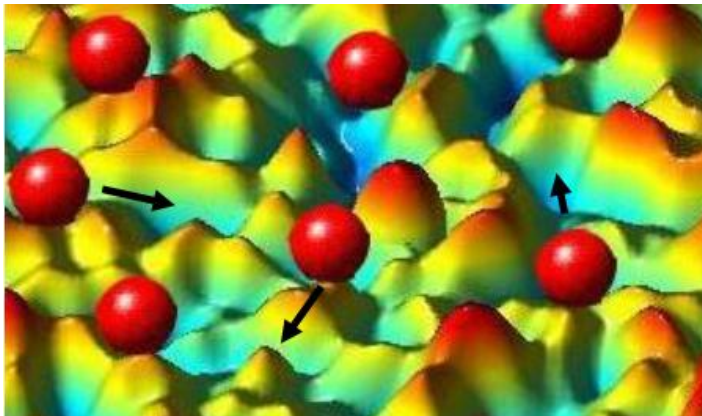


Localization,
anomalous diffusion and
slow relaxations
in disordered systems

Ariel Amir , work with
Yuval Oreg and Yoseph Imry



Outline

- Electron glass model

- slow relaxations and aging
- statics (Coulomb gap)

Dynamics:

- mapping to a new class of Random Matrix Theory (RMT)

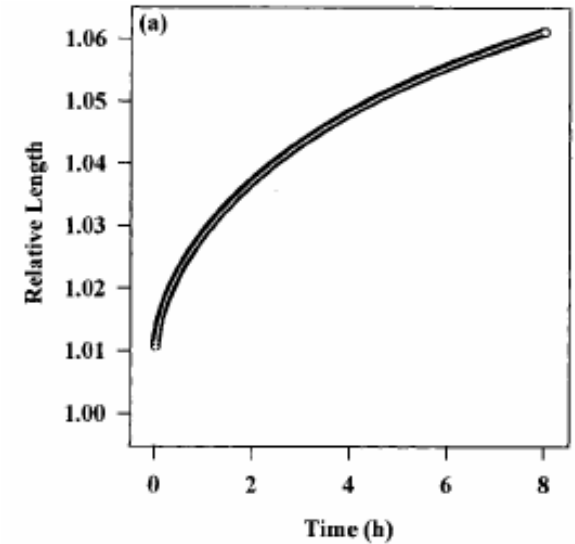
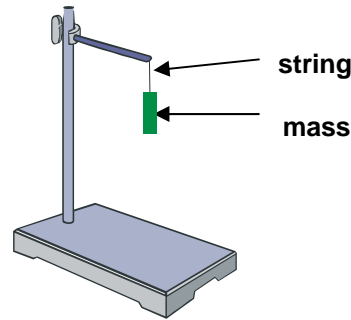
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- Solution of model: *eigenvalue distribution* through **moment calculation**
- Solution of model : **RG approach** → *localization properties*
- Implications: **anomalous diffusion**, **slow relaxations**, **localized phonons**
- Conclusions

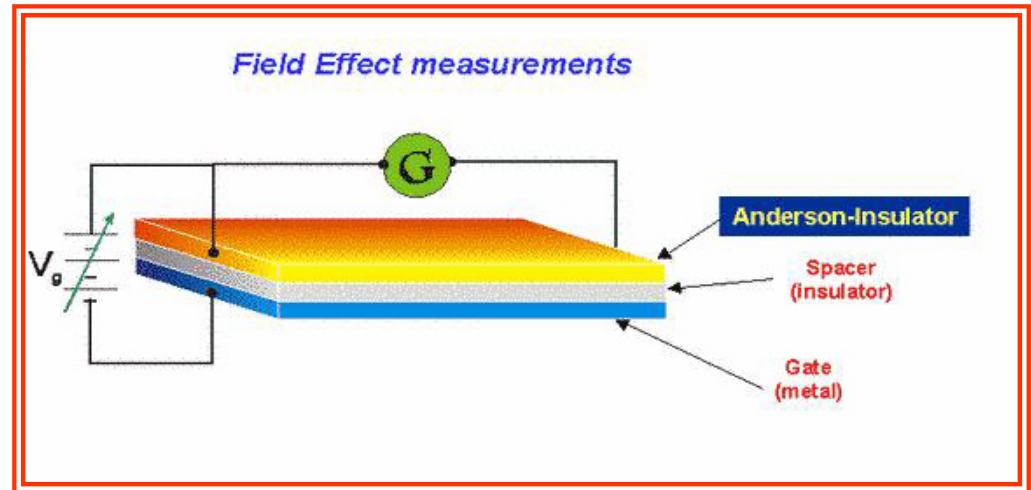
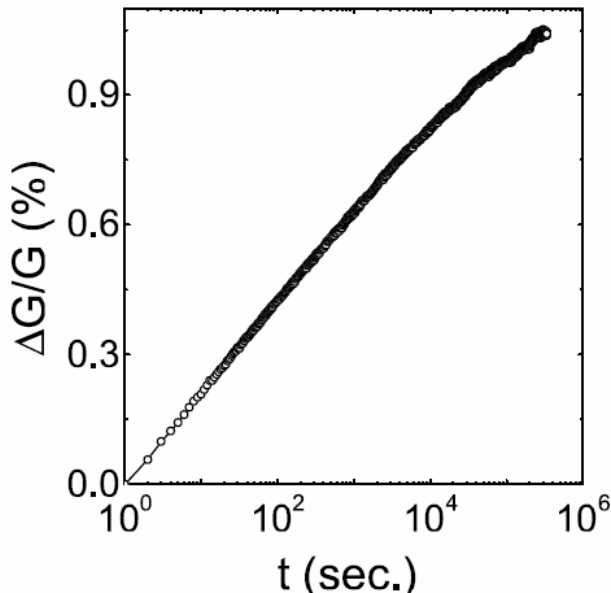
Slow relaxations in nature

W. Weber, *Ann. Phys.* (1835)

D. S. Thompson, *J. Exp. Bot.* (2001)



Electron glass- Experimental system



What are the ingredients leading to slow relaxations?

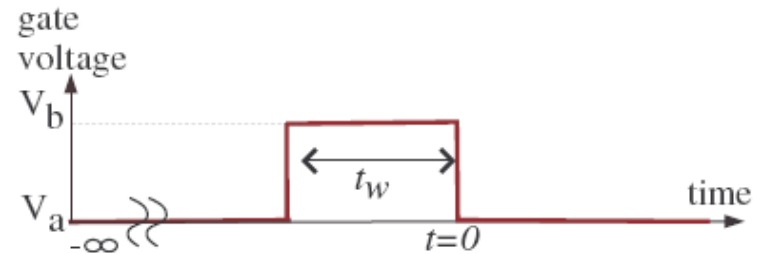
Logarithmic relaxations for 5 days!

Electron glass aging– experimental protocol

A. Vaknin and Z. Ovadyahu and M. Pollak, PRL 2000

Step I

System equilibrates for long time

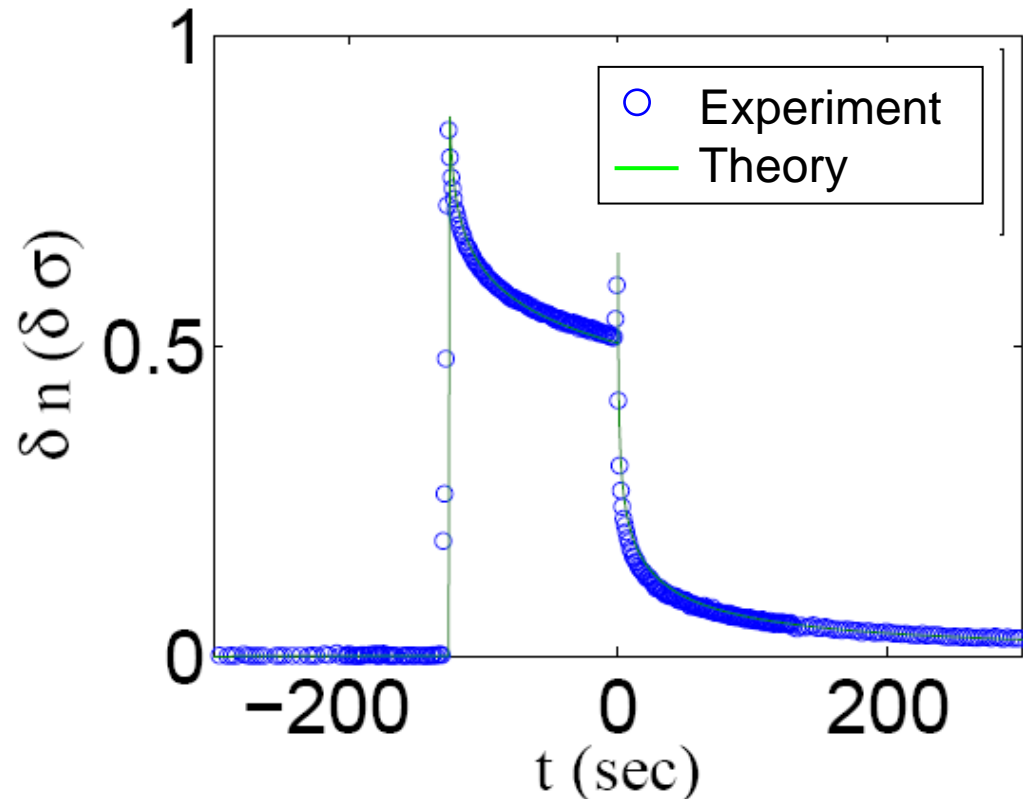


Step II

V_g is changed, for a time of t_w .

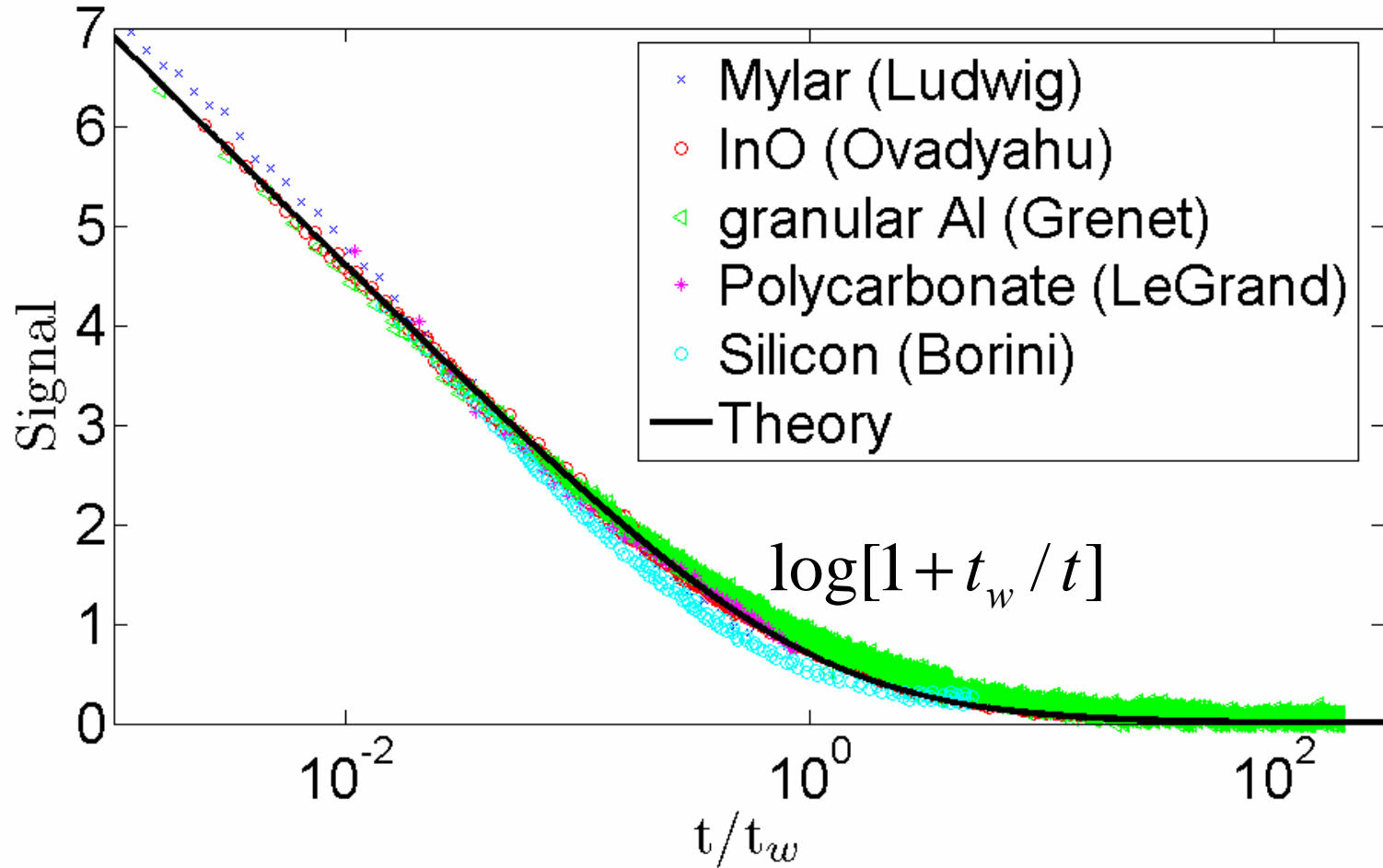
Throughout the experiment

Conductance is measured as a function of time.



