

# Magnetism of a Frustrated Four-Spin-Tube

M. Arlego, WB, Phys. Rev. B **84**, 134426 (2011)  
M. Arlego, et al. preprint



Marcelo Arlego  
Diego Rosales  
Gerardo Rossini

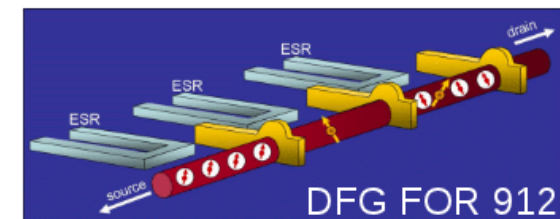
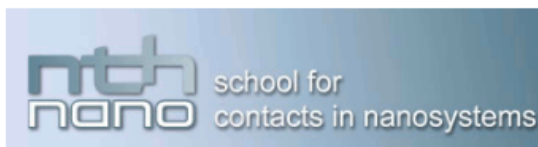
UNLP, La Plata, Argentina



Youssef Rahnavard  
Björn Willenberg

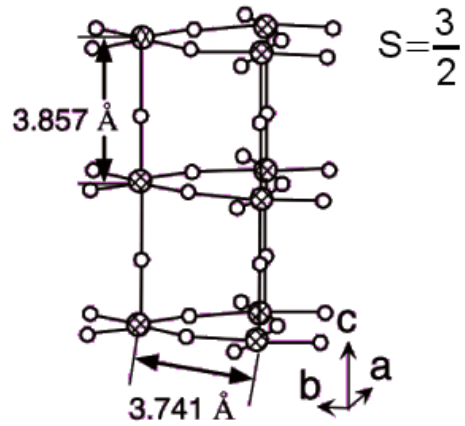
TUBS, Braunschweig Germany

## Funding

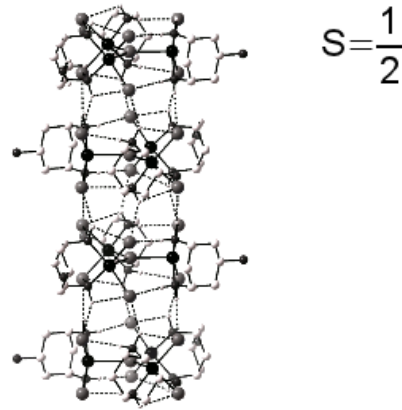
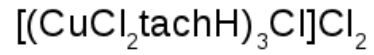


# Spin Tubes

 N=3



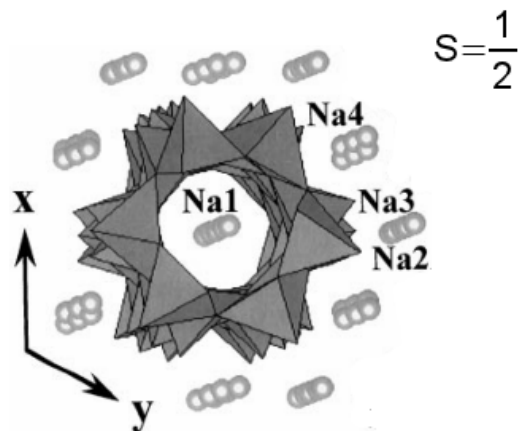
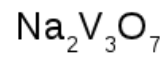
Manaka, et al.,  
JPSJ 78, 93701 (09)



Schnack, et al.,  
PRB 70, 17442 (04)

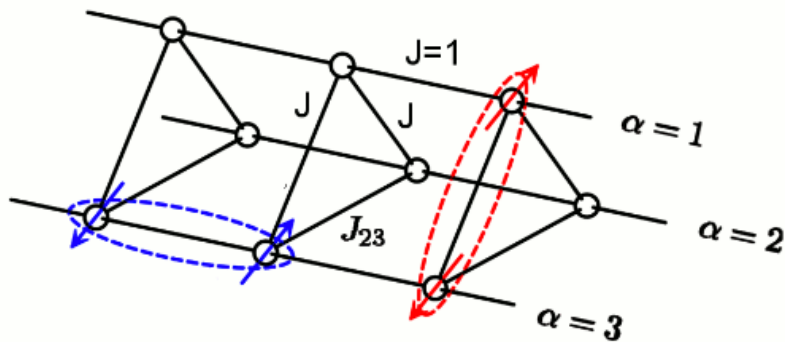
Geometric frustration

 N=9



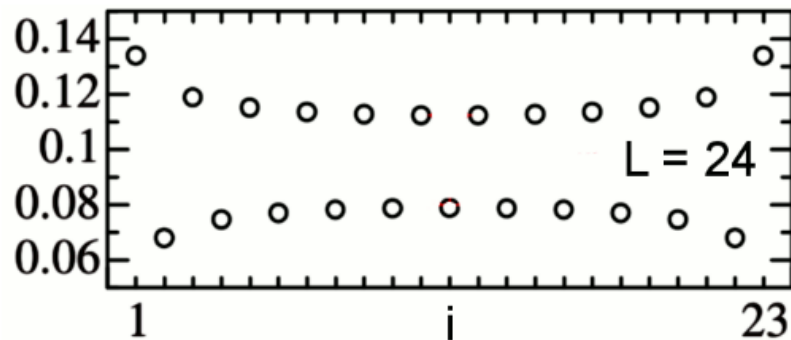
Millet, et al.,  
JSSC 147, 676 (99)

• Three leg tube

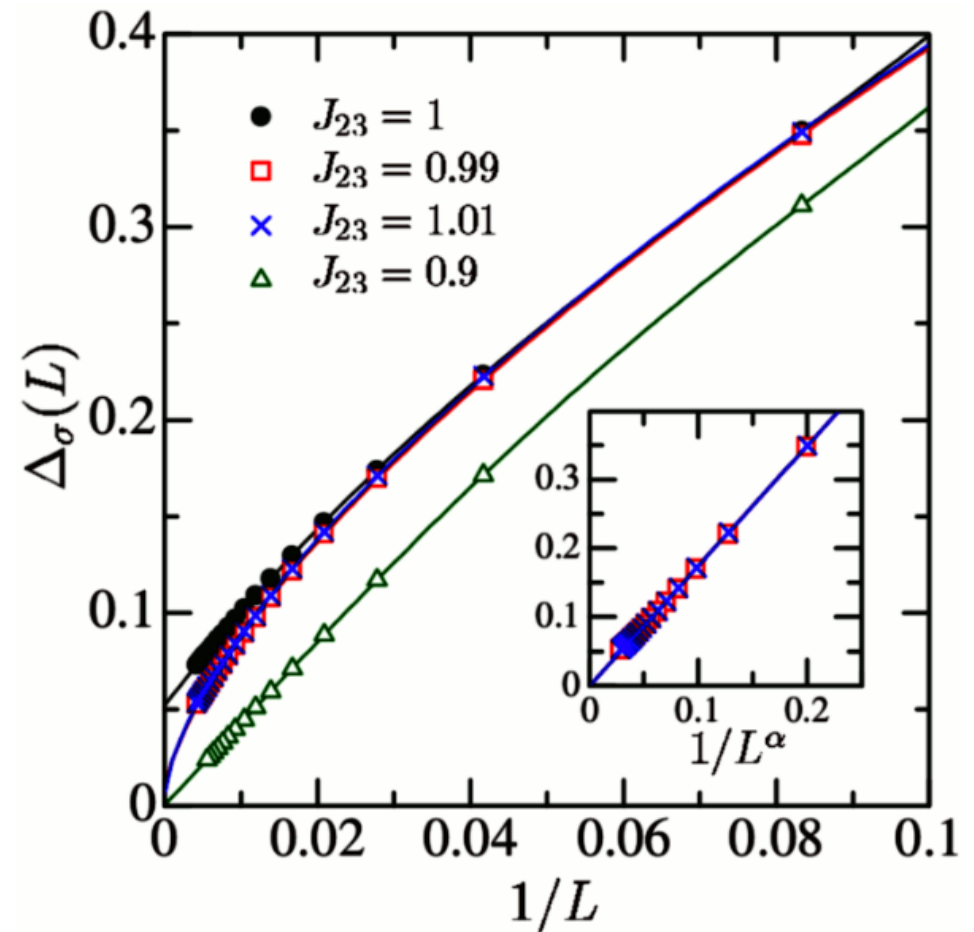


• consistent with LSM theorem

$$\langle \vec{S}_{\alpha,i} \cdot \vec{S}_{\alpha,i+1} \rangle$$



• ... gapped at isotropic point

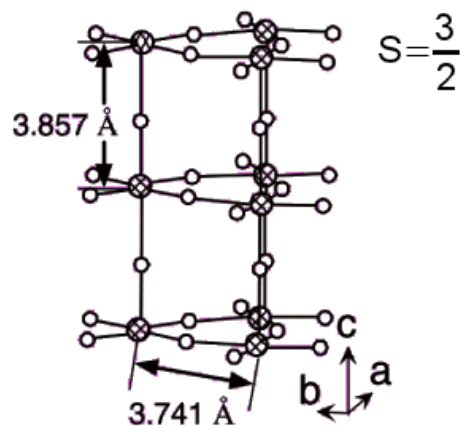


Kawano, Takahashi, JPSJ 66, 4001 (97)  
 Sakai, et al., Physica E 29, 633 (05)  
 Nishimoto, Arikawa, PRB 78, 054421 (08)

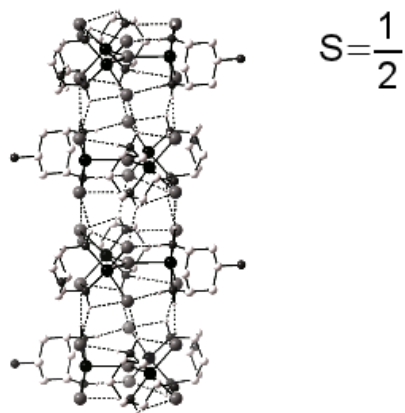
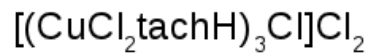
Sakai, et al., JPhys CondMat 22, 403201 (10)

# Spin Tubes

**N=3**

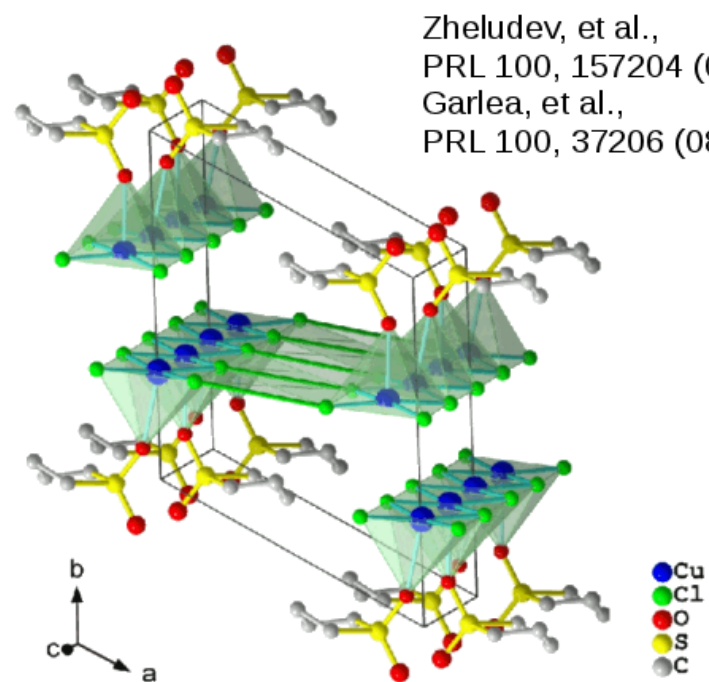
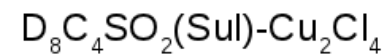


Manaka, et al.,  
JPSJ 78, 93701 (09)



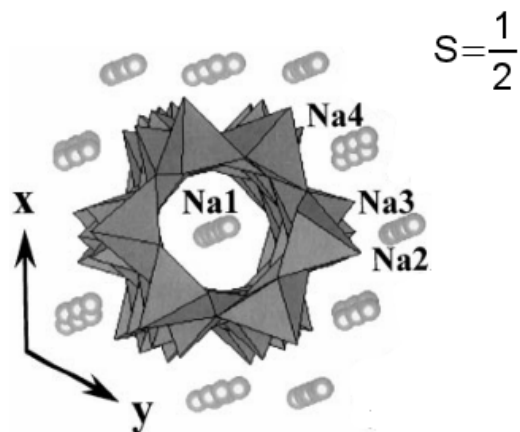
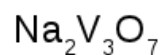
Schnack, et al.,  
PRB 70, 17442 (04)

**N=4**



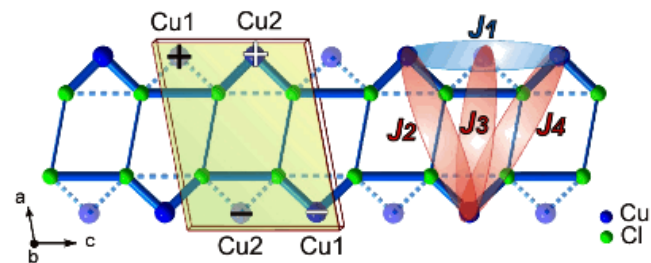
Zheludev, et al.,  
PRL 100, 157204 (08)  
Garlea, et al.,  
PRL 100, 37206 (08)

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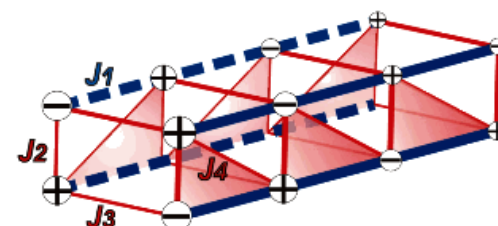


Millet, et al.,  
JSSC 147, 676 (99)

