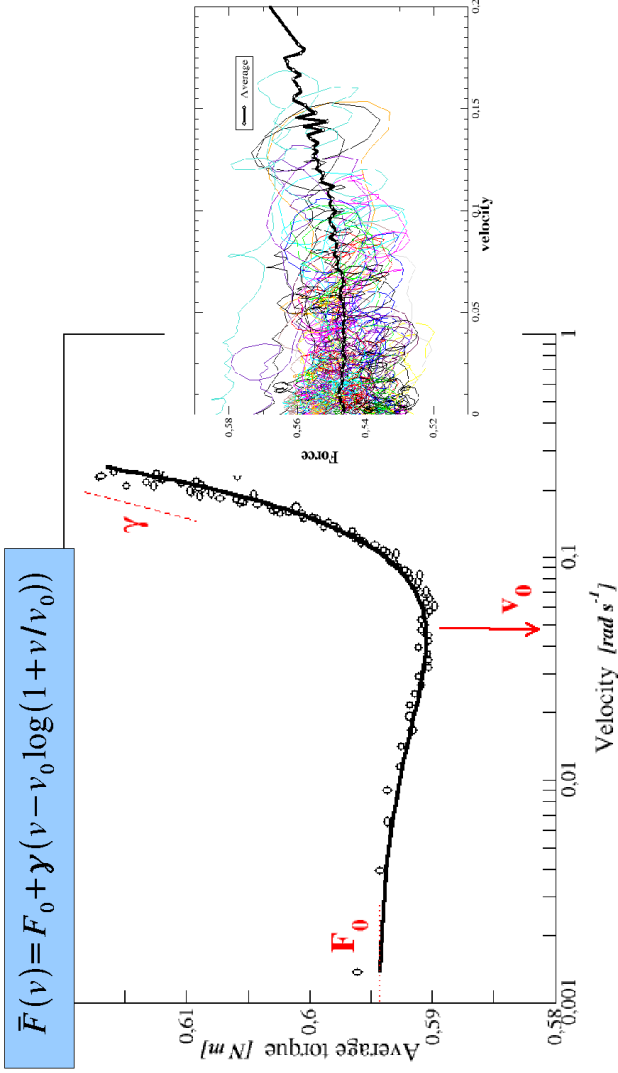


Hypothesis:

- Average viscous force
- Slip dependent random force

$$F(\theta, \dot{\theta}) = \bar{F}(\dot{\theta}) + F_f(\theta)$$

Viscous force



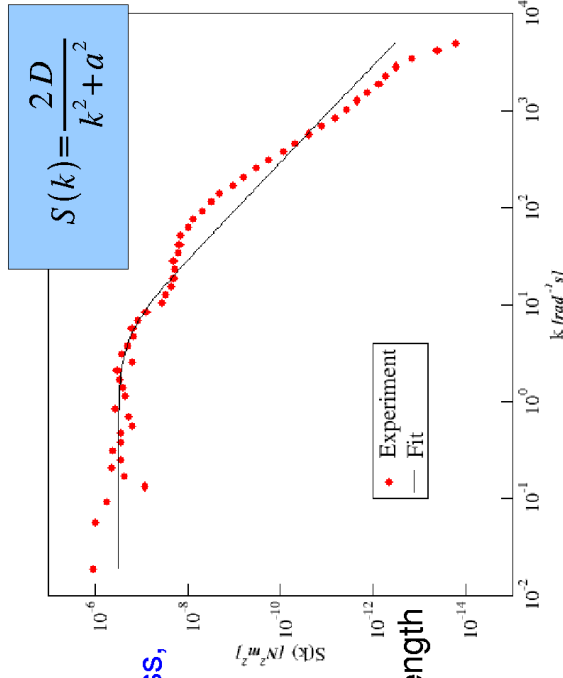
Random forces

The random frictional term is modeled by a confined Brownian process,

$$\frac{dF_f}{d\theta} = \eta(\theta) - aF_f$$

a inverse correlation length

$D = \langle \eta^2 \rangle$ noise variance



The model

$$I\ddot{\theta} = k(Vt - \theta) - F_0 + \gamma(\dot{\theta} - v_0 \log(1 + \dot{\theta}/v_0)) + F_f(\theta)$$

$$\frac{dF_f}{d\theta} = \eta(\theta) - aF_f \quad \langle \eta(\theta)\eta(\theta') \rangle = D\delta(\theta - \theta')$$

