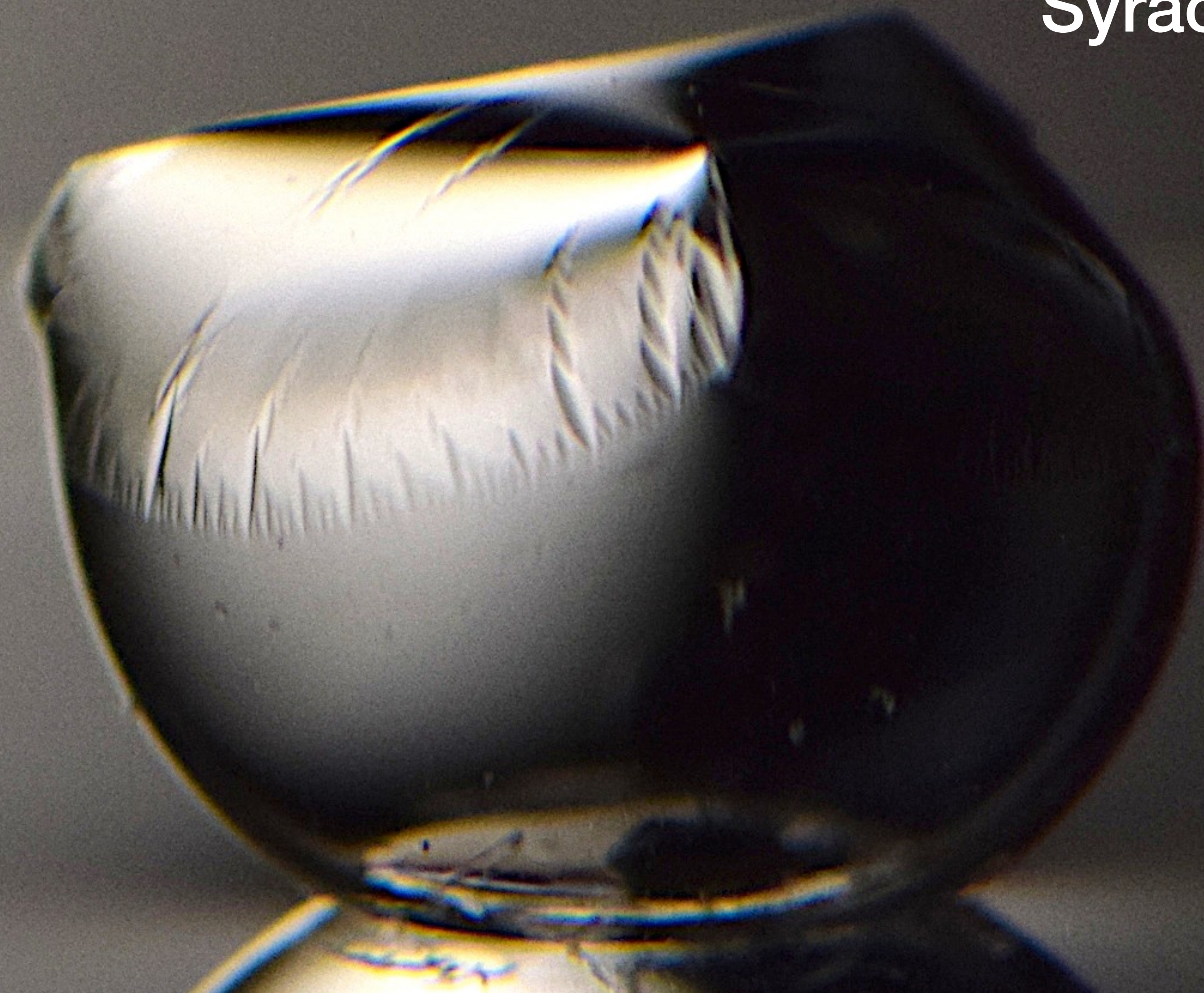


Sheets shaping liquids and liquids shaping sheets

Joseph D Paulsen
Syracuse University



It takes a village...

Paulsen Group

Yousra Timounay

Monica Ripp

Graham Leggat

Mengfei He

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UMass Amherst

Deepak Kumar

Chris Santangelo*

Benny Davidovitch

Thomas Russell

Narayanan Menon

K. Bugra Toga

Jooyoung Chang

Oxford

Dominic Vella

simulations



Teng Zhang



Monica Ripp



Yousra Timounay



Deepak Kumar



Tom Russell



Narayanan Menon



Vincent Démery



Benny Davidovitch



Chris Santangelo



Dominic Vella

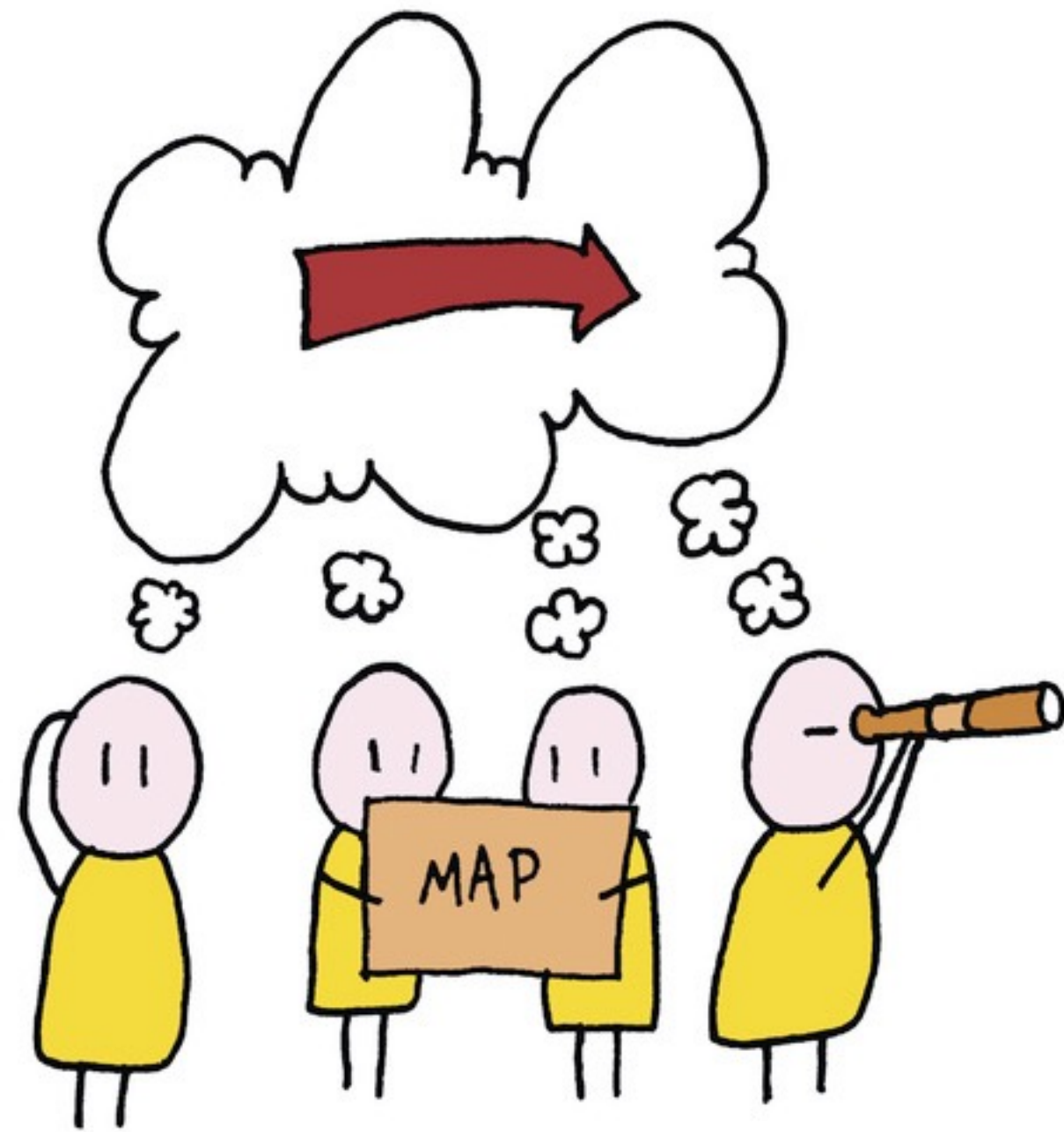
experiments

theory

Funding: W. M. Keck Foundation, NSF DMR-CAREER-1654102

NSF REU DMR-1460784, NSF REU DMR-1757749

*(now Syracuse)

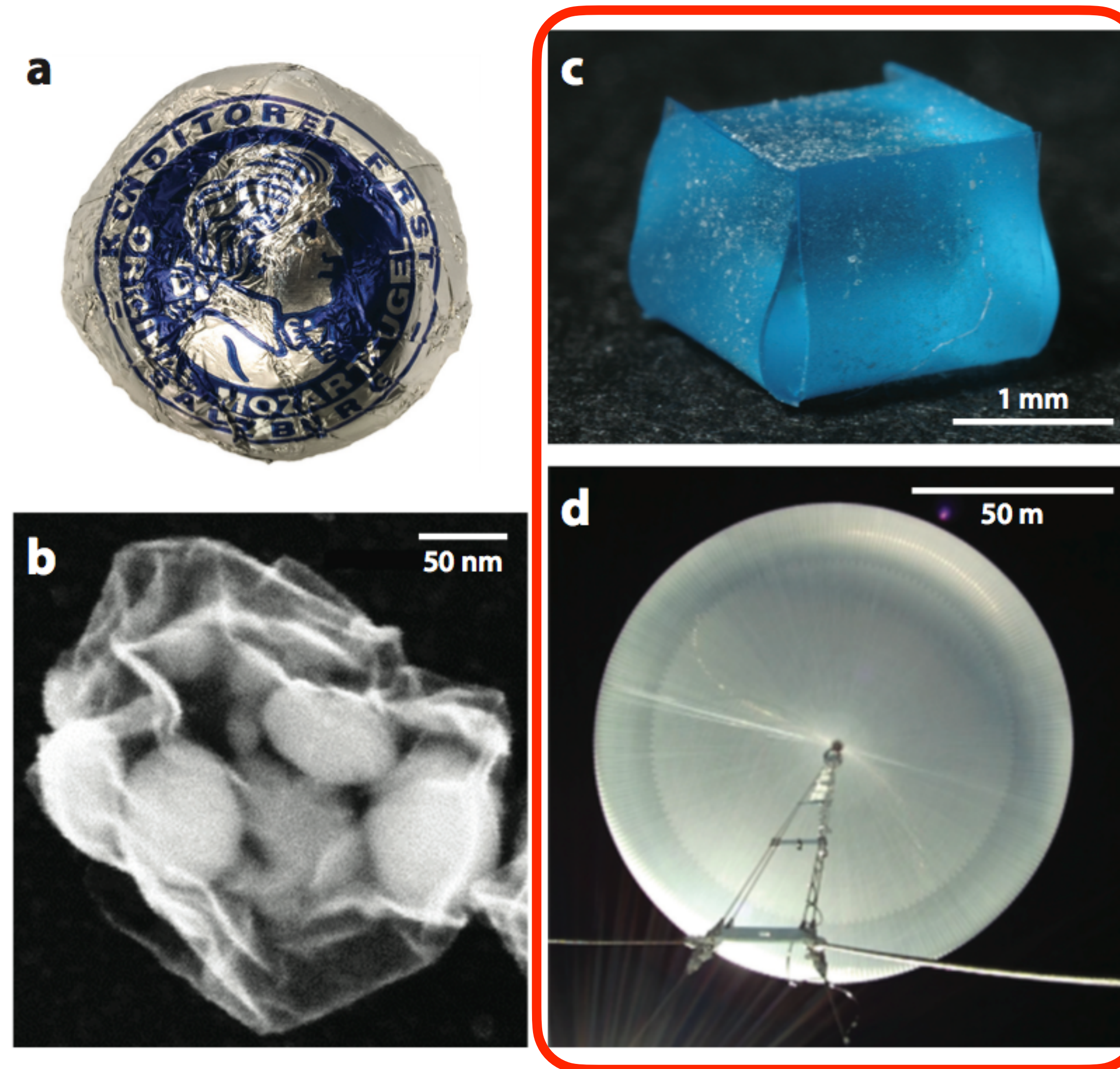


0. Wrapping a droplet with a thin solid
1. Indenting a floating film
2. Beyond the “gross shape”
3. Towards solid surfactants

we think we know what's going on

bleeding edge

We wrap objects in thin sheets to conceal, protect, or enhance them



Wrapper & contents
both easily deformed

- (a) Wikipedia, cf. Demaine et al. (2009)
- (b) Chen et al., Nano Lett. 12 (2012)
- (c) Py, Bico, Roman, et al. (2007)
- (d) Image courtesy of NASA

Figure from: “Wrapping liquids, solids, and gases in thin sheets”
JDP, Annu. Rev. Condens. Matter Phys. Vol. 10 (2019)

“Capillary Origami”

