

Entanglement in Strongly-Correlated Quantum Matter @ KITP, 2015

# Aspects of Holographic Quenches

## Part2: Locally Excited States

Tadashi Takayanagi

Yukawa Institute for Theoretical Physics (YITP), Kyoto University



This talk is a summary of our following works:

arXiv:1302.5703 [JHEP05(2013)080])

arXiv:1401.0539 [PRL 112(2014)111602]

arXiv:1403.0702 [PRD 90(2014)041701]

arXiv:1405.5946 [PTEP 2014 (2014) 9, 093B06]

arXiv:1410.2287 [JHEP01(2015)102]

(arXiv: 1503.08161)

## **Collaborators**

Pawel Caputa (Nordita) , Song He (YITP),

Joan Simon (Edinburgh), Andrius Stikonas (Edinburgh)

Masahiro Nozaki (YITP->Chicago),

Tokihiro Numasawa (YITP),

Kento Watanabe (YITP).

① Summary of Main Results (Just 3 pages)

(1-1) Our setup

**Take a locally excited state in a given (d+1) dim. CFT:**

$$|O(x)\rangle \equiv \underline{e^{-\varepsilon H}} \cdot \underline{O(x)} |0\rangle.$$

**UV regularization  
of local operator**

**(Note:  $\varepsilon \neq$  lattice spacing)**

**A primary state with dim.  $\Delta_O$**

$$\Rightarrow \text{Total energy} : \int T_{tt}(x) dx^d \approx \frac{\Delta_O}{\varepsilon}$$

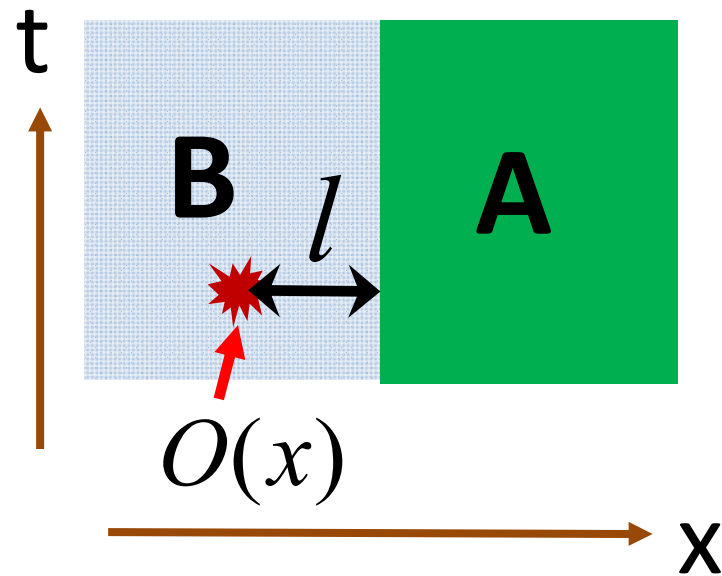
**Then we consider its time evolution:**

$$|O(x, t)\rangle = e^{-iHt} |O(x)\rangle.$$

(1-2) What to Compute

## The growth of (n-th Renyi) entanglement entropy

$$\Delta S_A^{(n)} \equiv S_A^{(n)} [ |O(x)\rangle ] - S_A^{(n)} [ |0\rangle ] .$$

