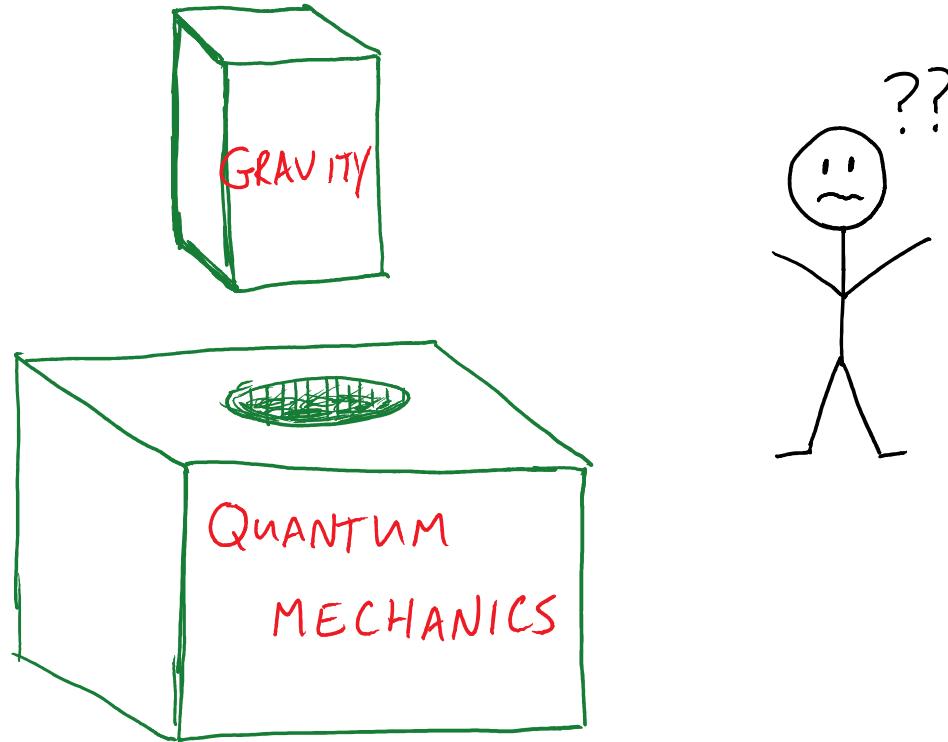


# GRAVITY AND ENTANGLEMENT

Mark Van Raamsdonk, UBC

KITP, April 2015

One of the greatest challenges for  
theoretical physics:



**GRAVITY** describes the dynamics of spacetime geometry and its interaction with matter.

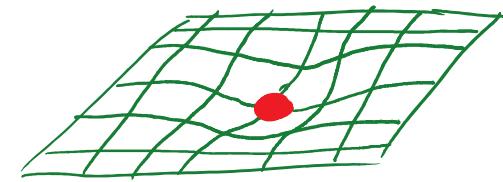
Central result:

EINSTEIN'S

EQUATION:-

$$G_{\mu\nu} = 8\pi G_N T_{\mu\nu}$$

measures  
intrinsic curvature  
of spacetime



measures energy &  
momentum density & flow

Fitting this into the framework of quantum mechanics has been very difficult!

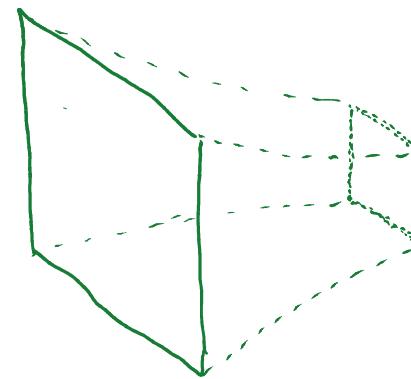
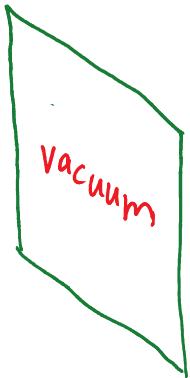
1997 : Remarkable progress via AdS/CFT  
correspondence in string theory (Maldacena)

QUANTUM  
GRAVITY  
(certain  
examples)

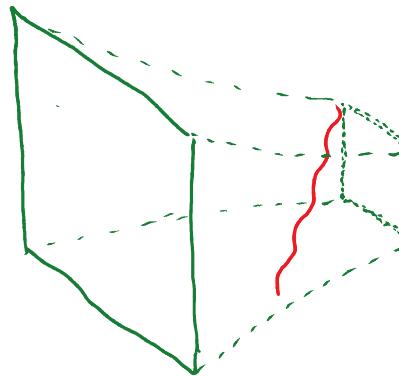
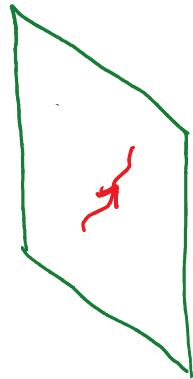
=  
exactly  
equivalent

ORDINARY  
QUANTUM  
SYSTEM  
(e.g. Quantum Field Th.  
on fixed spacetime)

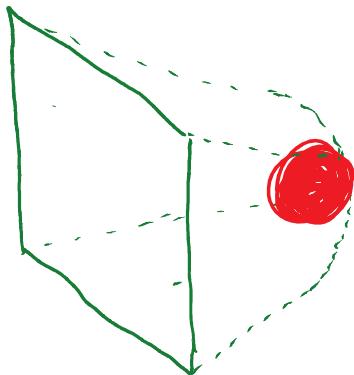
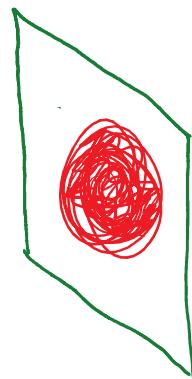
Different QFT states  $\longleftrightarrow$  Different spacetimes



empty  
spacetime



gravity  
wave



black  
hole

BIG QUESTION:

How/why do spacetime/gravity  
emerge from QFT physics?

Which states  $| \psi \rangle$  correspond to classical  
geometry?