Data: Ovadyahu et al.

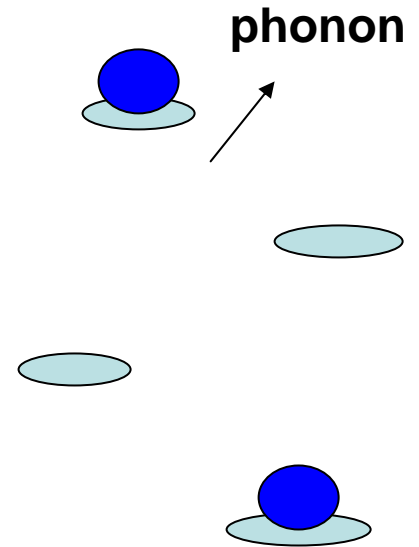
Aging and universality



Amir, Oreg and Imry, to be published

The model

- Strong localization due to disorder
→ randomly positioned sites, on-site disorder.
- Coulomb interactions are included
- “Phonons” induce transitions between configurations.
- Interference (quantum) effects neglected.



e.g:

Pollak (1970)

Shklovskii and Efros (1975)

Ovadyahu and Pollak (2003)

Muller and Ioffe (2004)

“Local mean-field” approximation - Dynamics

AA, Oreg and Imry, PRB (2008)

$$n_i \rightarrow \langle n_i \rangle, \quad \frac{dn_i}{dt} = \sum_j -\gamma_{i,j} + \gamma_{j,i}$$

$$\gamma_{i,j} = \exp(-2r_{ij} / \xi) n_i (1 - n_j) [N(|\Delta E|) + \theta(\Delta E)]$$

- ΔE includes the interactions
- N is the Bose-Einstein distribution
- ξ - the localization length

$$\left(E_i = \varepsilon_i + \sum_j \frac{n_j}{r_{ij}} \right)$$

At long times:

- The system reaches a locally stable point (metastable state).
- Many metastable states

(“Pseudo-ground-states”, *Baranovski et al., J. Phys. C, 1979*)

“Local mean-field” approximation - Equilibrium

- Detailed balance leads to Fermi-Dirac statistics ($n_j \rightarrow f_j$).
- Self-consistent set of equations for the energies:

$$E_i = \varepsilon_i - \sum_j \frac{1}{2} \frac{e^2}{r_{ij}} \tanh\left(\frac{E_j}{2T}\right) \quad (\text{assuming half filling}).$$

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Yields linear **Coulomb gap**

(+ temperature dependence)

M. Pollak, Discuss. Faraday Soc (1970)

*A.L. Efros and B.I. Shklovskii,
J. Phys. C (1975)*

A. L. Efros, J. Phys. C (1976).

M. Grunewald et al., J. Phys. C. (1982)

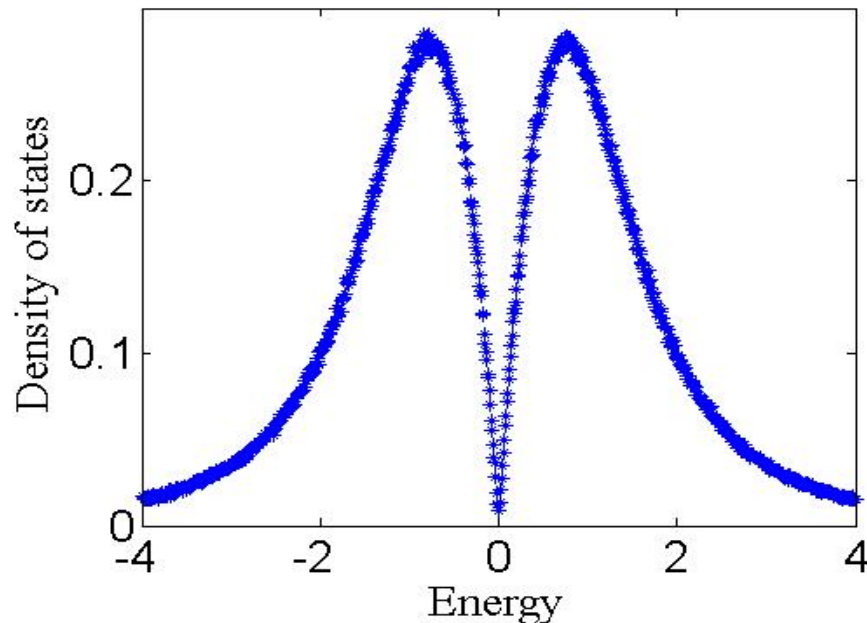
A.A. Mogilyanskii and M.E. Raikh (1989)

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Surer et al., PRL (2009)

Goethe et al., PRL (2009)



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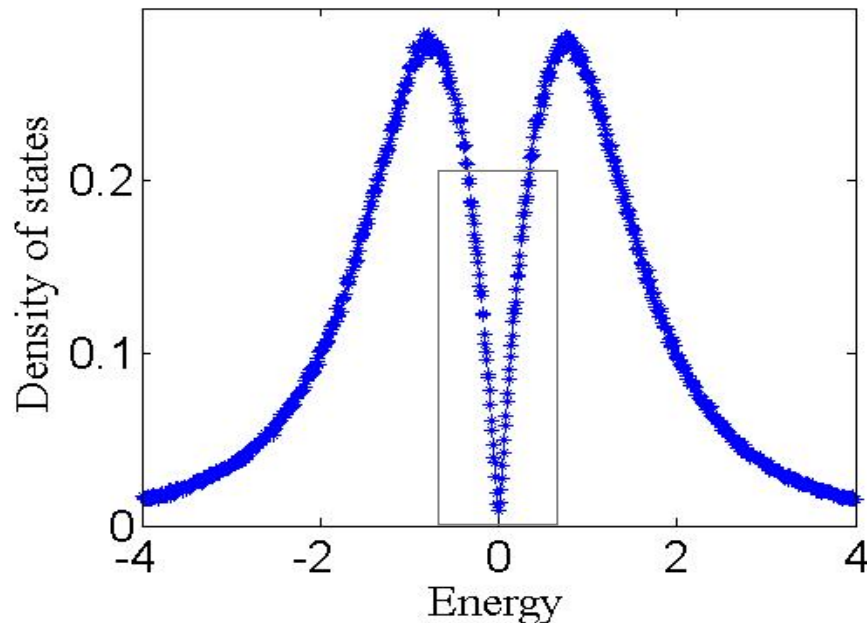
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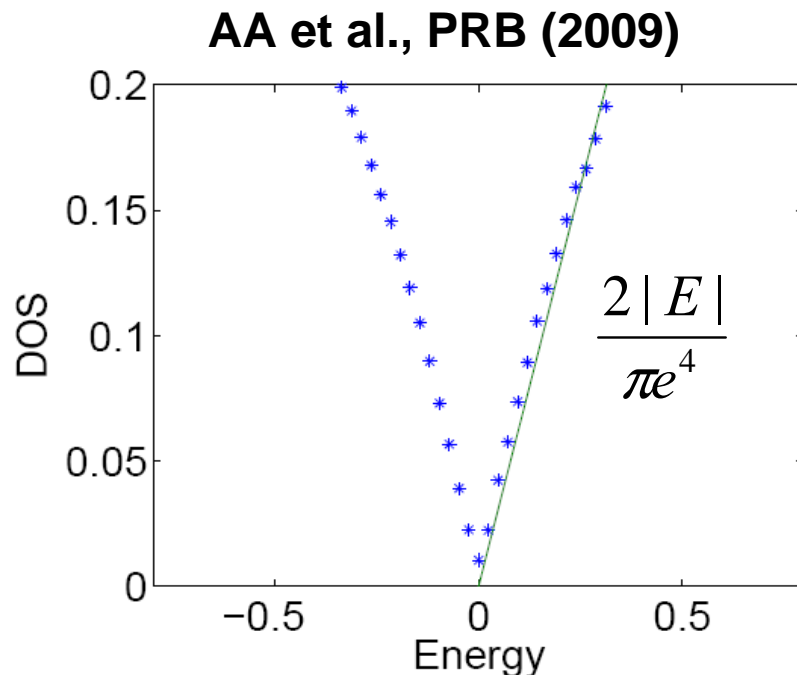
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$$\gamma_{i,j} = \exp(-2r_{ij} / \xi) n_i (1 - n_j) [N(|\Delta E|) + \theta(\Delta E)]$$

We saw: approach works well for statics

Moving on to dynamics...

Solution near locally stable point

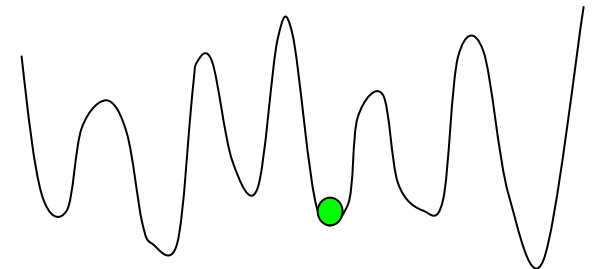
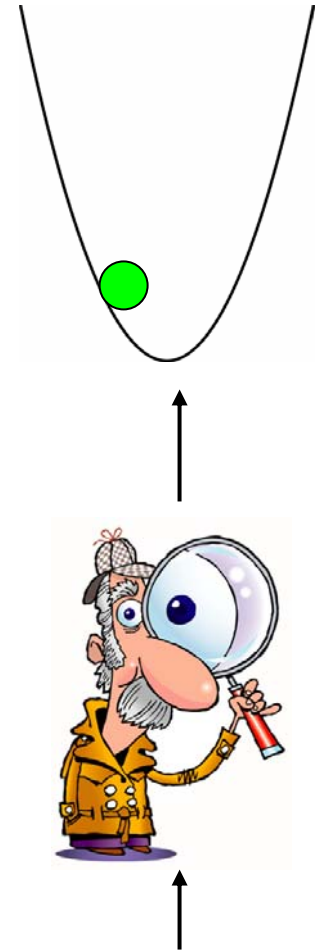
Close enough to the equilibrium (locally) stable point, one can linearize the equations, leading to the equation:

$$\frac{d\delta\vec{n}}{dt} = A \cdot \delta\vec{n}$$
$$A_{i,j} = \frac{\gamma_{i,j}^0}{n_j^0(1-n_j^0)} - \frac{e^2}{T} \sum_{l \neq i,j} \gamma_{i,k}^0 \left(\frac{1}{r_{i,j}} - \frac{1}{r_{i,k}} \right), \quad (i \neq j)$$
$$\gamma_{i,j}^0 \sim e^{-\frac{2r_{ij}}{\xi}} \quad (\text{Anderson Localization})$$

Sum of columns vanishes (particle conservation number)

For low temperatures, near a local minimum, second term is negligible \rightarrow

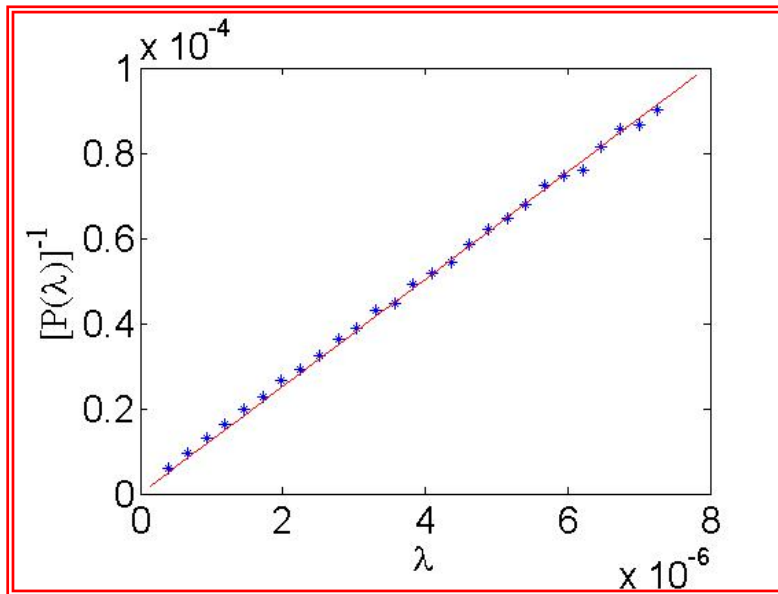
All eigenvalues are real and negative



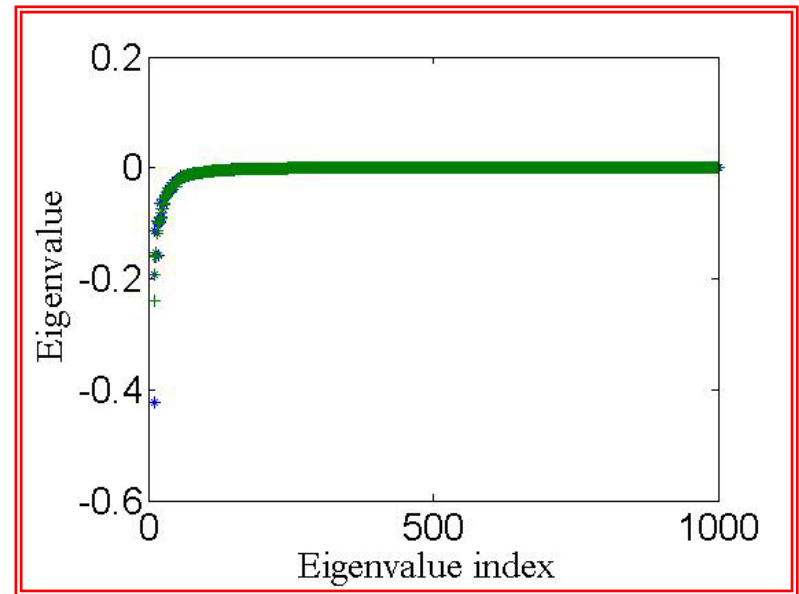
Eigenvalue Distribution

Solving numerically shows a distribution proportional to $\frac{1}{\lambda}$:

Eigenvalue distribution

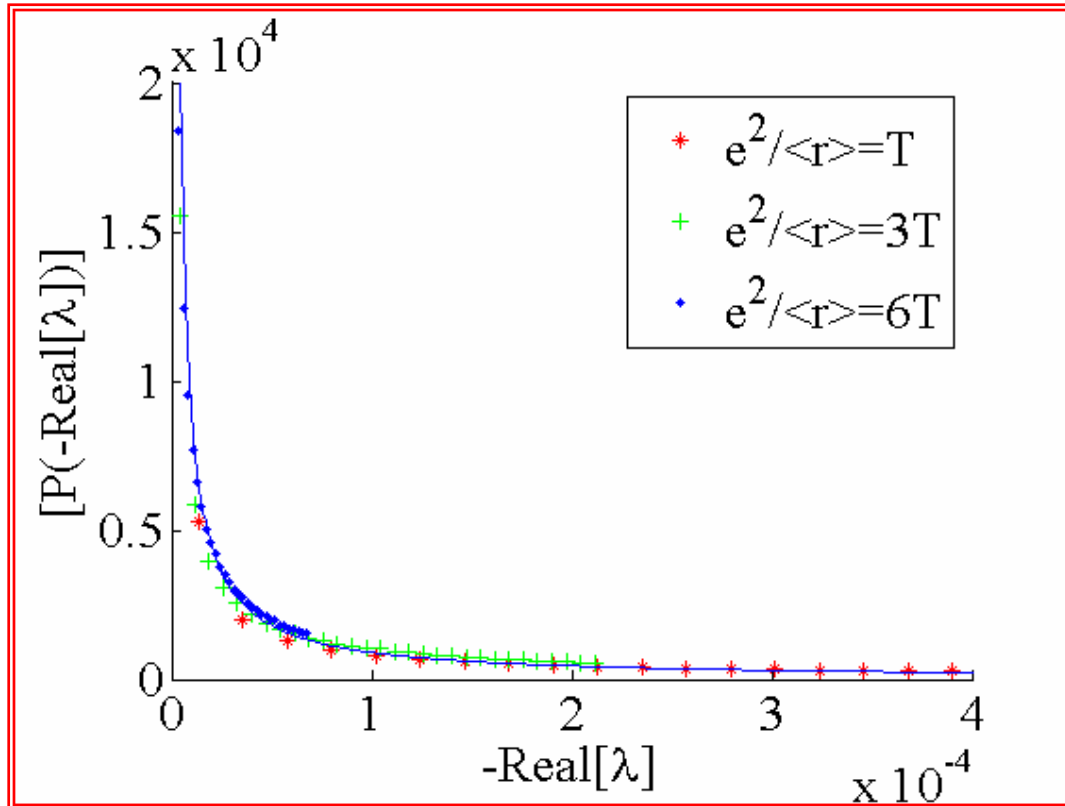


Eigenvalues



$$\sum_{\lambda} e^{-\lambda t} \longrightarrow \int P(\lambda) e^{-\lambda t} d\lambda \sim -\gamma_E - \log(\lambda_{\min} t)$$

Changing the interaction strength



- Slow relaxations occur also without interactions.
- Interactions push the distribution to (s)lower values.

Outline

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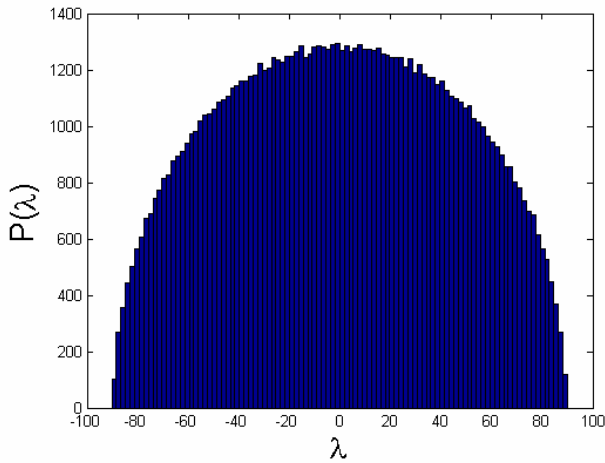
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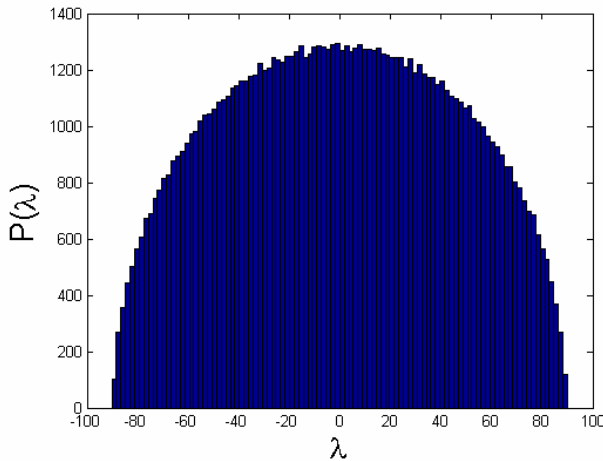
Digression: Random Matrix Theory in a nutshell

- Define a matrix with A_{ij} independent, Gaussian variables.
- Distribution of eigenvalues follows the “Wigner semicircle law”:



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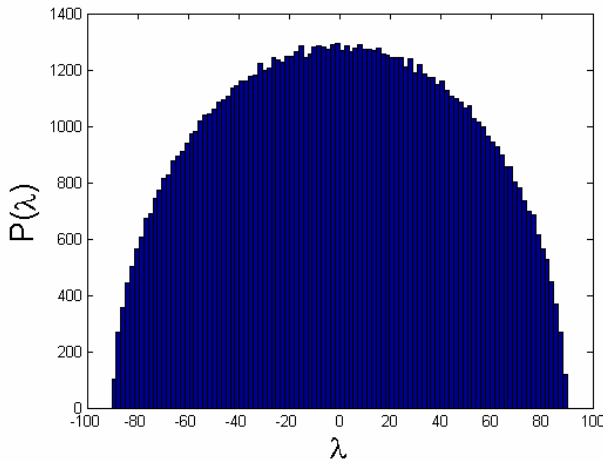


- Why is this interesting? Wide applications include:

Nuclear physics Wigner and Dirac (1951) Atomic physics Camarda and Georgopoulos *PRL* (1983) Mesoscopics Imry, *EPL* (1986), Beenakker, *PRL* (1993) Chaotic systems Bohigas, Giannoni and Schmit, *PRL* (1984) ...And many more.

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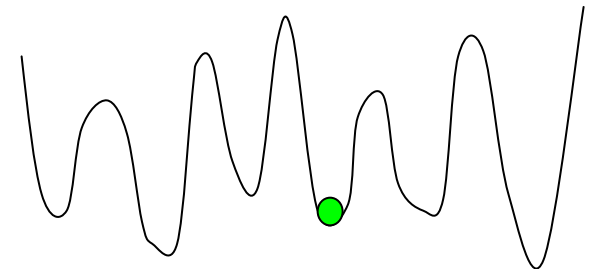
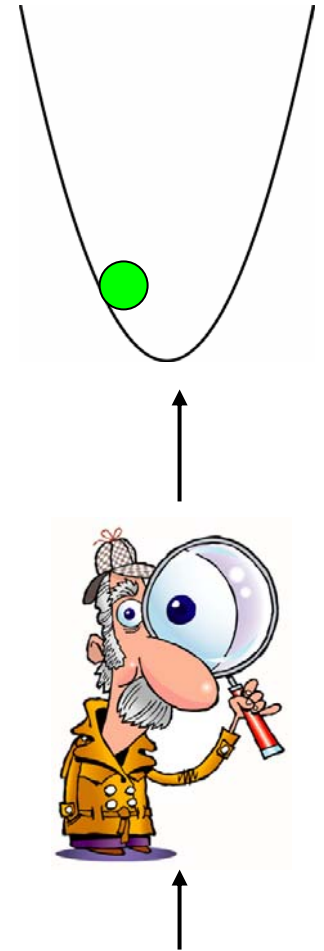
In the following we present a different class of RMT, also broadly applicable

Solution near locally stable point

Close enough to the equilibrium (locally) stable point, one can linearize the equations, leading to the equation:

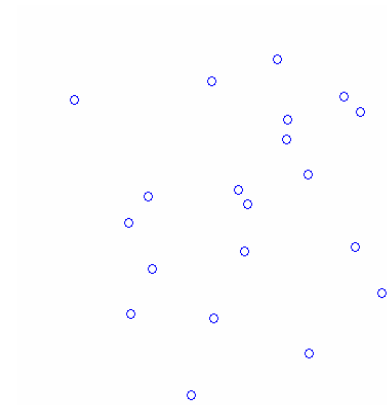
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Digression: What are Random Distance Matrices?

- 1) Choose N points randomly and uniformly in a d -dimensional cube.



I. M. Lifshitz, *Adv. Phys* (1964).

Mezard, Parisi and Zee, *Nucl. Phys.* (1999)

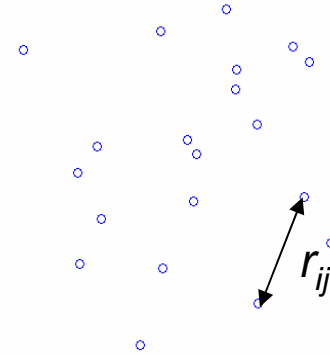
Bogomolny, Bohigas, and Schmit, *J. Phys. A: Math. Gen.* (2003).

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$$A_{i,j} = f(r_{ij}) , \quad f(r) = e^{-r/\xi}$$

(Euclidean distance) $\quad \varepsilon = \xi / \langle r \rangle$



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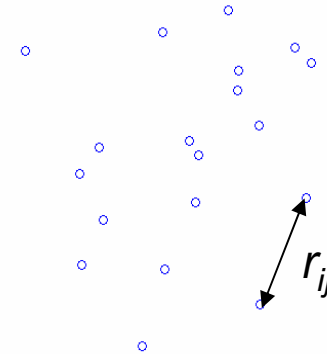
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$$A_{i,i} = -\sum_{j \neq i} A_{i,j} \quad \text{sum of every column vanishes}$$

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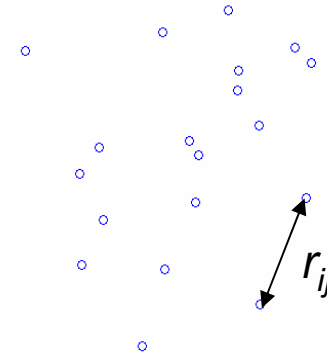
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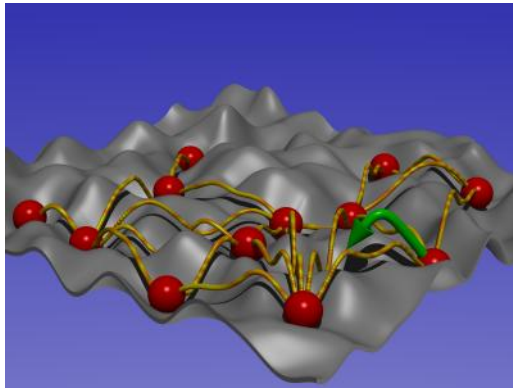
Q: What is the eigenvalue distribution?

What are the eigenmodes?

