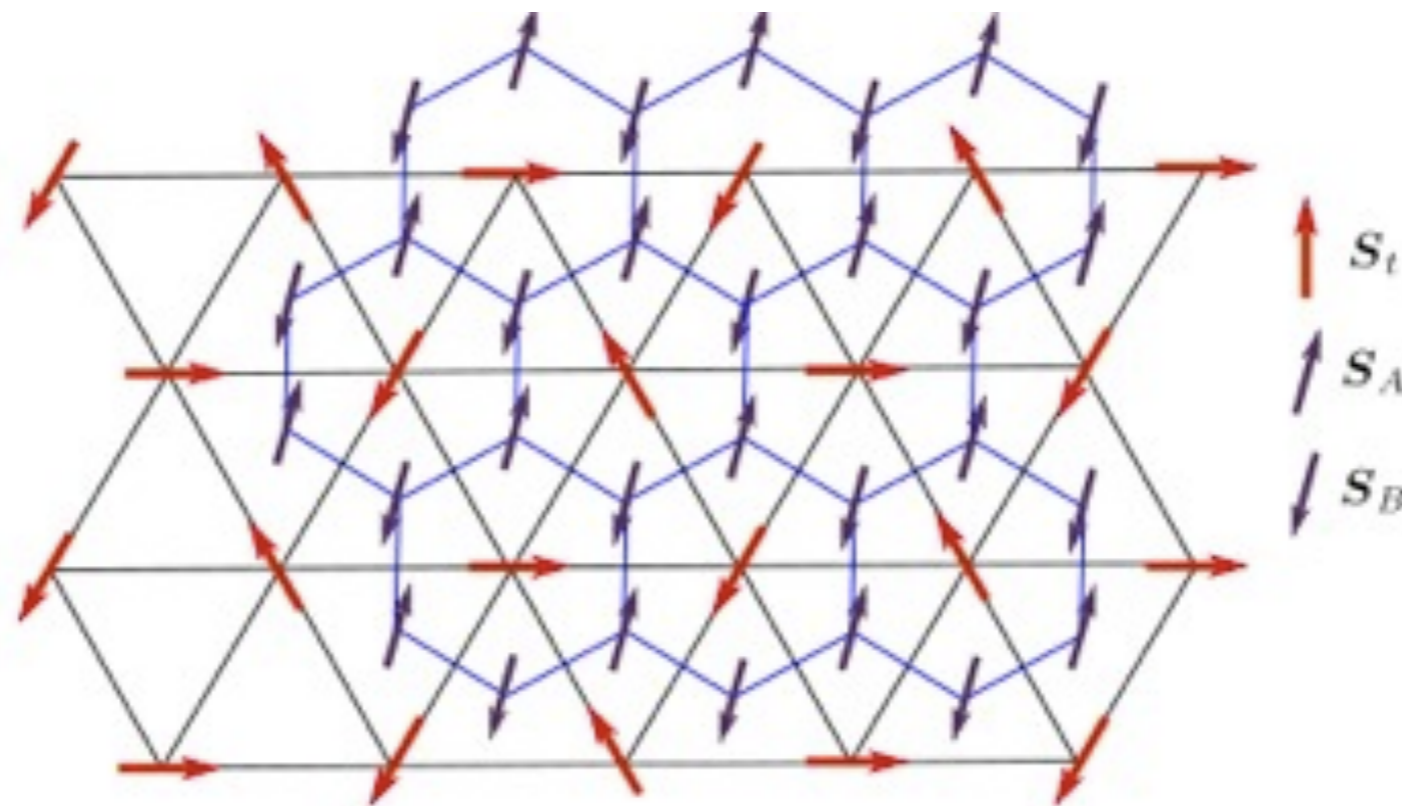


Emergent Critical Phase & Ricci Flow in a 2D Frustrated Heisenberg AFM

[arXiv:1206.5740](https://arxiv.org/abs/1206.5740)

Peter Orth
Premi Chandra
Piers Coleman
Joerg Schmalian



discussions +
Daniel Friedan



Exotic Phases of Frustrated
Magnets, KITP, Oct 8, 2012.



2D Heisenberg Antiferromagnets at Finite Temperature

2D Heisenberg Antiferromagnets at Finite Temperature

Hohenberg-Mermin-Wagner
Theorem (1966)



No Long-Range Order
at Finite Temperatures

$$\xi \sim a e^{\frac{2\pi J}{kT}}$$

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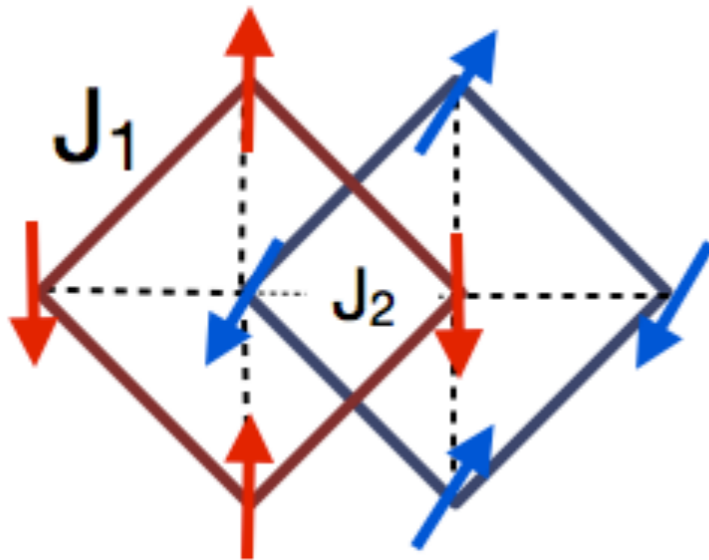


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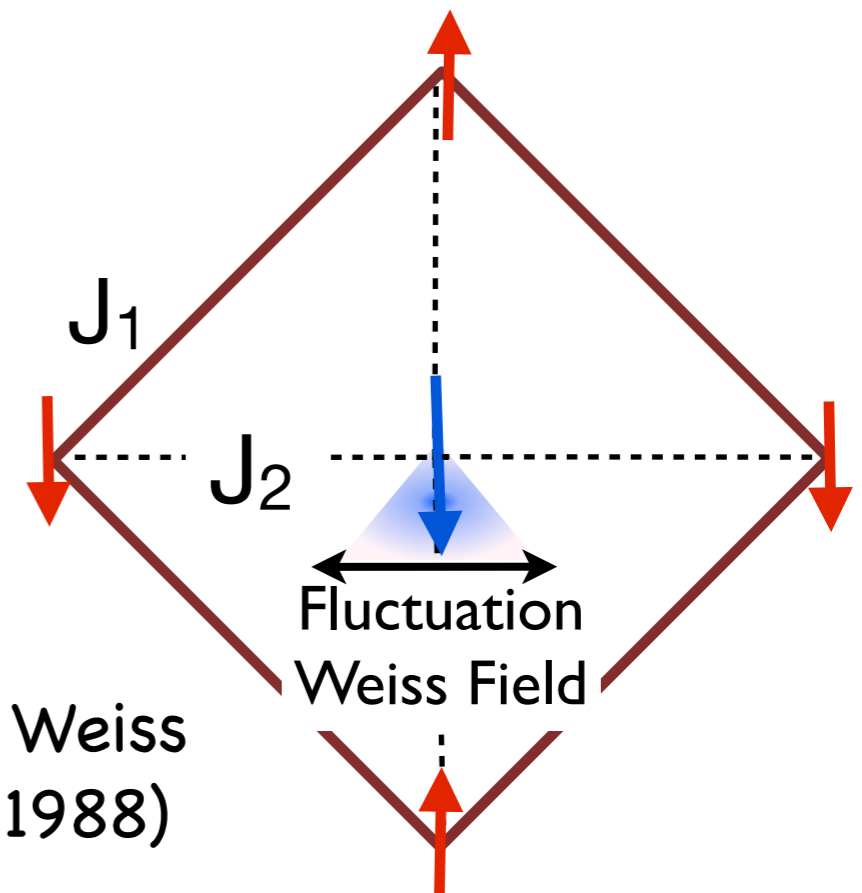
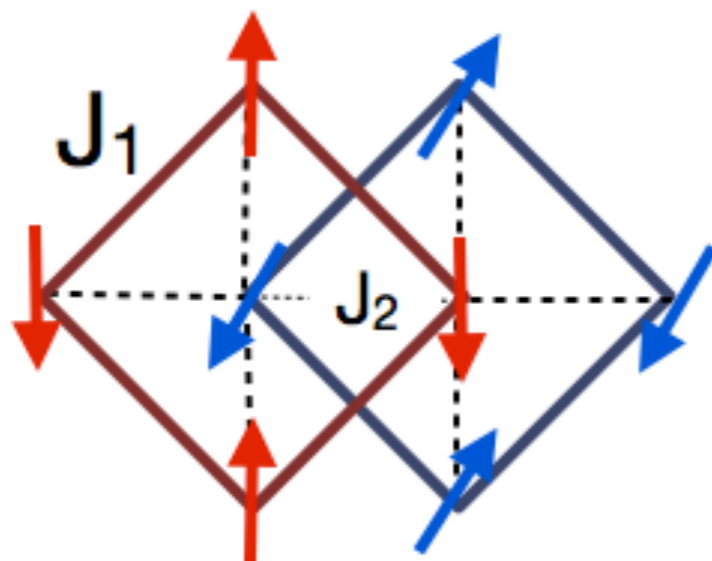


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Spins like to align the **fluctuation** Weiss field to their "easy plane". (Henley 1988)

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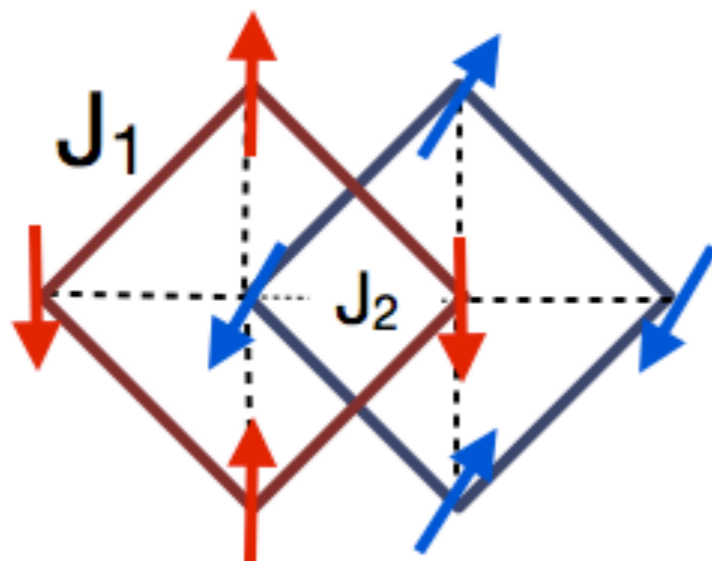


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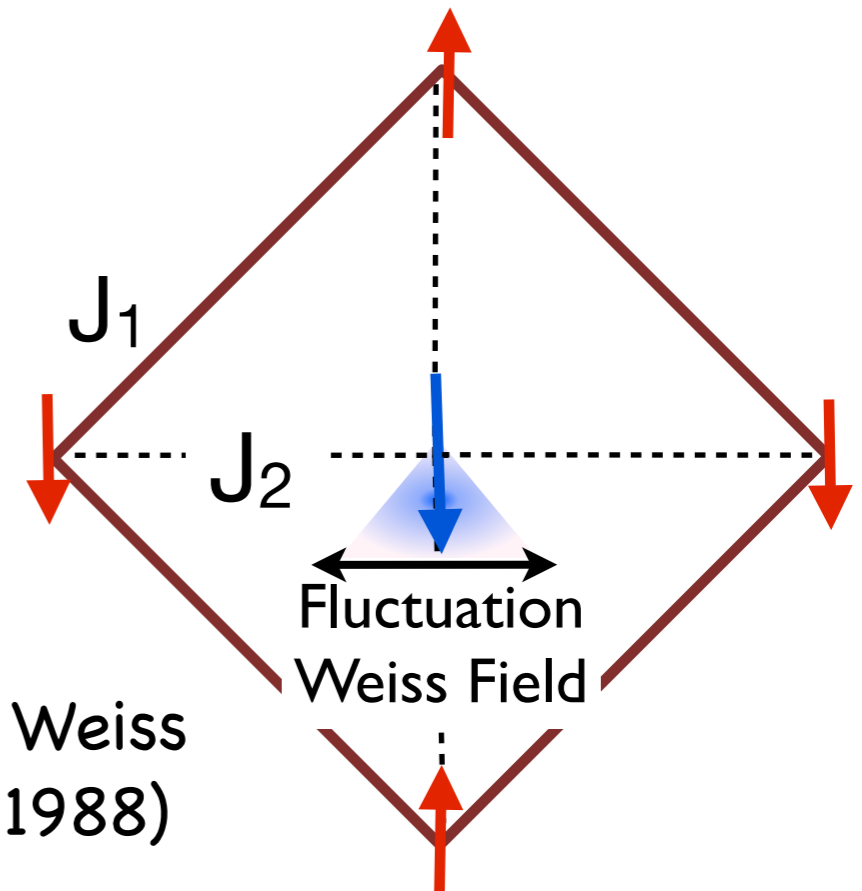
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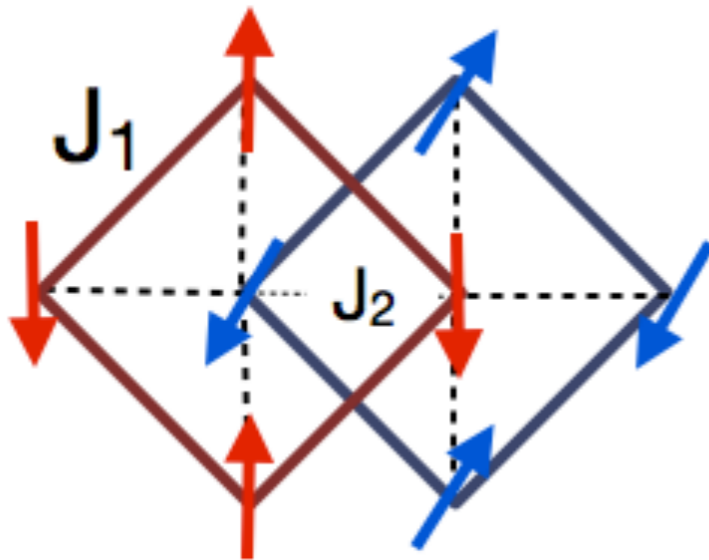


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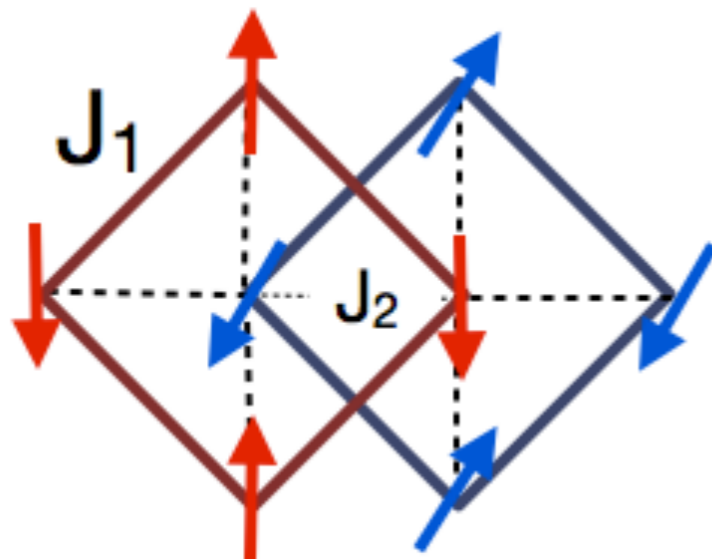


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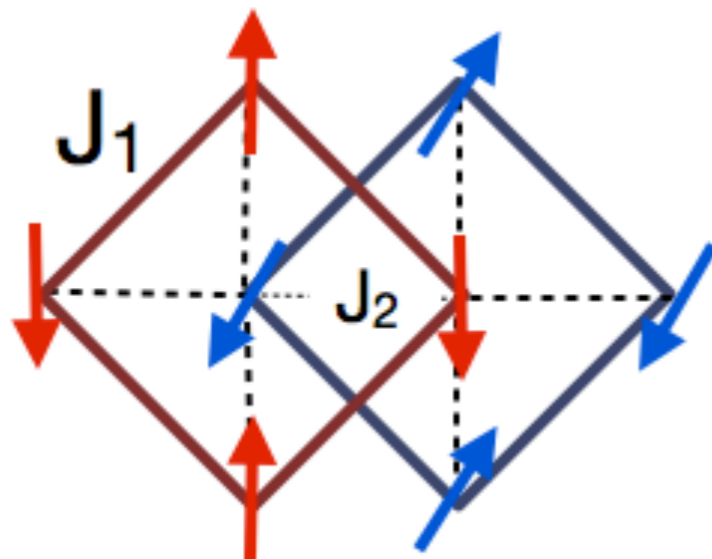


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$$\xi_{Z_2} \rightarrow \infty$$

$$T_{Z_2}$$

Chandra, Coleman and Larkin (90)

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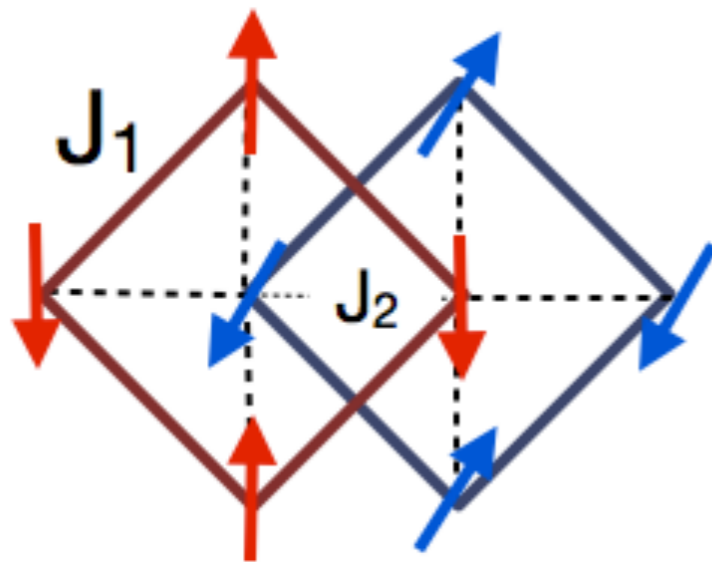


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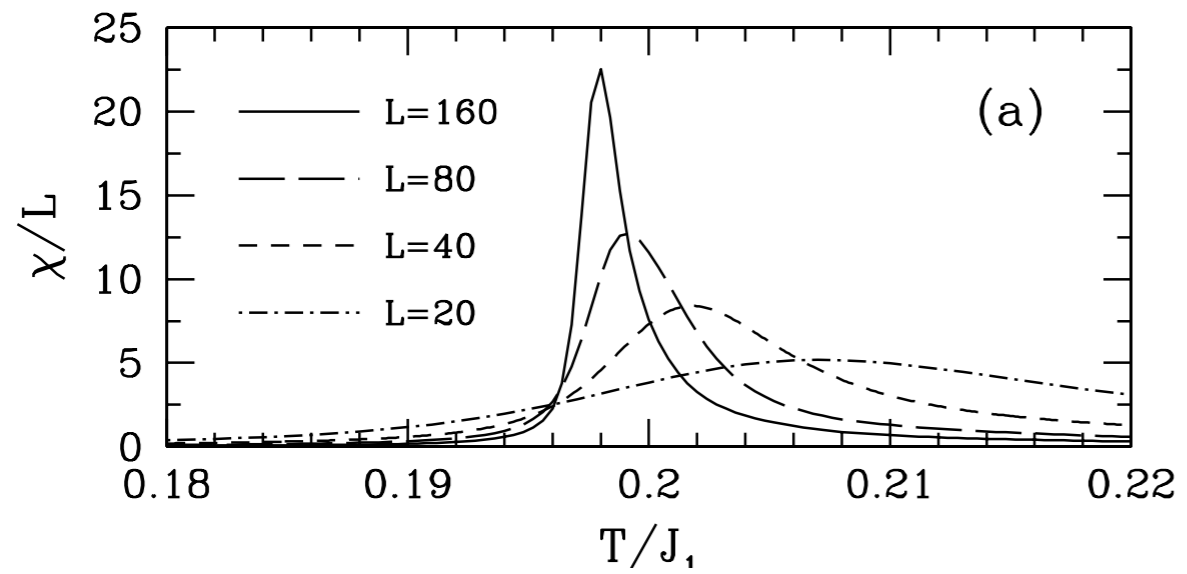
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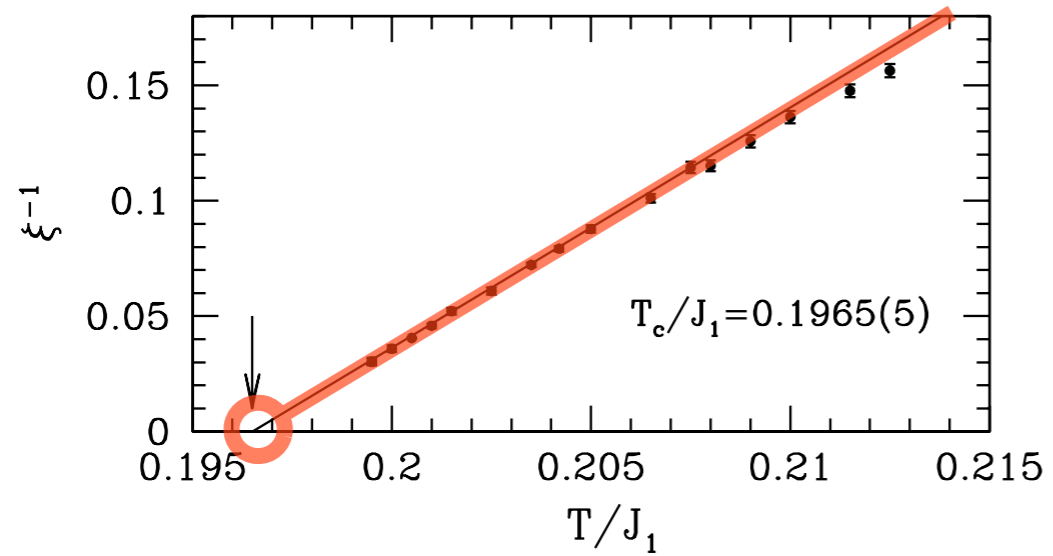
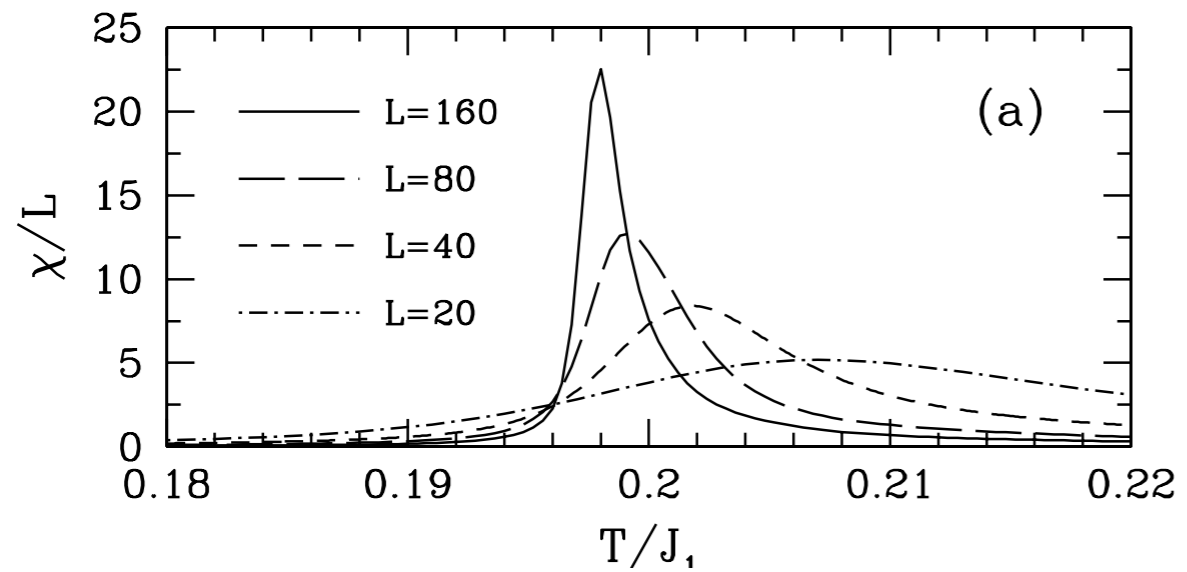
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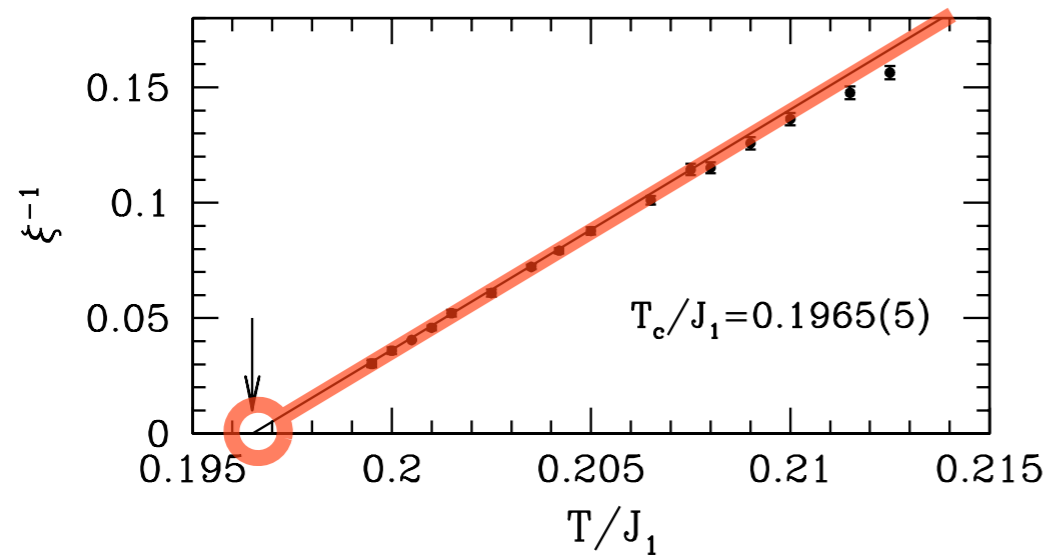
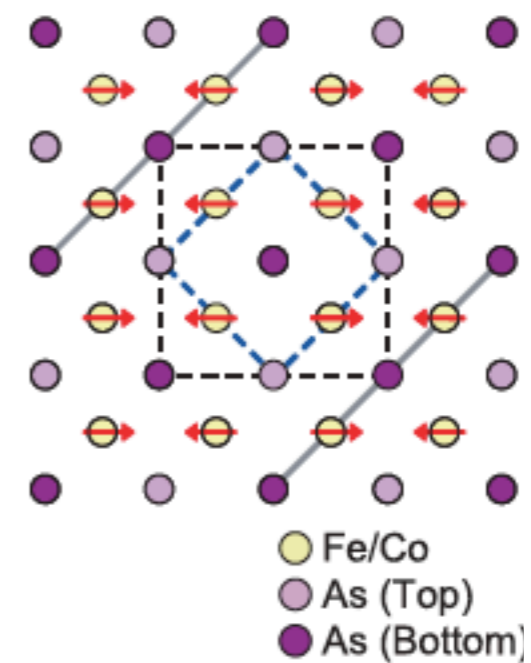
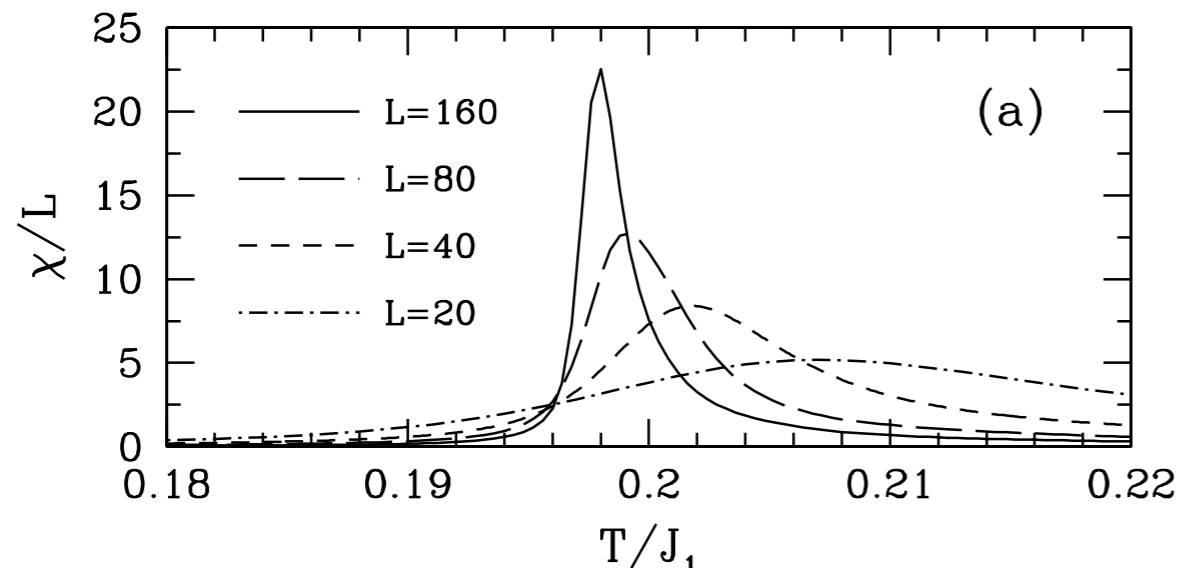
Emergent Z_2 Phase Transition in a disordered Heisenberg System.