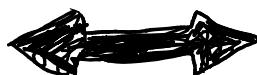
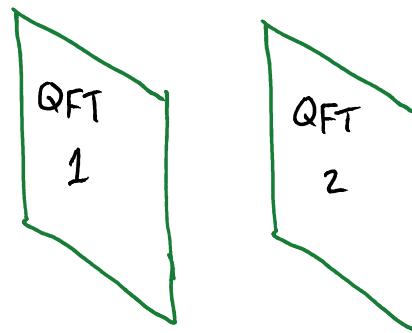
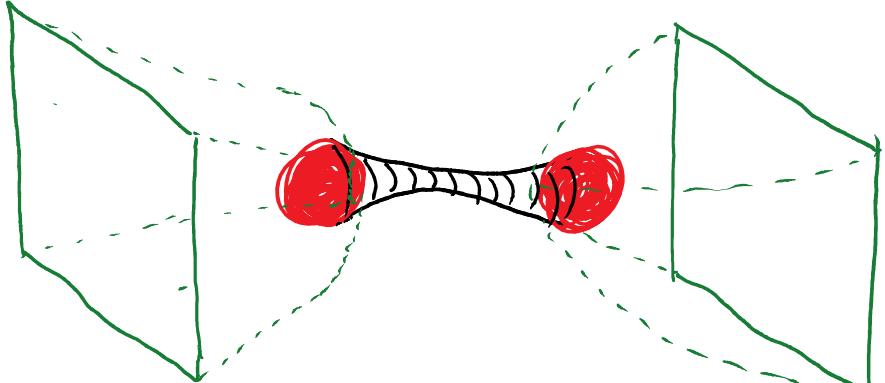
What can we learn about gravity?

Recent work: physics of entanglement + quantum  
information is crucial!

Maldacena 2001:

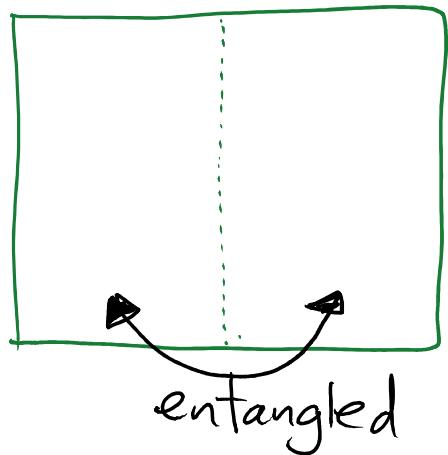
2 separate spacetimes  
connected by wormhole

2 separate QFTs  
entangled with  
one another



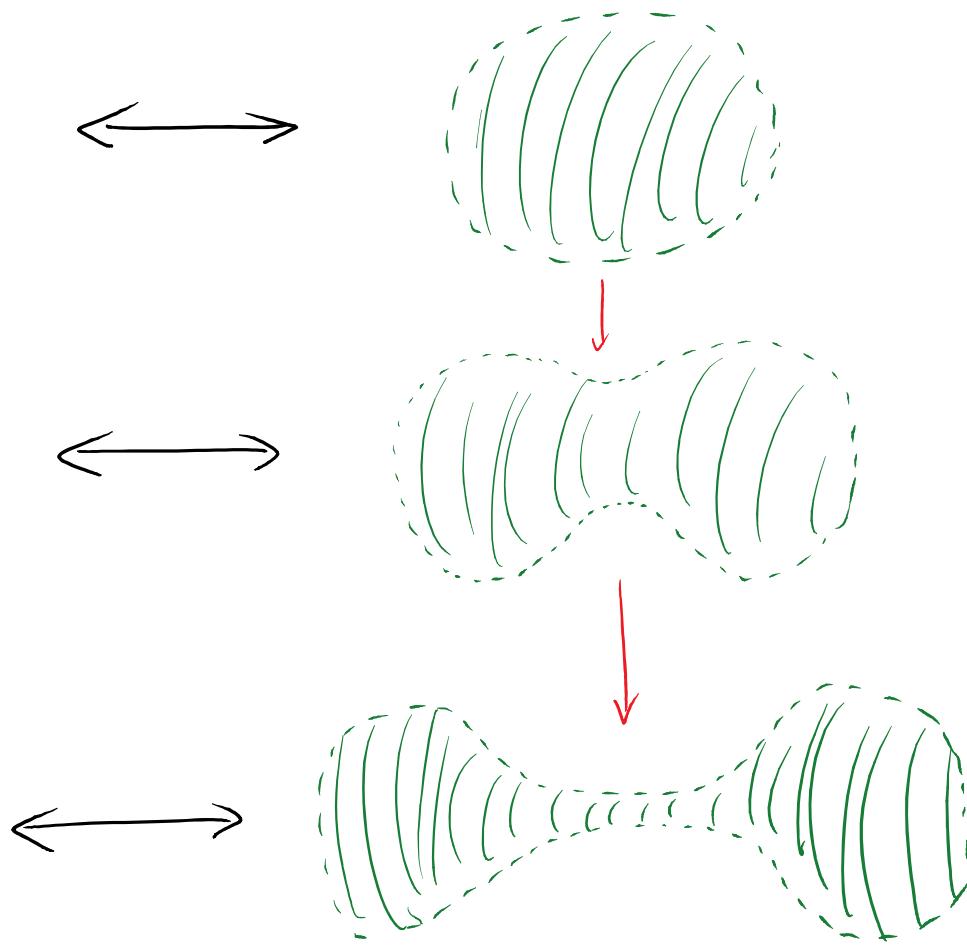
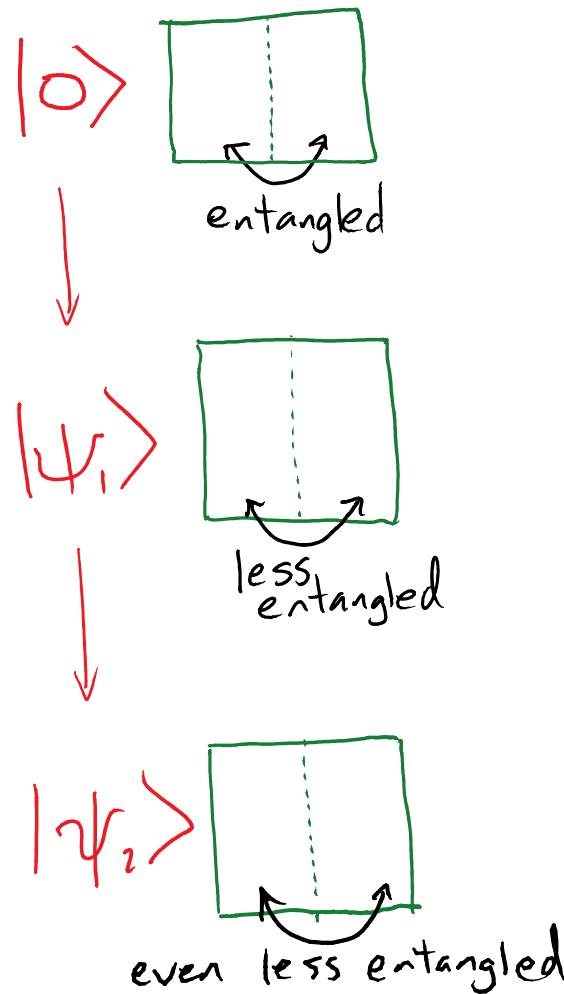
$$|\psi\rangle = |\text{green box}_1 \text{ red box}_2\rangle + |\text{red box}_1 \text{ green box}_2\rangle + |\text{green box}_1 \text{ red wavy box}_2\rangle + |\text{red box}_1 \text{ green wavy box}_2\rangle + \dots$$