Distance matrices – Motivation

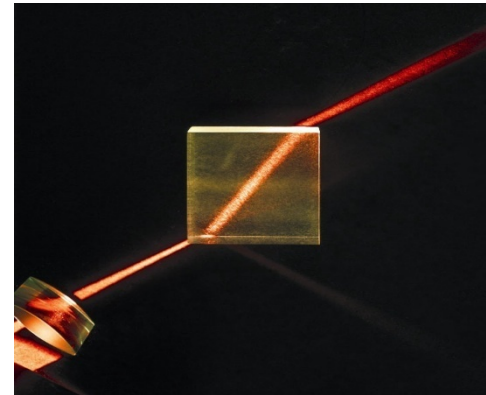
Relaxation in electron glasses

Amir, Oreg and Imry, PRB 2008



Photon propagation in a gas

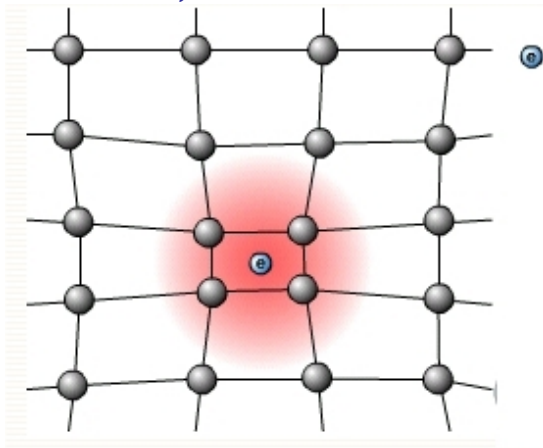
Akkermans, Gero and Kaiser, PRL 2008



Localization of phonons

Ziman, PRL 1982

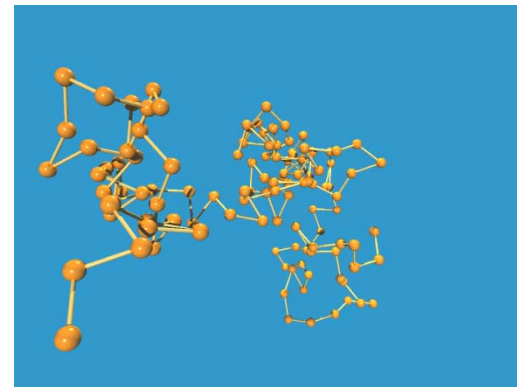
Vitelli et al., PRE 2010



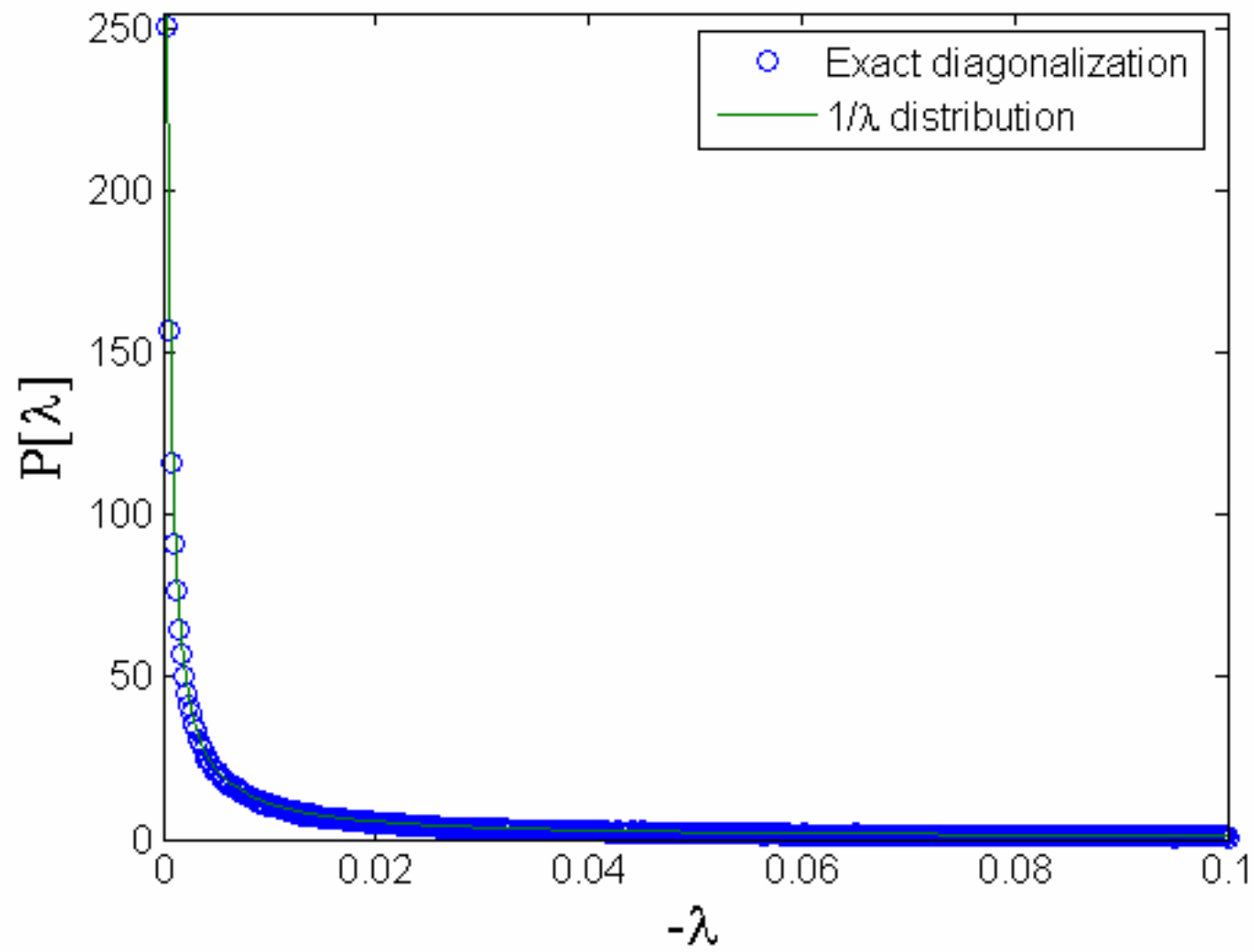
Anomalous diffusion

Scher and Montroll, PRB 1975

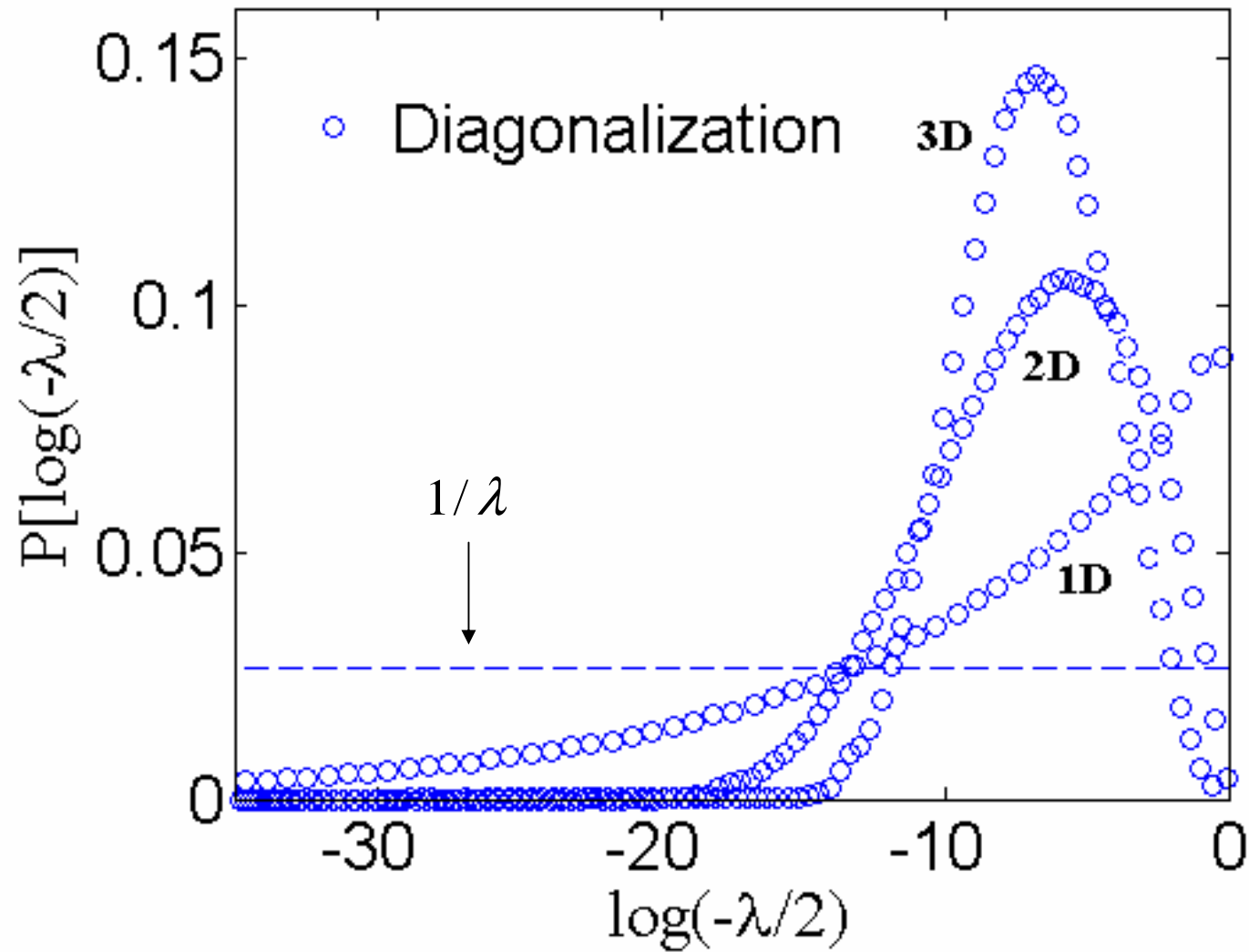
Metzler, Barkai and Klafter, PRL 1999



Results – 2D

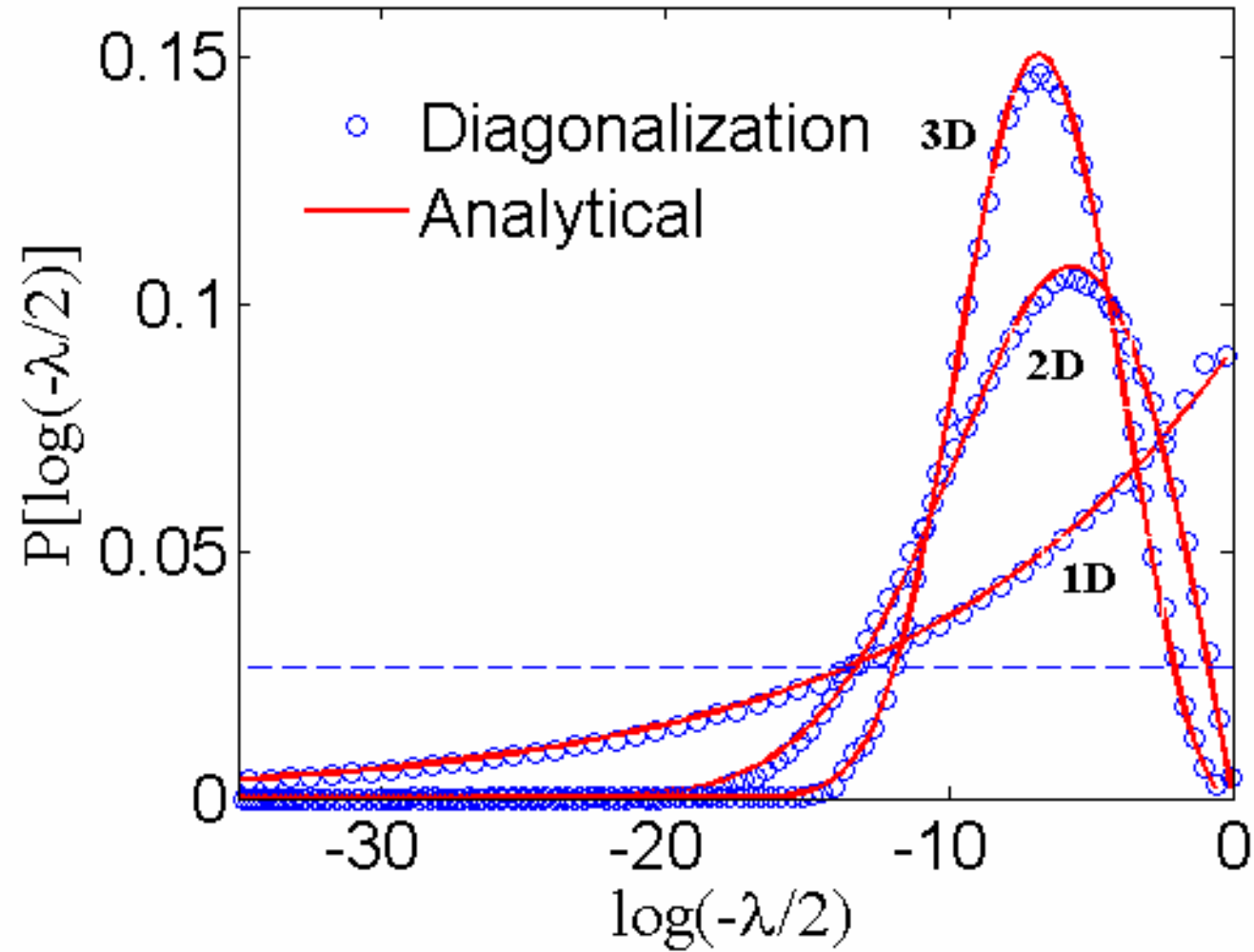


Results



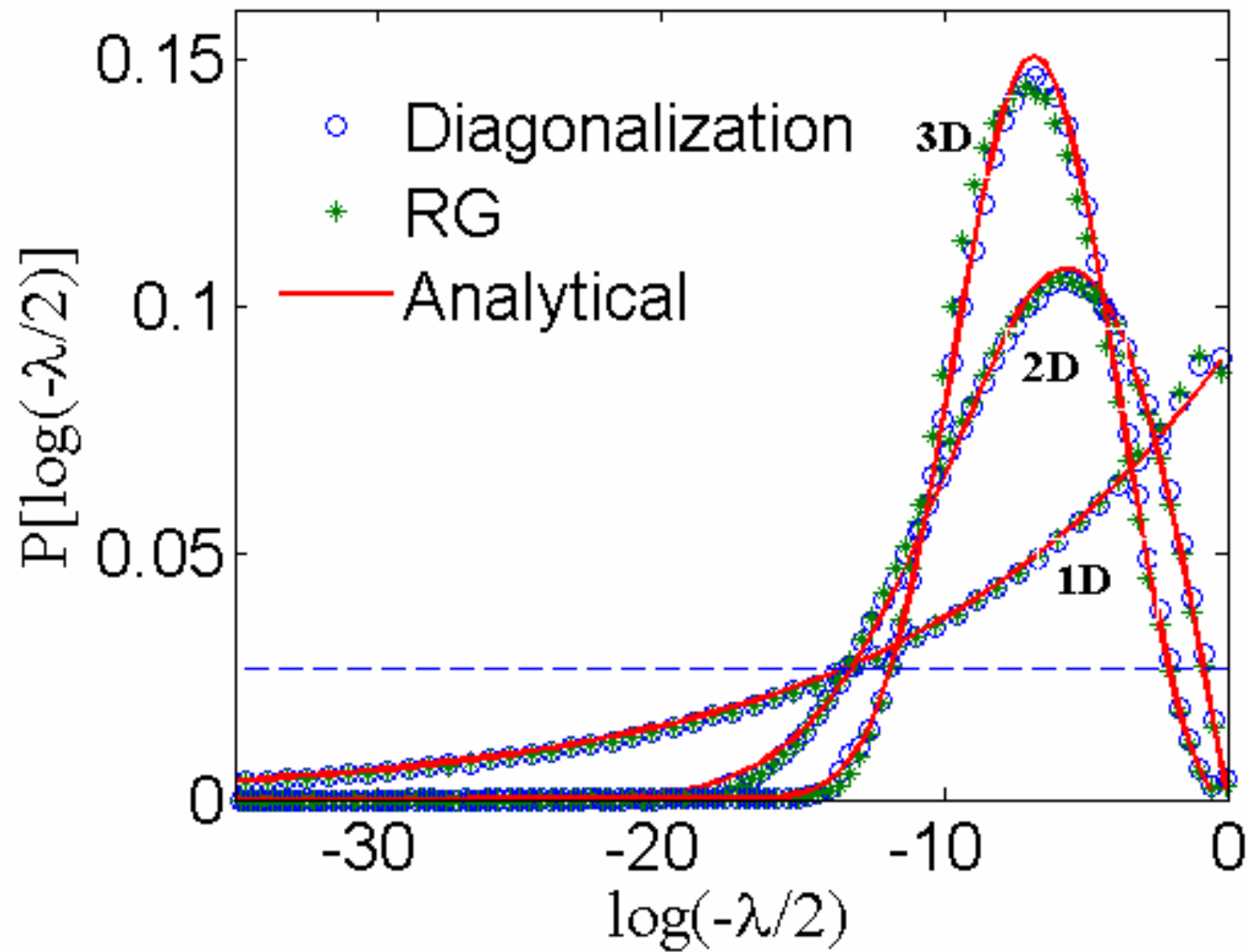
Results

(no fitting parameters)



Results

(no fitting parameters)



Exponential Distance Matrices- results

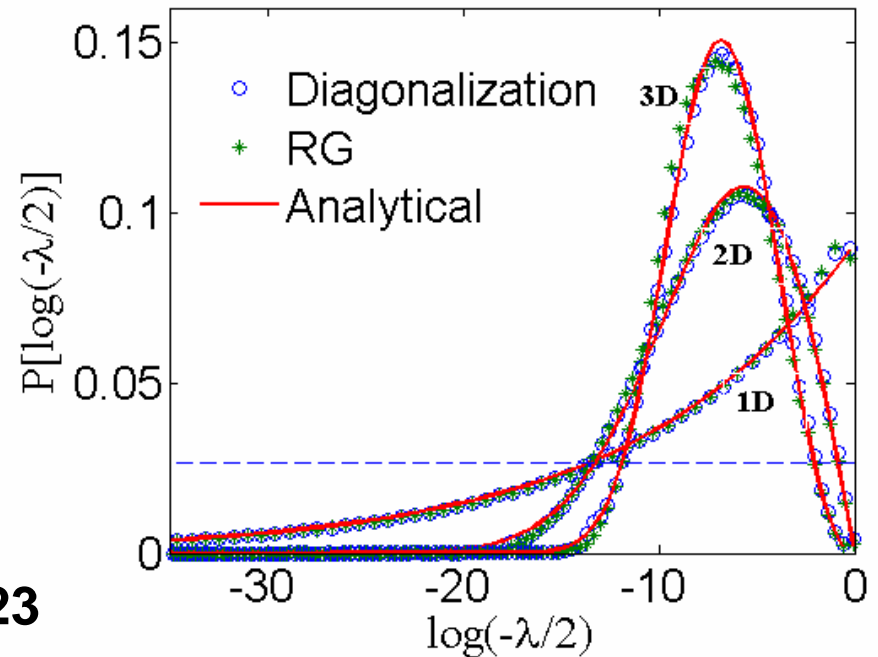