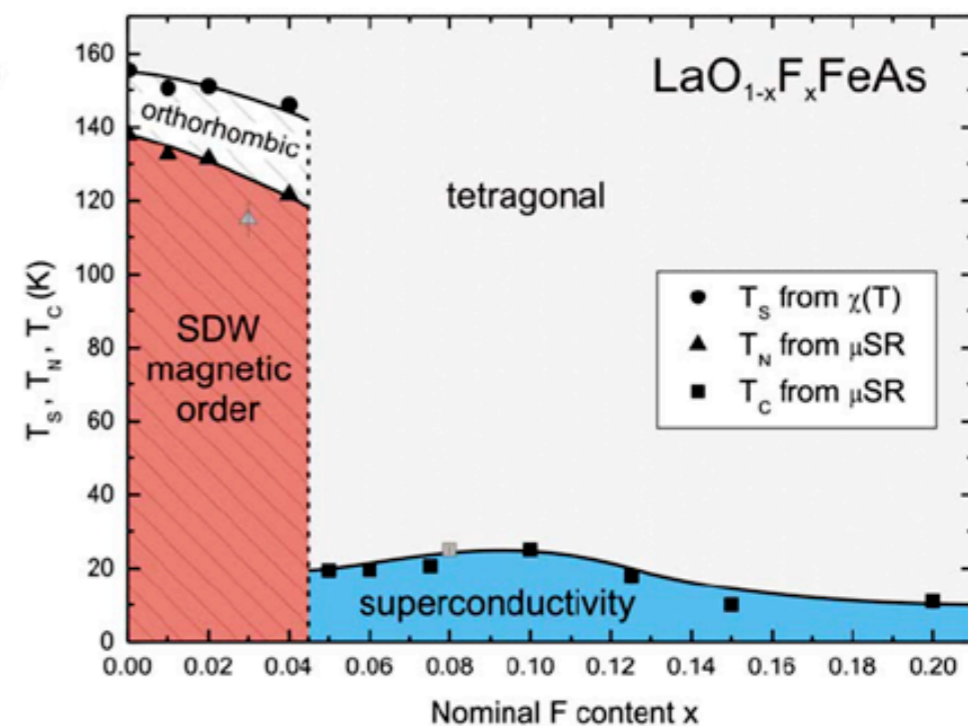
Weber et al (2003)



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Iron based superconductors (2008).

Z_2 generalized to Z_p ??

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$p \geq 5$ \longrightarrow Kosterlitz-Thouless Transition

Jose et al (77)

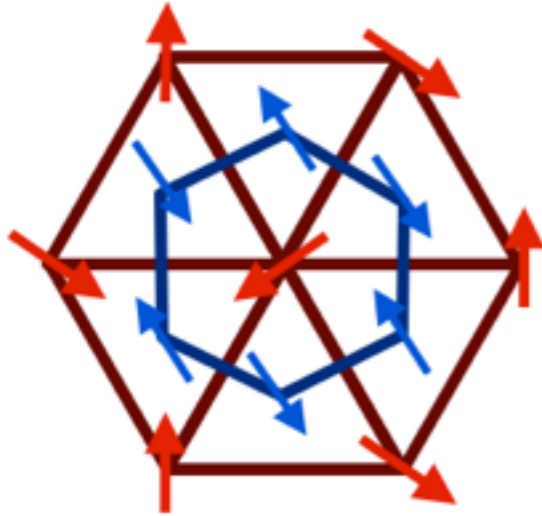
Z_2 generalized to Z_p ??

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Jose et al (77)

Can we find a model that has an emergent critical phase even though its underlying Heisenberg degrees of freedom have a finite correlation length?

2D Heisenberg AFM Hamiltonian



$$H = H_{hh} + H_{tt} + H_{th}$$

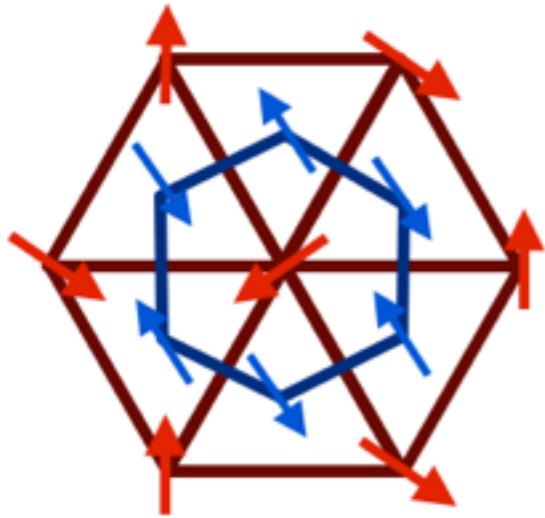
$$H_{\alpha\beta} = J_{\alpha\beta} \sum_{j=1}^{N_L} \sum_{\delta_{\alpha\beta}} S_{\alpha}(j) S_{\beta}(j + \delta_{\alpha\beta})$$

$$J_{tt} = J_{hh} = 1$$

$$J_{th} \ll 1$$

$$\alpha, \beta \in \{t, A, B\}$$

2D Heisenberg AFM Hamiltonian



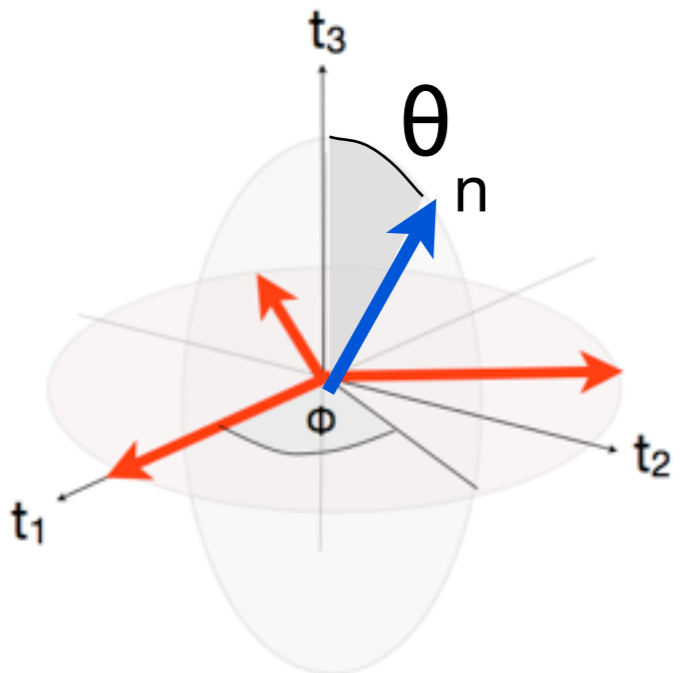
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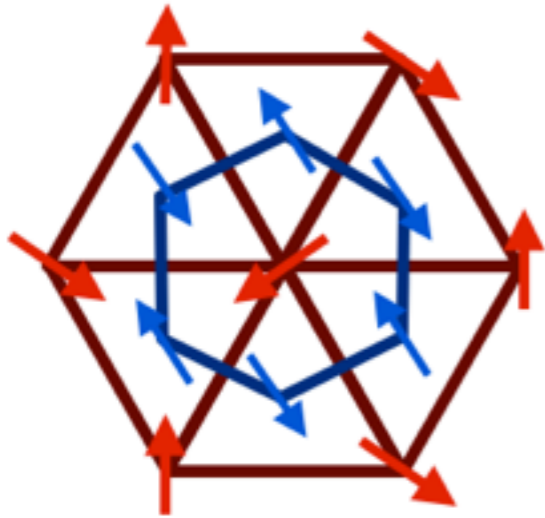
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Classically: two decoupled sublattices.

2D Heisenberg AFM Hamiltonian



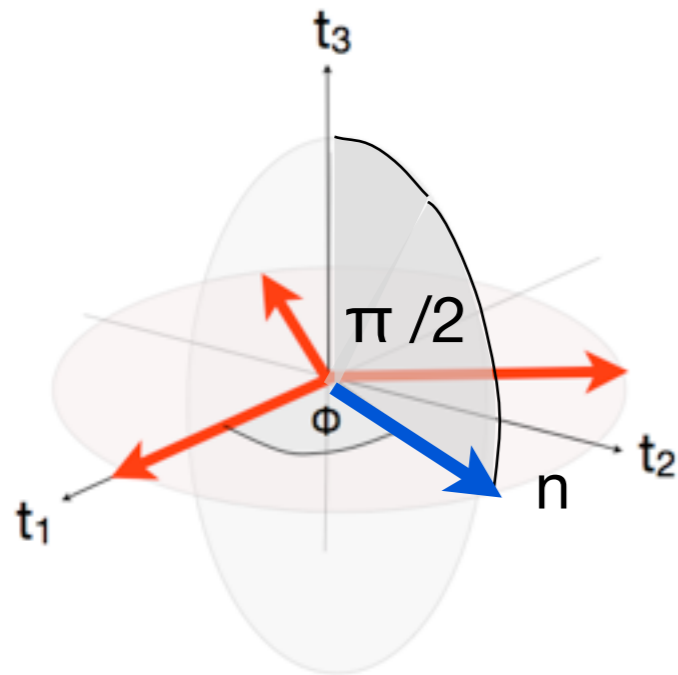
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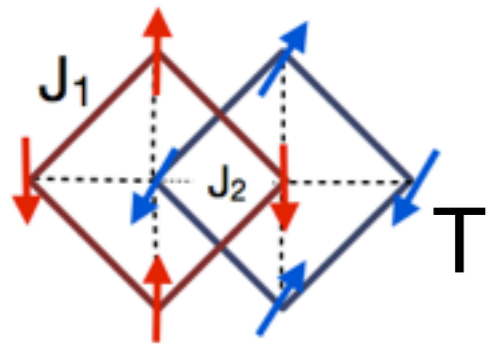
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Order from disorder drives coplanarity.



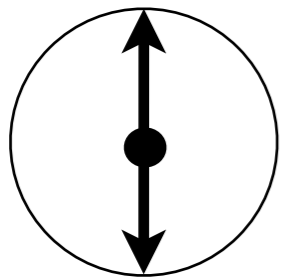
free moments

J_1

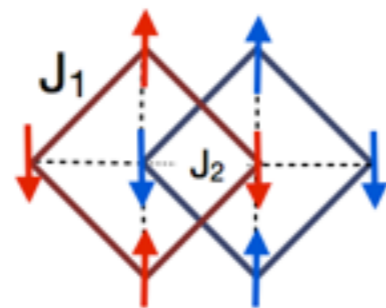
decoupled sublattices

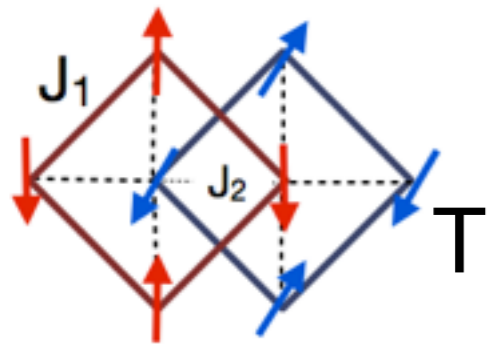
$J_1/\ln(J_1/J_2)$

T_{Z2}
Collinearity



Ising Z_2 order





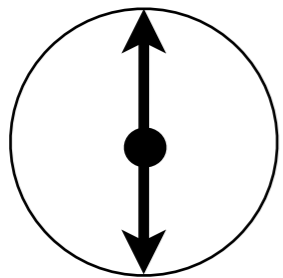
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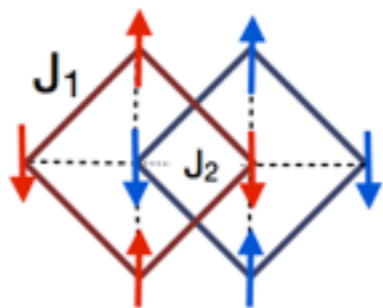
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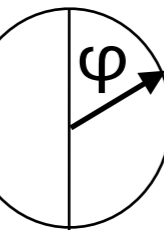
T_{BKT}

T_{Z6}

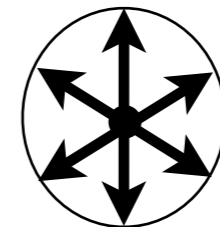
T

J_1

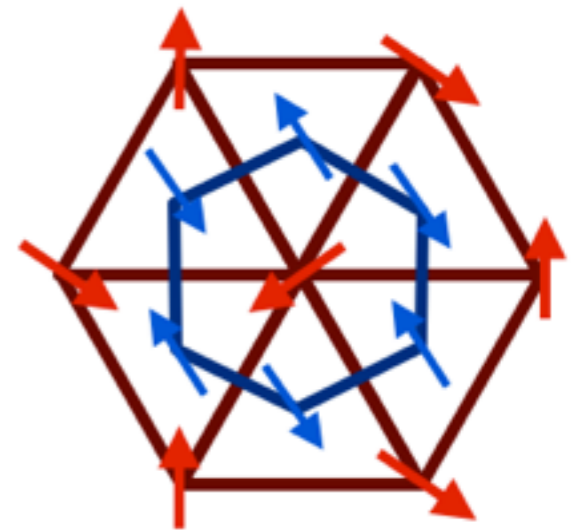
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Power law correlated xy Order



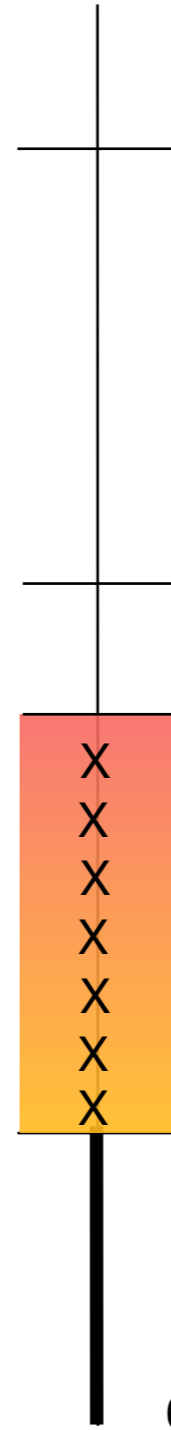
6-state Clock Order



(I)

T_{BKT}

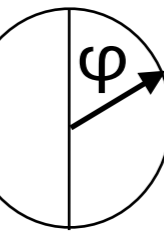
T_{Z6}



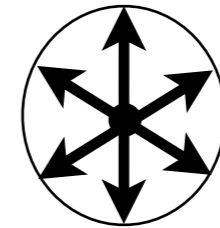
T

J_1

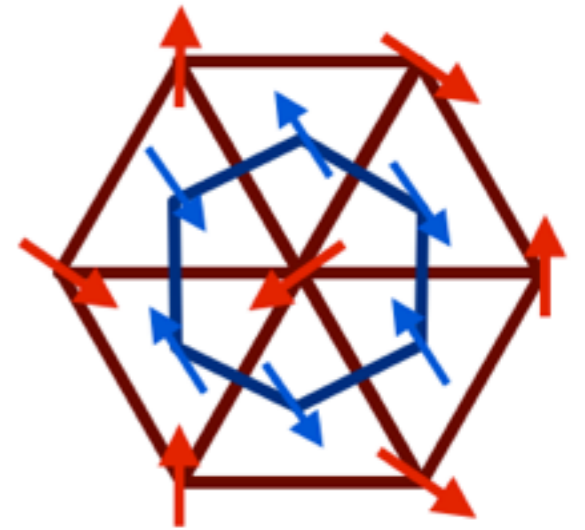
$J_1/\ln(J_1/J_2)$
Coplanarity

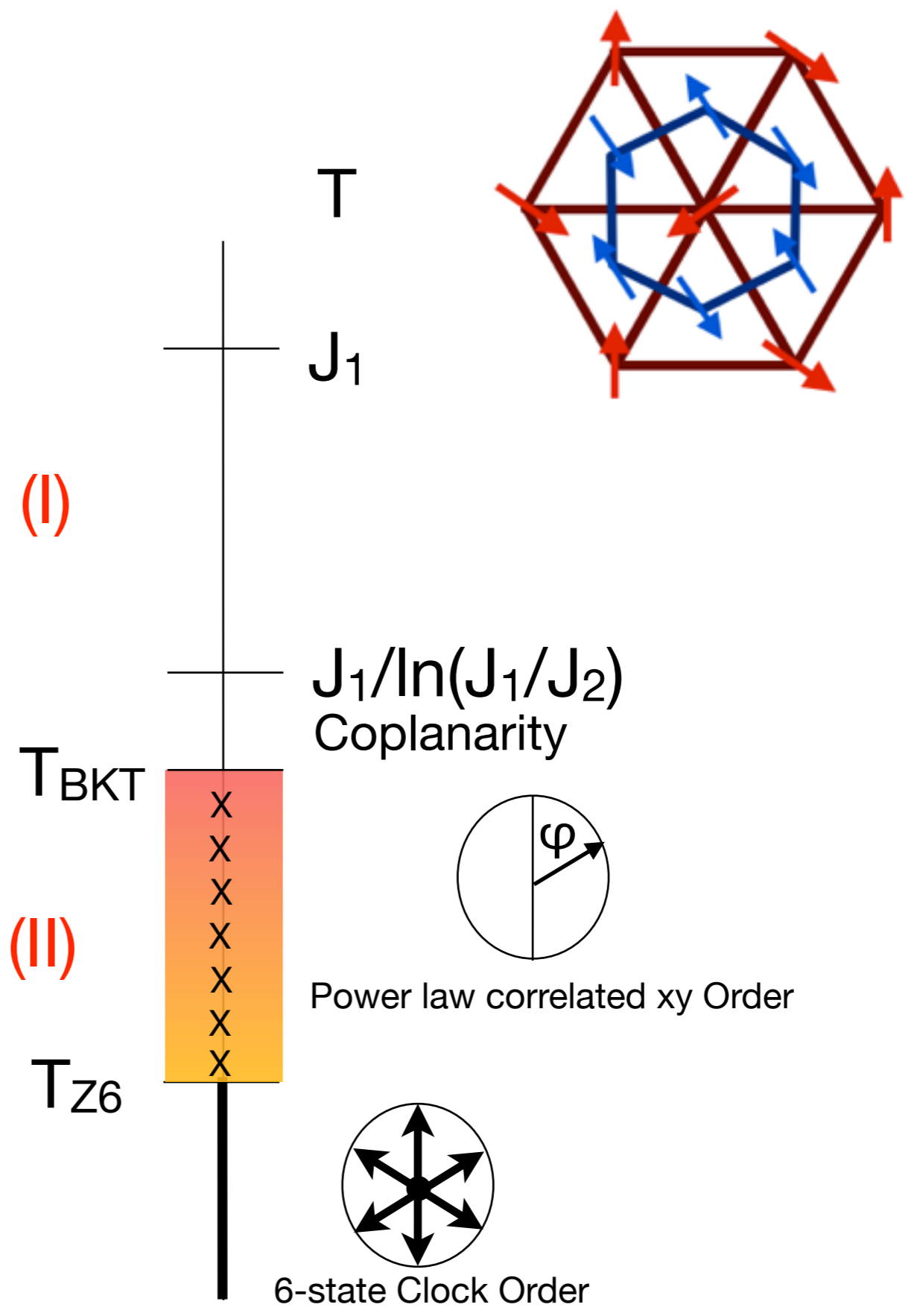


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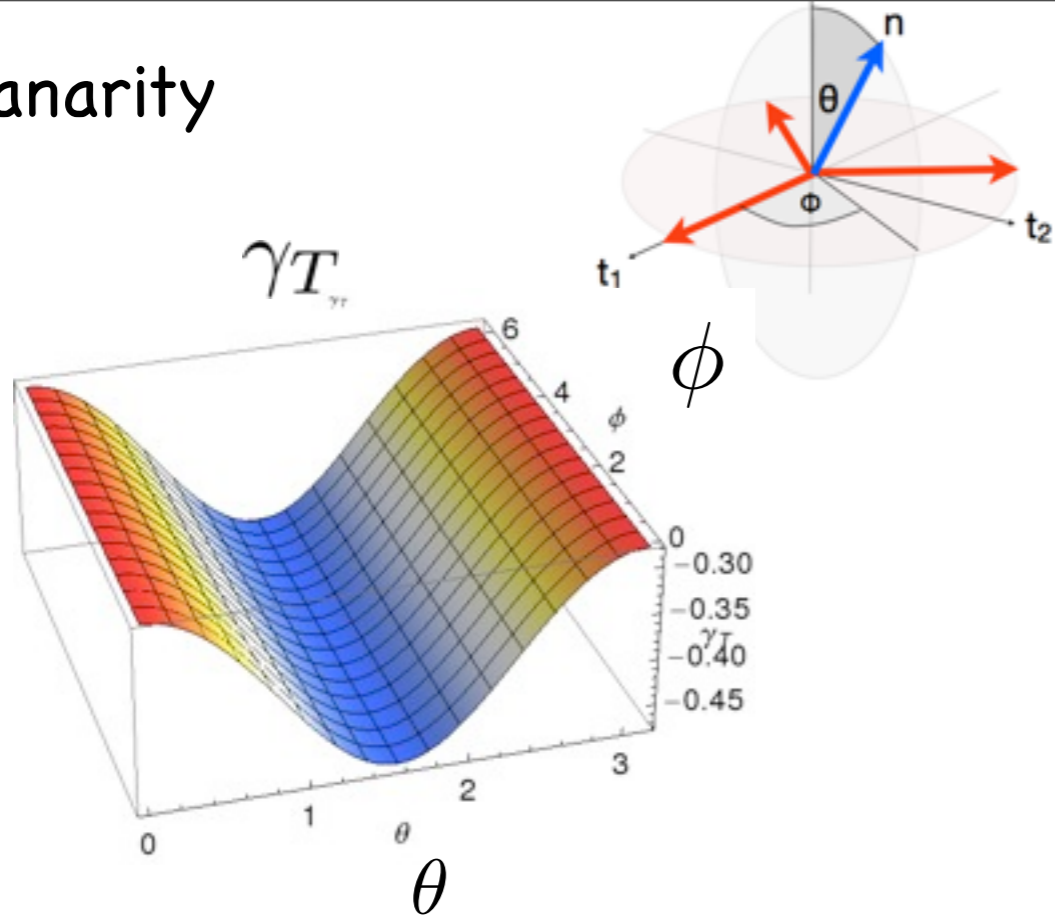


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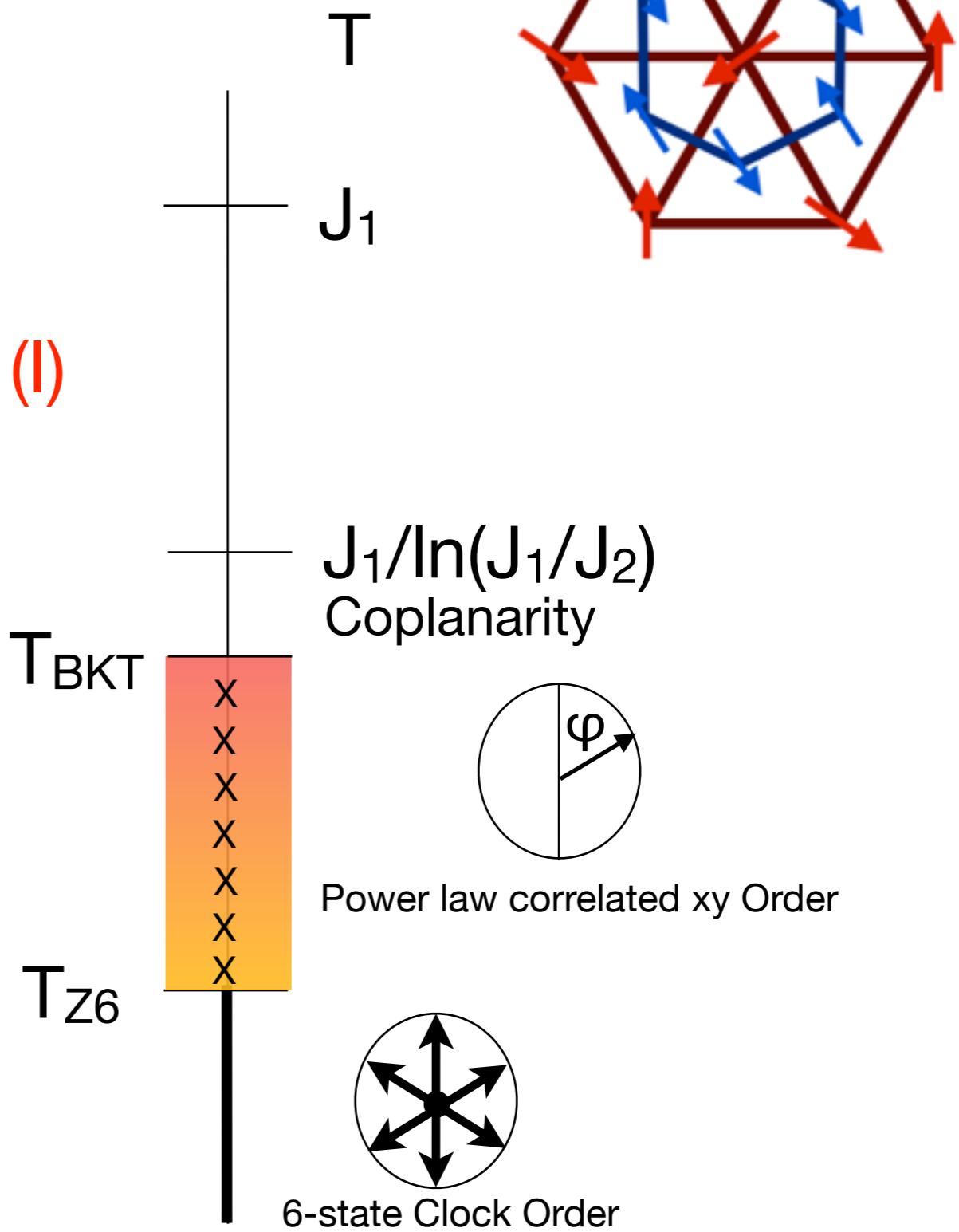
Coplanarity



