

Skyrmion Dynamics in Thin Films of Chiral Magnets

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Skyrmions and topological transport phenomena

Skyrmions in multiferroics toward E -control

Forced and spontaneous dynamics of skyrmions

Collaborators

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- ***PSI (Switzerland)***

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- ***Groningen Univ. (Holland)***

M. Mostovoy

- ***Fudan Univ. (Peoples R China)***

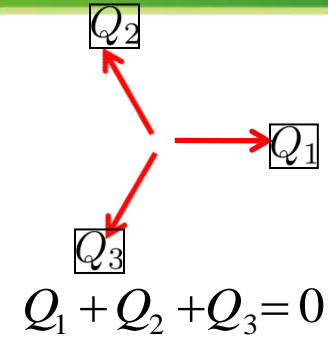
Yufan Li, Xiofeng Jin

Skyrmion Crystal

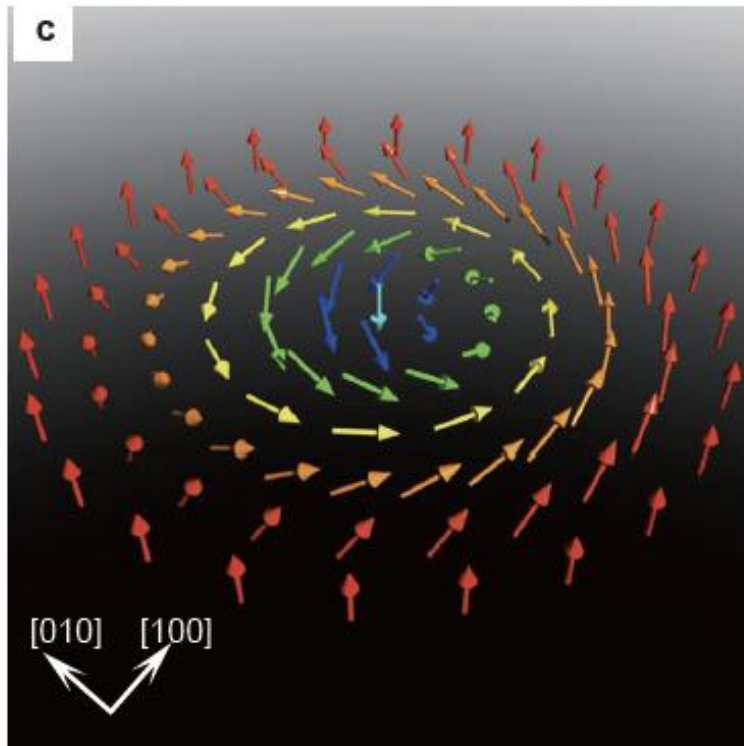
Superposition of three Helix without phase shift

$$M(r) \approx M_f + \sum_{i=1}^3 M_{Q_i}(r + \Delta r)$$

$$M_{Q_i}(r + \Delta r) = A[n_{i1}\cos(Q_i \cdot r) + n_{i2}\sin(Q_i \cdot r)]$$

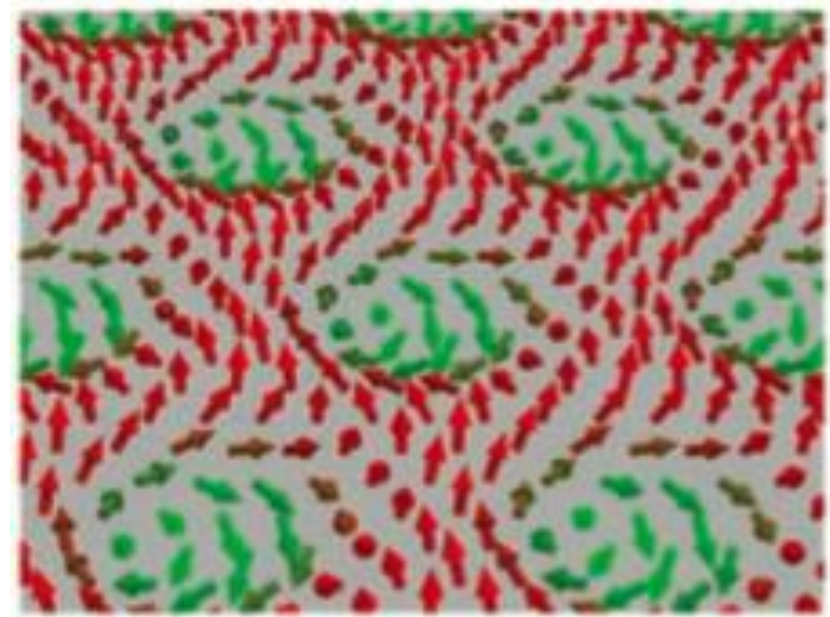


Skyrmion



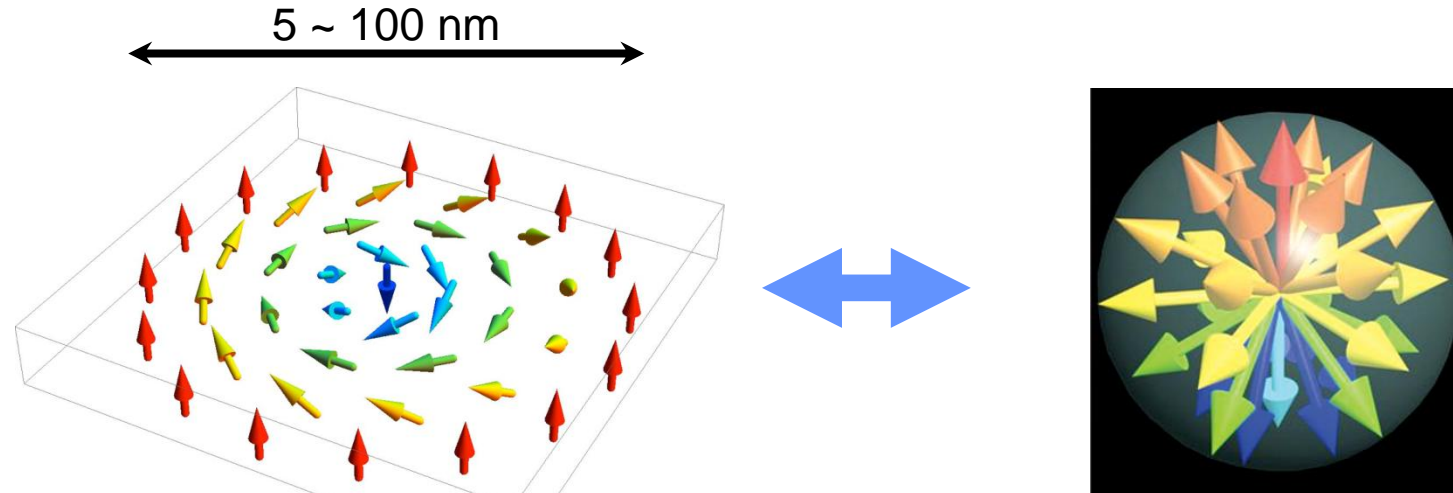
Skyrmion crystal

3-flod-Q



S. Muhlbauer et al. Science 323, 915 (2009).

What is magnetic skyrmion?

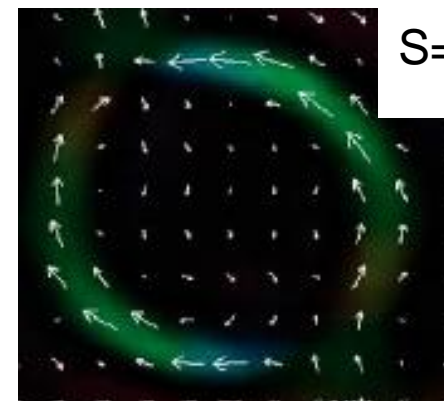
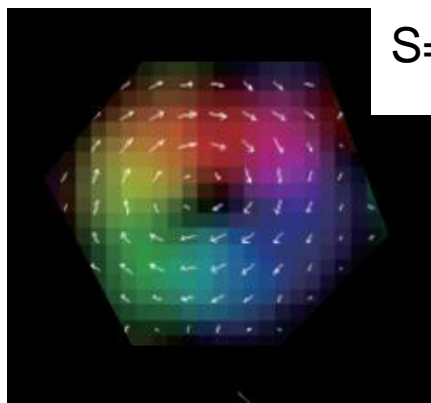


Topologically-stable spin vortex
with particle-like nature

“skyrmion number”

$$S = \frac{1}{4\pi} \int \vec{n} \cdot \frac{\partial \vec{n}}{\partial x} \times \frac{\partial \vec{n}}{\partial y} d\vec{r} = -1$$

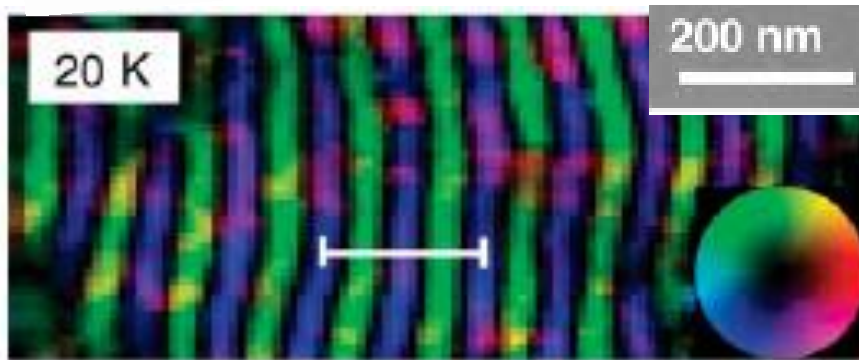
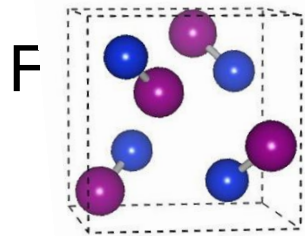
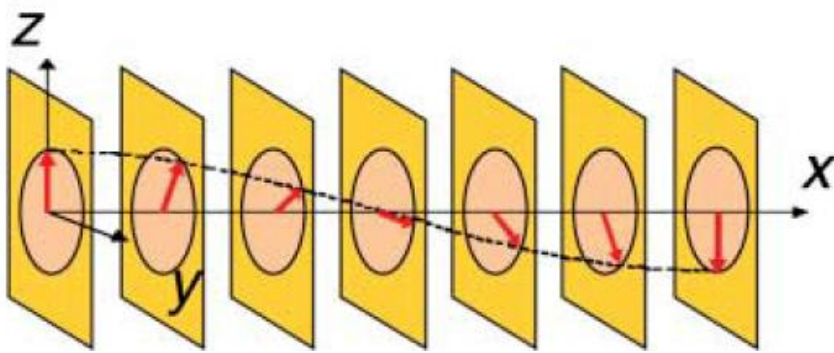
Lateral component of
M of some bubbles



a pair of
Bloch lines

Toward real space observation of Skyrmion structure

structure



M. Uchida, Y. Onose, Y. Matsui, Y. Tokura,
Science (2006)

$$H = \sum \left(\underbrace{-J\vec{S}_i \cdot \vec{S}_j}_{\text{Ferro}} + \underbrace{\vec{D}_{ij} \cdot (\vec{S}_i \times \vec{S}_j)}_{\text{DM}} \right)$$