BUT: lots of entanglement already in QFT ground state  
dual to empty spacetime.



QFT fields in any region entangled with  
fields outside

What happens if we remove this entanglement?



2 regions of space pinch off from each other!

MIR, MNR, Czech, Noguera, Karzmarek

Suggests that classical spacetime geometry  
emerges via entanglement of degrees of freedom  
in dual QFT!

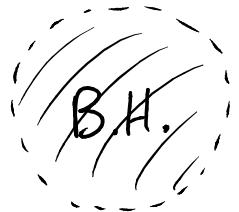
MVR, Swingle 2009

No classical spacetime without quantum entanglement.

Can we be more quantitative?

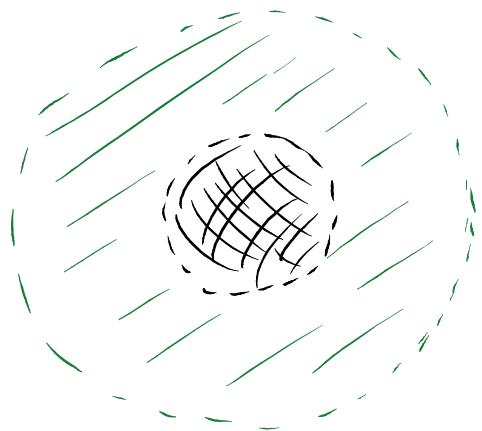
First hints: black holes + thermodynamics

1970s



$\frac{\text{Area}}{4G_N}$  behaves like entropy.

AdS/CFT:



in thermal ensemble

$$\{|E_i\rangle, p_i = \frac{1}{Z} e^{-\beta E_i}\}$$

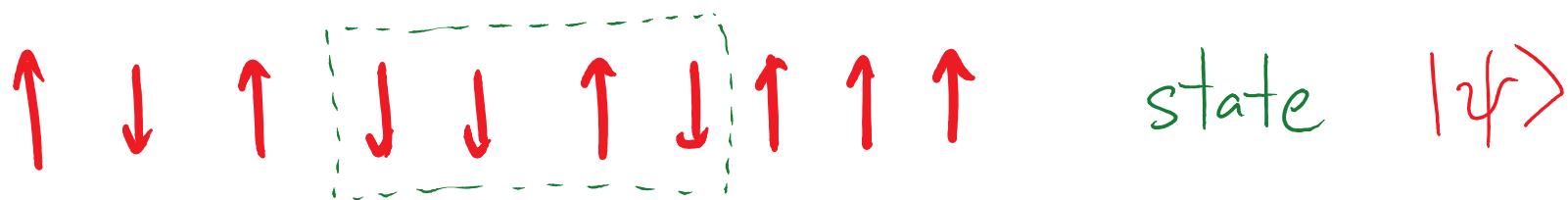
$$\frac{\text{Area}}{4G_N} = S_{\text{CFT}} = -\sum_i p_i \log p_i$$

Key point (Ryu + Takayanagi)

Entropy of subsystems also  
has a geometric interpretation.

# Entropy of subsystems (= Entanglement Entropy)

Consider any quantum system w. subsystem A



state  $|\psi\rangle$   
subsystem A → described by ENSEMBLE  $\{P_i, |\psi_i^A\rangle\}$

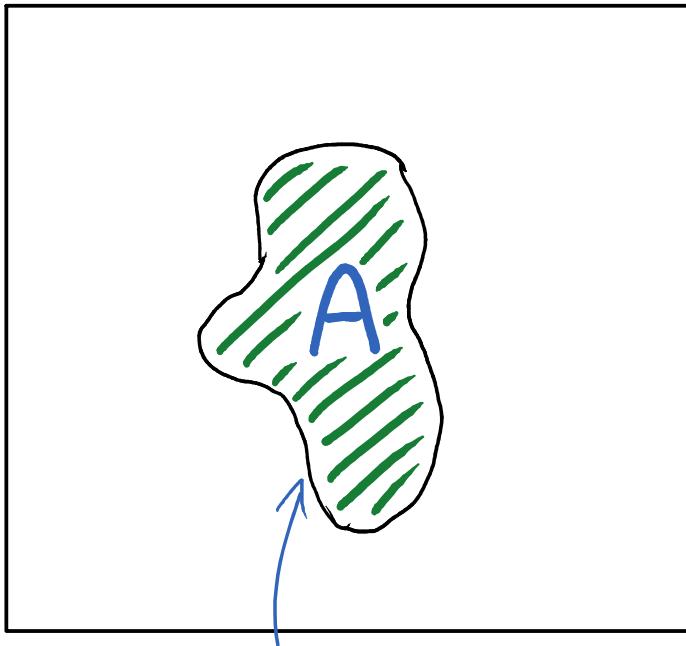
$\{P_i\} \neq \{1\}$  → A is ENTANGLED w. rest  
↓ classical uncertainty about state of A