Py, Bico, Roman, Baroud, et. al., 2007



“Thick” films ($t \sim 60 \mu\text{m}$): long wavelength bending

What happens for a film that is **1000 times thinner?**

Bending energies ($\sim t^3$) are then **$\times 10^9$ smaller!**

$t = 80 \text{ nm}$ circular polystyrene sheet



1 mm

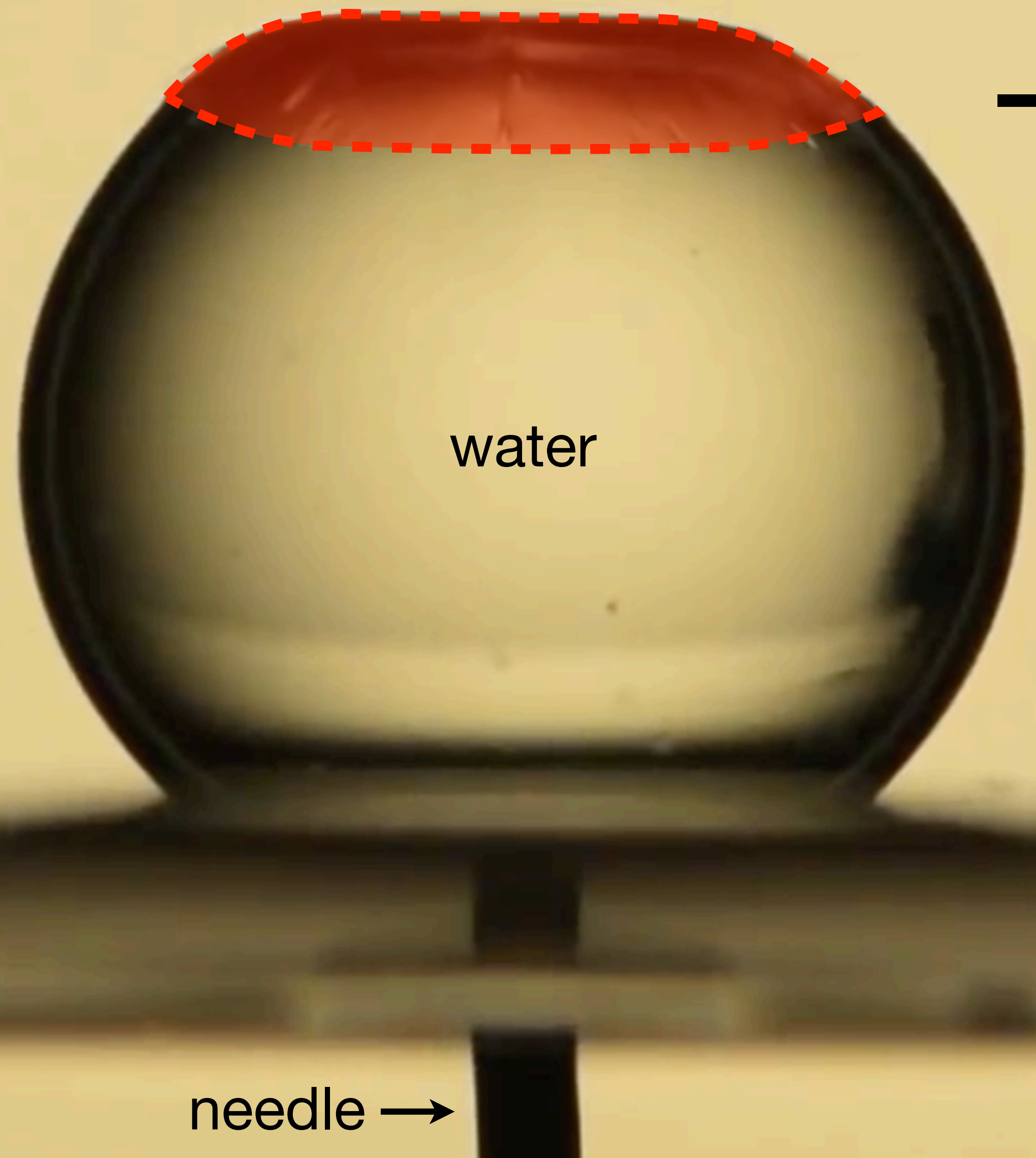


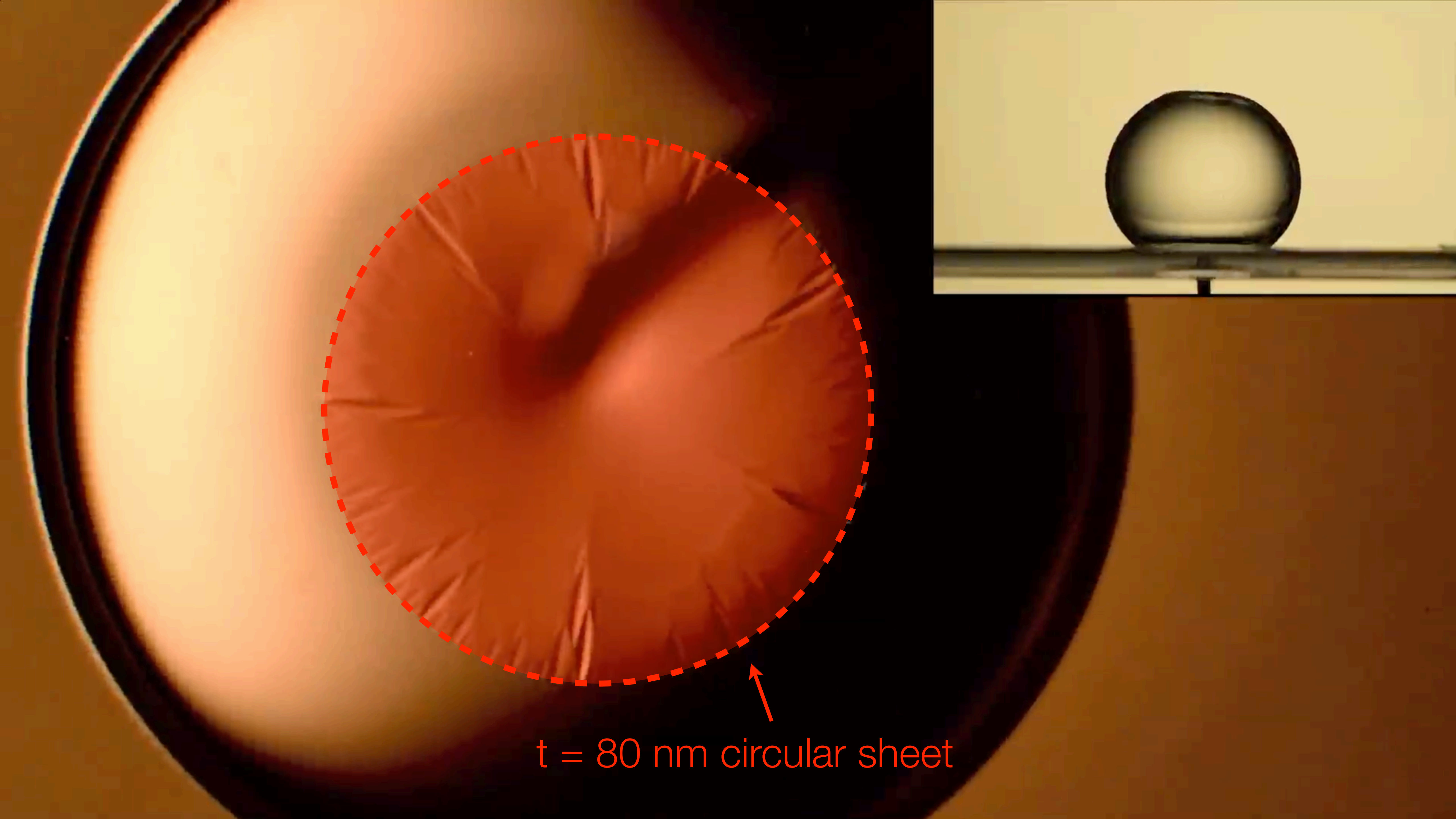
water

silicone oil

fluorinated oil

needle →



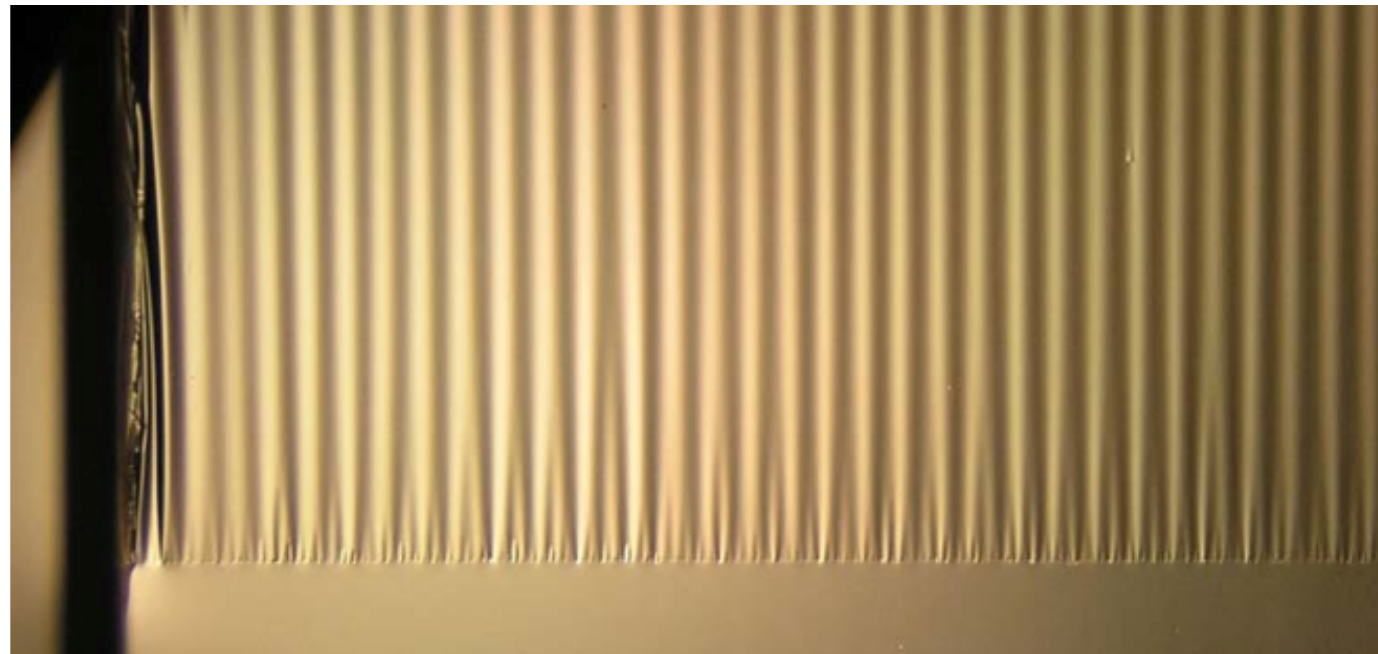


$t = 80 \text{ nm}$ circular sheet

How do thin sheets respond to confinement?

