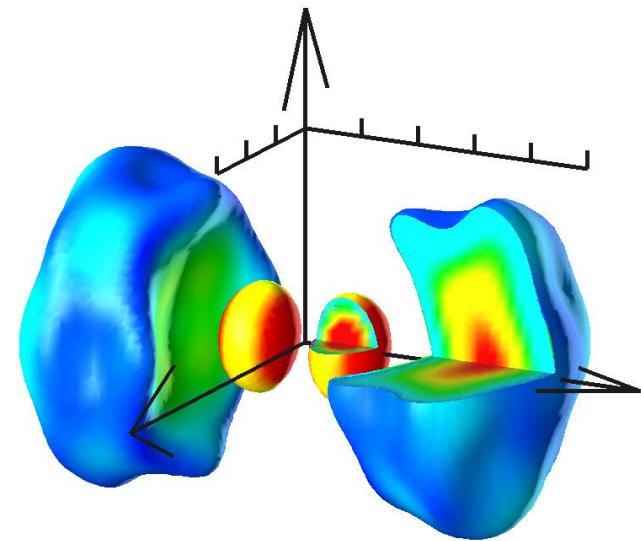
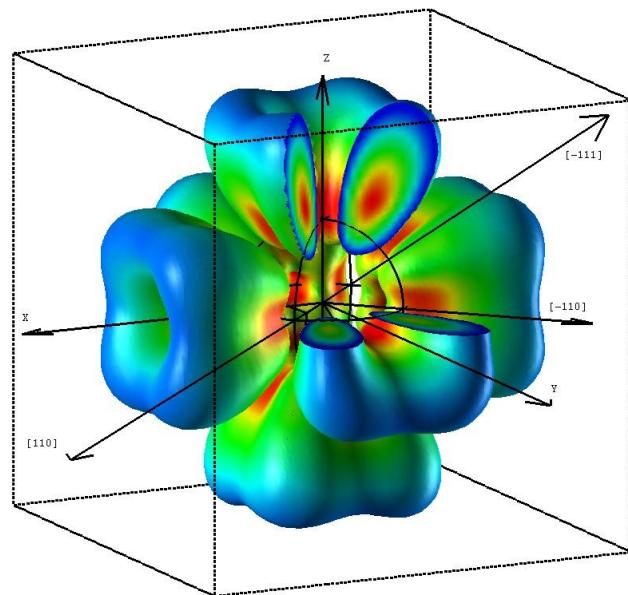


Simplifying Excitations in Charge-Transfer Insulators



Wei Ku

Brookhaven National Lab & SUNY Stony Brook



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- DOE-CMSN
- Taiwanese Student Fellowship

Theoretical works

- Chi-Cheng Lee (BNL & Tamkang Univ.) ← about to move on
- Chen-Lin Yeh (BNL & Tamkang Univ.)
- Hung-Chung Hsueh (Tamkang Univ.)

Experiments

- Ben C. Larson (ORNL) – PRL 2007
- Igor Zaliznyak (BNL) – Nature Physics 2009
- Peter Abbamonte (UIUC) – PNAS 2008

Discussions

- George Sawatzky (UBC)
- Adolfo G. Eguiluz (UT-Knoxville / ORNL)



Outline

- Recent X-ray experiments
 - Short wave length → sensitivity to local excitations
 - Strong anisotropy → sensitivity to orbitals orientation
- Strong local interaction & short-range correlation
 - Local approach
- Fundamental difficulty of charge-transfer insulators
 - Defining “local” in a crystal with “natural” symmetry
- Simple understanding of anisotropy from local picture
- 70% missing spectral weight in INS
- Treating local problem via TDDFT: TD-LDA+ U
- Treating local problem beyond perturbation
- Propagation of local excitations



For the Purists & the Rest

- In a many-body system, G_1, G_2, \dots, G_N not apparently related
 - Band structure (or DOS) is not sufficient
 - At least linear response is necessary
- TD-DFT: time evolution of density is legally accessible



Non-Resonant Inelastic X-Ray Scattering (NIXS)

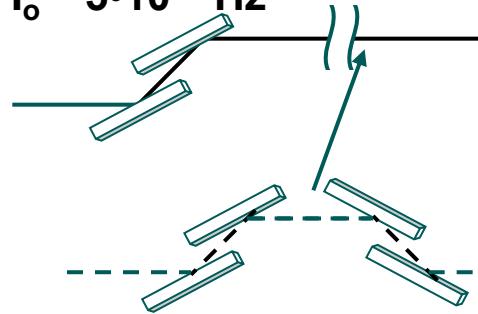
Spherically Bent
Analyzer Crystal

UNI-CAT ID-33

7.59 keV

$\Delta E \sim 1.1$ eV

$I_0 \sim 5 \cdot 10^{12}$ Hz

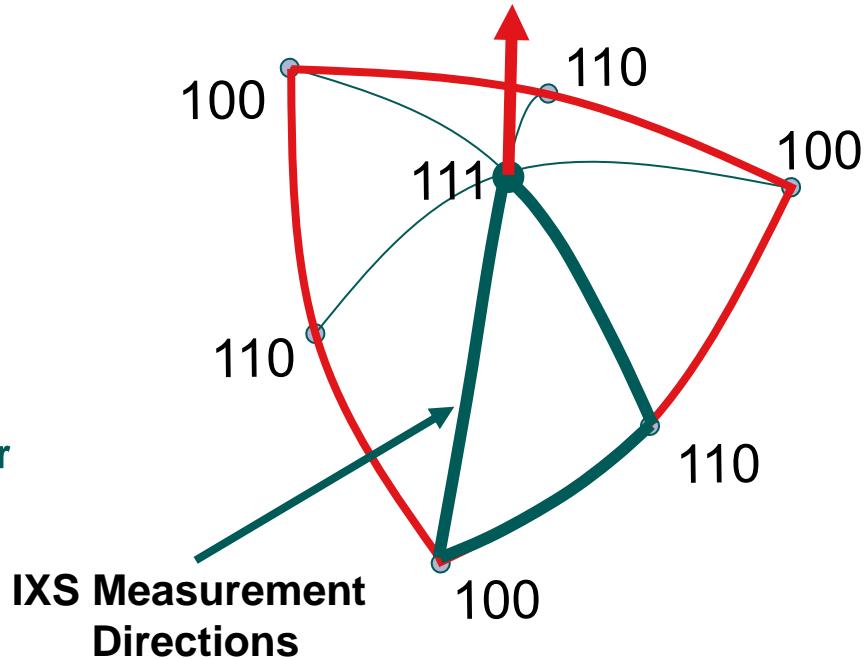


UNI-CAT ID-33

CHESS C-Line

$\Delta E \sim 0.3$ eV

$I_0 \sim 10^{11}$ Hz



$$s(\vec{q}, \omega) = \frac{1}{V} S(\vec{q}, \omega) = -2\hbar \text{Im} \chi_{\vec{G}\vec{G}'}(\vec{q} - \vec{G}, \omega)$$

Absolute IXS Measurements

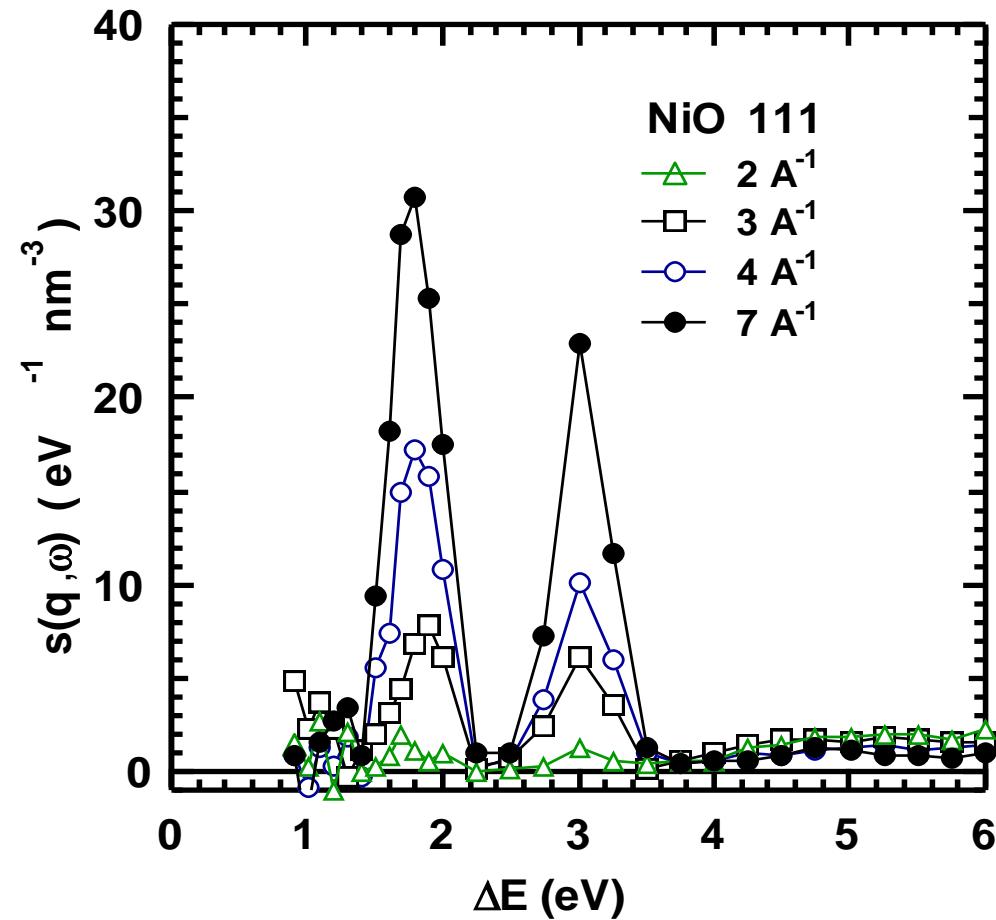
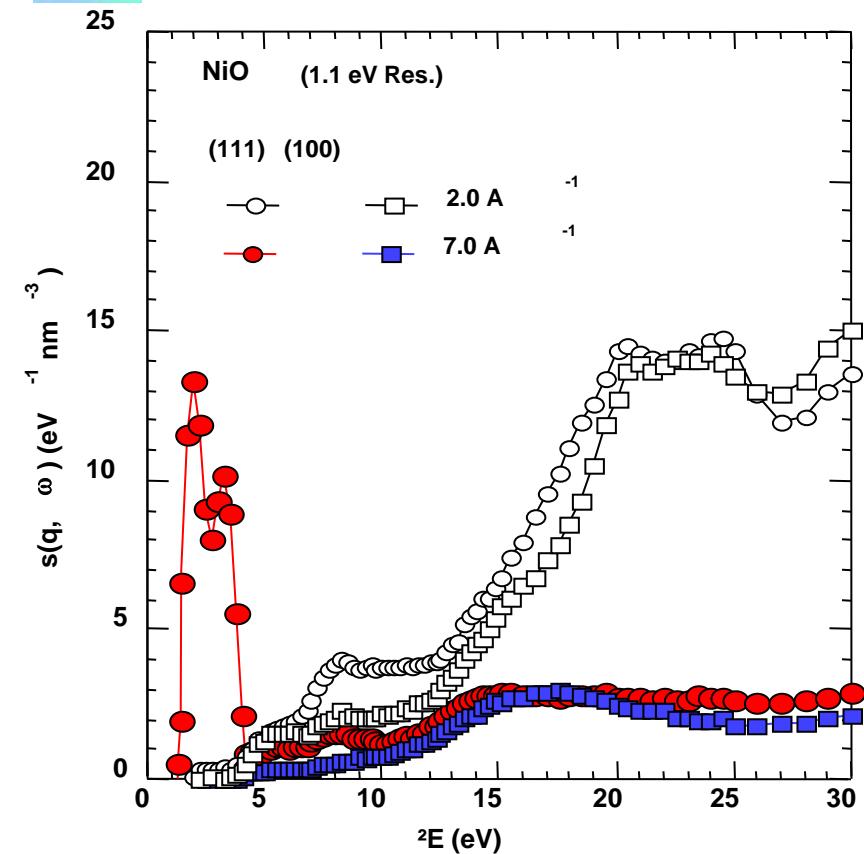
$$\frac{d^2\sigma}{d\Omega d\omega} = r_0^2 (\vec{e}_i \cdot \vec{e}_f)^2 \left(\frac{\omega_f}{\omega_i} \right)^2 S(\vec{q}, \omega)$$