Quantify via entropy:

$$S(A) = - \sum_i P_i \log P_i$$

# QFT entanglement entropy:

CFT state  $| \psi \rangle$



general region

$S(A)$  divergent but can define sensible quantities:

$$S_{|\psi\rangle}(A) - S_{|vac\rangle}(A)$$

$$S(A) + S(B) - S(A \cup B)$$

$$\equiv I(A, B)$$

↑ mutual information

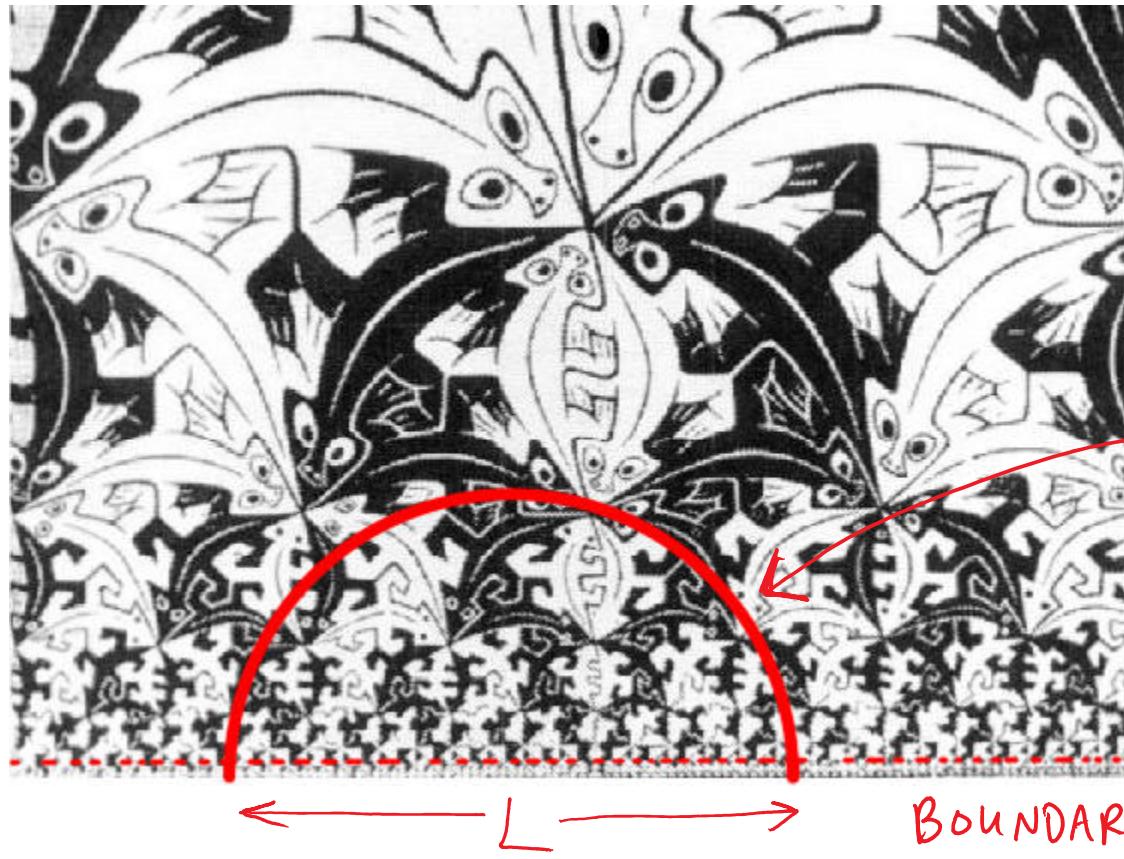
Example:

$$S = \frac{c}{3} \log \left( \frac{L}{\epsilon} \right)$$

"central charge" = measure of  
number of d.o.f.

for ANY 1+1D conformal field theory.

Ryu + Takayanagi: this result can be represented geometrically.



spatial slice of  
anti-de Sitter  
space

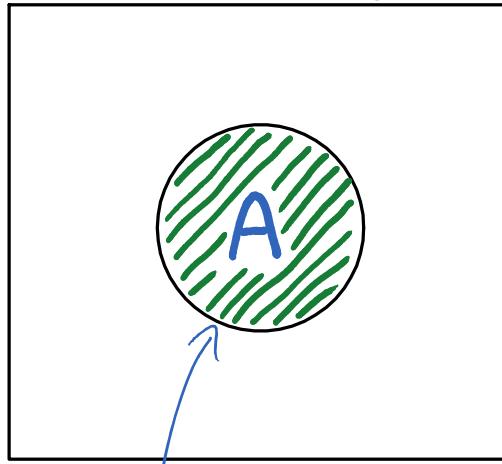
shortest length  
curve that has  
same boundary  
as interval  $L$

$$\text{length} = \frac{\epsilon}{3} \ln\left(\frac{L}{\epsilon}\right)$$

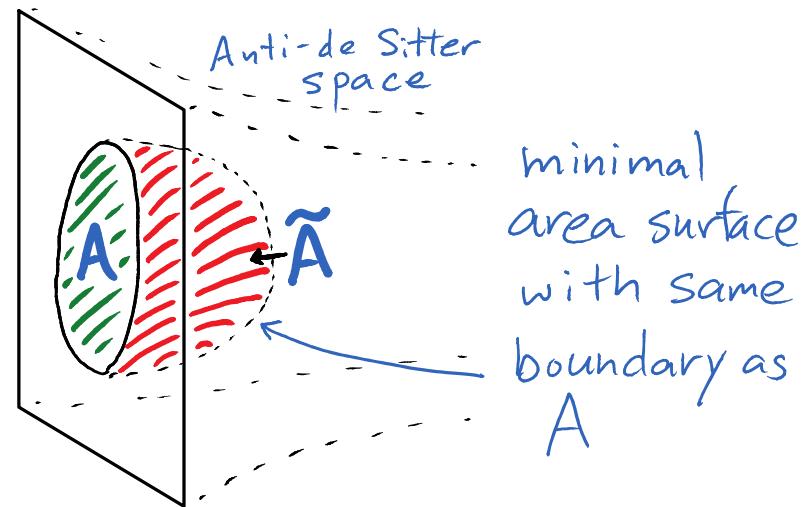
Higher dimensions:

for any CFT:

vacuum state



ball shaped region



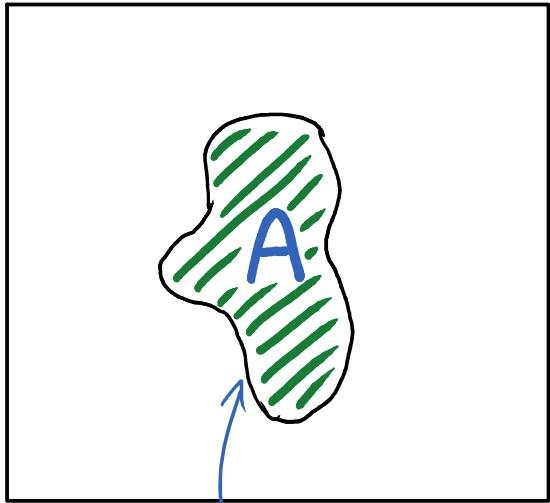
$$S(A) \propto \text{Area}(\tilde{A})$$

cf.  $S = \frac{\text{Area}}{4G_N}$  for  
black holes!

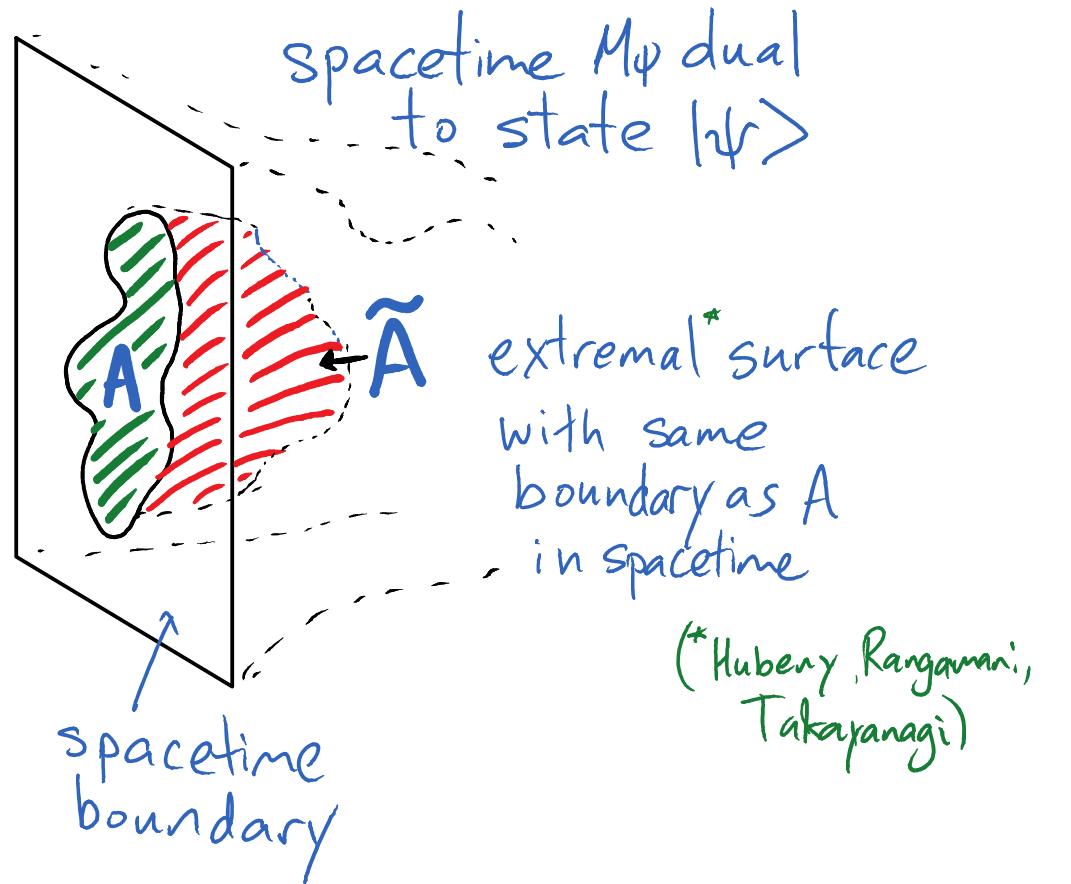
Ryu-Takayanagi: conjectured this to hold for general regions  $A$  + general states  $|\psi\rangle$  in "holographic" CFTs

# GEOMETRY FROM ENTANGLEMENT

CFT state  $| \Psi \rangle$



general region



spacetime  $M_4$  dual  
to state  $| \Psi \rangle$

extremal\* surface  
with same  
boundary as A  
in spacetime

(\*Hubeny, Rangamani,  
Takayanagi)

$| \Psi \rangle \rightarrow$  calculate  
 $S_A$  for  
many A  $\rightarrow$  find  $M_4$  s.t.

$$\text{Area}(\tilde{A}) = S_A$$

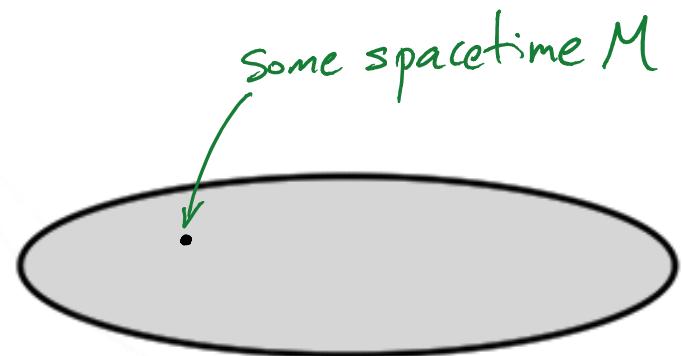
Can (plausibly) reconstruct geometry from entanglement!