$$P(\lambda) = \frac{dC_d \varepsilon^d \log^{d-1}(\lambda/2) e^{-\frac{C_d}{2} \varepsilon^d \log^d(\lambda/2)}}{2\lambda}$$

(arbitrary dimension d)

$$\varepsilon = \xi / \langle r \rangle$$

C_d = volume of a d -dimensional sphere

- Logarithmic corrections to $1/\lambda$
- In dimensions > 1 : cutoff at $e^{-C/\varepsilon^{d/(d-1)}}$



Amir, Oreg and Imry, arxiv: 1002.2123
To appear shortly in PRL

Analytical approach in a nutshell

Part I – Moment calculation

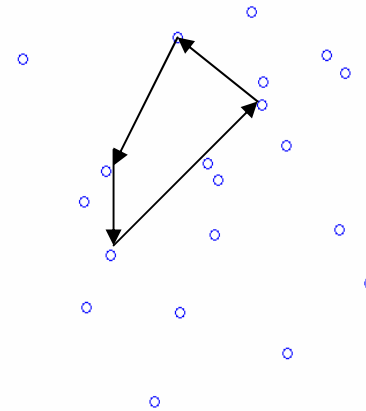
$$I_k = \int \lambda^k P(\lambda) d\lambda = \frac{1}{N} \langle A_{i_1, i_2} A_{i_2, i_3} \dots A_{i_k, i_1} \rangle$$

The k 'th moment:

$$I_k = 2^{k-1} d! C_d (\epsilon/k)^d$$

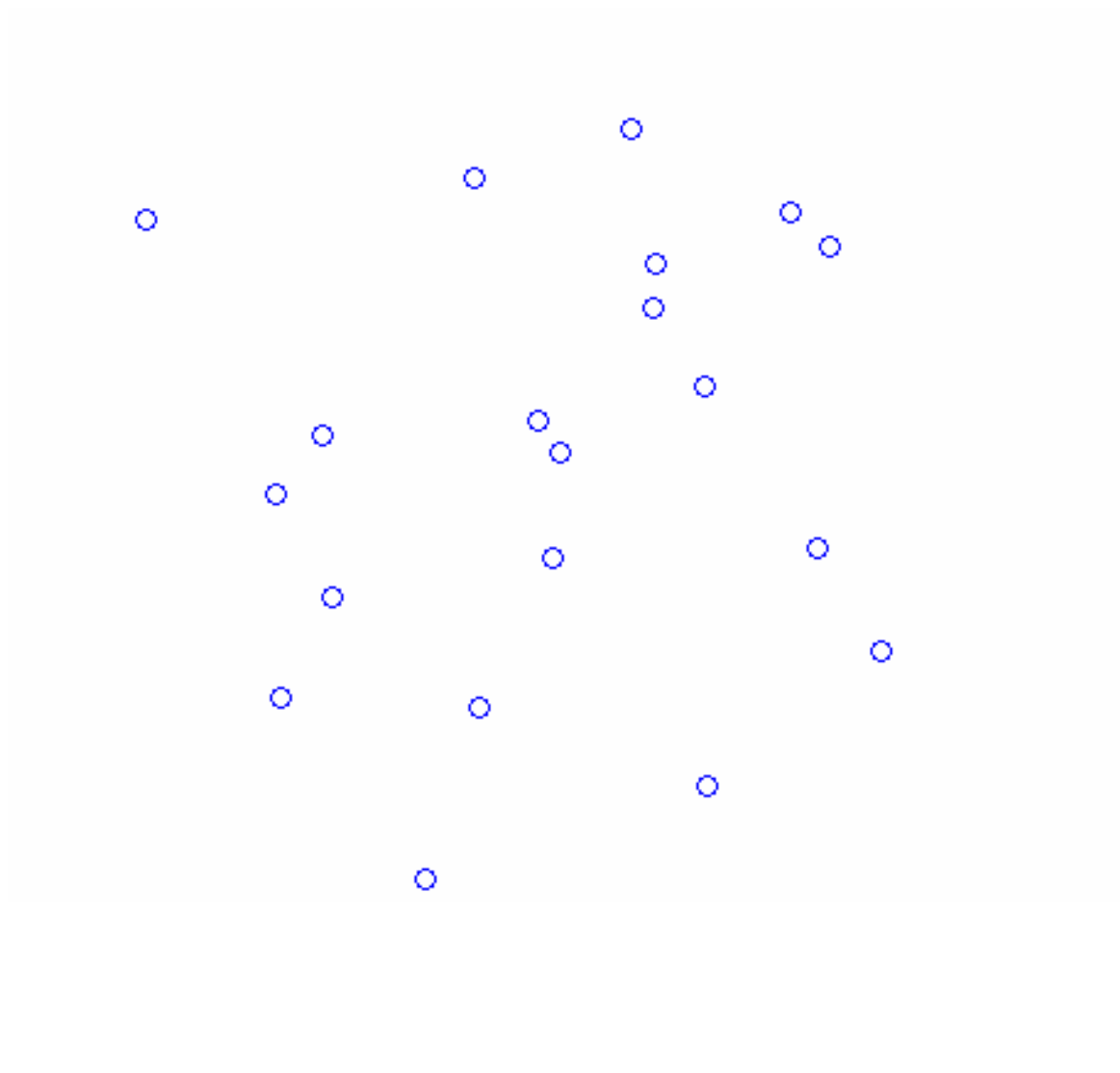


leads to the distribution function.

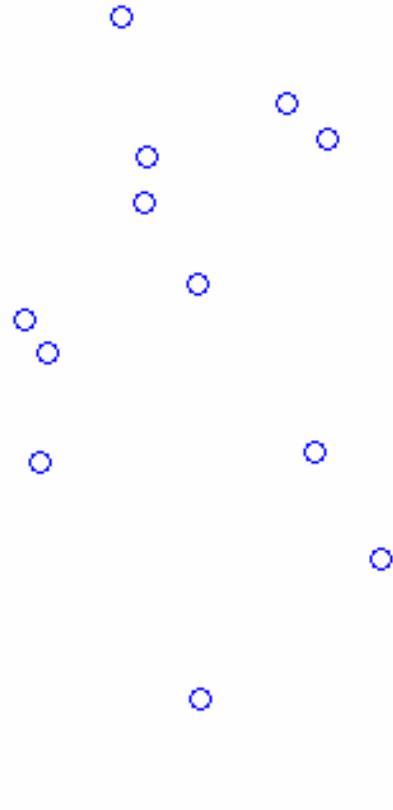


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Renormalization group approach