Absolute Response Calculations

$$S(\vec{q}, \omega) = -2\hbar V \text{Im} \chi_{\vec{G}\vec{G}}(\vec{q} - \vec{G}, \omega)$$



Local Excitations in the Mott Gap in NiO



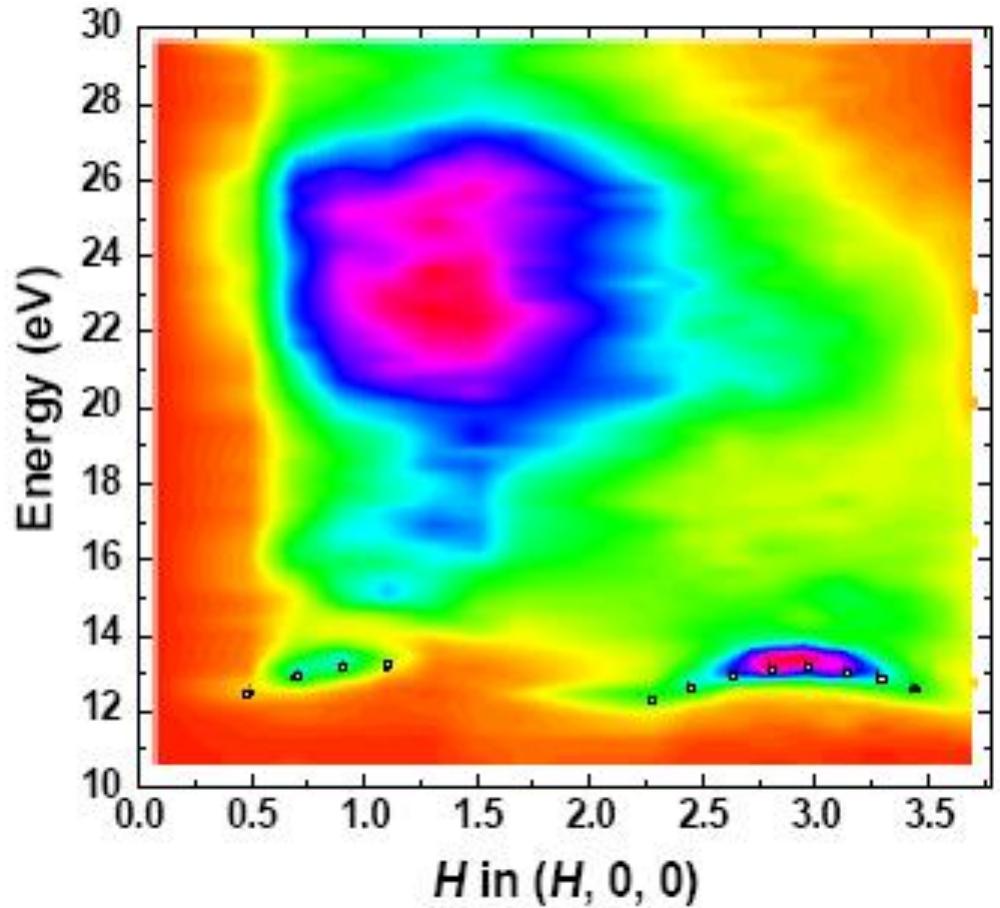
- New features at large q in the Mott gap! B. C. Larson et al., PRL 2007
- Short wave length → local excitations (small exciton)
- Strong angular dependence: (100) \neq (111)

Tightly-Bound Excitons in Charge Transfer Insulators: case study of LiF

- Tightly bound exciton
- Charge transfer insulator
→ p-h in different atoms
- Frenkel or Wannier exciton ?
- Simple dispersion
→ propagation in space/time

Inelastic X-ray scattering

- Structured spectral weight
- Clear dispersion at large q !
- observe fs dynamics



P. Abbamonte et al., PNAS 2008



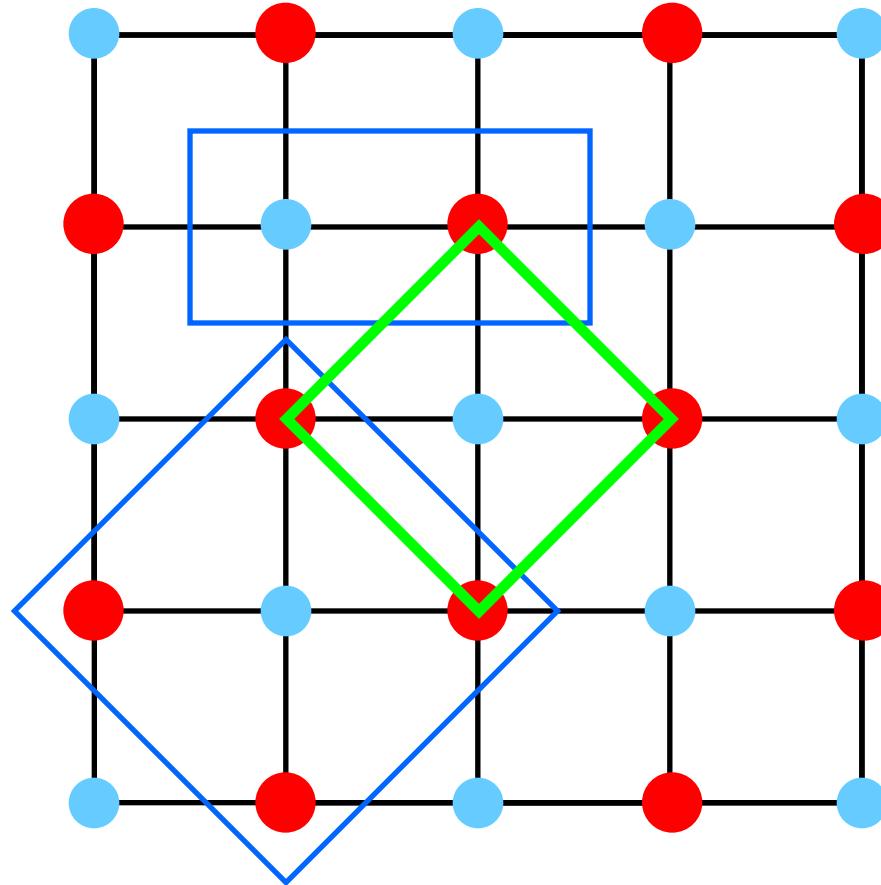
Local Picture for Strongly Correlated Systems

$$H = H_0 + V = \boxed{H_{local}} + H_{nonlocal}$$

- V too big for perturbation
- Maximize the terms in the “local” part
 - symmetric Wannier Representation → defines “local”
- Treat local part “accurately” and leave non-local part as modification



How to define “local” in CT-insulators?



- Periodic symmetry
- Point group symmetry
- Simultaneously keep both? How to split the Hilbert space?