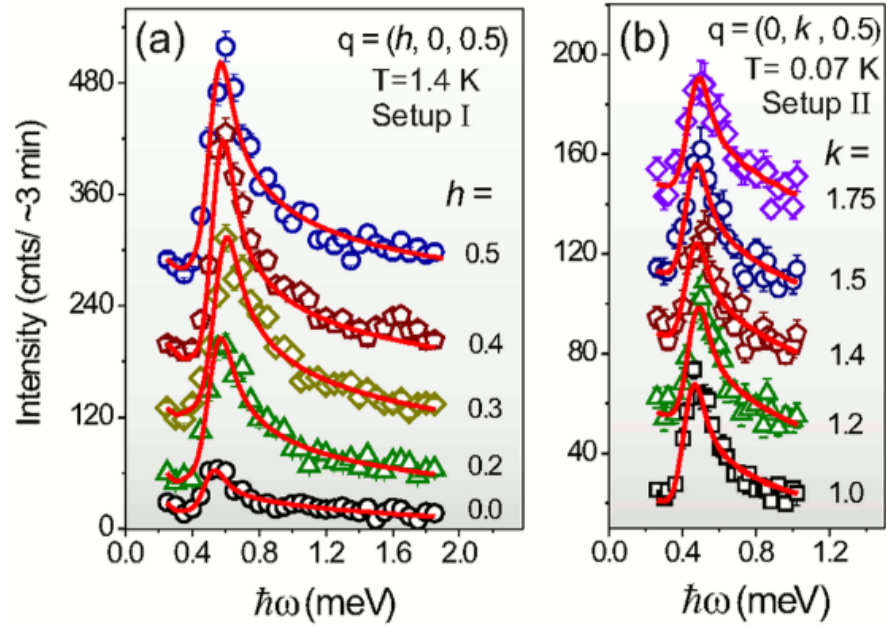
Frustration



$S = \frac{1}{2}$

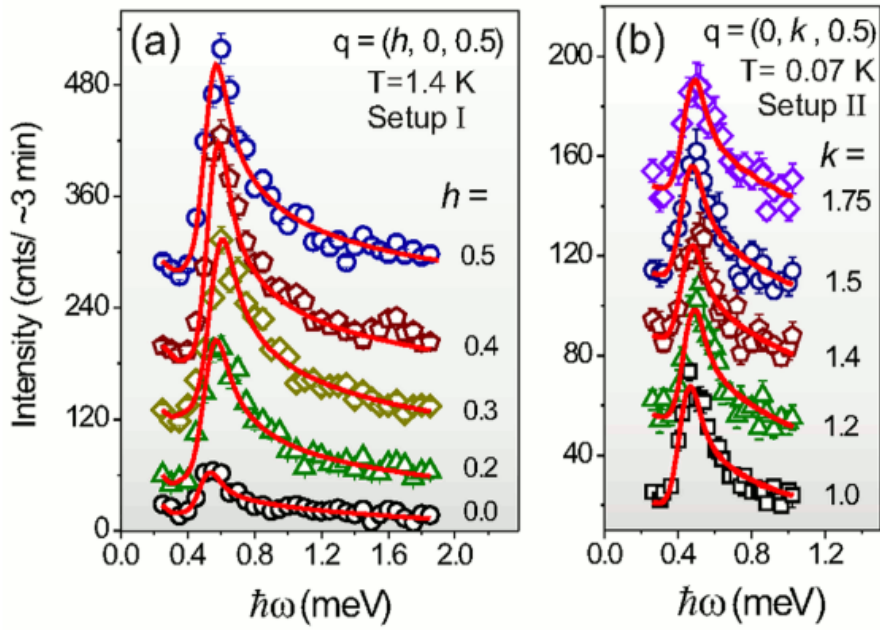


INS: strongly 1D

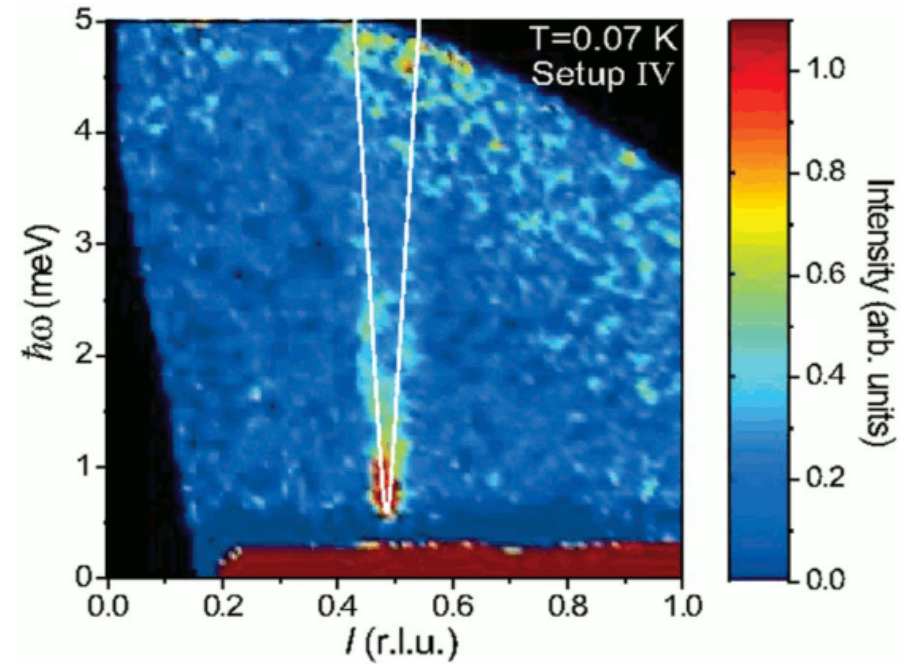


Sul-Cu<sub>2</sub>Cl<sub>4</sub>

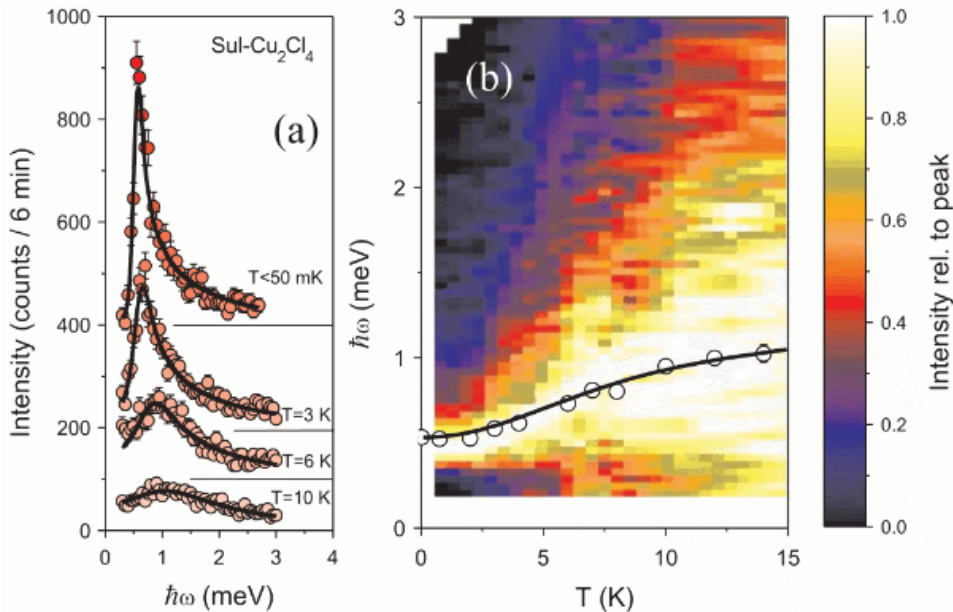
INS: strongly 1D



INS: c-axis dispersion



Const. q-scan vs. T @ Q



gapped

incommensurate  $\rightarrow$  frustration

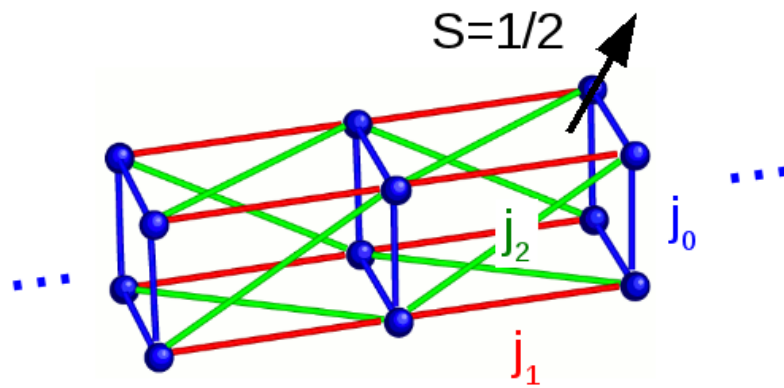
$$Q = \pi(1 - 0.044)$$

velocity

$$\partial \epsilon / \partial k \sim 100 \text{ K} \sim J_{\text{avg}}$$

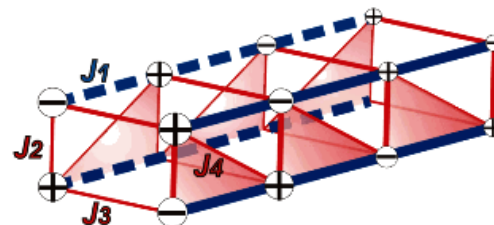
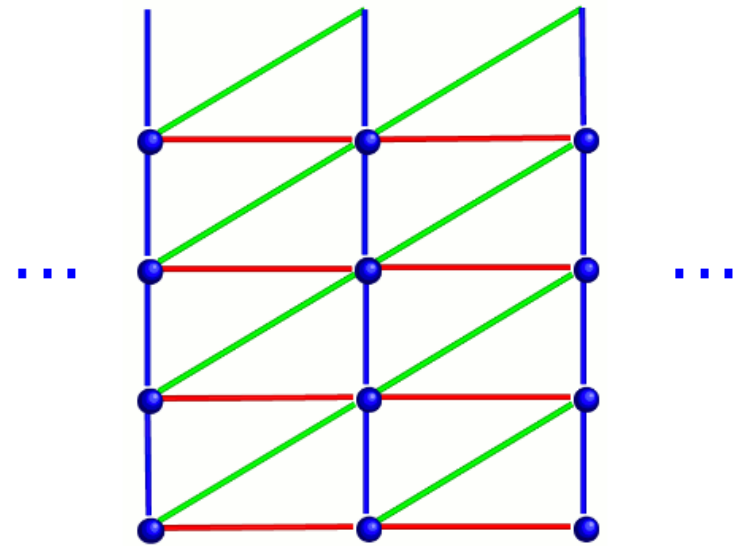
# Tube Model

- Geometrically frustrated four-spin tube



$$H = \sum_{\langle lm \rangle} J_{lm} \vec{S}_l \cdot \vec{S}_m \quad \text{all } j_i \text{ AFM}$$

- Anisotropic triangular lattice on a torus



1) Tube Materials & Model

2) Phases and Static Correlations

3) Plaquette Regime

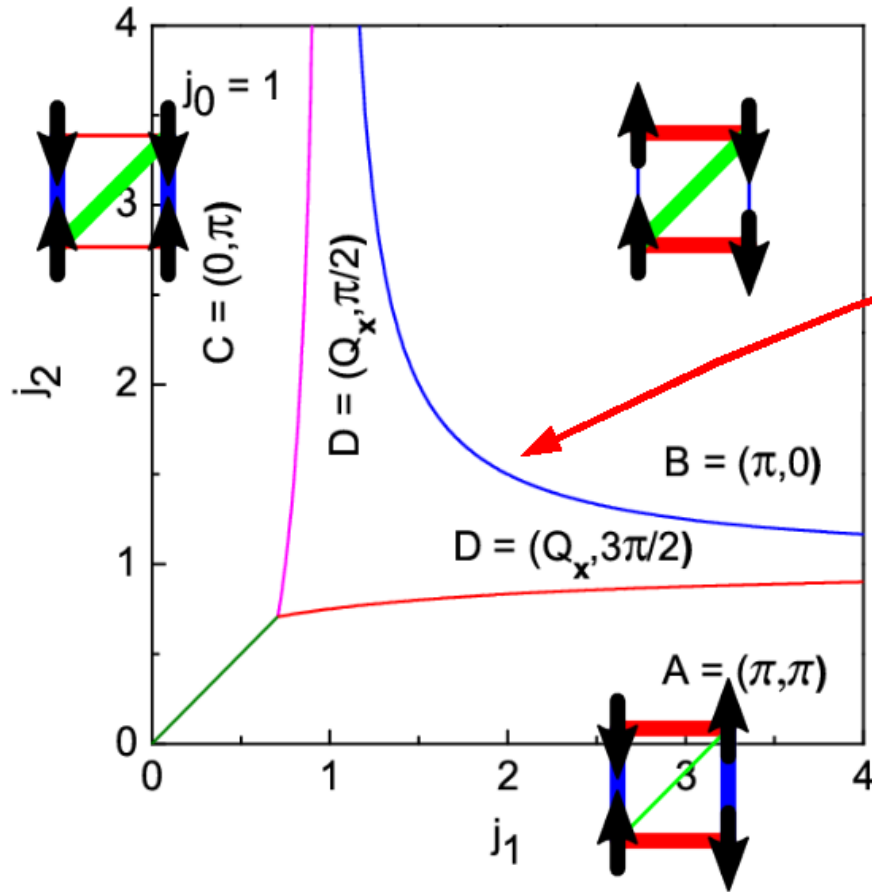
4) Strong Leg Coupling



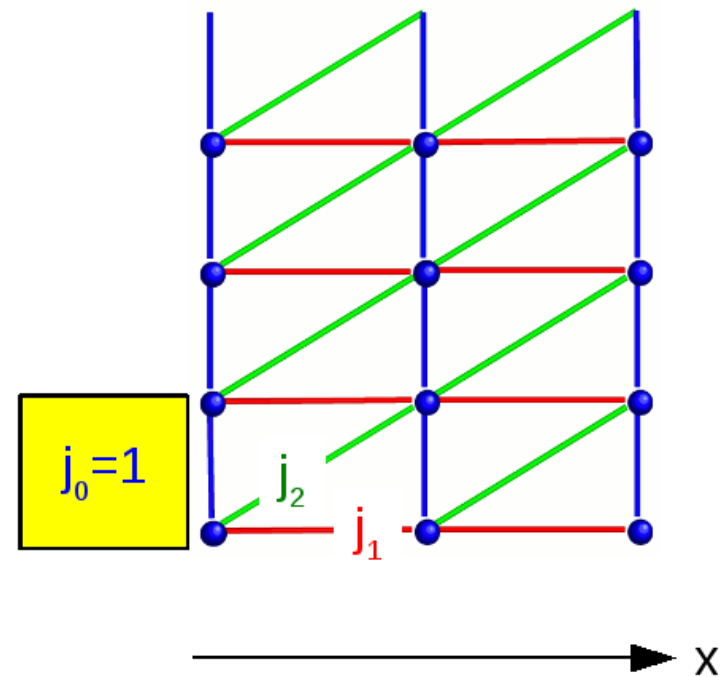


# Phases

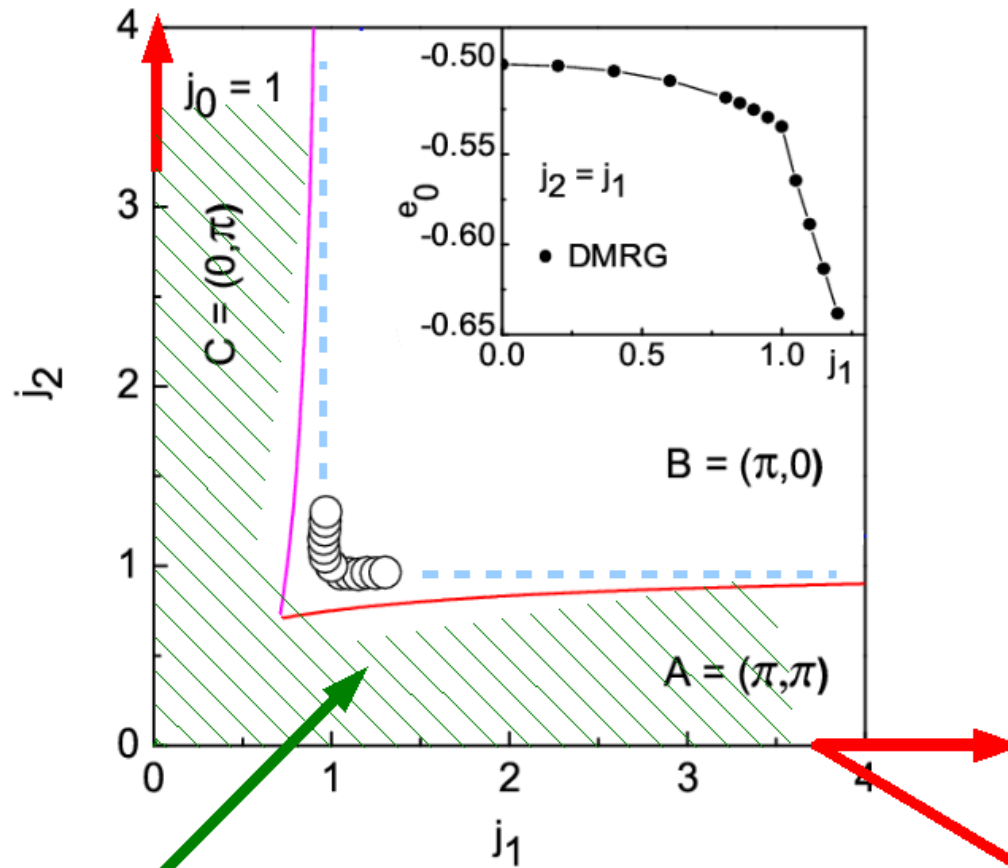
## Classical



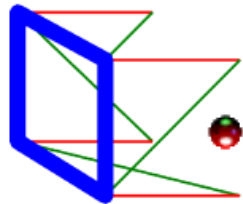
All 1<sup>st</sup> order



# Phases



- DMRG  $L=50$ :
  - One line of 1<sup>st</sup> order transitions terminating plaquette state
  - Only moderate  $c \rightarrow q$  renormalization
  - Very weak beyond  $j_1(j_2) \sim 1.5$ ,  $j_2(j_1) \sim 1$
  - No transition into  $(\pi, 0)$  state



coupled 4-spin singlet-plaquettes

non-frustrated ladder-case  $j_2=0$

Cabra, Honecker, Pujol, PRB 58, 6241 (98)  
Kim and Sólyom, PRB 60, 15230 (99)

$j_1 = \infty, j_2 = 0$  Luttinger liquid

$j_1 \leftrightarrow j_2$  symmetry

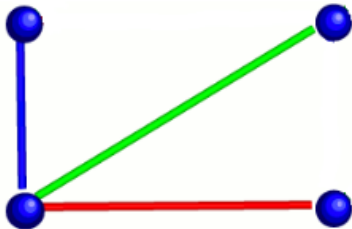
# Phases: Schwinger Bosons

## • SU(2) invariant decoupling

$$\mathbf{S}_l \cdot \mathbf{S}_m =: \hat{B}_{lm}^+ \hat{B}_{lm} : - \hat{A}_{lm}^+ \hat{A}_{lm}$$

$$\text{FM} \quad \hat{B}_{lm}^+ = \frac{1}{2} \sum_{\sigma} b_{l\sigma}^+ b_{m\sigma}$$

$$\text{AFM} \quad \hat{A}_{lm} = \frac{1}{2} \sum_{\sigma} \sigma b_{l\sigma} b_{m-\sigma}$$



$$\left. \begin{aligned} A[B]_{lm} &= \langle \hat{A}[\hat{B}]_{lm} \rangle \\ \sum_{\sigma} \langle b_{l\sigma}^+ b_{l\sigma} \rangle - 2S &= 0 \end{aligned} \right\} \begin{array}{l} \text{6 bond mean fields} \\ \text{1 Lagrange params.} \end{array}$$

Ceccatto, Gazza, Trumper, PRB 47, 12329 (93)  
Flint, Coleman, PRB 79, 014424 (09)

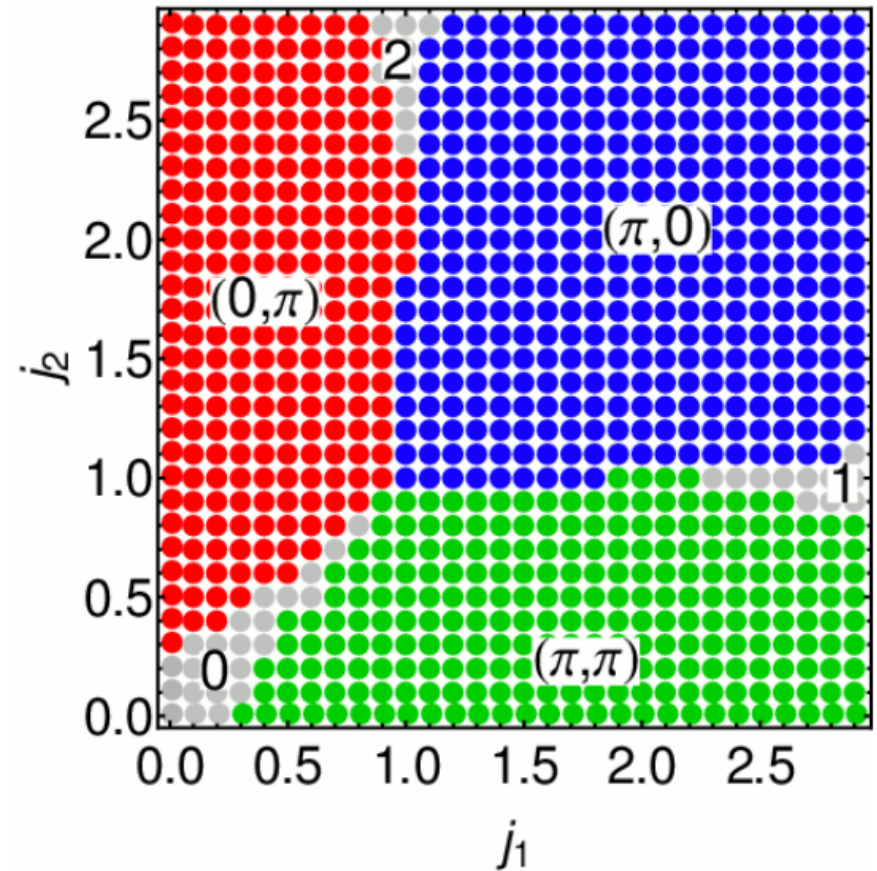
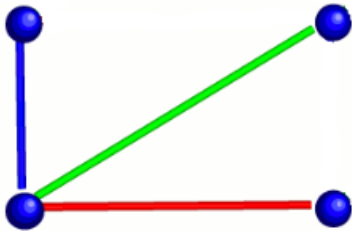
# Phases: Schwinger Bosons

