Spin-density wave in the electron gas in Hartree-Fock, Reduced Density-Matrix Functional Theory, and Exact-Exchange Spin-DFT

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Outline

- Overhauser's spiral spin density wave (SSDW)
- Reduced Density Matrix Functional Theory (RDMFT)
- Numerical results: HF and RDMFT
- SSDW in exact-exchange spin-DFT

Uniform electron gas in Hartree-Fock approximation

classic papers by Overhauser (PRL (1960), PR (1962)):

Overhauser's theorem (analytical proof)

In the HF approximation, the paramagnetic state of the uniform electron gas is unstable w.r.t. formation of spin- or charge-density waves for all electron densities

ansatz for HF orbitals in SSDW state:

$$\Phi_{1\mathbf{k}}(\mathbf{r}) = rac{1}{\sqrt{\Omega}} \exp(i\mathbf{k} \cdot \mathbf{r}) \left(egin{array}{c} \cos(heta_{\mathbf{k}}) \\ \sin(heta_{\mathbf{k}}) \exp(i\mathbf{q} \cdot \mathbf{r}) \end{array}
ight)$$

$$\Phi_{2\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{\Omega}} \exp(i\mathbf{k} \cdot \mathbf{r}) \left(\begin{array}{c} -\sin(\theta_{\mathbf{k}}) \\ \cos(\theta_{\mathbf{k}}) \exp(i\mathbf{q} \cdot \mathbf{r}) \end{array} \right)$$



SSDW for uniform electron gas in Hartree-Fock

ansatz provides HF solution with total HF energy lower than paramagnetic state if following self-consistency are satisfied

HF self-consistency conditions for SSDW

$$\tan(2\theta_{\mathbf{k}}) = \frac{2g_{\mathbf{k}}}{\varepsilon_{\uparrow \mathbf{k}} - \varepsilon_{\downarrow \mathbf{k} + \mathbf{q}}} \qquad \varepsilon_{\sigma \mathbf{k}} = \frac{\mathbf{k}^{2}}{2} - V_{\sigma}(\mathbf{k})$$

$$V_{\uparrow}(\mathbf{k}) = \int \frac{d^{3}k'}{(2\pi)^{3}} \frac{4\pi}{|\mathbf{k} - \mathbf{k}'|} \left(n_{1\mathbf{k}} \cos^{2}(\theta_{\mathbf{k}}) + n_{2\mathbf{k}} \sin^{2}(\theta_{\mathbf{k}}) \right)$$

$$V_{\downarrow}(\mathbf{k} + \mathbf{q}) = \int \frac{d^{3}k'}{(2\pi)^{3}} \frac{4\pi}{|\mathbf{k} - \mathbf{k}'|} \left(n_{1\mathbf{k}} \sin^{2}(\theta_{\mathbf{k}}) + n_{2\mathbf{k}} \cos^{2}(\theta_{\mathbf{k}}) \right)$$

$$2g_{\mathbf{k}} = \int \frac{d^{3}k'}{(2\pi)^{3}} \frac{4\pi}{|\mathbf{k} - \mathbf{k}'|} \left(n_{1\mathbf{k}} - n_{2\mathbf{k}} \right) \sin(2\theta_{\mathbf{k}})$$

SSDW for uniform electron gas in Hartree-Fock

ansatz leads to constant density n and spin-spiral density wave for magnetization density $\mathbf{m}(\mathbf{r})$ (use $\mathbf{q} = (0, 0, q)$)

$$\mathbf{m}(\mathbf{r}) = \begin{pmatrix} m_0 \cos(qz) \\ m_0 \sin(qz) \\ 0 \end{pmatrix}$$

and

$$m_0 = -rac{1}{2} \int rac{\mathsf{d}^3 k}{(2\pi)^3} (n_{1\mathbf{k}} - n_{2\mathbf{k}}) \sin(2 heta_{\mathbf{k}})$$

although a simple model, no numerical solution of SSDW in HF for 3-D electron gas has been given !!

often assumed: optimal wavevector for SSDW $q\stackrel{<}{\sim} 2k_F$



Reduced Density Matrix Functional Theory

One-particle reduced density matrix (1-RDM)

$$\gamma(\mathbf{r},\mathbf{r}')=N\int\!\!\mathsf{d}^3r_2\ldots\int\!\!\mathsf{d}^3r_N\Psi^*(\mathbf{r}',\mathbf{r}_2,\ldots,\mathbf{r}_N)\Psi(\mathbf{r},\mathbf{r}_2,\ldots,\mathbf{r}_N)$$

spectral decomposition:

$$\gamma(\mathbf{r}, \mathbf{r}') = \sum_{i} n_{i} \Phi_{i}(\mathbf{r}) \Phi_{i}^{\dagger}(\mathbf{r}')$$

 n_i : occupation numbers

 $\Phi_i(\mathbf{r})$: natural orbitals (Pauli spinors)