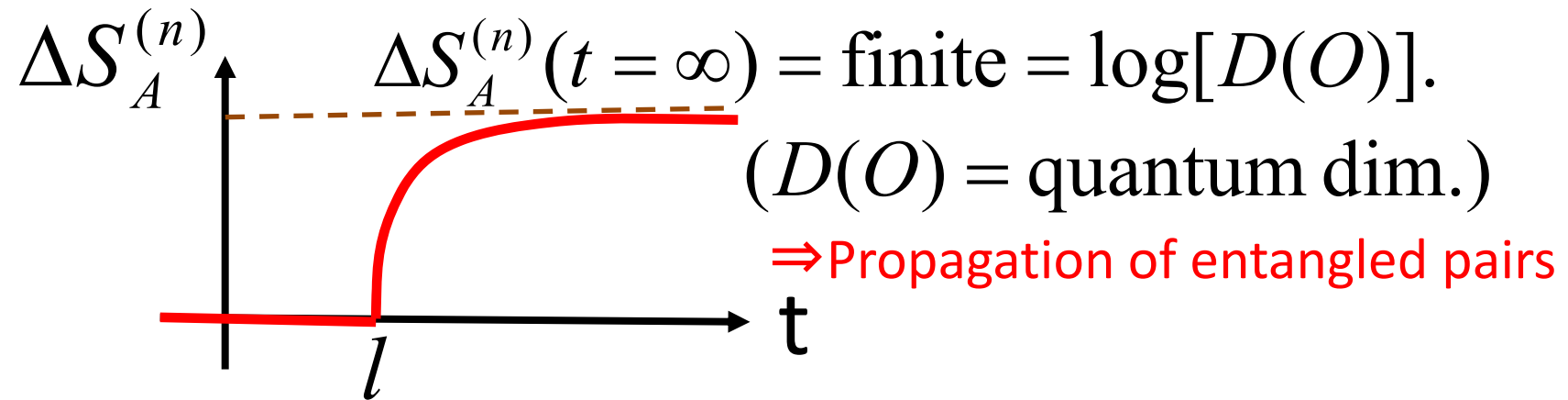


For simplicity, we choose  
**A = a half space** .

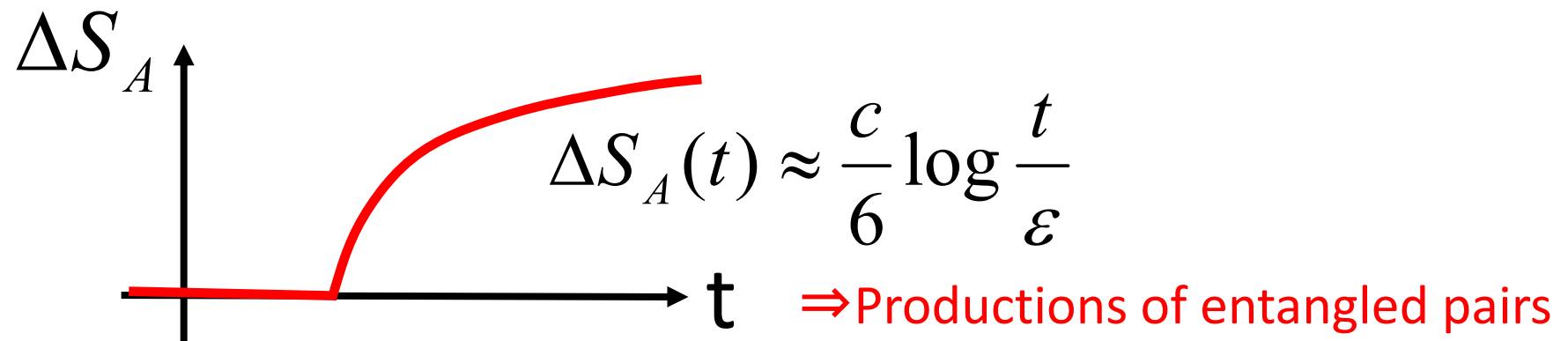
This calculation will show  
propagations and generations  
of quantum entanglement.