No disorder limit:

$$D = 0 \quad a = 0$$

$V < v_0$: periodic stick-slip

$$T^* \sim (I/k)^{1/2}$$

$V > v_0$: steady sliding

Overdamped limit:

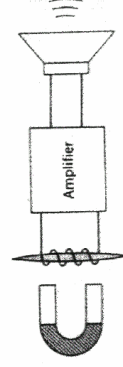
$$\gamma V \gg I \quad V \gg v_0$$

disorder dominated dynamics

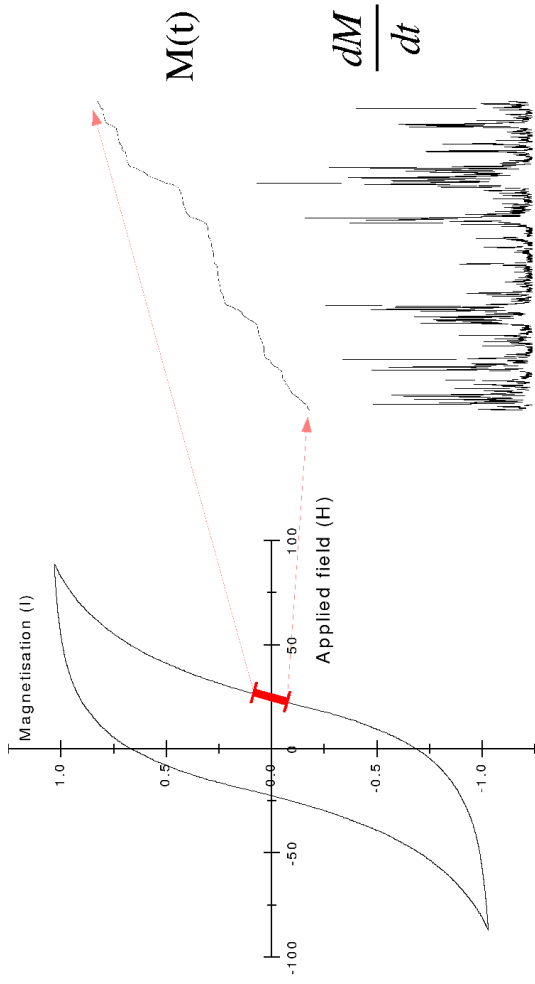
power law scaling

Domain wall model!

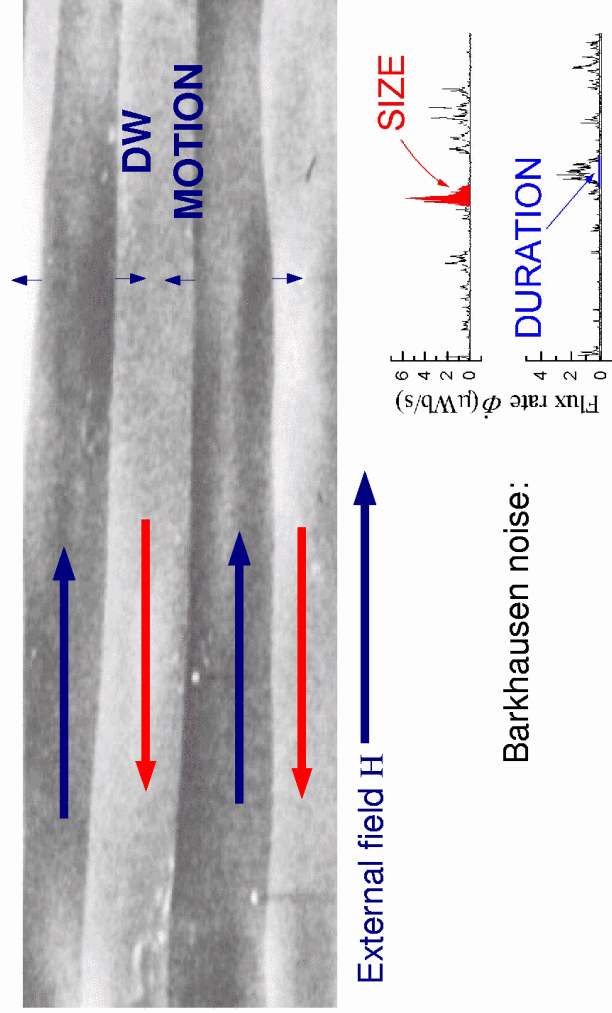
Intermezzo the Barkhausen effect



The Barkhausen effect



Domain wall motion



Domain wall model (ABBM)