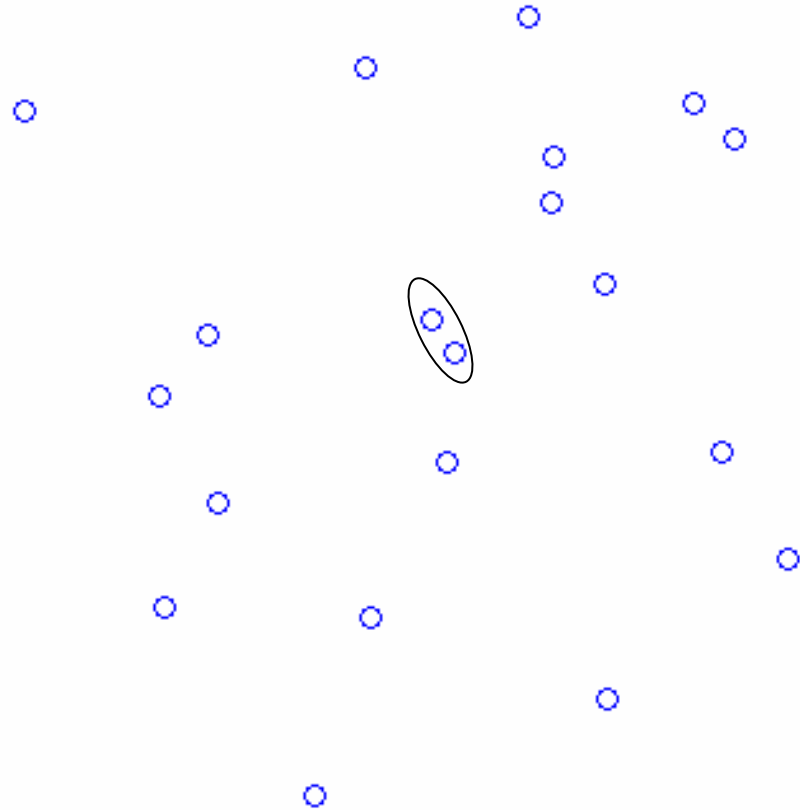


Renormalization group approach



Mechanical intuition: *network of masses and springs*

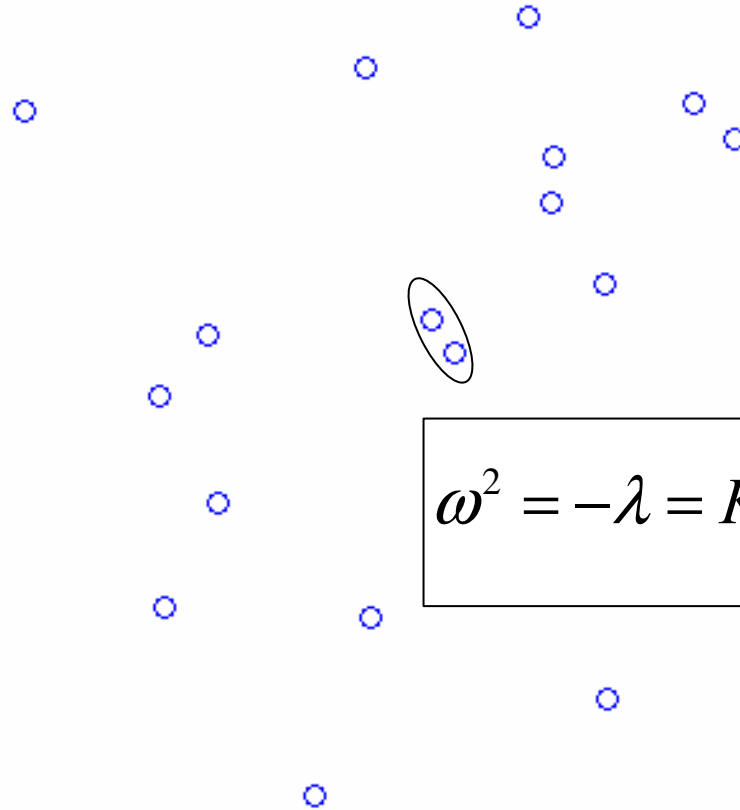
Renormalization group approach



Dasgupta and Ma, PRB 1980

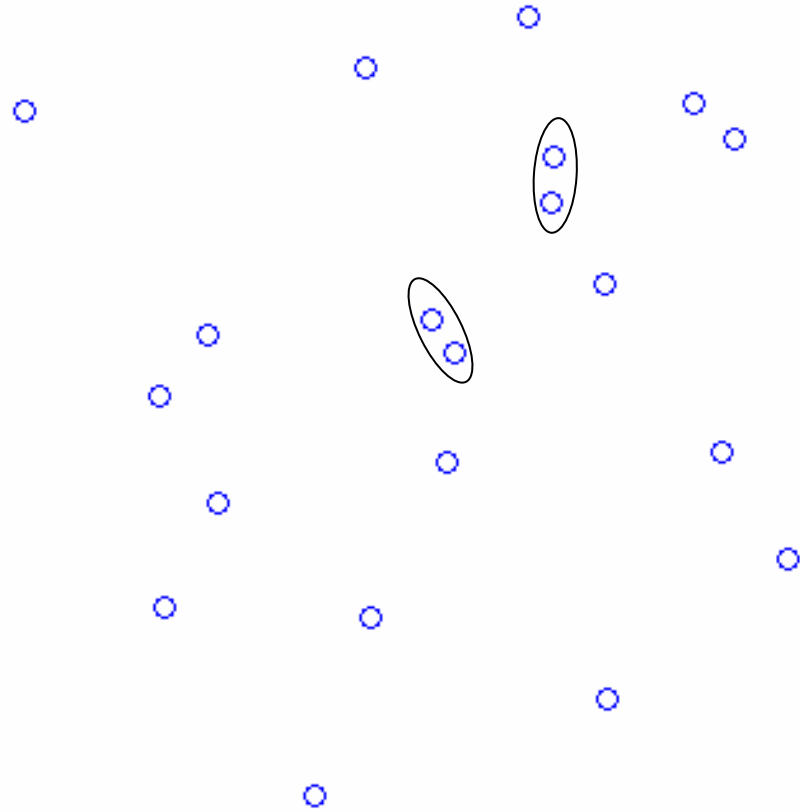
Fisher, PRL (1992) etc.

Renormalization group approach

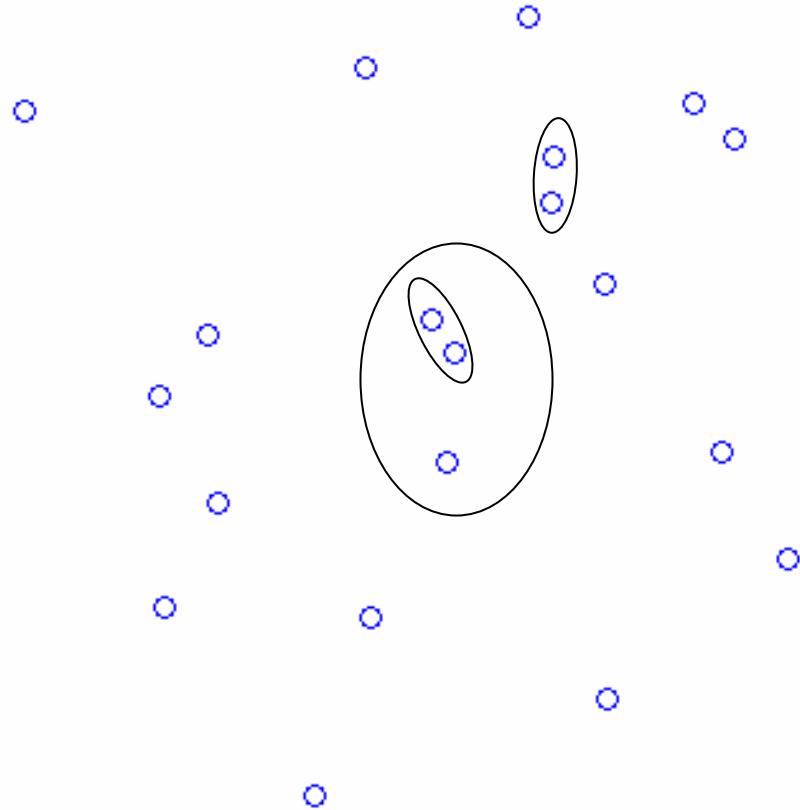


$$\omega^2 = -\lambda = K / \mu , \quad \frac{1}{\mu} = \frac{1}{m_1} + \frac{1}{m_2}$$

Renormalization group approach

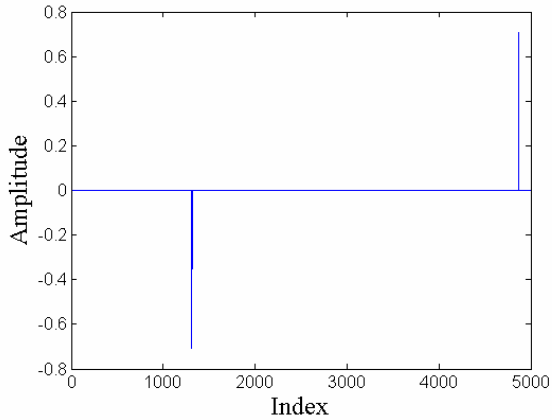


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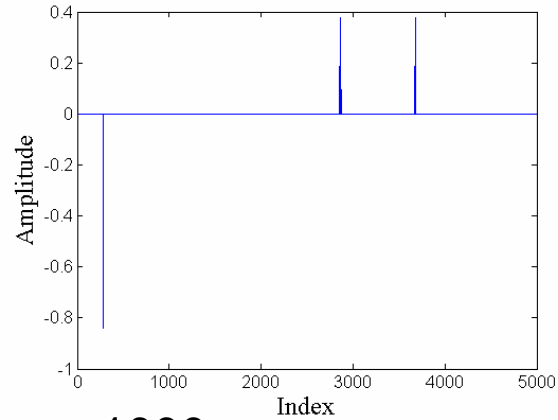


RG procedure yields growing clusters

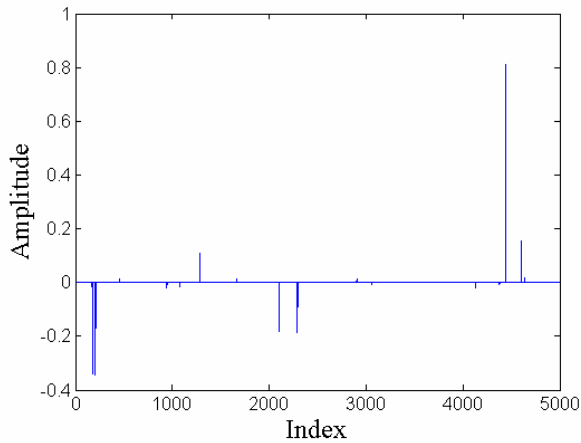
Renormalization group approach



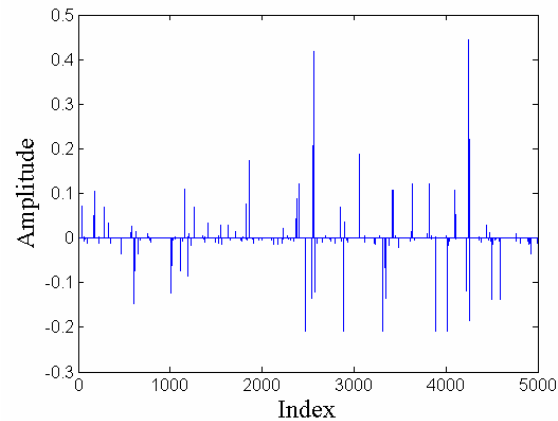
3'rd $\lambda \sim -1.86$



1000 $\lambda \sim -0.05$



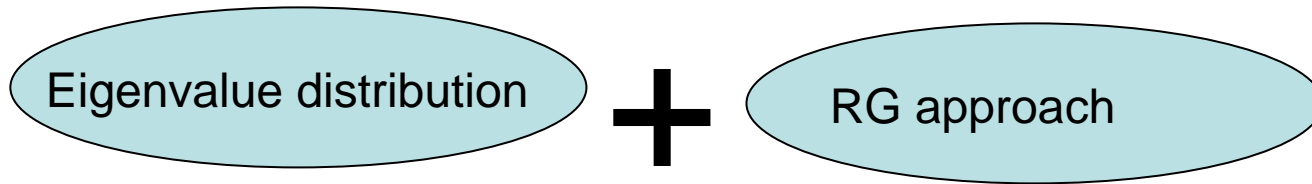
3000 $\lambda \sim -9.6 \cdot 10^{-4}$



4000 $\lambda \sim -8.5 \cdot 10^{-5}$

Examples of eigenmodes of a 5000X5000 matrix

Renormalization group approach



→

$$n_c \sim e^{\frac{C_d}{2} \varepsilon^d |\log^d(-\lambda/2)|}$$

Number of points in a cluster of a given eigenvalue

Renormalization group approach

Eigenvalue distribution

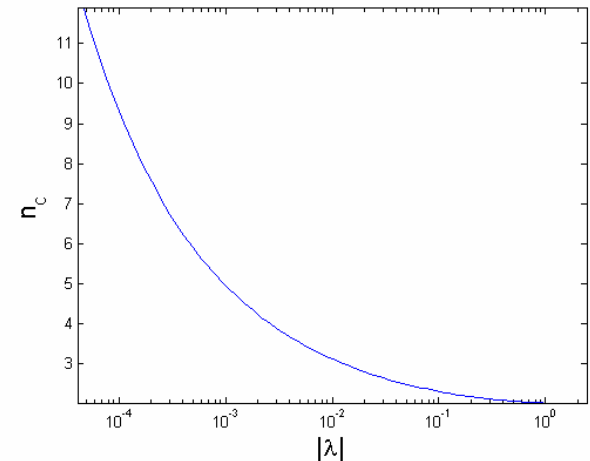


RG approach



$$n_c \sim e^{\frac{C_d}{2} \epsilon^d |\log^d(-\lambda/2)|}$$

Number of points in a cluster of a given eigenvalue



- Eigenmodes are localized clusters (“phonon localization”)
- Size of clusters diverges at low frequencies

Amir, Oreg and Imry, *Localization, anomalous diffusion and slow relaxations: a random distance matrix approach*, **arxiv: 1002.2123**

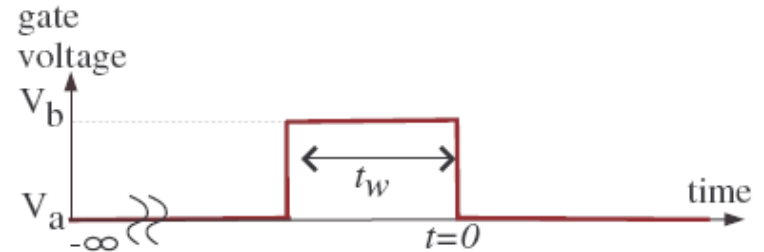
Intermediate Summary

- Statics: Coulomb gap
- Dynamics: Microscopic model reduces to RMT problem of a different class.
- Eigenvalue distribution (relaxation times) follow approximately $1/\lambda$ distribution.
- Same RMT class relevant for various other problems.