## (1-3) Summary of Main Results

### (i) Integrable CFTs [Massless Free Fields, Minimal Models etc.]



### (ii) Holographic CFTs [AdS3/CFT2]

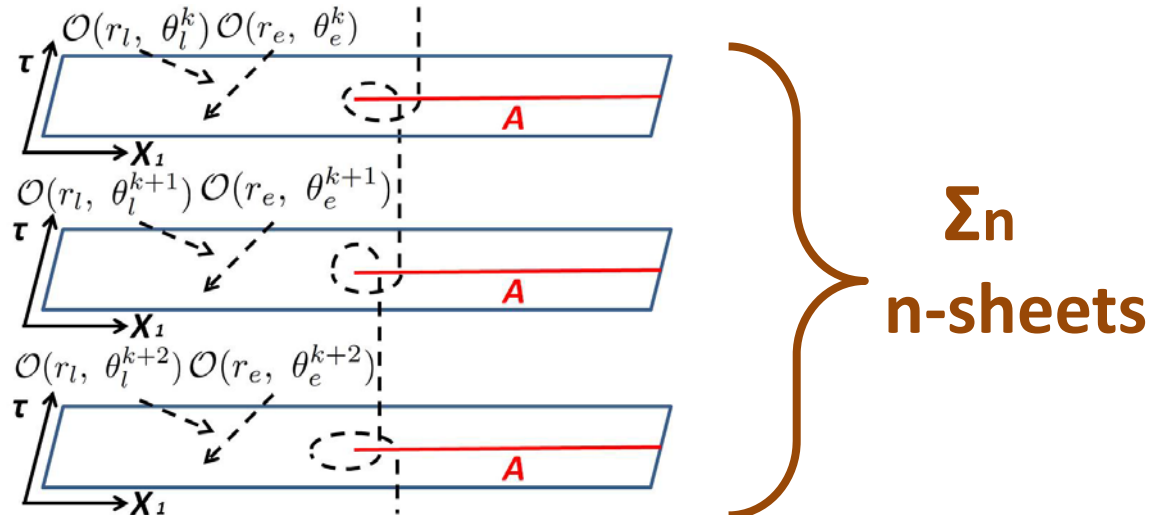


## ② Free Field Theory Calculations [Nozaki-Numasawa-TT 14]

### (2-1) Replica formulation

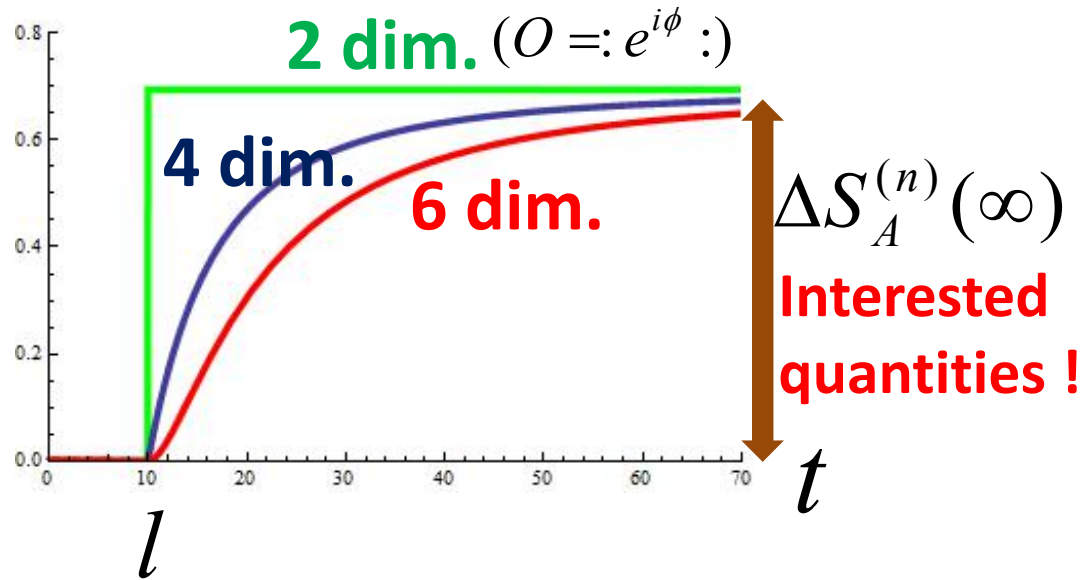
The n-th Renyi EE can be expressed in terms of 2n-point correlation functions on  $\Sigma_n$  :

$$\Delta S_A^{(n)} = \frac{1}{1-n} \cdot \left[ \log \left\langle O(r_l, \theta_l^n) O(r_e, \theta_e^n) \cdots O(r_l, \theta_l^1) O(r_e, \theta_e^1) \right\rangle_{\Sigma_n} - n \cdot \log \left\langle O(r_l, \theta_l) O(r_e, \theta_e) \right\rangle_{\Sigma_1} \right].$$