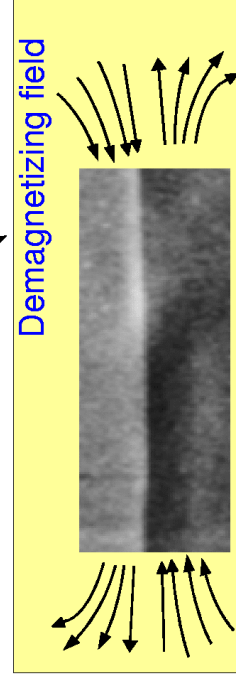
B Alessandro, C Beatrice, G Bertotti, A Montorsi
Journal of Applied Physics 68, 2901 (1990).

$$\Gamma \frac{dx}{dt} = ct - kx - H_{pin}(x)$$

Damping
eddy currents

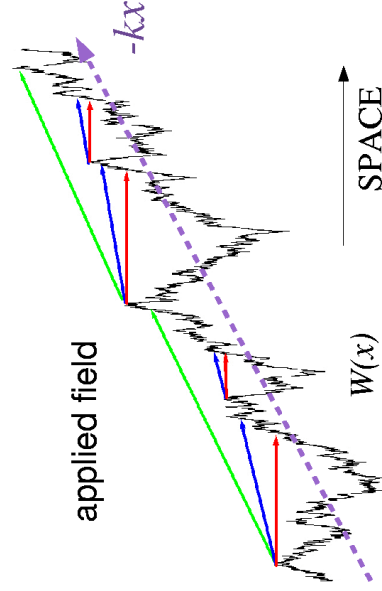
Applied field
constant rate

Pinning field
Brownian correlations



$$\frac{dH_{pin}}{dx} = \eta(x)$$

Avalanches in the ABBM model



Exact results:
velocities

$$P(v) = v^{c-1} \exp(-kv) / N$$

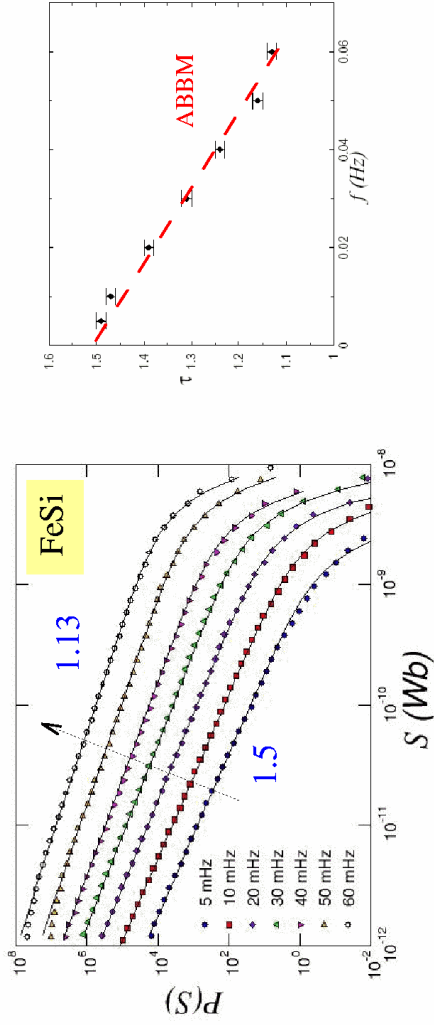
sizes

$$P(s) = s^{-3/2+c/2} f(sk^2)$$

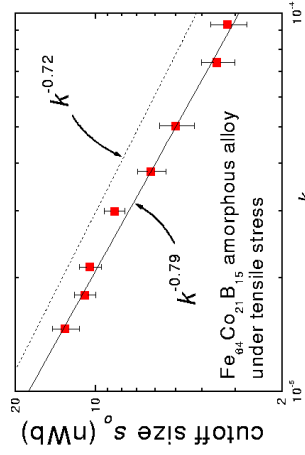
durations

$$P(T) = T^{-2+c} g(Tk)$$

Field rate dependence

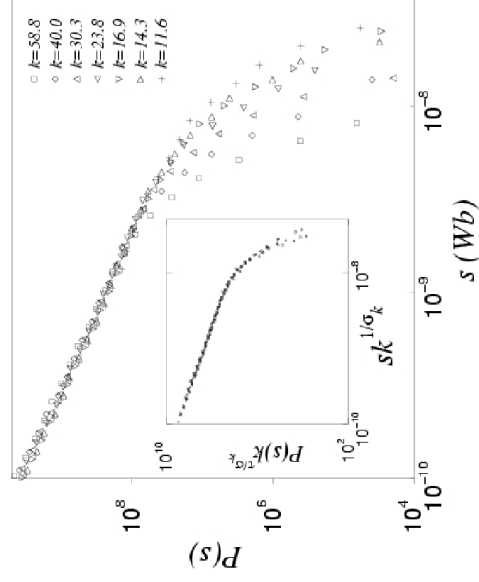


Effect of the demagnetizing field