...they wrinkle

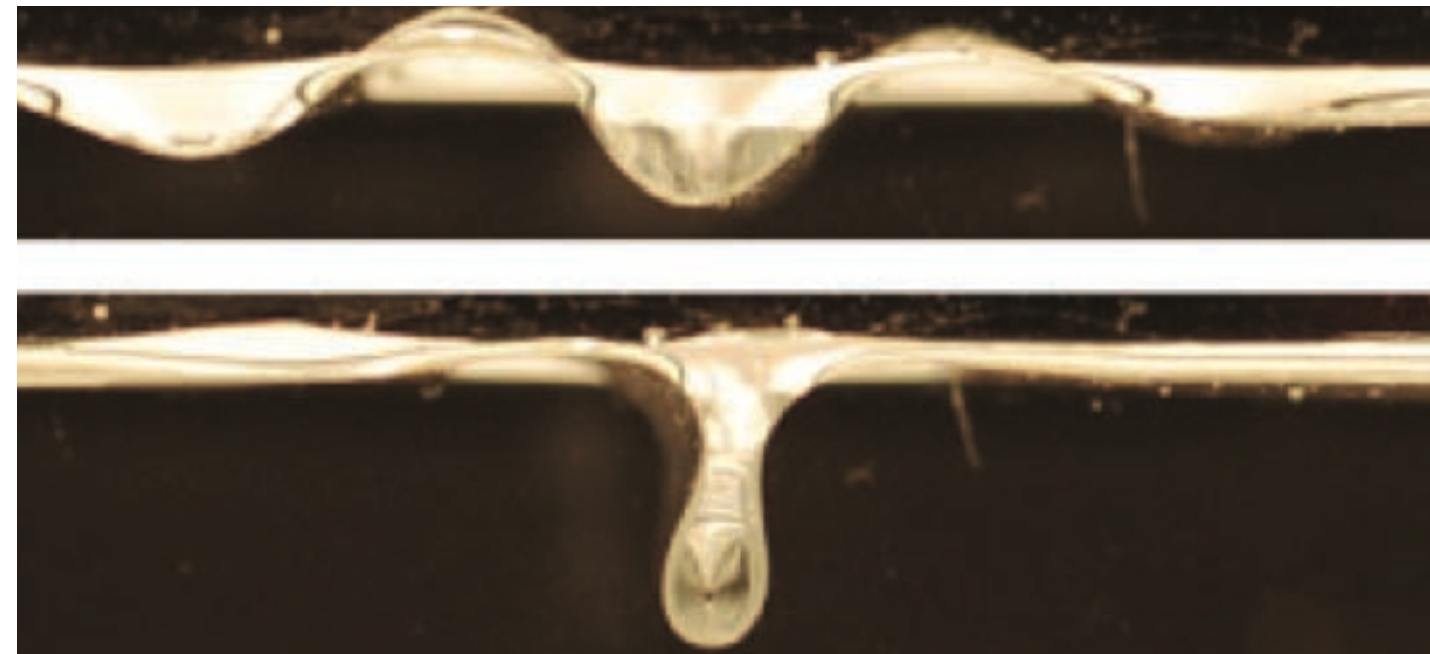
Smooth deformation



Huang et al., *PRL* 2010

...they fold

Localized deformation



Pocivavsek et al., *Science* 2008

...they crumple

Localized deformation and stress

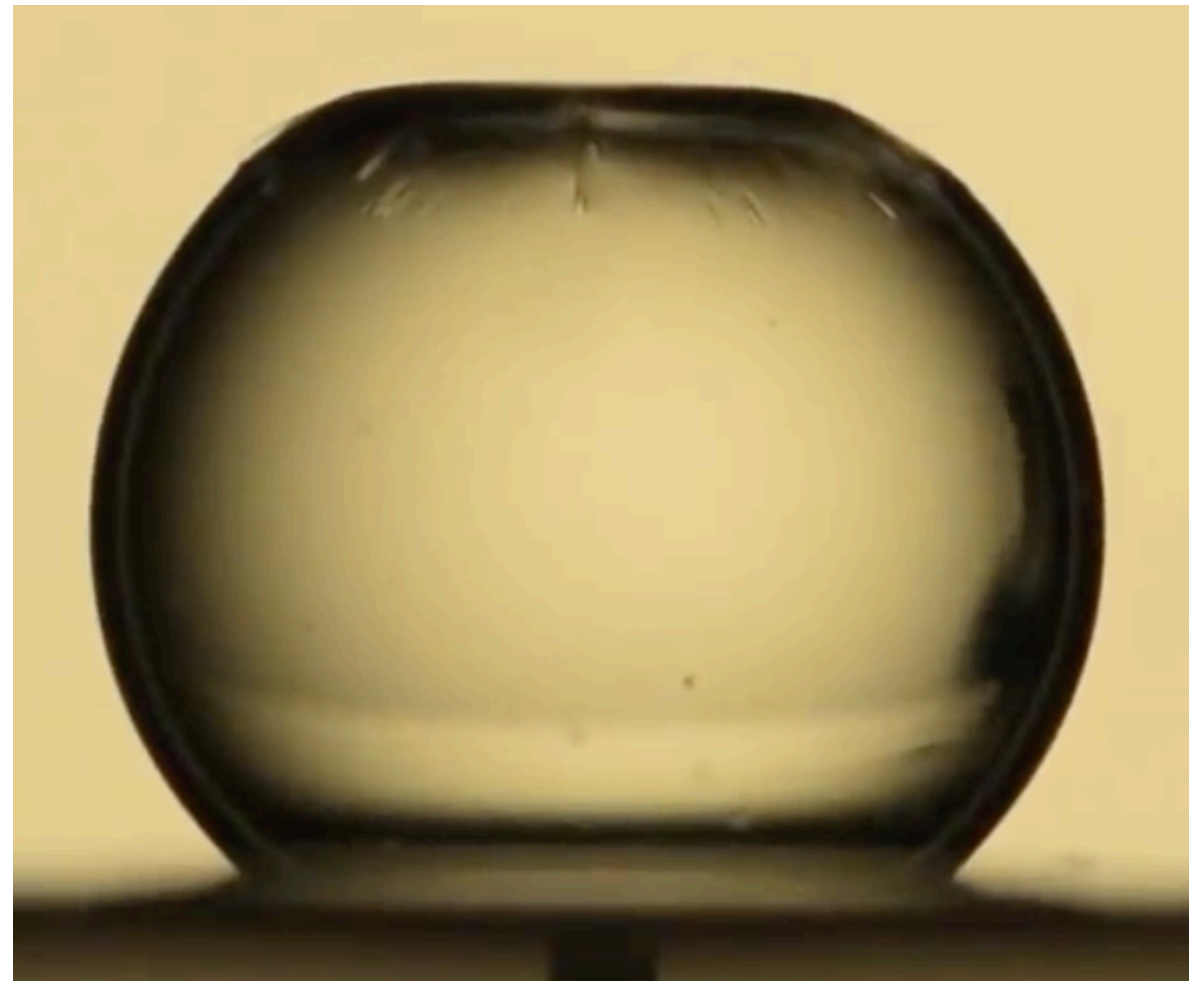


Subtle physics, nonlinear sheet equations

Wrapping: all of these in highly curved geometry!

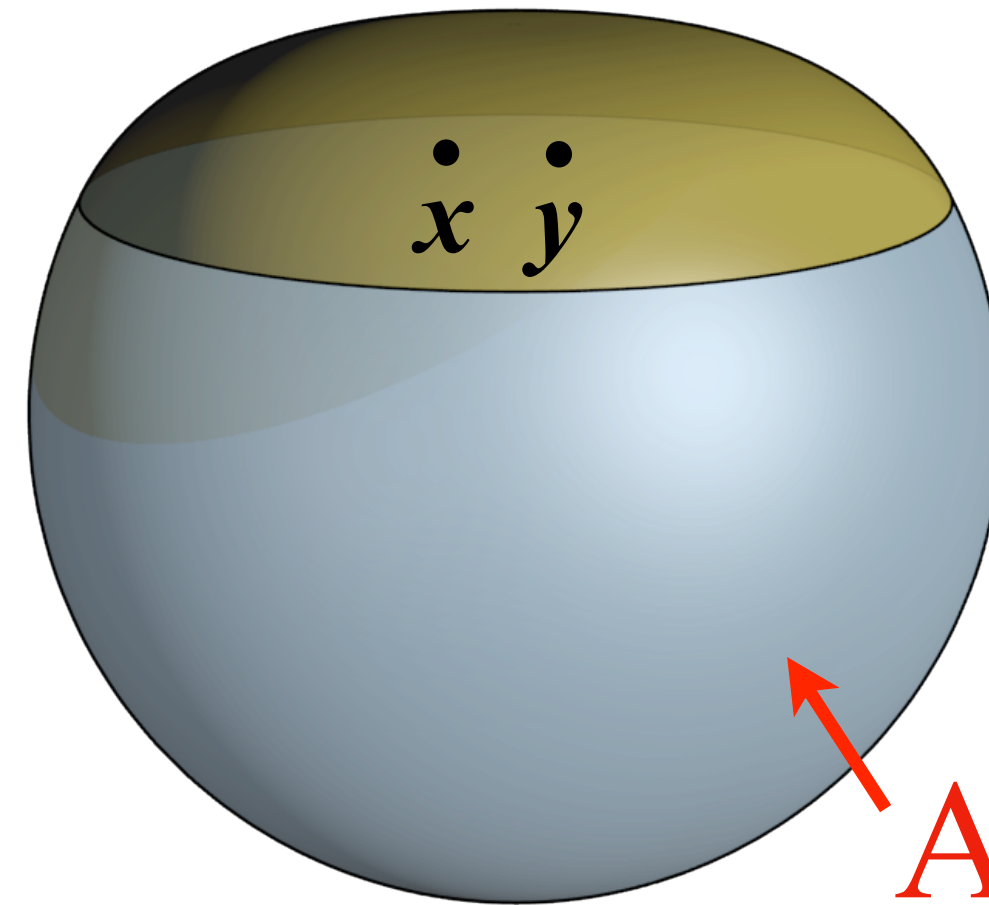
The Plan:

actual shape



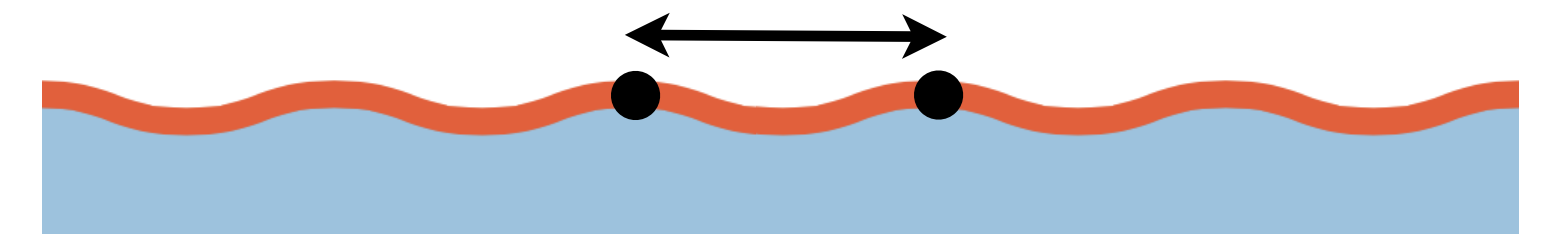
=

“gross shape”



+

“small-scale features”
(wrinkles, crumples, folds)



...which allow in-plane
compression at negligible cost

Gross shape **can compress**
but not stretch:

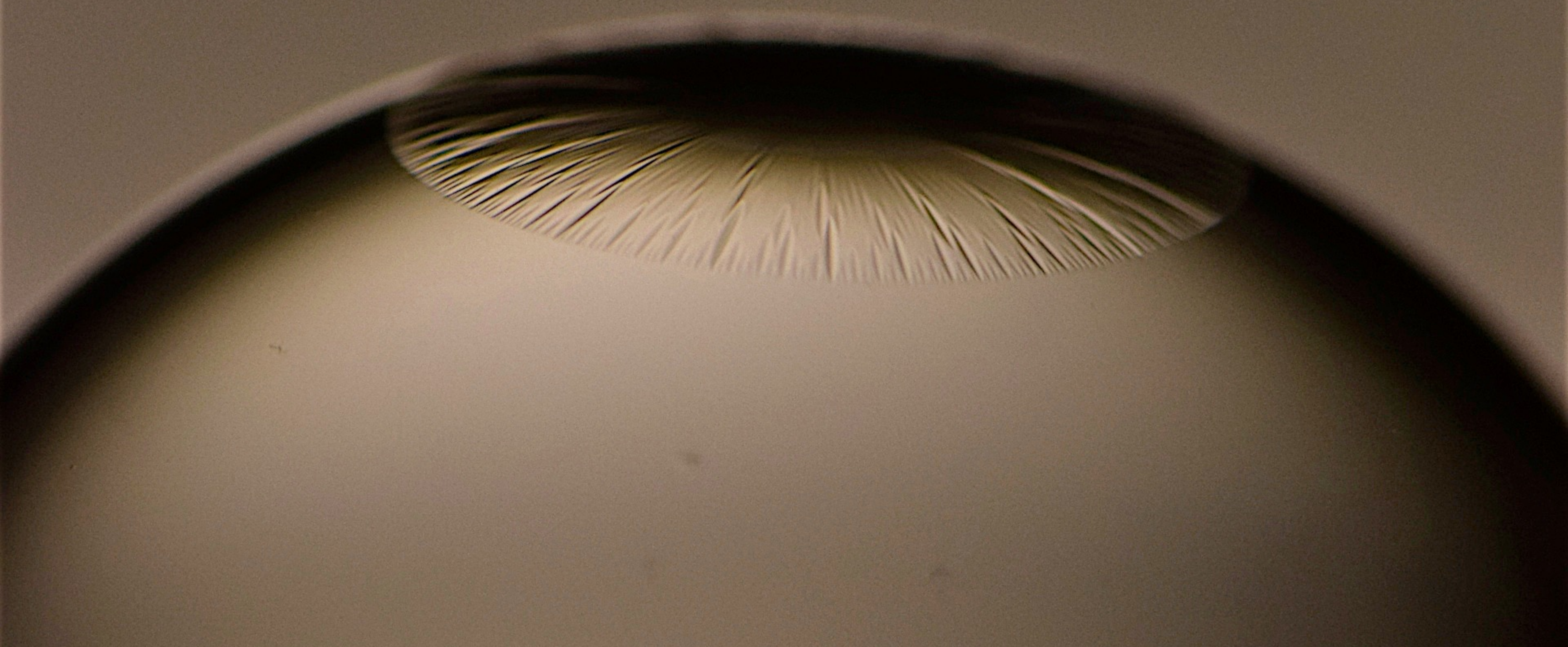
$$|f(x) - f(y)| \leq |x - y|$$

where f : planar sheet \rightarrow gross shape

Geometric model: $U = \gamma A_{\text{free}}$ (System seeks **gross shape** that minimizes A_{free})

JDP, Démery, et al., Nat. Mater. 2015

Test case 1: Large droplet (gross shape is axisymmetric)

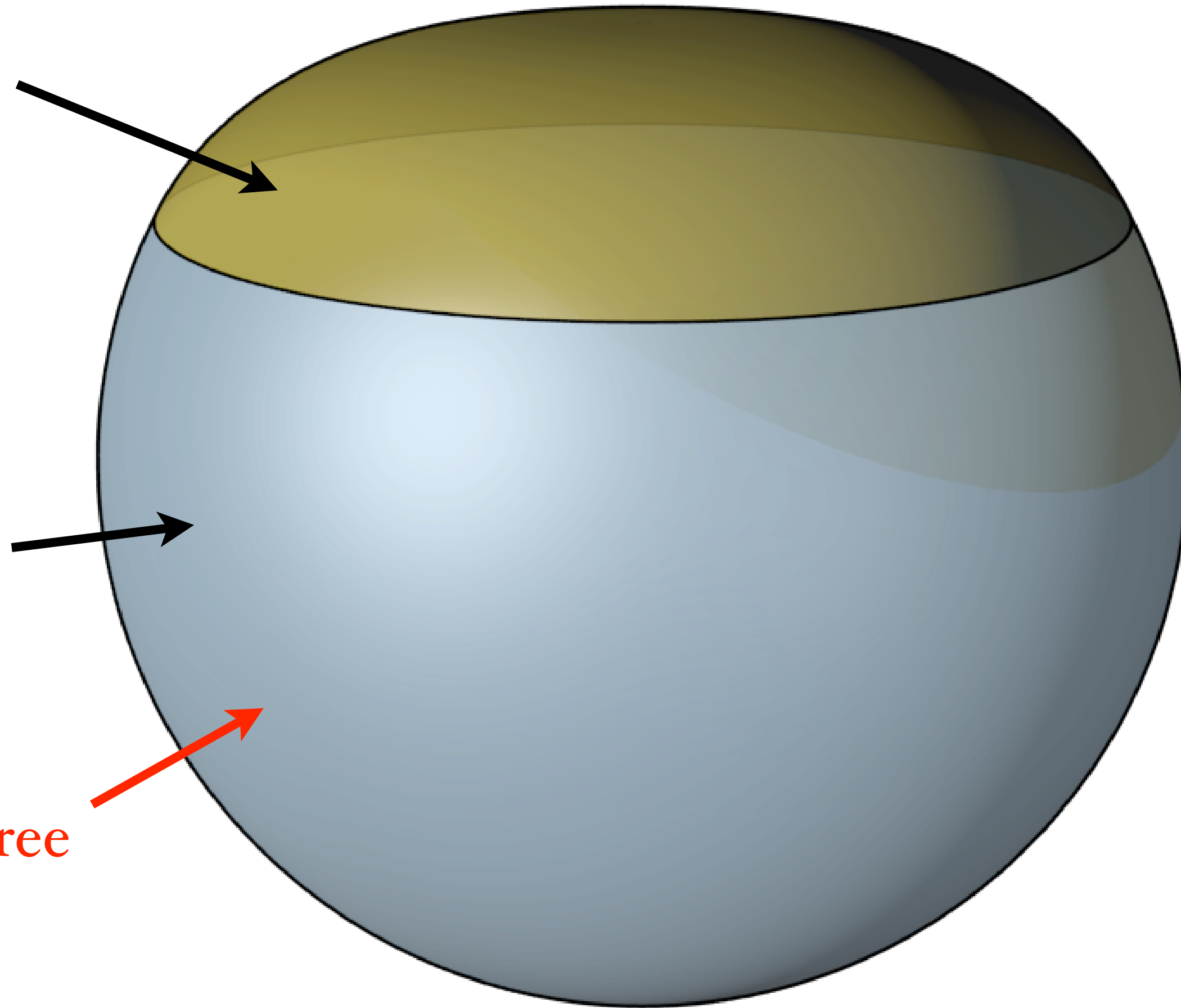


Global minimum of $U = \gamma A_{\text{free}}$ among all axisymmetric configurations

