Fluctuation induced ordering on the kagome lattice

Gia-wei Chern

Los Alamos

Roderich Moessner MPI-PKS

mpipks



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Outline

Order by disorder

- degeneracy and frustration
- zero, soft and hard modes
- Phase diagram of KHAFM
 - coplanar order by disorder and dynamical symmetry
 - numerical freezing

New algorithm

stable low-T regime

Effective descriptions

- effective Potts model
- two-component height model
 - ⇒ dipolar ordering via KT transition



Kagome lattice = corner-sharing triangles

$$\mathcal{H} = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j = \frac{J}{2} \sum_{\alpha} \vec{\ell}_{\alpha}^2$$

 \Rightarrow Each triangle: $\vec{\ell}_{\alpha} = \vec{S}_{\alpha 1} + \vec{S}_{\alpha 2} + \vec{S}_{\alpha 3} = 0$

Triangle-based network: *local* "weathervane mode" at zero energy Chandra/Harris/Chalker et al. 1992

- ► local d.o.f.!
- extensive ground-state dimension of kagome magnet





Degeneracies are intrinsically unstable

Extensive ground-state degeneracy

- huge low-energy d.o.s.
- \Rightarrow any perturbation is strong!
 - subleading terms determine low-T behavior

Frustrated magnets have three regimes

- high-temperature paramagnet
- cooperative paramagnet ("universal")
- non-generic low-T

 Θ_{cw}/T_F parametrises strength of frustration Obradors, Ramirez



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(Thermal) order by disorder

At T > 0, minimise $F = U - TS \approx -TS$

maximise entropy!



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Ordered states tend to have highest entropy!

Ground states: x = 0 or y = 0

► can in principle explore both axes

Ground states weighted by fluctuations

 $\lim_{T\to 0} A_{\epsilon}/A_{T} \to 1!$

 \Rightarrow System localised near origin Mechanism: two 'soft' directions

'crossing' points of ground-states



All coplanar states have extensive number of soft (harmonic zero) modes! Chalker et al., 1992

- ► quadru-, octupolar order: (e^{3iΘ}) Zhitomirsky 2008
- perfect extensive degeneracy of free energy at harmonic level
- maps onto 3-state Potts model

 $\mathsf{Huse} + \mathsf{Rutenberg} \ 90\mathsf{s}$

Order is cut off by T at exp. large distance Mermin + Wagner



Analytics: unknown (purely anharmonic problem) Holdsworth + Shender

 unweighted coplanar Potts ensemble is critical Baxter

Numerics: unknown

 algorithms freeze as coplanar obdo sets in Reimers, Huse+Rutenberg, Zhitomirsky (1992-2012)

 not even small systems can be solved exactly



Cluster algorithm for Heisenberg model

Exploit dynamical symmetry: $\eta = \pm 1$ maps between coplanar states Hassan+R.M.

$$\begin{split} \dot{\ell}^{1}_{\alpha} + \dot{\ell}^{2}_{\alpha} &= \left(\eta S^{2}_{\alpha\beta}\right) \left(\eta \ell^{3}_{\beta}\right) - S^{3}_{\alpha\beta} \ell^{2}_{\beta} \\ \dot{\ell}^{2}_{\alpha} - \dot{\ell}^{1}_{\alpha} &= S^{3}_{\alpha\beta} \ell^{1}_{\beta} - \left(\eta S^{1}_{\alpha\beta}\right) \left(\eta \ell^{3}_{\beta}\right) \\ \left(\eta \dot{\ell}^{3}_{\alpha}\right) &= \left(\eta S^{1}_{\alpha\beta}\right) \ell^{2}_{\beta} - \left(\eta S^{2}_{\alpha\beta}\right) \ell^{1}_{\beta} \end{split}$$

Conventional cluster algorithm fails

small fluctuations around coplanar states

Write $\vec{S} = \vec{S}^{(0)} + \vec{s}$, cluster algo acts on $\vec{S}^{(0)}$: colour exchange around loop

• Boltzmann factor only for 'anharmonic part' \vec{s}_i

Can equilibrate systems with $O(10^3)$ spins for $T \rightarrow 0$.



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Low-T regime reached for $T \lesssim 10^{-4} J$

• "entropic" Boltzmann factors: $\exp\left(-\frac{TS}{T}\right) = \exp(-S)$ is T-independent _{cf Henley}

Enhanced correlations saturate

thermodyn. limit still not reachable!



Coplanar states map onto (2-component) Gaussian height model: Holdsworth + Shender

 $ec{z}_eta = ec{z}_lpha + ec{S}_{lphaeta}$

Harmonic (XY=Potts) problem at exactly solvable critical point Baxter; Huse+Rutenberg



Role of anharmonic fluctuations?

Rough phase \Rightarrow algebraic spin correlations

Kosterlitz-Thouless physics

 $\mathsf{Flat}\ \mathsf{phase} \Rightarrow \mathsf{long}\mathsf{-range}\ \mathsf{order}$

can classify all candidate orders

Height-space is periodic

► histogram in unit cell ⇒ nature of ordering



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Enhanced stiffness \implies dipolar order

Gaussian model: $S = \frac{K}{2} \int |\nabla z|^2 d^2 r$



K grows upon inclusion of fluctuations? Yes!

$$\Theta(q) = \mathrm{K}_{\mathrm{H'berg}}/\mathrm{K}_{\mathrm{potts}} = \left|z_{\mathrm{potts}}(q)/z_{\mathrm{H'berg}}(q)
ight|^2$$

Enhanced but weak maximum at Γ point suggests $\sqrt{3}\times\sqrt{3}$ order

thermodynamic limit?



Try to fit entropic weights by additional interaction

simulate effective model

Single parameter \mathcal{J}_2 suffices: n.n.n. interactions only



 $\mathcal{J}_2 \approx 0.019$ is quite small

Data collapse and KT transition

Can simulate up to $10^{6}\ \rm{spins}$

- consider broader phase diagram
- Data collapse for KT transition with $\mathcal{J}_2^c=0$
 - ▶ small ordered moment, $< 10\% m_{\rm SAT}$

Hard to discern because of long correlation length

- rather than smallness: cf. chirality
- further phase transitions?

Castelnovo et al.



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The Heisenberg kagome afm

Complex instance of obdo

- harmonic: coplanar ('octupolar') order
- ► anharmonic: ('dipolar') spin order
- ordering is delicate
 - small T = 0 order parameter (classically!)
 - large crossover length

New efficient algorithm cf. Schnabel+Landau

► can access low-*T* regime

Effective field theory + Potts model

access long wavelengths+large-scales

Beyond n.n. Heisenberg model

- ► semiclassics Henley; perturbations
- connection to experiment

Also: pyro. XY; hyperkagome/garnet; ...