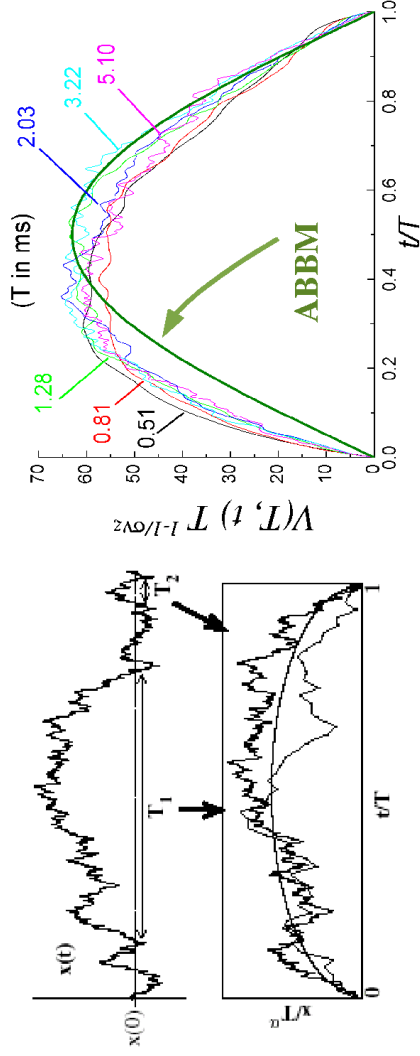


cutoff to the power law increases as k decreases but **not** as $1/k^2$

The exponent can be understood considering DW fluctuations Durin & Zapperi Phys. Rev. Lett. 84, 4705 (2000).



Average pulse shape



Theoretical shape is symmetric also true for RFIM, DW depinning...

Experiments are asymmetric!

Eddy current damping

Damping comes from eddy currents:

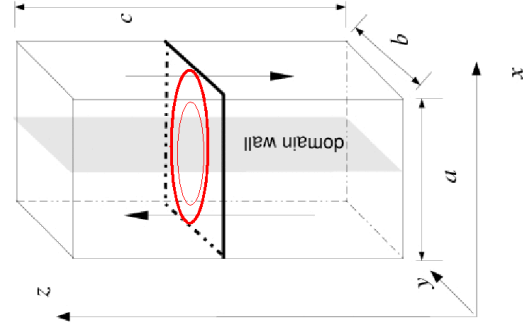
$$\partial_t H_e - \sigma \mu \nabla^2 H_e = 0$$

$$\partial_x H_e^x - \partial_x H_e^y = 2\sigma I_s v(t)$$

In the **quasistatic** approximation we get:

$$I_s v dx = H_{eff} dx$$

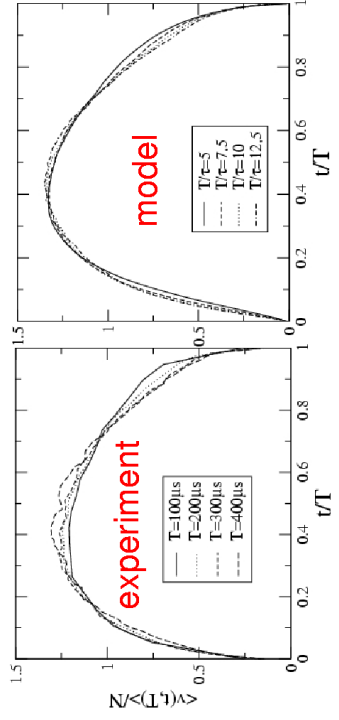
Energy dissipated by eddy currents Work done by the field