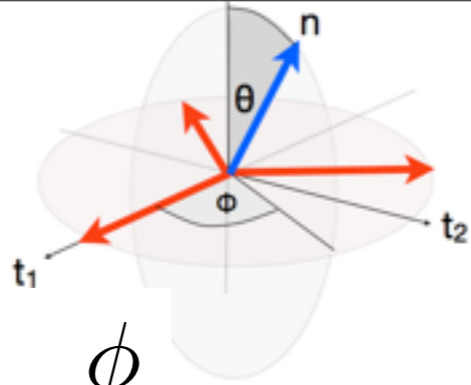
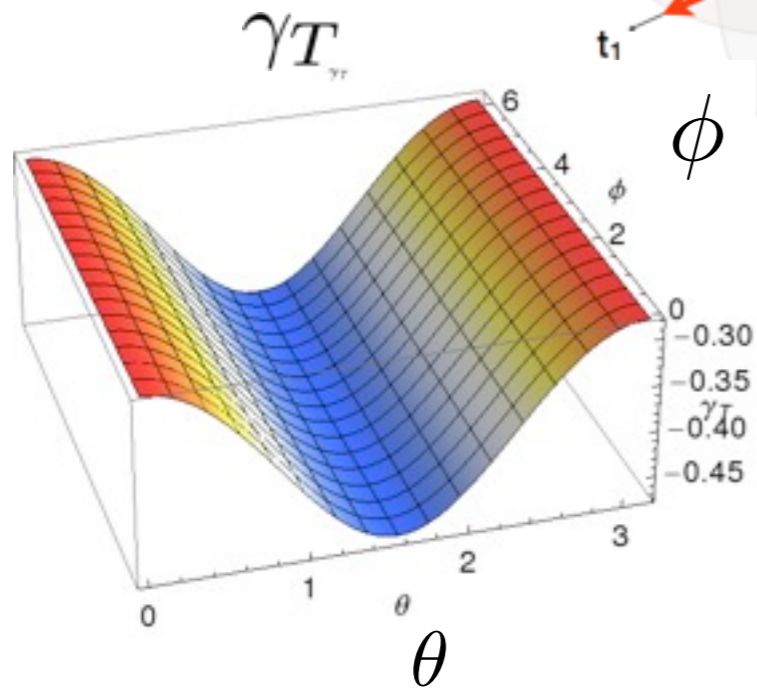
$$\delta F(\theta, \phi) = \sum_{\alpha, p \in MBZ} E_{\alpha}(p) \left(\langle B_{\alpha, p}^{\dagger} B_{\alpha, p} \rangle + \frac{1}{2} \right)$$

$$\delta F(\theta, \phi) = \gamma_T(\theta, \phi) T$$

(I)



Coplanarity



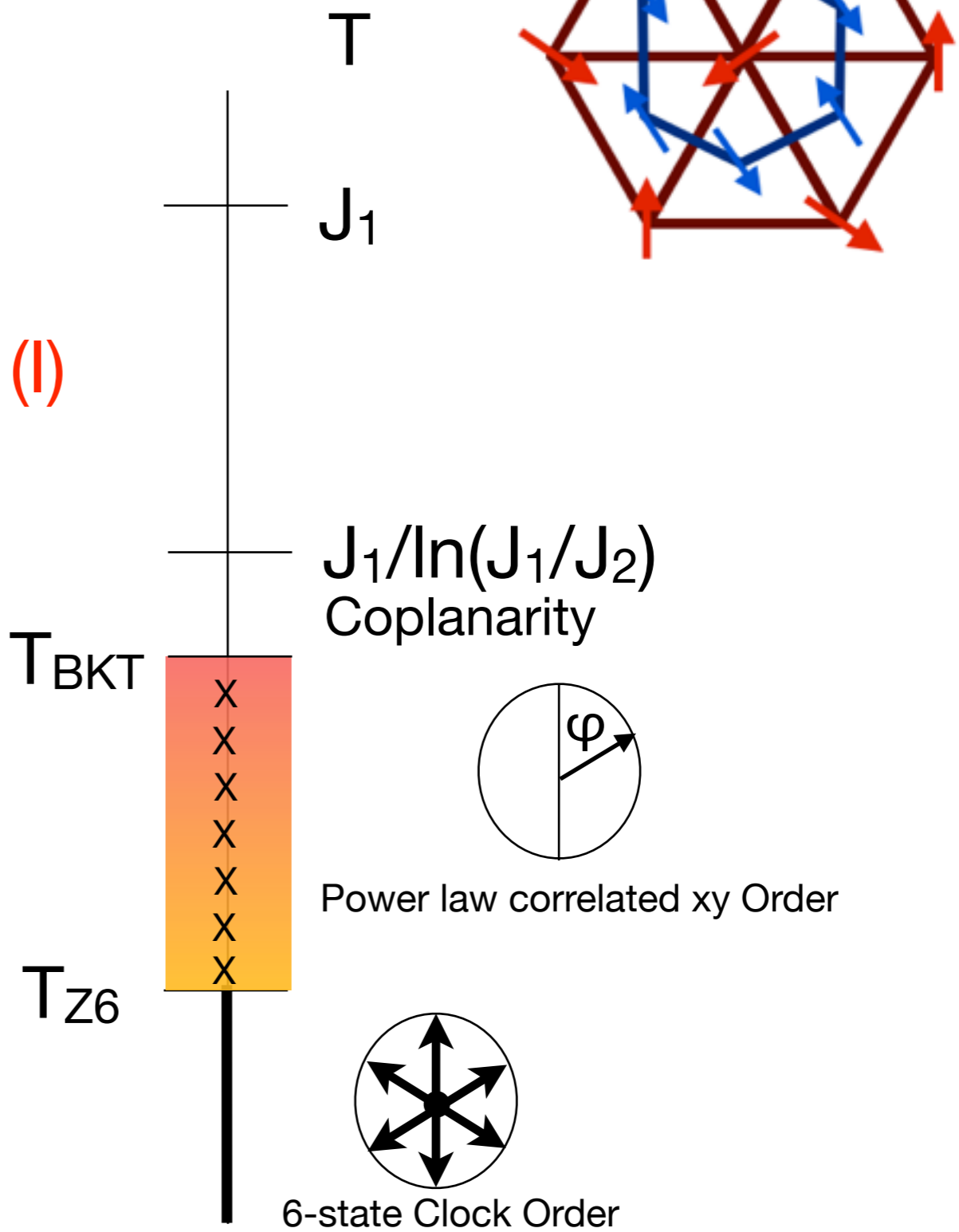
Minimum occurs at $\theta = \frac{\pi}{2}$

→ Coplanar Spin Ordering

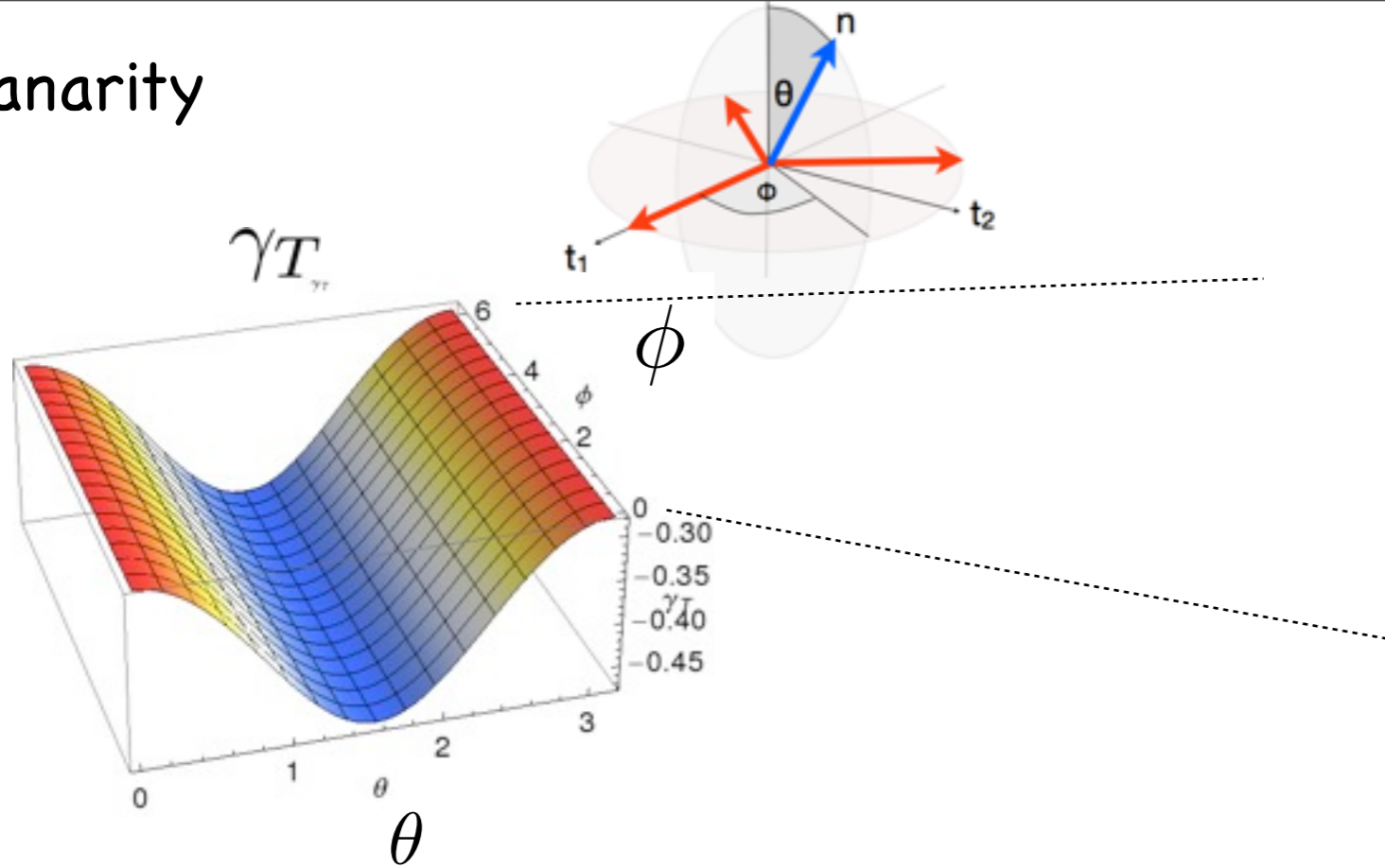
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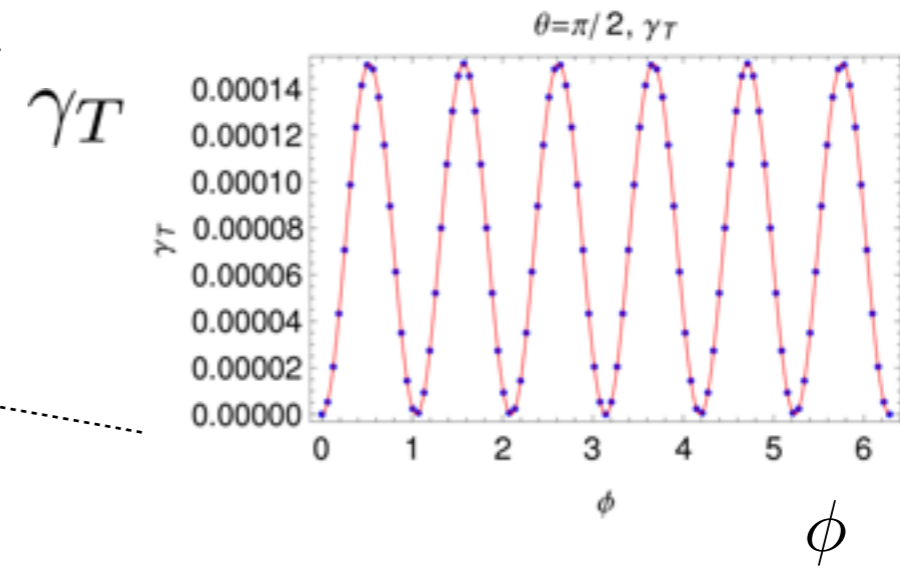
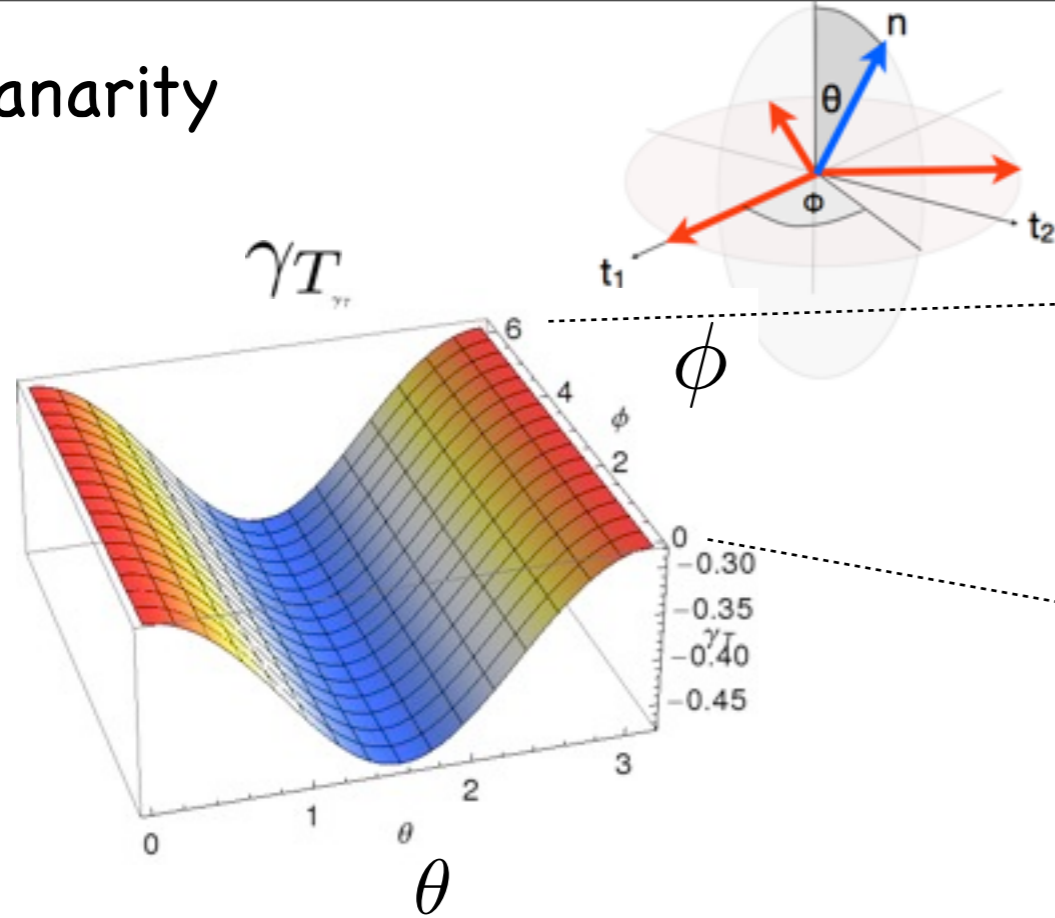
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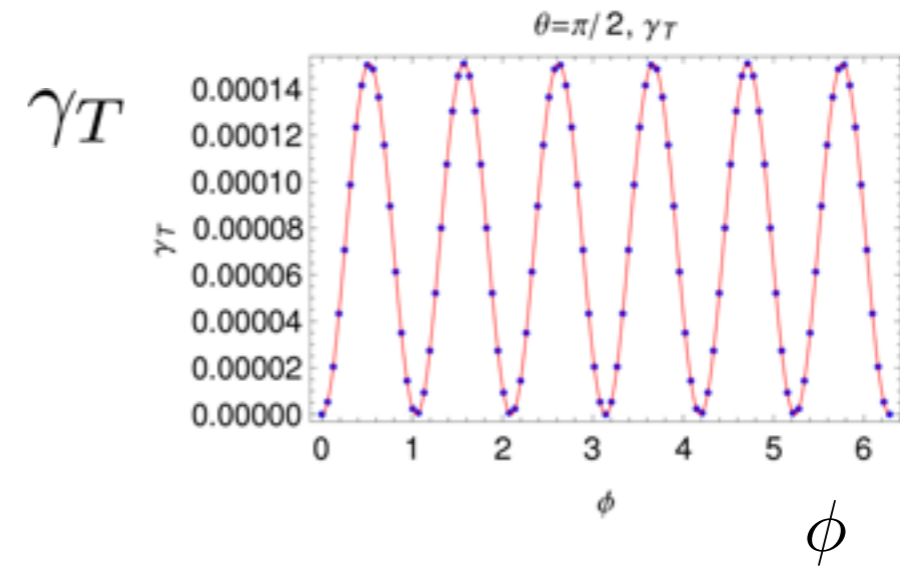
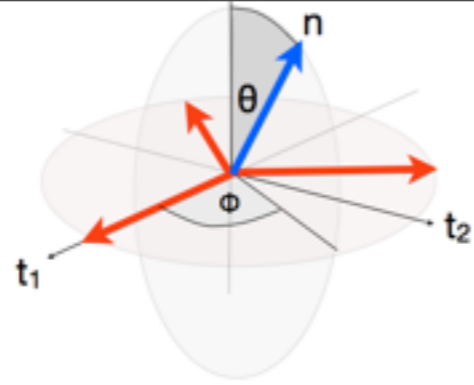
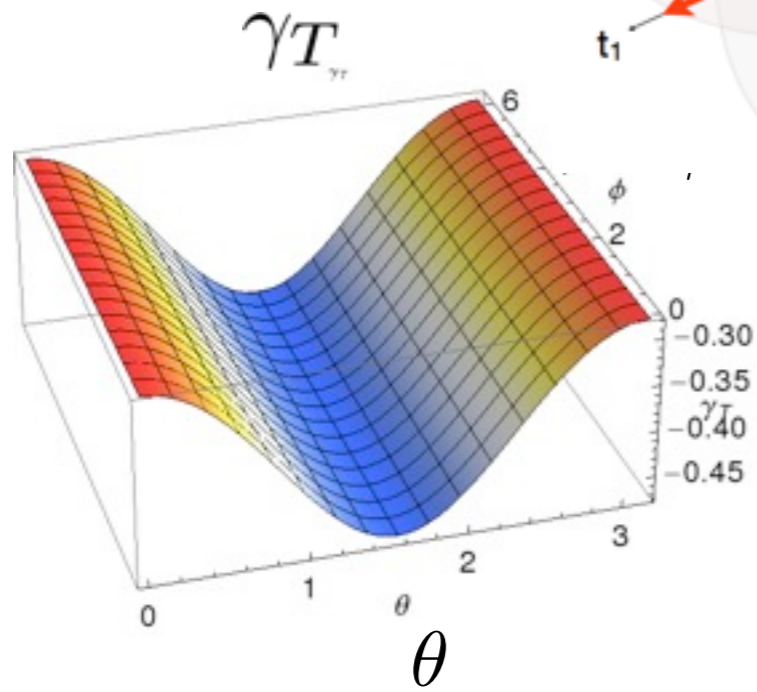
Six Degenerate Minima at $\phi = \frac{2\pi n}{6}$

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Coplanarity

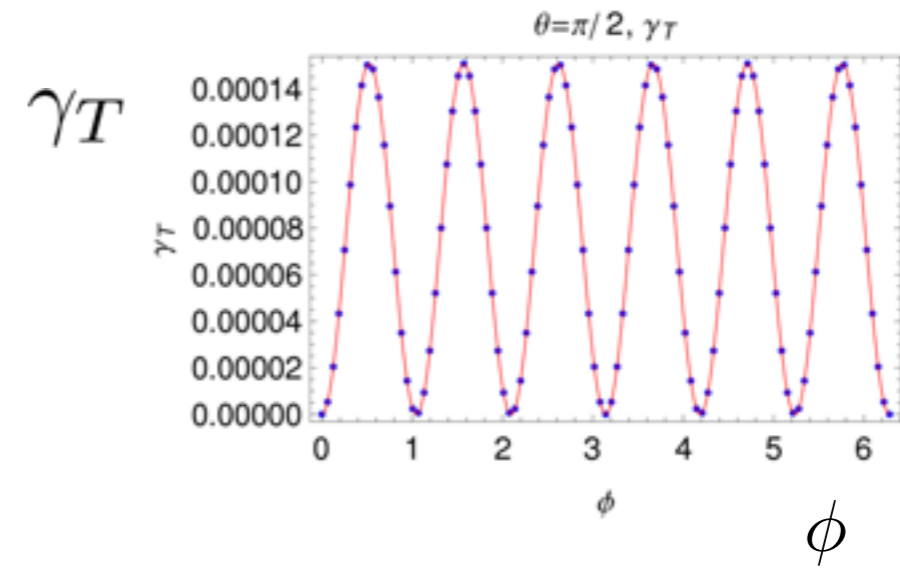
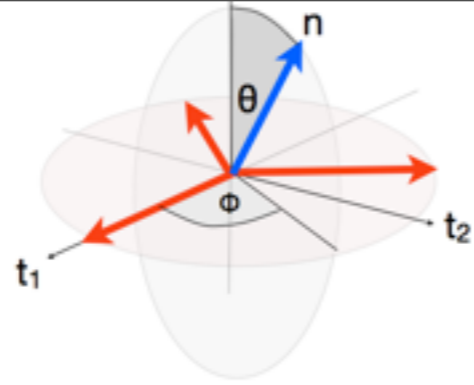
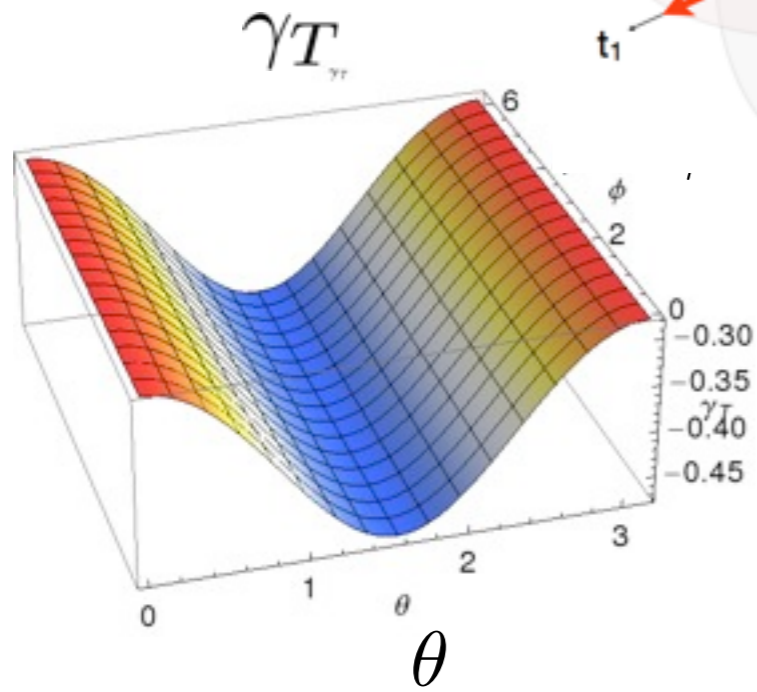


