

s sourcery

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References

- s source framework, gapped phases w/ JM - 1407.8203
- entropy bound w/ JM - 1505.07106
- gapless “square root” states w/JM and SX - *coming soon*
- Important other work: White’s DMRG, MPS, Vidal’s MERA, other tensor networks, long history of real space RG, see 1407.8203 for an extensive list of references

The ground state problem

Given a local Hamiltonian H , “determine” its ground state and compute physical properties

Is there any way this problem could have a general solution? Very hard to imagine,

provably false in some cases, yet I hope to

convince you that for a very broad class of H the answer is YES!



[QMA hard: Gottesman et al., also glassy states, etc.]

Families of Hamiltonians/states

Throughout we consider **families of Hamiltonians indexed by system size**: $\{H_L\}$

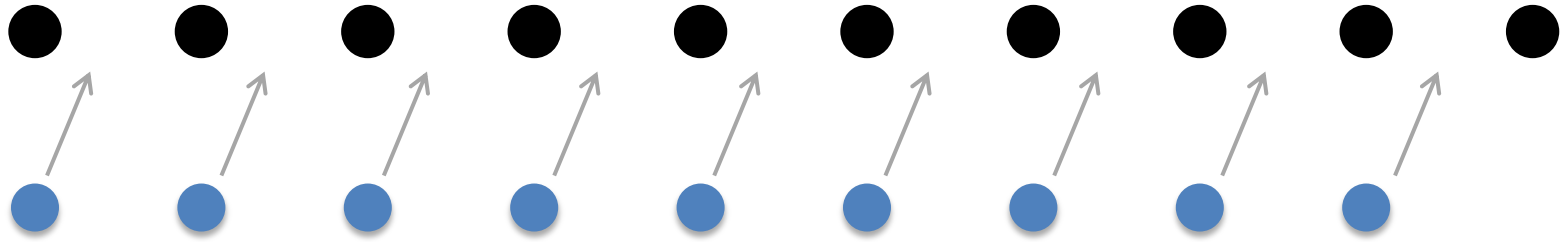
These Hamiltonians have **corresponding ground states**: $\{|\psi_L\rangle\}$

We will study transformations between states at different L ; **product states can always be subtracted or added at will**: $|0\rangle^M, \forall M$

DEFINITION AND PROPERTIES

“RG” construction of wavefunction

L sites



L sites



L black sites are interleaved with L blue sites using a **quasi-local unitary**. The output is the black state on $2L$ sites.



$2L$ sites

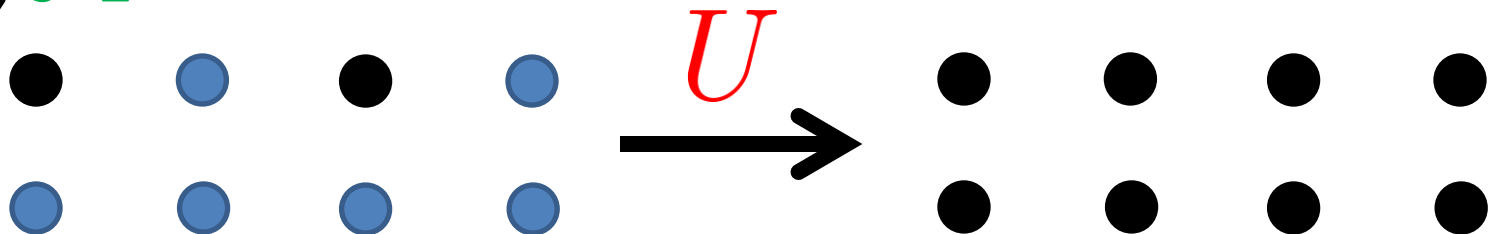
$$|\psi_{2L}\rangle = U(|\psi_L\rangle |0\rangle^L)$$

slight abuse of notation: “fixed point”
[BGS-McGreevy ‘14]

s source RG fixed point

A d -dimensional s source RG fixed point is a system where a ground state on $(2L)^d$ sites can be constructed from s copies of ground states on L^d sites times some unentangled degrees of freedom by acting with a quasi-local unitary

$d=2, s=1$



s source RG definition again

A family of states is s source fixed point if
(for large enough L):

$$\begin{array}{ccc}
 \text{d-torus} & & \text{d-torus} \\
 \underbrace{L \times \dots \times L}_d & \longleftrightarrow & \underbrace{2L \times \dots \times 2L}_d \\
 U_L \left(\underbrace{|\psi_L\rangle \dots |\psi_L\rangle}_s |0\rangle^{(2^d - s)L^d} \right) & = & |\psi_{2L}\rangle
 \end{array}$$