Beyond the quasi-static approximation: non-local damping

$$\Gamma v(t) \rightarrow \int_0^t dt' f(t-t')v(t')$$

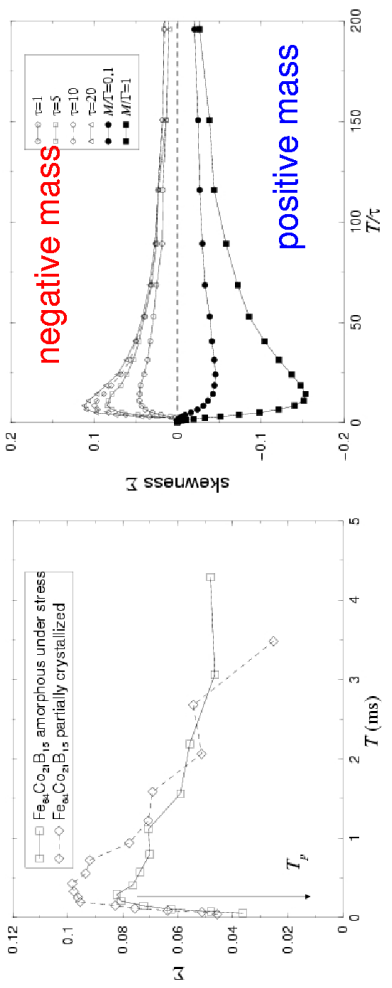
$$f(t) = -\frac{32 I_s^2}{\mu \sigma \pi} \theta_2 [e^{-4\pi^2 t(\alpha^2 \mu \sigma)}] \sum_n \frac{e^{-t/\tau_n}}{(2n+1)^2}$$



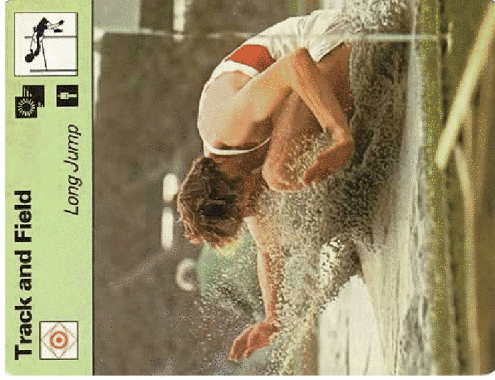
Negative domain wall mass

In Fourier space: $f(\omega)v(\omega) = (\Gamma + i\omega M(\omega))v(\omega)$
 the effective mass M is negative at all frequencies:

$$M(\omega=0) = -8 I_s^2 b^3 \mu \sigma / \pi^5 \approx 7 \cdot 10^{-5} \text{ Kg/m}^2$$



back to..... Granular friction



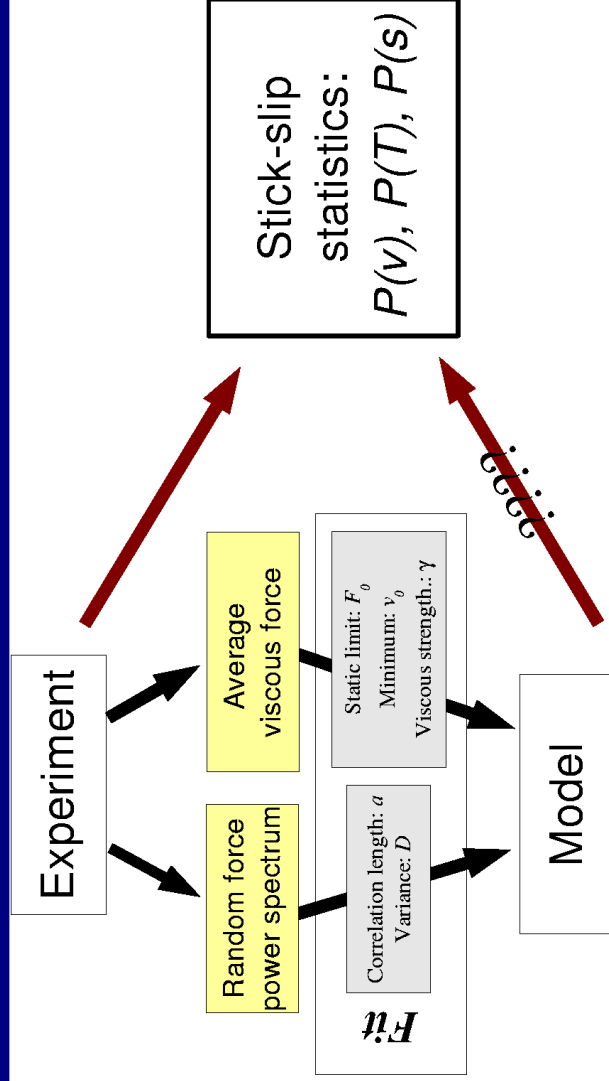
$$I\ddot{\theta} = k(Vt - \theta) - F(\theta, \dot{\theta})$$

$$F = F_0 + \gamma(\dot{\theta} - v_0 \log(1 + \dot{\theta}/v_0)) + F_f(\theta)$$