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Coplanarity



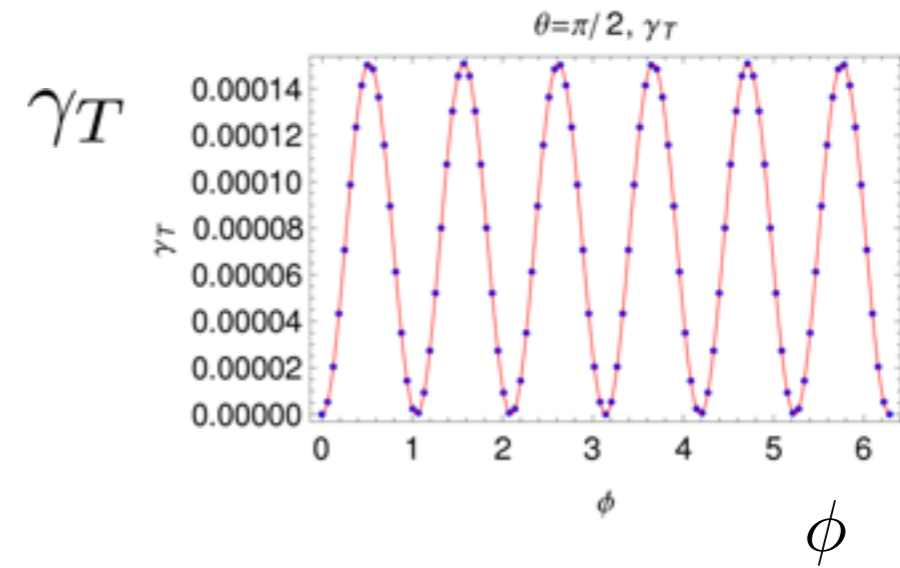
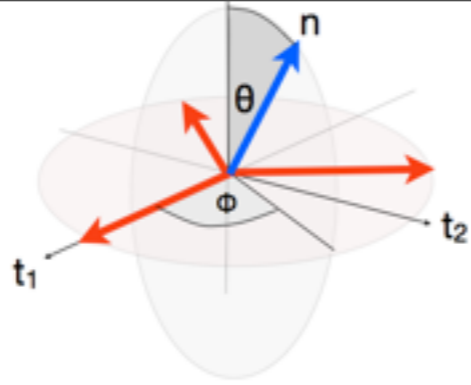
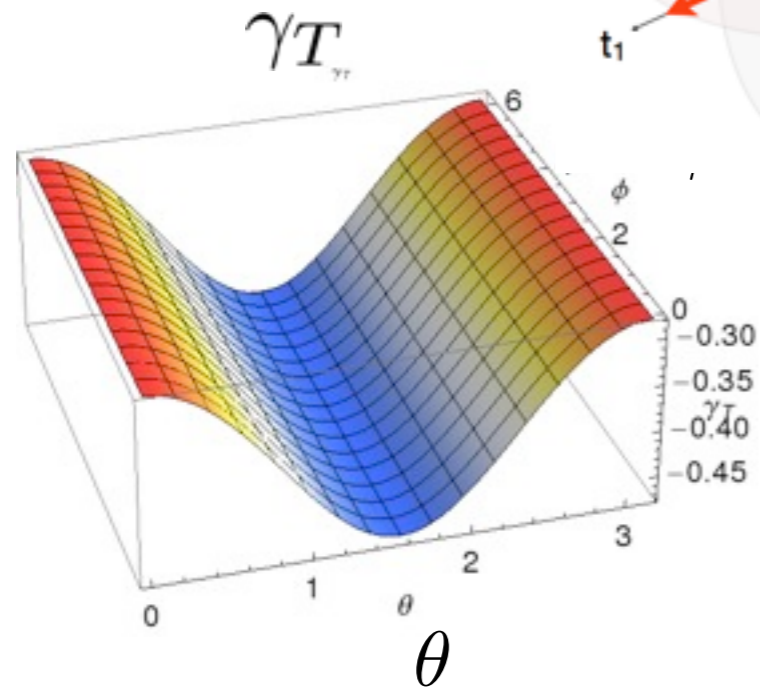
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→ Coplanar Spin Ordering

$$S = S_h + S_t + S_{th}$$

Coplanarity



Minimum occurs at $\theta = \frac{\pi}{2}$

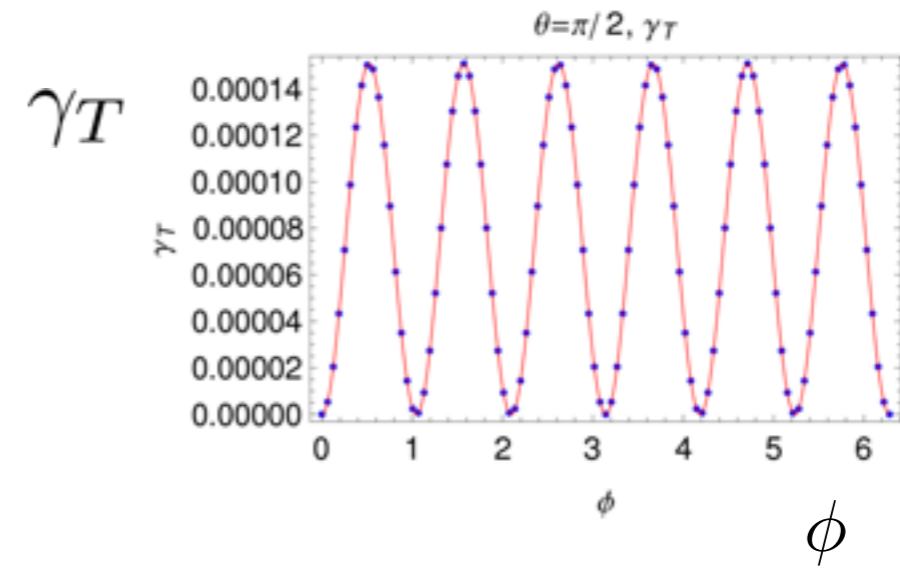
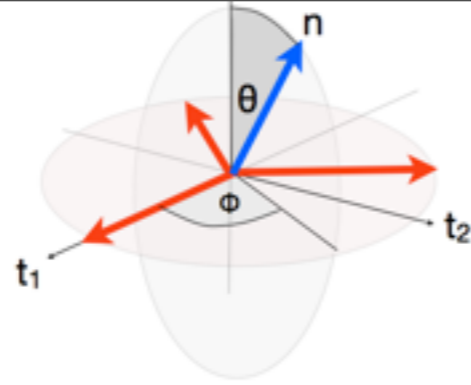
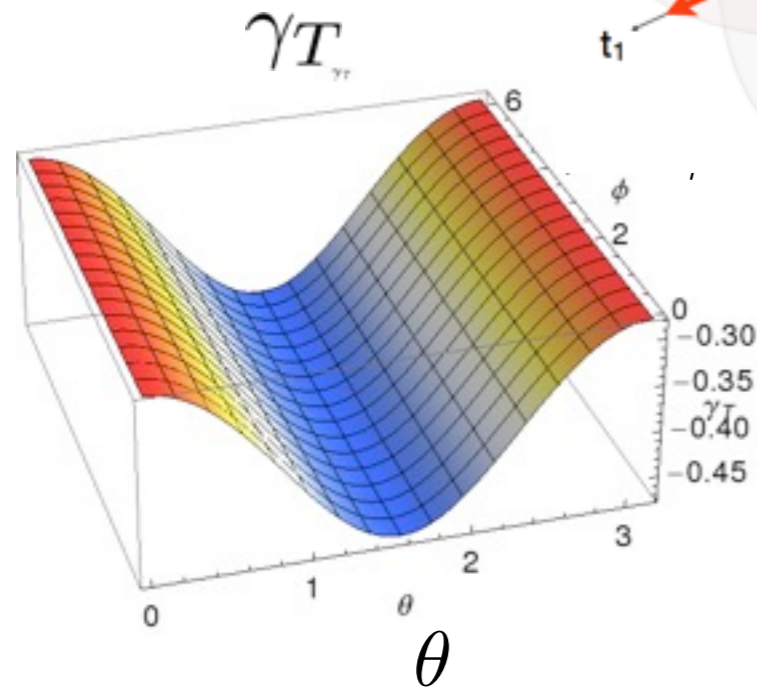
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Six Degenerate Minima at $\phi = \frac{2\pi n}{6}$

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$$S_h = \frac{K_h}{2} \int d^2x (\nabla n)^2$$

Coplanarity



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Six Degenerate Minima at $\phi = \frac{2\pi n}{6}$

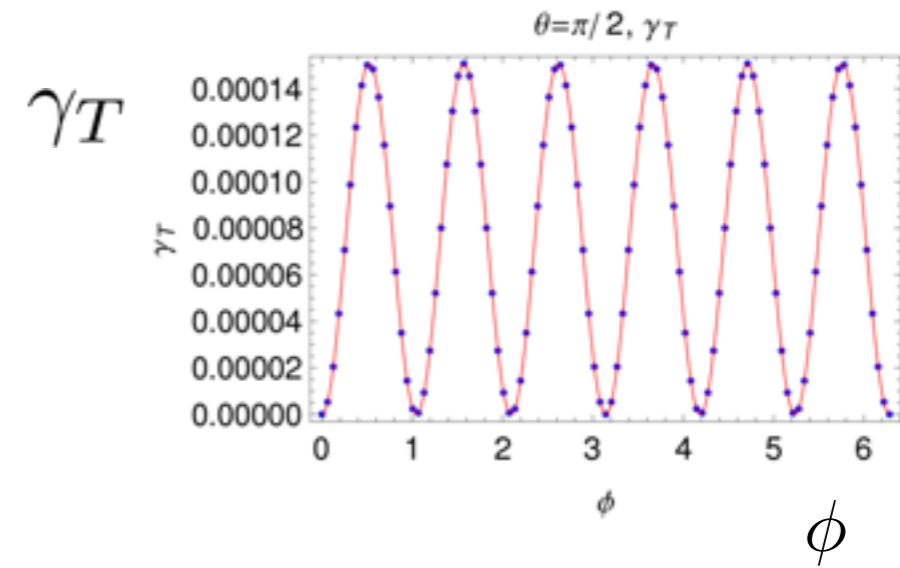
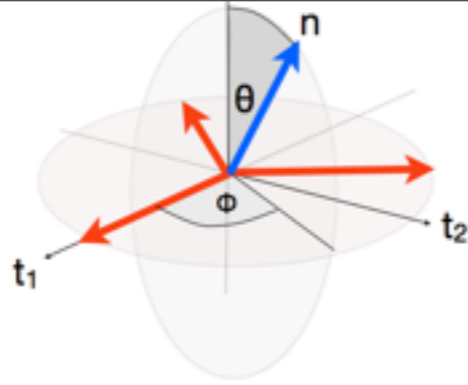
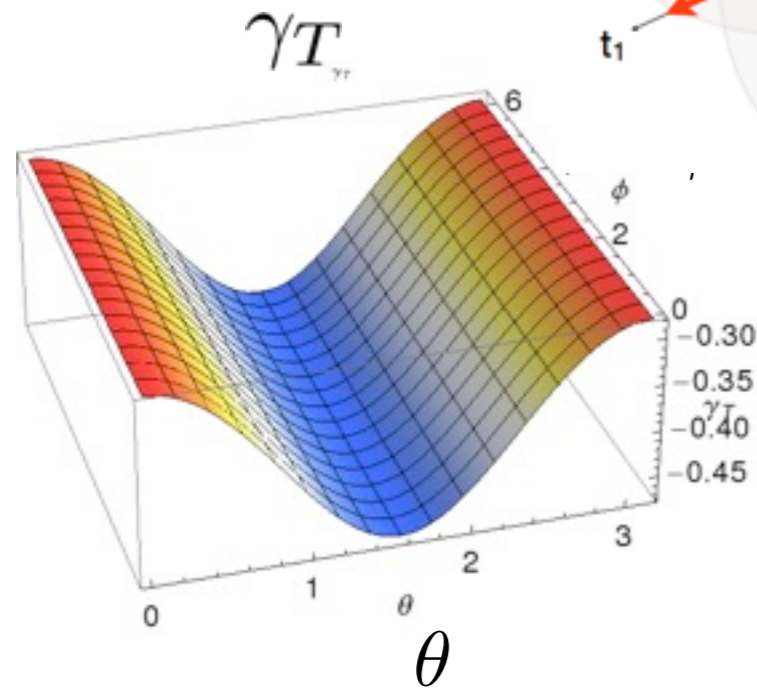
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$$S_{th} = \int d^2x (\alpha \cos^2 \theta + \lambda \cos 6\phi)$$

$$O(J_{th}^2)$$

$$O(J_{th}^6)$$

Both terms relevant
at high temperatures

The Coplanarity Cross-over Temperature

ξ = coherence length of coplanar fluxes

$$ae^{\frac{2\pi J_1}{T}} = a \left(\frac{J_1}{J_2} \right)$$

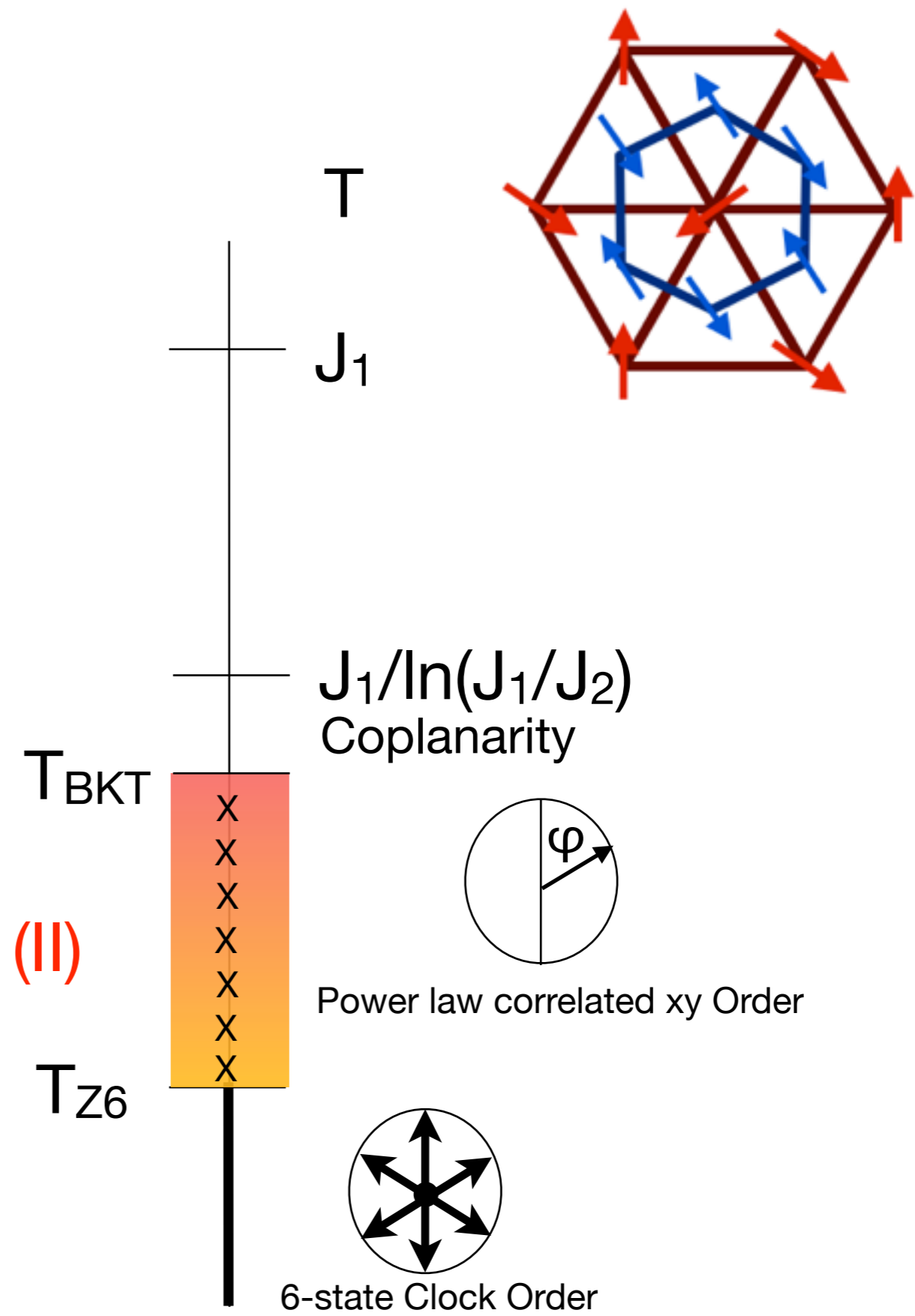
The Coplanarity Cross-over Temperature

ξ = coherence length of coplanar fluxes

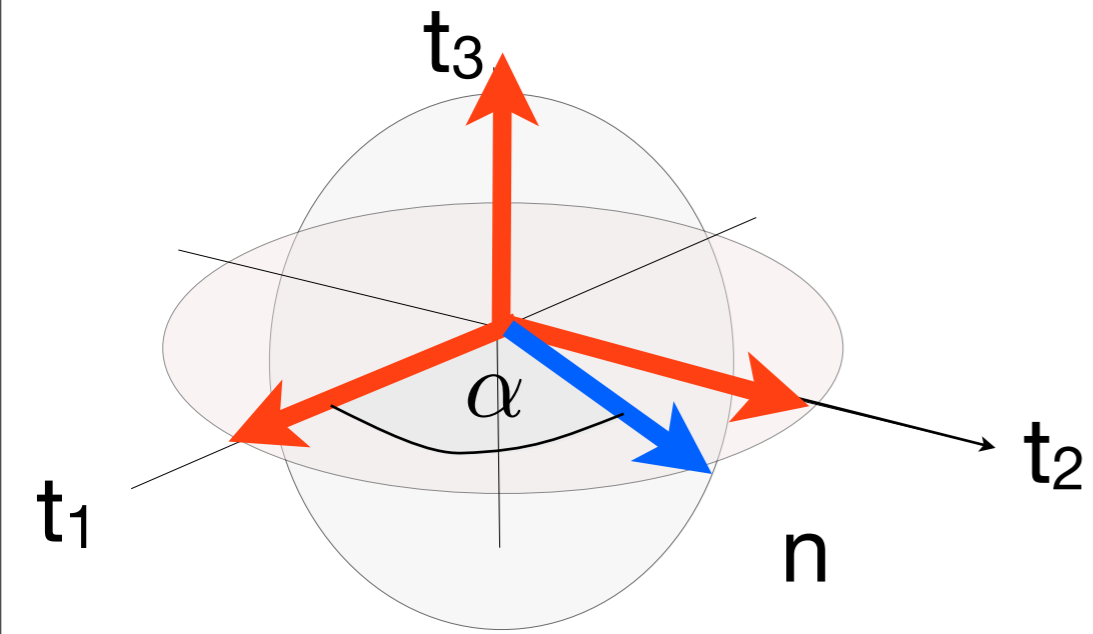
$$ae^{\frac{2\pi J_1}{T}} = a \left(\frac{J_1}{J_2} \right)$$



$$T_{\text{coplanarity}} \sim \frac{J_1}{\ln \left(\frac{J_1}{J_2} \right)}$$

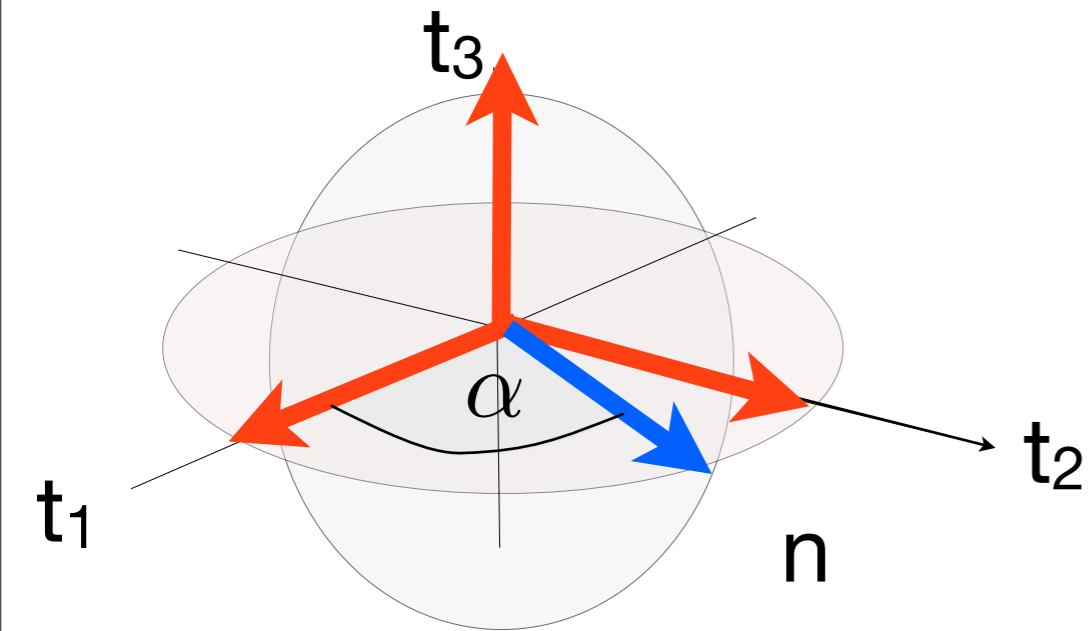


Scaling in the Coplanar State



Scaling in the Coplanar State

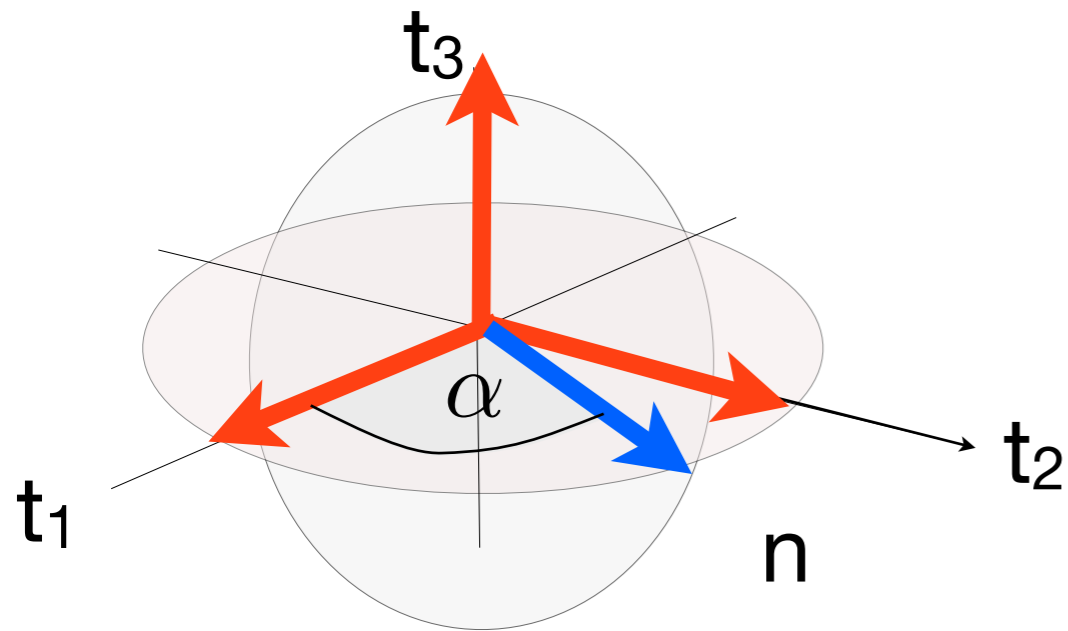
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Scaling in the Coplanar State

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$SO(3)$ "center of mass" motion of underlying spin fluid.

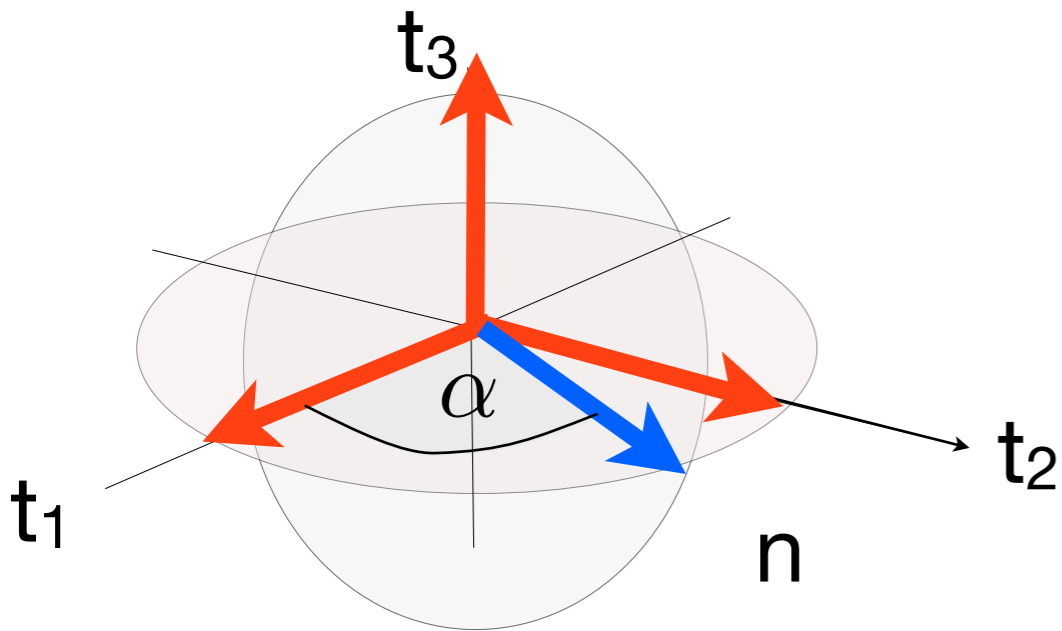


Scaling in the Coplanar State

Order parameter: $SO(3) \times U(1)$:

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$U(1)$: "relative phase" α between two sublattices.

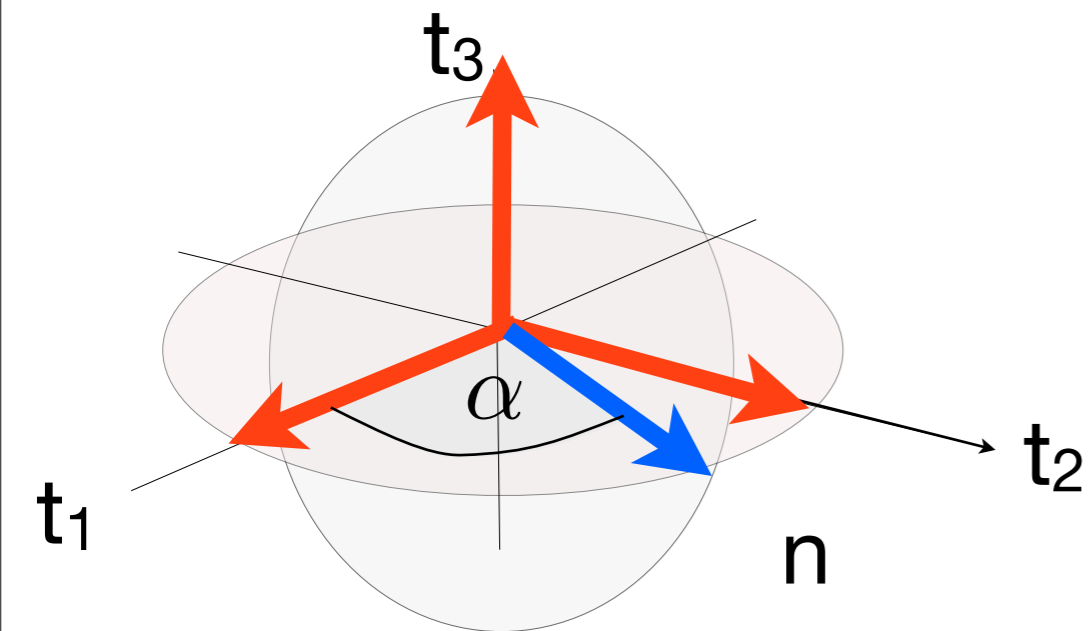


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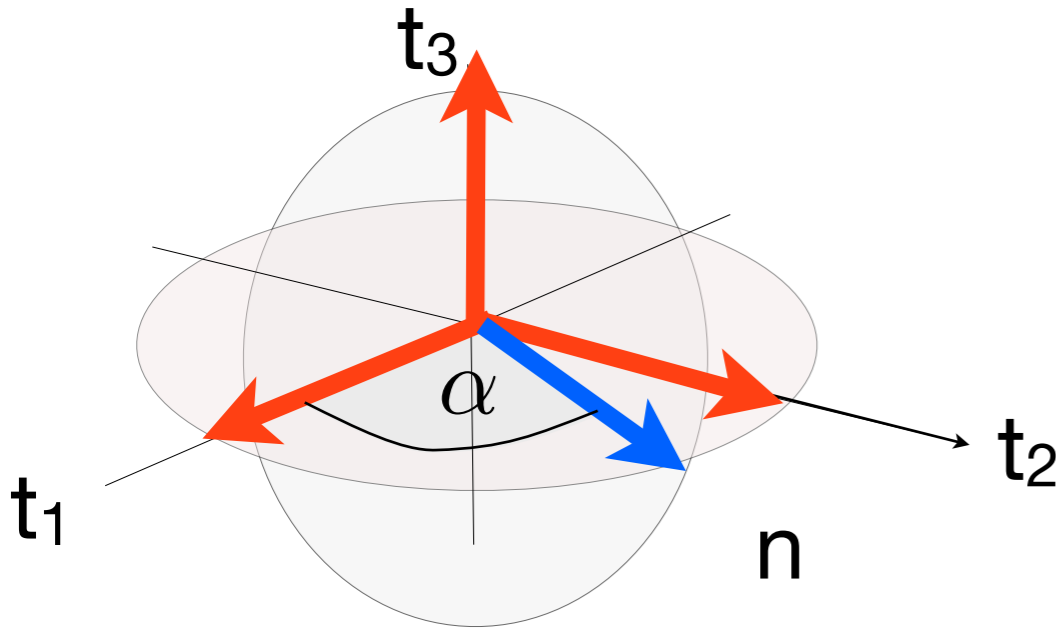
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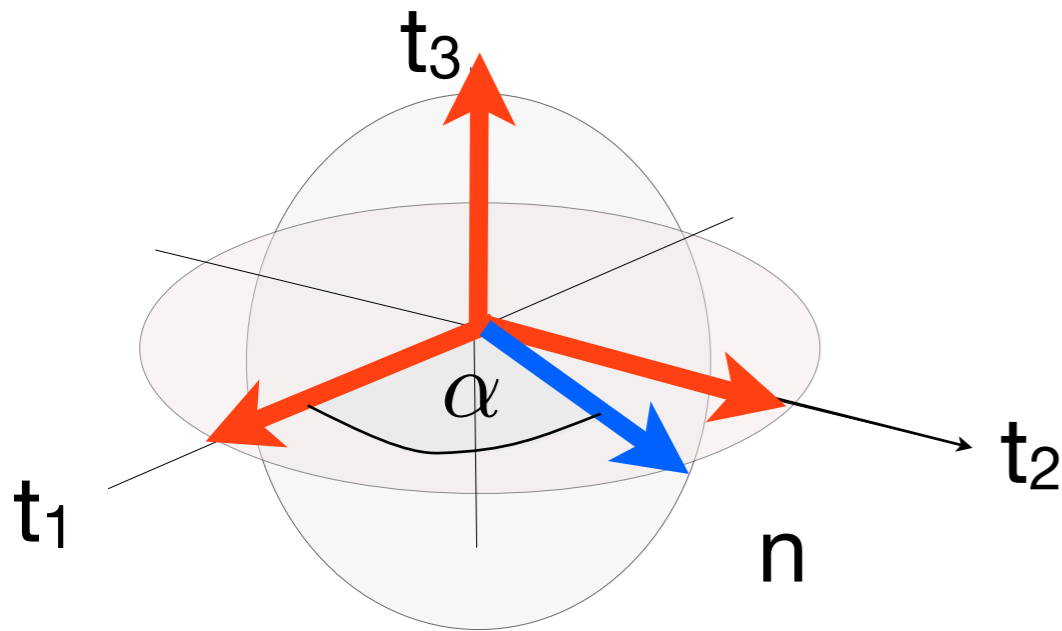
Scaling: confirms decoupling of $U(1)$ Relative Degree of Freedom to form a phase with topologically stable vortices. Binding of the vortices leads to a power law phase in which the 6-fold anisotropy is irrelevant.