Some properties

$$S(A) = -\text{tr}(\rho_A \log \rho_A)$$

Recursive entropy
bounds:

$$S(2R) \leq sS(R) + kR^{d-1}$$

$$S(2R) \geq sS(R) - k'R^{d-1}$$

result uses [Van Acoleyen-Marien-Verstraete]

$G(L)$ = ground state degeneracy

Ground state
degeneracy lemma
(for gapped case):

$$G(2L) = G(L)^s$$

Local operators \rightarrow local operators

$$|\psi_{2L}\rangle = U(|\psi_L\rangle|0\rangle^L)$$

$$\langle\psi_{2L}|O_{loc}|\psi_{2L}\rangle = \langle\psi_L|\tilde{O}_{loc}|\psi_L\rangle$$

iterate: can compute local expectation value in $O(\log(L))$ steps

$$\tilde{O}_{loc} = \langle 0|^L U^\dagger O_{loc} U |0\rangle^L$$

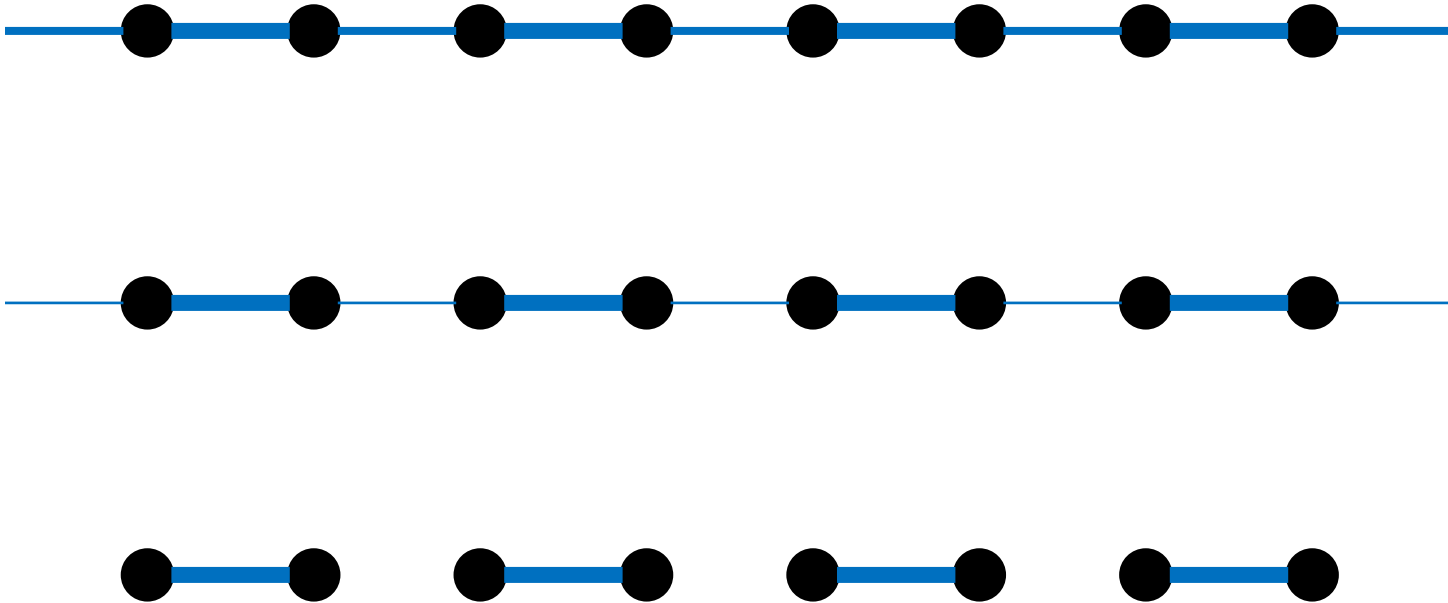
O remains local because:

1. U spreads O by at most the speed of light times a time of order one
2. The number of sites is halved at every step

EXISTENCE RESULTS

Example: trivial insulator, $s=0$

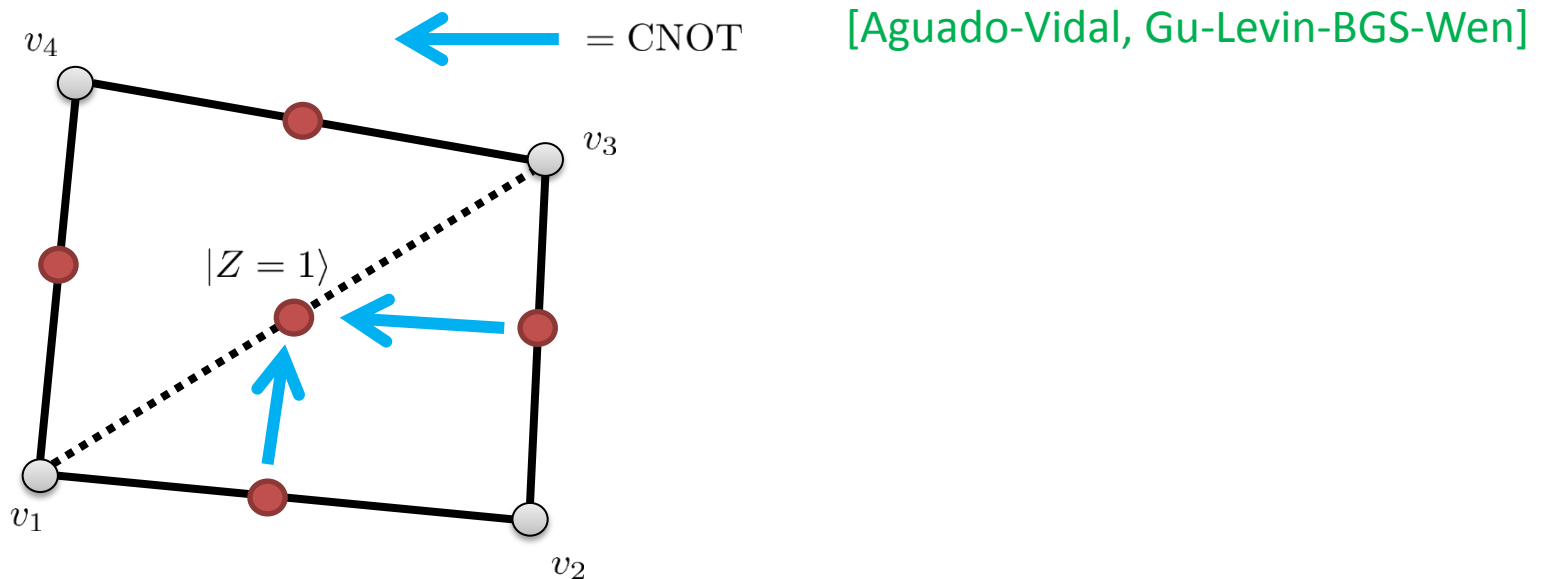
One particle per unit cell, alternating weak bonds:



Deform weak
bonds to zero

Example: (gapped) gauge theory, $s=1$

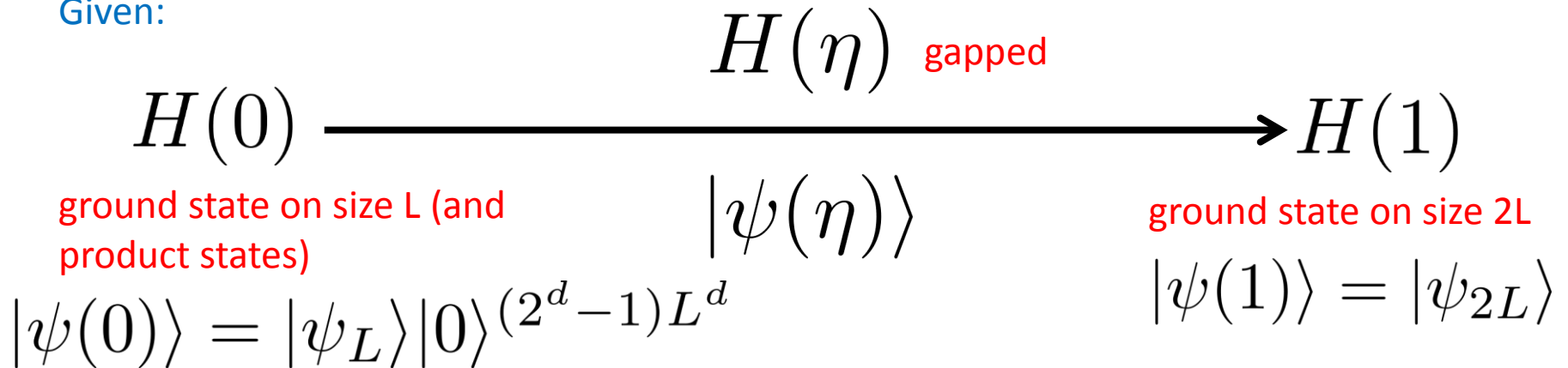
Toric code, discrete gauge theory, $d>2$, ...