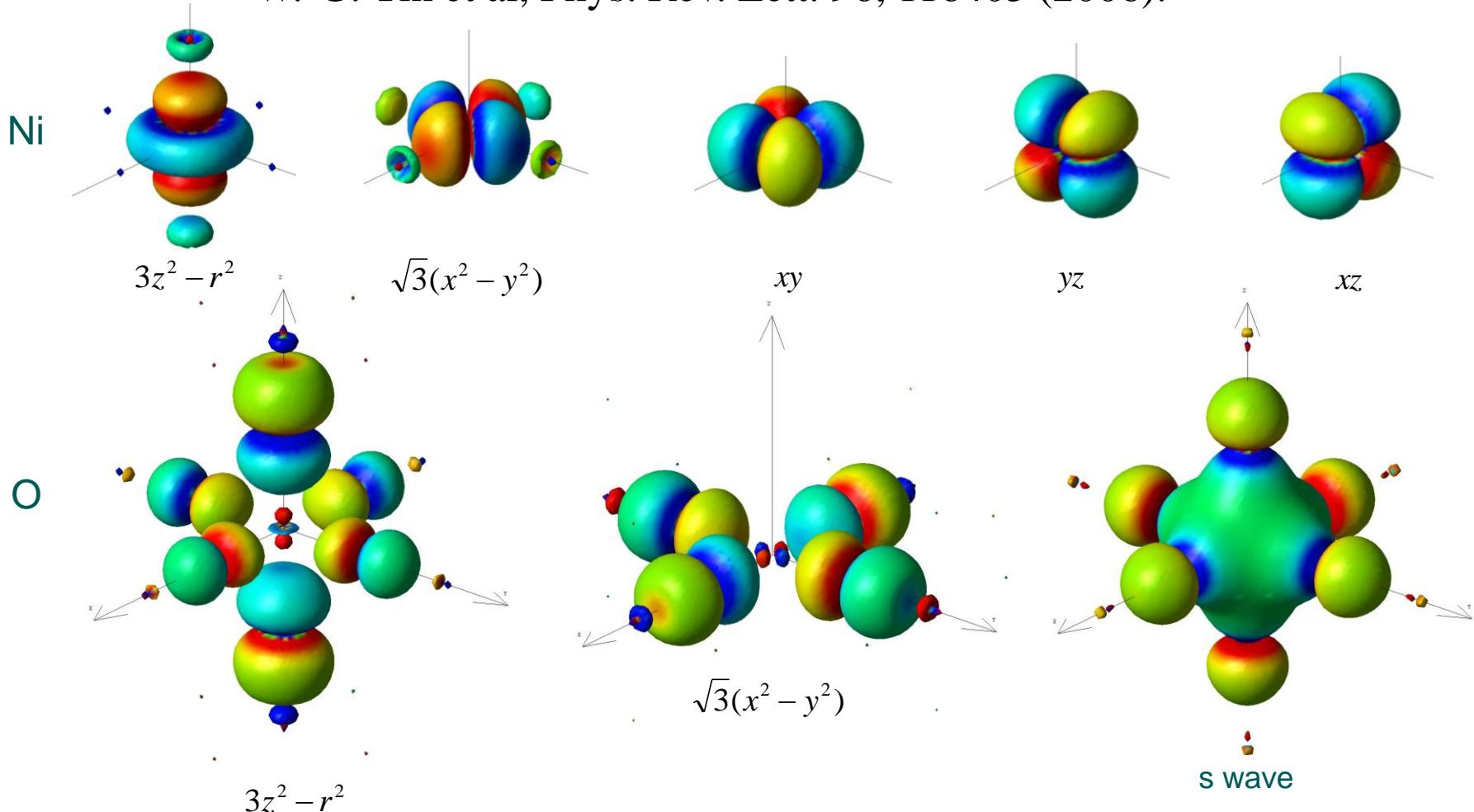


Symmetric Wannier Functions for CT-Insulators

W. Ku et al., Phys. Rev. Lett. **89**, 167204 (2002).

R. L. Barnett et al., Phys. Rev. Lett. **96**, 026406 (2006).

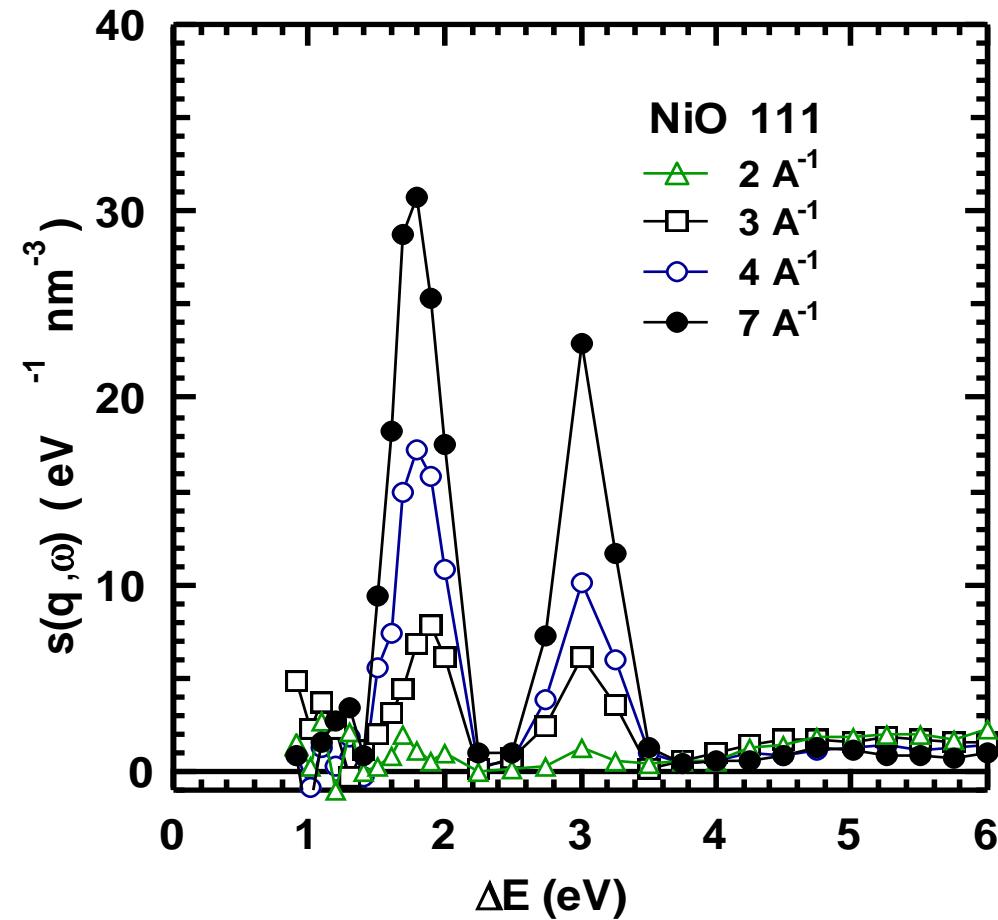
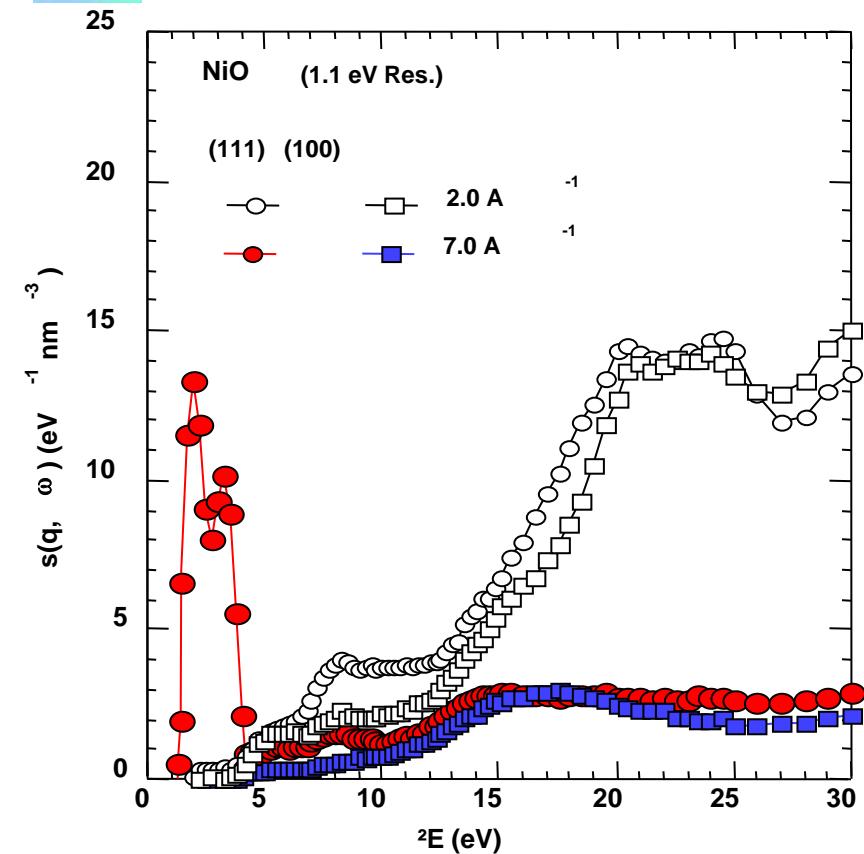
W.-G. Yin et al, Phys. Rev. Lett. **96**, 116405 (2006).



- O-*p* orbitals → additional Ni-*d* orbitals (no double counting of O orbitals)
- “local” is now defined by this “super-atom”



Local Excitations in the Mott Gap in NiO



- New features at large q in the Mott gap! B. C. Larson et al., PRL 2007
- Short wave length → local excitations
- Strong angular dependence: (100) \neq (111)



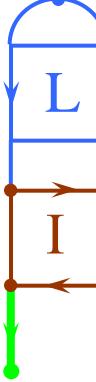
q-dependence of Localized Excitons

Energy-resolved Wannier states

$$\chi(\mathbf{q}, \mathbf{q}'; w) = \sum_{11'} \langle 1 | e^{-i\mathbf{q} \cdot \hat{\mathbf{x}}} | 1' \rangle \langle 1' | e^{i\mathbf{q}' \cdot \hat{\mathbf{x}}} | 1 \rangle L(1, 1'; 1', 1; w)$$

Diagram illustrating the decomposition of the susceptibility χ into energy-resolved Wannier states:

χ =  + 

 + 

Decomposition into energy-resolved Wannier states:

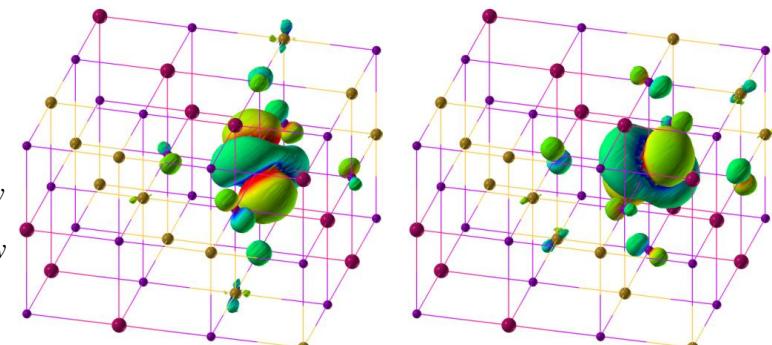
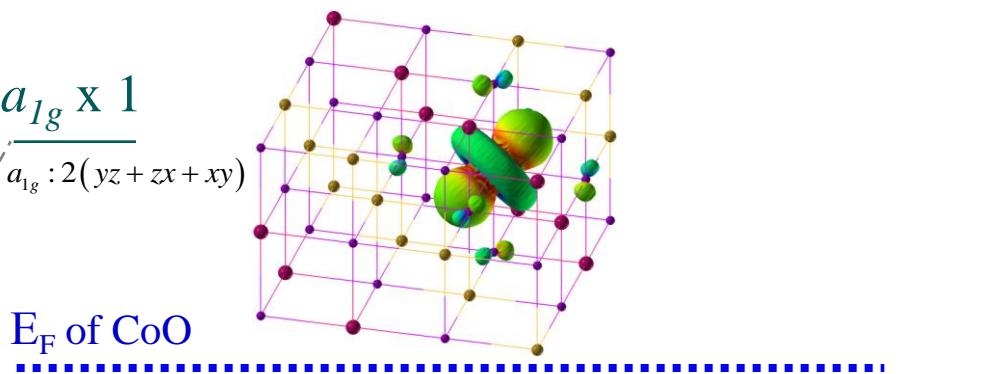
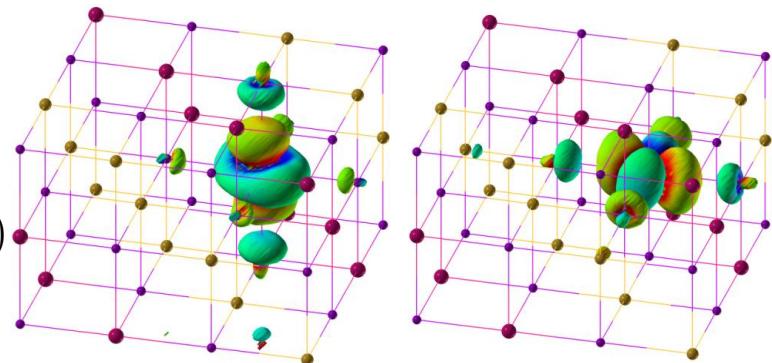
$d \times 5$