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$$\psi \rightarrow \psi' = \psi + r\alpha$$

$$r = \kappa/2I_3$$

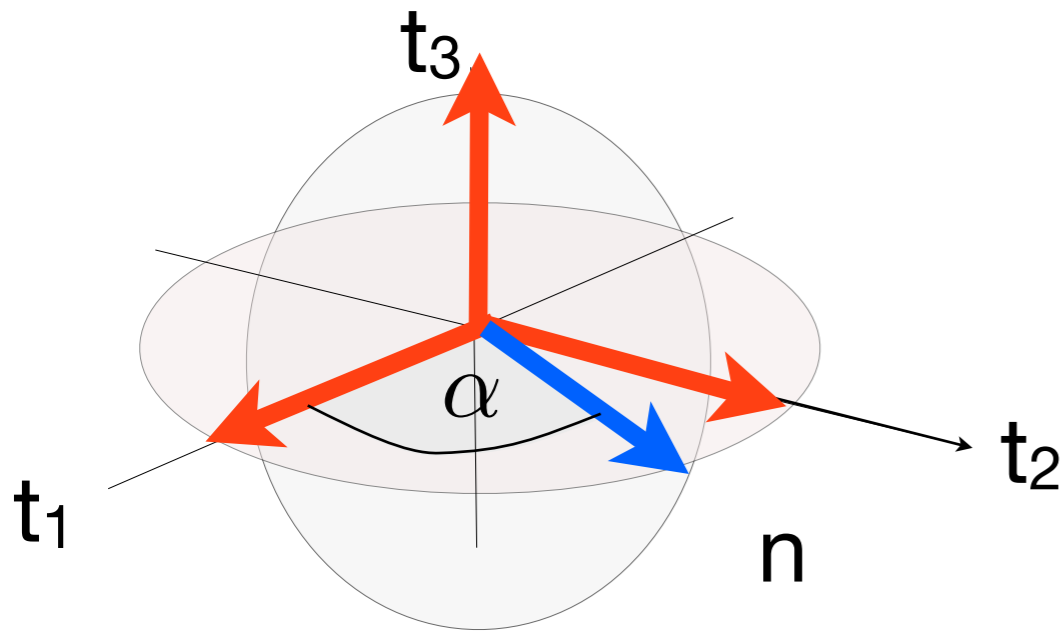
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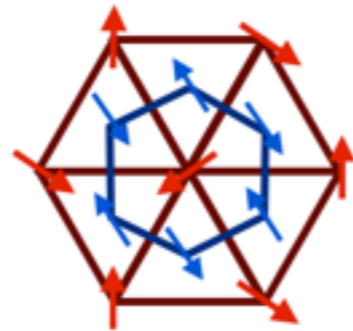
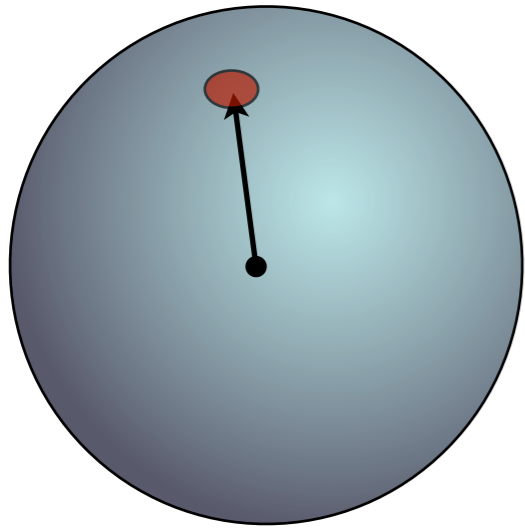
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Magnetism, Gravity and String Theory.

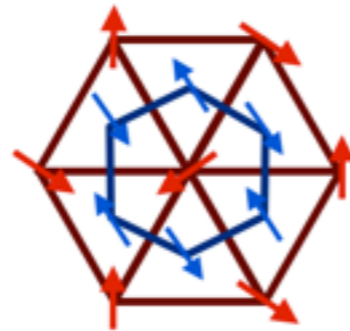
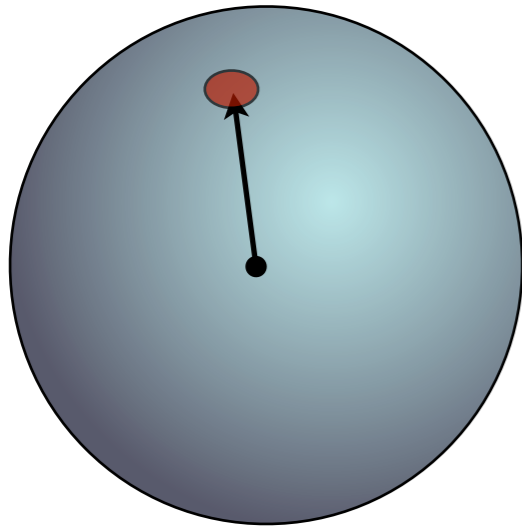
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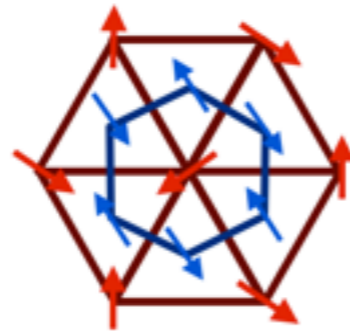
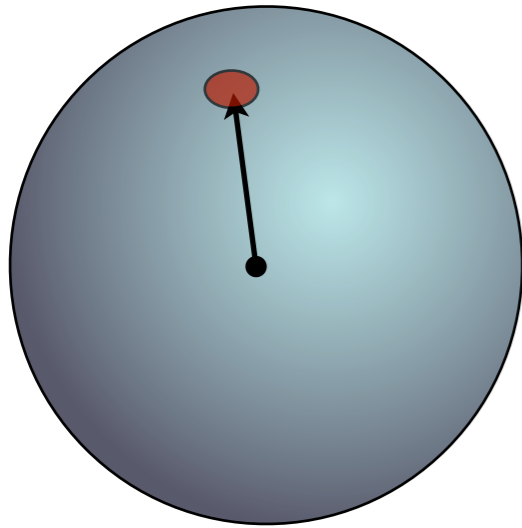
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Magnetization = 4D vector

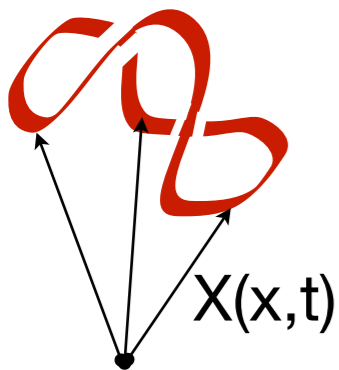
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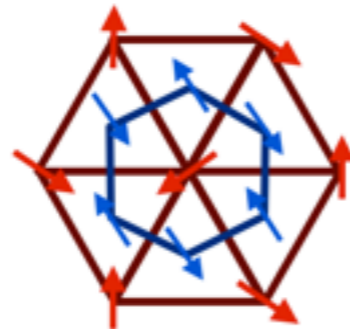
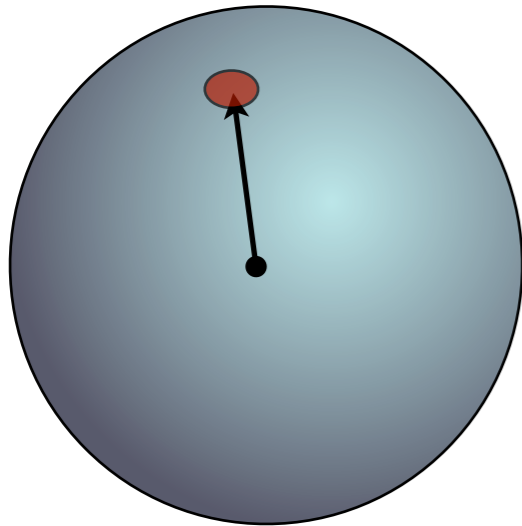


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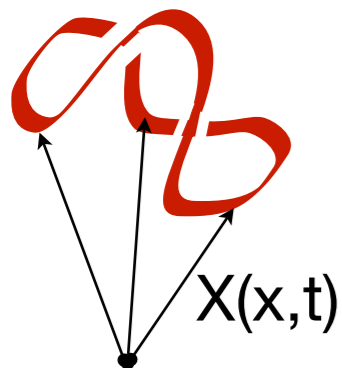
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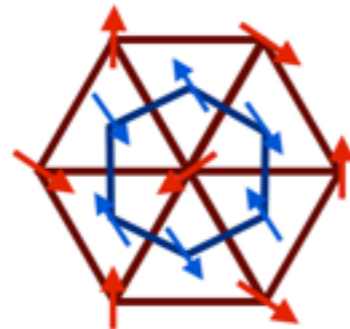
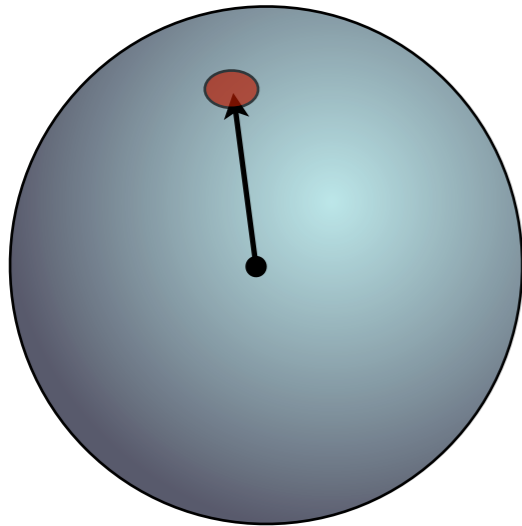
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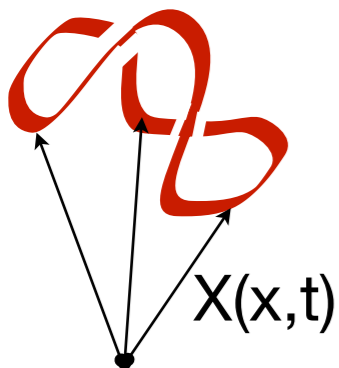
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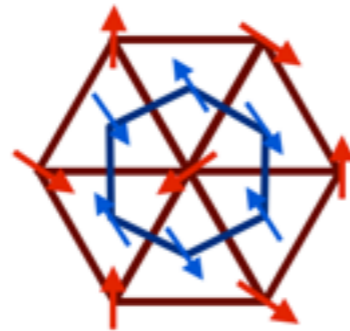
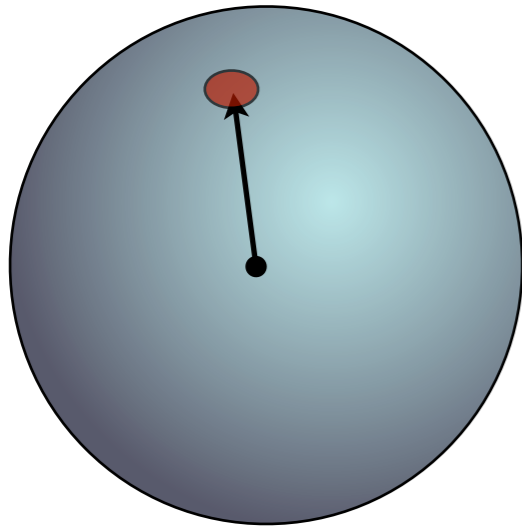


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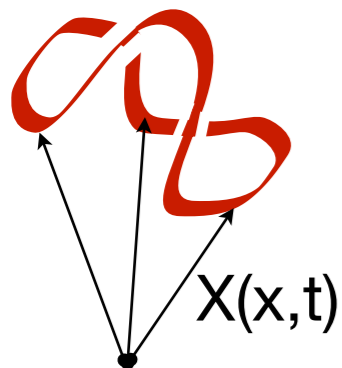
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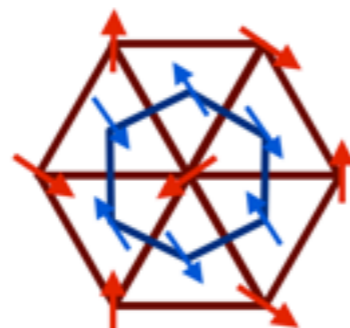
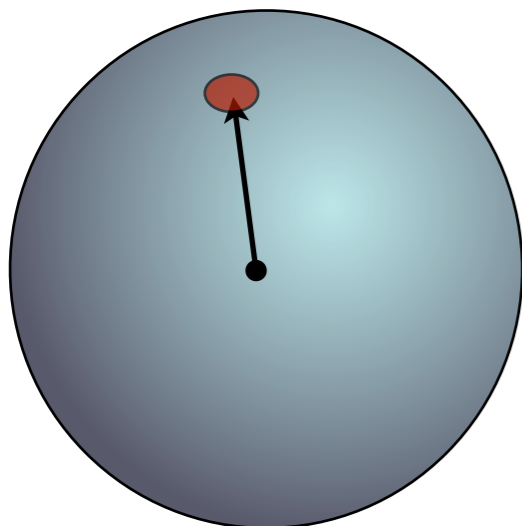
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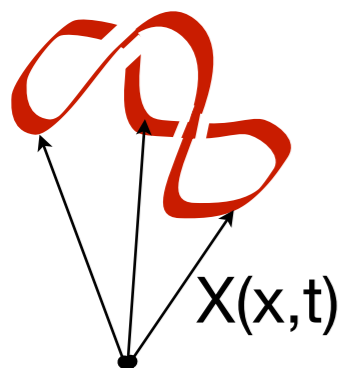
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Friedan '80, Hamilton '81, Perelmann '06

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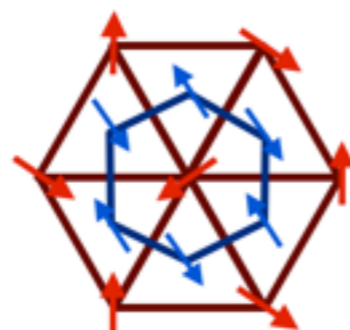
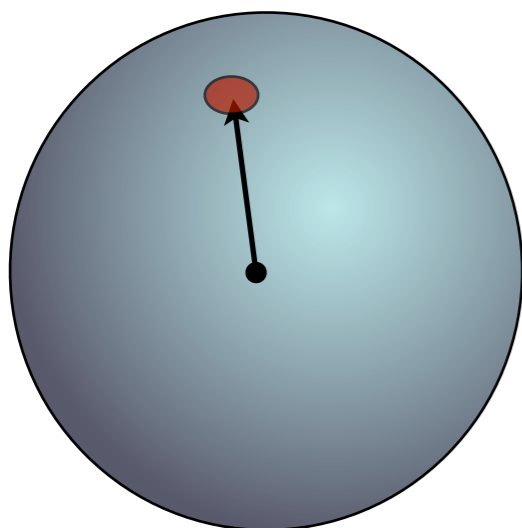
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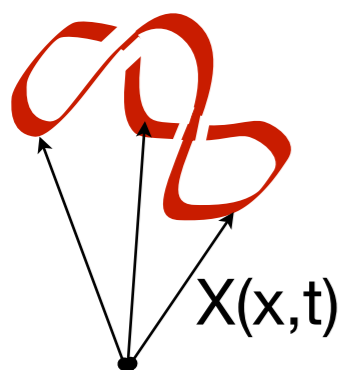
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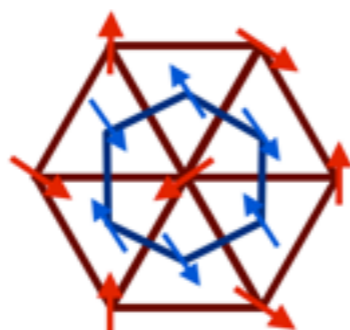
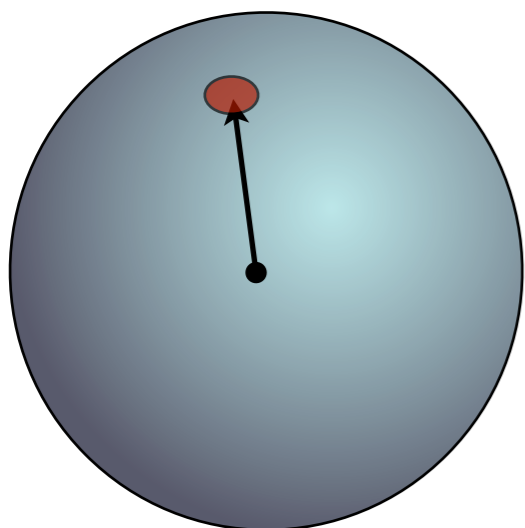
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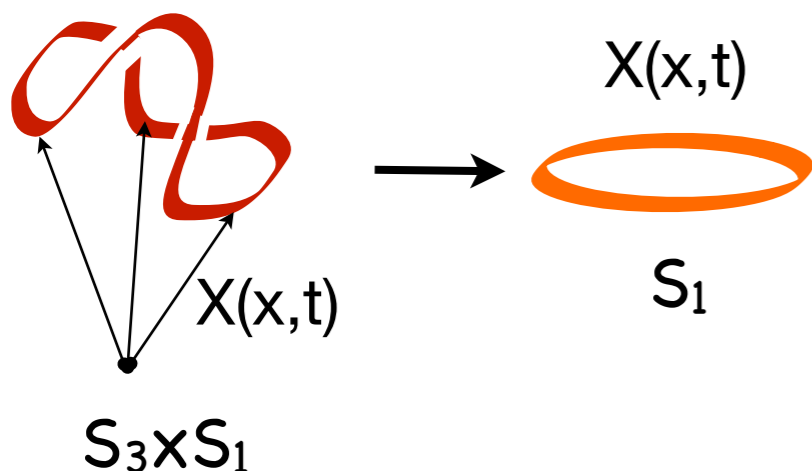
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The decoupling of the U(1) degrees of freedom from the SO(3) degrees of freedom is a kind of compactification from a four to a one dimensional universe.

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Small mathematica code to calculate the Ricci tensor

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Small mathematica code to calculate the Ricci tensor

Metric Tensor

$$g_{\mathbf{l}} = \begin{pmatrix} \sin^2(\theta) (I_1 \sin^2(\psi) + I_2 \cos^2(\psi)) + I_3 \cos^2(\theta) & \sin(\theta) (I_1 - I_2) \sin(\psi) \cos(\psi) & I_3 \cos(\theta) & \frac{1}{2} \kappa \cos(\theta) \\ \sin(\theta) (I_1 - I_2) \sin(\psi) \cos(\psi) & I_1 \cos^2(\psi) + I_2 \sin^2(\psi) & 0 & 0 \\ I_3 \cos(\theta) & 0 & I_3 & \frac{\kappa}{2} \\ \frac{1}{2} \kappa \cos(\theta) & 0 & \frac{\kappa}{2} & I_\varphi \end{pmatrix};$$

Cristoffel Symbol

```

For[i = 1, i ≤ 4, i++,
  For[k = 1, k ≤ 4, k++,
    For[l = 1, l ≤ 4, l++,
      Γ[[i, k, l]] =
        Sum[1/2 * g_u[[i, j]] *
          (D[g_l[[j, k]], x[[l]] + D[g_l[[j, l]], x[[k]]] -
            D[g_l[[k, l]], x[[j]]))]]]]

```

Small mathematica code to calculate the Ricci tensor

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$$g_l = \begin{pmatrix} \sin^2(\theta) (I1 \sin^2(\psi) + I2 \cos^2(\psi)) + I3 \cos^2(\theta) & \sin(\theta) (I1 - I2) \sin(\psi) \cos(\psi) & I3 \cos(\theta) & \frac{1}{2} \kappa \cos(\theta) \\ \sin(\theta) (I1 - I2) \sin(\psi) \cos(\psi) & I1 \cos^2(\psi) + I2 \sin^2(\psi) & 0 & 0 \\ I3 \cos(\theta) & 0 & I3 & \frac{\kappa}{2} \\ \frac{1}{2} \kappa \cos(\theta) & 0 & \frac{\kappa}{2} & I\varphi \end{pmatrix};$$

$$(\Gamma^i)_{kl} = \frac{1}{2} g^{ij} (\nabla_l g_{jk} + \nabla_k g_{jl} - \nabla_j g_{kl})$$

Ricci Tensor

For $i = 1, i \leq 4, i++,$

For $k = 1, k \leq 4, k++,$

Riccill[[i, k]] =

$$\sum_{l=1}^4 \left(D[\Gamma[[1, i, k]], x[[1]]] - D[\Gamma[[1, i, 1]], x[[k]]] + \sum_{m=1}^4 (\Gamma[[m, 1, m]] * \Gamma[[1, i, k]] - \Gamma[[m, i, 1]] * \Gamma[[1, k, m]]) \right);$$

$$R^k_{ijl} = \Gamma^k_{i1,j} - \Gamma^k_{ij,1} + \Gamma^k_{jn} \Gamma^n_{li} - \Gamma^k_{ln} \Gamma^n_{ij}$$

$$R_{ij} = R^k_{ikj}$$

$$\frac{dg_{ij}}{dl} = -\frac{1}{2\pi} R_{ij}$$

Small mathematica code to calculate the Ricci tensor

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$$g_l = \begin{pmatrix} \sin^2(\theta) (I1 \sin^2(\psi) + I2 \cos^2(\psi)) + I3 \cos^2(\theta) & \sin(\theta) (I1 - I2) \sin(\psi) \cos(\psi) & I3 \cos(\theta) & \frac{1}{2} \kappa \cos(\theta) \\ \sin(\theta) (I1 - I2) \sin(\psi) \cos(\psi) & I1 \cos^2(\psi) + I2 \sin^2(\psi) & 0 & 0 \\ I3 \cos(\theta) & 0 & I3 & \frac{\kappa}{2} \\ \frac{1}{2} \kappa \cos(\theta) & 0 & \frac{\kappa}{2} & I\varphi \end{pmatrix};$$

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$$R_{ij} = R^k_{ikj}$$

$$\frac{dg_{ij}}{dl} = -\frac{1}{2\pi} R_{ij}$$

Small mathematica code to calculate the Ricci tensor

This is the renormalization of I3

$$-\text{FullSimplify} \left[\frac{1}{2\pi} \text{Riccill}[[3, 3]] \right]$$

$$-\frac{I3^2 - (I1 - I2)^2}{4\pi I1 I2}$$

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$$g_l = \begin{pmatrix} \sin^2(\theta) (I1 \sin^2(\psi) + I2 \cos^2(\psi)) + I3 \cos^2(\theta) & \sin(\theta) (I1 - I2) \sin(\psi) \cos(\psi) & I3 \cos(\theta) & \frac{1}{2} \kappa \cos(\theta) \\ \sin(\theta) (I1 - I2) \sin(\psi) \cos(\psi) & I1 \cos^2(\psi) + I2 \sin^2(\psi) & 0 & 0 \\ I3 \cos(\theta) & 0 & I3 & \frac{\kappa}{2} \\ \frac{1}{2} \kappa \cos(\theta) & 0 & \frac{\kappa}{2} & I\varphi \end{pmatrix};$$

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For $i = 1, i \leq 4, i++,$

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`Riccil1[[i, k]] =`

`$\sum_{l=1}^4 \left(D[\Gamma[[1, i, k]], x[[l]]] - D[\Gamma[[1, i, l]], x[[k]]] +$`

`$\sum_{m=1}^4 (\Gamma[[m, l, m]] * \Gamma[[1, i, k]] - \Gamma[[m, i, l]] * \Gamma[[1, k, m]]) \right) \right]$` ;

$$\frac{dg^{ab}}{dl} = -\frac{1}{2\pi} R^{ab}$$

Small mathematica code to calculate the Ricci tensor

This is the renormalization of I3

`-FullSimplify $\left[\frac{1}{2\pi} \text{Riccil1}[[3, 3]] \right]$`

$$-\frac{I3^2 - (I1 - I2)^2}{4\pi I1 I2}$$

Ricci Tensor

For $i = 1, i \leq 4, i++,$

For $k = 1, k \leq 4, k++,$

`Riccill[[i, k]] =`

$$\sum_{l=1}^4 \left(\mathbf{D}[\Gamma[[1, i, k]], \mathbf{x}[[1]]] - \mathbf{D}[\Gamma[[1, i, 1]], \mathbf{x}[[k]]] + \sum_{m=1}^4 (\Gamma[[m, 1, m]] * \Gamma[[1, i, k]] - \Gamma[[m, i, 1]] * \Gamma[[1, k, m]]) \right)]] ;$$

$$\frac{dg^{ab}}{dl} = -\frac{1}{2\pi} R^{ab}$$

Small mathematica code to calculate the Ricci tensor

This is the renormalization of I3

`-FullSimplify [$\frac{1}{2\pi}$ Riccill[[3, 3]]]`

$$-\frac{I_3^2 - (I_1 - I_2)^2}{4\pi I_1 I_2}$$

$$\frac{dI_1}{dl} = \frac{(I_2 - I_3)^2 - I_1^2}{4\pi I_2 I_3} - \frac{(I_1^2 - I_2^2) \kappa^2}{16\pi I_2 I_3^2 \left(I_\varphi - \frac{\kappa^2}{4I_3} \right)}$$