Sheet: elliptic integral
(enabled by small-scale wrinkling, not shown)

Drop: sphere

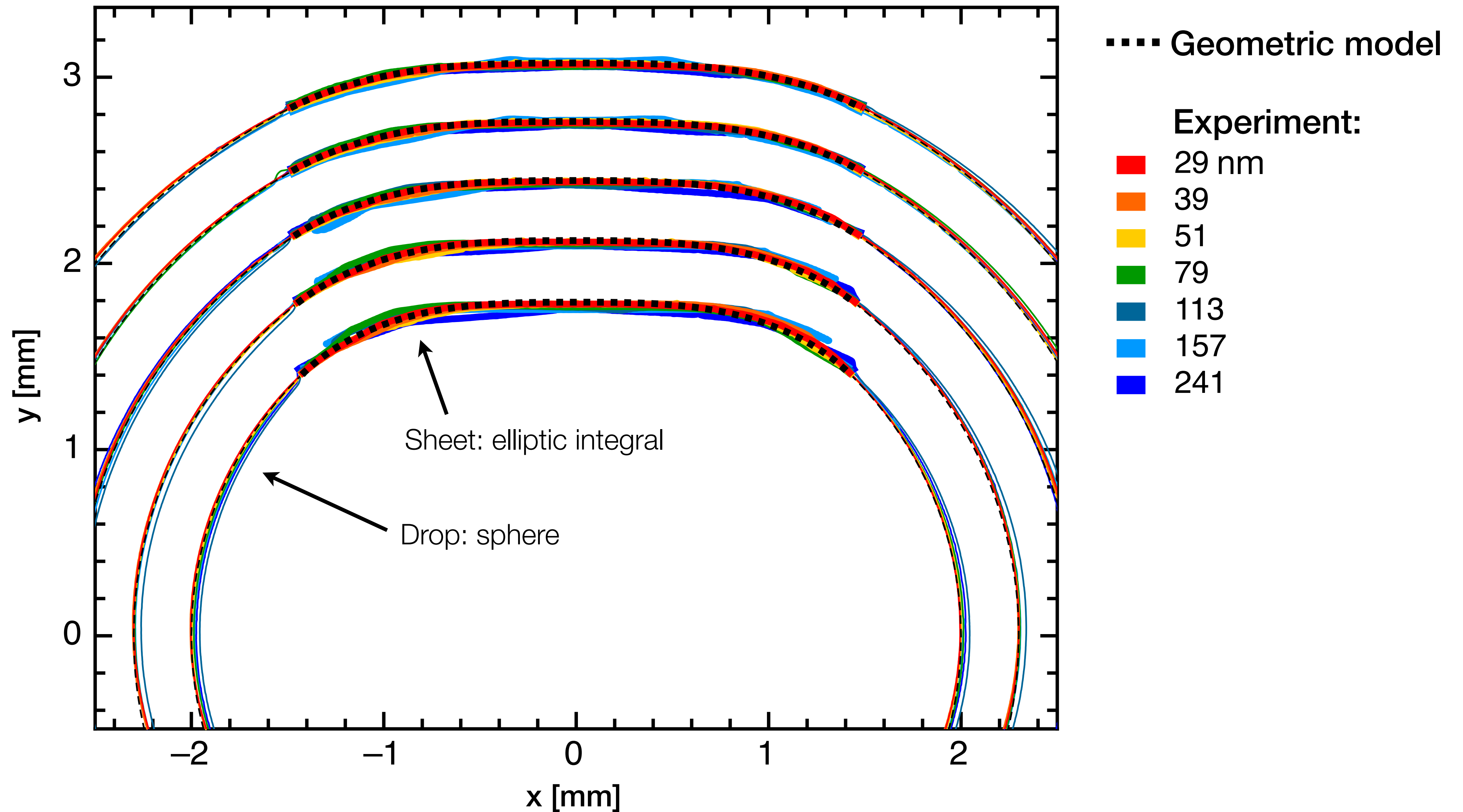
A_{free}

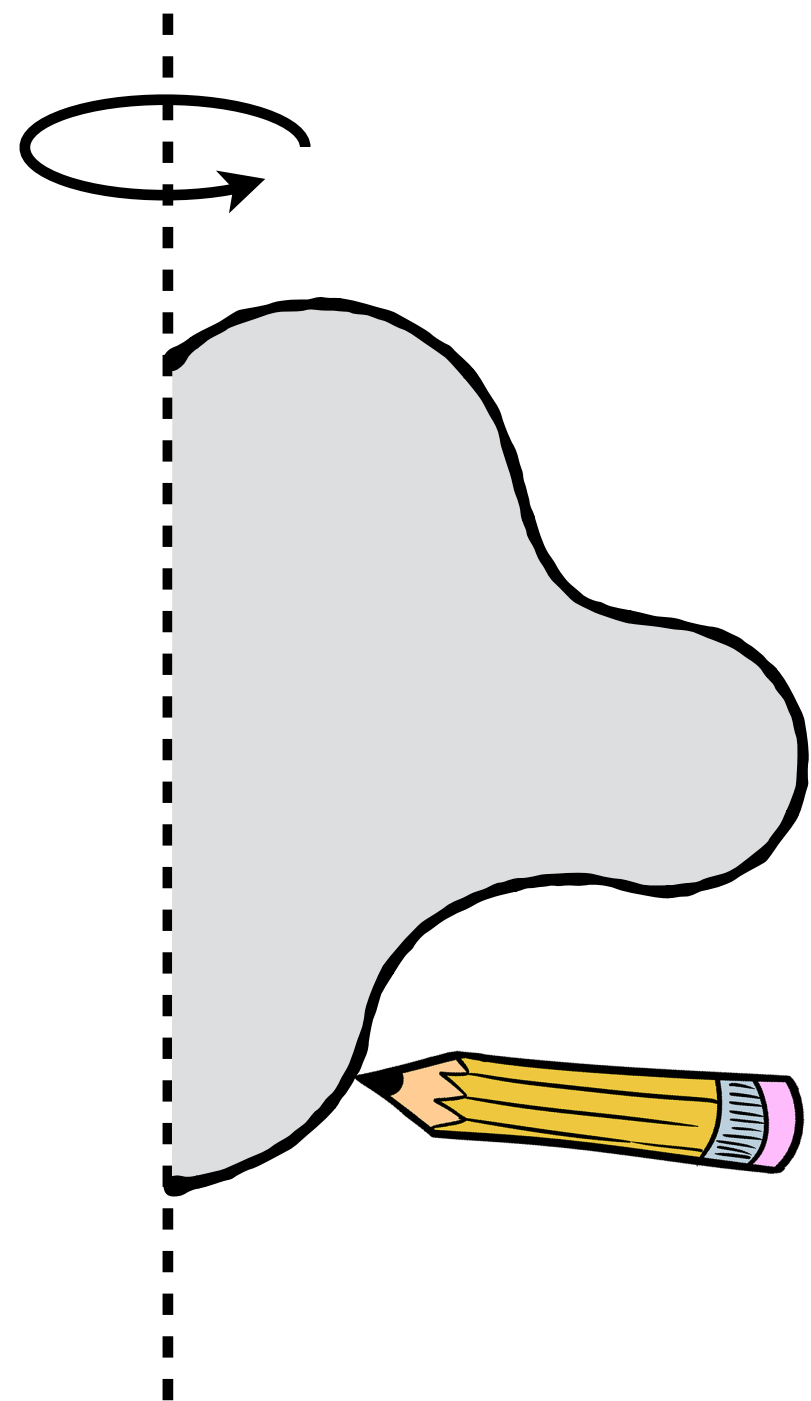


Single parameter: W/R

W : sheet radius
 $(4\pi/3) R^3$: drop volume

Comparing geometric model ($U = \gamma A_{\text{free}}$) with experiment





1. Draw a line starting and ending on z-axis
2. Consider its surface of revolution

Optimization:

Maximize volume for given surface area

Maximize volume for given arclength

Solution:

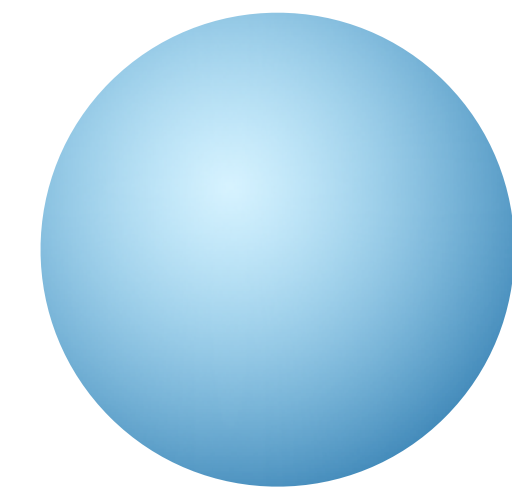
Sphere

Elliptic integral:

$$f(x) = \int_x^a \frac{t^2}{\sqrt{a^4 - t^4}} dt$$

where $a = \frac{4r\sqrt{2\pi}}{\Gamma(1/4)^2}$

Example:

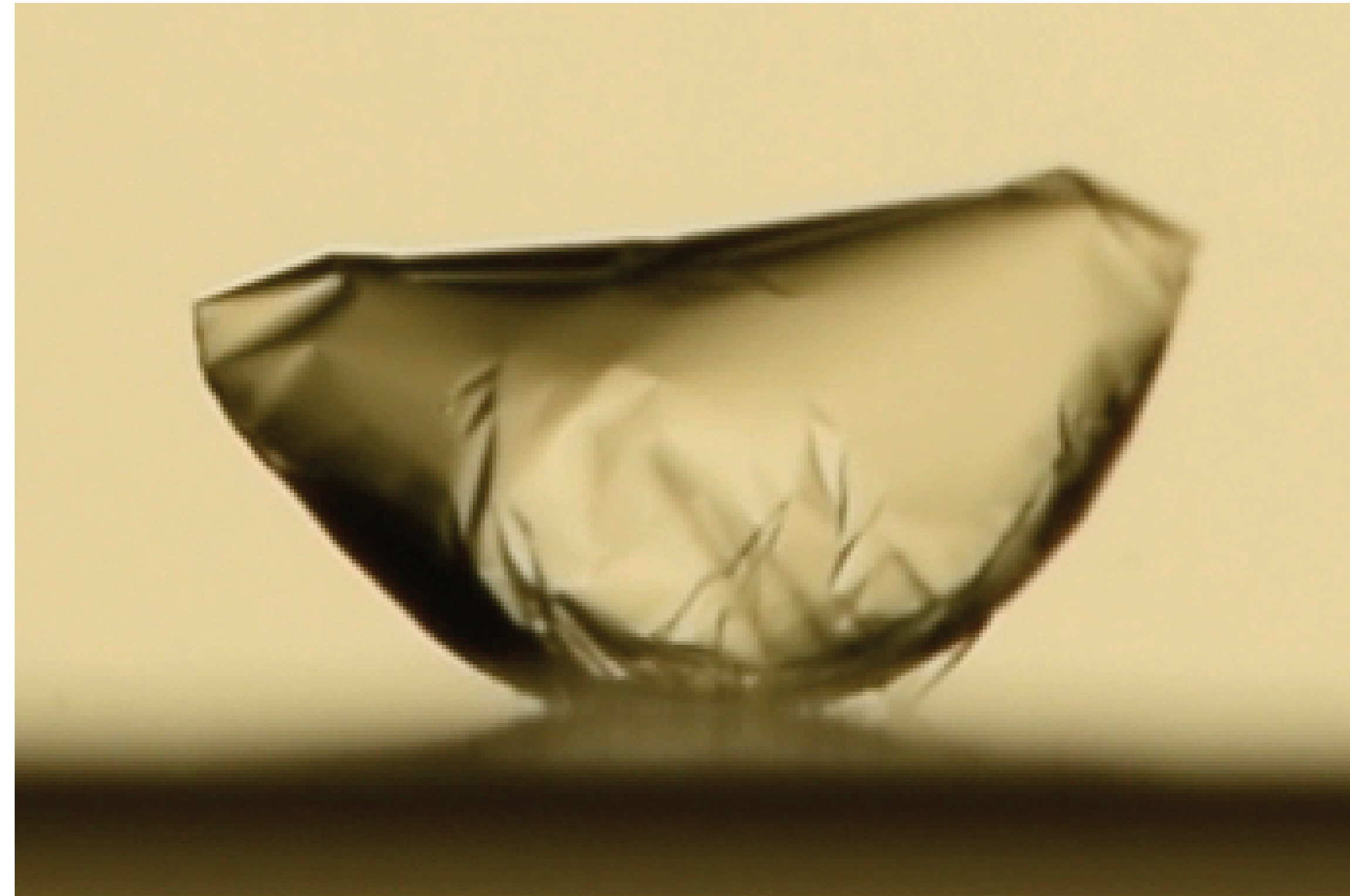


liquid droplet, soap bubble

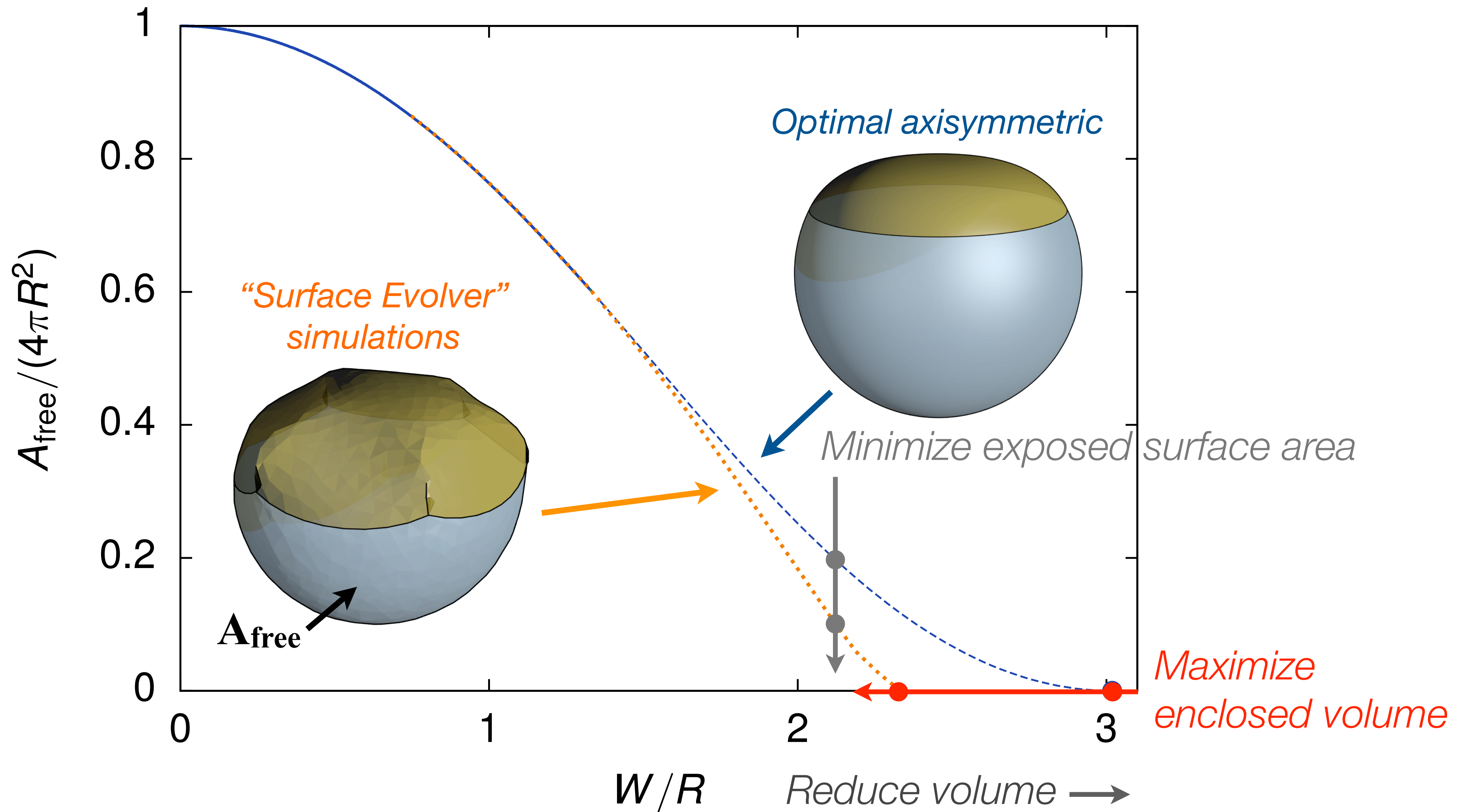


G.I. Taylor 1919: Parachute
Paulsen (not me) 1994: Mylar balloon

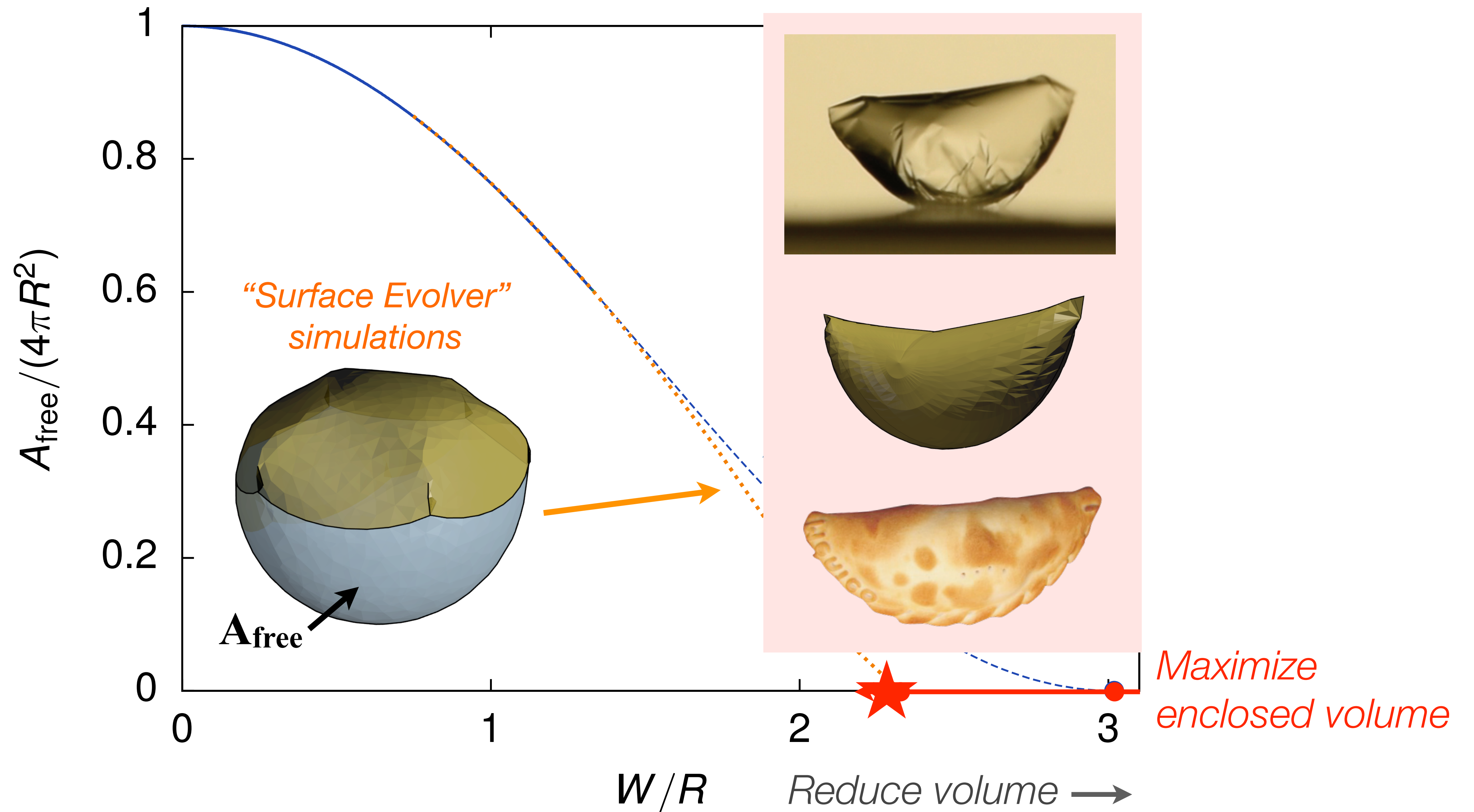
What about small drops?
(no longer axisymmetric)

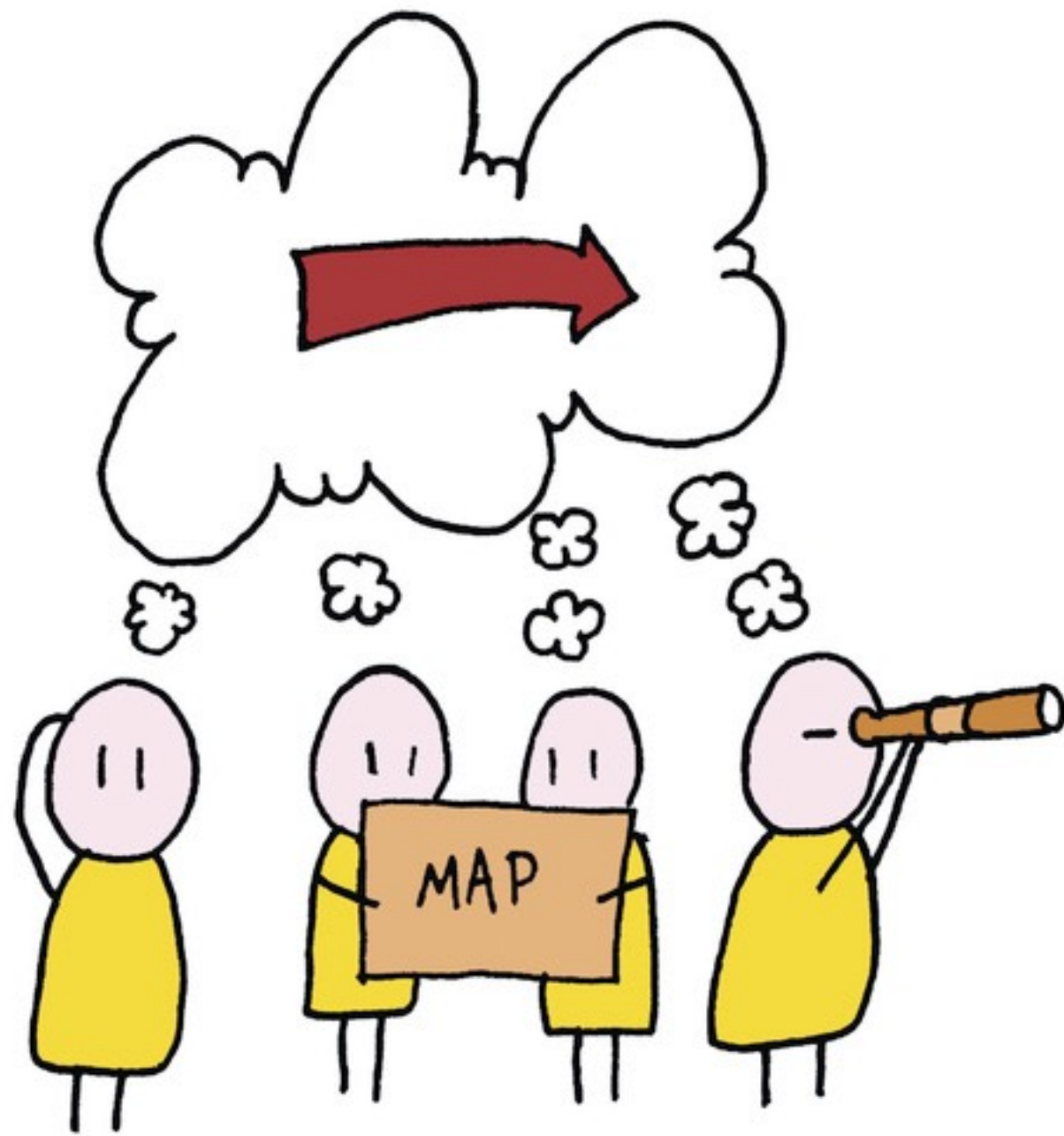


Breaking axial symmetry improves coverage



Breaking axial symmetry improves coverage





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Indentation:

Holmes & Crosby, PRL 2010

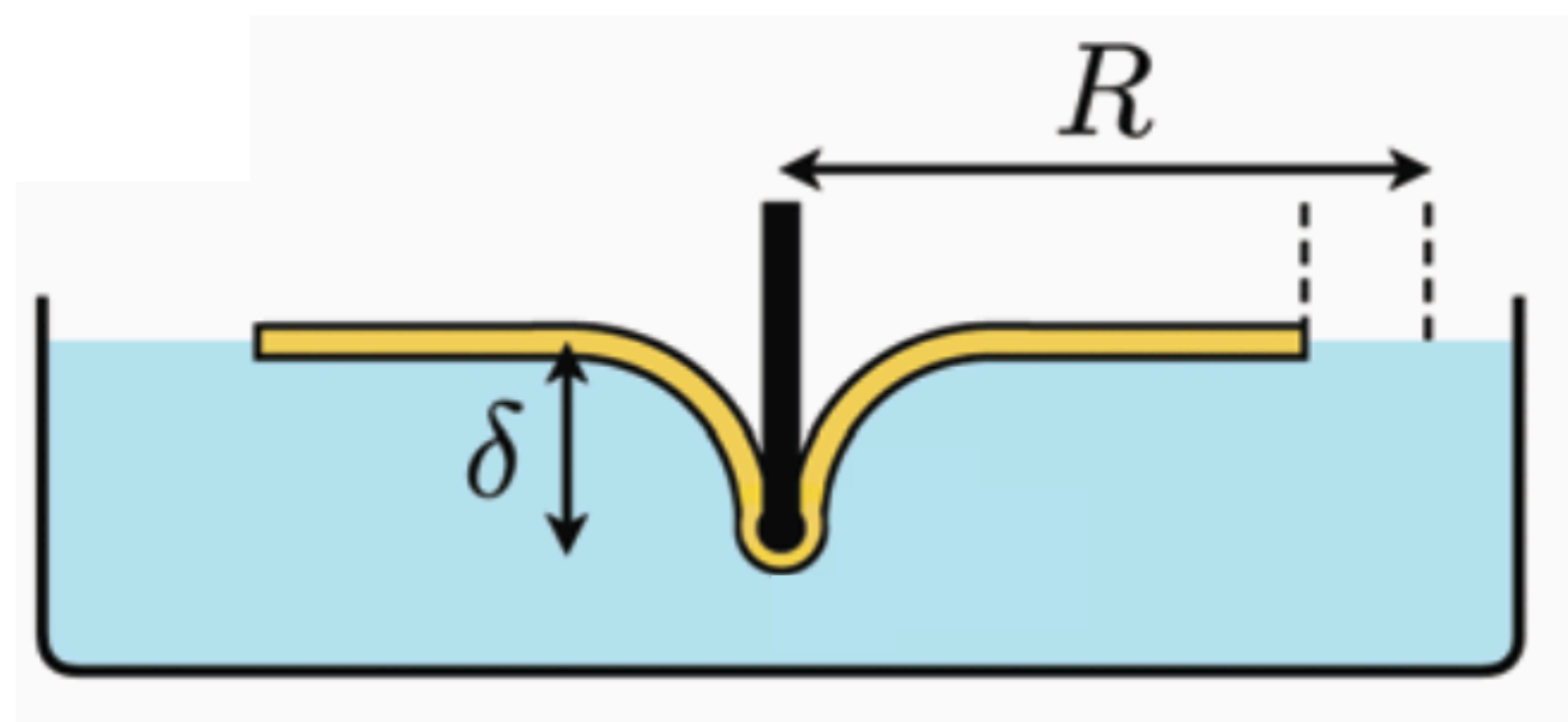
Vella, Davidovitch, et al., PRL 2015, PRE 2018

