

# Precision EW Measurements at the Large Hadron Collider at 8 TeV and 13-14 TeV

( $\sin^2\theta_W^{\text{eff}}$ ,  $\sin^2\theta_W^{\text{on-shell}}$ ,  $M_W^{\text{indirect}}$ ,  $M_W^{\text{direct}}$  )

Arie Bodek, University of Rochester

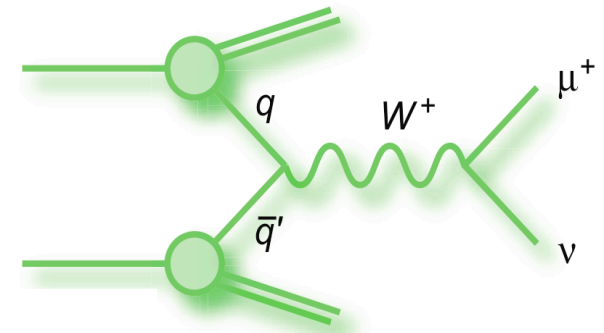
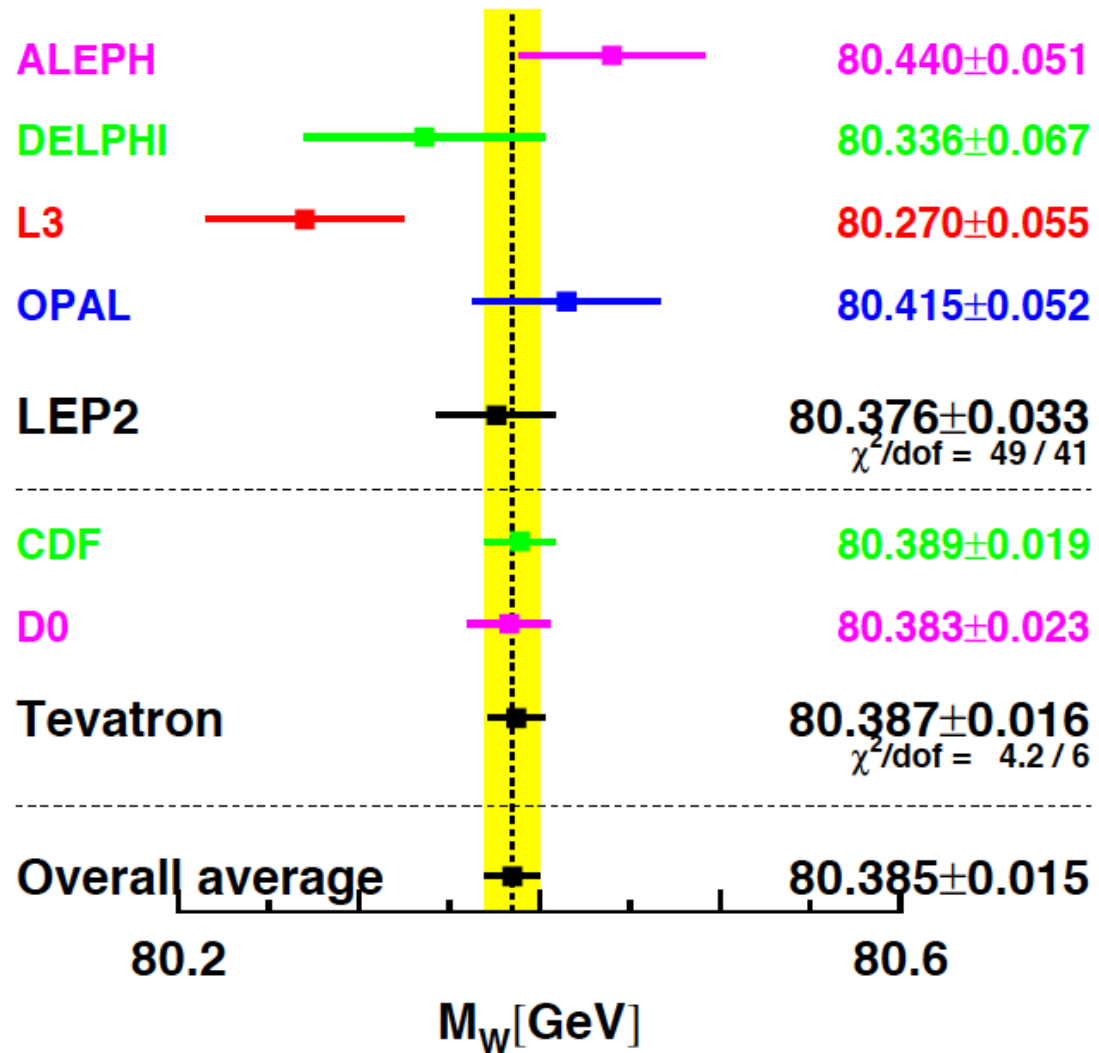


SM session of the

“Experimental Challenges for the LHC Run II” KITP workshop  
**Monday March 28, 2016**

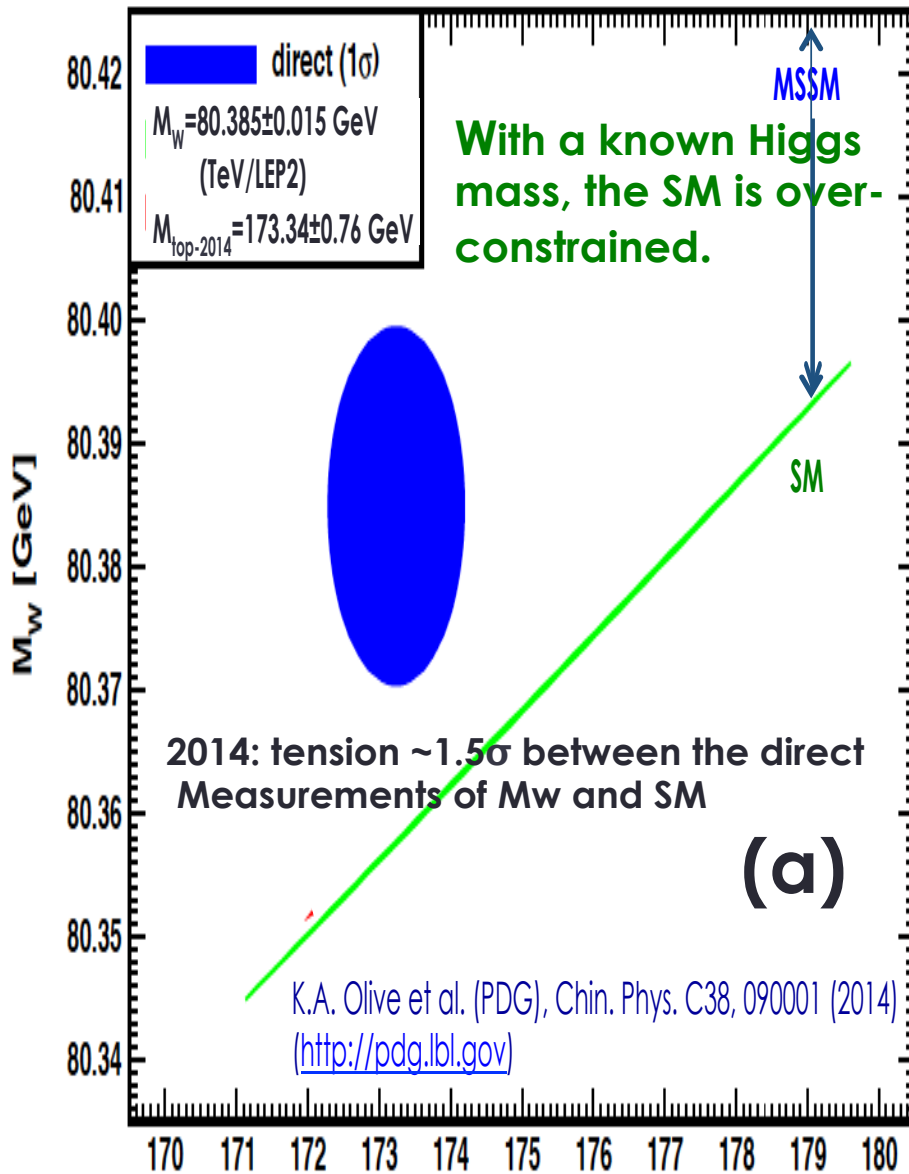
# Direct measurement of W mass LEP & Tevatron 2

<http://pdg.lbl.gov/2014/reviews/rpp2014-rev-w-mass.pdf>



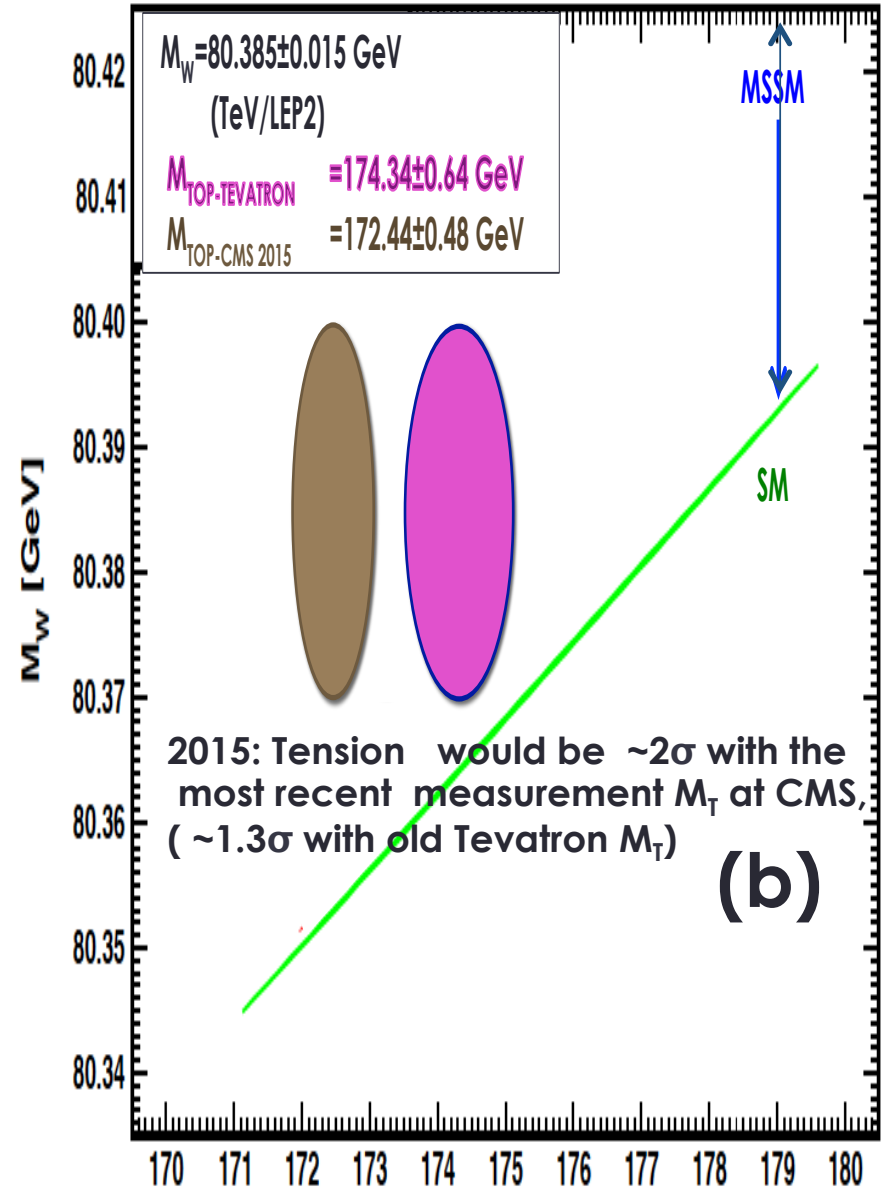
Most recent Tevatron  $M_W^{\text{direct}}$  measurements (CDF and Dzero)

$$\Delta M_W^{\text{direct}} \sim 20 \text{ MeV}$$



$m_t$  [GeV]