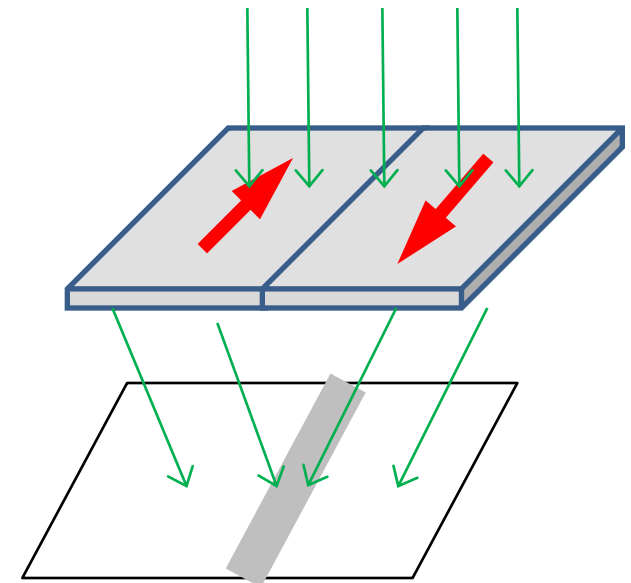
Ferro + DM



Helical spin structure

Long period $\sim aJ/D \sim 10\text{nm}-300\text{nm}$


Lorentz microscope
electrons

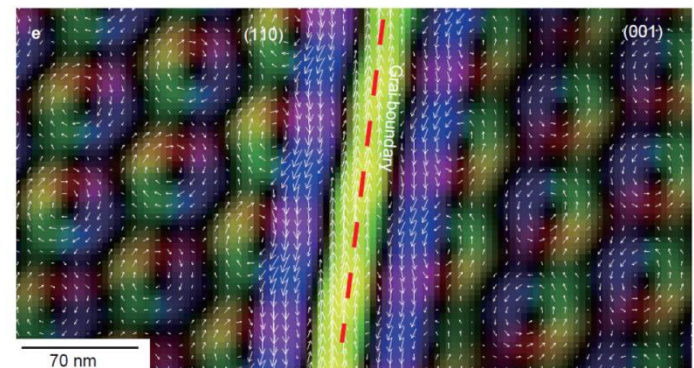
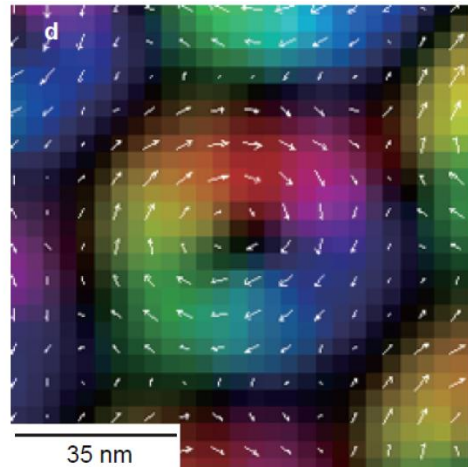
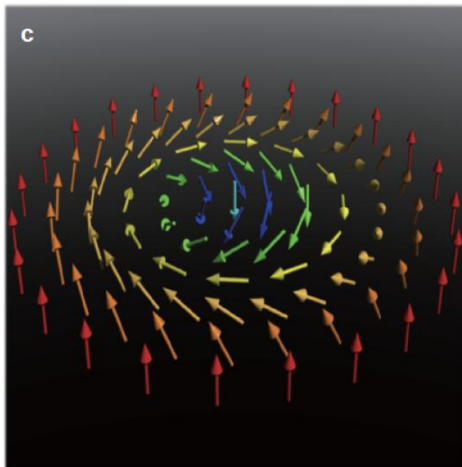
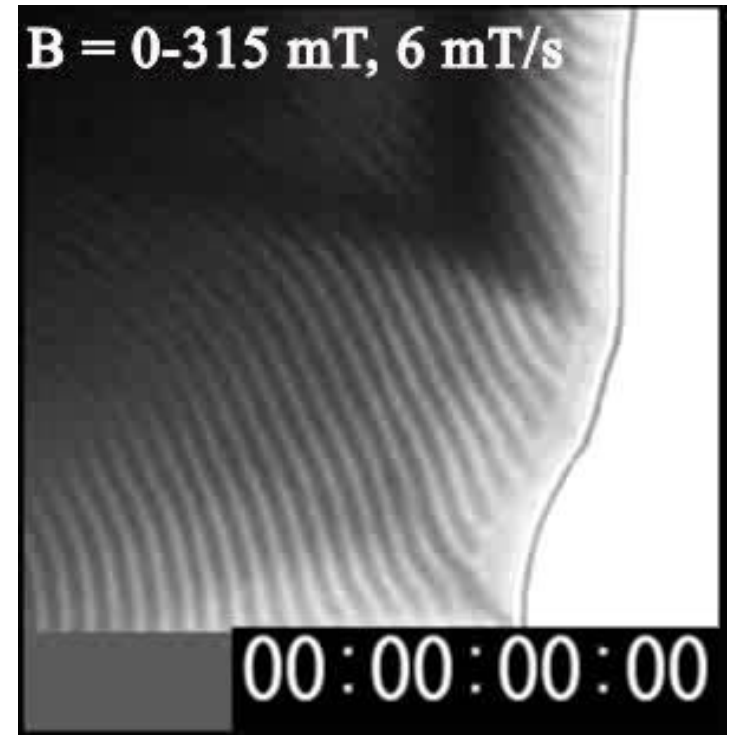
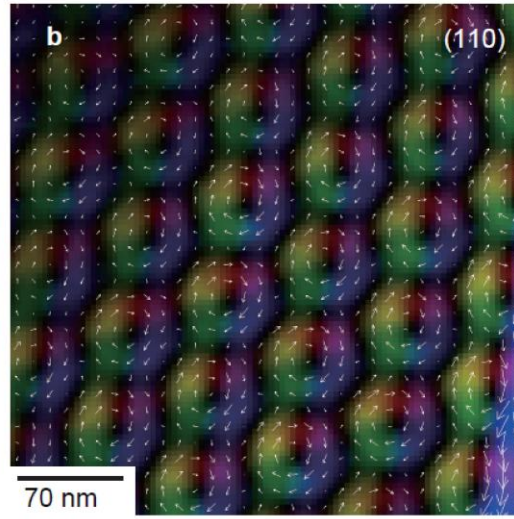
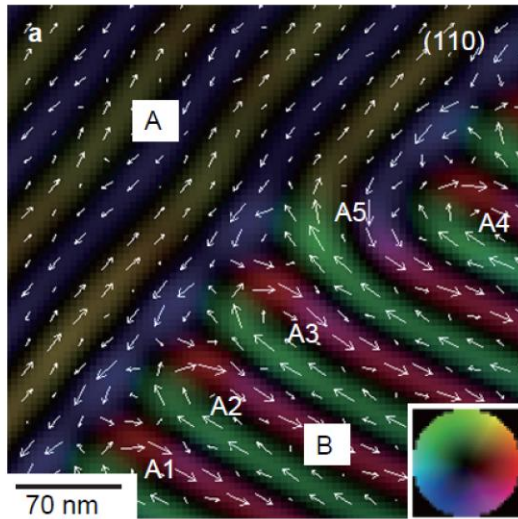


FeGe: from helical to skyrmion crystal at 260K

X.Z. Yu et al. Nat. Mater.(2011)

H=0

 H=0.1T



Near room-temperature formation of SkX in FeGe⁷

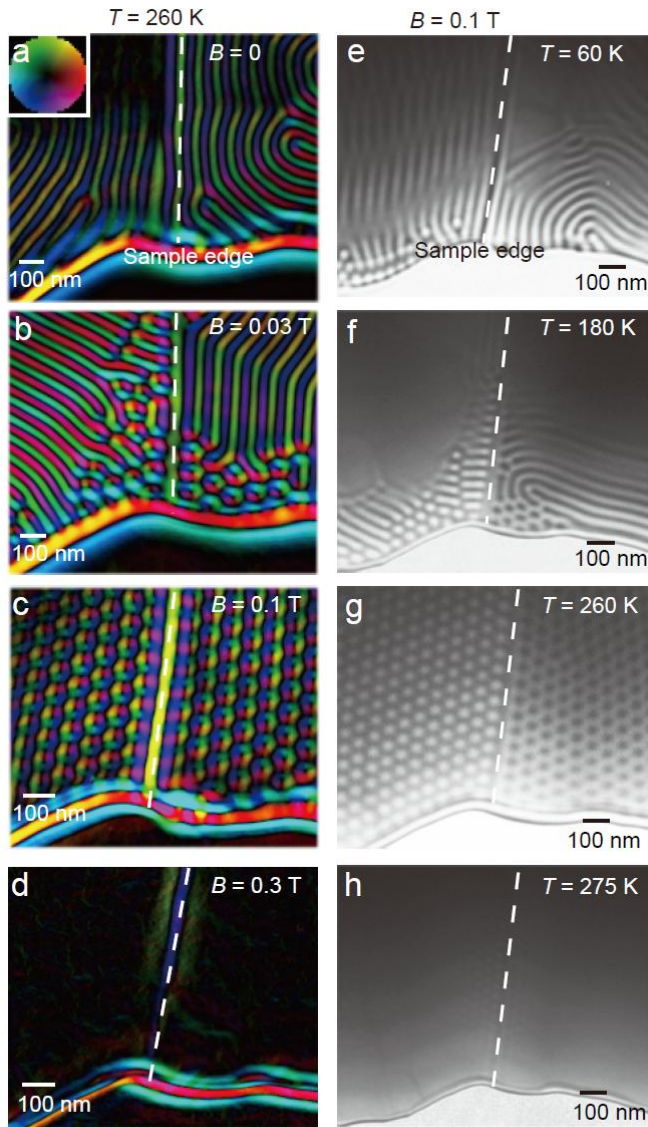


FIG. 2 Yu *et al.*

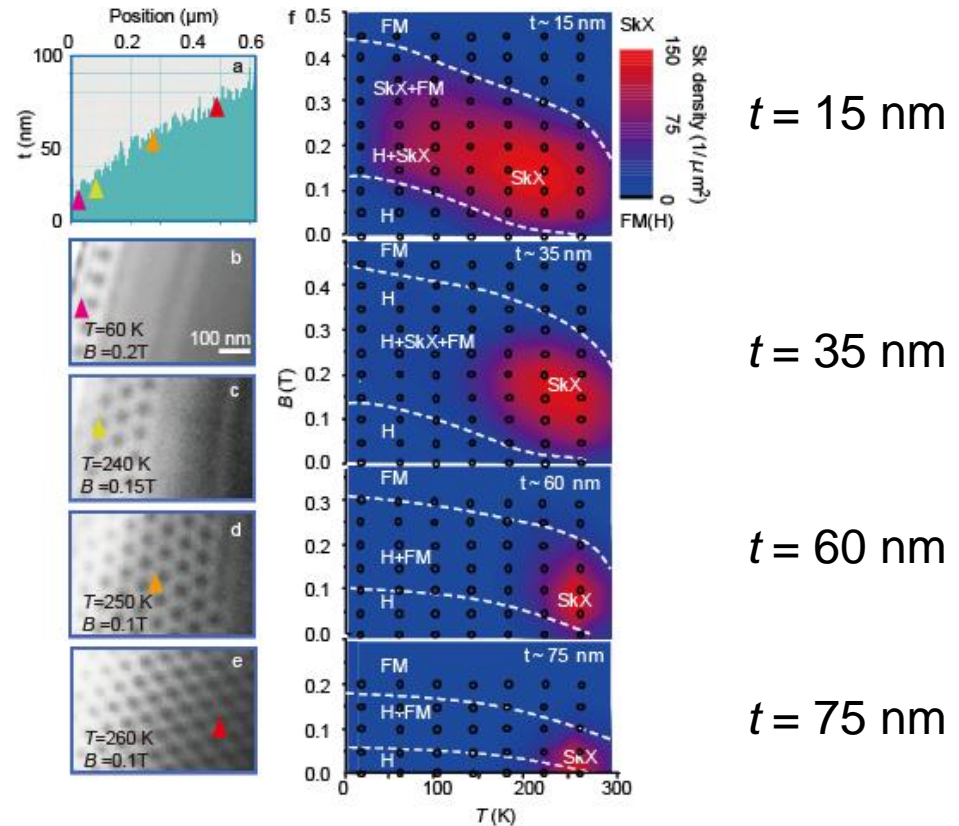
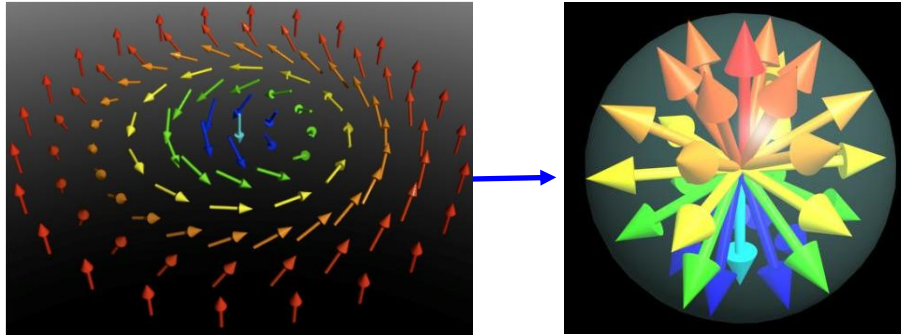


FIG. 3 Yu *et al.*

-Near room-temperature formation of SkX
 -Stability of SkX depend largely on the thickness.

X. Z. Yu, N. Kanazawa, Y. Onose, K. Kimoto, W.Z. Zhang, S. Ishiwata, Y. Matsui, and Y. Tokura, *Nature Mater.* 10 106 (2011)