Question: Couldn't I do this for  
any state in any field theory?

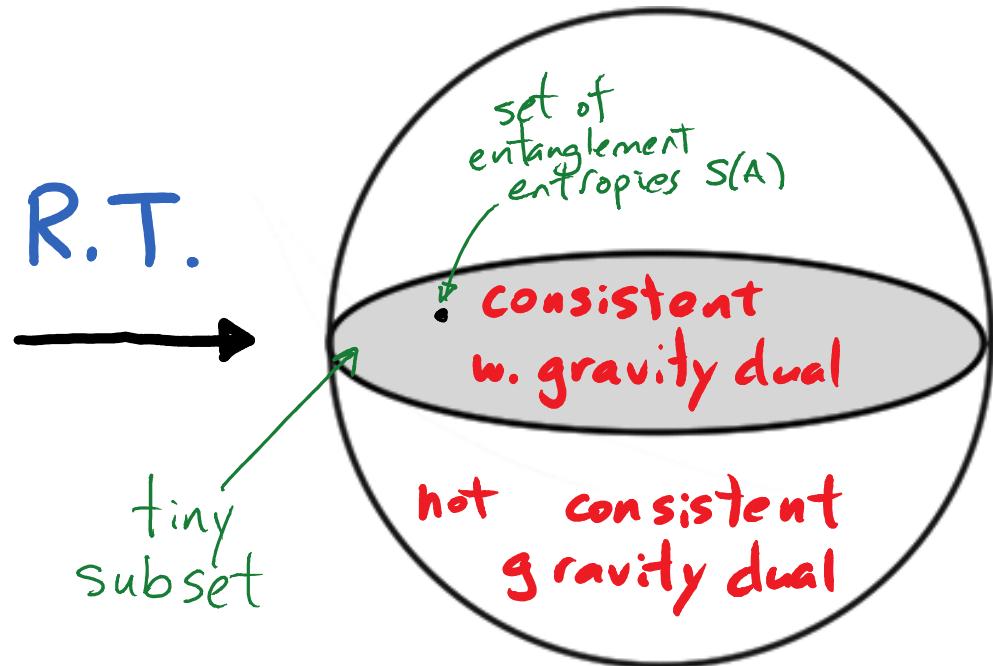
Question: Couldn't I do this for  
any state in any field theory?

A: No! Usually no  $M$  will reproduce  
entanglement structure  $S(A)$

Entanglement structure of states w.  
geometrical dual is highly constrained

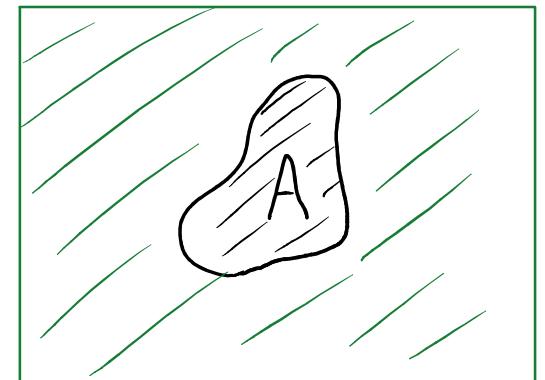


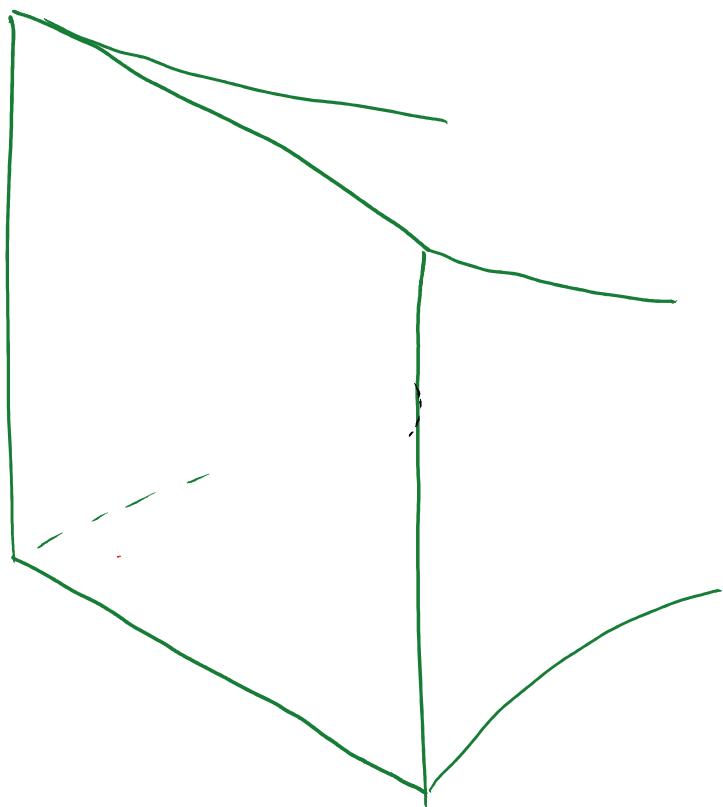
Asymptotically  
AdS spacetimes



functions  $S(A)$

||  
potential entanglement  
structures





Example: states dual to  
vacuum solutions of Einstein's  
equations

$S(A)$  for infinitesimal  
regions (equivalent to  $\langle T_{\mu\nu} \rangle$ )

↓ R.T.

asymptotic metric

↓ Einst. Eqn.

interior metric (to some distance)

↓ R.T.

$S(A)$  for any region  $A$  (not too large)

Entanglement structure determined by local data!