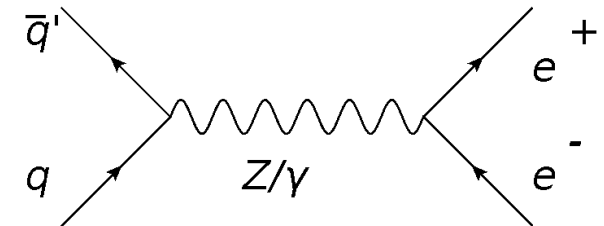
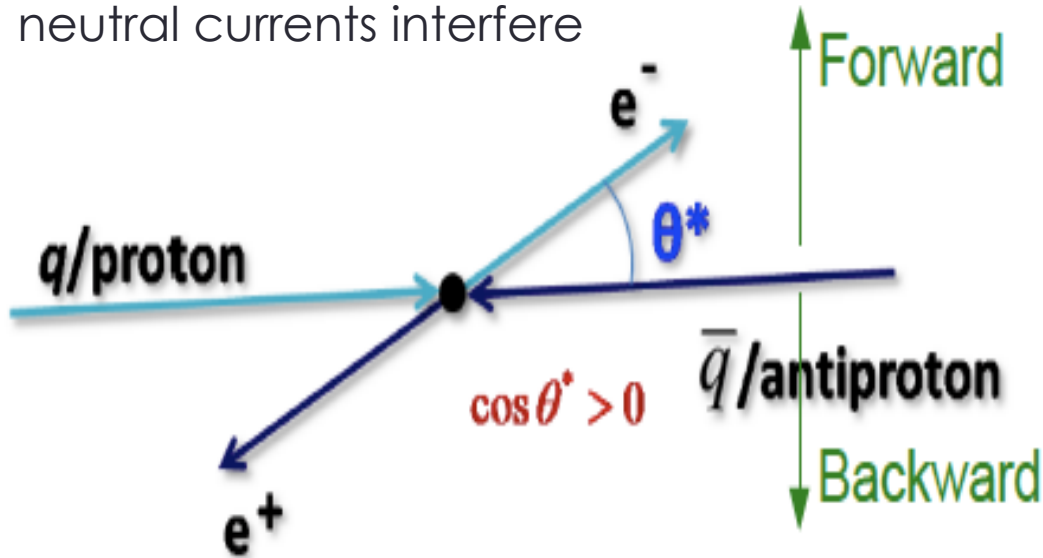
15 MeV error in W mass



$m_t$  [GeV]

$A_{FB}$  for  $e^+e^-$  or  $\mu^+\mu^-$  pairs in the Z boson Region is sensitive to the effective EW mixing angle  $\sin^2\theta_{eff}$

The axial and vector neutral currents interfere



Define Forward-Backward asymmetry:

$$A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B}$$

$$\sin^2\theta_{eff}^{lept} \approx 1.037 \cdot \sin^2\theta_w \quad [ \text{ZFITTER } \kappa_e(\sin^2\theta_w, M_Z) \text{ form factor} ]$$

(above relation is approximate) one needs to include complex EW radiative correction form factors in the theory predictions for  $A_{FB}$  to extract the on-shell  $\sin^2\theta_w$

$$\sin^2\theta_w = 1 - M_w^2 / M_Z^2$$

$M_W$  can also be determined indirectly via the relation

$$\sin^2\theta_W^{\text{on-shell}} = 1 - M_W^2 / M_Z^2$$

$\pm 0.00040$  error in  $\sin^2\theta_W$  is equiv. to  $\pm 20$  MeV error in  $M_W$  (indirect)

Both  $\sin^2\theta_W^{\text{on-shell}}$  and  $\sin^2\theta_{\text{eff}}^{\text{leptonic}} (M_Z)$  can be extracted from Drell-Yan forward-backward asymmetry (Afb) if we include EW radiative corrections.  $M_W^{\text{indirect}}$  can be extracted from  $\sin^2\theta_W^{\text{on-shell}}$

- If the SM is correct, then measurements of  $M_W^{\text{direct}}$  and  $M_W^{\text{indirect}}$  should agree. Deviations may imply the possibility of new physics.
- Similarly different measurements of  $\sin^2\theta_{\text{eff}}^{\text{leptonic}} (M_Z)$  should also agree and deviations may imply new physics.

We are now able to minimize experimental systematic errors, as shown in the next few slides

- Therefore, PDF uncertainties limit the precision of the measurements of  $\sin^2\theta_W^{\text{eff}}$ ,  $\sin^2\theta_W^{\text{on-shell}}$ ,  $M_W^{\text{indirect}}$  and  $M_W^{\text{direct}}$  at the LHC.
- We will show that Drell-Yan Afb provide further constraints on PDFs

# Tevatron $ee$ & $\mu\mu$ $9 \text{ fb}^{-1}$ : $\sin^2\theta_{\text{eff}}(M_Z)$ 6

$D\phi ee$

$$\sin^2\theta_{\text{eff}} = 0.23146 \pm 0.00043(\text{stat})$$

$$\pm 0.00008(\text{syst})$$

$$\pm 0.00017(\text{NNPDF2.3 PDFs NLO})$$

$$\sin^2\theta_{\text{eff}} = \mathbf{0.23146 \pm 0.00047 \text{ (total)}}$$

**CDF  $ee$  &  $\mu\mu$**

$$\sin^2\theta_{\text{eff}} = 0.23221 \pm 0.00043(\text{stat})$$

$$\pm 0.00005(\text{syst})$$

$$\pm 0.00016(\text{NNPDF3.0 PDFs NNLO})$$

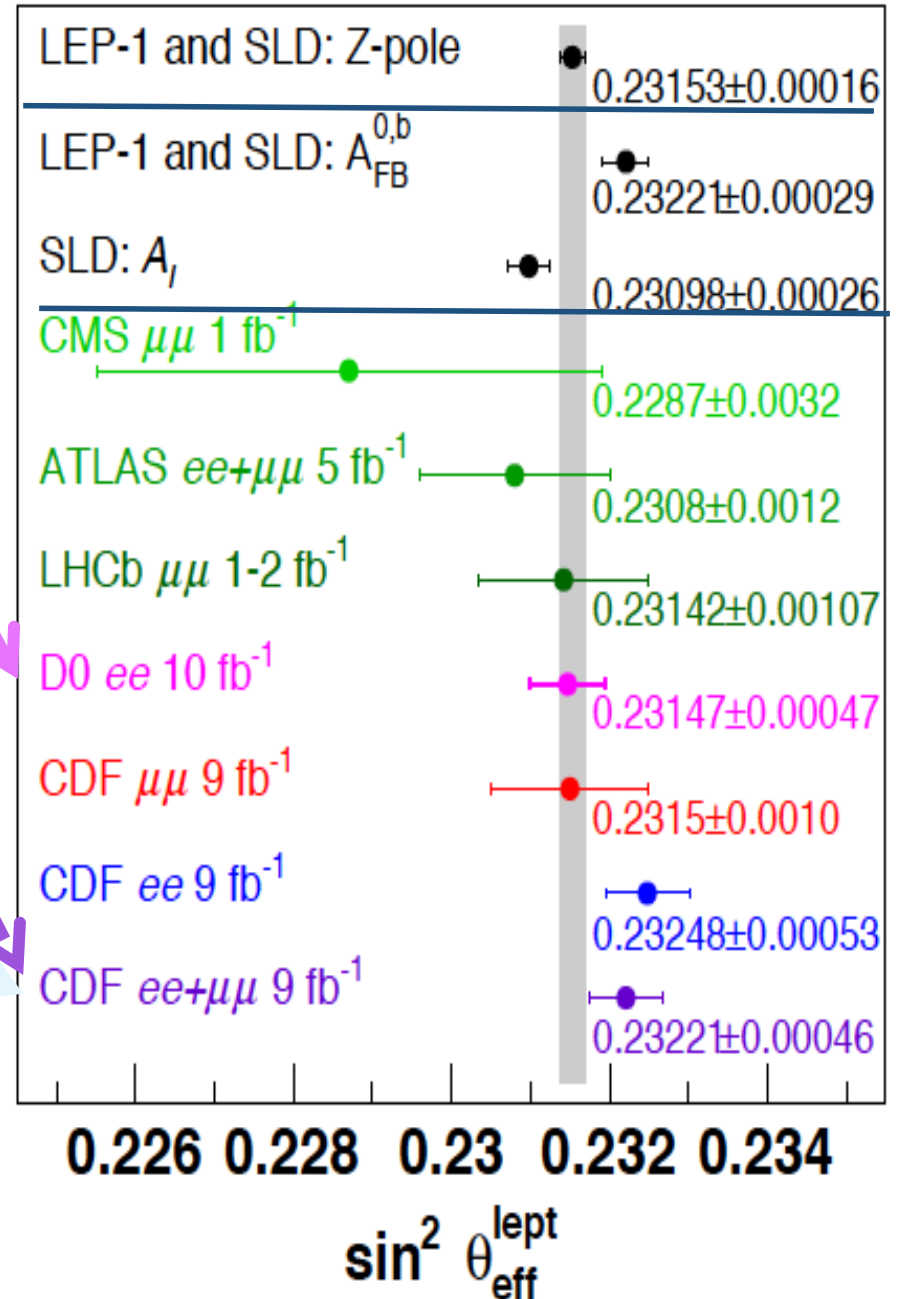
$$\sin^2\theta_{\text{eff}} = \mathbf{0.23221 \pm 0.00046 \text{ (total)}}$$

PDF errors smaller at the Tevatron than at the LHC.

$$\mathbf{ATLAS \sin^2\theta_{\text{eff}} = 0.23080 \pm 0.00050(\text{stat})}$$

$$\pm 0.00060(\text{syst}) \pm \mathbf{0.00090(\text{pdf})}$$

Need to greatly reduce PDF errors at LHC (by a factor of 4 with respect to ATLAS)



# CDF $ee$ & $\mu\mu$ $9 \text{ fb}^{-1}$ Indirect $M_W$ measurement

7

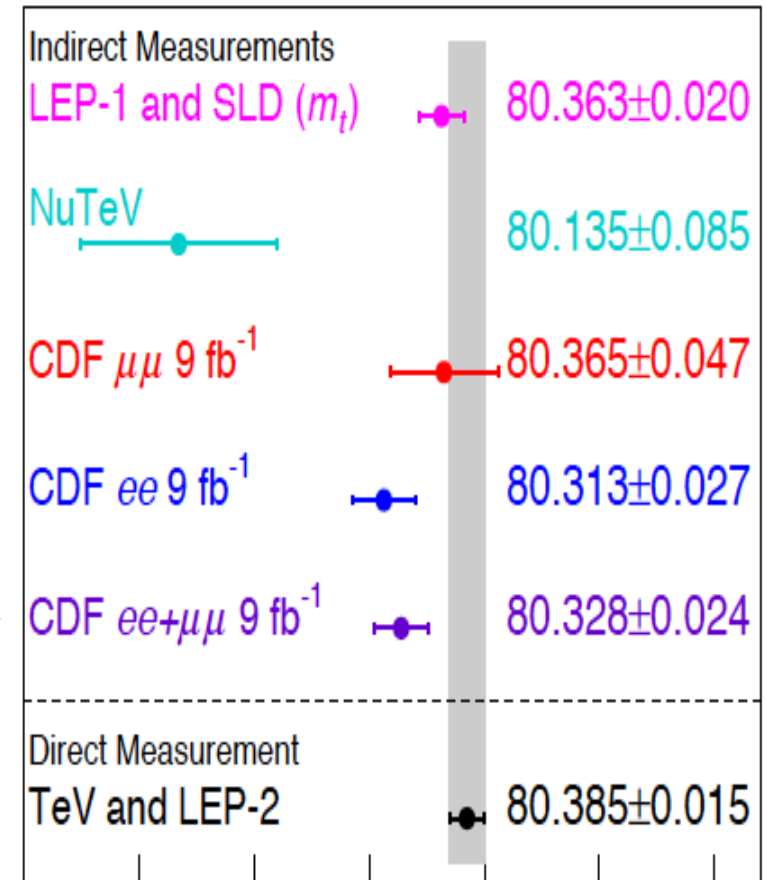
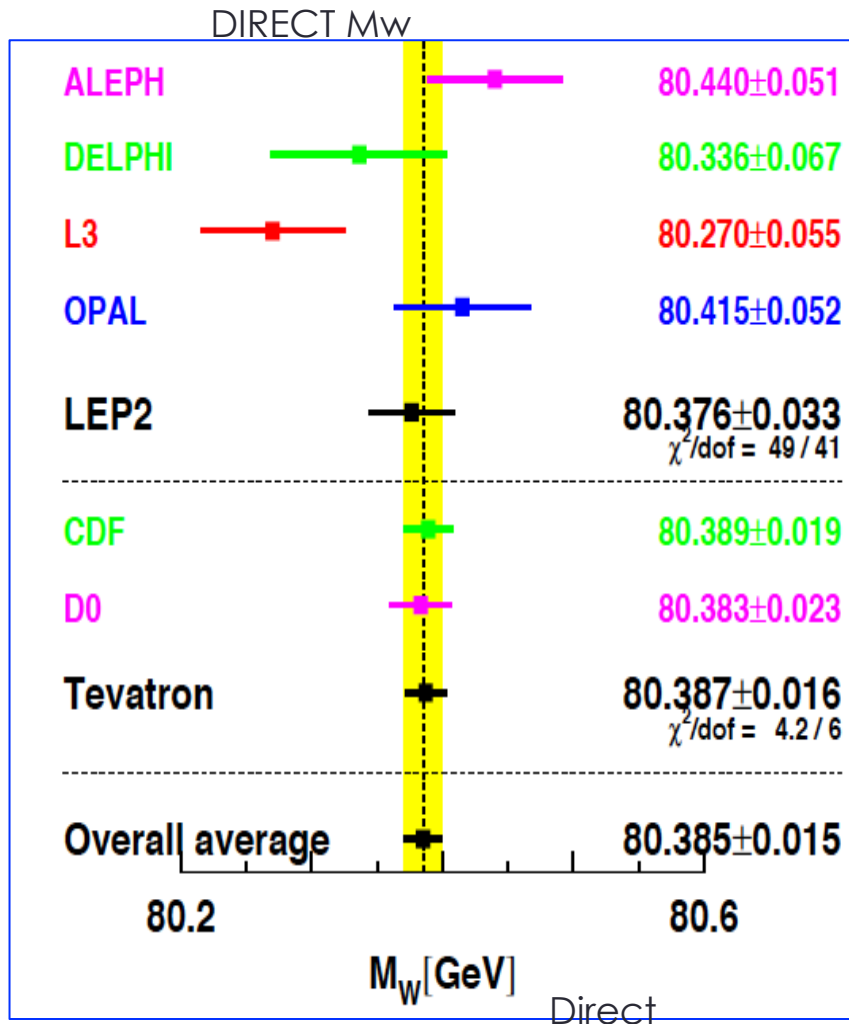
Aim at 10 MeV error in both direct and indirect  $M_W$  measurement at the LHC.