Real-space fictitious magnetic field in a skyrmion spin texture

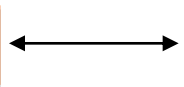


scalar spin chirality

Solid angle $\Omega = 4\pi$

In strong coupling case

One skyrmion



One magnetic flux ϕ_0

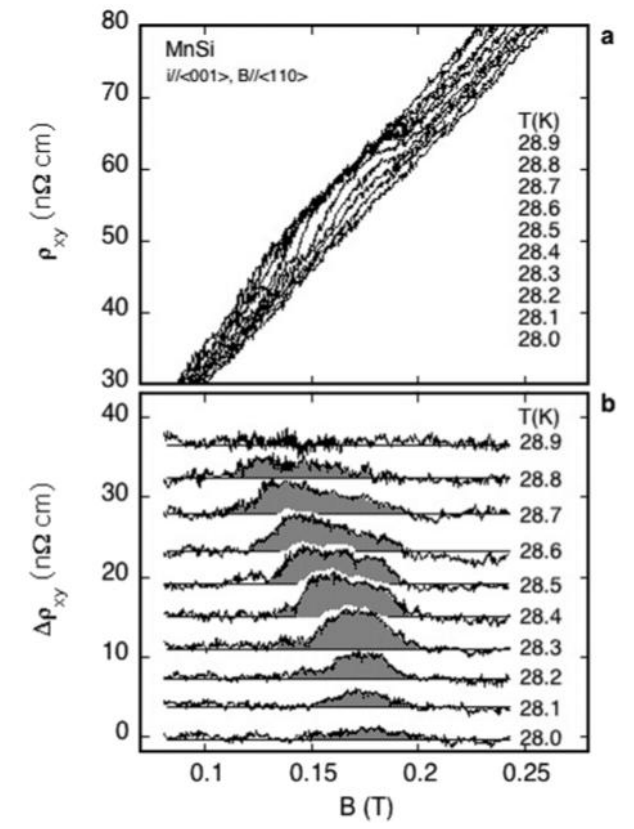
$$\phi_0 = h/e$$

Emergent magnetic field

$$\mathbf{B}_{\text{eff}}^z = -\phi_0/A$$

A: skyrmion size

High skyrmion density \Leftrightarrow Large topological Hall Effect

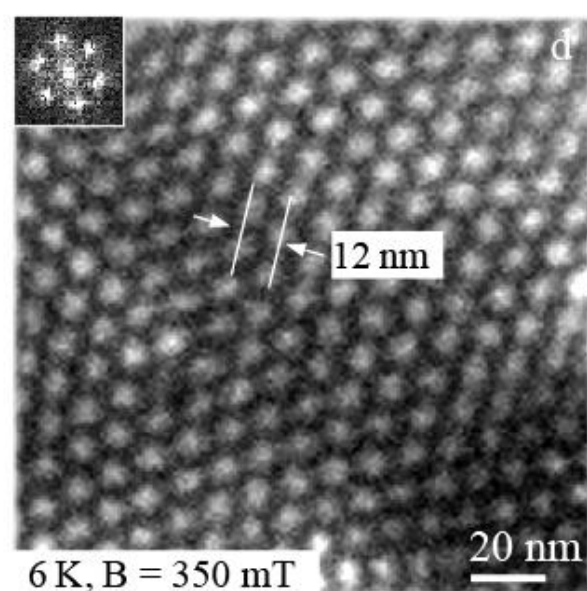
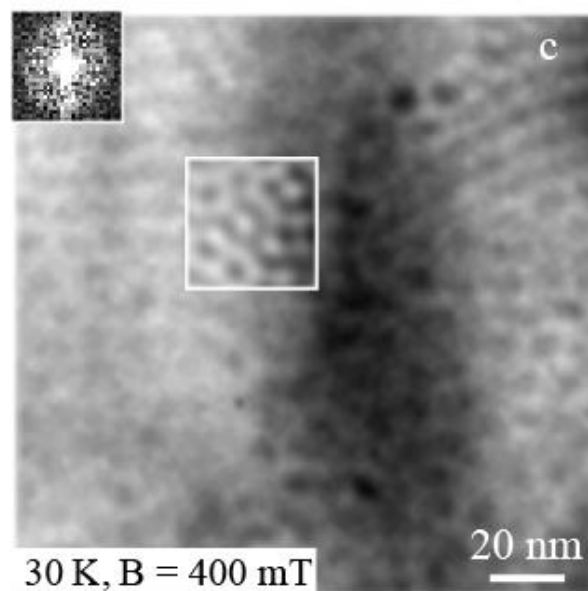
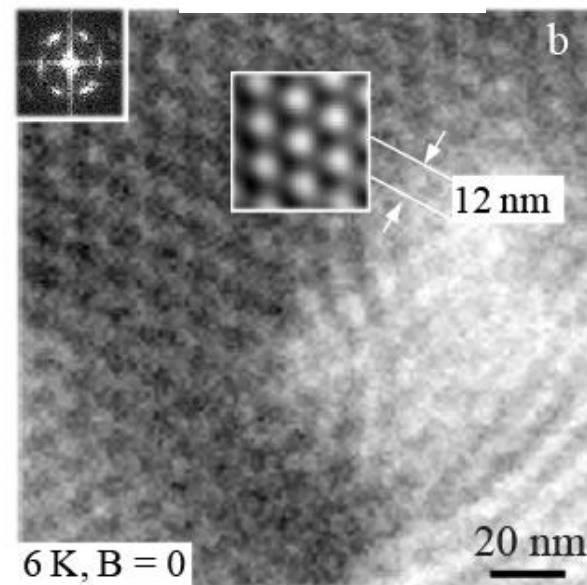
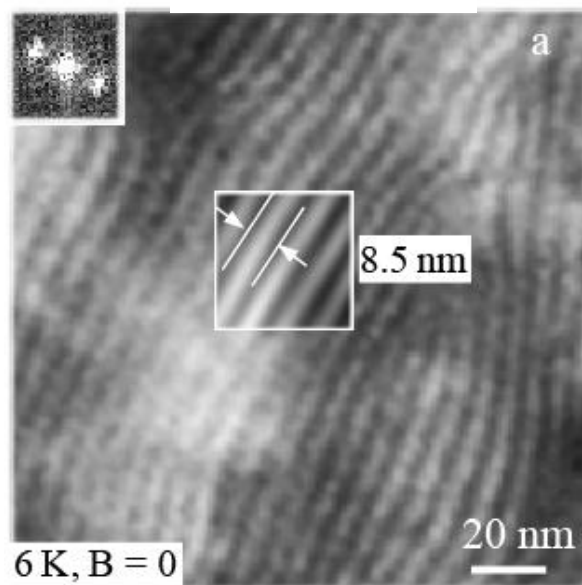
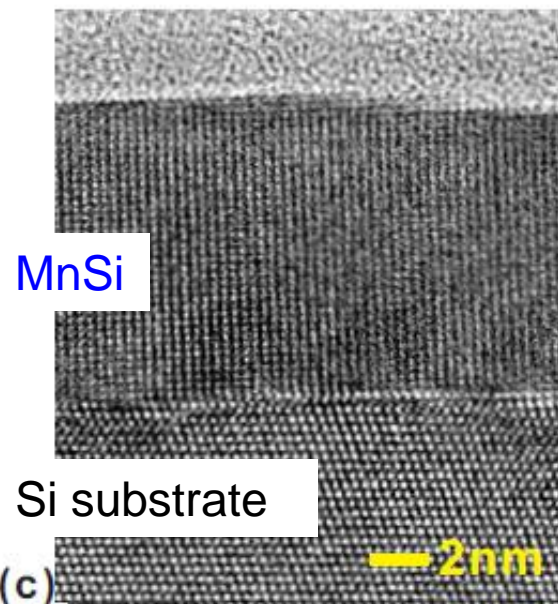
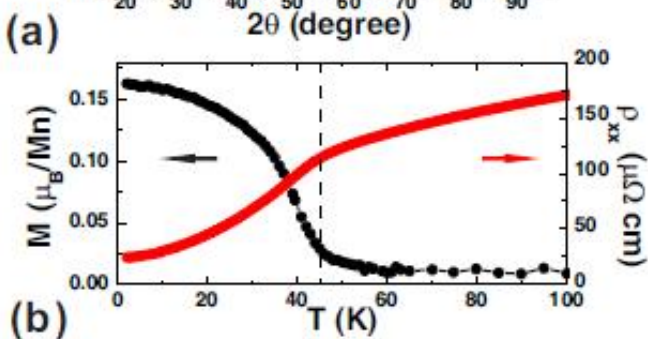
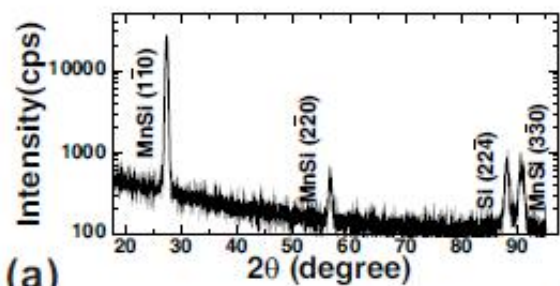


A. Neubauer *et al*, PRL 102 186602 (2009)

Ultrathin epitaxial thin films of MnSi

10nm-thick

20nm-thick



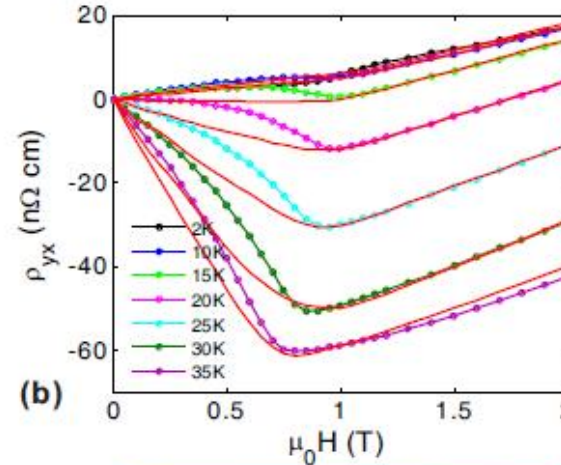
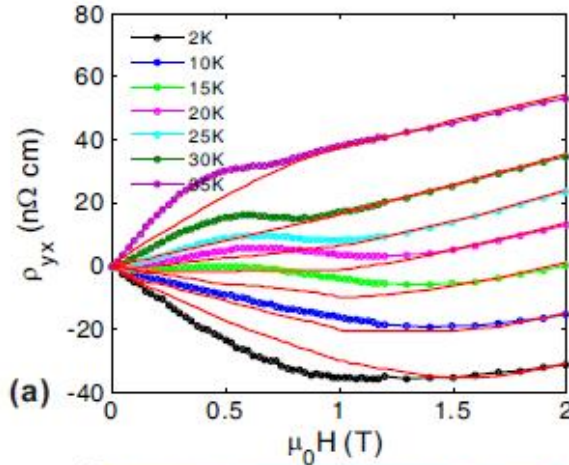
Skyrmion phase mapping by topological Hall resistivity

Yufan Li, Kanazawa, Kagawa

Conventional **anomalous** + **normal** Hall effects

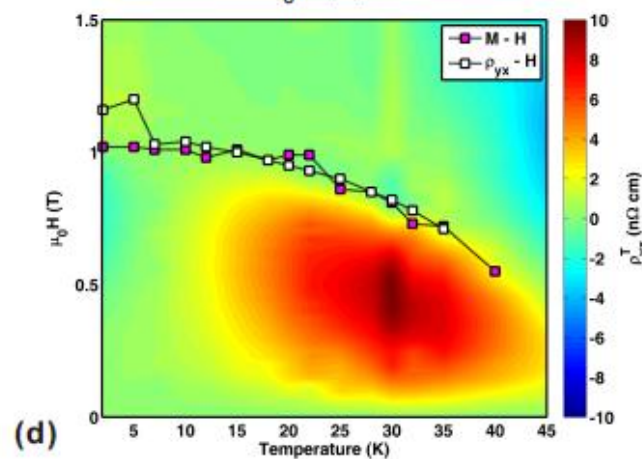
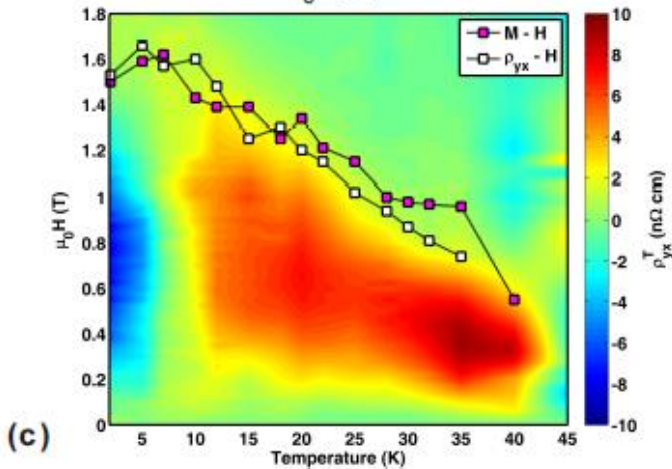
10nm-thick

50nm-thick



$$\rho_{yx}^A = \underbrace{R_0 B_z}_{\text{normal}} + \underbrace{\mu_0 R_S M_z}_{\text{anomalous}}$$

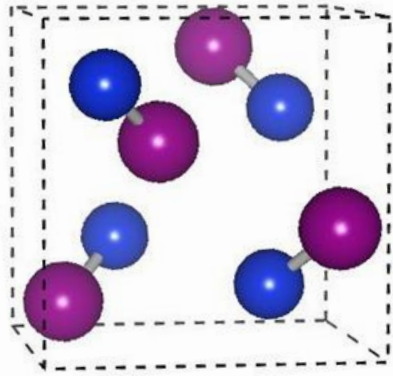
$$\mu_0 R_S = S_A \rho_{xx}^2$$



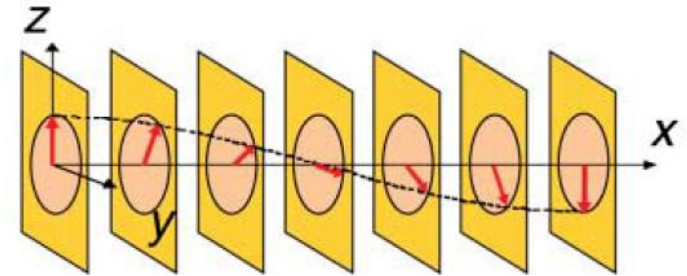
See also the late paper on FeGe thin film;
S. X. Huang and C. L. Chien, Phys. Rev. Lett. **108**, 267201 (2012)

Magnetic phase diagrams of B20 TMSi, TMGe

B20 structure



Cubic but noncentros
(Chiral)



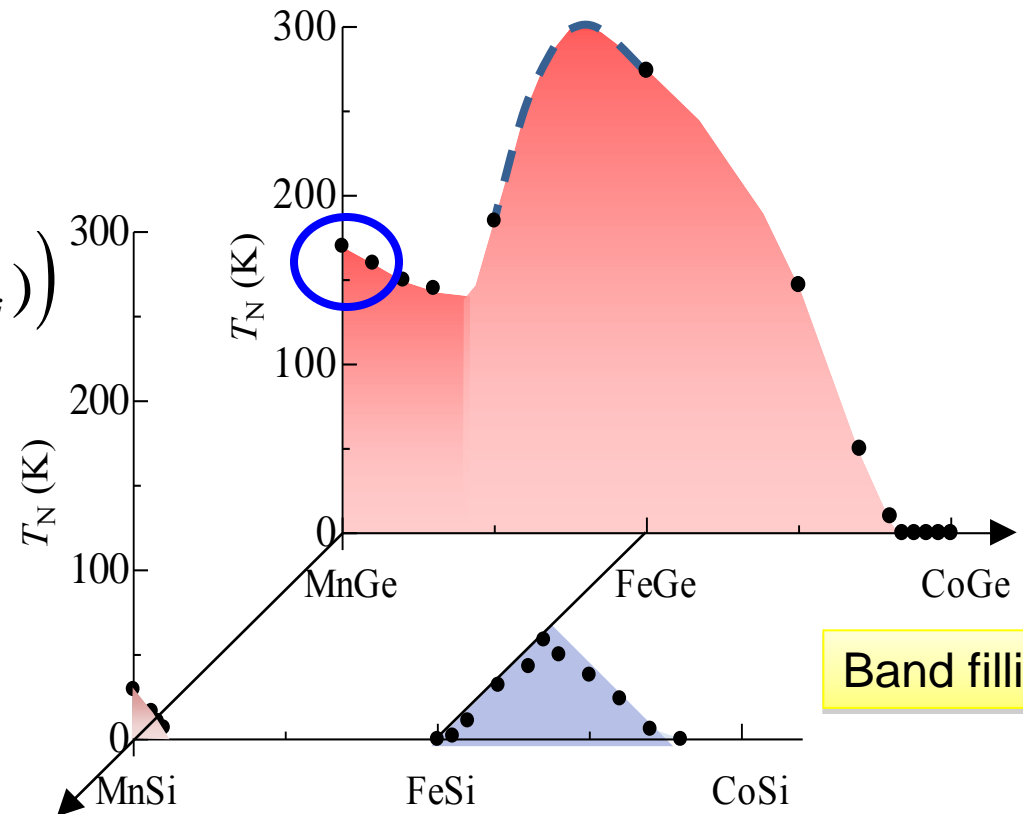
$$H = \sum \left(\underbrace{-J \vec{S}_i \cdot \vec{S}_j}_{\text{Ferro}} + \underbrace{\vec{D}_{ij} \cdot (\vec{S}_i \times \vec{S}_j)}_{\text{DM}} \right)$$