## SU(2) invariant decoupling

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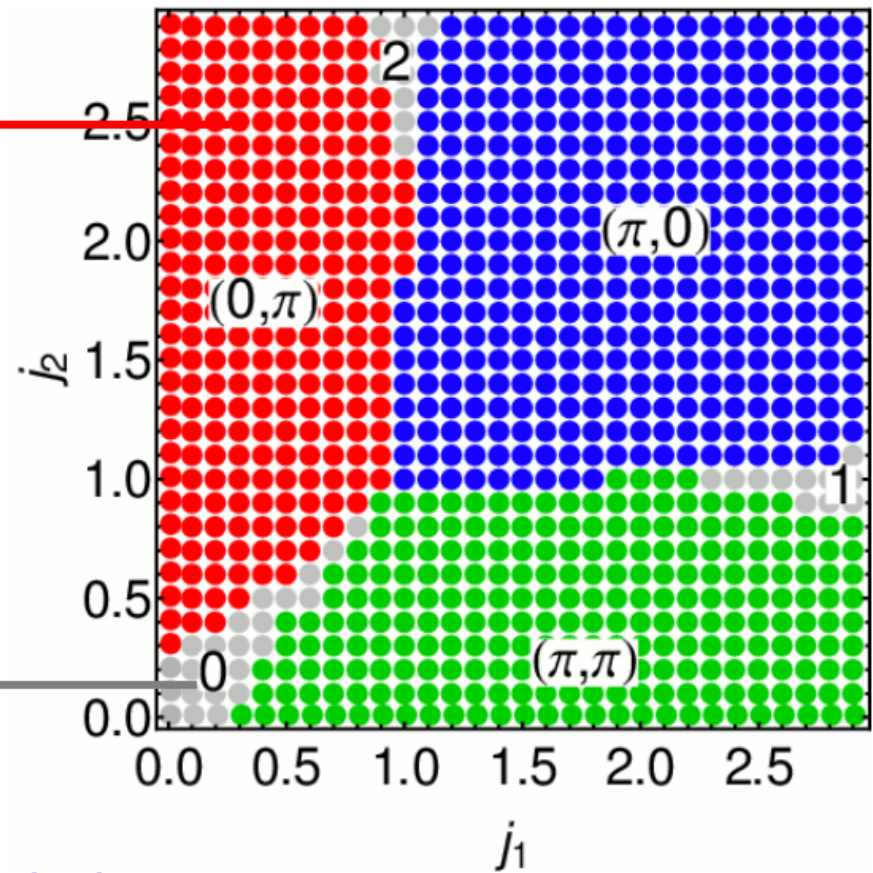
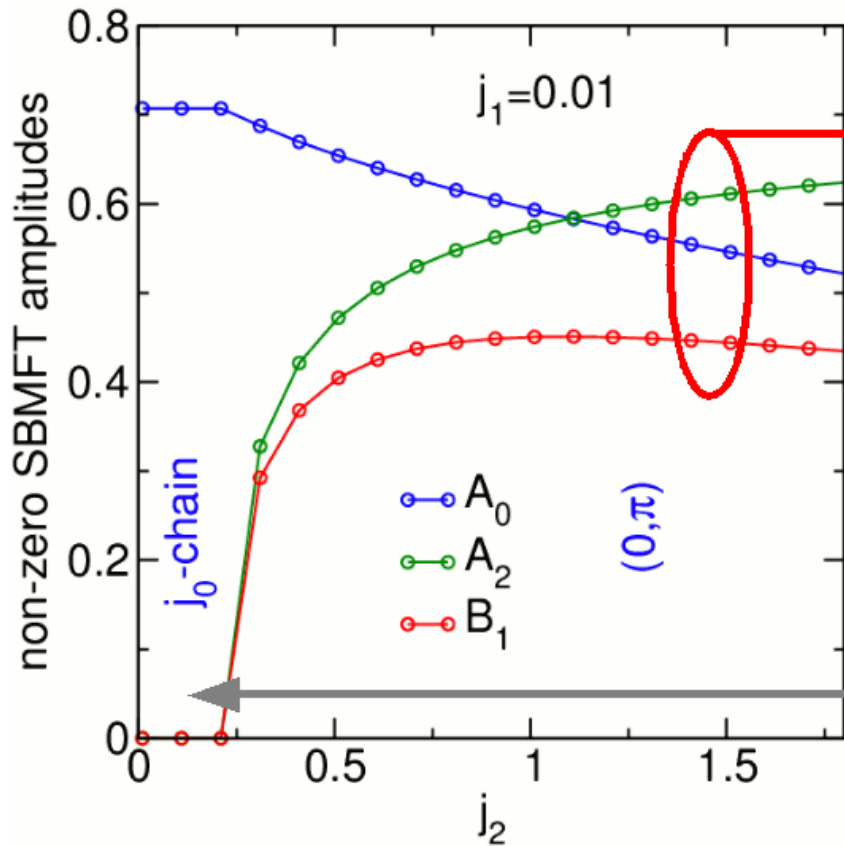
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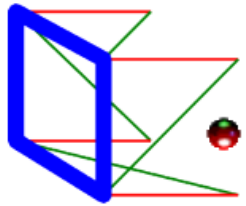
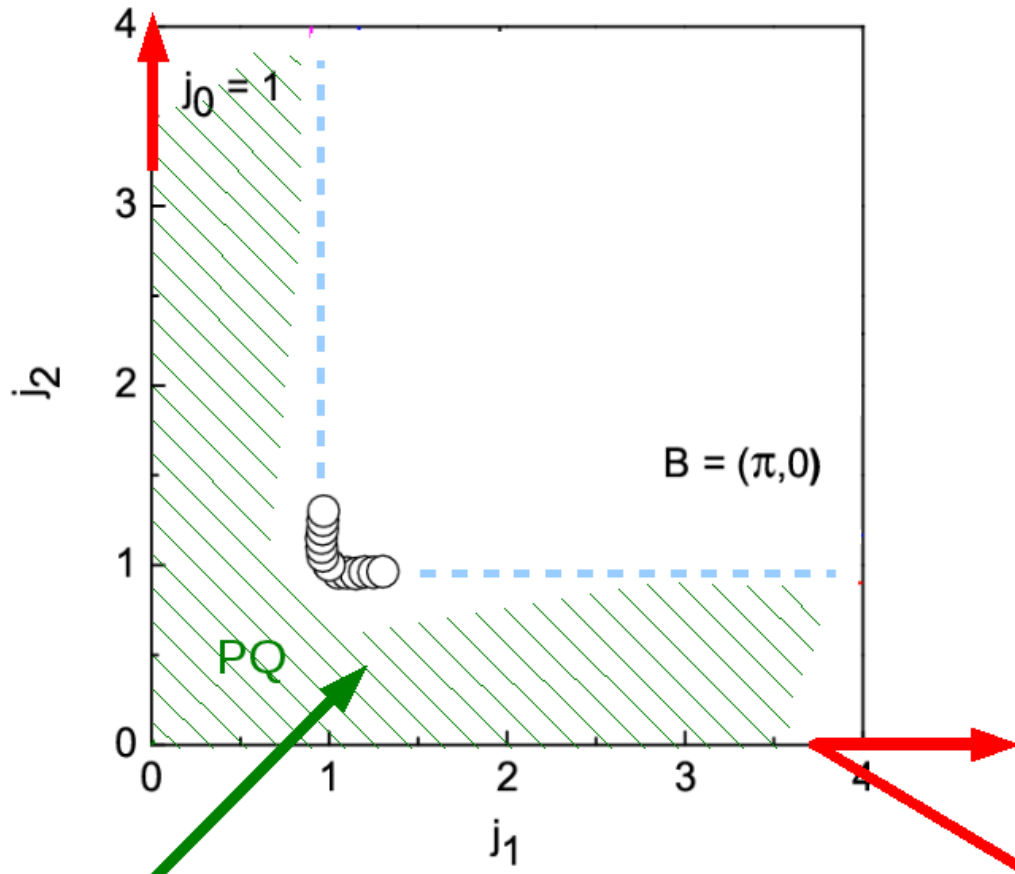
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# Phases: Schwinger Bosons



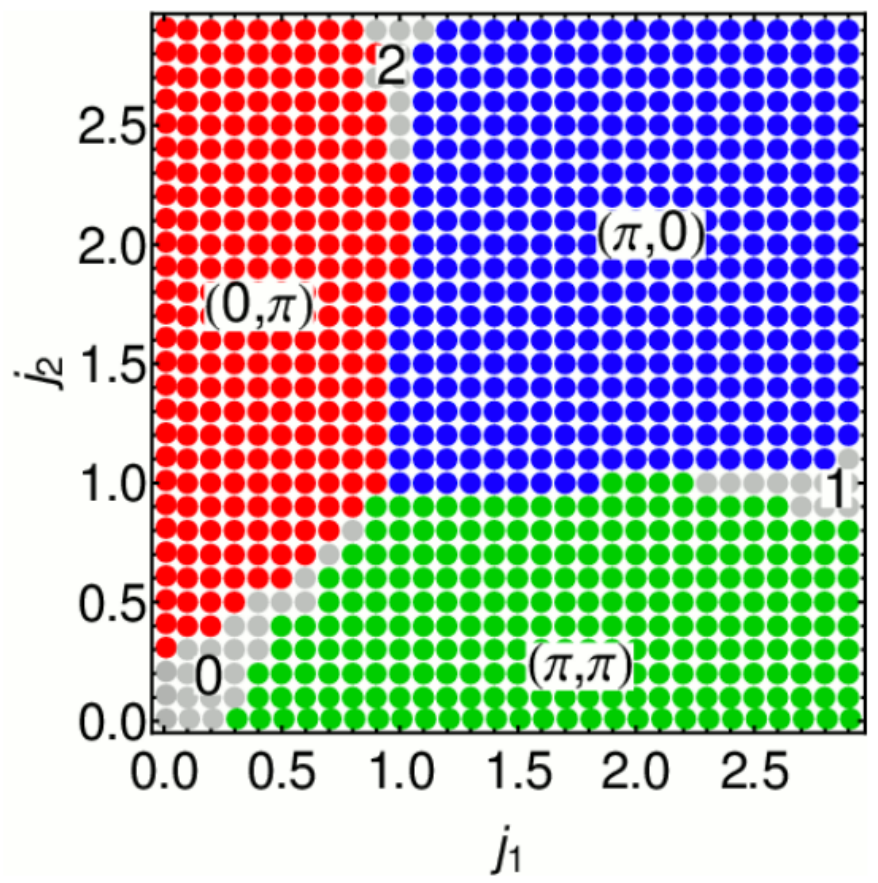
- SR0 similar to classical LRO but no boson condensate
- no IC phase
- red ↔ green ↔ blue: 1<sup>st</sup> order
- RGB ↔ gray: 2<sup>nd</sup> order
- gray = unphysical ⚡

# Phases



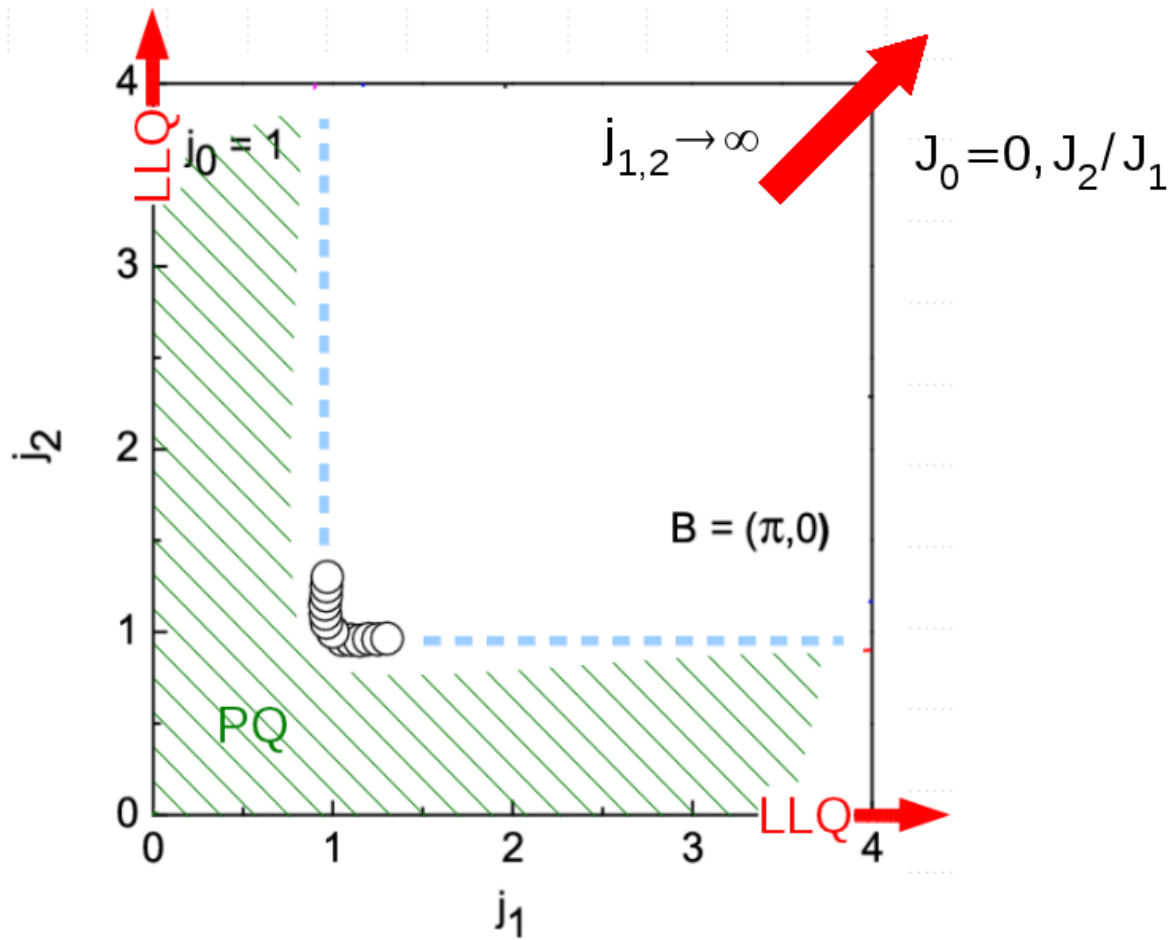
● coupled 4-spin plaquettes

● Schwinger Bosons



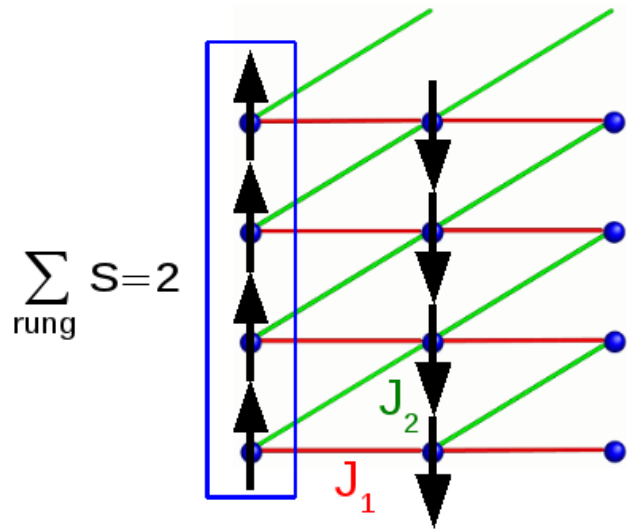
●  $j_1 = \infty, j_2 = 0$  Luttinger liquid