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I₁→I₂

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As isotropy develops, r stops renormalizing, U(1) phase decouples with

finite stiffness

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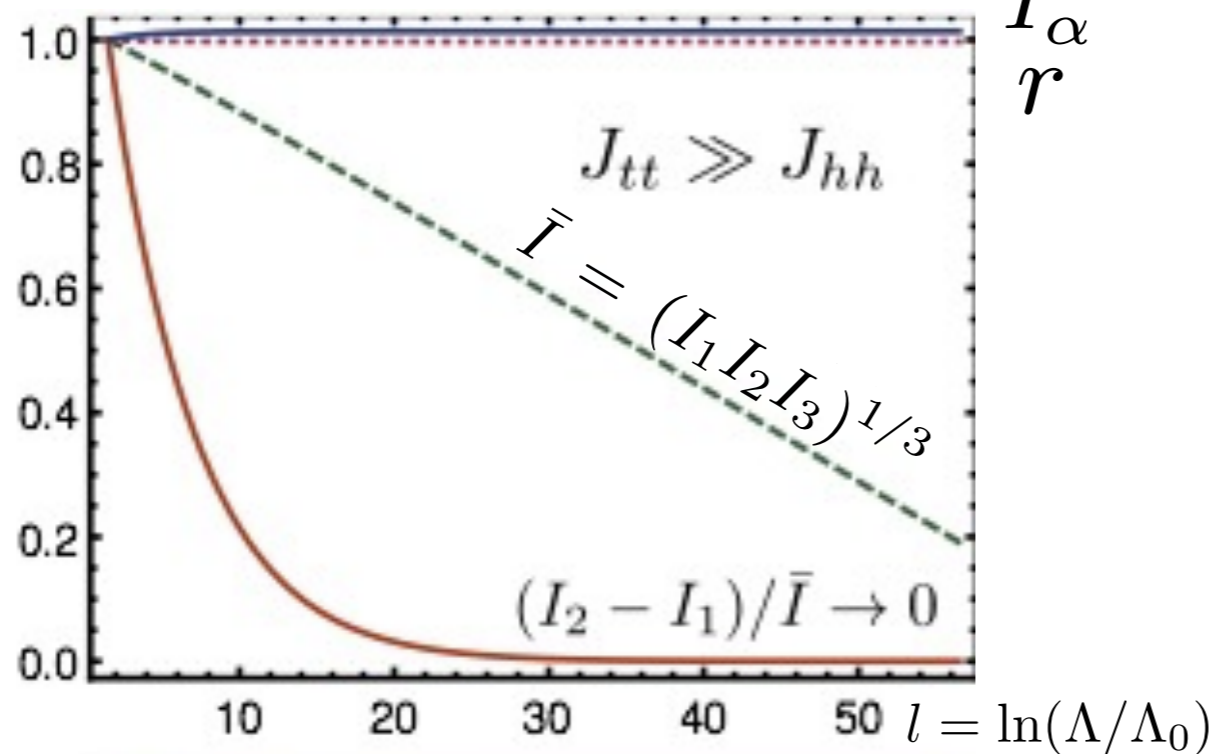
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$$r = \kappa/2I_3$$

$$I_\alpha \rightarrow I'_\alpha = I_\alpha - \kappa^2/4I_3$$

Decoupling becomes complete as $I_1 \rightarrow I_2$

$$\bar{I} = (I_1 I_2 I_3)^{1/3} \quad \text{---} \quad (I_2 - I_1)/\bar{I} \quad \text{---} \quad r$$



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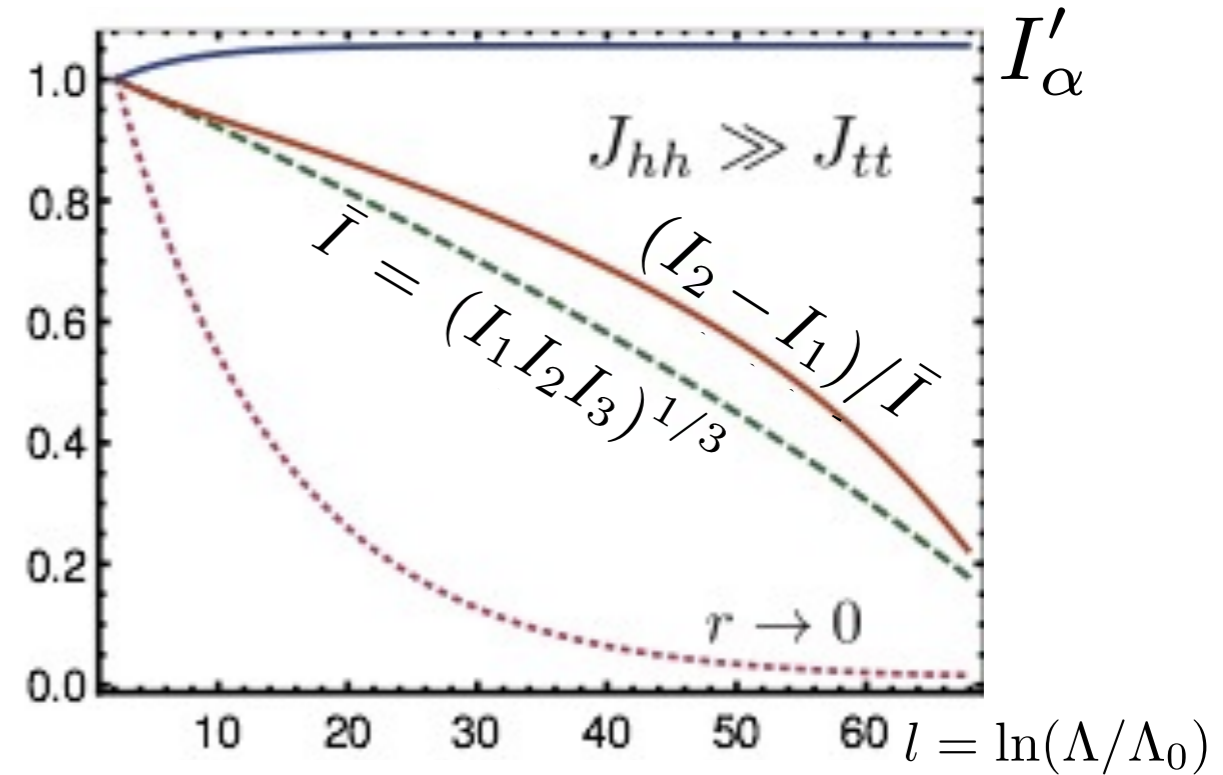
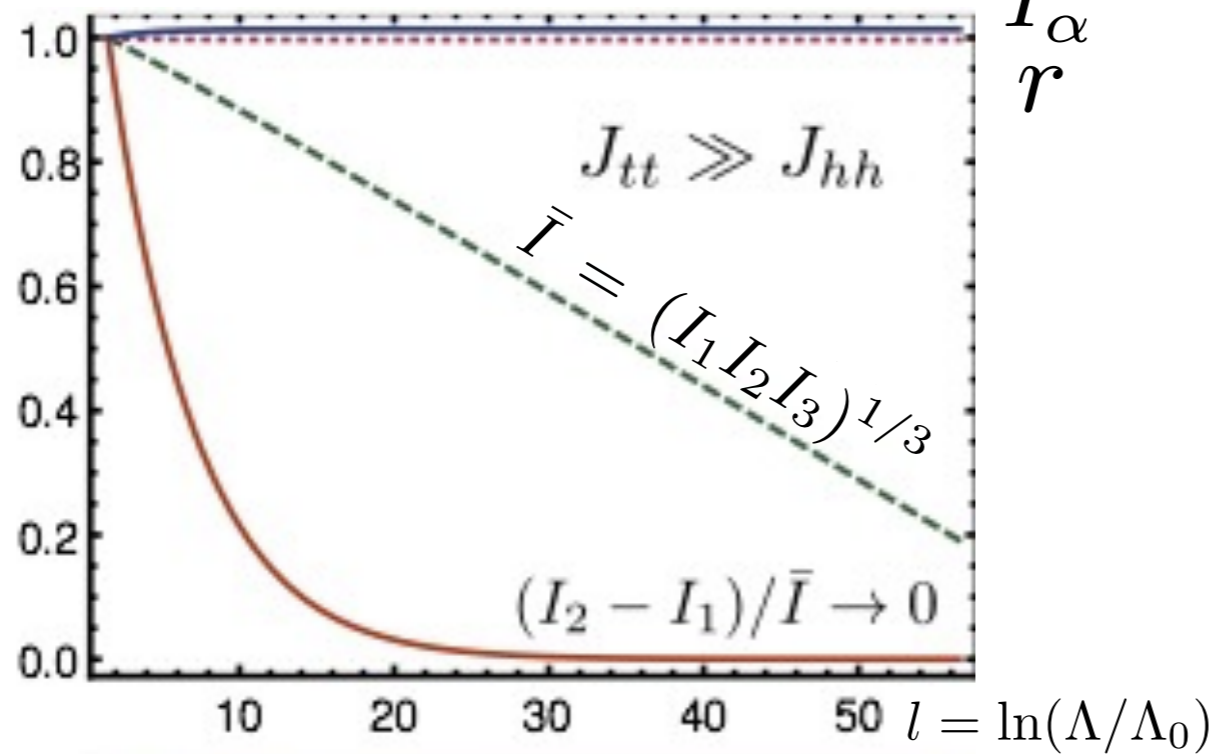
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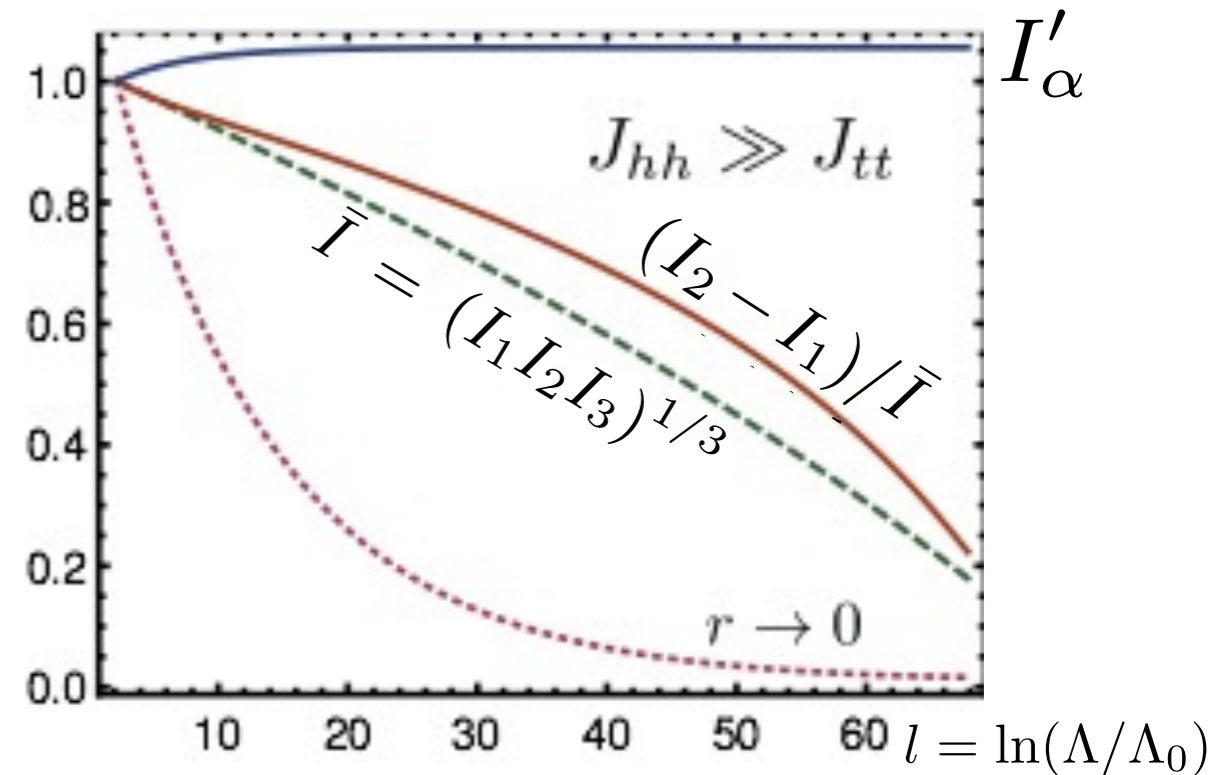
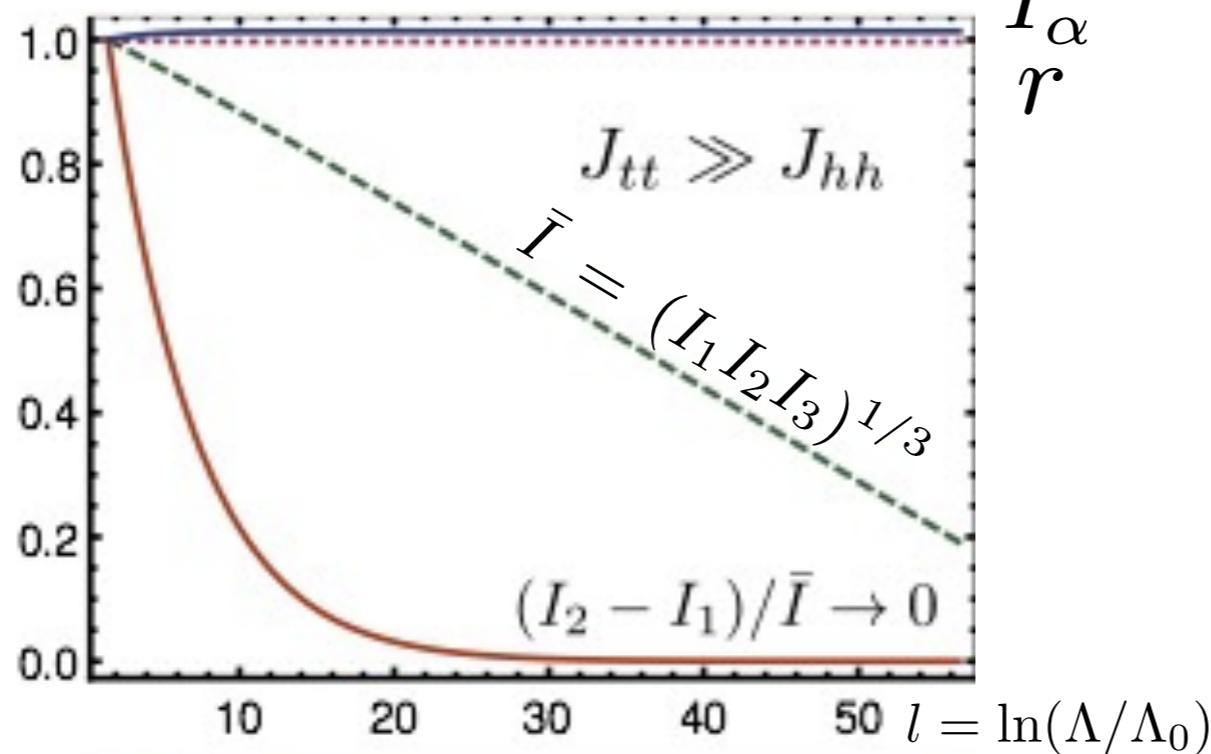
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Decoupling of U(1) Degrees of Freedom in Both Parameter Regimes

$$S_{\mathbb{Z}_6} = \frac{1}{2} \int d^2x \left[(I'_\alpha (\partial_\mu \alpha))^2 + \lambda \cos(6\alpha) \right].$$

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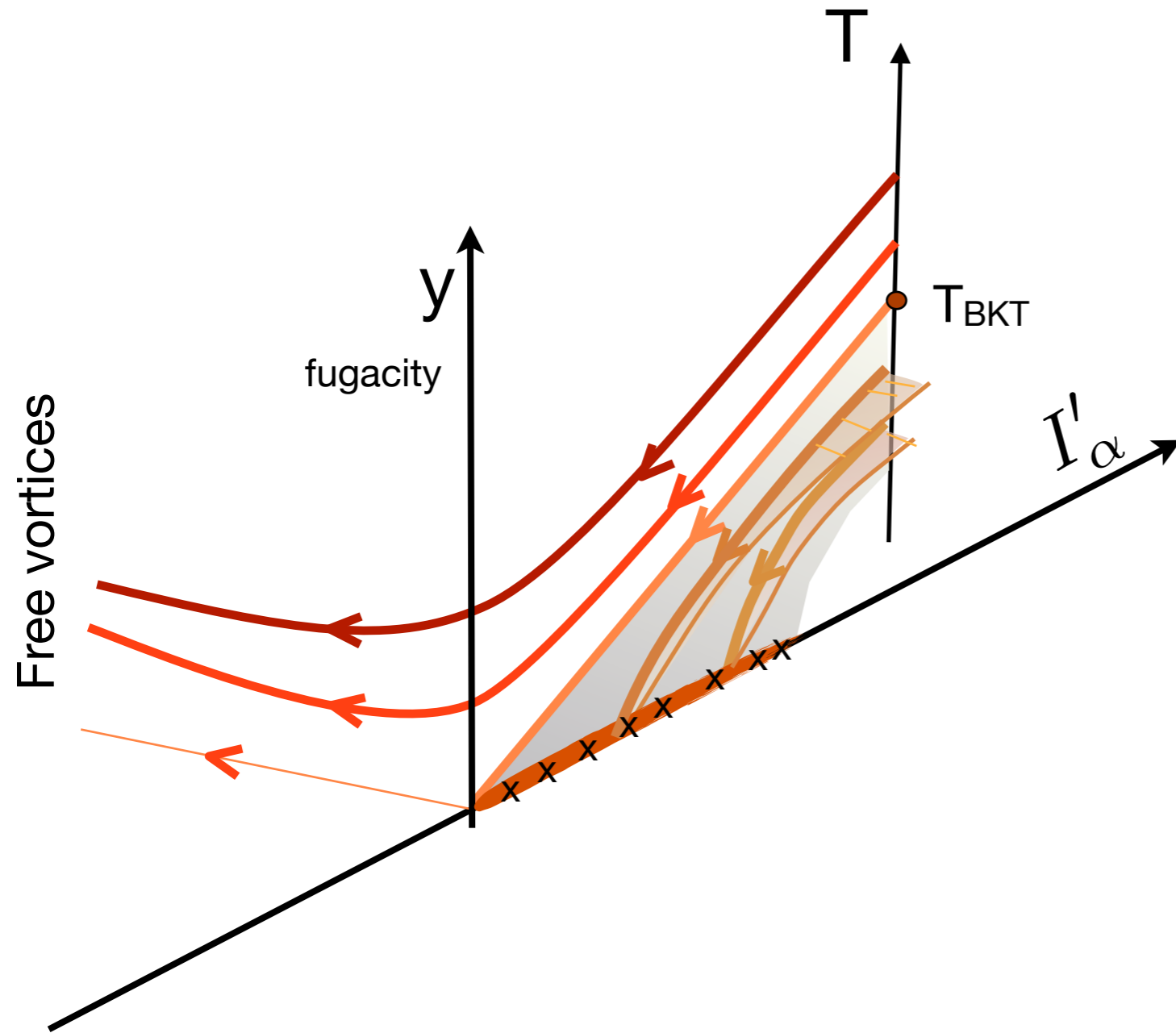
$$\frac{d \ln(\lambda)}{dl} = 2 - \frac{n^2/2}{2\pi I'_\alpha} = 2 - \frac{n^2}{8}$$

$n > 4$, @ T_{BKT}
Anisotropy Irrelevant
(Jose et al, 1977)