Ferro + DM



Helical spin structure

Long period $\sim aJ/D$
 $\sim 10\text{nm}-300\text{nm}$



Bandwidth

Band filling

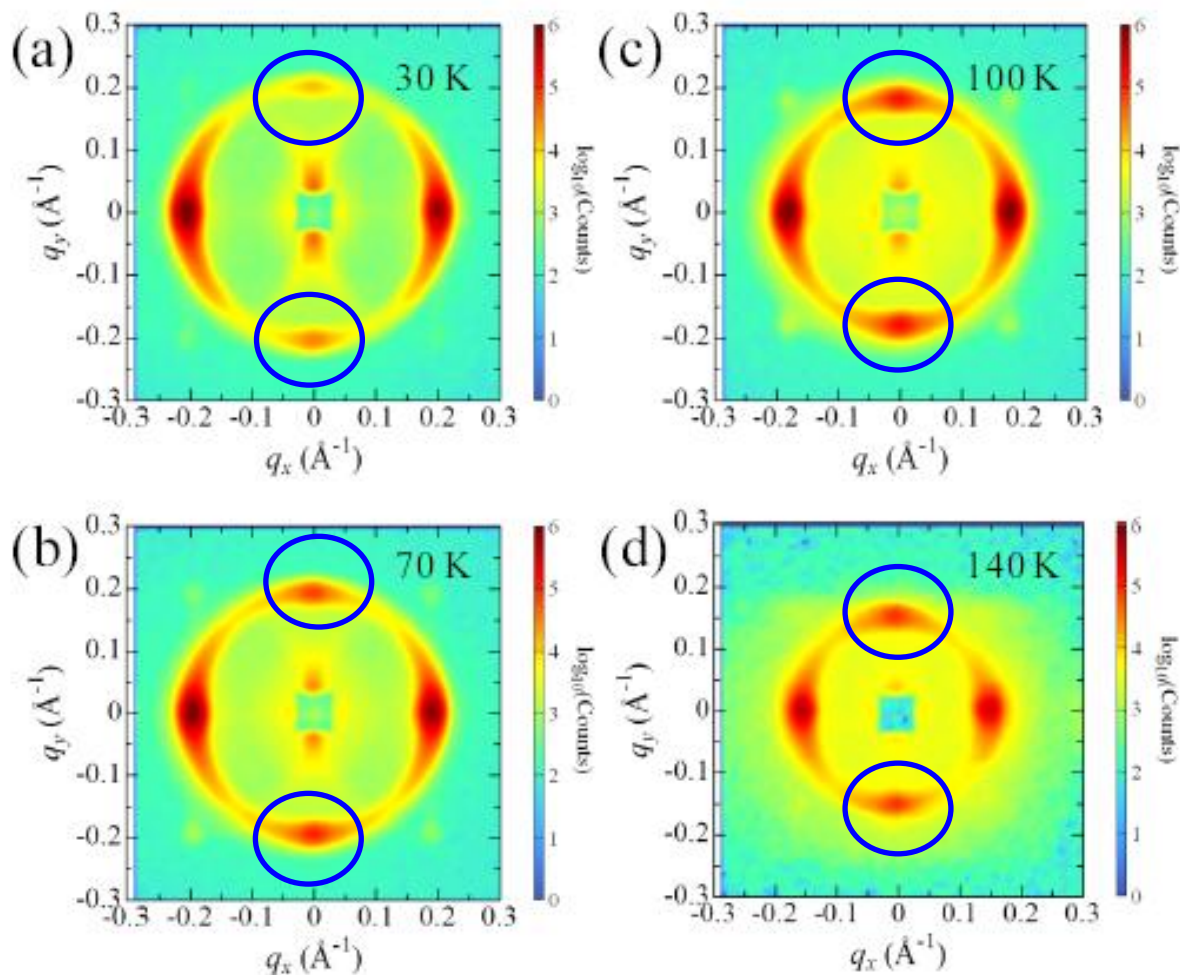
Small angle neutron scattering on MnGe (polyXtal)

B (10T) then B=0



After application of high magnetic field

in collaboration
with Keimer group

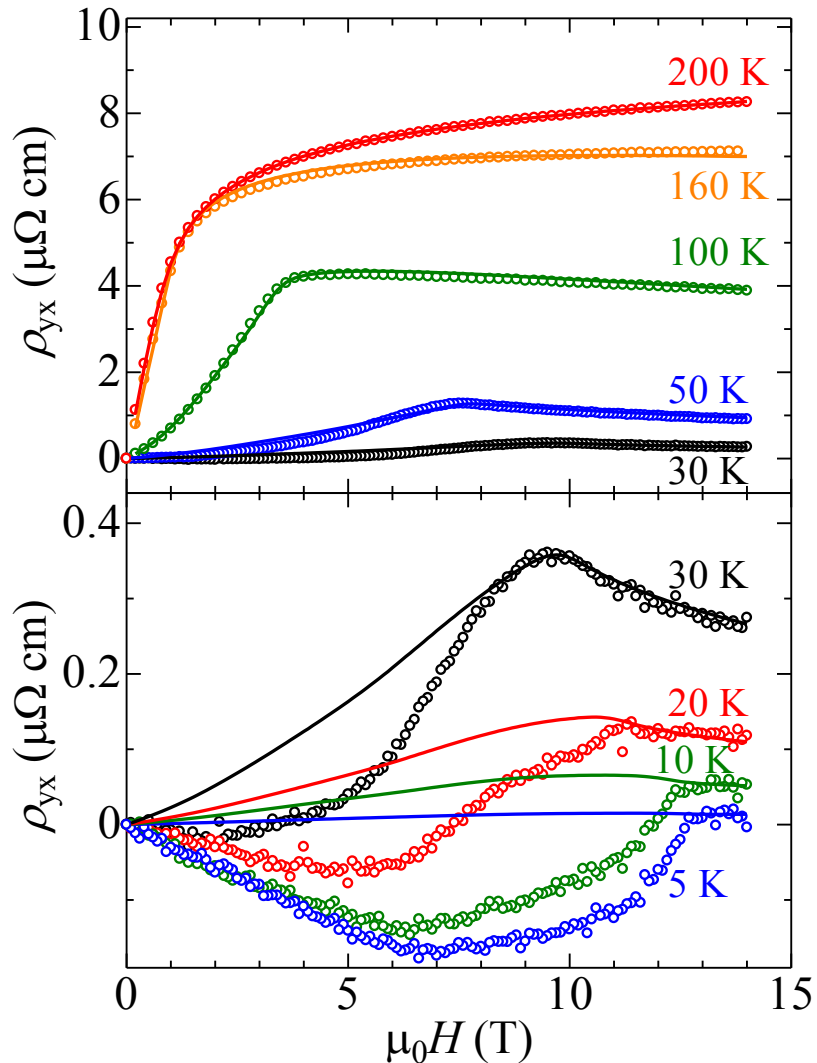


$\lambda = 3 - 6 \text{ nm}$

$q's // \langle 100 \rangle$

Evidence for multiple-q structure even at B=0

Topological Hall effect in MnGe



$H > H_C$

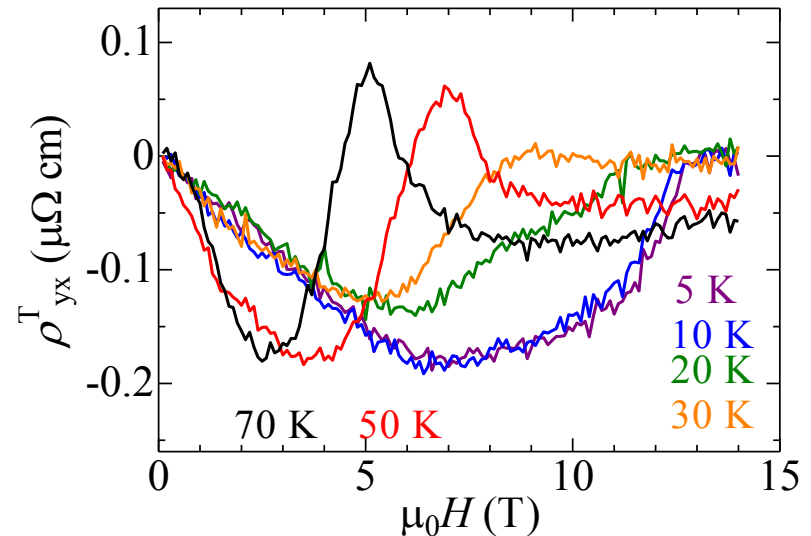
Induced ferromagnetic state
 → “Conventional” anomalous Hall effect

Solid lines: estimate of

$$\rho_{yx}^A = R_0 B_z + \mu_0 R_S M_z$$

$$\mu_0 R_S = S_A \rho_{xx}^2$$

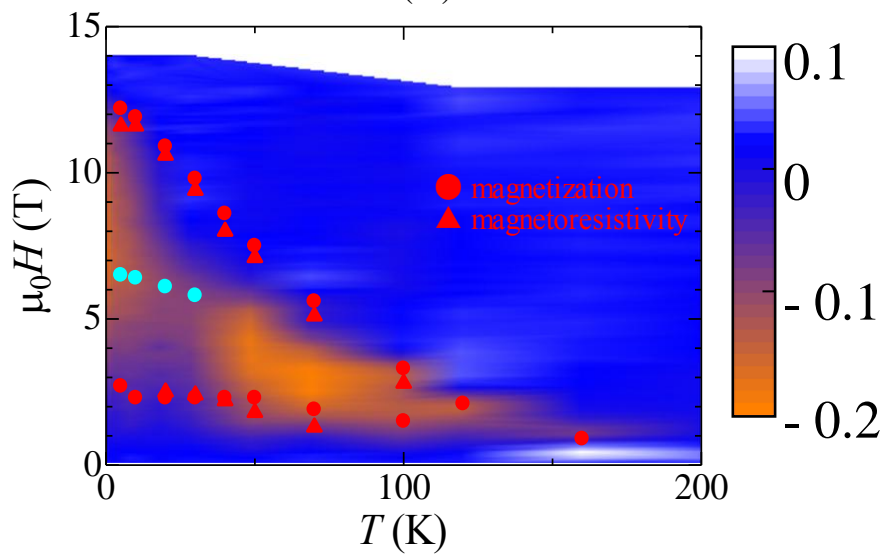
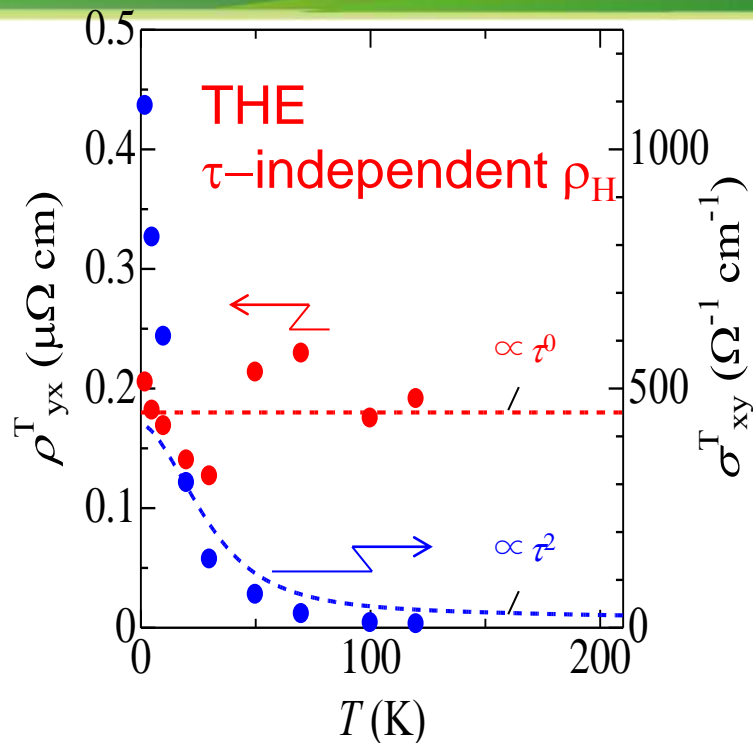
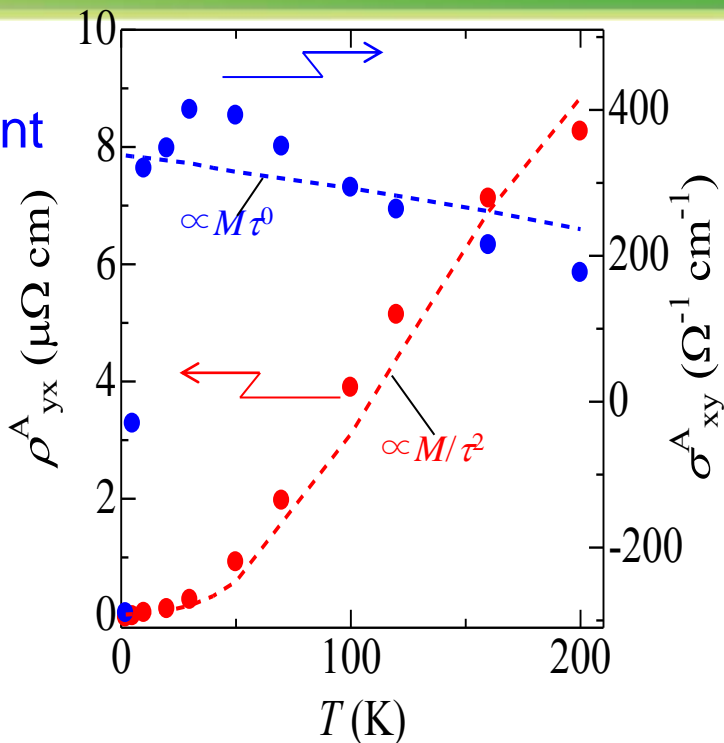
Components of THE