$$\sin^2 \theta_{\text{eff}}^{\text{lept}} = 0.23221 \pm 0.00046$$

$$\sin^2 \theta_W = 0.22400 \pm 0.00045 \quad \text{On-shell}$$

$$M_W(\text{indirect}) = 80.328 \pm 0.024 \text{ GeV}/c^2.$$

INDIRECT  $M_W$



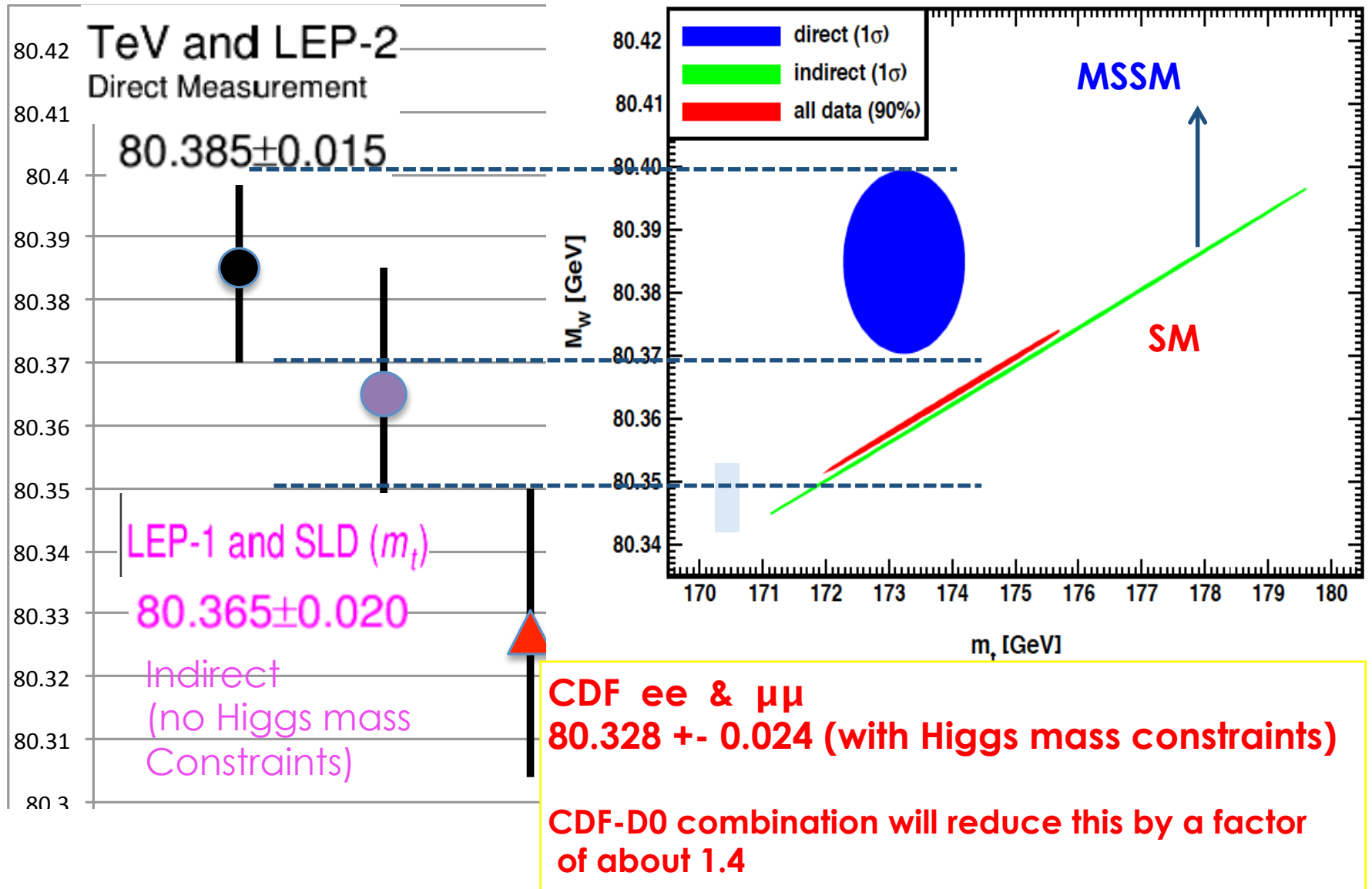
80 80.1 80.2 80.3 80.4 80.5 80.6  
W-boson mass ( $\text{GeV}/c^2$ )

<http://pdg.lbl.gov/2014/reviews/rpp2014-rev-w-mass.pdf>

# CDF $ee$ & $\mu\mu$ $9 \text{ fb}^{-1}$ Indirect $M_W$ measurement 8

<http://pdg.lbl.gov/2014/reviews/rpp2014-rev-standard-model.pdf>

K.A. Olive et al. (PDG), Chin. Phys. C38, 090001 (2014) (<http://pdg.lbl.gov>)





**Systematic errors from momentum and energy calibration are now very small**

New technique used for both  $\mu^+\mu^-$  and  $e^+e^-$  for both data and hit level MC.  
( Ref A. Bodek et al. Euro. Phys. J. C72, 2194 (2012))

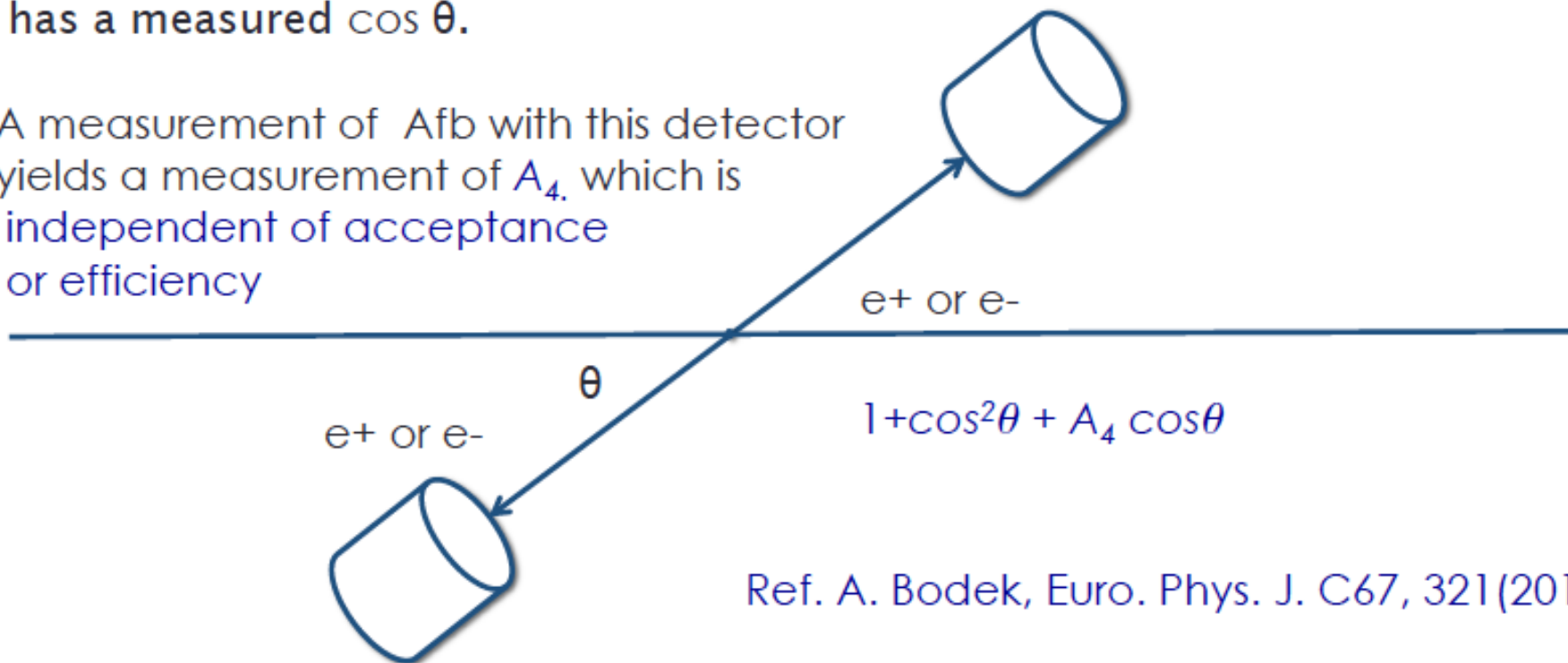
**Step 1 : Remove the correlations between the scale for the two leptons** by getting an initial calibration using Z events and requiring that the **mean  $\langle 1/P_T \rangle$**  of each lepton in bins of  $\eta, \Phi$  and charge be correct.

**Step2: The Z mass used as a calibration.** The Z mass as a function of  $\eta, \Phi$ , (and charge for  $\mu^+\mu^-$ ) of each lepton be correct

- **Reference for muons:** Expected Z mass (post FSR) smeared by resolution (with acceptance cuts).
- **Reference for electrons:** Expected Z mass (post FSR + clustered FSR photons), smeared by resolution (with acceptance cuts).
- Now Used in CDF, CMS. A similar technique is used in Dzero.

Imagine a detector with acceptance for only one value of  $\cos \theta$ . Each event has a measured  $\cos \theta$ .

A measurement of  $A_{fb}$  with this detector yields a measurement of  $A_4$ , which is independent of acceptance or efficiency



Ref. A. Bodek, Euro. Phys. J. C67, 321(2010)

$\cos \theta = 1$  yields best measurement of  $A_4$ .  $\cos \theta = 0$  yields no measurement of  $A_4$

We can combine measurements of  $A_4$  with different detectors at different  $\cos(\theta)$  by weighting events. Events with  $\cos(\theta) = 0$  have zero weight.

Events with  $\cos \theta = 1$  have maximum weight.  $\rightarrow$  obtain smaller statistical error.

$A_{fb}(\text{all } \cos \theta) = (3/8) A_4 \rightarrow$  No acceptance corrections needed.

**Systematic errors from acceptance and efficiencies are now very small**

New technique: event weighting method for  $A_{\text{FB}}$  analyses

Ref. A. Bodek, Euro. Phys. J. C67, 321 (2010)

$$dN/d\cos\theta = 1 + \cos^2\theta + A_0(M, P_T) (1 - 3\cos^2\theta)/2 + A_4(M) \cos\theta$$