$t_{2g} \times 3$

$e_g \times 2$
 $e_g 1 \square 3z^2 - r^2$
 $e_g 2 \square \sqrt{3}(x^2 - y^2)$

$a_{1g} \times 1$
 $a_{1g} : 2(yz + zx + xy)$

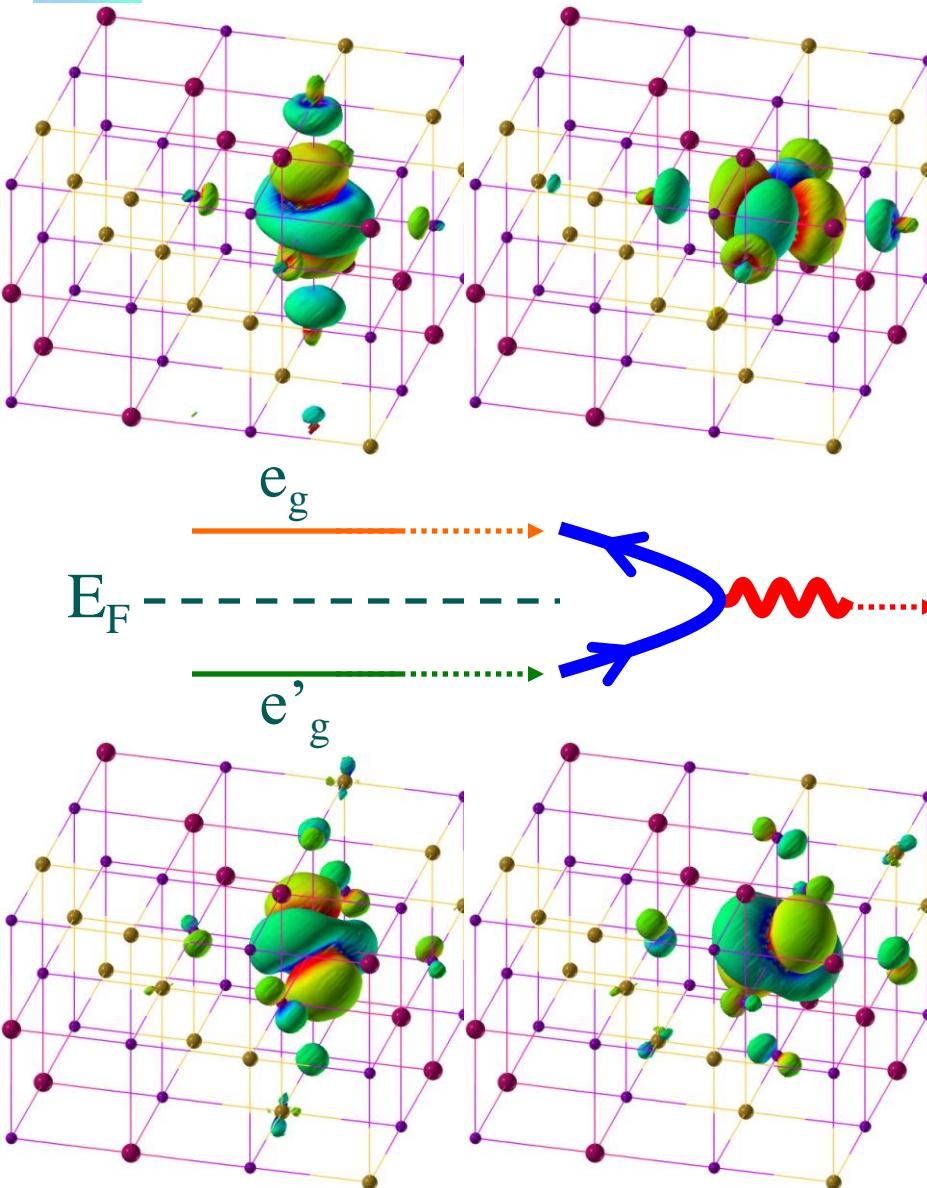
$e'_g \times 2$
 $e'_g 1 \square (1 + \sqrt{3})yz + (1 - \sqrt{3})zx - 2xy$
 $e'_g 2 \square (1 - \sqrt{3})yz + (1 + \sqrt{3})zx - 2xy$



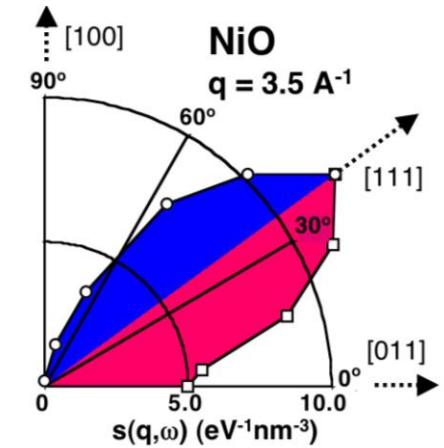
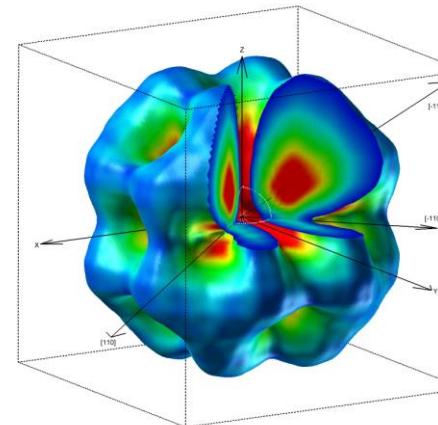