Nearly temperature independent

topological Hall effects via Skyrmion lattice

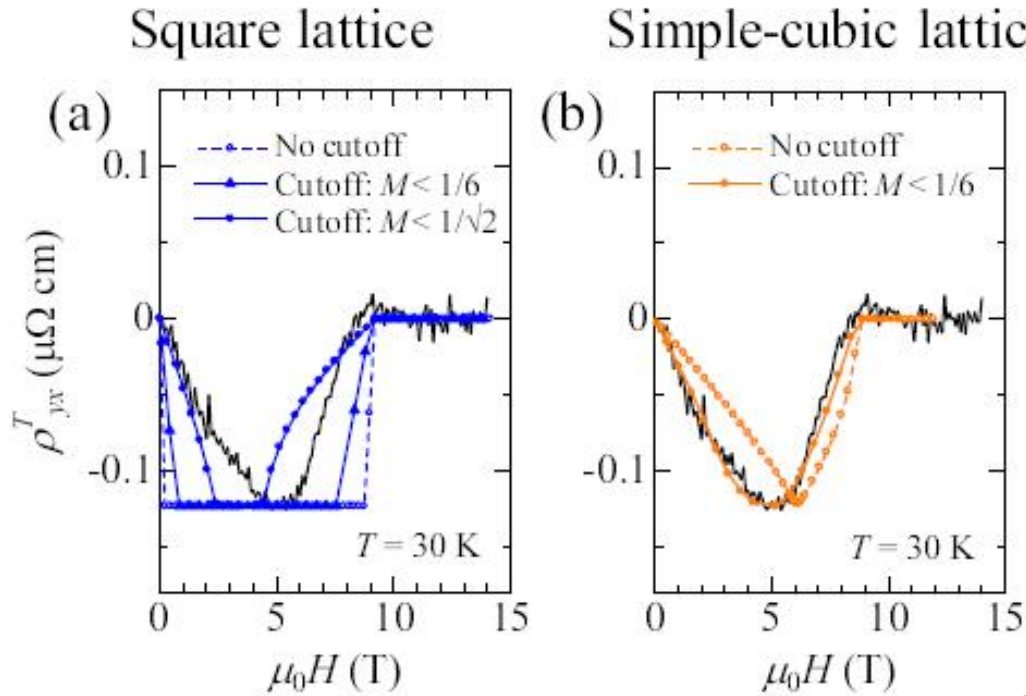
AHE
 σ_H constant



MnGe

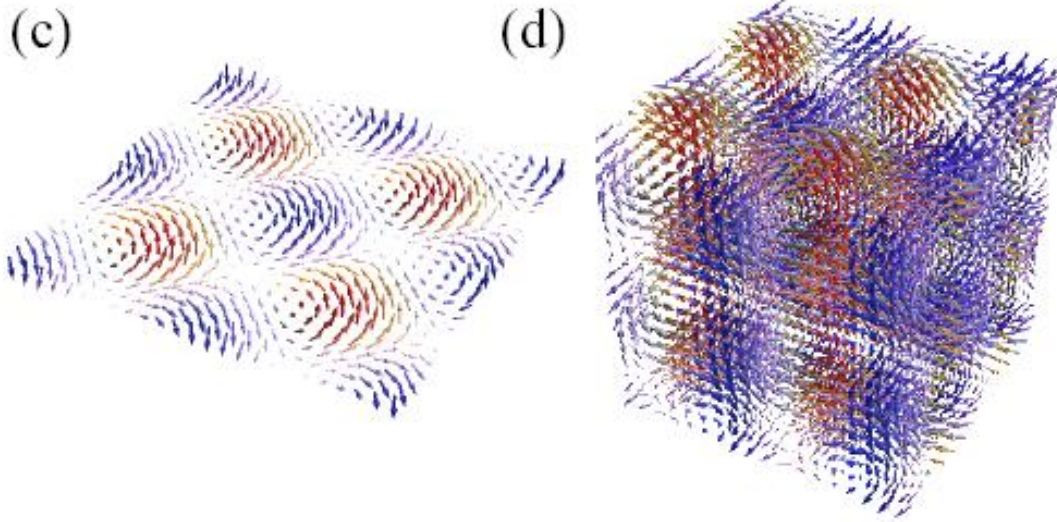
SkX subsistent
 to ground state

Possible 2D (meron) or 3D (hedgehog) Skyrmion Xtal at B=0



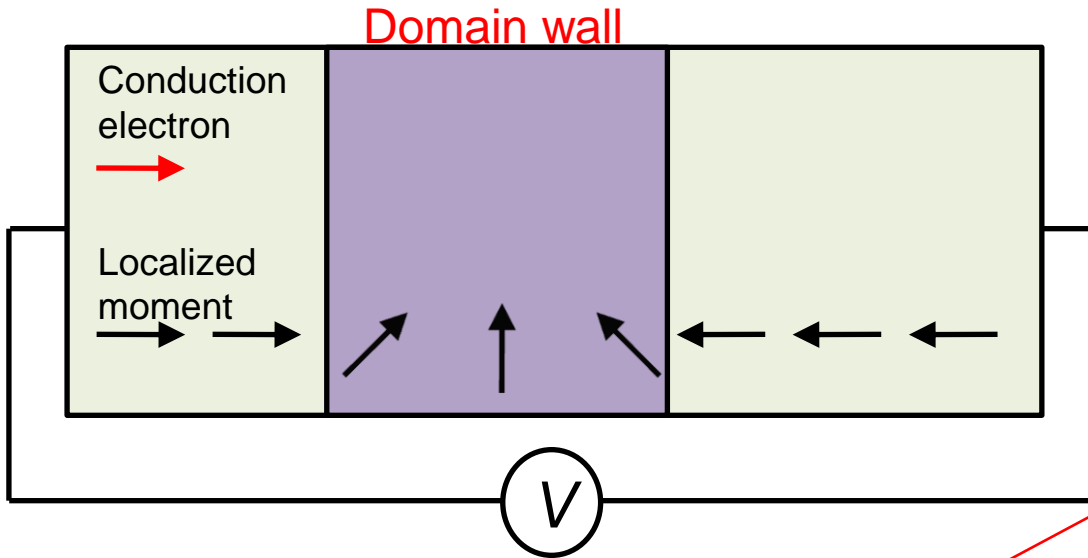
B. Binz and A. Vishwanath,
Physica B 403, 1336 (2008).

skyrmion-antiskyrmion



Current drive of skyrmions and emergent EM field

Domain wall motion by spin transfer torque



Topological Hall effect
Emergent magnetic field
 h

-Emergent electric field

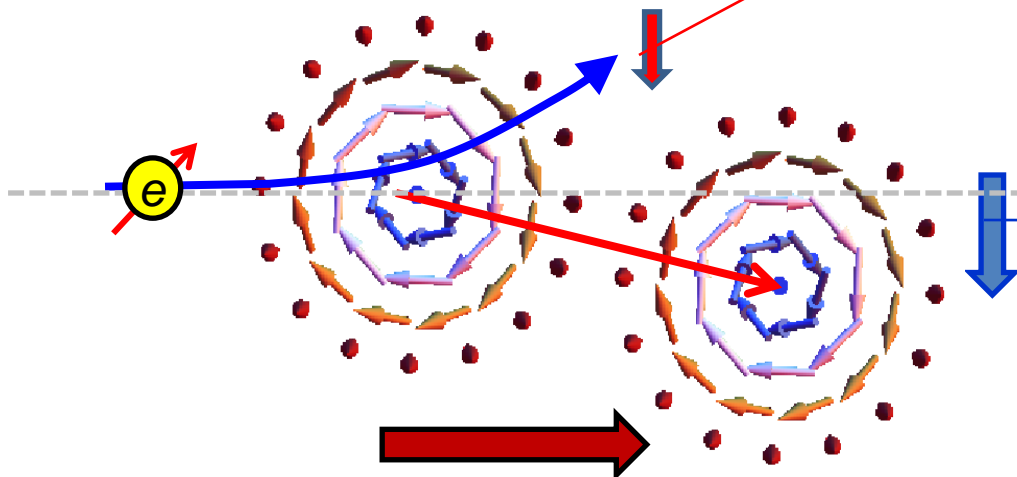
$$\mathbf{e} = -\frac{1}{c} [\mathbf{V}_{\parallel} \times \mathbf{h}]$$

→ appears as reduction of topological Hall effect

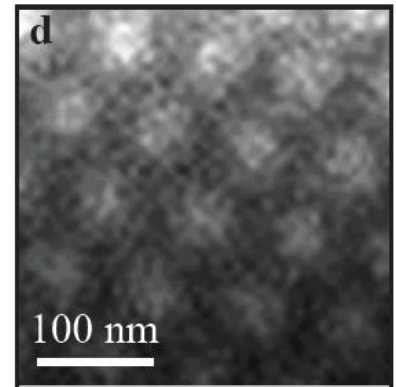
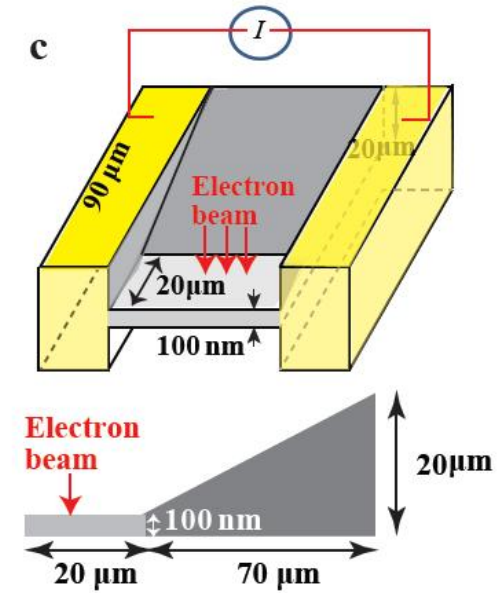
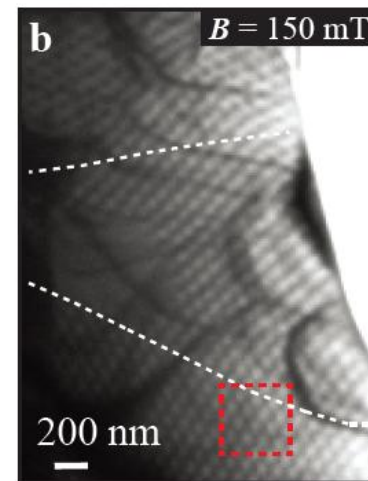
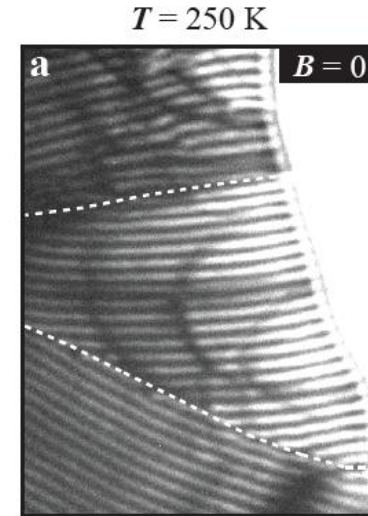
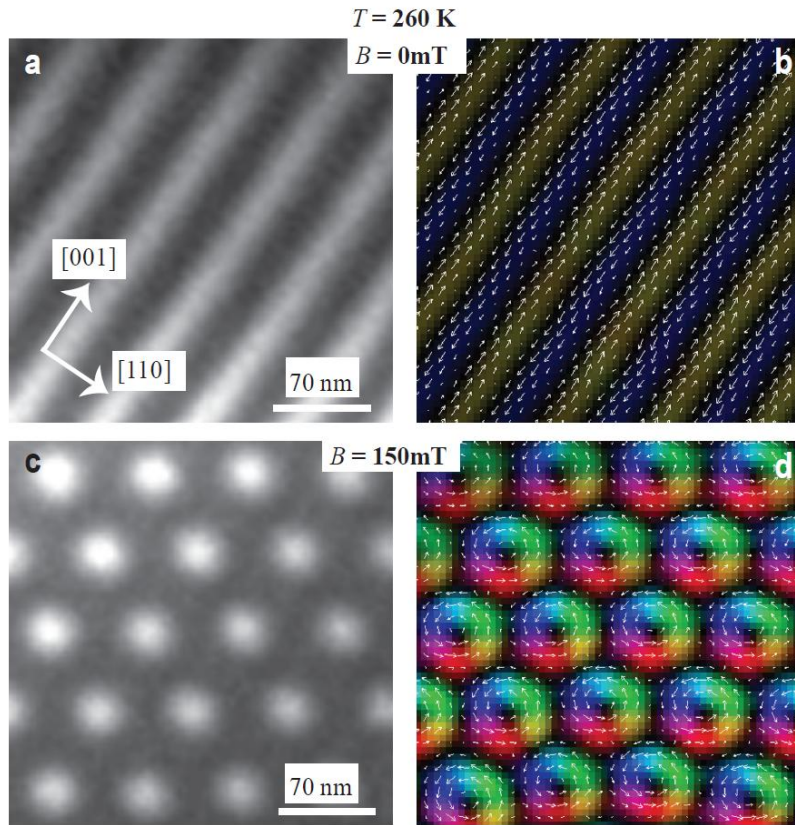
-counteraction of topological Hall effect (THE)

→ skyrmion Hall effect

Skyrmion motion by spin transfer torque

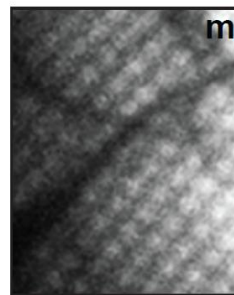
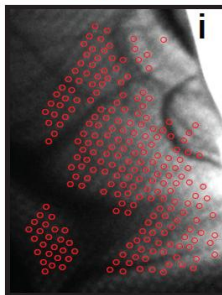
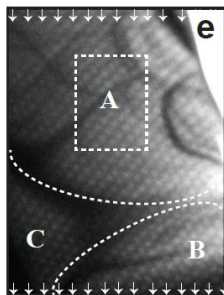
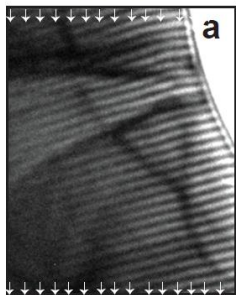


Current driven skyrmion flow in FeGe film

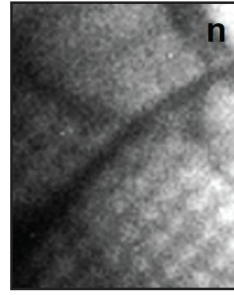
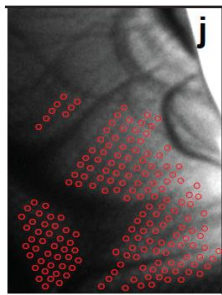
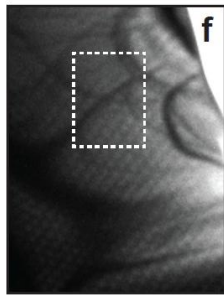
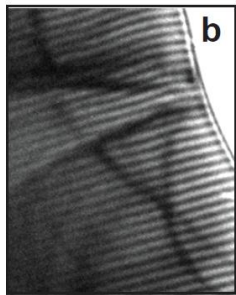


$T = 250 \text{ K}$ $B = 0$

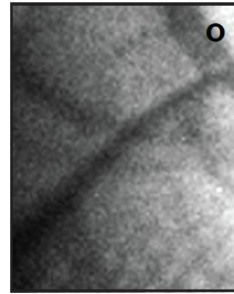
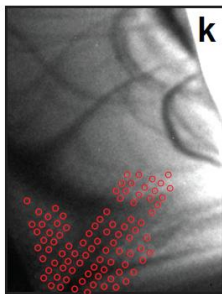
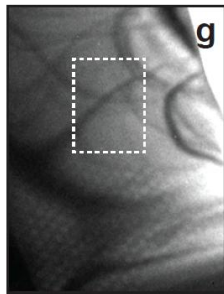
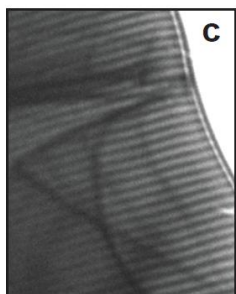
$t = 0$
 $I = 0$