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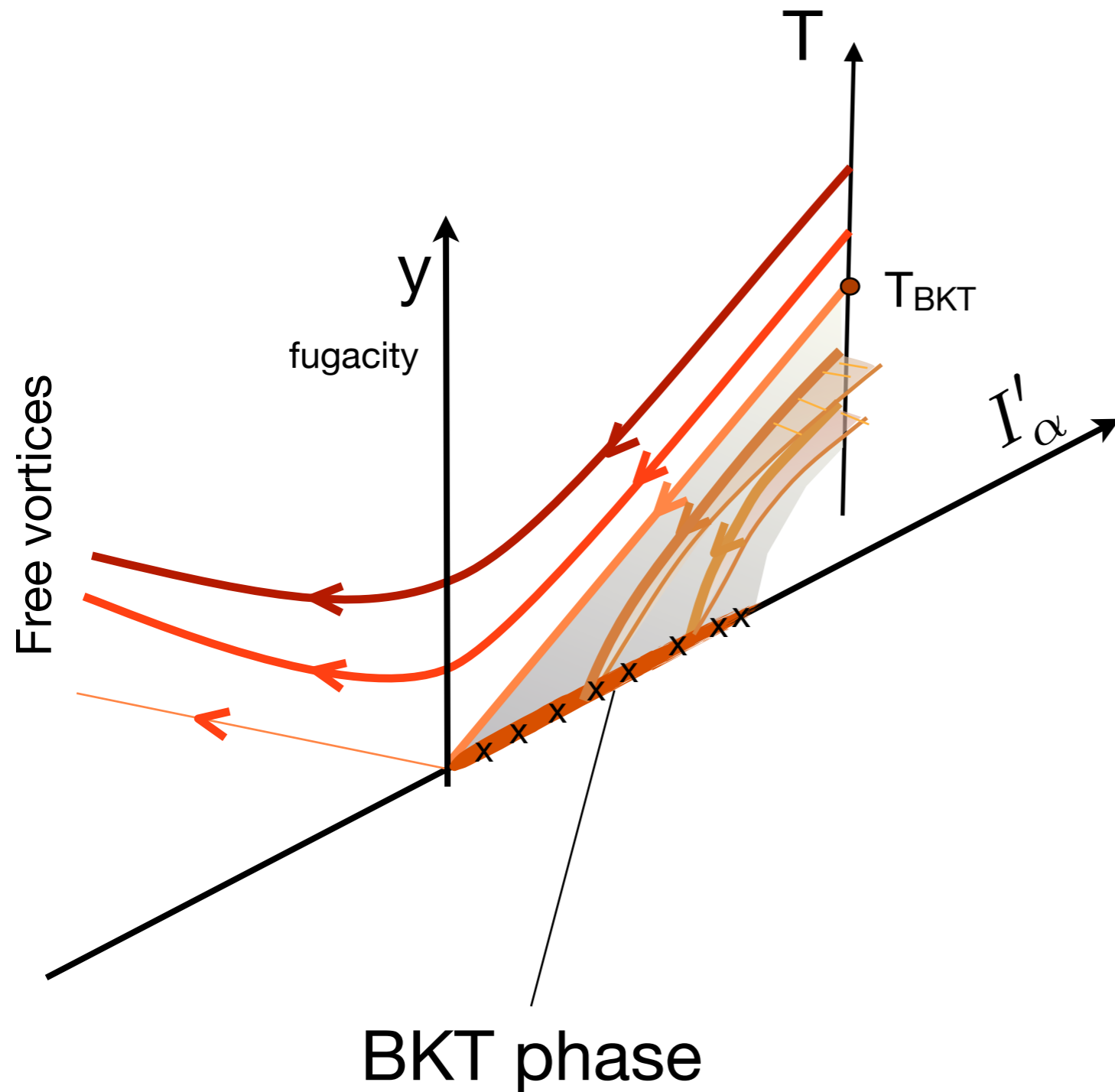
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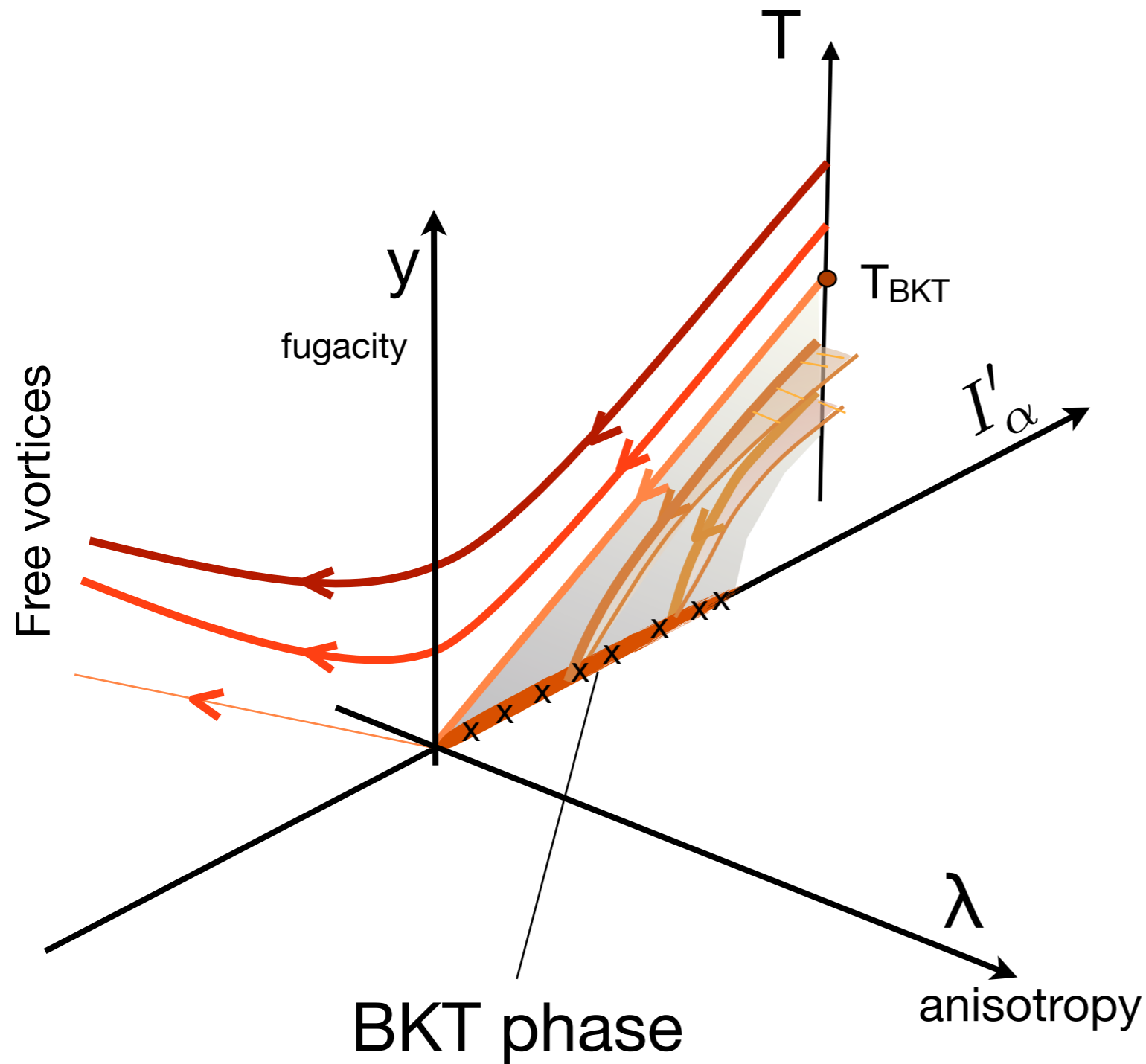
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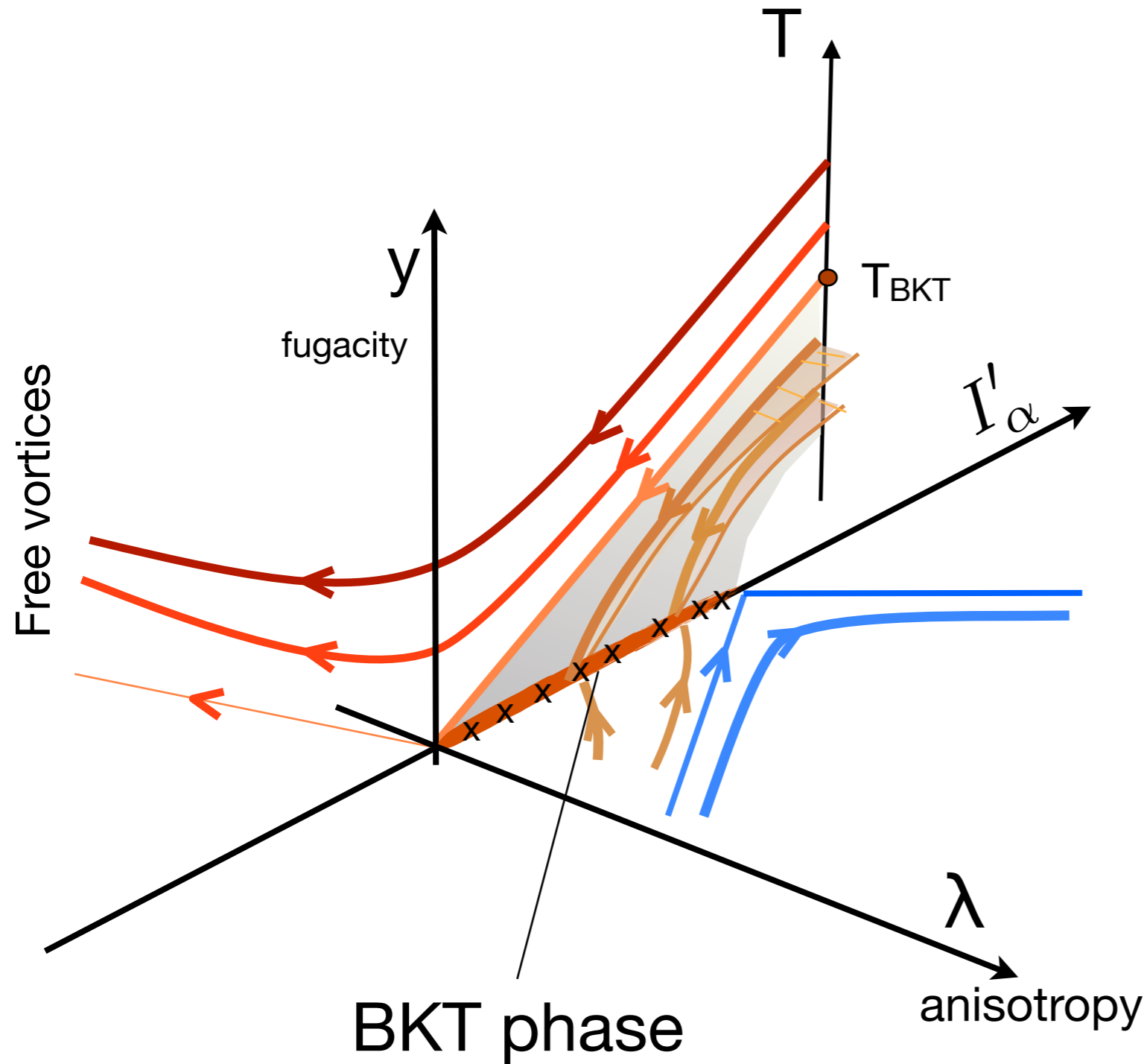
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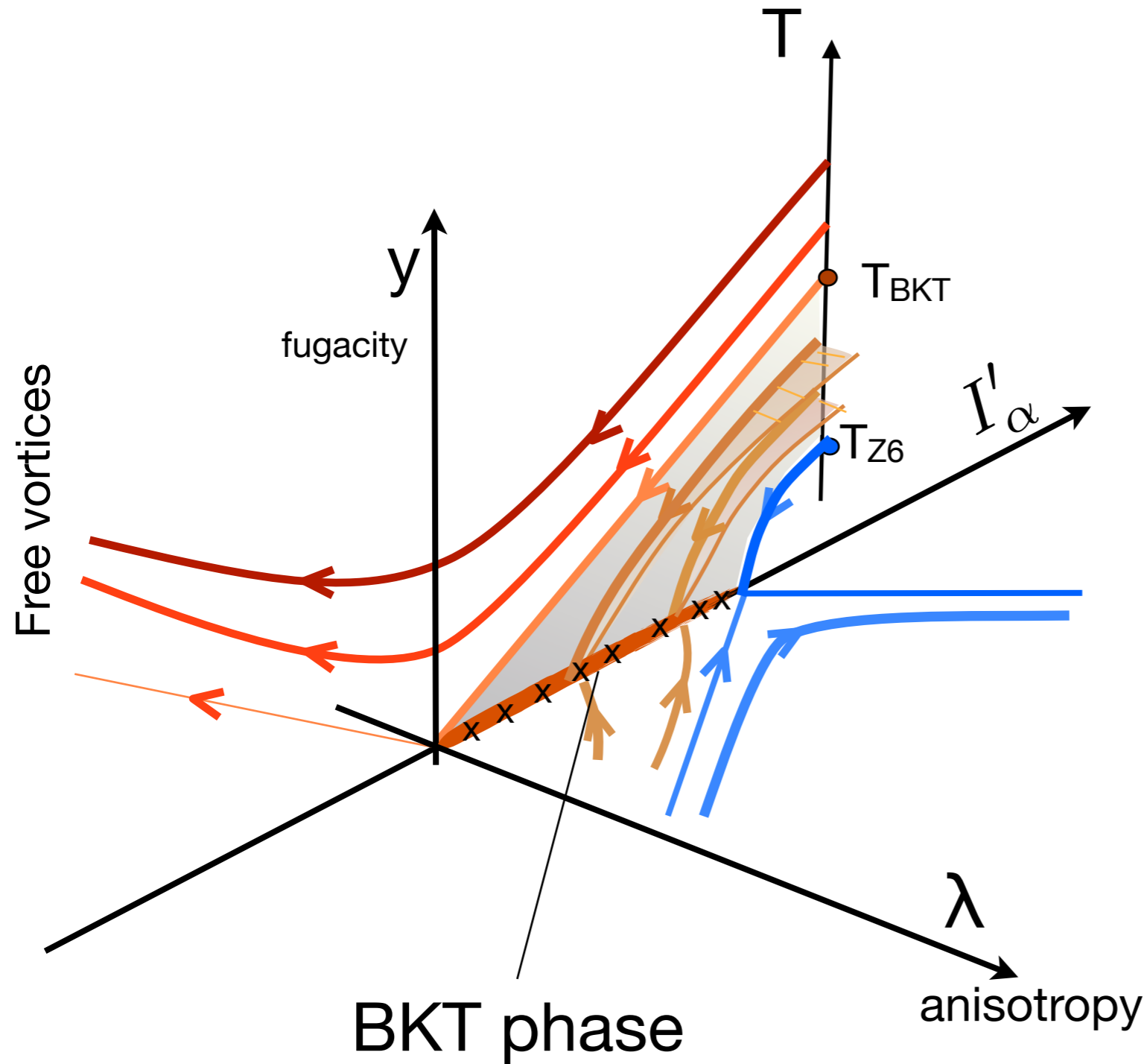
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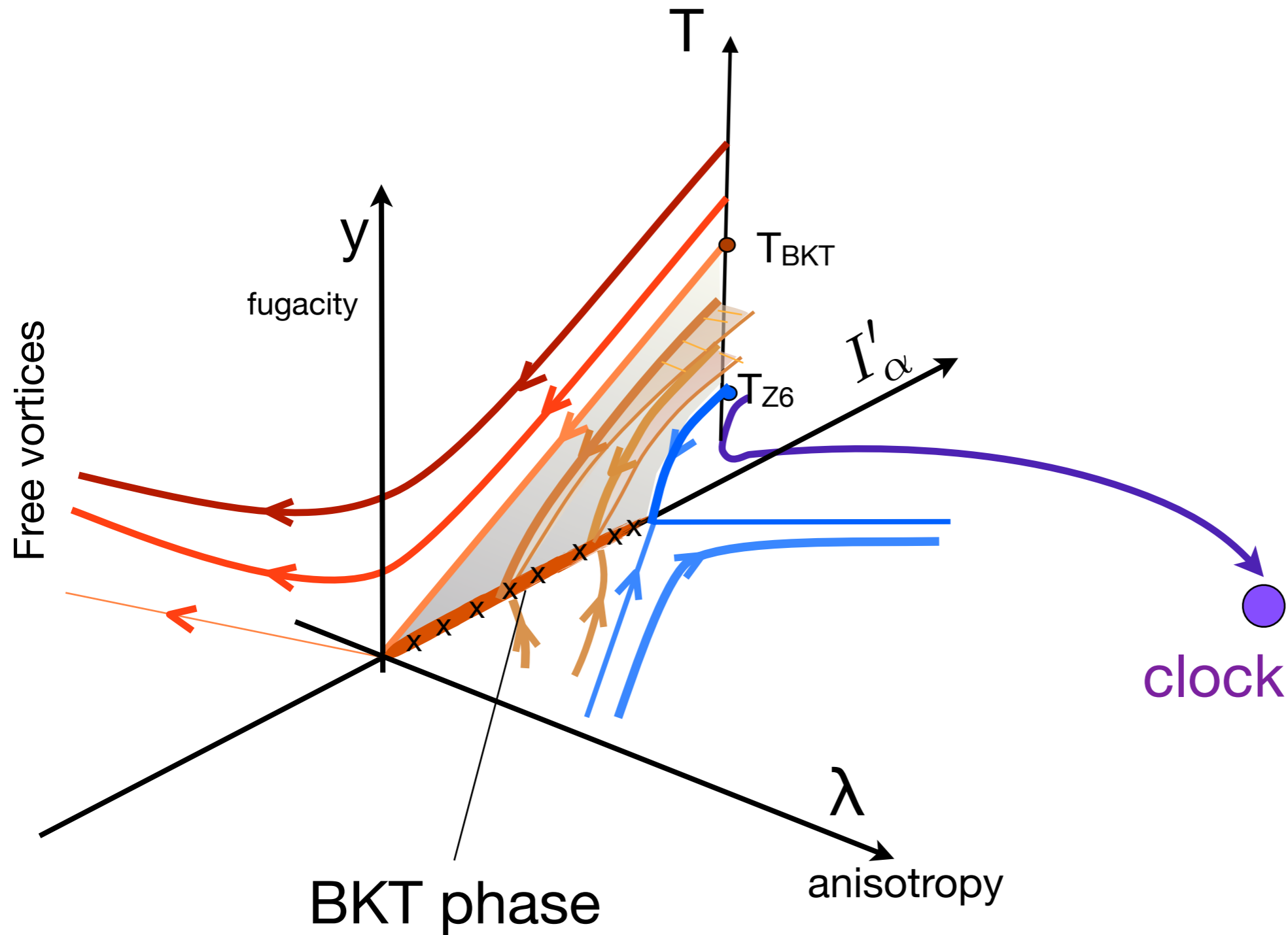
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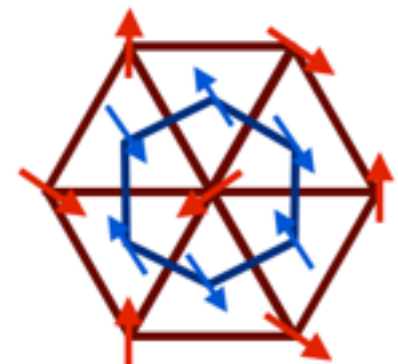
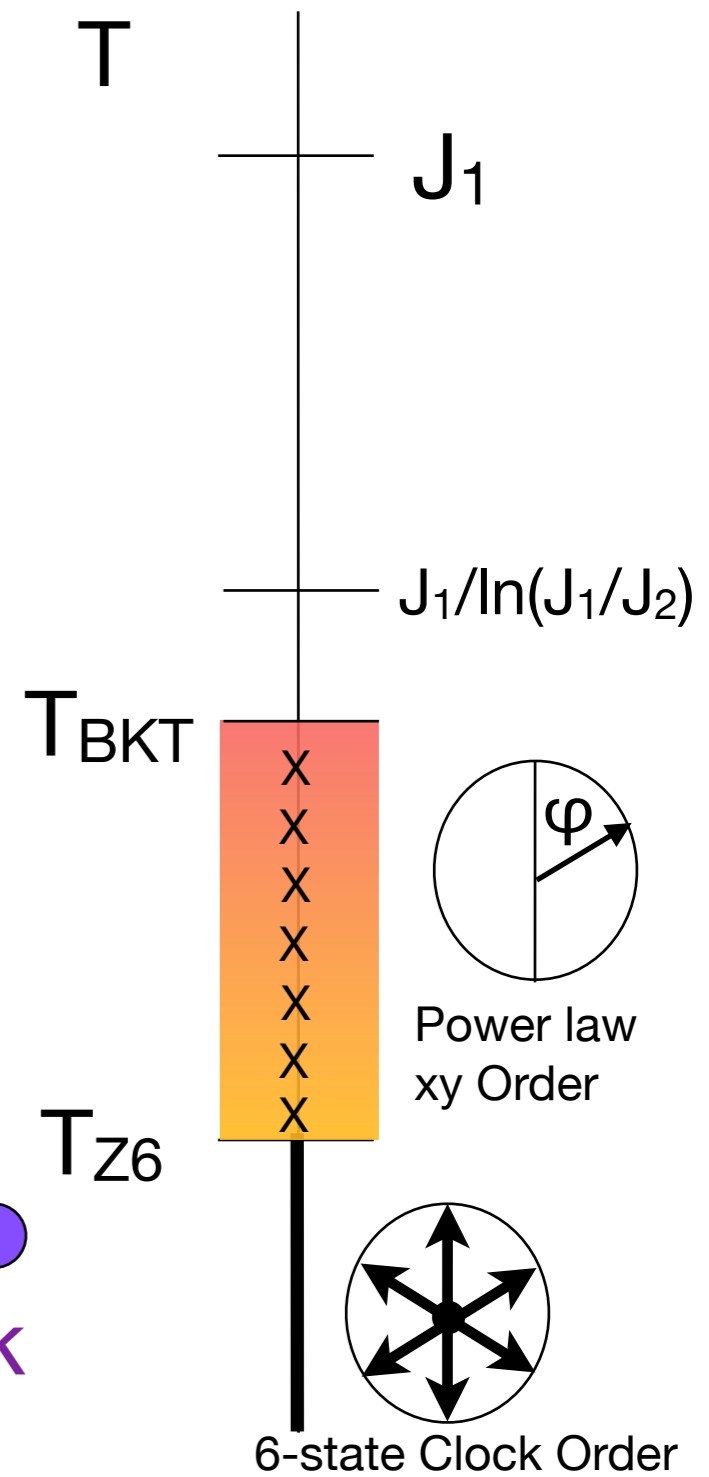
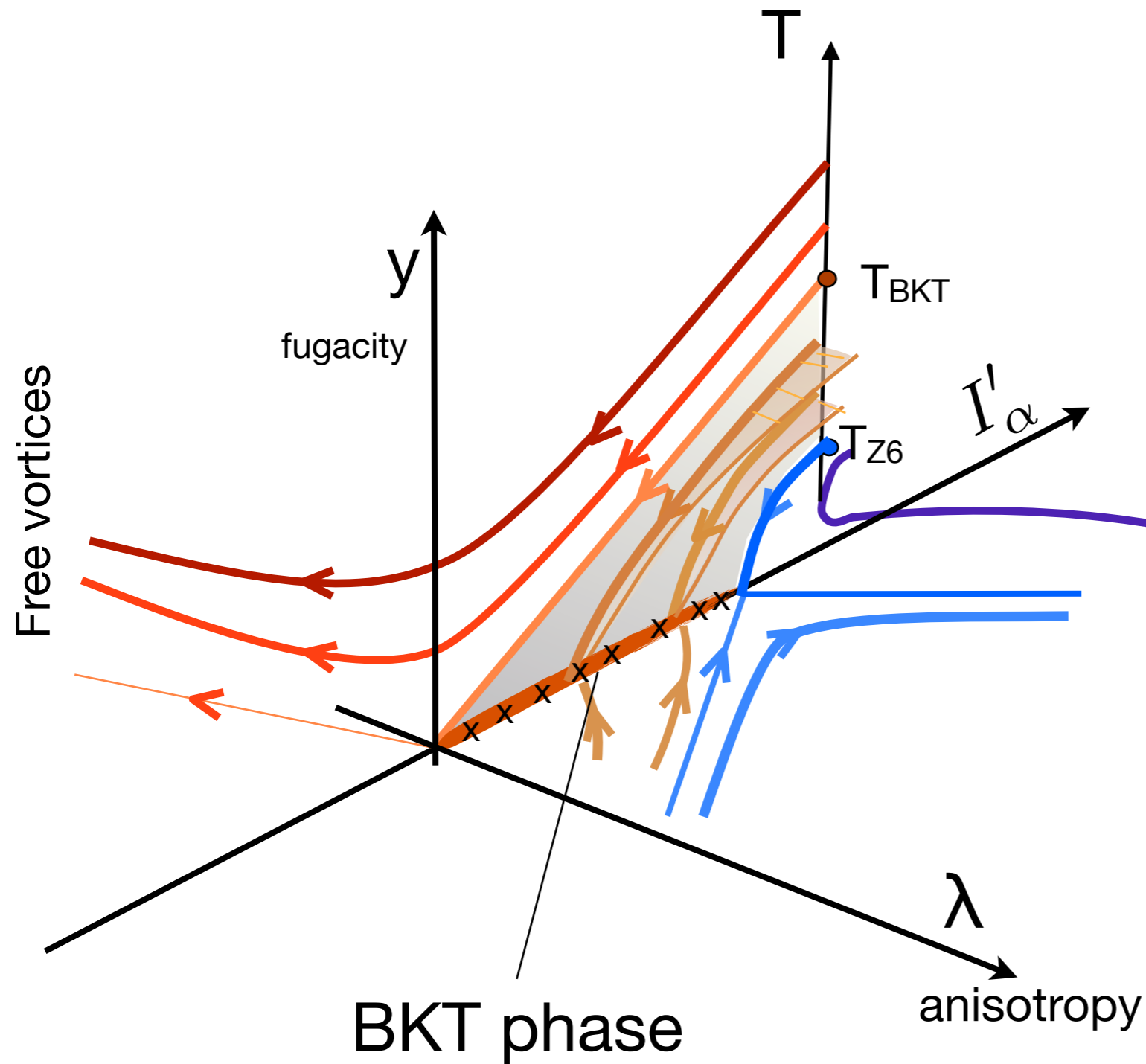
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Real space observation of spin frustration in Cr on a triangular lattice

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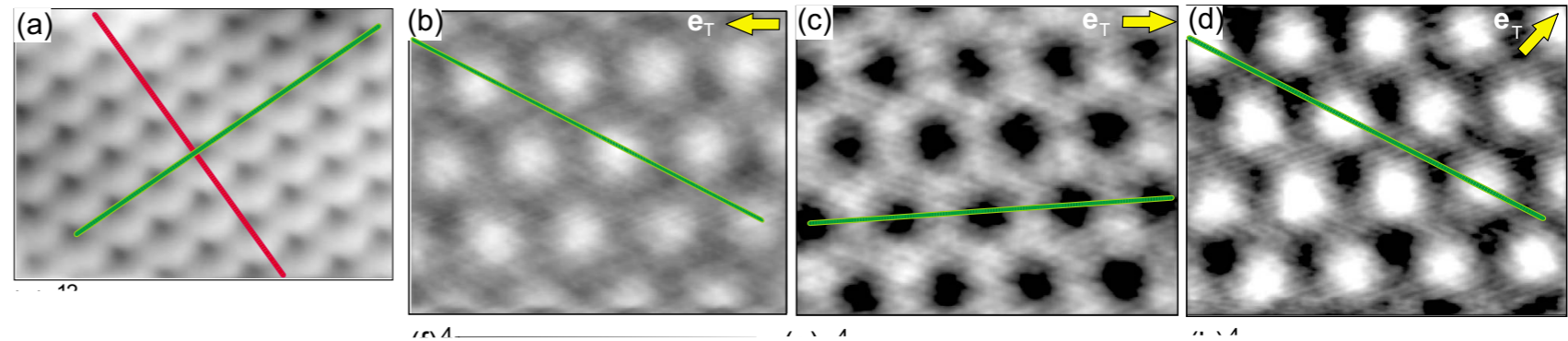
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Single layer Cr on Pd(111) surface
Spin Polarized STM (SP-STM)



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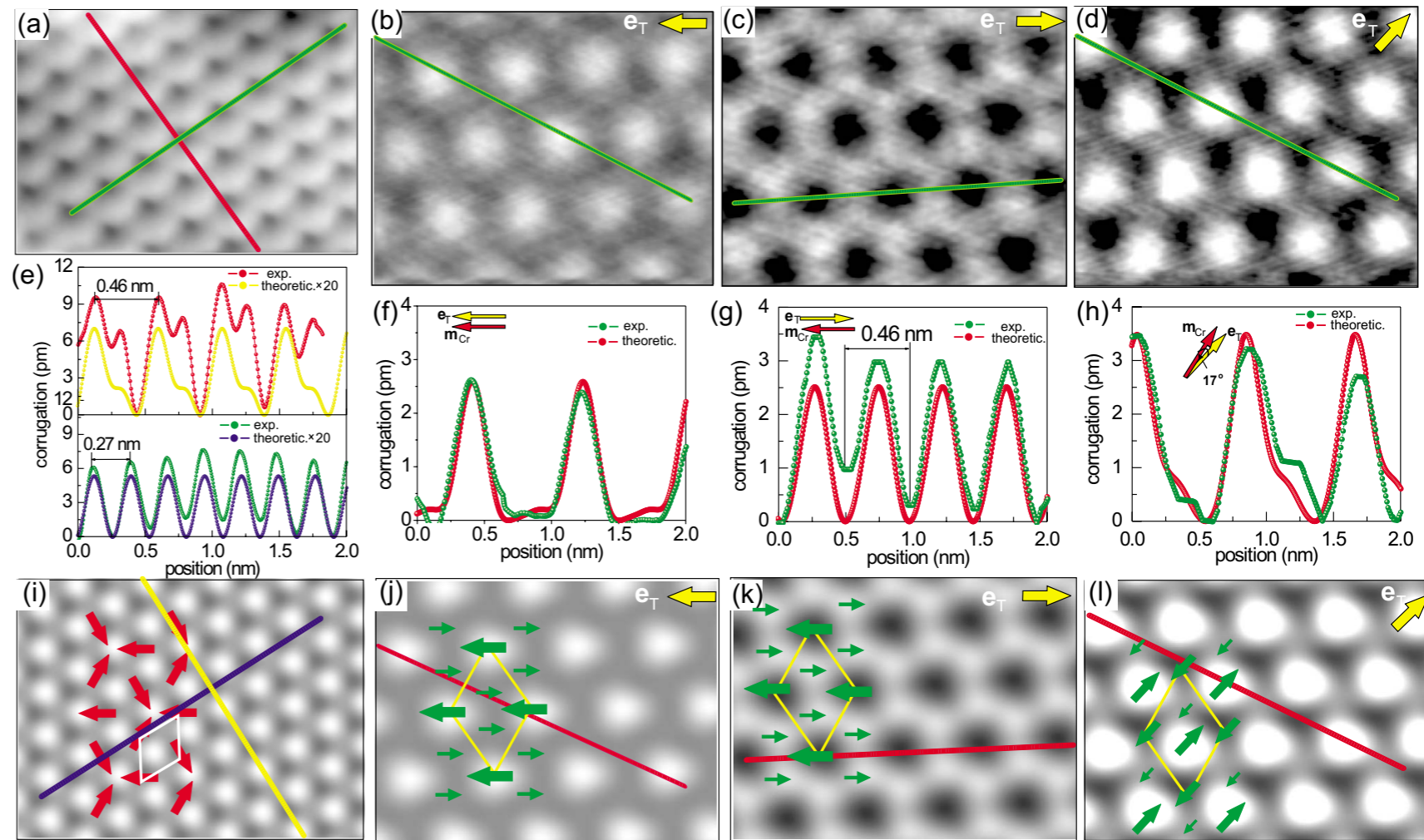
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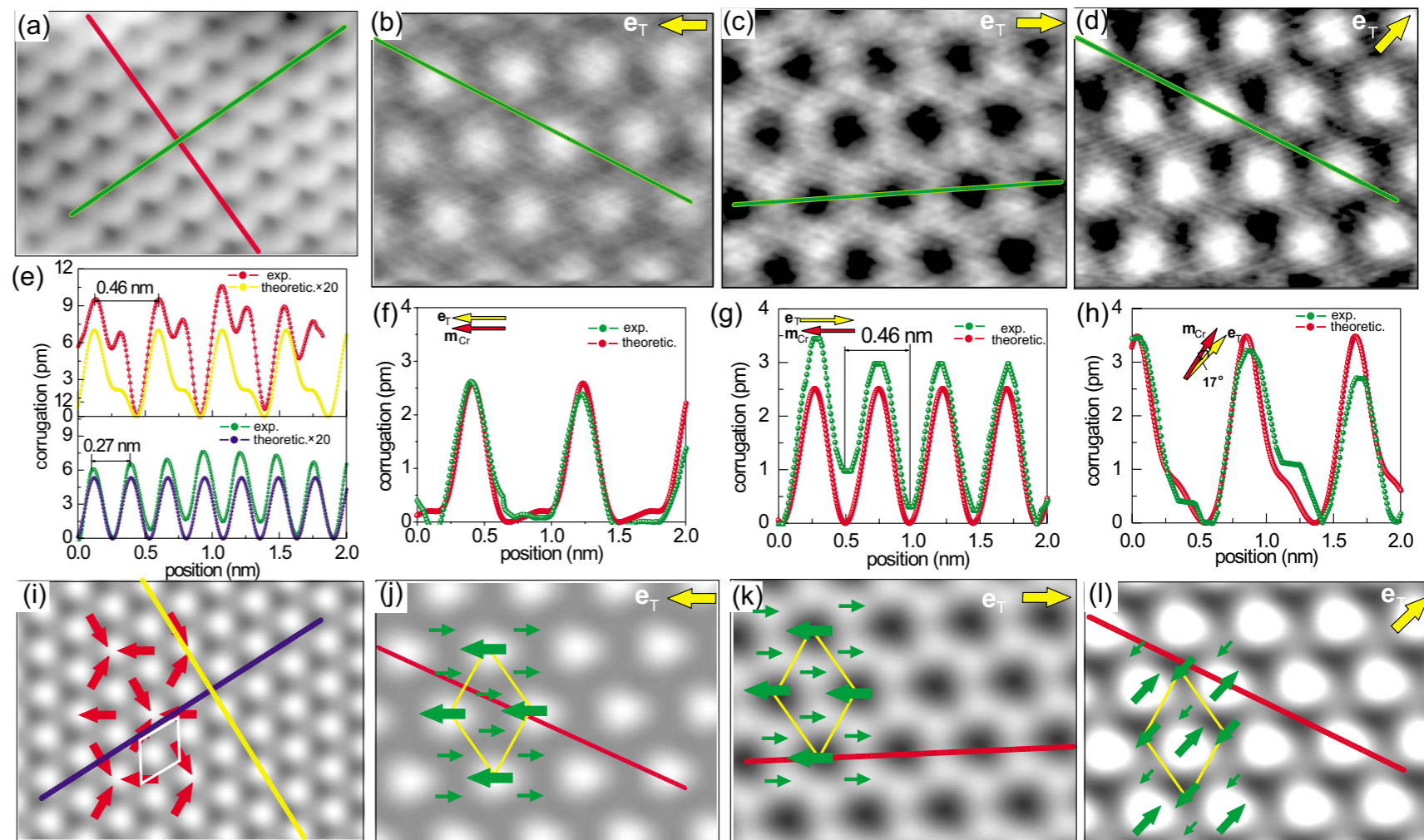
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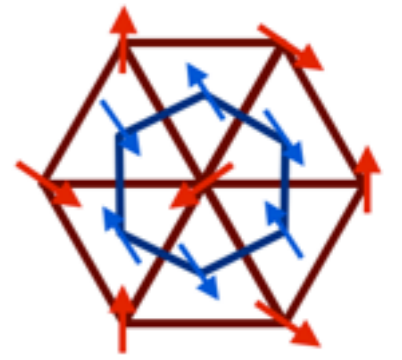
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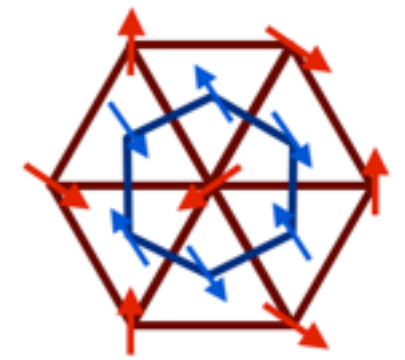
Bilayer: possible candidate for powerlaw phase?