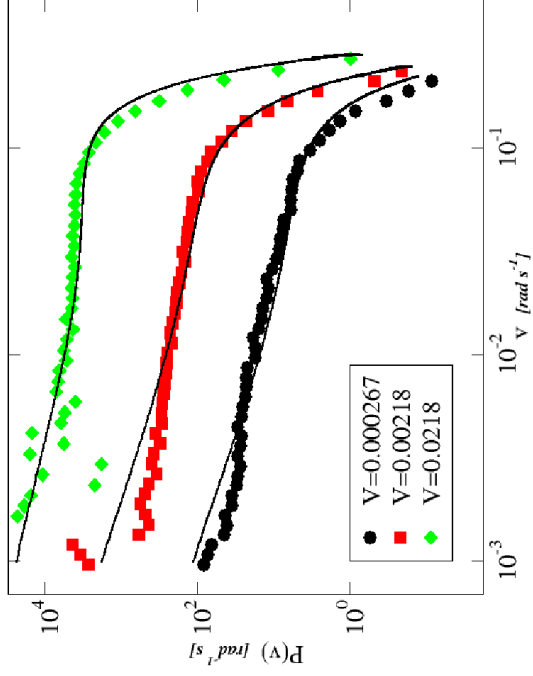
$$\frac{dF_f}{d\theta} = \eta(\theta) - aF_f$$

