Quantum criticality in diluted 2D antiferromagnets

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Simulations of dimer-diluted S=1/2 Heisenberg models



C Quantum criticality in the presence of disorder

Geometric/quantum criticality at percolation point

Background: Random antiferromagnets ID random J>0 Heisenberg chain: $H = \sum J_i \vec{S}_i \cdot \vec{S}_{i+1}$ Infinite-randomness fixed point; random-singlet phase: Spins form singlets on all length-scales at T=0Ma, Dasgupta, Hu, 1979; D. Fisher, 1994 **Renormalization procedure:** successive decimation of strongest-coupled pairs $\chi \sim \frac{1}{T \ln^2(T)}$ Dynamic exponent: $Z = \infty$ $\omega = k^Z, \Delta_L = L^{-Z}$

 $\begin{array}{ll} \mbox{Dimerized chain:} & H = \sum_i [J_i + \delta(-1)^i] \vec{S}_i \cdot \vec{S}_{i+1} \\ \mbox{Griffiths fixed point;} & Z \propto 1/\delta \\ \mbox{Igloi, Juhasz, Rieger (2000)} \end{array}$

2D random transverse Ising model: $Z = \infty$ Motrunich, Mau, Huse, D. Fisher (2000); Pich, Young, Rieger, Kawashima (1998)

2D random J Heisenberg models: RG procedure \Rightarrow conventional scaling Lin, Melin, Rieger, and Igloi (2003)







Site-diluted 2D Heisenberg model

Cuprates: $Cu \rightarrow Zn$ substitution

Fraction p of non-magnetic sites

Sublattice magnetization vs p

Is p_c less than percolation point $p^* \approx 0.41$?

Kato et al. (2000); p_c=p*

Sandvik (2001); percolating cluster ordered

Classical percolation transition





Comparison with neutron scattering; $La_2(Zn,Mg)_xCu_{1-x}O_4$



Vajk, Greven, Mang, Lynn, Gehring (2001)

Order slightly suppressed; likely other interactions present

Quantum-criticality in dimerized 2D systems



Singlet-formation at strong bonds ("dimers") ⇒ antiferromagnetic to spin-gapped transition



 $g_c \approx 2.5$ in both bilayer and single-layer model

Quantum-critical diluted systems

Unpaired spins in site-diluted gapped system ⇒ localized moments form; order antiferromagnetically ⇒ no quantum-critical point



Dimer-diluted models:



Bilayer model

Sandvik (2002);Vajk, Greven (2002) Sknepnek,Vojta,Vojta (2004)



Geometric percolation transition expected for g < g^*







Check of scaling (Z) using antiferromagnetic quantities $q = (\pi, \pi)$ $S(q) = \langle S_{-q}^{z} S_{q}^{z} \rangle = \int d\omega S(q, \omega) \sim L^{1-\eta}$ $\chi(q) = \int_{0}^{1/T} d\tau \langle S_{-q}^{z}(\tau) S_{q}^{z}(0) \rangle = \int d\omega \frac{1}{\omega} S(q, \omega) \sim L^{1+z-\eta}$



$$1 - \eta = \gamma_S$$
$$1 + z - \eta = \gamma_\chi$$

Data at g close to g^* consistent with

$$z = \gamma_{\chi} - \gamma_S \approx 1.36$$



Correlation length measured by neutron scattering in $La_2(Zn,Mg)_xCu_{1-x}O_4$

Vajk, Greven, Mang, Lynn, Gehring (2001)

Close to percolation: $\xi \sim 1/T^z, z \approx 1.4$

Could this be due to multi-critical point (g*, p*) in extended parameter space?

How universal is the behavior found for the diluted bilayer antiferromagnet?

Can similar behavior be observed in a single layer?

Single layer with staggered dimers





Dimers form triangular lattice; $p^*=1/2$

Brick-wall lattice when g=0; different percolation point



$$P^*(g=0) \approx 0.29$$

Uniform susceptibility at $p^*(g>0)$

- Indicates critical point; $g_c \approx 1.247$
- Finite susceptibility at critical point
- Divergent susceptibility for $g < g_c$



- Divergence follows closely $T^{-1/2}$ for g > I
- Faster divergence for smaller g
- Curie form for g=0, due to "broken clusters"







Line of critical points with continuously varying exponents



$S(\pi,\pi), \ \chi(\pi,\pi)$ Finite-size scaling at g=1.0 Z \approx 3.0

Disagrees with Z≈4.0 from uniform susceptibility None-asymptotic behavior?



Conclusions

Novel quantum-critical behavior discovered numerically in dimer-diluted systems



Conventional scaling for the $p < p^*$ transition; $z \approx 1.31$

Why does the single layer have a line of critical points? Conditions under which the percolating cluster orders?