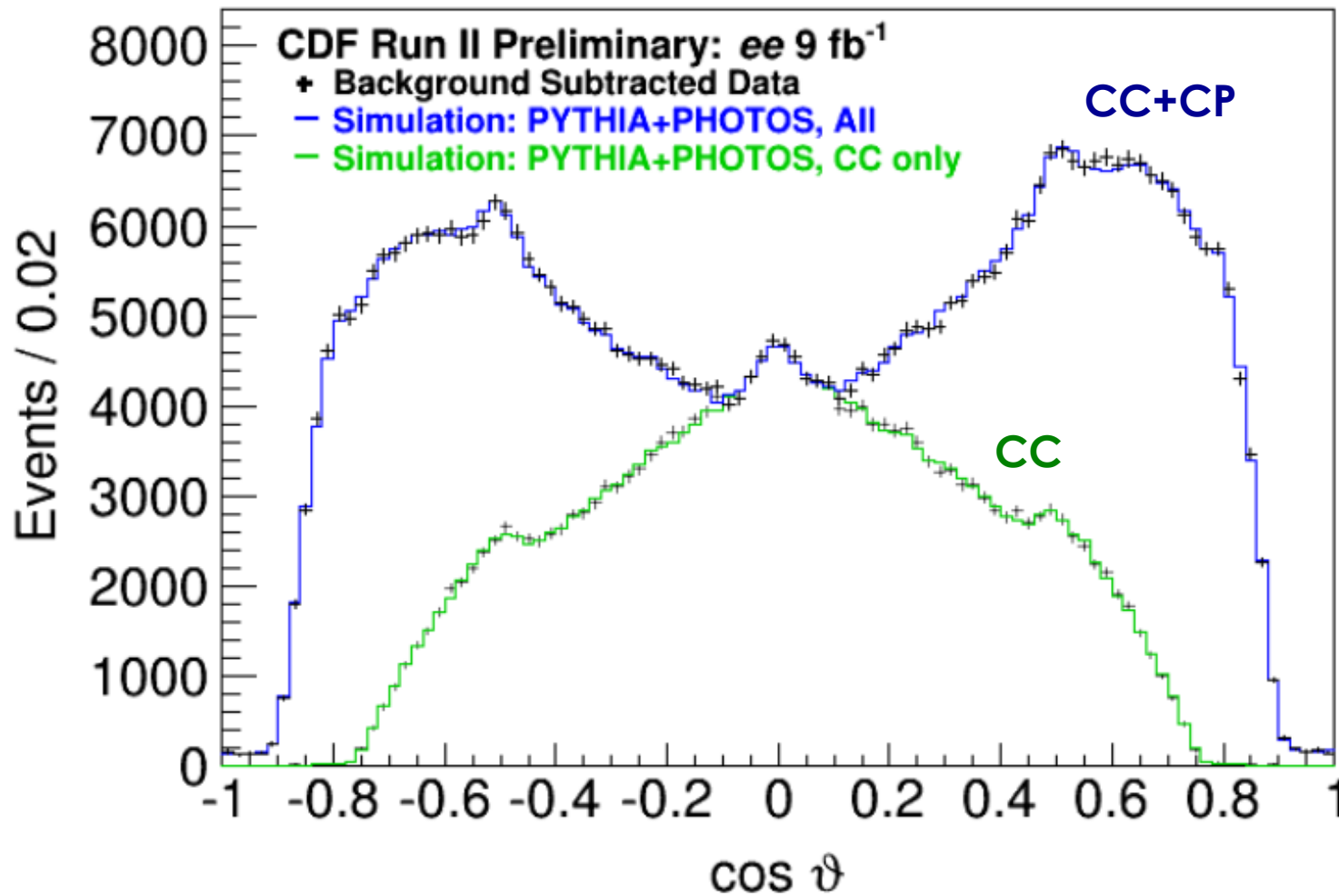
Angular event weighting is equivalent to extraction of  $A_4(M)$  in bins of  $\cos\theta$ , and averaging the results.

Events at large  $\cos\theta$  provide better determination of  $A_4$ , so they are weighted more than events at small  $\cos\theta$ .

For each  $\cos\theta$  acceptance and efficiencies cancel to first order in the measurement of  $A_4$ . Then set  $A_{\text{fb}} = (3/8)A_4$

Statistical errors with event weighting are 20% smaller. Systematic errors largely cancel.

Event weighting does not correct for resolution smearing and final state radiation, which are included later in the unfolding using  $M$ .



The error in  $A_{FB}$  is reduced if we have more acceptance at large  $\cos \theta$ ,  
 Standard  $A_{FB}$  method requires precise knowledge of acceptance and efficiencies.

Measure  $A_4 \rightarrow A_{FB}$

PDF fits have a limited number of parameters and therefore cannot precisely describe PDFs in all regions of  $x$  and  $Q^2$  at the same time. In addition, sometimes experimental errors are inflated to account for inconsistencies between experiments at different regions of  $x$  and  $Q^2$ .

- For  $M_W$  extraction from  $W$  data and  $\sin^2\theta_{\text{eff}}$  extraction from  $Z$  data we are mainly interested in precise PDFs in the  $x$  and  $Q^2$  regions relevant for the production of  $W$  and  $Z$  for a particular CM energy.

For the Tevatron only PDFs which are consistent (within errors) with  $W$  and  $Z$  production data at CDF or D0 are relevant.

For LHC only PDFs which are consistent (within measurement errors) with  $W$  and  $Z$  production data at the LHC are relevant.

- We can further constrain the PDF using new precise LHC data that has not been used in the published PDF fits. (e.g. 8 and 14 TeV  $W$  asymmetry data).
- In addition, we have developed a way to include Drell-Yan Afb data to further constrain PDFs in this important  $x$  and  $Q^2$  region. (Afb data have never been included in any PDF fits before).

From: Parton Distribution Uncertainty in the Measurement of  $M_W$  in Proton - Anti-proton Collisions, W.James Stirling, Alan D. Martin (Durham U.). Phys.Lett. B237 (1990) 551, DOI: [10.1016/0370-2693\(90\)91223-X](https://doi.org/10.1016/0370-2693(90)91223-X)

**At the Tevatron**, for valence-valence W production (for  $P_T=0$  W's), when integrated over all rapidity, the  $P_T$  distribution of the lepton from W decay is given by an expression which is independent of PDFs.

$$\frac{dN^{e\tau}}{d \cos \theta^*} = \frac{3}{8} (1 \pm \cos \theta^*)^2 \quad \cos \theta^* = \left(1 - \frac{4p_T^2}{M_W^2}\right)^{1/2}$$

$$\frac{1}{\sigma} \frac{d\sigma}{dp_T^2} = \frac{3}{2M_W^2} \left(1 - \frac{4p_T^2}{M_W^2}\right)^{-1/2} \left(2 - \frac{4p_T^2}{M_W^2}\right) \quad \rightarrow \text{f independent of PDFs}$$

For full acceptance

However, if the detector has limited lepton rapidity coverage ( $R$ ) the expression is modified:

$$\frac{1}{\sigma} \frac{d\sigma}{dp_T^2} = \frac{3}{2M_W^2} \left(1 - \frac{4p_T^2}{M_W^2}\right)^{-1/2} \left(2 - \frac{4p_T^2}{M_W^2}\right) F(p_T) :$$

Where  $F(P_T)$  is PDF dependent

$$\frac{1}{\sigma} \frac{d\sigma}{dp_T^2} = \frac{3}{2M_W^2} \left(1 - \frac{4p_T^2}{M_W^2}\right)^{-1/2} \left(2 - \frac{4p_T^2}{M_W^2}\right) F(p_T),$$

$$F(p_T) = \frac{1}{\sigma} \int_R dy \frac{d\sigma}{dy} + \left(1 - \frac{4p_T^2}{M_W^2}\right)^{1/2} \left(1 - \frac{2p_T^2}{M_W^2}\right)^{-1} \frac{1}{\sigma} \int_R dy \left(\frac{d\sigma^+}{dy} - \frac{d\sigma^-}{dy}\right)$$

At the Tevatron, for the valence-valence process at  $PT=0$ , precision measurements of the  $W$  lepton charge asymmetry reduce the PDF error in the predicted transverse momentum distributions

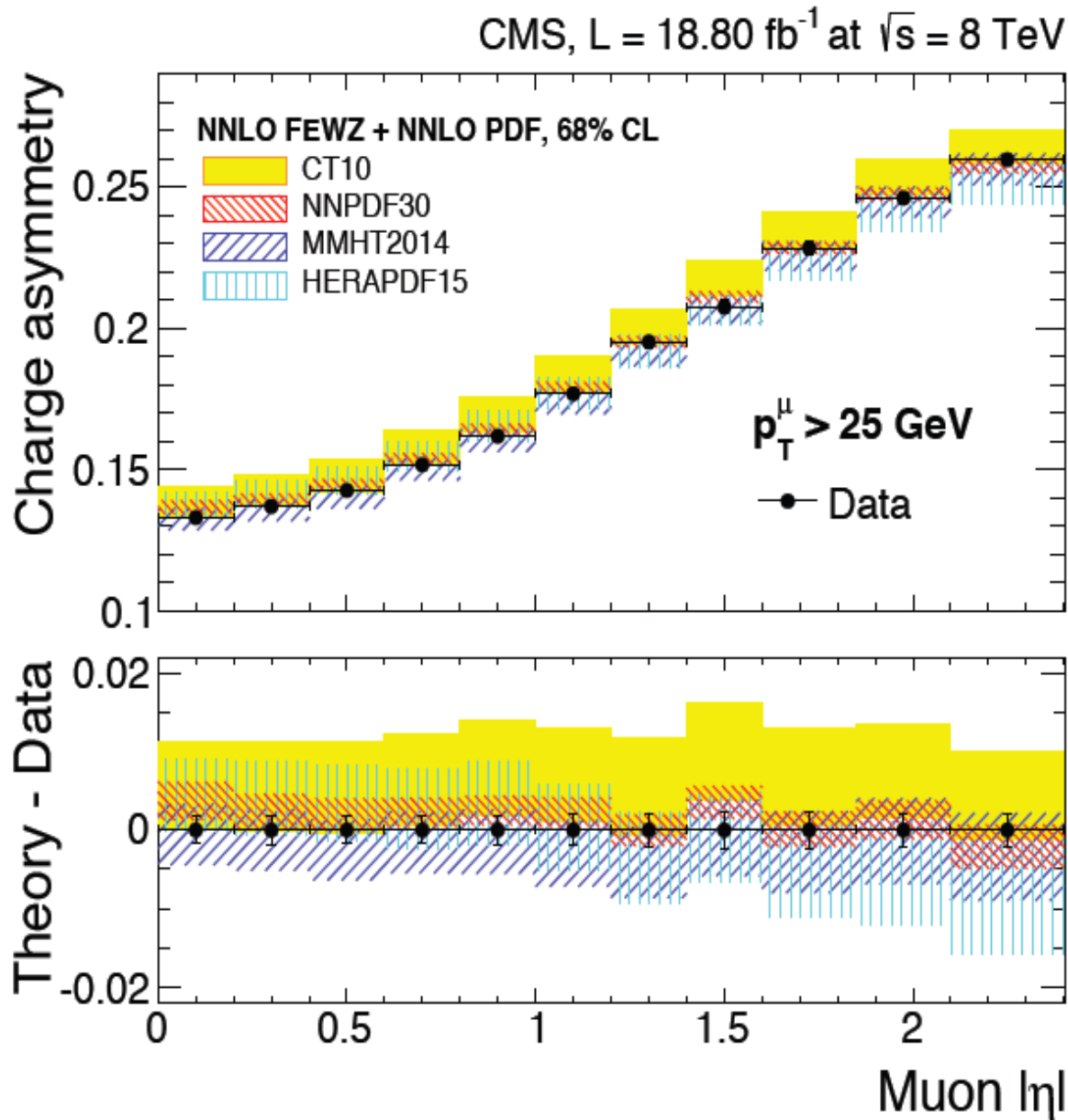
-> consequently,  $A_W$  helps to reduce PDF errors in measurement of  $M_W$ .

Only PDFs which are consistent with **W lepton charge asymmetry measurements at the Tevatron should be considered.**

$$\left(\frac{W^-}{W^+}\right)_{Tev} \approx \frac{d(x_1) u(x_2) + s(x_1) c(x_2)}{u(x_1) d(x_2) + c(x_1) s(x_2)} \approx \frac{d}{u}(x_1) / \frac{d}{u}(x_2)$$

Any other data which constrains the quark PDFs in the  $W$  and  $Z$  region would also help eg. Drell-Yan Afb, (and to a lesser extent  $W$  decay lepton differential cross sections  $d\sigma/d\eta$ ,  $Z$  rapidity distributions,  $W/Z$  cross section ratios etc).

However, asymmetries can have more impact because of smaller experimental systematic errors (and less sensitivity to QCD scale and higher order terms).



8 TeV CMS W asym  
Data further constrains  
NNPDF 3.0 PDFs (red)  
which have included  
CMS 7 TeV data.

Very few CT10 replicas  
would be consistent  
with the CMS 8 TeV W  
Asymmetry data  
(no LHC data included  
in CT10).



Source uncertainty in CDF measurement with 9.4 fb <sup>-1</sup> (electron+muons)	Uncertainty in $\sin^2 \theta_{\text{eff}}^{\text{lept}}$ ( $M_Z$ )	Uncertainty in $M_W^{\text{indirect}}$ (GeV)	Uncertainty in $M_W$ direct (2012) 2.2 fb <sup>-1</sup> (CDF). MeV
Data: Statistics	$\pm 0.00042$ (stat)	0.020	0.012 (stat)
Data: Energy scale	$\pm 0.00003$ (syst)	0.001	0.007 (lepton scale)*
Data: Backgrounds	$\pm 0.00002$ (syst)	0.001	0.003 (Background)
Prediction: PDFs	$\pm 0.00016$ (syst)	0.008	0.010 (PDFs)**
Prediction: QCD EBA (NLO minus LO)	$\pm 0.00007$ (syst)	0.003	0.006 (Recoil energy)
Prediction: QCD scales	$\pm 0.00002$ (syst)	0.001	0.005 (PT W)
			0.004 (QED rad)
All systematics	$\pm 0.00018$ (syst)	0.009	0.015
Total: (stat+syst)	$\pm 0.00046$ (total)	0.023	0.019

$$\sin^2 \theta_{\text{eff}}^{\text{lept}} = 0.23221 \pm 0.00046$$

On-shell

$$\sin^2 \theta_W = 0.22400 \pm 0.00045$$

$$M_W(\text{indirect}) = 80.328 \pm 0.024 \text{ GeV}/c^2 .$$

The PDF errors include constraints from ATb data

(which reduce PDF errors in  $\sin^2 \theta_W$  from 0.00020 to 0.00015)

\*Will be reduced using new lepton energy calibrations  
 \*\* Will be reduced with LHC based PDFs

# Reducing PDF errors for $\sin^2\theta_{\text{eff}}$ at CDF and LHC 18

We use the CDF combined  $e+e$   $\mu^+\mu$  Afb CDF data to constrain PDFs

[Ref : A. Bodek, J. Han, A. Khukhunaishvili, W. Sakumoto:](#) "Using Drell-Yan forward-backward asymmetry to constrain parton distribution functions"  
arXiv:1507.02470. (**Eur.Phys.J. C76 (2016) no.3, 115**)

All PDF groups provide a default (central) PDF set. There are two methods that are used for the determination of PDF uncertainties.