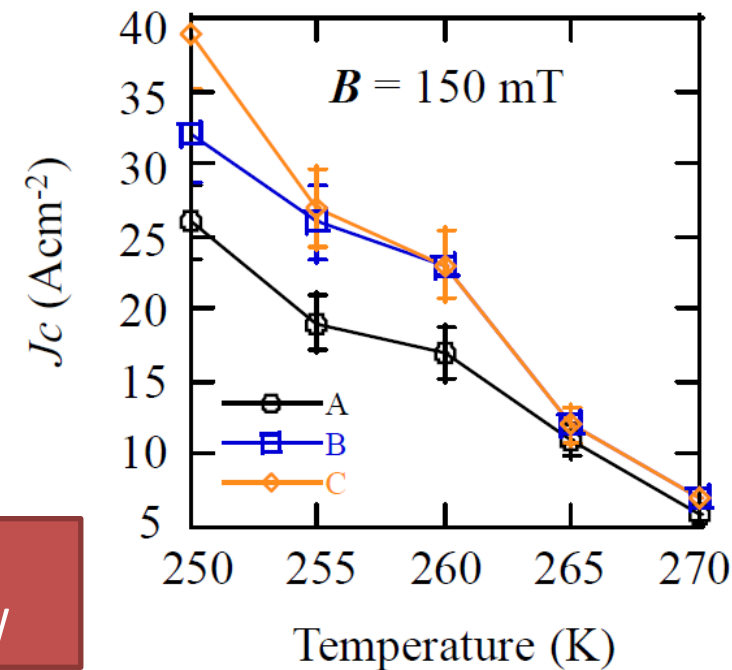
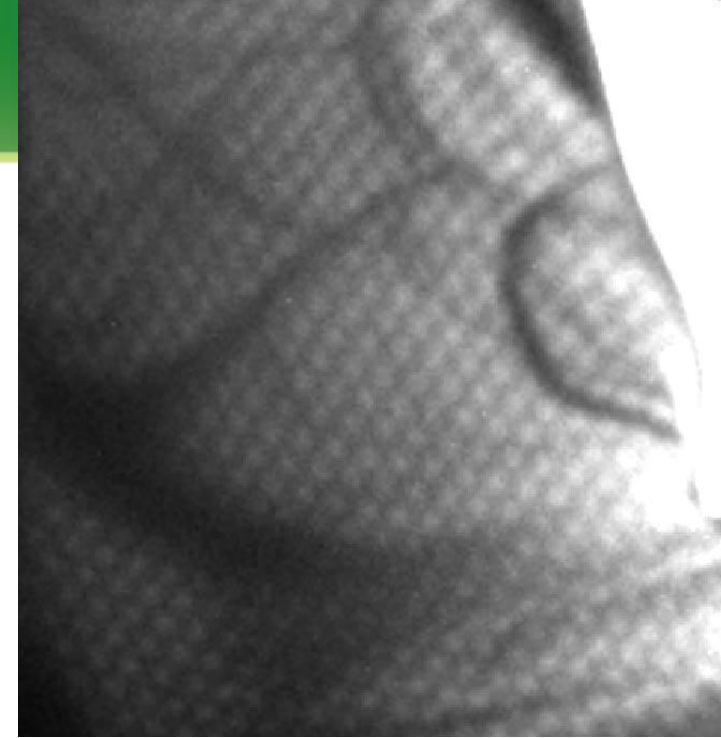
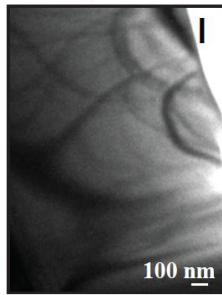
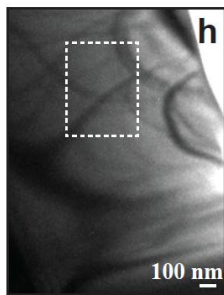
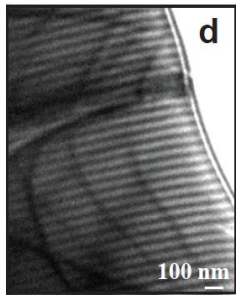
$t = 48 \text{ s}$
 0.41 mA



$t = 52 \text{ s}$
 0.50 mA



$t = 55 \text{ s}$
 0.61 mA



$J_c < 40 \text{ A/cm}^2$!!for skyrmion (not skew)
>10⁵ orders of magnitude smaller than J_c for DW

No pinning effect on skyrmion motion

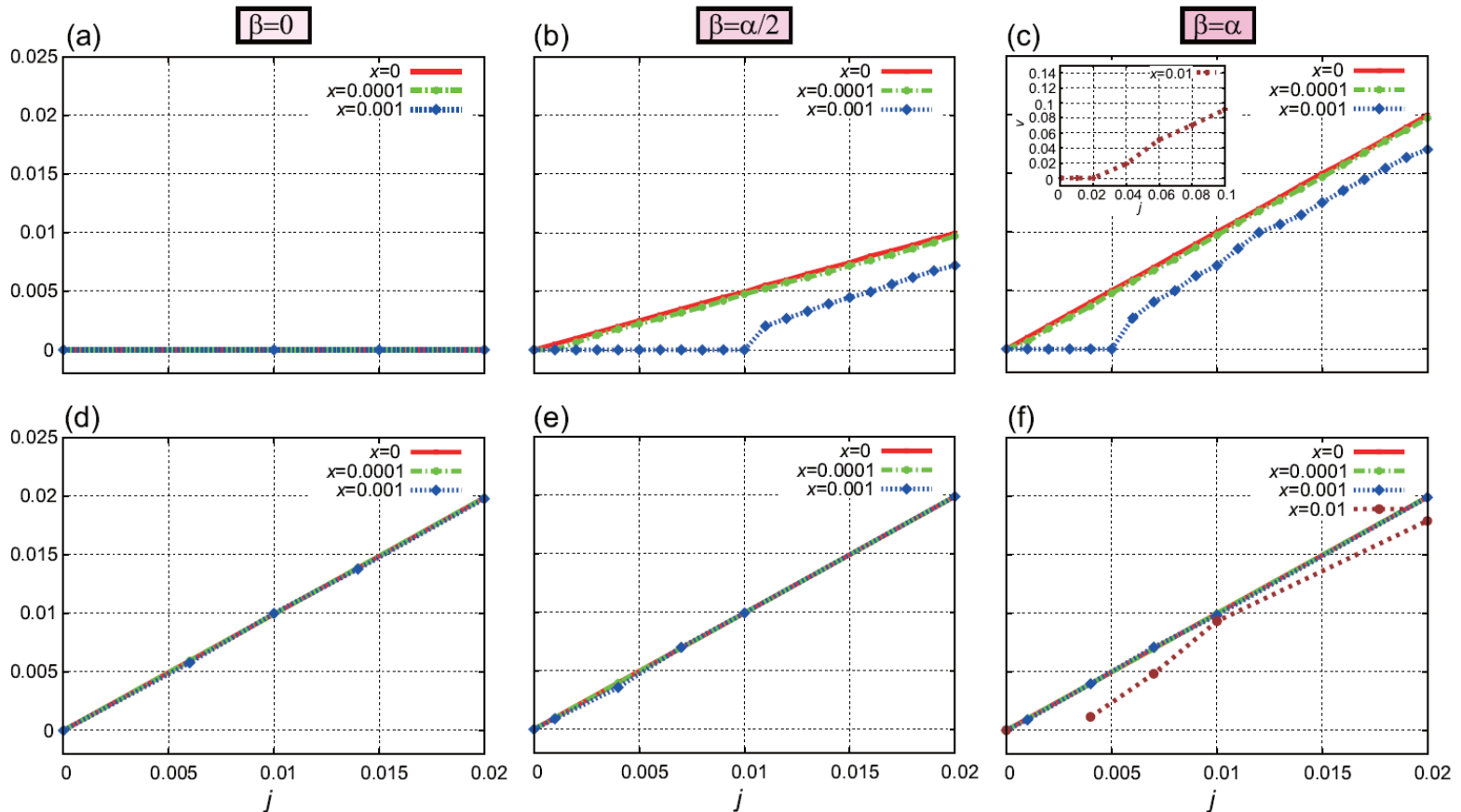
LLG

$$\frac{d\vec{M}_r}{dt} = \gamma \vec{M}_r \times B_r^{\text{eff}} - \frac{\alpha}{M} \vec{M}_r \times \frac{d\vec{M}_r}{dt} - \frac{pa^3}{2eM} (\vec{j} \cdot \vec{\nabla}) \vec{M}_r - \frac{pa^3\beta}{2eM^2} \left[\vec{M}_r \times (\vec{j} \cdot \vec{\nabla}) \vec{M}_r \right],$$

Iwasaki-Mochizuki-Nagaosa (2012)

screw

SS

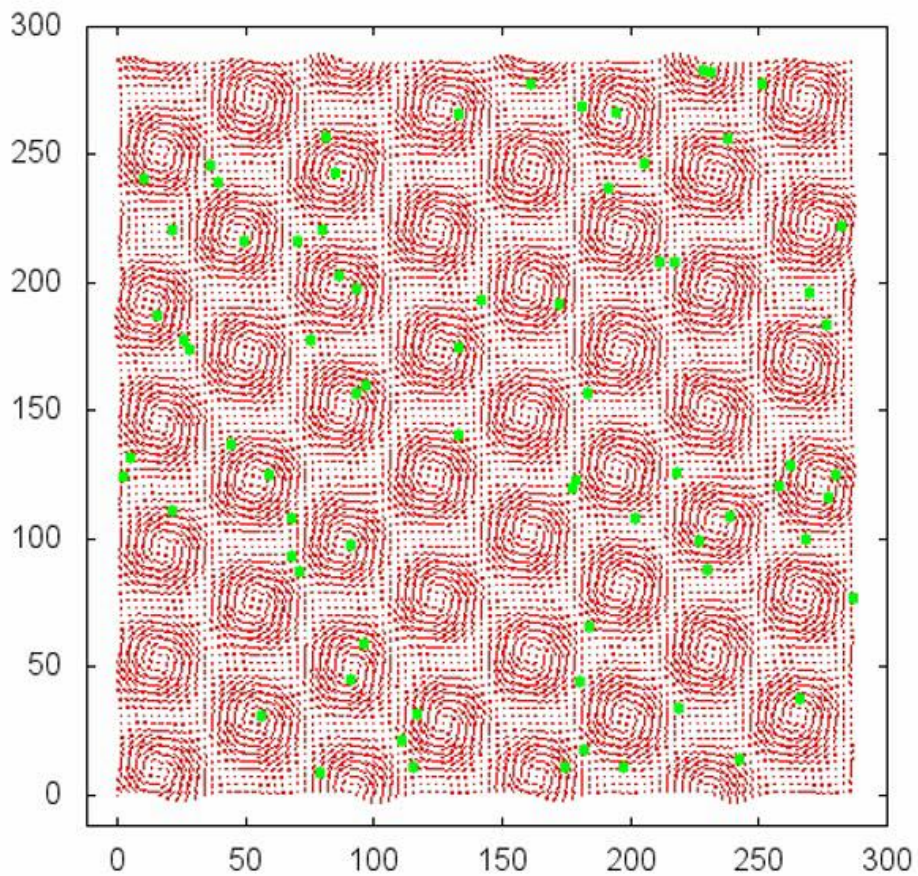


skyrmion

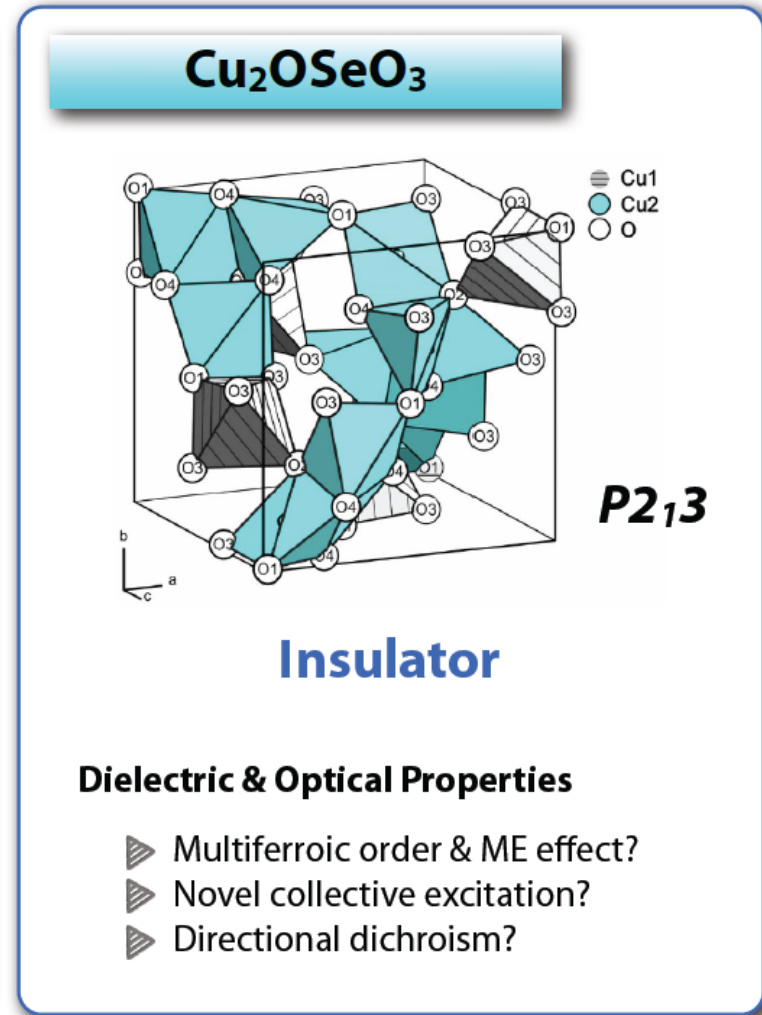
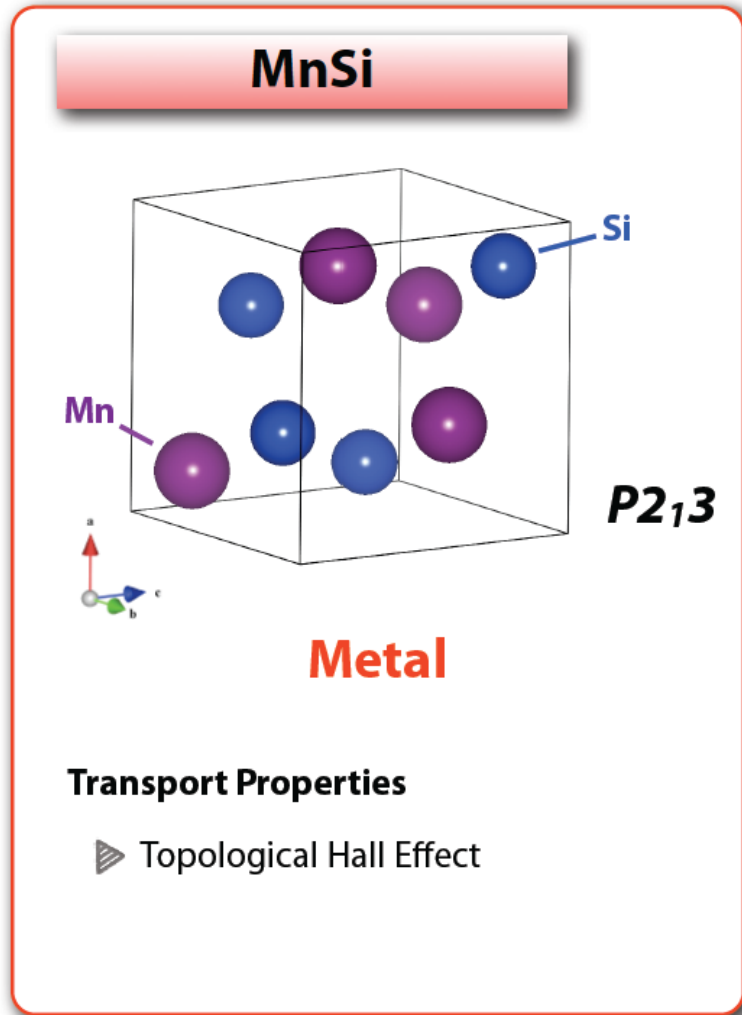
SkX