Local Excitations in NiO and CoO

Point group symmetry and new selection rules

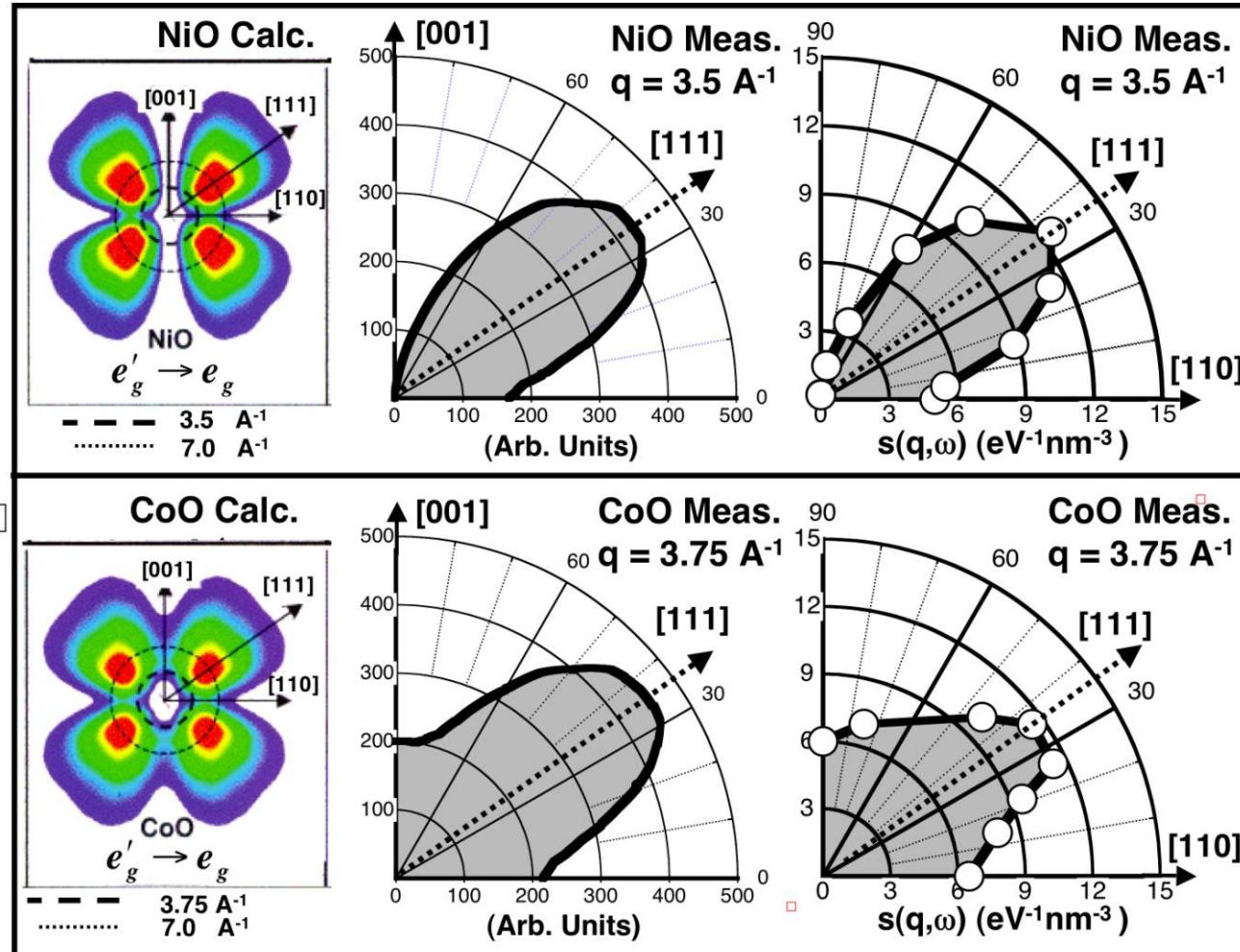


- Local cubic point group symmetry
→ nodal directions
→ new selection rules



Decoupling 6D exciton WF into sum over small number of pair of 3D WF

Angular Dependence of Local Excitations

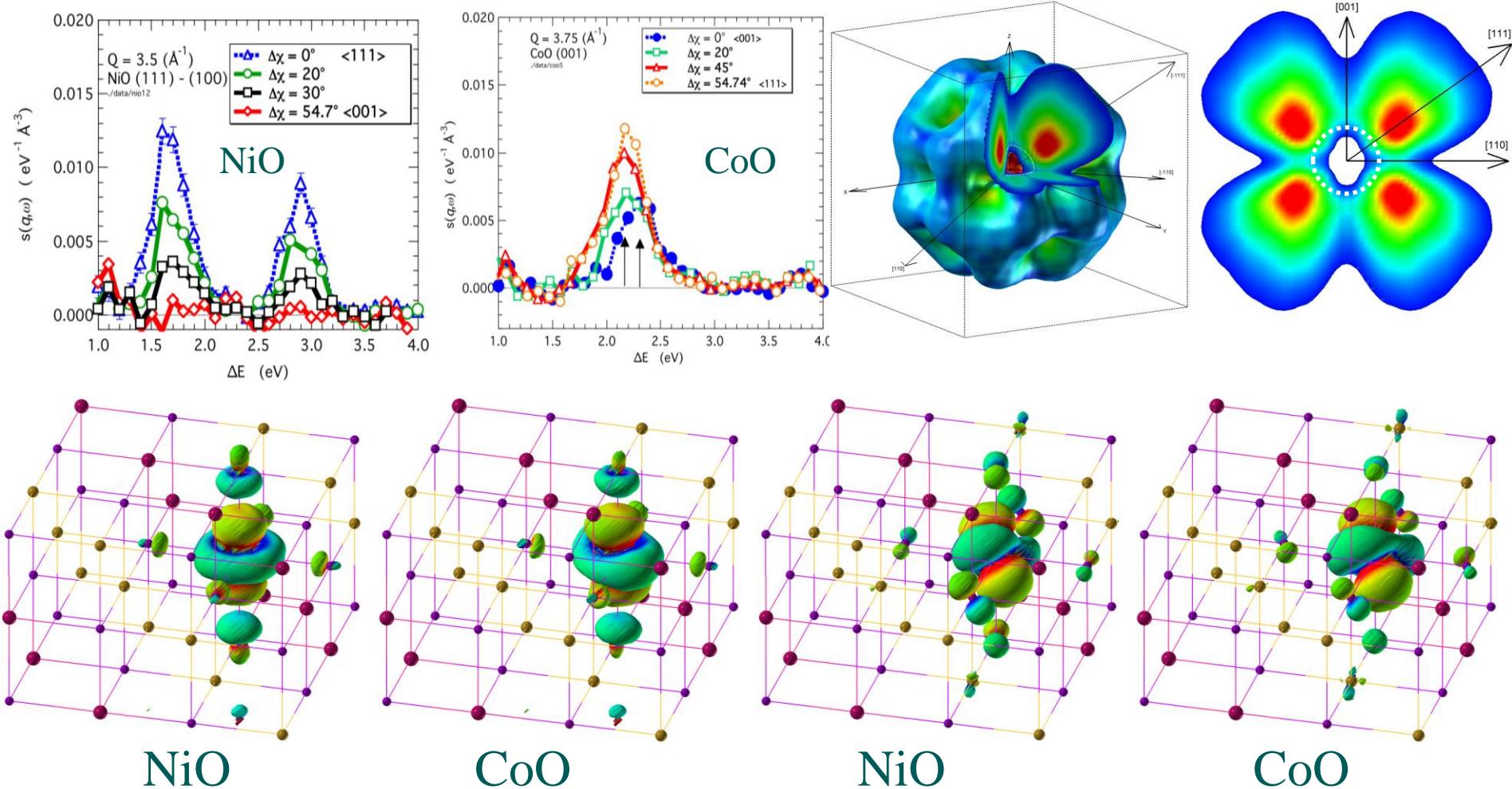


- Nodal direction \leftarrow point group symmetry
- Lack of [100] node in CoO \rightarrow weak symmetry breaking



Local Excitations in NiO and CoO

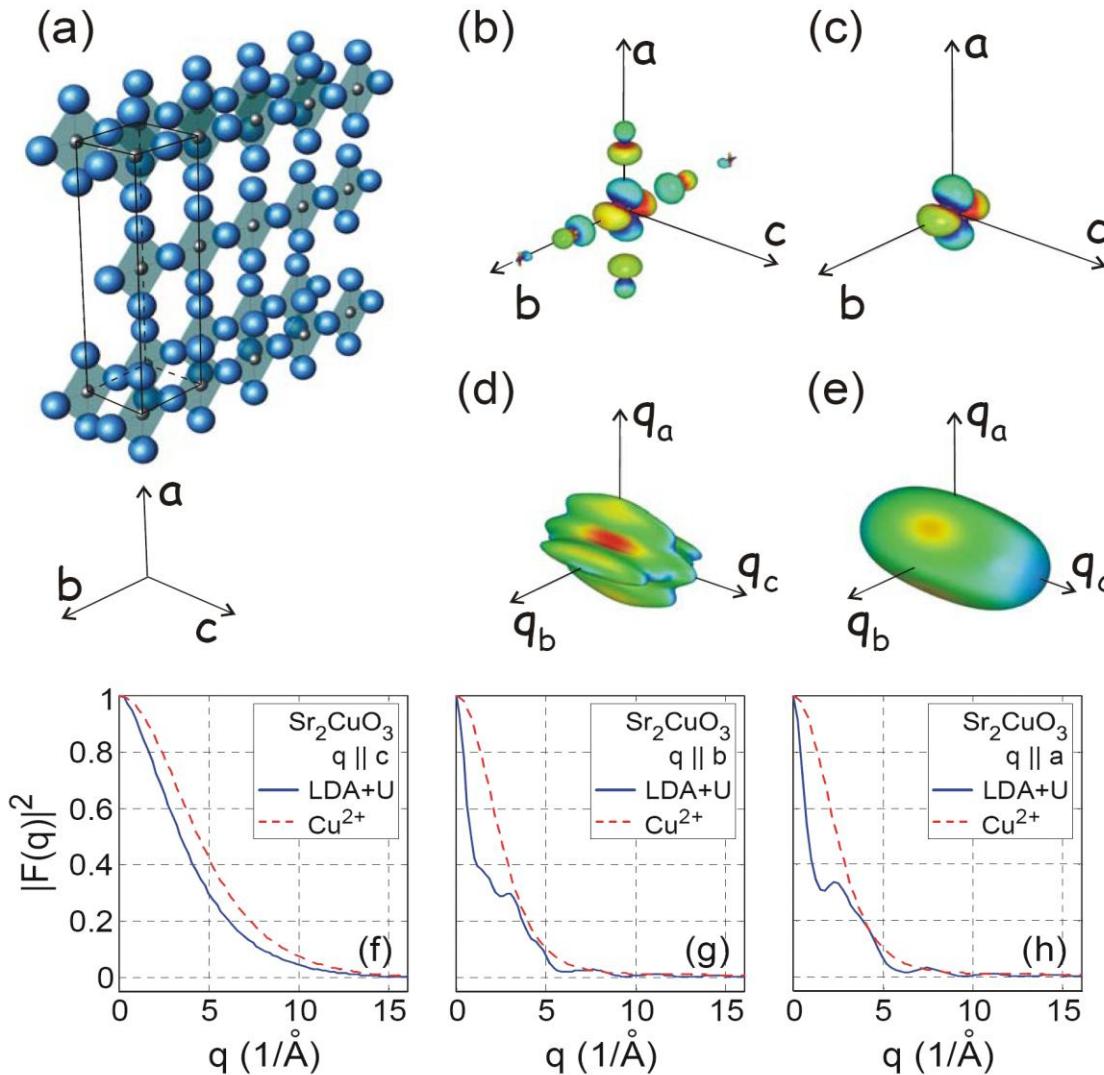
Sensitive probe of weak symmetry breaking



- Lost of nodal directions : extremely sensitive to weak symmetry breaking
- Visualization of symmetry breaking via Wannier functions



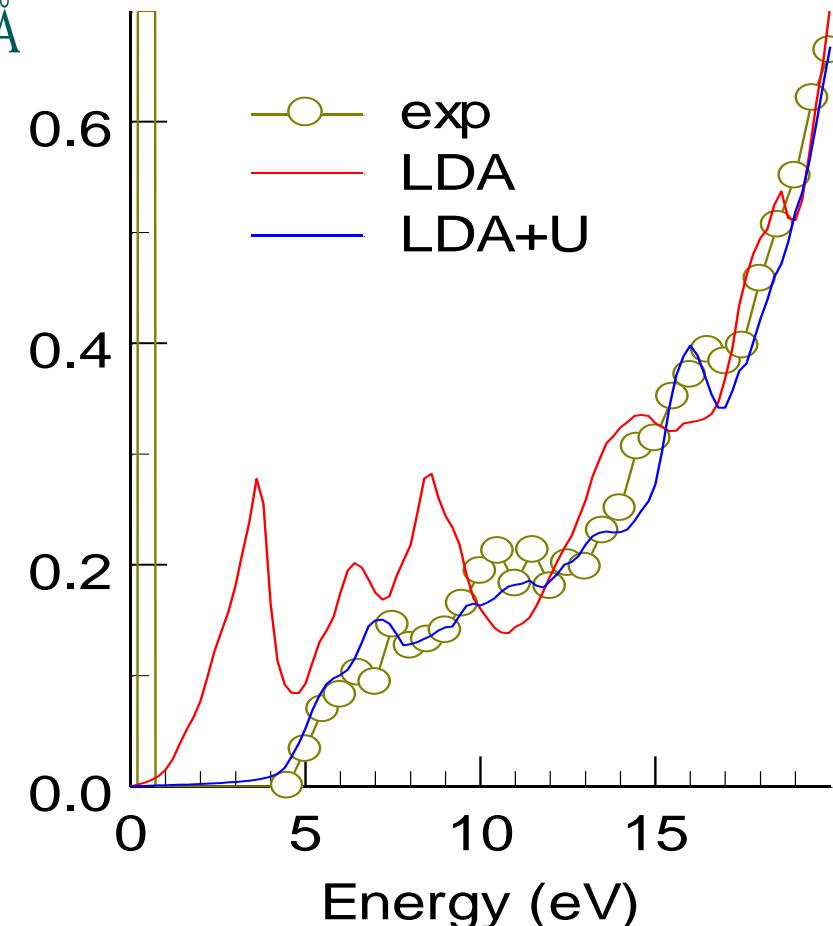
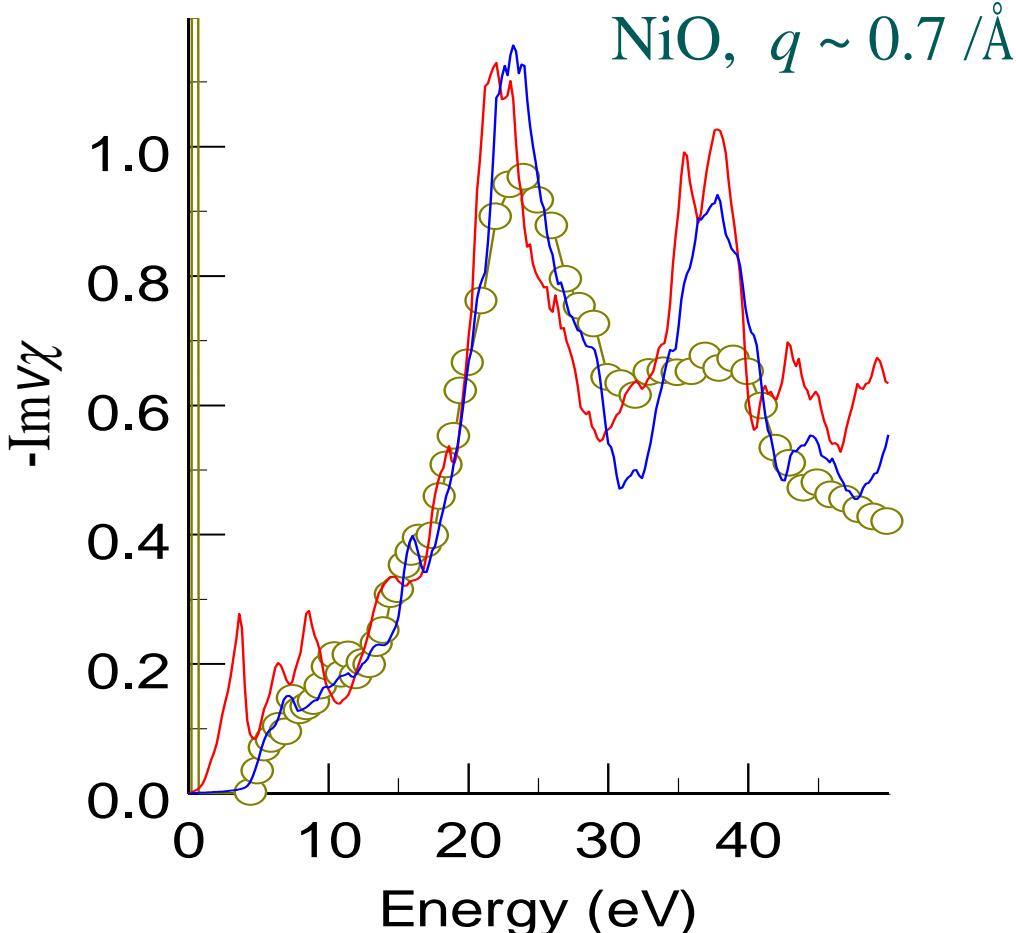
70% Missing Spectral Weight of INX in Cuprates