Gilbert Theorem

Ground state Ψ^N_0 and ground state energy of system of N interacting electrons is functional of 1-RDM

Ground state energy

$$E_V[\gamma] = T[\gamma] + V[\gamma] + W[\gamma]$$

kinetic energy (exact):

$$T[\gamma] = \frac{1}{2} \sum_{\sigma} \int d^3r \lim_{\mathbf{r} \to \mathbf{r}'} \nabla' \nabla \gamma_{\sigma\sigma}(\mathbf{r}, \mathbf{r}')$$

potential energy (exact):
$$V[\gamma] = \sum_{\sigma} \int \!\! \mathrm{d}^3 r V(\mathbf{r}) \gamma_{\sigma\sigma}(\mathbf{r},\mathbf{r})$$

Interaction energy (approximation needed):

$$W[\gamma] = \sum_{\sigma_1\sigma_2} \int\!\!\mathrm{d}^3r_1 \int\!\!\mathrm{d}^3r_2 \frac{P_{\sigma_1\sigma_2}[\gamma](\mathbf{r}_1,\mathbf{r}_2)}{|\mathbf{r}_1-\mathbf{r}_2|}$$

with ground state pair density $P_{\sigma_1\sigma_2}[\gamma](\mathbf{r}_1,\mathbf{r}_2)$

Approximation for energy functional

use density matrix power functional (Sharma et al, PRB **78**, 201103(R) (2008))

$$\begin{split} P_{\sigma\sigma'}[\gamma](\mathbf{r},\mathbf{r}') &= \frac{1}{2}\gamma_{\sigma\sigma}(\mathbf{r},\mathbf{r})\gamma_{\sigma'\sigma'}(\mathbf{r}',\mathbf{r}') - \frac{1}{2}\gamma_{\sigma\sigma'}^{\alpha}(\mathbf{r},\mathbf{r}')\gamma_{\sigma'\sigma}^{\alpha}(\mathbf{r}',\mathbf{r}) \\ \gamma^{\alpha}(\mathbf{r},\mathbf{r}') &= \sum_{i} n_{i}^{\alpha}\Phi_{i}(\mathbf{r})\Phi_{i}^{\dagger}(\mathbf{r}') \quad \text{ and } \quad 0.5 \leq \alpha < 1 \end{split}$$

limiting cases:

 $\alpha = 1$: Hartree-Fock

lpha = 0.5: Müller or Buijse-Baerends functional

(Müller, Phys. Lett. (1984), Buijse, Baerends, Mol. Phys. (2002))

for SSDW: use spinors of the form of Overhauser's HF spinors



Numerical Procedure

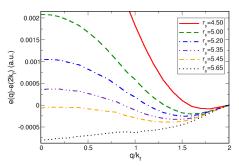
total energy per particle is functional of occupation numbers $n_{1\mathbf{k}}$, $n_{2\mathbf{k}}$ and of angle $\theta_{\mathbf{k}}$ discretize k-space with points $k_i \longrightarrow$ total energy becomes high-dimensional function of $n_{1\mathbf{k}_i}$, $n_{2\mathbf{k}_i}$ and $\theta_{\mathbf{k}_i} \longrightarrow$ optimziation with steepest descent

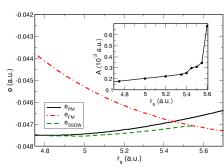
details in: F.G. Eich et al, cond-mat/0910.0534

HF total energies and phase diagram

HF total energy per electron as function of q for various r_{s}

HF total energy of PM, FM, and SSDW phases inset: SSDW amplitude

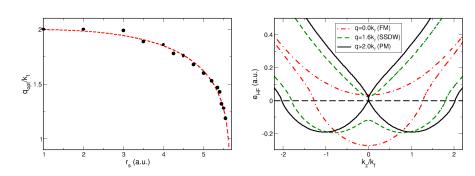




Optimal SSDW wavevector and HF single-particle bands

optimal SSDW wavevector

HF energy bands at $r_s = 5$



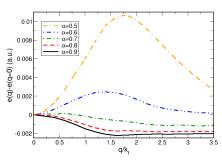
note: q not necessarily close to $2k_F!$

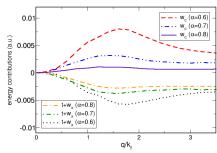


RDMFT total energy and energy contributions

RDMFT total energy per electron for $r_s = 5.0$ and different values of α

energy contributions for $r_s=$ 5.0 and different lpha





corelation destroys SSDW!