no (minimal) pinning effect on SkX

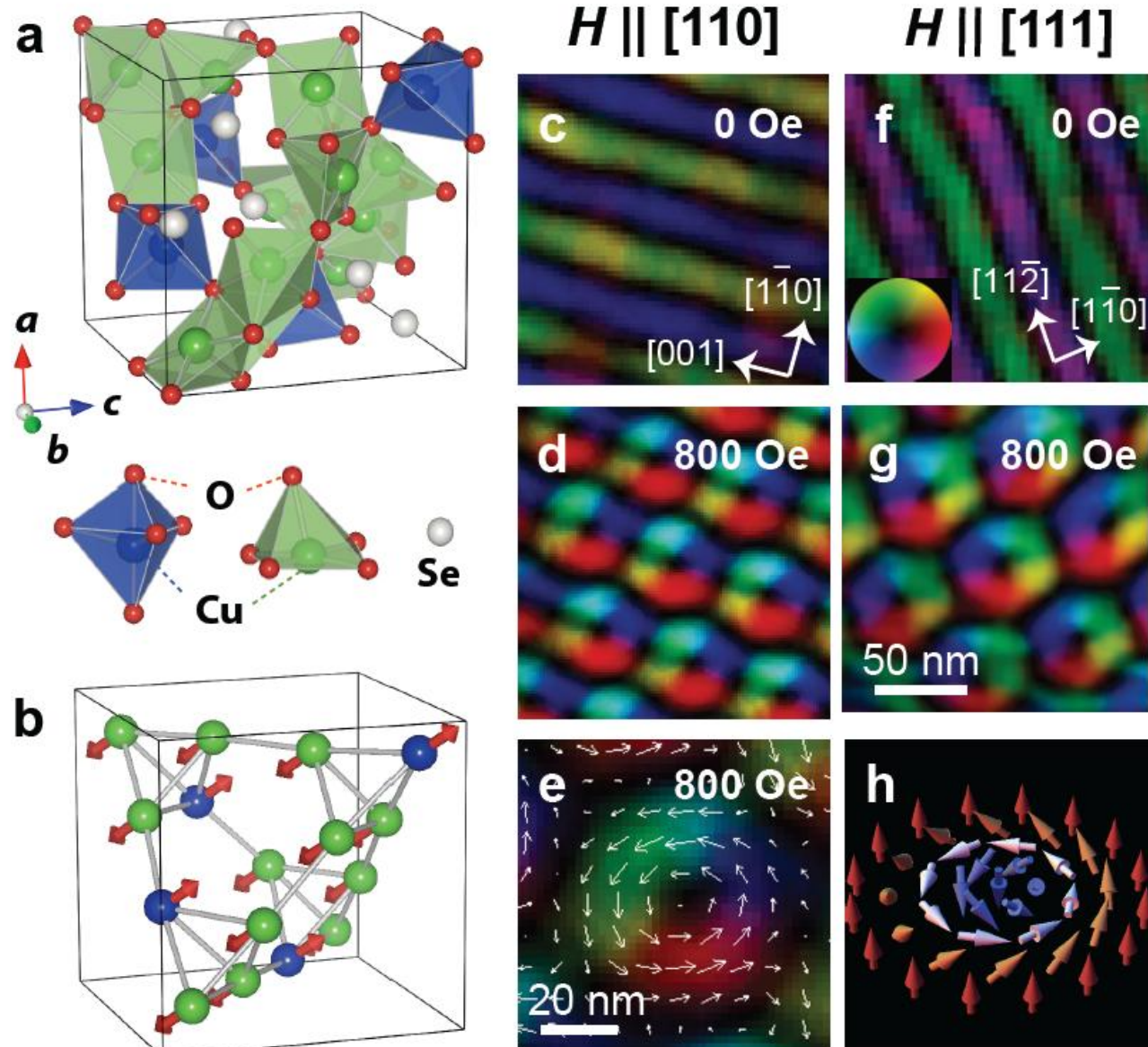
Simulation of current-driven skyrmions under pinning sites



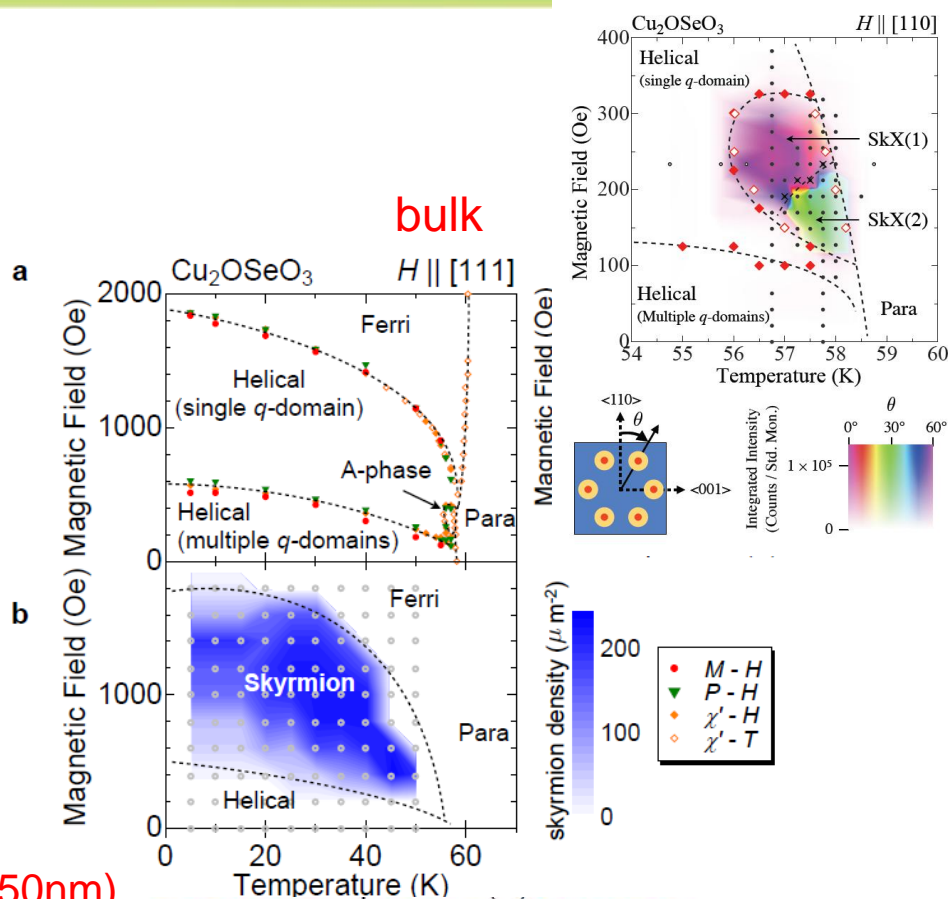
Cu_2OSeO_3 : Chiral Magnetic Insulator



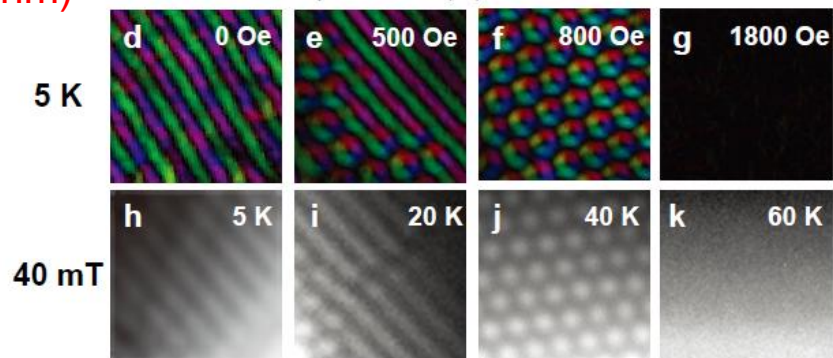
Lorentz TEM observation of thin flake of Cu_2OSeO_3



Skyrmion crystal phase: bulk vs. thin film



film (50nm)

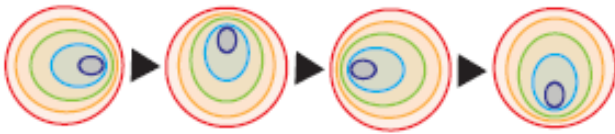


Seki et al. PRB (RC), 85, 220406(2012).

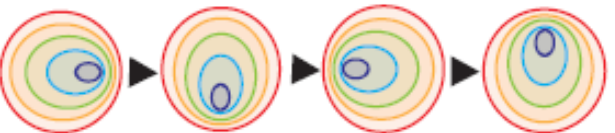
See also Adams et al. PRL, 108, 237204 (2012).

$H//x,y$ on skyrmions

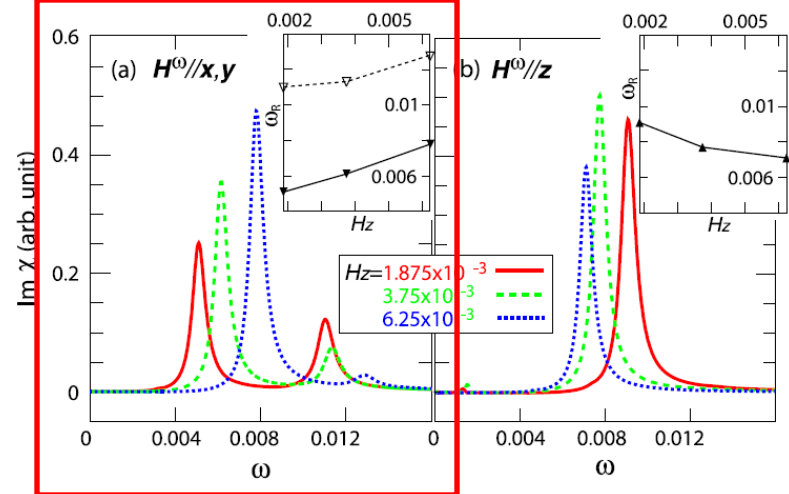
d Counter Clock Wise mode



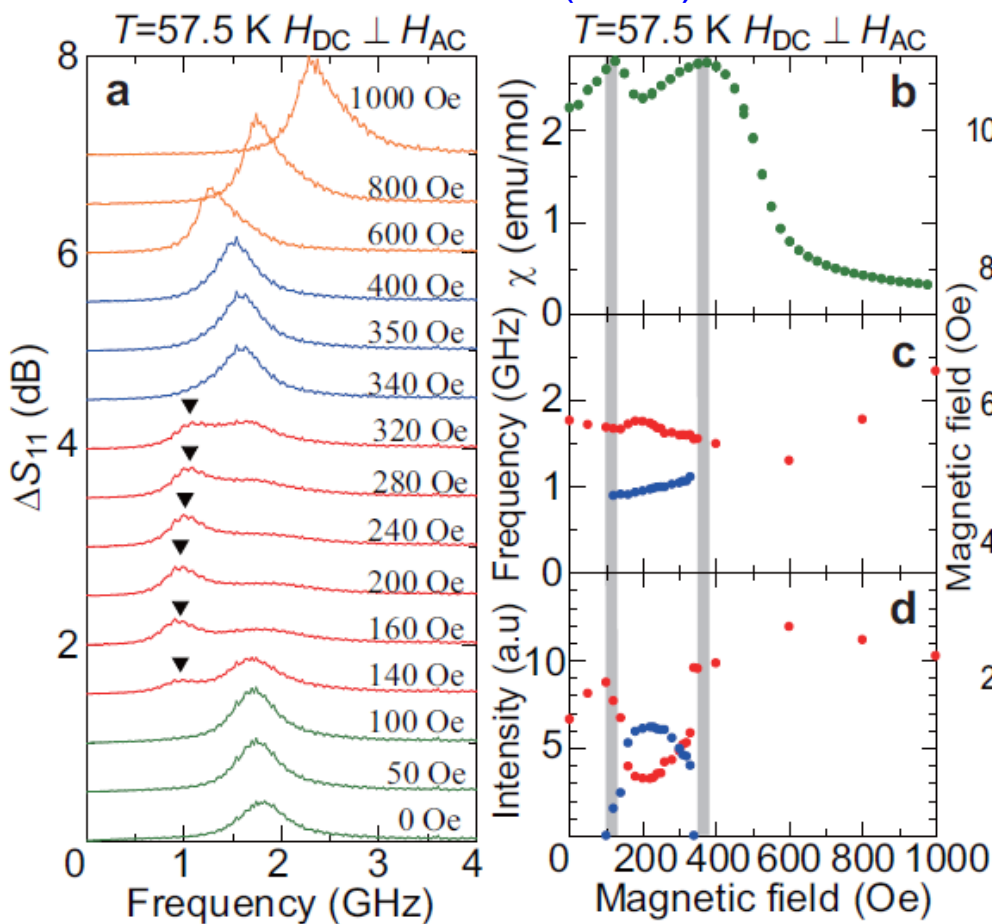
e Clock Wise mode



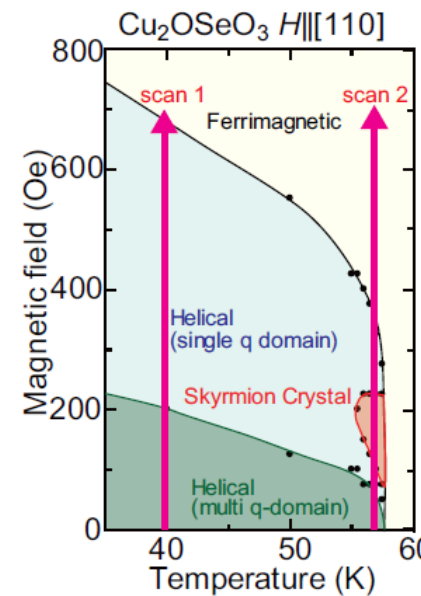
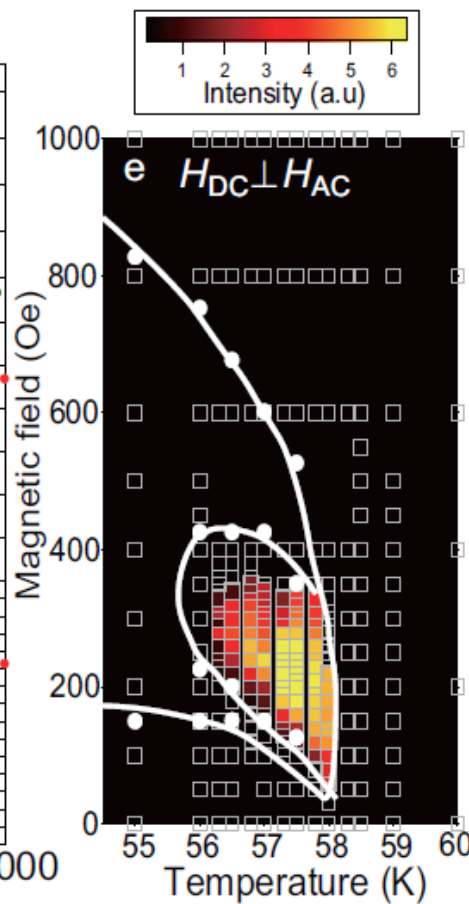
$H^\omega//xy$