## (2-2) Results in free massless scalar theory

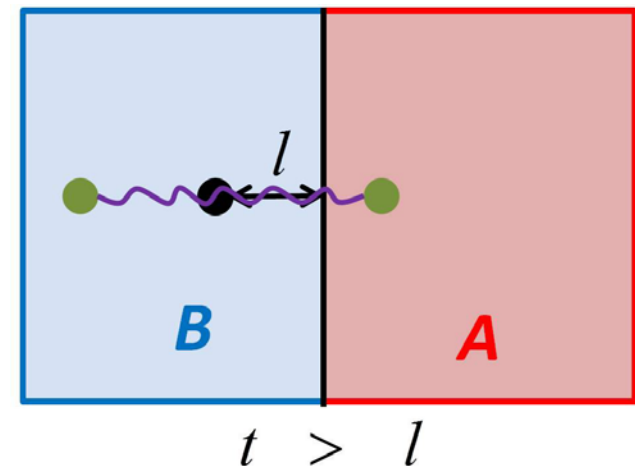
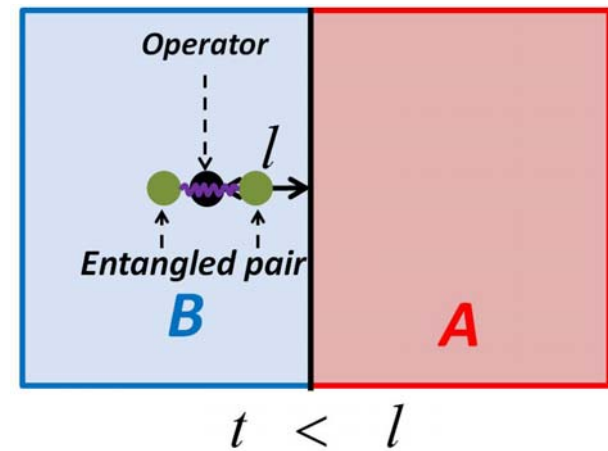
$\Delta S_A^{(2)}$  for  $O =: \phi :$  (i.e.  $k = 1$ )



E.g.  $\Delta S_{A(4\text{dim})}^{(2)} = \log\left(\frac{2t^2}{t^2 + l^2}\right)$ .

Note:

$\Delta S_A^{(n)f}$  is 'topologically invariant'  
under deformations of A.



$$\Delta S_A^{(n)f} \text{ for } O = \phi^k \text{ in } d+1 > 2 \text{ dim.}$$

TABLE I.  $\Delta S_A^{(n)f}$  and  $\Delta S_A^f (= \Delta S_A^{(1)f})$  for free massless scalar field theories in dimensions higher than two ( $d > 1$ ).

	$n$	$k = 1$	$k = 2$	$\dots$	$k = l$
$\Delta S_A^{(n)f}$ Renyi Entropy	2	$\log 2$	$\log \frac{8}{3}$	$\dots$	$-\log \left( \frac{1}{2^{2l}} \sum_{j=0}^l ({}^l C_j)^2 \right)$
	3	$\log 2$	$\frac{1}{2} \log \frac{32}{5}$	$\dots$	$-\frac{1}{2} \log \left( \frac{1}{2^{3l}} \sum_{j=0}^l ({}^l C_j)^3 \right)$
	$\vdots$	$\vdots$	$\vdots$	$\vdots$	$\vdots$
	$m$	$\log 2$	$\frac{1}{m-1} \log \frac{2^{2m-1}}{2^{m-1}+1}$	$\dots$	$\frac{1}{1-m} \log \left( \frac{1}{2^{ml}} \sum_{j=0}^l ({}^l C_j)^m \right)$
$\Delta S_A^f$	1	$\log 2$	$\frac{3}{2} \log 2$	$\dots$	$l \log 2 - \frac{1}{2^l} \sum_{j=0}^l {}^l C_j \log {}^l C_j$

von-Numann EE

**EPR state !**

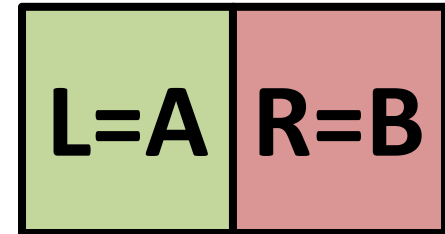
[For a proof: Nozaki, arXiv:1405.58754]