Paulsen, Hohlfeld, et al., PNAS 2016

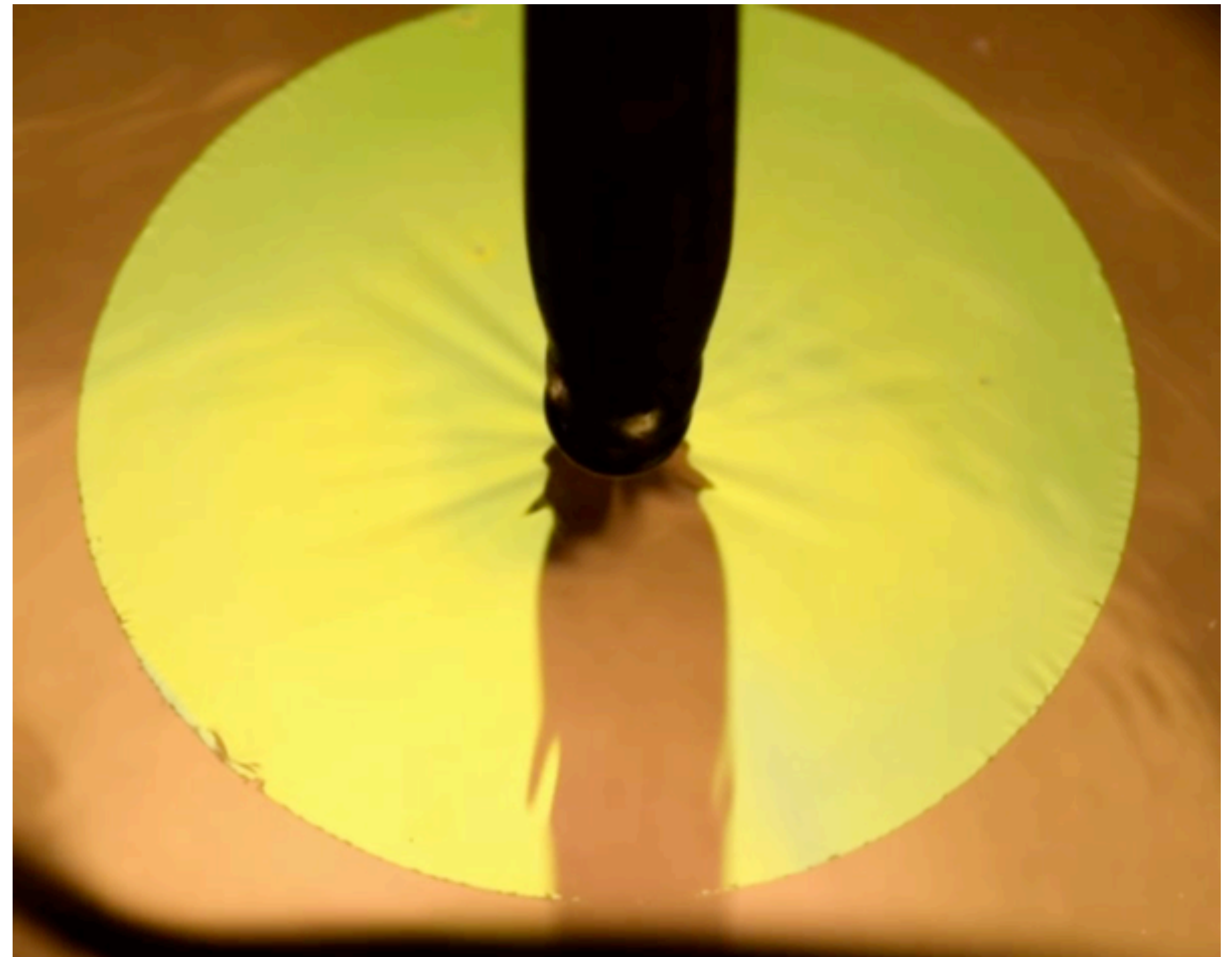
Ripp, Démery, Zhang, JDP, Soft Matter 2020



Monica Ripp



side view



top view

Geometric model for poking:

- Energy functional, now with *gravity*

$$U = U_{\text{gravity}} + \gamma(\Delta A_{\text{free}}) \quad \leftarrow \text{fluid energies only}$$

$$= \pi \int_0^\infty \left[\rho g r \zeta(r)^2 + 2\gamma R \left(\sqrt{1 + \zeta'(r)^2} - 1 \right) \right] dr$$

- Optimal $\zeta(r)$ satisfies **Euler-Lagrange equation:**

$$\zeta''(r) = \frac{r\zeta(r)}{\ell_{\text{curv}}^3} [1 + \zeta'(r)^2]^{3/2} \quad \text{where} \quad \ell_{\text{curv}} = \ell_c^{2/3} R^{1/3}$$

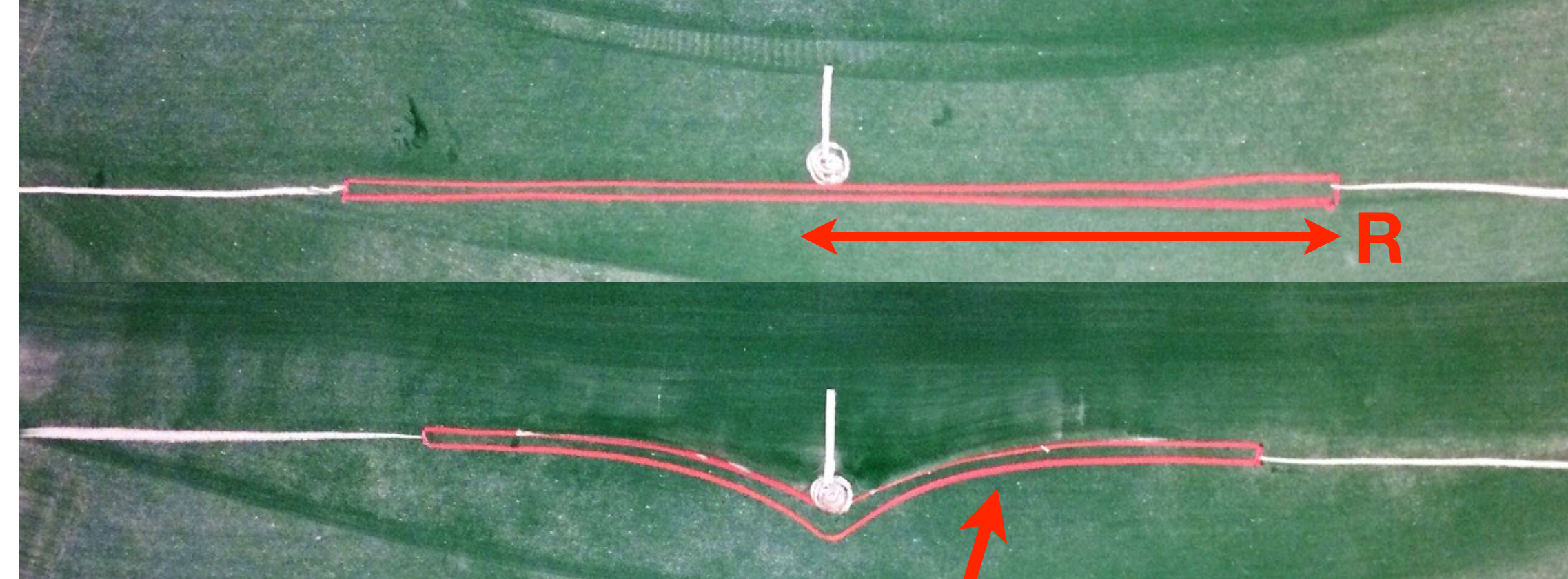
- **Small slopes:** Reduces to $\zeta''(r) = r\zeta(r)/\ell_{\text{curv}}^3$ (with same Airy function solution from Dominic's talk)

- **BONUS:** Geometric model gives access to **large slopes behavior:**

$$\text{Profile: } \zeta(r) \sim -[3 \log(1/r)/2]^{2/3} \text{ as } r \rightarrow 0$$

$$\text{Energy: } U \sim 2\pi\gamma R\delta$$

$$\text{Force: } F \sim 2\pi\gamma R$$

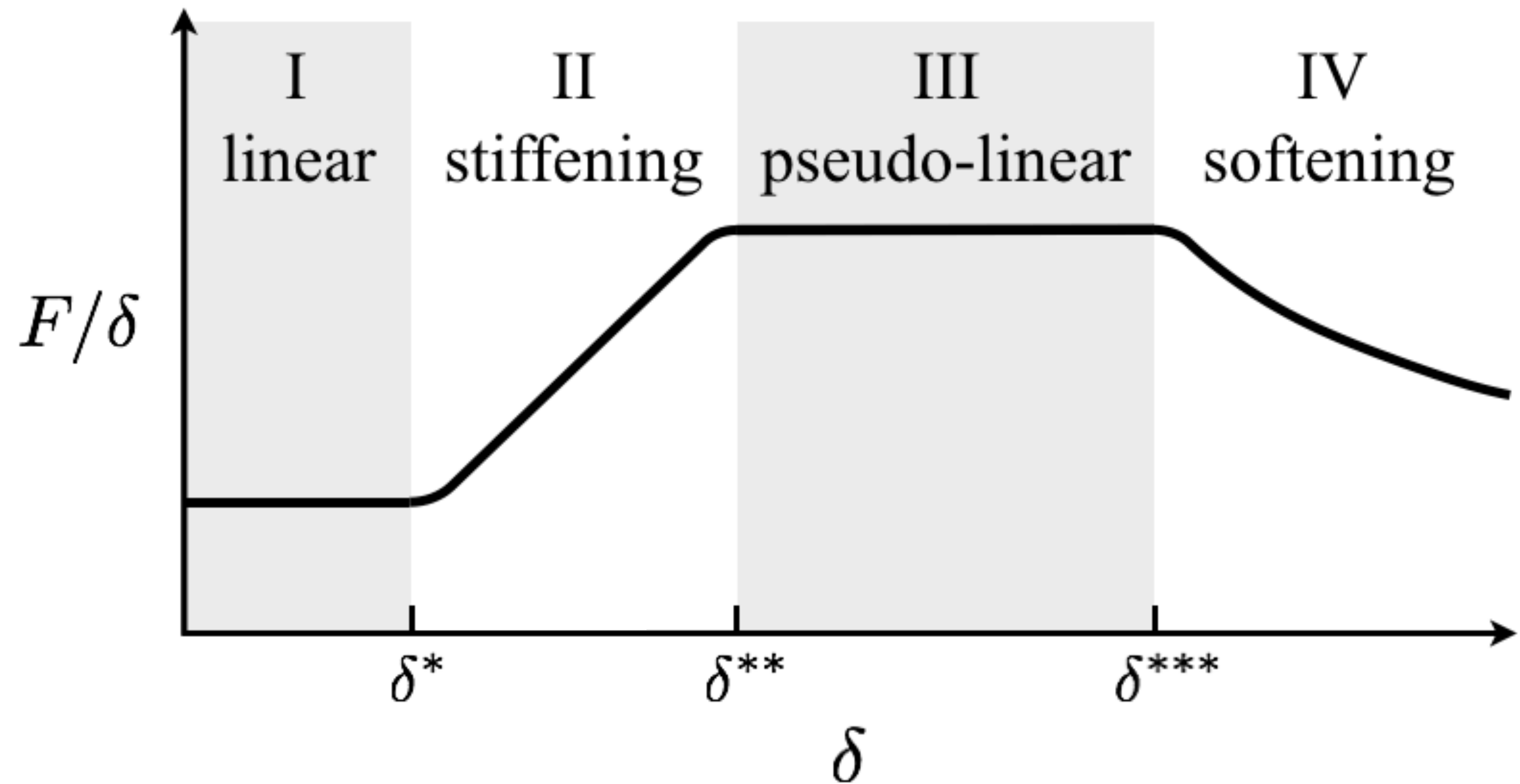


Problem boils down to finding axisymmetric height function: $\zeta(r)$



Vincent Démery

“Spring constant” F/δ , felt by indenter



Force:

I: $F \sim \gamma\delta$

II: $F \sim \sqrt{Y\rho g}\delta^2$

III: $F \approx 4.581(\gamma R_{\text{film}})^{2/3}(\rho g)^{1/3}\delta$

IV: $F \sim 2\pi\gamma R_{\text{film}}$

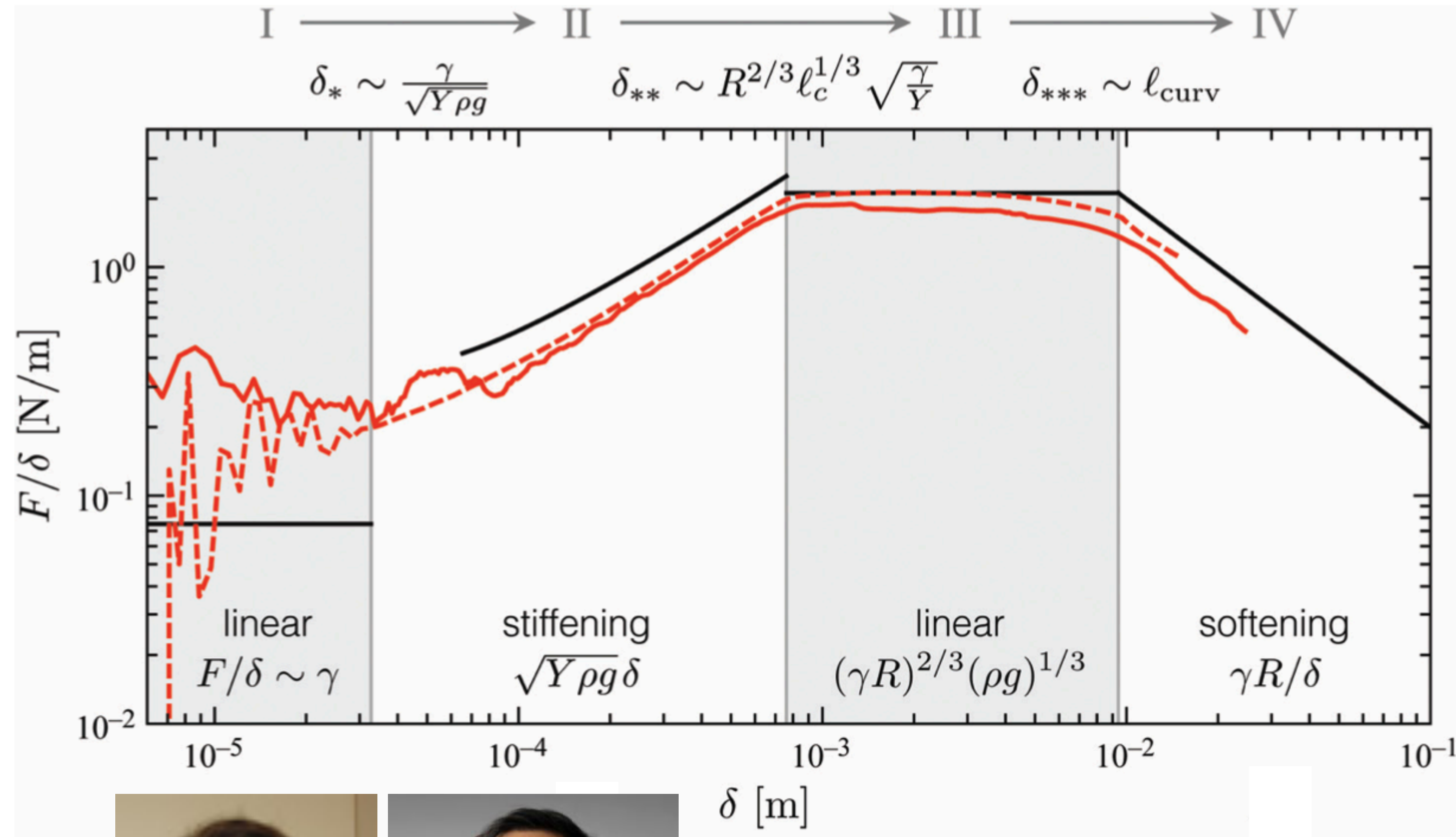
■ Föppl–von Kármán (cf. Benny, Dominic’s talks)