Onose et al. PRL (2012)



Mochizuki PRL (2012)

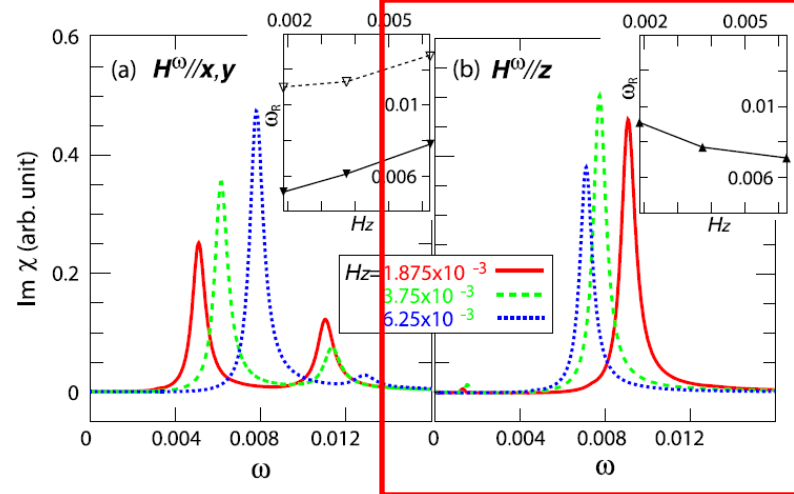
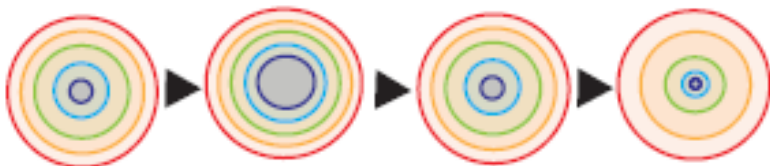


H//z on skyrmions

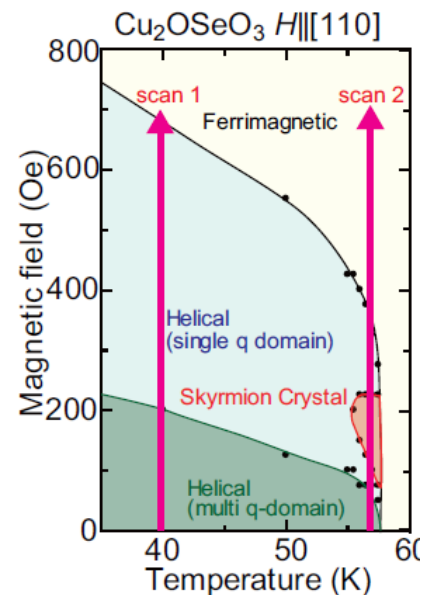
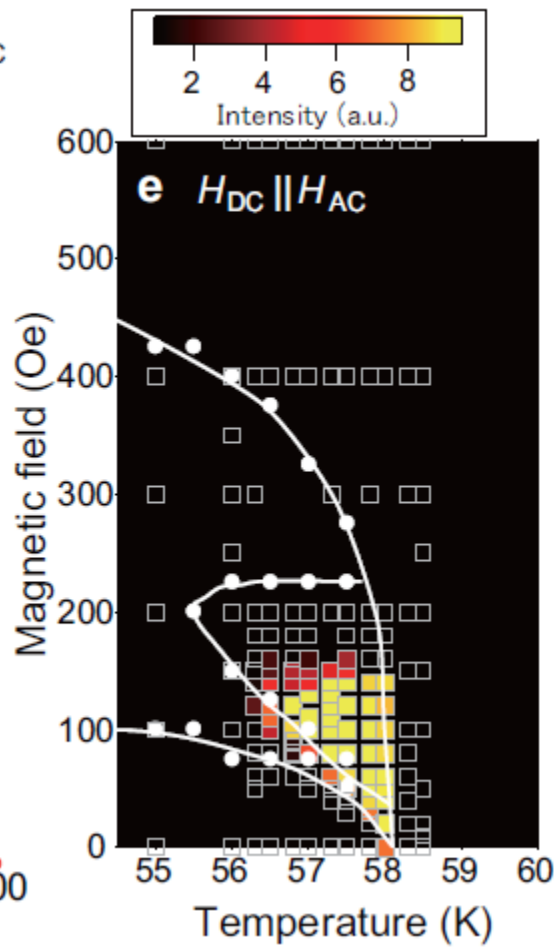
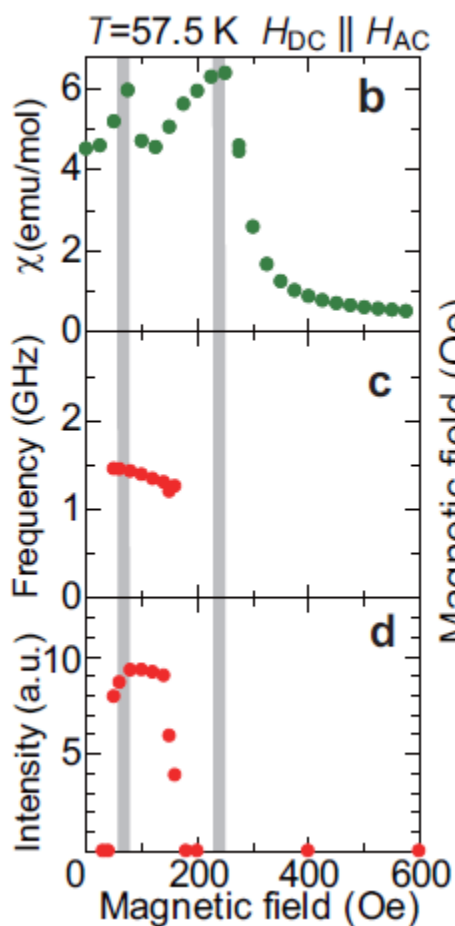
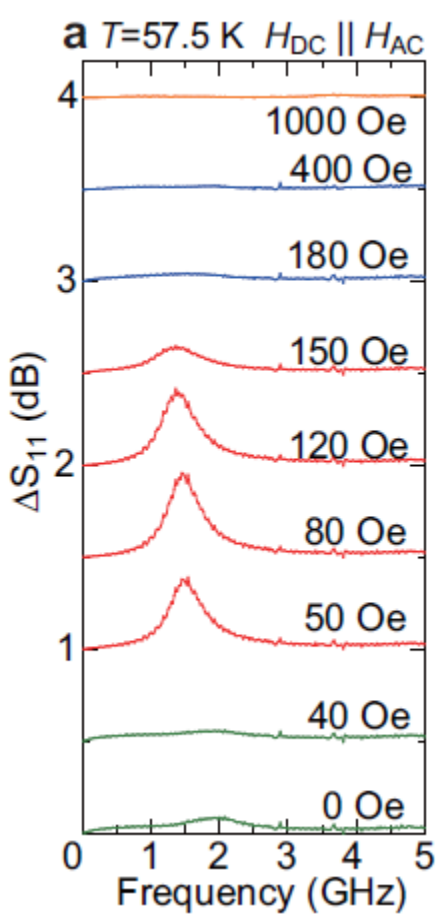
Onose et al. PRL (2012)

f Breathing mode

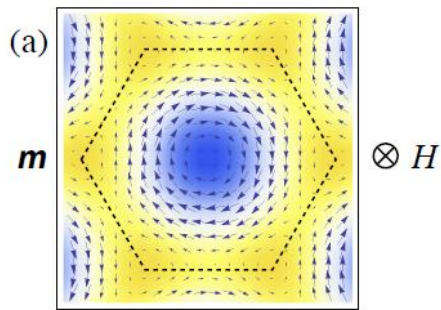
$H^\omega // c$



Mochizuki PRL (2012)

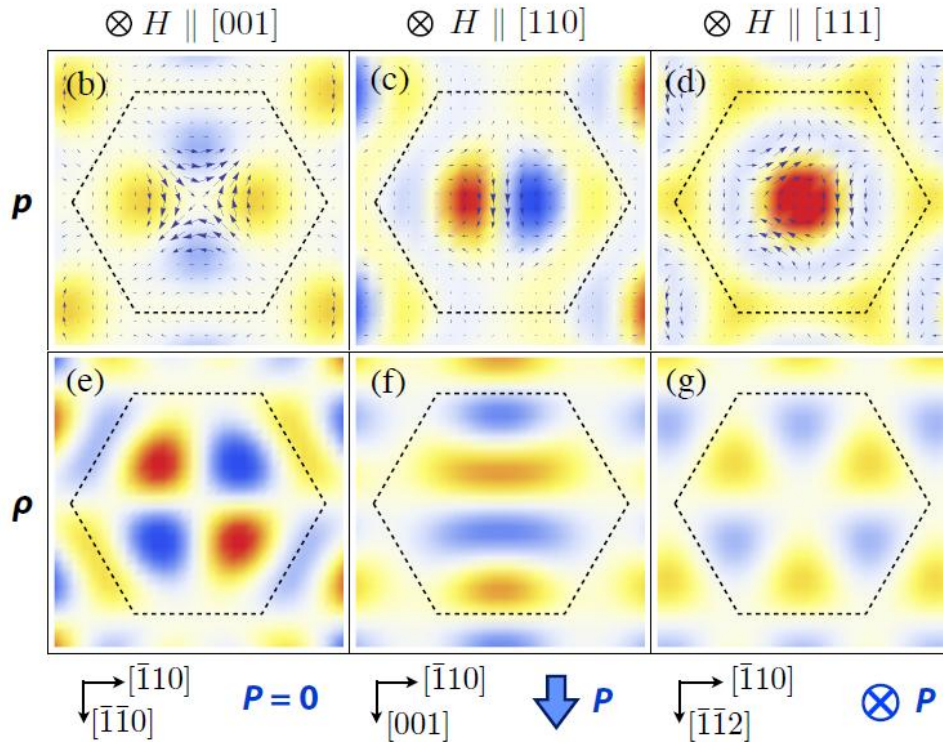
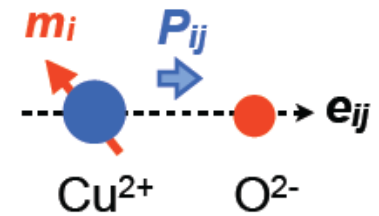


Cu₂OSeO₃: P and ρ distributions in skyrmion



d - p hybridization model

$$\vec{p}_{ij} \propto (\vec{e}_{ij} \cdot \vec{m}_i)^2 \vec{e}_{ij}$$



quadrupole

dipole

polarization P



E -field drive

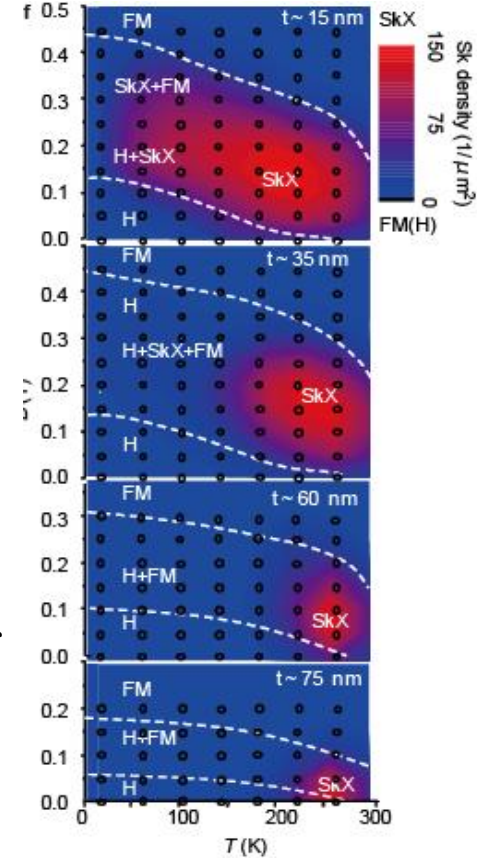
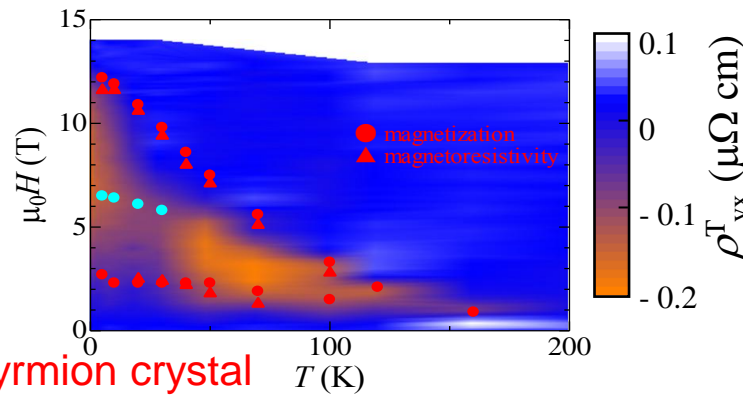
charge ρ

Stabilization Skyrmions and Skyrmion Xtal in form of thin films

Topological Hall effect as probe for SkX

emergent EM fields

MnGe
zero-field skyrmion crystal



Skyrmion transport and dynamics

toward skyrmionics

- ▶ current-drive of Skyrmions ($\sim 10 \text{ A/cm}^2$)
- ▶ multiferroic skyrmions, E-drive/optical-control of Skyrmion
- ▶ spontaneous ratchet motion in thermal equilibrium