**Hessian Matrix:** Use a set of eigenvector error PDFs The PDF uncertainties in a measurement are determined by repeating the analysis for all of the error PDF sets, and adding in quadrature the difference in the result obtained with the error PDFs and the result obtained with the default PDF.

**Monte Carlo Replicas:** Use a set of  $N$  (e.g. 100 or 1000) replica PDFs. Each of the PDF replicas has equal probability of being correct. The central value of any observable is the average of the values of  $\sin^2\theta_{\text{eff}}$  extracted with each one of the  $N$  PDF replicas. The PDF error is the RMS of the values extracted using all  $N$  replicas.

- For every one of the 100 NNPDF3.0 replicas find the  $\sin^2\theta_{\text{eff}}$  that gives the best fit. (Replicas can also be generated for any PDF set).
- Without any additional constraints, all 100 PDF sets of NNPDF3.0 (or any other set) are equally likely.
- Therefore, the **average** value of the 100 determination of  $\sin^2\theta_{\text{eff}}$  is the measurement.
- The RMS of the  $\sin^2\theta_{\text{eff}}$  values from all 100 PDF replicas is the PDF error.

The same procedure can be used to determine the PDF errors in the measurement of  $M_w$ .

**Monte Replica Method:**

$$\langle s \rangle = \frac{1}{N} \sum_{i=1}^N s_i \quad (12) \quad s = \sin^2\theta \text{ or } M_w.$$

$$\sigma_{pdf} = \sqrt{\frac{\sum_{i=1}^N (s_i - \langle s \rangle)^2}{N - 1}} \quad (13)$$

and the uncertainty in the PDF error is  $\Delta\sigma_{pdf} = \frac{\sigma_{pdf}}{\sqrt{2(N-1)}}$

The calculated standard PDF errors will be the same for both Hessian and Bayesian (replica) methods

For any given a set of Hessian eigenvector PDFs there is a prescription to generate an arbitrary number of PDF replicas.

The replica method is preferable for two reasons:

1. We can easily add constraints from new data.
2. We can easily find if the new data is consistent or inconsistent with the PDFs.

We use 100 NNPDF3.0 NNLO PDFs (these included LHC data) in the CDF Analysis of  $\sin^2\theta_w$ .

For these 100 replicas: CDF gets  $\sin^2\theta_w = 0.22401 \pm 0.00042$  (Statistical error)  
RMS is the PDF error is  $\pm 0.00020$  (PDF)

- New measurements can be incorporated into the ensemble without refits
  - Ensemble PDFs are reweighted

$$W_k = \frac{\exp(-\chi_k^2/2)}{\sum_{l=1}^N \exp(-\chi_l^2/2)}$$

The new result = weighted mean.  
The new weighted RMS is a reduced PDF error

$\chi_{k,l}^2$ : between new measurement and prediction with ensemble PDF k

**It is clear how to do this for new results for processes that have been used in previous PDF fits. (e.g. new LHC W asymmetry data)**

We now show how we can get both  $\sin^2\theta_w$  AND also further constrain PDFs from the same  $A_{FB}(M)$  data

A. [Bodek, J. Han, A. Khukhunaishvili, W. Sakumoto](#): arXiv:1507.02470 (Eur.Phys.J. C76 (2016) no.3, 115)

18. G. Watt and R. S. Thorne (MRST), JHEP 08:052 (2012) (arXiv:1205.4024)
19. <https://mstwpdf.hepforge.org/random/>
20. Walter T. Giele, and Stephane Keller, Phys.Rev. D58 (1998) 094023 (arXiv:hep-ph/9803393).
21. Nobuo Sato, J. F. Owens, Harrison Prosper, Phys. Rev. D 89, 114020 (2014) (arXiv:1310.1089)
22. Hannu Paukkunen, Pia Zurita, "PDF reweighting in the Hessian matrix approach", <http://arxiv.org/abs/1402.6623>
23. Richard D. Ball, Valerio Bertone, Francesco Cerutti, Luigi Del Debbio, Stefano Forte, Alberto Guffanti, Jose I. Latorre, Juan Rojo, Maria Ubiali, Nucl.Phys.B849, 112 (2011) arXiv:1012.0836.

Here  $s$  is a measurement of  $\sin^2\theta_w$  (or  $M_w$ )

$$w_i = \frac{e^{-\frac{1}{2}\chi_i^2}}{\sum_{i=1}^N e^{-\frac{1}{2}\chi_i^2}} \quad (14)$$

$w_i$  is the weight for each PDF.

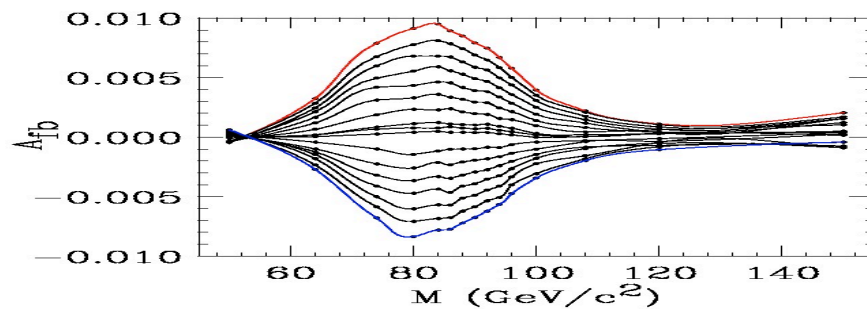
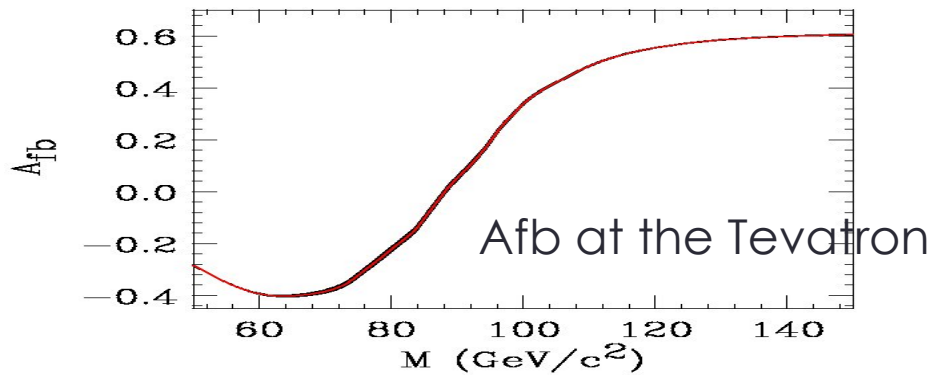
$$\langle s \rangle = \sum_{i=1}^N w_i s_i \quad (15)$$

$$\sigma_{pdf} = \sqrt{\frac{\sum_{i=1}^N w_i (s_i - \langle s \rangle)^2}{1 - 1/N_{eff}}} \quad (16)$$

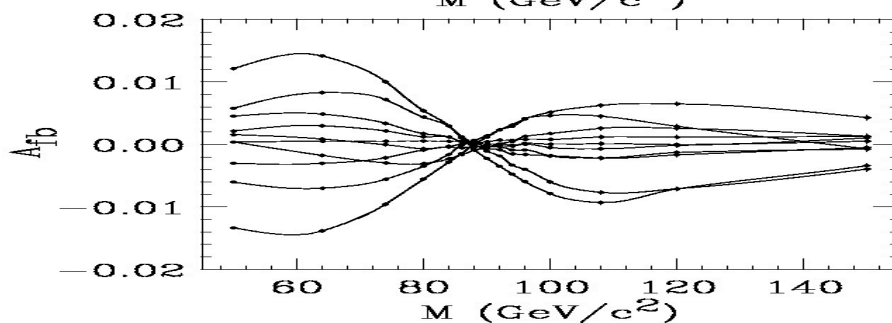
The weights reduce the effective number of replicas from  $N$  to  $N_{eff}$  where

$$N_{eff} = \frac{1}{\sum_{i=1}^N w_i^2} \quad (17)$$

and the uncertainty in the PDF error is  $\Delta\sigma_{pdf} \approx \frac{\sigma_{pdf}}{\sqrt{2(N_{eff}-1)}}$ .



lines are variation from different  $\sin^2\theta_{\text{eff}}$ .  $\Delta\sin^2\theta_{\text{eff}}$  changes  $A_{fb}$  mostly in the region of the Z pole.



Lines are variation from different PDFs (NNPDF3.0).

Different PDFs change  $A_{fb}$  mostly above and below the region of the Z peak. (different  $d/u$ , different antiquarks).

To extract  $\sin^2\theta_{\text{eff}}$  we change  $\sin^2\theta_{\text{eff}}$  till we get the best fit to  $A_{fb}(M)$ . However, if we use an incorrect PDF it will give a poor fit.



**Constraints from W lepton Charge asymmetry at the Tevatron**  
*sensitive to slope of d/u* (x range at the Tevatron)

$$\left(\frac{W^-}{W^+}\right)_{Tev} \approx \frac{d(x_1) u(x_2) + s(x_1) c(x_2)}{u(x_1) d(x_2) + c(x_1) s(x_2)} \approx \frac{d}{u}(x_1) / \frac{d}{u}(x_2) \quad (11)$$

**Constraints from Afb at the Tevatron (x range at the Tevatron)**  
 (Afb is diluted by the fraction of antiquarks and the d/u ratio)  
*Sensitive to the absolute values of d/u and ubar/u,*

$$D_{AFB}^{Tev}(d) \propto \left[\frac{d}{u}(x_1)\right]^2 \quad (9)$$

$$D_{AFB}^{Tev}(\bar{q}) \propto \left[\frac{\bar{u}}{u}(x_1)\right]^2 \quad (10)$$

W asym data at CDF reduces PDF errors for Mw measurement but do not reduce PDF errors in the measurement of  $\sin^2\theta_{eff}$ .

**However,** Afb data reduces errors in both Mw and  $\sin^2\theta_{eff}$  measurements.



$\chi^2_{\min}$  versus  $\sin^2\theta_W$  for each ensemble PDF:

Weighted Mean  $\rightarrow \sin^2\theta_W = 0.22401 \pm 0.00042$   
 Weighted RMS = **reduced PDF error =  $\pm 0.00015$**   
 (from  $\pm 0.00020$ )

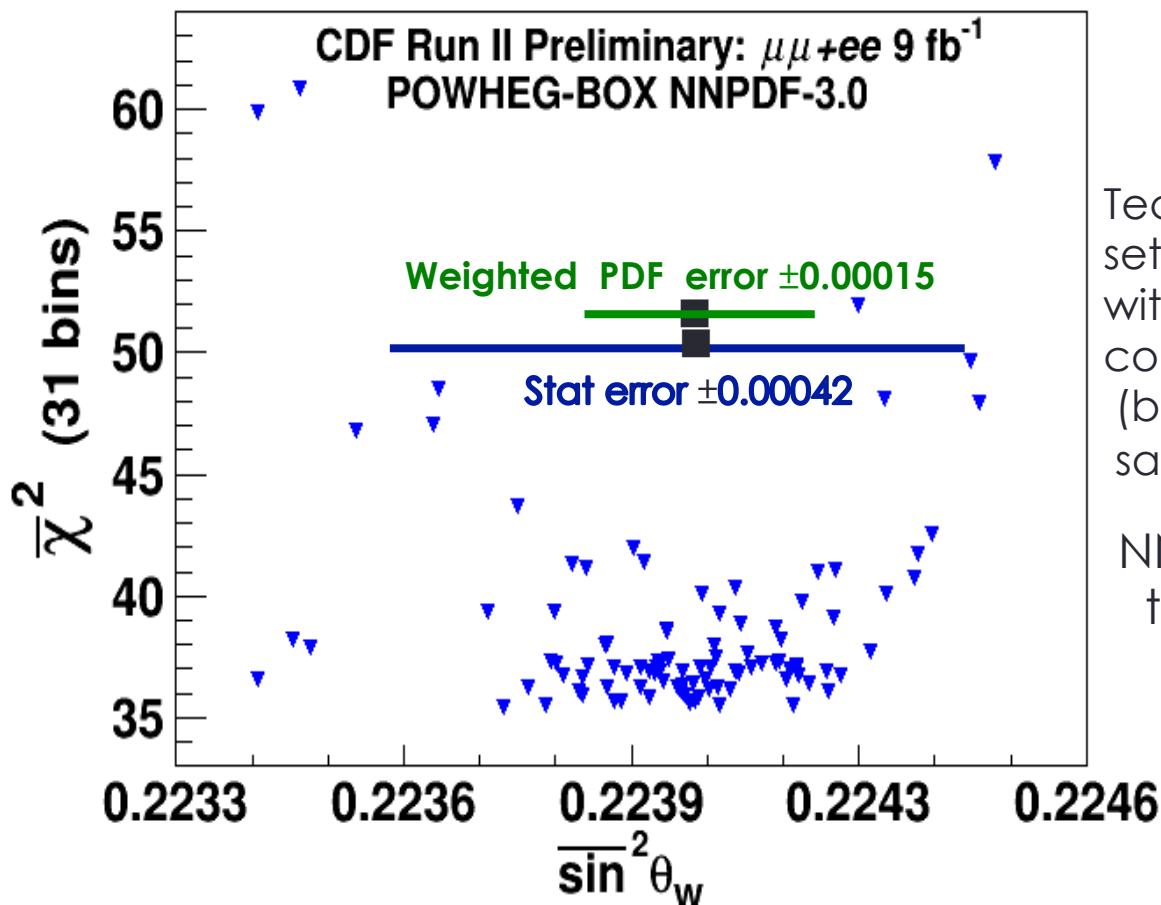
Ensemble PDFs are constrained by reweighting

$$W_k = \frac{\exp(-\chi_k^2/2)}{\sum_{l=1}^N \exp(-\chi_l^2/2)}$$

Technique can be used with any PDF set provided the PDF set is consistent with the new data. If the PDF sets are consistent with each other the result (but not the PDF error) will be the same.