## Heuristic Explanation

First , notice that in free CFTs, there are definite **particles moving at the speed of light.**

$$\Rightarrow \phi \approx \underbrace{\phi_L}_{\text{left-moving}} + \underbrace{\phi_R}_{\text{right-moving}} \cdot$$



$$\begin{aligned} \phi^k |\text{vac}\rangle &\approx \sum_{j=0}^k \binom{k}{j} C_j \cdot (\phi_L)^j \cdot (\phi_R)^{k-j} |\text{vac}\rangle \\ &= 2^{-k/2} \sum_{j=0}^k \sqrt{\binom{k}{j} C_j} |j\rangle_L |k-j\rangle_R. \end{aligned}$$

$$\Rightarrow \Delta S_A^{(n)f} = \frac{1}{1-n} \log \left[ 2^{-nk} \sum_{j=0}^k \binom{k}{j} C_j^n \right]$$

$$\Delta S_A^f = k \log 2 - 2^{-k} \sum_{j=0}^k \binom{k}{j} C_j \cdot \log[\binom{k}{j} C_j].$$

Agree with  
replica  
Calculations !

### ③ Rational 2d CFTs [He-Numasawa-Watanabe-TT 14]

#### (3-1) Free Scalar CFT in 2d

Consider following two operators in the free scalar CFT:

$$(i) \quad O_1 = e^{i\alpha\phi}, \quad \Rightarrow \quad \Delta S_A^{(n)f} = 0.$$

$$|O_1\rangle = e^{i\alpha\phi_L} |0\rangle_L \otimes e^{i\alpha\phi_R} |0\rangle_R \Rightarrow \text{Direct product state}$$


$$(ii) \quad O_2 = e^{i\alpha\phi} + e^{-i\alpha\phi}, \quad \Rightarrow \quad \Delta S_A^{(n)f} = \log 2.$$

$$\begin{aligned} |O_2\rangle &= e^{i\alpha\phi_L} |0\rangle_L \otimes e^{i\alpha\phi_R} |0\rangle_R + e^{-i\alpha\phi_L} |0\rangle_L \otimes e^{-i\alpha\phi_R} |0\rangle_R \\ &\approx |\uparrow\rangle_L |\uparrow\rangle_R + |\downarrow\rangle_L |\downarrow\rangle_R \Rightarrow \text{EPR state} \end{aligned}$$

(3-2) Rational 2d CFTs (e.g. minimal models, WZW models)

2<sup>nd</sup> Renyi EE  $\Rightarrow$  **4 pt. function:**  $\langle O(\infty)O(1)O(z)O(0) \rangle$ .

(i) **Early time** ( $0 < t < l$ ):  $(z, \bar{z}) \rightarrow (0,0)$ .

(ii) **Late time** ( $t \geq l$ ):  $(z, \bar{z}) \rightarrow (1,0)$ .  **Chiral Fusion Transformation**  
 $z \rightarrow 1-z$

This allows us to prove  $\Delta S_A^{(n)} = \log \underline{D(O)}$  for any n.  
**quantum dim.**