## General Questions:

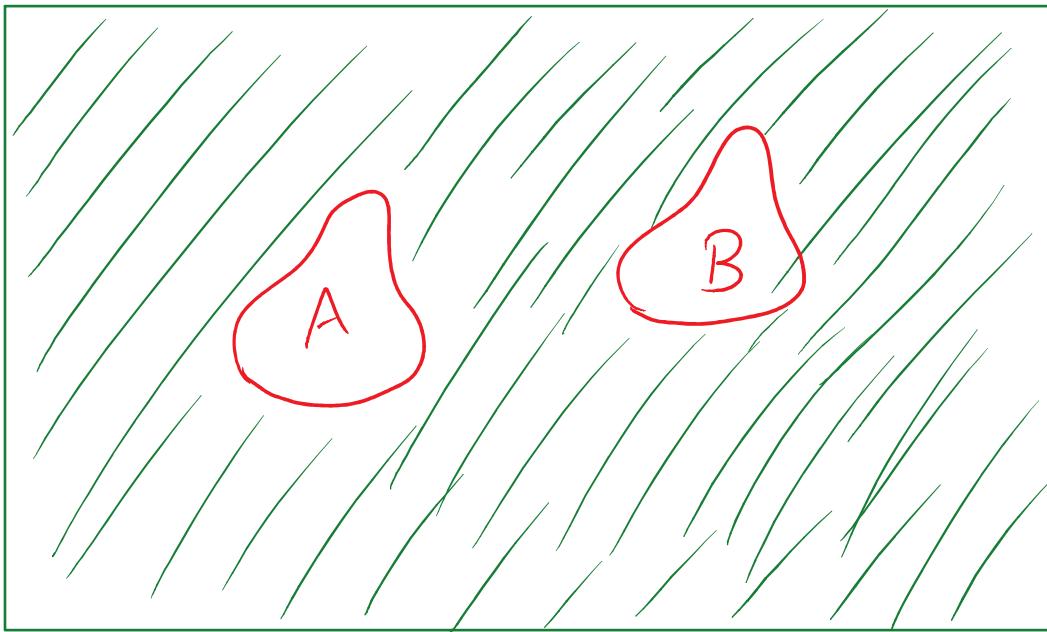
- \* Which entanglement structures are consistent w. geometrical description? \*
- \* Why do holographic Hamiltonians prefer states w. this structure? \*

Can we learn anything about gravity?

Can we learn anything about gravity?

Yes! Constraints on entanglement structure  
restrict possible spacetimes.

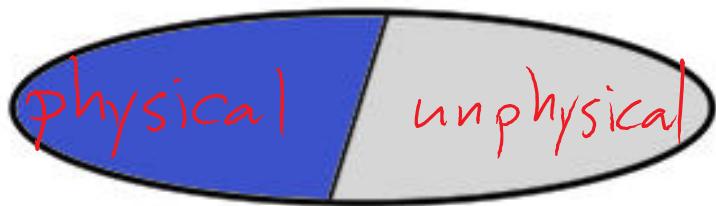
\* Not all functions  $S(A)$  represent consistent entanglement structures \*



e.g. must have:

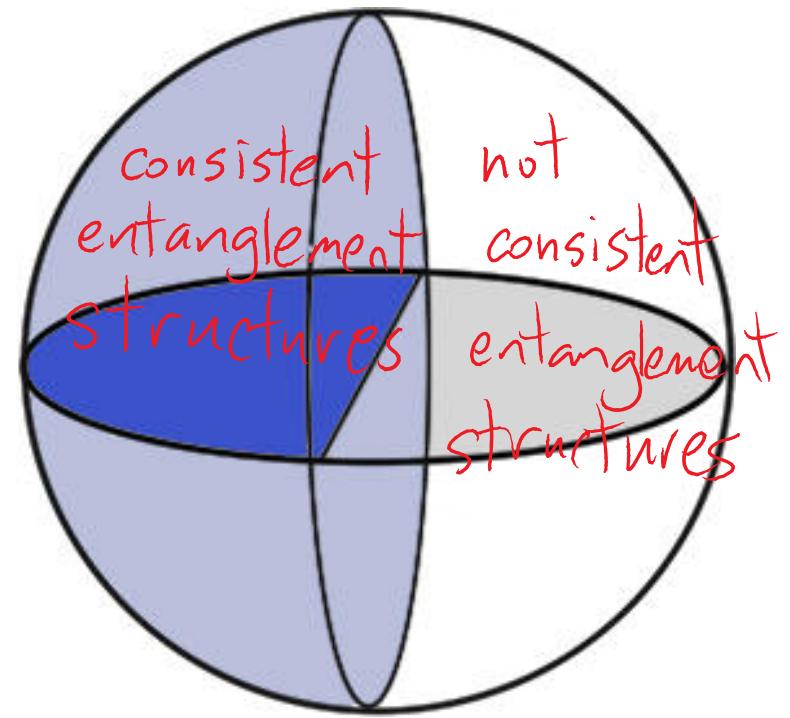
$$S(A) + S(B) - S(A \cup B) \geq 0$$

Question : which geometries give rise  
to consistent entanglement structures?



a sympt. AdS geometries

R.T.  
→



functions  $S(A)$

# ENTANGLEMENT INEQUALITIES

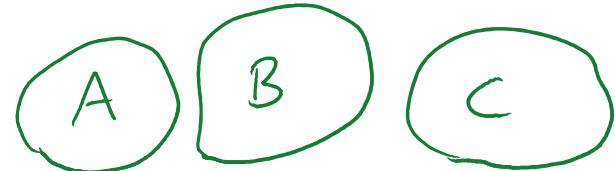
SUBADDITIVITY:

$$I(A,B) = S(A) + S(B) - S(A \cup B) \geq 0$$

MUTUAL INFORMATION: measure of entanglement/correlations between  $A \circ B$

STRONG SUBADDITIVITY:

$$I(ABC) \geq I(BC)$$



$$S(AB) + S(BC) \geq S(ABC) + S(B)$$

# POSITIVITY + MONOTONICITY OF RELATIVE ENTROPY.

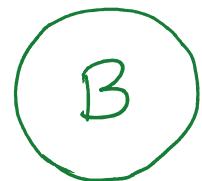
$$S(\rho \parallel \sigma) = \text{tr}(\rho \log \rho) - \text{tr}(\rho \log \sigma)$$

- measure of distinguishability of  $\rho, \sigma$
- positive, larger for  $\rho, \sigma$  corresponding to larger subsystems.