NNPDF3.0 NNLO is consistent with the CDF  $A_{FB}(M)$  data

These constrained PDFs can be used for other analyses (e.g. direct measurement of the W mass)



# Toy Study of PDF Systematics with CMS pseudo-data 26

- PDFs are the dominant uncertainty in the measurement of  $\sin^2\theta_W$  at the LHC
- Mainly come in as the uncertainty in the dilution of Afb.

**ATLAS result:**  $0.23080 \pm 0.00050(\text{stat}) \pm 0.00060(\text{syst}) \pm 0.00090(\text{pdf})$

(Now: Do a Toy MC study with CMS like detector)

**Pseudo data with CMS like detector 8 TeV 15 fb<sup>-1</sup> sample  
(7 M reconstructed  $\mu\mu$  events)  $\rightarrow$  expected error in  $\sin^2\theta_{\text{eff}}$ :**

Expected Stat Error:  $\pm 0.00043$

PDF errors with CT10:  $\pm 0.00077$

PDF errors with NNPDF3.0:  $\pm 0.00051 \rightarrow \pm 0.00026$  the new approach

[A. Bodek, et al arXiv:1507.02470](#) asymmetry to constrain parton distribution functions" arXiv:1507.02470. (**Eur.Phys.J. C76 (2016) no.3, 115**)

Actual CMS sample at 8 TeV is 19 fb<sup>-1</sup> sample

(8.2 M reconstructed  $\mu\mu$  events and 6.8 M reconstructed  $ee$  events)

Constraints from W lepton Charge asymmetry at the LHC  
*(sensitive to absolute value of d/u)*

$$\begin{aligned} \left(\frac{W^-}{W^+}\right)_{LHC} &\approx \frac{d(x_1) \bar{u}(x_2) + s(x_1) \bar{c}(x_2)}{u(x_1) \bar{d}(x_2) + c(x_1) \bar{s}(x_2)} \\ &\approx \frac{d/u(x_1)}{\bar{d}/\bar{u}(x_2)} \approx \frac{d}{u}(x_1) \end{aligned}$$

Constraints from Afb at the LHC (x range of LHC)

Afb is diluted by the fraction of antiquarks and the d/u ratio)

**( sensitive to the absolute values of d/u and ubar/u)**

$$D_{AFB}^{LHC}(d) \propto \frac{d(x_1)\bar{d}(x_2)}{u(x_1)\bar{u}(x_2)} \approx \frac{d}{u}(x_1) \quad (20)$$

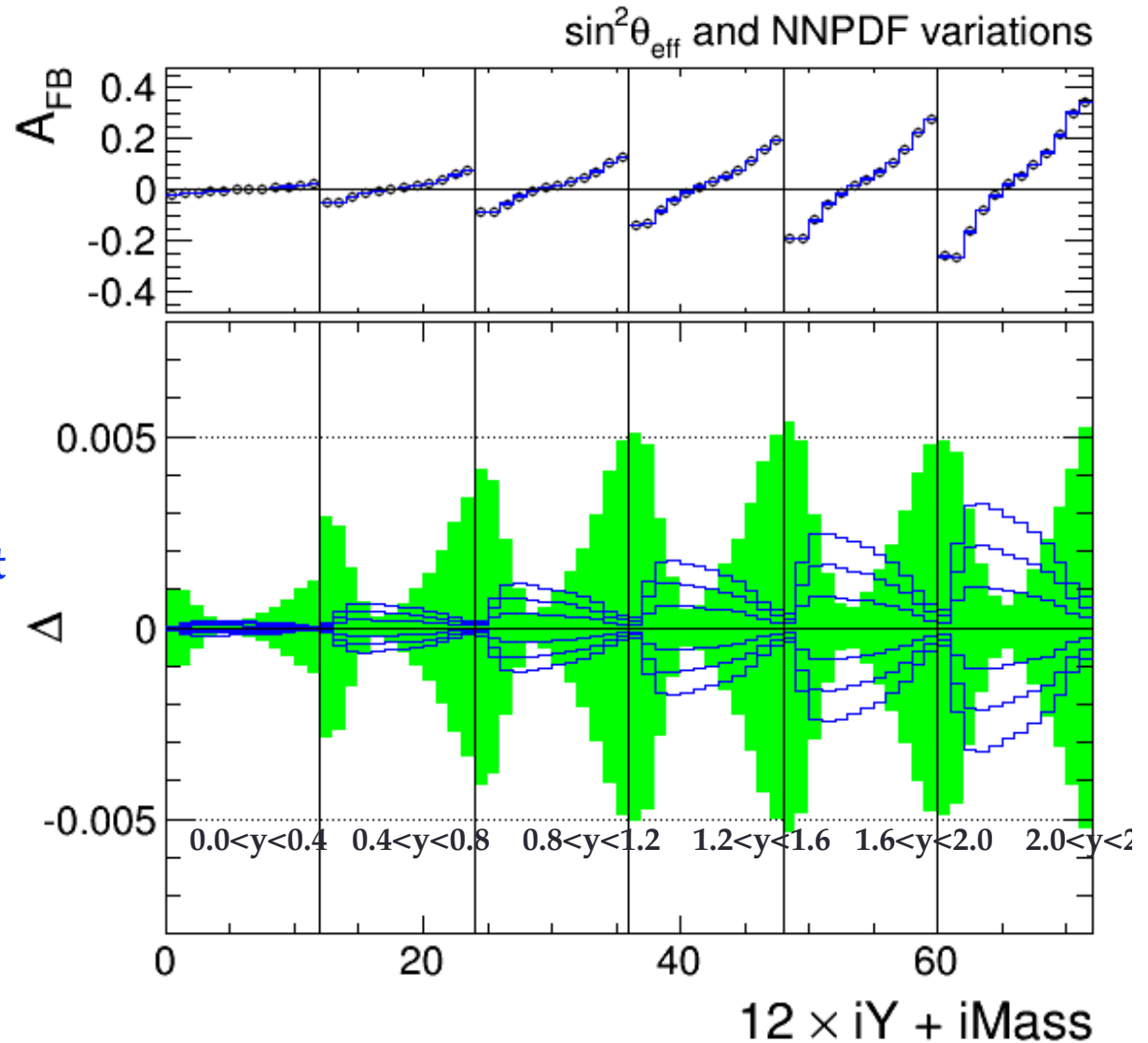
$$D_{AFB}^{LHC}(\bar{q}) \propto \frac{\bar{u}(x_1)u(x_2)}{u(x_1)\bar{u}(x_2)} \approx \frac{\bar{u}}{u}(x_1) \quad (21)$$

Both Waym and Drell-Yan Afb at the LHC can be used to reduce PDF errors on both Mw and  $\sin^2\theta_{\text{eff}}$

Green band is RMS variation from different PDFs (NNPDF3.0). Different PDFs change  $A_{\text{FB}}$  above and below the region of the Z peak.

Blue lines are variation from different  $\sin^2\theta_{\text{eff}}$  (steps of 0.00040).  $\sin^2\theta_{\text{eff}}$  changes  $A_{\text{FB}}$  in the region of the Z pole. Does not change  $A_{\text{FB}}$  at high or low mass.

Note that dilution decreases for higher rapidity.  $\sin^2\theta_{\text{eff}}$  should be the same – independent of rapidity. This provides additional information on PDFs



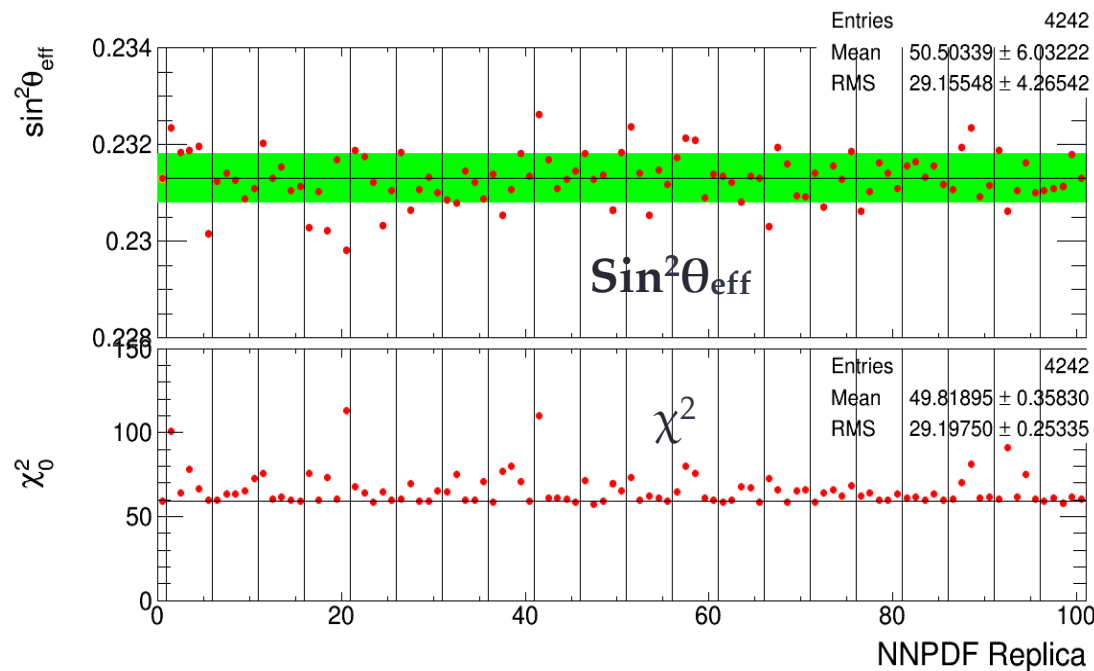
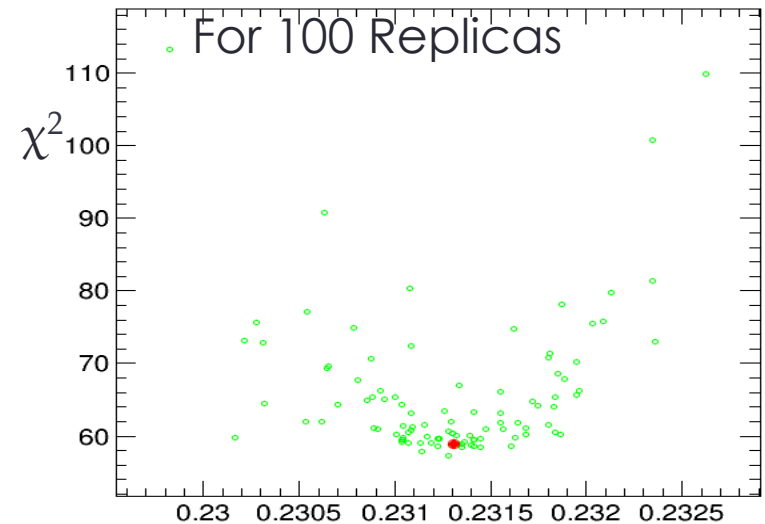
Pseudo-data 7M  $\mu\mu$  reconstructed events CMS like detector

# Constraining PDFs with Bayesian re-weighting.

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Extract  $\text{Sin}^2\theta_{\text{eff}}$  for each PDF replica. Plots shown for one pseudo experiment

- We now assign NNPDF replicas weights, based on how well they describe Afb data.  $\text{weight} \sim \exp(-\chi^2/2)$
- PDFs with bad  $\chi^2$  will have small weights