Phases: AFM 'Spin-2 Chain' Limit  $J_0=0$

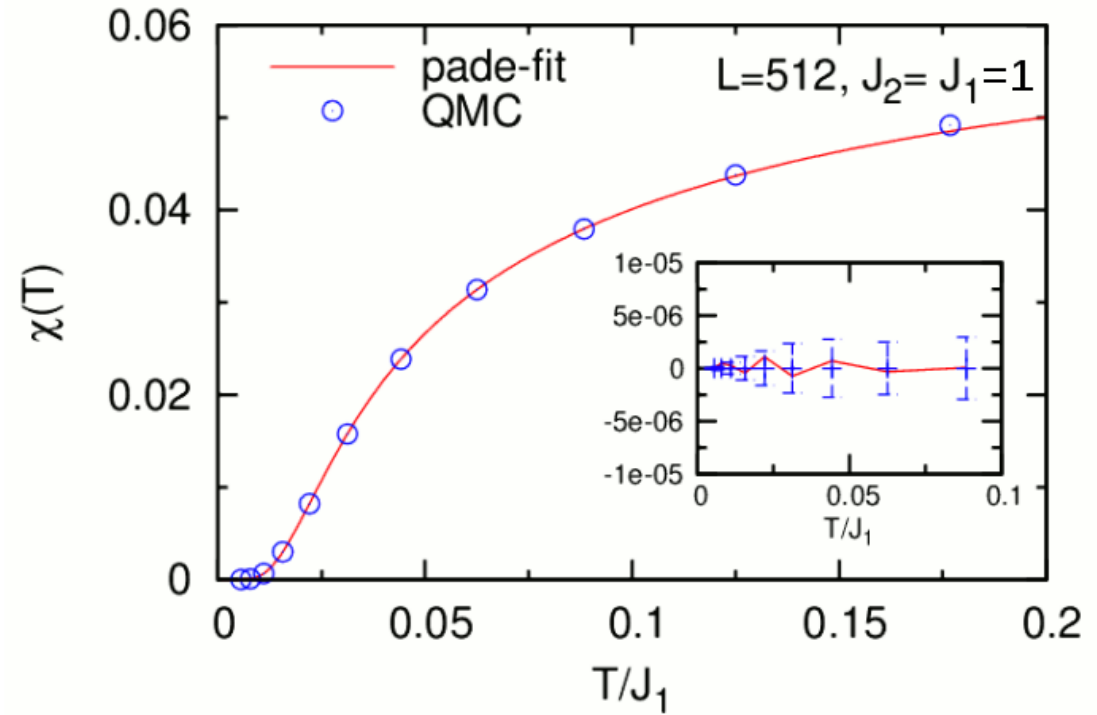


# Phases: AFM 'Spin-2 Chain' Limit $J_0=0$

● Unfrustrated. Classical: 1D  $S=2$  AFM



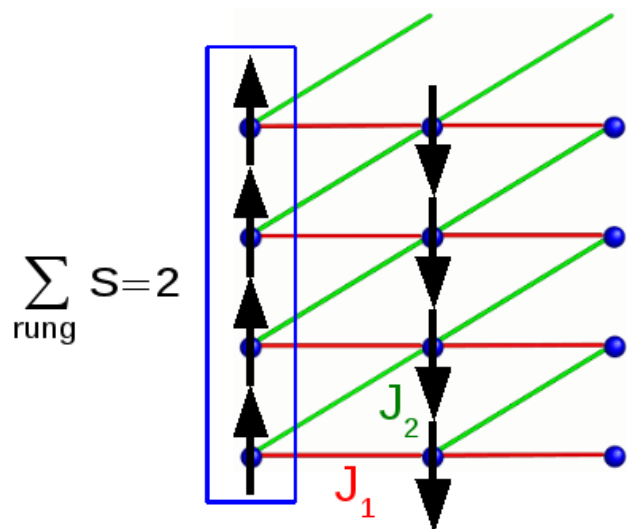
● QMC uniform susceptibility: spin gap



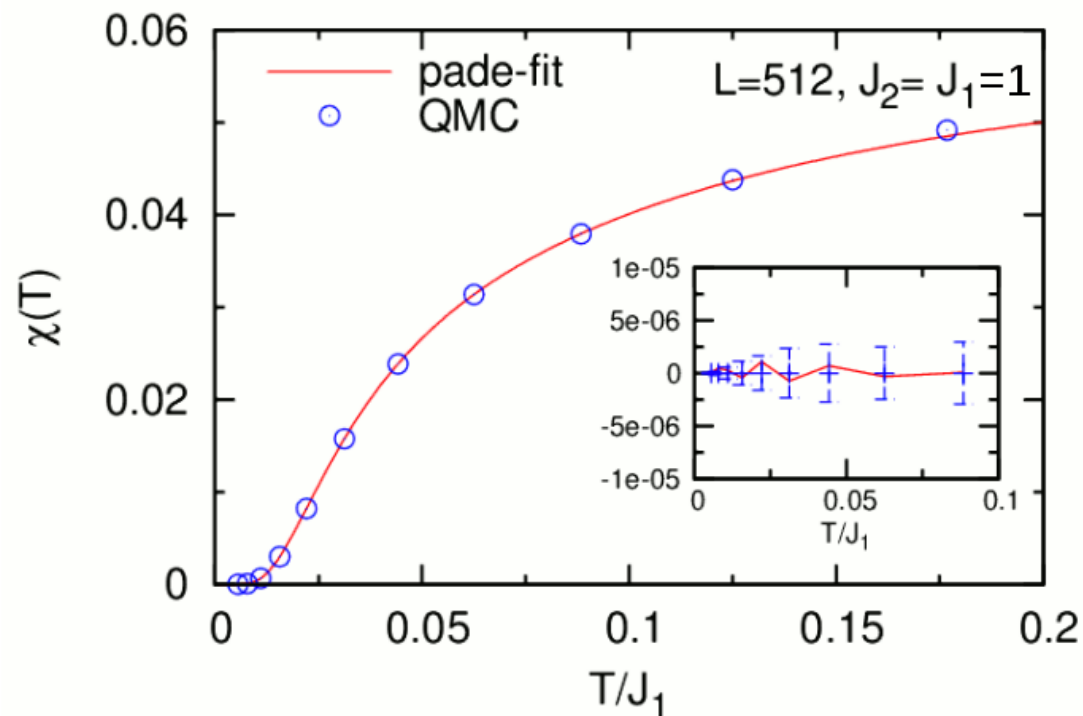


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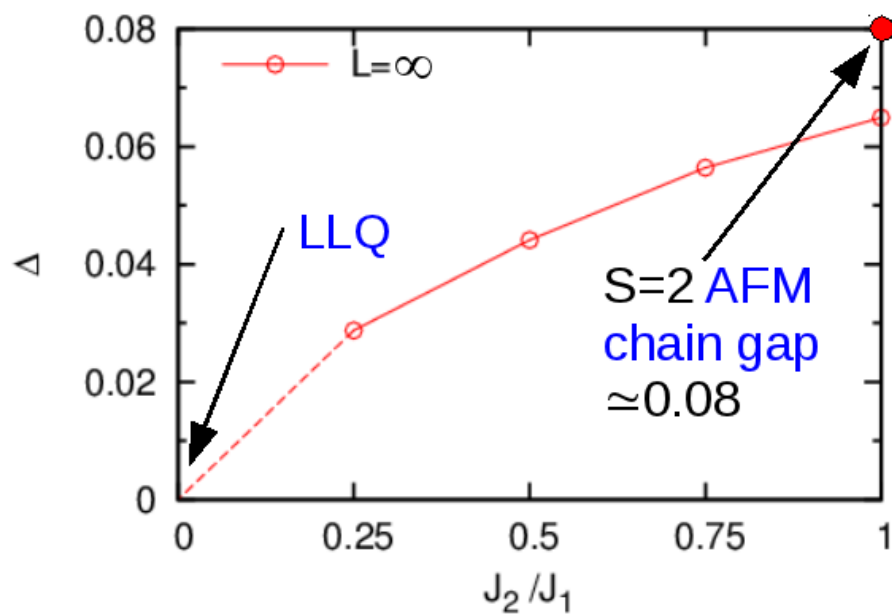
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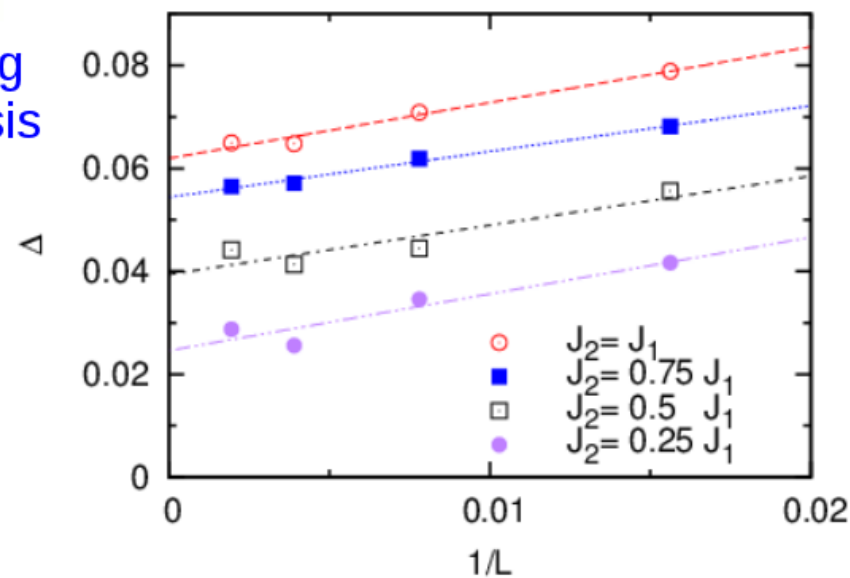
● QMC uniform susceptibility: spin gap



● Gap vs.  $J_2/J_1$

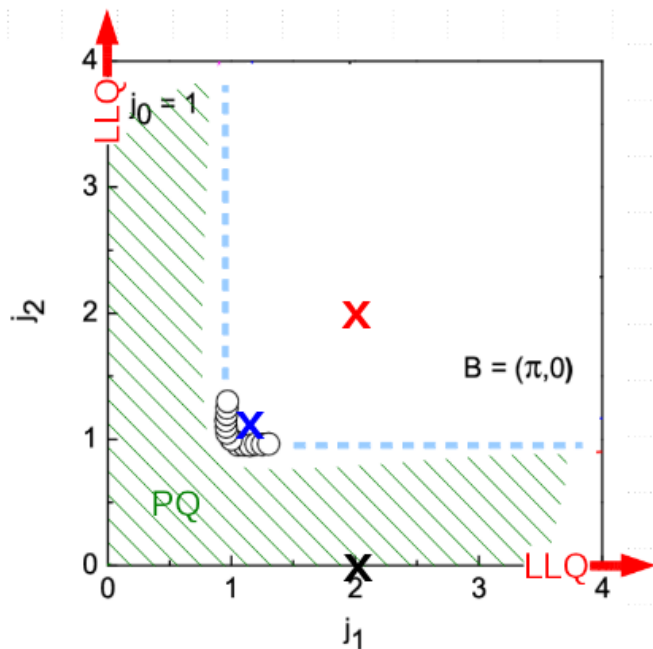


● Scaling analysis

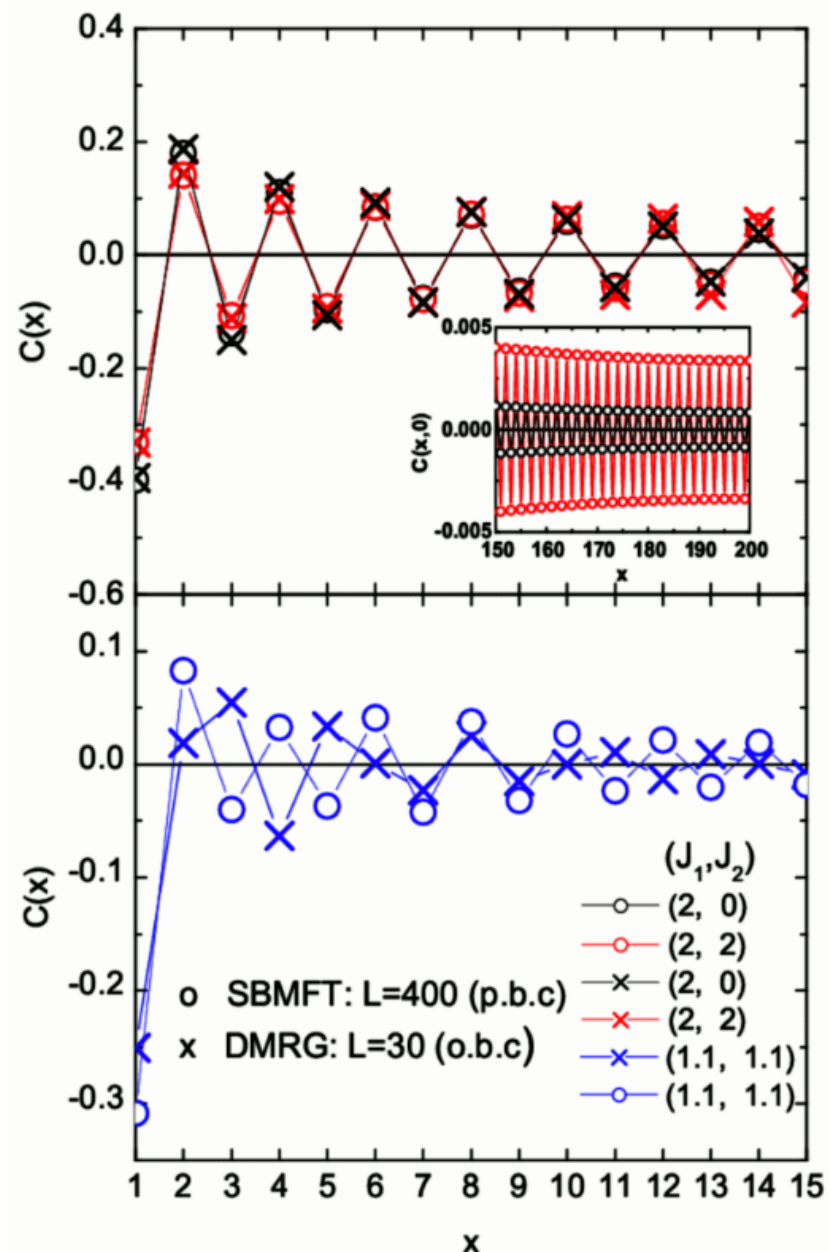


# Static Correlation Functions & Structure Factor

$C(x) = \langle S_{r=(x,0)}^z S_{(0,0)}^z \rangle$  from DMRG & SBMFT

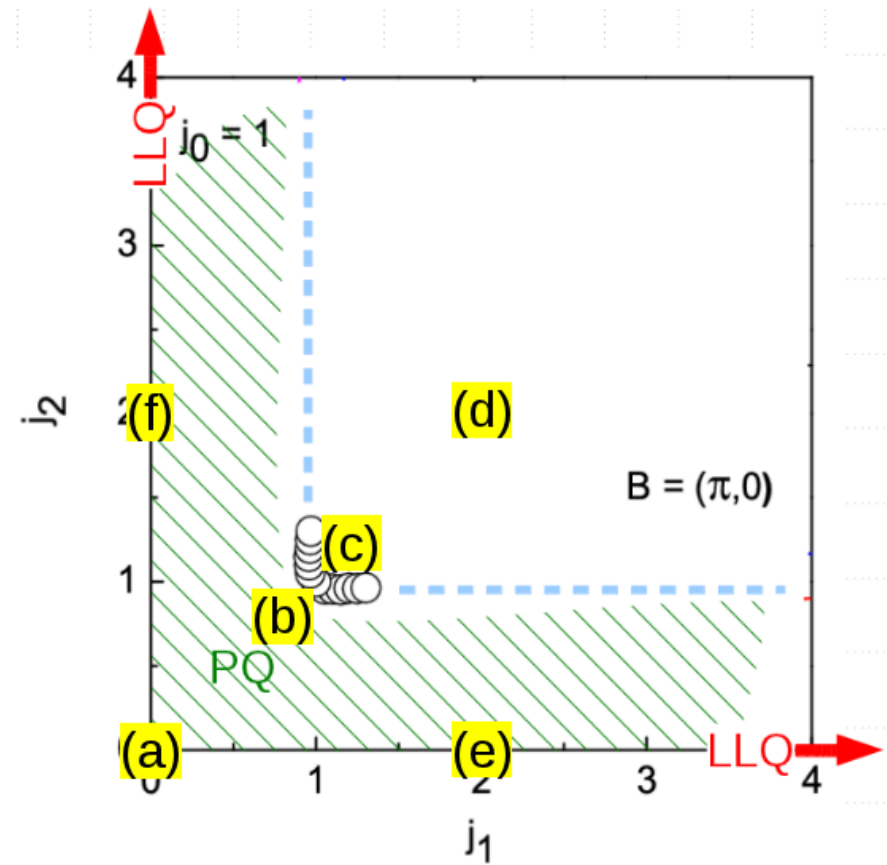
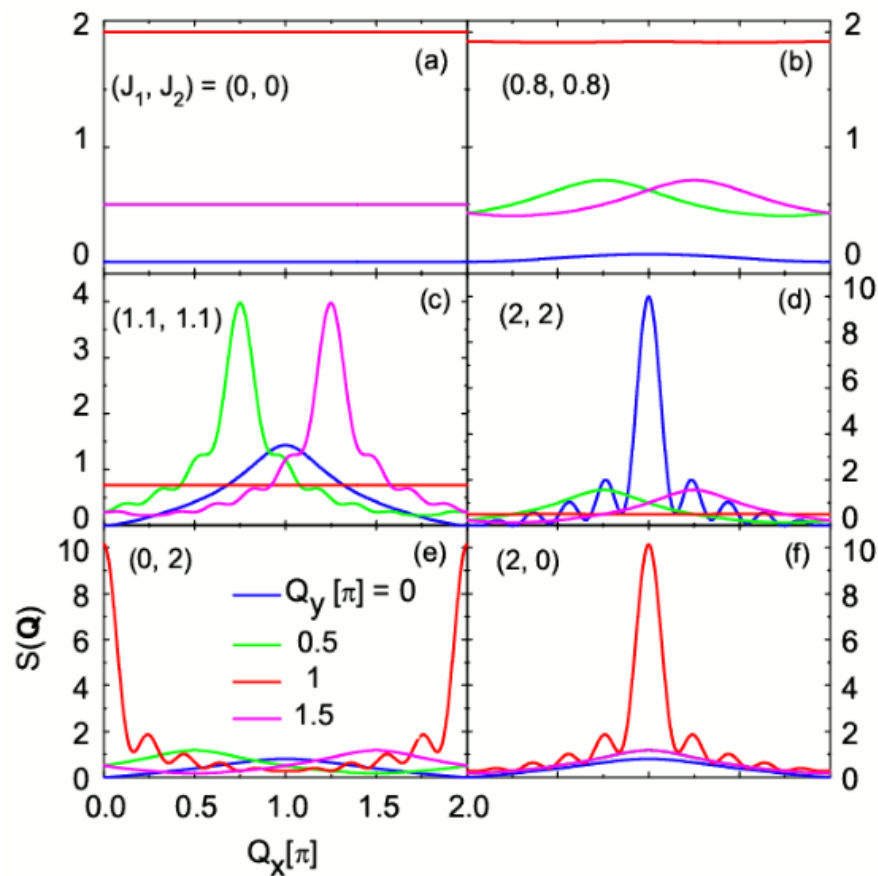