Charge Excitations in NiO and CoO

Small momentum transfer



- Small $q \rightarrow$ inter-site excitations
- LDA+ U approximation greatly improves the gap and line shape
- Good agreement at small q in absolute unit



TDDFT via LDA+ U Functional (TD-LDA+ U)

The diagram illustrates the decomposition of the Green's function G and the self-energy v_s . The Green's function G is shown as a horizontal arrow pointing left, equal to the sum of two terms: G_0 (dotted arrow) and G_0 (dotted arrow). The self-energy v_s is shown as a blue circle, equal to the sum of several components: v_{ext} (orange circle), $v_{Hartree}$ (green circle with dashed vertical line), v_{XC} (green circle with dashed vertical line), $v_{local\ Hartree}$ (green circle with dashed vertical line), $v_{local\ Fock}$ (red dashed circle with green arrow), and a correction term $-[U(\text{green circle } -1/2) - J(\text{green circle } -1/2)] \cdots$.

response function

$$L = \frac{\delta G}{\delta v_{ext}} = -G \frac{\delta G^{-1}}{\delta v_{ext}} G$$

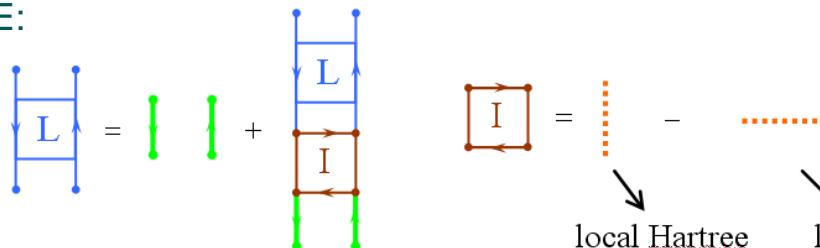
The diagram shows the Bethe-Salpeter equation for the response function L . It is represented by a blue circle with a double-headed arrow, equal to the sum of two terms: a blue circle with a vertical line and a blue circle with a horizontal line. The horizontal line is labeled L . This is further decomposed into a vertical line with a green arrow and a horizontal line with a red arrow, plus a term labeled I (a red square with arrows) and a green vertical line. The entire equation is enclosed in a red bracket with the label "(Bethe-Salpeter equation)".

The diagram shows the decomposition of the Hartree term. A red square with arrows is equal to a vertical dashed line plus a vertical dotted line minus a term labeled $(U-J)$ (two green dots). Arrows point from the terms to labels: "Hartree" under the vertical dashed line and "long-range screening" under the vertical dotted line.

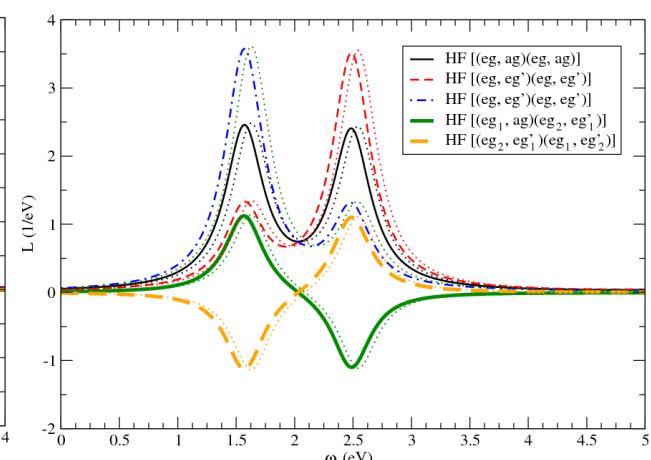
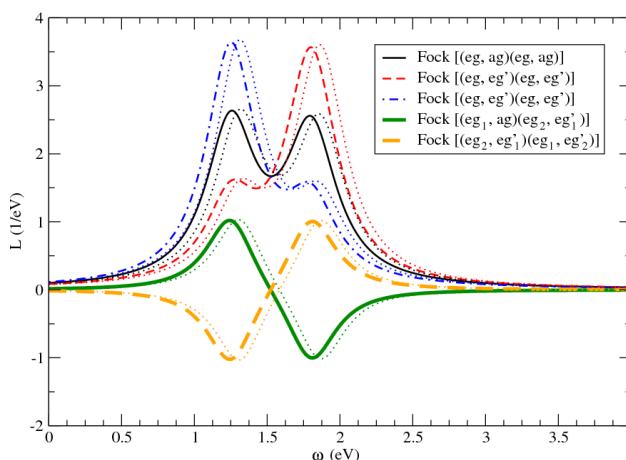
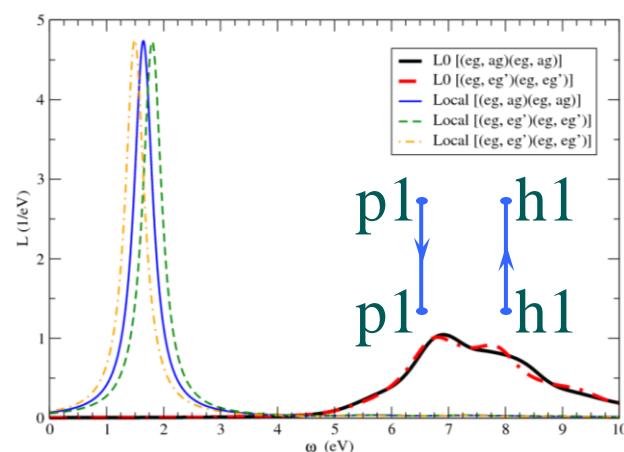
The diagram shows the decomposition of the local Hartree and local Fock terms. A blue box contains a vertical dotted line minus a horizontal dotted line. Arrows point from the terms to labels: "local Hartree" under the vertical dotted line and "local Fock p-h attraction" under the horizontal dotted line.

Formation of Frenkel Excitons in Local Picture

BSE:



(Wannier basis, laptop done all the job)



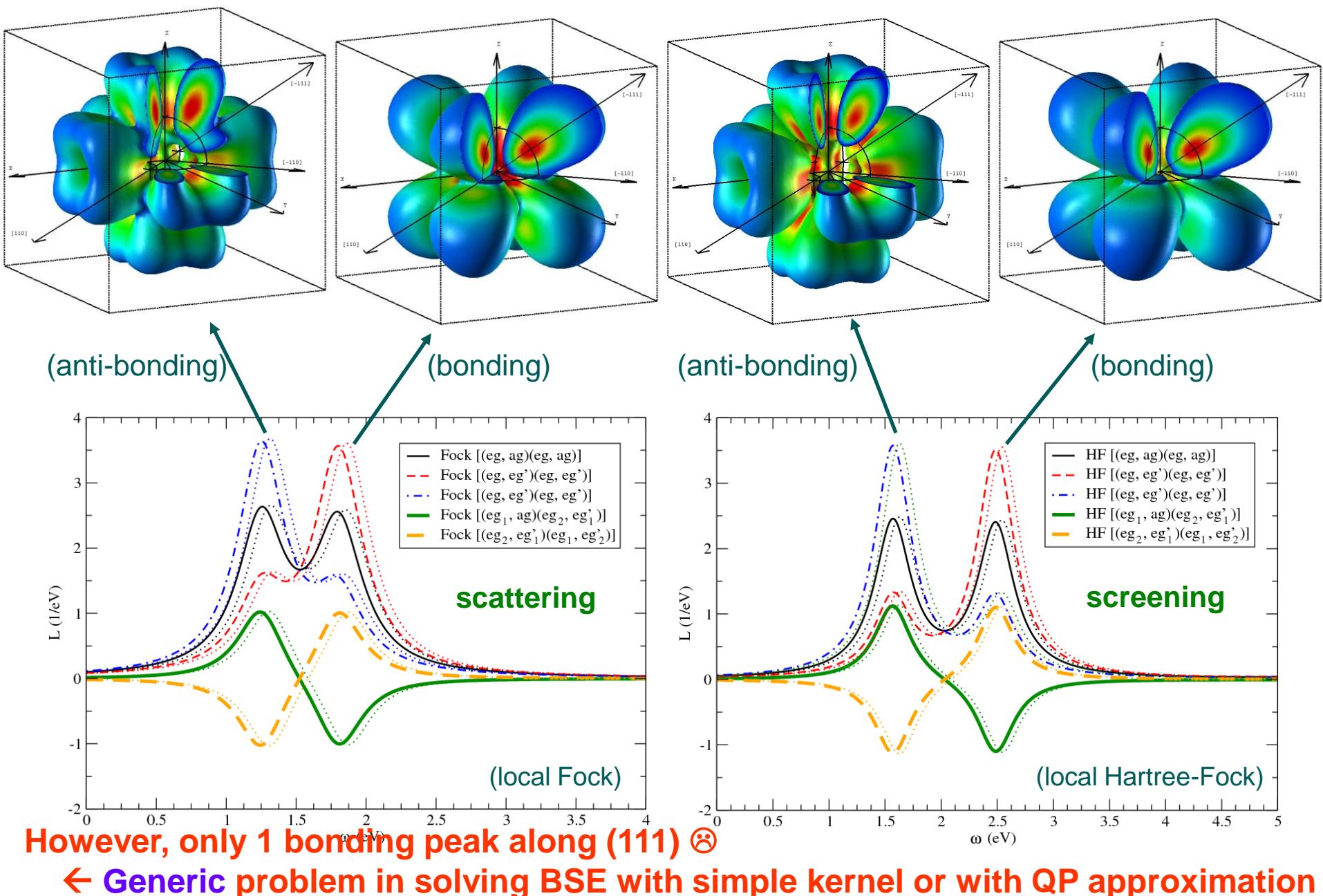
+ $p1$ $h1$ **binding**
 $p1$ $h1$
 same pair
 p-h attraction

+ $p2$ $h2$ **scattering**
 $p1$ $h1$
 local Fock

+ $p1$ $h1$ **screening**
 $p1$ $h1$
 local Hartree

- Exciton energy in reasonable agreement with experiment
- Strongly hybridized Frenkel excitons

Capability and Limitation of the Approximate Functional





Solving Local Interacting Problem Accurately

1. Map to an interacting Hamiltonian

W.-G. Yin, D. Volja and Wei Ku, Phys. Rev. Lett. **96**, 116405 (2006).