$$\langle \eta(\theta)\eta(\theta') \rangle = D\delta(\theta - \theta')$$

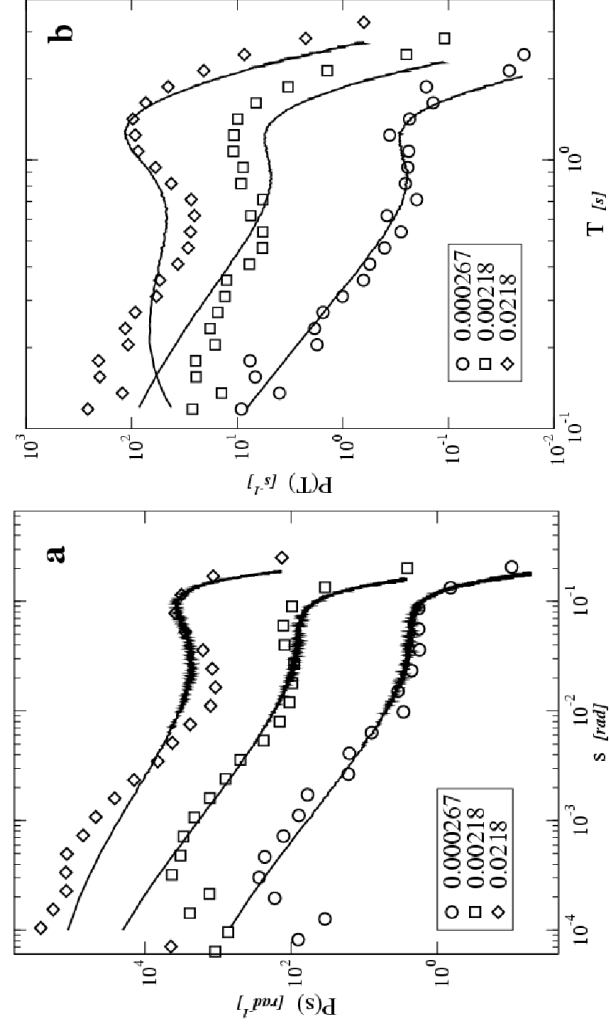
Comparison Experiments & Model



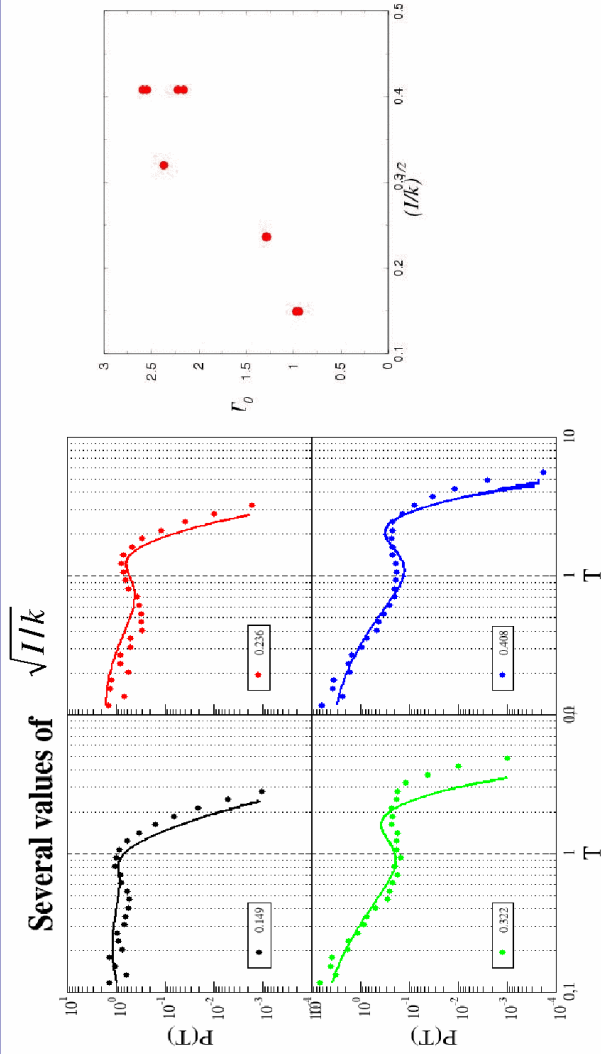
Simulations: velocity distributions



Simulation results: avalanches

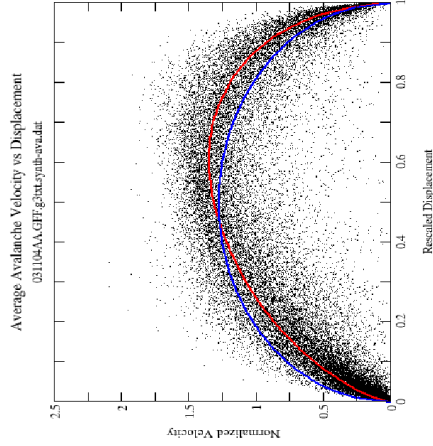


Effect of spring and inertia

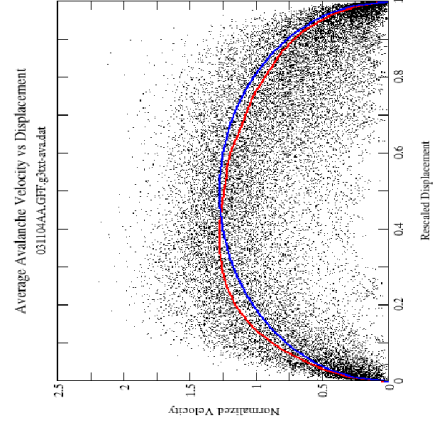


Granular pulse shapes (...in progress...)

Model:
positive mass

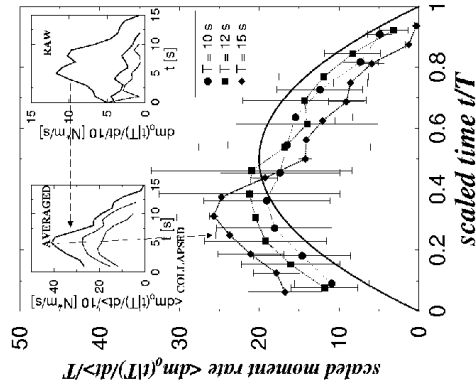


Experiments:
negative mass?

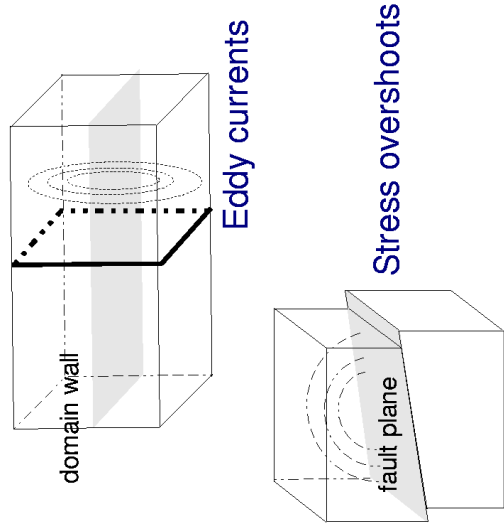


....and Earthquakes?