Topological quantum liquid: insensitive to arbitrary smooth deformations of space

Tool: adiabatic expansion

Given:



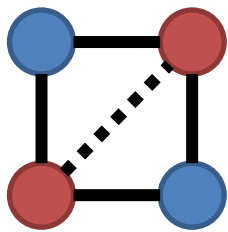
→ Exact: use Hastings-Wen quasi-adiabatic tech

Note: this may not be the most efficient U. However, this is a **non-variational** way to construct the ground state!

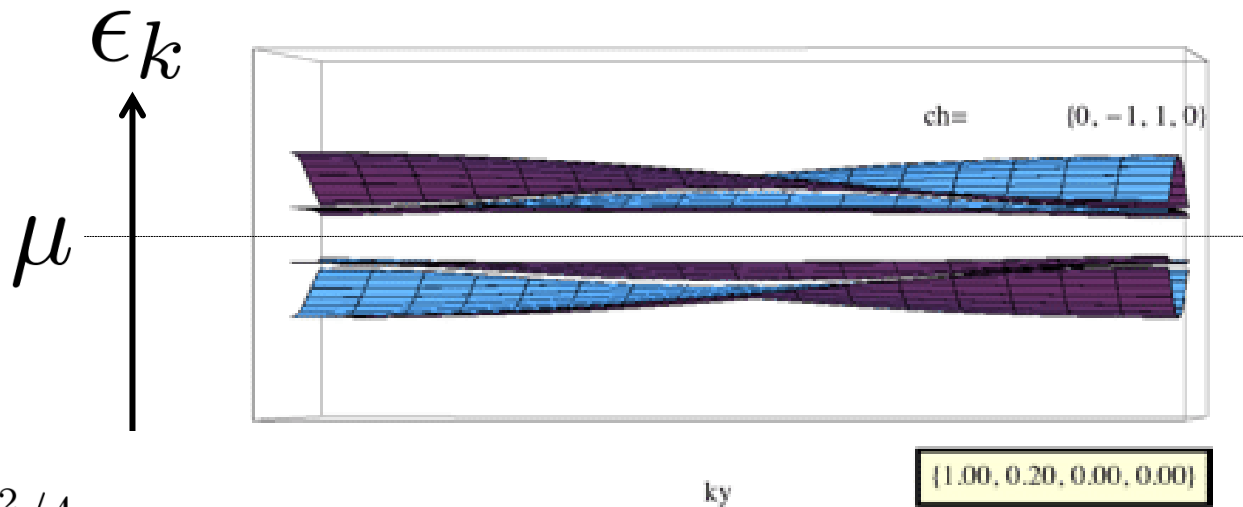
Example: chiral insulators, $s=1$

Examples:

1. Integer quantum Hall, Chern insulators
2. Massive Dirac fermion, $d=2$



Sites: $L^2 \rightarrow L^2/2 \rightarrow L^2/4$



Example: CFTs, $s=1$ [CONJECTURE]

Some evidence:

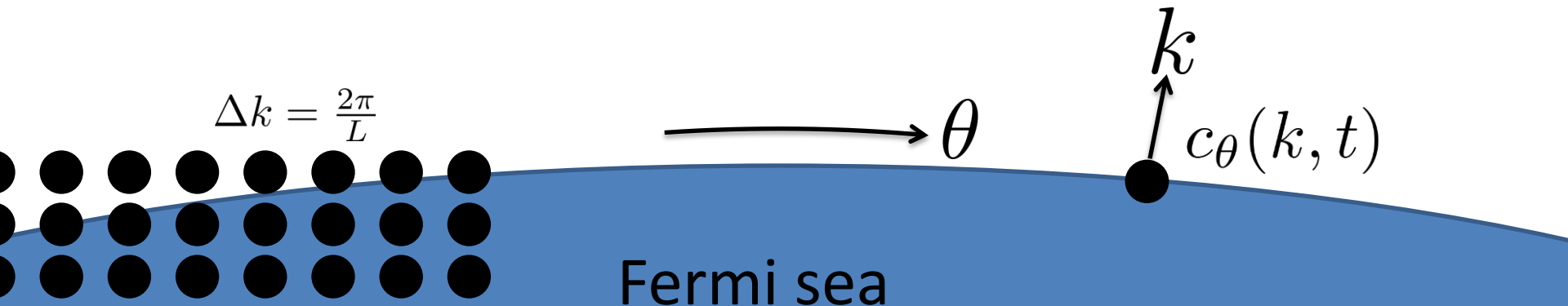
1. Consistent with structure of entanglement and correlations
2. Correlations easy to include
3. MPS approximation results in 1d [Verstraete-Cirac]