Ex. Ising model :  $\Delta S_A^{(n)} [I] = \Delta S_A^{(n)} [\varepsilon] = 0,$

$$\Delta S_A^{(n)} [\sigma] = \log \sqrt{2}.$$

## ④ Holographic Analysis

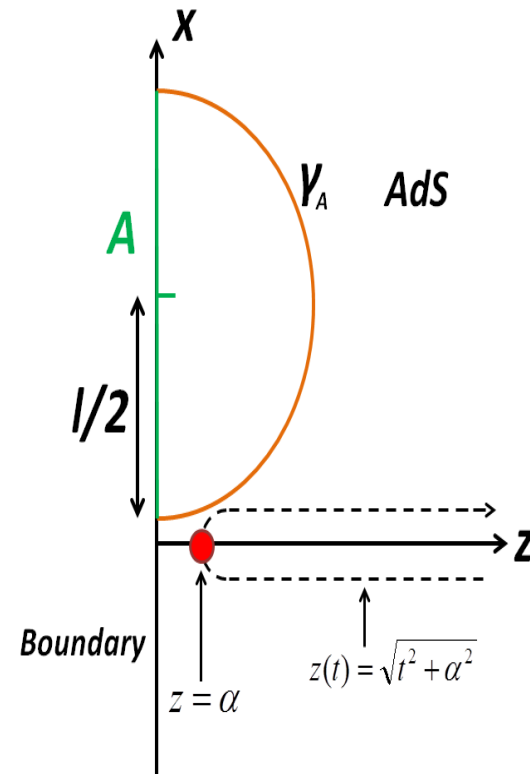
A locally excited state

~ A falling particle in AdS.

$$\Delta_o \approx mR$$

[e.g. stress tensors agree with the CFT.]

We can find an analytical metric using the Horowitz-Itzhaki map.

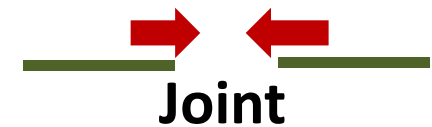


$$\Delta S_A \approx \frac{c}{6} \log \left( \frac{t}{\varepsilon} \right).$$

[Holographic Calculations: Numasawa-Nozaki-TT 13,  
Caputa-Nozaki-TT 14]

[Large c CFT computations: Asplund-Bernamonti-  
Galli-Hartman 14]

cf.  $\Delta S_A \approx \frac{c}{3} \log t$  , for local quenches



in the sense of Calabrese-Cardy 2007. [Holographic Calculation: Ugajin 13]

⑤ Finite Temperature Analysis [Caputa-Simon-Stikonas-TT 14]

Consider a local operator excited state at finite temp.

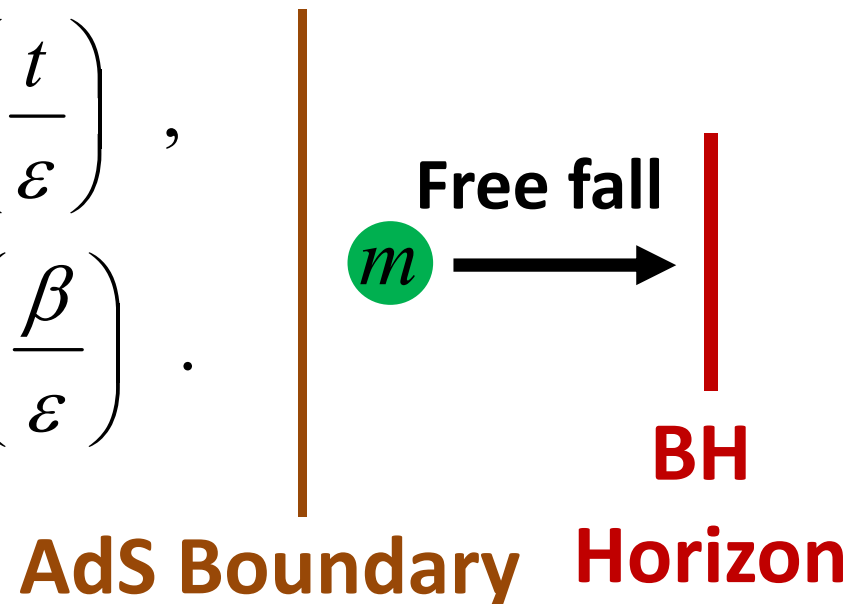
(i) Integrable CFTs  $\Rightarrow$  Only  $O(\epsilon/\beta)$  differences.

(ii) Holographic CFTs (=AdS Blackholes)

$\Rightarrow$  EE saturates at late time !

$$0 < t \ll \beta : \quad \Delta S_A^{(1)} \approx \frac{c}{6} \log\left(\frac{t}{\epsilon}\right) ,$$

$$t \gg \beta : \quad \Delta S_A^{(1)} \approx \frac{c}{6} \log\left(\frac{\beta}{\epsilon}\right) .$$



# A Sketch of Time Evolution