Spin-density wave in non-collinear spin-DFT

Ref: S. Kurth, F.G. Eich, PRB 80, 125120 (2009)

Kohn-Sham equation of non-collinear spin-DFT

$$\left(-\frac{\nabla^2}{2} + v_s(\mathbf{r}) + \mu_B \boldsymbol{\sigma} \mathbf{B}_s(\mathbf{r})\right) \Phi_i(\mathbf{r}) = \varepsilon_i \Phi_i(\mathbf{r})$$

here: assume form of KS potentials:

$$v_s(\mathbf{r}) = 0$$

$$\mathbf{B}_s(\mathbf{r}) = (B\cos(qz), B\sin(qz), 0)$$



analytic solution: Kohn-Sham orbitals and orbital energies

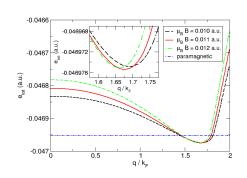
$$\begin{split} & \Phi_{1\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{\Omega}} \exp(i\mathbf{k}\mathbf{r}) \begin{pmatrix} \cos(\theta_k) \\ \sin(\theta_k) \exp(iqz) \end{pmatrix} \\ & \Phi_{2\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{\Omega}} \exp(i\mathbf{k}\mathbf{r}) \begin{pmatrix} -\sin(\theta_k) \\ \cos(\theta_k) \exp(iqz) \end{pmatrix} \\ & \varepsilon_{1\mathbf{k}} = \frac{k_x^2 + k_y^2}{2} + \varepsilon_\kappa^{(-)} \qquad \varepsilon_{2\mathbf{k}} = \frac{k_x^2 + k_y^2}{2} + \varepsilon_\kappa^{(+)} \\ & \varepsilon_\kappa^{(\pm)} = \frac{\kappa^2}{2} + \frac{q^2}{8} \pm \sqrt{\frac{q^2}{4}\kappa^2 + \mu_B^2 B^2}, \qquad \kappa = k_z + \frac{q}{2} \\ & \tan(\theta_\kappa) = \frac{1}{2\alpha} (1 - \sqrt{1 + 4\alpha^2}), \qquad \alpha = \frac{\mu_B B}{q\kappa} \end{split}$$

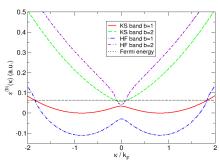
total energy per particle $e_{EXX}^{tot}(q,B) \longrightarrow \text{minimize w.r.t. } q \text{ and } B$

Energy minimization: occupied states in two KS bands

total energy per electron as function of q for various B for $r_s = 5.4$

HF and KS bands for optimal $q=1.68~k_F$ and $\mu_B B=0.011$ a.u.

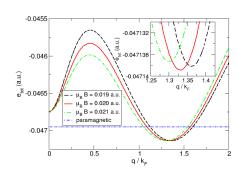


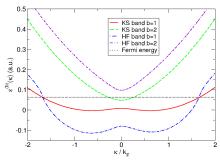


Energy minimization: occupied states in one KS band

total energy per electron as function of q for various B for $r_s=5.4\,$

HF and KS bands for optimal $q=1.33~k_F$ and $\mu_B B=0.020$ a.u.

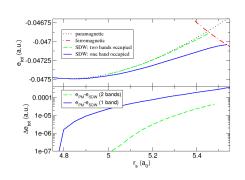


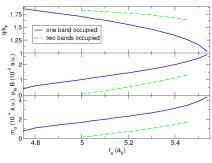


Phase diagram and optimal parameters: 1 and 2 bands

Phase diagram: PM, FM and SSDW (occupations in 1 and 2 bands)

optimal values of q and B and SSDW amplitude





OEP equations in non-collinear spin-DFT

four OEP equations

$$\sum_{i}^{occ} \left(\Phi_{i}^{\dagger}(\mathbf{r}) \Psi_{i}(\mathbf{r}) + h.c. \right) = 0$$

$$-\mu_B \sum_{i}^{occ} \left(\Phi_i^{\dagger}(\mathbf{r}) \boldsymbol{\sigma} \Psi_i(\mathbf{r}) + h.c. \right) = 0$$

 $\Psi_i(\mathbf{r})$: orbital shifts (see Txema's talk last week)

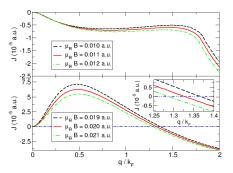
for SSDW: first and last OEP eqs. exactly satisfied 2nd and 3rd OEP eq. equivalent

$$J(q,B)\cos(qz) = 0 J(q,B)\sin(qz) = 0$$



is OEP equation satisfied?

prefactor of OEP eq. for $r_s=$ 5.4 as function of q for various B upper panel: two-band case, lower panel: one-band case



only for one-band case OEP eq. is satisfied for the parameter values minimizing the total energy!

Collaborators:

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- C.R. Proetto, Free Univ. Berlin, Germany, and Centro Atomico Bariloche, Argentina
- S. Sharma and E.K.U. Gross, MPI Halle, Germany

Summary

- SSDW instability in the uniform electron gas in HF, RDMFT, EXX-SDFT
- ullet HF: optimal wavevector can be far from $2k_f$
- RDMFT: correlation destroys SSDW
- ullet EXX-SDFT: SSDW stable over smaller range of r_s than HF
- EXX-SDFT: occupation in one band with holes below the Fermi energy gives lower energy than two-band case and is consistent with OEP equations