Later: provably true for some other gapless (but non-relativistic) scale invariant states ...

Example: FS, $s=2^{d-1}$ [CONJECTURE]

- **Conjecture:** Metals (Fermi liquids) in d dimensions are fixed points but require multiple copies of size L to make size $2L$

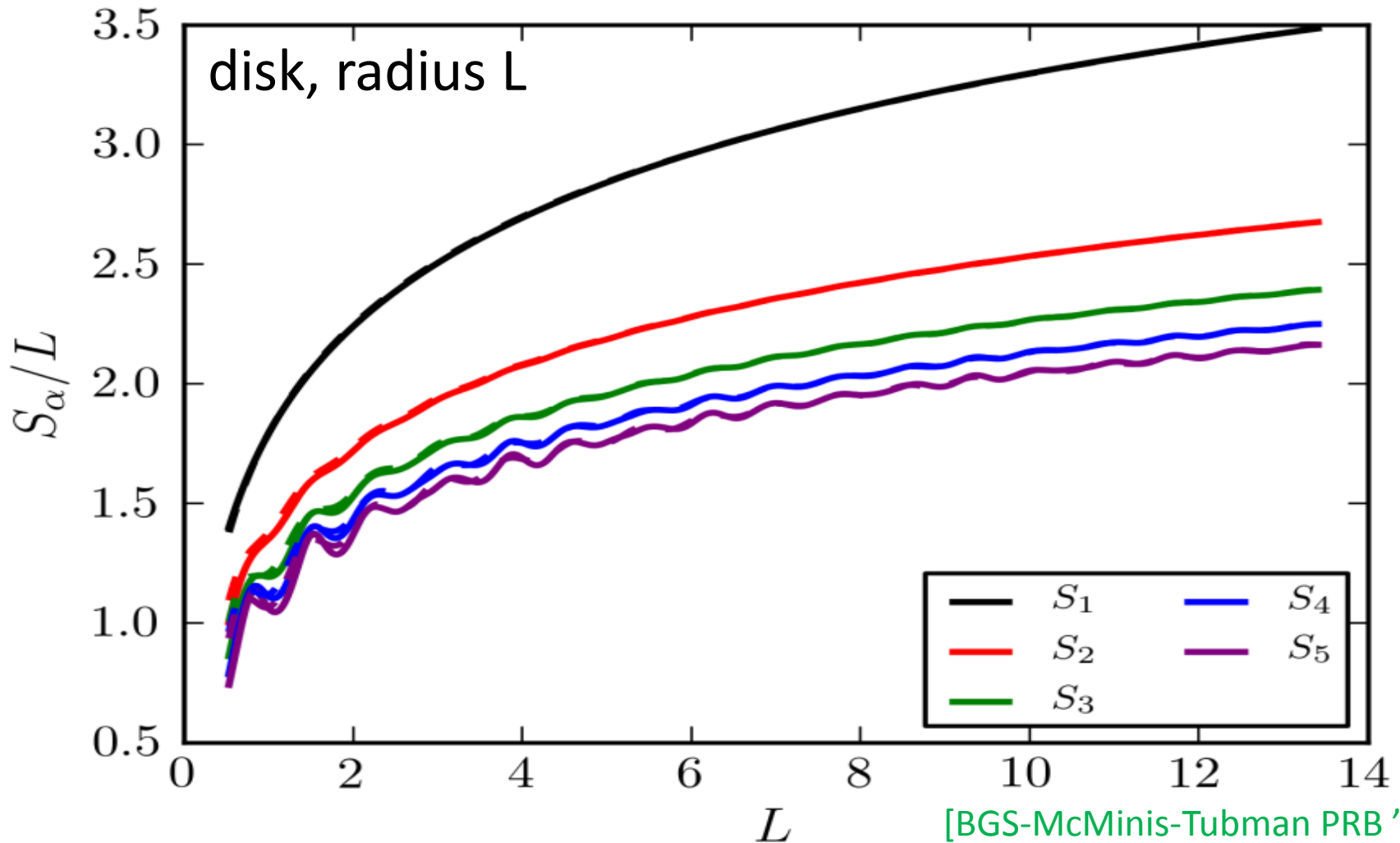
$$\mathcal{S} = \int d\theta \int dt dk c_{\theta}^{\dagger}(k, t) (i\partial_t - v_F k) c_{\theta}(k, t)$$