S and Q



S and Q

- Emergent Z_6 Clock order in 2D triangular-hexagonal AFM.



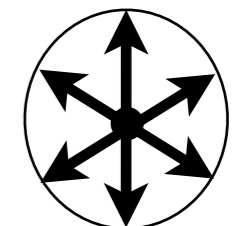
T

J_1

$J_1/\ln(J_1/J_2)$

T_{BKT}

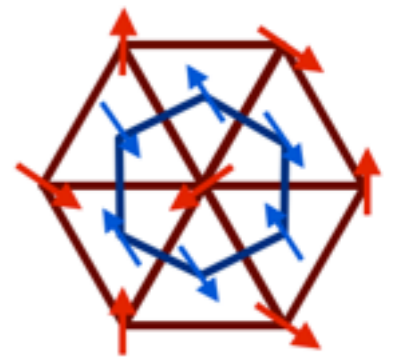
T_{Z_6}



6-state Clock Order

S and Q

- Emergent Z_6 Clock order in 2D triangular-hexagonal AFM.
- 2D Heisenberg model with emergent BKT transition.

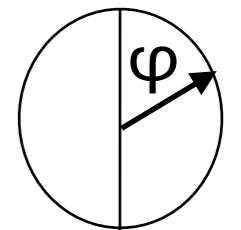
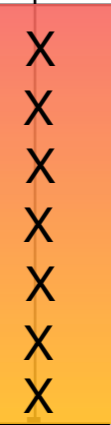


T

J_1

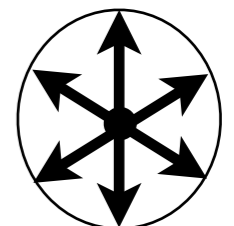
$J_1/\ln(J_1/J_2)$

T_{BKT}



Powerlaw phase

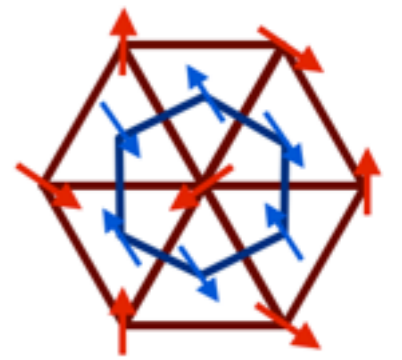
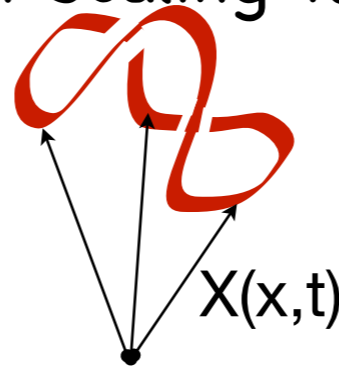
T_{Z6}



6-state Clock Order

S and Q

- Emergent Z_6 Clock order in 2D triangular-hexagonal AFM.
- 2D Heisenberg model with emergent BKT transition.
- Practical Application of Friedan Scaling to Magnetism



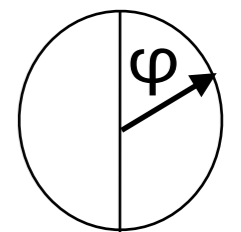
T

J_1

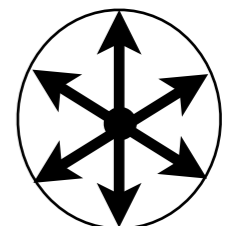
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T_{BKT}

T_{Z6}



Powerlaw phase



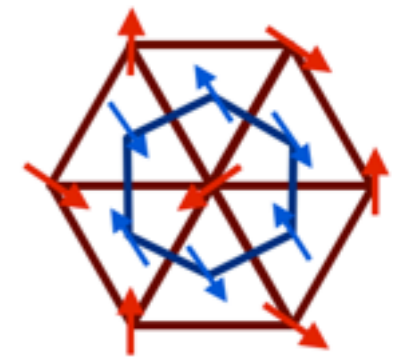
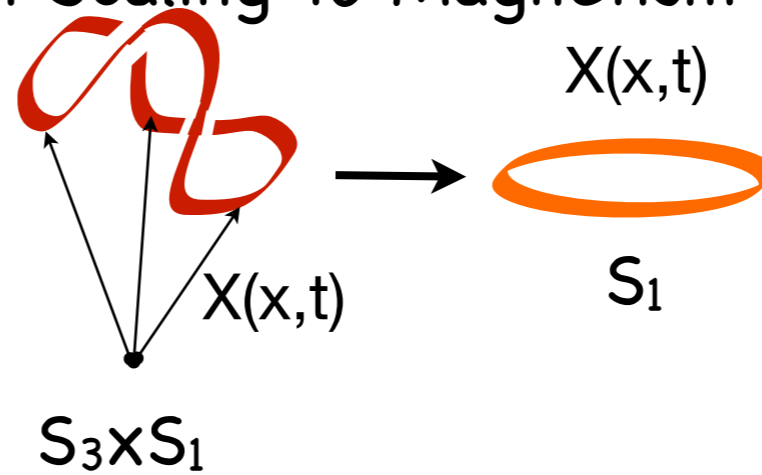
6-state Clock Order

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• Toy model for compactification



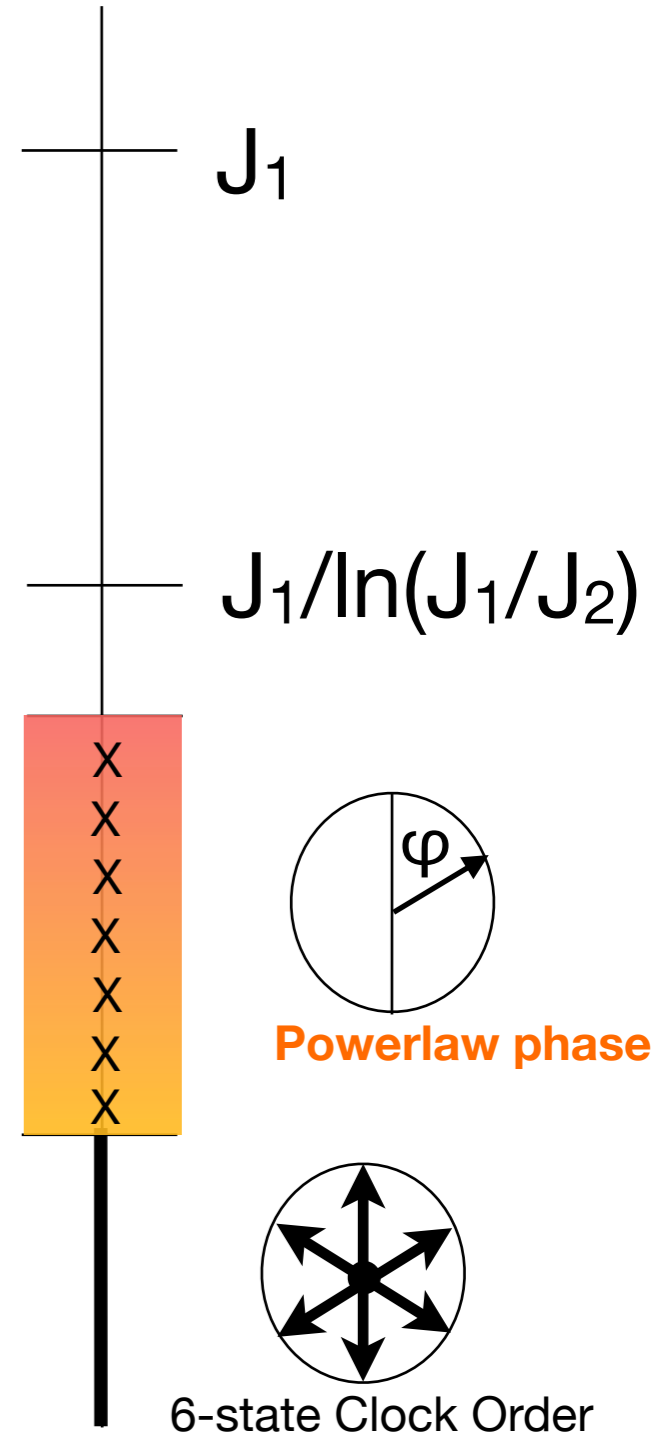
T

J_1

$J_1/\ln(J_1/J_2)$

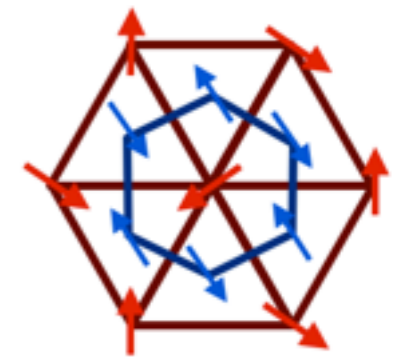
T_{BKT}

T_{Z6}



S and Q

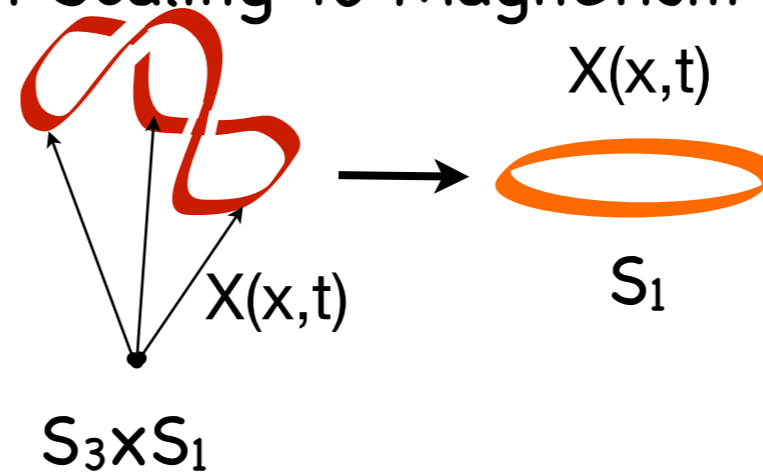
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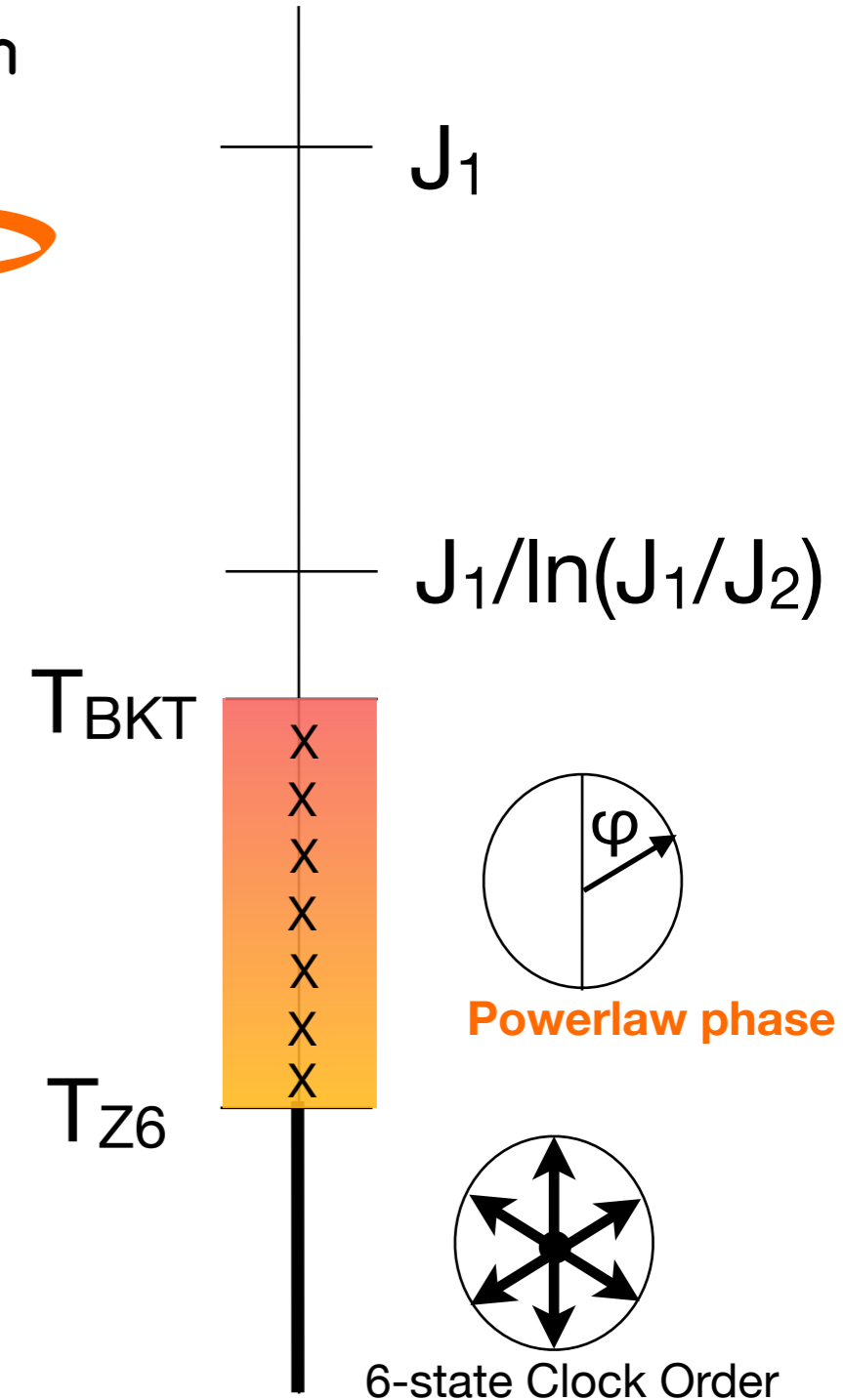
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J_1

- Realization of mapping of RG into time. (cf AdSCFT)

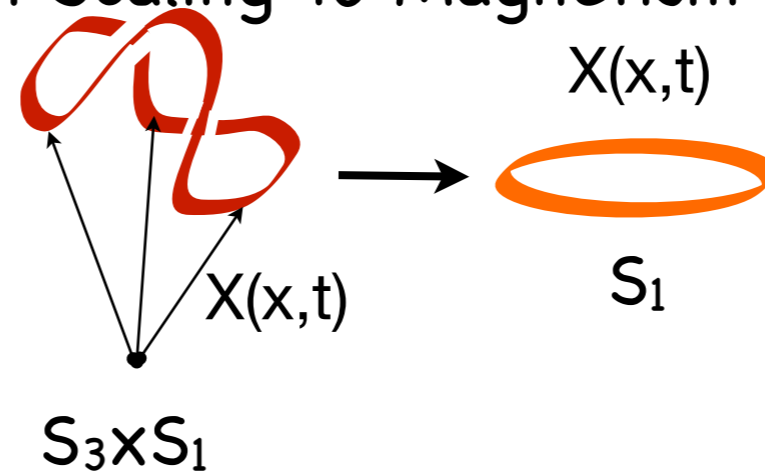


6-state Clock Order

S and Q

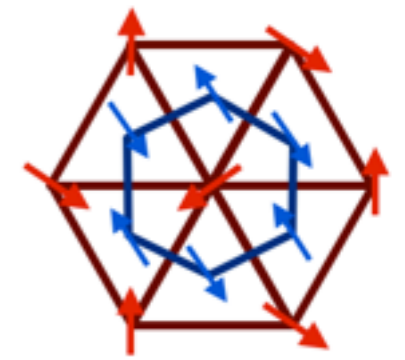
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- Realization of mapping of RG into time. (cf AdSCFT)

$$S = \int d\tau \left[\frac{1}{2} g^{ab} \dot{g}_{ab} - \frac{1}{2\pi} R \right]$$



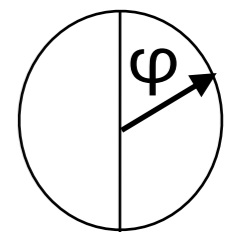
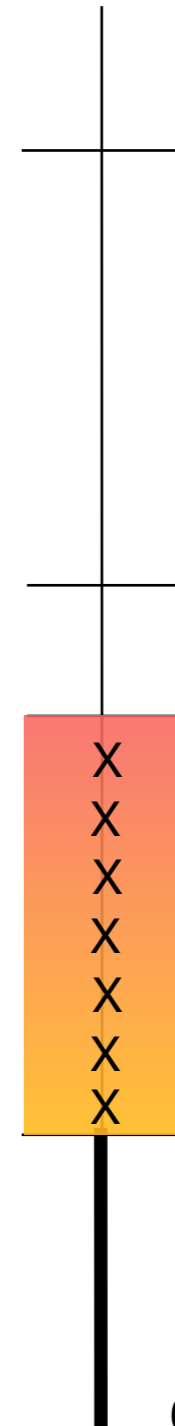
T

J₁

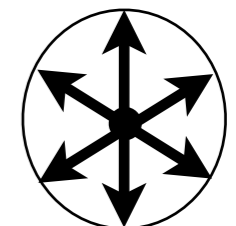
J₁/ln(J₁/J₂)

T_{BKT}

T_{Z6}



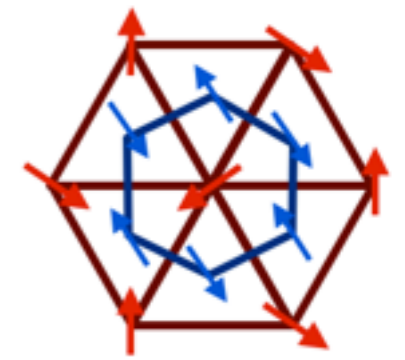
Powerlaw phase



6-state Clock Order

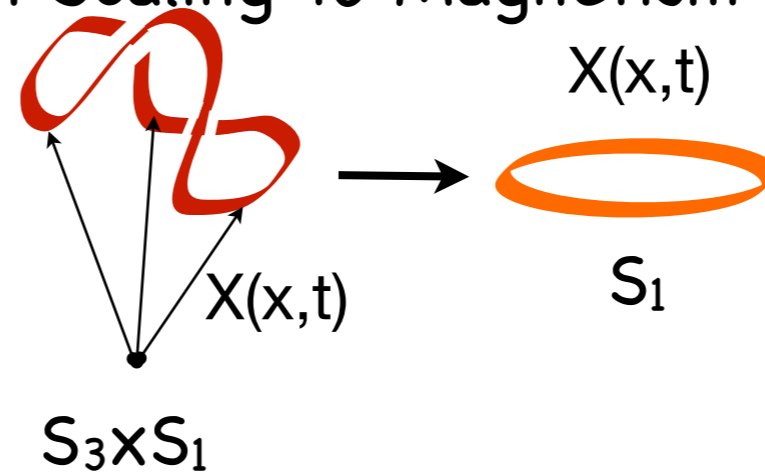
S and Q

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J₁

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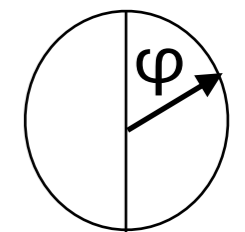
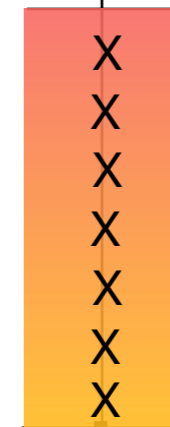
$$S = \int d\tau \left[\frac{1}{2} g^{ab} \dot{g}_{ab} - \frac{1}{2\pi} R \right]$$

$$\frac{\delta S}{\delta g^{ab}} = 0 = \dot{g}_{ab} - \frac{1}{2\pi} R_{ab}$$

Thanks: D. Friedan

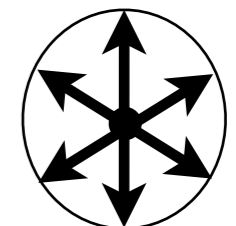
T_{BKT}

J₁/ln(J₁/J₂)



Powerlaw phase

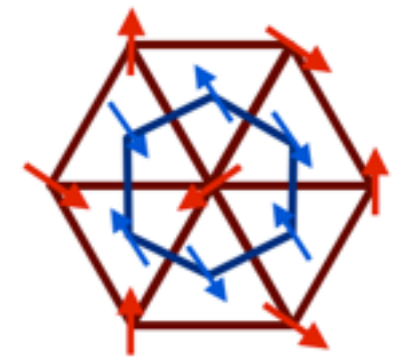
T_{Z6}



6-state Clock Order

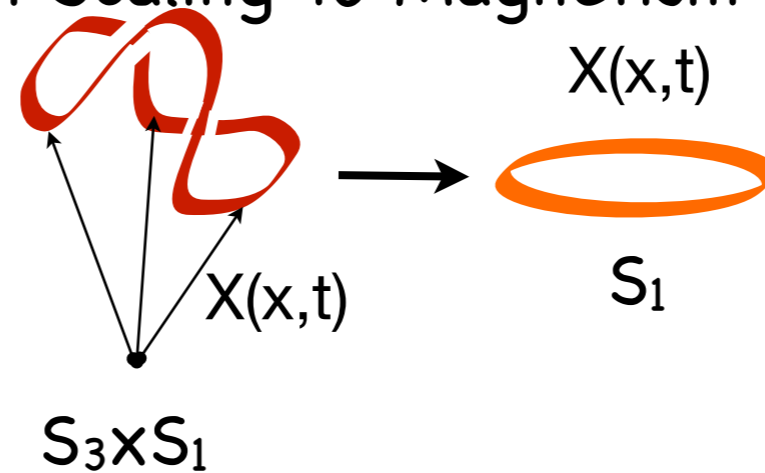
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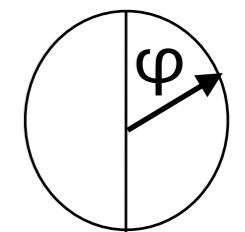
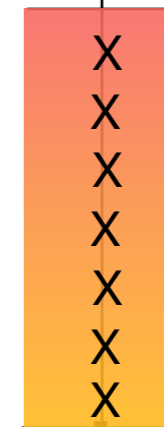
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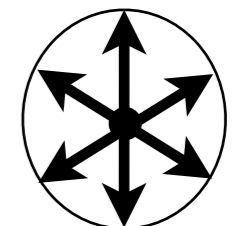
T_{BKT}

J₁/ln(J₁/J₂)



Powerlaw phase

T_{Z6}

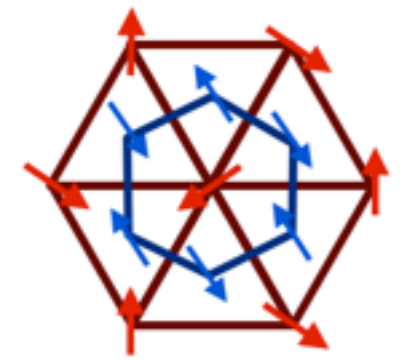


6-state Clock Order

- Quantum version?

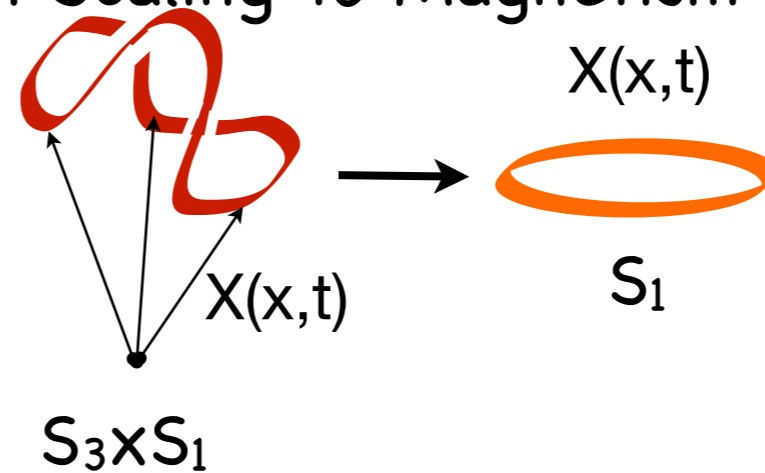
S and Q

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J₁

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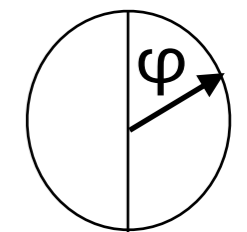
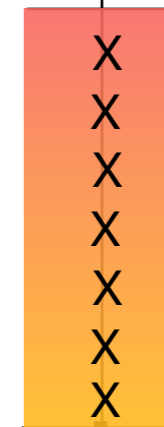
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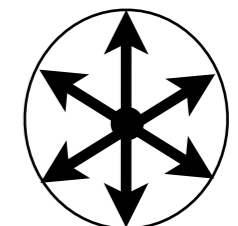
T_{BKT}

J₁/ln(J₁/J₂)



Powerlaw phase

T_{Z6}



6-state Clock Order

- Quantum version?

Can one suppress T_{Z6} to zero : power law spin-liquid?

Thank you!