- commensurate cases  
DMRG and SBMFT agree ✓
- however: incommensurate  
region in DMRG beyond 1<sup>st</sup>  
order transition
- $C(x) \sim \exp(-x/l)$  with  $l \rightarrow 0$  as  $j_{1,2} \rightarrow 0$



# Static Correlation Functions & Structure Factor

$$S(\mathbf{Q}) = \sum_{\mathbf{r}} \exp(i\mathbf{Q} \cdot \mathbf{r})$$

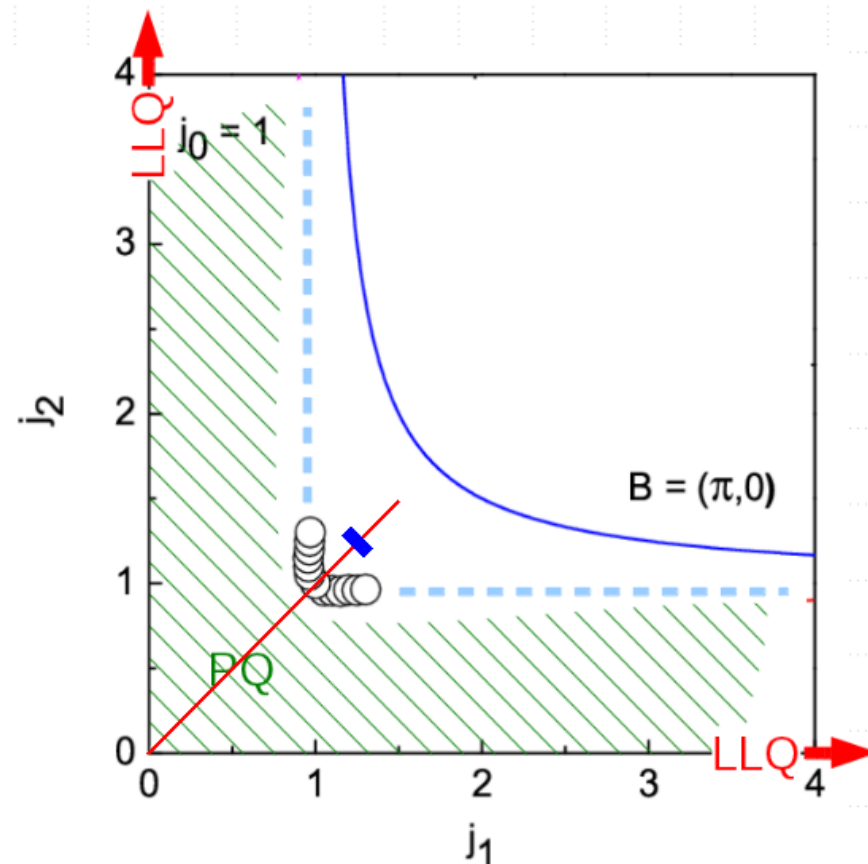
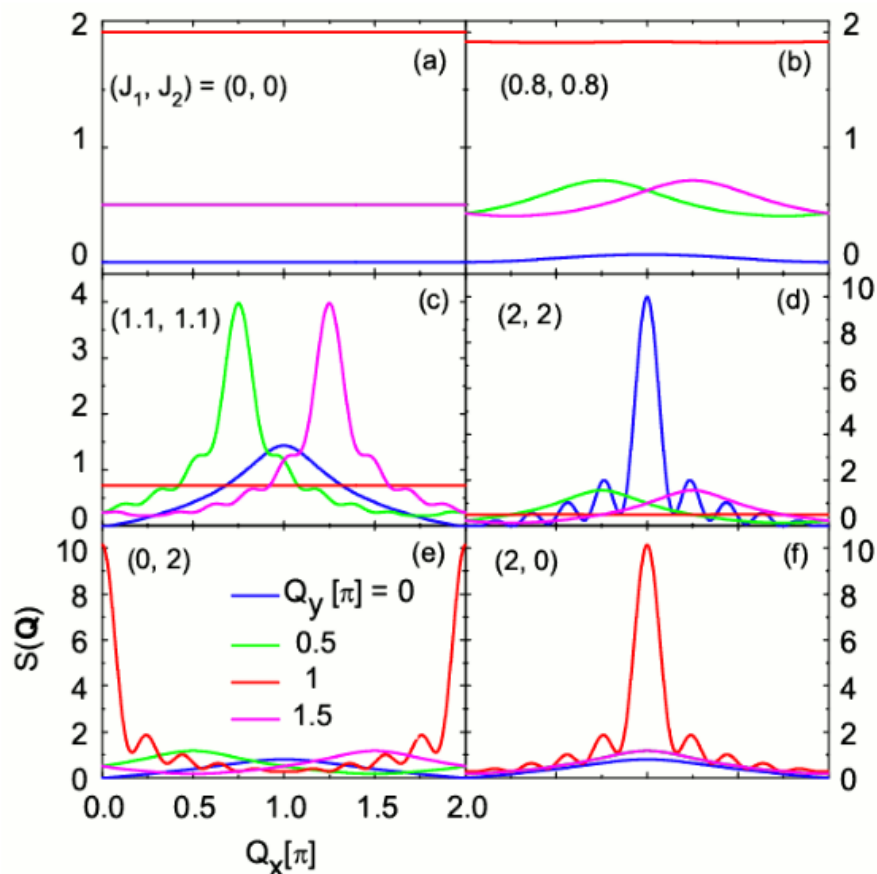


$j_{1,2} \ll 1$ : plaquette regime,  $Q_y \rightarrow \pi$ ,  $Q_x \rightarrow \text{const}$

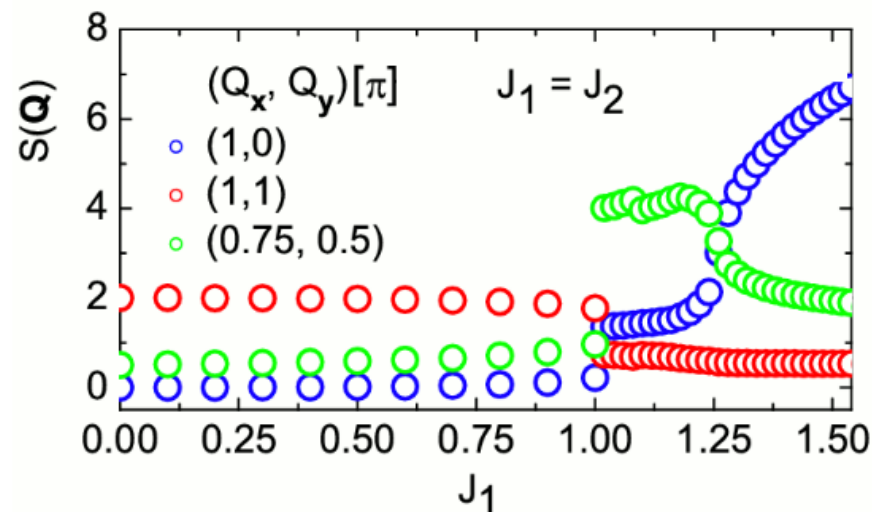
$j_{1,2} > \sim 1$ : incommensurate,  $Q_y, Q_x \cong$  classical

# Static Correlation Functions & Structure Factor

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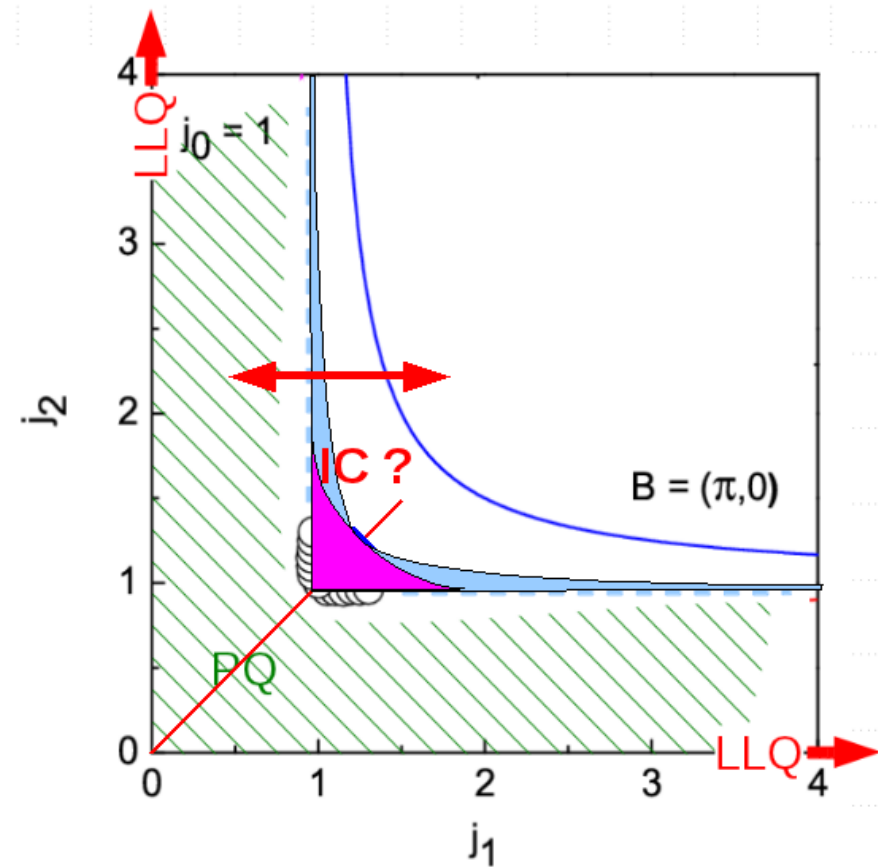
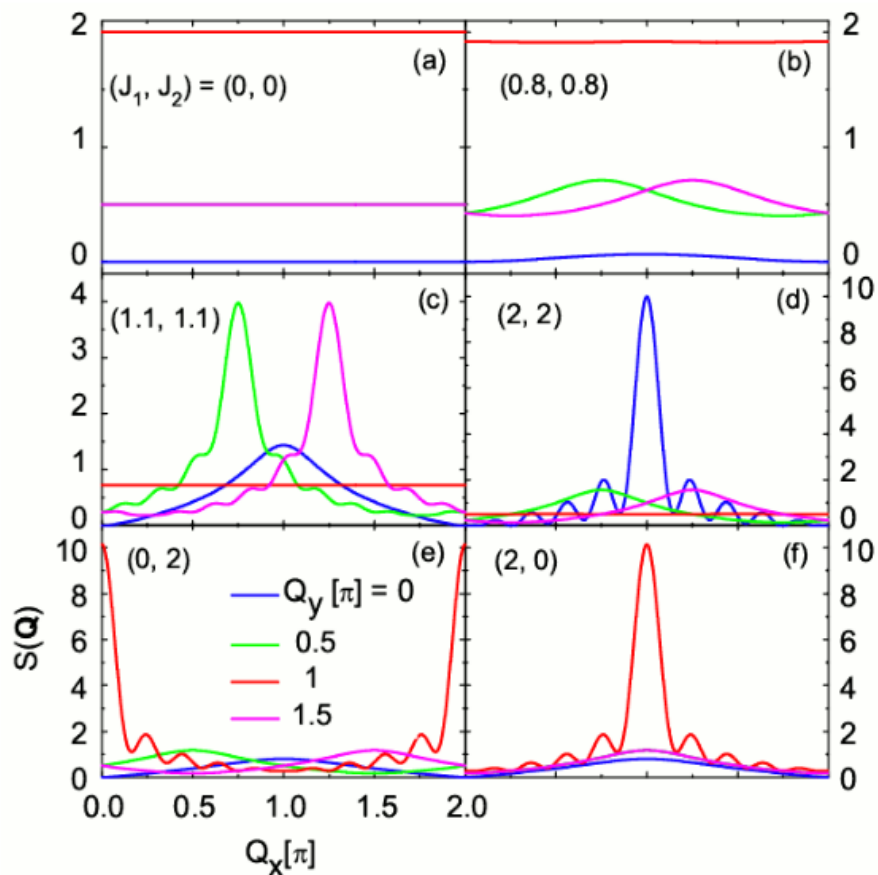


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- strongly renormalized 'transition' into  $(\pi, 0)$
- small IC region likely

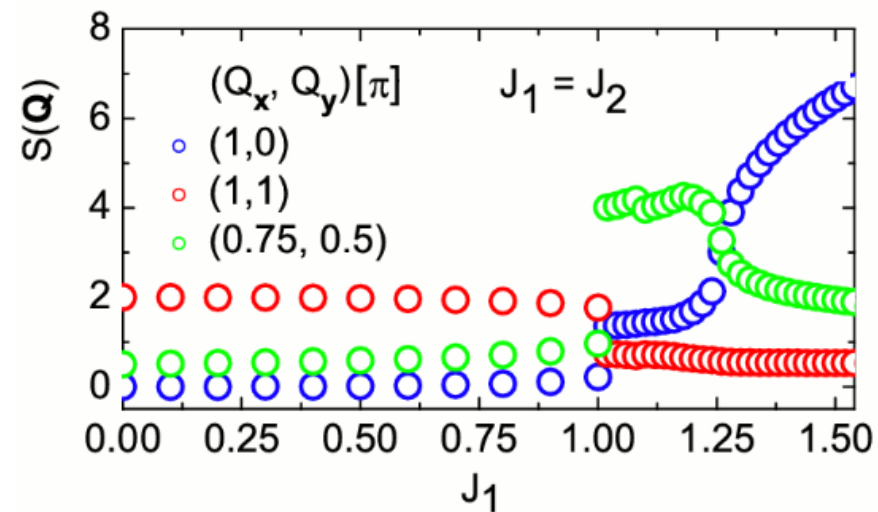


# Static Correlation Functions & Structure Factor

$$S(\mathbf{Q}) = \sum_{\mathbf{r}} \exp(i\mathbf{Q} \cdot \mathbf{r})$$



- $j_{1,2} \ll 1$ : plaquette regime,  $Q_y \rightarrow \pi$ ,  $Q_x \rightarrow \text{const}$
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1) Tube Materials & Model

2) Phases and Static Correlations

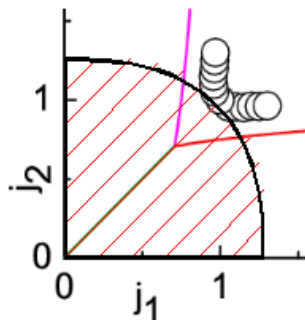
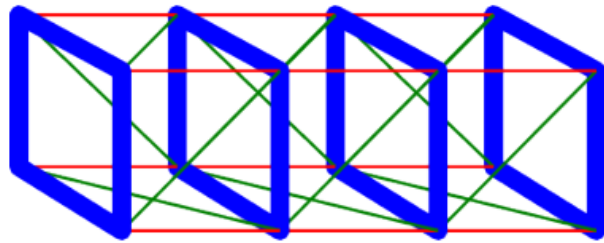
3) Plaquette Regime