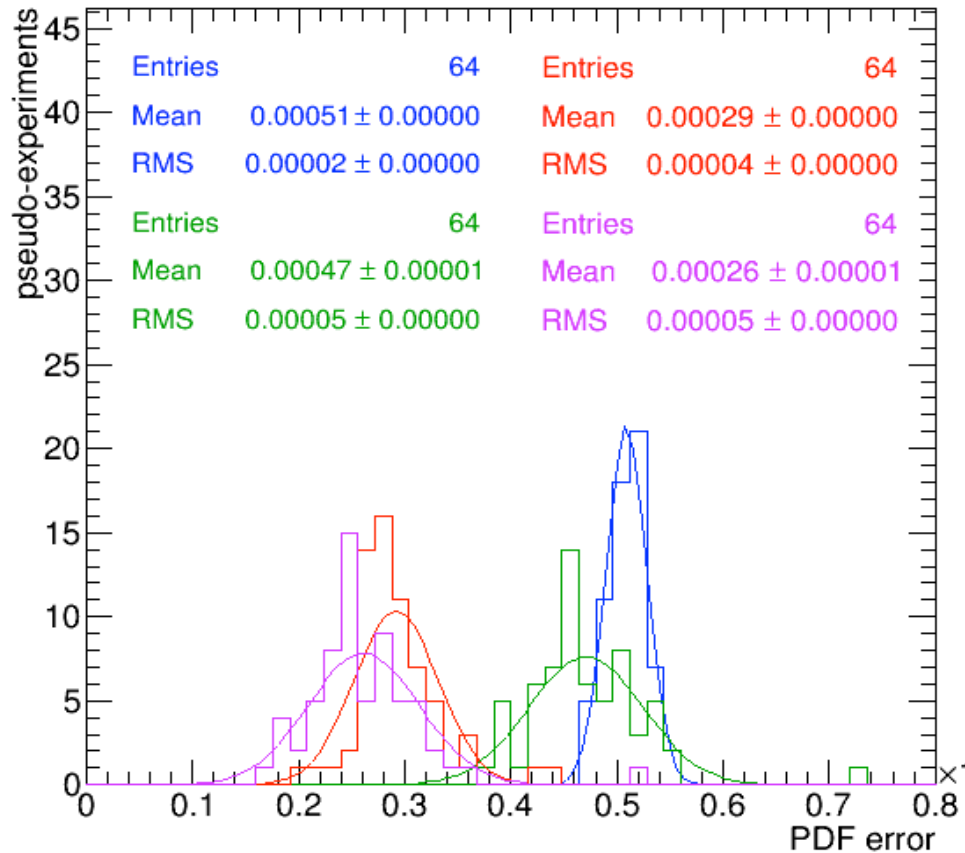
$\text{Sin}^2\theta_{\text{eff}}$

Example of one of the 64 pseudo experiments

$\chi^2$

PDF errors for **unconstrained** and constrained NNPDF3.0 For the 64 pseudo-experiments CMS-like detector (7M  $\mu\mu$  reconstructed events) Errors in  $\sin^2\theta_{\text{eff}}$ .

**Stat error** 0.00043  
**CT10 PDF error** 0.00077



**Method and PDF error**

NNPDF3.0 Unconstrained

100 replicas  $\sigma_{\text{PDF}} = \underline{\pm 0.00051}$

NNPDF3.0 8 TeV  $\mu$  W asym

constrained weight  $\sim \exp(-\chi^2/2)$

$\sigma_{\text{PDF}} = \underline{\pm 0.00047}$

NNPDF3.0 8 TeV  $\mu\mu$  AFB

constrained weight  $\sim \exp(-\chi^2/2)$

$\sigma_{\text{PDF}} = \underline{\pm 0.00029}$

NNPDF 3.0 8 TeV combined

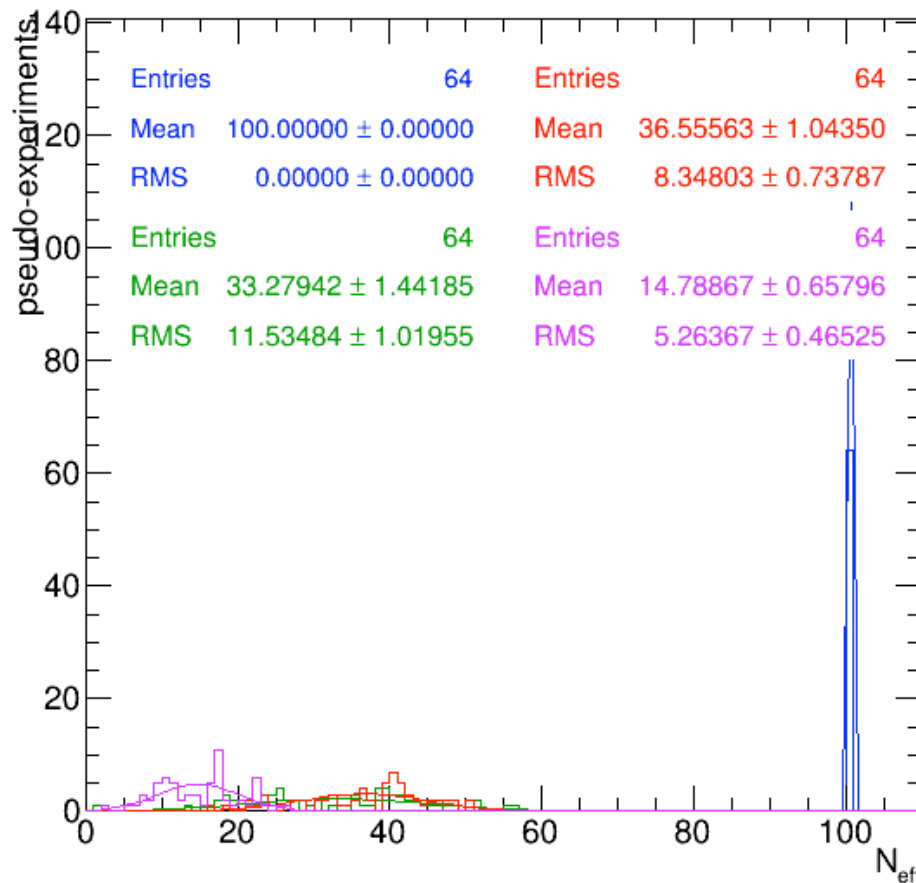
constraint from Afb AND

8 TeV W asym: weight  $\sim \exp(-\chi^2/2)$

$\sigma_{\text{PDF}} = \underline{\pm 0.00026}$

At 8 TeV Afb and 8 TeV W asym data reduce PDF errors to  $\pm \underline{0.00026}$ . **Factor of 3 smaller than CT10. Will become even smaller with 13-14 TeV data)**

What is the number of effective replicas for **unconstrained** and constrained NNPDF3.0 for various levels of constraints.



Method # of effective replicas  
No constraints

$\sigma_{\text{PDF}} = \pm 0.00051$  100 replicas

W asym data weight  $\sim \exp(-\chi^2/2)$

$\sigma_{\text{PDF}} = \pm 0.00047$  33 replicas

AFB data weight  $\sim \exp(-\chi^2/2)$

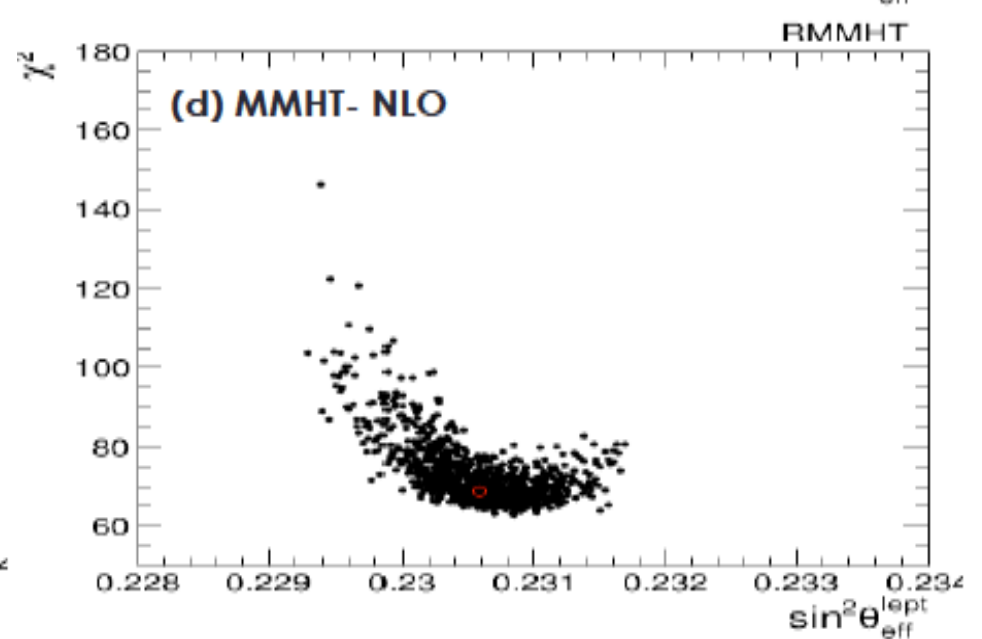
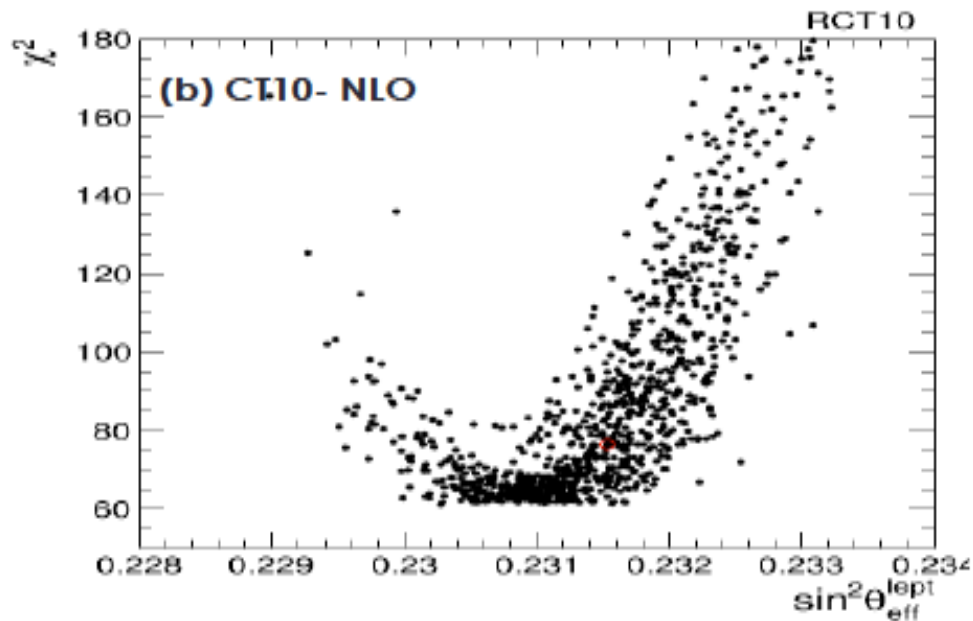
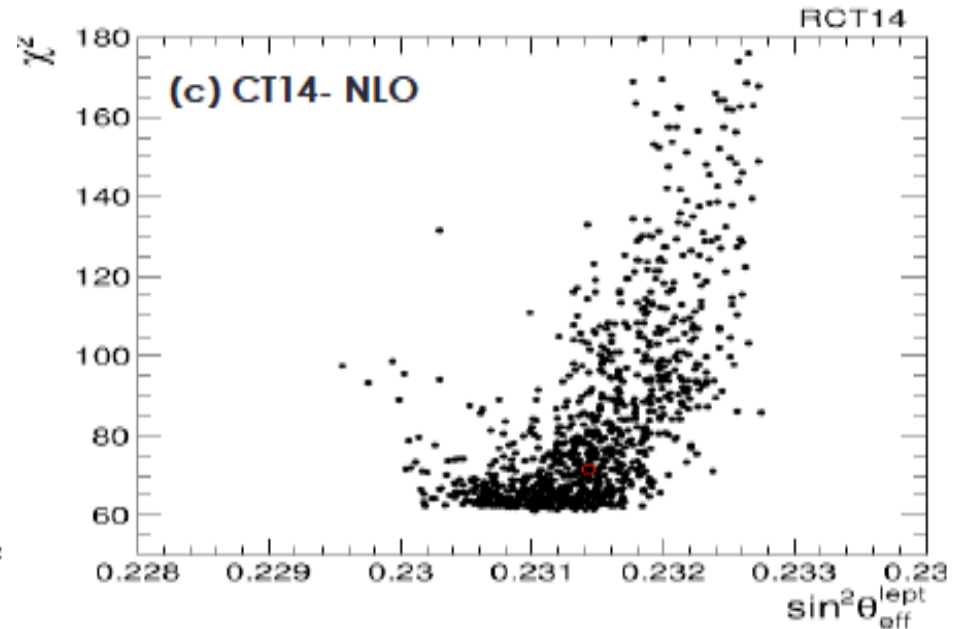
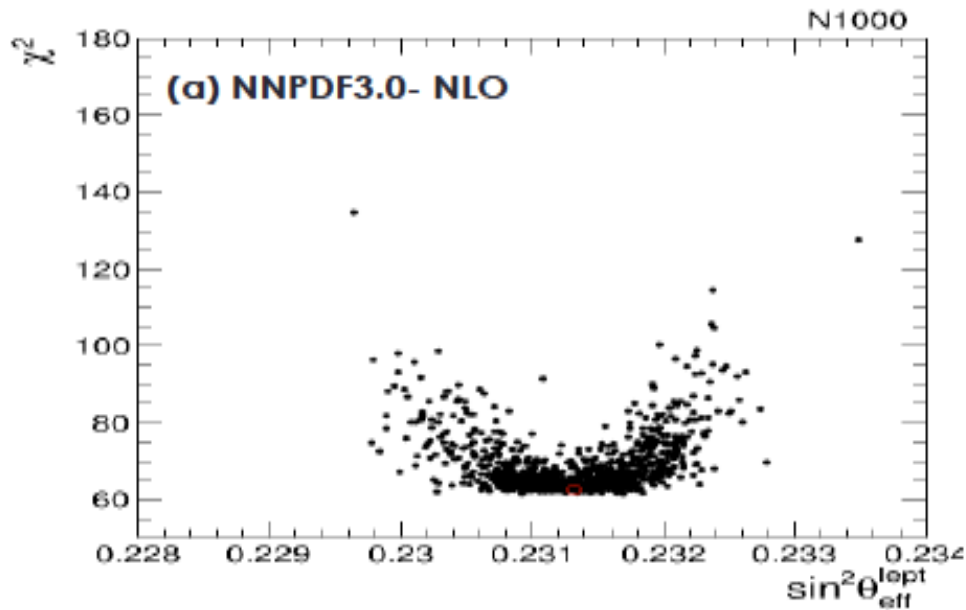
$\sigma_{\text{PDF}} = \pm 0.00029$  37 replicas