2. Solve local part exactly

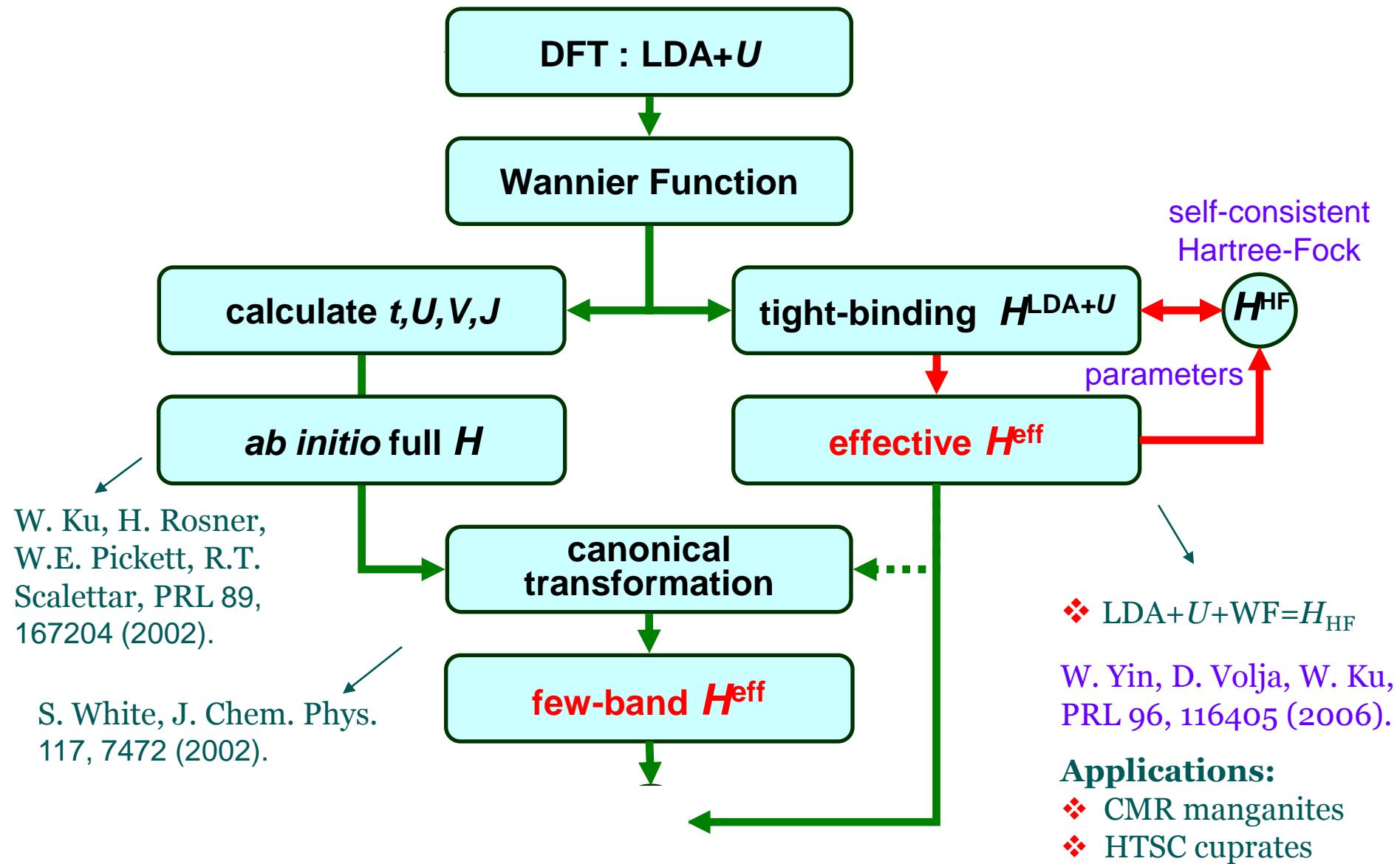
3. Add non-local part via

1) effective kinetic energy for the excitons

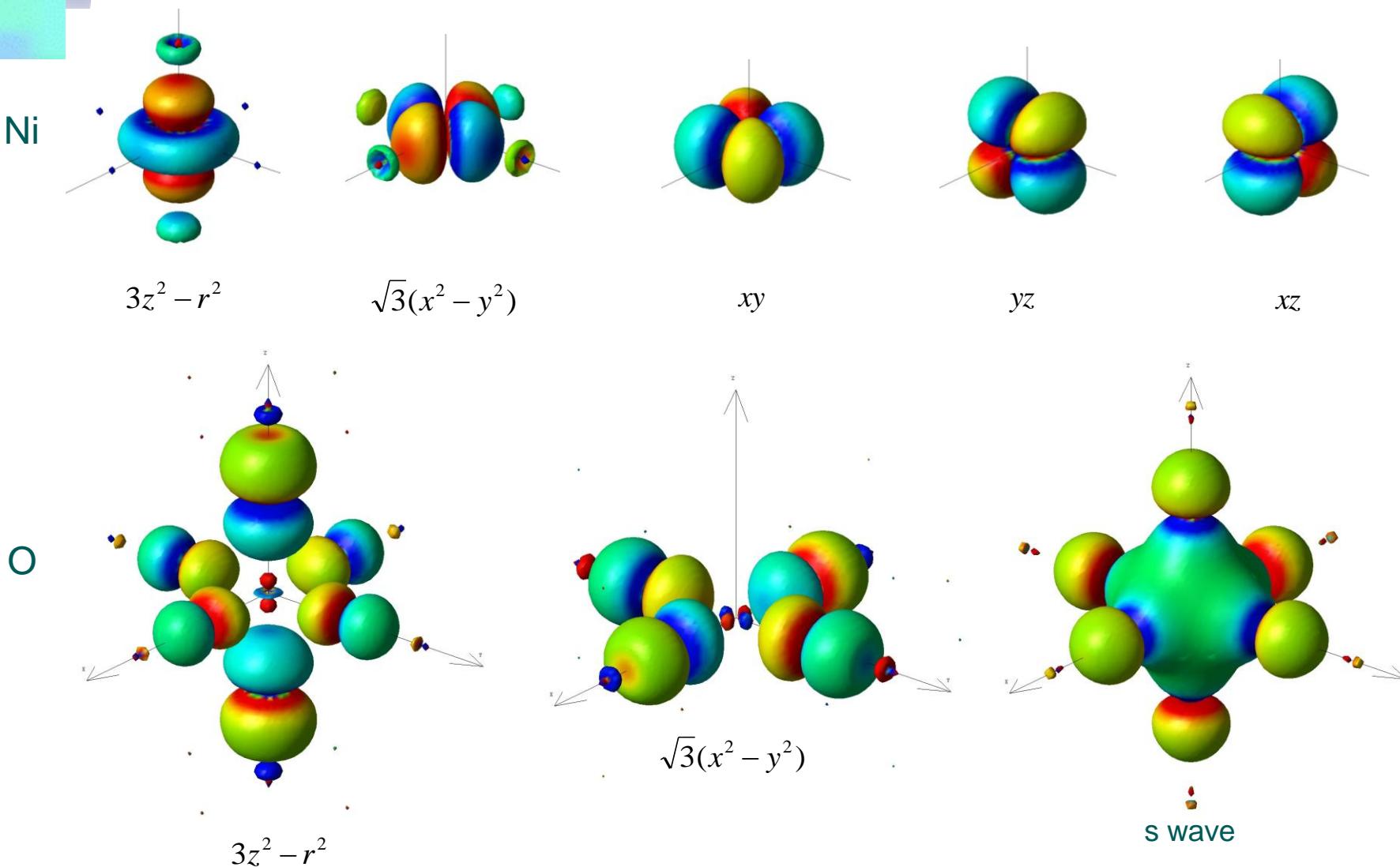
2) extended Hubbard X-operator (powerful but demanding)



Direct mapping to effective Hamiltonian



Super Atom for Charge Transfer Insulator

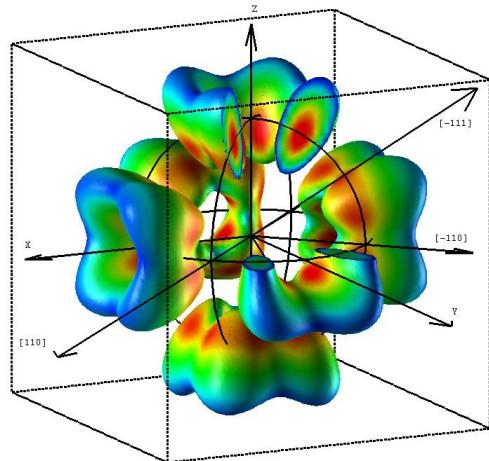


Maximize the contributions of “local atom”