4) Strong Leg Coupling



# Spectrum in the Plaquette Regime: Series expansion analysis

## ● Plaquette basis



## plaquette spectrum

E	S	q <sub>i</sub>
1	2	3
0	0 ⊕ 1 ⊕ 1	2
-1	1	1
-2	0	0

# quanta

equidistant

## ● Continuous unitary trafo $H_{\text{eff}} = e^{-\eta(I)} H e^{\eta(I)}$

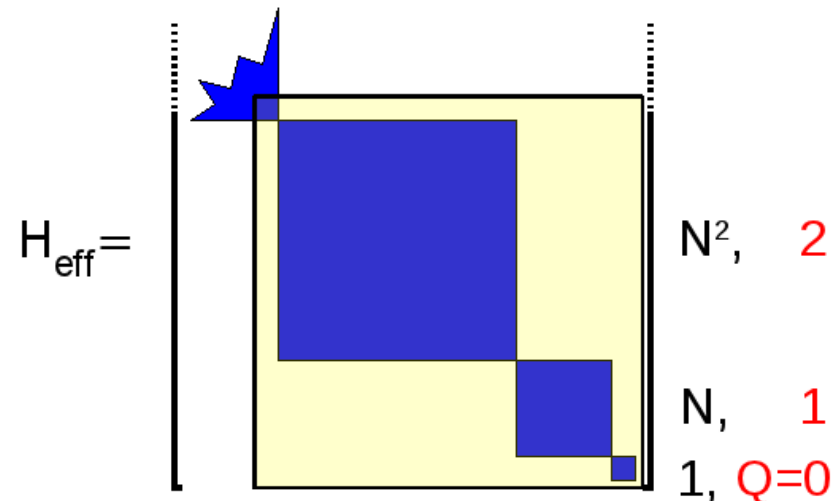
Wegner, Ann. Phys. (94)  
Knetter, et al., EPJB (00)

$$H = H_0 + \sum_{i=1}^2 \sum_{n=-N}^N j_i T_{i,n}$$

$$[H_{\text{eff}} |_{l \rightarrow \infty}, \sum_l q_l] = 0$$

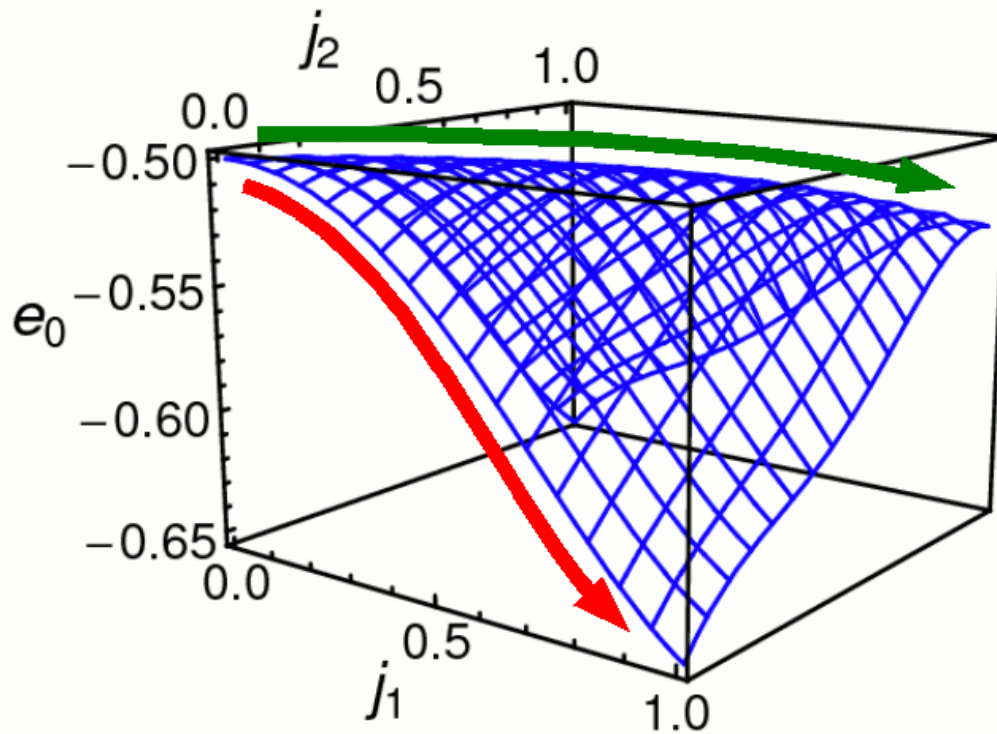
$$H_{\text{eff}} = H_0 + \sum_{ij} \boxed{j_1^i j_2^j} (\dots \Pi \boxed{C} \dots T \dots)$$

from DEQN



# Ground State Energy

$e_0 = \langle Q=0 | H_{\text{eff}} | Q=0 \rangle_0 / L$  up to  $O(7)$  in  $j_{1,2}$

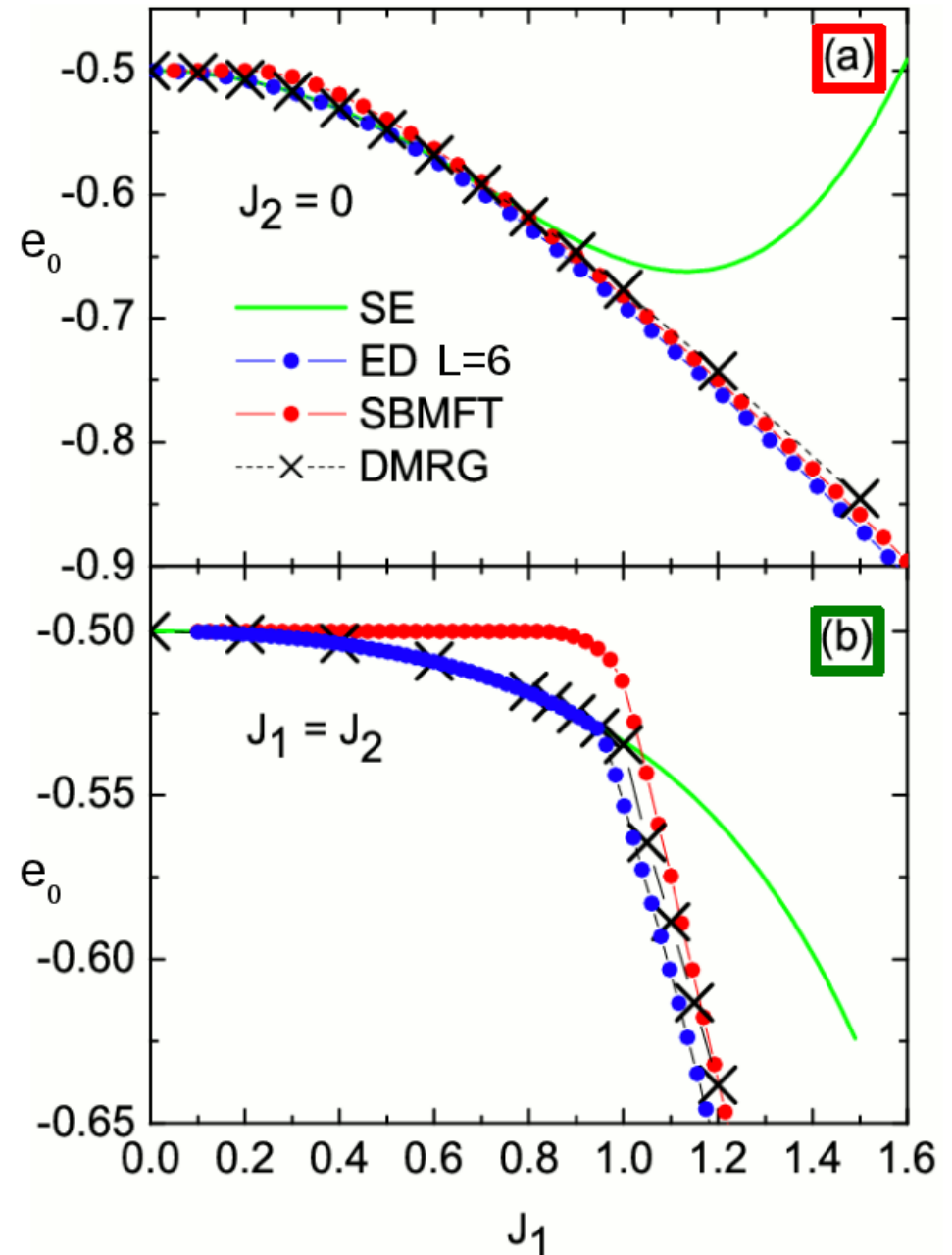


frustration  $\uparrow$  energy gain  $\downarrow$

up to 1<sup>st</sup> order QPT agrees with DMRG ✓

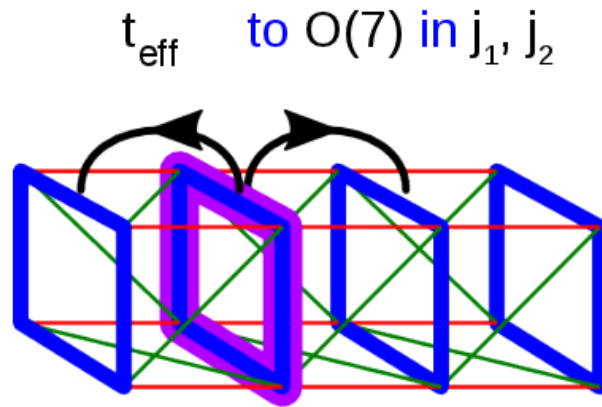
convergence radius  $\sim 1$

contrast ...





# Elementary One-Particle Excitations



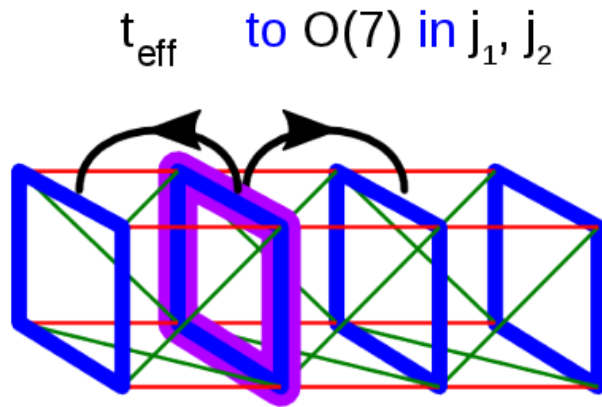
- From  $Q=1$  sector

$$H_{\text{eff}} |Q=1, r=0\rangle_0 = \sum_l t_l |Q=1, r=l\rangle_0 \quad \text{to } O(7)$$

- From  $Q=2$  sector?

$$\langle Q=2, 2\text{pt} | H_{\text{eff}} | Q=2, 1\text{pt} \rangle_0 = 0 \quad \text{to } O(7)$$

# Elementary One-Particle Excitations



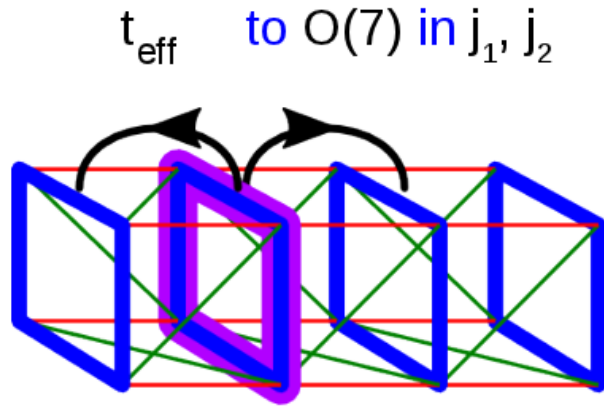
- From  $Q=1$  sector

$$H_{\text{eff}} |Q=1, r=0\rangle_0 = \sum_l t_l |Q=1, r=l\rangle_0 \quad \text{to } O(7)$$

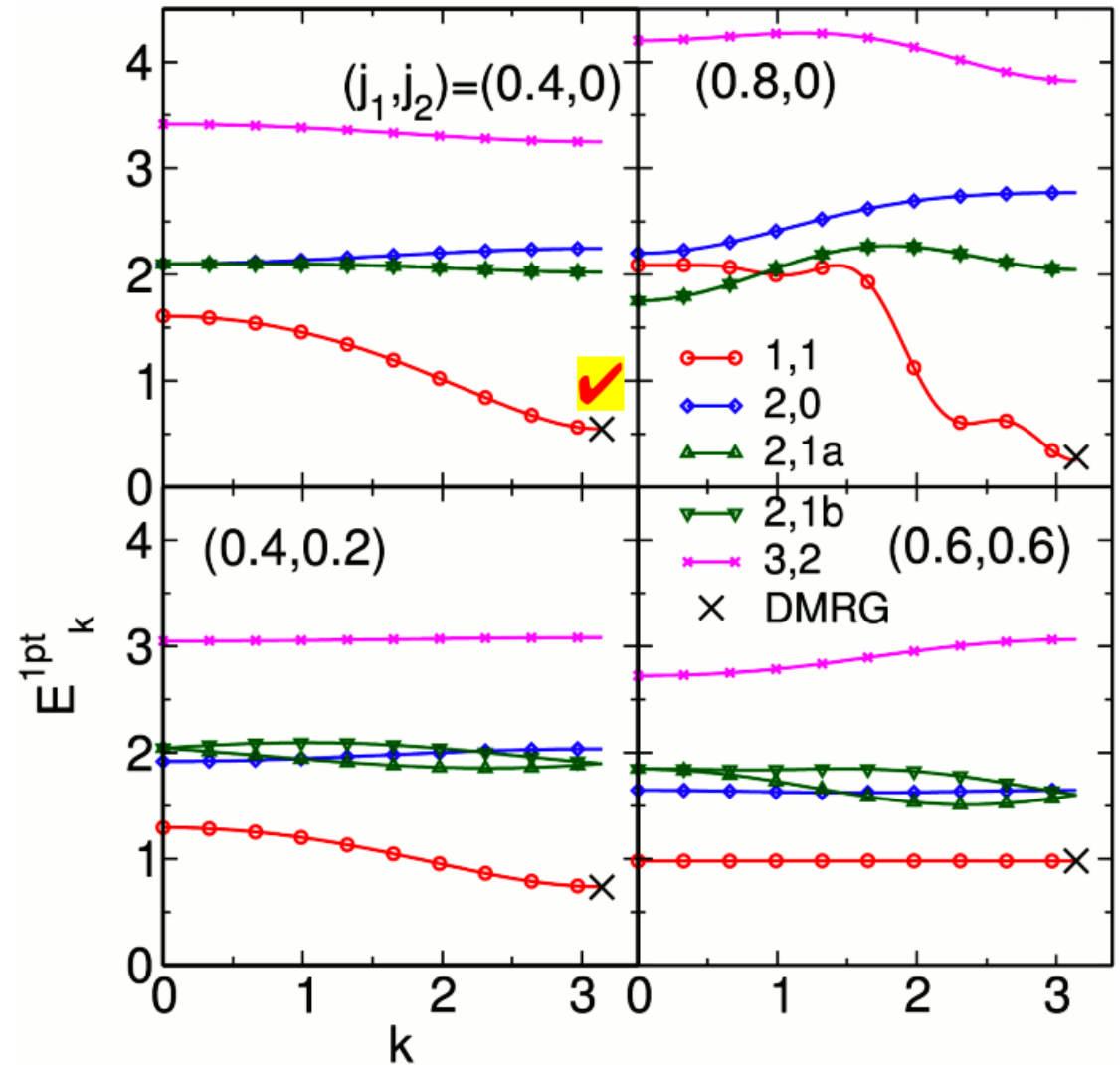
- From  $Q=2$  sector?

$$H_{\text{eff}} |Q=2, \text{1pt}, r=0\rangle_0 = \sum_l t_l |Q=2, \text{1pt}, r=l\rangle_0 \quad \text{to } O(7)$$

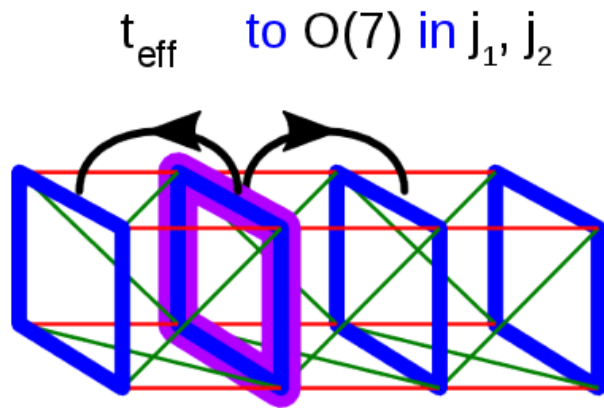
# Elementary One-Particle Excitations



Q=1	1 triplon	un-protected
2	2 triplons + 1 singlon	
3	1 quinton (?)	



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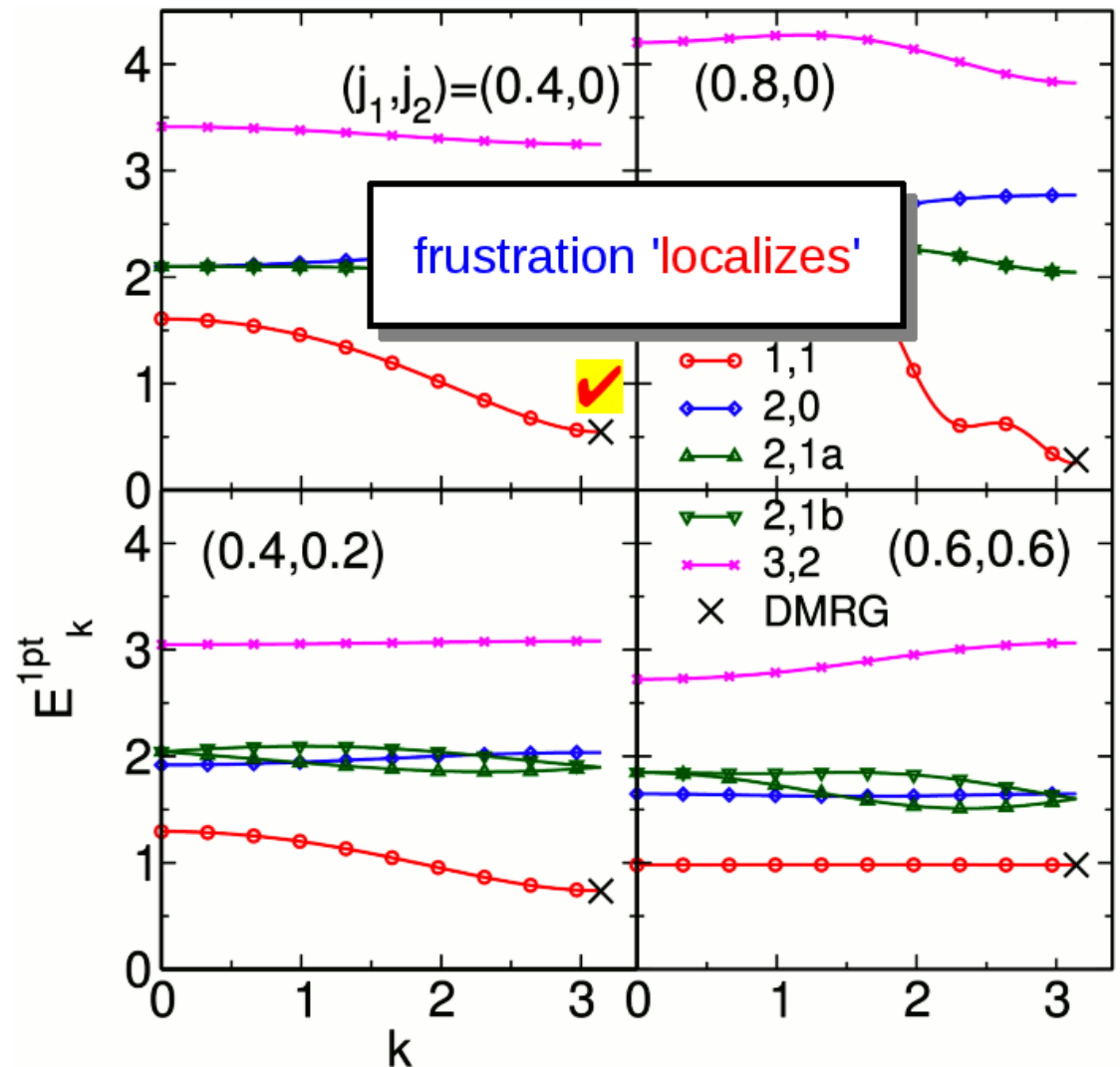
•  $j_1 \leftrightarrow j_2$  from symmetries eg.:

$$E_{1,1}^{1\text{pt}}(\mathbf{k}, j_1, j_2) = E_{1,1}^{1\text{pt}}(\mathbf{k} + \pi, j_2, j_1), \dots$$