$\rho$ : density matrix for ball B in state  $| \psi \rangle$

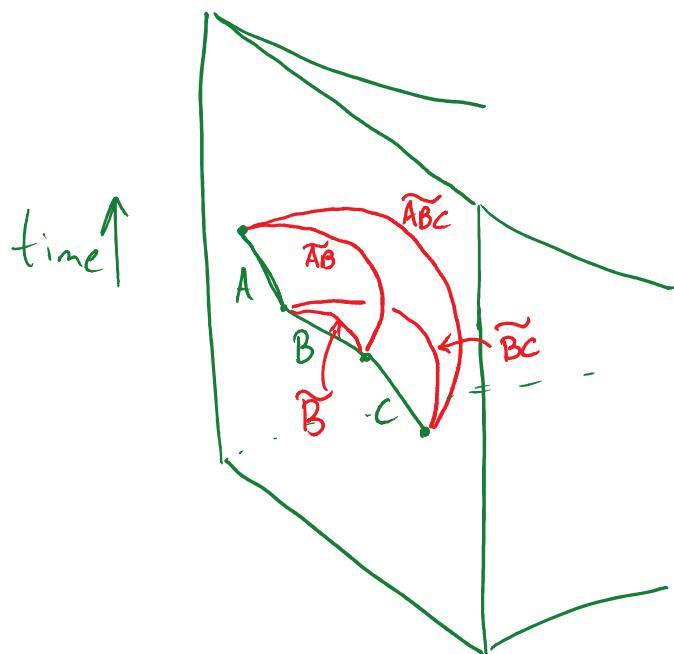
$\sigma$ : density matrix for ball B in vacuum

$$S(\rho \parallel \sigma) = \Delta \langle H_\sigma \rangle - \Delta S$$



$$H_\sigma = -\ln \sigma = 2\pi \int \frac{R^2 - r^2}{2R} \cdot T_{oo}$$

Using R.T., translate these  
to statements about geometry



e.g.

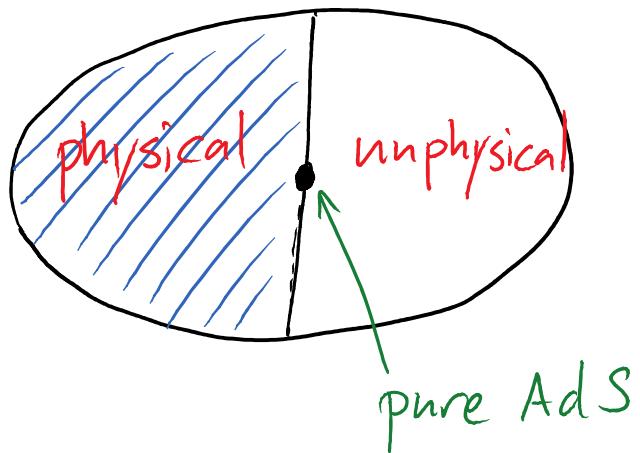
Strong subadditivity



$$\text{Area}(\tilde{AB}) + \text{Area}(\tilde{BC})$$

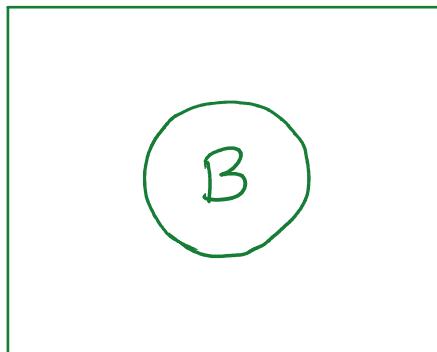
$$\geq \text{Area}(\tilde{ABC}) + \text{Area}(\tilde{B})$$

Not true for all geometries



Start w. geometries

$M$  near pure AdS.  
 $(|\psi\rangle = |\text{vac}\rangle + \delta|\psi\rangle)$

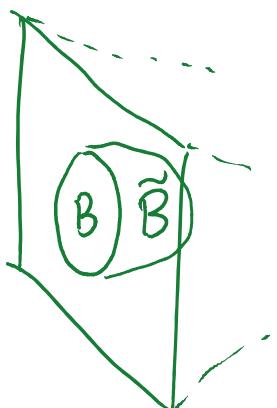


Positivity of relative entropy

$$\Rightarrow \delta S_B = 2\pi \int_B d^4x \frac{R^2 - r^2}{2R} \delta \langle T_{00} \rangle$$

Blanco  
Casini  
Hung  
Myers

True for all balls in all  
Lorentz frames.



Translates to constraint on  
metric perturbation.

Result:

\* Physical perturbations to pure AdS must satisfy Einstein's Equations to linear order \*

- Lashkari, McDermott, MVR

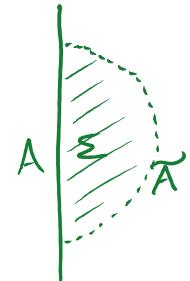
- Faulkner, Guica, Hartman, Myers, MVR

- Swingle, MVR

More recently:

w Swingle: Refined Ryu-Takayanagi formula:  $S_{CFT} = \frac{\text{Area}}{4G_N} + S_\Sigma$   
 $\Downarrow SS = 8E$

$\delta\langle T_{\mu\nu}\rangle$  is source for linearized Einst. Eqns.



w. Lashkari

Rabideau

Sabella-Garnier

also: Rangamani, Takayanagi  
Hubeny, Bhattacharya

Ooguri, Lin, Marcoli, Stoica

Constraints at non-linear level

(strong subadditivity + positivity/monotonicity  
of relative entropy)

give rise to energy conditions!

(GRAVITY PROGRAM TALK NEXT WEEK)

# SUMMARY

structure of entanglement  
in QFT



spacetime  
geometry

constraints on QFT  
entanglement



Einstein Equations  
(so far: linearized)  
+ Energy Conditions