$$S_\alpha = \frac{1}{1-\alpha} \log(\text{tr}(\rho^\alpha)) \quad \text{general definition, probes spectrum}$$

$$S_\alpha = a_1 L \ln(L) + a_2 L + a_3 \frac{1}{L^{a_4}} \cos(a_5 L) \quad \text{form of } S \text{ for a metal, } d=2$$

[Klich-Gioev PRL '06, BGS PRL '09]



Example: “square root states”, $s=1$

$$H = \sum_r \left(-X_r + e^{-\beta J Z_r \sum_{r' \in \text{nn}(r)} Z_{r'}} \right)$$

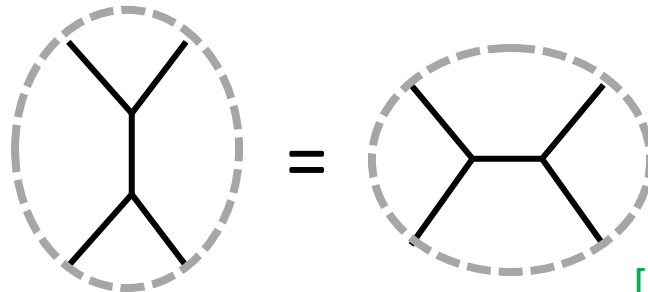
$$|\text{Ising}\rangle = \frac{1}{\sqrt{Z}} \sum_{\sigma} e^{\frac{\beta J}{2} \sum_{rr'} \sigma_r \sigma_{r'}} |\sigma\rangle$$

Hamiltonian is positive

$$H = \sum_r Q_r, \quad Q_r \geq 0$$

Use invariance of statistical partition function:

[Levin-Nave]



→ $s=1$

[BGS-Xu-McGreevy soon]

EXTENSIONS, COMMENTS, AND WRAP-UP

What is s ?

- Gapped systems: usually $s=0$ or $s=1$
- Gapless systems, scaling to a point: likely $s=1$
- Metals and NFLs: likely $s=2^{d-1}$

Is there a principled way of determining s ?

Intuition: lots of entanglement \rightarrow lots of low lying states

Scaling theory of critical states

$$\xi(T) \sim T^{-1/z} \quad \text{correlation length}$$

$$s(T) \sim \left(\frac{1}{\xi}\right)^{d-\theta} \sim T^{\frac{d-\theta}{z}} \quad \text{thermal entropy density}$$

$$\text{Key idea: } -\log(\rho_{A,\text{gs}}) \sim \sum_x \frac{H_x}{T(x)} \quad \text{half space gs density matrix (max ent)}$$

$$\begin{aligned} \theta < d - 1 &\rightarrow S_{EE} \sim \text{area} \\ \theta = d - 1 &\rightarrow S_{EE} \sim \text{area}^* \log \end{aligned}$$

$$s = 2^\theta$$

Evidence for generality of s sourcery

- Rigorous constructions
- Sufficient to capture correlations and entanglement entropy
- Some numerical evidence
- Even sufficient to capture effective Schmidt rank ...

Effective Schmidt rank

Q: What is the physical meaning of entanglement entropy?

A1: $nS(\rho)$ is the cost to compress $\rho^{\otimes n}$ in the limit of a large n number of copies

A2: For a single copy the answer is

$$H_{\max}^{\epsilon}(\rho) = \min_{\|\sigma - \rho\|_1 < \epsilon} \log(\text{rank}(\sigma))$$

$$H_{\max}^{\epsilon}(A, \text{CFT gs}) = S(A) \left(1 + \mathcal{O} \left(\sqrt{\frac{\log(\frac{1}{\epsilon})}{S(A)}} \right) \right)$$

Summary

- An RG inspired framework for the exact description of ground states – “efficient” calculation of physical properties, mounting evidence of generality
- Stay tuned for extensions to thermal states, etc. and results about CFTs, etc.
- A possibly hard case is FS+gauge field ...
- Foundation for holography via tensor networks?