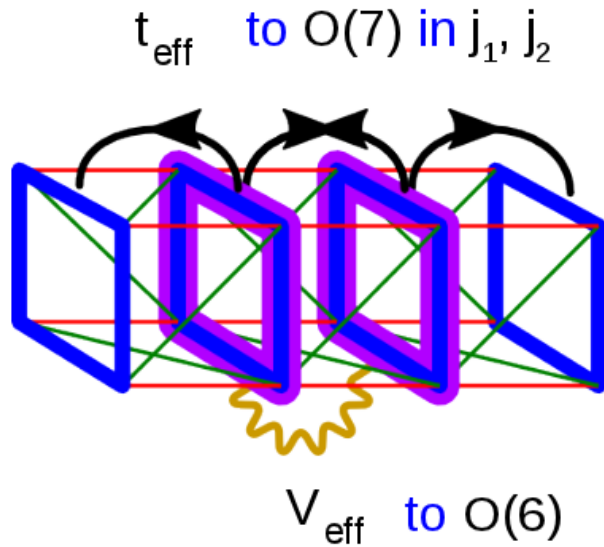
• Gap @  $j_2=0$

$$E_{1,1}^{1\text{pt}}(\pi, j_1, 0) = 1 - \frac{4}{3}j_1 + \frac{41}{108}j_1^2 + \frac{349}{1296}j_1^3 + \frac{4596401}{39191040}j_1^4 - \frac{169497997}{4702924800}j_1^5 - \frac{689874137639377}{2986545364992000}j_1^6 - \frac{8430165345498432721}{45156565918679040000}j_1^7$$

Cabra, et al. PRB 58, 6241 (98) ←



# Two-Particle Excitations



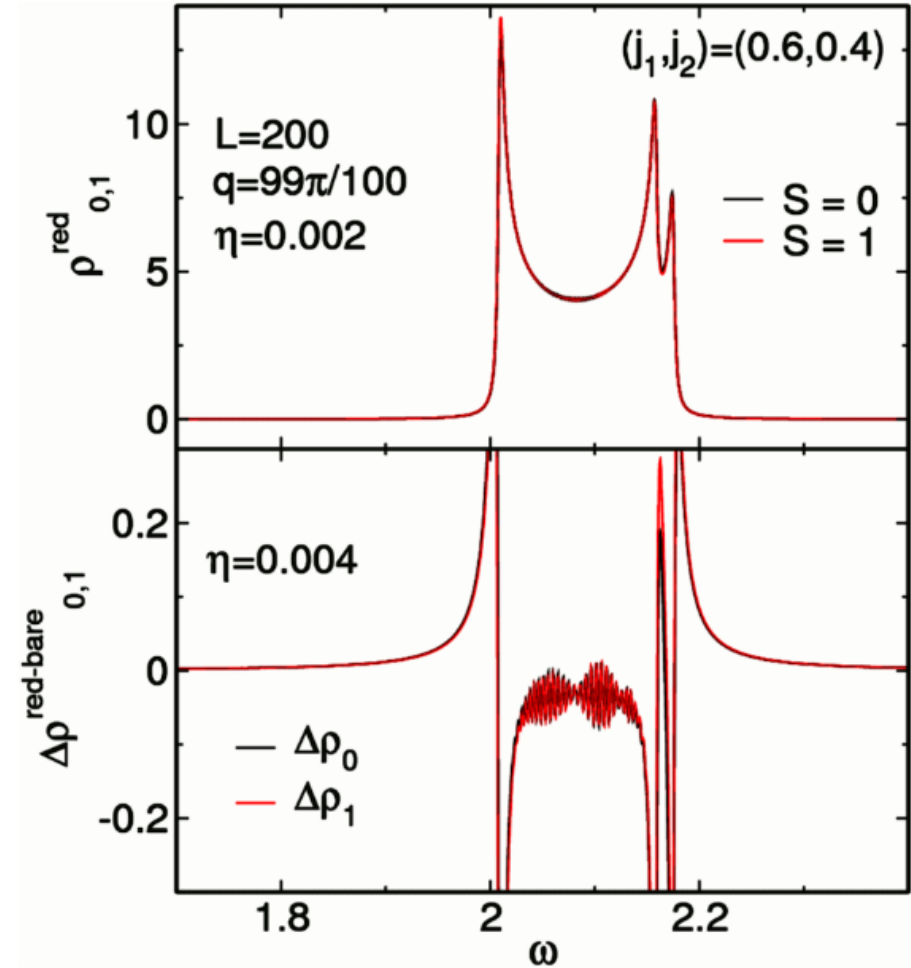
$$|q, d\rangle^{\text{Sm}} = \frac{1}{\sqrt{L}} \sum_r e^{iq(r+d/2)} |Q=2, r, r+d\rangle^{\text{Sm}}$$

$\text{Sm} \langle q, d | H_1 + H_2 | q, d \rangle^{\text{Sm}}$  on finite  $L$

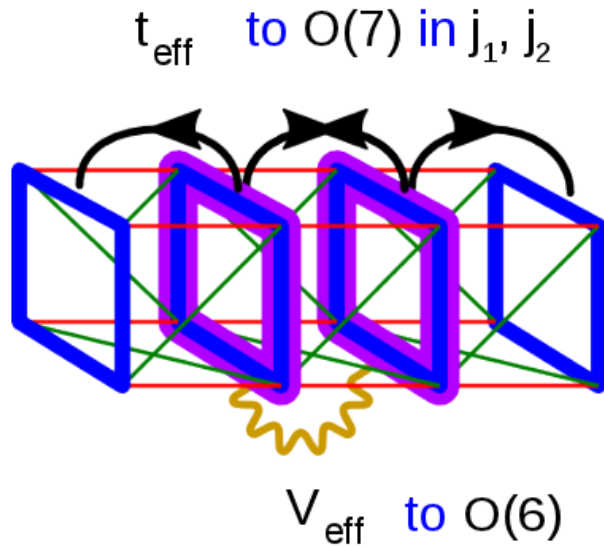
one-kinetic erg. hard-core

two-particle irreducible 'real' interactions

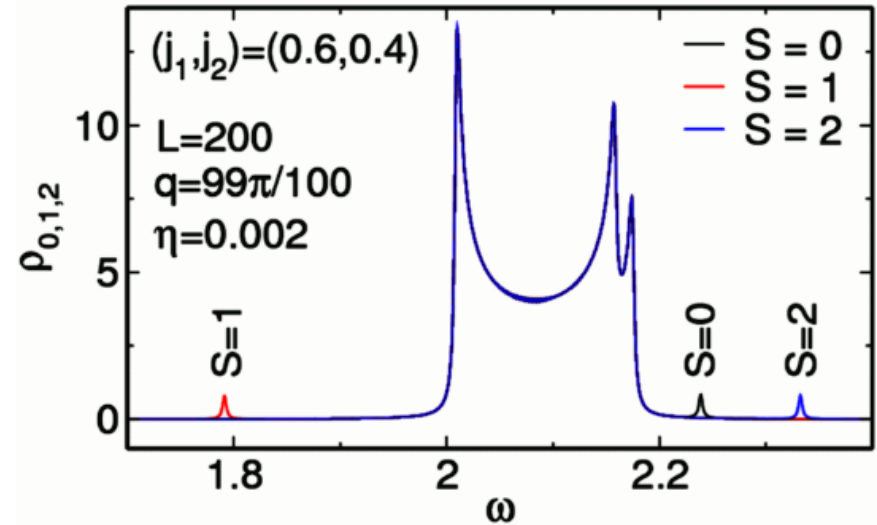
one-particle irreducible



# Two-Particle Excitations

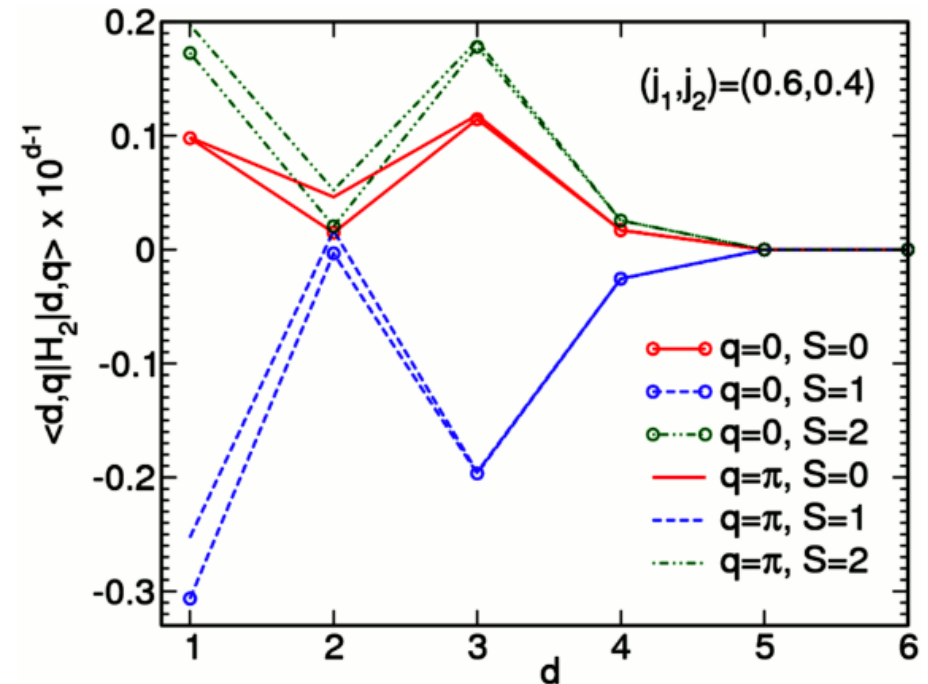


one- and two-particle irreducible

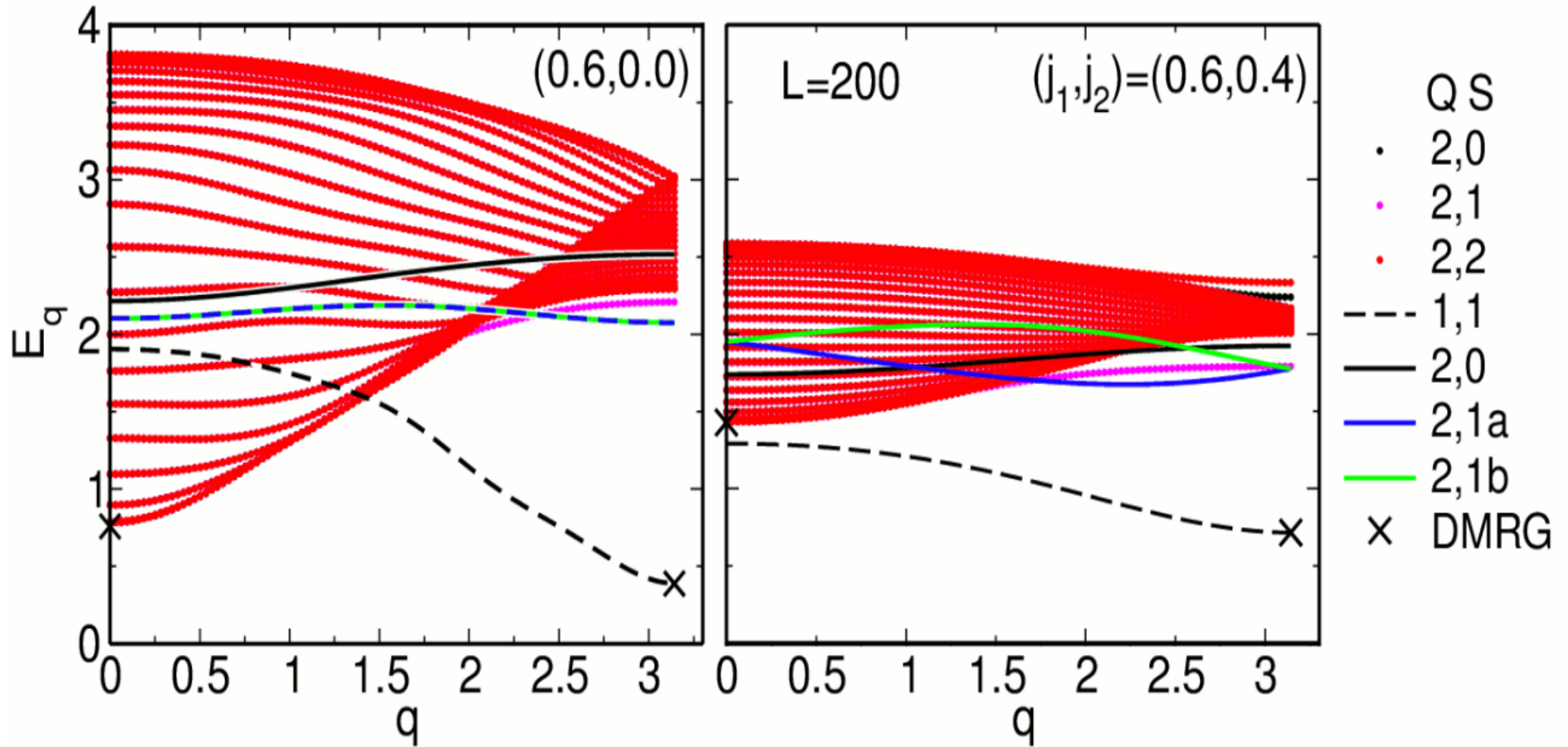


Scattering continua + bound triplet + anti-bound singlet + quintet

two-particle potential as  $f(d)$  rapidly decaying

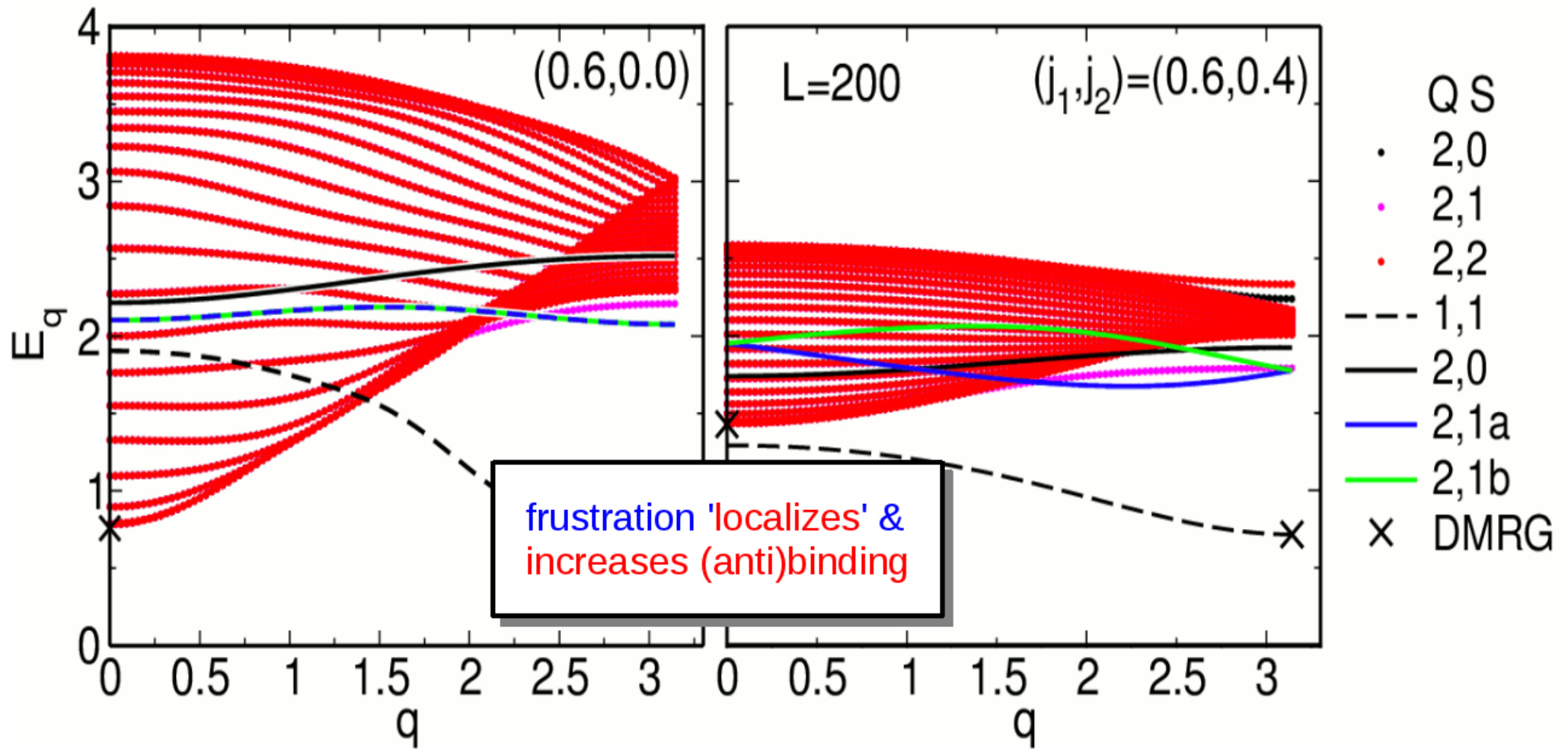


# One- & Two-Particle Dispersion



- several one-particle states within the two-particle continuum
- three collective (anti)bound states
- excellent agreement with DMRG