Vella, Davidovitch, et al., *PRL* 2015, *PRE* 2018

■ Geometric model: $U = U_{\text{gravity}} + \gamma(\Delta A_{\text{free}})$

Ripp, Démery, Zhang, JDP, *Soft Matter* 2020

“Spring constant” F/δ , felt by indenter



Force:

- I: $F \sim \gamma\delta$
- II: $F \sim \sqrt{Y\rho g}\delta^2$
- III: $F \approx 4.581(\gamma R_{\text{film}})^{2/3}(\rho g)^{1/3}\delta$
- IV: $F \sim 2\pi\gamma R_{\text{film}}$



Monica Ripp



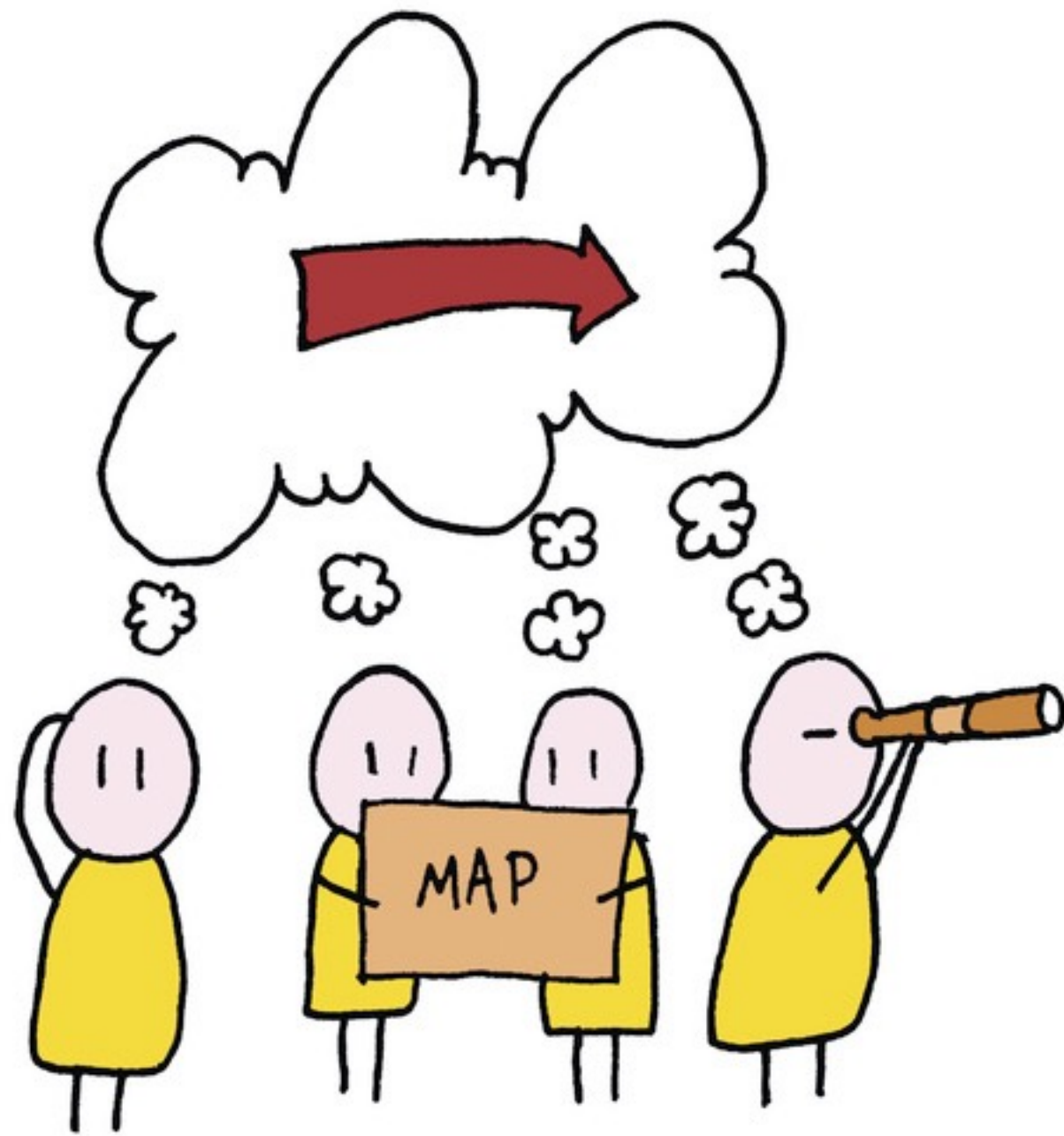
Teng Zhang

■ Föppl–von Kármán (cf. Benny, Dominic’s talks)

Vella, Davidovitch, et al., *PRL* 2015, *PRE* 2018

■ Geometric model: $U = U_{\text{gravity}} + \gamma(\Delta A_{\text{free}})$

Ripp, Démery, Zhang, JDP, *Soft Matter* 2020



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1. Indenting a floating film

Geometric model can predict forces (in wrinkly regimes)

2. Beyond the “gross shape”

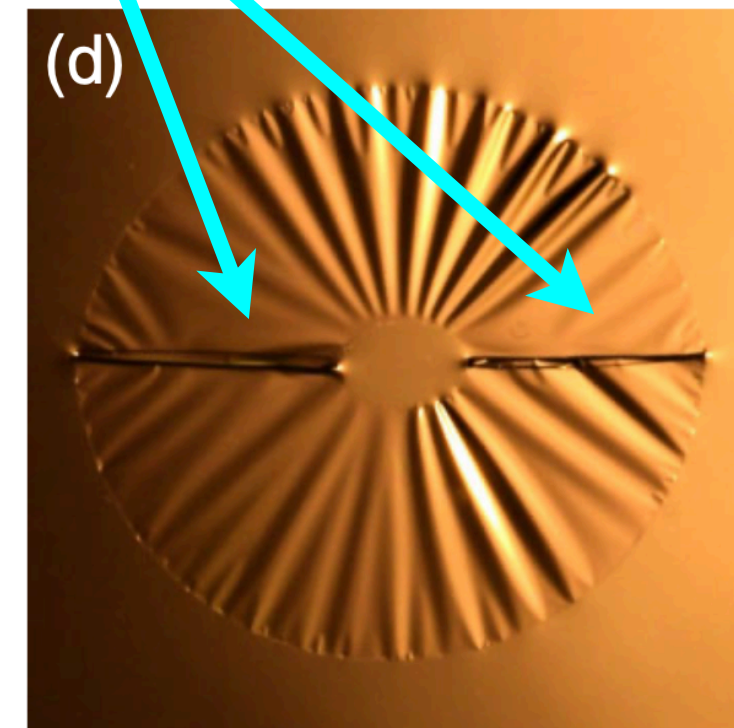
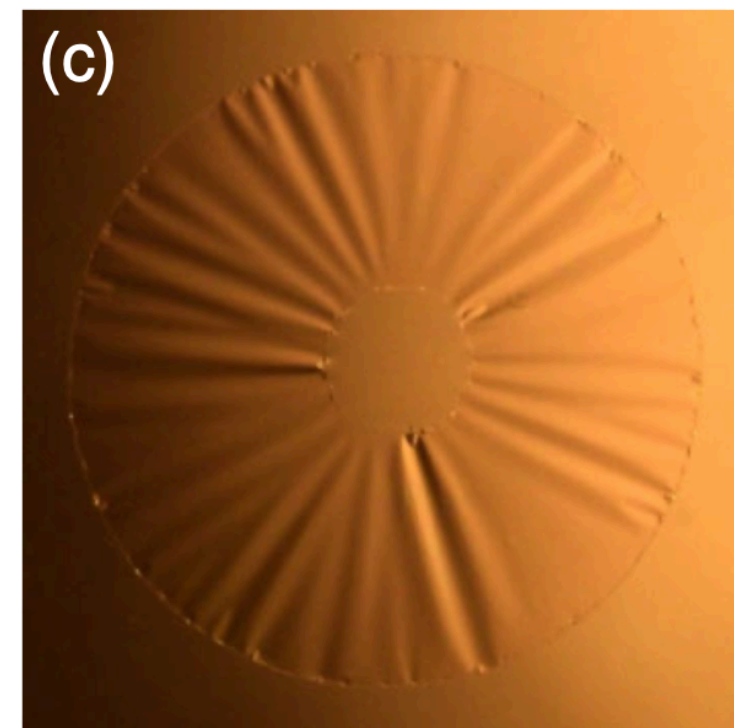
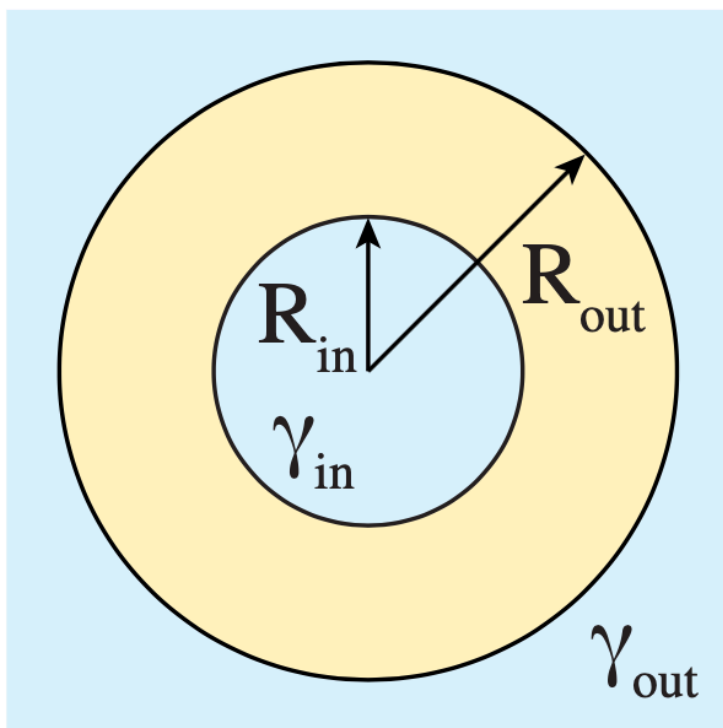
3. Towards solid surfactants

A very reasonable outlook:

The geometric model $U = \gamma A$ is great for explaining gross shapes, forces,
BUT it doesn't predict small-scale features (wrinkles, crumples, folds),
i.e., how excess length is stored

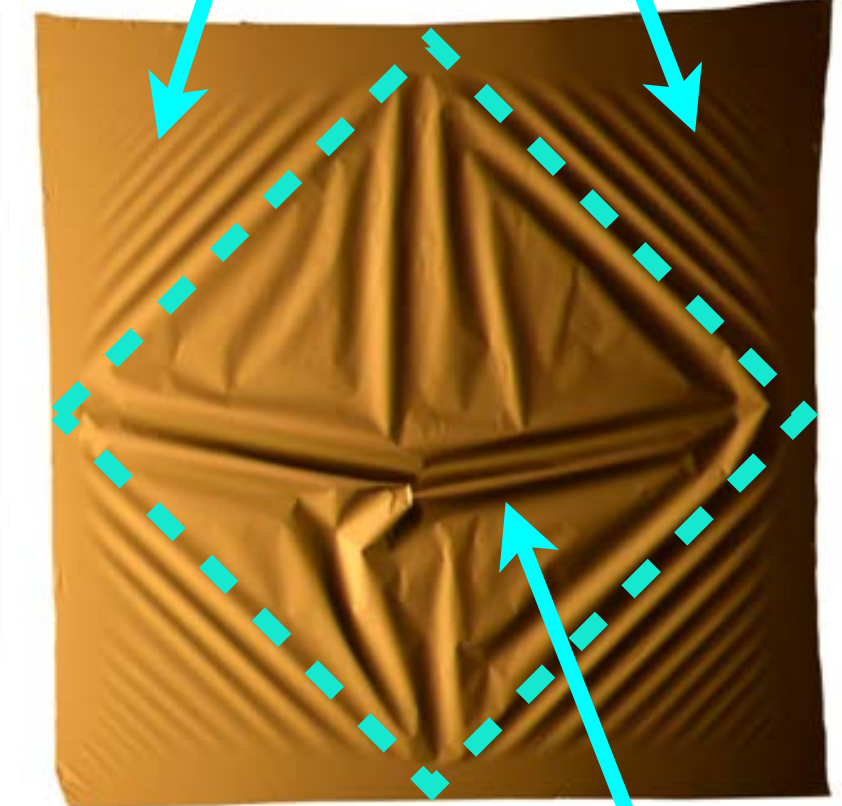
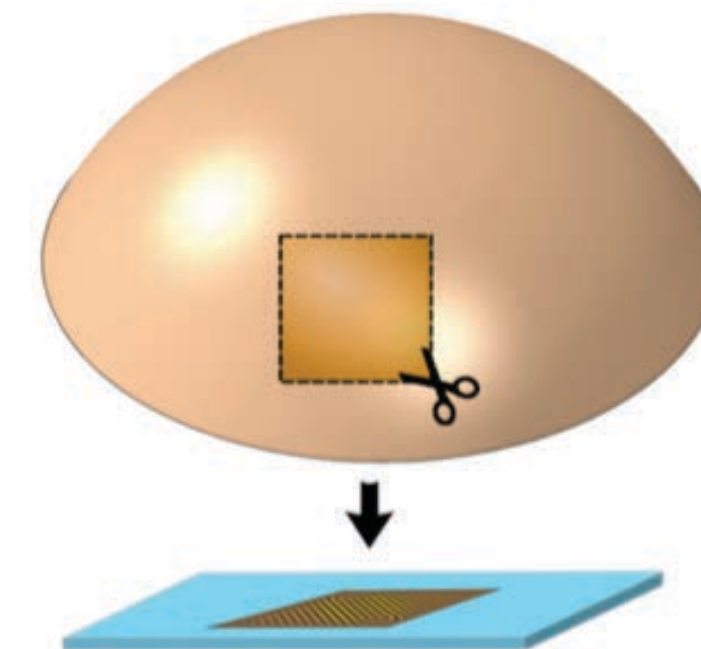
Recent progress:

Annular film on a flat liquid bath:



optimal gross shape
implies folds (vs. wrinkles)

Stamping a curved shell onto a plane:

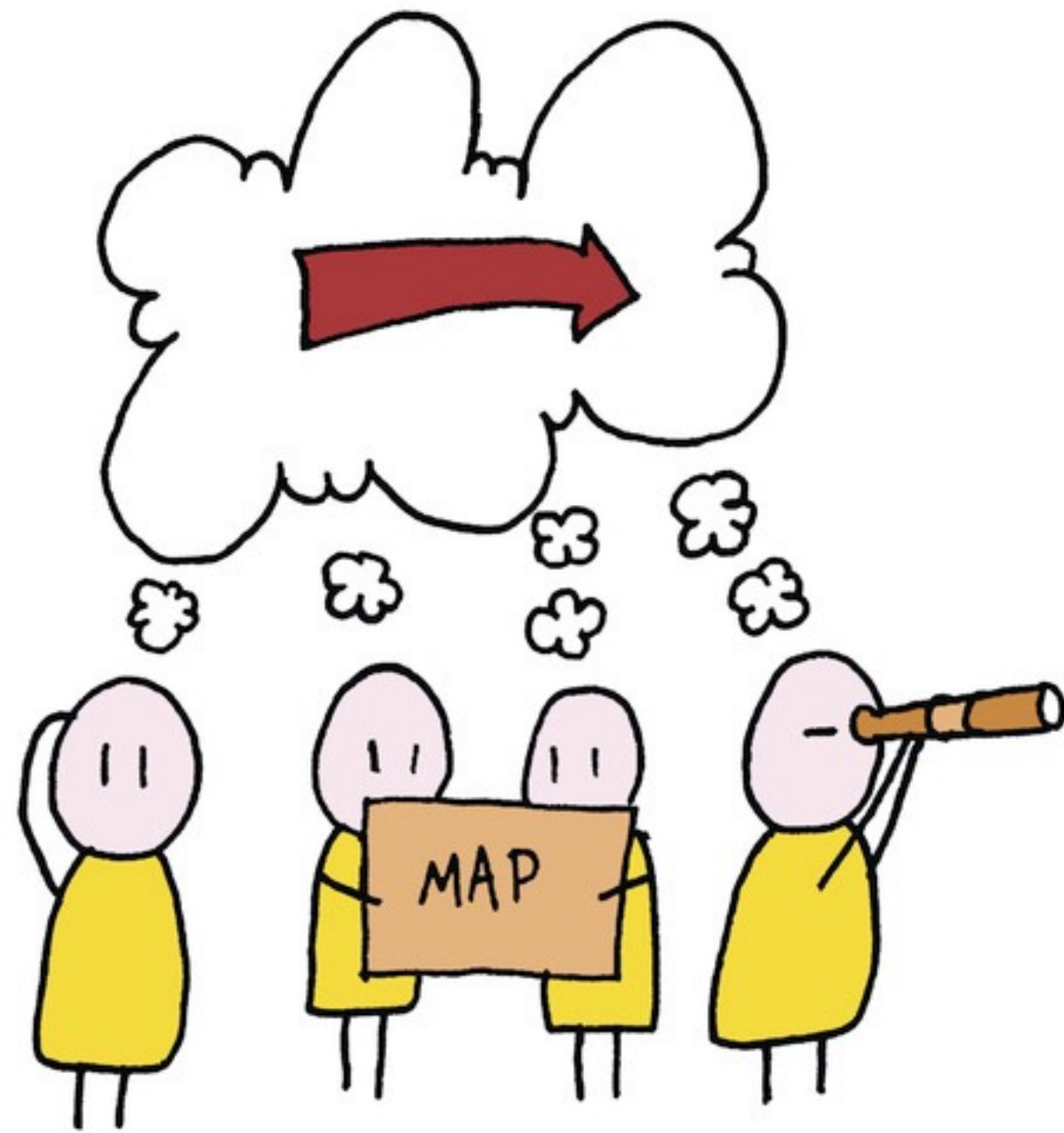


wrinkle direction **fixed** by
optimal gross shape

wrinkles **free**

JDP, Démercy, Davidovitch, Menon, et al., PRL 118 (2017)

Tobasco, JDP, Katifori, et al. (arXiv 2004.02839)



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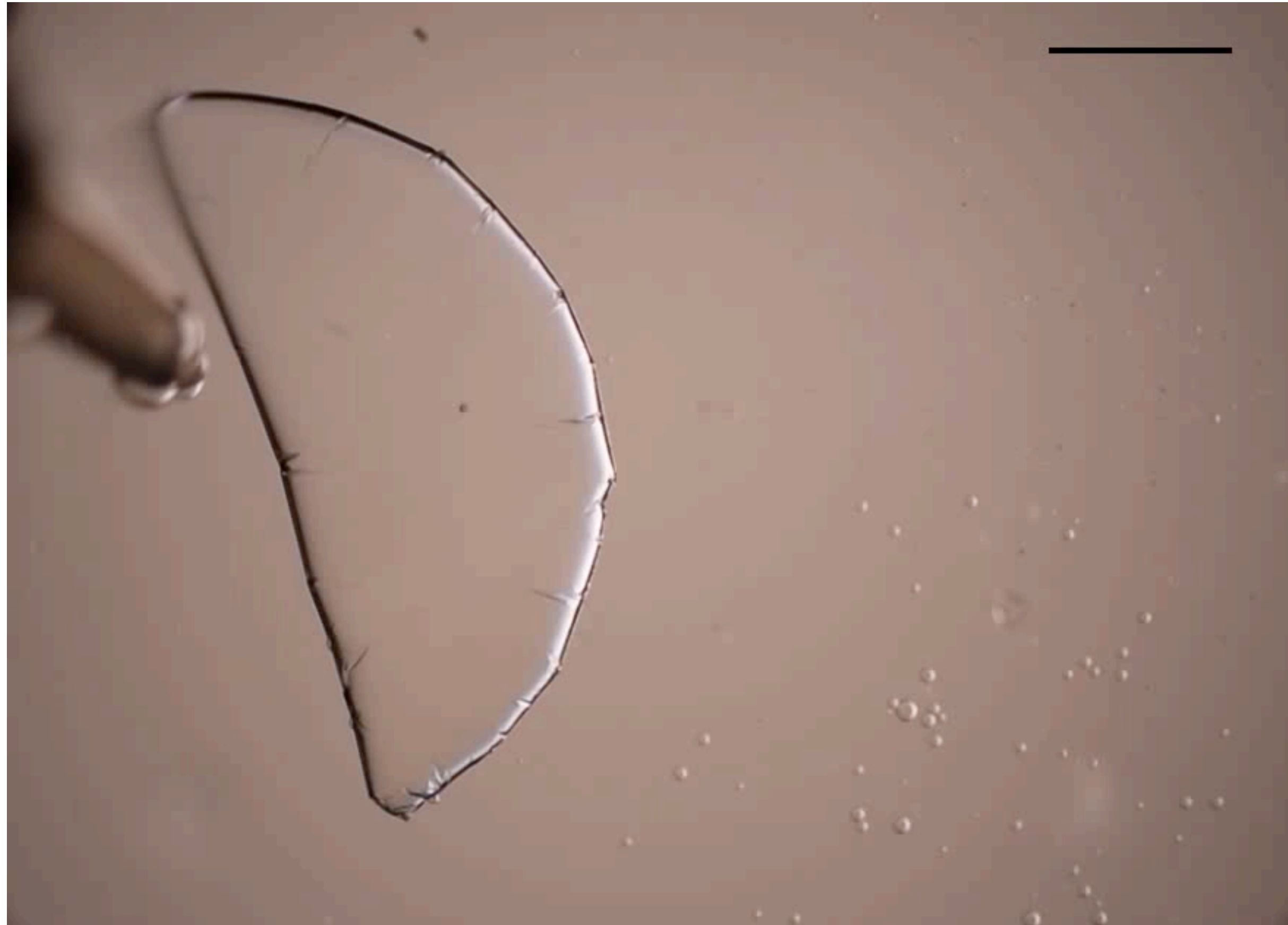
Geometric model can predict forces (in wrinkly regimes)

2. Beyond the “gross shape”

Some folding transitions and wrinkle layouts may be understood as a geometric optimization

3. Towards solid surfactants

Elastic films as “**solid surfactants**”



Complementary to molecular & particle surfactants

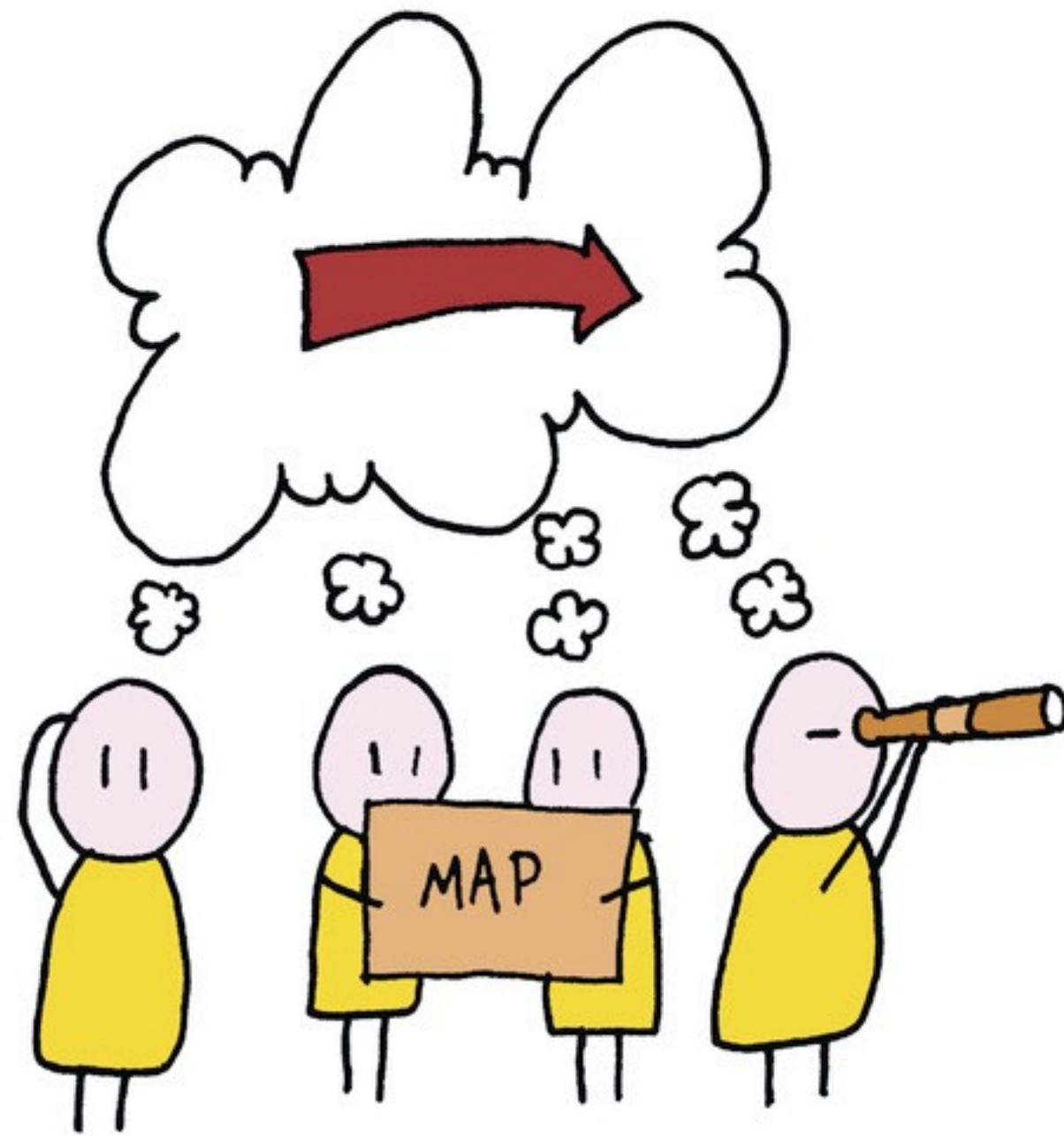
Interesting properties for applications:

- Achieve non-spherical shapes
- Tailor mechanical, optical properties
- Platform for surface patterning (physical, chemical)
- Sequester/protect liquid cargo

Conclusions

0. Wrapping a droplet with a thin solid

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1. Indenting a floating film

Geometric model can predict forces (in wrinkly regimes)

2. Beyond the “gross shape”

Some folding transitions and wrinkle layouts may be understood as a geometric optimization

3. Towards solid surfactants

Ultrathin sheets: A tool for tailoring droplets, emulsions, interfaces (shape, rheology, surface chemistry, ...)

“Wrapping liquids, solids, and gases in thin sheets”

JDP, Annu. Rev. Condens. Matter Phys. Vol. 10 (2019), **arXiv:1804.07425**