AFB + W asym weight  $\sim \exp(-\chi^2/2)$

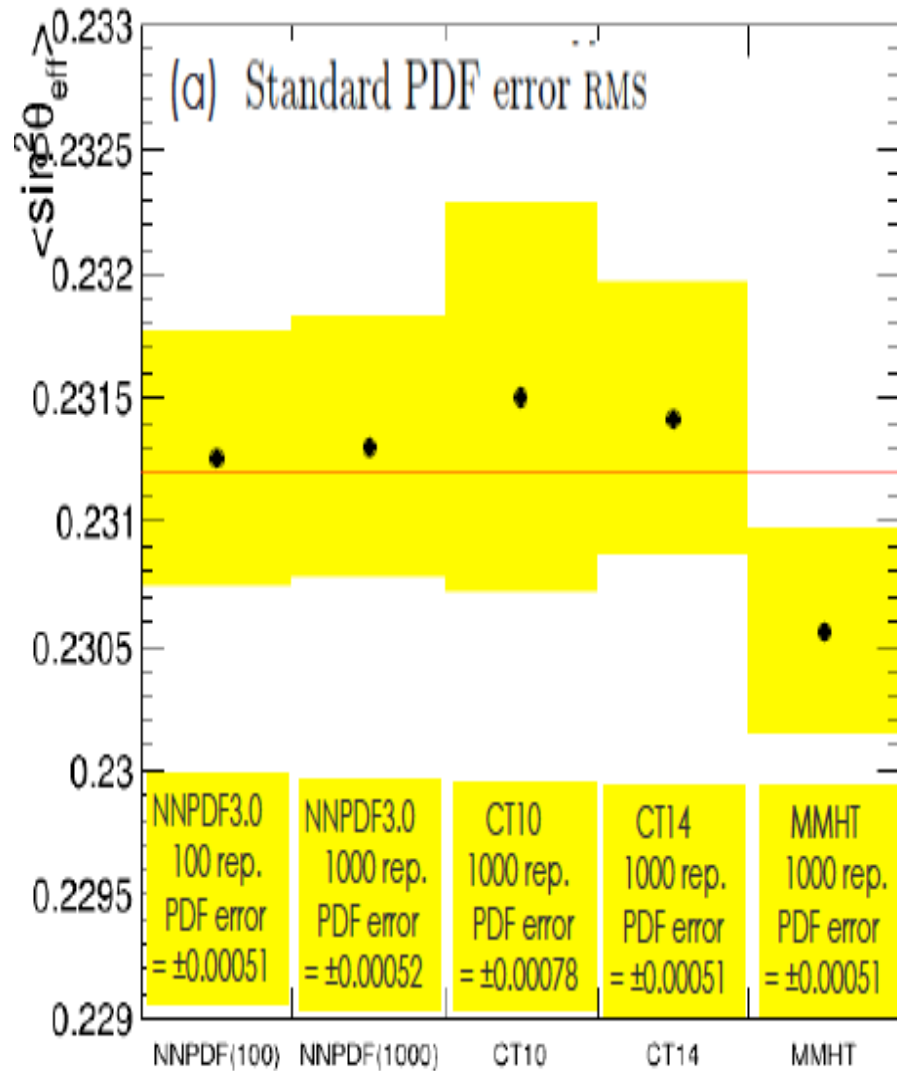
$\sigma_{\text{PDF}} = \pm 0.00026$  15 replicas

*We are repeating the study with 1000 replicas*

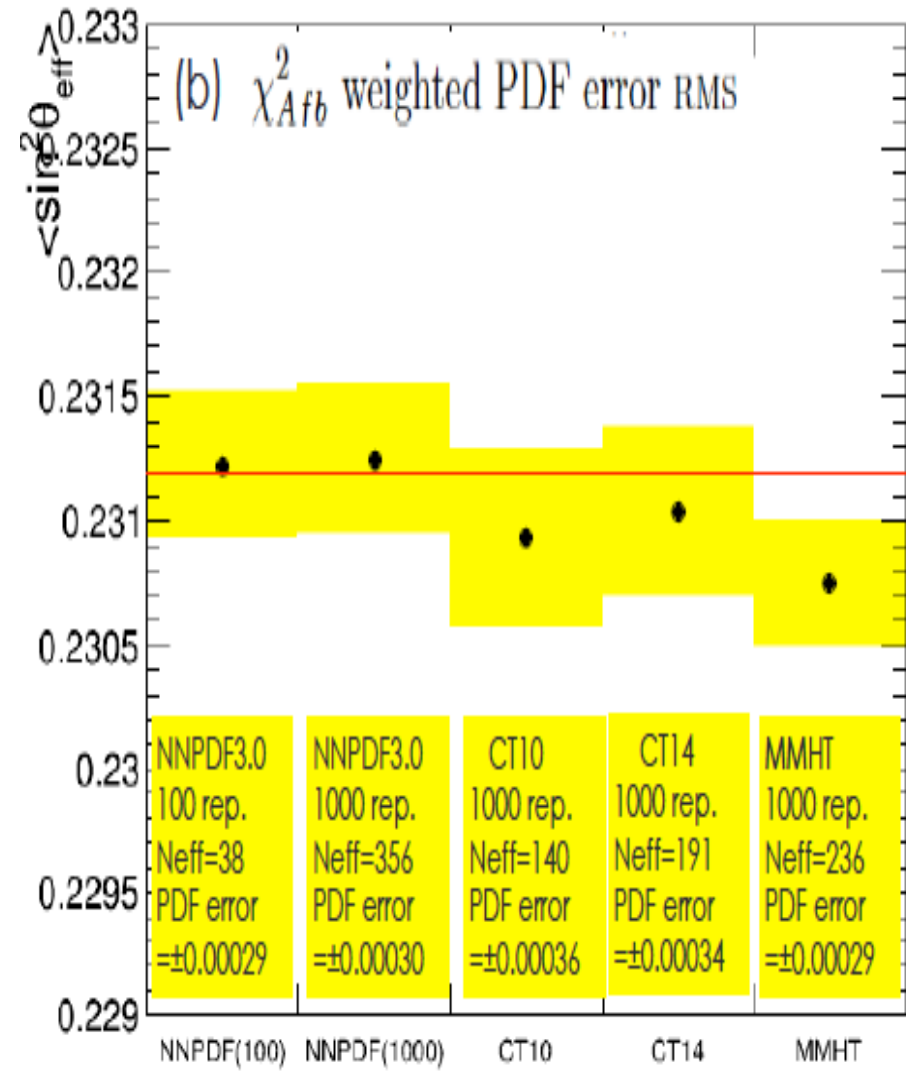
# Afb pseudo 8 TeV CMS data (NNPDF3.0 central) 32



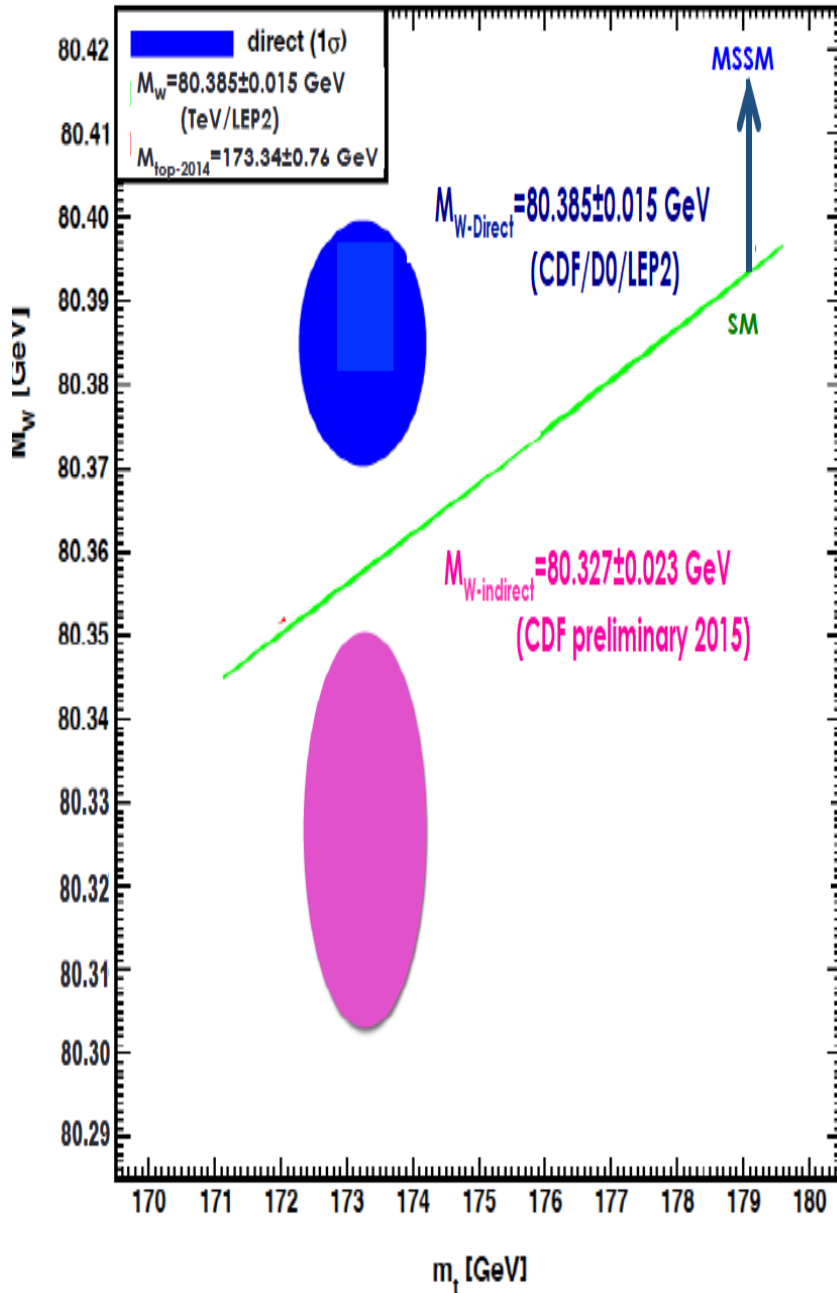




Model	Bias
NNPDF3.0	+0.00005
NNPDF3.0	+0.00008
CT10	+0.00031
CT14	+0.00022
MMHT	-0.00063



Model	Bias
NNPDF3.0	+0.00003
NNPDF3.0	+0.00006
CT10	-0.00026
CT14	-0.00016
MMHT	-0.00044



CMS like detector	2016 sample	2017-18 sample
Energy	8 TeV	13-14 TeV
Number of reconstructed events	8.2M $\mu^+\mu^-$ 6.8M $e^+e^-$	120M $\mu^+\mu^-$ -
$\Delta \sin^2 \theta_W$		
Statistical error	$\pm 0.00034$	$\pm 0.00011$
Weighted PDF error	$\pm 0.00022$	$\pm 0.00014$
(Stat+PDF) error	$\pm 0.00040$	$\pm 0.00018$
$\Delta M_W^{indirect}$	MeV	MeV
Statistical error	$\pm 17$	$\pm 5$
weighted PDF error	$\pm 11$	$\pm 7$
(Stat+PDF) error	$\pm 20$	$\pm 9$

CMS  
End of 2015  
Analysis in  
Progress  
Match CDF

CMS  
end of 2017  
A factor of  
2 reduction  
in error