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1) Tube Materials & Model

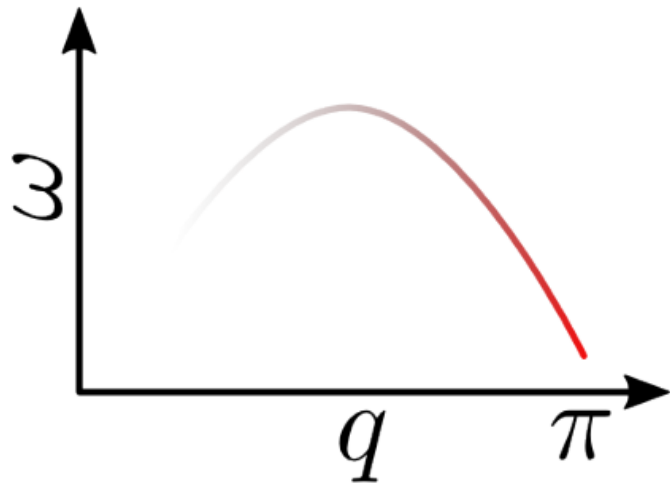
2) Phases and Static Correlations

3) Plaquette Regime

4) Strong Leg Coupling



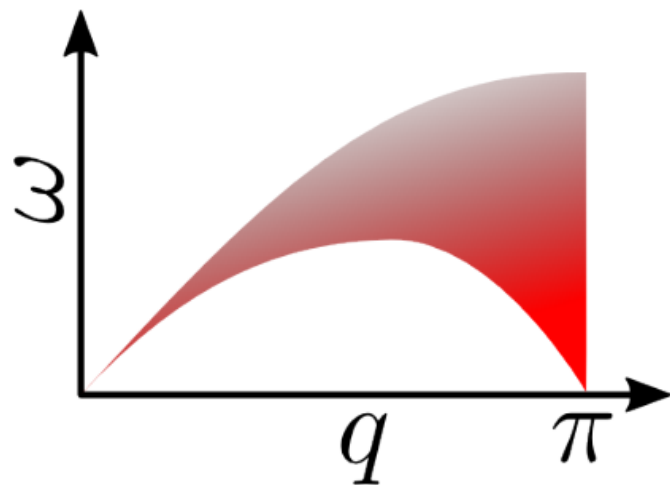
# Spinon-Magnon Crossover



$J_0=0$

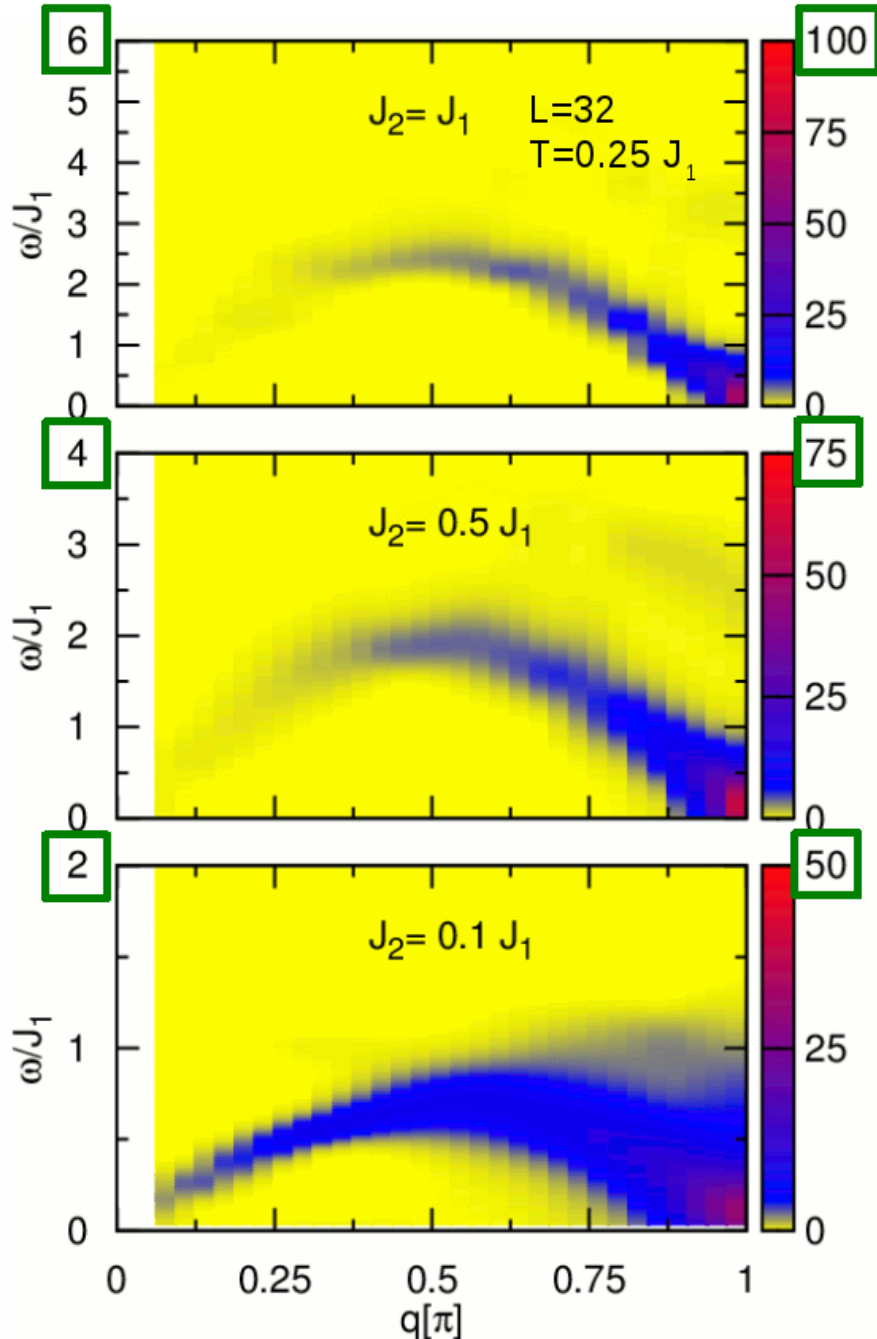
$J_2/J_1=1$ ,  $(\pi,0)$   $S \sim 2$ -chain: magnons

$J_2/J_1=0$ , 4  $S \sim 1/2$ -chain: spinons



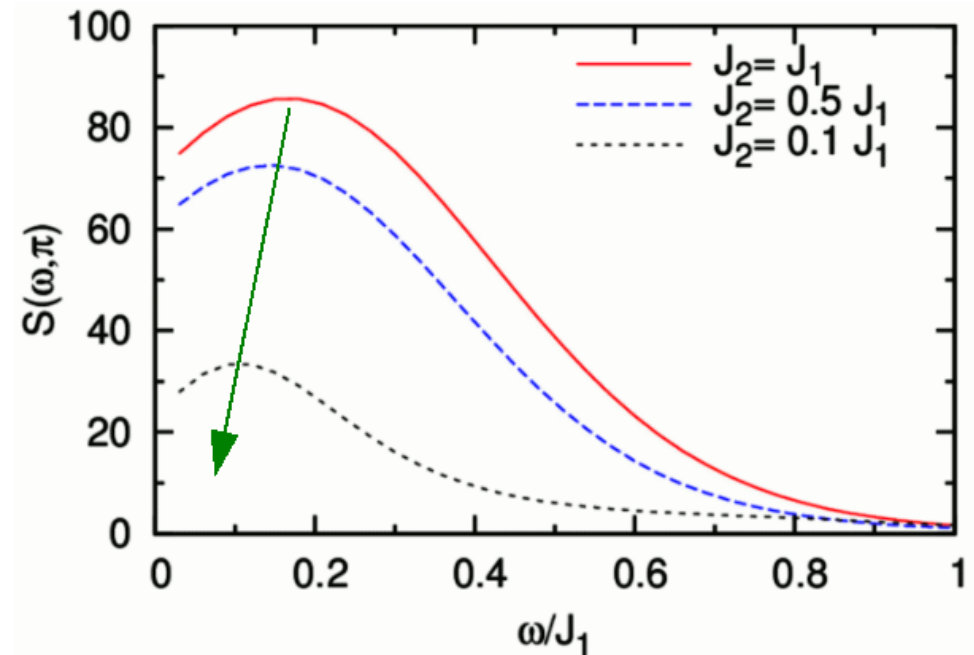
# Spinon-Magnon Crossover

## Dynamic structure factor vs $J_2/J_1$



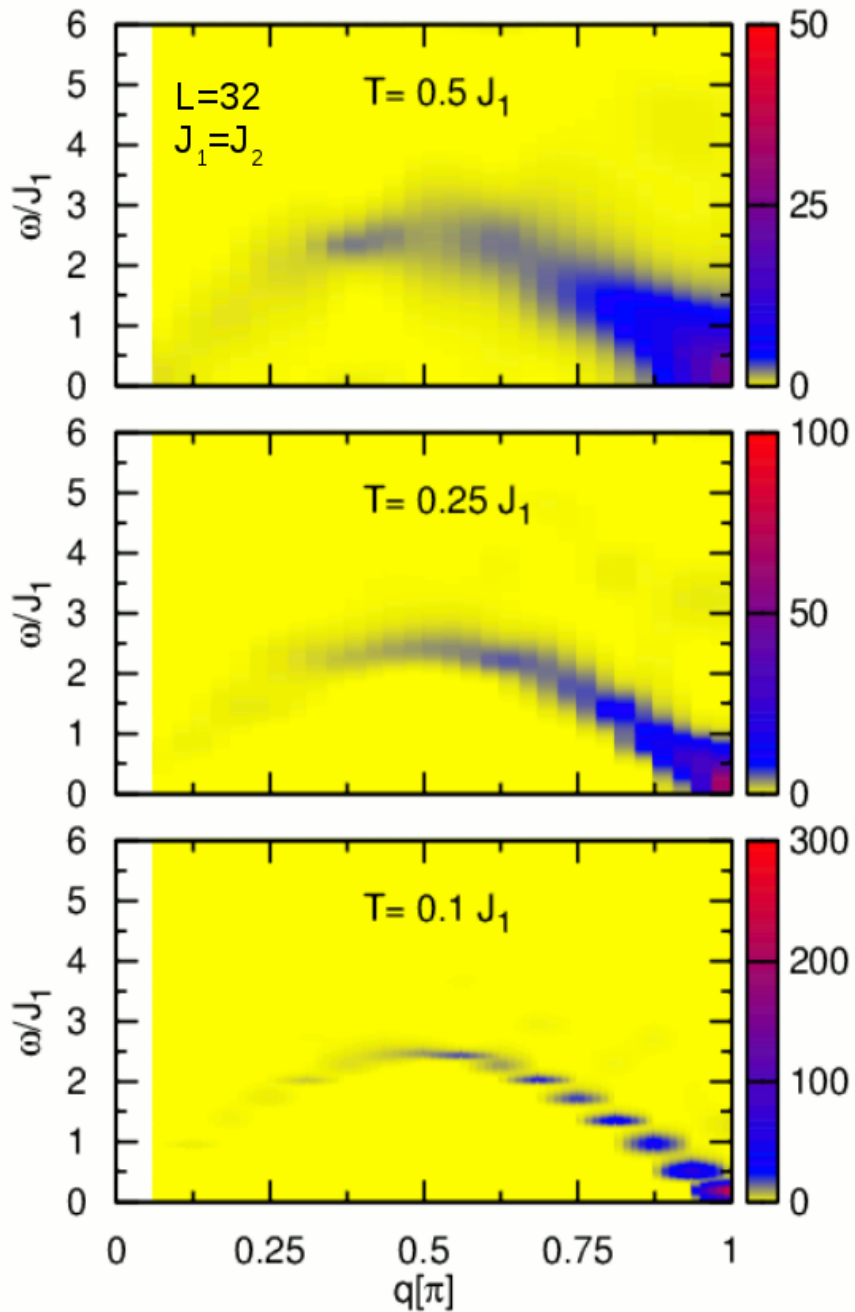
$J_0=0$   
 $J_2/J_1=1$ ,  $(\pi, 0)$  S~2-chain: magnons  
 $J_2/J_1=0$ , 4 S-1/2-chain: spinons

## Peak shift ~ gap closure

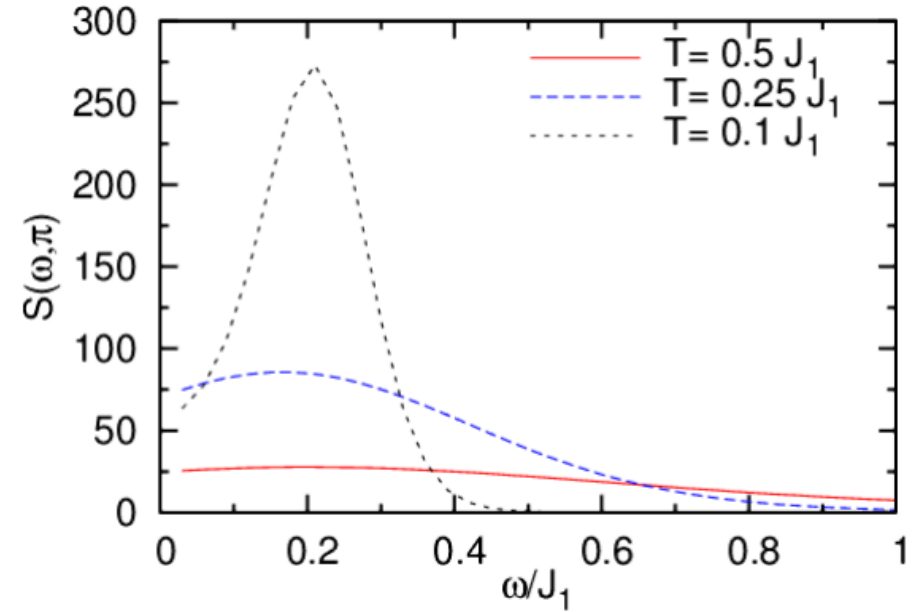


# Lineshape vs Temperature

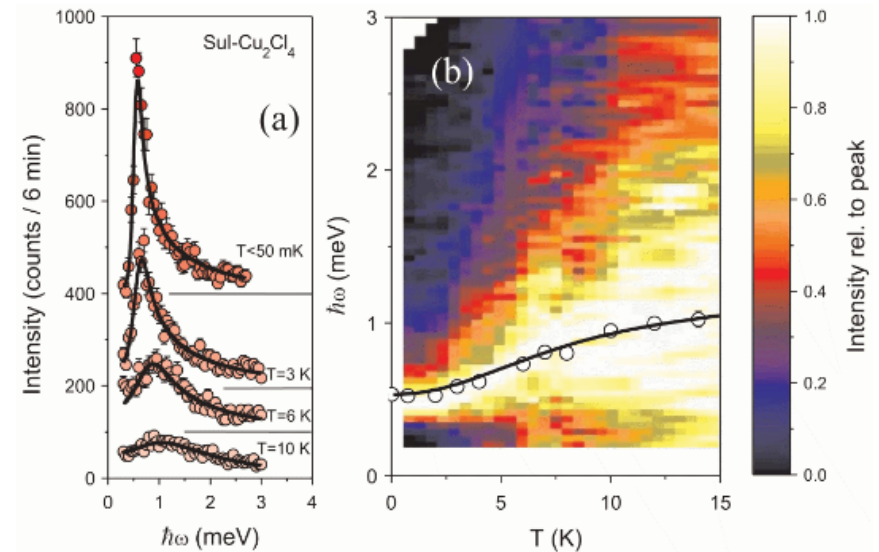
## Dynamic structure



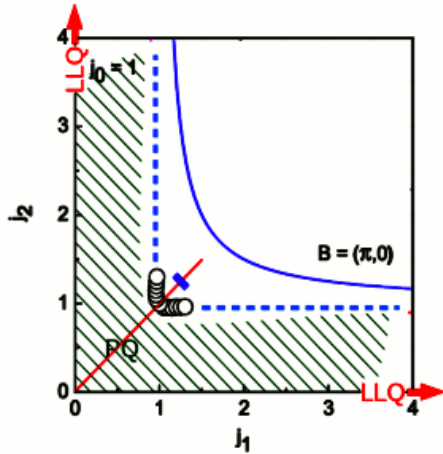
## Cut at $q=\pi$



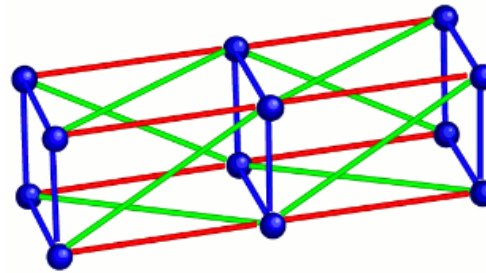
## ... tempting



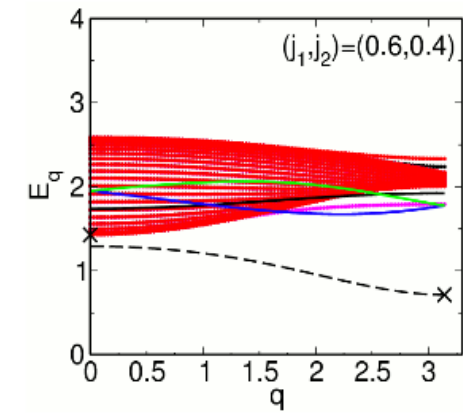
# Conclusions



- Plaqueette, incommensurate, Luttinger, and Haldane phases



- Novel elementary excitations and bound states



- Spinon-magnon crossover