Electron glass aging– experimental protocol

Step I

System equilibrates for long time

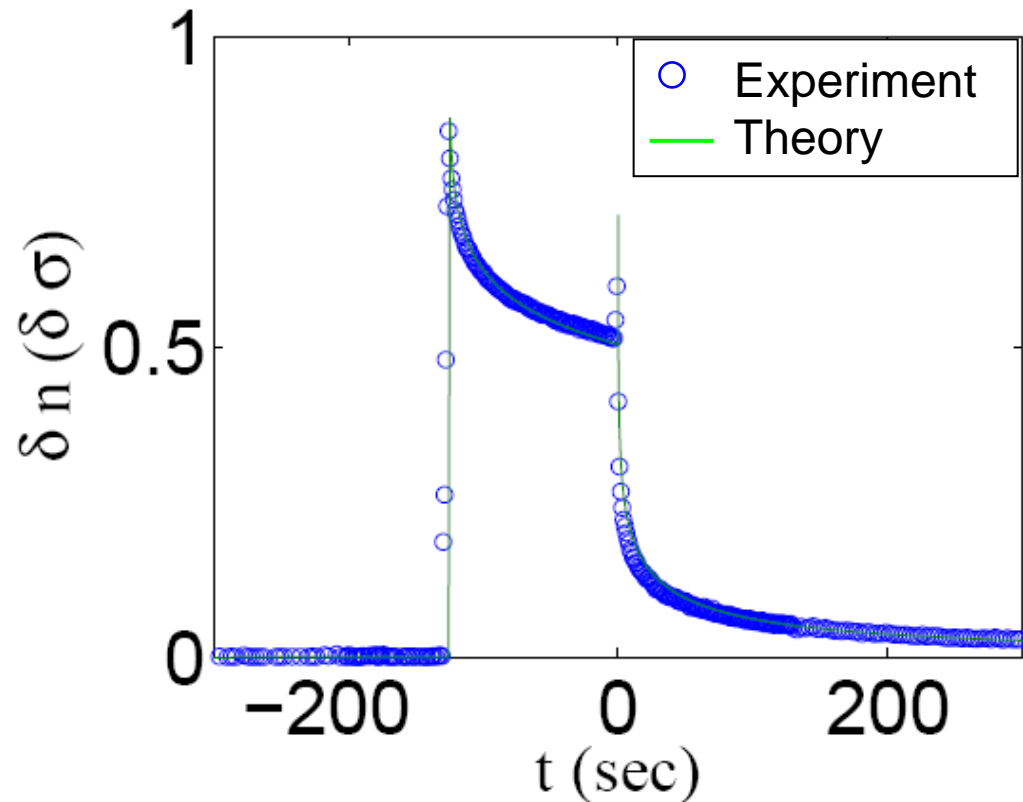


Step II

V_g is changed, for a time of t_w .

Throughout the experiment

Conductance is measured as a function of time.

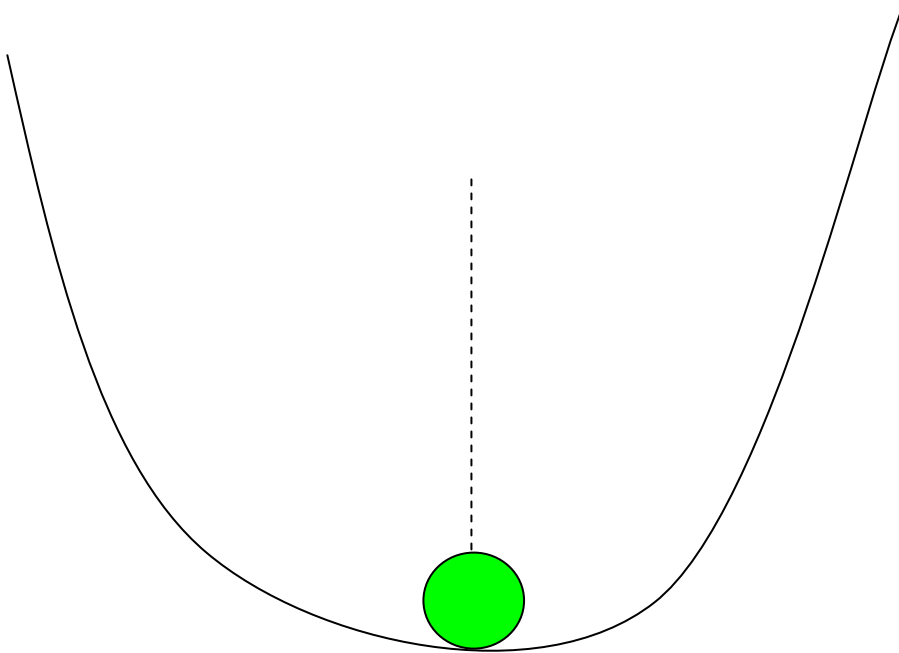


Data: *Ovadyahu et al.*

Aging – physical picture

Assume a parameter of the system is slightly modified (e.g: V_g)

After time t_w it is changed back. What is the response?

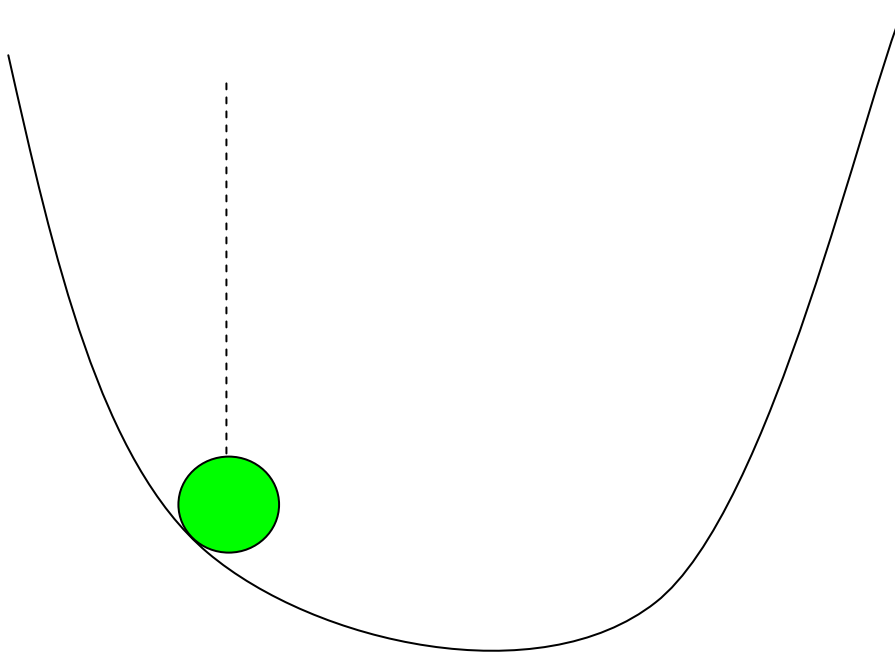


Initially, system is at some local minimum

Aging – physical picture

Assume a parameter of the system is slightly modified (e.g: V_g)

After time t_w it is changed back. What is the response?



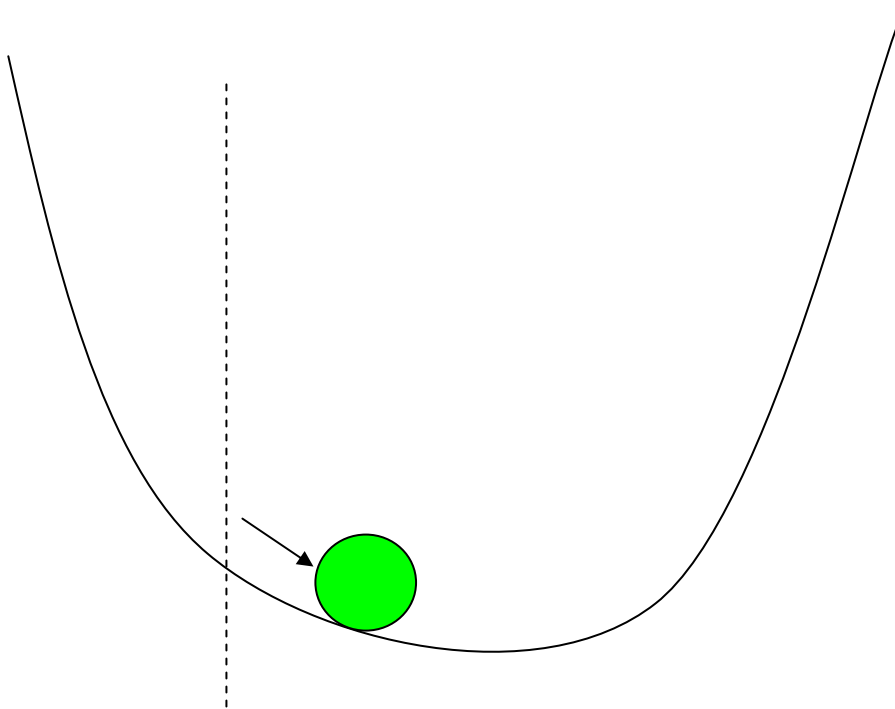
At time $t=0$ the potential changes,

and the system begins to roll towards the new minimum

Aging – physical picture

Assume a parameter of the system is slightly modified (e.g: V_g)

After time t_w it is changed back. What is the response?

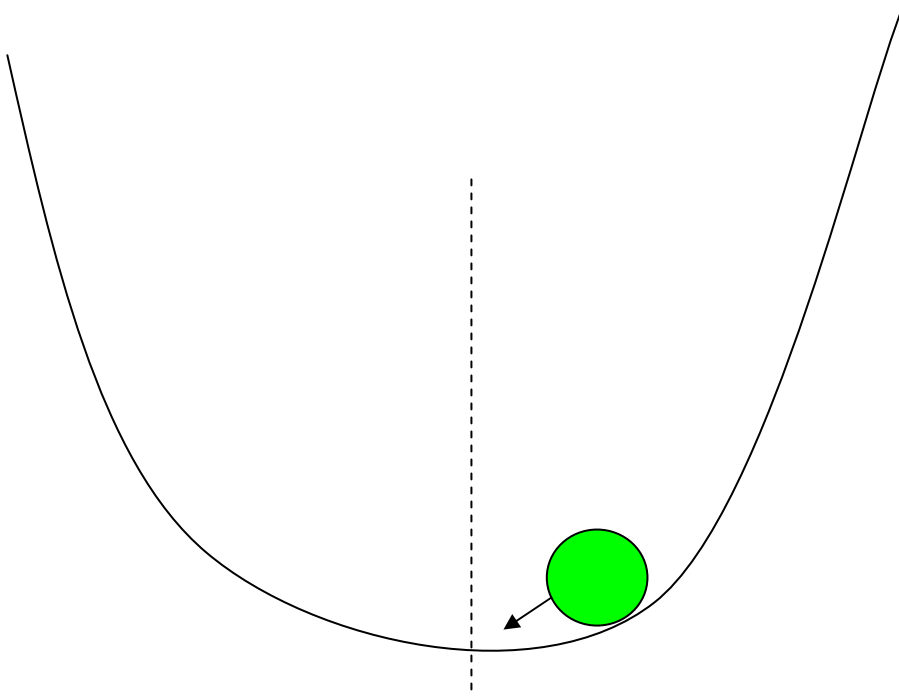


At time t_w the system reached some new configuration

Aging – physical picture

Assume a parameter of the system is slightly modified (e.g: V_g)

After time t_w it is changed back. What is the response?



**Now the potential is changed back to the initial form-
the particle is not in the minimum!**

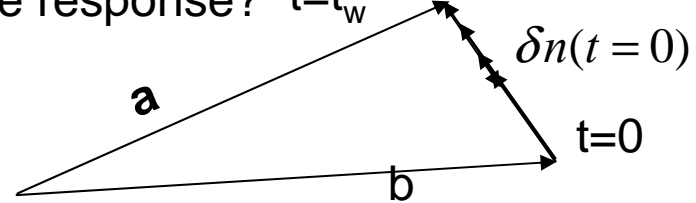
The longer t_w , the further it got away from it.

It will begin to roll down the hill.