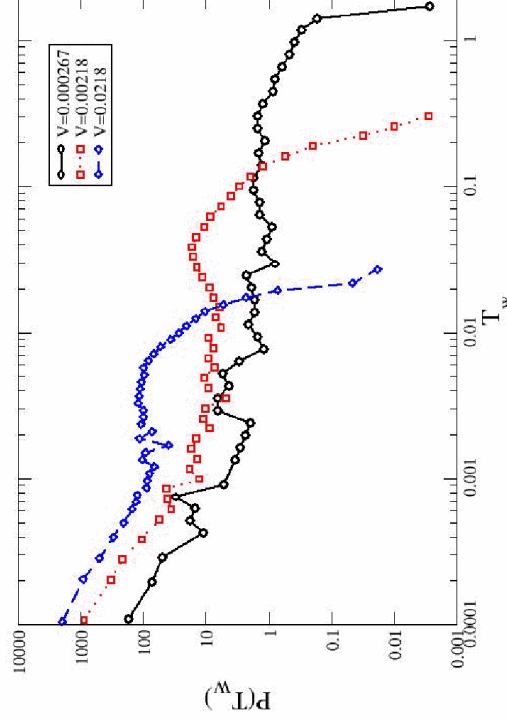
asymmetric shape maybe due to similar mechanism?



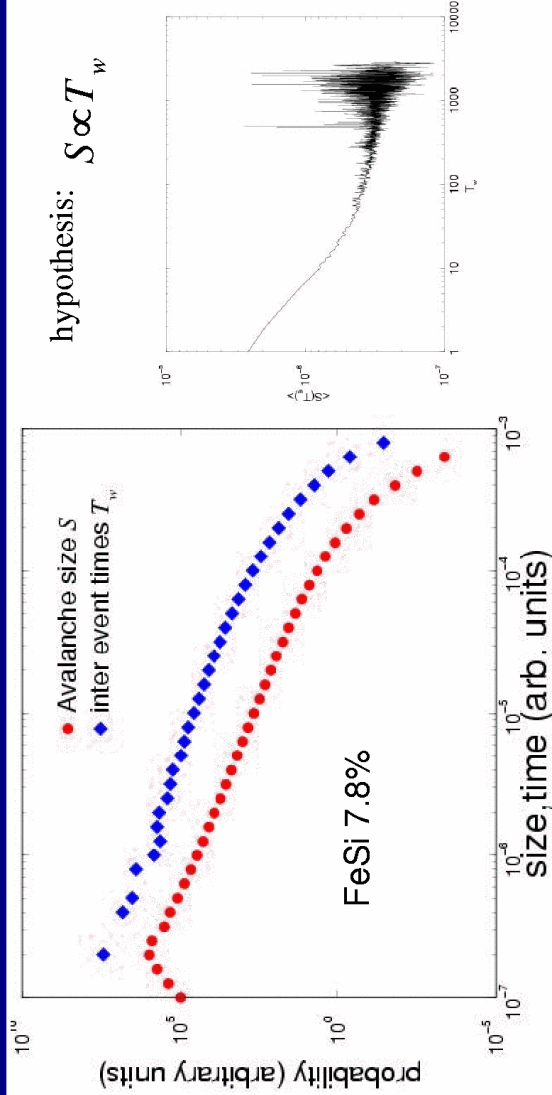
cond-mat/0509226



Inter-event time distribution for granular friction



Inter-event time distribution for Barkhausen noise



Domain walls and Granular media

$$I \ddot{\theta} = K(Vt - \theta) - F_0 - \gamma(\dot{\theta} - v_0 \log(1 + \dot{\theta}/v_0)) - \delta F(\theta)$$

slip	$\theta \rightarrow x$	domain wall jump
shear rate	$KV \rightarrow c$	magnetic field rate
spring constant	$K \rightarrow k$	demagnetizing factor
frictional damping	$\gamma \rightarrow \Gamma$	eddy current dissipation
random forces	$\delta F \rightarrow W$	pinning forces
inertia	$I \rightarrow -M$	eddy current mass

$$M \ddot{x} = ct - kx - \Gamma \dot{x} - W(x)$$