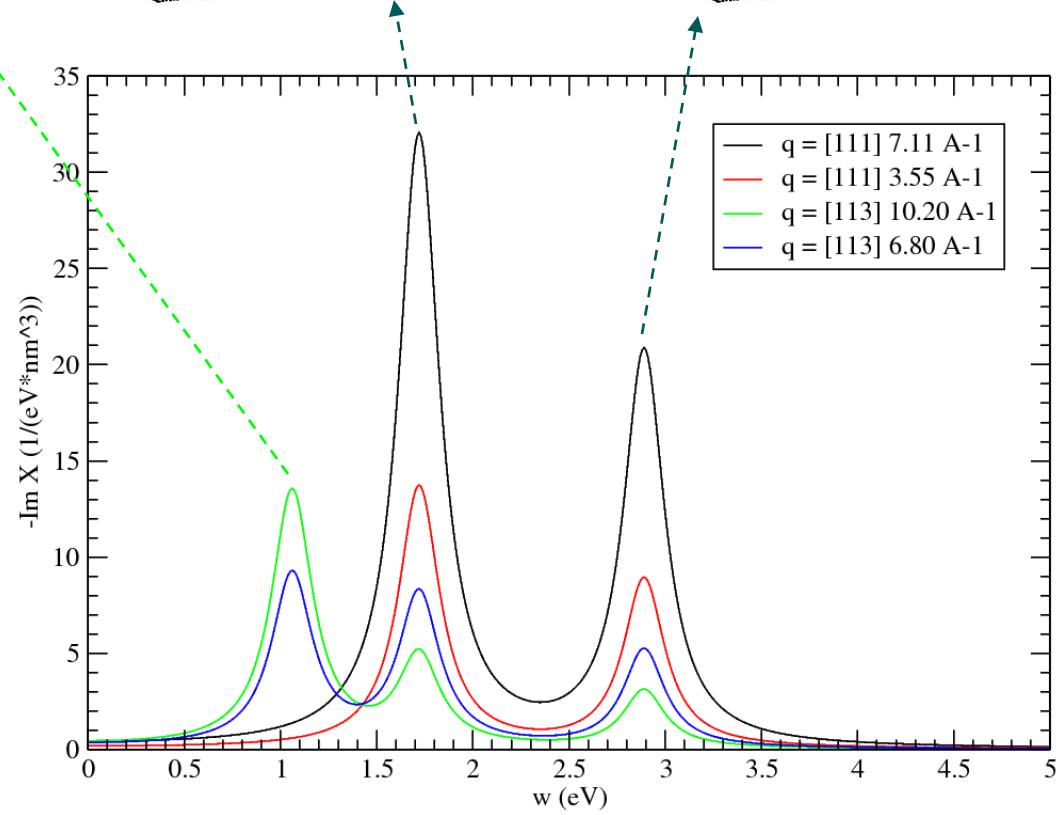
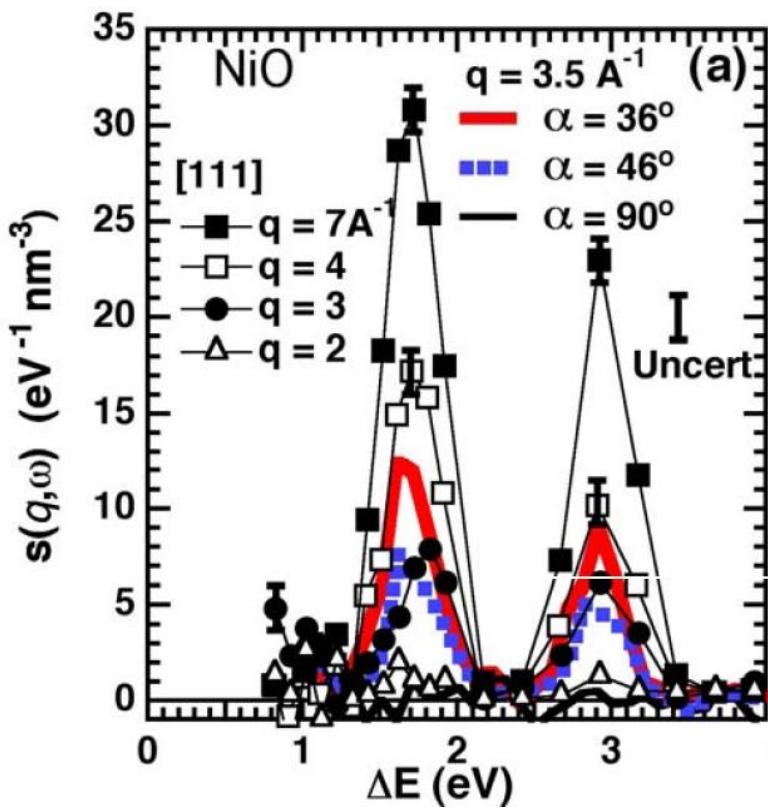
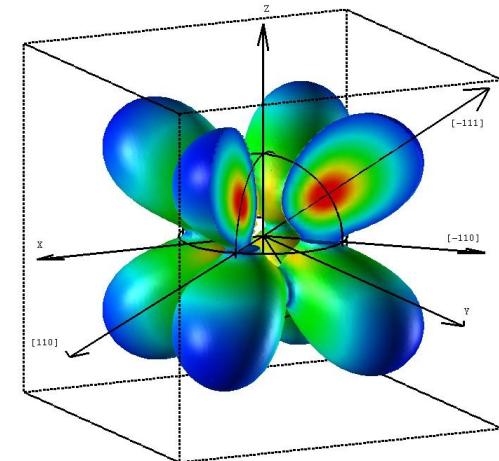
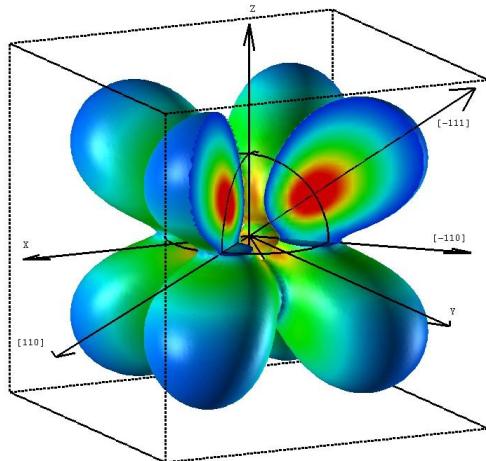


Density response for super atom

(eg-t2g=0.65eV)



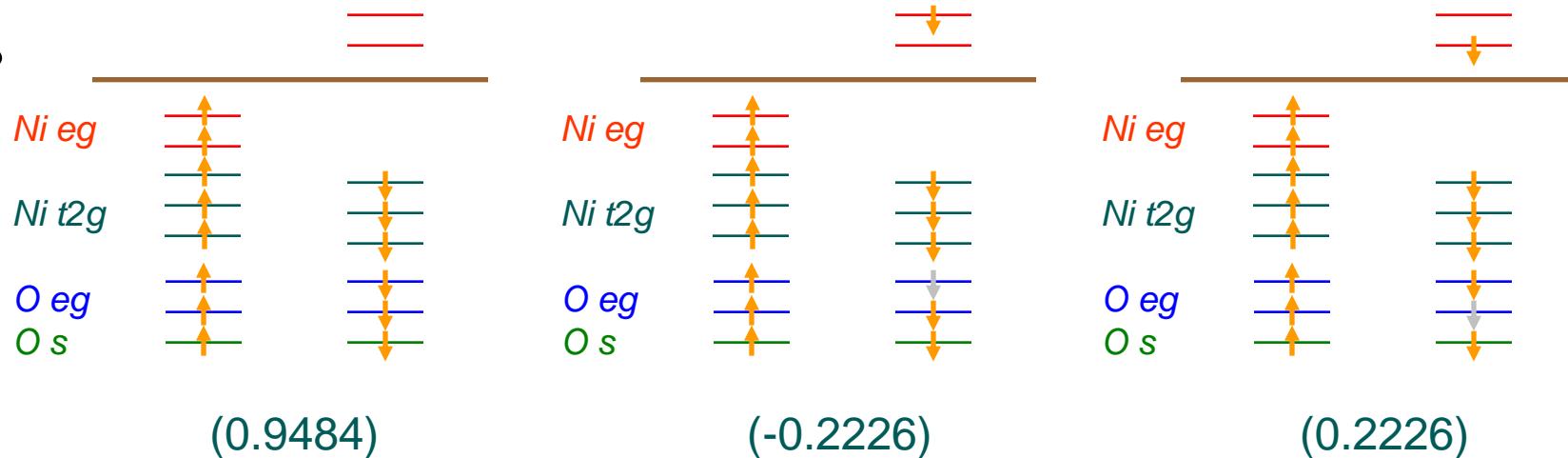
(antibonding-type)



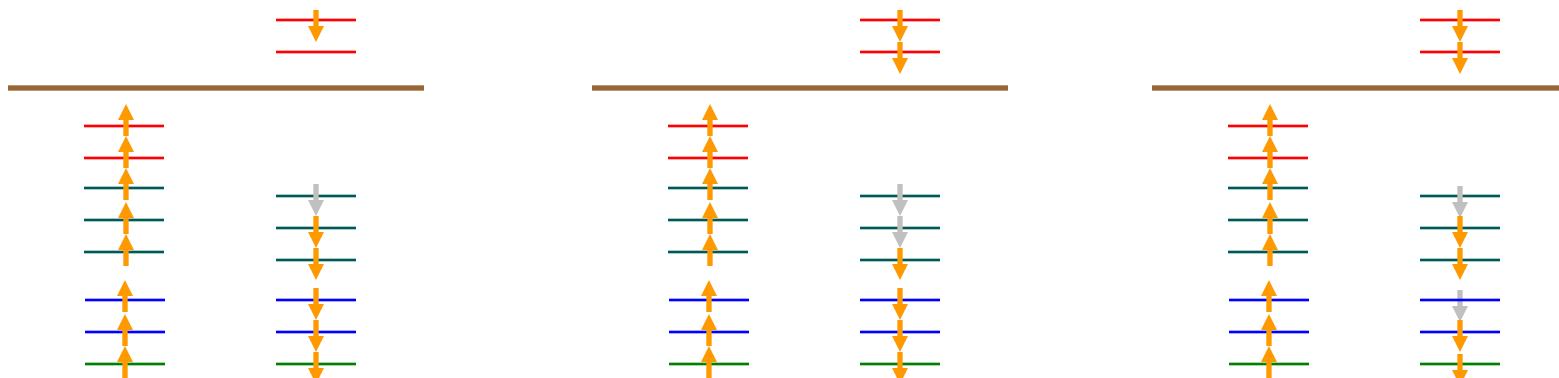


Multiplet Splitting Made Possible with MB Hilbert Space

GS



T_{1g}





Summary

- Recent X-ray experiments
 - Short wave length → sensitivity to local excitations
 - Strong anisotropy → sensitivity to orbitals orientation
- Strong local interaction & short-range correlation
 - Local approach
- Defining “local” in a crystal with “natural” symmetry
 - Symmetric Wannier functions
- Simple understanding of anisotropy from local picture
 - Single p-h pair (Frenkel exciton within super atom)
- 70% missing spectral weight in INS
- Treating local problem via TDDFT: TD-LDA+ U
 - Easy visualization of physics, good energy scale
 - Generic limitation of adiabatic approximation of current *ab initio* approx.
- Treating local problem beyond perturbation
- Propagation of local excitations
 - p-h kinetic kernel