Aging – Analysis

Assume a parameter of the system is slightly modified (e.g: V_g)

After time t_w it is changed back. What is the response? $t=t_w$



Sketch of calculation

If a and b configurations are close enough in phase space:

$$\delta n(t = t_w) \sim \sum_{\text{eigenmodes } \alpha} \chi_{\alpha} e^{-\lambda_{\alpha} t_w} |V_{\alpha}\rangle \Rightarrow \sum_{\text{eigenmodes } \alpha} e^{-\lambda_{\alpha} t_w} =$$

modes are independent and contribute uniformly

Logarithmic relaxation during step II

Time t after the perturbation is switched off:

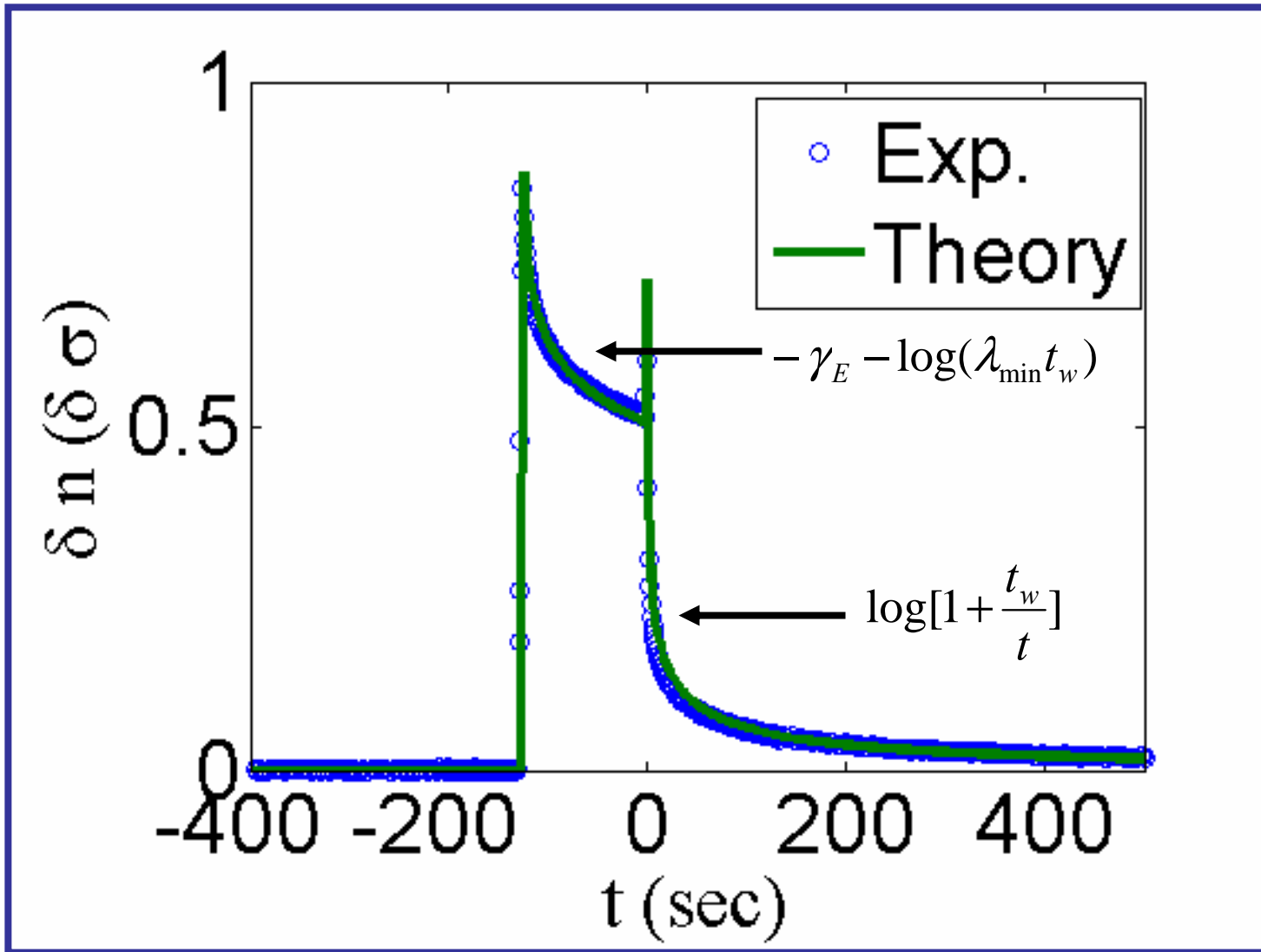
$$\delta n(t) \sim \sum_{\text{eigenmodes } \alpha} \chi_{\alpha} (1 - e^{-\lambda_{\alpha} t_w}) e^{-\lambda_{\alpha} t} |V_{\alpha}\rangle = f(t + t_w) - f(t)$$

Full aging

Only $1/\lambda$ distribution yields full aging!

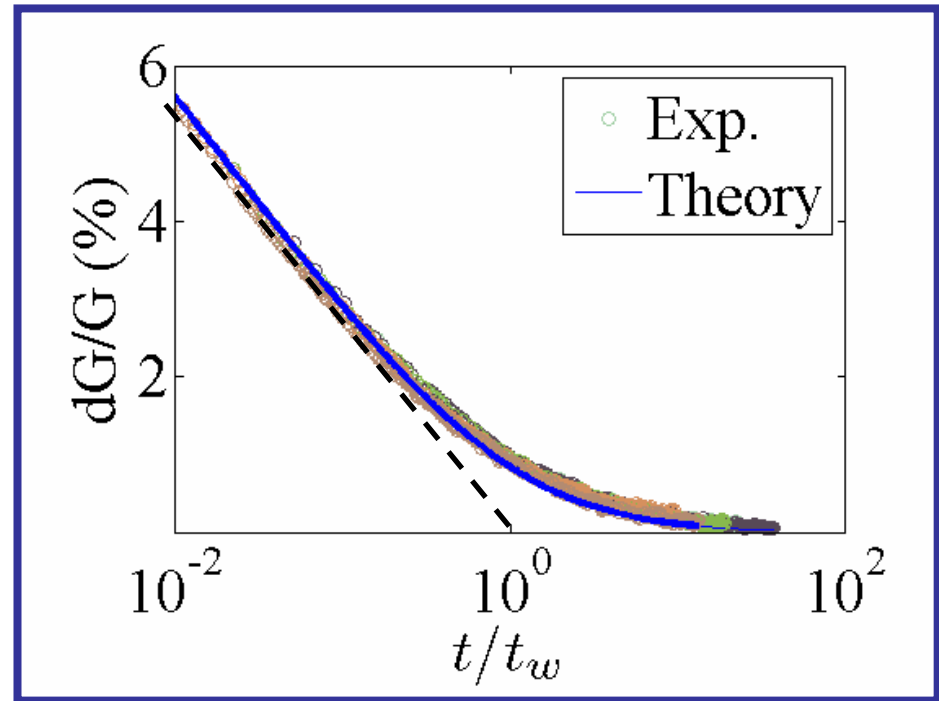
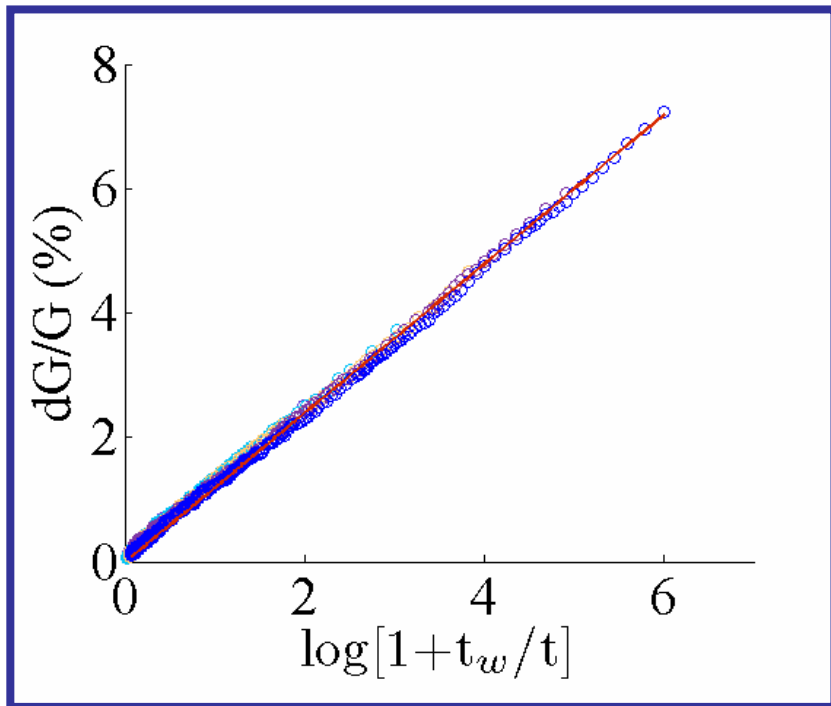
See also: T. Grenet et al. Eur. Phys. J B 56, 183 (2007)

Aging Protocol - Results



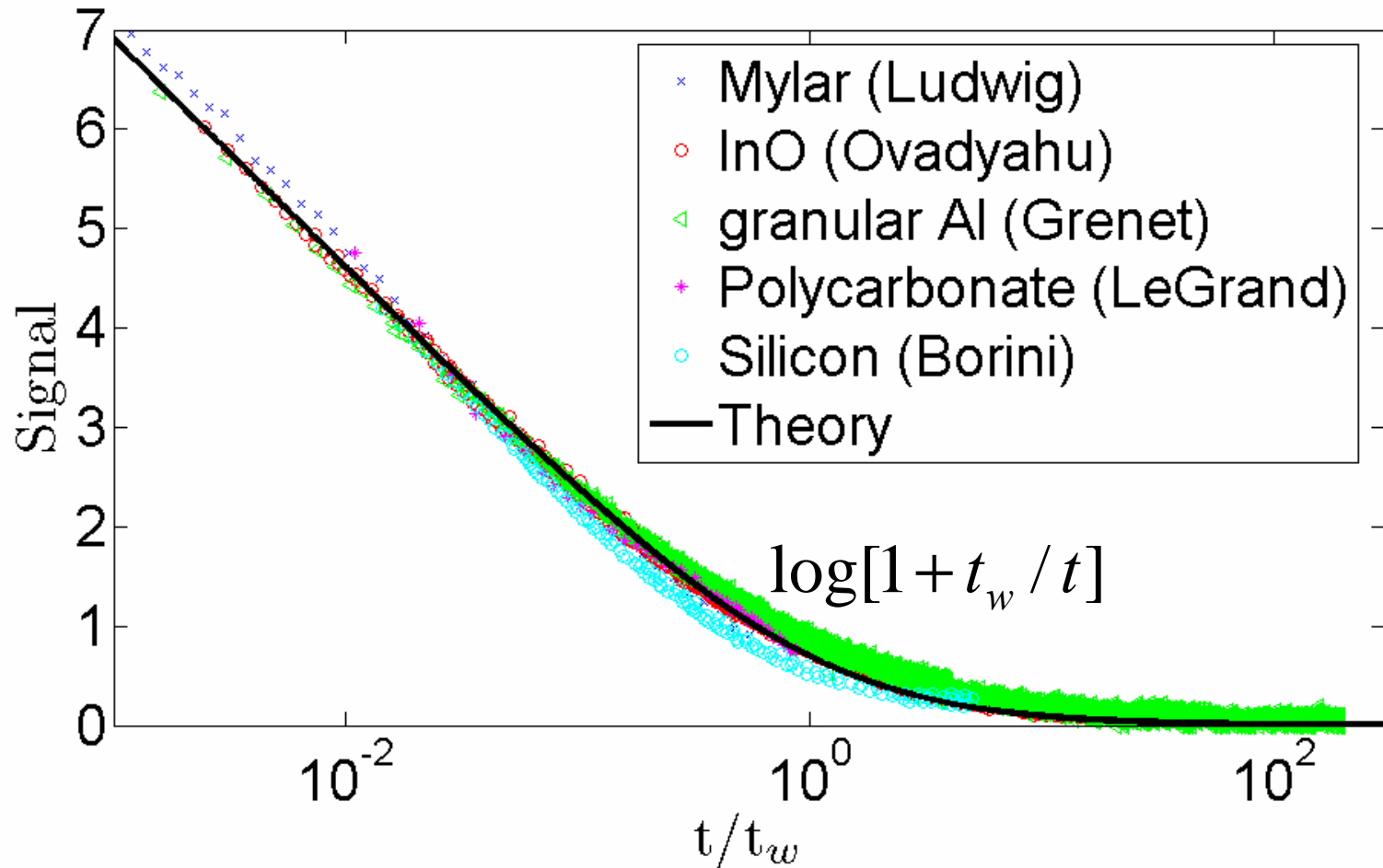
Detailed fit to experimental data

- Full aging
- Deviations from logarithm start at t/t_w



Amir, Oreg and Imry, *PRL* (2009)

Full aging and universality



Amir, Oreg and Imry, to be published

Conclusions

- Statics: Coulomb gap
- Dynamics near locally stable point: many slow *localized* modes, $\sim \frac{1}{\lambda}$ distribution.

How universal? **We believe: a very relevant RMT class.**

- One obtains **full aging**, with relaxation approximately of the form :

$$\delta\sigma \sim \log\left[1 + \frac{t_w}{t}\right]$$

More details:

Phys. Rev. B 77, 1, 2008 (local mean-field model)

Phys. Rev. Lett. 103, 126403 (2009) (aging properties)

Phys. Rev. B 80, 245214 2009 (variable-range hopping)

Ann. Phys. 18, 12, 836 (2009) ($1/f$ noise)

arxiv: 1002.2123, soon in Phys. Rev. Lett. (exponential matrices – solution)

Electron glass dynamics – Review (soon online)